Title: The imprint of supernovae and AGN feedback on the SZ sky

Date: Apr 27, 2009 09:45 AM

URL: http://pirsa.org/09040032

Abstract: I discuss whether supernovae at high-redshift can result in a detectable SZ signal at small angular scales. I also discuss various aspects of

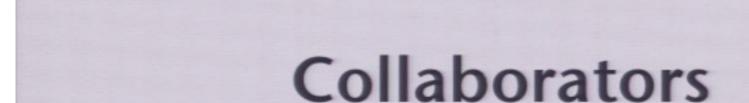
AGN feedback on galaxy clusters.



AGN Feedback Heating in Galaxy Clusters

Peng Oh, UCSB

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Fulaí Guo (UCSB grad student, soon UC Santa Cruz postdoc)

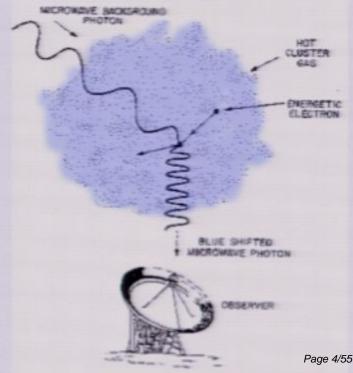
Mateusz Ruszkowskí (Míchigan)

Papers:

Guo g Oh, 2008, MNRAS, 384, 251 Guo, Oh g Ruszkowski, 2008, ApJ, 688, 859 Pirsa: 09040032

For cosmology, we'd like clusters to be nice spherical cows...





But they often contain a

beating heart...

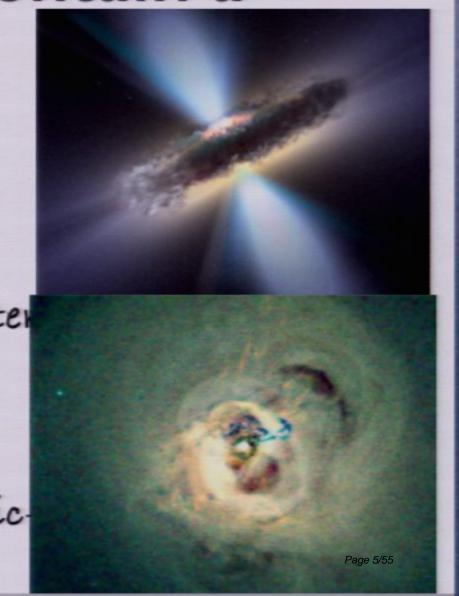
which could affect:

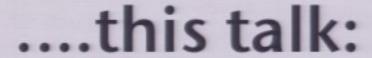
and pressure profiles

O SZ decrement, y-parameter

O Gas fractions

□ Non-thermal pressure support: turbulence, cosmic-





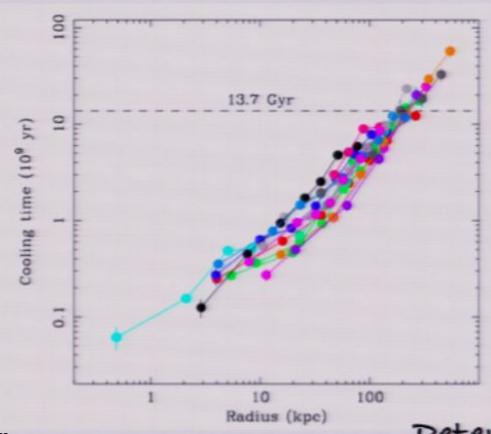
- ☐ A model of AGN heating: cosmic ray heating (but see C. Pfrommer's talk)
- A possible explanation for the dichotomy between cool core and non-cool core clusters (many speakers today, incld Babul, Bode, Nagai)



Cosmic-Ray Heating

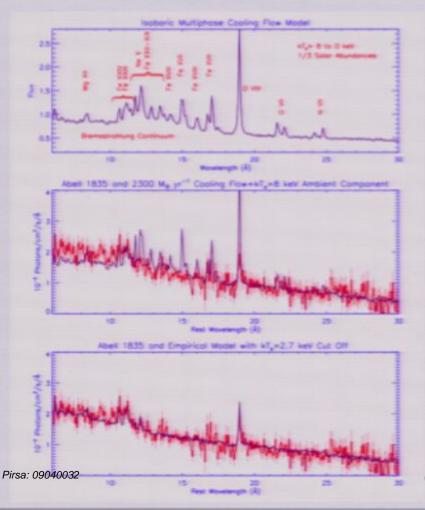
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Although gas cooling times in clusters are often short...



Peterson & Fabian (2006)

Gas does not appear to cool below ~1/3 of T_vir



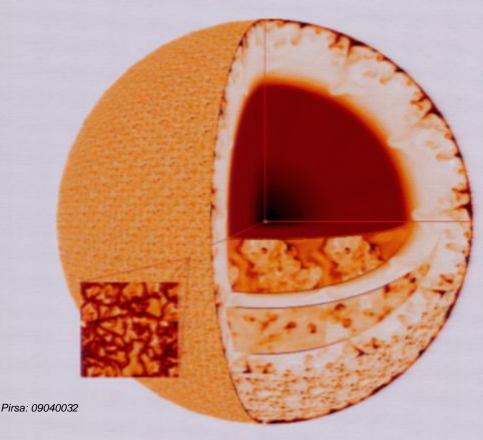
can only fit spectra if prevent gas from cooling below ~1/3 of ambient temperature

universal across different cluster temperatures

Peterson et al (2001)

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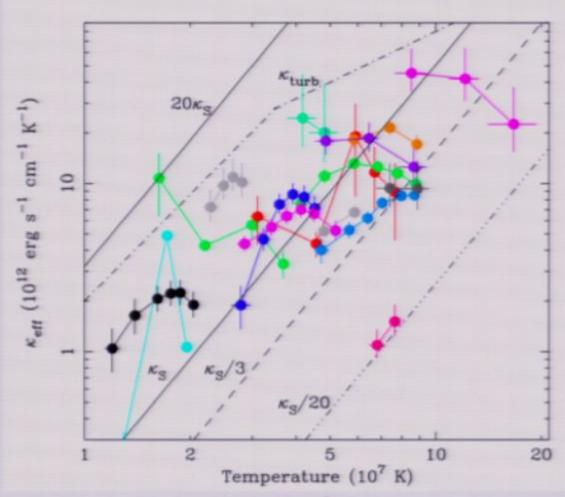
Heating is an attractive solution



Cluster sits in quasithermal equilibrium: just like a star!

