

Title: On the analogy between thermodynamics and quantum entanglement

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Abstract: According to the second law of thermodynamics the entropy of a system cannot decrease by adiabatic state transformations. In quantum mechanics, the '\degree of entanglement\' of a state cannot increase under state transformations of a certain kind (local operations assisted by classical communication) In this talk I will explore the significance of the analogy between these two statements.

# On the analogy between thermodynamics and quantum entanglement

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# Classical Thermodynamics and its Second Law

Clausius (1850), Kelvin (1851), Planck (1897), Carathéodory (1909):

- ▶ concerned with macroscopic systems (e.g.: gas or mixture of vapor, liquid and solids) in thermal equilibrium states.
- ▶ Each such system has a state space  $\Gamma$  consisting of all possible equilibrium states  $X$ . (E.g.: the  $(p, V)$  diagram for a fluid.)
- ▶ There exists a real-valued function  $S$  on  $\Gamma$ .  
 $S(X)$  is called the *entropy* of the system in state  $X$ .
- ▶ Second Law of thermodynamics:  
*The entropy of a system can never decrease as a result of an adiabatic process.*
- ▶ Here, an adiabatic process is a state transformation which involves no exchange of heat with the environment.

# Classical Thermodynamics and its Second Law

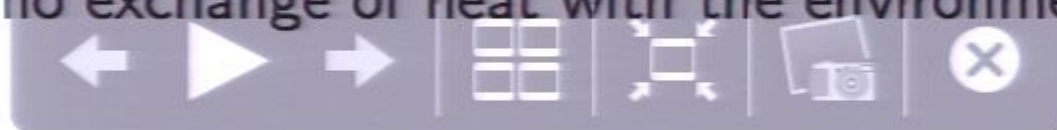
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# Problems with classical formulations of the Second Law

- ▶ These treatments of the Second Law are not mathematically sophisticated. They make abundant use of silent assumptions.
- ▶ Does the Second Law hold for all TD systems? (Or only for "simple" systems?)
- ▶ Is there a unique entropy function?
- ▶ Exactly what state transformations for a given system can be realized by adiabatic processes?



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# Axiomatic approach of Lieb and Yngvason (1999)

## Notation

$X \preceq Y$  :  $\exists$  an adiabatic process that transforms state  $X$  into  $Y$

$X + Y$  : the composition of two systems in states  $X$  and  $Y$ .

$tX$  : a "scaled version" of state  $X$

## Axioms:

1.  $X \preceq X$
2. If  $X \preceq Y$  and  $Y \preceq Z$  then  $X \preceq Z$ .
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For all  $X, Y$  in the same state space:  $X \preceq Y$  or  $Y \preceq X$ .

## Second Law à la L&Y:

### Theorem

*If Axioms 1–6 and the CH hold, then there is a function  $S$  such that:*

$$X \preceq Y \text{ in the same state space} \iff S(X) \leq S(Y), \quad (1)$$

*and which is additive and extensive, i.e.:*

$$S(t_1 X_1 + \cdots + t_n X_n) = \sum_{i=1}^n t_i S(X_i) \quad (2)$$

*Moreover this function is **unique** upto affine transformations.*

## Remarks

- ▶ No differentiability of  $S$  or assumptions on the topology of  $\Gamma$  needed. This is **much** more general than classical treatments.
- ▶ LY adopt a liberal reinterpretation of "adiabatic":  
*"State  $Y$  is adiabatically accessible from state  $X$ , in symbols  $X \lesssim Y$ , if it is possible to change the state  $X$  to  $Y$  by means of an interaction with some auxiliary device ... and a weight, in such a way that the device returns to its initial state at the end of the process whereas the weight may have changed its position in a gravitational field."*
- ▶ LY provide an appealing fresh starting point for attempts to reduce the Second Law to statistical mechanics. (Since there is a concrete realisation of these axioms in terms of probability measures on another phase space)
- ▶ Main weakness is the Comparability Hypothesis. This is only claimed to hold for "simple" systems.
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Consider quantum system composed of two subsystems (far from each other) in a pure state

