

Title: Unconventional Pairing and Impurities in Superfluid Helium-3

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URL: <http://pirsa.org/09040026>

Abstract: The growing fascination with unconventional pairing is driven in part by continuing discoveries of exotic superconductors. The first of these, superfluid ^3He , was found by Osheroff, Richardson, and Lee in 1971. This was followed soon thereafter by superconductivity in the heavy fermion compound, UPt_3 . And then an explosion of interest accompanied the observation of superconductivity in cuprates, Sr_2RuO_4 , and organic materials. The newest discoveries are superconducting compounds of FeAs. These systems have been demonstrated (or in some cases it is just suspected) that they have pairing condensates with non-zero angular momentum, $L= 1, 2$, and even 3. But all of them have the common hallmark of a high degree of sensitivity to impurities. In this talk I will discuss impurity effects in the best known of these unconventionally paired systems, ^3He , a paradigm for the other unconventional superconductors. Impurity scattering is deftly controlled in superfluid ^3He by imbibing it into high porosity silica aerogel. We can understand the suppression of its superfluid state (the transition temperature), the effect on its order parameter (the pairing energy), the appearance of quasiparticle bound states (gaplessness), and possibly new phases, in the context of current theory. I will discuss experiments from many laboratories and their theoretical interpretation leading to the topical question of the day, "Can anisotropic scattering stabilize new anisotropic states?"



NORTHWESTERN
UNIVERSITY

Unconventional Pairing and Impurities in Superfluid ^3He

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Northwestern University
National Science Foundation DMR-0703656
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Outline:

Superfluid ^3He in silica-aerogel

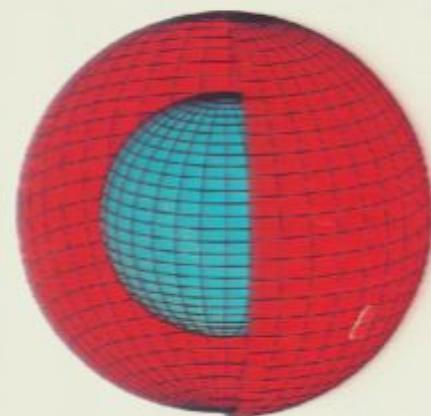
Impurity scattering models:
homogeneous or inhomogeneous?

Control of phase stability and nucleation with
global anisotropy

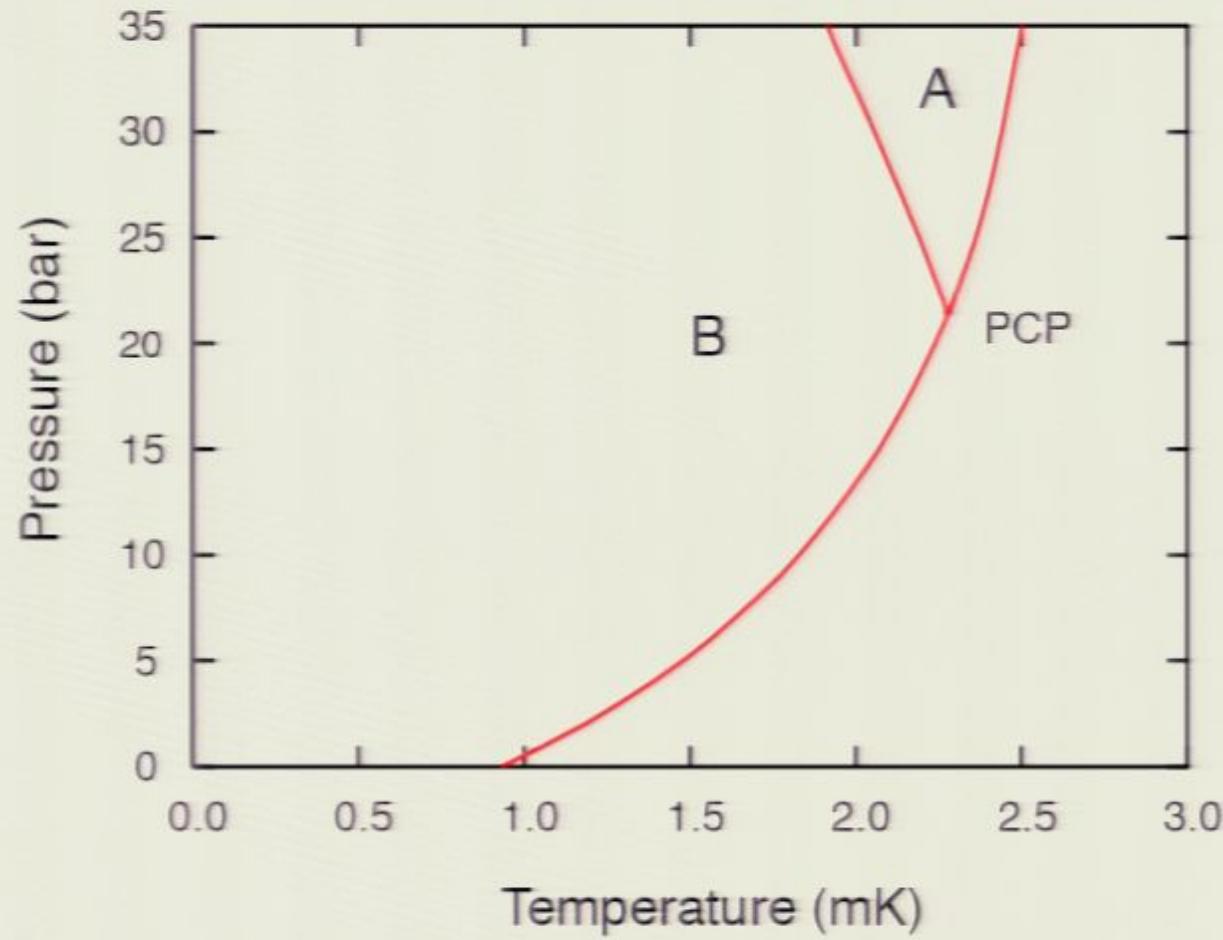
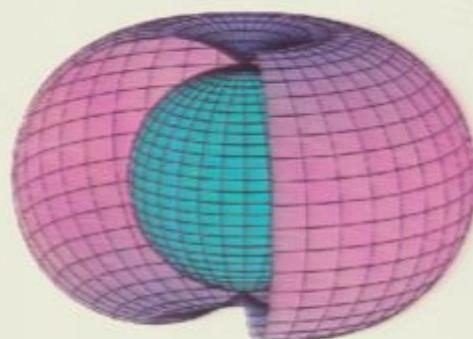
New phases?

Superfluid ^3He : phase diagram

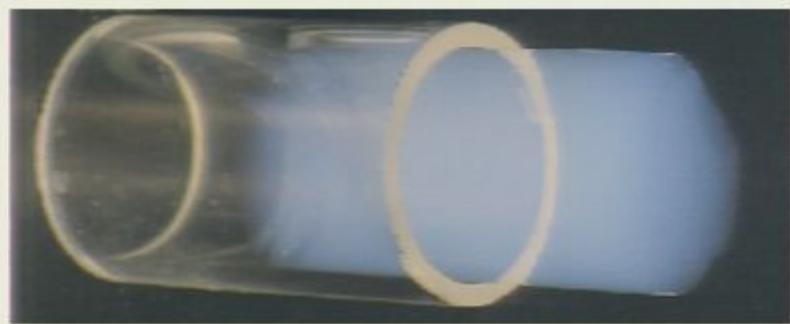
B - phase



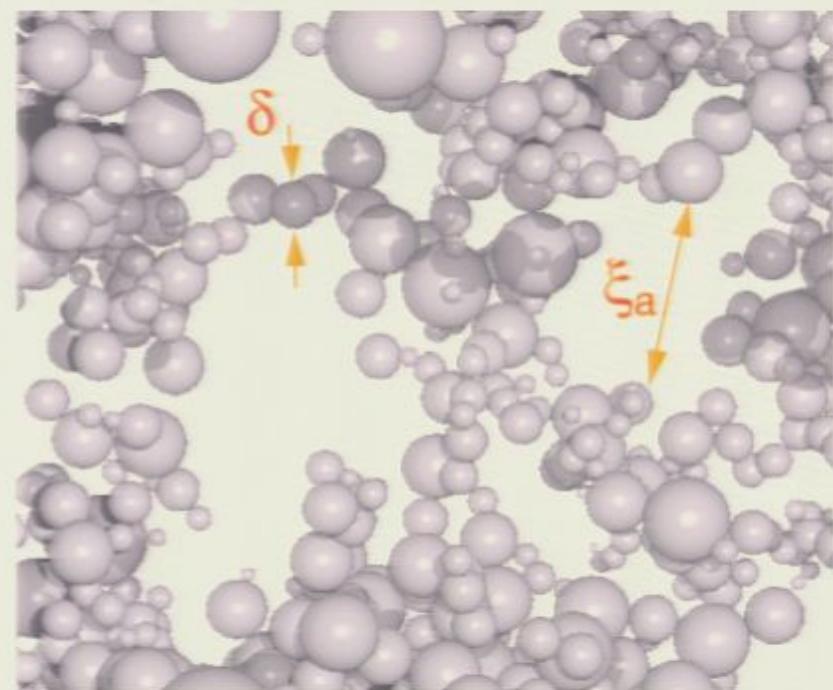
A - phase



Silica aerogel:



DLCA Calculation,
porosity, $\phi = 98\%$, Tom Lippman



particle diameter

$$\delta \approx 3 \text{ nm}$$

mean free path

$$\lambda \approx 150 \text{ nm}$$

correlation length

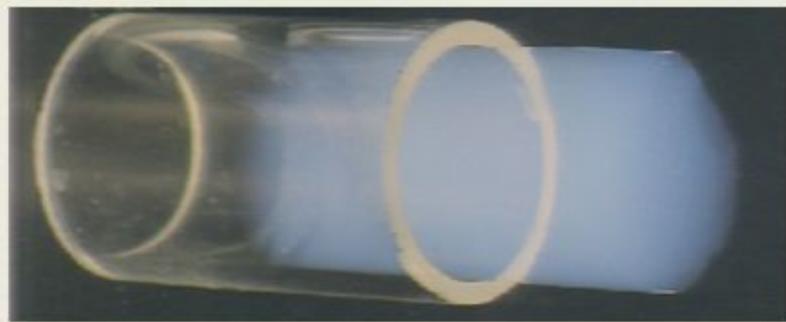
$$\xi_a \approx 20 \text{ nm}$$

superfluid

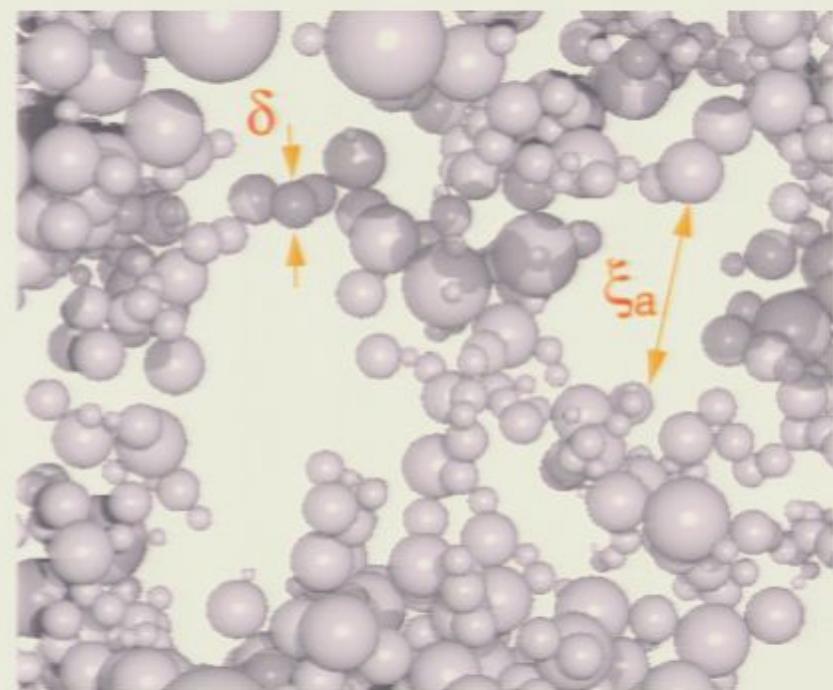
coherence length:

$$\xi \approx 20 \text{ to } 80 \text{ nm} \quad (34 \text{ to } 0 \text{ bar})$$

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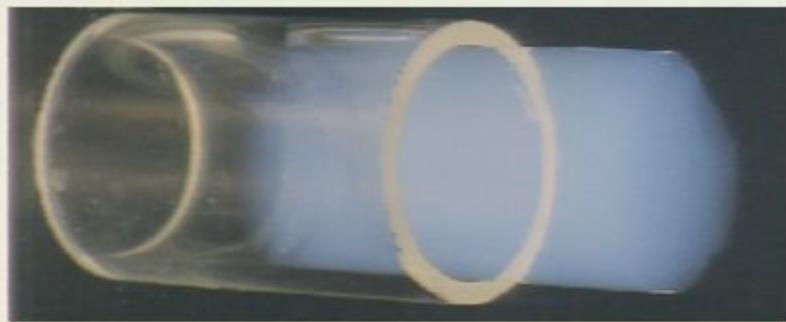
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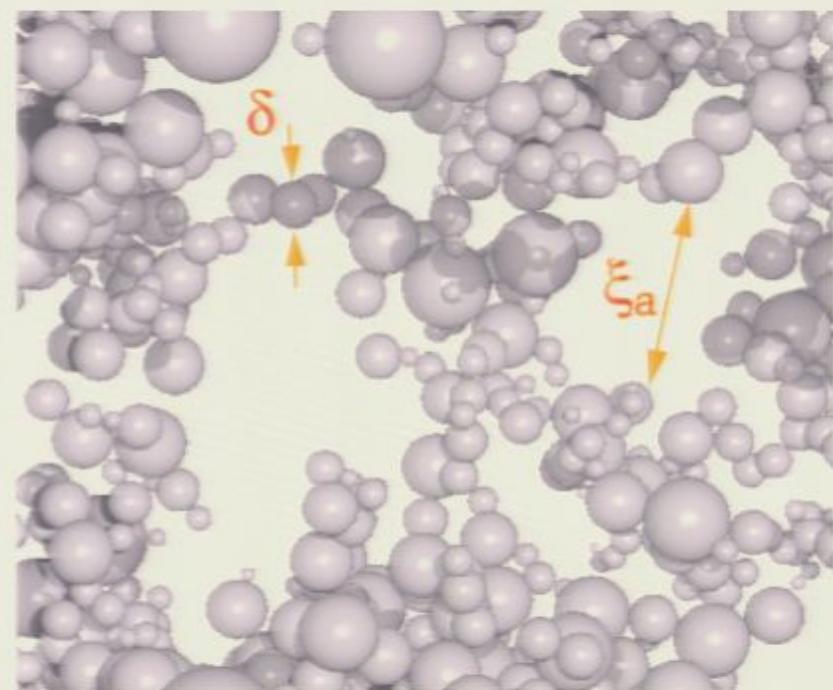
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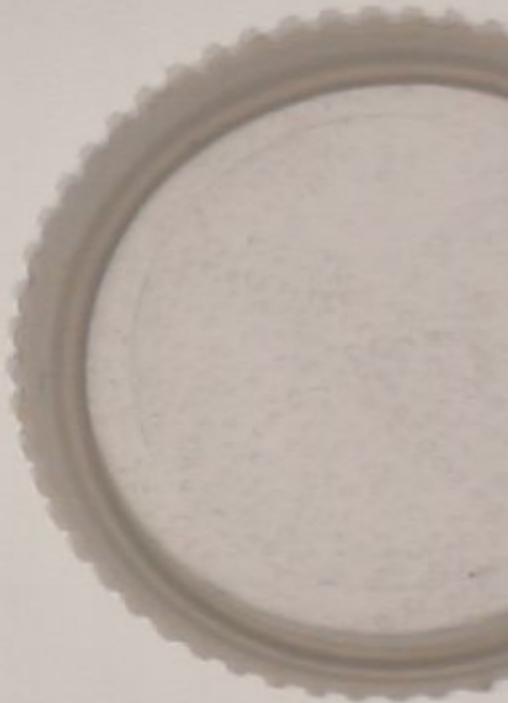
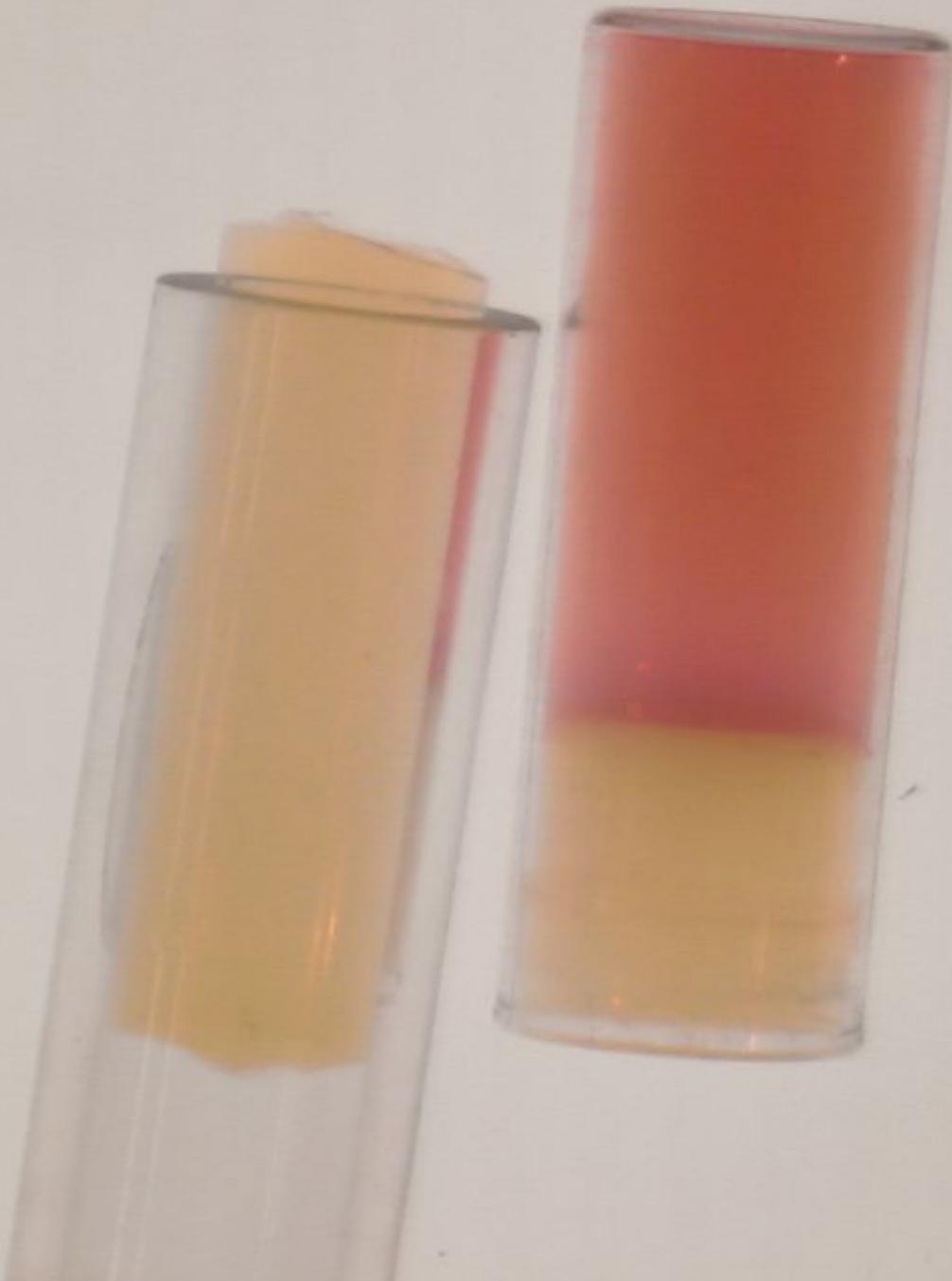
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superfluid

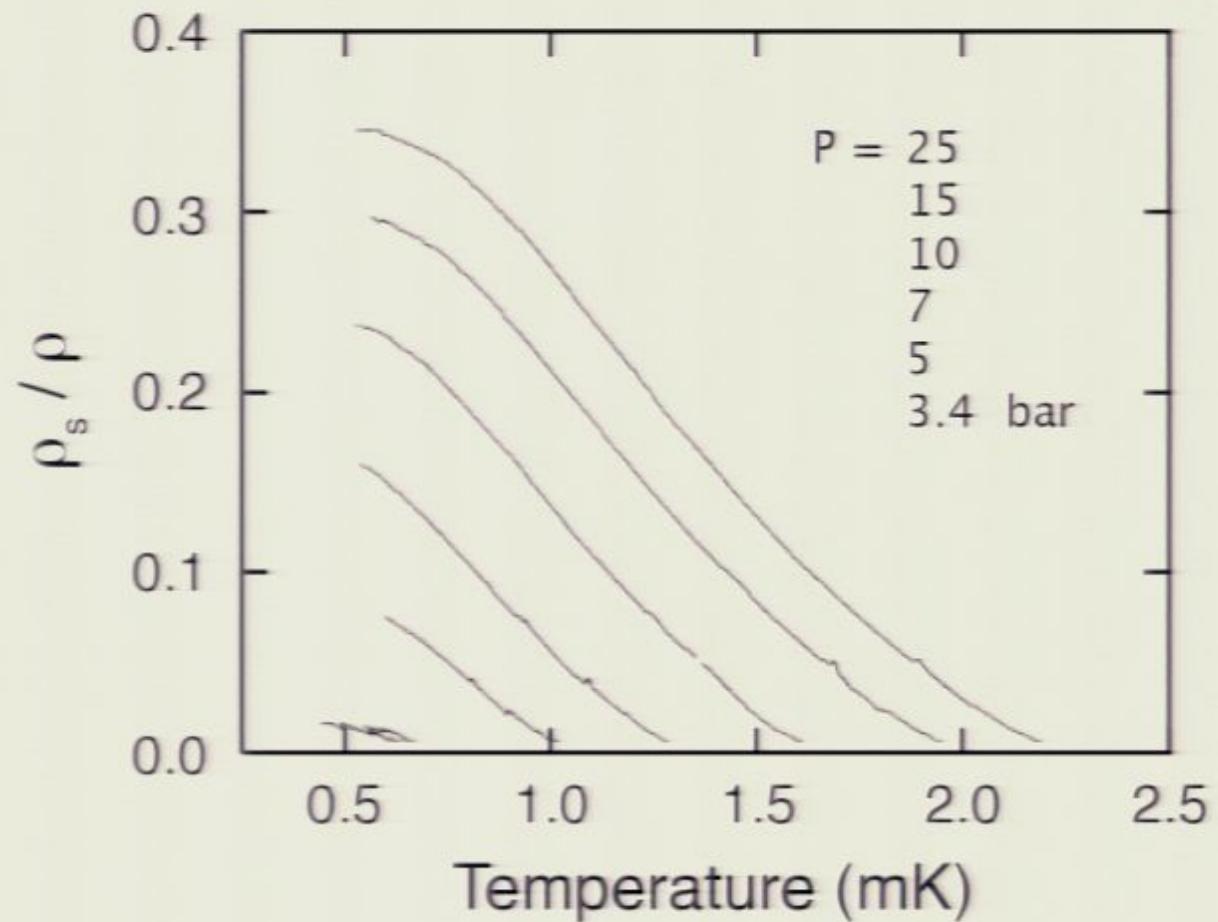
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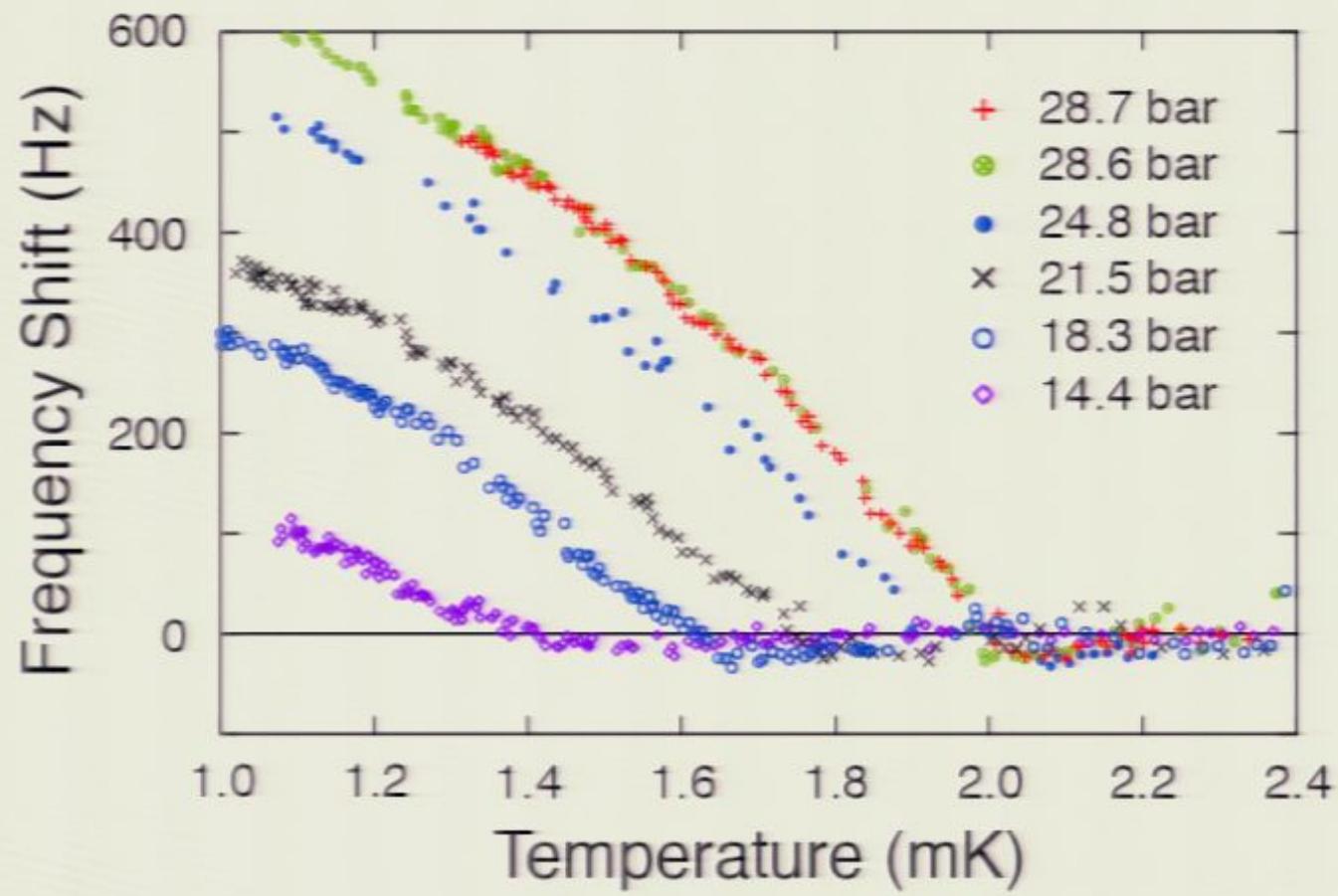
Superfluid density

