

Title: Energy transfer in warped string compactifications and warped Kaluza-Klein modes as a dark matter candidate

Date: Jun 16, 2009 02:00 PM

URL: <http://pirsa.org/09060005>

Abstract: Strongly warped regions, or throats, are a common feature of string theory compactifications. In the early, hot universe, energy will be transferred between these throats and between throats and the standard model. Using the gauge-gravity duality, we calculate the rate of this energy transfer. Due to the warping, the resulting decay rate of throat-localized Kaluza-Klein states to other throats or the standard model is strongly suppressed. If their lifetime is longer than the current age of the universe, these states are an interesting dark matter candidate. We discuss a scenario along these lines.

ENERGY TRANSFER IN WARPED STRING COMPACTIFICATIONS AND WARPED KALUZA-KLEIN MODES AS A DARK MATTER CANDIDATE

Benedict v. Harling

School of Physics
University of Melbourne
Australia

June 16, 2009

Based on work
with A. Hebecker,
T. Noguchi
and S. Halverson

arXiv:0705.3641
arXiv:0801.4013
arXiv:0902.1188

ENERGY TRANSFER IN WARPED STRING COMPACTIFICATIONS AND WARPED KALUZA-KLEIN MODES AS A DARK MATTER CANDIDATE

Benedict v. Harling

School of Physics
University of Melbourne
Australia

June 16, 2009

Based on work
with A. Hebecker,
T. Noguchi
and S. Halter:

arXiv:0705.3648

arXiv:0801.4015

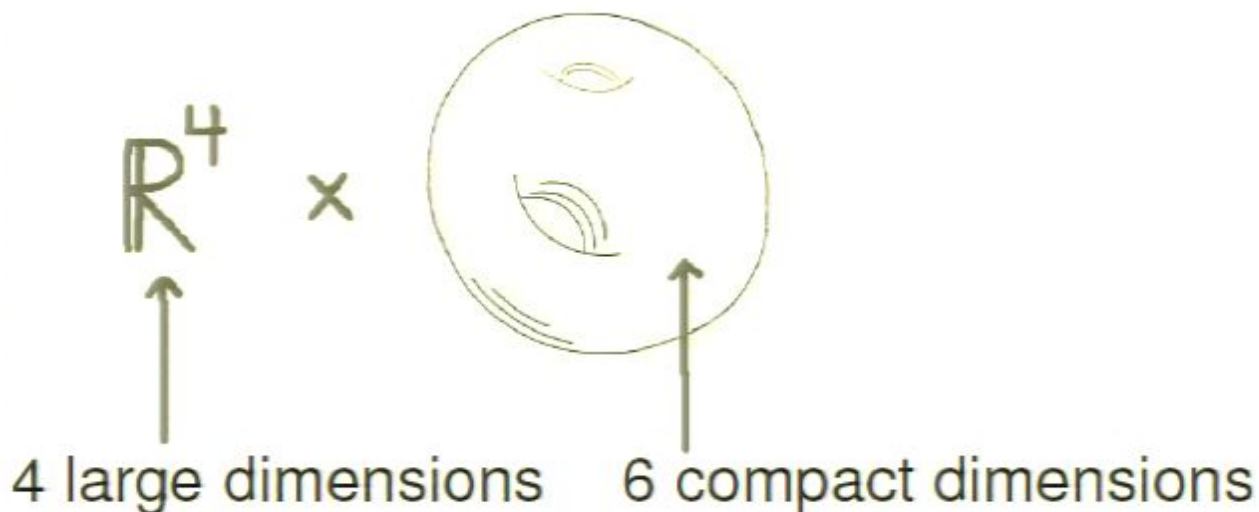
arXiv:0907.1111

INTRODUCTION: THROATS FROM FLUX

- We consider **type IIB string theory**. Lives in 10 spacetime dimensions.
 - The 6 extra dimensions are thought to be compactified.
 - Type IIB string theory has form fields (generalizations of U(1) gauge field) which can carry energy density or flux.
 - Backreaction of flux on geometry can create strongly warped regions, so-called throats.

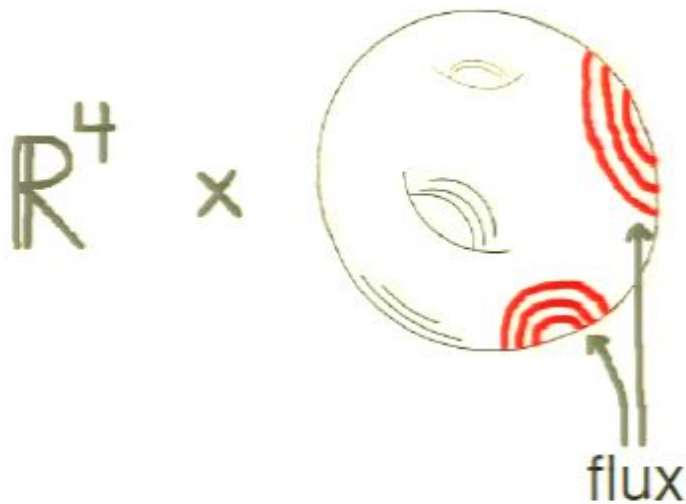
INTRODUCTION: THROATS FROM FLUX

- We consider **type IIB string theory**. Lives in 10 spacetime dimensions.
- The 6 extra dimensions are thought to be **compactified**.
- Type IIB string theory has form fields (generalizations of U(1) gauge field) which can carry energy density or flux.
- Backreaction of flux on geometry can create strongly warped regions, so-called throats.



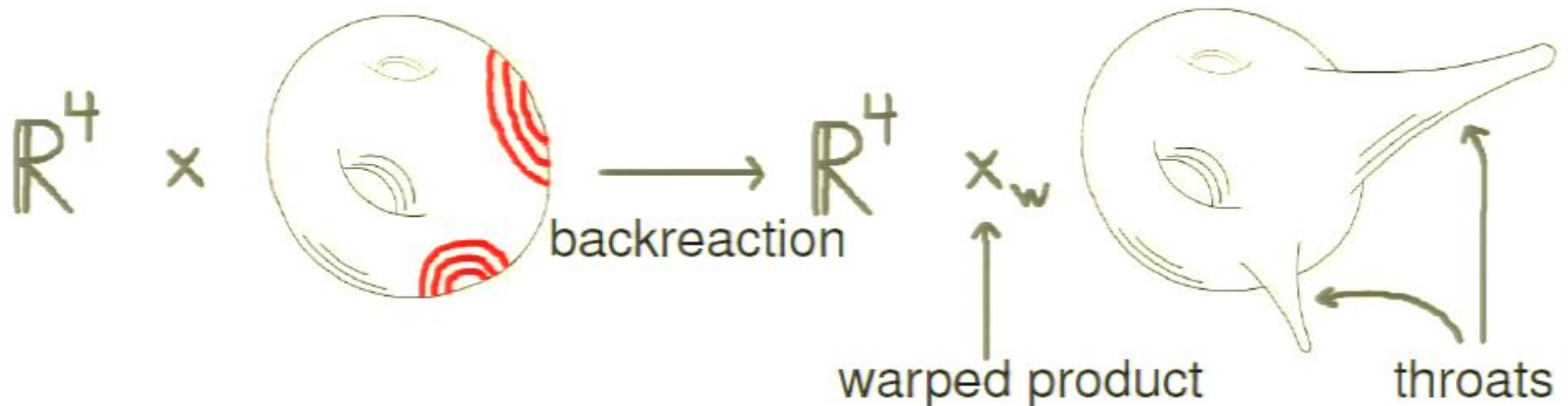
INTRODUCTION: THROATS FROM FLUX

- We consider **type IIB string theory**. Lives in 10 spacetime dimensions.
- The 6 extra dimensions are thought to be **compactified**.
- Type IIB string theory has **form fields** (generalizations of U(1) gauge field) which can carry energy density or **flux**.
- Backreaction of flux on geometry can create strongly warped regions, so-called throats.

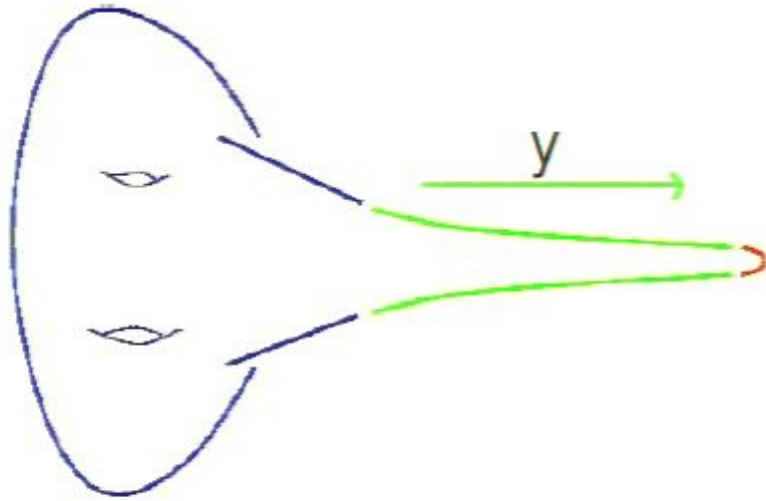


INTRODUCTION: THROATS FROM FLUX

- We consider **type IIB string theory**. Lives in 10 spacetime dimensions.
- The 6 extra dimensions are thought to be **compactified**.
- Type IIB string theory has **form fields** (generalizations of U(1) gauge field) which can carry energy density or **flux**.
- Backreaction of flux on geometry can create strongly warped regions, so-called **throats**.



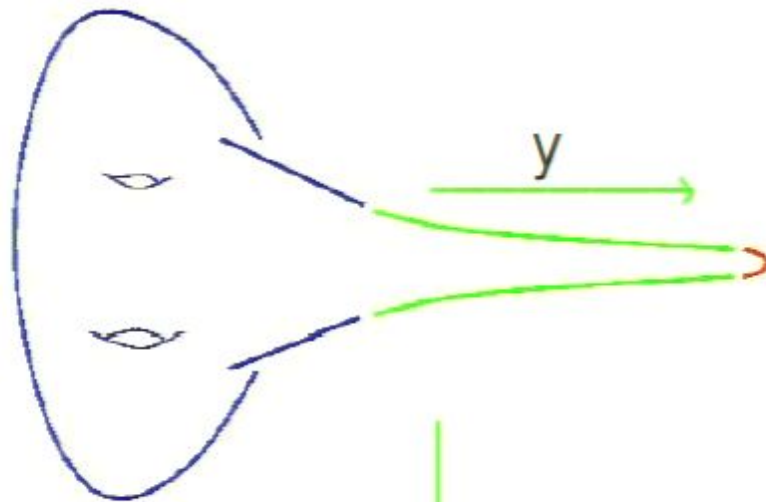
INTRODUCTION: THROATS AS RS MODELS



Metric in the throat region:

$$ds^2 = e^{-2\pi y/R} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 + .$$

INTRODUCTION: THROATS AS RS MODELS



Metric in the throat region:

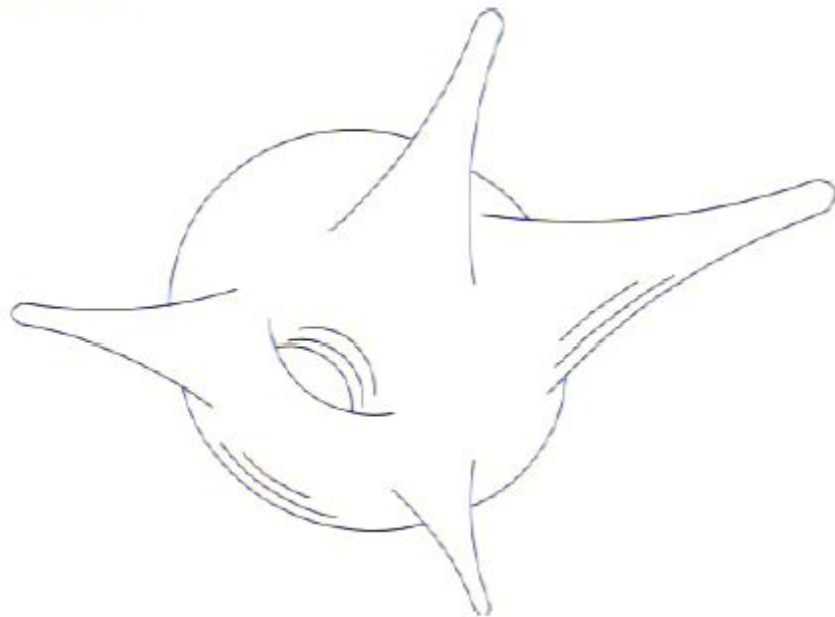
$$ds^2 = e^{-2\pi y/R} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 + .$$



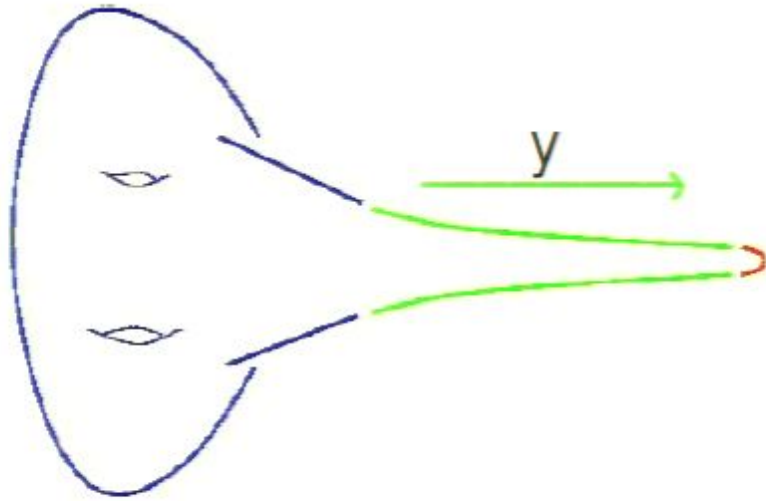
Throats are string realizations of the **Randall-Sundrum model**.

MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on D-branes (hyperplanes in 10d space on which open strings end) e.g. in a throat.
- SM is not in the early universe. \Rightarrow Throats will be heated up.
- The so-produced KK particles in throats decay back to the SM and other throats.



INTRODUCTION: THROATS AS RS MODELS

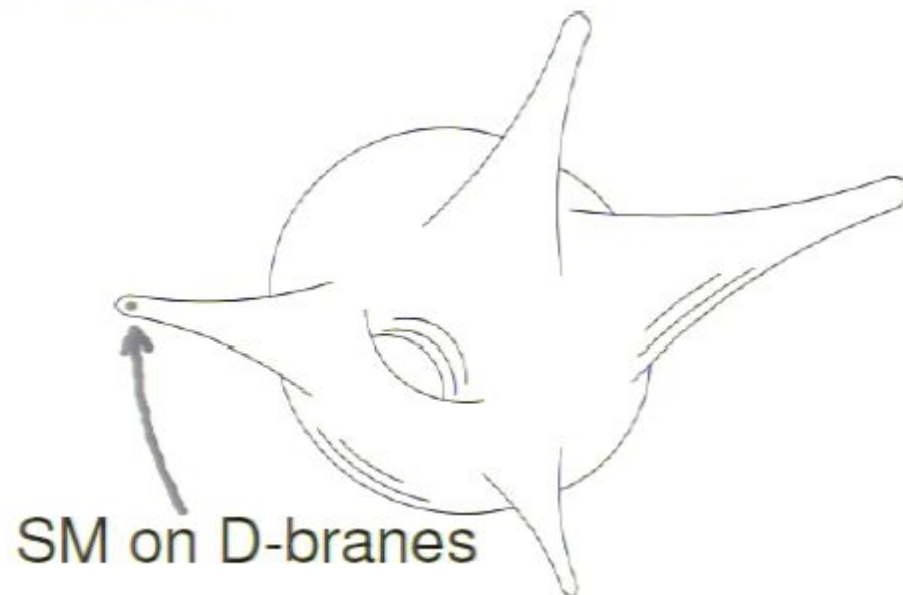


Metric in the throat region:

$$ds^2 = e^{-2\pi y/R} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 + .$$

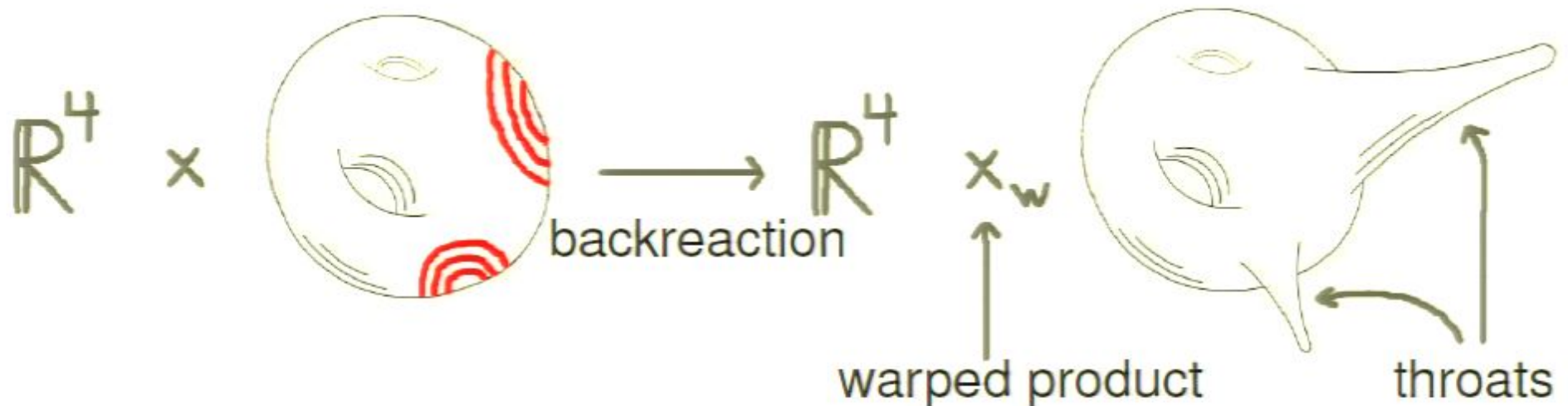
MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on **D-branes** (hyperplanes in 10d space on which open strings end) e.g. in a throat.
 - SM is not in the early universe. \Rightarrow Throats will be heated up.
 - The so-produced KK particles in throats decay back to the SM and other throats.



