

Title: Introduction to the Bosonic String

Date: Feb 20, 2009 10:00 AM

URL: <http://pirsa.org/09020012>

Abstract: This course provides a thorough introduction to the bosonic string based on the Polyakov path integral and conformal field theory. We introduce central ideas of string theory, the tools of conformal field theory, the Polyakov path integral, and the covariant quantization of the string. We discuss string interactions and cover the tree-level and one loop amplitudes. More advanced topics such as T-duality and D-branes will be taught as part of the course. The course is geared for M.Sc. and Ph.D. students enrolled in Collaborative Ph.D. Program in Theoretical Physics. Required previous course work: Quantum Field Theory (AM516 or equivalent). The course evaluation will be based on regular problem sets that will be handed in during the term. The primary text is the book: 'String theory. Vol. 1: An introduction to the bosonic string. J. Polchinski (Santa Barbara, KITP) . 1998. 402pp. Cambridge, UK: Univ. Pr. (1998) 402 p.' All interested students should contact Alex Buchel at abuchel@uwo.ca as soon as possible.

CFT (continuation).

CFT (continuation)

$$-\delta A(\epsilon) - \frac{\epsilon}{2\pi i} \int_R d^2\sigma \sqrt{g} \nabla_\mu J^\mu \cdot A(\omega) = 0$$

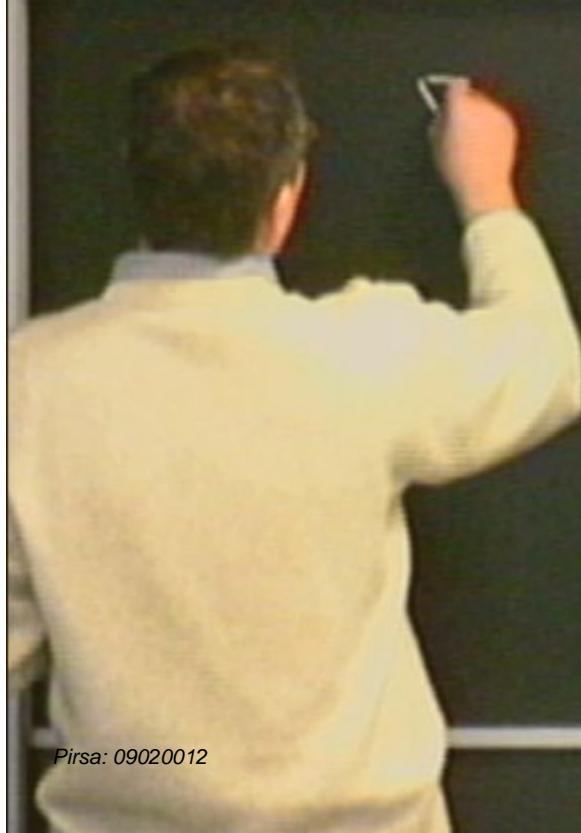


CFT (continuation).

$$-\delta A(\zeta_0) - \frac{e}{2\pi i} \int_R d^2\sigma \sqrt{g} \nabla_\mu J^\mu \cdot A(\zeta) = 0$$

ζ_0 must be inside R

$$G_0 \text{ must be inside } R$$
$$\int_{\mathcal{K}} J^a F g \nabla_a J^a A(G_0) = \sum \delta A(G_0)$$



σ_0 must be inside D

$$\int\limits_K J(\sigma) \nabla_\sigma J^* A(\sigma_0) = \delta A(\sigma_0)$$

$$\nabla_\sigma J^* A(\sigma) = \frac{2\pi}{\epsilon} \frac{1}{r-g} \delta^2(\sigma - \sigma_0) \nabla A(r).$$

σ_0 must be inside D

$$\int_{\mathbb{R}^n} \delta(\sigma - \sigma_0) \nabla_{\sigma} \left[\int_{\mathbb{R}^n} \delta(\sigma - \sigma_0) A(\sigma_0) \right] = \delta A(\sigma_0)$$

$$\nabla_{\sigma} \int_{\mathbb{R}^n} \delta(\sigma - \sigma_0) A(\sigma_0) = \frac{2\pi}{\gamma} \frac{1}{\sqrt{1-\gamma}} \delta^2(\sigma - \sigma_0) \nabla A(\sigma).$$

σ_0 must be inside D

$$\int_{\mathbb{C}} \delta \sigma \cdot \nabla_{\bar{\sigma}} \left[\int_{\mathbb{D}} \delta A(\sigma_0) \right] = - \oint_{\partial D} \delta A(\sigma_0)$$

$$\nabla_{\sigma} \Im^* A(\sigma_0) = \frac{2\pi i}{\varepsilon} \frac{1}{r-g} \delta^2(\sigma - \sigma_0) \cdot \nabla A(r).$$

$$\int_{\partial D} n_{\sigma} \Im^* A(\sigma_0) = \frac{2\pi}{\varepsilon} \delta A(\sigma_0)$$

$$-\nabla A(\sigma_0) + \frac{i\pi}{2n} \int_R d^2\sigma \bar{f}_j \nabla_j \bar{J}^a A(a) = 0$$

σ_0 must be inside R

$$\int_R d^2\sigma \bar{f}_j \nabla_j \left[\bar{J}^a(\sigma) A(a) \right] = -\frac{2\pi}{\xi} \delta A(\sigma_0) \quad \left\{ \langle \bar{\psi}_0 \bar{J}^a \rangle = 0 \right.$$

$$\int_{\partial R} n_a \bar{J}^a \cdot A(a) = \frac{2\pi}{\xi} \delta A(\sigma_0)$$

$$\oint_{\partial R} \left(i_z dz - i_{\bar{z}} d\bar{z} \right) A(z_0, \bar{z}_0) = \frac{2\pi}{\epsilon} \delta A(z_0, \bar{z}_0)$$

$$\oint_{\partial R} \left(j_z dz - i \bar{z} d\bar{z} \right) A(z_0, \bar{z}_0) = \frac{2\pi}{\xi} \delta A(z_0, \bar{z}_0)$$

\Rightarrow LHS simplifies when,

$$j_z \text{ is holomorphic} \quad \bar{\partial} j_z = 0$$

$$\oint_{\partial R} \left(j_z dz - i \bar{z} d\bar{z} \right) A(z_0, \bar{z}_0) = \frac{2\pi}{\epsilon} \delta_A(z_0, \bar{z}_0)$$

\Rightarrow LHS simplifies when.

$$j_z \text{ is holomorphic } \bar{\partial} j_z = 0$$

$$j_{\bar{z}} \text{ is antiholomorphic } \partial j_{\bar{z}} = 0$$

$$\oint_{\partial R} \left(j_z dz - i \bar{z} d\bar{z} \right) A(z_0, \bar{z}_0) = \frac{2\pi}{\epsilon} \delta A(z_0, \bar{z}_0)$$

\Rightarrow LHS simplifies when.

$$2\pi i \left[\operatorname{Re} j_z A(z_0, \bar{z}_0) \right]$$

j_z — is holomorphic $\bar{\partial} j_z = 0$

$j_{\bar{z}}$ — is antiholomorphic $\partial j_{\bar{z}} = 0$

$$\oint_{\partial R} (j_z dz - i \bar{z} d\bar{z}) A(z_0, \bar{z}_0) = \frac{2\pi}{\epsilon} \delta_A(z_0, \bar{z}_0)$$

\Rightarrow LHS simplifies when.

$$j_z \text{ is holomorphic } \bar{\partial} j_z = 0$$

$$j_{\bar{z}} \text{ is antiholomorphic } \partial j_{\bar{z}} = 0$$

$$2\pi i \left[\text{Res } j_z A(z, \bar{z}_0) \right]$$

$$+ 2\pi i \left[\text{Res } j_{\bar{z}} A \right]$$

$$\oint_{\partial R} (j_z dz - i_{\bar{z}} d\bar{z}) A(z_0, \bar{z}_0) = \frac{2\pi}{\epsilon} \delta A(z_0, \bar{z}_0)$$

\Rightarrow LHS simplifies when.

$$j_z \text{ is holomorphic } \bar{\partial} j_z = 0$$

$$j_{\bar{z}} \text{ is antiholomorphic } \partial j_{\bar{z}} = 0$$

$$2\pi i \left[\text{Res } j_z A(z, \bar{z}_0) \right]$$

$$+ 2\pi i \left[\text{Res } j_{\bar{z}} A \right]$$

$$= \frac{2\pi}{\epsilon} \delta A(z_0, \bar{z}_0)$$

$j_z \sim i s$ holomorphic $\partial j_z = 0$
 $j_{\bar{z}} \sim i s$ antiholomorphic $\partial j_{\bar{z}} = 0$

$$+ 2\pi i [Res j_{\bar{z}} A] \\ = \frac{i\pi}{\epsilon} \delta A(z_0, \bar{z}_0)$$

$$f(z)$$
$$Res f(z) = \alpha \beta f$$

$\int_C f(z) dz$ is holomorphic $\Rightarrow \int_C f(z) dz = 0$
 $\int_{\bar{C}} f(z) dz$ is antiholomorphic $\Rightarrow \int_{\bar{C}} f(z) dz = 0$

$$+ 2\pi i [\operatorname{Res} \int_C A] \\ = \frac{2\pi i}{\varepsilon} \delta A(z_0, \bar{z}_0)$$

$f(z)$

$$\operatorname{Res} f(z) = \text{coeff of } \frac{1}{z}$$

$$f(z) = \sum_{-\infty}^{+\infty} \frac{p_n}{z^{n+1}} \Rightarrow \operatorname{Res} f(z) = p_1$$

$$\begin{array}{l|l}
\begin{array}{l}
j_z \sim i s \text{ holomorphic} \quad \partial j_z = 0 \\
j_{\bar{z}} \sim i s \text{ anti-holomorphic} \quad \partial j_{\bar{z}} = 0
\end{array} & + 2\pi i \left[\operatorname{Res}_{j_z} A \right] \\
& = \frac{2\pi i}{\varepsilon} \delta A(z_0, \bar{z}_0)
\end{array}$$

$f(z)$

$\operatorname{Res} f(z) = \text{coeff of } \frac{1}{z}$

$$f(z) = \sum_{-\infty}^{+\infty} \frac{p_n}{z^n} \Rightarrow \operatorname{Res} f(z) = p_1$$

$$\operatorname{Res}_{z \rightarrow z_0} j_z A(z_0, \bar{z}_0) + \overline{\operatorname{Res}_{\bar{z} \rightarrow \bar{z}_0} j_{\bar{z}} A(z_0, \bar{z})} = \sum_{i \in \mathcal{E}} \delta A(z_0, \bar{z}_0)$$

$\gamma \alpha x = 0.$

Example 1

$$\delta x^m = \varepsilon a^m \rightarrow \text{action is invariant}$$

Example 1 $\delta x^m = \epsilon a^m g(r) \rightarrow$ action is invariant

$$S = \frac{1}{2\pi d} \int d^2 r (\partial x)^2$$

$$\delta S = \frac{\epsilon a^m}{2\pi d} \int d^2 r$$



$$S = \frac{1}{4\pi r} \int d^2r (\partial x)^2$$

$$\delta S = \frac{q\bar{q}}{2\pi r} \int d^2r \partial^a X^m \partial_a \beta$$

$$\delta S = \frac{ie}{2\pi d} \int d^2\sigma \ j^a X^{\mu} \partial_a P = -\bar{w}$$

$$[d\tau] e^{-S[\tau]} = [d\tau] e^{-S[\tau]} \left[1 + \frac{i\varepsilon}{2\pi} \int d^2\sigma \ j^a \nabla_a P + O(\varepsilon) \right]$$

$$\delta S = \frac{ie}{2\pi d} \int d^2\sigma \partial^a X^m \partial_a P = -\bar{\psi}$$

$$[d\tau] e^{-S[\tau]} = [d\tau] e^{-S[\tau]} \left[1 + \frac{i\varepsilon}{2\pi} \int d^2\sigma j^a \nabla_a P + O(\varepsilon) \right]$$

$$\frac{e^{i\omega t}}{\omega} = \frac{i\varepsilon}{2\pi}$$

$$SS = \frac{q_0}{2\pi d} \int d^2\sigma \ j^a X^{ab} \partial_b P$$

$$[d\varphi] e^{-S[\varphi]} = [d\tau] e^{-S[\tau]} \left[1 + \frac{i\varepsilon}{2\pi} \int d^2\sigma \ j^a \nabla_a P + O(\varepsilon) \right]$$

$$\vec{J} \cdot \vec{X} \frac{\varepsilon q_m}{2\pi d} = - \frac{i\varepsilon}{2\pi} J^a$$

