Title: The null energy condition and its violators

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Abstract: TBA

Alberto Nicolis Columbia University

The null energy condition and its violators

w/ Rattazzi and Trincherini, to appear

also: w/ Dubovsky, Gregoire, Rattazzi, '05 w/ Creminelli, Luty, Senatore, '06

Energy conditions in GR

- " E > 0"
- several ways to make it covariant: weak, strong, dominant, null (...?)
- ullet different contractions of $T_{\mu
 u}$

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@ CTC's

positive energy theorem

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O CTC's

positive energy theorem

2nd law of black-hole thermodynamics

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- O CTC's
- positive energy theorem
- 2nd law of black-hole thermodynamics

... and bad ones:

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- @ CTC's
- positive energy theorem
- 2nd law of black-hole thermodynamics

... and bad ones:

singularity theorems

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The "null" one (NEC) stands out as the most robust

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 $T_{\mu\nu} n^{\mu} n^{\nu} \geq 0$ for all null n^{μ} 's

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The "null" one (NEC) stands out as the most robust

 $\sigma T_{\mu\nu} \, n^\mu n^\nu \geq 0 \;\; {
m for \; all \; null \;} n^\mu \, {
m s} \, {
m for \;} {
m all \;} {
m for \;} {
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m for \;}$

ullet saturated by a c.c. $T_{\mu
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all the others violated or fixed by a suitable c.c.

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- The "null" one (NEC) stands out as the most robust
- $\sigma T_{\mu\nu} \, n^\mu n^\nu \geq 0$ for all null n^μ 's
- $m{o}$ saturated by a c.c. $T_{\mu
 u} \propto g_{\mu
 u}$
- all the others violated or fixed by a suitable c.c. ambiguous somewhat

NEC closely related to the "dominant" one (DEC)

$$\mathsf{DEC} = \mathsf{NEC} + \ \rho \geq 0$$

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more covariantly

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no superluminal flow of energy-momentum for any observer.

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 \bullet For cosmology: NEC $\hfill(\rho+p)\geq 0$

$$(\rho + p) \ge 0$$

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$$\Rightarrow$$

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$$egin{aligned} egin{aligned} egin{aligned} egin{aligned} \dot{H} \propto -(
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NEC



Expansion?

 $m{\circ}$ For cosmology: NEC $\qquad \qquad (
ho+p)\geq 0$



$$(\rho + p) \ge 0$$

Friedmann eqs.

$$\dot{H} \propto -(\rho + p)$$

@ NEC



Expansion? Big Bang Need UV-completion

$$\bullet$$
 For cosmology: NEC $\qquad \qquad (\rho+p) \geq 0$



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Expansion?

Big Bang

Need UV-completion

For black holes:

NEC/DEC



2nd law



$$\dot{H} \propto -(\rho + p)$$

@ NEC



Expansion?

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For black holes:

NEC/DEC



crucial for thermodynamic/holographic interpretation

holographic cosmology

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Difficult! Usually:

stability NEC

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Difficult! Usually:

Simplest example

$$\mathcal{L} = \pm \frac{1}{2} (\partial \phi)^2 - V(\phi)$$

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Simplest example

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$$\phi = \phi(t) \qquad (\rho + p) = \pm \dot{\phi}^2$$

Difficult! Usually:



Simplest example

$$\mathcal{L} = \pm \frac{1}{2} (\partial \phi)^2 - V(\phi)$$

$$\phi = \phi(t)$$
 $(\rho + p) = \pm \dot{\phi}^2$

I want to qualify the "usually"

Neglect gravity for the moment

Well defined QFT question

Whatever we get, will translate into an "Einstein frame" statement

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Fact

NEC needs spontaneous Lorentz breaking

$$T_{\mu\nu} \neq \eta_{\mu\nu}$$

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- There are light Golstones!
- Their dynamics largely model-independent

Fact

NEC needs spontaneous Lorentz breaking

$$T_{\mu\nu} \neq \eta_{\mu\nu}$$

- There are light Golstones!
- Their dynamics largely model-independent
- Those are the guys to worry about

Consider a system of scalars

$$\mathcal{L} = F(\phi^I, \partial \phi^I, \partial \partial \phi^I, \dots)$$
$$I, J, \dots = 1, \dots, N$$

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- Includes any mixture of fluids and solids
- ullet At low energies stop at the $\partial\phi^I$ level
- Enough to see Lorentz breaking

$$\mathcal{L} \to F(\phi^I, \partial \phi^I) \to F(\phi^I, B^{IJ})$$
$$B^{IJ} \equiv \partial_\mu \phi^I \partial^\mu \phi^J$$

