

Title: A Mechanism for Asymptotic Safety of Chiral Yukawa Systems

Date: Nov 07, 2009 05:00 PM

URL: <http://pirsa.org/09110126>

Abstract:

A mechanism for asymptotic safety of chiral Yukawa systems

Michael Scherer

Collaboration with Holger Gies and Stefan Rechenberger

ITP, Jena University



Presented at Perimeter Institute, Waterloo, Canada, November 7, 2009

Outline

- 1 *Triviality and the hierarchy problem*
- 2 *Asymptotic safety and the flow equation*
- 3 *Truncation for chiral Yukawa systems*
- 4 *Fixed-points of chiral Yukawa systems*
- 5 *Numerical example for Higgs/top mass prediction*

An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).

An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).
- Recently, the AS scenario for QG has received strong support thanks to new techniques in non-perturbative quantum field theory.

An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).
- Recently, the AS scenario for QG has received strong support thanks to new techniques in non-perturbative quantum field theory.
- However, the setting of the AS scenario is more general and might also be applied to other QFTs that have problems with non-renormalizability.



An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).
- Recently, the AS scenario for QG has received strong support thanks to new techniques in non-perturbative quantum field theory.
- However, the setting of the AS scenario is more general and might also be applied to other QFTs that have problems with non-renormalizability.
- In the standard model of particle physics the Higgs sector is plagued by two problems:

An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).
- Recently, the AS scenario for QG has received strong support thanks to new techniques in non-perturbative quantum field theory.
- However, the setting of the AS scenario is more general and might also be applied to other QFTs that have problems with non-renormalizability.
- In the standard model of particle physics the Higgs sector is plagued by two problems:

triviality & hierarchy problem

An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).
- Recently, the AS scenario for QG has received strong support thanks to new techniques in non-perturbative quantum field theory.
- However, the setting of the AS scenario is more general and might also be applied to other QFTs that have problems with non-renormalizability.
- In the standard model of particle physics the Higgs sector is plagued by two problems:

triviality & hierarchy problem

- Both problems could be solved within the AS scenario.

An asymptotic safety scenario for the standard model

- The asymptotic safety (AS) scenario is mainly discussed in the context of a quantum theory for gravity (QG).
- Recently, the AS scenario for QG has received strong support thanks to new techniques in non-perturbative quantum field theory.
- However, the setting of the AS scenario is more general and might also be applied to other QFTs that have problems with non-renormalizability.
- In the standard model of particle physics the Higgs sector is plagued by two problems:

triviality & hierarchy problem

- Both problems could be solved within the AS scenario.
- As a toy model for the SM we will investigate a class of chiral Yukawa systems

Triviality

- Higgs field is parametrized in terms of a bosonic field ϕ with a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 + \frac{m^2}{2}\phi^2 + \frac{\lambda}{8}\phi^4.$$

Triviality

- Higgs field is parametrized in terms of a bosonic field ϕ with a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 + \frac{m^2}{2}\phi^2 + \frac{\lambda}{8}\phi^4.$$

- Perturbative computation of the one-loop correction to the four-Higgs-boson coupling yields relation between the bare and the renormalized coupling λ

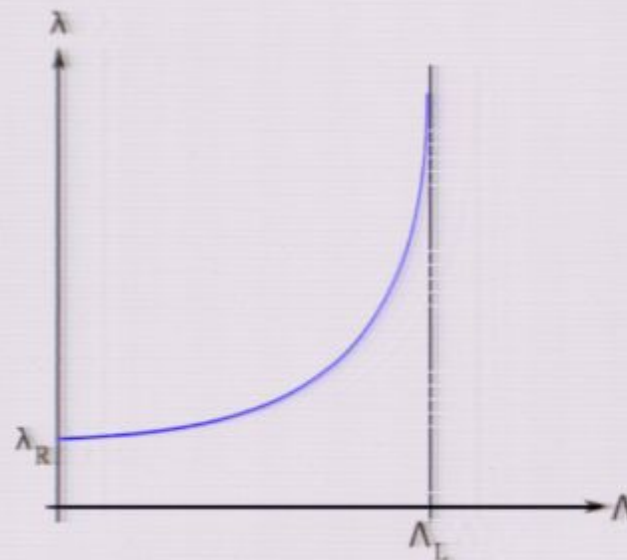


Triviality

- Higgs field is parametrized in terms of a bosonic field ϕ with a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 + \frac{m^2}{2}\phi^2 + \frac{\lambda}{8}\phi^4.$$

- Perturbative computation of the one-loop correction to the four-Higgs-boson coupling yields relation between the bare and the renormalized coupling λ



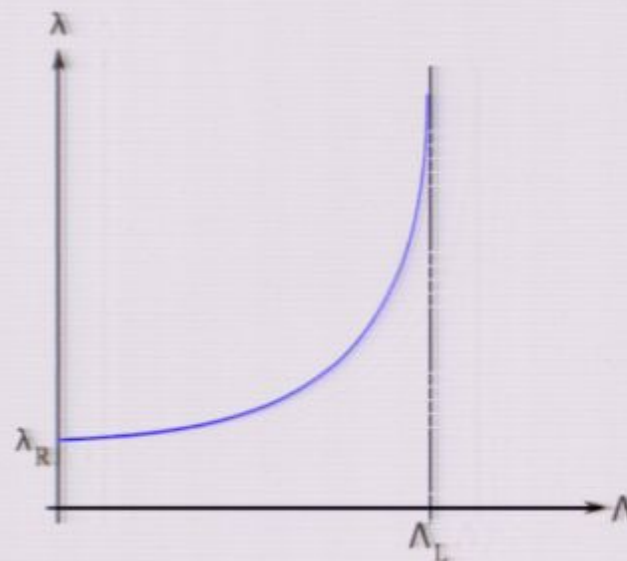
- Landau-pole indicates breakdown of perturbative QFT \rightarrow new d.o.f.?

Triviality

- Higgs field is parametrized in terms of a bosonic field ϕ with a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 + \frac{m^2}{2}\phi^2 + \frac{\lambda}{8}\phi^4.$$

- Perturbative computation of the one-loop correction to the four-Higgs-boson coupling yields relation between the bare and the renormalized coupling λ



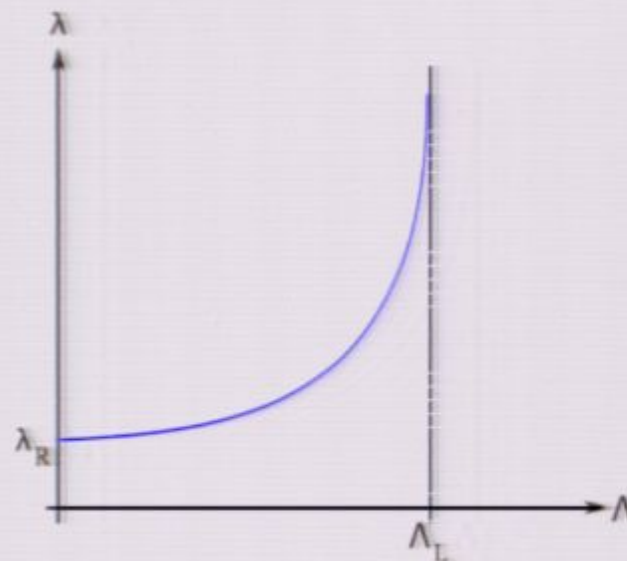
- Landau-pole indicates breakdown of perturbative QFT \rightarrow new d.o.f.?
- Perturbation theory relies on an expansion around zero coupling.

Triviality

- Higgs field is parametrized in terms of a bosonic field ϕ with a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 + \frac{m^2}{2}\phi^2 + \frac{\lambda}{8}\phi^4.$$

- Perturbative computation of the one-loop correction to the four-Higgs-boson coupling yields relation between the bare and the renormalized coupling λ



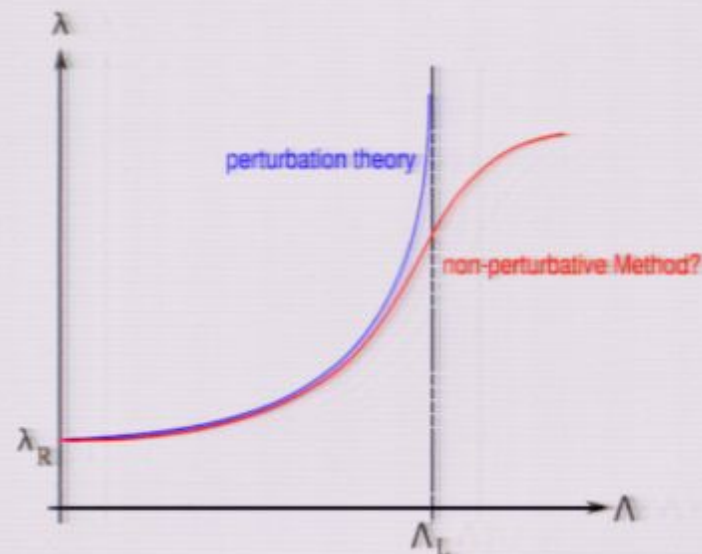
- Landau-pole indicates breakdown of perturbative QFT \rightarrow new d.o.f.?
- Perturbation theory relies on an expansion around zero coupling.
- Near the Landau pole perturbation theory will lose its validity since λ grows large

Triviality

- Higgs field is parametrized in terms of a bosonic field ϕ with a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 + \frac{m^2}{2}\phi^2 + \frac{\lambda}{8}\phi^4.$$

- Perturbative computation of the one-loop correction to the four-Higgs-boson coupling yields relation between the bare and the renormalized coupling λ



- Landau-pole indicates breakdown of perturbative QFT \rightarrow new d.o.f.?
- Perturbation theory relies on an expansion around zero coupling.
- Near the Landau pole perturbation theory will lose its validity since λ grows large
- Need non-perturbative tool to study triviality & take into account fermions!

