

Title: The Clock Ambiguity and its Implications

Date: Sep 04, 2008 10:30 AM

URL: <http://pirsa.org/08090045>

Abstract:

The clock ambiguity and its implications

- 1) The clock ambiguity
- 2) Physics in context of C.A.
- 3) Field theory (without gravity)
- 4) Specific heat ("with gravity")
- 5) Conclusions + next steps

AA:gr-qc 9408023

AA+Iglesias Aug '07 (PRD '08)

AA+Iglesias May '08 (NYAS preprint)

1)

1) The clock ambiguity

$$H|\psi\rangle_s = 0$$

1) The clock ambiguity

$$H|\psi\rangle_s = 0$$

$$S = C \otimes R$$

clock $\{ |t_i\rangle_c \}$

rest

1) The clock ambiguity

$$H|\psi\rangle_s = 0$$

$$S = C \otimes R$$

clock $\uparrow$ rest $\uparrow$

$$\{|t_i\rangle_c\} \quad \{|j\rangle_R\}$$

1) The clock ambiguity

$$H|\psi\rangle_s = 0$$

$$S = \underset{\substack{\uparrow \\ \text{clock} \\ \{ |t_i\rangle_c \}}}{C} \otimes \underset{\substack{\uparrow \\ \text{rest} \\ \{ |j\rangle_R \}}}{R}$$

1) The clock ambiguity

$$H|\psi\rangle_s = 0$$

$$S = \underbrace{\mathbb{C}}_{\text{clock}} \otimes \underbrace{\mathbb{R}}_{\text{rest}}$$

$\{ |t_i\rangle \}$ $\{ |j\rangle_R \}$

$$|\psi\rangle_s = \sum_{ij} \alpha_{ij} |t_i\rangle |j\rangle_R$$

$$|\phi(t)\rangle_R = \sum_j \alpha_{ij} |j\rangle_R$$

Prob of $|x\rangle_R$ at time t_3

$$\frac{|\langle x | \phi(t_3) \rangle_R|^2}{|\langle \phi(t_3) | \phi(t_3) \rangle|}$$

Prob of $|x\rangle_R$ at time t_3

$$\frac{|\langle x | \phi(t_3) \rangle_R|^2}{|\langle \phi(t_3) | \phi(t_3) \rangle_R|}$$

Consider new state in $\mathbb{R}$ with new time
evolution: Q'_{ij}

Consider new state in $\mathcal{R}$ with new time evolution: α'_{ij}

$$|\psi'\rangle_S = \sum_{ij} \alpha'_{ij} |t_i\rangle_C |j\rangle_R$$

$$M|\psi\rangle_S = |\psi'\rangle_S$$

Consider new state in R with new time evolution: α'_{ij}

$$|\psi'\rangle_S = \sum_{ij} \alpha'_{ij} |t_i\rangle_C |j\rangle_R$$

$$M|\psi\rangle_S = |\psi'\rangle_S$$

define

$$|k\rangle_S = |t_i\rangle_C |j\rangle_R \\ k(i,j)$$

$$|\psi'\rangle = \sum_k \alpha'_k |k\rangle_s$$

$$M'|\psi'\rangle = \sum_k \alpha'_k m'_k |k\rangle_s$$

$$|\psi'\rangle_S = \sum_k \alpha'_k |k\rangle_S$$

$$M'|\psi'\rangle_S = \sum_k \alpha'_k \underbrace{m'_k |k\rangle_S}_{|k'\rangle_S}$$

$$|\psi\rangle_S = \sum_k \alpha'_k |k'\rangle_S$$

$$|\psi'\rangle_S = \sum_k \alpha'_k |k\rangle_S$$

$$M'|\psi'\rangle_S = \sum_k \alpha'_k \underbrace{m'_k |k\rangle_S}_{|k'\rangle_S}$$

$$|\psi\rangle_S = \sum_k \alpha'_k |k'\rangle_S$$

$$= \sum_k \alpha'_{ij} |e_i\rangle_{C'} |j\rangle_{K'}$$

2) Physics w/ C.A.

Quasiseparability

2) Physics w/ C.A.

