

Title: Linear Optics Quantum Process Tomography

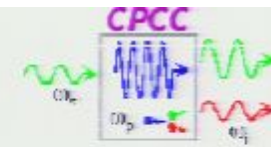
Date: Aug 27, 2008 02:30 PM

URL: <http://pirsa.org/08080047>

Abstract: The field of linear optics quantum computing (LOQC) allows the construction of conditional gates using only linear optics and measurement. This quantum computing paradigm bypasses a seemingly serious problem in optical quantum computing: it appears to be very hard to produce a meaningful interaction between two single photons. But what if this obstacle were instead an advantage? By assuming that none of the physical components that make up an LOQC gate produce a direct photon-photon interaction, we dramatically reduce the space of gates which are possible for a given number of input and output qubits. In fact, by parametrizing a gate according to its action on single photons, instead of on multiple photons, it is possible to exponentially reduce the number of measurements necessary to fully characterize an LOQC gate. In addition, this approach to LOQC process tomography may have additional experimental advantages when non-ideal input states are used for this characterization.



Quantum Estimation: Theory and Practice, 27 August 2008
Perimeter Institute, Waterloo



Linear Optics Quantum Process Tomography

Joe Altepeter, Milja Medic, and Prem Kumar

Center for Photonic Communication and Computing
EECS and Physics Departments, Northwestern University

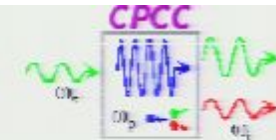
Evan Jeffrey

Department of Physics
University of Leiden, The Netherlands



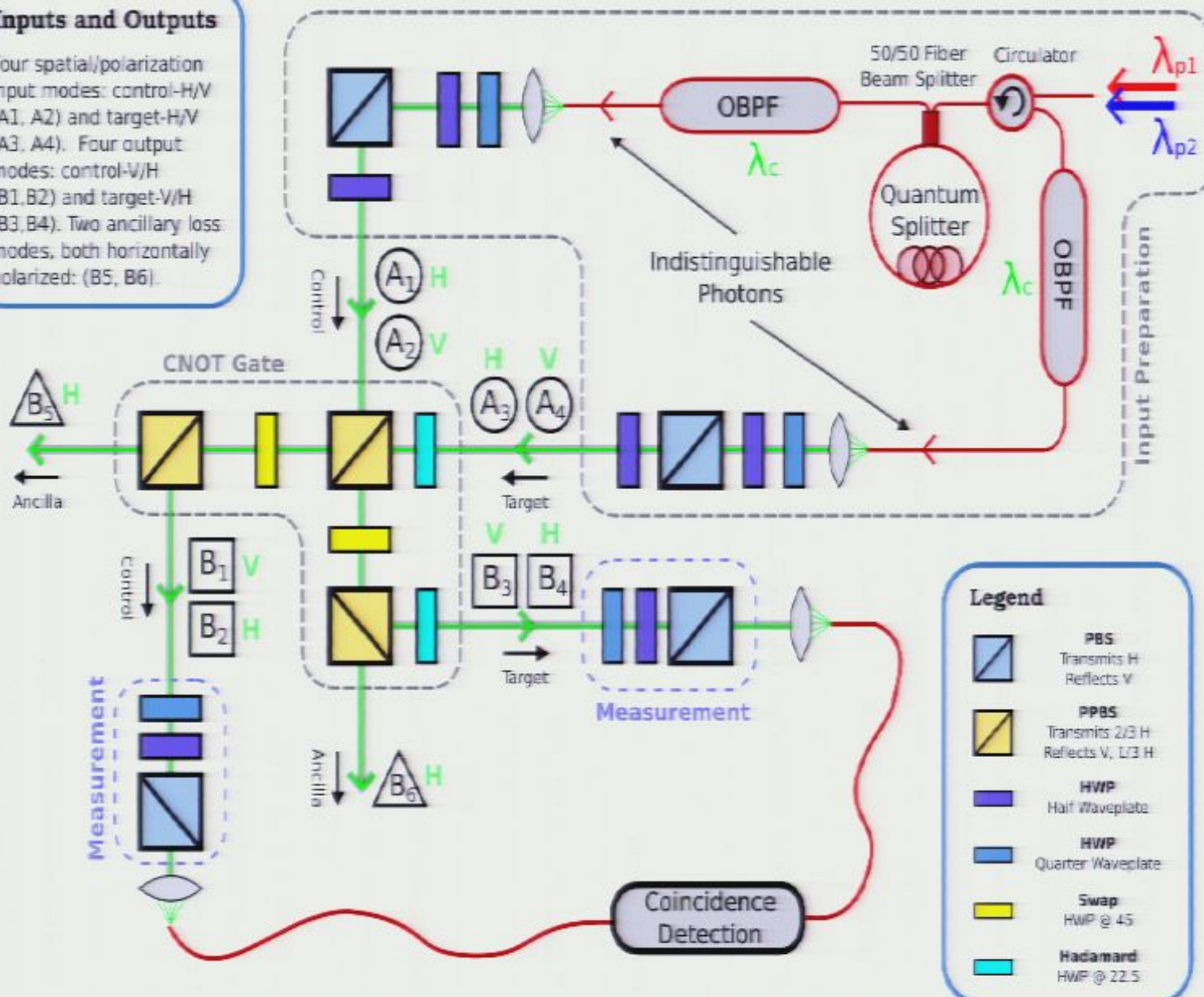


A Telecom-band CNOT Gate



Inputs and Outputs

Four spatial/polarization input modes: control-H/V (A1, A2) and target-H/V (A3, A4). Four output modes: control-V/H (B1, B2) and target-V/H (B3, B4). Two ancillary loss modes, both horizontally polarized: (B5, B6).