$$S = \frac{q_m}{2\pi d} \int d^2\sigma \partial^a X^m \partial_a P$$

$$\langle d\varphi \rangle e^{-S[\varphi]} = \langle d\varphi \rangle e^{-S[\varphi]} \left[1 + \frac{i\varepsilon}{2\pi} \int d^2\sigma \partial^a \nabla_a P + O(\varepsilon) \right]$$

$$\partial_X \frac{\varepsilon q_m}{2\pi d} = -\frac{i\varepsilon}{2\pi} J^a \Rightarrow T_q = q_m T_q^m$$

$$T_q^m = + \frac{i}{2\pi} \partial_a X^m$$

$$\delta S = \frac{e^S}{2\pi d} \int d^2\sigma \ j^\alpha X^\mu \partial_\alpha P$$

$$[d\varphi] e^{-S[\varphi]} = [d\tau] e^{-S[\tau]} \left[1 + \frac{i\varepsilon}{2\pi} \int d^2\sigma \ j^\alpha \nabla_\alpha P + O(\varepsilon) \right]$$

$$j_\lambda \frac{\varepsilon q_m}{2\pi d} = -\frac{i\varepsilon}{2\pi} J^\lambda \Rightarrow T_q^\lambda = q_m T^\mu q$$

$$\boxed{T_0^\mu = + \frac{i}{2\pi} \partial_\alpha X^\mu} \Rightarrow \begin{aligned} T_2^\mu &= \frac{i}{2\pi} \partial X^\mu \\ \tilde{T}_2^\mu &= \frac{i}{2\pi} \bar{\partial} X^\mu \end{aligned}$$

$$SS = \frac{e^{-S}}{2\pi d} \int d^2\sigma \partial^a X^m \partial_a P$$

$$[d\varphi] e^{-S[\varphi]} = [d\tau] e^{-S[\tau]} \left[1 + \frac{i\varepsilon}{2\pi} \int d^2\sigma \partial^a \nabla_a P + O(\varepsilon) \right]$$

$$\partial_X \frac{\varepsilon q_m}{2\pi d} = -\frac{i\varepsilon}{2\pi} J^a \Rightarrow T_q^a = q_m T_q^a \quad \text{holomorphic}$$

$$\boxed{T_q^a = +\frac{i}{2\pi} \partial_a X^m} \Rightarrow T_2^a = \frac{i}{2\pi} \partial_a X^m$$

$$T_{\bar{2}}^a = \frac{i}{2\pi} \bar{\partial}_a X^m \quad \text{antiholomorphic}$$

$$A = :e^{ikx_{\text{continuation}}}:$$

$$J_2 : e^{ikx} :$$



$$A = :e^{ikx}:$$

$$J_2 : e^{ikx(z, \bar{z}_0)} = : \bar{J}_2 e^{ikx} : + \text{cross-correlations.}$$

T
can be singular as $z \rightarrow z_0$

$J_2 = \text{constant} + T_2 \epsilon + \text{(Cross-contracting)}$
 T
can be singular as $\epsilon \rightarrow 0$.

$\Rightarrow \underline{\text{A rule}}:$ the most singular in $\frac{1}{\epsilon - \epsilon_0}$ term
comes from $\omega_0 \mapsto$ Cross-contracting

$$T^a A(\sigma_0) = \frac{2\pi}{i\varepsilon} \delta A(\sigma_0)$$

\Rightarrow A rule. the most singular is
comes from most can be singular if $z \rightarrow z_0$
 $z - z_0$, then cross-contract. 0.

(5) $\frac{d}{dx} : \mathcal{D}X : : C^{\text{left}} \rightarrow \mathcal{D}(X) \text{ (SAL))}$

\Rightarrow A rule: the most singular in $\frac{1}{z-z_0}$, turn
comes from most cross-contractions

(5) $\frac{i}{d} : \text{cikx} : C ; \tilde{p} \approx \frac{1}{d!} \left(-\frac{1}{2} \lambda' \ln |z-z_0|^2 \right)$

only singular pieces.

\Rightarrow A rule: the most singular in $\frac{1}{z-z_0}$, turn can be singular if $z \rightarrow z_0$.
comes from most cross-contractions.

(5) $\frac{i}{d} : \partial x^i : = C^{ikx} ; \tilde{P} \sim \frac{1}{d!} \left(-\frac{1}{2} \omega' \right) \ln |z - z_0|^2).$
only singular if $C^{ikx} : \tilde{P}$ is.

\Rightarrow A rule: the most singular in $\frac{1}{z-z_0}$, turn
 comes from most cross-contracting.

$$\textcircled{S} \quad \frac{\partial}{\partial x^\nu} : e^{ikx} : \sim \frac{1}{\omega^1} \left(-\frac{1}{2} \omega^1 \ln |z-z_0|^2 \right).$$

only singular pieces

$$= \frac{\omega^m}{2(z-z_0)} : e^{ikx} :$$

\Rightarrow A rule: the most singular in $\frac{1}{z-z_0}$, turn
 comes from most cross-contractions
 can be singular as $r \rightarrow r_0$

$$\textcircled{5} \quad \frac{i}{\lambda'} : \partial x^u : e^{ikx} ; \tilde{\rho} \sim \frac{1}{\lambda'} \left(-\frac{1}{2} \lambda' \ln |z-z_0|^2 \right).$$

$\frac{\delta}{\delta x^v}$ (only singular pieces) $i k^u : e^{ikx} :$

$$\frac{1}{\lambda'} : \bar{\partial} x^u : e^{ikx} = \frac{k^u}{2(z-z_0)} : e^{ikx} :$$

$$= \frac{k^u}{2(\bar{z}-\bar{z}_0)} : e^{ikx} .$$

$$EHS = \left(\frac{k^u}{2} + \frac{k^l}{2} \right) \cdot A$$



$$EHS = \left(\frac{k^w}{2} + \frac{k^u}{2} \right) A = k^w A =$$



$$EHS = \left(\frac{k^u}{2} + \frac{k^m}{2}\right) \cdot A = k^u A = \frac{1}{i\varepsilon} \delta A$$

$$\rightarrow \delta A = ik^u \varepsilon A$$

from
Ward identity

$$e^{ikx}$$



$$EHS = \left(\frac{k^u}{2} + \frac{k^v}{2}\right) A = k^u A = \frac{1}{i\varepsilon} \delta A$$

$$\rightarrow \delta A = ik^u \varepsilon A$$

from

Ward identity

$$A = e^{ikx} \rightarrow e^{ikx + i k^u \varepsilon}$$

$$X^u \rightarrow X^u + \varepsilon$$

Example 2 $\delta\sigma^a$

Example 2

$$\delta\sigma^a = \epsilon U^a$$

$\uparrow U^a$ is a constant.

Example 2 $\delta\sigma^a = \epsilon U^a$

$$\sigma^a \rightarrow \sigma^{a'} = \sigma^a + \epsilon U^a$$

$$X^m(\sigma') = X^m(\sigma)$$

Example 2

$$\delta\sigma^a = \varepsilon U^a$$

$\uparrow U^a$ is a constant

$$\sigma^a \rightarrow \sigma^{a'} = \sigma^a + \varepsilon U^a$$

$$X^u(s') = X^u(s)$$

$$\delta X^u(s) = X^u(s') - X^u(s)$$



Example 2

$$\delta\sigma^n = \epsilon U^n$$

$$\sigma^n \rightarrow \sigma^{n'} = \sigma^n + \epsilon U^n$$

$\uparrow U^n$ is a constant

$$X''(s') = X''(s) \Rightarrow X''(s + \delta s)$$

$$\delta X''(s) = X''(s) - X''(s)$$



Example 2

$$\delta\sigma^a = \varepsilon U^a \quad T_0^a \text{ is a constant.} \quad X'(s) = X(s-s_0)$$

$$\sigma^a \rightarrow \sigma^{a'} = \sigma^a + \varepsilon U^a$$

$$X''(s') = X''(s) \Rightarrow X''(s + \delta s) = X''(s) \Rightarrow$$

$$\delta X''(s) = X''(s) - X''(s)$$

Example 2

$$\delta\sigma^a = \epsilon U^a$$

\uparrow
 U^a is a constant.

$$\sigma^a \rightarrow \sigma^{a'} = \sigma^a + \epsilon U^a$$

$$X'(s) = X(s - \epsilon)$$

$$X''(s') = X''(s) \Rightarrow X''(\underbrace{s + \frac{\delta s}{s'}}_{s'}) = X''(s) \Rightarrow$$

$$\delta X''(s) = X''(s') - X''(s)$$

Example 2 $\delta\sigma^a = \epsilon U^a$

$$\sigma^a \rightarrow \sigma^{a'} = \sigma^a + \epsilon U^a$$

$$X^m'(\sigma') = X^m(\sigma) \Rightarrow X^m'(\sigma + \frac{\delta\sigma}{\epsilon}) = X^m(\sigma) \Rightarrow$$

$$\delta X^m(\sigma) = X^m'(\sigma) - X^m(\sigma)$$

$$= X^m(\sigma - \epsilon U^a) - X^m(\sigma) = -\epsilon U^a \partial_a X^m$$

Example 2

$$\sigma^a = \xi U^a \quad \uparrow U^a \text{ is a constant.} \quad X'(s) = X(s-t_0)$$

$$s \rightarrow \sigma^a = s + \xi U^a$$

$$X''(s') = X''(s) \Rightarrow X''(\underbrace{s + \xi U^a}_{s'}) = X''(s) =$$

$$\delta X(s) = X''(s) - X''(s)$$

$$= X''(s - \xi U^a) - X''(s) = \boxed{-\xi U^a \partial_s X'' = \delta X}$$

$$\Im S = \delta \left[\frac{1}{4\pi k'} \int d^2 \sigma \partial^a X^m \partial_a X_m \right]$$

$$= \frac{1}{2\pi \alpha'} \int d^2 \sigma \partial^a X^m \partial_a X_m$$

$$\delta S = \delta \left[\frac{1}{4\pi k'} \int d^2\theta \epsilon^{ij} \partial^0 X^i \partial_n X_j \right]$$

$$= \frac{1}{2\pi \alpha'} \int d^2\sigma \epsilon^{ij} X^i \partial_n \delta X_j$$

$$\epsilon(\vec{r}-\vec{s})$$

$$\delta S = \delta \left[\frac{1}{4\pi G} \int d^4x \partial^\mu X^\nu \partial_\mu X_\nu \right]$$

$$= \frac{1}{2\pi\alpha'} \int d^2\sigma \partial^b X^\mu \partial_b \delta X_\mu \Big] =$$

$$\delta X_\mu = - \mathcal{E}(\sigma) \cup^a \partial_a X_\mu$$

$$= \frac{k^a}{2(\vec{r} - \vec{r}_0)} \cdot \mathcal{E}^{i \times}$$

$$= \frac{1}{2\pi\alpha'} \int d^2\sigma \partial^b X^a \partial_b \delta X_a = -\frac{i}{2\pi r} \int d^2\zeta \partial^a X^b \partial_a \varepsilon \partial_b X_b$$

$$\delta X_a = -\underline{\varepsilon(\zeta)} \underline{\partial^a} \underline{\partial_a X_a}$$

$$\begin{aligned} \frac{1}{\lambda^2} : \bar{\partial} X^a : e^{i\kappa X} &= \frac{\kappa}{2(z-\bar{z}_0)} : e^{i\kappa X} : \\ &= \frac{\kappa}{2(\tilde{z}-\bar{z}_0)} : e^{i\kappa X} : \end{aligned}$$

$$= \frac{1}{2\pi\epsilon_0} \int d^2\sigma \partial^b X^a \partial_b \delta X_a = -\frac{1}{2\pi\epsilon_0} \int d^2\sigma \partial^a X^b \partial_b \delta X_a$$

$$\delta X_a = -\underline{\mathcal{E}(e) U^a \partial_a X_a}$$

$$-\frac{U^a}{2\pi\epsilon_0} \int d^2\sigma \partial^a X^b \partial_b \delta X_a$$

$$\frac{1}{\lambda}; \bar{\partial} X^a; e^{ikx} = \frac{K^a}{2(z - z_0)}; e^{ikx};$$

$$= \frac{K^a}{2(\bar{z} - \bar{z}_0)}; e^{ikx}.$$

$$\delta X_{\mu} = - \underline{\mathcal{E}}(\sigma) D^+ \partial_\sigma X_{\mu}$$

$$\left(\frac{U}{2\pi\sigma^4} \right) d^2\sigma \, D X \int d_\theta J_\theta \chi \, \mathcal{E}$$

$$\textcircled{(5)} \quad \frac{1}{\lambda} : \partial X^u : e^{ikx} ; \quad \nearrow \frac{1}{2\lambda} \left(-\frac{1}{2} \lambda' \right) \ln |z - z_0|^2 .$$