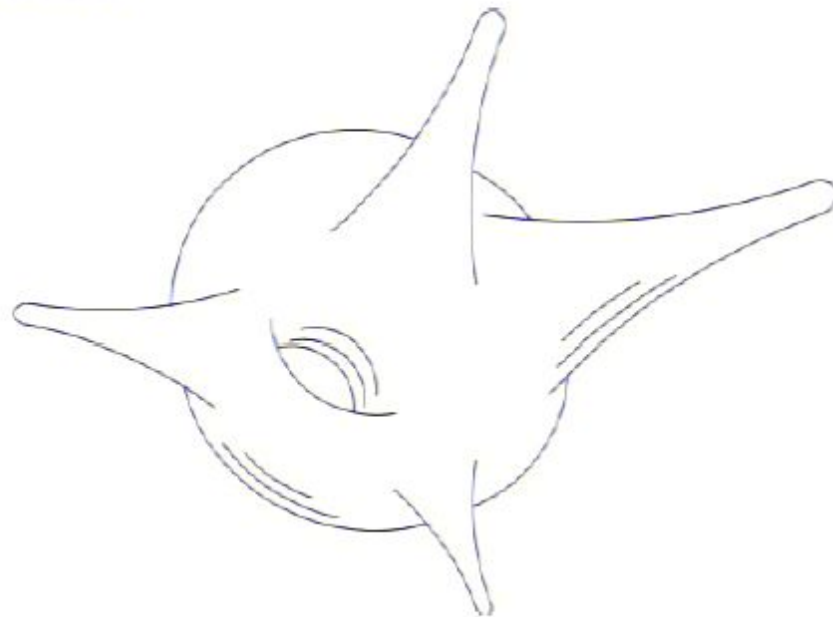
INTRODUCTION: THROATS FROM FLUX

- We consider **type IIB string theory**. Lives in 10 spacetime dimensions.
- The 6 extra dimensions are thought to be **compactified**.
- Type IIB string theory has **form fields** (generalizations of U(1) gauge field) which can carry energy density or **flux**.
- Backreaction of flux on geometry can create strongly warped regions, so-called **throats**.

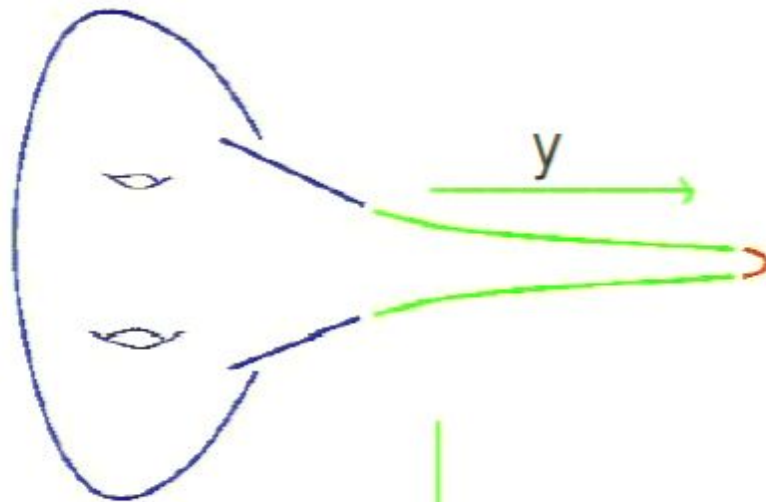


MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on D-branes (hyperplanes in 10d space on which open strings end) e.g. in a throat.
- SM is not in the early universe. \Rightarrow Throats will be heated up.
- The so-produced KK particles in throats decay back to the SM and other throats.



INTRODUCTION: THROATS AS RS MODELS



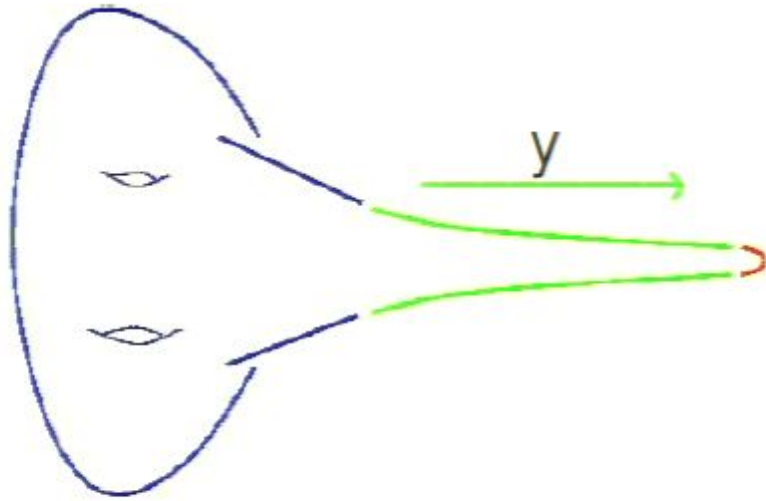
Metric in the throat region:

$$ds^2 = e^{-2\pi y/R} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 + .$$



Throats are string realizations of the **Randall-Sundrum model**.

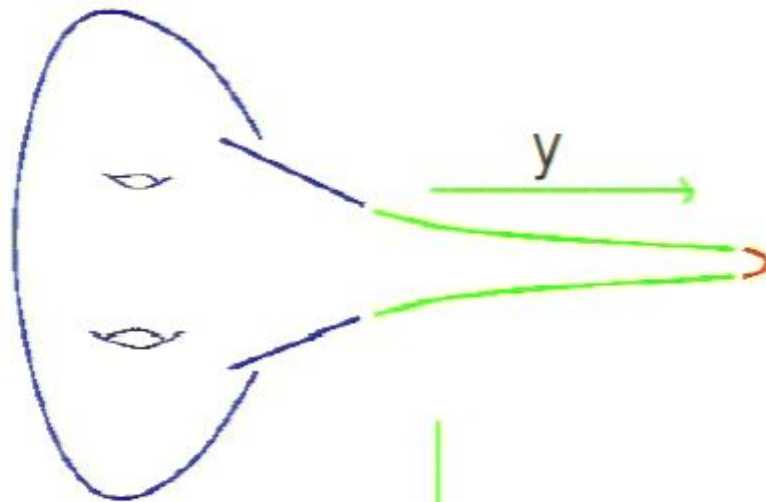
INTRODUCTION: THROATS AS RS MODELS



Metric in the throat region:

$$ds^2 = e^{-2\pi y/R} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 + .$$

INTRODUCTION: THROATS AS RS MODELS



Metric in the throat region:

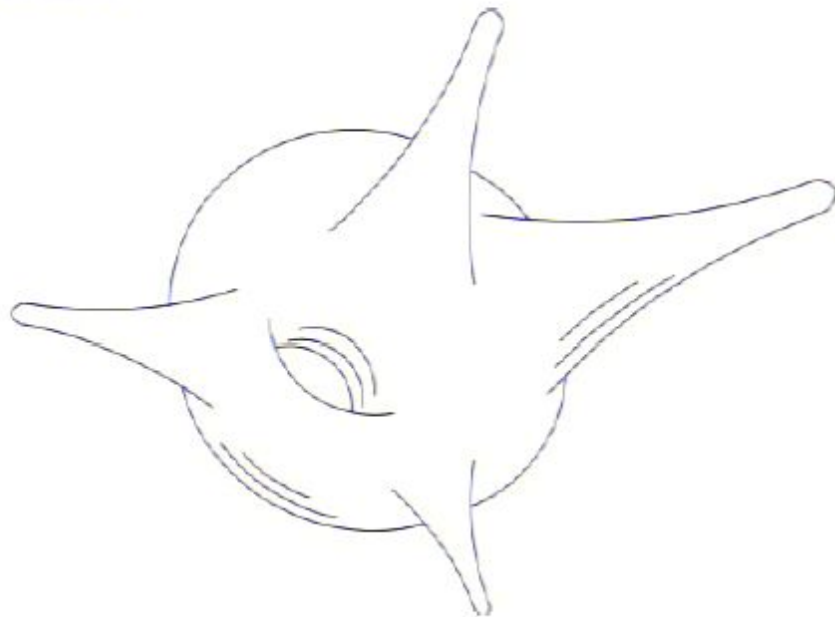
$$ds^2 = e^{-2\pi y/R} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2 + .$$



Throats are string realizations of the **Randall-Sundrum model**.

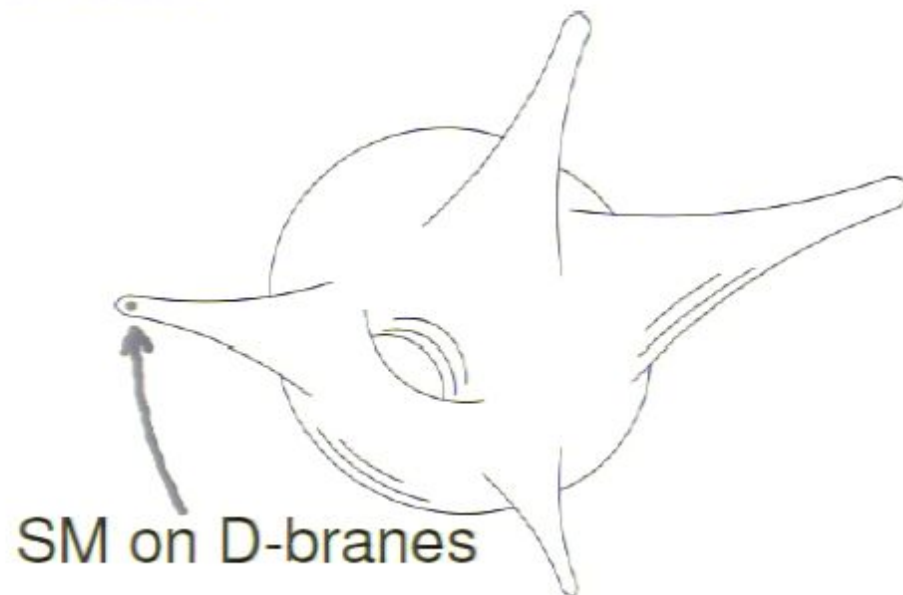
MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on D-branes (hyperplanes in 10d space on which open strings end) e.g. in a throat.
- SM is not in the early universe. \Rightarrow Throats will be heated up.
- The so-produced KK particles in throats decay back to the SM and other throats.



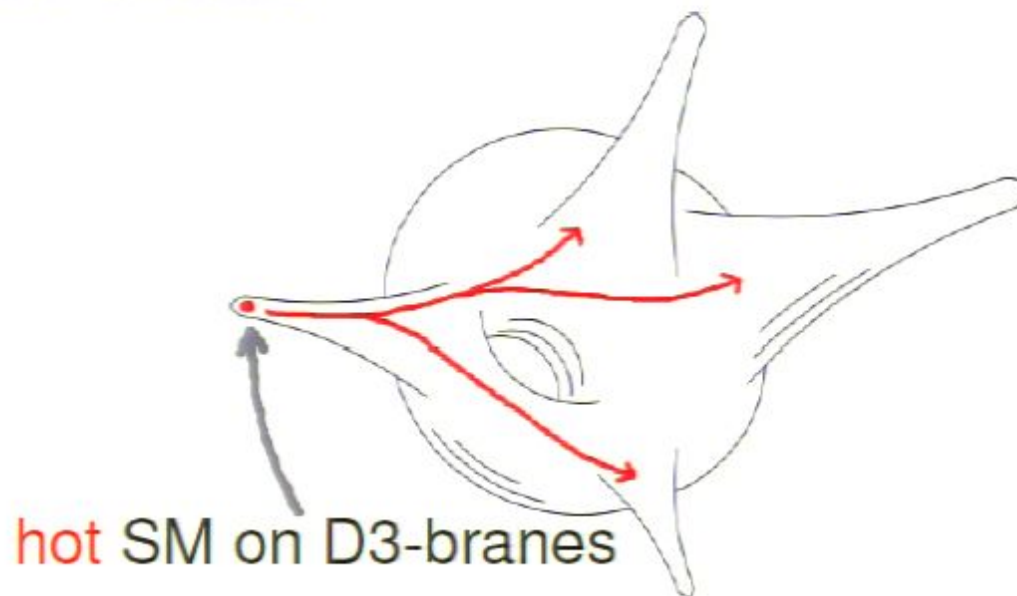
MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on **D-branes** (hyperplanes in 10d space on which open strings end) e.g. in a throat.
 - SM is not in the early universe. \Rightarrow Throats will be heated up.
 - The so-produced KK particles in throats decay back to the SM and other throats.



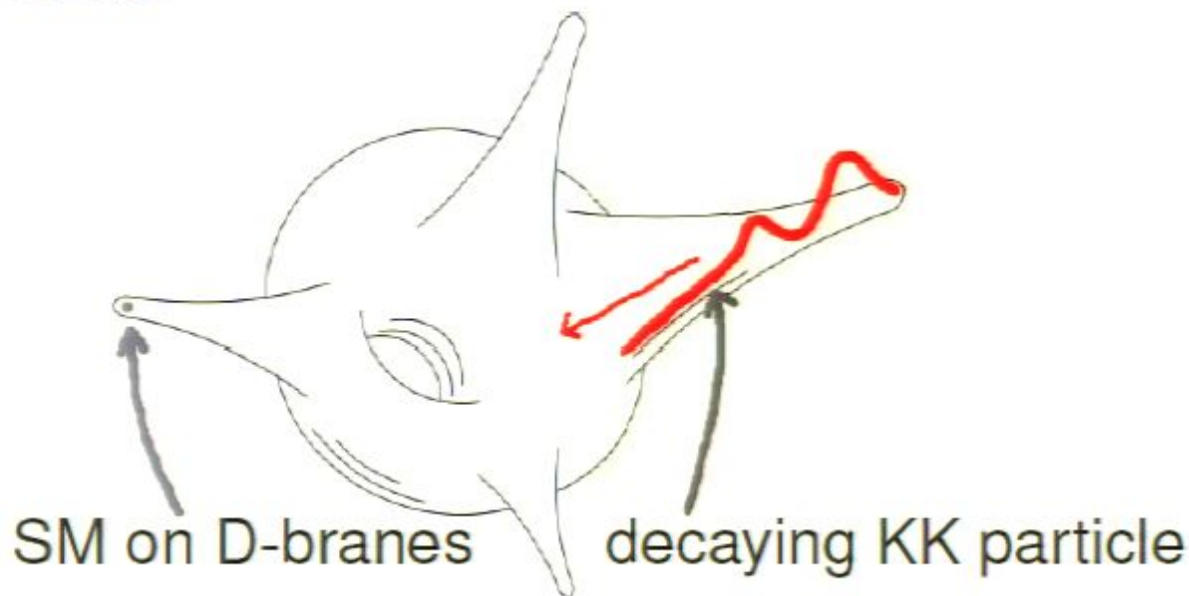
MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on **D-branes** (hyperplanes in 10d space on which open strings end) e.g. in a throat.
- SM is hot in the early universe. \Rightarrow Throats will be **heated up**.
- The so-produced KK particles in throats decay back to the SM and other throats.



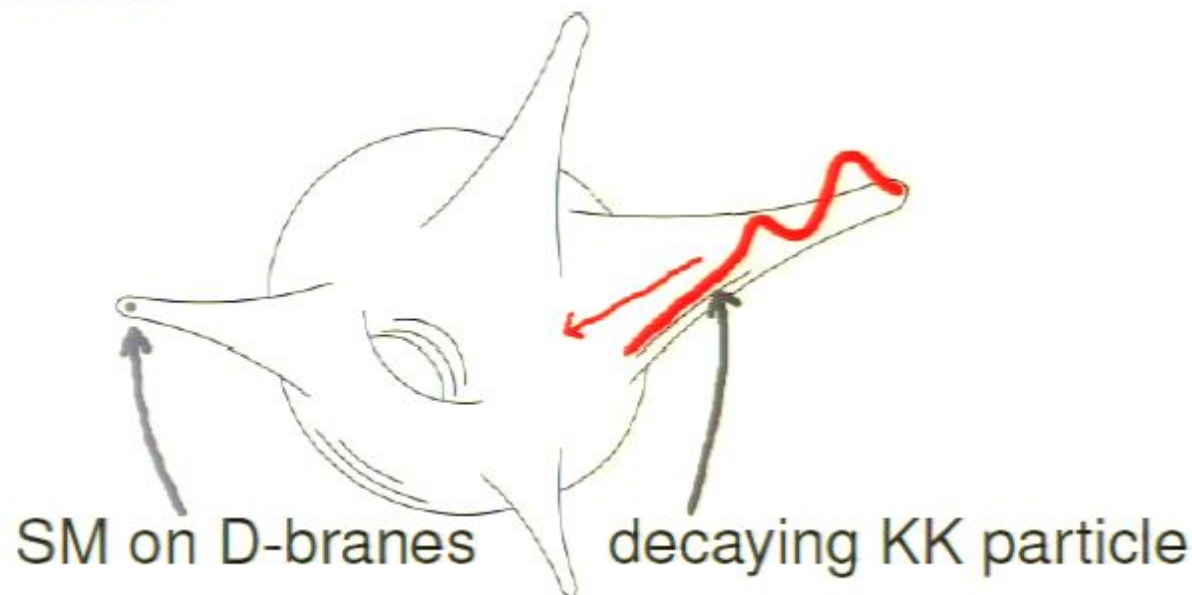
MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on **D-branes** (hyperplanes in 10d space on which open strings end) e.g. in a throat.
- SM is hot in the early universe. \Rightarrow Throats will be **heated up**.
- The so-produced KK particles in throats **decay back** to the SM and other throats.



