Title: Quantum foundations minus probability theory

Date: May 05, 2009 04:00 PM

URL: http://pirsa.org/09050010

Abstract: Researchers in quantum foundations claim (D'Ariano, Fuchs, ...):

Quantum = probability theory + x

and hence:

x = Quantum - probability theory

Guided by the metaphorical analogy:

probability theory / x = flesh / bones

we introduce a notion of quantum measurement within x, which, when flesing it with Hilbert spaces, provides orthodox quantum mechanical probability calculus.

Pirsa: 09050010 Page 1/108

How much QM can we recover without a priori assuming instrumentalist concepts such as measurement and probability but just 'compoundness'?

Pirsa: 09050010 Page 2/108

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and hence:

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Guided by the metaphorical analogy:

$$\frac{\text{flesh}}{\text{bones}} = \frac{\text{probability}}{x}$$

We start from x, (i) extract (the skeleton of) classicality from it, (ii) extract probability from it, (iii) extract quantum measurements from it, and (iv) describe quantum-classical interaction within it. Flesing it with FHilb yields QM probabilistic calculus.

Pirsa: 09050010 Page 3/108

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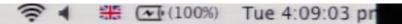
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Is there a conceptual picture behind this game?

### Classicizing quantumness

arXiv:0904.1997

Somewhere out there, in ontic reality, is a world.

That world is called the 'quantum universe'.

We would like to probe that world.

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Pirsa: 09050010 Page 6/108

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We refer to 'identifiable parts' of it as systems, and to their 'identifiable changes' as processes. To joint parts and processes we refer by  $- \otimes -$ , and to consecutive processes by  $- \circ -$ .

Pirsa: 09050010 Page 7/108

One could either think of this as structure of the interface or as structure of the system.

Our language := system, process,  $\otimes$ ,  $\circ$ 

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Pirsa: 09050010 Page 8/108

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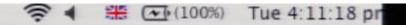
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Pirsa: 09050010 Page 9/108

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Pirsa: 09050010 Page 10/108

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What does such a rigid stance buy us?

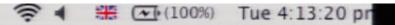
Pirsa: 09050010 Page 11/108

# Why does a tiger have stripes and a lion doesn't?





Pirsa: 09050010 Page 12/108



### Why does a tiger have stripes and a lion doesn't?





prey ⊗ predator ⊗ environment

hunt

dead prey ⊗ eating predator

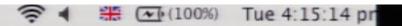
Pirsa: 09050010 Page 13/108

Processes and their compositions:

$$f \equiv \frac{1}{f}$$
  $1_A \equiv \left| \begin{array}{cc} g \circ f \end{array} \right| = \frac{1}{g}$   $f \otimes g \equiv \frac{1}{f}$ 

States and effects for systems, and resulting numbers:

$$I \xrightarrow{\psi} A \equiv \psi \qquad A \xrightarrow{\pi} I \equiv \Phi \qquad I \xrightarrow{\pi} I \equiv \Phi$$



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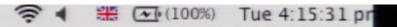
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Upside-down flipping i.e. adjoints:

pics:	1	1
cats:	$A \xrightarrow{f} B$	$B \xrightarrow{f^{\dagger}} A$
FHilb:	linear map	its adjoint



# **CLASSICAL INTERFACES**

Pirsa: 09050010 Page 16/108

#### An classical interface is:

$$A \xrightarrow{\delta} A \otimes A = \bigvee$$

$$A \xrightarrow{\varepsilon} I = \P$$

#### such that:

- 1.  $\varepsilon$  is a *unit* for  $\delta$ ;
- 2.  $\delta$  is coassociative;
- 3.  $\delta$  is cocommutative;
- 4.  $\delta$  is isometry;
- 5.  $\delta$  is Frobenius.

In FHilb we can encode all ONBs as

$$\delta: \mathcal{H} \to \mathcal{H} \otimes \mathcal{H} :: |i\rangle \mapsto |ii\rangle \qquad \varepsilon: \mathcal{H} \to \mathbb{C} :: |i\rangle \mapsto 1.$$

**Thm.** Interfaces  $(\delta, \varepsilon)$  in **FHilb** exactly correspond with orthonormal bases on the underlying Hilbert space.

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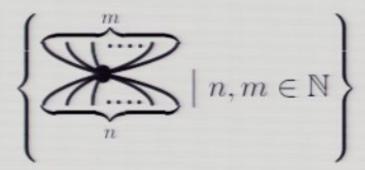
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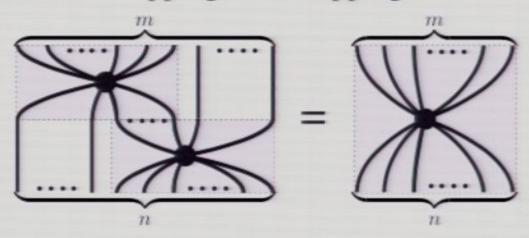
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Pirsa: 09050010 Page 20/108

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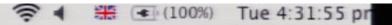
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States and effects for systems, and resulting numbers:

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### Why does a tiger have stripes and a lion doesn't?



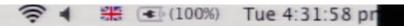


prey ⊗ predator ⊗ environment

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Pirsa: 09050010 Page 29/108



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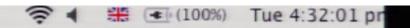
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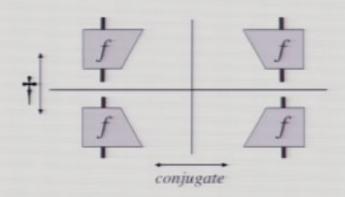
Upside-down flipping i.e. adjoints:

pics:	1	1
cats:	$A \xrightarrow{f} B$	$B \xrightarrow{f^{\dagger}} A$
FHilb:	linear map	its adjoint



pics:	1	<b>=</b>	
cats:	$A \xrightarrow{f} B$	$f^{\sharp} = A \xrightarrow{(e_A \otimes 1_B) \circ (1_A \otimes f^{\dagger} \otimes 1_B) \circ (1_A \otimes \eta_B)} B$	
FHilb:	linear map	its conjugate	

pics:	17	
cats:	$A \xrightarrow{f} B$	$f^* = B \xrightarrow{(\epsilon_B \otimes 1_A) \circ (1_B \otimes f \otimes 1_A) \circ (1_B \otimes \eta_A)} A$
FHilb:	Hilb: linear map its transposed	



# MEASUREMENT ↔ NO-SIGNALING

Pirsa: 09050010 Page 32/108

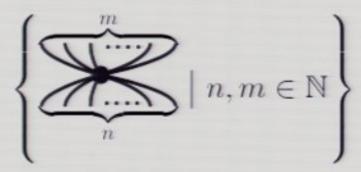
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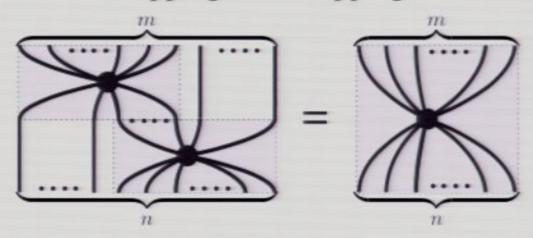
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Pirsa: 09050010 Page 33/108

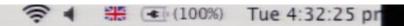
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Pirsa: 09050010 Page 34/108



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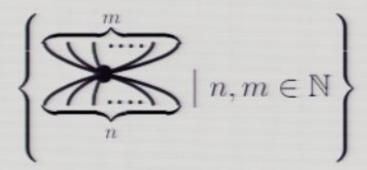
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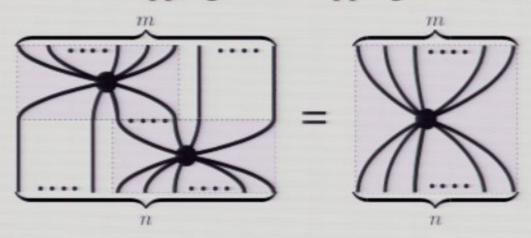
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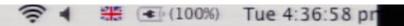
Pirsa: 09050010 Page 37/108

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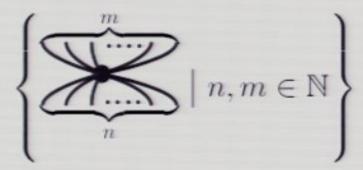
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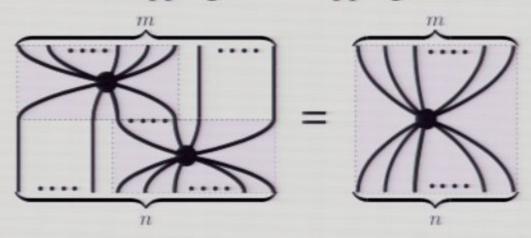
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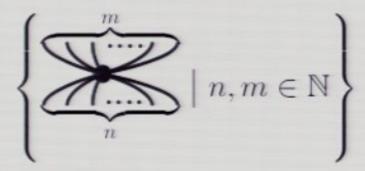


Pirsa: 09050010 Page 41/108

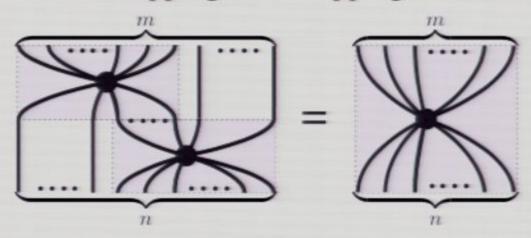
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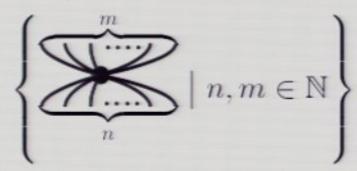
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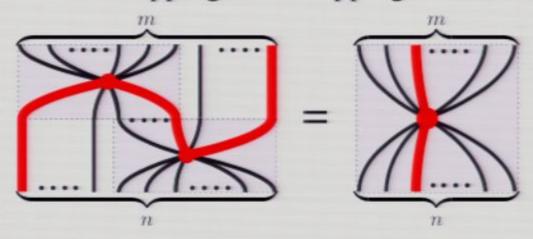
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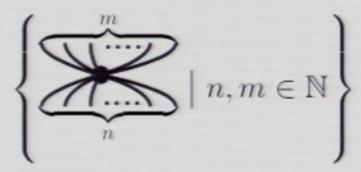
Pirsa: 09050010 Page 43/108



