

Title: Cosmic ray anomaly and metastable dark matter particles

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Abstract: TBA

Cosmic ray anomaly and metastable dark matter particles

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CR anomaly

- cosmic ray anomaly seen by PAMELA, ATIC, FERMI, HESS
- due to annihilating DM, decaying DM, astrophysics?
- if DM \Rightarrow several interesting features
 - large annihilation xsec
 - preferentially into leptons
- this talk: how the discussion changes if dark sector includes metastable particles

The scales

- χ DM, but dark sector also includes $\phi \Leftarrow$ metastable
- assume that ϕ decays through higher dim operators
- for guidance use "see-saw" scale

$$m_\nu = y^2 \frac{v^2}{\Lambda} \Rightarrow \Lambda = 6 \cdot 10^{12} \text{ GeV} \left(\frac{y}{0.1} \right)^2 \left(\frac{0.1 \text{ eV}}{m_\nu} \right)$$

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- if ϕ decays through dim 5 ops, $\frac{1}{\Lambda} O$

$$\tau \simeq 16\pi \frac{\Lambda^2}{m_\phi^3} = 0.1 \mu s \left(\frac{\Lambda}{10^{13} \text{ GeV}} \right)^2 \left(\frac{3 \text{ TeV}}{m_\phi} \right)^3$$

- possibly relevant for early universe, explain the large boost factors?

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- if ϕ decays through dim 6 ops, $\frac{1}{\Lambda^2} O$

$$\tau \simeq 16\pi \frac{\Lambda^4}{m_\phi^5} = 1.4 \cdot 10^{12} \text{ s} \left(\frac{\Lambda}{10^{13} \text{ GeV}} \right)^4 \left(\frac{3 \text{ TeV}}{m_\phi} \right)^5$$

- decays after it flies $\lambda \simeq 1.3 \text{ kpc} \left(\frac{v\gamma}{0.1} \right) \left(\frac{10^{13} \text{ GeV}}{\Lambda} \right)^4 \left(\frac{m_\phi}{3 \text{ TeV}} \right)^3$
- propagation on galactic distances

Outline

- large boost factors
- gamma rays from galactic center
- conclusions

Large cross sections

Boost factor

- if simple thermal relic then

$$\Omega_{\text{DM}} \propto 1/\langle\sigma_{Av}\rangle_F \Rightarrow \langle\sigma_{Av}\rangle_F \simeq 3 \times 10^{-26} \text{cm}^3/\text{s}$$

- cosmic ray flux

$$\Phi_{e^+} \propto \langle\sigma_{Av}\rangle \rho_{\text{DM}}^2 / m_{\text{DM}}^2$$

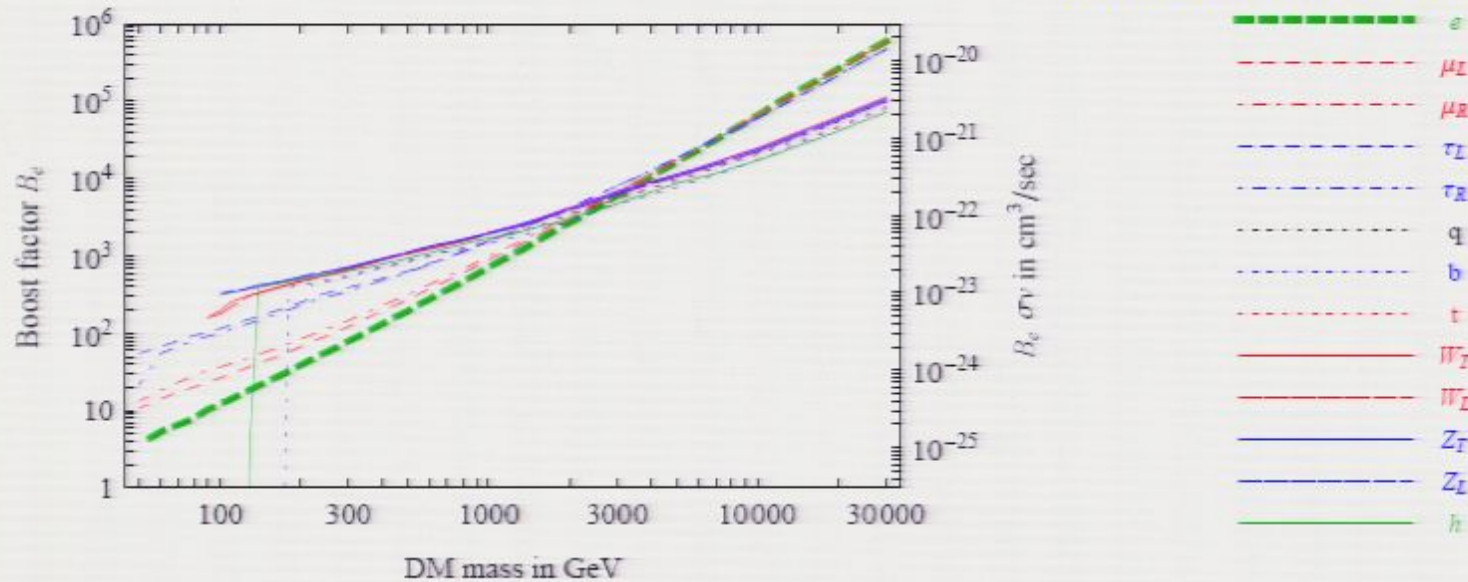
- define a "boost factor"

$$B \equiv \frac{\langle\sigma_{Av}\rangle \rho_{\text{DM}}^2}{\langle\sigma_{Av}\rangle_F (\bar{\rho}_{\text{DM}})^2}$$

Boost factors

model independent analysis for Pamela

Cirelli, Kadastik, Raidal, Strumia 2008



very large "boost factors"

Unitarity bound

- unitarity bound also imposes the upper bound on DM mass
- s-wave annihilation: $\sigma v \leq 4\pi/(vM^2)$
- this bounds the possible boost factor

$$B \leq \frac{4\pi}{vM^2 \langle \sigma v \rangle_F} = 5 \cdot 10^6 \left(\frac{1\text{TeV}}{M} \right)^2 \left(\frac{10^{-3}}{v} \right)$$