Also explains lack of cold gas/stars

Thermal Conduction

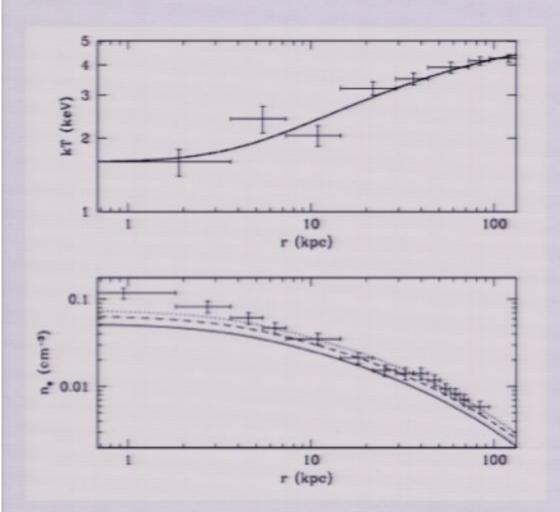


Conduction at fraction of classical Spitzer value close to what's needed.

Coincidence??

Pirsa: 09040002 eterson & Fabian (2006)

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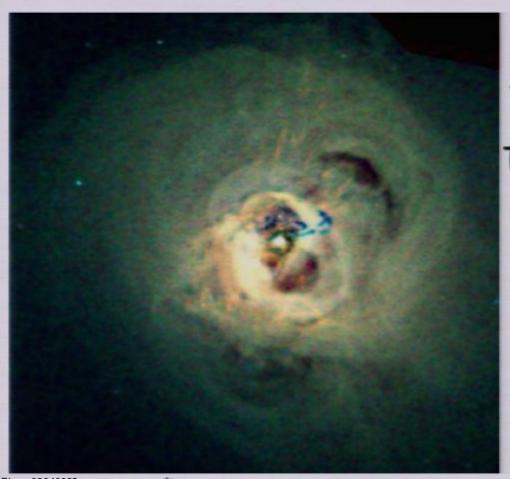


can build conduction-only models in hydrostatic and thermal equilibrium

But: suffer finetuning problems, tend to be globally unstable

Zakamska & Narayan (2002)

AGN/radio galaxy heating



Bubbles observed in ICM, filled with hot/relativistic plasma

Maybe: entrain cold gas pdv work

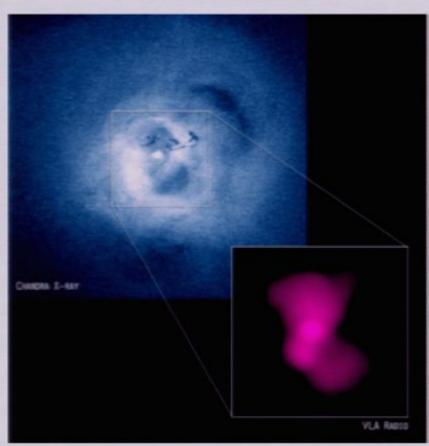
This talk: cosmic ray

heating

(quo g oh 2008)

Chandra image, Perseus cluster

Why cosmic rays?



We see radio synchrotron emission

Many sources: jets, accretion shock, SN Provide gentle, distributed heating



- □ Authors have considered dynamical and heating effects (via Coulomb, hadronic and Alfven wave interactions) (Boehringer & Morfill 1988, Loewenstein et al 1991, Repaheli & Silk 1995, Colafrancesco et al 2004, Jubelgas et al 2006, Pfrommer et al 2006)
- ☐ But CRS in previous models did not successfully stop cooling flow

A key problem: CR transport is slow

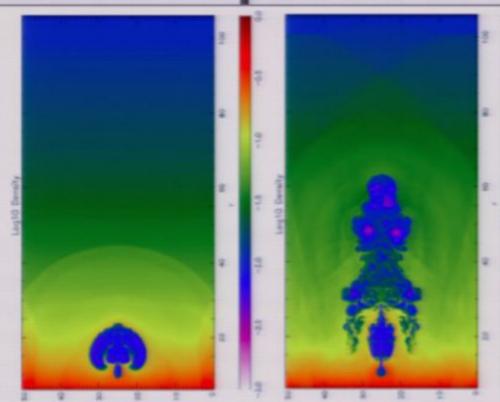
$$\boldsymbol{F}_{c} = \gamma_{c} E_{c} (\boldsymbol{u} + \boldsymbol{v}_{A}) - \boldsymbol{n} \kappa_{c} (\boldsymbol{n} \cdot \boldsymbol{\nabla} E_{c}),$$
 (A14)

$$\frac{\partial E_{c}}{\partial t} = (\gamma_{c} - 1)(\boldsymbol{u} + \boldsymbol{v}_{A}) \cdot \nabla E_{c} - \nabla \cdot \boldsymbol{F}_{c} + \bar{Q}. \quad (A15)$$

