

Title: Entanglement detection with bounded reference frames

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Abstract: Violation of local realism can be probed by theory-independent tests, such as Bell's inequality experiments. There, a common assumption is the existence of perfect, classical, reference frames, which allow for the specification of measurement settings with arbitrary precision. However, if the reference frames are "bounded", only limited precision can be attained. We expect then that the finiteness of the reference frames limits the observability of genuine quantum features. Using spin coherent states as reference frames, we determined their minimal size necessary to violate Bell's inequalities in entangled systems ranging from qubits to macroscopic dimensions. In the latter, the reference frame's size must be quadratically larger than that of the system. Lacking such large reference frames, precludes quantum phenomena from appearing in everyday experience.

Entanglement detection with bounded reference frames

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Reference:

F. Costa, N. Harrigan, T. Rudolph, Č. Brukner,
arXiv:0903.0035

Perimeter Institute,
August '09

Reference frames in quantum mechanics

Reference frames in quantum mechanics

- Physical states are defined relative to a reference frame
e.g. position with respect to a Cartesian frame, time with respect to a clock, etc...

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- What is possible to say in absence of a reference frame?

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- What is possible to say in absence of a reference frame?
- What if the reference frame itself is a quantum system?

Relative properties between quantum mechanical systems

Outline

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- Reference frames in quantum mechanics: a brief review.

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- Higher spins and classicality.
- Conclusions.

Communicating a direction without a shared reference frame

Comunicating a direction without a shared reference frame



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Formalizing reference frames



$|\psi\rangle$

Formalizing reference frames

Change of reference frame induces a passive transformation



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$$|\psi\rangle \longrightarrow R(g)|\psi\rangle$$

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e.g. rotations-SU(2)

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Transformations of a RF form a group

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Passive transformations are a representation of the group

e.g. spin

Quantum states in absence of a reference frame

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$$\mathcal{H} = \bigoplus_q \mathcal{M}_q \otimes \mathcal{N}_q$$

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Projects over this

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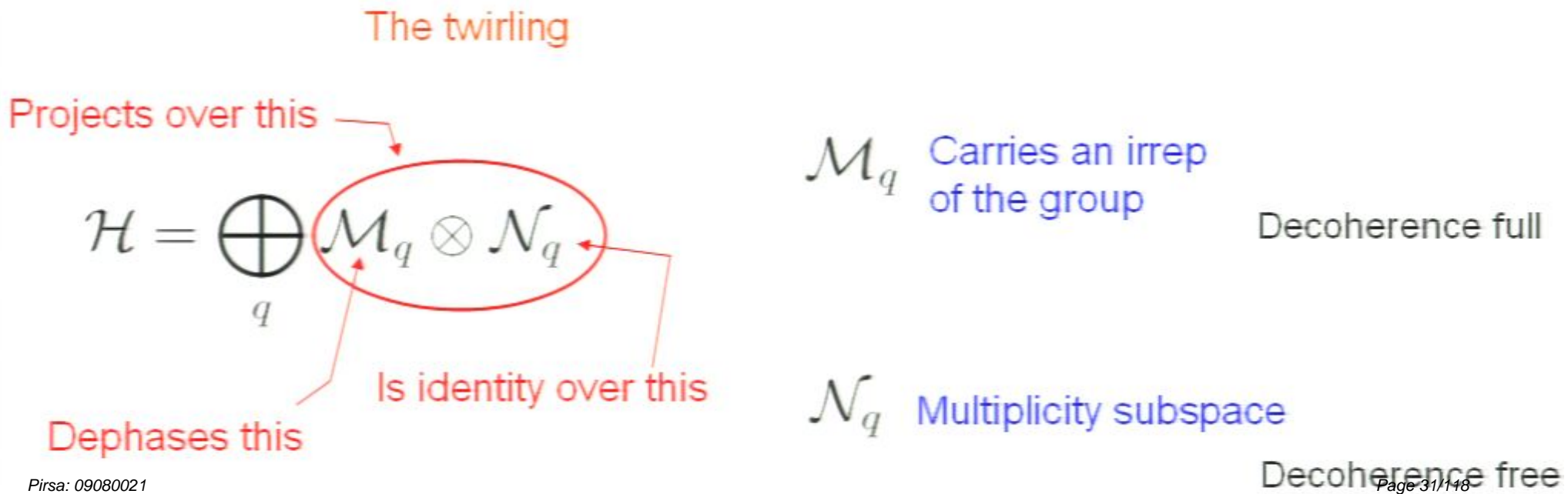
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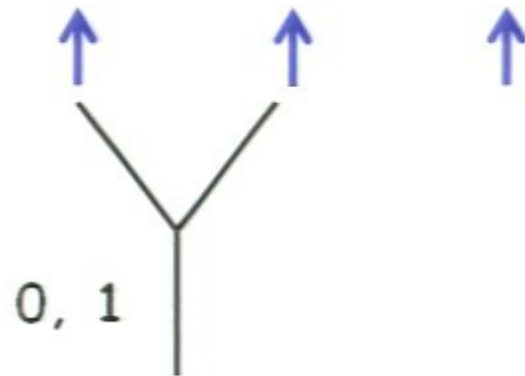
Group action induces decomposition of Hilbert space



