

Title: How do supermassive black holes get into galaxies?

Date: Dec 09, 2009 02:00 PM

URL: <http://pirsa.org/09120013>

Abstract: I will give an account of our current understanding of the formation and growth of the central supermassive black holes in galaxies from an astrophysical perspective.

# How do supermassive black holes get into galaxies

Martin Haehnelt



M,J



$$ds^2 = \frac{\Delta}{\rho^2} (cdt - a \sin^2 \theta d\phi)^2 - \frac{\sin^2 \theta}{\rho^2} [(r^2 + a^2) d\phi - a cdt]^2 - \frac{\rho^2}{\Delta} dr^2 - \rho^2 d\theta^2,$$

where  $\Delta$  and  $\rho$  are

$$\left. \begin{aligned} \Delta &= r^2 - 2\mu r + a^2, \\ \rho^2 &= r^2 + a^2 \cos^2 \theta. \end{aligned} \right\}$$



stellar mass black holes

1-100  $M_{\odot}$

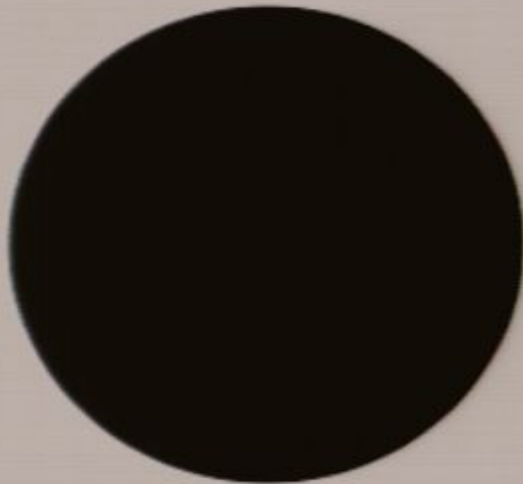


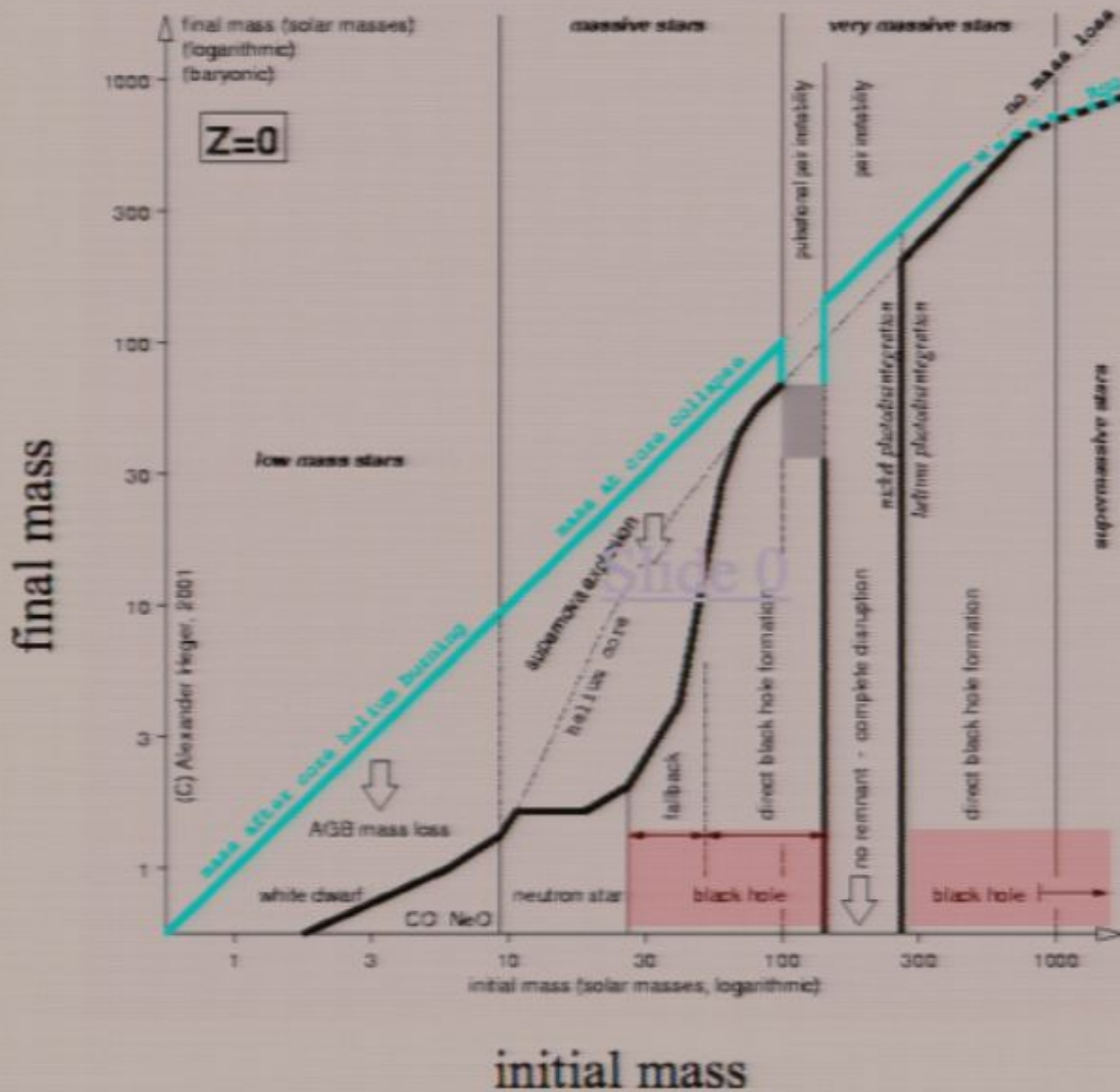
Do intermediate mass bh exist ?



supermassive black holes

$3 \times 10^5 - 3 \times 10^{10} M_{\odot}$





The end-state of massive stars is a black hole.





stellar mass black holes

$1-100 M_{\odot}$

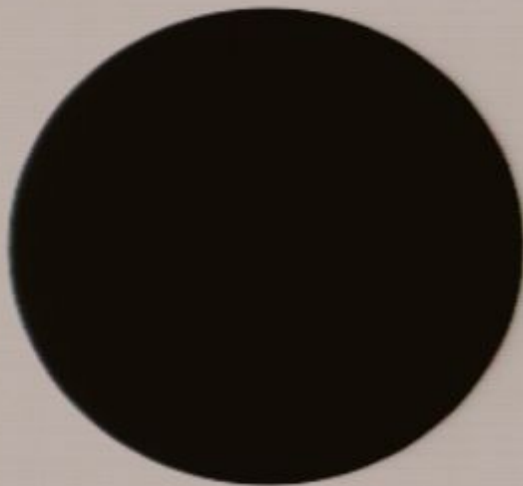


Do intermediate mass bh exist ?

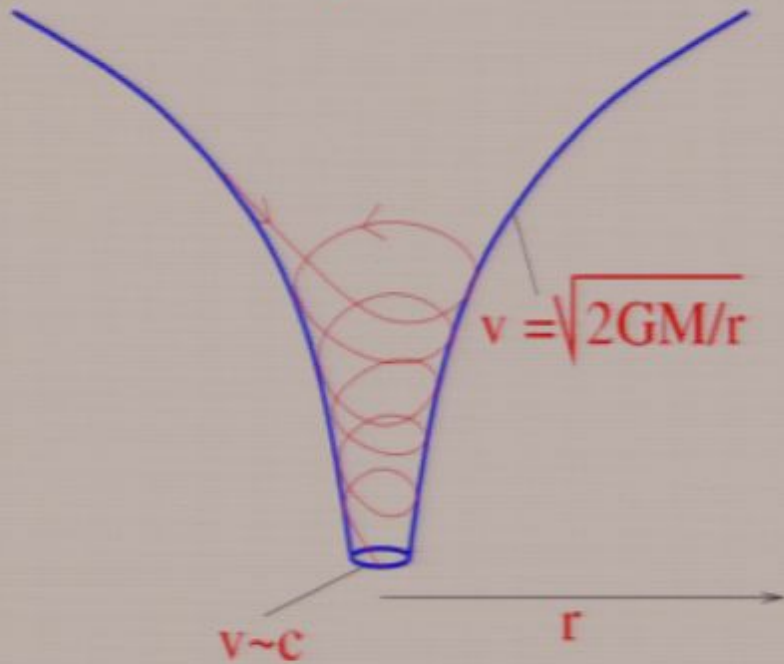
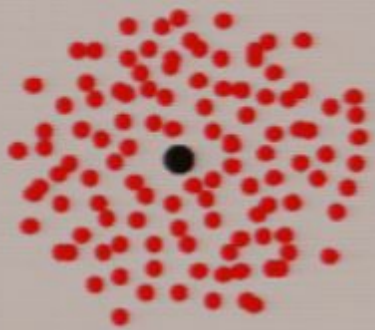


supermassive black holes

$3 \times 10^5 - 3 \times 10^{10} M_{\odot}$



# Nearby galaxies

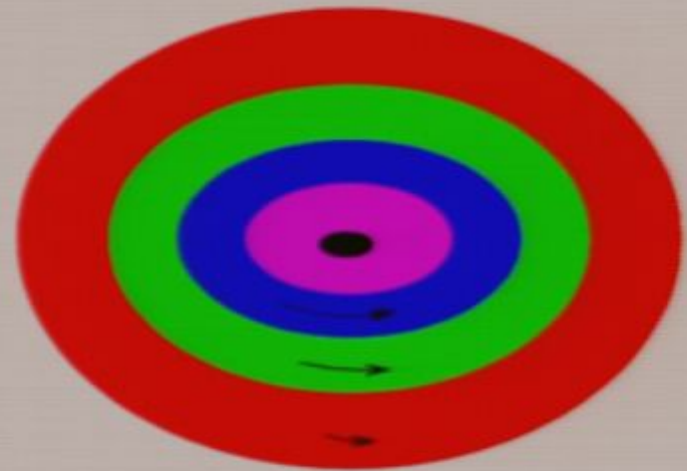


# QSOs

Accretion disc



Quasar



$$T \sim 100\,000\text{ K}$$

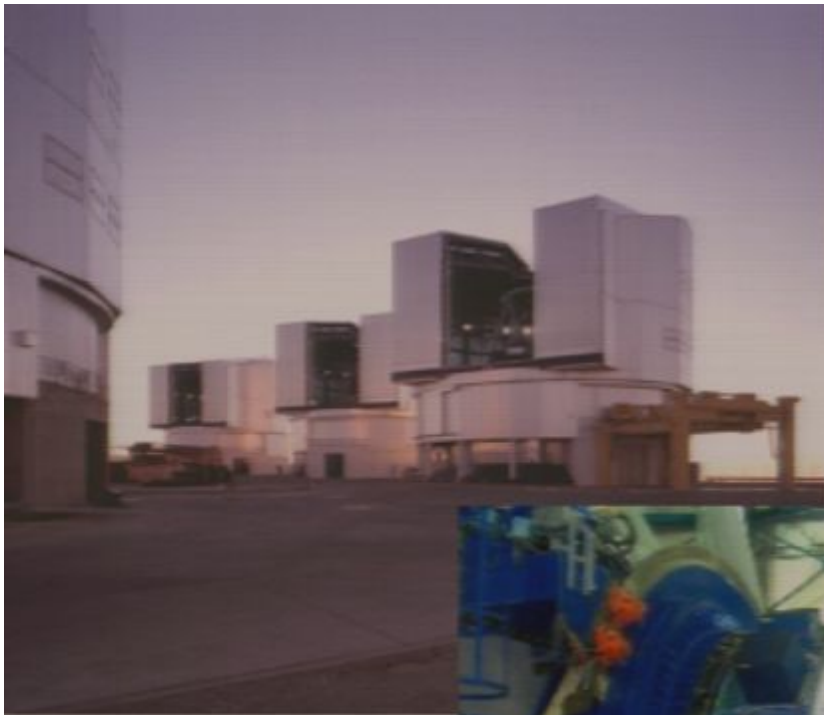
$$L \sim 0.1 \dot{M} c^2$$



# The supermassive black hole at the galactic centre

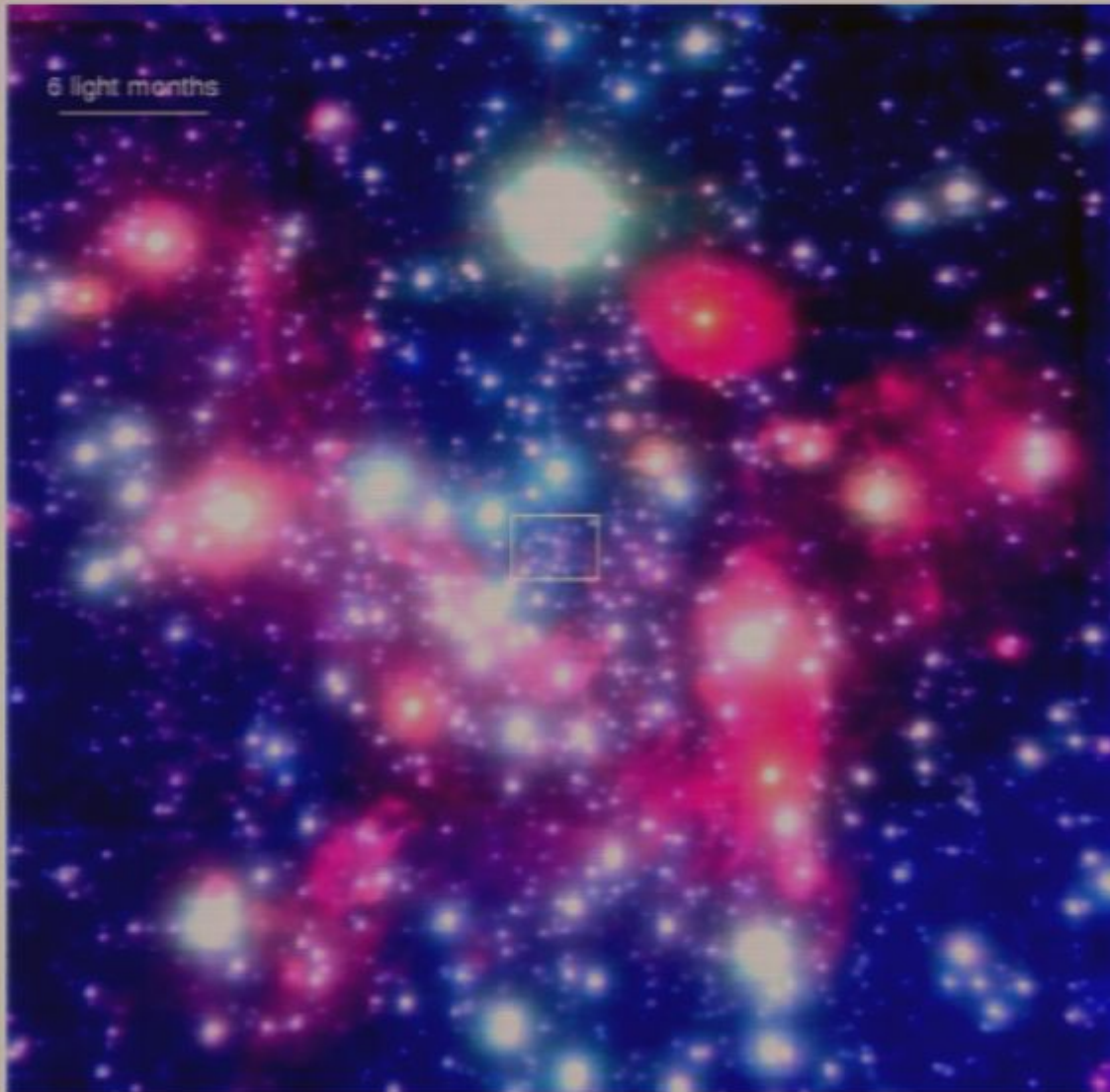






# NACO@VLT



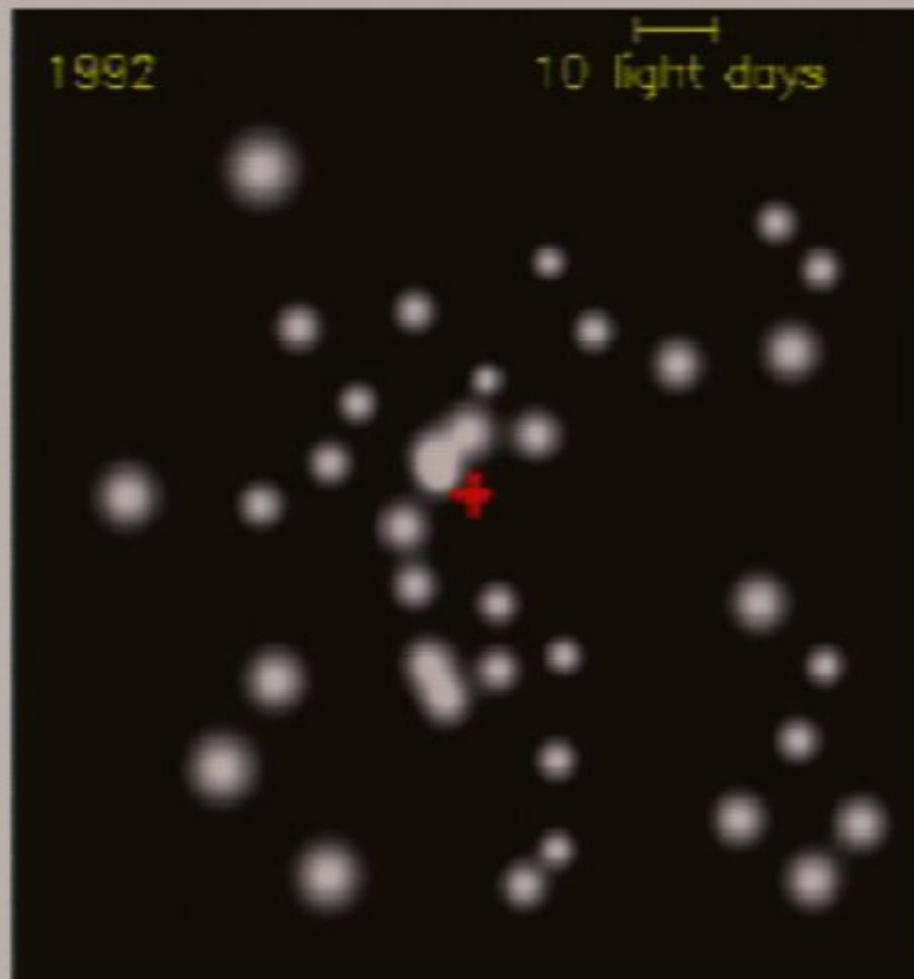


Piercing through the dust  
in the galactic centre at  
IR wavelengths.

# The black hole at the centre of our galaxy

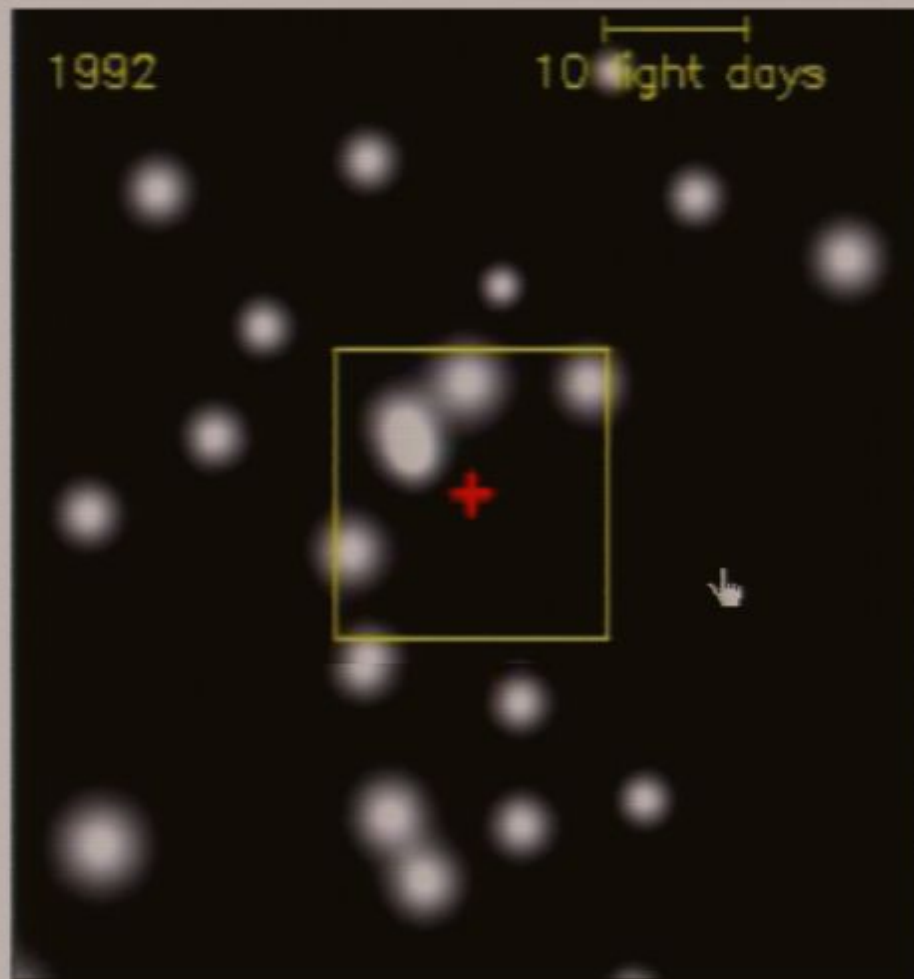


# The black hole at the centre of our galaxy



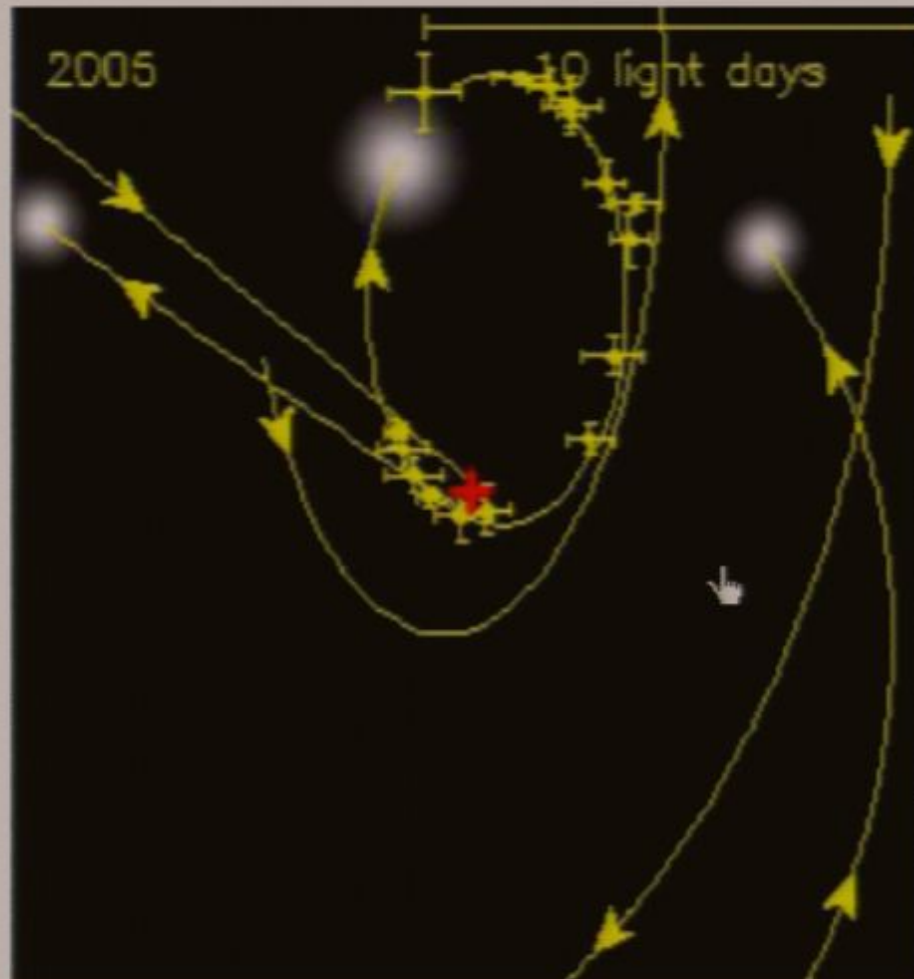


# The black hole at the centre of our galaxy

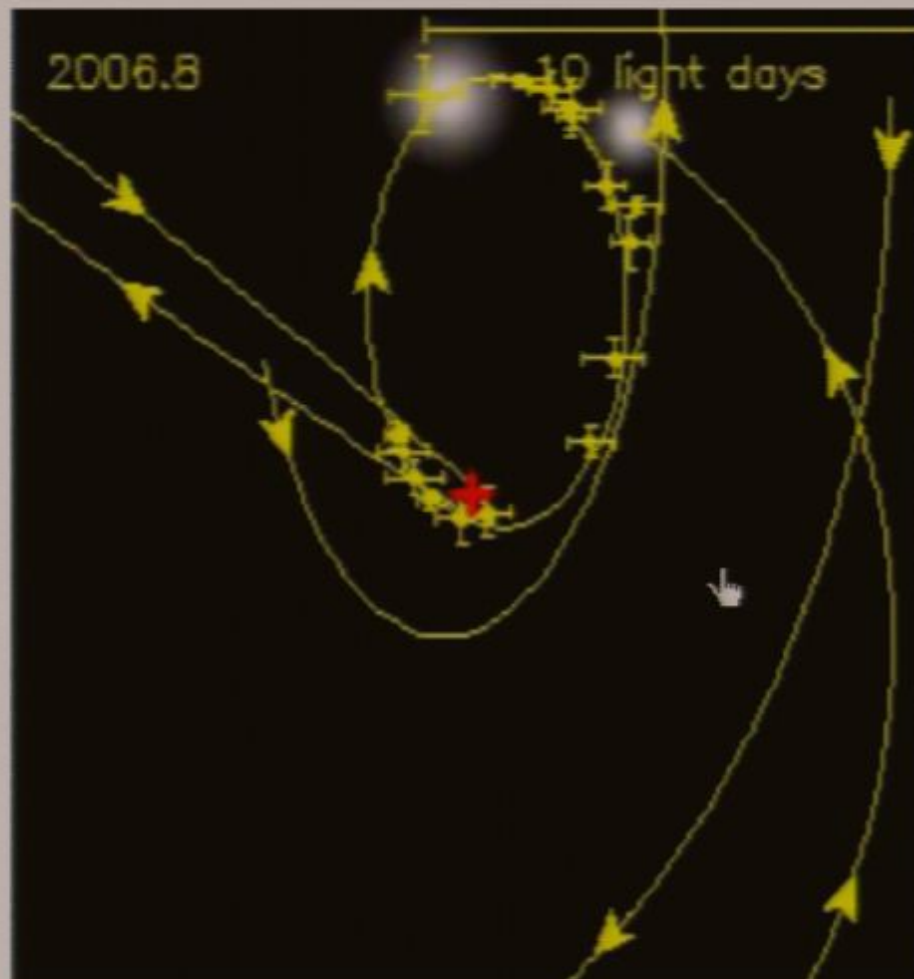


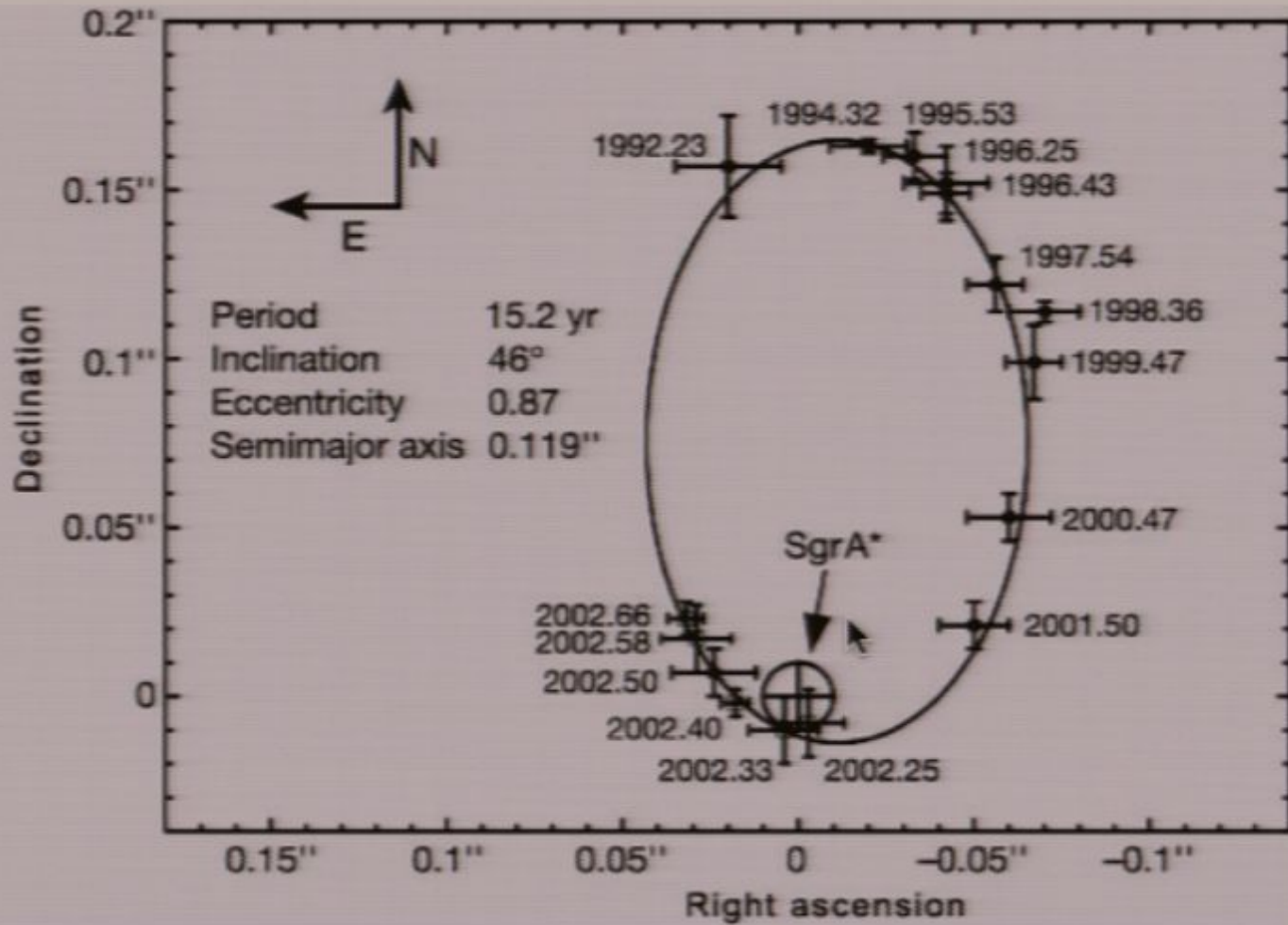


# The black hole at the centre of our galaxy

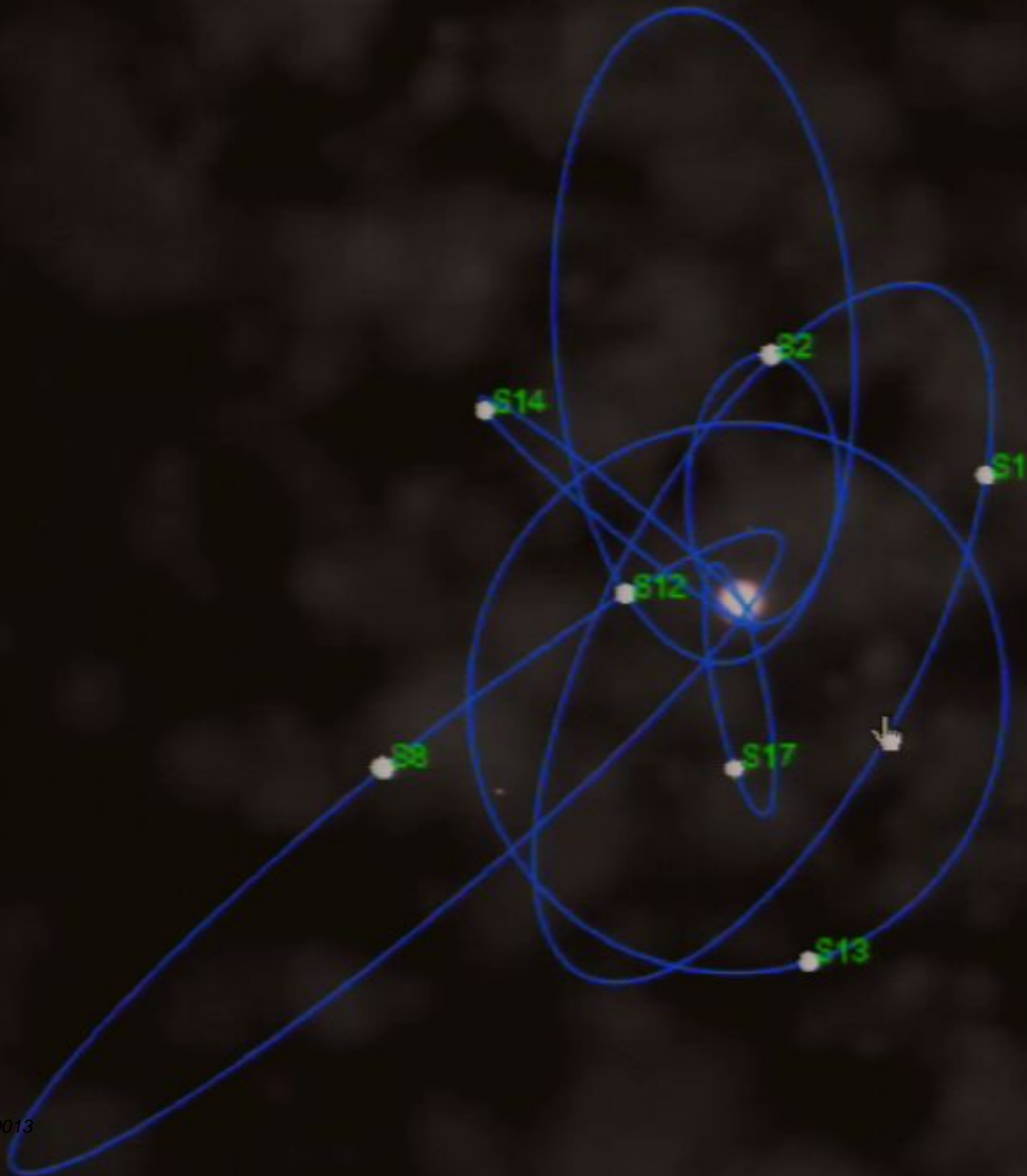


# The black hole at the centre of our galaxy

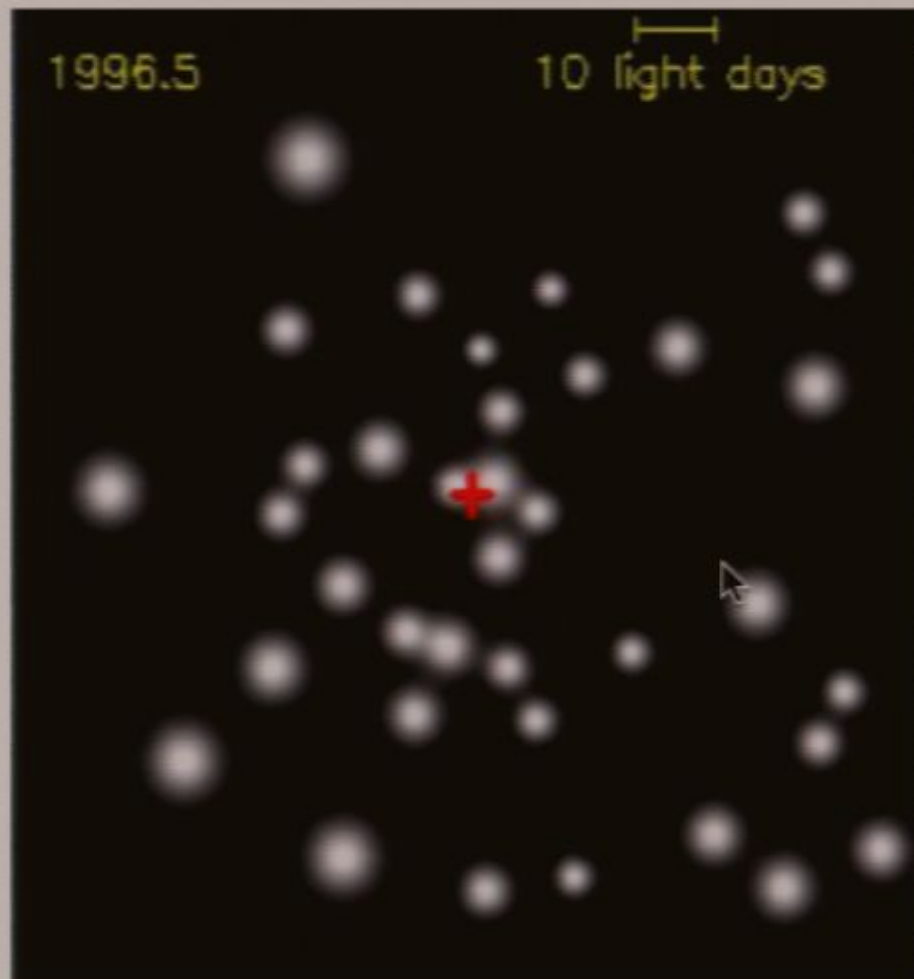




1993 09 09 13:58:59 UTC  
45000000x faster

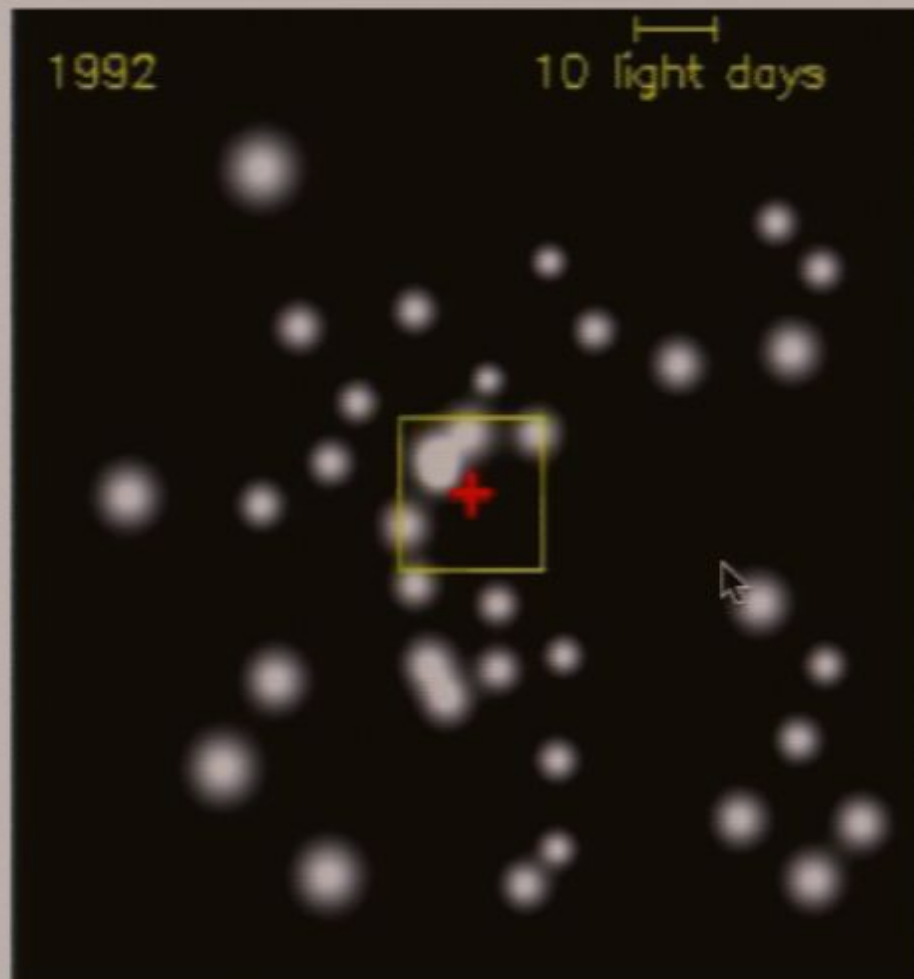


# The black hole at the centre of our galaxy

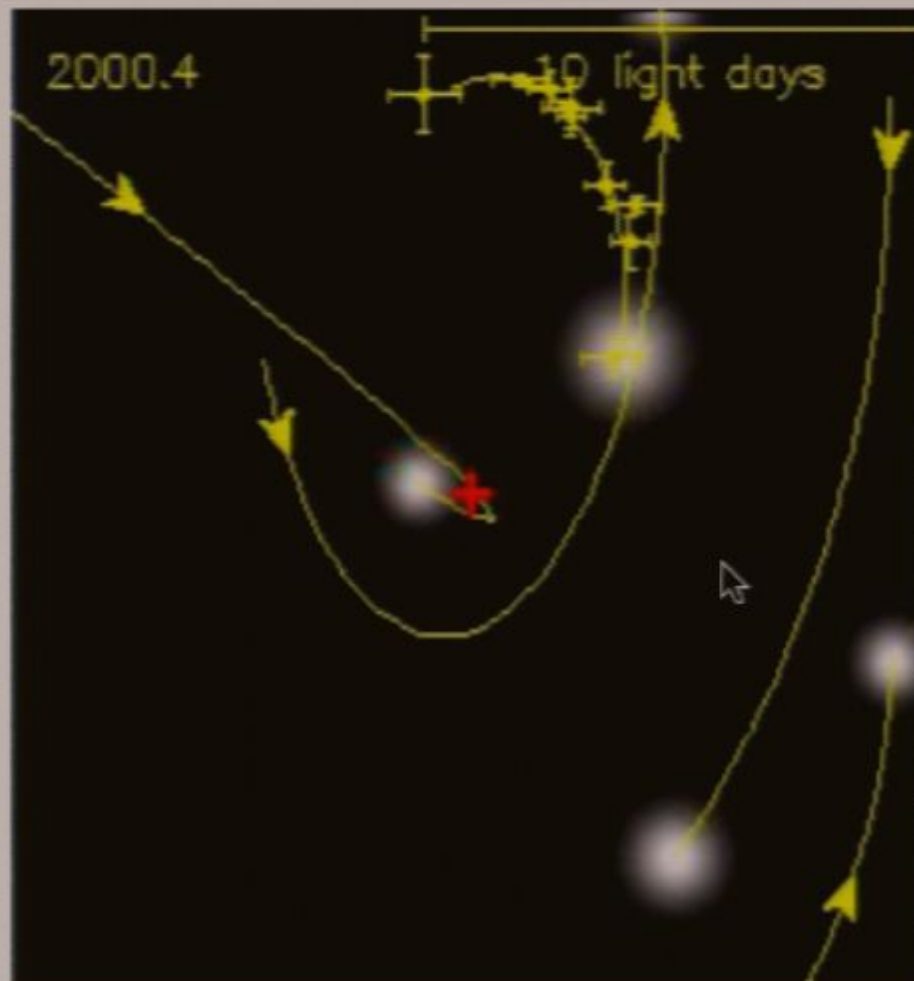




# The black hole at the centre of our galaxy

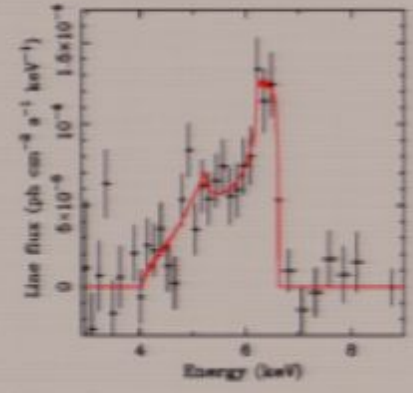


# The black hole at the centre of our galaxy

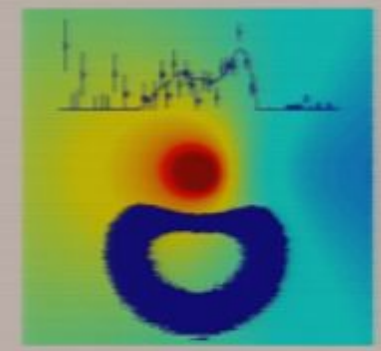


- ativistic redshifts
- SS survey
- A]
- A
- ical merging history in a bright  
mation and evolution of super  
ries
- binary black holes coalescenc
- le Interactions
- is the Rosetta stone?