Hierarchy problem

We observe a huge hierarchy in the standard model:

$$\Lambda_{\text{EW}} \sim 10^2 \text{ GeV} \ll \Lambda_{\text{GUT}} \sim 10^{16} \text{ GeV}.$$

Hierarchy problem

We observe a huge hierarchy in the standard model:

$$\Lambda_{EW} \sim 10^2 \text{ GeV} \ll \Lambda_{GUT} \sim 10^{16} \text{ GeV}.$$

The Higgs mass renormalizes quadratically ($\delta m^2 \sim \Lambda^2$). In perturbation theory the relation between bare and renormalized coupling is given by

$$m_R^2 \sim m_{\Lambda,UV}^2 - \delta m^2$$



Hierarchy problem

We observe a huge hierarchy in the standard model:

$$\Lambda_{EW} \sim 10^2 \text{ GeV} \ll \Lambda_{GUT} \sim 10^{16} \text{ GeV}.$$

The Higgs mass renormalizes quadratically ($\delta m^2 \sim \Lambda^2$). In perturbation theory the relation between bare and renormalized coupling is given by

$$\underbrace{m_R^2}_{\sim 10^4 \text{ GeV}^2} \sim \underbrace{m_{\Lambda,UV}^2}_{\sim 10^{32} (X + \dots 10^{-28}) \text{ GeV}^2} - \delta m^2$$

with a counterterm $\delta m^2 = X \cdot 10^{32} \text{ GeV}^2$.

Hierarchy problem

We observe a huge hierarchy in the standard model:

$$\Lambda_{EW} \sim 10^2 \text{ GeV} \ll \Lambda_{GUT} \sim 10^{16} \text{ GeV}.$$

The Higgs mass renormalizes quadratically ($\delta m^2 \sim \Lambda^2$). In perturbation theory the relation between bare and renormalized coupling is given by

$$\underbrace{m_R^2}_{\sim 10^4 \text{ GeV}^2} \sim \underbrace{m_{\Lambda,UV}^2}_{\sim 10^{32} (X + \dots 10^{-28}) \text{ GeV}^2} - \delta m^2$$

with a counterterm $\delta m^2 = X \cdot 10^{32} \text{ GeV}^2$.

- Perform a fine-tuning with a precision of $\Lambda_{EW}^2 / \Lambda_{GUT}^2 \sim 10^{-28}$.

Hierarchy problem

We observe a huge hierarchy in the standard model:

$$\Lambda_{EW} \sim 10^2 \text{ GeV} \ll \Lambda_{GUT} \sim 10^{16} \text{ GeV}.$$

The Higgs mass renormalizes quadratically ($\delta m^2 \sim \Lambda^2$). In perturbation theory the relation between bare and renormalized coupling is given by

$$\underbrace{m_R^2}_{\sim 10^4 \text{ GeV}^2} \sim \underbrace{m_{\Lambda,UV}^2}_{\sim 10^{32} (X + \dots 10^{-28}) \text{ GeV}^2} - \delta m^2$$

with a counterterm $\delta m^2 = X \cdot 10^{32} \text{ GeV}^2$.

- Perform a fine-tuning with a precision of $\Lambda_{EW}^2 / \Lambda_{GUT}^2 \sim 10^{-28}$.
- This seems to be “unnatural”.

A hierarchy problem corresponds to the existence of a large critical exponent $\Theta_I > 0$ at a fixed point, e.g. in ϕ^4 -theory we find at the GFP $\Theta = 2$.

Flow equation & asymptotic safety

Exact renormalization group equations (ERGE) derived from path-integral representation (Wetterich '93)

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \{ [\Gamma_k^{(2)}[\Phi] + R_k]^{-1} (\partial_t R_k) \}, \quad \partial_t = k \frac{d}{dk}$$

Flow equation & asymptotic safety

Exact renormalization group equations (ERGE) derived from path-integral representation (Wetterich '93)

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \{ [\Gamma_k^{(2)}[\Phi] + R_k]^{-1} (\partial_t R_k) \}, \quad \partial_t = k \frac{d}{dk}$$

Plugging in an effective average action $\Gamma_k[\Phi] = \sum_i g_{i,k} \mathcal{O}_i$, we obtain β -functions

$$\partial_t g_{i,k} = \beta_{i,k}(g_{1,k}, g_{2,k}, \dots)$$



Flow equation & asymptotic safety

Exact renormalization group equations (ERGE) derived from path-integral representation (Wetterich '93)

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \{ [\Gamma_k^{(2)}[\Phi] + R_k]^{-1} (\partial_t R_k) \}, \quad \partial_t = k \frac{d}{dk}$$

Plugging in an effective average action $\Gamma_k[\Phi] = \sum_i g_{i,k} \mathcal{O}_i$, we obtain β -functions

$$\partial_t g_{i,k} = \beta_{i,k}(g_{1,k}, g_{2,k}, \dots)$$

At a (possibly non-Gaussian) fixed-point we linearize the β -functions

$$\partial_t g_{i,k} = B_i^j (g_{j,k} - g_j^*), \quad B_i^j = \left. \frac{\partial \beta_i}{\partial g_{j,k}} \right|_{g^*} \quad \otimes$$

Flow equation & asymptotic safety

Exact renormalization group equations (ERGE) derived from path-integral representation (Wetterich '93)

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \{ [\Gamma_k^{(2)}[\Phi] + R_k]^{-1} (\partial_t R_k) \}, \quad \partial_t = k \frac{d}{dk}$$

Plugging in an effective average action $\Gamma_k[\Phi] = \sum_i g_{i,k} \mathcal{O}_i$, we obtain β -functions

$$\partial_t g_{i,k} = \beta_{i,k}(g_{1,k}, g_{2,k}, \dots)$$

At a (possibly non-Gaussian) fixed-point we linearize the β -functions

$$\partial_t g_{i,k} = B_i^j (g_{j,k} - g_j^*), \quad B_i^j = \left. \frac{\partial \beta_i}{\partial g_{j,k}} \right|_{g^*}$$

With the eigenvectors V^I and eigenvalues Θ^I of the stability matrix we give a general solution of the linearized fixed-point equation

$$g_{i,k} = g_i^* + \sum_I C_I V_i^I \left(\frac{k_0}{k} \right)^{\Theta_I}$$

Flow equation & asymptotic safety

Exact renormalization group equations (ERGE) derived from path-integral representation (Wetterich '93)

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \{ [\Gamma_k^{(2)}[\Phi] + R_k]^{-1} (\partial_t R_k) \}, \quad \partial_t = k \frac{d}{dk}$$

Plugging in an effective average action $\Gamma_k[\Phi] = \sum_i g_{i,k} \mathcal{O}_i$, we obtain β -functions

$$\partial_t g_{i,k} = \beta_{i,k}(g_{1,k}, g_{2,k}, \dots)$$

At a (possibly non-Gaussian) fixed-point we linearize the β -functions

$$\partial_t g_{i,k} = B_i^j (g_{j,k} - g_j^*), \quad B_i^j = \left. \frac{\partial \beta_i}{\partial g_{j,k}} \right|_{g^*}$$

With the eigenvectors V^I and eigenvalues Θ^I of the stability matrix we give a general solution of the linearized fixed-point equation

$$g_{i,k} = g_i^* + \sum_I C_I V_i^I \left(\frac{k_0}{k} \right)^{\Theta_I}$$

- $\text{Re } \Theta_I > 0$: relevant coupling (to be fixed by experiment)

Flow equation & asymptotic safety

Exact renormalization group equations (ERGE) derived from path-integral representation (Wetterich '93)

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \{ [\Gamma_k^{(2)}[\Phi] + R_k]^{-1} (\partial_t R_k) \}, \quad \partial_t = k \frac{d}{dk}$$

Plugging in an effective average action $\Gamma_k[\Phi] = \sum_i g_{i,k} \mathcal{O}_i$, we obtain β -functions

$$\partial_t g_{i,k} = \beta_{i,k}(g_{1,k}, g_{2,k}, \dots)$$

At a (possibly non-Gaussian) fixed-point we linearize the β -functions

$$\partial_t g_{i,k} = B_i^j (g_{j,k} - g_j^*), \quad B_i^j = \left. \frac{\partial \beta_i}{\partial g_{j,k}} \right|_{g^*}$$

