



Title: Entanglement and Secret-Key Distillation from a Complementary Information Tradeoff

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Abstract: One of the quintessential features of quantum information is its exclusivity, the inability of strong quantum correlations to be shared by many physical systems. Likewise, complementarity has a similar status in quantum mechanics as the sine qua non of quantum phenomena. We show that this is no coincidence, and that the central role of exclusivity in quantum information theory stems from the phenomenon of complementarity. We adopt an information-theoretic approach to complementarity, which leads to a new and simple definition of private states and new proofs of the achievable asymptotic rates of both secret key and entanglement distillation. From the latter follows a new proof of the direct part of the quantum noisy channel coding theorem.

Quantum Info as Complementary Classical Info:  
Secret Key and Entanglement Distillation via Processing  
Complementary Information

Joseph M. Renes  and Jean-Christian Boileau 

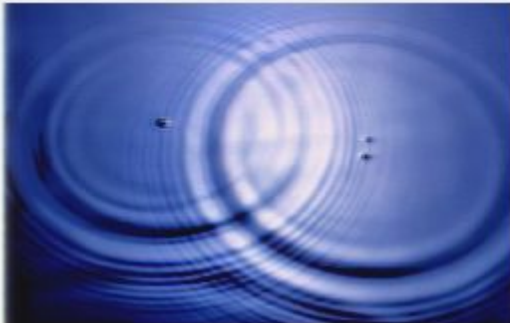


Theoretical Quantum Physics, Institut für Angewandte Physik  
Technische Universität Darmstadt

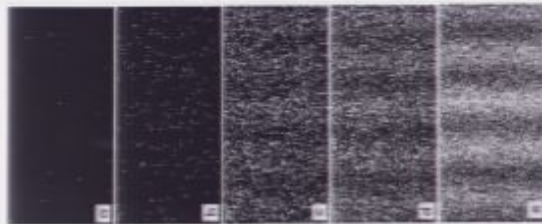


Center for Quantum Information and Quantum Control  
University of Toronto

## Complementarity is the essence of quantum mechanics...

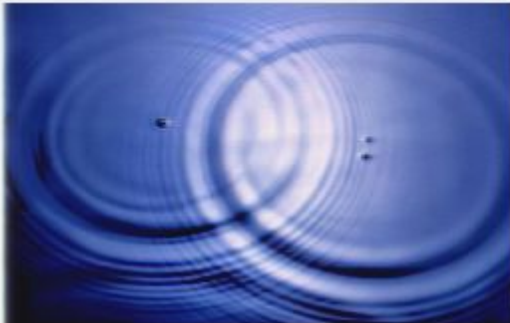


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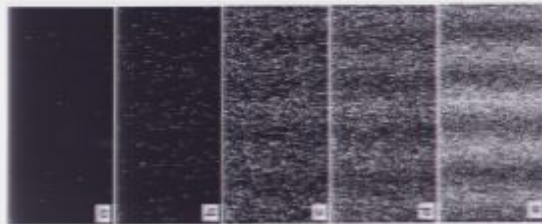


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But does it really *explain* anything?

YES! We can directly and concretely understand several tasks in quantum information theory in terms of complementarity.

## Outline

- 1 Define what we mean by complementarity. Use an information-theoretic approach. (Surprise!)
- 2 Consider two common tasks in qinfo:  
Distillation of secret keys or entangled pairs.  
Recall known results and examine the logic behind them.
- 3 Construct protocols by using the complementarity approach.

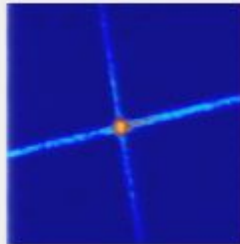
## Information-Theoretic Complementarity

## Entropic Uncertainty Relation

Maassen & Uffink  
PRL 60 1103 (1988)

$$H(O^A) + H(\tilde{O}^A) \geq -\log \max_{j\tilde{k}} |\langle j|\tilde{k}\rangle|^2$$

$O^A, \tilde{O}^A$  operators on system  $A$ ; eigenvectors  $|j\rangle$  and  $|\tilde{k}\rangle$ ;  $H(\cdot)$  entropy.