Example: 3 spin-1/2 particles with no external RF

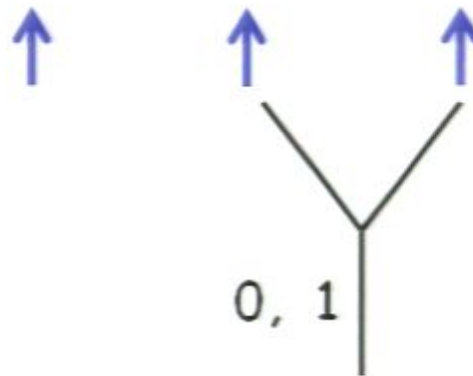


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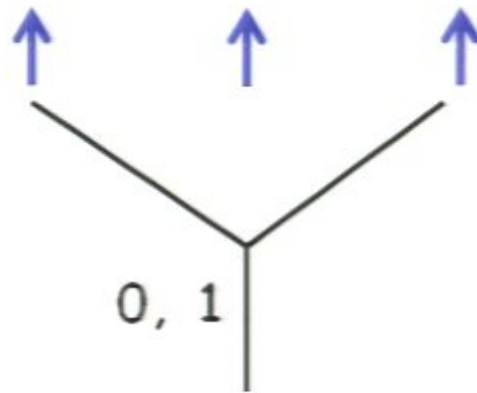
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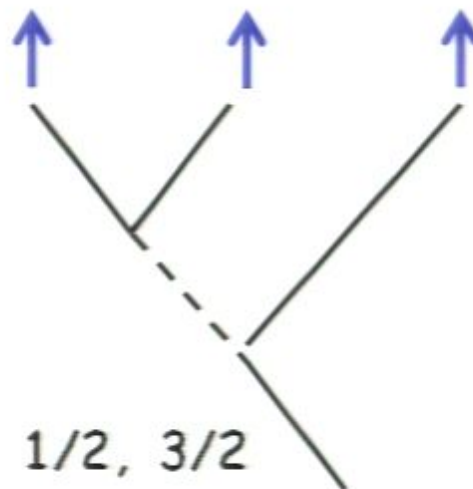
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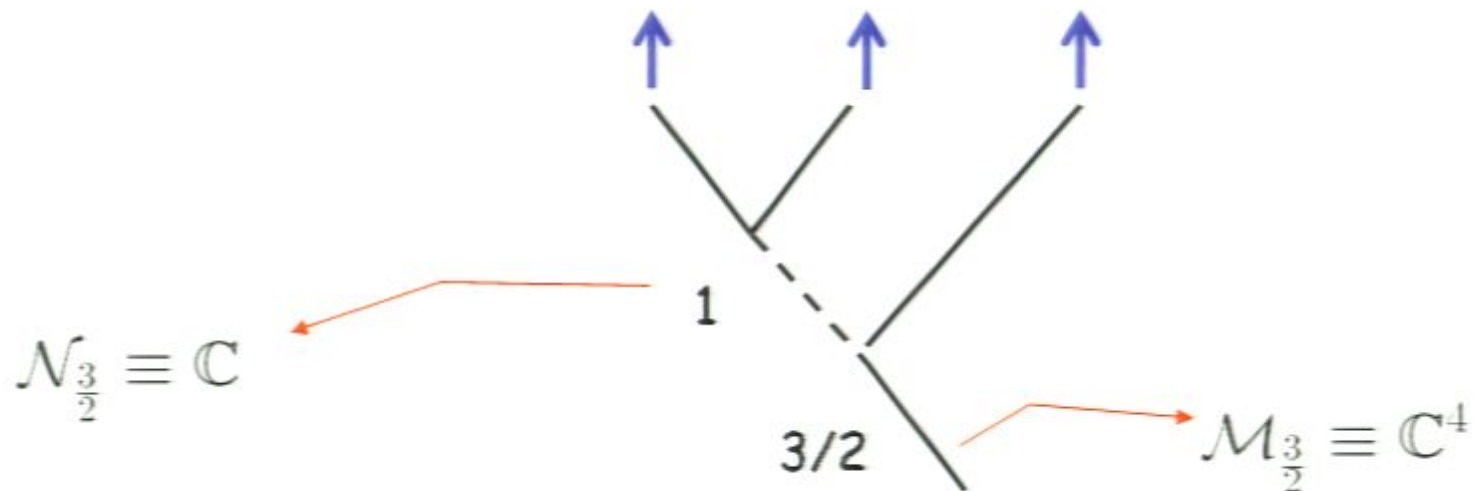