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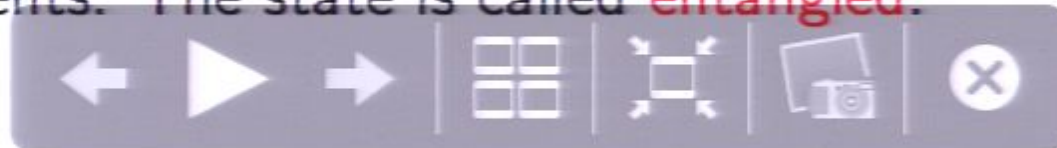
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one can also attribute a pure state to its components. The state is called **separable**.

- ▶ But if it takes the form

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Theorem ( Schmidt's biorthogonal decomposition theorem:)

For every state  $|\Psi\rangle$  there are orthonormal bases on  $\mathcal{H}_1$  and  $\mathcal{H}_2$  such that

$$|\Psi\rangle = \sqrt{p_1}|\alpha_1\rangle|\beta_1\rangle + \cdots + \sqrt{p_n}|\alpha_n\rangle|\beta_n\rangle \quad (5)$$

with  $p_i \geq 0$  and  $\sum p_i = 1$ . The values  $p_\Psi = (p_1, \dots, p_n)$  are called *Schmidt coefficients* of state  $|\Psi\rangle$ .

For a separable (i.e. not-entangled) state,  $p_\Psi = (1, 0, \dots, 0)$

For a 'maximally' entangled state  $p_\Psi = (1/n, \dots, 1/n)$ .

## Measures of entanglement and LOCC operations

- ▶ There have been many proposals of entropy-like measures of entanglement, or to prove the uniqueness of such a measure. (e.g.: Vedral & Plenio (1997), Rudolph (2001), Donald e.a. (2002), Vedral & Kashefi (2002))
- ▶ There is no consensus on this issue; in particular for mixed states.
- ▶ However, consensus exists on the following claim: Whatever one means by the "amount of entanglement" of a state  $|\Psi\rangle$ , it is not possible to increase the amount of entanglement by LOCC operations.

### Definition

LOCC operations: all state transformations that can be achieved by local operations (such as local unitary evolutions, local interactions with third systems, or local measurements), assisted by classical communication.



## Theorem (Nielsen's Theorem)

State  $|\Psi\rangle$  can be transformed (with certainty) into  $|\Psi'\rangle$  by LOCC operations if and only if  $p_{\Psi'}$  is **majorized** by  $p_{\Psi}$

Where

### Definition (Majorization)

probability distribution  $p = (p_1, \dots, p_n)$  is majorized by  $q = (q_1, \dots, q_n)$  (Notation:  $p \preceq q$ ) if  $p$  can be written as a convex mixture of  $q$  and permutations of  $q$ :

$$p_i = \sum_j \alpha_j (\Pi_j q)_i \quad \text{for permutation matrices } \Pi_j \text{ and } \alpha_j \geq 0, \sum \alpha_j = 1$$

- ▶ In words:  $p$  is more “uniform”, or more “disordered” than  $q$ .
- ▶ A minimal requirement on a putative measure of entanglement, say  $E(|\Psi\rangle)$  is thus that it should respect majorization in the sense that

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## A partial analogy between LOCC and adiabatic operations

- ▶ Majorization is a preorder (i.e. Axioms 1 and 2 above hold)
- ▶ Interpret combinations as products:

$$X + Y \iff p \otimes q = (p_1 q_1, p_1 q_2, \dots, p_n q_n)$$

Then Axiom 3 holds.

- ▶ Interpret scaling as taking multiple copies:

$$tX \iff p^{\otimes t} = p \otimes \dots \otimes p, \quad t \in \mathbb{N}$$

Then Axiom 5 holds in the integer form  $p^{\otimes(n+k)} \sim p^{\otimes n} \otimes p^{\otimes k}$ .