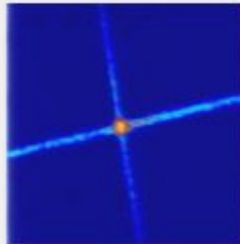
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- But how much info can we *simultaneously* extract?

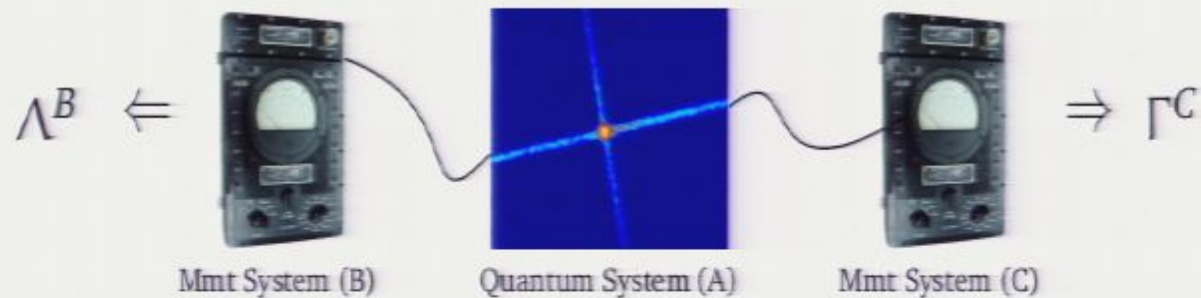


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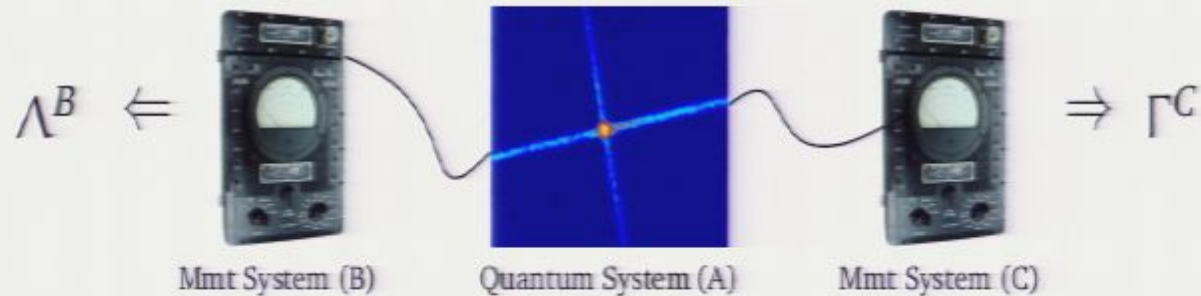
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## Information Exclusion Principle

Hall  
PRL 74 3307 (1995)

$$H(O^A|\Lambda^B) + H(\tilde{O}^A|\Gamma^C) \geq -\log \max_{j,k} |\langle j|\tilde{k}\rangle|^2$$

$\Lambda^B, \Gamma^C$  measurements on B, C;  $H(\cdot|\cdot)$  conditional entropy

- ☞ tradeoff in simultaneously available, complementary classical information

## Complementary Information Tradeoff

JMR & JCB  
arxiv:0806.3984 [quant-ph]

$$H(O^A|B) + H(\tilde{O}^A|C) \geq -\log \max_{j\tilde{k}} |\langle j|\tilde{k}\rangle|^2$$

- ☞ Replace classical mutual information with Holevo upper bound.  
 $H(O^A|B)$  is quantum conditional entropy after Alice measures  $O^A$
- ☞ Numerical support for conjecture up to dimension 12.
- ☞ Holds for generalized Paulis  $X^A$  and  $Z^A$ .

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Distillation of Secret Keys or Entangled Pairs

(And state merging, too.)