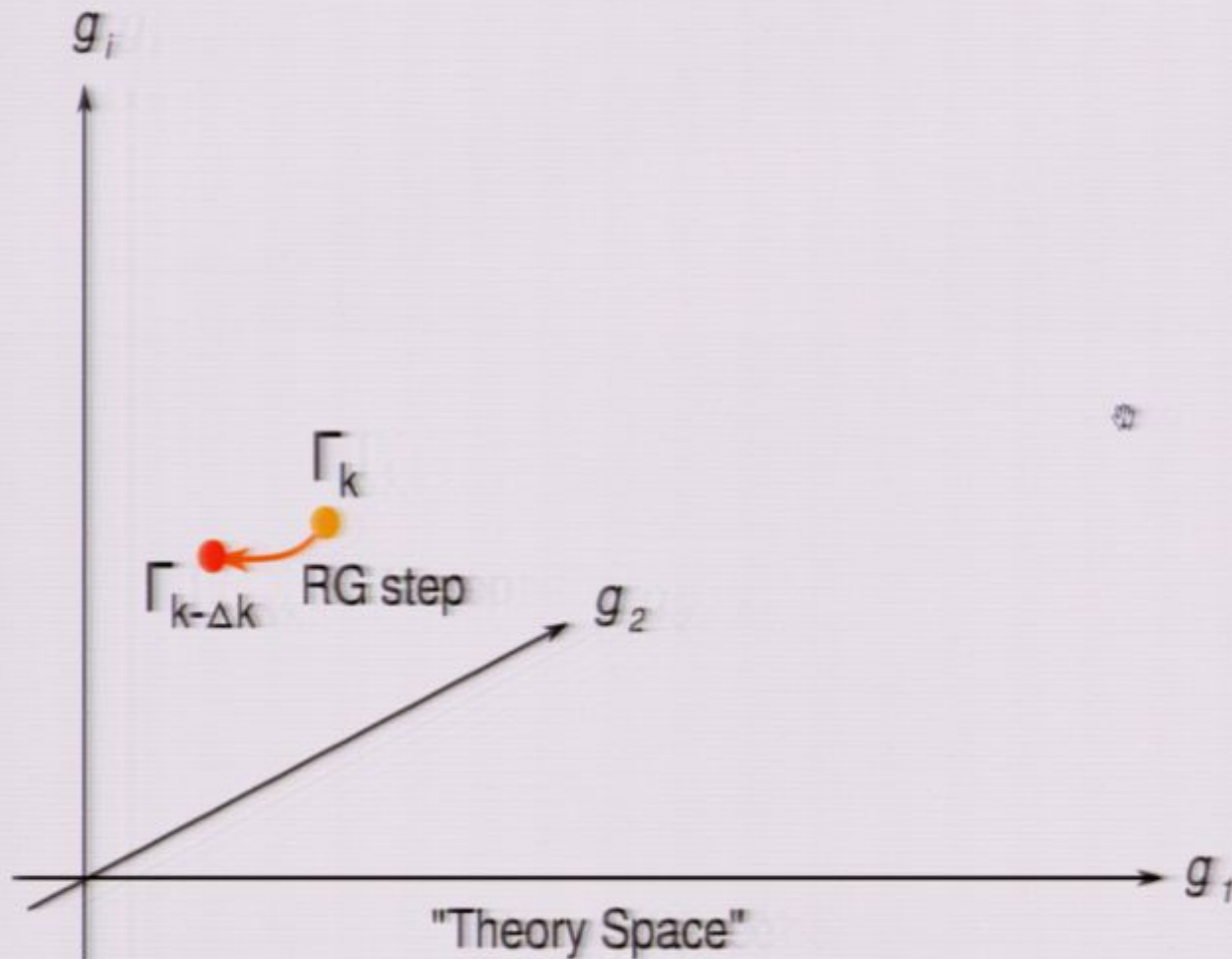
With the eigenvectors V^I and eigenvalues Θ^I of the stability matrix we give a general solution of the linearized fixed-point equation

$$g_{i,k} = g_i^* + \sum_I C_I V_i^I \left(\frac{k_0}{k} \right)^{\Theta_I}$$

- $\text{Re } \Theta_I > 0$: relevant coupling (to be fixed by experiment)
- $\text{Re } \Theta_I < 0$: irrelevant coupling (prediction for physical observable in the IR)

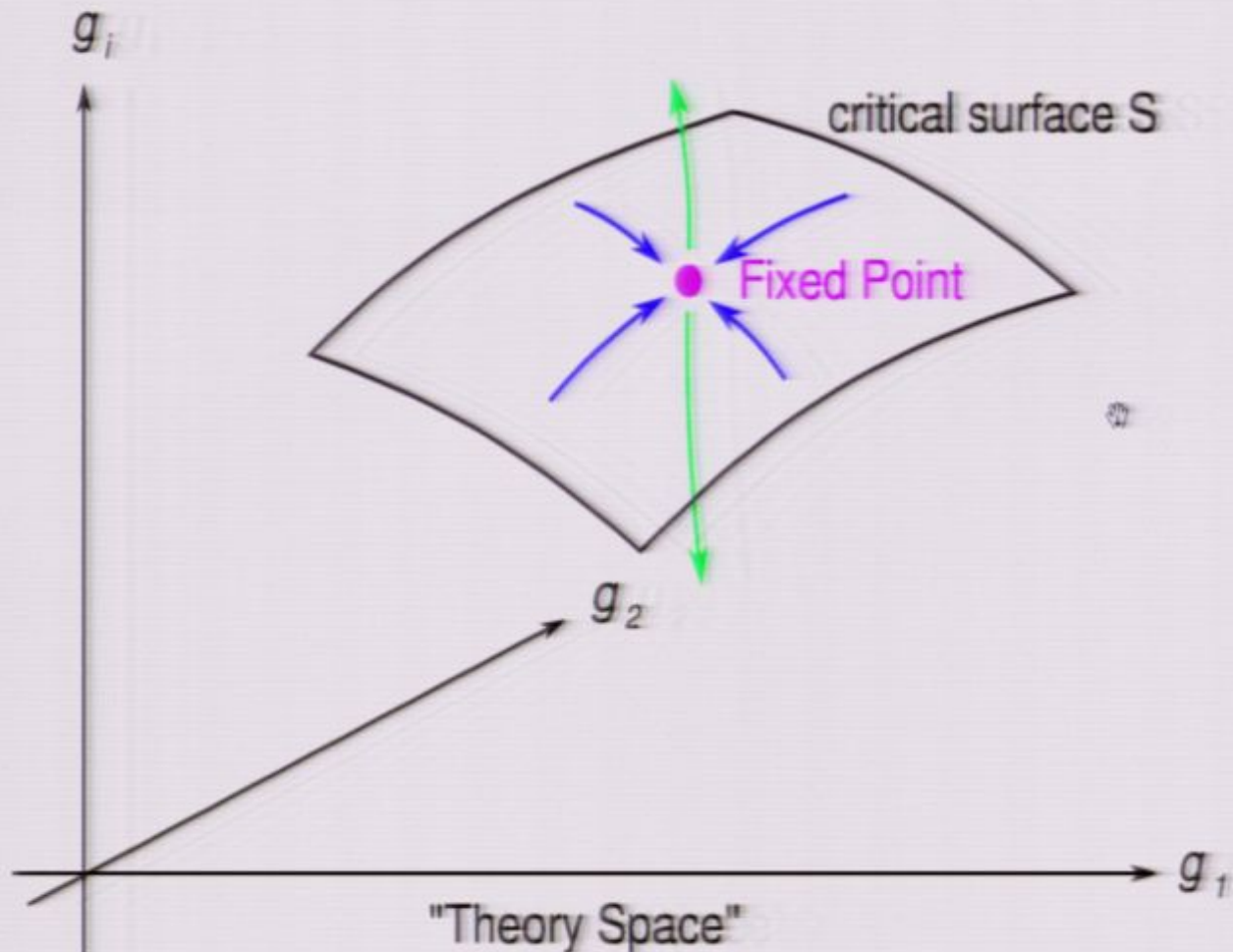
Asymptotic safety & theory space

Effective average action: $\Gamma_k[\chi] = \sum_i g_{i,k} \mathcal{O}_i$. Scale dependence: $\partial_t \Gamma_k[\chi] = \sum_i \beta_{i,k} \mathcal{O}_i$.