### Relativistic redshifts



Tanaka et al 1995



Reynolds et al 2000

Perimeter, 9<sup>th</sup> December 2009

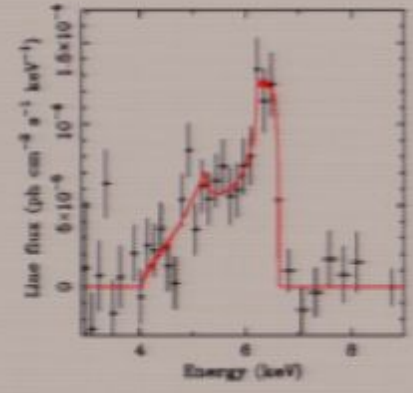


Click to add notes

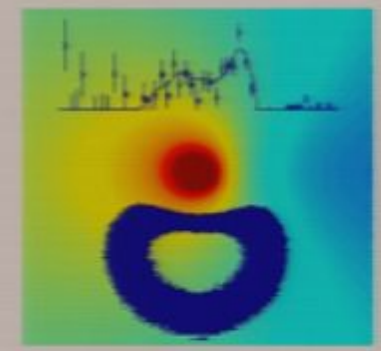


- ativistic redshifts
- 3S survey
- ical merging history in a ht elliptical
- mation and evolution of ermassive binaries
- binary black holes
- lescence?
- le Interactions
- is the Rosetta stone?

# Relativistic redshifts



Tanaka et al 1995

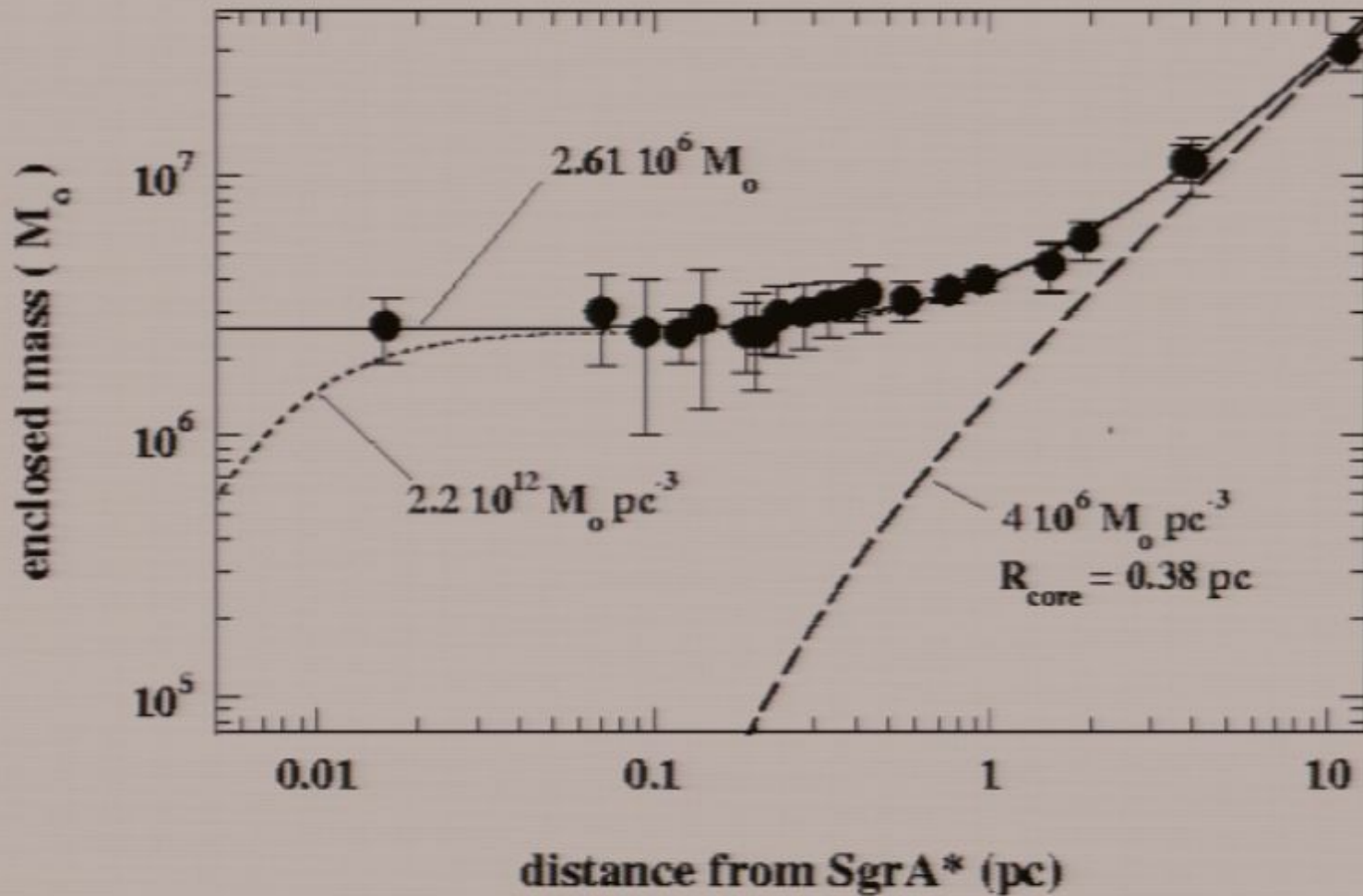


Reynolds et al 2000



Click to add notes



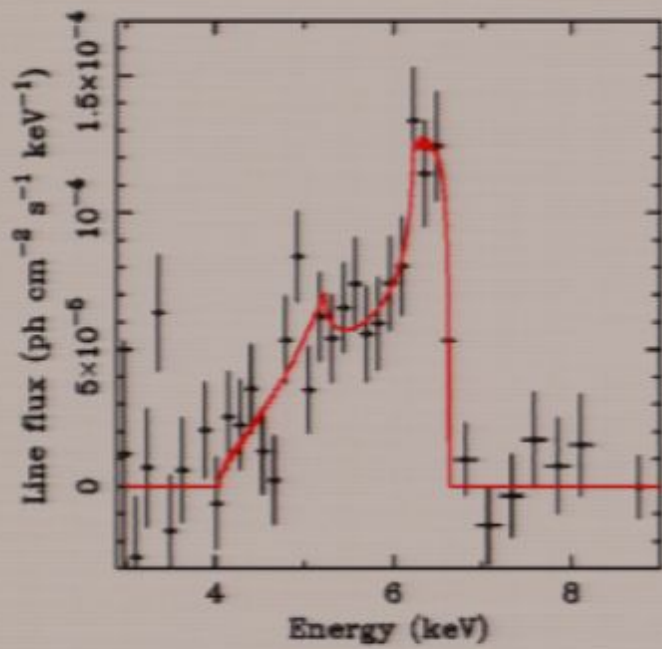


For astrophysicists the presence of a black hole with mass several million that of the sun has been established beyond reasonable doubt.

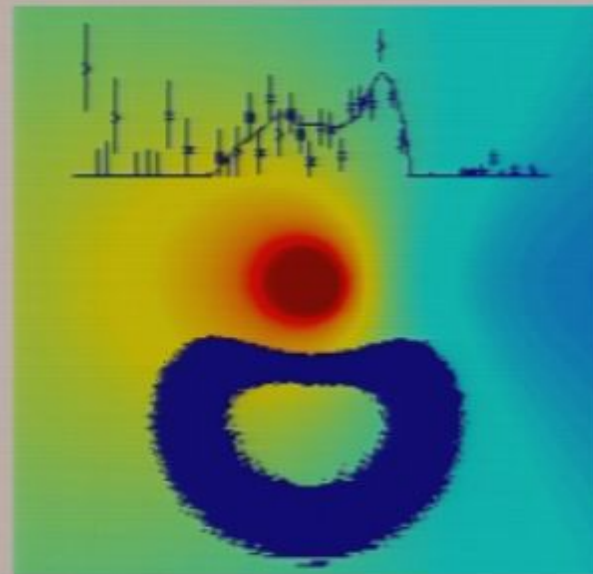




# Relativistic redshifts



Tanaka et al 1995

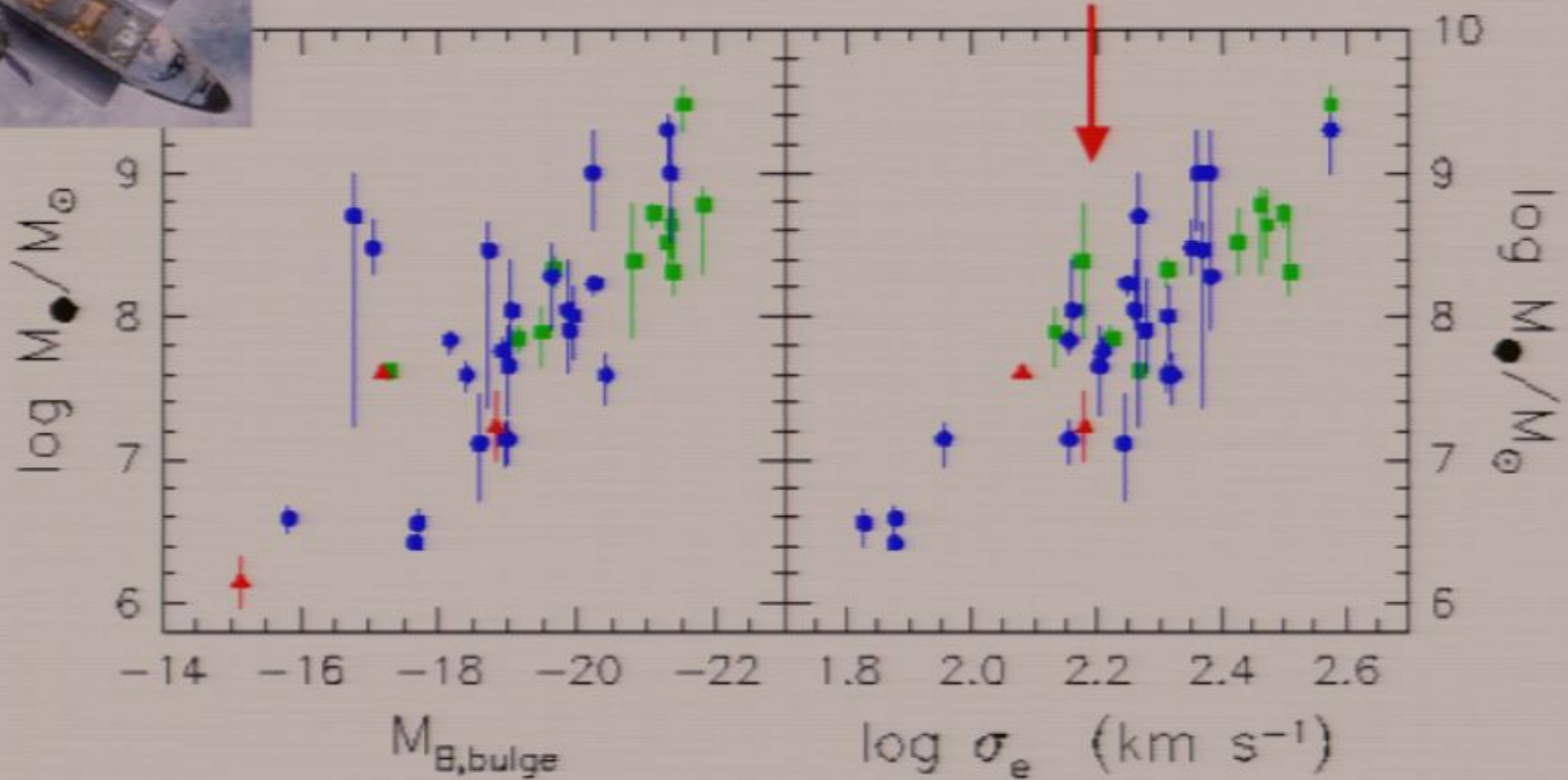


Reynolds et al 2000





Is this really tighter?



black hole mass scales with 
/ bulge mass Kormendy 2003  
\ stellar velocity dispersion of the bulge



# Massive black holes?

Giant Ellipticals/S0s



Yes

Spirals



Yes but black hole mass scales with bulge mass not total mass

Dwarfs



~~Possibly~~

A few

Globular Clusters



Maybe



# SDSS survey



# SDSS survey



1 million galaxy  
spectra

100 000 QSOs

<http://www.sdss.org>



A very large,  
homogeneous  
high-quality  
data set.



early-type



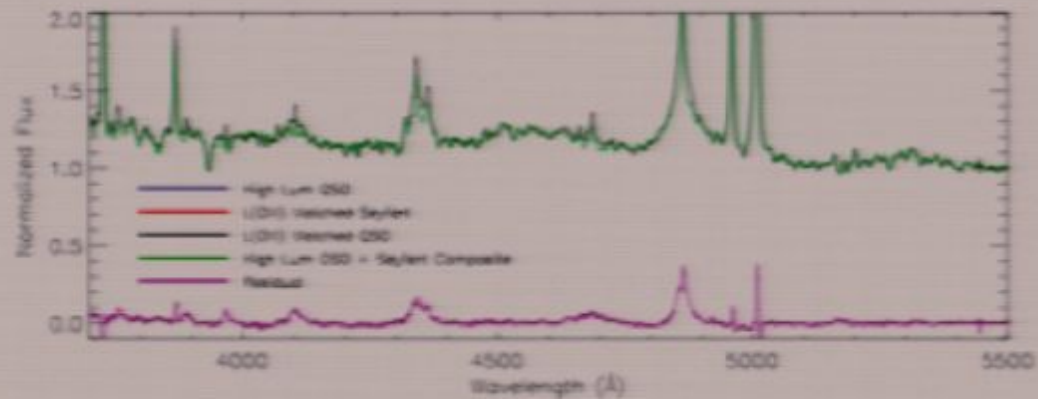
late-type

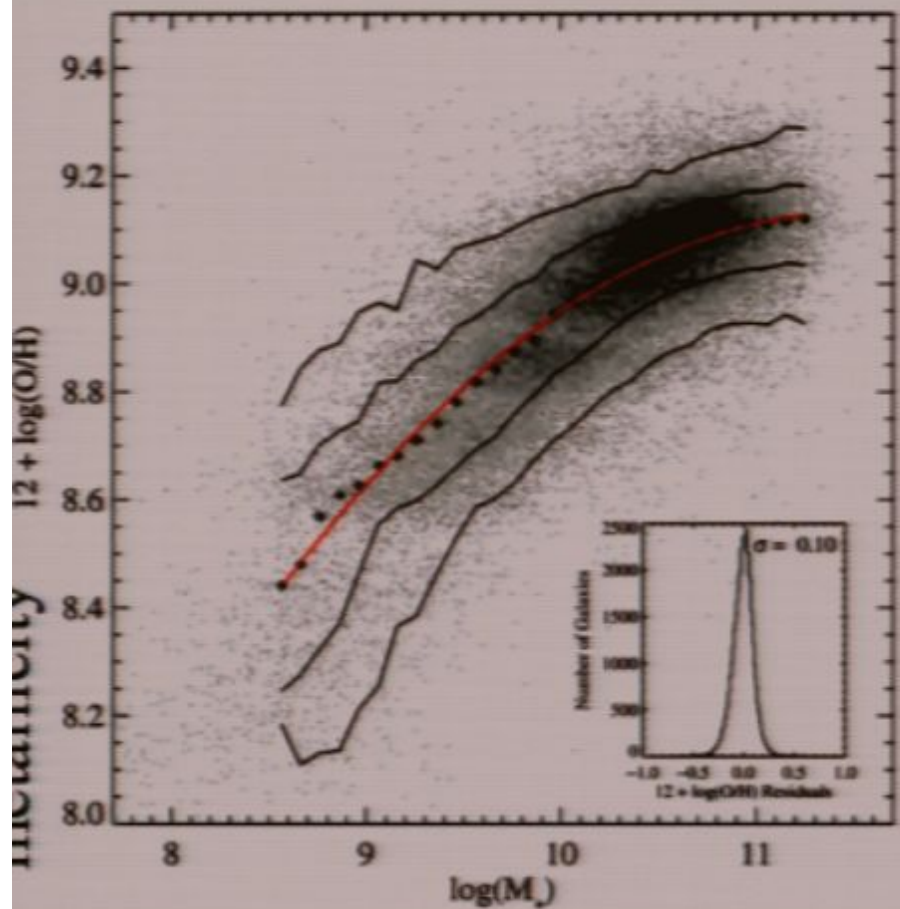


early-type with massive companion



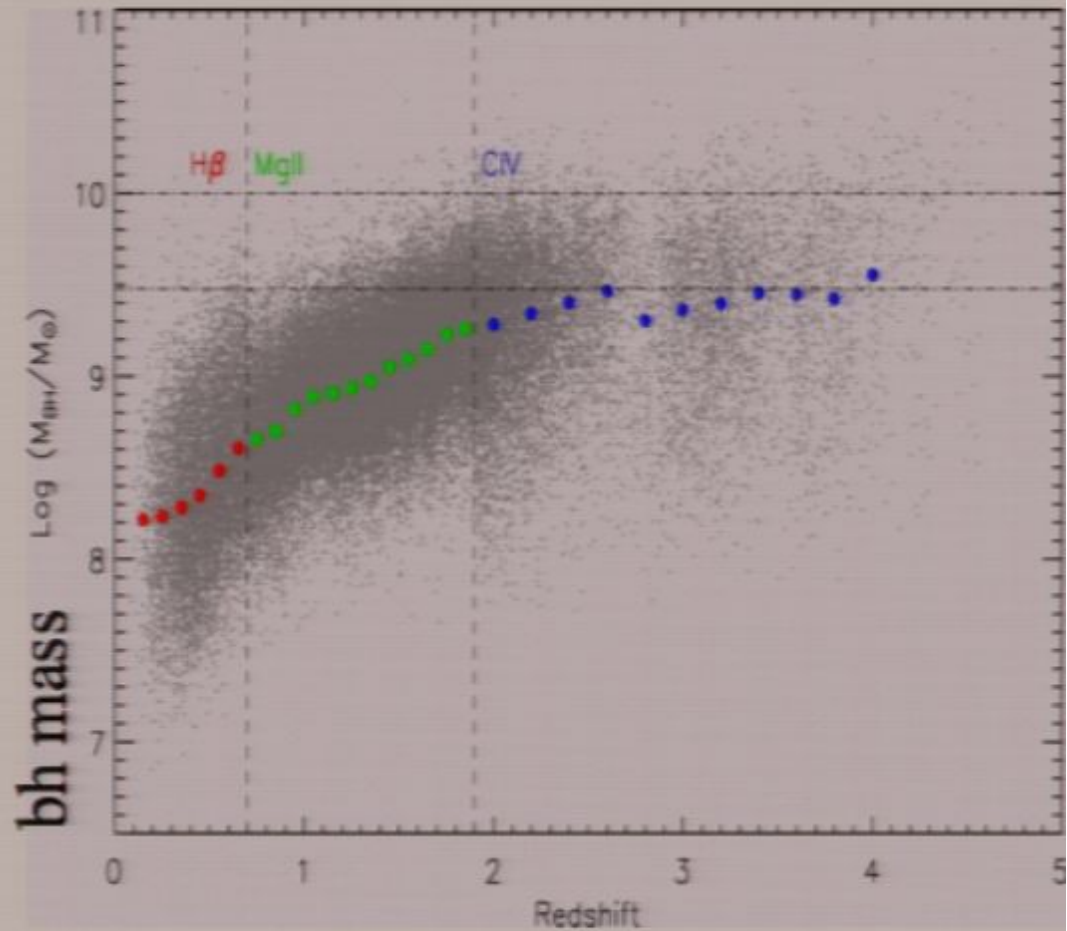
late-type with low mass companion





stellar mass

Tremonti et al.

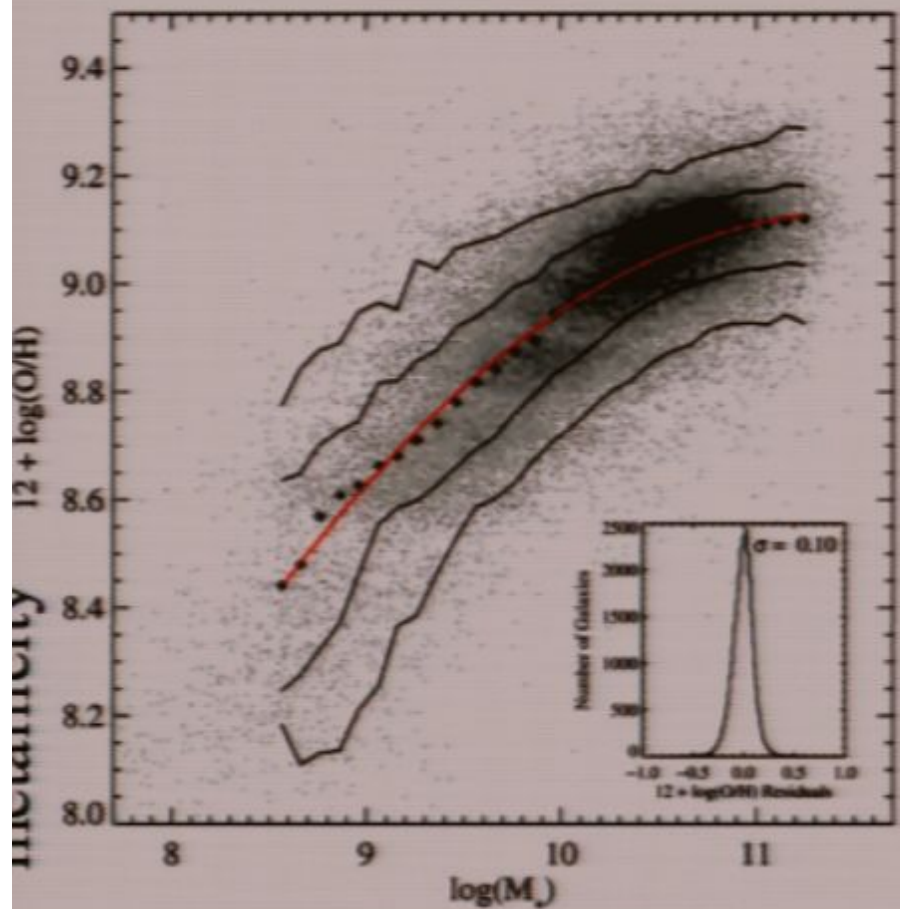


Shen et al.



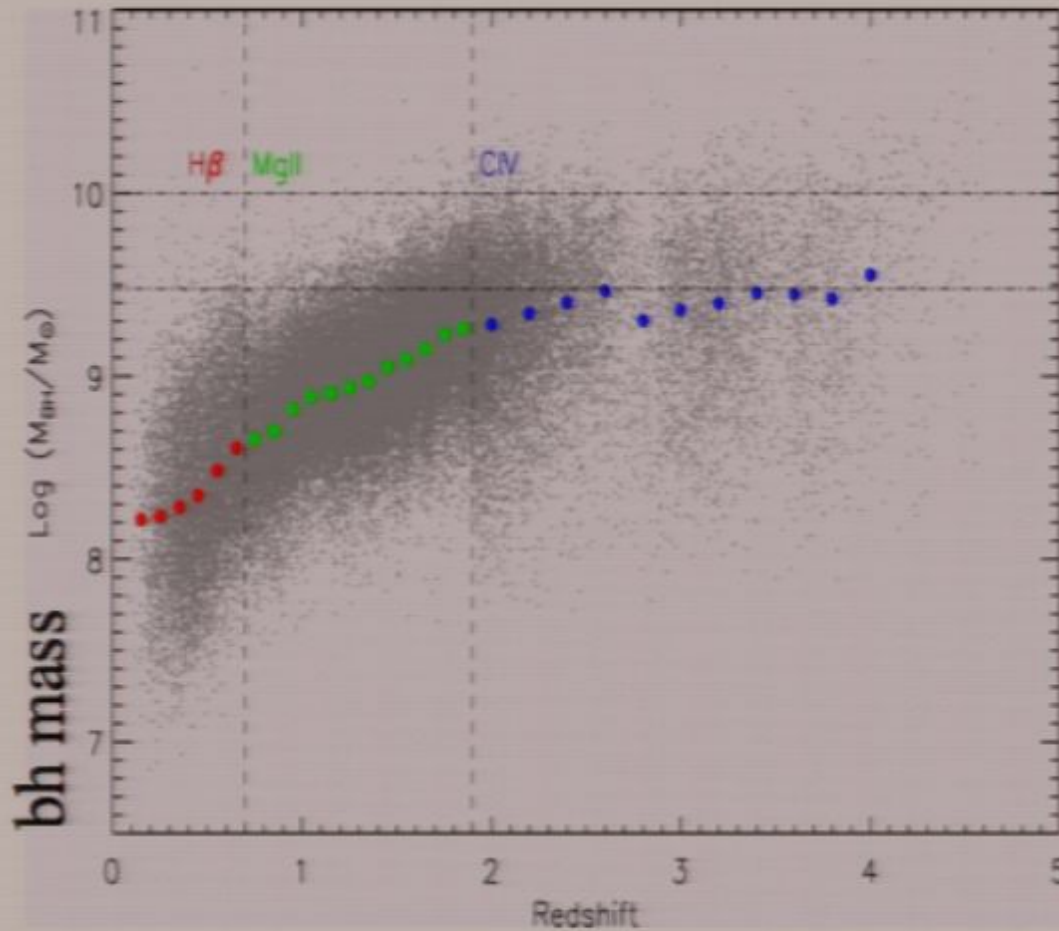
# Co-evolution of galaxies and black holes





stellar mass

Tremonti et al.

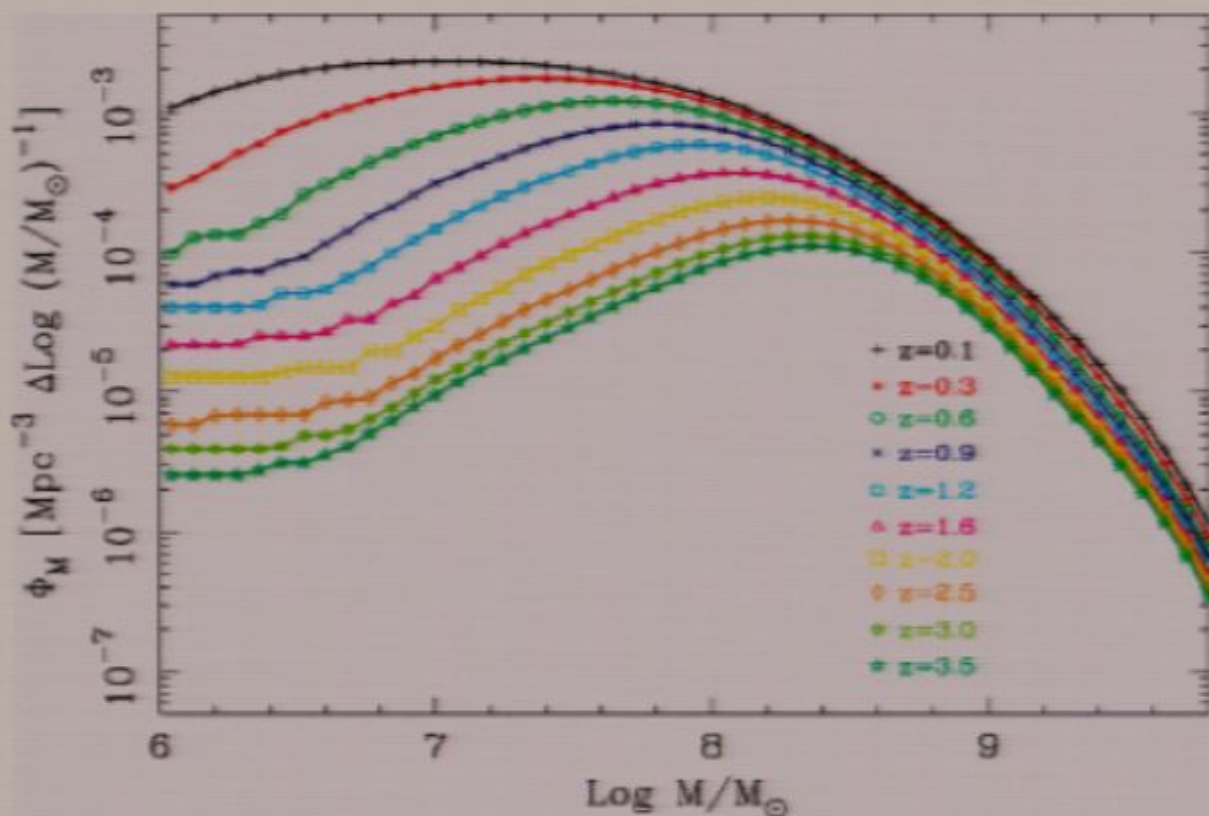


Shen et al.





## Evolution of the black hole mass function as reconstructed from accretion history

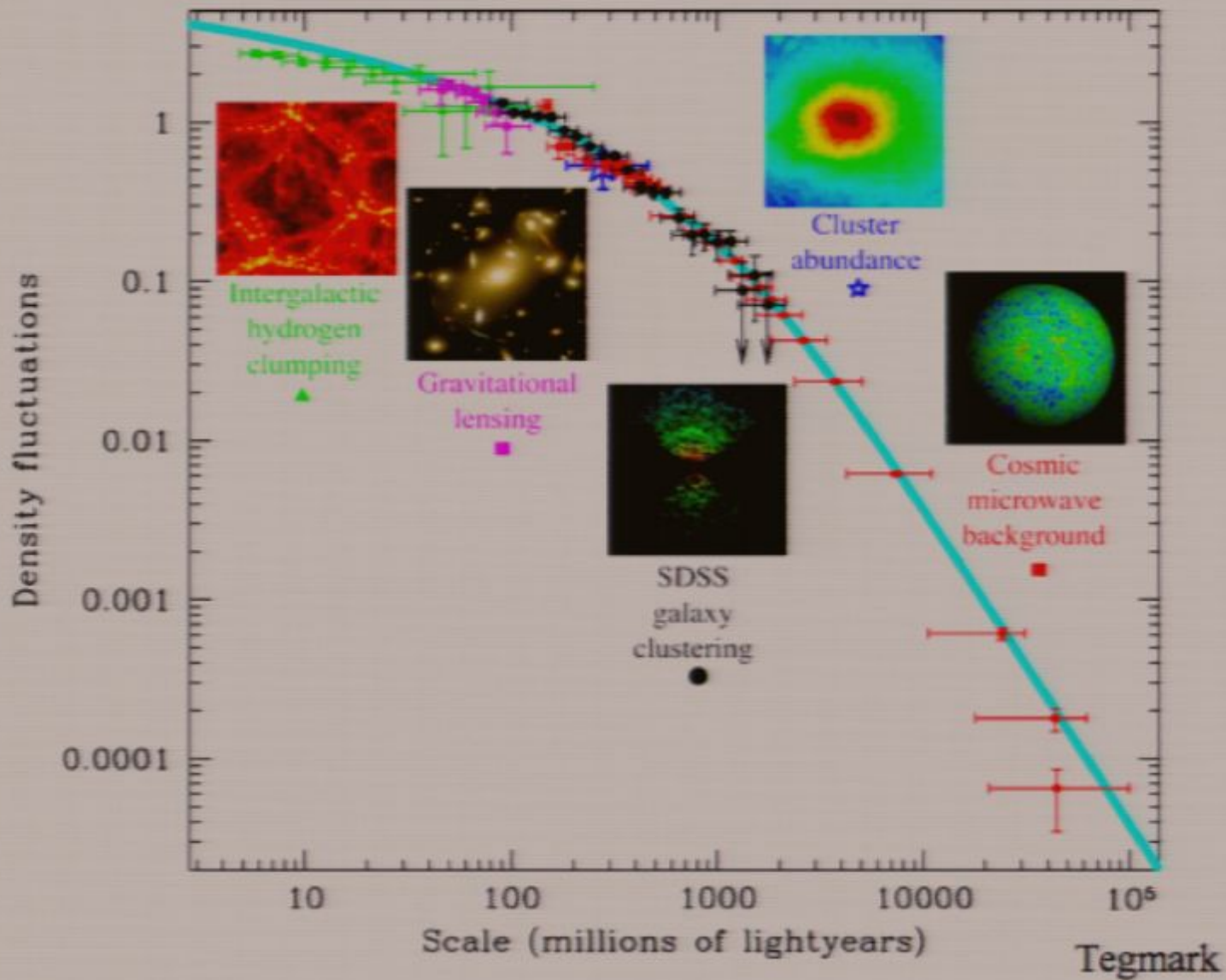


Merloni 2004

### Cosmic Downsizing

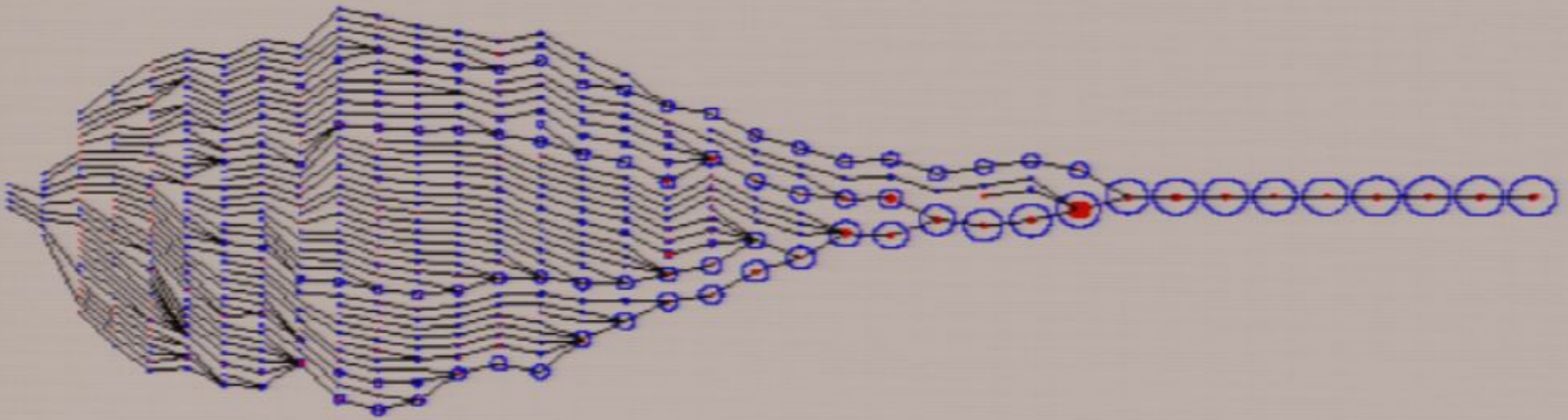
Most of the big black holes form at high redshift. Most of the black holes still growing at the present day are small. These small black holes have gained most of their mass recently.





**We know the DM power spectrum very well!**

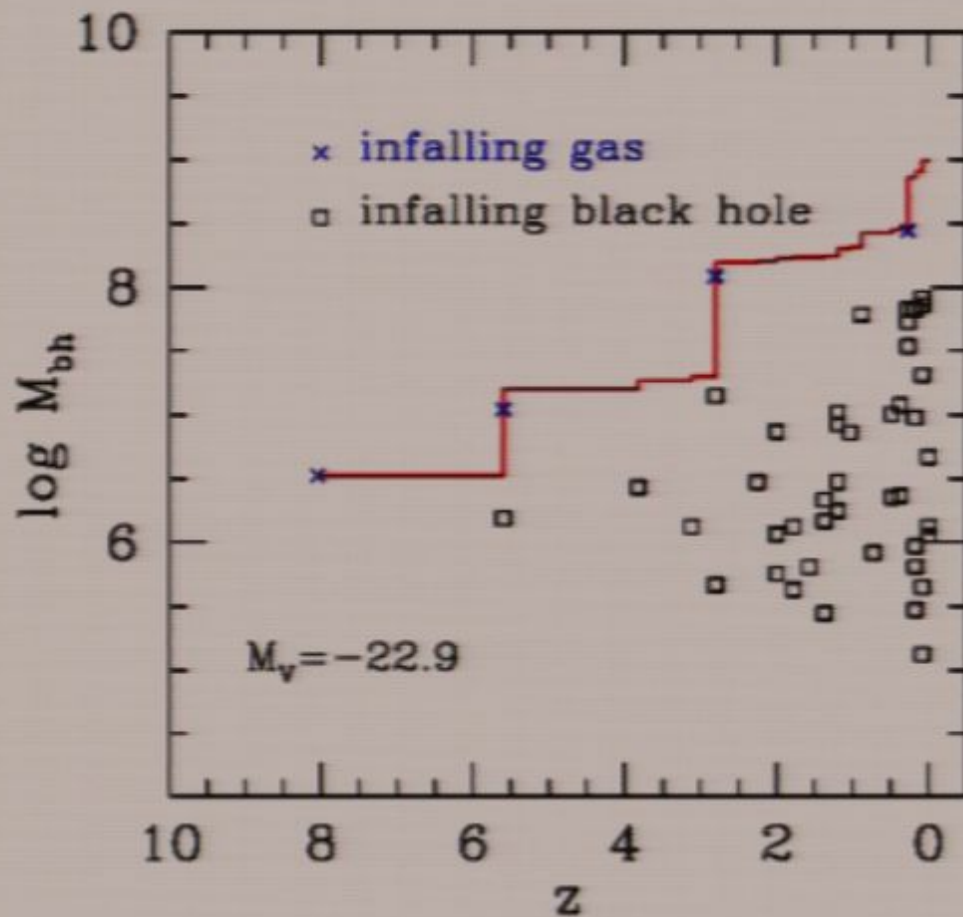




Bullock et al.

