

Title: Quantum Device Theory at the University of Waterloo

Date: Apr 23, 2009 12:15 PM

URL: <http://pirsa.org/09040025>

Abstract:

Quantum device theory research at the University of Waterloo

F.K. Wilhelm

Institute for Quantum Computing (IQC) and Department of Physics and
Astronomy, University of Waterloo, Canada

4-corners condmat meeting 2009, Perimeter Institute

The QDT group



PI: Frank
Wilhelm



PDF: Jay
Gambetta



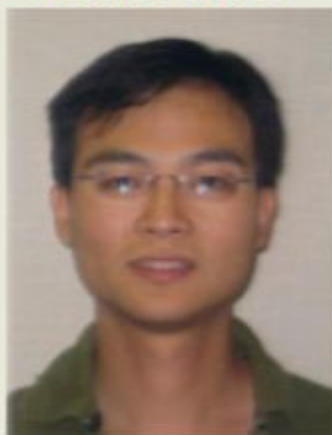
Brendan
Osberg



Felix Motzoi



PDF: Bill
Coish



Friend:
Tzu-Chieh
Wei



Peter
Groszkowski



Farzad
Qassemi

Research interests

Foundation: Condensed matter theory

Main interest: devices that can be used as qubits

- Coherent spins in semiconductor quantum dots
- Quantum noise in nanostructures
- Decoherence theory
- Circuit quantum electrodynamics
- Quantum measurement theory
- Optimal control theory
- Quantum error correction

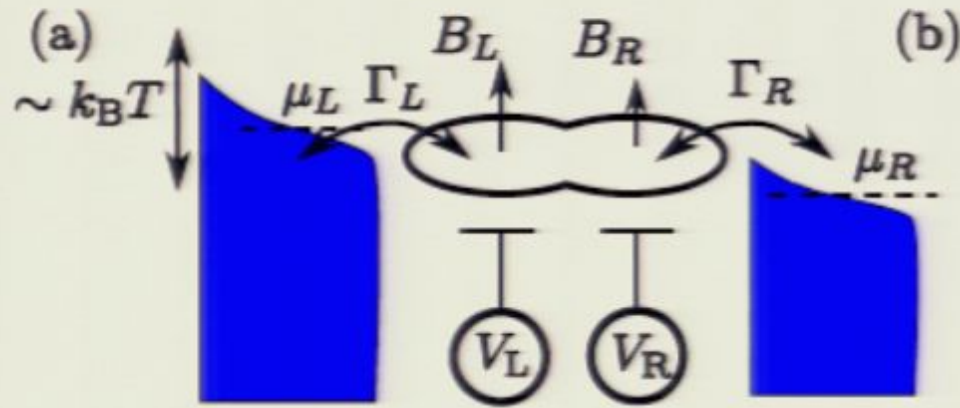
Few-body correlations far from equilibrium

Contents

- 1 Spin in quantum dots
- 2 Noise in tunnel junctions

Stationary and transient leakage current in the Pauli spin blockade

Lateral double quantum dot

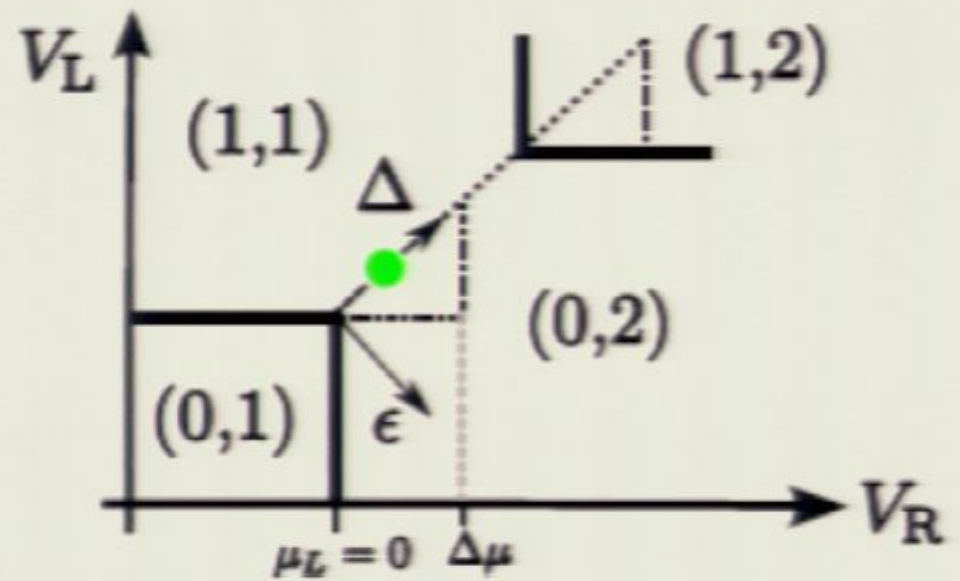


Hamiltonian

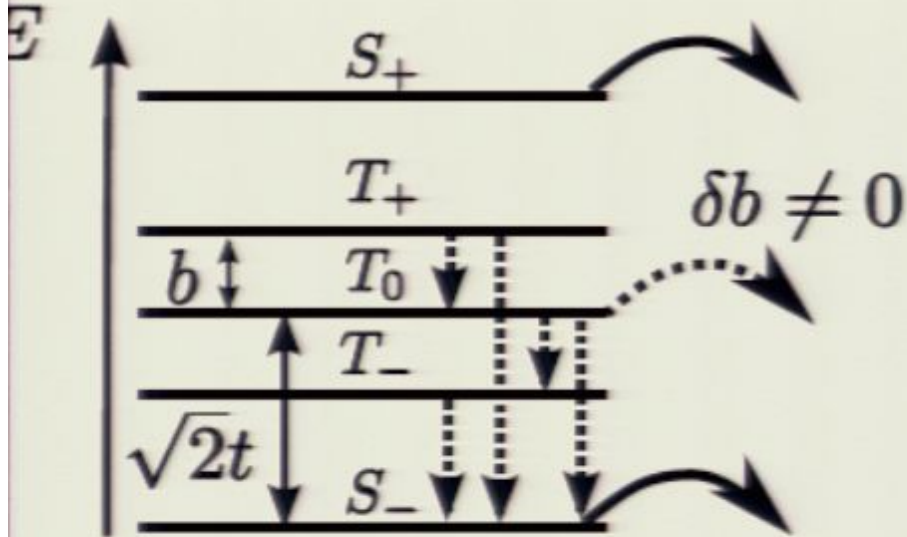
$$H = H_{DD} + H_{Lead} + H_{D-L}$$

$$H_{DD} = \sum_l \left(\frac{U}{2} n_l (n_l - 1) - V_l n_l \right) + U' n_L n_R - \sum_{\sigma} t \left(d_{L\sigma}^{\dagger} d_{R\sigma} + \text{h.c.} \right) + bS$$

Charge stability diagram



Spin structure



- T: (1,1) Triplet states
- S: Hybridized (2,0) and (1,1) singlet
Projected Hamiltonian

- Tunnel contacts *not* spin-active!
- Transport cycle:

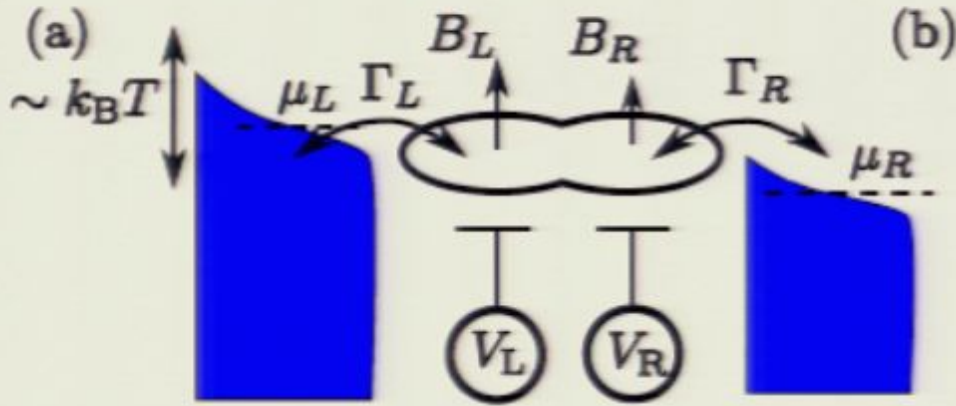
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- Can get stuck in triplet: Pauli spin blockade

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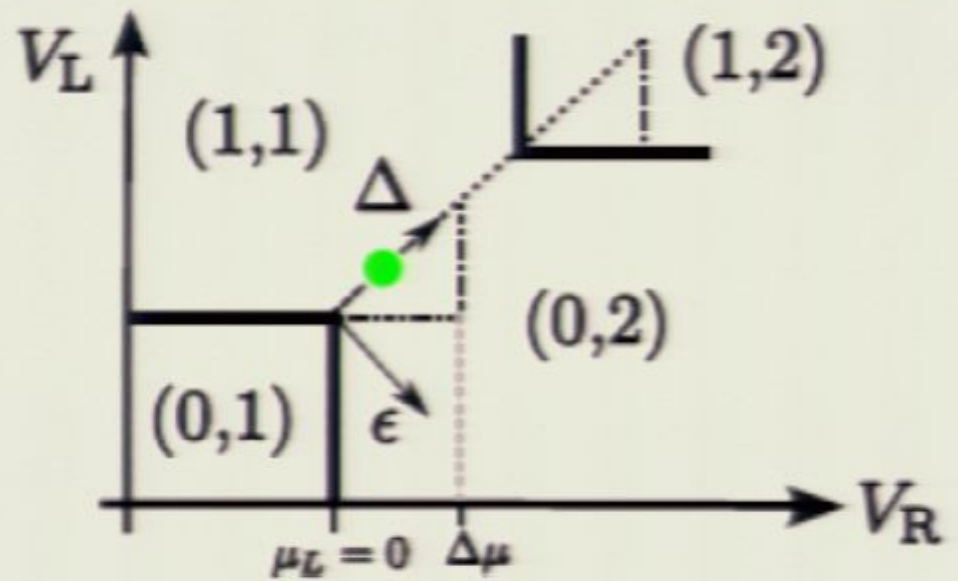


