

Title: Introduction to Effective Field Theory - Lecture 7B

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Abstract:

gain  $\left(\frac{m_W}{m_Z}\right)^2 \sim 10^{24} \cdot \left(\frac{E}{m_W}\right)^2 \sim 10^{24}$

2) EFT at  $M_W$ :

$\mathcal{L}_{\text{NR}} \supset \frac{G_F}{\sqrt{2}} (\bar{\nu} \gamma^\mu \gamma_\nu \nu) (\bar{e} \gamma^\mu \gamma_\nu e) + \dots$  all be suppressed by powers of  $M_W$ .

$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_W^2}$



$m_\nu = 0.05 \text{ GeV}$   $m_\nu = 5.11 \times 10^{-4} \text{ GeV}$   $A(\nu\gamma \rightarrow \nu\gamma) \approx \dots$

matching gives:

$$\mathcal{L}_{\text{eff}} = \frac{e(\frac{1}{2} + 2s_W^2)}{40\pi m_\nu^2} \left(\frac{G_F}{\sqrt{2}}\right) \left[ 5 N_{\mu\nu} F^{\mu\nu} (F_{\lambda\rho} F^{\lambda\rho}) - 14 N_{\mu\nu} F^{\nu\lambda} F_{\lambda\rho} F^{\rho\mu} \right]$$

$$N_{\mu\nu} = \partial_\mu (\bar{\nu} \gamma_\nu \gamma_L \nu) - (m \leftrightarrow \nu)$$

→ gives the leading result

$$\int \mathcal{D}Y_{,\mu\nu} \mathcal{D}Y_{,\mu\nu} e^{iS_0(A) + iS_0(\psi) + i \int J_{\mu\nu} A} \left[ \frac{\delta}{\delta Y_{,\mu\nu}} \left( \frac{1}{2} Y_{,\mu\nu} + b Y_{,\mu\nu} \right) \right]$$

but  $\int \mathcal{D}Y_{,\mu\nu} e^{iS_0(A, \psi) + i \int J_{\mu\nu} A} \sim a J_{\mu\nu}$

$$\varphi_{\mu\nu} \sim \frac{\varphi^2}{16\pi m^2} \left[ \left( \frac{\delta}{\delta Y_{,\mu\nu}} F^{\mu\nu} \right)^2 - 4 F_{\mu\nu} F^{\mu\nu} \right]$$

$$\int dY_{\mu\nu} e^{iS_0(A) + iS_0(\psi) + i\int J_{\mu\nu}^{\alpha} \left[ \bar{\psi} \gamma_{\mu} (\psi + b \gamma_5 \psi) \right]}$$

$$a(\bar{\psi} \gamma_{\mu} \psi) + b(\bar{\psi} \gamma_{\mu} \gamma_5 \psi)$$

but  $\int dY_{\mu\nu} e^{iS_0(A, \psi) + i\int J_{\mu\nu}^{\alpha} \sim a J_{\mu\nu}^{\text{em}} + b J_{\mu\nu}^{\text{ax}}$

$$\varphi_{\mu\nu}^2 \sim \frac{\varphi^2}{16\pi m^2} \left[ \left( \frac{1}{\Lambda} F_{\mu\nu} \right)^2 - 4 F_{\mu\nu} \tilde{F}^{\mu\nu} \right] + \mathcal{O}\left(\frac{1}{m^2}\right)$$

$m_W = 80 \text{ GeV}$   $m_Z = 91.1876 \text{ GeV}$   $A(\nu\bar{\nu} \rightarrow \nu\bar{\nu}) \approx \dots$

gives:

$$\frac{e(\frac{1}{2} + 2s_W^2)}{40\pi m_W^2} \left( \frac{G_F}{\sqrt{2}} \right) \left[ 5 N_{\mu\nu} F^{\mu\nu} F_{\lambda\rho} F^{\lambda\rho} - 14 N_{\mu\nu} F^{\nu\lambda} F_{\lambda\rho} F^{\rho\mu} \right]$$

$$N_{\mu\nu} = \partial_\mu (\bar{\nu} \gamma_\nu \nu) - (m \leftrightarrow \nu)$$

