

Title: Qu-transitions. Phase transitions in the quantum era.

Date: Apr 22, 2009 02:00 PM

URL: <http://pirsa.org/09040013>

Abstract: Physicists are often so awestruck by the lofty achievements of the past, we end up thinking all the big stuff is done, which blinds us to the revolutions ahead. We are still firmly in the throes of the quantum revolution that began a hundred years ago. Quantum gravity, quantum computers, qu-bits and quantum phase transitions, are manifestations of this ongoing revolution. Nowhere is this more so, than in the evolution of our understanding of the collective properties of quantum matter. Fifty years ago, physicists were profoundly shaken by the discovery of universal power-law correlations at classical second-order phase transitions. Today, interest has shifted to Quantum Phase Transitions: phase transitions at absolute zero driven by the violent jiggings of quantum zero-point motion. Quantum, or Qu-transitions have been observed in ferromagnets, helium-3, ferro-electrics, heavy electron and high temperature superconductors. Unlike its classical counterpart, a quantum critical point is a kind of 'black hole' in the materials phase diagram: a singularity at absolute zero that profoundly influences wide swaths of the material phase diagram at finite temperature. I'll talk about some of the novel ideas in this field including 'avoided criticality' - the idea that high temperature superconductivity nucleates about quantum critical points - and the growing indications that electron quasiparticles break up at a quantum critical point.

Qu-Transitions



P. Coleman
(CMT, Rutgers)

Perimeter Inst.
Waterloo, Apr 21
2009.

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Qu-Transitions

Phase transitions in the quantum era"



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- The Qu-era: from classical to quantum.
- Classical criticality
- Heavy Fermion Quantum Criticality.
- New Approaches and Ideas
- Avoided Criticality

1758 in Paris: 72 years after “Principia”

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Classical revolution is still in full sway.

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Return of Halley’s comet.

Publishers in Paris decide to bring out the 1st French Translation of Newton’s Principia, which they have held in proof form for ten years.

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PRINCIPES
MATHÉMATIQUES

DE LA

PHILOSOPHIE NATURELLE,

Par feu M^{onsieur} le Marquis DU CHASTELLET.

TOME PREMIER.



A PARIS,

DESSANT & SAILLANT, rue S. Jehu de Beuvain.
Chez LABRÉY, rue de la Comédie Française,
au Paroisse.

M. D. C. C. L. V. I.

AVEC APPROBATION, ET PRIVILEGE DU ROI.

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Par *Madame la Marquise du CHÂTELET.*

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Marquise Emilie du Châtelet
(1707-1749)

Mathematical Physicist:
Translator and interpreter of Principia.

"ce beau probleme astronomico-geometrique"

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Newtons Momentum

$$\sum_i m_i \vec{v}_i$$

Leibniz' "vis vivre"

$$\sum_i m_i (v_i)^2$$

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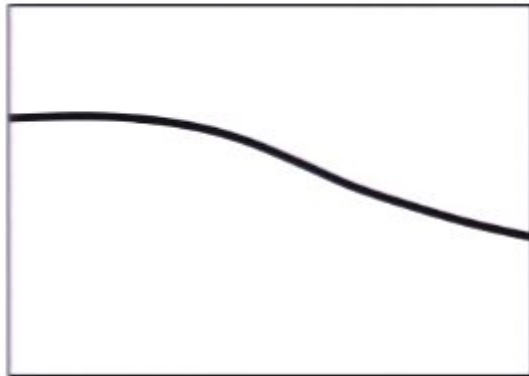
Resolution of the controversy (and the missing
factor of a half) required a further 60-80 years.

108 years after Planck, many
surprises later, the quantum era is in full sway.

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“With a heavy heart, I have been converted to the idea that Fermi -Dirac, not Einstein-Bose is the correct statistics. I wish to write a short note on its application to paramagnetism.”

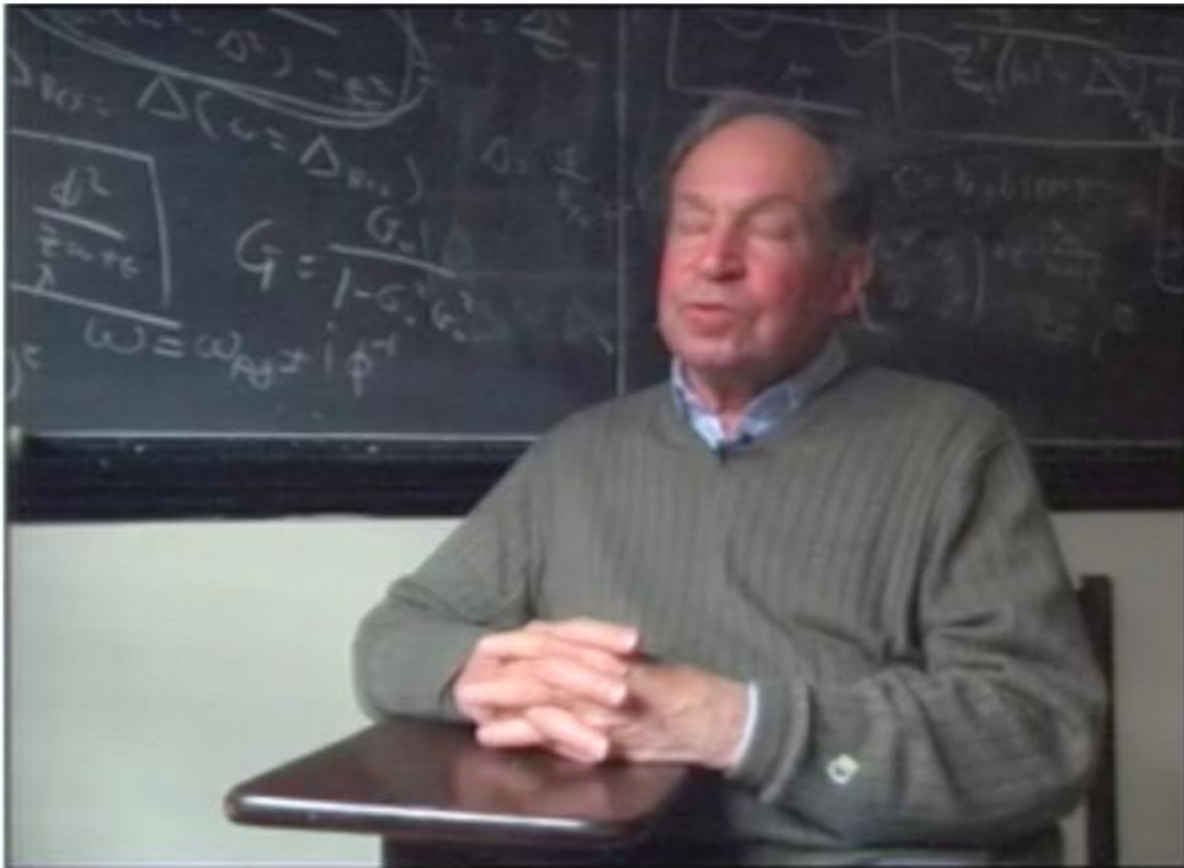
W. Pauli, in letter to Schrödinger, Dec 1926.

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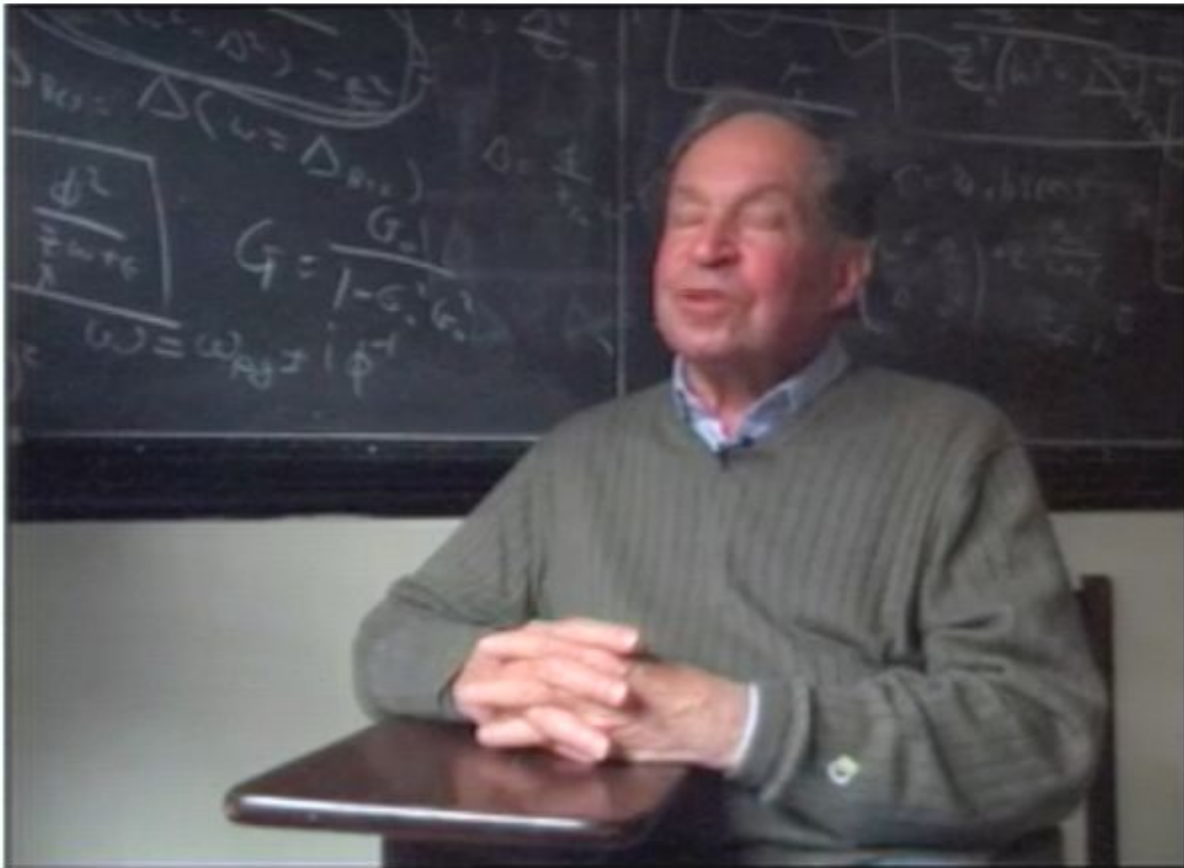
David Pines in
musicofthequantum.rutgers.edu

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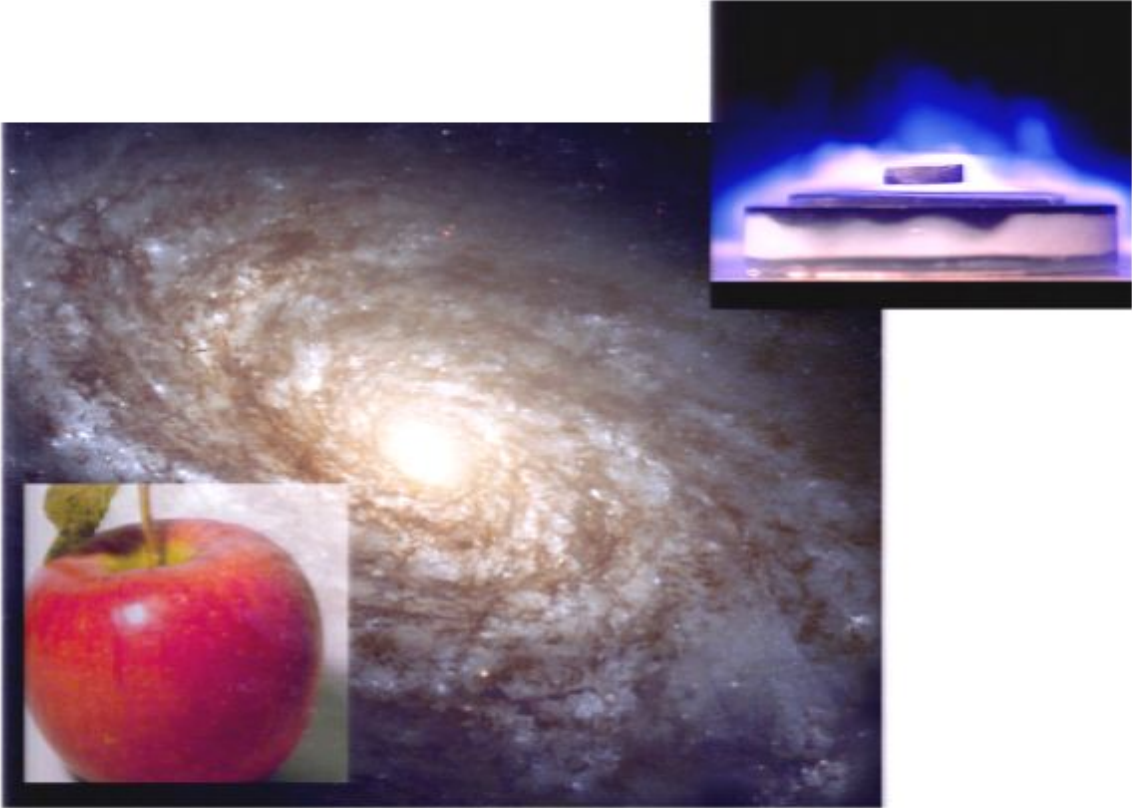
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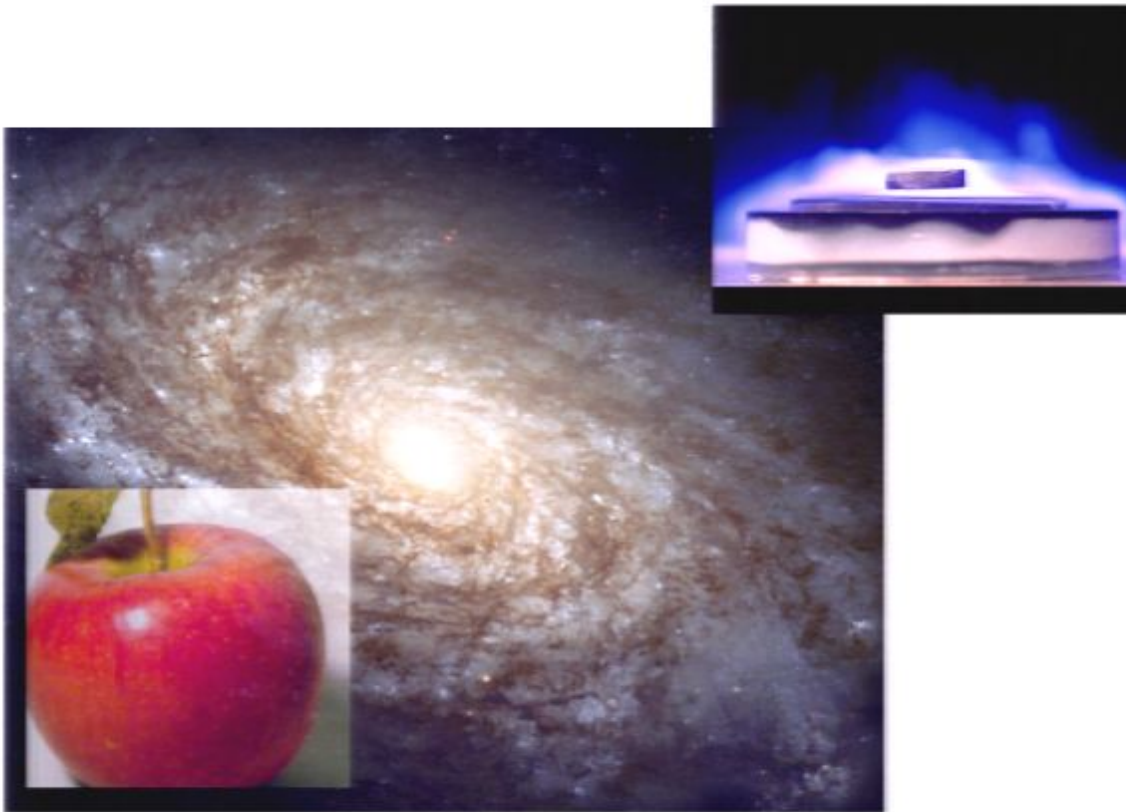


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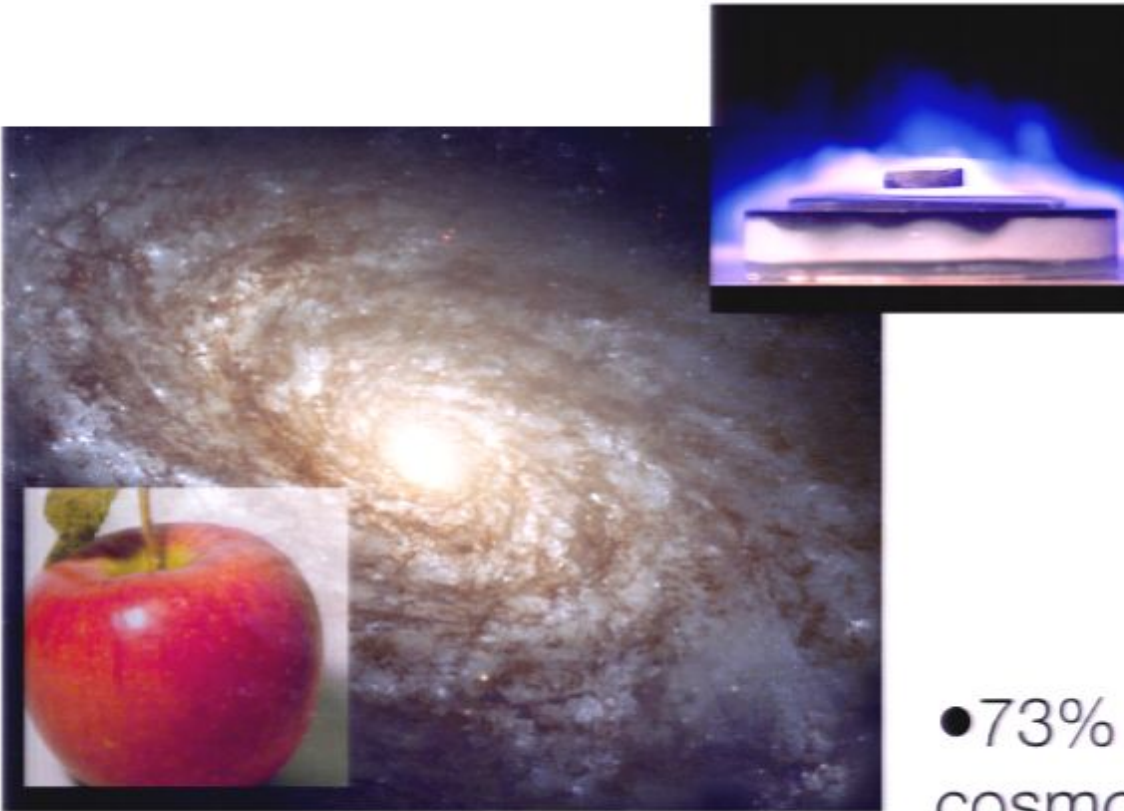
Quantum zero point fluctuations:



Quantum zero point fluctuations: major unsolved problem of the quantum era.

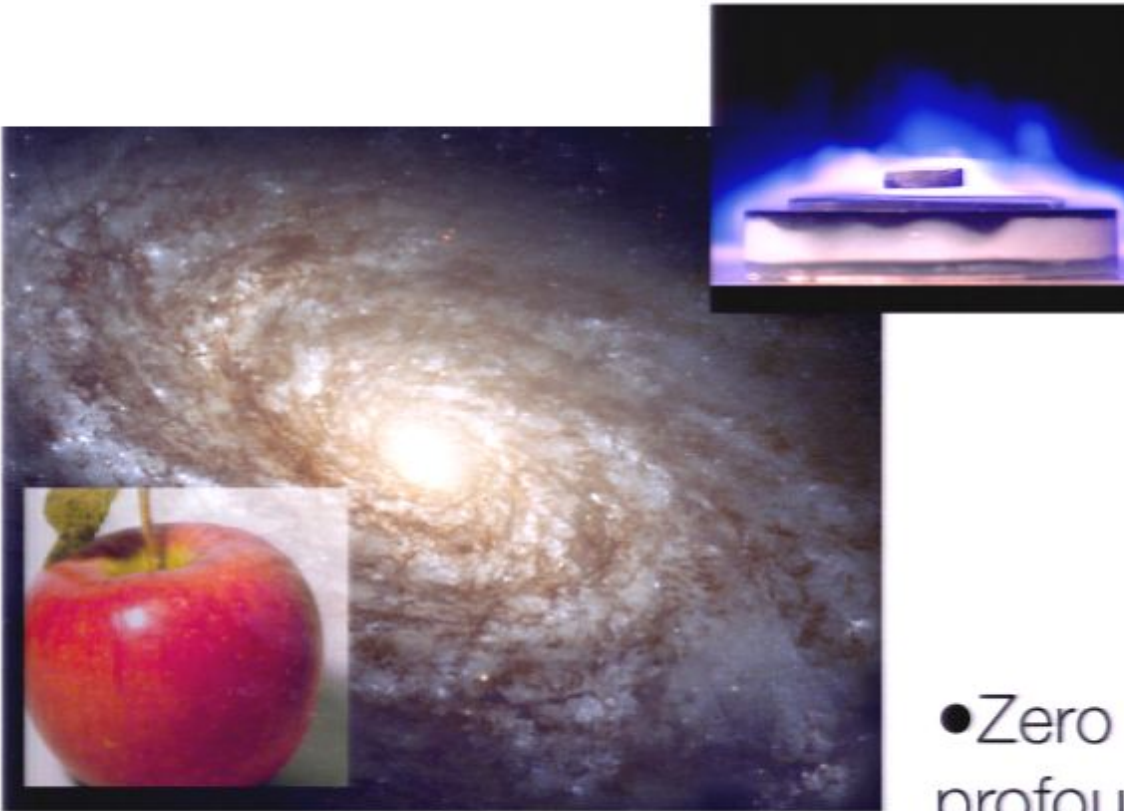


Quantum zero point fluctuations: major unsolved problem of the quantum era.



- 73% of the mass of the cosmos is “Dark Energy”: an unidentified form of zero point energy, causing the expansion to accelerate.

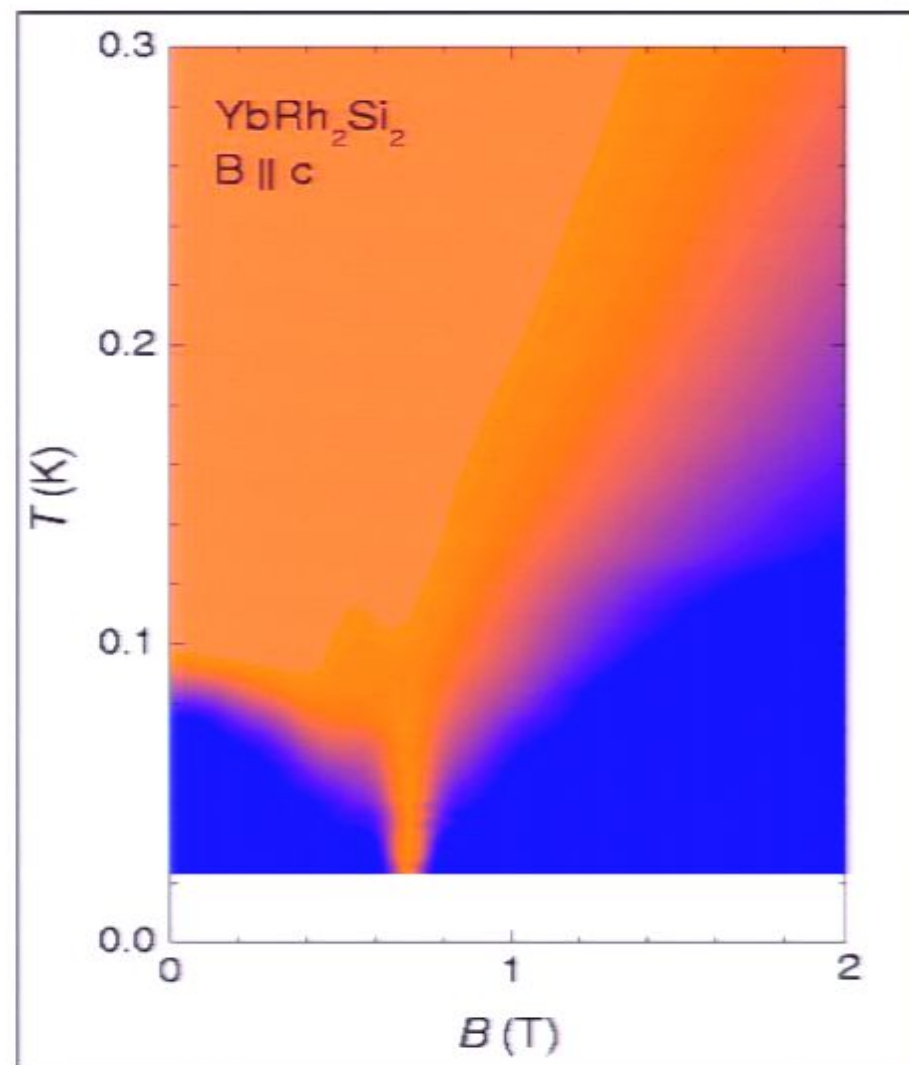
Quantum zero point fluctuations: major unsolved problem of the quantum era.



- Zero point fluctuations profoundly transform matter, endowing it with marked tendency to develop new forms of order.

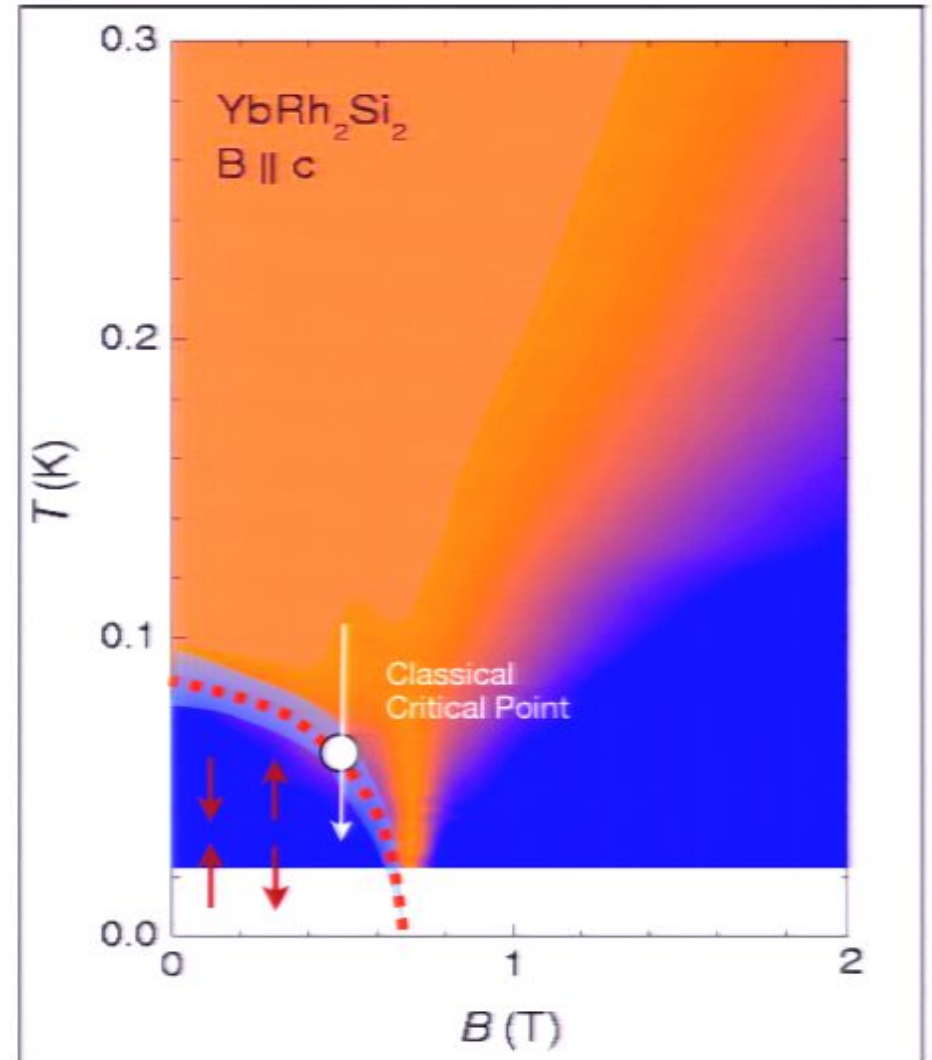
Classical Criticality

Custers et al (2002)

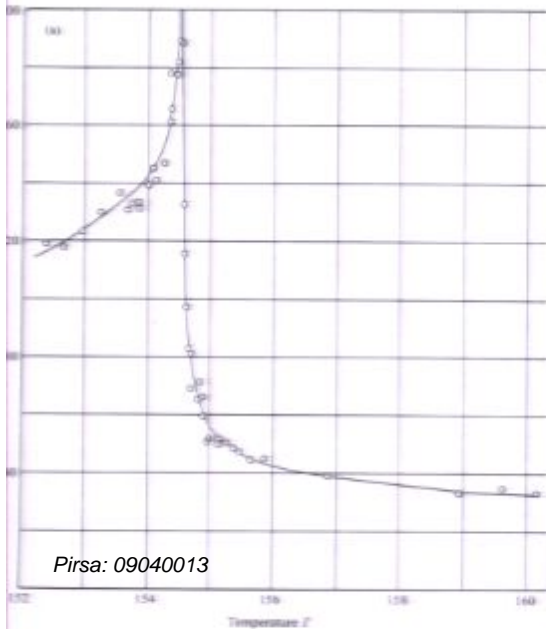


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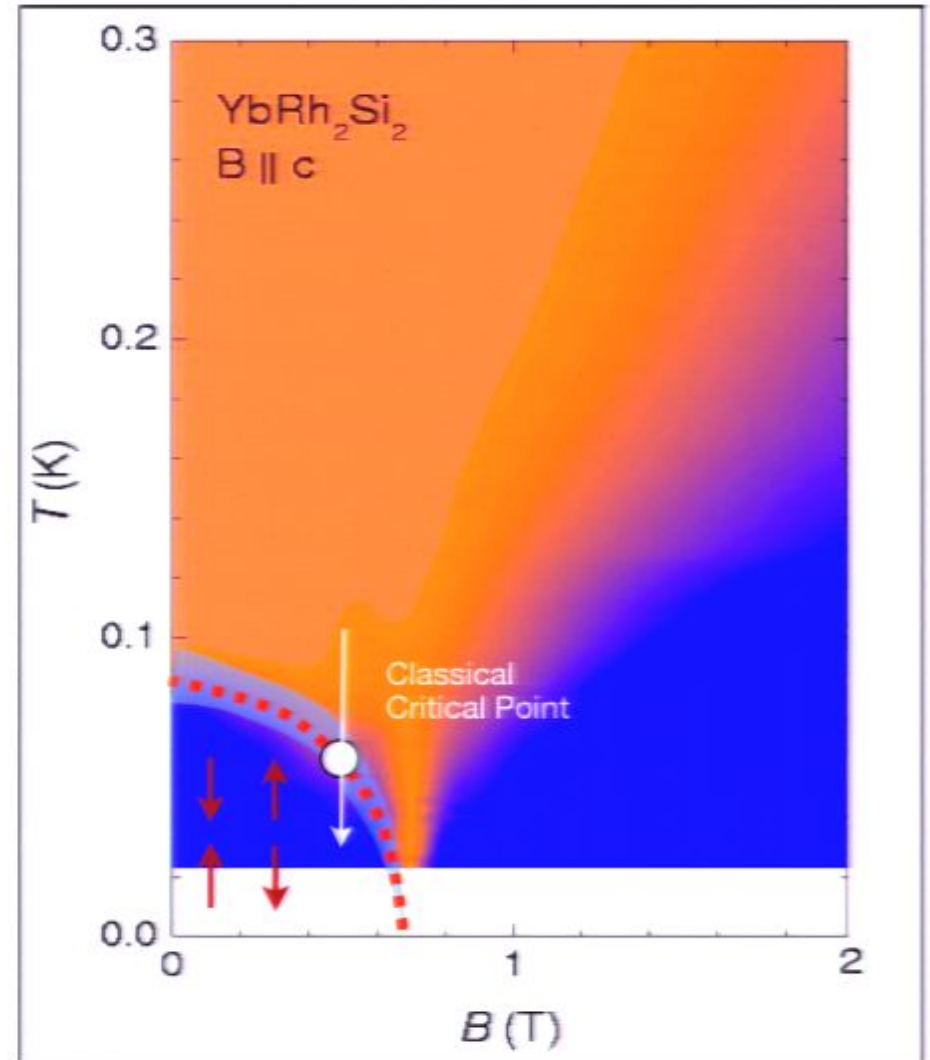
Classical Criticality



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Oxygen. (Voronel et al 1963).

Custers et al (2002)



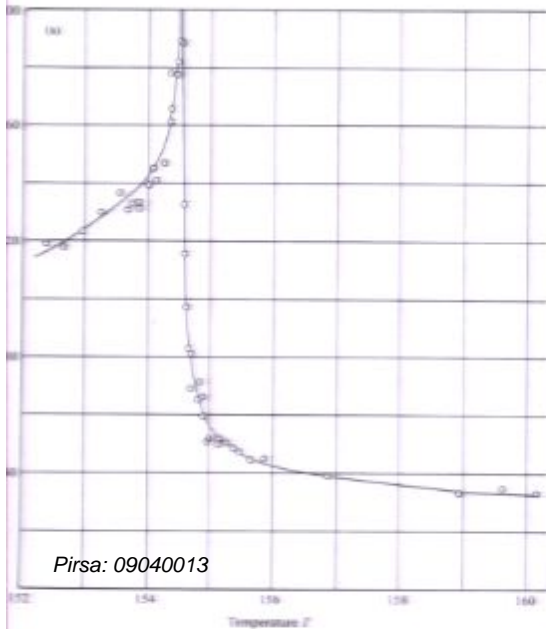
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Classical Criticality

Michael Fisher



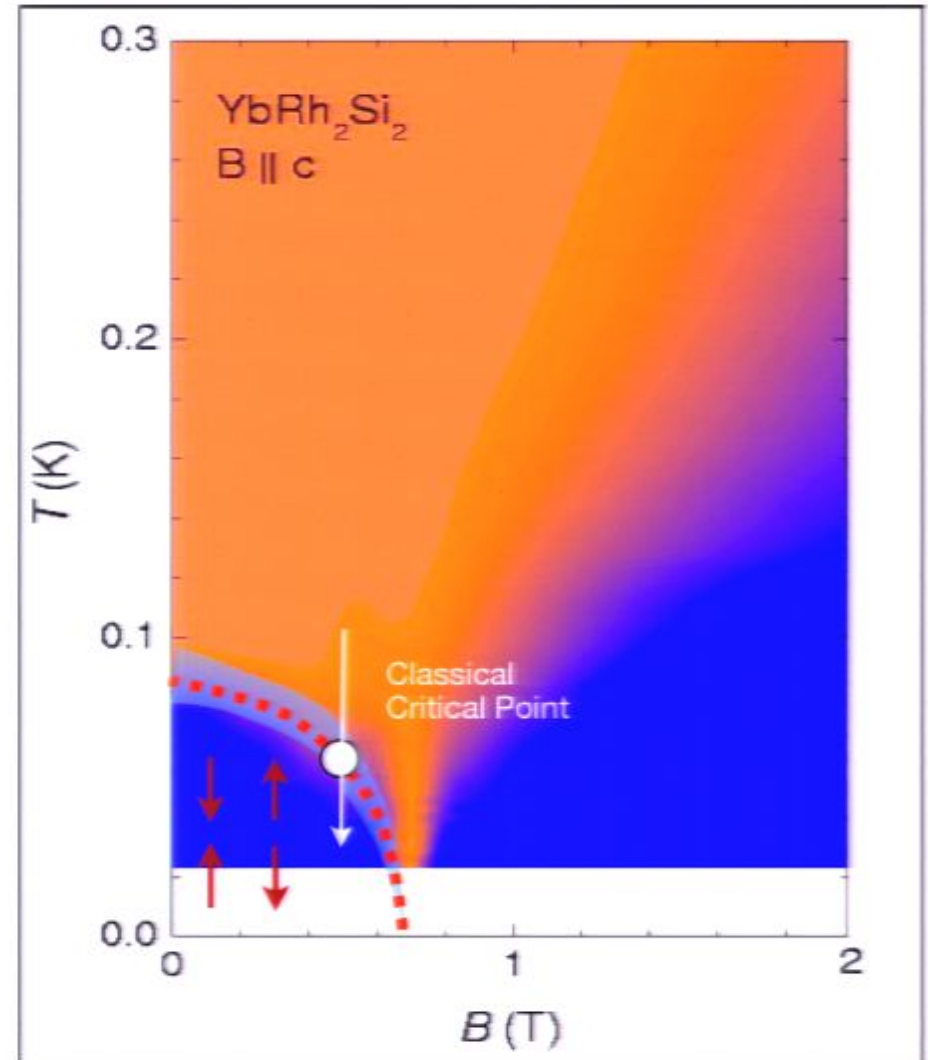
"New insights into physics often come from revisiting areas once thought to be closed." Michael Fisher.



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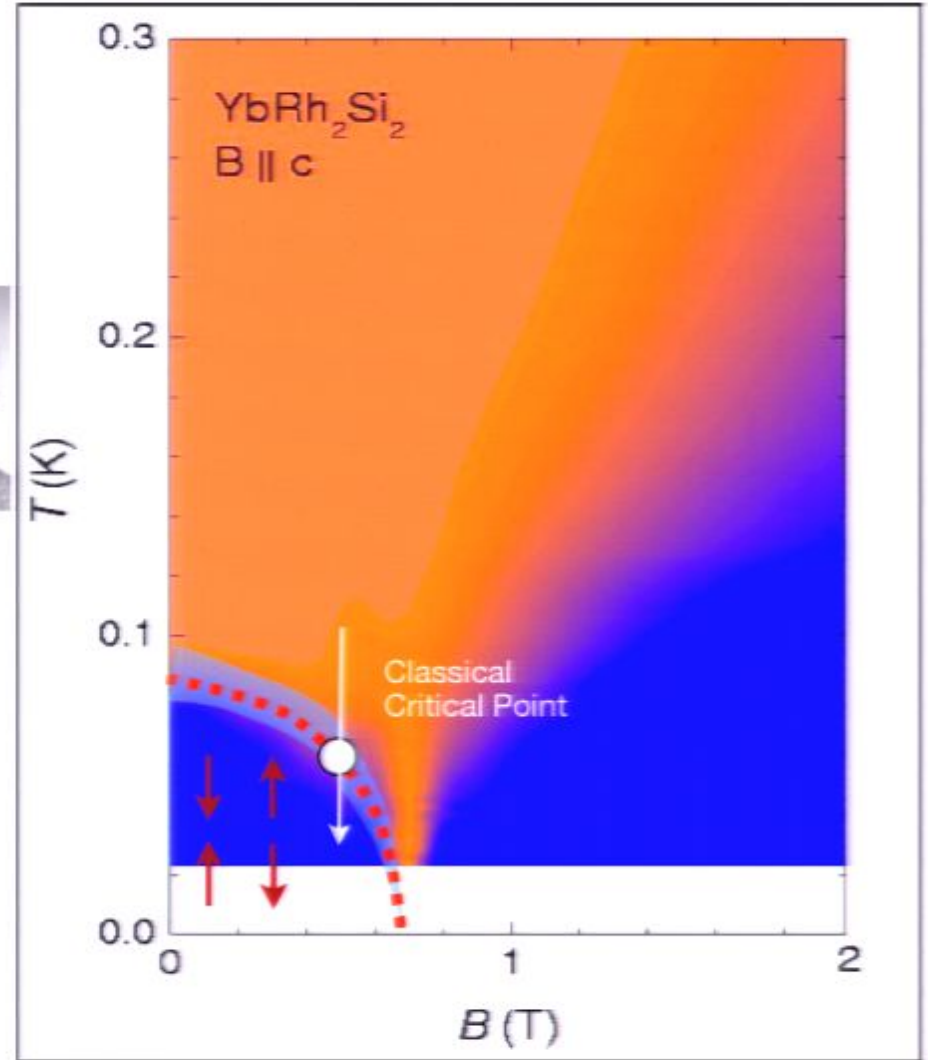
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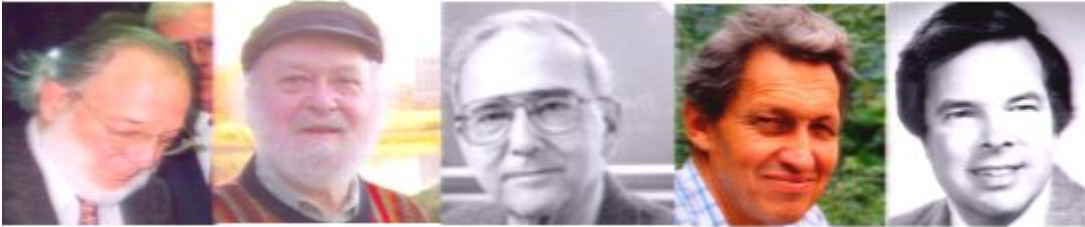
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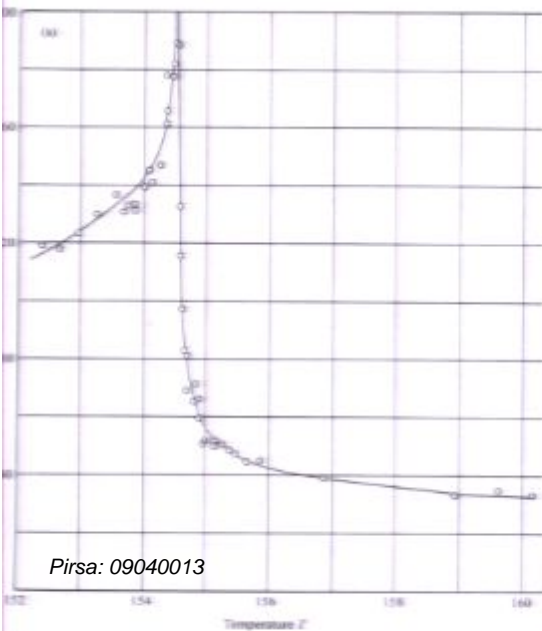
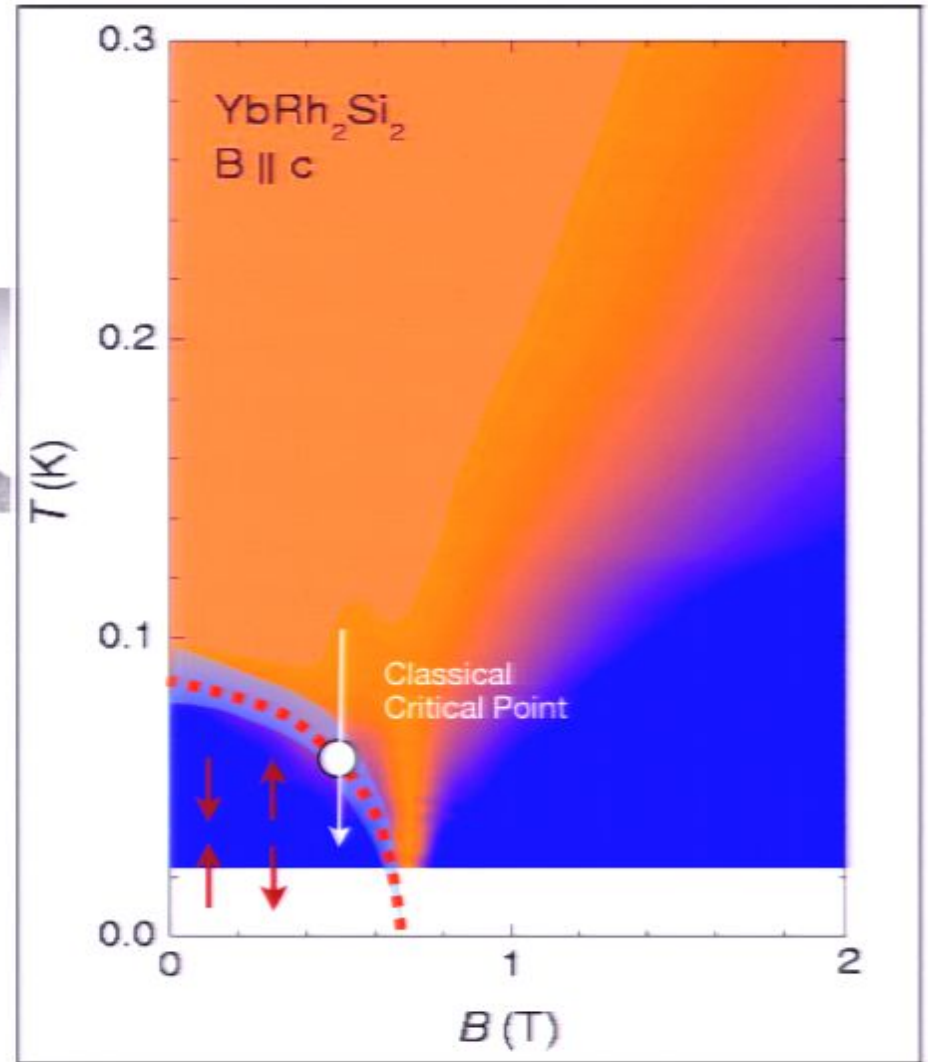
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Oxygen. (Voronel et al 1963).

