

Title: Conceptual problems in phenomenological interpretation in searches for variation of constants and violation of various invariances

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Abstract: At present a number of current or proposed experiments are directed towards a search for a `new physics\` by detecting variations of fundamental physical constants or violations of certain basic symmetries. Various problems related to the phenomenology of such experiments will be considered.

# Conceptual problems in phenomenology for searches for 'new physics'

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and Max-Planck-Institut für Quantenoptik (Garching)



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# Why new physics via precision atomic study?

## We observe at present:

A **fantastic** progress in precision atomic studies

- quantum optics with manipulations on single atoms and precision measurements is the most dynamic part of modern physics
- 3 Nobel prices in 9 years for recent results!

A **terrible** shortage of data and very slow progress on

- a possible extension of the Standard  model of electroweak and strong interactions
  - (even on details of the Standard model – still no Higgs particle has been seen up-to-date);
- quantum gravity.

# Why new physics via precision atomic study?

**We observe at present:**

A **fantastic** progress in precision atomic studies

A **terrible** shortage of data and very slow

- quantum manipulation of atoms
- measurement of most modern
- 3 Nobel prizes in the last 10 years

What is good for atomic physics should be useful for the rest of physics.

**Why not to go for 'new physics'?**

- quantum gravity.




# Outline

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- Variations of constants
- Physics at the Planck scale
- CPT
- Violation of symmetries
- General discussion
- Examples
  - Examples
    - Examples
    - Examples
    - Examples
    - ...






# Testing 'old physics' and looking for 'new physics'

## Two kinds of experiments

Some pretend to be model-independent.

The others are model-dependent. 

Both are not free off conceptual problems, 'hidden' arbitrary suggestions and missintepretations.



# Examples: Variations and LTI

---

Variation of a constant  
 $A$  means that

$$A(t_1) \neq A(t_2) .$$

LTI (Local Time Invariance), a part of `big` relativistic invariances, requires in particular that a result of a measurement of a fundamental scalar value for  $t_1$  and  $t_2$  should be the same.



# Examples: Variations and LTI

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LTI (Local Time Invariance), a part of 'big' relativistic invariances, requires in particular that a

Often it is believed that  
**LTI is** equivalent to non-variation  
of the fundamental constants.

**However, that is not.**



# Examples: Variations and LTI

---

- LTI suggests  
`fundamental  
experiments` :
- free particles,
  - free space,
  - **no environment,**
  - etc.

- Example of the  
environment:
- local gravity is  
changing in time  
because of the Sun  
and the Moon at a  
1-ppm level;
  - $g$  can be measured  
at 10-ppb level.

# Examples: Variations and LTI

LTI suggests  
`fundamental  
experiments`:

- free particles,
- free space
- no environment
- etc.

Example of the  
environment:

- local gravity is  
changing in time

In modern precision study  
the `environment` may be  
something far outside of the lab.

**Matter**, dark matter,  
bath of Cosmic Microwave  
Background ( $\gamma$ ,  $\nu$ ,  $g$ ).



# Examples: Variations and LTI

LTI suggest

Example of the

ment:

ivity is

g in time

The inflation model:

once and long long time ago

in a very remote past

(even before MPQ was established

and before hydrogen atoms were formed)

real changes in values of  $m_e$ ,  $\alpha$ ,  $m_p$  took place.

study

may be

of the lab.

etc.

**Matter**, dark matter,  
bath of Cosmic Microwave  
Background ( $\gamma$ ,  $\nu$ ,  $g$ ).

# Examples: Variations and LTI

LTI suggest

Example of the

Why?

Because of cooling of the Universe.

The very Universe acted  
as a piece of the environment  
and caused a phase transition  
changing  $m_e$ ,  $\alpha$ ,  $m_p$  etc:

no violation of LPT,  
but still the variations took place.





# Examples: EDM of an electron

## Common knowledge:

- Modern theory allows no EDM until some violations of basic symmetries are involved.
- A standard electron has the EDM because of CP violation in Standard model (SM) but it is below any detectable level.
- The motivation for a search of a bigger value of the EDM is due to extensions of the SM.
- Their allowed range is for about 5 orders of magnitude.
- The results relate to half of the range (in the logarithmic scale) and they are negative.

**The interpretation is model-independent.**



# Examples: EDM of an electron

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Model-dependent!

Correct!

Incorrect!

# Examples: EDM of an electron

What else can induce the EDM?

- Modern theory allows no EDM until **some violations of basic symmetries** are involved.
- If some relativistic invariances are violated (isotropy, the Lorentz invariance) that could produce the EDM.

A violation of the Lorentz invariance.

- The preferred frame.
  - `Local` - Dark Matter
  - `Global` - CMB
- The EDM can be directed not along the spin  $\mathbf{s}$ , but **along the velocity  $\mathbf{v}$**  in respect to the preferred frame or along  $[\mathbf{v} \times \mathbf{s}]$ .



# Examples: EDM of an electron

What else can induce the EDM?

- Modern theory allows a violation of the Lorentz symmetry. This involves a preferred frame. This can be the rest frame of the Dark Matter or the CMB. The EDM can be induced along the direction of the preferred frame or along  $[\mathbf{v} \times \mathbf{s}]$ .
- If some symmetry is broken, different experiments are needed. But still: the EDM is not along the direction of the preferred frame or along  $[\mathbf{v} \times \mathbf{s}]$ .

That is indeed a very different situation.

Different experiments are needed.

But still:

that is the EDM.





# `Variation of the constants`

---

- The very wording is inconsistent:  
`constants` cannot vary.
- That is not a problem of semantics.
- The very conception of `constants` assumes:
  - there are certain basic equations
  - and there are certain parameters there.
- `Variation of the constants` suggests that we can
  - keep the equations,
  - but vary their constants.
- **That is inconsistent.**
- If we want `to change` the constants, we **have** to re-consider the equations.



