

Title: Ekpyrotic/Cyclic Cosmology - Lecture 5

Date: Jul 01, 2009 09:00 AM

URL: <http://pirsa.org/09070001>

Abstract:

V Cyclic Universes

Single Field



V Cyclic Universes

Single Field

$V(\phi)$

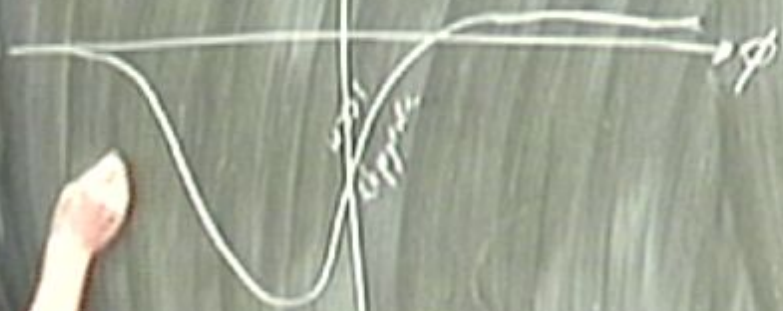


V Cyclic Universes

Single Field

$V(\phi)$

$$d = e^{Ht}$$

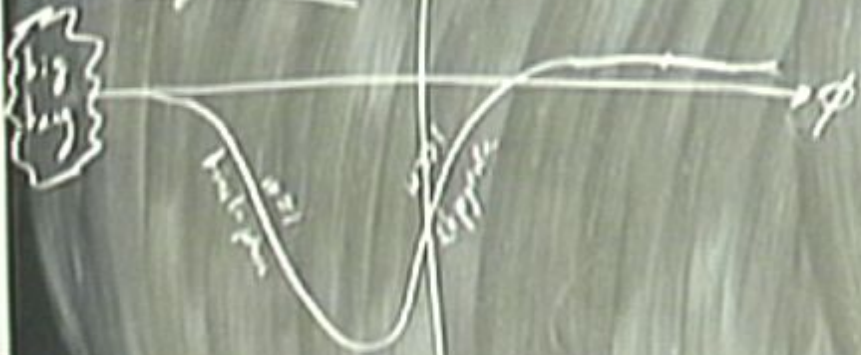


V Cyclic Universes

Single Field

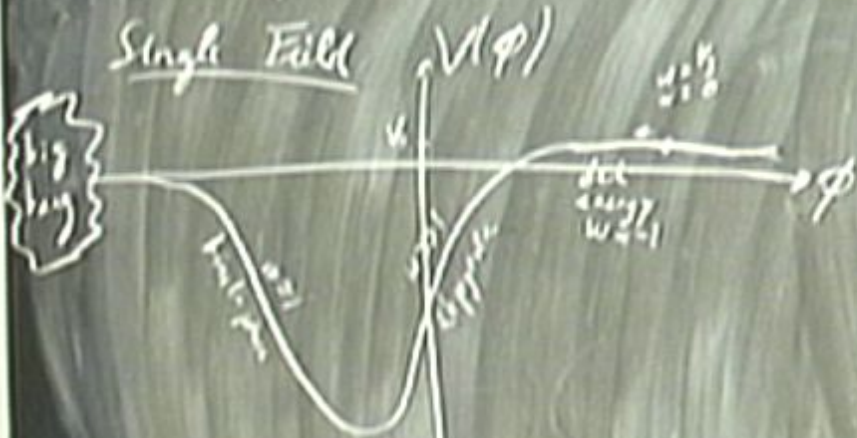
$V(\phi)$

$$d = c \cdot \phi$$



V Cyclic Universes

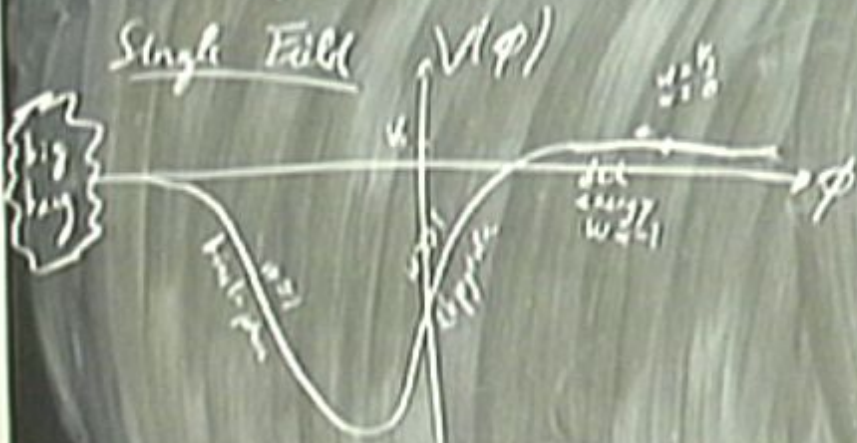
Single Field

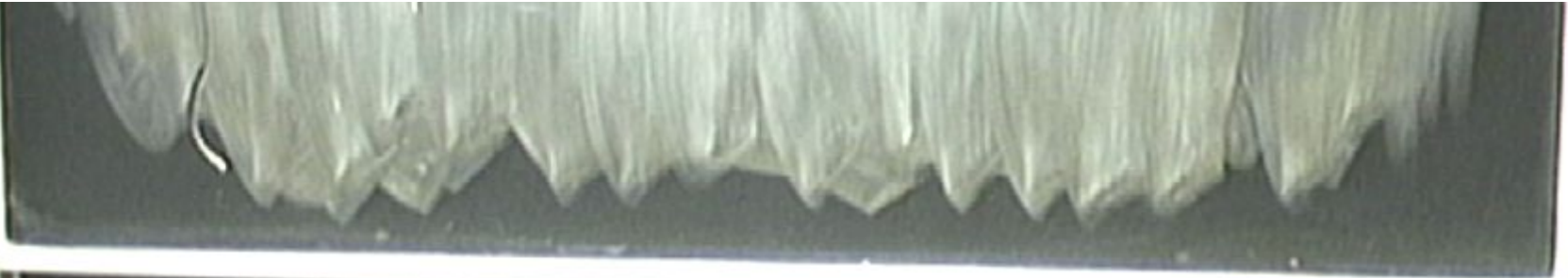


$$d = c$$

V Cyclic Universes

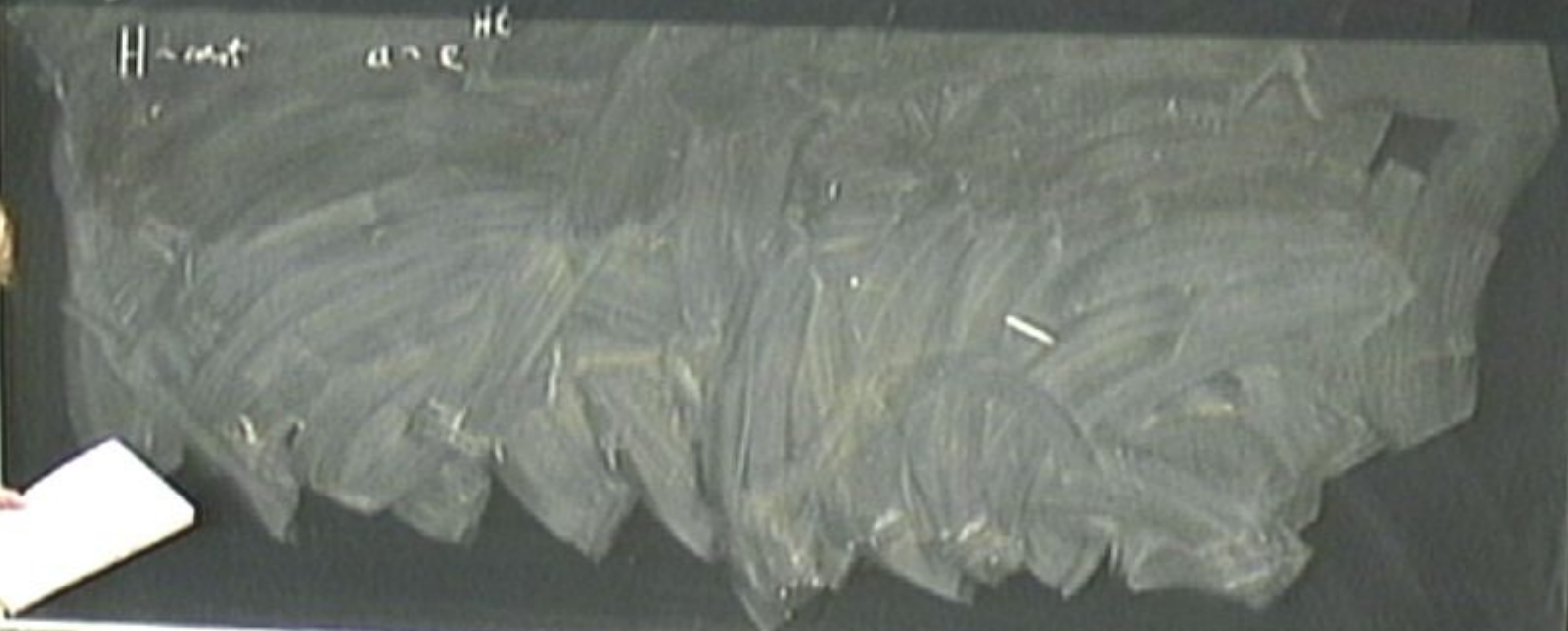
