

Title: Testing the m_p/m_e cosmological variation from H_2 lines in High-Redshift QSO spectra

Date: Jul 15, 2008 10:10 AM

URL: <http://pirsa.org/08070011>

Abstract:

*Testing $\mu = m_p/m_e$ cosmological variation
from H₂ lines in high-redshifted QSO spectra*

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W.Ubachs, E.Reinhold

Dirac

Nature, 1937

"Large Number Hypothesis"

$$G \sim t^{-1}$$

Contemporary theoretical models of unification of all the physical interactions predict variations of fundamental physical constants in the course of the cosmological evolution.

In particular, Superstring/M-theory predict time-variations of coupling constants and masses of elementary particles:

$$g_i = g_i[f(\phi)] \quad \text{and} \quad m_i = m_i[f(\phi)]$$

However, different theories predict different variations and different relations between constants (e.g. $\alpha = e^2/\hbar c$ and $\mu = m_p/m_e$; Calmet & Fritzsch, 2002, Phys.Lett.B540; Langacker et al., 2002, Phys.Lett.B528; Flambaum & Shuryak, 2002, Phys.Rev.D65):

$$\alpha(\mu) \iff \mu(\alpha)$$

Searching of μ -variation becomes more significant after possible finding of α -variation (Webb et. al.).

Experimental detection or even determination of rigid upper limit on μ -variation would be an important tool for selection of theoretical models of the fundamental physical interactions.

The last observational results

| Epoch | Reference | Constant | $\Delta x/x$ | $\dot{x}/x, \text{yr}^{-1}$ |
|----------------------|------------------------|----------------------------|---------------------------------|----------------------------------|
| “Now and Here” | 2007 Fortier et al. | α | $(3.6 \pm 6.0) \times 10^{-15}$ | $(-0.6 \pm 1.0) \times 10^{-15}$ |
| | 2006 Peik et al. | α | $(1.6 \pm 2.3) \times 10^{-15}$ | $(-0.3 \pm 0.4) \times 10^{-15}$ |
| | 2003 Bize et al. | α | $< 2.4 \times 10^{-15}$ | $< 1.2 \times 10^{-15}$ |
| | 1995 Prestage et al. | α | $< 3.7 \times 10^{-14}$ | $< 3.7 \times 10^{-14}$ |
| “Oklo” | 2006 Petrov et al. | α | $(0.1 \pm 0.7) \times 10^{-7}$ | $(-0.5 \pm 3.5) \times 10^{-17}$ |
| | 2004 Lamoreaux et al. | α | $(4.5 \pm 0.2) \times 10^{-8}$ | $(-2.3 \pm 0.1) \times 10^{-17}$ |
| | 2003 Fujii et al. | α | $(8.8 \pm 0.7) \times 10^{-8}$ | $(-4.4 \pm 0.4) \times 10^{-17}$ |
| “Q S O” | 2007 Levshokov et al. | α | $(5.4 \pm 2.5) \times 10^{-6}$ | $(-5.4 \pm 2.5) \times 10^{-16}$ |
| | 2004 Chand et al. | α | $(-0.6 \pm 0.6) \times 10^{-6}$ | $(0.6 \pm 0.6) \times 10^{-16}$ |
| | 2003 Murphy et al. | α | $(-0.5 \pm 0.1) \times 10^{-5}$ | $(6.4 \pm 1.4) \times 10^{-16}$ |
| | 1999 Webb et al. | α | $(-1.1 \pm 0.4) \times 10^{-5}$ | $(2.2 \pm 5.1) \times 10^{-16}$ |
| | 2006 Reinhold et al. | μ | $(2.0 \pm 0.6) \times 10^{-5}$ | $(2.0 \pm 0.6) \times 10^{-15}$ |
| | 2005 Ivanchik et al. | μ | $(1.7 \pm 0.7) \times 10^{-5}$ | $(1.7 \pm 0.7) \times 10^{-15}$ |
| | 2005 Kanekar et al. | μ | $< 1.4 \times 10^{-5}$ | $< 2.1 \times 10^{-15}$ |
| | 1995 Cowie & Songaila | μ | $(0.8 \pm 6.3) \times 10^{-4}$ | $(0.8 \pm 6.3) \times 10^{-14}$ |
| | 2007 Tzanavaris et al. | $\alpha^2 g_p \mu^{-1}$ | $(0.6 \pm 2.0) \times 10^{-5}$ | $(0.6 \pm 2.0) \times 10^{-15}$ |
| | 1995 Cowie & Songaila | $\alpha^2 g_{pp} \mu^{-1}$ | $(0.7 \pm 1.1) \times 10^{-5}$ | $(0.7 \pm 1.1) \times 10^{-15}$ |

What advantage have astrophysical methods in comparison with precise laboratory ones ?

