

Title: Polarisation tomography of macro- and mesoscopic quantum states of light

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Abstract: We report an experiment on reconstructing the quantum state of bright (macroscopic) polarization-squeezed light generated in a birefringent (polarization-maintaining) fibre due to the Kerr nonlinearity. The nonlinearity acts on both H and V polarization components, producing quadrature squeezing; by controlling the phase shift between the H and V components one can make the state squeezed in any Stokes observable. The tomography is performed by measuring histograms for a series of Stokes observables, and the resulting histograms (tomograms) are processed in a way similar to the classical 3D Radon transformation. At the output, we obtain the polarization Q-function, which in the case of large photon numbers coincides with the polarization W-function. An interesting extension of the performed experiment will be going down to lower photon numbers (mesoscopic quantum states), and we expect a different behaviour of polarization W and Q functions in this case. An experiment on producing such states is discussed.

POLARIZATION TOMOGRAPHY OF MACRO- AND MESOSCOPIC STATES OF LIGHT

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³ Departamento de Física, Universidad de Guadalajara

⁴ Departamento de Óptica, Facultad de Física, Universidad Complutense Madrid



OUTLINE

Polarization of light: quantum and classical description

Polarization tomography

Experiment made: polarization tomography of Kerr-squeezed light with bright polarized component

Experiment planned: polarization tomography of bright squeezed vacuum

POLARIZATION OF LIGHT

Classical Stokes parameters (1852)

$$S_0 \equiv I_H + I_V, \quad S_1 \equiv I_H - I_V,$$

$$S_2 \equiv I_{45} - I_{-45}, \quad S_3 \equiv I_R - I_L$$



Quantum description: Stokes operators

$$\hat{S}_0 \equiv \mathbf{a}_H^+ \mathbf{a}_H + \mathbf{a}_V^+ \mathbf{a}_V, \quad \hat{S}_1 \equiv \mathbf{a}_H^+ \mathbf{a}_H - \mathbf{a}_V^+ \mathbf{a}_V,$$

$$\hat{S}_2 \equiv \mathbf{a}_{45}^+ \mathbf{a}_{-45} - \mathbf{a}_{-45}^+ \mathbf{a}_{45}, \quad \hat{S}_3 \equiv \mathbf{a}_R^+ \mathbf{a}_R - \mathbf{a}_L^+ \mathbf{a}_L,$$

$$\vec{\hat{S}} = \{\hat{S}_1, \hat{S}_2, \hat{S}_3\}, \quad \vec{n} = \{\cos \vartheta, \sin \vartheta \cos \varphi, \sin \vartheta \sin \varphi\}$$

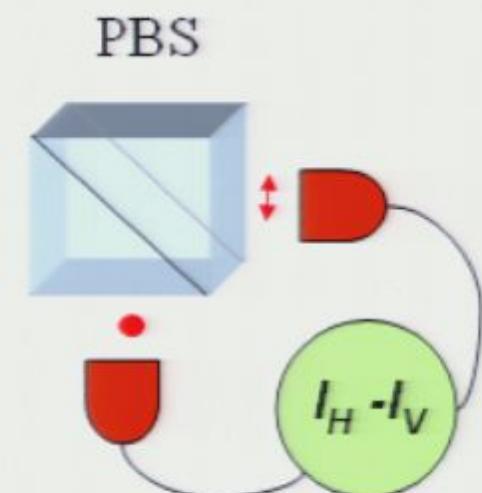
$$\hat{S}(\vartheta, \varphi) = \vec{\hat{S}} \cdot \vec{n}$$

'Stokes space', not a sphere

$$[\hat{S}_i, \hat{S}_j] = 2i\hat{S}_k \Rightarrow \Delta^2 S_i \Delta^2 S_j \geq |\langle S_k \rangle|^2$$

$$i, j, k = 1, 2, 3$$

$\Delta^2 S_i < |\langle S_k \rangle| < \Delta^2 S_j$ polarization squeezing

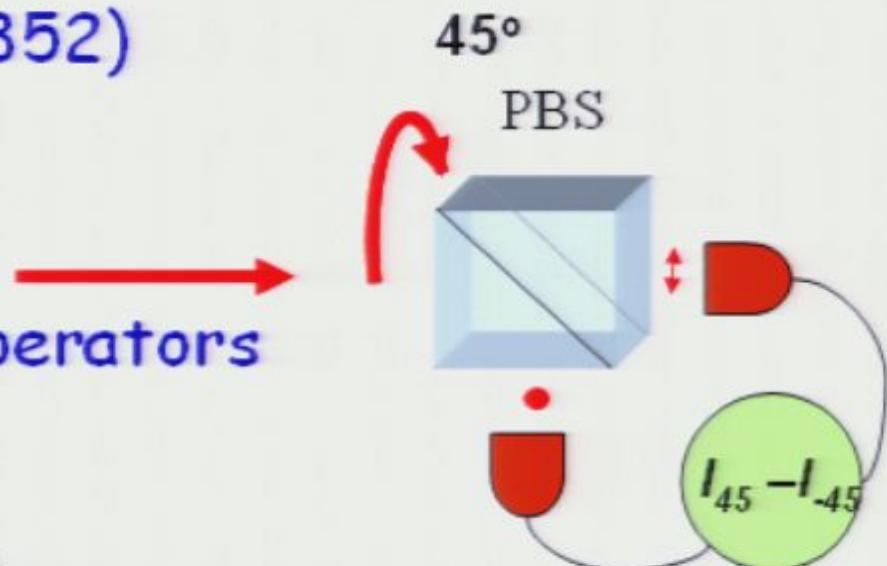


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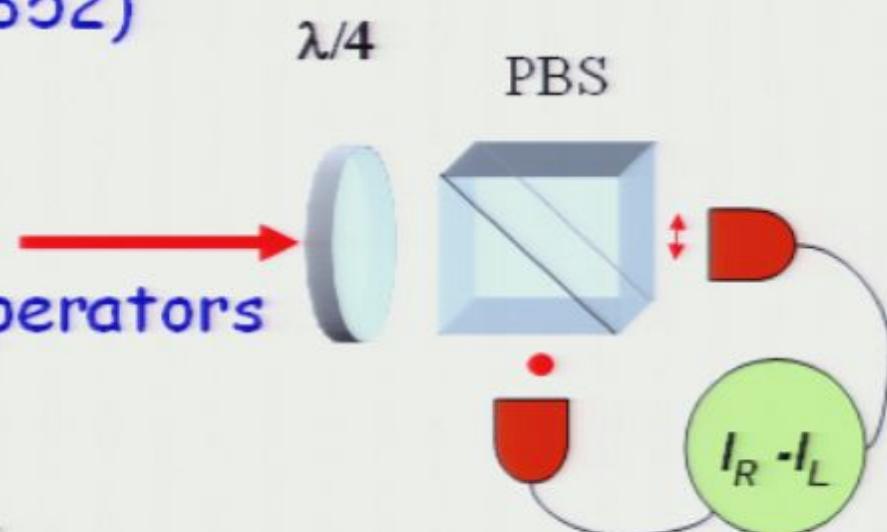
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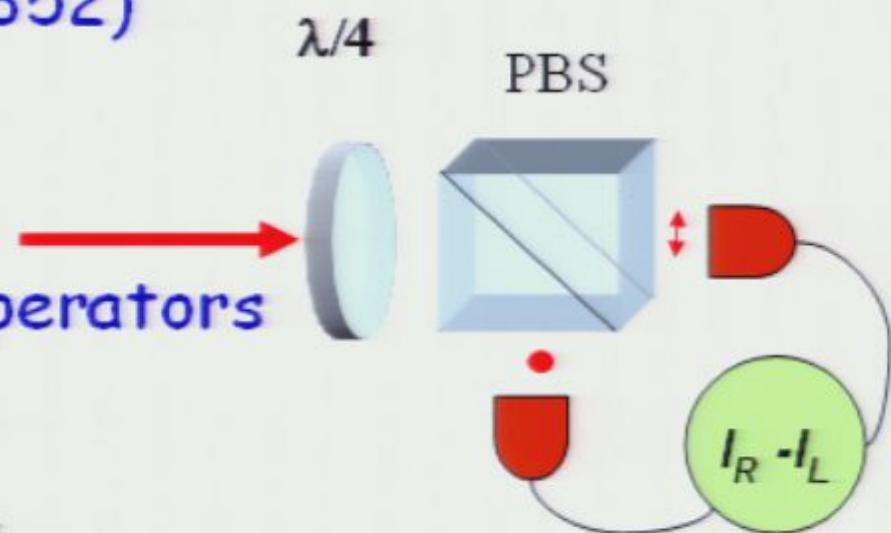
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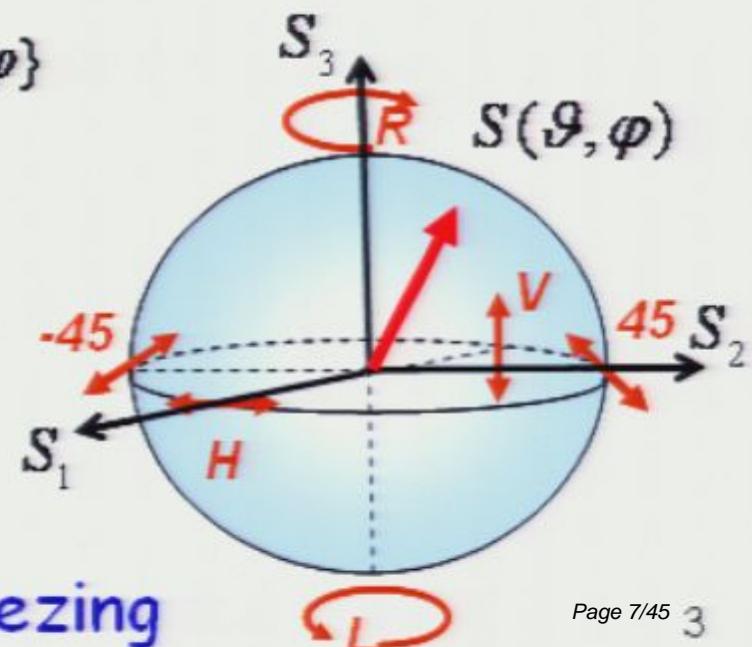
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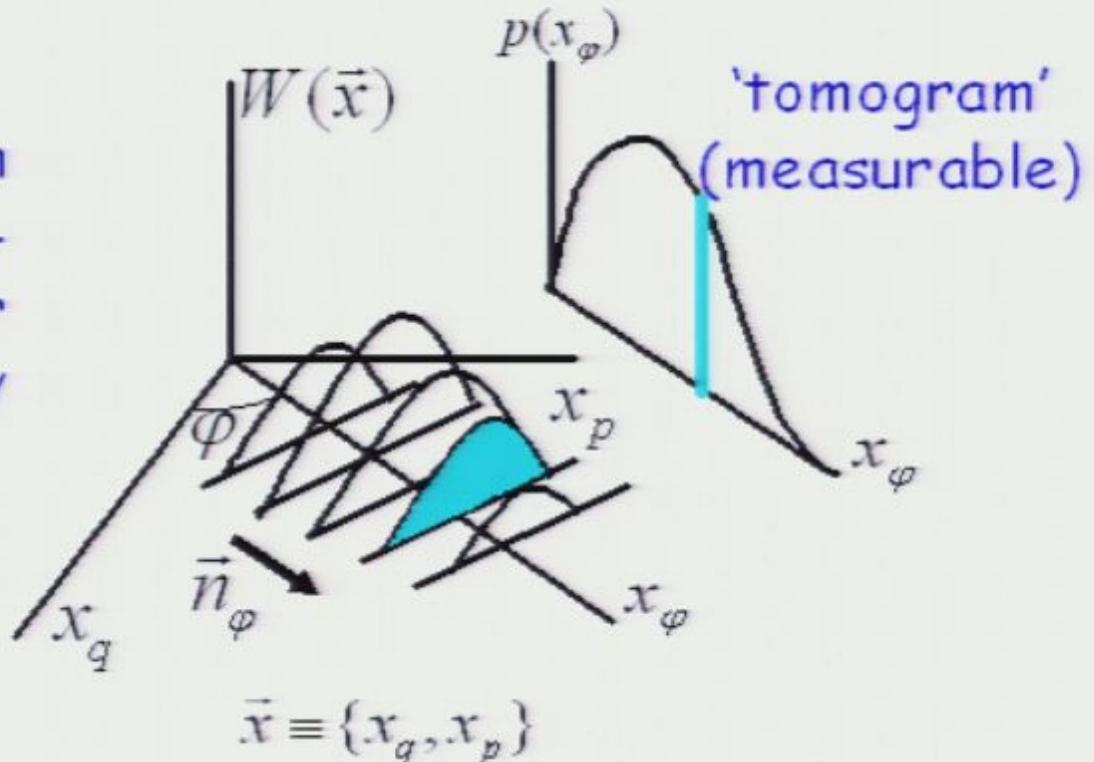