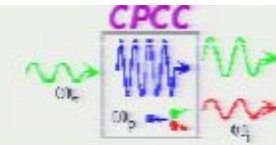


Legend

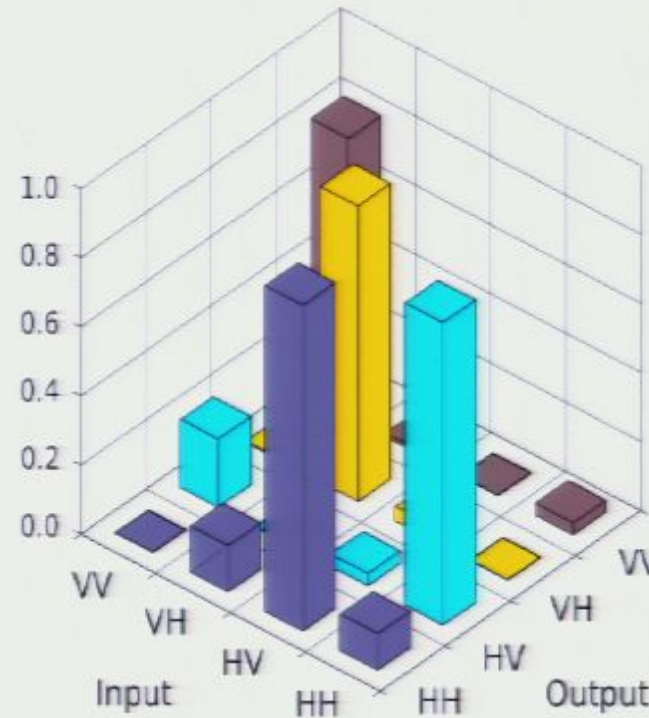
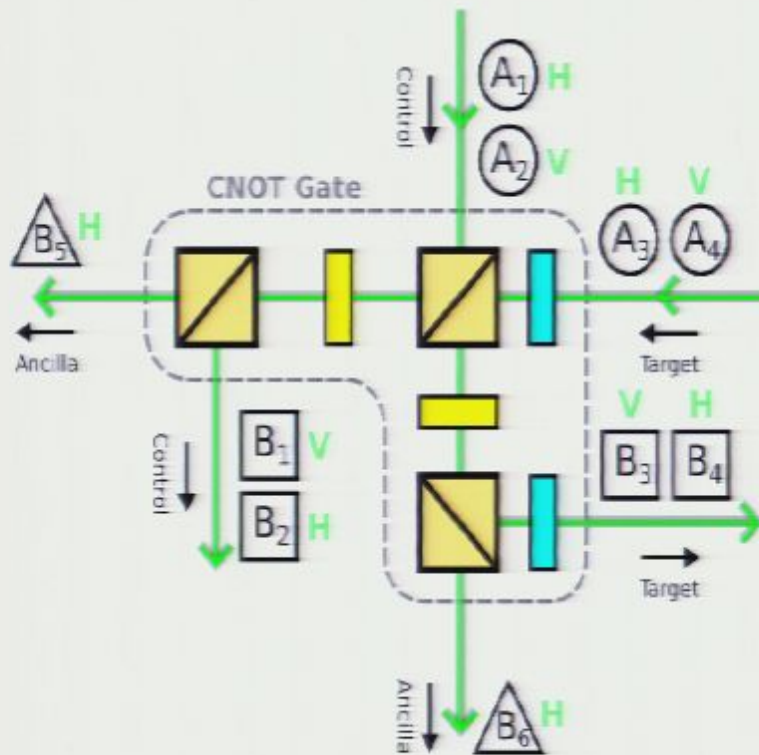
- PBS**
Transmits H
Reflects V
- PPBS**
Transmits 2/3 H
Reflects V, 1/3 H
- HWP**
Half Waveplate
- HWP**
Quarter Waveplate
- Swap**
HWP @ 45
- Hadamard**
HWP @ 22.5



First Experimental Results



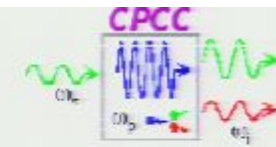
- Errors not from mode mismatch.



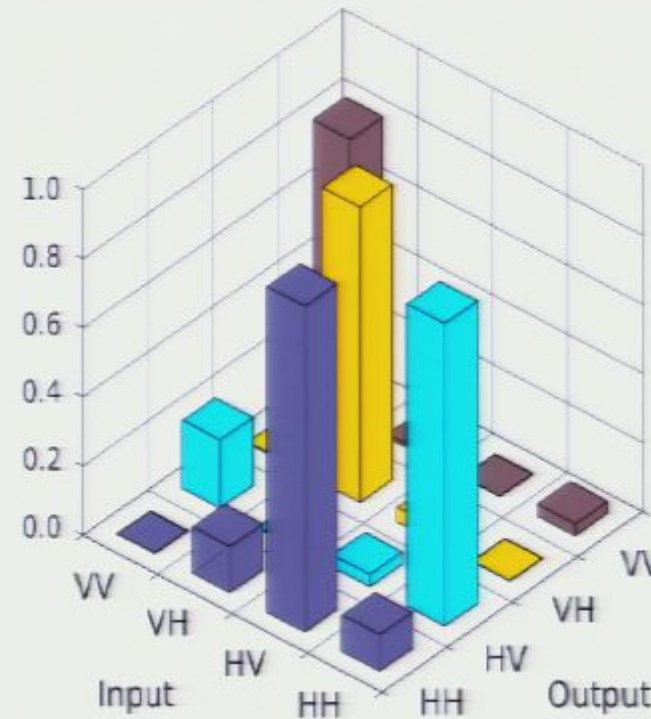
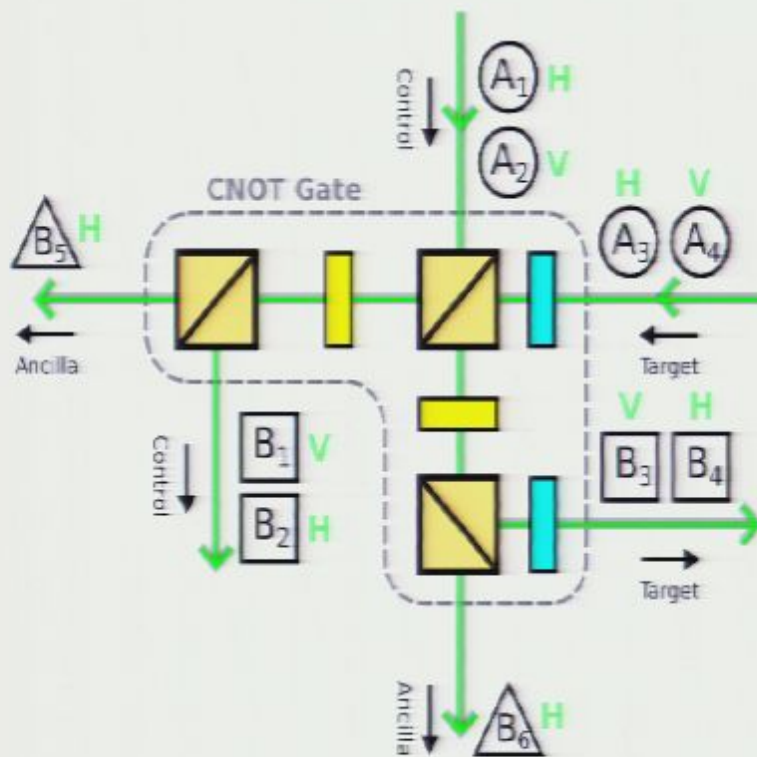
State	F	T	S_L
HH+VV	0.80(2)	0.43(6)	0.42(4)
HH-VV	0.75(3)	0.30(5)	0.53(4)
HV+VH	0.75(3)	0.34(6)	0.53(4)
HV-VH	0.80(2)	0.29(4)	0.49(3)