- ▶ Axiom 6 holds in the integer form

$$\text{If } p^{\otimes k} \otimes z \preceq_T q^{\otimes k} \otimes z' \text{ for } k \longrightarrow \infty \text{ then } p \preceq_T q.$$

- ▶ But Axiom 4 **fails** in this case. One may have

$$p \not\preceq q \text{ but } p \otimes p \preceq q \otimes q$$

- ▶ Also, a lemma following from Axiom 1–6 is **Cancellation**

$$\text{If } X + Z \preceq Y + Z \text{ then } X \preceq Y$$

The corresponding statement for majorization **fails**.



## Extending the analogy

- ▶ By Nielsen's theorem, this last failure means that there are quantum states  $|\Psi\rangle$  and  $|\Psi'\rangle$  such that it is impossible to transform  $|\Psi\rangle$  into  $|\Psi'\rangle$  by LOCC, but it is nevertheless possible to transform  $|\Psi\rangle|\Omega\rangle$  into  $|\Psi'\rangle|\Omega\rangle$  for a suitable state  $|\Omega\rangle$ ! This is known as **entanglement catalysis** (Jonathan and Plenio, 1999).
- ▶ Entanglement catalysis suggests introducing a preorder that extends majorization, called **trumping** (Nielsen).

### Definition (Trumping)

Distribution  $p$  is trumped by  $q$  ( $p \preceq_T q$ ) iff  $\exists$  a probability distribution  $z = (z_1, \dots, z_m)$  such that  $p \otimes z \preceq q \otimes z$ .

Clearly:  $p \preceq q \implies p \preceq_T q$ .

- ▶ Idea: Replace majorization by trumping. Then all the Axioms 1–6 hold!

## Where the analogy breaks down

- ▶ But there is **no way** the Comparability Hypothesis can hold in this case (except for qubits); if  $n > 2$  probability distributions are simply not completely ordered by trumping.
- ▶ However, we would like to get rid of this hypothesis anyway.

Recently, Turgut (2007) and Klimesh (2007) obtained the following:

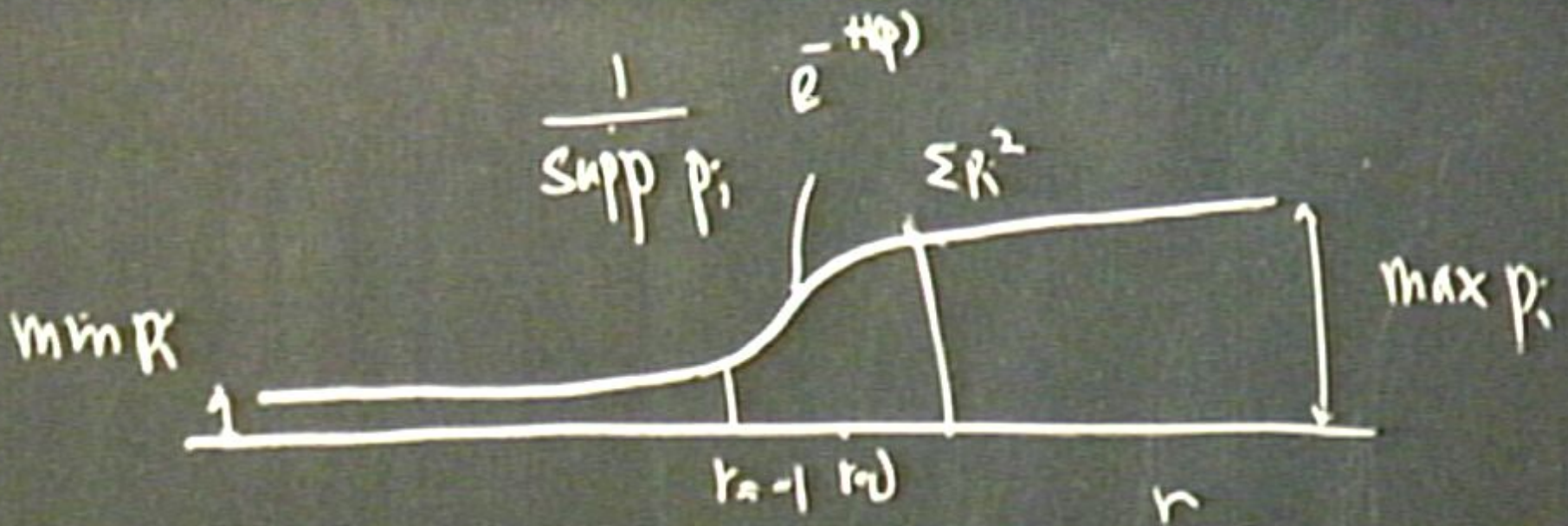
### Theorem (Turgut-Klimesh)