→ gives the leading result

$m_W = 80 \text{ GeV}$   $m_Z = 91.1876 \text{ GeV}$   $A(\nu_Y \rightarrow \nu_Y) \approx \dots$

matching gives:

$$\mathcal{L}_H = \frac{e(\frac{1}{2} + 2s_W^2) \alpha}{10\pi m_W^2} \left( \frac{G_F}{\sqrt{2}} \right) \left[ 5 N_{\mu\nu} F^{\mu\nu} F_{\lambda\rho} F^{\lambda\rho} - 14 N_{\mu\nu} F^{\nu\lambda} F_{\lambda\rho} F^{\rho\mu} \right]$$

$$N_{\mu\nu} = \partial_\mu (\bar{\nu}_L \gamma_\nu \nu_L) - (m \leftrightarrow \nu)$$

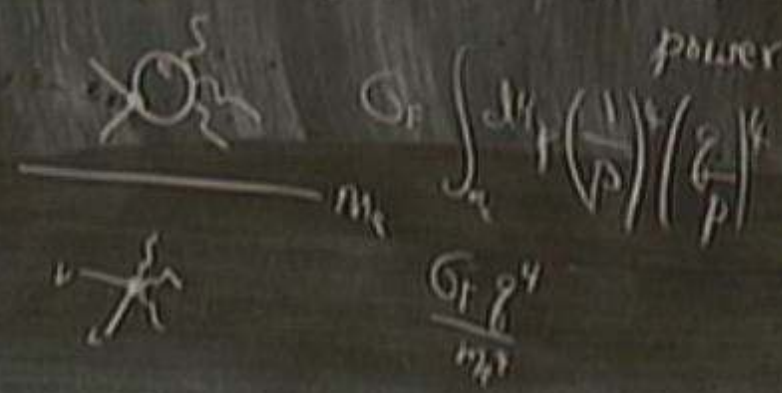
→ gives the leading result

gain  $\left(\frac{m_{\nu}}{M_{Pl}}\right)^2 \sim 10^{-24} \cdot \left(\frac{E}{m_e}\right)^4 \sim 10^{-24}$

2) EFT at  $M_{Pl}$ :

$L_{NR} \supset \frac{G_F}{2\sqrt{2}} (\bar{\nu}_\mu \gamma^\mu \gamma_\nu \nu) (\bar{e} \gamma^\mu \gamma_\mu e) + \dots$  all be suppressed by powers of  $M_{Pl}$ .

$\frac{G_F}{c^2} = \frac{g^2}{8M_W^2}$



$\frac{G_F g^4}{m_\nu^2}$

$$\mathcal{L} = -\frac{1}{2} g_{\mu\nu}(\phi) \partial_\mu \phi^\alpha \partial_\nu \phi^\alpha +$$

$$G \rightarrow H$$

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$$G \rightarrow H$$

$$\mathcal{L} = - (\dot{\phi})^2 \partial_\mu \phi^2 \gamma^\mu \phi^2 +$$

$$\phi^2 + \xi(\phi)$$

$$T_{\mu\nu} \rightarrow T_{\mu\nu} + \delta T_{\mu\nu} + \partial_\mu \xi^\nu T_{\mu\nu} + \partial_\nu \xi^\mu T_{\mu\nu}$$

$$\mathcal{L} = -\frac{1}{2} g_{\mu\nu}(\phi) \partial_\mu \phi^\alpha \partial_\nu \phi^\alpha +$$

$$G \rightarrow H \quad \phi^\alpha \rightarrow \phi^\alpha + \xi^\alpha(\phi)$$

$$J_{\mu\nu} \rightarrow J_{\mu\nu} + \partial_\mu \xi^\alpha \partial_\nu g_{\alpha\beta} + \partial_\nu \xi^\alpha \partial_\mu g_{\alpha\beta} + \partial_\mu \xi^\alpha \partial_\nu g_{\alpha\beta} + \mathcal{L}_\xi g_{\mu\nu}$$

