

Title: Introduction to the Bosonic String

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Abstract: This course provides a thorough introduction to the bosonic string based on the Polyakov path integral and conformal field theory. We introduce central ideas of string theory, the tools of conformal field theory, the Polyakov path integral, and the covariant quantization of the string. We discuss string interactions and cover the tree-level and one loop amplitudes. More advanced topics such as T-duality and D-branes will be taught as part of the course. The course is geared for M.Sc. and Ph.D. students enrolled in Collaborative Ph.D. Program in Theoretical Physics. Required previous course work: Quantum Field Theory (AM516 or equivalent). The course evaluation will be based on regular problem sets that will be handed in during the term. The primary text is the book: 'String theory. Vol. 1: An introduction to the bosonic string. J. Polchinski (Santa Barbara, KITP) . 1998. 402pp. Cambridge, UK: Univ. Pr. (1998) 402 p.' All interested students should contact Alex Buchel at abuchel@uwo.ca as soon as possible.

Conformal Field Theory (CFT)



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⇒ CFT CFT
 $X \rightarrow \lambda X$

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⇒ CFT → $\mathbb{R}^{1,1}$ CFT

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fix world-sheet diff + Weyl

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⇒ CFT $X \rightarrow \xi X$ CFT

→ for string theory
fix world-sheet diff + Weyl

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⇒ CFT $X \rightarrow \xi X$ CFT

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fix world-sheet diff + Weyl

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⇒ CFT $X \rightarrow \xi X$ CFT

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six world-sheet diff + Weyl

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⇒ CFT $X \rightarrow \xi X$ CFT

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⇒ CFT $X \rightarrow \xi X$ CFT (a scaling transformation part conformal)

\Rightarrow CFT

$$X \rightarrow \xi X$$

CFT

(a scaling transformation is part of Conformal)

CFTs show up in condensed matter @

\Rightarrow CFT

$$X \rightarrow \xi X$$

CFT

(a scaling transformation is part of Conformal)

CFTs show up in condensed matter @
2nd order phase transitions

\Rightarrow CFT

$$X \rightarrow \xi X$$

CFT

(a scaling transformation is part of Conformal)

CFTs show up in condensed matter @
2nd order phase transitions

Notes

"Applied CFT" Paul Ginsparg

D -massless scalars in 2-dim. (toy CFT).

D-massless scalars in 2-dim. (toy QFT)

$$S = \frac{1}{4\pi\alpha'} \int d^2\sigma \left[\partial_\mu X^\mu \partial^\mu X_\mu + \partial_2 X^\mu \partial_2 X_\mu \right]$$

D-massless scalars in 2-dim. (toy CFT)

$$S = \frac{1}{4\pi\alpha'} \int d^2\sigma \left[\partial_1 X^\mu \partial_1 X_\mu + \partial_2 X^\mu \partial_2 X_\mu \right]$$

↑ Polyakov action with $\eta_{ab} \Rightarrow \eta_{\sigma\tau}$

$$S = \frac{1}{4\pi\alpha'} \int d^2\sigma \left[\partial_1 X^\mu \partial_1 X_\mu + \partial_2 X^\mu \partial_2 X_\mu \right]$$

↑ Polyakov action with
(σ_0, σ_1)

$\eta_{ab} \Rightarrow \eta'_{ab} \rightarrow \delta_{ab}$
Wick rotation

↑ Polyakov action with

$$\begin{pmatrix} \sigma_0 & \sigma_1 \\ \tau & \sigma \end{pmatrix}$$

→ Wick rotation

$$\begin{pmatrix} \sigma_2 & \sigma_1 \end{pmatrix}$$

$$h_{ab} \Rightarrow \eta_{ab} \rightarrow \delta_{ab}$$

Wick rotation

$$\boxed{\sigma_2 \equiv i\sigma_0}$$

τ σ

Wick rotation (σ_2, σ_1)

$$|\sigma_2 \equiv i\sigma_0|$$

$$z = \sigma^1 + i\sigma^2$$

$$\partial_z =$$

$$\bar{z} = \sigma^1 - i\sigma^2$$

↑ Polyakov action with

$$\begin{pmatrix} \sigma_0 & \sigma_1 \\ \tau & \sigma \end{pmatrix}$$

Wick rotation

$$\begin{pmatrix} \sigma_2 & \sigma_1 \end{pmatrix}$$

$\gamma_{ab} \Rightarrow \eta_{ab} \rightarrow \delta_{ab}$
Wick rotation

$$\sigma_2 \equiv i\sigma_0$$

$$z = \sigma' + i\sigma^2$$

$$\bar{z} = \sigma' - i\sigma^2$$

$$\partial \equiv \partial_z = \frac{1}{2}(\partial_1 - i\partial_2)$$

$$\bar{\partial} \equiv \partial_{\bar{z}} = \frac{1}{2}(\partial_1 + i\partial_2)$$

$$0 = z\bar{z} = 0 \quad \dots$$

Wick rotation (σ_2, σ_1) $\left| \sigma_2 \equiv i\sigma_0 \right|$

$$z = \sigma_1 + i\sigma_2$$

$$\bar{z} = \sigma_1 - i\sigma_2$$

$$\partial \equiv \partial_z = \frac{1}{2}(\partial_1 - i\partial_2)$$

$$\bar{\partial} \equiv \partial_{\bar{z}} = \frac{1}{2}(\partial_1 + i\partial_2)$$

$$\partial \bar{z} = 1 \quad \bar{\partial} z = 0$$

z
↑
holomorphic

\bar{z}
↑
anti-holomorphic

$$v = (v^1, v^2)$$

$$v^2 = v^1 + i v_2$$

$$v^2 = v^1 - i v_2$$

$$v = (v^1, v^2)$$

$$v^2 = v^1 + i v_2$$

$$v^{\bar{2}} = v^1 - i v_2$$

$$\text{Sub} \rightarrow \begin{pmatrix} g_{z\bar{z}} & g_{z\bar{z}} \\ g_{z\bar{z}} & g_{\bar{z}\bar{z}} \end{pmatrix}$$

$$U^2 = U^1 - iU_2$$

$$\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$u^2 = u' - i u_2$$

$$g^{\bar{z}z} = g^{z\bar{z}} = 2$$

$$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$g^{zz} = g_{zz} = 2$$
$$\sqrt{g} = \frac{1}{2} \quad dz$$

$$\begin{pmatrix} -\frac{1}{2} & \\ & 0 \end{pmatrix}$$



$$\sqrt{g} = \sqrt{\det g}$$

$$d\sigma^1 d\sigma^2 = d^2\sigma = \sqrt{\det g} d\sigma^1 d\sigma^2$$

$$\sqrt{g} = \sqrt{\sum}$$

$$d^2x \sqrt{g}$$

$$d\sigma^1 d\sigma^2 = d^2\sigma = \sum d\sigma^1 d\sigma^2$$

$$\sqrt{g} = \frac{1}{2} \quad d\sigma^1 d\sigma^2 = d^2\sigma = 2 d\sigma^1 d\sigma^2$$
$$d^2z \sqrt{g} = d\sigma^2$$

$$\sqrt{g} = \frac{1}{2}$$

$$d\sigma^1 d\sigma^2 = d^2\sigma = \frac{1}{2} d\sigma^1 d\sigma^2$$

$$d^2z \sqrt{g} = d\sigma^2$$

$$\int d^2\sigma \delta^2(\sigma_1, \sigma_2) = \int d^2\sigma \delta(\sigma_1) \delta(\sigma_2) = 1$$

$$\int d^2z \delta^2(\bar{z}, \bar{z}) = 1 \Rightarrow \delta^2(\bar{z}, \bar{z}) = \frac{1}{2} \delta^2(\sigma_1, \sigma_2)$$

$$|\bar{g}| = \frac{1}{2}$$

$$d\sigma^1 d\sigma^2 = d^2\sigma = \frac{1}{2} d\sigma^1 d\sigma^2$$

$$d^2z |\bar{g}| = d\sigma^2$$

$$\int d^2\sigma \delta^2(\sigma_1, \sigma_2) = \int d^2\sigma \delta(\sigma_1) \delta(\sigma_2) = 1$$

$$\int d^2z \delta^2(\bar{z}, \bar{z}) = 1 \Rightarrow \delta^2(\bar{z}, \bar{z}) = \frac{1}{2} \delta^2(\sigma_1, \sigma_2)$$

Divergence theorem.

