

Title: Quantum computational phases of matter: measurement-based quantum computing in the Haldane phase

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Abstract: A recent breakthrough in quantum computing has been the realization that quantum computation can proceed solely through single-qubit measurements on an appropriate quantum state. One exciting prospect is that the ground or low-temperature thermal state of an interacting quantum many-body system can serve as such a resource state for quantum computation. The system would simply need to be cooled sufficiently and then subjected to local measurements. It would be unfortunate, however, if the usefulness of a ground or low-temperature thermal state for quantum computation was critically dependent on the details of the system's Hamiltonian; if so, engineering such systems would be difficult or even impossible. A much more powerful result would be the existence of a robust ordered phase which is characterized by the ability to perform measurement-based quantum computation. I'll discuss some recent results on the existence of such a computational phase of matter. I'll first outline some positive results on a phase of a toy model that contains the cluster state. Then, in a realistic model of coupled spin-1 particles, I'll demonstrate the existence of a computational phase. This result is obtained by using a local measurement sequence to "renormalize" the state to a computationally-universal fixed point. Together, these results reveal that the characterization of computational phases of matter has a rich, complex structure — one which is still poorly understood. Joint work with Gavin Brennen, Akimasa Miyake, and Joseph Renes.

# Quantum computational phases of matter: Measurement-based quantum computing in the Haldane phase

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**Stephen Bartlett**



The University of Sydney  
AUSTRALIA

School of Physics

*in collaboration with:*

Gavin Brennen (Macquarie)

Akimasa Miyake (Perimeter Institute)

Joe Renes (TU Darmstadt)

# Measurement-based quantum computing

VOLUME 86, NUMBER 22

PHYSICAL REVIEW LETTERS

28 MAY 20

- Quantum computing can proceed through *measurements* rather than unitary evolution
- Measurements are strong and incoherent: easier?

Uses a resource state such as the *cluster state*:

- a universal circuit board
- a lattice of spins in a specific *entangled* state

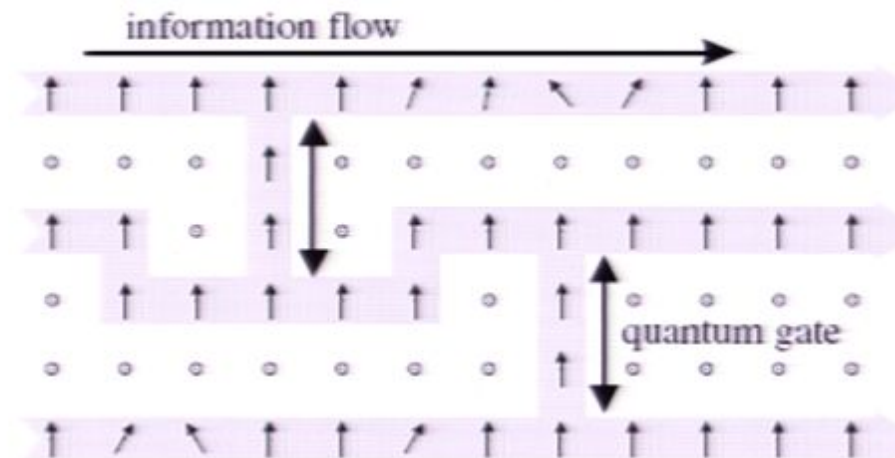
## A One-Way Quantum Computer

Robert Raussendorf and Hans J. Briegel

*Theoretische Physik, Ludwig-Maximilians-Universität München, Germany*

(Received 25 October 2000)

We present a scheme of quantum computation that consists entirely of one-qubit measurements on a particular class of entangled states, the cluster states. The measurements are used to imprint a quantum logic circuit on the state, thereby destroying its entanglement at the same time. Cluster states are thus one-way quantum computers and the measurements form the program.



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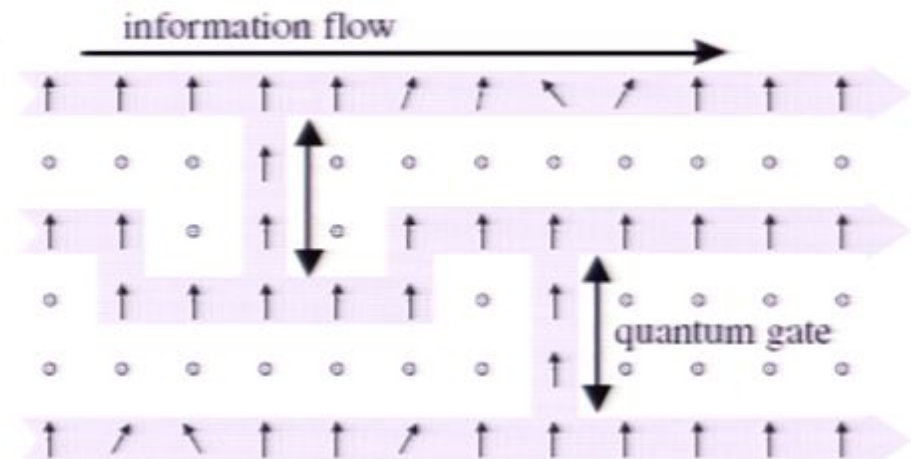
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Computational properties → properties of states

**Q:** What makes the cluster state work?

**Q:** What properties of a state are necessary?

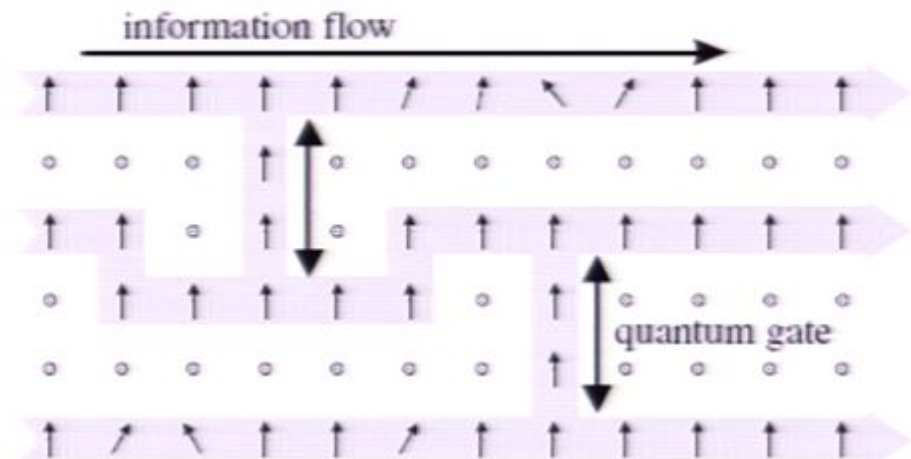
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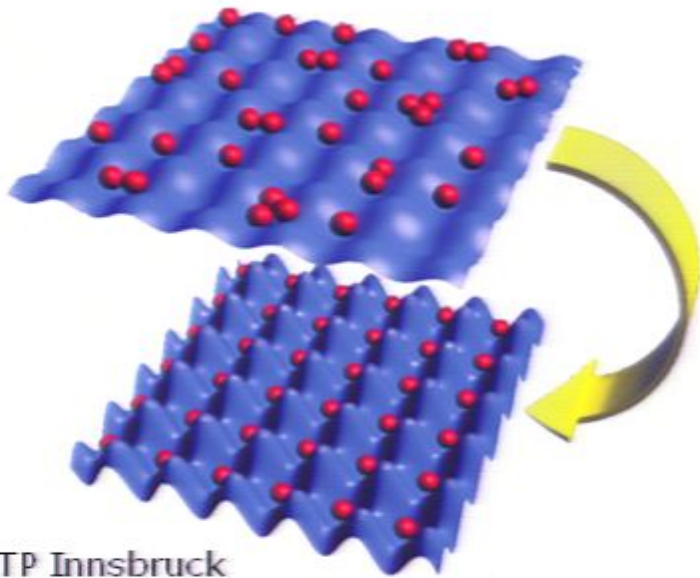
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# Resource states & Hamiltonians

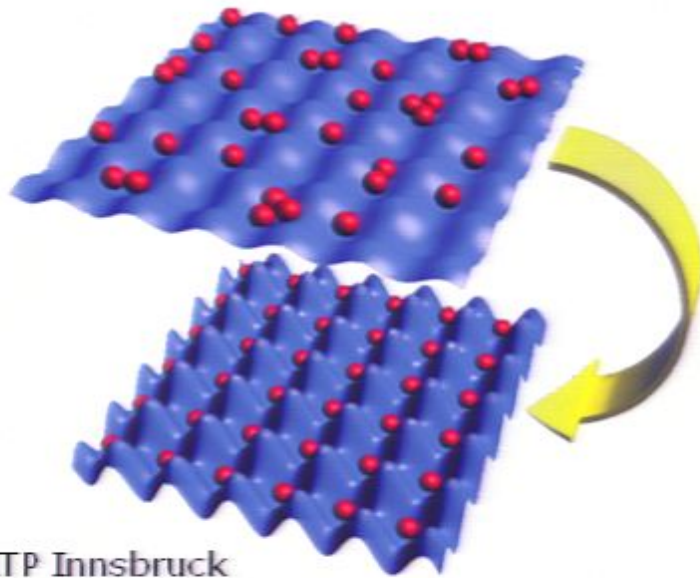


- Resource states can be:
  - constructed with unitary gates
  - the ground state of a coupled quantum many-body system (Raussendorf *et al.*)

ITP Innsbruck  
[www.uibk.ac.at/th-physik/qo/research/](http://www.uibk.ac.at/th-physik/qo/research/)

$$H_{\text{res}} = -\sum_i \sigma_i^x - \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z$$

# Resource states & Hamiltonians



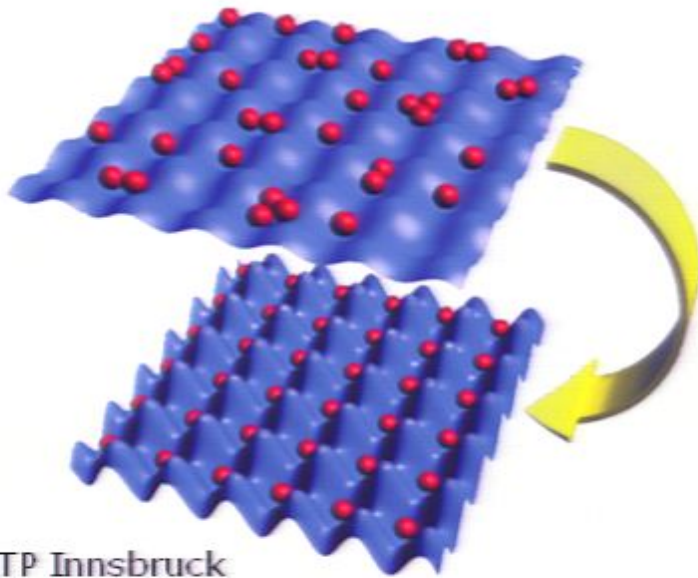
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$$H_{\text{cluster}} = - \sum_{\text{sites}} \begin{array}{c} Z \\ | \\ Z-X-Z \\ | \\ Z \end{array}$$



# Resource states & Hamiltonians



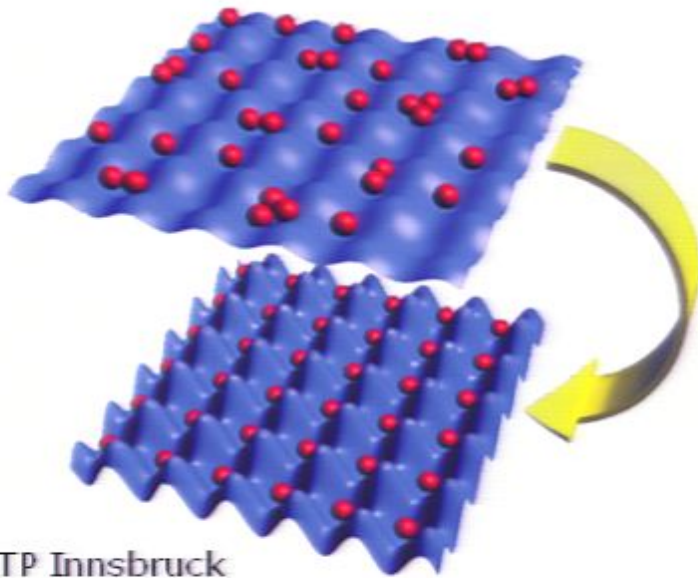
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- Is the system **fragile** or **robust** to local perturbations, or finite temperature?