Porto and Parpia
PRL 74, 4667 (1995)



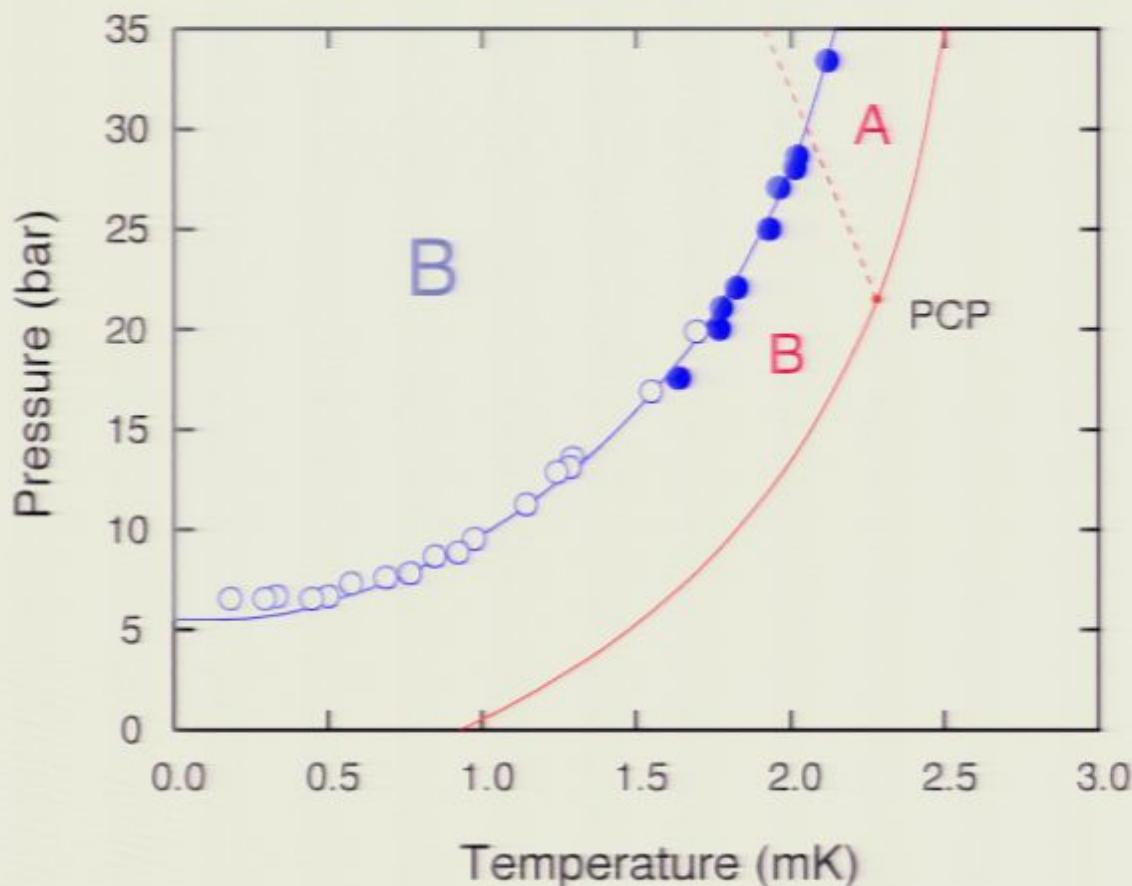
NMR frequency shift:

$$\Delta\omega \propto \Delta^2$$



Sprague et al.
PRL 75, 661 (1995);

Superfluid ^3He in aerogel: phase diagram



Torsional Oscillator

Matsumoto et al.

PRL 79, 253 (1997);

Acoustics:

*Gervais et al. PRL 87,
035701 (2001);*

PRB 66, 054528 (2002).

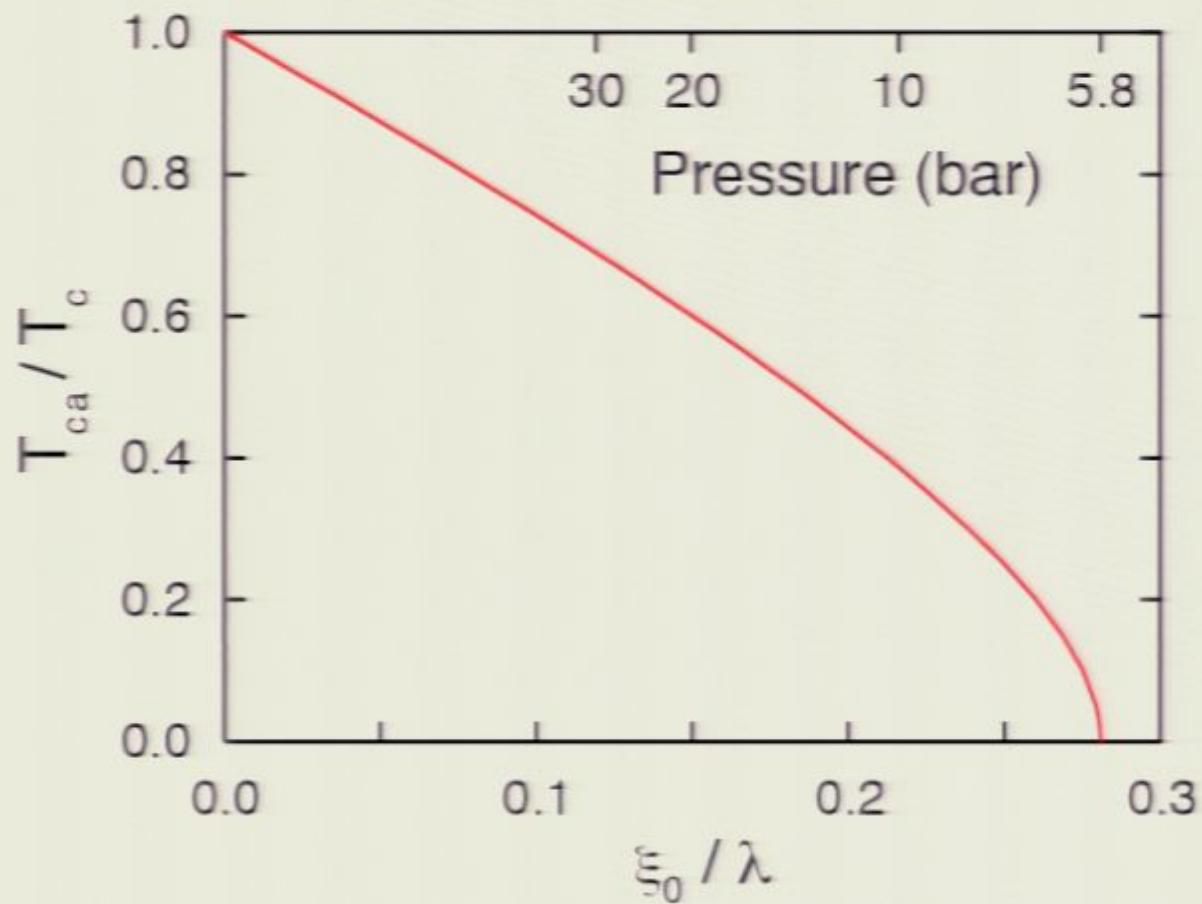
cw-NMR:

*Barker et al. PRL 85,
2148 (2000).*

pulsed-NMR:

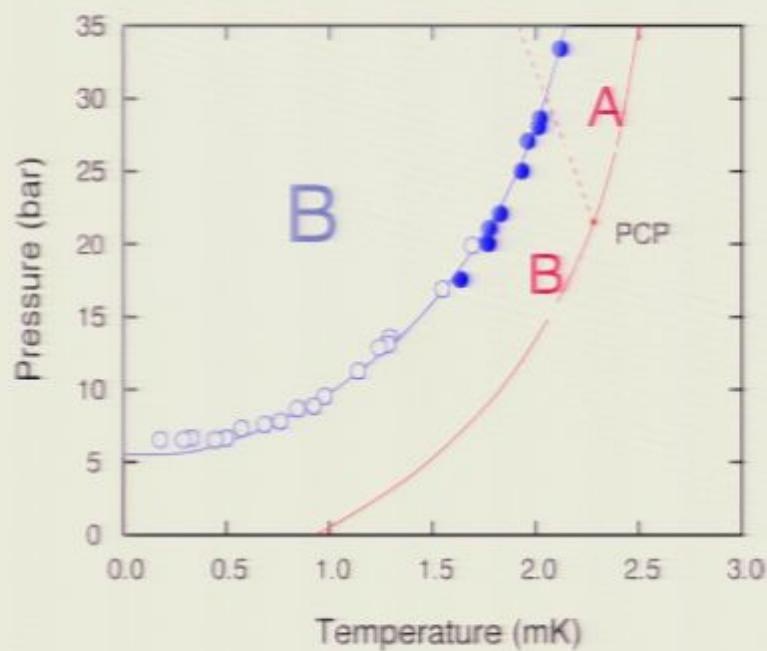
*Dmitriev et al. JETP
Lett. 76, 371 (2002).*

Abrikosov-Gorkov theory:

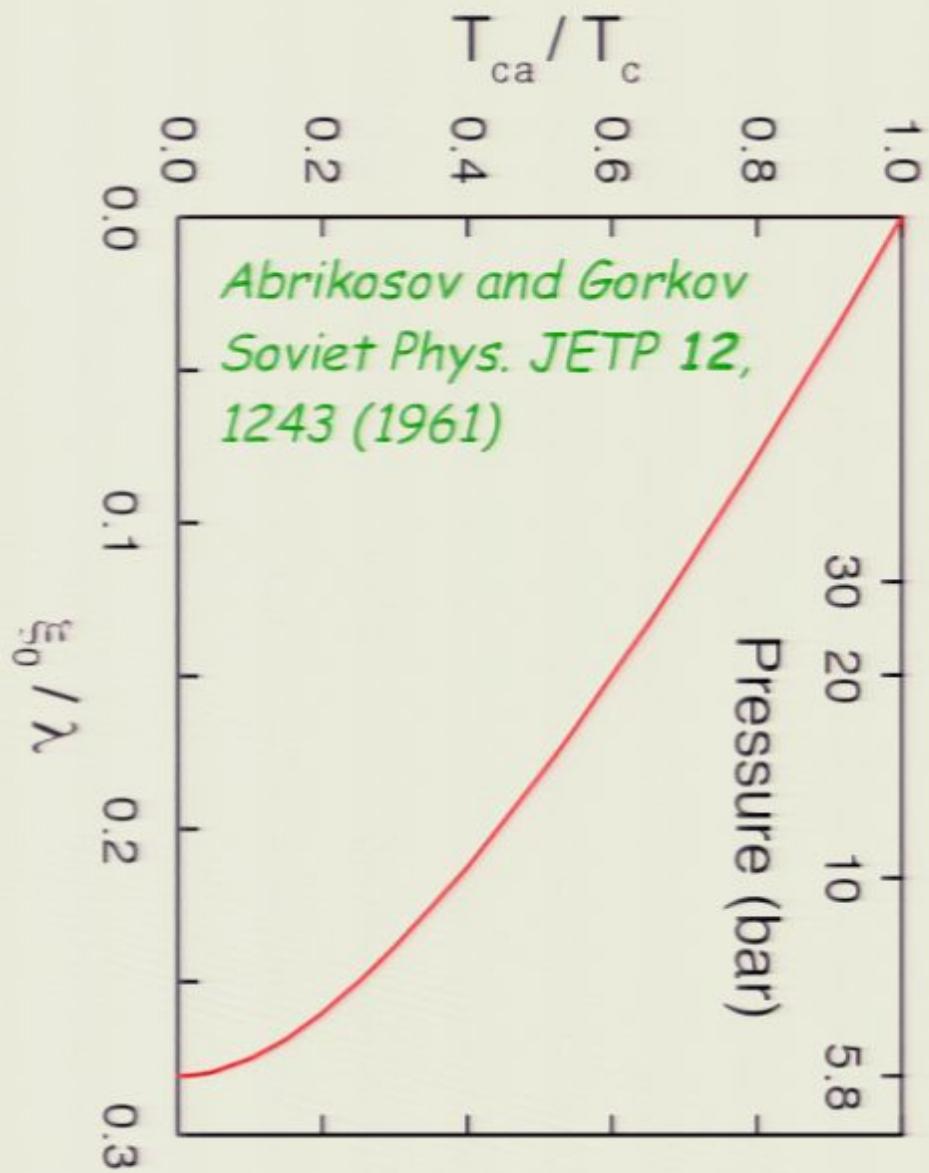


Abrikosov and Gorkov
Soviet Phys. JETP 12,
1243 (1961)

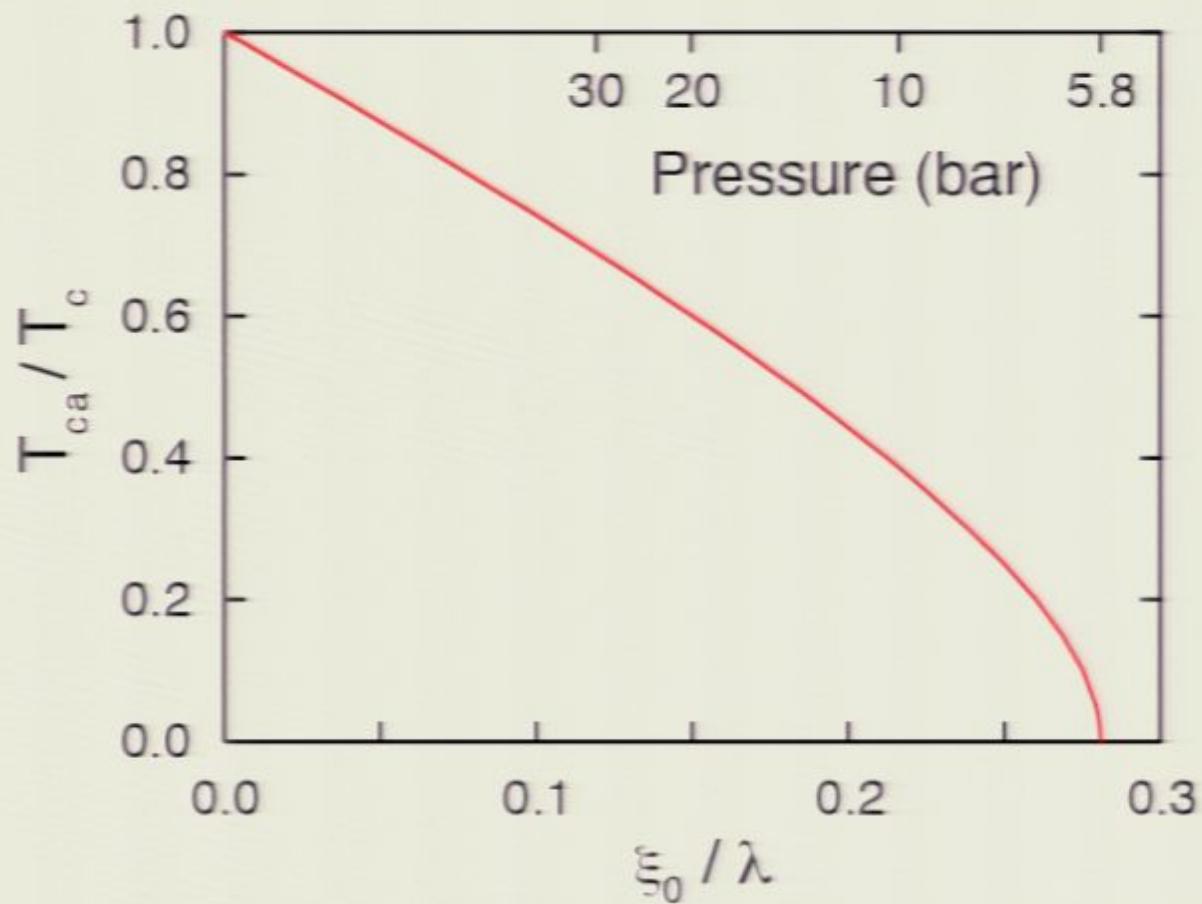
Abrikosov-Gorkov rotated:



Phase diagram
critical pressure

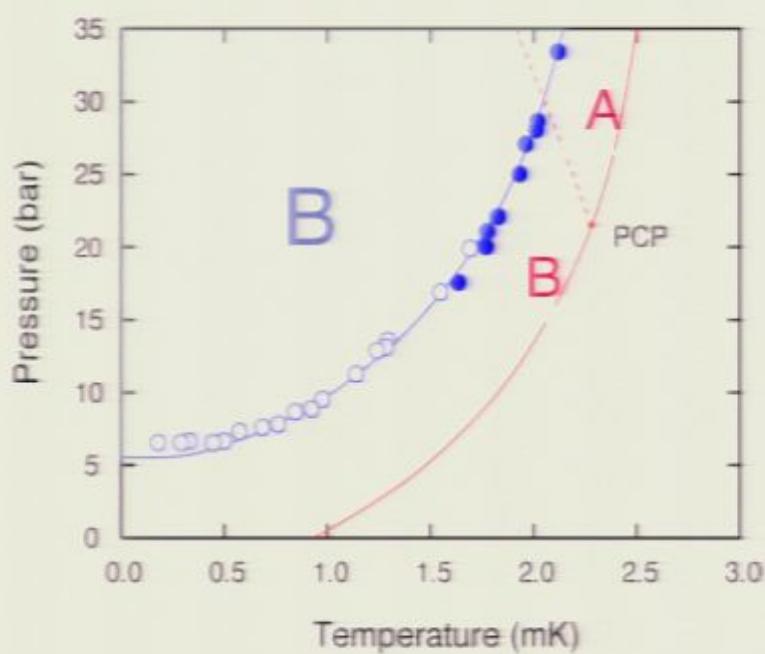


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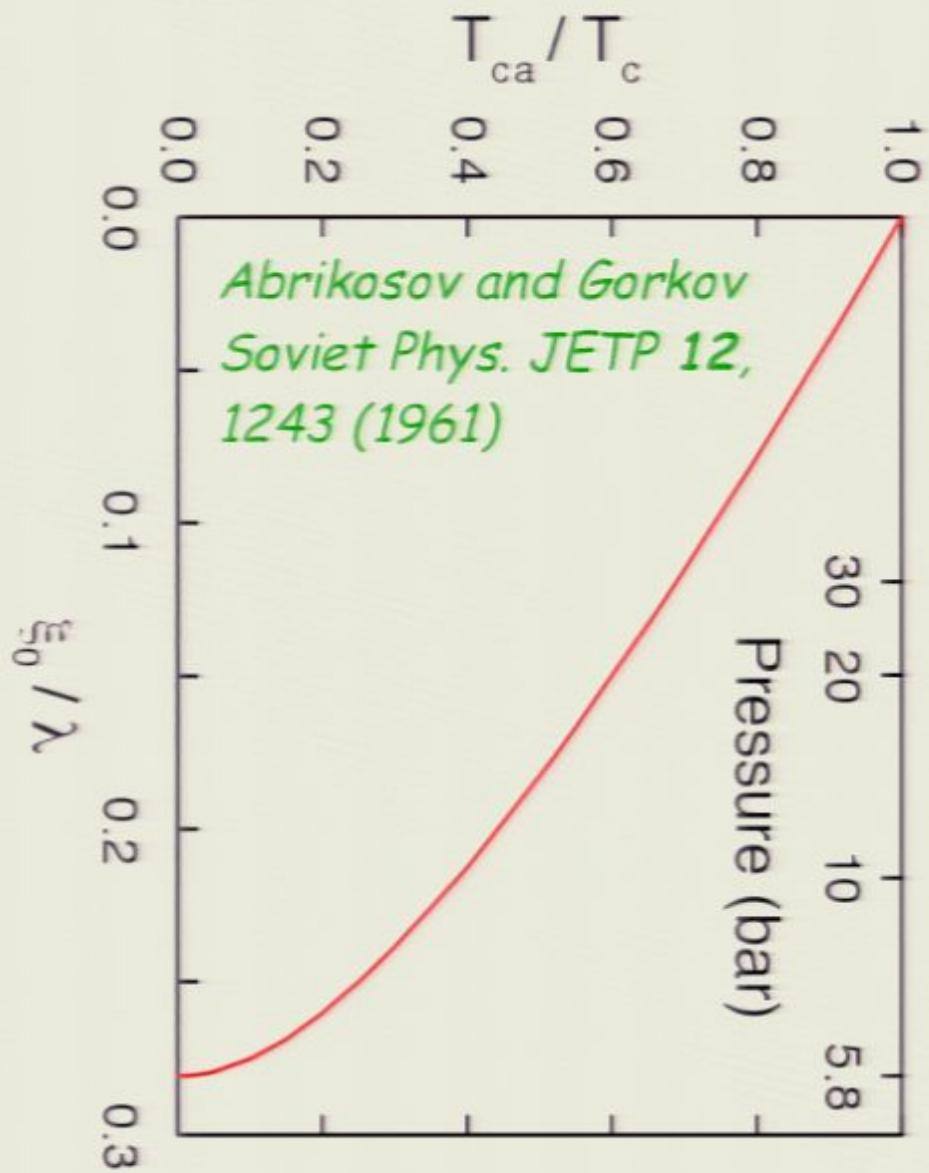