First Experimental Results



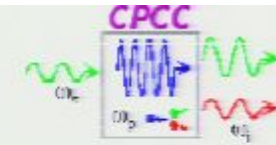
- Errors not from mode mismatch.
- We can't take more data.



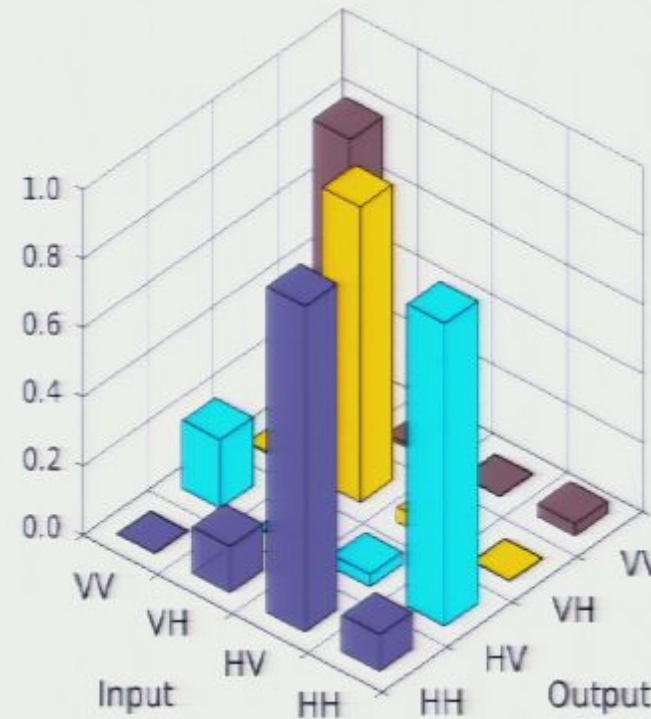
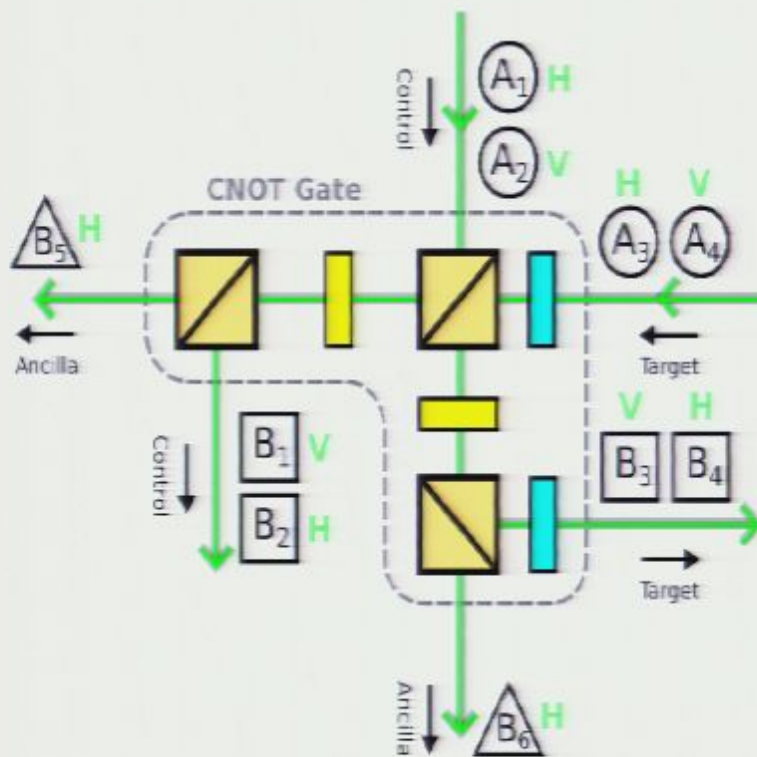
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First Experimental Results



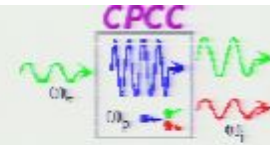
- Errors not from mode mismatch.
- We can't take more data.
- **What about multi-pairs?**



State	F	T	S_L
HH+VV	0.80(2)	0.43(6)	0.42(4)
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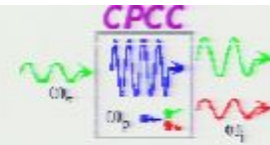
Explaining the Multipair Errors



How does a CNOT gate operate on two photons?



Explaining the Multipair Errors



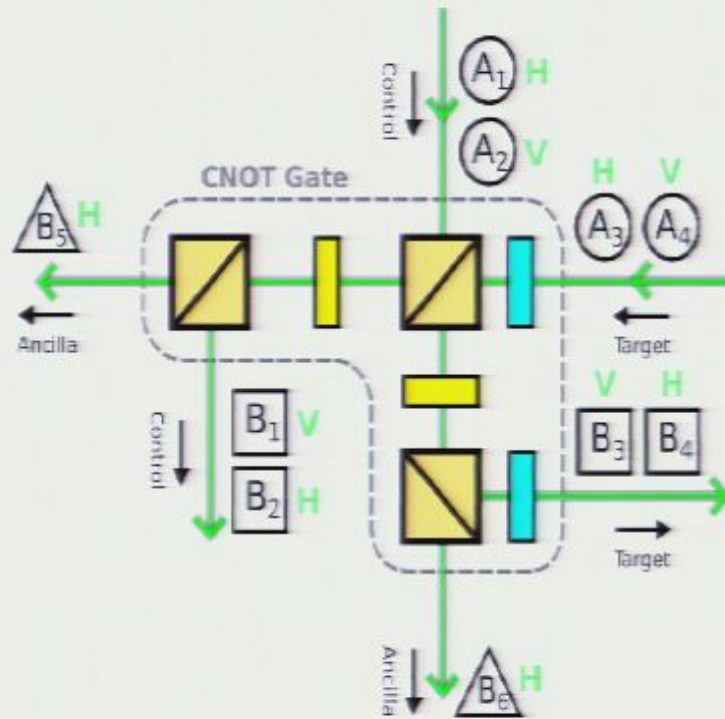
How does a CNOT gate operate on two photons?