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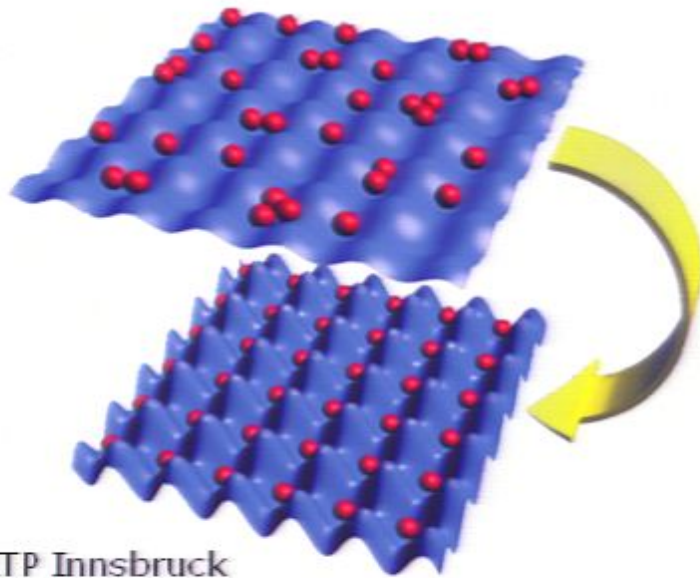


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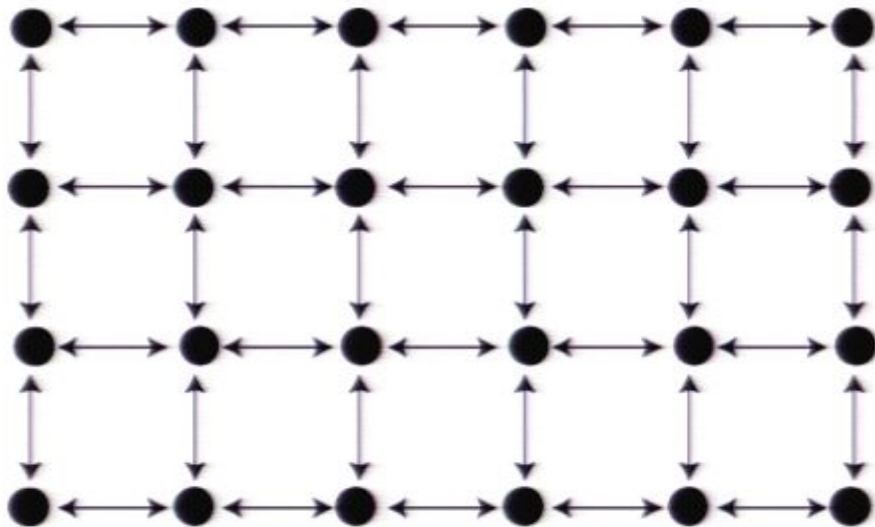
**Central idea:** identify suitable *order parameters* that characterize the ability to perform MBQC on a state

# Identifying phases, part 1

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# MBQC gates and correlation functions

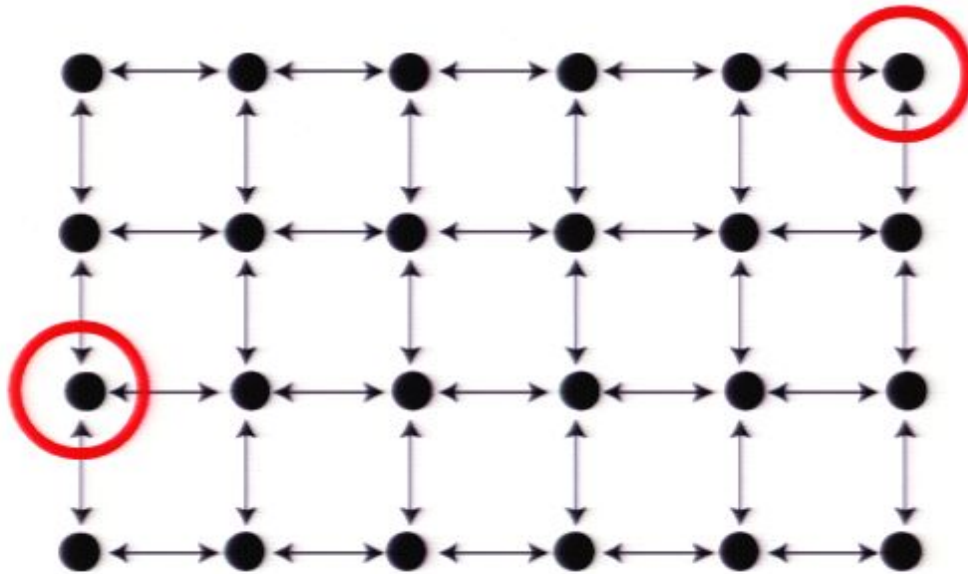
- With a cluster state, you can teleport a qubit (identity gate) between any two points on the lattice using local measurements on the qubits in between





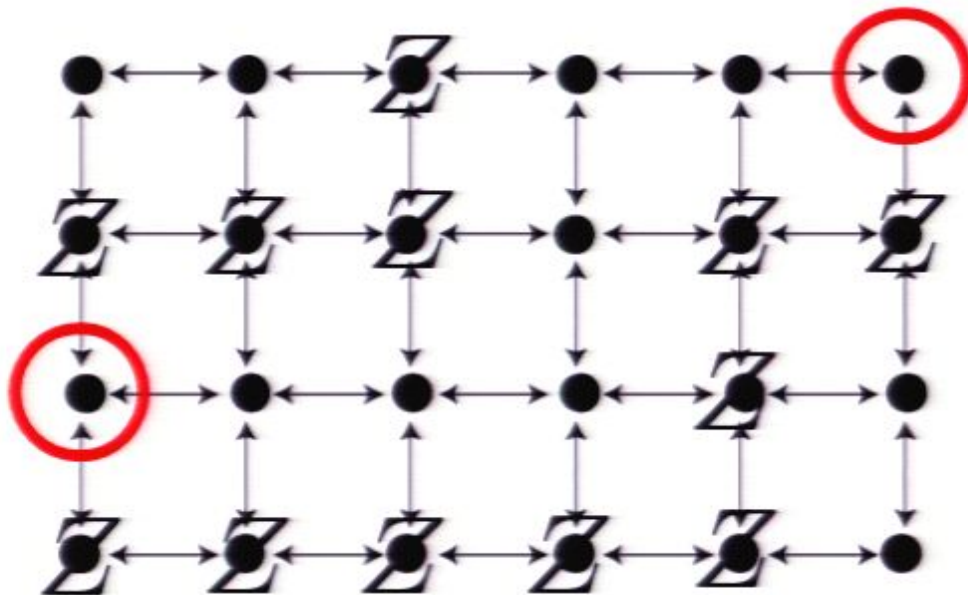
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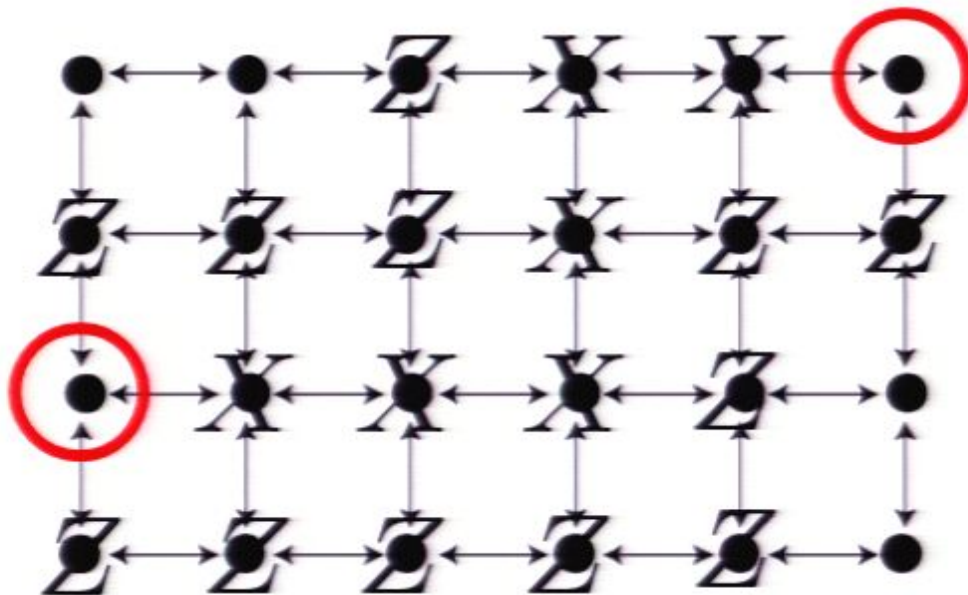
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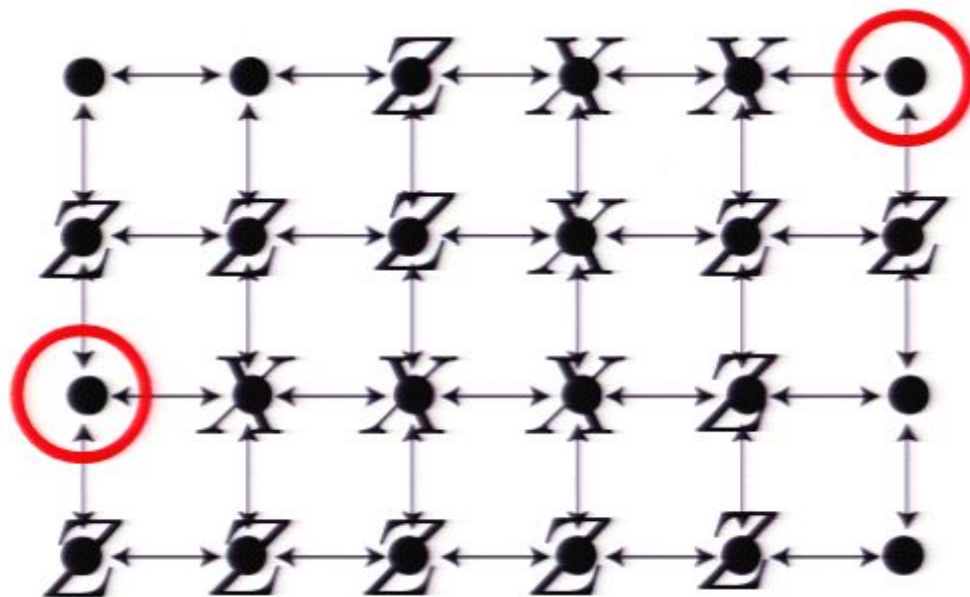
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- Measurement outcomes determine a Pauli correction to gate
- Other gates are given by different measurement patterns
- Performance of a gate is quantified by the fidelity of the resulting entangled state



# MBQC gates and correlation functions

**Approach:** Quantum gates as correlation functions

- Order parameters to characterise *phases* which are universal for MBQC



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$$\rho_{\text{pm}} = \sum_J U_J P_J \rho_0 P_J U_J^\dagger$$

Measurement projector

Unitary correction

The equation is annotated with red arrows. One arrow points from the text "Measurement projector" to the  $P_J$  term in the equation. Another arrow points from the text "Unitary correction" to the  $U_J^\dagger$  term in the equation.