“20th Century Revolution”

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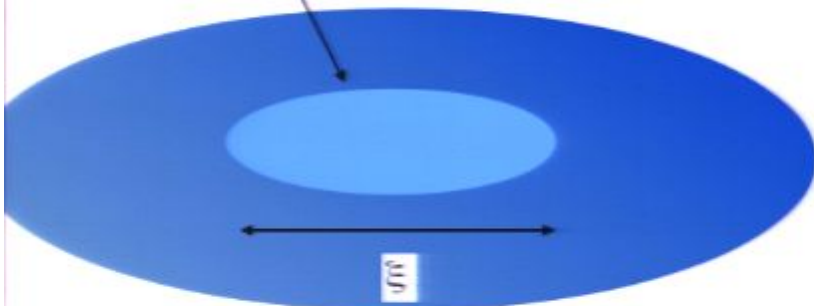
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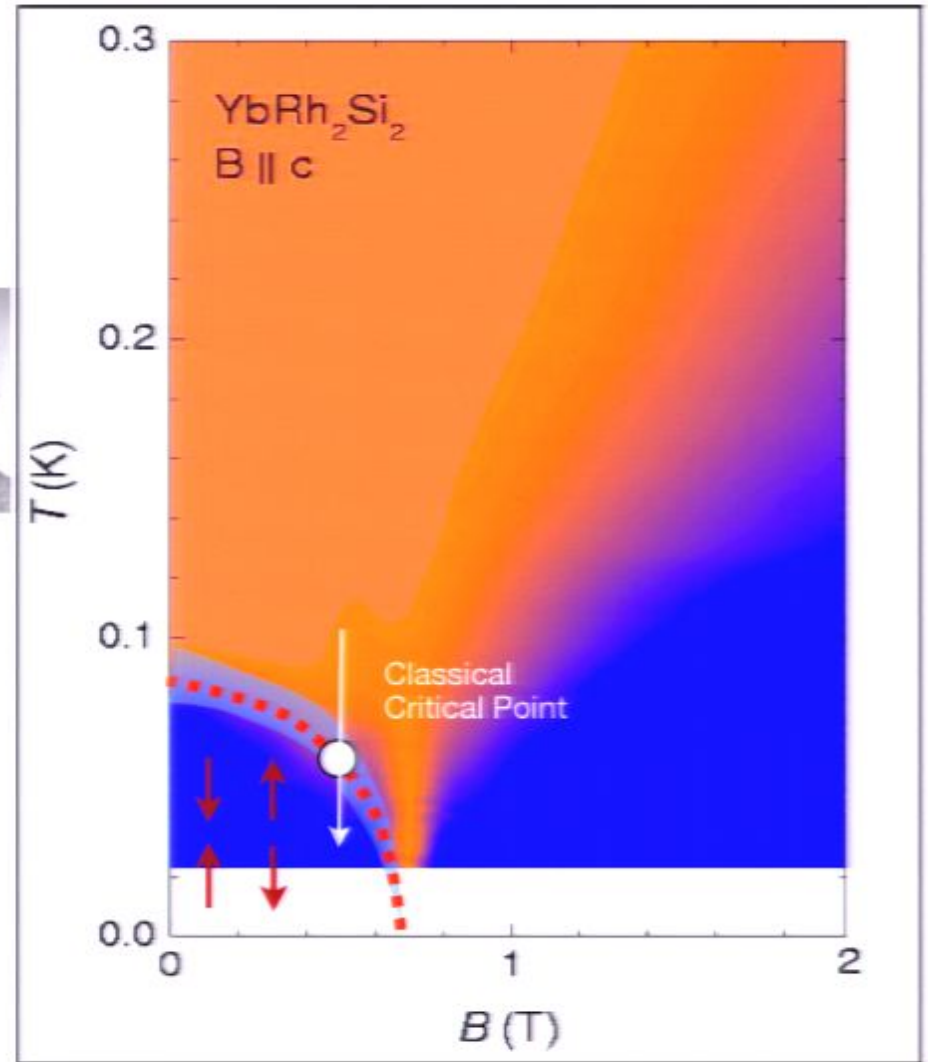
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$$\langle \psi(x)\psi(0) \rangle \sim \frac{1}{x^{d-2+\eta}}$$

Critical matter



Custers et al (2002)



Classical Criticality

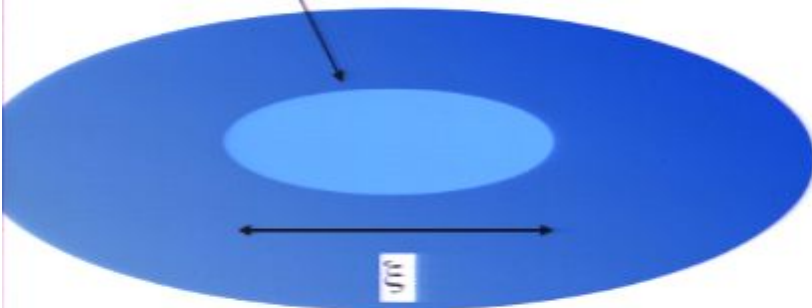
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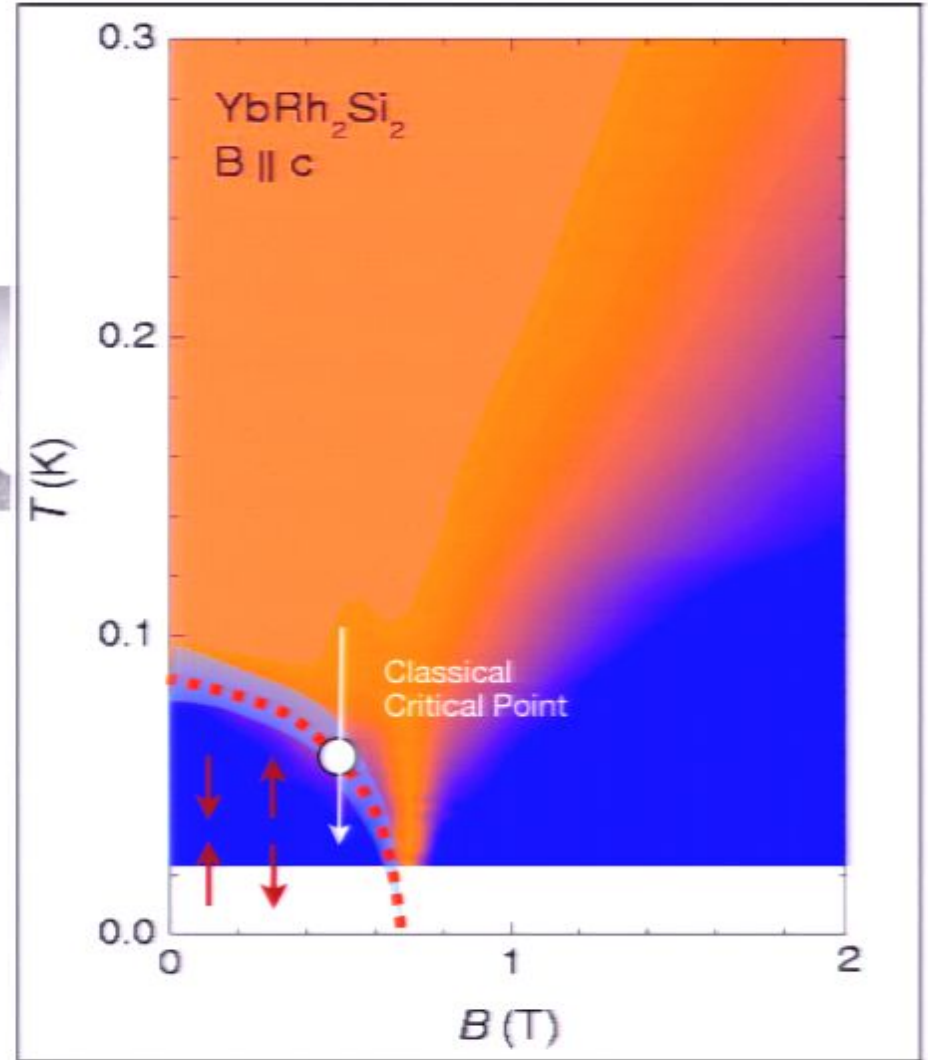
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Critical matter - universal



Custers et al (2002)

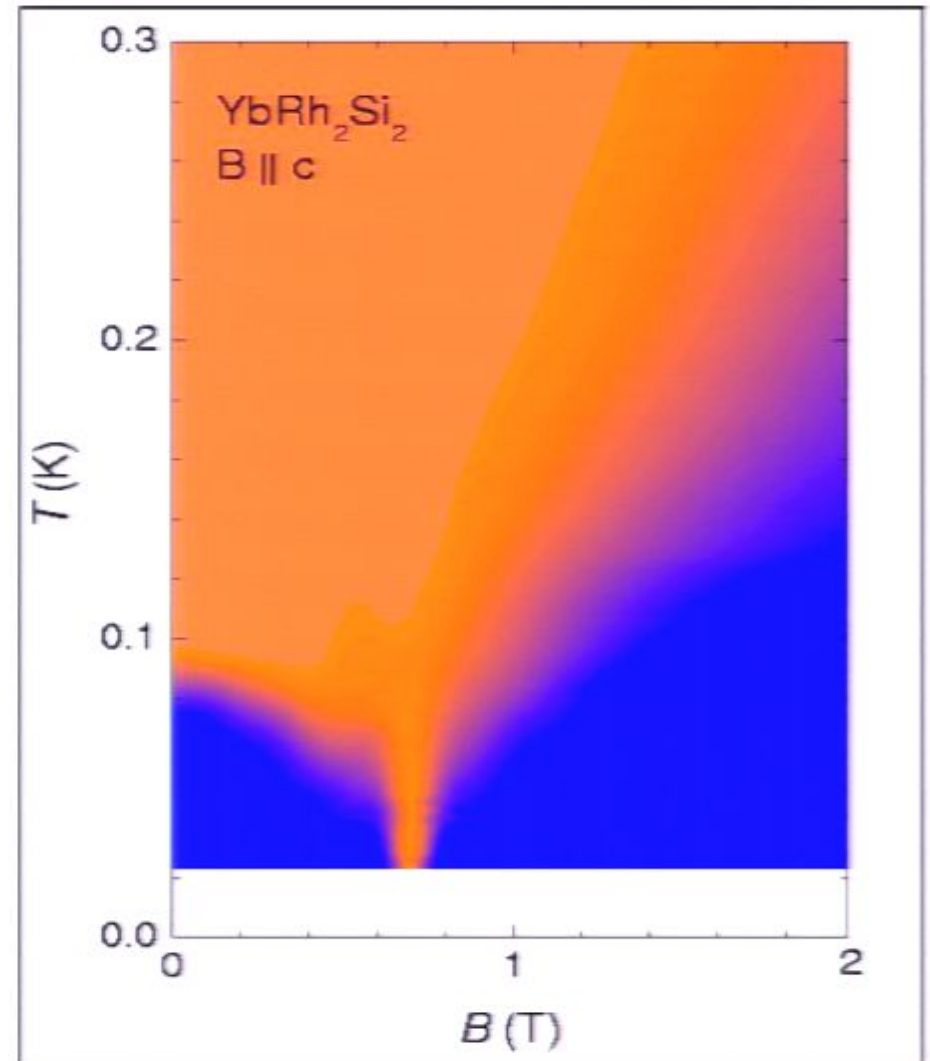


“20th Century Revolution”

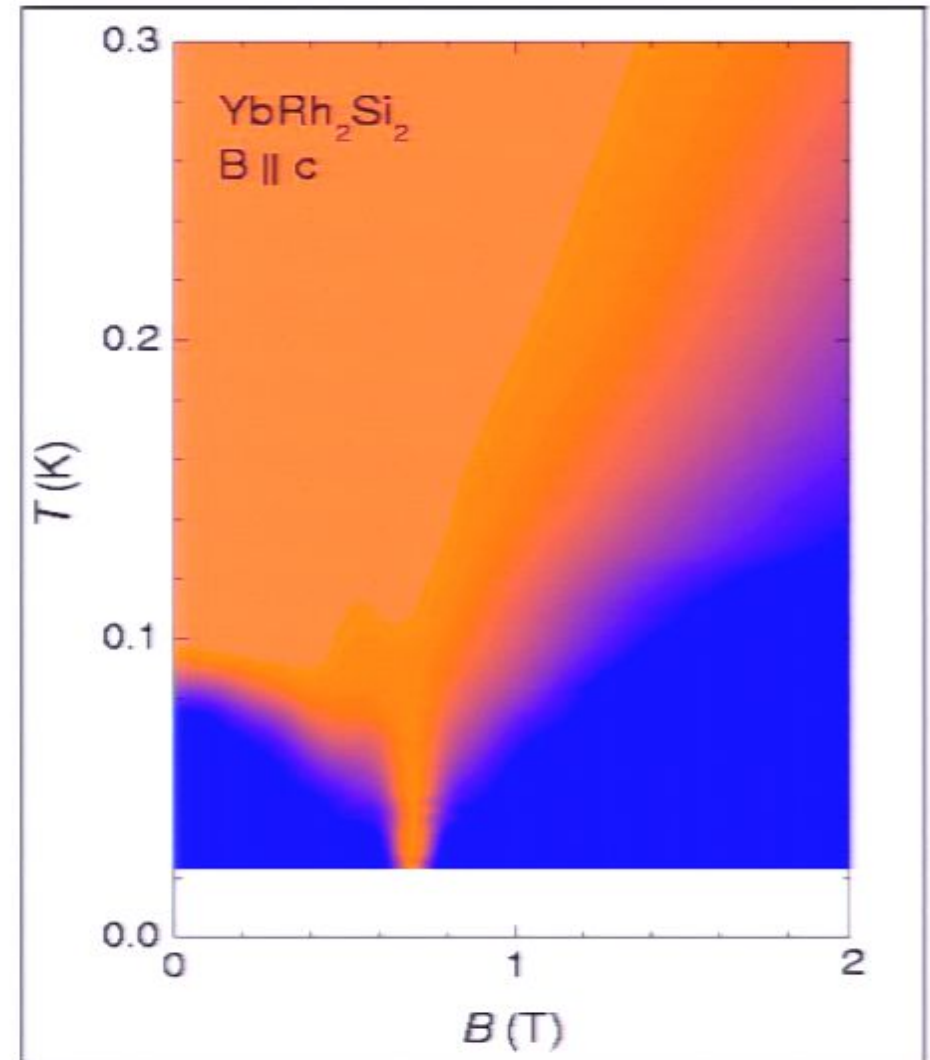
Qu-Transition

driven by zero point energy.

Custers et al (2002)



Q_y-Transition
Phase transition
driven by zero point energy.

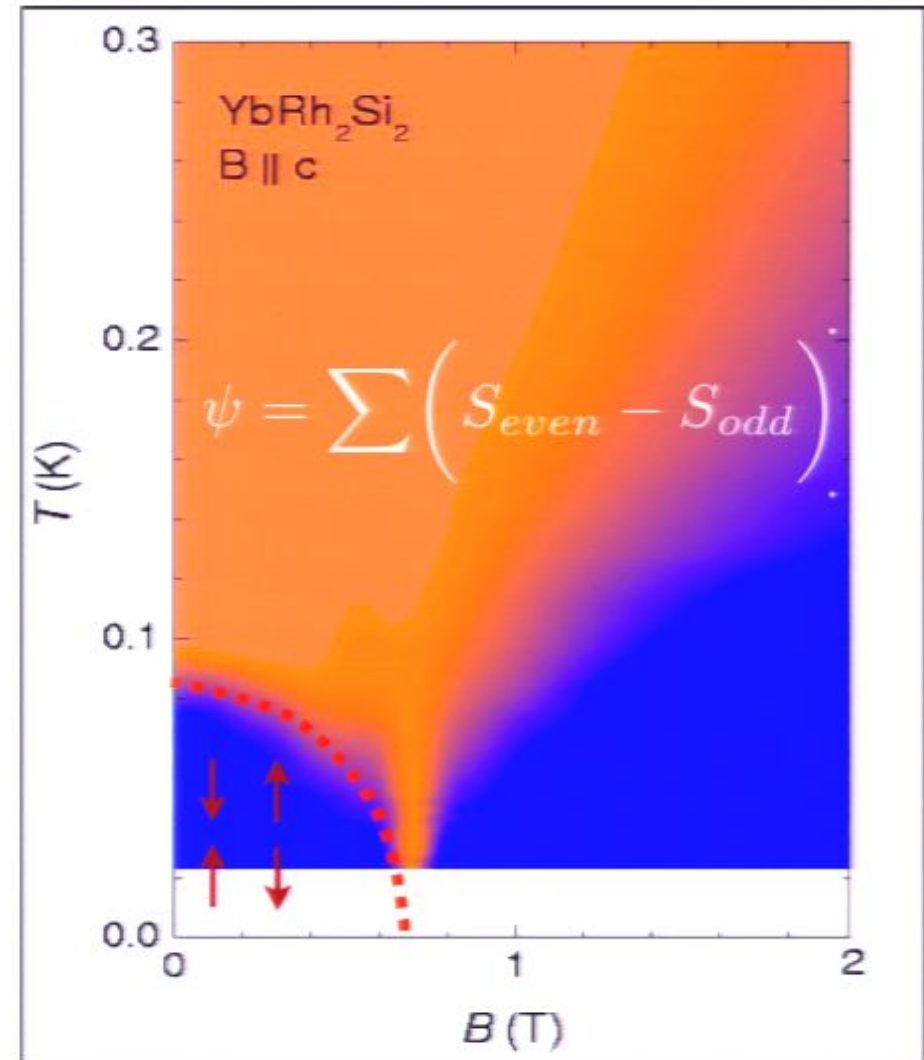


Q₀-Transition

Phase transition

driven by zero point energy.

$$[H, \psi] = 0$$

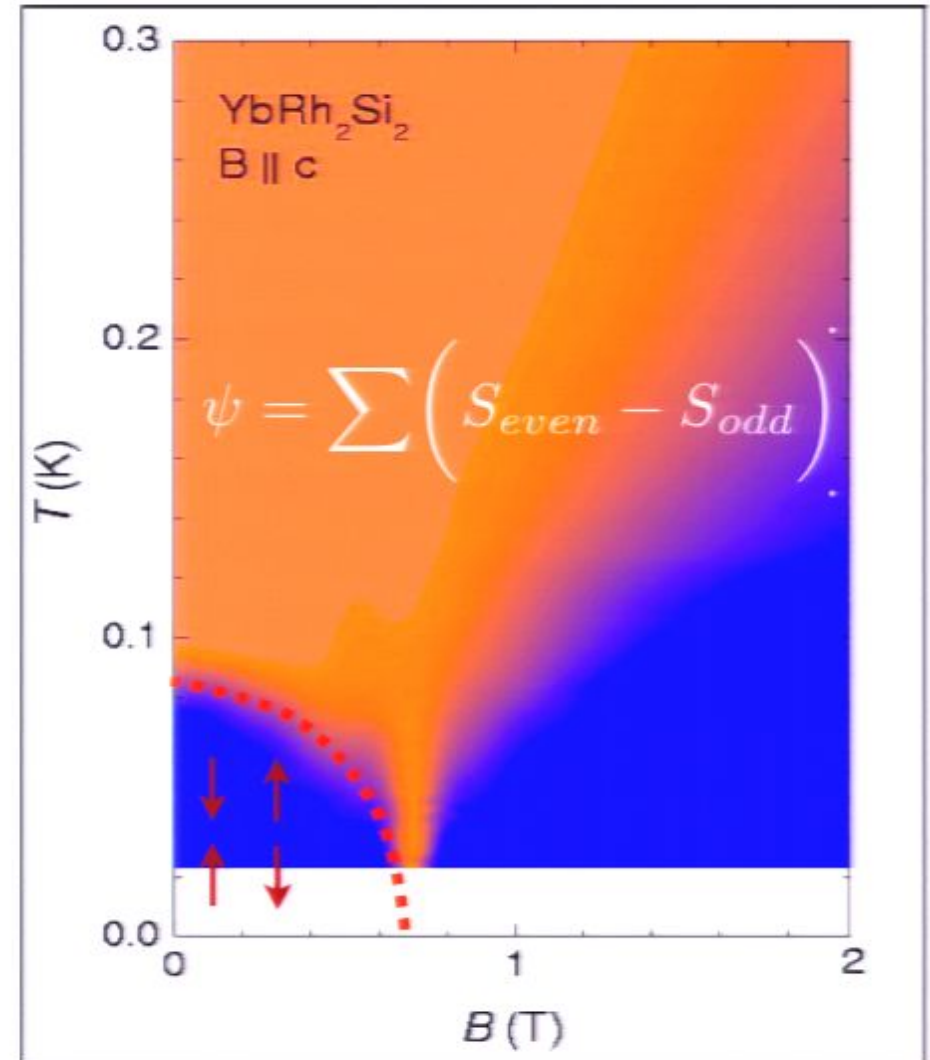


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Q_v-Transition

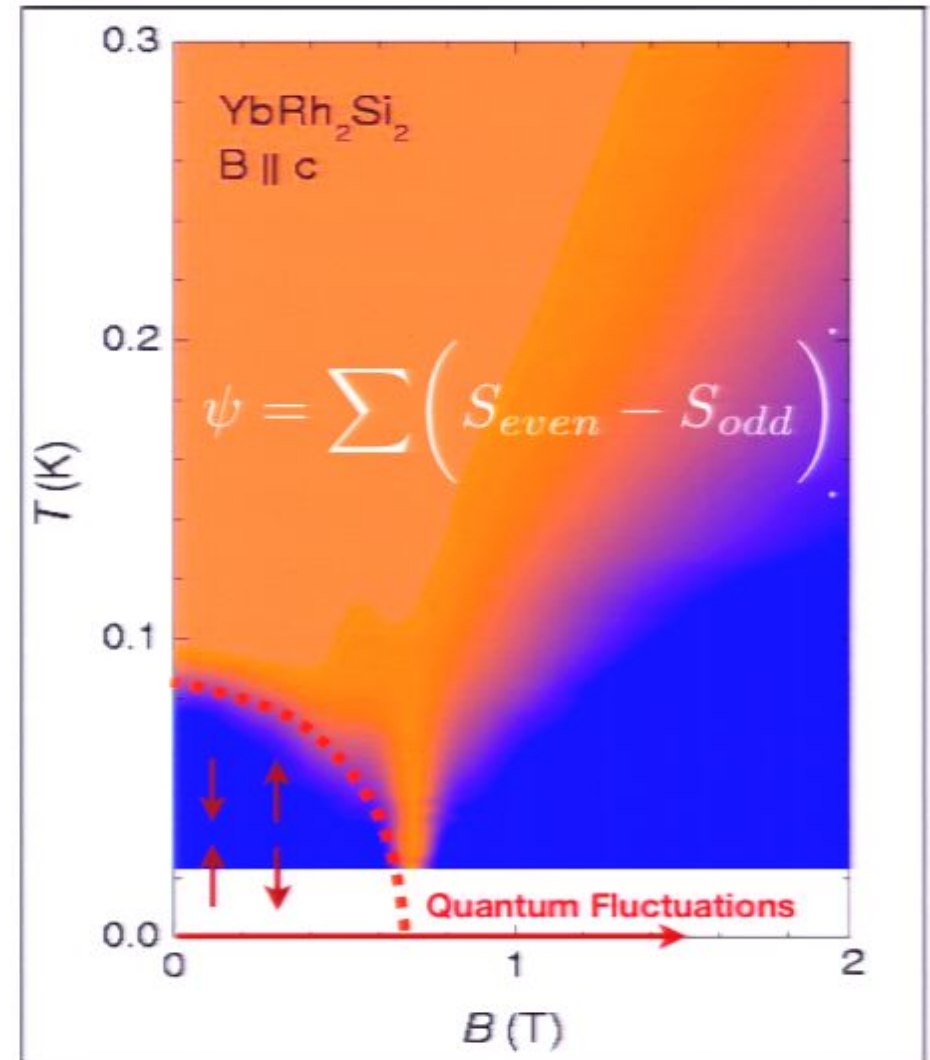
Phase transition

driven by zero point energy.

$$[H, \psi] \neq 0$$

What happens when the time and length scale of coherent fluctuations expands to fill the entire material?

Pisa: 09040073



Q_y-Transition

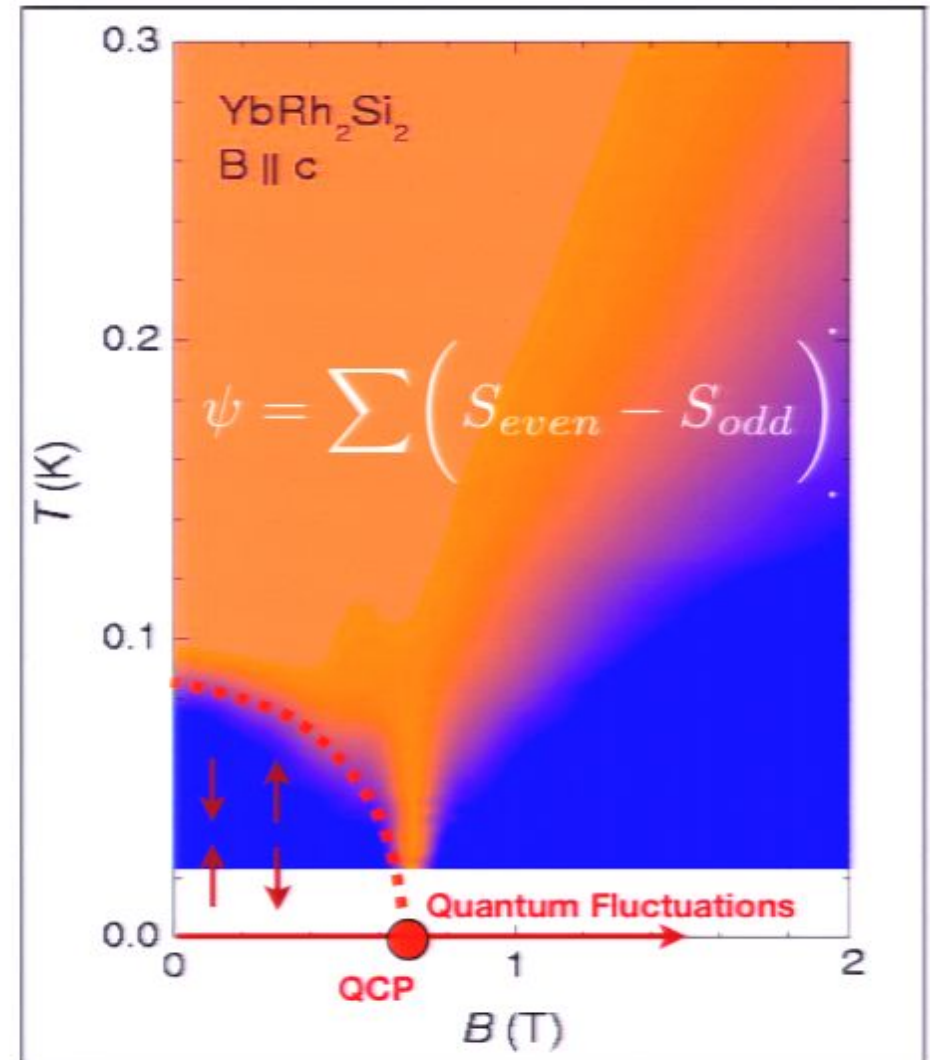
Phase transition

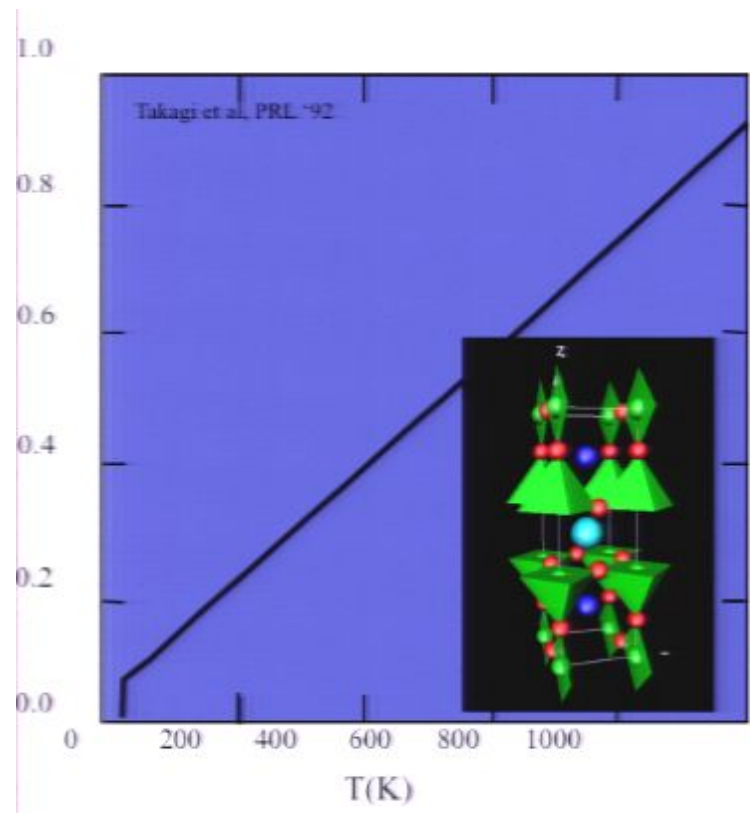
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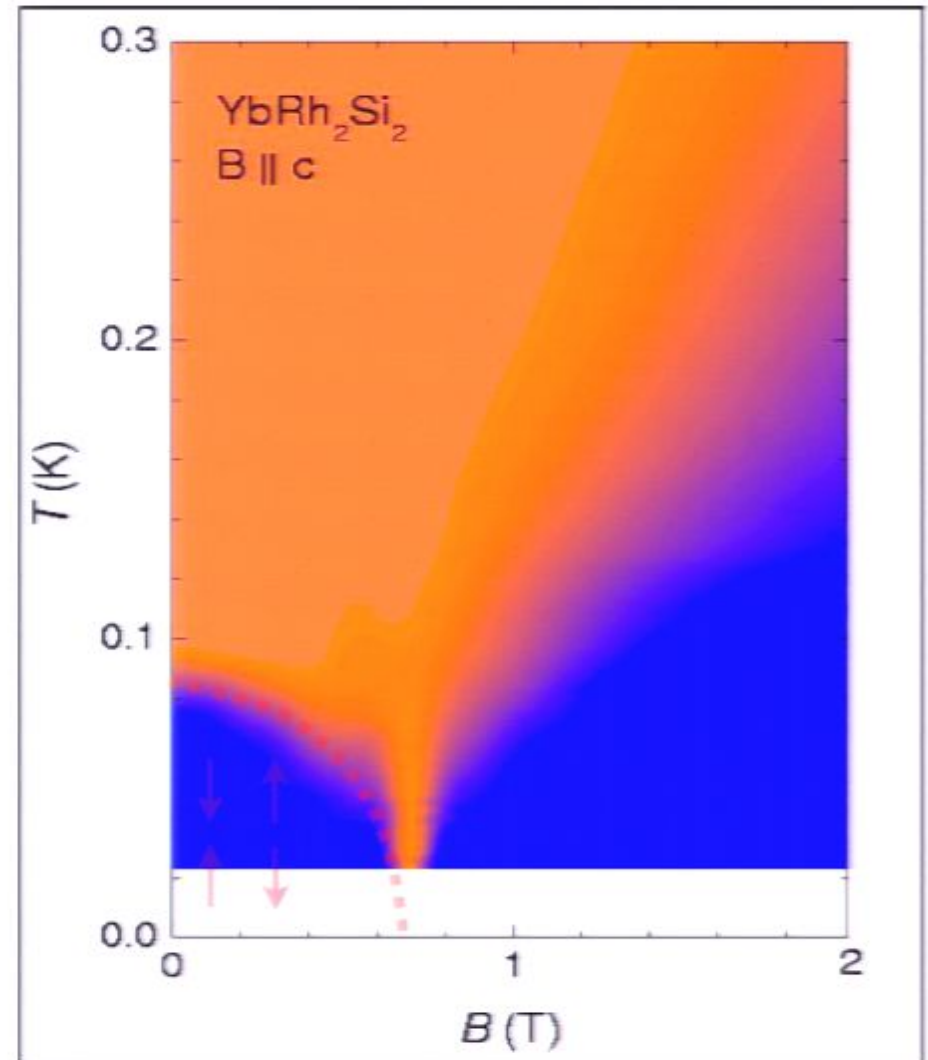




Cuprates
 $T_c = 11-92\text{K}$

Q_y-Transition

Phase transition
driven by zero point energy.

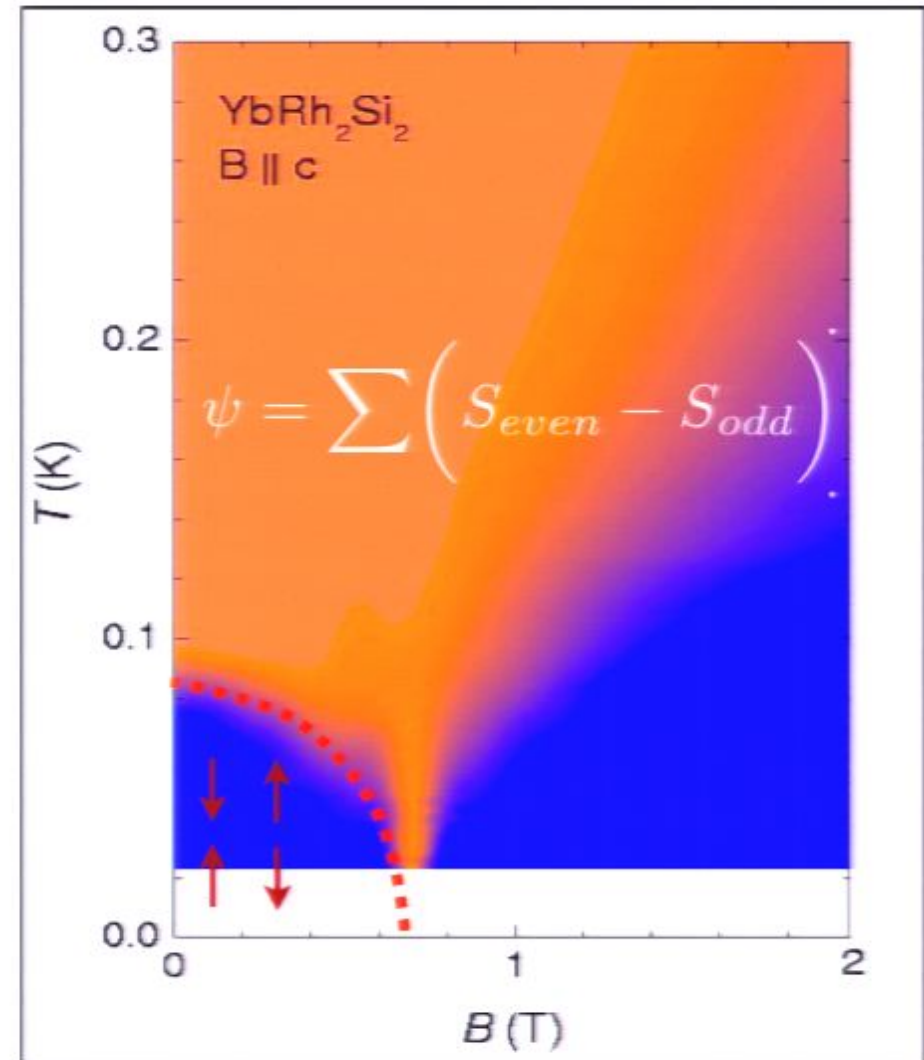


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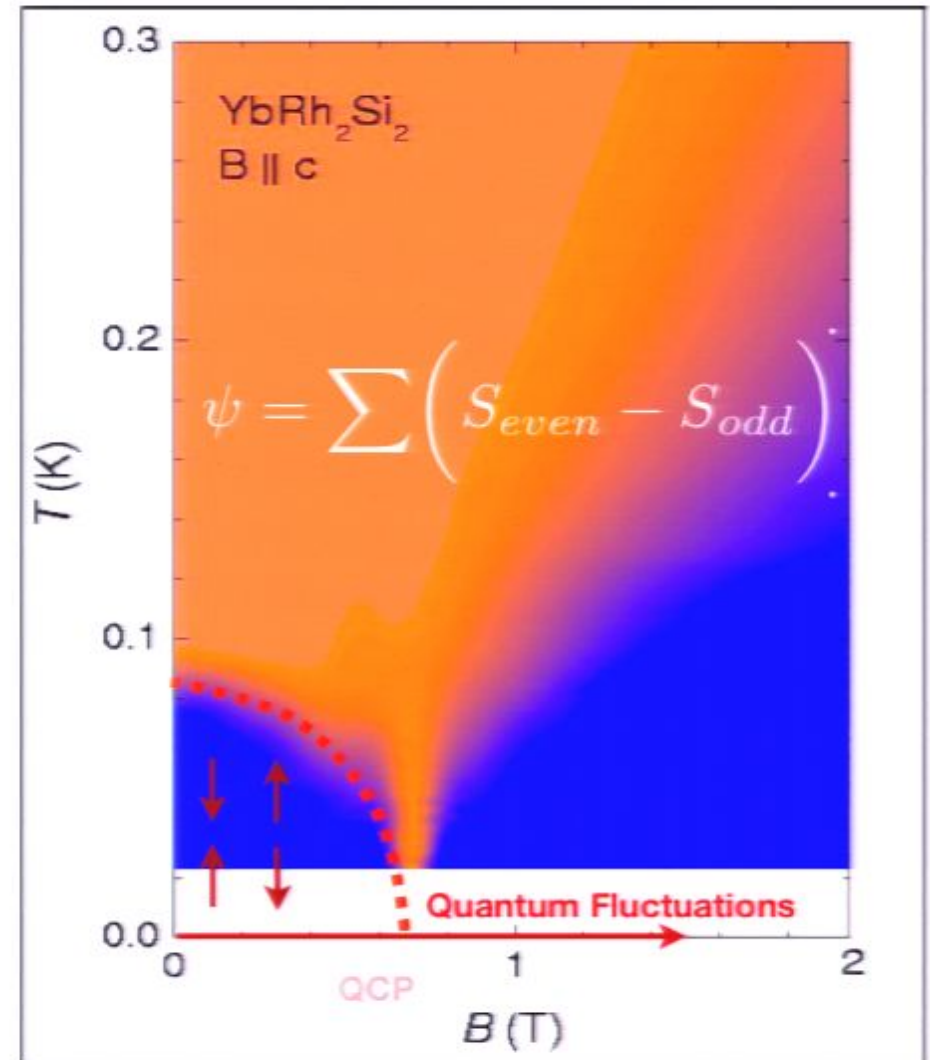
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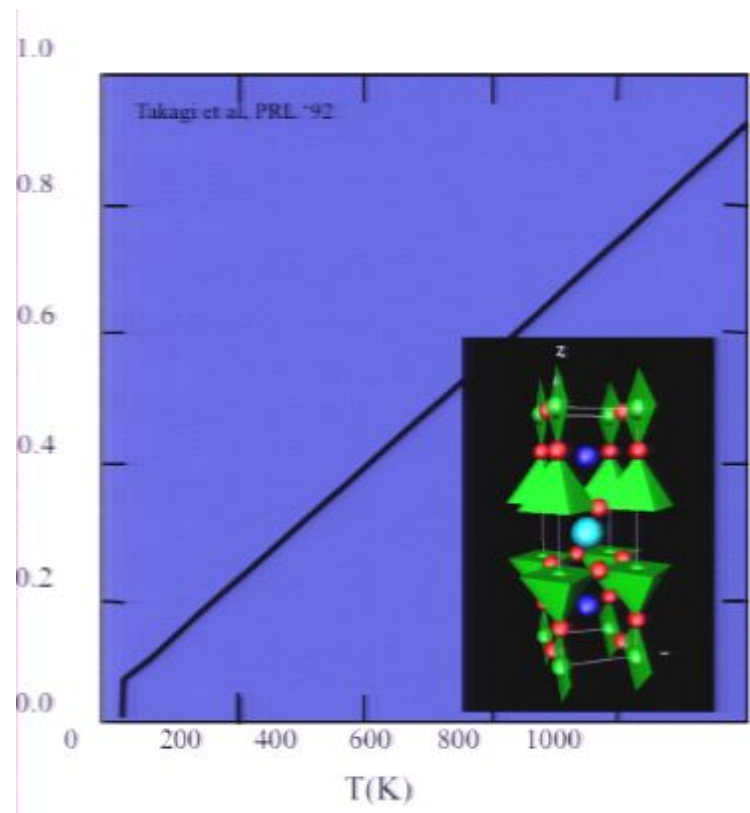
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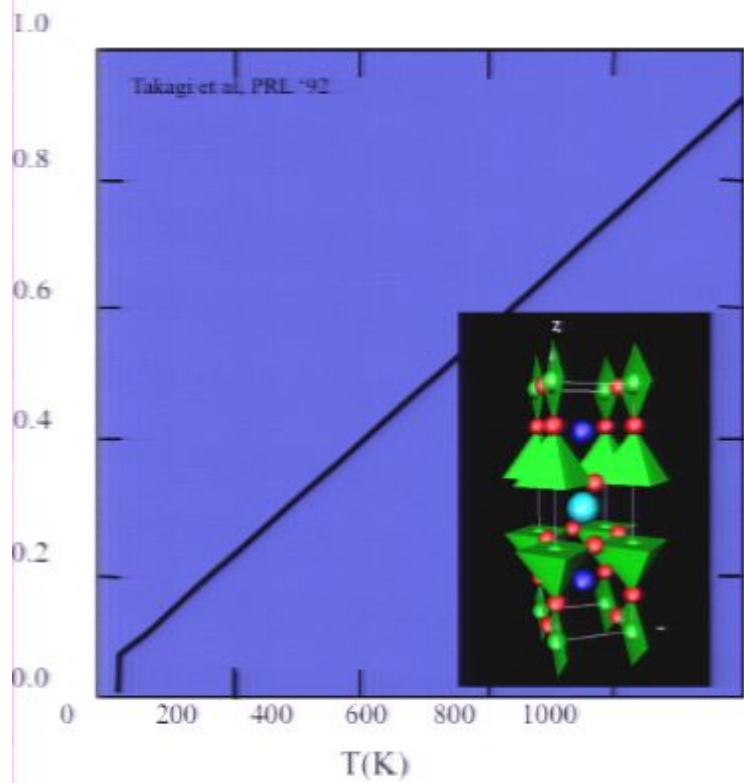