$$\frac{\delta}{\delta X^\nu} \quad \text{only singularities} \quad i F^u : e^{ikx} :$$

$$\begin{aligned} \frac{1}{\lambda} : \bar{\partial} X^u : e^{ikx} &= \frac{K^u}{2(z - z_0)} : e^{ikx} : \\ &= \frac{K^u}{2(\bar{z} - \bar{z}_0)} : e^{ikx} : \end{aligned}$$

$$\delta \chi_m = -\underline{\mathcal{E}(\sigma)} \circ \partial_a \chi_m$$

$$\left(\frac{v^*}{2\pi\sigma^4} \right) d^2\sigma \partial^* X^a \partial_a \chi_m$$

$$-\delta S = \frac{i}{2\pi} \int d^2\sigma \sqrt{-g} j^a \partial_a \mathcal{E}$$

$$\delta X_u = - \underline{\varepsilon(\sigma)} \circ \underline{\partial_\alpha} \underline{X_{u_1}} - \left(\frac{U}{2\pi\sigma^4} \right) d^2\sigma \circ \underline{\partial X}^\alpha \underline{\partial_\alpha} \underline{X}^\beta \underline{\varepsilon}$$

$$-\delta S = \frac{i e}{2\pi} \int d^2\sigma \sqrt{g} \ j^\alpha \underline{\partial_\alpha \varepsilon}$$

$$\textcircled{2}: \partial^\alpha X^\mu \partial_\alpha \partial_\beta X^\nu \cdot \varepsilon = \frac{1}{2} \partial_\nu [\partial^\alpha X^\mu \partial_\alpha X_\mu]$$

$$\partial X^a = - \underline{\varepsilon}(\sigma) \partial^\alpha \partial_\alpha X^a$$

$$\left(\frac{e^{\sigma}}{2\pi\omega} \right) d^2\sigma \partial^\alpha \partial_\alpha X^a \underline{\varepsilon}$$

$$-\delta S = \frac{i}{2\pi} \int d^2\sigma \sqrt{g} j^a \partial_a \underline{\varepsilon}$$

$$\textcircled{2}: \partial^\alpha X^a \partial_\alpha \partial_\beta X^a \cdot \varepsilon = \frac{1}{2} \partial_\beta [\partial^\alpha X^a \partial_\alpha X_a] \cdot \varepsilon$$
$$\hookrightarrow -\frac{1}{2} (\partial X)^2 \cdot \partial_\beta \varepsilon$$

$$T_v = \frac{v^b}{i\omega} \left(\gamma_a X^m \partial_b X_m - \frac{1}{2} \delta_{ab} (\partial X)^2 \right)$$

$$J_a = \frac{e^b}{i\omega} \left(\partial_a X^m \partial_b X_m - \frac{1}{2} \delta_{ab} (\partial X)^2 \right)$$

$$\Pi_{ab} = -\frac{1}{2} \left(\partial_a X^m \partial_b X_m - \frac{1}{2} \delta_{ab} (\partial X)^2 \right)$$

$$J_a = \frac{\partial^b}{\partial x^b} \left(\gamma_a X^m \partial_b X_m - \frac{1}{2} \delta_{ab} (\partial X)^2 \right)$$

$$\Pi_{ab} \stackrel{def}{=} -\frac{1}{2} \left(\gamma_a X^m \gamma_b X_m - \frac{1}{2} \delta_{ab} (\partial X)^2 \right)$$

$$J_a = -\frac{1}{i} \psi^\dagger \Gamma_{ab} = i \psi^b \Pi_{ab}$$

In 4CFT:

$$T_{ab} = -\frac{1}{2} : \partial_a X^c \partial_b X_c - \frac{1}{2} \delta_{ab} (\partial X)^c :$$

In 4CFT

$$T_{ab} = -\frac{1}{2} : \partial_a X^c \partial_b X_c - \frac{1}{2} \delta_{ab} (\partial X)^c :$$

Conformal inv

In 4CFT

$$T_{ab} = -\frac{1}{2} : \partial_a X^c \partial_b X_c - \frac{1}{2} \delta_{ab} (\partial X)^c :$$

Conformal inv

$$T_a^a = 0$$

In 4CFT

$$T_{ab} = -\frac{1}{2} : \partial_a X^\mu \partial_b X_\mu - \frac{1}{2} \delta_{ab} (\partial X)^\mu : \\ \Rightarrow \text{Conformal inv}$$

$$T_{a^q} = 0$$

$$g^{z\bar{z}} = 2, g^{\bar{z}\bar{z}} = 2$$

In CFT

$$T_{ab} = -\frac{1}{2} : \partial_a X^\mu \partial_b X_\mu - \frac{1}{2} \delta_{ab} (\partial X)^2 :$$

Conformal inv

$$T_{a\bar{a}} = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}\bar{z}}$$

$$g^{z\bar{z}} = 2, g^{\bar{z}\bar{z}} = 2$$



τ^a

T_{0^a} is a constant

$$\chi(g) = \chi(e^{-ig})$$

Only nonvanishing $T_{ab} : \{T_{\bar{z}\bar{z}}, T_{\bar{z}\bar{z}}\}$



$$\wedge (\omega - \epsilon \partial^a) \rightarrow (0) = \Gamma \epsilon \partial^a \partial_0 \wedge = 0 \wedge$$

Only nonvanishing T_{ab} : $\{T_{22}, T_{\bar{2}\bar{2}}\}$
⇒ use conservation of the current.



$$\wedge (\omega - \epsilon \omega^*) - X(\omega) = -\epsilon \omega \partial_{\omega} X = \omega \wedge$$

Only nonvanishing T_{ab} : $\{T_{22}, T_{\bar{2}\bar{2}}\}$
⇒ use conservation of the current.

$$\nabla_a^a T_{ab} = 0$$



$$\wedge (\mathcal{L} - \mathcal{E} J^a) - X (0) = -\mathcal{E} O_{0,1} X = 0 \wedge$$

Only nonvanishing T_{ab} : $\{T_{22}, T_{\bar{2}\bar{2}}\}$
⇒ use conservation of the current.

$$\nabla^a T_{a1} = 0 \Rightarrow \nabla^a T_{a1} = 0$$



$$\wedge (\omega - \epsilon \sigma^a) - X (0) = \Gamma \epsilon 0 \sigma_0 X = 0 \wedge$$

Only nonvanishing $T_{\alpha\bar{\beta}}$: $\{T_{22}, T_{\bar{2}\bar{2}}\}$
 \Rightarrow use conservation of the current.

$$\nabla^\mu T_{\alpha\bar{\beta}} = 0 \Rightarrow \nabla^\mu T_{\alpha\bar{\beta}} = 0$$

$\bar{\partial}$



$$\wedge (\mathcal{L} - \mathcal{E} J^\alpha) - X(\mathcal{O}) = \Gamma \mathcal{E} O \circ_{\mathcal{O}_0} X = \mathcal{O} X$$

Only nonvanishing T_{ab} : $\{T_{zz}, T_{\bar{z}\bar{z}}\}$

⇒ use conservation of the current.

$$\nabla_a^a T_{ab} = 0 \Rightarrow \nabla^a T_{ab} = 0$$

$$\bar{\partial} T_{zb} + \partial$$

$$\wedge (\omega - \epsilon \sigma^a) - X (0) = \Gamma \epsilon \sigma \sigma_0 X = \sigma \wedge$$

Only nonvanishing T_{ab} : $\{T_{zz}, T_{\bar{z}\bar{z}}\}$

\Rightarrow use conservation of the current.

$$\nabla_a T_{ab} = 0 \Rightarrow \nabla^a T_{ab} = 0$$

$$\bar{\partial} T_{zb} + \partial T_{\bar{z}b} = 0$$

$$\wedge (\sigma - \epsilon \omega^a) - X(\sigma) = -\epsilon \sigma \partial_0 X = \sigma \wedge$$

Only nonvanishing T_{ab} : $\{T_{zz}, T_{\bar{z}\bar{z}}\}$

\Rightarrow use conservation of the current.

$$\nabla_a T_{az} = 0 \Rightarrow \nabla^a T_{az} = 0$$

$$\bar{\partial} T_{z\bar{b}} + \partial T_{\bar{z}b} = 0 \quad (b=2) \Rightarrow \bar{\partial} T_{z\bar{z}} = 0$$

$$b=\bar{z} \Rightarrow \partial T_{\bar{z}\bar{z}} = 0$$

⇒ use conservation of the current.

$$\nabla_a^a J_{a\bar{a}} = 0 \Rightarrow \nabla_a^a T_{a\bar{a}L} = 0$$

$$\bar{\partial} T_{z\bar{z}L} + \partial \bar{T}_{\bar{z}\bar{z}L} = 0 \quad l=2 \Rightarrow \bar{\partial} T_{z\bar{z}} = 0$$

$$T_{z\bar{z}} = T(z)$$

$$l=2 \Rightarrow \partial T_{\bar{z}\bar{z}} = 0$$

$$\bar{\partial} T_{\bar{z}\bar{z}} = \frac{1}{2} \bar{\partial} X^a$$

anti holomorph.

⇒ use conservation of the current.

$$\nabla^a J_{a\bar{b}} = 0 \Rightarrow \nabla^a T_{a\bar{b}} = 0$$

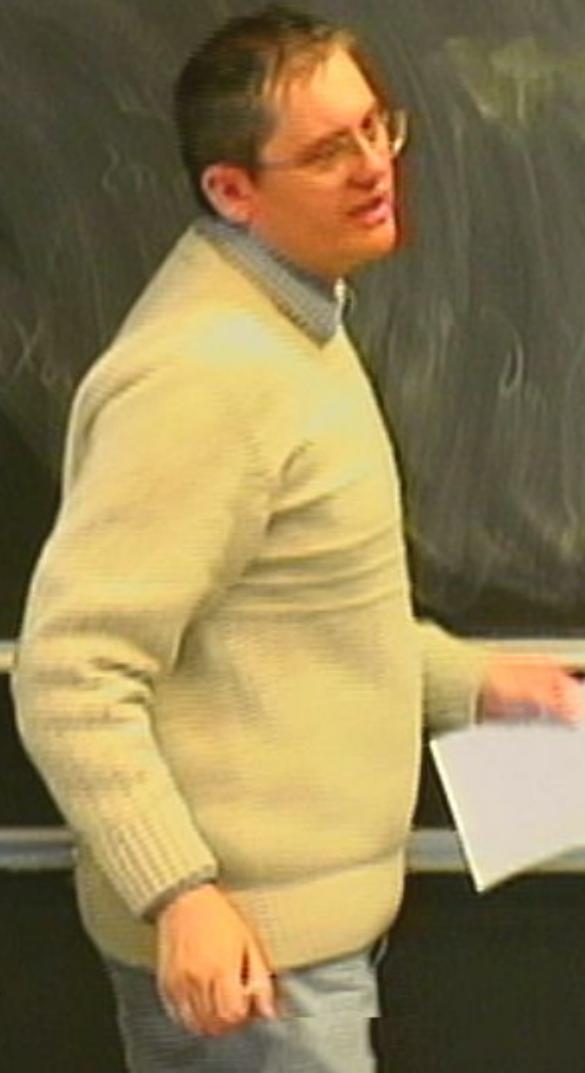
$$\bar{\partial} T_{z\bar{b}} + \partial \bar{T}_{\bar{z}b} = 0 \quad b=\bar{z} \Rightarrow \bar{\partial} T_{z\bar{z}} = 0$$

$$T_{z\bar{z}} = T(z) \quad \& \quad T_{\bar{z}\bar{z}} = \tilde{T}(\bar{z}) \quad b=\bar{z} \Rightarrow \bar{\partial} T_{\bar{z}\bar{z}} = 0$$

$$J_{\bar{z}}^a = \frac{1}{2\pi} \bar{\partial} X^a$$

anti holomorph.