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$$\langle AB \rangle_{\text{pm}} = \sum_J \text{Tr}[(AB)U_J P_J \rho_0 P_J U_J^\dagger] = \sum_J \text{Tr}[(AB_J)P_J \rho_0 P_J]$$

$$B_J = U_J B U_J^\dagger$$

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Measurement projector (pointing to  $P_J$ )  
 Unitary correction (pointing to  $U_J^\dagger$ )

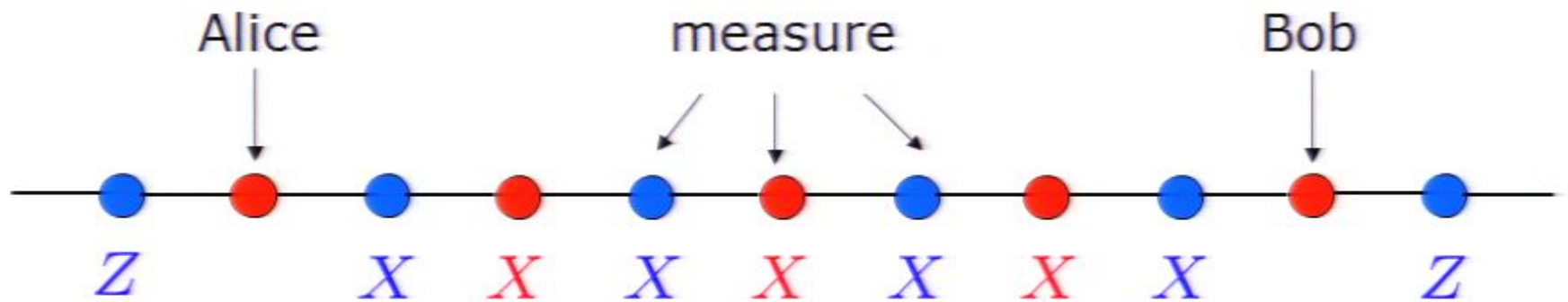
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# Example: Cluster-state gates



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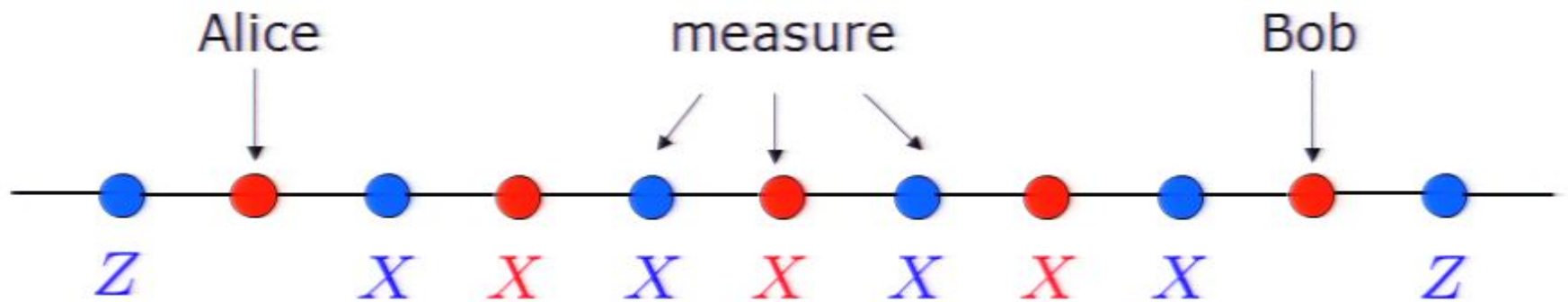
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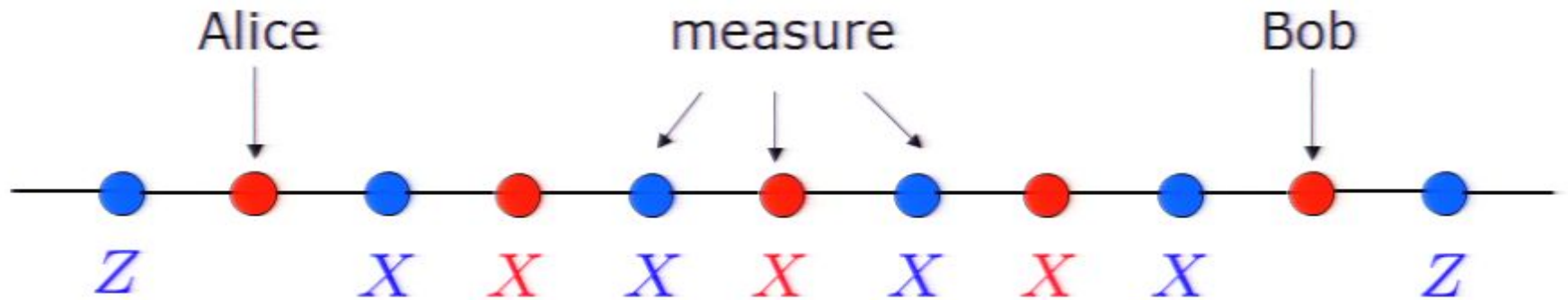
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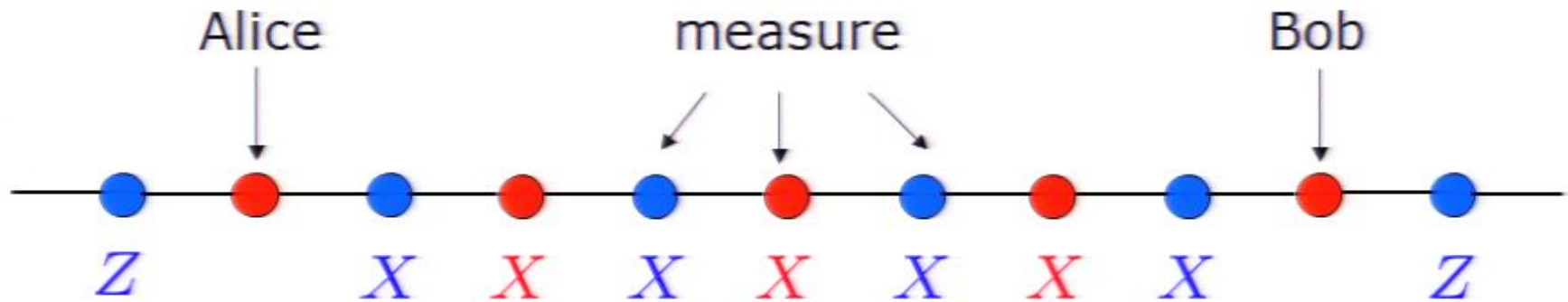


String-like order parameter  
Quantifies fidelity of 1-  
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$$\langle XX \rangle_{\text{pm}} = \langle Z_{2i-1} \left( \sum_{k=i}^{j-1} X_{2k} \right) Z_{2j-1} \rangle$$

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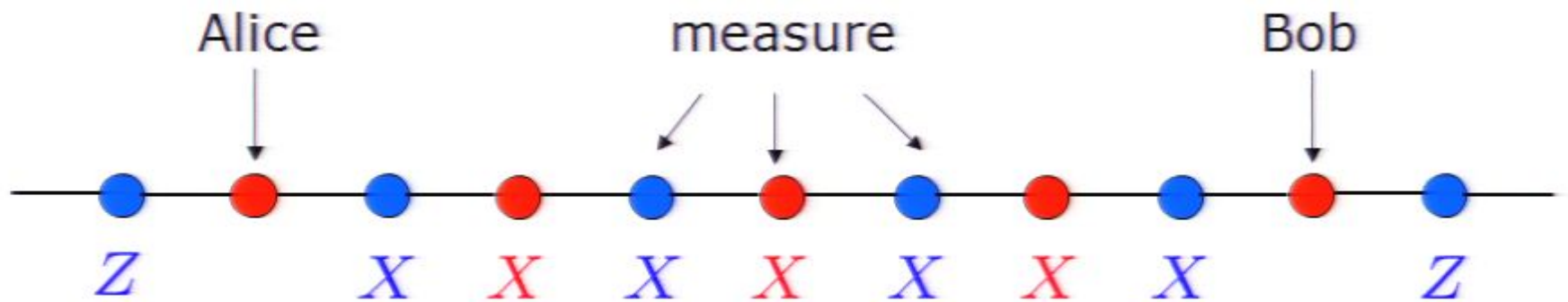
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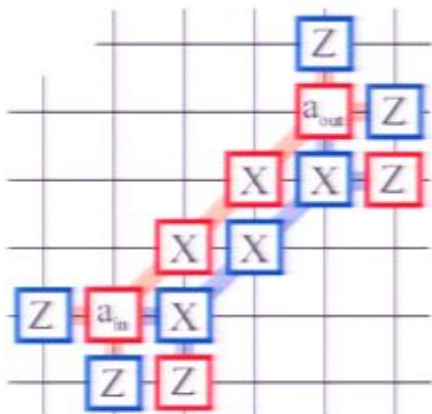
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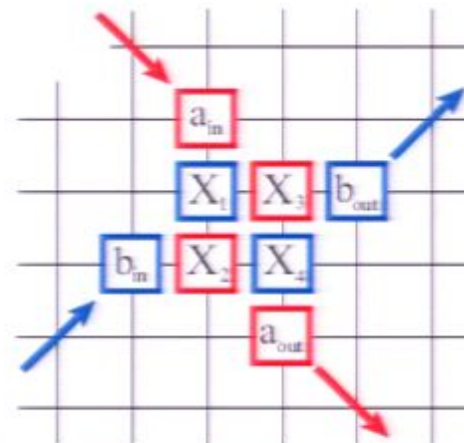
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2D:



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**Interacting strings order parameter**  
Quantifies CNOT gate fidelity

## Cluster Hamiltonian with transverse field

$$H = - \sum_{\text{sites } i} K_i + BX_i$$



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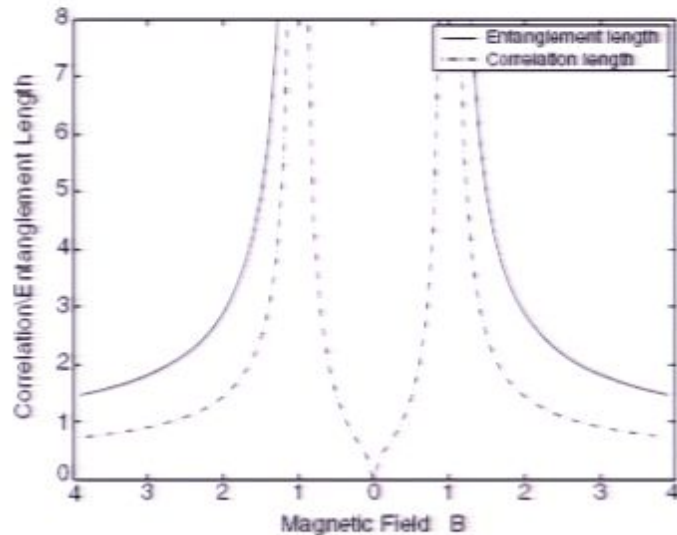
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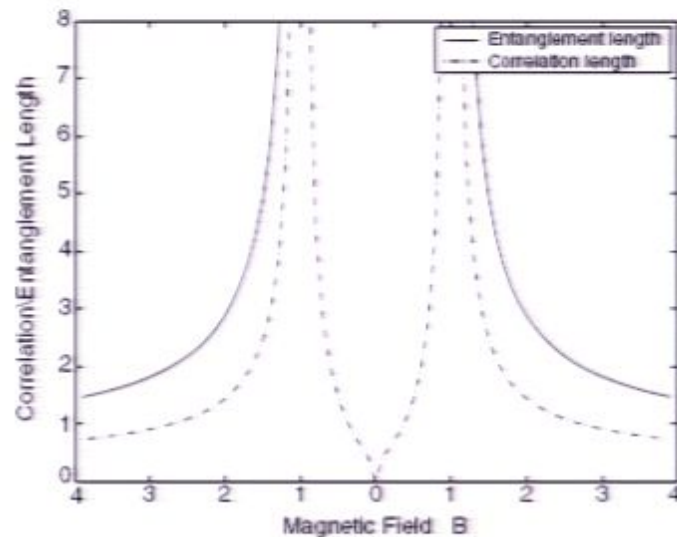
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2D: Universal gate set for MBQC is long-ranged in  $B < 1$  phase

Doherty and Bartlett

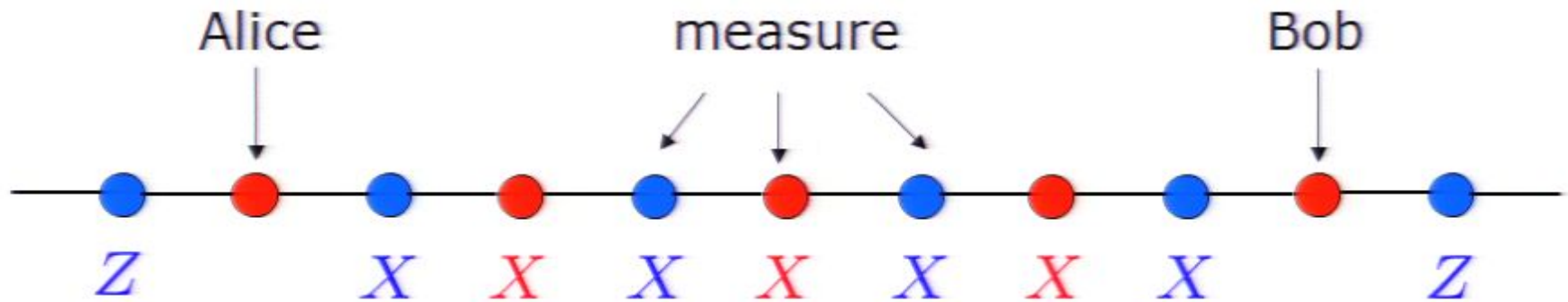
*PRL 103, 020506 (2009)*



# Conclusions 1.0

- Quantum gates – correlation functions – order parameters to identify **universal phases** for MBQC
- MBQC is a new type of long range “string” order
- **Problem:** the ‘cluster Hamiltonian’ is a toy model that does not describe a realistic system
- Can we apply these results to a realistic spin system?