DM haloes grow by merging

# Typical merging history in a bright elliptical



Kauffmann & Haehnelt 2002

Growth of black holes dominated by accretion of gas and feedback regulated

At late time also growth by infall and merging of black holes

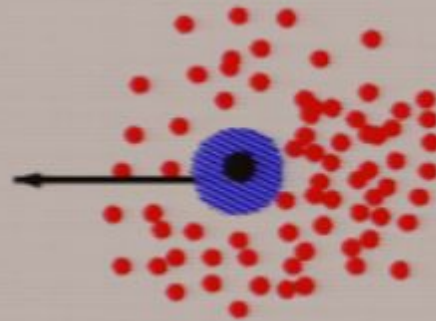
How do  $10^6 M_{\text{sol}}$  black holes form?



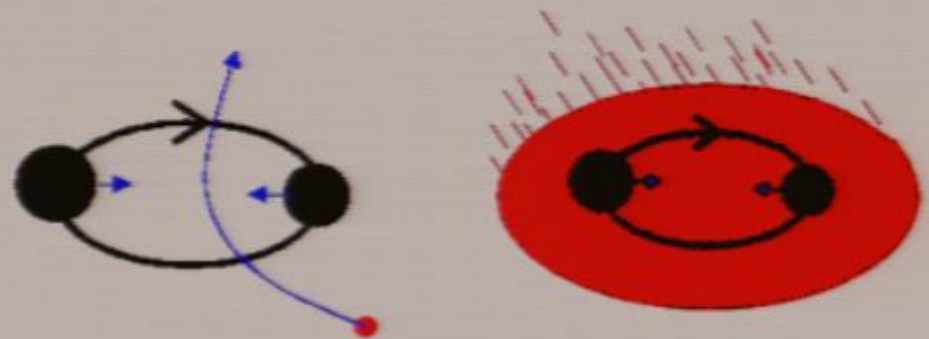
# Formation and evolution of supermassive binaries

## 1. Dynamical friction

$$t \propto a$$

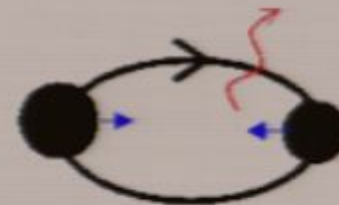


## 2. Binary hardening due to stars or accretion of gas



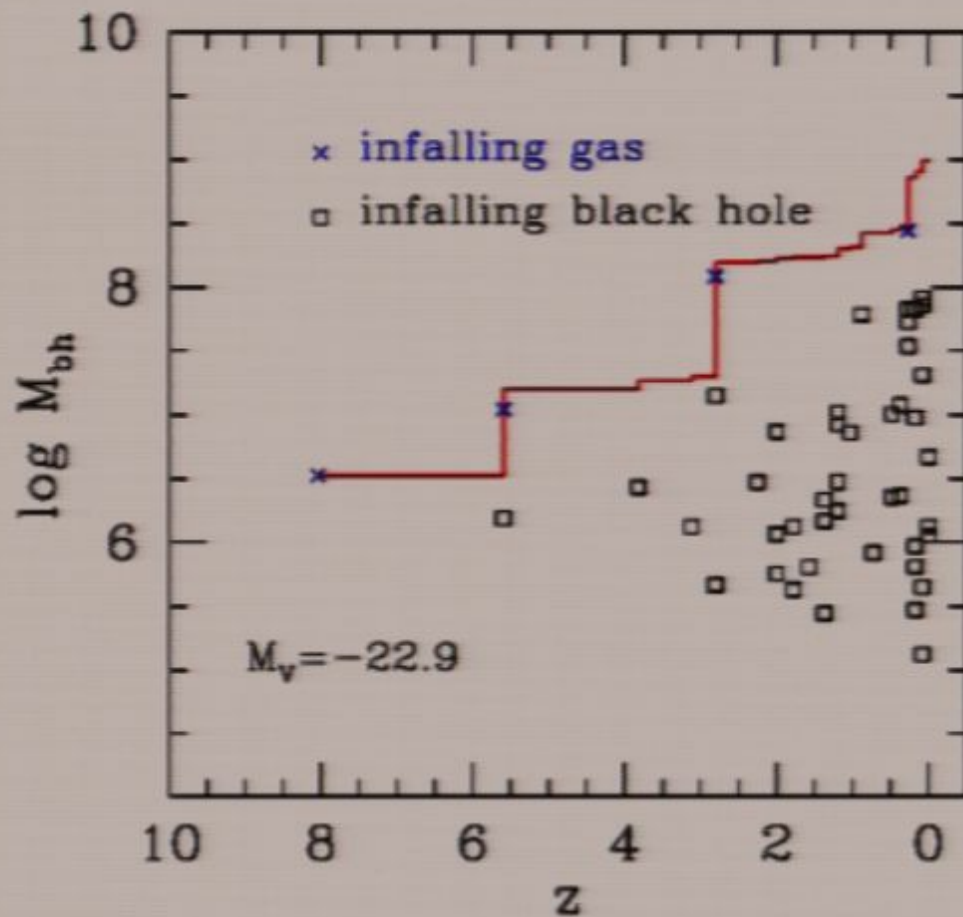
## 3. Gravitational radiation

$$t \propto a^4$$





# Typical merging history in a bright elliptical

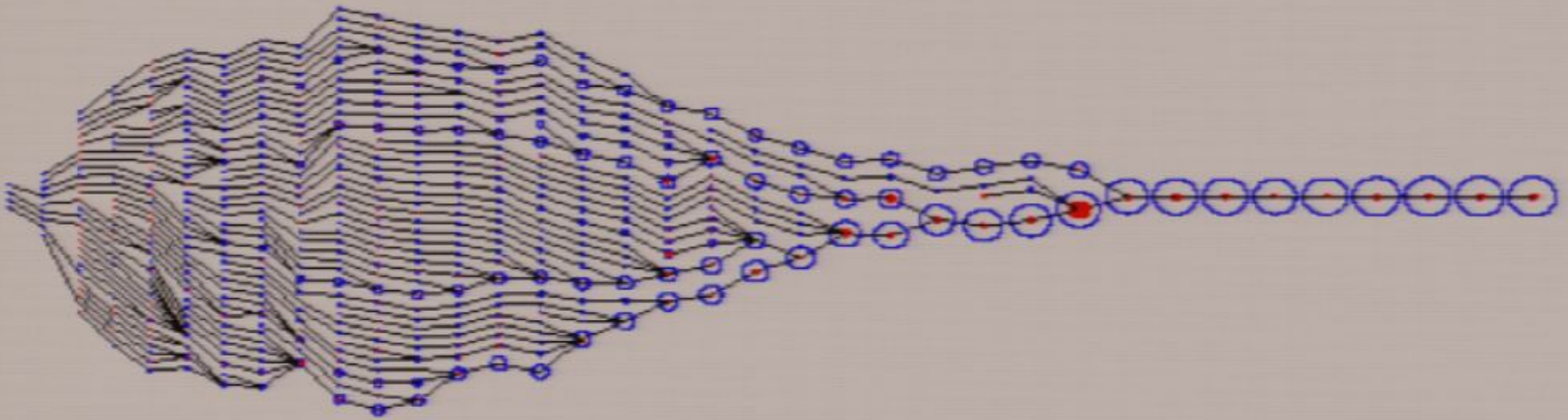


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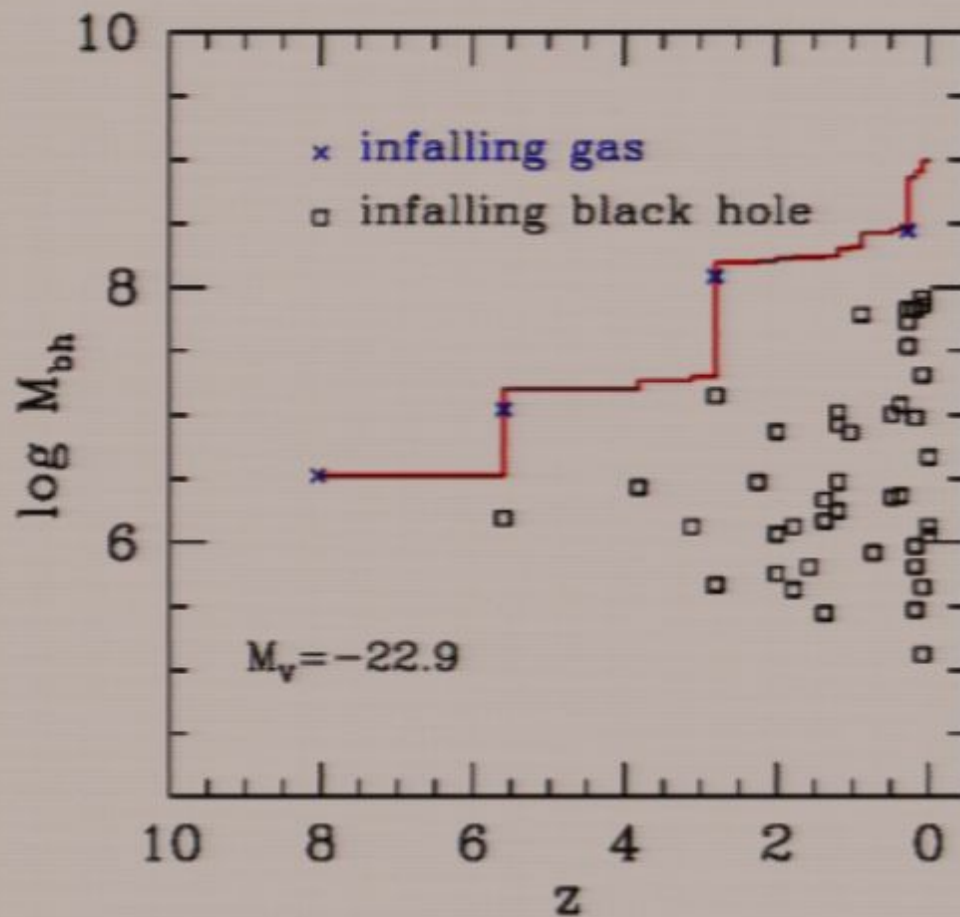
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Bullock et al.

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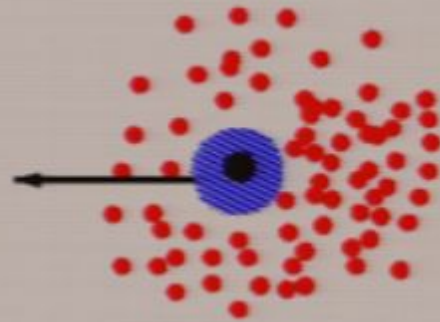
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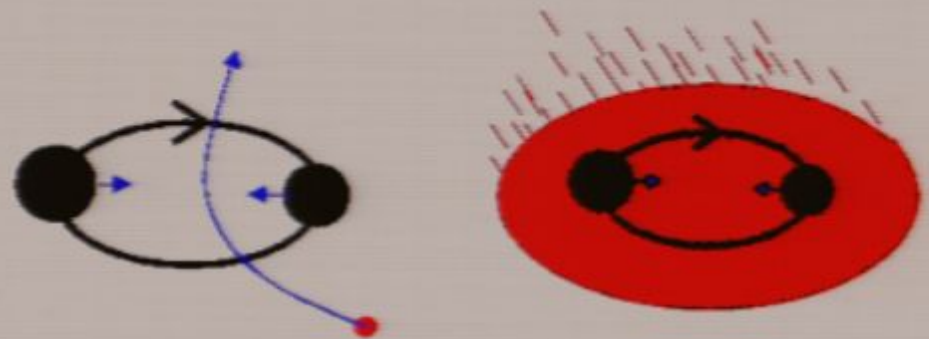
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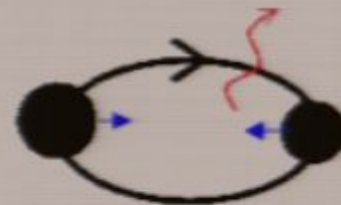


## 2. Binary hardening due to stars or accretion of gas



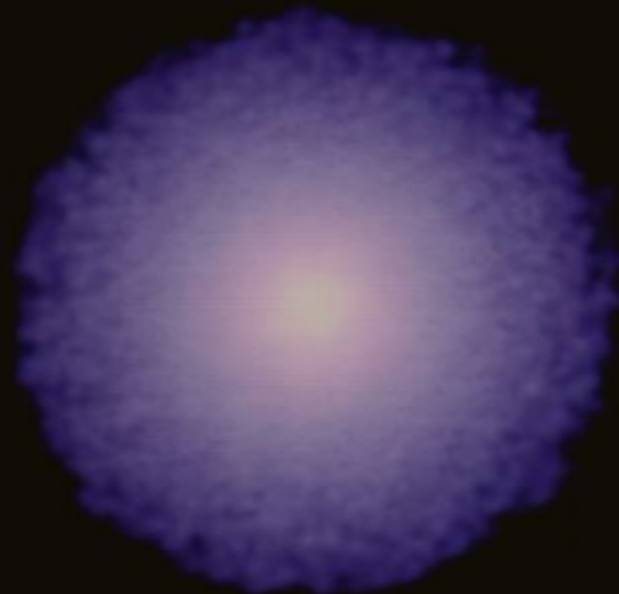
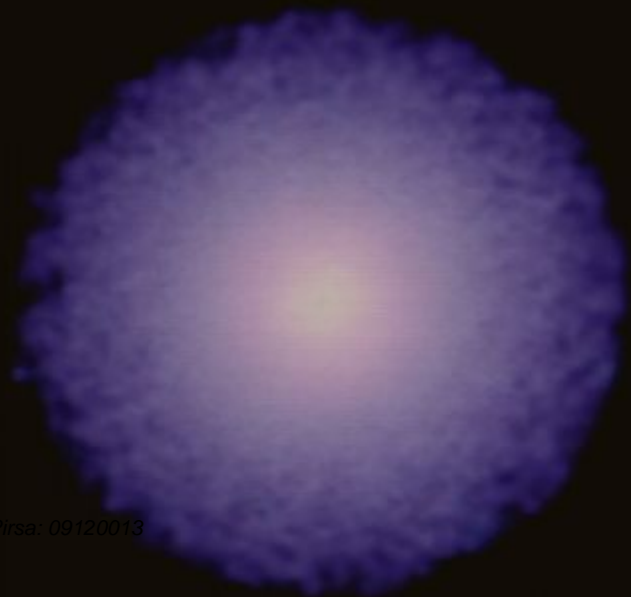
## 3. Gravitational radiation

$$t \propto a^4$$





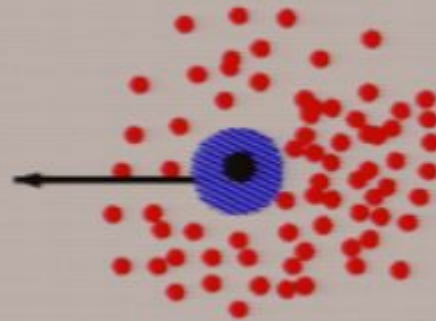
$\Gamma = 10 \text{ Myr}$



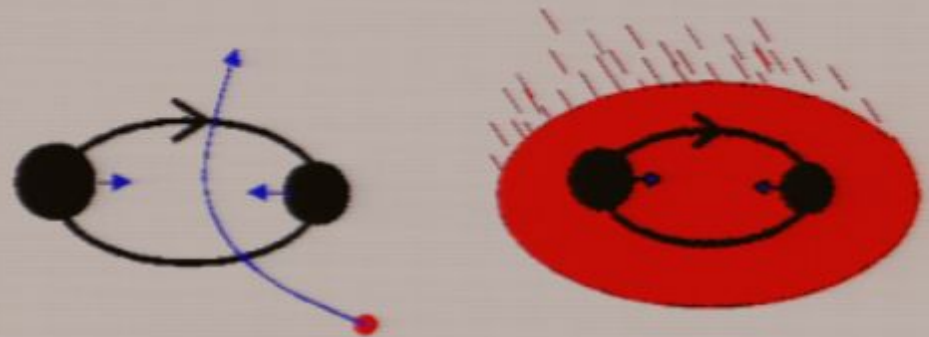
# Formation and evolution of supermassive binaries

## 1. Dynamical friction

$$t \propto a$$



## 2. Binary hardening due to stars or accretion of gas



## 3. Gravitational radiation

$$t \propto a^4$$



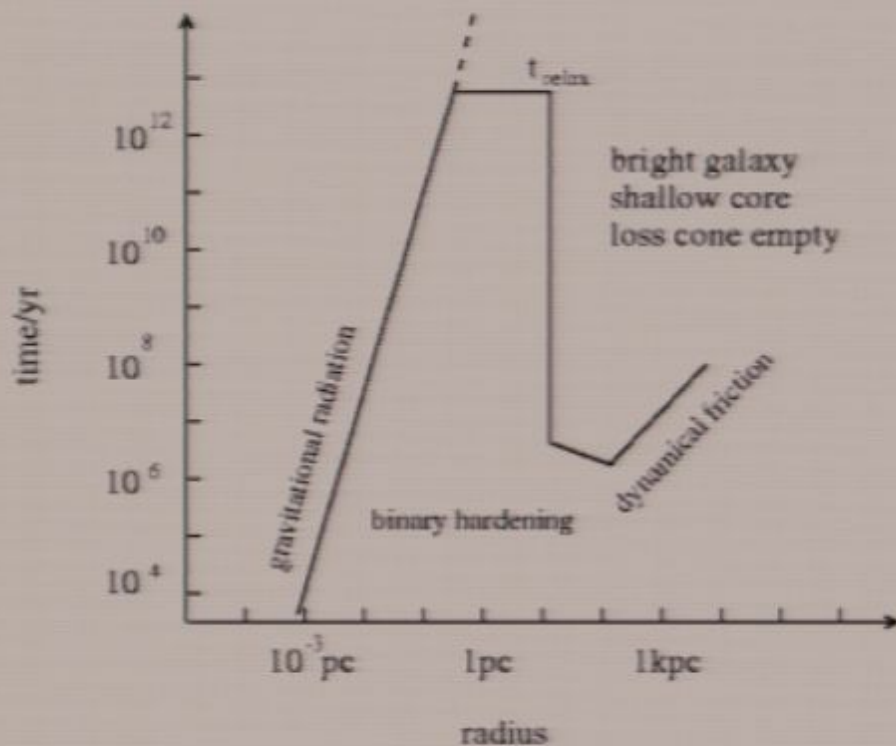
# Do binary black holes coalesce?

Need either:

accretion of gas

or

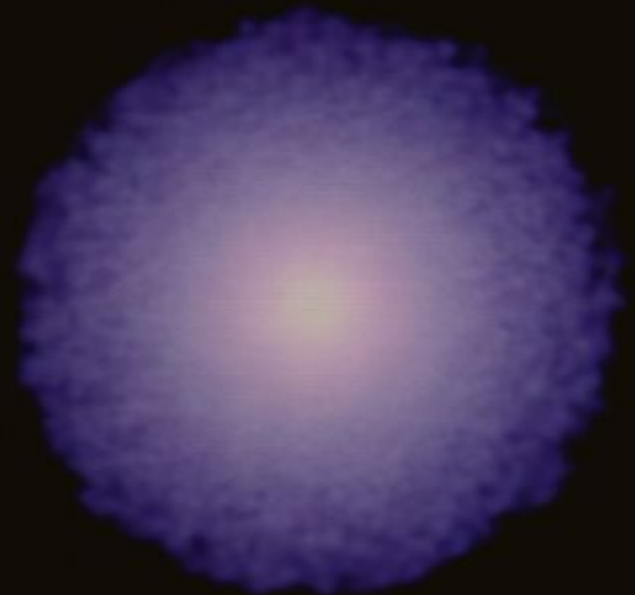
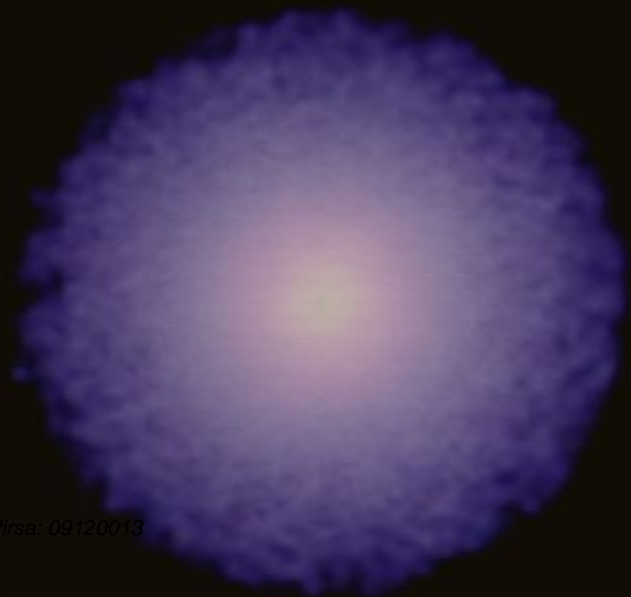
efficient refilling of  
loss cone (only in  
small galaxies and  
maybe at high  
redshift)



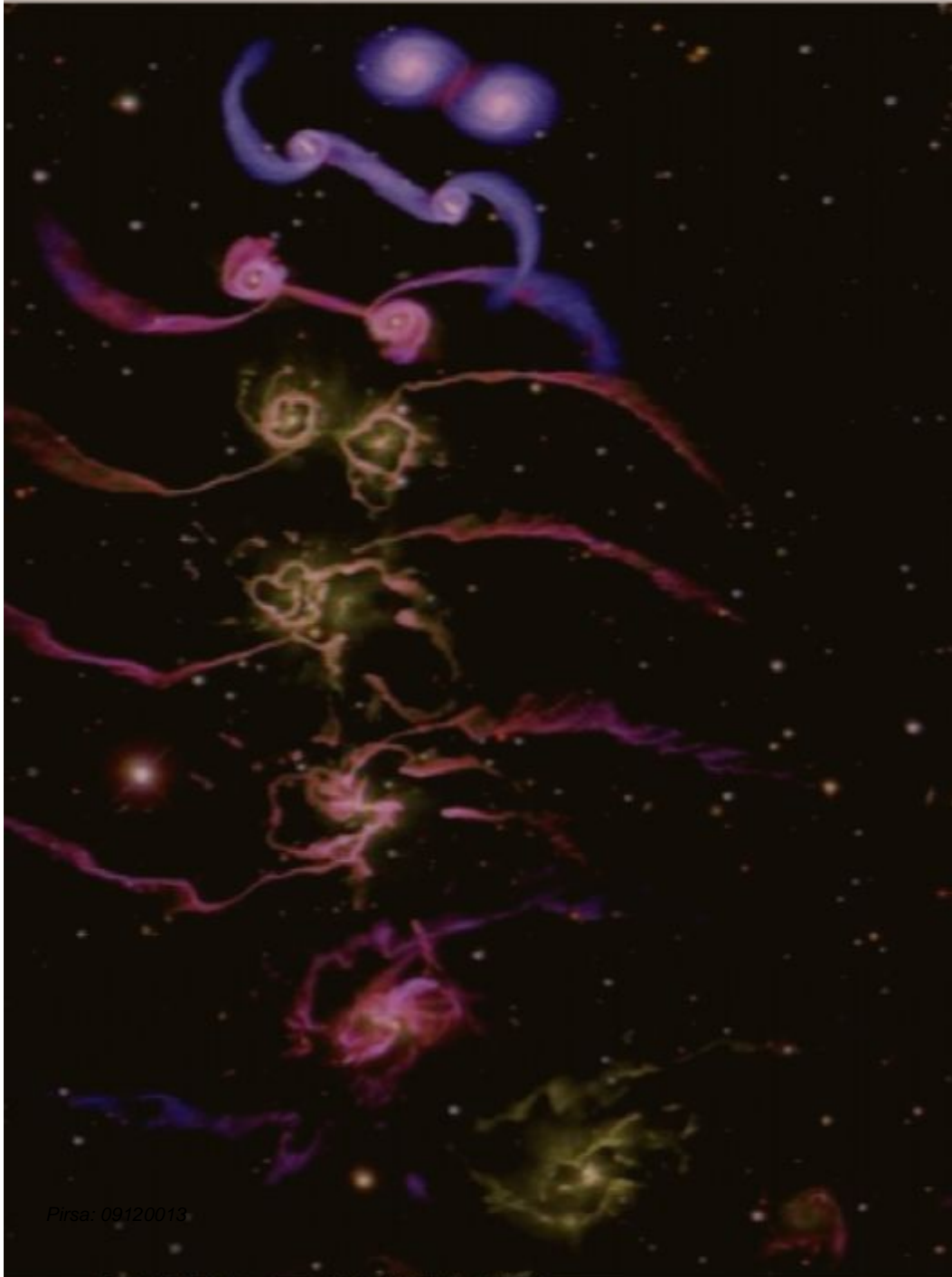
Begelman Blandford & Rees 1980  
Yu 2002



$\Gamma = 10 \text{ Myr}$



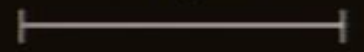
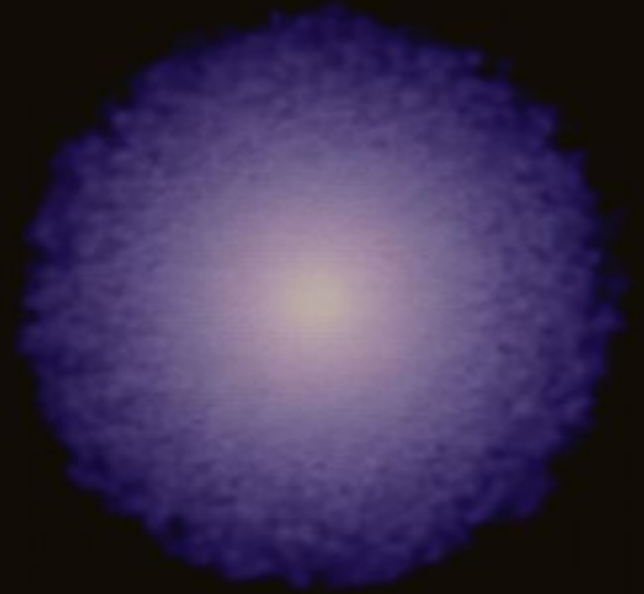
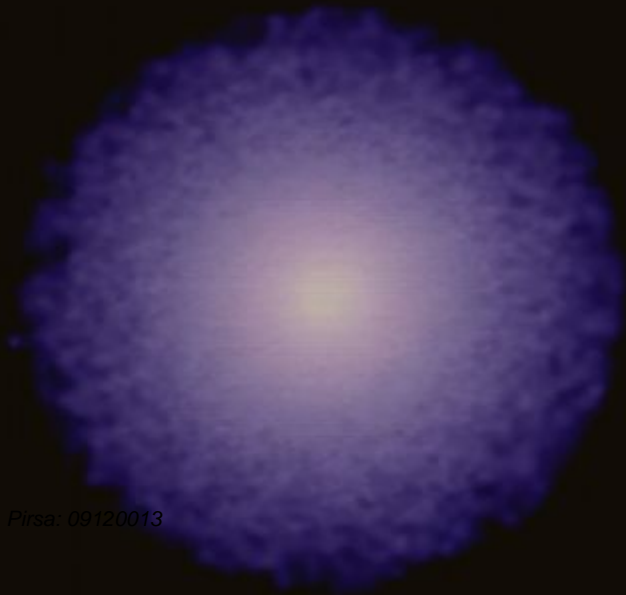