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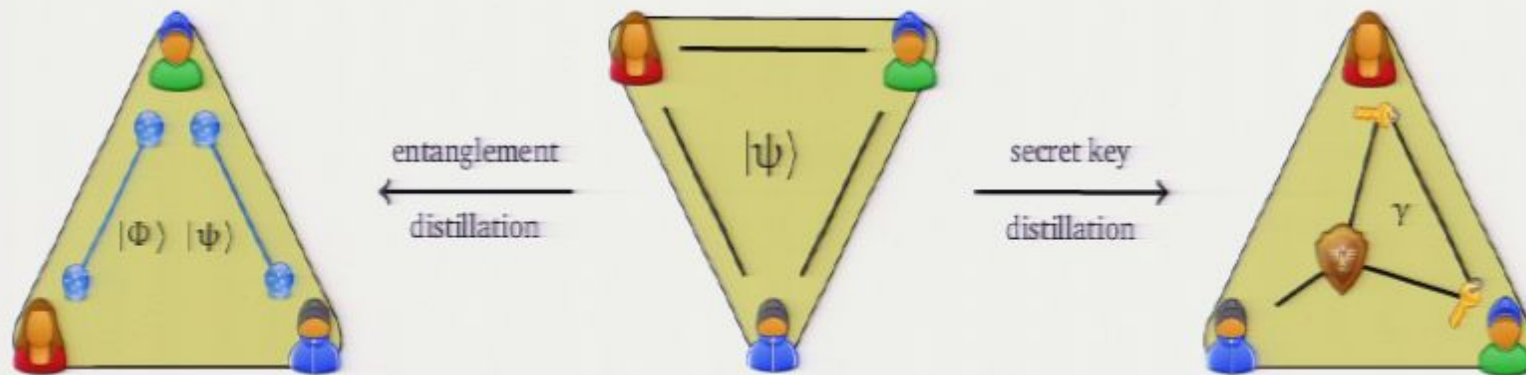
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- Initially Alice, Bob, and Eve share (many copies of) a pure state  $|\psi\rangle$ .
- In *entanglement distillation* Alice and Bob create entangled states  $|\Phi\rangle$  using only one-way classical communication.
- Related is the *state merging* protocol in which Alice transfers her part of the purification of Eve's system to Bob. Both can be performed simultaneously.
- In *key distillation* Alice and Bob use one-way classical communication to create a shared private random classical string. In QM this is described by a *private state*  $\gamma$ , comprising the two (identical) key strings and a *shield* which deflects Eve's correlations from it.



- ☞ Key distillation rate Devetak & Winter  
ProcRoySocA 461 207 (2005)

$$K_{\rightarrow}(\psi) \geq I(Z^A : B) - I(Z^A : E) = H(Z^A|E) - H(Z^A|B)$$

- ☞ Entanglement distillation rate

$$E_{\rightarrow}(\psi) \geq H(B) - H(AB) = -H(A|B).$$

- ☞ Communication cost  $I(A : E)$  in either case.

- ☞ State merging at the same rates/costs Horodecki, Oppenheim, Winter  
Nature 436 7051 (2005)

- ☞ Channel coding from entanglement distillation via teleportation



Decouple  
Alice & Eve



☞ Sensible approach for key generation...

☞ Use Uhlmann's Theorem for entanglement generation: Schumacher & Westmoreland QIP 15 (2002).

$$\psi^A \simeq 1^A/d^A \quad \text{and} \quad \psi^{AE} \simeq \psi^A \otimes \psi^E$$

$$\Downarrow \quad \exists U^B \quad \text{s.t.}$$

$$U^B |\psi\rangle^{ABE} \simeq \sum_{jk} \sqrt{q_k/d^A} |j\rangle^A |j, k\rangle^B |k\rangle^E$$

- $d^A$  is the dimension of  $A$
- eigenvalues/vectors of  $\psi^A$  ( $\psi^E$ ) are  $1/d^A$ ,  $|j\rangle^A$  ( $q_k, |k\rangle^E$ ),
- $|j, k\rangle^B$  are orthonormal.

☞ By coherently measuring  $B$ ,  $\psi^{AB}$  can be (nearly) transformed into  $\Phi^{AB}$ .

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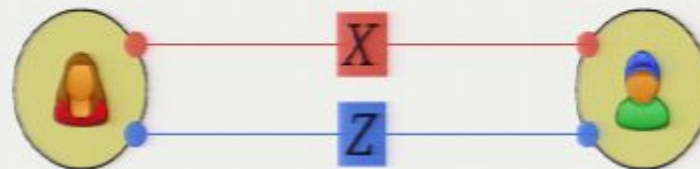
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## Complementarity-based Approach

Couple Alice and Bob classically, twice.