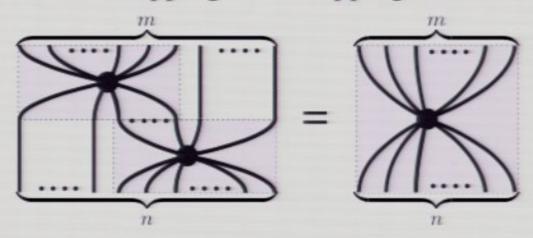
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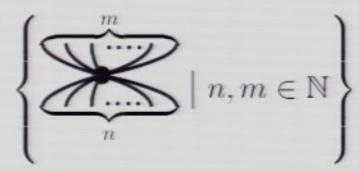
for instances 
$$\delta_0^2 = \bigvee$$
 and  $\delta_2^0 = \bigvee$ .



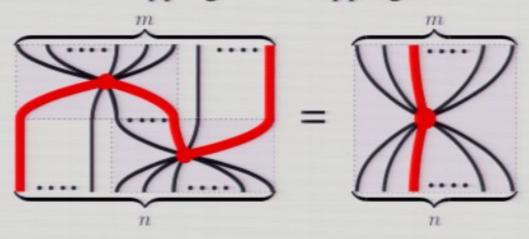
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Pirsa: 09050010 Page 45/108



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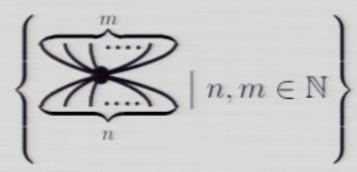
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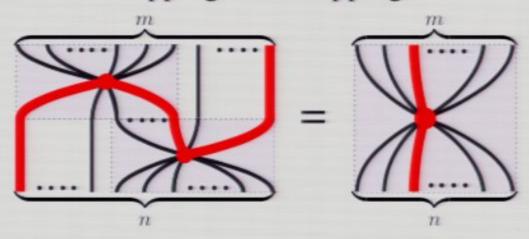
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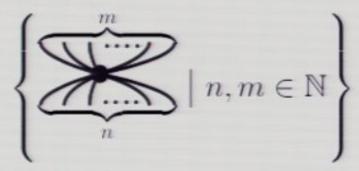
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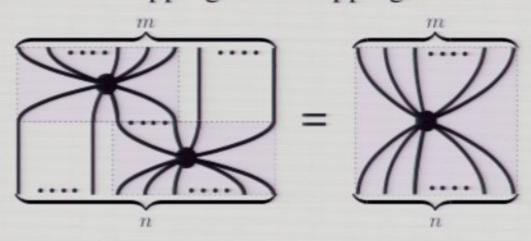
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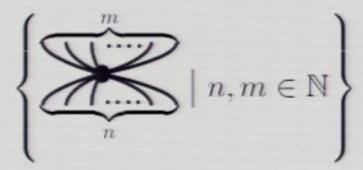
Pirsa: 09050010 Page 49/108

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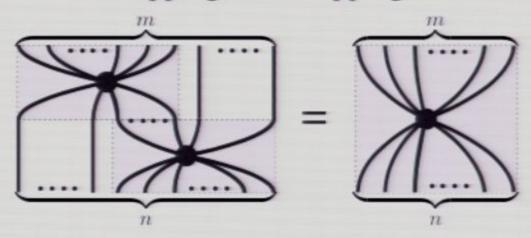
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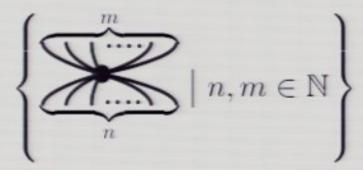
Pirsa: 09050010 Page 51/108

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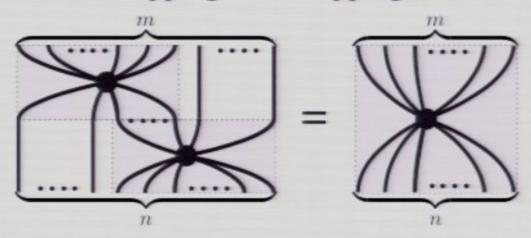
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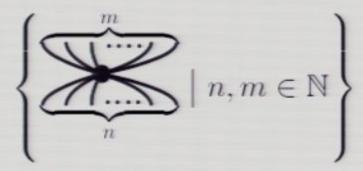
Pirsa: 09050010 Page 53/108

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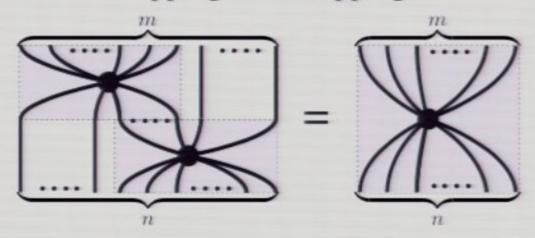
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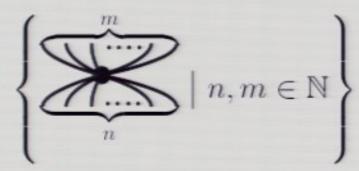
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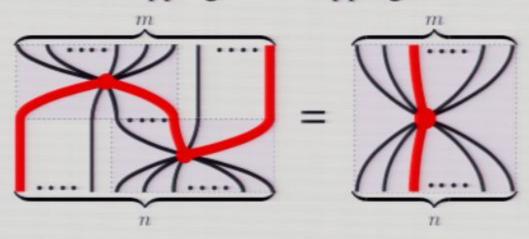
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Pirsa: 09050010 Page 55/108



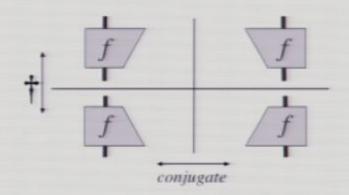
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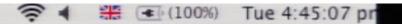
pics:	1	=
cats:	$A \xrightarrow{f} B$	$f^{\sharp} = A \xrightarrow{(\epsilon_A \otimes 1_B) \circ (1_A \otimes f^{\dagger} \otimes 1_B) \circ (1_A \otimes \eta_B)} B$
FHilb:	linear map	its conjugate

pics:	1	
cats:	$A \xrightarrow{f} B$	$f^* = B \xrightarrow{(\epsilon_B \otimes 1_A) \circ (1_B \otimes f \otimes 1_A) \circ (1_B \otimes \eta_A)} A$
FHilb:	linear map	its transposed

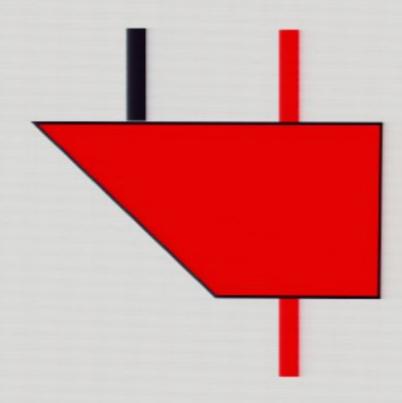


## MEASUREMENT ↔ NO-SIGNALING

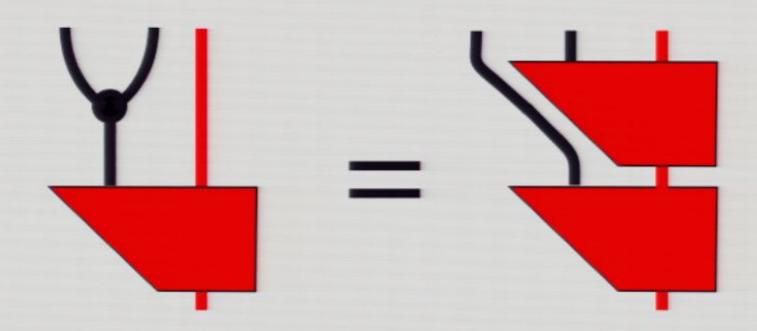
Pirsa: 09050010 Page 58/108



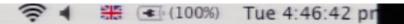
 $\mathcal{M}:A \to X \otimes A$ 



 $\mathcal{M}:A\to X\otimes A$ 

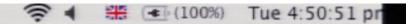


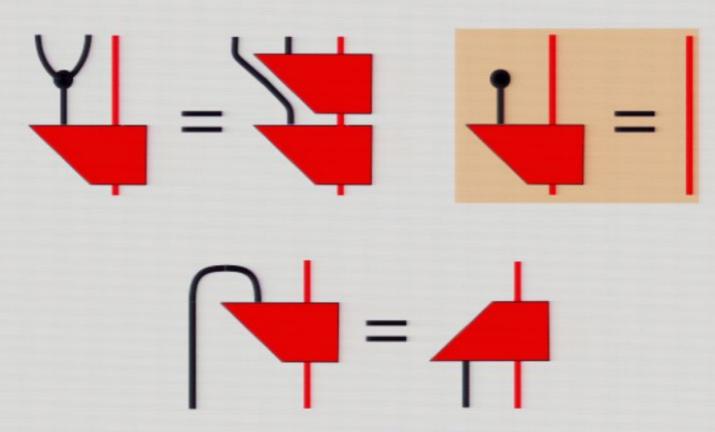
⇒ von Neumann projection postulate.



