

Title: Inflationary Cosmology - Lecture 3

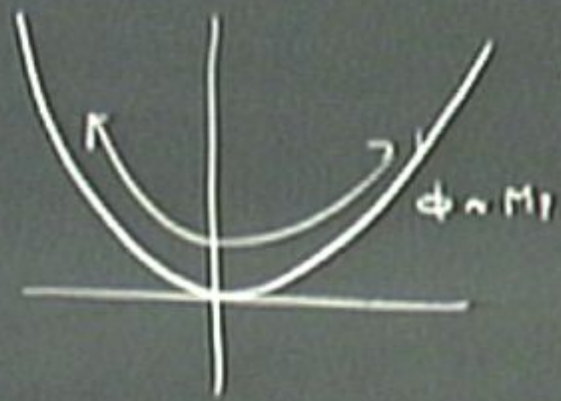
Date: Jun 30, 2009 11:30 AM

URL: <http://pirsa.org/09060087>

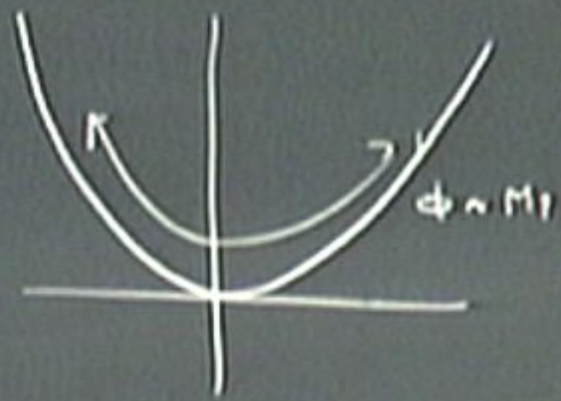
Abstract:

Reheating

Reheating

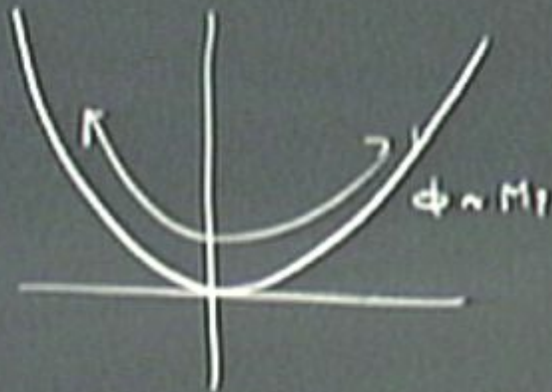


Reheating



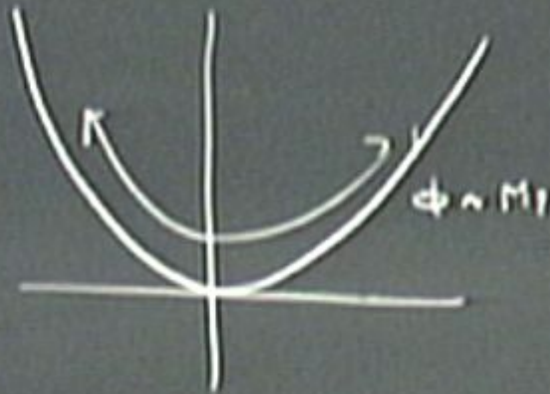
$$m_{\phi} \rightarrow \frac{\epsilon_{\phi}}{m} \rightarrow \frac{1}{2m} \left(\dot{\phi}^2 + m^2 \phi^2 \right) \approx \frac{1}{2} m_{\text{eff}}^2$$

Reheating



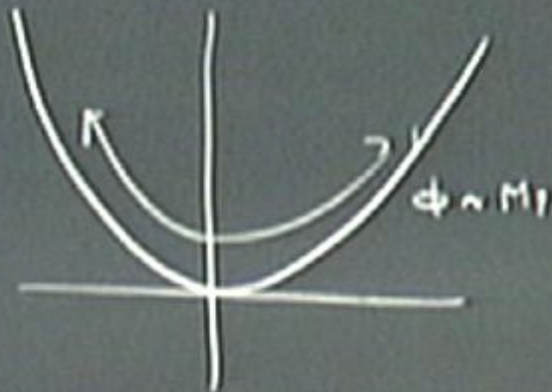
$$m_{\phi} = \frac{\epsilon_{\phi}}{m} = \frac{1}{2m} \left(\dot{\phi}^2 + m^2 \phi^2 \right) \approx \frac{1}{2} m \bar{\phi}^2 \quad 10^{92} \text{ cm}^{-3}$$

Reheating



$$m_{\phi} = \frac{\epsilon_{\phi}}{m} \rightarrow \frac{1}{2m} \left(\dot{\phi}^2 + m^2 \phi^2 \right) \approx \frac{1}{2} m \frac{H^2}{m^2} \approx 10^{92} \text{ cm}^{-3}$$

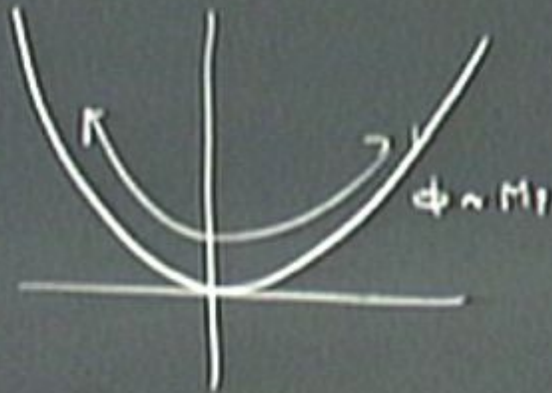
Reheating



$$m_{\phi} = \frac{E_{\phi}}{m} = \frac{1}{2m} \left(\dot{\phi}^2 + m^2 \phi^2 \right) \approx \frac{1}{2} m \bar{\phi}^2 \quad 10^{92} \text{ cm}^{-3}$$

