

Title: XRay observations normalized ICM models: SZ scaling relations and cosmological implications

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Abstract: We build simple, 'top-down', models for the gas density and temperature profiles for clusters of galaxies based on current high precision XRay observations so as to 'exactly' satisfy observed XRay scaling relation between temperature and mass. The gas is assumed to be in hydrostatic equilibrium along with a component of non-thermal pressure due to dispersion and the gas fraction reaches universal value only at or beyond the virial radius. For these models, we calculate the Sunyaev-Zel'dovich Effect (SZE) scaling relations. We show that all the predicted SZE scaling relations between the integrated SZE flux and the gas temperature, the gas mass, the total mass, as well as, the gas fraction are in excellent agreement with recent SZE observations by Bonamente et al (2008). The consistency between the global properties of clusters detected in X-Ray's and in SZE hints that we are looking at the same population of clusters as a whole. Implications for SZE power spectrum, SZE flux-M200 scaling relation and number counts are discussed

# X-Ray Observations Normalized ICM Models- SZ Scaling Relations and Cosmological Implications

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## SCALING RELATIONS and REASON TO STUDY THEM -

The relations between the different integrated properties of a cluster and with its mass are known as scaling relations.

Ex – Xray Flux – Mass / SZ-flux – Mass / GasTemp – Mass/ Xray Flux – Temp etc

- A) Cosmology – cluster connection is mainly through mass. For example, the mass function of  
▪haloes (i.e clusters) depends sensitively on cosmology (like  $\sigma_8$ , DE etc).  
▪Many cluster surveys are planned/ongoing use clusters as cosmological probes.  
▪Similarly, new CMB anisotropies due to clusters depend on how these clusters of diff masses  
▪are distributed in the sky.
- B) Unfortunately, for most of these surveys/CMB anisotropies clusters masses are not directly  
▪measured. They have to be inferred from cluster observables (like flux).
- So, **KNOWLEDGE OF SCALING RELATION IS CRUCIAL FOR USING CLUSTERS IN COSMOLOGY.**  
▪ **INCORRECT SCALINGS ----> INCORRECT COSMOLOGY.**
- C) Theoretical understanding needed because observationally one cannot go to large cluster  
▪radii to cover the entire mass of the cluster.



## SCALING RELATION IN OBSERVATIONS AND THEORY -

### OBSERVATIONS -

1. X-Ray/SZ observations can give the flux over a volume of the cluster. The volume is mainly limited by the sensitivity of the observations (ex- for xrays  $r_{2500}$ , sometimes  $r_{500}$ , SZ can be observed further away).
2. XRay spectroscopy gives temperature. Older observations assumed isothermality. Now, we can get temp profiles. Temperature quoted are weighted over range in radii by the Xray emission.
3. Masses can be inferred in several ways – a) directly from lensing, b) indirectly from dynamics of galaxies, c) from assuming the gas to be in hydrostatic equilibrium. Each has its advantages/disadvantages.

### THEORY ( Analytical or Simulations) -

1. In simulations, gas and dark matter particles are evolved from tiny perturbations at high  $z$  to form clusters at lower  $z$ 's. The distribution of DM + gas, directly gives the different density and temp profiles. The DM follows a universal profile known as NFW profile. Analytic models assume spatial symmetry.
2. From the temperature profiles, one can construct the 'emission weighted or 'spectroscopic-like' weighted average temperature for a cluster.
3. Different fluxes (say Xray/SZ) or masses within a cluster radius can be obtained from summing over particles (in simulations) or by integrating the analytical density profiles.

## ON SCALING RELATIONS AND SELF-SIMILARITY -

**Assumption:** Gravity ONLY decides the thermodynamic state of the ICM. No preferred scale  $\rightarrow$  hence self-similar (first proposed by Kaiser)\_

Mass  $M_{\Delta_c}$  is enclosed in radius

$$\begin{aligned} R_{\Delta} &\rightarrow M_{\Delta} \sim \rho_c(z) \Delta_c R_{\Delta}^3 \\ \rho_c(z) &\sim E^2(z)_- \\ R &\sim M^{1/3} E^{-2/3}(z)_- \end{aligned}$$

**1) M-T reln:**

$$M_{\Delta_c} \propto T^{3/2} E^{-1}(z)$$

**2)  $L_X$ -T reln:**

$$L_X = \int_V \left( \frac{\rho_{\text{gas}}}{\mu m_p} \right)^2 \Lambda(T) dV$$

$$L_X \propto M_{\Delta_c} \rho_c T^{1/2} \propto T^2 E(z)$$

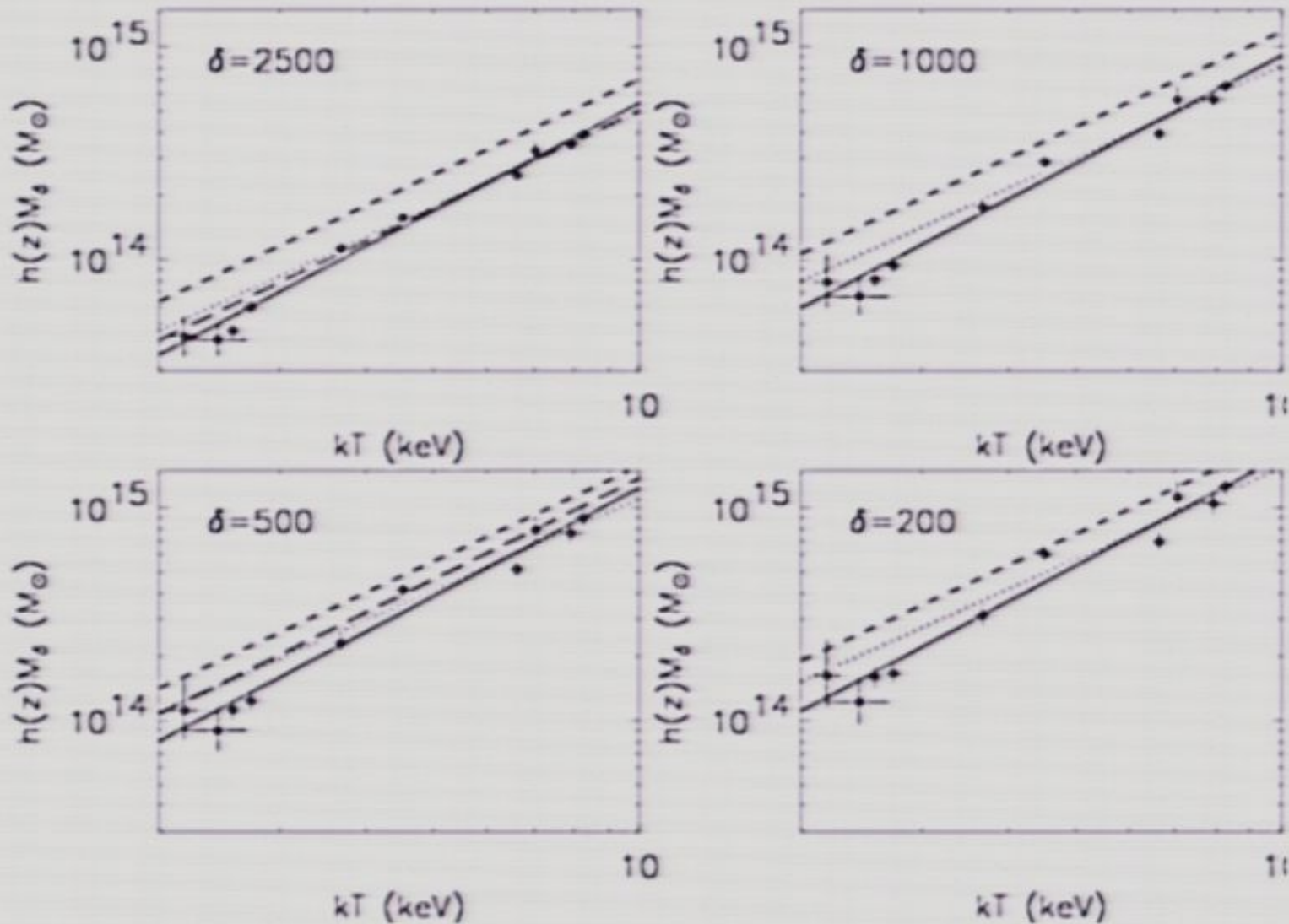
**3) SZ relns:**

$$\Delta S \propto \int y(\theta) d\Omega \propto d_A^{-2} \int T n_e d^3r \propto d_A^{-2} T^{5/2} E^{-1}(z)$$

$$y_0 \propto T^{3/2} E(z) \propto L_X^{3/4} E^{1/4}(z)$$



## SIMULATIONS-VS-OBSERVATIONS -



(Observations by Arnaud et al, Simulations by Evrard et al)

Non radiative sims way off.

Radiative is better but has much room for improvement (ask Daisuke, Christoph)

**Q) So What is Going On ?**

**A) Our lack of full understanding of the physics !**

Many experts work on, has spoken today, will speak after me, will keep working on...

I'm not going to focus on the micro physics

I'll something much naïve, something more phenomenological.

...and see what we get out of constructing simple model and observations!

Bottomline - Let X-Ray observations (much precise) aid SZ science



A bit of motivation -

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Most analytical or simulation models are '**bottom-up**' models.  
Simulations have similar problems to KS model. Both have many similarities.

Present cluster models faces tension :

1. It fails to reproduce scaling relations at low masses (Xrays)
2. Used as SZ cluster templates -vs- "CBI-excess" .

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Caveats – 1) Need good observations at higher redshifts also

- 2) There are no observations at virial radius.
- 3) SZ surveys have started but only two SZ observational scaling papers.
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**We build a very simple 'top-down' cluster model that fulfils these requirements.  
Easier to run MCMC with analytic models**



## ANALYTIC MODELLING OF GALAXY CLUSTER - THE DARK MATTER

1. The primary component of a cluster is the dark matter halo (NFW).
2. The size / mass of a halo - depends on our definition.
  - a) They are connected by a constant over-density
  - b) Choice of overdensity -
    - i) xray observations typically go upto  $\Delta = 2500, 500$
    - ii) Nbody sims use  $\Delta = 200, 180\Omega_M$
    - iii) Virial radius is  $\Delta = 170$  (for sCDM),  $\sim 100$  (LCDM)

$$r_{vir} = \left[ \frac{M_{vir}}{\frac{4\pi}{3} \rho_{crit}(z) \Delta_c(z)} \right]^{1/3}$$

- c) Given NFW profile, one can convert between definitions.
- d) Clusters do not have a well-defined outer boundary. A 'good' boundary is the shock radius \_ typically 2-3 virial radius.

Will be important later for normalization

The underlying variation is in halo concentration with cluster mass and redshift.  
We use observed c-M relation from Comerford & Natarajan (2008)



## ANALYTIC MODELLING OF GALAXY CLUSTER - THE BARYONIC MATTER

1. Baryonic matter is dominated by Intra-cluster gas.
2. Following early X-ray observations, the most common and still used by many observers is the so called  $\beta$ -model, first proposed by Cavaliere & Fusco-Femiano (1976)\_

$$n_e(r) = n_{e0} \left( 1 + \frac{r^2}{r_c^2} \right)^{-3\beta/2}$$

$$M_{\text{tot}}(r) = \frac{3\beta k T_X}{G\mu m_p} \frac{r^3}{r_c^2 + r^2}$$

Recent observations and simulations for long time shows model to be incorrect at small and large radii. Especially, gas profiles tend to be steeper at outer region

--> **model underestimates total mass.**

3. Later improvement is incorporating non-isothermality and analytically solving for gas in equilibrium in NFW haloes (spherical/triaxial) by Makino & Suto. But has free parameters.
4. Komatsu & Seljak (2001) has proposed a 'Universal model for cluster gas' based on simple assumptions (with no free params). This is the **favoured model** at present.

## THE TEMPERATURE PROFILE -

### Cool Core Clusters -

Due to increased density at cluster cores, cooling time is less than Hubble time \_ gas can cool easily and forms cool core. Gas flows in to form 'cooling-flows'.

We motivate our temp profiles from X-ray observation :

- 1) within the cool cores (Sanderson et al 2007)

$$T(r) \propto r^{-4}$$

- 2) Outside the cool cores, we take the temp profile to be 'polytropic'  
(Vikhlinin et al 06, Sanderson 07, Arnaud et al 07, Sun et al 08)

$$T(r) \propto \rho(r)^{\gamma-1}$$

### Non Cool-core Clusters -

The temperature profile is taken to be polytropic through out. It is almost flat in the inner radii.

(Sanderson et al 2007, O'Hara et al 2007)

Our final temp profiles matches very well the trend and extent seen in observations of falling temperatures at large cluster radii.



## GAS DYNAMICAL EQUATION AND THE DENSITY PROFILE -

Solve

$$\frac{d\phi(x)}{dx} = \underbrace{\frac{1}{\rho(x)} \frac{dP(x)}{dx}}_{\text{THERMAL PRESSURE}} + \underbrace{\frac{1}{\rho(x)} \frac{d[\rho(x)\sigma_r^2(x)]}{dx}}_{\text{DISPERSION PRESSURE}} + \underbrace{2\beta(x) \frac{\sigma_r^2(x)}{x}}_{\text{ANISOTROPIC PRESSURE}}$$

$$x = r/r_c$$

We neglect the anisotropic term (typically small)

In general, most analytic models neglect non-thermal terms and only use term 1 (example, Komatsu-Seljak)

$$\frac{d\phi(x)}{dx} = \frac{1}{\rho(x)} \frac{dP(x)}{dx}$$

We know that neglecting non-thermal pressure biases mass-estimate  
--> and hence scaling relations.

Bias is as bad as working with a beta-model mass estimate.

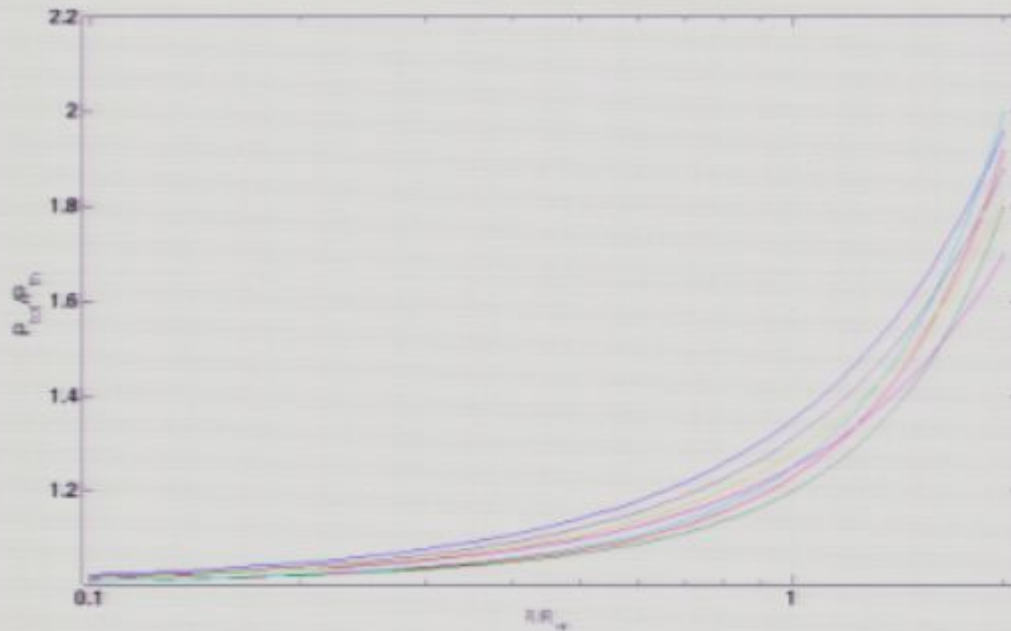


## NON THERMAL PRESSURE CONTRIBUTION -

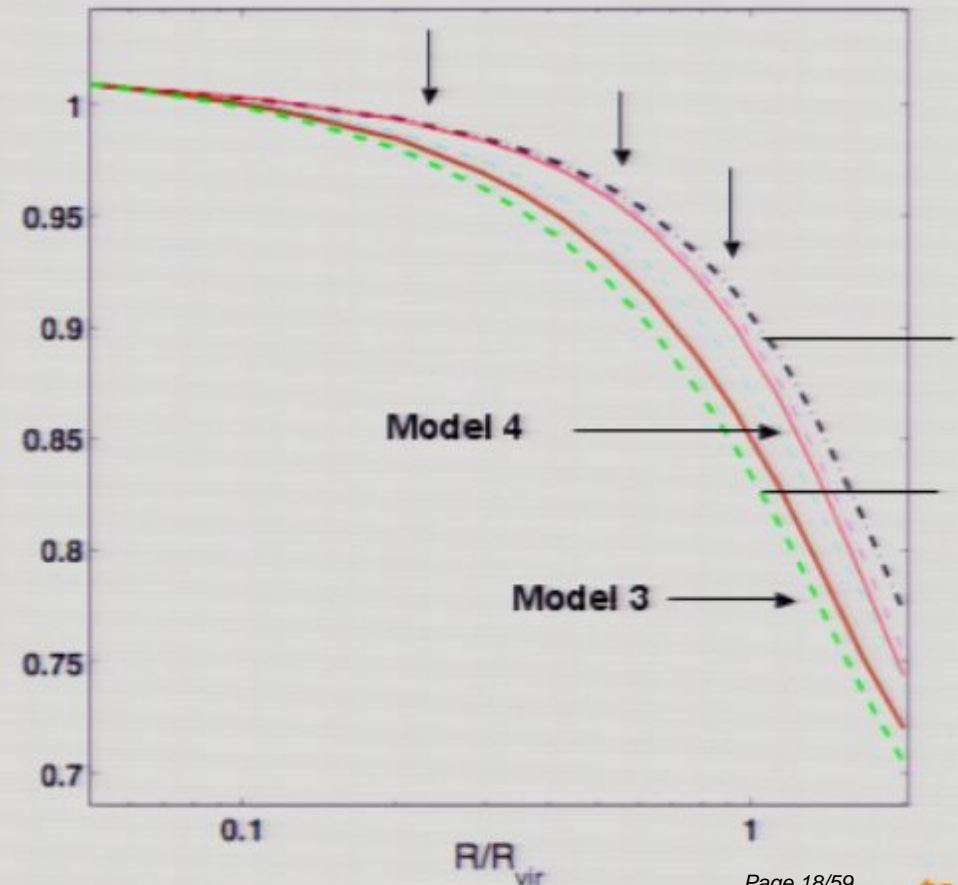
Xray observations do not give us any direct inputs into the non-thermal pressure support.