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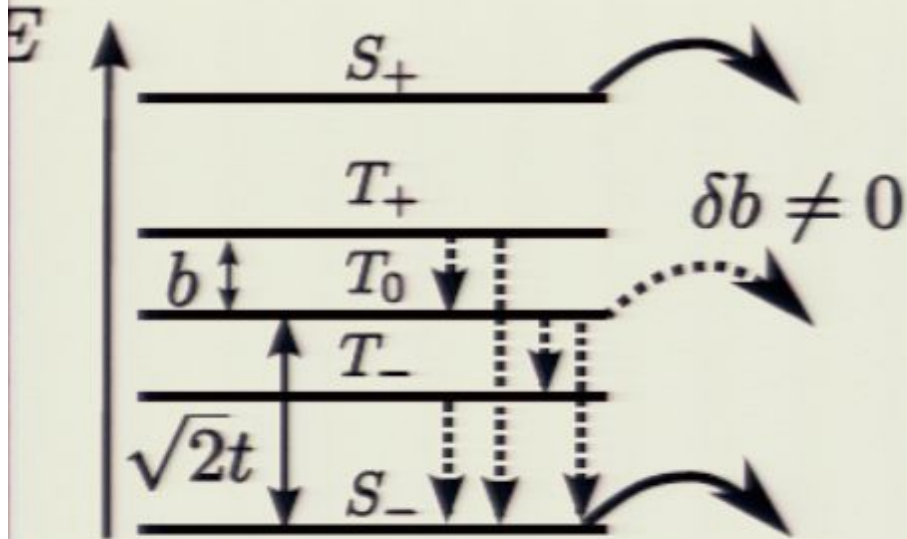
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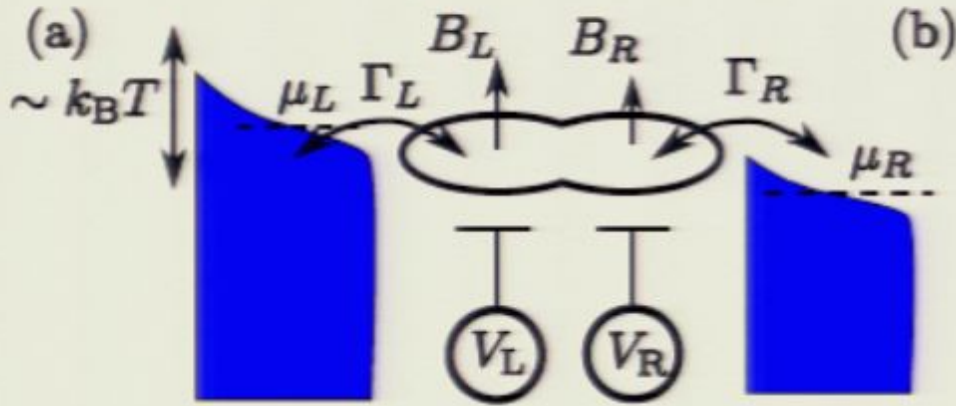
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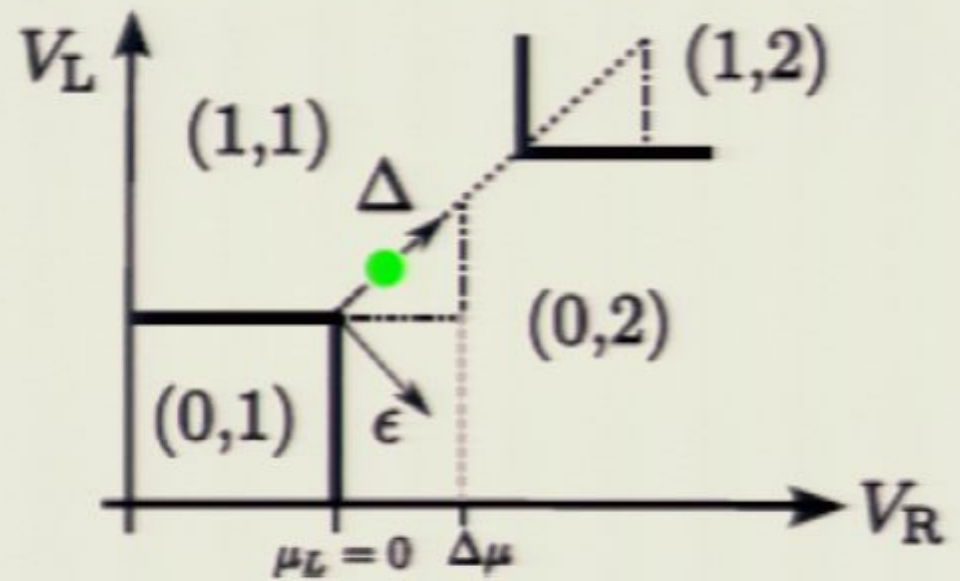


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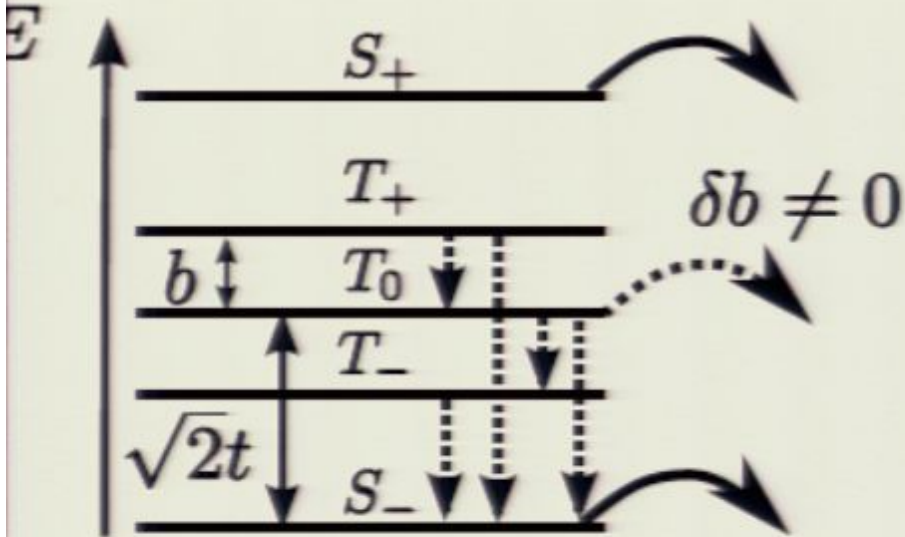
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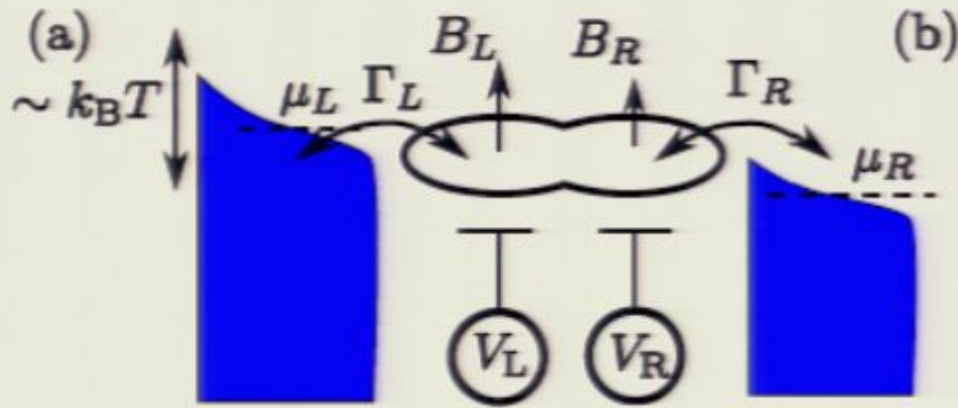
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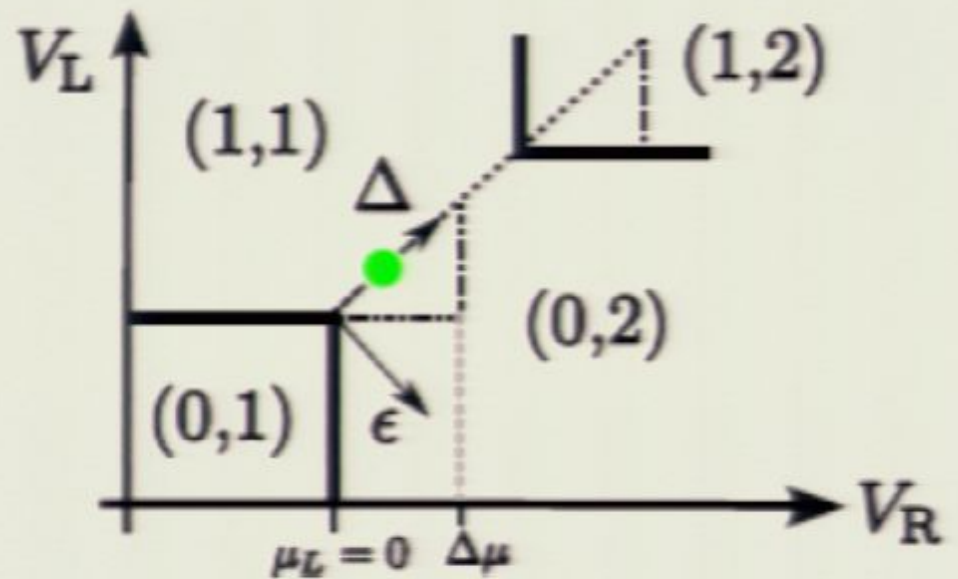


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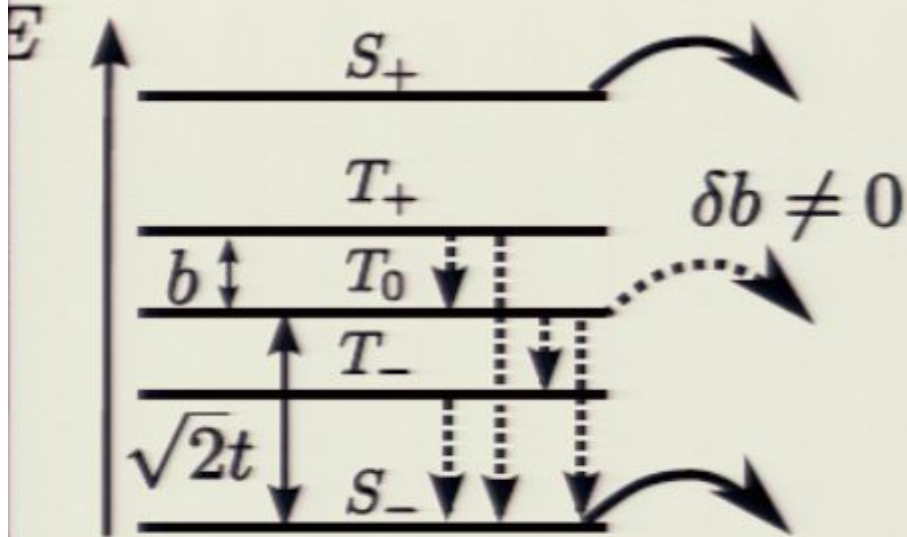
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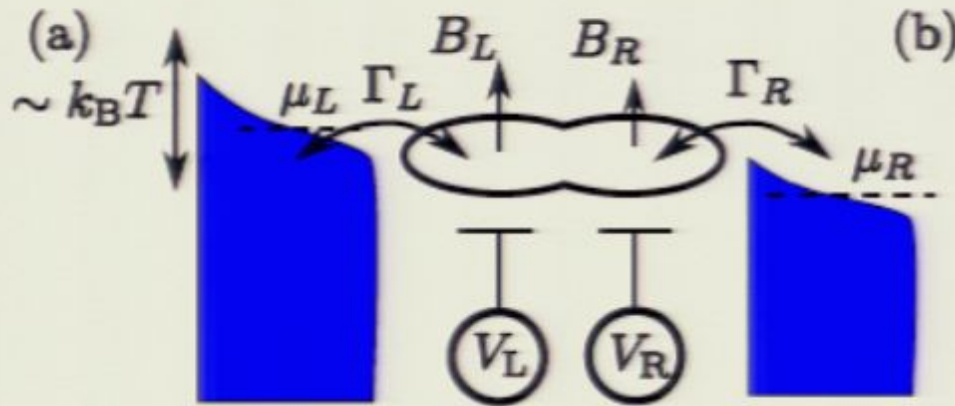
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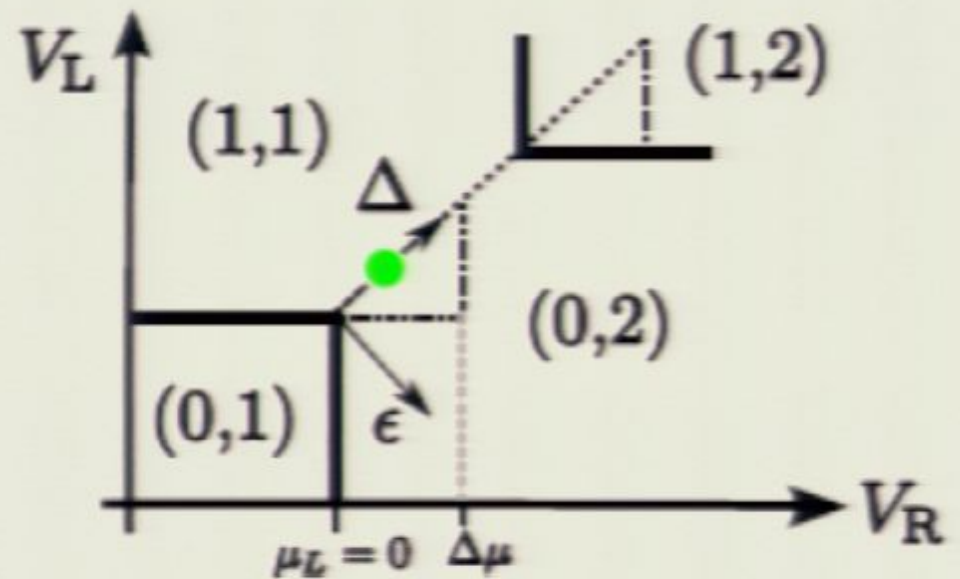