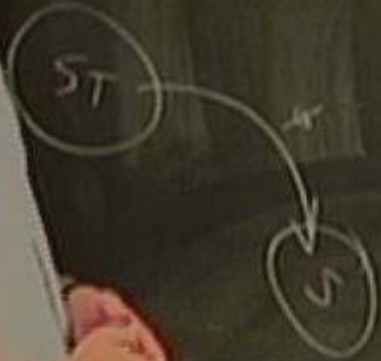
$$\mathcal{L} = -\frac{1}{2} g_{\mu\nu}(\phi) \partial_\mu \phi^\alpha \partial_\nu \phi^\alpha +$$

$$G \rightarrow H \quad \phi^\alpha \rightarrow \phi^\alpha + \xi^\alpha(\phi)$$

$$g_{\mu\nu} \rightarrow g_{\mu\nu} + \partial_\mu \xi^\alpha g_{\alpha\nu} + \partial_\nu \xi^\alpha g_{\mu\alpha} + \partial_\mu \xi^\alpha \partial_\nu \xi^\beta g_{\alpha\beta}$$

$$g_{\mu\nu} + \mathcal{L}_\xi g_{\mu\nu}$$

$$\mathcal{L}_\phi = -\frac{1}{2} g_{\mu\nu}(\phi) \partial_\mu \phi^\alpha \partial_\nu \phi^\alpha +$$



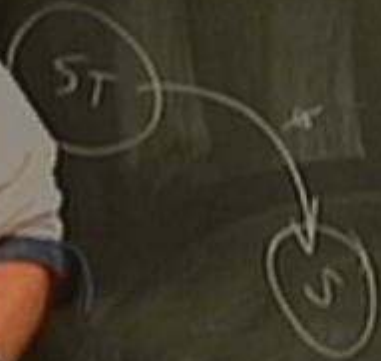
$$G \rightarrow H$$

$$\phi \rightarrow \phi + \xi(\phi)$$

$$G/H$$

$$J_{\mu\nu} \rightarrow J_{\mu\nu} - \partial_\nu \xi^\alpha \partial_\mu g_{\alpha\beta} + \partial_\nu \xi^\alpha g_{\mu\beta} + \partial_\mu \xi^\alpha g_{\alpha\nu} - J_{\mu\nu} + \mathcal{L}_\xi g_{\mu\nu}$$

$$\mathcal{L} = -\frac{1}{2} g_{\mu\nu}(\phi) \partial_\mu \phi^\alpha \partial_\nu \phi^\alpha +$$



$$G \rightarrow H \quad \phi \rightarrow \phi + \xi(\phi)$$

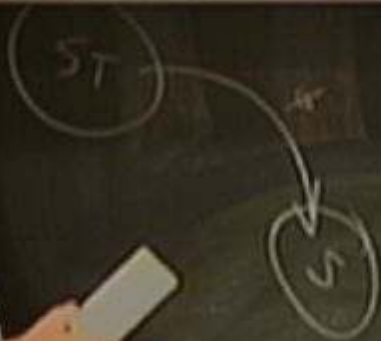
$$g_{\mu\nu} \rightarrow g_{\mu\nu} - \partial_\mu \xi^\alpha \partial_\nu \xi^\alpha + \partial_\mu \xi^\alpha \partial_\nu \xi^\alpha + \partial_\mu \xi^\alpha \partial_\nu \xi^\alpha$$

$$g_{\mu\nu} + \mathcal{L}_\xi g_{\mu\nu}$$

# Goldstone Bosons

SM as an EFT

GR " " "



$$G \rightarrow H$$

$$\phi \rightarrow \phi + \xi(b)$$

$$G/H$$

$$\mathcal{L}_M \rightarrow \mathcal{L}_M + \partial_\mu \xi^a \mathcal{L}_M + \partial_\mu \xi^a \mathcal{L}_M + \partial_\mu \xi^a \mathcal{L}_M$$

$$\mathcal{L}_M + \mathcal{L}_S \mathcal{L}_M$$

# Goldstone Bosons

SM as an EFT

GR " " "