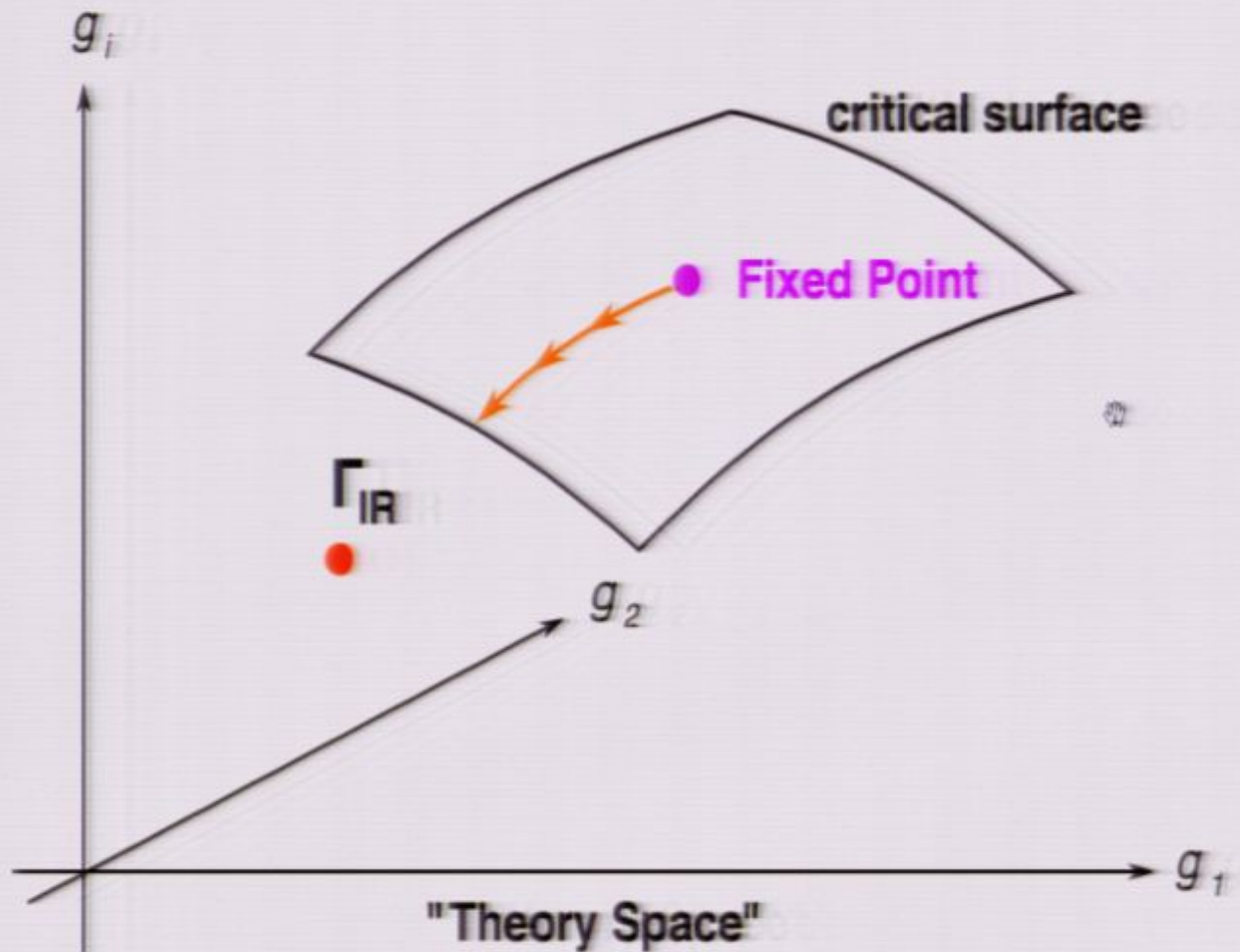
Asymptotic safety & theory space

Effective average action: $\Gamma_k[\chi] = \sum_i g_{i,k} \mathcal{O}_i$, Scale dependence: $\partial_t \Gamma_k[\chi] = \sum_i \beta_{i,k} \mathcal{O}_i$.



Asymptotic safety & theory space

Effective average action: $\Gamma_k[\chi] = \sum_i g_{i,k} \mathcal{O}_i$, Scale dependence: $\partial_t \Gamma_k[\chi] = \sum_i \beta_{i,k} \mathcal{O}_i$.

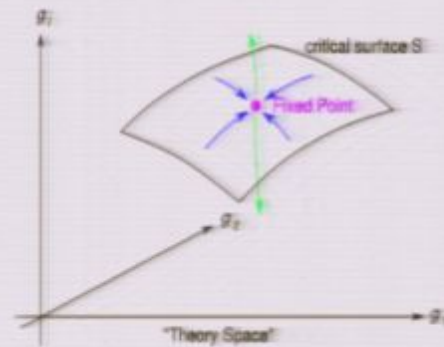


Triviality & hierarchy problem in the asymptotic safety scenario



- Dimension of the critical surface: $\Delta = \dim S = \text{number of relevant directions}$.

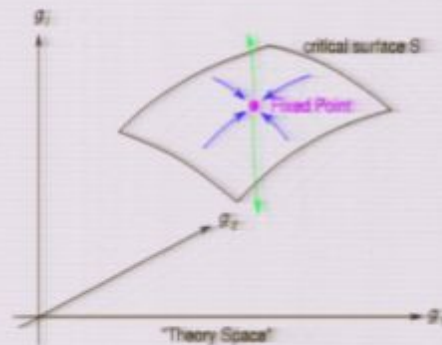
Triviality & hierarchy problem in the asymptotic safety scenario



- Dimension of the critical surface: $\Delta = \dim S = \text{number of relevant directions}$.
- If a non-Gaussian fixed point (NGFP) exists we can draw the limit $\Lambda \rightarrow \infty$ and the system is independent from the UV cutoff.

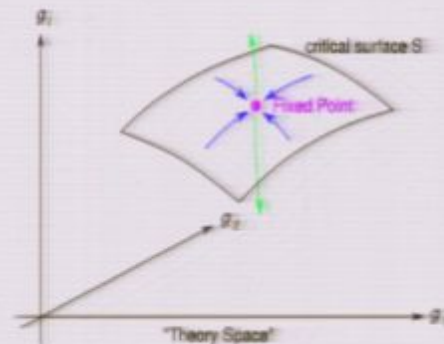


Triviality & hierarchy problem in the asymptotic safety scenario



- Dimension of the critical surface: $\Delta = \dim S = \text{number of relevant directions}$.
- If a non-Gaussian fixed point (NGFP) exists we can draw the limit $\Lambda \rightarrow \infty$ and the system is independent from the UV cutoff.
- NGFP solves the triviality problem, because the UV limit is well-defined.

Triviality & hierarchy problem in the asymptotic safety scenario



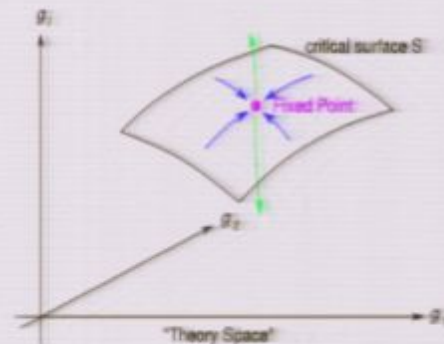
- Dimension of the critical surface: $\Delta = \dim S =$ number of relevant directions.
- If a non-Gaussian fixed point (NGFP) exists we can draw the limit $\Lambda \rightarrow \infty$ and the system is independent from the UV cutoff.
- NGFP solves the triviality problem, because the UV limit is well-defined.
- If $\Delta < \infty \rightarrow$ system is predictive, because there is only a finite number of parameters to be fixed (by experiment).

Triviality & hierarchy problem in the asymptotic safety scenario



- Dimension of the critical surface: $\Delta = \dim S = \text{number of relevant directions}$.
- If a non-Gaussian fixed point (NGFP) exists we can draw the limit $\Lambda \rightarrow \infty$ and the system is independent from the UV cutoff.
- NGFP solves the triviality problem, because the UV limit is well-defined.
- If $\Delta < \infty \rightarrow$ system is predictive, because there is only a finite number of parameters to be fixed (by experiment).
- We find a hierarchy problem if there exist large critical exponents $\Theta_I > 0$.

Triviality & hierarchy problem in the asymptotic safety scenario



- Dimension of the critical surface: $\Delta = \dim S = \text{number of relevant directions}$.
- If a non-Gaussian fixed point (NGFP) exists we can draw the limit $\Lambda \rightarrow \infty$ and the system is independent from the UV cutoff.
- NGFP solves the triviality problem, because the UV limit is well-defined.
- If $\Delta < \infty \rightarrow$ system is predictive, because there is only a finite number of parameters to be fixed (by experiment).
- We find a hierarchy problem if there exist large critical exponents $\Theta_I > 0$.
- RG computation will show how large the Θ_I are at a NGFP.

Triviality & hierarchy problem in the asymptotic safety scenario



- Dimension of the critical surface: $\Delta = \dim S = \text{number of relevant directions}$.
- If a non-Gaussian fixed point (NGFP) exists we can draw the limit $\Lambda \rightarrow \infty$ and the system is independent from the UV cutoff.
- NGFP solves the triviality problem, because the UV limit is well-defined.
- If $\Delta < \infty \rightarrow$ system is predictive, because there is only a finite number of parameters to be fixed (by experiment).
- We find a hierarchy problem if there exist large critical exponents $\Theta_I > 0$.
- RG computation will show how large the Θ_I are at a NGFP.
- If all of them are small $\ll 1$ then the hierarchy problem is solved.

Toy model - Chiral Yukawa system without gauge bosons

Derivative expansion, leading-order truncation

$$\Gamma_k = \int d^d x \left\{ i(\bar{\psi}_L^a \not{\partial} \psi_L^a + \bar{\psi}_R \not{\partial} \psi_R) + (\partial_\mu \phi^{a\dagger})(\partial^\mu \phi^a) \right. \\ \left. + U_k(\phi^{a\dagger} \phi^a) + \bar{h}_k \bar{\psi}_R \phi^{a\dagger} \psi_L^a - \bar{h}_k \bar{\psi}_L^a \phi^a \psi_R \right\}$$

Toy model - Chiral Yukawa system without gauge bosons

Derivative expansion, leading-order truncation

$$\Gamma_k = \int d^d x \left\{ i(\bar{\psi}_L^a \not{\partial} \psi_L^a + \bar{\psi}_R \not{\partial} \psi_R) + (\partial_\mu \phi^{a\dagger})(\partial^\mu \phi^a) \right. \\ \left. + U_k(\phi^{a\dagger} \phi^a) + \bar{h}_k \bar{\psi}_R \phi^{a\dagger} \psi_L^a - \bar{h}_k \bar{\psi}_L^a \phi^a \psi_R \right\}$$

- N_L left-handed fermions ψ_L^a
- one right-handed fermion ψ_R

Toy model - Chiral Yukawa system without gauge bosons

Derivative expansion, leading-order truncation

$$\Gamma_k = \int d^d x \left\{ i(\bar{\psi}_L^a \not{\partial} \psi_L^a + \bar{\psi}_R \not{\partial} \psi_R) + (\partial_\mu \phi^{a\dagger})(\partial^\mu \phi^a) \right. \\ \left. + U_k(\phi^{a\dagger} \phi^a) + \bar{h}_k \bar{\psi}_R \phi^{a\dagger} \psi_L^a - \bar{h}_k \bar{\psi}_L^a \phi^a \psi_R \right\}$$

- N_L left-handed fermions ψ_L^a
- one right-handed fermion ψ_R
- N_L complex bosons ϕ^a
- invariant under chiral $U(N_L)_L \otimes U(1)_R$ transformations



