

Title: General Relativity on Trial: Gravitational Waves and the Parameterized Post-Einsteinian Framework

Date: Oct 08, 2009 01:40 PM

URL: <http://pirsa.org/09100082>

Abstract: With the imminent detection of gravitational waves by ground-based interferometers, such as LIGO, VIRGO and TAMA, pulsar timing observations, and proposed space-borne detectors, such as LISA, we must ask ourselves: how much do we trust general relativity? The confirmation of general relativity through Solar System experiments and binary pulsar observations has proved its validity in the weak-field, where velocities are small and gravity is weak, but no such tests exist in the strong, dynamical regime, precisely the regime of most interest to gravitational wave observations. Unfortunately, because of their inherent feebleness, the extraction of gravitational waves from detector noise relies heavily on the technique of matched filtering, where one constructs waveform filters or templates to clean the data. Currently, all such waveforms are constructed with the implicit assumption that general relativity is correct both in the weak and strong, dynamical regimes. Such an assumption constitutes a fundamental bias that will introduce a systematic error in the detection and parameter estimation of signals, and in turn can lead to a mischaracterization of the universe through incorrect inferences about source event rates and populations. In this talk, I will define this bias, explain its possible consequences and propose a remedy through a new scheme: the parameterized post-Einsteinian framework. In this framework one enhances waveforms via the inclusion of post-Einsteinian parameters that both interpolate between general relativity and well-motivated alternative theories, but also extrapolate to unknown theories, following sound theoretical principles, such as consistency with conservation laws and symmetries. The parameterized post-Einsteinian framework should allow matched filtered data to select a specific set of post-Einsteinian parameters without *a priori* assuming the validity of the former, thus allowing the data to either verify general relativity or point to possible dynamical strong-field deviations.

# the Parameterized Post-Quantum Cryptography

Nico Yunes  
Princeton University

October '09,  
Perimeter Institute

(with Frans Preforius, arxiv:0909.3636)

# General Relativity on Trial

The Case for Gravitational Waves and  
the Parameterized Post-Einsteinian Framework

Nico Yunes  
Princeton University

October '09,  
Perimeter Institute

(with Frans Pretorius, arxiv:0909.3636)

# General Relativity on Trial

The Case for Gravitational Waves and  
the Parameterized Post-Einsteinian Framework

Nico Yunes  
Princeton University

October '09,  
Perimeter Institute

(with Frans Pretorius, arxiv:0909.3636)

# The Evidence

Solar System Experiments: light deflection, perihelion precession, frame-dragging (LAGEOS, spacecraft tracking)

# The Evidence

**Solar System Experiments:** light deflection, perihelion precession, frame-dragging (LAGEOS, spacecraft tracking)

**Binary Pulsar Observations:** Quadrupolar emission, frame-dragging (Hulse-Taylor, double binary pulsar)

# The Evidence

Solar System Experiments: light deflection, perihelion precession, frame-dragging ( Lageos, spacecraft tracking)

Binary Pulsar Observations: Quadrupolar emission, frame-dragging (Hulse-Taylor, double binary pulsar)

Other Tests: WEP tests (laboratory), EM Lorentz Invariance and preferred frames (Michelson-Morley, etc.)

# The Evidence

**Solar System Experiments:** light deflection, perihelion precession, frame-dragging (LAGEOS, spacecraft tracking)

**Binary Pulsar Observations:** Quadrupolar emission, frame-dragging (Hulse-Taylor, double binary pulsar)

**Other Tests:** WEP tests (laboratory), EM Lorentz Invariance and preferred frames (Michelson-Morley, etc.)

**No Dynamical and Strong Field tests of GR !**

# The Evidence

**Solar System Experiments:** light deflection, perihelion precession, frame-dragging (LAGEOS, spacecraft tracking)

**Binary Pulsar Observations:** Quadrupolar emission, frame-dragging (Hulse-Taylor, double binary pulsar)

**Other Tests:** WEP tests (laboratory), EM Lorentz Invariance and preferred frames (Michelson-Morley, etc.)

**No Dynamical and Strong Field tests of GR !**

**Cosmological Tests:** Well...we need dark energy and dark matter...but otherwise we are doing great! (WMAP, SN)

# What is a Strong/Dynamical Source?

Binary Black Hole Late Inspiral, Merger and Ringdown

# What is a Strong/Dynamical Source?

Binary Black Hole Late Inspiral, Merger and Ringdown

Yes. Gravitational waves are red/blue  
and black holes are cyan in GR

But is GR  
right here?



# What is a Strong/Dynamical Source?

Binary Black Hole Late Inspiral, Merger and Ringdown

Yes. Gravitational waves are red/blue  
and black holes are cyan in GR

But is GR  
right here?