CAN WE MEASURE MORE THAN THE STOKES PARAMETERS AND THEIR VARIANCES?

QUANTUM TOMOGRAPHY

'Traditional' quantum tomography, or Wigner-function tomography, or homodyne tomography

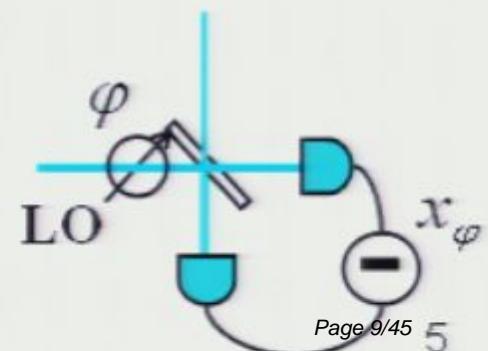
D.T.Smithey, M.Beck,
 M.Raymer, and A.Faridani,
PRL 70, 1244 (1993)



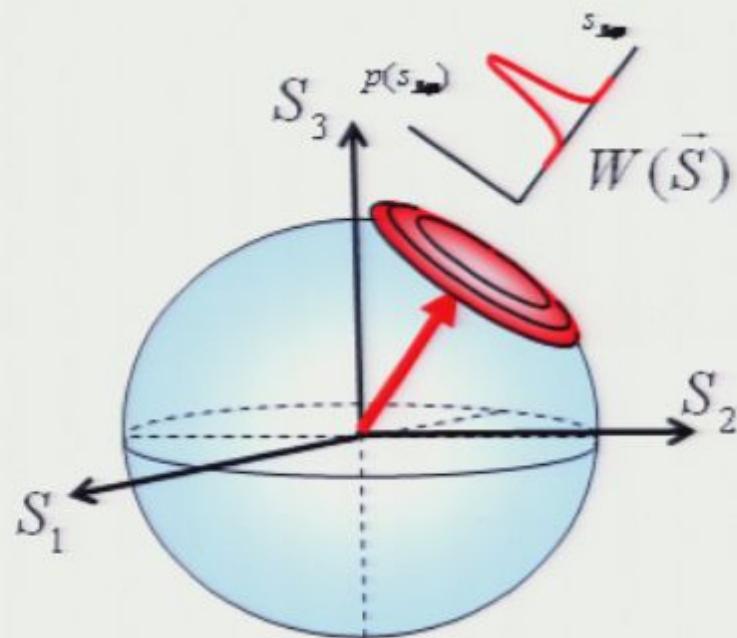
Direct Radon transform: $p(x_\varphi) = \int W(\vec{x}) \delta(\vec{x} \cdot \vec{n}_\varphi - x_\varphi) d^2 \vec{x}$

Inverse Radon transform: $p(x_\varphi) \rightarrow W(\vec{x})$

Can we do this for polarization variables?



POLARIZATION TOMOGRAPHY



(some quasiprobability function in
the Stokes space)

$$\vec{n} = \{\cos \vartheta, \sin \vartheta \cos \varphi, \sin \vartheta \sin \varphi\}$$

$$p(s_{\vartheta\varphi}) = \int W(\vec{S}) \delta(\vec{S} \cdot \vec{n} - s_{\vartheta\varphi}) d^3 \vec{S}$$

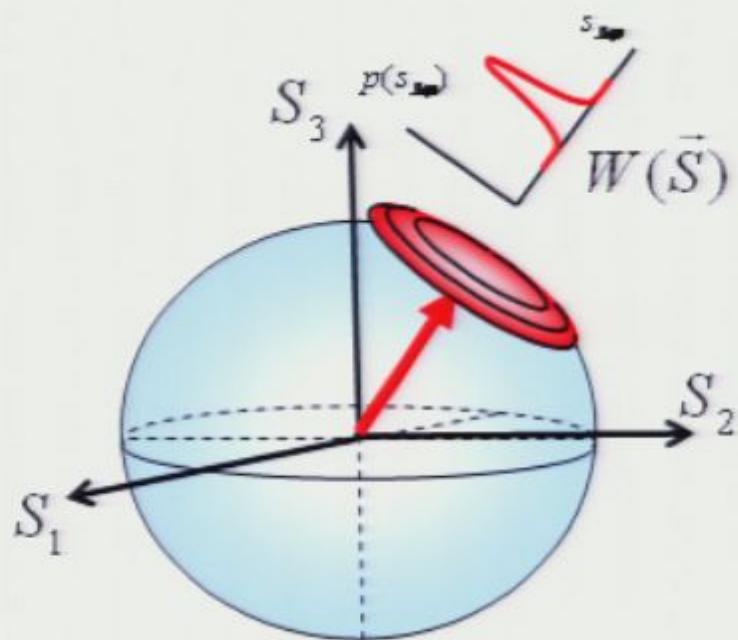
tomogram of a Stokes measurement

P.A. Bushev, V.P. Karassiov,
 A.V. Masalov, and A.A. Putilin,
Optics and Spectroscopy 91,
 526 (2001)