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 $m{\circ}$ Lorentz breaking solution: $\partial_{\mu}\phi^{I} \neq 0$

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$$\mathcal{L} \sim \left[F_{IJ} \eta_{\mu\nu} + 2 F_{IK,JL} \, \partial_{\mu} \phi^K \partial_{\nu} \phi^L \right] \partial^{\mu} \pi^I \partial^{\nu} \pi^J$$

(Dubovsky, Gregoire, Nicolis, Rattazzi 2005)

Theorem:

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(Dubovsky, Gregoire, Nicolis, Rattazzi 2005)

Theorem:

Stability and

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(Dubovsky, Gregoire, Nicolis, Rattazzi 2005)

Theorem:

Stability and isotropy

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(Dubovsky, Gregoire, Nicolis, Rattazzi 2005)

Theorem:

Stability and

isotropy or subluminality

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Theorem:

Stability and

isotropy or subluminality



NEC !

Conversely, there are stable systems with anistropic and superluminal NEC-violating solutions

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Conversely, there are stable systems with anistropic and superluminal NEC-violating solutions

useless for cosmology

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Conversely, there are stable systems with anistropic and superluminal NEC-violating solutions

useless for cosmology

 to evade the theorem, try to evade the assumptions

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Rewind...

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Rewind... Consider a system of scalars

$$\mathcal{L} = F(\phi^I, \partial \phi^I, \partial \partial \phi^I, \dots)$$
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Most relevant at low energies

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Higher derivatives problematic (when important):

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- Most relevant at low energies
- Higher derivatives problematic (when important):
- classically...

$$(\partial\phi)^2 + \frac{1}{M^2}(\Box\phi)^2 \to (\partial\phi)^2 - (\partial\chi)^2 + M^2\chi^2$$

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- Higher derivatives problematic (when important):
- classically...

$$(\partial \phi)^2 + \frac{1}{M^2} (\Box \phi)^2 \to (\partial \phi)^2 - (\partial \chi)^2 + M^2 \chi^2$$

... and quantum-mechanically:



(Arkani-Hamed, Cheng, Luty, Mukohyama 2003)

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(Arkani-Hamed, Cheng, Luty, Mukohyama 2003)

 Suppose the system is degenerate at lowest order in spatial derivatives

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(Arkani-Hamed, Cheng, Luty, Mukohyama 2003)

 Suppose the system is degenerate at lowest order in spatial derivatives

$$\mathcal{L} \sim \dot{\pi}^2 - 0 \cdot (\vec{\nabla}\pi)^2$$

Higher derivative terms

$$(\Box \phi)^2 \rightarrow \ddot{\pi}^2, \ (\vec{\nabla} \dot{\pi})^2, \ (\nabla^2 \pi)^2$$

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negligible at low energies

(Arkani-Hamed, Cheng, Luty, Mukohyama 2003)

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leading gradient energy

@ Consistent derivative expansion (w/ non-relativistic scaling $\omega \sim \vec{k}^2$)

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no ghost-like new d.o.f.

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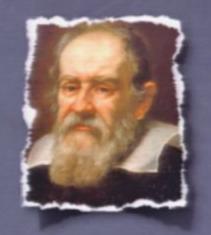
no ghost-like new d.o.f.

allows consistent NEC-violating cosmological scenarios: bounce, "starting the universe", w<-1 now</p>

> (Creminelli, Luty, Nicolis, Senatore 2006, Creminelli, Senatore 2007, Creminelli's talk, yesteeday)

2nd Caveat: the Galileon

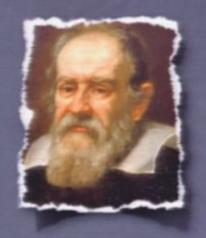
(Nicolis, Rattazzi, Trincherini 2008)



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2nd Caveat: the Galileon

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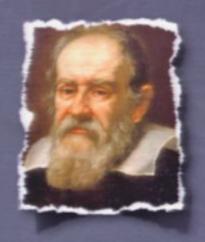


The (classical) problem is having higherderivative eom

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2nd Caveat: the Galileon

(Nicolis, Rattazzi, Trincherini 2008)



 The (classical) problem is having higherderivative eom

Is there a higher-derivative Lagrangian that yields two-derivative eom?