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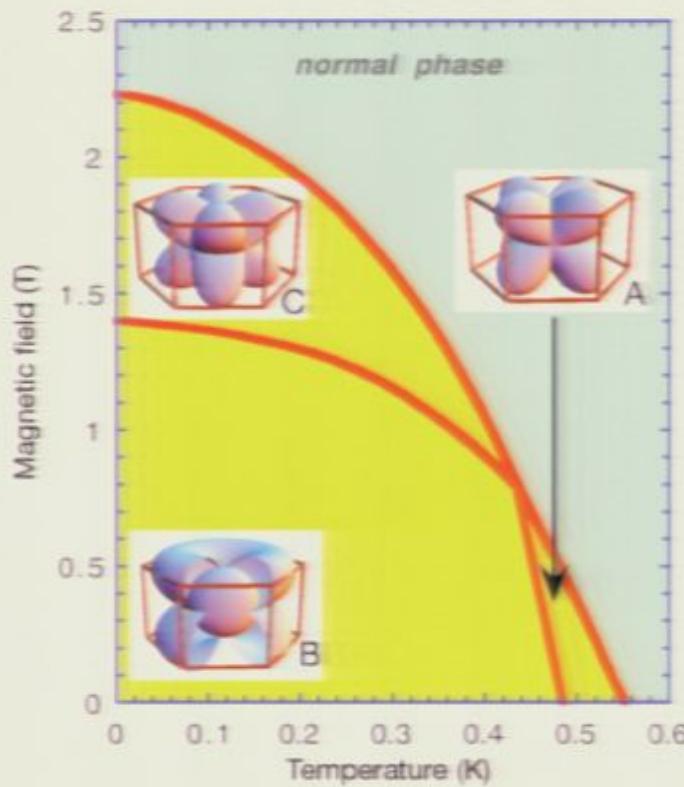
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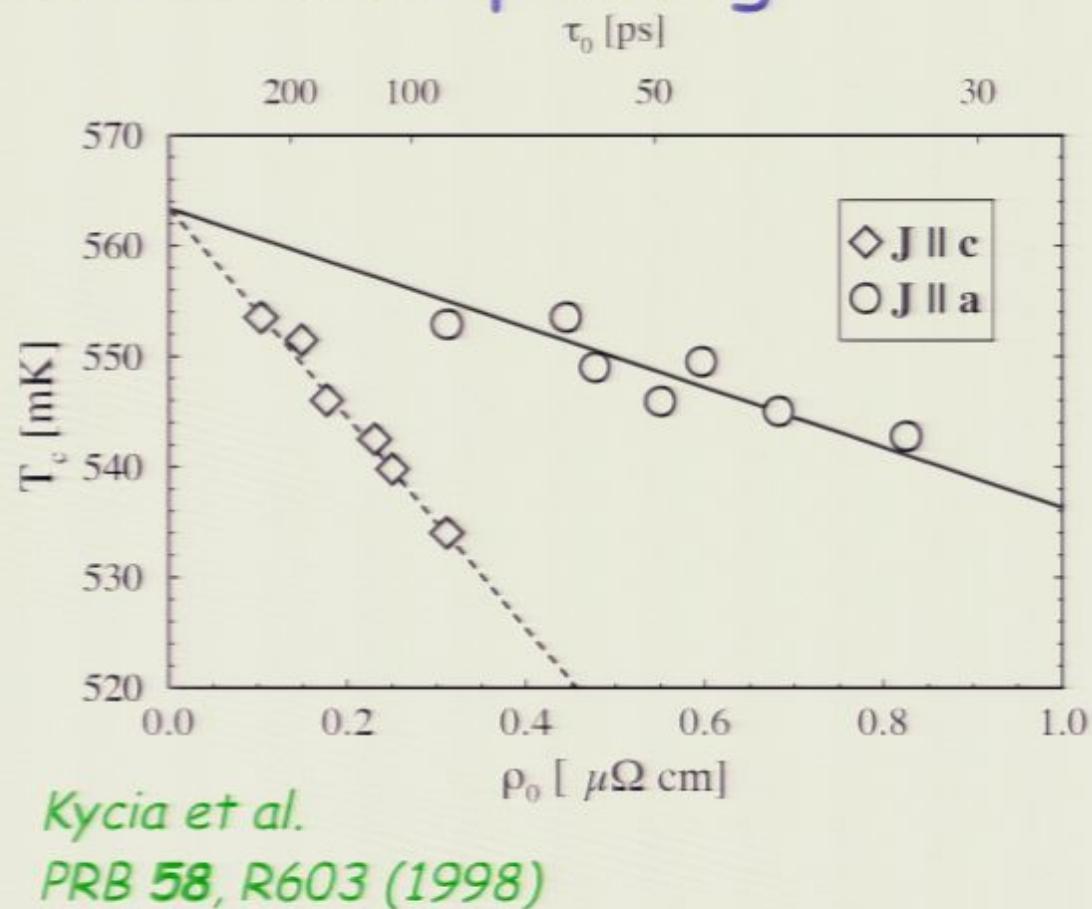
Phase diagram
critical pressure



UPt₃ another unconventional pairing system:



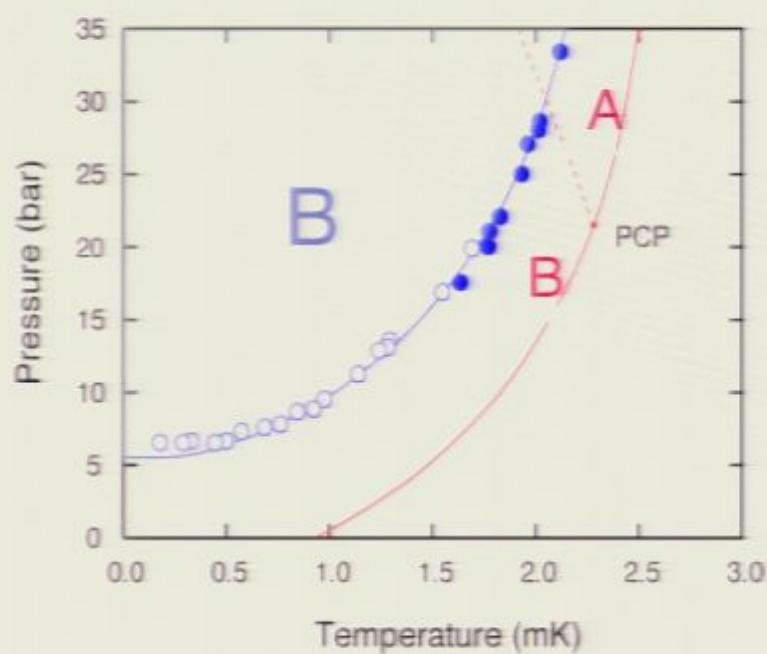
Phase diagram



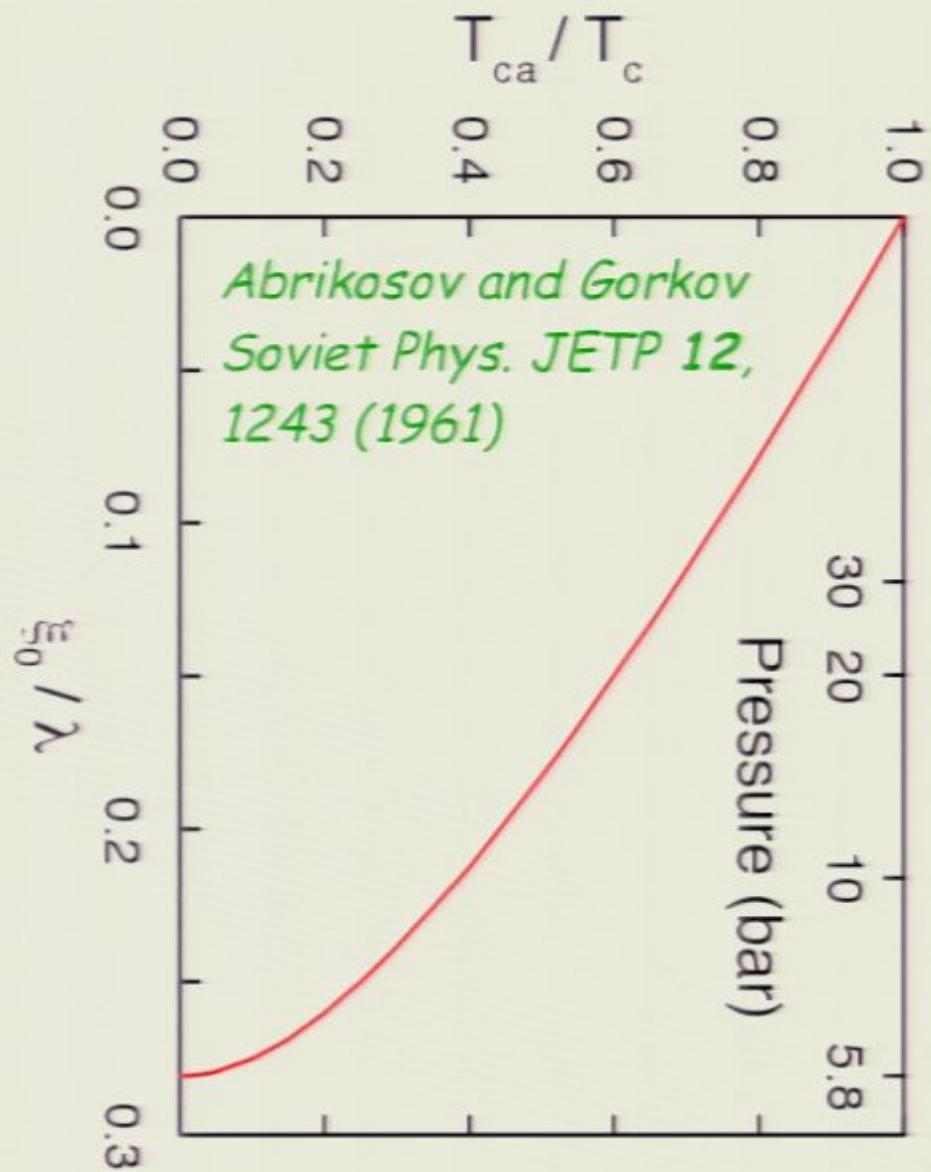
Kycia et al.
PRB 58, R603 (1998)



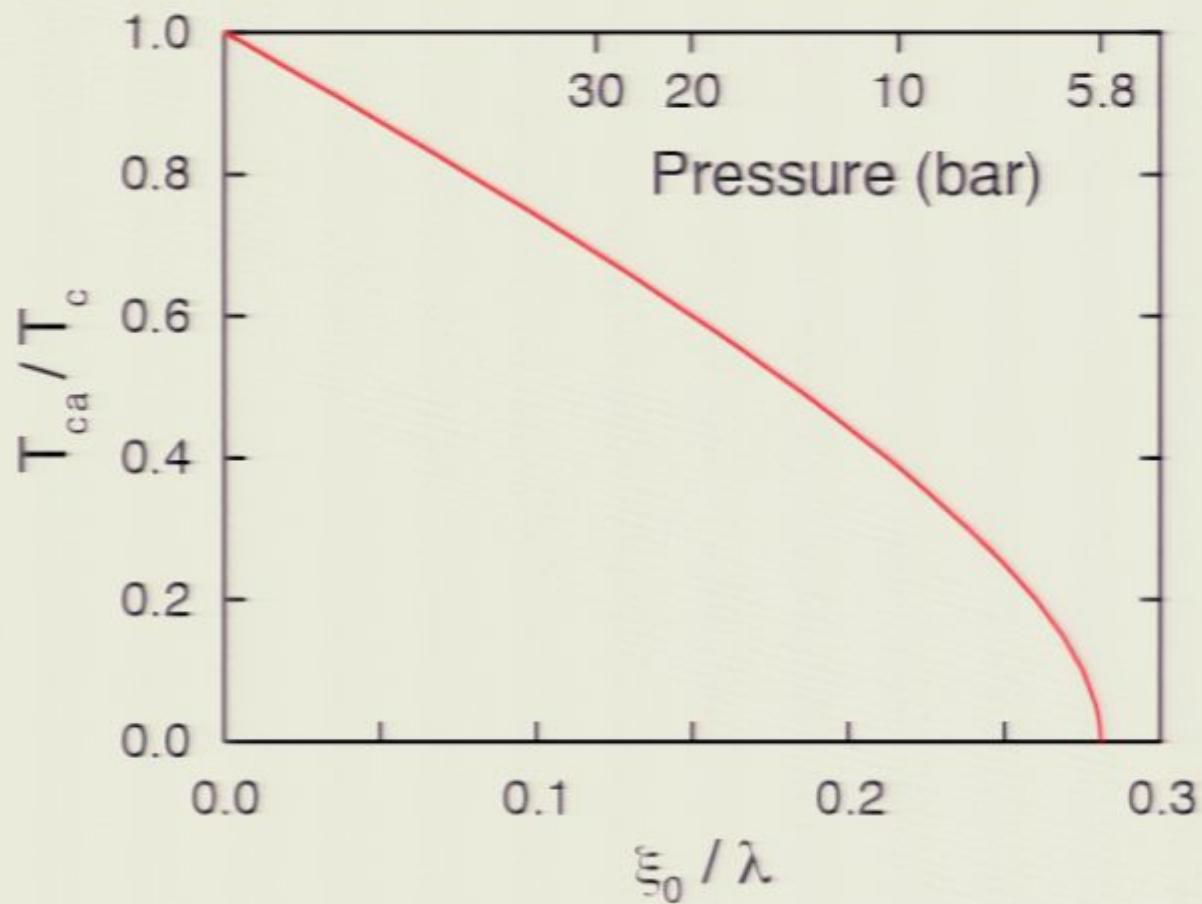
Abrikosov-Gorkov rotated:



Phase diagram
critical pressure

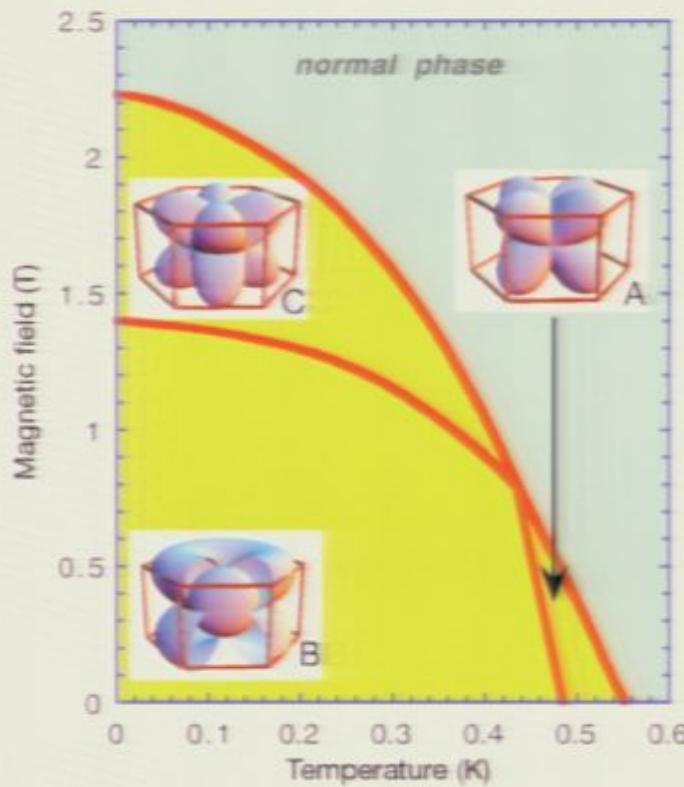


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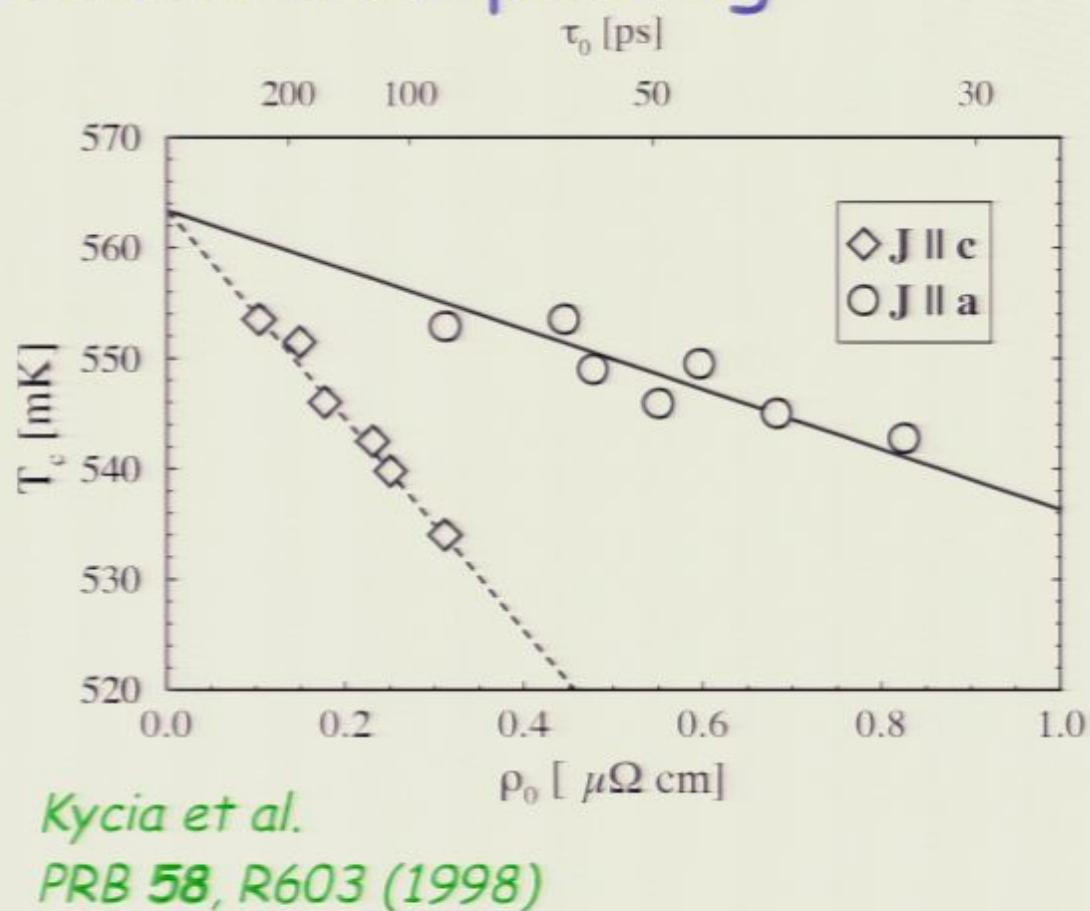


Abrikosov and Gorkov
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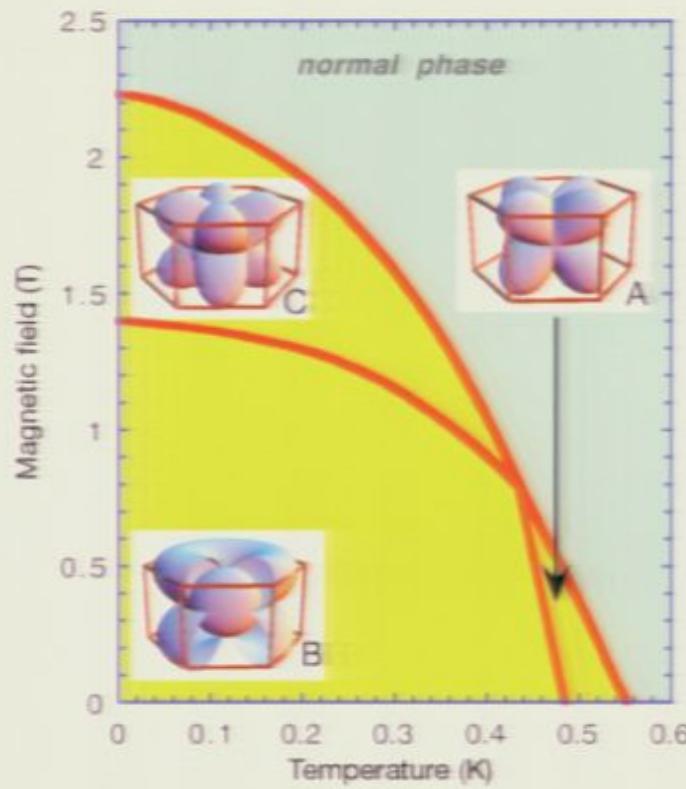
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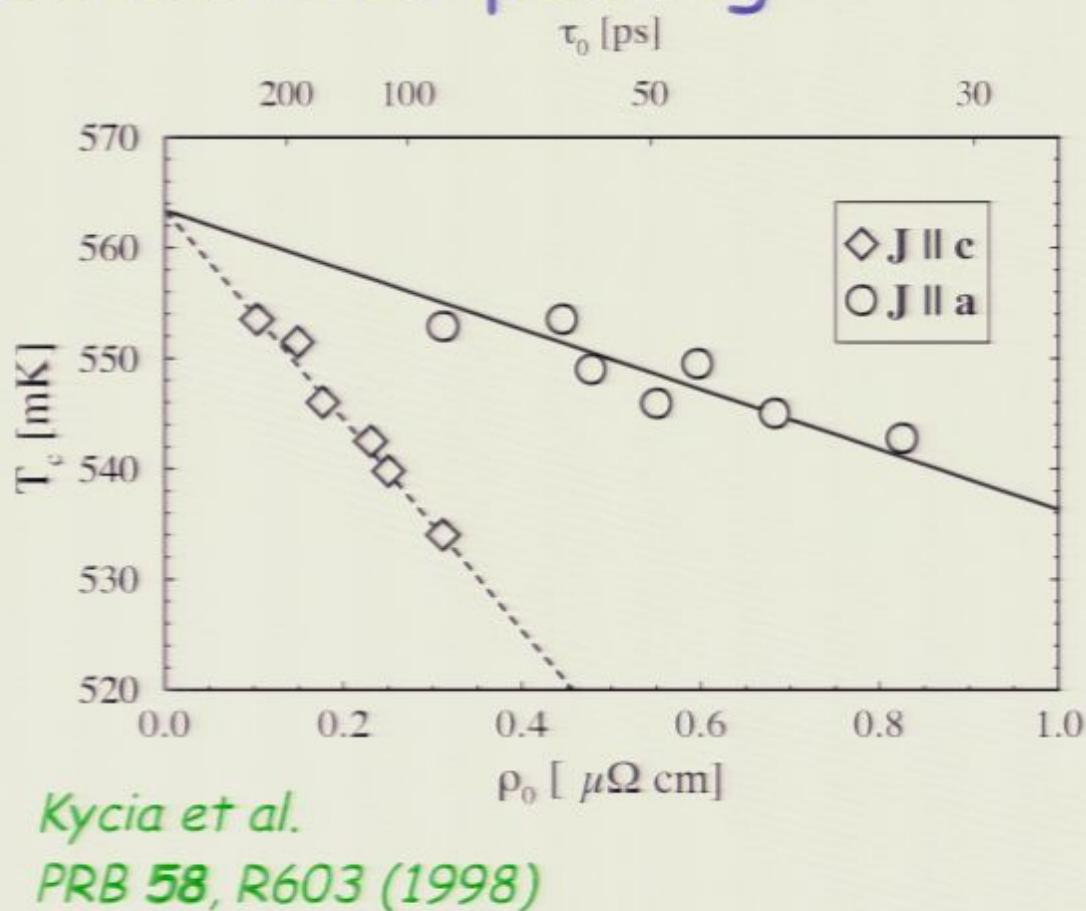
Phase diagram



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Phase diagram



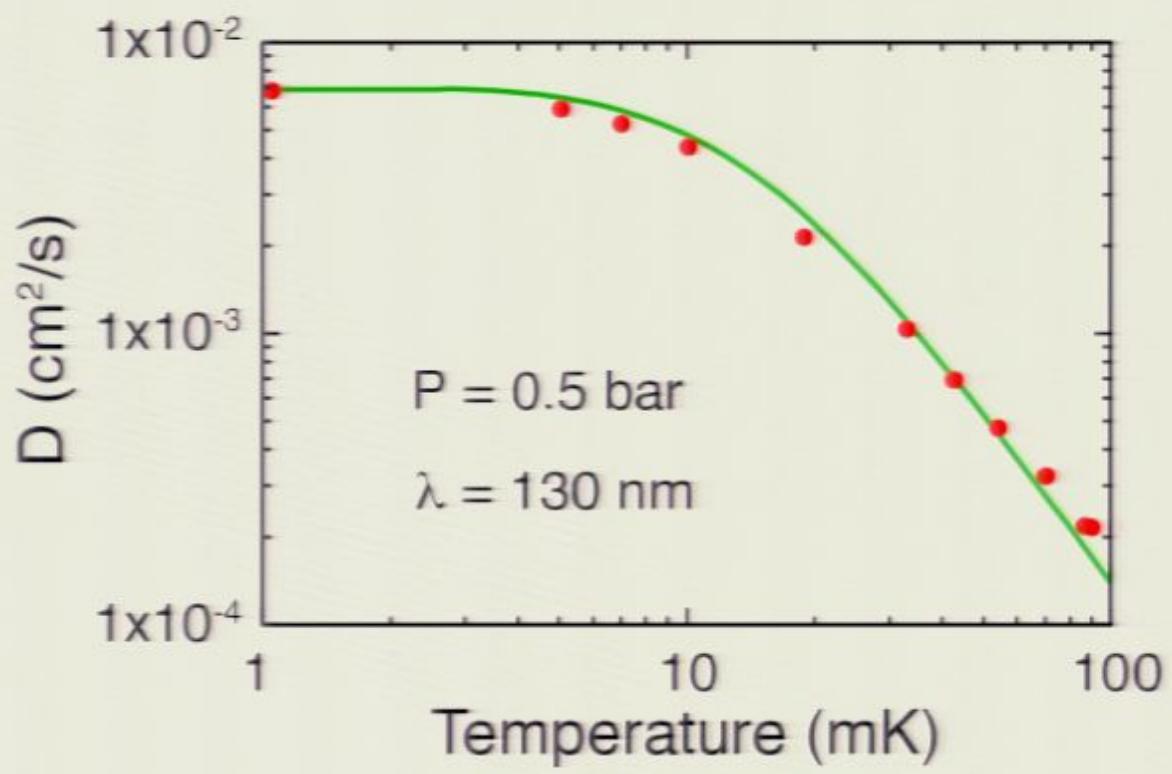
Kycia et al.
PRB 58, R603 (1998)