In above notation,

$$S = \frac{1}{2\pi\alpha'} \int d^2z \partial X^\mu \bar{\partial} X_\mu$$

$$S = \frac{1}{2\pi\alpha'} \int d^2\sigma \partial X^\mu \bar{\partial} X_\mu$$

" $d\sigma d\bar{\sigma}$

→ EOM

$$0 = \frac{\delta S}{\delta \bar{\sigma}} = -\frac{2}{2\pi\alpha'} \partial \bar{\partial} X^\mu = 0$$

∂

Paul Ginsparg

$$0 = \frac{\partial \mathcal{L}}{\partial x^m} = -\frac{2}{2\pi\alpha'} \sqrt{-g} g^{mn} x^m = 0$$

$$\boxed{\partial \mathcal{L} = 0}$$

CFT is show u ... matter

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Notes

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yl Ginsparg

$$\frac{\partial \bar{\partial} X^m}{\partial z} = 0 \quad X^m = X(z, \bar{z})$$

$$\bar{\partial} [\partial X^m] = 0$$

"Applied QFT" Paul Ginsparg

$$\partial \bar{\partial} X^m = 0 \quad X^m = X^m(z, \bar{z})$$

$$\bar{\partial} [\partial X^m] = 0$$

$$\partial X^m = \partial X^m(z)$$

$$\bar{\partial} X^m = \bar{\partial} X^m(\bar{z})$$

or holomorphic function.

$$\partial \bar{\partial} X^m = 0 \quad X^m = X^m(z, \bar{z})$$

$$\bar{\partial} [\partial X^m] = 0$$

$$\partial X^m = \partial X^m(z)$$

↙ a holomorphic function.

$$\partial [\bar{\partial} X^m] = 0$$

$$\bar{\partial} X^m = \bar{\partial} X^m(\bar{z})$$

Some relations

$$\int [dx] e^{-s}$$

Some relations

$$\langle \mathcal{F}[x] \rangle = \int [dx] e^{-S}$$

Some relations

$$\langle \mathcal{F}[x] \rangle = \int [dx] e^{-S} \mathcal{F}[x]$$

Some relations

$$\langle \mathcal{F}[x] \rangle = \frac{\int [dx] e^{-S} \mathcal{F}[x]}{\int [dx] e^{-S}}$$

$$\int [dx] e^{-S}$$

$$\langle f(x) \rangle = \int dx e^{-f(x)}$$

$$\langle 0| \int dx e^{-S} | \rangle$$

holomorphic

↑
antiholomorphic

$$\langle F[x] \rangle = \int [dx] e^{-S} F[x(\vec{t}, \vec{x})]$$

||

$$\int [dx] e^{-S}$$

$$\langle 0|F|0 \rangle$$

||

$\langle 0 | F | 0 \rangle$

$$F_1[X(z_1, \bar{z}_1)] \quad F_2[X(z_2, \bar{z}_2)] \quad F_2[X(z_2, \bar{z}_2)]$$

$\langle 0 | \hat{T} | 0 \rangle$

$$\int dx e^{iS[x(z, \bar{z})]} \hat{T}_1[x(z, \bar{z})] \hat{T}_2[x(z, \bar{z})]$$

$$\langle 0 | F | 0 \rangle$$

$$\int dx e^{iF_1[X(z_1, \bar{z}_1)] + iF_2[X(z_2, \bar{z}_2)] + iF_3[X(z_3, \bar{z}_3)]}$$

$$\langle 0 | F_1 F_2 | 0 \rangle$$

$$\langle 0 | F | 0 \rangle$$

$$\int dx e^{iS[x]} F_1[x(z_1, \bar{z}_1)] F_2[x(z_2, \bar{z}_2)] F_3[x(z_3, \bar{z}_3)]$$

$$\langle 0 | F_1 F_2 F_3 | 0 \rangle$$

$$\langle 0 | \hat{T} | 0 \rangle$$

$$\int dx e^{iS[x(z, \bar{z})]} \hat{T}_1[x(z, \bar{z})] \hat{T}_2[x(z, \bar{z})]$$

$$\|$$

$$\langle 0 | \hat{T}_1 \hat{T}_2 | 0 \rangle$$

$$\|$$

$$\langle 1 | \hat{T} | 2 \rangle$$

$$\langle 0 | \hat{T} | 0 \rangle$$

$$\int dx e^{\hat{T}_1[x(z_1, \bar{z}_1)]} \hat{T}_2[x(z_2, \bar{z}_2)] \hat{T}_2[x(z_2, \bar{z}_2)]$$

$$\| \langle 0 | \hat{T}_1 \hat{T}_2 | 0 \rangle$$

$$\| \langle 1 | \hat{T} | 2 \rangle$$

$$| 2 \rangle = \hat{T}_2 | 0 \rangle$$

Statement:

$$\int dx \partial_x [\text{something}] = 0$$

$$\int dx \partial_x [\text{something}] = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X^m}$$

$$\int dx \partial_x [\text{something}] = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X^m} [\text{something}]$$

$$\int dx \partial_x [\text{something}] = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X^m} [\text{something}]$$



$$0 = \int [dx] \frac{\delta}{\delta X^m} [\text{something}]$$

$$\int d^2z (\gamma_z \partial_z \phi + \partial_{\bar{z}} \phi) = i \oint (\phi d\bar{z} - \phi dz)$$

$$0 = \int [dx] \frac{\delta}{\delta x^m} [\bar{e}^s]$$

$$\delta \bar{e}^s = 1 \Rightarrow \delta^2 \bar{e}^s = 0$$

Example: $\delta \bar{e}^s = 1$

$$\delta \bar{e}^s = 1 \Rightarrow \delta^2 \bar{e}^s = 0$$

$$\delta^2 \bar{e}^s = 0$$

$$\delta \bar{e}^s = 1 \Rightarrow \delta^2 \bar{e}^s = 0$$



$$0 = \int [dx] \frac{\delta}{\delta x^m} [e^{-S}] = - \int [dx] e^{-S} \frac{\delta S}{\delta x^m}$$

~~$$\int [dx] \frac{\delta}{\delta x^m} [e^{-S}] = - \int [dx] e^{-S} \frac{\delta S}{\delta x^m}$$~~

$$0 = \int [dx] \frac{\delta}{\delta x^m} [e^{-S}] = - \int [dx] e^{-S} \frac{\delta S}{\delta x^m}$$

$$= - \int [dx] e^{-S} \dots$$

$$0 = \int [dx] \frac{\delta}{\delta x^m} [e^{-S}] = - \int [dx] e^{-S} \frac{\delta S}{\delta x^m}$$

$$= - \int [dx] e^{-S} \left(\frac{2}{2\pi\alpha'} \partial \bar{x}^m \right)$$

$$0 = \int [dx] \frac{\delta}{\delta X^{\mu}(\bar{t}, \bar{x})} \left[e^{-S} \right] = - \int [dx] e^{-S} \frac{\delta S}{\delta X^{\mu}}$$

$$= - \int [dx] e^{-S} \left(\frac{2}{2\pi\alpha'} \partial \bar{X}^{\mu} \right) = \frac{1}{\pi\alpha'} \langle \partial \bar{X}^{\mu} | 0 \rangle$$