Single Field





DE

H₂CO₃ a₂e^{HC}



DE

$H = \text{const}$ $u \sim e^{Ht}$

$\frac{H_{DE} \cdot u}{a_{DE} \cdot L} = e^{N_{DE}}$

$$3H^2 = \frac{1}{2} \dot{\eta}^2 + V$$



DE

$$H \sim \text{const} \quad a \sim e^{Ht} \quad \frac{H_{DE, \text{end}}}{a_{DE, \text{ly}}} \equiv e^{N_{DE}}$$

$$3H^2 = \frac{1}{2} \dot{\eta}^2 + V \quad \text{set to } H=0 \text{ as } V < 0$$

DE

$$H \sim \text{const} \quad a \sim e^{Ht} \quad \frac{H_{DE, \text{end}}}{a_{DE, \text{end}}} \equiv e^{N_{DE}}$$

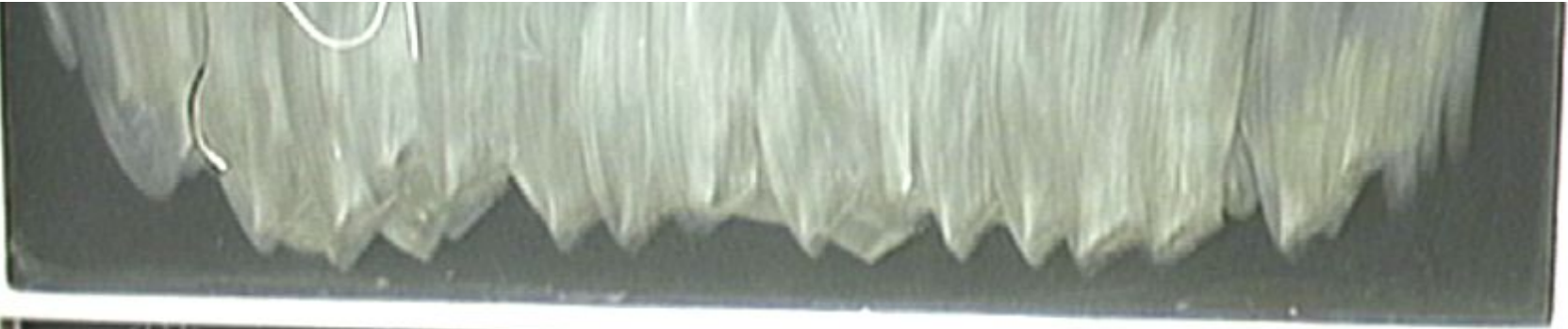
$$3H^2 = \frac{1}{2} \dot{\phi}^2 + V \quad \text{set to } H=0 \text{ as } V \ll 0$$