$$\textcircled{A_1}^\dagger \textcircled{A_3}^\dagger \rightarrow \boxed{B_1}^\dagger \boxed{B_3}^\dagger$$

$$\textcircled{A_1}^\dagger \textcircled{A_4}^\dagger \rightarrow \boxed{B_1}^\dagger \boxed{B_4}^\dagger$$

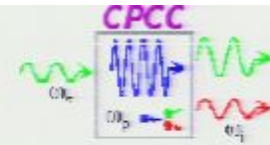
$$\textcircled{A_2}^\dagger \textcircled{A_3}^\dagger \rightarrow \boxed{B_2}^\dagger \boxed{B_4}^\dagger$$

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Explaining the Multipair Errors



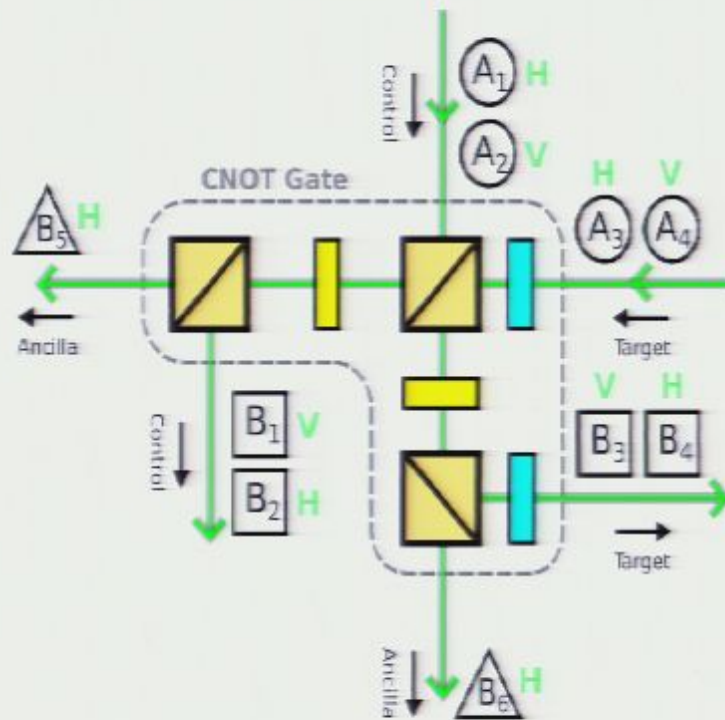
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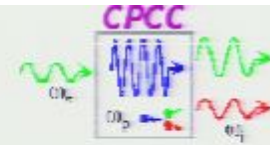
$$\textcircled{A_2}^\dagger \textcircled{A_4}^\dagger \rightarrow \boxed{B_2}^\dagger \boxed{B_3}^\dagger$$



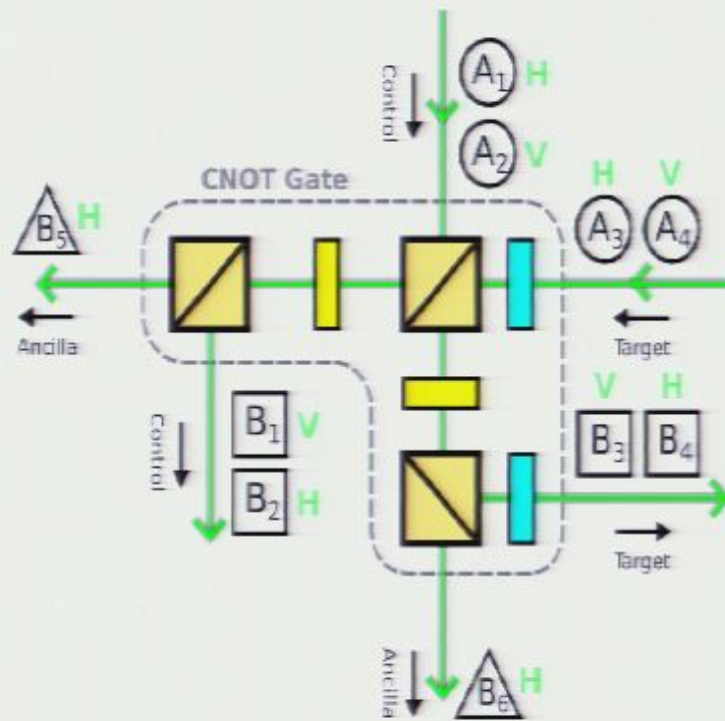
How do we use this to predict the CNOT's effect on four input photons? **Seems like a hard problem...**



Back to the Drawing Board

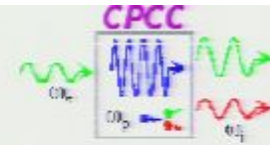


What about a single photon model?





Back to the Drawing Board



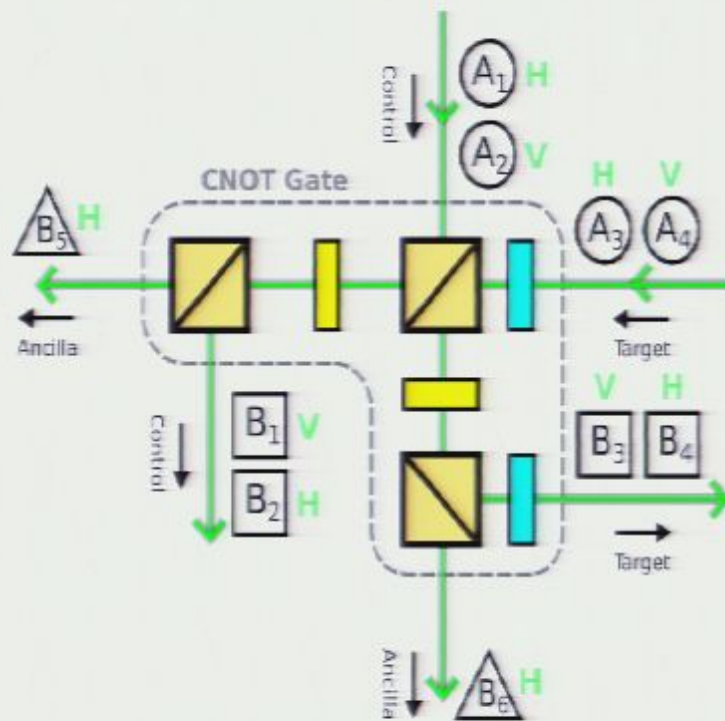
What about a single photon model?