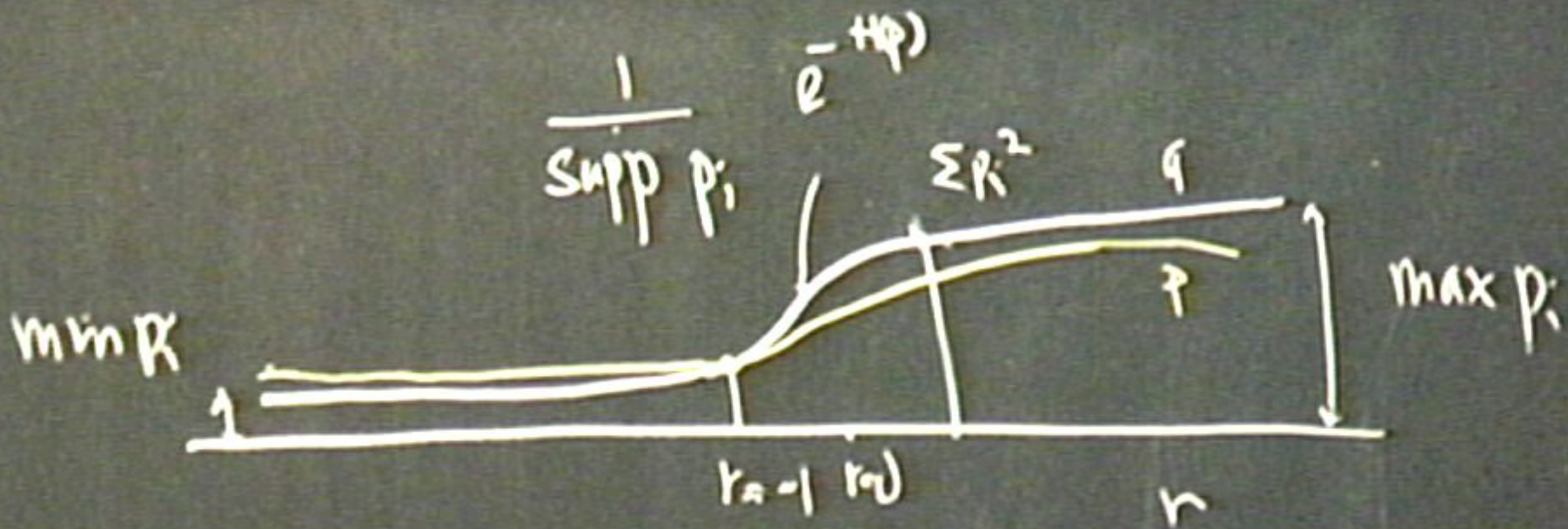
*Suppose  $p$  and  $q$  only have non-zero elements. Then  $p \prec_T q$  holds if and only if for all  $r \in \mathbb{R}$*

$$\begin{aligned} M_r(p) &< M_r(q) && \text{if } r > -1 \text{ and} \\ M_r(p) &> M_r(q) && \text{if } r < -1 \end{aligned}$$