Elliptic

$a \sim \text{const}$

$$H \sim \frac{1}{t} \quad V \sim \frac{1}{t^2}$$

$$H \text{ grav } \ll \left| \frac{V_{\text{vac}}}{V_{\text{eff}}(\phi)} \right|^{1/2}$$



$$\frac{DE}{a_{DE} \cdot t} = c \quad H_0 \quad \frac{a_{DE=red}}{a_{DE=ly}} = e^{N_{DE}}$$

$$3H^2 = \frac{1}{2} \dot{q}^2 + V \quad \text{set to } H=0 \text{ as } V < 0$$

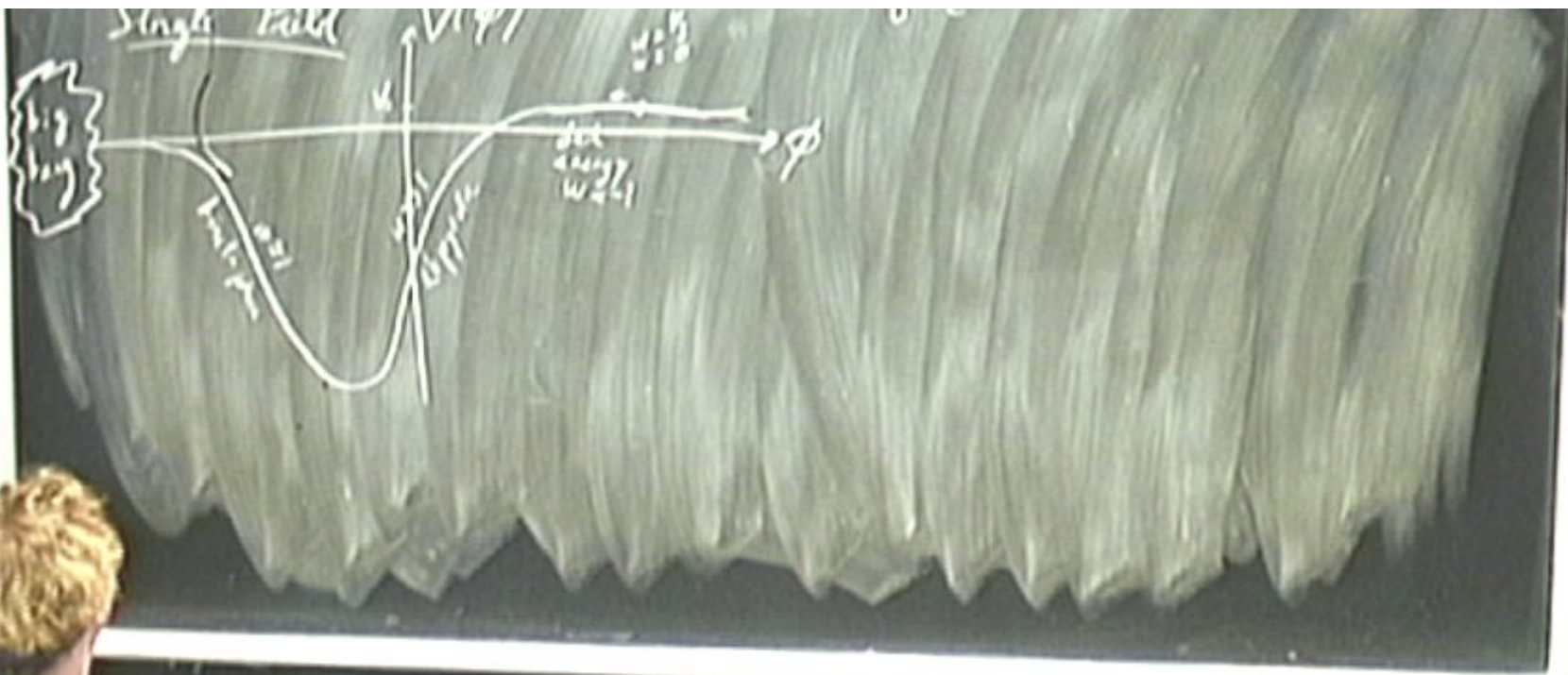
Elliptical

$a = \text{const}$

$$H = \frac{1}{t} \quad \dot{V} = \frac{1}{t^2}$$

$$|H|_{\text{grav}} \ll \left| \frac{V_{\text{kin}}}{V_{\text{pot}}} \right|^{1/2}$$

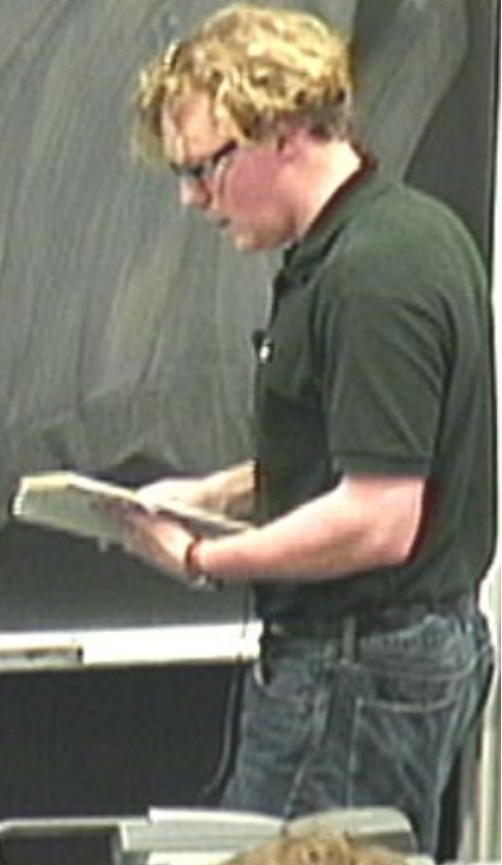




$\psi = \frac{1}{\sqrt{2\pi}}$
 $\psi = \frac{1}{\sqrt{2\pi}}$
 $|\psi|_{\text{right}} \propto \left| \frac{\psi_{\text{left}}}{\sqrt{2\pi}} \right|^{1/2}$

$$\frac{KE}{T_r}$$

begin at scalar KE - radiation equality



$\frac{KE}{T_r}$

T_r

begin at scalar $\vec{k} \cdot \vec{E}$ - radiation equality

$\rho \propto \omega^{-6}$

KE

T_r

begin at scalar $\vec{k} \cdot \vec{E}$ - radiation equality

$$\rho \propto \omega^{-6}$$

$$\frac{\alpha_{KE \cdot \text{rad}}}{\alpha_{KE \cdot \text{by}}} = \left(\frac{\rho_{\text{rad}}}{\rho_{\text{by}}} \right)^{1/2}$$