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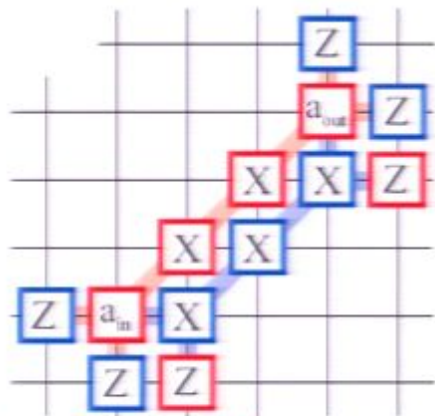


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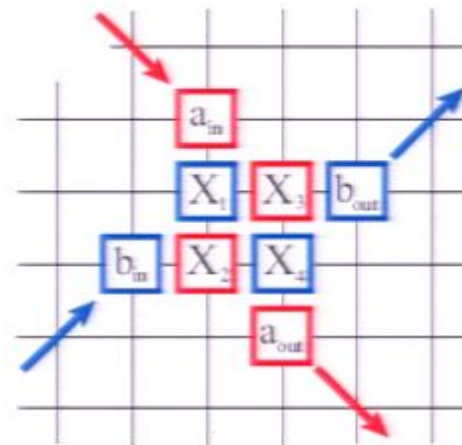
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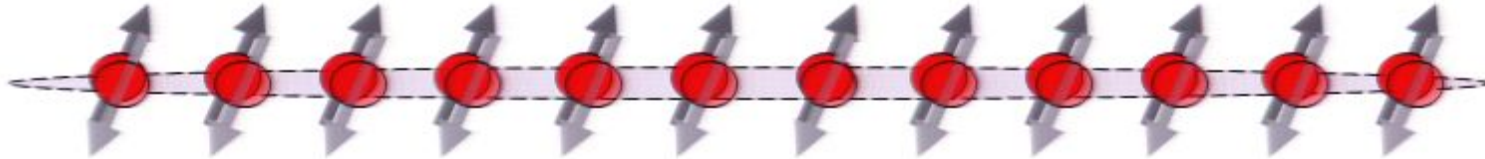
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# Identifying phases, part 2

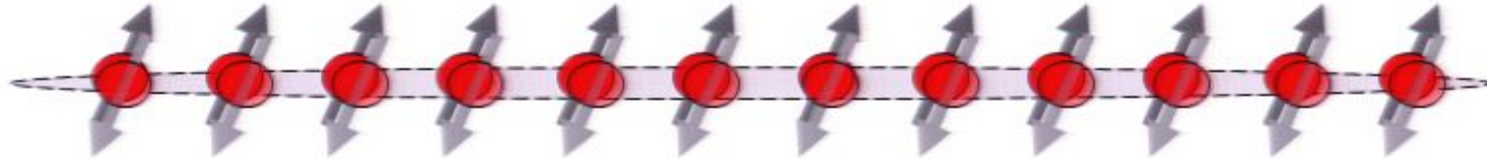
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# Affleck-Kennedy-Lieb-Tasaki (AKLT)



Rotationally-invariant spin chains

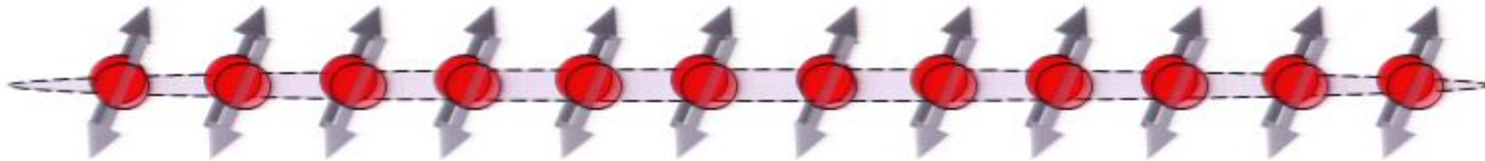
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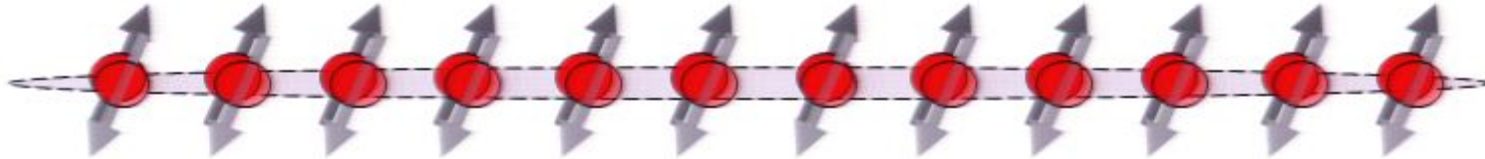
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- **Ground state:** valence bond solid / projected entangled pair state (PEPS)



$P =$  triplet projection (isometry)

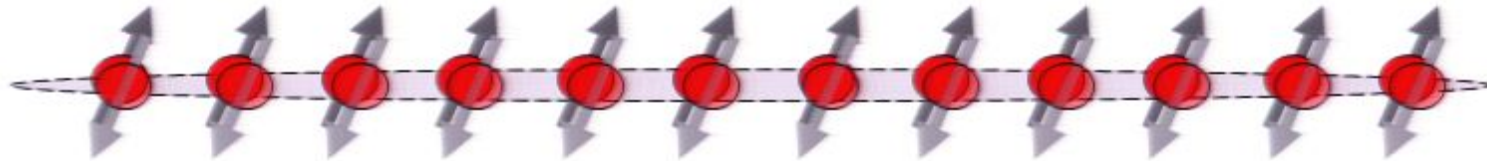


# AKLT as a Matrix Product State

- PEPS yields a Matrix Product State (MPS) description of the AKLT ground state

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$$|x\rangle = -|J_x = 0\rangle = -\frac{1}{\sqrt{2}}(|J_z = 1\rangle - |J_z = -1\rangle) \quad \sigma_1 = X$$

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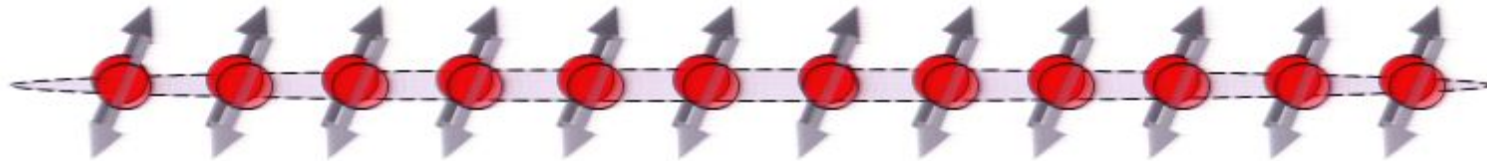
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# AKLT as a Matrix Product State

- PEPS yields a Matrix Product State (MPS) description of the AKLT ground state

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# Affleck-Kennedy-Lieb-Tasaki (AKLT)



## Rotationally-invariant spin chains

- **Haldane conjecture (1983):** ground state properties of Heisenberg spin chains depend on the spin. Integer spin chains are *gapped*.
- **AKLT state (1987):** solution for a spin-1 Heisenberg chain with a gap.

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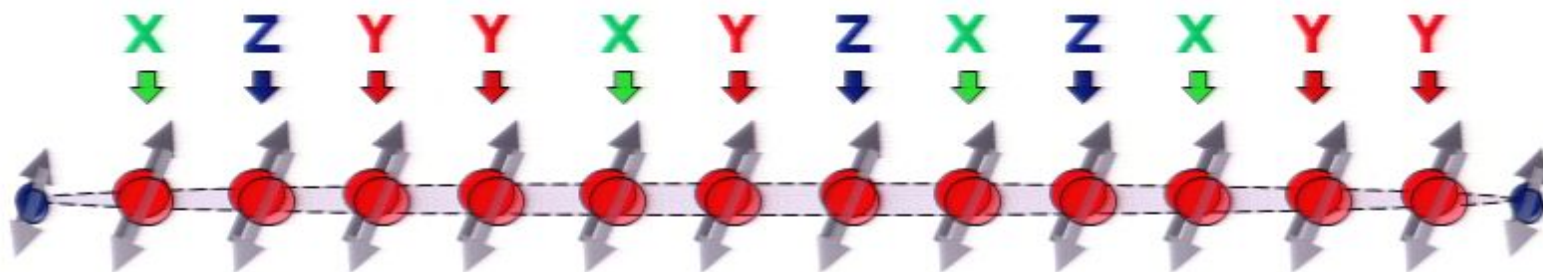
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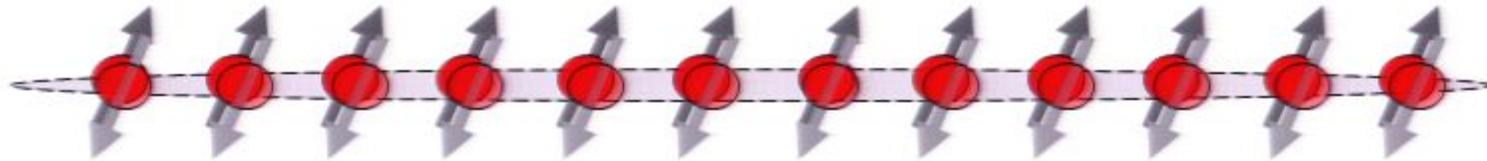
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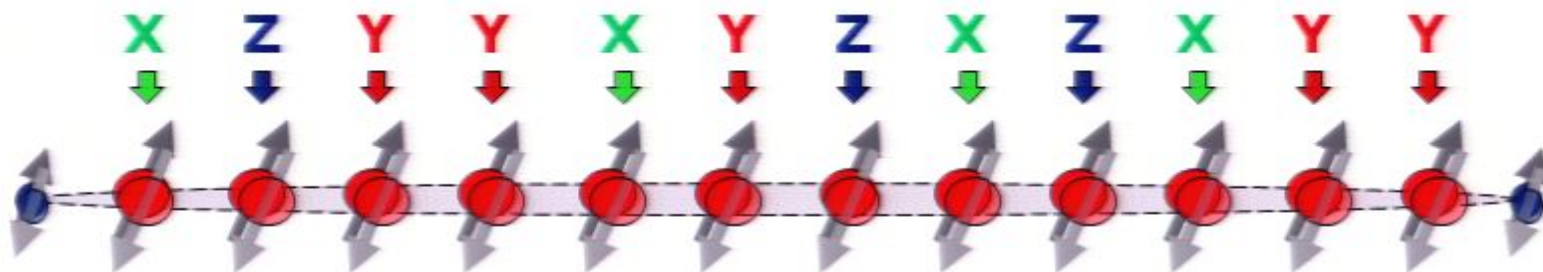
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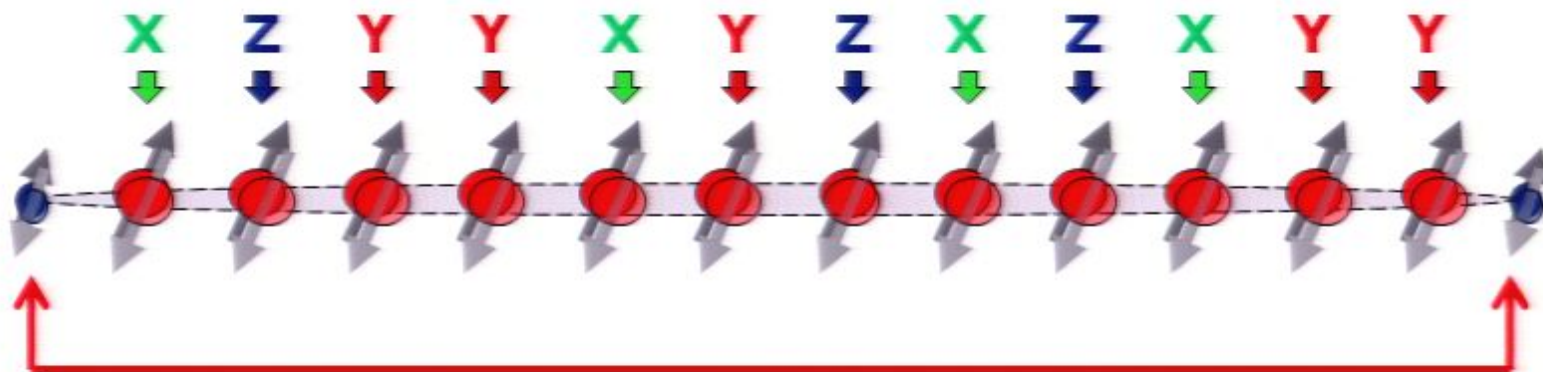
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$$(X^{\#x,y} Z^{\#y,z})_B |\Psi^-\rangle_{AB}$$