Simulations give a handle on non-thermal pressure. However, instead of adopting any particular simulation results, we only adopt the relative thermal/non-thermal pressure. This quantity is more 'robust' across different simulations

Rasia, Pfrommer pvt communication



$$P_{NT}/P_T = \frac{\rho(x)\sigma_r^2(x)}{P(x)}$$



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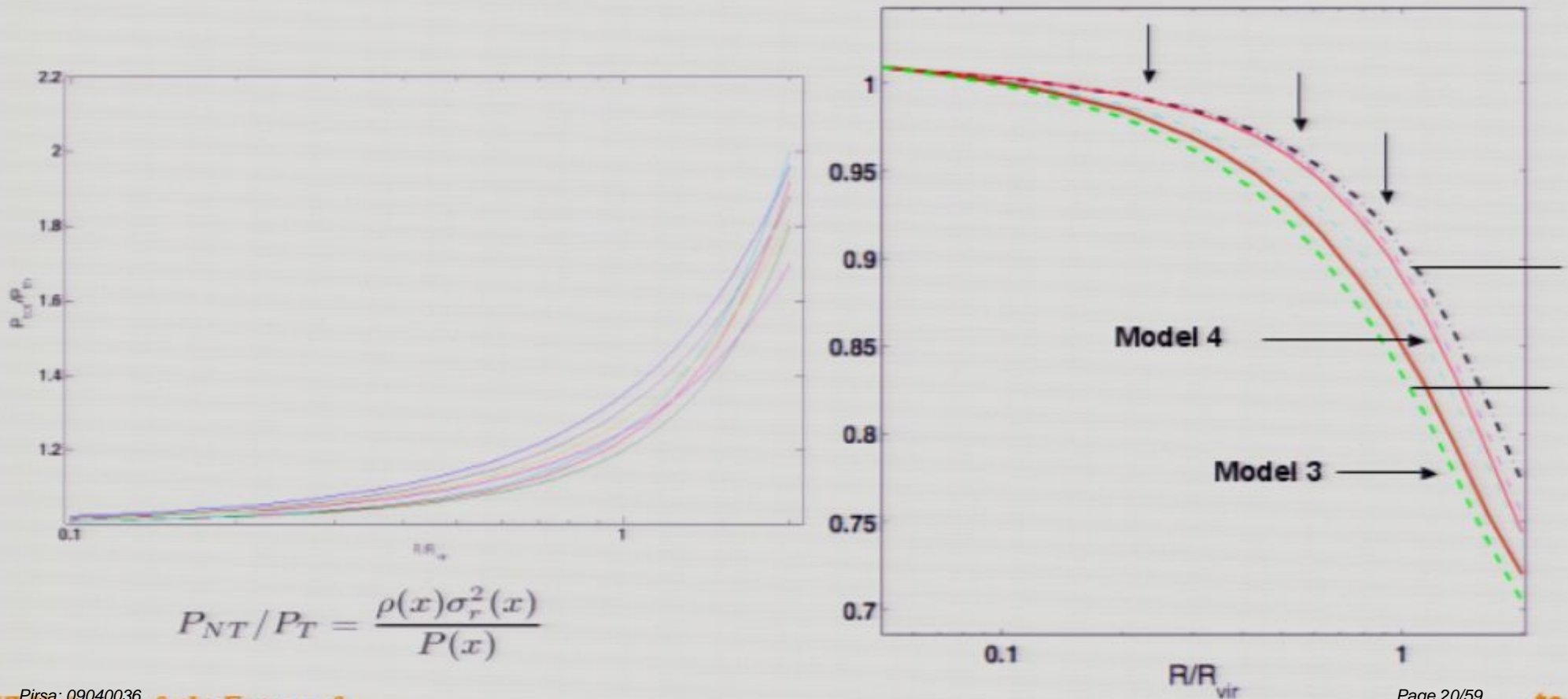
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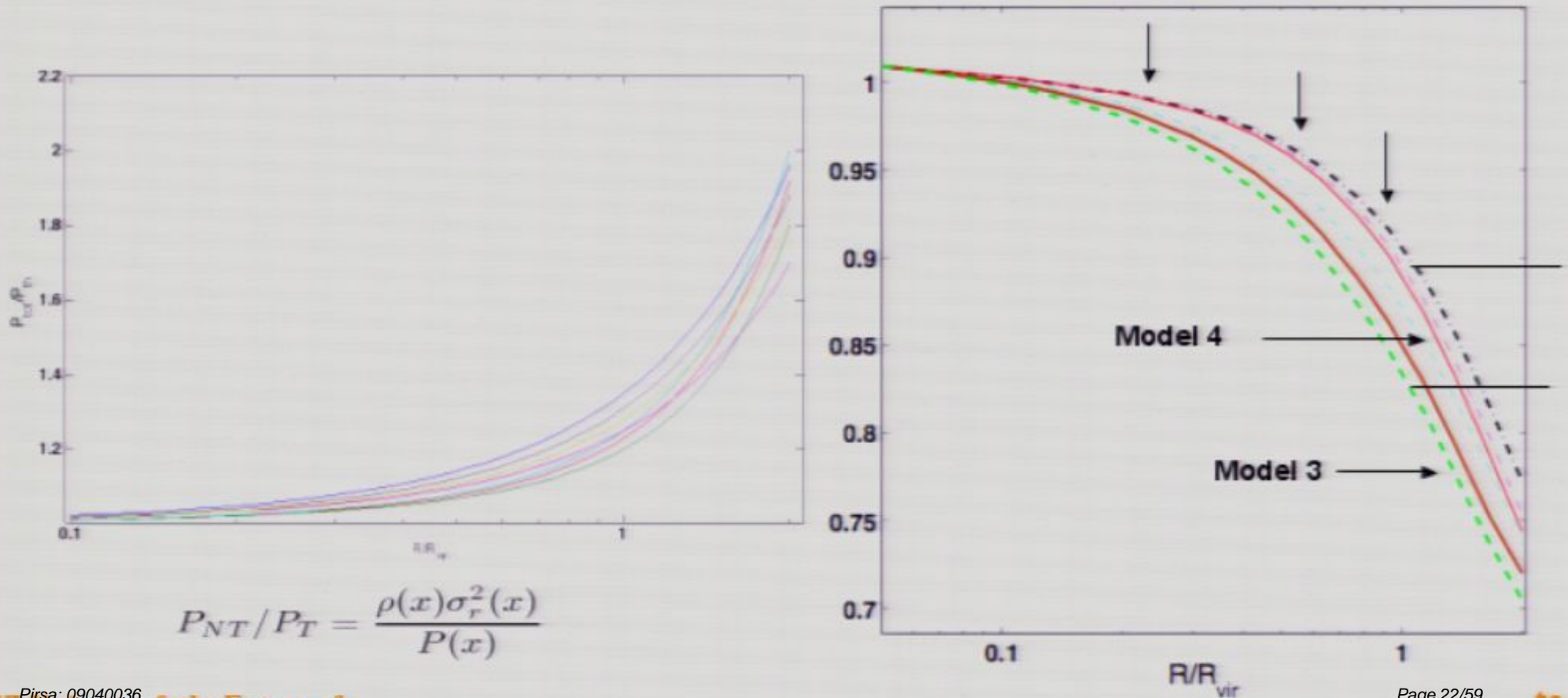
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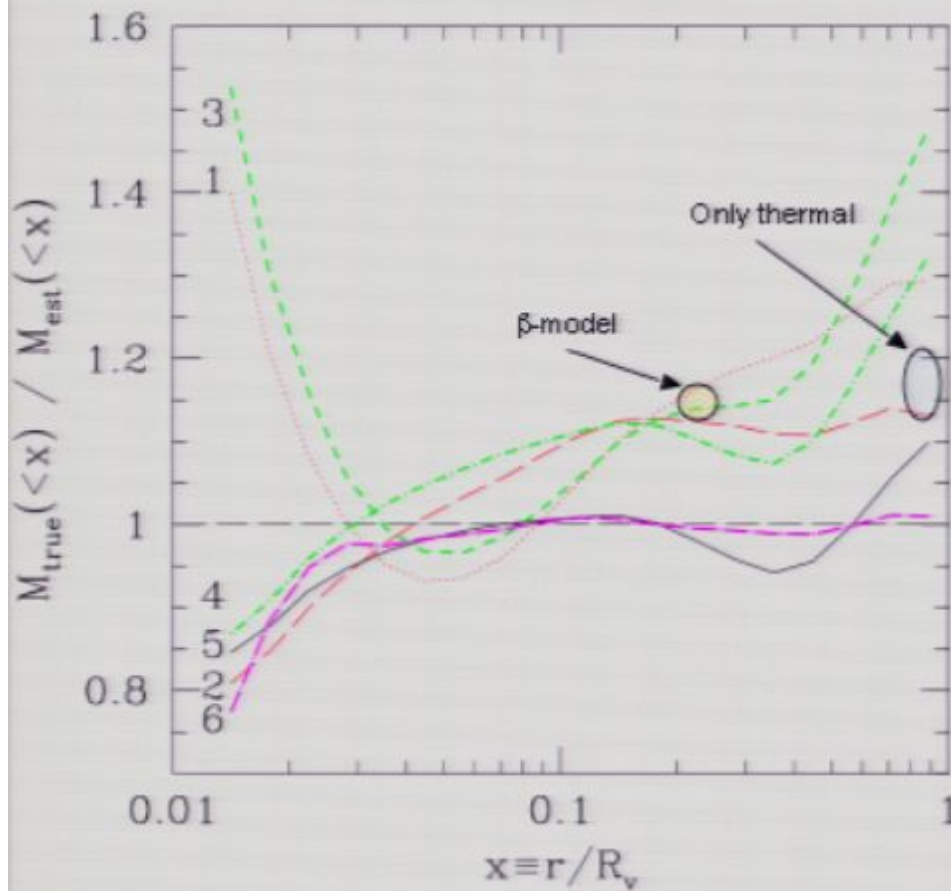
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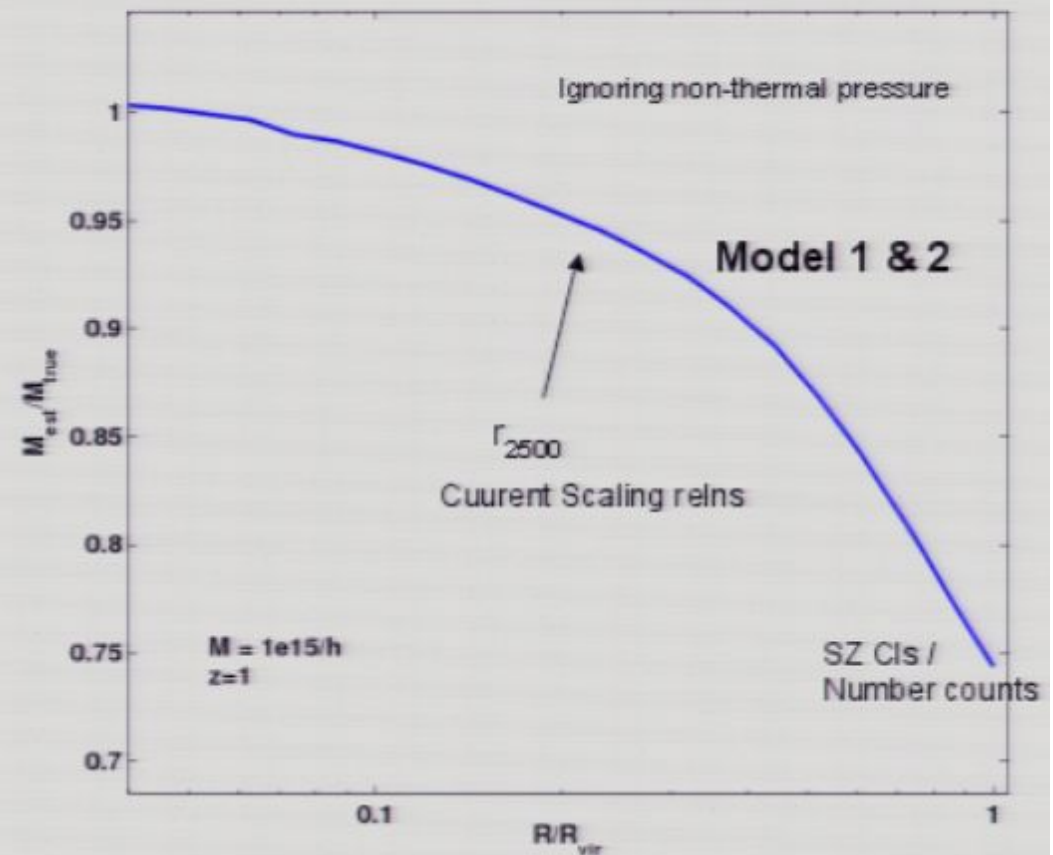


## ACTUAL CLUSTER MASS – VS – RECOVERED MASS -

Simple hydrostatic eqm or ignoring non-thermal pressure, assumption of isothermality as well as commonly used beta model biases recovered cluster mass.



Simulations (Rasia et al 2004)