Using quasar spectra Astrophysicists can study physical conditions which were existing Gyr ago.

aboratory measurements

We measure $\Delta\mu$ with high accuracy (!) during
days
months
a few years

e. $\Delta t = \text{day} - \text{some years}$

$$\frac{\Delta\mu}{\Delta t} = ?$$

Astrophysical measurements

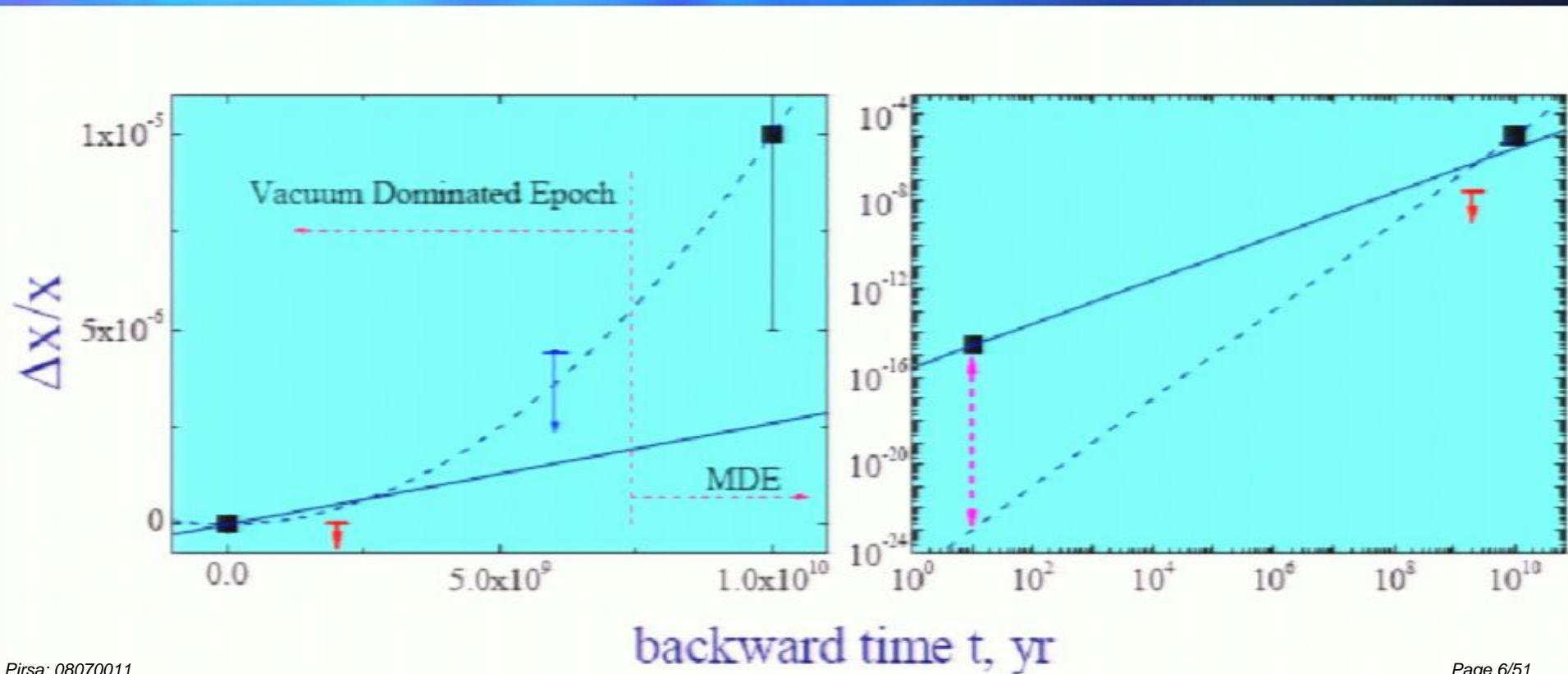
You measure $\Delta\mu$ not so precise

but

the Universe gives us

$\Delta t = 10 - 14 \text{ Gyr}$

$$\Delta\mu/\mu = f(t) = a_0 (=0) + a_1 t + a_2 t^2 + \dots$$

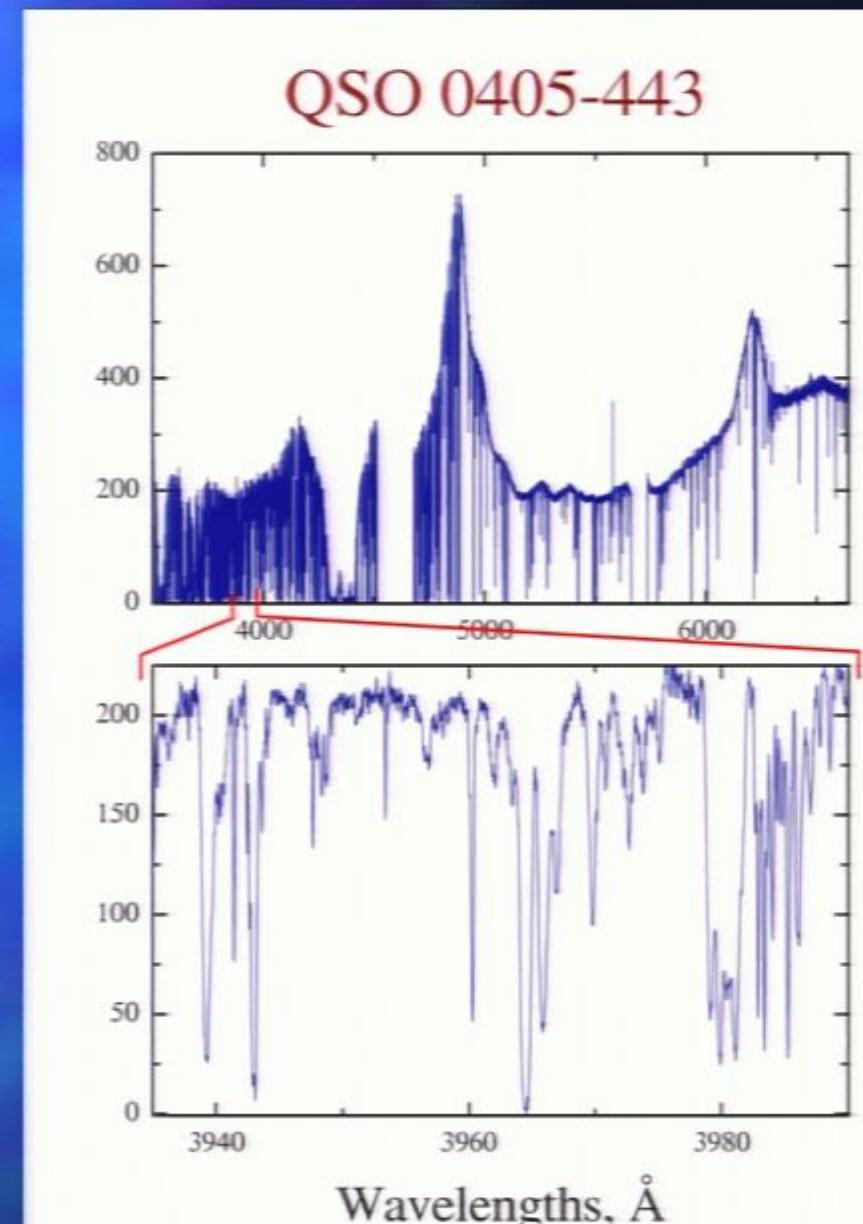


Why quasar spectra ?

Quasars are the most luminous and distant visible objects in the universe.

Therefore, the light traveling from QSO to observer bring us information about earlier epochs of the Universe (10-14 yr ago).

In studying of absorption systems in QSO spectra we obtained information about physical conditions at the epochs of the spectrum formation.



The real possibility of experimentally testing the cosmological variation of μ appeared only after the discovery of H₂ molecule clouds at high redshift (1985).

Today more than 100 000 quasars are identified and only in 17 of them H₂ absorption systems were observed because for detecting such systems we need a high-resolution spectrograph and very large optical telescopes (10m Keck or 8m VLT)

And (2+2) of 17 have H₂ absorption systems which are suitable for our analyses (number of lines is about 40 for each systems, well-determined line profile, ...)

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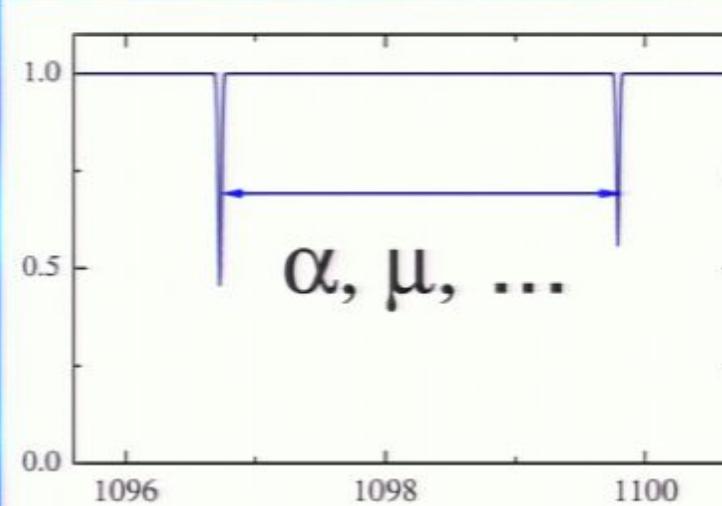
And (2+2) of 1
which are suitable
lines is about 4
determined line pr



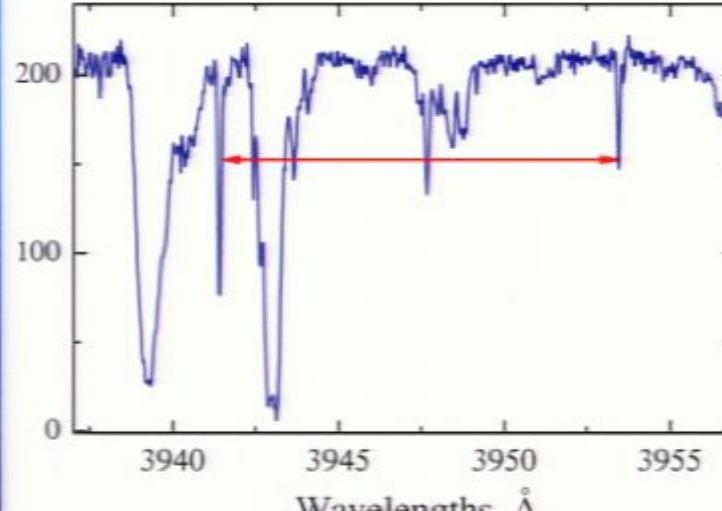
| | | |
|--------------------|-------------|-----------------------|
| Q 0515 -441 | 1.15 | HST |
| Q 1331+170 | 1.78 | HST |
| Q 0551 -336 | 1.96 | VLT |
| Q 0013 -004 | 1.97 | MMT, VLT |
| Q 1444+014 | 2.09 | VLT |
| Q 1232+082 | 2.34 | VLT |
| Q 0027 -183 | 2.40 | VLT |
| Q 2343+125 | 2.43 | VLT |
| Q 2348 -011 | 2.44 | VLT |
| Q 0405 -443 | 2.59 | VLT |
| Q 0528 -250 | 2.81 | AAT, Keck, VLT |
| Q 0347 -383 | 3.02 | VLT |
| Q 1443+272 | 4.22 | VLT |

Astrophysical methods of determination of possible fundamental constant changes are based on comparison of wavelengths measured in quasar spectra with ones measured in laboratory

"Here"
and
"Now"



"There"
and
"12 Gyr ago"



$$\lambda_{obs} = \lambda_{lab} (1 + z)$$

**So, for solving the problem of
possible cosmological variation
of $\mu = m_p/m_e$
we need to have**

1. Observed wavelengths of H_2 lines formed at the earlier epochs
2. Laboratory wavelengths of these lines
3. A function $\lambda(\alpha, \mu, \dots)$ that shows us how wavelengths depend on fundamental constants

L. New astronomical observations of H₂ lines

/ A. Ivanchik, P. Petitjean, D. Varshalovich, B. Aracil,
R. Srianand, H. Chand, C. Ledoux, and P. Boissé

Astronomy & Astrophysics 404, 2005 /

2. New laboratory measurements of H₂ lines

/ E. Reinhold, R. Buning, U. Hollenstein,
A. Ivanchik, P. Petitjean, and W. Ubachs

PRL 96, 2006 /

3. New calculations of Sensitivity Coefficients

/ V. Meshkov, A. Stolyarov,
A. Ivanchik, and D. Varshalovich

JETP Letters 83, 2006 /

1.

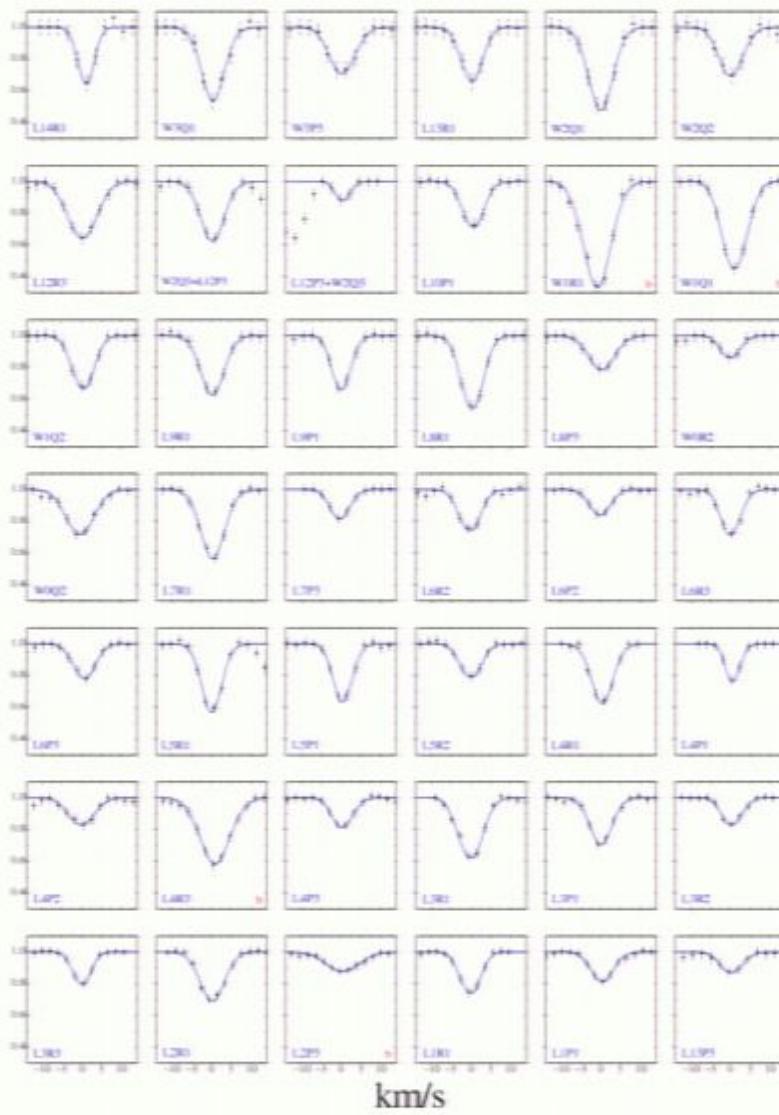
We have high-quality optical spectrum ($S/N \sim 40$, $R \sim 53\,000$)
for two quasars
(Q 0347-383 $z_{em} = 3.22$ and Q 0405-443 $z_{em} = 3.02$)
obtained by 8m-VLT/UVES of the ESO

In each of the QSO spectra there is an H_2 absorption system
at $z_{abs} = 3.0249$ and 2.5947

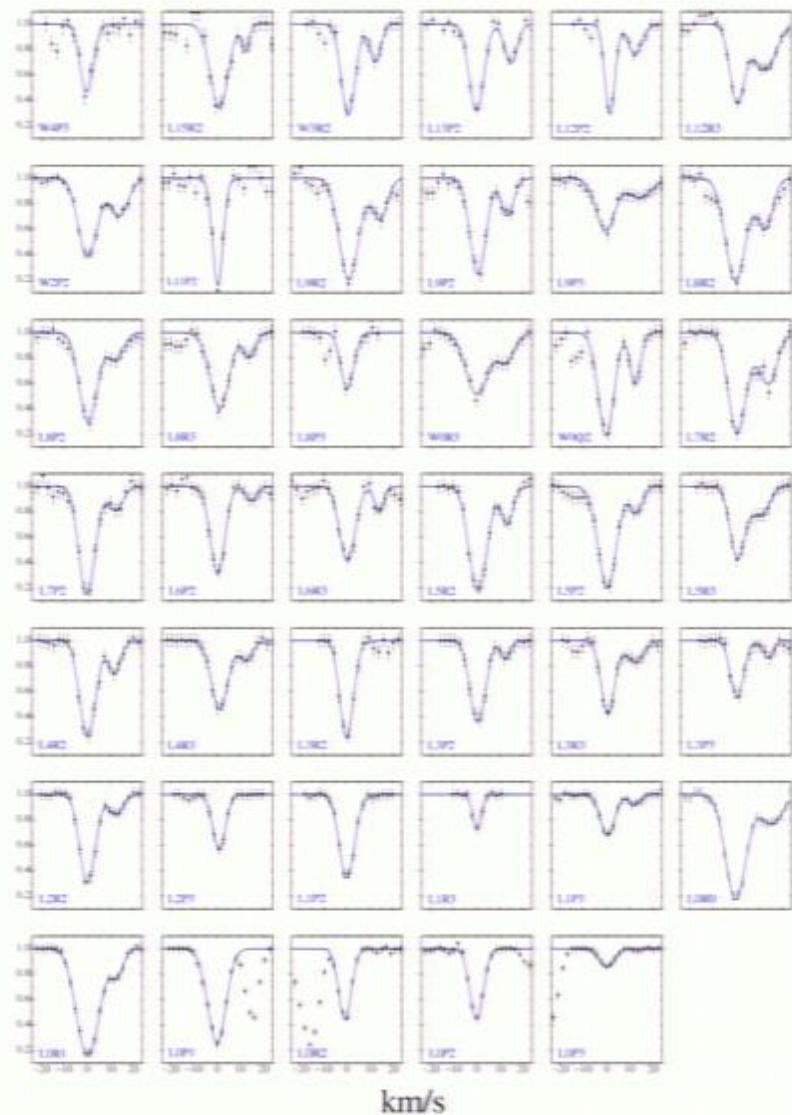
and

We have two additional spectra Q 1232+082 and HE 0027-183
with absorption systems at $z_{abs} = 2.3377$ and 2.4018

Q 0347-383



Q 0405-443

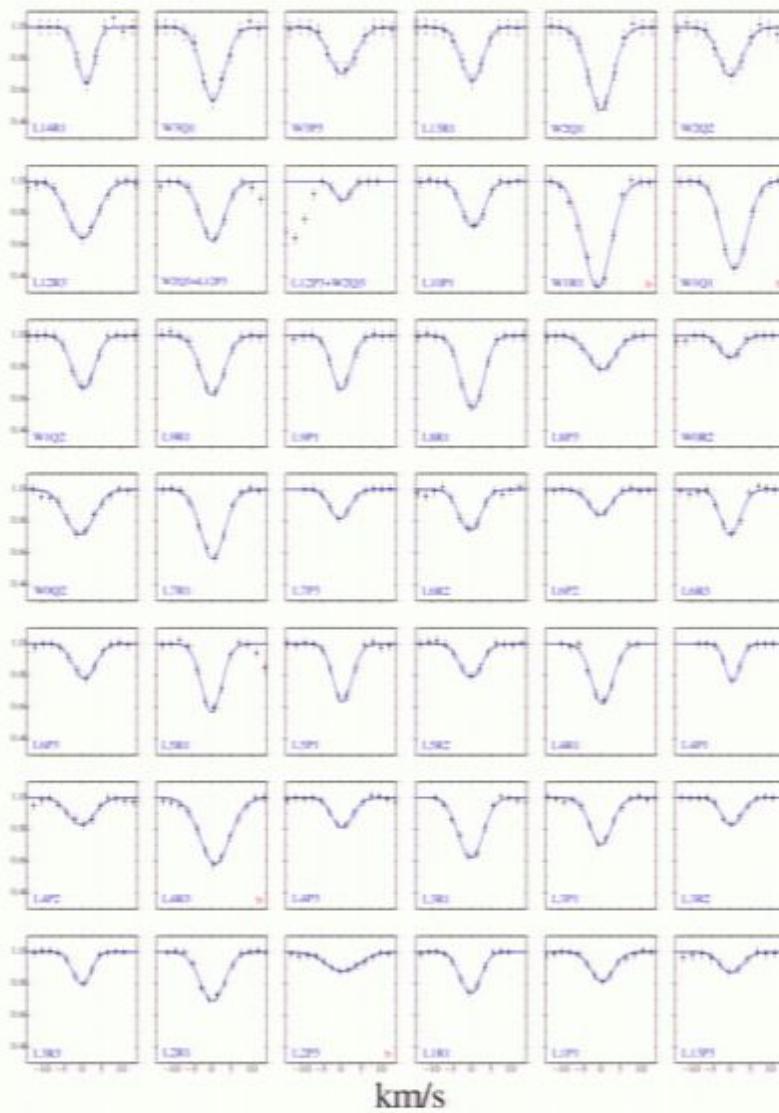


2.

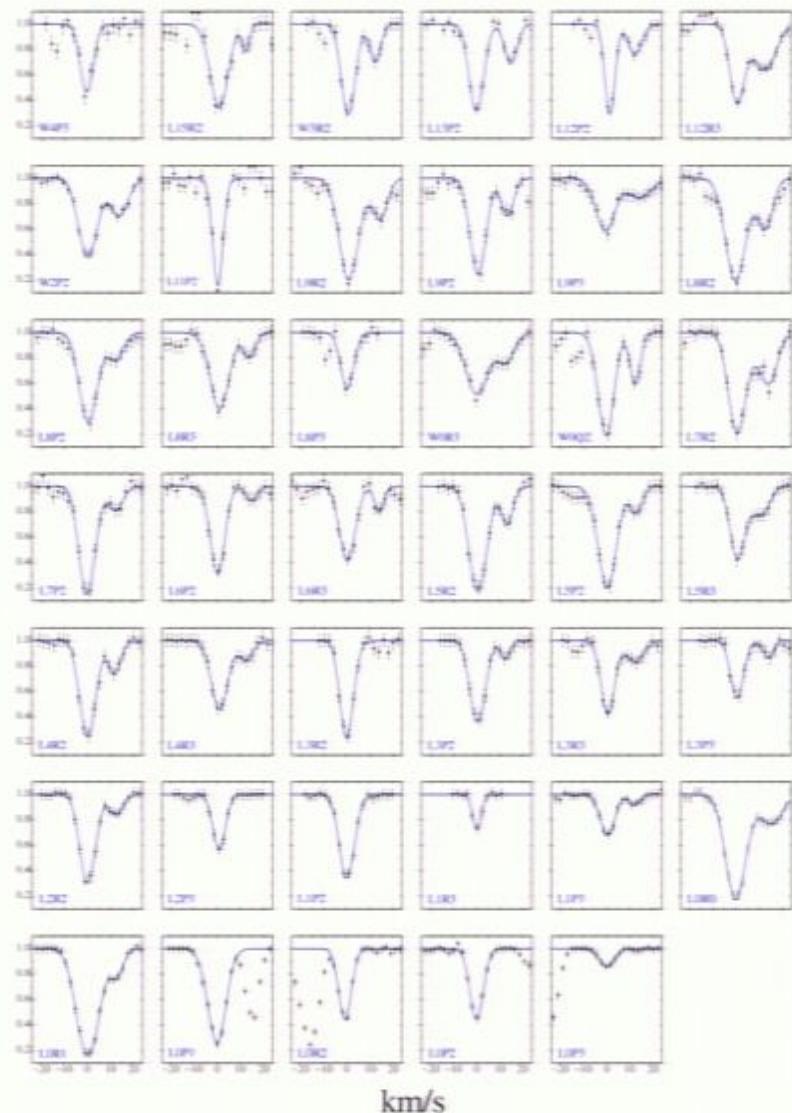
Previously we used Abgrall's atlas (1993) of H_2 laboratory wavelengths which gives errors $\sigma_\lambda \sim 1.5 \text{ m}\text{\AA}$

but now observational accuracy becomes comparable with laboratory one and we need to have more precise H_2 laboratory wavelengths.

Q 0347-383



Q 0405-443

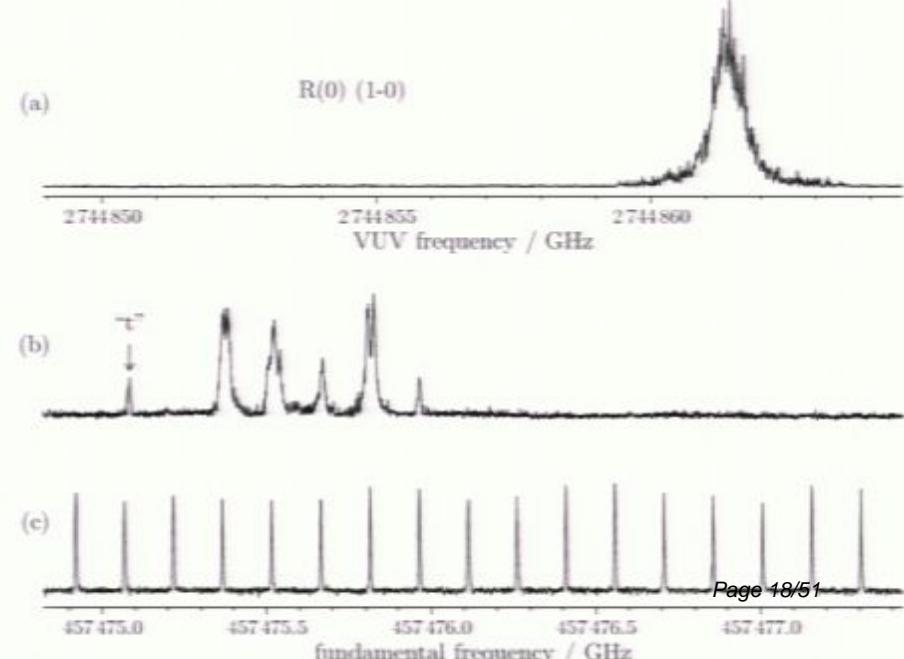


2.

Specially for our H_2 lines observed in the QSO spectra new extremely accurate wavelengths were measured using ultraviolet laser spectroscopy performed in Amsterdam)

- Philip et al., Can. J. Chem., **82**, 713, 2004
- Reinhold et al., PRL, **96**, 151101, 2006

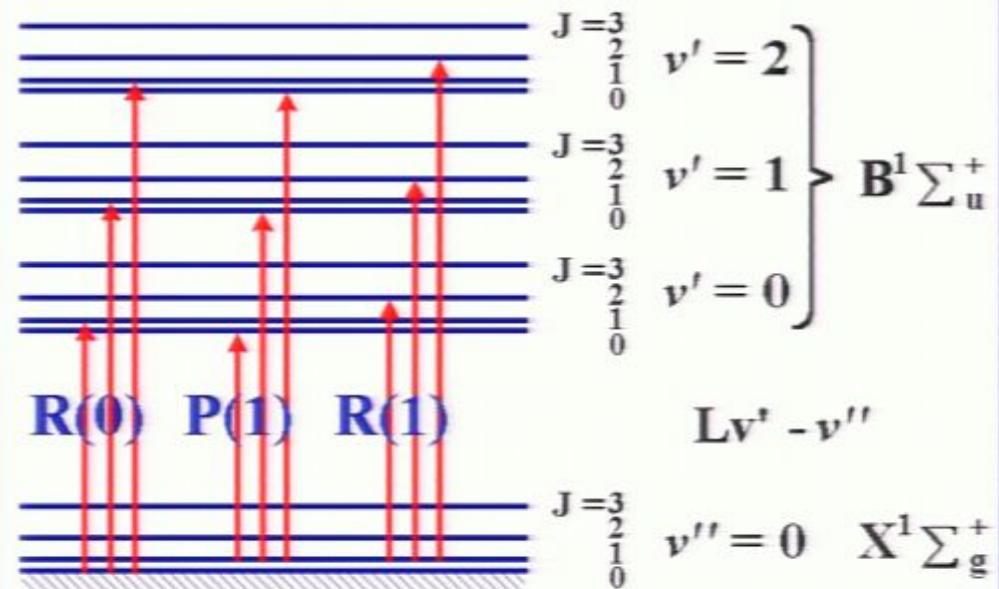
It gives errors
 $\sigma_\lambda \sim 0.07 \text{ m}\text{\AA}$
(i.e. more than
20 times better)



3. How wavelengths change if μ vary ?

$$K_i = \frac{\mu}{\lambda_i} \frac{d\lambda_i}{d\mu}$$

Sensitivity Coefficients



$$E_{\text{vib}} \sim \sqrt{m_e/m_p}, \quad E_{\text{rot}} \sim (m_e/m_p).$$

Wavelengths of electron-vibro-rotational lines depend on the reduced mass of the molecule. The dependence differs for different transitions. Thus, a measured wavelength λ_i of a line formed in an absorption system at redshift z_{abs} can be written as

$$\lambda_i = \lambda_i^0 \cdot (1 + K_i \cdot \Delta\mu/\mu) \cdot (1 + z_{\text{abs}})$$

where K_i defined as

$$K_i = \frac{\mu}{\lambda_i} \frac{d\lambda_i}{d\mu}$$

3.