□ RH

Reheating

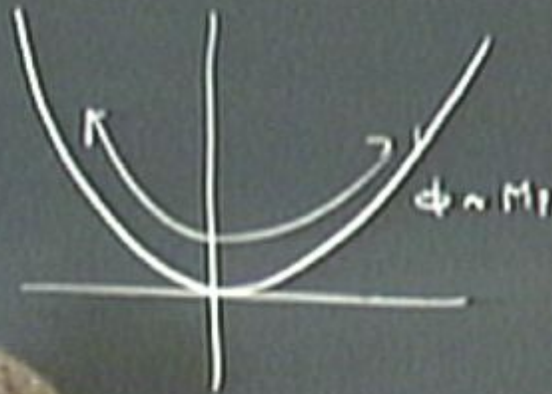


$$m_\phi = \frac{\epsilon_\phi}{m} = \frac{1}{2m} (\dot{\phi}^2 + m^2 \phi^2) \approx \frac{1}{2} m \Phi^2 \quad 10^{92} \text{ cm}^{-3}$$

$$\Gamma \approx H$$

$$g_* T_{\text{re}}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_p^2$$

Reheating



$$m_\phi = \frac{c}{\lambda}$$

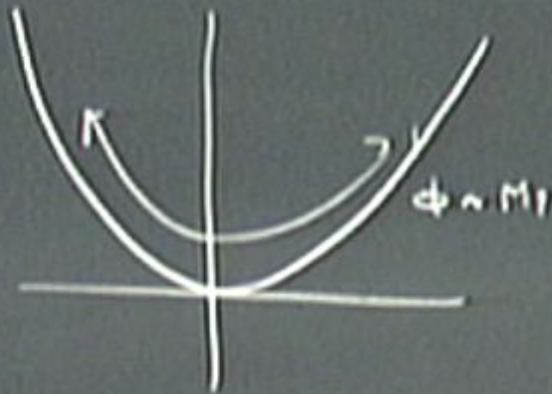
$$\left(\dot{\phi}^2 + m^2 \phi^2 \right) \geq \frac{1}{2} m \Phi^2$$

$$10^{92} \text{ cm}^{-3}$$

$$T_{\text{re}} \sim H^2 M_p^2 \sim \Gamma^2 M_p^2$$

$$T_{\text{re}} = \sqrt{M_p \Gamma g_*^{-1/2}}$$

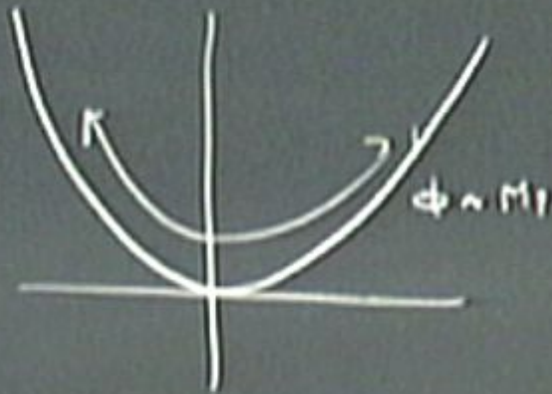
Reheating



$$m_\phi = \frac{\epsilon_\phi}{m} = \frac{1}{2m} (\dot{\phi}^2 + m^2 \phi^2) \approx \frac{1}{2} m \bar{\phi}^2 \quad 10^{9e} \text{ cm}^{-3}$$

$$g_* T_{re}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_r^2 \quad T_{re} = \sqrt{M_r \Gamma g_*^{-1/4}}$$

(P) Reheating



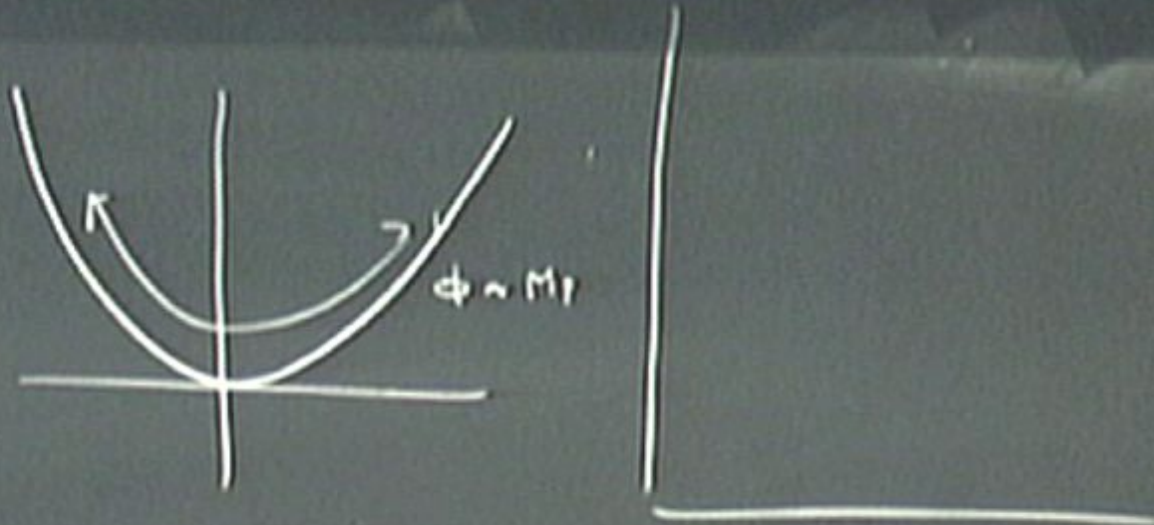
$$m_\phi = \frac{\epsilon_\phi}{m} = \frac{1}{2m} (\dot{\phi}^2 + m^2 \phi^2) \approx \frac{1}{2} m \dot{\phi}^2 \quad 10^{92} \text{ cm}^{-3}$$