Model 1 non-thermal component

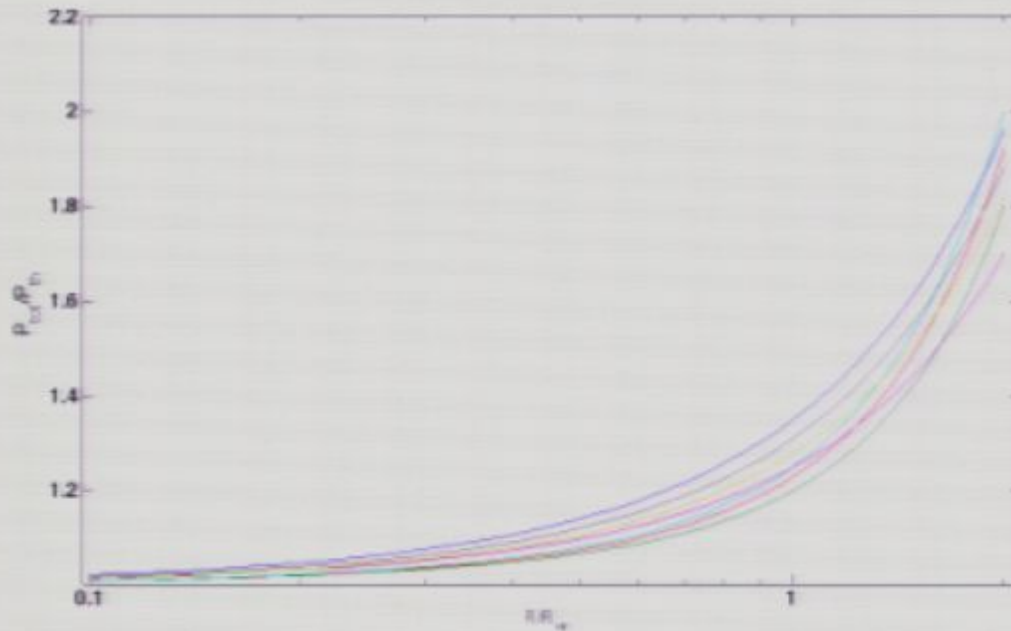


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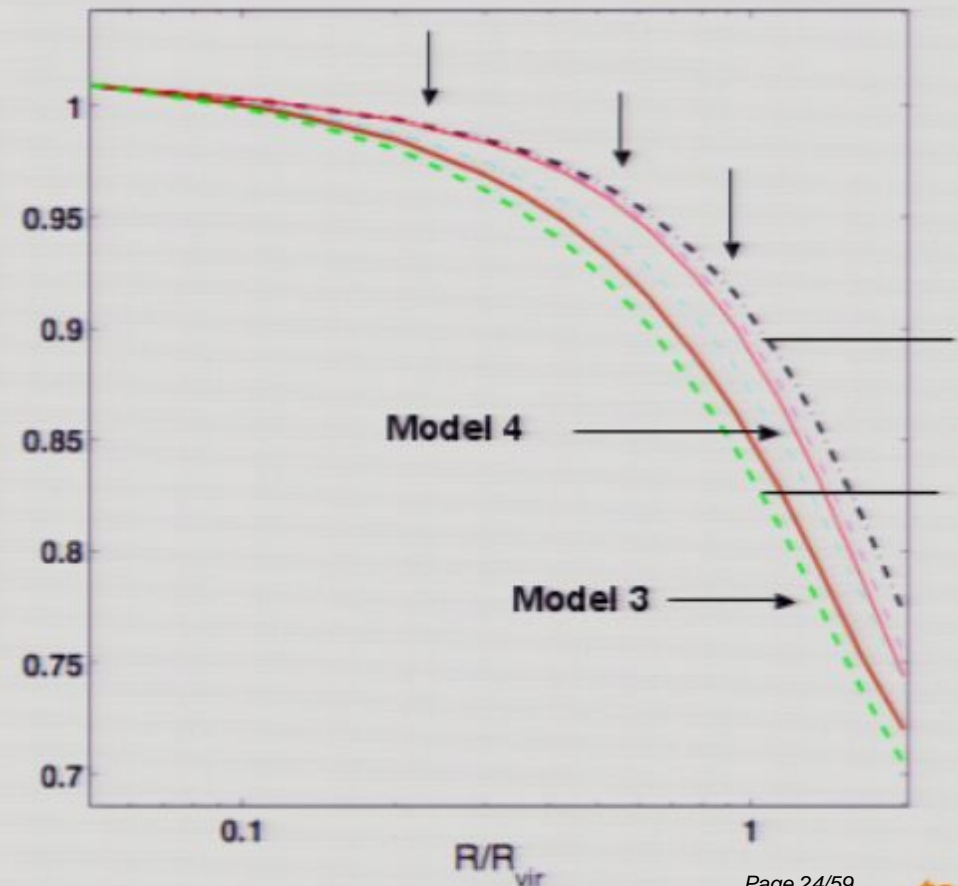
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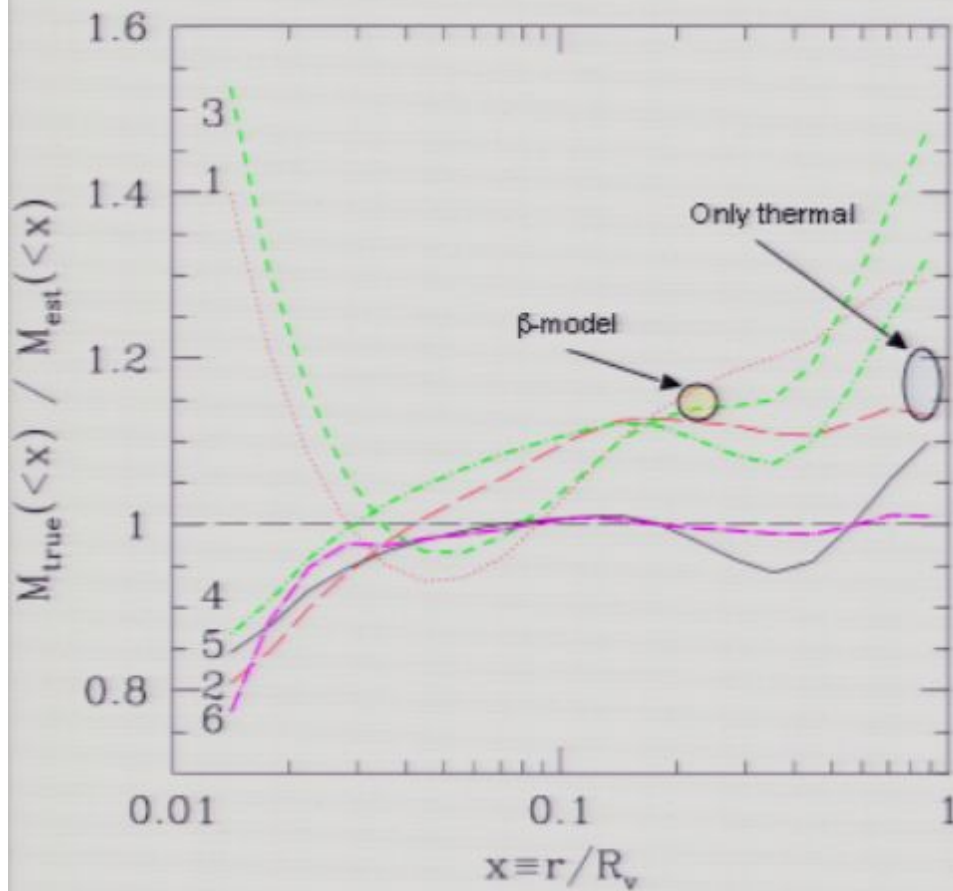


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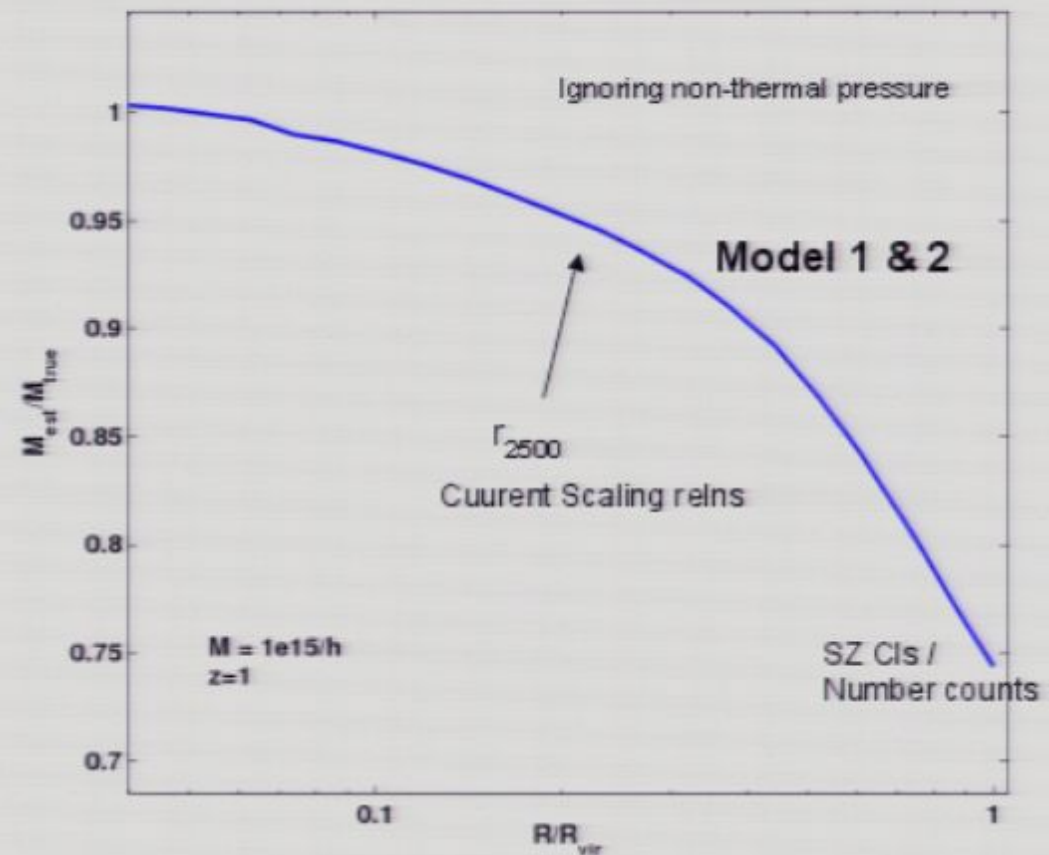


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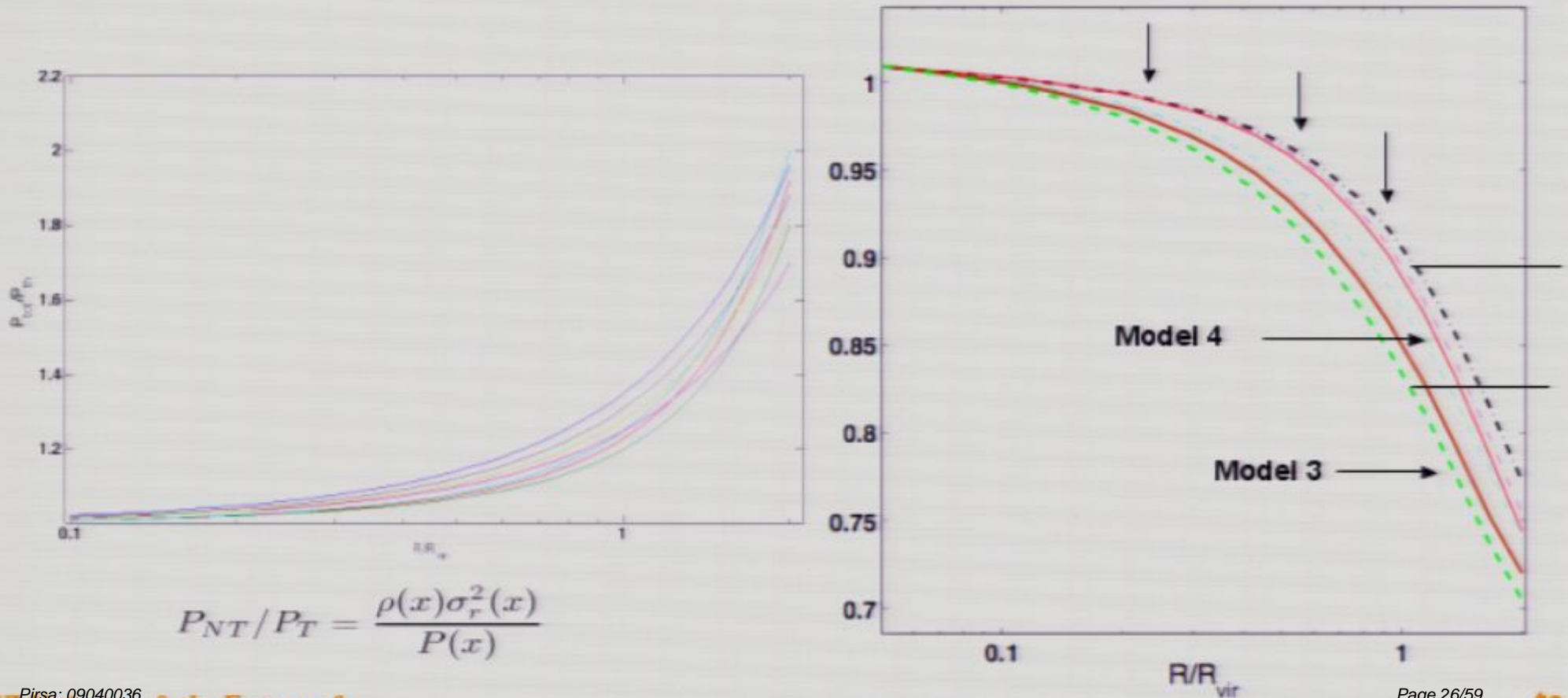


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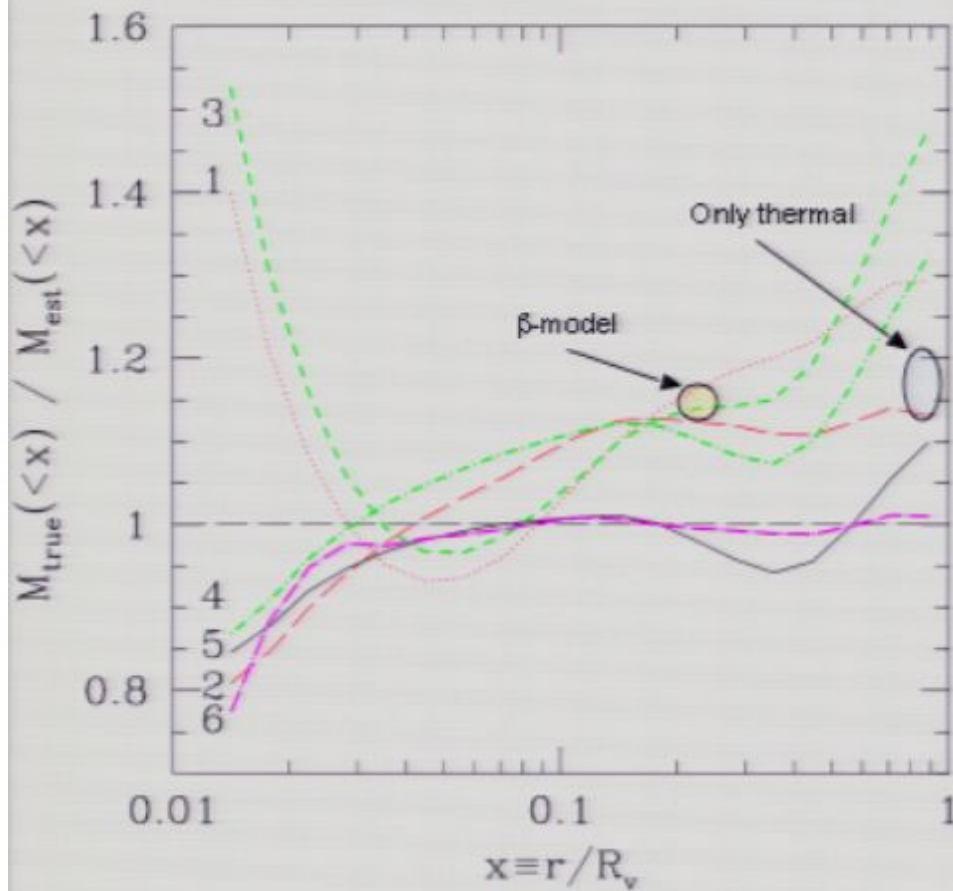
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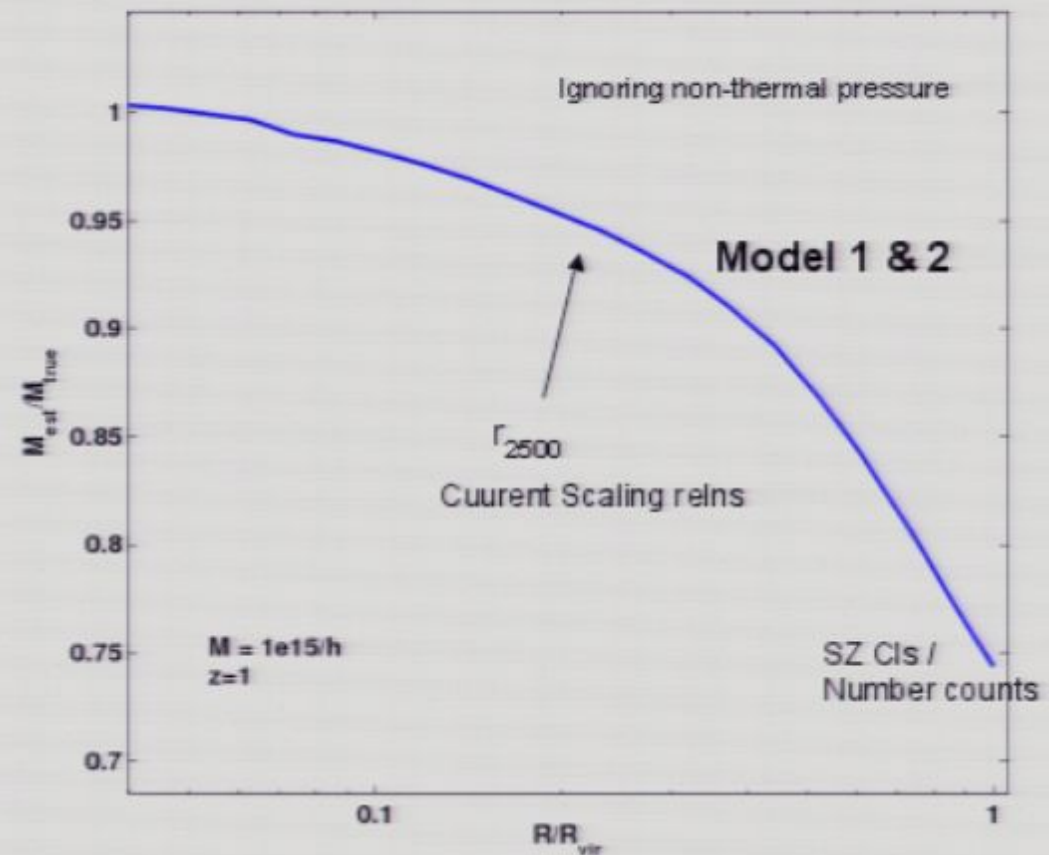


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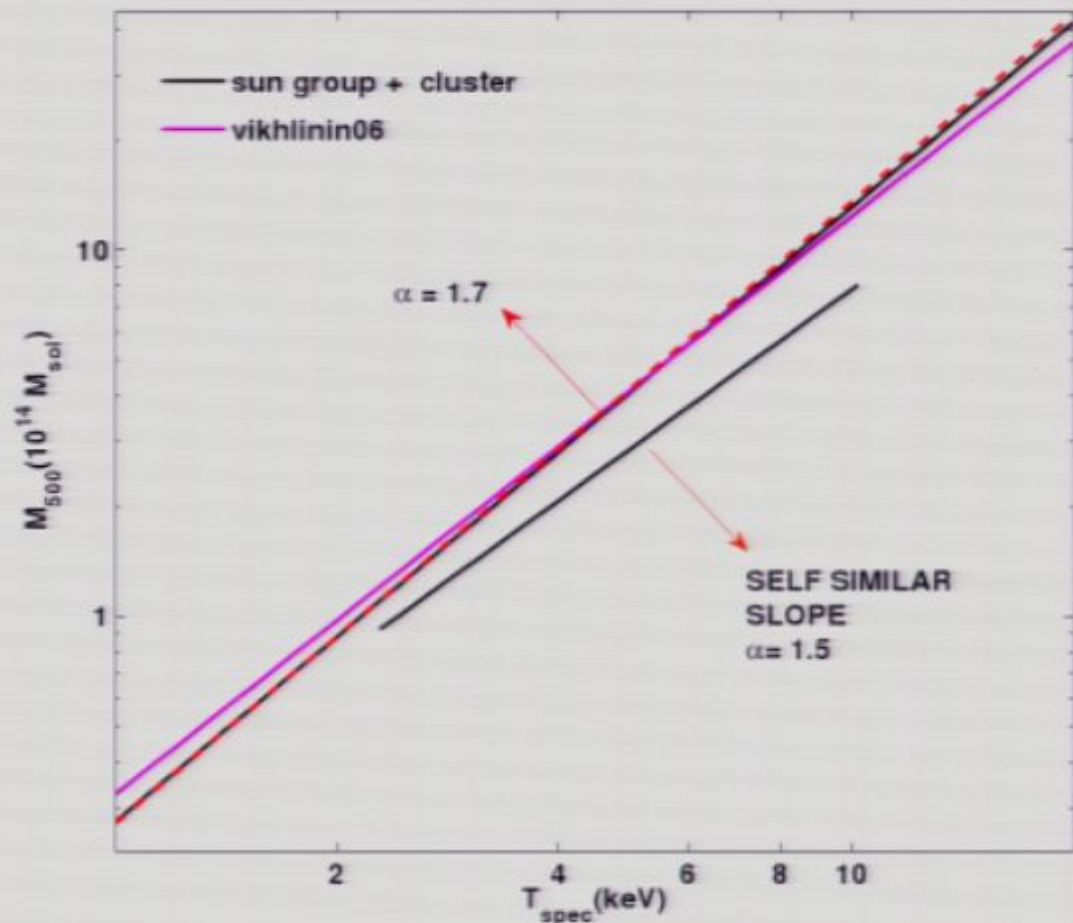
Model 1 non-thermal component

## NORMALIZING TEMPERATURE -

We hide our lack of knowledge of cluster gas physics by --> forcing the temperature to follow Xray observations 'exactly'.