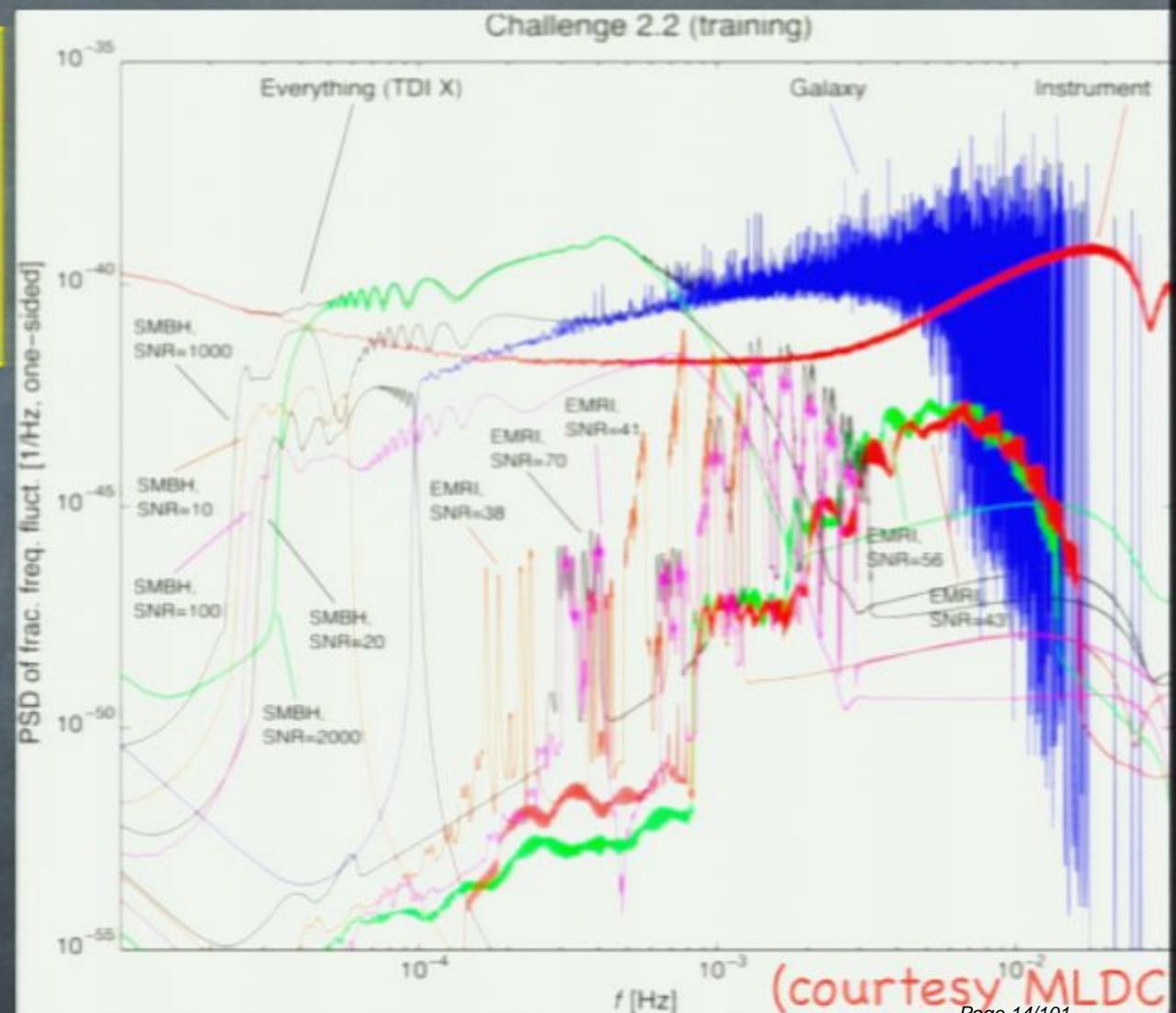


# Does it matter whether GR is right?

YES! You will see  
only what you  
expect to see

# Does it matter whether GR is right?

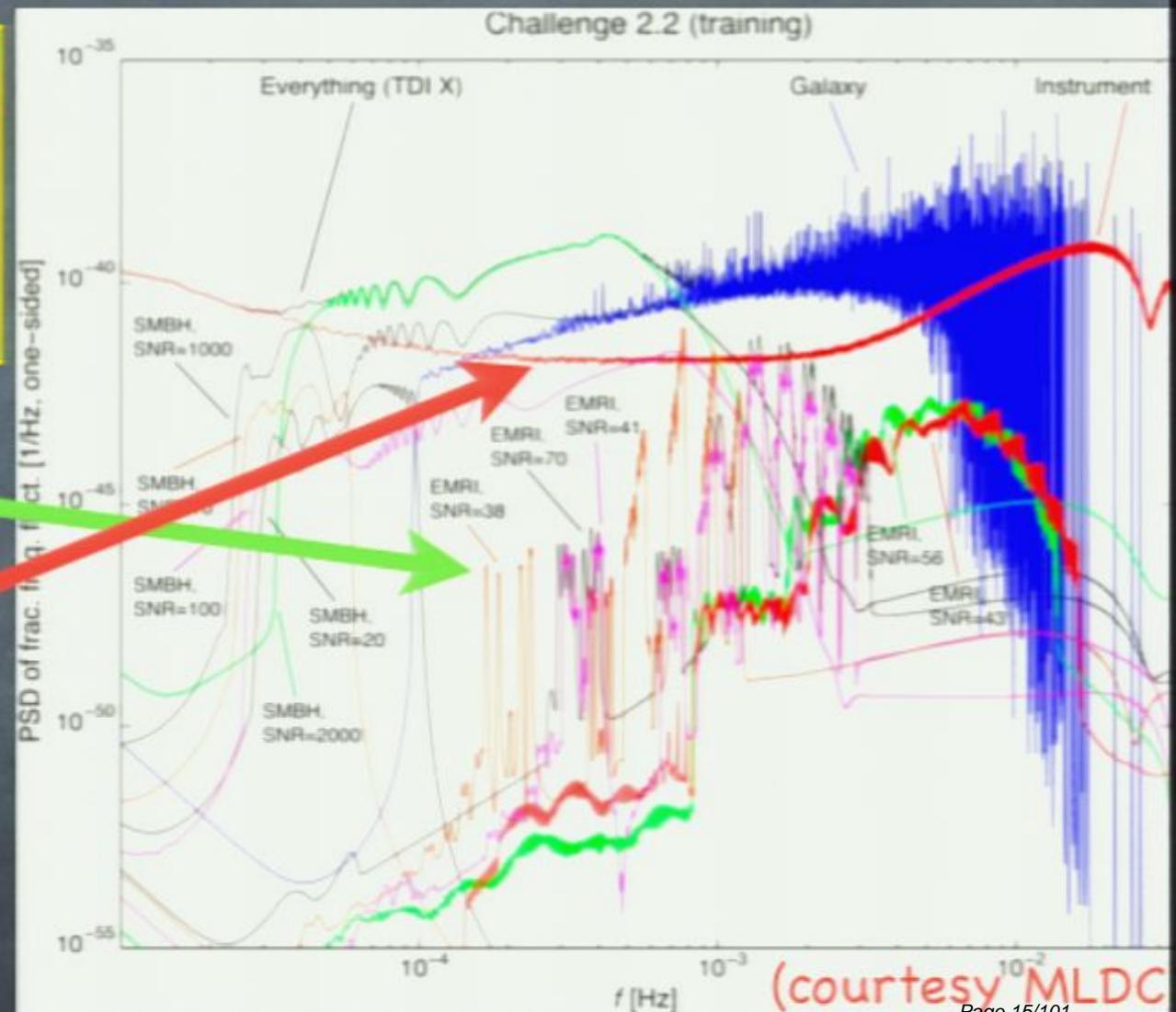
YES! You will see  
only what you  
expect to see



# Does it matter whether GR is right?

YES! You will see  
only what you  
expect to see

A lot of sources  
buried well under  
the noise curve

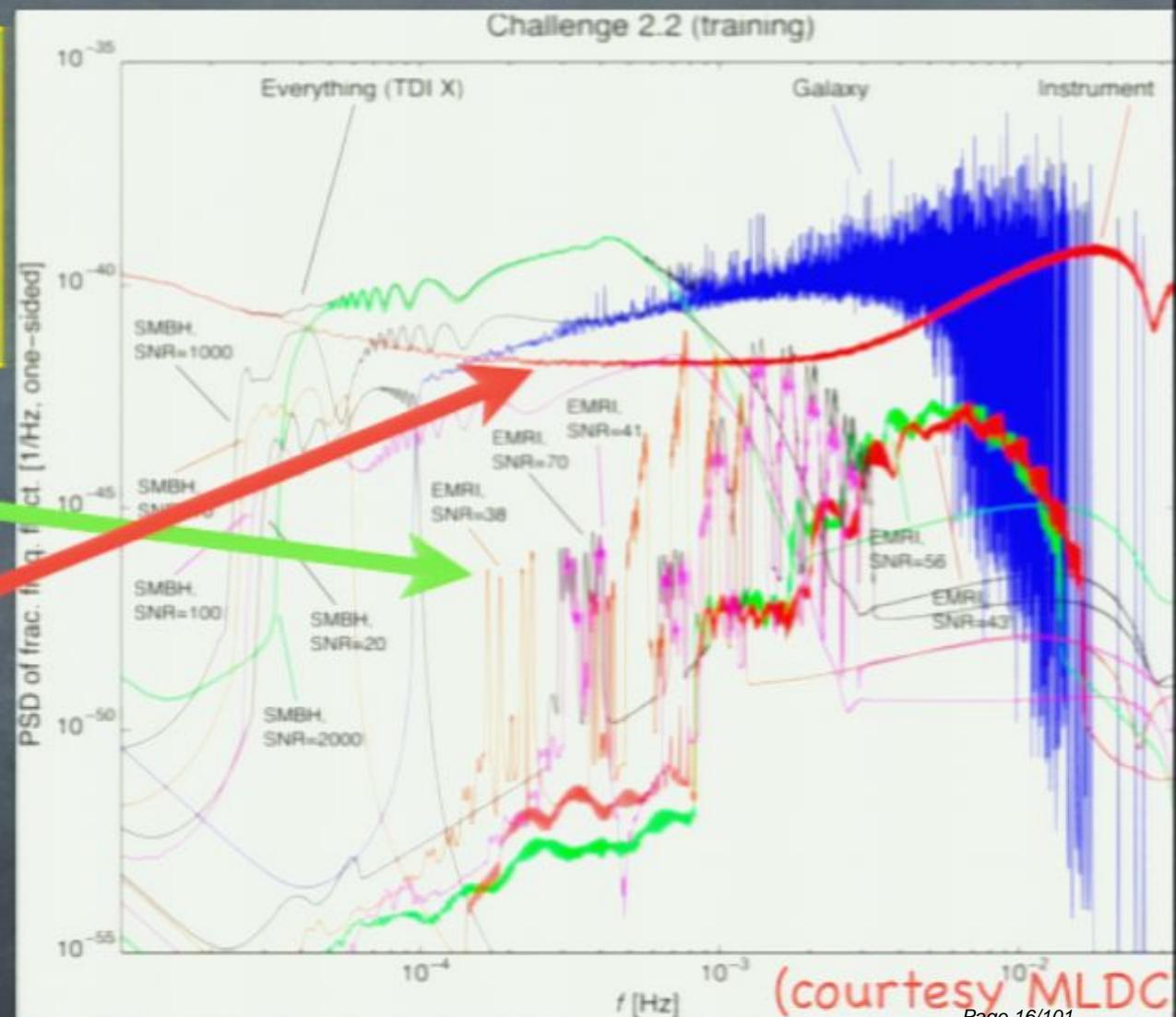


# Does it matter whether GR is right?

YES! You will see  
only what you  
expect to see

A lot of sources  
buried well under  
the noise curve

(Maximize cross-  
correlation between a  
template and data as a  
function of sys. par.)



