Title: Towards a generally covariant averaging process for metrics in general relativity

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Abstract: The speculation that Dark Energy can be explained by the backreaction of present inhomogeneities on the evolution of the background cosmology has been increasingly debated in the recent literature. We demonstrate quantitively that the backreaction of linear perturbations on the Friedmann equations is small but is nevertheless non-vanishing. This indicates the need for an improved averaging procedure capable of averaging tensor quantities in a generally covariant way. We present an averaging process which decomposes the metric into Vielbeins selected employing a variational principle, and parallel-transports them to a single point at which they can be averaged. The functionality of the process is discussed in specific 2-d examples, and its application to 3-surfaces and metric recovery in cosmology is outlined.

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# Towards a generally covariant averaging process for metrics in general relativity

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Perimeter Institute, Canada August 14, 2008

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Introduction

- 2 Averaging Problem
- Perturbation Theory
- 4 Generally Covariant Averaging
- Conclusions

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Universe is homogeneous and isotropic on large scales

Introduction

⇒ Use exact solution to Einstein equations (FLRW metric) to model the universe

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Conclusions

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- ⇒ Use exact solution to Einstein equations (FLRW metric) to model the universe
  - CMB is isotropic with only small anisotropies
- ⇒ Describe by linear perturbations about the FLRW solution

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- ⇒ Describe by linear perturbations about the FLRW solution
  - Astronomical observations (galaxy clustering and motions, gravitational lensing, CMB, type la supernovae, Lyman  $\alpha$ , etc.)
- ⇒ ACDM model with 76% dark energy, 20% dark matter, and 4% baryonic matter

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- ⇒ Is dark energy the backreaction of the generation and evolution of inhomogeneities on the evolution of the background Cosmology? Page 7/90

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#### The Averaging Problem

Standard Cosmology based on

$$G_{\mu\nu}(\langle g_{\mu\nu}\rangle) = 8\pi G \langle T_{\mu\nu}\rangle + \Lambda \langle g_{\mu\nu}\rangle$$

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for some average  $\langle A \rangle$  in a domain  $\mathcal D$ 

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→ Modifications can in principle act as a dark enegy

$$G_{\mu\nu}(\langle g_{\mu\nu}\rangle) = 8\pi G \langle T_{\mu\nu}\rangle + 8\pi G T_{\mu\nu}^g + \Lambda \langle g_{\mu\nu}\rangle$$

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#### Einstein Equations in 3+1 Form

Foliate spacetime with family of spacelike hypersurfaces

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## Einstein Equations in 3+1 Form

- Foliate spacetime with family of spacelike hypersurfaces
- Projection operators  $n^{\mu}=rac{1}{lpha}(1,eta^{i})$  and  $h_{\mu\nu}=g_{\mu\nu}+n_{\mu}n_{\nu}$
- Line element  $ds^2 = -(\alpha^2 + \beta_i \beta^i)dt^2 + 2\beta^i dt dx^i + h_{ij} dx^i dx^j$
- Extrinsic curvature  $2K_{ii} = -\mathcal{L}_n h_{ii} = -\dot{h}_{ii}/\alpha$  (for  $\beta^i = 0$ )

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- → Hamiltonian constraint:

$$\mathcal{R} + K^2 - K_j^i K_i^j = 16\pi G\rho + 2\Lambda$$

⇒ Evolution equation:

$$\frac{1}{\alpha}\dot{K}_{ij} =$$

$$\mathcal{R}_{ij} - 2K_i^n K_{nj} + KK_{ij} - 8\pi GS_{ij} + 4\pi Gh_{ij} (S-\rho) - \Lambda h_{ij} - \frac{1}{\alpha} D_i D_j \alpha$$

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- Define average  $\langle A \rangle = \frac{1}{V} \int_{\mathcal{D}} A \sqrt{h} d^3 \mathbf{x}$
- Define Hubble rate  $3H_{\mathcal{D}} = 3\dot{a}_{\mathcal{D}}/a_{\mathcal{D}} = \dot{V}/V$