$$0 = \int [dx] \frac{\delta}{\delta X_{(z, \bar{z})}^m} \left[e^{-S} F(z, \bar{z}) \right]_{z \neq z'} = - \int [dx] e^{-S} \frac{\delta S}{\delta X^m}$$

$$= - \int [dx] e^{-S} \left(\frac{z}{2\pi\alpha'} \partial \bar{X}^m \right) = \frac{1}{\pi\alpha'} \langle \partial \bar{X}^m | 0 \rangle$$

$$0 = \int [dx] \frac{\delta}{\delta X_{(z, \bar{z})}^m} \left[e^{-S} F(z, \bar{z}) \right]_{z \neq z'} = - \int [dx] e^{-S} \frac{\delta S}{\delta X^m}$$

$$= - \int [dx] e^{-S} \left(\frac{z}{2\pi\alpha'} \partial \bar{\partial} X^m \right) = \frac{1}{\pi\alpha'} \langle \partial \bar{\partial} X^m | 0 \rangle$$

$$\langle \partial \bar{\partial} X \rangle_F = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X_{(z, \bar{z})}^m} \left[e^{-S} \right]_{z \neq z'} = - \int [dx] e^{-S} \frac{\delta S}{\delta X^m}$$

$$= - \int [dx] e^{-S} \left(\frac{z}{2\pi\alpha'} \partial \bar{\partial} X^m \right) = \frac{1}{\pi\alpha'} \langle 0 | \partial \bar{\partial} X^m | 0 \rangle$$

$$\langle \partial \bar{\partial} X \rangle_F = 0$$

$\bar{\partial} X = 0$ as an operator statement

$$0 = \int [dx] \frac{\delta}{\delta X_{(z, \bar{z})}^m} \left[e^{-S} \mathcal{F}(z, \bar{z}) \right]_{z \neq z'} = - \int [dx] e^{-S} \frac{\delta S}{\delta X^m}$$

$\mathcal{F} = 0$

$$= - \int [dx] e^{-S} \left(\frac{z}{2\pi\alpha'} \partial \bar{\partial} X^m \right) = \frac{1}{\pi\alpha'} \langle \partial \bar{\partial} X^m | 0 \rangle$$

$$\langle \bar{\partial} X \rangle_{\mathcal{F}} = 0$$

$\bar{\partial} X = 0$ as an operator statement.

$$\int dx \partial_x [\text{something}] = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X^m} [\text{something}]$$

Some/

is an operator statement

$$\int_a^b dx \partial_x [\text{something}] = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X^m} [\text{something}]$$

Some

operator statement

$$\int_a^b dx \partial_x [\text{something}] = 0$$

$$0 = \int [dx] \frac{\delta}{\delta X^m} [\text{something}]$$

$$S = \frac{1}{2m} \int dx |\partial x|^2$$

Some

as an operator statement

$$0 = \int [dx] \frac{\delta}{\delta X^{\mu}(z, \bar{z})} \left[e^{-S} \right]$$



$$0 = \int [dx] \frac{\delta}{\delta X^{\mu}(z, \bar{z})} \left[e^{-S} X^{\nu}(z', \bar{z}') \right]$$



$$0 = \int [dx] \frac{\partial}{\partial X^m} (\bar{z}, \bar{z}) \left[\bar{e}^s X^m (\bar{z}', \bar{z}') \right]$$

$$= \int d[x] \left\{ \bar{e}^s \left(\frac{\partial}{\partial X^m} \right)^2 \frac{1}{\pi d'} \partial \bar{z} X^m \right.$$

$$0 = \int [dx] \frac{\delta}{\delta X^m(z, \bar{z})} \left[e^{-S} X^0(z', \bar{z}') \right]$$

$$= \int [dx] \left\{ e^{-S} \left(-\frac{1}{\pi \alpha'} \partial \bar{\partial} X^m + e^{-S} \frac{\delta X^0(z', \bar{z}')}{\delta X^m(z, \bar{z})} \right) \right\}$$



$$= \int d[X] \left\{ e^{-S} \left(\frac{1}{\pi \alpha'} \right)^2 \partial \bar{\partial} X^m + e^{-S} \frac{\delta X^m(z, \bar{z})}{\delta X^m(z', \bar{z}')} \delta_{m\nu}^{\nu} \delta^2(z - z', \bar{z} - \bar{z}') \right\}$$

$$= \int d[X] \left\{ e^{-S_0} X^2 \left(\frac{1}{\pi d'} \right)^2 \partial \bar{\partial} X^m + e^{-S_0} \frac{\delta X^m(z', \bar{z}')}{S X^m(z, \bar{z})} \right\}$$

$$\partial | \bar{\partial} X^m | = 0 \quad \partial X^m = \bar{\partial} X^m$$



$$= \int d[X] \left\{ e^{-S_0} X^2 \frac{1}{\pi \alpha'} \partial \bar{\partial} X^m(z, \bar{z}) + e^{-S_0} \frac{\delta X^m(z', \bar{z}')}{\delta X^m(z, \bar{z})} \delta_{\mu\nu}^2 \delta^2(z' - z, \bar{z}' - \bar{z}) \right\}$$



$$\partial X^m = \partial X^m(\frac{\tau}{\alpha'})$$

$$0 = \int [dx] \frac{\delta}{\delta X^m(z, \bar{z})} \left[e^{-S} X(z', \bar{z}') \right]$$

$$d[X] \left\{ e^{-S} X(z', \bar{z}') \frac{\delta X^m(z', \bar{z}')}{\delta X^m(z, \bar{z})} + \frac{1}{\pi \alpha'} \partial_{\bar{z}} X(z, \bar{z}) \frac{\delta}{\delta X^m(z, \bar{z})} \left[\frac{1}{\delta_{m\mu}} \delta^2(z' - z, \bar{z}' - \bar{z}) \right] \right\}$$

$$= \frac{1}{\pi \alpha'} \partial_{\bar{z}} \partial_z X = 0$$

$$0 = \int [dX] \frac{\delta}{\delta X^m(z, \bar{z})} \left[e^{-S} X^m(z', \bar{z}') \right]$$

$$= \int [dX] \left\{ e^{-S} X^m(z, \bar{z}) \frac{\delta}{\delta X^m(z, \bar{z})} + e^{-S} \frac{\delta X^m(z', \bar{z}')}{\delta X^m(z, \bar{z})} \right\}$$

$$= \frac{1}{\pi} \partial_{\bar{z}} \bar{\partial}_z X^m = 0$$

$$= \int d[X] \left\{ e^{-S} \frac{1}{\pi \alpha'} \partial \bar{\partial} X^m(z, \bar{z}) + e^{-S} \frac{\delta X^\nu(z', \bar{z}')}{\delta X^m(z, \bar{z})} \delta^2(z' - z, \bar{z}' - \bar{z}) \right\}$$

$$= \frac{1}{\pi \alpha'} \partial \bar{\partial} \langle X^m(z, \bar{z}) X^\nu(z', \bar{z}') \rangle$$

$$\partial \bar{\partial} X = 0 \Rightarrow \partial X = 0$$

$$= \int d[X] \left\{ e^{-S_0} X^m(-1) \frac{1}{\pi d'} \partial \bar{\partial} X^m(z, \bar{z}) + e^{-S} \frac{\delta X^m(z, \bar{z})}{\delta X^m(z, \bar{z})} \right\}$$

$$= \frac{1}{\pi d'} \partial_z \bar{\partial}_{\bar{z}} \langle X^m(z, \bar{z}) X^m(z', \bar{z}') \rangle + \langle \delta^2(z, \bar{z}) \rangle$$

$$\partial \bar{\partial} X^m = 0 \quad \bar{\partial} X^m = \bar{\partial} X^m(\bar{z})$$