Reconstruction:

$$p(s_{\vartheta\varphi}) \rightarrow W(\vec{S})$$

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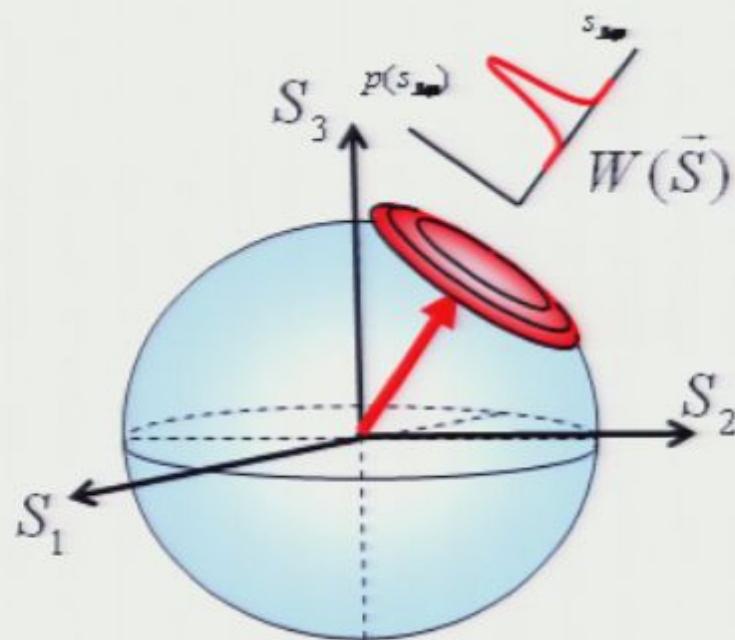
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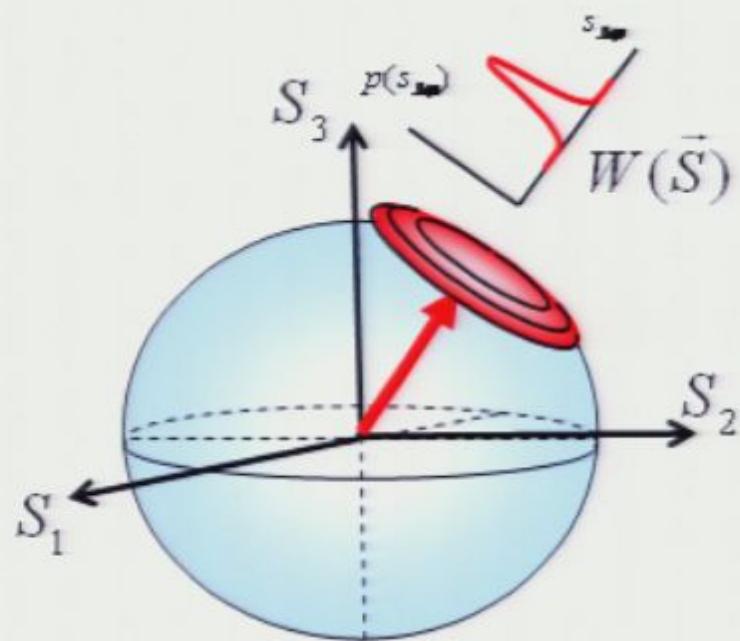
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$$\langle S(\theta, \varphi)^n \rangle$$



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RECONSTRUCTION OF QUASIPROBABILITY FUNCTIONS

W-function $W(\vec{S}) \equiv \int_{\mathbb{R}^3} \frac{d^3 \vec{\omega}}{(2\pi)^3} \chi(\vec{\omega}) e^{-i \vec{\omega} \cdot \vec{S}}, \quad \chi(\vec{\omega}) \equiv \left\langle e^{i \vec{\omega} \cdot \vec{S}} \right\rangle$

Reconstruction $W(\vec{S}) \propto \int_{-\infty}^{\infty} ds_{g\varphi} \int_{S_2} d^2 \vec{n} \frac{d^2 p(s_{g\varphi})}{ds_{g\varphi}^2} \delta(s_{g\varphi} - \vec{S} \cdot \vec{n})$
 (for $S_0 \gg 1$)

K.B. Wolf, Opt. Commun. 132, 343 (1996);

V.P. Karassiov and A.V. Masalov, J. Opt. B 4, 366 (2002)

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Q-function $Q(S_0, \vec{n}) = \langle S_0, \vec{n} | \hat{\rho}_{S_0} | S_0, \vec{n} \rangle$

Reconstruction

for $S_0 \gg 1$: $Q(S_0, \vec{n}) \propto (2S_0 + 1) \int_{-\infty}^{\infty} ds_{\mathcal{S}\varphi} \int_{S_2} d^2 \vec{n} \frac{d^2 p(s_{\mathcal{S}\varphi})}{ds_{\mathcal{S}\varphi}^2} \delta(s_{\mathcal{S}\varphi} - \vec{S} \cdot \vec{n})$



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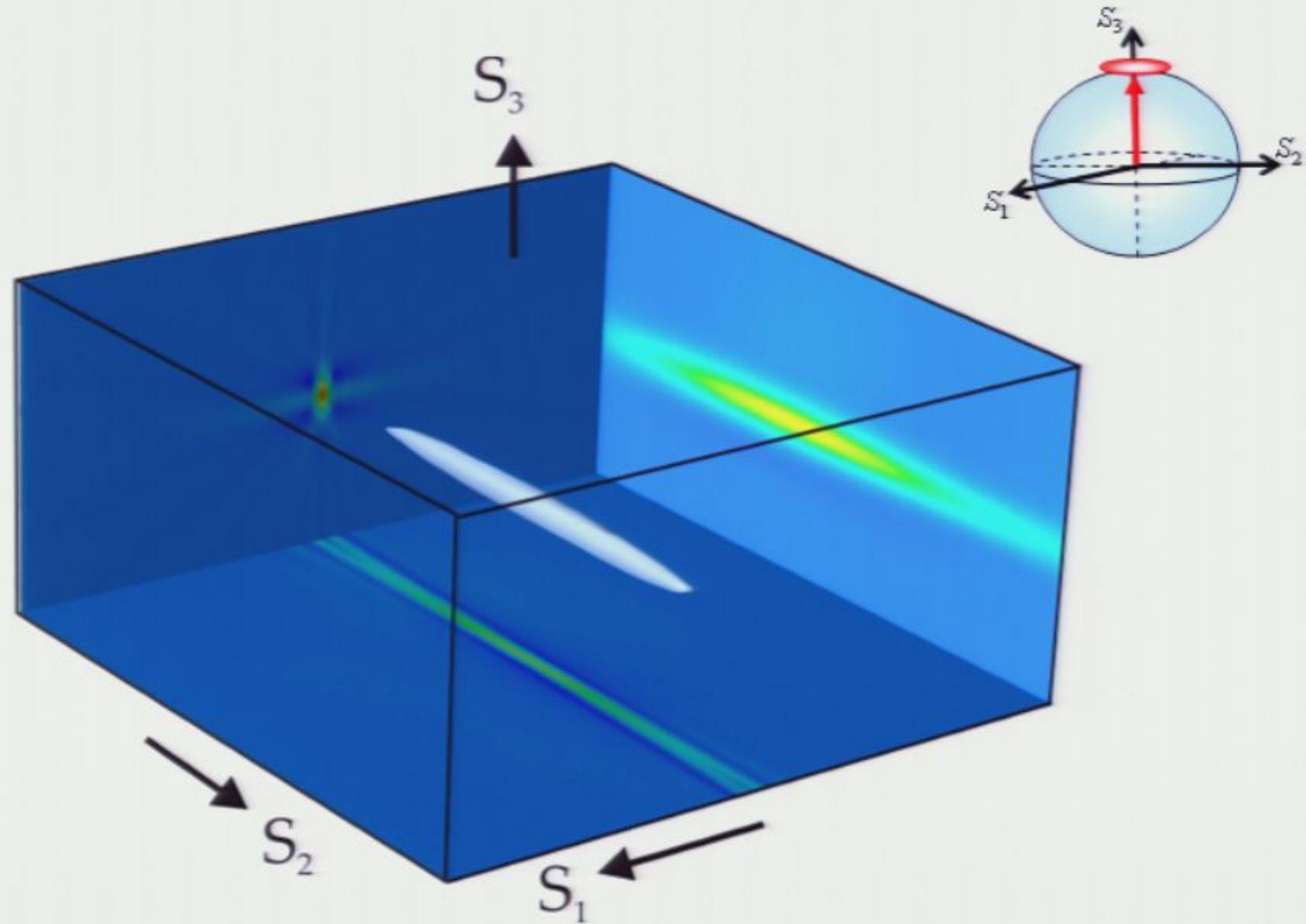
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Q-function $Q(S_0, \vec{n}) = \langle S_0, \vec{n} | \hat{\rho}_{S_0} | S_0, \vec{n} \rangle$

Reconstruction for $S_0 \gg 1$: $Q(S_0, \vec{n}) \propto (2S_0 + 1) \int_{-\infty}^{\infty} ds_{\mathcal{S}\varphi} \int_{S_2} d^2 \vec{n} \frac{d^2 p(s_{\mathcal{S}\varphi})}{ds_{\mathcal{S}\varphi}^2} \delta(s_{\mathcal{S}\varphi} - \vec{S} \cdot \vec{n})$

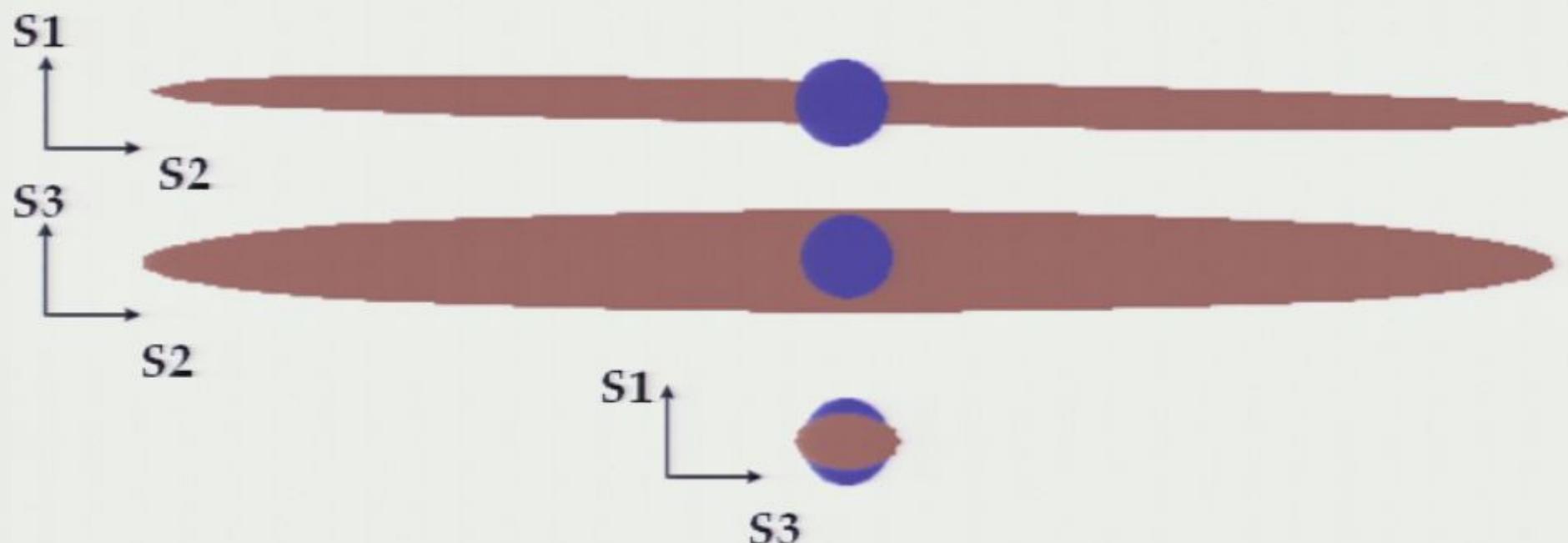
$S_0 = N \sim 10^7 \Rightarrow$ W-reconstruction ~ Q-reconstruction ~ classical 3D inverse Radon transformation