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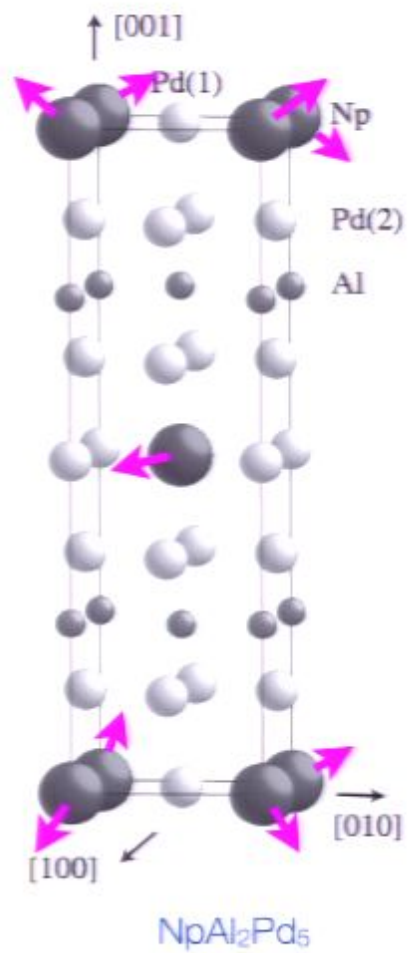




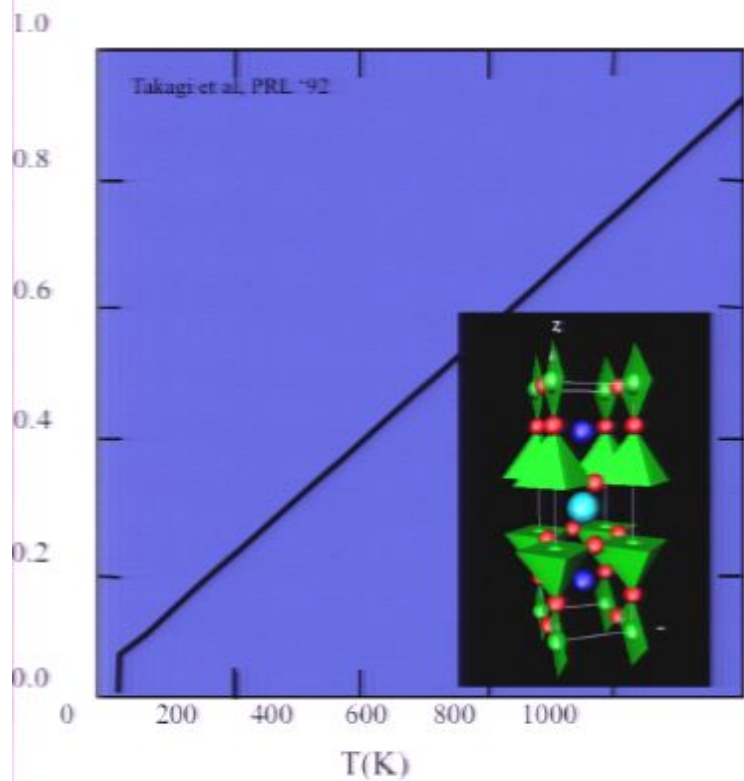
Cuprates
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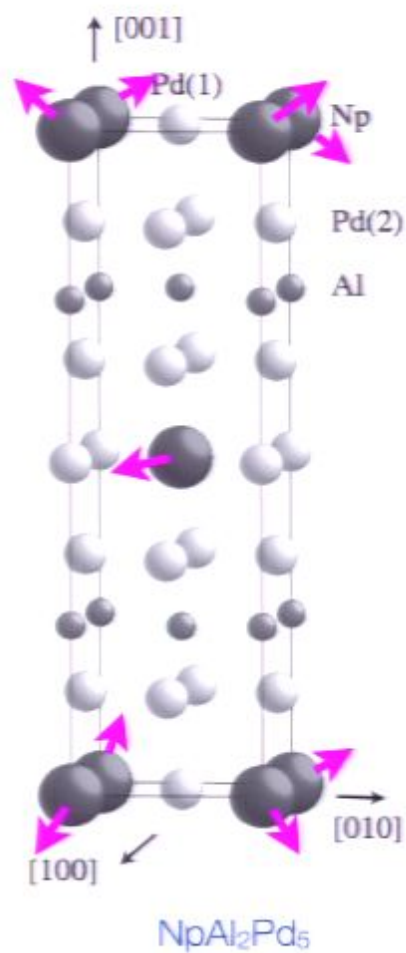
Cuprates
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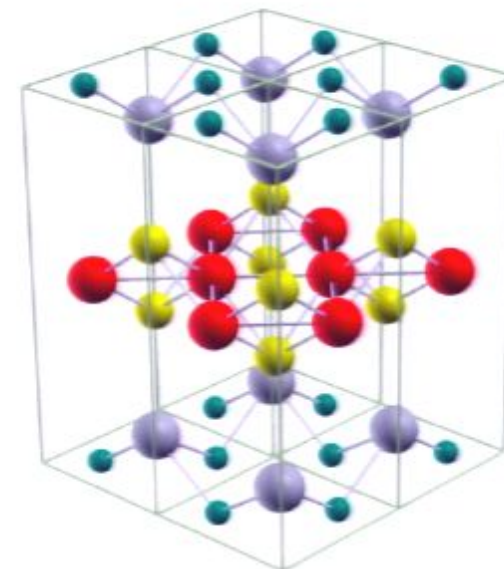
HF 115s
 $T_c = 0.2 - 18.5\text{K}$



Cuprates
 $T_c = 11 - 92 \text{ K}$

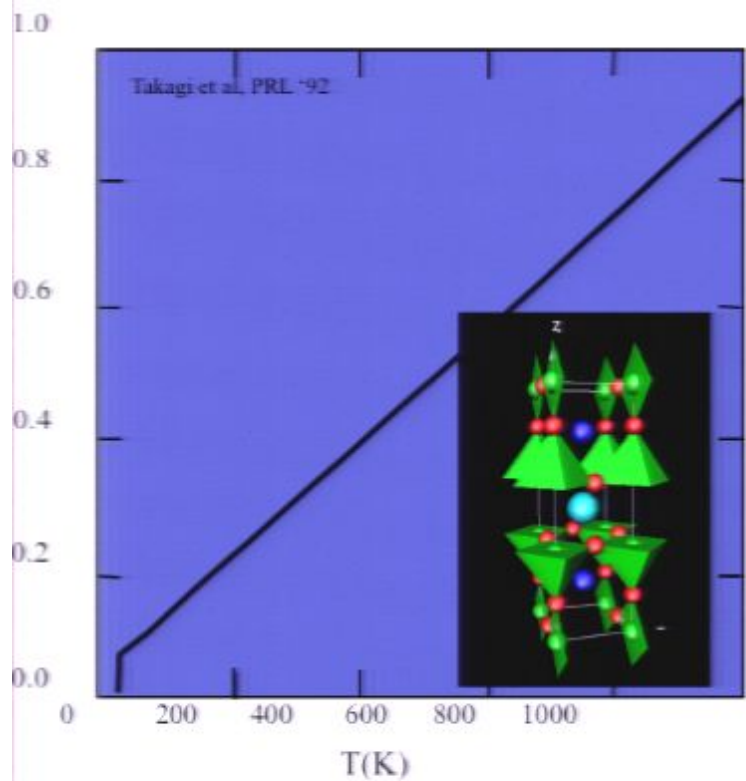


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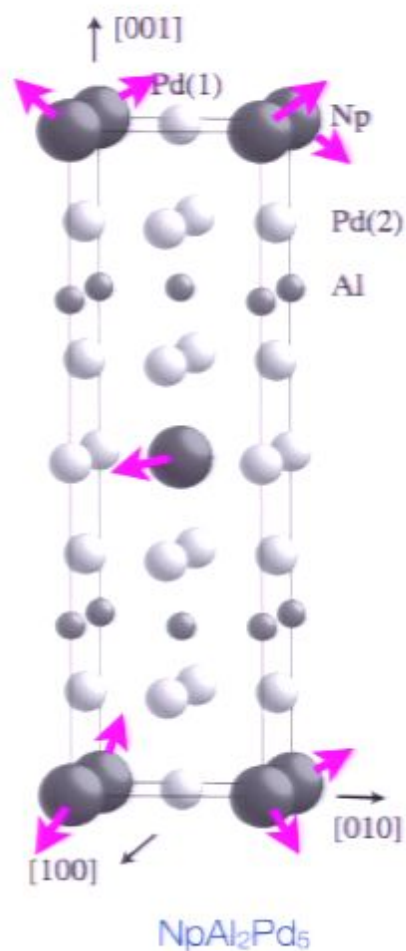


$\text{PrF}_x\text{O}_{1-x}\text{FeAs}$
 Z.A. Ren et.al. Beijing. (08)

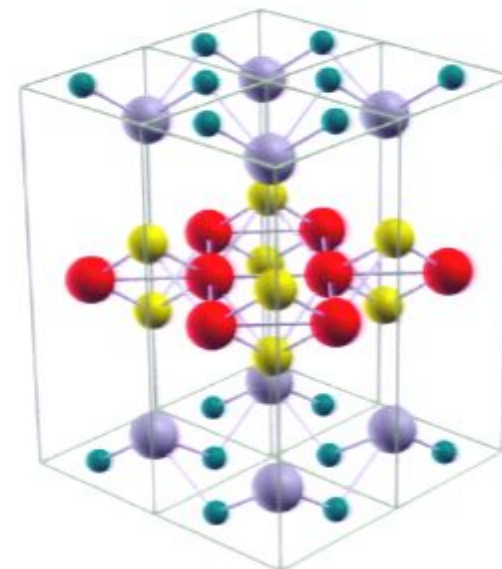
Iron oxypnictides
 $T_c = 6 - 53 \text{ ++ ? K}$



Cuprates
 $T_c = 11 - 92\text{K}$



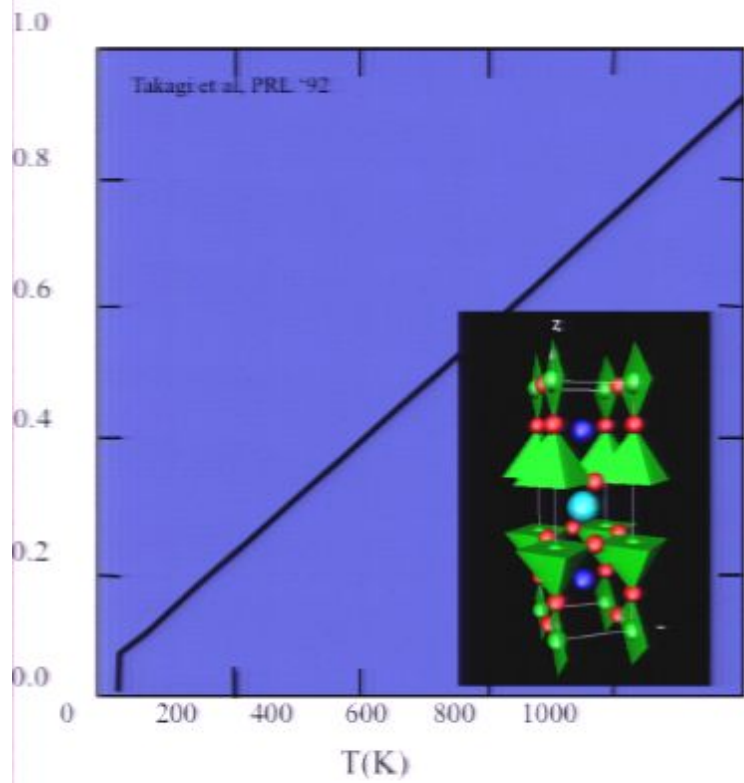
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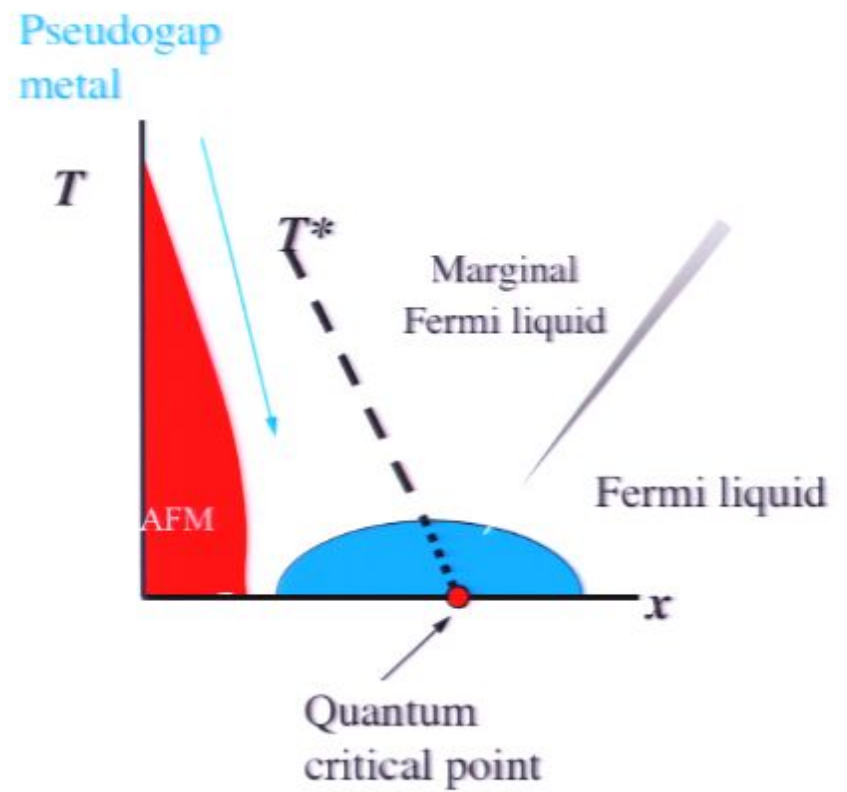
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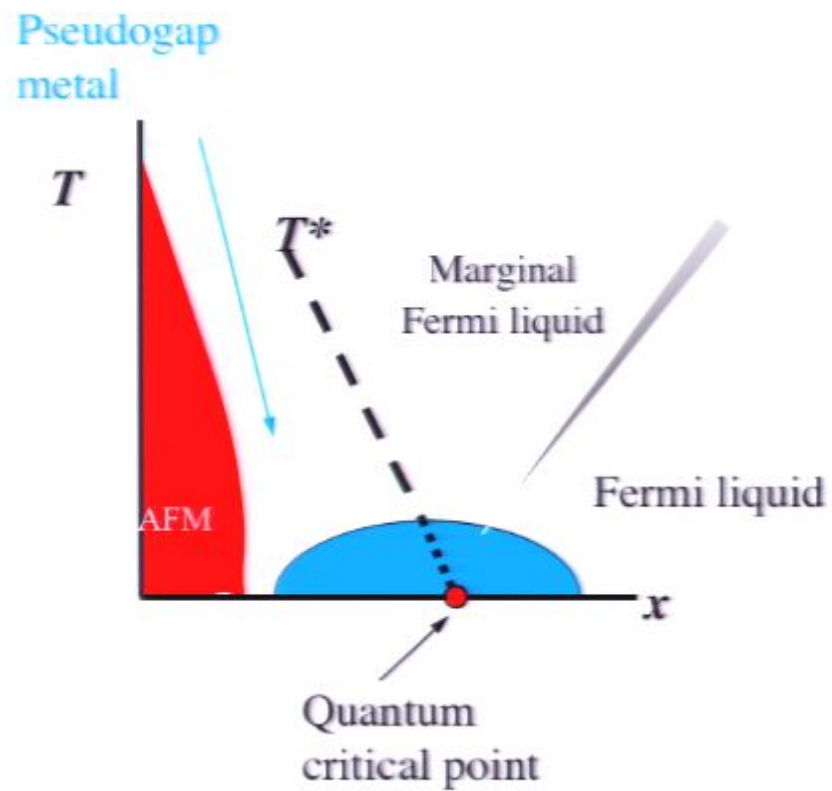
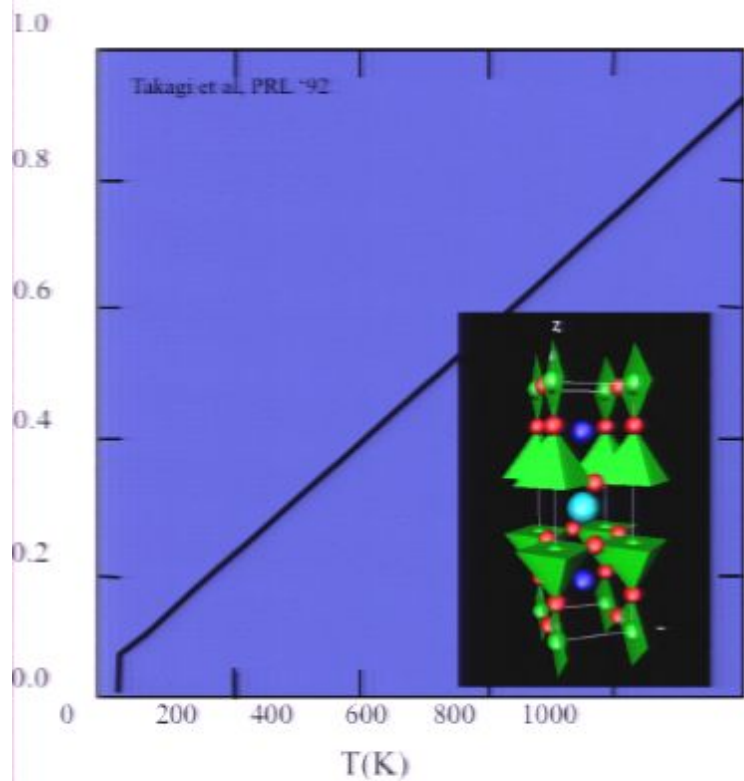
Iron oxypnictides
 $T_c = 6 - 53 \text{ ++ ? K}$

What kind of electronic fluid yields
 high T_c superconductivity?



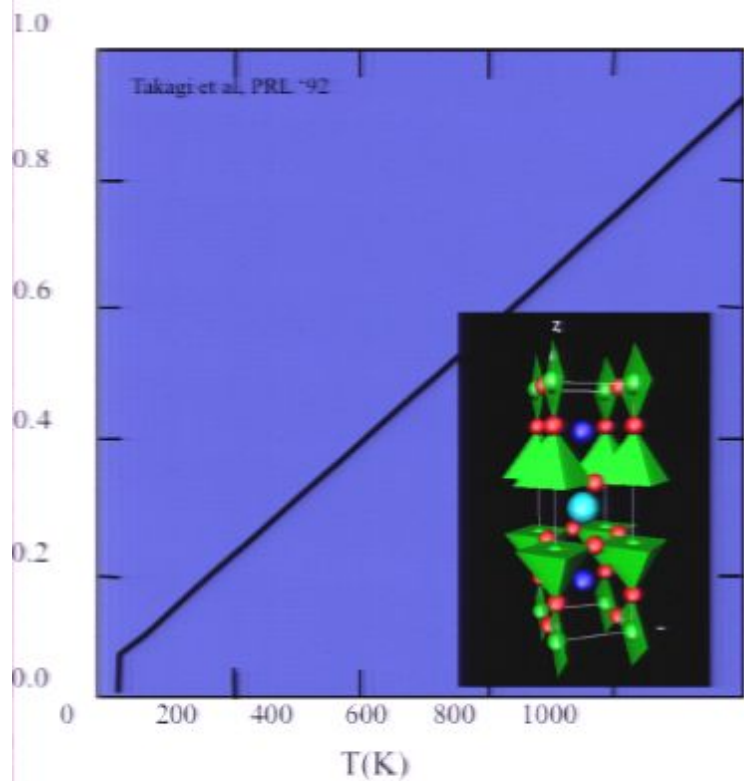
Cuprates
 $T_c = 11-92\text{K}$





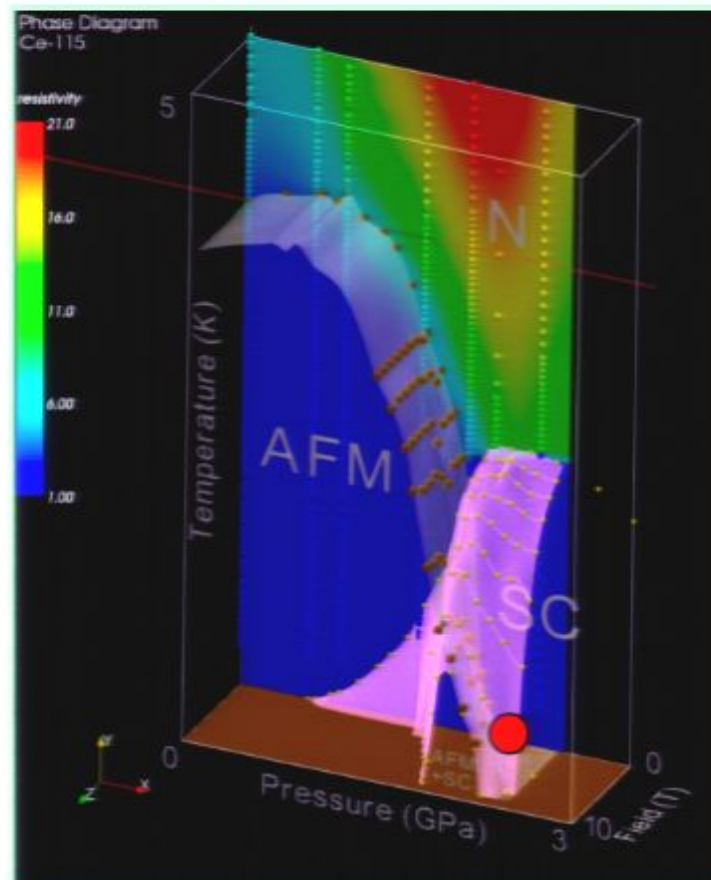
Cuprates
 $T_c = 11-92K$

“Avoided criticality”

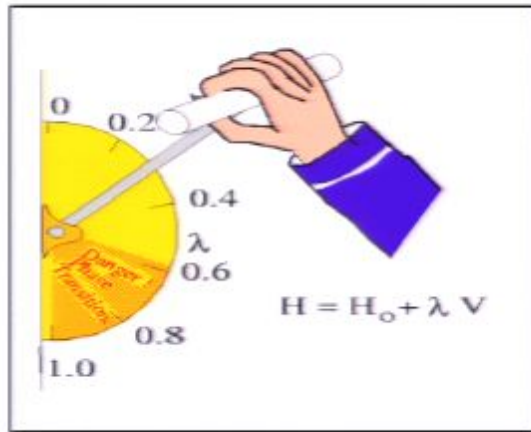


Cuprates
 $T_c = 11-92\text{K}$

“Avoided criticality”

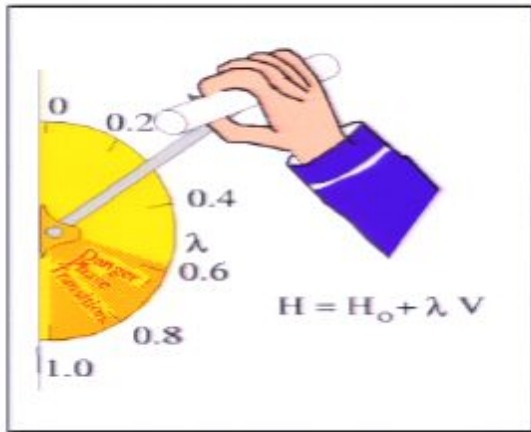


Tuson Park, (2007).



Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

Landau, JETP 3, 920 (1957)



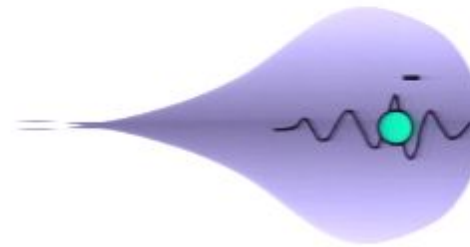
Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

Landau, JETP 3, 920 (1957)



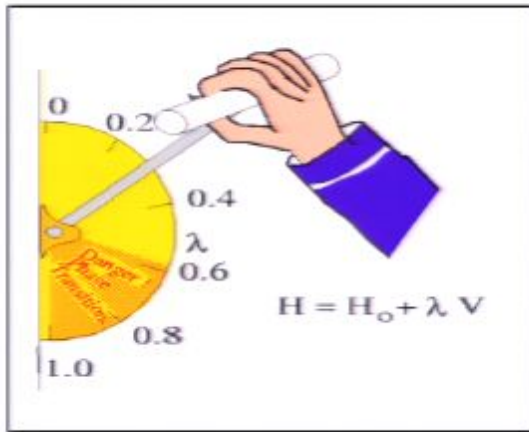
$|e^- \rangle$

Interactions
adiabatically



$|qp^- \rangle$

“Quasiparticle”



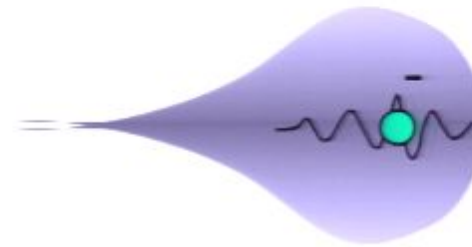
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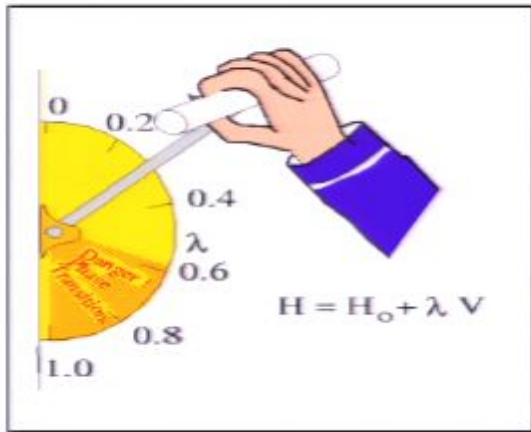
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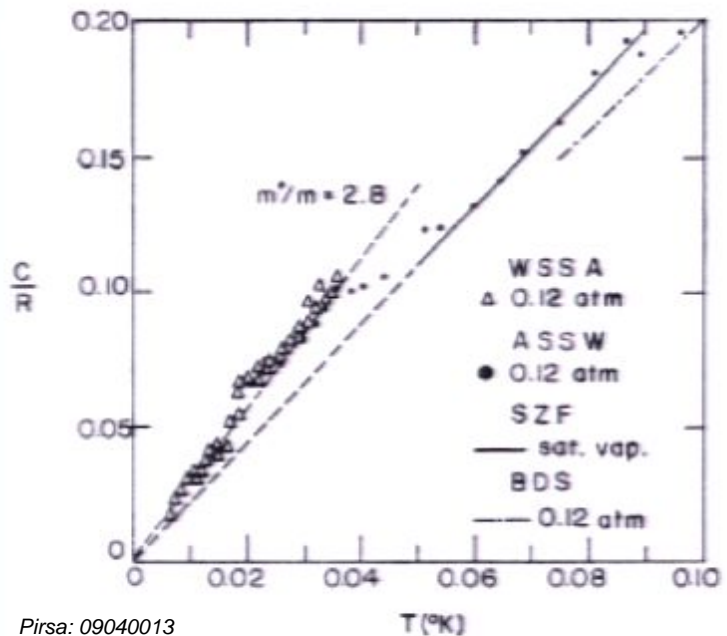
$$\frac{m^*}{m} = \frac{N(0)^*}{N(0)} = 1 + \frac{F_1^s}{3}$$



Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

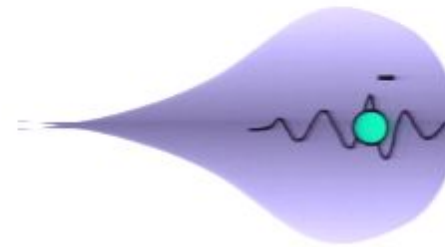
Landau, JETP 3, 920 (1957)

He-3 (1950/60s)
(Fairbanks, many others)



Pirsa: 09040013

“Quasiparticle”



$$|qp^-\rangle$$

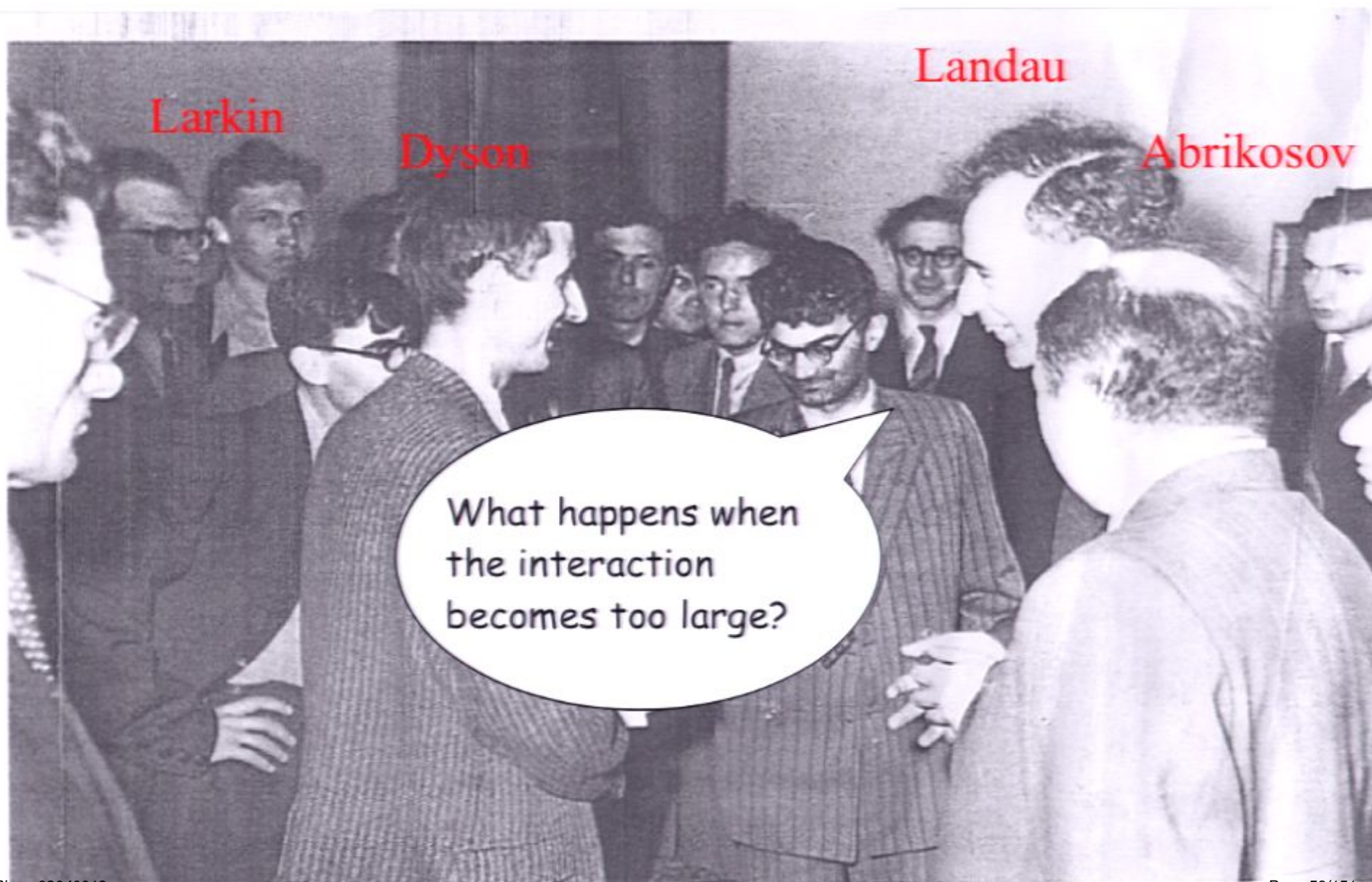
$$\frac{m^*}{m} = \frac{N(0)^*}{N(0)} = 1 + \frac{F_1^s}{3}$$

Page 51/151

Landau's Question.



Landau's Question.





What happens when the interaction becomes too large?



Landau 1936



"Electrons order"

What happens when the interaction becomes too large?



What happens when the interaction becomes too large?

Landau 1936

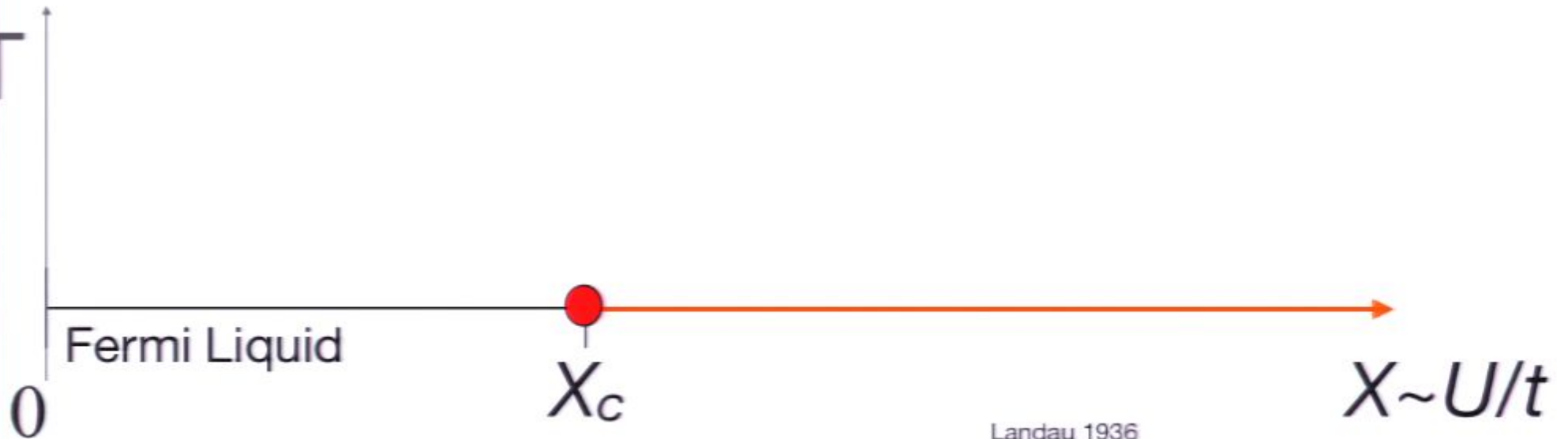


"Electrons order"

Mott 1947



"Electrons localize"



What happens when the interaction becomes too large?

Landau 1936



"Electrons order"

Mott 1947

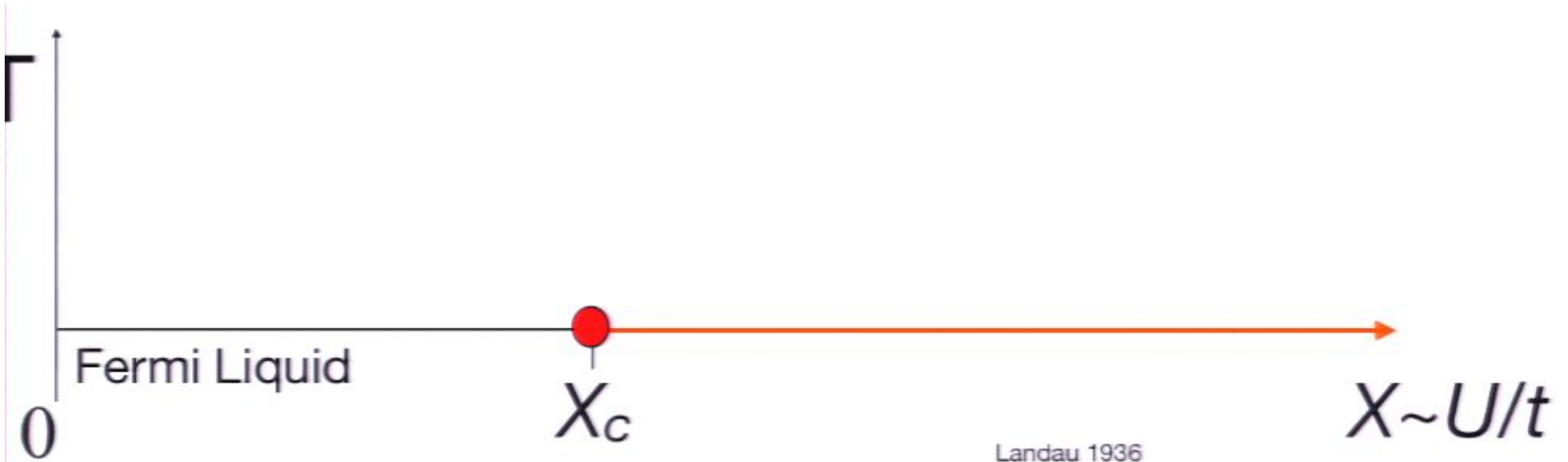


"Electrons localize"

Anderson 1961



"Local moments form"



What happens when the interaction becomes too large?

Landau 1936



"Electrons order"

Mott 1947



"Electrons localize"

Anderson 1961



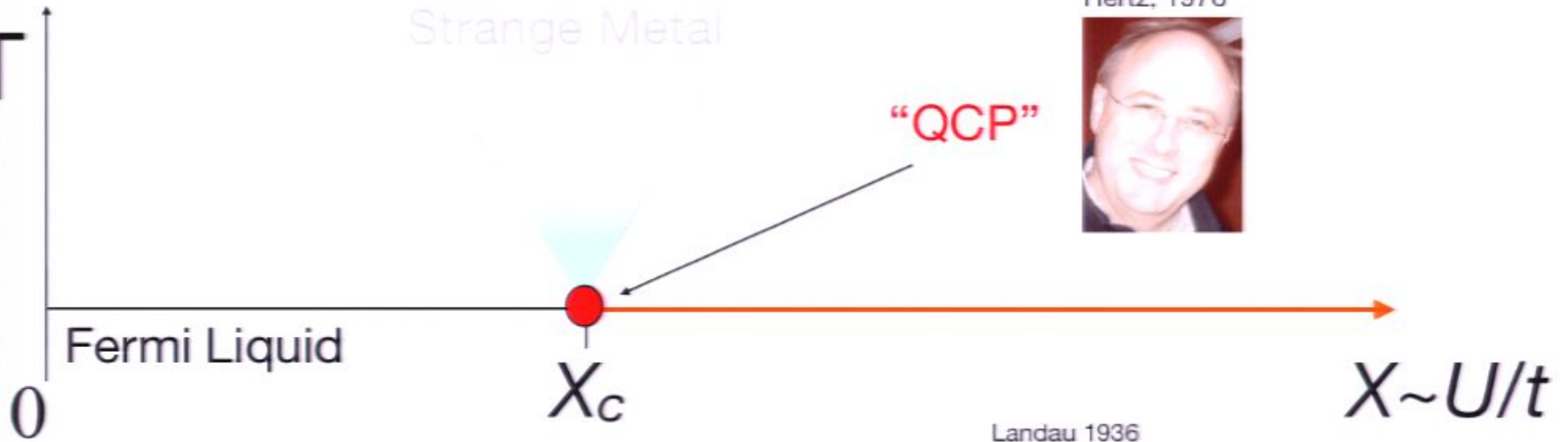
"Local moments form"

Strange Metal

Hertz, 1976



“QCP”



Landau 1936



“Electrons order”

Mott 1947



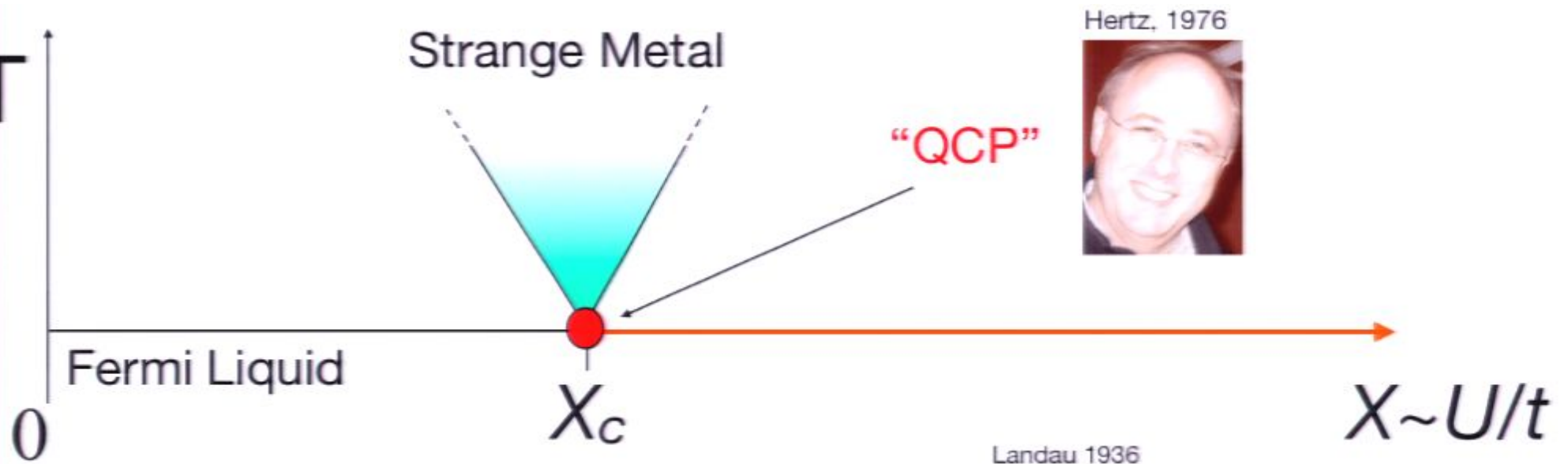
“Electrons localize”

Anderson 1961



“Local moments form”

What happens when the interaction becomes too large?



Hertz, 1976



Landau 1936



"Electrons order"

Mott 1947



"Electrons localize"

Anderson 1961



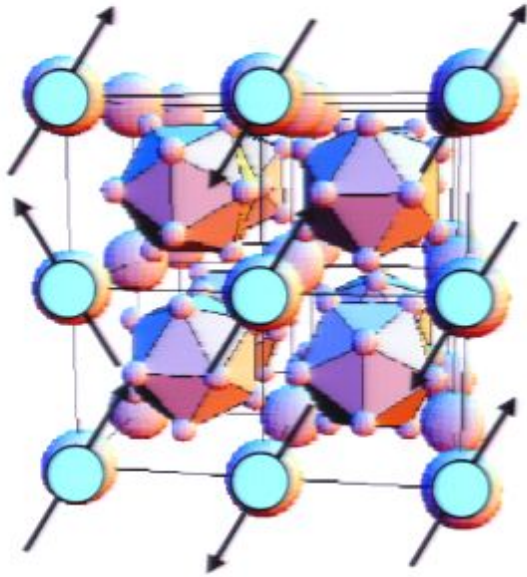
"Local moments form"

What happens when the interaction becomes too large?