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⇒ Extrinsic curvature evolution ⇒ Raychaudhuri equation:

$$\frac{\ddot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = -\frac{4\pi G}{3} \left\langle \alpha^{2} (\rho + S) \right\rangle + \frac{\Lambda}{3} \left\langle \alpha^{2} \right\rangle + \frac{1}{3} \left( \mathcal{Q}_{\mathcal{D}} + \mathcal{P}_{\mathcal{D}} \right)$$

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#### Modifications

Kinematical backreaction:

$$Q_{\mathcal{D}} = \left\langle \alpha^2 \left( K^2 - K_j^i K_i^j \right) \right\rangle - \frac{2}{3} \left\langle \alpha K \right\rangle^2$$

Dynamical backreaction:

$$\mathcal{P}_{\mathcal{D}} = \langle \dot{\alpha} K \rangle + \langle \alpha D^i D_i \alpha \rangle$$

Curvature contribution:

$$\mathcal{R}_{\mathcal{D}} = \langle \alpha^2 \mathcal{R} \rangle$$

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#### Perturbation Theory

Take the Newtonian metric in the form

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$$\mathcal{R}_{\mathcal{D}} = \frac{6}{a^2} \langle (\nabla \phi)^2 + 4\phi \nabla^2 \phi \rangle$$

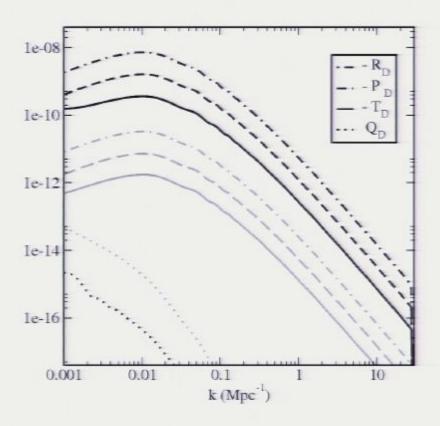
$$\mathcal{Q}_{\mathcal{D}} = 6 \langle \dot{\phi}^2 \rangle$$

$$\mathcal{P}_{\mathcal{D}} = -6 \frac{\dot{a}}{a} \langle \dot{\phi} \phi \rangle - 3 \langle \dot{\phi}^2 \rangle + \frac{2}{a^2} \langle \phi \nabla^2 \phi - (\nabla \phi)^2 \rangle)$$

$$\mathcal{T}_{\mathcal{D}} = \frac{8\pi G}{3} \rho \langle 2\phi \delta + a^2 v^2 \rangle$$

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#### Modifications at z = 10 and z = 0



Evaluate corrections with CMBEasy, e.g.

$$Q_{\mathcal{D}} = 6 \int \mathcal{P}_{\psi}(k) \left| \dot{\phi}(t,k) \right|^2 \frac{dk}{k}$$

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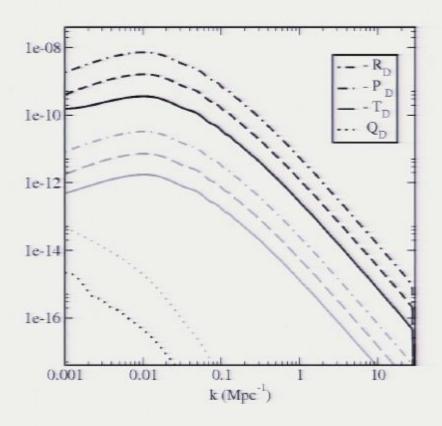
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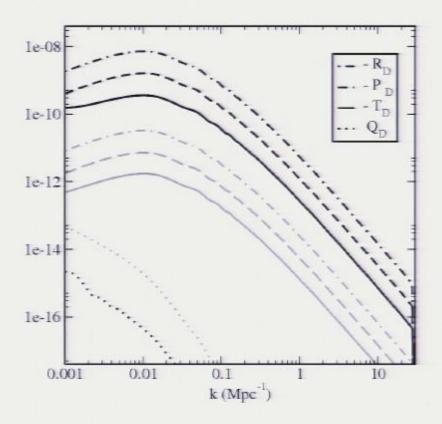
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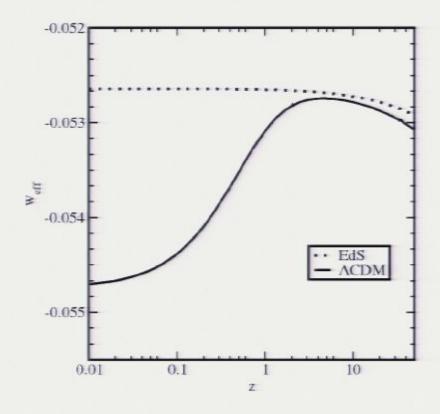
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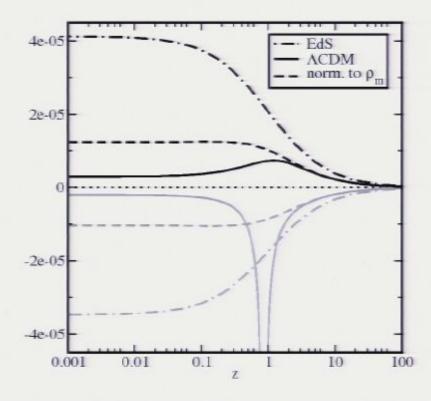
## Effective Equation of State