MOTIVATION: HEATED THROATS

- Throats are believed to be a **common feature** in the landscape of string theory vacua [Hebecker, March-Russell].
- Standard model can be realized on **D-branes** (hyperplanes in 10d space on which open strings end) e.g. in a throat.
- SM is hot in the early universe. \Rightarrow Throats will be **heated up**.
- The so-produced KK particles in throats **decay back** to the SM and other throats.



OUTLINE OF THE TALK

1 CALCULATION OF THE HEAT TRANSFER RATE



OUTLINE OF THE TALK

1 CALCULATION OF THE HEAT TRANSFER RATE

- 2.1.1 Heat transfer through a wall
- 2.1.2 Heat transfer through a window
- 2.1.3 Heat transfer through a door

OUTLINE OF THE TALK

1 CALCULATION OF THE HEAT TRANSFER RATE



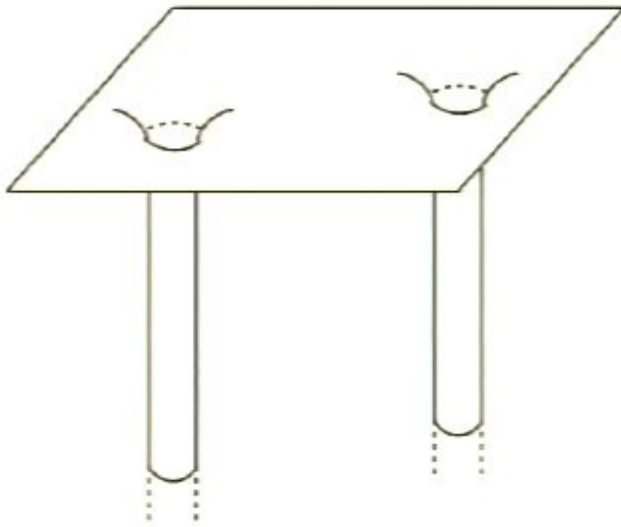
2 DARK MATTER IN A THROAT: THERMAL PRODUCTION



OUTLINE OF THE TALK

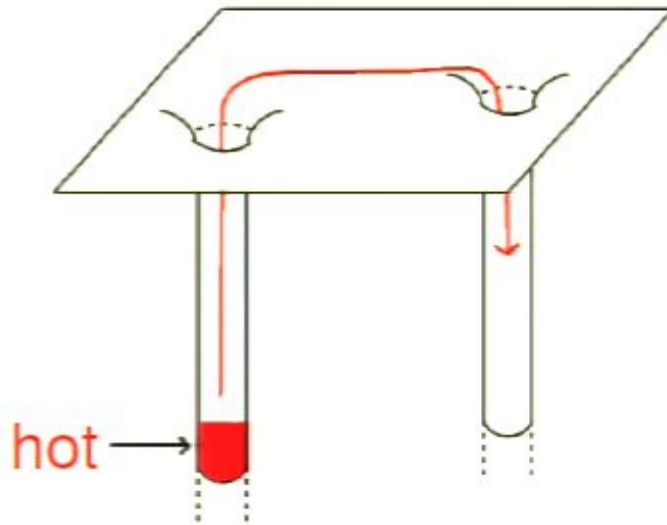
- 1 CALCULATION OF THE HEAT TRANSFER RATE
- 2 DARK MATTER IN A THROAT: THERMAL PRODUCTION
- 3 DARK MATTER IN A THROAT: DECAY

HEAT TRANSFER RATE



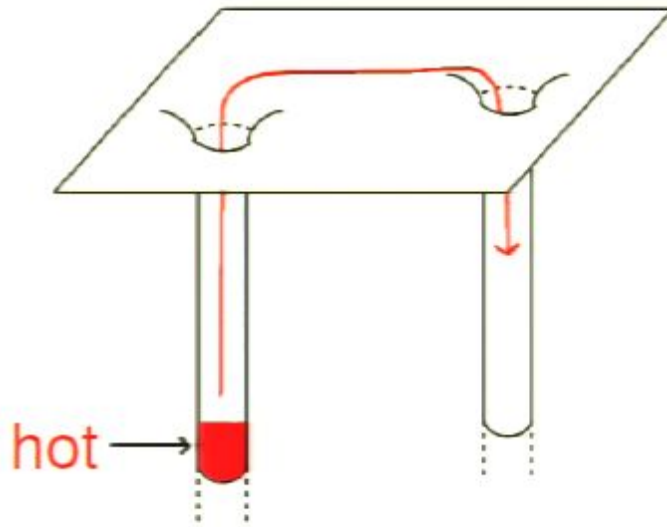
- We consider two $\text{AdS}_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by transition probability of particles between the two throats.
- Problem turns out to be a multi-dimensional tunneling problem. \Rightarrow Difficult to solve.
- Way out via AdS/CFT duality: String theory on $\text{AdS}_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

HEAT TRANSFER RATE



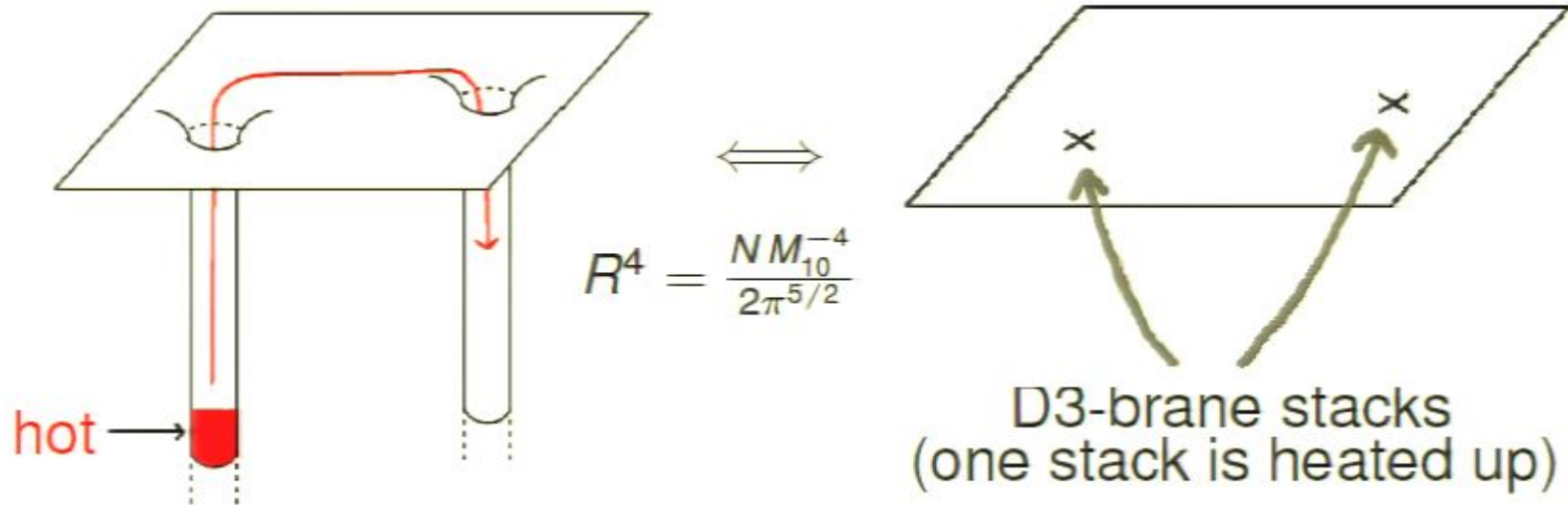
- We consider two $AdS_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by transition probability of particles between the two throats.
- Problem turns out to be a multi-dimensional tunneling problem. \Rightarrow Difficult to solve.
- Way out via AdS/CFT duality: String theory on $AdS_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

HEAT TRANSFER RATE



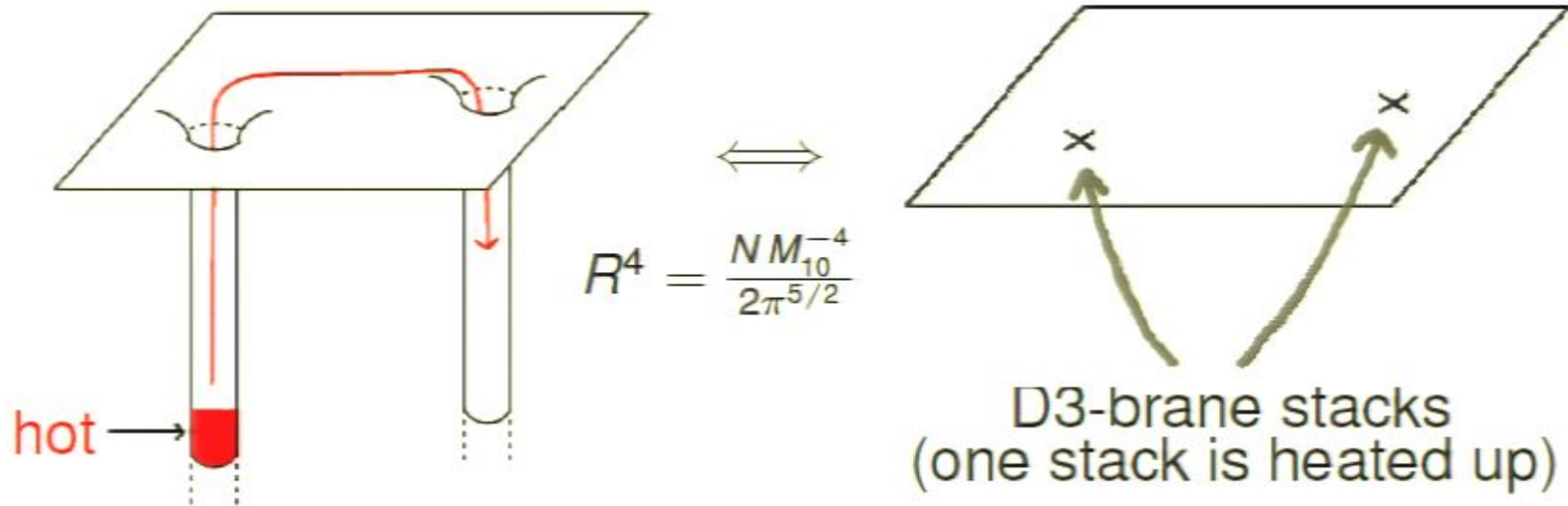
- We consider two $\text{AdS}_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by **transition probability** of particles between the two throats.
- Problem turns out to be a **multi-dimensional tunneling problem**. \Rightarrow Difficult to solve.
- Way out via AdS/CFT duality: String theory on $\text{AdS}_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

HEAT TRANSFER RATE



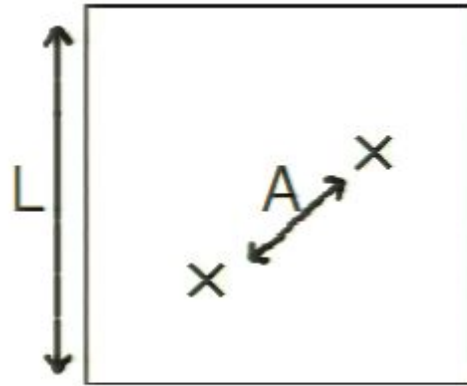
- We consider two $\text{AdS}_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by **transition probability** of particles between the two throats.
- Problem turns out to be a **multi-dimensional tunneling problem**. \Rightarrow Difficult to solve.
- Way out via **AdS/CFT duality**: String theory on $\text{AdS}_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

HEAT TRANSFER RATE



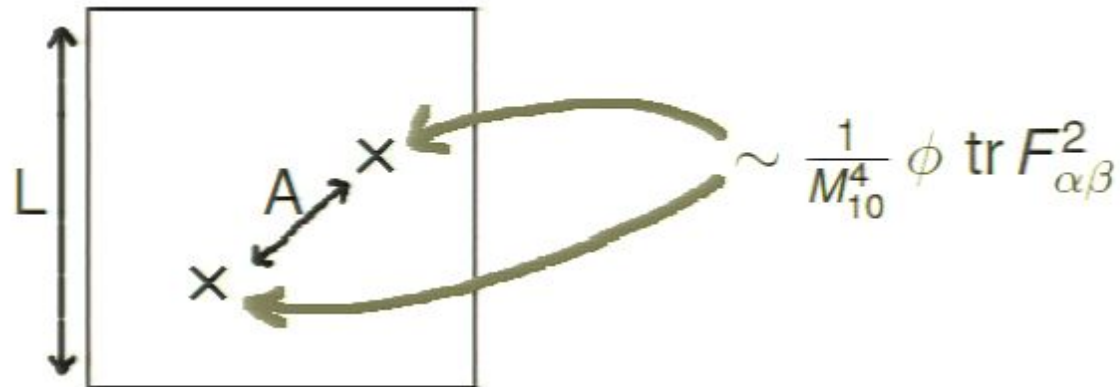
- We consider two $\text{AdS}_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by **transition probability** of particles between the two throats.
- Problem turns out to be a **multi-dimensional tunneling problem**. \Rightarrow Difficult to solve.
- Way out via **AdS/CFT duality**: String theory on $\text{AdS}_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

HEAT TRANSFER RATE



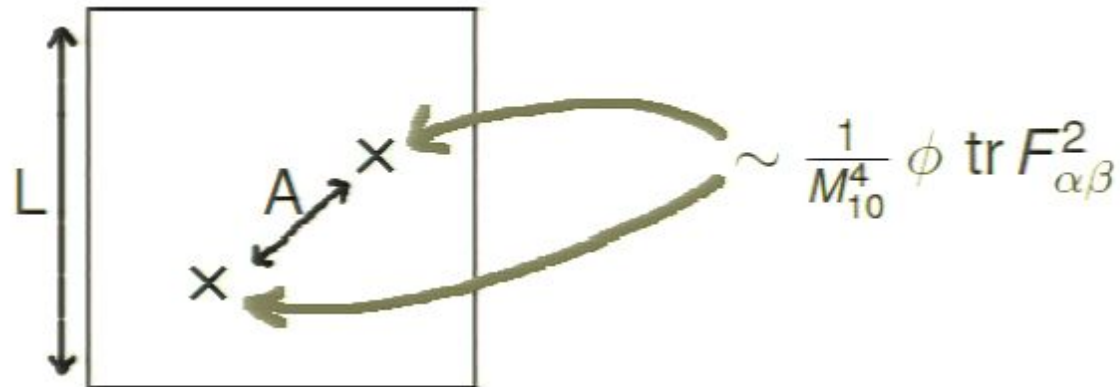
- World-volume theory on stack of N D3-branes is $\mathcal{N} = 4$ $U(N)$ super Yang-Mills theory.
 - The two brane stacks are coupled by SUGRA fields in the embedding space.
 - Gauge theories are strongly coupled, but relevant process is protected by a non-renormalization theorem (Dixon, Freedman).

HEAT TRANSFER RATE



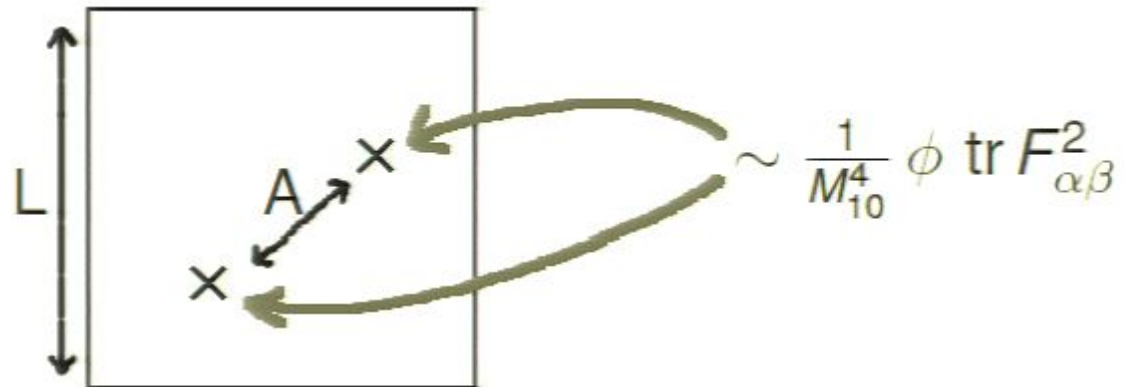
- World-volume theory on stack of N D3-branes is $\mathcal{N} = 4$ $U(N)$ super Yang-Mills theory.
- The two brane stacks are coupled by SUGRA fields in the embedding space.
- Gauge theories are strongly coupled, but relevant process is protected by a non-renormalization theorem (Seiberg-Witten).

HEAT TRANSFER RATE



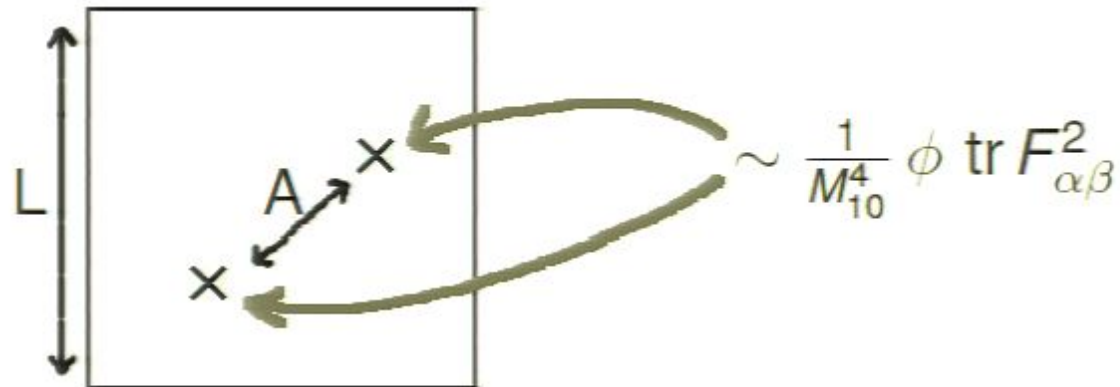
- World-volume theory on stack of N D3-branes is $\mathcal{N} = 4$ $U(N)$ super Yang-Mills theory.
- The two brane stacks are coupled by SUGRA fields in the embedding space.
- Gauge theories are strongly coupled, but relevant process is protected by a non-renormalization theorem (Seiberg-Witten).