### OBSERVED M-T RELATIONS

$$M_{\delta} E(z) = A \left[ \frac{T}{5 \text{ keV}} \right]^{\alpha}$$



We use M-T relation from Sun et al 2008 (more like a superset of previous data, like Vikhlinin)

$$\alpha = 1.68 \pm 0.04 \quad A = (2.85 \pm 0.18) h^{-1} 10^{14} M_{\odot}$$

## NORMALIZING GAS DENSITY -

Cooling/heating influences gas fraction as a function of radius.  
However, at sufficiently large scale (virial radius and beyond), the integrated gas mass fraction attains Universal value.

$$M_{\text{gas}}(r > R_{\text{vir}}) = \Omega(1) \Omega_b / \Omega_m M_{\text{halo}}(r > R_{\text{vir}})$$

Different cases:

Model 1 - Universal at  $R_{\text{vir}}$ ;  $P_{\text{tot}}/P_{\text{thermal}}$  from Rasia et al (2004).

Model 2 -  $0.9 * \text{Universal}$  and  $R_{\text{vir}}$ ;  $P_{\text{tot}}/P_{\text{thermal}}$  from Rasia et al (2004).

Model 3 - Universal at  $2R_{200}$ ;  $P_{\text{tot}}/P_{\text{thermal}}$  for highest non-therm press

Model 4 - Universal at  $2R_{200}$ ;  $P_{\text{tot}}/P_{\text{thermal}}$  for intermediate non-therm press

Model 5 - Universal at  $2R_{200}$ ; only thermal pressure

Model 6 - Universal at  $R_{\text{vir}}$ ; only thermal pressure



Temperature influences density through dynamical equation. Once mass normalised, the density influences calculation of average temperature, needed for normalization.

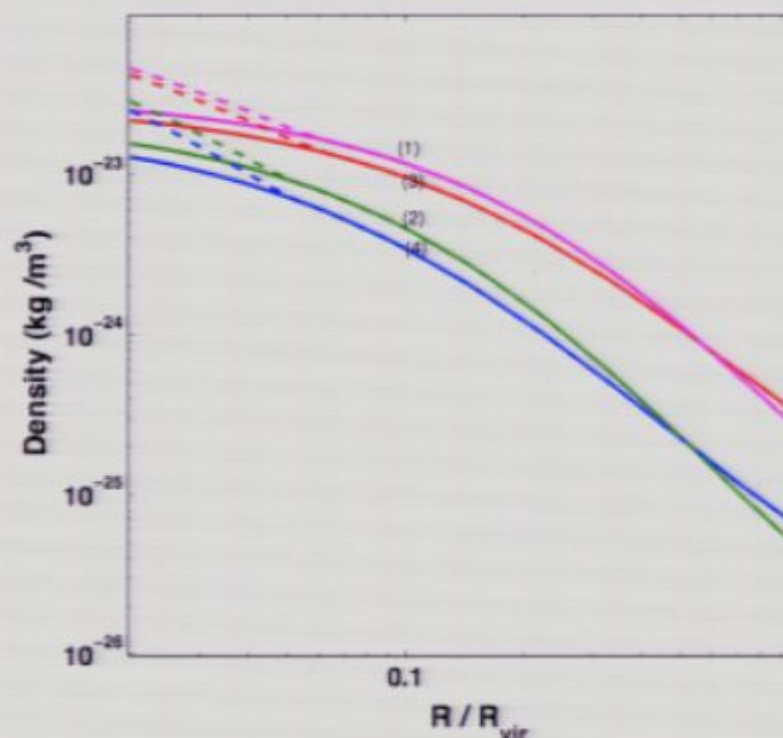
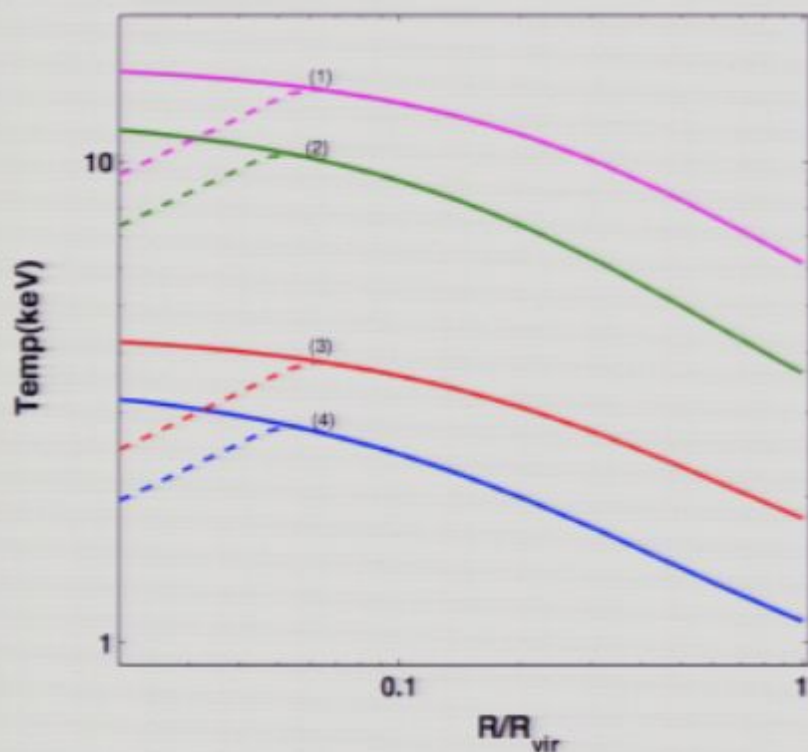
$$T_{sl} = \frac{\int n^2 T^{.75} / T^{1/2} dV}{\int n^2 T^{.75} / T^{3/2} dV}$$

Thus the set of equations need to be solved iteratively, for each model.

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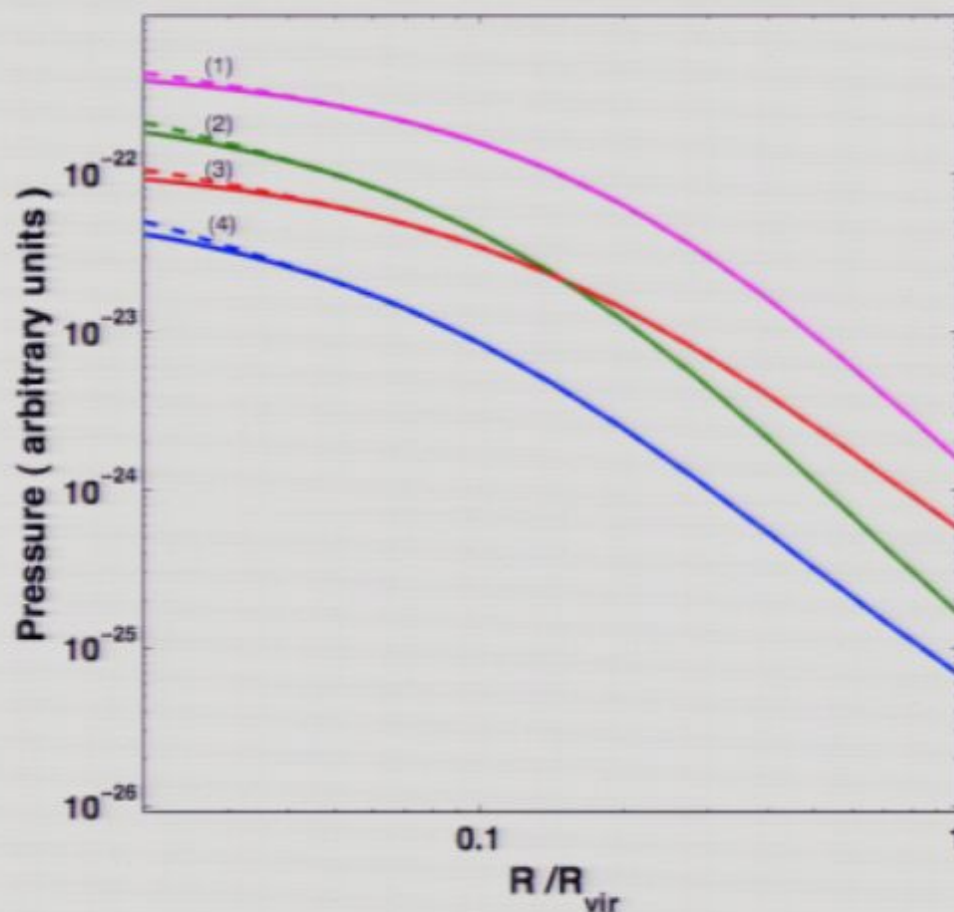
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## Recap -

Till now, our entire modelling of cluster gas density and temperature is based upon

1. Recent high resolution xray observations for temp
2. universality of baryon fraction at largest radii
3. non-thermal component from simulations.

With these, we construct SZ properties of clusters, i.e

1a) SZ cluster scaling relations (and compare with recent Bonamente et al 2008 obs)

$$\log Y = A + B \log X$$

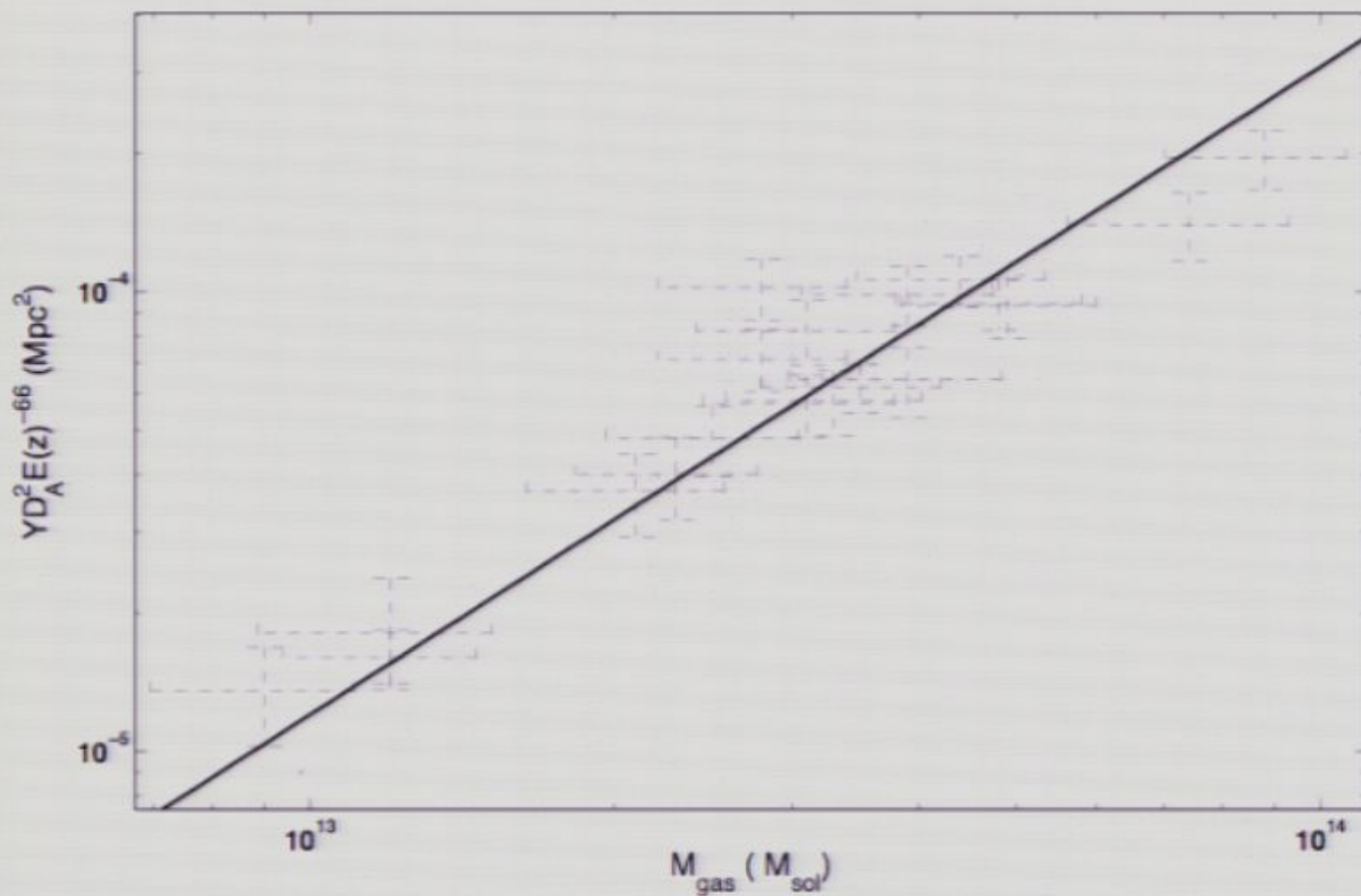
Y = SZ Flux

X = Gas Mass, Total Mass, Temperature

1b) SZ  $M_{200}$  scaling reln needed for  $dndz$  ( more later by Christoph)

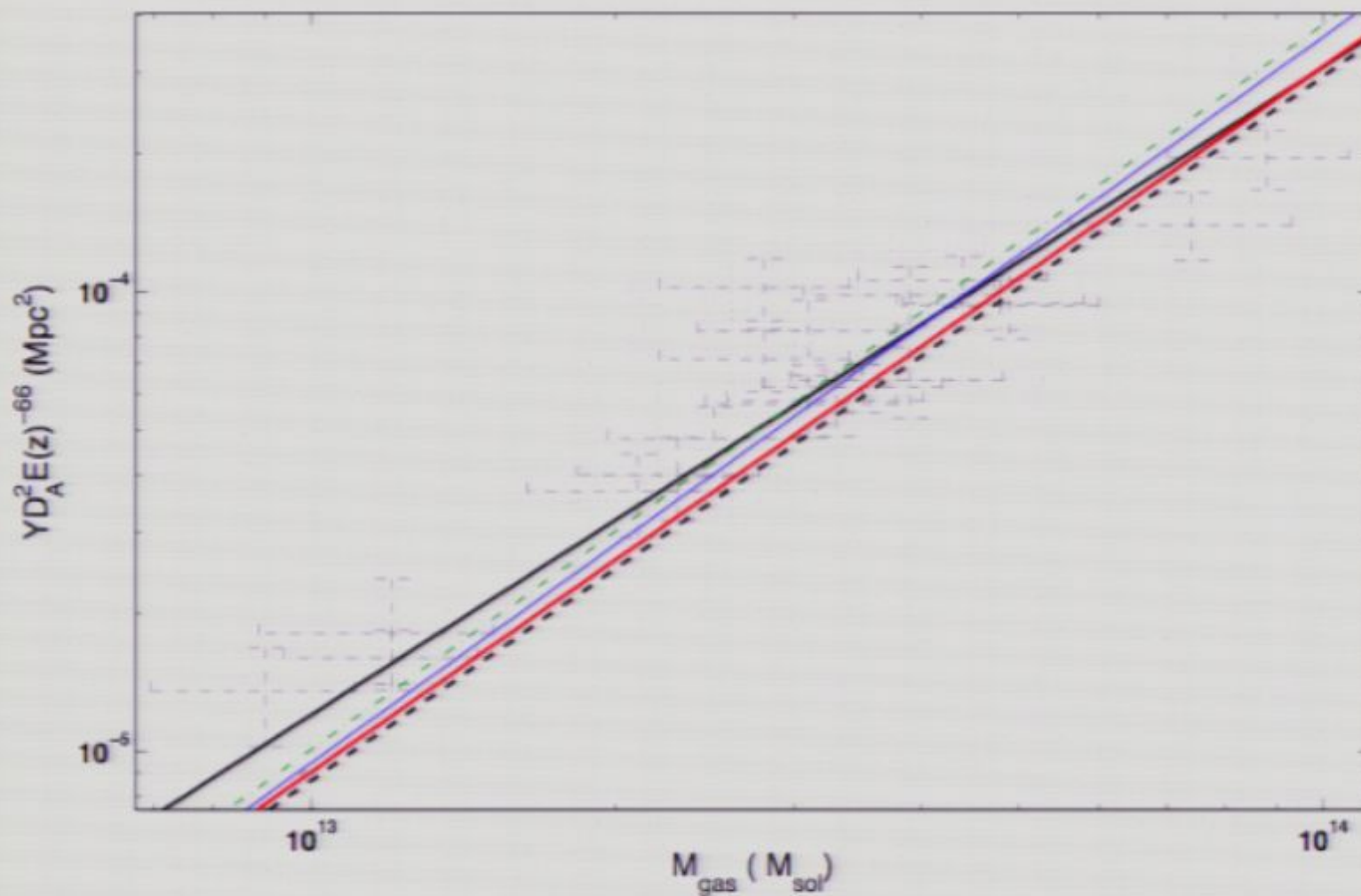
2) SZ power spectrum and compare with CBI+ observations (more later by Jonathan)