Applied CFT
 Paul Ginsparg

$$= \int_{\mathbb{R}^2} \delta^2(z-z') X^m(z, \bar{z}) X^0(z', \bar{z}') dz d\bar{z} + \int_{\mathbb{R}^2} \delta^2(z-z') X^m(z, \bar{z}) X^0(z', \bar{z}') dz d\bar{z}$$

$$= \int_{\mathbb{R}^2} \delta^2(z-z') X^m(z, \bar{z}) X^0(z', \bar{z}') dz d\bar{z} = \int_{\mathbb{R}^2} \delta^2(z-z') X^m(z, \bar{z}) X^0(z', \bar{z}') dz d\bar{z}$$

$$\frac{1}{\pi} \int_{\mathbb{C}} \bar{z} \langle X^m(z, \bar{z}) X^n(z', \bar{z}') \rangle + \gamma^2 \langle \delta^2(z' - z) \rangle$$

$$\frac{1}{\pi} \int_{\mathbb{C}} \bar{z} \langle X^m(z, \bar{z}) X^n(z', \bar{z}') \rangle = \int_{\mathbb{C}} \delta^2(z - z')$$

Normal ordering

$$= \frac{1}{\mathcal{N}} \langle \partial_{\bar{z}} \bar{\psi}_2 | X^m(z, \bar{z}) X(\bar{z}', \bar{z}') \rangle + \gamma^2 \langle \delta^2(z' - z) \rangle$$

Normal ordering

$$\begin{aligned} \partial \bar{\psi} X^m |_{\text{classically}} &= 0 \\ \partial \bar{\psi} X^m |_q &\neq 0 \end{aligned}$$

$$\partial_{\bar{z}} \chi^{\text{hol}}(z_1, \bar{z}_1) \chi^{\text{ant}}(z_2, \bar{z}_2)$$



$$\partial_{\bar{z}} \chi^m(z_1, \bar{z}_1) \chi^m(z_2, \bar{z}_2) = - \pi \delta^2(z_{12})$$

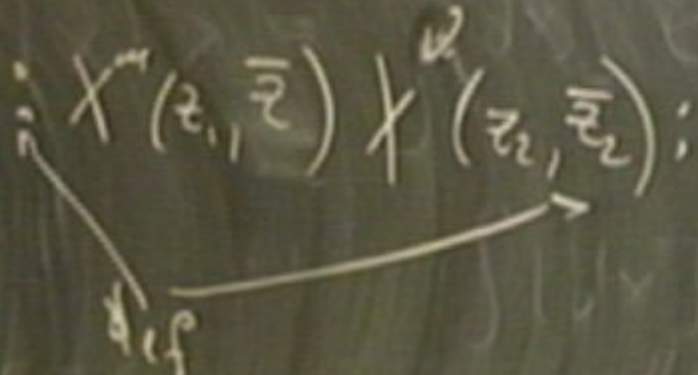
$$z_{12} \equiv z_1 - z_2$$

$$\partial_{\bar{z}_1} \bar{\chi}^m(z_1, \bar{z}_1) \chi^m(z_2, \bar{z}_2) = - \underbrace{\frac{\pi \alpha'}{4} \eta^{\mu\nu} \delta^2(z_{12})}_{\text{Quantum correction to classical FOM}} \quad z_{12} \equiv z_1 - z_2$$

$$\partial_{\bar{z}_1} \bar{\chi}^m(z_1, \bar{z}_1) \cdot \chi^m(z_2, \bar{z}_2) = - \underbrace{\frac{\pi \alpha' \hbar^2}{2} \delta^2(z_{12})}_{\text{Quantum correction to classical FOM}}$$

$z_{12} \equiv z_1 - z_2$

$$\partial_{\bar{z}_1} \bar{X}^m(z_1, \bar{z}_1) \cdot X^{n, \psi}(z_2, \bar{z}_2) = - \underbrace{\pi \alpha' \hbar^{-1} \delta^2(z_{12})}_{\text{Quantum correction to classical FOM}} \quad z_{12} \equiv z_1 - z_2$$



$$\partial_{\bar{z}_1} \bar{\psi}(z_1, \bar{z}_1) \cdot \psi(z_2, \bar{z}_2) = - \underbrace{\frac{\pi \hbar^2}{\hbar^2} \delta^2(z_{12})}_{\text{Quantum correction}}$$

$z_{12} \equiv z_1 - z_2$

$$\partial_{\bar{z}_1} \bar{\psi}(z_1, \bar{z}_1) \psi(z_2, \bar{z}_2) = 0$$

to classical FOM

def

Def

$$\vdash X^{\text{any}} \vdash = X^{\text{any}}$$

Def

$$: X^{m_1} : = X^{m_1}$$

$$: X^{m_1} X^{m_2} : = X^{m_1} X^{m_2}$$

Def

$$: X^{\alpha} : = X^{\alpha}$$

$$: X^{\alpha} X^{\beta} : = X^{\alpha} X^{\beta} + \frac{\alpha!}{2} \eta^{\alpha\beta} p_n |z, z|^2$$

Def

$$: X^{\mu} : = X^{\mu}$$

$$: X^{\mu} X^{\nu} : = X^{\mu} X^{\nu} + \frac{\alpha'}{2} \eta^{\mu\nu} \ln |z_1, z_2|^2$$

$$\partial_{\bar{z}_1} : X^{\mu} X^{\nu} : = -\frac{\alpha'}{2} \eta^{\mu\nu}$$

Def

$$: X^m : = X^m$$

$$: X^m X^n : = X^m X^n + \frac{d^1}{2} \eta^{mn} |z, z|^2$$

$$z, \bar{z} : X^m X^n : = -\frac{d^1}{2} \eta^{mn} \delta^{(z, z)}$$

Def

$$: X^m : = X^m$$

$$: X^m X^n : = X^m X^n + \frac{\alpha'}{2} \eta^{mn} h |z, z|^2$$

$$\partial_{\bar{z}_1} : X^m X^n : = -\frac{\alpha'}{2} \eta^{mn} \delta^{(z, z)} + \frac{\alpha'}{2} \eta^{mn} \partial_{\bar{z}_1} h |z, z|^2$$

Because : $\oint_{\gamma} \ln |z|^2 = 2\pi \delta^2(z)$

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Proof

If $z \neq 0$

$$\oint_{\gamma} [\ln z + \ln \bar{z}] =$$

Because : $\partial_{\bar{z}} \ln |z|^2 = 2\pi \delta^2(z)$

Proof

If $z \neq 0$

$$\partial_{\bar{z}} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial_{\bar{z}} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

Because : $\partial\bar{\partial} \ln |z|^2 = 2\pi \delta^2(z)$

Proof

If $z \neq 0$

$$\partial\bar{\partial} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial\bar{\partial} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

$$\int d^2z \partial\bar{\partial} \ln |z|^2 =$$

Because: $\partial\bar{\partial} \ln |z|^2 = 2\pi \delta^2(z)$

Proof

If $z \neq 0$

$$\partial\bar{\partial} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial\bar{\partial} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

$$\text{LHS} \int d^2z \partial\bar{\partial} \ln |z|^2 =$$

$$\text{RHS} \int 2\pi \delta^2 - d^2z = 2\pi$$

Because: $\partial\bar{\partial} \ln |z|^2 = 2\pi \delta^2(z)$

Proof

If $z \neq 0$

$$\partial\bar{\partial} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial\bar{\partial} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

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Proof

If $z \neq 0$

$$\partial_{\bar{z}} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial_{\bar{z}} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

$$\text{LHS} \int d^2z \partial_{\bar{z}} \ln |z|^2 = \int d^2z \partial_{\bar{z}} \ln z$$