Minimal requirements for reasonable notion of measurement

Pirsa: 09050010 Page 61/108





Asserts no-signaling

Pirsa: 09050010 Page 62/108

**Thm.** In **FHilb** the above rules yield exactly all projector spectra arising from self-adjoint operators.

#### proof.

Projection postulate ⇒

idempotence

mutual orthogonality

No signaling ⇒

Completeness of spectrum

Minimal requirement II ⇒

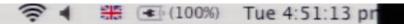
Orthogonality of projectors

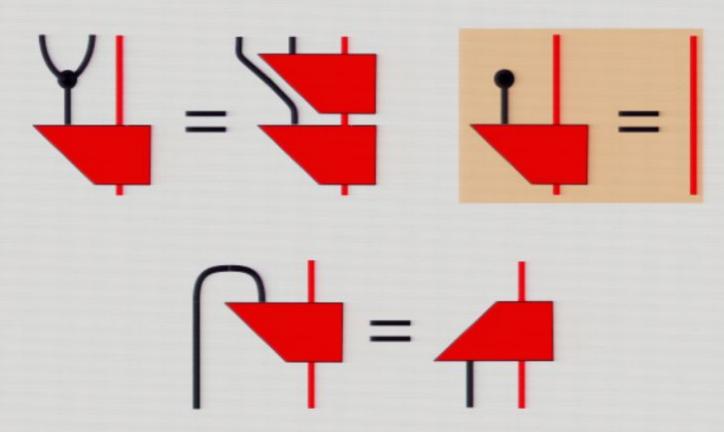
$$P_i^2 = P_i$$

$$P_i \circ P_{j \neq i} = \mathbf{0}$$

$$\sum_{i} P_{i} = 1_{\mathcal{H}}$$

$$\frac{P_i^{\dagger} = P_i}{\text{PROJECTOR}}$$
SPECTRUM





Asserts no-signaling

Pirsa: 09050010 Page 64/108

**Thm.** In **FHilb** the above rules yield exactly all projector spectra arising from self-adjoint operators.

#### proof.

Projection postulate ⇒

idempotence

mutual orthogonality

No signaling ⇒

Completeness of spectrum

Minimal requirement II ⇒

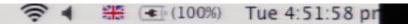
Orthogonality of projectors

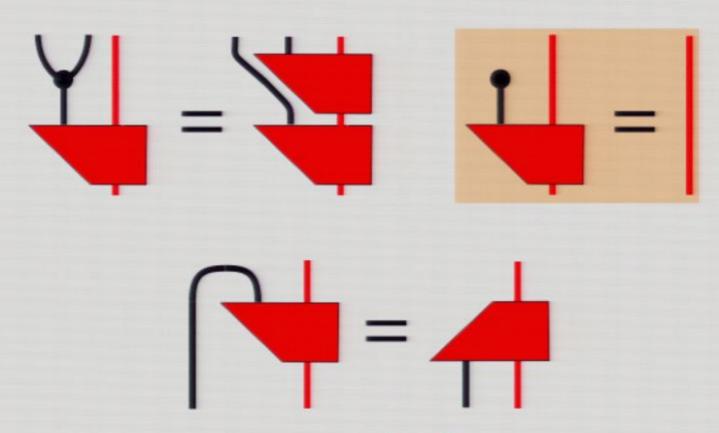
$$P_i^2 = P_i$$

$$P_i \circ P_{j \neq i} = \mathbf{0}$$

$$\sum_{i} P_{i} = 1_{\mathcal{H}}$$

$$\frac{P_i^{\dagger} = P_i}{\text{PROJECTOR}}$$
SPECTRUM





Asserts no-signaling

Pirsa: 09050010 Page 66/108

**Thm.** In **FHilb** the above rules yield exactly all projector spectra arising from self-adjoint operators.

#### proof.

Projection postulate ⇒

idempotence

mutual orthogonality

No signaling ⇒

Completeness of spectrum

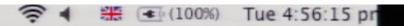
Minimal requirement II ⇒

Orthogonality of projectors

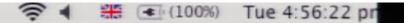
$$P_i^2 = P_i$$
$$P_i \circ P_{j \neq i} = \mathbf{0}$$