# Two examples:

The Maxwell equations of propagation of a photon at media (which obviously violate LPT, LTI, isotropy and relativity)

- not just  $c(t,x)$ ,
- but:
  - $\varepsilon(t,x)$
  - $\mu(t,x)$
  - and **their time and space gradients** are present in basic equations!

A model:

- a space-time with  $\hbar(t)$ , but not with  $\hbar(x)$  – that an isotropic situation!
- the electron's spin angular moment,  $L$ , is conserved, but

$$J_z = \pm 1/2 \hbar(t)$$

which is inconsistent.



# The Feynman's continual integral and the interference

The soul of quantum theory is interference and a way to include it as a base principle as a continual integral.

- Quantum mechanics is based on the path integral (over possible trajectories).
- Quantum field theory is based an integral over all evolving field configurations.

**The action is the phase.**  
The least action is an enhanced trajectory leading to classical mechanics.







# Modification of 'old physics'

- We should take the most fundamental level (quantum field theory).
- We should write the proper confinual integral.
- We should introduce various dimensionless factors  $\zeta(x,t)$ .
- A chage  $c \rightarrow c \zeta(x,t)$  is a kind of description of  $c(t)$ .

- Example:  
in classical  
electrodymanic  
Lagrangian a  
substitute

$$\mathbf{E}^2 + \mathbf{B}^2 \rightarrow \epsilon \mathbf{E}^2 + \mathbf{B}^2/\mu$$

leads to the Maxwell equations in media, properly introducing gradients of  $\epsilon(x,t)$  and  $\mu(x,t)$ .



# Modification of 'old physics'

- We should take the most fundamental level (quantum field theory).
- We should write the proper Lagrangian

**E** and **B** are components of  $F_{\mu\nu}$ ,  
the same quantities used  
for the vacuum case.

- A charge  $c \rightarrow c \zeta(x,t)$  is a kind of description of  $c(t)$ .

- Example:  
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# Modification of 'old physics'

- We should take the most fundamental level
- Example: in classical

In fact even in Lagrangian some gradients can appear.

If we start from  $dX/d(ct)$  we can substitute it for

$(1/c)(dX/dt)$  or  $(d/dt)(X/c)$  or  $dX/d(ct)$ .

Going from Lagrangian to equations we produce more gradients.





# Gradient terms (GT)

---

- GT allow to look for variations of dimensional quantities:
  - detecting gradient rates
  - directly comparing **quantities** (not the results!) for different cases (like the MM experiments)
- Gradients can be introduced only on base a model and the results are most likely **model-dependent**.



# Gradients and different kinds of searches for variations

There are three kinds of searches for variations of fundamental constants, **which differently deal with GT**

- fast measurements – long separation
- monitoring
- differential measurements

# Gradients and different kinds of searches for variations

There are three kinds of searches for variations of fundamental constants, which differently deal with  $GT$

- fast measurements – long separation
- monitoring
- differential measurements

- Fast measurement – gradient effects  $\sim dX/dt \times \tau$ .
- Long separation – difference  
 $X(t+T) - X(t) \sim dX/dt \times T$ .
- $T \gg \tau$ . The gradient effects can be neglected.
- These measurements deal with dimensionless results.





# Gradients and different kinds of searches for variations

There are three kinds of searches for variations of fundamental constants, which

- Fast measurement – gradient effects  $\sim dX/dt \times \tau$ .
- Long separation –

In atomic spectroscopy:

the time of the measurement is a combination of the interaction time (photon + atom) and the coherence time.

measurements

these measurements deal with dimensionless results.



# Gradients and different kinds of searches for variations

There are three kinds of searches for variations of fundamental constants, which differently deal with GT

- fast measurements – long separation
- **monitoring**
- differential measurements

That can be

- a continuous measurement
- long measurements – long separation
- phase-coherent separated measurements.
- For any of them  $T \sim \tau$ . The gradient effects cannot be neglected.
- **Any interpretation need a 'complete' theory.**



# Gradients and different kinds of searches for variations

What is 'the phase coherent separated measurements'?

That is a kind of a classical Ramsey-scheme experiment.

Dealing with an atom we do not 'know' the position of an electron and we study a kind of average radius.

Dealing with a planet we see the phase of its rotation.

Even doing 'fast measurements' with 'long separation', but 'remembering' the phase, we should still deal with GT.

That can be

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- phase-coherent

Likely all tests of General Relativity and searches for variations of  $G$  are a kind of monitoring.

**No particular constraint on  $dG/dt$  obtained without any assumption of GT cannot be correct!**



# Gradients and different kinds of searches for variations

What is the phase coherent

Another example: measuring  $R$  in two body system.

For simplicity:  $M \gg m$  and a circular orbit

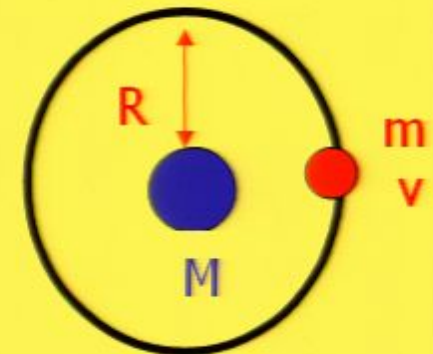
Everything is stable:

Kepler's law:  $v^2 R = GM$ ,  
a conservation of  $v = 2\pi R/T$   
and no dependence on  $m$ .

If  $G$ ,  $M$  and  $m$  are changing there are two effects:  
changing  $GM$  shifts  $R$  directly,  
changing  $m$  but keeping  $p = mv$  shifts  $v$  and thus  $R$ .

