

Title: Gravitoelectromagnetism

Date: Nov 21, 2008 11:00 AM

URL: <http://pirsa.org/08110041>

Abstract: Gravitomagnetism is a subtle concept. Adding Lorentz invariance to Newtonian gravity leads to magnetism, but Einsteinian gravitomagnetism differs from Maxwell's electromagnetism. The differences lead to confusion when Lense-Thirring precession is wrongly ascribed to gyroscopes, and when authors disagree about whether lunar laser ranging has measured gravitomagnetism. To clarify these issues, we analyze electric and magnetic effects in local Lorentz frames using the tetrad formalism.

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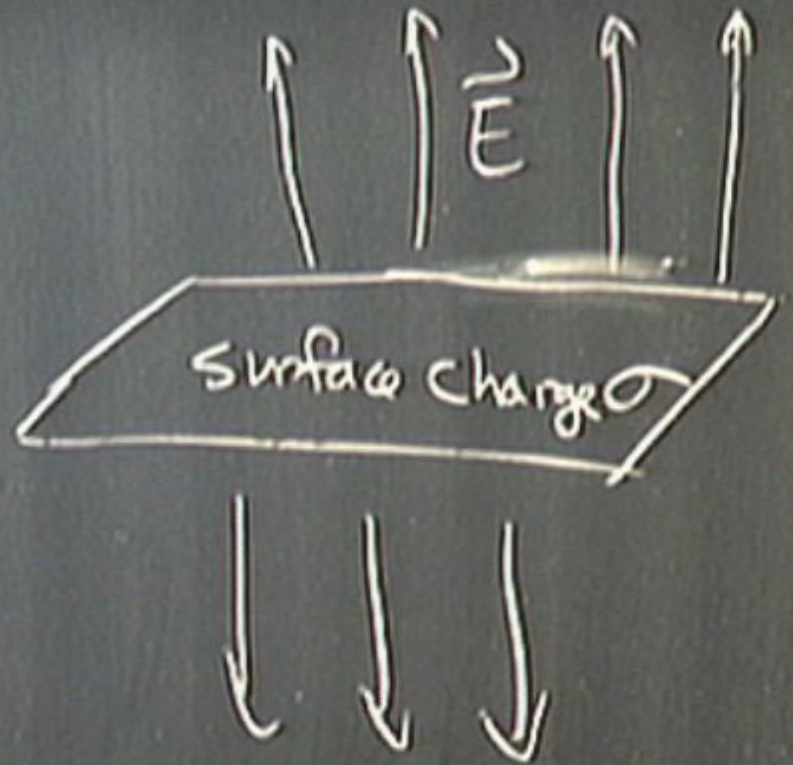
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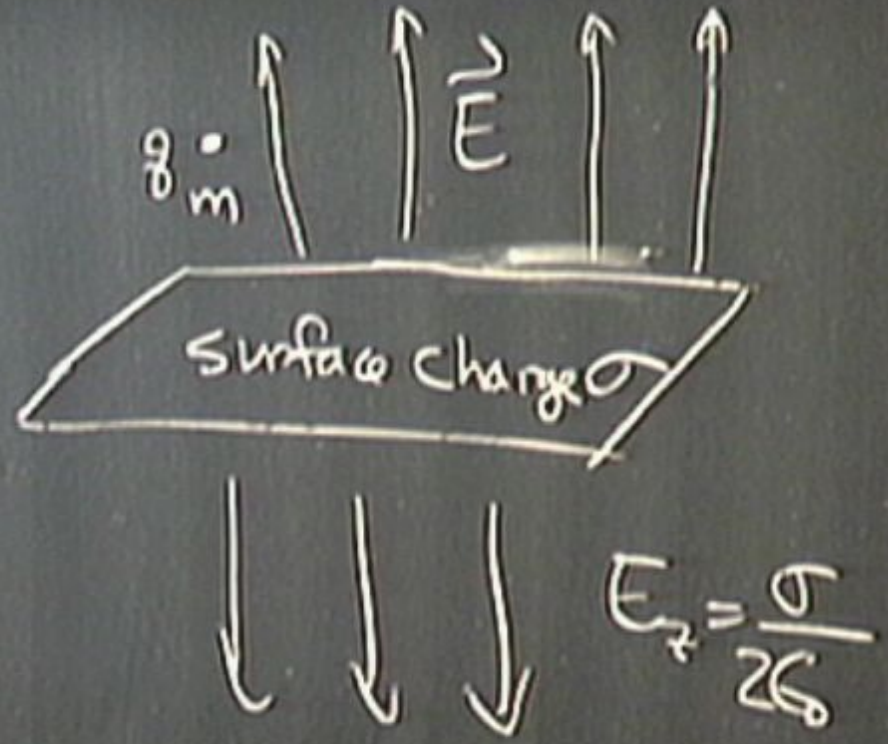
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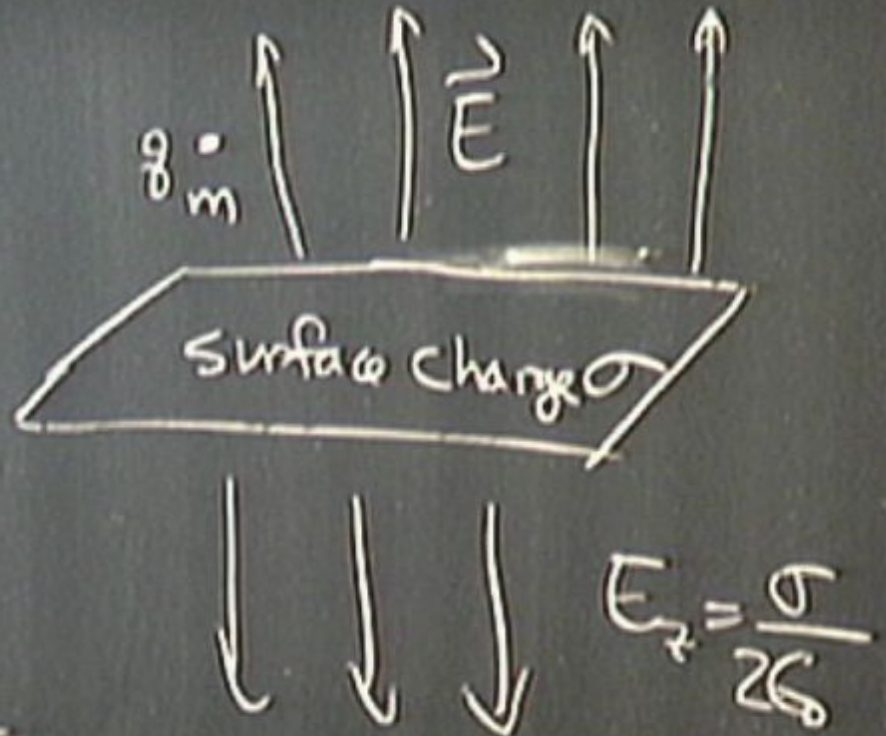