ST

S

$G \rightarrow H$

$G/H$

$$\phi \rightarrow \phi + \xi(\phi)$$

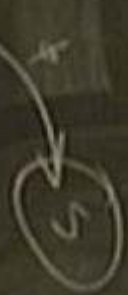
$$J_{\mu\nu} \rightarrow J_{\mu\nu} + \partial_\mu \xi^\nu g_{\mu\nu} + \partial_\nu \xi^\mu g_{\mu\nu} + \partial_\mu \xi^\mu g_{\nu\nu}$$

$$J_{\mu\nu} + \mathcal{L}_S g_{\mu\nu}$$

$$\mathcal{L} = -\frac{1}{2} g_{\mu\nu}(\phi) \partial_\mu \phi^\alpha \partial_\nu \phi^\alpha +$$

$$G \rightarrow H \quad \phi \rightarrow \phi + \xi(\phi)$$

$$g_{\mu\nu} \rightarrow g_{\mu\nu} + \partial_\mu \xi^\alpha g_{\alpha\nu} + \partial_\nu \xi^\alpha g_{\mu\alpha} + \partial_\mu \xi^\alpha \partial_\nu \xi^\beta g_{\alpha\beta} + \mathcal{L}_\xi g_{\mu\nu}$$



$G/H$



Weinberg (for  $SU_2 \times SU_2 / SU_2$ )  
Callan, Coleman, Wess + Zumino (CCWZ formalism)  
~ 1968

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a global symmetry is spontaneously broken, <sup>in a field theory</sup> there  
are Goldstone bosons:

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a global symmetry is spontaneously broken, <sup>in a field theory</sup> there are Goldstone bosons:

$$Q = \int d^3x j^0, \quad \partial_\mu j^\mu = 0 \text{ (conserved)}$$

If a global symmetry is spontaneously broken, <sup>in a field theory</sup> there are Goldstone bosons:

$$\partial_\mu j^{\mu A} = 0 \text{ (conserved)}$$

$$Q = \int d^3x \mathcal{J}^0, \quad \partial_\mu \mathcal{J}^\mu = 0 \text{ (conserved)}$$

Goldstone  
there must be a state  
 $|Q\rangle$  for which  
 $\langle 0 | \mathcal{J}^\mu(x) | Q \rangle \neq 0$

where  $|Q\rangle$  - ground state

$$Q = \int d^3x j^0, \quad \partial_\mu j^\mu = 0 \text{ (conserved)}$$

Goldstone  
 there must be a state  
 $|\phi\rangle$  for which  
 $\langle 0 | j^\mu(x) | \Omega \rangle \neq 0$

where  $|\Omega\rangle$  is the ground state.

$\phi(x)$  for which  $\langle \phi \rangle \neq 0$ .

but  $\phi \rightarrow g\phi$  is a symmetry.

$$\langle \phi \rangle \rightarrow g\langle \phi \rangle \neq \langle \phi \rangle$$

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$$\neq 0 = g(x)\langle \phi \rangle$$



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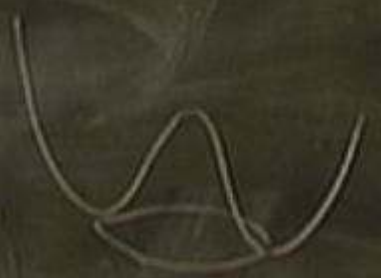
$$\neq \int d^3x \langle \phi \rangle$$

$$Q = \int d^3x f^0, \quad \partial_\mu f^\mu = 0 \text{ (conserved)}$$

Goldstone  
 there must be a state  
 $|G\rangle$  for which  
 $\langle G | \phi(x) | G \rangle \neq 0$

$|G\rangle$  is the ground state.