# So what if we use the wrong filters?

Wrong filters could:

- 1) Filter out wild observations.
- 2) Induce a systematic error in parameter estimation.

# So what if we use the wrong filters?

Wrong filters could:

- 1) Filter out wild observations.
- 2) Induce a systematic error in parameter estimation.

A wild observation is saying: "Here is something from which we may learn a lesson, perhaps of a kind not anticipated beforehand and perhaps more important than the main object of the study." (Kruskal)

# So what if we use the wrong filters?

Wrong filters could:

- 1) Filter out wild observations.
- 2) Induce a systematic error in parameter estimation.

A wild observation is saying: "Here is something from which we may learn a lesson, perhaps of a kind not anticipated beforehand and perhaps more important than the main object of the study." (Kruskal)

This is bad:

For Cosmology: Wild observations could inform Dark Energy models. Do we need to modify GR in the infrared?

# So what if we use the wrong filters?

Wrong filters could:

- 1) Filter out wild observations.
- 2) Induce a systematic error in parameter estimation.

A wild observation is saying: "Here is something from which we may learn a lesson, perhaps of a kind not anticipated beforehand and perhaps more important than the main object of the study." (Kruskal)

This is bad:

For Cosmology: Wild observations could inform Dark Energy models. Do we need to modify GR in the infrared?

# Road Map

- I. A brief review of gravitational wave astrophysics
- II. Alternative Theories? Who Ordered That ?!
- III. Fundamental Bias and the ppE framework

## Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

---

GWs := name we give to oscillatory perturbations  
in the gravitational field or metric tensor.

---

## Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

↑      ↑      ↑  
metric   Minkowski   tensor

---

GWs := name we give to oscillatory perturbations  
in the gravitational field or metric tensor.

---

## Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + G_{\mu\nu} = 8\pi T_{\mu\nu}$$

↑      ↑      ↑  
metric   Minkowski   GW metric pert.

---

GWs := name we give to oscillatory perturbations  
in the gravitational field or metric tensor.

---

## Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + G_{\mu\nu} = 8\pi T_{\mu\nu} \rightarrow$$

GW metric pert.

metric tensor      Minkowski      GW metric pert.

Einstein tensor      Stress-Energy tensor

---

GWs := name we give to oscillatory perturbations  
in the gravitational field or metric tensor.

---

## Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + \square_\eta h_{\mu\nu} = \tau_{\mu\nu}[h^2]$$

metric tensor      Minkowski      GW metric pert.

Einstein tensor      Stress-Energy tensor

---

GWs := name we give to oscillatory perturbations  
in the gravitational field or metric tensor.

---

# Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + G_{\mu\nu} = 8\pi T_{\mu\nu} \rightarrow \square_\eta h_{\mu\nu} = \tau_{\mu\nu}[h^2]$$

metric tensor      Minkowski      GW metric pert.

Einstein tensor      Stress-Energy tensor

Flat-space, diff. wave op.

Annoying non-linearities

GWs := name we give to oscillatory perturbations in the gravitational field or metric tensor.

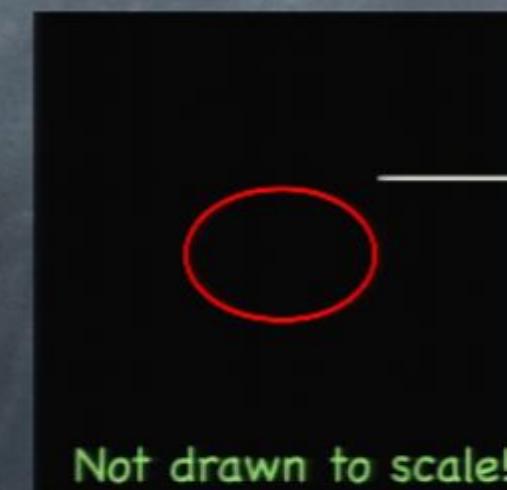
# Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + \square_\eta h_{\mu\nu} = \tau_{\mu\nu}[h^2]$$

metric tensor      Minkowski      GW metric pert.      Einstein tensor      Stress-Energy tensor      Flat-space, diff. wave op.      Annoying non-linearities

GWs := name we give to oscillatory perturbations in the gravitational field or metric tensor.

Given a ring of test-particles, the plus polarization induces:



Change in length is about 1e-18 meters (for a 4km interferometer and a typical GW). That's about 1/1000 nucleus diameter

# Ia. What are Gravitational Waves?

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + G_{\mu\nu} = 8\pi T_{\mu\nu} \rightarrow \square_\eta h_{\mu\nu} = \tau_{\mu\nu}[h^2]$$

metric tensor      Minkowski      GW metric pert.

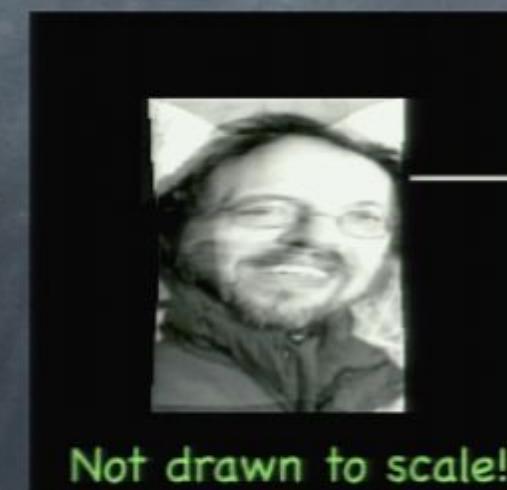
Einstein tensor      Stress-Energy tensor

Flat-space, diff. wave op.