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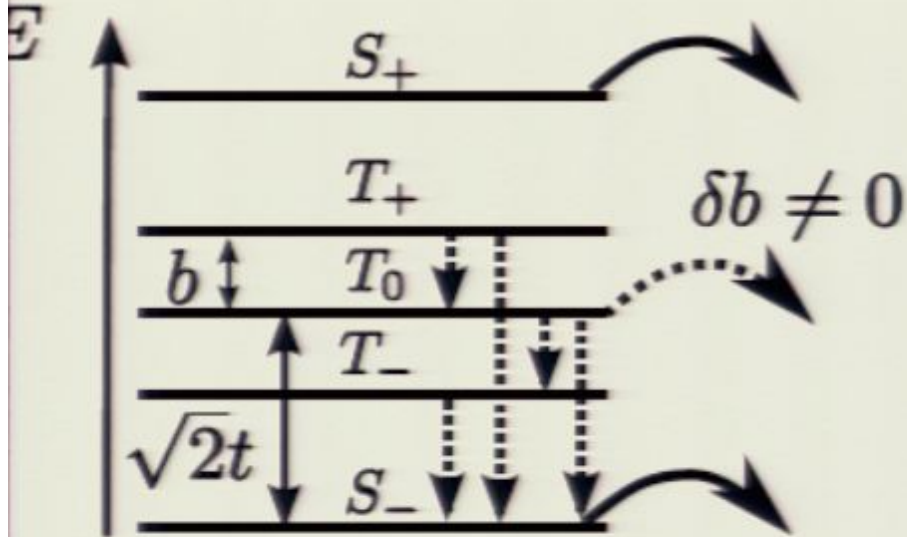
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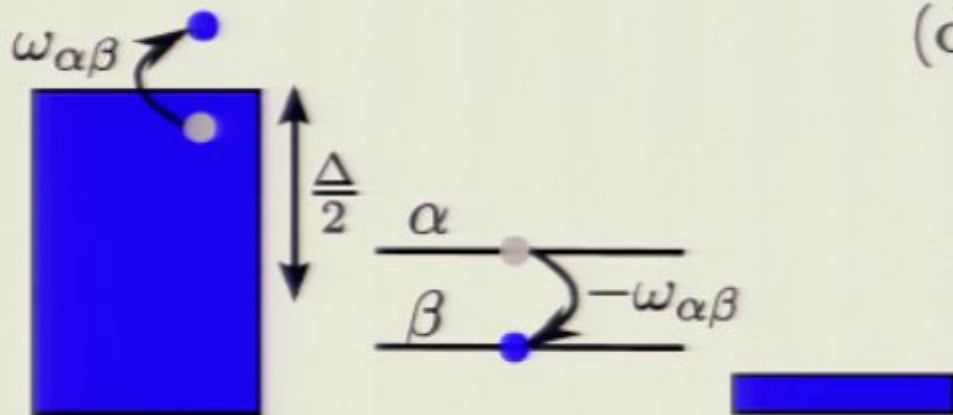
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Spin-flip cotunneling



(d)

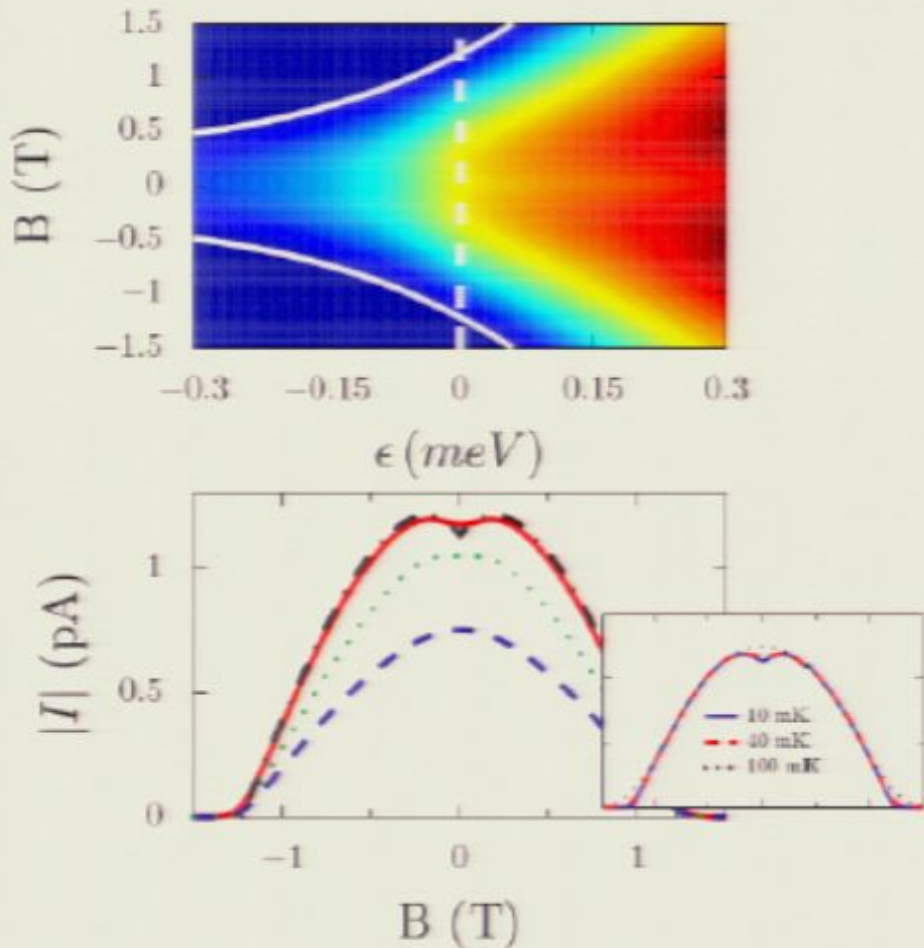
- second-order process
- second electron can have any spin
- High bias \rightarrow no Kondo effect
- sf-cotunneling breaks Pauli blockade at long times

Mathematical procedure

- Schrieffer-Wolf transformation for higher-order processes
- Fermi's golden rule rates for processes
- Pauli master equation for probabilities, currents

Result: Fast spin initialization

Leakage current:

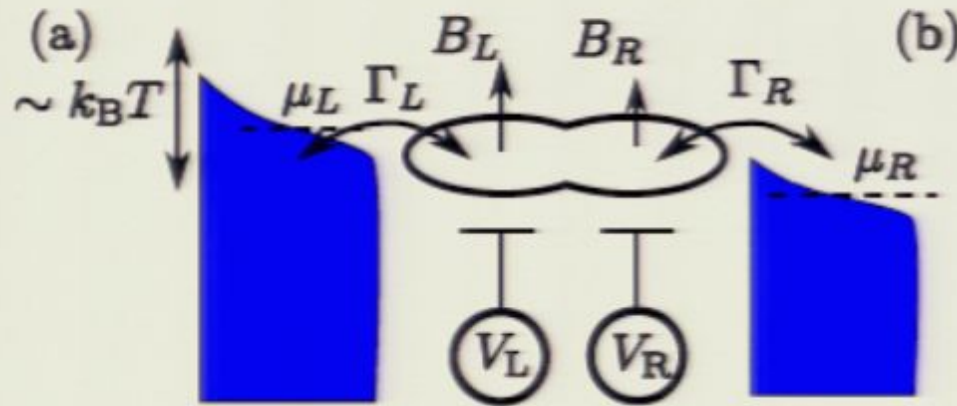


- cutoff at large B : $E_{T-} < E_{S-}$: spin blockade in presence at cotunneling at $T = 0$
- allows spin initialization — fast compared to inelastic spin relaxation
- Zero-field dip: At large δB , T_0 can directly escape via S_0

F. Qassemi, W.A. Coish, and FKW, PRL in press, arXiv:0812.2957

Stationary and transient leakage current in the Pauli spin blockade

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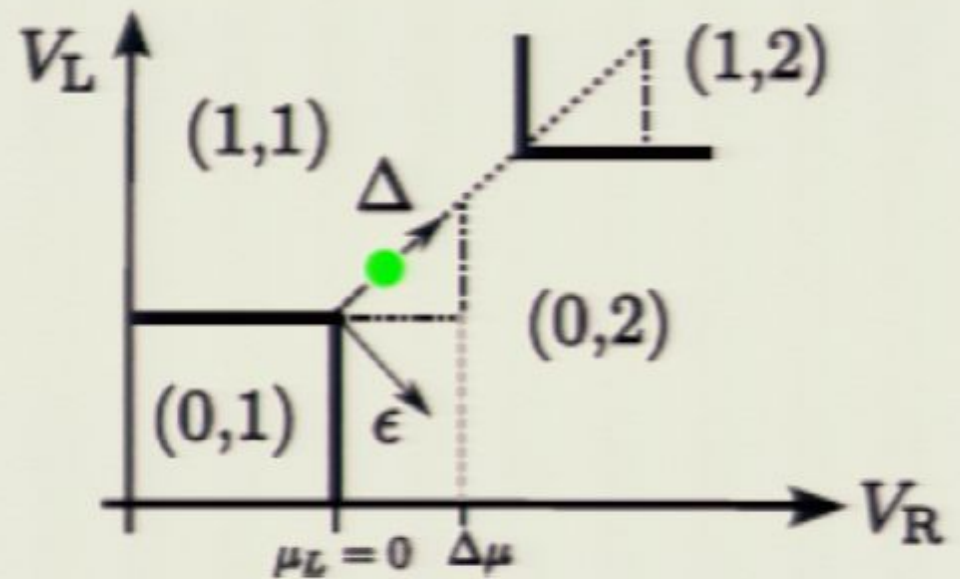