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$$\neq 0 \quad g(x) \langle \phi \rangle$$

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$$\langle \phi \rangle \rightarrow g \langle \phi \rangle \neq \langle \phi \rangle$$

$$|\phi\rangle = g(x)|\phi\rangle$$



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$$\neq \langle g(x) \phi \rangle$$

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$$\neq g(x) \langle \phi \rangle$$



Consider first a toy model.

$$\phi = \begin{pmatrix} \phi_1 \\ \vdots \\ \phi_n \end{pmatrix} \quad N \text{ real scalars.}$$

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$$\partial_\mu \phi^\dagger \partial^\mu \phi - V(\phi), \quad V(g\phi) = V(\phi) \text{ for all } \phi.$$

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$$\mathcal{L} = -\frac{1}{2} \partial_\mu \phi^T \partial^\mu \phi - V(\phi), \quad V(g\phi) = V(\phi) \text{ for all } g.$$

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2D-EFT at  $M_W$ :

$$\mathcal{L}_{\text{eff}} \supset \frac{G_F}{2\sqrt{2}} (\bar{\nu} \gamma^\mu \gamma_\nu) (\bar{e} \gamma^\mu \gamma_\nu e) + \dots \quad \text{all be suppressed by powers of } M_W.$$

$$\frac{G_F}{2\sqrt{2}} = \frac{g^2}{8M_W^2}$$



$$G_F \int d^4x \left( \frac{1}{\sqrt{2}} \right)^2 \left( \frac{g}{p} \right)^4$$

$$\frac{G_F g^4}{M_W^2}$$

$$\mathcal{L} = \frac{1}{c} \partial_\mu \phi^\top \partial^\mu \phi - V(\phi) \quad , \quad V(g\phi) = V(\phi) \text{ for all } g.$$

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$$g = e^{i\omega^a T_a}$$

$\swarrow$  real  
 $\searrow$  real

$g \leftrightarrow$  orthogonal  
 $\leftrightarrow$  unitary  
 $\leftrightarrow$  real

$T_a \leftrightarrow$  antisymmetric  
 $\leftrightarrow$  hermitian  
 $\leftrightarrow$  imaginary

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$$T_a = -T_a^\top = -T_a^* = T_a^\dagger \leftrightarrow \text{real}$$

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Define  $\mathcal{L}$

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$$g = e^{i\omega^a T_a} \quad \text{real}$$

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$$T_a = -T_a^T = -T_a^\dagger = T_a^T \quad \text{real}$$

Define

$$\chi_f = g(x) \chi(x) \quad \text{where}$$

$$\mathcal{L} = \frac{1}{c} \partial_\mu \phi^\dagger \partial^\mu \phi - V(\phi) \quad , \quad V(g\phi) = V(\phi) \text{ for all } g. \\ \text{and } \phi \neq 0 \text{ at min.}$$

$$g = e^{i\omega^a T_a} \quad \text{real}$$

$g \Leftrightarrow$  orthogonal  $T_a \Leftrightarrow$  antisymmetric  
 $\Leftrightarrow$  unitary  $\Leftrightarrow$  hermitian  
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$$T_a = -T_a^T = -T_a^\dagger = T_a^T \quad \Leftrightarrow \text{real}$$

Define  $\chi_F = g(\omega) \chi(\omega)$

$$\tilde{\mathcal{L}} = \frac{1}{c} \partial_\mu \phi^\dagger \partial^\mu \phi - V(\phi) \quad , \quad V(g\phi) = V(\phi) \text{ for all } g. \\ \text{and } \phi \neq 0 \text{ at min.}$$

$$g = e^{i\omega^\dagger T_a} \quad \text{real}$$

$g \leftrightarrow$  orthogonal  $T_a \leftrightarrow$  antisymmetric  
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$$T_a = -T_a^T = -T_a^\dagger = T_a^T \quad \leftrightarrow \text{real}$$

Define  $\chi_f = g(x) \chi(x)$  where  $g(x) = e^{i\omega^\dagger x T_a}$

Define  $\phi = \psi$  to be the config minimizing  $V(\phi)$ .