Annoying non-linearities

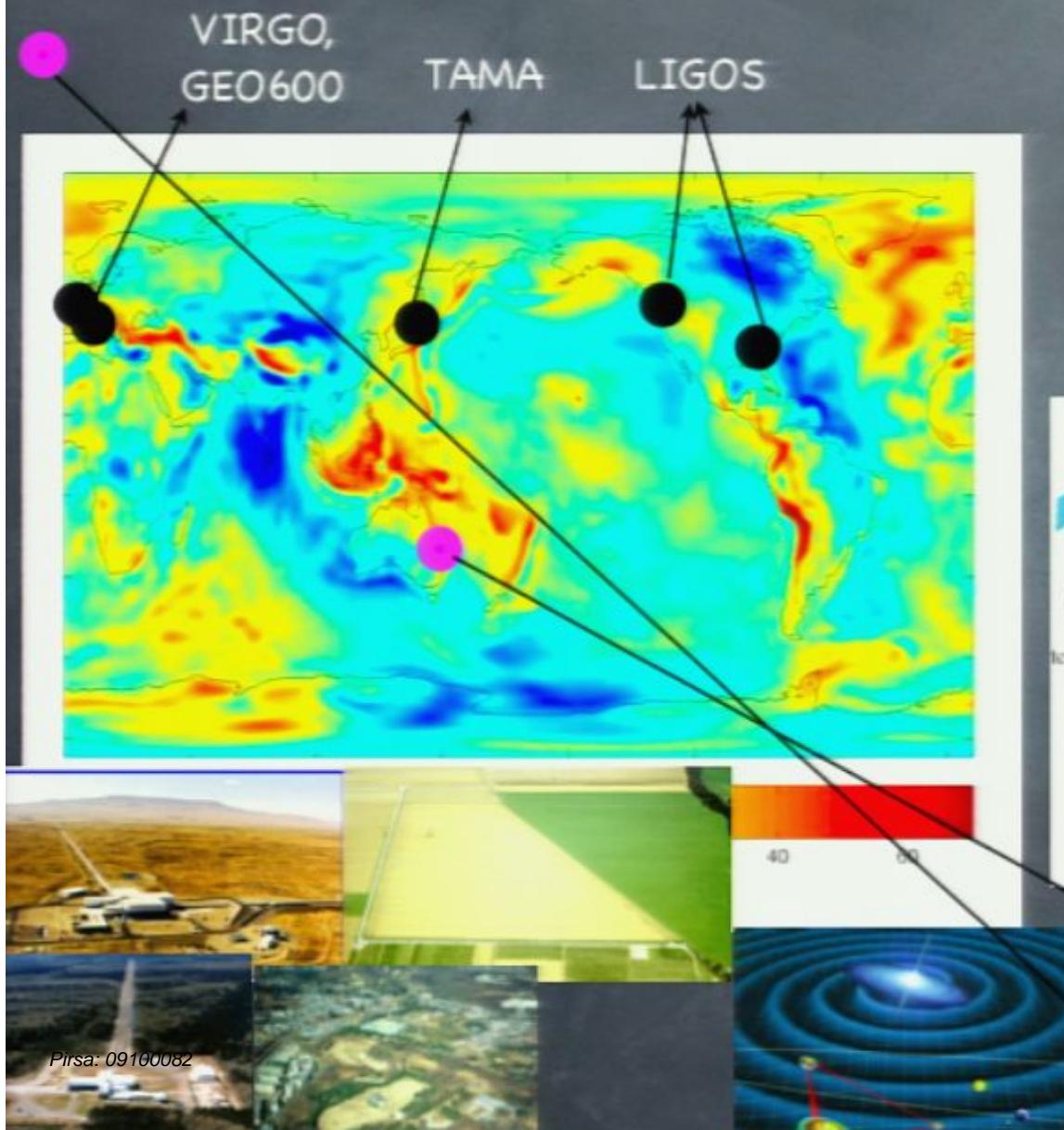
GWs := name we give to oscillatory perturbations in the gravitational field or metric tensor.

So for example, let's shoot Lee with a GW gun face on...

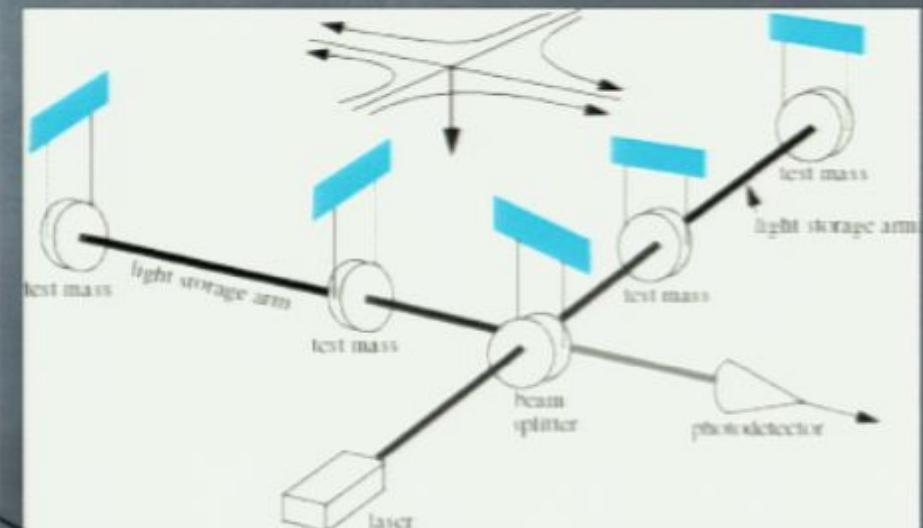


Change in length is about  $1e-18$  meters (for a 4km interferometer and a typical GW). That's about 1/1000 nucleus diameter

## Ib. How do the detectors work?



Bounce light off mirrors  
and look for interference  
pattern when the light  
recombines.



New Australian  
Detector ?

LISA

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f)\tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

signal-to-noise ratio →

$$\rho^2 \sim \int \frac{\tilde{d}(f)\tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

data

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\text{signal-to-noise ratio} \rightarrow \rho^2 \sim \int \frac{\tilde{d}(f)\tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

data                      a measure of detector noise  
(spectral noise density)

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template (projection of GW metric perturbation) →  $\tilde{h}(f; \lambda^\mu)$

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template parameters that characterize system →  $\tilde{h}(f; \lambda^\mu)$

template (projection of GW metric perturbation) →  $\tilde{d}(f) \tilde{h}(f; \lambda^\mu)$

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template parameters that characterize system →  $\tilde{h}(f; \lambda^\mu)$

template (projection of GW metric perturbation) →  $\tilde{h}(f; \lambda^\mu)$

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

(1) If  $h(f)$  is very wrong, you will filter out the signal.

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template parameters that characterize system →  $\tilde{h}(f; \lambda^\mu)$

template (projection of GW metric perturbation) →  $\tilde{h}(f; \lambda^\mu)$

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

(1) If  $h(f)$  is very wrong, you will filter out the signal.

(2) If  $h(f)$  is slightly wrong, you get wrong best-fit pars.

## IIa. Why would the Waveform be wrong?

- (1) Because we use approximations to solve the Einstein Equations -> Modeling error.
- 
-

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f)\tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f)\tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  
data → a measure of detector noise  
(spectral noise density)

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template (projection of GW metric perturbation) →  $\tilde{h}(f; \lambda^\mu)$

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template parameters that characterize system →  $\tilde{h}(f; \lambda^\mu)$

template (projection of GW metric perturbation) →  $\tilde{h}(f; \lambda^\mu)$

---

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

---

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template parameters that characterize system →  $\tilde{h}(f; \lambda^\mu)$

template (projection of GW metric perturbation) →  $\tilde{h}(f; \lambda^\mu)$

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

(1) If  $h(f)$  is very wrong, you will filter out the signal.

## Ic. How does the Analysis Work?

$$\rho^2 \sim \int \frac{\tilde{d}(f) \tilde{h}(f; \lambda^\mu)}{S_n(f)} df$$

signal-to-noise ratio →  $\rho^2$

data →  $\tilde{d}(f)$

a measure of detector noise (spectral noise density) →  $S_n(f)$

template parameters that characterize system →  $\tilde{h}(f; \lambda^\mu)$

template (projection of GW metric perturbation) →  $\tilde{d}(f) \tilde{h}(f; \lambda^\mu) / S_n(f)$

Matched Filtering: Maximize the signal-to-noise ratio over all template parameters  $\lambda^\mu$

- (1) If  $h(f)$  is very wrong, you will filter out the signal.
- (2) If  $h(f)$  is slightly wrong, you get wrong best-fit pars.

## IIa. Why would the Waveform be wrong?

- (1) Because we use approximations to solve the Einstein Equations -> Modeling error.
- 
-

## IIa. Why would the Waveform be wrong?

- (1) Because we use approximations to solve the Einstein Equations -> Modeling error.
- (2) Because we use the Einstein Equations, instead of the M-Equations -> Fundamental bias.

## IIa. Why would the Waveform be wrong?

- (1) Because we use approximations to solve the Einstein Equations -> Modeling error.
- (2) Because we use the Einstein Equations, instead of the M-Equations -> Fundamental bias.

It's not easy to fool Mother Nature! (Wald)

## IIa. Why would the Waveform be wrong?

- (1) Because we use approximations to solve the Einstein Equations -> Modeling error.
- (2) Because we use the Einstein Equations, instead of the M-Equations -> Fundamental bias.