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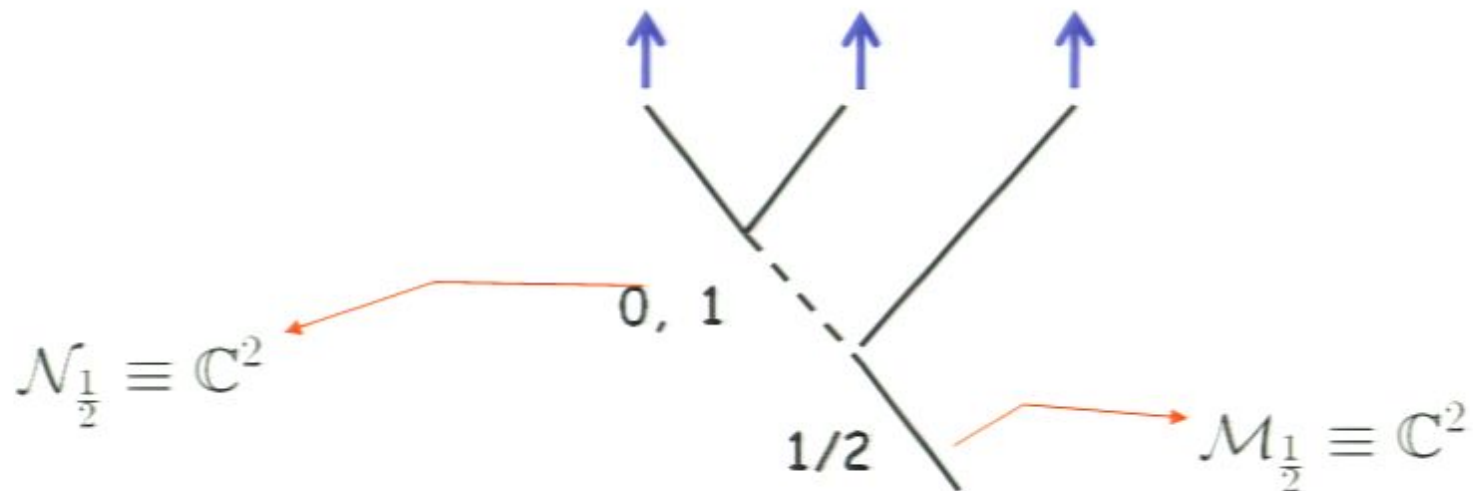
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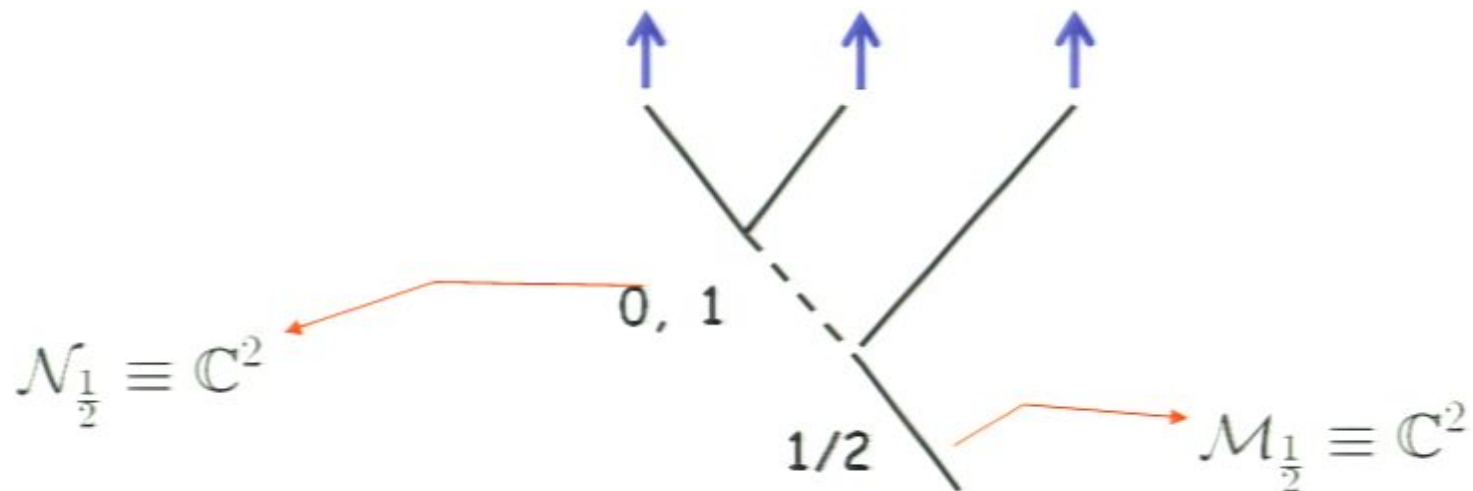
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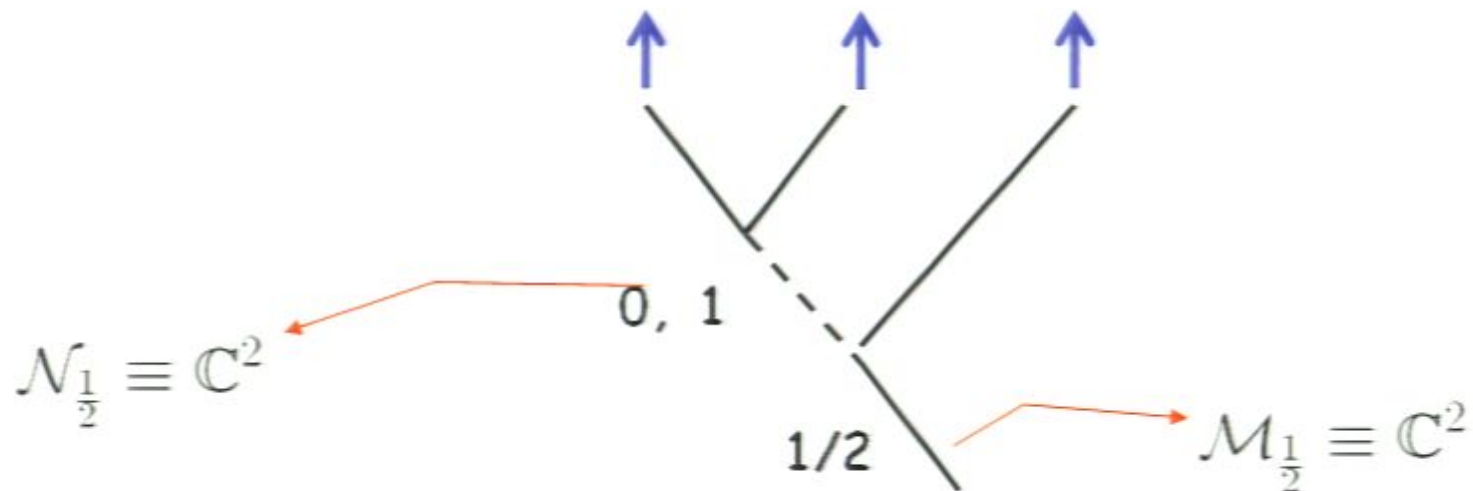