- from model indep. bound on previous slide for leptonic decays

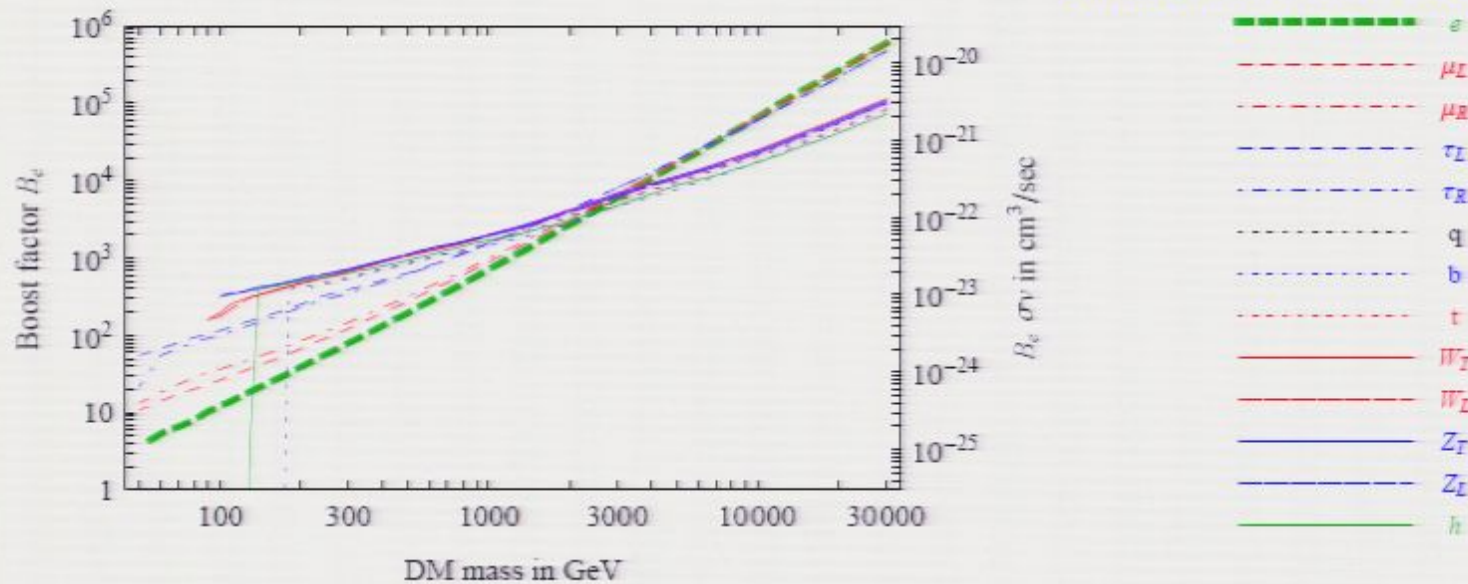
$$M \leq 8 - 9 \text{ TeV}$$

($M \leq 12\text{TeV}$ for hadronic decays)

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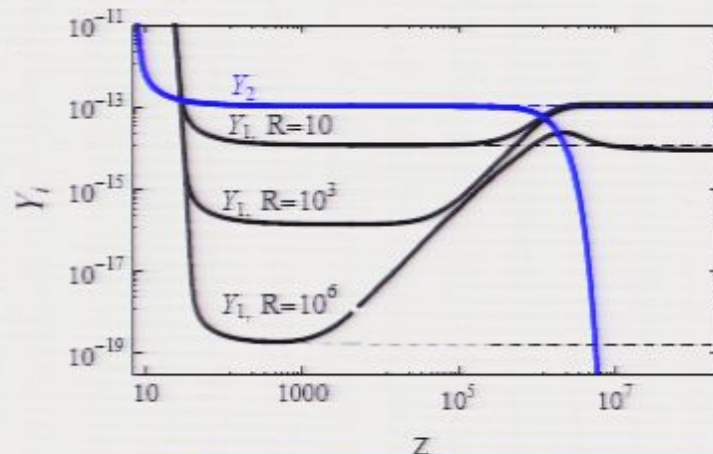
Large boost factors

- Sommerfeld enhancement due to a new long range attractive interaction between WIMPs
Hisano et al. 03-06; Cirelli et al. 07-08; Pospelov, Ratz 08; Arkani-Hamed et al. 08; Fox, Poppitz 08; Pospelov 08; Bai, Hambye 08; Nomura, Thaler 08; Ackerman et al 08;
- annihilation through resonances
- recombination through WIMP-onium Pospelov, Ratz 2008
- nonthermal DM: 2 component dark sector (2CDS) model Fairbairn, JZ, 2008
- decaying DM
Chen, Takahashi, Yanagida, 08; Ibarra, Tran 08; Hamaguchi, Nakamura, Shirai Yanagida 08; Yin, Yuan, Liu, hang, Bi, Zhu 08
- different temperature in the DM sector and in the visible sector Nelson and Spitzer

2-component dark sector

Fairbairn, JZ, 0810.4147

- an example of DM not a simple thermal relic
- two components: χ_2, χ_1
 - χ_2 is metastable, decays after freeze-out
 $\chi_2 \rightarrow \chi_1 + X_{SM}$
 - χ_1 is the DM that we observe now - and also gives Pamela/ATIC/FERMI/HESS signal
 - this setup decouples $\langle \sigma_{AV} \rangle_F$ from Ω_{DM}



$$B \simeq \frac{Y_1(\infty)}{Y_1^{\text{Th.rel.}}} \simeq N_{\text{dec}} R$$

$$R = \frac{m_1 \langle \sigma_{A1} v_1 \rangle}{m_2 \langle \sigma_{A2} v_2 \rangle}, \quad z = m_1/T, \quad Y_i(z) = n_i(z)/s(z)$$