RESULTS: Q-FUNCTION RECONSTRUCTION





RESULTS: COMPARISON WITH A COHERENT STATE



6.2 dB squeezing almost along S_1

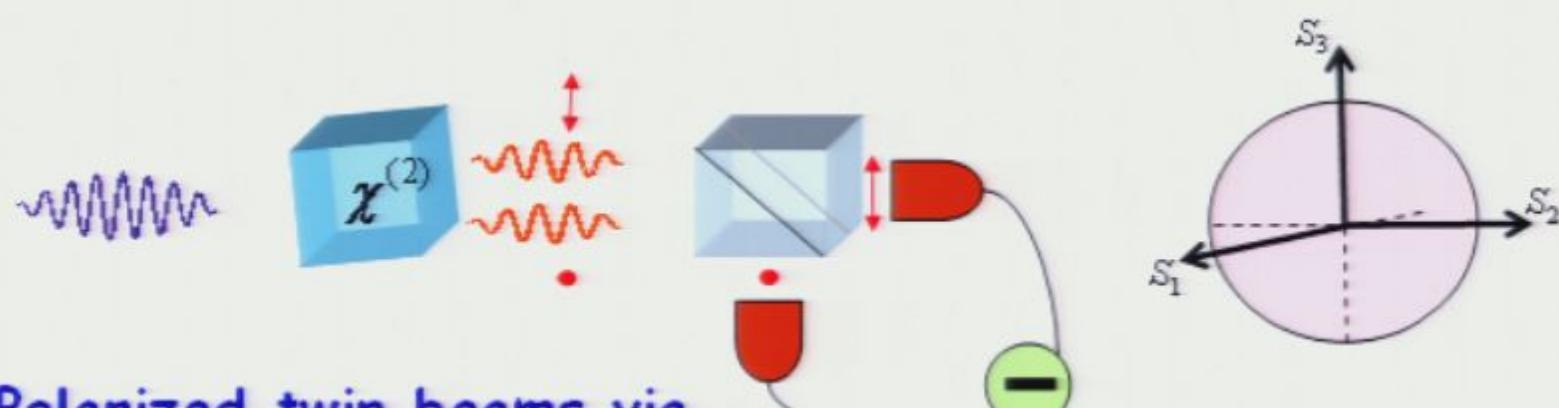
Much larger anti-squeezing, due to phase noise

Slightly above shot noise along S_3

FROM MACROSCOPIC TO MESOSCOPIC SCALE

Further questions:

1. What happens if S_0 reduces, so that $\text{Var}(S_i)$ gets closer to S_0 ?
2. What happens if the classical polarized component reduces, $\text{Var}(S_i)$ gets closer to S_3 ?
3. Can we prepare arbitrary shapes of $Q(S)$ and how can we use them?

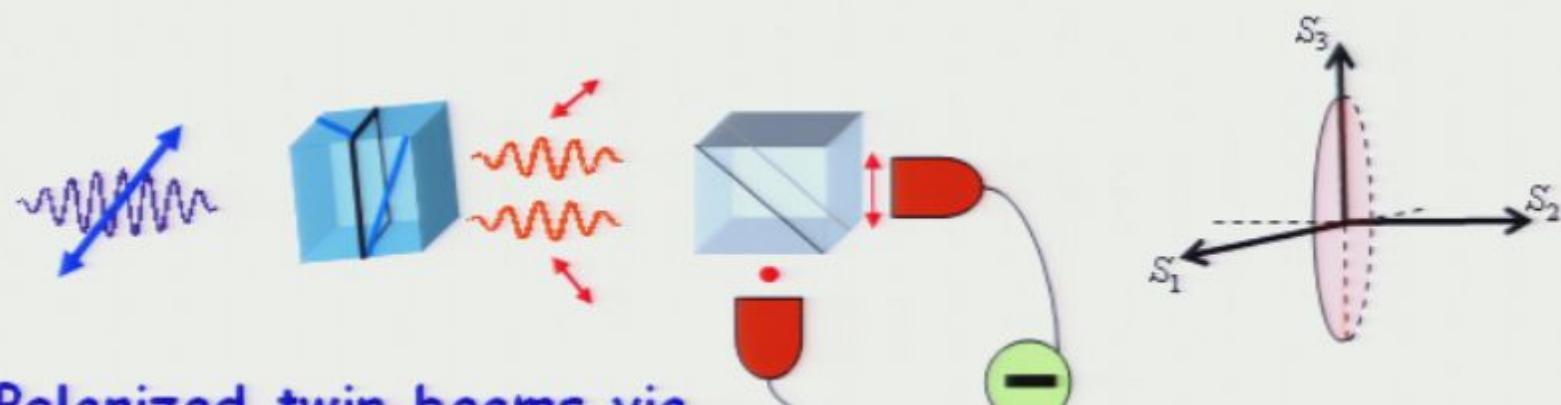


Polarized twin beams via travelling-wave OPA

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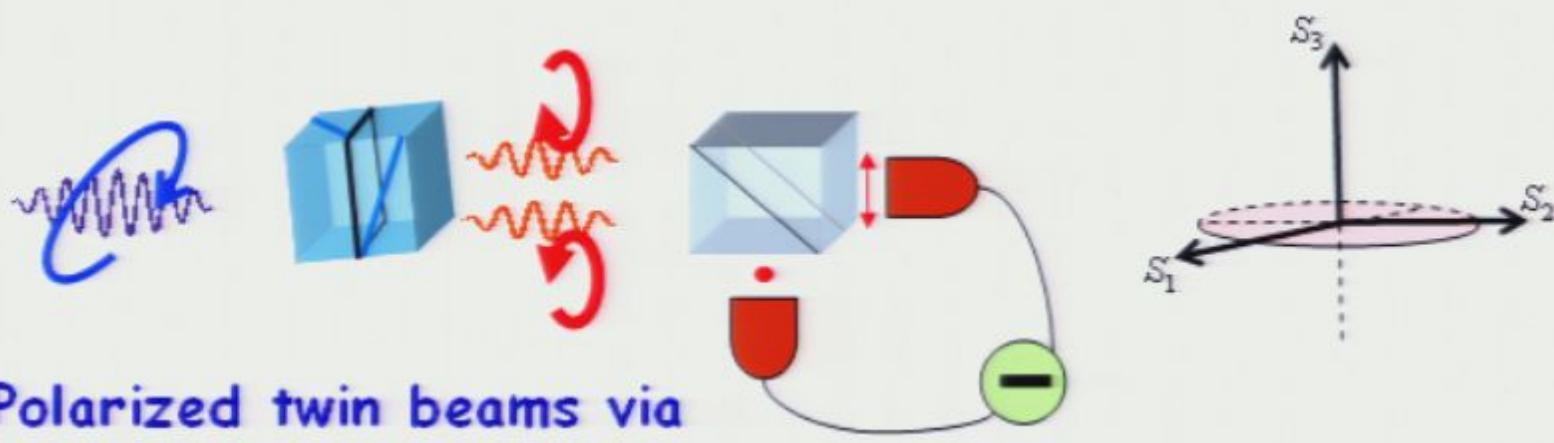