Toy model - Chiral Yukawa system without gauge bosons

Derivative expansion, leading-order truncation

$$\Gamma_k = \int d^d x \left\{ i(\bar{\psi}_L^a \not{\partial} \psi_L^a + \bar{\psi}_R \not{\partial} \psi_R) + (\partial_\mu \phi^{a\dagger})(\partial^\mu \phi^a) \right. \\ \left. + U_k(\phi^{a\dagger} \phi^a) + \bar{h}_k \bar{\psi}_R \phi^{a\dagger} \psi_L^a - \bar{h}_k \bar{\psi}_L^a \phi^a \psi_R \right\}$$

- N_L left-handed fermions ψ_L^a
- one right-handed fermion ψ_R
- N_L complex bosons ϕ^a
- invariant under chiral $U(N_L)_L \otimes U(1)_R$ transformations
- define $\rho = \phi^{a\dagger} \phi^a$.
- dimensionless quantities:

$$\tilde{\rho} = k^{2-d} \rho, \quad \tilde{h}^2 = k^{d-4} h_k^2, \quad u(\tilde{\rho}) = k^{-d} U_k(\rho)|_{\rho=k^{d-2}\tilde{\rho}}$$

Toy model - Chiral Yukawa system without gauge bosons

Derivative expansion, leading-order truncation

$$\Gamma_k = \int d^d x \left\{ i(\bar{\psi}_L^a \not{\partial} \psi_L^a + \bar{\psi}_R \not{\partial} \psi_R) + (\partial_\mu \phi^{a\dagger})(\partial^\mu \phi^a) \right. \\ \left. + U_k(\phi^{a\dagger} \phi^a) + \bar{h}_k \bar{\psi}_R \phi^{a\dagger} \psi_L^a - \bar{h}_k \bar{\psi}_L^a \phi^a \psi_R \right\}$$

- N_L left-handed fermions ψ_L^a
- one right-handed fermion ψ_R
- N_L complex bosons ϕ^a
- invariant under chiral $U(N_L)_L \otimes U(1)_R$ transformations
- define $\rho = \phi^{a\dagger} \phi^a$.
- dimensionless quantities:

$$\tilde{\rho} = k^{2-d} \rho, \quad \tilde{h}^2 = k^{d-4} h_k^2, \quad u(\tilde{\rho}) = k^{-d} U_k(\rho)|_{\rho=k^{d-2}\tilde{\rho}}$$

For the regime with spontaneously broken symmetry (SSB), we expand the effective potential about its minimum:

$$\kappa := \tilde{\rho}_{\min} > 0,$$

$$u = \frac{\lambda_2}{2!} (\tilde{\rho} - \kappa)^2 + \frac{\lambda_3}{3!} (\tilde{\rho} - \kappa)^3 + \dots$$

$\kappa, \lambda_{n_{\max}}, \lambda_2 > 0.$

Toy model - Chiral Yukawa system without gauge bosons

Derivative expansion, leading-order truncation

$$\Gamma_k = \int d^d x \left\{ i(\bar{\psi}_L^a \not{\partial} \psi_L^a + \bar{\psi}_R \not{\partial} \psi_R) + (\partial_\mu \phi^{a\dagger})(\partial^\mu \phi^a) \right. \\ \left. + U_k(\phi^{a\dagger} \phi^a) + \bar{h}_k \bar{\psi}_R \phi^{a\dagger} \psi_L^a - \bar{h}_k \bar{\psi}_L^a \phi^a \psi_R \right\}$$

- N_L left-handed fermions ψ_L^a
- one right-handed fermion ψ_R
- N_L complex bosons ϕ^a
- invariant under chiral $U(N_L)_L \otimes U(1)_R$ transformations
- define $\rho = \phi^{a\dagger} \phi^a$.
- dimensionless quantities:

$$\tilde{\rho} = k^{2-d} \rho, \quad \tilde{h}^2 = k^{d-4} h_k^2, \quad u(\tilde{\rho}) = k^{-d} U_k(\rho)|_{\rho=k^{d-2}\tilde{\rho}}$$

For the regime with spontaneously broken symmetry (SSB), we expand the effective potential about its minimum:

$$\kappa := \tilde{\rho}_{\min} > 0,$$

$$u = \frac{\lambda_2}{2!} (\tilde{\rho} - \kappa)^2 + \frac{\lambda_3}{3!} (\tilde{\rho} - \kappa)^3 + \dots$$

$$\kappa, \lambda_{n_{\max}}, \lambda_2 > 0.$$



Fixed-point mechanism

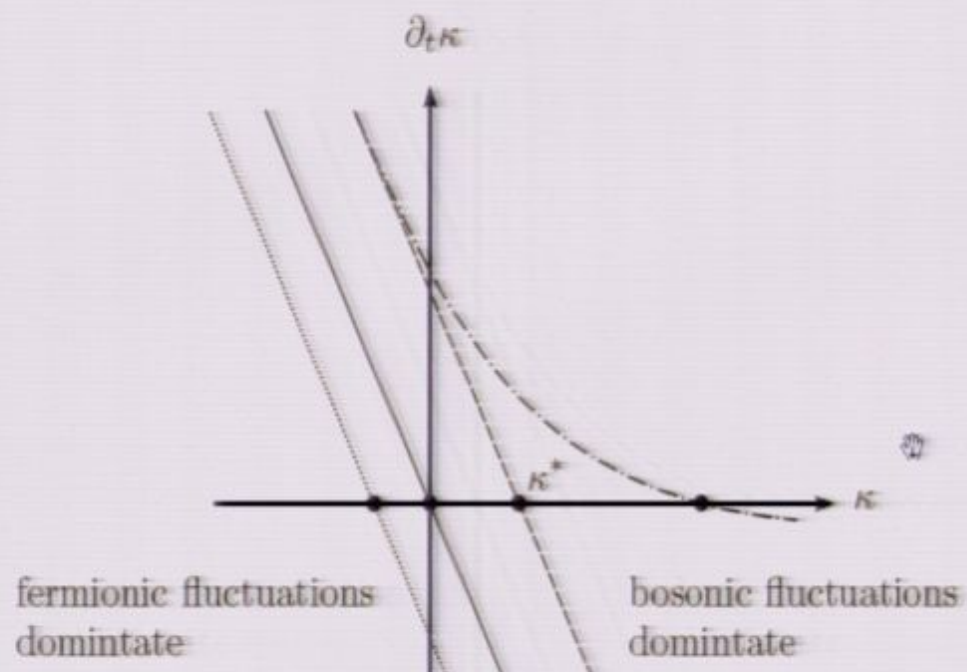
Loop contributions to the running of κ :

$$\partial_t \kappa = -2\kappa + \text{bosonic interactions} - \text{fermionic interactions.} \quad (1)$$

Fixed-point mechanism

Loop contributions to the running of κ :

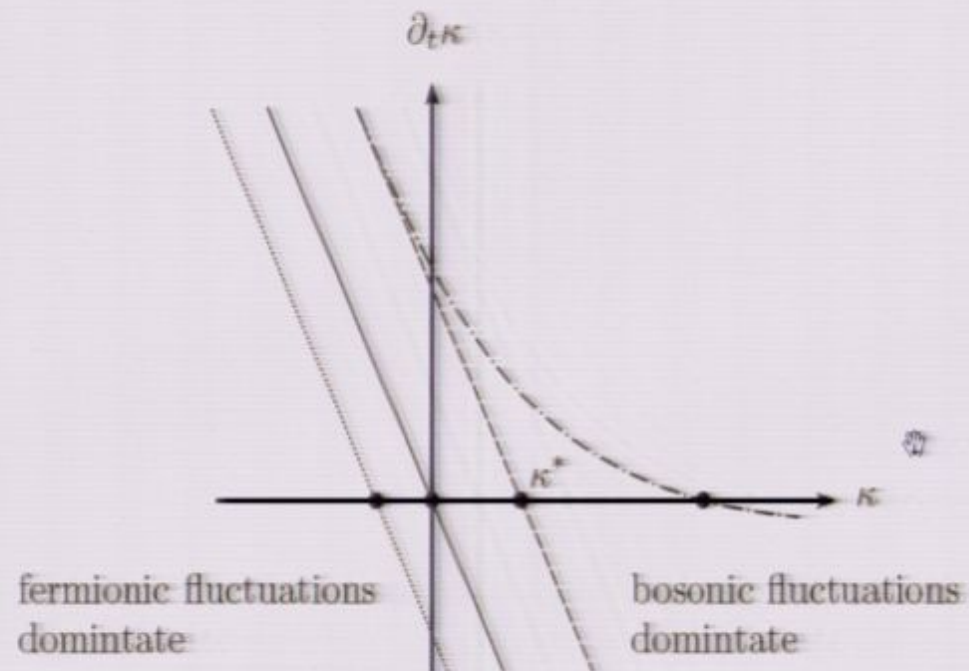
$$\partial_t \kappa = -2\kappa + \text{bosonic interactions} - \text{fermionic interactions.} \quad (1)$$



Fixed-point mechanism

Loop contributions to the running of κ :

$$\partial_t \kappa = -2\kappa + \text{bosonic interactions} - \text{fermionic interactions.} \quad (1)$$