**SZFlux -  $M_{\text{gas}}$**

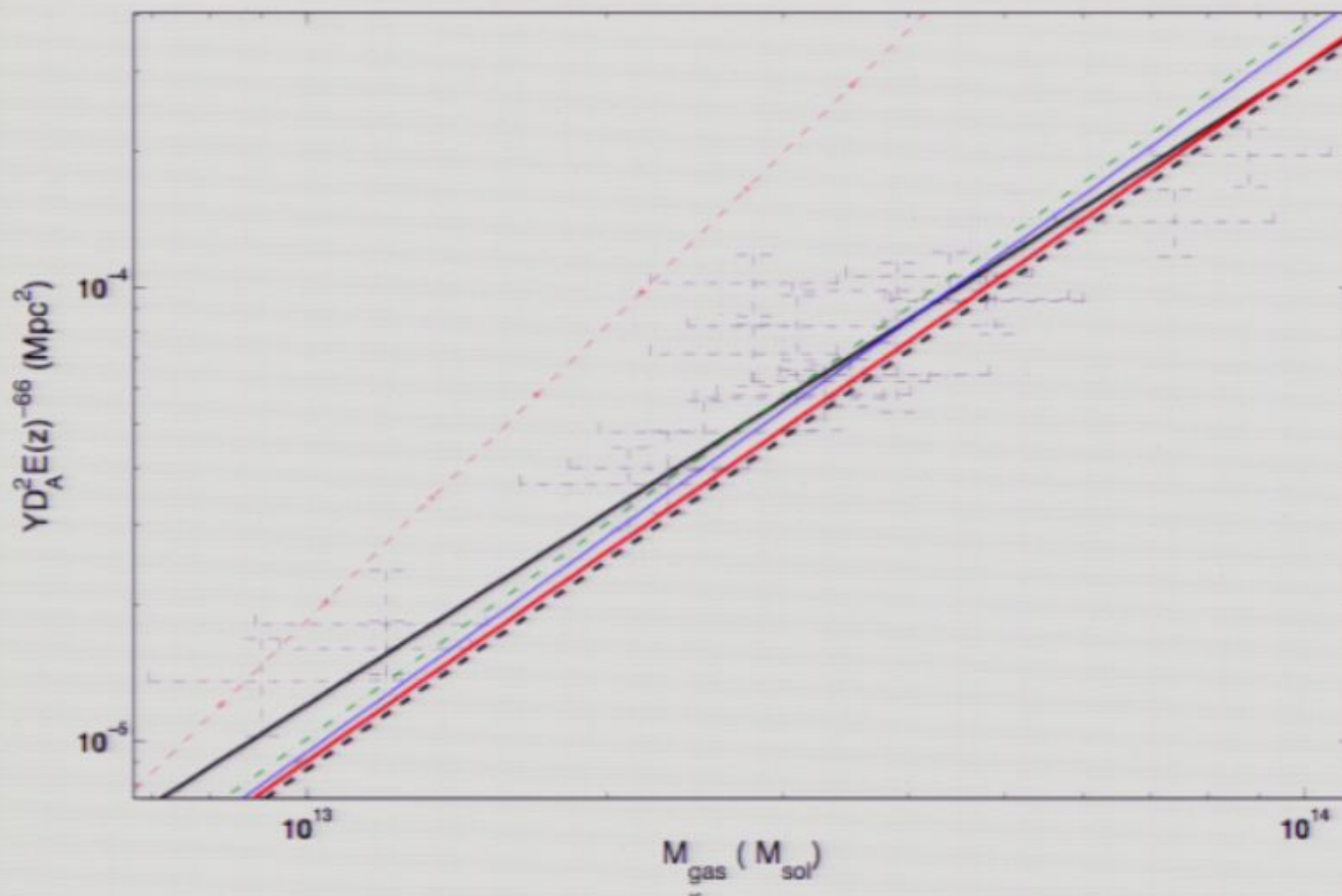


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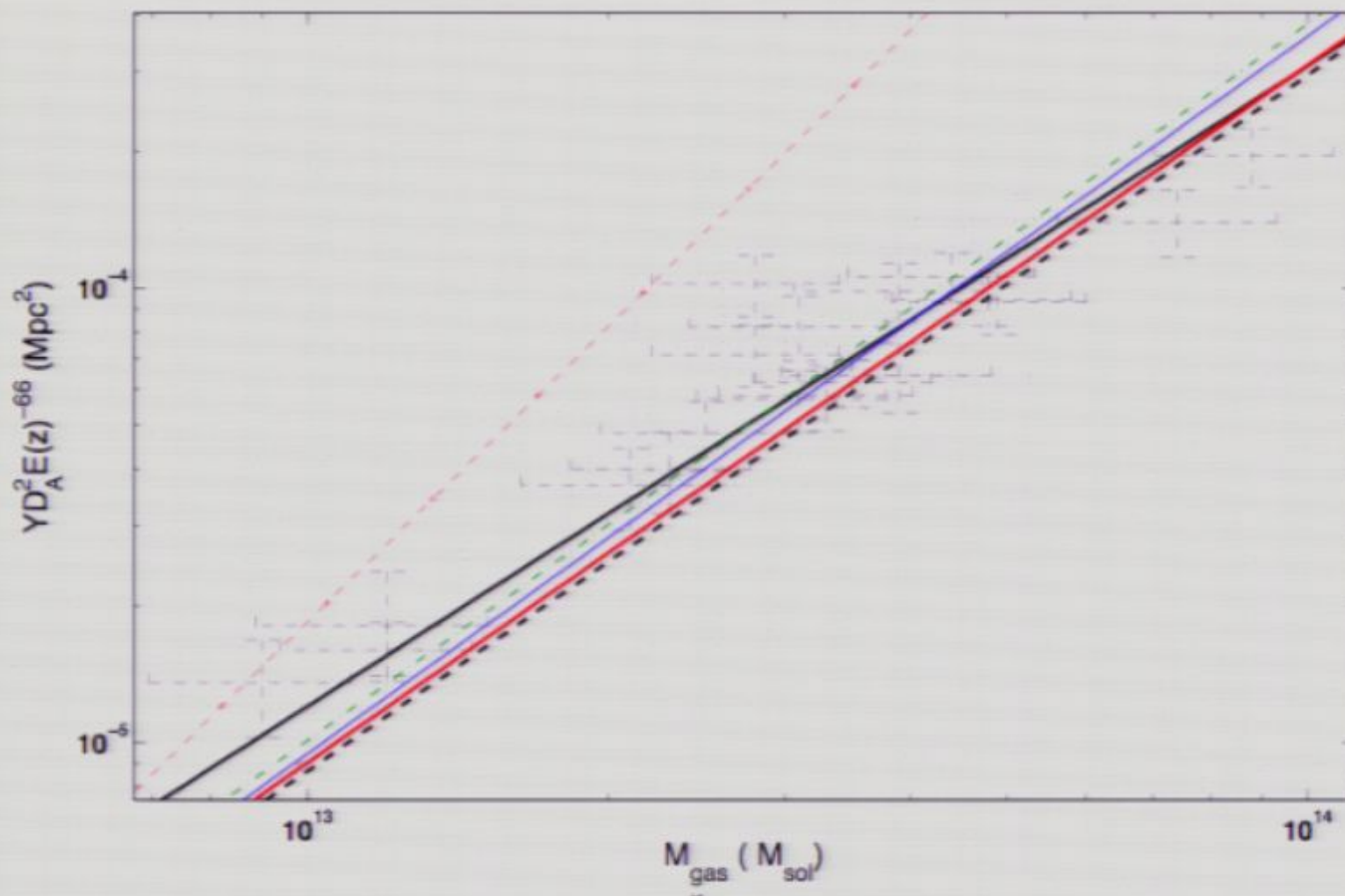




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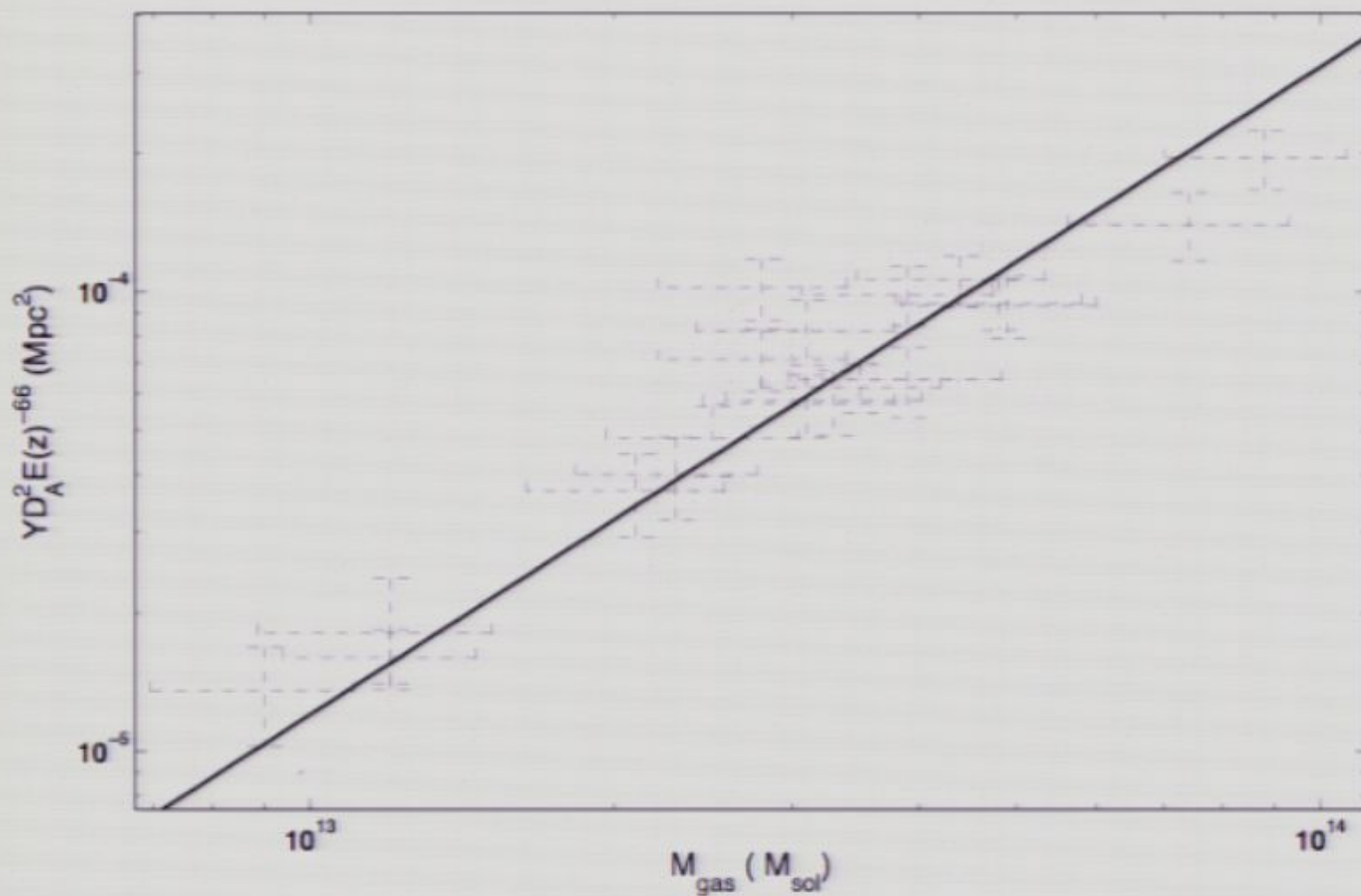
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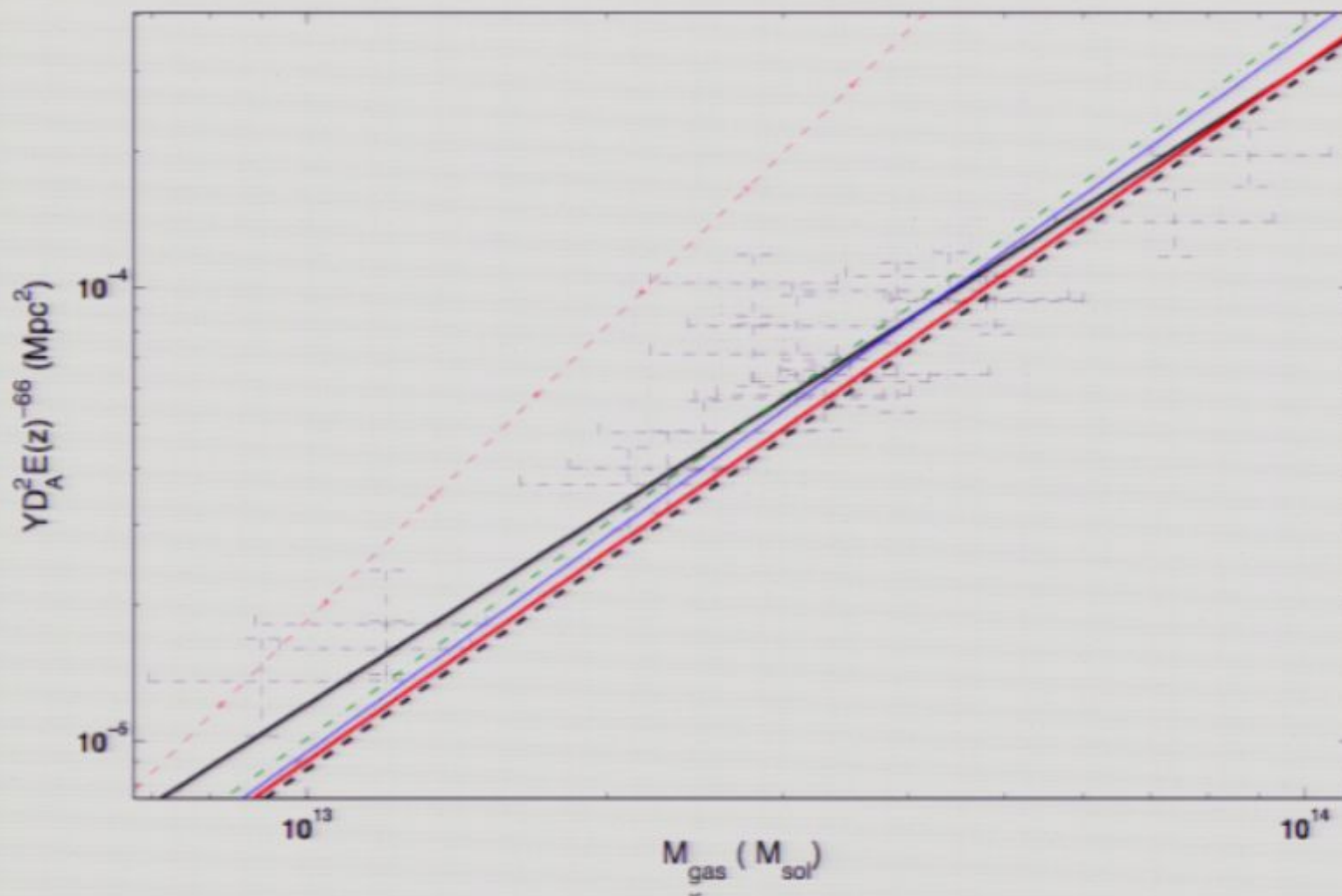
**SZFlux -  $M_{\text{gas}}$**

	A	B	$\chi^2$	A	B	$\chi^2$
Bonamente			.9252			.6713
model 1	-25.55	1.5794	1.0343			.5754
model 2	-25.51	1.5781	.9274			.4987
model 3	-25.0082	1.5357	1.3435			.7916
model 4	-24.98	1.5321	1.5741			.9812
thermal 1	-25.1567	1.5405	2.53			1.48
thermal 2	-25.577	1.5725	2.28			1.31
K S	23.0122	2.1758	10.50			5.2206