The results can be found from conservation laws:

$$R = GM \times m^2/p^2.$$





# Gradients and different kinds of searches for variations

What is the phase coherent

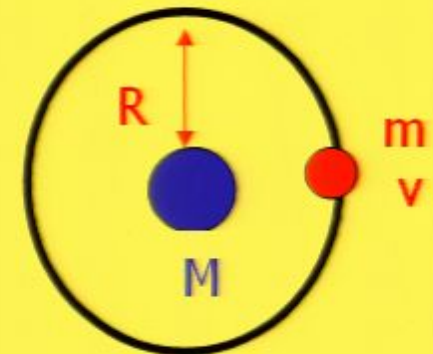
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 $R = GM \times m^2/p^2$ .



From equations of motion:  
 changing  $GM$  – no gradients,  
 changing  $m$  produces term  
 $m a = \dots - v dm/dt$ .



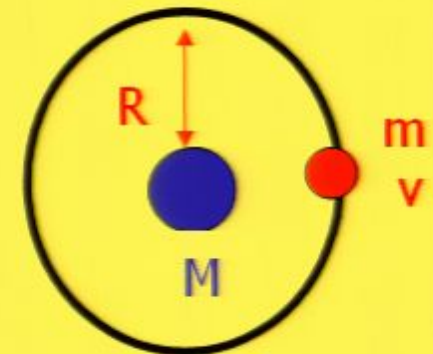


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And to be precise we got a result for a time variation of

$$GM \times m^2 m_e$$

which is in a quite realistic scenario

$$\sim m_p^3 m_e / M_{pl}^2.$$

That is dimensional which can be because of a gradient effect is involved.

From equations of motion:

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# Gradients and different kinds of searches for variations

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Another example: measuring  $R$  in two body system.

F Mass cannot change in conventional physics.

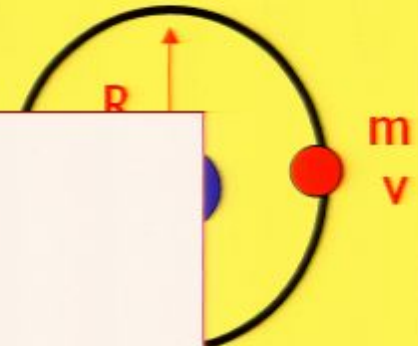
The choices:

1. Use a 'completely new physics'.

2. Use a modification of old physics:

e.g.  $m$  is varying because it picks up certain amount of dark matter or dust or because of evaporation and radiation. – That reads: we introduce a mechanism which may imply transfer of other values ( $E$ ,  $p$ ,  $L$  etc).

That is dimensional which can be because of a gradient effect is involved.



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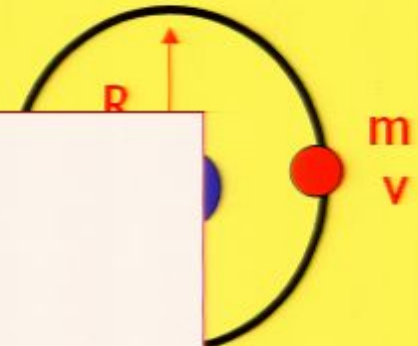
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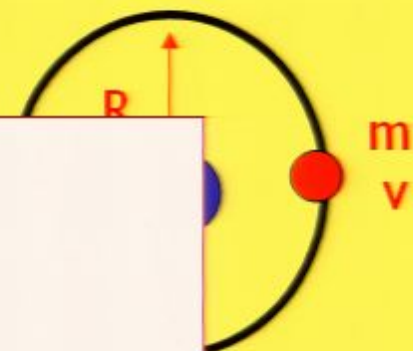
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# Gradients and different kinds of searches for variations

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There are three kinds of searches for variations of fundamental constants, which differently deal with GT

- fast measurements – long separation
- monitoring
- differential measurements
- Experiment directly sensitive to a gradient of a constant
  - e.g. for a transition which is proportional to the gradient
  - Those experiments are obviously model-dependent.
- MM-like experiments



# Gradients and different kinds of searches for variations

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- MM-like experiments

MM-experiments are also model dependent.





# The renormalization and the Planck-scale physics

---

- We believe that the most fundamental physical quantities are defined at the Planck scale.
- We observe various effects at energy scale much below the Planck scale.
- To build a really fundamental theory we should express observables in terms of the Planck-scale-determined values,
  - **which we cannot do.**

# The renormalization and the Planck-scale physics

- We believe that the most fundamental physical quantities are defined at the Planck scale.
- To build a really fundamental theory we should express observables in terms of

In a sense we have not a fundamental theory, but a **'fundamental constraint'**.

It is not a theory, but a **constraint**, in a sense that we **cannot predict** *ab initio* the experimental results, but only **relations** between them.

It is a **fundamental** constraint in a sense, that we need to perform **very few** measurements to predict **a lot**.





# The renormalization and the Planck-scale physics

---

The main problem is the Planck-scale physics.

We cannot derive any properties for `our` energies from there because of lack of knowledge about their physics.

Instead, we have developed the renormalization procedure which allows to avoid any involvement of that physics into low-energy calculations.

Such approach is successful and one can expect that we cannot see any effects from the Planck-scale physics at all.

**That is not correct.**

We are still able to see dynamics at the Planck scale if any.

Simple example: a compactification with  $R(t)$ .



# The renormalization and the Planck-scale physics

The example:

a multidimensional world. All space dimensions, but 3, are compact with compactification radius  $R$ , which is  $\sim$  Planck length, but  $R = R(t)$ .

In such a case,  $R(t)$  could enter the renormalized results, e.g.,

$$\alpha = \alpha_0 (1 + k \alpha_0 \ln(m_e R)) .$$

As long as  $\alpha_0$  is a constant, we can see no effects related to  $\alpha_0$ , because all results, such as the Lamb shift, are expressed in terms of  $\alpha$ .