$$\Gamma \approx H$$

$$g_* T_{re}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_p^2$$

$$T_{re} = \sqrt{M_p \Gamma g_*^{-1/4}}$$

(P) Reheating

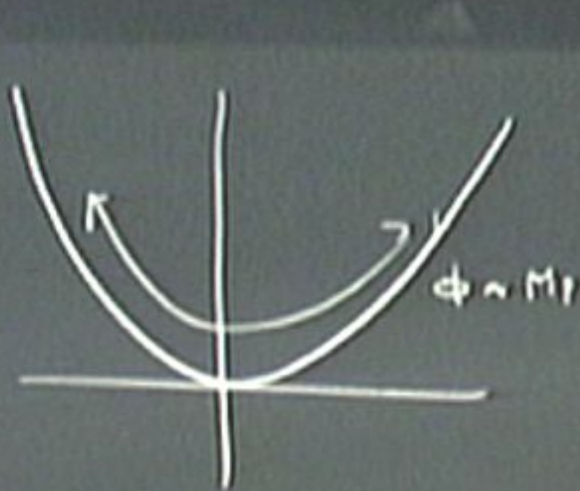


$$m_\phi = \frac{\epsilon_\phi}{m} = \frac{1}{2m} \left(\dot{\phi}^2 + m^2 \phi^2 \right) \approx \frac{1}{2} m \Phi^2 \quad 10^{92} \text{ cm}^{-3}$$

$\Gamma \approx H$

$$g_* T_{re}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_p^2 \quad T_{re} = \sqrt{M_p \Gamma g_*^{-1/4}}$$

(P) Reheating



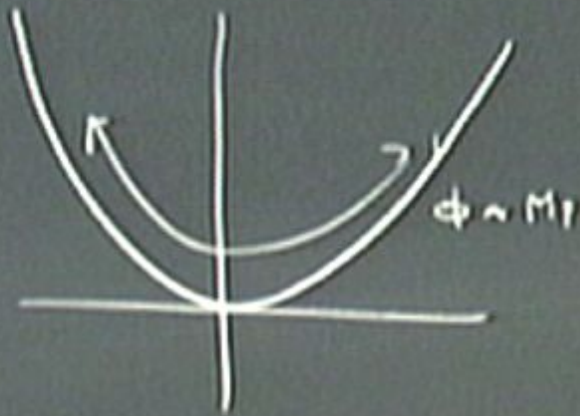
$$m_\phi = \frac{\epsilon_\phi}{m} = \frac{1}{2m} (\dot{\phi}^2 + m^2 \phi^2) \approx \frac{1}{2} m \Phi^2 \quad 10^{92} \text{ cm}^{-3}$$

$\Gamma \approx H$

$$g_* T_{re}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_p^2$$

$$T_{re} = \sqrt{M_p \Gamma g_*^{-1/4}}$$

(P) Reheating



- Conservation of ζ
-

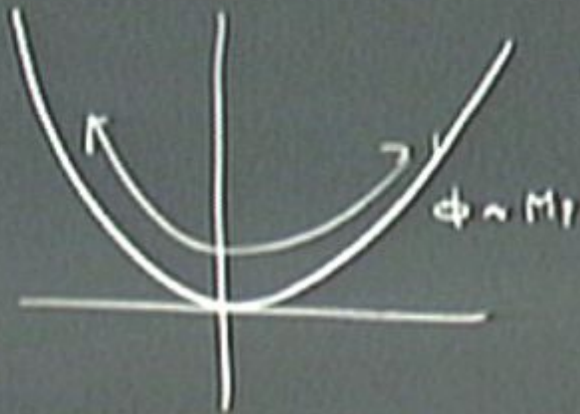
$$m_\phi = \frac{\epsilon_\phi}{m} \rightarrow \frac{1}{2m} (\dot{\phi}^2 + m^2 \phi^2) \approx \frac{1}{2} m \Phi^2 \quad 10^{92} \text{ cm}^{-3}$$

Γ_{RH}

$$g_* T_{RH}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_p^2$$

$$T_{RH} = \sqrt{M_p \Gamma g_*^{-1/4}}$$

(P) Reheating



• Conservation of ζ

$$10^{16} \text{ GeV} \longleftrightarrow 1 \text{ MeV}$$

$$m_\phi = \frac{\epsilon_\phi}{m} = \frac{1}{2m} (\dot{\phi}^2 + m^2 \phi^2) \approx \frac{1}{2} m \dot{\phi}^2 \quad 10^{92} \text{ cm}^{-3}$$