- Effective density:  $(8\pi G/3)\rho_{\rm eff} = T_{\mathcal D} (\mathcal Q_{\mathcal D} + \mathcal R_{\mathcal D})/6$
- Effective pressure:  $16\pi Gp_{\text{eff}} = \mathcal{R}_{\mathcal{D}}/3 \mathcal{Q}_{\mathcal{D}} 4\mathcal{P}_{\mathcal{D}}/3$
- ⇒ Effective equation of state:

 $w_{ ext{eff}} = -(1/3)(\mathcal{R}_{\mathcal{D}} - 4\mathcal{P}_{\mathcal{D}} - 3\mathcal{Q}_{\mathcal{D}})/(\mathcal{R}_{\mathcal{D}} - 6\mathcal{T}_{\mathcal{D}} + \mathcal{Q}_{\mathcal{D}}) \approx -1_{\text{Page-Loop}}$ 

## Impact on Large-Scale Evolution



- Einstein de-Sitter and ΛCDM (WMAPIII concordance)
- $\bullet \sim 10^{-5}$  impact as predicted, maxima at  $z \approx 1.4$  and  $z \approx 0.7$

⇒ Backreaction is a small but non-vanishing physical effect

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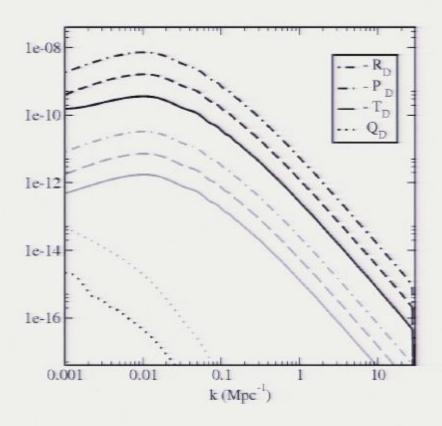
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- ⇒ Backreaction is a small but non-vanishing physical effect
  - Newtonian metric cannot be used
  - Backreaction is a non-linear effect
  - Averaging process

$$\langle A \rangle = \frac{1}{V} \int_{\mathcal{D}} A \sqrt{h} d^3 \mathbf{x}$$

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Perturbation Theory

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- depends on choice of coordinate system
- cannot be used to average vector and tensor quantities

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- Background free approach
- ⇒ Need a generally covariant averaging process

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Conclusions

Introduction

## Generally Covariant Averaging Process for the Metric

- Averaging Process must be independent of coordinate system
- ⇒ Parallel transport tensor quantities along geodesics to the same point before averaging

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# Generally Covariant Averaging Process for the Metric

- Averaging Process must be independent of coordinate system
- → Parallel transport tensor quantities along geodesics to the same point before averaging
  - Decompose metric into a right-handed orthochronous Minkowski tetrad

$$g_{\mu\nu}(x) = \eta_{\alpha\beta} E^{\alpha}{}_{\mu}(x) E^{\beta}{}_{\nu}(x)$$

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# Generally Covariant Averaging Process for the Metric