Transport mean-free-path λ :

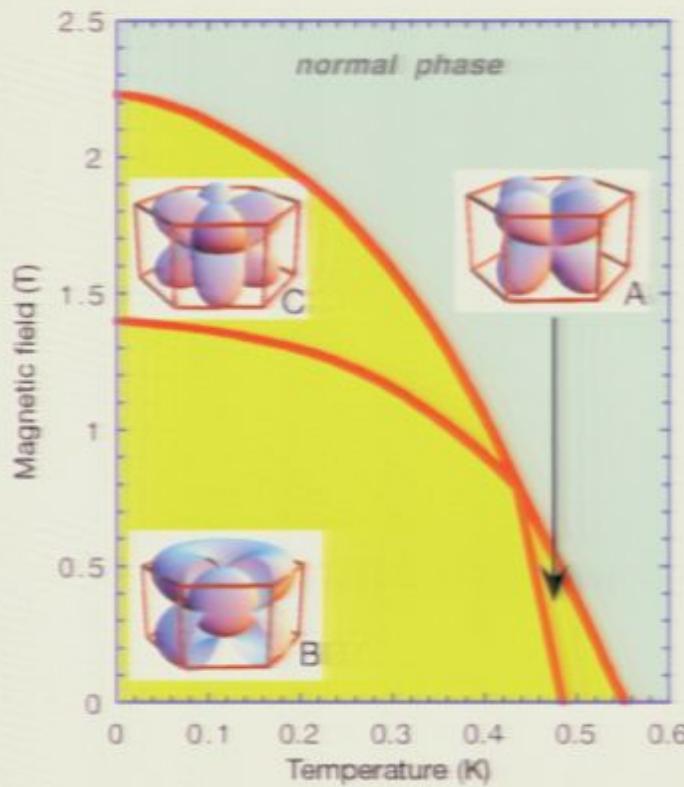
For pure ^3He
in the Fermi liquid

$$D \propto 1/T^2$$

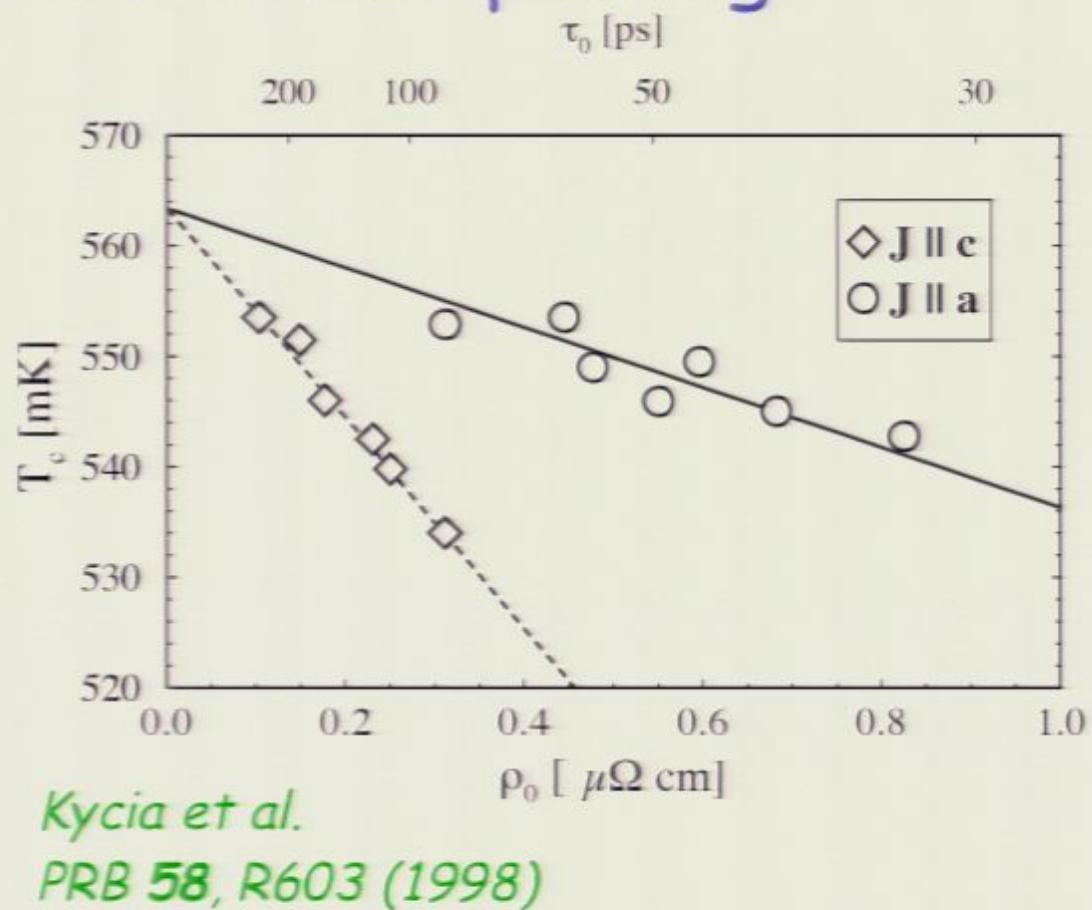


Collin, PhD thesis Univ.
J. Fourier (2002);
Bunkov et al JPCheM 66,
1325 (2005)

UPt₃ another unconventional pairing system:



Phase diagram



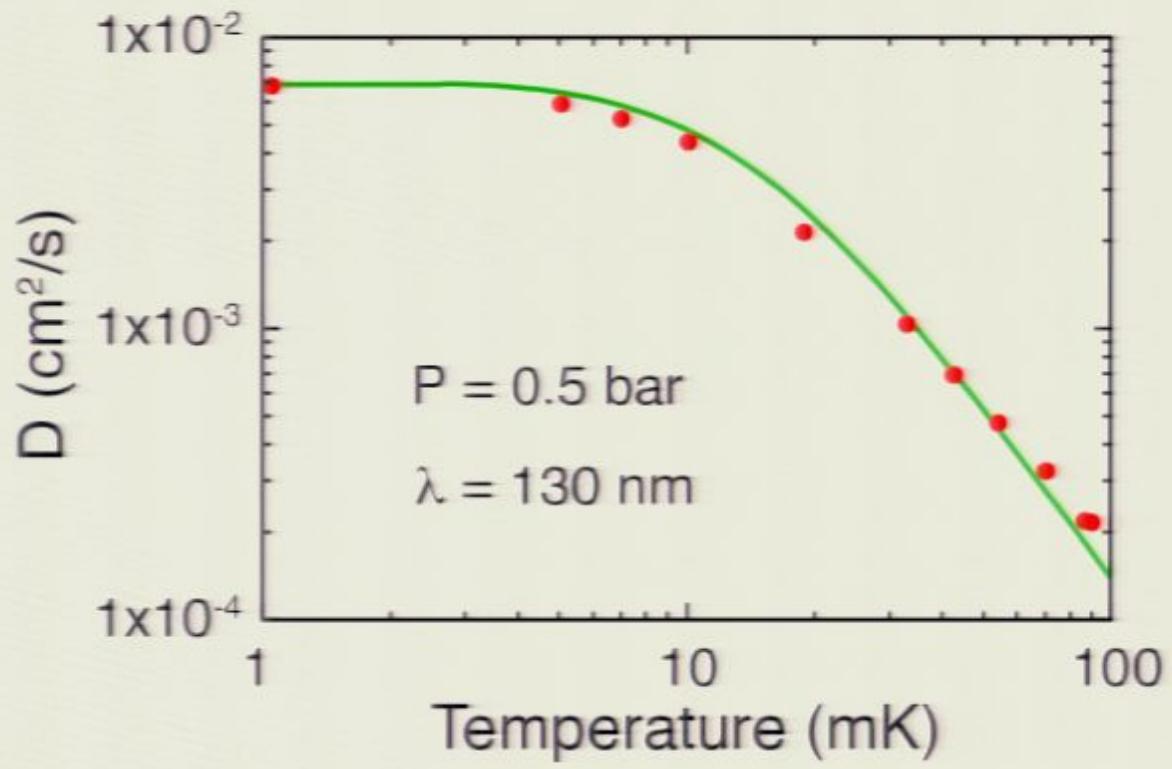
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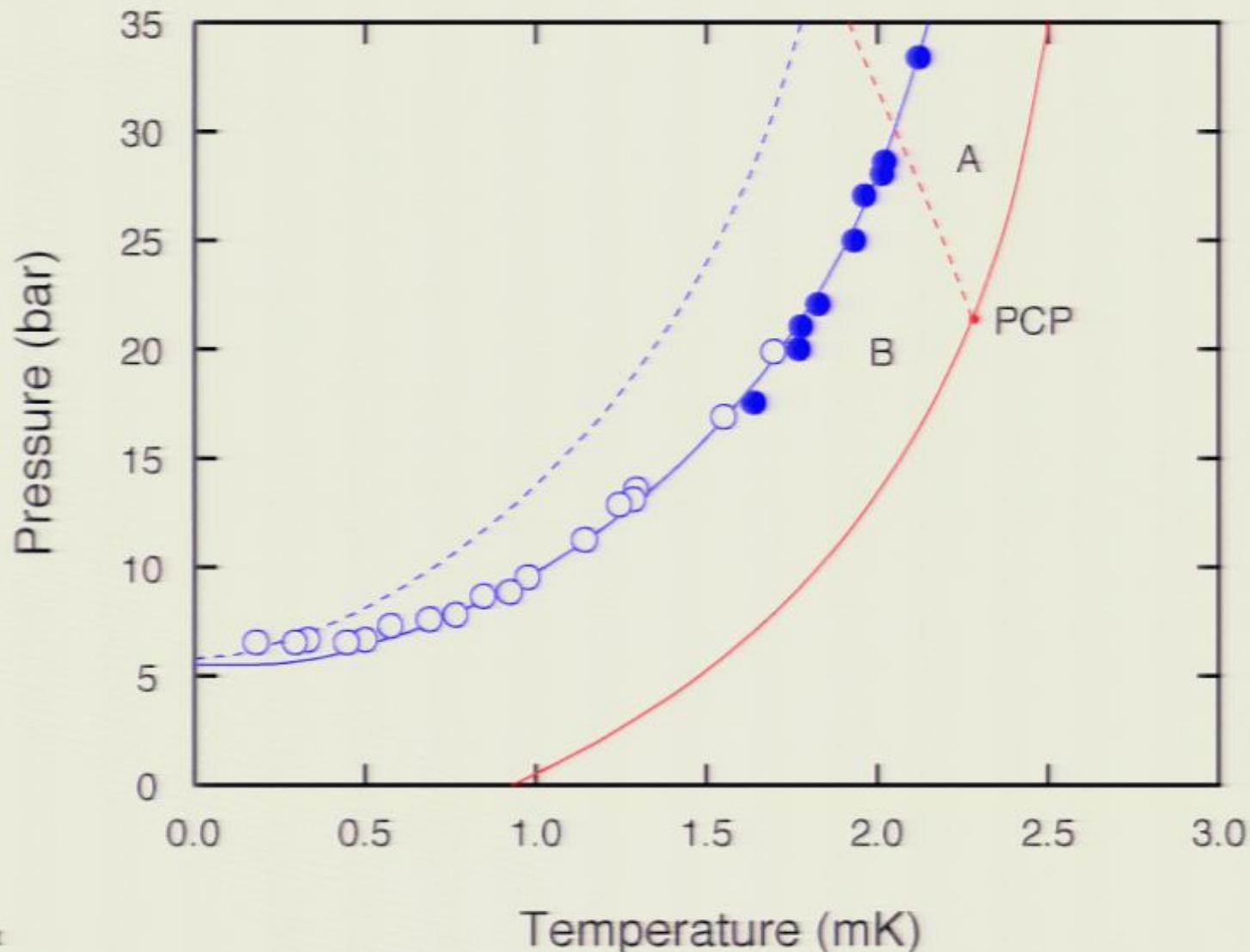
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Bunkov et al JPC Chem 66,
1325 (2005)

Phase diagram and AG theory:



Ginzburg-Landau free energy: (pure ^3He)

$$\begin{aligned} F = & -\alpha \text{Tr}(AA^\dagger) + g_z H_\mu (AA^\dagger)_{\mu\nu} H_\nu + \beta_1 |\text{Tr}(AA^T)|^2 \\ & + \beta_2 [\text{Tr}(AA^\dagger)]^2 + \beta_3 \text{Tr}(AA^T(AA^T)^*) \\ & + \beta_4 \text{Tr}((AA^\dagger)^2) + \beta_5 \text{Tr}(AA^\dagger(AA^\dagger)^*) \end{aligned}$$

$$\alpha = \frac{N(0)}{3} \left(\frac{T}{T_c} - 1 \right),$$

$$\frac{\beta_i}{\beta_0} = (-1, 2, 2, 2, -2), i = 1, \dots, 5,$$

$$\beta_0 = \frac{7\zeta(3)}{120\pi^2} \frac{N(0)}{(k_B T_c)^2},$$

Ginzburg-Landau theory for impurities:

Homogeneous Isotropic Scattering Model (HISM):

Silica aerogel is a system of point scatterers for ${}^3\text{He}$

One parameter: the transport mean free path, λ (unitary limit).

Determines the pair-breaking parameter: $x = \xi / \lambda$.

$$\ln \frac{T_c}{T_{c,a}} = 2 \sum_{n=1} \left(\frac{1}{2n-1} - \frac{1}{2n-1+x} \right)$$

$$\begin{pmatrix} \beta_1^a \\ \beta_2^a \\ \beta_3^a \\ \beta_4^a \\ \beta_5^a \end{pmatrix} = \beta_0^a \begin{pmatrix} -1 \\ 2 \\ 2 \\ 2 \\ -2 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \\ 0 \\ 1 \\ -1 \end{pmatrix} + \begin{pmatrix} \Delta\beta_1^{sc,a} \\ \Delta\beta_2^{sc,a} \\ \Delta\beta_3^{sc,a} \\ \Delta\beta_4^{sc,a} \\ \Delta\beta_5^{sc,a} \end{pmatrix}$$

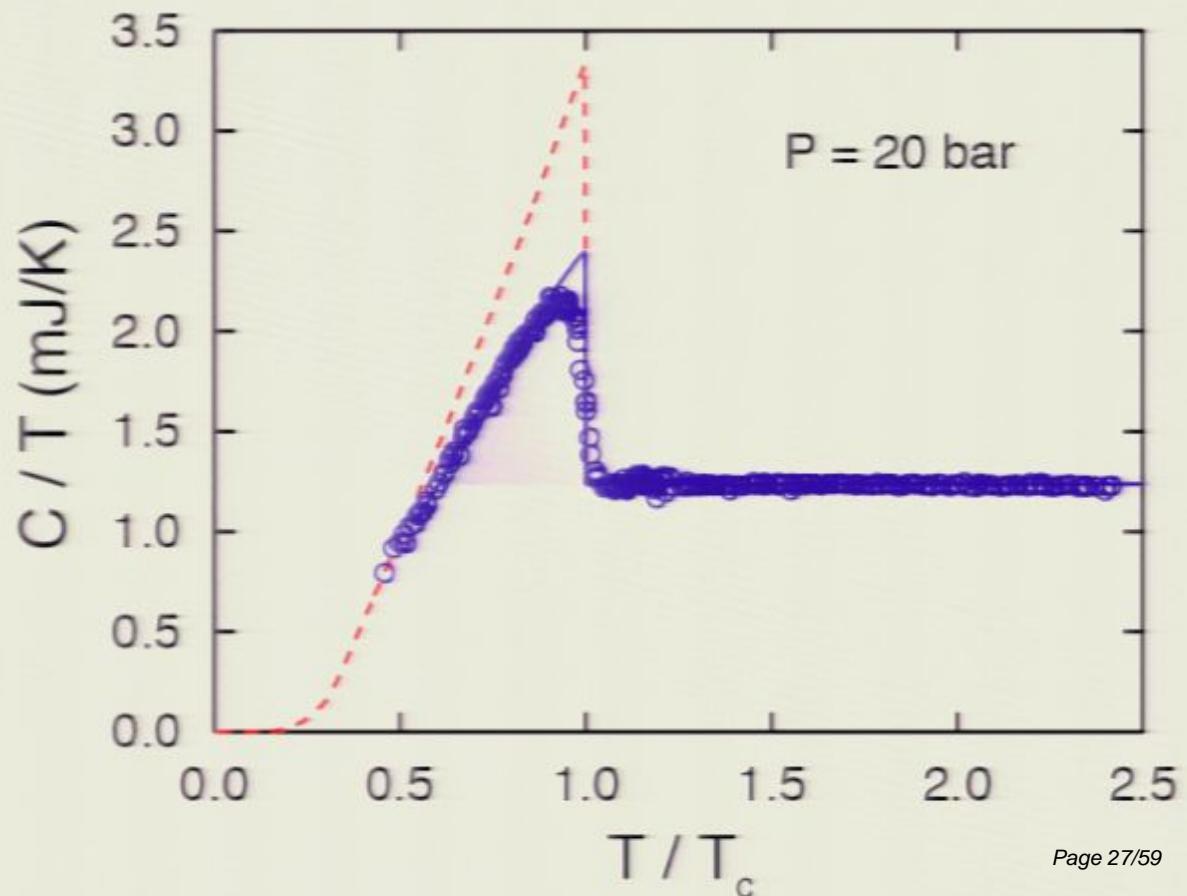
$$\beta_0^a = \frac{N(0)}{30(\pi k_B T_c)^2} \sum_{n=1} \frac{1}{(2n-1+x)^3},$$

$$b = \frac{N(0)}{9(\pi k_B T_c)^2} \left(\sin^2 \delta_0 - \frac{1}{2} \right) \sum_{n=1} \frac{x}{(2n-1+x)^4}.$$

Thuneberg et al.
PRL 80, 2861 (1998);
Page 26/59

Specific heat jump:

$$\Delta(T) = \pi k_B T_c \left[\frac{2}{3} \left(\frac{T_c}{T} - 1 \right) \frac{\Delta C}{C} \right]^{1/2}$$



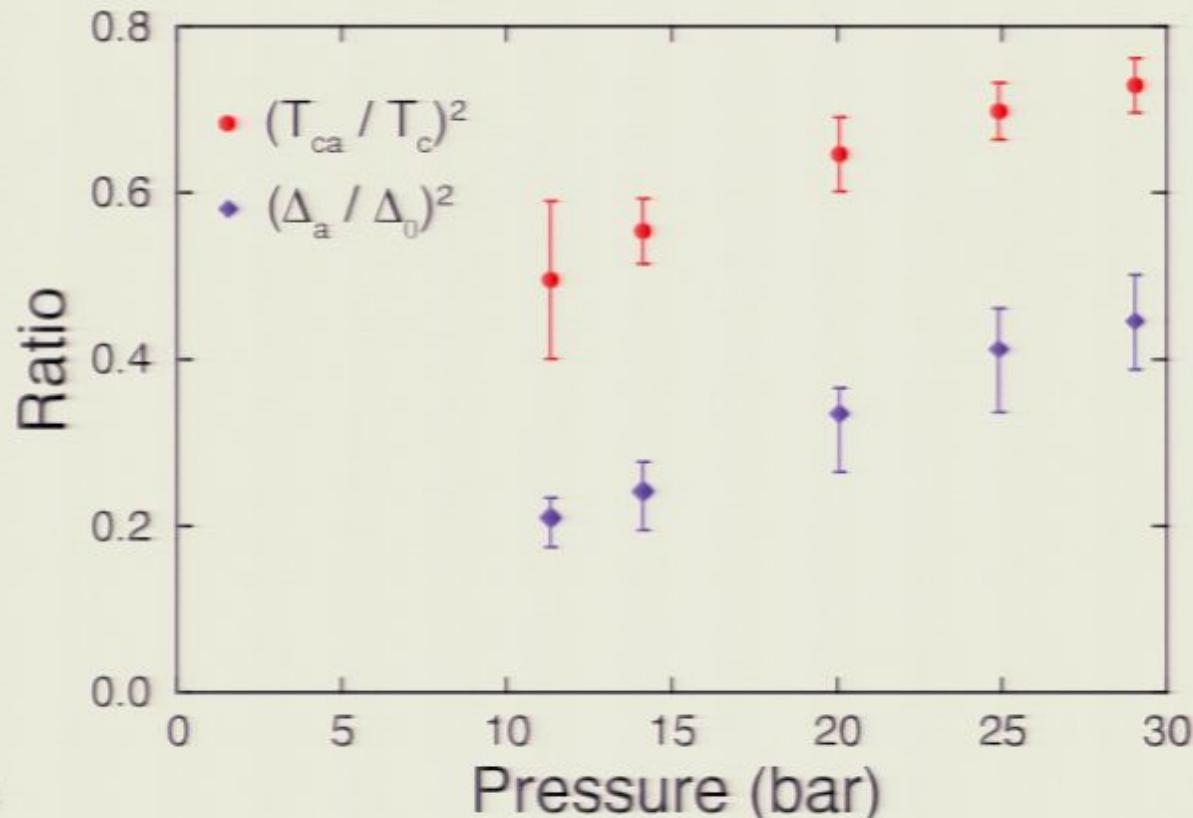
Choi et al.
PRL 93, 145301 (2004);

Failure of HISM:

Gap from specific heat jump:

$$\Delta^2(T) = (\pi k_B T_c)^2 \frac{2}{3} \frac{\Delta C}{C} \left(\frac{T_c}{T} - 1 \right)$$

Homogeneous Isotropic Scattering Model (HISM): suppression ratio must be the same for the energy gap and T_c (unitary scattering)



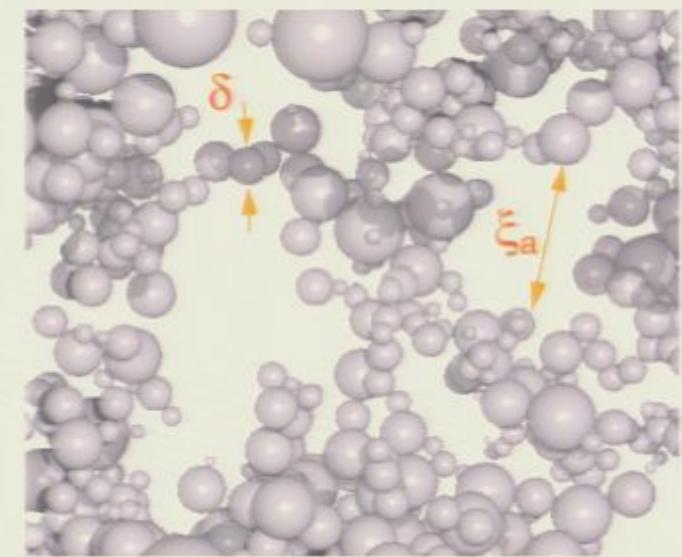
Ginzburg-Landau theory (cont'd):

Inhomogeneous Isotropic Scattering Model (IISM):

Define a new pair-breaking parameter:

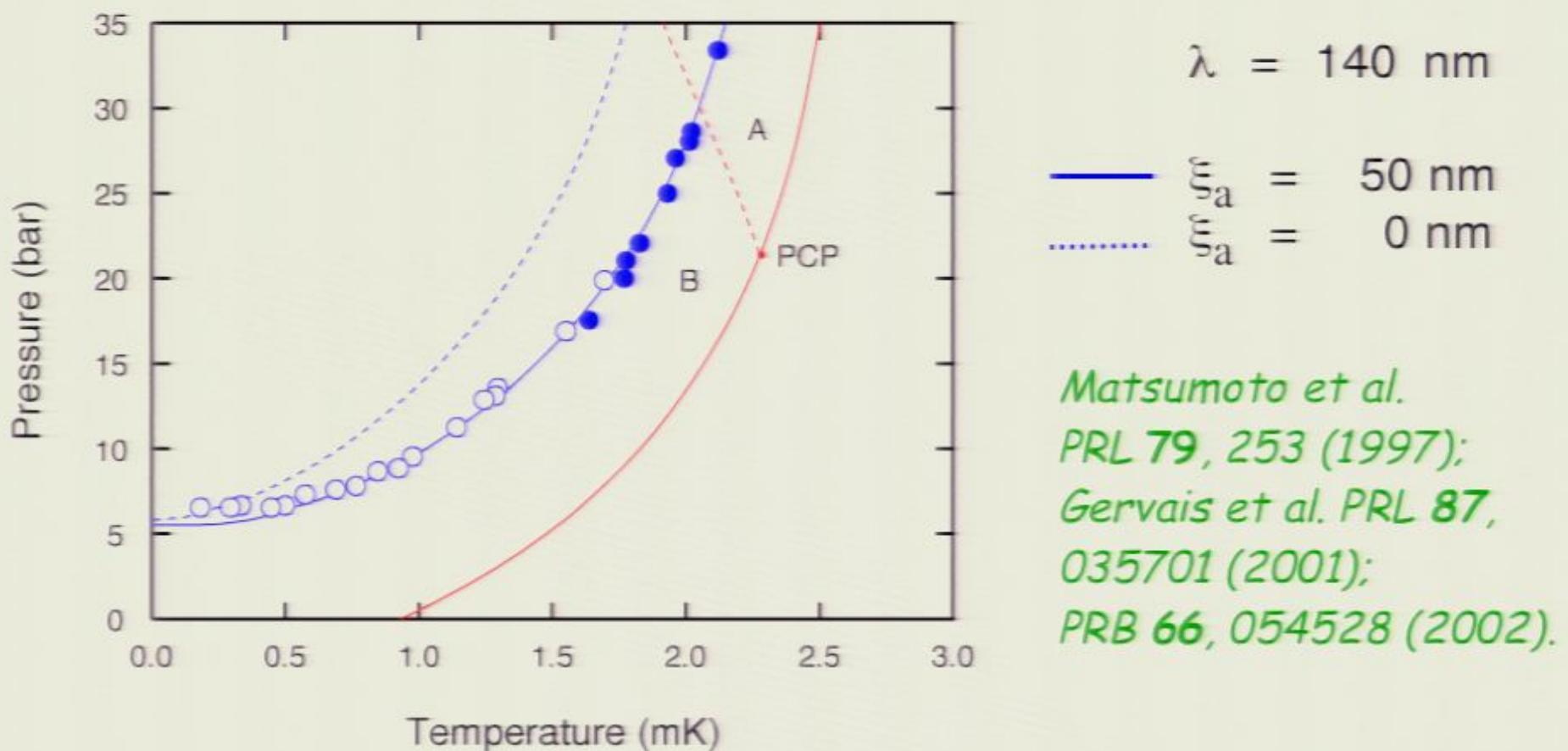
$$x = (\xi / \lambda) / [1 + \xi_a^2 / \lambda \xi]$$

in terms of a scattering correlation length ξ_a , proposed by Sauls and Sharma and calculated by Fomin.



