Title: Observational Constraints on Gravitational Degrees of Freedom

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Abstract: TBA

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The Phenomenology of Gravity



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12 August 2008

A Case For / Introduction To "Phenomenology" (with examples)











A Case For / Introduction To "Phenomenology" (with examples)

• Parsimony (gratia Alan Cooney, U. Arizona)











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- Parsimony (gratia Alan Cooney, U. Arizona)
- Consistency (gratia Chanda Prescod-Weinstein, P.I.)











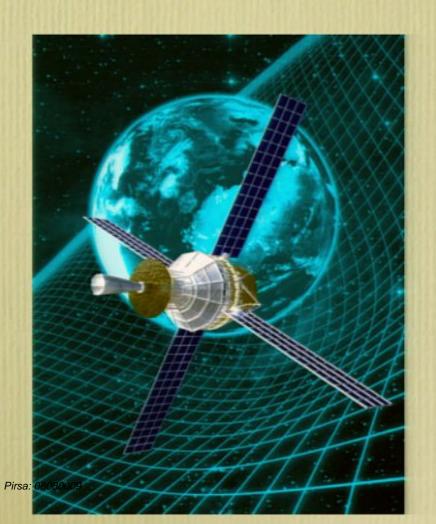
Example I.

Degrees of Freedom in Gravity

(work in collaboration with Alan Cooney & Dimitrios Psaltis @ U. Arizona)

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"Parametrized Post-Newtonian"



Pioneered by Will and Nordvedt in the 1970s a general "phenomenology" of solar system gravity

more than just "interoperability"

— a deep assertion of physical
principles

Newtonian gravity

$$g_{00} = -1 + 2U$$

$$U = \int \frac{\rho(x')}{|x - x'|} d^3x$$

"First order"
post-Newtonian
(in a particular
gauge)

$$g_{00} = -1 + 2U + 2U^2$$

possible modifications . . .

$$g_{00} = -1 + 2U + 2\beta U^2 - 2\xi \Phi_W \cdots$$

$$\Phi_W = \int \frac{\rho' \rho''(x - x')}{|x - x'|^3} \cdot \left(\frac{x' - x''}{|x - x''|} - \frac{x - x''}{|x' - x''|} \right) d^3 x' d^3 x''$$

"Matter Dictatorship"

All possible modifications to the metric are functions of the matter/energy distribution and its velocity [& preferred frames]

$$\Phi_W = \int \frac{\rho' \rho''(x - x')}{|x - x'|^3} \cdot \left(\frac{x' - x''}{|x - x''|} - \frac{x - x''}{|x' - x''|}\right) d^3x' d^3x''$$



$$A = \int \frac{\rho'[v' \cdot (x - x')]^2}{|x - x'|^3} d^3x'$$

&c., &c.

not quite "no new degrees of freedom" — e.g., Brans-Dickie theories can fit in, but not f(R) (in general) Page 9/58

PPN as Phenomenology

plus sides

$$\mathcal{A} = \int \frac{\rho'[v'\cdot(x-x')]^2}{|x-x'|^3} d^3x'$$

"matter dictatorship": parsimony without restriction forms of dictatorship easy to cover (but see Alexander, 2008)

downsides (for cosmology &c.)

requires "post-Newtonian" conditions: shallow potentials that fade at large distance (but can adapt the gauge to cosmology)

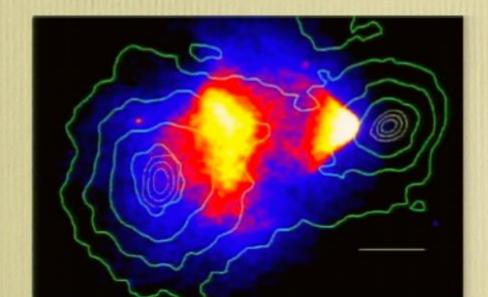
MOND vs. Dark Matter

MOND has parsimony — one parameter "acceleration scale," but not relativistic

TeVeS has phenomenological parsimony — matter dictatorship[*], but not fundamentally so (new vector and scalar fields)