Define  $t_i$  to be those  $T_n$  for which  $t_i \psi = 0$   
(ie  $e^{i\omega t_i} \psi = \psi$ )

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$e^{i\omega t_i} \in H \subset G$  is a subgroup of  $G$ .  
(ie  $e^{i\omega t_i} \psi = \psi$ )

Define  $X_a$  to be those  $T_a$ 's for which  $X_a \psi \neq 0$

$\{X_{a_i} t_i\}$

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Define  $X_\alpha$  to be those  $T_a$ 's for which  $X_\alpha \psi \neq 0$ .

$\{T_a\} = \{X_{\alpha_i} t_i\}$   $X_\alpha$ 's label the coset  $G/H$

(momentum conserved)  $\langle G | \phi(x) | Q \rangle \neq 0$

the directions  $X_{\alpha} U$  are the ones obtained  
by a  $G$  transf<sup>n</sup> from the vac., so  $\phi'$ , projected  
in these directions,  $\phi^T X_{\alpha} U$ , are the Goldstone  
directions

the directions  $X_{\alpha} \psi$  are the ones obtained  
 by a  $G$  transf<sup>n</sup> from the vac., so  $\phi'$ , projected  
 in these directions,  $\phi^T X_{\alpha} \psi$ , are the Goldstone  
 directions (and so are contained in  $\mathfrak{g}(x)$  if  $\phi = g(x)\psi$ ),  
 We demand  $X_{\alpha}(x)$  satisfy  $X_{\alpha}^T X_{\beta} \psi = 0$  for all  $\alpha, \beta$  and all  $X_{\alpha}$ .

the directions  $X_{\alpha} \psi$  are the ones obtained  
 by a  $G$  transf<sup>n</sup> from the vac., so  $\phi'$ , projected  
 in these directions,  $\phi^T X_{\alpha} \psi$ , are the Goldstone  
 directions (and so are contained in  $g(x)$  if  $\phi = g(x)\psi$ ),  
 We demand  $X(x)$  satisfy  $X(x)^T X_{\alpha} \psi = 0$  for all  $x$ , and all  $X_{\alpha}$ .

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 We demand  $\chi(x)$  satisfy  $\chi^T X_{\alpha} \psi = 0$  for all  $x$ , and all  $X_{\alpha}$ .

$g = e^{i\omega x}$  the directions  $X_{\omega} \psi$  are the ones obtained  
 $\phi^T g'_{\omega} \psi = f(\omega)$  by a  $G$  transf<sup>n</sup> from the vac., so  $\phi'$ , projected  
 in these directions,  $\phi^T X_{\omega} \psi$ , are the Goldstone  
 directions (and so are contained in  $g(x)$  if  $\phi = g(x)\psi$ ).  
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 in these directions,  $\phi^T X_{\omega} \psi$ , are the Goldstone  
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 We demand  $X(x)$  satisfy  $X(x)^T X_{\omega} \psi = 0$  for all  $x$ , and all  $X_{\omega}$ .

$g = e^{i\omega x}$  the directions  $X_{\omega} \psi$  are the ones obtained  
 $\phi^T g'_{\omega} \psi = f(\omega)$  by a  $G$  transf<sup>n</sup> from the vac., so  $\phi'$ , projected  
 $\frac{\partial}{\partial \omega} = 0$  in these directions,  $\phi^T X_{\omega} \psi$ , are the Goldstone  
 directions (and so are contained in  $g(x)$  if  $\psi = g(x)\psi_0$ ).