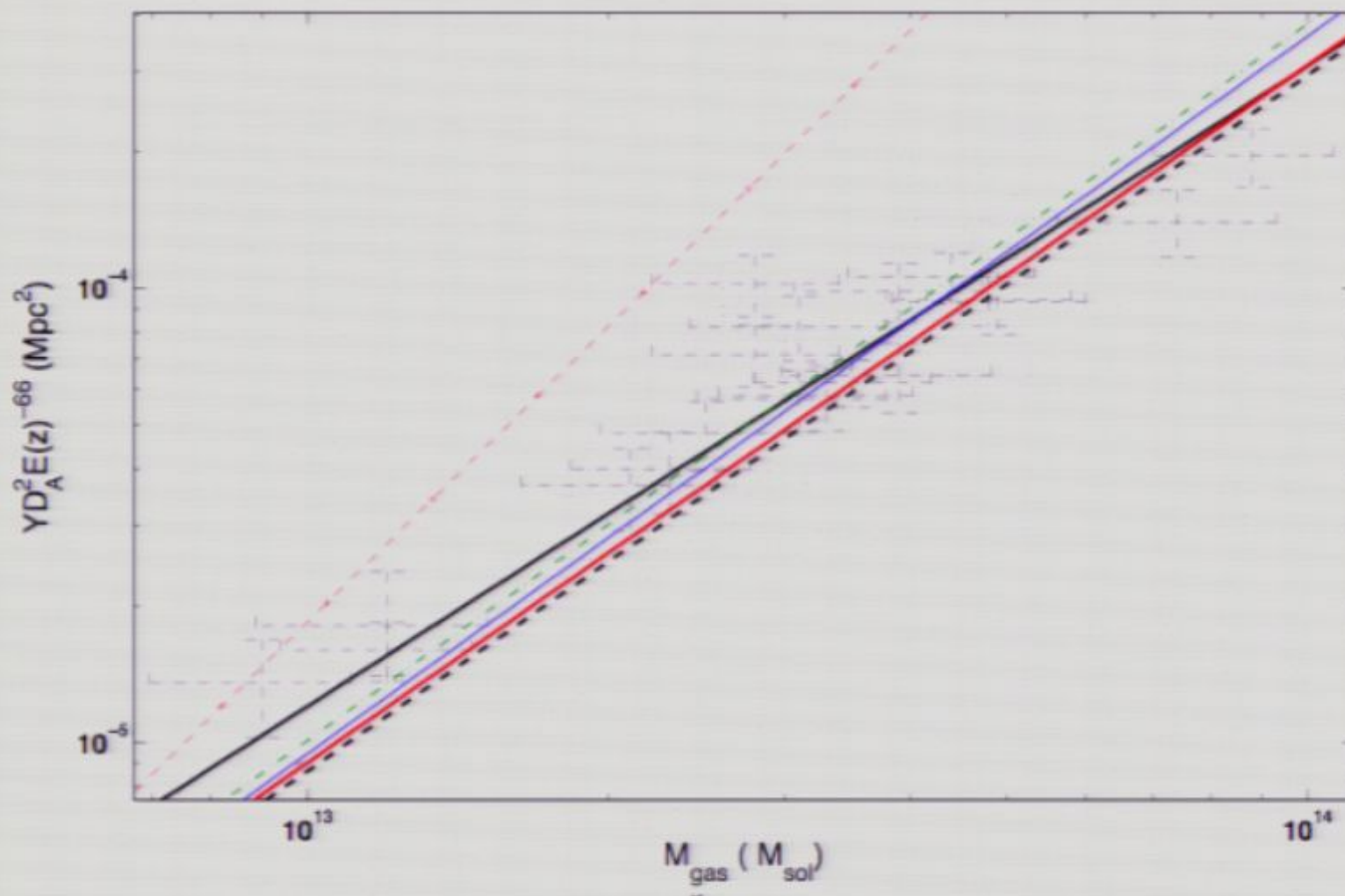




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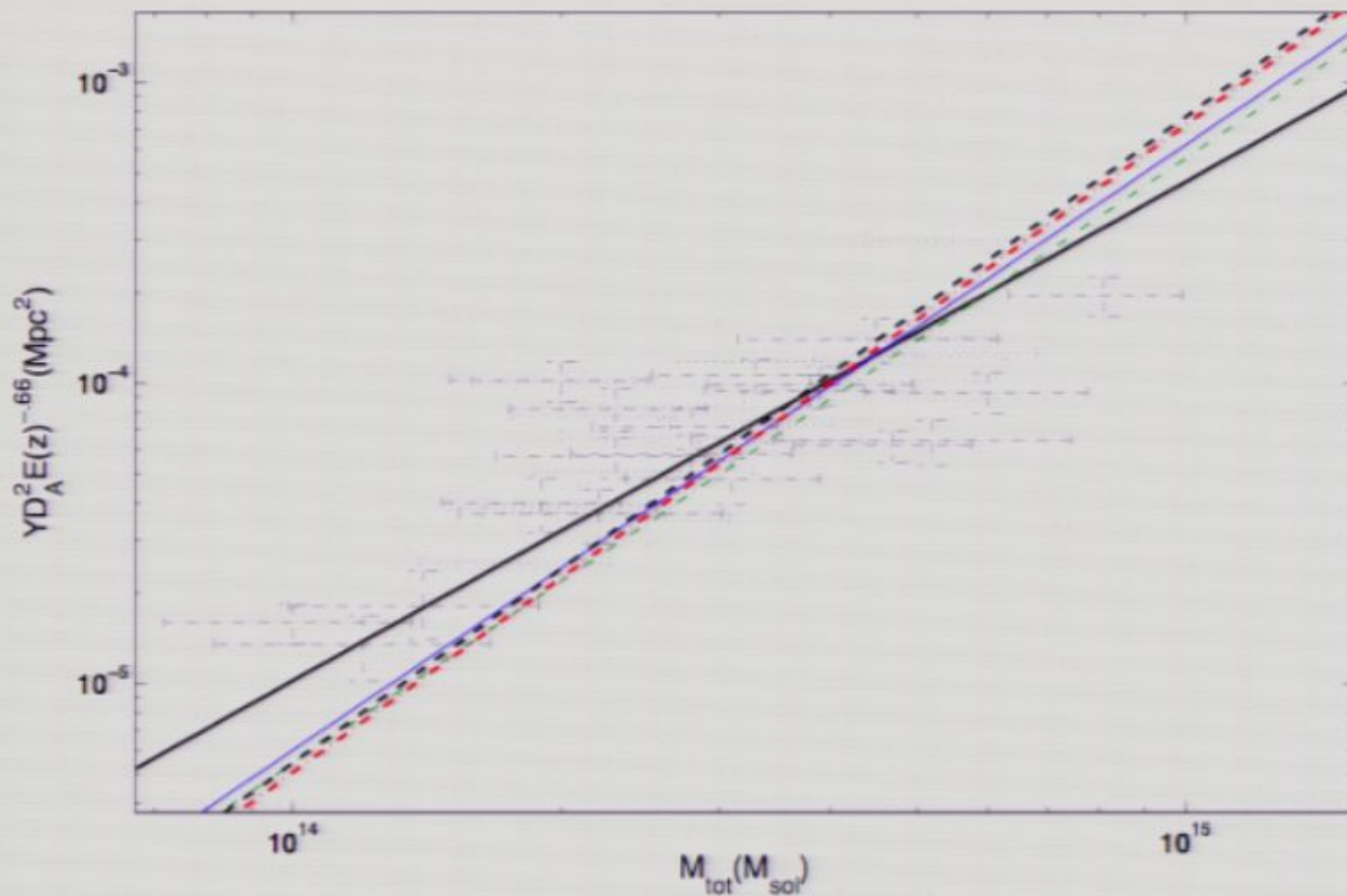


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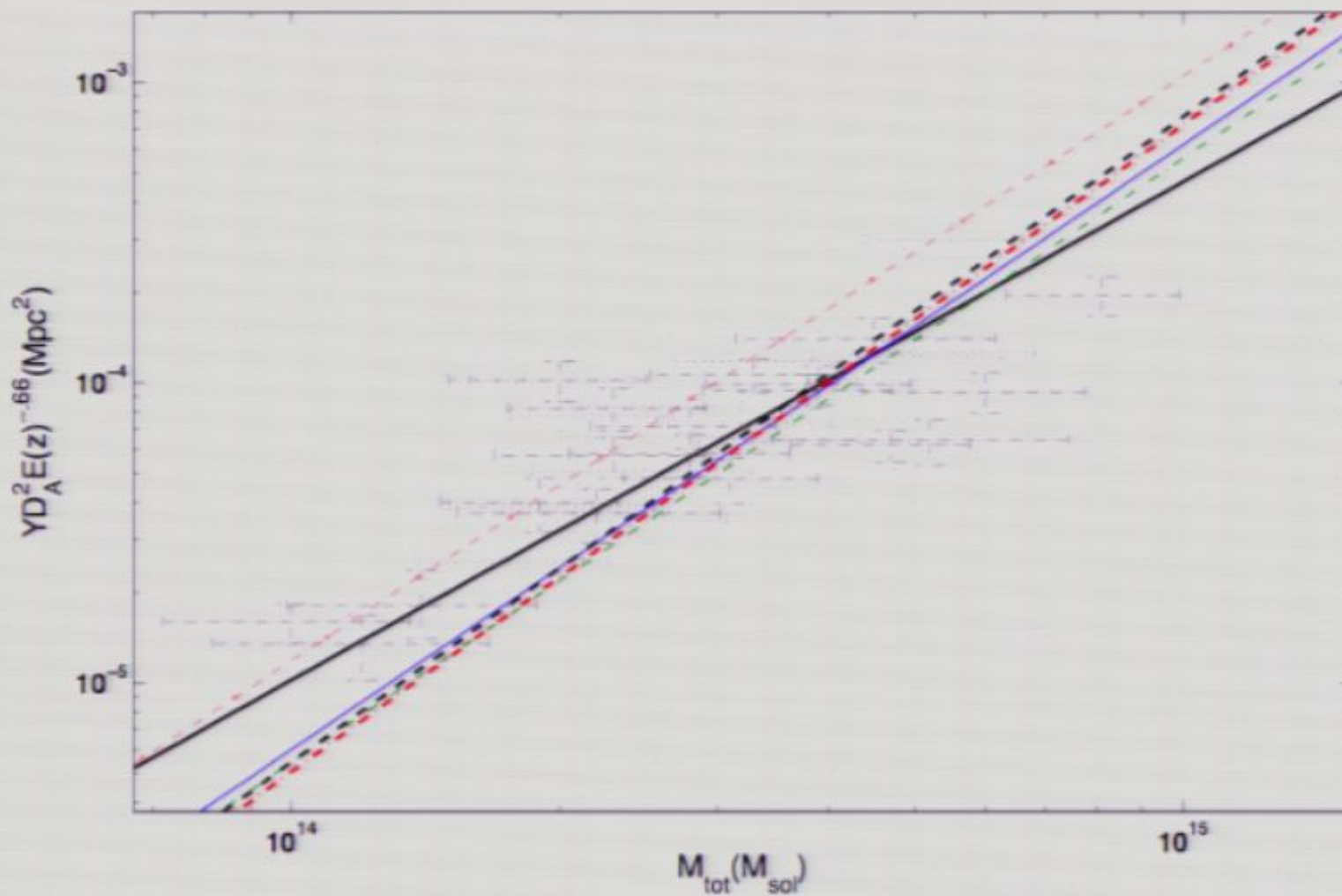
	A	B	$\chi^2$	A	B	$\chi^2$
Bonamente			.9252			.6713
model 1	-25.55	1.5794	1.0343			.5754
model 2	-25.51	1.5781	.9274			.4987
model 3	-25.0082	1.5357	1.3435			.7916
model 4	-24.98	1.5321	1.5741			.9812
thermal 1	-25.1567	1.5405	2.53			1.48
thermal 2	-25.577	1.5725	2.28			1.31
K S	23.0122	2.1758	10.50			5.2206



**SZFlux -  $M_{\text{tot}}$**

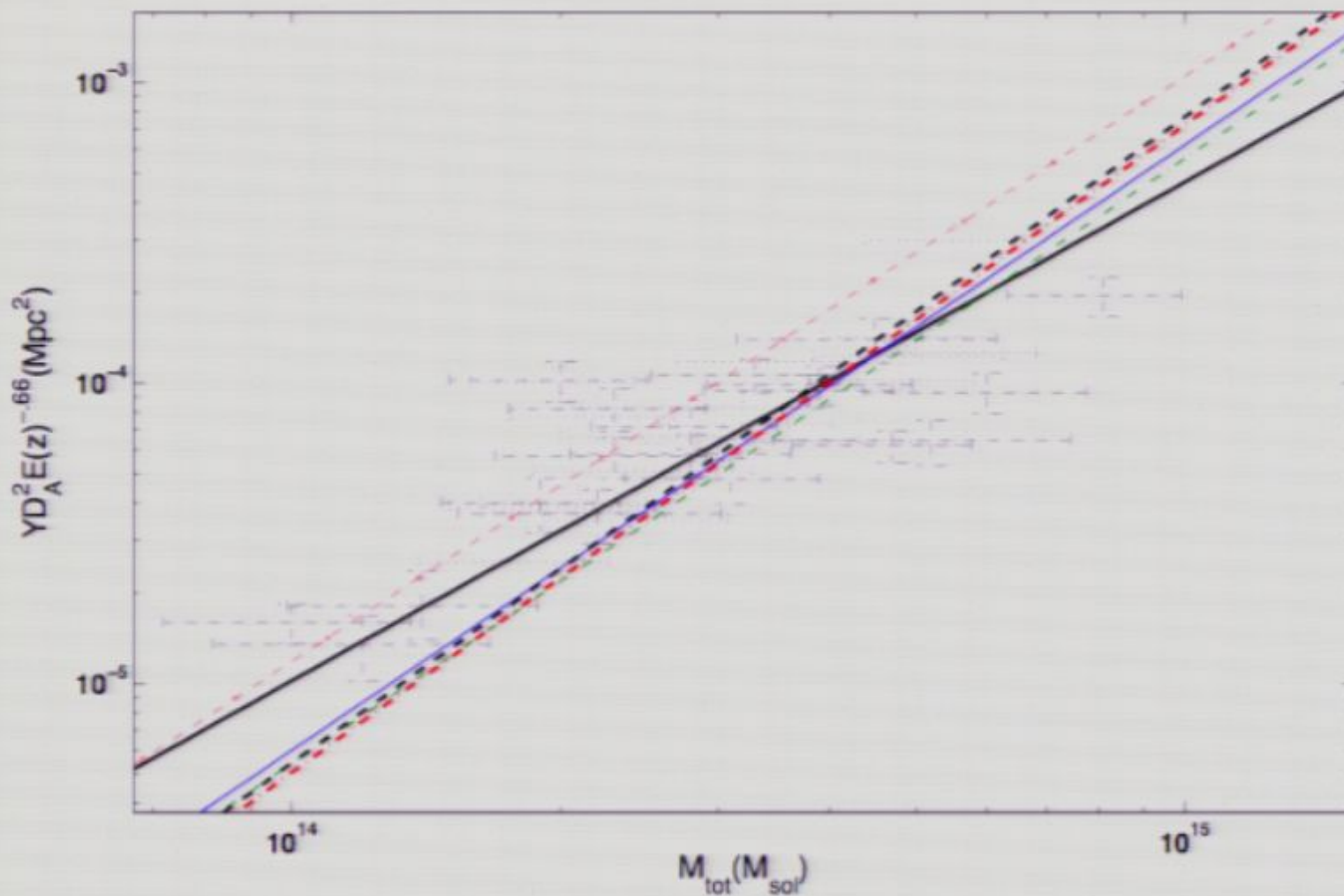


**SZ Flux -  $M_{\text{tot}}$**



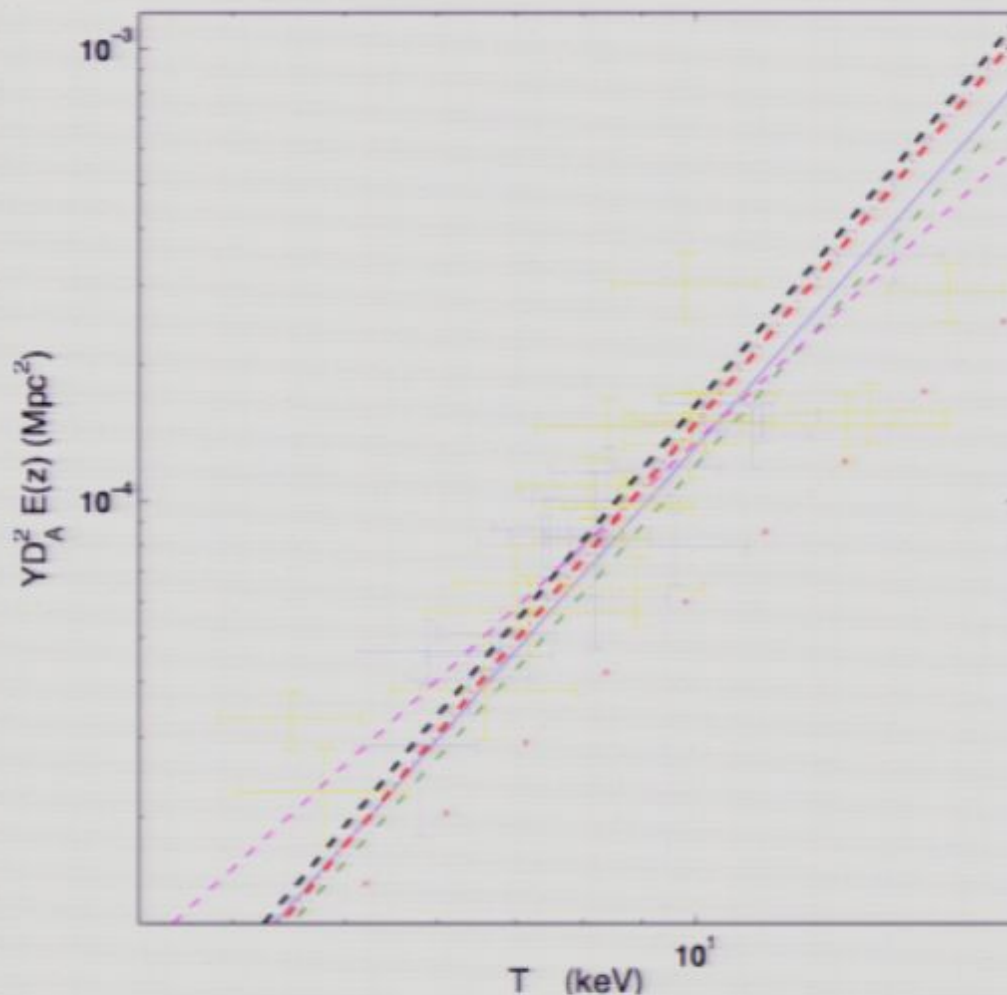
**SZ Flux -  $M_{\text{tot}}$**





**SZFlux -  $M_{\text{tot}}$**

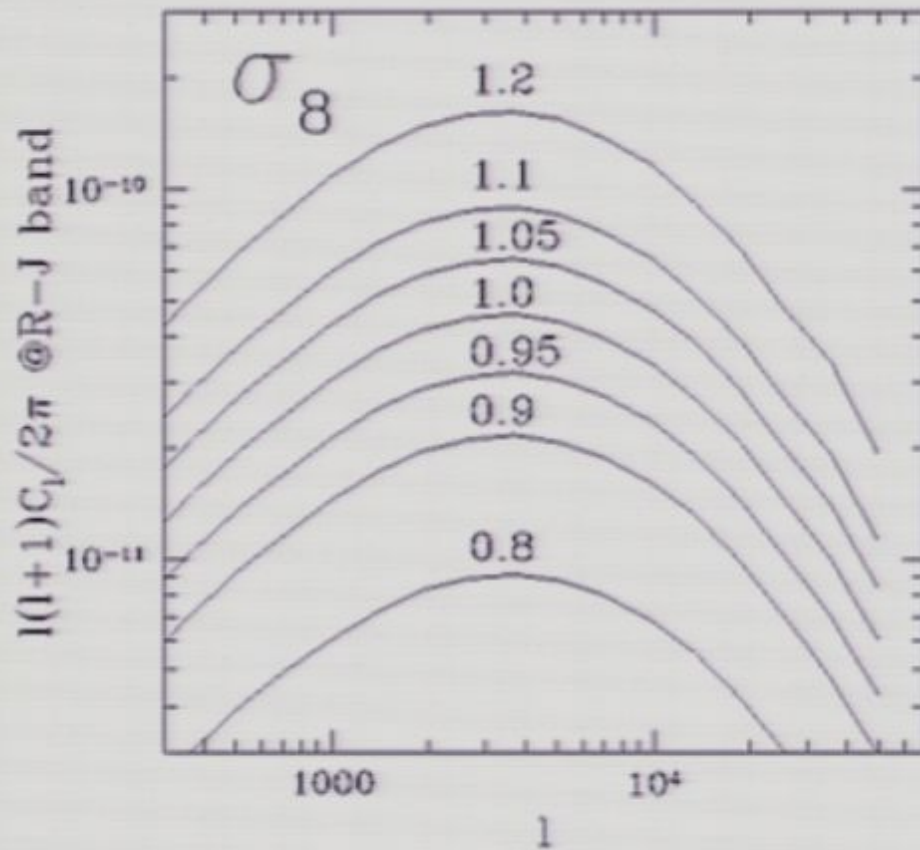
	A	B	$\chi^2_1$	A	B	$\chi^2_2$
Bonamente			1.5714			.7214
model 1	-33.43	2.015	1.8299			.8878
model 2	-33.45	2.013	2.0959			1.0584
model 3	-35.46	2.1546	1.8868			.8945
model 4	-35.37	2.1504	1.7617			.8255
thermal 1	-35.1914	2.1431	1.66			.7303
thermal 2	-34.05	2.0626	1.64			.74
K-S	-32.2853	1.9549	2.5651			1.1946