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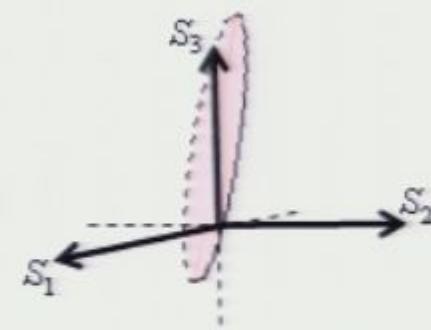
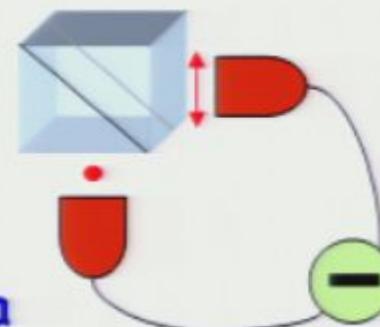
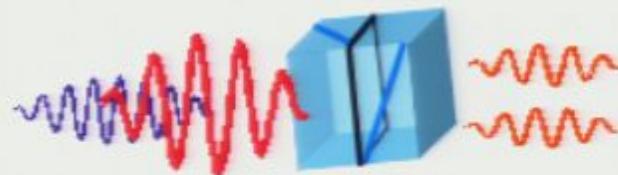
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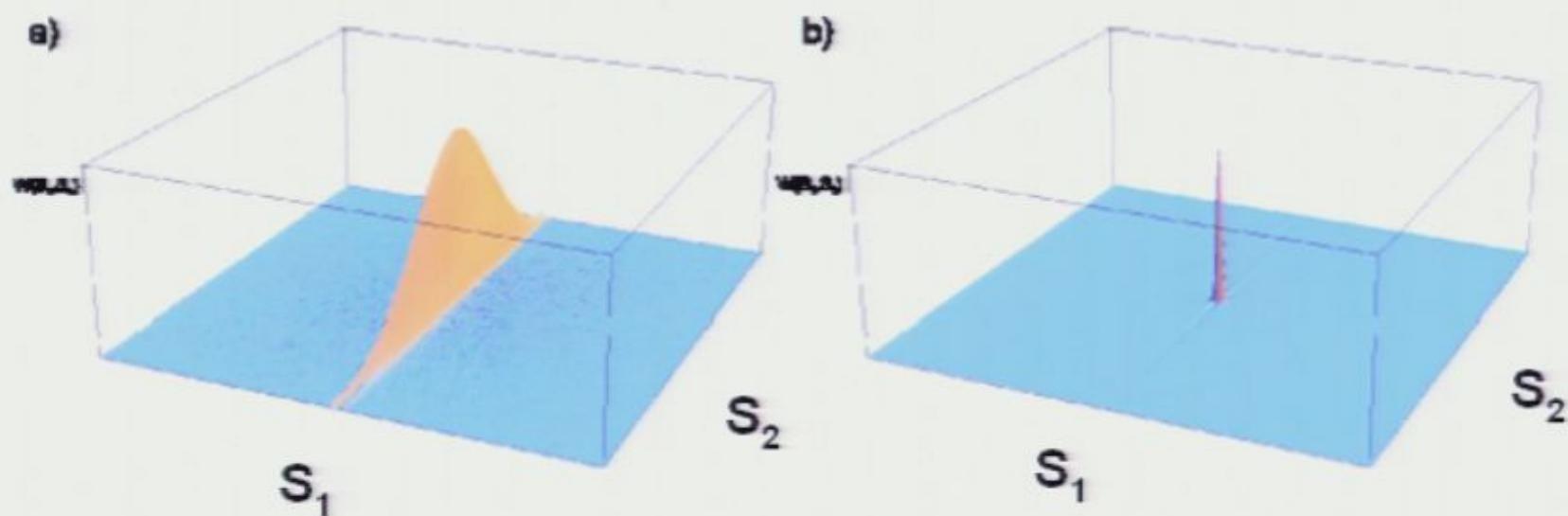
LATEST NEWS

from Zdenek Hradil and
Jaroslav Rehacek (Palacky
University, Olomouc,
Czech Republic)

Problem: negative 'ripples' created by
the inverse Radon transformation in
the Q-function (also in W-function)

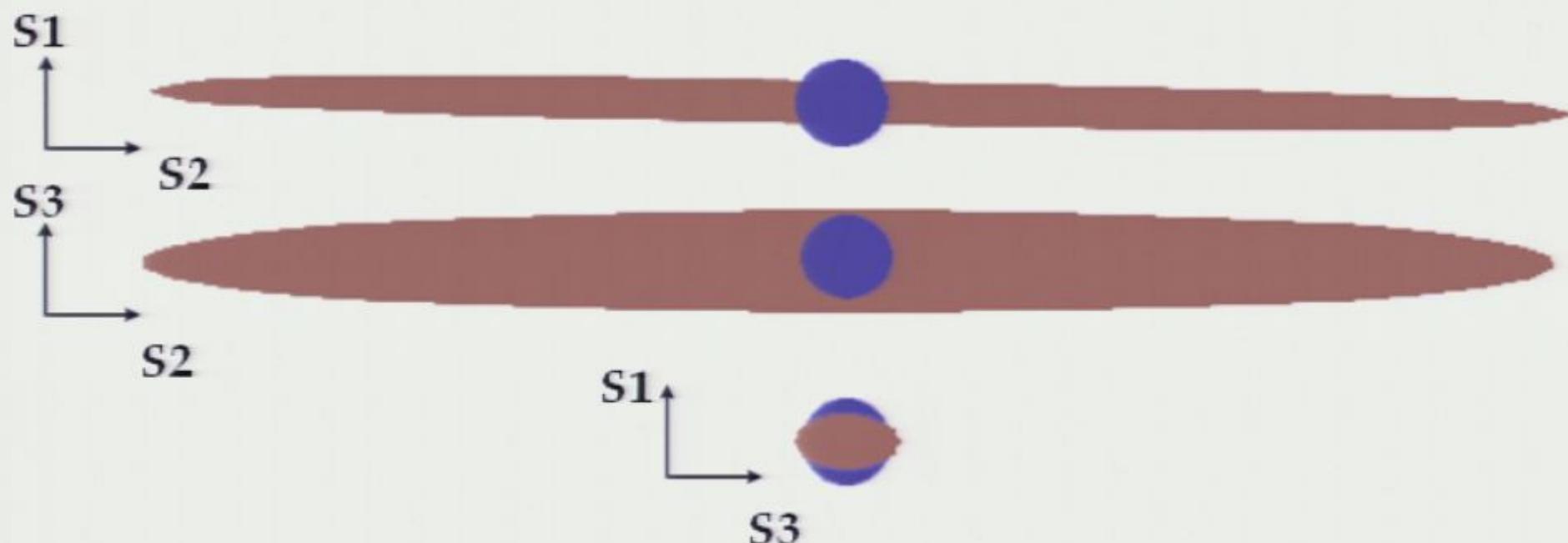
Solution: Maximum-likelihood approach
eliminates the Q-function negativities