HEAT TRANSFER RATE



- World-volume theory on stack of N D3-branes is $\mathcal{N} = 4$ $U(N)$ super Yang-Mills theory.
- The two brane stacks are coupled by SUGRA fields in the embedding space.
- Gauge theories are strongly coupled, but relevant process is protected by a non-renormalization theorem (Seiberg-Witten).

HEAT TRANSFER RATE



- World-volume theory on stack of N D3-branes is $\mathcal{N} = 4$ U(N) super Yang-Mills theory.
- The two brane stacks are coupled by SUGRA fields in the embedding space.
- Gauge theories are strongly coupled, but relevant process is protected by a non-renormalization theorem [Gubser, Klebanov].

HEAT TRANSFER RATE

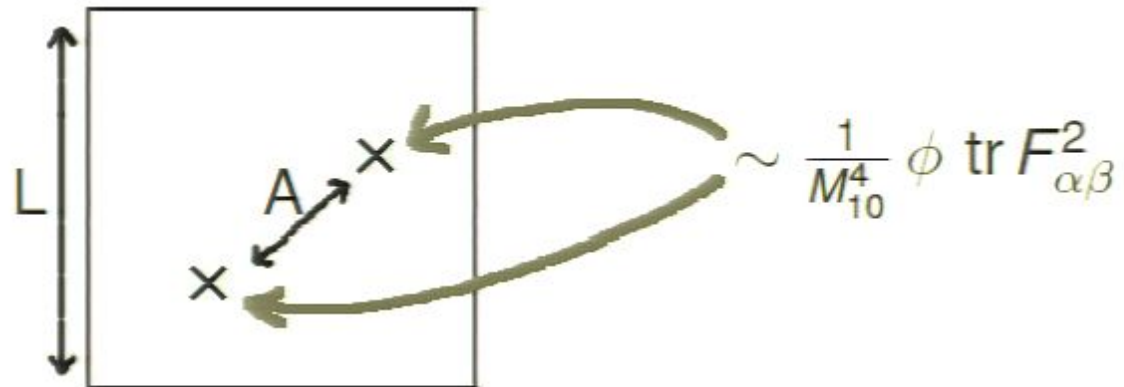
- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



- Thermally averaging the above process gives heat transfer rate (per 4d volume):

$$\dot{p} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE



- World-volume theory on stack of N D3-branes is $\mathcal{N} = 4$ $U(N)$ super Yang-Mills theory.
- The two brane stacks are coupled by SUGRA fields in the embedding space.
- Gauge theories are strongly coupled, but relevant process is protected by a non-renormalization theorem [Gubser, Klebanov].

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type

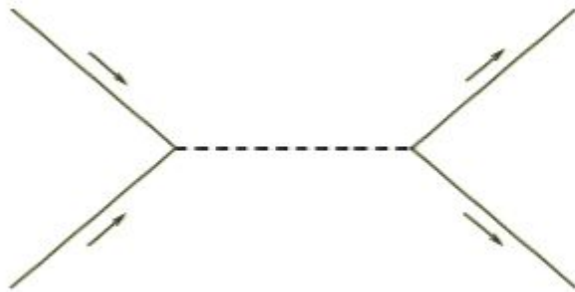


- Thermally averaging the above process gives heat transfer rate (per 4d volume):

$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



summing
over KK
modes

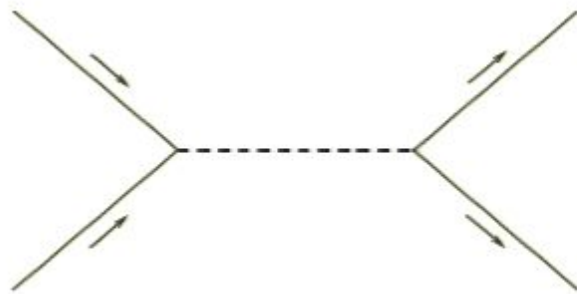
$$\begin{aligned}
 \mathcal{M} &\sim \frac{g^2}{M_{10}^6 L^6} \sum_{n \in \mathbb{Z}^6} \frac{e^{2i n \cdot \Delta / L}}{s - m_n^2 + i\epsilon} \\
 &\sim \frac{g^2}{M_{10}^6 A^2} + \frac{g^2}{M_{10}^6 L^6}
 \end{aligned}$$

- Thermally averaging the above process gives heat transfer rate (per 4d volume):

$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9$$

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

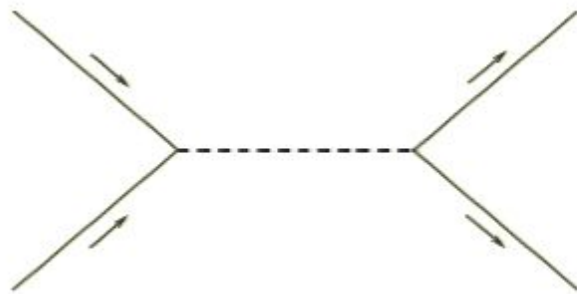
$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives heat transfer rate (per 4d volume):

$$\dot{p} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

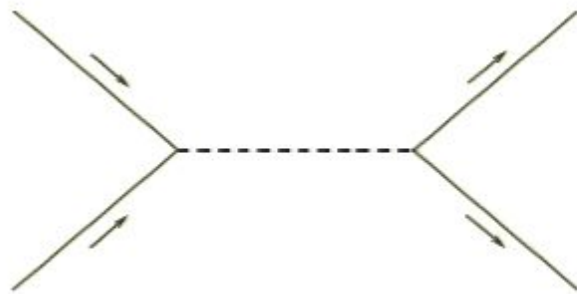
$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives **heat transfer rate** (per 4d volume):

$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

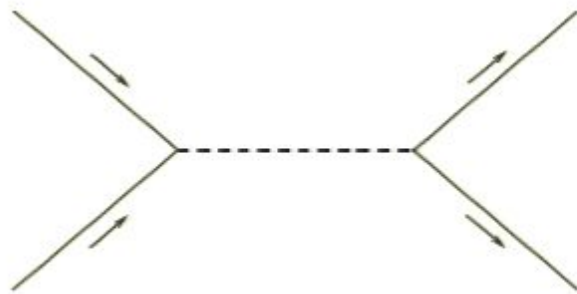
$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives heat transfer rate (per 4d volume):

$$\rho \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

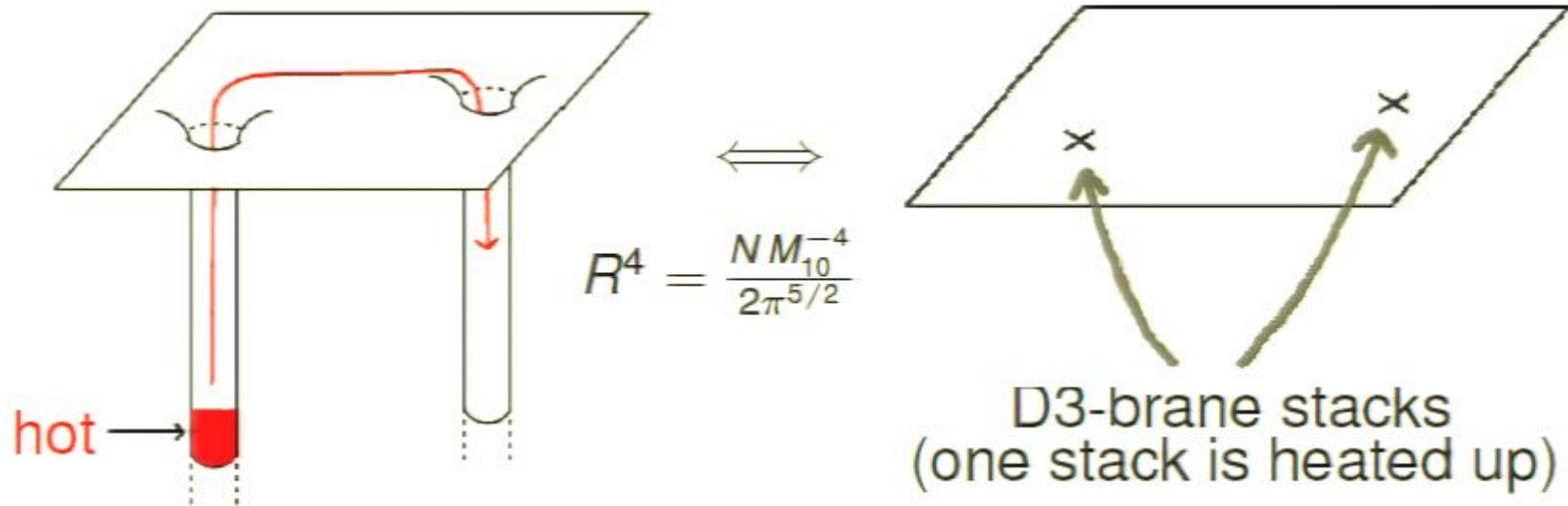
$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives **heat transfer rate** (per 4d volume):

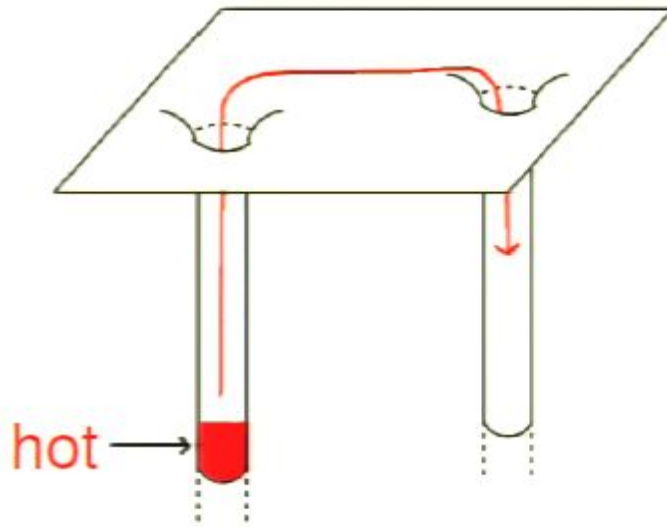
$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE



- We consider two $\text{AdS}_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by **transition probability** of particles between the two throats.
- Problem turns out to be a **multi-dimensional tunneling problem**. \Rightarrow Difficult to solve.
- Way out via **AdS/CFT duality**: String theory on $\text{AdS}_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

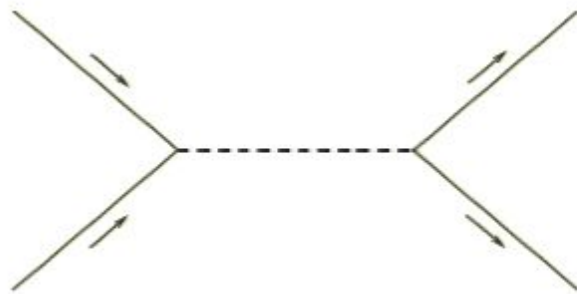
HEAT TRANSFER RATE



- We consider two $\text{AdS}_5 \times S^5$ throats embedded in a 6d torus.
- Heat transfer rate determined by **transition probability** of particles between the two throats.
- Problem turns out to be a **multi-dimensional tunneling problem**. \Rightarrow Difficult to solve.
- Way out via AdS/CFT duality: String theory on $\text{AdS}_5 \times S^5$ is dual to $\mathcal{N} = 4$ $U(N)$ gauge theory realized on stack of D3-branes.

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives **heat transfer rate** (per 4d volume):

$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

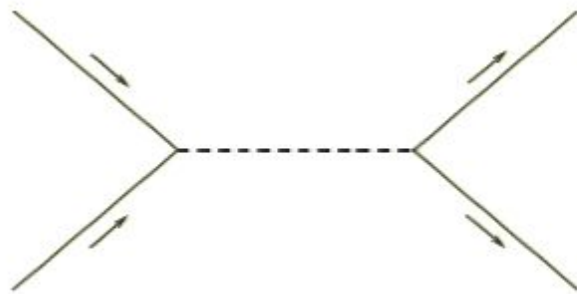
THROAT DARK MATTER: SETUP

- Consider a compactification with a **large number** of throats.
- Assume that standard model lives on some D-branes in the unwarped part of the compact space.
- Assume that only the standard model is heated up directly by the reheating mechanism.
- Throats are then heated up by heat transfer from the hot standard model.



HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

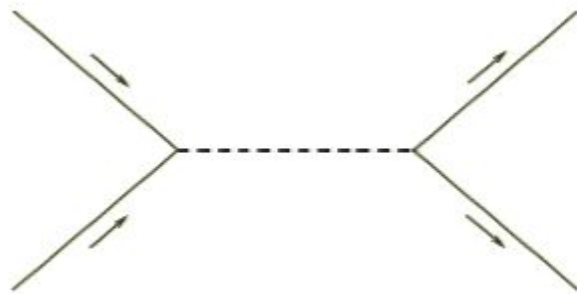
$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives **heat transfer rate** (per 4d volume):

$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

HEAT TRANSFER RATE

- KK expansion of SUGRA fields. \Rightarrow From 4d viewpoint one has two gauge theories coupled by tower of KK modes.
- Heated throat corresponds to heated gauge theory. Energy transfer due to processes of the type



\Rightarrow
summing
over KK
modes

$$\mathcal{M} \sim \frac{s^2}{M_{10}^8 L^6} \sum_{\vec{n} \in \mathbb{Z}^6} \frac{e^{2\pi i \vec{n} \vec{A}/L}}{s - m_{\vec{n}}^2 + i\epsilon}$$

$$\sim \frac{s^2}{M_{10}^8 A^4} + \frac{s}{M_{10}^8 L^6}.$$

- Thermally averaging the above process gives **heat transfer rate** (per 4d volume):

$$\dot{\rho} \sim \frac{R_1^8 R_2^8}{A^8} T^{13} + \frac{R_1^8 R_2^8}{L^{12}} T^9.$$

THROAT DARK MATTER: SETUP

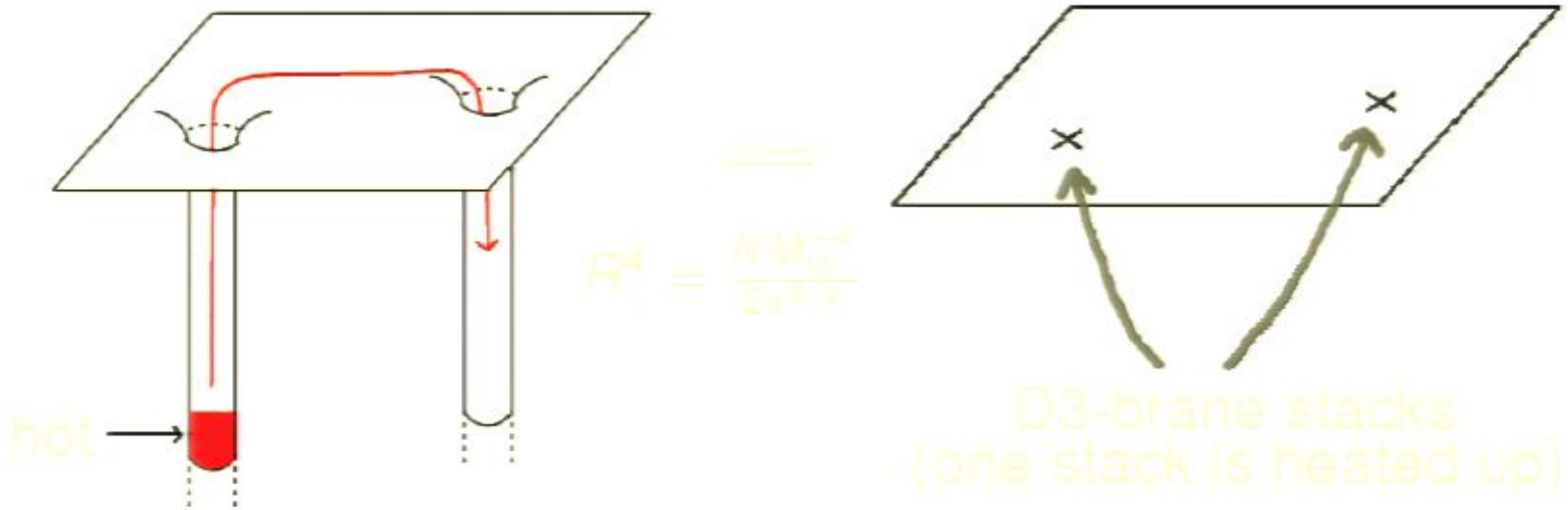
- Consider a compactification with a **large number** of throats.
- Assume that standard model lives on some D-branes in the unwarped part of the compact space.
- Assume that only the standard model is heated up directly by the reheating mechanism.
- Throats are then heated up by heat transfer from the hot standard model.



HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
 ⇒ **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

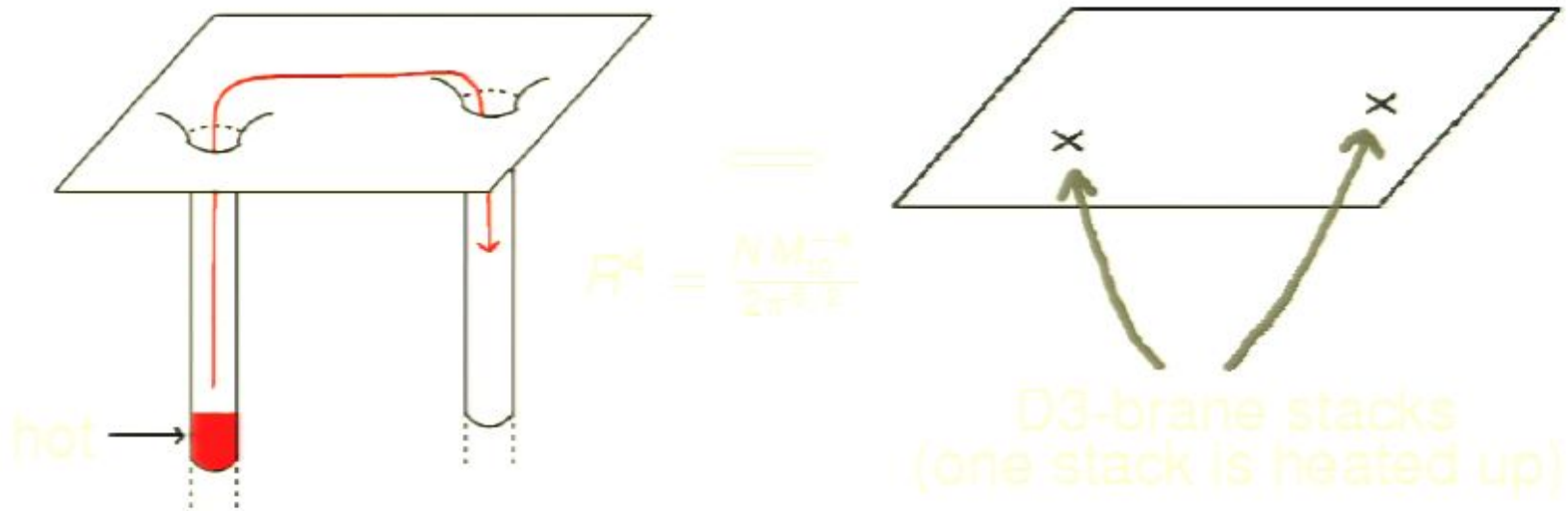
$$\dot{h} \sim \frac{g R^8}{M_{10}^8 A^5} T^{13} \sim \frac{g R^8}{M_{10}^8 L^{12}} T^9$$



HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
 ⇒ **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

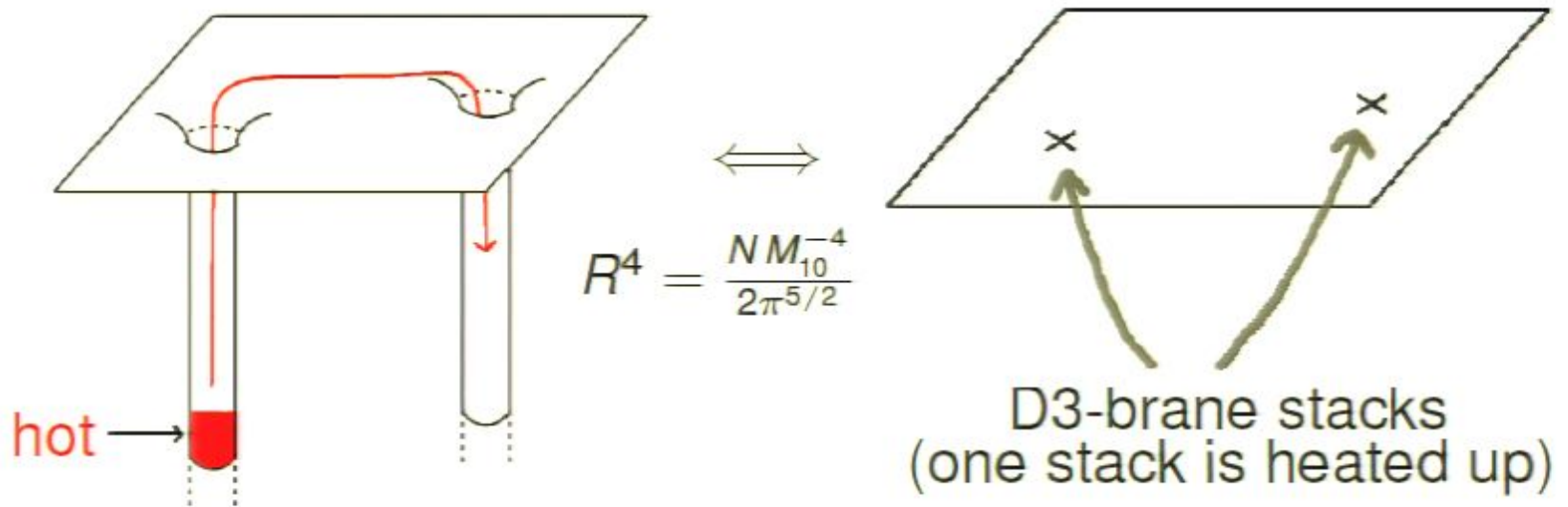
$$\dot{h} \sim \frac{g R^8}{M_{10}^8 A^8} T^{13} = \frac{g R^8}{M_{10}^8 L^{12}} T^9$$



HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
 \Rightarrow **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

$$\dot{\rho} \sim \frac{g R^8}{M_{10}^8 A^8} T^{13} + \frac{g R^8}{M_{10}^8 L^{12}} T^9.$$



HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
⇒ **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

$$\dot{\rho} \sim \frac{g R^8}{M_{10}^8 A^8} T^{13} + \frac{g R^8}{M_{10}^8 L^{12}} T^9.$$

- Use relations $R^8 = N^2 M_{10}^{-8}$ (N^2 number of dof of dual gauge theory) and $M_4^2 \sim L^6 M_{10}^8$.
- Restrict to setup with $A \sim L$ and $T < L^{-1}$. ⇒ Second term dominates. ⇒ Heat transfer rate from SM to throat is

$$\dot{\rho} \sim g N^2 \frac{T^9}{M_4^4}.$$

HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
⇒ **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

$$\dot{\rho} \sim \frac{g R^8}{M_{10}^8 A^8} T^{13} + \frac{g R^8}{M_{10}^8 L^{12}} T^9.$$

- Use relations $R^8 = N^2 M_{10}^{-8}$ (N^2 number of dof of dual gauge theory) and $M_4^2 \sim L^6 M_{10}^8$.
- Restrict to setup with $A \sim L$ and $T < L^{-1}$. ⇒ Second term dominates. ⇒ **Heat transfer rate from SM to throat is**

$$\dot{\rho} \sim g N^2 \frac{T^9}{M_4^4}.$$

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} scale like radiation ($\rho_{\text{throat}} \propto a^{-4}$) and how long like matter ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in gauge theory picture.
- Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
- Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
⇒ **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

$$\dot{\rho} \sim \frac{g R^8}{M_{10}^8 A^8} T^{13} + \frac{g R^8}{M_{10}^8 L^{12}} T^9.$$

- Use relations $R^8 = N^2 M_{10}^{-8}$ (N^2 number of dof of dual gauge theory) and $M_4^2 \sim L^6 M_{10}^8$.
- Restrict to setup with $A \sim L$ and $T < L^{-1}$. ⇒ Second term dominates. ⇒ **Heat transfer rate from SM to throat is**

$$\dot{\rho} \sim g N^2 \frac{T^9}{M_4^4}.$$

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} scale like radiation ($\rho_{\text{throat}} \propto a^{-4}$) and how long like matter ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in gauge theory picture.
- Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
- Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

HEAT TRANSFER TO THE THROAT

- Calculation of heat transfer rate used D-brane description.
⇒ **Result valid for SM on small number of D-branes.**
- This gives ($g \sim 100$ number of dof of SM):

$$\dot{\rho} \sim \frac{g R^8}{M_{10}^8 A^8} T^{13} + \frac{g R^8}{M_{10}^8 L^{12}} T^9.$$

- Use relations $R^8 = N^2 M_{10}^{-8}$ (N^2 number of dof of dual gauge theory) and $M_4^2 \sim L^6 M_{10}^8$.
- Restrict to setup with $A \sim L$ and $T < L^{-1}$. ⇒ Second term dominates. ⇒ **Heat transfer rate from SM to throat is**

$$\dot{\rho} \sim g N^2 \frac{T^9}{M_4^4}.$$

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} scale like radiation ($\rho_{\text{throat}} \propto a^{-4}$) and how long like matter ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in gauge theory picture.
- Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
- Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} scale like radiation ($\rho_{\text{throat}} \propto a^{-4}$) and how long like matter ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in gauge theory picture.
- Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
- Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} scale like radiation ($\rho_{\text{throat}} \propto a^{-4}$) and how long like matter ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in gauge theory picture.
- Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
- Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} **scale like radiation** ($\rho_{\text{throat}} \propto a^{-4}$) and how long **like matter** ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in gauge theory picture.
- Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
- Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} **scale like radiation** ($\rho_{\text{throat}} \propto a^{-4}$) and how long **like matter** ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in **gauge theory picture**.
 - Generically, throats have finite length (as in RS1). In dual gauge theory, this corresponds to existence of a confinement scale Λ .
 - Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} **scale like radiation** ($\rho_{\text{throat}} \propto a^{-4}$) and how long **like matter** ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in **gauge theory picture**.
- Generically, throats have **finite length** (as in RS1). In dual gauge theory, this corresponds to existence of a **confinement scale Λ** .
- **Another entry in AdS/CFT dictionary: KK particles in throat correspond to glueballs of dual gauge theory.**

ENERGY DENSITY IN THE THROAT SECTOR

- **Energy density** in throat directly after reheating then is

$$\rho_{\text{throat,RH}} \sim g^{1/2} N^2 \frac{T_{\text{RH}}^7}{M_4^3} \sim \dot{\rho} \cdot H_{\text{RH}}^{-1}.$$

Important for late-time abundance: How long does ρ_{throat} **scale like radiation** ($\rho_{\text{throat}} \propto a^{-4}$) and how long **like matter** ($\rho_{\text{throat}} \propto a^{-3}$) ?

- Convenient to discuss this in **gauge theory picture**.
- Generically, throats have **finite length** (as in RS1). In dual gauge theory, this corresponds to existence of a **confinement scale Λ** .
- Another entry in AdS/CFT dictionary: **KK particles** in throat correspond to **glueballs** of dual gauge theory.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

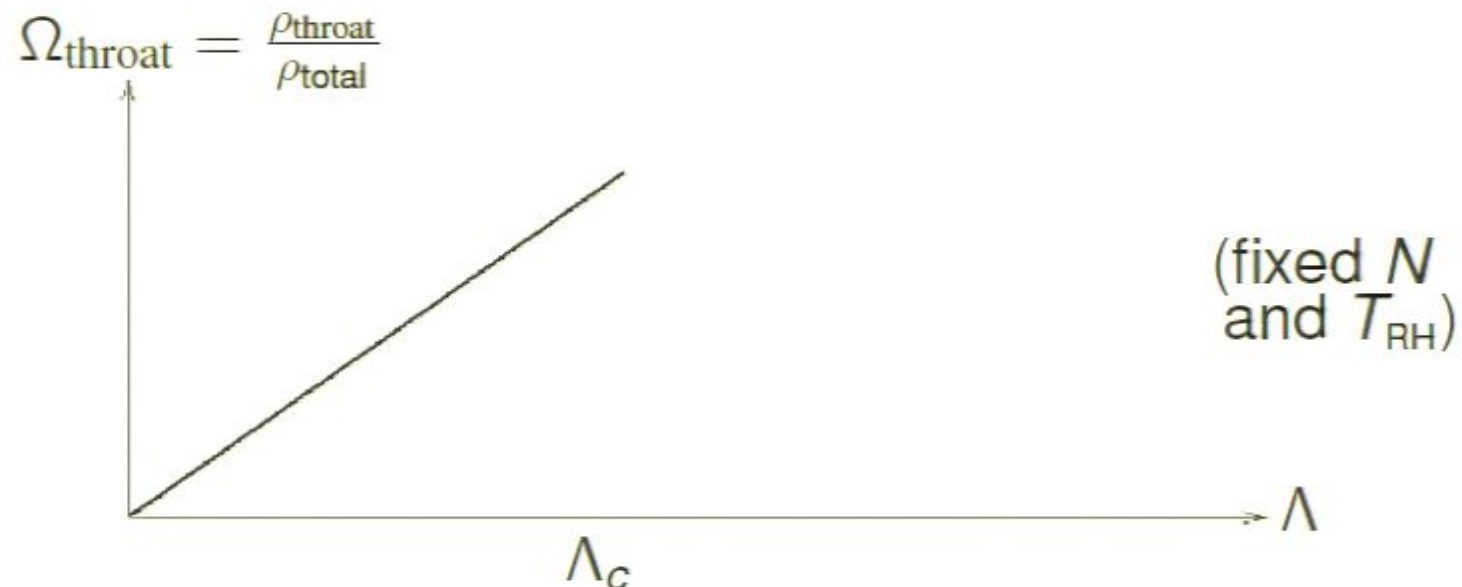
- If $\rho_{\text{throat,RH}}$ larger than critical energy density $\rho_{\text{cr}} \sim N^2 \Lambda^4$:
Gauge theory in deconfined phase.
 - $\rho_{\text{throat}} \propto a^{-4}$ until temperature drops to $\Lambda \Rightarrow$ confinement phase transition $\Rightarrow \rho_{\text{throat}} \propto a^{-3}$ afterwards.
 - Contribution of glueballs to total energy density Ω_{throat} as function of Λ (a throat with $\Lambda = \Lambda_c$ has $\rho_{\text{throat,RH}} = \rho_{\text{cr}}$).

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}}$ larger than critical energy density $\rho_{\text{cr}} \sim N^2 \Lambda^4$:
Gauge theory in deconfined phase.
- $\rho_{\text{throat}} \propto a^{-4}$ until temperature drops to $\Lambda \Rightarrow$ confinement phase transition $\Rightarrow \rho_{\text{throat}} \propto a^{-3}$ afterwards.
- Contribution of glueballs to total energy density Ω_{throat} as function of Λ (a throat with $\Lambda = \Lambda_c$ has $\rho_{\text{throat,RH}} = \rho_{\text{cr}}$):

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}}$ larger than critical energy density $\rho_{\text{cr}} \sim N^2 \Lambda^4$:
Gauge theory in deconfined phase.
- $\rho_{\text{throat}} \propto a^{-4}$ until temperature drops to $\Lambda \Rightarrow$ confinement phase transition $\Rightarrow \rho_{\text{throat}} \propto a^{-3}$ afterwards.
- **Contribution of glueballs to total energy density Ω_{throat} as function of Λ (a throat with $\Lambda = \Lambda_c$ has $\rho_{\text{throat,RH}} = \rho_{\text{cr}}$):**



LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):

Gauge theory in confined phase directly after reheating.

- Recall that heat transfer to throat is due to processes of the type:

- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.

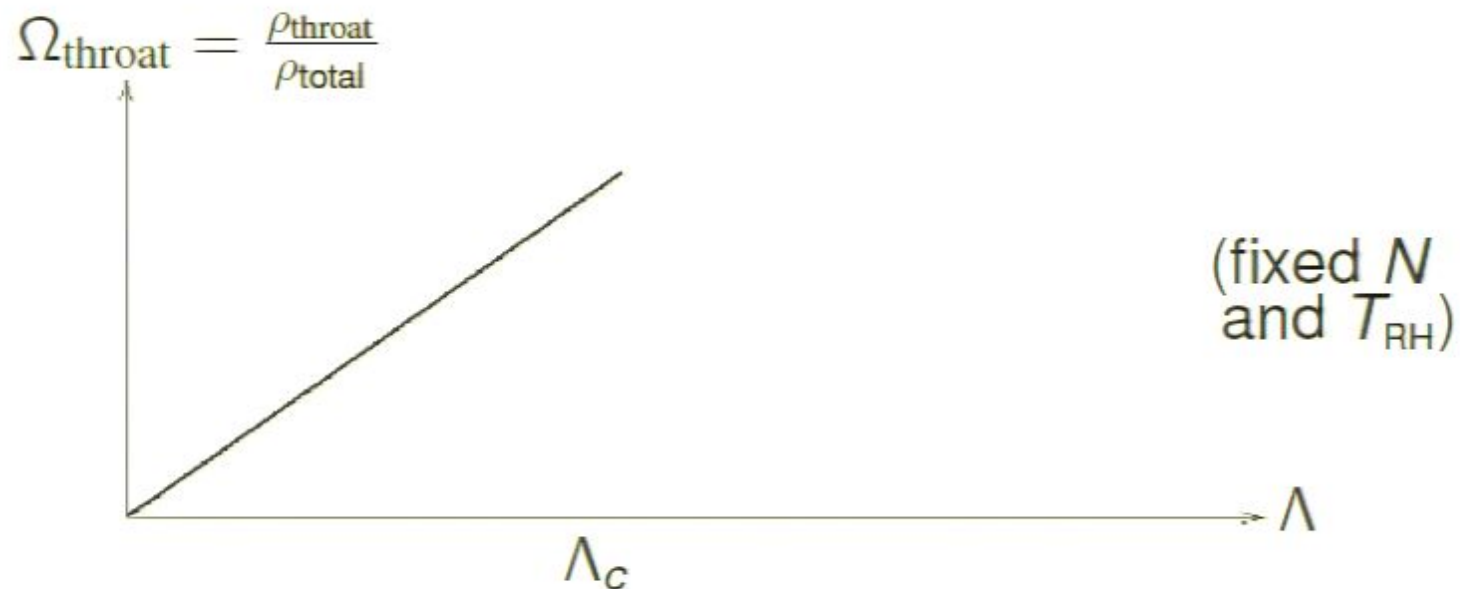
- Different in large- N gauge theories:

Decay leads to large number of nonrelativistic glueballs:

Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}}$ larger than critical energy density $\rho_{\text{cr}} \sim N^2 \Lambda^4$:
Gauge theory in deconfined phase.
- $\rho_{\text{throat}} \propto a^{-4}$ until temperature drops to $\Lambda \Rightarrow$ confinement phase transition $\Rightarrow \rho_{\text{throat}} \propto a^{-3}$ afterwards.
- **Contribution of glueballs to total energy density Ω_{throat} as function of Λ (a throat with $\Lambda = \Lambda_c$ has $\rho_{\text{throat,RH}} = \rho_{\text{cr}}$):**



LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):

Gauge theory in confined phase directly after reheating.