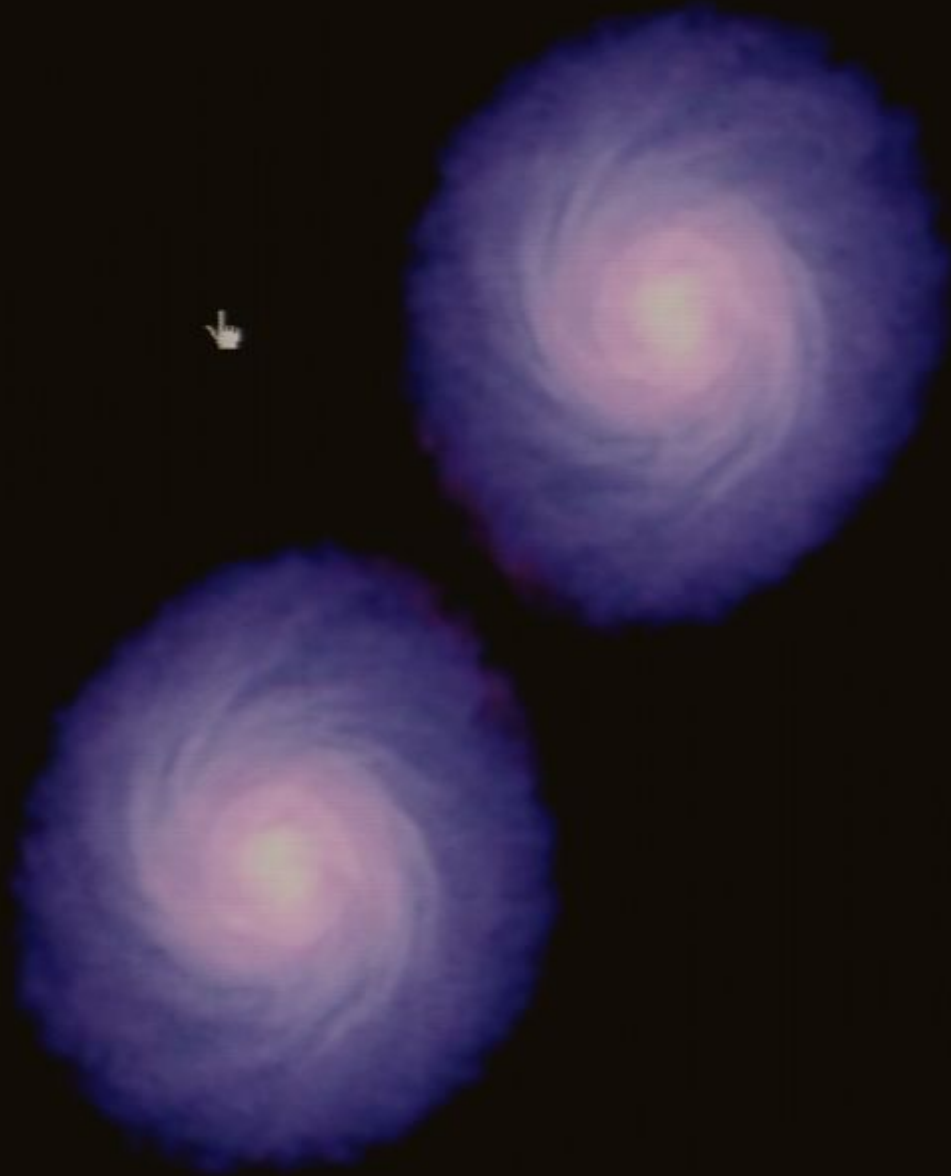
NGC2207



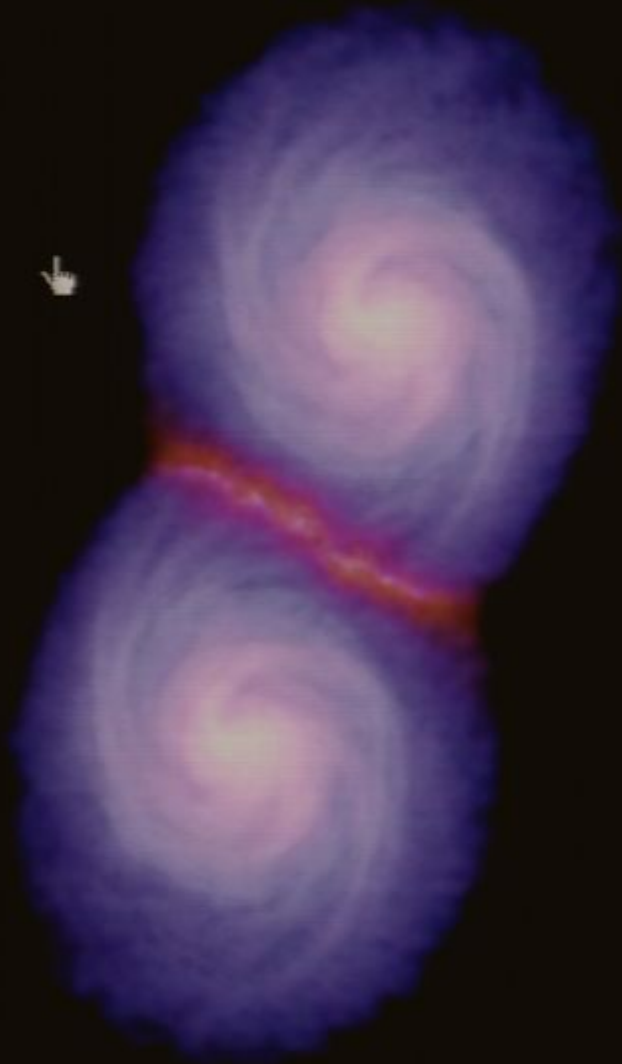
$\Gamma = 0$  Myr



$\Gamma = 230 \text{ Myr}$



$\Gamma = 290 \text{ Myr}$

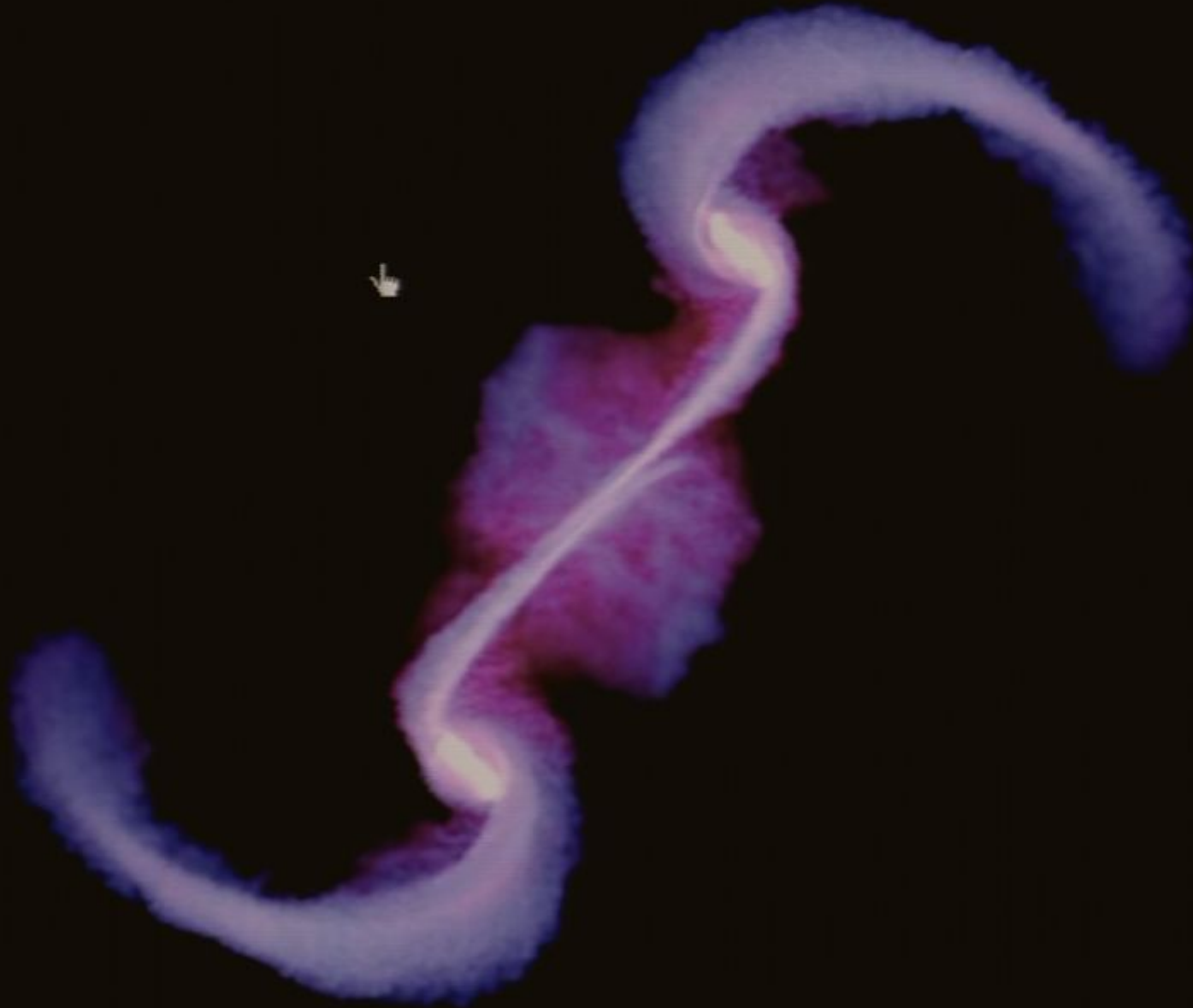




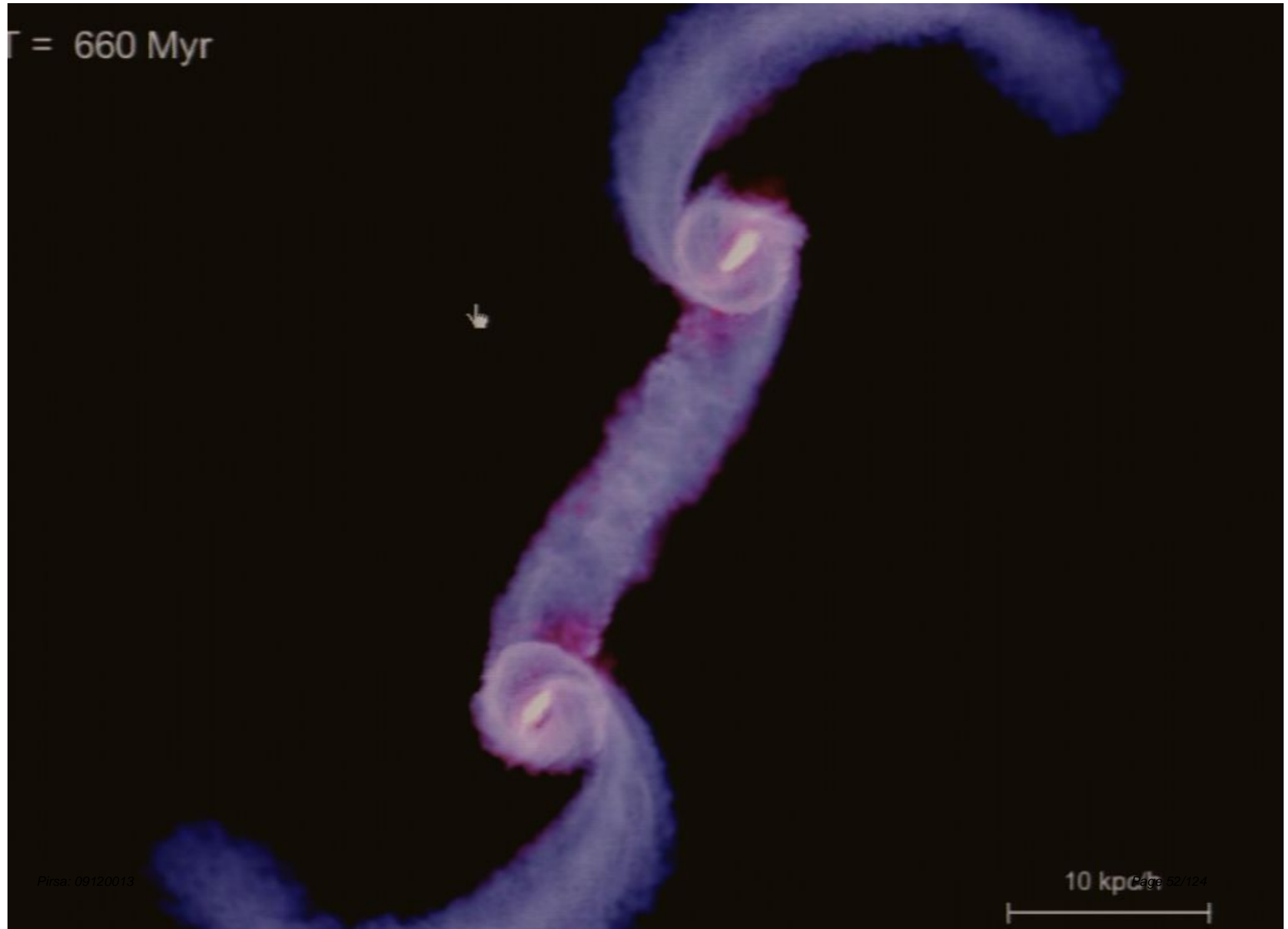
$\Gamma = 360 \text{ Myr}$



$\tau = 550 \text{ Myr}$



$\Gamma = 660 \text{ Myr}$



$\Gamma = 760 \text{ Myr}$

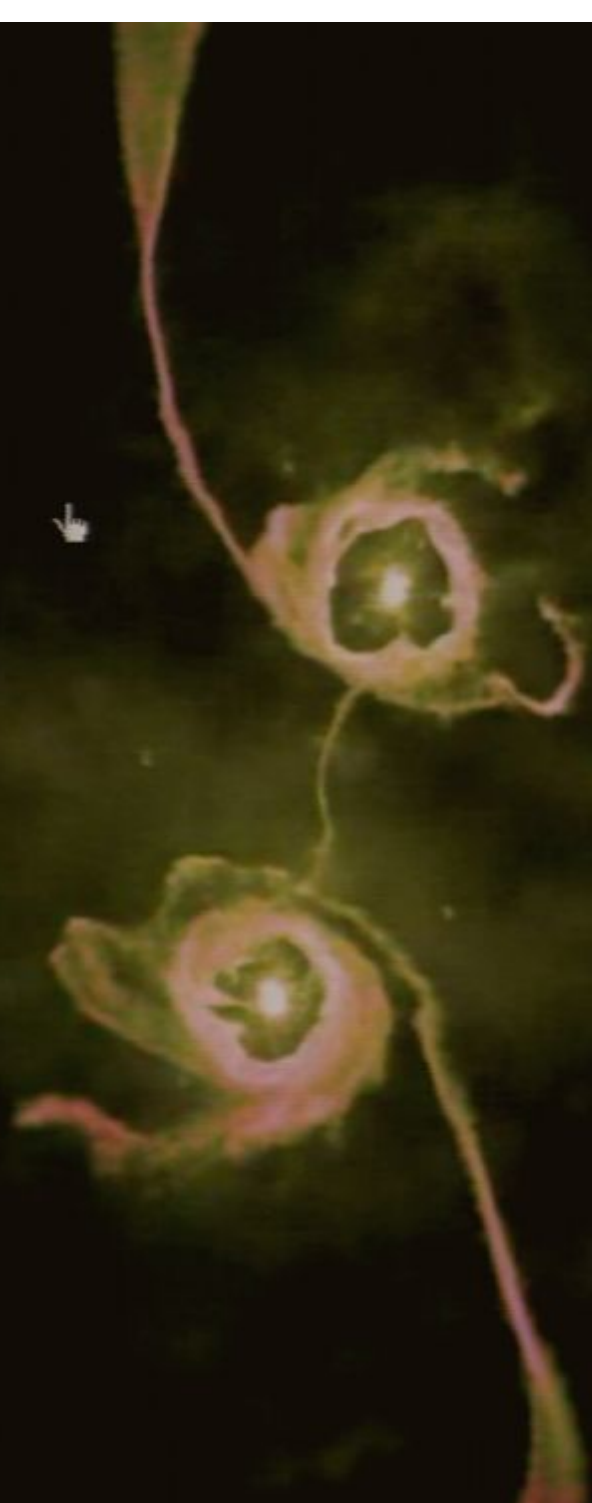




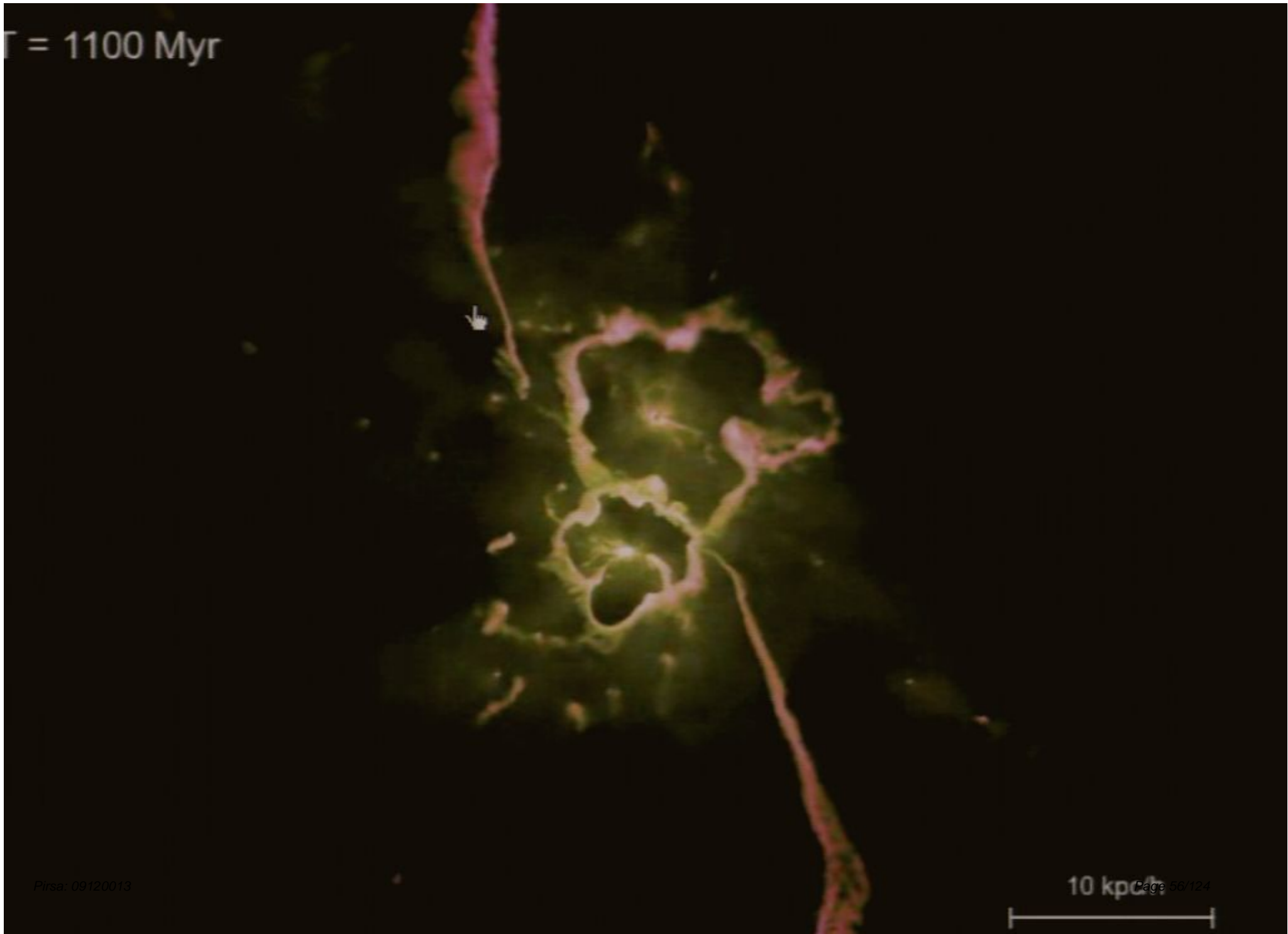
$\Gamma = 900 \text{ Myr}$



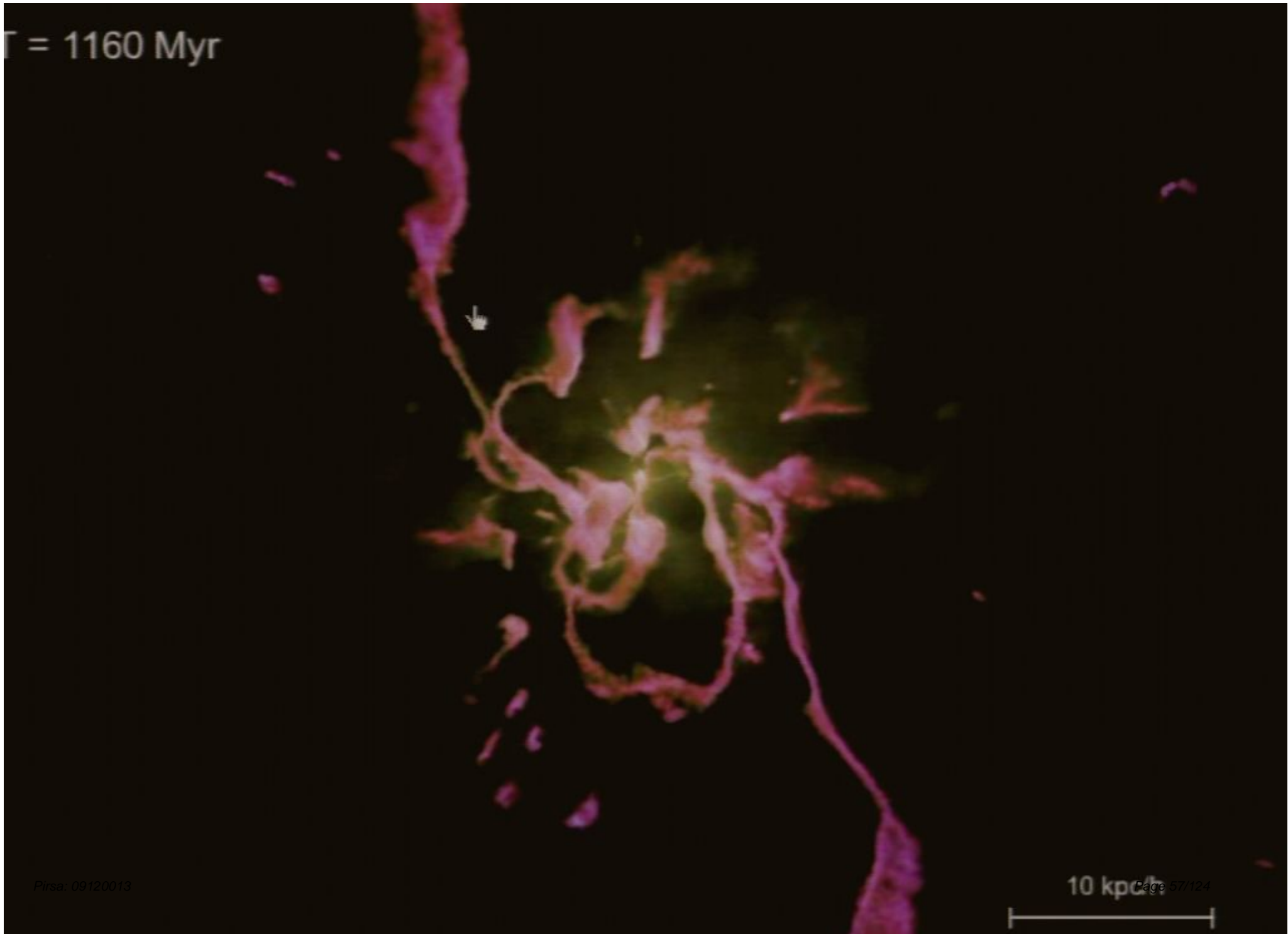
$\Gamma = 1030 \text{ Myr}$



$\Gamma = 1100 \text{ Myr}$

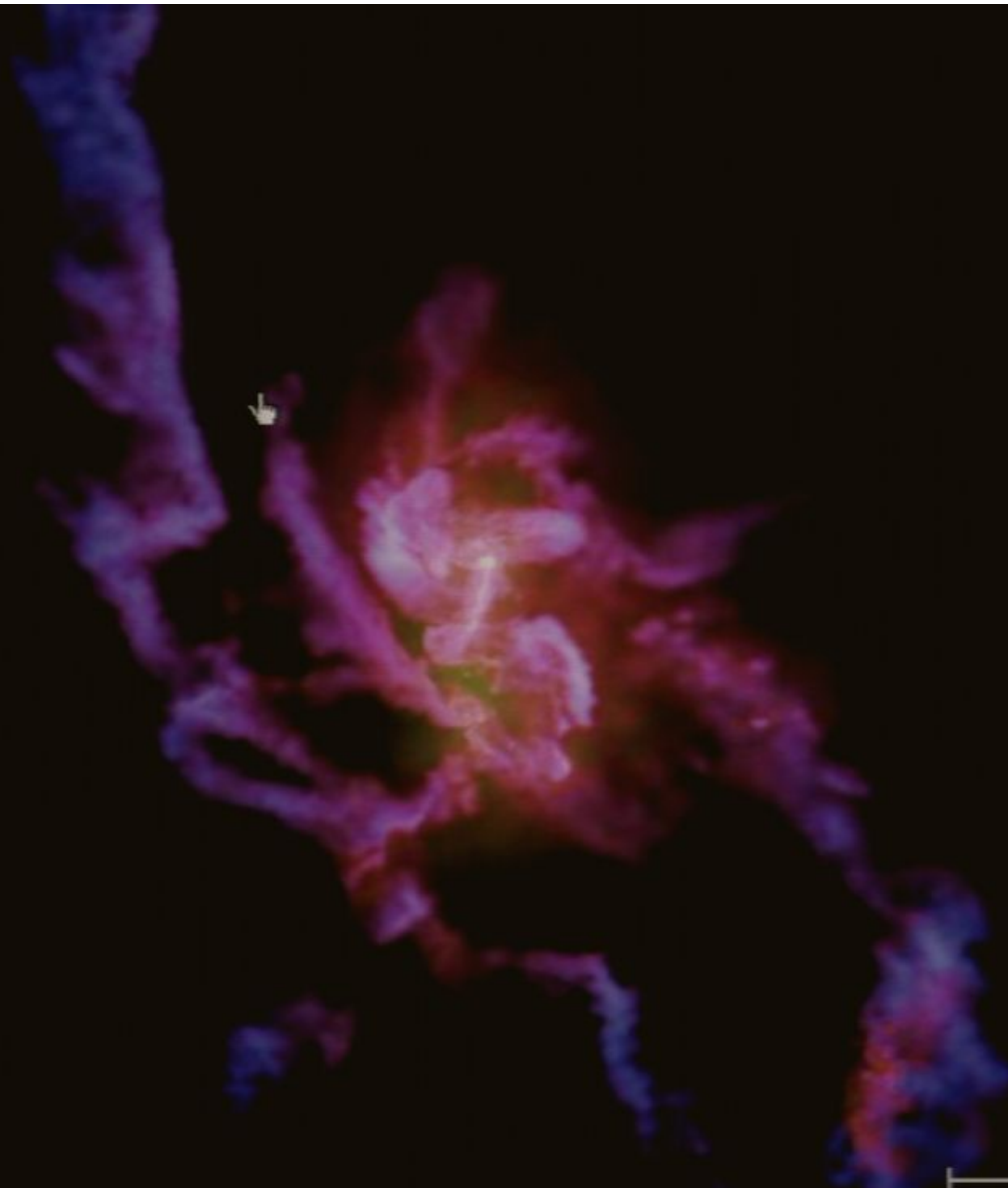


$\Gamma = 1160 \text{ Myr}$

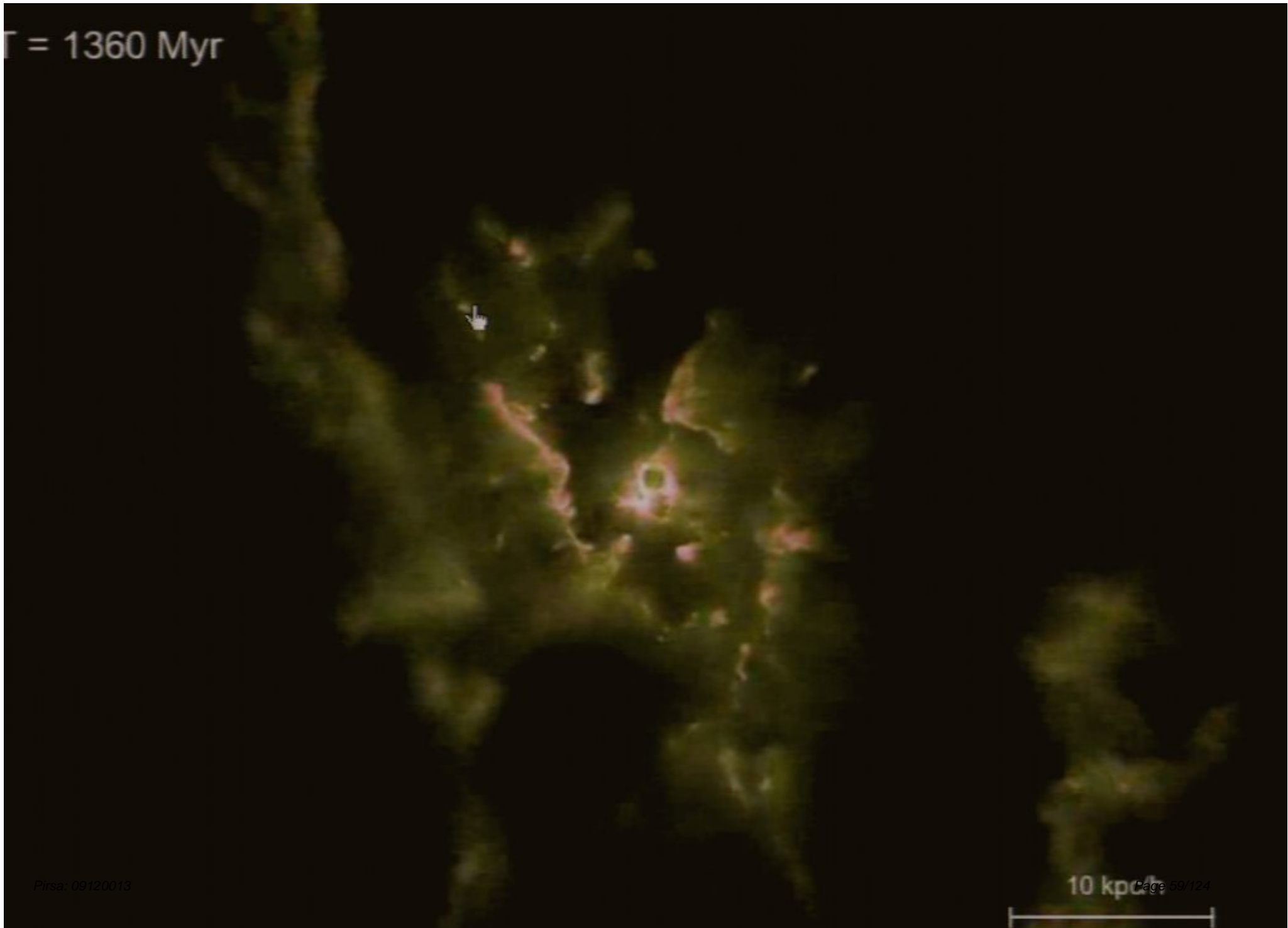




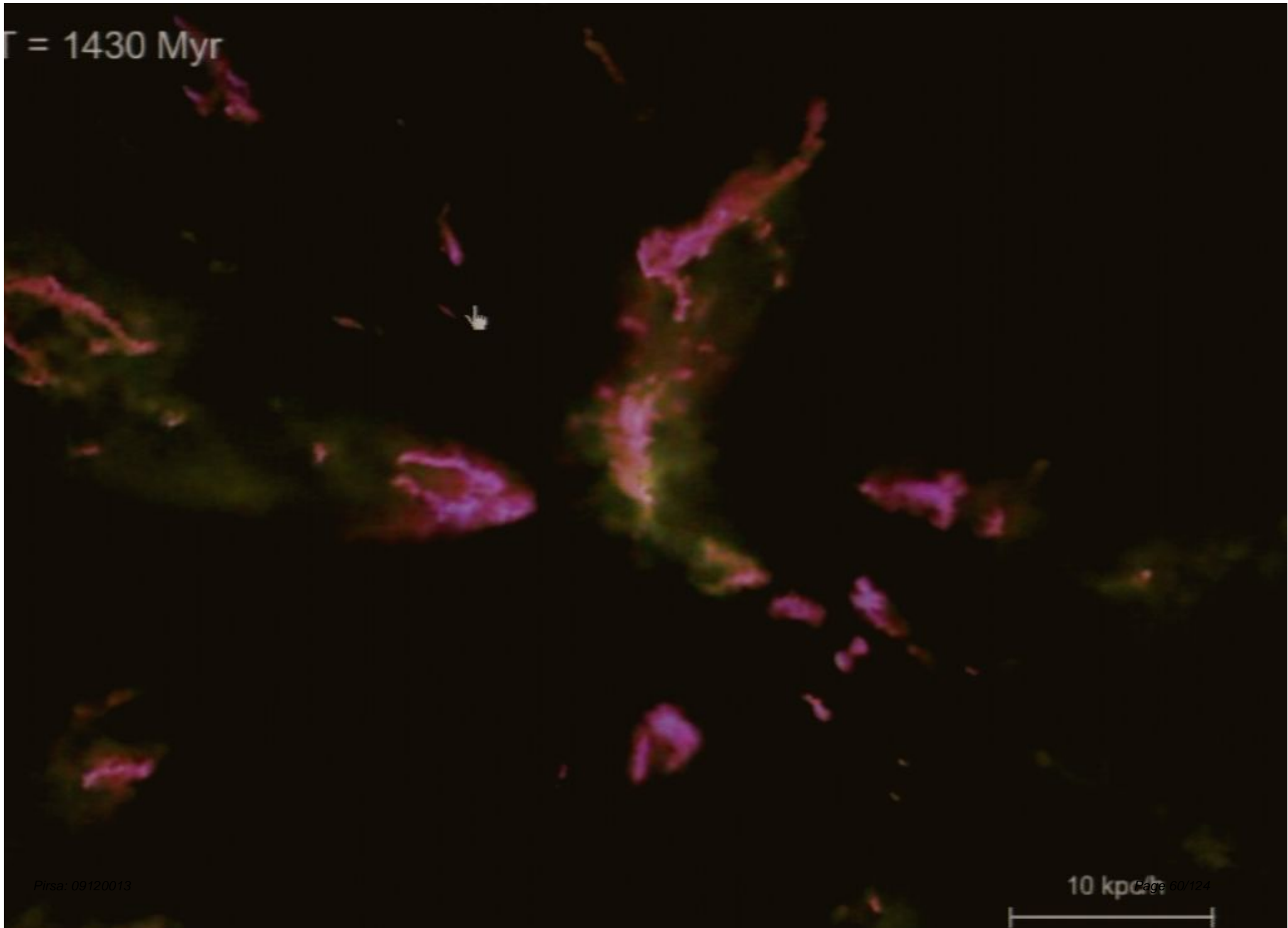
$\tau = 1260$  Myr



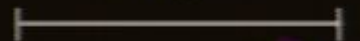
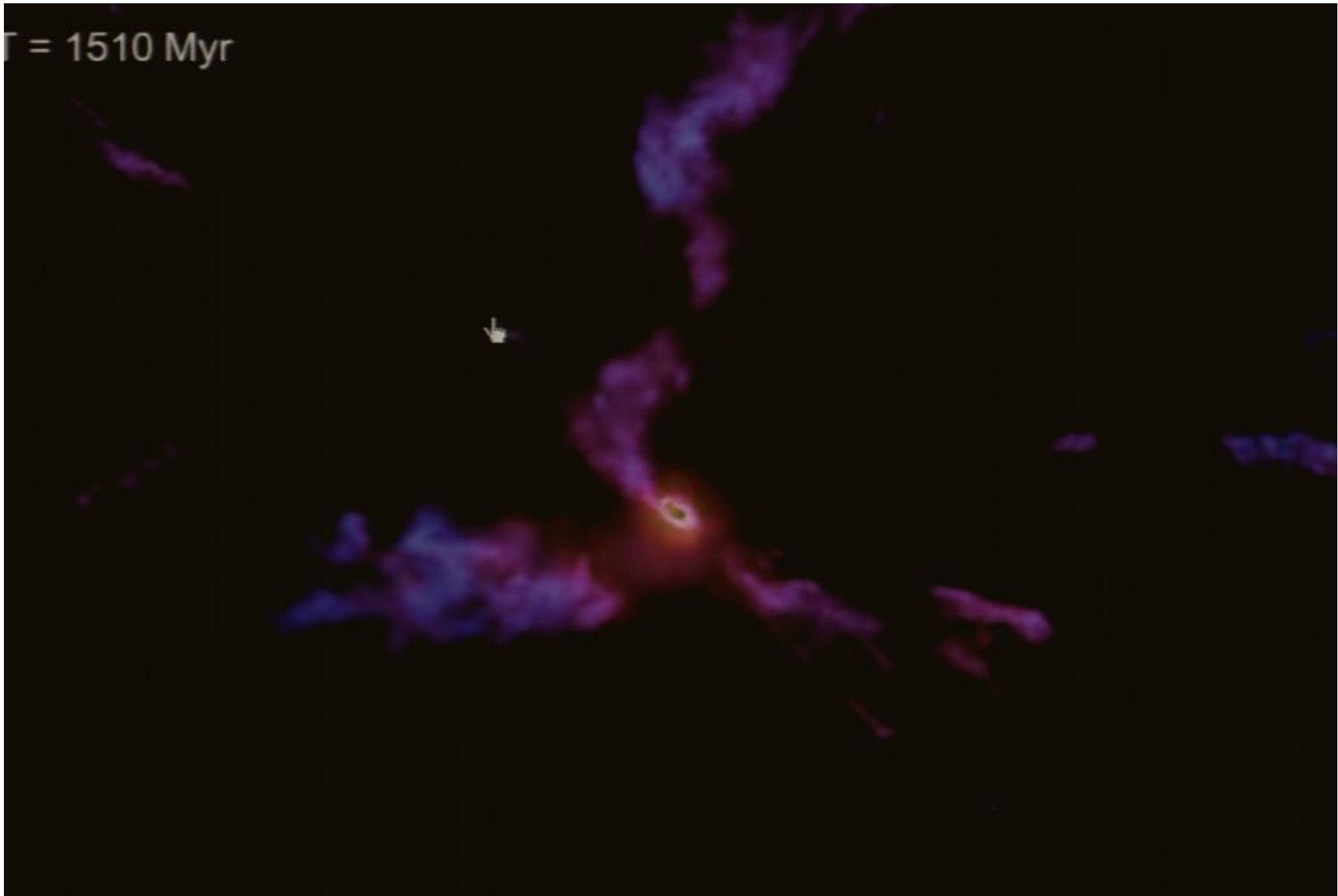
$\tau = 1360$  Myr



$\Gamma = 1430 \text{ Myr}$



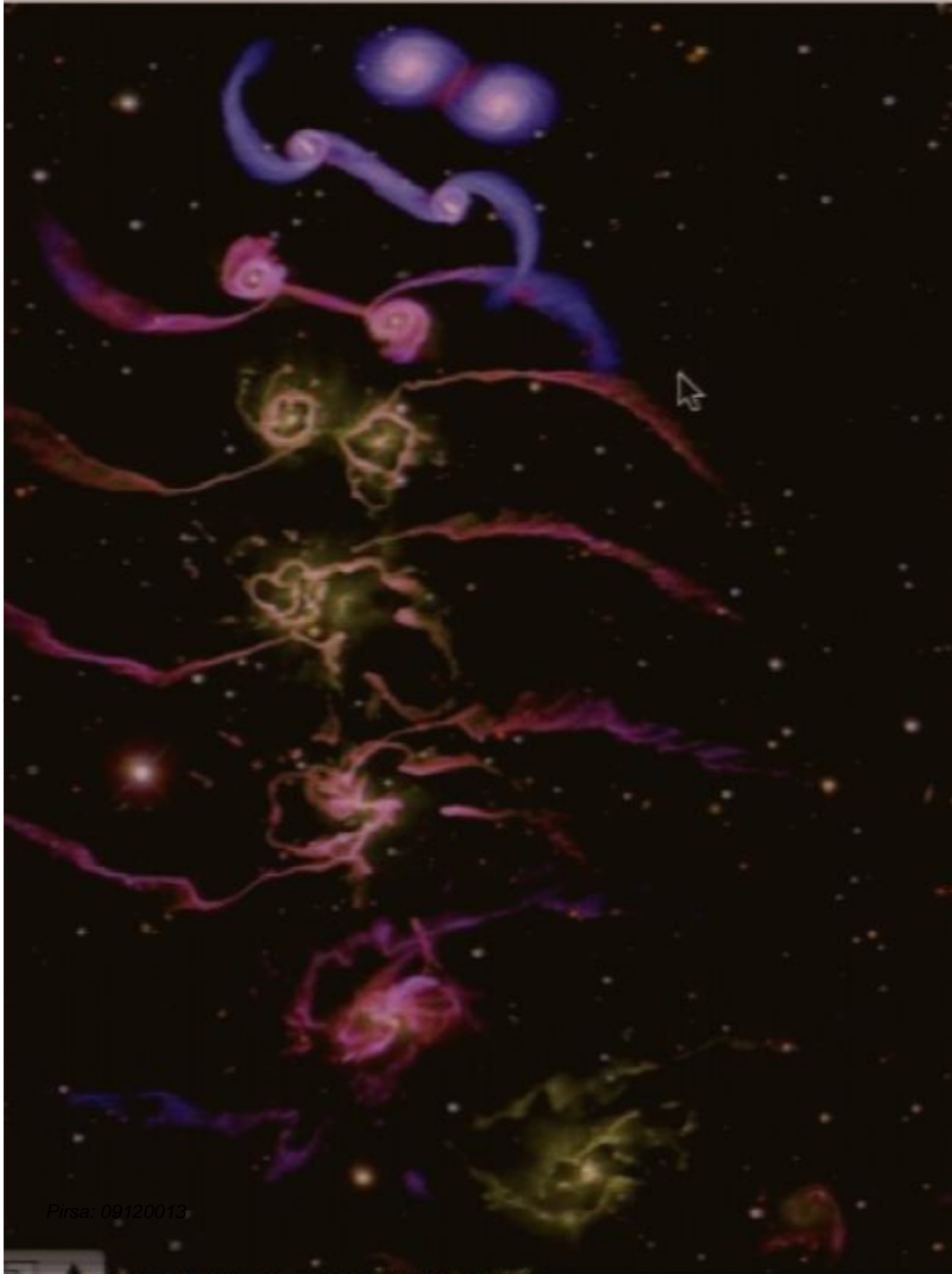
$\Gamma = 1510 \text{ Myr}$



$\Gamma = 1740 \text{ Myr}$





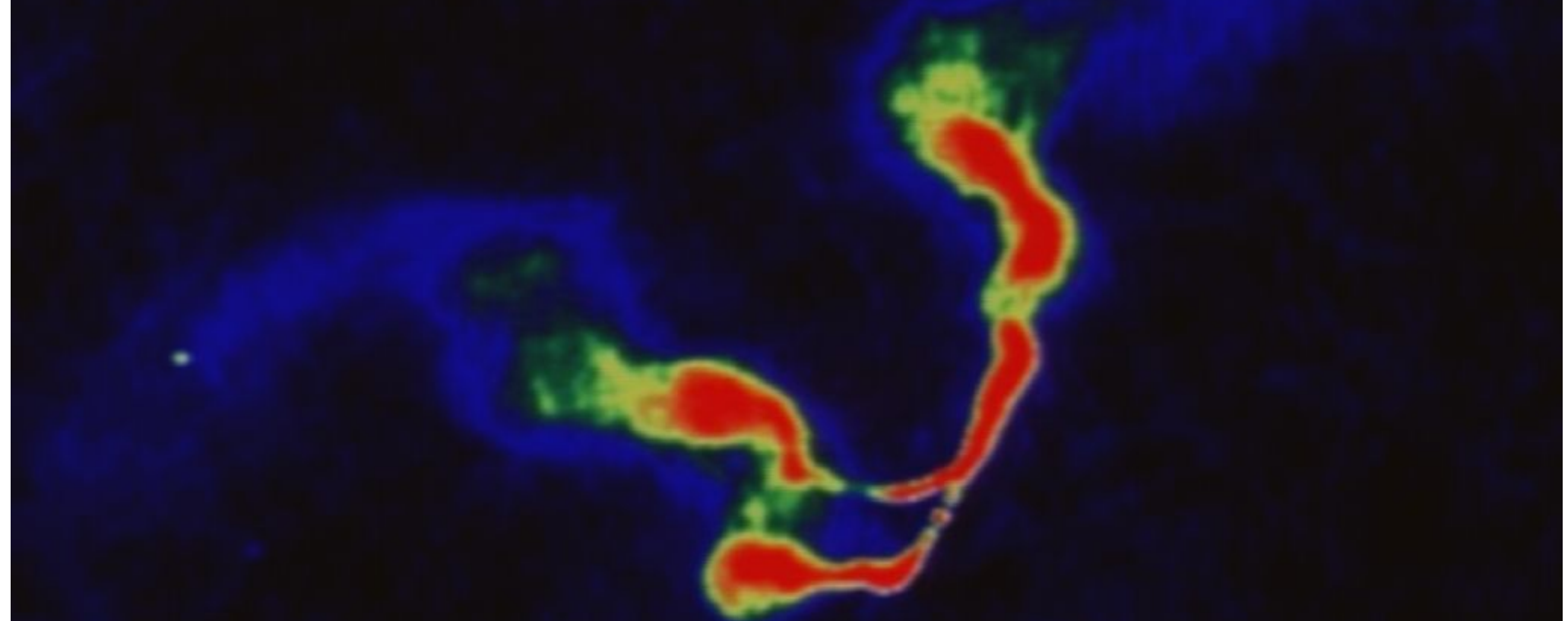


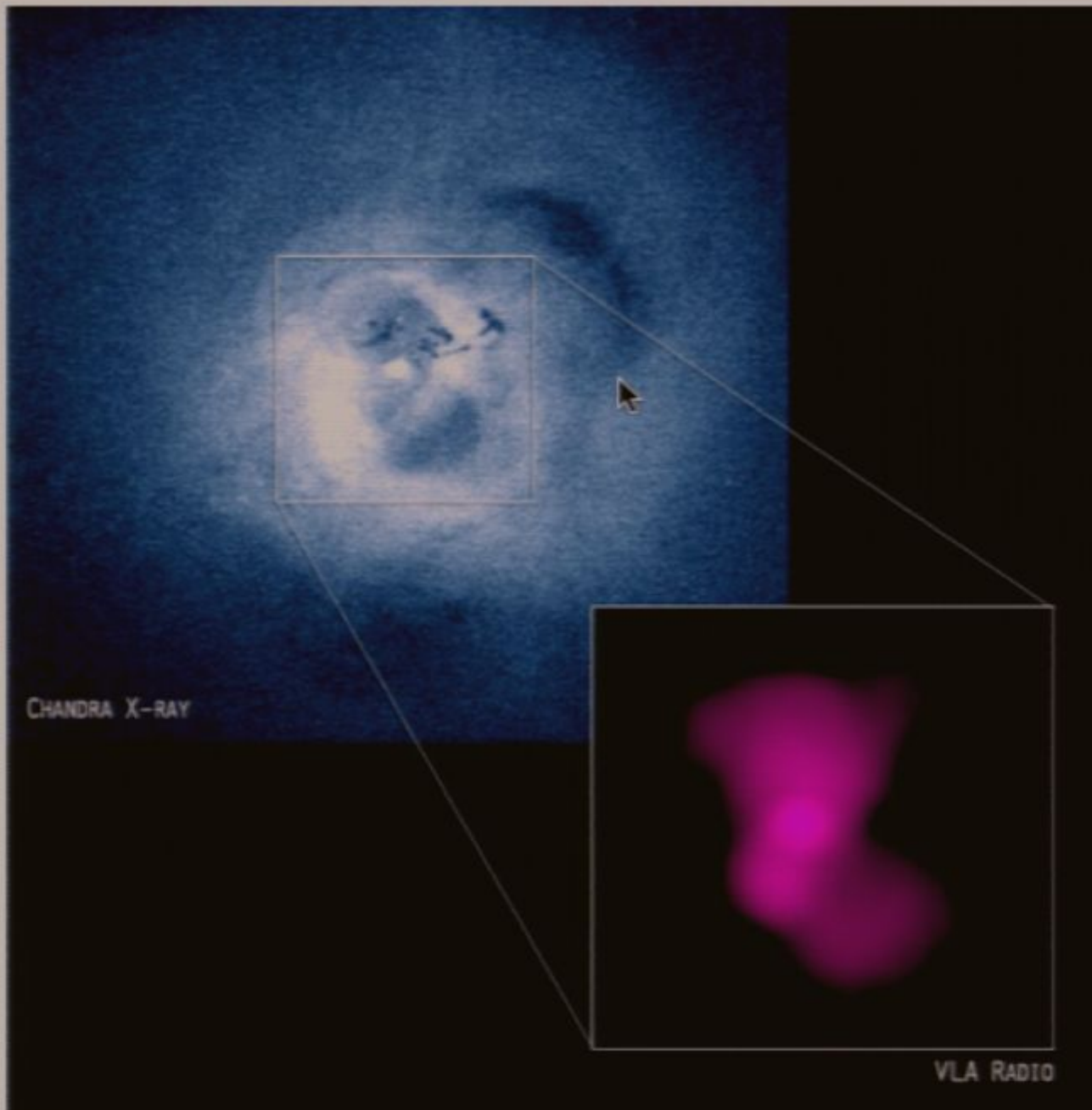
NGC2207



# A binary quasar

43



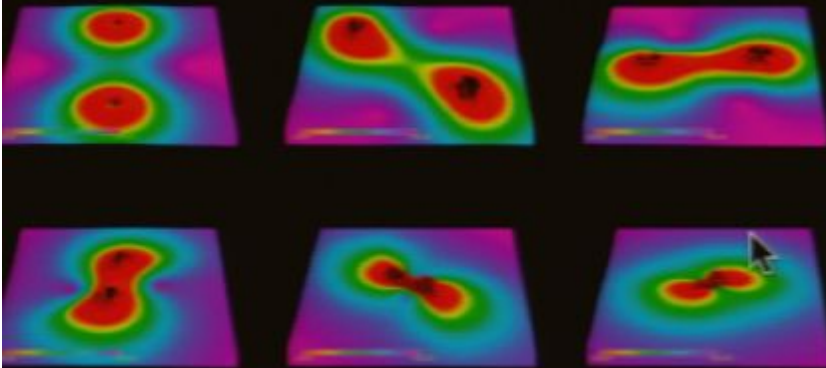


Supermassive black holes heat the gas in galaxy clusters .

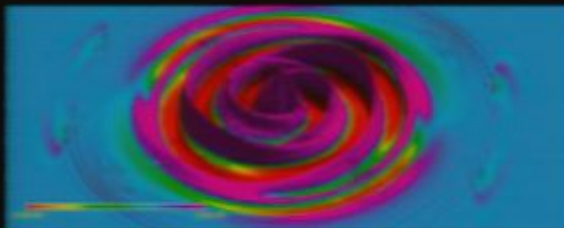
# A Revolution in Numerical Relativity







Snapshots showing a time sequence of the field representing the dynamics of two merging black holes.



A snapshot showing the field representing gravitational waves generated from two inspiralling black holes.