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# MBQC using AKLT states

PRL 101, 010502 (2008)

PHYSICAL REVIEW LETTERS

week ending  
4 JULY 2008

## Measurement-Based Quantum Computer in the Gapped Ground State of a Two-Body Hamiltonian

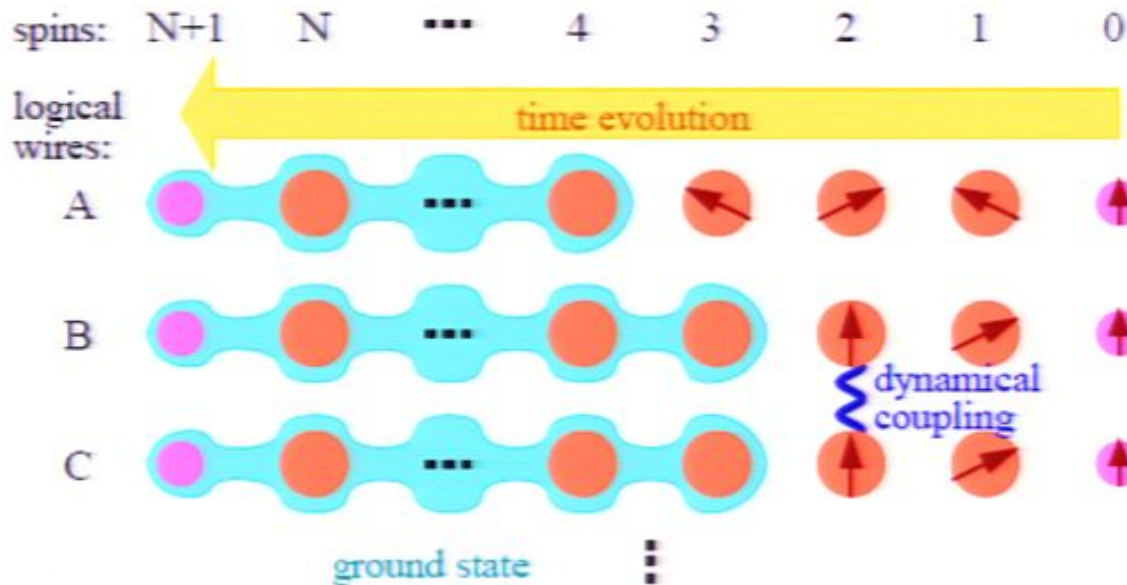
Gavin K. Brennen<sup>1,2</sup> and Akimasa Miyake<sup>2,3</sup>

<sup>1</sup>Department of Physics, Macquarie University, Sydney, NSW 2109, Australia

<sup>2</sup>Institute for Quantum Optics and Quantum Information, Austrian Academy of Sciences, Innsbruck, Austria

<sup>3</sup>Institute for Theoretical Physics, University of Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria

(Received 12 March 2008; published 2 July 2008)



- 2-D spin lattice
- Individual chains as quantum comp wires
- Two-qubit logic (CPHASE) by dynamical coupling
- Gap protects against thermal noise

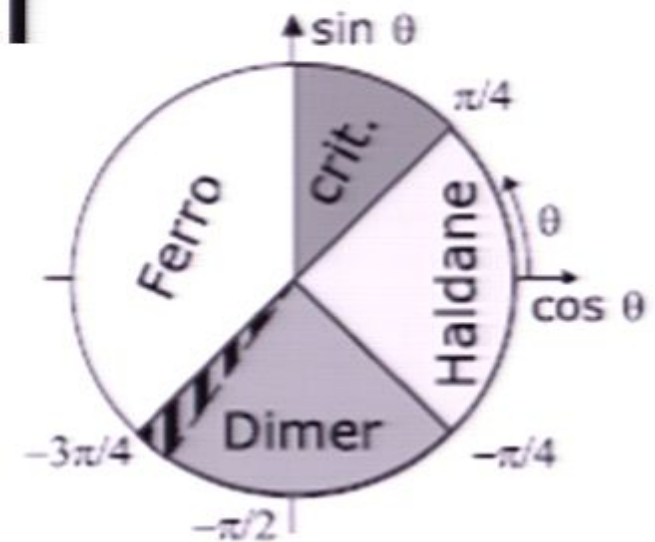
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Romero-Isart, Eckert, Sanpera  
PRA 75, 050303(R) (2007)

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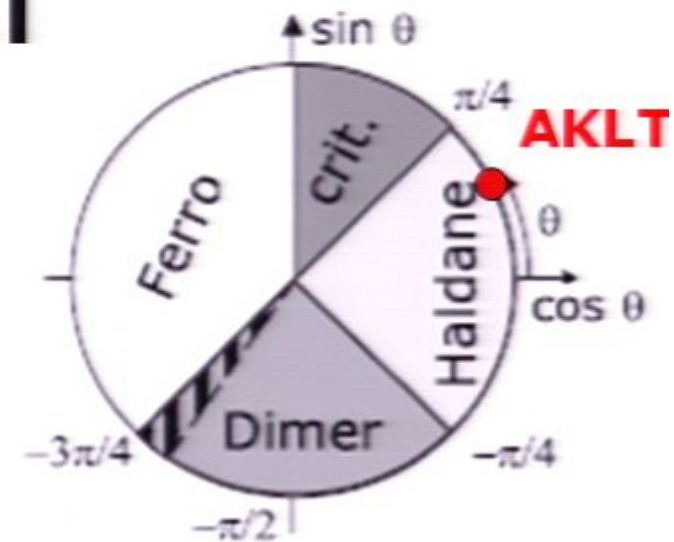
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Romero-Isart, Eckert, Sanpera  
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# MBQC gates and correlation functions

**Approach:** Quantum gates as correlation functions

- Order parameters to characterise *phases* which are universal for MBQC



$$\rho = \sum_J U_J P_J \rho_0 P_J U_J^\dagger$$

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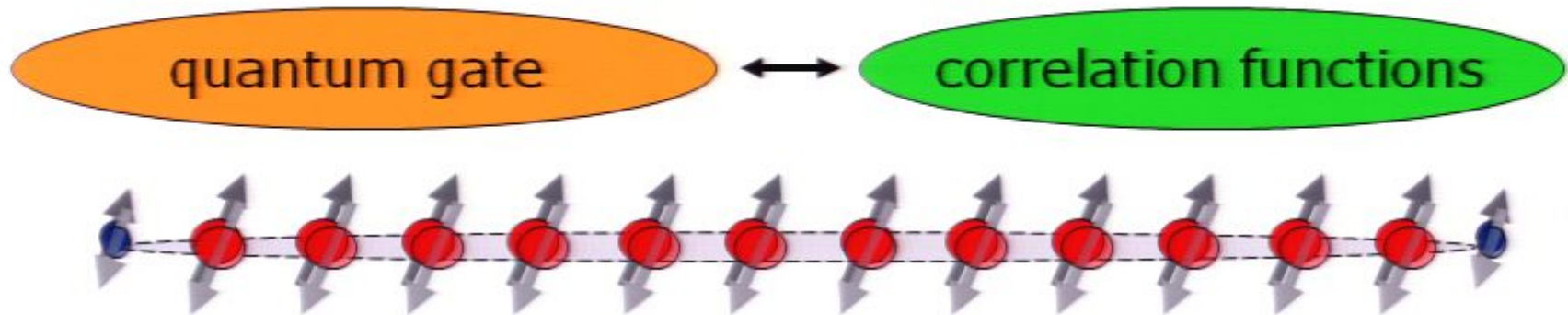
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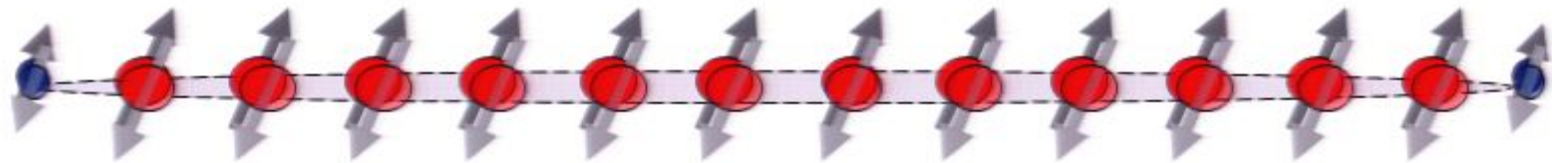
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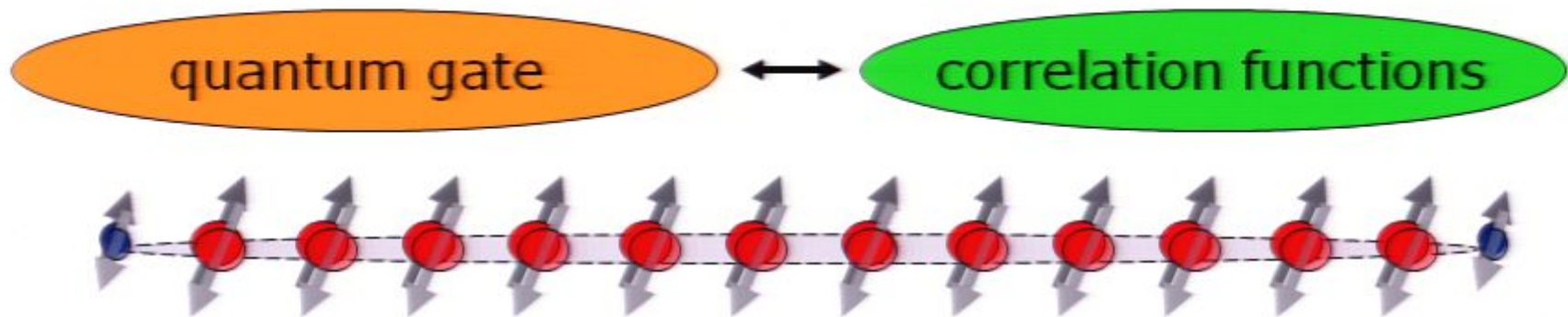
$$\langle XX \rangle_{\text{pm}} = \langle X \vec{V}_x X \rangle$$

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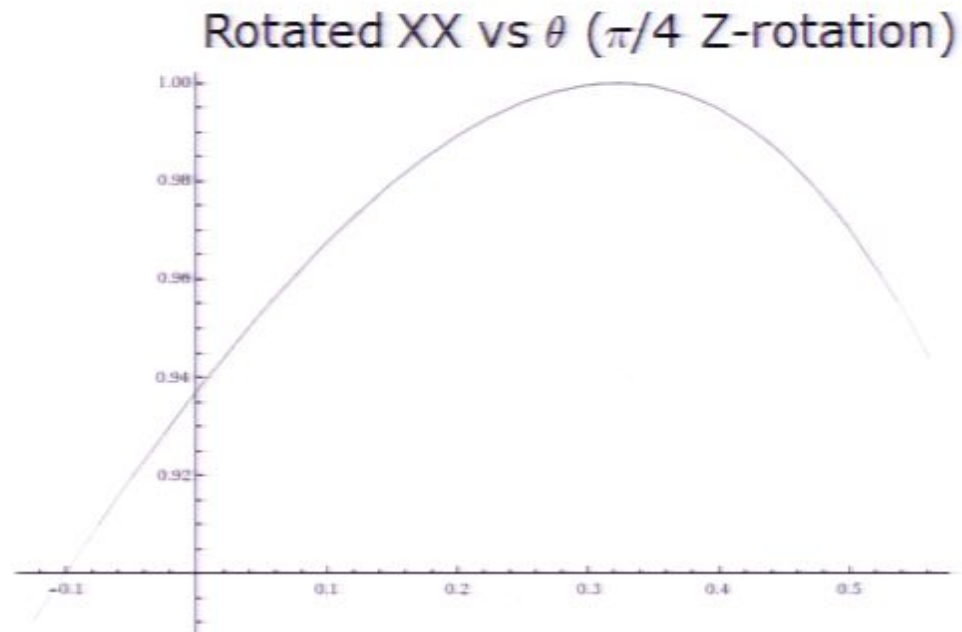
**Remarkable result:** due to rotational-invariance of our Hamiltonian, these order parameters are 1 at any length, for any ground state in the Haldane phase