RADON-TRANSFORMATION RESULTS





RESULTS: COMPARISON WITH A COHERENT STATE

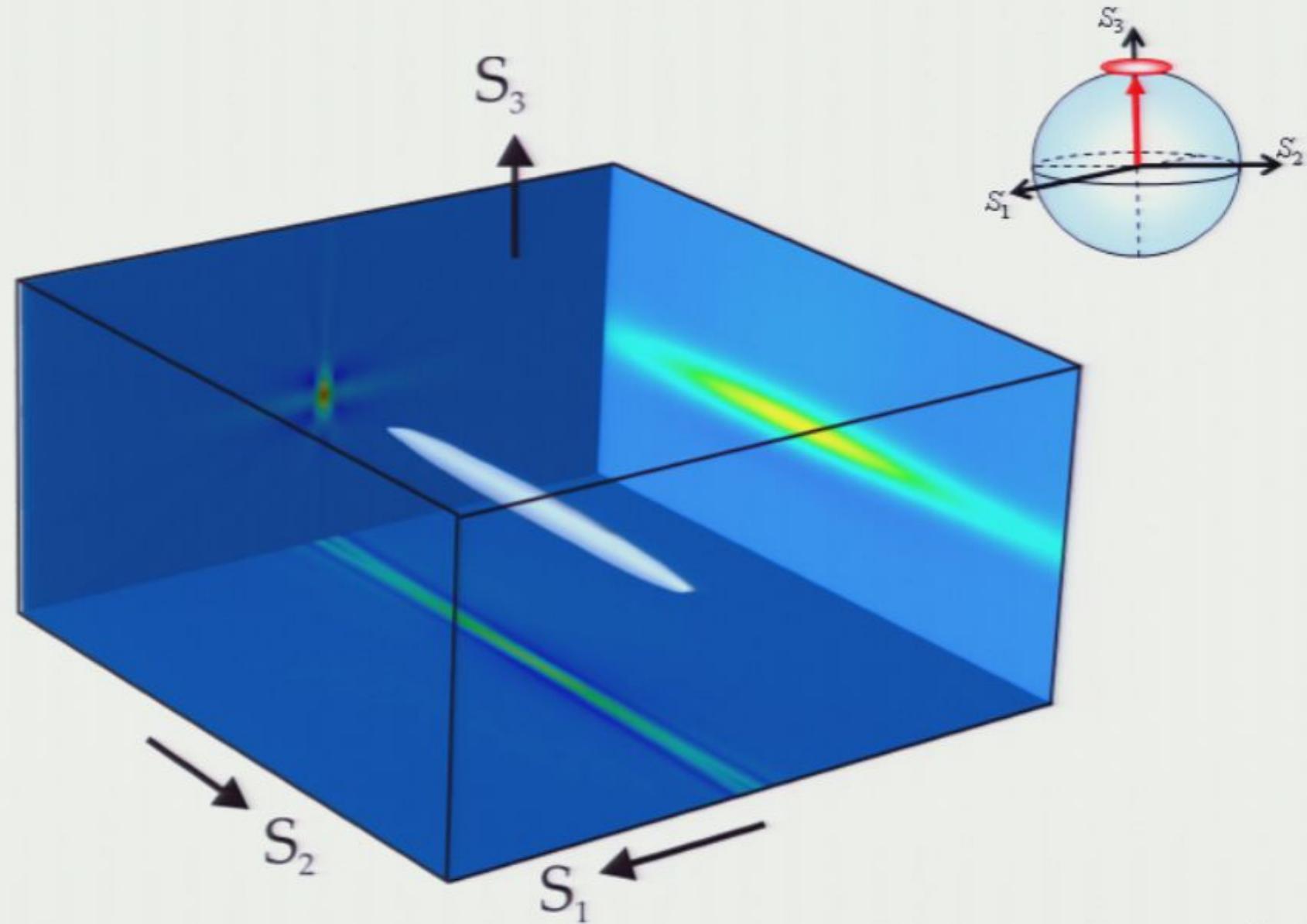


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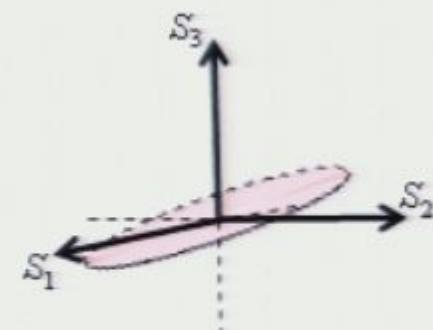
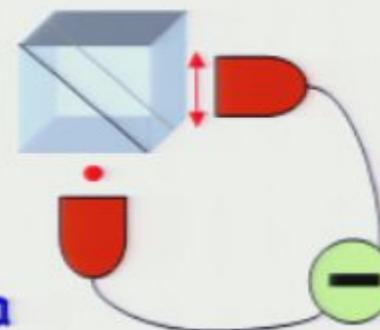
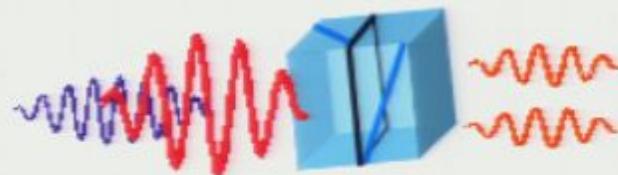
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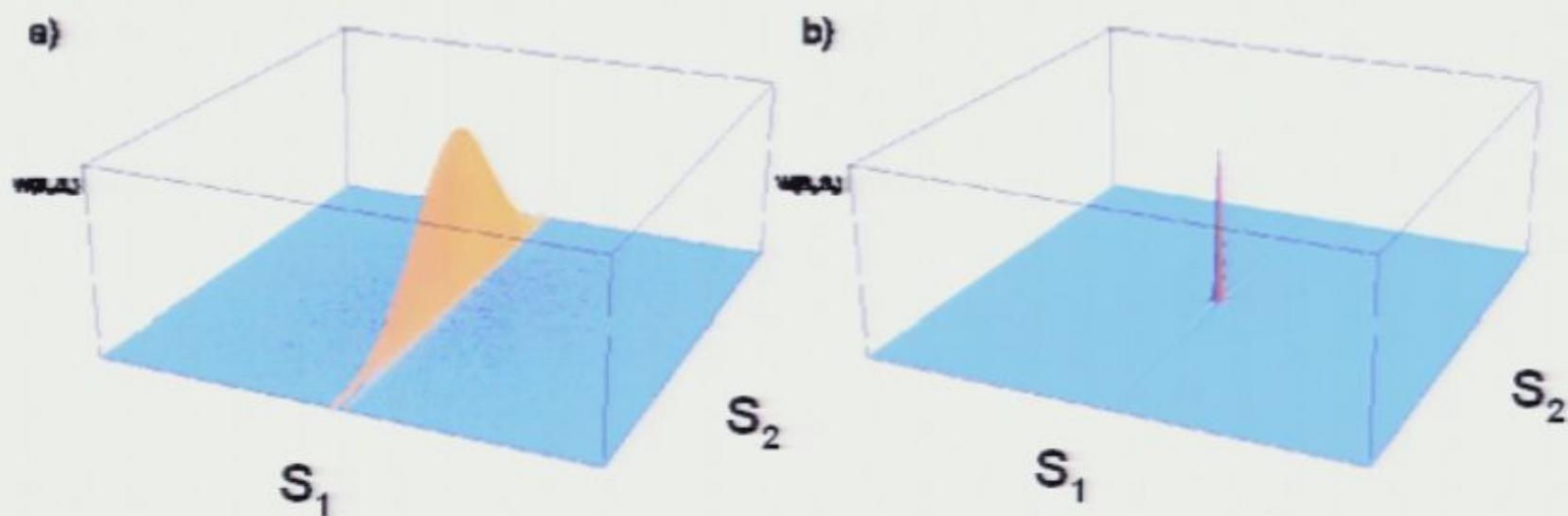
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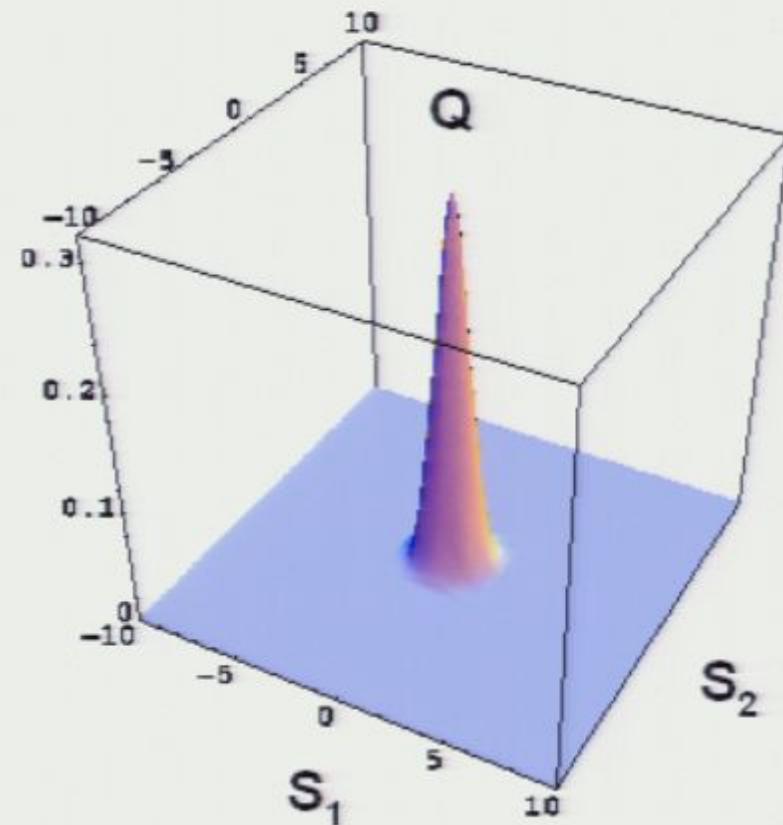
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the Q-function (also in W-function)