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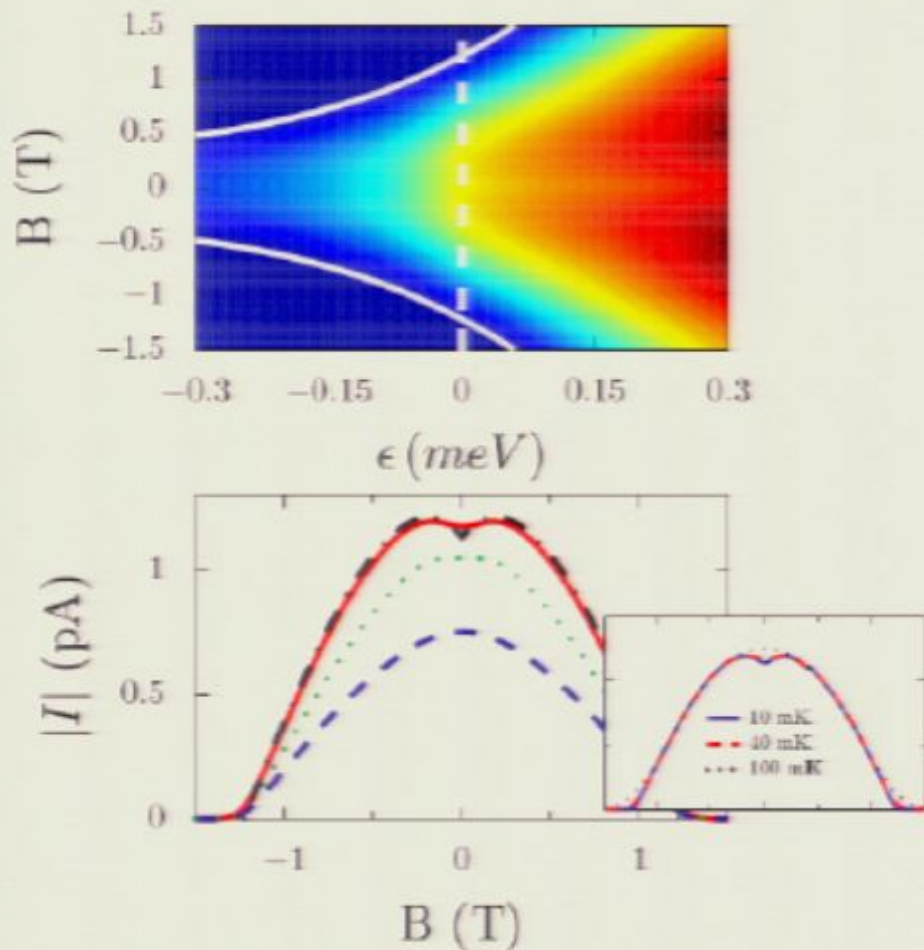
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Charge stability diagram



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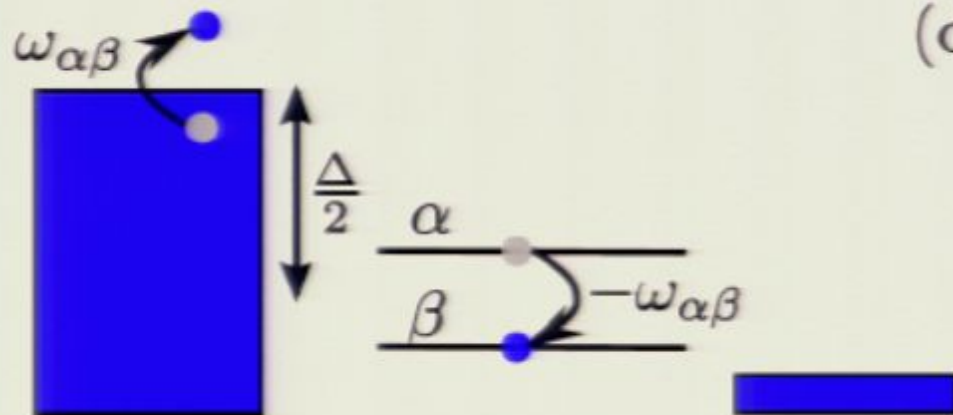
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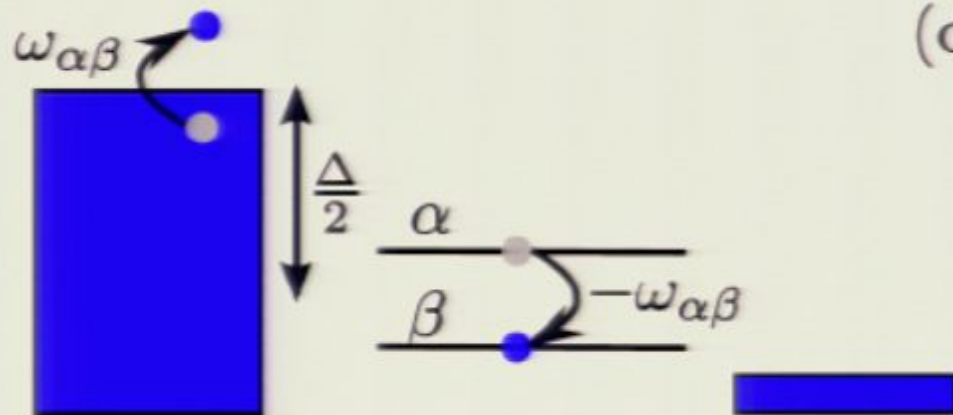
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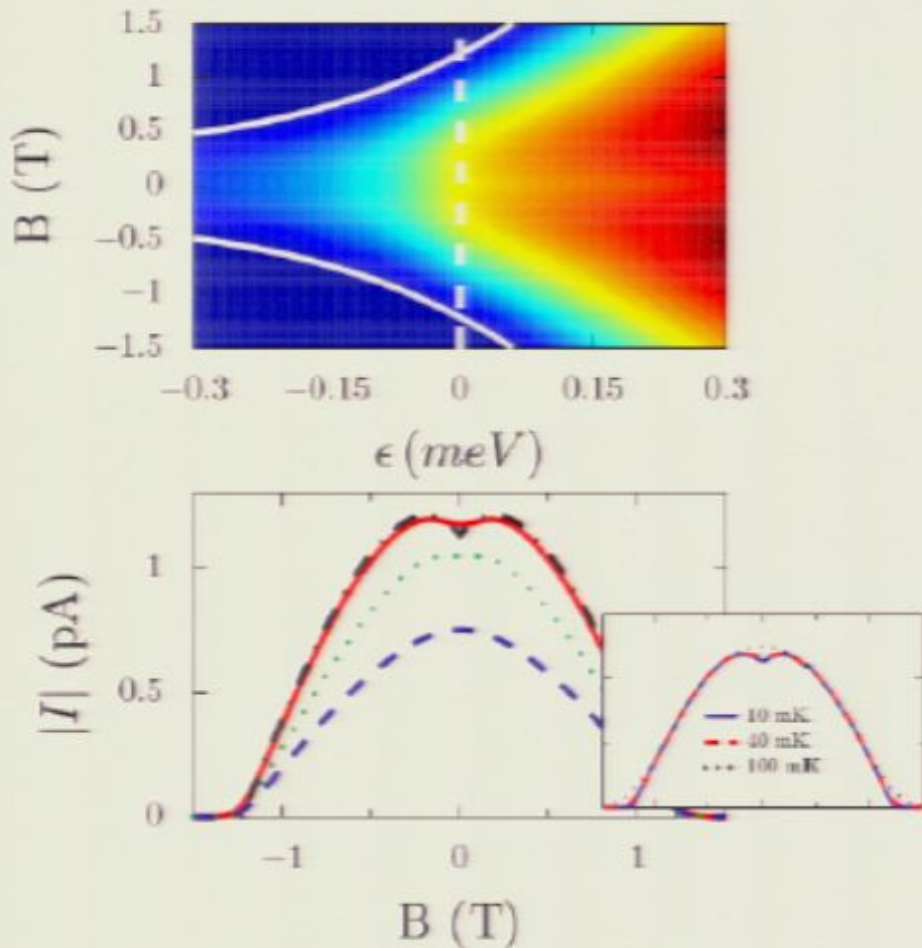
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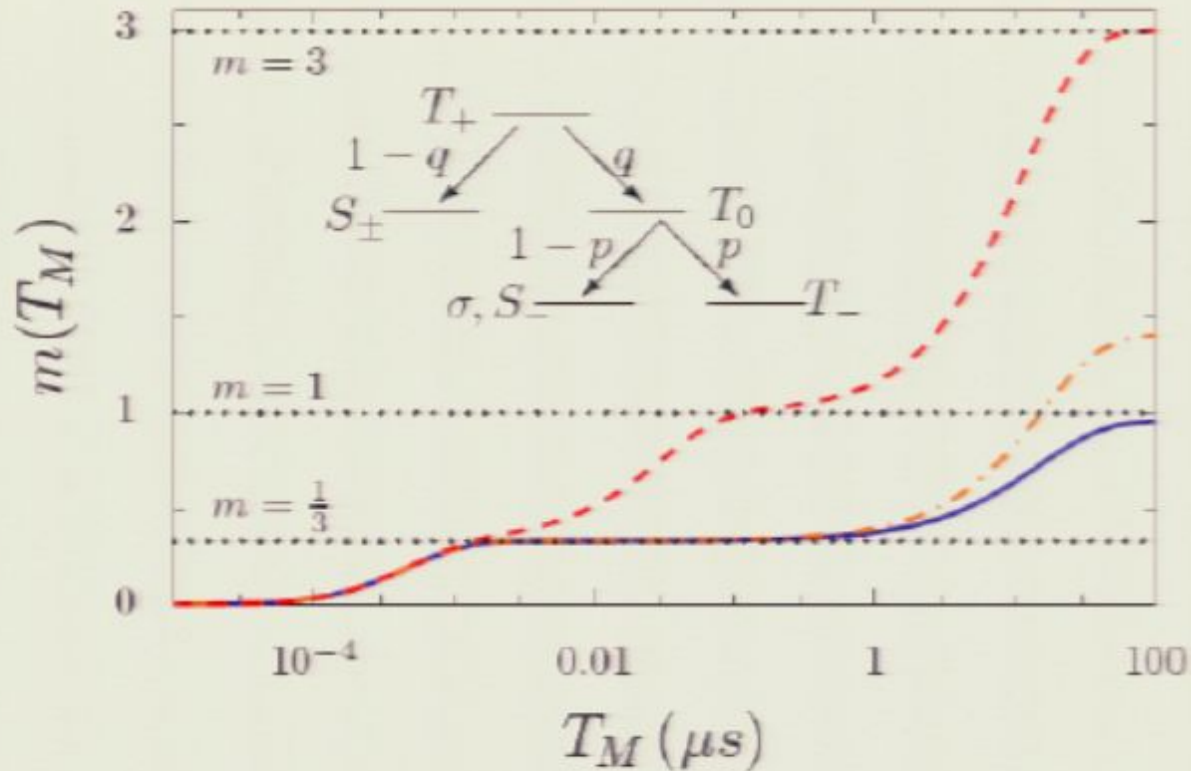
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Itinerant current statistics



- Transferred electron number (measurement time)
- Fractional values from rate branching ratio