- Quantum correlations between Alice and Bob are equivalent to two complementary classical correlations, say  $X$  and  $Z$  correlations.
- Quantum error correction strategy. But can we achieve the known rates?  
Answer: YES!
- Gives an explicit construction of the decoder

## Complementarity of private/entangled states

## Entangled Pairs

$$\Rightarrow H(Z^A|B) = H(X^A|B) = 0 \Rightarrow \psi^{AE} = \psi^A \otimes \psi^E;$$

now follow decoupling argument from before.

$$\Rightarrow \psi^{AB} \text{ has maximal entanglement when } H(Z^A|B) = H(X^A|B) = 0$$

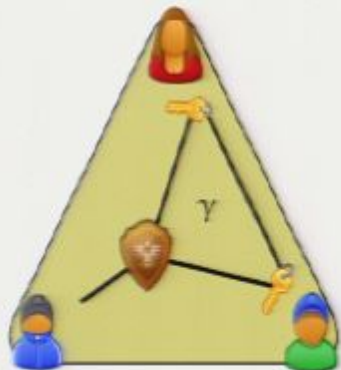
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### Private States Horodecki<sup>3</sup>, Oppenheim PRL 94 160502 (2005)



$Z^A$  generates the key.

☞  $\exists \Lambda^B$  such that  $H(Z^A|\Lambda^B) = 0$

➔ Key is correlated

☞  $\exists \tilde{\Lambda}^{BS}$  such that  $H(X^A|\tilde{\Lambda}^{BS}) = 0$

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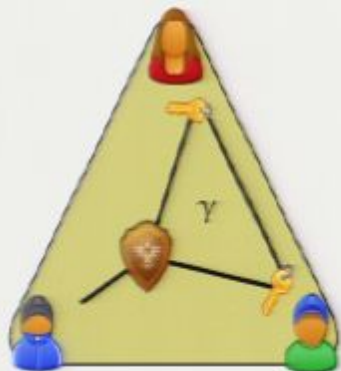
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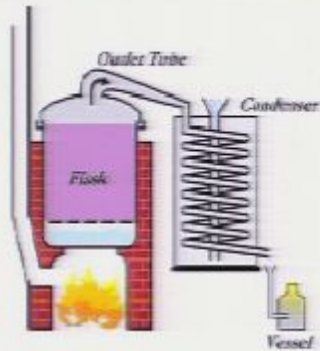
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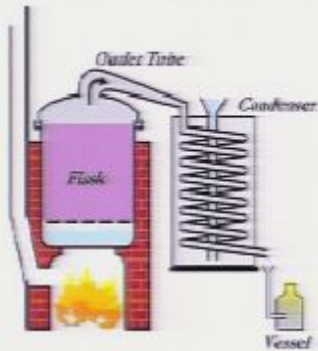
☞ Also works approximately

## Private State/Entanglement Distillation...



Given many copies of an arbitrary resource state,  $\Psi^{AB(S)} = (\psi^{AB(S)})^{\otimes n}$ , convert into (approximate) private/entangled states using local operations and (one-way) classical communication (with high probability).

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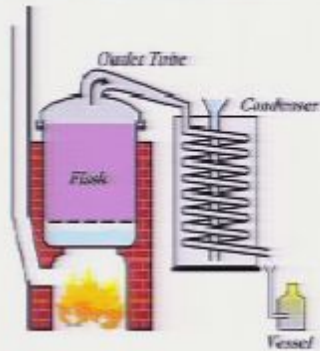
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X  $H(X^A)$