Still: we should see variation of  $\alpha$ , caused by the Planck scale via  $R(t)$ .

involvement of that physics  
into low-energy calculations.

Simple example: a  
compactification with  $R(t)$ .

# The renormalization and the Planck-scale physics

A classical/quantum analogy:

Usually, quantum effects are small and dealing with classical macroscopic bodies we can hardly see effects from the quantum scale.


However, we can see non-conservation effects.

In classical physics the mass is conserved.  
In quantum it does not.

Measuring classical mass we can see that.

Another example: no problem to see radioactivity.





# The MM experiment and the Planck scale physics

---

- The experiment

The classical MM experiment suggested that the path of a light, determined by a ruler of a solid-state bulk matter remains the same, while the velocity of the light could change if we change the direction of the propagation.

- A suggestion:

In the leading approximation the length of a substance ruler is determined by the Coulomb force in the atoms.

- The question:

Is the Coulomb force relativistic or not?

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# The MM experiment etc: two scenarios

The basic question is what is the fundamental quantity?

- The charge.
- Its value is a result of, e.g., spontaneous breakdown, not related to relativity.
- The Coulomb law is

$$\sim Z_1 Z_2 e^2 / r$$

- We can violate the relativity not affecting the charge and the Coulomb law.

The path does not change.

- The fine structure constant  $\alpha$  is calculable (at the Planck scale).
- The charge is calculable via  $\alpha, \hbar, c$ . The Coulomb law is

$$\sim Z_1 Z_2 \alpha \hbar c / r$$

- If we violate the relativity, the Coulomb law is sensitive to that.

The path changes.

# The MM experiment etc: two scenarios

The basic question is what is the fundamental quantity?

- The charge.
- Its value is a result of, e.g.,
- The fine structure constant  $\alpha$  is calculable (at the Planck

A suggestion related to the Planck scale is crucial for the interpretation of the classical MM experiment.

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The path does not change.

The path changes.





# How to violate a symmetry?

---

- A direct violation:  
the mass as a violation  
of the left-right  
symmetry
- Anomaly  
 $\pi^0$  decay:
  - a singularity
  - the regularization is  
possible only for a  
part of symmetries
- Spontaneous breakdown  
the Standard model;  
magnetic properties:
  - the potential is  
symmetric
  - the symmetric states are  
not stable
  - the least energy states  
are not symmetric

# How to violate a symmetry?

- A direct violation:  
the mass as a violation

Analogy: when we measure  $L_x$  and  $L_y$ , that is well defined in classical mechanics, but needs additional 'regularizations' in quantum mechanics, by e.g., saying that  $L_x$  is measured at  $t - \epsilon$  and  $L_y$  is at  $t + \epsilon$  with  $\epsilon \rightarrow 0$

The result depends on sign of  $\epsilon$ .

After 'regularizing' the measurements, both  $L_x$  and  $L_y$  are still conserved,

but the conservations cannot be seen at one time.

- Spontaneous breakdown

the Standard model;  
magnetic properties:  
the potential is symmetric  
the symmetric states are not stable  
the least energy states are not symmetric



# How to violate a symmetry?

- A direct violation of the mass as a function of the left-right symmetry
- Anomaly  
 $\pi^0$  decay:
  - a singularity
  - the regularization is possible only as a part of symmetry

Example: **left**-handed and **right**-handed organic molecules.

Electrodynamics is **isotropic**, but there is no stable state symmetric in respect to reflection  $x \rightarrow -x$ .

Usual symmetric combination of LH and RH molecules does not work because the tunnel transition time is too large.

The LH asymmetric molecule is destroyed much much faster than the oscillation period.



# How to violate a symmetry?

---

## **Observational problems:**

- we cannot look for a symmetry, but only for its consequences;
- the symmetry may be conserved, but its naive consequences would be incorrect.
- Non-commutativity of the coordinates in quantum gravity.

An analogy:

- **Classical mechanics:**  
For isotropy we can measure  $L_x$ ,  $L_y$ ,  $L_z$  and repeat the measurements.  
The results stand the same.
- **Quantum mechanics:**  
the results are not the same even for an isotropic problem.





# How to violate a symmetry?

---

The most natural language for a direct violation is external field:

- classical (matter, dark matter etc);
- quantum (condensate).

Such a field may be very similar to conventional electromagnetic or gravitational field, but is is `selective`.



# How to violate a symmetry?

The most natural  
language for a direct

Such a field may be  
very similar to

What does it mean 'selective'?

The electric field is universal in a sense, that  
its interaction is of the same form for  
any elementary particle.

The difference is only a charge of a particle,  
which is 0 or  $\pm 1$  for free particles.

ic or  
eld,  
tive'.



# CPT violation and a 'selective' field: an example

What experimentalists  
want for CPT:

- different charges –  
not good because of  
gauge invariance  
and charge  
conservation;

- masses:

$$m_{\pm} = m_0 \pm \delta m$$

We can rewrite that as

$$m_{\pm} = m_0 + qeU$$

where  $q = \pm 1$

and  $U = \delta m/e$ .

$U$  is an effective  
potential.

As long as we deal  
only with electrons  
we cannot see any  
effects due to  $U$

# CPT violation and a 'selective' field: an example

What experimentalists

If we deal with few kinds  
of particles ( $e, \mu, \dots$ )  
and  $U_e \neq U_{\mu'}$   
i.e., if the field is selective,  
we can observe effects  
proportional to  $U_e - U_{\mu'}$   
e.g. the muon decay.

We can rewrite that as

$$m_{\pm} = m_0 + qeU$$

where  $q = \pm 1$

and  $U = \delta m/e$ .

$U$  is an effective  
potential.

As long as we deal  
only with electrons  
we cannot see any  
effects due to  $U$

The result is **not** CPT asymmetry of the mass.