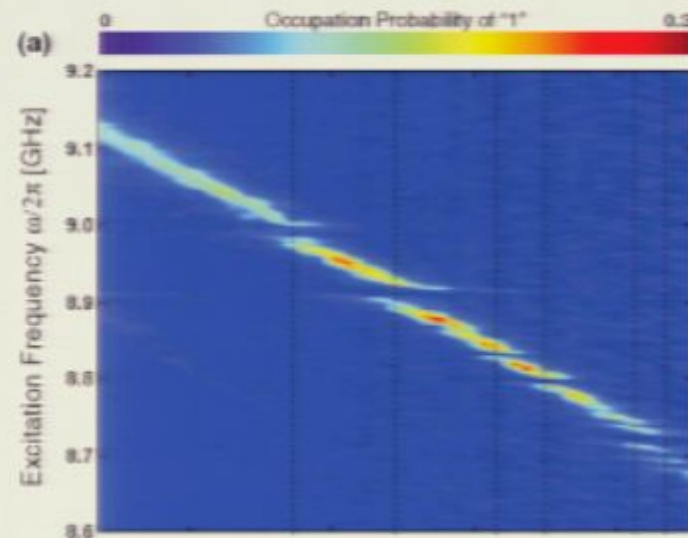
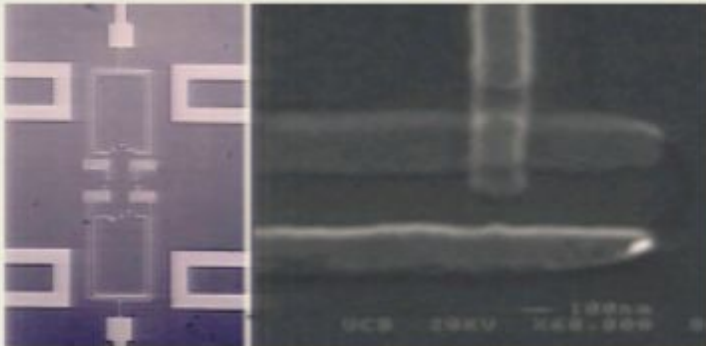
- $m = 1/3$: 3 of 4 levels blocking
- $m = 1$: 2 of 4 levels blocking
- $m = 3$: 1 of 4 levels blocking

Outline

- 1 Spin in quantum dots
- 2 Noise in tunnel junctions

Superconducting tunnel junctions

Tool for applications: Bolometers, SQUIDs, qubits



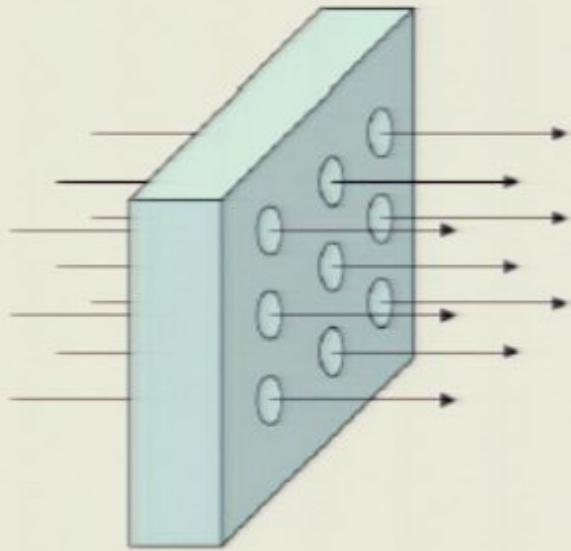
Excess noise in junctions limits sensitive experiments:

Qubit: R.W. Simmonds *et al.*, PRL 2004.

Bolometer / Mixer: P. Dieleman *et al.*, PRL 1997.

SQUIDs: F. Wellstood *et al.*, APL 2004.

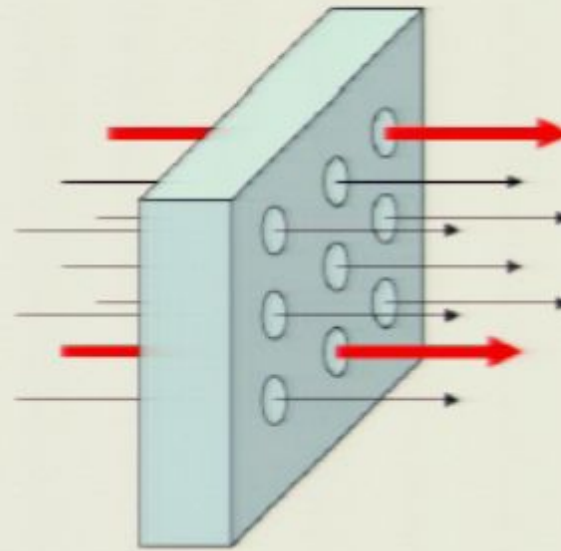
Rough junctions



Homogenous tunnel junction:

Transmission $T \ll 1$, M modes

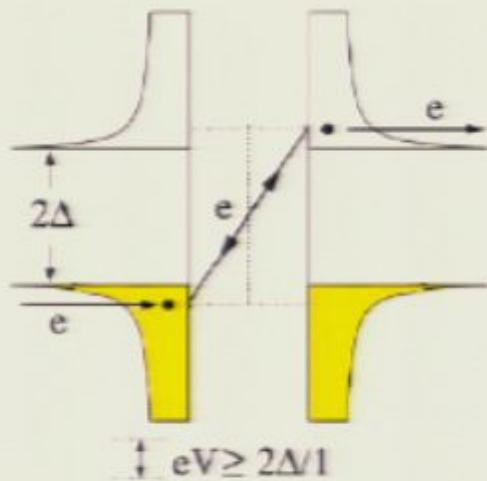
Phase qubit: $M \simeq 10^6$, $T \simeq 0.003$



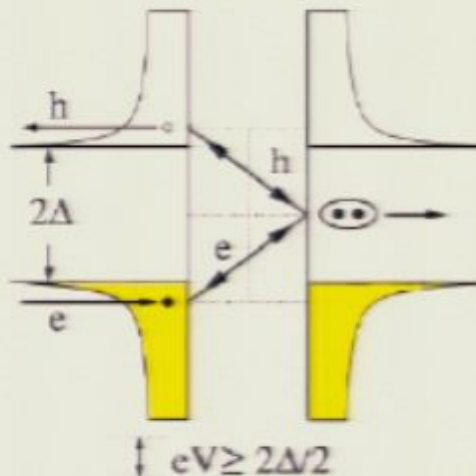
Rough junction: $T \propto e^{-d\sqrt{2m(U-E)}}$

Pinhole(s), $T \simeq 1$.

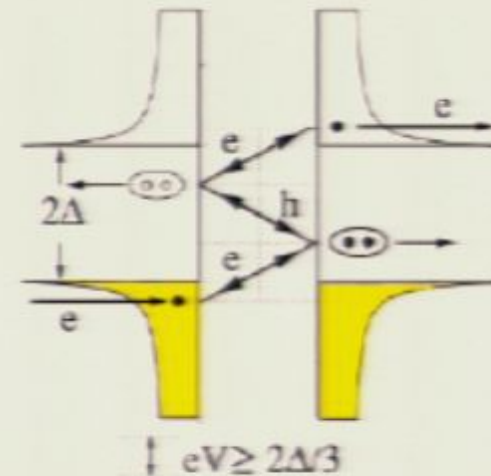
Multiple Andreev reflection



transferred charge $1e$



transferred charge $2e$



transferred charge $3e$

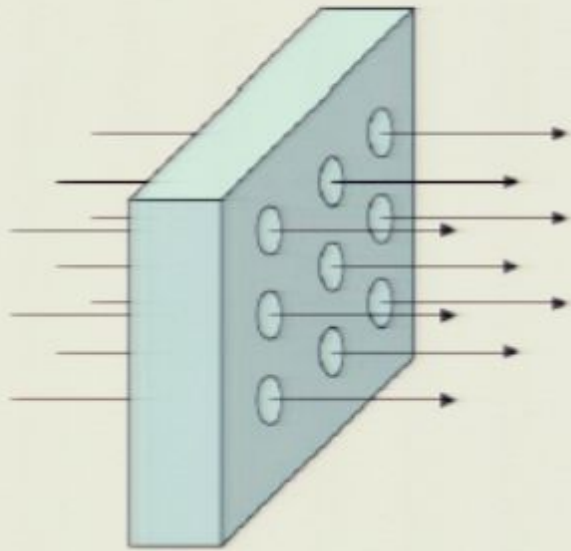
Transferred charge $n \geq e \frac{2\Delta}{eV}$

Reflection at off-diagonal (electron-hole potential)

\Rightarrow can take place above gap.

$V \rightarrow 0$: Supercurrent-carrying bound state

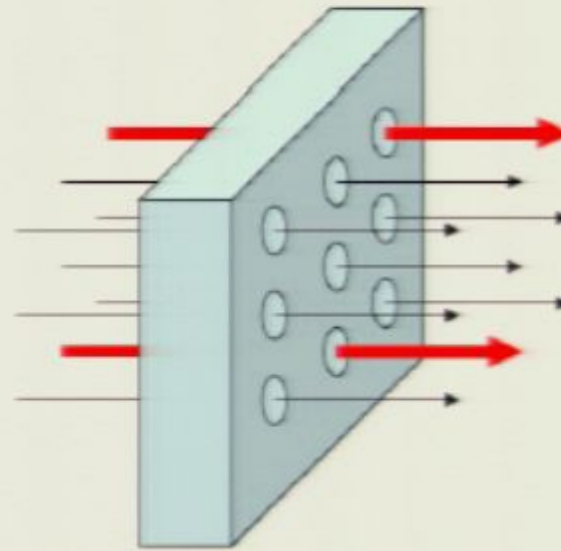
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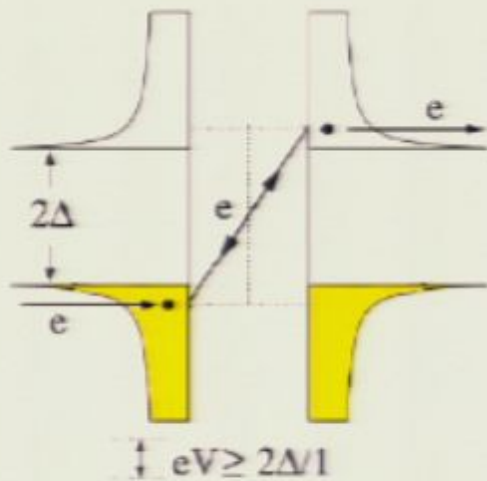
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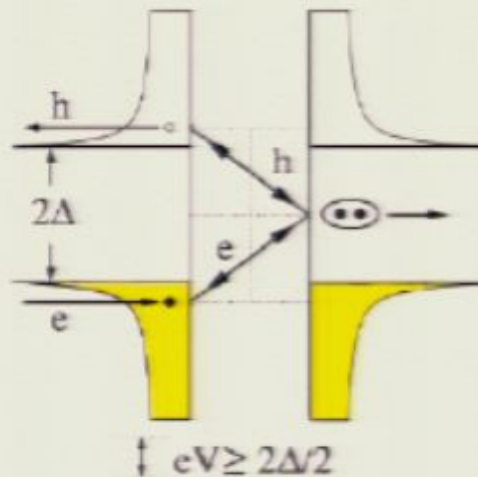
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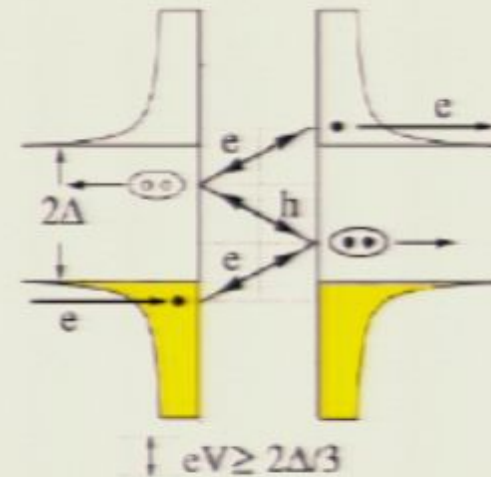
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Transport and shot noise

$T = 0$ quantum transport, no thermal noise.