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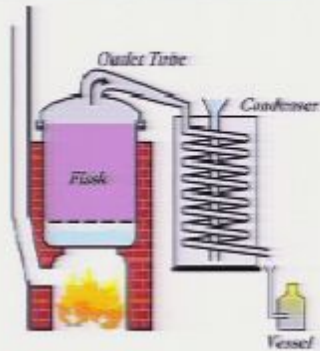
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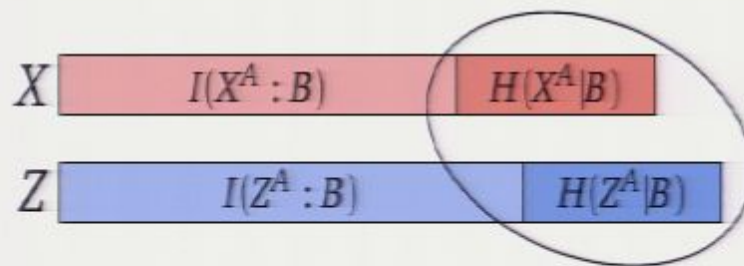
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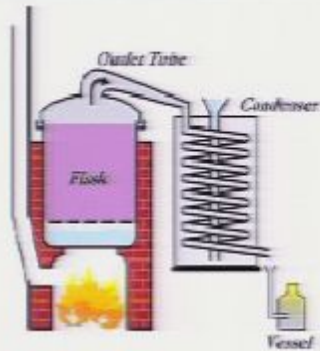
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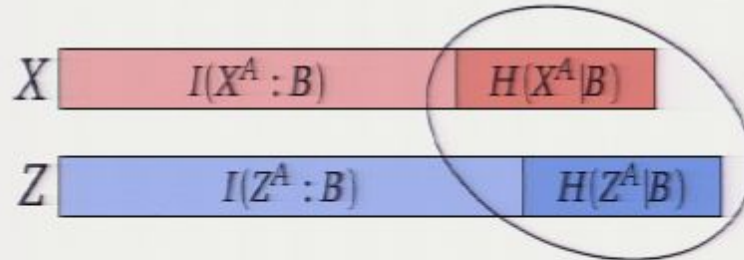
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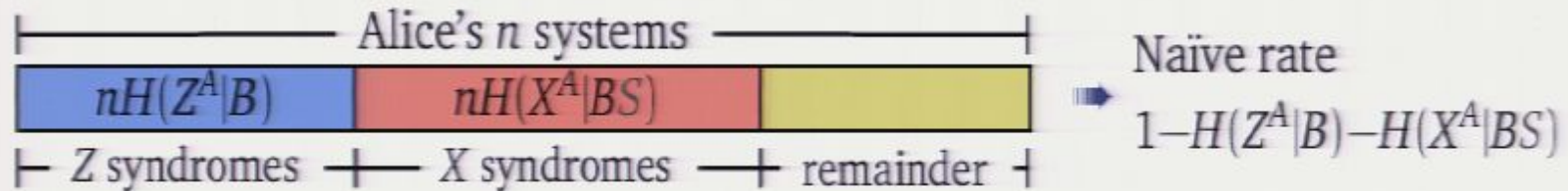
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- Finally, the (static) HSW theorem sets the size of the hints  $\Rightarrow H(Z^A|B)$ .

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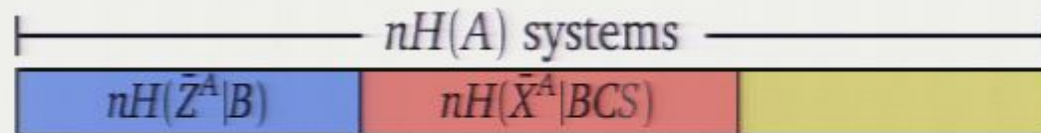