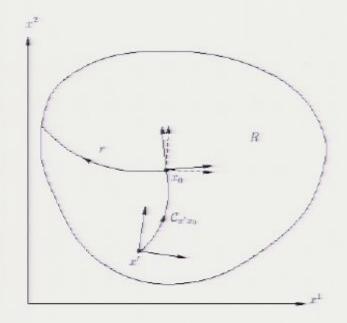
- Averaging Process must be independent of coordinate system
- ⇒ Parallel transport tensor quantities along geodesics to the same point before averaging
  - Decompose metric into a right-handed orthochronous Minkowski tetrad

$$g_{\mu\nu}(x) = \eta_{\alpha\beta} E^{\alpha}{}_{\mu}(x) E^{\beta}{}_{\nu}(x)$$

 Find (up to global Lorentz-transformations) unique tetrad field, the maximally smooth tetrad field, by following Lagrangian

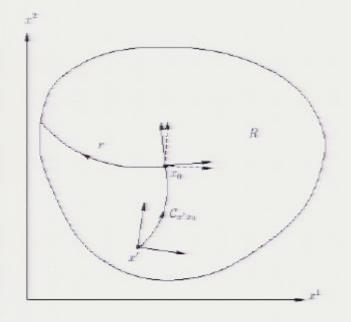
$$\mathcal{L}_{\mathrm{MS}} = (D_{\mu} E^{\alpha}{}_{\rho}) (D_{\nu} E^{\beta}{}_{\lambda}) g^{\mu\nu} g^{\rho\lambda} \eta_{\alpha\beta}$$

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$$V(x', x_0; \mathcal{C}_{x_0x'}) = \mathcal{P} \exp\left[-\int_{\mathcal{C}_{x_0x'}} dz^{\mu} \Gamma_{\mu}(z)\right]$$

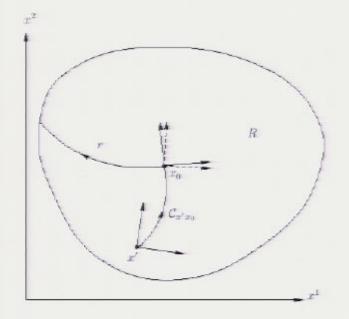
where  $\Gamma_{\mu}(x)$  are four matrices with components  $(\Gamma_{\mu}(x))^{\lambda}_{\nu} = \Gamma^{\lambda}_{\mu\nu}(x)$ 



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$$\begin{split} \langle E^{\alpha}{}_{\mu}(x_0) \rangle \\ &= \int_{R} f(x_0, x'; \mathcal{C}_{x'x_0}) \widehat{V}_{\mu}{}^{\nu}(x_0, x'; \mathcal{C}_{x'x_0}) E^{\alpha}{}_{\nu}(x') \sqrt{-g(x')} \ d^4x' \end{split}$$



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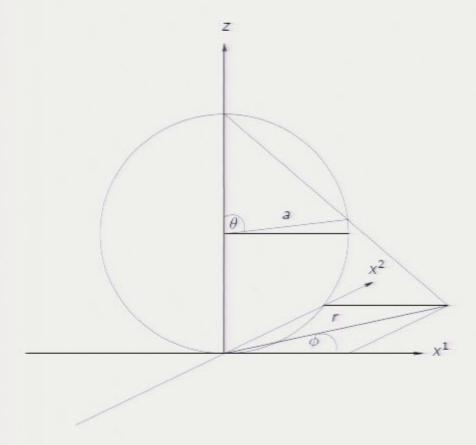
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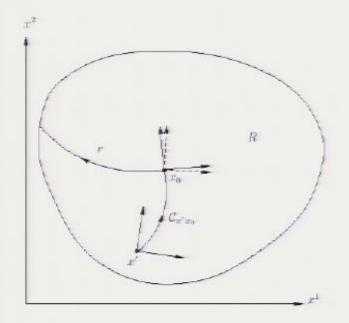
⇒ Averaged metric:

$$\langle g_{\mu\nu}(x)\rangle = \eta_{\alpha\beta} \langle E^{\alpha}{}_{\mu}(x)\rangle \langle E^{\beta}{}_{\nu}(x)\rangle$$