Conductance:

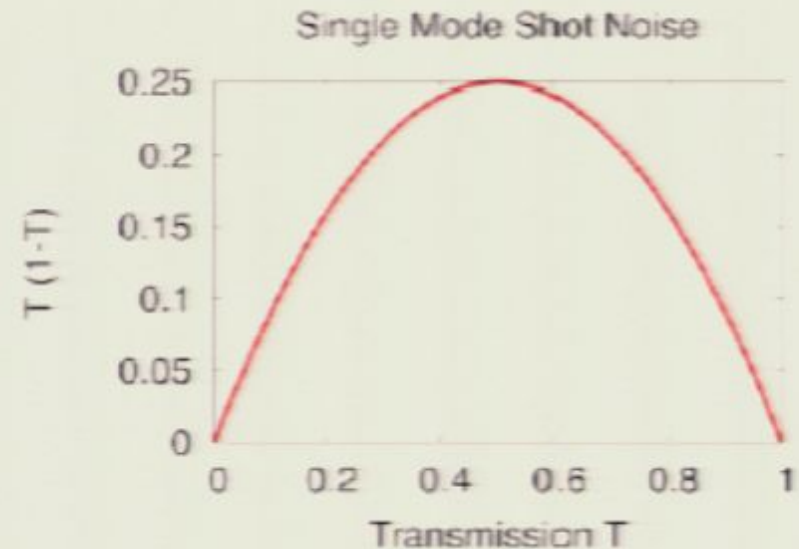
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Shot noise:

$$S_I = 2eVG_Q \sum_n T_n(1 - T_n)$$

Fano factor: $F = S_I/2eI$,

Tunnel junction: $F = 1$



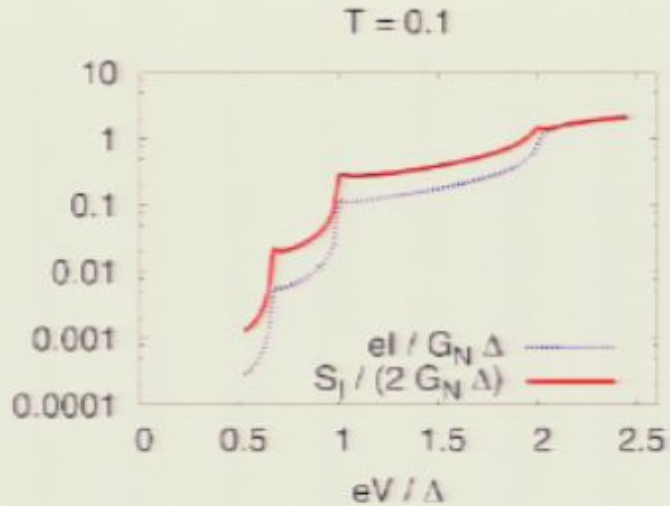
Technique: Full counting statistics of charge transfer

Cumulant generating function $S(\chi)$:

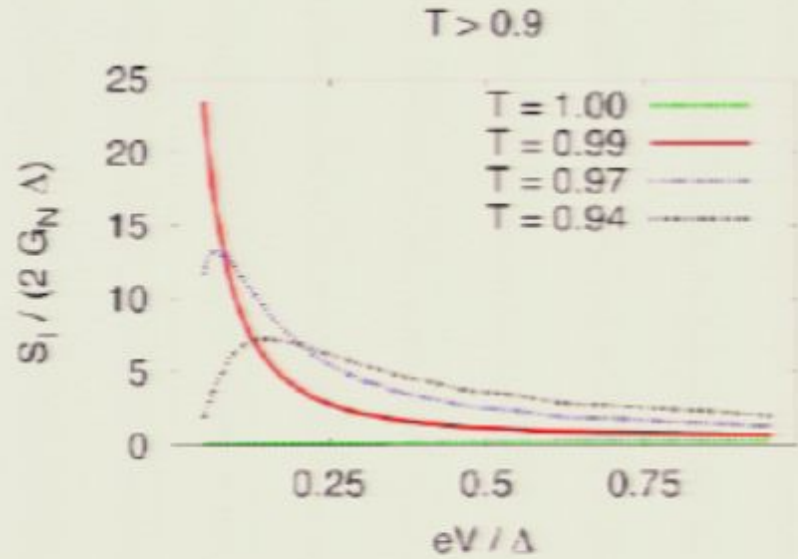
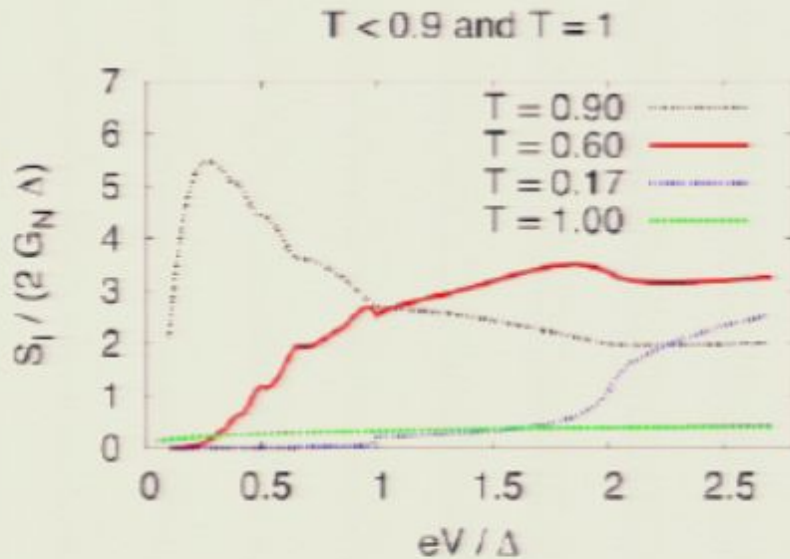
$$e^{-S(\chi)} = \sum_N P(N) e^{iN\chi}$$

Calculated from Keldysh Green's function including counting field (Levitov, Lesovik, Nazarov, Belzig, 1996 —)

Single mode contact



Subharmonic gap structure
 Crossover to $S \propto 2eN \propto 1/V$
 Residual noise at $T = 1$: Certain about transport happening, uncertain about number of cycles (supergap AR!)



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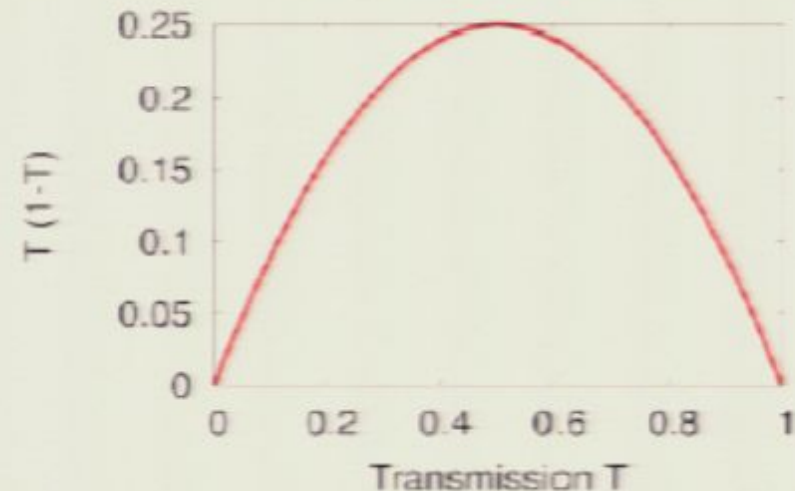
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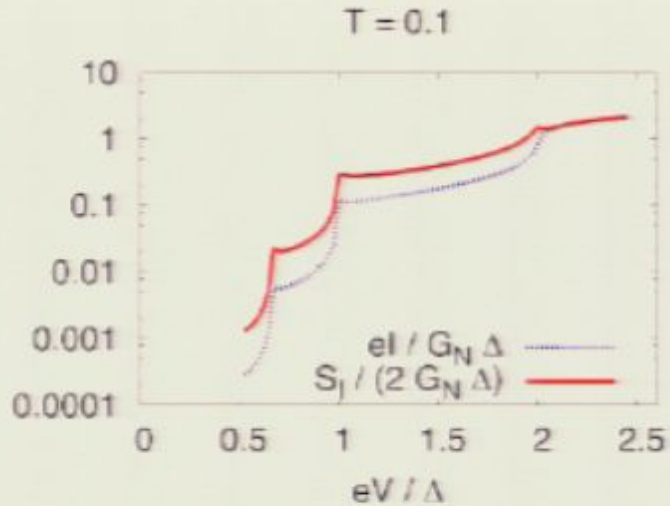
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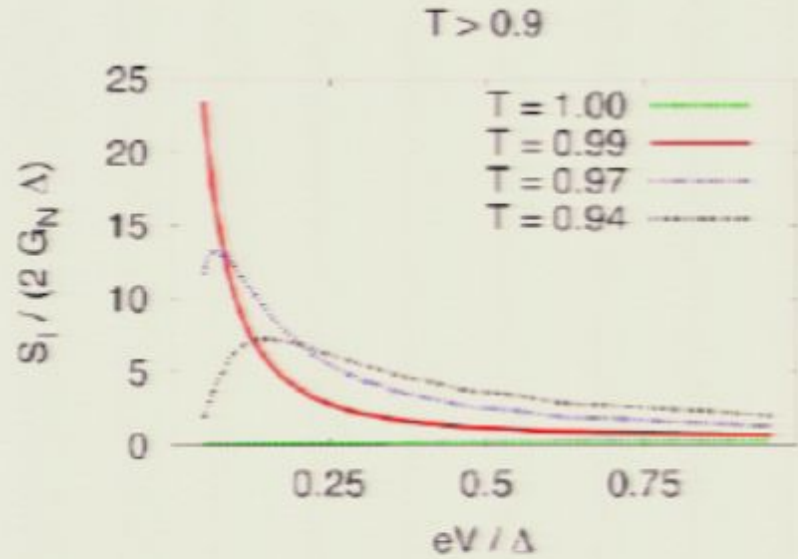
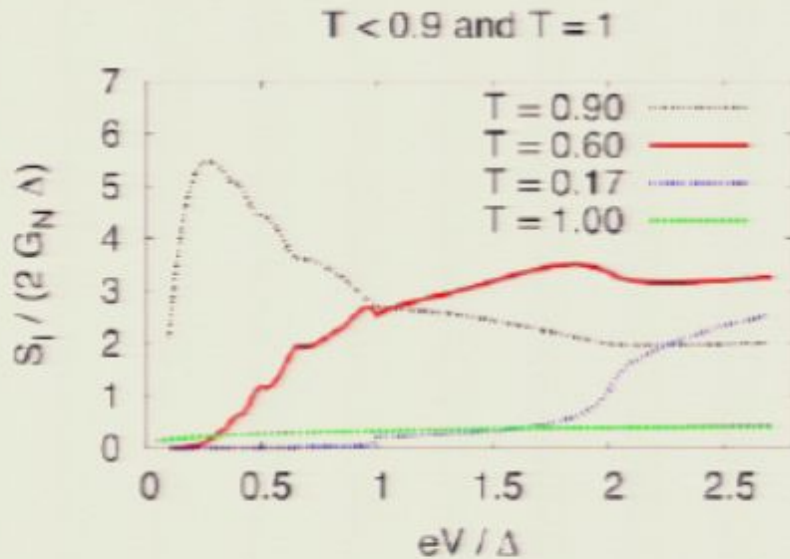
Single Mode Shot Noise