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TORONTO

$\frac{KE}{T_r}$

T_r

begin at scalar KE - radiation equality

$$\rho \propto a^{-6}$$

$$\frac{\alpha_{KE-rad}}{\alpha_{KE-by}} = \left(\frac{\rho_{KE-by}}{\rho_{KE-rad}} \right)^{1/2}$$

$$a \propto t^{1/3}$$

$$\rho \propto t^{-2}$$

$\frac{KE}{T_r}$

T_r

try at scalar KE - radiation equality

$$\rho \propto \alpha^{-6}$$

$$\frac{\alpha_{KE-rad}}{\alpha_{sc-by}} = \left(\frac{\rho_{sc-by}}{\rho_{KE-rad}} \right)^{1/6}$$

$$\rho_{sc-by} \sim |N_{rad}|$$

$$\rho_{sc-rad} \sim T_r^6$$

$$\alpha \sim t^{-1/3}$$

$$\rho \sim t^{-2}$$

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$\frac{KE}{Tr}$

Tr

temp of scalar KE - radiation equality

$$\rho \propto \alpha^{-6}$$

$$\frac{\alpha_{KE-out}}{\alpha_{KE-in}} = \left(\frac{\rho_{KE-in}}{\rho_{KE-out}} \right)^{1/6}$$

$$\rho_{KE-in} \sim |V_{in}|^4$$

$$\rho_{KE-out} \sim Tr^6$$

$$a \sim (1/\lambda)^3$$

$$\rho \sim \frac{1}{\lambda^6}$$

$$\gamma_{KE} \approx \ln \frac{|V_{in}|^4}{Tr}$$



EXIT
EVACUATION
DIRECTION
ARROWS

KE

T_r

temp of scalar KE-radiation equality

$$\rho \propto a^{-3}$$

$$\frac{\alpha_{KE-out}}{\alpha_{sc-by}} = \left(\frac{\rho_{sc-by}}{\rho_{KE-out}} \right)^{1/2}$$

$$\rho_{sc-by} \sim |V_{in}|$$

$$\rho_{sc-out} \sim T_r^4$$

$$a \sim (t)^{2/3}$$

$$\rho \sim t^{-6}$$

$$\gamma_{KE} \approx \ln \frac{|V_{in}|}{T_r}$$

$$a \text{ given by } e^{2/3 \gamma_{sc}}$$



$$\begin{aligned}
 & \rho = \frac{1}{T} \\
 & H = \frac{1}{T} \ln \dots \\
 & \gamma_{Ka} = \dots \\
 & \text{a gram by } e^{4\gamma_{Ka}} \\
 & \frac{H_{\text{out}}}{H_{\text{in}}} = e^{-2\gamma_{Ka}}
 \end{aligned}$$



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radiation (2 matter)

$$a \propto \frac{1}{T^4}$$

a grows by $\frac{T_r}{T_0} \approx e^{\text{Hrad}}$
temp of CMB today

redshift (& matter)

$$a \propto \frac{1}{z}$$

$$a \sim t^{2/3}$$

$$H = \frac{1}{2t} \propto a^{-2}$$

a grows by $\frac{T_r}{T_0} = e^{H_{\text{rad}} t}$
temp of CMB today

H shrinks by $e^{-2H_{\text{rad}} t}$

radiation (2 matter)

$$a \propto \frac{1}{T}$$

$$a \sim t^{1/2}$$

$$H = \frac{1}{2t} \propto a^{-2}$$

a grows by $\frac{T_r}{T_0} \approx e^{H_{rad} t}$
temp of CMB today

H shifts by $e^{-2H_{rad} t}$

net evolution

a grows by $e^{H_{SE} t}$

radiation (& matter)

$$a \propto \frac{1}{T}$$

$$a \sim t^{1/2}$$

$$H = \frac{1}{2t} \propto a^{-2}$$

$$a \text{ grows by } \frac{T_r}{T_0} \approx e^{H_{\text{rad}} t}$$

temp of CMB today

$$H \text{ shrinks by } e^{-2H_{\text{rad}} t}$$

net evolution

$$a \text{ grows by } e^{H_{SE} + \frac{2}{3} H_{IC} + H_{\text{rad}} t}$$

$$\frac{H_1}{H_0} = e^{H_{\text{ref}} - 2\gamma_{\text{ref}} - 2H_{\text{rad}}}$$

$$= \left| \frac{V_{\text{min}}}{V_{\text{ref-ty}}} \right|^{1/2} \frac{T_r^2}{|V_{\text{min}}|^{1/2}}$$

$$\frac{H_1}{H_0} = e^{H_{\text{el}} - 2\gamma_{\text{re}} - 2H_{\text{rad}}}$$

$$= \left| \frac{V_{\text{min}}}{V_{\text{el-by}}} \right|^{1/2} \frac{T_r^2}{|V_{\text{min}}|^{1/2}} \frac{T_0^2}{T_r^2} = \frac{T_0^2}{V_{\text{el-by}}^{1/2}}$$

$V_{\text{el-by}} \approx V_0$

4

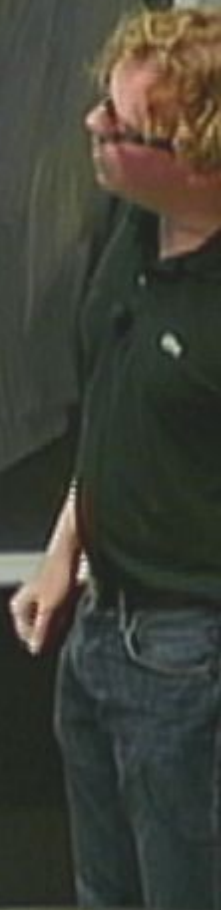
$$\frac{H_1}{H_0} = e^{H_{el} - 2\gamma_{ke} - 2H_{ad}}$$

$$= \left| \frac{V_{min}}{V_{el-by}} \right|^{1/2} \frac{T_r^2}{|V_{min}|^{1/2}} \frac{T_0^2}{T_r^2} = \frac{T_0^2}{V_{el-by}^{1/2}}$$

$$= 1$$

$$V_{el-by} \approx V_0 \approx T_0^4$$

4



$$\frac{H_i}{H_0} = e^{N_{ek} - 2\gamma_{kc} - 2N_{red}}$$

$$= \left| \frac{V_{min}}{V_{ek-by}} \right|^{1/2} \frac{T_r^2}{|V_{min}|^{1/2}} \frac{T_0^2}{T_r^2} = \frac{T_0^2}{V_{ek-by}^{1/2}}$$

$$= 1 \quad \text{cyclic}$$

$$V_{ek-by} \approx V_0 \approx T_0^4$$

$$N_{ek} = 2(\gamma_{kc} + N_{red})$$

4

$$H_0 = \left(\frac{V_{\text{min}}}{V_{\text{ek-by}}} \right)^{1/2} \frac{T_r^2}{|V_{\text{min}}|^{1/2}} \frac{T_0^2}{T_r^2} = \frac{T_0^2}{V_{\text{ek-by}}^{1/2}}$$