$$T(z) = -\frac{1}{\lambda} : \partial X^a \partial X_a : Y$$

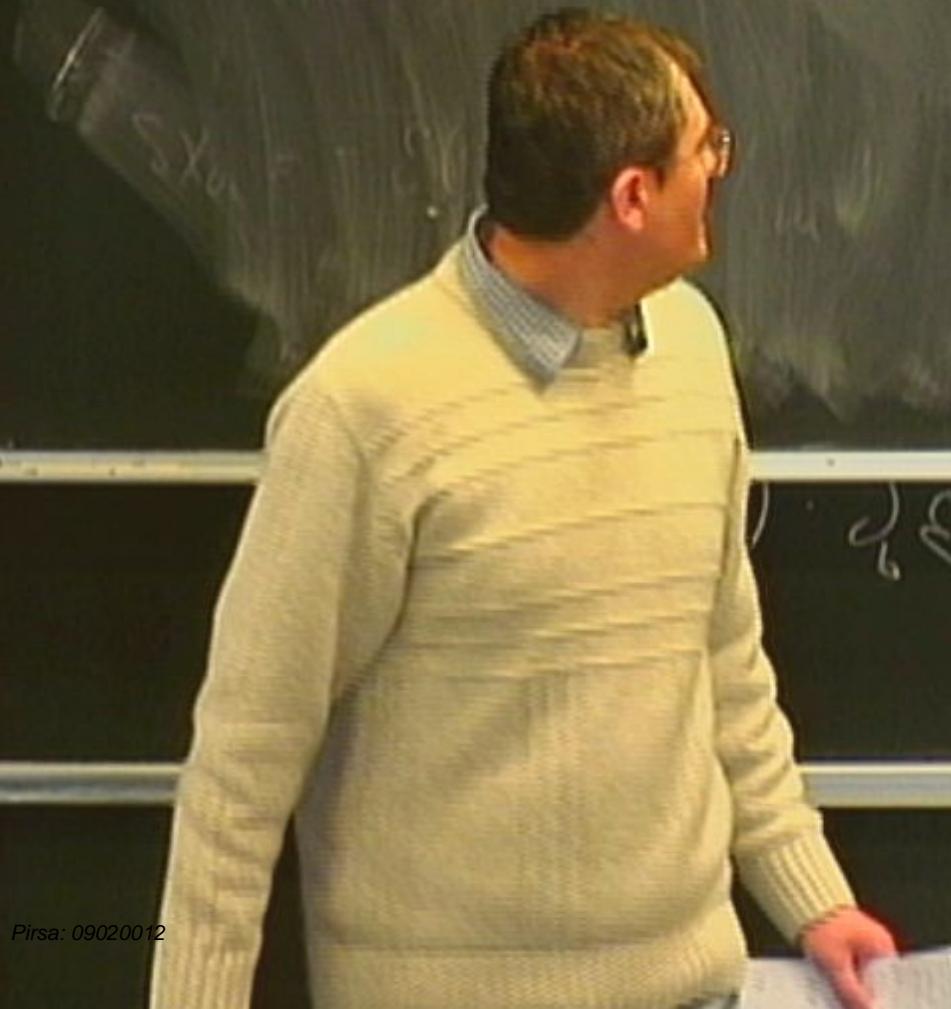


$$T(\tau) = -\frac{1}{\lambda} : \partial X^a \partial X_a : \quad ; \quad \tilde{T}(\xi) = -\frac{1}{\alpha} : \bar{\partial} X^a \bar{\partial} X^a :$$

$$T(z) = -\frac{1}{2} : \partial X^a \bar{\partial} X_a : \quad)$$

$$\tilde{T}(z) = -\frac{1}{2} : \bar{\partial} X^a \partial X_a :$$

⇒ A good exercise is to check that Ward identities
are satisfied for $\delta \phi^a = -\epsilon \eta \omega^a$



$$T(z) = -\frac{1}{2} : \partial X^a \partial X_a : \quad \tilde{T}(z) = -\frac{1}{2} : \bar{\partial} \bar{X}^a \bar{\partial} \bar{X}^a :$$

⇒ A good exercise is to check that Ward identities
satisfied for $\delta \phi^a = + \epsilon \partial \phi^a$



$$\partial \phi^a$$

$$\nabla(\tau) = -\frac{1}{2} : \partial X^a \partial X_a : \quad ; \quad \tilde{\nabla}(\xi) = -\frac{1}{2} : \bar{\partial} \bar{X}^a \bar{\partial} \bar{X}^a :$$

\Rightarrow A good exercise is to check that Ward identities are satisfied for $\delta g^\alpha = +\epsilon \alpha \omega^\alpha$

$$j = i \omega^\alpha \nabla_\alpha L$$

$$-\frac{1}{2} (\partial X)^a \cdot \partial_a \xi$$

$$\mathcal{T}(\tau) = -\frac{1}{2} : \partial X^a \partial X_a : ; \quad \tilde{\mathcal{T}}(\xi) = -\frac{1}{2} : \bar{\partial} X^a \bar{\partial} X^a :$$

⇒ A good exercise is to check that Ward identities
satisfied for $\delta \zeta^a = +\epsilon \theta \omega^a$

$$j = i \omega^a \mathcal{T}_{ab}$$

$$\bar{\partial} \mathcal{T} = 0$$

$$-\frac{1}{2} (\partial X)^a \cdot \partial_b \xi$$



⇒ A good exercise is to check that Ward identities
are satisfied for $\delta \phi^a = +\epsilon \theta \omega^a$

$$\left. \begin{aligned} j &= i \bar{\psi}^\alpha \Gamma_\alpha \psi \\ \bar{\partial} \Gamma &= 0 \end{aligned} \right\} \quad \bar{\psi}^\alpha \text{ was a constant.}$$



\Rightarrow A good exercise is to check that Ward identities
are satisfied for $\delta \phi^a = +\epsilon a \psi^a$

$\downarrow \begin{cases} j = i \psi^a \nabla_a \phi \end{cases}$ ψ^a was a constant.

$$\partial \bar{\nabla} = 0$$

generalization

$$j_z = i \psi(z) \nabla(z)$$

$$\tilde{j}_{\bar{z}} = i \bar{\psi}(\bar{z}) \bar{\nabla}(\bar{z})$$

$$\nabla \cdot \partial \chi \cdot \Sigma = \frac{1}{2} \partial_\nu \left(\partial^a \chi^a \partial_\mu \chi_{\mu} \right) \cdot \Sigma$$

$$+ (\partial \chi)^2 \cdot \partial_\nu \Sigma$$

$$\begin{aligned} & \left. \begin{aligned} J = i \langle O^a \bar{T}_{ab} \rangle \Big\} \quad O^a \text{ being a constant.} \\ \bar{\partial} T = 0 \end{aligned} \right\} \text{ general solution} \\ & \rightarrow J_z = i \langle O(z) \bar{T}(z) \rangle \quad \bar{J}_{\bar{z}} = i \langle \bar{O}(\bar{z}) \bar{T}(\bar{z}) \rangle \end{aligned}$$

②

$$X^\mu \cdot \varepsilon = \frac{1}{2} \partial_\nu \left[\partial^\mu X^\nu \partial_\lambda X_{\lambda} \right] \cdot \varepsilon$$

$$(\partial X)^2 \cdot \partial_\nu \varepsilon$$

generalization

$$\partial T = 0 \quad \rightarrow \quad J_z = i \langle \psi(z) | T(z) \rangle$$

$$\tilde{J}_z = i \langle \overline{\psi}(\bar{z}) | \tilde{T}(\bar{z}) \rangle$$

$$\nabla_a T^a = \bar{\delta} J_z + \partial T_{\bar{z}} = 0$$

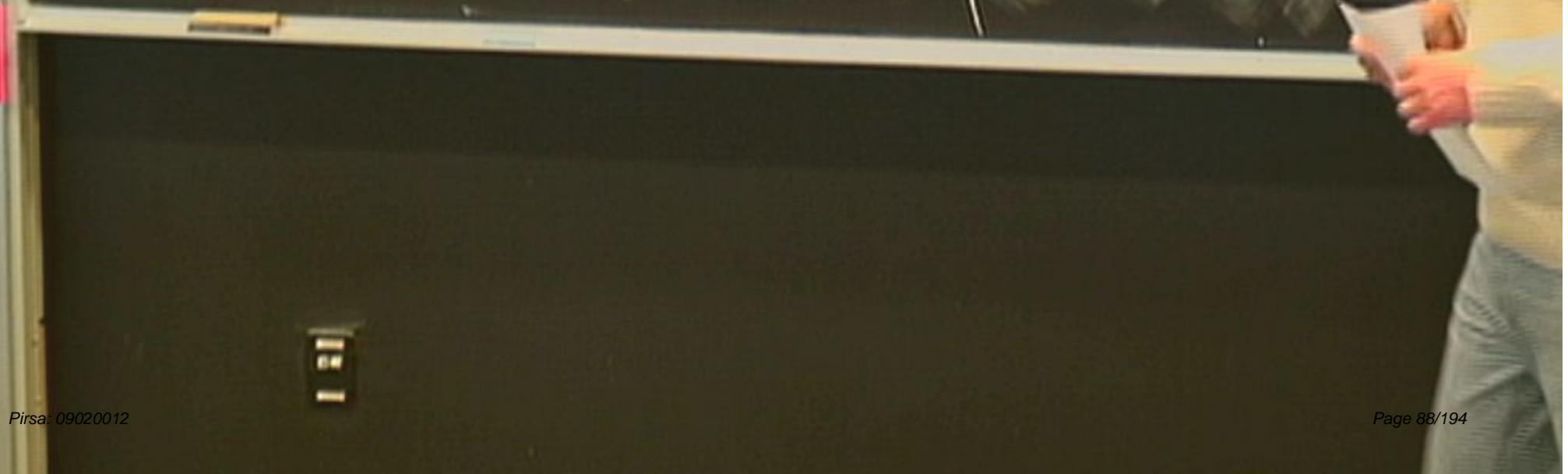


$$\begin{array}{l} \partial T = 0 \\ \text{generalization} \\ \rightarrow J_z = i \sigma(z) T(z) \quad \tilde{J}_{\bar{z}} = (\overline{\sigma(z)}) \tilde{T}(\bar{z}) \end{array}$$

$$\nabla_a T^a = \bar{\partial} J_z + \partial \tilde{J}_{\bar{z}} = 0$$

\Rightarrow The new currents are conserved provided
 $\sigma(z)$ is holomorphic

We want to find δX



We want to find δX due to symmetries $T_2 = 0$

2



We want to find δX due to symmetries $J_2 = 0 \text{ T}$

$$\frac{P_{ez}}{z \rightarrow z_0} j_z A + \frac{\widehat{P}_{es}}{\bar{z} \rightarrow \bar{z}_0} \tilde{j}_z A = \frac{1}{\epsilon} \delta A$$

We want to find δX due to symmetries $T_2 = 10T$

$$\frac{R_{xz}}{z \rightarrow z} j_z A + \frac{\widehat{R}_{xz}}{\bar{z} \rightarrow \bar{z}} \tilde{j}_z A = \frac{1}{i\varepsilon} (\delta A)$$

We want to find δX due to symmetries $J_2 = 0 \text{ T}$

$$\overset{P_{xz}}{\underset{z \rightarrow z_0}{\mathcal{J}_2}} A + \overset{\widetilde{P}_{xz}}{\underset{\bar{z} \rightarrow \bar{z}_0}{\mathcal{J}_2}} A = \frac{1}{i\varepsilon} \delta A$$

$$\Rightarrow A = X^m$$

$$\sum_{z \rightarrow z_0} j_z A + \sum_{\bar{z} \rightarrow \bar{z}_0} \tilde{j}_{\bar{z}} A = \frac{1}{i\varepsilon} \delta A$$

$$\Rightarrow A = X^m$$

$$j_z = i \partial^m T$$

$$\Rightarrow A = X$$

$$j_* = \cup T$$

$$T(r) : X^*(r_0, \bar{r}_0)$$

$$T_a = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}\bar{z}}$$

$$\Rightarrow A = X$$

$$j_\infty = i \partial \nabla$$

$$:\nabla(z) : : X''(z_0, \bar{z}_0) : = -\gamma \frac{1}{z} : \partial z \partial V(z) : X''(z_0, \bar{z}_0) :$$

$$T_{\bar{z}} = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}\bar{z}}$$

$$j_\mu = i \bar{\psi} \Gamma$$

$$\langle T(z) : X^\mu(z_0, \bar{z}_0) \rangle = -\frac{1}{2} \partial z \partial V(z) : X^\mu(z_0, \bar{z}_0) : \\ \sim -\left(\frac{1}{z}\right).$$

$$\Rightarrow \text{Conformal inv} \quad \delta_{ab} - \frac{1}{2} \partial_a \chi_{ab} - \frac{1}{2} \partial_b \chi_{ab} (\partial \chi) :$$

$$T_{a^q} = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}\bar{z}}$$

$$r_0 = 10^9$$

$$\begin{aligned} T(z) &:= \chi''(z_0, \bar{z}_0) = -\frac{1}{2} : \partial z \partial \chi(z) : \chi'''(z_0, \bar{z}_0) : \\ &\sim -\left(\frac{1}{2}\right) \cdot \left(-\frac{1}{2} \ln(r - r_0)^2\right) \end{aligned}$$

$$T_{ab} = -\frac{1}{2} : \partial_a \chi''' \partial_b \chi''' - \frac{1}{2} \delta_{ab} (\partial \chi)' :$$