Stereographic Projection:



• Metric:  $g_{ij} = (\frac{2a}{L})^4 \delta_{ij}$  where  $L^2 = 4a^2 + (x^1)^2 + (x^2)^2$ 



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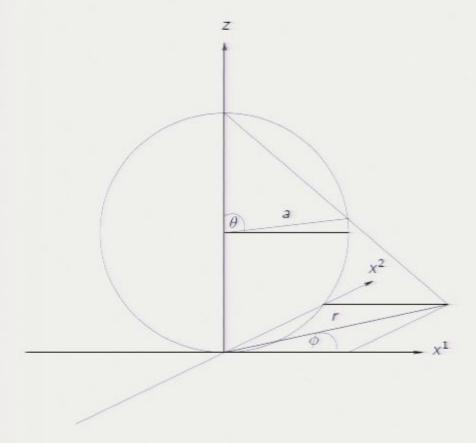
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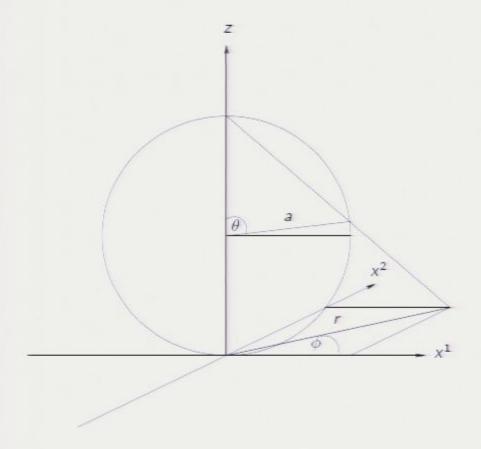
#### Stereographic Projection:



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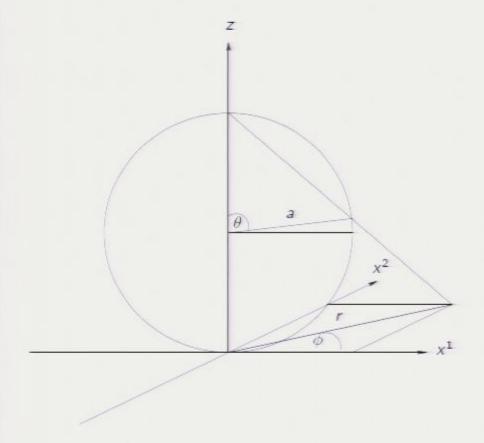
Averaging Problem



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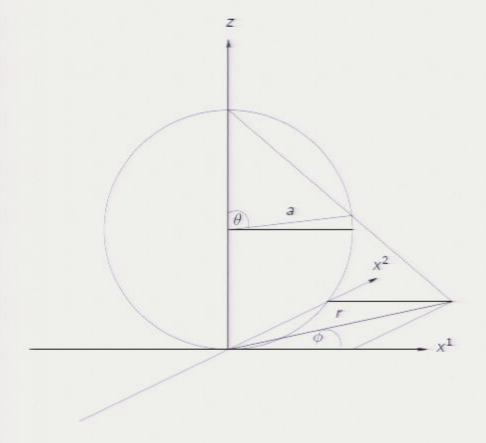
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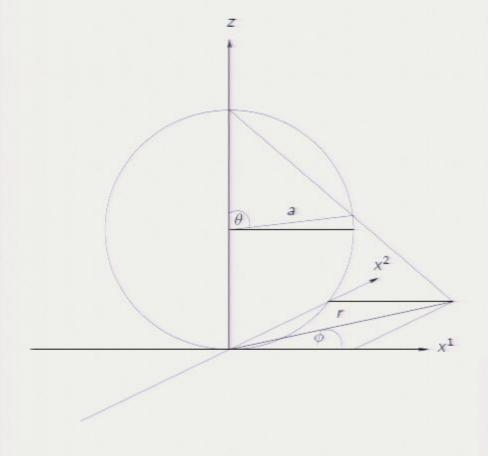
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- Averaged metric:

$$\langle g_{ij}(x)\rangle = g_{ij}(x)$$

Generally Covariant Averaging

## Averaging the Metric of a Perturbed Two-Sphere

- Perturb spherical coordinates with function f(x, y, z)
- Stereographic projection leads to perturbed metric:

$$(g_P)_{ij} = (1 + 2\eta f)(\frac{2a}{L})^4 \delta_{ij}$$

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  - Reference dyad field  $\widetilde{E}^{a}_{i}$  with  $\widetilde{E}^{a}_{i}\widetilde{E}^{b}_{j}$   $\delta_{ab}=g_{ij}$
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Averaging Problem