# CPT violation and a 'selective' field

- The selective potential  $U$  is similar to Kostelecky's  $a$  term.
- Some other terms are similar to the magnetic field.
- The parametrization is related to the vacuum.
- Often in experiments we should deal with residual electromagnetic fields.
- We shield from them.
- The shielding may be based on re-distribution of the electron
  - charges
  - currents
  - spins

The true magnetic field universally interacts with spin and orbital motion (of the same particle) and with different particles.

# CPT violation and a 'selective' field: example

- The selective potential  $U$  is similar to Kostelecky's  $a$ -term.
- Some other terms are

$$e \mathbf{j}_i \mathbf{A}_i + \mu_i \mathbf{s}_i \mathbf{B}_i$$

For true residual magnetic field:

$$\mathbf{B} = \mathbf{B}_i = \text{rot } \mathbf{A}_i$$

For violated symmetry they are small and independent.

- The parametrization is related to the vacuum.
- Often in experiments we should deal with residual electromagnetic fields.
- We shield from them.
- The shielding may be based on re-distribution of the electron
  - charges
  - currents
  - spins



# CPT violation and a 'selective' field: example

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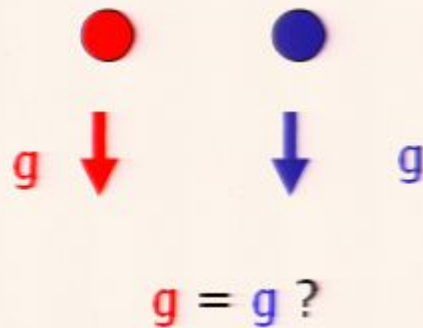
For violated symmetry they are separate  
and independent.

When shielding residual magnetic field  
by rearranging 'electron spins we change  
the effective hamiltonian to

$$e \mathbf{j}_i (\mathbf{A}_i - [\mathbf{r} \times \mathbf{B}_e]) + \mu_i \mathbf{s}_i (\mathbf{B}_i - \mathbf{B}_e)$$

- currents
- spins

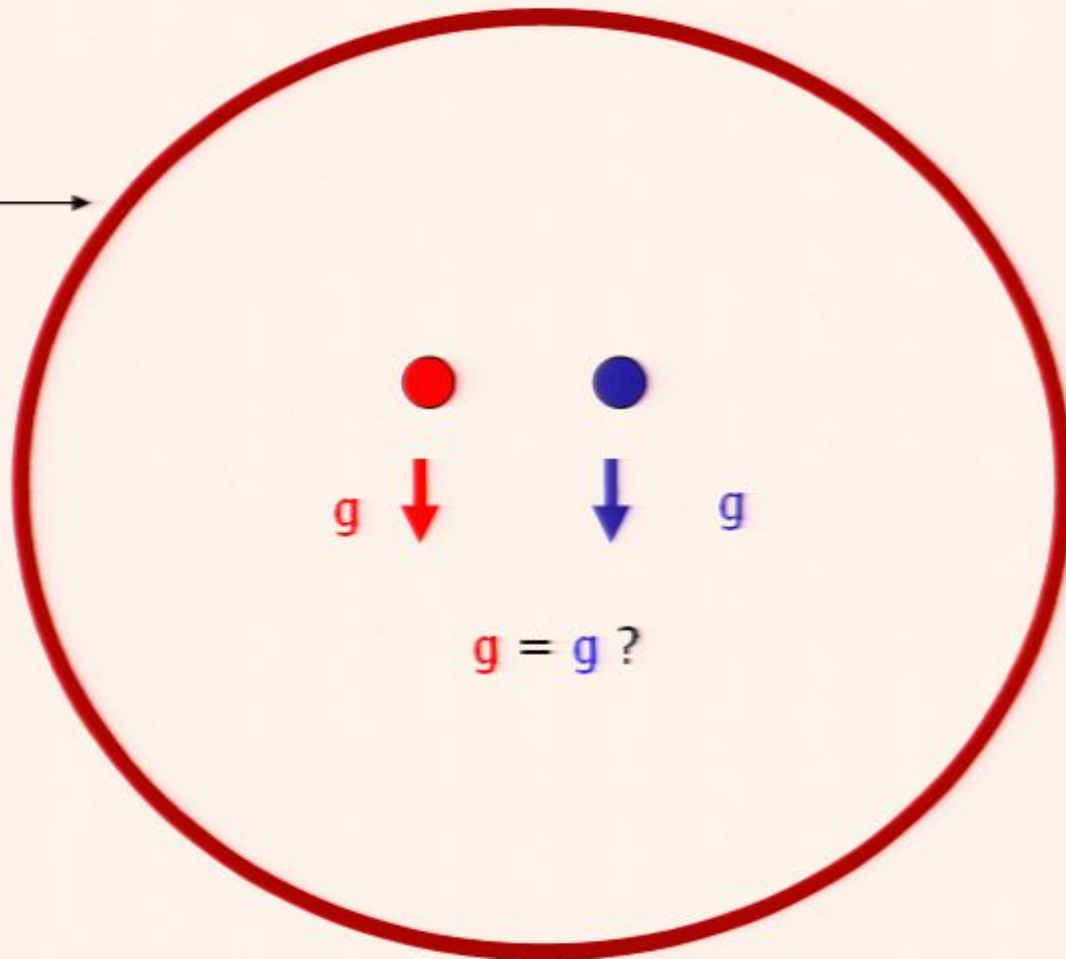
# CPT: gravity of **electron** and positron