$\Gamma \gg H$

$$g_* T_{re}^4 \approx H^2 M_p^2 \approx \Gamma^2 M_p^2$$

$$T_{re} = \sqrt{M_p \Gamma g_*^{-1/4}}$$

RUCB

• Classification of infl. models

$$\frac{1}{\sqrt{\epsilon}} \frac{H}{M_p} \approx 10^{-5}$$

$$\frac{H}{M_p}$$

GW signal

① High energy models

$$V \propto \phi^m$$

$$\epsilon, \eta = \left(\frac{H}{\dot{\phi}} \right)^2$$

$$\delta \phi$$

DL/CB

• Classification of infl. models

$$\frac{H}{M_p} \sim 10^{-5}$$

$$\frac{H}{M_p}$$

GW signal

① High energy modes

$$V \propto \phi^m \quad \epsilon, \eta = \left(\frac{H}{\dot{\phi}}\right)^2 \sim \frac{1}{2}$$

$$\Delta \phi \gg M_H$$

DLICB

• Classification of infl. models

$$\frac{H}{M_p} \sim 10^{-5}$$

$\frac{H}{M_p}$ GW signal

① High energy models

$$V \propto \phi^n \quad \epsilon, \eta = \left(\frac{M_p}{\phi}\right)^2 \sim \frac{1}{N}$$

$$\Delta \phi \gg M_H$$

$$V \sim 10^4 \text{ GeV} \sim M_{\text{out}}$$

• Classification of infl. models

$$\frac{H}{M_p} \sim 10^{-5}$$

$$\frac{H}{M_p} \text{ GW signal}$$

① High energy models

$$V \propto \phi^m \quad \epsilon, \eta = \left(\frac{H}{M_p}\right)^2 \sim \frac{1}{N}$$

$$\Delta \phi \gg M_H$$

$$V \sim 10^V \text{ GeV} \sim M_{\text{out}}$$

② Hybrid models

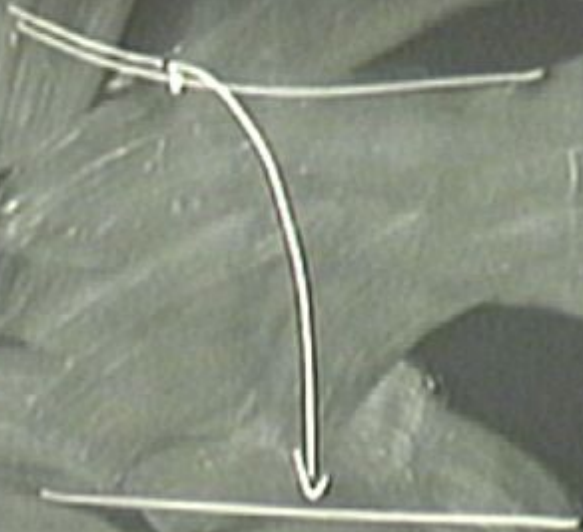
$$V = \frac{1}{2} m^2 \phi^2 + \frac{1}{4} \lambda (\gamma^2 - M^2)^2 + \frac{1}{2} \lambda' \gamma^2 \phi^2$$

$$V = \frac{1}{2} m^2 \phi^2 + \frac{1}{4} \lambda (\gamma^2 - M^2)^2 + \frac{1}{2} \lambda' \gamma^2 \phi^2$$

D3



D3



$$V = \frac{1}{2} m^2 \phi^2 + \frac{1}{4} \lambda (\gamma^2 - M^2)^2 + \frac{1}{2} \lambda' \gamma^2 \phi^2$$

$E \lll 1$



• Classification of infl. models

$$\frac{1}{\sqrt{\epsilon}} \frac{H}{M_p} \approx 10^{-5}$$

$$\frac{H}{M_p} \text{ GW signal}$$

① High energy models

$$V \propto \phi^n \quad \epsilon, \eta = \left(\frac{H}{\dot{\phi}}\right)^2 \approx \frac{1}{2}$$

$$\Delta \phi \gg M_H$$

$$V \sim 10^4 \text{ GeV} \sim M_{\text{out}}$$

② Hybrid models

$$\Delta \phi \approx \frac{\phi}{H} \approx \frac{v}{H^2} \approx \sqrt{\epsilon} M_p$$

• Classification of infl. models

① High energy models

$V \propto \phi^n$ $\epsilon, \eta = \left(\frac{M_{pl}}{\phi}\right)^2 \sim \frac{1}{N}$

$\Delta \phi \gg M_{pl}$

$V \sim 10^4 \text{ GeV} \sim M_{out}$

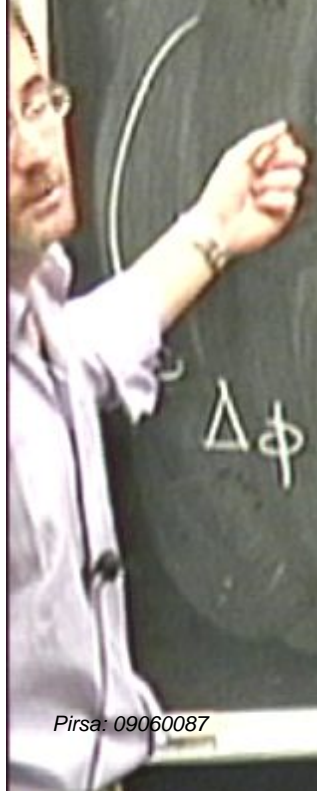
② Hybrid models

$\frac{1}{\sqrt{\epsilon}} \frac{H}{M_{pl}} \sim 10^{-5}$

$\frac{H}{M_{pl}}$ GW signal

$\Delta \phi \sim \frac{\phi}{H} \sim \frac{v}{H^2} \sim \sqrt{\epsilon} M_{pl}$

Lyth bound



DUCB

Classification of infl. models

① High energy models

$$V \propto \phi^n \quad \epsilon, \eta = \left(\frac{M_{Pl}}{\phi}\right)^2 \sim \frac{1}{2}$$

$$\Delta \phi \gg M_{Pl}$$

$$V \sim 10^4 \text{ GeV}^4 \sim M_{out}^4$$

② Hybrid models

$$\frac{1}{\sqrt{\epsilon}} \frac{H}{M_{Pl}} \sim 10^{-5}$$

$$\frac{H}{M_{Pl}}$$

GW signal

$$\Delta \phi \sim \frac{\phi}{H} \sim \frac{c}{H^2} \sim \sqrt{\epsilon} M_{Pl}$$