2CDS II

- for large boost factors $\langle \sigma_{A1} v_1 \rangle \gg \langle \sigma_{A2} v_2 \rangle$
- to avoid wash-out χ_2 should decay after freeze-out

$$\Gamma_2 \ll 10^{-17} \text{GeV} \cdot \left(\frac{10^3}{B} \right) \cdot \left(\frac{m_1}{1 \text{TeV}} \right)^2$$

$$\text{or } \tau_2 \gg (10^{-7} \text{s}) \times (B/10^3) \cdot (1 \text{TeV}/m_1)^2$$

- from nucleosynthesis $\tau_2 < 1 \text{s}$
- if through dim 5 operators: $\Gamma_2 \simeq m_2^3 / (16\pi\Lambda^2)$, then

$$5 \cdot 10^{15} \text{GeV} \left(\frac{m_2}{1 \text{TeV}} \right)^{3/2} > \Lambda \gg 10^{12} \text{GeV} \cdot \left(\frac{B}{10^3} \right)^{1/2} \left(\frac{m_{1,2}}{1 \text{TeV}} \right)^{1/2}$$

- note: the allowed range includes see-saw scale
- if produced at LHC would fly macroscopic distance

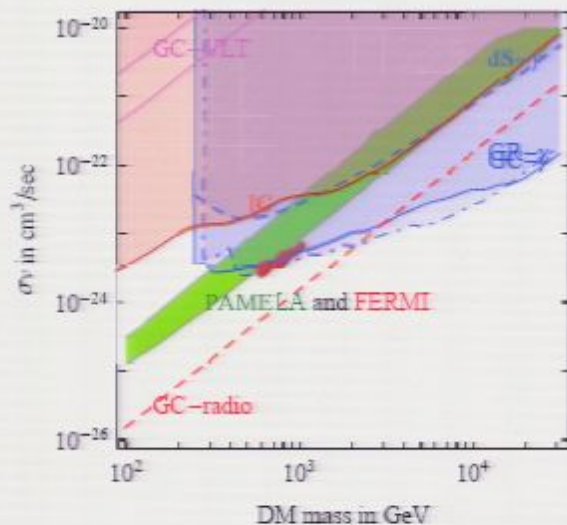
Gammas from galactic center

High energy photons

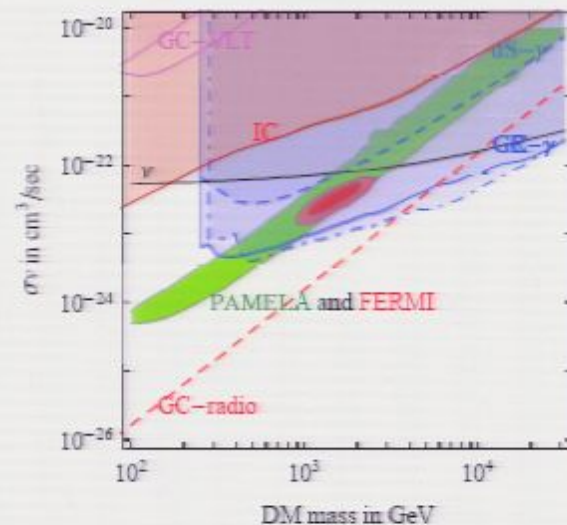
- annihilation into charged particles \Rightarrow inevitably high energy gamma from bremsstrahlung
- a curse and a blessing
 - a signal of DM (especially if a line)
 - right now there is a tension between cosmic ray anomaly and gamma from galactic center

Meade et al. 0905.0480

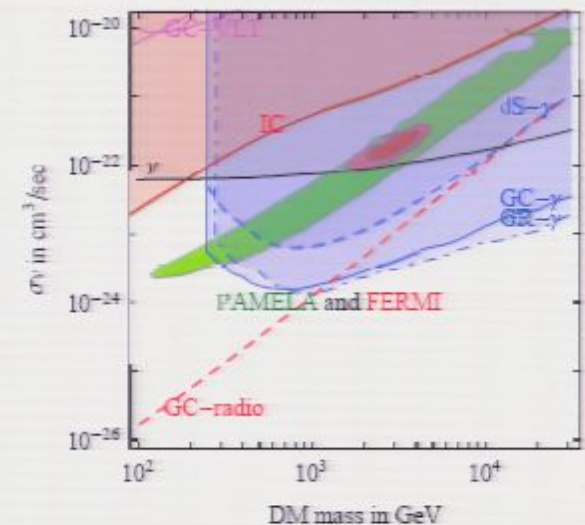
DM DM $\rightarrow e^+e^-$, NFW profile



DM DM $\rightarrow \mu^+\mu^-$, NFW profile



DM DM $\rightarrow \tau^+\tau^-$, NFW profile



Long lived intermediate states

ways out:

- less cuspy DM profiles in the galactic center?
- there is no problem if

$$\chi\chi \rightarrow \phi\phi \rightarrow 2SM 2\overline{SM}, \quad \text{with } \phi \text{ long lived}$$

Rothstein, Schwetz, JZ 0903.3116

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Rothstein, Schwetz, JZ 0903.3116

- if ϕ decays through dim 6 operators, $\tau \simeq 16\pi\Lambda^4/m_\phi^5$

$$\Lambda \sim \frac{2 \times 10^{12} \text{ GeV}}{(\beta\gamma)^{1/4}} \left(\frac{\lambda}{10 \text{ kpc}} \right)^{1/4} \left(\frac{m_\phi}{1 \text{ TeV}} \right)^{5/4}$$

- for a see-saw scale \Rightarrow propagation on galactic distances

Effective density

- for annihilating DM, the photon flux seen on earth

$$\frac{d\Phi_\gamma}{dE_\gamma} = \frac{\langle\sigma v\rangle}{4\pi} \frac{r_\odot \rho_\odot^2}{m_\chi^2} \frac{dN_\gamma}{dE_\gamma} J \Delta\Omega$$