Conformal inv

$$g^{z\bar{z}} = 2, g^{\bar{z}\bar{z}} = 2$$

$$T_{a^q} = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}\bar{z}}$$

$$T(z) := \chi''(z_0, \bar{z}_0) := -\frac{1}{2} \partial z \partial \chi(z) = \chi''(z_0, \bar{z}_0) \\ \sim -\left(\frac{1}{2}\right) \cdot \left(-\frac{1}{2} \ln(z - z_0)^2\right)$$

$$T_{ab} = -\frac{1}{2} : \partial_a \chi'' \partial_b \chi'' - \frac{1}{2} \delta_{ab} (\partial \chi)^2 : \\ \Rightarrow \text{Conformal inv}$$

$$g^{\bar{z}\bar{z}} = 2, g^{\bar{z}, \bar{z}} = 2 \\ T_{\bar{z}\bar{z}} = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}z}$$

$$T(z) := X''(z_0, \bar{z}_0) := -\frac{1}{2} : \partial z \partial \bar{z} : X''(z_0, \bar{z}_0) :=$$

$$\sim -\left(\frac{1}{2}\right) \cdot \left(-\frac{\pi^2}{2} \ln |z - z_0|^2\right) \cdot 2$$

$$T_{ab} = -\frac{1}{2} : \partial_a X'' \partial_b X'' - \frac{1}{2} \delta_{ab} (\partial X)^2 :$$

\Rightarrow Conformal inv

$$g^{z\bar{z}} = 2, g^{\bar{z}\bar{z}} = 2$$

$$T_{a\bar{a}} = 0 \Rightarrow T_{z\bar{z}} = 0 = T_{\bar{z}\bar{z}}$$

$$T(z) := X(z_0, \bar{z}_0) := -\frac{1}{2} : \partial X \partial V(z) : X''(z_0, \bar{z}_0) :$$

$$\sim -\left(\frac{1}{2}\right) \cdot \left(-\frac{1}{2} \ln((z-z_0)^2)\right) \cdot 2 : \partial X(z_0) :$$

$$T_{ab} = - \partial_a X'' \partial_b X'' - \frac{1}{2} \delta_{ab} (\partial X)^2$$

Confinement

$$T_q^q$$

$$g^{z\bar{z}} = 2, g^{\bar{z}\bar{z}} = 2$$

$$T_{z\bar{z}} = 0 = T_{\bar{z}z}$$

$$\sim -\left(\frac{1}{2}\right) \left(-\frac{1}{2} \ln(k^2 - z_0)^2\right) \cdot 2 : \partial X(z_0)$$

$\pi' X'' \sim$

$$\sim -\left(\frac{1}{z}\right) \left(-\frac{1}{2} \ln(z-z_0)^2\right) \cdot 2 : \partial X(z_0)$$

$$\pi' X'' \sim \frac{1}{z-z_0} \partial X(z_0)$$

$$\sim -\left(\frac{1}{z}\right) \left(-\frac{1}{2} \ln(z-z_0)^2\right) \cdot 2 : \partial X(z_0)$$

$$\pi^* X'' \sim \frac{1}{z-z_0} \partial X(z_0)$$

$\tilde{\pi}^* X''$

$$\sim -\left(\frac{1}{z}\right) \left(-\frac{1}{2} \ln(r_0 - z_0)^2\right) \cdot 2 \cdot \partial X(z_0).$$

$$\pi^* X'' \sim \frac{1}{z-z_0} \partial X(z_0)$$

$$\tilde{\pi}^* X'' \sim \frac{1}{\bar{z}-\bar{z}_0} \bar{\partial} X(\bar{z}_0)$$

$$\sim -\left(\frac{1}{\pi}\right) \left(-\frac{\pi i}{2} \ln(z-z_0)\right)^2 \cdot 2 = \partial X(z_0).$$

$$\Pi^+ X'' \sim \frac{1}{z-z_0} \partial X(z_0)$$

$$\tilde{\Pi}^- X'' \sim \frac{1}{\bar{z}-\bar{z}_0} \bar{\partial} X(\bar{z}_0)$$

$$\partial X(z_0)$$

$$\sim -\left(\frac{1}{z}\right) \left(-\frac{\pi i}{2} \ln(z-z_0)\right)^2 \cdot 2 = \partial X(z_0).$$

$$T' X'' \sim \frac{1}{z-z_0} \partial X(z_0)$$

$$\bar{T} X'' \sim \frac{1}{\bar{z}-\bar{z}_0} \bar{\partial} X(\bar{z}_0)$$

$$\therefore \partial X(z_0) + \bar{\partial} X(\bar{z}_0)$$

$$\sim \left(\frac{1}{z} \right) \left(-\frac{\pi i}{2} \ln(z - z_0)^2 \right) \cdot 2 : \partial X(z_0).$$

$$T' X'' \sim \frac{1}{z - z_0} \partial X(z_0)$$

$$\bar{T} X'' \sim \frac{1}{\bar{z} - \bar{z}_0} \bar{\partial} X(\bar{z}_0)$$

$$10 \partial X(z_0) + 10 \bar{\partial} X(\bar{z}_0)$$

$$\sim -\left(\frac{1}{z}\right) \left(-\frac{\pi^2}{2} \ln(z-z_0)^2\right) \cdot 2 = \partial X(z_0)$$

$$\text{Tr } X'' \sim \frac{1}{z-z_0} \partial X(z_0)$$

$$\tilde{\text{Tr}} X'' \sim \frac{1}{\bar{z}-\bar{z}_0} \bar{\partial} X(\bar{z}_0)$$

$$(\text{Tr } X(z_0) + i \tilde{\text{Tr}} \bar{\partial} X(\bar{z}_0)) = \frac{1}{2\pi} \delta X$$

$$\delta X^{\mu} = -\varepsilon \Omega(z_0)$$

$$\sim -\left(\frac{1}{z}\right) \cdot \left(-\frac{i}{2} \ln(z-z_0)\right)^2 \cdot 2 = \Im X(z_0)$$

$$\Re X \sim \frac{1}{z-z_0} \Im X(z_0)$$

$$\tilde{\Re} X \sim \frac{1}{\bar{z}-\bar{z}_0} \Im X(\bar{z}_0)$$

$$(\Re z) \Im X(z) + (\Im \bar{z}) \Im X(\bar{z}) = \frac{1}{8} \Delta X$$

$$\delta X^{\mu} = -\varepsilon \Omega(z_0) \partial X - \varepsilon \bar{\Omega}(z_0) \bar{\partial} X$$

$$\delta X^{\mu} = -\epsilon \Omega(z_0) \partial X - \epsilon \bar{\Omega}(z_0) \bar{\partial} X$$

??

$$\delta X^m = -\varepsilon \Omega(z_0) \partial X - \varepsilon \bar{\Omega}(z_0) \bar{\partial} X$$

??

\Rightarrow

$$\delta X^{\mu} = -\varepsilon \mathcal{O}(z_0) \partial X - \varepsilon \bar{\mathcal{O}}(z_0) \bar{\partial} X$$

$$z_0 \rightarrow z_0 + \varepsilon \mathcal{O}(z_0)$$

$$\delta X(z) = -\epsilon \psi(z_p) \bar{\partial} X - \epsilon \bar{\psi}(z_p) \bar{\partial} X$$

??

$$z \rightarrow z_p + \epsilon \psi(z_p)$$

$$\delta X^m(z) = -\varepsilon \psi(z_r) \partial X - \varepsilon \bar{\psi}(z) \bar{\partial} X$$

??

$$z \rightarrow z_r + \varepsilon \psi(z_r)$$

$$\delta X^m = X^{m'}(z', \bar{z}') - X^m(z, \bar{z}) = X^m(z - \varepsilon \psi, z - \varepsilon \bar{\psi}) - X^m$$

anti holomorphic.

$$\delta X^{\mu} = -\varepsilon \psi(z_p) \partial X - \varepsilon \bar{\psi}(z) \bar{\partial} X$$

??

$$z \rightarrow z_0 + \varepsilon \psi(z_p)$$

$$\begin{aligned}\delta X^{\mu} &= X^{\mu'}(z', \bar{z}') - X^{\mu}(z, \bar{z}) = X^{\mu}(z - \varepsilon \psi, z - \varepsilon \bar{\psi}) - X^{\mu} \\ &= -\varepsilon \psi \partial X - \varepsilon \bar{\psi} \bar{\partial} X\end{aligned}$$

holomorphic.

$$\delta X^{\mu}(z) = -\varepsilon \mathcal{U}(z_p) \partial X - \varepsilon \bar{\mathcal{U}}(\bar{z}) \bar{\partial} X$$

??

$$z \rightarrow z + \varepsilon \mathcal{U}(z_p)$$

$$\begin{aligned}\delta X^{\mu} &= X^{\mu'}(z', \bar{z}') - X^{\mu}(z, \bar{z}) = X^{\mu}(z - \varepsilon \mathcal{U}, \bar{z} - \varepsilon \bar{\mathcal{U}}) - X^{\mu}(z, \bar{z}) \\ &= -\varepsilon \mathcal{U} \partial X - \varepsilon \bar{\mathcal{U}} \bar{\partial} X\end{aligned}$$

holomorphic.

$$\delta X^m(z) = -\varepsilon \partial \psi(z_p) \partial X - \varepsilon \bar{\psi}(z) \bar{\partial} X$$

$$z \rightarrow z_0 + \varepsilon \psi(z_p) \quad \xrightarrow{\text{infinitesimal transform}} \quad z \rightarrow z' = f(z)$$

$$\begin{aligned}\delta X^m &= X^{m'}(z', \bar{z}') - X^m(z, \bar{z}) = X^m(z - \varepsilon \psi, z - \varepsilon \bar{\psi}) - X^m \\ &= -\varepsilon \psi \partial X - \varepsilon \bar{\psi} \bar{\partial} X\end{aligned}$$

holomorphic.

$$(z \rightarrow z_0 + \varepsilon \cup z_0) \dots \dots \dots (z)$$

$$\delta X^m = X^{m'}(z, \bar{z}') - X^m(z, \bar{z}) = X^m(z - \varepsilon \cup, z - \varepsilon \bar{\cup}) - X^m(z, \bar{z})$$

$$= -\varepsilon \cup \partial X - \varepsilon \bar{\cup} \bar{\partial} X$$

$J^a \Rightarrow J_a^* = a_m J_m^m$ holomorphic

$\left. + \frac{i}{\pi} \partial_a X^m \right] \Rightarrow J_{\bar{a}}^m = \frac{i}{\pi} \bar{\partial} X^m$

$\check{J}_{\bar{a}}^m = \frac{i}{\pi} \bar{\partial} X^m$ antiholomorphic.

$$\wedge((c,c) - \wedge(c,c)) = \wedge(c(c/c) - x)$$
$$= -\epsilon \cup \partial x - \epsilon \bar{\partial} \bar{\partial} x$$

$$f(\tau) = \zeta \tau$$

$$\partial_\alpha X^\mu$$

$$\nabla((c, \epsilon)) = \nabla((\tau, \epsilon)) = \nabla((c \cdot \tau / (\tau - c))) - X$$
$$= -\epsilon \cup \partial X - \epsilon \bar{\cup} \bar{\partial} X$$

$$f(\tau) = \zeta \tau$$

$|\zeta| = 1 \rightarrow$ just a rotation

$$\nabla \cdot (\epsilon, \mathbf{E}) = \nabla \cdot (\epsilon_0 \mathbf{E}) = \nabla \cdot (\epsilon_0 \epsilon \mathbf{E}) - \nabla \cdot (\epsilon \mathbf{E})$$

$$f(\tau) = \zeta \tau$$

$|\zeta| = 1 \rightarrow$ just a ratio

But in general: not scaling

$$\wedge((c,c) - \wedge(r,r)) = \wedge(c-c/r, c/r - r)$$
$$= -\epsilon \cup \partial X - \epsilon \overline{\cup} \bar{\partial} X$$

$$f(z) = \zeta z$$

$|z|, |\zeta| = 1 \rightarrow$ just a rotation

But in general : rot + scaling

$$\frac{y'^1}{y'^2} = \frac{y^1}{y^2}$$

$$\nabla((c, \epsilon)) = \nabla((\bar{c}, \bar{\epsilon})) = \nabla(c - \bar{c}) + \nabla(\epsilon - \bar{\epsilon})$$

$$= -\epsilon \cup \partial X - \bar{\epsilon} \bar{\cup} \bar{\partial} X)$$

$$f(z) = \zeta z$$

$|z| = 1 \rightarrow$ just a rotation

But in general : not scaling } $y^m \rightarrow \lambda \cdot y^m$

$$\frac{y'^1}{y^1} = \frac{y'^2}{y^2}$$

$$\nabla((\zeta, \zeta)) = \nabla(\zeta, \zeta) = -\zeta \cup \partial X - \zeta \bar{\cup} \bar{\partial} X$$

$$f(z) = \zeta z$$

$\Im z \mid \zeta \mid = 1 \rightarrow$ just a ratio..