Lyth bound

$$\hbar \neq 0 = 1$$

- Harmonic oscillator
- Free field in dS

$$\hbar \neq 0 = 1$$

- Harmonic oscillator
- Free field in dS
- Scalar + tensor spectrum

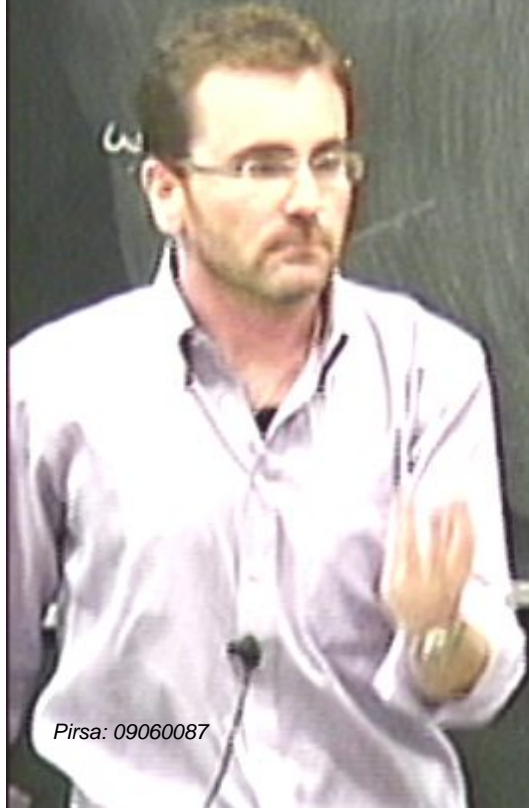
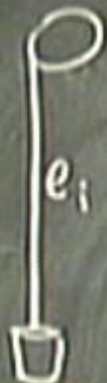
$$\hbar \neq 0 = 1$$

- Harmonic oscillator
- Free field in dS
- Scalar & tensor spectrum

$\hbar \neq 0 = 1$ Free field theory

- Harmonic oscillator
- Free field in dS
- Scalar + tensor spectrum



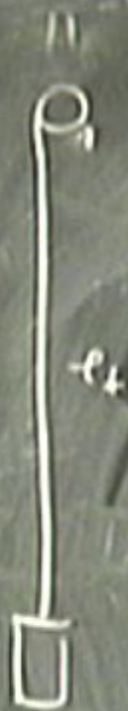
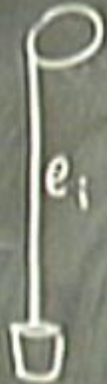




Adiabatic:

$$T \gg \omega_i^{-1}, \omega_t^{-1}$$

$$\omega = \sqrt{\frac{g}{l}}$$



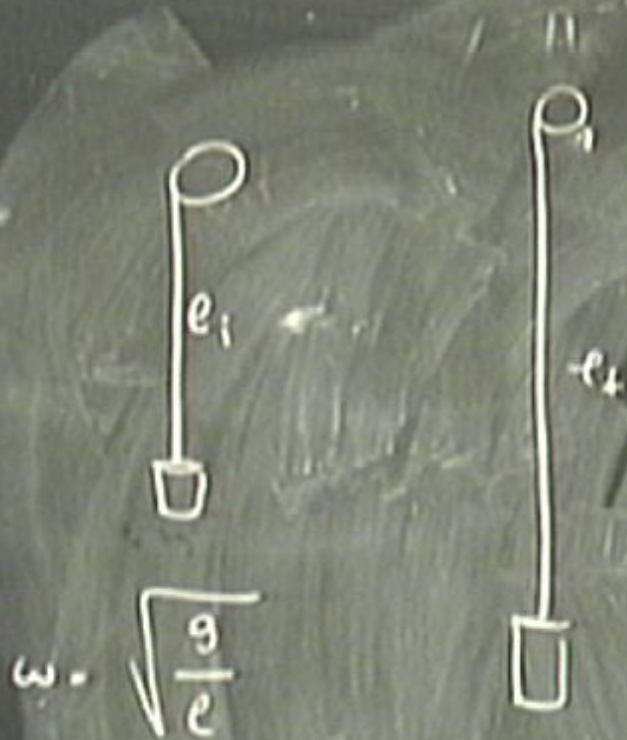
• Adiabatic

$$T \gg \omega_i^{-1}, \omega_f^{-1}$$

Vacuum stays there

• Sudden limit

$$T \ll \omega_i^{-1}, \omega_f^{-1}$$



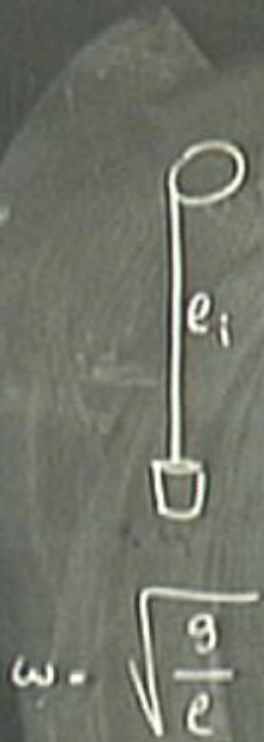
• Adiabatic:

$$T \gg \omega_i^{-1}, \omega_f^{-1}$$

Vacuum stays there

• Sudden limit

$$T \ll \omega_i^{-1}, \omega_f^{-1}$$



• Adiabatic

$$T \gg \omega_i^{-1}, \omega_f^{-1}$$

Vacuum stays there

Sudden limit

$$T \ll \omega_i^{-1}, \omega_f^{-1}$$





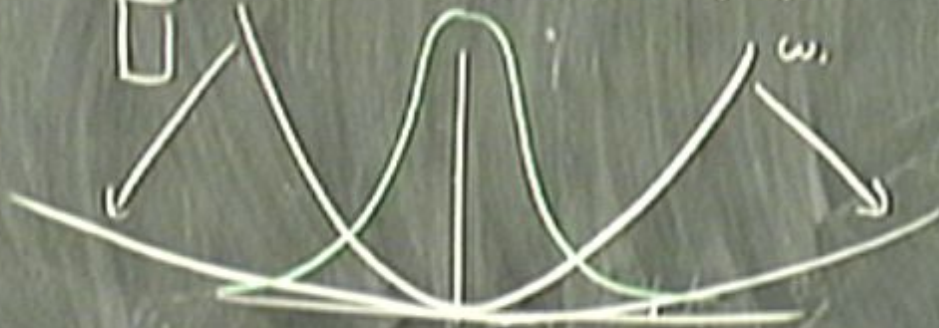
• Adiabatic:

$$T \gg \omega_i^{-1}, \omega_f^{-1}$$

Vacuum stays there

• Sudden limit

$$T \ll \omega_i^{-1}, \omega_f^{-1}$$



$$\Psi_m(x) = N_m H_m(\sqrt{\alpha} x) e^{-\frac{\alpha}{2} x^2}$$

$$\alpha = \frac{2m}{\hbar^2 E}$$

$$\Psi_m(x) = N_m H_m(\sqrt{\kappa} x) e^{-\frac{\kappa}{2} x^2} \quad \kappa = \frac{2mE}{\hbar^2}$$

Old Vacuum $\Psi^0(x) = \left(\frac{\kappa}{\pi}\right)^{1/4} e^{-\frac{\kappa}{2} x^2}$



$$\psi_m(x) = N_m H_m(\sqrt{\alpha} x) e^{-\frac{\alpha}{2} x^2} \quad \alpha = \frac{m}{\hbar^2}$$

Old Vacuum: $\psi(x) = \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-\frac{\alpha}{2} x^2} = \sum_{m=0}^{\infty} c_m \psi_m(x)$

$$\psi_m(x) = N_m H_m(\sqrt{\alpha} x) e^{-\frac{\alpha}{2} x^2} \quad \alpha = \frac{m\omega}{\hbar}$$

Old vacuum $\psi(x) = \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-\frac{\alpha}{2} x^2} = \sum_{m=0}^{\infty} c_m \psi_m(x)$

$$E_f = \sum_{m=0}^{\infty} \left(2m + \frac{1}{2}\right) \hbar \omega_f |c_m|^2 = \hbar \omega_f \left(\frac{1}{2} + \frac{(\omega_f - \omega_i)^2}{4\omega_i\omega_f} \right)$$

$$\psi_m(x) = N_m H_m(\sqrt{\alpha} x) e^{-\frac{\alpha}{2} x^2} \quad \alpha = \frac{m\omega}{\hbar v}$$

Old vacuum $\psi(x) = \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-\frac{\alpha}{2} x^2} = \sum_{m=0}^{\infty} c_m \psi_m(x)$

$$E_f = \sum_{m=0}^{\infty} \left(m + \frac{1}{2}\right) \hbar \omega_f |c_m|^2 = \hbar \omega_f \left(\frac{1}{2} + \frac{(\omega_f - \omega_i)^2}{4\omega_f \omega_i} \right)$$

$$\psi_m(x) = N_m H_m(\sqrt{k}x) e^{-\frac{\lambda}{2}x^2} \quad \lambda = \frac{m\omega}{\hbar}$$

Old Vacuum: $\psi(x) = \left(\frac{\lambda}{\pi}\right)^{1/4} e^{-\frac{\lambda}{2}x^2} = \sum_{m=0}^{\infty} c_m \psi_m(x)$

$$E_f = \sum_{m=0}^{\infty} \left(2m + \frac{1}{2}\right) \hbar \omega_f |c_m|^2 = \hbar \omega_f \left(\frac{1}{2} + \frac{(\omega_f - \omega_i)^2}{4\omega_f \omega_i} \right)$$

$\omega_f \ll \omega_i$

$$= \frac{\hbar \omega_i}{4}$$

$$\psi_m(x) = N_m H_m(\sqrt{k}x) e^{-\frac{1}{2}\alpha x^2} \quad \alpha = \frac{m\omega}{\hbar v}$$

Old Vacuum $\psi(x) = \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-\frac{\alpha}{2}x^2} = \sum_{m=0}^{\infty} c_m \psi_m(x)$

$$E_f = \sum_{m=0}^{\infty} \left(2m + \frac{1}{2}\right) \hbar \omega_f |c_m|^2 = \hbar \omega_f \left(\frac{1}{2} + \frac{(\omega_f - \omega_i)^2}{4\omega_f \omega_i} \right)$$

$\omega_f \ll \omega_i$

$$= \frac{\hbar \omega_i}{4} \gg \frac{\hbar \omega_f}{2}$$

$$\psi_m(x) = N_m H_m(\sqrt{k}x) e^{-\frac{1}{2}\alpha x^2} \quad \alpha = \frac{m\omega}{\hbar v}$$

Old Vacuum: $\psi(x) = \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-\frac{\alpha}{2}x^2} = \sum_{m=0}^{\infty} c_m \psi_m(x)$

$$E_f = \sum_{m=0}^{\infty} \left(2m + \frac{1}{2}\right) \hbar \omega_f |c_m|^2 = \hbar \omega_f \left(\frac{1}{2} + \frac{(\omega_f - \omega_i)^2}{4\omega_f \omega_i} \right)$$