- Recall that heat transfer to throat is due to processes of the type:

- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.

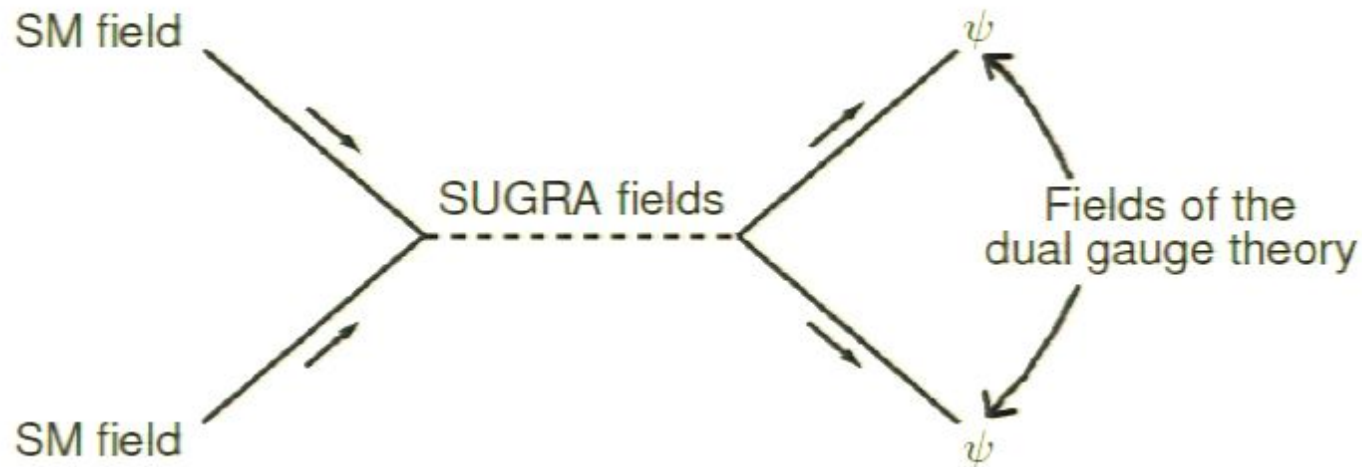
- Different in large- N gauge theories:

Decay leads to large number of nonrelativistic glueballs.

Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

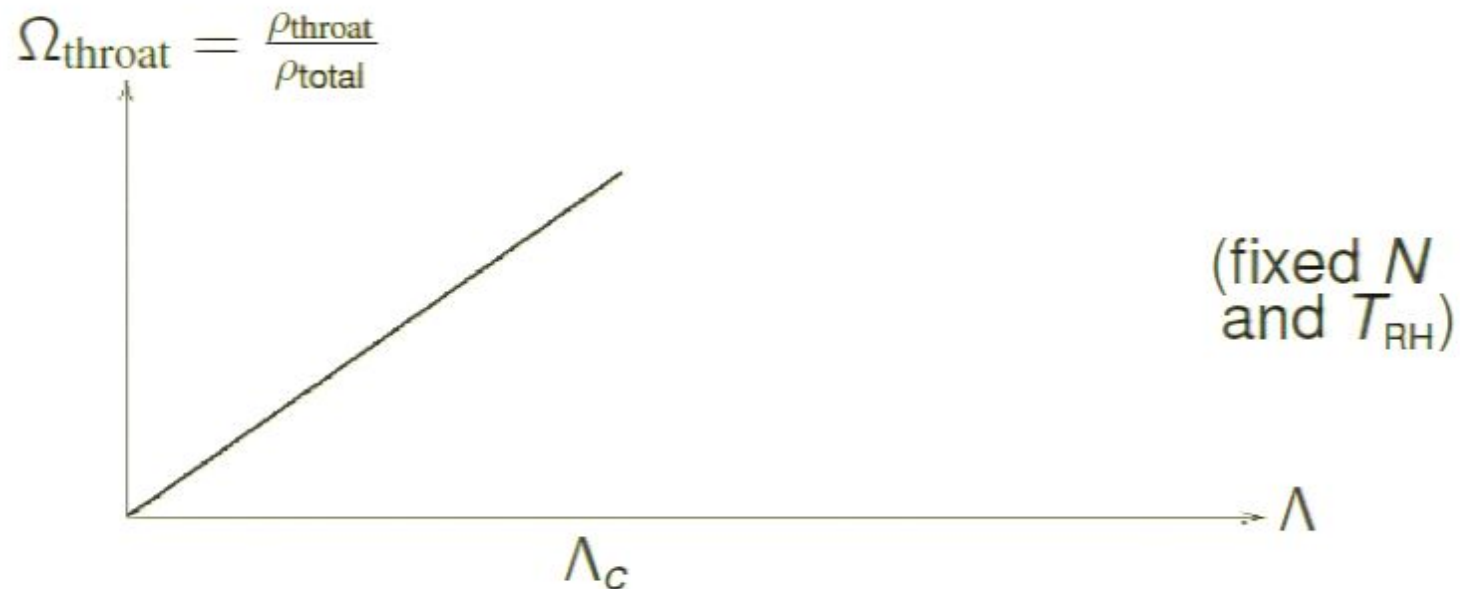
- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):
Gauge theory in confined phase directly after reheating.
- Recall that heat transfer to throat is due to processes of the type:



- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.
- Different in large- N gauge theories:
Decay leads to large number of nonrelativistic glueballs.
Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}}$ larger than critical energy density $\rho_{\text{cr}} \sim N^2 \Lambda^4$:
Gauge theory in deconfined phase.
- $\rho_{\text{throat}} \propto a^{-4}$ until temperature drops to $\Lambda \Rightarrow$ confinement phase transition $\Rightarrow \rho_{\text{throat}} \propto a^{-3}$ afterwards.
- **Contribution of glueballs to total energy density Ω_{throat} as function of Λ (a throat with $\Lambda = \Lambda_c$ has $\rho_{\text{throat,RH}} = \rho_{\text{cr}}$):**



LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):

Gauge theory in confined phase directly after reheating.

- Recall that heat transfer to throat is due to processes of the type:

- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.

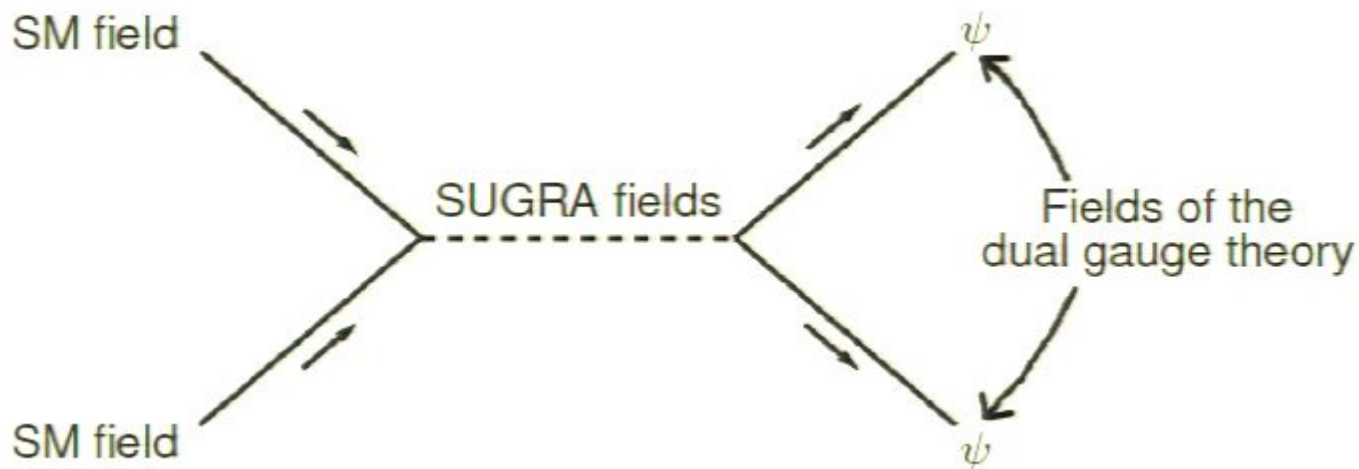
- Different in large- N gauge theories:

Decay leads to large number of nonrelativistic glueballs.

Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):
Gauge theory in confined phase directly after reheating.
- Recall that heat transfer to throat is due to processes of the type:



- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.
- Different in large- N gauge theories:
Decay leads to large number of nonrelativistic glueballs.
Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):

Gauge theory in confined phase directly after reheating.

- Recall that heat transfer to throat is due to processes of the type:

- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.

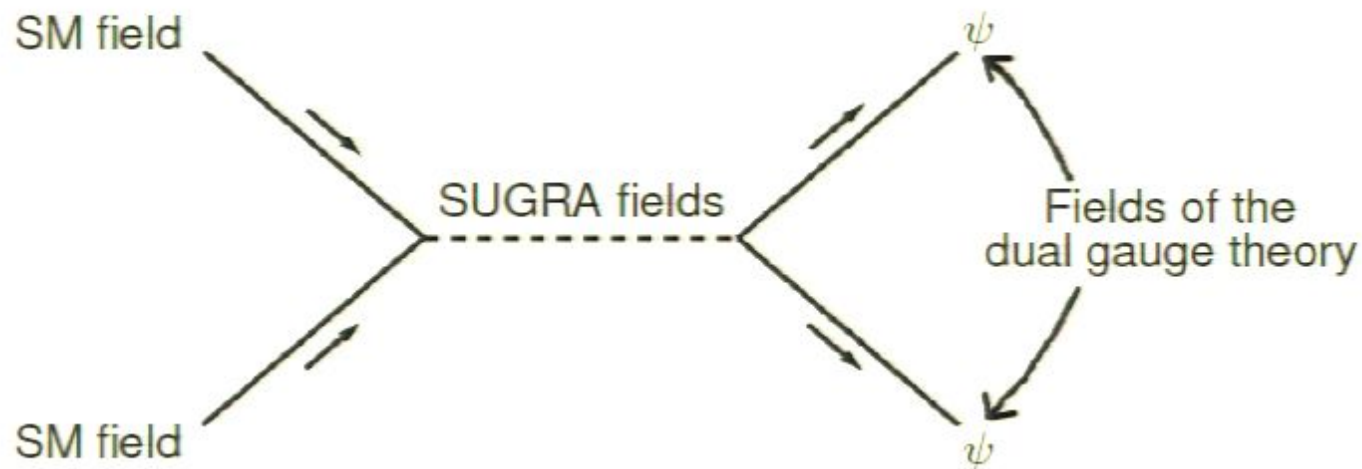
- Different in large- N gauge theories:

Decay leads to large number of nonrelativistic glueballs.

Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

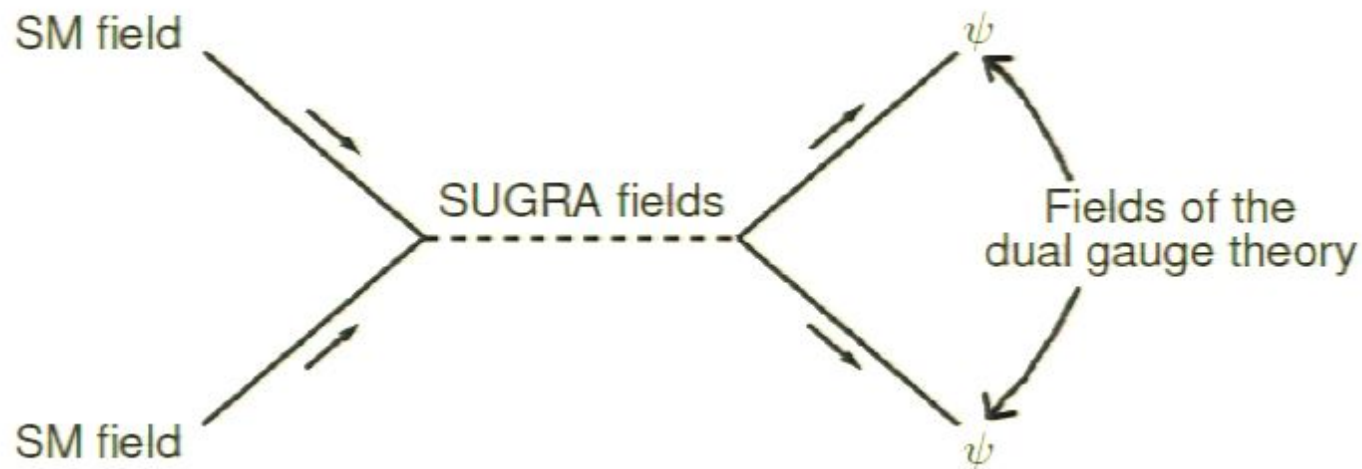
- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):
Gauge theory in confined phase directly after reheating.
- Recall that heat transfer to throat is due to processes of the type:



- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.
- Different in large- N gauge theories:
Decay leads to large number of nonrelativistic glueballs.
Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

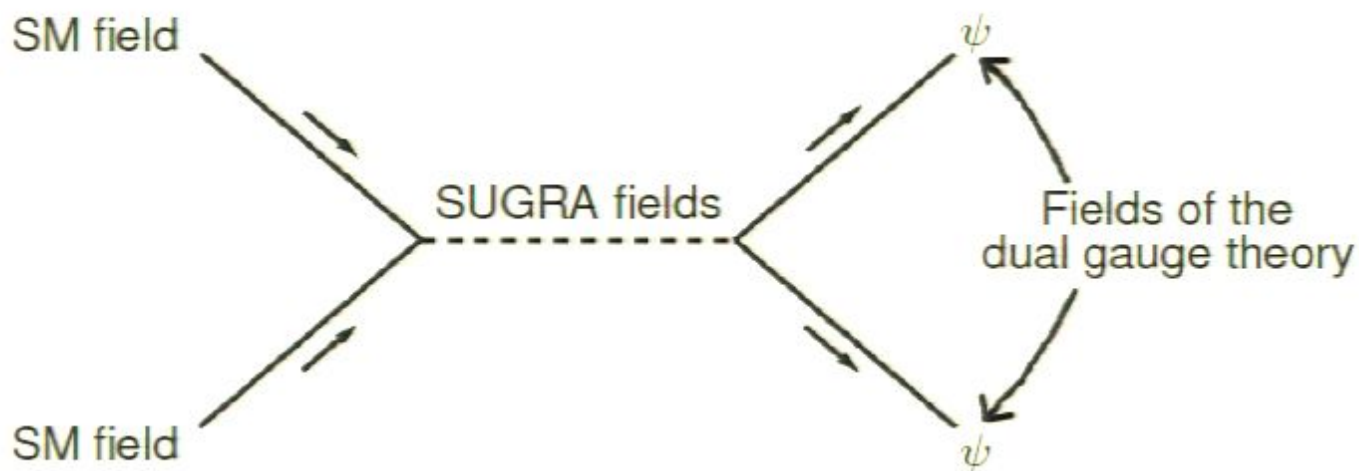
- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):
Gauge theory in confined phase directly after reheating.
- Recall that heat transfer to throat is due to processes of the type:



- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.
- Different in large-N gauge theories:
Decay leads to large number of nonrelativistic glueballs.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

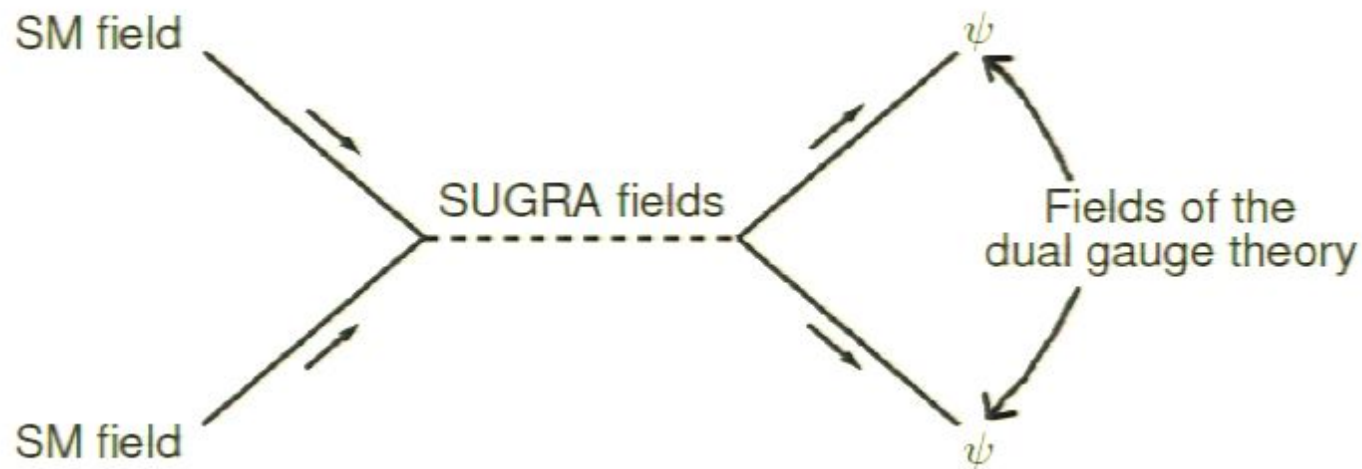
- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):
Gauge theory in confined phase directly after reheating.
- Recall that heat transfer to throat is due to processes of the type:



- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.
- Different in large-N gauge theories:
Decay leads to large number of nonrelativistic glueballs.