It's not easy to fool Mother Nature! (Wald)

A Minimal (?) Set of Criteria:

1. Weak-Field Consistency (existence and stability of physical solutions, satisfaction of precision tests).
2. Strong-Field Inconsistency (deviations only where experiments cannot currently rule out modifications)
3. Metric theory (gravity is curved spacetime, controlled by metric tensor that is WEP consistent).

## IIb. What theories satisfy that ?!

Some examples include:

- 
- (i) Scalar-Tensor theories: Promote Newton's constant to a scalar field,  $G = \phi(x^\mu)$  and give it dynamics. Passes weak-field tests for non-dynamical coupling. Spontaneous scalarization can lead to strong deviation in NS inspirals and mergers.
-

## IIb. What theories satisfy that ?!

Some examples include:

- 
- (i) Scalar-Tensor theories: Promote Newton's constant to a scalar field,  $G = \phi(x^\mu)$  and give it dynamics. Passes weak-field tests for non-dynamical coupling. Spontaneous scalarization can lead to strong deviation in NS inspirals and mergers.
  - (ii) Chern-Simons Modified Gravity: Enhance the action with the Pontryagin density  $\phi(x^\mu)R \wedge R$ . Passes weak-field tests. It could lead to strong-field deviations (parity violation?).
-

## IIb. What theories satisfy that ?!

Some examples include:

- 
- (i) Scalar-Tensor theories: Promote Newton's constant to a scalar field,  $G = \phi(x^\mu)$  and give it dynamics. Passes weak-field tests for non-dynamical coupling. Spontaneous scalarization can lead to strong deviation in NS inspirals and mergers.
  - (ii) Chern-Simons Modified Gravity: Enhance the action with the Pontryagin density  $\phi(x^\mu)R \wedge R$ . Passes weak-field tests. It could lead to strong-field deviations (parity violation?).
  - (iii) Gauss-Bonnet Modified Gravity: Enhance the action with the Gauss-Bonnet term  $\phi(x^\mu)(a_1R^2 + a_2R_{\mu\nu}R^{\mu\nu} + a_3R_{\mu\nu\delta\gamma}R^{\mu\nu\delta\gamma})$

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

## IIb. What theories satisfy that ?!

Some examples include:

- 
- (i) Scalar-Tensor theories: Promote Newton's constant to a scalar field,  $G = \phi(x^\mu)$  and give it dynamics. Passes weak-field tests for non-dynamical coupling. Spontaneous scalarization can lead to strong deviation in NS inspirals and mergers.
  - (ii) Chern-Simons Modified Gravity: Enhance the action with the Pontryagin density  $\phi(x^\mu)R \wedge R$ . Passes weak-field tests. It could lead to strong-field deviations (parity violation?).
  - (iii) Gauss-Bonnet Modified Gravity: Enhance the action with the Gauss-Bonnet term  $\phi(x^\mu)(a_1 R^2 + a_2 R_{\mu\nu} R^{\mu\nu} + a_3 R_{\mu\nu\delta\gamma} R^{\mu\nu\delta\gamma})$

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

$$h(t) = F_+ h_+ + F_\times h_\times$$

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

response  
function

$$h(t) = F_+ h_+ + F_\times h_\times$$

Beam-pattern functions (depend  
on geometry of detector only)

---

---

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

response  
function

$$h(t) = F_+ h_+ + F_\times h_\times$$

Beam-pattern functions (depend  
on geometry of detector only)

plus and cross  
polarizations of  $h_{\mu\nu}$

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

$$h(t) = F_+ h_+ + F_\times h_\times$$

response function →

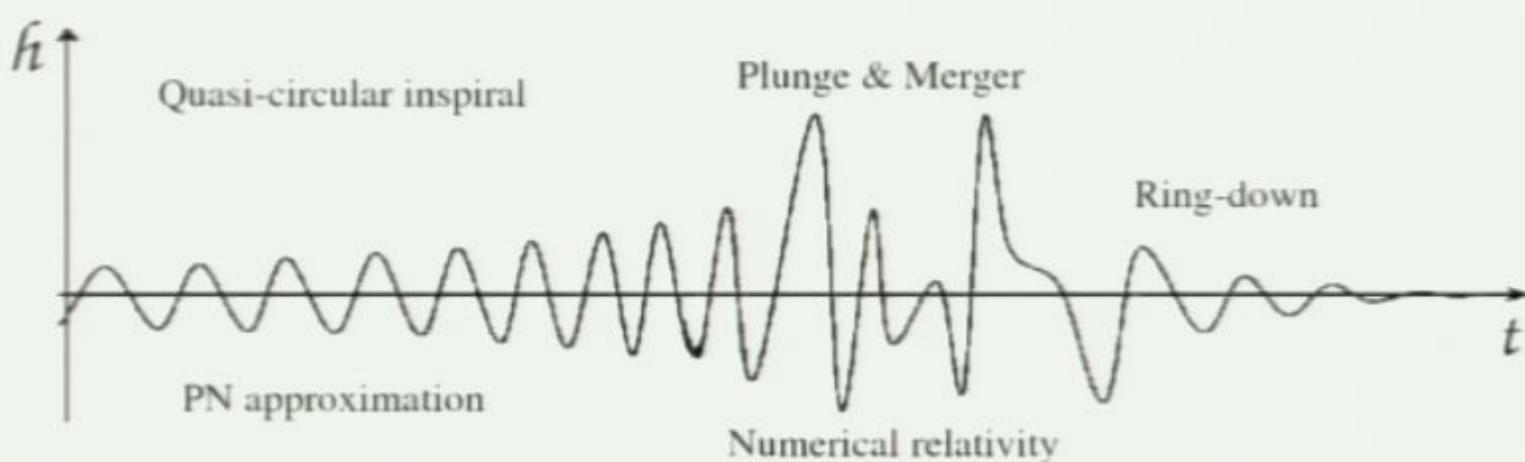
Beam-pattern functions (depend on geometry of detector only)

plus and cross polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---



## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

response  
function

$$h(t) = F_+ h_+ + F_\times h_\times$$

Beam-pattern functions (depend  
on geometry of detector only)

plus and cross  
polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

response  
function

$$h(t) = F_+ h_+ + F_\times h_\times$$

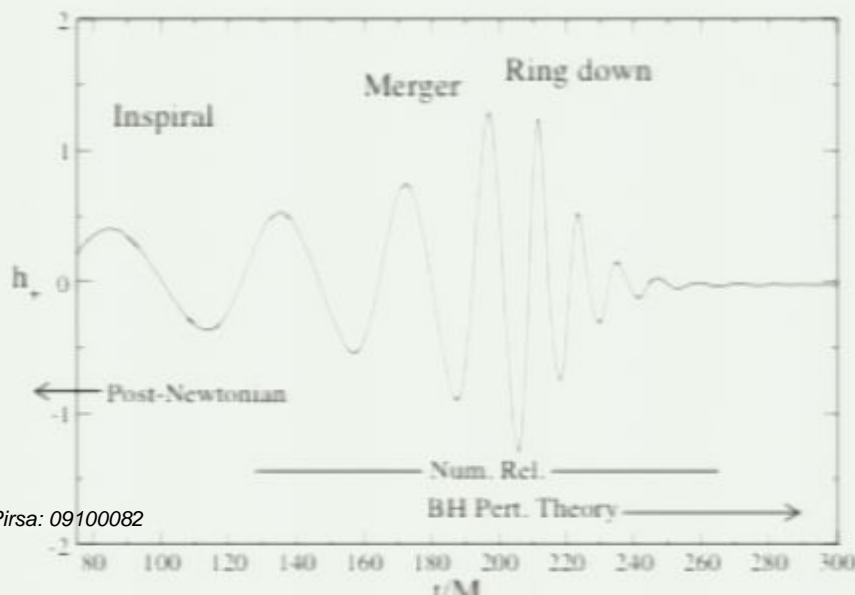
Beam-pattern functions (depend  
on geometry of detector only)