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- Solve  $\delta S = 0$  with  $S = \int_{\mathcal{P}} d^2x \sqrt{g} (D_i E^a_i) (D_k E^b_i) g^{ik} g^{jl} \delta_{ab}$
- Introduce vector field  $u^k = (D_i \widetilde{E}^c{}_i) \widetilde{E}^d{}_l \epsilon_{cd} g^{ik} g^{jl}$

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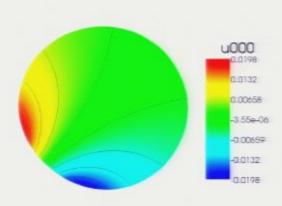
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$$g^{ik}(\partial_i \partial_k \phi) = -\frac{1}{2} D_k u^k \text{ on } R$$
$$\frac{\partial \phi}{\partial n} = -\frac{1}{2} n_k u^k \text{ on } \partial R$$

 $\Rightarrow \frac{\partial \phi}{\partial n} = -\frac{1}{2} n_k u^k \text{ on } \partial R$ 

# The Gaussian Shaped Perturbation

Gascoigne3D



$$\left(\frac{\partial^2}{(\partial x^1)^2} + \frac{\partial^2}{(\partial x^2)^2}\right)\phi(x^1, x^2) = 0$$
 on  $R$ 

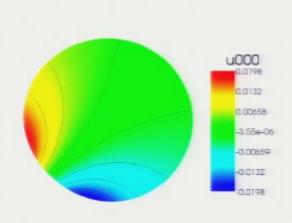
Neumann boundary conditions on  $\partial R$ 

$$\frac{\partial \phi}{\partial n} = \eta \cos^{-2}(\frac{r}{2a})(h(r,\gamma) + \frac{1}{2a^2}\partial_{\gamma} \int_{0}^{r} f(s',\gamma)ds')$$

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#### Gascoigne3D

ntroduction



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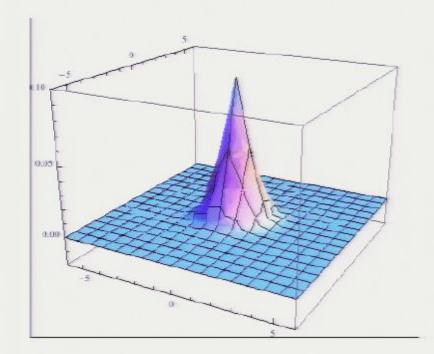
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R is the area inside  $\partial R$  given by

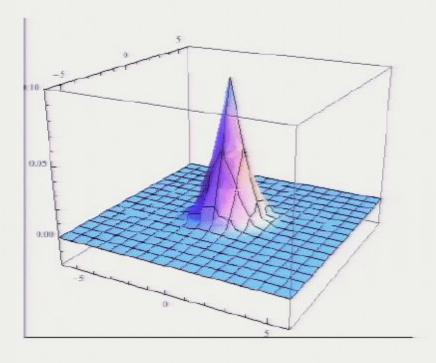
$$\alpha^{1}(\gamma) = 2a \tan(\frac{r}{2a}) \cos \gamma + \eta \frac{v(r,\gamma)}{\cos^{2}(\frac{r}{2a})} \sin \gamma$$
$$-\eta \frac{\cos \gamma}{\cos^{2}(\frac{r}{2a})} \int_{0}^{r} f(s',\gamma) ds'$$
$$\alpha^{2}(\gamma) = 2a \tan(\frac{r}{2a}) \sin \gamma - \eta \frac{v(r,\gamma)}{2^{2}(r,\gamma)} \cos \gamma$$