Sauls and Sharma et al. PRB 68, 224502 (2003); Fomin, Pis'ma ZhETF 88, 65 (2008); Surovtsev, arXiv:0810.1102; Hanninen and Thuneberg PRB 67, 214507 (2003).

Phase diagram with correlated disorder:



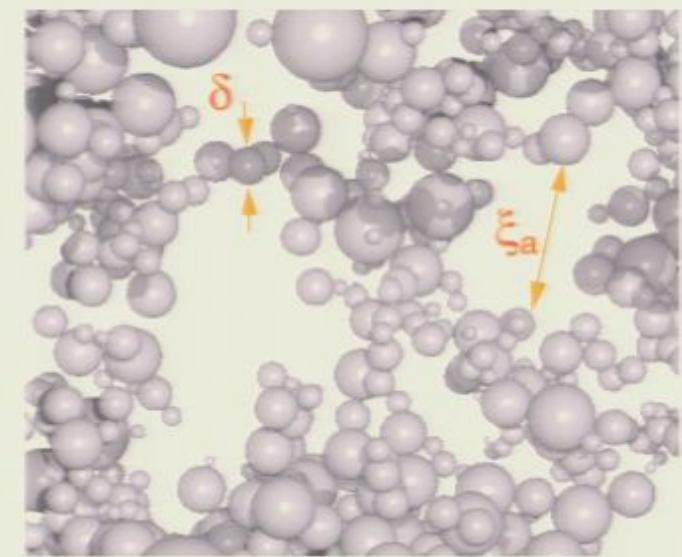
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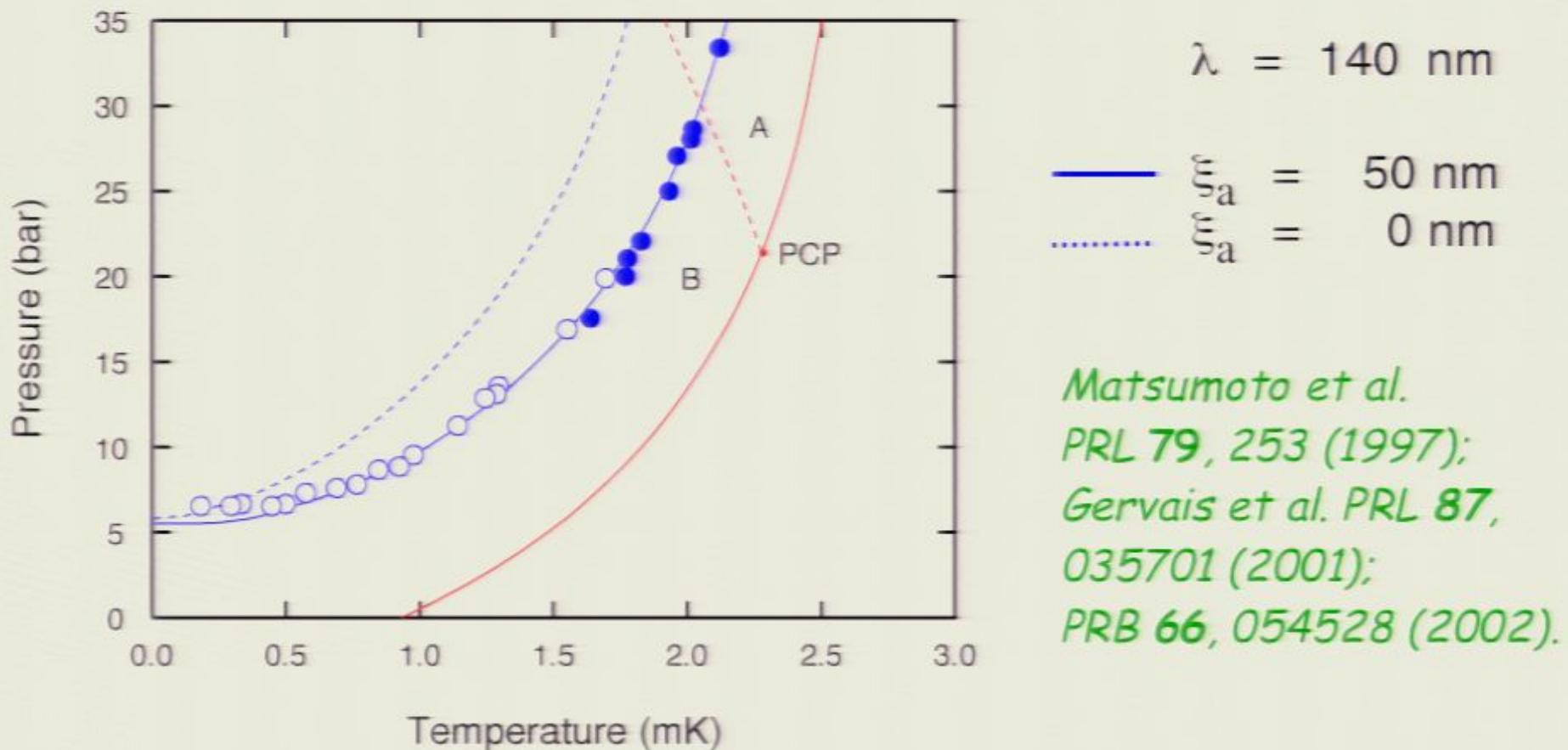
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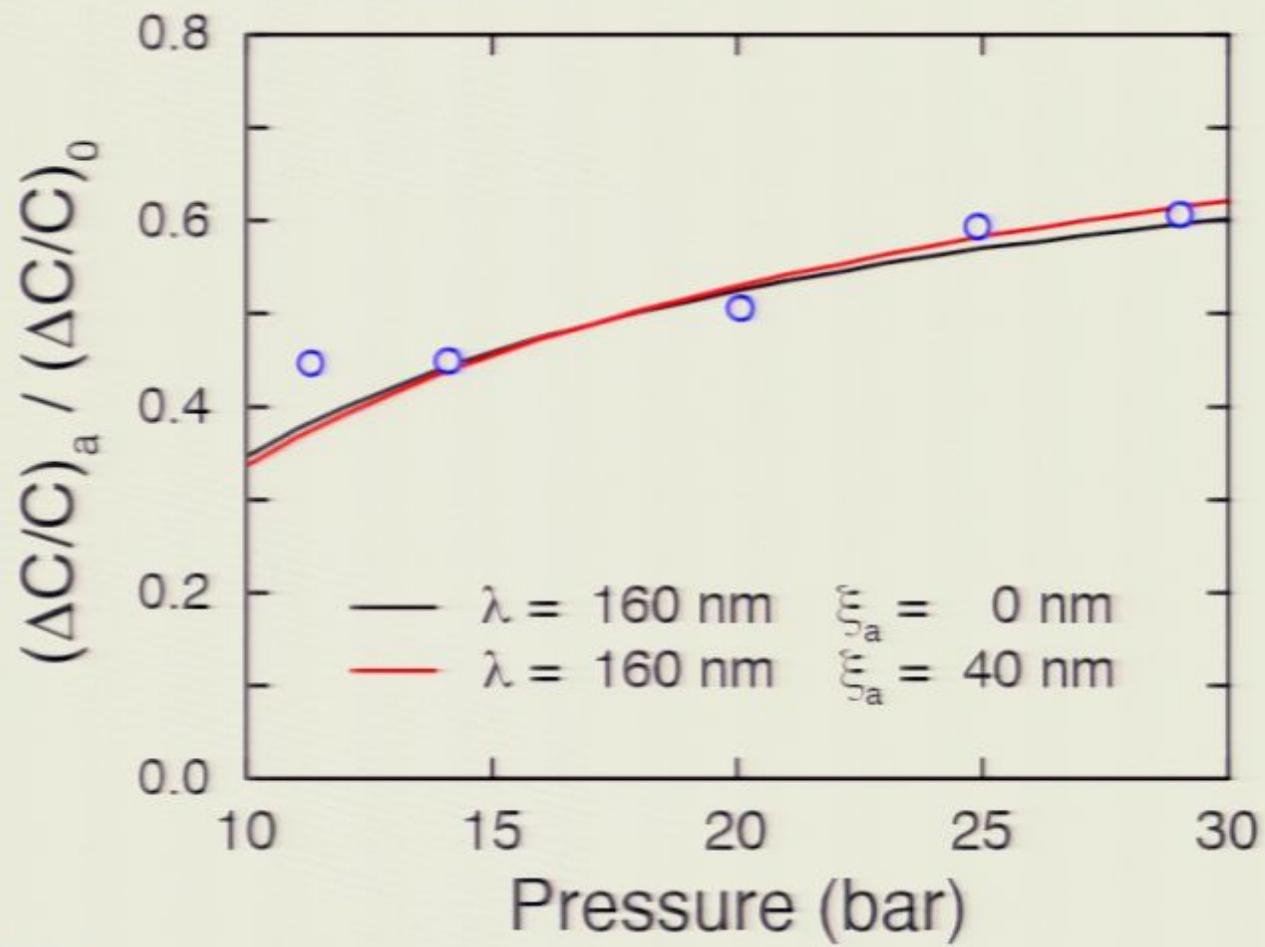
Phase diagram with correlated disorder:



$\lambda = 140 \text{ nm}$
— $\xi_a = 50 \text{ nm}$
- - - $\xi_a = 0 \text{ nm}$

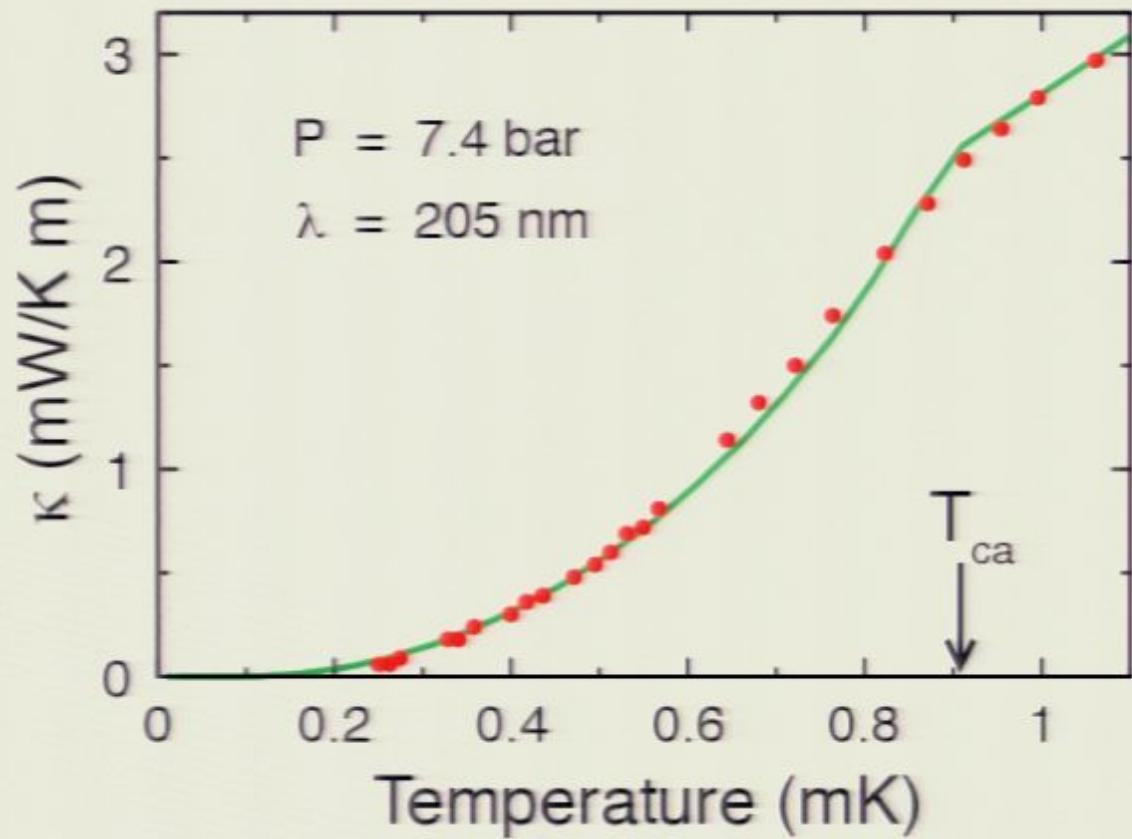
*Matsumoto et al.
PRL 79, 253 (1997);
Gervais et al. PRL 87,
035701 (2001);
PRB 66, 054528 (2002).*

Specific heat jump:



Choi et al.
PRL 93, 145301 (2004);

Thermal conductivity:



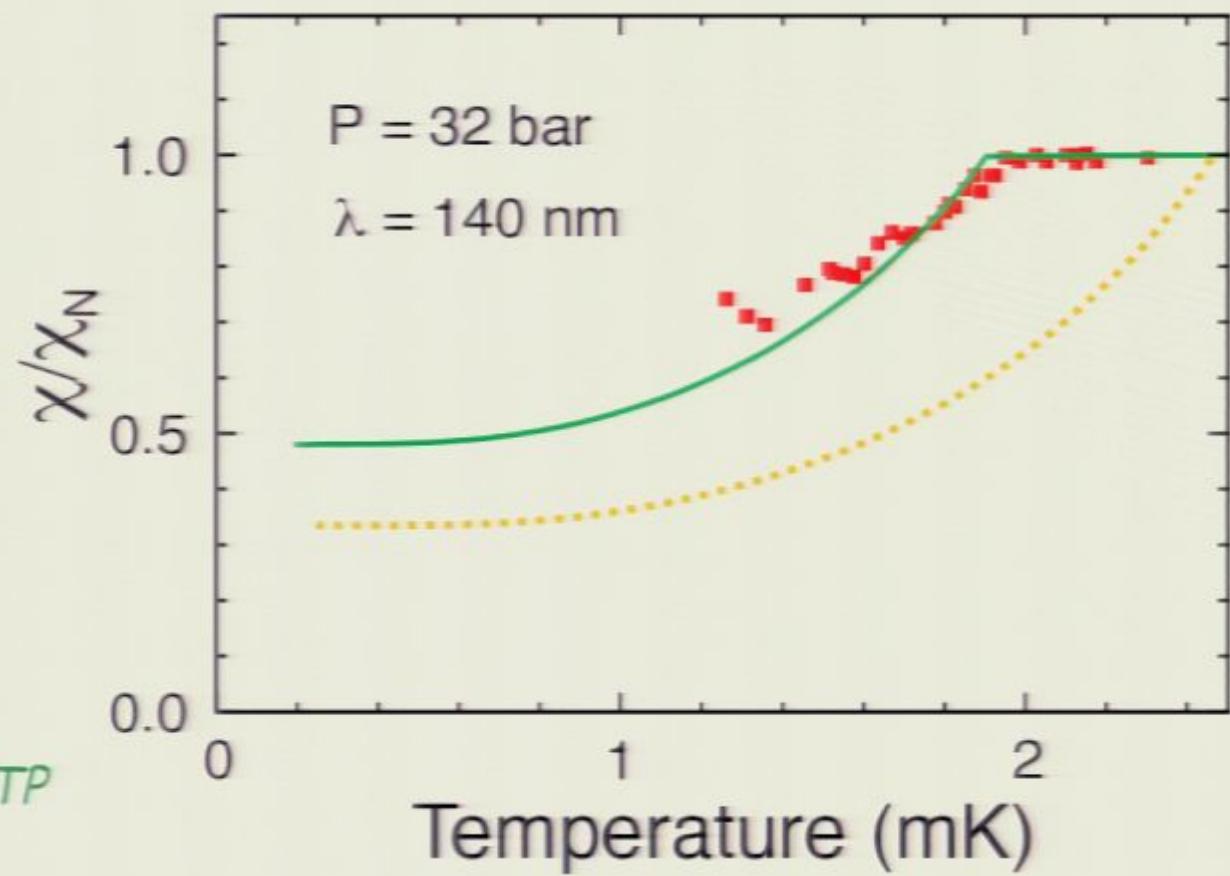
Fisher et al.

JLTP 126, 673 (2001);

Sharma and Sauls,

PhysicaB 329, 313 (2003).

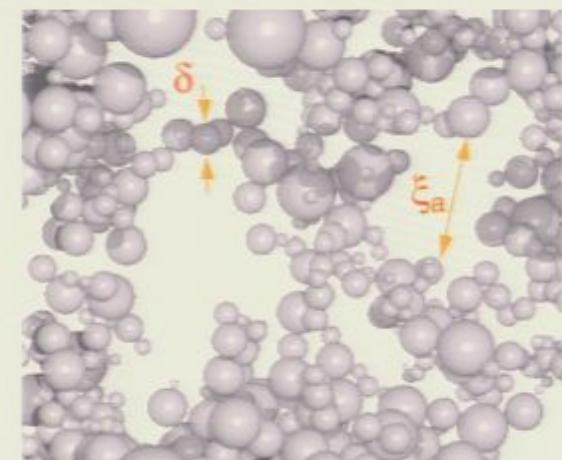
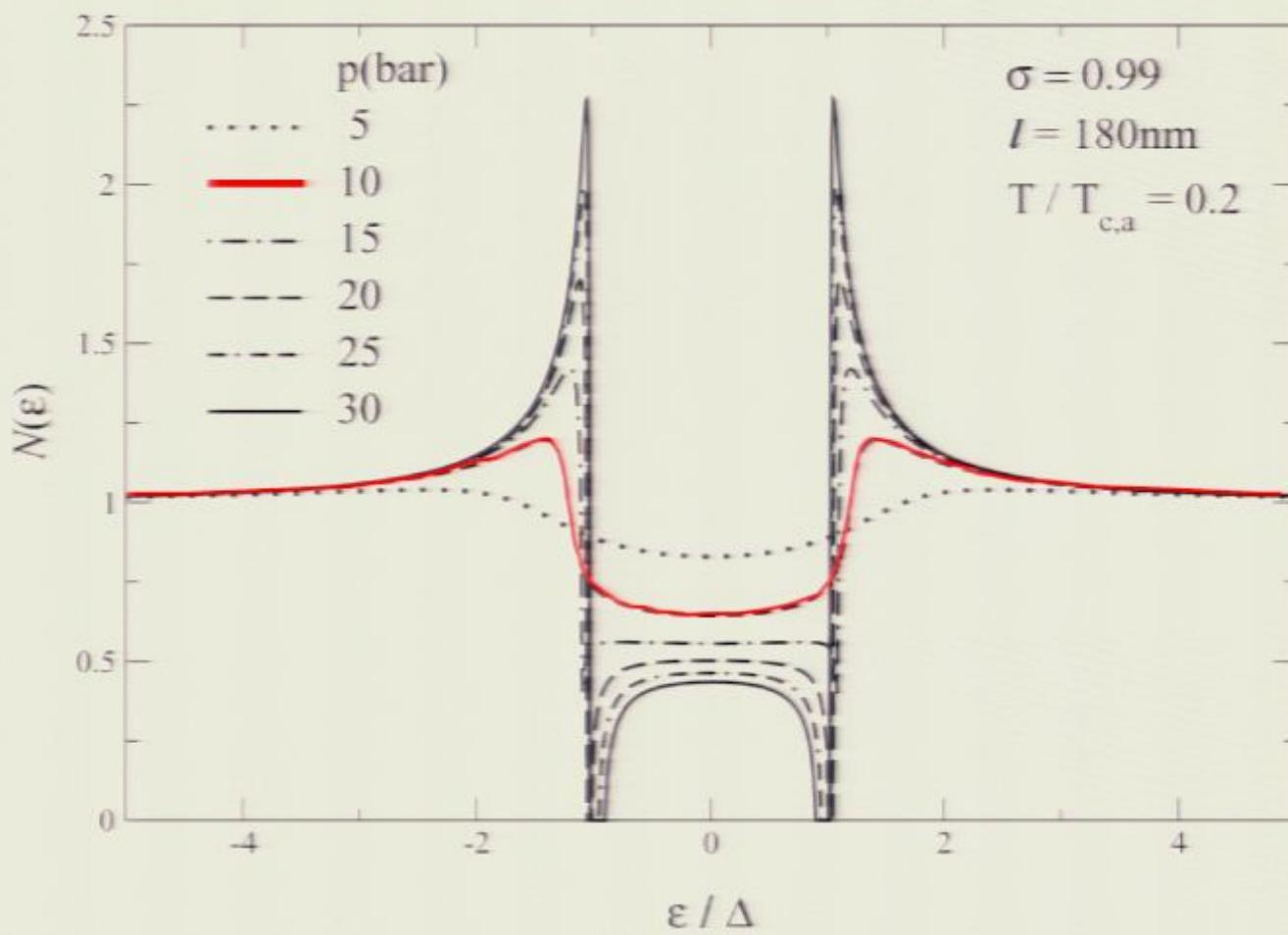
Magnetic susceptibility:



Barker *et al.*
PRL 85, 2148 (2000);
Sharma and Sauls, JLTP
125, 115 (2001).