See Popp et al.,  
PRA 71, 042306 (2005)

Order parameters exist to quantify the identity gate, and are long-ranged in the Haldane phase

# Numerical results

- Find ground states of the open/fixed chain using DMRG
- Localizable entanglement = 1 throughout Haldane phase
- Rotation fidelity drops off slowly with theta



# Why are rotations failing?

- Re-express the Hamiltonian as

$$H(\theta_{\text{AKLT}} + \varepsilon) = H_{\text{AKLT}} + \frac{5}{6}\varepsilon \sum_{j=1}^n \text{SWAP}_{j,j+1}$$

- Ansatz:

$$|\text{gs}(\theta)\rangle \simeq \alpha|\text{AKLT}\rangle + \sum_{j=1}^n \beta_j \text{SWAP}_{j,j+1}|\text{AKLT}\rangle$$

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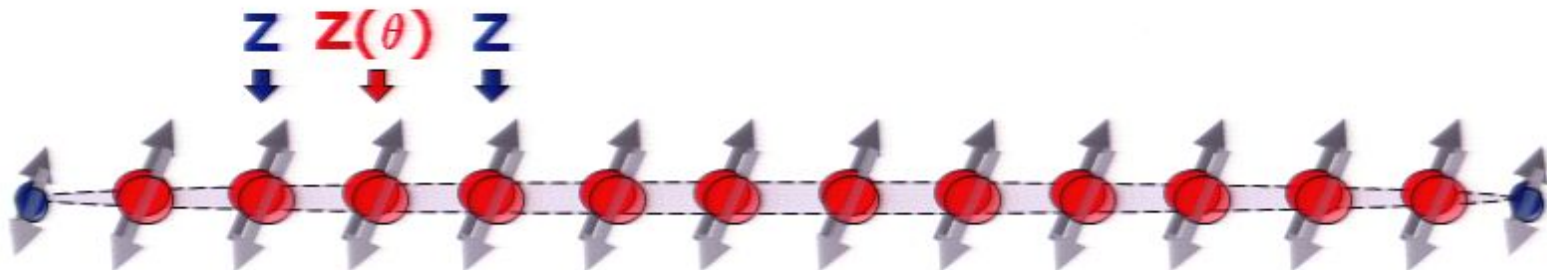
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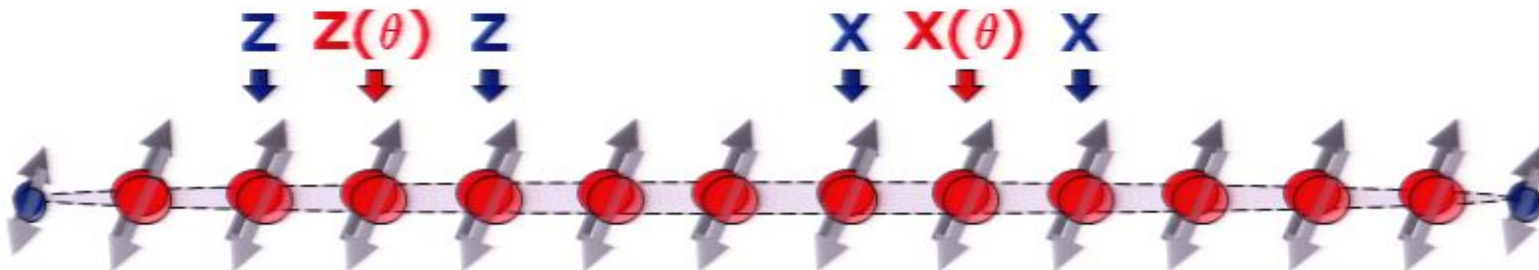
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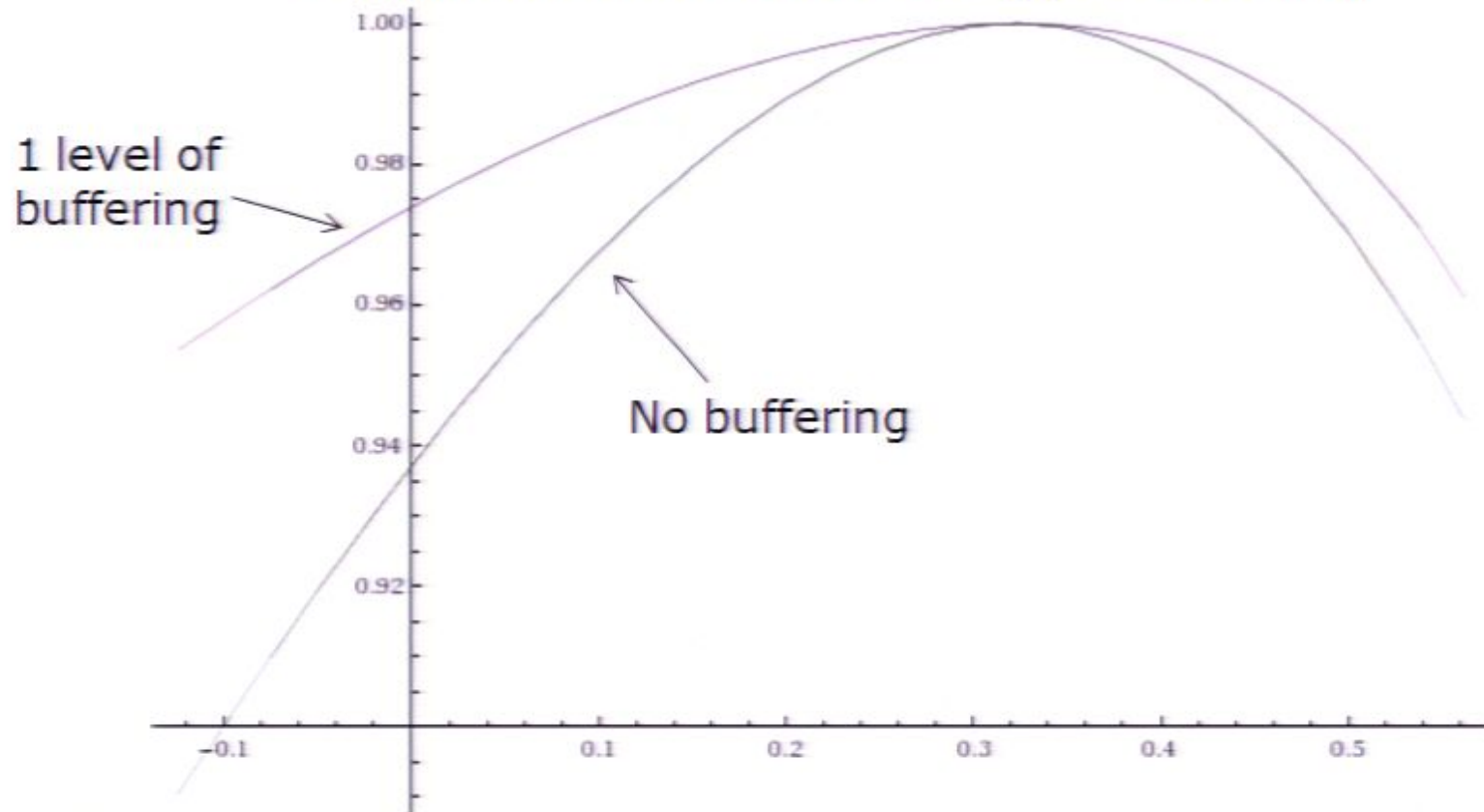
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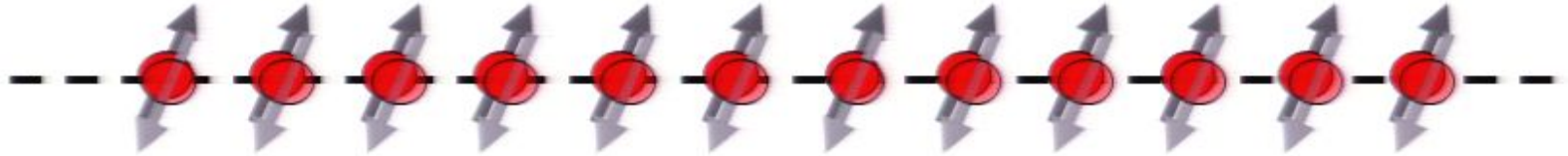


# Buffered rotations

Comparison of rotated XX vs  $\theta$  ( $\pi/4$  Z-rotation)

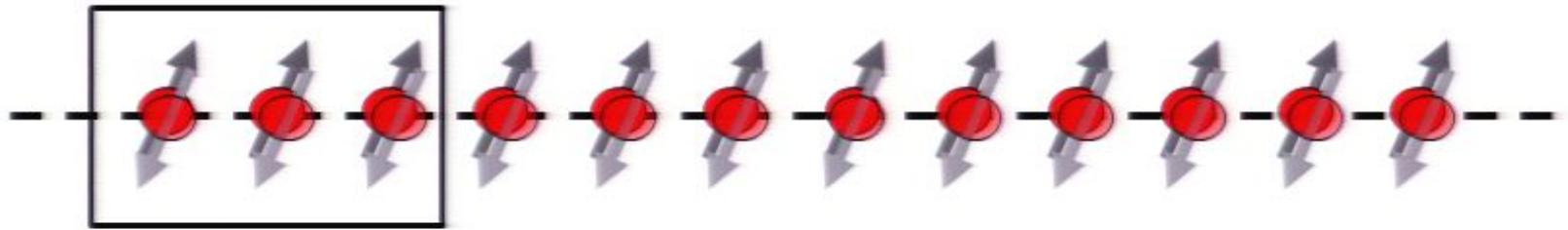


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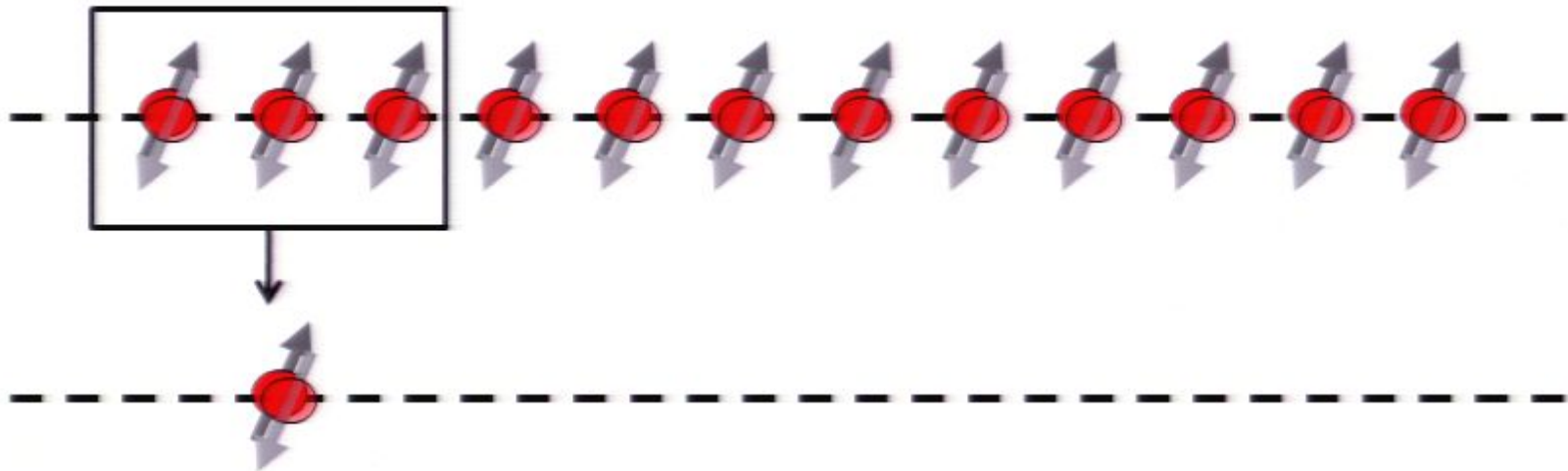