We demand  $\chi(x)$  satisfy  $\chi^T(x) X_{\omega} \psi = 0$  for all  $x$ , and all  $X_{\omega}$ .

$$g(x) = e^{i\omega(x)X_a}$$

↳ goldstone fields.

$$g(x) = e^{i\omega^a(x) X_a} \rightarrow \text{goldstone fields.}$$

$$\phi = e^{i\omega^a X_a} \chi$$

$$\text{under symmetry } \phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^a X_a} \tilde{\chi}$$

$$g(x) = e^{i\omega^*(x)X_a} \rightarrow \text{goldstone fields.}$$

$$\phi = e^{i\omega^* X_a} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^* X_a} \tilde{\chi}$

Claim:  $\omega^* \rightarrow \tilde{\omega}^*$  where:  $g e^{i\omega^* X_a} = e^{i\tilde{\omega}^* X_a}$

$$g(x) = e^{i\omega^*(x)X_a} \rightarrow \text{goldstone fields.}$$

$$\phi = e^{i\omega^* X_a} \chi$$

symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^* X_a} \tilde{\chi}$

Claim:  $\omega^* \rightarrow \tilde{\omega}^*$  where:  $g e^{i\omega^* X_a} = e^{i\tilde{\omega}^* X_a} e^{i\omega^* X_a}$

$$\chi \rightarrow \tilde{\chi} = h\chi \quad h \in H$$

$$g(x) = e^{i\omega^*(x) X_a} \rightarrow \text{goldstone fields.}$$

$$\phi = e^{i\omega^* X_a} \chi$$

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$$\chi \rightarrow \tilde{\chi} = h\chi \quad h \in H$$

$$g(x) = e^{i\omega^*(x)X_a} \rightarrow \text{goldstone fields.}$$

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$$\chi \rightarrow \tilde{\chi} = h\chi \quad h \in H$$

$$g(x) = e^{i\omega^*(x)X_{\omega}}$$

↳ goldstone fields.

$$\phi = e^{i\omega^*X_{\omega}} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^*X_{\omega}} \tilde{\chi}$

Claim:  $\omega^* \rightarrow \tilde{\omega}^*$  where:

$$ge^{i\omega^*X_{\omega}} = e^{i\tilde{\omega}^*X_{\omega}} \underbrace{e^{i\omega^*(\omega)\omega}}_{h \in H}$$

$\chi \rightarrow \tilde{\chi} = h\chi$   $h \in H$

$$g(x) = e^{i\omega^*(x)X_a} \rightarrow \text{goldstone fields.}$$

$$\phi = e^{i\omega^* X_a} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^* X_a} \tilde{\chi}$

Claim:  $\omega^* \rightarrow \tilde{\omega}^*$  where:

$$g e^{i\omega^* X_a} = e^{i\tilde{\omega}^* X_a} e^{i\omega^*(x) t}$$

$\chi \rightarrow \tilde{\chi} = h\chi$   $h \in H$

nonlinear  
real



$$g(x) = e^{i\omega^*(x) X_a} \rightarrow \text{goldstone fields.}$$

$$\phi = e^{i\omega^* X_a} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^* X_a} \tilde{\chi}$

Claim:  $\omega^* \rightarrow \tilde{\omega}^*$

where:

$$g e^{i\omega^* X_a} = e^{i\tilde{\omega}^* X_a} e^{i\text{rot}(a) \cdot t}$$

$$\tilde{\chi} \rightarrow \chi = h\chi$$

$$h \in H$$

nonlinear realization of  $G$

if  $g \in H$  then  $e^{i\omega^* X_a} = e^{i\tilde{\omega}^* X_a}$

$$g(x) = e^{i\omega^*(x) X_a} \rightarrow \text{goldstone fields.}$$

$$\tilde{\omega}^* X_a = g^{-1} \omega^* X_a$$

$$\phi = e^{i\omega^* X_a} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^* X_a} \tilde{\chi}$

nonlinear realization

Claim:  $\omega^* \rightarrow \tilde{\omega}^*$  where:

$$g e^{i\omega^* X_a} = e^{i\tilde{\omega}^* X_a} e^{i\omega^*(g^{-1}x) X_a}$$

$\chi \rightarrow \tilde{\chi} = h\chi$   $h \in H$

$G$   
if  $g \in H$   
then

$$g(x) = e^{i\omega^*(x)X_x}$$

one fields.