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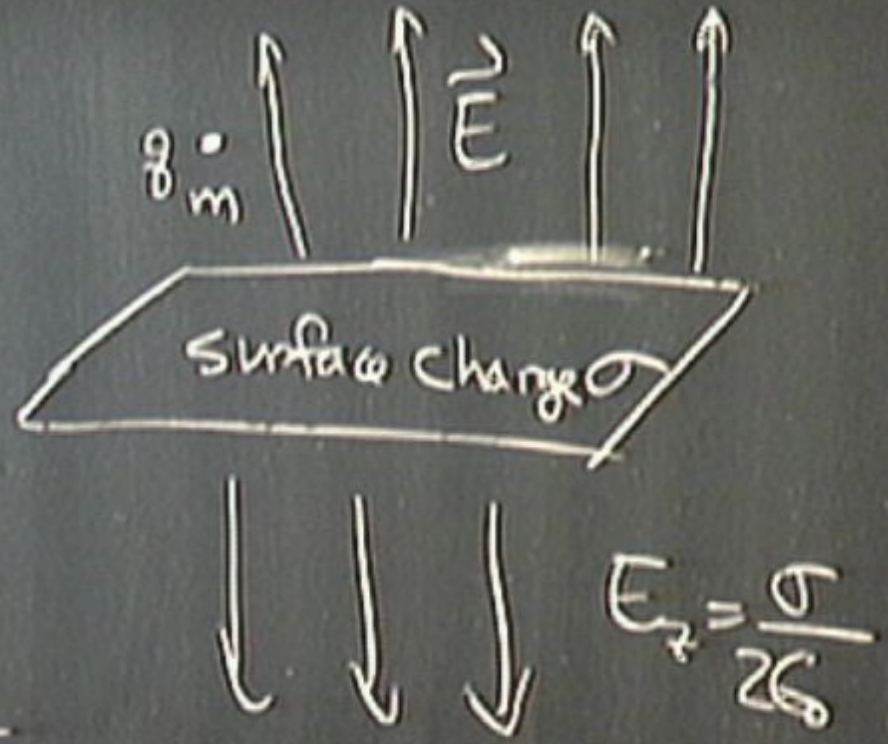
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$$a = \frac{q E_z}{m}$$