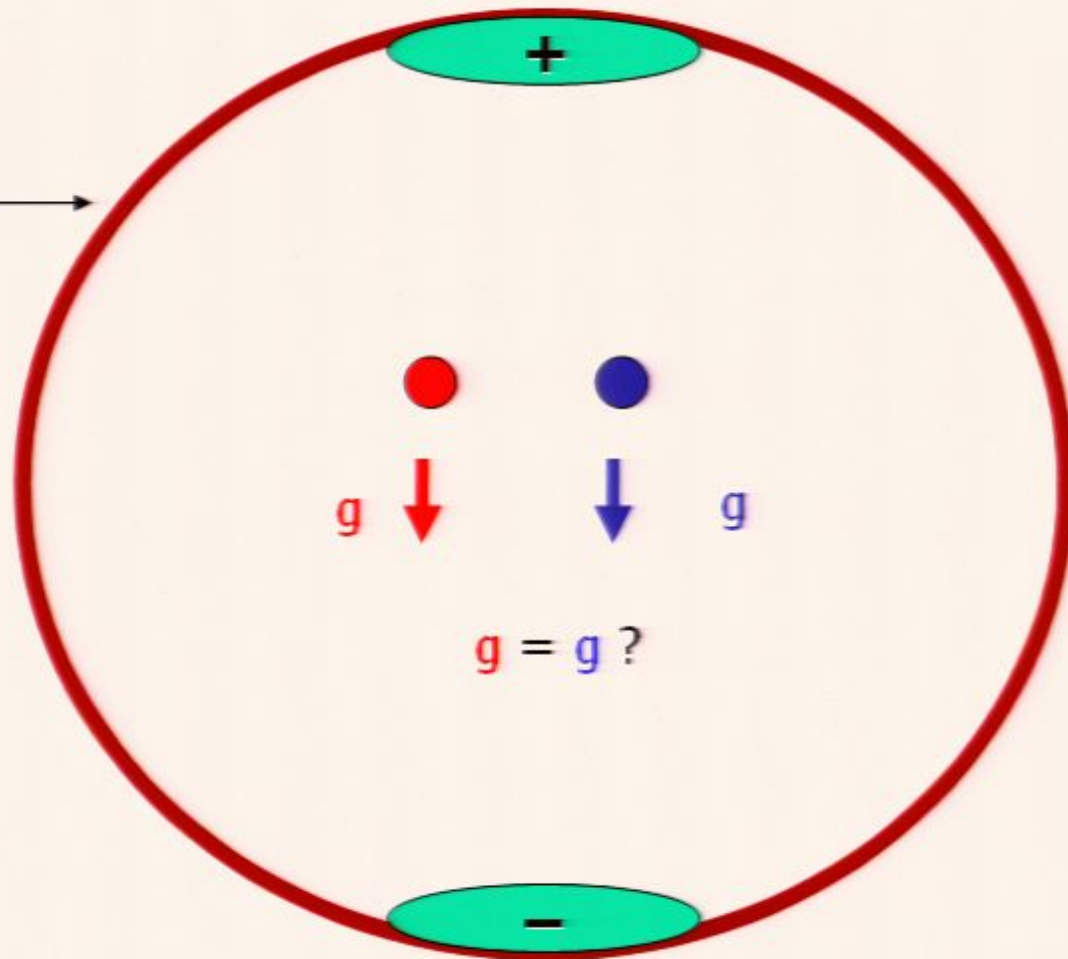
# CPT: gravity of **electron** and **positron**

Shielding  
from electric  
field



# CPT: gravity of **electron** and **positron**

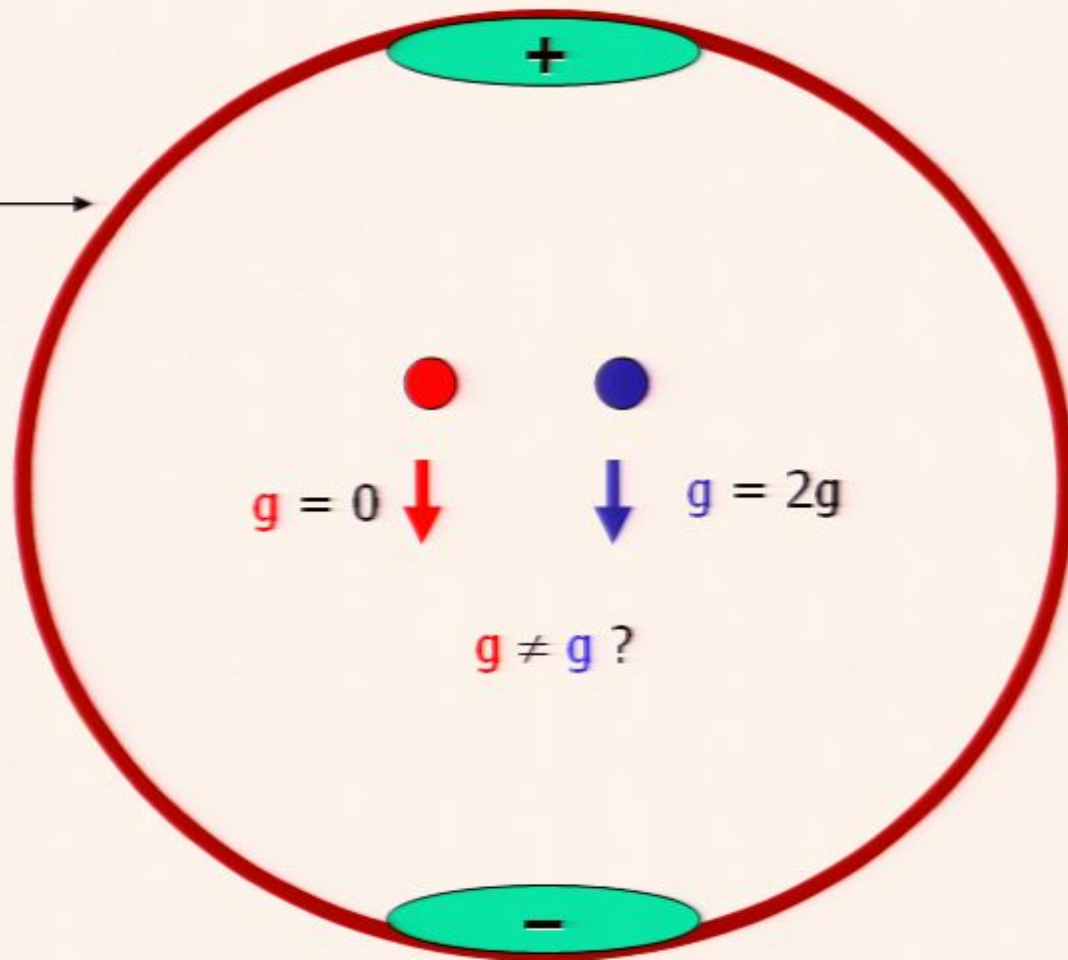
Shielding  
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field