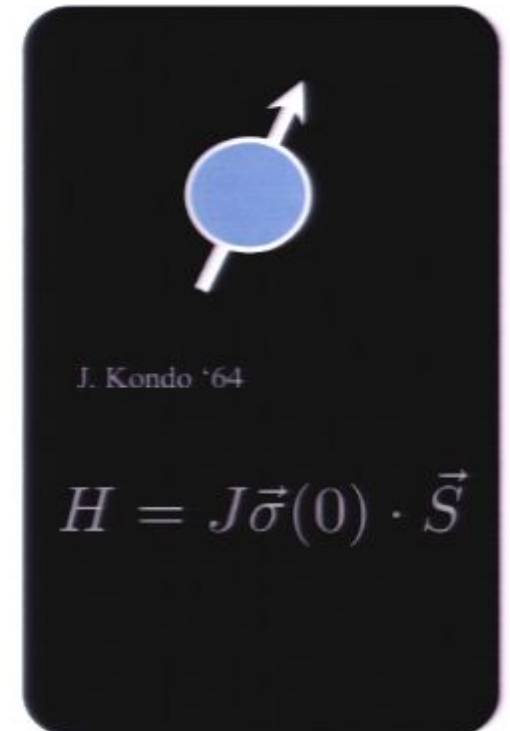
Heavy Electron Quantum Criticality.

Heavy Fermion Metals

[Review: cond-mat/0612006](#)



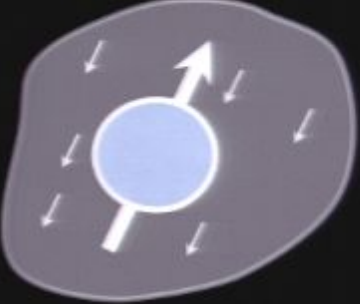
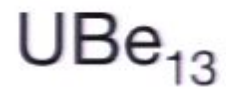
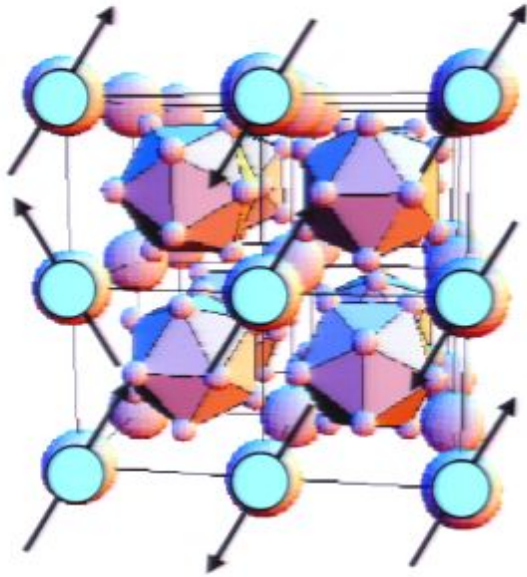
UBe₁₃



A diagram illustrating the Kondo effect. It features a blue circle representing a conduction electron with a white arrow pointing upwards and to the right, indicating its spin. Below the circle, the text "J. Kondo '64" is written. At the bottom, the Kondo Hamiltonian is given as $H = J\vec{\sigma}(0) \cdot \vec{S}$.

Heavy Fermion Metals

[Review: cond-mat/0612006](#)



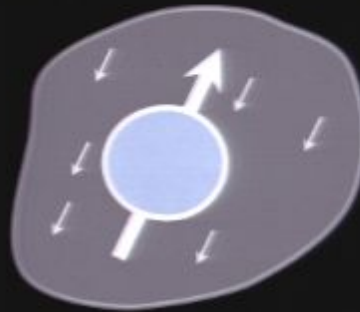
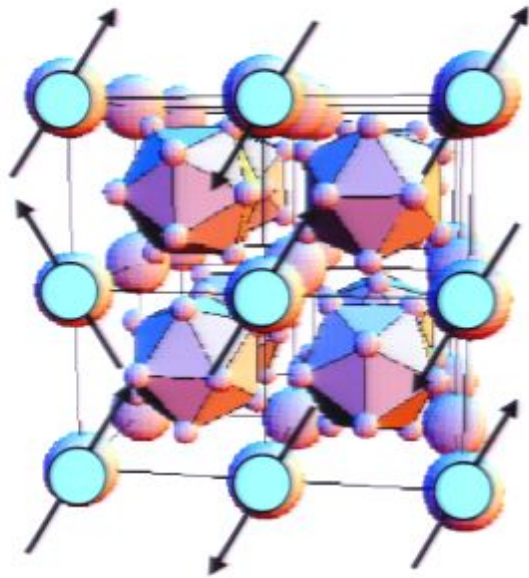
A diagram illustrating the Kondo effect. It shows a central blue circle representing a magnetic impurity with a large upward-pointing arrow indicating its spin. This impurity is surrounded by a grey, irregularly shaped region representing the conduction electron sea. Several smaller arrows within this region point towards the impurity, representing the screening of the impurity's spin by the conduction electrons.

J. Kondo '64

$$H = J\vec{\sigma}(0) \cdot \vec{S}$$

Heavy Fermion Metals

[Review: cond-mat/0612006](#)



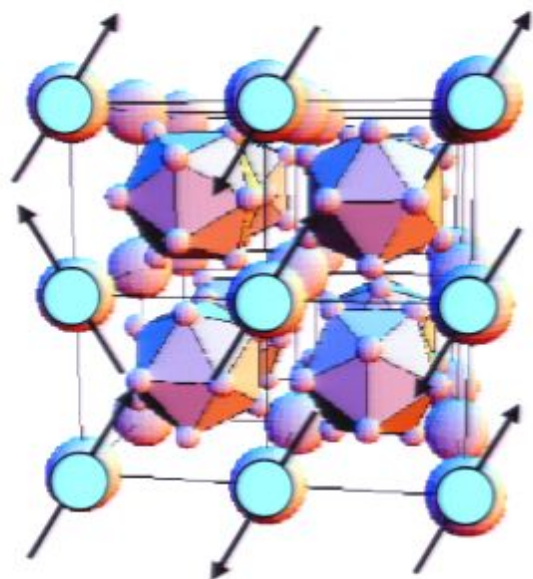
A diagram illustrating the Kondo effect. It shows a central blue circle representing a magnetic impurity with a white arrow pointing upwards, indicating its spin. This impurity is surrounded by a grey, irregularly shaped region representing the conduction electron sea. Several white arrows point towards the impurity from the surrounding region, representing the scattering of conduction electrons by the impurity's spin.

J. Kondo '64

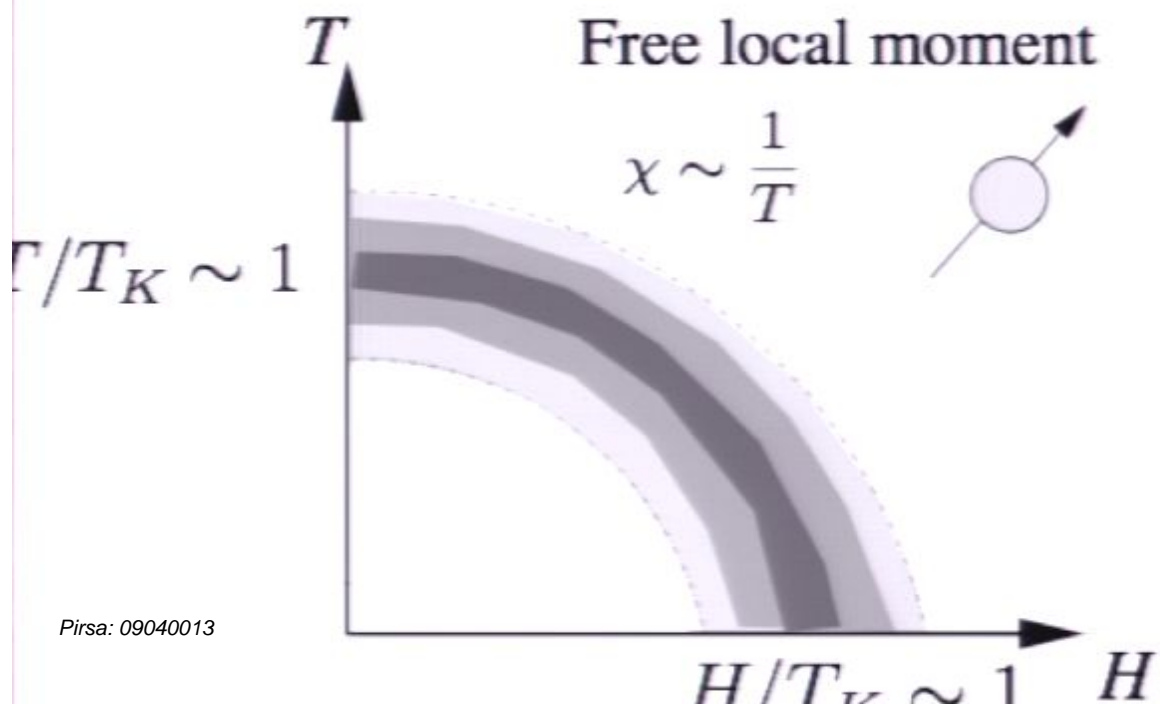
$$H = J\vec{\sigma}(0) \cdot \vec{S}$$
$$T_K \sim \epsilon_F e^{-\frac{\epsilon_F}{J}}$$

Heavy Fermion Metals

[Review: cond-mat/0612006](#)



UBe₁₃



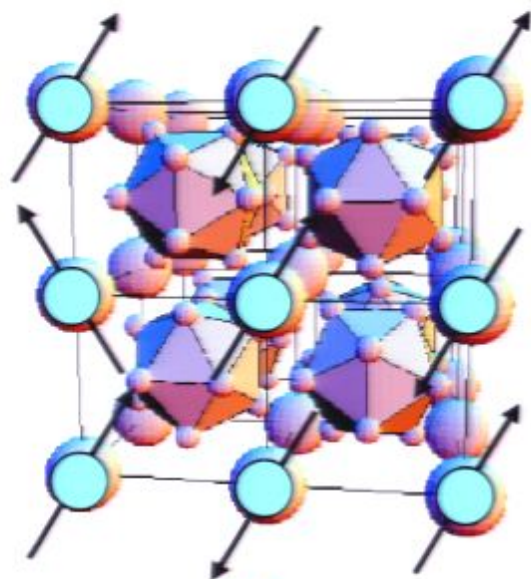
J. Kondo '64

$$H = J\vec{\sigma}(0) \cdot \vec{S}$$

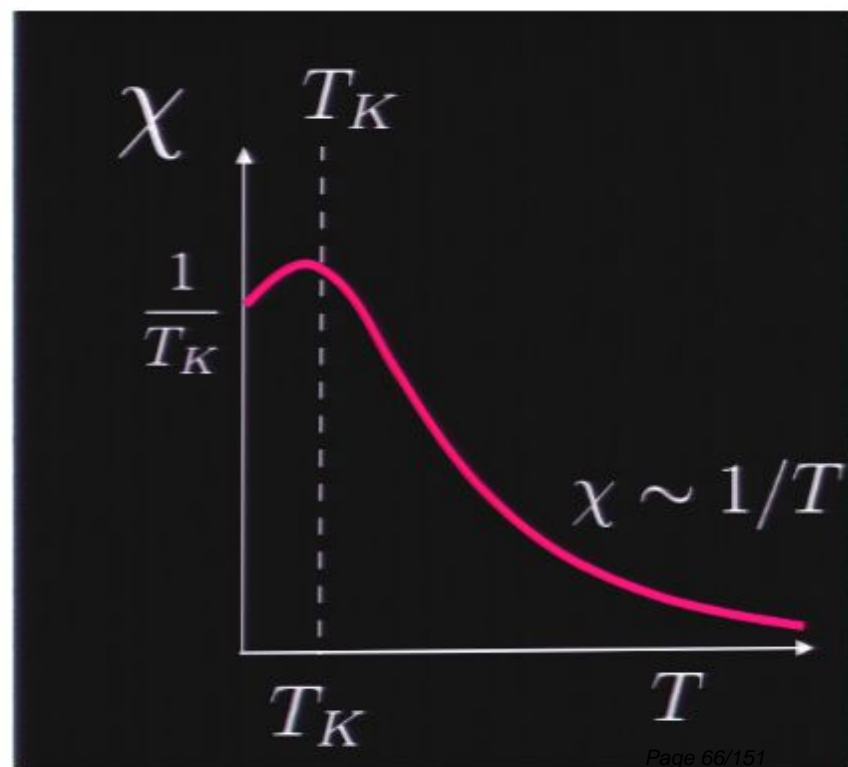
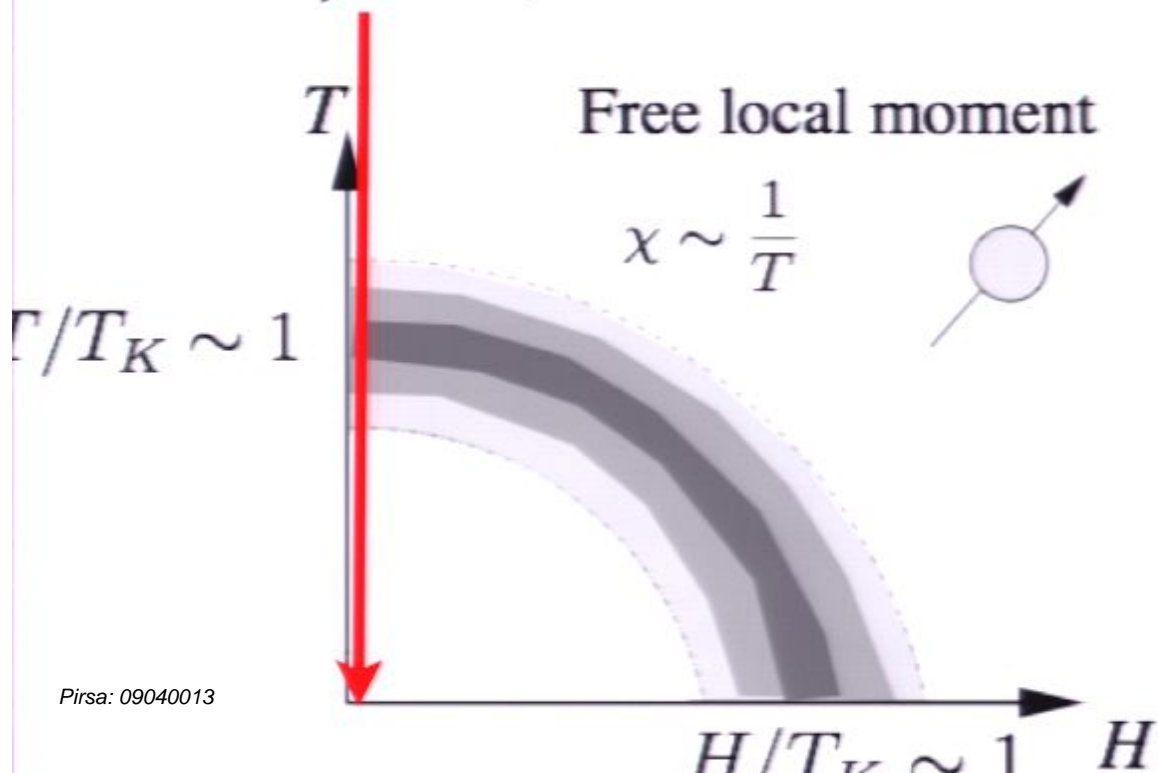
$$T_K \sim \epsilon_F e^{-\frac{\epsilon_F}{J}}$$

Heavy Fermion Metals

[Review: cond-mat/0612006](#)

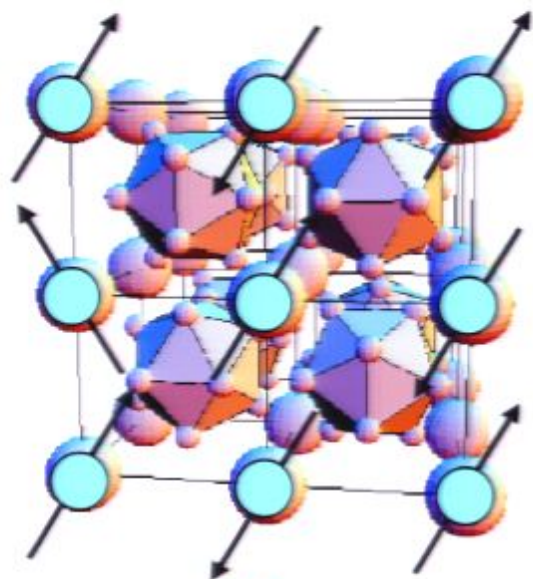


UBe_{13}

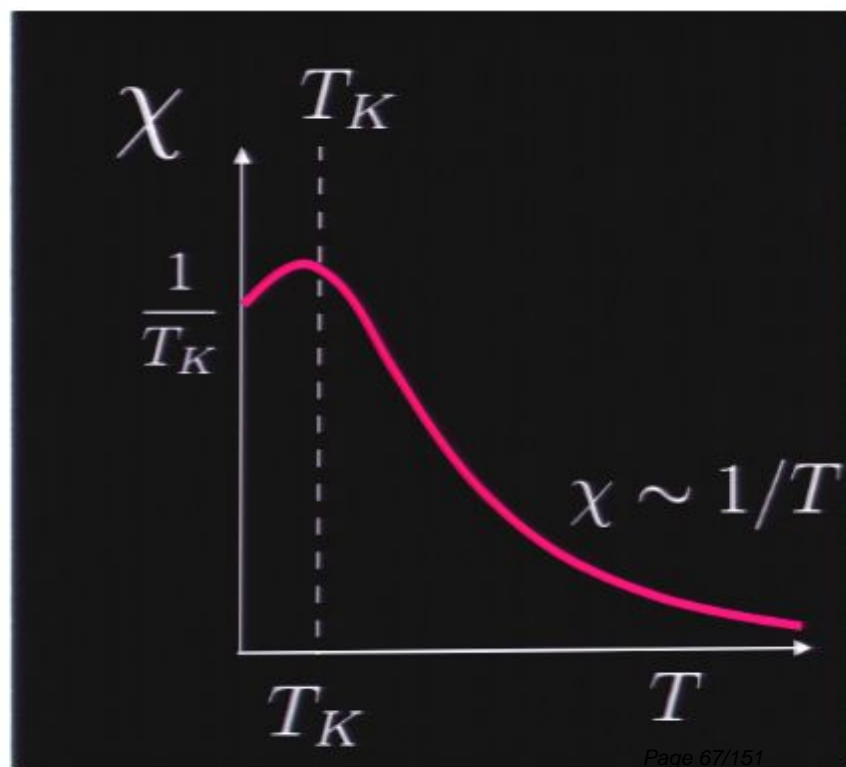
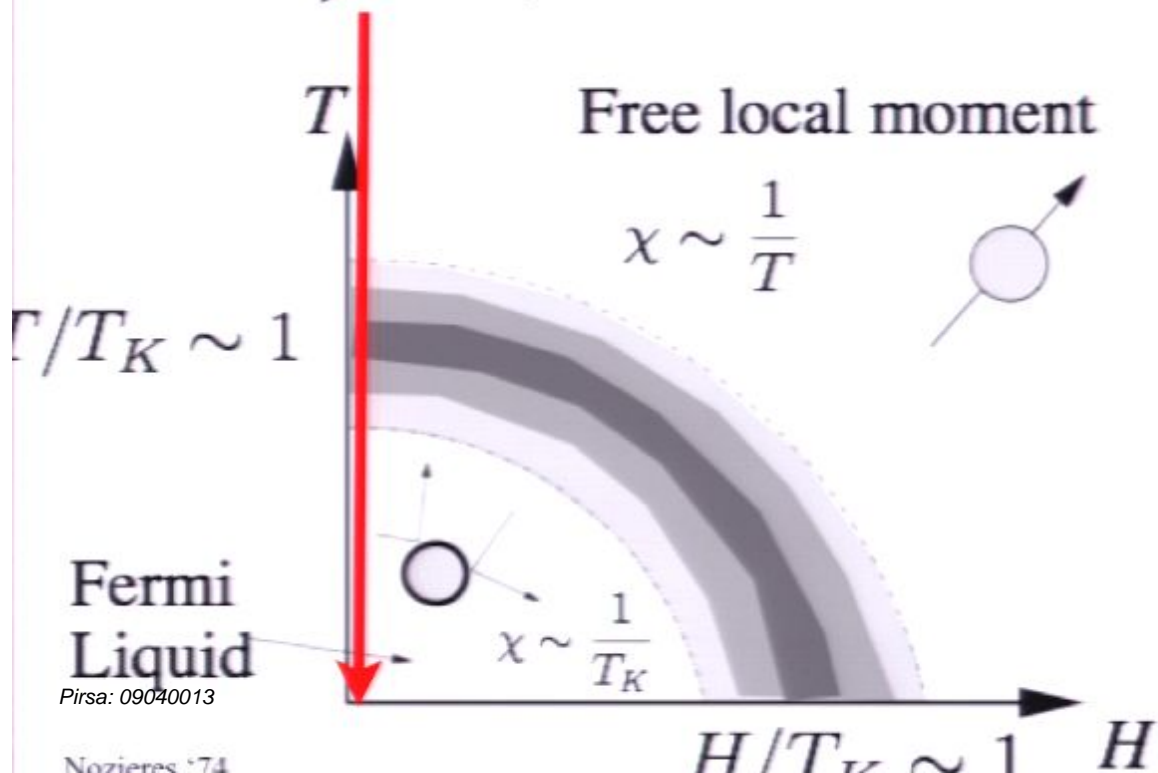


Heavy Fermion Metals

[Review: cond-mat/0612006](#)

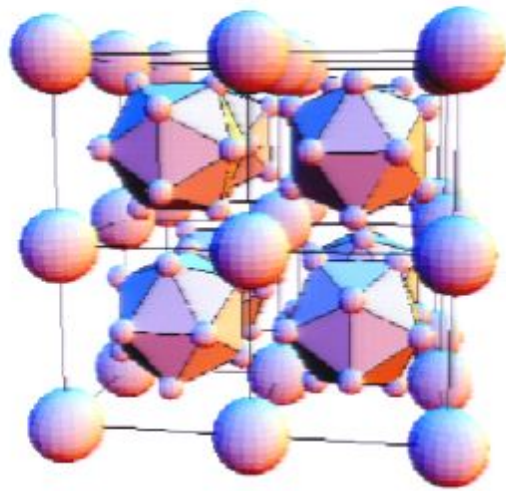


UBe₁₃

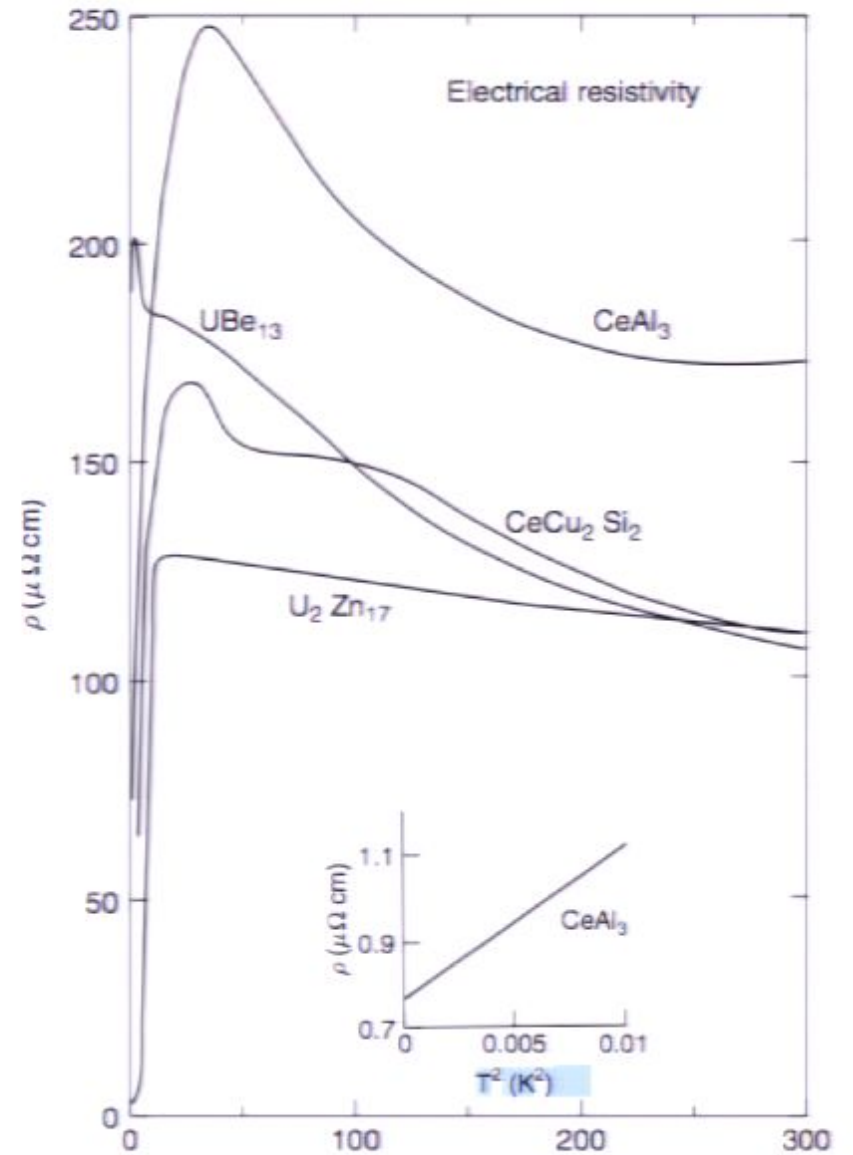
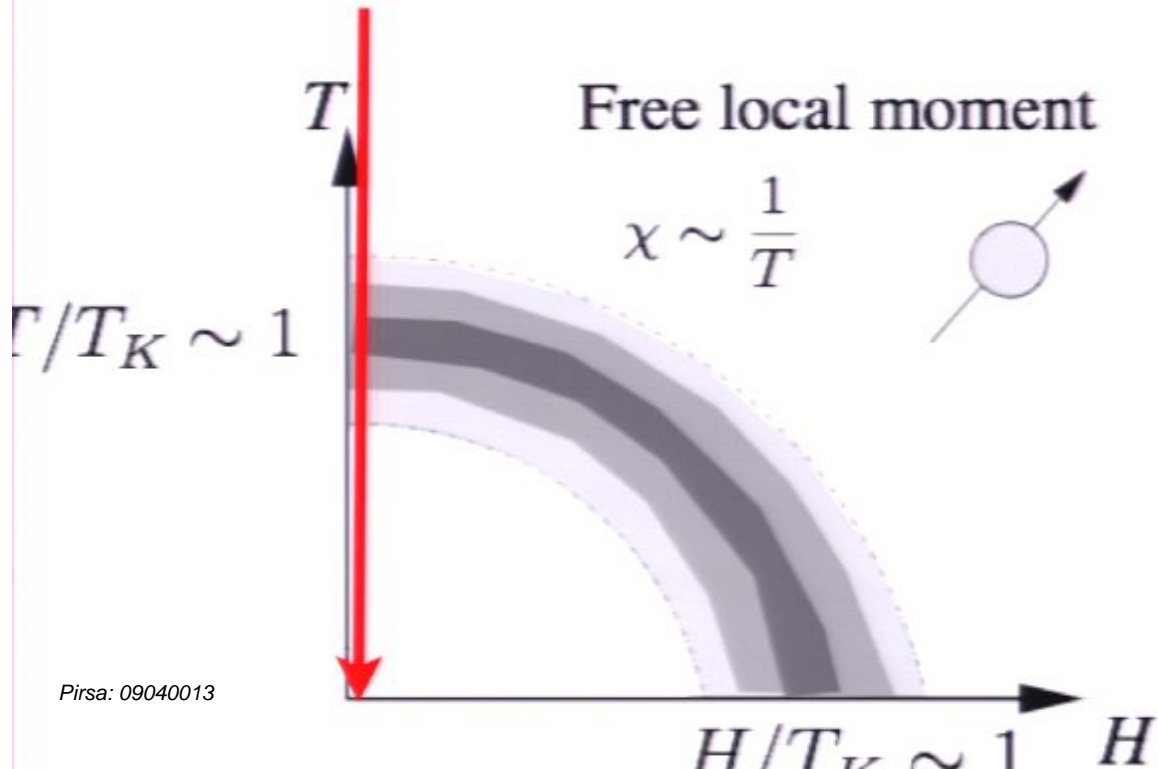


Heavy Fermion Metals

[Review: cond-mat/0612006](#)

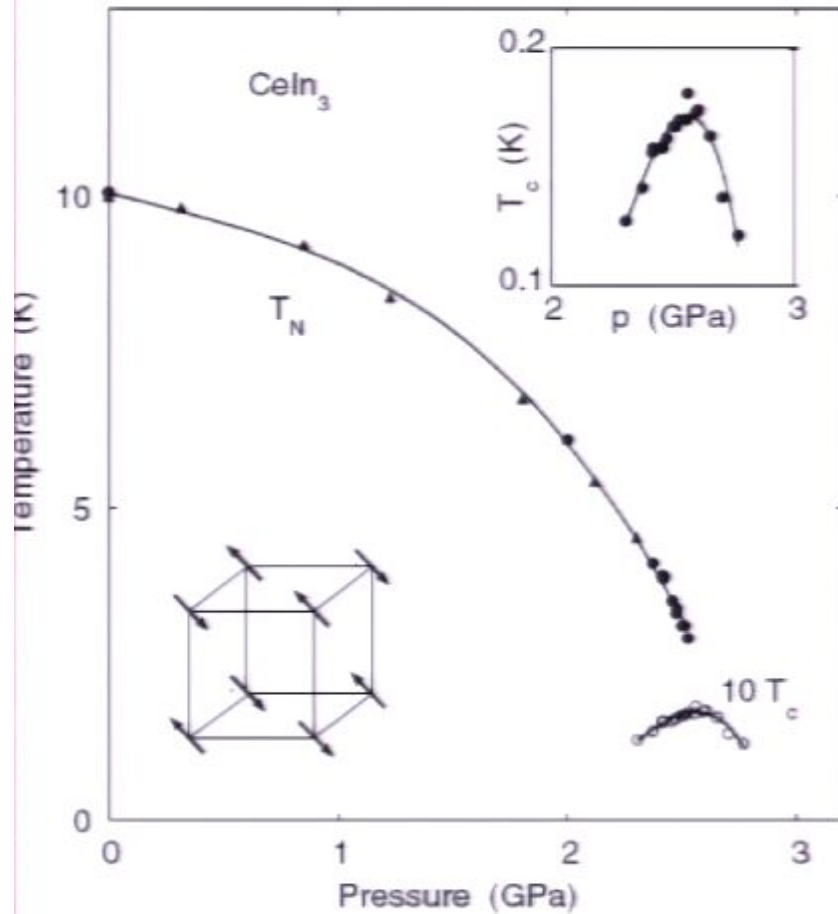


UBe₁₃



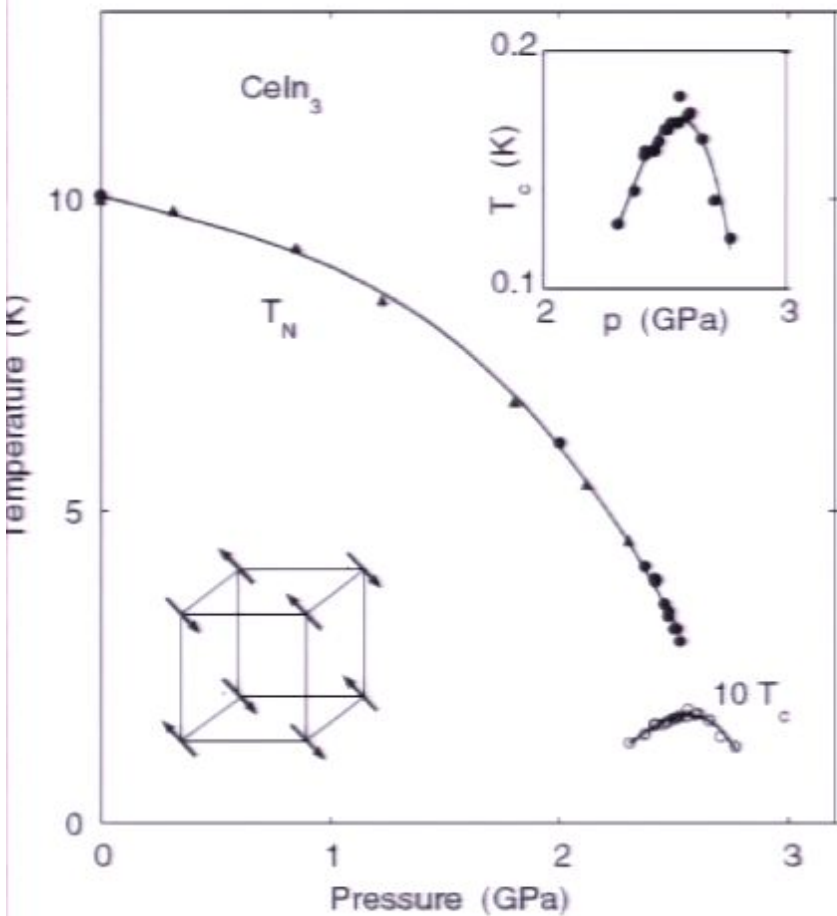
Experiments

“Avoided Criticality”

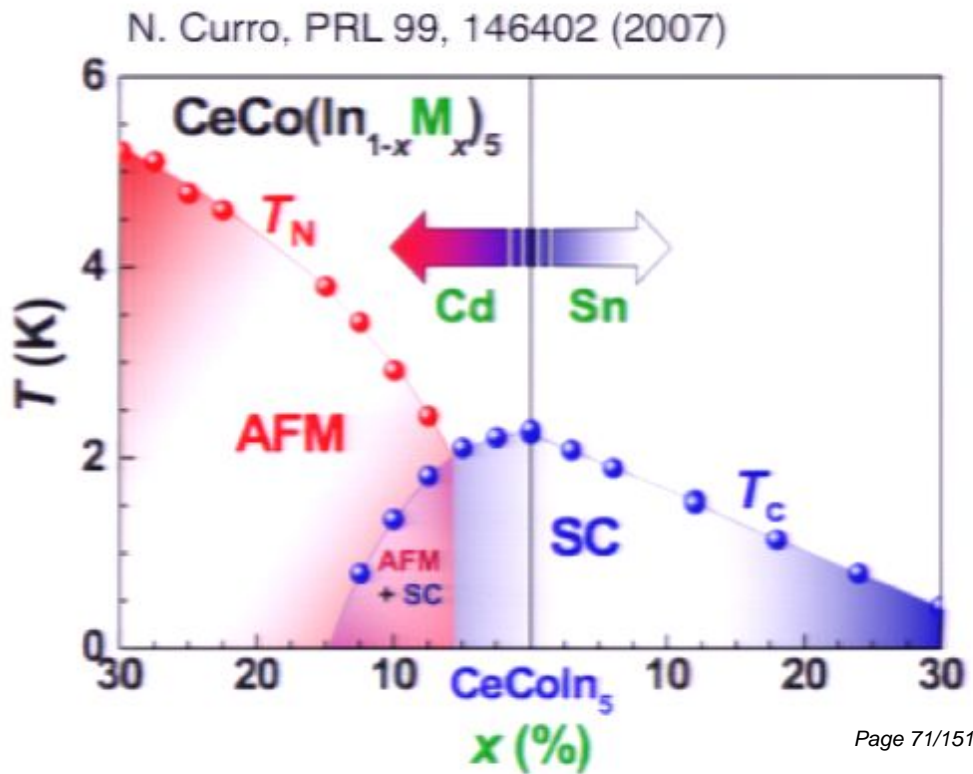


Mathur et al, Nature 394, 39 (1998)

“Avoided Criticality”

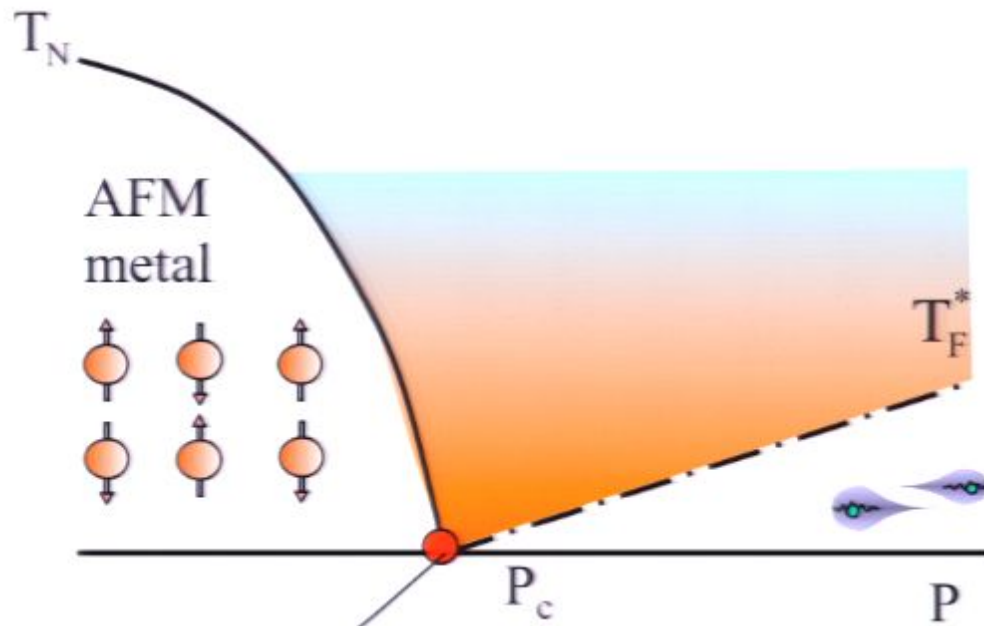


Mathur et al, Nature 394, 39 (1998)