$V_{\text{ek-by}} = V_0 = \sqrt{T_0^4}$

$$= 1 \quad \text{cyclic}$$

$$N_{\text{ek}} = 2 (y_{\text{hc}} + N_{\text{rad}})$$

N_{DE} free parameter

4

typical values $|V_{\text{min}}| \sim (10^{-3} M_{\text{pl}})^4 = (10^{15} \text{ GeV})^4$

$T_0 \approx 10^{-3} \text{ eV}$

typical values $N_{\text{min}} \sim (10^{-3} M_{\text{pl}})^4 \approx (10^{15} \text{ GeV})^4$

$T_0 \approx 10^{-3} \text{ eV}$ $T_r \sim (10^{-6} M_{\text{pl}}) = 10^{14} \text{ GeV}$



typical values $N_{\text{min}} \sim (10^{-3} M_{\text{pl}})^4 = (10^{15} \text{ GeV})^4$

$T_0 \sim 10^{-3} \text{ eV}$ $T_r \sim (10^{-6} M_{\text{pl}}) = 10^{12} \text{ GeV}$

$\gamma_{\text{RE}} \sim 7 \rightarrow N_{\text{rel}} \sim 124$

$N_{\text{rad}} \sim 55$

$N_{\text{DE}} \approx 0 \Rightarrow$ a grav by $e^{60} \approx 10^{26}$ per cycle

typical values $N_{\text{min}} \sim (10^{-3} M_{\text{pl}})^4 = (10^{15} \text{ GeV})^4$

$T_0 = 10^{-3} \text{ eV}$ $T_r \sim (10^{-6} M_{\text{pl}}) = 10^{12} \text{ GeV}$

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current size of observable universe $\sim 10^{26} \text{ m}$

typical values $N_{\text{min}} \sim (10^{-3} M_{\text{pl}})^4 = (10^{15} \text{ GeV})^4$

$T_0 \sim 10^{-3} \text{ eV}$ $T_r \sim (10^{-6} M_{\text{pl}}) = 10^{12} \text{ GeV}$

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$N_{\text{DE}} \approx 0 \Rightarrow$ a gross by $e^{60} \approx 10^{26}$ per cycle

current size of observable universe $\sim 10^{26} \text{ m}$

one cycle ago $\sim 1 \text{ m}$



how to evade problems of oscillating model :

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- $\Omega_{m,0} < \Omega_{crit}$ potential causes recollapse

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↳ also is DE

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- Chaotic mixmaster avoided due to ekpyrosis

how to evade problems of oscillating model :

- $\Omega_m < \Omega_{crit}$ potential causes recollapse
↳ also is DE
- chaotic mixmaster avoided due to ekpyrosis
- entropy diluted each cycle

- $\Omega_m < \Omega_{crit}$ potential curves recollapse

↳ also is DE

- chaotic mixmaster avoided due to ekpyrosis

- entropy diluted each cycle

Single $\{ \text{cyclic model} \}$ → attractor → stable



$H \sim \text{const}$

$a \sim e$

$$\frac{\dot{a}}{a} = H = \text{const} \Rightarrow a \sim e^{Ht}$$

$$3H^2 = \frac{1}{2} \dot{\phi}^2 + V$$

set to $H=0 \Rightarrow V < 0$

Ekpyrotic

$a \sim \text{const}$

$$H = \frac{1}{t}$$

$$V = \frac{1}{t^2}$$

$$\|H\|_{\text{grav}} \ll \left| \frac{V_{\text{kin}}}{V_{\text{eff}}} \right|^{1/2}$$

Two fields & the Phoenix universe

$$\delta S \propto \frac{1}{T}$$

Two fields & the Phoenix universe

$$S \propto \frac{1}{t}$$



Two fields & the Phoenix universe

$$\delta S \propto \frac{1}{t}$$

Work at t_{end} $\delta S^2 \propto \delta t^2$



Two fields & the Phoenix universe

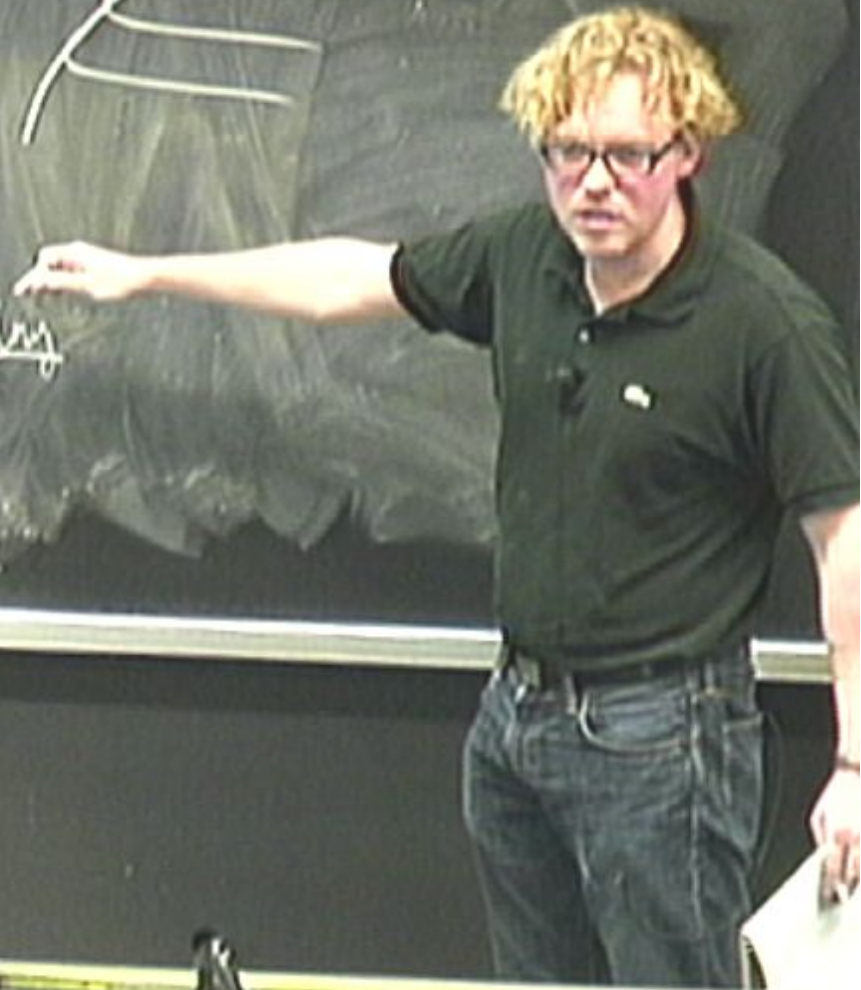
$$\delta S \propto \frac{1}{t}$$

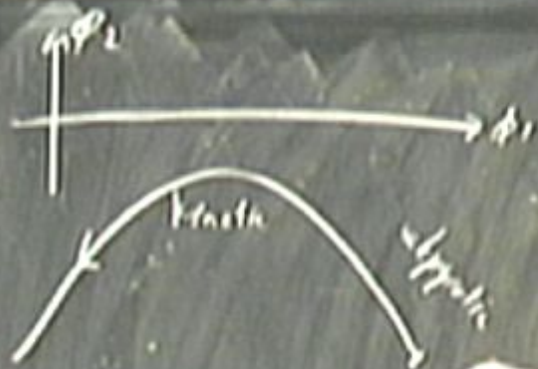
What at t_{end} $\delta S^2 \propto \sigma^2$