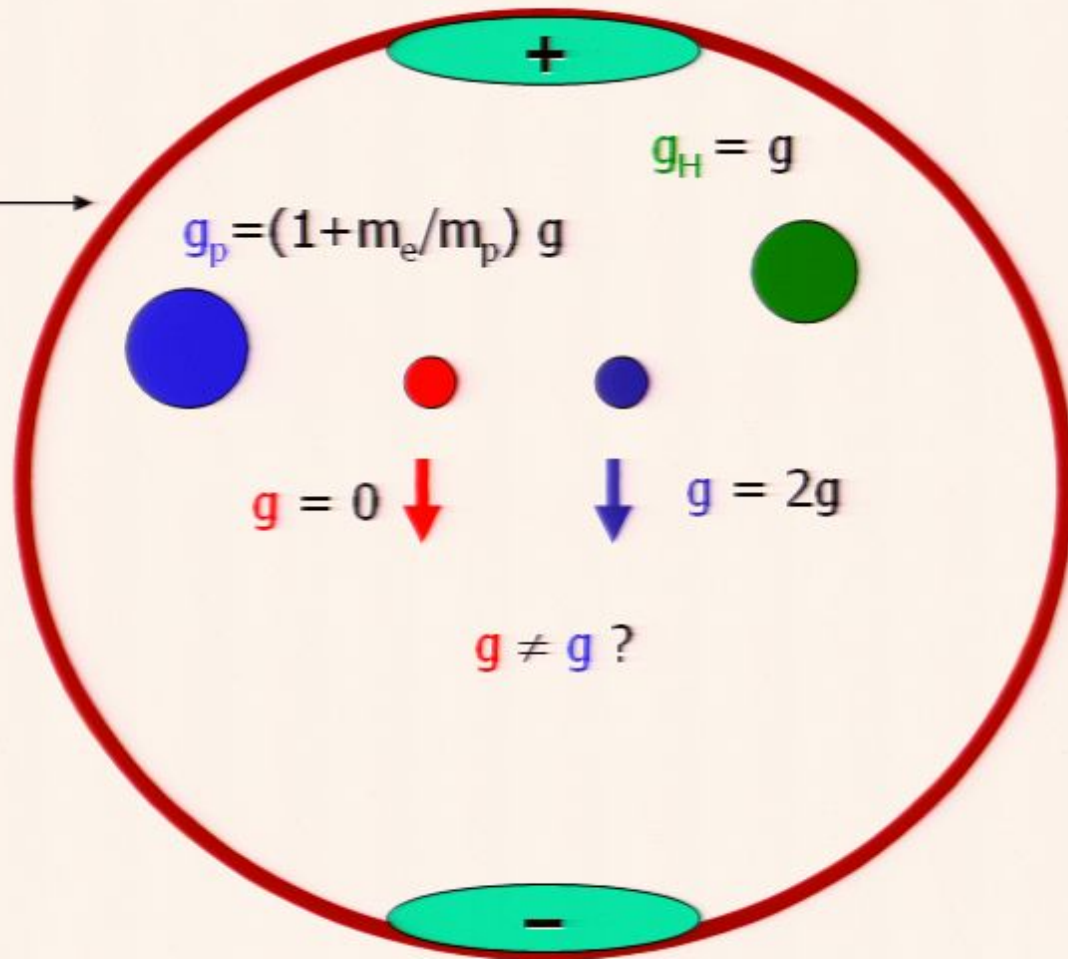
# CPT: gravity of **electron** and **positron**

Shielding  
from electric  
field



# Gravity of electron, positron, proton, and hydrogen

Schielding  
from electric  
field







# Why these fancy `violating terms` looks as the fields?

When we want to build a theory we are very limited in what we can write. The result should be relativistic invariant, local etc.

Violating fields and true fields (EM or G) satisfy similar conditions.

However, when we have a true fundamental theory (EM or G) we have certain constraints on relation between different terms (spin/orbit; Newton/PostNewton) and a kind of universality (conservation of charge, mass etc).

# Microscopic and macroscopic description

- One more issue about effective fields and equations is due to their scale.
- An example: a kind of classical-field effect can be caused by dark matter.

Experiments of macroscopic scale would show a departure from conventional relativity (e.g., the MM experiment).

Experiments of microscopic scale would confirm relativity (e.g., the  $g-2$ ,  $E_0=mc^2$ , and the hydrogen spectrum).





# Summary

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To introduce into phenomenology a certain piece of 'new physics' in a model independent way is not as easy and simple as some expect.

**Don't panic, but be careful!**

Gravity of electron, positron, proton, and hydrogen

Why these fancy 'violating terms' looks as the fields?

When we want to talk about the violation of the symmetry, we should use the language of the fields. The symmetry is broken by the insertion of the fields. The fields are the source of the violation. The fields are the source of the violation. The fields are the source of the violation.

Microscopic and macroscopic description

The microscopic description is the one that is most useful. The microscopic description is the one that is most useful. The microscopic description is the one that is most useful.