Quasiseparability $\Rightarrow$ locality

$$H = \int d^3x \mathcal{H}(x)$$

3 Field Theory (without gravity)

3 Field Theory (without gravity)

→ Random H

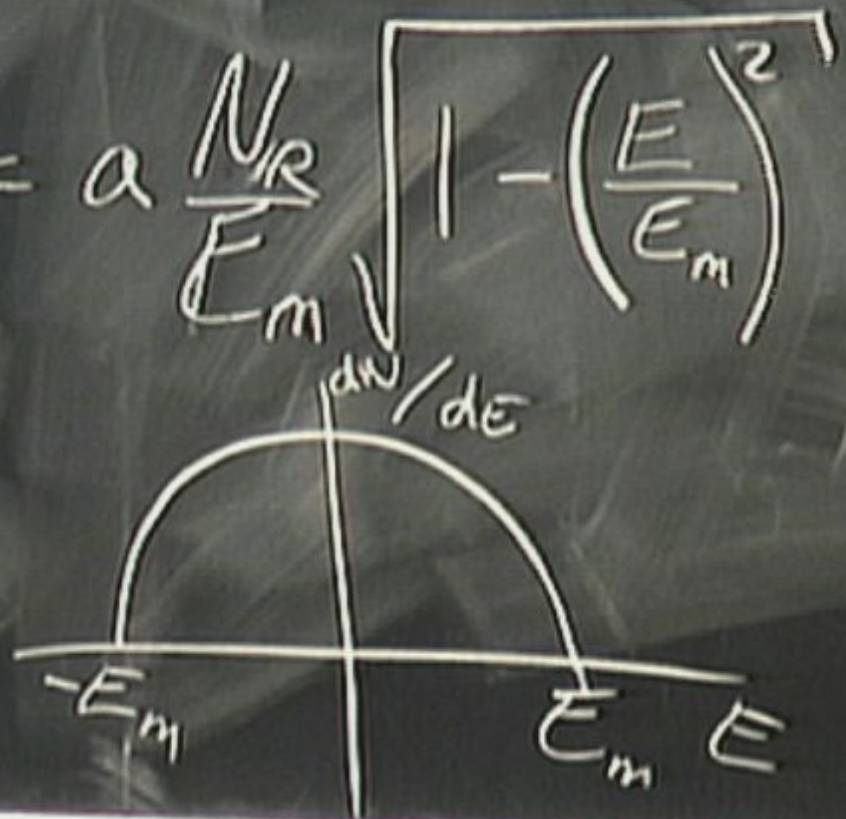
$$\frac{dN}{dE} = a \frac{N_R}{E_m} \sqrt{1 - \left(\frac{E}{E_m}\right)^2}$$

3 Field Theory (without gravity)

→ Random H

$$\frac{dN}{dE} = a \frac{N_R}{E_m} \sqrt{1 - \left(\frac{E}{E_m}\right)^2}$$

Wigner semicircle

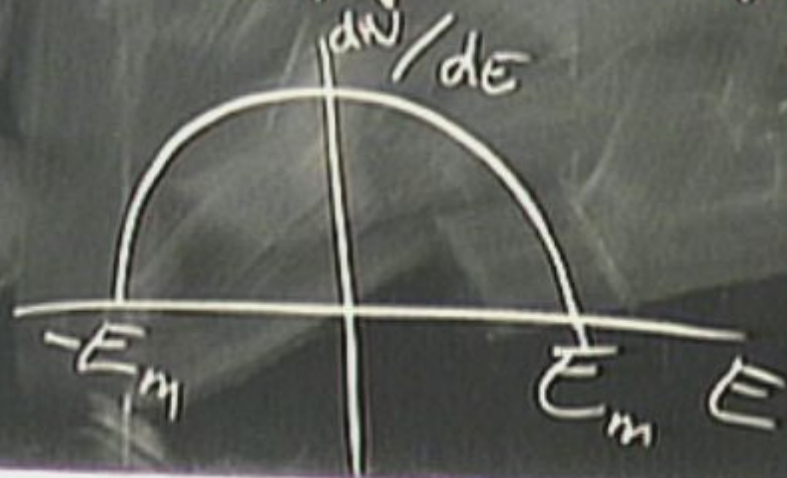


3 Field Theory (without gravity)

→ Random H

$$\frac{dN}{dE} = a \frac{N_R}{E_m} \sqrt{1 - \left(\frac{E}{E_m}\right)^2}$$

Wigner semicircle



$$\left. \frac{dN}{dE} \right|_{FT} = \frac{\beta}{E} \exp \left\{ b \left(\frac{E}{\Delta R} \right)^2 \right\}$$

$$FT \text{ regulator} = \Delta k = \frac{2\pi}{L}$$

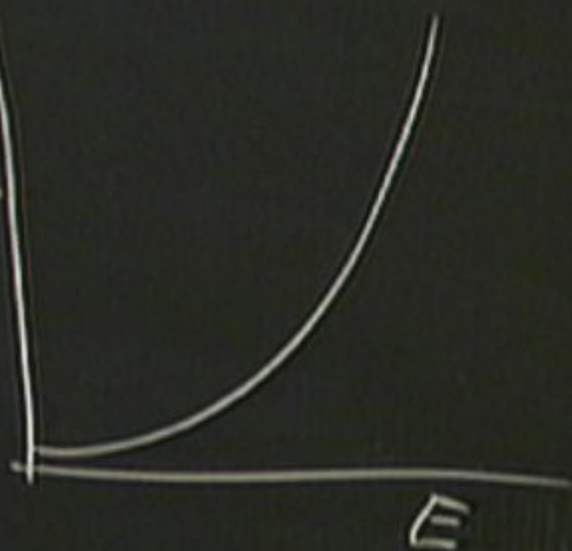
$$\left. \frac{dN}{dE} \right|_{FT} = \frac{B}{E} \exp \left\{ b \left(\frac{E}{\Delta R} \right)^4 \right\}$$