Diffusive and other CR transport timescales are long

Leads to overpressured center with insufficient heating at outskirts (though may drive turbulent convection: Chandran & collaborators)

Our model: use bubbles to transport CRs Bubbles disrupted by



Bubbles disrupted by Rayleigh-Taylor & Kelvin-Helmholtz instabilities as rise

(Also: CRs diffuse out)

Fast way of

transporting CRs: rise

time ~ sound crossing

Bruggen & Kaiser (2002)

time

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- D 1D Zeus code: solve time-dependent hydrodynamic equations + CR heating § transport equations
- calculate steady steady CR spectrum, assuming Coulomb, hadronic and Alfven-wave energy loses (latter dominates):

$$\Gamma_{wave} = v_A \frac{dP_c}{dr}$$

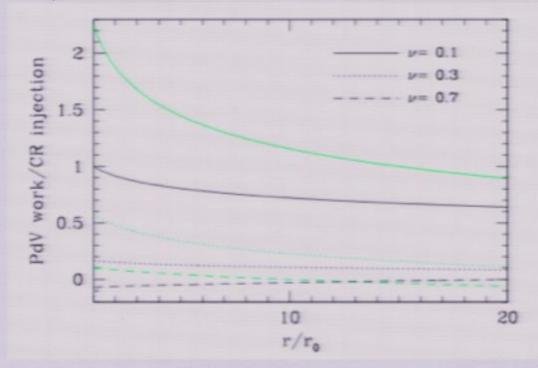
Assume energy density in bubbles is a power-law with radius (note: CR injection rate depends on gas cooling---feedback effect)

$$L_{
m bubble} \sim -\epsilon \dot{M}_{
m in} c^2 \left(rac{r}{r_0}
ight)^{-
u} \quad {
m for} \,\, r > r_0,$$

$$Q_{c} = \nabla \cdot \mathbf{F}_{\text{bubble}} \sim -\frac{1}{4\pi r^{2}} \frac{\partial L_{\text{bubble}}}{\partial r} \left[1 - e^{-(r/r_{0})^{2}} \right]$$
$$\sim -\frac{\nu \epsilon \dot{M}_{\text{in}} c^{2}}{4\pi r_{0}^{3}} \left(\frac{r}{r_{0}} \right)^{-3-\nu} \left[1 - e^{-(r/r_{0})^{2}} \right],$$
(19)

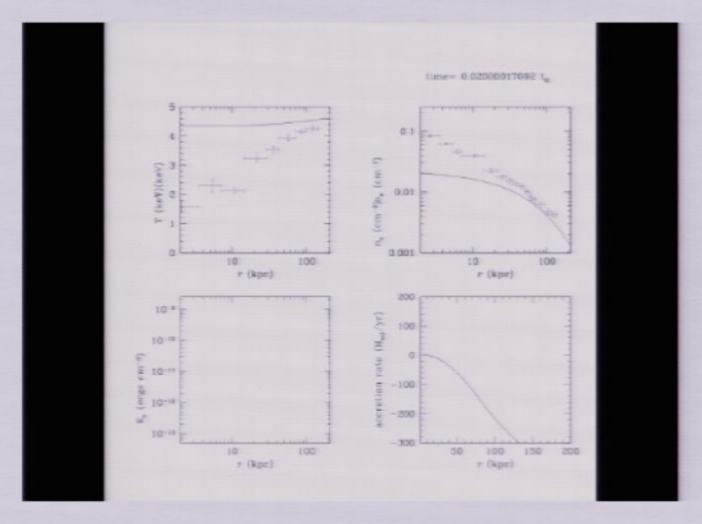
Slope is free parameter, implicitly specifies bubble disruption rate

Amount of energy lost to pdv work is small, at most comparable to the bubble disruption rate

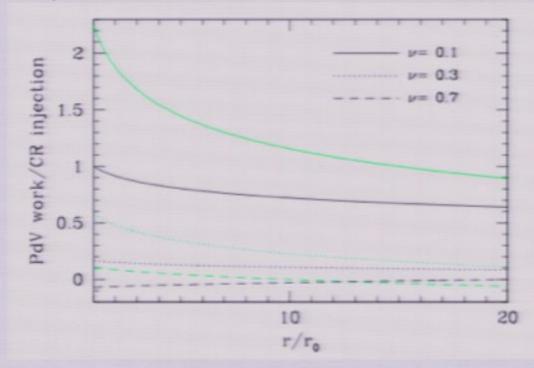


(pdv work can also heat ICM, we ignore it)

Bottom line: it works!

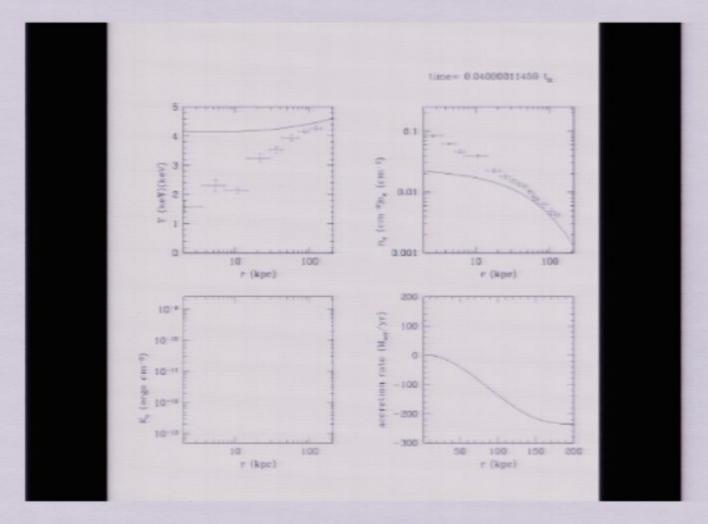


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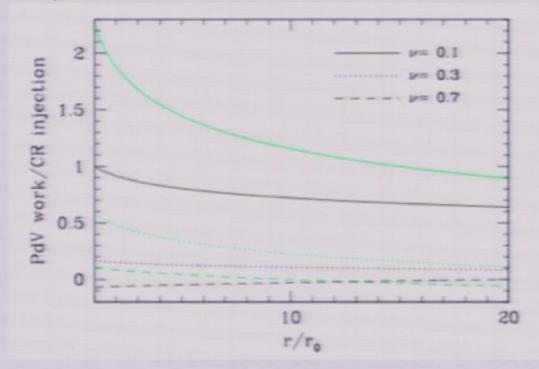