$$A_1^\dagger \rightarrow \sqrt{\frac{1}{3}} (B_1^\dagger + B_4^\dagger - B_3^\dagger)$$

$$A_2^\dagger \rightarrow \sqrt{\frac{1}{3}} (-\sqrt{2} B_5^\dagger - B_2^\dagger)$$

$$A_3^\dagger \rightarrow \sqrt{\frac{1}{3}} (B_1^\dagger - B_4^\dagger - B_6^\dagger)$$

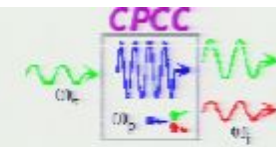
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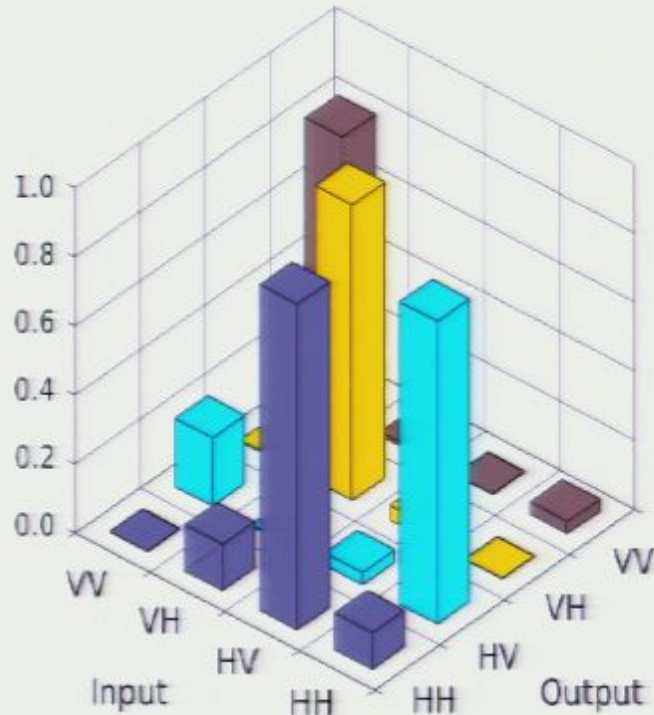
Now we can calculate how **any number of photons** will be affected by the gate.



Corrected Experimental Results



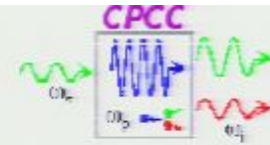
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Back to the Drawing Board