$\omega_f \ll \omega_i$

$$= \frac{\hbar \omega_i}{4} \gg \frac{\hbar \omega_f}{2}$$

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \right)$$

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \sqrt{\frac{1}{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{p}_1) / \sqrt{2}$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\hat{x}_1 - i\hat{p}_1 \right)$$

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \frac{1}{\sqrt{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\dots \right) = \dots$$

Square $e^{-\frac{1}{2}\xi^2 \hat{a}^2} - \frac{1}{2}\xi^2 \hat{a}^{+2}$

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \frac{1}{\sqrt{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\dots \right) = \dots$$

Squeeze $e^{\frac{1}{2}\xi \hat{a}^\dagger - \frac{1}{2}\xi \hat{a}} = S(\xi) \quad \xi = r e^{i\theta}$

$$S = \hat{a} \cosh r - \hat{a}^\dagger e^{i\theta} \sinh r$$

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \frac{1}{\sqrt{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\quad \quad \quad \right) = \quad \quad \quad$$

Squeeze operator: $e^{\frac{1}{2}\xi^* \hat{a}^2 - \frac{1}{2}\xi \hat{a}^{\dagger 2}} = S(\xi) \quad \xi = r e^{i\theta}$

$$S^\dagger \hat{a} S = \hat{a} \cosh r - \hat{a}^\dagger e^{i\theta} \sinh r$$

$$S^\dagger \hat{a}^\dagger S = \hat{a}^\dagger \cosh r - \hat{a} e^{-i\theta} \sinh r$$

Bogolubov transformation

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \frac{1}{\sqrt{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$|\xi\rangle = S(\xi) |0\rangle$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\quad \quad \quad \right) = \quad \quad \quad$$

Squeeze operator: $e^{\frac{1}{2}\xi^* \hat{a}^2 - \frac{1}{2}\xi \hat{a}^{\dagger 2}} = S(\xi) \quad \xi = r e^{i\theta}$

$$S^\dagger \hat{a} S = \hat{a} \cosh r - \hat{a}^\dagger e^{i\theta} \sinh r$$

$$S^\dagger \hat{a}^\dagger S = \hat{a}^\dagger \cosh r - \hat{a} e^{-i\theta} \sinh r$$

Bogolubov transformation

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \frac{1}{\sqrt{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\dots \right) = \dots$$

Squeeze operator: $e^{\frac{1}{2}\zeta^* \hat{a}^2 - \frac{1}{2}\zeta \hat{a}^{\dagger 2}} = S(\zeta)$

$$S^\dagger \hat{a} S = \hat{a} \cosh \tau - \hat{a}^\dagger e^{i\theta} \sinh \tau$$

$$S^\dagger \hat{a}^\dagger S = \hat{a}^\dagger \cosh \tau - \hat{a} e^{-i\theta} \sinh \tau$$

$$|\zeta\rangle = S(\zeta) |0\rangle$$

$$\underbrace{S \hat{a} S^\dagger}_{\hat{a}_{\text{eff}}} S |0\rangle = 0$$

$$\zeta = \dots$$

Bogolubov transformation

$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \frac{1}{\sqrt{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$|\xi\rangle = S(\xi) |0\rangle$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(- \right) = -$$

$$\underbrace{S \hat{a} S^\dagger}_{\hat{a}_d} S |0\rangle = 0$$

Squeeze operator. $e^{\frac{1}{2}\xi^* \hat{a}^2 - \frac{1}{2}\xi \hat{a}^{\dagger 2}} = S(\xi)$

$$\xi = r e^{i\theta}$$

$$S^\dagger \hat{a} S = \hat{a} \cosh r - \hat{a}^\dagger e^{i\theta} \sinh r$$

$$S^\dagger \hat{a}^\dagger S = \hat{a}^\dagger \cosh r - \hat{a} e^{-i\theta} \sinh r$$

Bogolubov transformation

$$\hat{y}_1 + i\hat{y}_2 = (\hat{x}_1 + i\hat{x}_2)e^{-i\theta/2}$$

$$\hat{y}_1 + i\hat{y}_2 = (\hat{x}_1 + i\hat{x}_2)e^{-i\theta/2}$$

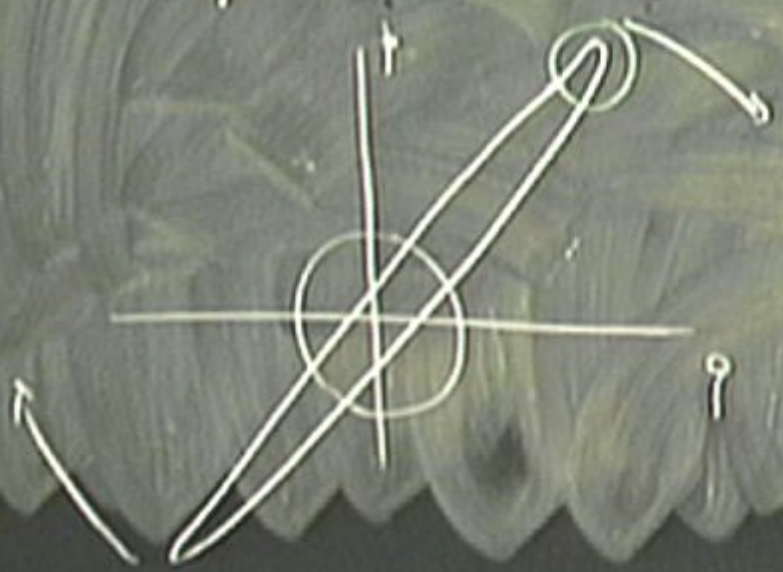
$$S^\dagger (\hat{y}_1 + i\hat{y}_2) S = \hat{y}_1 e^{-z} + i\hat{y}_2 e^z$$

$$\hat{y}_1 + i\hat{y}_2 = (\hat{x}_1 + i\hat{x}_2)e^{-i\theta/2}$$

$$S^+ \begin{pmatrix} \hat{y}_1 \\ \hat{y}_2 \end{pmatrix} S = \hat{y}_1 e^{-\tau} + i\hat{y}_2 e^{\tau}$$