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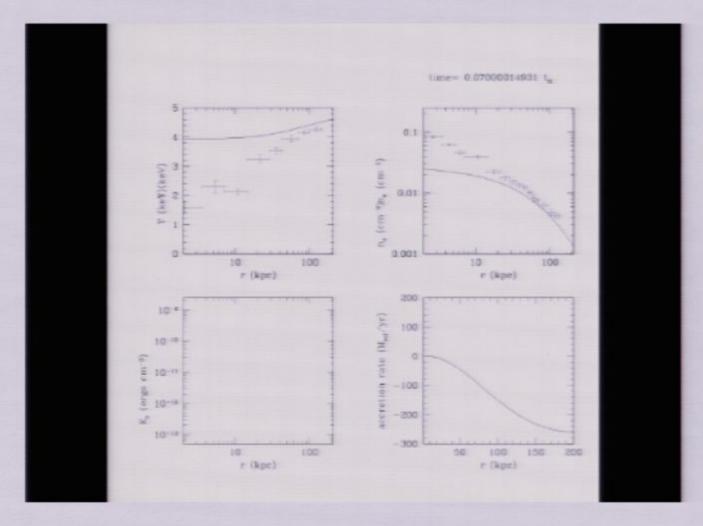


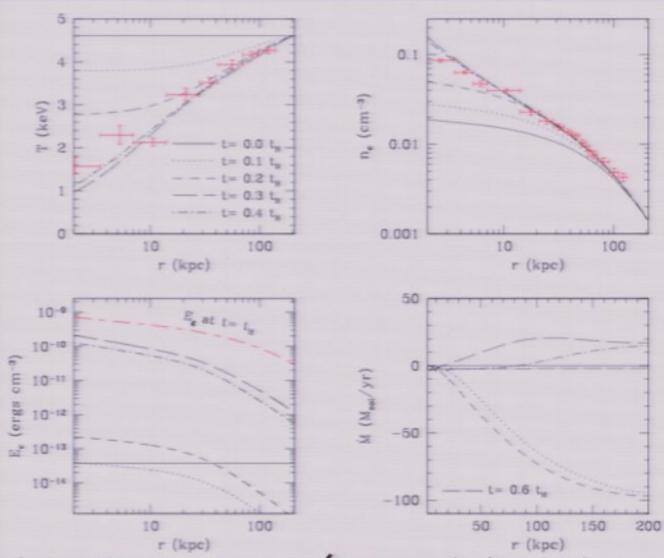
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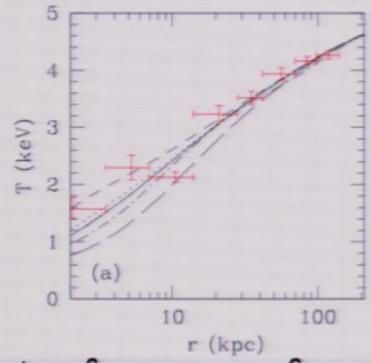
Note that CR pressure is much less than thermal

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No fine tuning

Works (i.e., no massive cooling flow) starting from arbitrary initial conditions (unlike other models...)

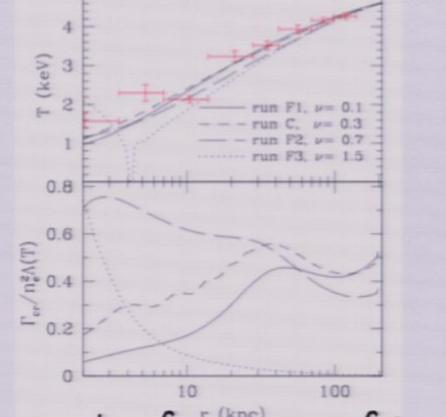


works for range of AGN +

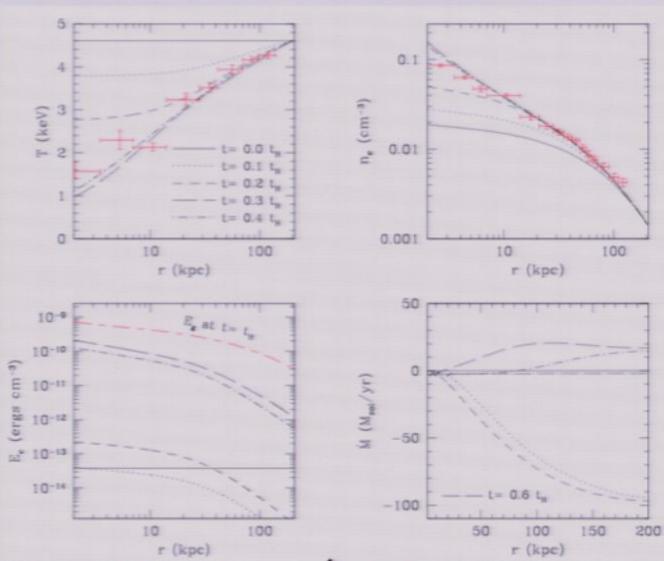
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AGN +