[*]. caveat. — only sort of.

observation ruled out (?) both — Clowe et al. 2006



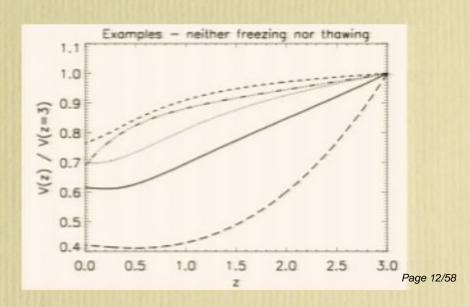
The Age of Quintessence

 "in for a penny, in for a pound" — a newly cavalier attitude towards new degrees of freedom (but a time of great hope.)

• pick your $V(\varphi)$, find your E.O.M., discover great

fundamental physics!

· where we are today:



$$\int \sqrt{-g} [R + (\partial \phi)^2 + V(\phi) + \mathcal{L}_{\text{matt}}]$$

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$$\int \sqrt{-g}[R + (\partial \phi)^2 + V(\phi) + \mathcal{L}_{\text{matt}}] - \int \sqrt{-g}[R + f(R) + \mathcal{L}_{\text{matt}}]$$

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but! — as an "exact" theory, can be rewritten:

$$\int \sqrt{-g} [R + (\partial \phi)^2 + \tilde{V}(\phi) + \mathcal{L}_{\text{matt}}(\phi)]$$

f(R) is the new $V(\varphi)$?

 Certainly new behaviors, with various clustering scales — looks in many ways like k-essence, which has non-trivial sound speeds. (See, e.g., work by Sawicki & Hu.)

• One of the most interesting things in "classical

cosmology" today

• But not conceptually distinct.

Still new degrees of freedom!



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One "escape route"

Forget about it! — just parametrize w, the E.O.S.: pick one you like, e.g.,

$$w(a) = w_0 + (1-a)w_a$$

$$w(z) = w_0 + zw_1$$

nice for doing observational tests given the current data—but "so open-minded your brain falls out"? You are never sure what physics you are holding constant.

 $w_a \leq 0.1$ means what, exactly? is it inconsistent with structure formation (e.g.)?

Pirsa: 08080009 you'll never know without a "real" model behind Page 19/58

Should we bite the bullet?

- Classical, local f(R) theories that are IR complete generically have new degrees of freedom, and are massless in cosmologically relevant regimes.
- Can we "do more with less"? Do the observations require new degrees? Is there a dark energy bullet cluster? [Not yet!]

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- Stay with Lagrangian formalism retain, for free, general covariance (so no nasty MOND ⇒ TeVeS surprises.)
- But assume it is a perturbative approximation to a larger theory "we know not what" that fixes "we know not how" fictional extra degrees that appear perturbatively.

$$\mathcal{L} = \int \frac{1}{2}\dot{q}^2 + \alpha q(x + \delta x)q(x + \delta x')d^3x \qquad \begin{array}{l} \textit{second order UV non-local} \\ \\ = \int \frac{1}{2}\dot{q}^2 + \alpha \left[q + \frac{dq}{dx}\delta x + \frac{1}{2}\frac{d^2q}{dx^2}(\delta x)^2 + \cdots \right] \left[q + \frac{dq}{dx}\delta x' + \frac{1}{2}\frac{d^2q}{dx^2}(\delta x')^2 + \cdots \right] \\ \\ \overset{\text{Pirsa: 08080009}}{\approx} \int \frac{1}{2}\dot{q}^2 + \alpha \left[q^2 + q\frac{dq}{dx}(\delta x + \delta x' + \left(\frac{dq}{dx}\right)^2 \delta x \delta x' + \cdots \right] \qquad \begin{array}{l} \text{fictitious bigher order local} \\ \\ & \text{fictitious bigher order local} \end{array}$$

The Formalism

 Adapt the "perturbative localization" of Woodard, Simon et al. — developed for the classic case of (UV) non-local theories.

