Title: Mapping classical fields to quantum states

Date: Oct 01, 2009 04:00 PM

URL: http://pirsa.org/09100093

Abstract: Efforts to extrapolate non-relativistic (NR) quantum mechanics to a covariant framework encounter well-known problems, implying that an alternate view of quantum states might be more compatible with relativity. This talk will reverse the usual extrapolation, and examine the NR limit of a real, classical scalar field. A complex scalar \psi that obeys the Schrodinger equation naturally falls out of the analysis. One can also replace the usual operator-based measurement theory with classical measurement theory on the scalar field, and examine the NR limit of this as well. In this limit, the local energy density of the field scales as \psi \psi \p2, adding credibility to this approach. With the added postulate that " all measurements correspond to boundary conditions that extremize the classical action" (see arXiv:0906.5409), additional quantitative comparisons emerge between this \psi and the standard quantum wavefunction. Uncertainty can then be introduced (along with a " collapse" due to Bayesian updating) by simply giving the classical scalar field two components instead of one, leading to an intriguing \psi-epistemic model.

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- Easily extrapolatable to curved spacetime

$$\mathcal{L} = \frac{1}{2} \left(g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi - \frac{m^2 c^2}{\hbar^2} \phi^2 \right)$$

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 Euler-Lagrange equation in Minkowski spacetime is Klein-Gordon Equation (KGE);

$$\left(\hbar^2 \frac{\partial^2}{\partial t^2} - \hbar^2 c^2 \nabla^2 + m^2 c^4\right) \phi = 0$$

(or, why ϕ isn't complex)

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Define:
$$\omega(k) = \sqrt{c^2 k^2 + m^2 c^4 / \hbar^2}$$
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Solutions to complex KGE:

$$\phi(\boldsymbol{x},t) = \int \left[a(\boldsymbol{k})e^{i(\boldsymbol{k}\cdot\boldsymbol{x}-\omega t)} + b(\boldsymbol{k})e^{i(\boldsymbol{k}\cdot\boldsymbol{x}+\omega t)} \right] d\boldsymbol{k}$$

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Single particle Schr. Eqn:
$$\psi(\boldsymbol{x},t) = \int \left[\Psi(k)e^{i[\boldsymbol{k}\cdot\boldsymbol{x}-(\omega-\omega_0)t]}\right]d\boldsymbol{k}$$

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In Minkowski spacetime, NR limit, one easily sees

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- Adding scalar potential can be done without adding a new term; just adjust g_{μν} in weak-field limit.

(WKB solutions match S.E.'s in NR limit)

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 $\psi_R = Re(\psi)$ $\frac{\partial}{\partial x^i} \Rightarrow \partial_i$ $\psi_I = Im(\psi)$

 $T_{0i} \approx (\psi_I \partial_i \psi_R) \cos^2(\omega_o t) - (\psi_R \partial_i \psi_I) \sin^2(\omega_o t) - 2Re(\psi \partial_i \psi) \sin(2\omega_0 t)$

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Lack of a conserved 4-current for the real KGE does not mean that it can't contain NRQM.

What if action extremization constrains measurements?

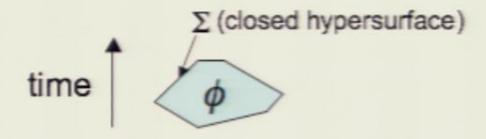
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$$\text{time} \qquad \qquad \sum \text{(closed hypersurface)}$$

$$\delta S = \int [\Box + m^2 c^2/\hbar^2] \phi \, \delta \phi \, d^4 \Omega + \oint_{\Sigma} \partial_{\mu} \phi \, \delta \phi \, d^{\mu} \Sigma$$

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Unlike QFT, constraints on ϕ can't be expressed via ϕ

Define "non-relativistic" field via stress energy tensor:

$$\frac{T_{0i}}{T_{00}} = \frac{v_i}{c} \ll 1$$

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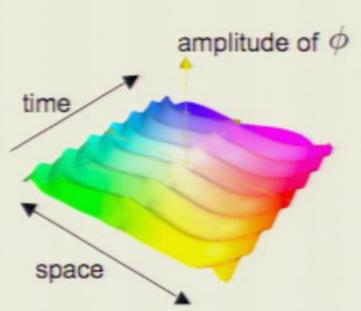
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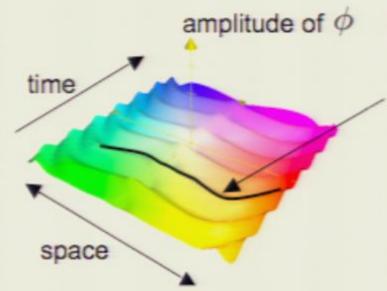
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Hypersurface requirements: Very near crest or trough of ϕ