Quantum Criticality: divergent specific heat capacity

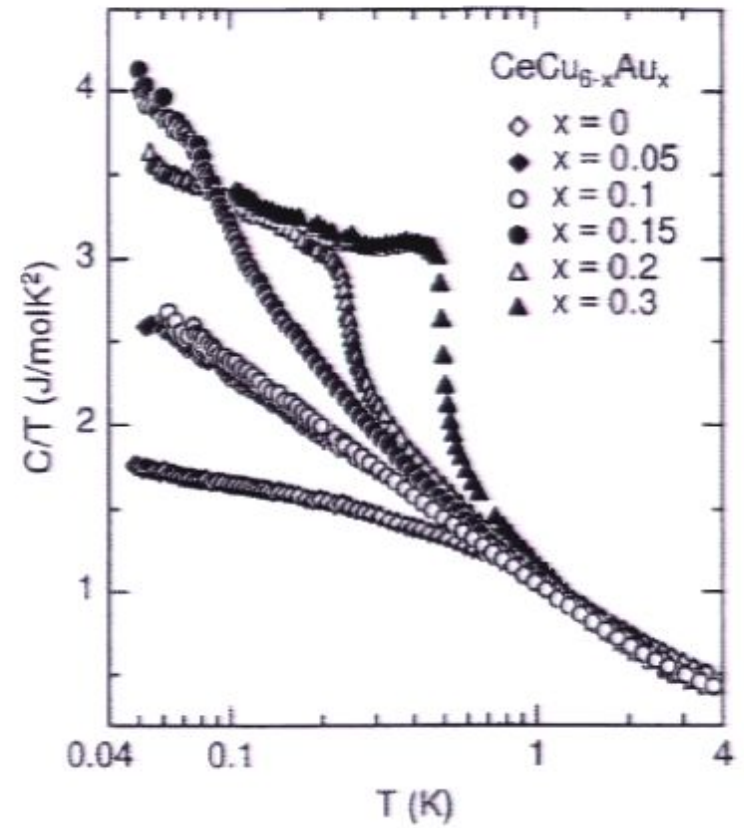
Heavy Fermion
Materials



Quantum Critical
point

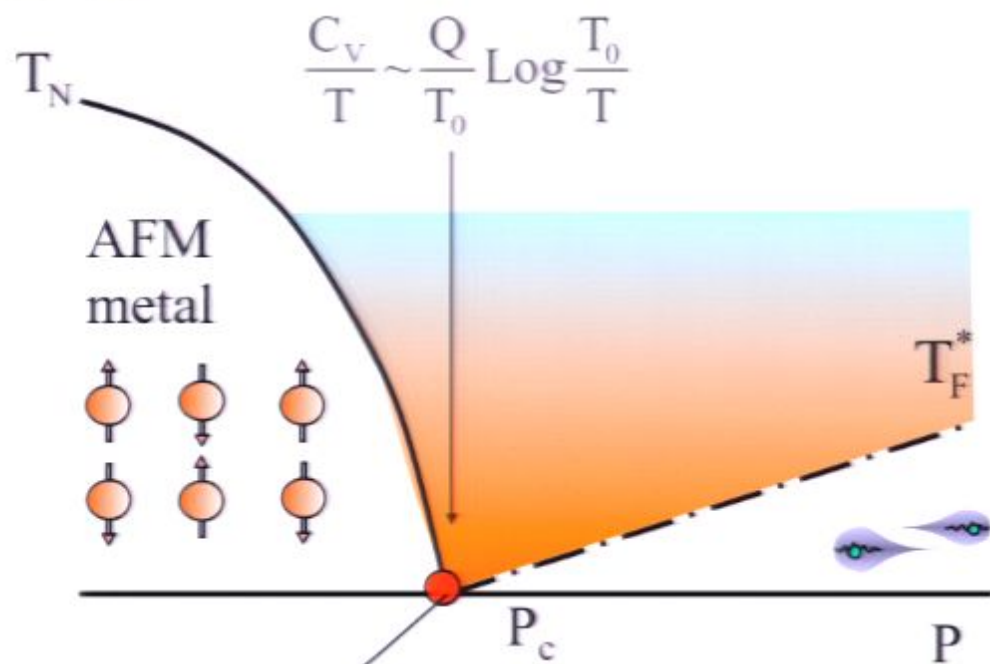
Pirsa: 09040013

H. Von Lohneyson (1996)



Quantum Criticality: divergent specific heat capacity

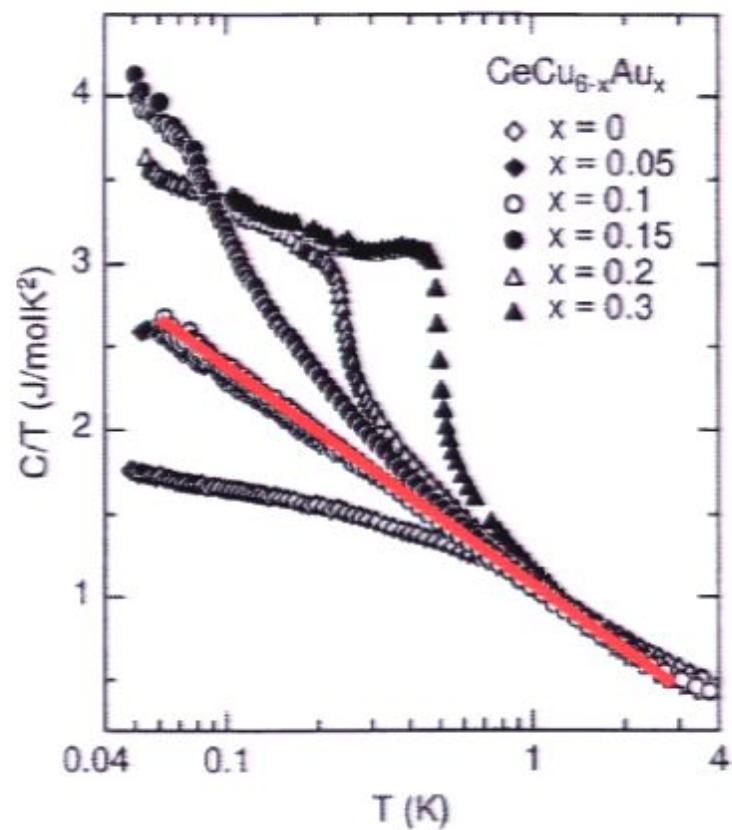
Heavy Fermion
Materials



Quantum Critical
point

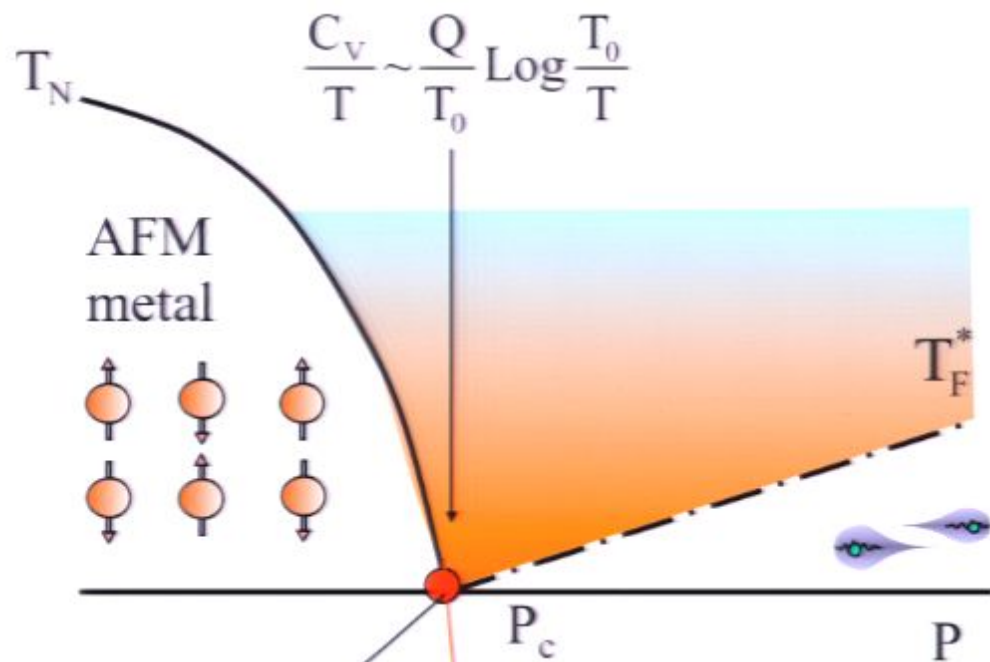
Pirsa: 09040013

H. Von Lohneyson (1996)



Quantum Criticality: divergent specific heat capacity

Heavy Fermion
Materials

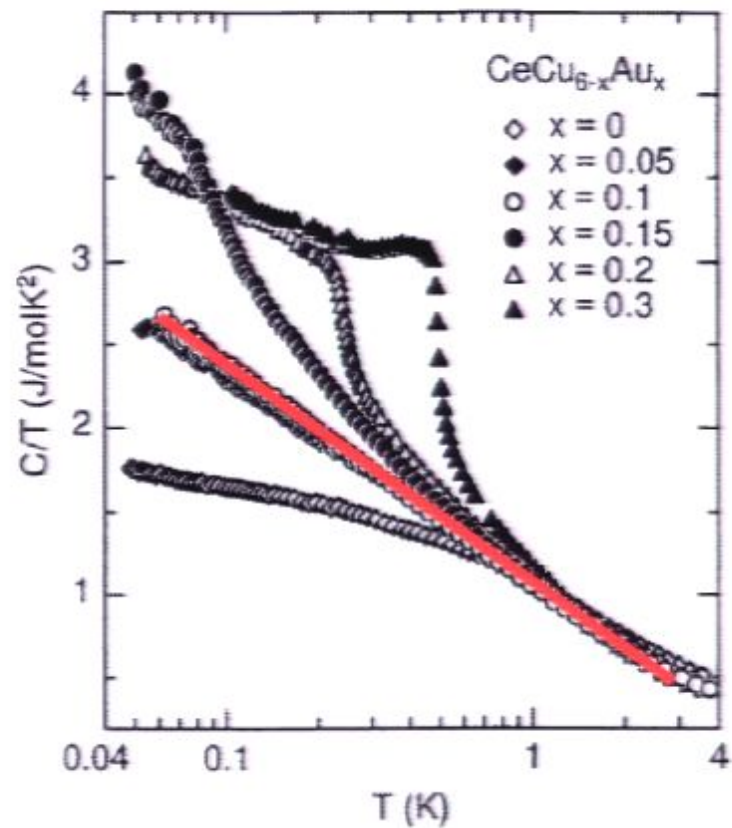


Quantum Critical
Point

$$\frac{m^*}{m} \rightarrow \infty$$

Pirsa: 09040013

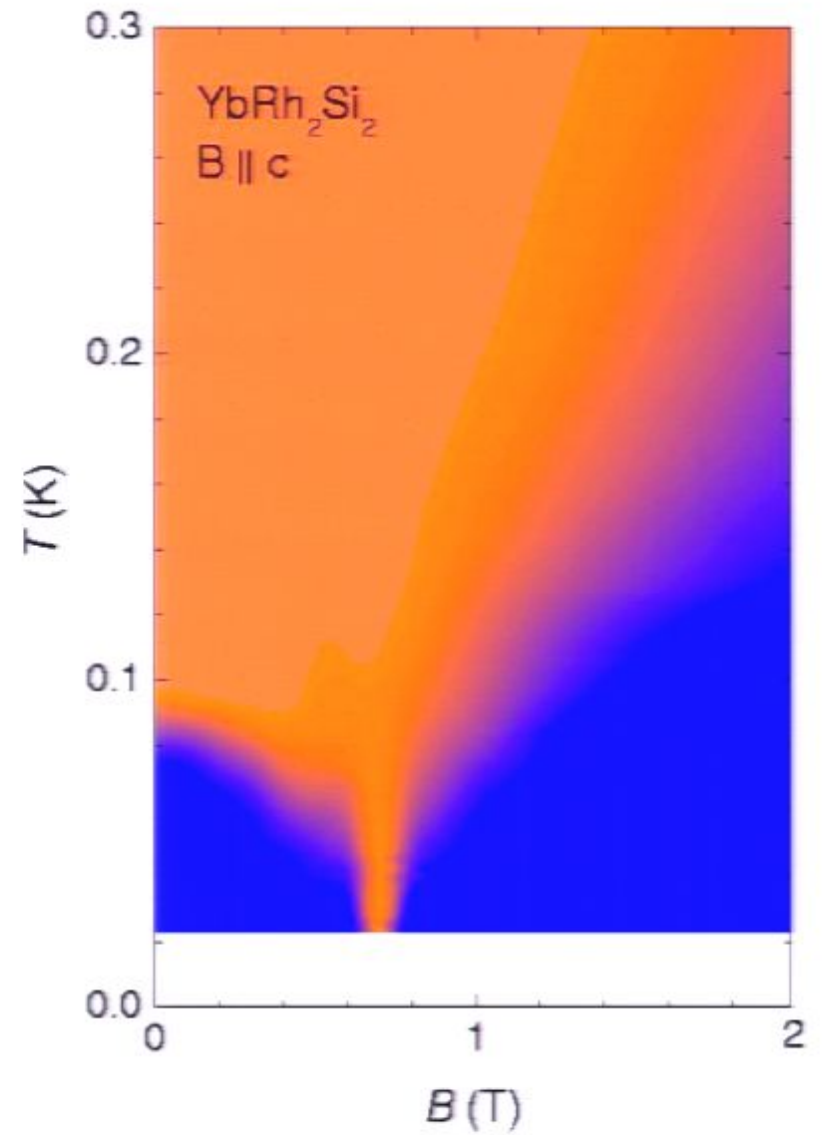
H. Von Lohneyson (1996)



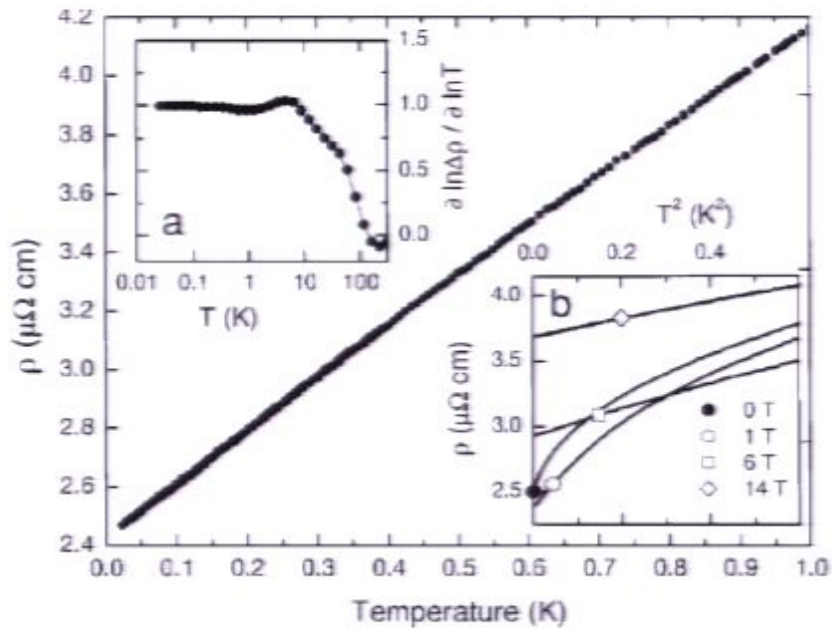
Field Tuned Criticality in YbRh₂Si₂.

YbRh₂Si₂

Trovarelli et al (2000).

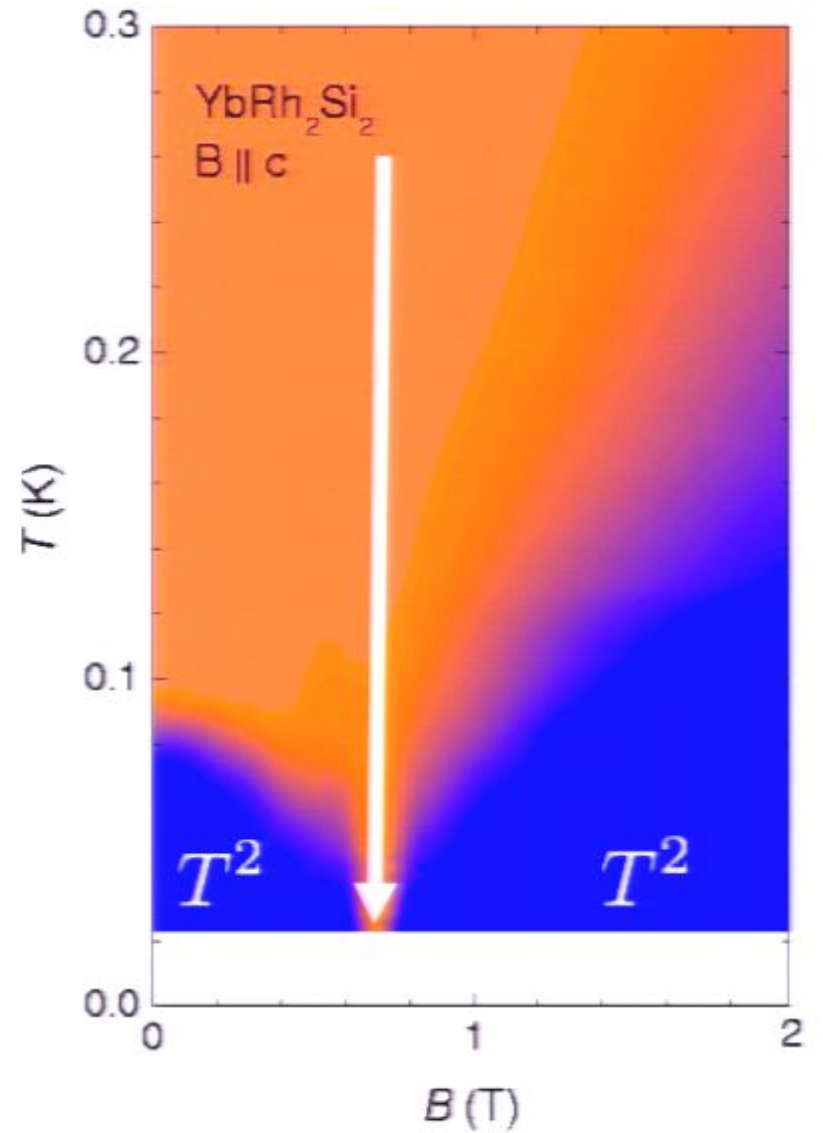


Field Tuned Criticality in YbRh₂Si₂.

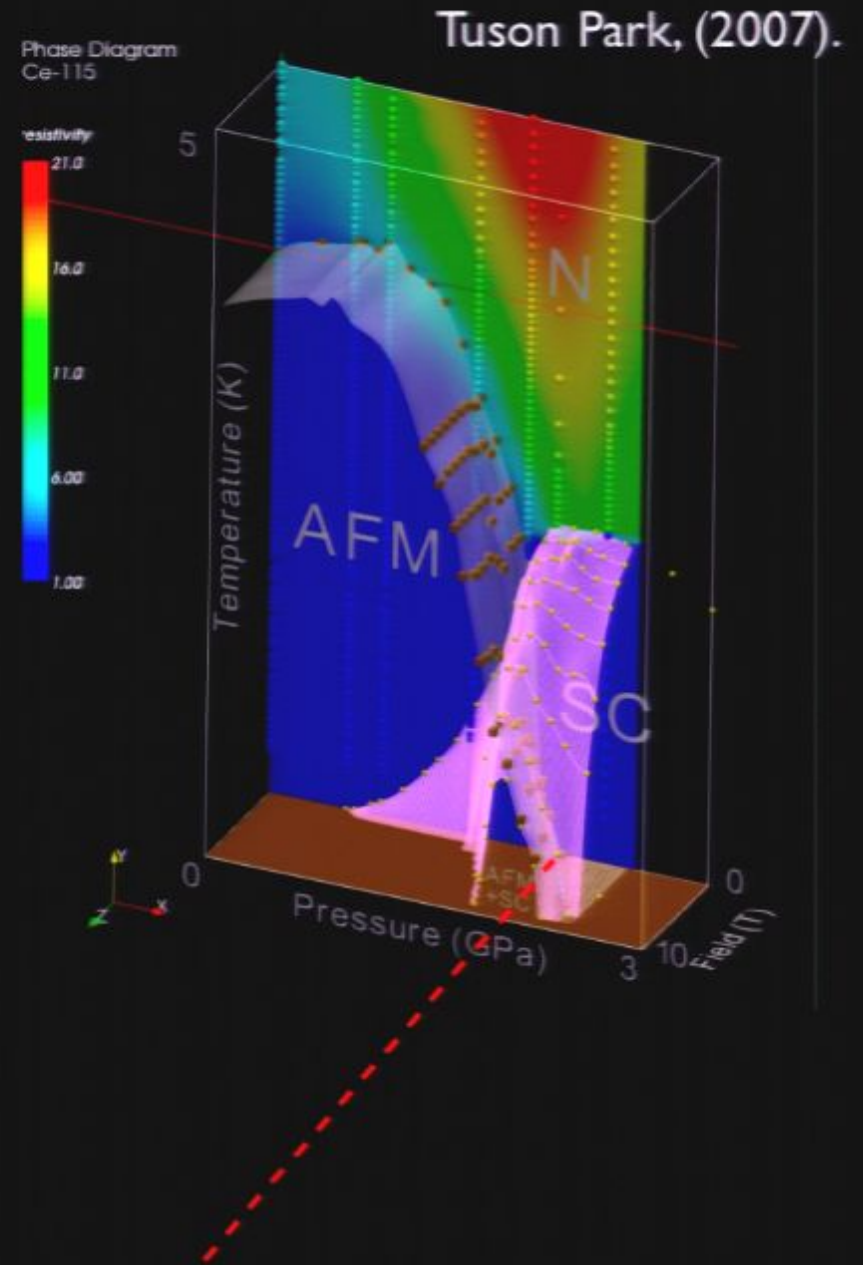


YbRh₂Si₂

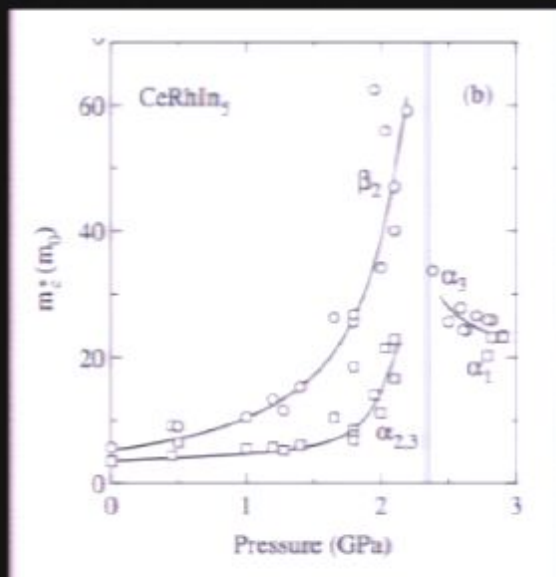
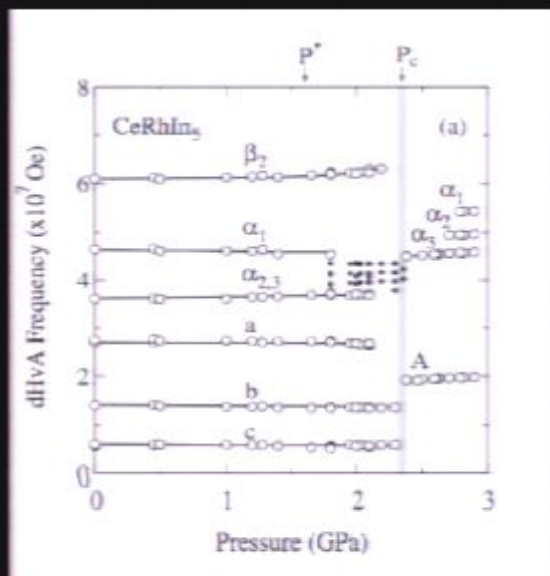
Trovarelli et al (2000).



Reconstruction of the Fermi Surface and mass divergence

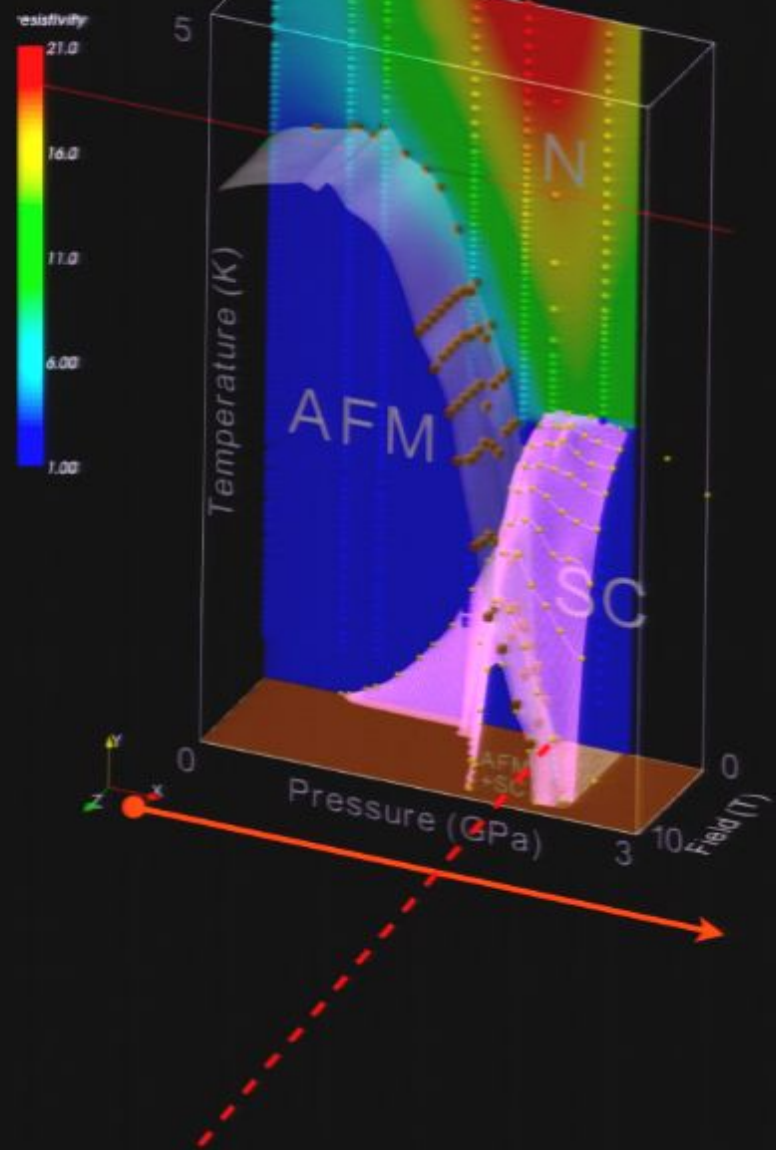


Reconstruction of the Fermi Surface and mass divergence



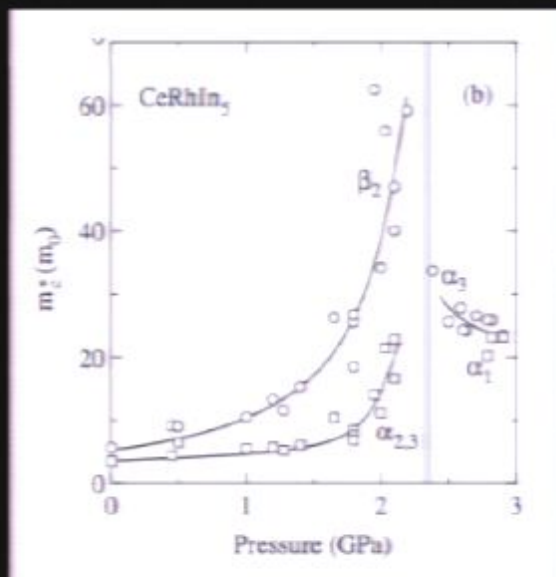
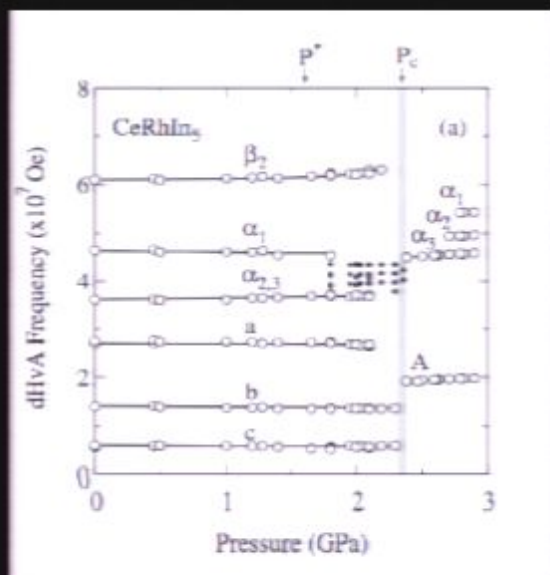
Phase Diagram
Ce-115

Tuson Park, (2007).



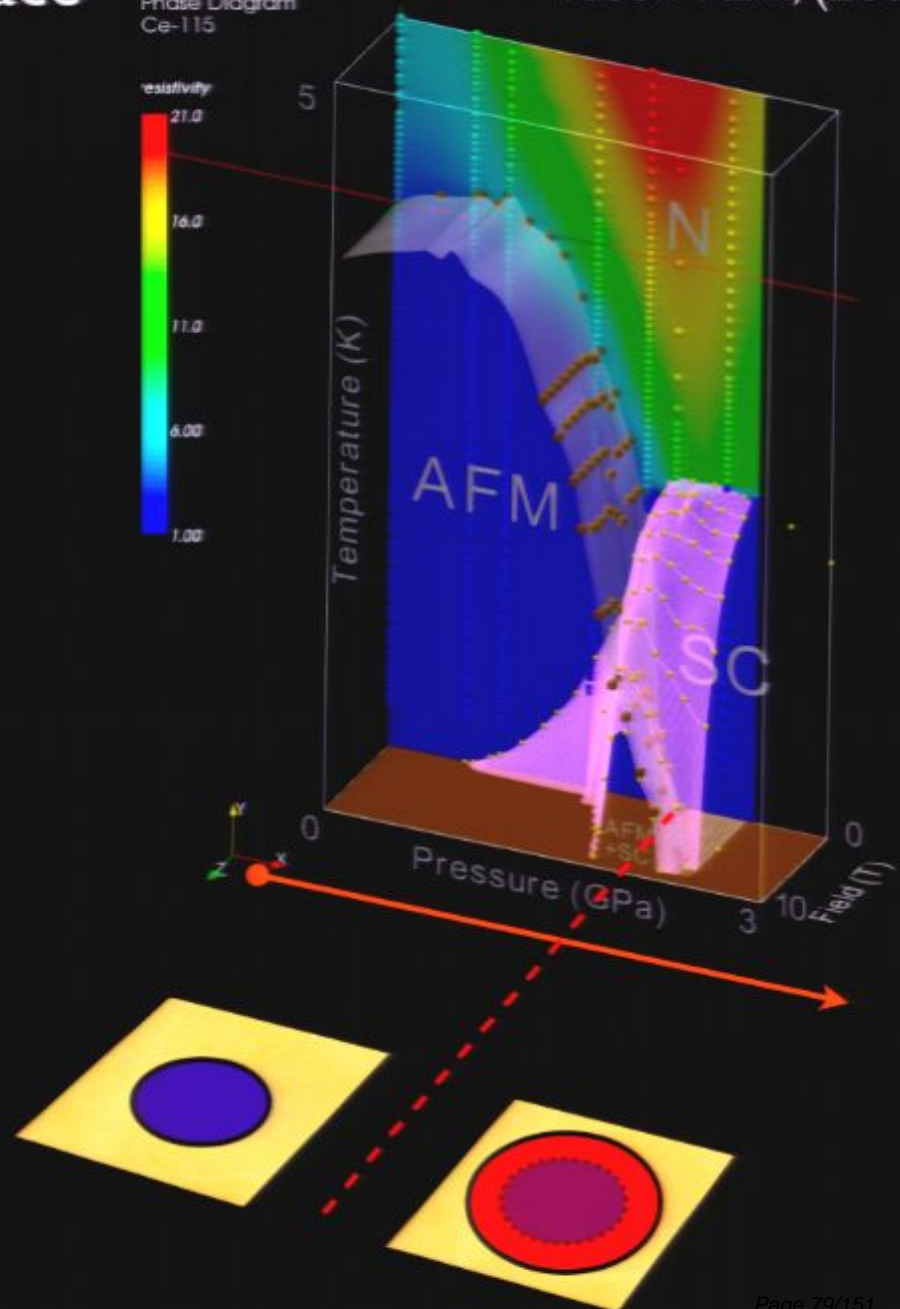
Reconstruction of the Fermi Surface and mass divergence

Tuson Park, (2007).

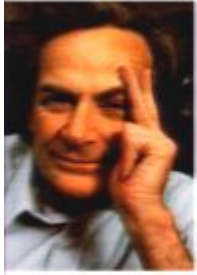


Phase Diagram
Ce-115

resistivity



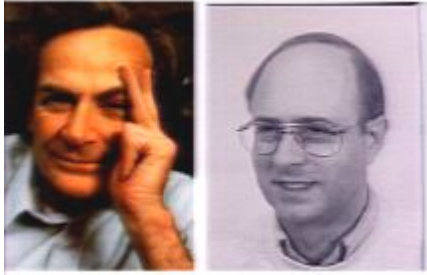
Black Hole in the Phase Diagram.



Feynman


Hertz

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

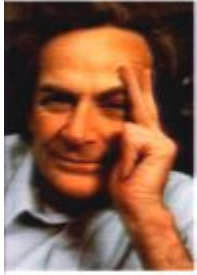


Feynman Hertz

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

$$\frac{\hbar}{k_B T}$$






Feynman

Hertz

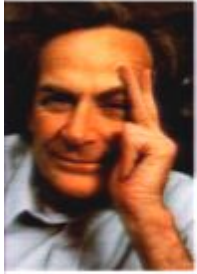
Temperature: "Casimir effect" in time.

(PC, L. Palova, P. Chandra (08))

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

$$\frac{\hbar}{k_B T}$$





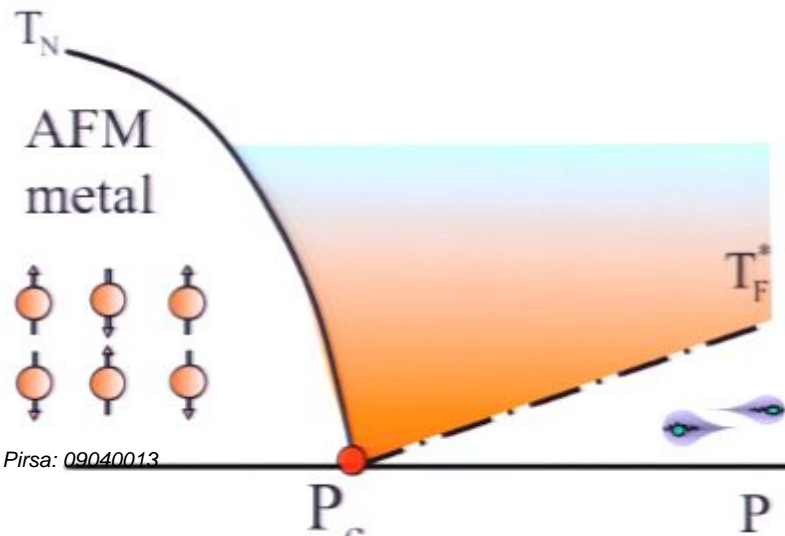
Feynman

Hertz

Temperature: "Casimir effect" in time.

(PC, L. Palova, P. Chandra (08))

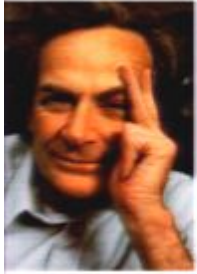
$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$



Pirsa: 09040013



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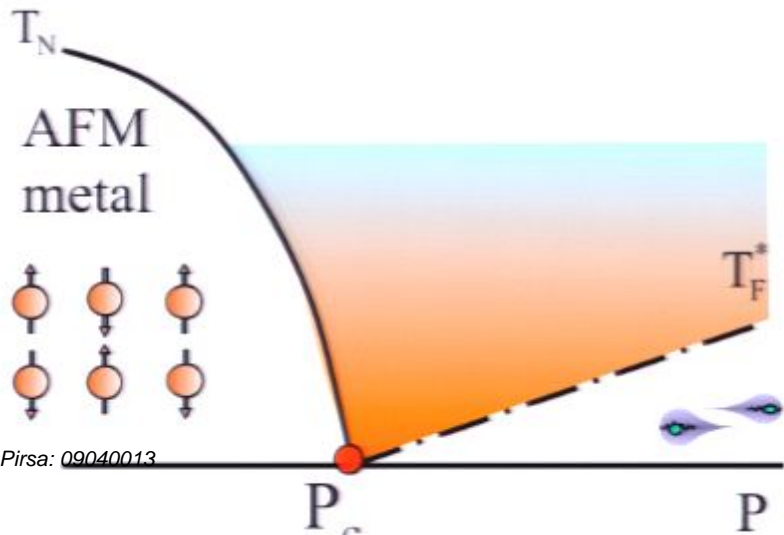
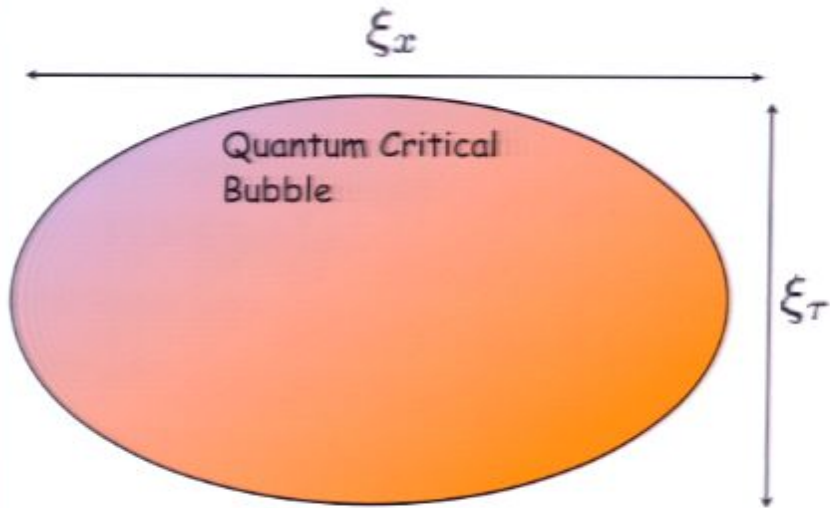
Feynman

Hertz

Temperature: "Casimir effect" in time.

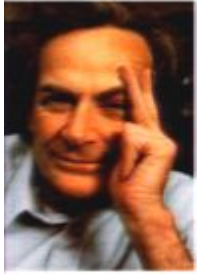
(PC, L. Palova, P. Chandra (08))

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$



Pirsa: 09040013



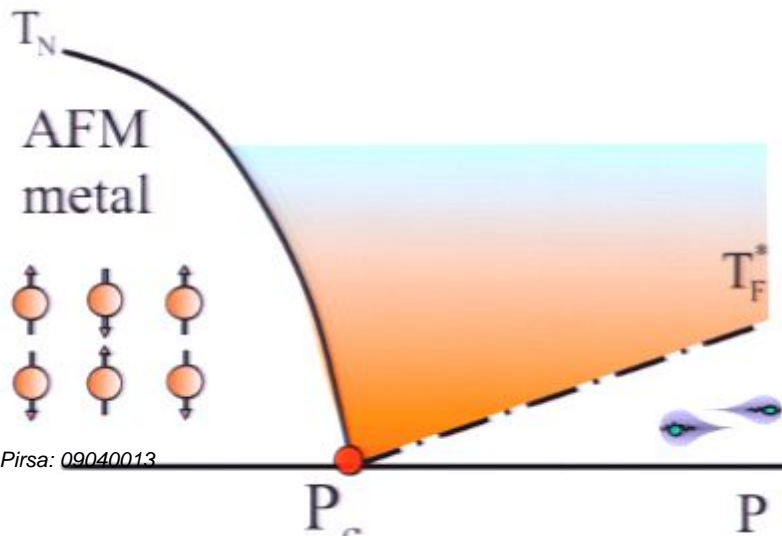
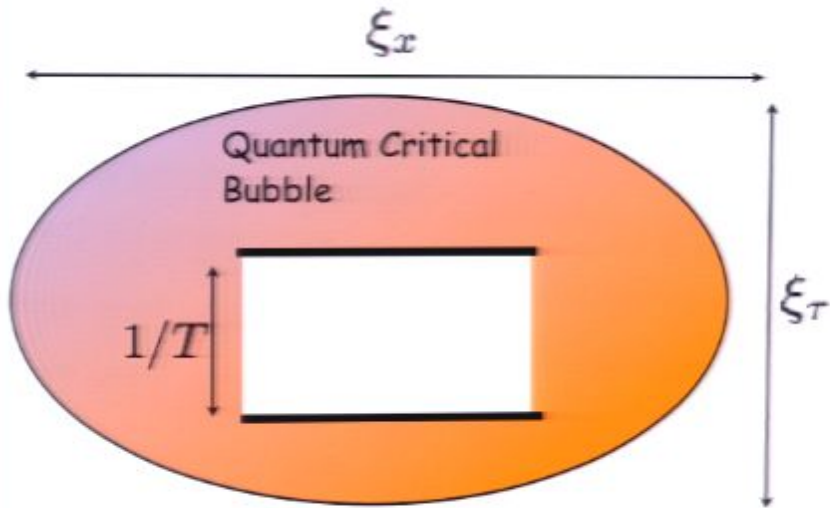


Feynman Hertz

Temperature: "Casimir effect" in time.

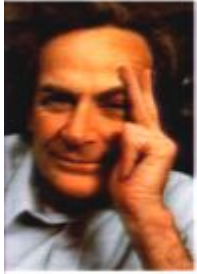
(PC, L. Palova, P. Chandra (08))

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$



Pirsa: 09040013



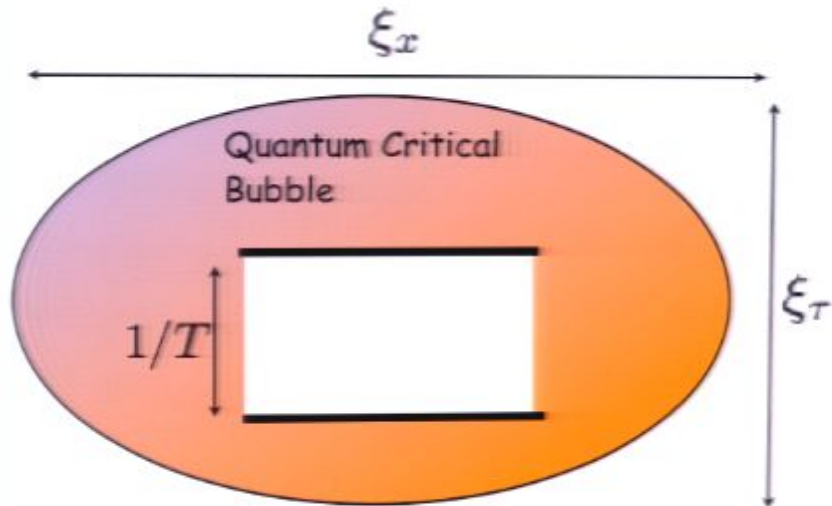


Feynman Hertz

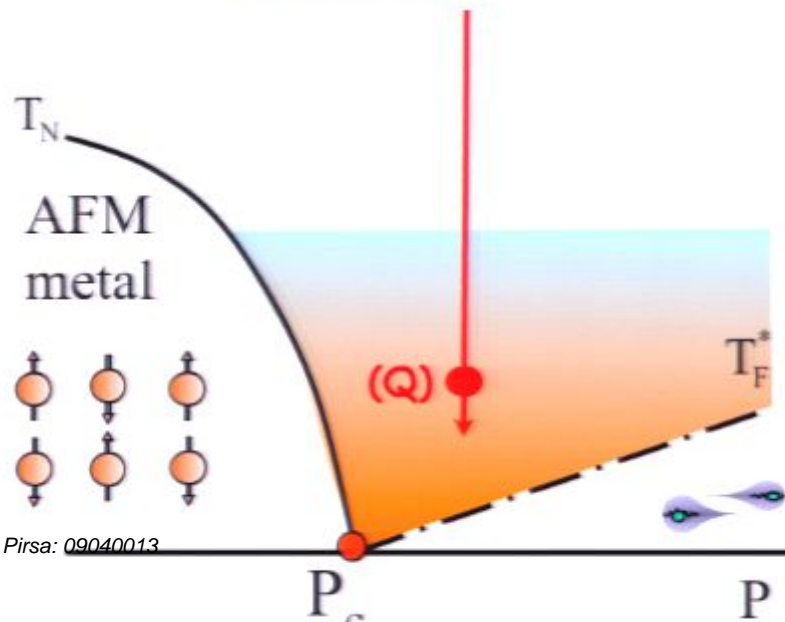
Temperature: "Casimir effect" in time.

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$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

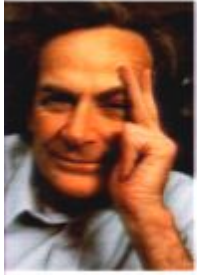


(Q) Quantum critical region:
interior of correlation bubble.



Pirsa: 09040013



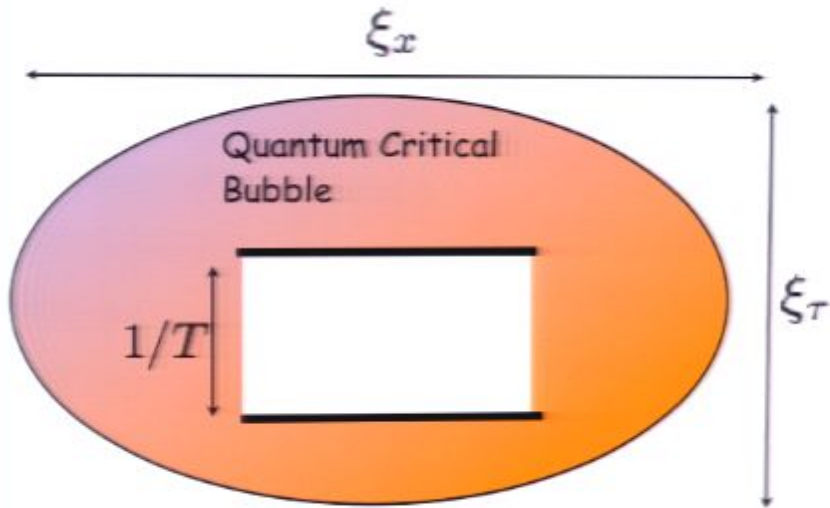


Feynman Hertz

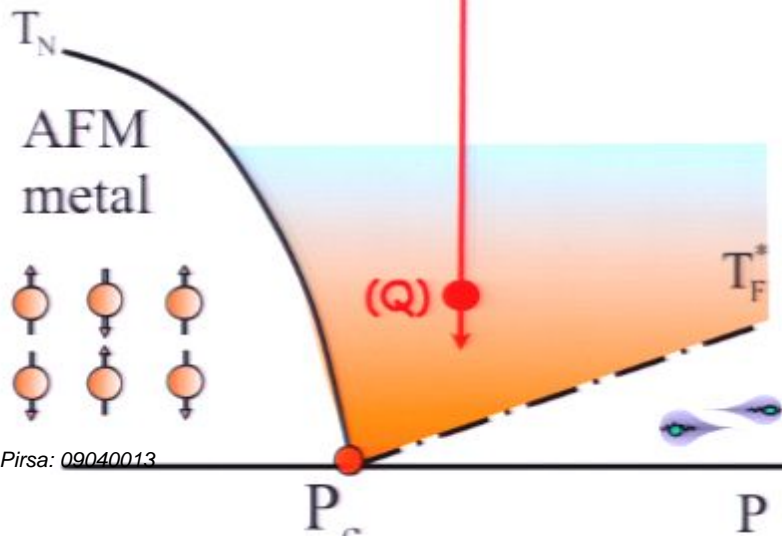
Temperature: "Casimir effect" in time.

(PC, L. Palova, P. Chandra (08))

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

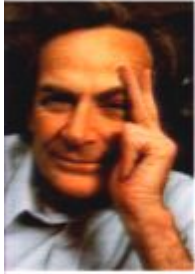


(Q) Quantum critical region:
interior of correlation bubble.