Single mode contact



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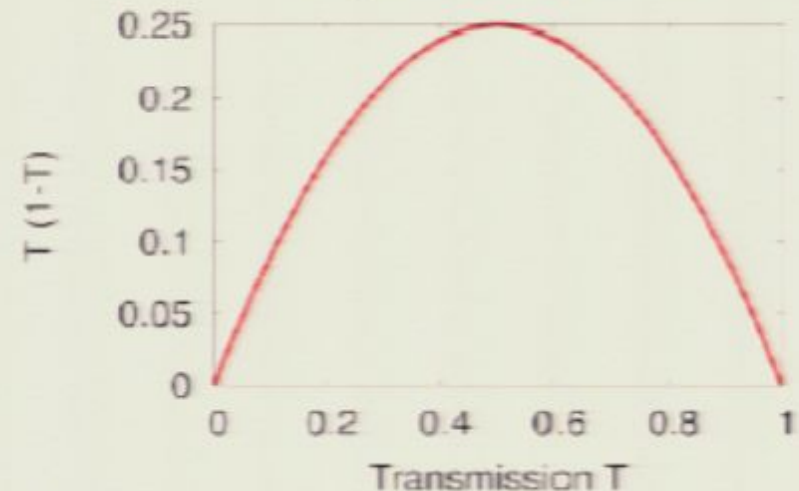
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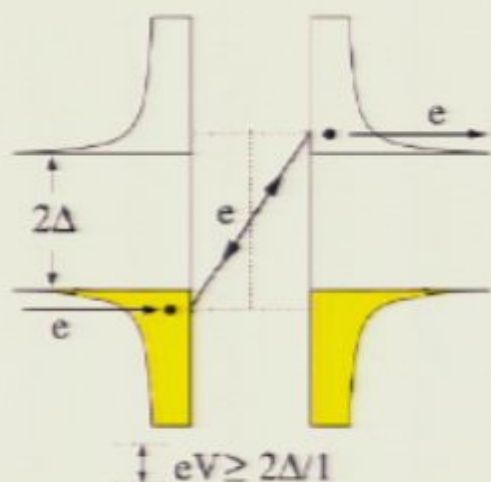
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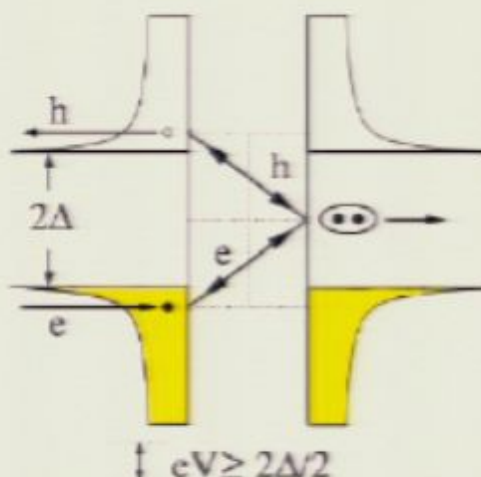
Single Mode Shot Noise



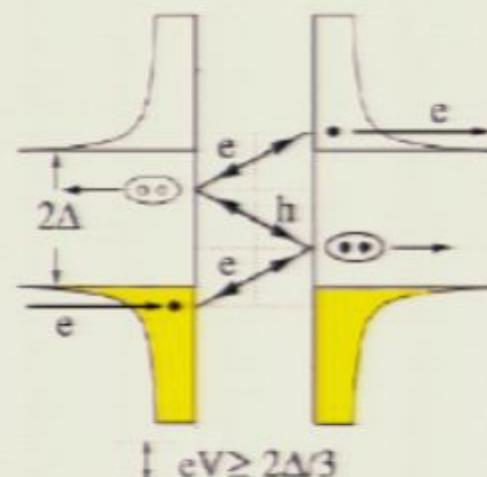
Multiple Andreev reflection



transferred charge $1e$



transferred charge $2e$



transferred charge $3e$

Transferred charge $n \geq e \frac{2\Delta}{eV}$

Reflection at off-diagonal (electron-hole potential)

⇒ can take place above gap.

$V \rightarrow 0$: Supercurrent-carrying bound state

Transport and shot noise

$T = 0$ quantum transport, no thermal noise.

Conductance:

$$G_N = G_Q \sum_n T_n$$

Shot noise:

$$S_I = 2eVG_Q \sum_n T_n(1 - T_n)$$

Fano factor: $F = S_I/2el$,

Tunnel junction: $F = 1$

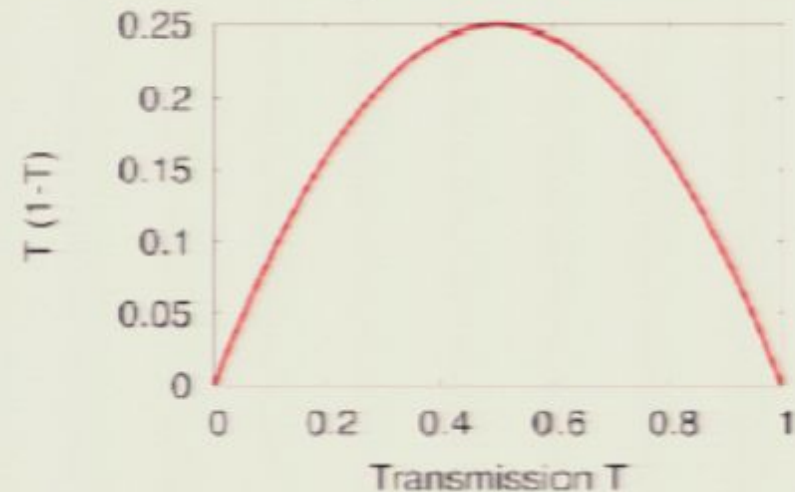
Technique: Full counting statistics of charge transfer

Cumulant generating function $S(\chi)$:

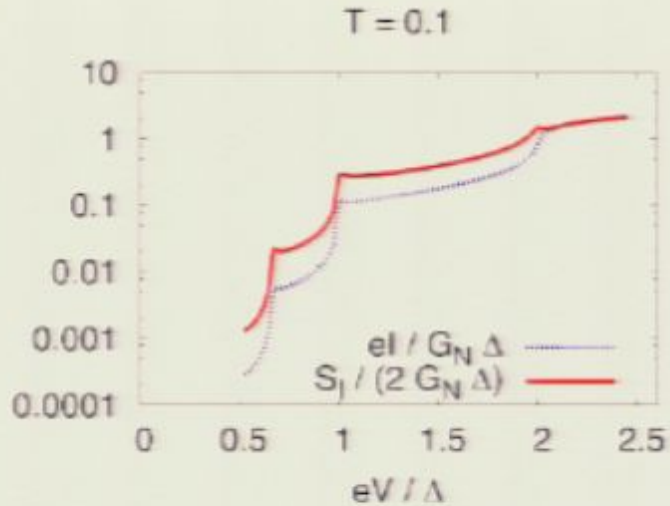
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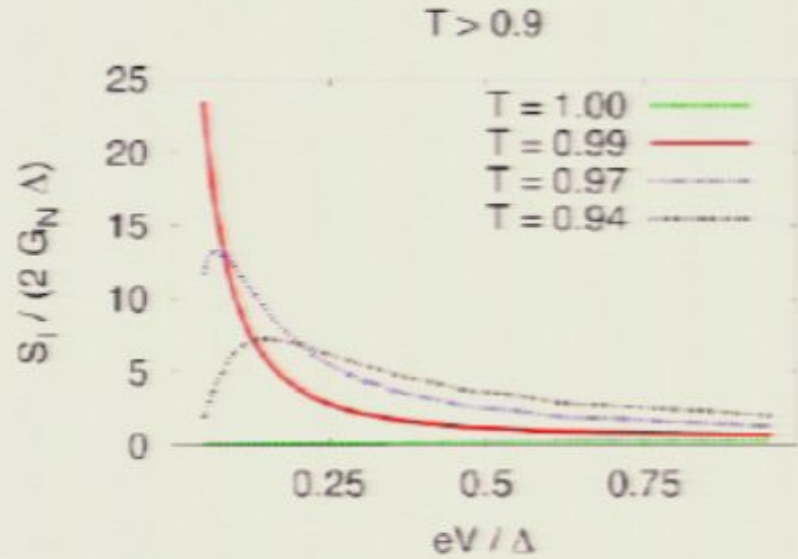
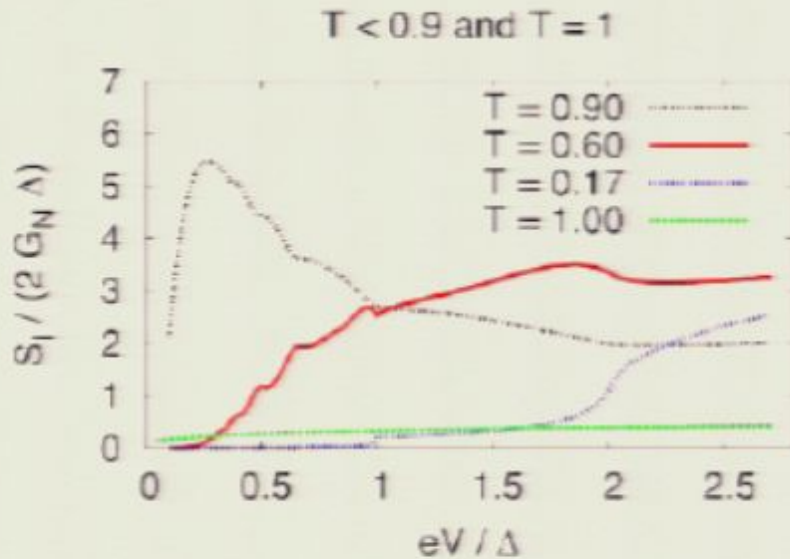
Single Mode Shot Noise



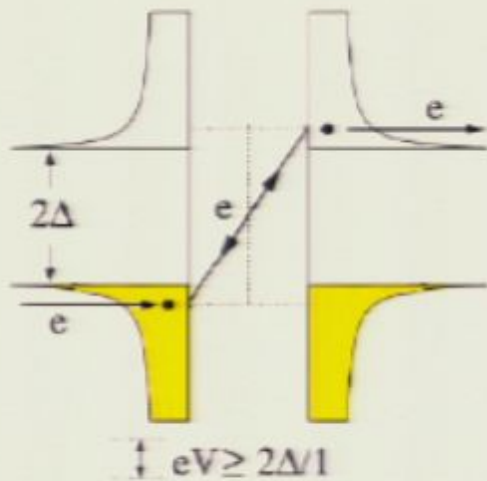
Single mode contact



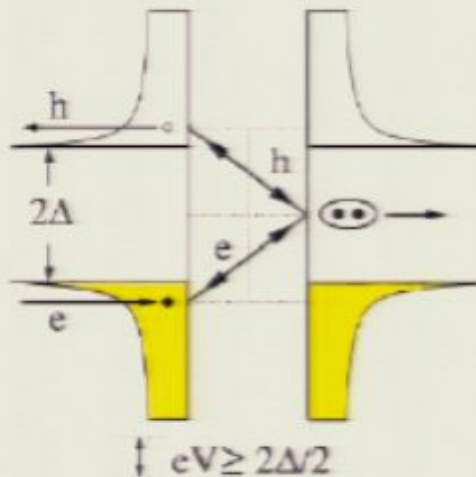
Subharmonic gap structure
 Crossover to $S \propto 2eN \propto 1/V$
 Residual noise at $T = 1$: Certain about transport happening, uncertain about number of cycles (supergap AR!)