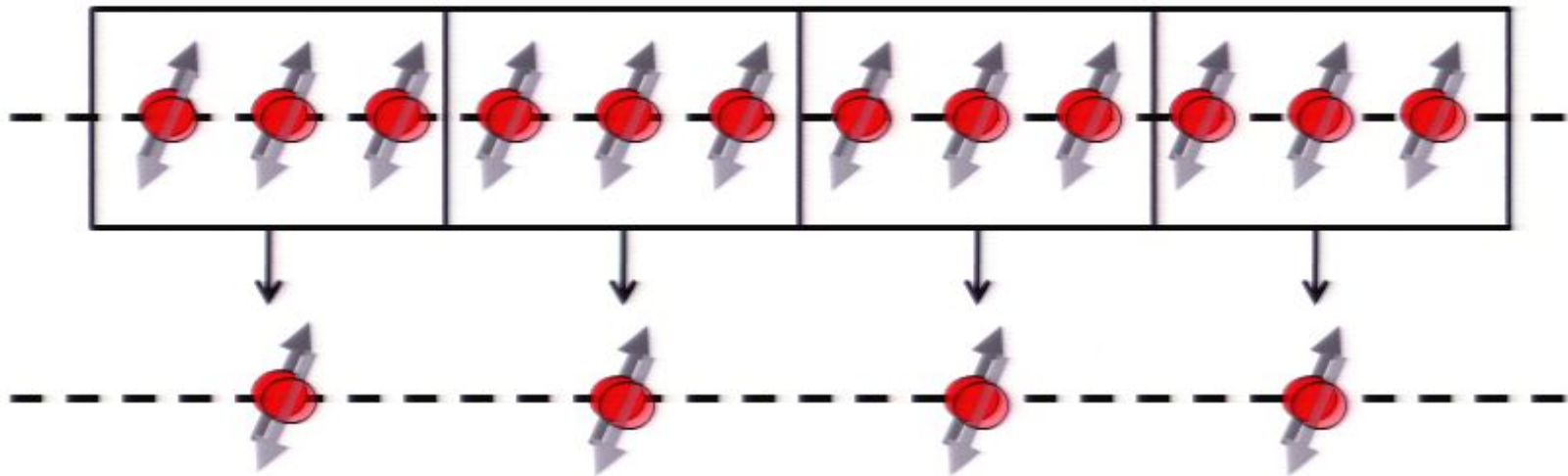
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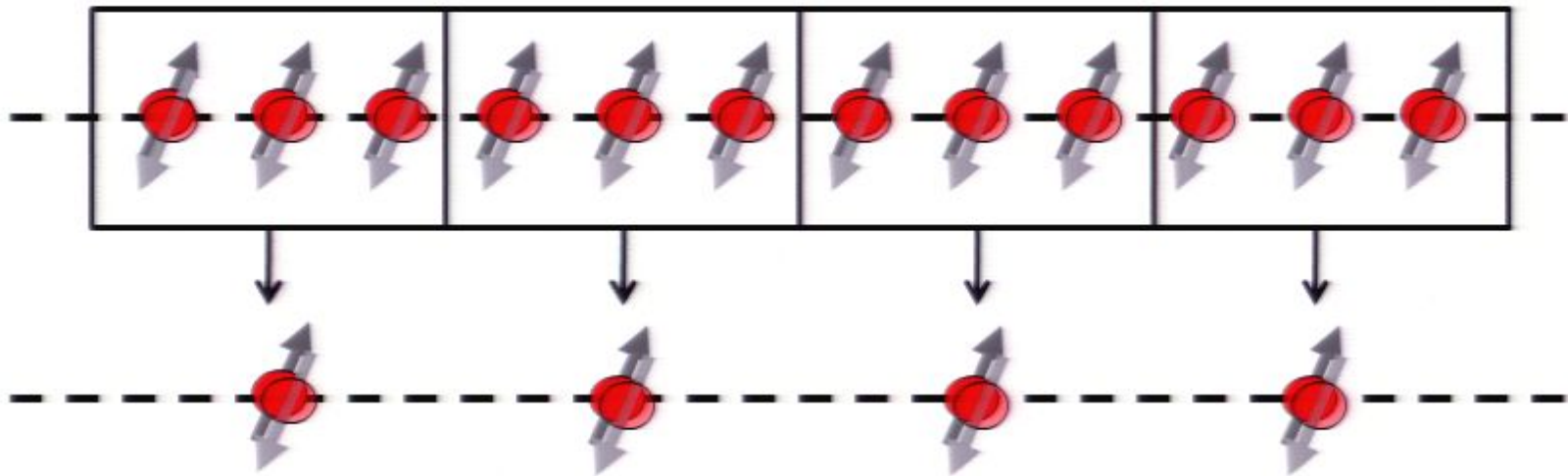
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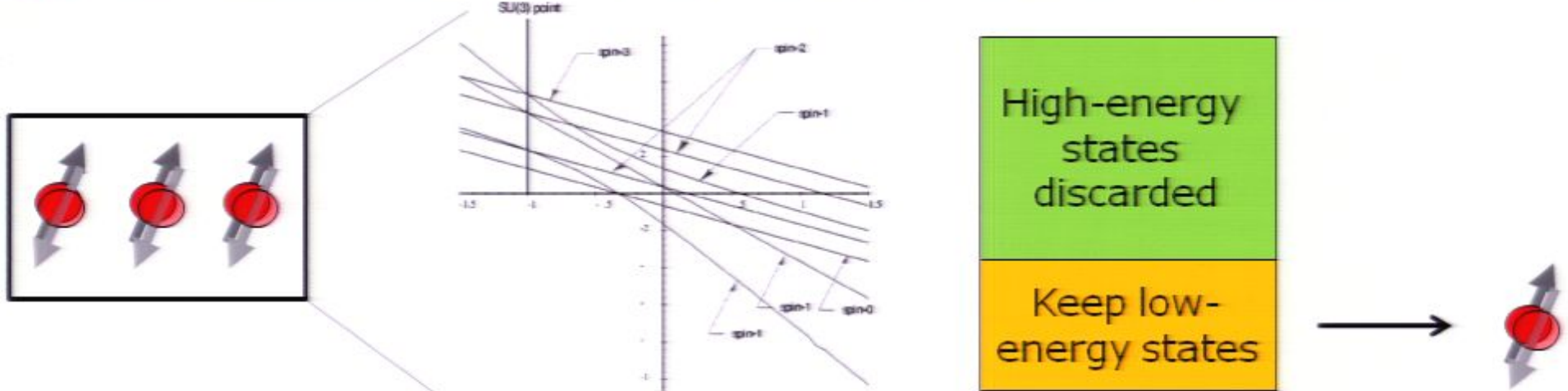
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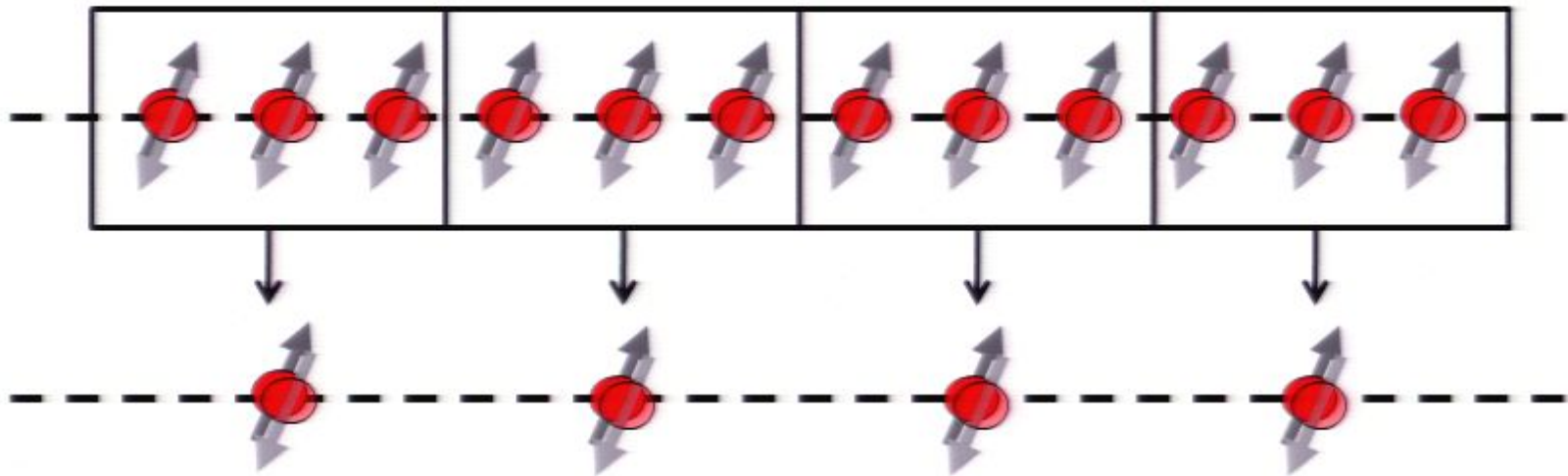
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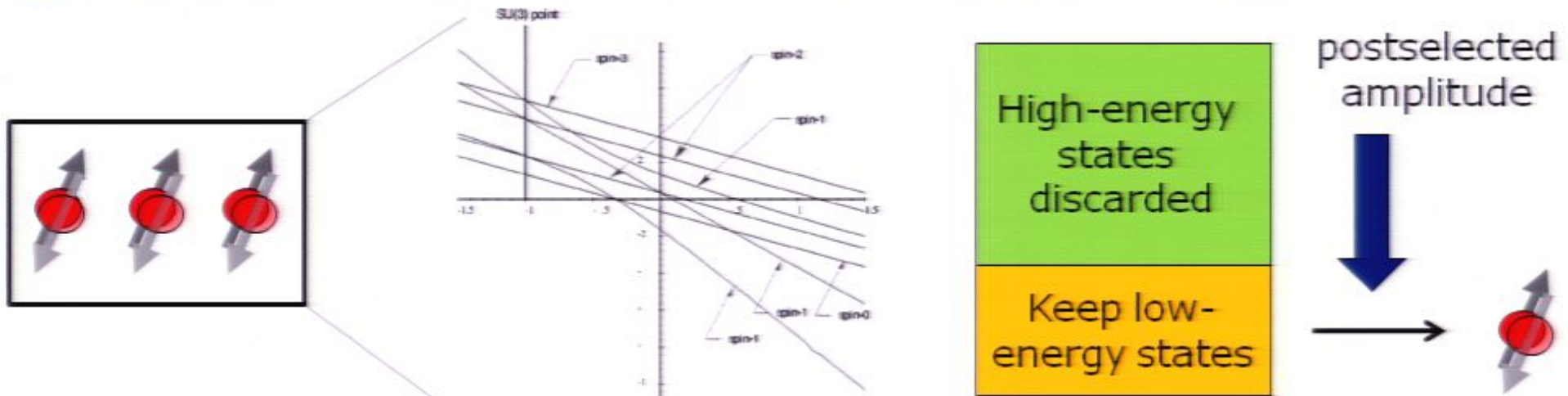
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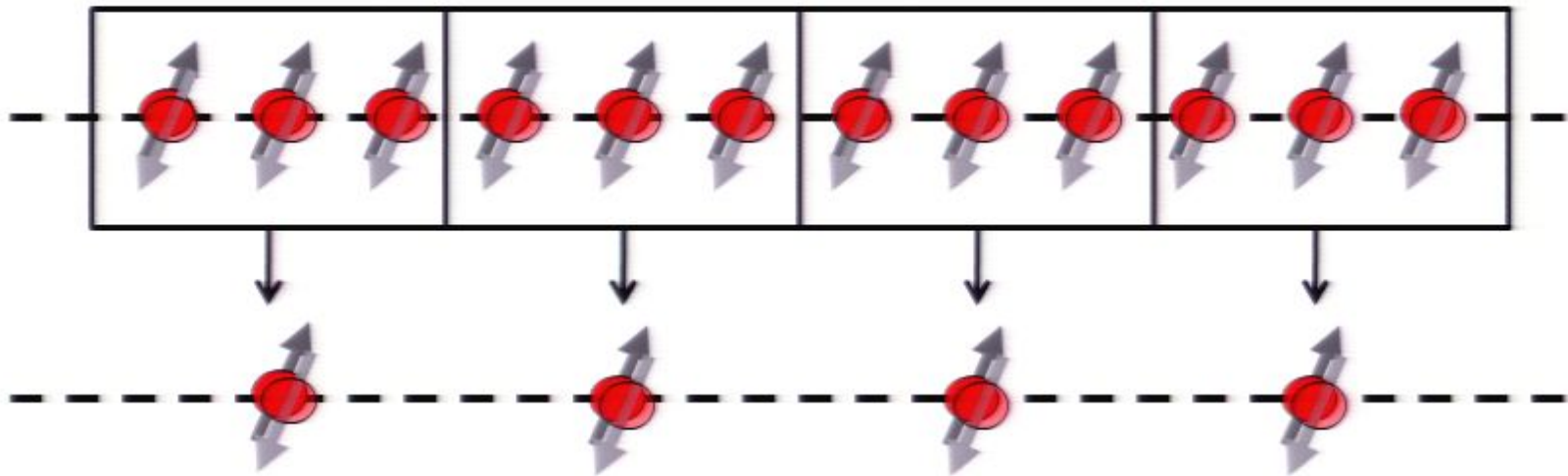
## Conclusions 2.0