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After the twirling only these degrees of freedom are accessible

Quantum reference frames

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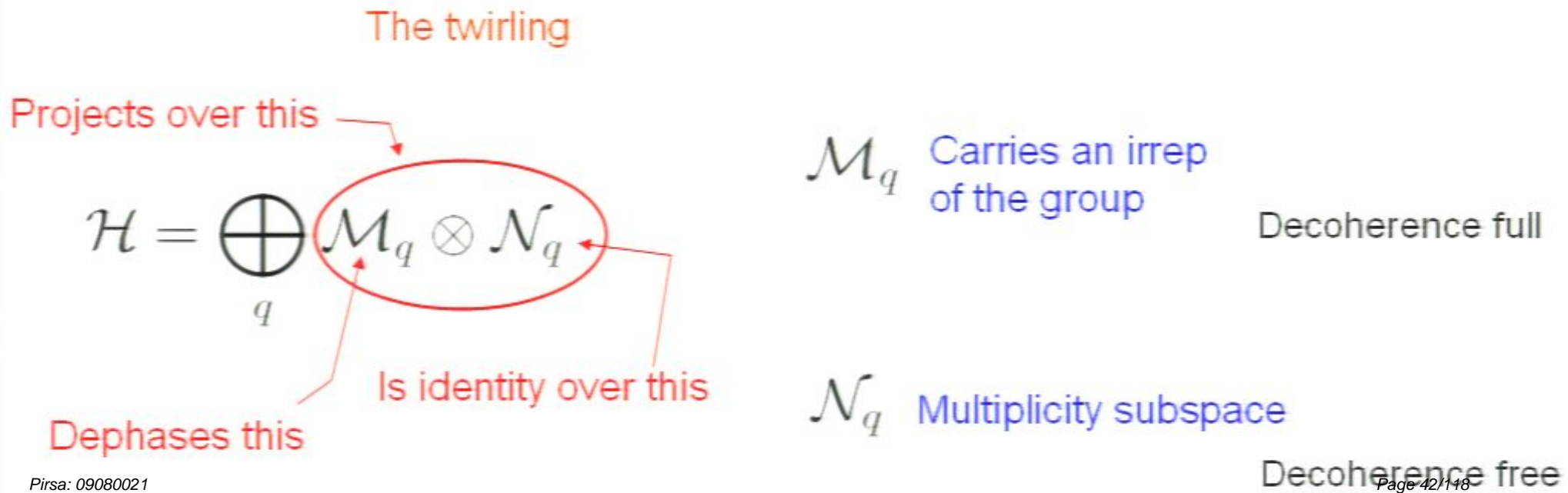
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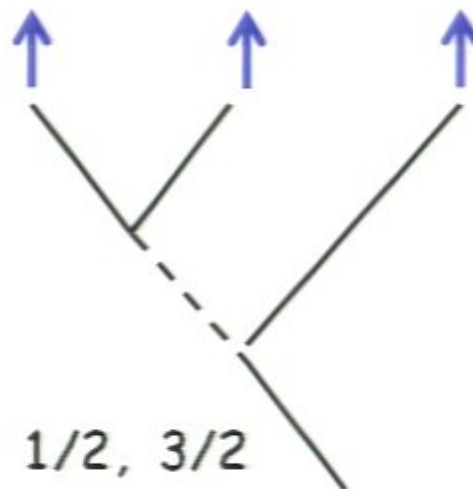
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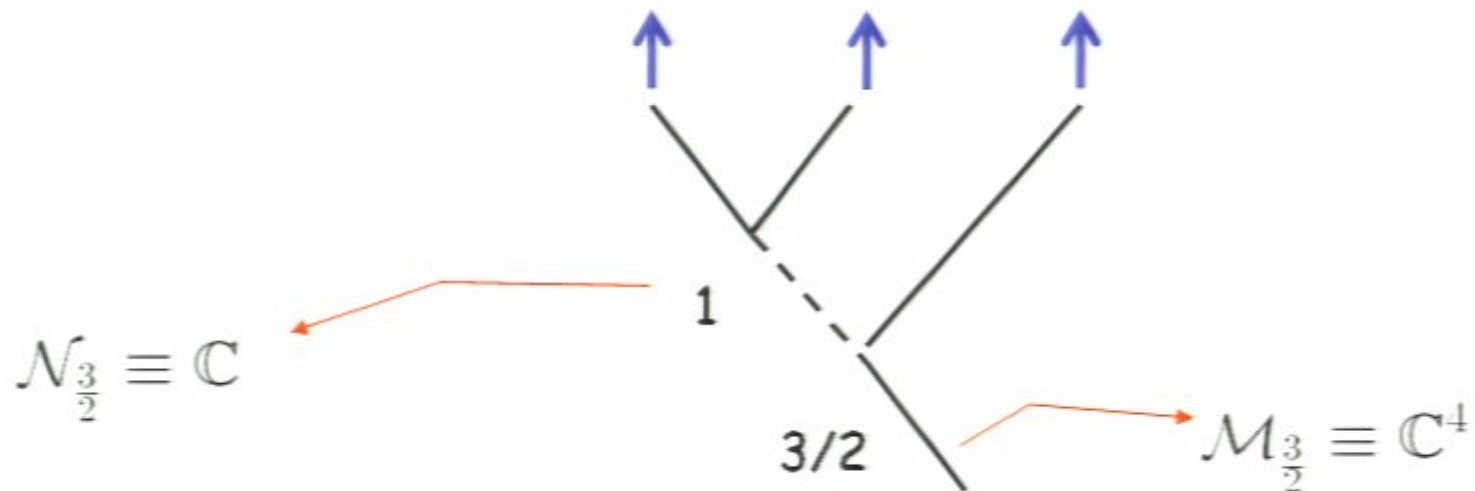