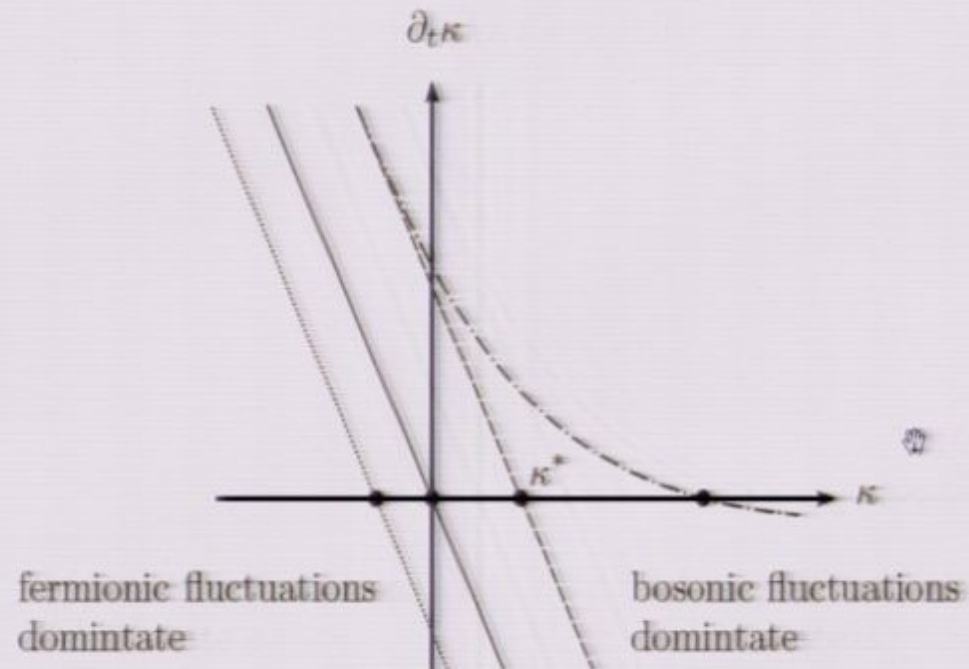


- dominating fluctuations of the boson field allow for a positive κ^*

Fixed-point mechanism

Loop contributions to the running of κ :

$$\partial_t \kappa = -2\kappa + \text{bosonic interactions} - \text{fermionic interactions.} \quad (1)$$



- dominating fluctuations of the boson field allow for a positive κ^*
- a suitable κ -dependence flattens the β -function near the fixed-point, which reduces the hierarchy problem
- near the FP the vev exhibits a conformal behaviour $v \sim k$ (cf. talk by H. Gies)

Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.

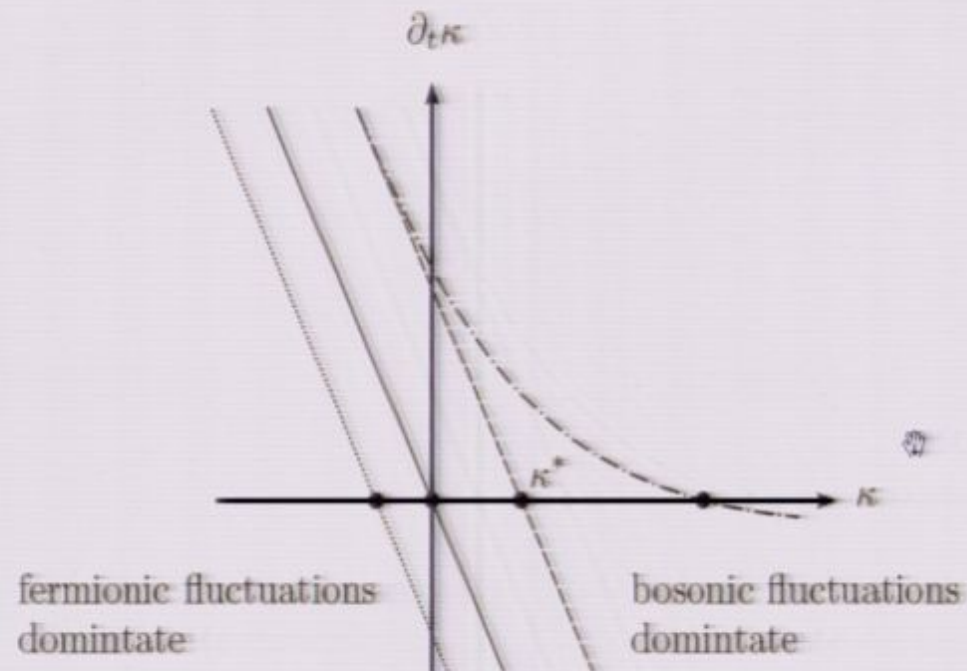
Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.

Fixed-point mechanism

Loop contributions to the running of κ :

$$\partial_t \kappa = -2\kappa + \text{bosonic interactions} - \text{fermionic interactions.} \quad (1)$$



- dominating fluctuations of the boson field allow for a positive κ^*
- a suitable κ -dependence flattens the β -function near the fixed-point, which reduces the hierarchy problem
- near the FP the vev exhibits a conformal behaviour $v \sim k$ (cf. talk by H. Gies)

Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.

Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.
- A leading order truncation can be parametrized by three couplings: h^2, λ, κ .

Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.
- A leading order truncation can be parametrized by three couplings: h^2, λ, κ .

$$\partial_t h^2 = \beta_h(h^2, \lambda, \kappa) = 0,$$

$$\partial_t \lambda = \beta_\lambda(h^2, \lambda, \kappa) = 0.$$

\Rightarrow we obtain a conditional fixed-point

$$\partial_t \kappa = \beta_\kappa(h^{2*}, \lambda^*, \kappa) = 0. \quad \text{⊗}$$

Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.
- A leading order truncation can be parametrized by three couplings: h^2, λ, κ .

$$\partial_t h^2 = \beta_h(h^2, \lambda, \kappa) = 0,$$

$$\partial_t \lambda = \beta_\lambda(h^2, \lambda, \kappa) = 0.$$

\Rightarrow we obtain a conditional fixed-point

$$\partial_t \kappa = \beta_\kappa(h^{2*}, \lambda^*, \kappa) = 0. \quad \oplus$$

The β_κ -function receives the contributions

$$\beta_\kappa = -2\kappa + \text{bosonic interactions} - \text{fermionic interactions}$$

Fixed-point analysis for our toy model

- Whether or not the balancing is possible crucially depends on the d.o.f. of the model.
- A leading order truncation can be parametrized by three couplings: h^2, λ, κ .

$$\partial_t h^2 = \beta_h(h^2, \lambda, \kappa) = 0,$$

$$\partial_t \lambda = \beta_\lambda(h^2, \lambda, \kappa) = 0.$$

\Rightarrow we obtain a conditional fixed-point

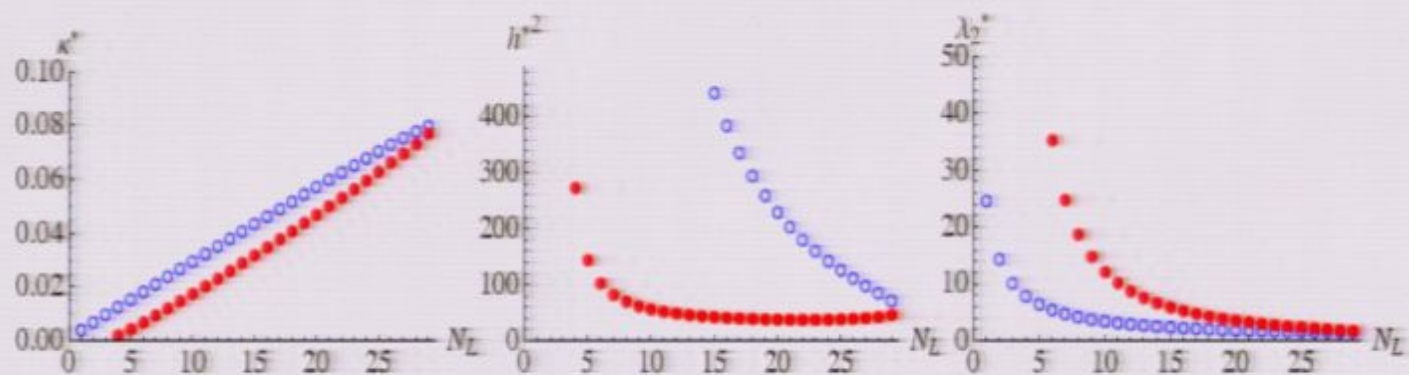
$$\partial_t \kappa = \beta_\kappa(h^{2*}, \lambda^*, \kappa) = 0. \quad \text{⊗}$$