But in general : not scaling } $y^m \rightarrow \lambda \cdot y^m$

$$\frac{y^{m_1}}{y^{m_2}} = \frac{y^{m_1}}{y^{m_2}}$$

special
confor-
mial
transf.

$$= -\mathcal{E} \cup \partial X - \mathcal{E} \bar{\cup} \bar{\partial} X$$

$$f(\vec{r}) = \zeta \vec{r}$$

$\nabla |\zeta| = 1 \rightarrow$ just a rotation

But in general : rot + scalings } $y^m \rightarrow \lambda \cdot y^m$

$$\frac{y'^1}{y^1} = \frac{y'^2}{y^2} \dots \frac{y'^n}{y^n}$$

← special
confor-
transf.

$$\delta X^{(n)}(z) = -\varepsilon \cup(z_p) \partial X - \varepsilon \bar{\cup}(z) \bar{\partial} X$$

??

$$z_p \rightarrow z_p + \varepsilon \cup(z_p)$$

infinitesimal transform $\xrightarrow{f(z)}$ is holomorphic

$$\begin{aligned}\delta X^{(n)} &= X^{(n)}(z', \bar{z}') - X^{(n)}(z, \bar{z}) = X^{(n)}(z - \varepsilon \cup, z - \varepsilon \bar{\cup}) - X^{(n)} \\ &= -\varepsilon \cup \partial X - \varepsilon \bar{\cup} \bar{\partial} X\end{aligned}$$

Conformal inv and DPF's



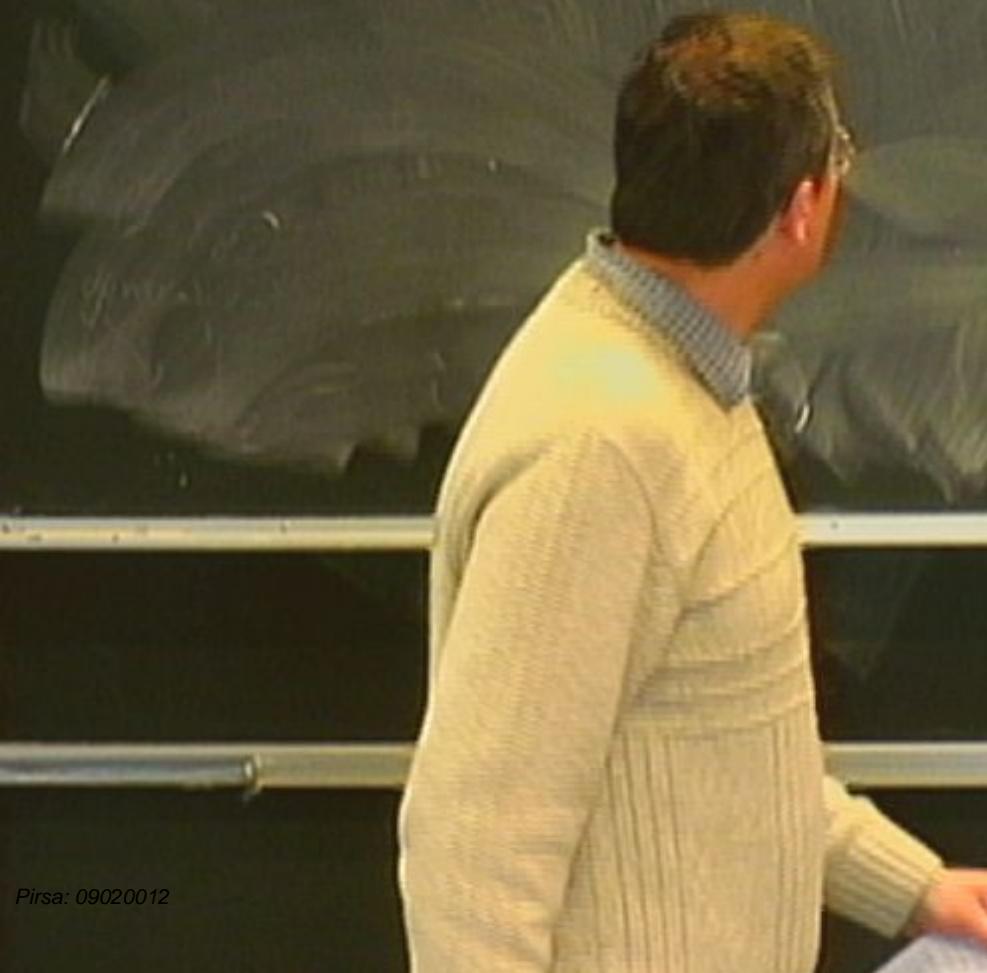
Conformal Inv and DPF's

$$T(z) A(0,0)$$



Conformal Inv and DPF's

$$\Psi(z) A(0,0) \sim \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0,0) + \text{non-singular}$$



Conformal inv and OPE's

$$\Psi(z) A(0,0) \sim \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0,0) + \text{non-singular}$$

are completely determined
by conformal transf



$$j_2 = \vec{v} \cdot \vec{\sigma}(r) T'(r)$$

$$\frac{p_{es}}{2\pi D} j_2 A(\sigma, \phi) = \frac{1}{\epsilon} \delta A$$



• Informal (transl.)

$$j_{\alpha} = \delta \phi(\alpha) T'(\alpha)$$

$$\int_D j_{\alpha} A(\alpha, \vec{\alpha}) = \sum_{\alpha} \delta A$$

042

$$j_z = i \mathcal{O}(z) T(z)$$

$$\underset{z \rightarrow D}{\text{Res}} j_z A(0, \bar{z}) = \frac{1}{\varepsilon} \delta A$$

$$\underset{z \rightarrow 0}{\text{Res}} i \mathcal{O}(z) \left[\sum_{n=0}^{\infty} \frac{1}{z^{n+1}} B^{(n)}(0, \bar{z}) \right] + \text{higher}$$

$$j_z = i \mathcal{O}(z) T'(z)$$

$$\rightarrow_D j_z A(0, \bar{z}) = \frac{i}{\epsilon} \delta A$$

$$\begin{aligned} & \text{provided } \\ & \lim_{z \rightarrow 0} i \mathcal{O}(z) \left[\sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, \bar{z}) + \text{higher-order terms} \right] = \frac{i}{\epsilon} \delta A \end{aligned}$$



$$j_z = i \mathcal{O}(z) T(z)$$

$$\lim_{\epsilon \rightarrow 0} j_z A(0, \bar{z}) = \frac{1}{i\epsilon} \delta A$$
$$\lim_{\epsilon \rightarrow 0} i \mathcal{O}(z) \left[\sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, \bar{z}) + \cancel{\text{higher order terms}} \right] = \frac{1}{i\epsilon} \delta A$$

$$U(t) = \sum_{k=0}^{\infty} \frac{t^k}{k!} U^k$$

$$C(t) = \sum_{k=0}^{\infty} \frac{c_k e^{kt}}{k!}$$

$$U(t) = \sum_{k=0}^{\infty} \frac{U_k(t)}{k!}$$

$$\frac{U_k}{k!}$$



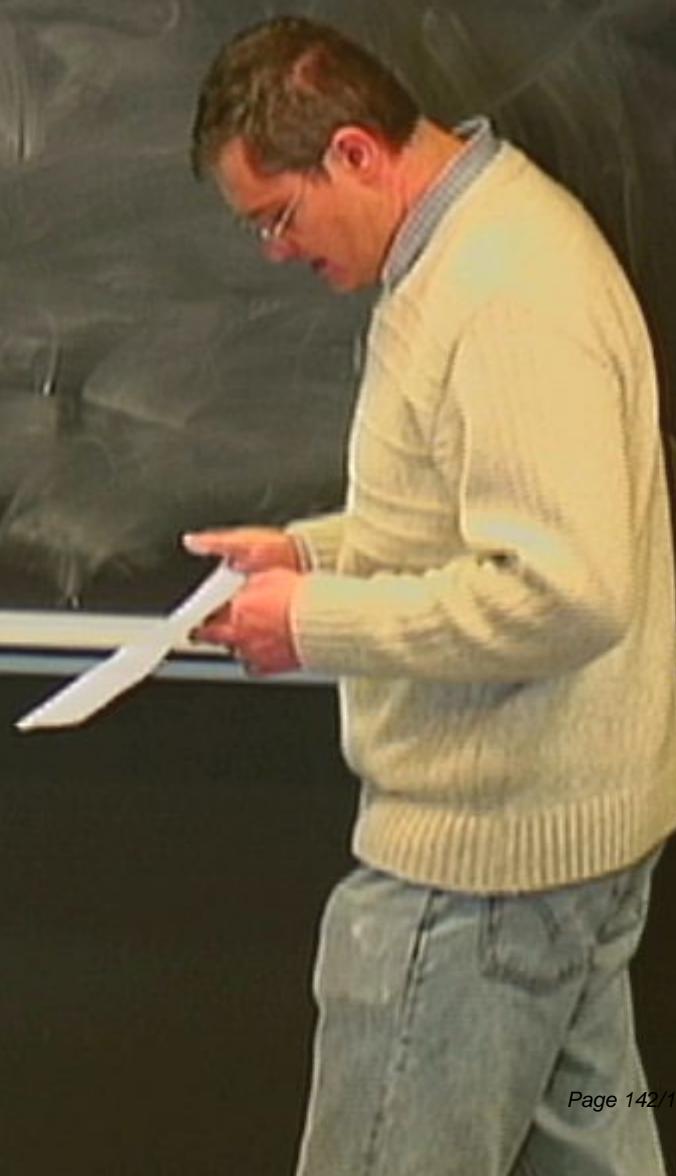
$$U(t) = \sum_{k=0}^{\infty} \frac{t^k}{k!} U^{(k)}$$

$$\therefore \frac{d^5 U(t)}{dt^5} \neq U^{(5)}(0, \vec{0})$$



$$U(t) = \sum_{k=0}^{\infty} \frac{U_0}{k!} t^k$$

$$\sum_{n=0}^{\infty} \frac{\partial^n U(t)}{n!} A^{(n)}(0, \vec{0}) =$$



$$U(t) = \sum_{k=0}^{\infty} \frac{U_k}{k!} t^k$$

$$\sum_{n=0}^{\infty} \frac{\delta U(t)}{n!} A^{(n)}(0, \bar{0}) = \frac{1}{e} \delta A$$

$$U(t) = \sum_{k=0}^{\infty} \frac{U_k}{k!} e^{ik\theta}$$

$$\sum_{n=0}^{\infty} \left(\frac{\partial^n U(t)}{n!} A^{(n)}(0, \bar{0}) \right) = \delta A$$

$$+ i \sum_{h=0}^{\infty} \frac{\overline{\partial^n U}}{n!} A^{(n)}(0, \bar{0})$$



$$U(t) = \sum_{k=0}^{\infty} \frac{\partial^k U}{\partial t^k}(0)$$

$$z \rightarrow z + \delta U(z)^T$$

$$\tilde{U} = \sum_{n=0}^{\infty} \frac{\partial^n U}{\partial t^n}(0) A^{(n)}(0, \bar{0}) = \delta A$$

$$+ i \sum_{n=0}^{\infty} \frac{\partial^n U}{\partial t^n}(0) A^{(n)}(0, \bar{0})$$

Consider a basis of operators

$$= -\partial \phi \partial X - \bar{\psi} \bar{\psi} \partial X$$

$$f(z) = \zeta z$$

$$\text{Is } |\zeta| = 1 \rightarrow$$

But in general

$$\left. \begin{aligned} \frac{y'}{y''} &= \frac{y''}{y'''} && \text{spiral} \\ &&& \text{conformal} \\ &&& \text{transf.} \\ y''' &\rightarrow \lambda \cdot y''' && \end{aligned} \right\} \text{a ratio, scaling}$$

$$f(z) = \zeta z$$

$|z| = 1 \rightarrow$ just a rotation

But in general: rot + scaling

$$\left. \begin{array}{l} y' = \frac{y^1}{y^2} \\ y' = \lambda y^1 \end{array} \right\} y^1 \rightarrow \text{special conformal transform}$$

Consider a basis of operators which are eigenvectors
of right scale transformation.