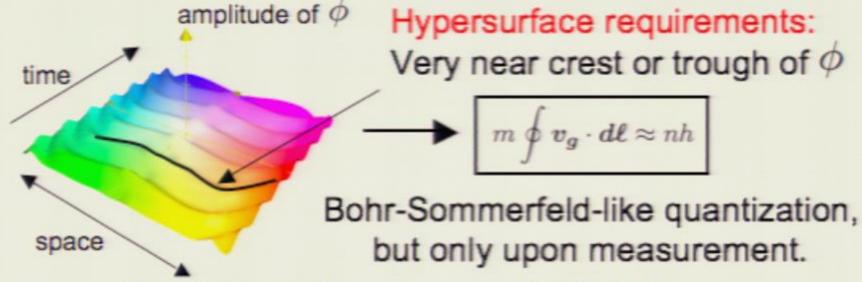
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arXiv:0906.5409; pirsa.org/09060031

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 How to "whittle down" acceptable boundary conditions for the case of only a few particles?
- Whence Probabilistic Measurement Outcomes?
 (not to mention apparent collapse, contextuality, Bell-inequality violations, etc., etc.)

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Due to nonlinearity, ψ_{meas} doesn't solve the S.E.!

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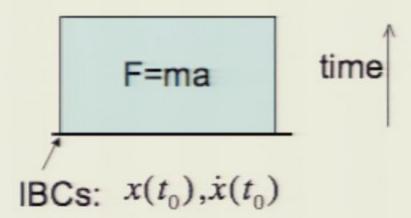
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Action Extremization!

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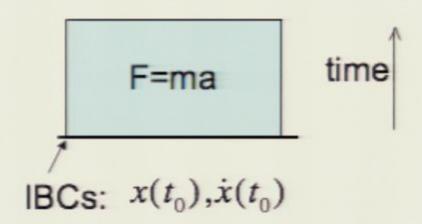
Action Extremization!

Newtonian Schema

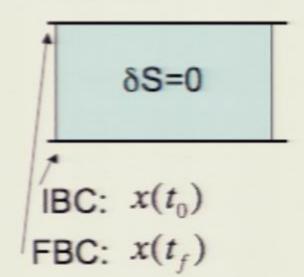


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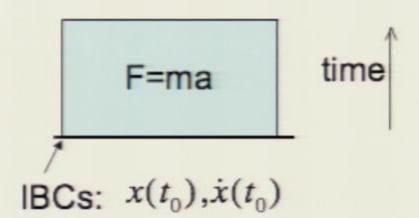


"Lagrangian Schema"

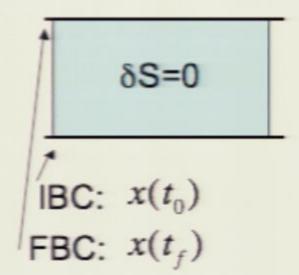


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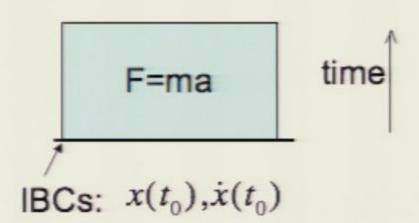


These pictures are not equivalent! (Same laws, different BCs)

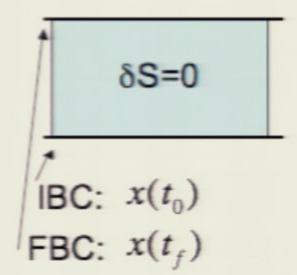
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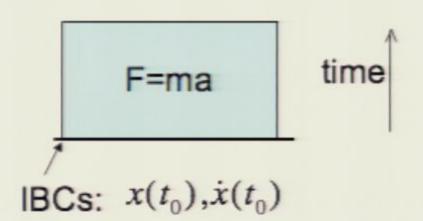


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BCs of physical systems must be physical constraints.