- the dimensionless J -factor is

$$J = \frac{1}{\Delta\Omega} \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} \frac{ds}{r_\odot} \frac{\rho_{\text{DM}}^2(r)}{\rho_\odot^2}$$

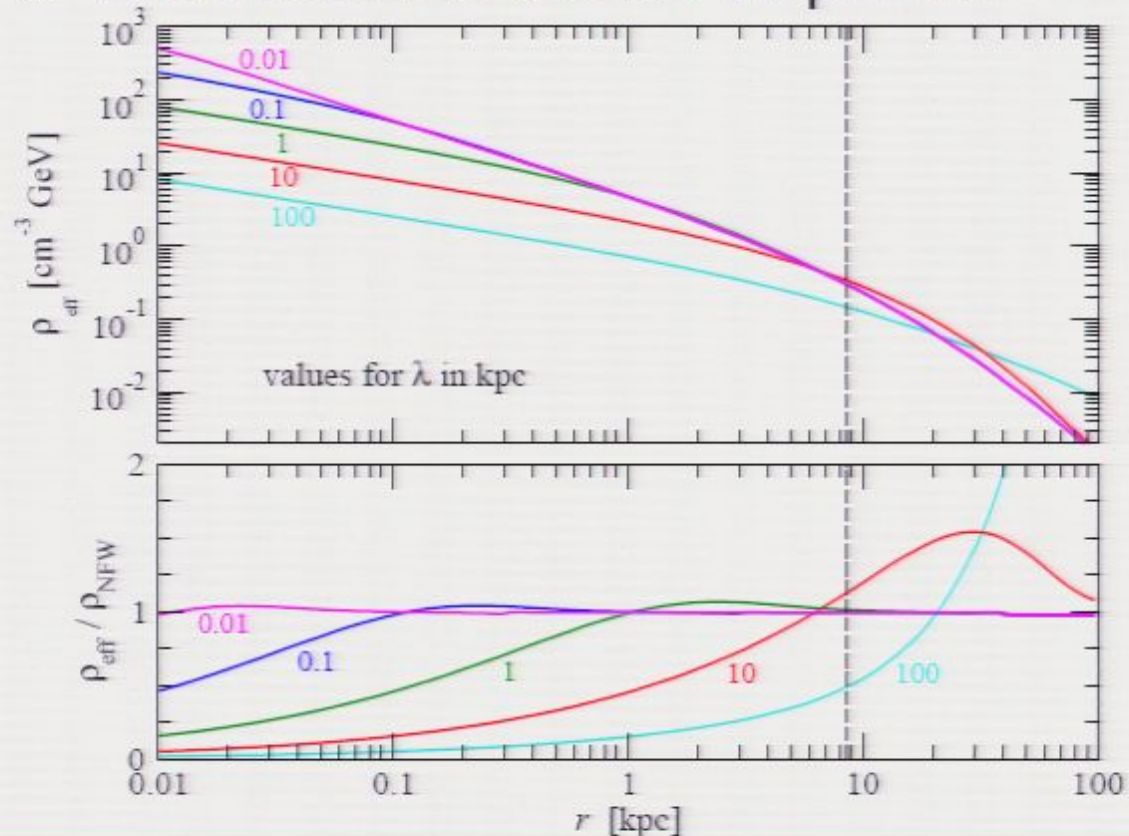
$$\rho_\odot = 0.3 \text{ GeV cm}^{-3} \text{ and } r_\odot = 8.5 \text{ kpc}$$

- for simplicity assume that ϕ nonrelat.
- then the effect of ϕ : $\rho \rightarrow \rho_{\text{eff}}$

$$\rho_{\text{eff}}^2(r) = \int d^3r' \rho^2(r') \frac{1}{4\pi\lambda} \frac{e^{-|\vec{r}-\vec{r}'|/\lambda}}{|\vec{r}-\vec{r}'|^2}$$

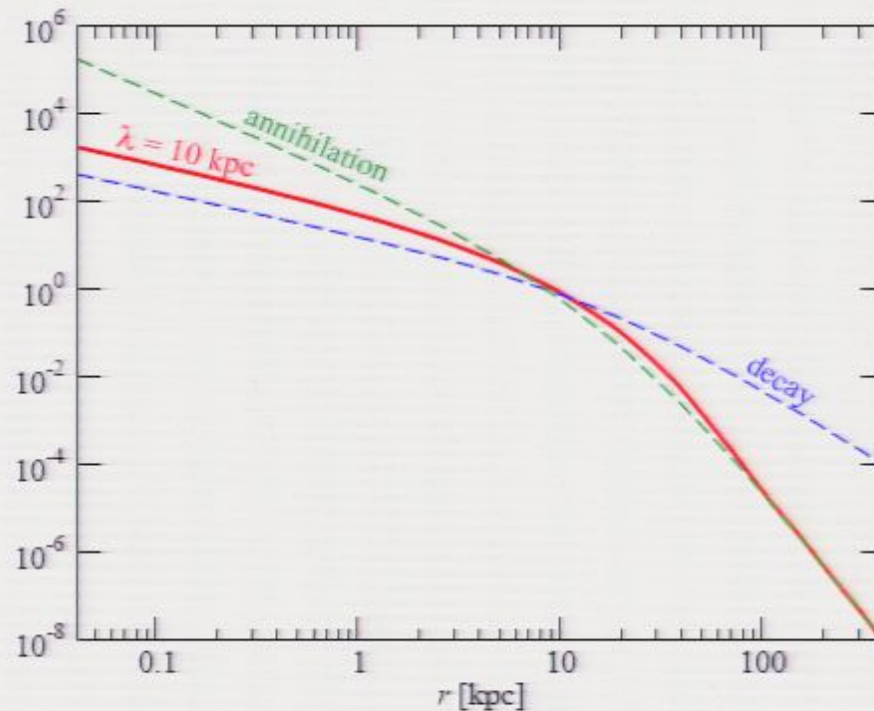
Effective density

- use NFW profile
 - for $r \lesssim \lambda$ suppressed $\rho_{\text{eff}}(r)$
 - for $r \gtrsim \lambda$ ϕ decays, so overproduction
 - for $r \gg \lambda$ we recover the NFW profile