Consider a basis of operators which are eigenvalues
of right scale transformation.

$$A(z, \bar{z}) \rightarrow$$

Consider a basis of operators which are eigenfunctions
of right scale transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') =$$

Consider a basis of operators which are eigenfunctions
of right scale transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{^{-h} \right\} \left\{-\bar{h}\right\} A(z, \bar{z})$$

Consider a basis of operators which are eigenvalues
of right scale transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{^{-1} \right\} \left\{^{-1} \right\} A(z, \bar{z})$$
$$z \rightarrow z' = \left\{ z \right\}$$

Consider a basis of operators which are eigenfunctions
of right scale transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{ \begin{array}{c} -h \\ -\bar{h} \end{array} \right\} A(z, \bar{z})$$
$$z \rightarrow z' = \left\{ \begin{array}{c} z \\ (h, \bar{h}) \end{array} \right\} - \text{conformal weights}$$

Consider a basis of operators which are eigenfunctions
of right scale transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{ \begin{matrix} -h \\ -\bar{h} \end{matrix} \right\} A(z, \bar{z})$$

h, \bar{h} $z \rightarrow z' = \left\{ \begin{matrix} z \\ h \end{matrix} \right\}$ - conformal weights

Consider a basis of operators which are eigenvalues
of right scale transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{ \begin{matrix} -h \\ -\bar{h} \end{matrix} \right\} A(z, \bar{z})$$

$h + \bar{h}$ - dim of A
 $h - \bar{h}$ - spin

$z' \rightarrow z = \left\{ \begin{matrix} h \\ \bar{h} \end{matrix} \right\}$ - conformal weights

$$\wedge (\mathcal{C} - \mathcal{E} \cup^*) - X(\mathcal{C}) = \Gamma \mathcal{E} \cup \mathcal{O}_0 X = \emptyset \wedge$$

$$D_2 : (h, l) \rightarrow (h+1, l)$$



$$\wedge (\mathcal{O} - \mathcal{E}\mathcal{O}^*) \rightarrow (\mathcal{O}) = \mathcal{E}\mathcal{O} \cap \mathcal{O}_0 \times = \mathcal{O} \wedge$$

$$D_2 : (h, \tilde{h}) \rightarrow (h+1, \tilde{h})$$

$$A \rightarrow \begin{cases} h & \tilde{h} \\ \tilde{h} & A \end{cases}$$



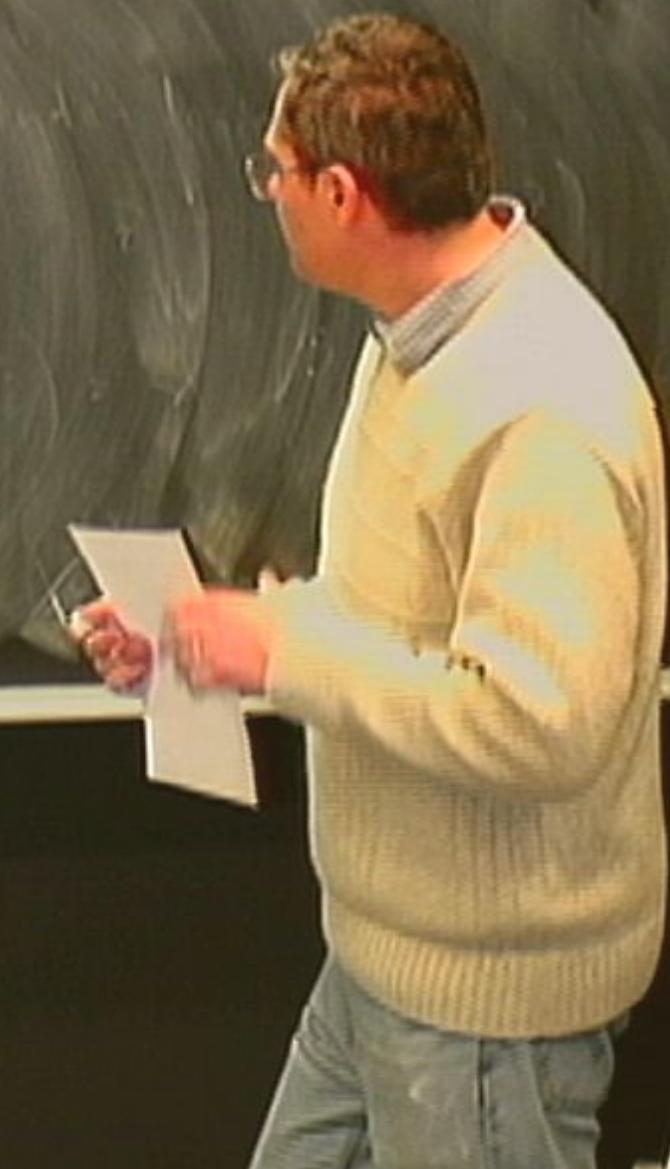
$$\wedge (\mathcal{O} - \mathcal{E}\mathcal{O}^*) \rightarrow (0) = \Gamma \mathcal{E} \mathcal{O} / \mathcal{O}_0, \wedge = \mathcal{O} \wedge$$

$$D_0 : (h, \tilde{h}) \rightarrow (h+1, \tilde{h}).$$

$$A \rightarrow \begin{cases} h \\ \tilde{h} \end{cases} \begin{cases} \tilde{h} \\ A \end{cases}$$

$$\gamma_2 A =$$

$$\delta \downarrow \frac{\partial}{\partial z} =$$



$$\wedge (\mathcal{O} - \mathcal{E}\mathcal{O}^*) \rightarrow (\mathcal{O}) = \Gamma \mathcal{E} \mathcal{O} / \mathcal{O}_0, \wedge = \mathcal{O} \wedge$$

$$D_0 : (h, \tilde{h}) \rightarrow (h+1, \tilde{h}).$$

$$A \rightarrow \begin{cases} -h & \tilde{h} \\ \tilde{h} & A \end{cases}$$

$$\gamma_2 A =$$

$$\downarrow \frac{\partial}{\partial z} \rightarrow \frac{\partial}{\partial z'} = \tilde{\xi}^1 \frac{\partial}{\partial z}$$



$$\wedge (\mathcal{O} - \mathcal{E}\mathcal{O}^*) \rightarrow (0) = \Gamma \mathcal{E} \mathcal{O} / \mathcal{O}_0, \wedge = 0 \wedge$$

$$\gamma_2 : (h, \tilde{h}) \rightarrow (h+1, \tilde{h})$$

$$A \rightarrow \begin{cases} -h \\ -\tilde{h} \end{cases} A$$

$$\gamma_2 A = \begin{cases} -1 \\ -h \end{cases} \begin{cases} -\tilde{h} \end{cases} \delta A$$

$$\delta \downarrow \frac{\partial}{\partial z} \rightarrow \frac{\partial}{\partial z'} = \begin{cases} -1 \\ -2 \end{cases}$$

$$\wedge (\mathcal{O} - \mathcal{E}\mathcal{O}^*) \rightarrow (\mathcal{O}) = \mathcal{E}\mathcal{O} \cap \mathcal{O}_0, \wedge = \mathcal{O} \wedge$$

$$\gamma_2 : (h, \tilde{h}) \rightarrow (h+1, \tilde{h}).$$

$$A \rightarrow \left\{ \begin{array}{c} h \\ \text{not } h \end{array} \right\} \left\{ \begin{array}{c} \tilde{h} \\ \text{not } \tilde{h} \end{array} \right\} A$$

$$\bar{\gamma}_2 : (h, \tilde{h})$$

$$\gamma_2 A = \bar{\gamma}^{-1} (\{h\}) \bar{\gamma}^{-1} (\{\tilde{h}\}) \gamma_2 A$$

$$\frac{\partial}{\partial z} \rightarrow \frac{\partial}{\partial z'} = \bar{\gamma}^{-1} \frac{\partial}{\partial z}$$



$$\wedge (\mathcal{C} - \mathcal{E} \cup^*) - X(C) = \Gamma \mathcal{E} \cup \mathcal{O}_0 X = C \wedge$$

$$\gamma_2 : (h, \tilde{h}) \rightarrow (h+1, \tilde{h})$$

$$A \rightarrow \left\{ \begin{array}{c} h \\ \downarrow \\ h+1 \end{array} \right\} \left\{ \begin{array}{c} \tilde{h} \\ \downarrow \\ \tilde{h} \end{array} \right\} A$$

$$\bar{\gamma}_2 : (h, \tilde{h}) \rightarrow (h, \tilde{h}+1)$$

$$\gamma_2 A = \bar{\gamma}^{-1} (\bar{\gamma}^{-h}) \bar{\gamma}^{-\tilde{h}} \gamma_2 A$$

$$\delta \frac{\partial}{\partial z} \rightarrow \frac{\partial}{\partial z'} = \bar{\gamma}^{-1} \frac{\partial}{\partial z}$$

Consider infinitesimal rigid scale trans.

$$z \rightarrow z' = z + \epsilon z$$



Consider infinitesimal rigid scale trans.

$$z \rightarrow z'' = z + \varepsilon z \Rightarrow \xi = (1 + \varepsilon)$$

Consider infinitesimal rigid scale trans.

$$z \rightarrow z' = z + \varepsilon \bar{z} \Rightarrow \xi = (1 + \varepsilon)$$

$$\cup(\bar{z})$$

Consider infinitesimal rigid scale trans.

$$z \rightarrow z' = z + \varepsilon z \Rightarrow \xi = (1 + \varepsilon)$$

$$\psi(\xi) = z$$

Consider infinitesimal rigid scale trans.

$$z \rightarrow z' = z + \varepsilon \psi(z) \Rightarrow \zeta = (1 + \varepsilon)$$

$$\psi(\bar{z}) = \bar{z}$$

$$A'(z', \bar{z}') =$$

Consider infinitesimal rigid scale trans.

$$z \rightarrow z' = z + \varepsilon \bar{z} \Rightarrow \bar{z}' = (1 + \varepsilon) \bar{z}$$

$$\cup(\bar{z}') = \bar{z}$$

$$A'(z', \bar{z}') = (1 + \varepsilon)^{-1} A(z, \bar{z})$$



Consider infinitesimal rigid scale trans.

$$z + \epsilon \psi(z)$$

$$z \rightarrow z' = z + \epsilon z \Rightarrow \xi = (1 + \epsilon)$$

$$\psi(\bar{z}) = \bar{z}$$

$$A'(z', \bar{z}') = (1 + \epsilon)^{-h} A(z, \bar{z}) + (1 - \epsilon^{-h}) A(z, \bar{z})$$

Consider infinitesimal rigid^a scale trans.

$$z \rightarrow z' = z + \varepsilon \bar{z} \Rightarrow \xi = (1 + \varepsilon)$$

$$\psi(\bar{z}') = \bar{z}$$

$$A'(z', \bar{z}') = (1 + \varepsilon)^{-h} A(z, \bar{z}) + (-\varepsilon h) A(z, \bar{z})$$

$$\delta A = A'(z, \bar{z}) - A'(z, \bar{z})$$

$$= A$$

$$\delta A = A'(z, \bar{z}) - A'(z, \bar{z})$$

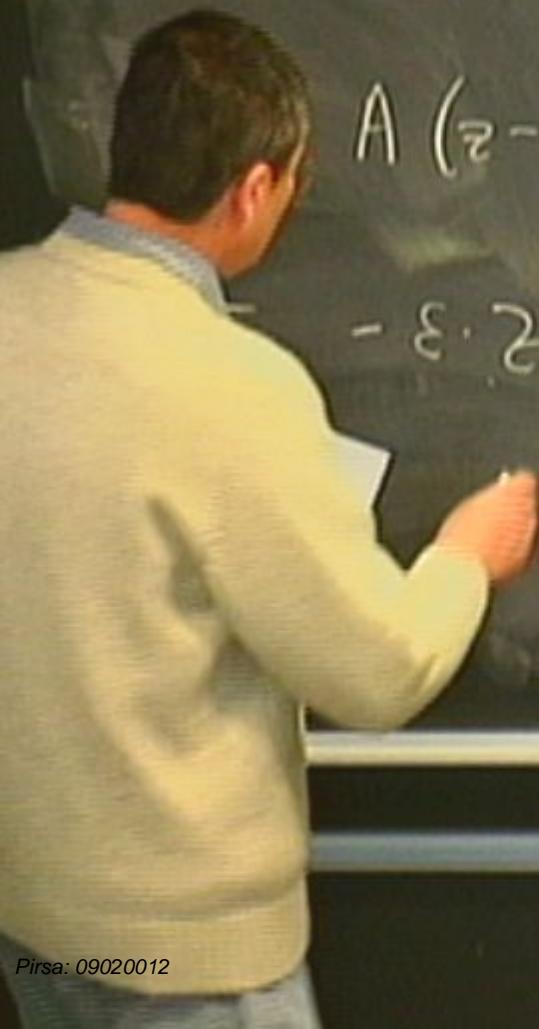
$$= A(z - \varepsilon z, \bar{z}) - h \varepsilon A(z, \bar{z}) - A(z, \bar{z})$$