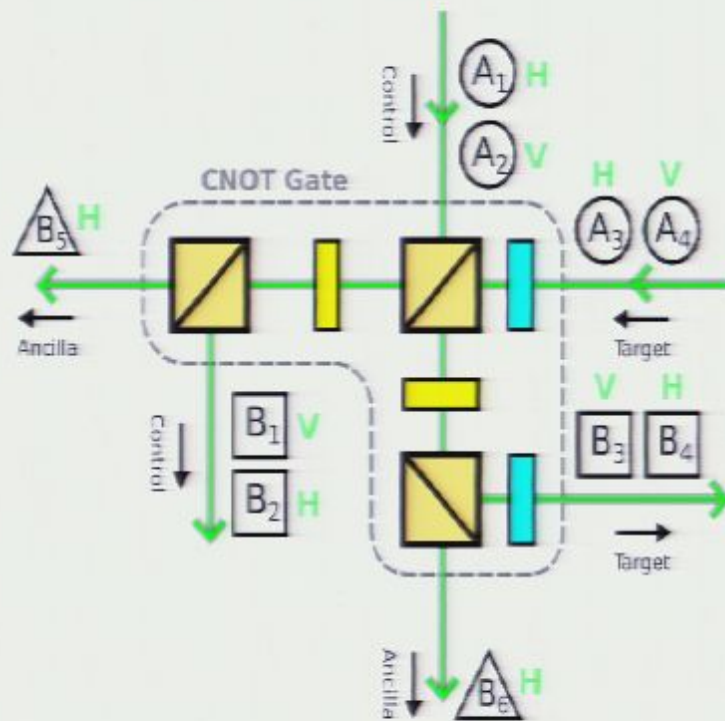
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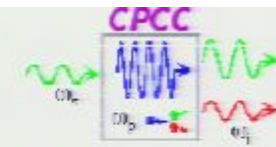
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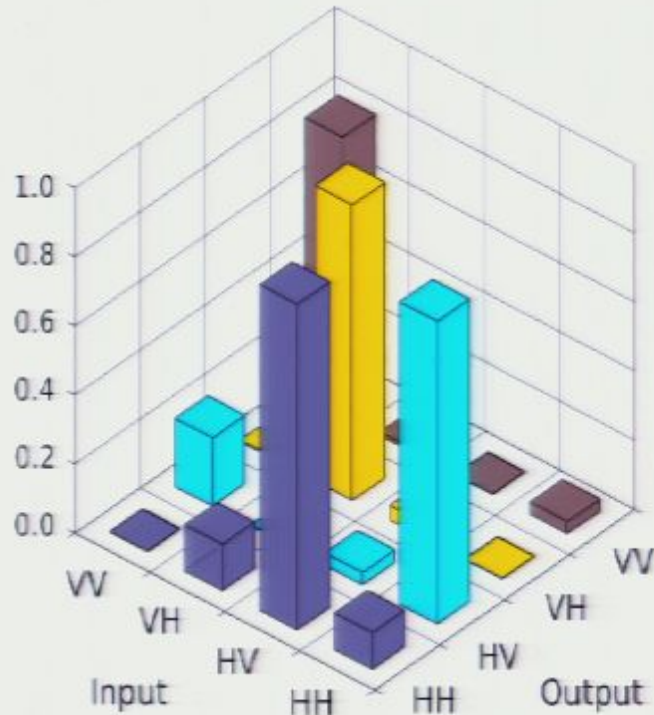
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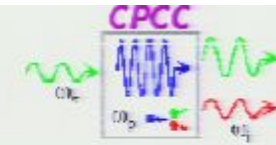
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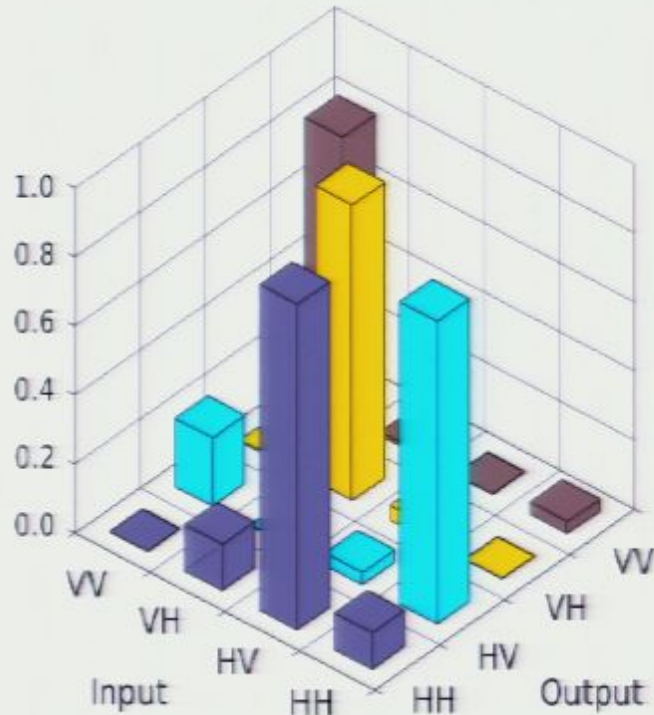
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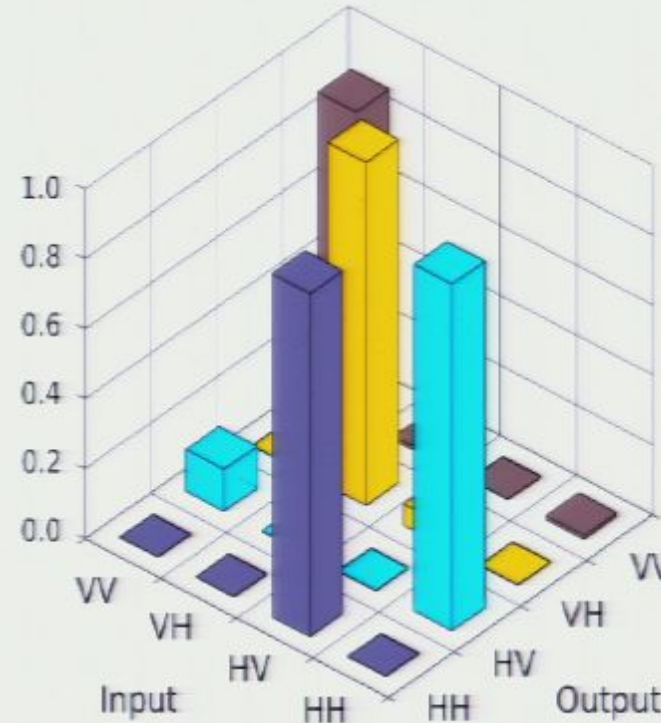
Corrected Experimental Results



Original Results



Multi-pairs Subtracted

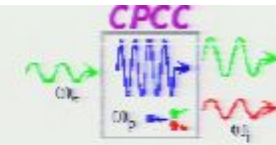


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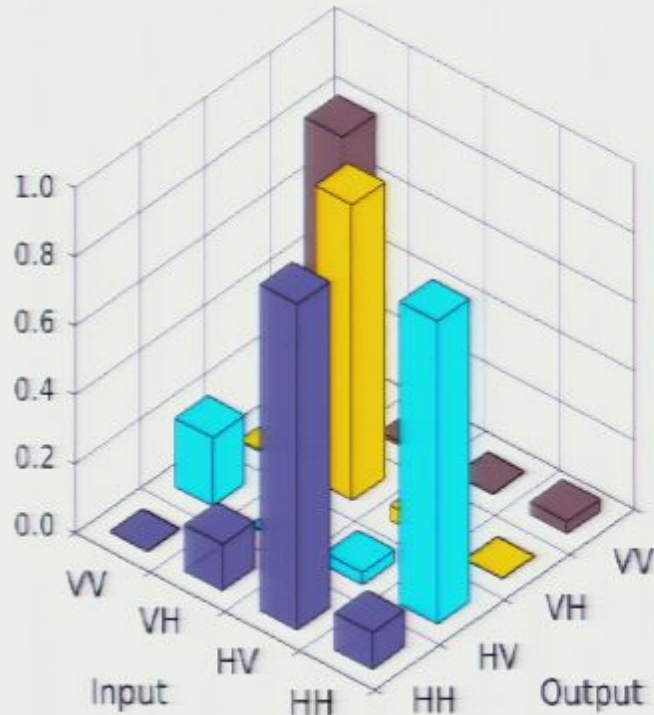
State	F	T	S_L
HH+VV	0.94(2)	0.88(6)	0.07(5)
HH-VV	0.92(3)	0.75(10)	0.17(7)
HV+VH	0.95(3)	0.85(10)	0.10(7)
HV-VH	0.92(3)	0.75(10)	0.17(7)