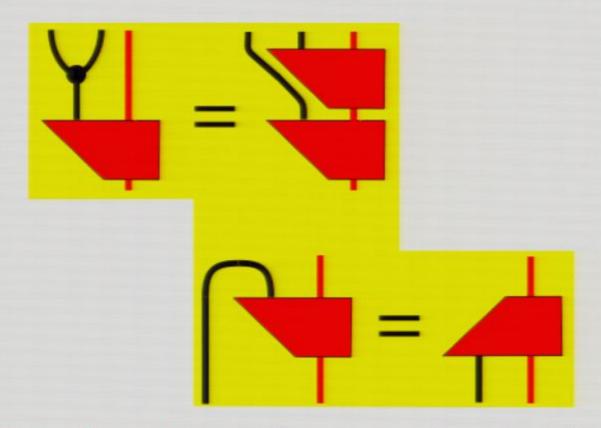
$$\sum_{i} P_{i} = 1_{\mathcal{H}}$$

$$\frac{P_i^{\dagger} = P_i}{\text{PROJECTOR}}$$
SPECTRUM



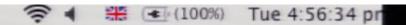
## Canonical observable :=



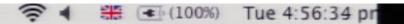


Minimal requirements for reasonable notion of measurement

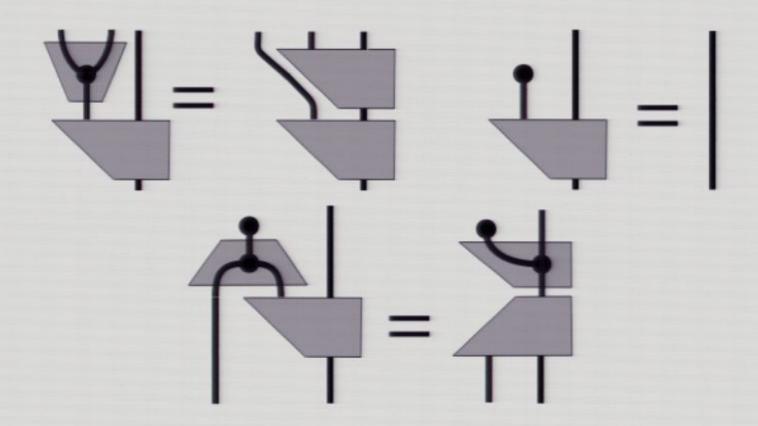
Pirsa: 09050010 Page 69/108

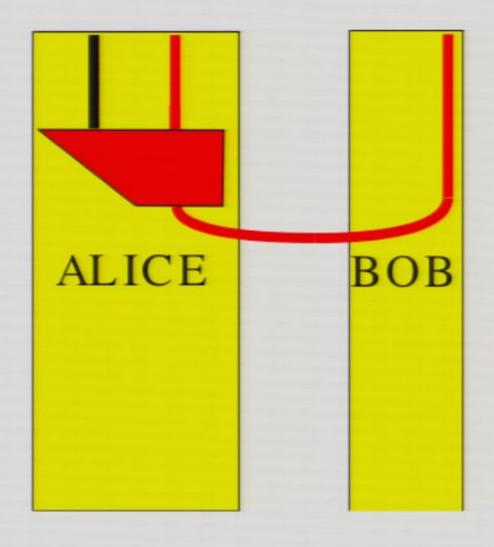


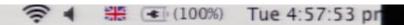
# Canonical observable := Copying



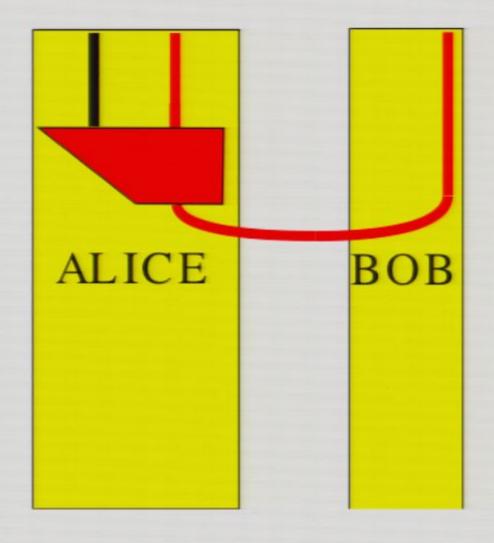