$$|\delta S(t_{\text{end}})| \leq \frac{1}{t}$$

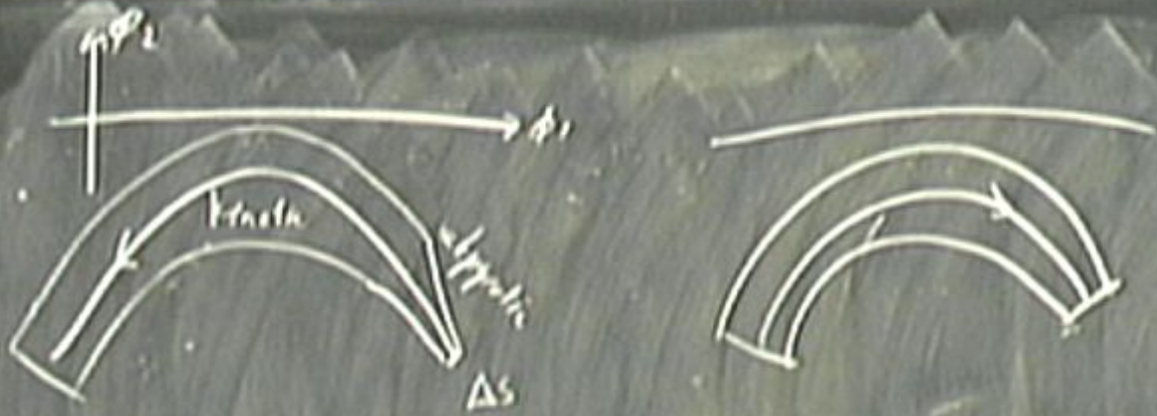
$$\rightarrow |\delta S(t_{\text{end}})| \leq \frac{1}{t} e^{-H_{\text{eff}} t}$$

tiny

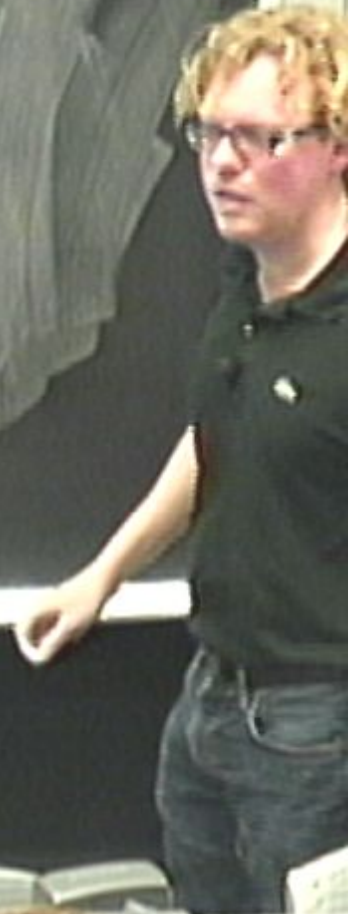








initial range $\Delta s \rightarrow$ final range $\sim e^{-Hx} \Delta s$



CAUTION
 SIGNIFICANT
 JERKING
 MOVEMENT



initial range $\Delta s \rightarrow$ final range $\sim e^{H\Delta s} \Delta s$

if outside of allowed range





initial range Δs \rightarrow final range $\sim e^{Ht} \Delta s$

if outside of allowed range
don't flatten enough

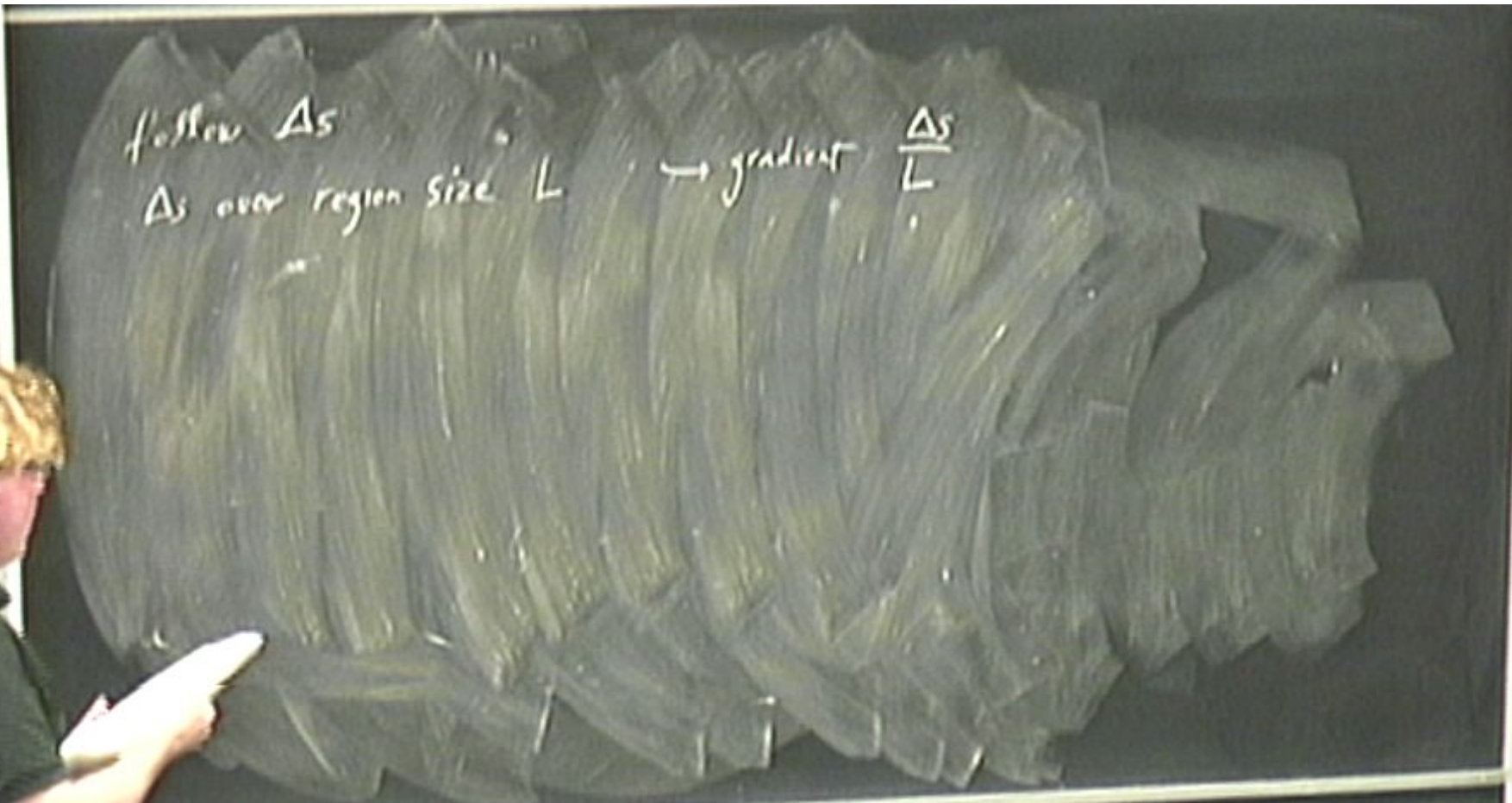
\rightarrow chaotic, mixmaster \rightarrow black holes / cycling stops