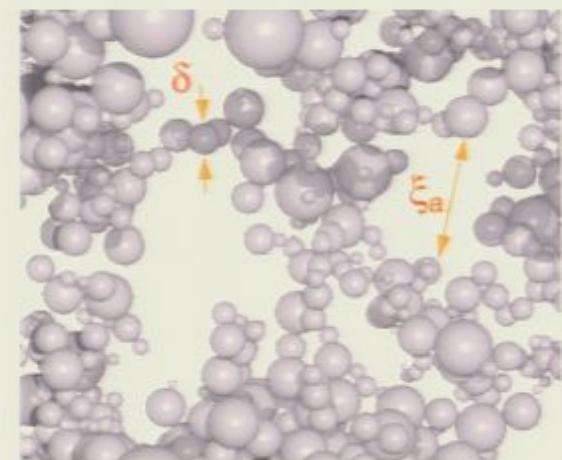
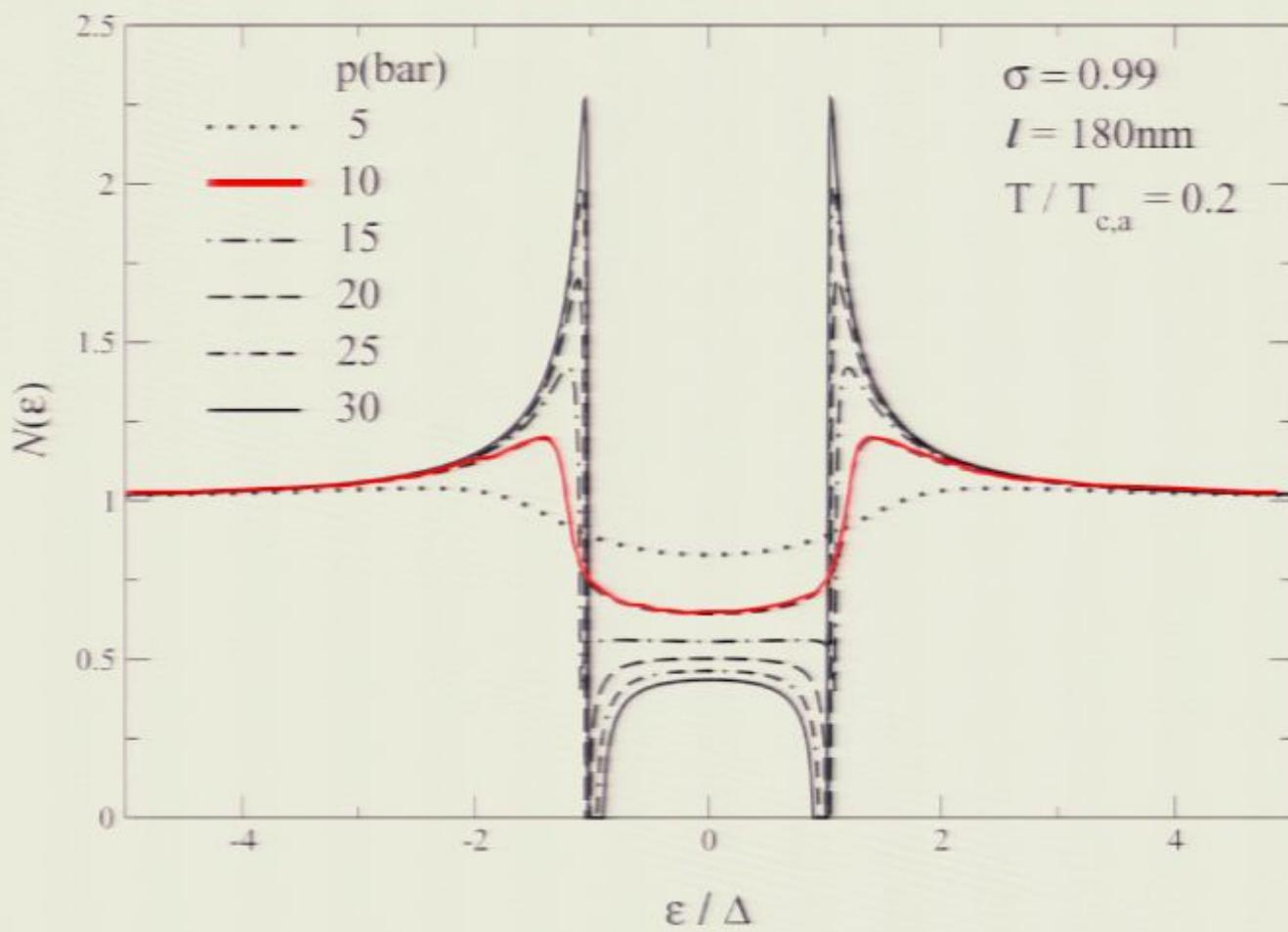
See also: Sprague *et al.* (1996); Collin (2002); Bunkov *et al.* (2005)

Gapless superfluidity: theory



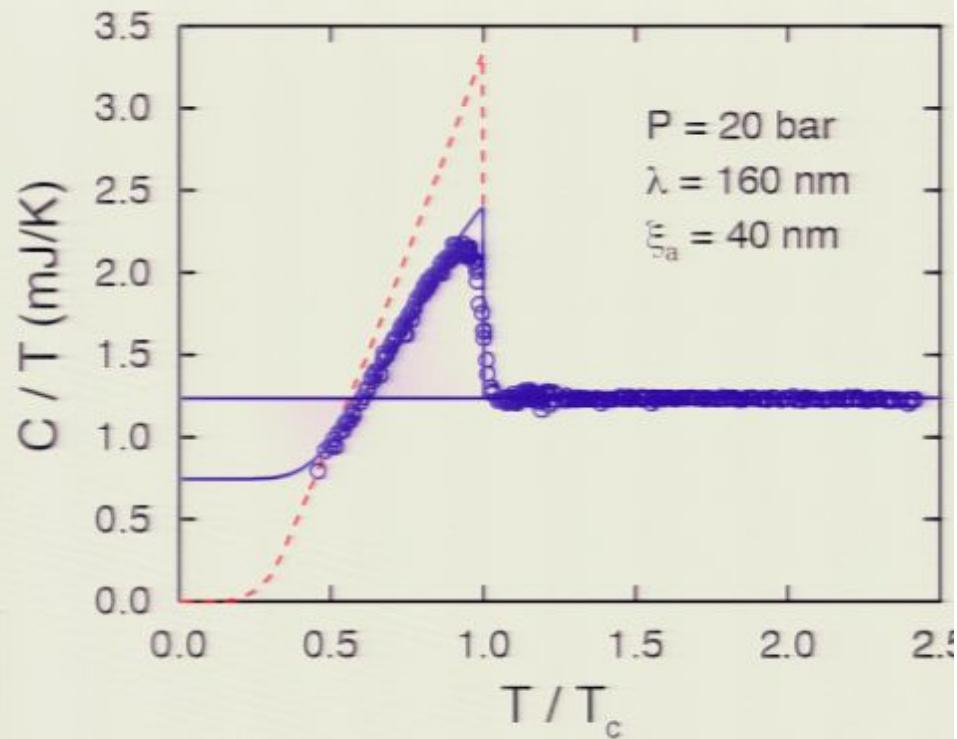
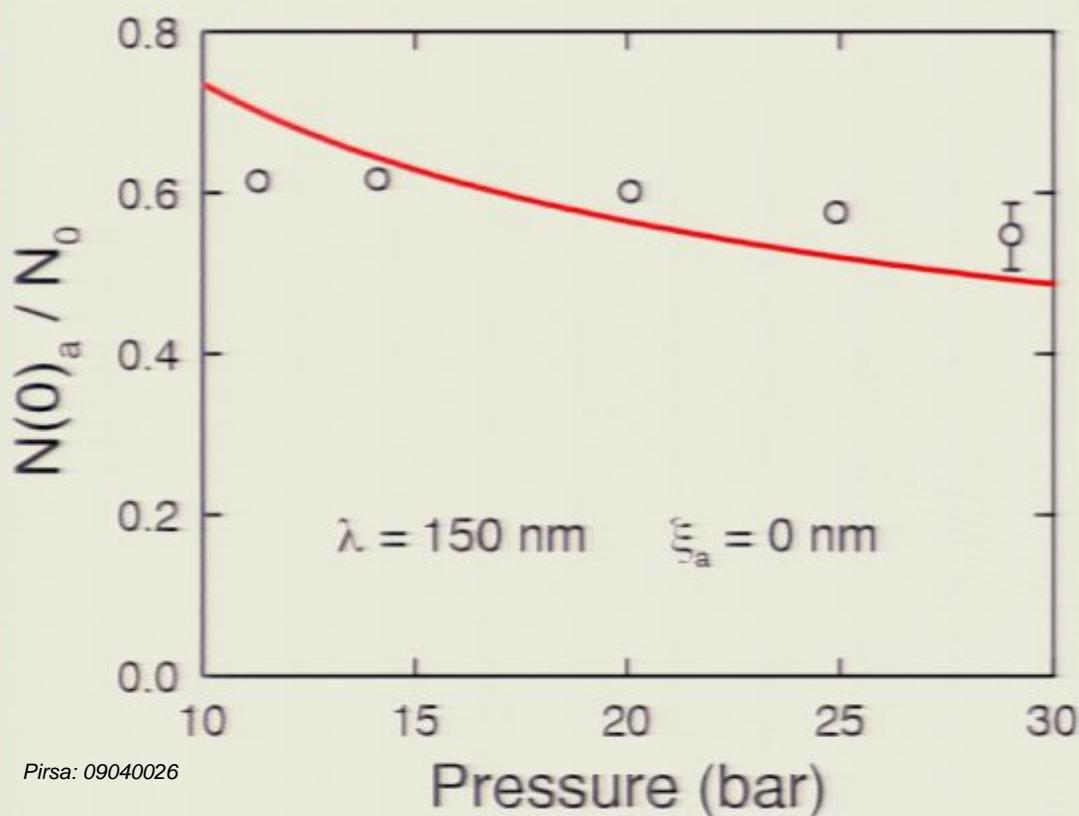
Sharma and Sauls.
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Gapless superfluidity: theory



Sharma and Sauls.
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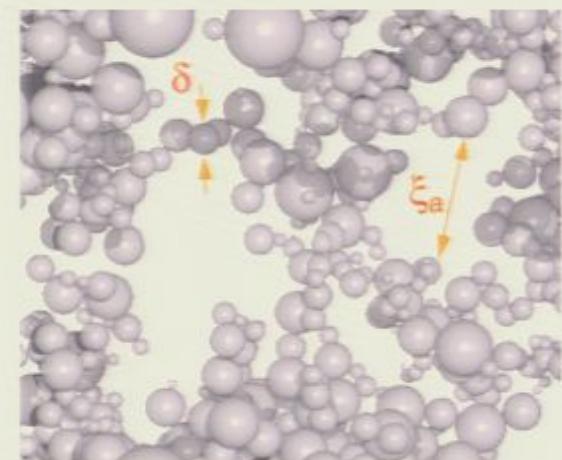
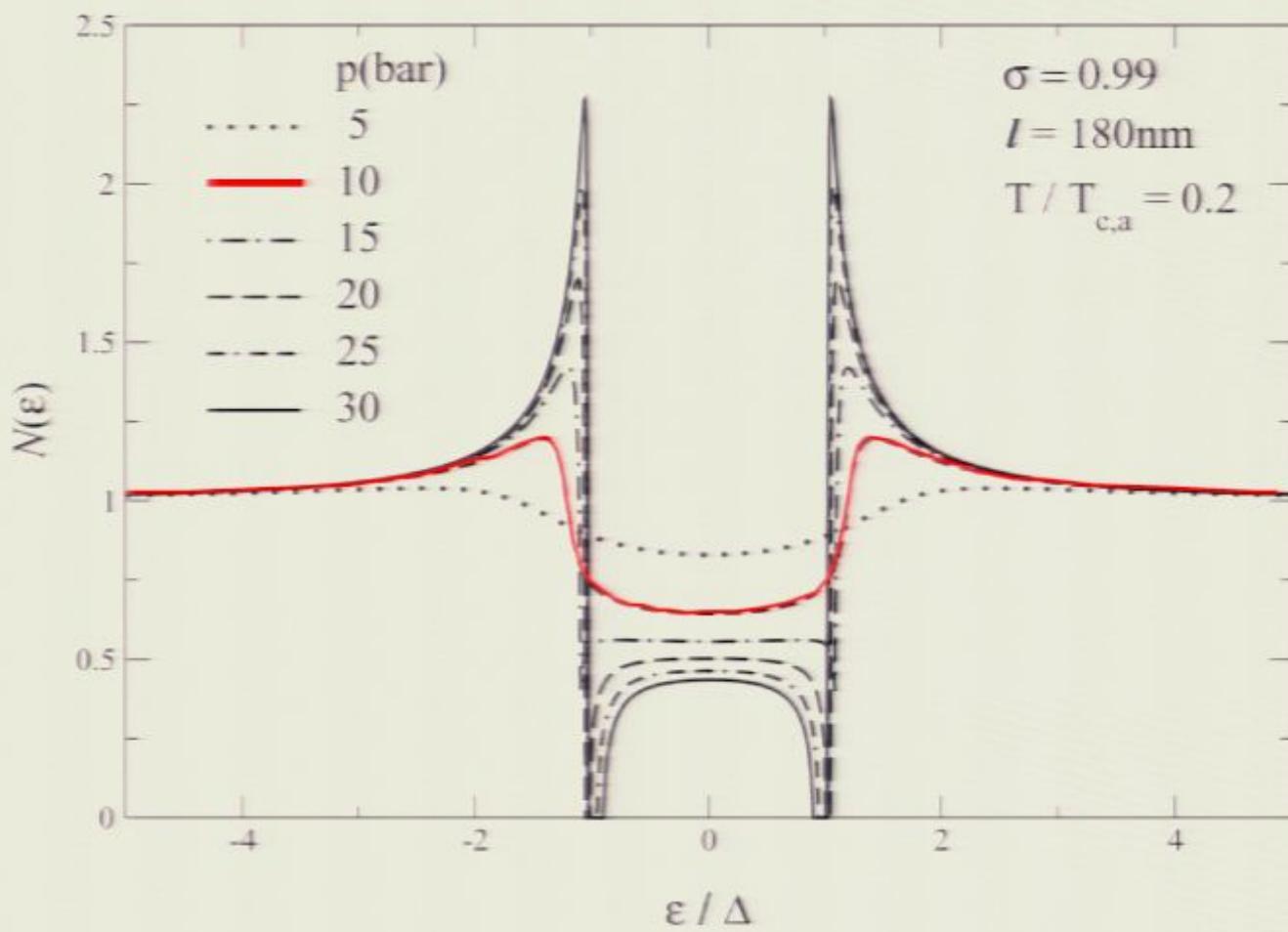
Gapless superfluidity $C(T)$:



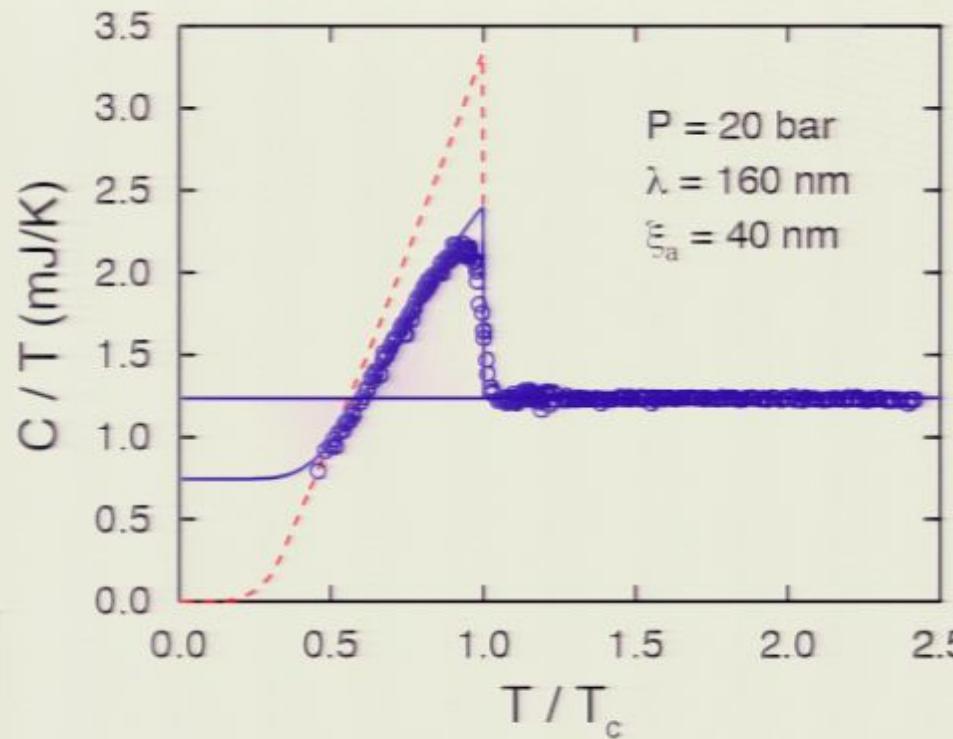
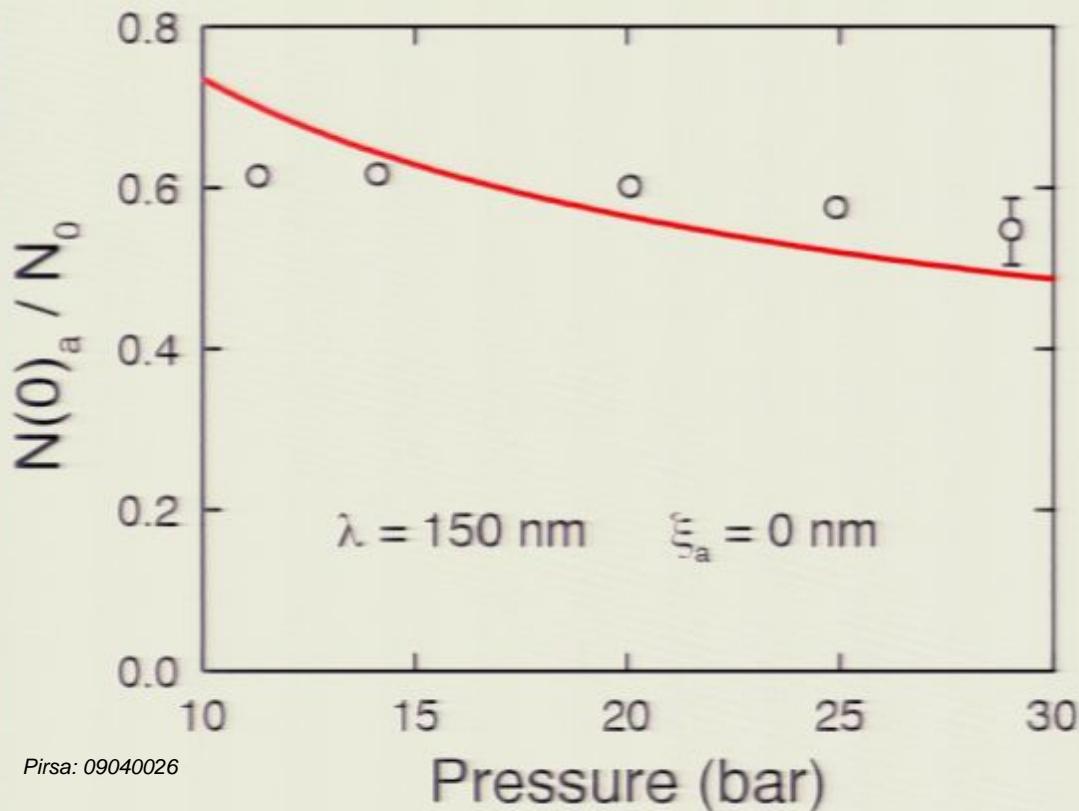
*Choi et al. PRL 93, 145301
(2004);*

*Also, attenuation of sound,
 $a(T=0)$, H.C. Choi et al. PRL
(2007)*

Gapless superfluidity: theory



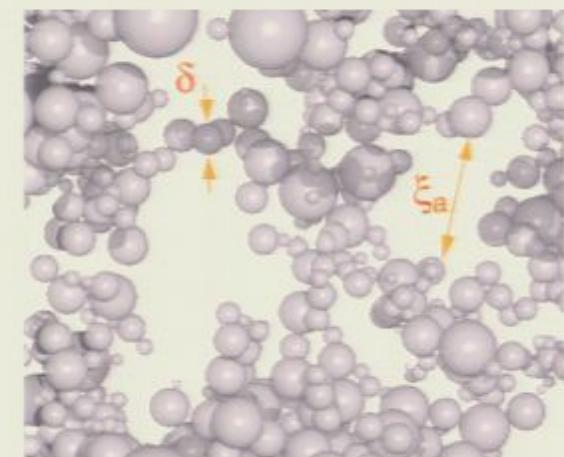
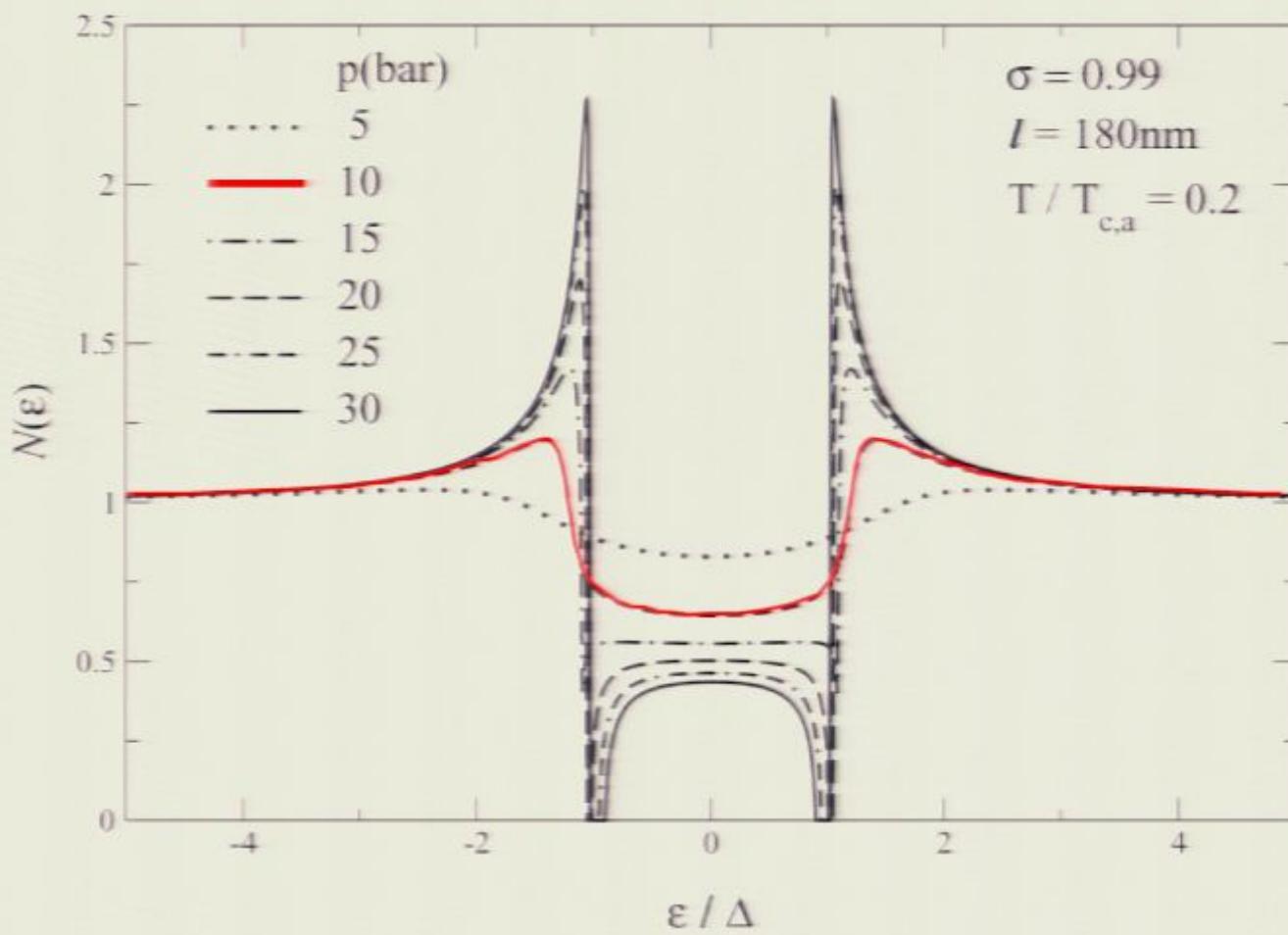
Gapless superfluidity $C(T)$:



*Choi et al. PRL 93, 145301
(2004);*

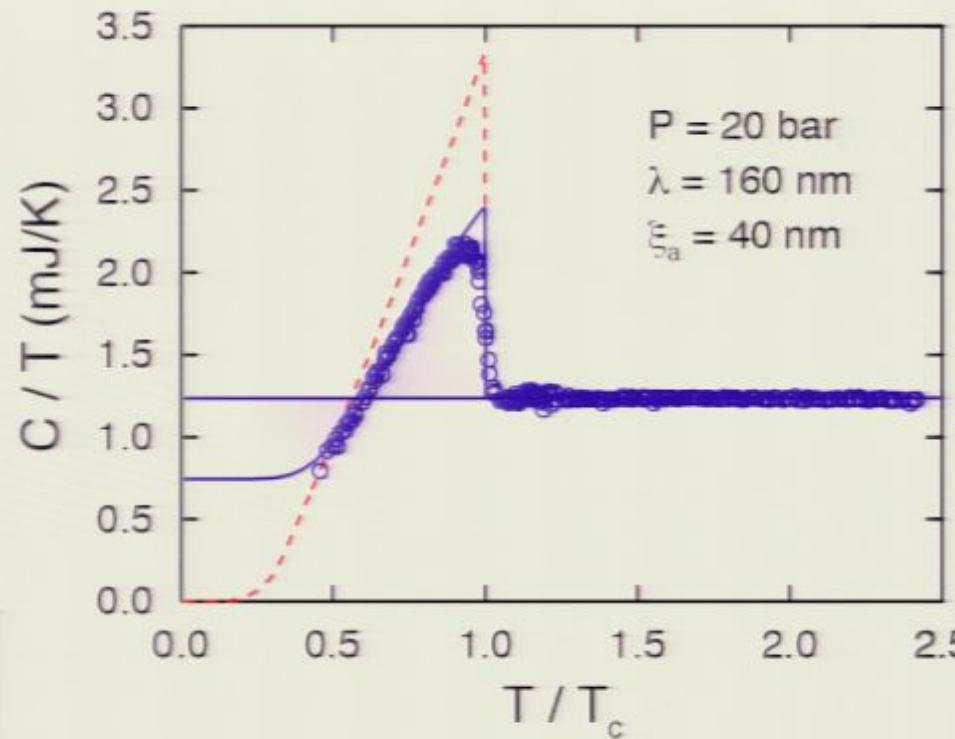
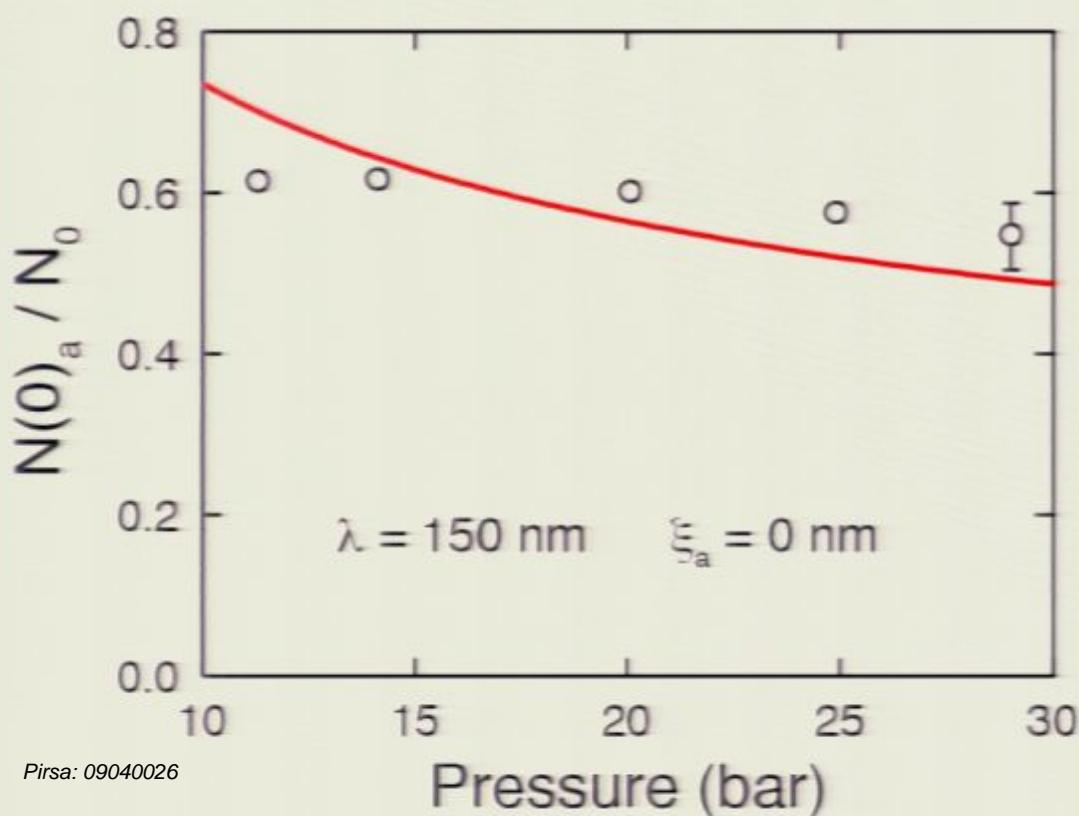
*Also, attenuation of sound,
 $a(T=0)$, H.C. Choi et al. PRL
(2007)*

Gapless superfluidity: theory



Sharma and Sauls.
JLTP 125, 115 (2001);

Gapless superfluidity $C(T)$:



*Choi et al. PRL 93, 145301
(2004);*

*Also, attenuation of sound,
 $a(T=0)$, H.C. Choi et al. PRL
(2007)*

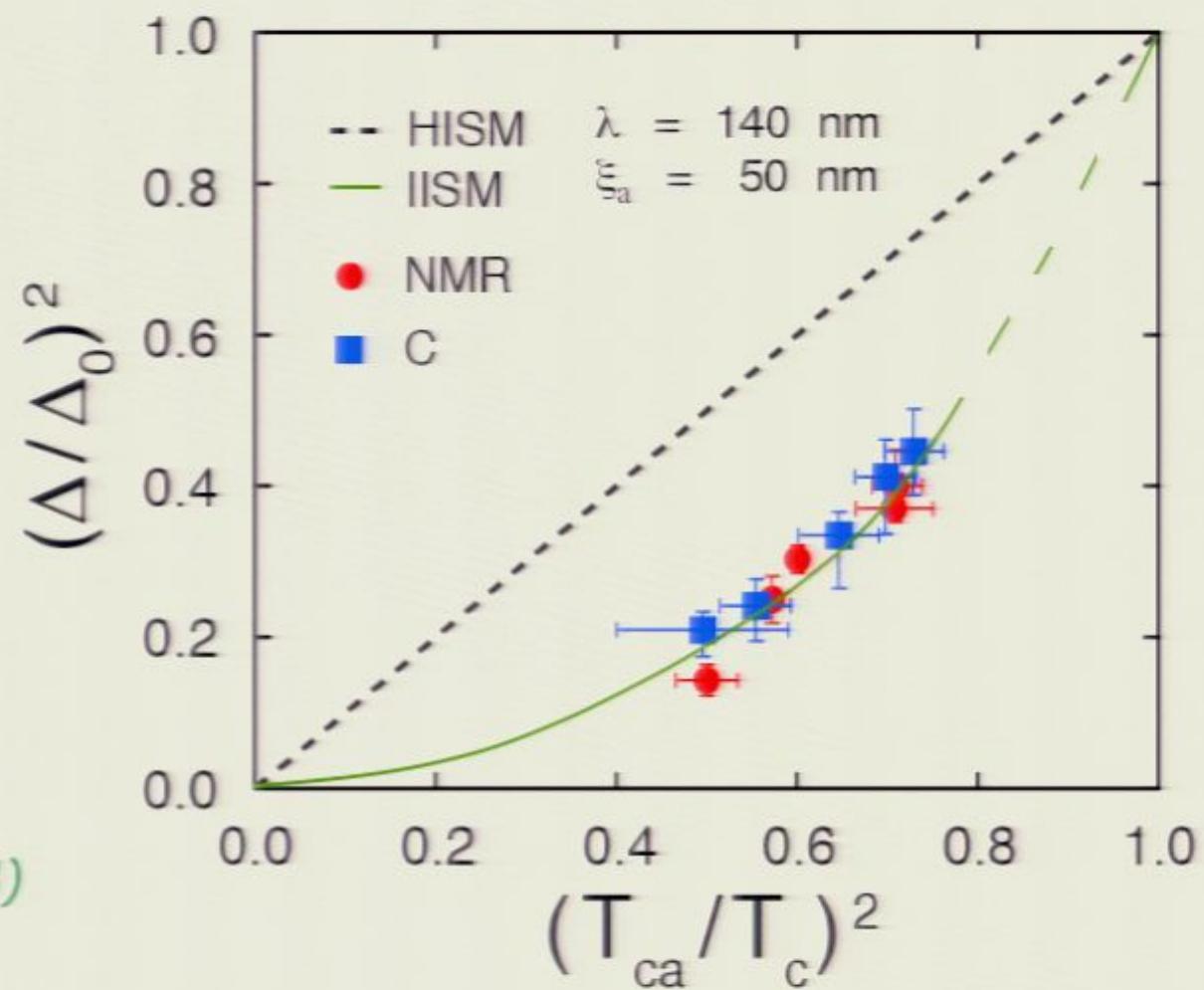
Success of IISM:

$$\Delta\omega_{\text{NMR}} \propto \Delta^2$$

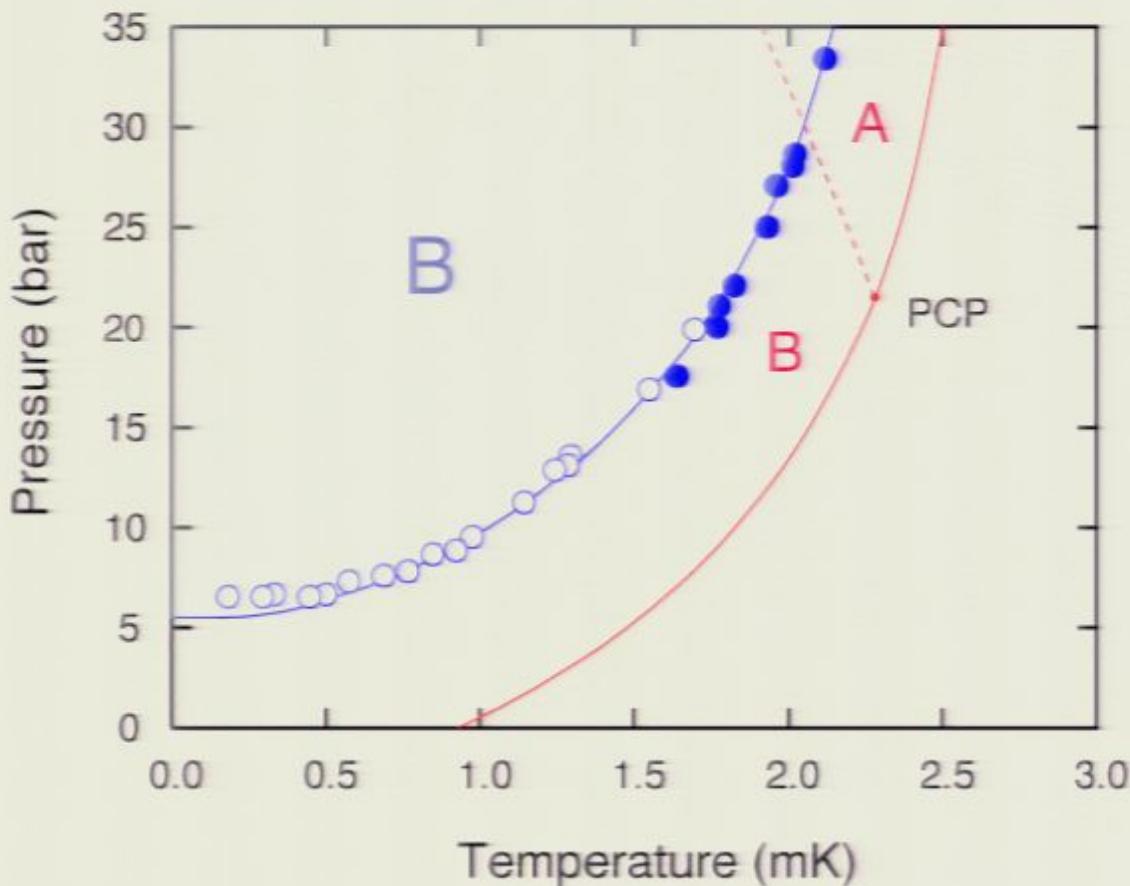
$$\frac{\Delta C}{C} \propto \Delta^2$$

Halperin et al.

JPSJ 77 111002 (2008)



Superfluid ^3He in aerogel phase diagram

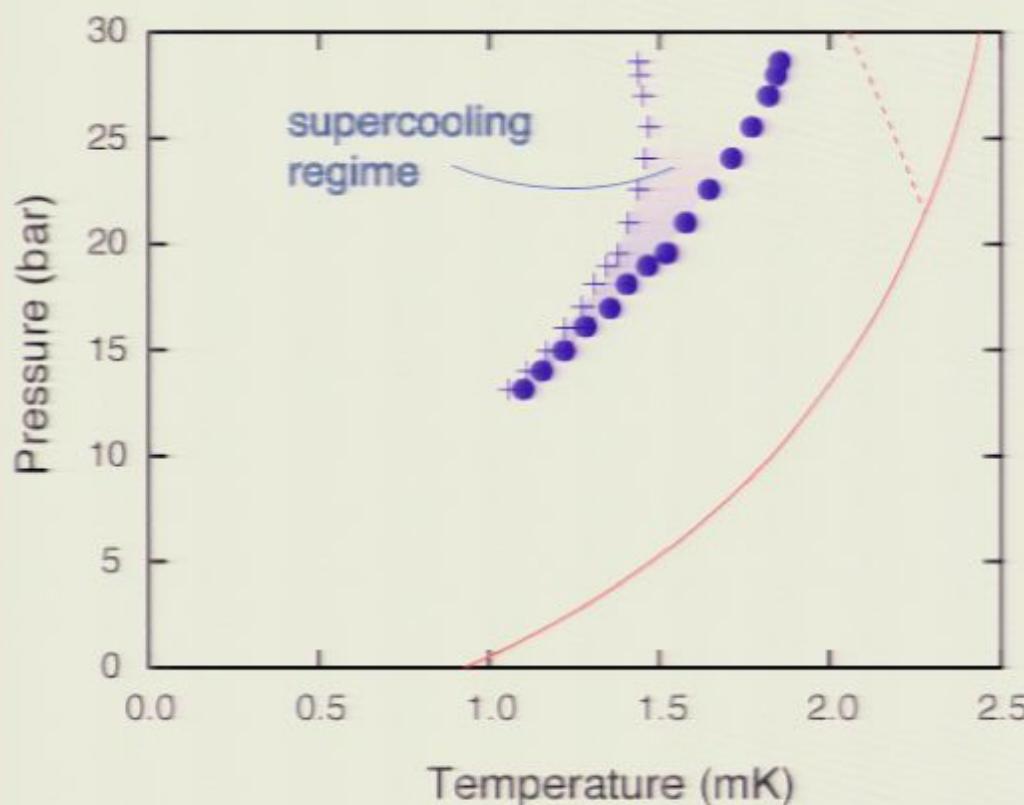


PCP ?

What happened to
the polycritical
point?

It's pushed up to
high pressure.

Metastable phase:



Nazaretski et al., JETP Lett.
79, 383 (2004).

The metastable A-like phase is an equal spin pairing state (ESP). But its orbital structure is not known.

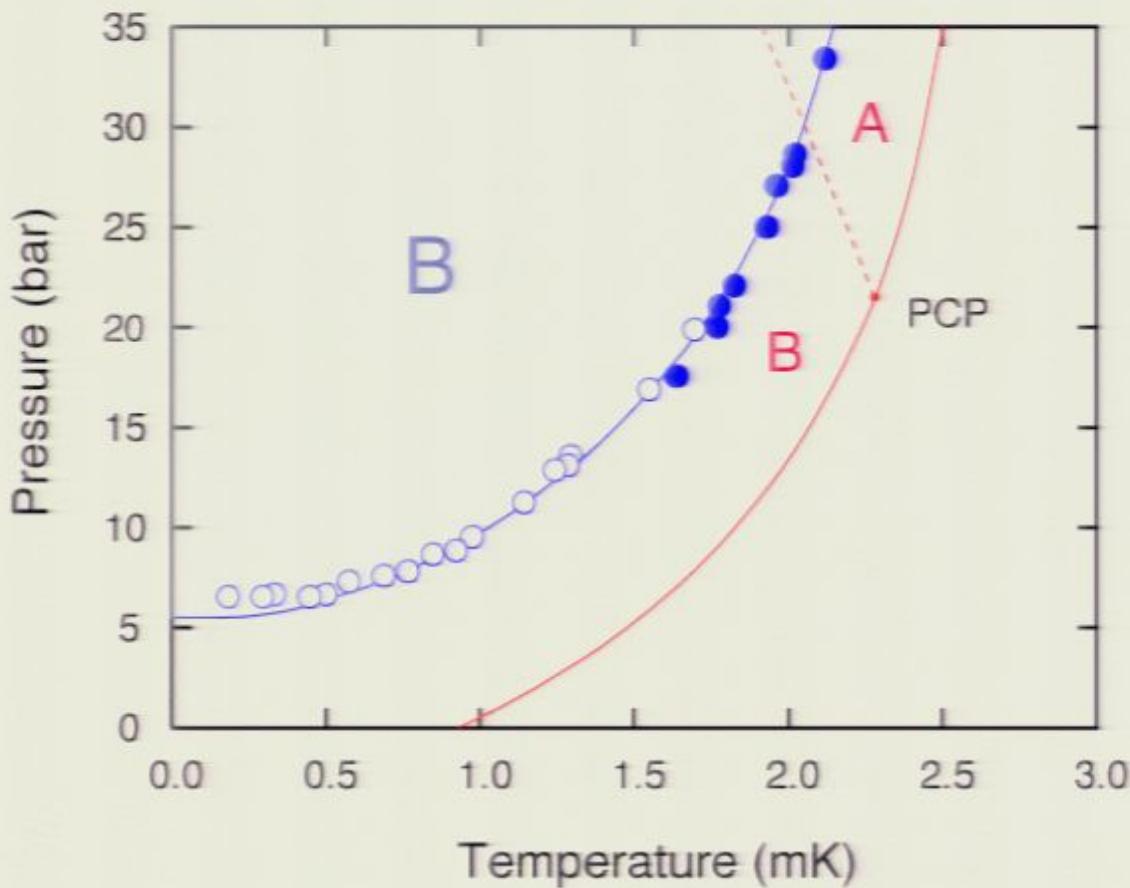
It could be the axial state like the A-phase (anisotropic) without long range orientational order, a Larkin-Imry-Ma state.

Volovik JETP Lett 84, 455 (2006).

It could be a quasi-isotropic (robust) state with long range order

Fomin JETP Lett 84, 624 (2006); Baramidze and Kharadze JLTP 150, 710 (2008)

Superfluid ^3He in aerogel phase diagram

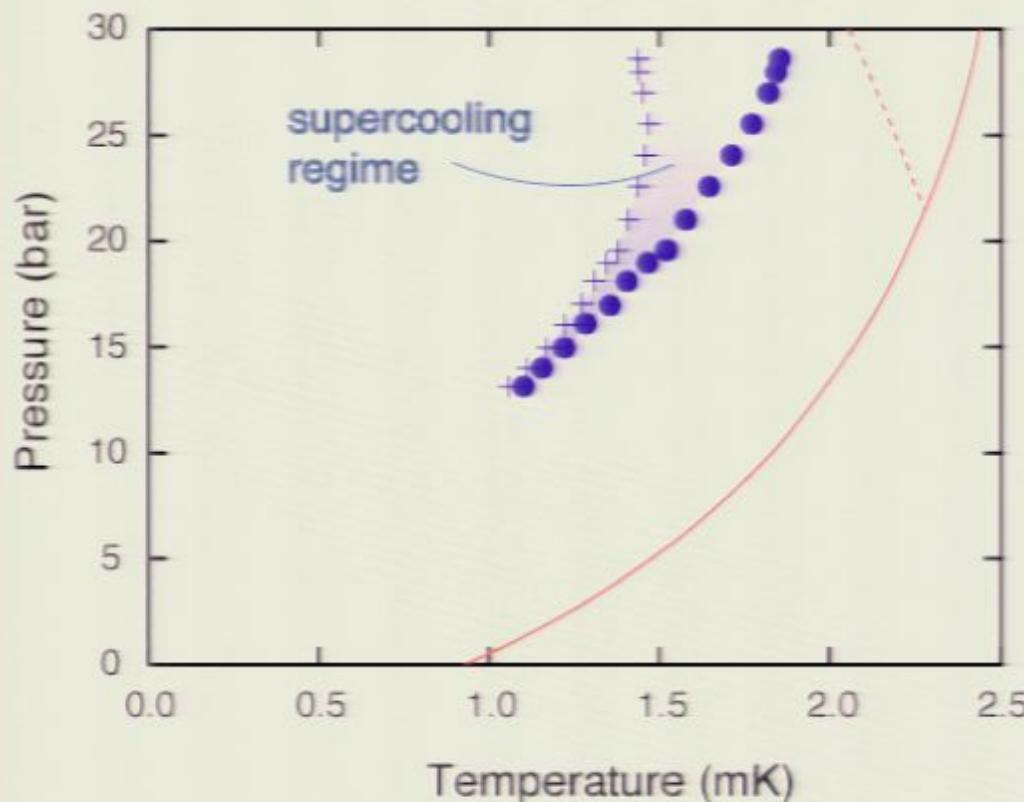


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Anisotropy:

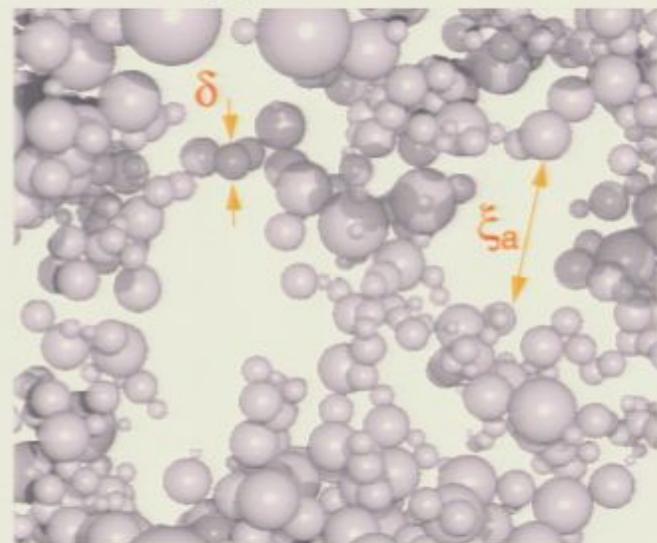
Anisotropic scattering favors anisotropic states (axial). Isotropic scattering favors isotropic phases:
Thuneberg et al. PRL 80, 2861 (1998).

The axial (LIM) state couples strongly to the anisotropic aerogel structure.

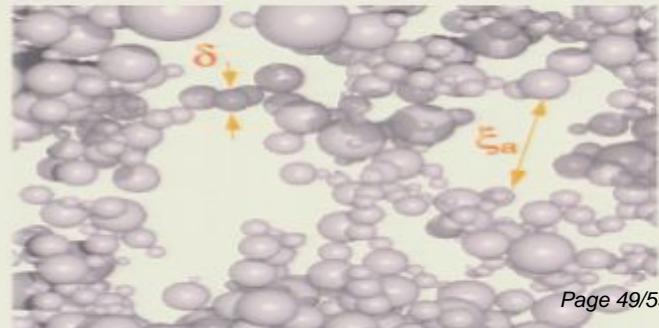
Can global anisotropy stabilize the A-like phase?