The β_κ -function receives the contributions



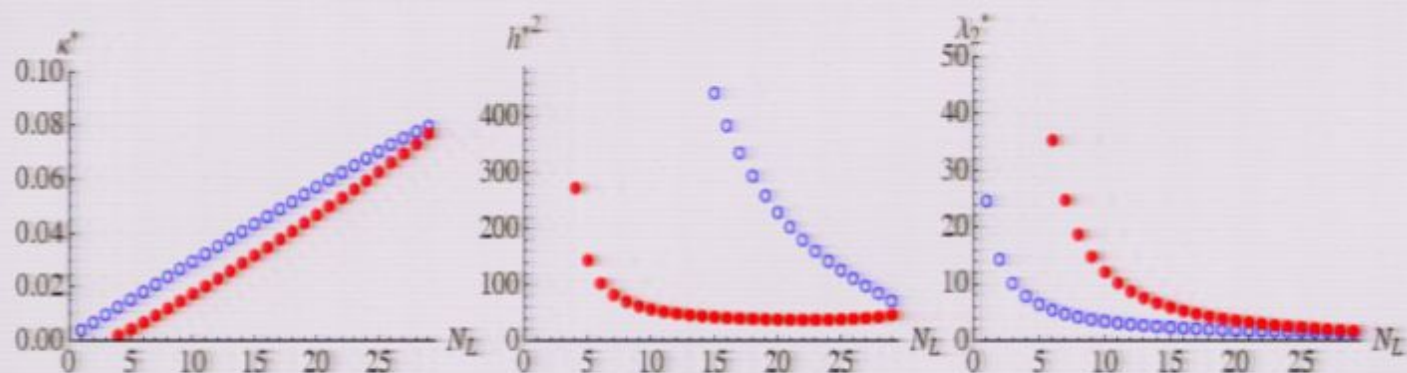
Fixed-points and critical exponents

We find a NGFPs for $1 \leq N_L \leq 57$

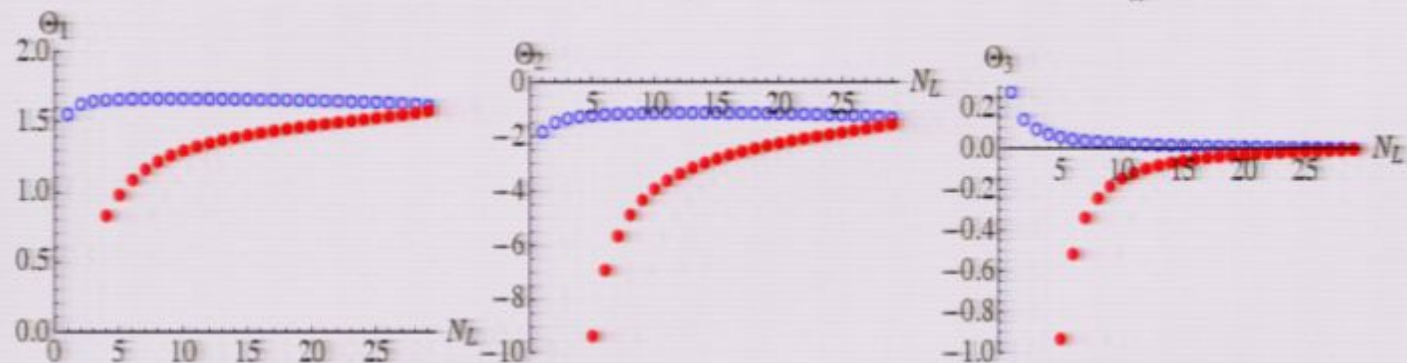


Fixed-points and critical exponents

We find a NGFPs for $1 \leq N_L \leq 57$



and the critical exponents:

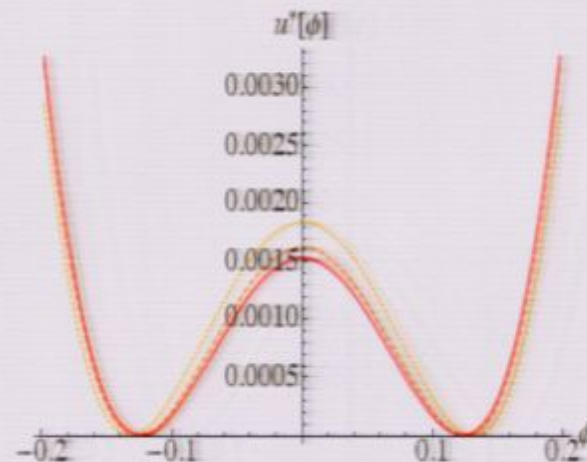


UV fixed-point regime for $N_L = 10$

- Example for a leading-order truncation expanded up to $\frac{\lambda_6}{6!} \rho^6$ in the effective potential and $N_L = 10$:

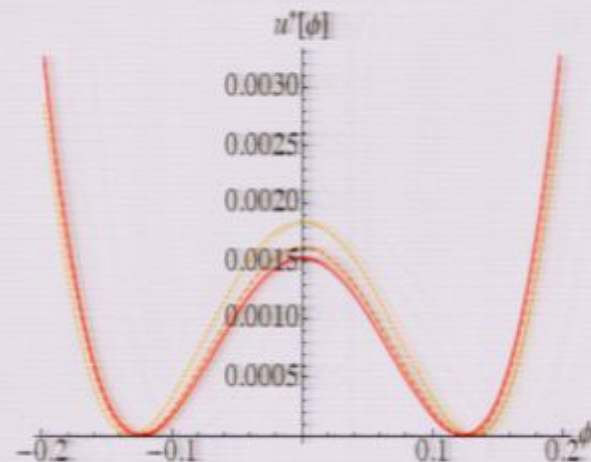
UV fixed-point regime for $N_L = 10$

- Example for a leading-order truncation expanded up to $\frac{\lambda_6}{6!} \rho^6$ in the effective potential and $N_L = 10$:
- Convergence of the fixed-point potential u^* at LO:



UV fixed-point regime for $N_L = 10$

- Example for a leading-order truncation expanded up to $\frac{\lambda_6}{6!} \rho^6$ in the effective potential and $N_L = 10$:
- Convergence of the fixed-point potential u^* at LO:

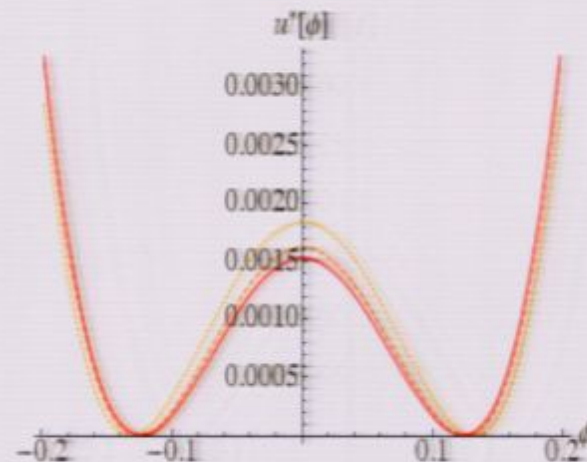


$$\text{FP : } \quad \kappa^* = 0.0152, \quad \lambda^* = 12.13, \quad h^{*2} = 57.41,$$

$$\text{Critical exponents : } \quad \Theta_1 = 1.056, \quad \Theta_2 = -0.175, \quad \Theta_3 = -2.350$$

UV fixed-point regime for $N_L = 10$

- Example for a leading-order truncation expanded up to $\frac{\lambda_6}{6!} \rho^6$ in the effective potential and $N_L = 10$:
- Convergence of the fixed-point potential u^* at LO:



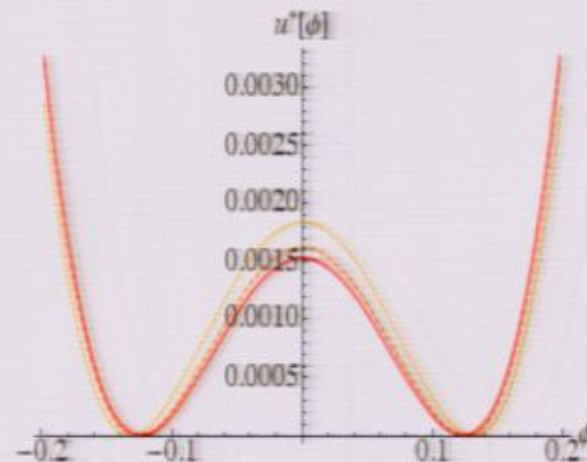
$$\text{FP : } \quad \kappa^* = 0.0152, \quad \lambda^* = 12.13, \quad h^{*2} = 57.41,$$

$$\text{Critical exponents : } \quad \Theta_1 = 1.056, \quad \Theta_2 = -0.175, \quad \Theta_3 = -2.350$$

- One relevant direction, corresponding to one physical parameter to be fixed.
- All other parameters are predictions from the theory.

UV fixed-point regime for $N_L = 10$

- Example for a leading-order truncation expanded up to $\frac{\lambda_6}{6!} \rho^6$ in the effective potential and $N_L = 10$:
- Convergence of the fixed-point potential u^* at LO:



$$\text{FP : } \quad \kappa^* = 0.0152, \quad \lambda^* = 12.13, \quad h^{*2} = 57.41,$$

$$\text{Critical exponents : } \quad \Theta_1 = 1.056, \quad \Theta_2 = -0.175, \quad \Theta_3 = -2.350$$

- One relevant direction, corresponding to one physical parameter to be fixed.
- All other parameters are predictions from the theory.
- The real part of the relevant direction is 1.056 and not anymore 2 \rightarrow Hierarchy problem weaker

(Toy-)Higgs mass and (Toy-)Top mass from asymptotic safety

- The flow can be fixed by one parameter, e.g. the IR value of κ .

(Toy-)Higgs mass and (Toy-)Top mass from asymptotic safety

- The flow can be fixed by one parameter, e.g. the IR value of κ .
- In a realistic model this would correspond to the vev which can be determined from the Z/W-boson masses:

$$v = \lim_{k \rightarrow 0} \sqrt{2\kappa k}$$

(Toy-)Higgs mass and (Toy-)Top mass from asymptotic safety

- The flow can be fixed by one parameter, e.g. the IR value of κ .
- In a realistic model this would correspond to the vev which can be determined from the Z/W-boson masses:

$$v = \lim_{k \rightarrow 0} \sqrt{2\kappa k}$$

- IR values of the other two parameters are predictions and are related to the Higgs and the top mass:

$$m_{\text{Higgs}} = \sqrt{\lambda_2} v, \quad m_{\text{top}} = \sqrt{h^2} v.$$



(Toy-)Higgs mass and (Toy-)Top mass from asymptotic safety

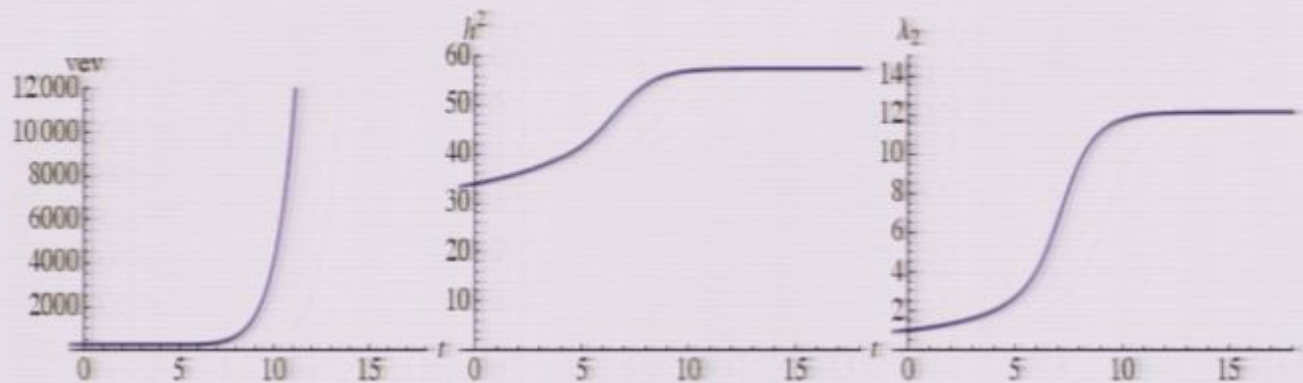
- The flow can be fixed by one parameter, e.g. the IR value of κ .
- In a realistic model this would correspond to the vev which can be determined from the Z/W-boson masses:

$$v = \lim_{k \rightarrow 0} \sqrt{2\kappa k}$$

- IR values of the other two parameters are predictions and are related to the Higgs and the top mass:

$$m_{\text{Higgs}} = \sqrt{\lambda_2} v, \quad m_{\text{top}} = \sqrt{h^2} v.$$

- Choosing $v = 246\text{GeV}$ and $N_L = 10$ as an example, we find



$$m_{\text{Higgs}} = 0.97v = 239\text{GeV}, \quad m_{\text{top}} = 5.56v = 1422\text{GeV}.$$

Discussion and Outlook



- The present theory reveals a possible AS mechanism for the standard model.

Discussion and Outlook



- The present theory reveals a possible AS mechanism for the standard model.
- Next to leading order in derivative expansion: NGFP might be destabilized due to Goldstone fluctuations and large values for the Yukawa coupling.



Discussion and Outlook



- The present theory reveals a possible AS mechanism for the standard model.
- Next to leading order in derivative expansion: NGFP might be destabilized due to Goldstone fluctuations and large values for the Yukawa coupling.
- We have massless Goldstone and fermion fluctuations, which are not present in the standard model.

Discussion and Outlook



- The present theory reveals a possible AS mechanism for the standard model.
- Next to leading order in derivative expansion: NGFP might be destabilized due to Goldstone fluctuations and large values for the Yukawa coupling.
- We have massless Goldstone and fermion fluctuations, which are not present in the standard model.
- In a simple Z_2 -symmetric Yukawa model (without Goldstone fluctuations) we observe a NLO fixed-point for $N_f < 1/3$.

Discussion and Outlook



- The present theory reveals a possible AS mechanism for the standard model.
- Next to leading order in derivative expansion: NGFP might be destabilized due to Goldstone fluctuations and large values for the Yukawa coupling.
- We have massless Goldstone and fermion fluctuations, which are not present in the standard model.
- In a simple Z_2 -symmetric Yukawa model (without Goldstone fluctuations) we observe a NLO fixed-point for $N_f < 1/3$.
- Also gravitational effects can be included: O. Zanusso, L. Zambelli, G. P. Vacca & R. Percacci

Discussion and Outlook



- The present theory reveals a possible AS mechanism for the standard model.
- Next to leading order in derivative expansion: NGFP might be destabilized due to Goldstone fluctuations and large values for the Yukawa coupling.
- We have massless Goldstone and fermion fluctuations, which are not present in the standard model.
- In a simple Z_2 -symmetric Yukawa model (without Goldstone fluctuations) we observe a NLO fixed-point for $N_f < 1/3$.
- Also gravitational effects can be included: O. Zanusso, L. Zambelli, G. P. Vacca & R. Percacci
- Next step: Include $SU(N_L)$ gauge bosons (work in progress with H. Gies and S. Rechenberger)

(Toy-)Higgs mass and (Toy-)Top mass from asymptotic safety

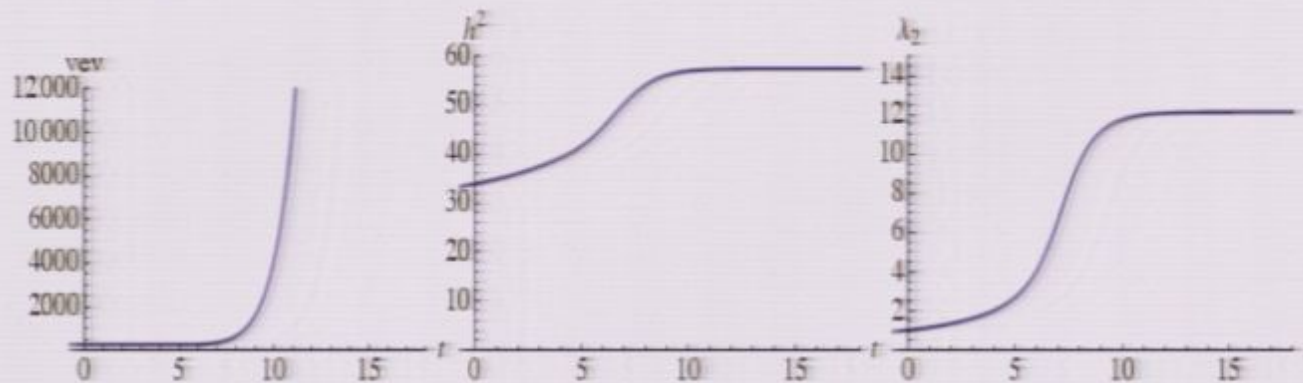
- The flow can be fixed by one parameter, e.g. the IR value of κ .
- In a realistic model this would correspond to the vev which can be determined from the Z/W-boson masses

$$v = \lim_{k \rightarrow 0} \sqrt{2\kappa k}$$

- IR values of the other two parameters are predictions and are related to the Higgs and the top mass

$$m_{\text{Higgs}} = \sqrt{\lambda_2} v, \quad m_{\text{top}} = \sqrt{h^2} v.$$

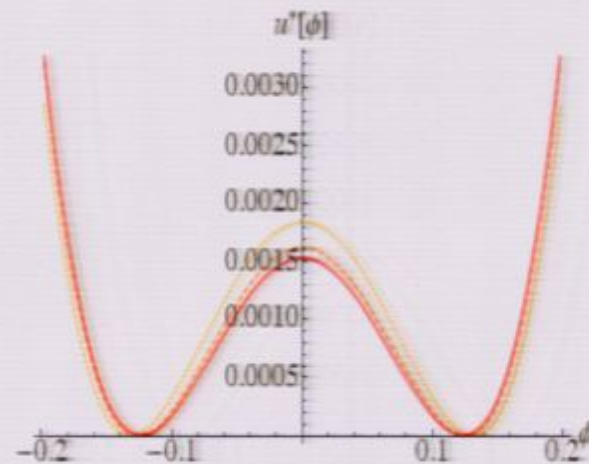
- Choosing $v = 246\text{GeV}$ and $N_L = 10$ as an example, we find



$$m_{\text{Higgs}} = 0.97v = 239\text{GeV}, \quad m_{\text{top}} = 5.56v = 1422\text{GeV}.$$

UV fixed-point regime for $N_L = 10$

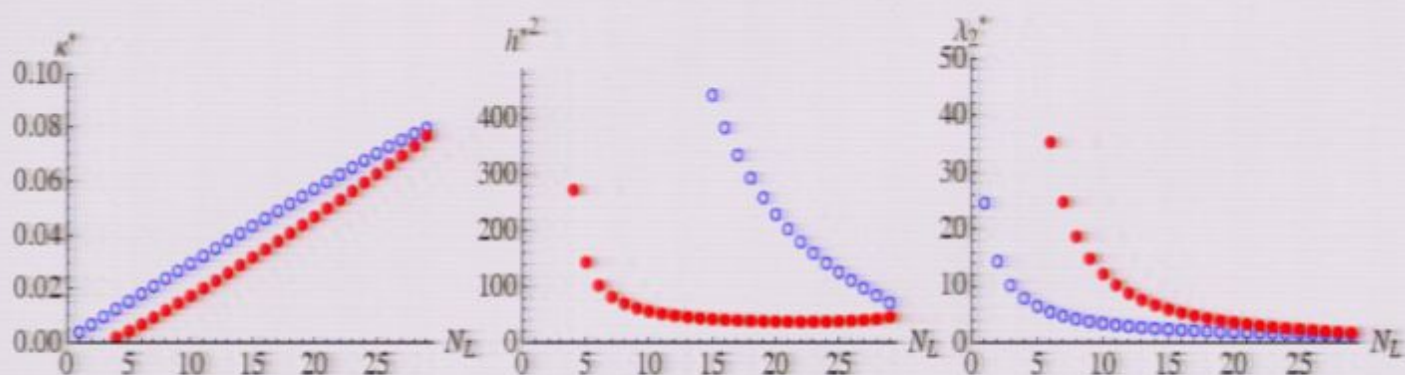
- Example for a leading-order truncation expanded up to $\frac{\lambda_6}{6!}\rho^6$ in the effective potential and $N_L = 10$:
- Convergence of the fixed-point potential u^* at LO:



$$\text{FP : } \quad \kappa^* = 0.0152, \quad \lambda^* = 12.13, \quad h^{*2} = 57.41,$$

Fixed-points and critical exponents

We find a NGFPs for $1 \leq N_L \leq 57$



and the critical exponents:

