Title: Why there is no information loss

Date: Dec 04, 2008 11:00 AM

URL: http://pirsa.org/08120016

Abstract: Using 2-dimensional CGHS black holes, I will argue that information is not lost in the Hawking evaporation because the quantum space-time is significantly larger than the classical one. I will begin with a discussion of the conceptual underpinnings of problem and then introduce a general, non-perturbative framework to describe quantum CGHS black holes. I will show that the Hawking effect emerges from it in the first approximation. Finally, I will introduce a mean field approximation to argue that, when the back reaction is included, future null infinity is `long enough\' to capture full information contained in pure states at past null infinity and that the S-matrix is unitary. There are no macroscopic remnants.

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Why there is no information loss

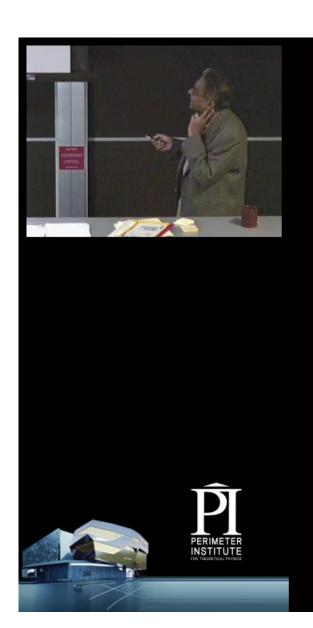
Because the quantum space-time is sufficiently larger than the classical.

Goal: A space-time resolution of the information-loss issue. A general paradigm and realization in CGHS BHs

organization:

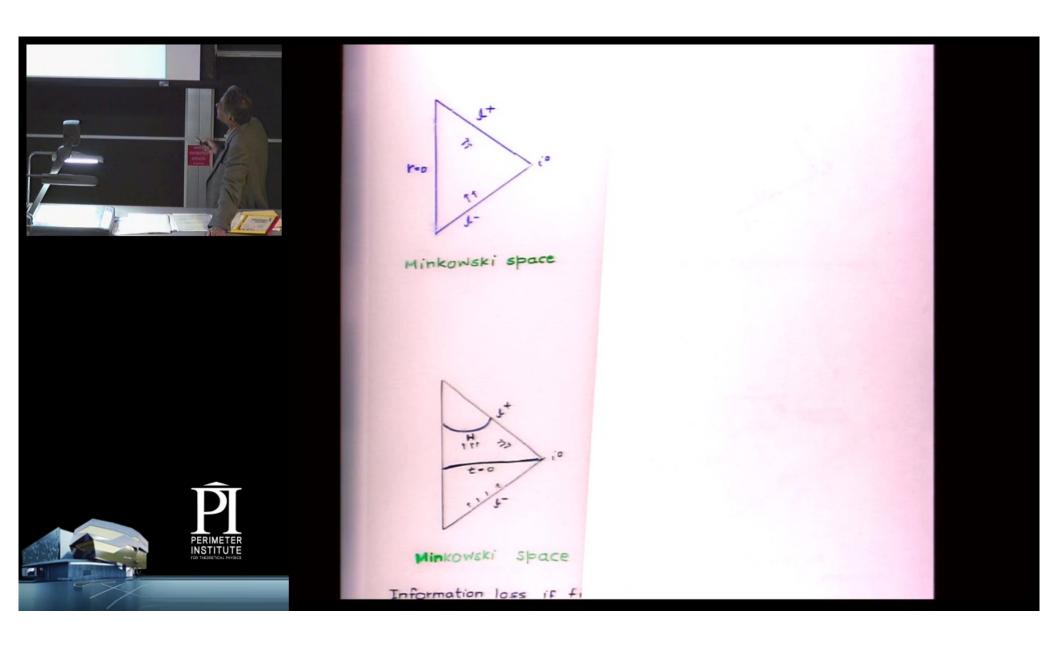
1. Introduction & Conceptual Setting

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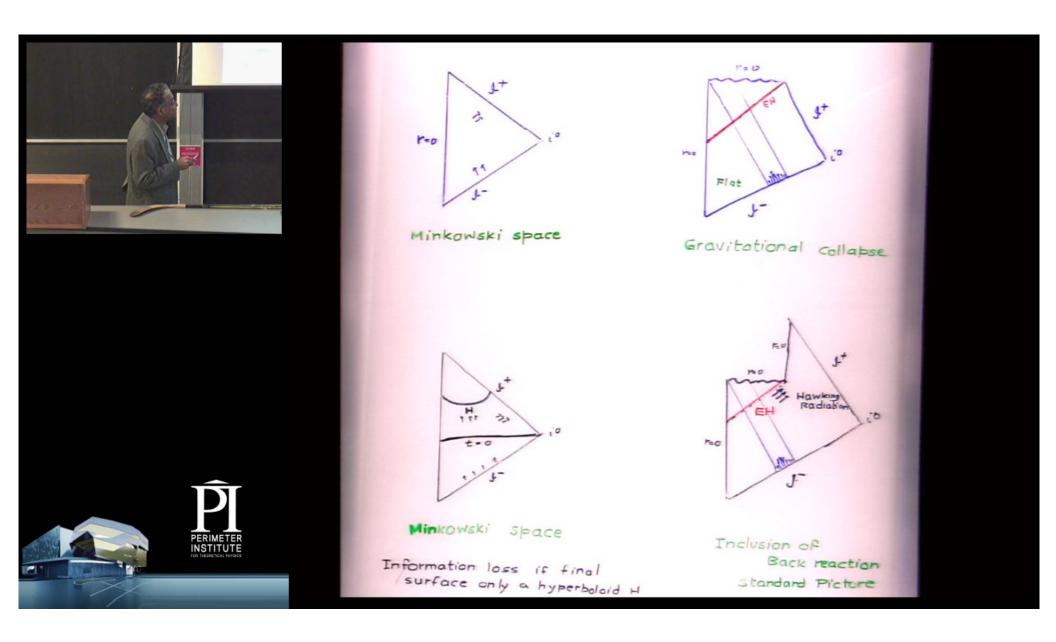


Goal: A space-time resolution of the information-loss issue · A general paradigm and realization in CGHS BHs organization: 1. Introduction & Conceptual Setting 2. CGHS BHs : Information Recovery 3 Summary a Outlook Joint Work With Victor Tavares & Madhavan Vandarajan (PRL '08); Motivation from AAA Bojowald (CQG '05)