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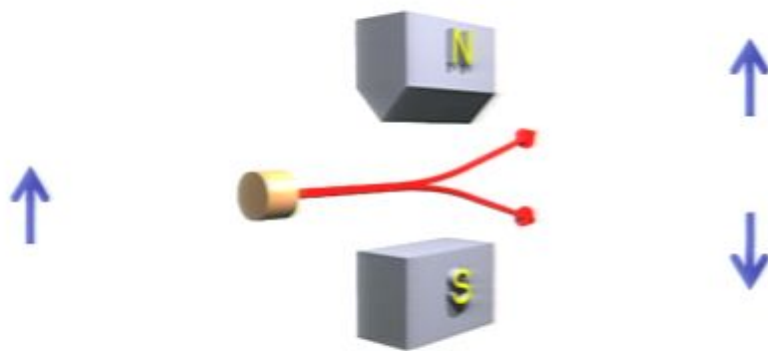


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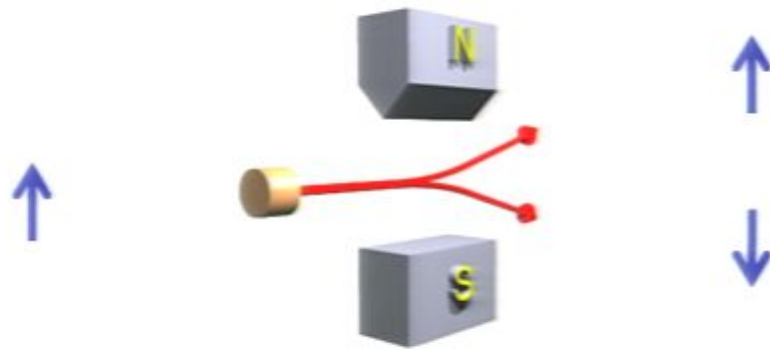
Quantum reference frames

Measure the projection of a spin $\frac{1}{2}$ particle along a given axis



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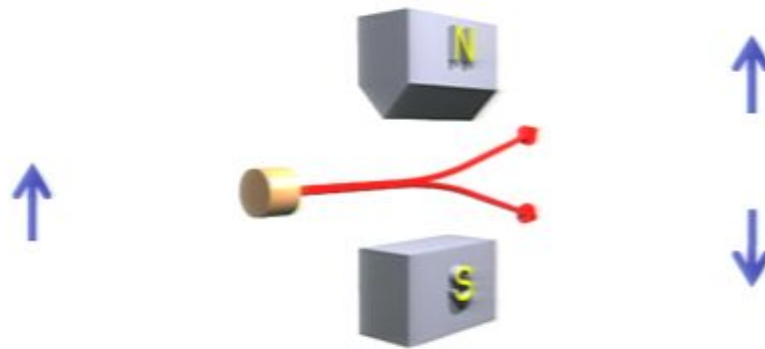
Use a spin j coherent state as reference frame



$|j, j\rangle$

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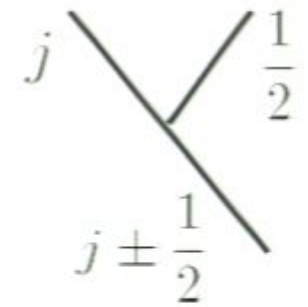


$$|j + \frac{1}{2}, j + \frac{1}{2}\rangle$$

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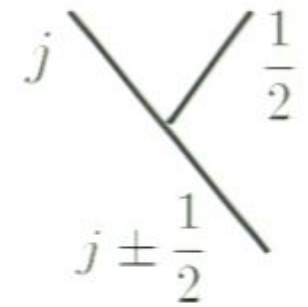
We can interpret:



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$j + \frac{1}{2}$ \longrightarrow Spin was up
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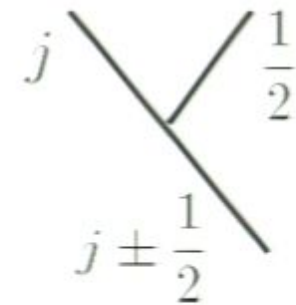


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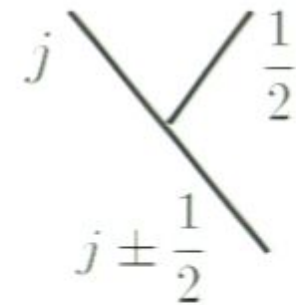


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But we have three possible states for the joint system:

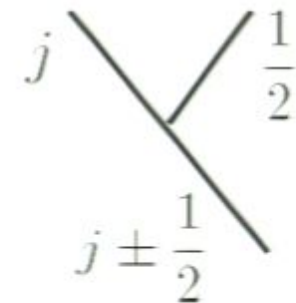
$$\left| \frac{1}{2} \right\rangle \otimes \left| j, j \right\rangle$$

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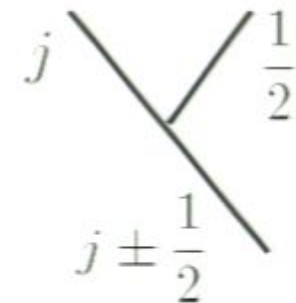
$$|\frac{1}{2}\rangle \otimes |j, j\rangle \longrightarrow \begin{aligned} &|j + \frac{1}{2}, j + \frac{1}{2}\rangle \\ &|j + \frac{1}{2}, j - \frac{1}{2}\rangle \\ &|j - \frac{1}{2}, j - \frac{1}{2}\rangle \end{aligned}$$

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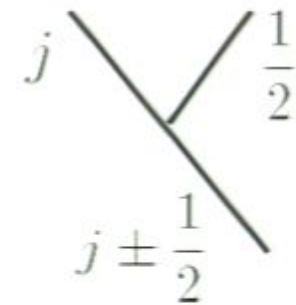
In one case we guess wrong!