Lorentz boost $\vec{u} = u \vec{e}_z$

Along the worldline, $\gamma_u \equiv \frac{1}{\sqrt{1-u^2}}$

$$t' = \gamma_u t$$

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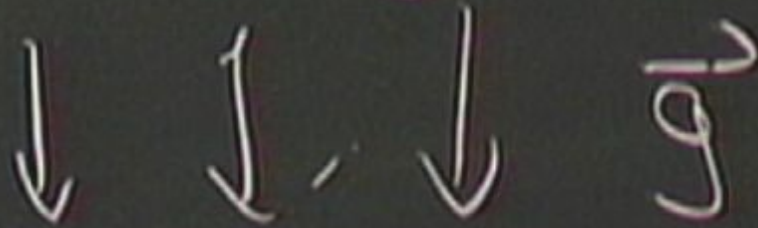
Substituting into $\vec{v}' = \frac{d\vec{x}'}{dt'}$, $\vec{p}' = \frac{m\vec{v}'}{\sqrt{1-v'^2}}$

$$\frac{d\vec{p}'}{dt'} = q(\vec{E}' + \vec{v}' \times \vec{B}')$$

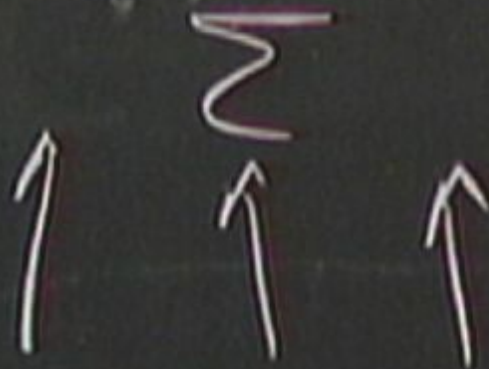
$$\vec{E}' = \gamma_u (\vec{E} + \vec{u} \times \vec{B})$$

$$\vec{B}' = \gamma_u (\vec{B} - \vec{u} \times \vec{E})$$

Now gravity:



Surface mass density



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↓ ↓ ↓ $\vec{g} = -g \hat{z}$

Surface mass density

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surface mass density

follow steps
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$$\frac{d\vec{p}}{dt} = m (\vec{g} + \vec{v} \times \vec{H})$$

gravitomagnetic field

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gravitomagnetic field

$$\vec{g}' = \gamma_u (\vec{g} - \vec{u} \times \vec{H}),$$

$$\vec{H}' = \gamma_u (\vec{H} + \vec{u} \times \vec{g})$$

$$g \Rightarrow \frac{G \Sigma}{2\pi}$$

$$\text{In GR, } \frac{dp^\mu}{d\tau} = -\frac{1}{m} \Gamma_{\nu\lambda}^{\mu} p^\nu p^\lambda \quad | \quad p^\mu = m \frac{dx^\mu}{d\tau}$$