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works for range of CR



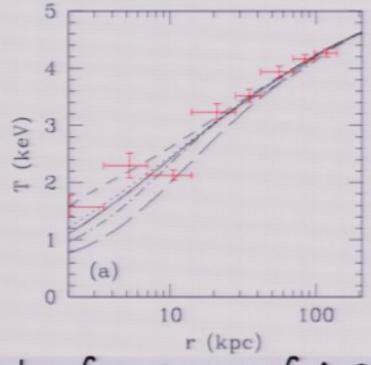
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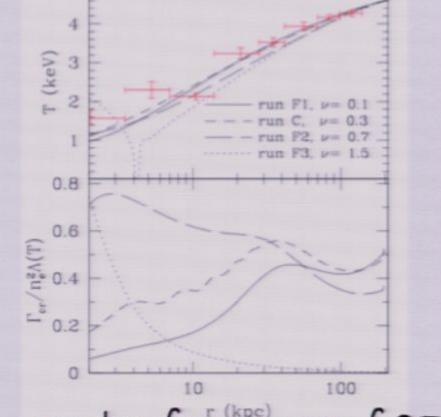


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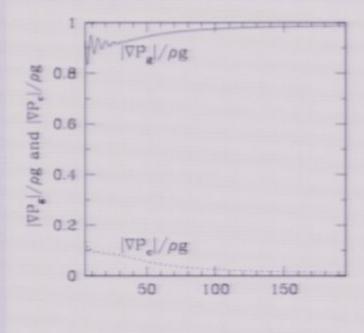
Pirsa: 09040032

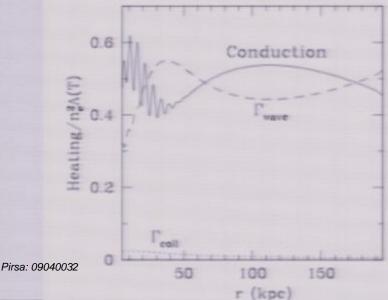
AGN +

Pirsa: 09040032



works for range of CR



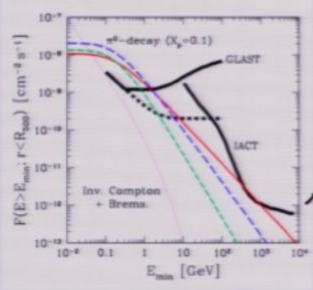


Required CR pressure gradients OK

Small fraction of thermal pressure gradient

most heating is wave heating

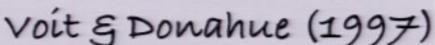
Observational tests

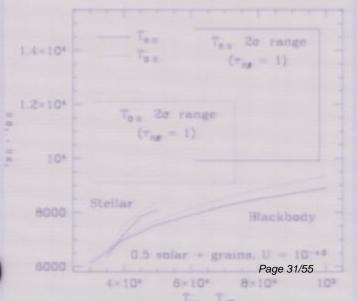


See gamma-rays from pion-decay with GLAST

Ando & Nagai (2007)

optical filaments: need source of anomolous heating?

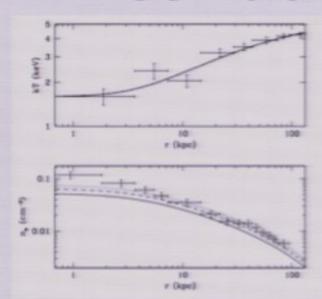






Global Stability in Cool Core vs. Non-Cool Core Clusters

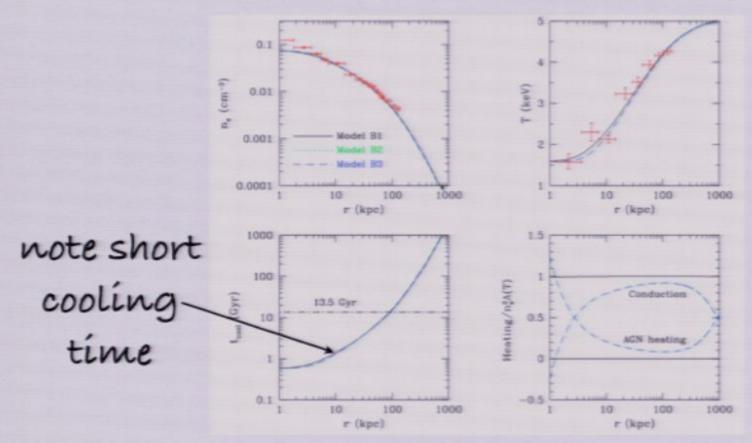
Let's look more closely at fine-tuning issues for conduction models...



can have equilibrium model which fits observations (solve eigenvalue problem)

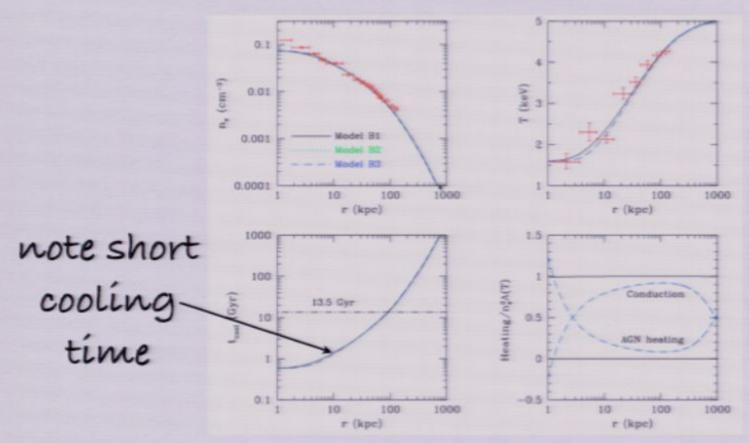
But it won't =0.4.0.6.0.8 evolve toward this state in general...