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But we have three possible states for the joint system:

$$|\frac{1}{2}\rangle \otimes |j, j\rangle \longrightarrow \begin{matrix} |j + \frac{1}{2}, j + \frac{1}{2}\rangle \\ |j + \frac{1}{2}, j - \frac{1}{2}\rangle \\ |j - \frac{1}{2}, j - \frac{1}{2}\rangle \end{matrix}$$

The middle state $|j + \frac{1}{2}, j - \frac{1}{2}\rangle$ is circled in blue.

In one case we guess wrong!

The probability for the wrong guess vanishes for $j \rightarrow \infty$

(Classical reference frame)

Spin measurement along „direction“

- Possible
Outcomes: $j = |j_{RF} - j_S|, \dots, j_{RF} + j_S$
- Outcome: $j = j_{RF} + m$
- Projector: $\hat{\Pi}_{j_{RF}+m} = \sum_{m'=-j_{RF}-m}^{j_{RF}+m} |j_{RF} + m, m'\rangle \langle j_{RF} + m, m'|$

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Quantum Reference Frame in a coherent state

Effective POVM:

$$\hat{P}_{j_S m}^{j_{RF}} = \langle j_{RF} | \hat{\Pi}_{j_{RF}+m} | j_{RF} \rangle \xrightarrow{j_{RF} \rightarrow \infty} |j_S, m\rangle \langle j_S, m| \equiv \hat{\Pi}_m$$

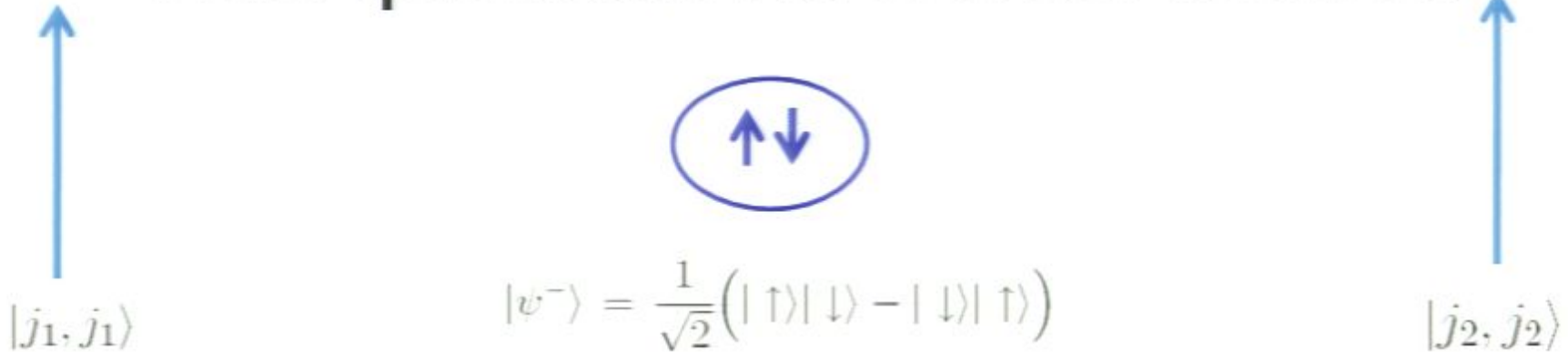


Von Neumann
spin projection

Violation of Bell's inequalities with quantum reference frames

$$S := \left| E(\vec{\alpha}_1, \vec{\beta}_1) + E(\vec{\alpha}_1, \vec{\beta}_2) + E(\vec{\alpha}_2, \vec{\beta}_1) - E(\vec{\alpha}_2, \vec{\beta}_2) \right| \leq 2,$$

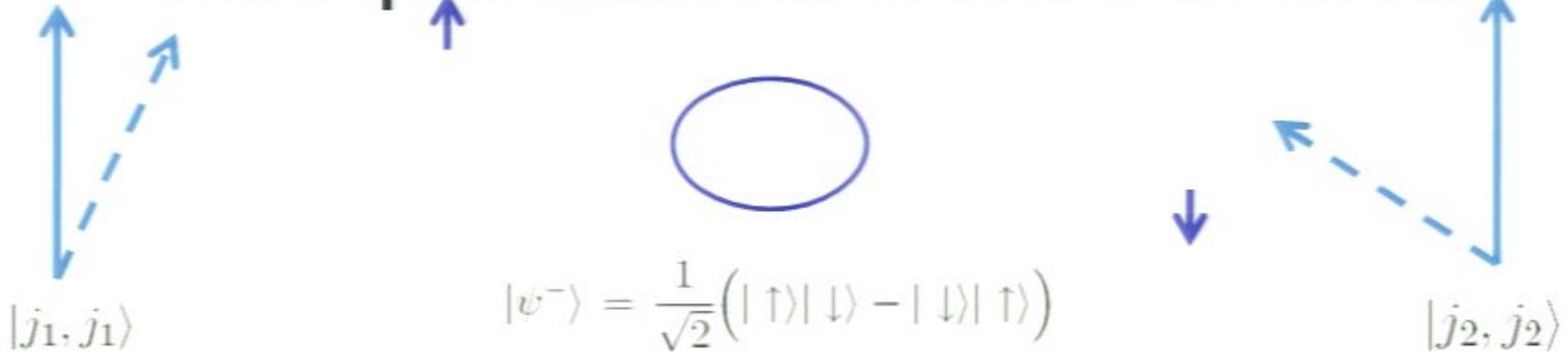
Violation of Bell's inequalities with quantum reference frames