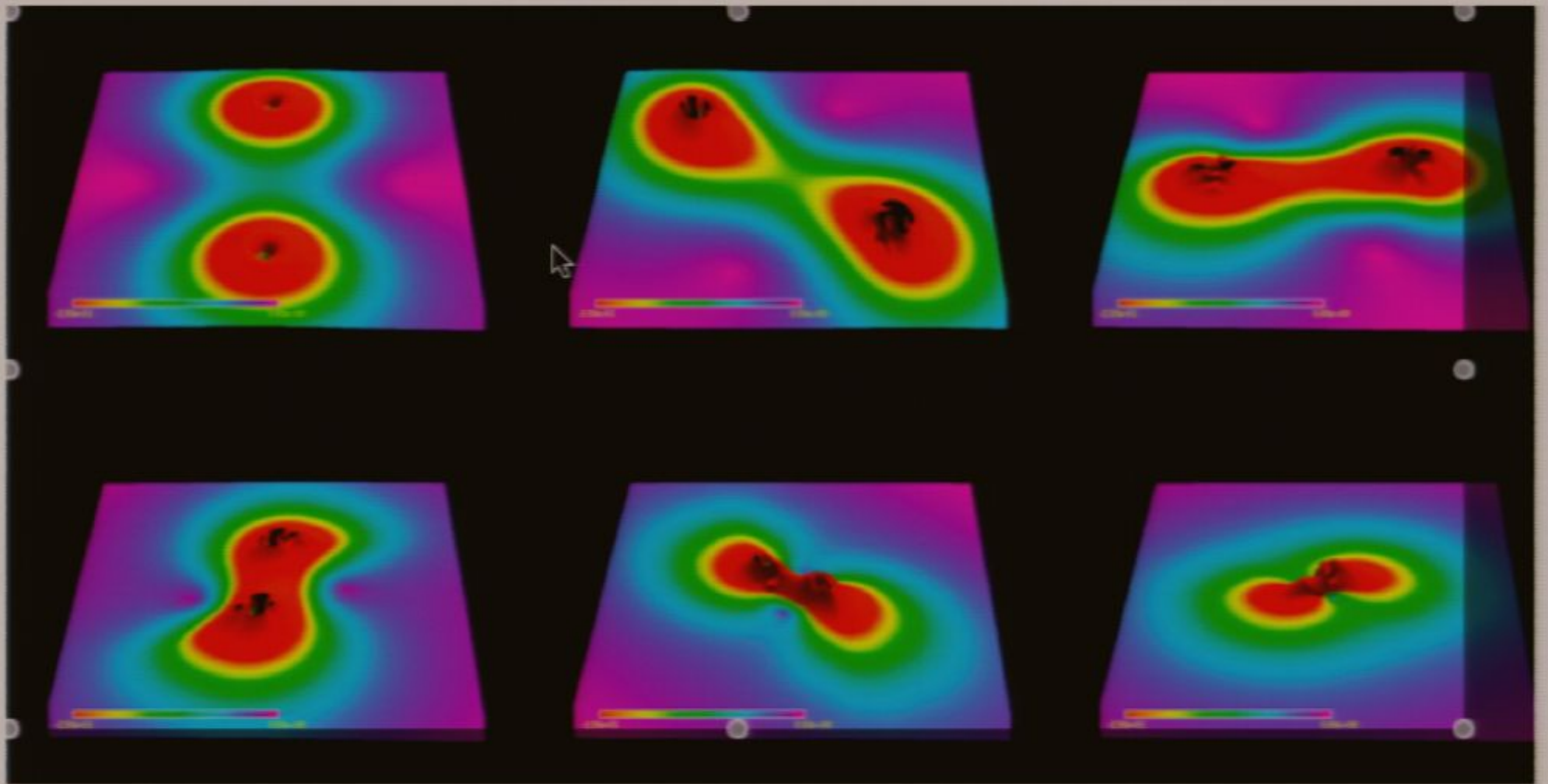
## Modeling Gravitational Wave Sources for LISA

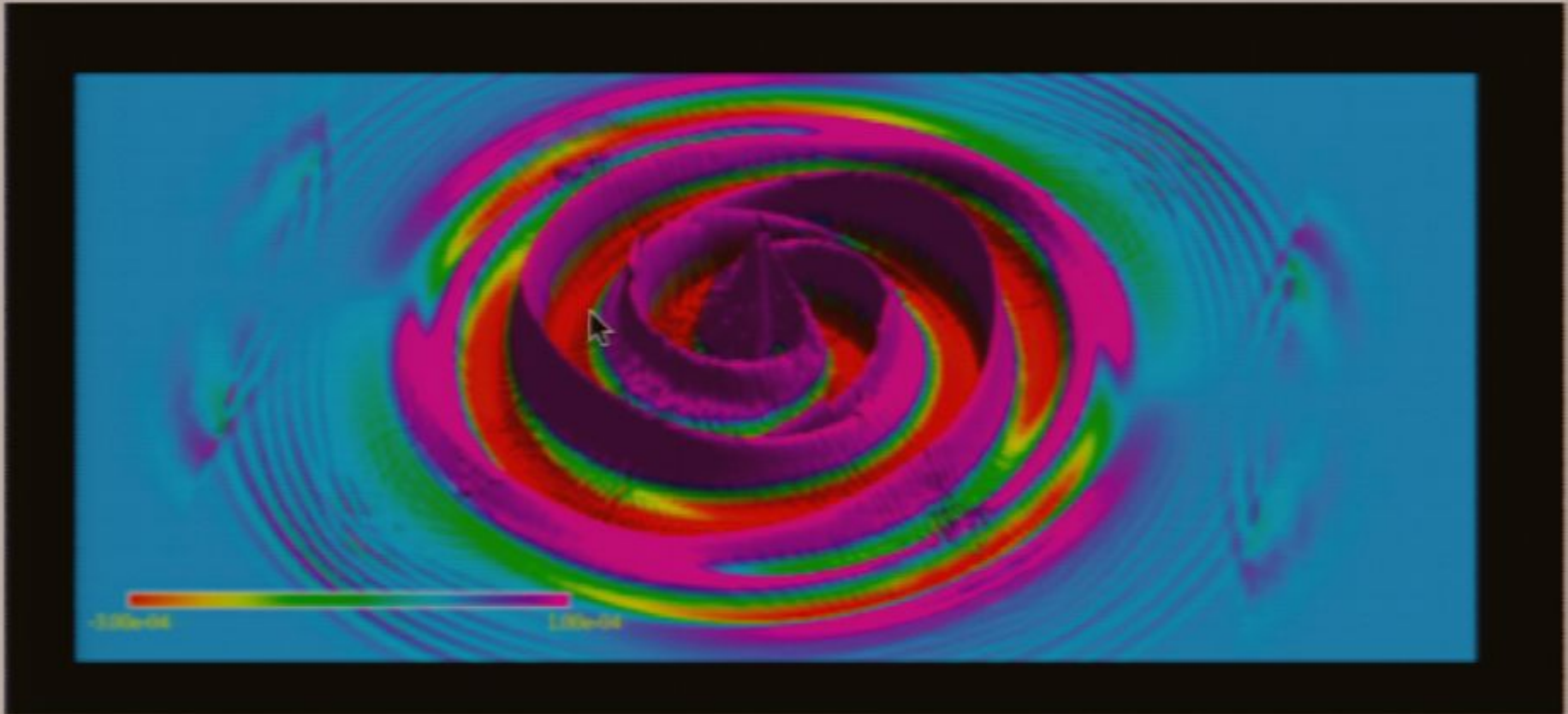
We model the astrophysical coalescence of comparable-mass, massive black hole binaries for different mass ratios and spins and calculate the resulting gravitational wave (GW) signatures. A key feature of our work is the use of mesh refinement techniques to handle the wide range of physical scales involved, from the black holes ( $\sim 1M$ ) to the gravitational waves ( $\sim 10\text{-}100M$ ), and to enable extraction of gravitational waveforms in the wave-zones.

Our methodology involves solving a closely coupled system of partial differential equations with many variables on a very large computational domain. Moreover, we use highly structured component-grids that are distributed across many processors, with a significant amount of communication between processors. Our simulations typically require hundreds of gigabytes and run on several hundred processors for hundreds of hours.

Astrophysical motivation guides our simulations. The waveforms determined by these simulations will be applied to analyzing and interpreting observed GW data from the Laser Interferometer Space Antenna (LISA) mission. Linear momentum loss due to asymmetrical radiation of GW in the unequal mass mergers imparts "kicks" to the merger remnant. High kick velocities from such mergers have the potential to strongly impact our understanding of how massive black holes have developed over cosmological time-scales.



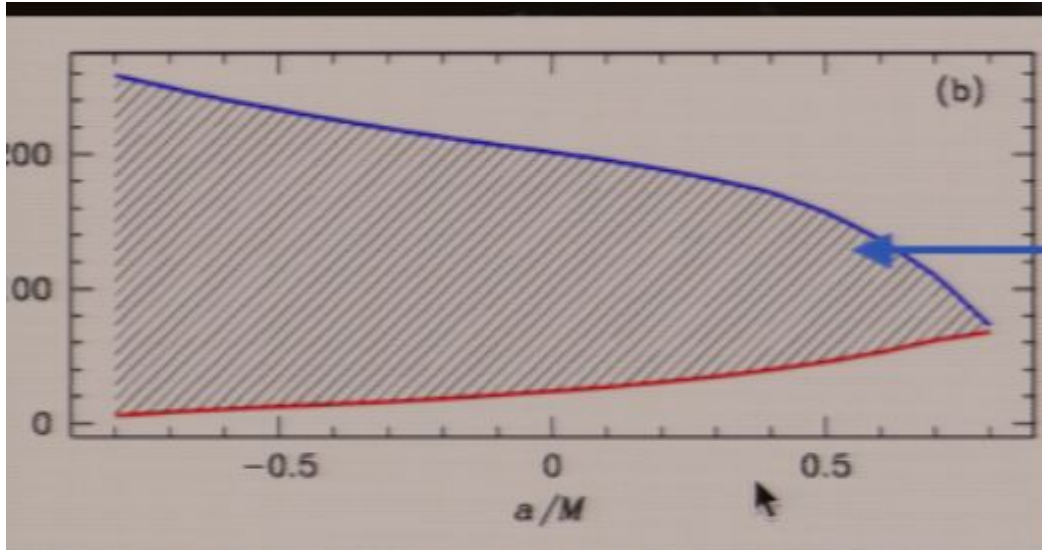




The merged black hole receives a recoil “kick” due to asymmetric gravitational wave emission.



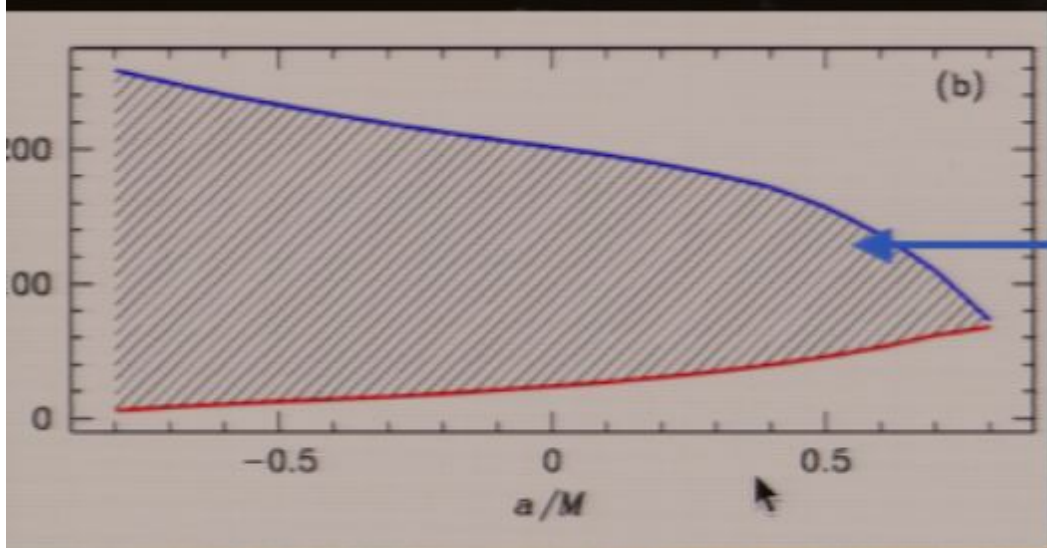




estimated kick velocity  
after binary merger

can be  $> 2000 \text{ km s}^{-1}$   
for anti-aligned spins!

Favata et al. 2004

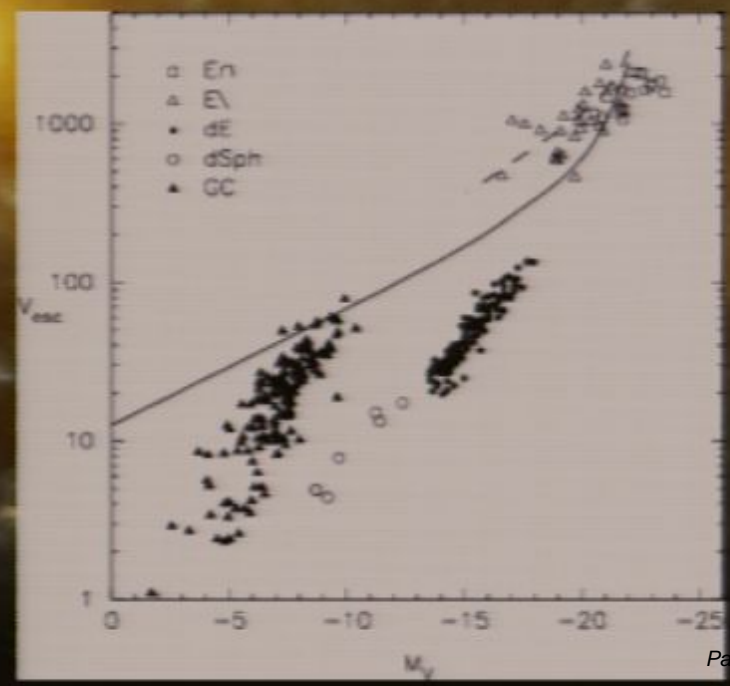


estimated kick velocity  
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can be  $> 2000 \text{ km s}^{-1}$   
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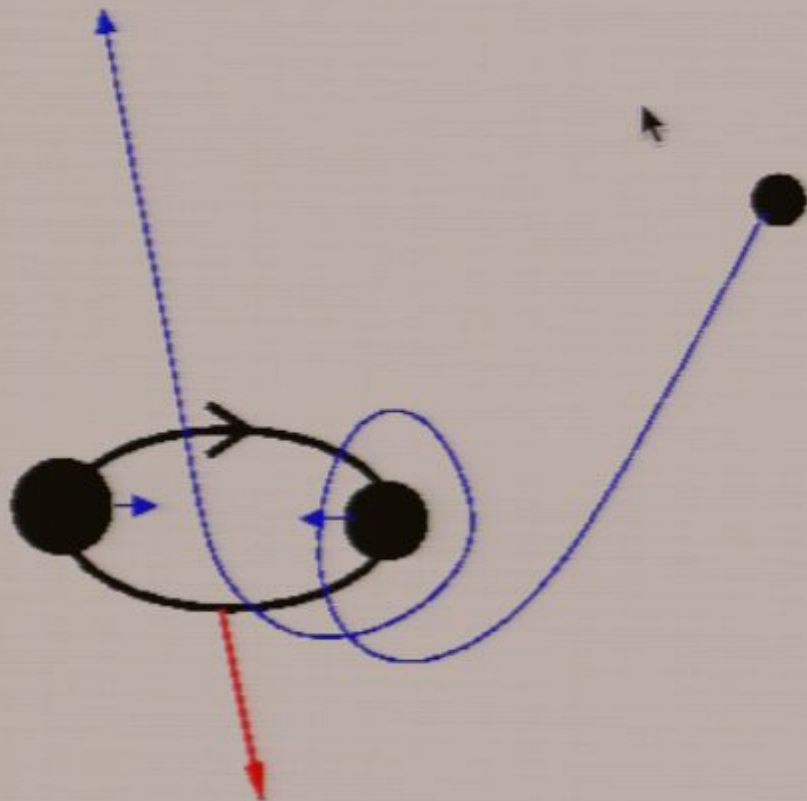
Favata et al. 2004

escape velocity from  
galaxies



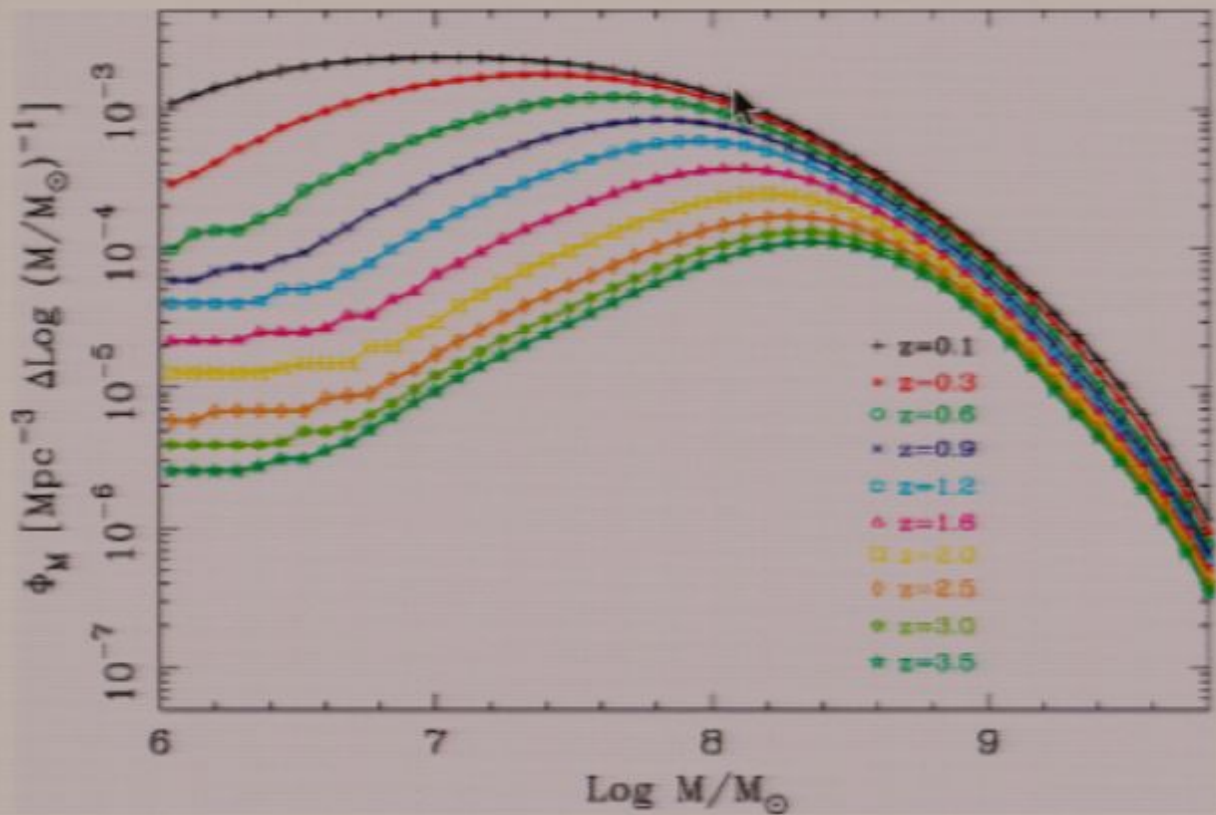


# Triple Interactions



If the binary does not coalesce before a third black hole falls in **sling shot ejection** of the lightest or in some cases all three black holes will occur

Saslaw, Valtonen & Aarseth 1974



The co-evolution of galaxies and their central supermassive galaxies is now reasonably well Understood.

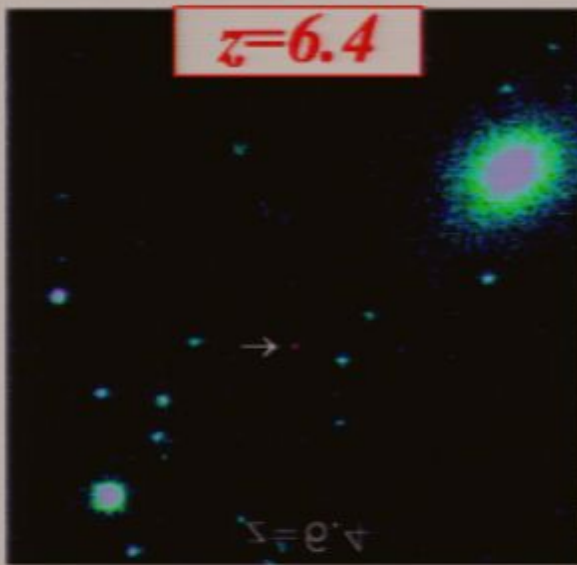
Merloni 2004



But how do supermassive black holes form  
in the first place?



# Is this the Rosetta stone?



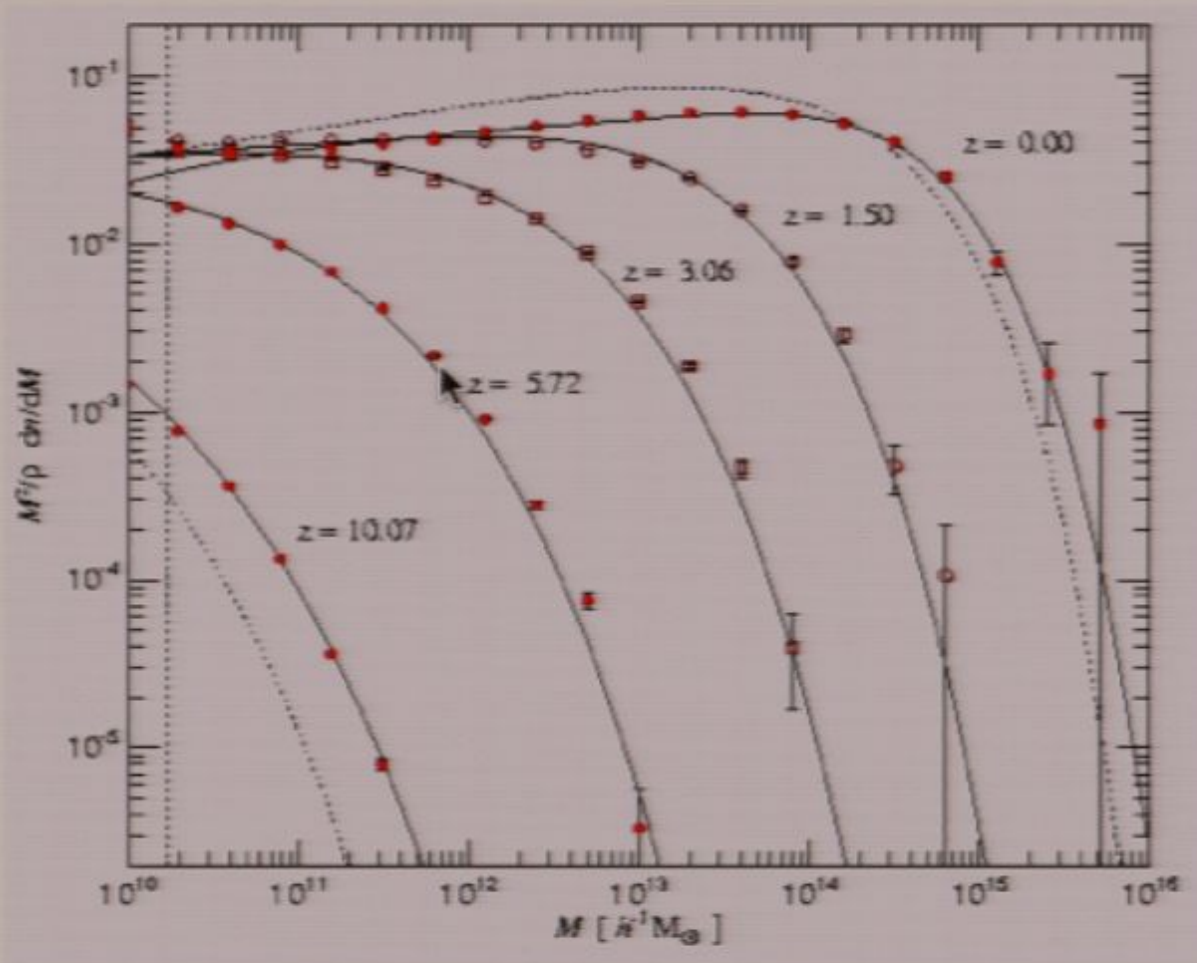
Black holes as massive as the most massive black holes today have already formed at  $z=6.4$ .

Estimated mass:  $3 \times 10^9 M_{\text{sol}}$

Fan et al.







Springel et al 2005

Sufficiently massive haloes do exist at  $z=6$ .





# Eddington-limited accretion



Radiation pressure due to Thomson scattering on electrons equals gravitational attraction if

$$\frac{GMm_p}{r^2} = \frac{L_{\text{Edd}} \sigma_T}{c 4\pi r^2}$$

The luminosity due to accretion can be written as

$$L_{\text{acc}} = \epsilon \dot{M} c^2$$

The Eddington limit for (spherical) accretion is

$$\dot{M}_{\text{Edd}} = \epsilon^{-1} \frac{Gm_p}{c\sigma_T} M$$

The characteristic timescale for Eddington limited exponential growth is

$$t_{\text{Salp}} = \epsilon \frac{c\sigma_T}{4\pi Gm_p} \approx 4.5 \epsilon_{0.1} 10^7 \text{ yr}$$





Age of Universe at  $z=6.4$ : 0.8-0.9 Gyr



For Eddington limited accretion only  $20 \epsilon_{0.1}^{-1}$   
e-foldings possible since Big Bang!





## Growth from stellar mass seeds requires

Eddington-limited accretion with duty cycle close to one  
and  
efficient growth in shallow potential wells  
and  
("fine tuning" of space density of stellar mass black hole  
seeds to avoid excessive ejection by black hole recoils  
in hierarchically merging proto-galaxies)

or

super-Eddington accretion



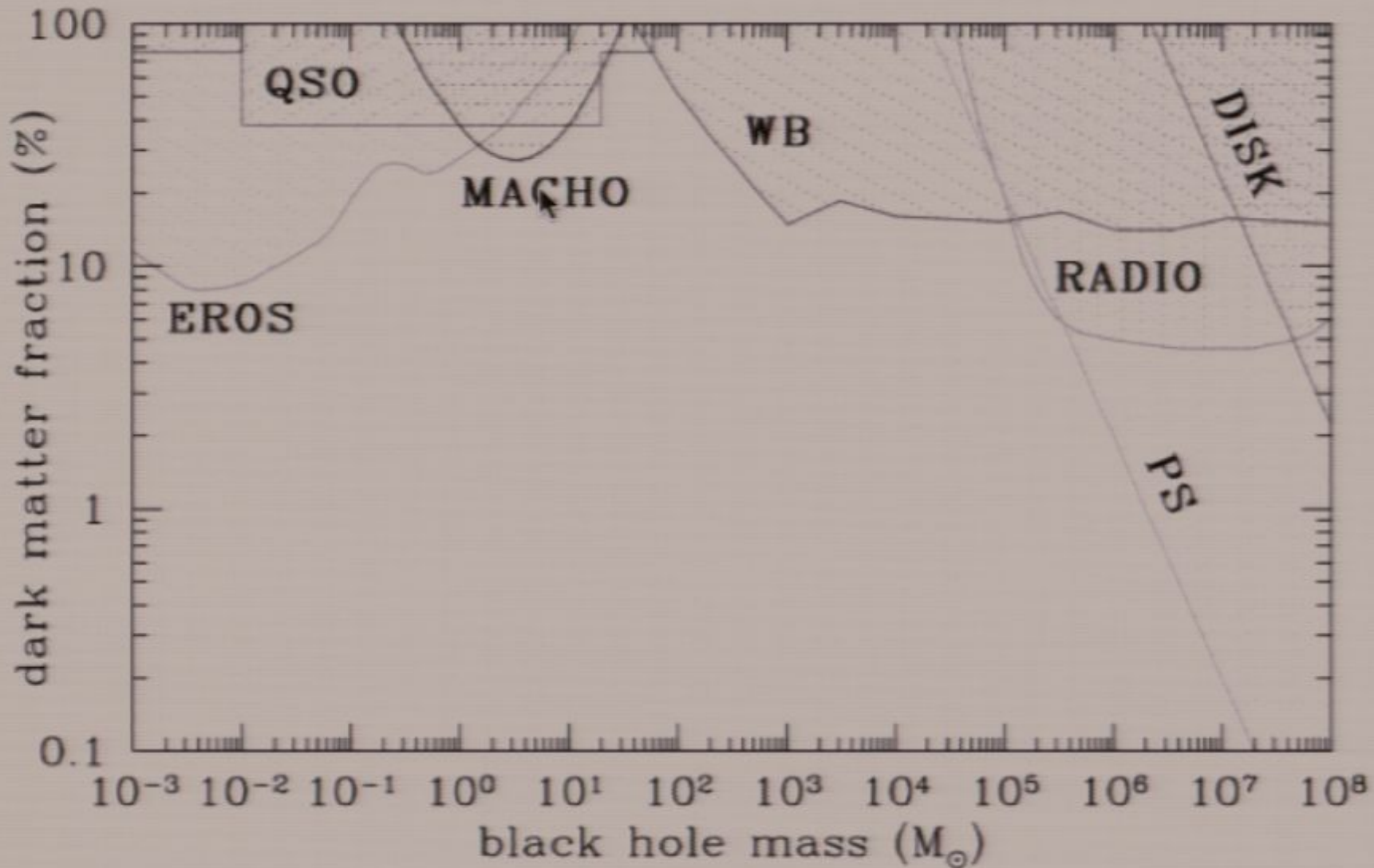


# A): Primordial Black Holes

- Proposed by many authors as artifacts of phase changes at, for example, quark-hadron transition. No consensus.
- Fractional mass is small fraction of DM.
- At origin, some fraction of horizon mass.
- “Compensated” at birth but soon uncompensated.
- Distribution Poisson like.



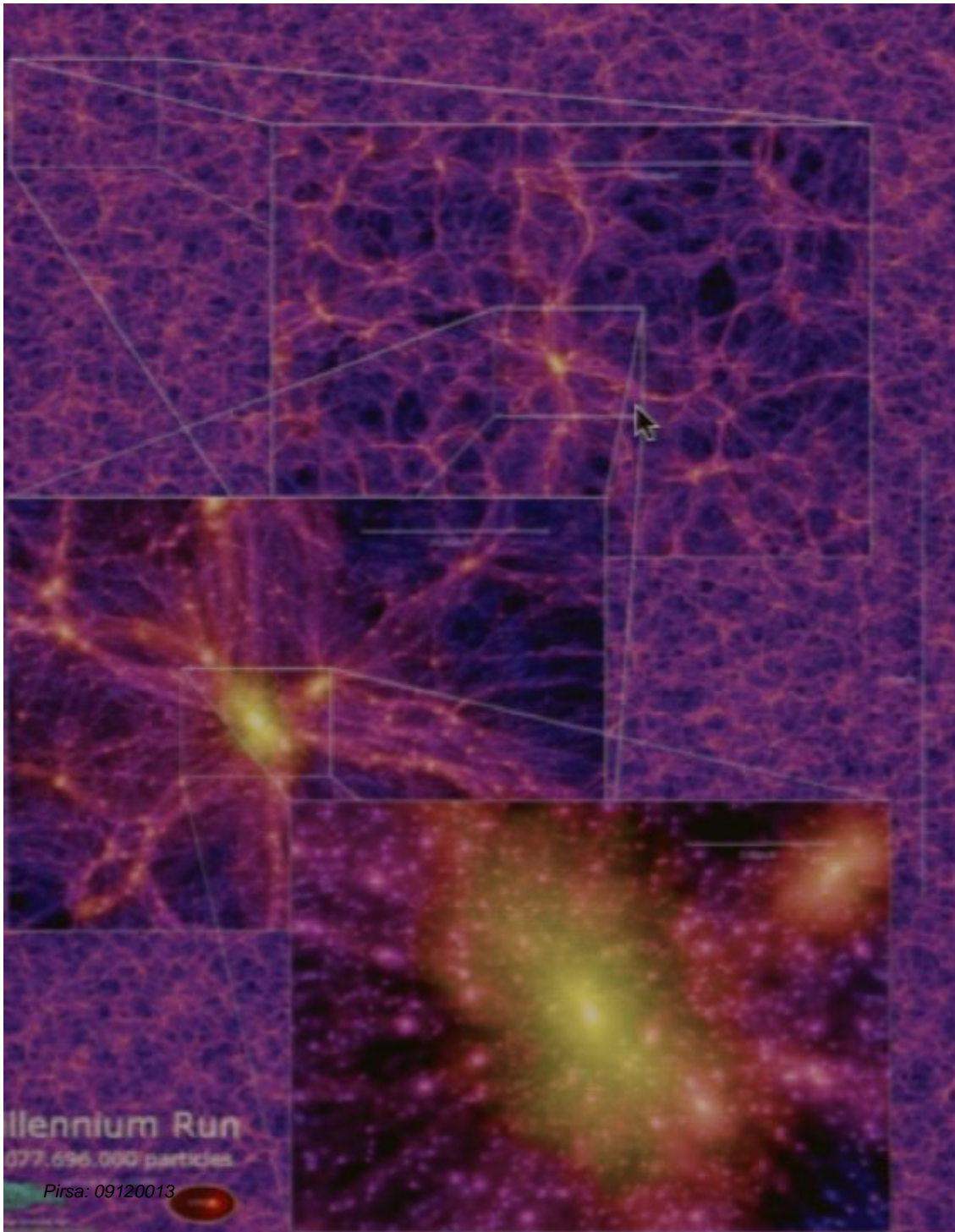
# Limits on primordial black holes





# The early growth of supermassive black holes





Volker Springel    Deborah Sijacki

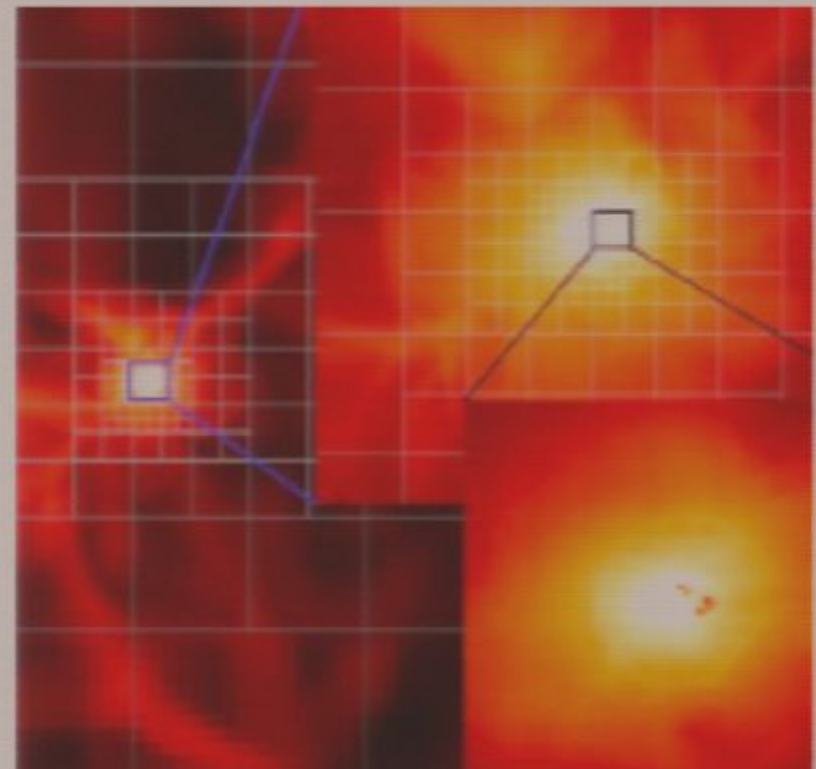
Resimulating the build-up of galaxies and black holes (including spin history and kicks due to gravitational wave re-coil) in the most massive halo at  $z=6$  with higher resolution.