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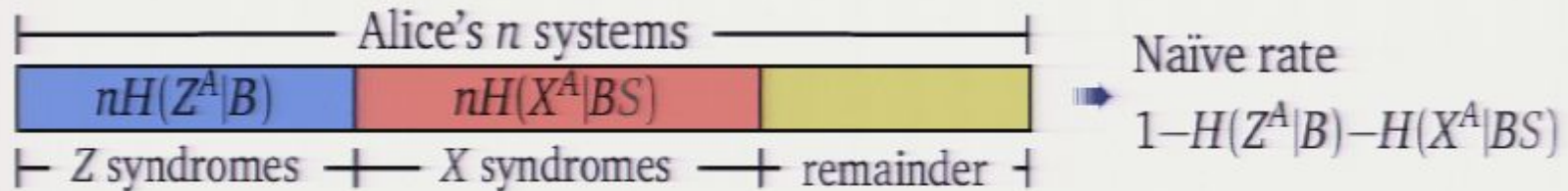
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- ➔ Which screws up the  $X^A$  HSW measurement on the  $BS$  system...
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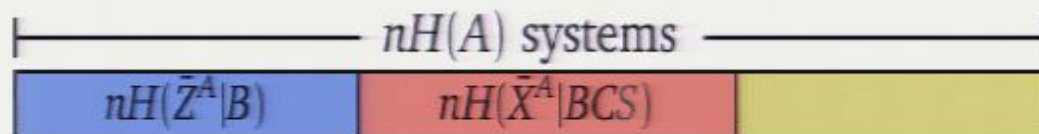
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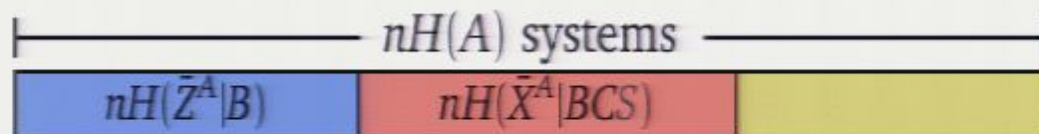
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Fundamental results of quantum information theory rest on the phenomenon of complementarity



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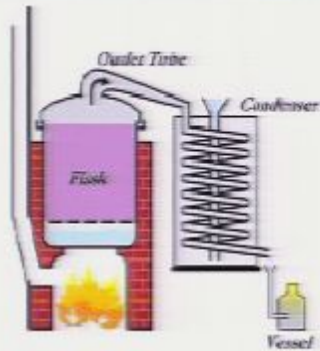


Possible  $z$ s



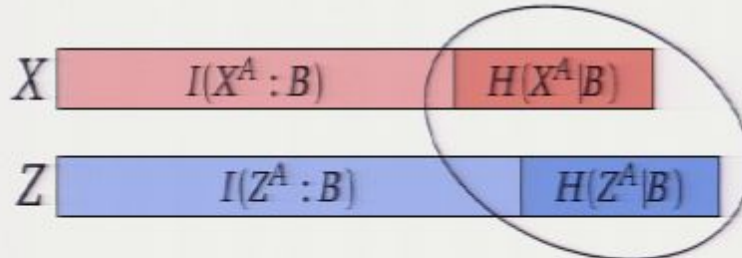
Support of  $\varphi_z^B$

## Private State/Entanglement Distillation...



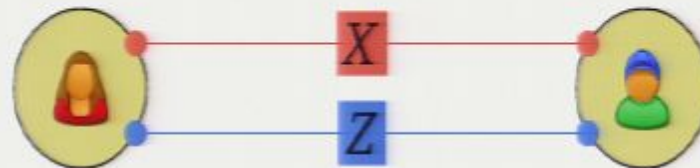
Given many copies of an arbitrary resource state,  
 $\Psi^{AB(S)} = (\psi^{AB(S)})^{\otimes n}$ , convert into (approximate)  
 private/entangled states using local operations and  
 (one-way) classical communication (with high probability).

... by transmitting “missing” information



send this info to Bob.

Couple Alice and Bob classically, twice.



- Quantum correlations between Alice and Bob are equivalent to two complementary classical correlations, say  $X$  and  $Z$  correlations.
- Quantum error correction strategy. But can we achieve the known rates?  
Answer: YES!
- Gives an explicit construction of the decoder



## Complementarity-based Approach

- ☞ Key distillation rate Devetak & Winter  
ProcRoySocA 461 207 (2005)

$$K_{\rightarrow}(\psi) \geq I(Z^A : B) - I(Z^A : E) = H(Z^A|E) - H(Z^A|B)$$

- ☞ Entanglement distillation rate

$$E_{\rightarrow}(\psi) \geq H(B) - H(AB) = -H(A|B).$$

- ☞ Communication cost  $I(A : E)$  in either case.