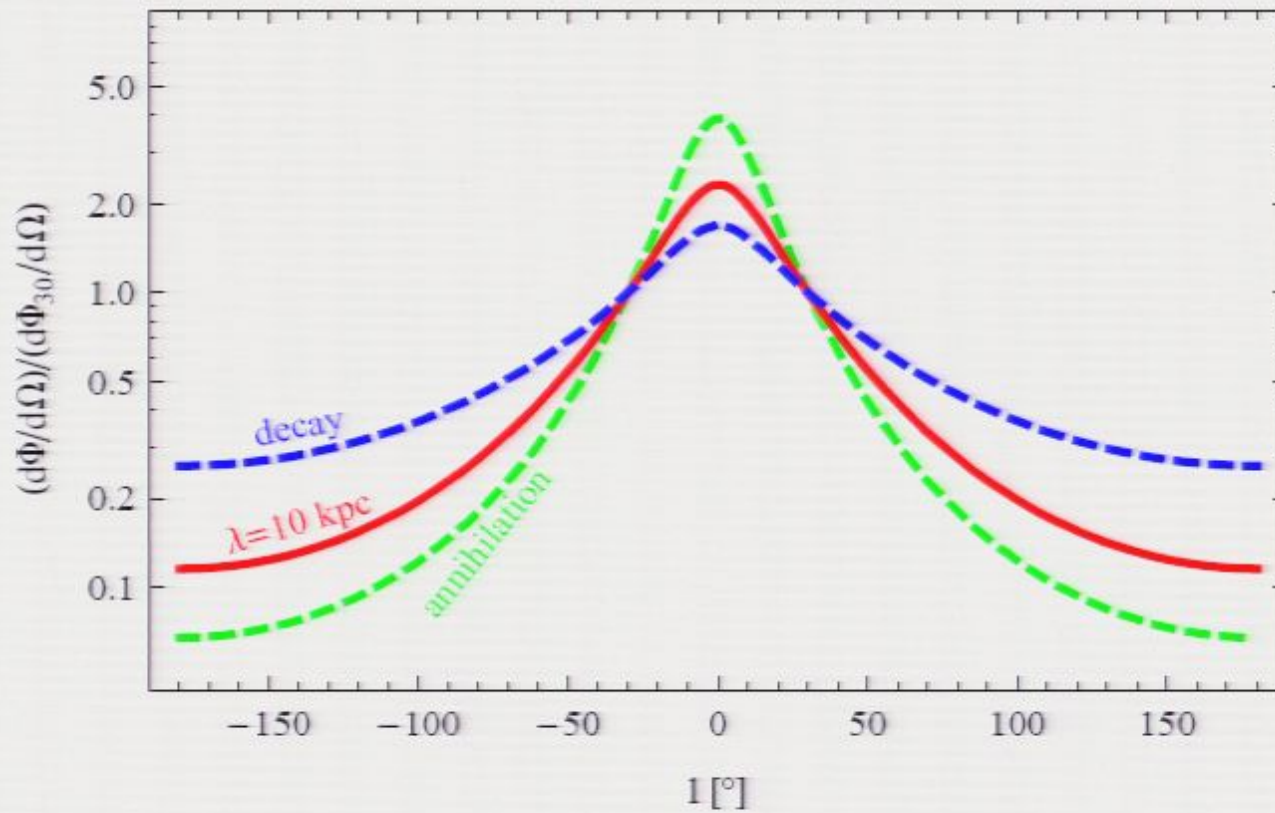
annih. vs decay

- the size of λ dialys the model effects from \sim annihilating DM to \sim decaying DM
- eg., for $r \ll \lambda$ and profiles $\rho(r) \propto r^{-\gamma} \Rightarrow \rho_{\text{eff}}^2(r) \propto 1/r^{2\gamma-1}$
 - for NFW ($\gamma = 1$): $\rho_{\text{eff}}^2(r) \propto 1/\rho(r)$



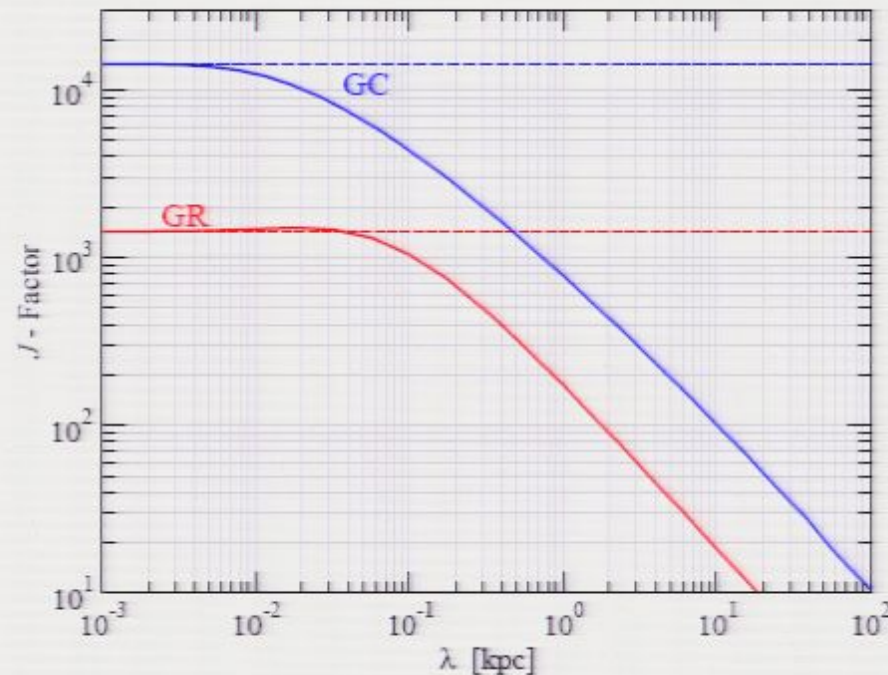
annih. vs decay

- comparing angular dependence



High energy photons

- the ϕ 's decay outside the galactic center
- photon flux from galactic center suppressed