# Canonical observable := Copying

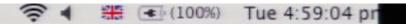




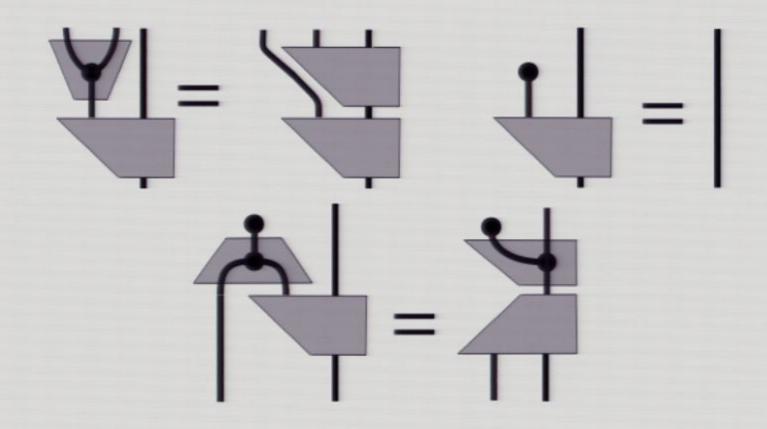


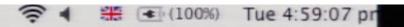
## Canonical observable :=



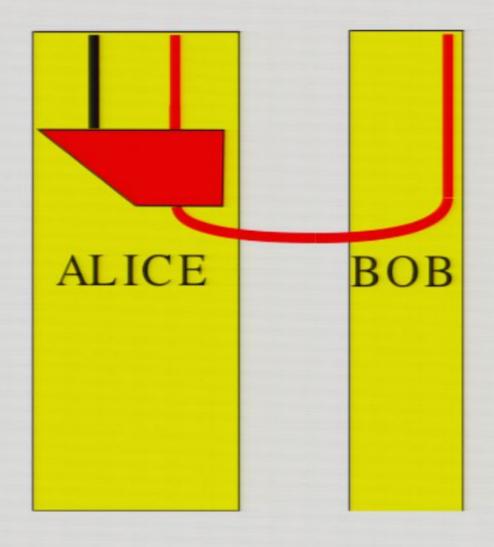


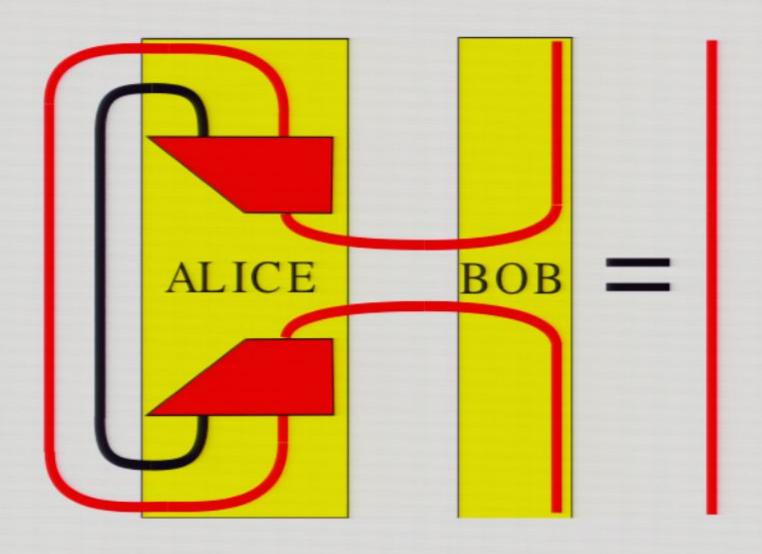
# Canonical observable := Copying

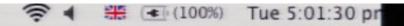


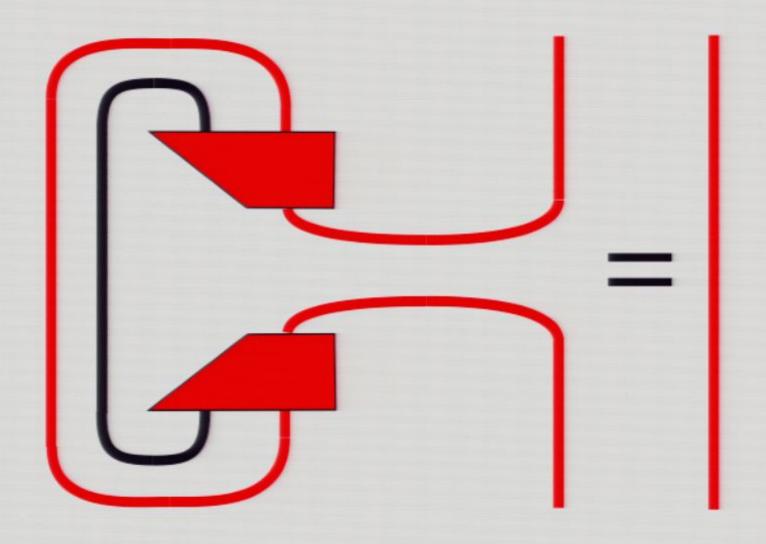


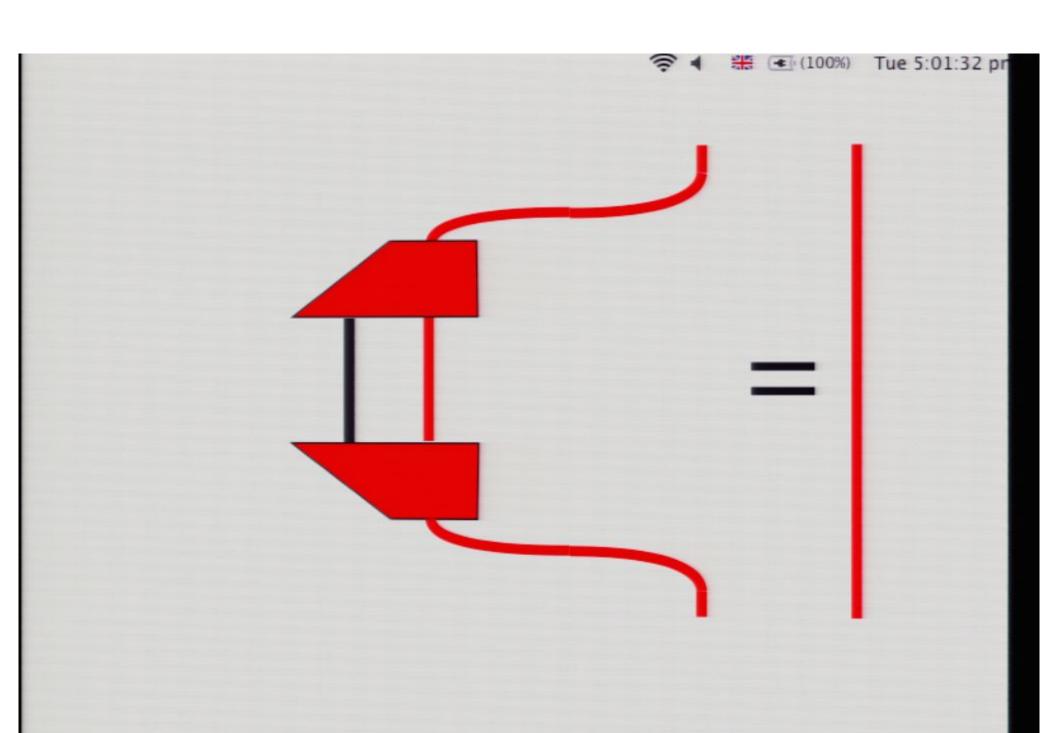
### Canonical observable :=

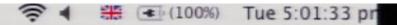


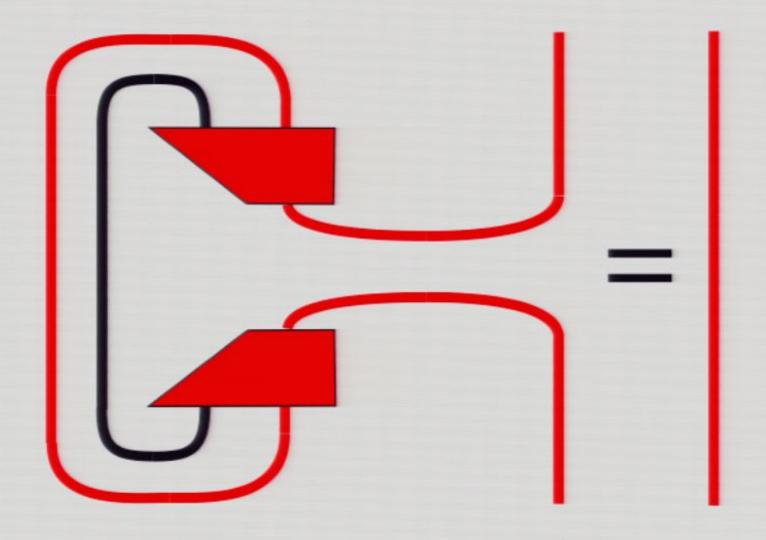


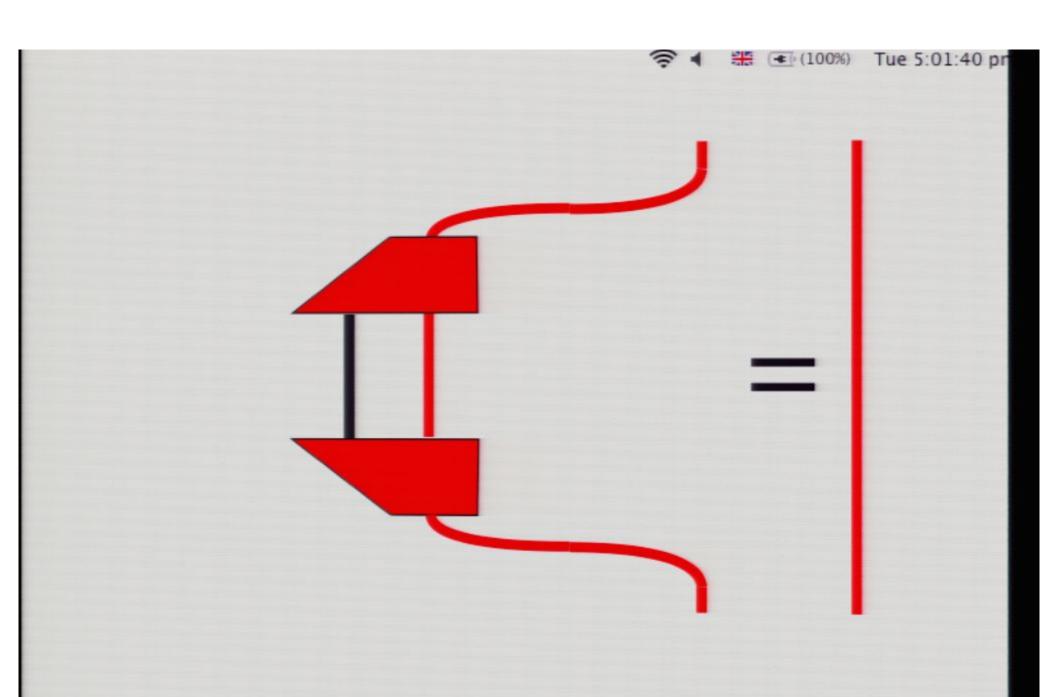