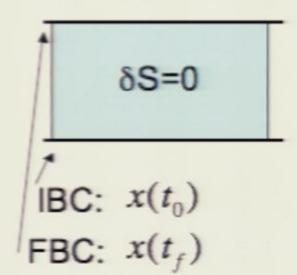
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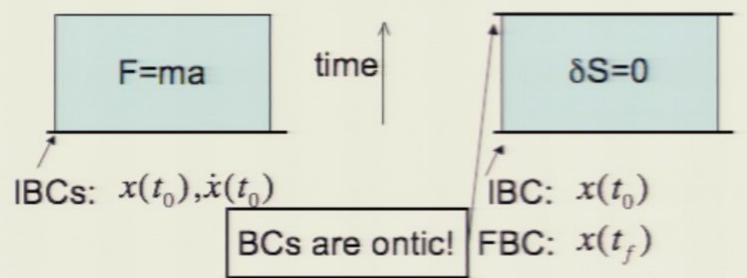
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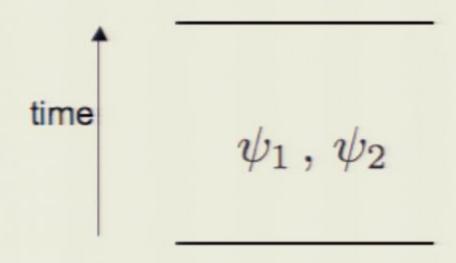
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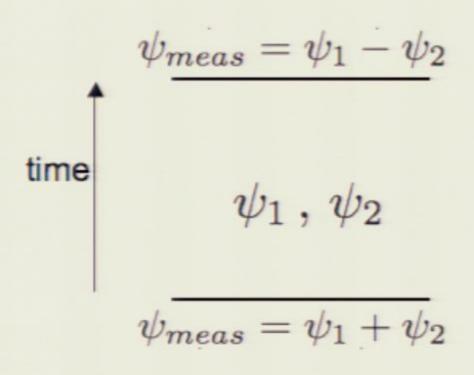


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time $\psi_{meas}=\psi_2$ $\psi_1\,,\,\psi_2$ $\psi_{meas}=\psi_1$

1) Aharanov/Vaidman (1991)

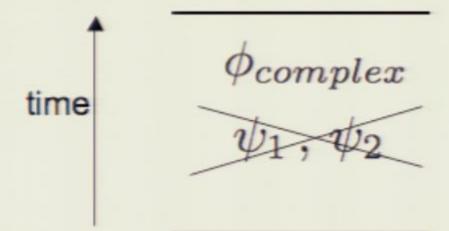
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- 1) Aharanov/Vaidman (1991)
- Wharton, Found. Phys. v37, p.159 (2007)

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Same rule as below



- 1) Aharanov/Vaidman (1991)
- Wharton, Found. Phys. v37, p.159 (2007)
- 3) Wharton, arXiv:0706.4075

Time-even measurements:

$$\psi_{meas} = \phi_{complex}$$

Time-odd measurements:

$$\psi_{meas} \Leftrightarrow \dot{\phi}_{complex}$$

Same rule as below

 ψ_1, ψ_2

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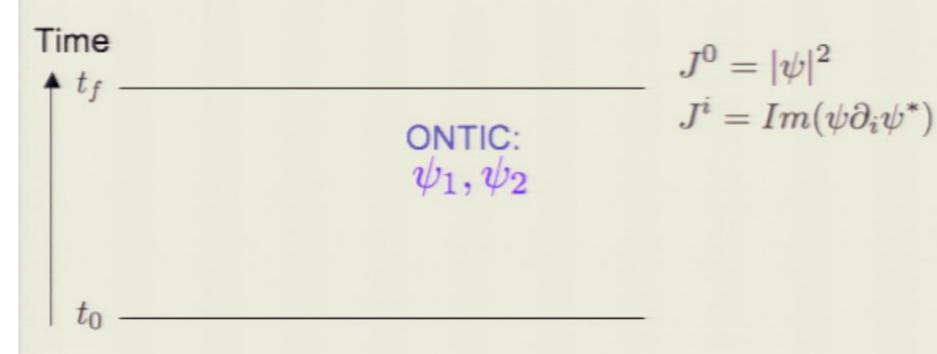
4) This Talk

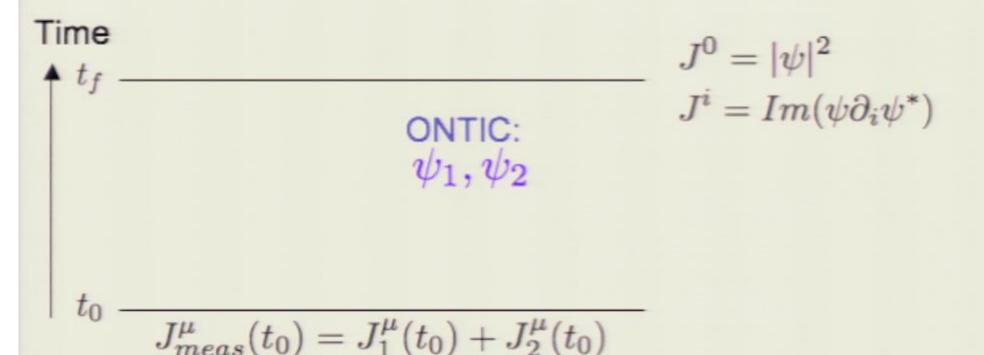
time

Time

ontic: ψ_1, ψ_2

 t_0





Time

 t_f

EPISTEMIC:

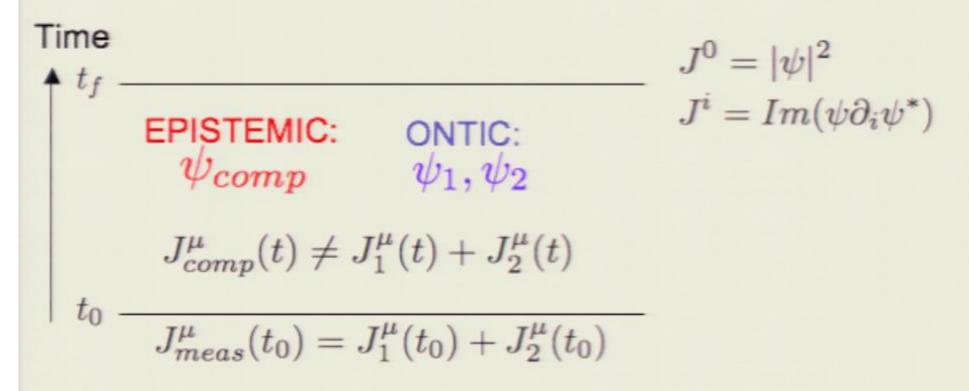
 ψ_{comp}

ONTIC:

 ψ_1, ψ_2

$$J^0 = |\psi|^2$$
$$J^i = Im(\psi \partial_i \psi^*)$$

$$J_{meas}^{\mu}(t_0) = J_1^{\mu}(t_0) + J_2^{\mu}(t_0)$$



Time
$$J_{meas}^{\mu}(t_f) = J_1^{\mu}(t_f) + J_2^{\mu}(t_f)$$
 $J^0 = |\psi|^2$ $J^i = Im(\psi \partial_i \psi^*)$ $J_{comp}^{\mu}(t) \neq J_1^{\mu}(t) + J_2^{\mu}(t)$ $J_{meas}^{\mu}(t_0) = J_1^{\mu}(t_0) + J_2^{\mu}(t_0)$

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Type of measurement at $t = t_f$ (and result) determine:

- parameters that are constrained in action extremization...
- one in turn determine ontic field values. (Natural contextuality!) Page 84/96

Time
$$J_{meas}^{\mu}(t_f) = J_1^{\mu}(t_f) + J_2^{\mu}(t_f)$$

EPISTEMIC: ONTIC: ψ_{comp} ψ_1, ψ_2
 $J_{comp}^{\mu}(t) \neq J_1^{\mu}(t) + J_2^{\mu}(t)$
 t_0
 $J_{meas}^{\mu}(t_0) = J_1^{\mu}(t_0) + J_2^{\mu}(t_0)$

$$J^0 = |\psi|^2$$
$$J^i = Im(\psi \partial_i \psi^*)$$

Open question:

What algorithm can one use to reconstruct ψ_1, ψ_2 from $J^{\mu}_{meas}(t_0), J^{\mu}_{meas}(t_f)$?

Type of measurement at $t = t_f$ (and result) determine:

- parameters that are constrained in action extremization...
- in turn determine ontic field values. (Natural contextuality!) Page 85/96

Pirsa: 09100093 Page 86/96

- Still need a probability rule
 - Should be joint probability distribution

$$P[\psi_{meas}(t_0), \psi_{meas}(t_f)]$$

Candidates:

$$P = |e^{iS}|^2$$
 $P = [\Delta F(\psi_1, \psi_2)]^2$

Pirsa: 09100093 Page 87/96

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- Need other (identical) particles
 - Don't expand into configuration space (Ontic state needn't bother encoding unperformed measurements)

Pirsa: 09100093 Page 88/96

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- Testable Experimental Consequences?
 - Key parameter: accurate time duration between preparation, measurement.

Pirsa: 09100093 Page 90/96

1) GR is a valuable clue: we should consider it!

Pirsa: 09100093 Page 91/96

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Pirsa: 09100093 Page 92/96

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Pirsa: 09100093 Page 93/96

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- 4) We need a consistent, spacetime view of measurements and boundary conditions. (Ideally without operators!)

Pirsa: 09100093 Page 94/96

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