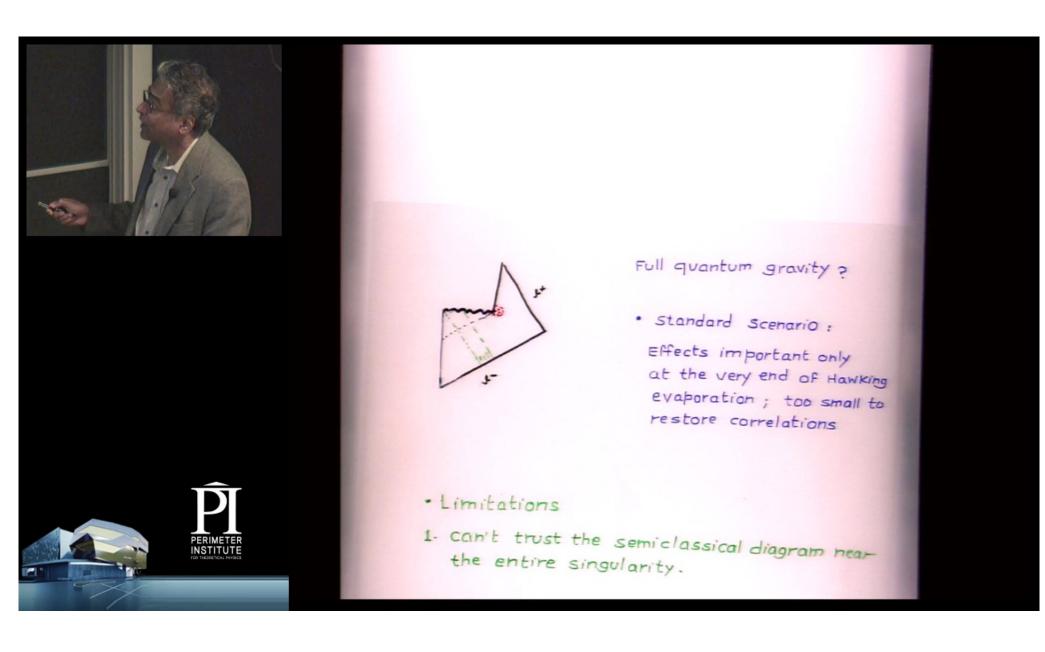
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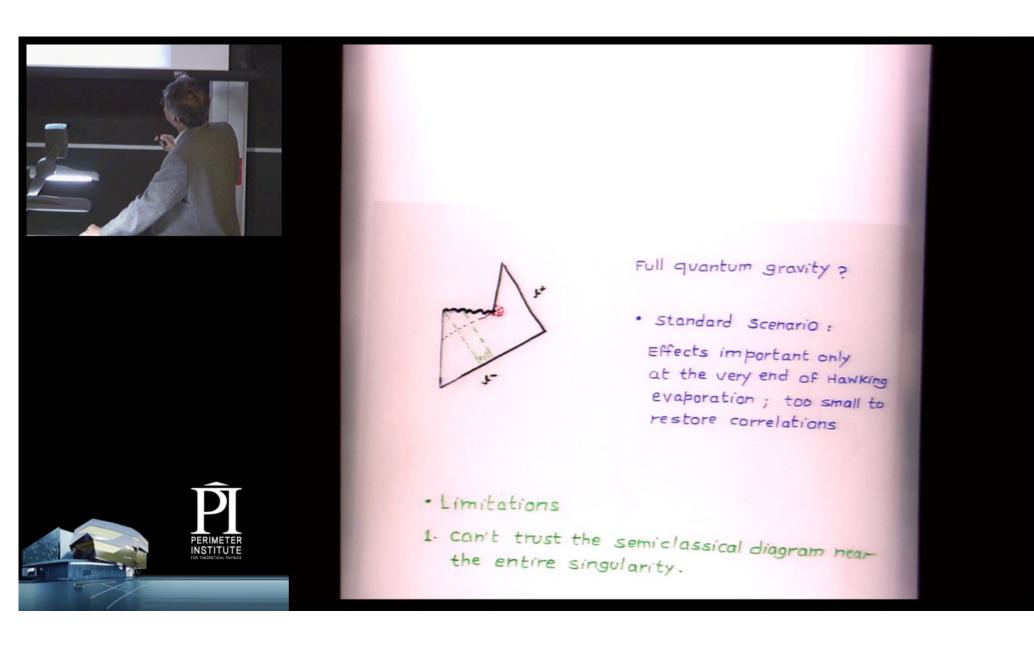
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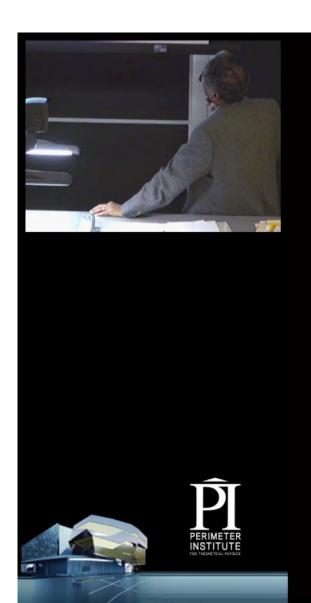


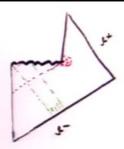
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Standard Scenario:
 Effects important only
 at the very end of Hawking
 evaporation; too small to
 restore correlations

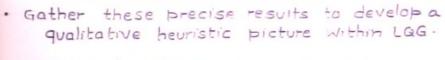
- · Limitations
- 1. can't trust the semiclassical diagram near the entire singularity.
- 2: Event horizon too global a concept.
 (e.g. Hajicék argument)

If singularity is resolved, event horizon may disappear completely or 'transcended'.

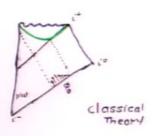
What then is a BH? What evaporates?

Well-defined quasi-local notion of Dynamical

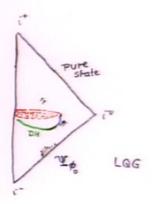




· spherical collapse of a scalar field: Large BH



LQG: strong Indications that quantum space-time does not end at singularity. > NO Sink of information!



Pure states on J evolve to pure states on Jt of the larger, quantum Space-time.







a strominger)

• Spherically Symmetric GR with massless f $49_{ab} = 9_{ab} + e^{2\phi} k^{-2} S_{ab}$ [k] = [1]

$$S^{(4)} = \frac{1}{G} \int d^2x \sqrt{g_1} \left[e^{-2\phi} \left(R + 2 \sqrt{a\phi} \nabla_a \phi + 2 e^{+2\phi} \kappa^2 \right) \right]$$
$$-\frac{1}{2} \int d^2x \sqrt{g_1} e^{-2\phi} \sqrt{af} \nabla_a f$$