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For simplicity assume purely two-derivative eom

$$\frac{\delta \mathcal{L}}{\delta \pi} = f(\partial_{\mu} \partial_{\nu} \pi)$$

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$$\frac{\delta \mathcal{L}}{\delta \pi} = f(\partial_{\mu} \partial_{\nu} \pi)$$

Invariant under

$$\pi(x) \to \pi(x) + c + b_{\mu} x^{\mu}$$

"Galilean invariance"

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For simplicity assume purely two-derivative eom

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Invariant under

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"Galilean invariance"

 $m{\circ}$ Analogous to $x(t)
ightarrow x(t) + x_0 + v_0 \, t$

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$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \pi^n$$

Simplest:

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \pi^n$$

 \circ Simplest: $\mathcal{L}^{(1)}=\pi$

$$\mathcal{L}^{(1)} = \pi$$

Next to simplest:

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \pi^n$$

$$\circ$$
 Simplest: $\mathcal{L}^{(1)}=\pi$

$$\mathcal{L}^{(1)} = \pi$$

• Next to simplest: $\mathcal{L}^{(2)} = (\partial \pi)^2$

$$\mathcal{L}^{(2)} = (\partial \pi)^2$$

Less trivial (DGP):

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \pi^n$$

• Simplest:
$$\mathcal{L}^{(1)} = \pi$$

• Next to simplest:
$$\mathcal{L}^{(2)} = (\partial \pi)^2$$

• Less trivial (DGP):
$$\mathcal{L}^{(3)} = (\partial \pi)^2 \, \Box \pi$$

Invariance:

Exactly one invariant at each order in π

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Exactly one invariant at each order in π

Galilean invariant Lagragian term

 \leftrightarrow

Cayley-invariant of matrix $\partial_{\mu}\partial_{\nu}\pi$

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Exactly one invariant at each order in π

Galilean invariant Lagragian term



Cayley-invariant of matrix $\partial_{\mu}\partial_{\nu}\pi$

Rank(
$$\partial_{\mu}\partial_{\nu}\pi$$
) = D



D invariants (+tadpole)

Exactly one invariant at each order in π

Galilean invariant Lagragian term

 \longleftrightarrow

Cayley-invariant of matrix $\partial_{\mu}\partial_{\nu}\pi$

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D invariants (+tadpole)

$$x, \dot{x}^2$$

Exactly one invariant at each order in π

Galilean invariant Lagragian term

 \longleftrightarrow

Cayley-invariant of matrix $\partial_{\mu}\partial_{\nu}\pi$

Rank($\partial_{\mu}\partial_{\nu}\pi$) = D

D invariants (+tadpole)

D=1 (mechanics):

$$x, \dot{x}^2$$

D=4 (us):
$$\pi, (\partial \pi)^2, \ldots \partial \pi \partial \pi (\partial^2 \pi)^3$$

Galilean Invariants

$$\mathcal{L}_{1} = \pi$$

$$\mathcal{L}_{2} = -\frac{1}{2}\partial\pi \cdot \partial\pi$$

$$\mathcal{L}_{3} = -\frac{1}{2}[\Pi]\partial\pi \cdot \partial\pi$$

$$\mathcal{L}_{4} = -\frac{1}{4}([\Pi]^{2}\partial\pi \cdot \partial\pi - 2[\Pi]\partial\pi \cdot \Pi \cdot \partial\pi - [\Pi^{2}]\partial\pi \cdot \partial\pi + 2\partial\pi \cdot \Pi^{2} \cdot \partial\pi)$$

$$\mathcal{L}_{5} = -\frac{1}{5}([\Pi]^{3}\partial\pi \cdot \partial\pi - 3[\Pi]^{2}\partial\pi \cdot \Pi \cdot \partial\pi - 3[\Pi][\Pi^{2}]\partial\pi \cdot \partial\pi + 6[\Pi]\partial\pi \cdot \Pi^{2} \cdot \partial\pi$$

$$+2[\Pi^{3}]\partial\pi \cdot \partial\pi + 3[\Pi^{2}]\partial\pi \cdot \Pi \cdot \partial\pi - 6\partial\pi \cdot \Pi^{3} \cdot \partial\pi)$$
(34)
(35)

$$\Pi^{\mu}{}_{\nu} \equiv \partial^{\mu}\partial_{\nu}\pi \qquad [\cdots] \equiv \text{Tr}\{\cdots\}$$

Quantum mechanically

- Galilean invariance protects the structure of the Lagrangian
- large classical non-linearities possible within EFT
- i.e., small radiative corrections and fluctuations perturbative

(Nicolis, Rattazzi, Trincherini, in preparation)