GC=Galactic center, GR=Galactic Ridge

effect on e^+e^- fits

- fix decay $\phi \rightarrow \mu^+\mu^-$, NFW profile, MED propagation model for el.
- injections spectra for photons, e^+ , e^- from pythia
- background from Galprop + free scale, slope allowed to vary by ± 0.05
- split Fermi and ATIC data

A technical detail

- solution of diff. eq., source $q(\vec{x}, E) = \kappa \left(\frac{\rho(\vec{x})}{\rho_0} \right)^2 f(\epsilon)$

$$\Psi(r, z, \epsilon) = \kappa \frac{\tau_E}{\epsilon^2} \int_{\epsilon}^{\infty} d\epsilon_S f(\epsilon_S) \tilde{I}(\lambda_D, r, z),$$

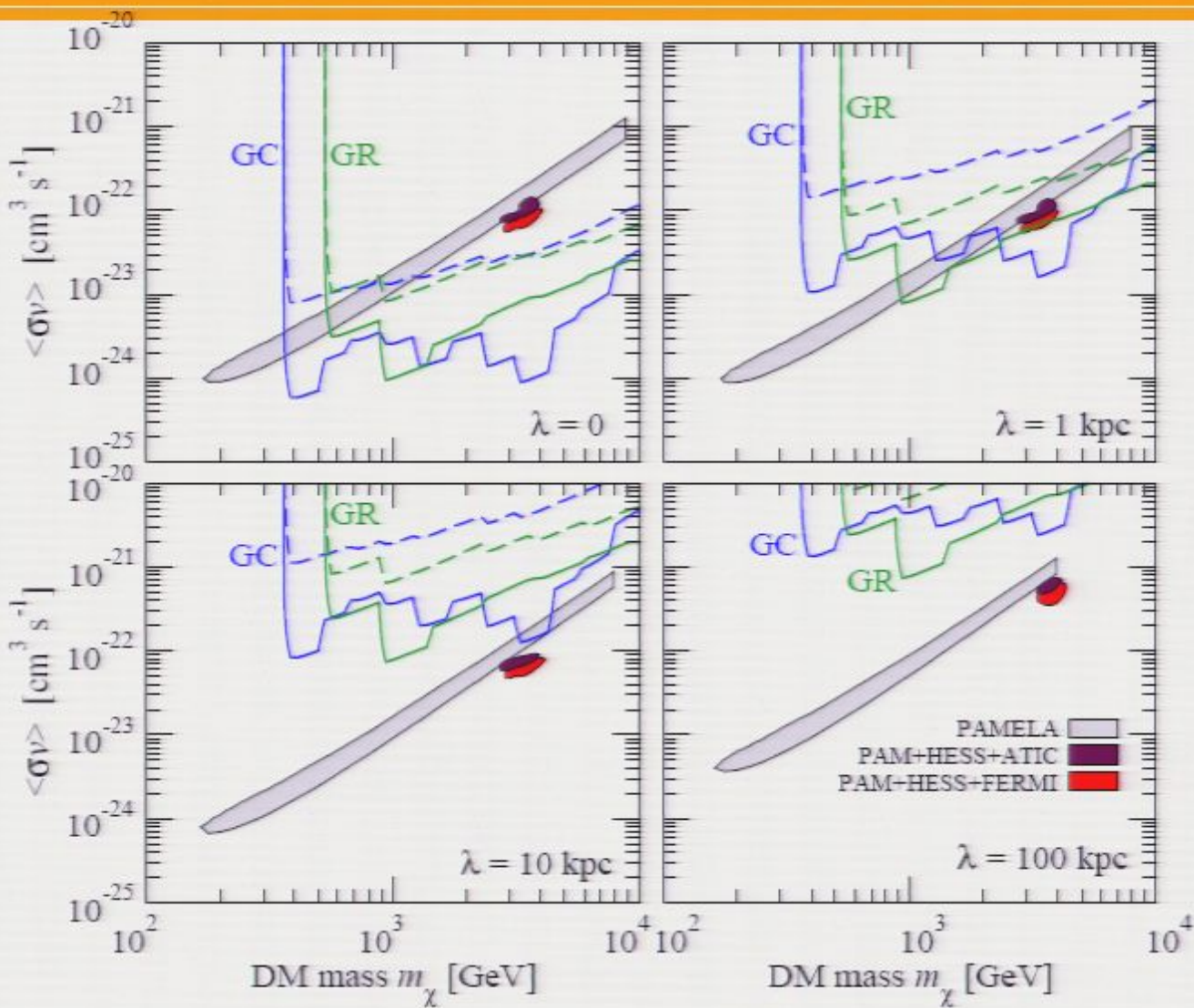
- astrophysics in $I(\lambda) \Rightarrow$ we need to vary it
 - if expanded in Bessel funcs., very slow converg.
 - can write a diff. eq., solved num. in ~ 10 s on pc

$$\nabla^2 \tilde{I}(\lambda_D, r, z) - \frac{2}{\lambda_D} \partial_{\lambda_D} \tilde{I}(\lambda_D, r, z) = 0,$$

boundary conditions:

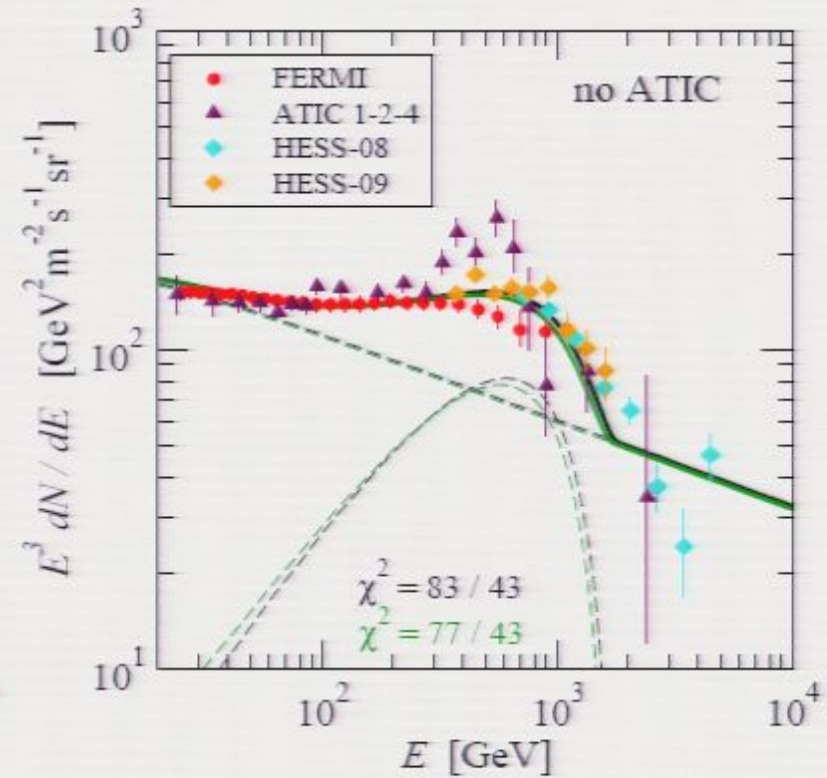
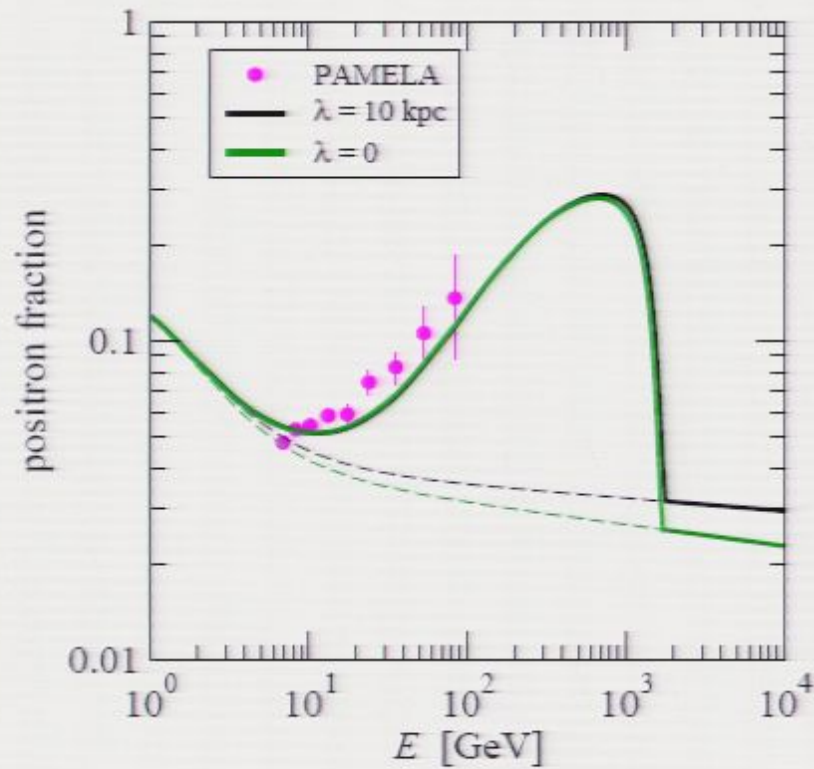
$$\begin{aligned} \tilde{I}(0, r, z) &= \left(\frac{\rho(r, z)}{\rho_0} \right)^2, & \tilde{I}(\infty, r, z) &= 0, \\ \tilde{I}(\lambda_D, R_{\text{gal}}, z) &= 0, & \tilde{I}(\lambda_D, r, \pm L) &= 0. \end{aligned}$$

e^+e^- fits



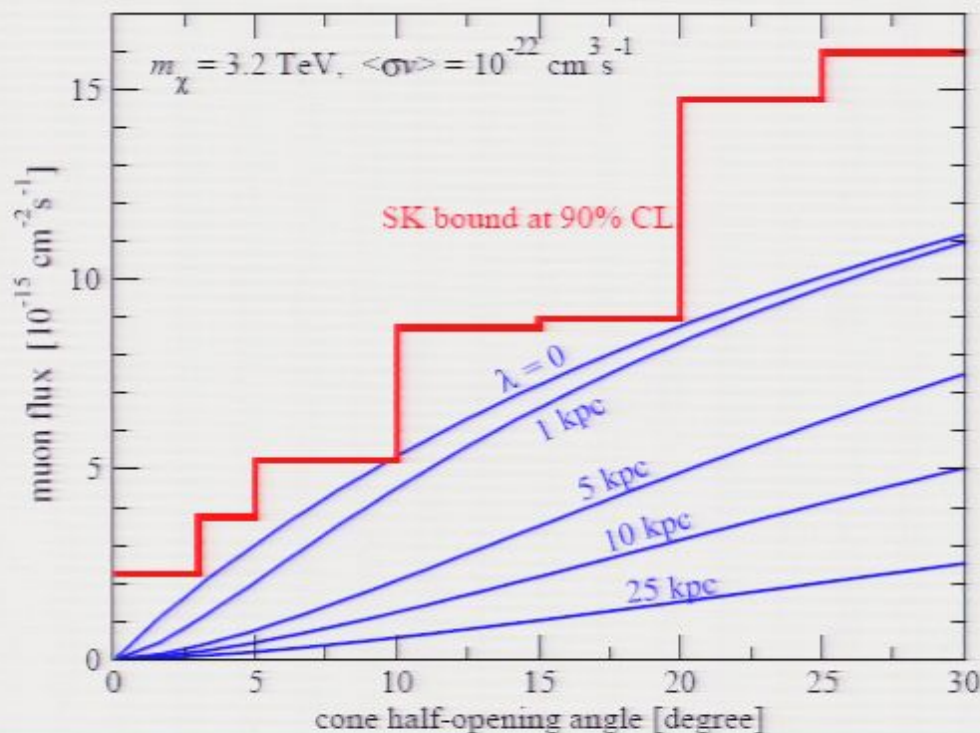
minimum $m_\chi \simeq 3.6\text{TeV}$, $\langle\sigma v\rangle \sim 10^{-22}\text{cm}^3\text{s}^{-1}$

comparison with data



Neutrino bounds

- Super-Kamiokande: upper limit on ν induced upward going muon flux from various extra-terrestrial sources
- potentially relevant for LLP scenario since 8 neutrinos/DM annihilation ($\chi\chi \rightarrow \phi\phi \rightarrow 2\mu^+2\mu^-$)



Can χ, ϕ be thermal relics?