follow Δs

Δs over region size L

→ gradient $\frac{\Delta s}{L}$



follow ΔS

ΔS over region size L

→ gradient $\frac{\Delta S}{L}$

over one cycle

$$\Delta S \rightarrow e^{N_{\text{el}}} \Delta S$$

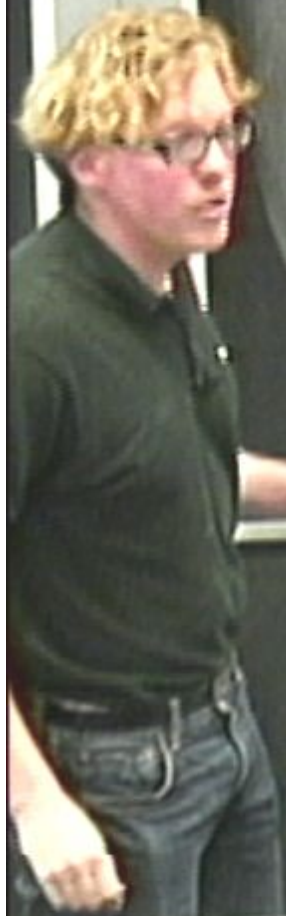
$$L \rightarrow L e^{N_{\text{DO}} + \frac{2}{3} N_{\text{O}} + N_{\text{wall}}}$$

follow Δs

Δs over region size $L \rightarrow$ gradient $\frac{\Delta s}{L}$

over one cycle $\Delta s \rightarrow e^{N_{\text{cell}}} \Delta s$
 $L \rightarrow L e^{N_{\text{cell}} + \frac{1}{2} \gamma_{\text{cell}} + N_{\text{cell}}}$

final gradient $\frac{\Delta s}{L} e$



follow ΔS

ΔS over region size $L \rightarrow$ gradient $\frac{\Delta S}{L}$

over one cycle $\Delta S \rightarrow e^{N_{DE}} \Delta S$
 $L \rightarrow L e^{N_{DE} + \frac{2}{3} N_{DE} + N_{DE}}$

final gradient = $\frac{\Delta S}{L} e^{60 - N_{DE}}$

follow ΔS

ΔS over region size $L \rightarrow$ gradient $\frac{\Delta S}{L}$

over one cycle $\Delta S \rightarrow e^{N_{DE}} \Delta S$

$L \rightarrow L e^{N_{DE} + \frac{2}{3} \gamma_{DE} + N_{null}}$

final gradient = $\frac{\Delta S}{L} e^{60 - N_{DE}}$

if $N_{DE} > 60$



follow ΔS

ΔS over region size $L \rightarrow$ gradient $\frac{\Delta S}{L}$

over one cycle $\Delta S \rightarrow e^{N_{DE}} \Delta S$
 $L \rightarrow L e^{N_{DE} + \frac{2}{3} N_{DE} + N_{DE}}$