# Cosmos



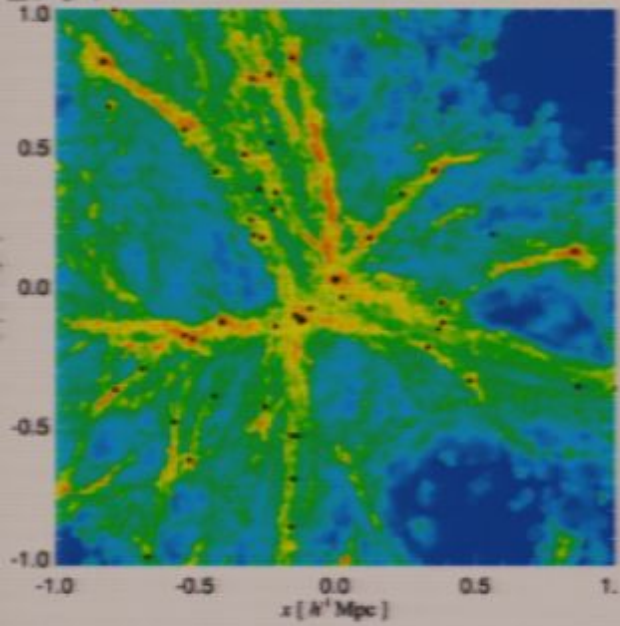
# ENZO AMR



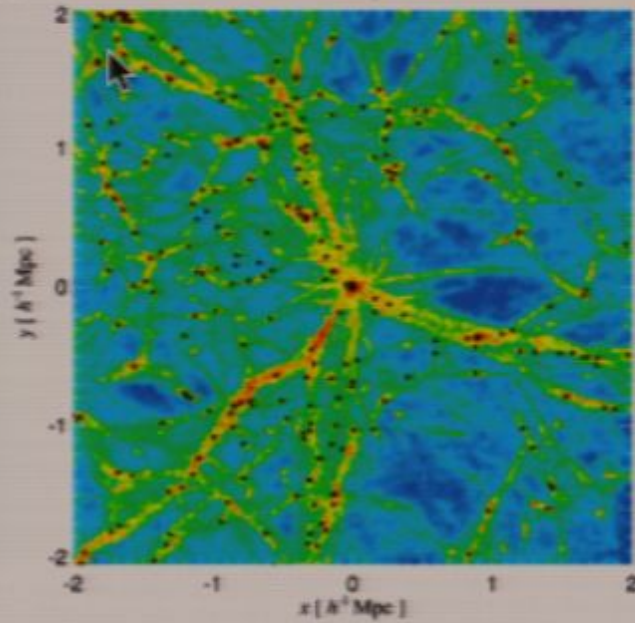
# Darwin



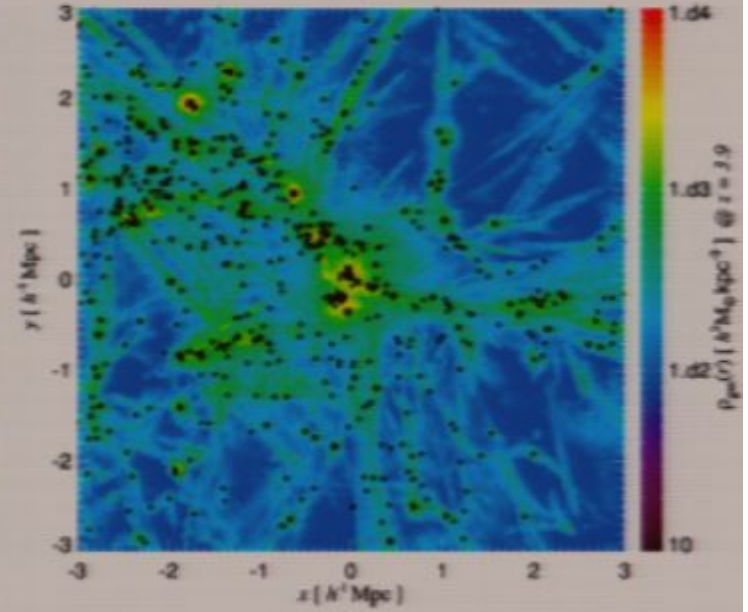
$\kappa_{\text{max}} [M_{\odot}/h]$ :  $\bullet$   $10^2-10^4$   $\bullet$   $10^4-10^7$   $\bullet$   $10^7-10^8$   $\bullet$   $10^8-10^9$   $\bullet$   $> 10^9$



$z=7.9$

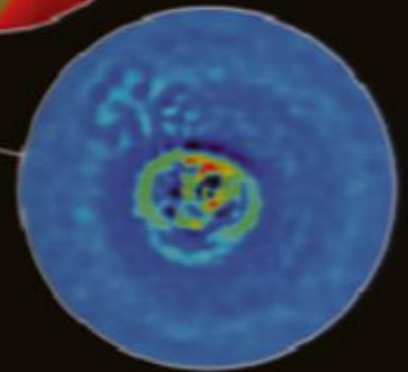
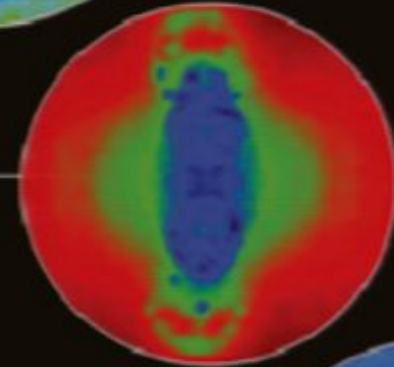
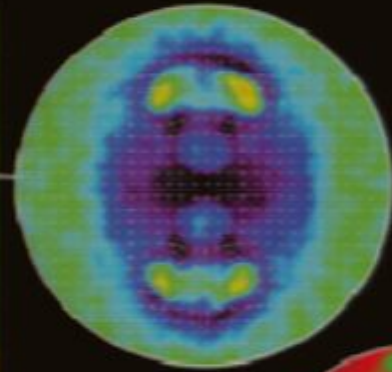
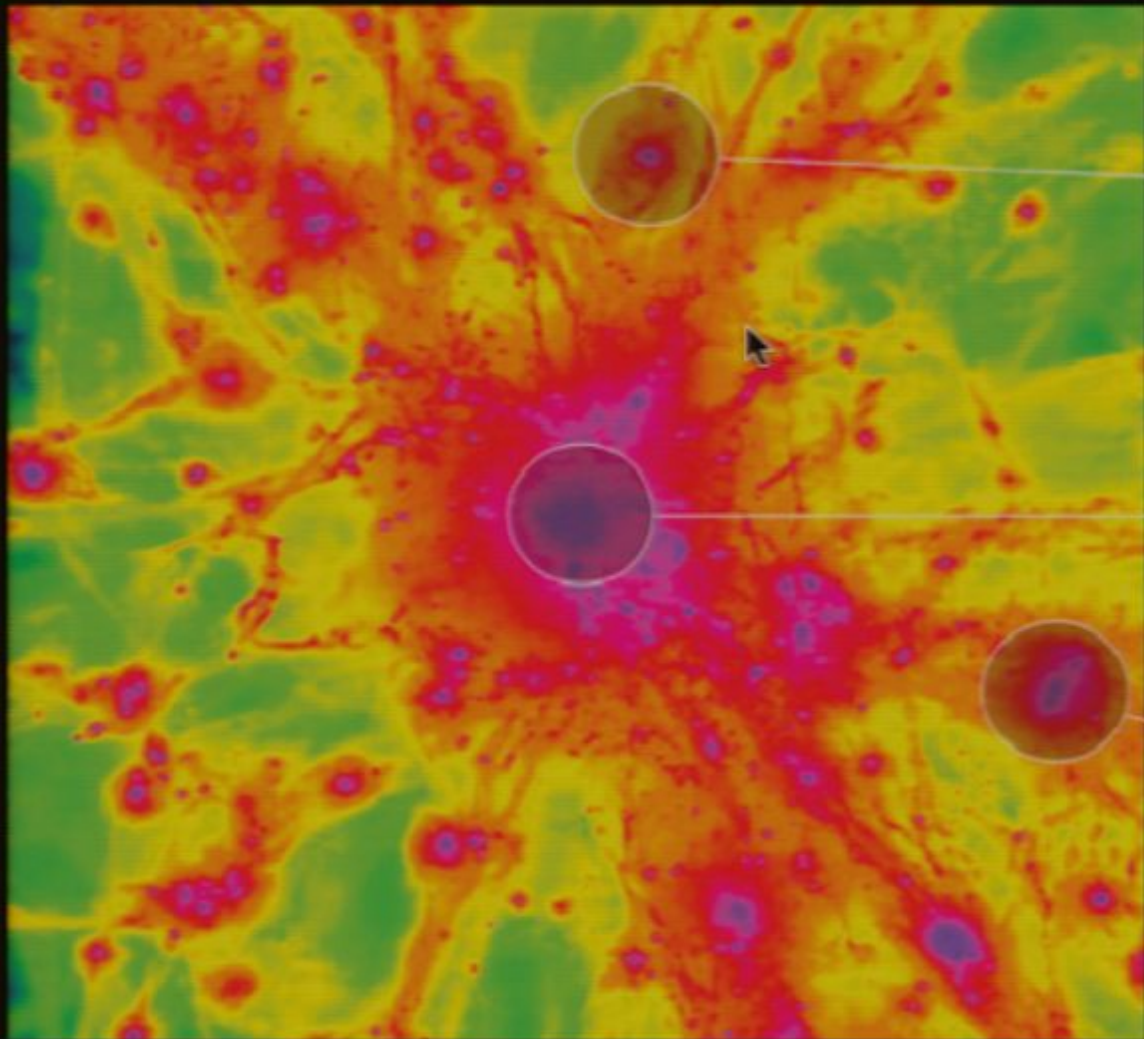


$z=6.2$



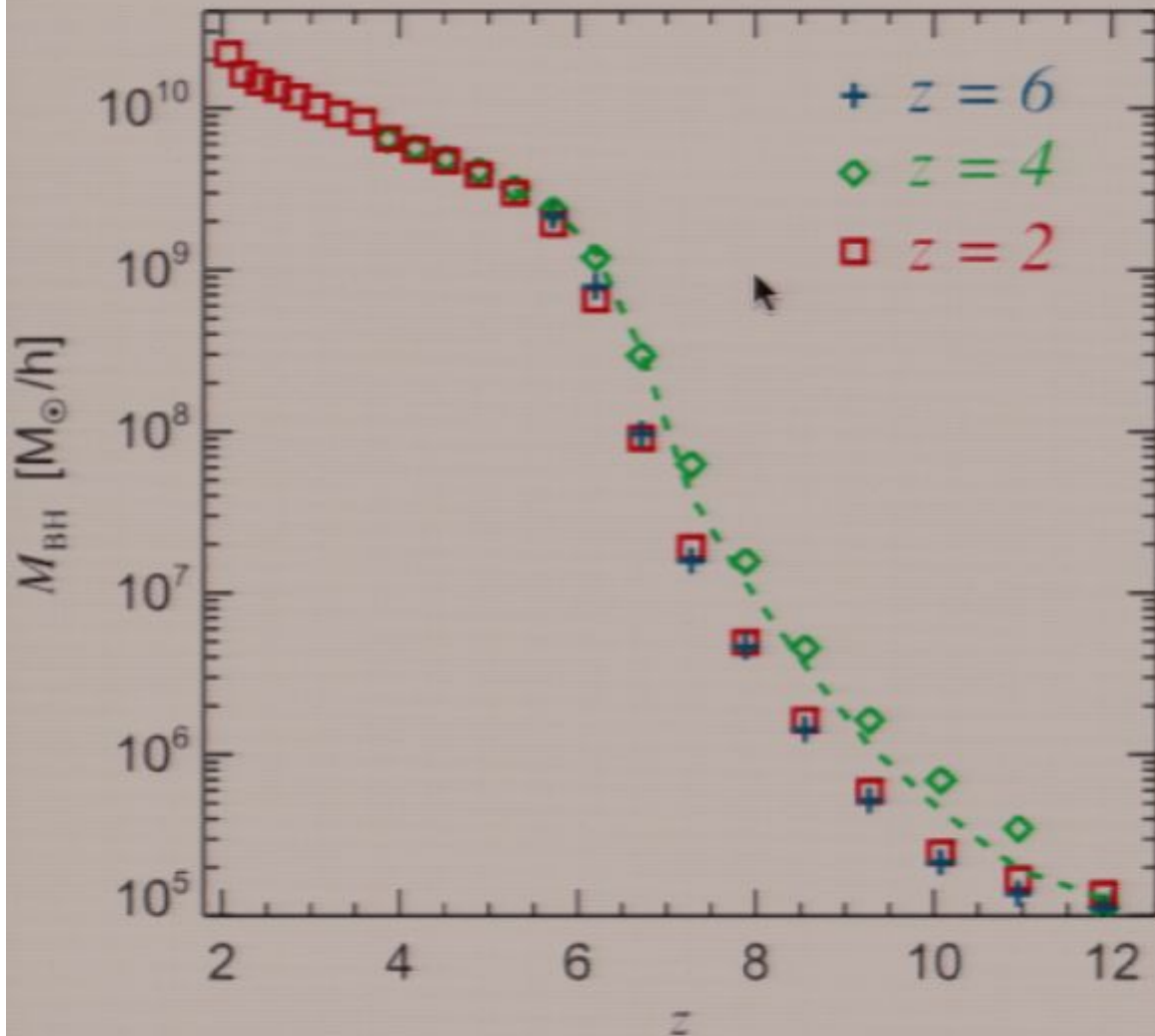
$z=3.9$





Sijacki

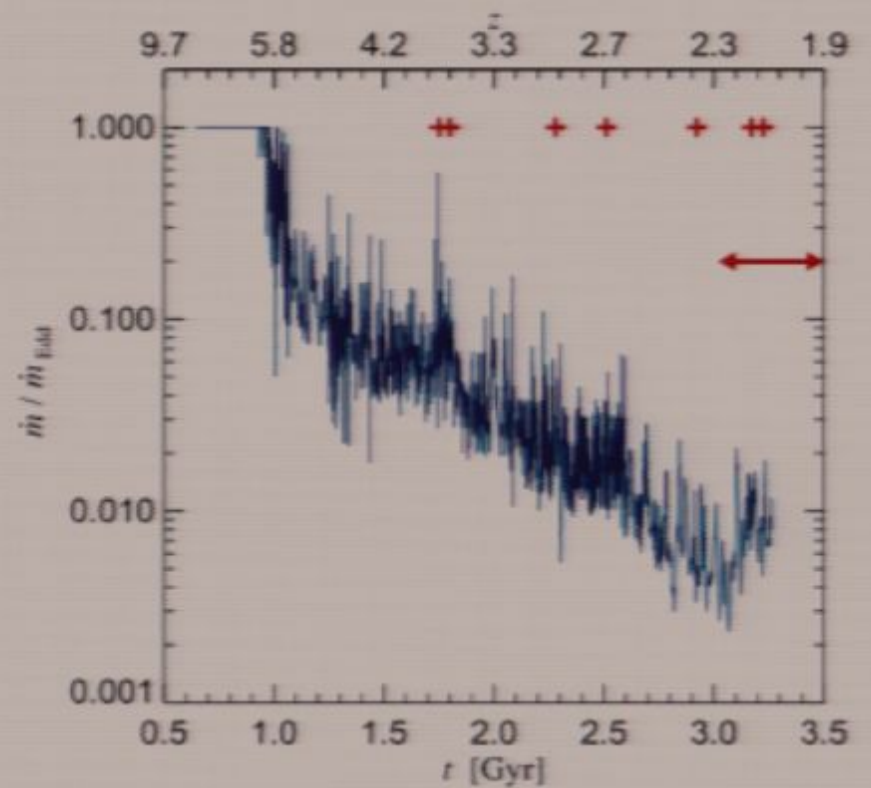
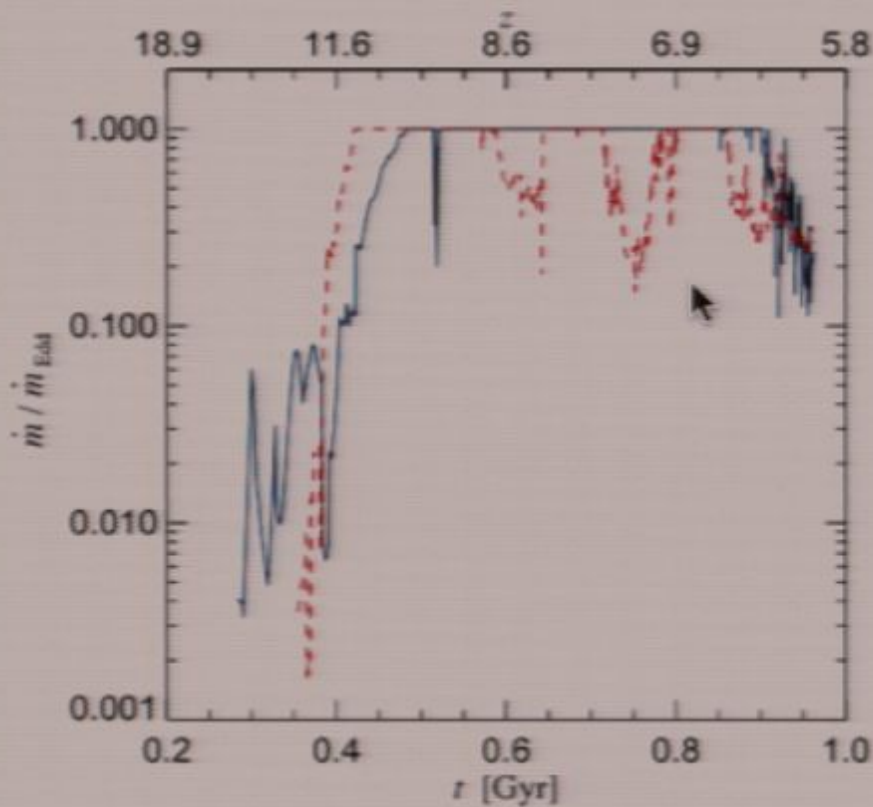
It is now possible to include most of the relevant physics.



By  $z=6$  the  $10^5 M_{\odot}$  seed black hole has grown into a supermassive black hole with mass several billion times that of the sun.

Sijacki, Springel & Haehnelt 2009



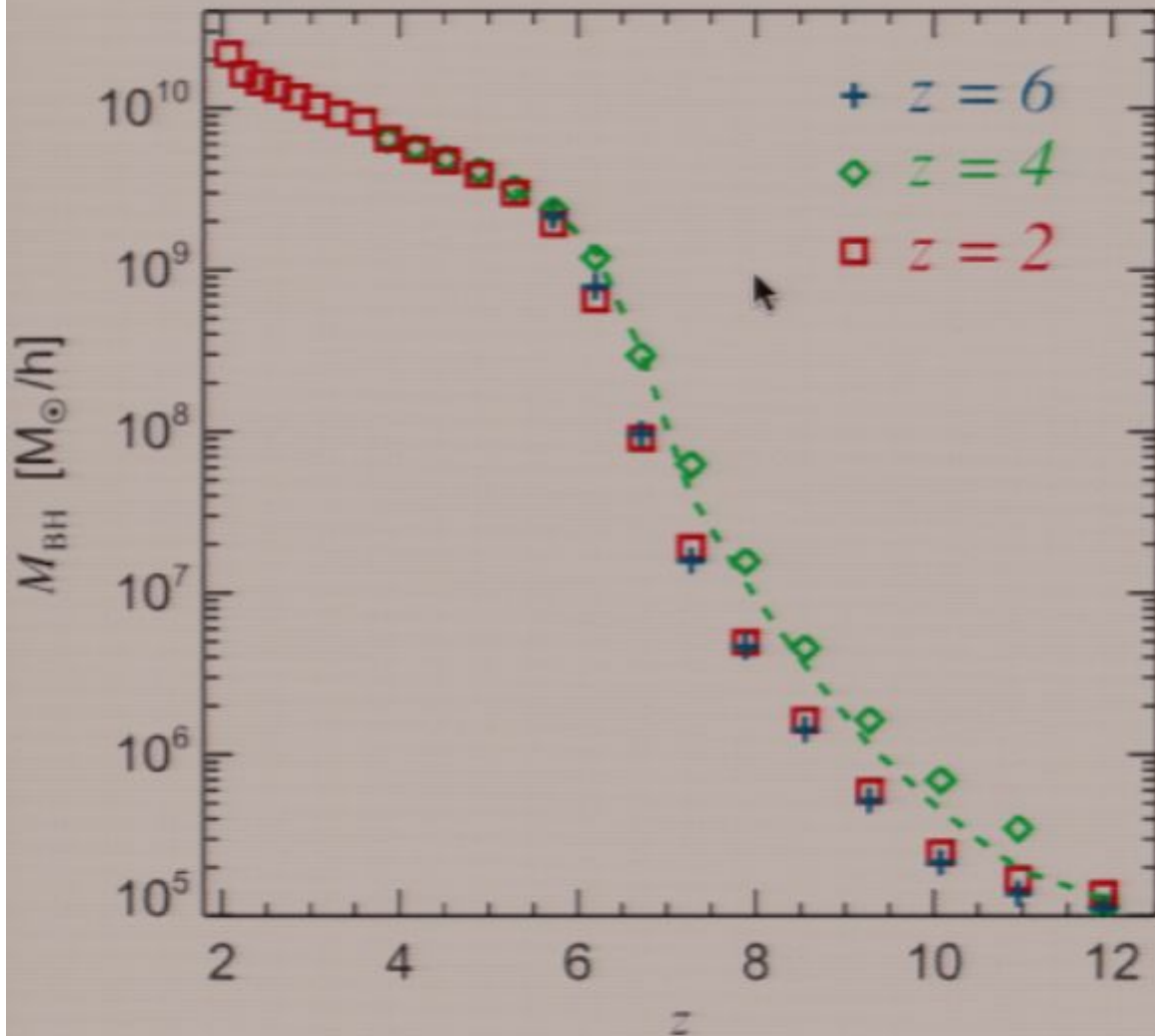


Sijacki, Springel & Haehnelt 2009

Early on the accretion rate is at the Eddington limit for more than 50% of the time.





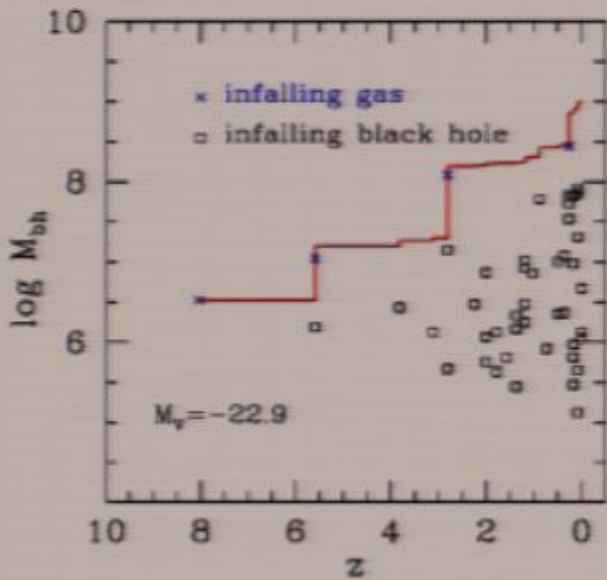


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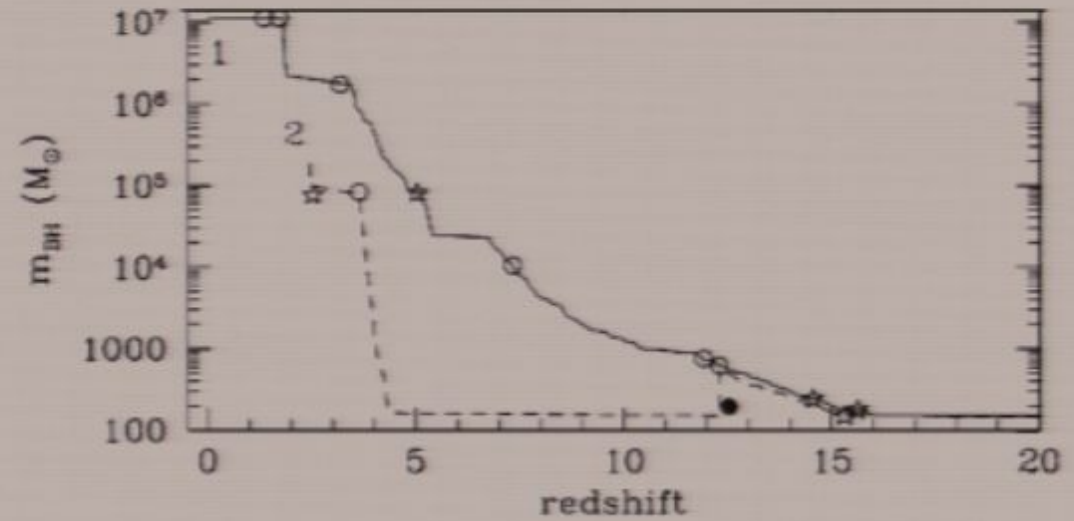
Sijacki, Springel & Haehnelt 2009







Kauffmann & Haehnelt



Volonteri et al.

Does the hierarchical growth start with a minimum seed mass?

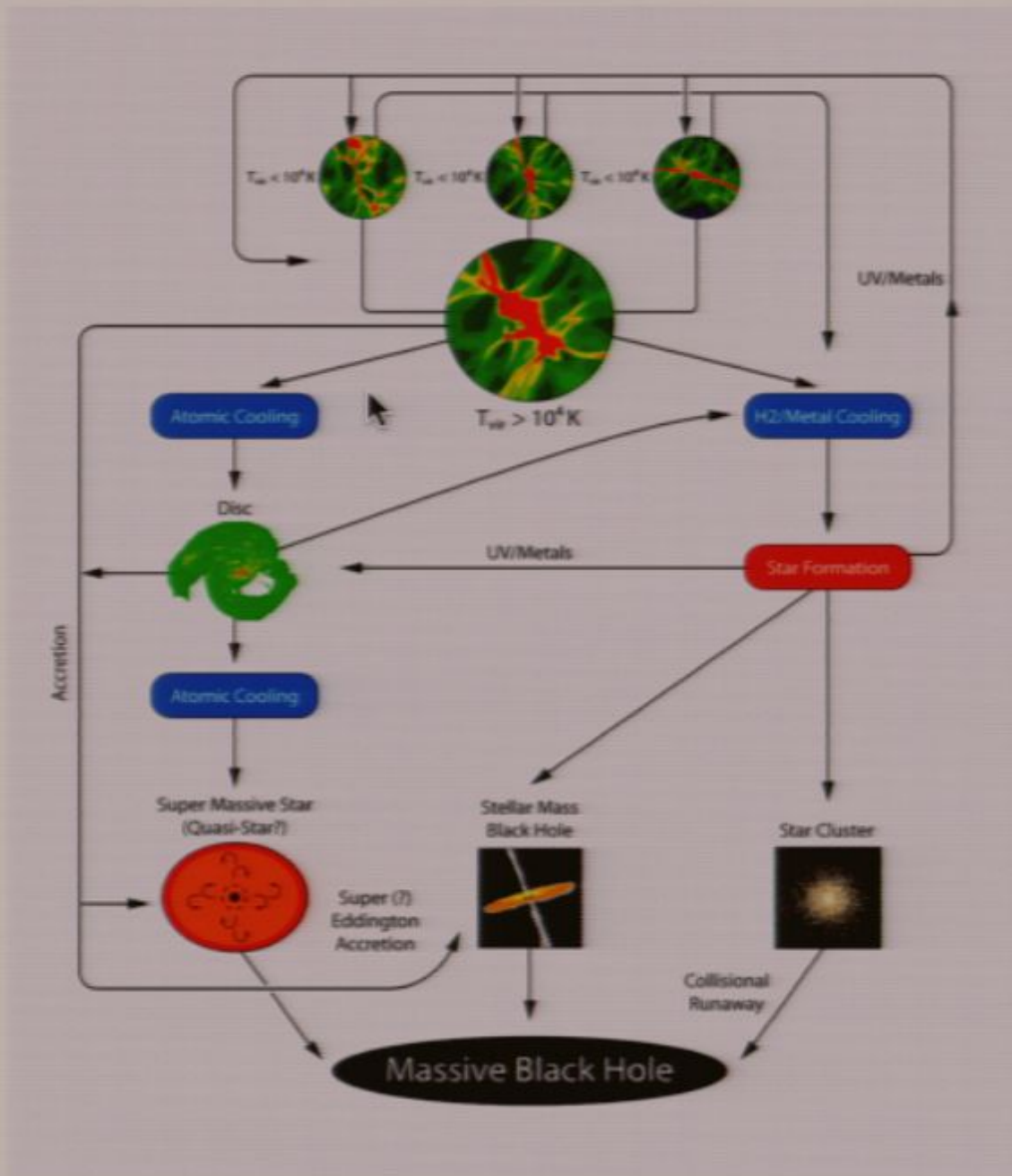
or

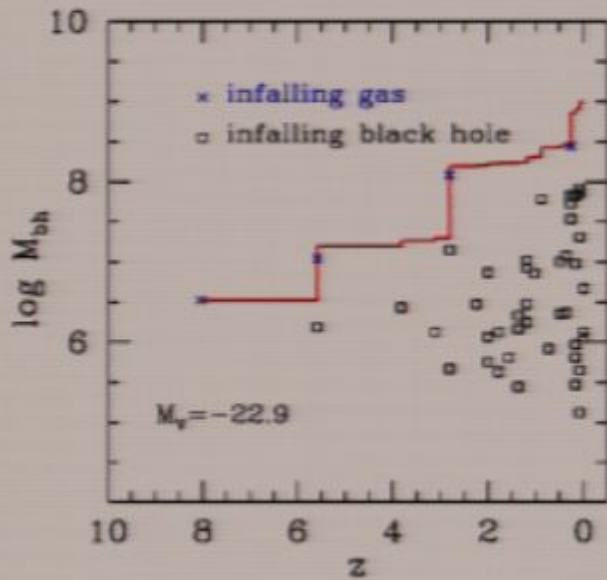
Does the hierarchical assembly extend all the way to stellar mass black holes?



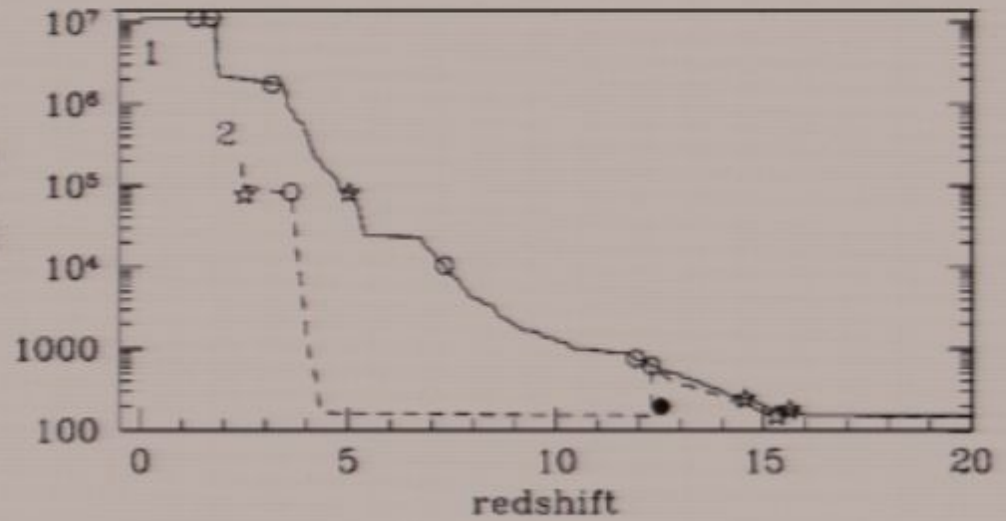
We most probably need massive seed black holes.  
How do massive seed black holes form?







Kauffmann & Haehnelt



Volonteri et al.

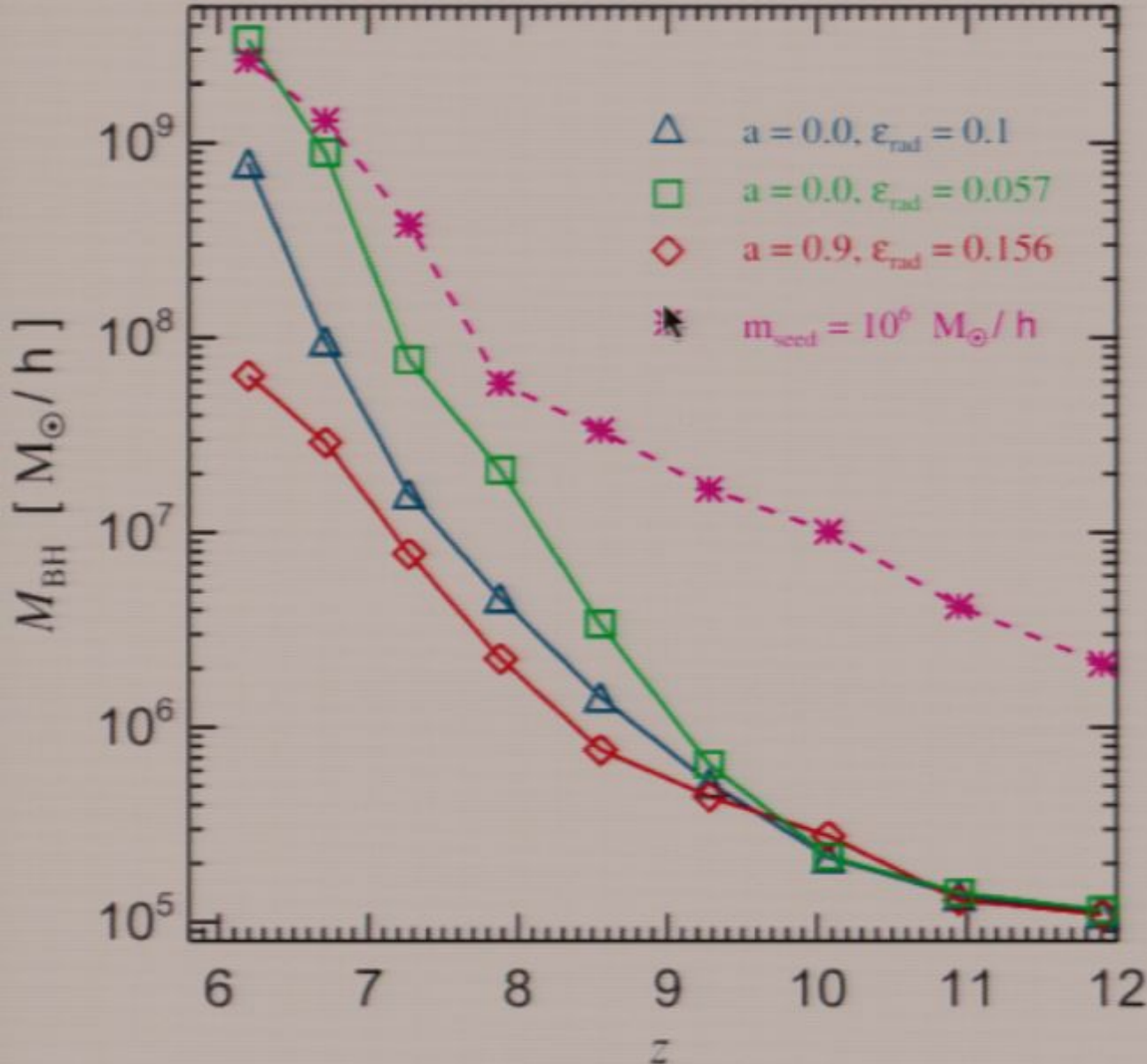
Does the hierarchical growth start with a minimum seed mass?

or

Does the hierarchical assembly extend all the way to stellar mass black holes?

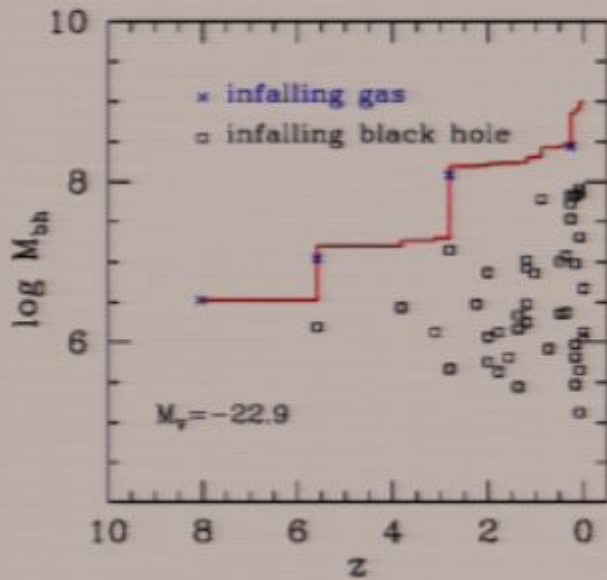




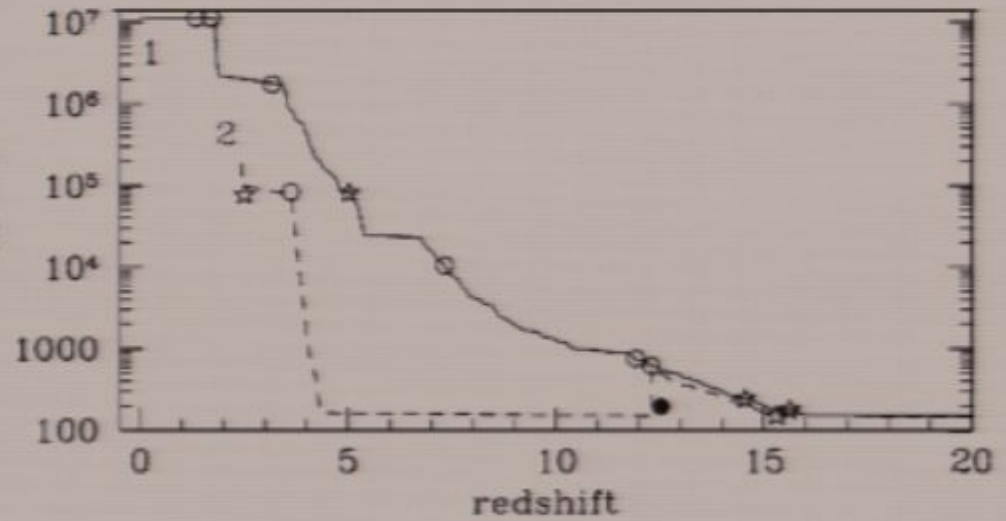


The mass at  $z=6$  depends strongly on the radiative efficiency of the accretion and therefore on the spin of the black hole.





Kauffmann & Haehnelt



Volonteri et al.

Does the hierarchical growth start with a minimum seed mass?

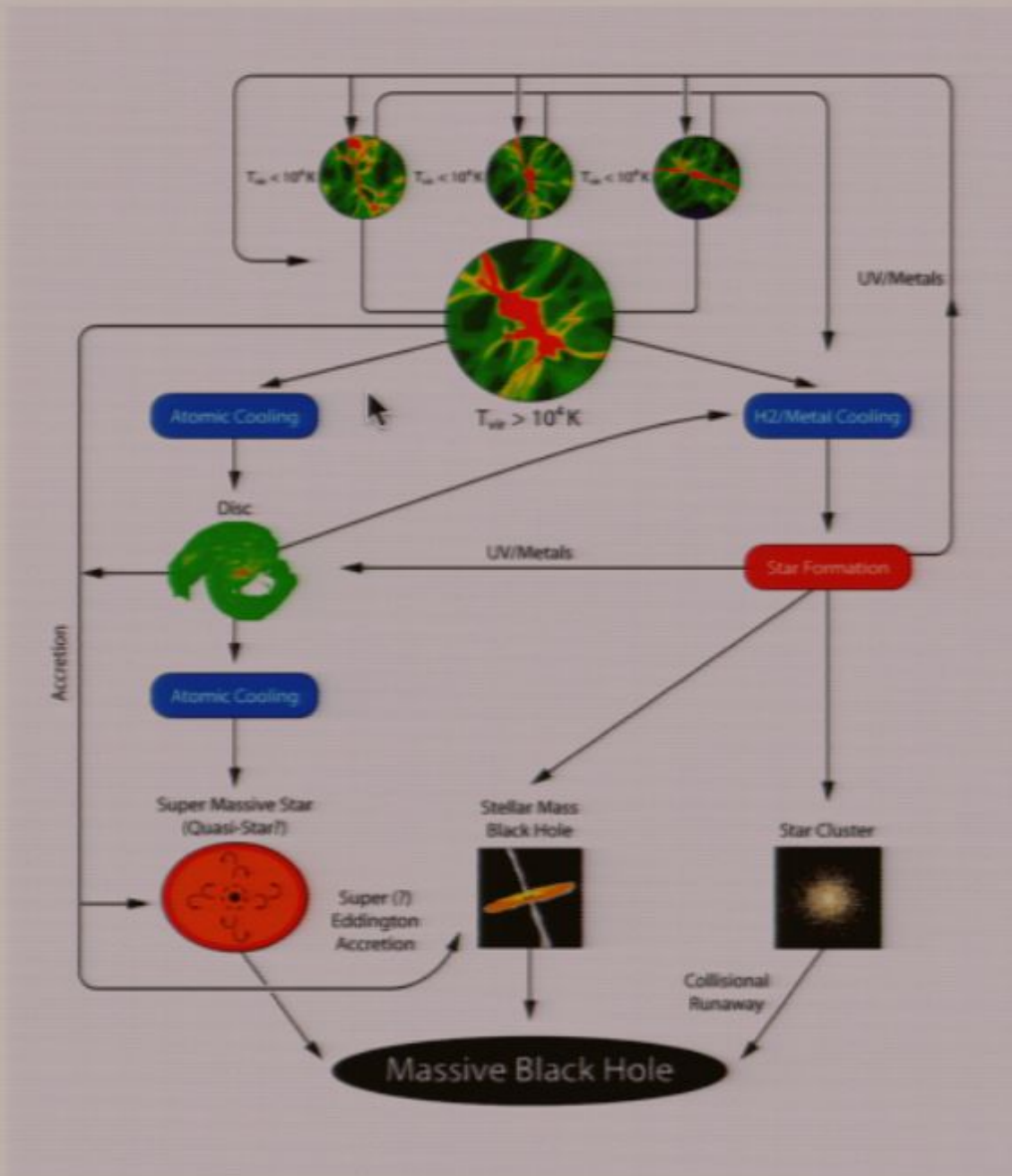
or

Does the hierarchical assembly extend all the way to stellar mass black holes?



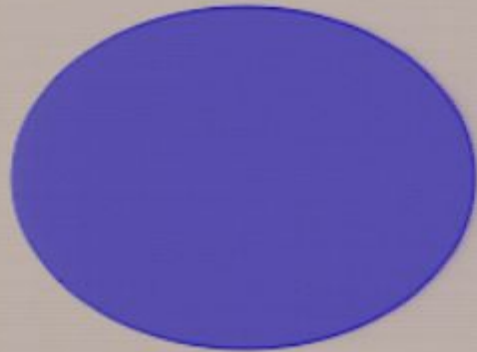
We most probably need massive seed black holes.  
How do massive seed black holes form?







# Direct collapse into a compact massive self-gravitating disc



haloes with  $T_{\text{vir}} \geq 10000\text{K}$   
with no metals (and H<sub>2</sub> suppression)  
are least prone to fragmentation



## FORMATION OF THE FIRST SUPERMASSIVE BLACK HOLES

VOLKER BRUNN<sup>1</sup> AND ABRAHAM LOEB<sup>1,2,3</sup>

Received 2002 December 18; accepted 2003 June 16

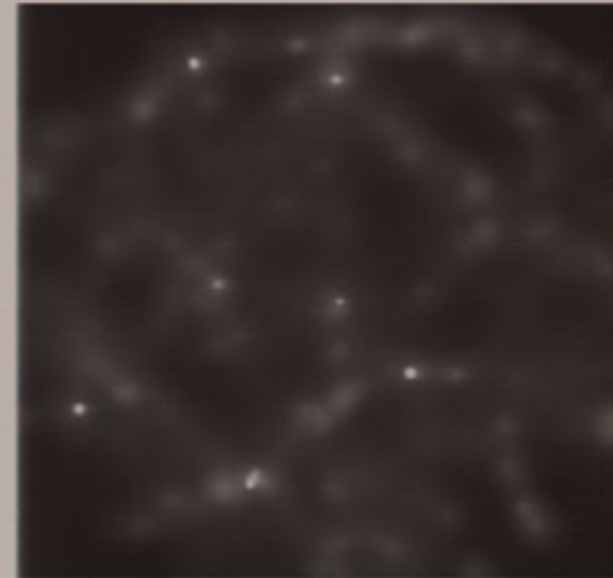
### ABSTRACT

We consider the physical conditions under which supermassive black holes could have formed inside the first galaxies. Our smoothed particle hydrodynamics simulations indicate that metal-free galaxies with a virial temperature of  $\sim 10^4$  K and suppressed  $H_2$  formation (due to an intergalactic UV background) tend to form a binary black hole system that contains a substantial fraction ( $\gtrsim 10\%$ ) of the total baryonic mass of the host galaxy. Fragmentation into stars is suppressed without substantial  $H_2$  cooling. Our simulations follow the condensation of  $\sim 5 \times 10^6 M_\odot$  around the two centers of the binary system to a scale of  $\lesssim 0.1$  pc. Low-spin galaxies form a single black hole instead. These early black holes lead to quasar activity before the epoch of reionization. Primordial black hole binaries lead to gravitational radiation emission at redshifts  $z \gtrsim 10$  that would be detectable by *Laser Interferometer Space Antenna*.