- ❑ Quantum gates – correlation functions – order parameters to identify **universal phases** for MBQC
- ❑ MBQC is a new type of long range “string” order
- ❑ Local filtering can **renormalize** the system – **general?**
- ❑ Short-ranged “errors” removed (probabilistically)
- ❑ Long-ranged MBQC order retained

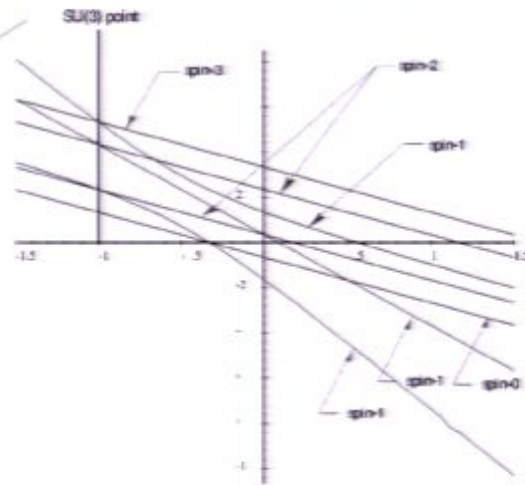
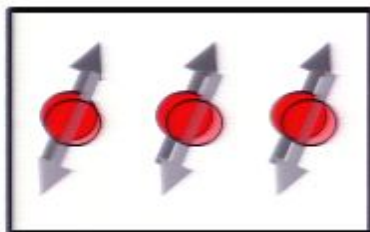
## Outlook

- ❑ Finite temperature
- ❑ Higher dimensions
- ❑ Adiabatic gate teleportation vs MBQC

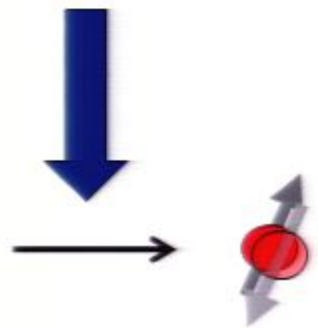
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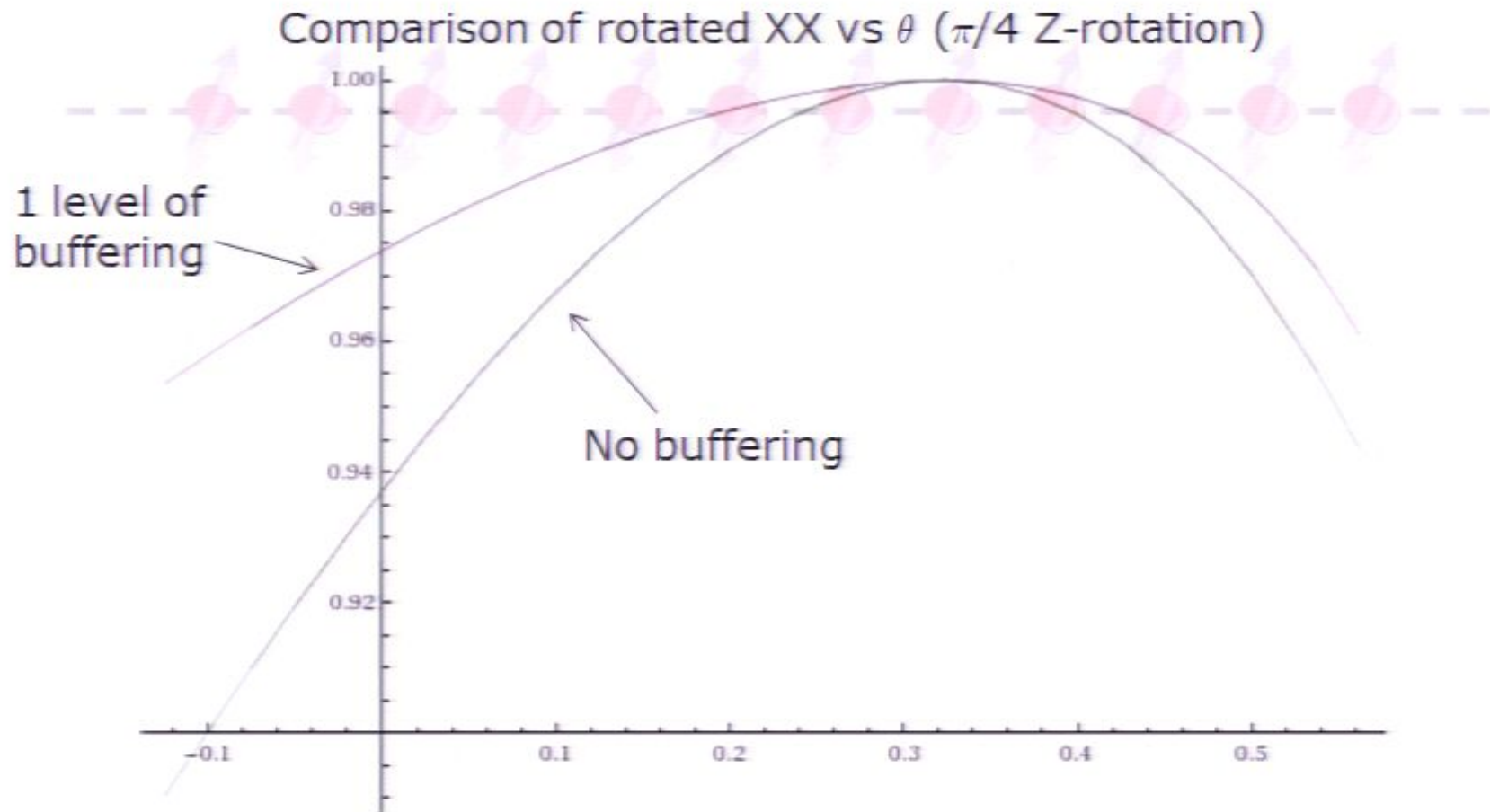


postselected amplitude



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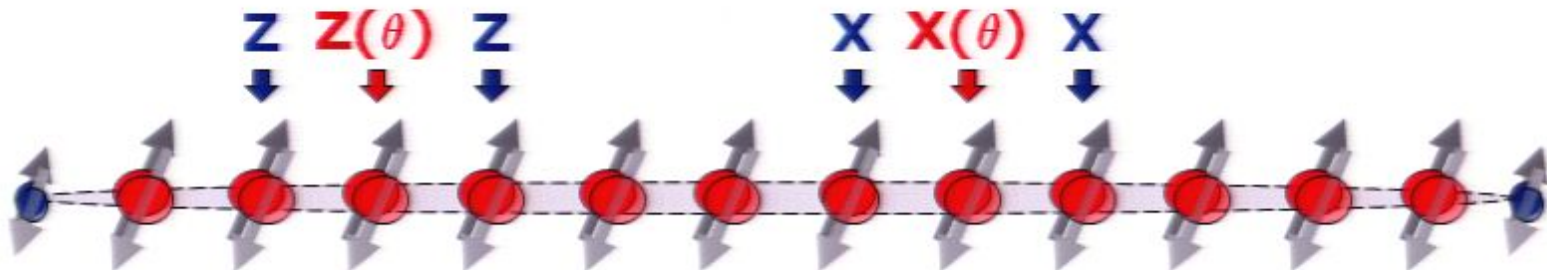
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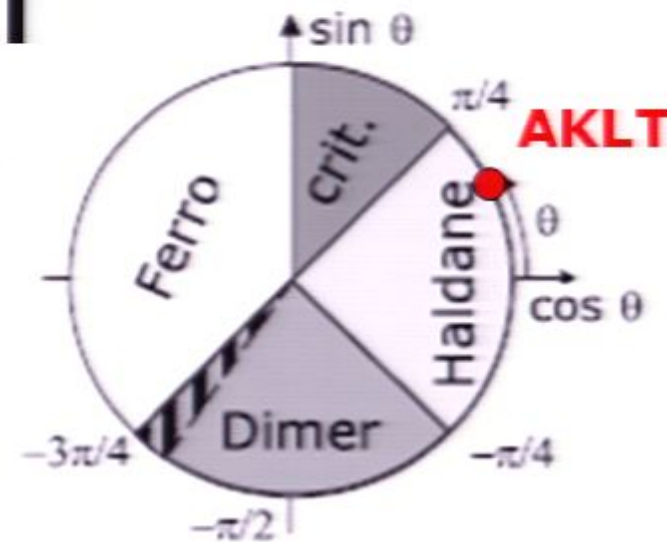
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- Rotation is probabilistic, but heralded. (First outcome: try again.)

# MBQC gates and correlation functions

**Approach:** Quantum gates as correlation functions

- Order parameters to characterise *phases* which are universal for MBQC



$$\rho = \sum_J U_J P_J \rho_0 P_J U_J^\dagger$$

Measurement projector  $P_J$  (indicated by a red arrow pointing to  $P_J$ )  
 Unitary correction  $U_J$  (indicated by a red arrow pointing to  $U_J$ )

$$\langle AB \rangle = \sum_J \text{Tr}[(AB)U_J P_J \rho_0 P_J U_J^\dagger] = \sum_J \text{Tr}[(AB_J)P_J \rho_0 P_J]$$

**Magic:**  $(AB_J)P_J = P_J A S B'$

$$B_J = U_J B U_J^\dagger$$

$$\langle AB \rangle = \text{Tr}[(A S B')\rho_0]$$

# Why are rotations failing?

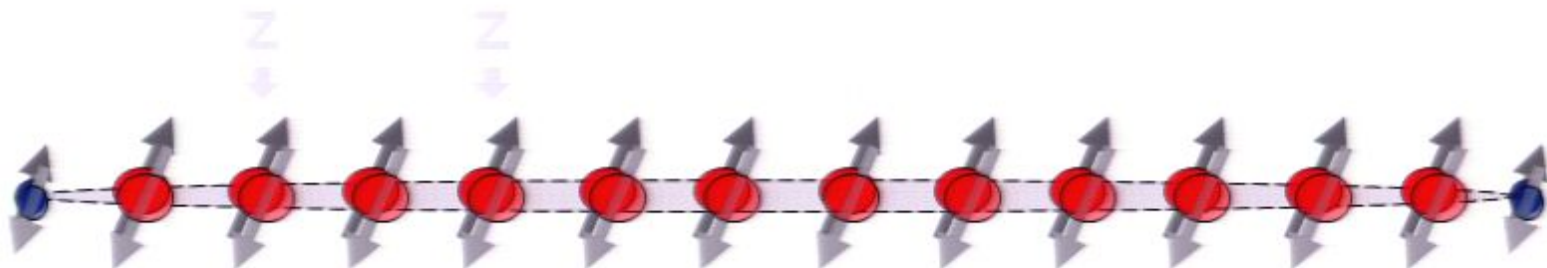
- Re-express the Hamiltonian as

$$H(\theta_{\text{AKLT}} + \varepsilon) = H_{\text{AKLT}} + \frac{5}{6}\varepsilon \sum_{j=1}^n \text{SWAP}_{j,j+1}$$

- Ansatz:

$$|\text{gs}(\theta)\rangle \simeq \alpha|\text{AKLT}\rangle + \sum_{j=1}^n \beta_j \text{SWAP}_{j,j+1}|\text{AKLT}\rangle$$

- Identity gate: SWAP-invariant, only total number of X, Y, Z's matter
- Rotation gates: buffer them against SWAP errors





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