final gradient = $\frac{\Delta S}{L} e^{60 - N_{DE}}$

if $N_{DE} > 60$ enlarge region with the right ICs

follow ΔS
 ΔS over region size $L \rightarrow$ gradient $\frac{\Delta S}{L}$

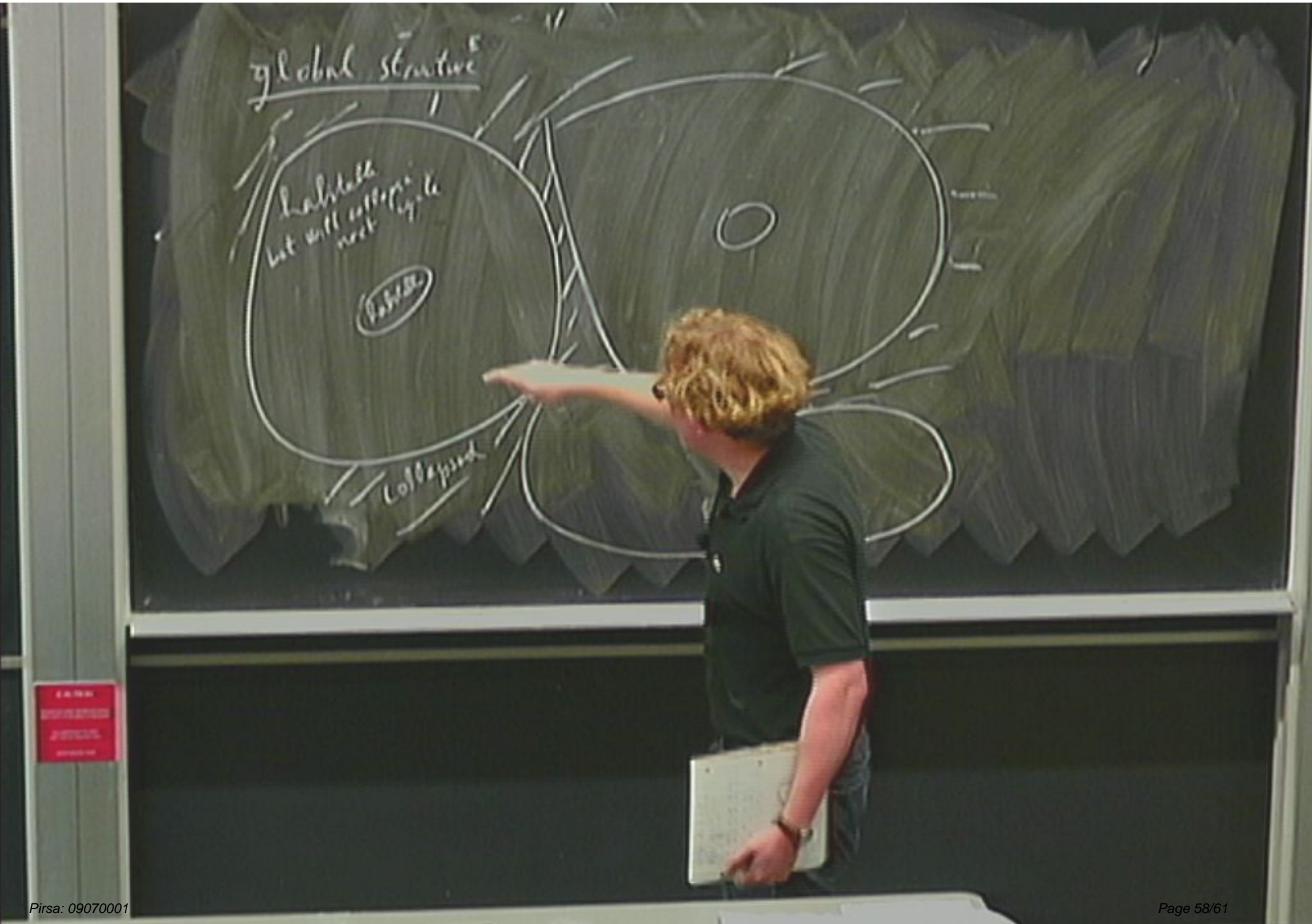
over one cycle $\Delta S \rightarrow e^{N_{DE}} \Delta S$
 $L \rightarrow L e^{N_{DE} + \frac{2}{3} \gamma_{DE} + N_{void}}$

final gradient = $\frac{\Delta S}{L} e^{60 - N_{DE}}$

if $N_{DE} > 60$ enlarge region with the right ICs

DE must last > 600 billion yrs





global structure

habitate
but will collapse &
reek

habitate

collapse

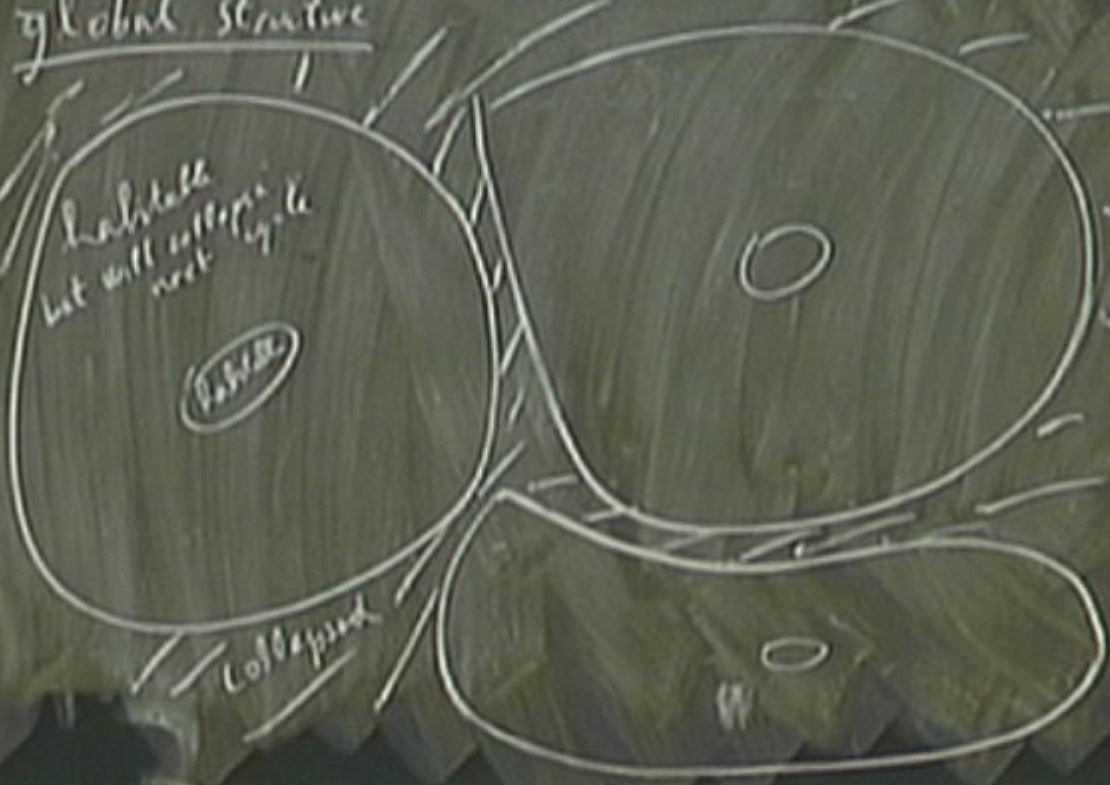
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global structure

habitate
but will collapse &
revert to

habitate

collapse



$$10^{-3} \text{ eV}$$

$$T_r \sim (10^{-6} \text{ MeV}) = 10^{12} \text{ GeV}$$

$$g_{\text{rel}} \sim 7$$

$$\rightarrow N_{\text{rel}} \sim 124$$

$$N_{\text{rad}} \sim 55$$

$$\frac{\Delta V}{V} = \frac{V_{\text{ISS}} (\delta S)^2}{V} = \epsilon (\delta S)^2$$

$$N_{\text{DE}} \approx 0 \Rightarrow \alpha \text{ grows by } e^{60} \approx 10^{26} \text{ per cycle}$$

current size of observable universe $\sim 10^{26} \text{ m}$

one cycle ago $\sim 1 \text{ m}^3$

$$(10^{-6} \text{ Mpc}) = 10^{12} \text{ GeV}$$

$$\frac{\Delta V}{V} = \frac{V_{,SS}(\delta S)^2}{V} = \epsilon (\delta S)^2 \leq \frac{1}{3}$$

124

$$\delta S < \frac{1}{\sqrt{3}}$$

by $e^{60} = 10^{26}$ per cycle

size universe $\sim 10^{26}$ m