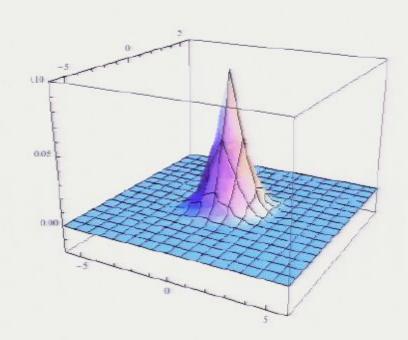
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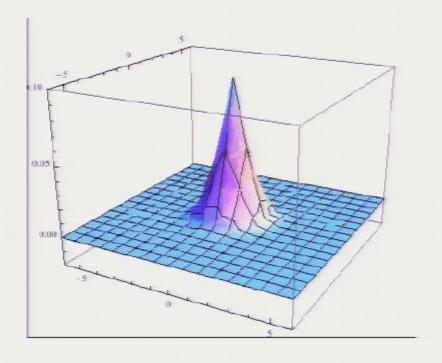
where 
$$h(\tau) = \frac{\partial f}{\partial x^2}\Big|_{(x^1, x^2) = (z^1(\tau), z^2(\tau))} \frac{dz^1}{d\tau}(0) - \frac{\partial f}{\partial x^1}\Big|_{(x^1, x^2) = (z^1(\tau), z^2(\tau))} \frac{dz^2}{d\tau}(0)$$
 and  $v(\tau, \gamma)$  fulfills the differential equation 
$$\frac{d^2v(\tau, \gamma)}{d\tau^2} + \frac{v(\tau, \gamma)}{z^2} = \frac{h(\tau, \gamma)}{\cos^2(\frac{\tau}{2s})}$$

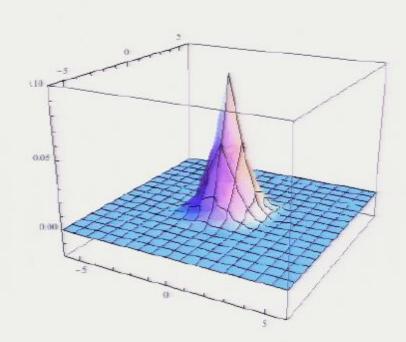


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Averaging effect in investigated example is too small

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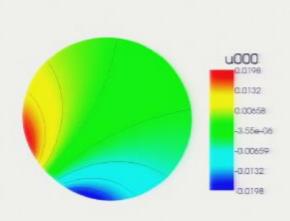
$$\mathcal{L} = (L_t E^a{}_i) \left( L_t E^b{}_j \right) \delta_{ab} t^i t^j$$

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# The Gaussian Shaped Perturbation

#### Gascoigne3D

ntroduction



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 on  $R$ 

Neumann boundary conditions on  $\partial R$ 

$$\frac{\partial \phi}{\partial n} = \eta \cos^{-2}(\frac{r}{2a})(h(r,\gamma) + \frac{1}{2a^2}\partial_{\gamma} \int_{0}^{r} f(s',\gamma)ds')$$

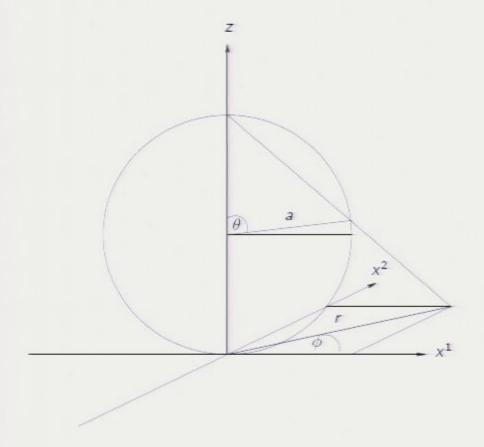
R is the area inside  $\partial R$  given by

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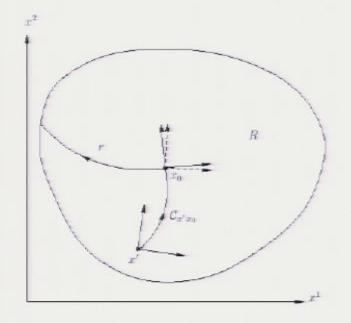
# Averaging the Metric of a Two-Sphere