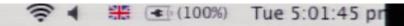


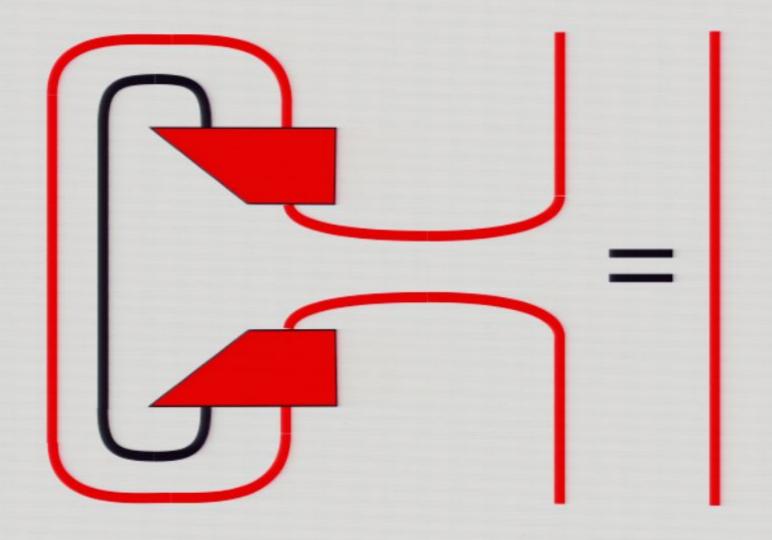


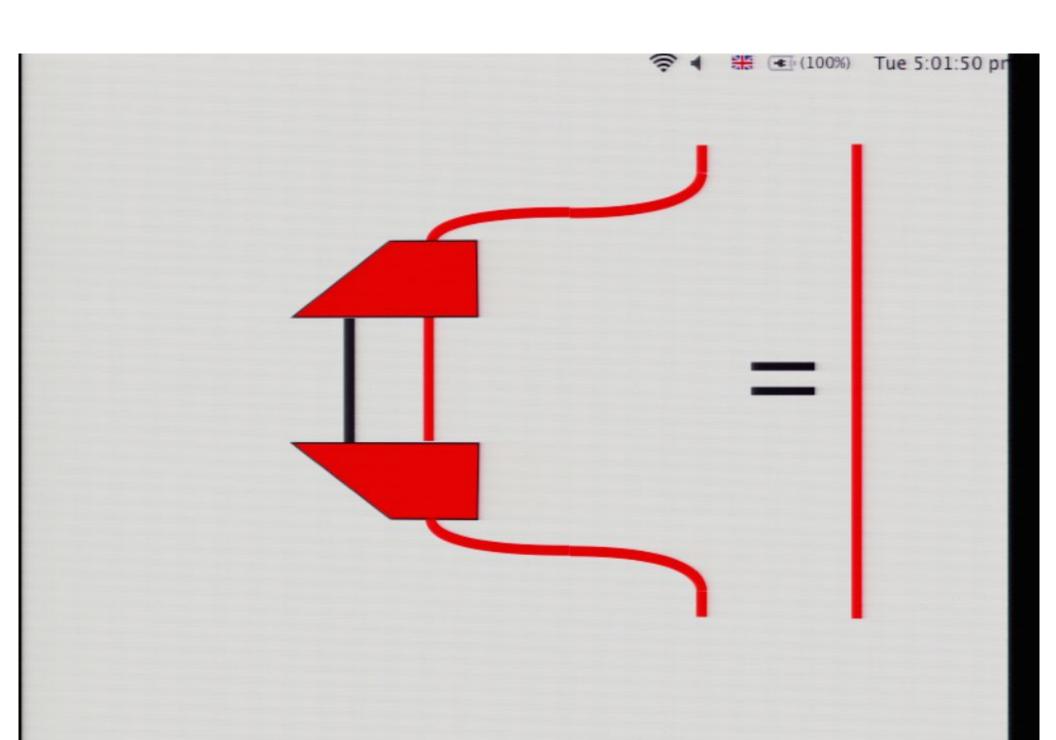


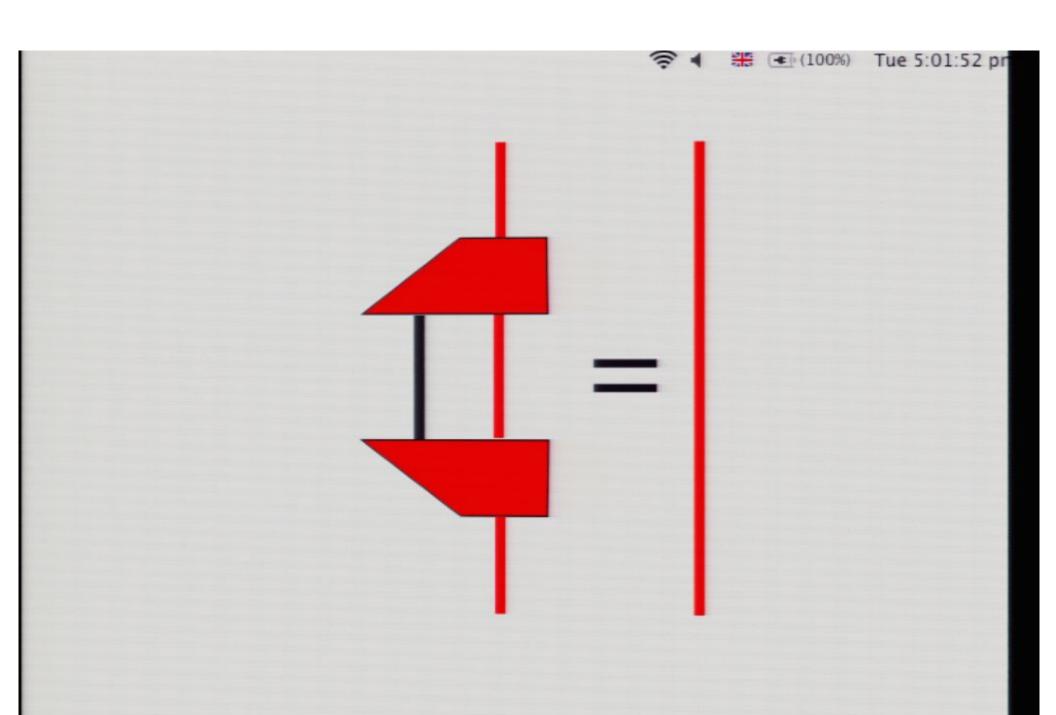


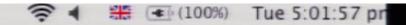


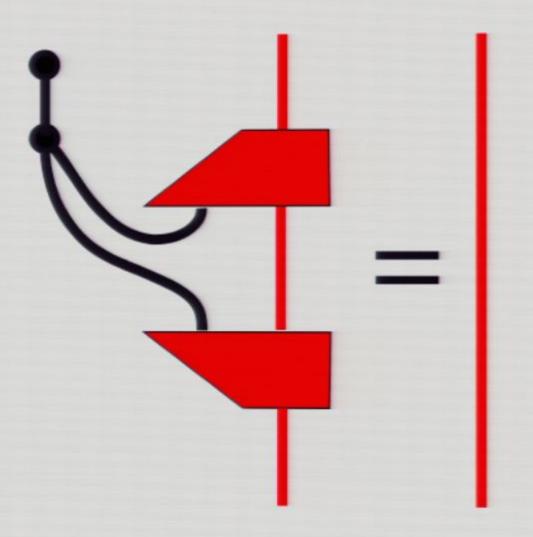


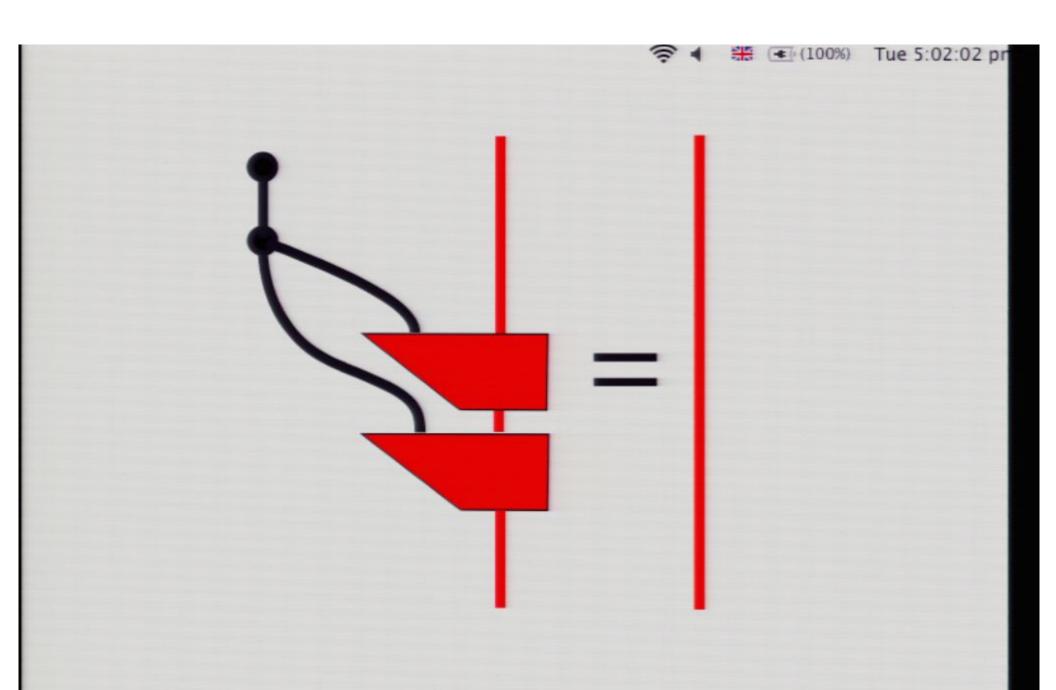


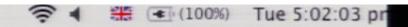


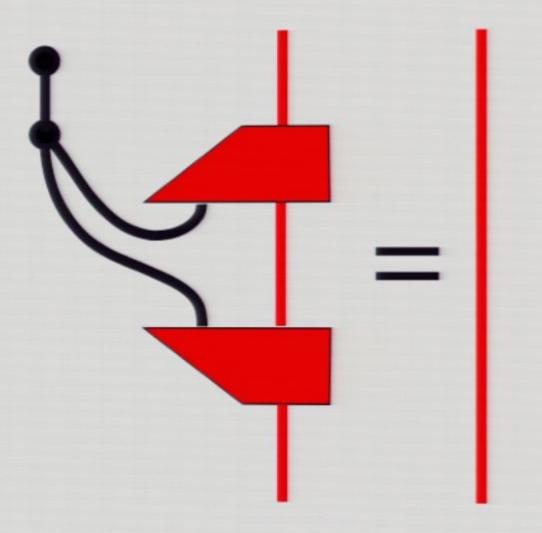


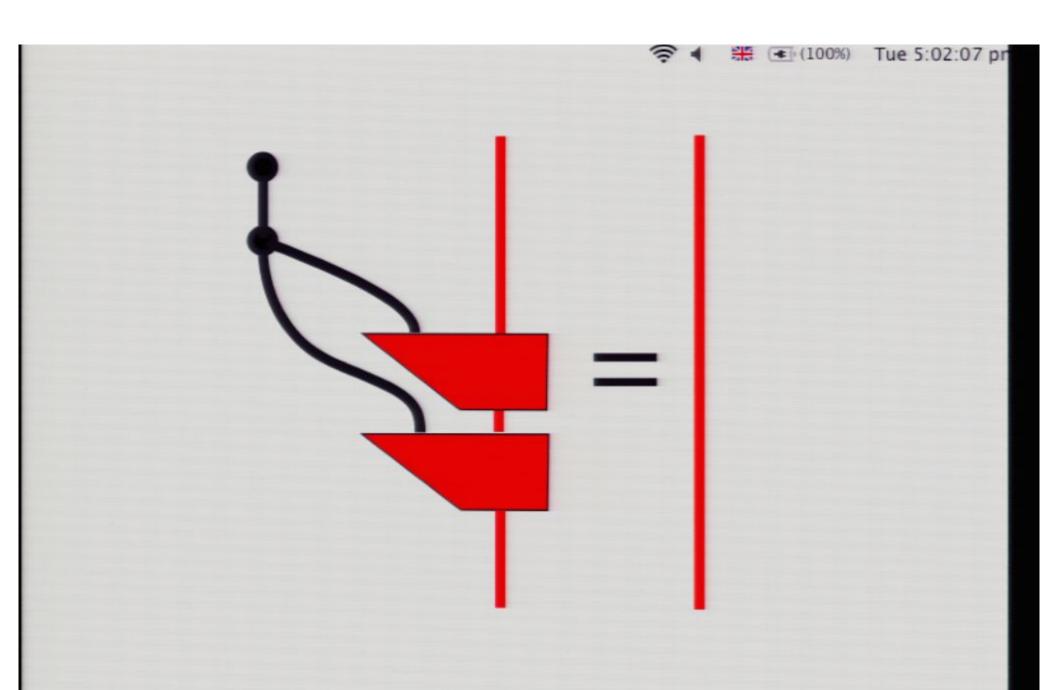


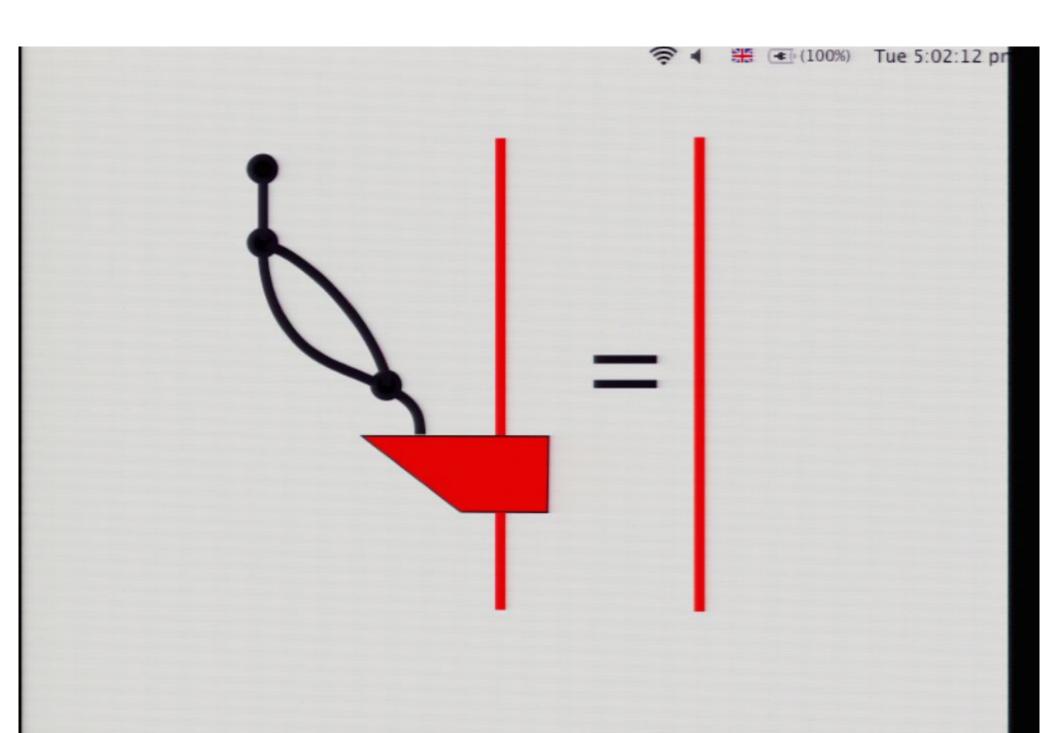


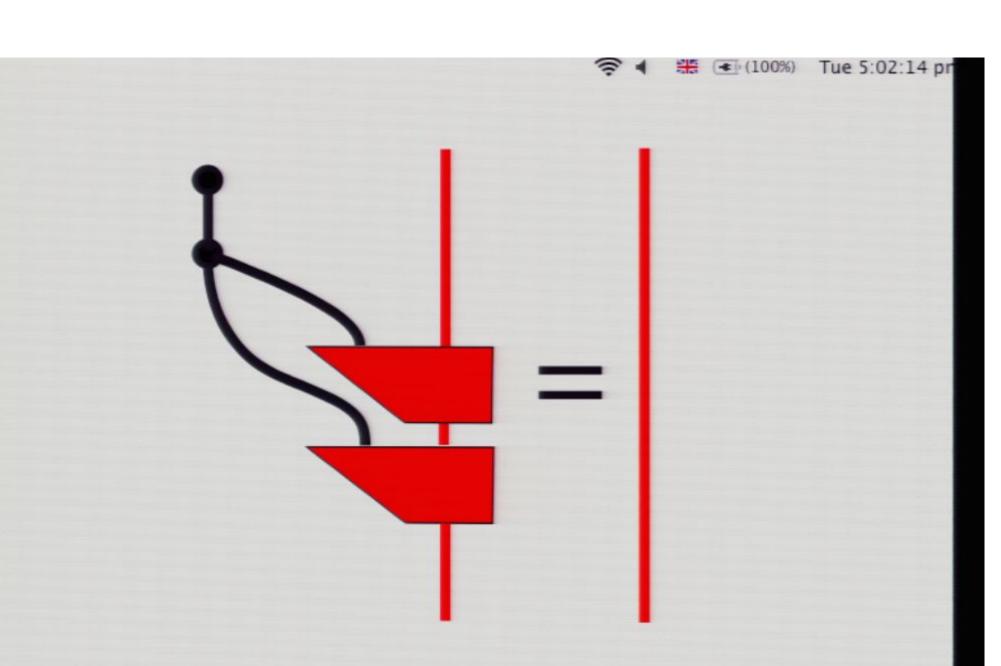


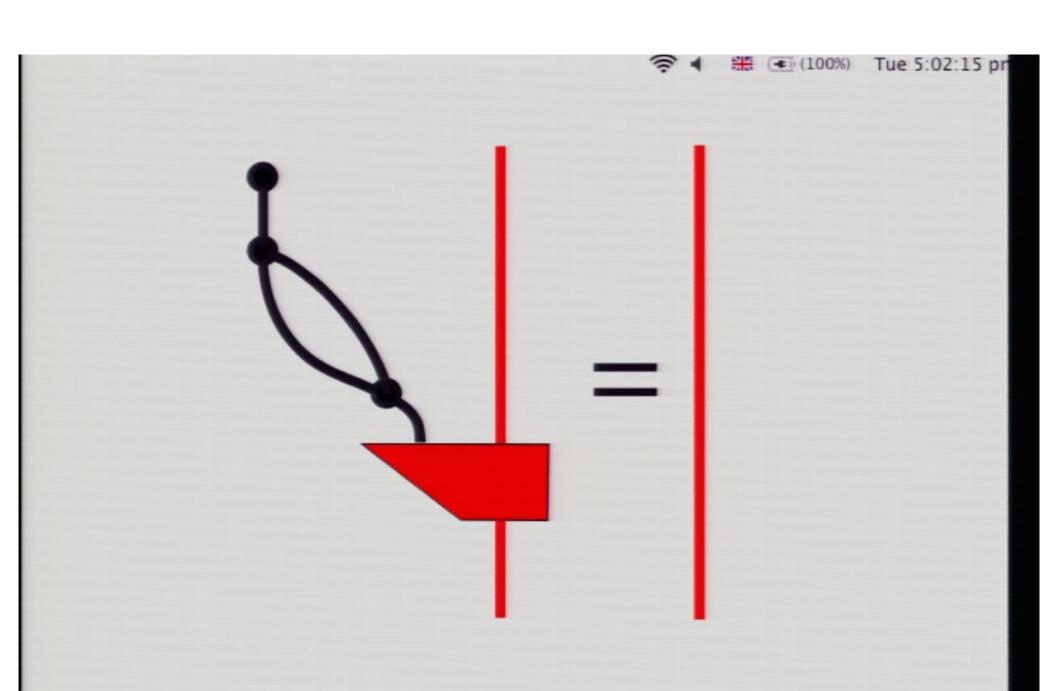


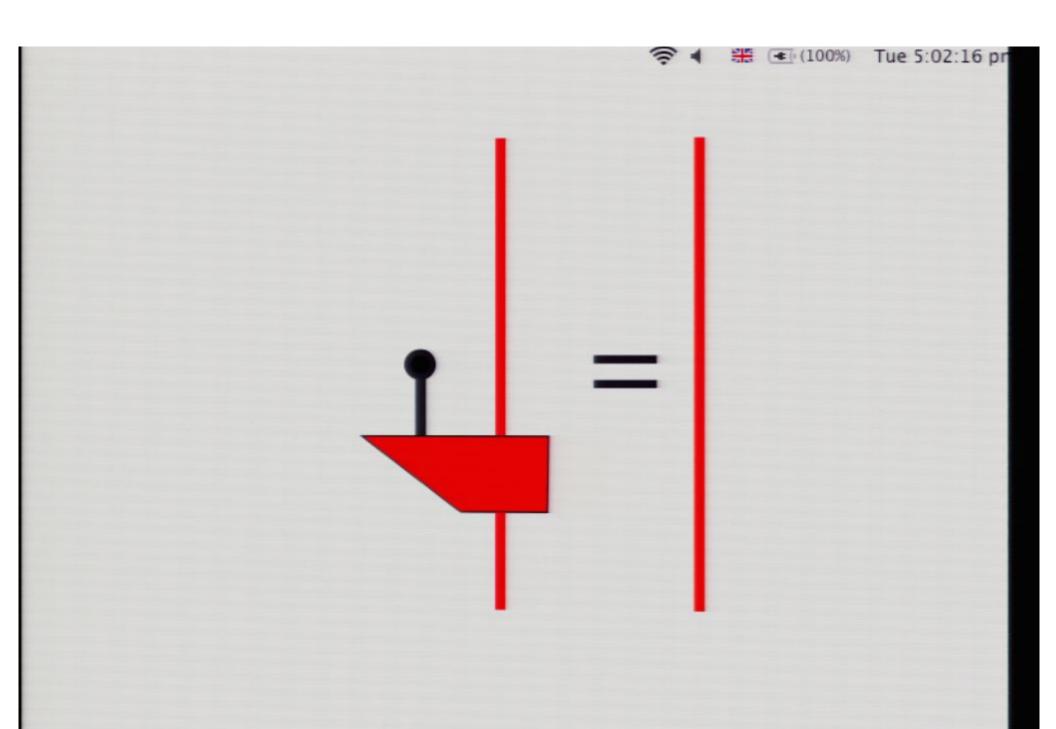