- ☞ State merging at the same rates/costs Horodecki, Oppenheim, Winter  
Nature 436 7051 (2005)

- ☞ Channel coding from entanglement distillation via teleportation

## Complementary Information Tradeoff

JMR & JCB  
arxiv:0806.3984 [quant-ph]

$$H(O^A|B) + H(\tilde{O}^A|C) \geq -\log \max_{j\tilde{k}} |\langle j|\tilde{k} \rangle|^2$$

- ☞ Replace classical mutual information with Holevo upper bound.  
 $H(O^A|B)$  is quantum conditional entropy after Alice measures  $O^A$
- ☞ Numerical support for conjecture up to dimension 12.
- ☞ Holds for generalized Paulis  $X^A$  and  $Z^A$ .

proof from strong subadditivity Christandl & Winter  
IEEE TIT 51 3159 (2005)

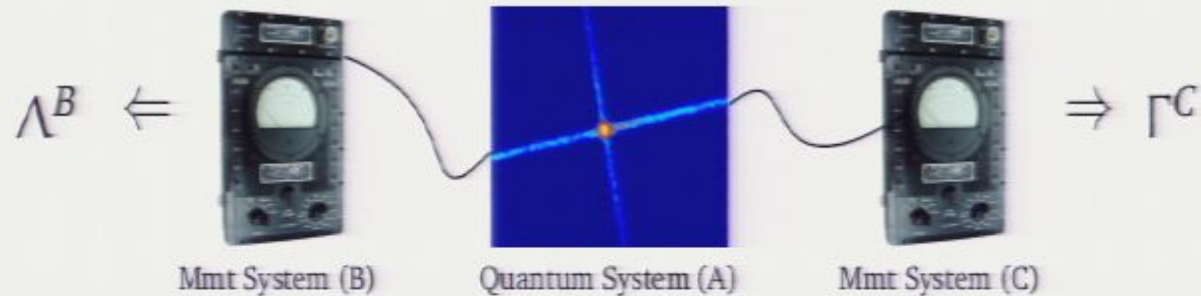


## Entropic Uncertainty Relation

Maassen & Uffink  
PRL 60 1103 (1988)

$$H(O^A) + H(\tilde{O}^A) \geq -\log \max_{j,k} |\langle j|\tilde{k}\rangle|^2$$

$O^A, \tilde{O}^A$  operators on system A; eigenvectors  $|j\rangle$  and  $|\tilde{k}\rangle$ ;  $H(\cdot)$  entropy.



- For generalized Paulis  $X$  and  $Z$ , any state  $|\Psi\rangle$  satisfies  $H(X) + H(Z) \geq \log d$
- But how much info can we *simultaneously* extract? Include mmt systems!

## Information Exclusion Principle

Hall  
PRL 74 3307 (1995)

$$H(O^A|\Lambda^B) + H(\tilde{O}^A|\Gamma^C) \geq -\log \max_{j,k} |\langle j|\tilde{k}\rangle|^2$$

$\Lambda^B, \Gamma^C$  measurements on B, C;  $H(\cdot|\cdot)$  conditional entropy

- tradeoff in simultaneously available, complementary classical information

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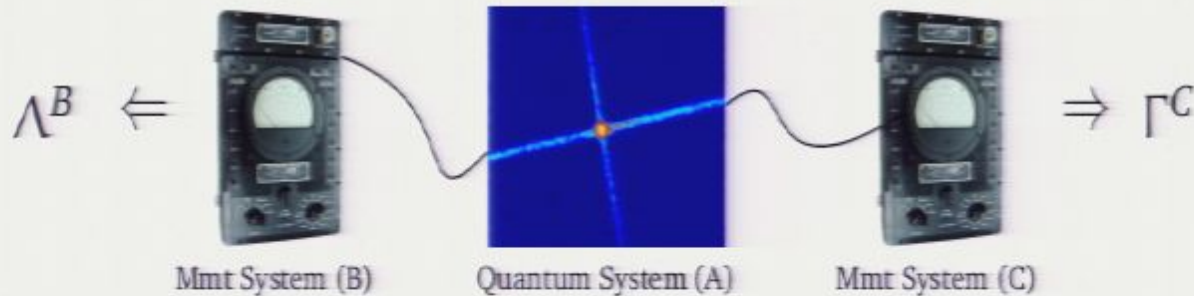


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