$\partial_{\bar{z}} \ln z = \partial_{\bar{z}} \ln |z|^2$

$$\text{RHS} \int 2\pi \delta^2(z) d^2z = 2\pi$$

Because : $\partial_{\bar{z}} \ln |z|^2 = 2\pi \delta^2(z)$

Proof

If $z \neq 0$

$$\partial_{\bar{z}} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial_{\bar{z}} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

$$\text{LHS} \int d^2z \partial_{\bar{z}} \ln |z|^2 = \int d^2z \partial_{\bar{z}} \ln z$$

$$\partial_{\bar{z}} \ln z = \partial_{\bar{z}} [\ln z + \ln \bar{z}] = \frac{1}{z}$$

$$\text{RHS} \int 2\pi \delta^2(z) d^2z = 2\pi$$

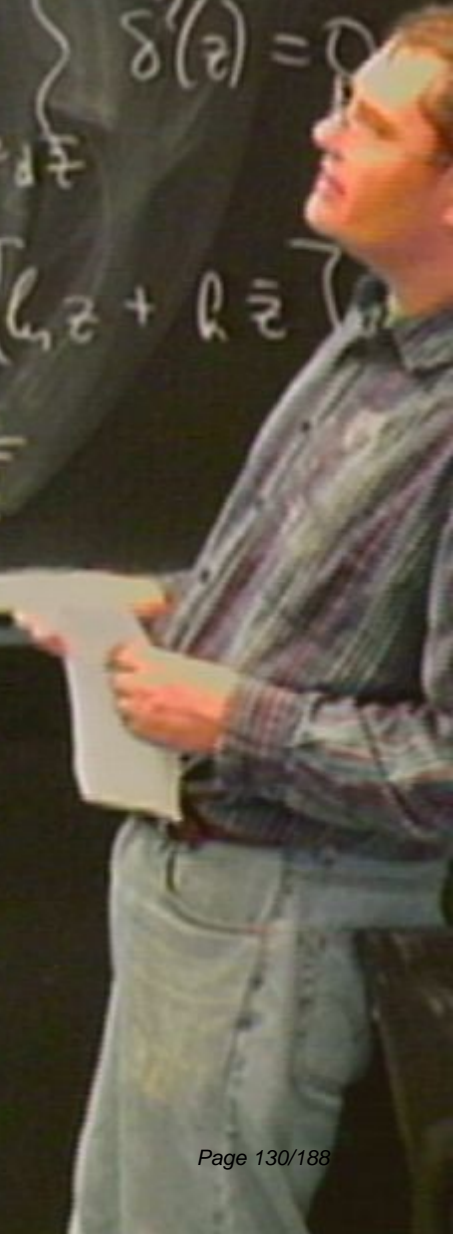
Because: $\partial_{\bar{z}} \ln |z|^2 = 2\pi \delta^2(z)$

Proof If $z \neq 0$

$$\partial_{\bar{z}} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial_{\bar{z}} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

$$\text{LHS} \int_{|z| < 1} d^2z \partial_{\bar{z}} \ln |z|^2 = \int_{|z| < 1} d^2z \partial_{\bar{z}} \circlearrowleft = \int_{|z|=1} \frac{1}{|z|} \partial_{\bar{z}} [\ln z + \ln \bar{z}] = \int_{|z|=1} \frac{1}{z} dz = \frac{1}{2\pi i} \int_{|z|=1} \frac{dz}{z} = 1$$

$$\text{RHS} \int d^2z 2\pi \delta^2(z) = 2\pi$$



Because : $\partial_{\bar{z}} \ln |z|^2 = 2\pi \delta^2(z)$

Proof

If $z \neq 0$

$$\partial_{\bar{z}} [\ln z + \ln \bar{z}] = 0 \quad \left. \vphantom{\partial_{\bar{z}} [\ln z + \ln \bar{z}]} \right\} \delta^2(z) = 0$$

$$\text{LHS} \int_{|z| < 1} d^2z \partial_{\bar{z}} \ln |z|^2 = \int_{|z| < 1} d^2z \partial_{\bar{z}} \ln z = \int_{|z|=1} \frac{d\bar{z}}{z} = 2\pi i$$

$$\text{RHS} \int 2\pi \delta^2(z) d^2z = 2\pi$$

$$0 = \int [dx] \frac{\delta}{\delta X^m} \left[e^{-S} F(z, \bar{z}) \right] = - \int [dx] e^{-S} \frac{\delta S}{\delta X^m}$$

$F = 0$

$$= - \int [dx] e^{-S} \left(\frac{z}{2\pi\alpha'} \partial \bar{X}^m \right) = \frac{1}{\pi\alpha'} \langle \partial \bar{X}^m | 0 \rangle$$

$$\langle \partial \bar{X}^m \rangle_F = 0$$

$\partial \bar{X}^m = 0$ as an operator statement

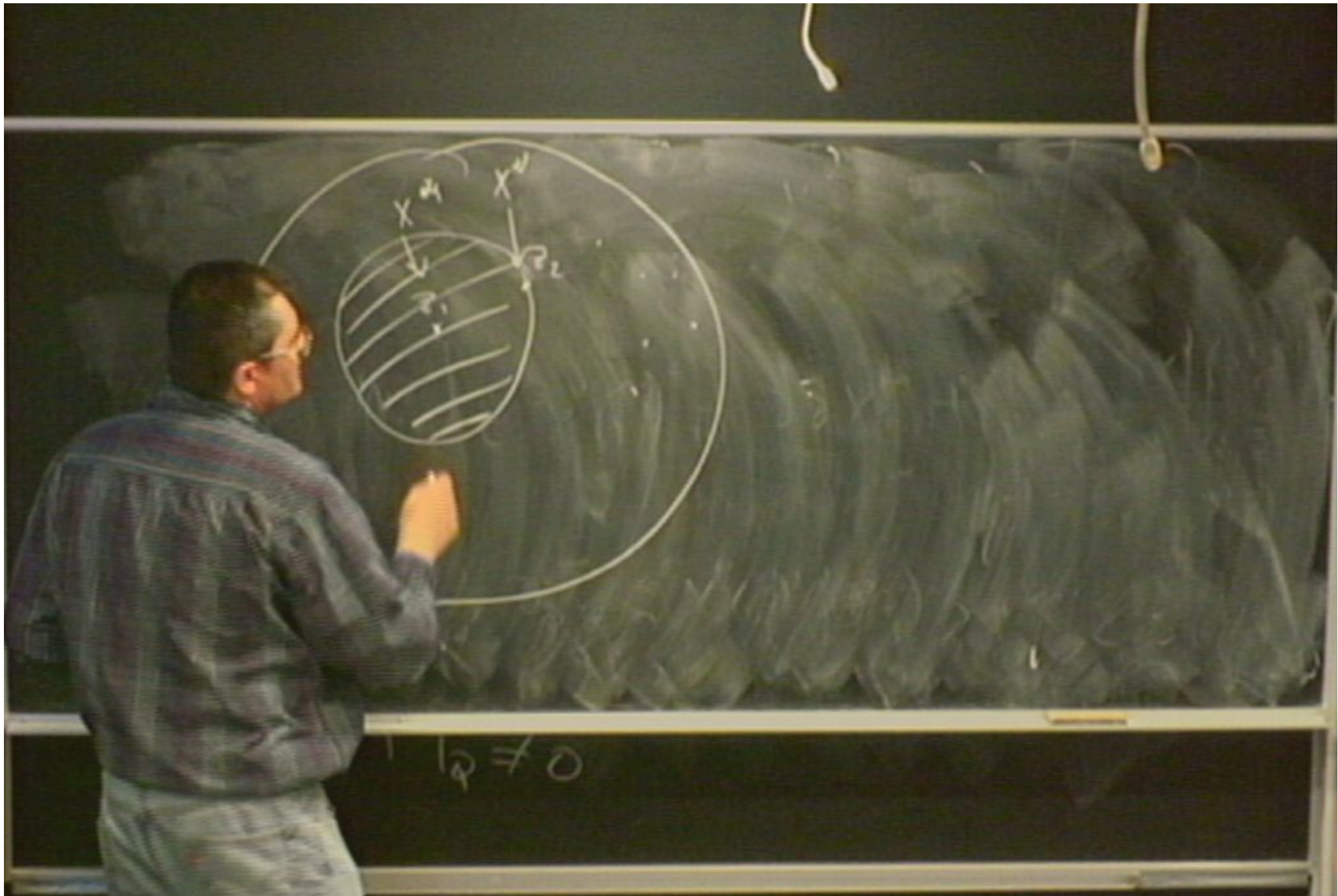
Def

$$:X^m: = X^m$$

$$:X^m X^j: = X^m X^j + \frac{\alpha'}{2} \eta^{mj} \ln |z_1, z_2|^2$$

$$\partial_i \bar{\partial}_i :X^m X^j: = -\frac{\alpha'}{2} \eta^{mj} \delta^i(z_1, z_2) + \frac{\alpha'}{2} \eta^{mj} \partial_i \bar{\partial}_i \ln |z_1, z_2|^2 = 0$$