*Vicente et al.
PRB 72, 075301 (2005)*

local
anisotropy



global
anisotropy



Global Anisotropy- new phases :

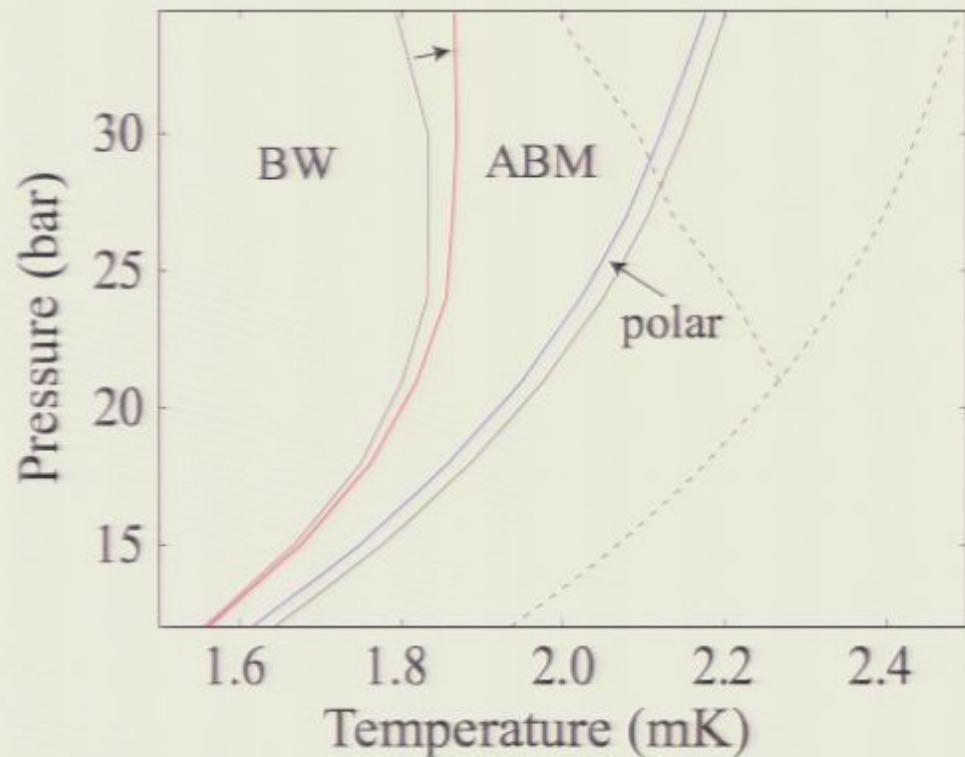
*Aoyama and Ikeda PRB 76
104512 (2007).*

Anisotropic scattering will stabilize (new) anisotropic states

Axial compression stabilizes the axial state

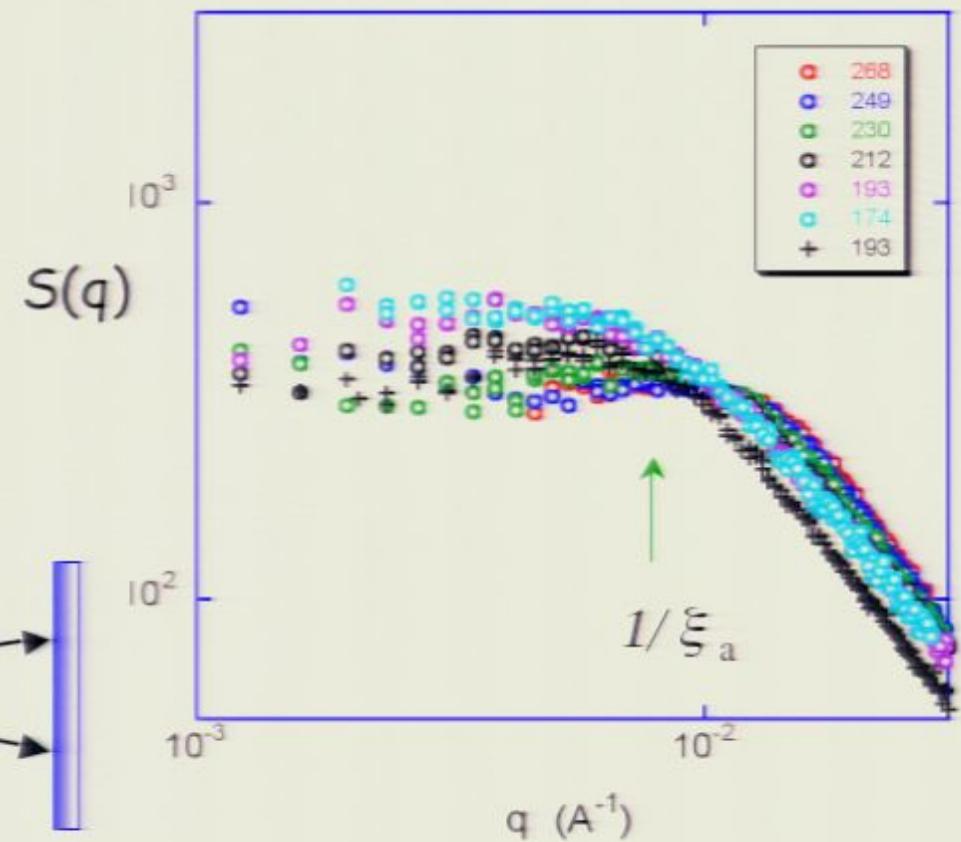
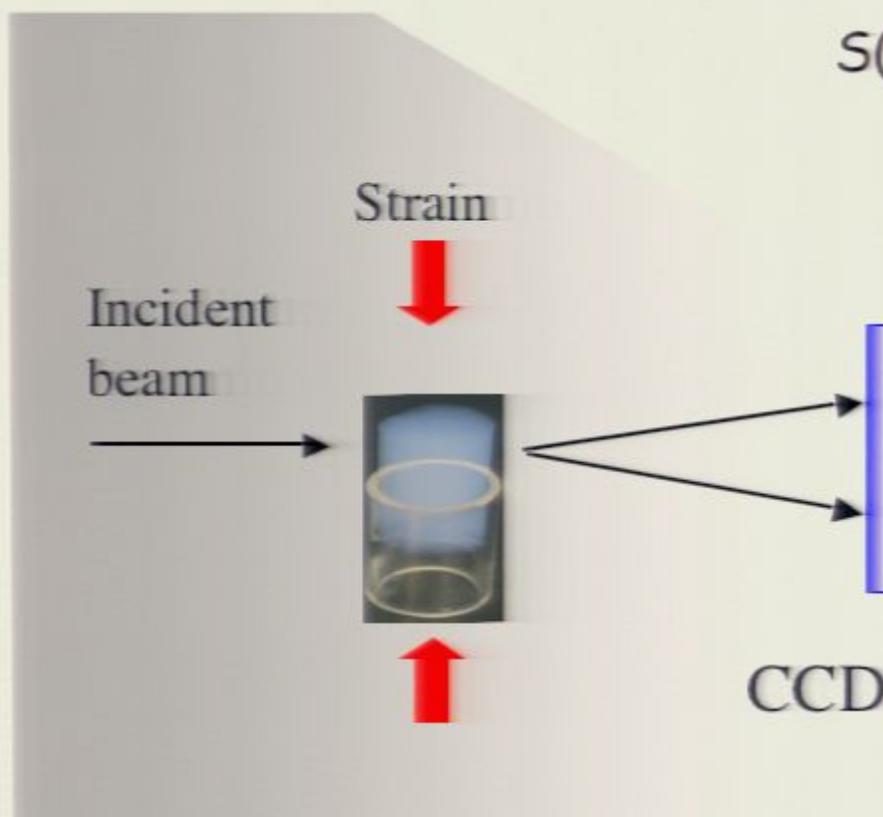
Radial compression stabilizes the polar state

radially compressed



Controlled homogeneous anisotropic scattering:

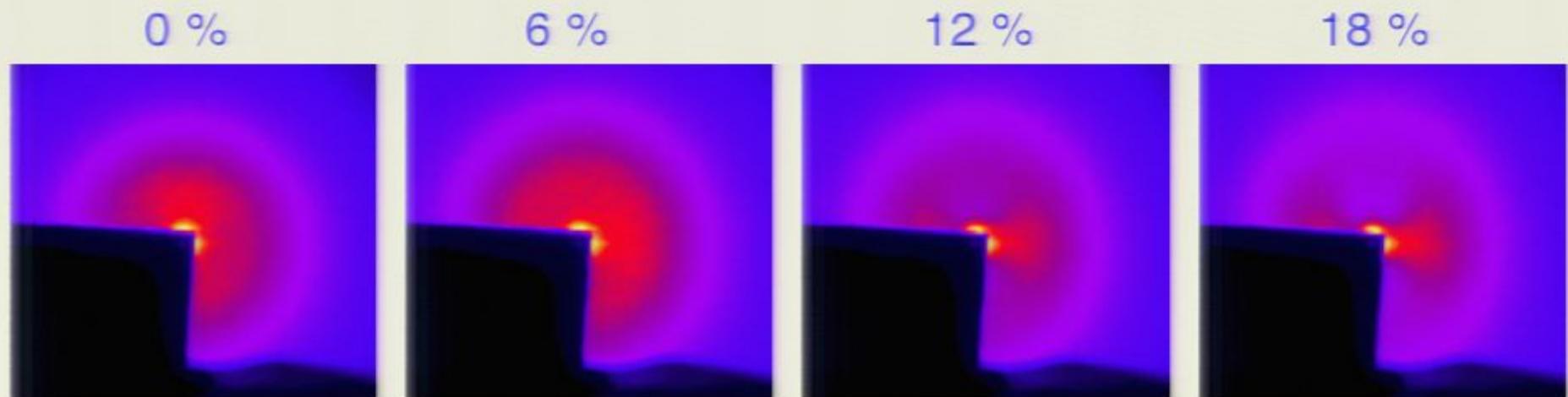
SAXS: Advanced Photon Source (ANL)



$$\frac{\lambda_{\perp}}{\lambda_{\parallel}} = \frac{\xi_{a\perp}}{\xi_{a\parallel}} \frac{2}{(1 + \xi_{a\perp}/\xi_{a\parallel})}$$

Anisotropy (axial compression):

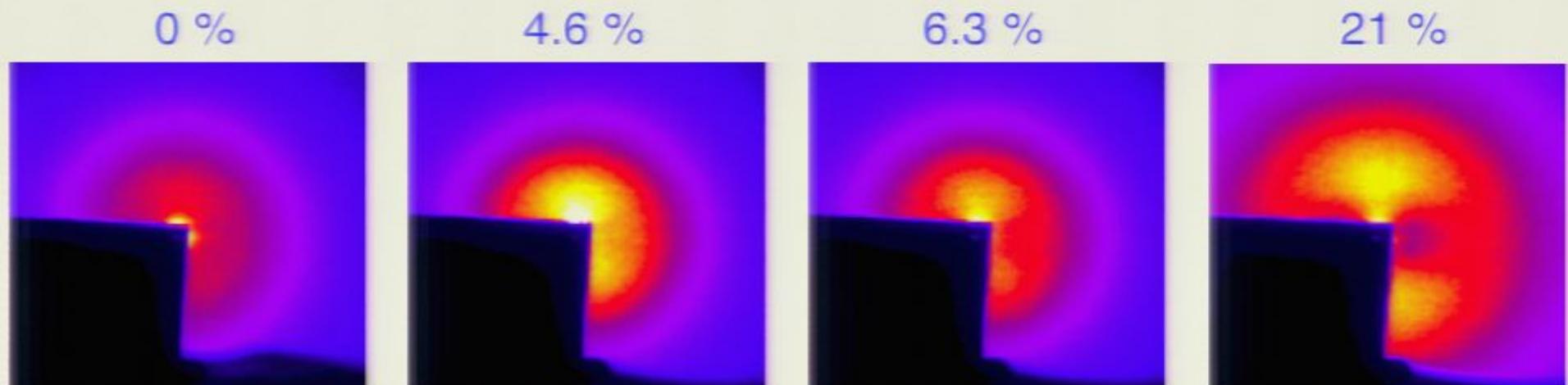
SAXS



Pollanen et al. JNCS 354 4668 (2008).

Anisotropy (radial compression):

SAXS

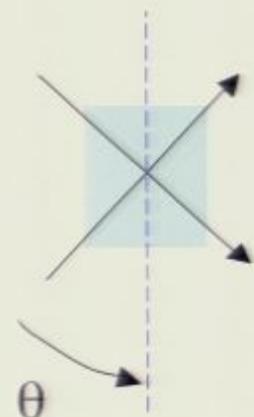


Pollanen et al. JNCS 354 4668 (2008).

Aerogel anisotropy characterization by optical birefringence:

crossed polarizer
rotation

20.9% axial strain



12.7% radial compression

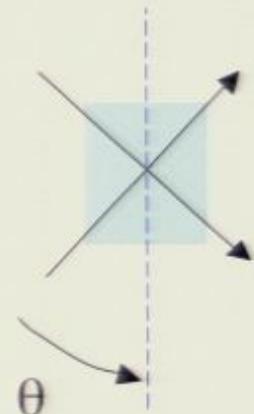


Aerogel anisotropy characterization by optical birefringence:

Sample (N.Mulders) from
ISSP NMR experiment

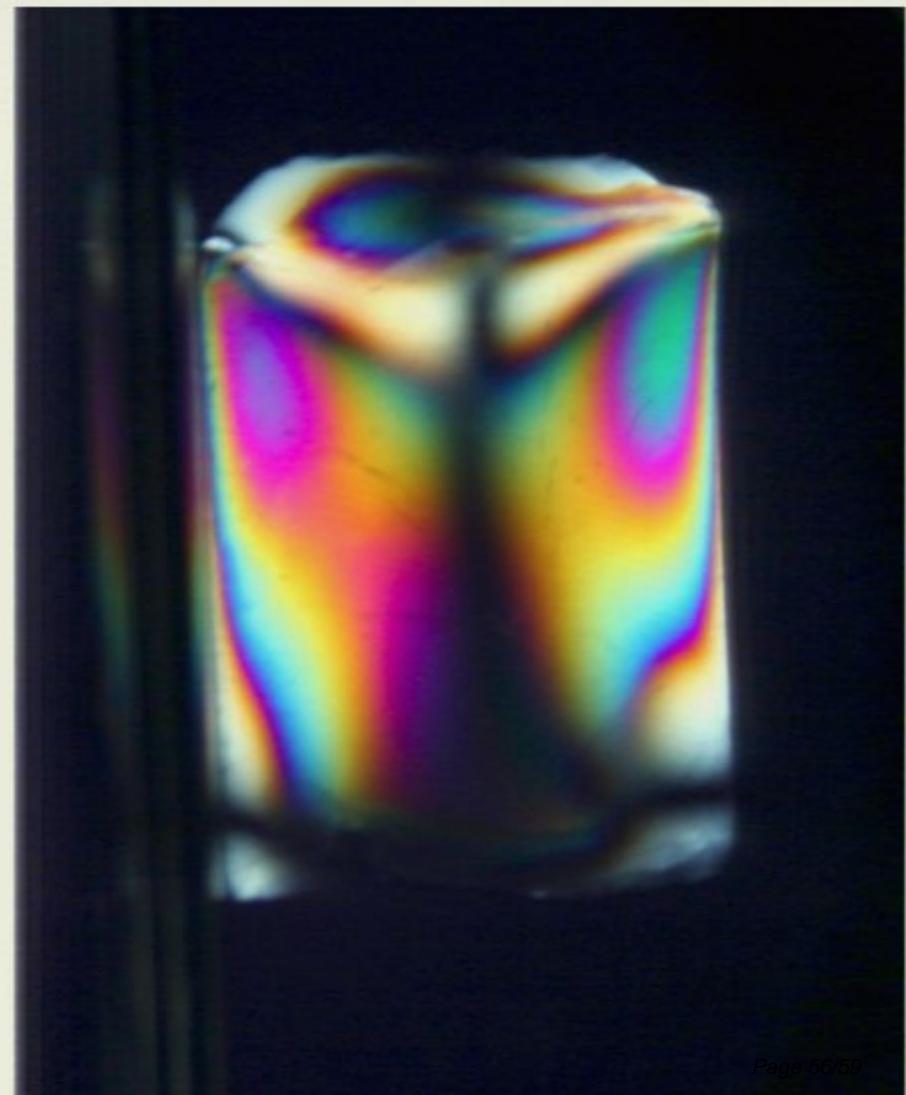
crossed polarizer at
45° to cylinder axis

Rotation about the
cylinder axis



Inhomogeneous aerogel:

*....inhomogeneous
98% aerogel*

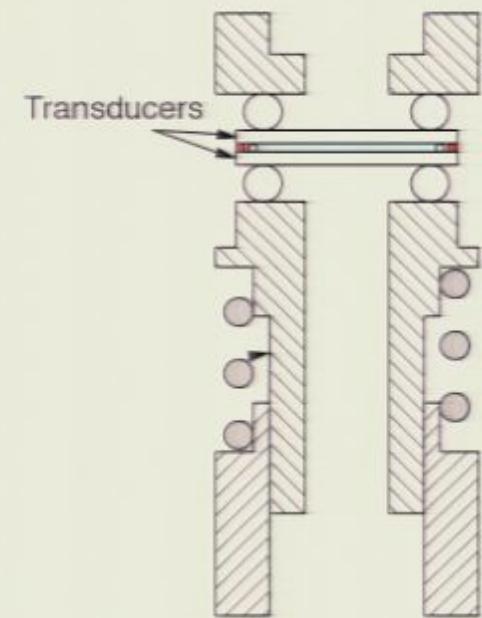
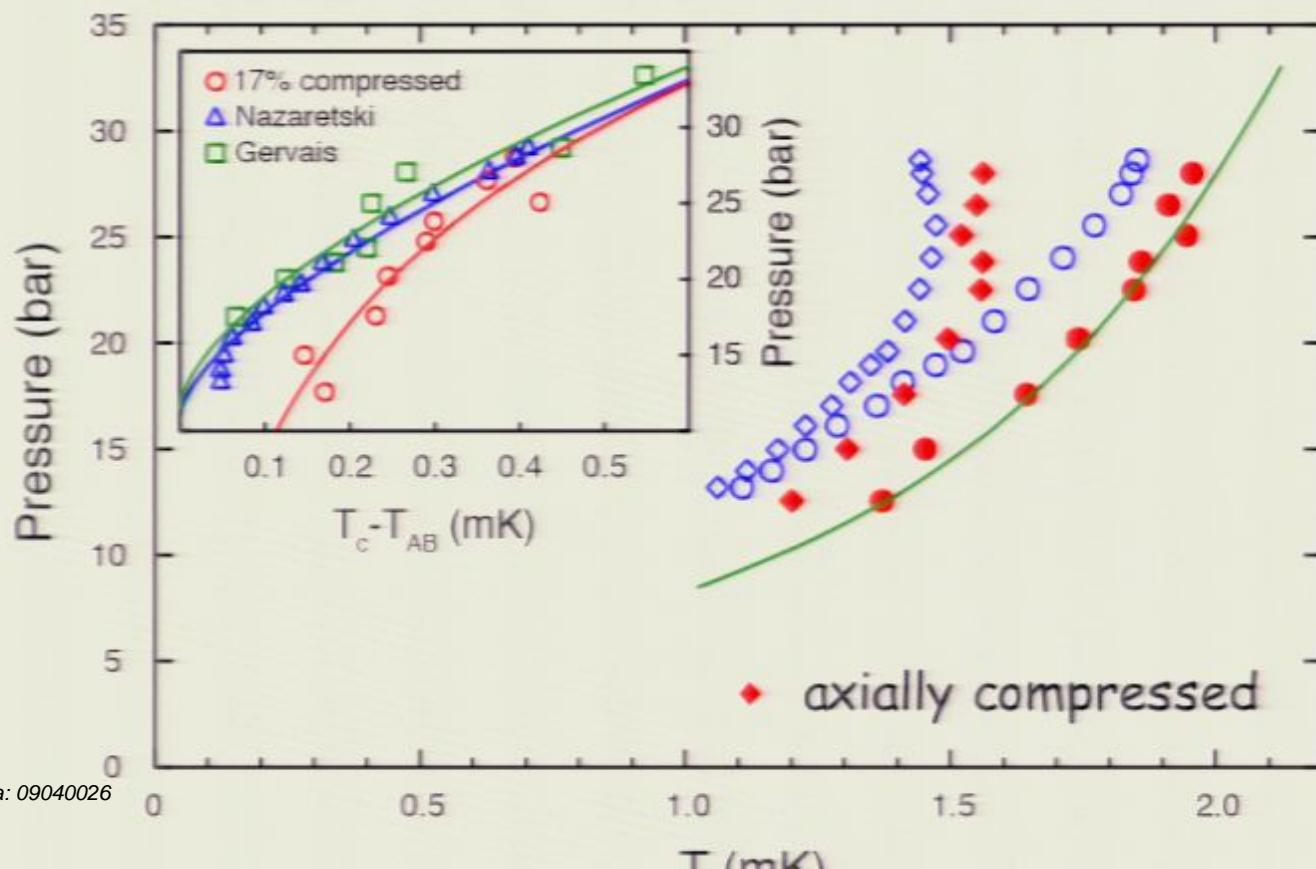


A-like phase stability: axial compression

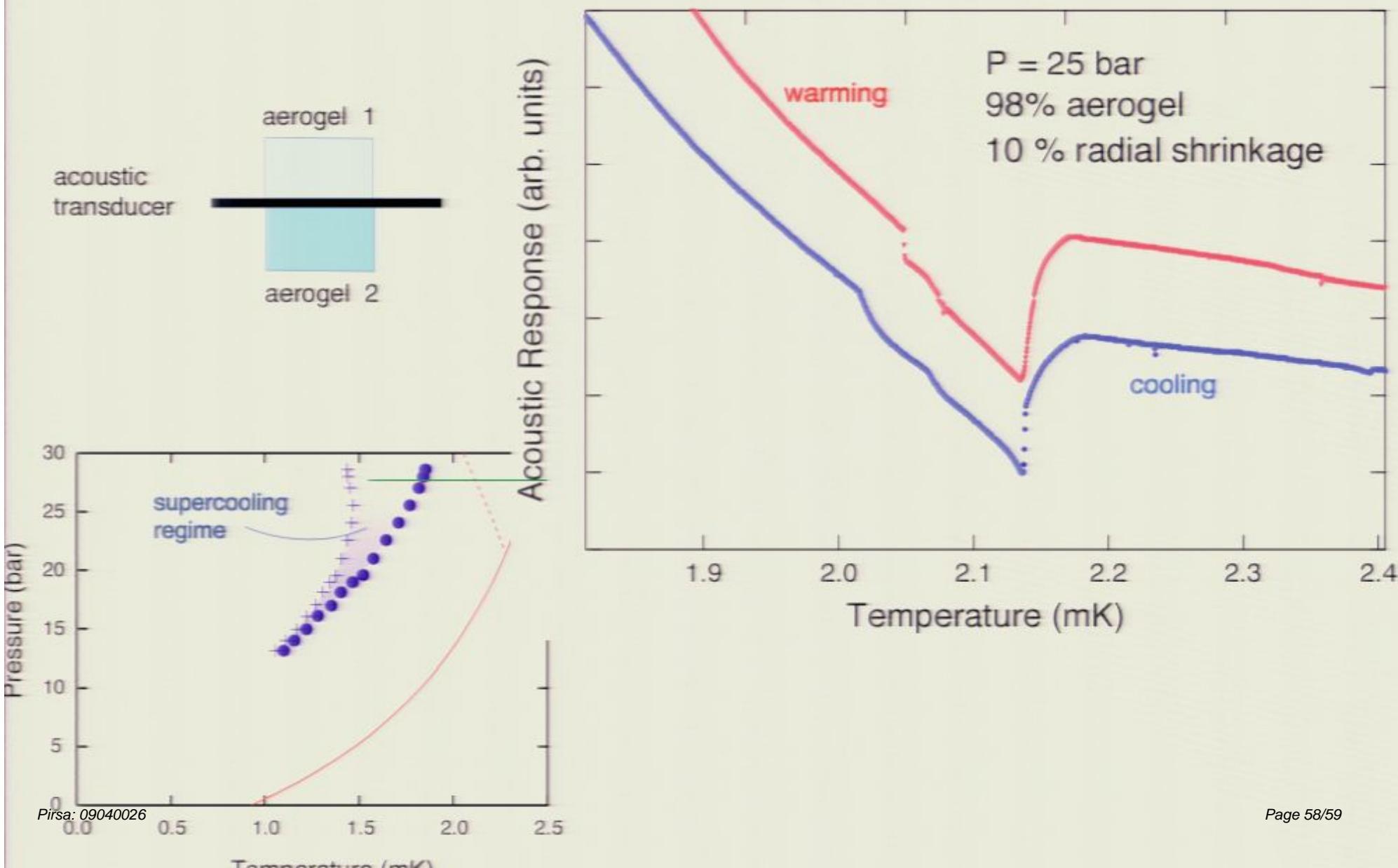
Test of the Vicente *et al.* idea.....

17% axial compression

- *Davis et al.*
PRB 100, 215304 (2008).



A-like phase stability: radial compression



Summary of ^3He in aerogel:

- Correlated impurity scattering models (IISM) work well
- Open question about the A-like phase orbital symmetry
- New phases?
- Possible control of phase stability and nucleation with global scattering anisotropy