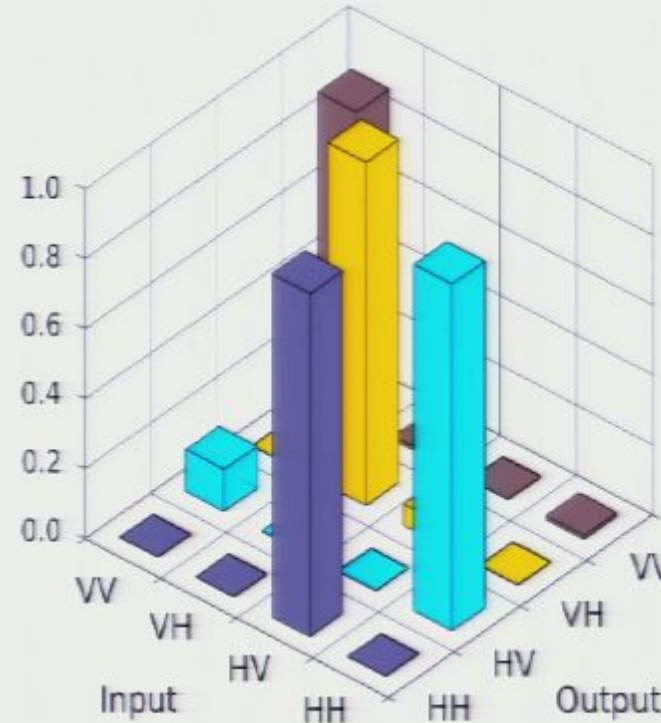
Corrected Experimental Results



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Multi-pairs Subtracted

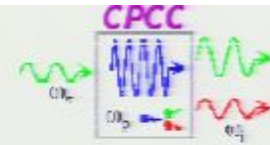


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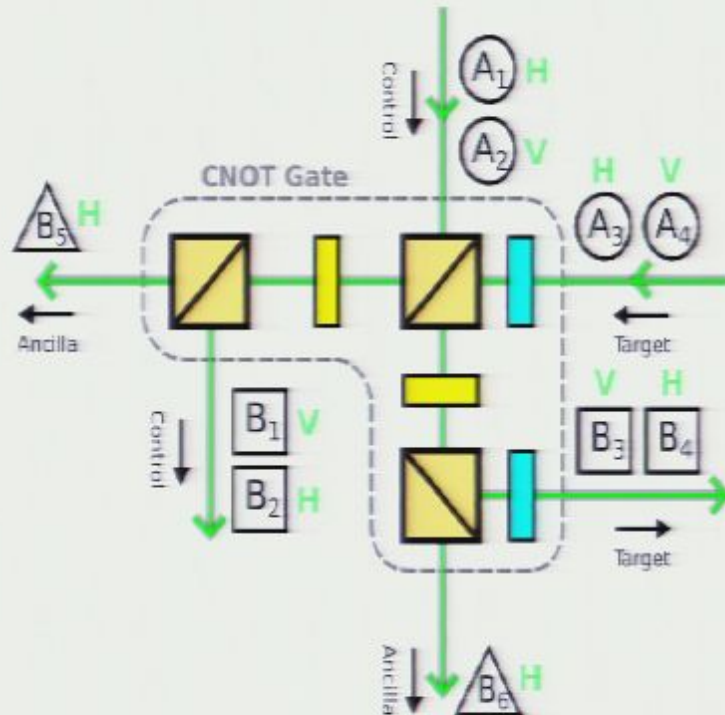
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Is there more to this technique?

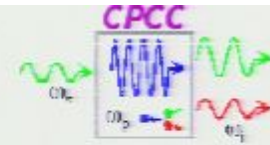


The Assumption:





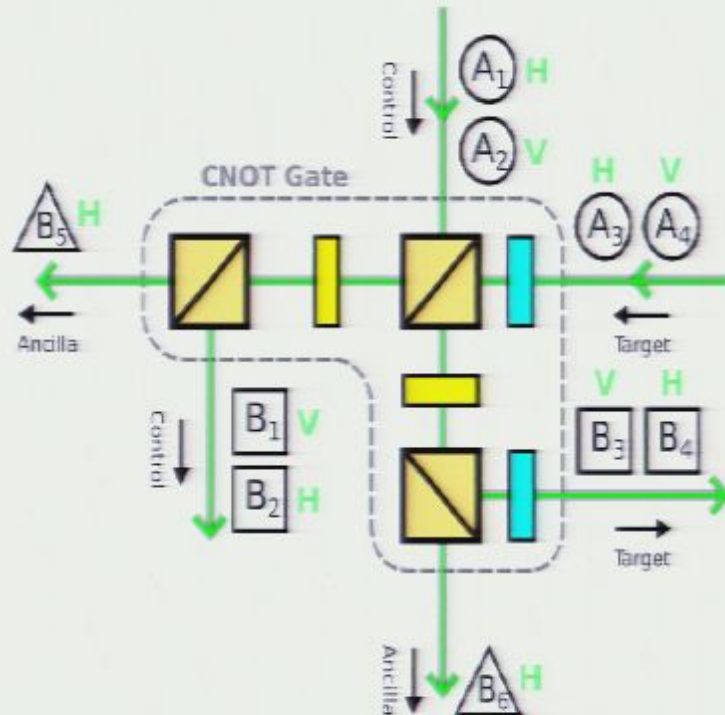
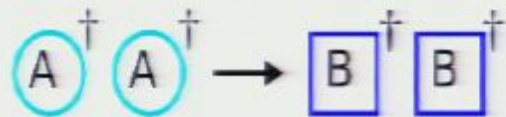
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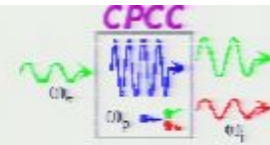


The Advantage:





No Loss, No Decoherence



The output is a coherent superposition of outputs:

$$A_1^\dagger \rightarrow \beta_{11} B_1^\dagger + \beta_{12} B_2^\dagger + \beta_{13} B_3^\dagger + \beta_{14} B_4^\dagger$$

$$A_2^\dagger \rightarrow \beta_{21} B_1^\dagger + \beta_{22} B_2^\dagger + \beta_{23} B_3^\dagger + \beta_{24} B_4^\dagger$$

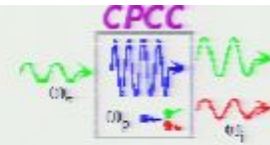
$$A_3^\dagger \rightarrow \beta_{31} B_1^\dagger + \beta_{32} B_2^\dagger + \beta_{33} B_3^\dagger + \beta_{34} B_4^\dagger$$

$$A_4^\dagger \rightarrow \beta_{41} B_1^\dagger + \beta_{42} B_2^\dagger + \beta_{43} B_3^\dagger + \beta_{44} B_4^\dagger$$

In other words, it's a unitary operation.