$$|Q| \neq 0$$



$$\langle A_i | A_j(z_n, \bar{z}_n) \rangle = ?$$

$$\langle A_i(z_1, \bar{z}_1) \dots A_i(z_n, \bar{z}_n) \rangle = ?$$

↑
correlation functions are observables in CFT

$$\langle A_i(z_i, \bar{z}_i) \cdots A_j(z_j, \bar{z}_j) \rangle = ?$$

↑
correlation functions are observables in CFT

$$z_{ij} \rightarrow 0$$

$$\langle A_i(z_1, \bar{z}_1) \dots A_i(z_n, \bar{z}_n) \rangle = ?$$

↑ correlation functions are observables in CFT

$z_{ij} \rightarrow 0$

↑ interested in singularities

Def

$$:X^m: = X^m$$

$$:X^m X^d: = X^m X^d + \frac{\alpha'}{2} \eta^{md} \ln |z, \bar{z}|^2$$

$$i\partial_{\bar{1}} :X^m X^d: = -\frac{\alpha'}{2} \eta^{md} \delta^{\langle z, \bar{z} \rangle} + \frac{\alpha'}{2} \eta^{md} \partial_{\bar{1}} \bar{\partial}_{\bar{1}} \ln |z, \bar{z}|^2 = 0$$

OPE

$$A_1(\sigma_1) A_2(\sigma_2) =$$

OPE

$$\langle A_i(\sigma_1) A_j(\sigma_2) \rangle \sim \sum_k c_{ij}^k(\sigma_1 - \sigma_2) \langle A_k(\sigma_1) \rangle$$

OPE

$$\langle A_1(\sigma_1) A_2(\sigma_2) \dots \rangle \sim \sum_k C_{ij}^k(\sigma_1 - \sigma_2) \langle A_k(\sigma_1) \rangle$$

OPE

$$\langle A_1(\sigma_1) A_2(\sigma_2) \dots \rangle \sim \sum_k c_{ij}^k(\sigma_1 - \sigma_2) \langle A_k(\sigma_1) \rangle$$

Typically

OPE

$$\langle A_1(\sigma_1) A_2(\sigma_2) \dots \rangle \sim \sum_k c_{ij}^k (\sigma_1 - \sigma_2) \langle A_k(\sigma_1) \rangle$$

Typically $c_{ij}^k \propto (r_{12})^k$

QFT

$$\langle A_1(\sigma_1) A_2(\sigma_2) \dots \rangle$$

$$\sum_k c_{ij}^k(\sigma_1 - \sigma_2) \langle A_k(\sigma_1) \rangle$$

Typically $c_{ij}^k \propto (r_{12})^k$

In normal QFT

QFT

$$\langle A_1(\sigma_1) A_2(\sigma_2) \dots \rangle$$

$$\sum_k c_{ij}^k (\sigma_1 - \sigma_2)^k$$

$$\langle A_k(\sigma_1) \rangle$$

Typically

$$c_{ij}^k \propto (r_{12})^k$$

In normal QFT

radius of convergence $\neq 0$

Typically, $\sum_{k=0}^{\infty} \varphi(z_1 z_2)^k$ | In normal QFT
 In CFT's OPEs have finite radius of convergence! $\neq 0$
convergence!

$$\partial_{z_1} : X^{\sim} X^{\vee} : = -\frac{1}{4\pi\alpha'} \eta^{\mu\nu} \delta(z_1, z_2) + \frac{\alpha'}{2} \eta^{\mu\nu} \partial_{z_1} h_{\mu\nu} |z_1, z_2|^2 = 0$$

Typically: $\sum_{k=0}^{\infty} \varphi(r_{12})^k$

In normal QFT
radius of convergence $\neq 0$

In CFT's OPEs have finite radius of

convergence!
 $O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \dots$

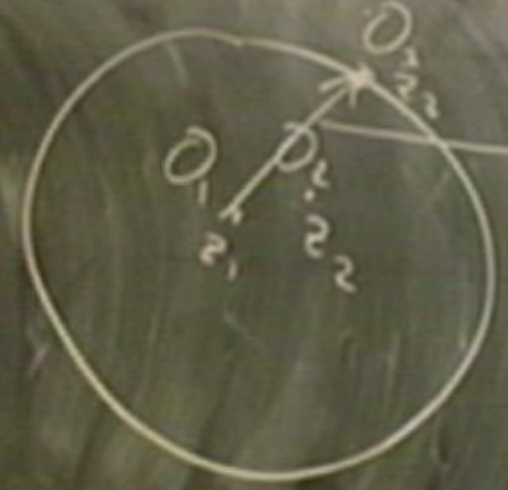
$$\partial_1 \bar{\partial}_1 : X^{\mu} X^{\nu} : = -\frac{1}{\pi \alpha'} \eta^{\mu\nu} \delta(z_1, z_2) + \frac{\alpha'}{2} \eta^{\mu\nu} \partial_1 \bar{\partial}_1 \ln |z_1 - z_2|^2 = 0$$

$$\sum_{k=1}^n \binom{n}{k} \binom{n-k}{k} z^k$$

$$\sum_k \begin{pmatrix} 1 & \\ & 2 \end{pmatrix} \mathcal{O}_{1k}(z_1) \mathcal{O}_{1k}(z_2)$$



$$\sum_k \left(\frac{z_1}{z_2} \right)^k \mathcal{O}_k(z_2)$$



radius of convergence

of z_1, z_2 OPE

$$\langle A_1(\sigma_1) A_2(\sigma_2) \dots \rangle \sim \sum c_{ij}^k (\sigma_1 - \sigma_2) \langle A_k(\sigma_1) \rangle$$

Typically

$$c_{ij}^k \propto \mathcal{Z}(\sigma_{12})^k$$

In normal QFT
radius of convergence $\neq 0$
radius of

In CFT's OPEs have

$$O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \frac{\text{const}}{\dots}$$

$$\langle A_1(\sigma_1) A_2(\sigma_2) \cdot B \rangle \left(\sum_k \binom{k}{j} (\sigma_1 - \sigma_2)^k \langle A_k(\sigma_1) B \rangle \right)$$

Typically, $\sum_{j=0}^k \binom{k}{j} \rho(r_{12})^k$

In normal QFT radius of convergence is ∞

In CFT's OPEs have finite radius of

$$O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \frac{\text{convergent!}}{}$$

$$\partial_1 \partial_1 : X''(1) X''(2) = 0$$

$$\alpha, \beta : X^{\alpha}(1) X^{\beta}(2) \dots = 0$$

$$\alpha \int \beta$$

$$\alpha_1 \alpha_2 : x''(1) x'(2) = 0$$

$$\alpha_1 [\alpha_1 : x'' x' :] = 0$$

$$\alpha_2 [: \alpha_2 x'' x' :] = 0$$

$$\partial, \bar{\partial}_1 : X''(z) X'(z) = 0$$

$$\partial_1 [\bar{\partial}_1 : X'' X'] = 0$$

$$\textcircled{\partial_1} [: \partial_1 X'' X'] = 0$$

$\partial_1 X''(z, \bar{z}) X'(z, \bar{z})$ is a holomorphic function

$$\partial_1 [: \partial_1 X^\mu X^\nu :] = 0$$

$$\textcircled{2} [: \partial_1 X^\mu X^\nu :] = 0$$

$: \partial_1 X^\mu(z_1, \bar{z}_1) X^\nu(z_2, \bar{z}_2) :$ is a holomorphic function