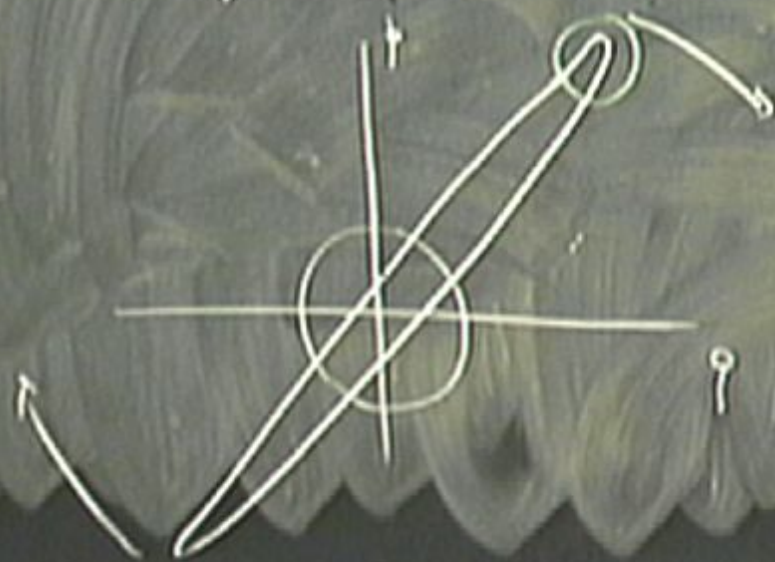
$$\hat{y}_1 + i\hat{y}_2 = (\hat{x}_1 + i\hat{x}_2) e^{-i\theta/2}$$

$$S^\dagger (\hat{y}_1 + i\hat{y}_2) S = \hat{y}_1 e^{-\tau} + i\hat{y}_2 e^{\tau}$$



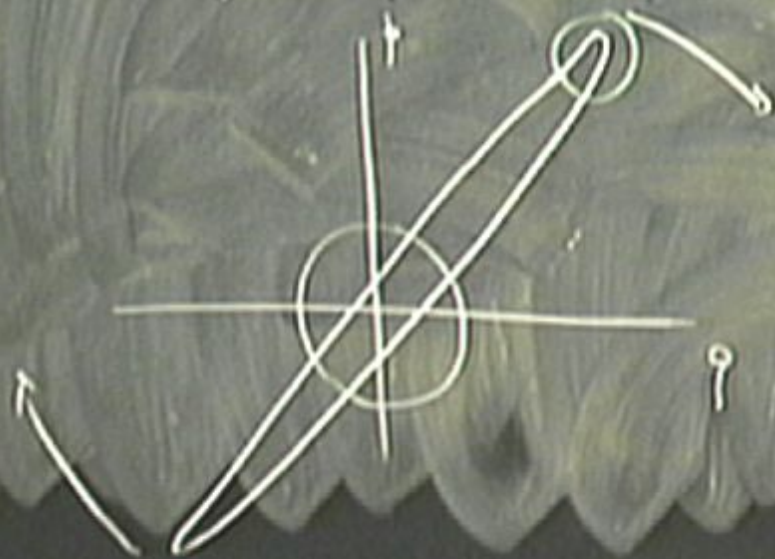
$$\hat{Y}_1 + i\hat{Y}_2 = (\hat{X}_1 + i\hat{X}_2)e^{-i\theta/2}$$

$$S^\dagger (\hat{Y}_1 + i\hat{Y}_2) S = \hat{Y}_1 e^{-\tau} + i\hat{Y}_2 e^{\tau}$$



$$\hat{y}_1 + i\hat{y}_2 = (\hat{x}_1 + i\hat{x}_2)e^{-i\theta/2}$$

$$S^+(\hat{y}_1 + i\hat{y}_2)S = \hat{y}_1 e^{-\tau} + i\hat{y}_2 e^{\tau}$$



$$\hat{a} = \frac{1}{\sqrt{2}} \left(\sqrt{\frac{m\omega}{\hbar}} \hat{x} + i \sqrt{\frac{1}{m\omega\hbar}} \hat{p} \right) = (\hat{x}_1 + i\hat{x}_2) / \sqrt{2}$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(\quad \quad \quad \right) = \quad \quad \quad$$

Squeeze operator: $e^{\frac{1}{2}\xi^* \hat{a}^2 - \frac{1}{2}\xi \hat{a}^{\dagger 2}} = S(\xi)$

$$S^\dagger \hat{a} S = \hat{a} \cosh \tau - \hat{a}^\dagger e^{i\theta} \sinh \tau$$

$$S^\dagger \hat{a}^\dagger S = \hat{a}^\dagger \cosh \tau - \hat{a} e^{-i\theta} \sinh \tau$$

$$|\xi\rangle = S(\xi) |0\rangle$$

$$\underbrace{S \hat{a} S^\dagger}_{\hat{a}} S |0\rangle = 0$$

Bogolubov transformation