Pirsa: 09040013



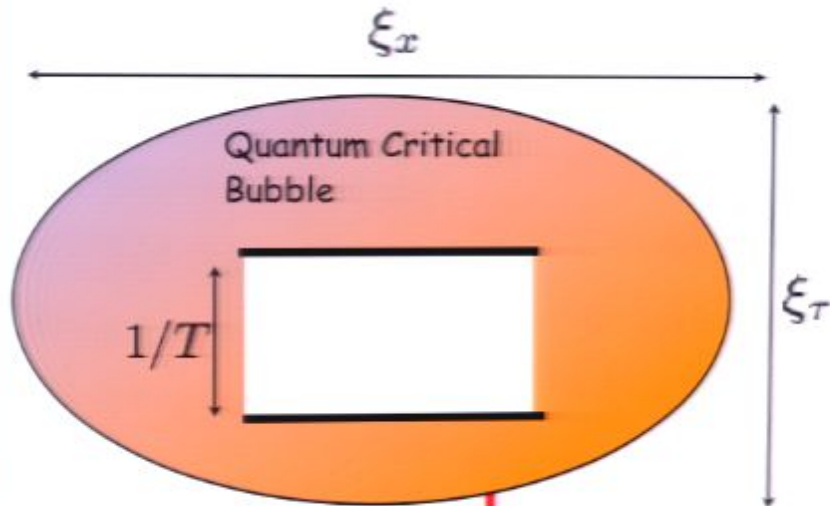


Feynman Hertz

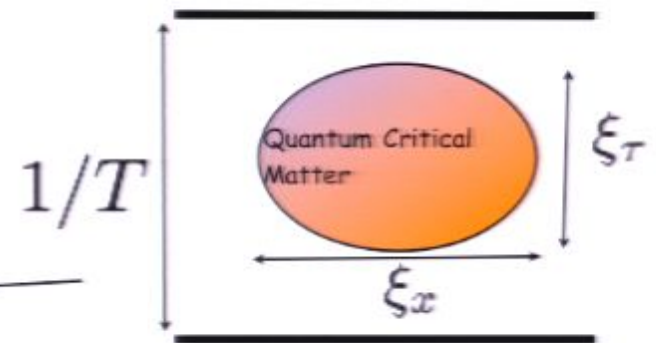
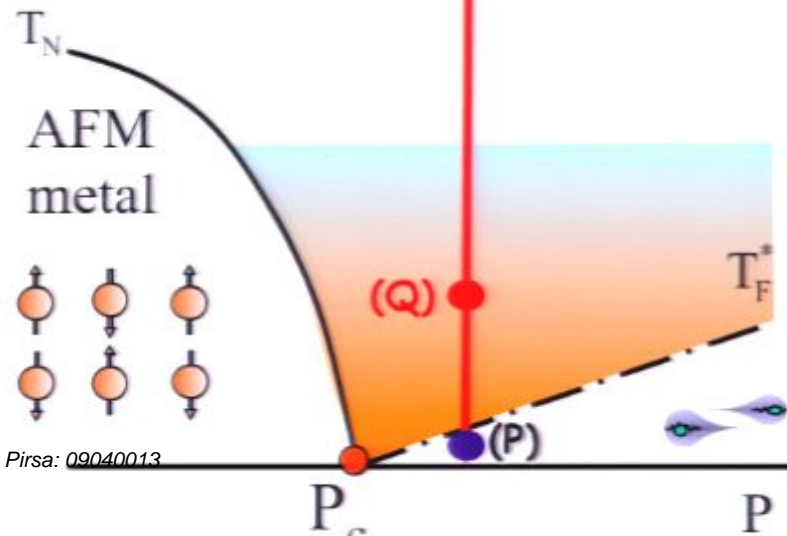
Temperature: "Casimir effect" in time.

(PC, L. Palova, P. Chandra (08))

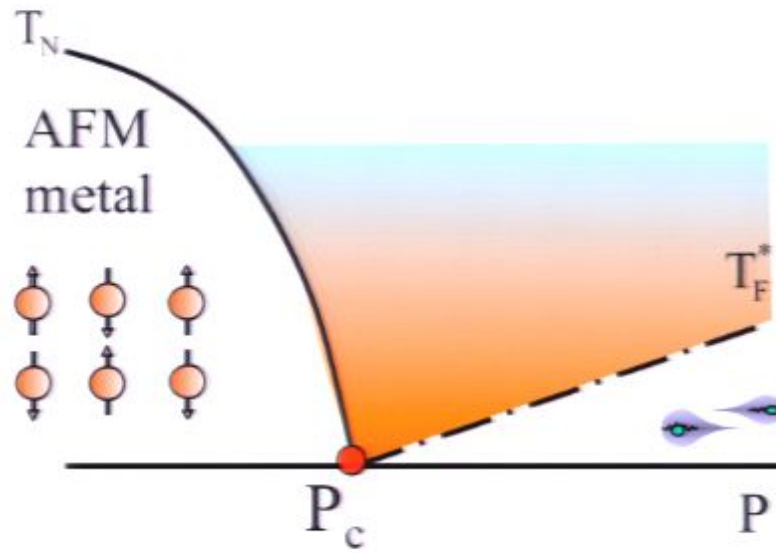
$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$



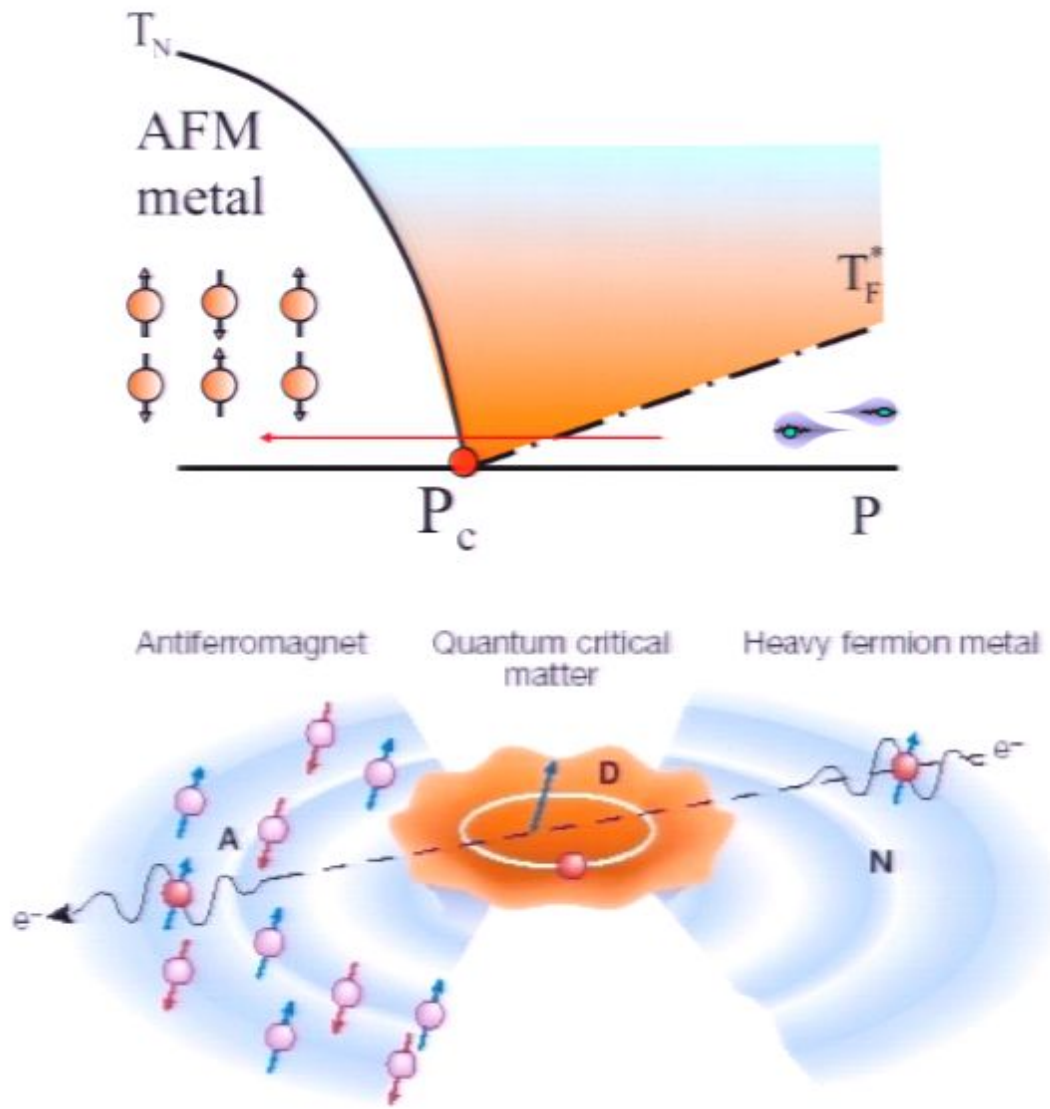
- (Q) Quantum critical region: interior of correlation bubble.
- (P) Paramagnet: probes exterior of correlation bubble



"Black Hole in the Phase Diagram".

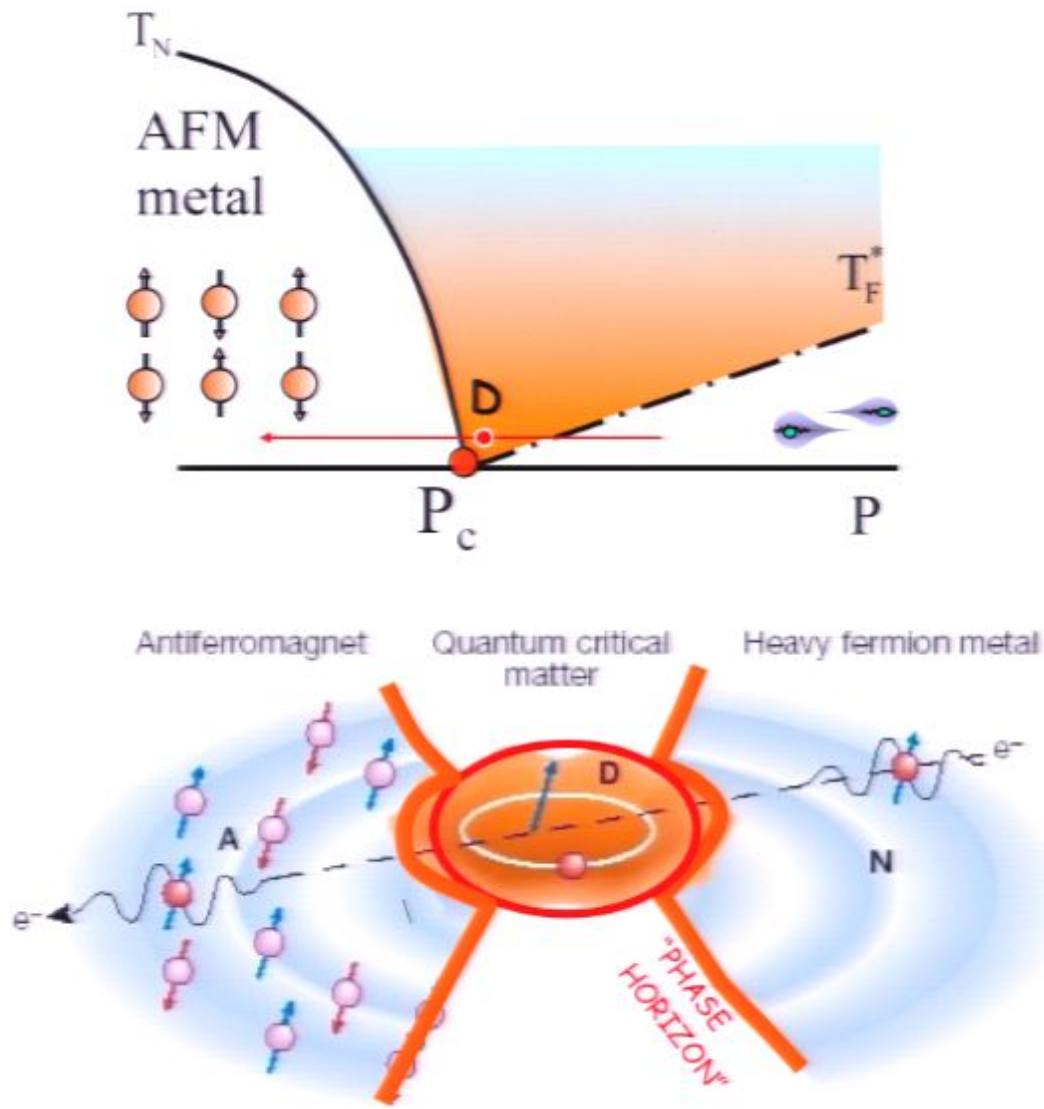


"Black Hole in the Phase Diagram".



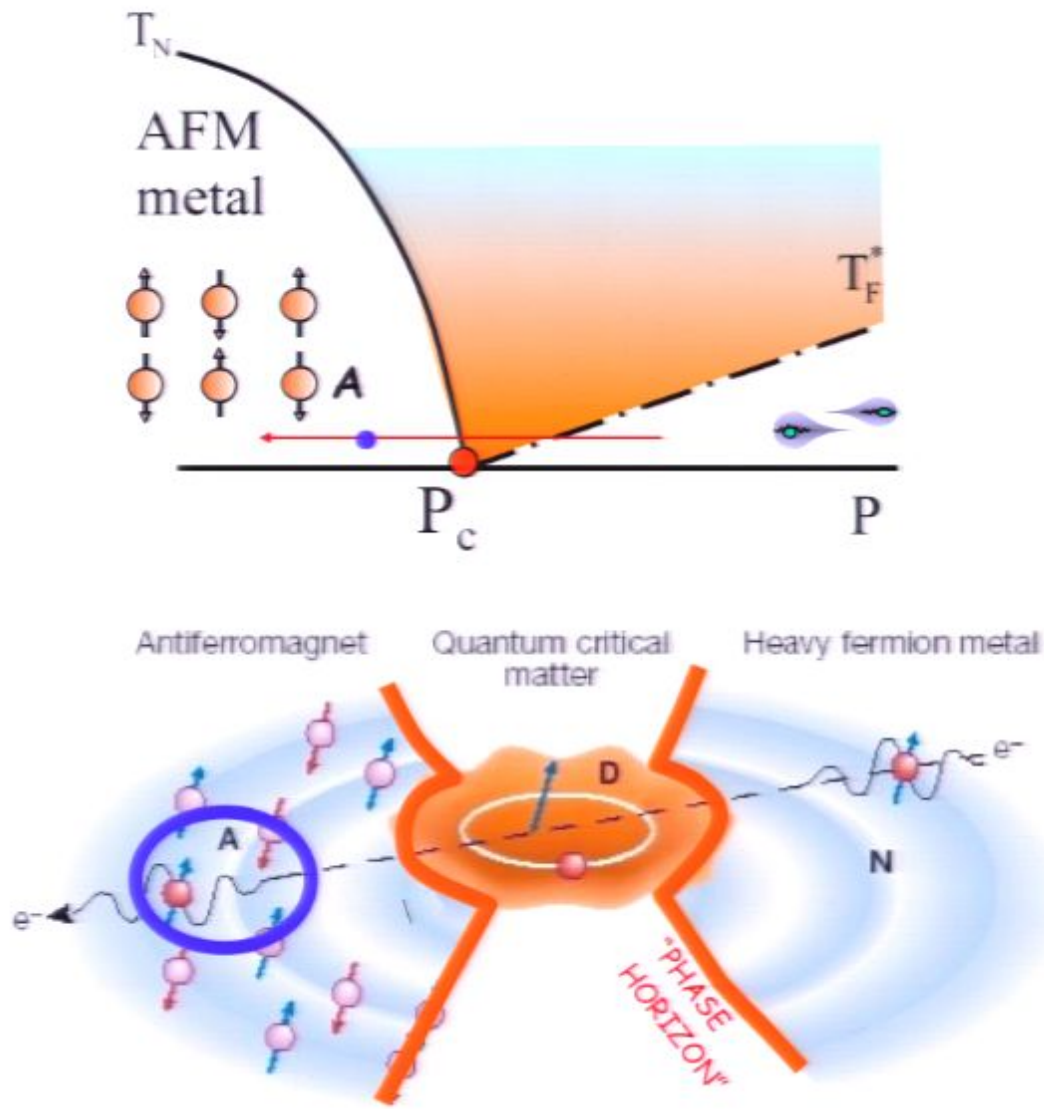
P.C and A. Schofield, Nature (2005)

"Black Hole in the Phase Diagram".



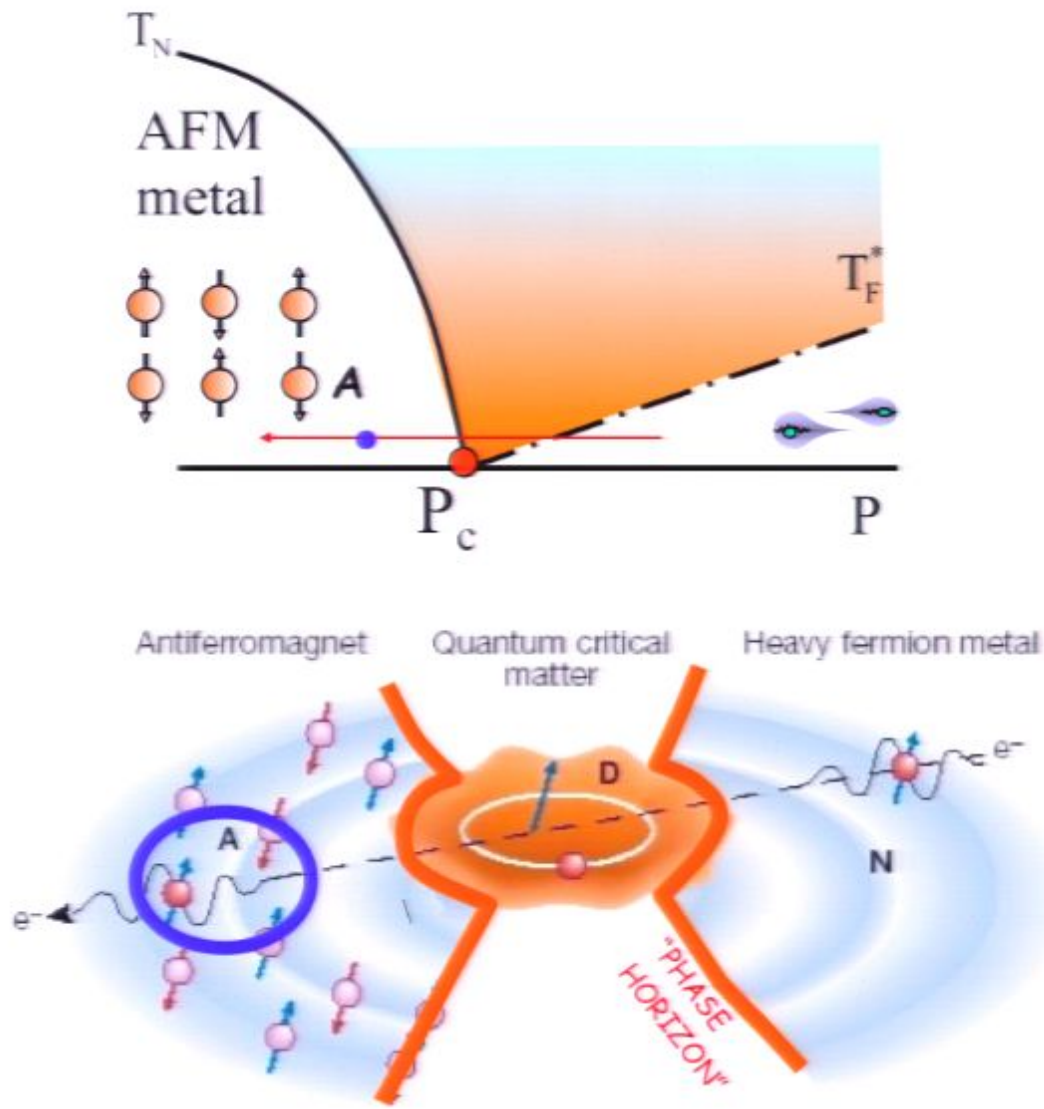
P.C and A. Schofield, Nature (2005)

"Black Hole in the Phase Diagram".



P.C and A. Schofield, Nature (2005)

"Black Hole in the Phase Diagram".

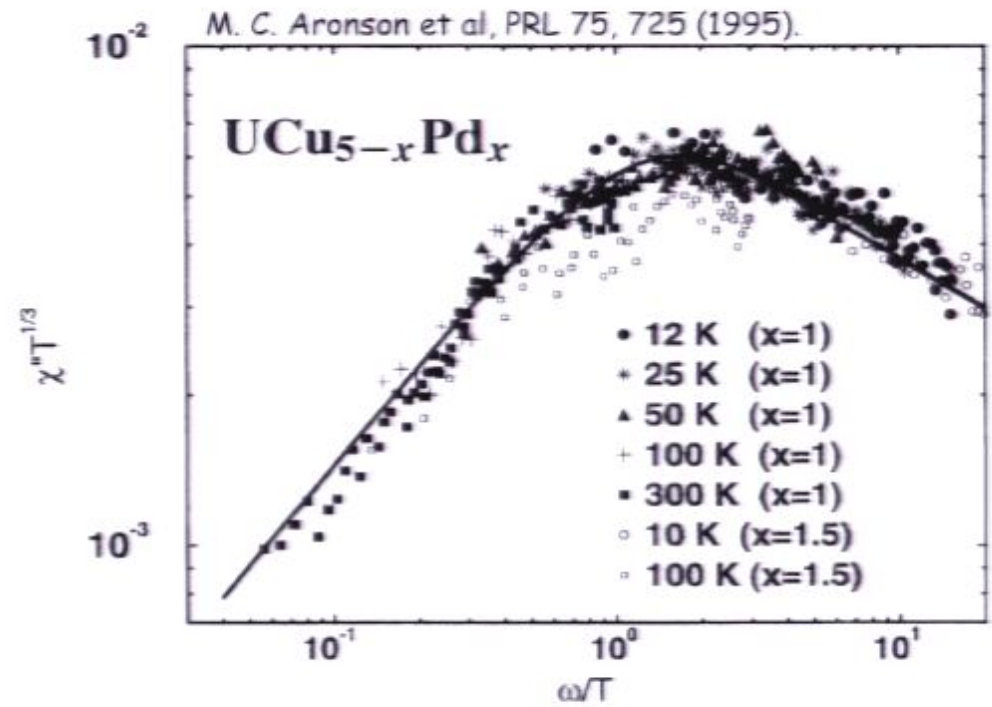


P.C and A. Schofield, Nature (2005)

E/T Scaling:



Meigan Aronson

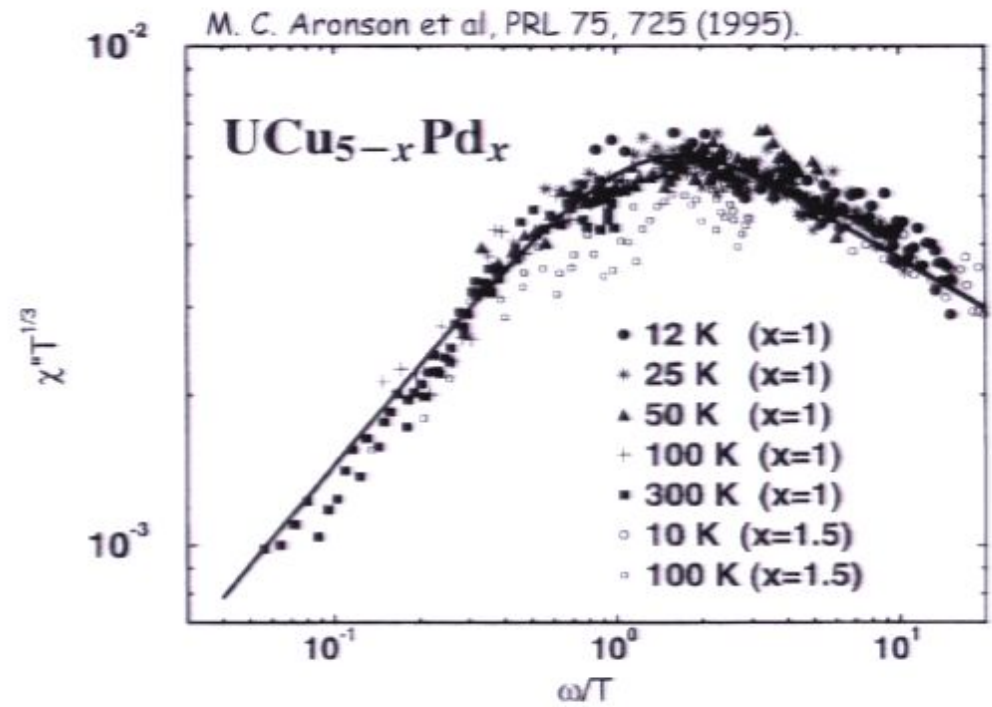


E/T Scaling:



Meigan Aronson

$$\chi''(E) = \frac{1}{E^{1-\alpha}} G\left(\frac{E}{T}\right)$$



E/T Scaling:

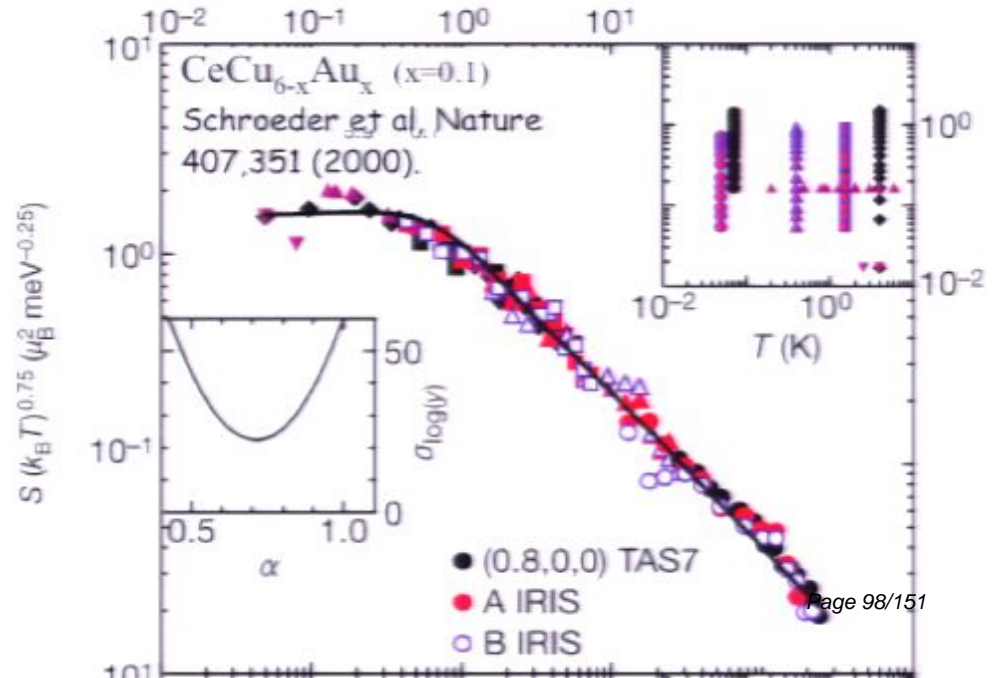
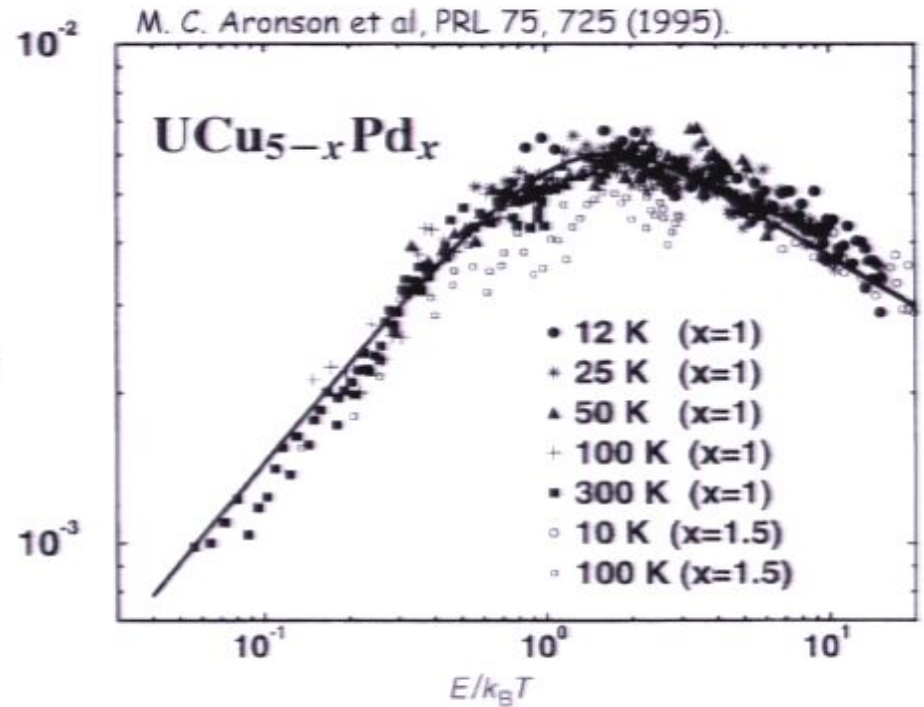


Meigan Aronson



Almut Schroeder

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E/T Scaling:



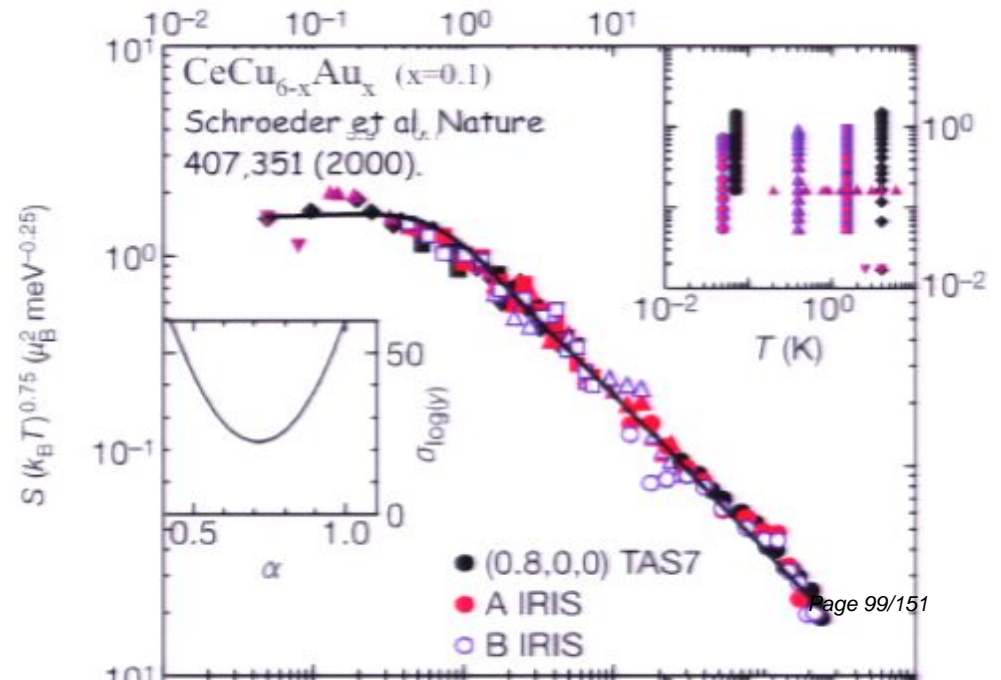
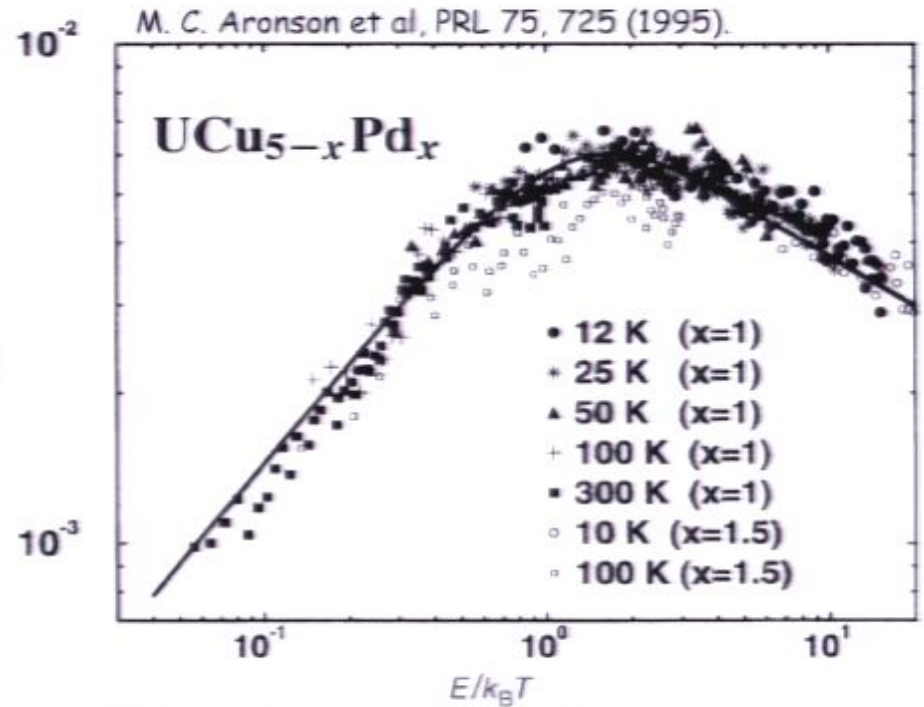
Meigan Aronson



Almut Schroeder

$$\chi''(E) = \frac{1}{E^{1-\alpha}} G\left(\frac{E}{T}\right)$$

Physics Below the upper
Critical Dimension.



$$\chi(q) = \frac{1}{q^{2 \cdot n}} f \left[\frac{q}{V} \right]$$

CeRhIn₅

E/T Scaling:



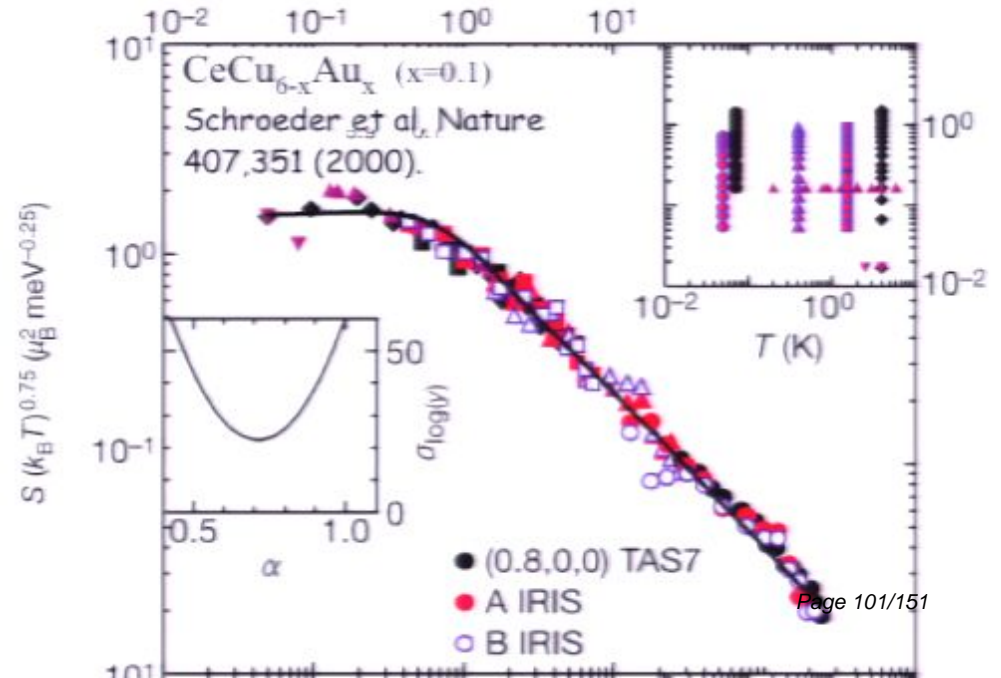
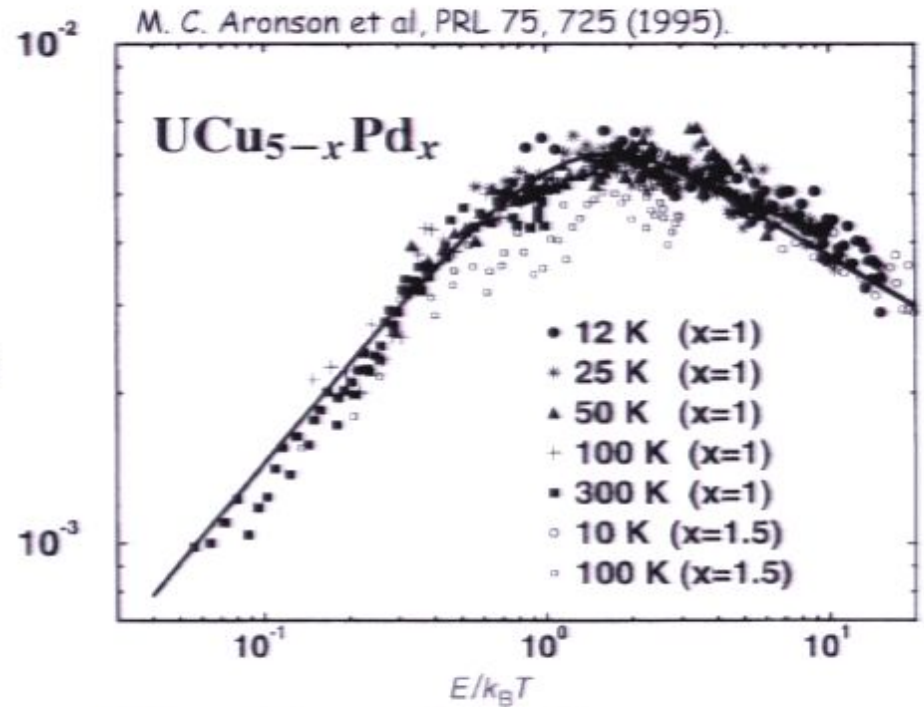
Meigan Aronson



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Physics Below the upper
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The Standard Model

Standard Model: Quantum SDW?



Doniach



Schrieffer



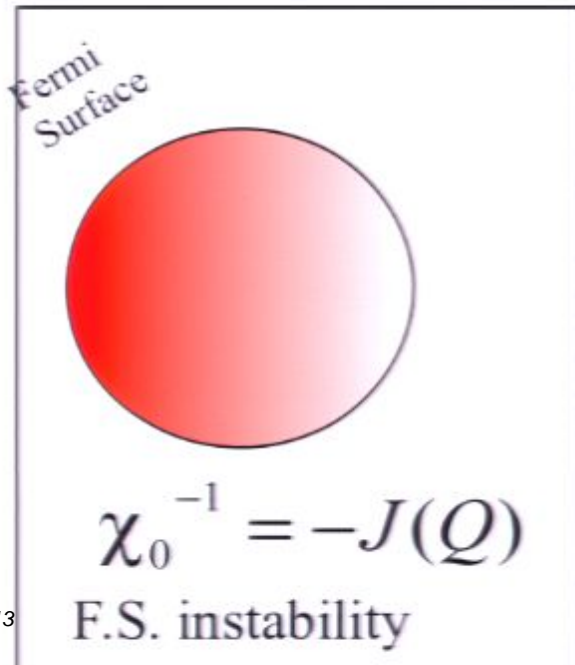
Hertz



Millis

- Moriya, Doniach, Schrieffer (60s)
- Hertz (76)
- Millis (93)

$$d_{\text{eff}} = d + z$$

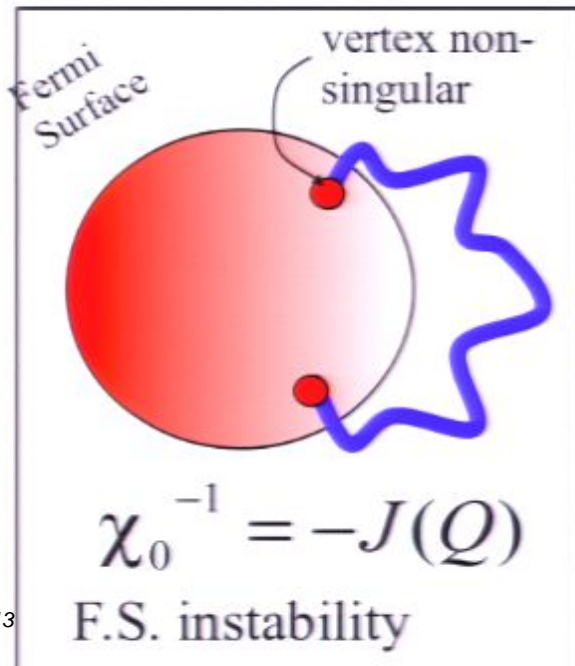


Standard Model: Quantum SDW?



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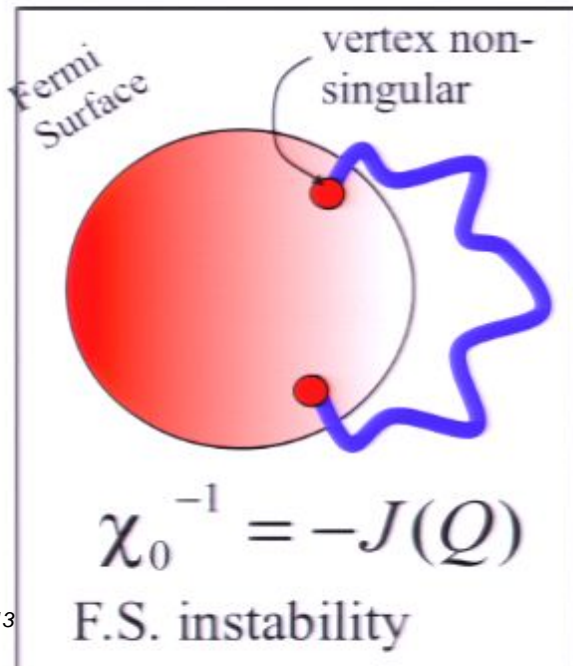


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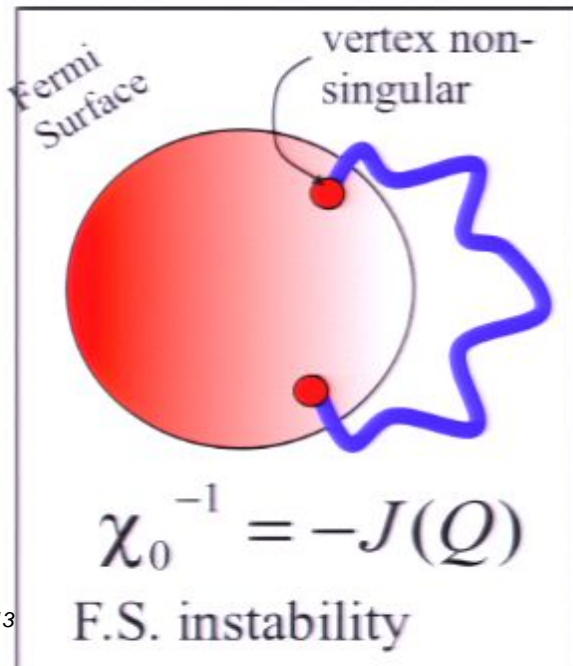
$$\chi^{-1}(q, \omega) \propto (\xi^{-2} + (q - Q)^2 - i\omega / \Gamma)$$

Standard Model: Quantum SDW?