$: \partial_3 X^\mu(z_1 + \xi, \bar{z}_1 + \bar{\xi}) X^\nu(z_2, \bar{z}_2) :$

$\xi = z_{12}$

correlation functions are observables in CFT

$z_{ij} \rightarrow 0$

↑ interested in singularities

$$\textcircled{2} \quad \langle \partial_1 X^m(z_1) X^j(z_2) \rangle = 0$$

$$\langle \partial_1 X^m(z_1, \bar{z}_1) X^j(z_2, \bar{z}_2) \rangle \text{ is a holomorphic function}$$

$$\langle \partial_1 X^m(z_1 + \epsilon, \bar{z}_1 + \bar{\epsilon}) X^j(z_2, \bar{z}_2) \rangle = \sum_{k=1}^{\infty} \frac{\epsilon^k}{k!} \partial_1^k X^m(z_1, \bar{z}_1) X^j(z_2, \bar{z}_2)$$

$\epsilon = z_{12}$

correlation functions are observables in CFT

$z_{ij} \rightarrow 0$
 ↑ interested in singularities

$\partial_1 X(z_1, \bar{z}_1) X'(z_2, \bar{z}_2)$ is a holomorphic function
 $\partial_3 X''(z_1 + \xi, \bar{z}_1 + \bar{\xi}) X'(z_1, \bar{z}_1) = \sum_{k=1}^{\infty} \frac{\xi^k}{k!} \partial^k X''(z_1, \bar{z}_1)$

$= \sum_{k=1}^{\infty} \frac{z_{12}^k}{k!} \partial^k X''(z_1, \bar{z}_1) X'(z_1, \bar{z}_1)$

$\partial_1 X(z_1, \bar{z}_1) X'(z_2, \bar{z}_2)$ is a holomorphic function
 $\partial_3 X''(z_1 + \xi, \bar{z}_1 + \bar{\xi}) X'(z_1, \bar{z}_1) = \sum_{k=1}^{\infty} \frac{\xi^k}{k!} : X' \partial^k X'' :$
 $\xi = z_{12}$

$$= \sum_{k=1}^{\infty} \frac{z_{12}^k}{k!} : X'(z_1, \bar{z}_1) \partial^k X''(z_1, \bar{z}_1) :$$

$$: X(z_1, \bar{z}_1) X'(z_1, \bar{z}_1) : = : X(z_1, \bar{z}_1) X''(z_1, \bar{z}_1) :$$

$$\textcircled{2} \quad \left[: \partial, X^m X^j : \right] = 0 \quad \left\{ \quad \right. \quad \left[: \bar{\partial}, X^m X^j : \right] = 0$$

$: \partial, X^m(z_1, \bar{z}_1) X^j(z_2, \bar{z}_2) :$ is a holomorphic func
 $: \partial, X^m(z_1 + \epsilon, \bar{z}_1 + \bar{\epsilon}) X^j(z_2, \bar{z}_2) : = \sum_{k=0}^{\infty} \frac{\epsilon^k}{k!} : X^j \partial^k X^m :$
 $\epsilon = z_{12}$

$$= \sum_{k=0}^{\infty} \frac{z_{12}^k}{k!} : X^j(z_2, \bar{z}_2) \partial^k X^m(z_1, \bar{z}_1) :$$

$$: X^j(z_1, \bar{z}_1) X^m(z_2, \bar{z}_2) : = : X^j(z_2, \bar{z}_2) X^m(z_1, \bar{z}_1) :$$

$\partial_1 X(z_1, \bar{z}_1) X'(z_2, \bar{z}_2)$ is a holomorphic function
 $\partial_3 X''(z_1 + \xi, \bar{z}_1 + \bar{\xi}) X'(z_2, \bar{z}_2) = \sum_{k=1}^{\infty} \frac{\xi^k}{k!} : X' \partial^k X'' :$
 $\xi = z_{12}$

$$= \sum_{k=1}^{\infty} \frac{z_{12}^k}{k!} : X'(z_1, \bar{z}_1) \partial^k X''(z_2, \bar{z}_2) :$$

$$: X'(z_1, \bar{z}_1) X''(z_2, \bar{z}_2) = : X'(z_2, \bar{z}_2) X''(z_1, \bar{z}_1) :$$

$$+ \sum_{k=1}^{\infty} \left\{ \frac{z_{12}^k}{k!} : X'' \partial^k X' : + \frac{\bar{z}_{12}^k}{k!} : X'' \partial^k X' : \right.$$

$$O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \sum c_k A_k(z_2, \bar{z}_2)$$

$$O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \sum c^k A_k(z_2, \bar{z}_2)$$

\updownarrow
 x_1

\updownarrow
 x_2

\updownarrow

1 2 1

$$O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \sum c^k A_k(z_2, \bar{z}_2)$$

\updownarrow
 x^m

\updownarrow
 x^m

\updownarrow

$$\left\{ x^2 \partial^k x^m; x^{2-k} x^m \right\}$$

$$O_1(z_1, \bar{z}_1) O_2(z_2, \bar{z}_2) = \sum c^k A_k(z_2, \bar{z}_2)$$

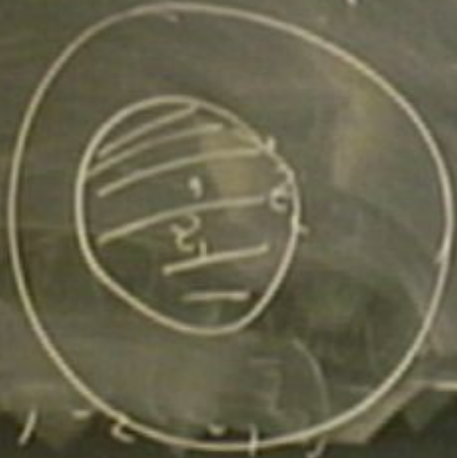
\Downarrow
 X^m

\Downarrow
 X^m

$\frac{z_{12}^k}{k!}$

\Downarrow

$$\left\{ X^j \partial^k X^m; X^j \partial^{-k} X^m \right\}$$



$$: X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n) :$$

$$: X^{u_1}(z_1, \bar{z}_1) \cdot X^{u_n}(z_n, \bar{z}_n) : = X^{u_1}(z_1, \bar{z}_1) \cdot X^{u_n}(z_n, \bar{z}_n) + \sum$$

$$: X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n) : = X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n) +$$

$$+ \sum \text{subtractions}$$

subtraction =

$$: X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n) : = X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n) +$$

+ \sum subtractions

subtraction = replace a pair of fields

$$\frac{1}{z} \alpha' \eta^{\mu\nu} p_\mu |z_{ij}|^2 \dots$$

$$0 = \int d[X] \frac{\delta}{\delta X_n(z_n, \bar{z}_n)} \left[e^{-S} X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n) \right]$$

$$0 = \int d[X] \frac{\delta}{\delta X_n(z_n, \bar{z}_n)} \left[e^{-S} \underbrace{X^{m_1}(z_1, \bar{z}_1) \cdots X^{m_n}(z_n, \bar{z}_n)}_{n-1 \text{ field}} \right]$$

$$:X^{m_1} X^{m_2} \cdots X^{m_n}: = X^{m_1} X^{m_2} \cdots X^{m_n} - \frac{1}{2} \sum_{i < j} \eta_{m_i m_j} \ln |z_i - z_j|^2$$

$$: X^{\mu_1}(z_1, \bar{z}_1) \cdots X^{\mu_n}(z_n, \bar{z}_n) : = \underbrace{X^{\mu_1}(z_1, \bar{z}_1) \cdots X^{\mu_n}(z_n, \bar{z}_n)} +$$