$$\tilde{\omega}^* X_x = g^{-1} \omega^* X_x g$$

$$\phi = e^{i\omega^* x}$$

under sym

$$\tilde{\phi} = g\phi = e^{i\tilde{\omega}^* X_x} \tilde{X}$$

nonlinear realization of  $G$ .

Claim:

$$e^{i\tilde{\omega}^* X_x} = e^{i\omega^* X_x} e^{i\omega^*(g^{-1}x)} t$$

if  $g \in H$  then  $e^{i\omega^* t} = g$ .

$$\tilde{X} = hX$$

$$h \in H$$

$$g(x) = e^{i\omega^\alpha(x) X_\alpha} \rightarrow \text{goldstone fields.}$$

$$\tilde{\omega}^\alpha = \text{linear plus higher order} \quad \tilde{\omega}^\alpha X_\alpha = g^{-1} \omega^\alpha X_\alpha g$$

inhomogeneous

$$\phi = e^{i\omega^\alpha X_\alpha} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^\alpha X_\alpha} \tilde{\chi}$

nonlinear realization of  $G$ .

Claim:  $\omega^\alpha \rightarrow \tilde{\omega}^\alpha$

where:

$$g e^{i\omega^\alpha X_\alpha} = e^{i\tilde{\omega}^\alpha X_\alpha} e^{i\omega^\alpha X_\alpha} g$$

$$\chi \rightarrow \tilde{\chi} = h\chi$$

$$h \in H$$

if  $g \in H$  then  $e^{i\omega^\alpha X_\alpha} = g$

$$g(x) = e^{i\omega^\alpha(x) X_\alpha}$$

↳ goldstone fields.

$$\tilde{\omega}^\alpha = \omega^\alpha + \dots$$

plus higher order  
inhomogeneous

$$\tilde{\omega}^\alpha X_\alpha = g^{-1} \omega^\alpha X_\alpha g$$

$$\phi = e^{i\omega^\alpha X_\alpha} \chi$$

under symmetry  $\phi \rightarrow \tilde{\phi} = g\phi = e^{i\tilde{\omega}^\alpha X_\alpha} \tilde{\chi}$

nonlinear realization of  $G$ .

Claim:  $\omega^\alpha \rightarrow \tilde{\omega}^\alpha$  where:

$g e^{i\omega^\alpha X_\alpha} = e^{i\tilde{\omega}^\alpha X_\alpha} e^{iU(h)T}$

$\chi \rightarrow \tilde{\chi} = h\chi$        $h \in H$

if  $g \in H$   
then  $e^{iU(h)T} = g$

$$\mathcal{L} = -\frac{1}{2} \partial_\mu \phi^T \partial^\mu \phi - V(\phi) \quad , \quad V(g\phi) = V(\phi) \text{ for all } \phi$$

matrix

and  $\phi \neq 0$  at min.

$$g = e^{i\omega^a T_a}$$

↖ real

G

$g \leftrightarrow$  orthogonal  $T_a \leftrightarrow$  antisymmetric  
 $\leftrightarrow$  unitary  $\leftrightarrow$  hermitian

$$T_a = -T_a^T = -T_a^\dagger = T_a^T$$

↖ real

$\leftrightarrow$  imaginary

Define  $\mathcal{L}_g = g(x) \mathcal{L}(x)$  where  $g(x) = e^{i\omega^a(x) T_a}$