plus and cross  
polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---



## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

response function

$$h(t) = F_+ h_+ + F_\times h_\times$$

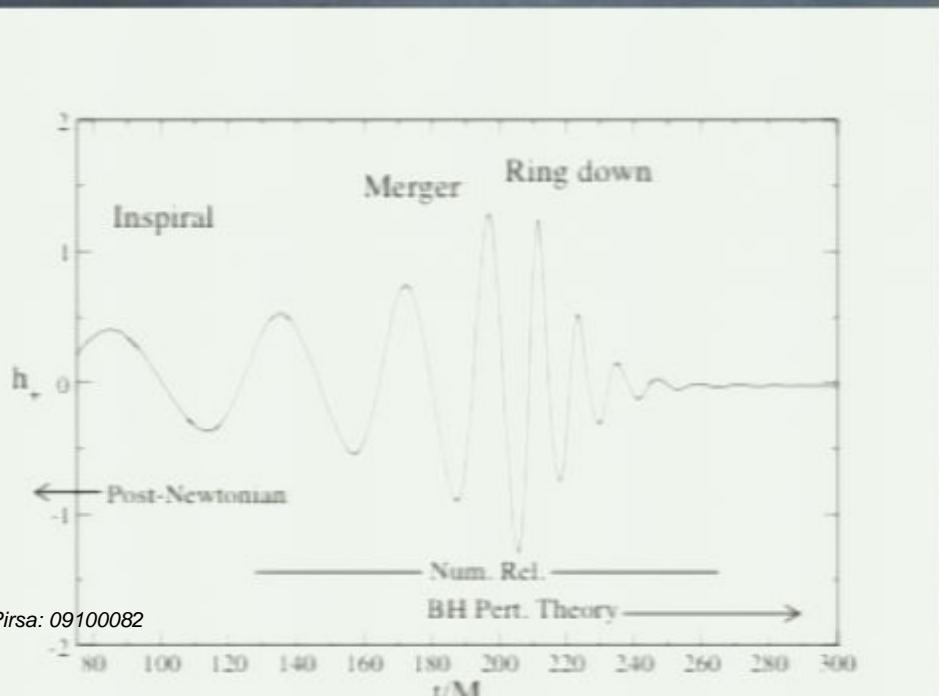
Beam-pattern functions (depend on geometry of detector only)

plus and cross polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---



## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

$$h(t) = F_+ h_+ + F_\times h_\times$$

response function →

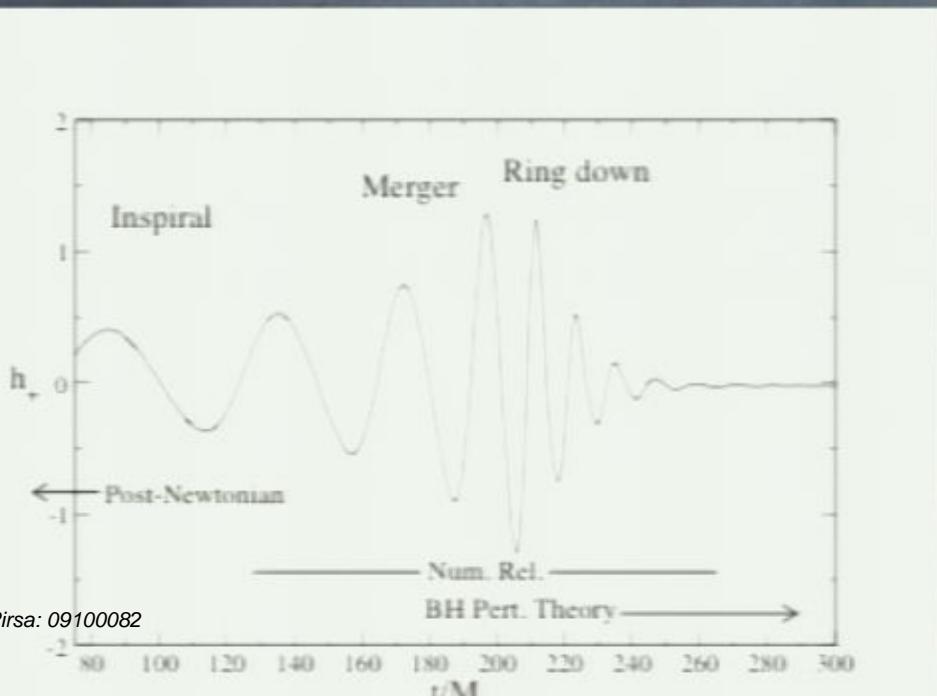
Beam-pattern functions (depend on geometry of detector only)

plus and cross polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---



$$\tilde{h}_I \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$$

Phase depends on  $F_{dot} = E_{dot} (dE/df)^{-1} \rightarrow$  Quadrupole radiation

$$\tilde{h}_M \propto f^{-2/3}$$
 Interpolation

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

$$h(t) = F_+ h_+ + F_\times h_\times$$

response function →  $h(t) = F_+ h_+ + F_\times h_\times$

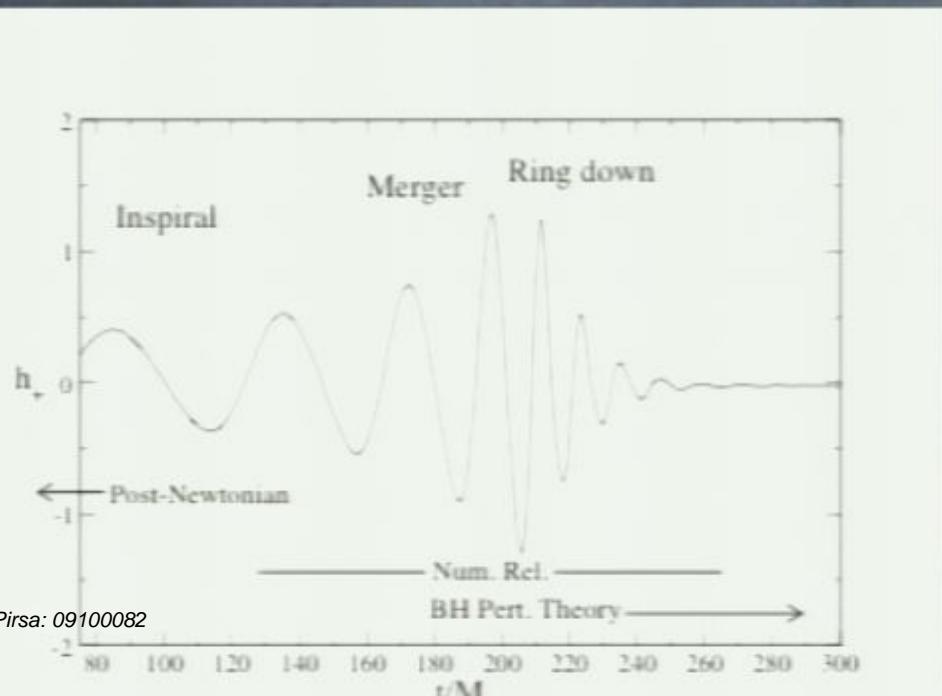
Beam-pattern functions (depend on geometry of detector only)

plus and cross polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---



$$\tilde{h}_I \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$$

Phase depends on  $F_{dot} = E_{dot} (dE/df)^{-1} \rightarrow$  Quadrupole radiation

$$\tilde{h}_M \propto f^{-2/3} \quad \text{Interpolation}$$

$$\tilde{h}_{RD} \propto \frac{\tau}{1+\tau^2(f-f_{RD})^2}$$

## IIc. What can be measured?

Interferometers measure a projection of the GW pert.