=

$$\delta A = A'(z, \bar{z}) - A'(z_0, \bar{z}_0)$$

$$A(z - \varepsilon z_0, \bar{z}) - h \in A(z_0, \bar{z}_0) - A(z_0, \bar{z})$$

$$- \varepsilon \cdot z$$



$$\begin{aligned}
 \Im A &= A(z, \bar{z}) - A(\bar{z}, z) \\
 &= A(z - \varepsilon z, \bar{z}) - h \varepsilon A(z, \bar{z}) - A(z, \bar{z}) \\
 &= -\varepsilon \cdot z \cdot \Im A(0) - \varepsilon \cdot h \cdot A(0)
 \end{aligned}$$

z → 0 O(z)

$$\left[\sum_{n \geq 0} \frac{1}{z^{n+1}} A^{(n)}(0, \bar{0}) + \cancel{\text{higher terms}} \right] = \frac{1}{i\varepsilon} \Im A$$

$$\begin{aligned}
 \Im A &= A(z, \bar{z}) - A(\bar{z}, z) \\
 &= A(z - \varepsilon z, \bar{z}) - h \varepsilon A(z, \bar{z}) - A(z, \bar{z}) \\
 &= -\varepsilon \cdot z \cdot \Im A(0) - \varepsilon \cdot h \cdot A(0) \\
 &\xrightarrow{\text{general}} A(0, \bar{0}) =
 \end{aligned}$$

$$\lim_{\varepsilon \rightarrow 0} \frac{(A(z))}{i\varepsilon} \left[\sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, \bar{0}) + \cancel{\text{higher order}} \right] = \frac{1}{i\varepsilon} \delta A$$

$$= A(z - \varepsilon z, \bar{z}) - h \in A(z, \bar{z}) - A(z, \bar{z})$$

$$= -\varepsilon \cdot z \cdot \partial A(0) - \varepsilon \cdot h \cdot A(0)$$

In general

$$T(z) A(0, \bar{0}) \approx \sum_{h=0}^{\infty} \frac{1}{z^{h+1}} A^{(h)}(0, 0) \Leftrightarrow$$

\square $\lim_{z \rightarrow 0} i \partial(z) \left[\sum_{h=0}^{\infty} \frac{1}{z^{h+1}} A^{(h)}(0, \bar{0}) + \cancel{\text{higher terms}} \right] = \frac{1}{i\varepsilon} \delta A$

$$= A(z - \varepsilon z, \bar{z}) - h \in A(z, \bar{z}) - A(z, \bar{z})$$

$$= -\varepsilon \cdot z \cdot \partial A(0) - \varepsilon \cdot h \cdot A(0)$$

In general

$$T(z) A(0, \bar{0}) \approx \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, 0) \Leftrightarrow \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{h^n} \left[z^n \partial_z A^{(n)} \right]$$

\hookrightarrow $\lim_{z \rightarrow 0} \frac{i \omega(z)}{z} \left[\sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, 0) + \cancel{\text{higher order}} \right] = \frac{1}{i\varepsilon} \delta A$

$$= A(z - \varepsilon z, \bar{z}) - h \in A(z, \bar{z}) - A(z, \bar{z})$$

$$= -\varepsilon \cdot 2 \cdot \partial A(0) - \varepsilon \cdot h \cdot A(0)$$

In general

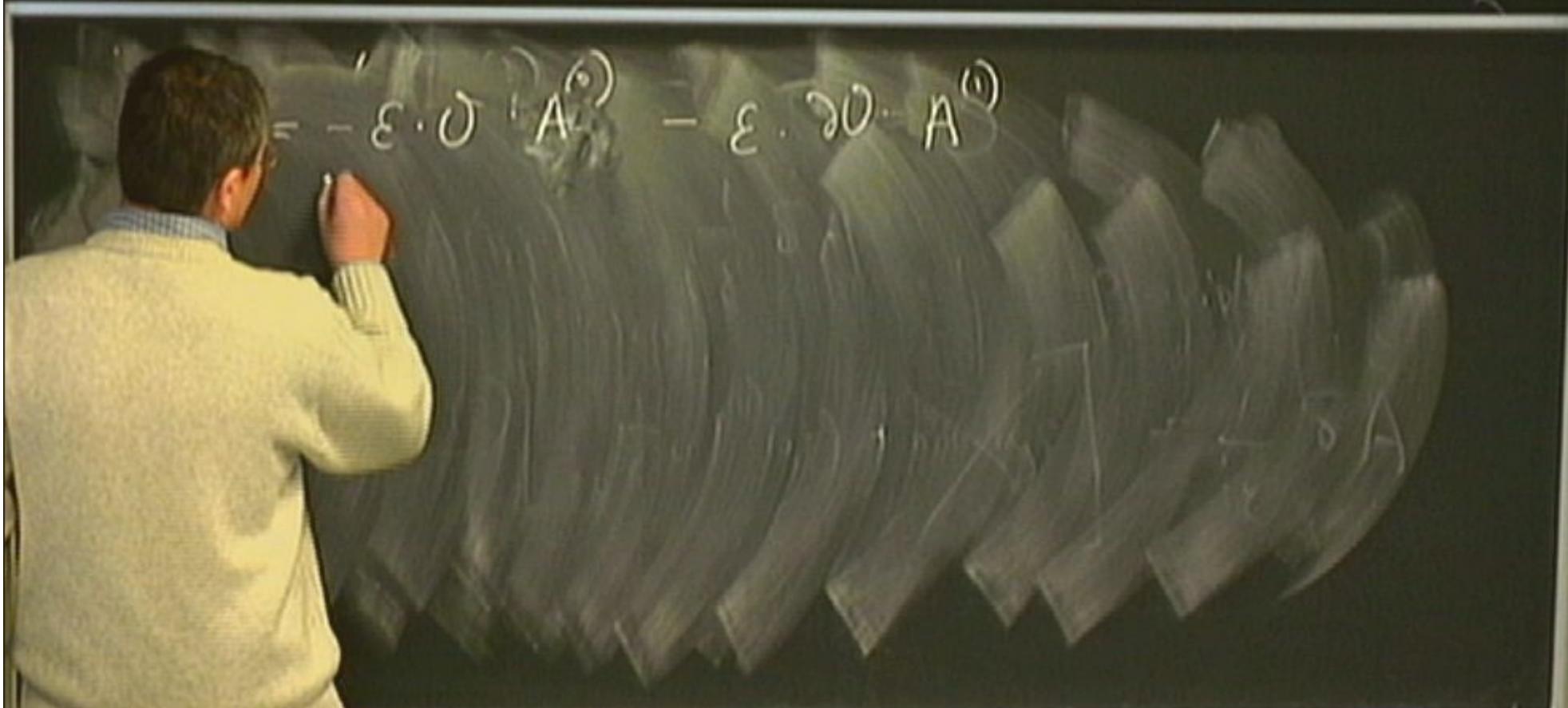
$$T(z) A(0, \bar{0}) \approx \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, 0) \Leftrightarrow \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{h^n} \left[i \partial A^{(n)} \right]$$

\hookrightarrow $\lim_{z \rightarrow 0} \frac{i \partial(z)}{z} \left[\sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, \bar{0}) + \text{higher terms} \right] = \frac{1}{i\varepsilon} \delta A$

$$T(z) A(0, \bar{z}) \sim \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, \bar{z}) \Leftrightarrow \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{n!} \left[\partial_z \partial_{\bar{z}}^n A \right],$$

$$\delta A =$$

$$T(z) A(0, \bar{z}) \approx \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0, 0) \Leftrightarrow \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{n!} \left[\partial_z \partial_{\bar{z}}^n A^{(n)} \right]$$



$$\overline{T(z)A(0,0)} \sim \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0,0) \Leftrightarrow \left\{ \begin{array}{l} \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{n!} \left[\partial_z \partial_z^n A \right] \end{array} \right.$$

$$\begin{aligned} A' &= -\varepsilon \cdot \partial_z A^{(0)} - \varepsilon \cdot \partial_z A^{(0)} \\ &\quad \underbrace{- \varepsilon \cdot z A^{(0)} - \varepsilon \cdot z A^{(0)}} \end{aligned}$$



$$\overline{T(z)A(0,0)} \approx \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0,0) \Leftrightarrow \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{n!} \left[\partial^n A \right],$$

$$= -\varepsilon \cdot \partial \cdot A^{(0)} - \varepsilon \cdot \partial \cdot A^{(0)}$$

$$\hookrightarrow -\varepsilon \cdot \partial A^{(0)} - \varepsilon \cdot \partial A^{(0)}$$

$$(0) =$$

$$\overline{T(z)A(0,0)} \approx \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} A^{(n)}(0,0) \Leftrightarrow \delta A = -\varepsilon \sum_{n=0}^{\infty} \frac{1}{h^n} \left[\partial_z A^{(n)} \right]$$

$$\delta A = -\varepsilon \cdot \partial_z A^{(0)} - \varepsilon \cdot \partial_0 A^{(0)}$$

$$\frac{-\varepsilon \cdot \partial_z A^{(0)} - \varepsilon \cdot \partial_0 A^{(0)}}{}$$

$$A^{(0)} = \partial A$$

$$A^{(n)} = h \cdot A$$

$$\mathbb{T}(z) A = \dots \frac{1}{z}$$



$$T(z) A = \dots + \frac{1}{z^2} A^{(2)} + \frac{1}{z} A^{(1)} + \dots$$

$$\mathbb{T}(z) A = \dots + \frac{1}{z^2} A^{(2)} + \frac{1}{z} A^{(1)} + \dots +$$
$$+ \frac{1}{z^2} A + \frac{1}{z} \partial A$$

$$\mathcal{T}(z) A = \dots + \frac{1}{z^2} A^{(2)} + \frac{1}{z} A^{(1)} + \dots - - - \mathcal{T}(z)$$

$$+ \frac{h}{z^2} \circled{A} + \frac{1}{z} \mathcal{D}A +$$



$$= \frac{h}{z^2} \circled{A} + \frac{1}{z} \partial A + \dots$$

Primary fields or tensor operators

$$\begin{aligned} h + \bar{h} & - \dim \text{of } A \\ h - \bar{h} & = \text{spin} \end{aligned} \quad \leftarrow \mathcal{R} = \{(h, \bar{h})\} - \text{conformal weights}$$

$$\dots - \frac{h}{z^2} \circled{A} + \frac{1}{z} \partial A + \dots$$

Primary fields or tensor operators

$$O'(z', \bar{z}') =$$

$$\begin{aligned} h + \tilde{h} & - \text{dim of } A \\ h - \tilde{h} & - \text{spin} \end{aligned} \quad \leftarrow z' = \left\{ \begin{array}{l} (h, \tilde{h}) - \text{conformal weights} \end{array} \right.$$

$$= \dots + \frac{h}{z^2} \circled{A} + \frac{1}{z} \partial A + \dots$$

Primary fields or tensor operators

$$O'(z', \bar{z}') = (z')^h (\bar{z}')^{\bar{h}}$$

$$\begin{aligned} h - \bar{h} &= \dim A \\ z' &\rightarrow \bar{z}' = \{ z \mid (h, \bar{h}) \text{- conformal weights} \} \end{aligned}$$

$$= \dots + \frac{h}{z^2} \langle A \rangle + \frac{1}{z} \partial A + \dots$$

Primary fields or tensor operators

$$\mathcal{O}'(z', \bar{z}') = (z')^h (\bar{z}')^{\bar{h}} \mathcal{O}(z, \bar{z})$$

$$\begin{cases} h + \bar{h} - \dim A \\ h - \bar{h} \end{cases} \rightarrow r = \left\{ \begin{array}{l} \in \mathbb{R} \\ (h, \bar{h}) \end{array} \right\} \text{ - conformal weights}$$

$$O(z, \bar{z}) = (z')^h (\bar{z}')^{\bar{h}} O(z, \bar{z})$$

$$T(z) O =$$

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{ \begin{matrix} -h \\ -\bar{h} \end{matrix} \right\} A(z, \bar{z})$$

$h + \bar{h}$ - dim of A
 $h - \bar{h}$ - 8916

$z \rightarrow z' = \{ z \in (h, \bar{h}) \text{ - conformal weights}$

$$\hookrightarrow O(z, \bar{z}) = (Oz) \begin{pmatrix} Oz \\ O\bar{z} \end{pmatrix}^T O(\bar{z}, \bar{\bar{z}})$$

$$\hookrightarrow T(z)O = \frac{h}{z^2}O + \frac{1}{z}\partial O$$

or right side transformation.

$$A(z, \bar{z}) \rightarrow A'(z', \bar{z}') = \left\{ \begin{matrix} -h \\ \bar{z}' \\ -\bar{h} \end{matrix} \right\} A(z, \bar{z})$$

$$\begin{matrix} h+\bar{h} & z \rightarrow z' = \{ z \in (h, \bar{h}) \text{ - conformal weights} \\ h-\bar{h} & \text{dim of } A \end{matrix}$$