- Moriya, Doniach, Schrieffer (60s)
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$$d_{eff} = d + z$$



$$\chi^{-1}(q, \omega) \propto (\xi^{-2} + (q - Q)^2 - i\omega / \Gamma)$$

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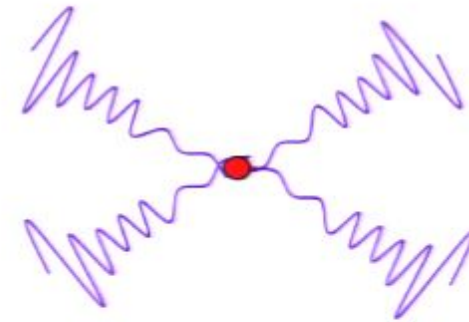
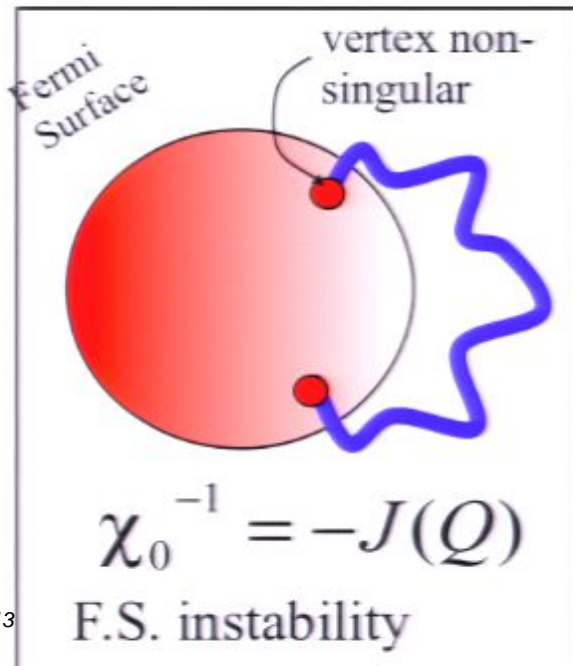
Time counts as $z=2$ scaling dimensions

Standard Model: Quantum SDW?



- Moriya, Doniach, Schrieffer (60s)
- Hertz (76)
- Millis (93)

$$d_{eff} = d + z$$



If $d + z = d + 2 > 4$:
 ϕ^4 terms “irrelevant”
 Critical modes are Gaussian.
 T is not the only energy scale.

$$\chi^{-1}(q, \omega) \propto (\xi^{-2} + (q - Q)^2 - i\omega / \Gamma)$$

$$\tau^{-1} \propto \xi^{-2}$$

Time counts as $z = 2$ scaling dimensions

New Ideas:

Break up of the electron.

$$H = \sum_k \varepsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J \sum_j (\psi_j^\dagger \vec{\sigma} \psi_j) \cdot \vec{S}_j$$

"THE BATTLEGROUND"

Kondo Lattice Model

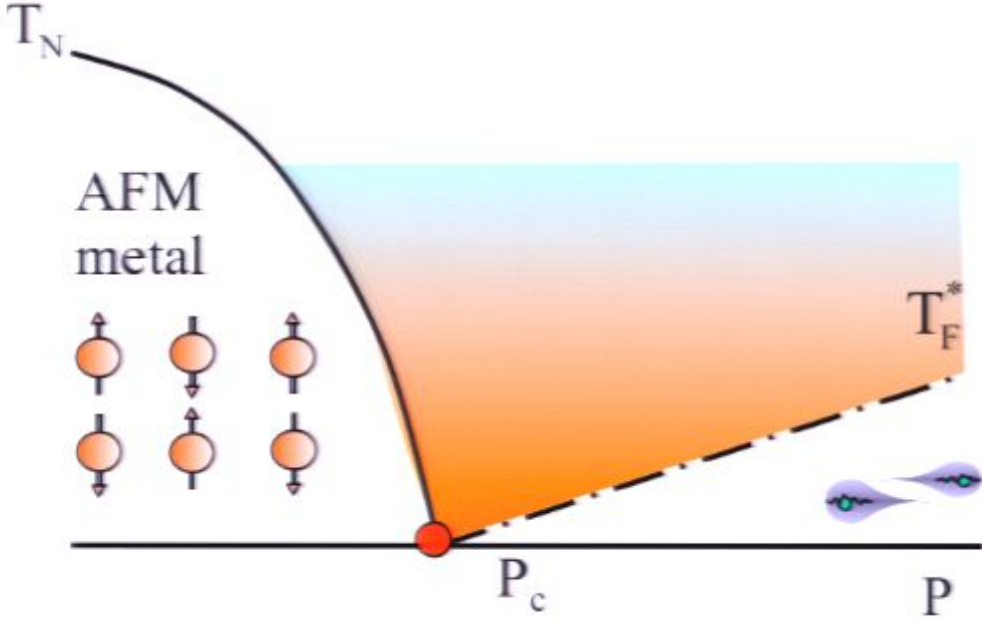
(Kasuya, 1951)

$$H = \sum_k \epsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J \sum_j (\psi_j^\dagger \vec{\sigma} \psi_j) \cdot \vec{S}_j$$

"THE BATTLEGROUND"

Kondo Lattice Model
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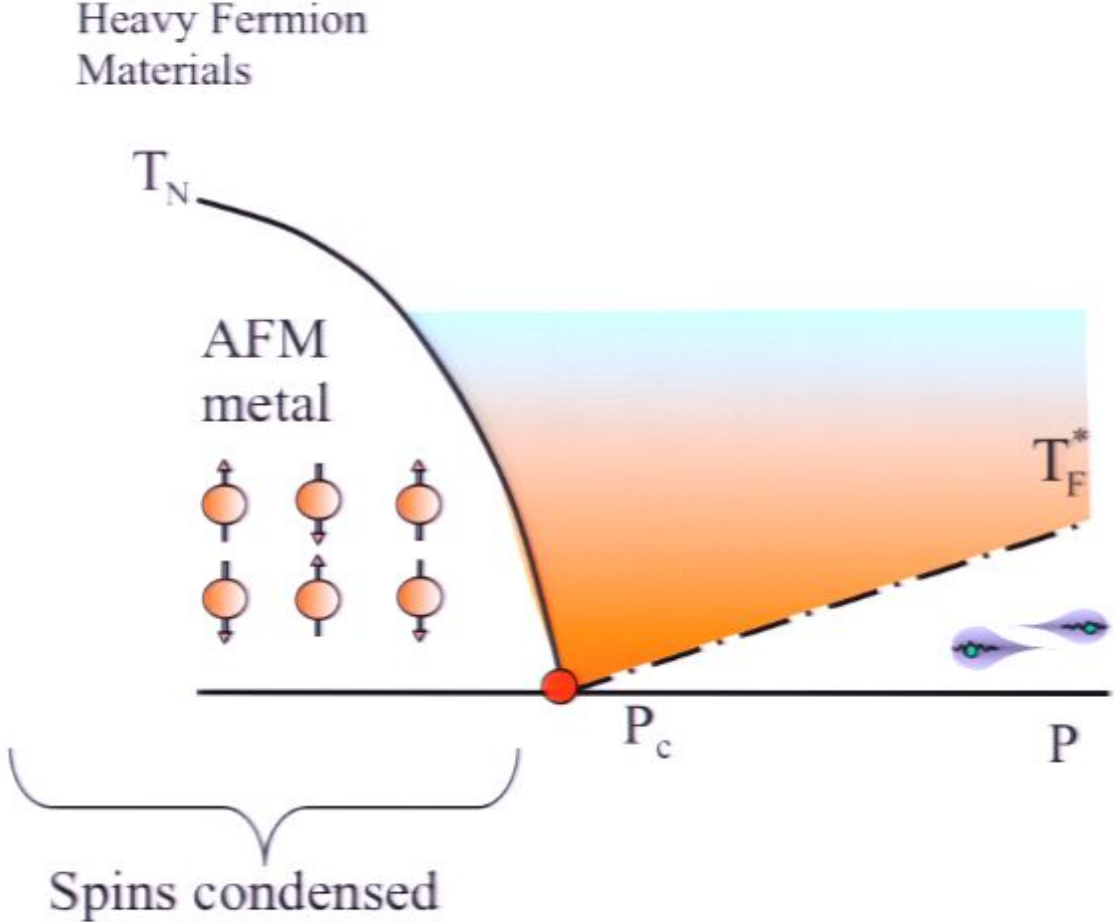
Heavy Fermion
Materials



$$H = \sum_k \epsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J \sum_j (\psi_j^\dagger \vec{\sigma} \psi_j) \cdot \vec{S}_j$$

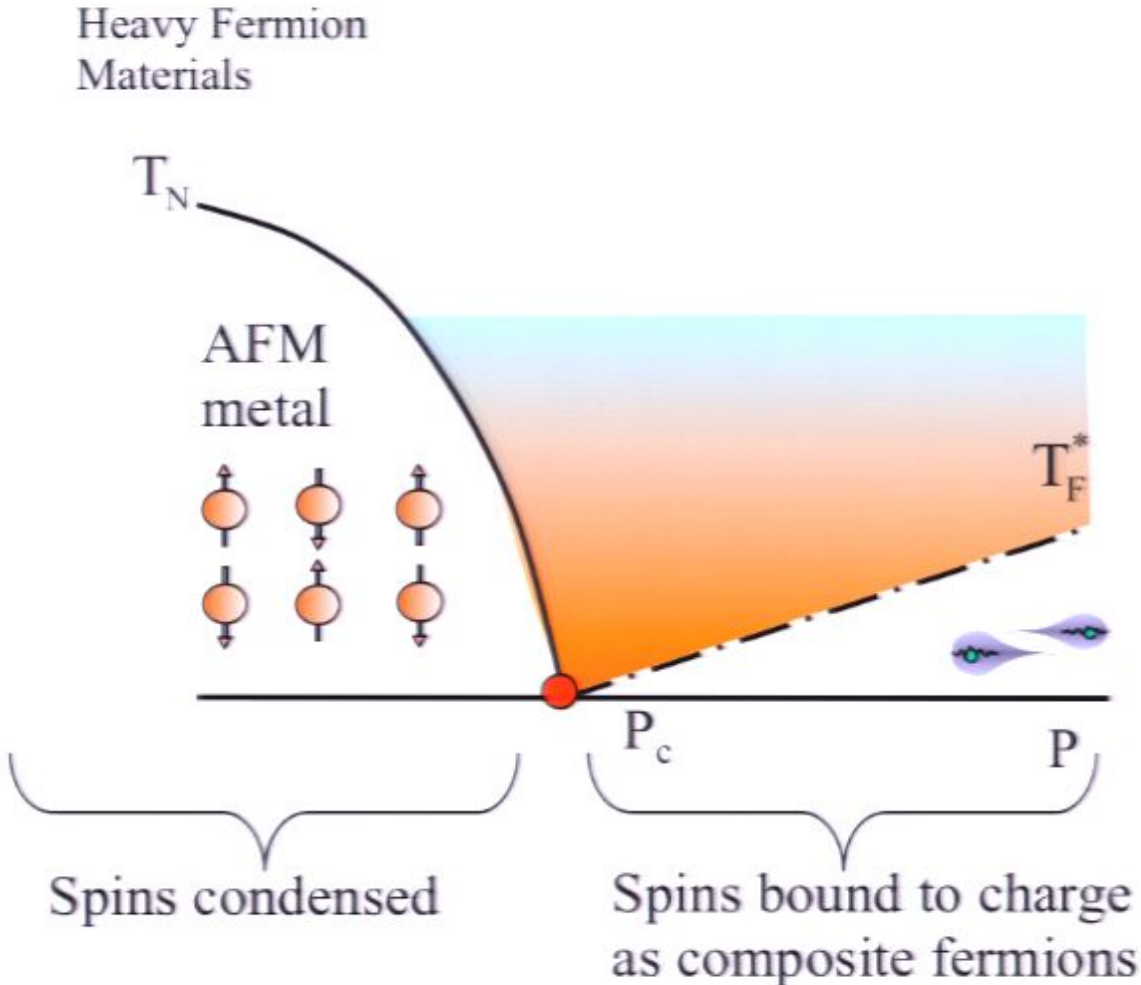
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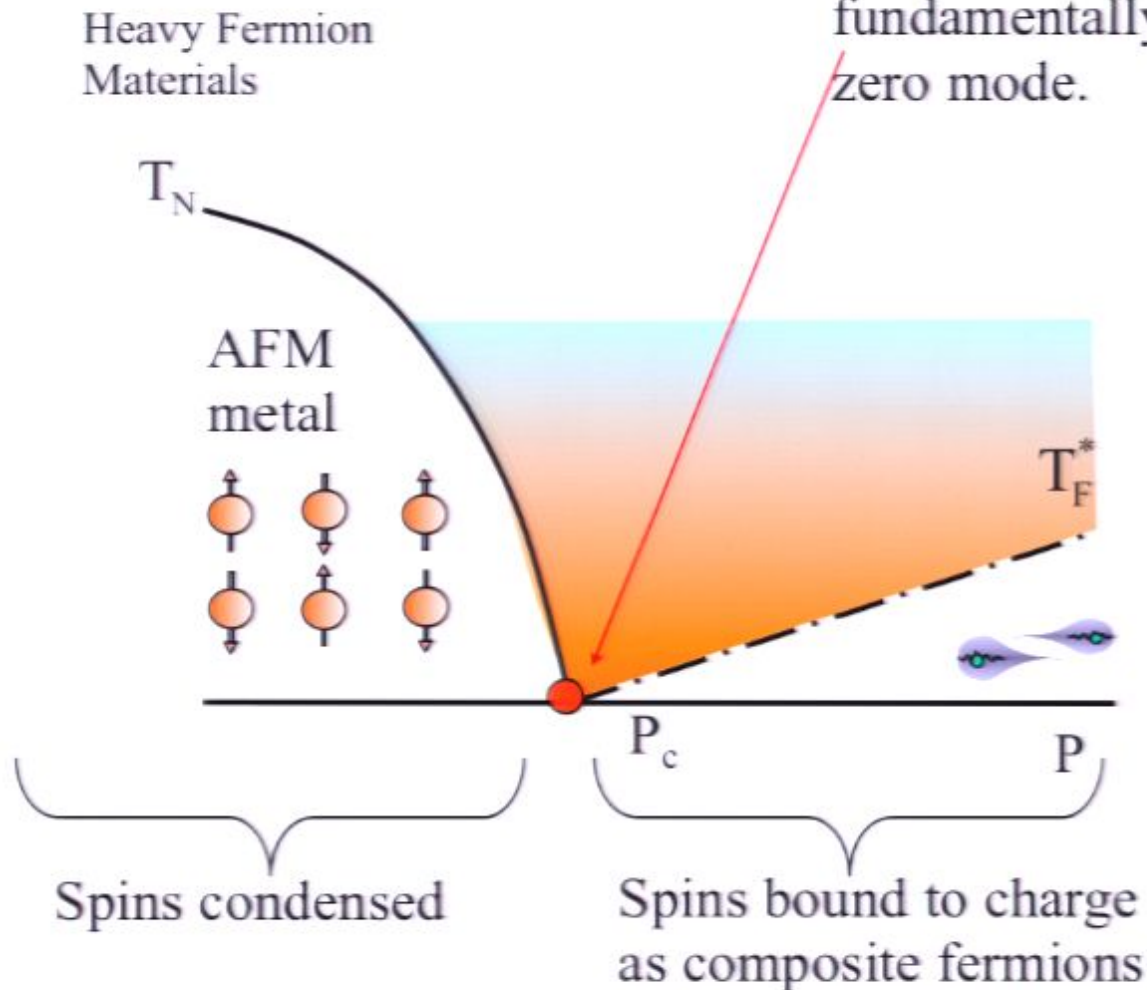
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"THE BATTLEGROUND"

Kondo Lattice Model

(Kasuya, 1951)

Deconfinement of spin:
fundamentally new kind of
zero mode.



New Methods



e^-

New Methods



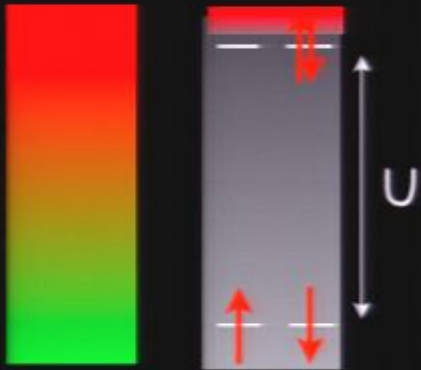
New Methods



Elimination of States
implies Gauge Fields.

(Read News, PC, Millis Lee... 80's)

New Methods



e^-

spin

Φ

W

$$\vec{S} = b^\dagger_\alpha \left(\frac{\vec{\sigma}}{2} \right)_{\alpha\beta} b_\beta$$

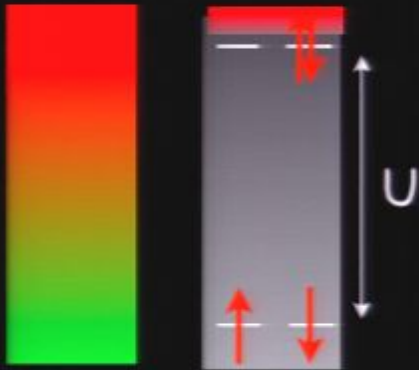
J. Schwinger '55

$$U(1) : b_j \rightarrow e^{i\theta_j} b_j$$

Elimination of States
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New Methods



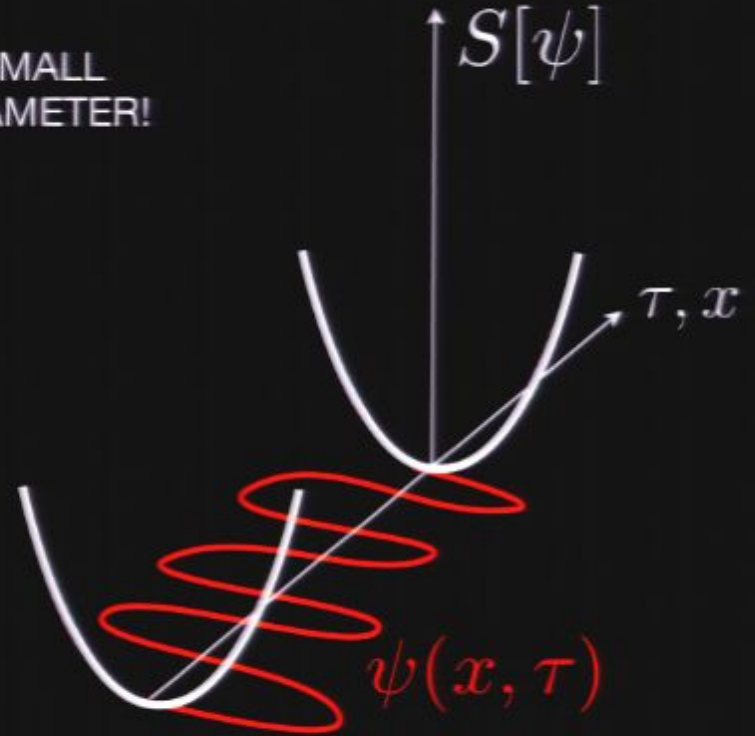
e^- spin
 Φ w

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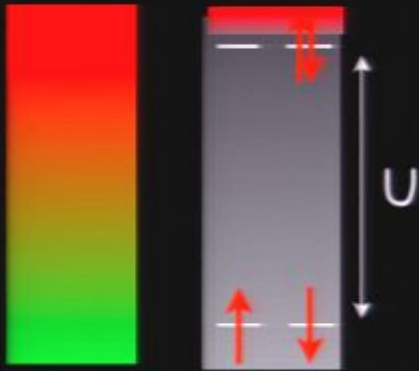
NO SMALL
PARAMETER!



Elimination of States
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New Methods



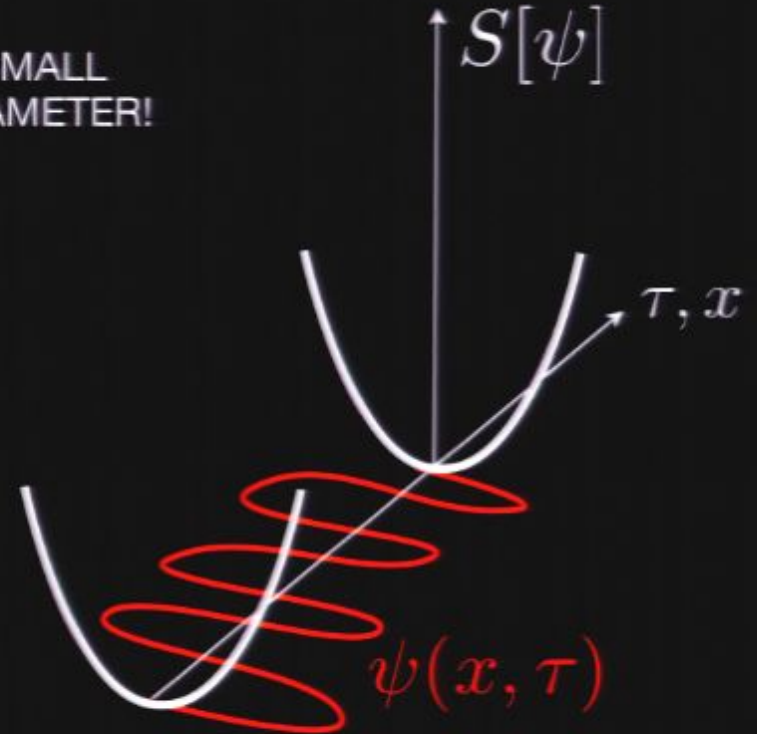
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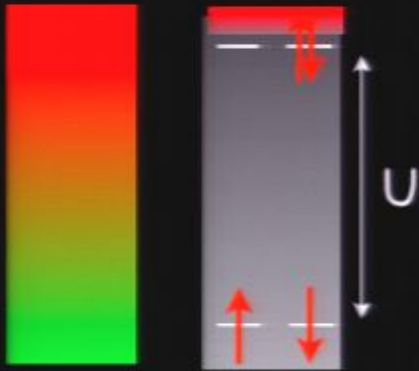


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Large N : family of models with “N” spin components, which retain the key physics and can be solved in the large N limit.

New Methods



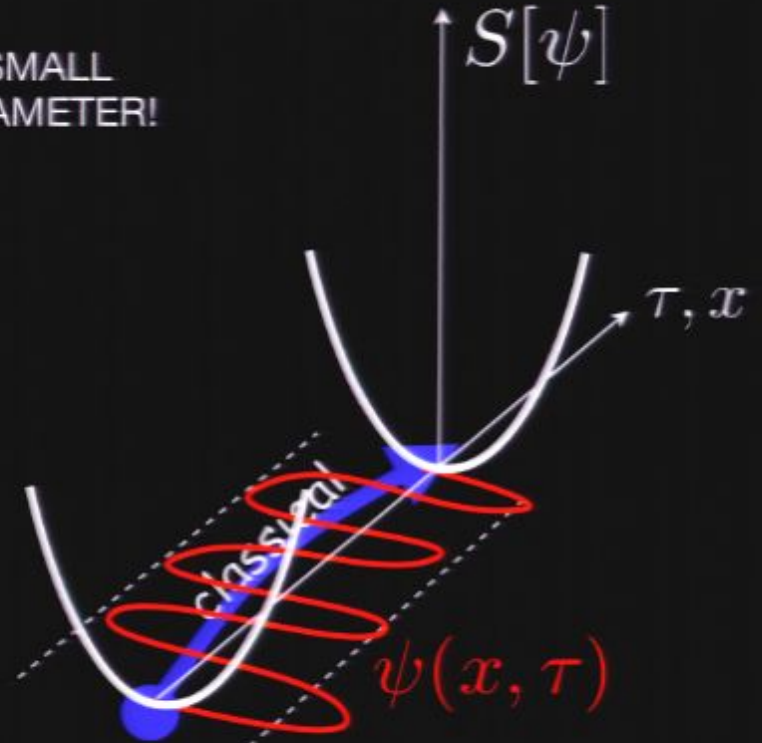
e^- spin
 Φ w

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J. Schwinger '55

$$U(1) : b_j \rightarrow e^{i\theta_j} b_j$$

NO SMALL
PARAMETER!



$$\frac{1}{N} \sim \hbar_{eff}$$

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(Read News, PC, Millis Lee... 80s)

Large N : family of models with “N” spin components, which retain the key physics and can be solved in the large N limit.

New Ideas

New Ideas

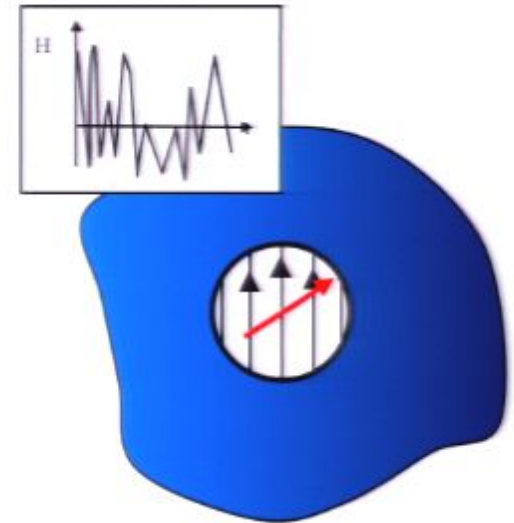
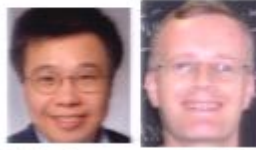
Si, Ingersent

- **Local quantum criticality**

(Si, Ingersent, Smith, Rabello, Nature 2001):

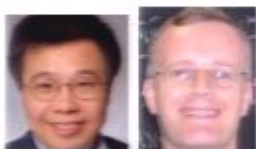
Spin is the critical mode,
Fluctuations critical in time.

Requires a two dimensional spin fluid



New Ideas

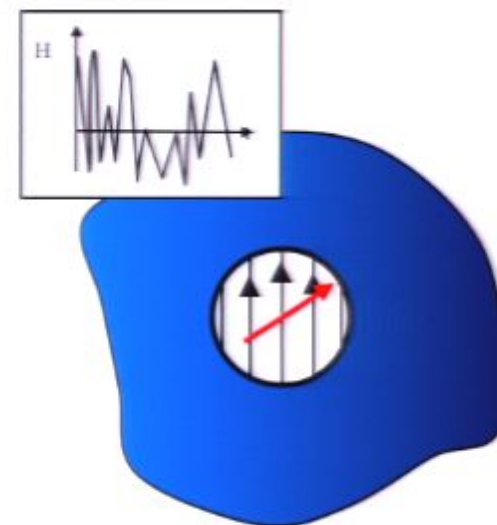
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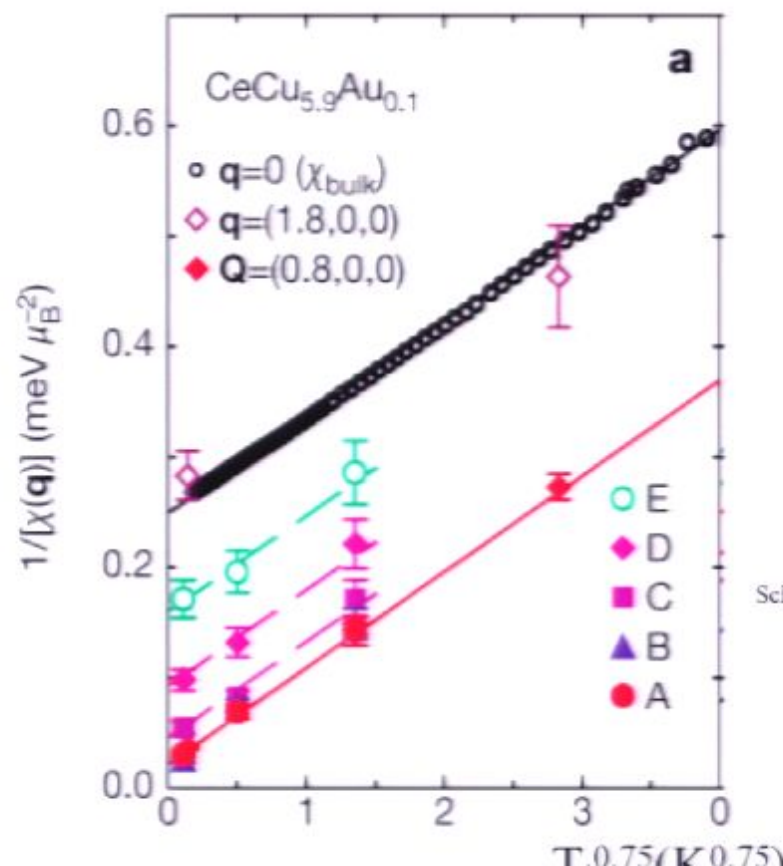
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Requires a two dimensional spin fluid

Locality of critical fluctuations

$$\chi^{-1} = \chi_0^{-1} + AT^\alpha$$



Schroeder et al. Nature 407,351(2000).

New Ideas

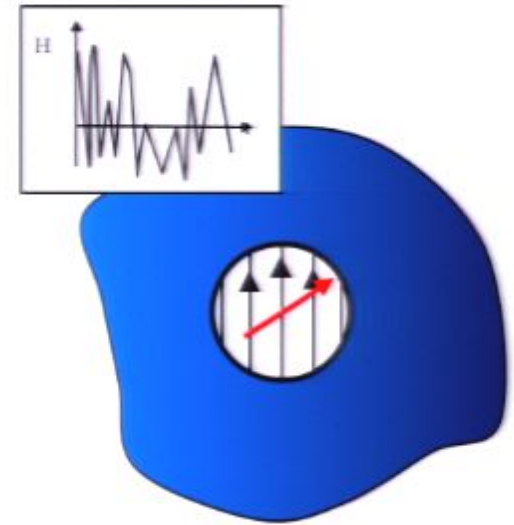
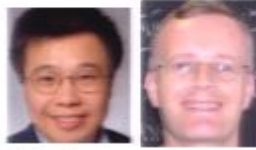
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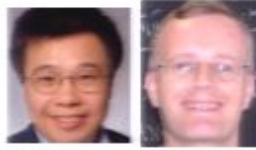
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New Ideas

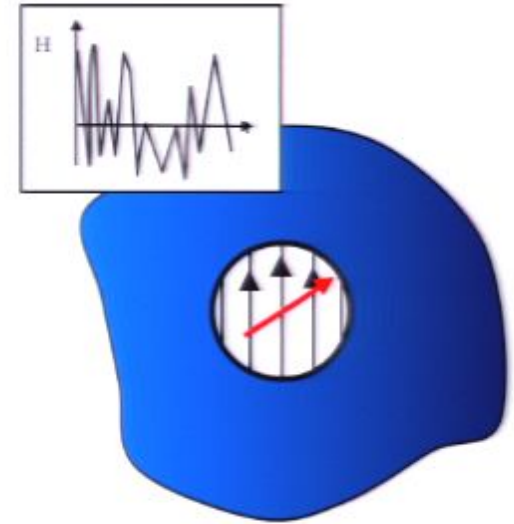
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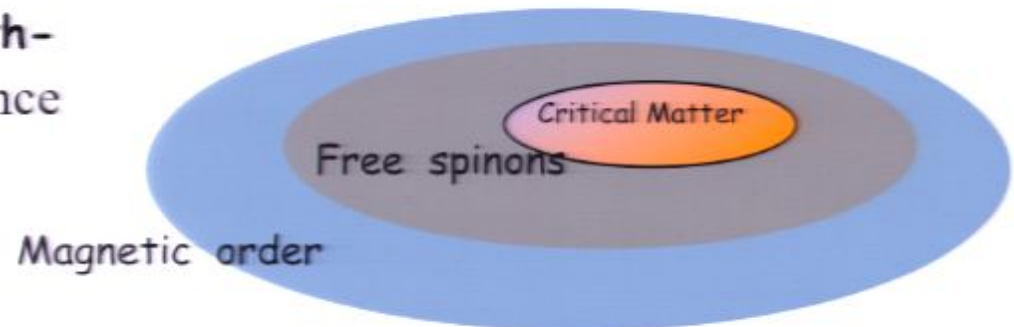


Requires a two dimensional spin fluid

- **Deconfined Criticality: Two diverging length-scales.** (Hermele et al 2004; Senthil et al, Science 2004).

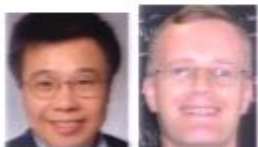


Senthil Sachdev Vishwanath



New Ideas

Si, Ingersent

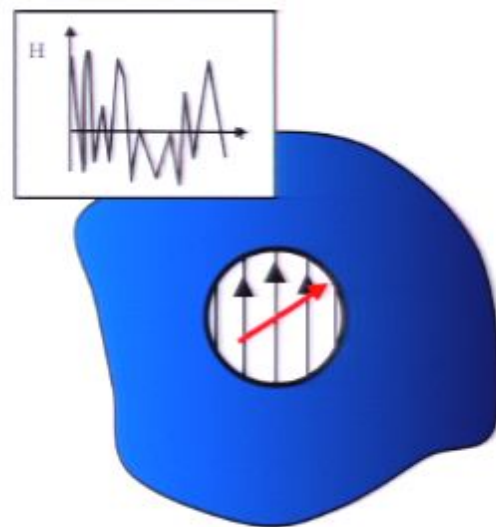


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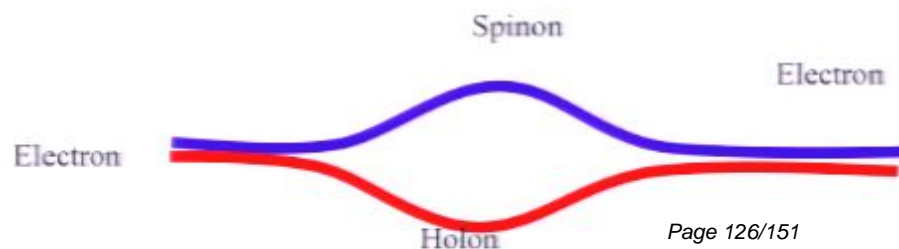
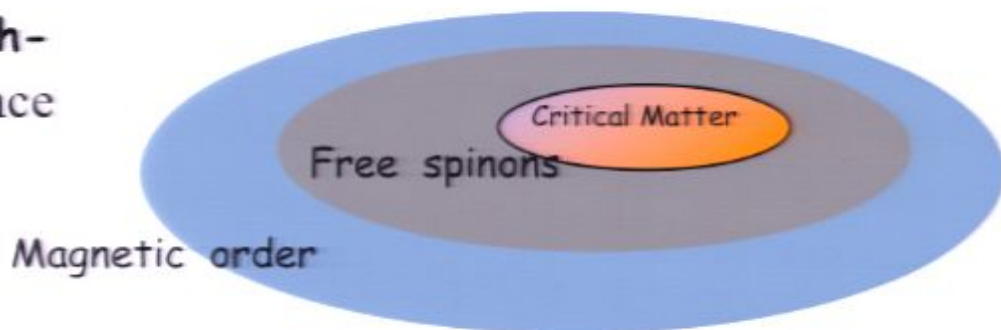
Senthil Sachdev Vishwanath

• Search for a new mean-field theory.

(PC, Pepin et al JCM, 2001, Rech et al 2005, Lebanon et al 2006.).



Pepin



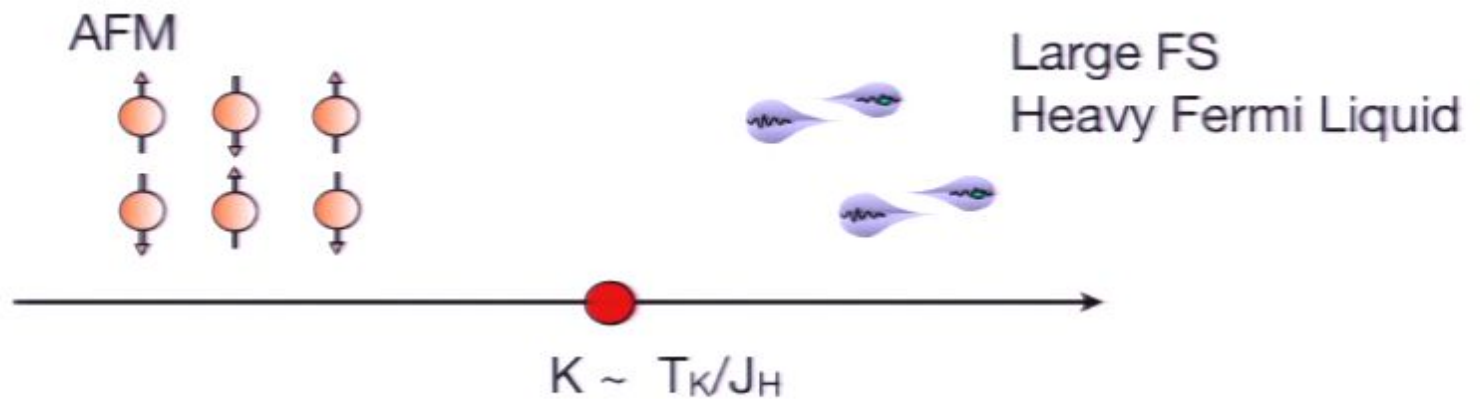
Q-Frustration

Kondo meets frustration

- Frustration and Kondo have different effects.

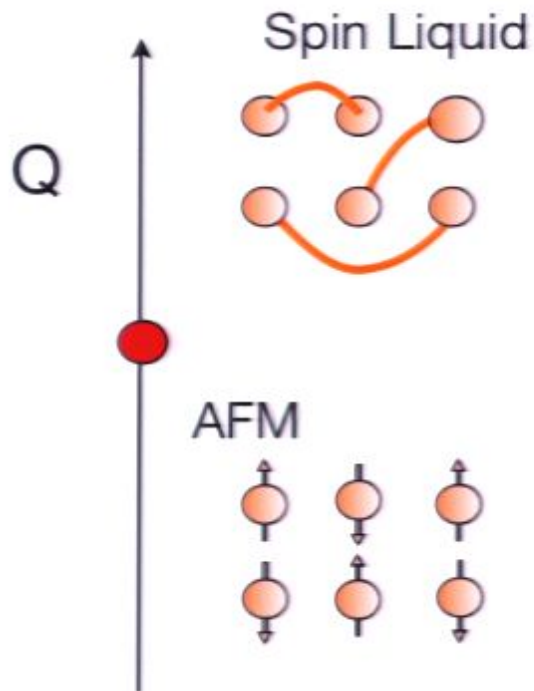
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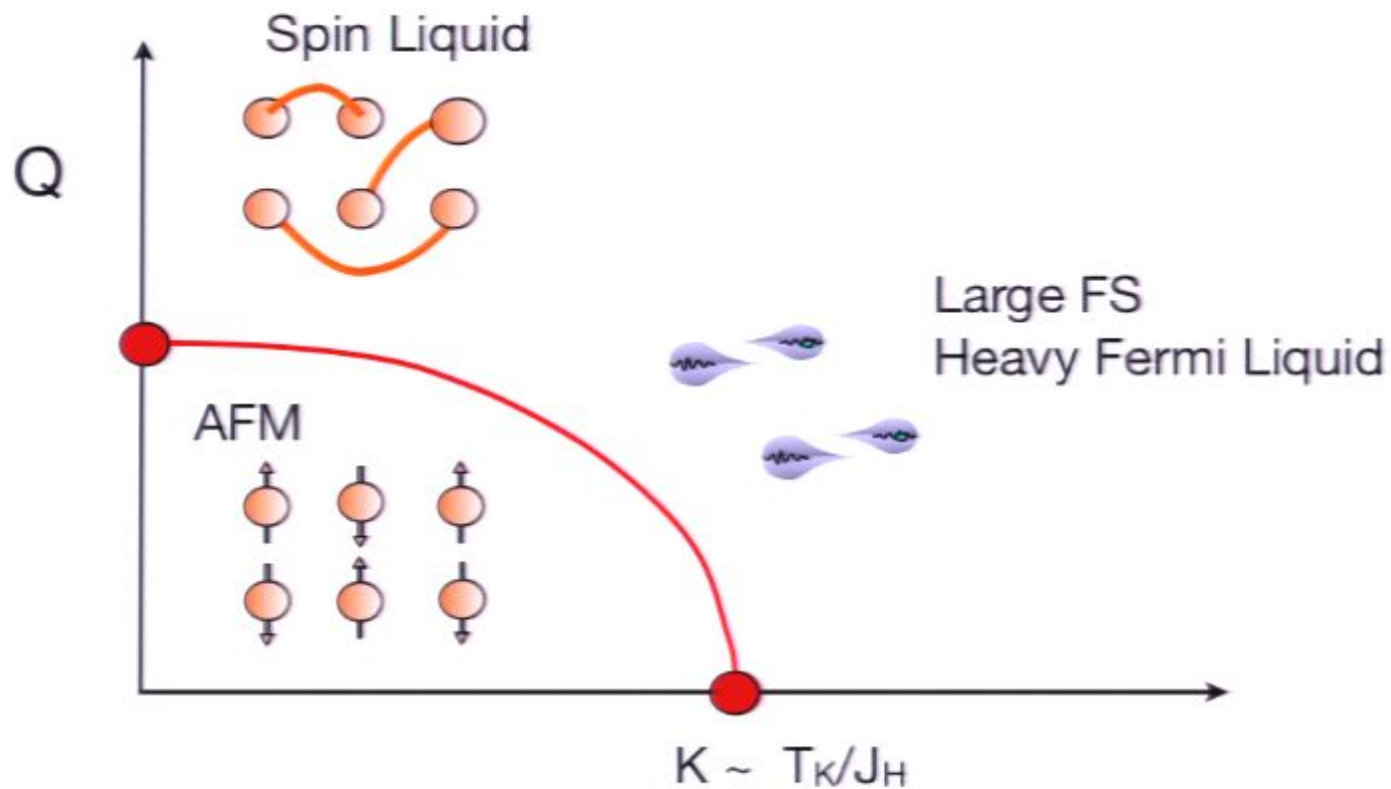
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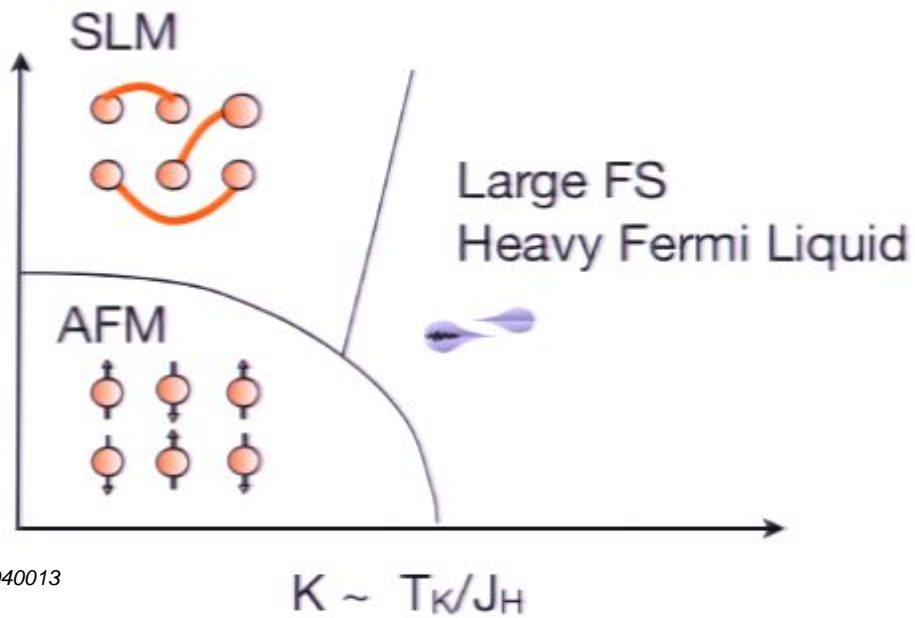