...look at differences with AGN feedback model



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...look at differences with AGN feedback model



Pirsa: Praíoserst, build a background equilibrium solution solution

...and perform a global stability

WKB

analysis

$$\left(\frac{P}{\rho} - v^2\right) \frac{d}{dr} (\nabla \cdot \boldsymbol{\xi}) = \left(r\sigma^2 + r\frac{d^2\Phi}{dr^2}\right) \frac{\xi}{r} + \frac{1}{\rho} \frac{d}{dr} \left(P\frac{\Delta T}{T}\right)$$

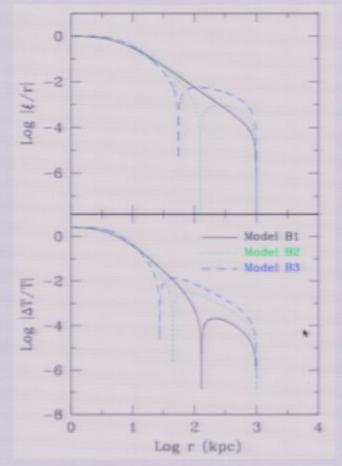
$$-2v^2 \frac{d}{dr} \left(\frac{\xi}{r}\right) + \left(2\sigma v + v\frac{dv}{dr} - \frac{1}{\rho} \frac{dP}{dr}\right) \frac{d\xi}{dr}$$

$$\kappa T \frac{d}{dr} \left(\frac{\Delta T}{T}\right) = F\left[\frac{7}{2} \frac{\Delta T}{T} - r\frac{d}{dr} \left(\frac{\xi}{r}\right) + \frac{\xi}{r}\right] + \frac{\Delta L_{\eta}}{4\pi r^2} 34)$$

$$\frac{1}{4\pi r^2} \frac{d}{dr} \Delta L_r = (P\sigma - \rho^2 \mathcal{L}_{\rho} - \mathcal{H})(\nabla \cdot \boldsymbol{\xi}) - \Delta \mathcal{H}$$

$$+ \left(\frac{P\sigma}{\gamma - 1} + \rho T \mathcal{L}_T + \frac{v}{\gamma - 1} \frac{dP}{dr} - \frac{\gamma v}{\gamma - 1} \frac{P}{\rho} \frac{d\rho}{dr}\right) \frac{\Delta T}{T}$$

$$+ Pv \frac{d}{dr} (\nabla \cdot \boldsymbol{\xi}) + \frac{Pv}{\gamma - 1} \frac{d}{dr} \left(\frac{\Delta T}{T}\right) (35)$$

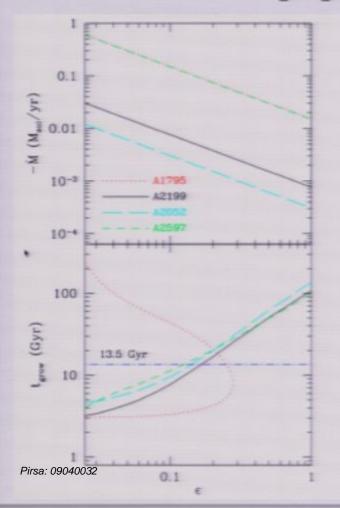


Growth rate is an eigenvalue of analysis Explore parameter space rapidly!

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Global unstable modes suppressed with AGN!



Suppression depends on efficiency

$$L_{\rm agn} = -\epsilon \dot{M}_{\rm in} c^2,$$

The crucial term: feedback

$$\Delta \mathcal{H}_{\text{feed}} \equiv \mathcal{H} \Delta \dot{M}(r_{\text{in}}) / \dot{M}_{\text{in}} = \frac{\mathcal{H} \sigma}{v_0} \xi(r_{\text{in}}) ,$$

TIMESCALES FOR THE CLUSTERS SHOWN IN TABLE II

Name	tcool,0 a (Gyr)	$t_{\infty,0}$ b (Gyr)	ϵ_{\min}^{d}	Model	tgrow c (Gyr)
A1795	0.9	0.6	0.28	A1	3.8
				A2	3.3, 43.3 °
				A3	stable
A2199	0.6	0.4	0.17	B1	2.8
				B2	4.4
				B3	16.9
A2052	1.1	0.7	0.14	CI	6.2
				C2	5.9
				C3	20.0
A2597	0.4	0.3	0.12	D1	1.8
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Observations:

 $\epsilon \sim 0.3$ (Heinz et al 2007)

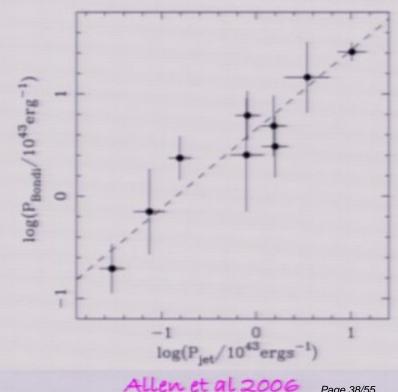
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some observations suggest