Solution: Maximum-likelihood approach
eliminates the Q-function negativities

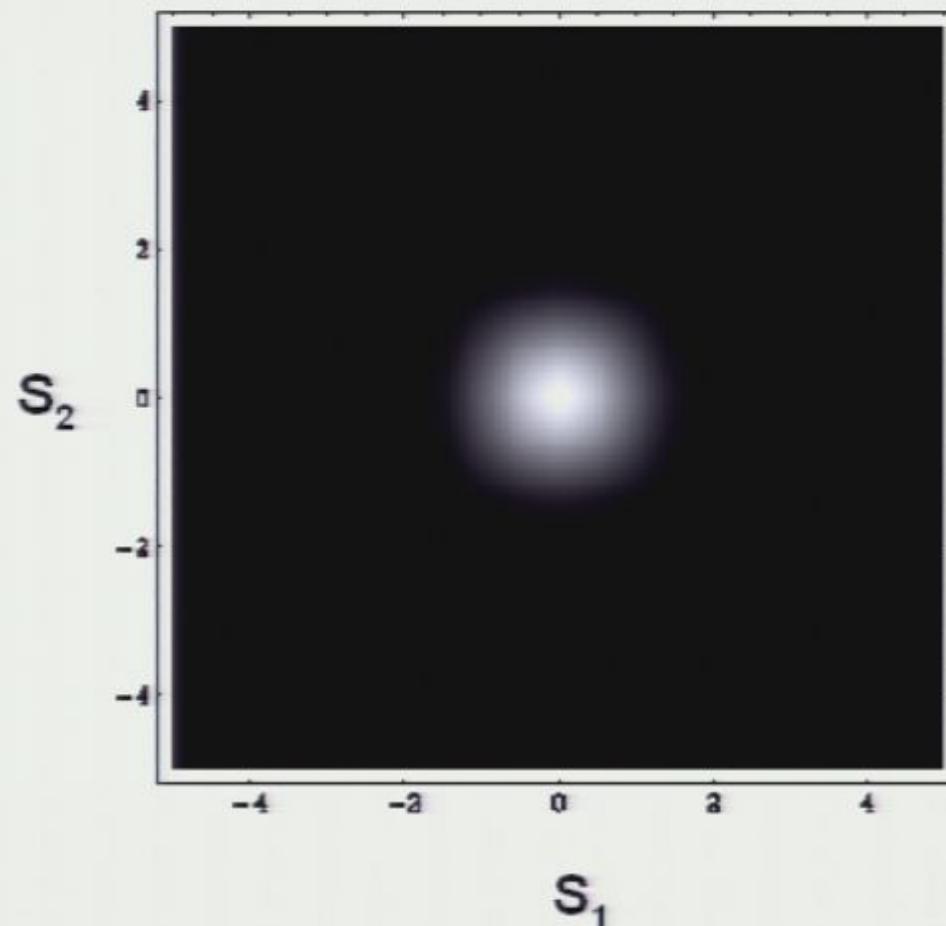
RADON-TRANSFORMATION RESULTS



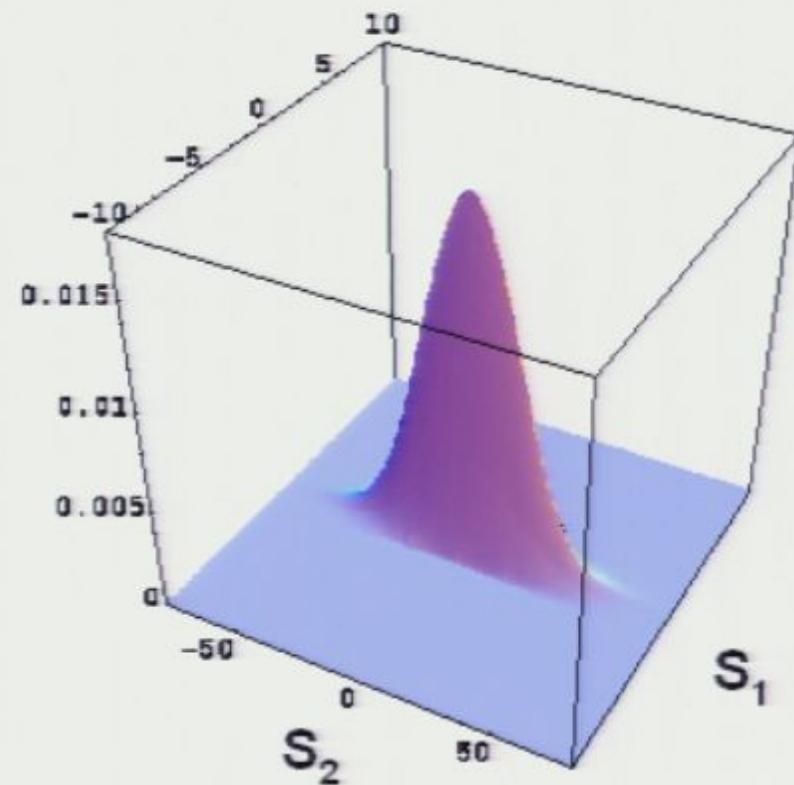
'DARK-PLANE TOMOGRAPHY': COHERENT STATE



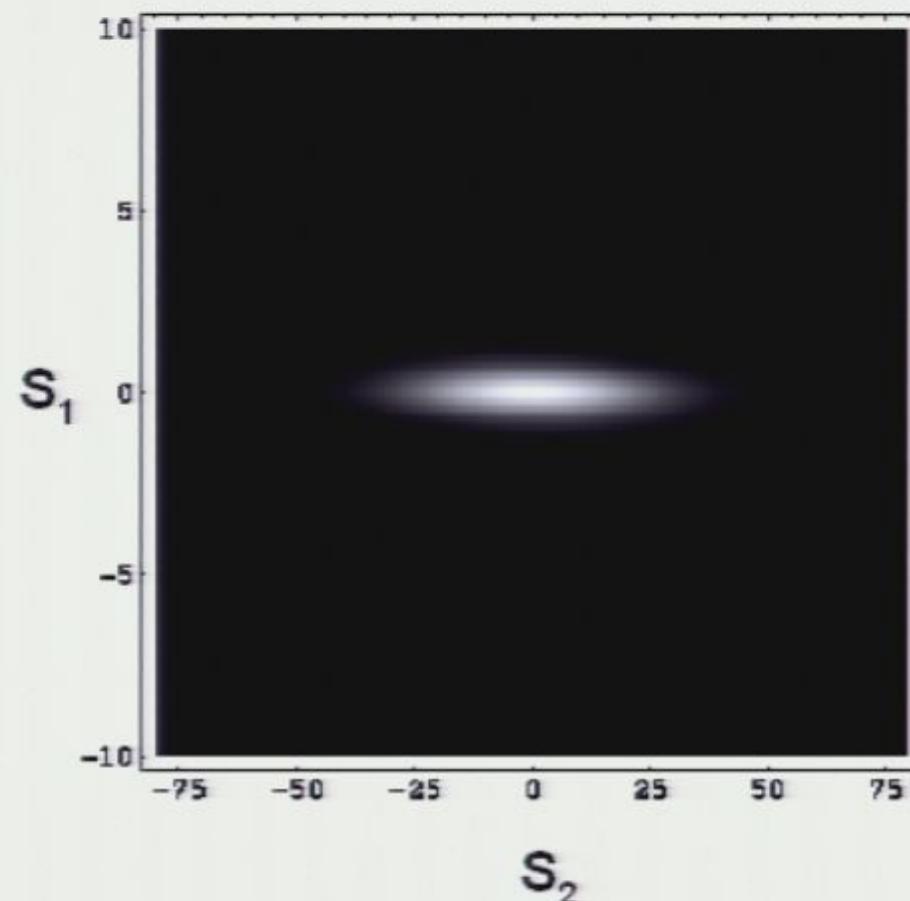
'DARK-PLANE TOMOGRAPHY': COHERENT STATE



'DARK-PLANE TOMOGRAPHY': SQUEEZED STATE



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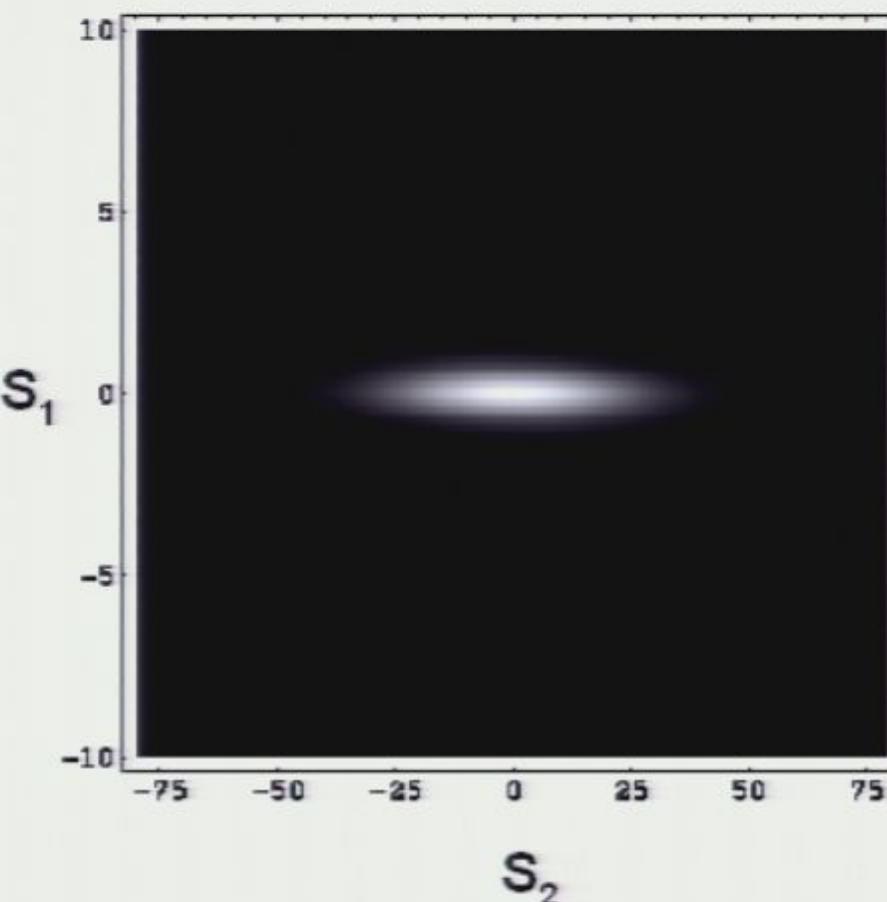




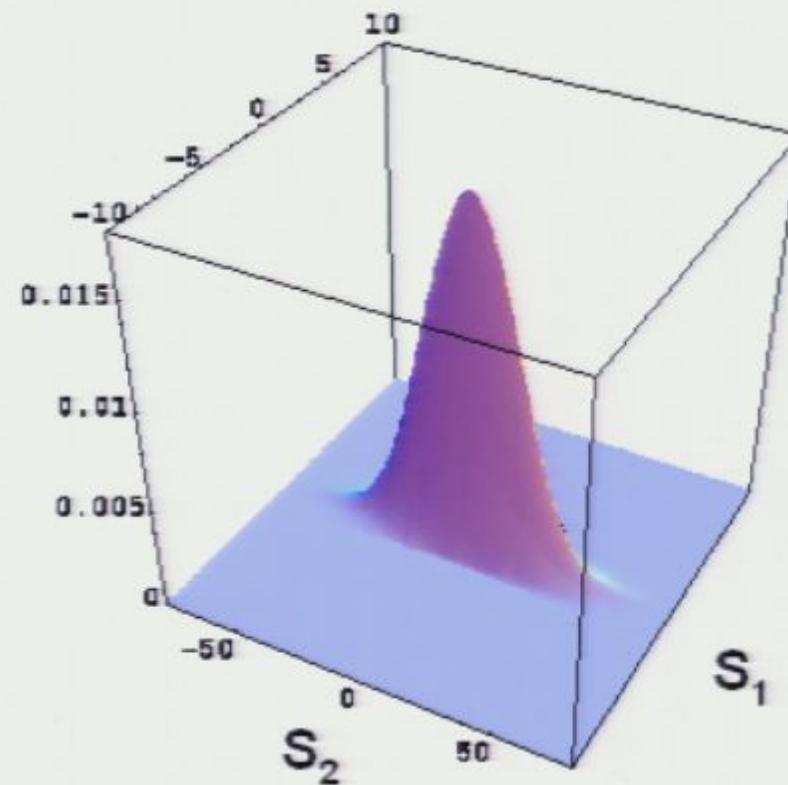
CONCLUSIONS

1. Performed: full polarization tomography of an intense squeezed state generated via third-order nonlinearity in an optical fibre. In the 'macroscopic' limit, reconstruction of W- and Q-quasiprobability functions is similar and coincides with the classical 3D inverse Radon transform. We now have a tool for polarization state reconstruction.
2. Planned: generation and full polarization tomography of 'mesoscopic' states of polarization-squeezed vacuum with and without polarized coherent component.