*Subject headings:* black hole physics — cosmology: theory — galaxies: formation — hydrodynamics — quasars: general

*On-line material:* color figures

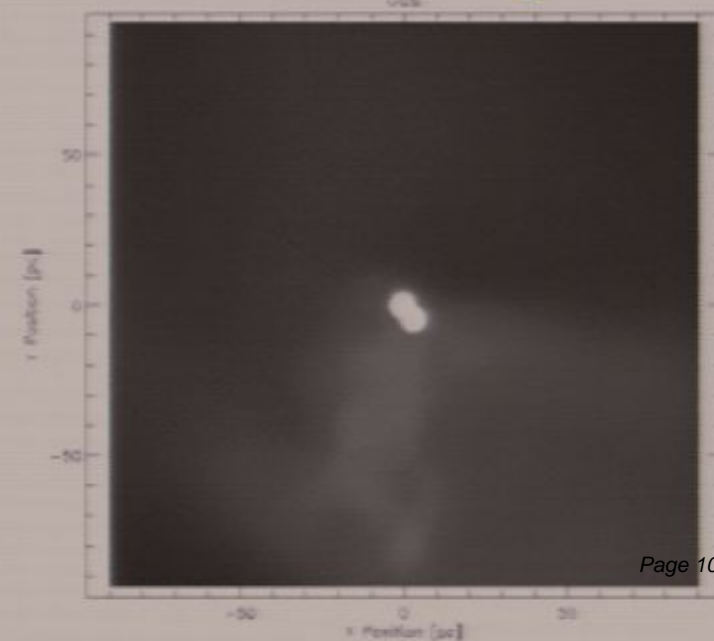
with  $H_2$



collapsing  $10^4$  K halo:

strongly suppressed  
fragmentation without  
 $H_2$  cooling

without  $H_2$



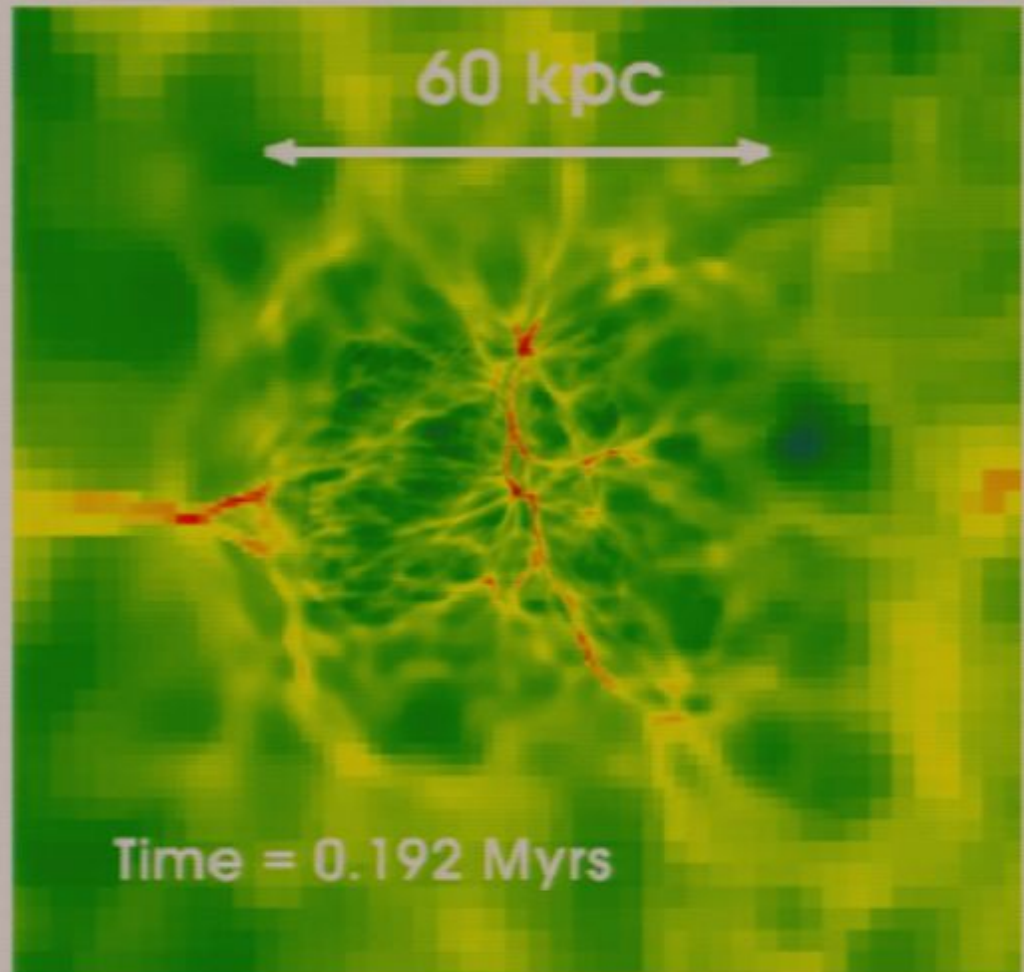
$$T_{\text{vir}} \sim 35000\text{K}$$

$$V_{\text{vir}} \sim 30\text{kms}^{-1}$$

$$M_{\text{tot}} \sim 3.3 \times 10^8 M_{\odot}$$

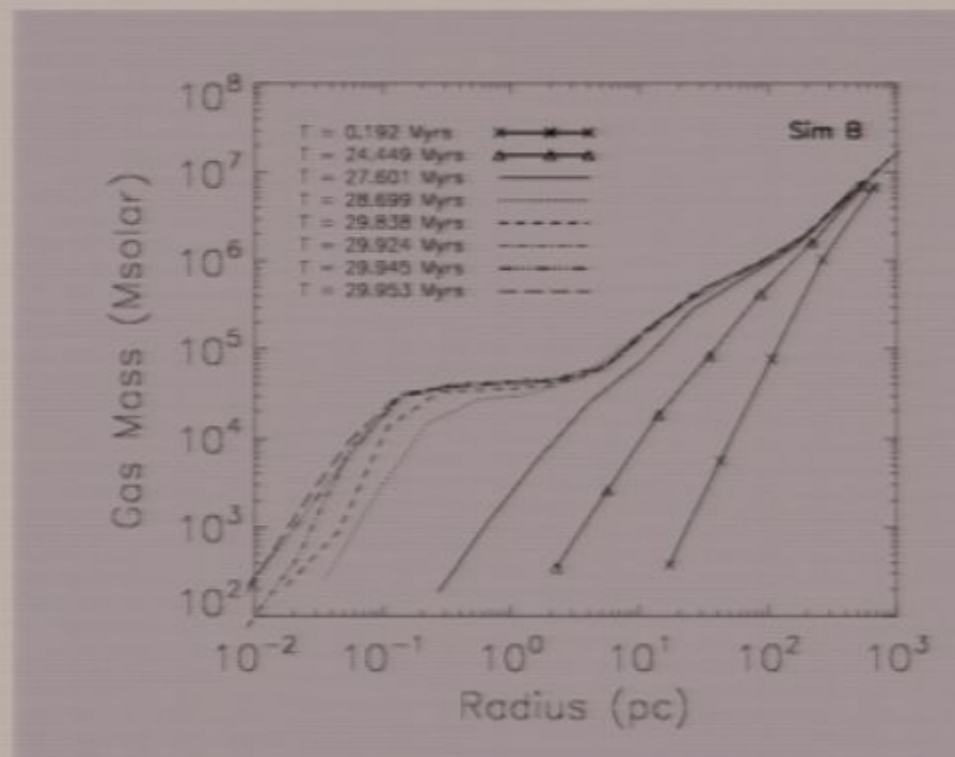
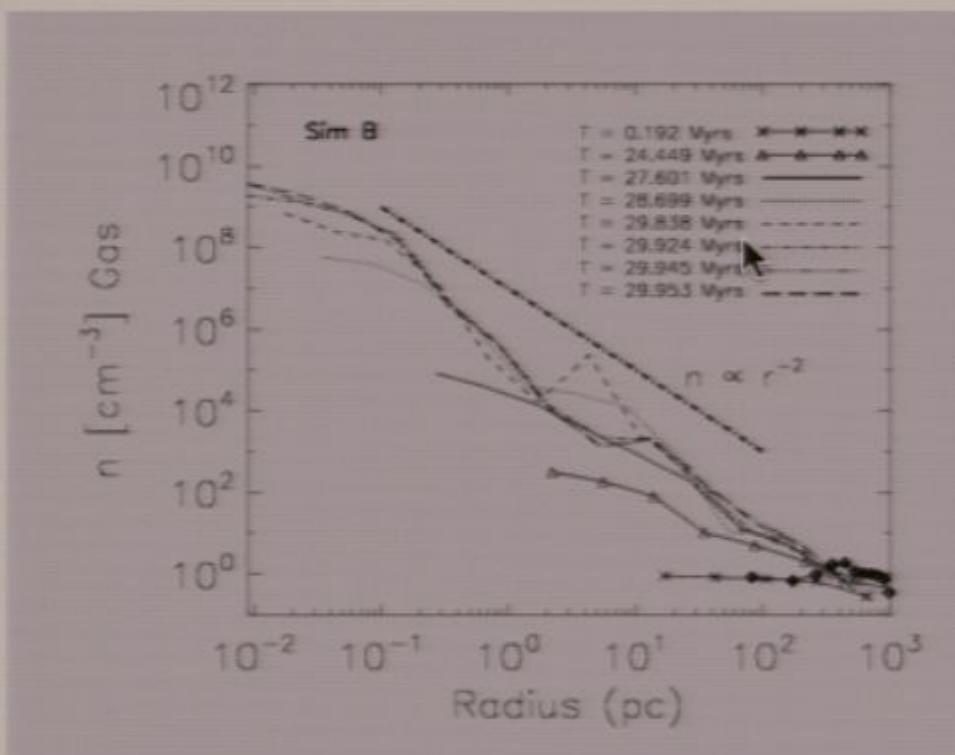
$$z \sim 15$$

no  $\text{H}_2$  cooling



Regan & Haehnelt 2009

## Isothermal collapse at $T \sim 7000\text{-}8000\text{K}$

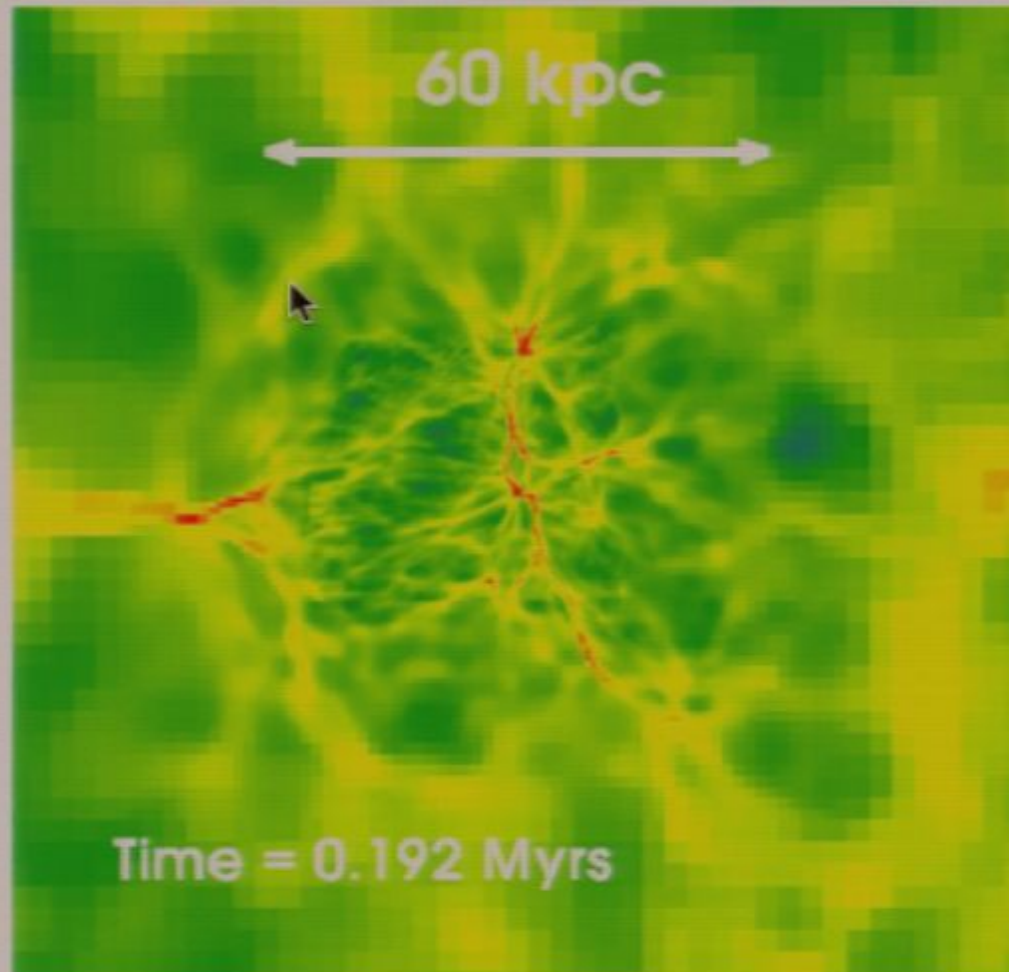


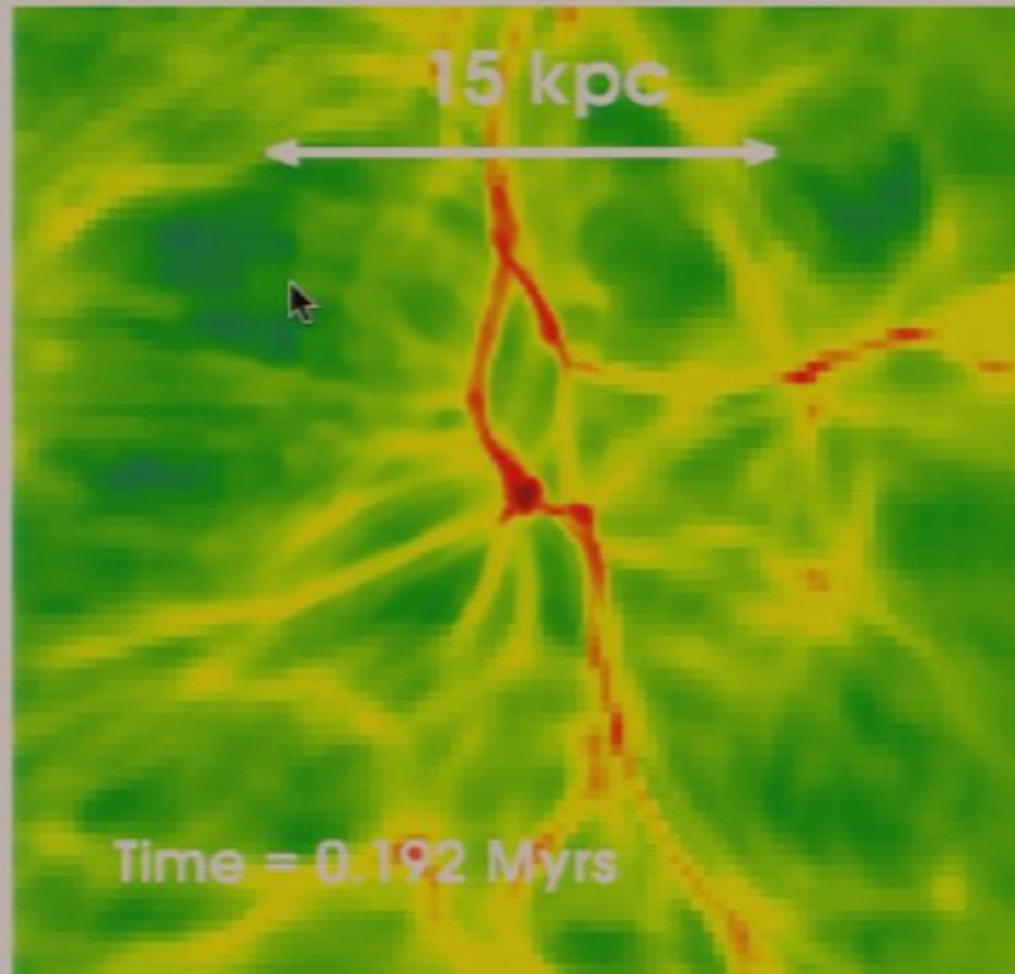
Regan & Haehnelt 2008

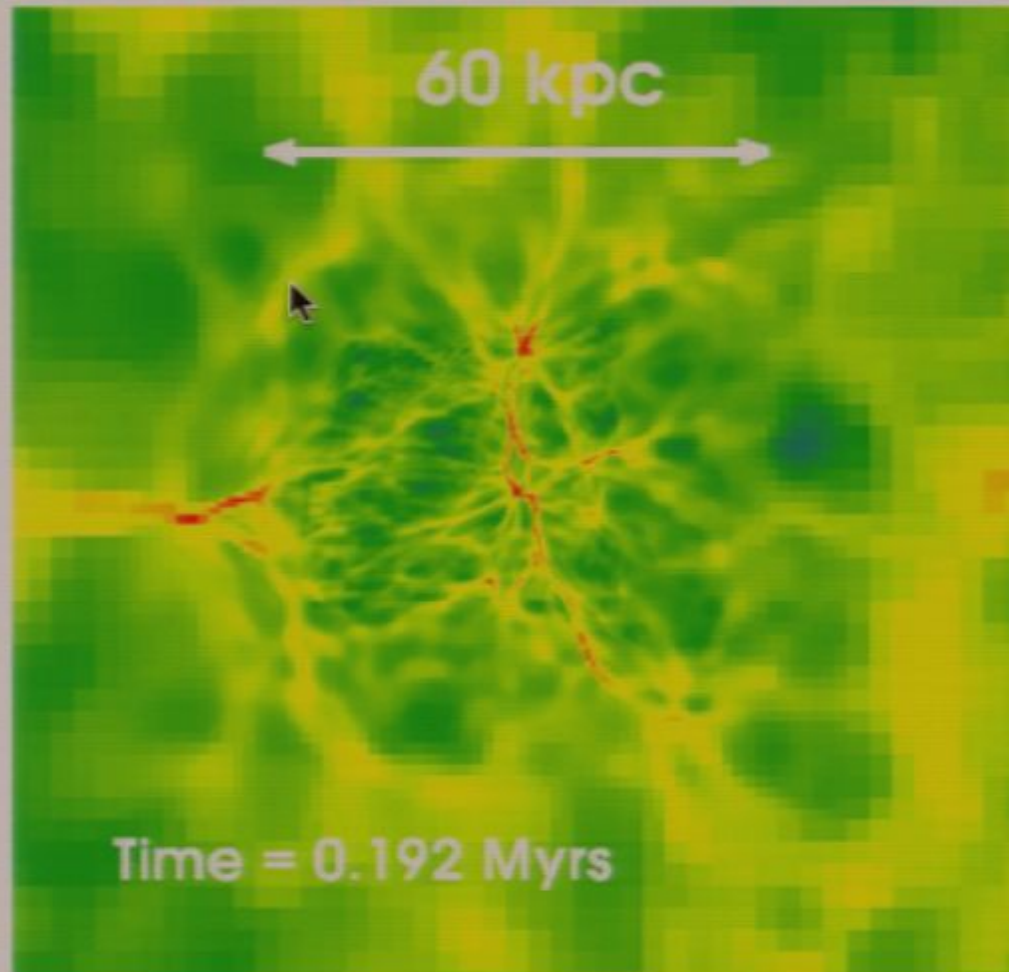
The inner  $2 \times 10^4 M_{\odot}$  collapse by a factor 1000 in radius before they settle into rotational support!