Here:

$$M_r(p) = \left( \sum_{i=1}^n p_i^{(r+1)} \right)^{1/r} \quad r \in \mathbb{R} \tag{6}$$





# Remarks

- ▶ The expressions  $M_r(p)$  for  $r \geq -1$  provide well-known generalizations of entropy;  $-\log M_{r-1}(p)$  is the so-called Rényi entropy of  $p$ .
- ▶ remarkably, the Turgut-Klimesh theorem also requires inequality for these expressions with  $r < -1$ .
- ▶ Main point: trumping holds **iff** there is inequality in a family of entropy-like expressions. Moreover, one has additivity and extensivity since

$$\log M_r(p_1^{\otimes t_1} \otimes \cdots \otimes p_n^{\otimes t_n}) = \sum_i t_i \log M_r(p_i) \quad \forall r \in \mathbb{R}$$

## Exploiting the analogy: I. From QM to TD

The Turgut-Klimesh theorem suggests the problem: if we drop the Comparability Hypothesis, can one still derive the Lieb-Yngvason second law for a **family** of entropies? Yes!

Theorem ( Dubra, Maccheroni & Ok (2004))

Let  $\Gamma$  be a (compact) mixture space, i.e.

if  $X, Y \in \Gamma$  then also  $tX + (1 - t)Y \in \Gamma$  for  $0 < t < 1$ .

► let  $\preceq$  be a preorder on  $\Gamma$  that satisfies

$$X \preceq Y \iff tX + (1 - t)Z \preceq tY + (1 - t)Z \quad (\forall t \in [0, 1])$$

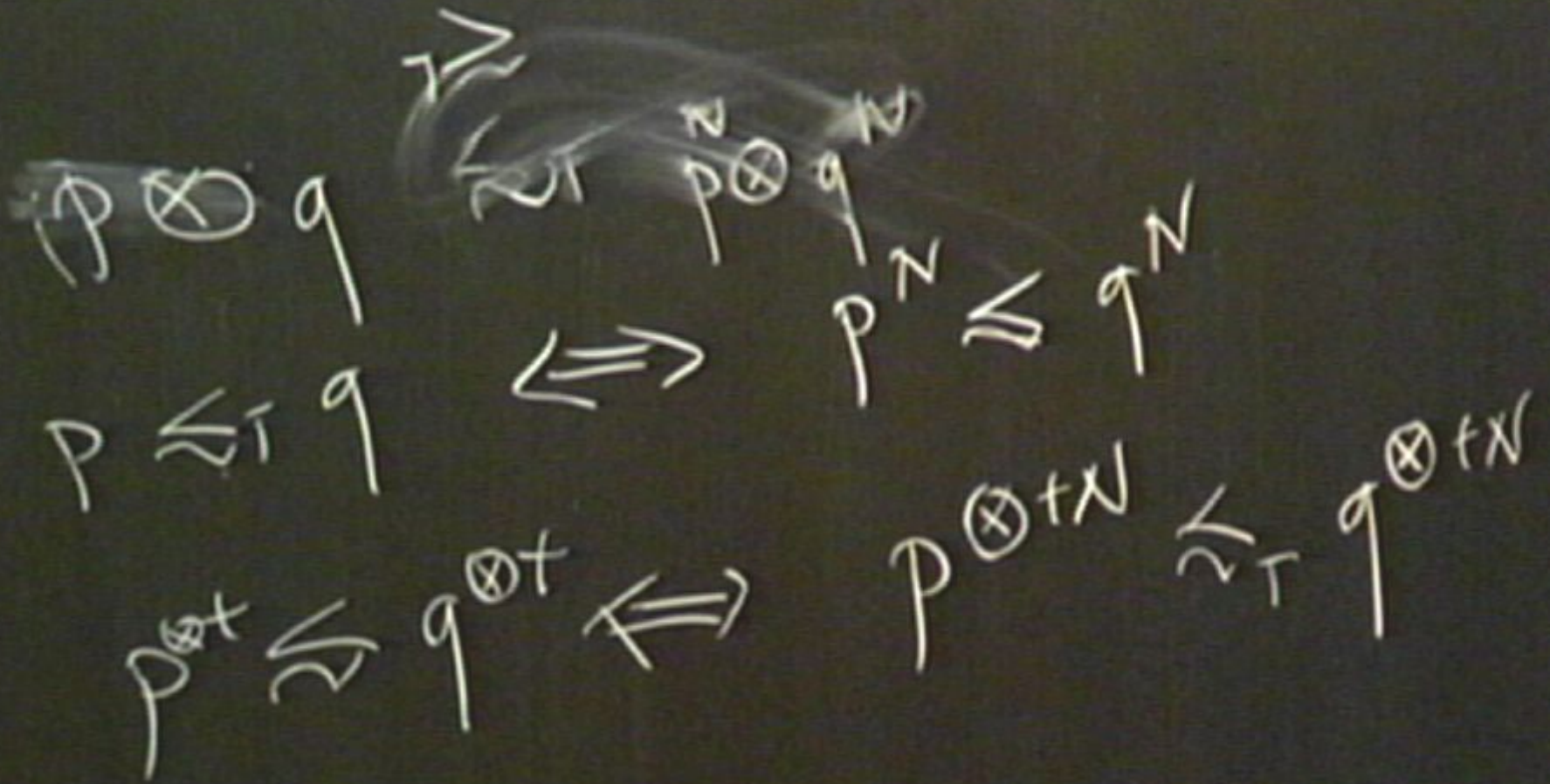
$\{t \in [0, 1] : tX + (1 - t)Z \preceq tY + (1 - t)Z'\}$  is closed