$$h(t) = F_+ h_+ + F_\times h_\times$$

response function →  $h(t) = F_+ h_+ + F_\times h_\times$

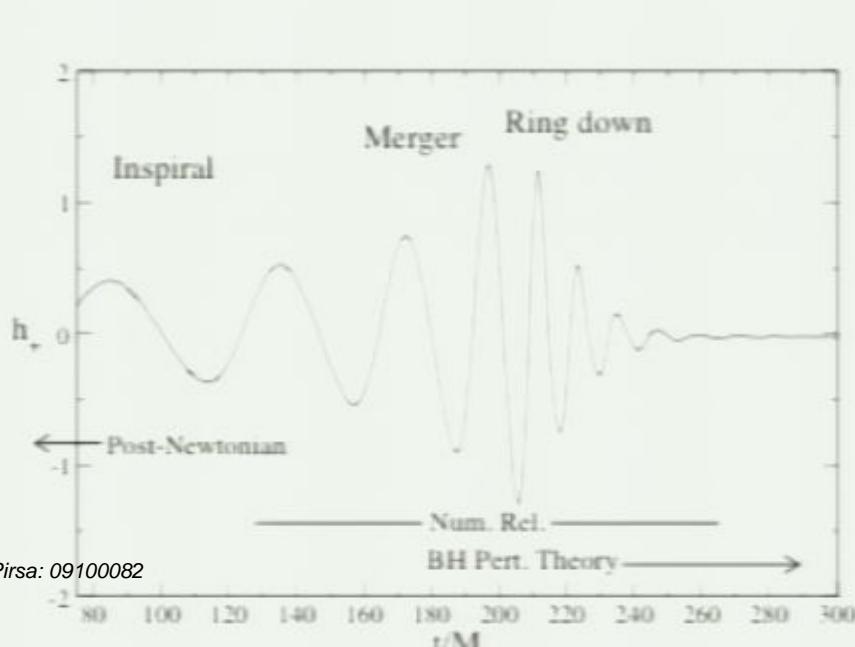
Beam-pattern functions (depend on geometry of detector only)

plus and cross polarizations of  $h_{\mu\nu}$

---

Concentrate on Binary Coalescences

---



$$\tilde{h}_I \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$$

Phase depends on  $F_{dot} = E_{dot} (dE/df)^{-1} \rightarrow$  Quadrupole radiation

$$\tilde{h}_M \propto f^{-2/3} \quad \text{Interpolation}$$

$$\tilde{h}_{RD} \propto \frac{\tau}{1+\tau^2(f-f_{RD})^2}$$

A Lorentzian is the FT of a exponentially damped sinusoid

## IIId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$

---

---

## IIId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$

---

(i) Scalar-Tensor theories:  
(in the inspiral phase)

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS}} f^{-7/3}$$

## IIId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0 f^{-5/3}}$

---

(i) Scalar-Tensor theories:  
(in the inspiral phase)

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS} f^{-7/3}}$$

because of dipolar  
energy emission

## IIId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$

(i) Scalar-Tensor theories:  
(in the inspiral phase)

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS} f^{-7/3}}$$

because of dipolar energy emission  
GW frequency related to scalar field

## IIId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$

(i) Scalar-Tensor theories:  
(in the inspiral phase)

Fourier Transform of  
response function

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS} f^{-7/3}}$$

because of dipolar  
energy emission  
GW frequency  
related to scalar field

## IIId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0} f^{-5/3}$

(i) Scalar-Tensor theories:  
(in the inspiral phase)

Fourier Transform of  
response function

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS} f^{-7/3}}$$

GR part

because of dipolar  
energy emission

GW frequency

related to scalar field

(ii) Chern-Simons Modified Gravity:  
(in the inspiral phase)

## IId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0 f^{-5/3}}$

(i) Scalar-Tensor theories:  
(in the inspiral phase)

Fourier Transform of  
response function

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS} f^{-7/3}}$$

GR part

because of dipolar  
energy emission  
GW frequency  
related to scalar field

(ii) Chern-Simons Modified Gravity:  
(in the inspiral phase)

$$\tilde{h} = \tilde{h}_{GR} (1 + \alpha_{CS} f)$$

## IId. How can the observable be modified?

Concentrate on the Inspiral Phase for now

(0) In GR:  $\tilde{h}_{GR} \propto f^{-7/6} e^{i\psi_0 f^{-5/3}}$

(i) Scalar-Tensor theories:  
(in the inspiral phase)

Fourier Transform of  
response function

$$\tilde{h} = \tilde{h}_{GR} e^{i\beta_{CS} f^{-7/3}}$$

GR part

because of dipolar  
energy emission  
GW frequency  
related to scalar field

(ii) Chern-Simons Modified Gravity:  
(in the inspiral phase)

$$\tilde{h} = \tilde{h}_{GR} (1 + \alpha_{CS} f)$$

related to CS coupling

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

ABC of ppN:

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

ABC of ppN:

- A) Expand the field equations about Minkowski
- B) Assuming a perfect fluid source and a PN expansion, solve the field equations in terms of Green func potentials
- C) Construct a generalization of the metric (a "super-metric") in terms of ppN potentials and ppN parameters.

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

ABC of ppN:

- A) Expand the field equations about Minkowski
- B) Assuming a perfect fluid source and a PN expansion, solve the field equations in terms of Green func potentials
- C) Construct a generalization of the metric (a "super-metric") in terms of ppN potentials and ppN parameters.

$$g_{ij} = \delta_{ij}(1 + 2\gamma U)$$

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

ABC of ppN:

- A) Expand the field equations about Minkowski
- B) Assuming a perfect fluid source and a PN expansion, solve the field equations in terms of Green func potentials
- C) Construct a generalization of the metric (a "super-metric") in terms of ppN potentials and ppN parameters.

$$g_{ij} = \delta_{ij}(1 + 2\gamma U)$$

Newtonian potential

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

ABC of ppN:

- A) Expand the field equations about Minkowski
- B) Assuming a perfect fluid source and a PN expansion, solve the field equations in terms of Green func potentials
- C) Construct a generalization of the metric (a "super-metric") in terms of ppN potentials and ppN parameters.

$$g_{ij} = \delta_{ij}(1 + 2\gamma U)$$

Newtonian potential  
ppN parameter

The diagram shows the equation  $g_{ij} = \delta_{ij}(1 + 2\gamma U)$  enclosed in a yellow rectangular box. Two arrows point from the text "Newtonian potential" and "ppN parameter" to the terms  $(1 + 2\gamma U)$  and  $\delta_{ij}$  respectively, indicating their components in the metric tensor.

## IIIa. Testing Alternative Theories

In the 1970's, the parameterized post-Newtonian framework was developed to cure an outbreak of alt. theories

ABC of ppN:

- A) Expand the field equations about Minkowski
- B) Assuming a perfect fluid source and a PN expansion, solve the field equations in terms of Green func potentials
- C) Construct a generalization of the metric (a "super-metric") in terms of ppN potentials and ppN parameters.

$$g_{ij} = \delta_{ij}(1 + 2\gamma U)$$

GR:  $\gamma = 1$

BD:  $\gamma = \frac{1 + \omega_{BD}}{2 + \omega_{BD}}$

Newtonian potential

ppN parameter

Pisa: 09100082

## IIIb. The parameterized post-Einsteinian scheme

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)
- B) Extract the gravitational perturbation far from the source and construct the Fourier Transform of the response function.
- C) Construct a generalization of the Fourier Transform of the response function (a “super-response”) in terms of a frequency series and ppE parameters.