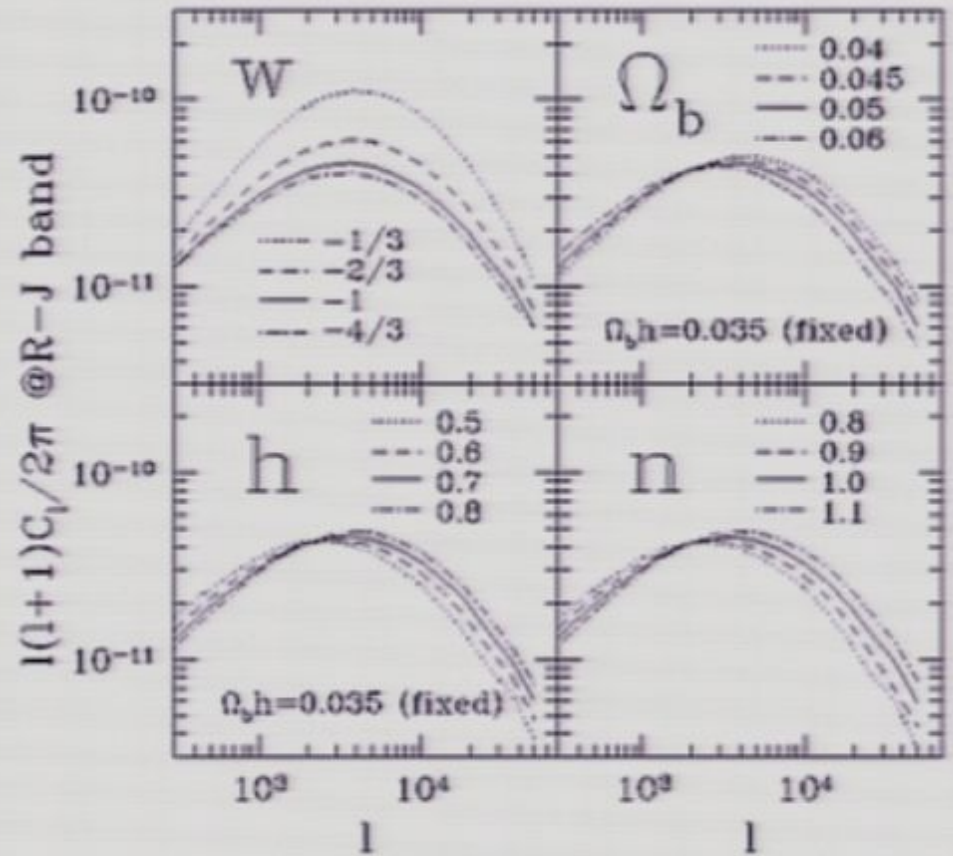
**SZFlux - Temp**

	T (keV)					
	A	B	$\chi^2_1$	A	B	$\chi^2_2$
Bonamente			.8362			.2034
model 1	-6.7967	2.9158	1.0769			.2969
model 2	-6.8358	2.9134	1.3164			.4582
model 3	-6.9282	3.0997	1.0768			.2903
model 4	-6.8790	3.0895	1.1217			.3172
thermal 1	-6.7716	3.0696	1.68			.64
thermal 2	-6.702	2.954	1.31			.4420
K S	6.47	2.98	7.47			2.7200

## On To SZ Power Spectrum -



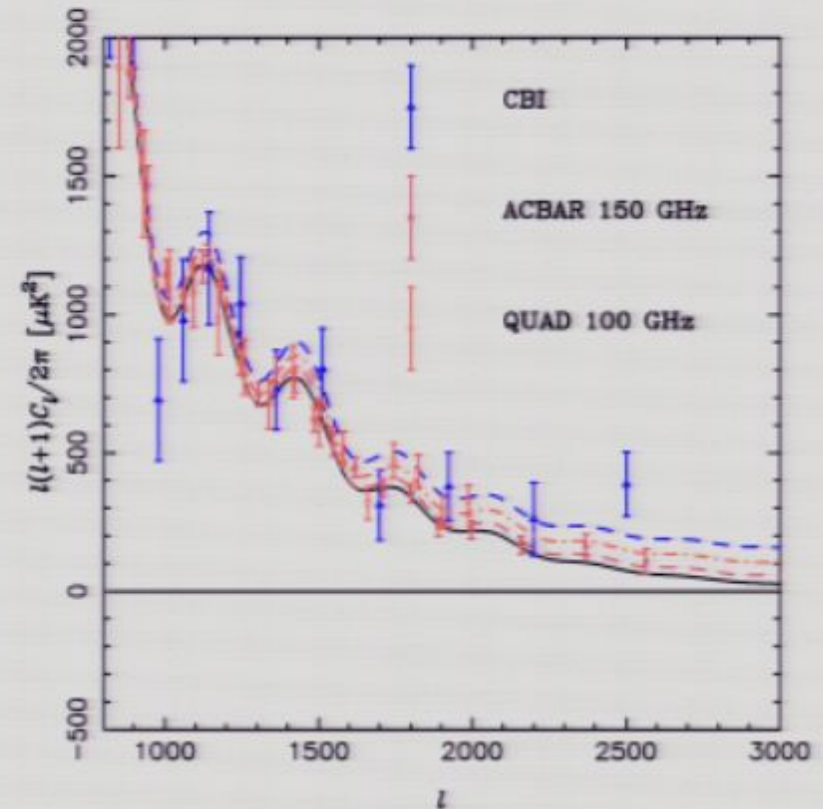
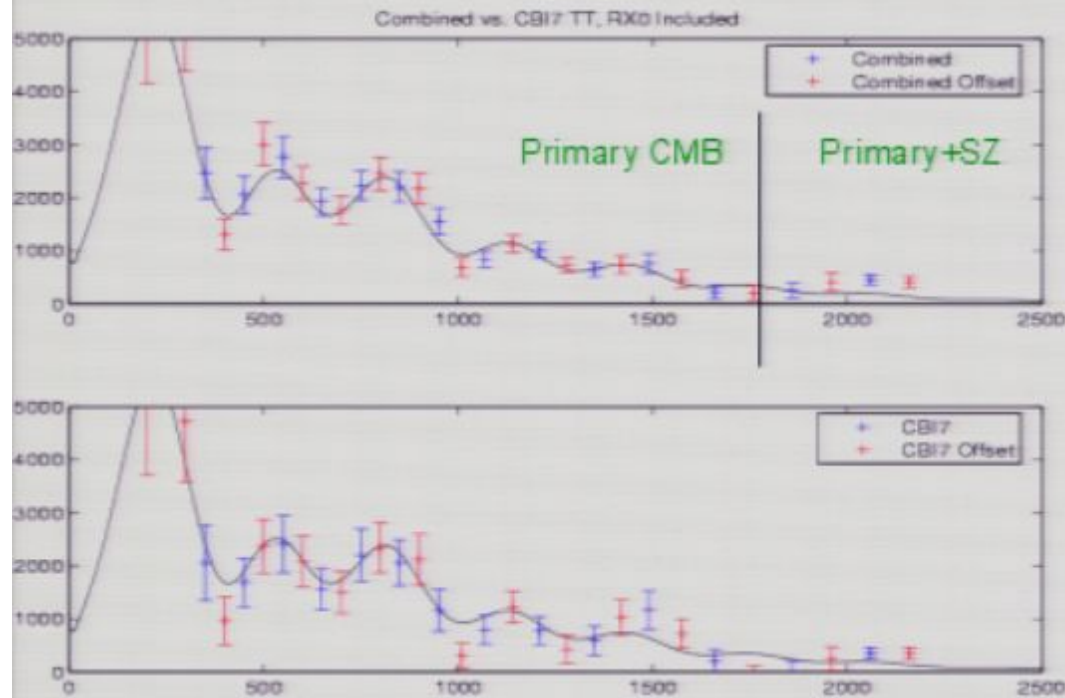
$$C_l \sim \sigma_8^{6-7}$$



Effect of other parameters are subdued under the effect of  $\sigma_8$ .



## EXCESS OF CMB ANISOTROPY OBSERVED AT SMALL ANGULAR SCALES -



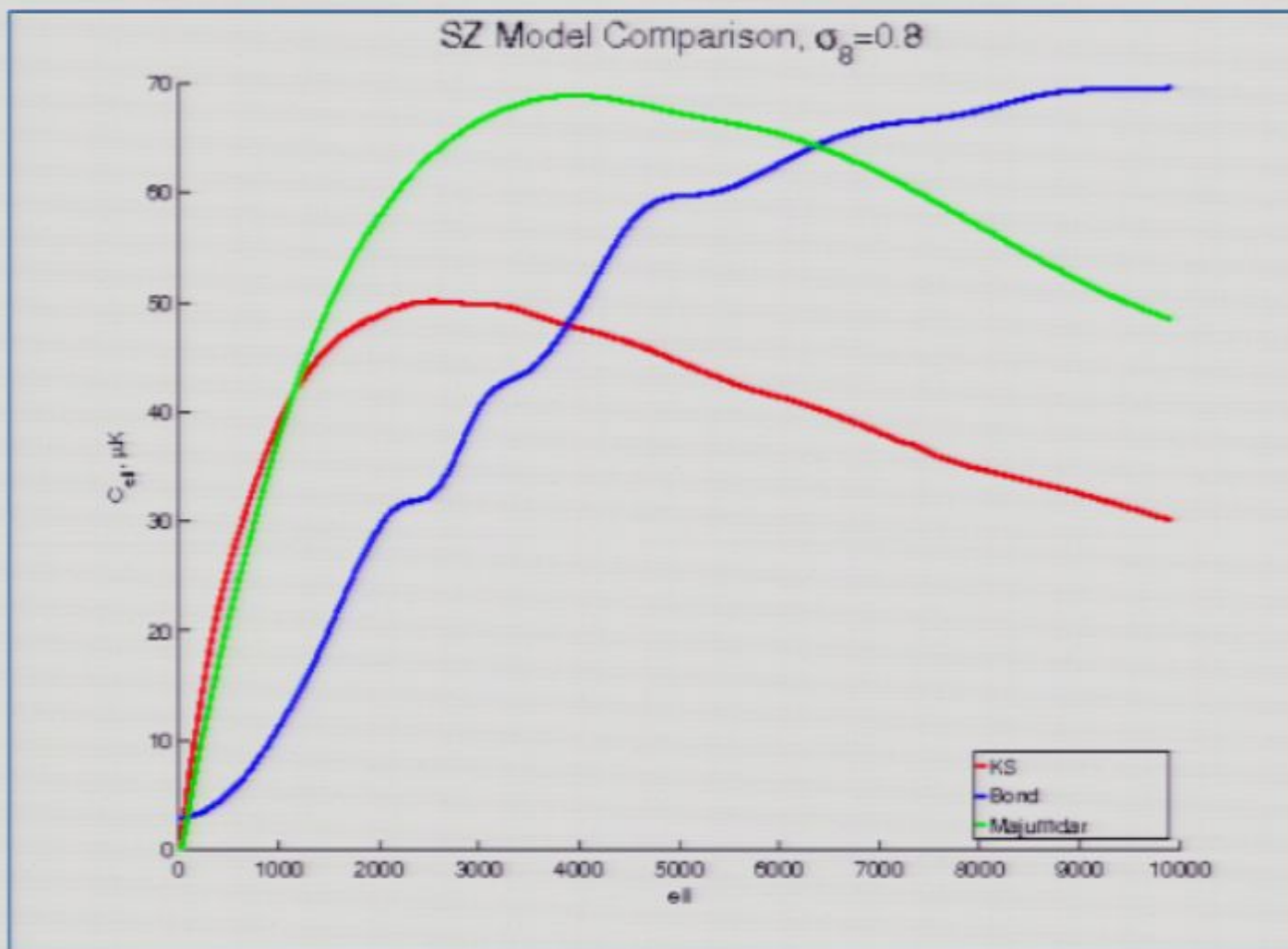
Sievers et al (2009) [and previously Bond et al (2006)] analyzed CBI data + radio sources at 30 GHz and found “excess” CMB anisotropy at  $l > 1500$ . Moreover, ACBAR & BIMA data also agree with CBI excess, if excess has SZ freq dependence.

They use two different SZ cluster templates: 1) Analytic Komatsu Seljak and 2) SPH simulation. They find  $\sigma_8 = 0.922 \pm 0.047$  and  $\sigma_8 = 0.988 \pm 0.049$ . Previous result was  $\sigma_8 = 1.0 - 1.05$ .

SZ CI WITH NEW SZ TEMPLATE AND  $\sigma_8$  -

More in Jon Siever's talk

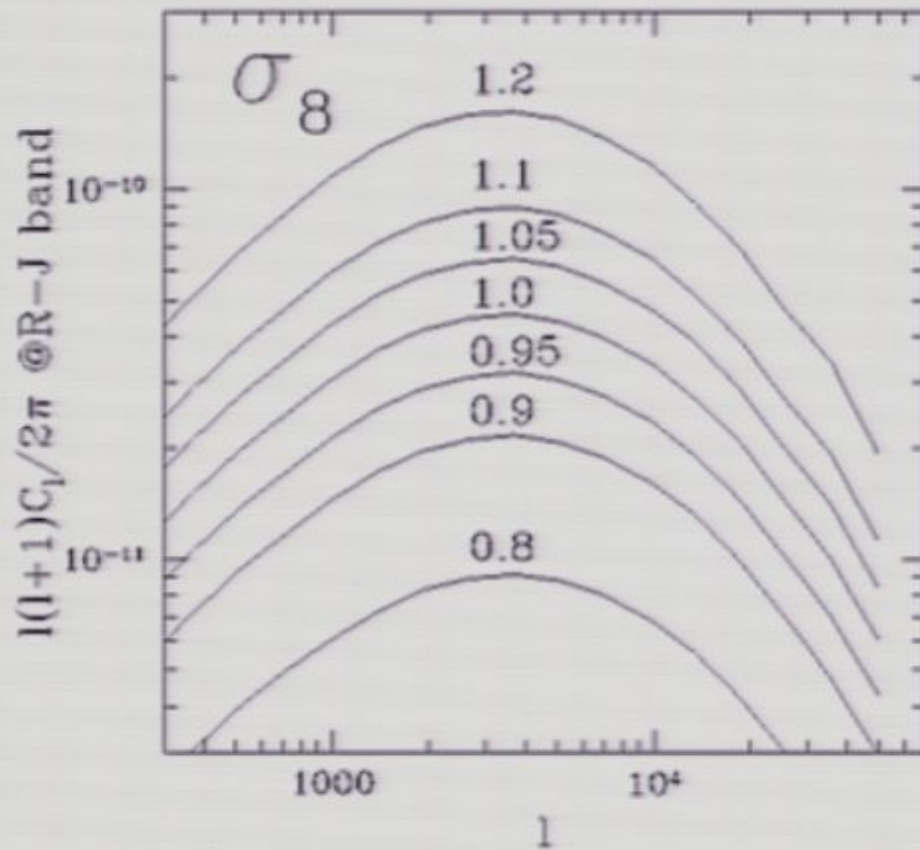
## SZ CI WITH NEW SZ TEMPLATE AND $\sigma_8$ -