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A, B: Lorentz indices

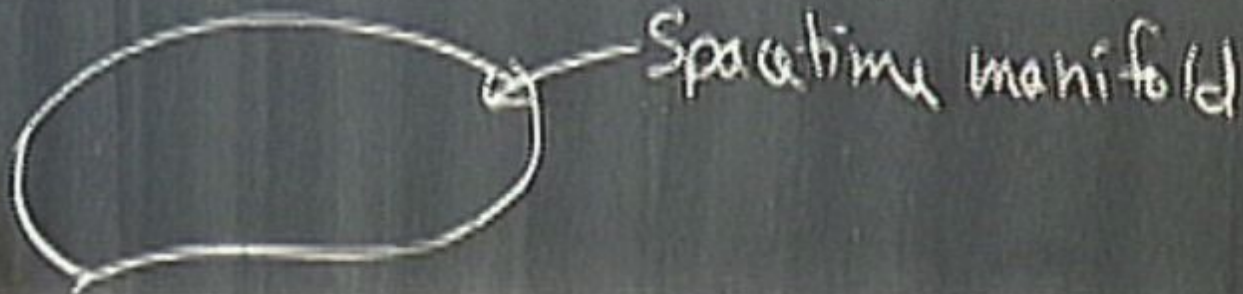
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Combine Lorentz frames + curved spacetime

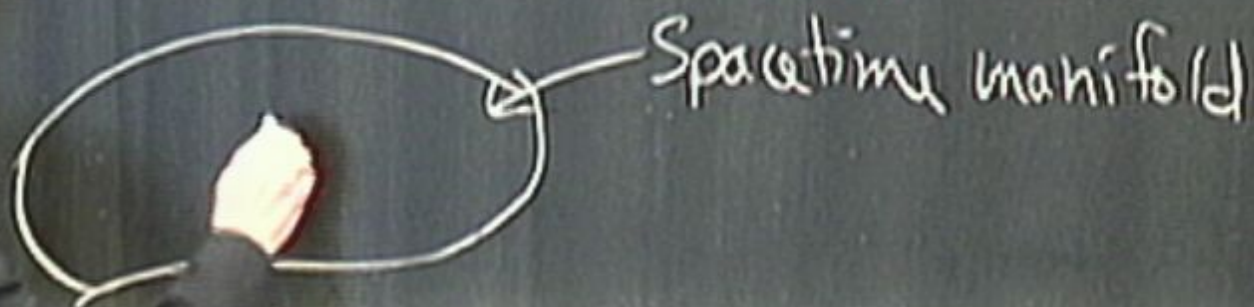


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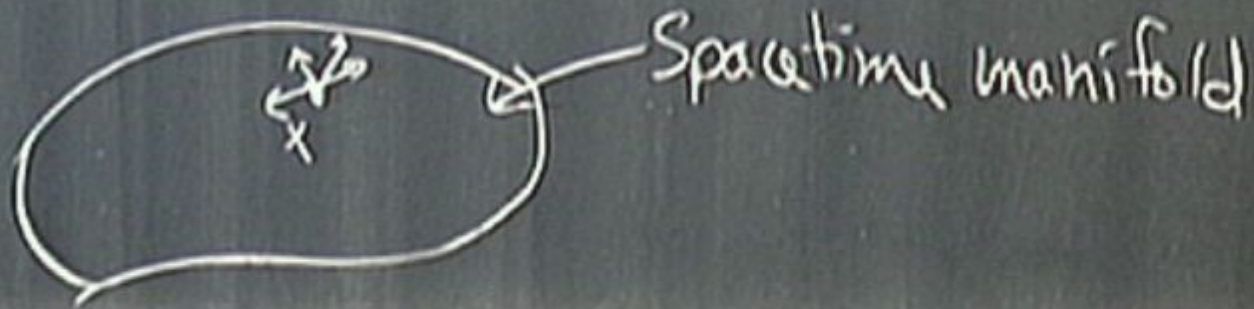
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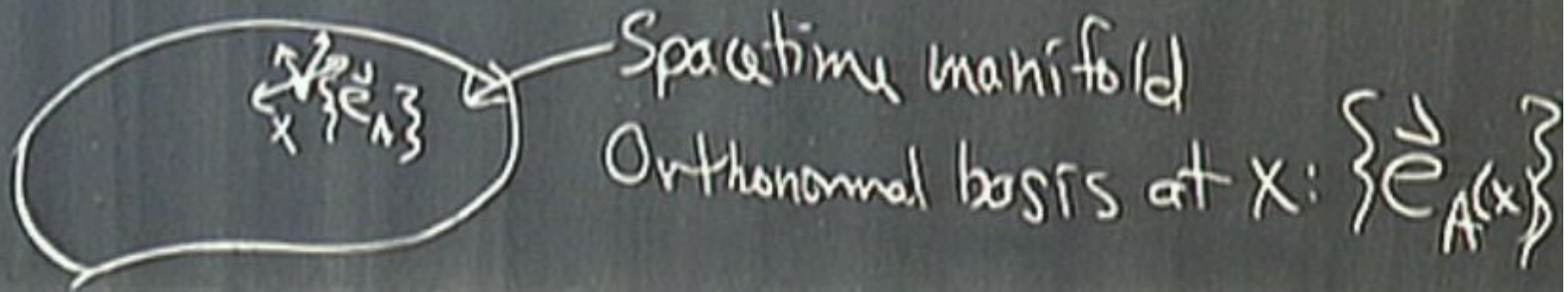
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Relating coordinate