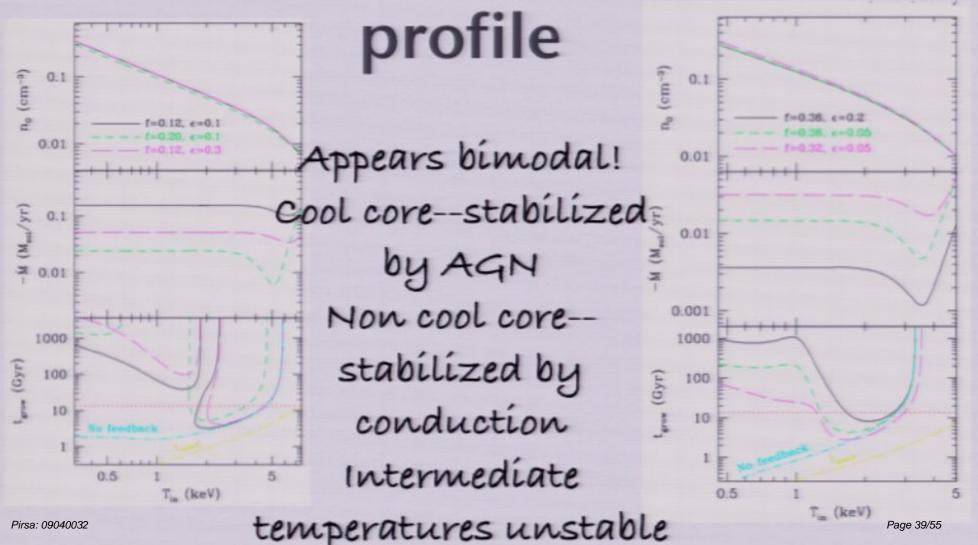
 $\epsilon \propto \dot{M}^{0.3-0.6}$

Stronger feedback,

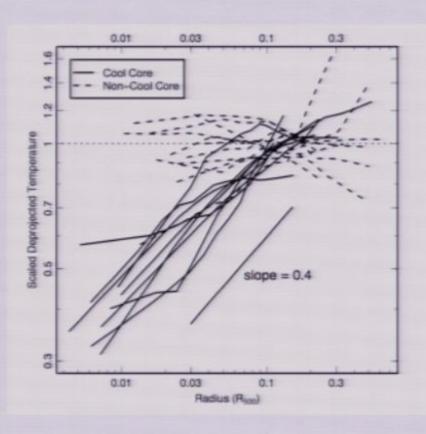
 ϵ_{\min} reduced

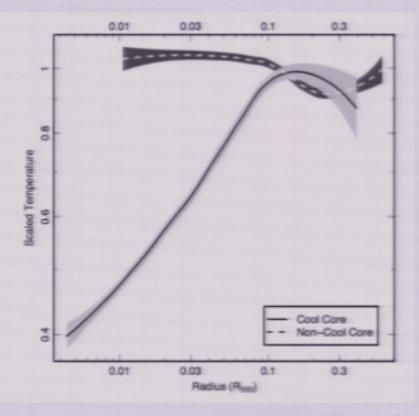


Dependence on background



Consistent with observations



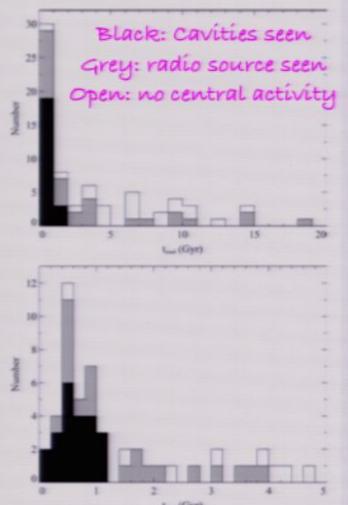


Sanderson et al (2006)

Bimodal cool-core/non-cool core population

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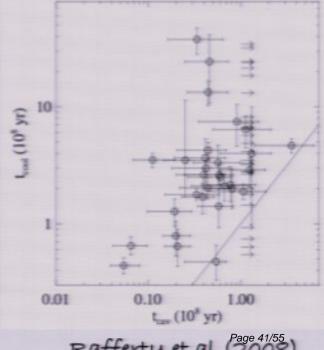
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Dunn & Fabian (2008)

More AGN activity is seen in systems with cool core systems (shorter central cooling time)

Shorter central cooling times correlate with younger cavities

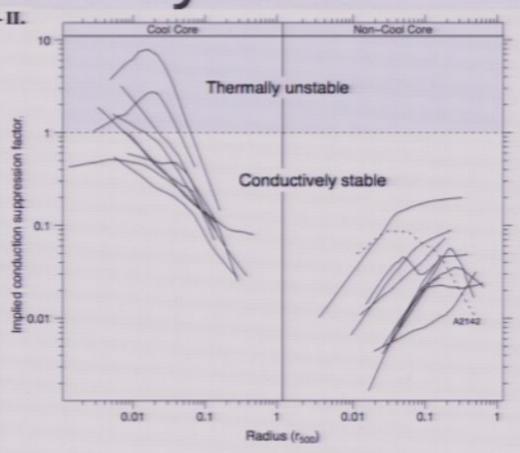


More recently...

A statistically-selected Chandra sample of 20 galaxy clusters – II. Gas properties and cool-core/non-cool core bimodality

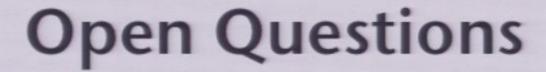
Alastair J. R. Sanderson^{1*}, Ewan O'Sullivan² and Trevor J. Ponman¹

Some more observational support for this picture!