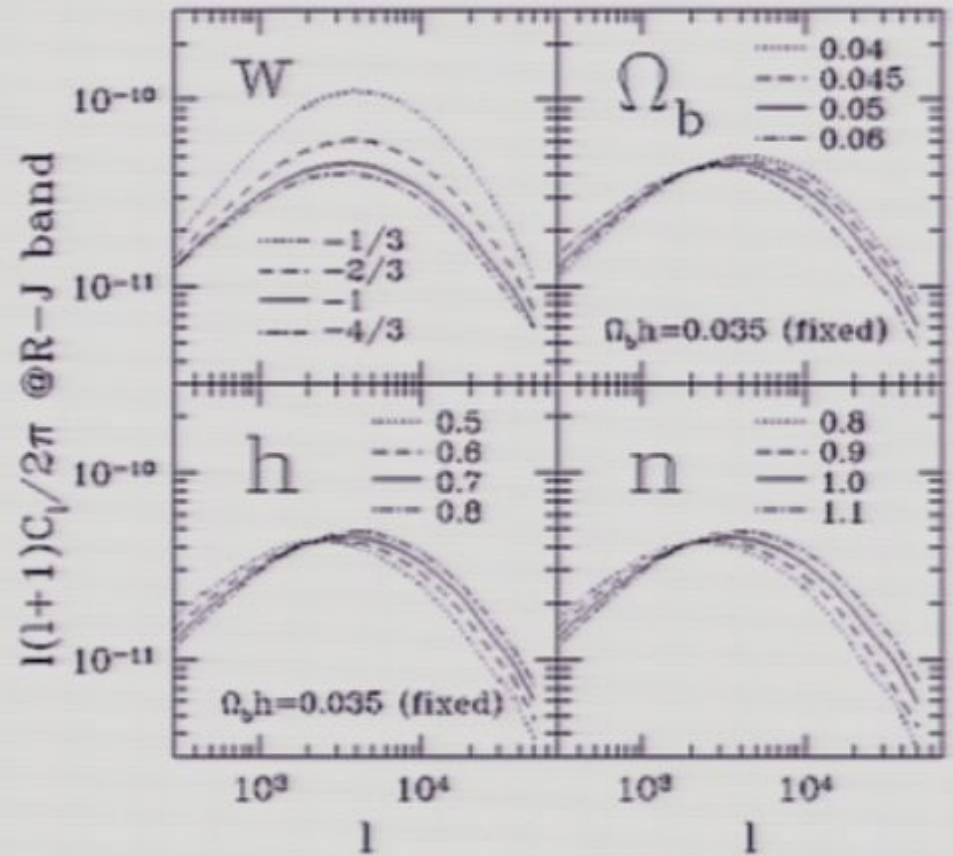
More in Jon Siever's talk



## On To SZ Power Spectrum -

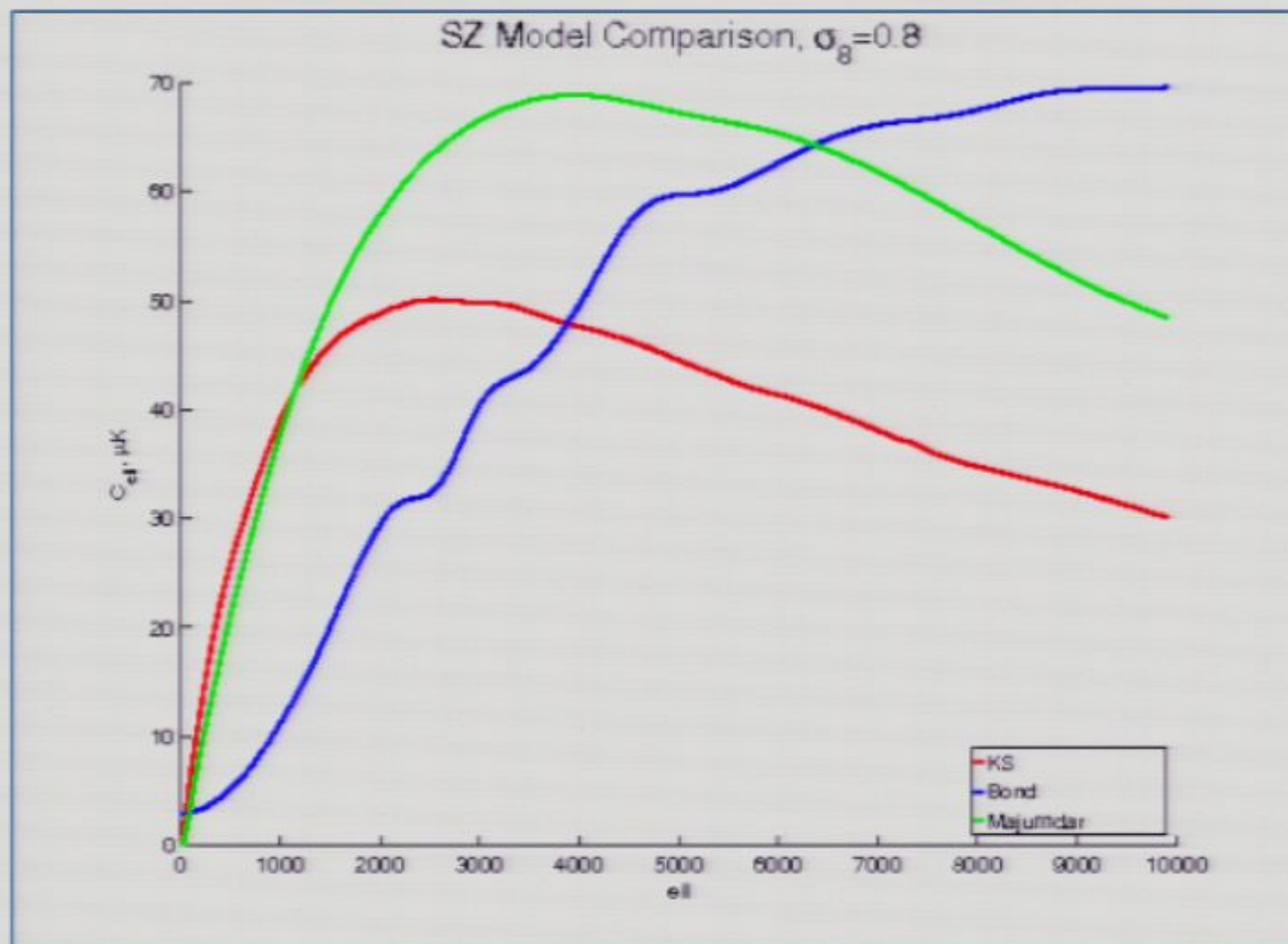


$$C_l \sim \sigma_8^{6-7}$$



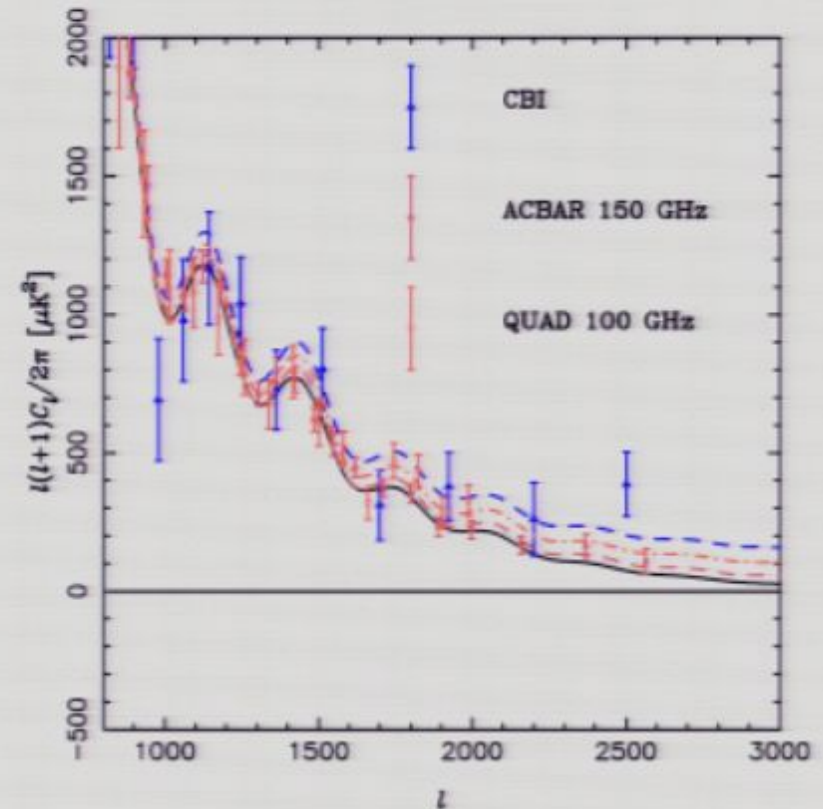
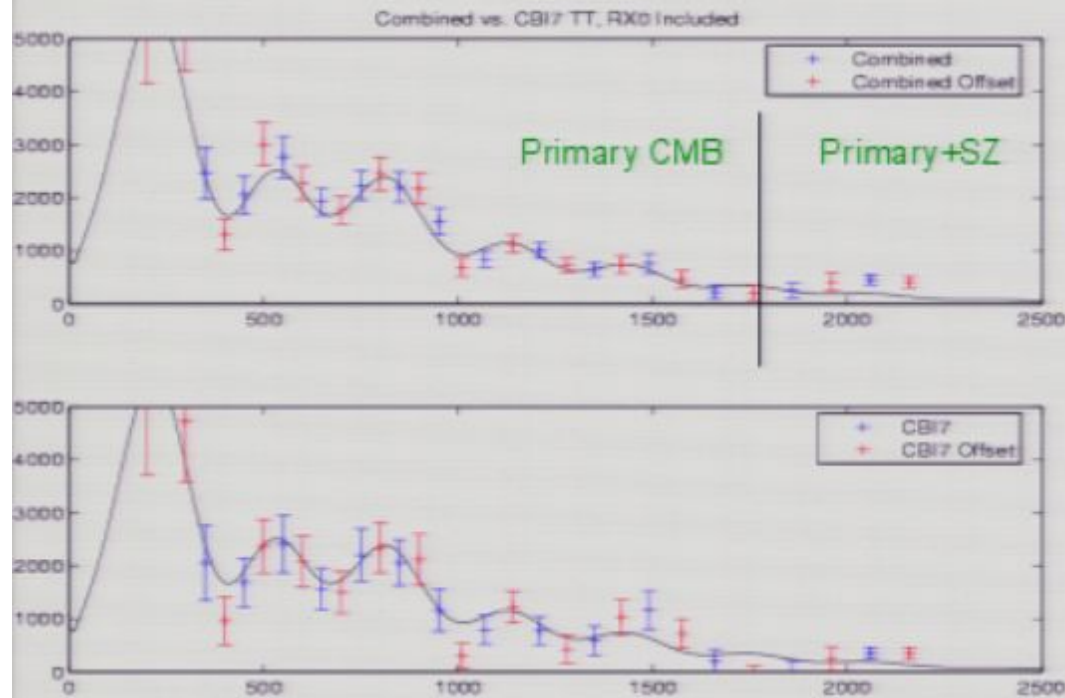
Effect of other parameters are subdued under the effect of  $\sigma_8$ .

## SZ CI WITH NEW SZ TEMPLATE AND $\sigma_8$ -



More in Jon Siever's talk

## EXCESS OF CMB ANISOTROPY OBSERVED AT SMALL ANGULAR SCALES -



Sievers et al (2009) [and previously Bond et al (2006)] analyzed CBI data + radio sources at 30 GHz and found “excess” CMB anisotropy at  $l > 1500$ . Moreover, ACBAR & BIMA data also agree with CBI excess, if excess has SZ freq dependence.

They use two different SZ cluster templates: 1) Analytic Komatsu Seljak and 2) SPH simulation. They find  $\sigma_8 = 0.922 \pm 0.047$  and  $\sigma_8 = 0.988 \pm 0.049$ . Previous result was  $\sigma_8 = 1.0 - 1.05$ .

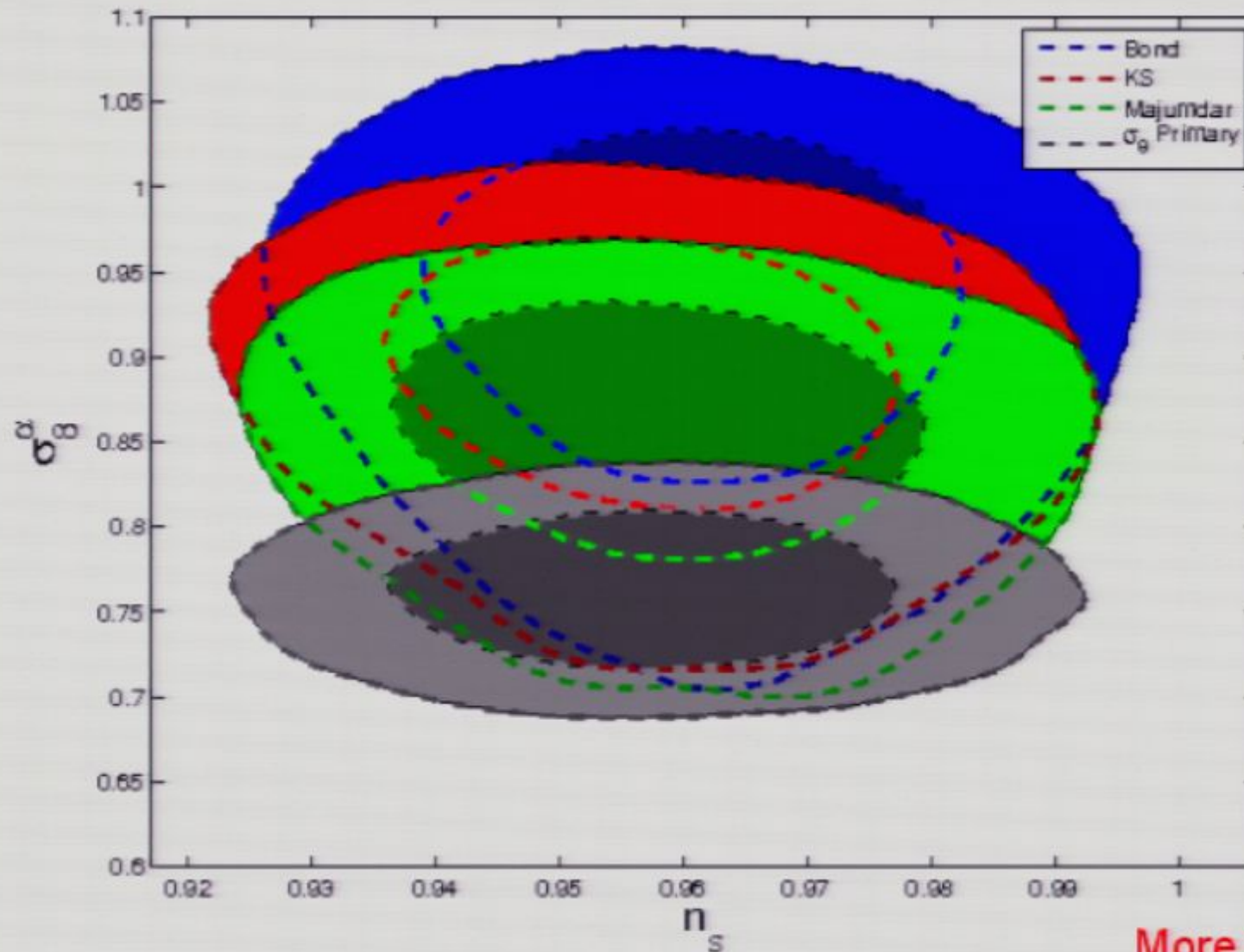


SZ CI WITH NEW SZ TEMPLATE AND  $\sigma_8$  -

More in Jon Siever's talk

## SZ CI WITH NEW SZ TEMPLATE AND $\sigma_8$ -

$\sigma_8 = 0.851 \pm 0.055$  w/out BIMA (including quad 2yr, acbar, cbi, sza, wmap5)

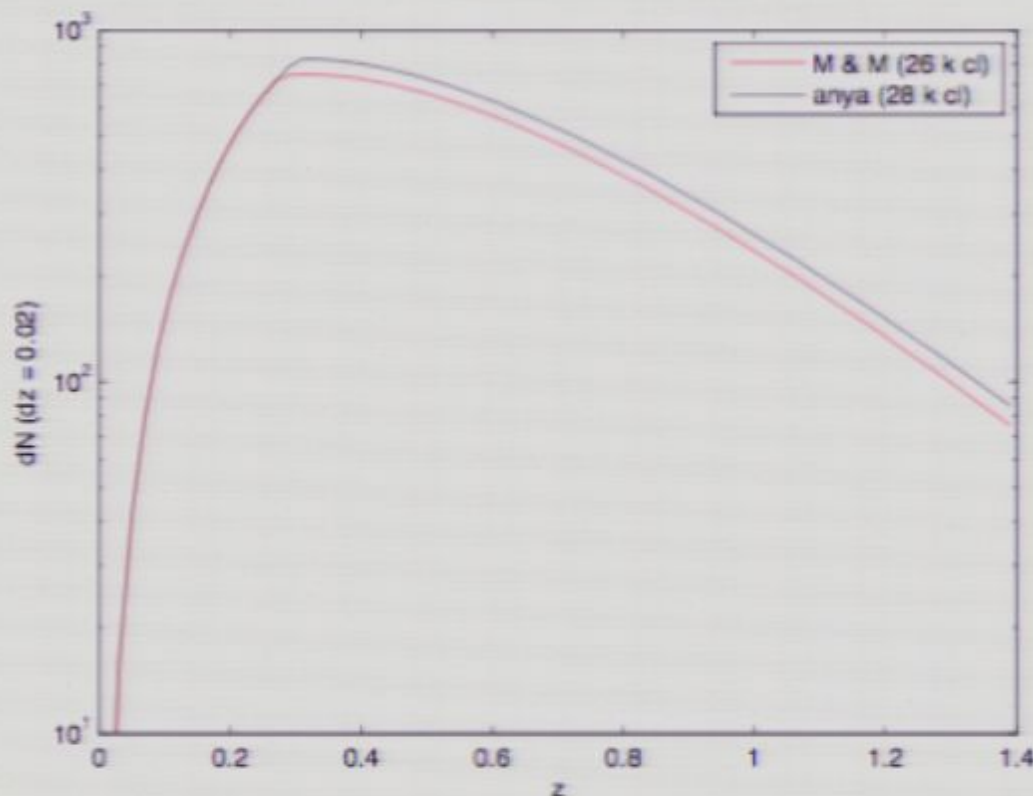


More in Jon Siever's talk

## SZ Flux - M200 relation and number counts -

$Y_{200} - M_{tot}$  SCALING RELATION(NON THERMAL-THERMAL)

	$A_{200}$	$B_{200}$
model 1(fid)	-28.0669	1.6123
non th 7	-28.1709	1.6197
non th 4	-28.1946	1.6216
non th 1	-28.2073	1.6221
model 1 (thermal)	-28.2302	1.6247



Imp if you want to do  
'precision cosmology'

More in  
Christoph's talk



## SUMMARY AND CONCLUSIONS -

**We have build a 'bottom-up' model of cluster gas density and temperature profiles, normalized so as to give observed Xray mass-temp scaling relation.** The resultant pressure profiles have been used to construct SZ cluster templates.

The non self-similar Xray scaling results in non self-similar SZ scaling relations. **We compare our results with the recently observed SZ scaling relations (Y-T, Y-Mgas, Y-Mtot) by Bonamente et al and find excellent agreement.**

There are two imp implications of this result -

1) We can say that we are observing the same 'family' of clusters in XRay and SZ, thus making modeling of SZ selection of cluster for surveys easier.

2) We can more confidently predict Y-M200 scaling reln used for cluster surveys. This is important if we want to do so called 'precision cosmology'

**We use the SZ templates to calculate the expected SZ power spectrum as arc-min scales. We compare our results to CBI observed anisotropy to constrain  $\gamma_8$ .**

**Our best fit  $\gamma_8 \sim 0.84$ . This is within the 1-sigma error bar of the best fit  $\gamma_8 \sim 0.817$  from WMAP. In contrast, previous and recent studies, using other cluster templates have got  $\gamma_8 \sim 0.93 - 1.0$  (i.e 4-7 sigma away from WMAP value)**

No Signal

VGA-1

No Signal

VGA-1