#### Stereographic Projection:



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- Connector:

$$\widehat{V}_{j}^{i}(0,\tau;\mathcal{C}_{\tau 0}) = \cos^{-2}(\frac{\tau}{2a})\delta_{ij}$$



Parallel transport along geodesics  $C_{x_0x'}$  realized by Wegner-Wilson line operator

$$V(x', x_0; \mathcal{C}_{x_0x'}) = \mathcal{P} \exp \left[ - \int_{\mathcal{C}_{x_0x'}} dz^{\mu} \; \Gamma_{\mu}(z) \right]$$

where  $\Gamma_{\mu}(x)$  are four matrices with components  $(\Gamma_{\mu}(x))^{\lambda}_{\nu} = \Gamma^{\lambda}_{\mu\nu}(x)$ 

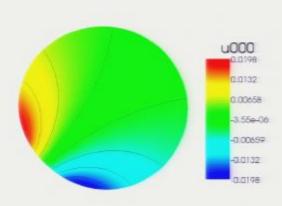
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⇒ Averaged metric:

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⇒ Use different Lagrangian to define initial tetrad field:

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• 
$$\mathcal{L} = (L_t E^a_i) (L_t E^b_j) \delta_{ab} (t^i t^j + g^{ij} Rs^2)$$

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Perturbation Theory

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• 
$$\mathcal{L} = (L_{\xi}E^{a}_{i})(L_{\xi}E^{b}_{j})\delta_{ab}\xi^{i}\xi^{j}$$

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ntroduction

 Problem: Used Lagrangian is similar to the Lagrangian which defines the geodetic induced parallel field

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$$\mathcal{L} = (L_{\xi}E^{a}_{i})(L_{\xi}E^{b}_{j})\delta_{ab}\xi^{i}\xi^{j}$$

o ...

⇒ They all fail to define a dyad field that can be used to average the perturbed plane in the desired way

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### Conclusions and Outlook

 Once we have a suitable Lagrangian we have a generally covariant averaging process which can be used to smooth metrics in the framework of GR

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$$\mathcal{L} = (L_t E^a{}_i) \left( L_t E^b{}_j \right) \delta_{ab} t^i t^j$$

- ⇒ Use different Lagrangian to define initial tetrad field:
  - $\mathcal{L} = (L_t E^a_i) (L_t E^b_j) \delta_{ab} (t^i t^j + g^{ij} Rs^2)$
  - $\mathcal{L} = (L_n E^a_i) (L_n E^b_j) \delta_{ab} n^i n^j$
  - $\mathcal{L} = (L_{\xi}E^{a}_{i})(L_{\xi}E^{b}_{j})\delta_{ab}\xi^{i}\xi^{j}$

•

ntroduction

 Problem: Used Lagrangian is similar to the Lagrangian which defines the geodetic induced parallel field

$$\mathcal{L} = (L_t E^a{}_i) \left( L_t E^b{}_j \right) \delta_{ab} t^i t^j$$

⇒ Use different Lagrangian to define initial tetrad field:

• 
$$\mathcal{L} = (L_t E^a_i) (L_t E^b_j) \delta_{ab} (t^i t^j + g^{ij} Rs^2)$$

• 
$$\mathcal{L} = (L_n E^a_i) (L_n E^b_j) \delta_{ab} n^i n^j$$

• 
$$\mathcal{L} = (L_{\xi}E^{a}_{i})(L_{\xi}E^{b}_{j})\delta_{ab}\xi^{i}\xi^{j}$$

. . . .

⇒ They all fail to define a dyad field that can be used to average the perturbed plane in the desired way

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### Conclusions and Outlook

- Once we have a suitable Lagrangian we have a generally covariant averaging process which can be used to smooth metrics in the framework of GR
- Apply it to different perturbation functions to study their interaction with each other and with the background sphere
- Apply it to three-sphere and three-plane corresponding to hypersurfaces of closed and flat FLRW models

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- Apply it to four-dimensional example which involves choice of boundary conditions on the congruence of light-like geodesics
- Apply averaging process to Cosmology and combine the two lines of research

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