then  $\exists$  a convex closed family  $\mathcal{S}$  of additive and extensive real-valued functions  $S$  on  $\Gamma$  such that:

$$X \preceq Y \iff S(X) \leq S(Y) \quad \forall S \in \mathcal{S}$$

## Exploiting the analogy: II. From TD to QM

- ▶ Nielsen's Theorem applies only to pure quantum states.
- ▶ For mixed states  $\rho$ , we may still use the axiomatic approach and interpret  $\rho \preceq \sigma$  as 'ELOCC-accessibility'. (ELOCC for LOCC assisted by entanglement catalysis).
- ▶ Then the result of Dobra, Maccheroni & Ok will also apply to mixed states.
- ▶ Of course, the question remains to identify this class of functions  $\mathcal{S}$  in this case.





## A partial analogy between LOCC and adiabatic operations

- ▶ Majorization is a preorder (i.e. Axioms 1 and 2 above hold)
- ▶ Interpret combinations as products:

$$X + Y \iff p \otimes q = (p_1 q_1, p_1 q_2, \dots, p_n q_n)$$

Then Axiom 3 holds.

- ▶ Interpret scaling as taking multiple copies:

$$tX \iff p^{\otimes t} = p \otimes \dots \otimes p, \quad t \in \mathbb{N}$$

Then Axiom 5 holds in the integer form  $p^{\otimes(n+k)} \overset{\uparrow}{\sim} p^{\otimes n} \otimes p^{\otimes k}$ .

- ▶ Axiom 6 holds in the integer form

$$\text{If } p^{\otimes k} \otimes z \preceq_T q^{\otimes k} \otimes z' \text{ for } k \longrightarrow \infty \text{ then } p \preceq_T q.$$

- ▶ But Axiom 4 **fails** in this case. One may have

$$p \not\preceq q \text{ but } p \otimes p \preceq q \otimes q$$

- ▶ Also, a lemma following from Axiom 1–6 is **Cancellation**

$$\text{If } X + Z \preceq Y + Z \text{ then } X \preceq Y$$

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