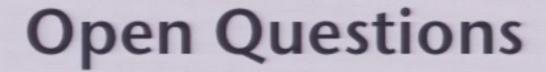


found in CC clusters whereas the flatter slope population are all non-CC clusters. We explore the role of thermal conduction in stabilizing the ICM and conclude that this mechanism alone is sufficient to balance cooling in non-CC clusters. However, CC clusters appear to form a

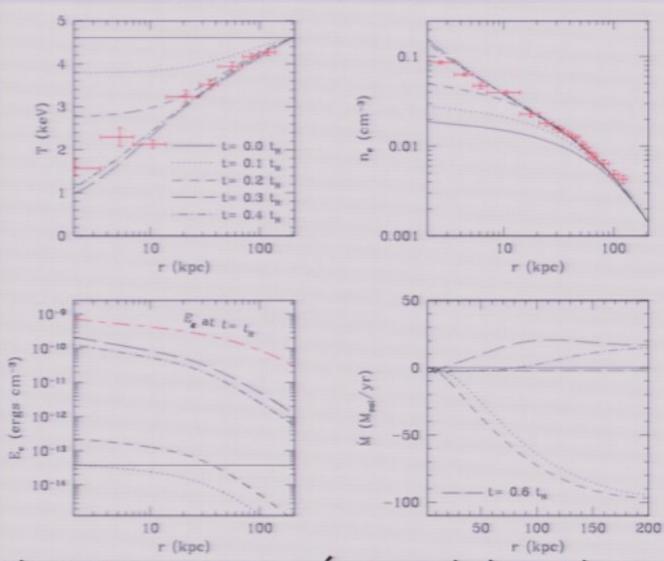
Pirsa: 09040032 stinct population in which heating from feedback is required in addition to conduction.



- What determines the final state a cluster relaxes toward (fastest decaying eigenfunction)?
- □ 3D simulations
- How to get gas to black hole—is Bondi accretion the whole story? (outflows, angular momentum, hot vs. cold accretion)



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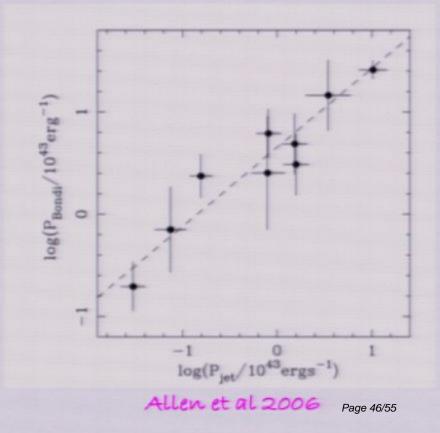


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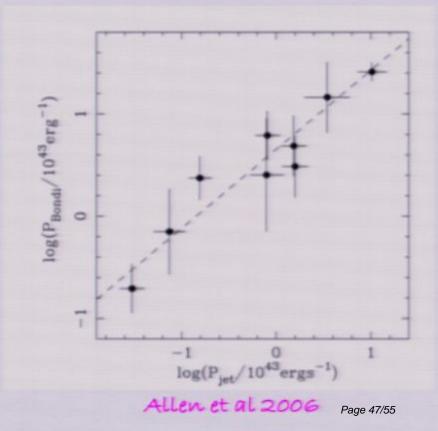
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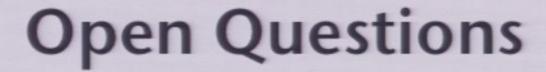
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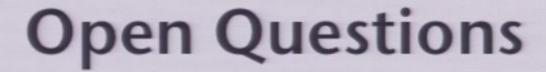
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- Bubble stability: at what rate are bubbles disrupted? 'Magnetic shielding', CR diffusivity, etc....
- □ Topology of magetic field lines: could hot/cold core clusters be two aspects of the same phenomenon, viewed at different times?

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- Cosmic ray heating can be important in clusters—rising bubbles (eventally disrupted) provide a fast means of transporting them
- Global stability analysis fast way of exploring parameter space. Predict: (i) minimal level of heating efficiency, (ii) bimodal central temperatures



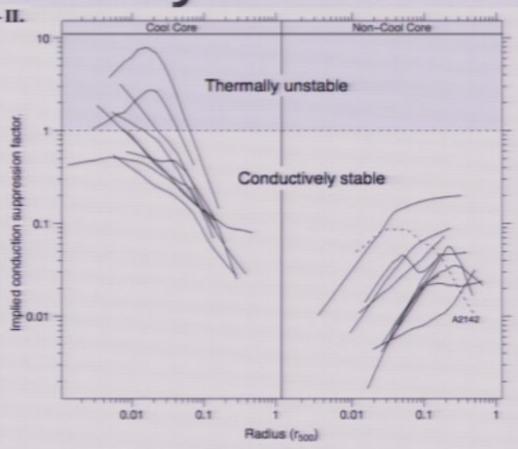
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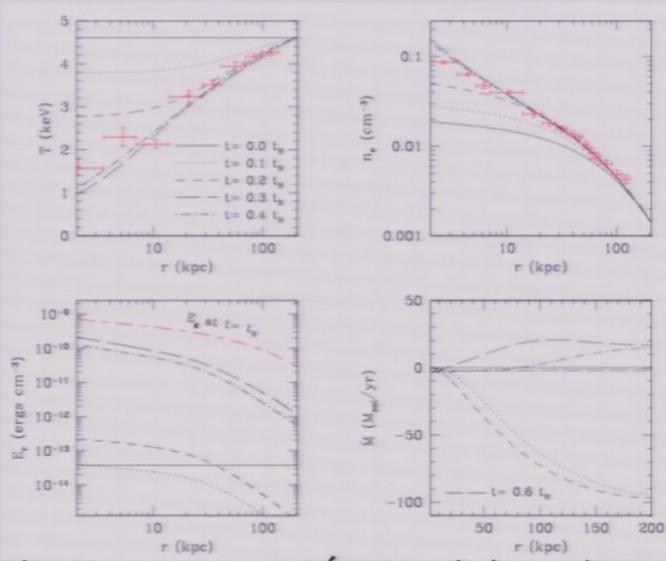
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School of Physics and Autonomy, Ontorropy of Biomorghum, Edghause, directingham 815 207, UK

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O SZ decrement, y-parameter

O Gas fractions

□ Non-thermal pressure support: turbulence, cosmíc

support: turbulence, cosmíd