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)
- B) Extract the gravitational perturbation far from the source and construct the Fourier Transform of the response function.
- C) Construct a generalization of the Fourier Transform of the response function (a “super-response”) in terms of a frequency series and ppE parameters.

$$\tilde{h}(f) = \begin{cases} \tilde{h}_I^{(GR)} \cdot (1 + \alpha u^a) e^{i\beta u^b} & f < f_M, \\ \gamma u^c e^{i(\delta + \epsilon u)} & f_M < f < f_{MRD}, \\ \zeta \frac{\tau}{1 + 4\pi^2 \tau^2 \kappa (f - f_{RD})^d} & f > f_{MRD}. \end{cases}$$

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)
- B) Extract the gravitational perturbation far from the source and construct the Fourier Transform of the response function.
- C) Construct a generalization of the Fourier Transform of the response function (a “super-response”) in terms of a frequency series and ppE parameters.

$$\bar{h}(f) = \begin{cases} \bar{h}_I^{(GR)} \cdot (1 + \alpha u^a) e^{i\beta u^b} & f < f_M, \\ \gamma u^c e^{i(\delta + \epsilon u)} & f_M < f < f_{MRD}, \\ \zeta \frac{\tau}{1 + 4\pi^2 \tau^2 \kappa (f - f_{RD})^d} & f > f_{MRD}. \end{cases}$$

ppE parameters

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)
- B) Extract the gravitational perturbation far from the source and construct the Fourier Transform of the response function.
- C) Construct a generalization of the Fourier Transform of the response function (a “super-response”) in terms of a frequency series and ppE parameters.

GR:  $(\alpha, a, \beta, b) = (0, a, 0, b)$

$$\bar{h}(f) = \begin{cases} \bar{h}_I^{(GR)} \cdot (1 + \alpha u^a) e^{i\beta u^b} & f < f_M, \\ \gamma u^c e^{i(\delta + \epsilon u)} & f_M < f < f_{MRD}, \\ \zeta \frac{\tau}{1 + 4\pi^2 \tau^2 \kappa (f - f_{RD})^d} & f > f_{MRD}. \end{cases}$$

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)
- B) Extract the gravitational perturbation far from the source and construct the Fourier Transform of the response function.
- C) Construct a generalization of the Fourier Transform of the response function (a “super-response”) in terms of a frequency series and ppE parameters.

GR:  $(\alpha, a, \beta, b) = (0, a, 0, b)$

BD:  $(\alpha, a, \beta, b) = (0, a, \beta_{BD}, -7/3)$

$$\tilde{h}(f) = \begin{cases} \tilde{h}_I^{(GR)} \cdot (1 + \alpha u^a) e^{i\beta u^b} & f < f_M, \\ \gamma u^c e^{i(\delta + \epsilon u)} & f_M < f < f_{MRD}, \\ \zeta \frac{\tau}{1 + 4\pi^2 \tau^2 \kappa (f - f_{RD})^d} & f > f_{MRD}. \end{cases}$$

## IIIb. The parameterized post-Einsteinian scheme

ABC of ppE:

- A) Given a physical scenario, solve the field equations ;)
- B) Extract the gravitational perturbation far from the source and construct the Fourier Transform of the response function.
- C) Construct a generalization of the Fourier Transform of the response function (a “super-response”) in terms of a frequency series and ppE parameters.

GR:  $(\alpha, a, \beta, b) = (0, a, 0, b)$

$$\tilde{h} = \tilde{h}_{\text{GR}} (1 + \alpha f^a) e^{i\beta f^b}$$

BD:  $(\alpha, a, \beta, b) = (0, a, \beta_{BD}, -7/3)$

CS:  $(\alpha, a, \beta, b) = (\alpha_{CS}, 1, 0, b)$

ppE parameters

## IIIc. Interpolation vs Extrapolation

## IIIc. Interpolation vs Extrapolation

Interpolation: To reproduce predictions from known alt. theories when ppE parameters acquire certain values.

## IIIc. Interpolation vs Extrapolation

Interpolation: To reproduce predictions from known alt. theories when ppE parameters acquire certain values.

Extrapolation: To reproduce predictions from unknown modifications (guided by symmetry and conservation principles) when ppE parameters acquire other values.

---

---

## IIIc. Interpolation vs Extrapolation

Interpolation: To reproduce predictions from known alt. theories when ppE parameters acquire certain values.

Extrapolation: To reproduce predictions from unknown modifications (guided by symmetry and conservation principles) when ppE parameters acquire other values.

$$\tilde{h} = \tilde{h}_{\text{GR}} (1 + \alpha f^a) e^{i\beta f^b}$$

## IIIc. Interpolation vs Extrapolation

Interpolation: To reproduce predictions from known alt. theories when ppE parameters acquire certain values.

Extrapolation: To reproduce predictions from unknown modifications (guided by symmetry and conservation principles) when ppE parameters acquire other values.

$$\tilde{h} = \tilde{h}_{\text{GR}} (1 + \alpha f^a) e^{i\beta f^b}$$

---

Use ppE-enhanced templates to filter data and let the data choose Nature's ppE parameters.

---

## IIIc. Interpolation vs Extrapolation

Interpolation: To reproduce predictions from known alt. theories when ppE parameters acquire certain values.

Extrapolation: To reproduce predictions from unknown modifications (guided by symmetry and conservation principles) when ppE parameters acquire other values.

$$\tilde{h} = \tilde{h}_{\text{GR}} (1 + \alpha f^a) e^{i\beta f^b}$$

---

Use ppE-enhanced templates to filter data and let the data choose Nature's ppE parameters.

---

You have constructed more general ppE templates that contain less fundamental bias at the cost of introducing more ppE parameters.

## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

Testing GR will require collaboration between waveform modelers, alternative theory architects and data analysts.

## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

Testing GR will require collaboration between waveform modelers, alternative theory architects and data analysts.

Difficult to see.

Always in motion the future is.

## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

Testing GR will require collaboration between waveform modelers, alternative theory architects and data analysts.

Difficult to see.

Always in motion the future is.

ppE templates for spinning, non-equal mass binaries?

## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

Testing GR will require collaboration between waveform modelers, alternative theory architects and data analysts.

Difficult to see.

Always in motion the future is.

ppE templates for spinning, non-equal mass binaries?

Effectiveness of ppE templates at extracting signals?

Other altenative theories predictions to parameterize?

## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

Testing GR will require collaboration between waveform modelers, alternative theory architects and data analysts.

Difficult to see.

Always in motion the future is.

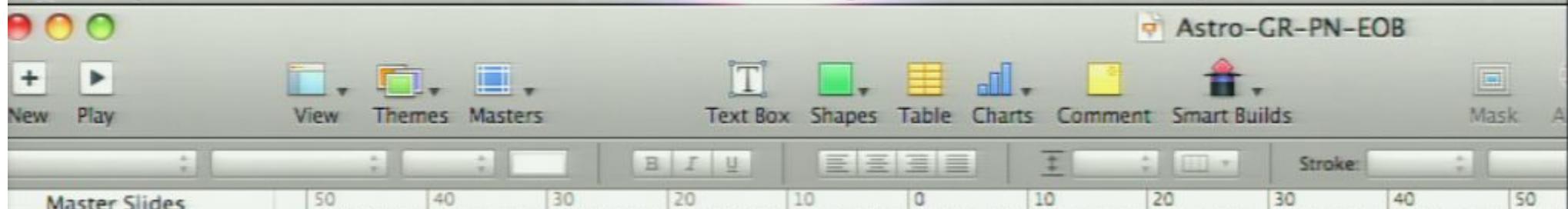
ppE templates for spinning, non-equal mass binaries?

Effectiveness of ppE templates at extracting signals?

Other altenative theories predictions to parameterize?

So how is GR found in this Trial?

Data is the jury and it will return a verdict soon.



## Conclusions

You cannot assume GR to test GR (Fundamental Bias).

New ppE framework is the first step to acknowledge fundamental bias and allow for true tests.

Testing GR will require collaboration between waveform modelers, alternative theory architects and data analysts.

Difficult to see.

Always in motion the future is.

ppE templates for spinning, non-equal mass binaries?

Effectiveness of ppE templates at extracting signals?

Other alternative theories predictions to parameterize?

So how is GR found in this Trial?



Finder File Edit View Go Window Help



Thu 2:53 PM

(59%)



NO NAME



Einstein-3



propEinstei  
guide.ftr  
n.pdf



PI-  
psswd.tex



statEinstei  
n.pdf