$$FT \text{ regulator} = \Delta k = \frac{2\pi}{L}$$

$$CFT: \alpha = (d-1)/d$$

$$\text{Verlinde: } \alpha(d) = 1/2$$

$$\left. \frac{dN}{dE} \right|_{FT}$$

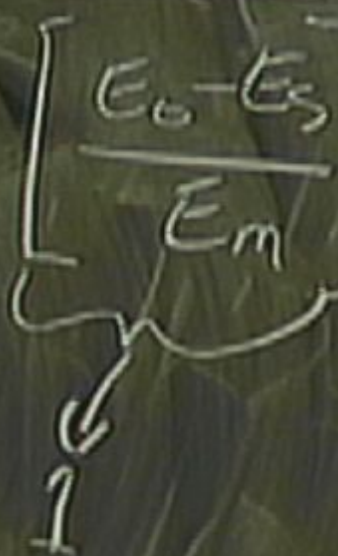


0th order equality:

$$N_R = \left[\frac{B}{\frac{2}{\pi} \left[1 - \left(\frac{E_0 - E_s}{E_m} \right)^{2/2} \right]^{1/2} \frac{E_m}{E_0}} \right] \exp \left\{ b \left(\frac{c E_0}{\Delta^*} \right)^2 \right\}$$

1st (and O^{-14}) order

$$\frac{E_0}{E_0 - E_S} \left(1 - \left[\frac{E_0 - E_S}{E_m} \right] \right)^2 = \alpha \phi \left(C \frac{E_0}{\Delta E} \right)^q$$



1st (and O^{th}) order

$$\frac{E_0}{E_0 - E_s} \left(\frac{E_0 - E_s}{E_m} \right)^2 = \alpha \left(C \frac{E_0}{\Delta k} \right)^2$$

$\downarrow$

$$E_s = E_0 + E_m(1 - \epsilon)$$

2nd Order

$$\left(\left. \frac{dN}{dE} \right|_{FT} - \left. \frac{dN}{dE} \right|_R \right)^2 \approx \left(\left(\frac{E_0}{\Delta R} \right)^2 \frac{\Delta E}{E_0} \right)^2 \equiv \Delta_e$$

$$\left. \frac{dN}{dE} \right|_{E_0}$$

α	$\Delta K(\text{GeV})$	$\Delta E(\text{GeV})$	Δ_2
$1/2$	10^{-25}	10^3	$10^{-24.5}$
$1/2$	10^{-25}	10^{11}	$10^{-16.5}$
$1/2$	10^{-42}	10^3	10^{-16}
$1/2$	10^{-42}	10^{11}	10^{-8}
$3/4$	10^{-25}	10^3	$10^{1.8}$
$3/4$	10^{-25}	10^{11}	$10^{9.8}$
$3/4$	10^{-42}	10^3	$10^{14.5}$
$3/4$	10^{-42}	10^{11}	$10^{22.5}$
1	10^{-25}	10^3	$10^{2.8}$
1	10^{-25}	10^{11}	10^{30}
1	10^{-42}	10^3	10^{45}
1	10^{-42}	10^{11}	10^{52}

$$E_0 = \rho_{\text{pl}} R_H^3 = 10^{80} \text{ GeV}$$

$$\Delta K = m_{\text{g}} = 10^{-25} \text{ (69)} = H_0 = 10^{-42} \text{ GeV}$$

$$\Delta E = E_{\text{ca}} = 10^3 \text{ GeV}$$

$$\text{(or)} \\ = E_{\text{CHECK}} = 10^{11} \text{ GeV}$$

α	$\Delta K(\text{GeV})$	$\Delta E(\text{GeV})$	Δ_2
$1/2$	10^{-25}	10^3	$10^{-24.5}$
$1/2$	10^{-25}	10^{11}	$10^{-16.5}$
$1/2$	10^{-42}	10^3	10^{-16}
$1/2$	10^{-42}	10^{11}	10^{-8}
$3/4$	10^{-25}	10^3	$10^{1.8}$
$3/4$	10^{-25}	10^{11}	$10^{9.8}$
$3/4$	10^{-42}	10^3	$10^{14.5}$
$3/4$	10^{-42}	10^{11}	$10^{22.5}$
1	10^{-25}	10^3	10^{28}
1	10^{-25}	10^{11}	10^{36}
1	10^{-42}	10^3	10^{45}
1	10^{-42}	10^{11}	10^{53}