$|j_1, j_1\rangle$ $|\psi^-\rangle = \frac{1}{\sqrt{2}}(|\uparrow\rangle|\downarrow\rangle - |\downarrow\rangle|\uparrow\rangle)$ $|j_2, j_2\rangle$

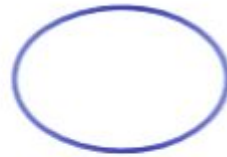
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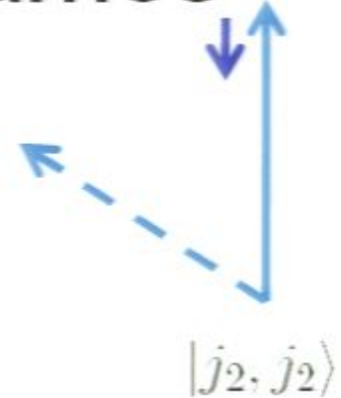


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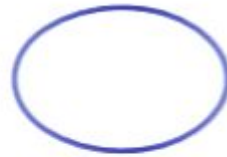


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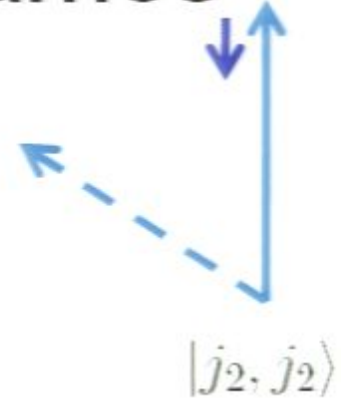
$$E^{j_1 j_2}(\vartheta) = \frac{1 - 4j_1 j_2 \cos \vartheta}{(2j_1 + 1)(2j_2 + 1)} \xrightarrow{j_1, j_2 \rightarrow \infty} E(\vartheta) = -\cos \vartheta$$

What is the minimum size of j_1 , j_2 for which Bell's inequalities are violated?

Violation of Bell's inequalities with quantum reference frames



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$$S_{j_1, j_2} = 2 \left| \frac{1 + 4\sqrt{2}j_1 j_2}{(2j_1 + 1)(2j_2 + 1)} \right|$$

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Interlude

Effective decoherence

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Effective decoherence map

$$\mathcal{E}_j(\rho) = \left(\frac{j}{j+1} \mathcal{I} + \frac{1}{j+1} \mathcal{G} \right) [\rho]$$

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Can we apply this to our case?

Yes, be we find a different result...

Bell inequalities under effective decoherence

Singlet state $|\psi^-\rangle\langle\psi^-| = \frac{1}{4}(\mathbb{1} - \vec{\sigma}^A \cdot \vec{\sigma}^B)$

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Apply an independent effective decoherence on each side

$$\mathcal{E}_{j_1} \otimes \mathcal{E}_{j_2}(|\psi^-\rangle\langle\psi^-|) = \frac{1}{4} \left[\mathbb{I} - \frac{j_1 j_2}{(j_1 + 1)(j_2 + 1)} \vec{\sigma}^A \cdot \vec{\sigma}^B \right]$$

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Correlation function $E_{j_1 j_2}(\vartheta) = -\frac{j_1 j_2}{(j_1 + 1)(j_2 + 1)} \cos \vartheta$

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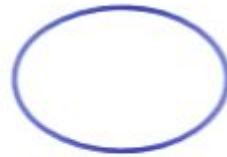
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Can we apply this to our case?

Yes, because we find a different result...

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$$S_{j_1, j_2} = 2 \left| \frac{1 + 4\sqrt{2}j_1 j_2}{(2j_1 + 1)(2j_2 + 1)} \right|$$

What is the minimum size of j_1, j_2 for which Bell's inequalities are violated?

Interlude

Effective decoherence

Is it possible to „trace out“ the reference frame and just work with the system?

YES.....

Bartlett, Rudolph, Spekkens, Turner '09

Effective decoherence map $\mathcal{E}_j(\rho) = \left(\frac{j}{j+1} \mathcal{I} + \frac{1}{j+1} \mathcal{G} \right) [\rho]$

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Can we apply this to our case?

Yes, be we find a different result...

Bell inequalities under effective decoherence

Singlet state $|\psi^-\rangle\langle\psi^-| = \frac{1}{4} (\mathbb{I} - \vec{\sigma}^A \cdot \vec{\sigma}^B)$

Apply an independent effective decoherence on each side

$$\mathcal{E}_{j_1} \otimes \mathcal{E}_{j_2}(|\psi^-\rangle\langle\psi^-|) = \frac{1}{4} \left[\mathbb{I} - \frac{j_1 j_2}{(j_1 + 1)(j_2 + 1)} \vec{\sigma}^A \cdot \vec{\sigma}^B \right]$$

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Larger than before

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Spin-0 State: $|\Psi_{j_S}^-\rangle := \frac{1}{\sqrt{2j_S + 1}} \sum_{m=-j_S}^{j_S} (-1)^{j_S-m} |j_S, m\rangle |j_S, -m\rangle$

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Angle θ must be given with precision better than $1/j$