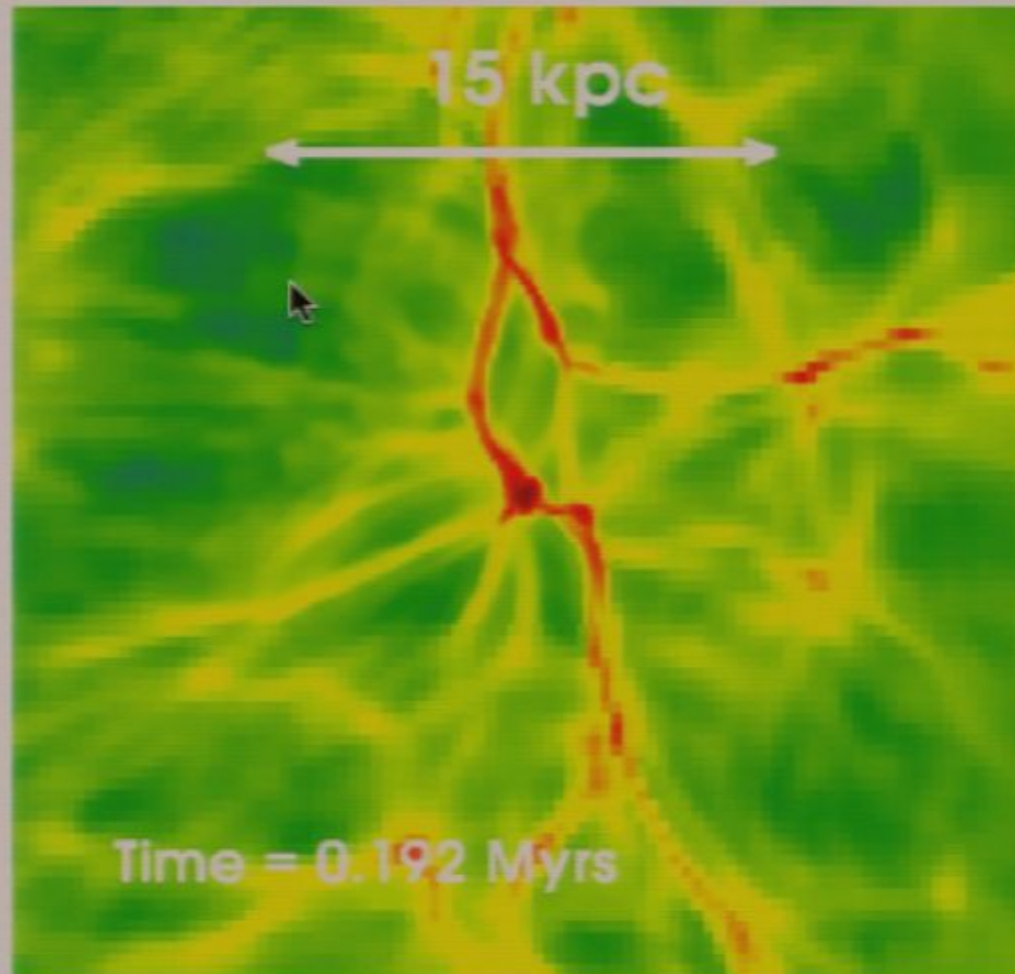




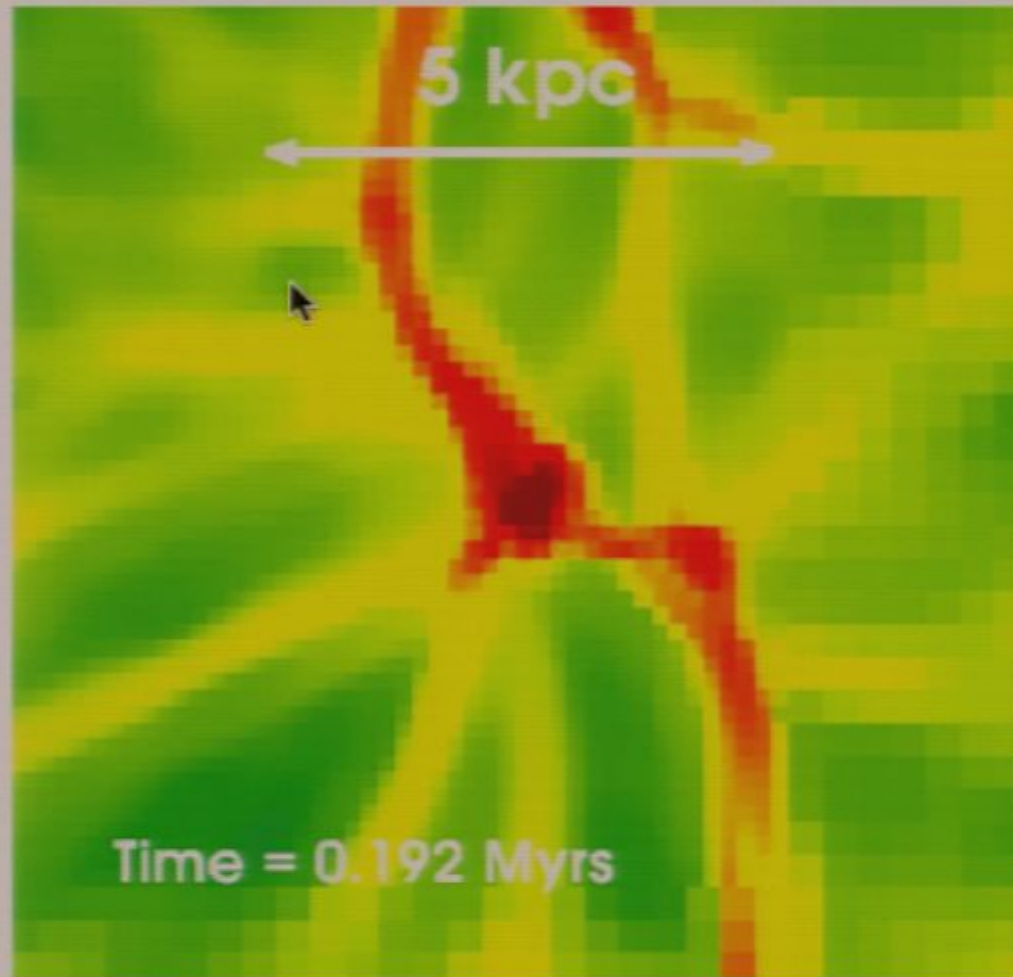


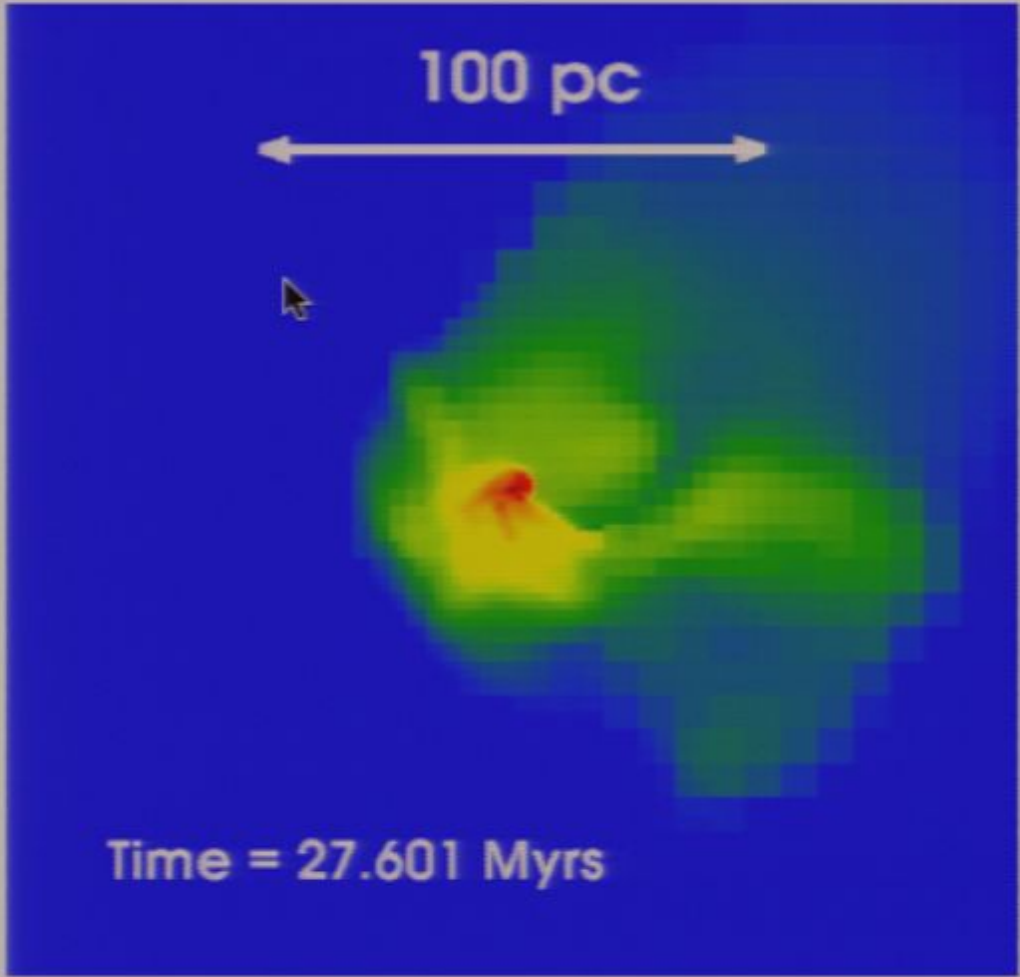


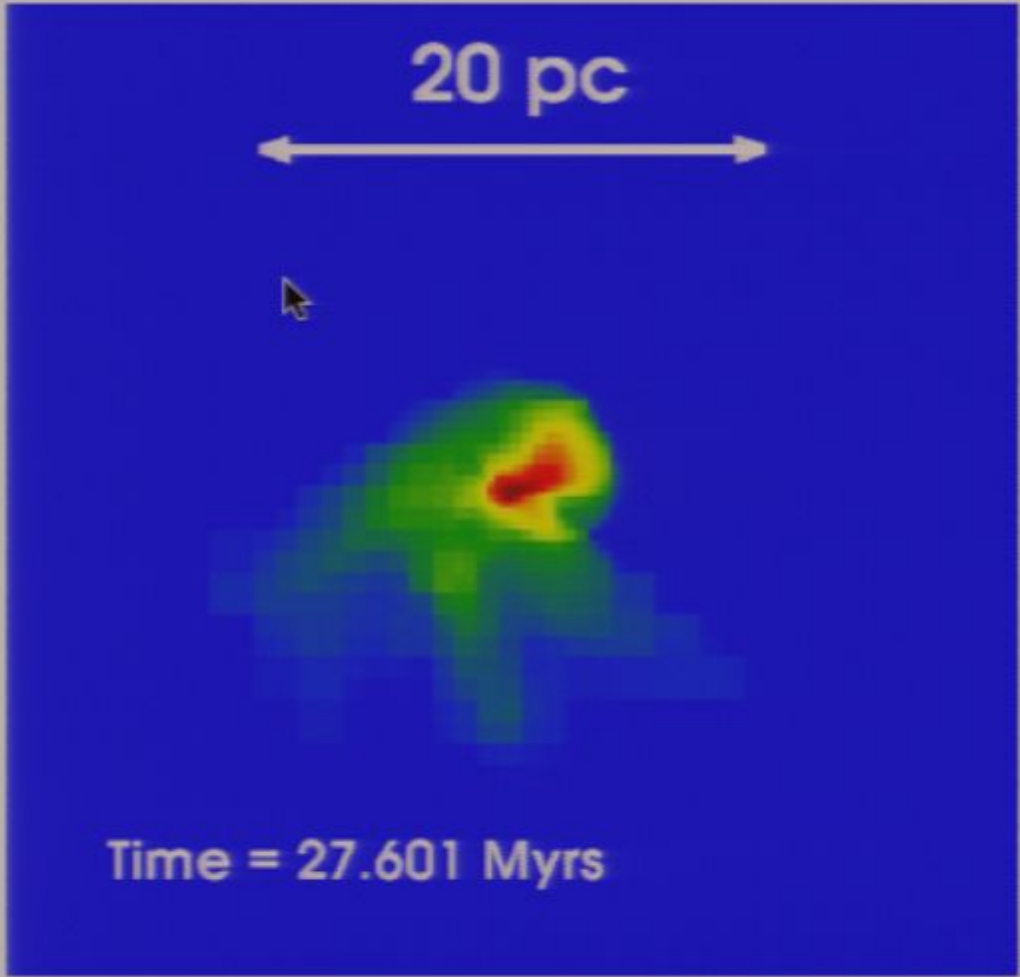


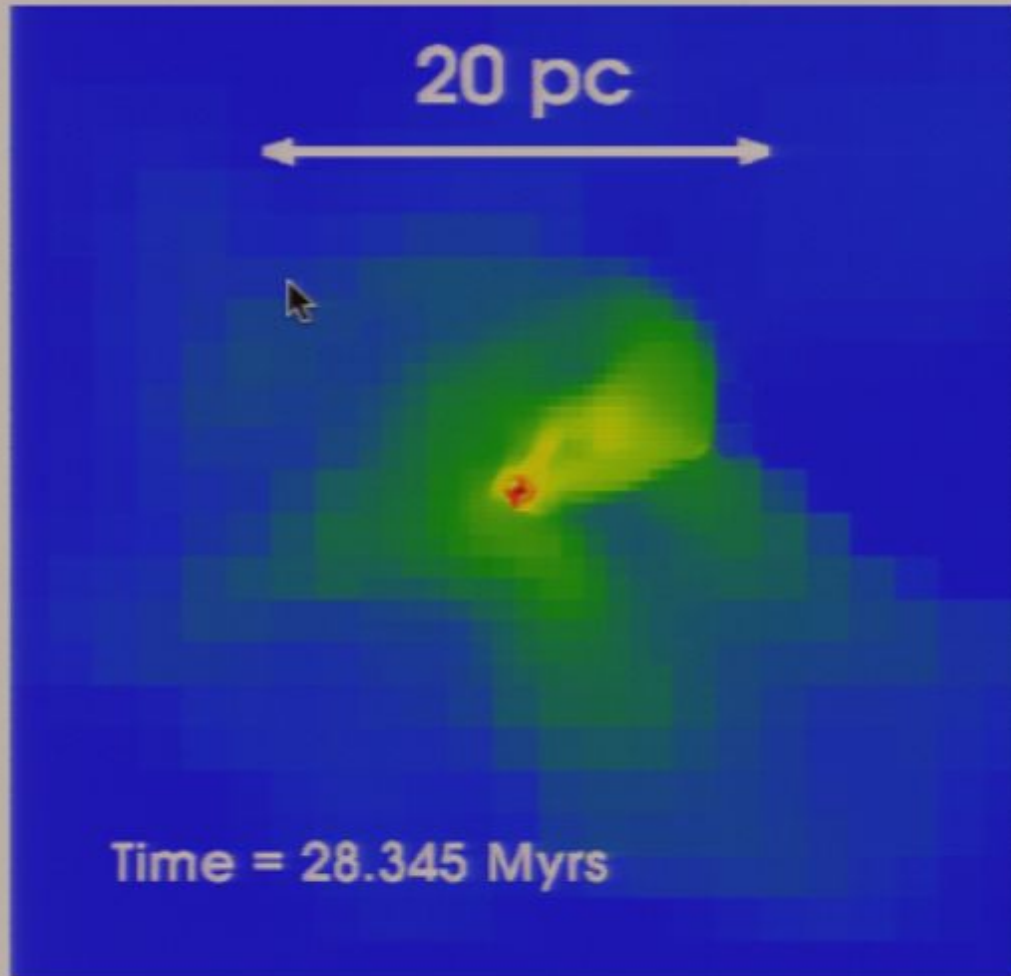




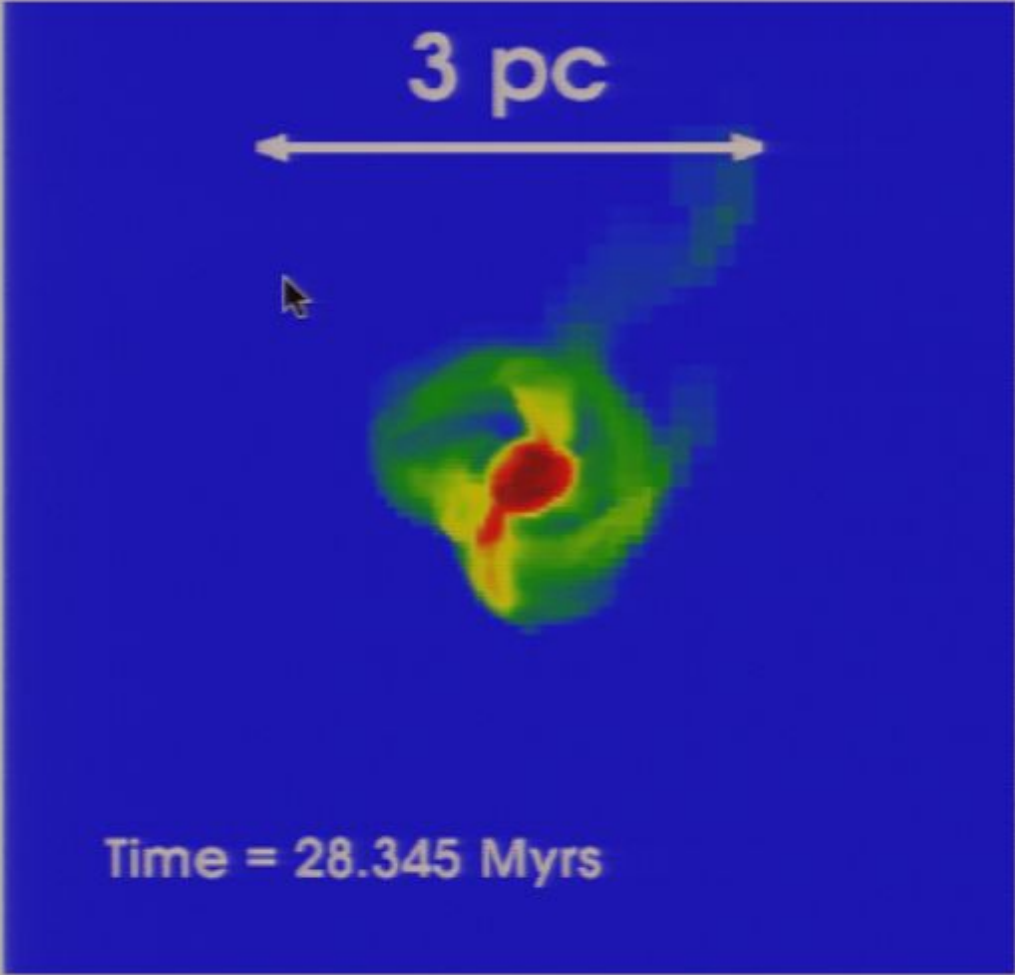


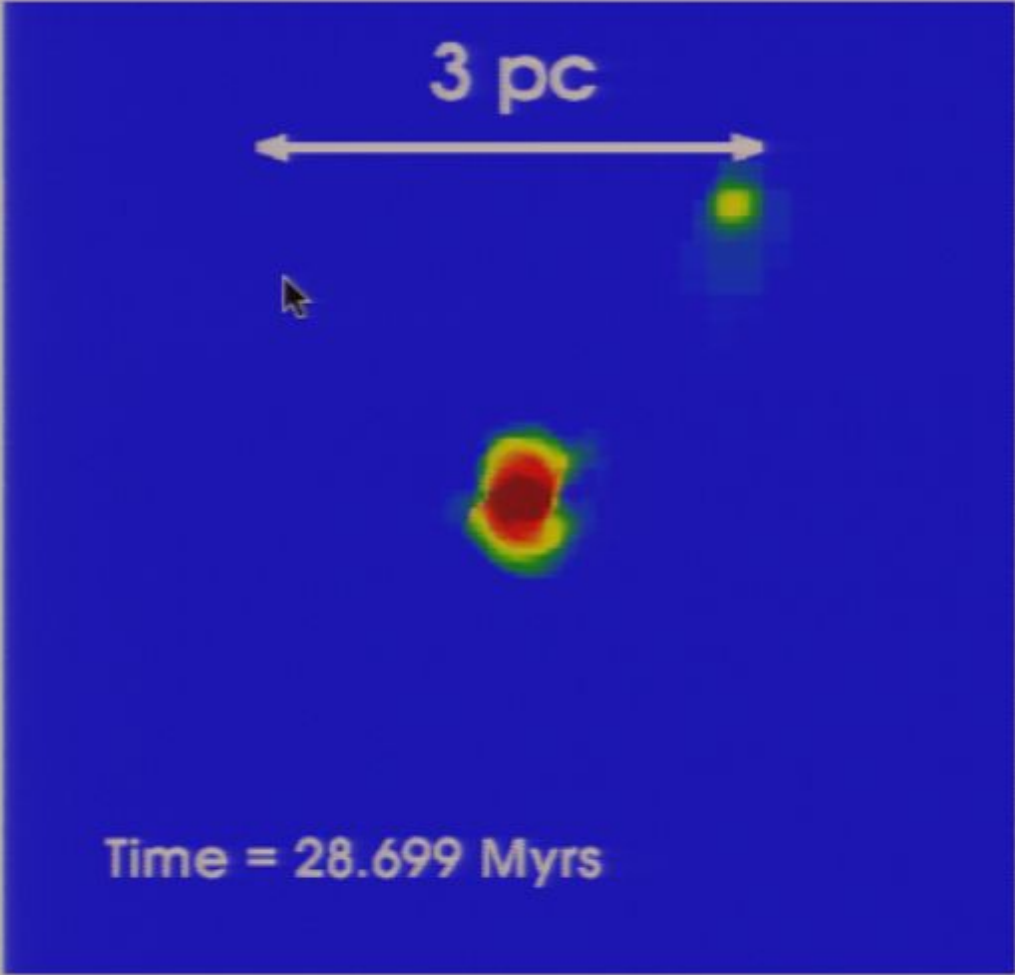


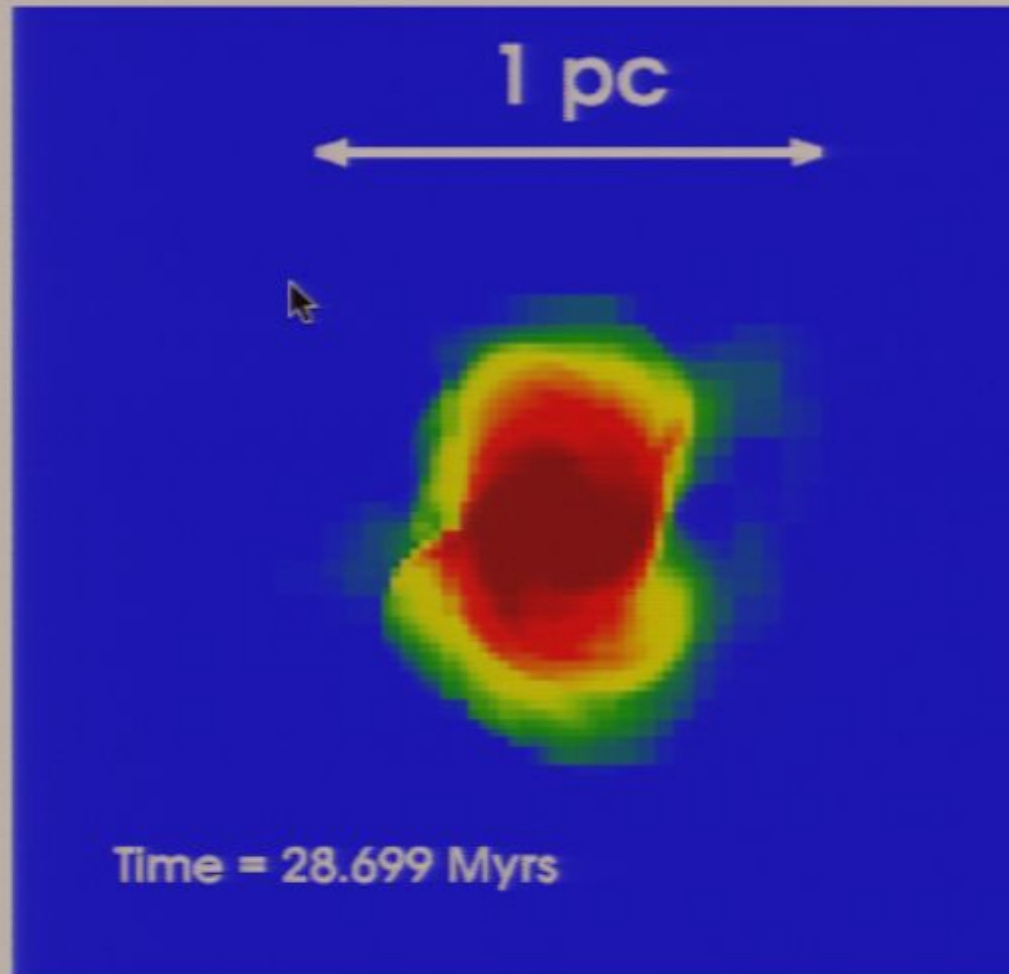


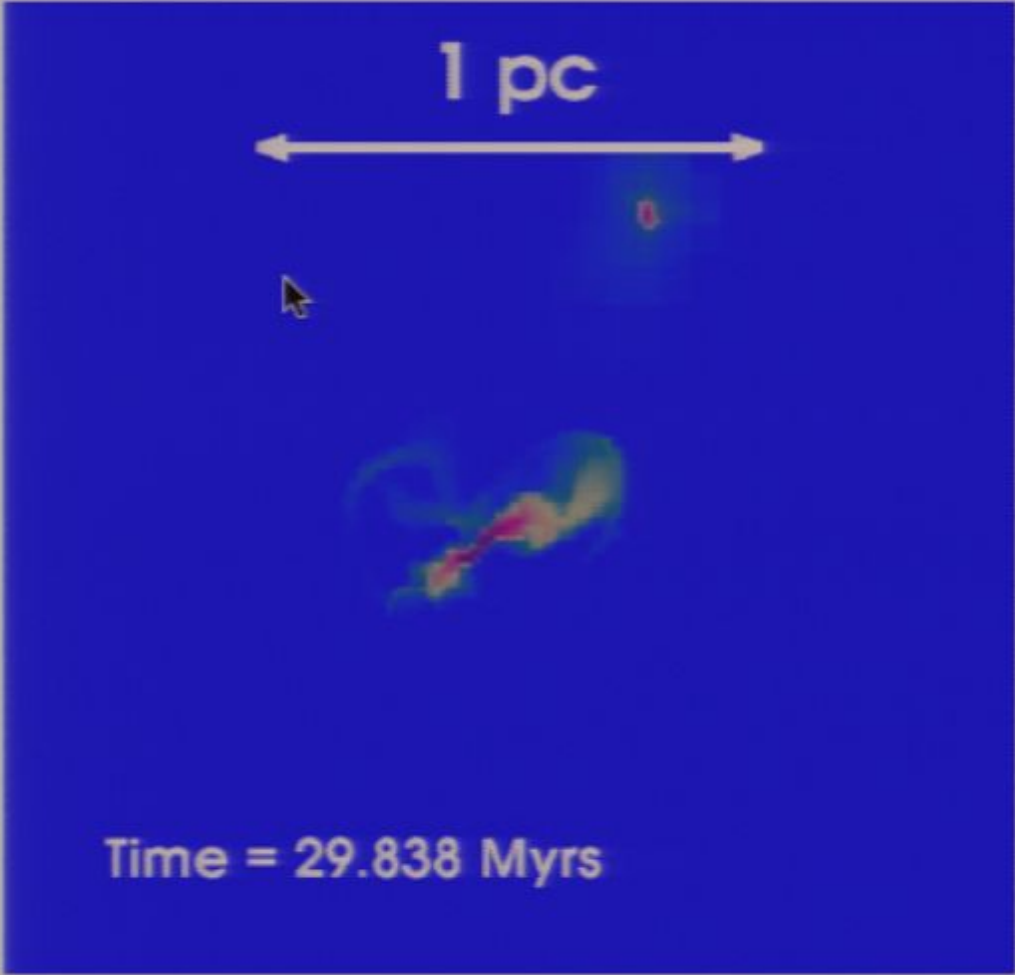




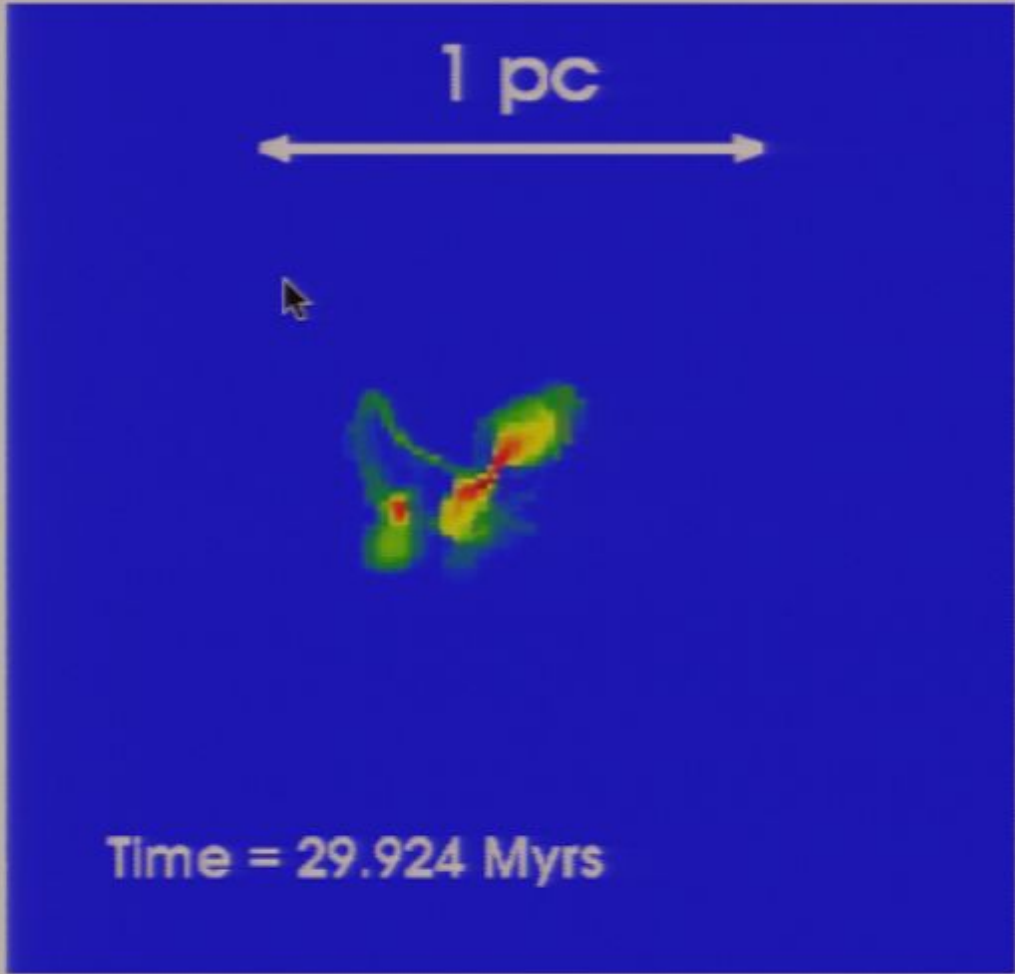


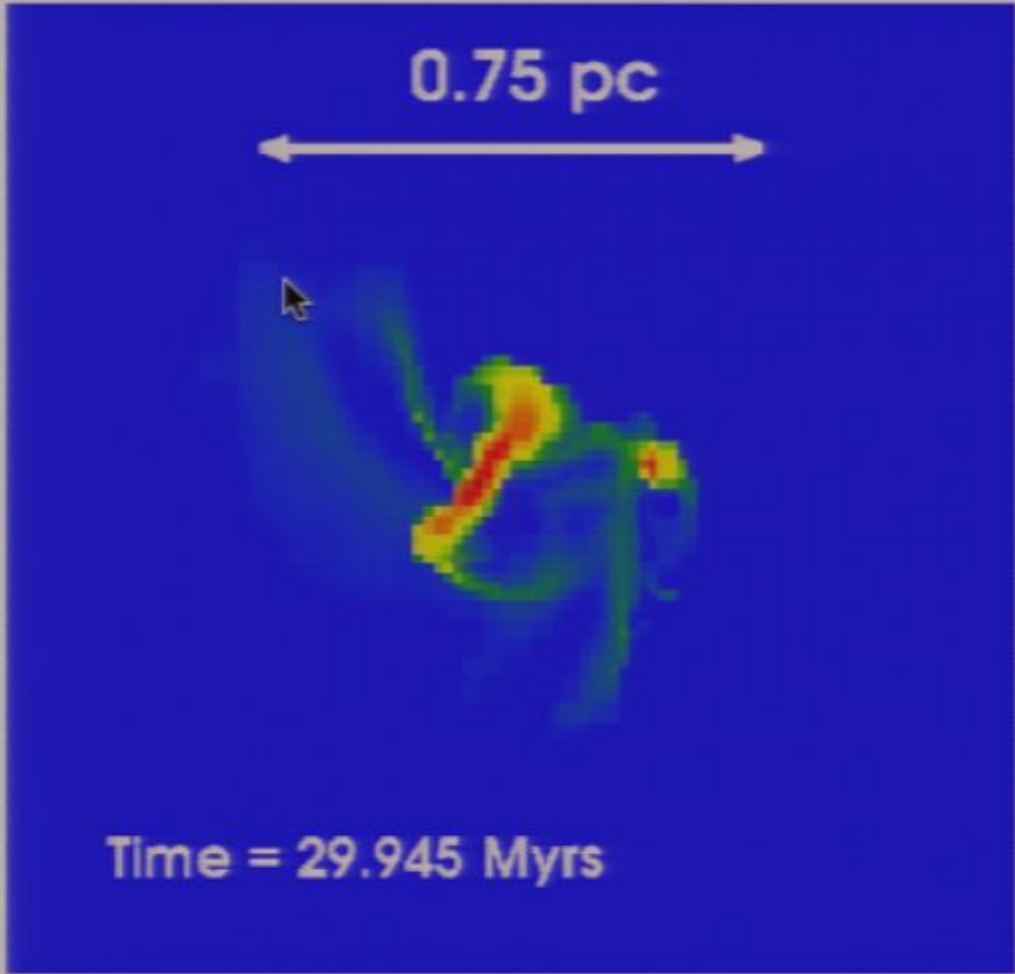


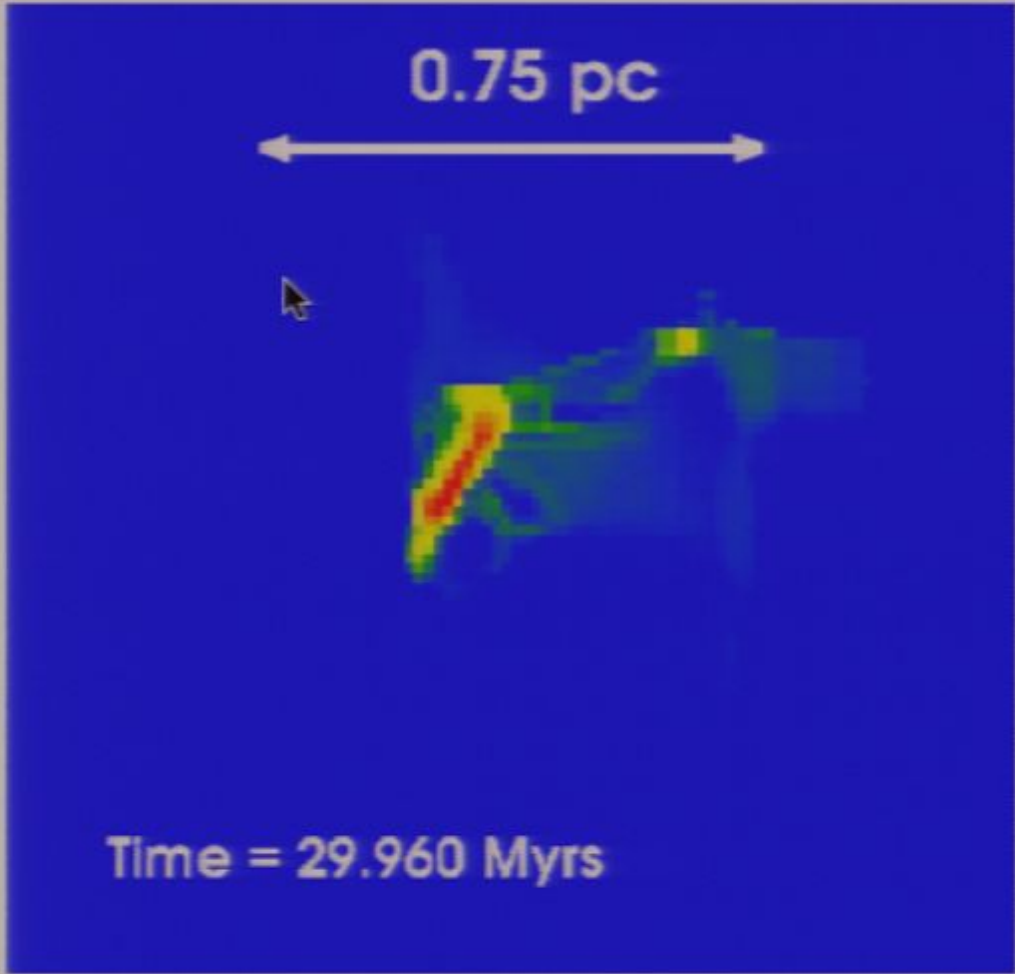


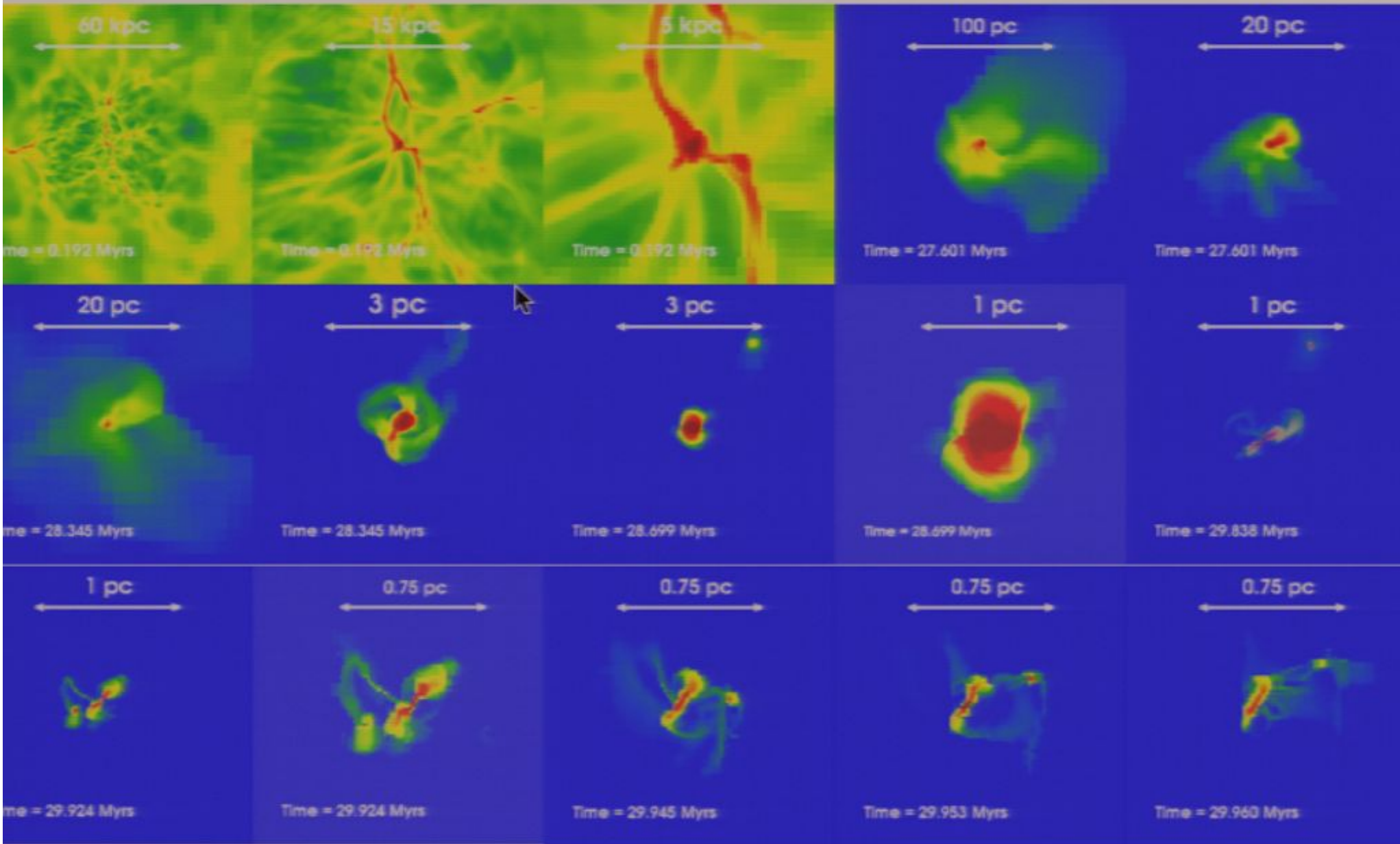




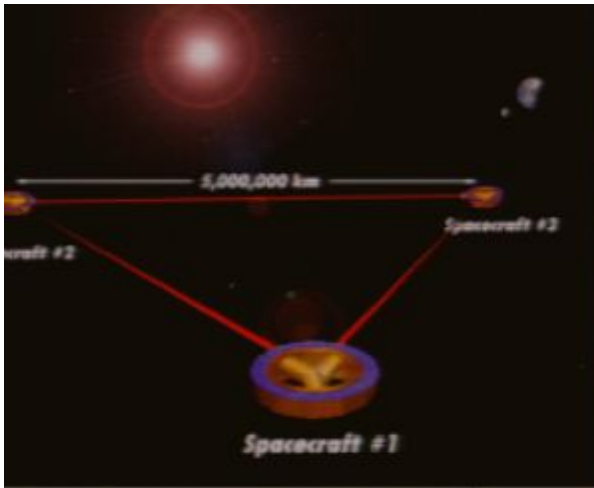




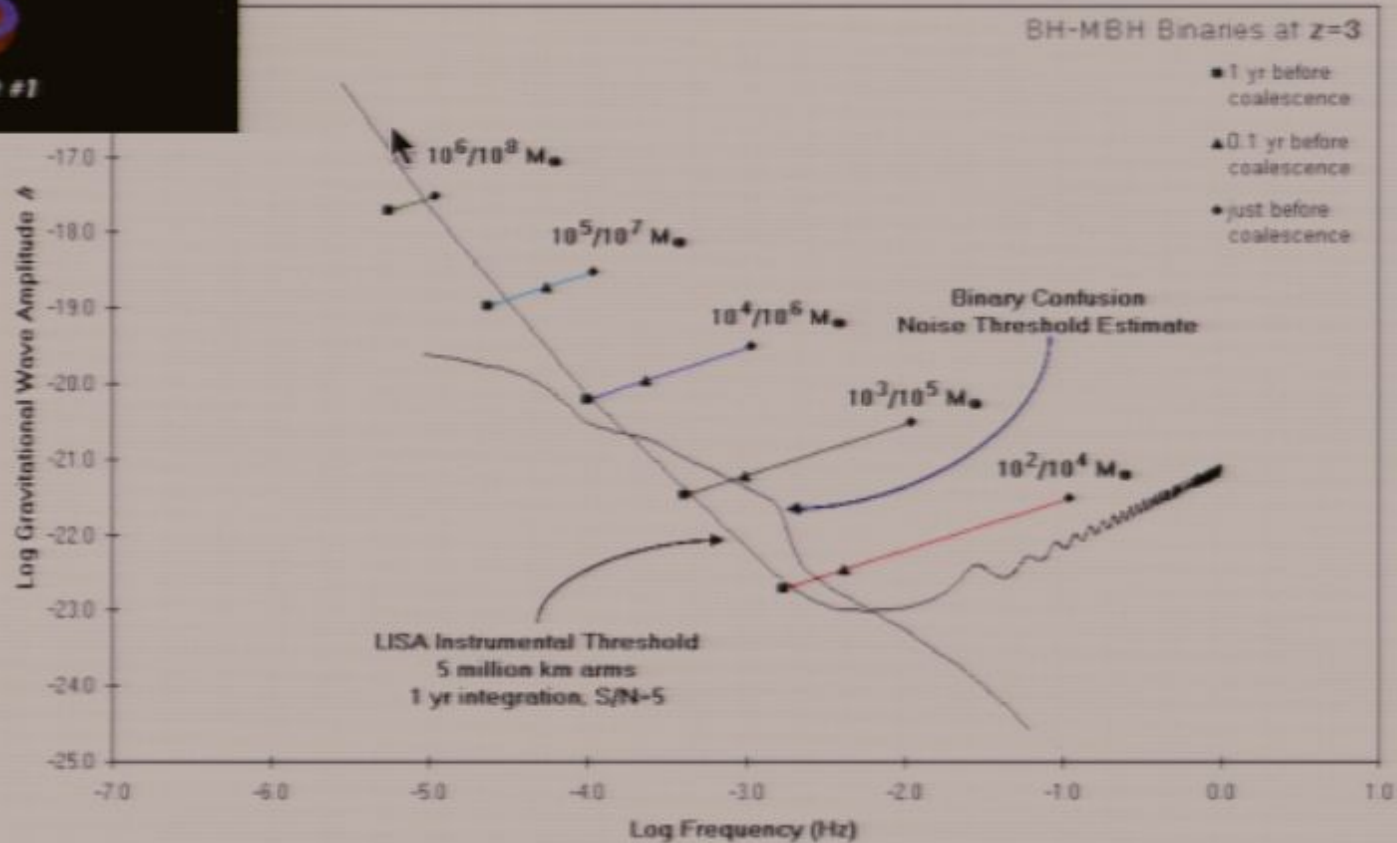








## Strain Amplitudes During Last Year Before BH-BH Coalescence



~~2004?~~

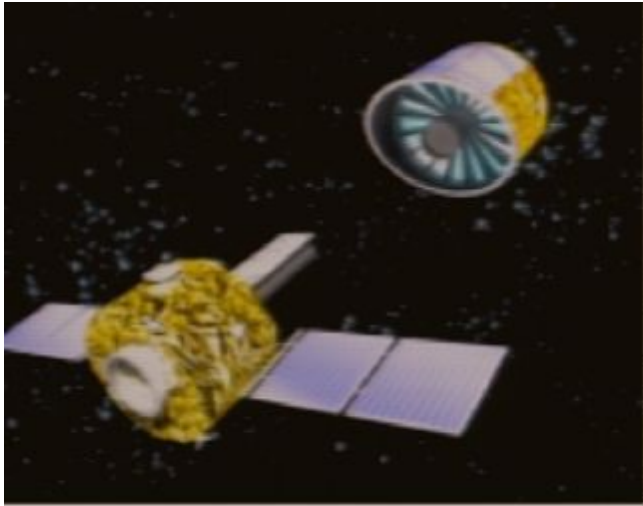
~~2011?~~

~~2017?~~

2020?

LISA will see mergers of  $10^5 - 10^7 M_\odot$   
binary black holes with high S/N

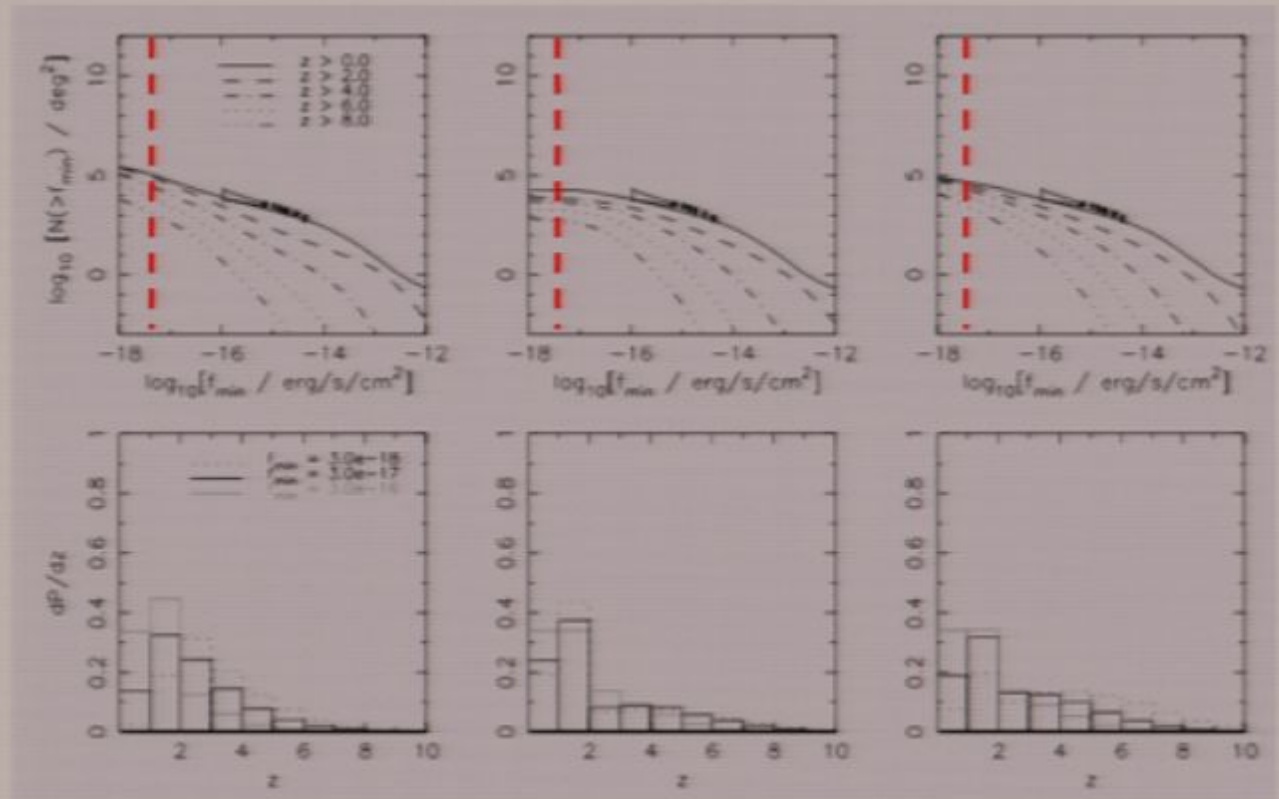


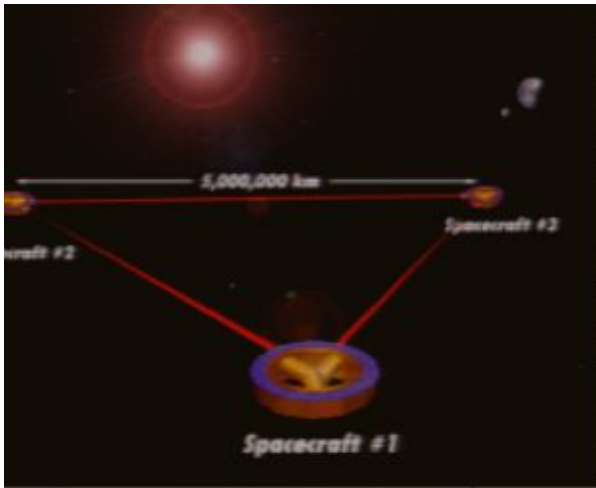


## Detecting quasars at very high redshift with next generation X-ray telescopes

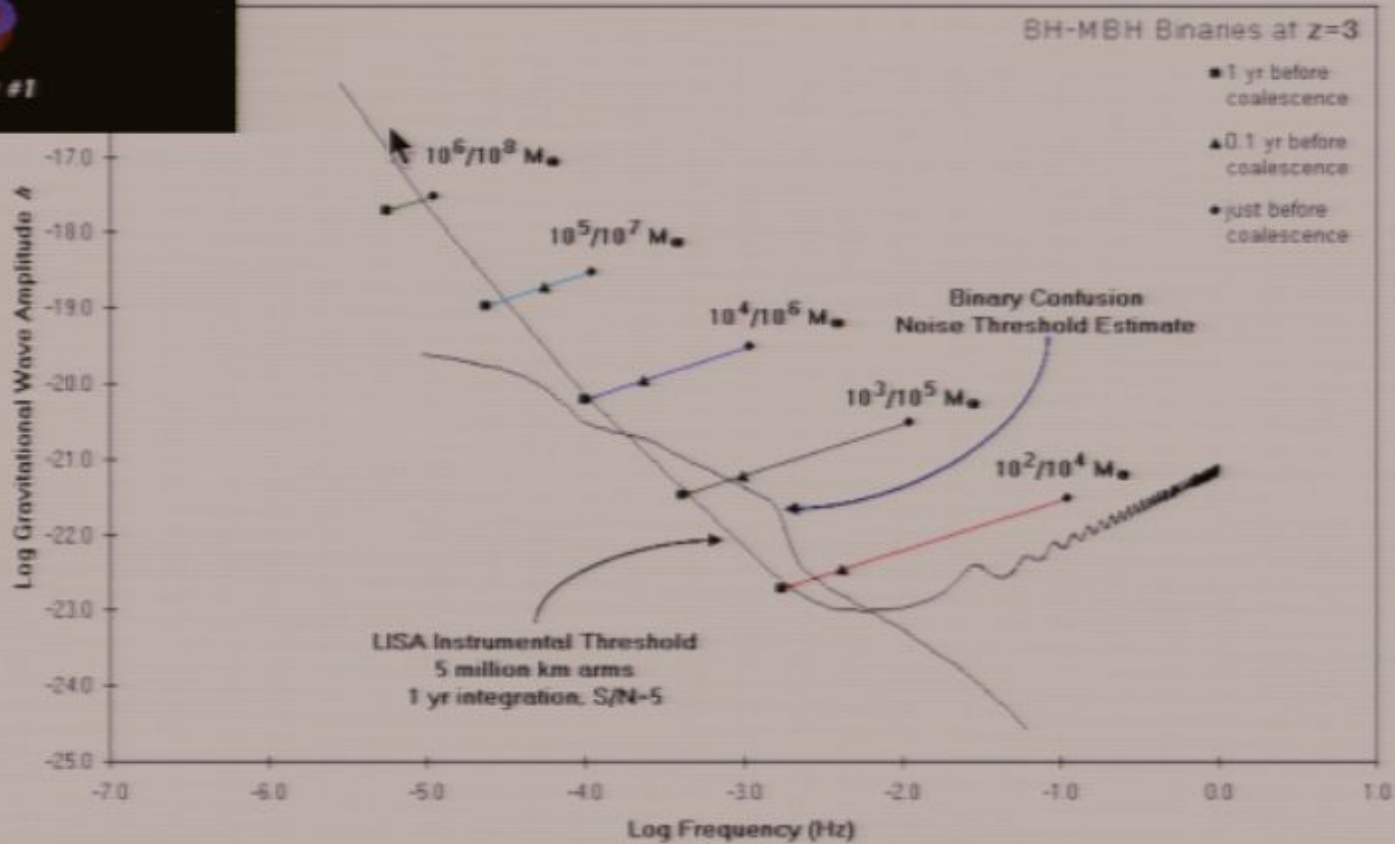
Kirsty J. Rhook\* & Martin G. Haehnelt†  
*Institute of Astronomy, Madingley Road, Cambridge CB3 0HA*

A significant fraction of the X-ray sources detected by XEUS should be at  $z > 6$ . Should be able to see black holes with masses as small as  $10^5 M_{\odot}$ .





### Strain Amplitudes During Last Year Before BH-BH Coalescence



~~2004?~~

~~2011?~~

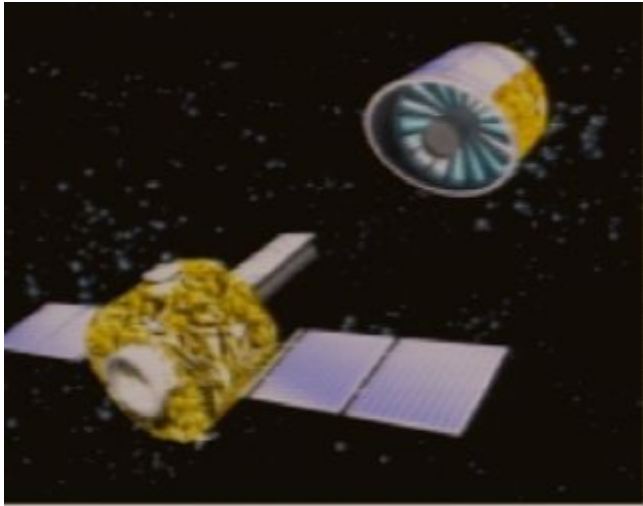
~~2017?~~

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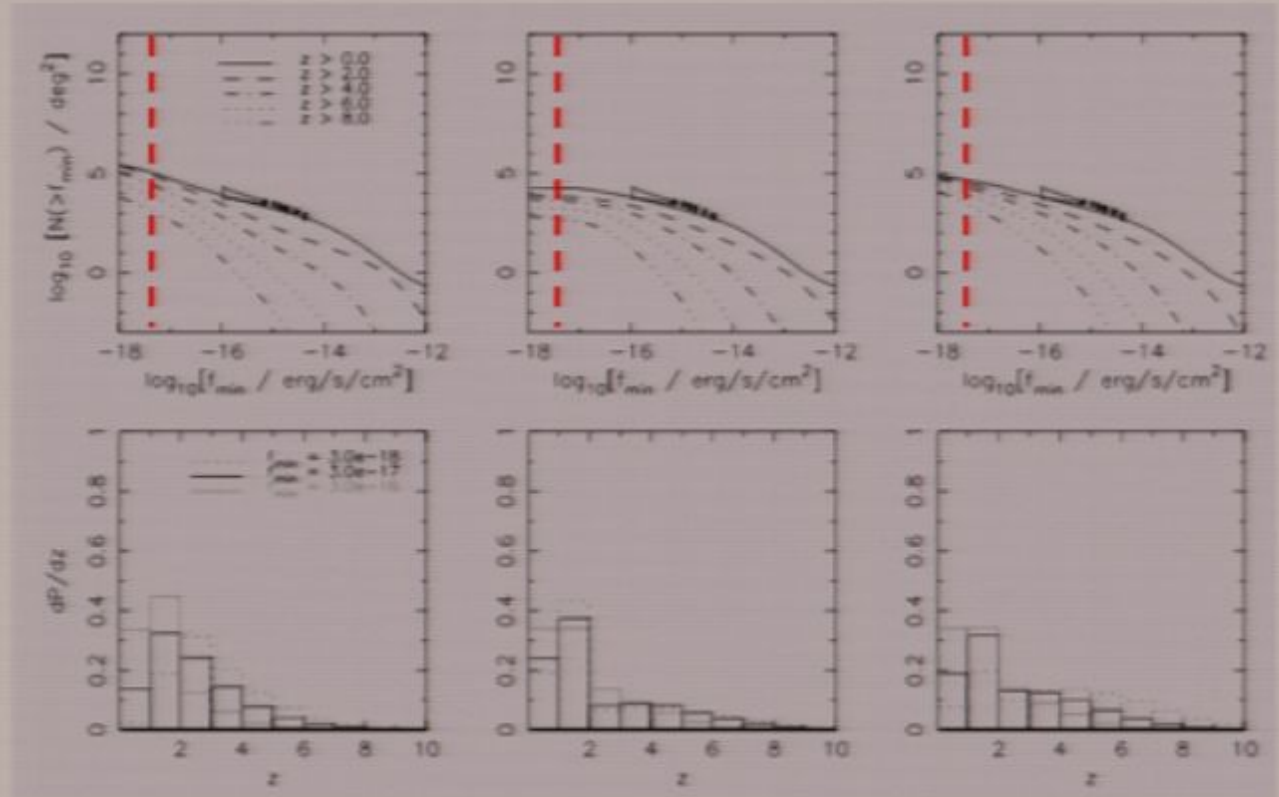




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# Summary

- Substantial progress on the observation and modelling of feedback regulated co-evolution of galaxies and their central black holes has been made which is now reasonably well understood.
- We still don't know how and when supermassive black holes form in the first place!
- Most probably require massive seed black holes. Direct collapse of gas in haloes with  $T_{\text{vir}} \geq 10000\text{K}$  with no metals (and  $\text{H}_2$  suppression) is least prone to fragmentation into stars. Further evolution via a massive Quasi-Star?
- Primordial massive black holes probably not required.
- LISA and future X-ray missions offer excellent prospects for unravelling the early build-up of supermassive black holes.



# Summary

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- Most probably require massive seed black holes. Direct collapse of gas in haloes with  $T_{\text{vir}} \geq 10000\text{K}$  with no metals (and  $\text{H}_2$  suppression) is least prone to fragmentation into stars. Further evolution via a massive Quasi-Star?
- Primordial massive black holes probably not required.
- LISA and future X-ray missions offer excellent prospects for unravelling the early build-up of supermassive black holes.