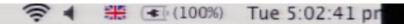




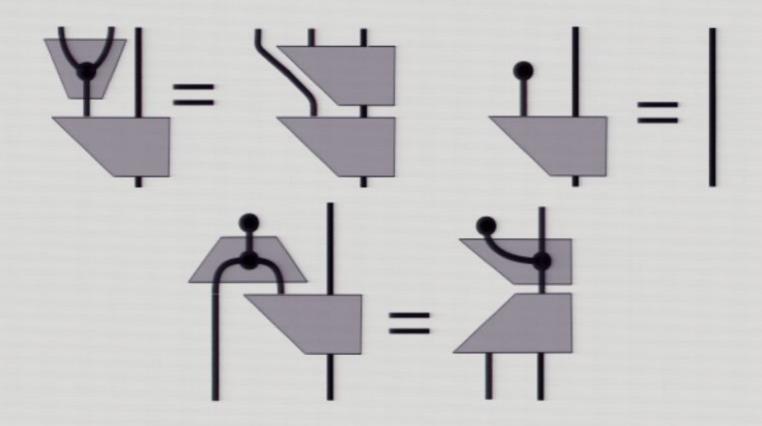


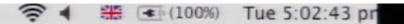




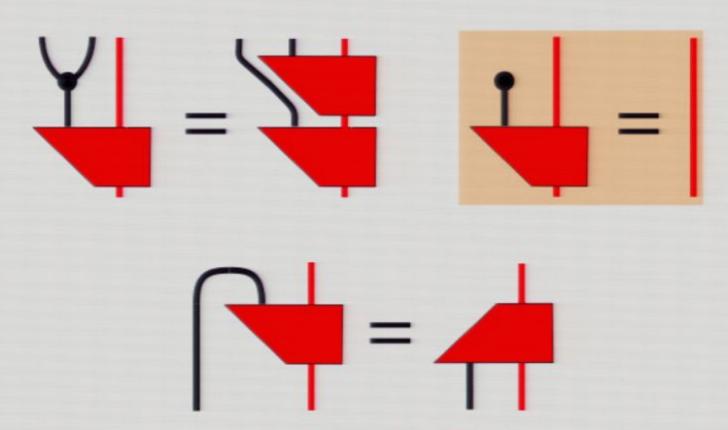


# **Canonical observable := Copying**



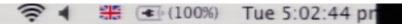


#### **Quantum measurement:**



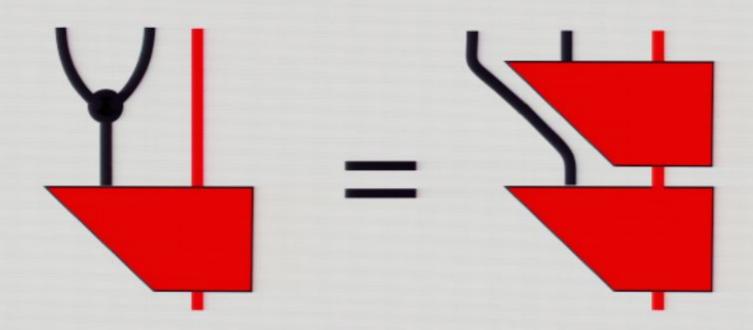
Asserts no-signaling

Pirsa: 09050010 Page 95/108

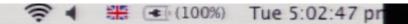


## Quantum measurement:

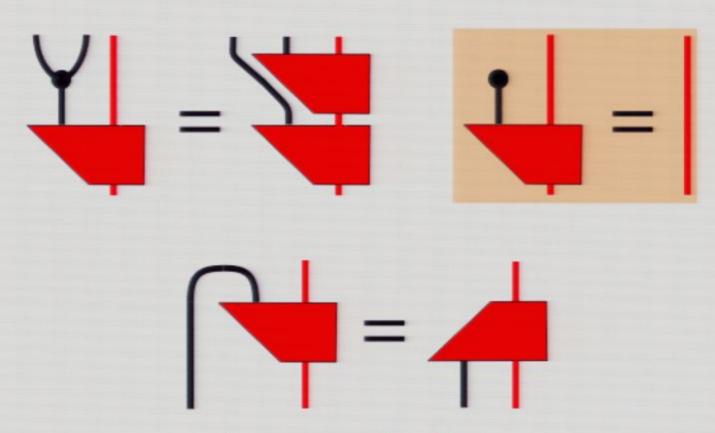
 $\mathcal{M}: A \to X \otimes A$ 



⇒ von Neumann projection postulate.