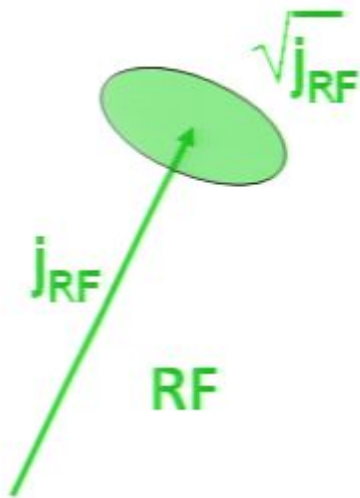
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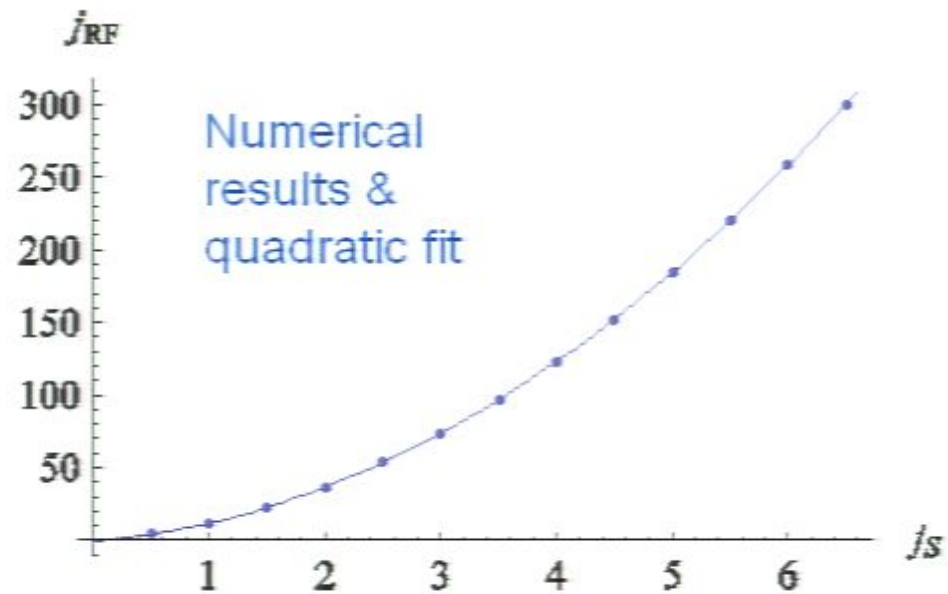
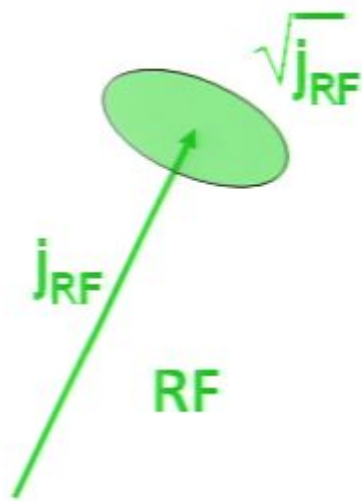
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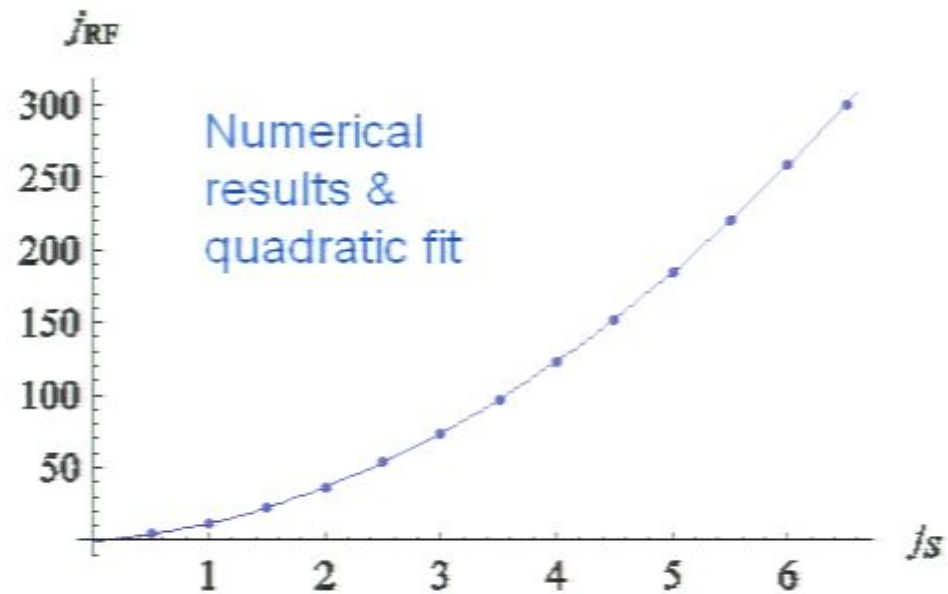
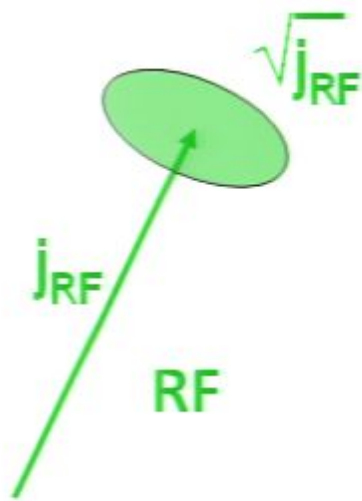
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Thank you for your attention!