- other constraints if χ a thermal relic and ϕ a meta-stable thermal relic
- LLP scenario interesting if $\lambda \gtrsim 10$ kpc, then

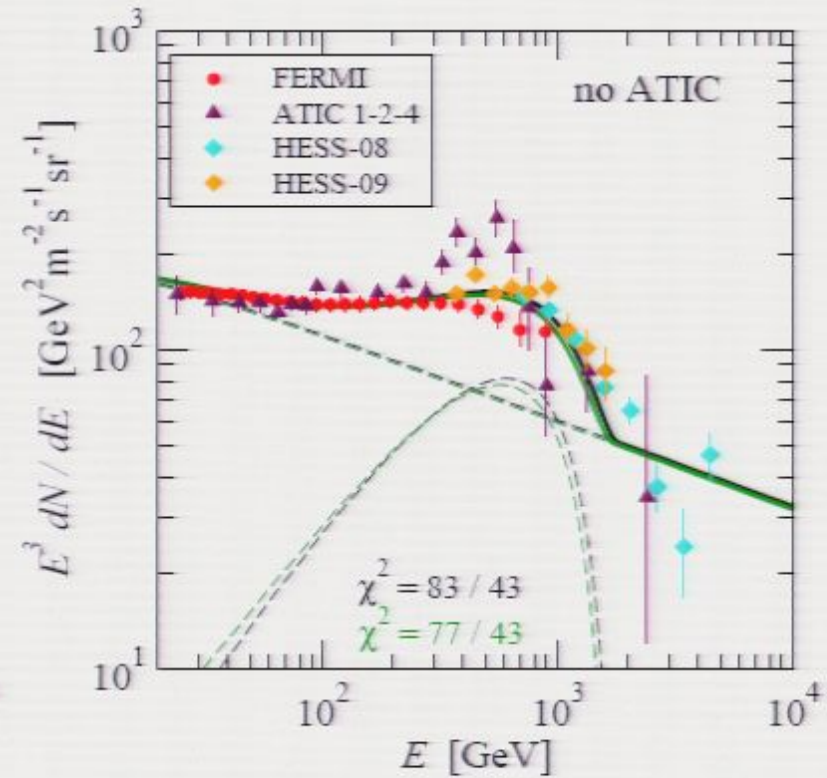
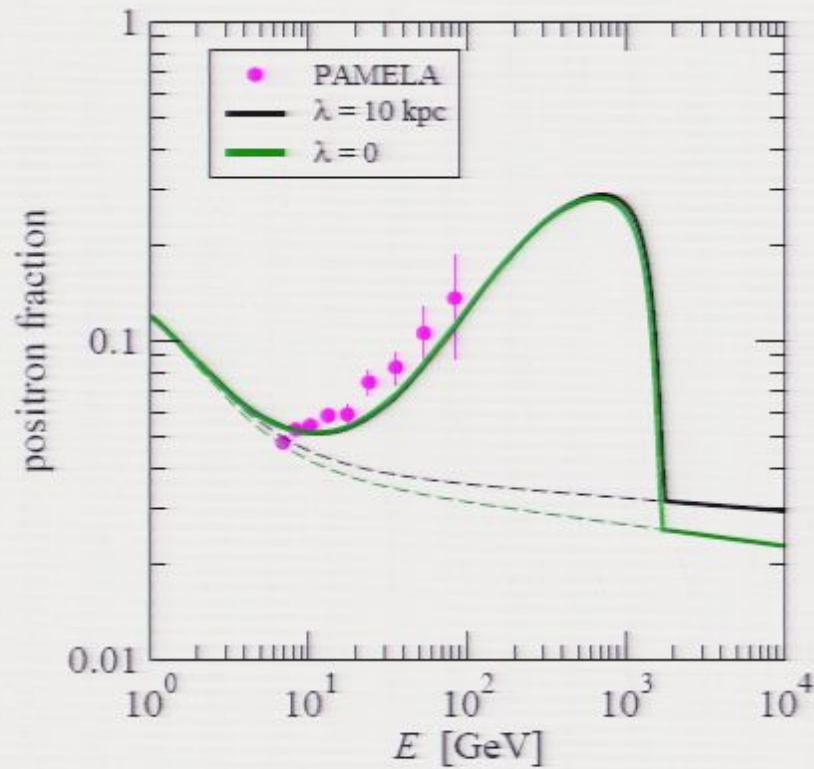
$$\tau_\phi = \frac{\lambda}{c\beta\gamma} \simeq \frac{10^{12} \text{ s}}{\beta\gamma} \left(\frac{\lambda}{10 \text{ kpc}} \right).$$

- such late decaying relics modify light element abundances $\Rightarrow \rho_\phi \lesssim 10^{-3} \rho_\chi$
- for $\tau \gtrsim 10^{13}$ s constr. from diffuse gamma ray bckg.
- ratio $\sigma_{\chi}^{\text{ann}} / \sigma_{\phi}^{\text{ann}}$ must be of order 10^{-3} or smaller
- an open exercise in model building

Conclusions

- have provided two scenarios with metastable particles
- they solve the problems for annihilating DM explanation of cosmic ray anomaly
 - the large boost factors
 - the lack of DM photon signal from galactic center

comparison with data



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