· Dilaton Gravity : 1+1 dim, 9ab, \$, \$

$$S = \frac{1}{G} \int d^2x \sqrt{191} \left[e^{2\phi} \left(R + 4 \nabla^2 \phi \nabla_0 \phi + 4 \kappa^2 \right) \right]$$
$$-\frac{1}{2} \int d^2x \sqrt{191} \nabla^2 f \nabla_0 f$$

· Both cases true degree of freedom : f

But equs simpler in CGHS

in particular:
$$\Box f = 0$$
 $\Leftrightarrow \Box f = 0$

$$g^{ab} = \Omega \eta^{ab}$$
refers to η^{ab}
Dilatonic geometry: Ω and $\overline{D} := e^{-2\phi}$

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* Spherically Symmetric GR with massless f

$$4g_{ab} = g_{ab} + e^{2\phi} \kappa^{2} S_{ab} \qquad [k] = \Gamma^{1}$$

$$S^{(k)} = \frac{1}{G} \int_{0}^{2} x \sqrt{s} \left[e^{2\phi} (R + 2 \sqrt{a} \phi \sqrt{a} \phi + 2 e^{2\phi} \sqrt{a}^{2}) \right]$$

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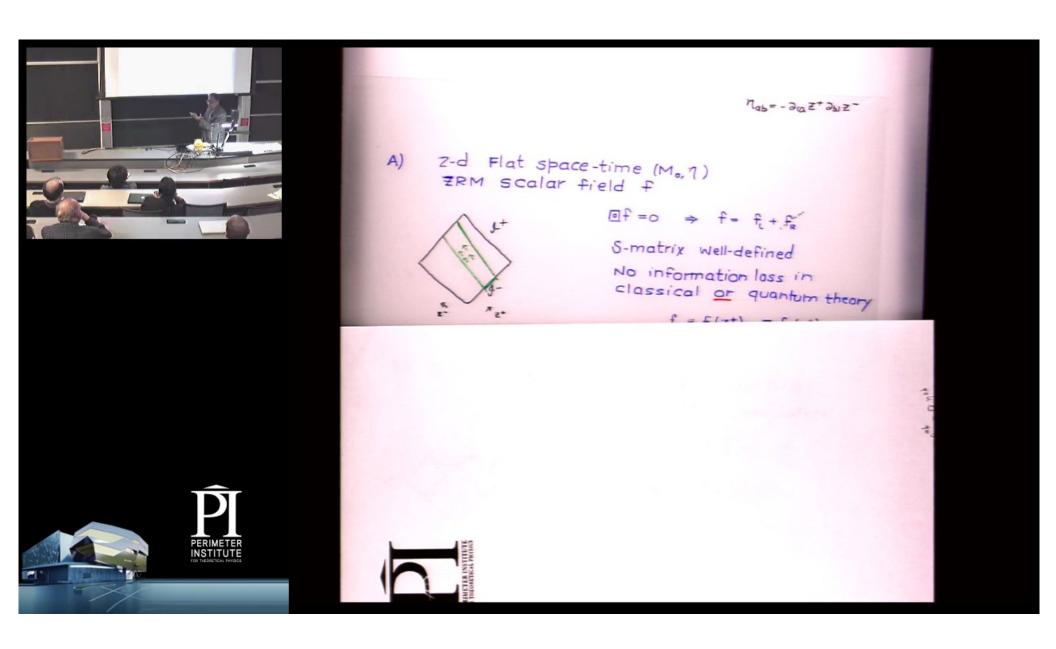
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$$determined in$$

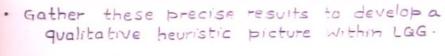
a closed form by (f, nab).

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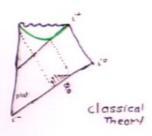


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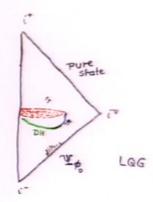




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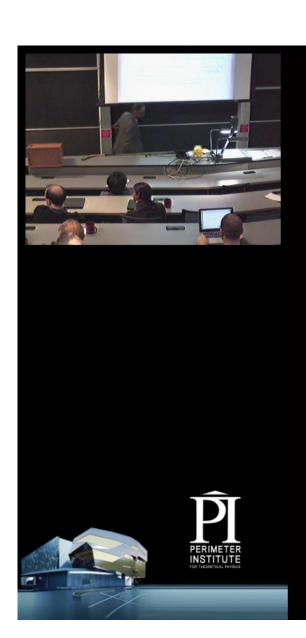


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