Summary

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savely

- Internet Explorer
- Outlook Express
- Skype
- FreeCell
- Spider Solitaire
- RPAD32.EXE
- Adobe Reader 7.0
- Solitaire
- Microsoft Office Word 2003
- WinEdt
- Microsoft Office PowerPoint 2003
- IrfanView - Thumbnails

All Programs

- My Documents
- My Recent Documents
- My Pictures
- My Music
- My Computer
- My Network Places
- Control Panel
- Set Program Access and Defaults
- Connect To
- Printers and Faxes
- Help and Support
- Search
- Run...

# Summary

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Don't panic, but be carreful!

Outline Slides

Gravity of electron, positron, proton, and hydrogen

Why these fancy "violating terms" looks as the fields?

When we want to talk about the interaction between a fermion and a boson, we usually use the term "interaction".

However, when we talk about the interaction between a fermion and a fermion, we usually use the term "exchange".

Microscopic and macroscopic description

The more we know about the microscopic world, the more we understand the macroscopic world.

An example: a kind of microscopic effect can be caused by dark matter.

Summary

To introduce into phenomenology a certain piece of new physics, in a model-independent way is not as easy and simple as some expect.

Don't panic, but be careful!

Turn off computer

Stand By Turn Off Restart

Cancel

phenomenology  
new

independent way is not as easy and simple as some expect.

Don't panic, but be careful!



Gravity of electron, positron, proton, and hydrogen

Why these fancy 'violating terms' looks as the fields?

When we want to talk about the interactions between the fields, we can write the Lagrangian as the sum of the kinetic terms and the interaction terms. The interaction terms are the ones that are 'violating'.

Microscopic and macroscopic description

The microscopic description is the one that is most fundamental. It is the one that is most difficult to understand. The macroscopic description is the one that is most intuitive. It is the one that is most easy to understand.

Summary

To introduce into phenomenology a certain piece of 'new physics' in a model independent way is not as easy and simple as some expect.

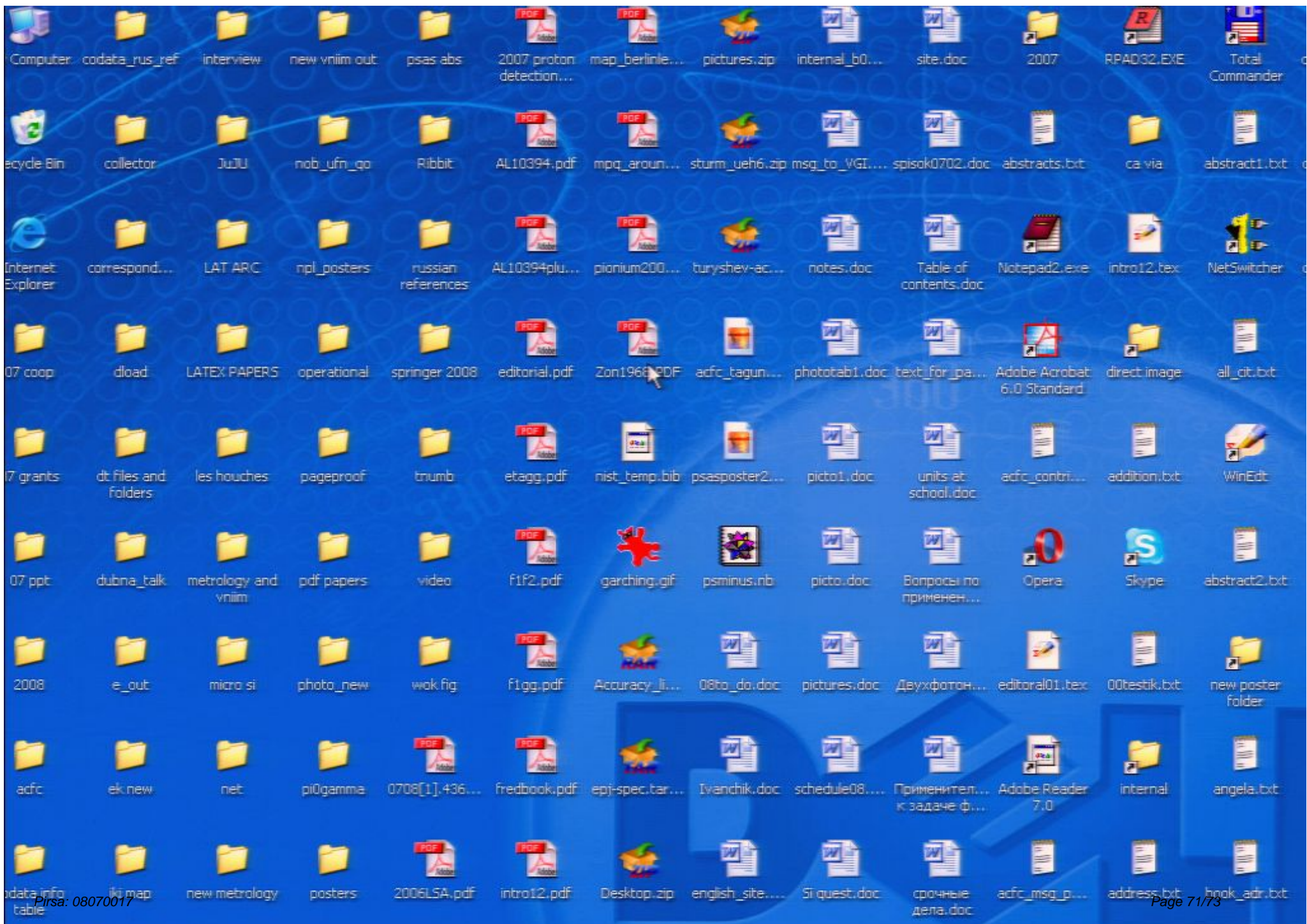
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Logging off...