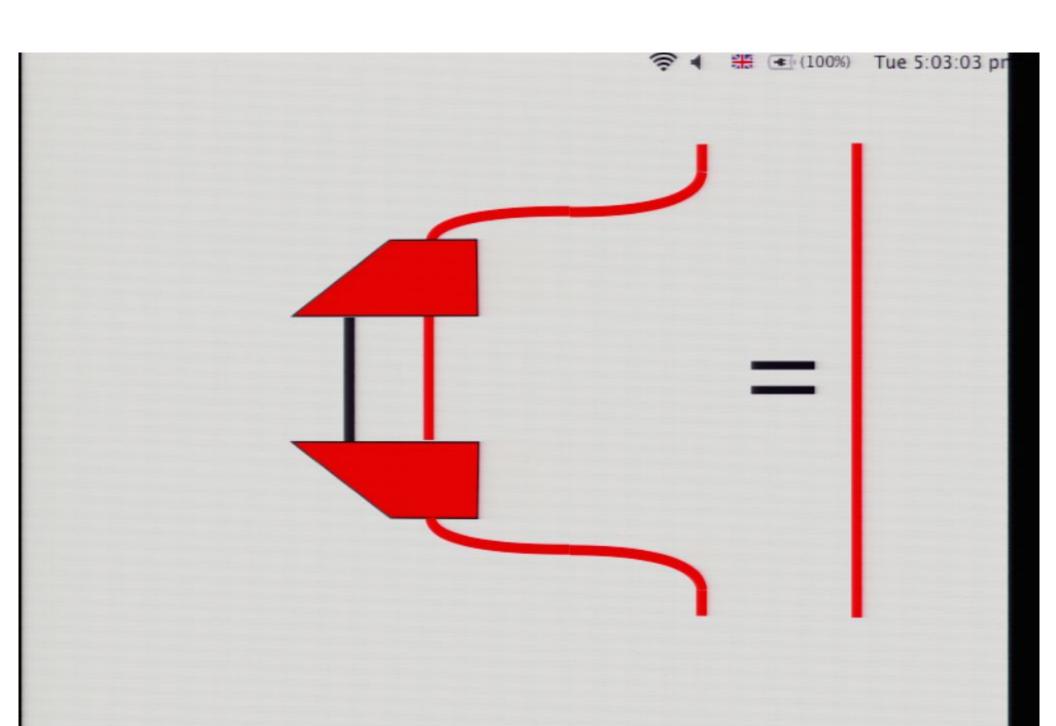


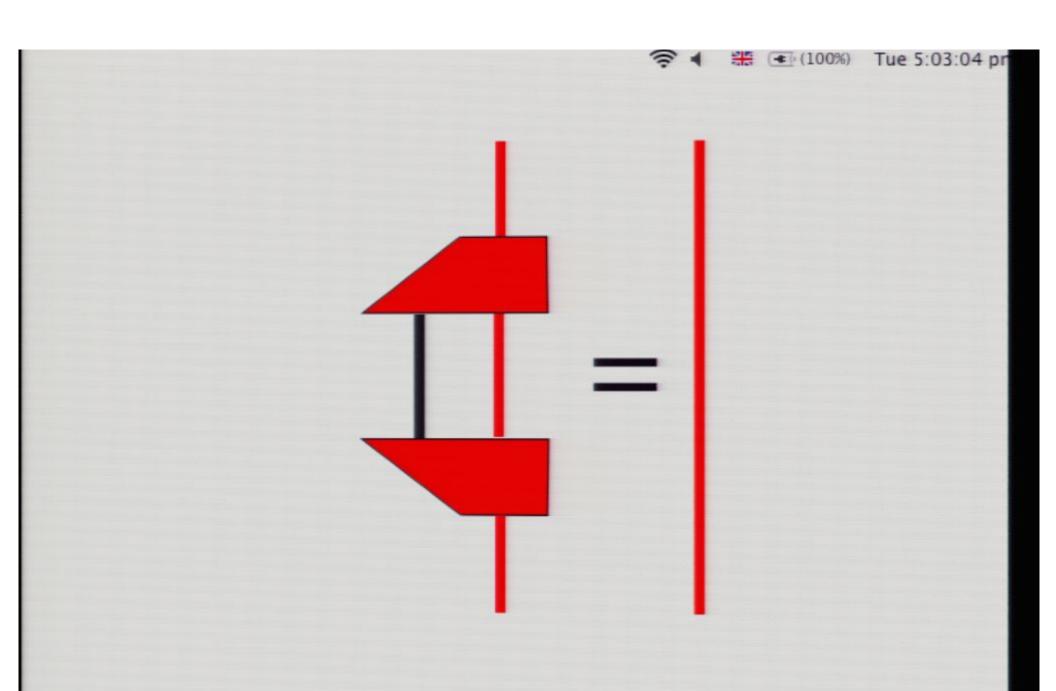
### **Quantum measurement:**

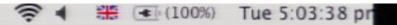


Asserts no-signaling

Pirsa: 09050010 Page 97/108

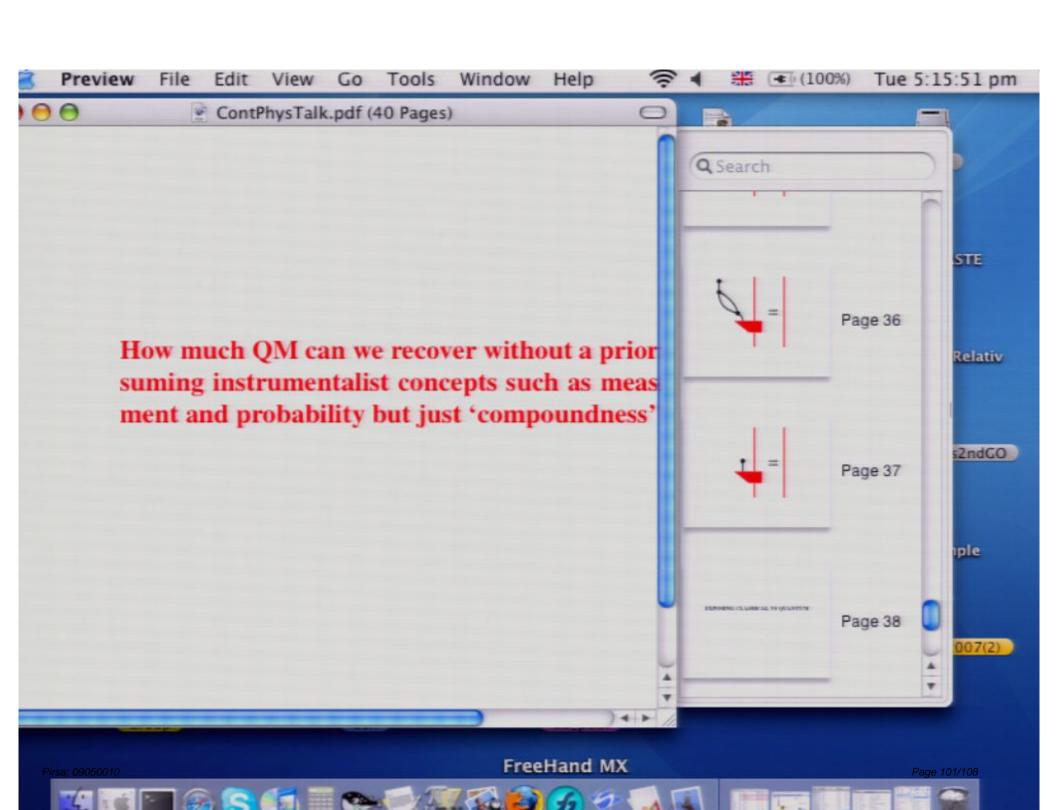






# EXPOSING CLASSICAL VS QUANTUM

Pirsa: 09050010 Page 100/108



An classical interface is:

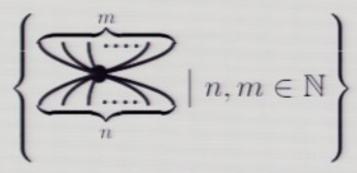
$$A \xrightarrow{\delta} A \otimes A = \bigvee$$

$$A \xrightarrow{\varepsilon} I = \P$$

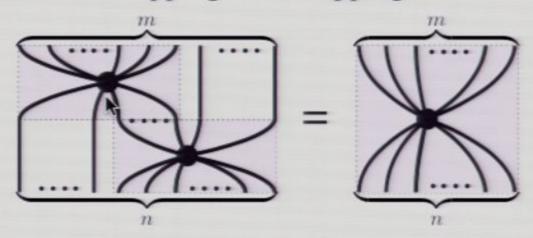
such that:

- 1.  $\varepsilon$  is a *unit* for  $\delta$ ;
- 2.  $\delta$  is coassociative;
- 3.  $\delta$  is cocommutative;
- 4.  $\delta$  is isometry;
- 5.  $\delta$  is *Frobenius*.

#### A classical interface is:



invariant under flipping and swapping, and such that:



An classical interface is:

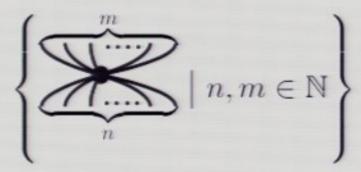
$$A \xrightarrow{\delta} A \otimes A = \bigvee$$

$$A \xrightarrow{\varepsilon} I = \P$$

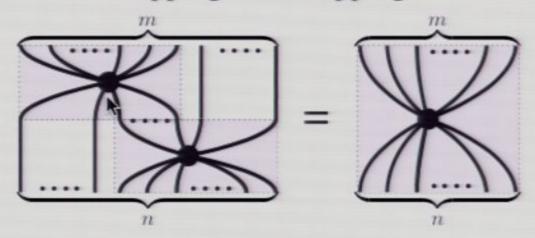
such that:

- 1.  $\varepsilon$  is a *unit* for  $\delta$ ;
- 2.  $\delta$  is coassociative;
- 3.  $\delta$  is cocommutative;
- 4.  $\delta$  is isometry;
- 5.  $\delta$  is Frobenius.

#### A classical interface is:



invariant under flipping and swapping, and such that:



An classical interface is:

$$A \xrightarrow{\delta} A \otimes A = \bigvee$$
$$A \xrightarrow{\varepsilon} I = \uparrow$$

such that:

- 1.  $\varepsilon$  is a *unit* for  $\delta$ ;
- 2.  $\delta$  is coassociative;
- 3.  $\delta$  is cocommutative;
- 4.  $\delta$  is isometry;
- 5.  $\delta$  is Frobenius.

How much QM can we recover without a priori assuming instrumentalist concepts such as measurement and probability but just 'compoundness'?

Pirsa: 09050010 Page 107/108