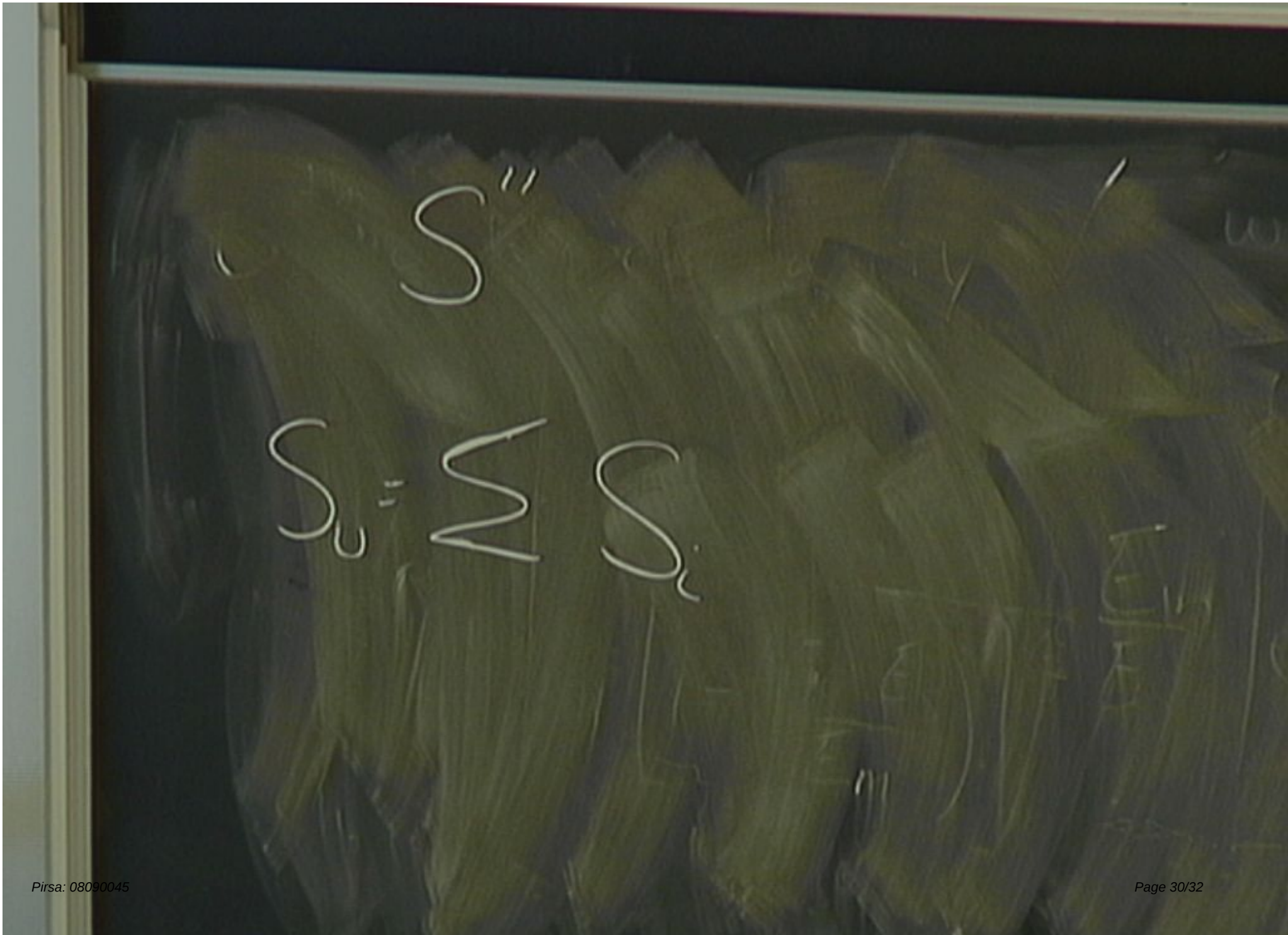
$$E_0 = \rho_u R_H^3 = 10^{80} \text{ GeV}$$

$$\Delta K = m_\gamma = 10^{-25} \text{ GeV} = H_0 = 10^{-42} \text{ GeV}$$

$$\Delta E = E_{\text{acc}} = 10^3 \text{ GeV}$$

$$\text{or } = E_{\text{CHECK}} = 10^{11} \text{ GeV}$$

$$\Delta E \sim \frac{1}{\delta t}$$



$$\frac{1}{2} \sqrt{\frac{1}{2}}$$

$$S''$$

$$S_c'' \sim S_i$$

$$S_{wig}'' = -\frac{1}{8}$$

S''

$$S_0 = \sum S_i$$

$$S''_{wig} = -\frac{1}{8}$$

Dominated by

$$S''_{0m} = \pm 10^{-2} \text{ GeV}$$