deceptive "formalism overlap"

$$\int \sqrt{-g} [R + \frac{\mu^4}{R} + \cdots]$$

looks like a "standard" f(R) theory — but appearances are deceptive; here we write it in terms of an action, but only because it keeps us honest in spacetime symmetries

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sa: 08080009
$$\int 1_{-2} + \left[q + \frac{dq}{dx}\delta x + \frac{1}{2}\frac{d^{2}q}{dx^{2}}(\delta x) + \frac{1}{2}\frac{d^{2}q}{dx^{2}}(\delta x')^{2} + \cdots\right]$$
Page 25/58

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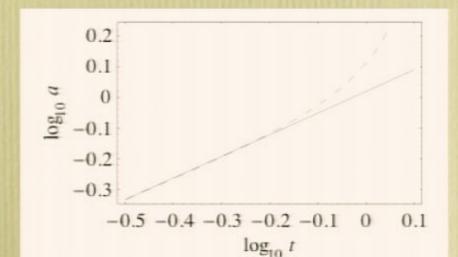
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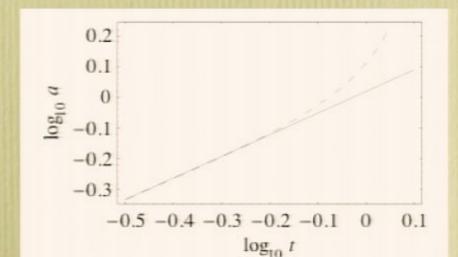
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Perturbative Validity?

control parameter for knowing when perturbative approximations break down : $R^2 \gg \mu^4$, or

$$H^4 \gg \frac{\mu^4}{36(1-q)^2}$$

i.e. — our solutions are valid even for "large" accelerations; (comes from the fact that H is not a scalar but rather gauge-dependent — Minkowski space can have large H!)

This Kind of Modified Gravity...

- has "interesting" behavior (e.g., can produce reliable cosmic acceleration.)
- may imply a "bizarre" fundamental theory but will respect many important *empirical* features such as general covariance, gauge independence, &c.
- is not as wildly unconstrained as you might think . . .

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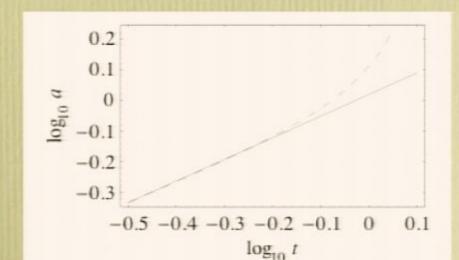
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Constraining Modified Gravity simpliciter

$$\int \sqrt{-g}[R+f(R,G_B)+\cdots]$$

(GB is the "Gauss-Bonnet" term — integral related to spacetime topology)

if f is an polynomial without a constant term,

$$w \leq -2$$

...ruled out by WMAP + Supernovae!

Taking the theory further

 homogenous expansion already a way to rule out entire classes of modified gravities quickly and powerfully.

$$\int \sqrt{-g} [R + f(R, R_{\mu\nu}R^{\mu\nu}, R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}) + \cdots]$$

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 produces reliable statements about, e.g., growth
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The Future of Modified Gravity

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 the simplicity of the mathematics means that large classes can be ruled out quickly; homogenous expansion + structure formation constraints may well rule out such a massive class that we can answer (to a string theorist's delight)

Pirsa: 08080009 are new gravitational degrees of freedom necessary Page 50/58

 what fundamental theories lie beneath our perturbative approximations? Non-local theories are one obvious place to look, but the relationship is non-trivial because of our UV/IR exchange.

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 Niayesh & Ghazal's work on Cuscutons & Gravitational Aethers — are their d.o.f.'s "frozen in the Cooney-DeDeo-Psaltis way"?

logically (but not physically?) independent ways to "freeze" apparent degrees of freedom:



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- homogenous solution dies away
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Pirsa: 08080009 are new gravitational degrees of freedom necessary Page 53/58

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Example II.

The Case of the Covariant Aether

(work in collaboration with Chanda Prescod-Weinstein @ P.I. & U. Waterloo)

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