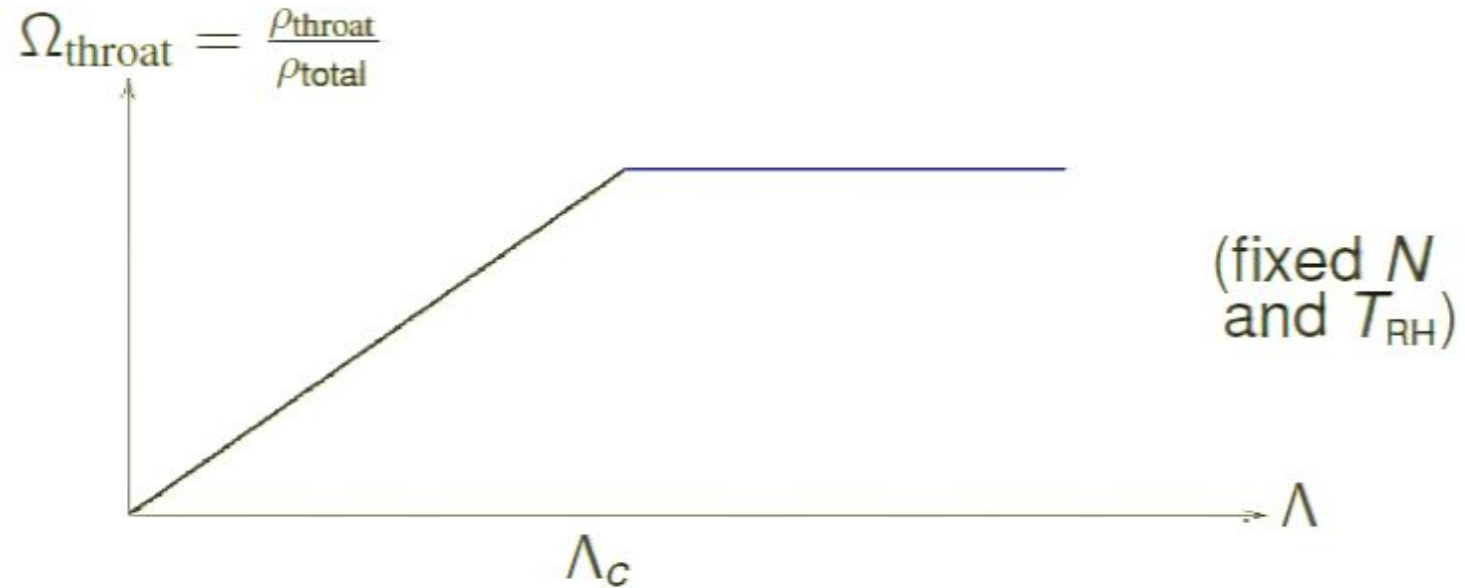
LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

- If $\rho_{\text{throat,RH}} < \rho_{\text{cr}}$ (or $\Lambda > \Lambda_c$):
Gauge theory in confined phase directly after reheating.
- Recall that heat transfer to throat is due to processes of the type:



- In QCD: Decay to two quarks leads to two jets and to small number of ultrarelativistic hadrons.
- Different in large-N gauge theories:
Decay leads to large number of nonrelativistic glueballs.
- Then, $\rho_{\text{throat}} \propto a^{-3}$ immediately after reheating.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS



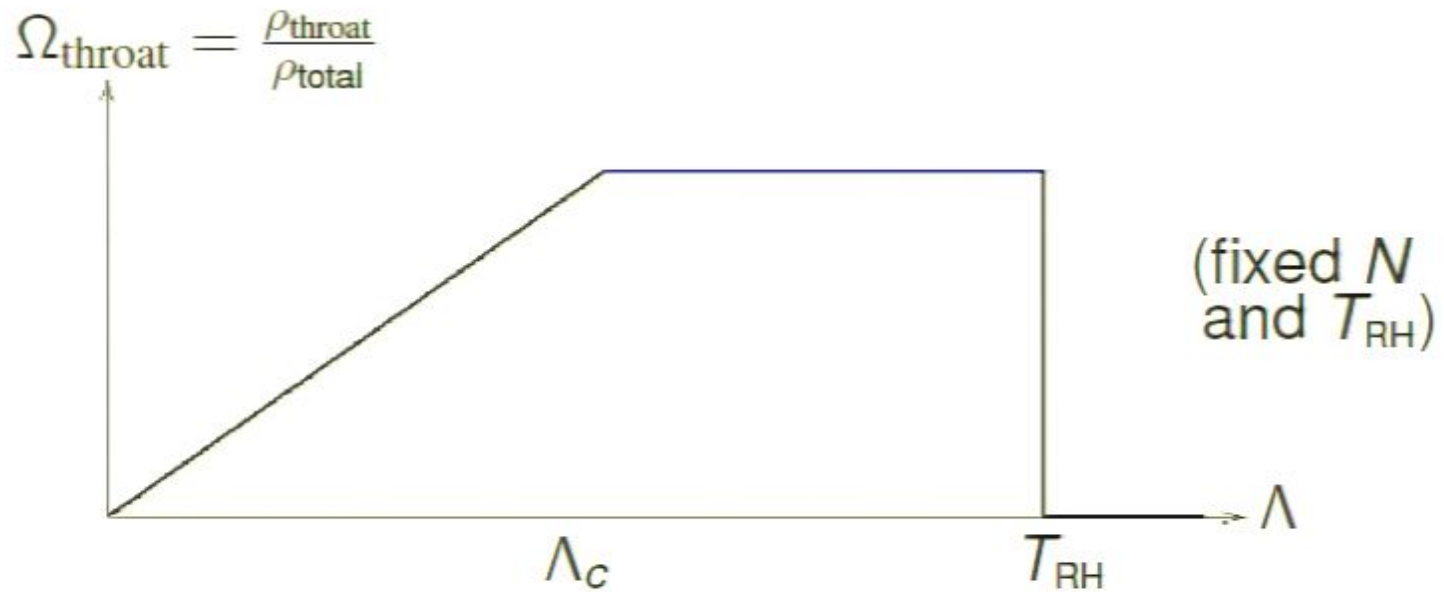
• For throats with $\Lambda_c < \Lambda < T_{RH}$ we find

$$\Omega_{\text{throat}} \sim \left(\frac{T_{RH} N^{1/2}}{6 \cdot 10^{11} \text{ GeV}} \right)^4$$

• For $N \sim 10$ (typical value) and $T_{RH} \sim 10^{11} \text{ GeV}$: $\Omega_{\text{throat}} = \mathcal{O}(1)$.

• These throats can account for the dark matter.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS



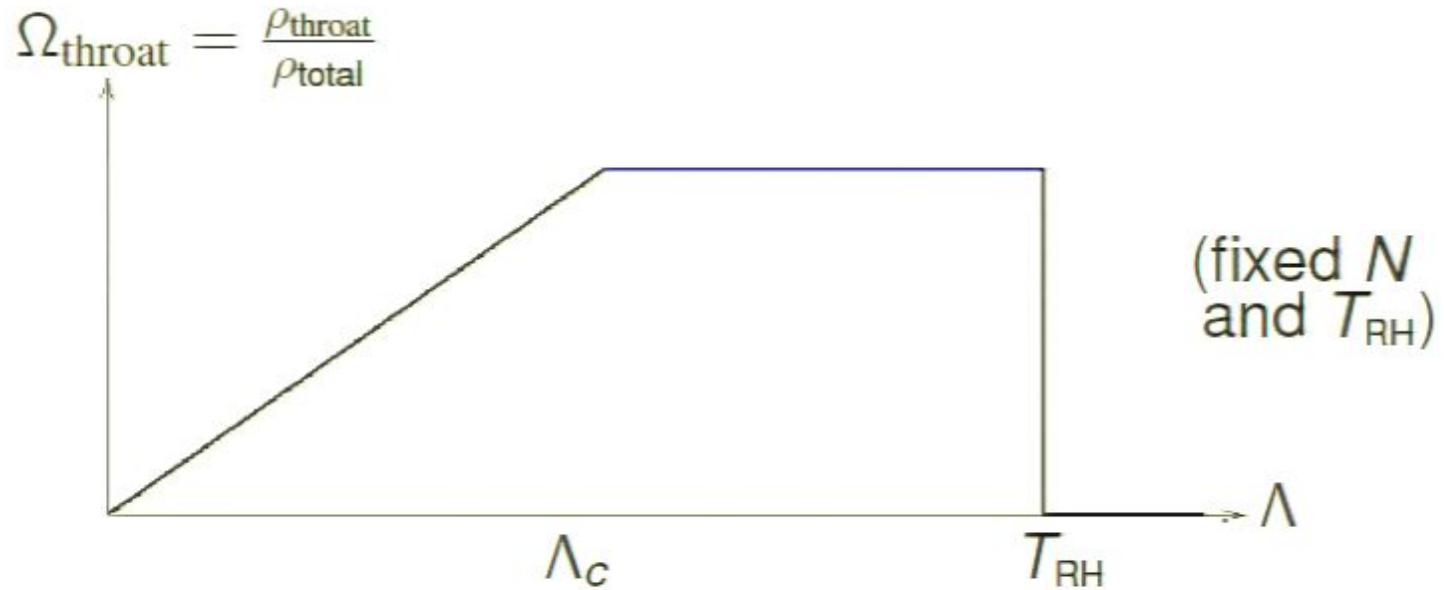
• For throats with $\Lambda_c < \Lambda < T_{\text{RH}}$ we find

$$\Omega_{\text{throat}} \sim \left(\frac{T_{\text{RH}} N^{1/2}}{6 \cdot 10^{11} \text{ GeV}} \right)^4$$

• For $N \sim 10$ (typical value) and $T_{\text{RH}} \sim 10^{11} \text{ GeV}$: $\Omega_{\text{throat}} = \mathcal{O}(1)$.

• These throats can account for the dark matter.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS



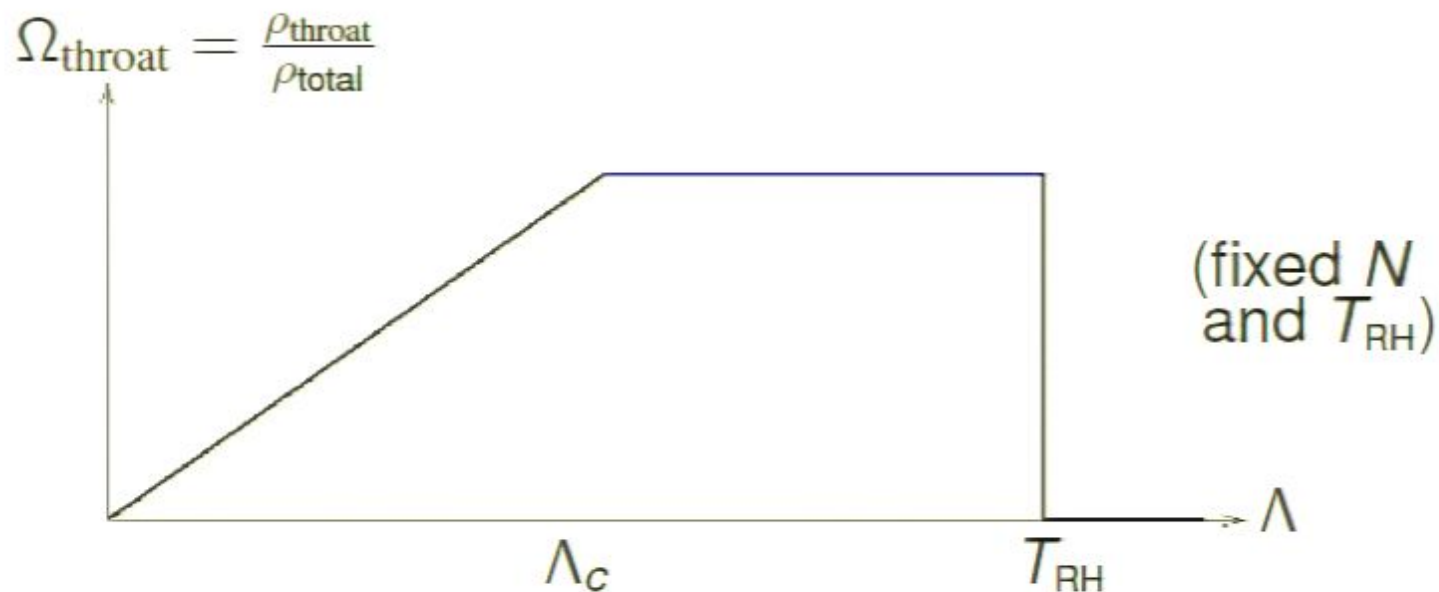
- For throats with $\Lambda_c < \Lambda < T_{\text{RH}}$ we find

$$\Omega_{\text{throat}} \sim \left(\frac{T_{\text{RH}} N^{1/2}}{6 \cdot 10^{11} \text{ GeV}} \right)^4 .$$

• For $N \sim 10$ (typical value) and $T_{\text{RH}} \sim 10^{11} \text{ GeV}$: $\Omega_{\text{throat}} = \mathcal{O}(1)$.

• These throats can account for the dark matter.

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS



- For throats with $\Lambda_c < \Lambda < T_{\text{RH}}$ we find

$$\Omega_{\text{throat}} \sim \left(\frac{T_{\text{RH}} N^{1/2}}{6 \cdot 10^{11} \text{ GeV}} \right)^4.$$

- For $N \sim 10$ (typical value) and $T_{\text{RH}} \sim 10^{11} \text{ GeV}$: $\Omega_{\text{throat}} = \mathcal{O}(1)$.

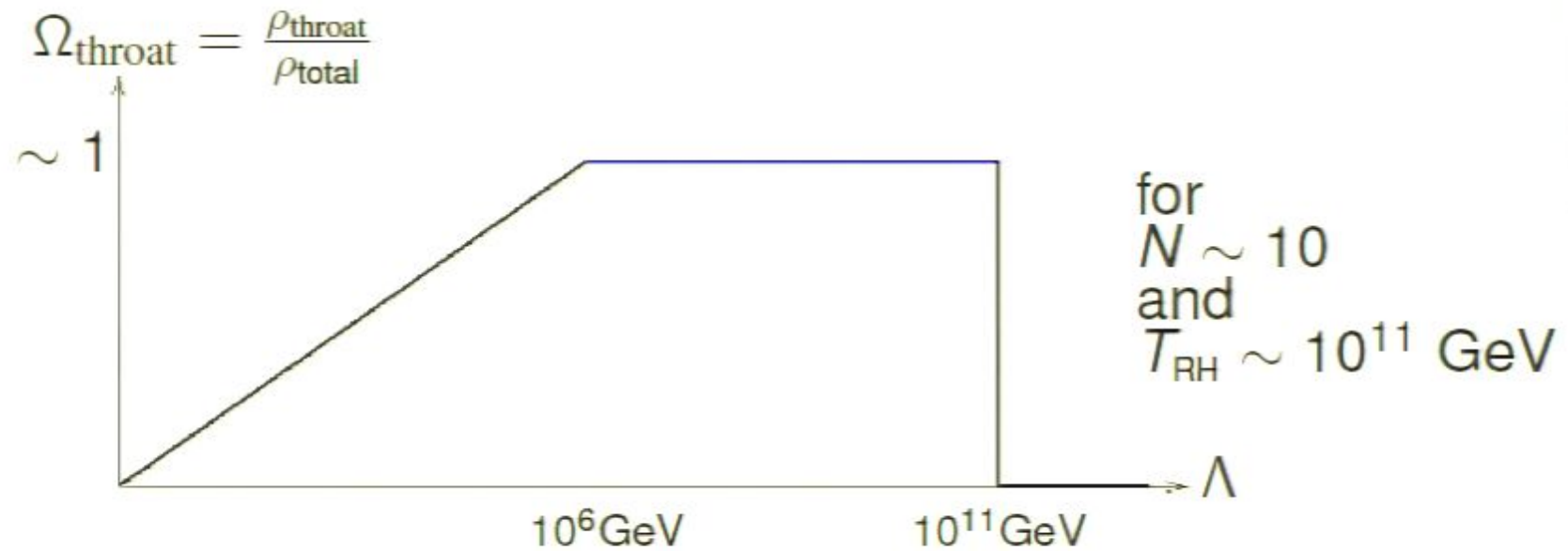
LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS



- Λ_c is a function of N and T_{RH} . Using $N \sim 10$ and $T_{\text{RH}} \sim 10^{11} \text{ GeV}$, we find $\Lambda_c \sim 10^6 \text{ GeV}$.

- String vacua with ~ 10 throats in this range probably not uncommon. But unfortunately, there is a problem ...

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

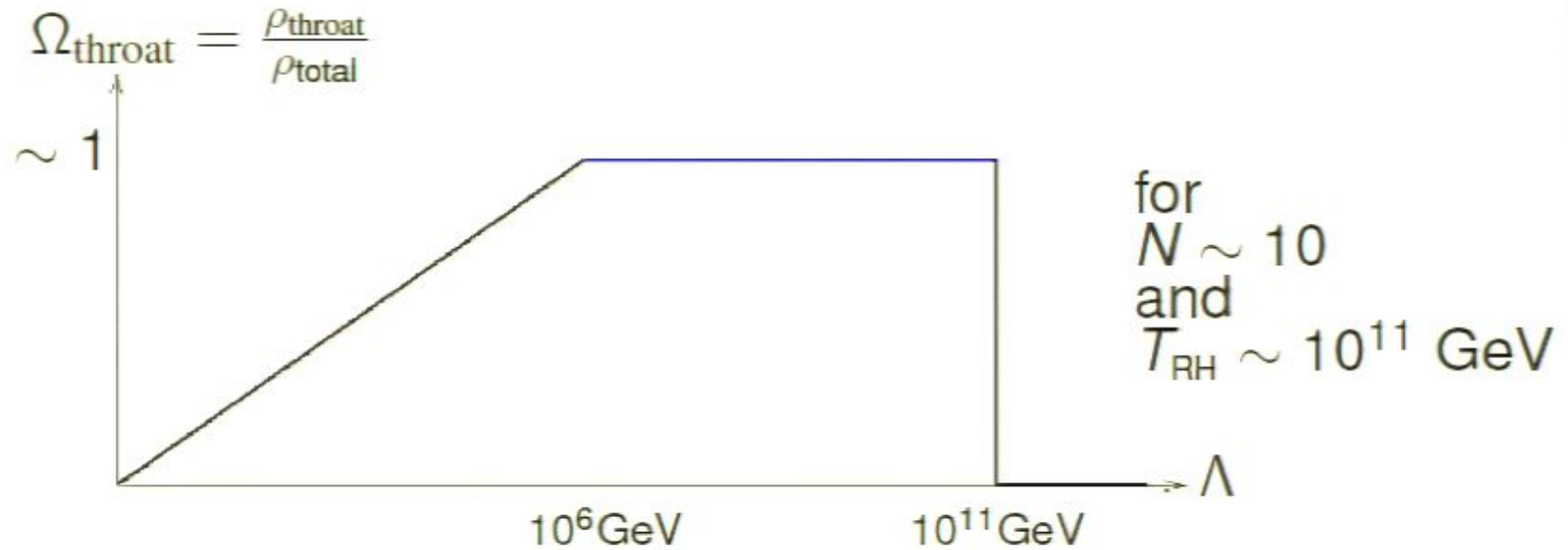


- Λ_c is a function of N and T_{RH} . Using $N \sim 10$ and $T_{\text{RH}} \sim 10^{11}$ GeV, we find $\Lambda_c \sim 10^6$ GeV.

Throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) could account for the dark matter if reheating temperature was $\sim 10^{11}$ GeV.

- String vacua with ~ 10 throats in this range probably not uncommon. But unfortunately, there is a problem ...

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS

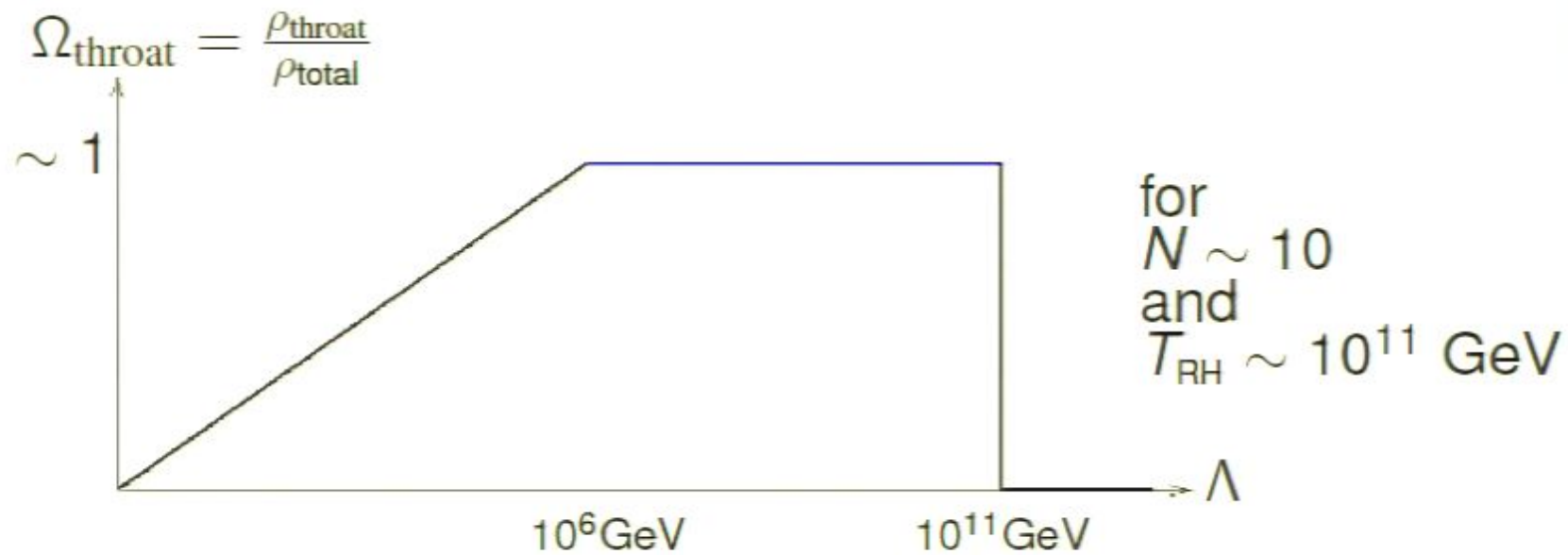


- Λ_c is a function of N and T_{RH} . Using $N \sim 10$ and $T_{\text{RH}} \sim 10^{11} \text{ GeV}$, we find $\Lambda_c \sim 10^6 \text{ GeV}$.

Throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) could account for the dark matter if reheating temperature was $\sim 10^{11} \text{ GeV}$.

• String vacua with ~ 10 throats in this range probably not uncommon. But unfortunately, there is a problem ...

LATE-TIME ABUNDANCE OF KK PARTICLES/GLUEBALLS



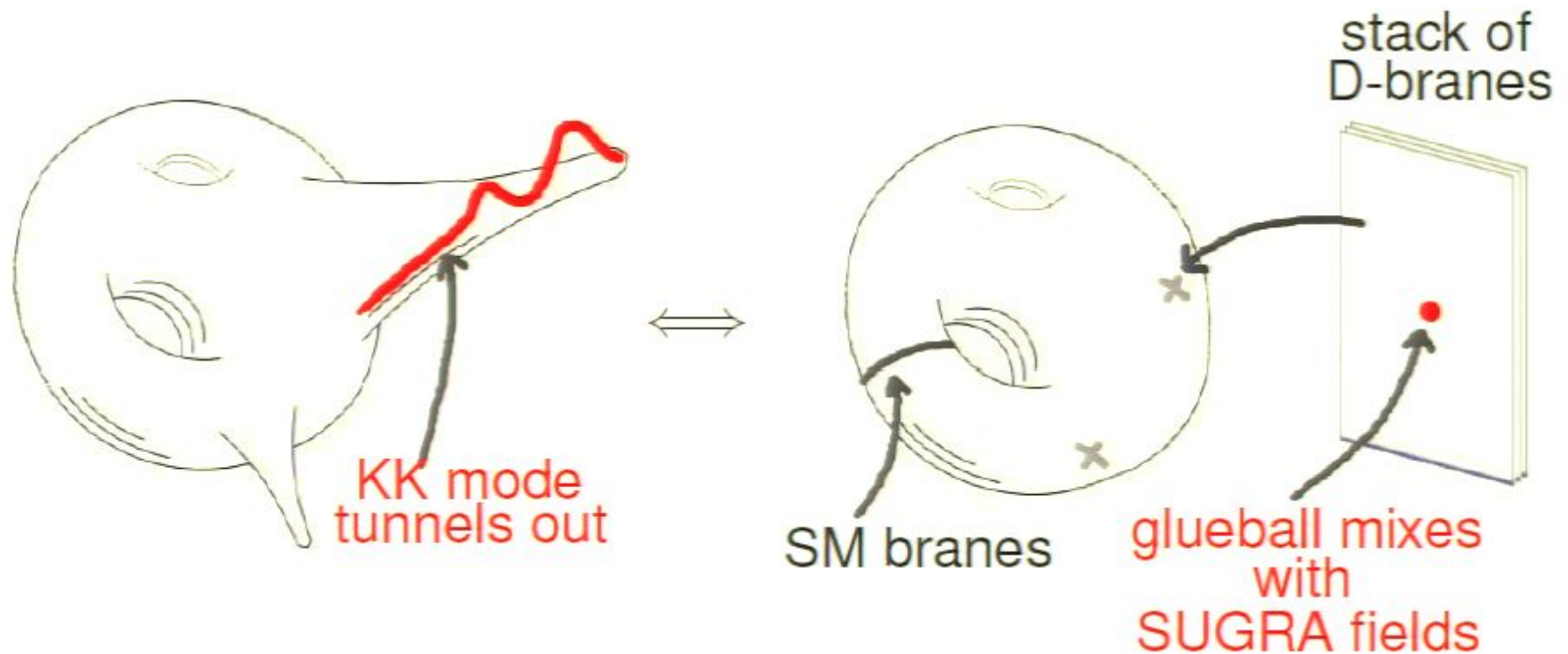
- Λ_c is a function of N and T_{RH} . Using $N \sim 10$ and $T_{\text{RH}} \sim 10^{11}$ GeV, we find $\Lambda_c \sim 10^6$ GeV.

Throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) could account for the dark matter if reheating temperature was $\sim 10^{11}$ GeV.

- String vacua with ~ 10 throats in this range probably not uncommon. **But unfortunately, there is a problem ...**

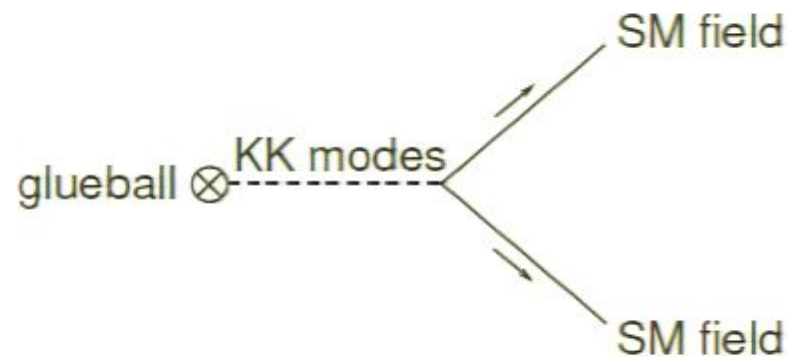
DECAYING DARK MATTER

- Glueballs mix with SUGRA fields in compact space which in turn couple to SM sector. \implies **Dark matter can decay to SM fields.**



DECAY OF SCALAR GLUEBALLS

- **Glueballs decay to SM** via KK modes of SUGRA fields in compact space:



- For $A \sim L$ and $\Lambda < L^{-1}$: Dominant contribution from zero mode.
- E.g. Spin-2-glueball mixing vertex $\propto N/M_4$ (from calculation in throat picture).
- Coupling to SM is $\propto M_4^{-1}$. **Decay rate** then is

$$\Gamma_{\text{glueball}} \sim g N^2 \frac{\Lambda^5}{M_4^4}.$$

OBSERVABLE SIGNATURES

- **Lifetime of glueballs** (for $N \sim 10$) is:

$$\tau_{\text{glueball}} \sim 10^{27} \left(\frac{7 \cdot 10^3 \text{ GeV}}{\Lambda} \right)^5 \text{ s.}$$

- Only for $\Lambda \lesssim 10^6 \text{ GeV}$ are the Spin-2-glueballs stable until our epoch ($\sim 10^{17} \text{ s}$).
- But glueballs decay to photons, hadrons, etc. \Rightarrow Need $\tau_{\text{glueball}} \gtrsim 10^{27} \text{ s}$ to avoid bounds from γ -ray observations.
- Thus, throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ may not give good dark matter candidates.
- On the other hand: Spin-2- (and other) glueballs probably first decay to a lightest glueball in throat sector. Maybe this state is sufficiently stable?

OBSERVABLE SIGNATURES

- **Lifetime of glueballs** (for $N \sim 10$) is:

$$\tau_{\text{glueball}} \sim 10^{27} \left(\frac{7 \cdot 10^3 \text{ GeV}}{\Lambda} \right)^5 \text{ s.}$$

- Only for $\Lambda \lesssim 10^6 \text{ GeV}$ are the Spin-2-glueballs stable until our epoch ($\sim 10^{17} \text{ s}$).
 - But glueballs decay to photons, hadrons, etc. \Rightarrow Need $\tau_{\text{glueball}} \gtrsim 10^{27} \text{ s}$ to avoid bounds from γ -ray observations.
 - Thus, throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ may not give good dark matter candidates.
 - On the other hand: Spin-2- (and other) glueballs probably first decay to a lightest glueball in throat sector. Maybe this state is sufficiently stable?

OBSERVABLE SIGNATURES

- **Lifetime of glueballs** (for $N \sim 10$) is:

$$\tau_{\text{glueball}} \sim 10^{27} \left(\frac{7 \cdot 10^3 \text{ GeV}}{\Lambda} \right)^5 \text{ s.}$$

- Only for $\Lambda \lesssim 10^6 \text{ GeV}$ are the Spin-2-glueballs stable until our epoch ($\sim 10^{17} \text{ s}$).
- But glueballs decay to photons, hadrons, etc. \Rightarrow Need $\tau_{\text{glueball}} \gtrsim 10^{27} \text{ s}$ to avoid bounds from γ -ray observations.
- Thus, throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ may not give good dark matter candidates.
- On the other hand: Spin-2- (and other) glueballs probably first decay to a lightest glueball in throat sector. Maybe this state is sufficiently stable?

OBSERVABLE SIGNATURES

- **Lifetime of glueballs** (for $N \sim 10$) is:

$$\tau_{\text{glueball}} \sim 10^{27} \left(\frac{7 \cdot 10^3 \text{ GeV}}{\Lambda} \right)^5 \text{ s.}$$

- Only for $\Lambda \lesssim 10^6 \text{ GeV}$ are the Spin-2-glueballs stable until our epoch ($\sim 10^{17} \text{ s}$).
- But glueballs decay to photons, hadrons, etc. \Rightarrow Need $\tau_{\text{glueball}} \gtrsim 10^{27} \text{ s}$ to avoid bounds from γ -ray observations.
- Thus, throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ may not give good dark matter candidates.
- On the other hand: Spin-2- (and other) glueballs probably first decay to a lightest glueball in throat sector. Maybe this state is sufficiently stable?

OBSERVABLE SIGNATURES

- Lifetime of glueballs (for $N \sim 10$) is:

$$\tau_{\text{glueball}} \sim 10^{27} \left(\frac{7 \cdot 10^3 \text{ GeV}}{\Lambda} \right)^5 \text{ s.}$$

- Only for $\Lambda \lesssim 10^6 \text{ GeV}$ are the Spin-2-glueballs stable until our epoch ($\sim 10^{17} \text{ s}$).
- But glueballs decay to photons, hadrons, etc. \Rightarrow Need $\tau_{\text{glueball}} \gtrsim 10^{27} \text{ s}$ to avoid bounds from γ -ray observations.
- Thus, throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ may not give good dark matter candidates.
- On the other hand: Spin-2- (and other) glueballs probably first decay to a lightest glueball in throat sector. **Maybe this state is sufficiently stable?**

CONCLUSIONS AND OUTLOOK

- We have calculated the **heat transfer rate** and glueball/KK particle **decay rate** between throats and between a throat and the SM.
- These rates are important in the early-universe cosmology of multi-throat compactifications (e.g. reheating after brane inflation in a throat).
- We have presented another application; KK modes in a throat/glueballs of the dual gauge theory are an interesting dark matter candidate.
- Thermal production of glueballs can be efficient if the glueballs become nonrelativistic quickly after reheating.
- Glueballs from throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) have the abundance of dark matter for $T_{RH} \sim 10^{11} \text{ GeV}$. But stability still an open issue!

CONCLUSIONS AND OUTLOOK

- We have calculated the **heat transfer rate** and glueball/KK particle **decay rate** between throats and between a throat and the SM.
- These rates are important in the early-universe cosmology of multi-throat compactifications (e.g. reheating after brane inflation in a throat).
- We have presented another application: KK modes in a throat/glueballs of the dual gauge theory are an interesting dark matter candidate.
- Thermal production of glueballs can be efficient if the glueballs become nonrelativistic quickly after reheating.
- Glueballs from throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) have the abundance of dark matter for $T_{RH} \sim 10^{11} \text{ GeV}$. But stability still an open issue!

CONCLUSIONS AND OUTLOOK

- We have calculated the **heat transfer rate** and glueball/KK particle **decay rate** between throats and between a throat and the SM.
- These rates are important in the early-universe cosmology of multi-throat compactifications (e.g. reheating after brane inflation in a throat).
- We have presented another application: KK modes in a throat/glueballs of the dual gauge theory are an **interesting dark matter candidate**.
- Thermal production of glueballs can be efficient if the glueballs become nonrelativistic quickly after reheating.
- Glueballs from throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) have the abundance of dark matter for $T_{\text{RH}} \sim 10^{11} \text{ GeV}$. But stability still an open issue!

CONCLUSIONS AND OUTLOOK

- We have calculated the **heat transfer rate** and glueball/KK particle **decay rate** between throats and between a throat and the SM.
- These rates are important in the early-universe cosmology of multi-throat compactifications (e.g. reheating after brane inflation in a throat).
- We have presented another application: KK modes in a throat/glueballs of the dual gauge theory are an **interesting dark matter candidate**.
- Thermal production of glueballs can be efficient if the glueballs become nonrelativistic quickly after reheating.
- Glueballs from throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) have the abundance of dark matter for $T_{\text{RH}} \sim 10^{11} \text{ GeV}$. But stability still an open issue!

CONCLUSIONS AND OUTLOOK

- We have calculated the **heat transfer rate** and glueball/KK particle **decay rate** between throats and between a throat and the SM.
- These rates are important in the early-universe cosmology of multi-throat compactifications (e.g. reheating after brane inflation in a throat).
- We have presented another application: KK modes in a throat/glueballs of the dual gauge theory are an **interesting dark matter candidate**.
- Thermal production of glueballs can be efficient if the glueballs become nonrelativistic quickly after reheating.
- Glueballs from throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) have the abundance of dark matter for $T_{RH} \sim 10^{11} \text{ GeV}$. **But stability still an open issue!**

CONCLUSIONS AND OUTLOOK

- We have calculated the **heat transfer rate** and glueball/KK particle **decay rate** between throats and between a throat and the SM.
- These rates are important in the early-universe cosmology of multi-throat compactifications (e.g. reheating after brane inflation in a throat).
- We have presented another application: KK modes in a throat/glueballs of the dual gauge theory are an **interesting dark matter candidate**.
- Thermal production of glueballs can be efficient if the glueballs become nonrelativistic quickly after reheating.
- Glueballs from throats with $10^6 \text{ GeV} < \Lambda < 10^{11} \text{ GeV}$ (and $N \sim 10$) have the abundance of dark matter for $T_{RH} \sim 10^{11} \text{ GeV}$. **But stability still an open issue!**