Modelling Loss

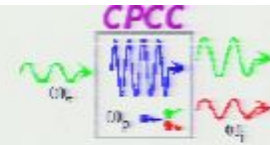


General loss:

$$\textcircled{A}^\dagger \rightarrow \beta_B \boxed{B}^\dagger + \beta_L \triangleleft L^\dagger$$



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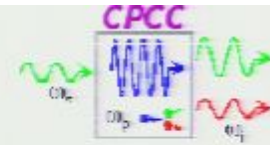
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Modelling Loss

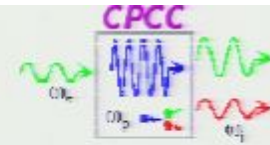


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Modelling Loss



General loss:

$$\textcircled{A}^\dagger \rightarrow \beta_B \boxed{B}^\dagger + \beta_L \triangle L^\dagger$$

Let's make a half of a bigger unitary:

$$\textcircled{A_1}^\dagger \rightarrow \beta_{11} \boxed{B_1}^\dagger + \beta_{12} \boxed{B_2}^\dagger + \beta_{13} \boxed{B_3}^\dagger + \beta_{14} \boxed{B_4}^\dagger + \beta_{15} \triangle B_5^\dagger + \beta_{16} \triangle B_6^\dagger + \beta_{17} \triangle B_7^\dagger + \beta_{18} \triangle B_8^\dagger$$

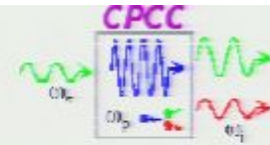
$$\textcircled{A_2}^\dagger \rightarrow \beta_{21} \boxed{B_1}^\dagger + \beta_{22} \boxed{B_2}^\dagger + \beta_{23} \boxed{B_3}^\dagger + \beta_{24} \boxed{B_4}^\dagger + \beta_{25} \triangle B_5^\dagger + \beta_{26} \triangle B_6^\dagger + \beta_{27} \triangle B_7^\dagger + \beta_{28} \triangle B_8^\dagger$$

$$\textcircled{A_3}^\dagger \rightarrow \beta_{31} \boxed{B_1}^\dagger + \beta_{32} \boxed{B_2}^\dagger + \beta_{33} \boxed{B_3}^\dagger + \beta_{34} \boxed{B_4}^\dagger + \beta_{35} \triangle B_5^\dagger + \beta_{36} \triangle B_6^\dagger + \beta_{37} \triangle B_7^\dagger + \beta_{38} \triangle B_8^\dagger$$

$$\textcircled{A_4}^\dagger \rightarrow \beta_{41} \boxed{B_1}^\dagger + \beta_{42} \boxed{B_2}^\dagger + \beta_{43} \boxed{B_3}^\dagger + \beta_{44} \boxed{B_4}^\dagger + \beta_{45} \triangle B_5^\dagger + \beta_{46} \triangle B_6^\dagger + \beta_{47} \triangle B_7^\dagger + \beta_{48} \triangle B_8^\dagger$$



Some Loss, No Decoherence



What about loss here?

$$A_1^\dagger \rightarrow \beta_{11} B_1^\dagger + \beta_{12} B_2^\dagger + \beta_{13} B_3^\dagger + \beta_{14} B_4^\dagger$$

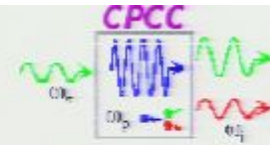
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$$A_4^\dagger \rightarrow \beta_{41} B_1^\dagger + \beta_{42} B_2^\dagger + \beta_{43} B_3^\dagger + \beta_{44} B_4^\dagger$$



Some Loss, No Decoherence



What about loss here?

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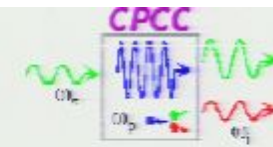
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Explicit loss coherences are needed when they will coherently interfere. In other words, for **low-order coincidences and singles.**



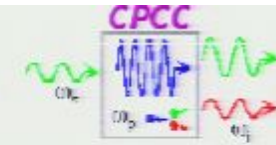
Modelling Decoherence



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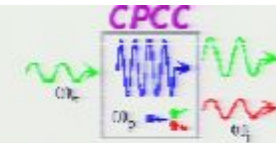
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• In general, there can be **any number of these decohering modes**, leading to any number of matrices.... hmm....

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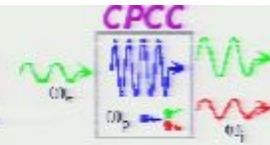
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A Brief Detour: Reviewing the χ -Matrix



$$\mathcal{E}(\rho) = \sum_k E_k \rho E_k^\dagger$$

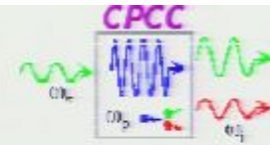
$$E_k = \sum_m e_{km} \tilde{E}_m$$

$$\chi_{mn} = \sum_k e_{km} e_{kn}^*$$

$$\mathcal{E}(\rho) = \sum_{mn} \tilde{E}_m \rho \tilde{E}_n^\dagger \chi_{mn}$$



Looks a lot Like Kraus Operators...



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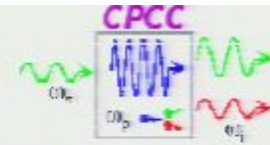
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The Point



Consider an N-qubit gate with 2N single-photon input modes. The 2N-element vector α defines an arbitrary input creation operator:

$$\mathbb{A}^\dagger \equiv \sum_i \alpha_i \mathbb{A}_i^\dagger \quad \rho' \equiv |\alpha\rangle\langle\alpha|$$

For any given LOQC process \mathcal{E} (where primed objects are in the single-photon basis):

$$\mathcal{E}(\rho') = \sum_k \beta_{ij}^{(k)} \rho' \beta_{ij}^{(k)\dagger}$$

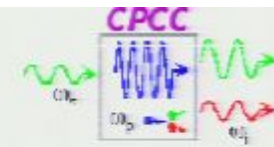
We can rewrite with the Hermitian χ_{mn} :

$$\mathcal{E}(\rho') = \sum_{mn} \tilde{E}'_m \rho' \tilde{E}'_n^\dagger \chi_{mn}$$

By diagonalizing χ_{mn} we recreate the beta map, and can predict the output state for any input.



In Conclusion



Scaling for
N qubit gates

Standard
Process

$$2^{4N}$$

LOQC Process

$$(2N)^4$$

LOQC Process
with Loss

$$(4N)^4$$



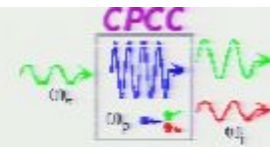
In Conclusion

Scaling looks good, but don't forget feed-forward.

Experimental Axioms:

- Errors are clues.
- Use all your data.

Ideas from smart theorists are always welcome.



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