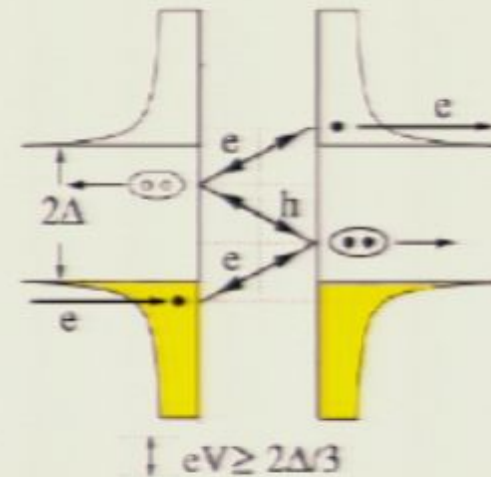
Multiple Andreev reflection



transferred charge $1e$



transferred charge $2e$



transferred charge $3e$

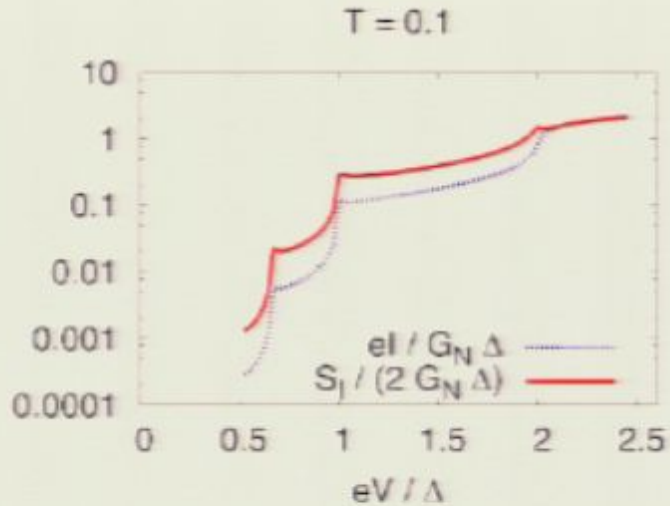
Transferred charge $n \geq e \frac{2\Delta}{eV}$

Reflection at off-diagonal (electron-hole potential)

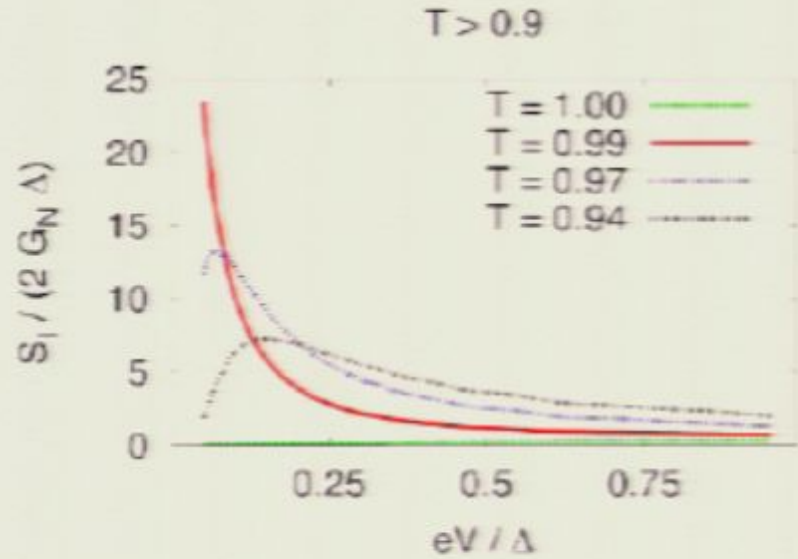
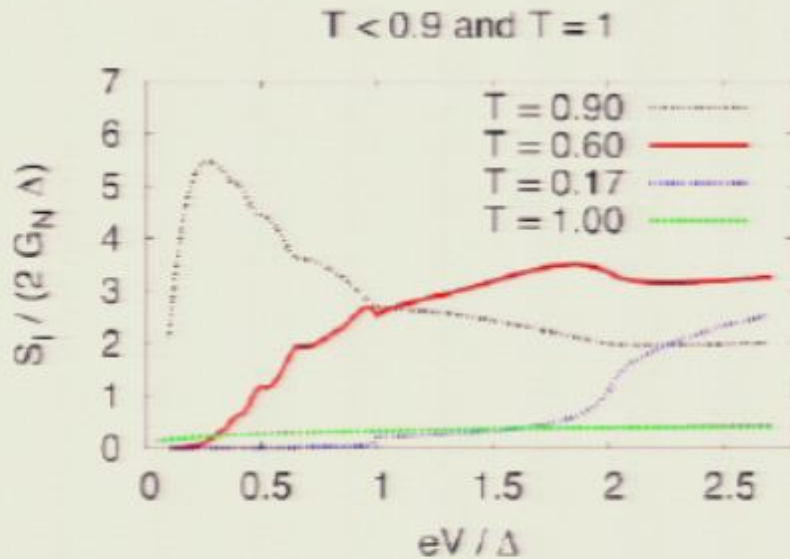
⇒ can take place above gap.

$V \rightarrow 0$: Supercurrent-carrying bound state

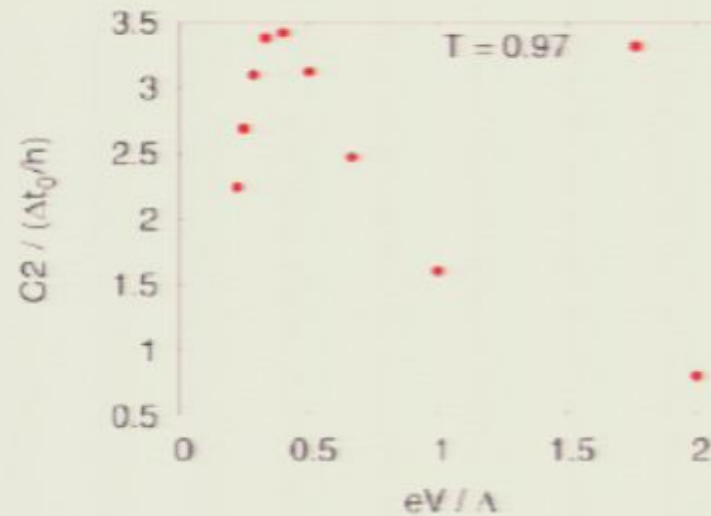
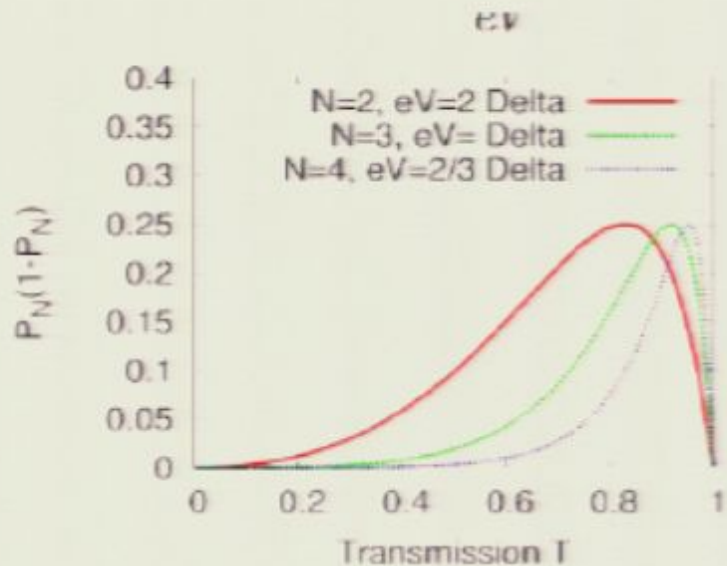
Single mode contact



Subharmonic gap structure
Crossover to $S \propto 2eN \propto 1/V$
Residual noise at $T = 1$: Certain about transport happening, uncertain about number of cycles (supergap AR!)



Counting Andreev cycles



Qualitative prediction $S \propto 2enT^n(1 - T^n)$ for n-th order Andreev Divergence as $1/V^\alpha$.

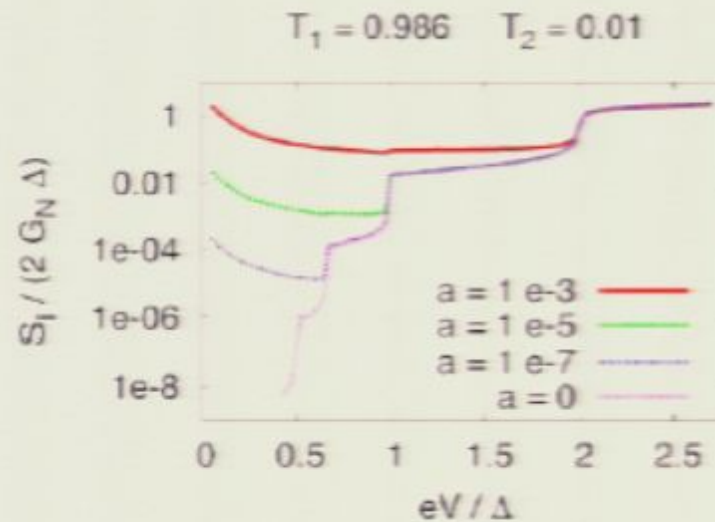
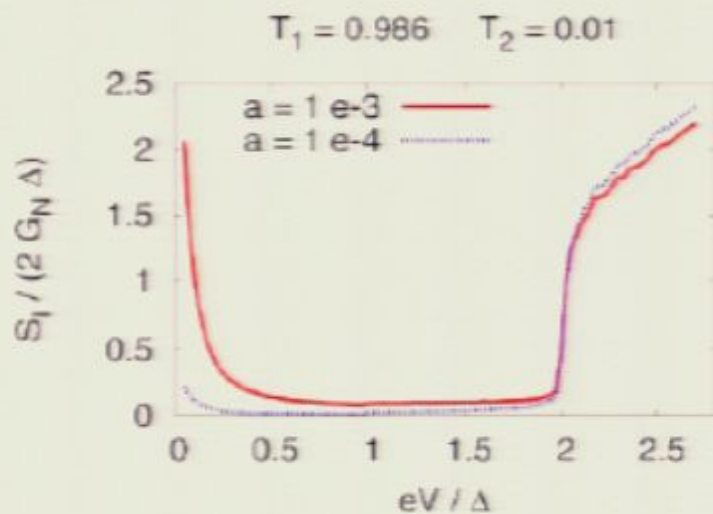
Qualitative: $\alpha = 1$, full model $\alpha \simeq 0.8$.

G. Heinrich and FKW, arXiv:0808.3705

Full tunnel junctions

Tunnel contact with fraction a pinholes

$$\rho(T) = a\delta(T - T_1) + (1 - a)\delta(T - T_2)$$



Pinhole can be hidden in current but dominate low- V -noise

Summary

- Spin-flip cotunneling
- Noise in superconducting tunnel junctions