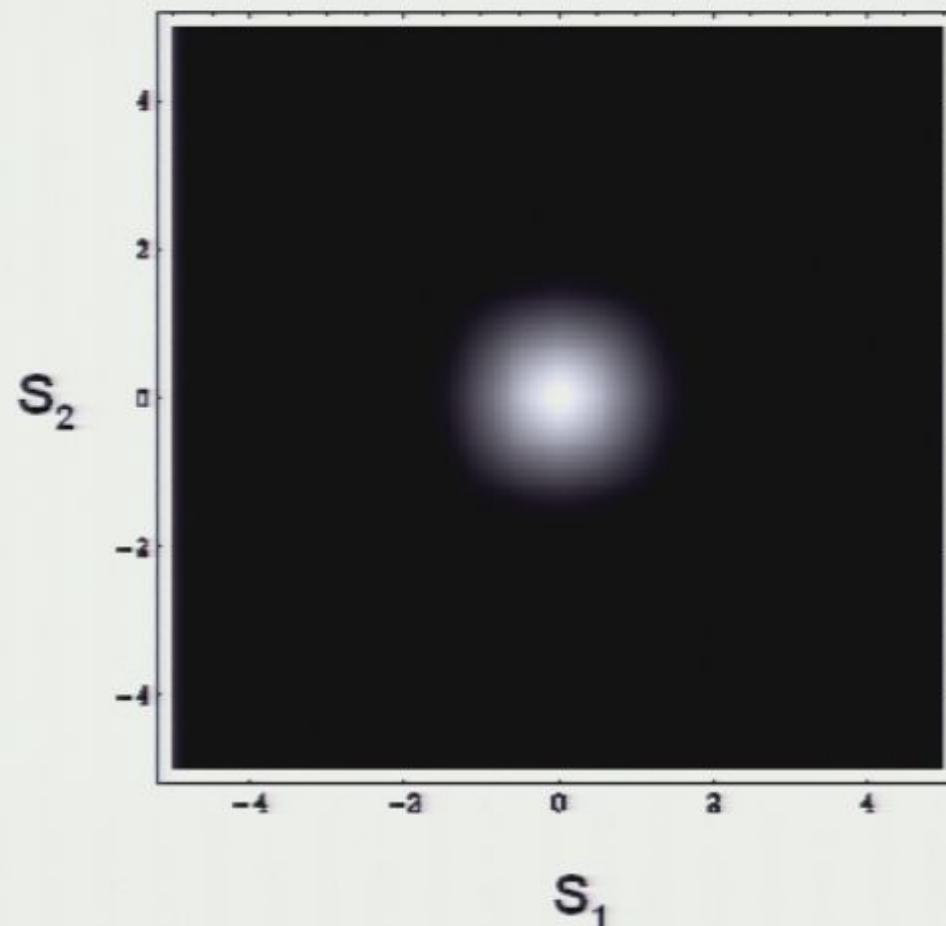
**THANK YOU FOR YOUR
ATTENTION!**



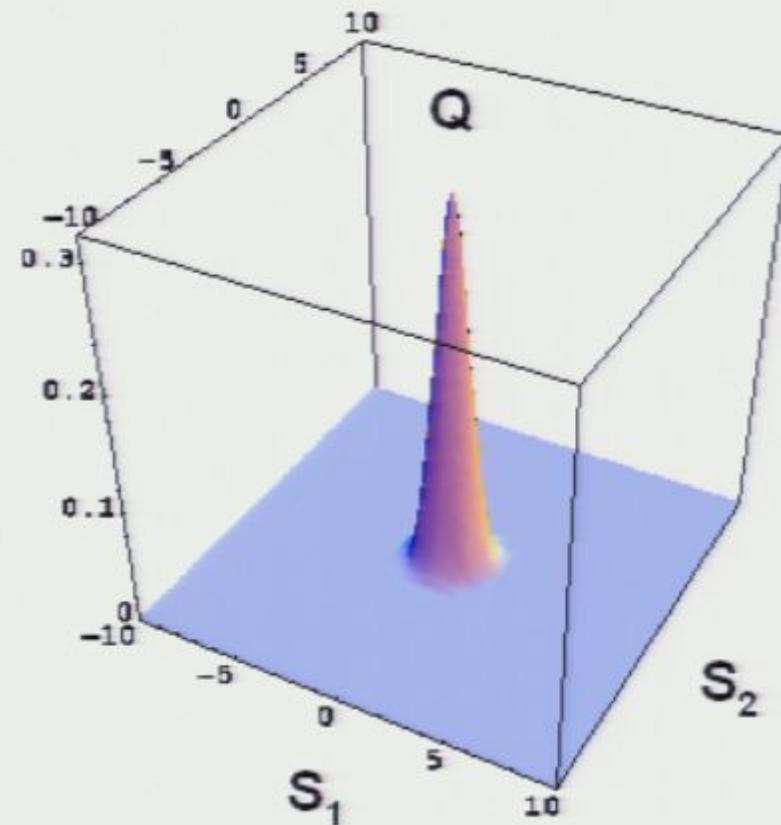
'DARK-PLANE TOMOGRAPHY': SQUEEZED STATE



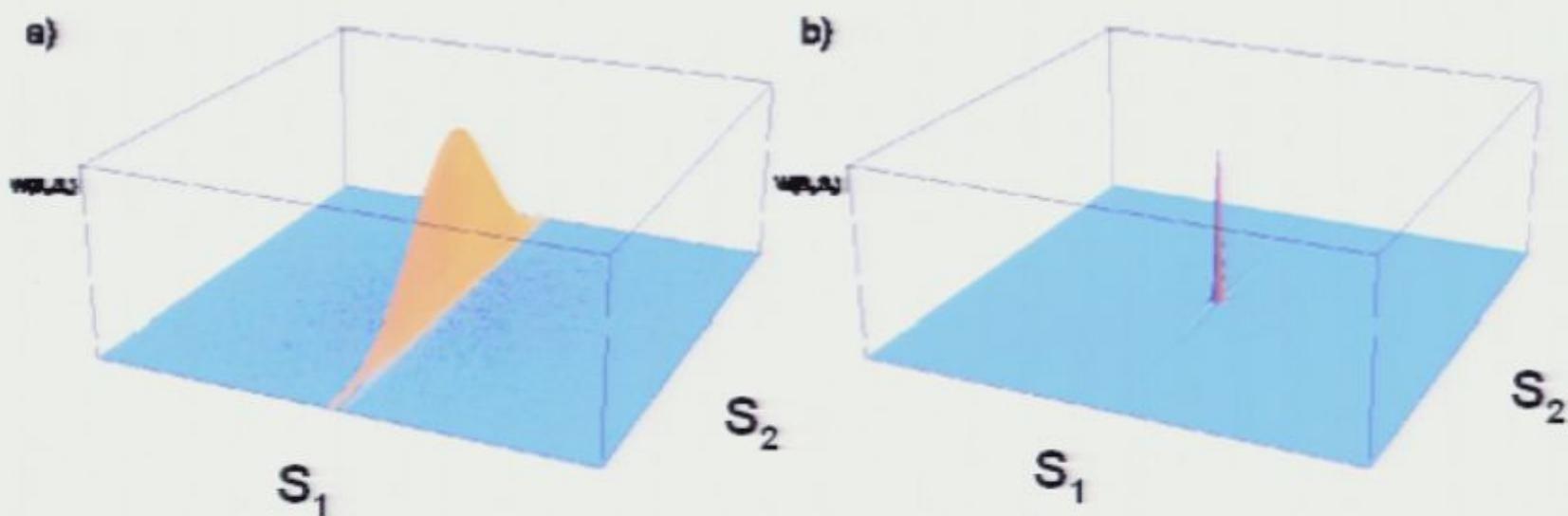
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RADON-TRANSFORMATION RESULTS



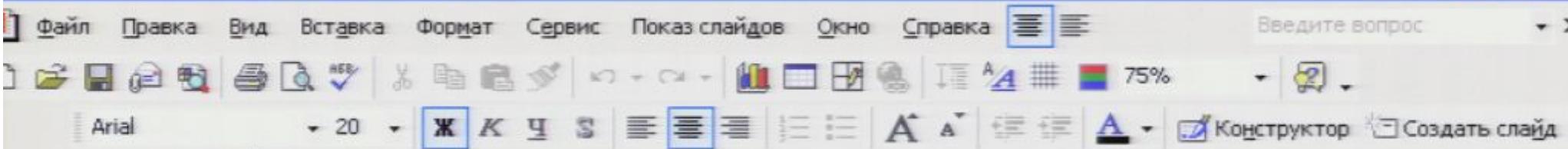


LATEST NEWS

from Zdenek Hradil and
Jaroslav Rehacek (Palacky
University, Olomouc,
Czech Republic)

Problem: negative 'ripples' created by
the inverse Radon transformation in
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POLARIZATION TOMOGRAPHY OF MACRO- AND MESOSCOPIC STATES OF LIGHT

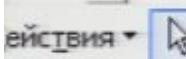
Christoph Marquardt¹, Joel Heersink¹, Ruifang Dong¹,
Maria V. Chekhova^{1,2}, Andrei B. Klimov³,
Luis L. Sánchez-Soto⁴, Ulrik L. Andersen¹,
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Заметки к слайду



Автофигуры