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(Nicolis, Rattazzi, Trincherini, in preparation)

higher-derivative Lagrangian with healthy two-derivative eom

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higher-derivative Lagrangian with healthy two-derivative eom

classical non-linear solutions inside the EFT $(\sim GR)$

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(Nicolis, Rattazzi, Trincherini, in preparation)

higher-derivative Lagrangian with healthy two-derivative eom

 \circ classical non-linear solutions inside the EFT (\sim GR)

possible exception to the no go theorem

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 Convenient to promote galilean transformation + Poincare' to conformal group

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 Convenient to promote galilean transformation + Poincare' to conformal group

$$\begin{cases} \pi \to \pi + c \\ \pi \to \pi + b_{\mu} x^{\mu} \end{cases}$$

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 Convenient to promote galilean transformation + Poincare' to conformal group

$$\begin{cases} \pi(x) \to \pi(\lambda x) + \log \lambda \\ \pi(x) \to \pi(x + bx^2 - (b \cdot x)x) - 2b_{\mu}x^{\mu} \end{cases}$$

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Convenient to promote galilean transformation + Poincare' to conformal group

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i.e., promote the galileon to a dilaton

$$g_{\mu\nu} = e^{2\pi} \eta_{\mu\nu}$$

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 Convenient to promote galilean transformation + Poincare' to conformal group

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i.e., promote the galileon to a dilaton

$$g_{\mu\nu} = e^{2\pi} \eta_{\mu\nu}$$

 galiean invariant terms become conformally invariant terms (upon straightforward modifications) -- same good features Just by symmetry, "de Sitter" solution

$$e^{\pi(x)} = -\frac{1}{H_0 t}$$

Just by symmetry, "de Sitter" solution

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Spontaneous breaking

$$SO(4,2) \rightarrow SO(4,1)$$

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Just by symmetry, "de Sitter" solution

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Spontaneous breaking

$$SO(4,2) \rightarrow SO(4,1)$$

scale invariance + conservation:

$$\begin{cases} \rho = 0 \\ p = \#\frac{1}{t^4} \end{cases}$$

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Fluctuations live in a fictitious deSitter space

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Fluctuations live in a fictitious deSitter space



exactly luminal

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Fluctuations live in a fictitious deSitter space



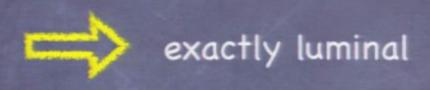
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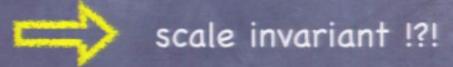


scale invariant !?!

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Fluctuations live in a fictitious deSitter space





No. They are massive $m \sim H_0$ (implied by broken time translation)

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$$\dot{H} \propto -(\rho + p) \sim 1/t^4$$

$$H \sim 1/|t|^3$$

Solution modified at late times

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$$\dot{H} \propto -(\rho + p) \sim 1/t^4$$

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Solution modified at late times

Derivatives grow... strong coupling

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$$\dot{H} \propto -(\rho + p) \sim 1/t^4$$

$$H \sim 1/|t|^3$$

Solution modified at late times

Derivatives grow... strong coupling

Reheating?

Final Thoughts

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Two consistent NEC-violating EFTs

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Two consistent NEC-violating EFTs

 Effective Lagrangian for cosmological adiabatic perturbations

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Two consistent NEC-violating EFTs

 Effective Lagrangian for cosmological adiabatic perturbations



only two possibilities (in cosmology, w/ one dof)

(Creminelli, Luty, Nicolis, Senatore 2006)

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Two consistent NEC-violating EFTs

Effective Lagrangian for cosmological adiabatic perturbations



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No superluminality, still...

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Lorentz-invariant UV completion?

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Lorentz-invariant UV completion?

 Galileon: superluminality for other solutions (w/ localized sources)

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Lorentz-invariant UV completion?

 Galileon: superluminality for other solutions (w/ localized sources)

also certain scattering amplitude too soft

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Lorentz-invariant UV completion?

 Galileon: superluminality for other solutions (w/ localized sources)

also certain scattering amplitude too soft



probably no standard Lorentz invariant UV completion

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 \circ GR: DEC (\sim NEC) = no superluminal flow

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 \circ GR: DEC (\sim NEC) = no superluminal flow

GR: NEC for matter implies CTC's

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 \circ GR: DEC (\sim NEC) = no superluminal flow

GR: NEC for matter implies CTC's

our no-go theorem: NEC (+ stability) implies superluminality for matter

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- \circ GR: DEC (\sim NEC) = no superluminal flow
- GR: NEC for matter implies CTC's
- our no-go theorem: NEC (+ stability) implies superluminality for matter
- Galileon: certain solutions violate NEC, others are superluminal

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- \circ GR: DEC (\sim NEC) = no superluminal flow
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Deep? Accidental?