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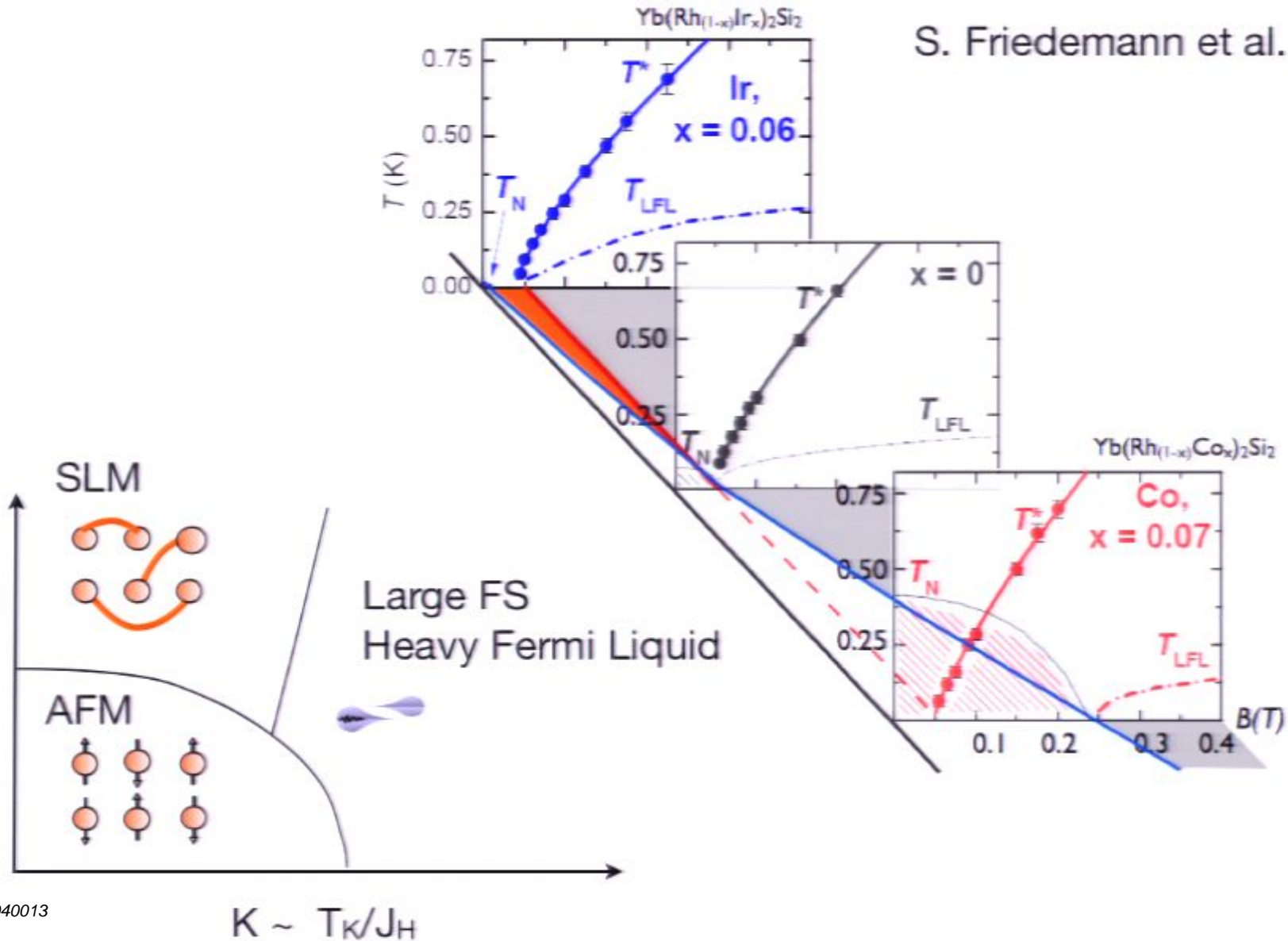


Experimental Support I



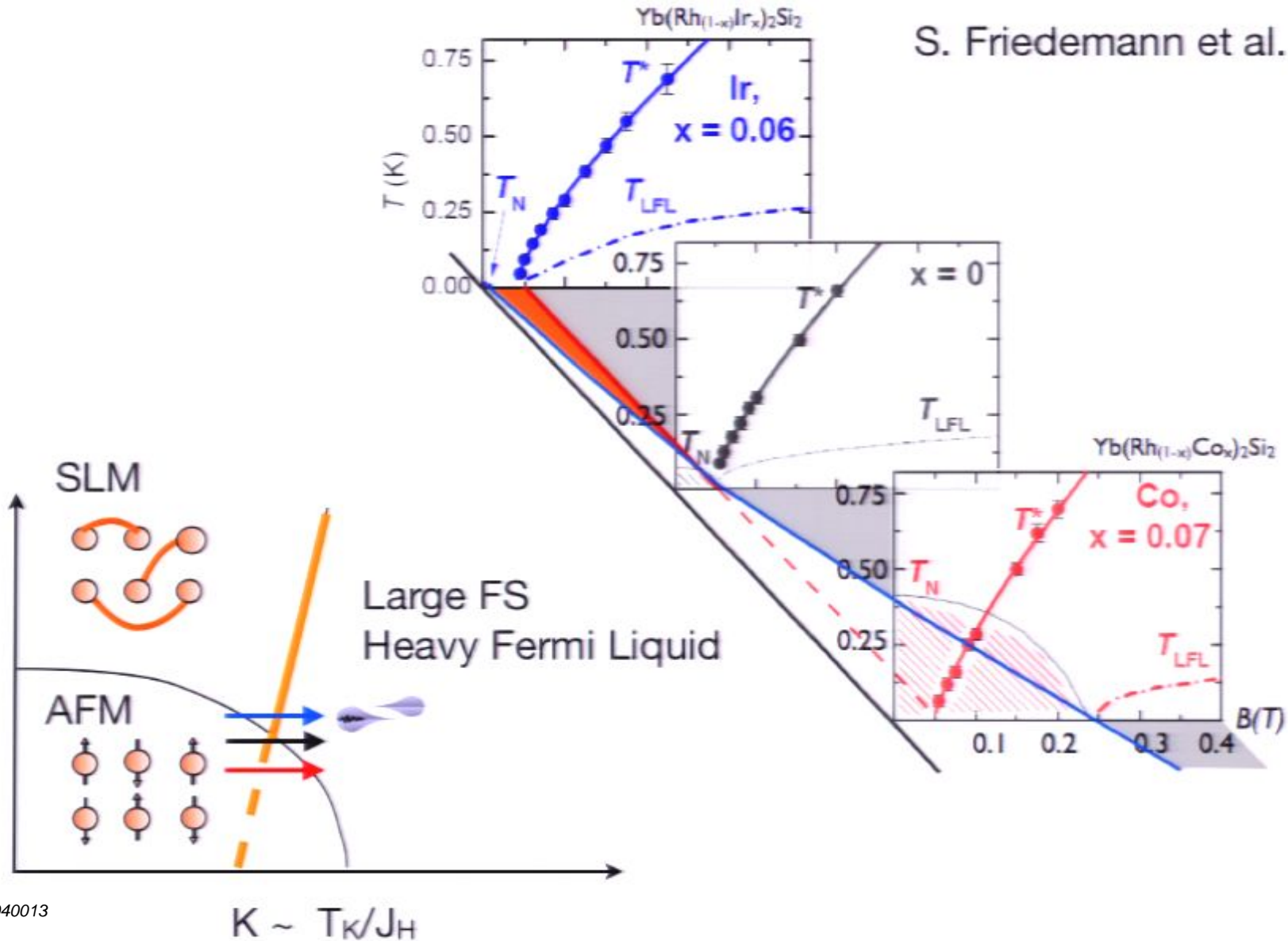
Experimental Support I

S. Friedemann et al. (2009).

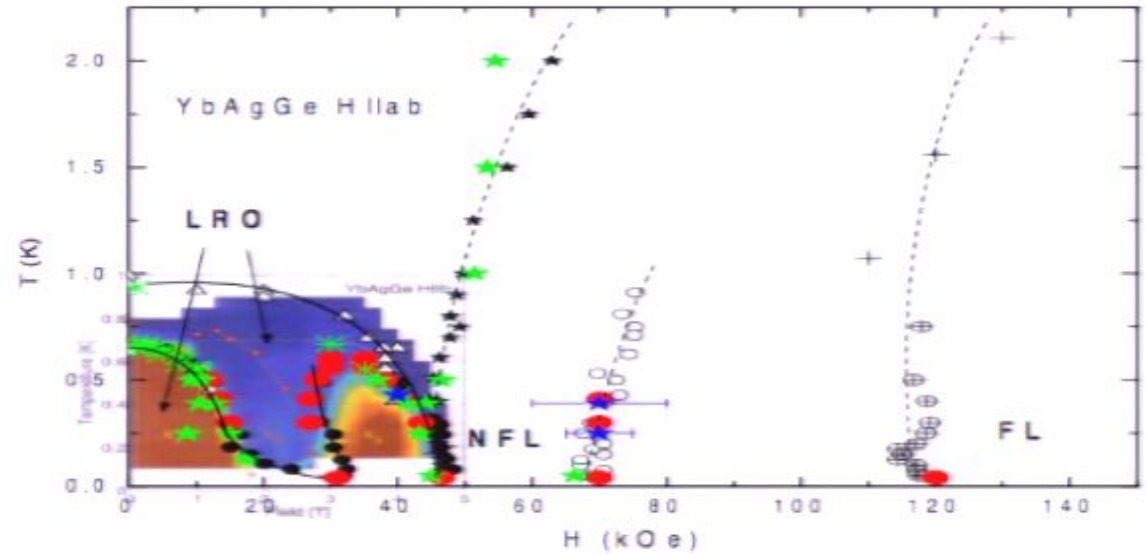


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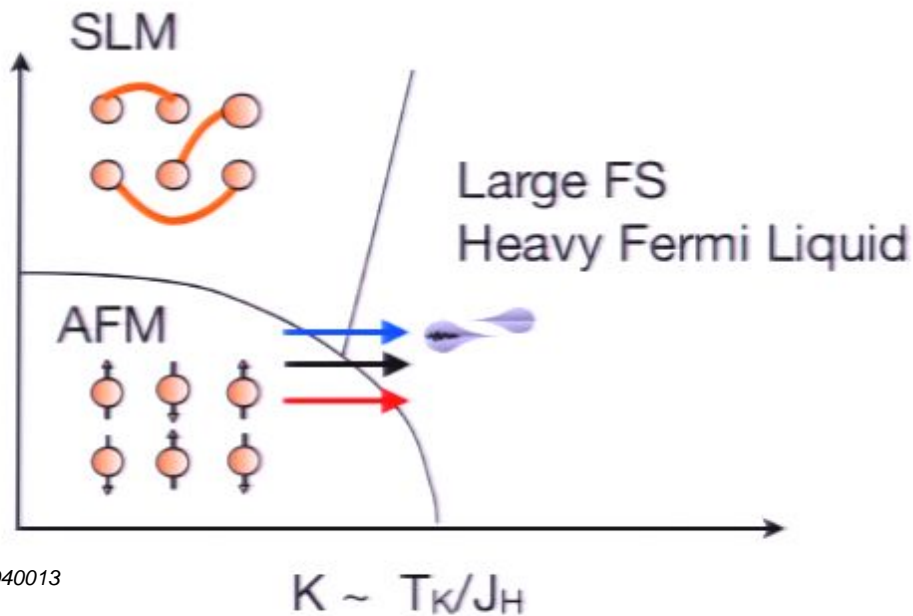
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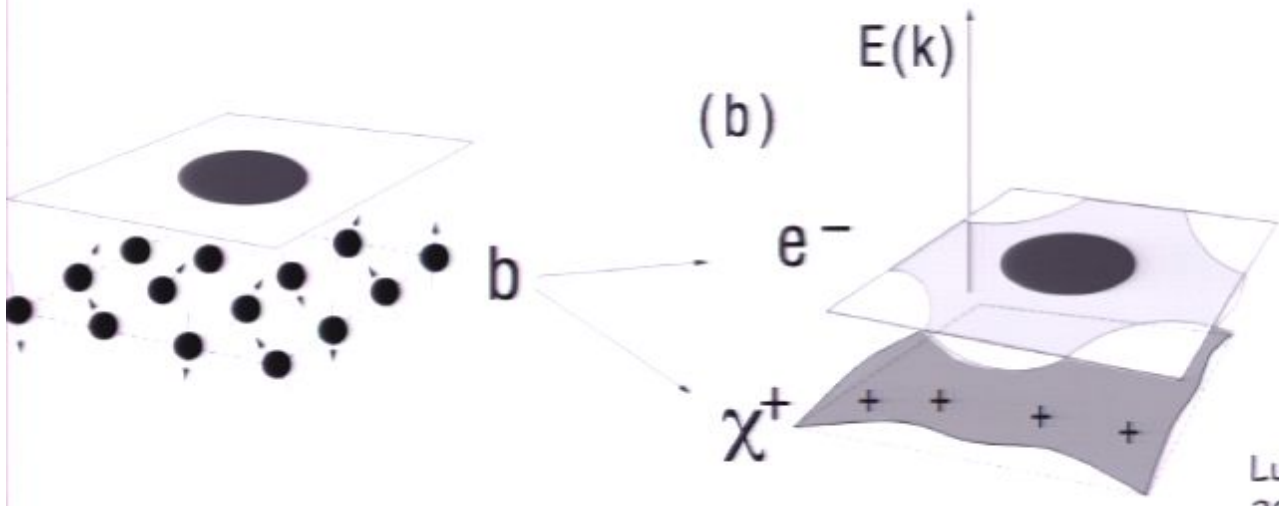


Experimental Support II



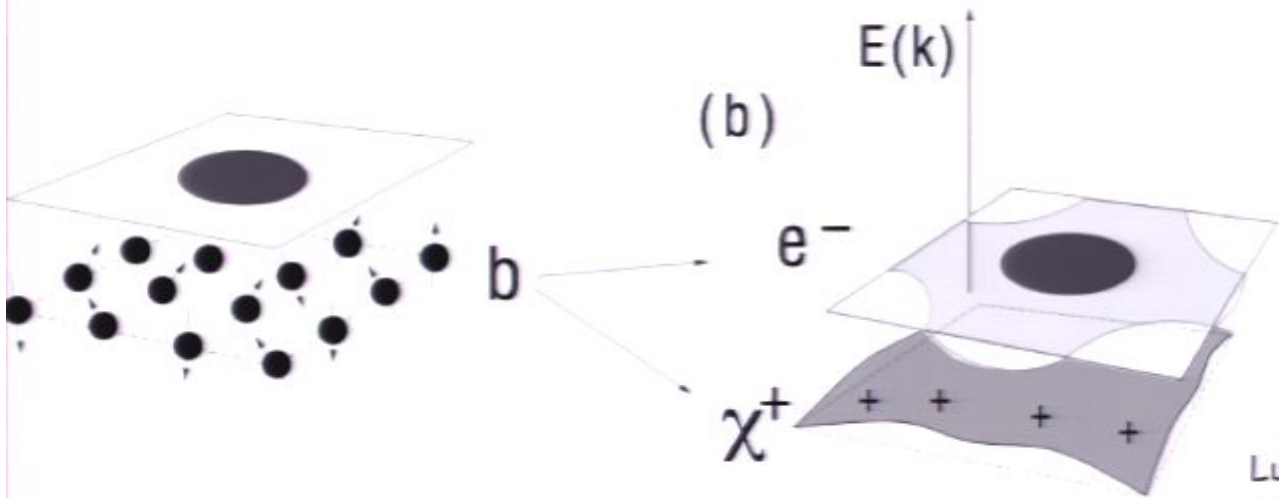
Canfield et al (unpublished)





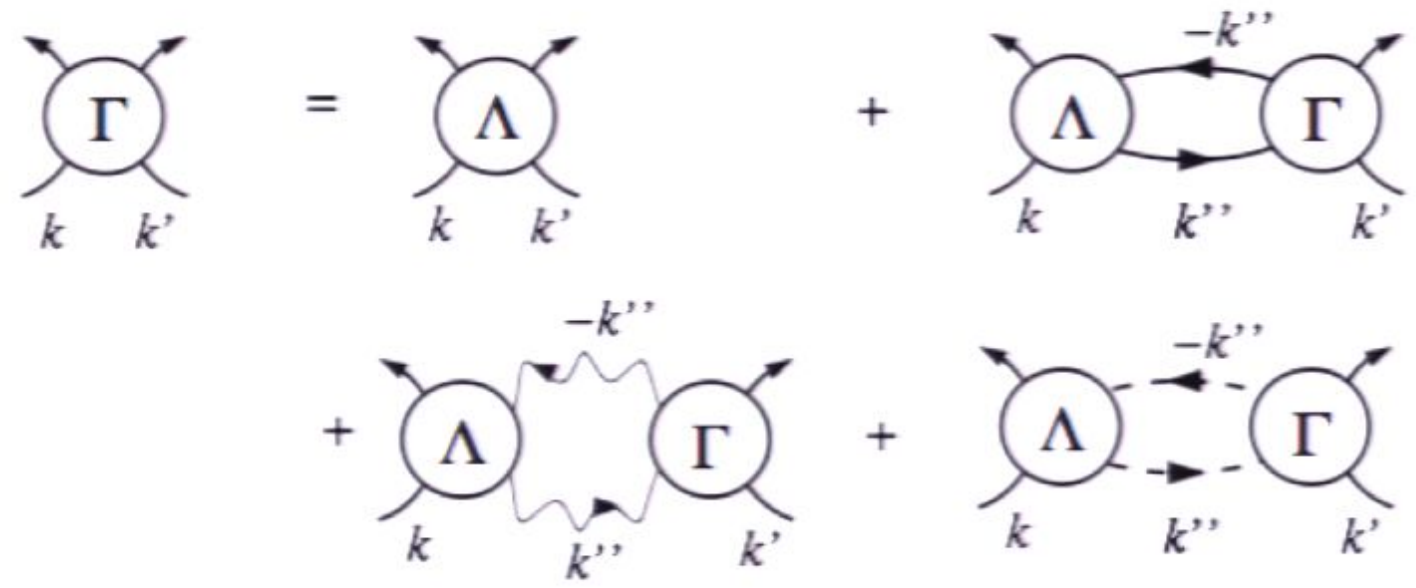
$$\frac{n_e}{K} = N \frac{v_{FS}}{(2\pi)^D} - \underbrace{\frac{v_\chi}{(2\pi)^D}}_{=1}$$

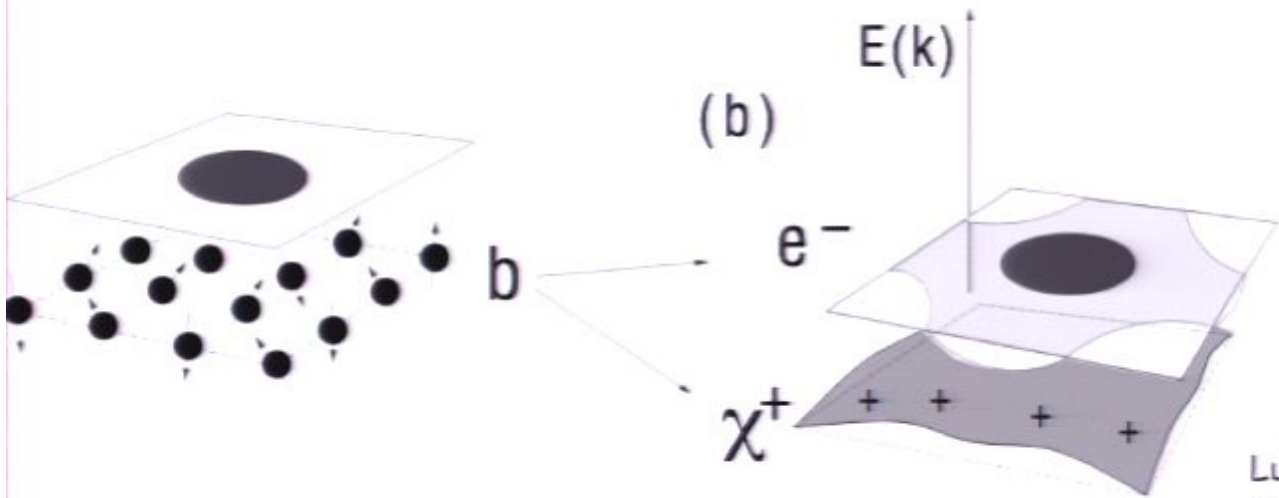
Luttinger sum rule for Kondo Lattice (Oshikawa 2000) P.C, I.Paul, J. Rech (05)



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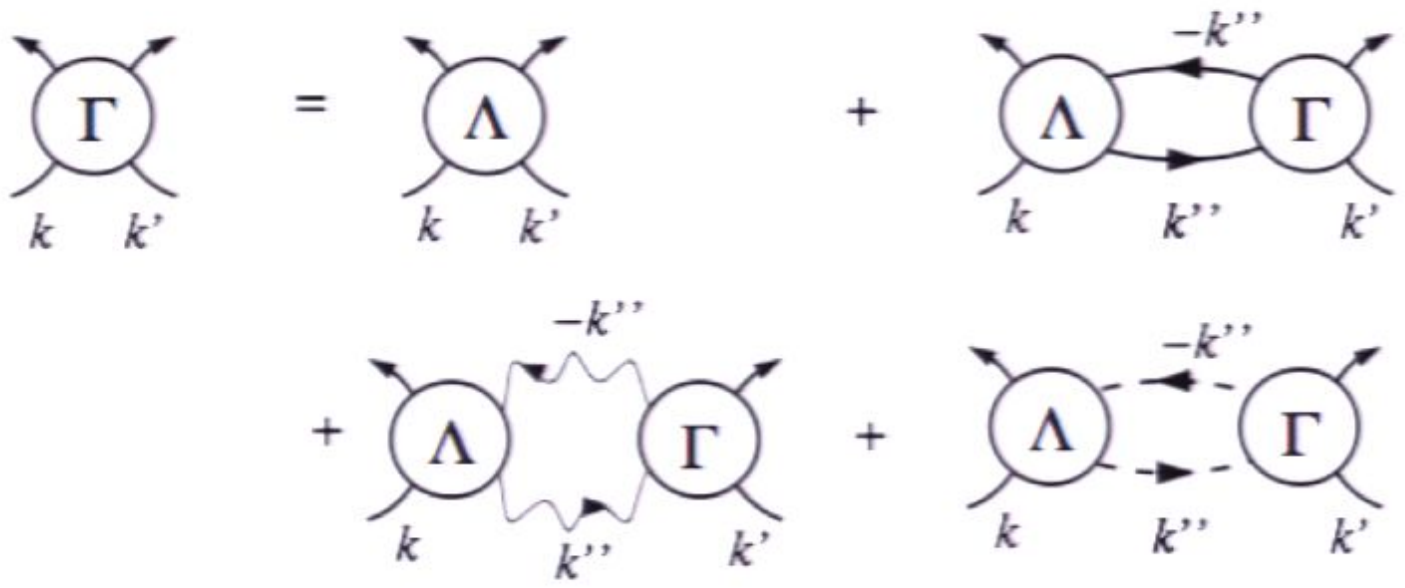
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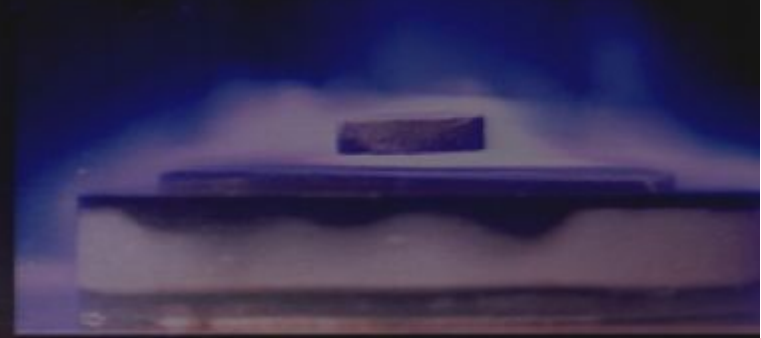


Fermi Liquid scattering parameters determined primarily by excitation of low-lying spinon and holon states..

Virtual spinons.

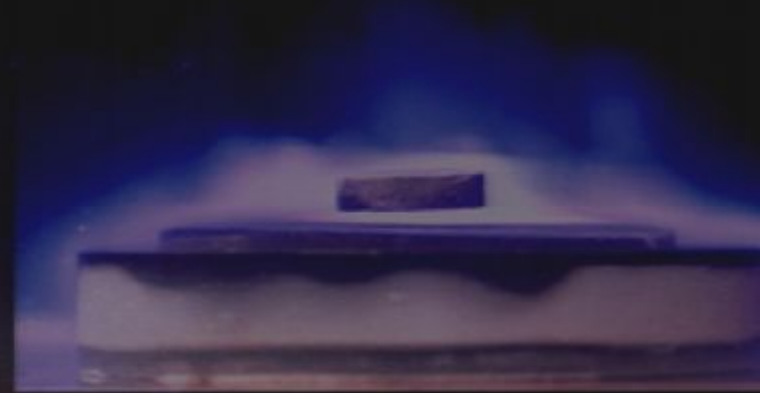
Virtual holons.

Towards a new Conjunction of ideas.

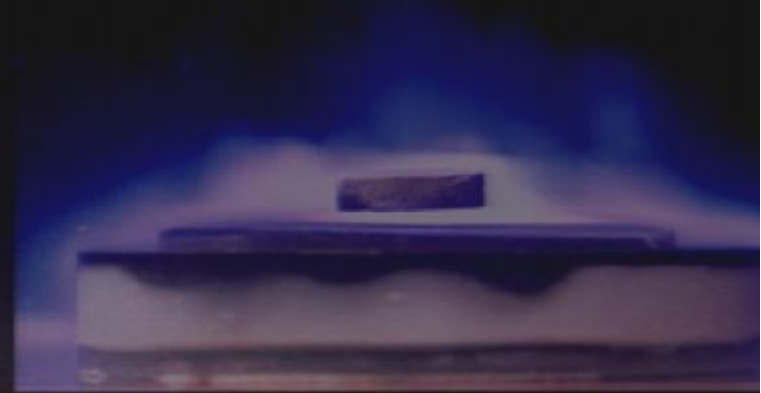


Towards a new Conjunction of ideas.

- Liu, McGreevy and Vegh arXiv:0903.2477
"Non-Fermi liquids from holography" (large N /AdsCFT).

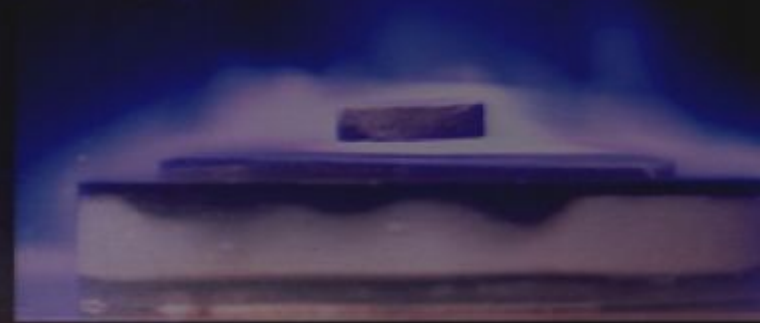


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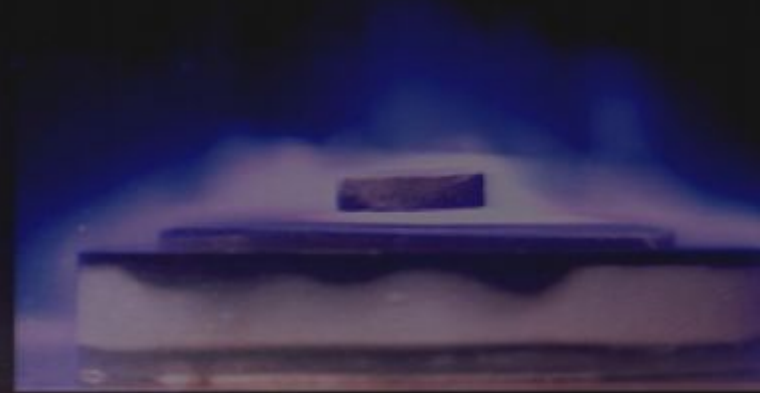
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All have in common the idea of a critical Fermi surface with branch-cuts replacing poles.

Collaborators



Flint



Nev.



Dzero



Rech

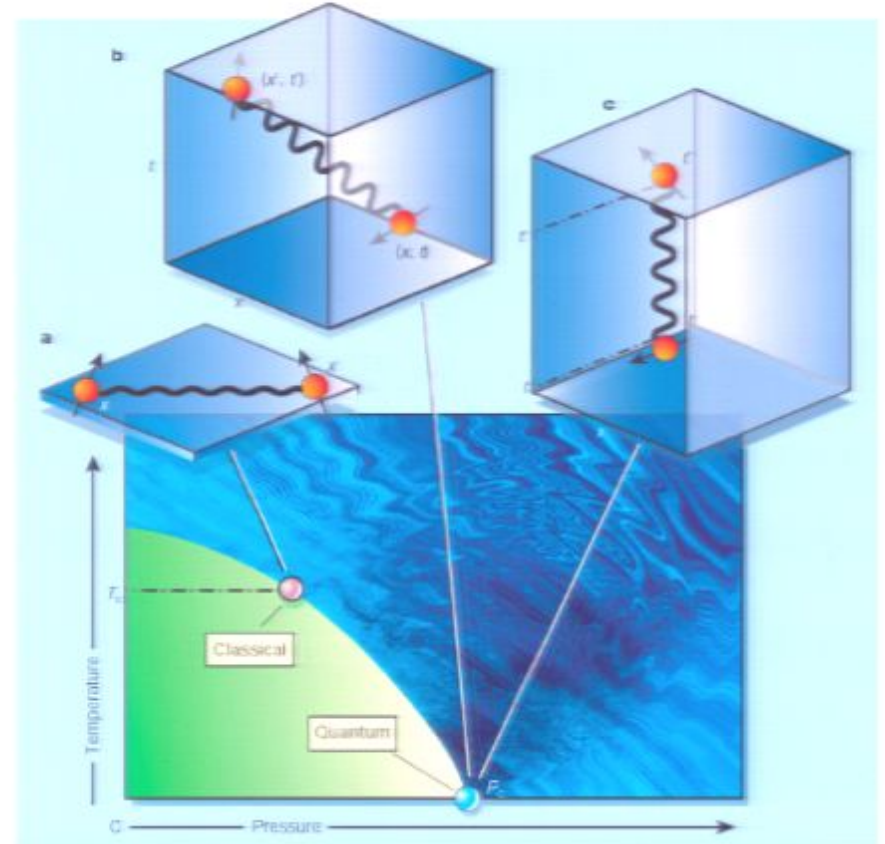


Lebanon

Rebecca Flint	Rutgers
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Maxim Dzero	Columbia/Rutgers
Jerome Rech,	ANL/Munich.
Eran Lebanon,	Israel.
Indranil Paul,	CNRS, Grenoble
Lucia Palova	Rutgers
Premi Chandra	Rutgers
Gergely Zarand	Budapest
Olivier Parcollet	SpHT Paris.
Andy Schofield	Birmingham
Qimiao Si	Rice, Houston
Catherine Pepin	SpHT Paris.
Almut Schroeder	Kent State
Gabriel Aeppli	LCN
Hilbert v. Lohneysen	Konlsruhe

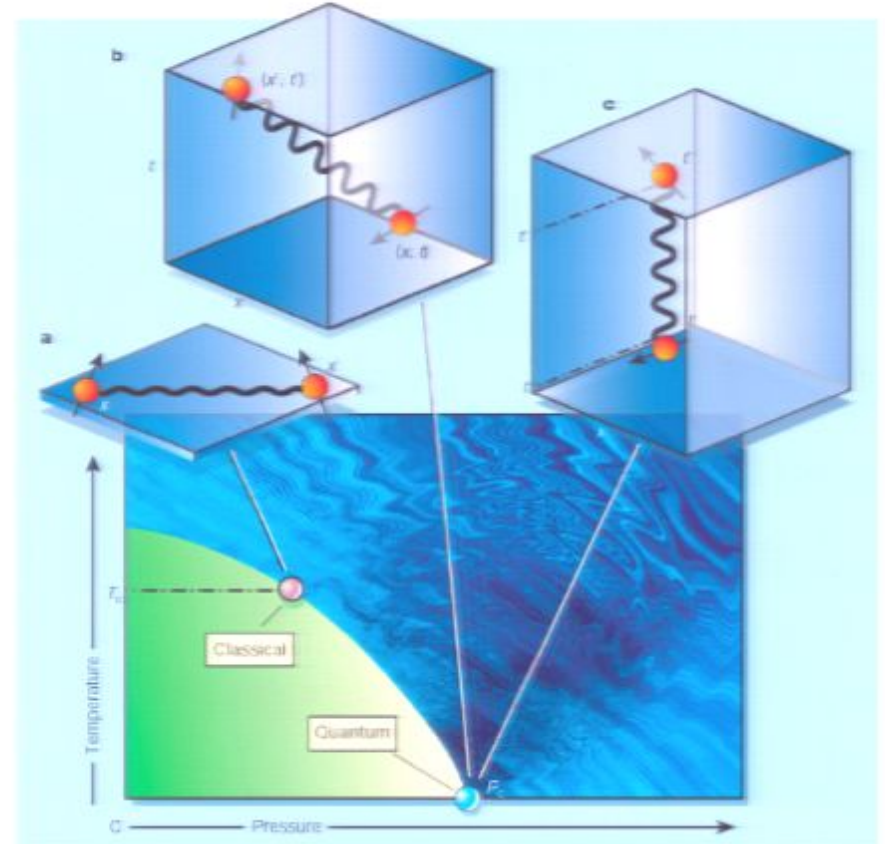


Conclusions



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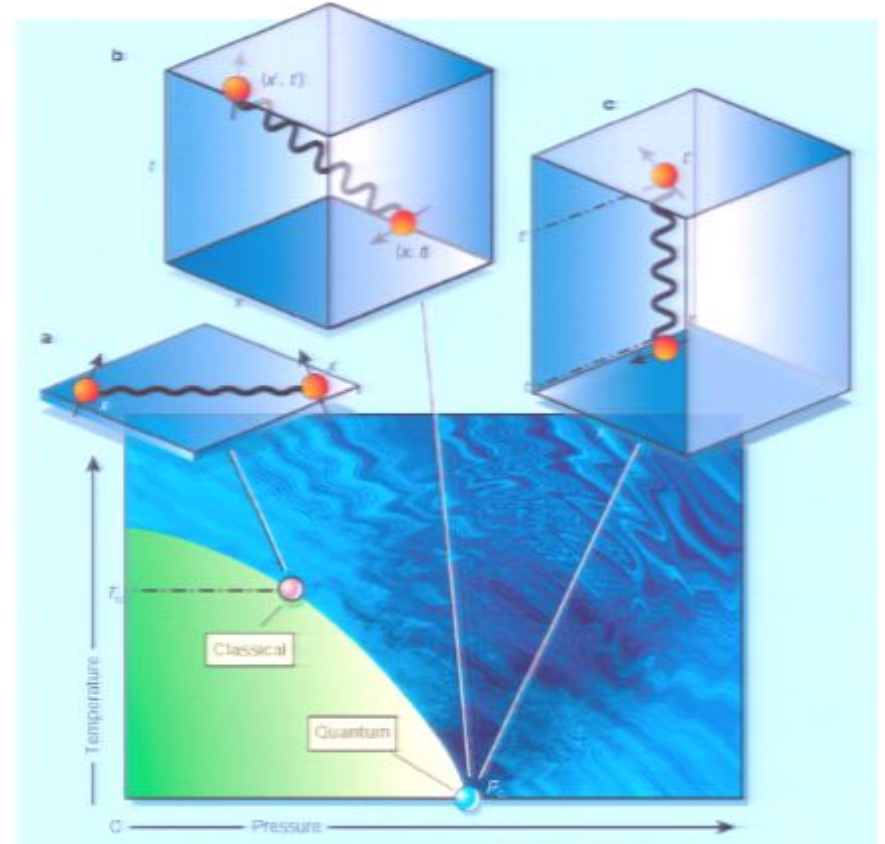
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we are just at the beginning
of the quantum era.



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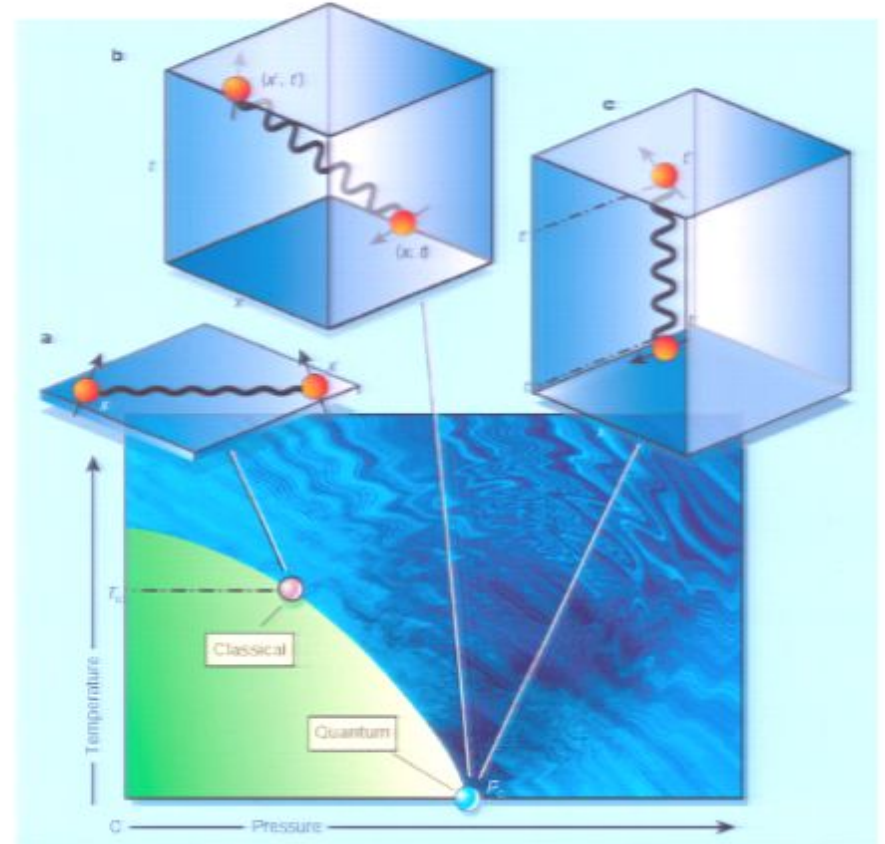


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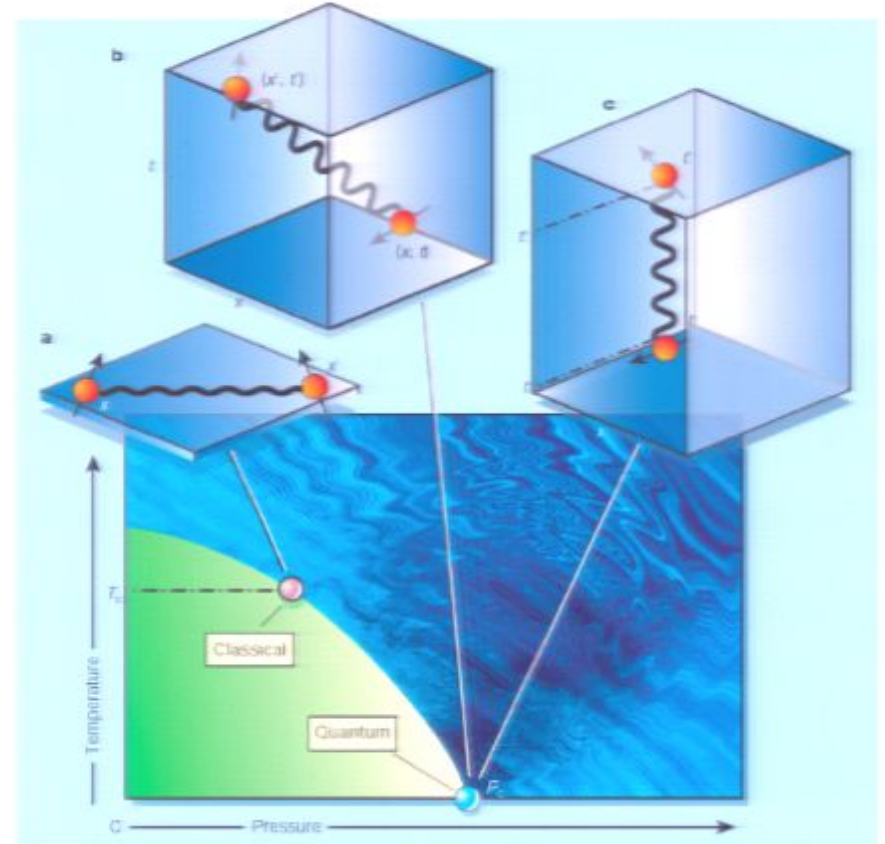
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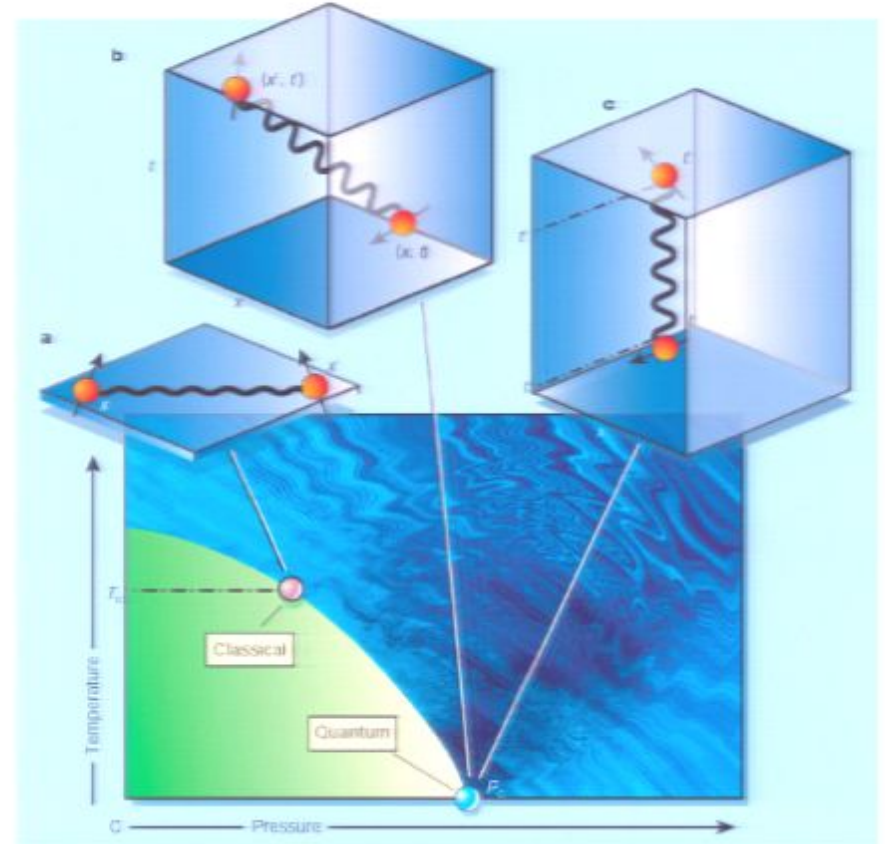


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New rule in material physics: avoided criticality.

New phases develop in order to avoid the
singular quantum critical point.