Relating coordinate and orthonormal bases:

$$\vec{e}_M(x) = \underbrace{e^A_M(x)}_{\text{tetrad or vierbein}} \vec{e}_A$$

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Spin connection
Ricci Rotation

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2 kinds of transformations:

Coordinate $x^\mu \rightarrow x'^M(x)$

Lorentz: $P^A \rightarrow \Lambda^A_{BP} P^B$

Given $V^A(x)$

$$D_m V^A \equiv \frac{\partial V^A}{\partial x^m} + \omega_{mB}^A V^B$$

Under L.T., $D_m V^A \rightarrow \Lambda^A_B (D_m V^B)$

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Given $V^A(x)$

ω_{mB}^A is a pseudo tensor

$$\mathbb{D}_m V^A \equiv \frac{\partial V^A}{\partial x^m} + \omega_{mB}^A V^B \quad (\text{cf. } \Lambda_m \rightarrow A_m - \partial_m \Phi) \quad (\text{eig } \Phi)$$

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$$P^A = E(1, v^I)$$

$$P^M = P^L(1, v^i)$$

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 P^A &= E(1, v^I) \\
 P^M &= P^L(1, v^i)
 \end{aligned}
 \left. \vphantom{\begin{aligned} P^A \\ P^M \end{aligned}} \right\}$$

$$\frac{dP^I}{dt} = E \left[(g^I + M^I_j W^j) + \epsilon^I_{BL} (\Omega^R + N^R_j v^j) v^L \right]$$

g^I

$$\left. \begin{aligned} P^A &= E(I, v^I) \\ P^M &= P^L(I, v^I) \end{aligned} \right\} \frac{dP^I}{dt} = E \left[(g^I + M^I_{jI}) + \epsilon^I_{KL} (\Omega^R + N^R_{jI}) v^L \right]$$

$$\omega^I_{t0} = \omega^0_{tI} = -g^I, \quad \omega^I_{tJ} = \epsilon^I_{JE} \Omega^R$$

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$$\omega_{t0}^I = \omega^0, \quad \omega_{tI}^I = -g^I, \quad \omega_{tJ}^I = \epsilon_{JK}^I \Omega^R$$

$$\omega_{j0}^I = \omega_{jI}^0 = -M_{jI}^I, \quad \omega_{jR}^I = \epsilon_{KL}^I N_{jI}^L$$

Weak-field limit:

[The rest of the page is heavily obscured by dark, overlapping brushstrokes, making the text illegible.]

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Choose Lorentz frames so $\vec{e}_{A=0} = \vec{v}$ far-velocity of a coordinate-stationary observer

3-vector $\vec{g} = -\nabla\Phi - \partial_t \underline{w}$, $\underline{\Omega} = -\frac{1}{2} \nabla \times \underline{w}$

$$M^I_j = -\frac{1}{2} \partial_t h^I_j - \epsilon^I_{jk} (\nabla \times \underline{w})^k, N_j^I = -\frac{1}{2} \epsilon^{Ikl} \partial_k h_{lj}$$

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3-vector

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 \omega^I_{t0} &= \omega^I_{tI} = -g^I, & \omega^I_{tJ} &= \epsilon^I_{JK} \Omega^R \\
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Spin Precession

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↑
(Coriolis!)

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(Coriolis!) \rightarrow Frame Dragging!

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2 Equivalence Principles
Acceleration, Rotation