+ \sum subtractions

subtraction = replace a pair of fields

$$\frac{1}{2} \alpha' \eta^{\mu\nu} \frac{p_\mu p_\nu}{|z_{ij}|^2} :$$

$$\vec{F} = \vec{F}(x_{w1}, \dots, x_{wk})$$



$$\vec{f} = \vec{f}(x_{w_1}, \dots, x_{w_k})$$

$$\vec{f} := \exp \left[\right]$$



$$\mathcal{F} = \mathcal{F}(X_{w_1}, \dots, X_{w_n})$$

$$\therefore \mathcal{F} := \exp \left[\frac{\delta}{4} \int d^2z_1 d^2z_2 \ln|z_1 - z_2|^2 \frac{\delta}{X^w(z_1, \bar{z}_1)} X_u(z_2, \bar{z}_2) \right]$$

$$\mathcal{F} = \mathcal{F} [X_{w_1}, \dots, X_{w_k}]$$

$$\therefore \mathcal{F} := \exp \left[\frac{\delta'}{\delta} \int d^2z_1 d^2z_2 \ln |z_1 z_2| \frac{\delta}{\delta X^{\mu\nu}(z_1, \bar{z}_1)} X_\nu(z_2, \bar{z}_2) \right] \mathcal{F}$$

$$\mathcal{F} \equiv \exp \left[-\frac{\mathcal{L}'}{4} \left(d^2 z_1, d^2 z_2, \sum \frac{\delta}{\delta X_{\mu}^{(1)}}, \sum \frac{\delta}{\delta X_{\mu}^{(2)}} \right) \right] \cdot \mathcal{F}$$



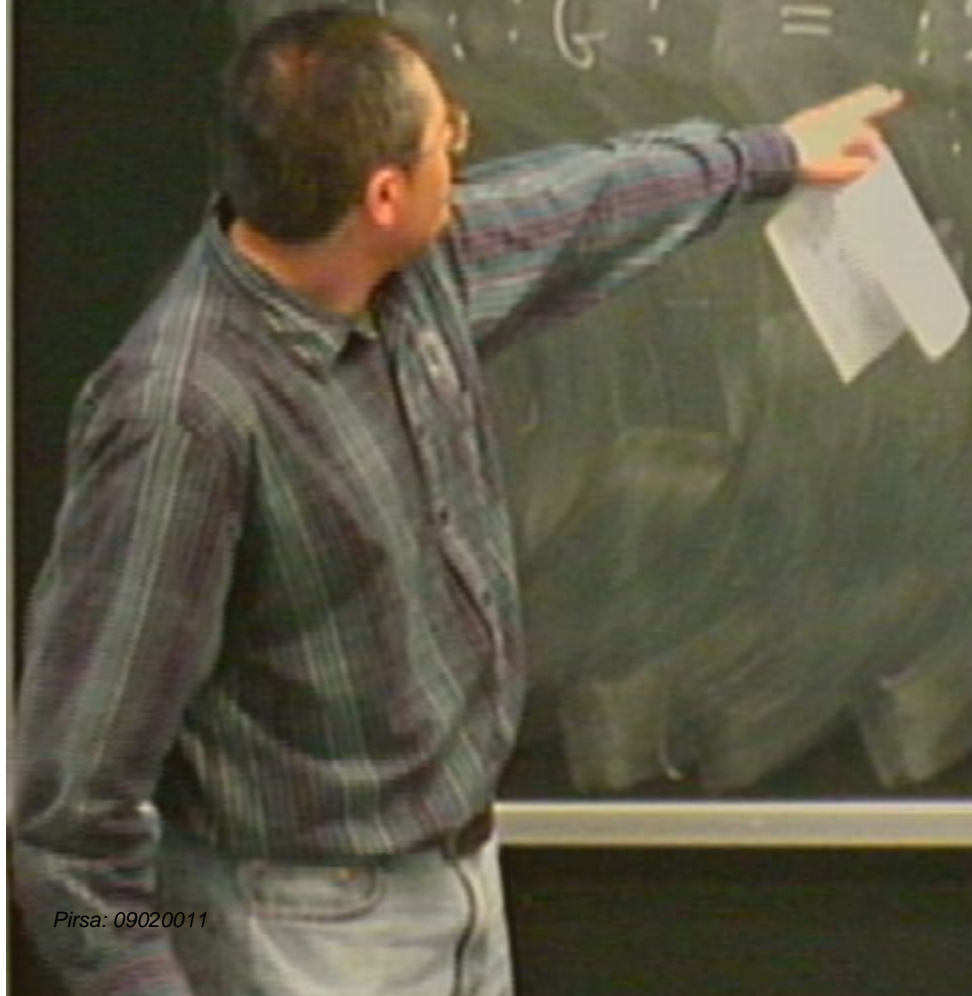
$$\mathcal{F} = \exp \left[-\frac{\alpha'}{4} \int d^2 z_1 d^2 z_2 \sum_{\mu} \sum_{\nu} \frac{\delta X_{\mu}^{\nu}(1)}{\delta X_{\nu}^{\mu}(2)} \right] \mathcal{F}$$

$$= : \mathcal{F} : + \sum \text{contractions}$$

$$\text{contraction} = -\frac{1}{2} \alpha' \eta^{\mu\nu} \ln |z_{12}|^2$$

contraction = $\bar{z} \cdot \frac{1}{|z|^2} \cdot \frac{1}{|z|^2}$

$$T : G \rightarrow G$$



contraction = $\bar{z}_2 \psi_{\alpha} \cdot \psi_{\beta} / |z_{12}|^2$

$$T : : G : = : T G : + \sum \text{cross-contractions}$$

$$\text{contraction} = \frac{1}{2} \int d^2z_1 d^2z_2 \ln |z_{12}|^2$$

$$:F::G: = :FG: + \sum \text{cross-contractions}$$

$$:F::G: = \exp \left[-\frac{\alpha'}{2} \int d^2z_1 d^2z_2 \ln |z_{12}|^2 \right]$$

$$\text{contraction} = \frac{1}{2} \int d^2z_1 d^2z_2 \ln |z_{12}|^2$$

$$:\mathcal{F}::\mathcal{G}: = :\mathcal{F}\mathcal{G}: + \sum \text{cross-contractions}$$

$$:\mathcal{F}::\mathcal{G}: = \exp \left[-\frac{\alpha'}{2} \int d^2z_1 d^2z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X_F^\mu(z_1, \bar{z}_1)} \frac{\delta}{\delta X_F^\nu(z_2, \bar{z}_2)} \right]$$

$$\text{contraction} = \bar{z}_1 \delta_{12} \cdot \ln |z_{12}|^2$$

$$:\mathbb{T}G: = :T\mathbb{G}: + \sum \text{cross-contractions}$$

$$:G: = \exp \left[-\frac{\alpha'}{2} \int d^2z_1 d^2z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X_F^{\mu}(z_1, \bar{z}_1)} \frac{\delta}{\delta X_G^{\nu}(z_2, \bar{z}_2)} \right]$$

$$x :T\mathbb{G}:$$

$$\text{contraction} = \frac{1}{2} \int d^2 z_1 d^2 z_2 \langle \psi(z_1) \psi(z_2) \rangle$$

$$:\mathcal{F}::\mathcal{G}: = :\mathcal{F}\mathcal{G}: + \sum \text{cross-contractions}$$

$$:\mathcal{F}::\mathcal{G}: = \exp \left[-\frac{\alpha'}{2} \int d^2 z_1 d^2 z_2 \langle \psi(z_1) \psi(z_2) \rangle \frac{\delta}{\delta X_F(z_1, \bar{z}_1)} \frac{\delta}{\delta X_G(z_2, \bar{z}_2)} \right] x : \mathcal{F}\mathcal{G} :$$