In previous work we used standard adiabatic approximation with energy level represented by Dunham formula:

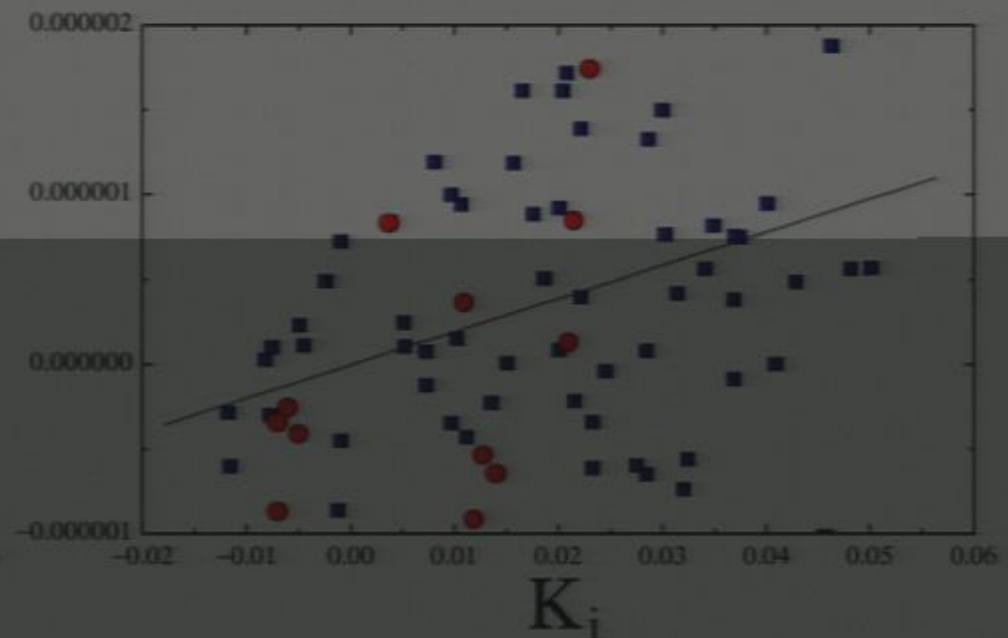
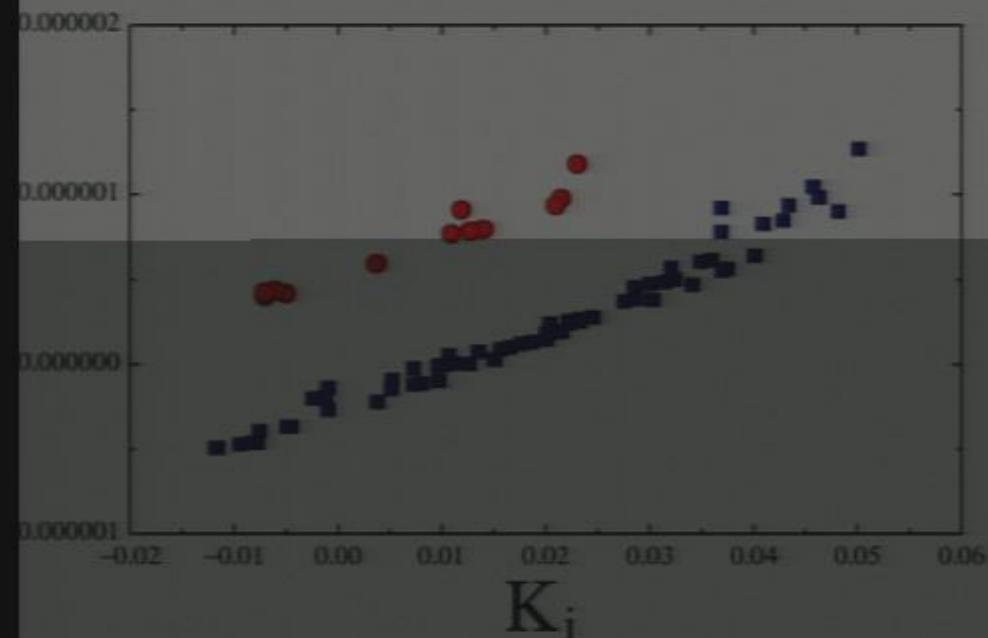
$$E(v, J) = \sum_{k,l} Y_{kl} \left(v + \frac{1}{2} \right)^k [J(J+1) - \Lambda^2]^l$$
$$Y_{kl} \propto \mu^{-l-k/2}$$

End of slide show, click to exit.

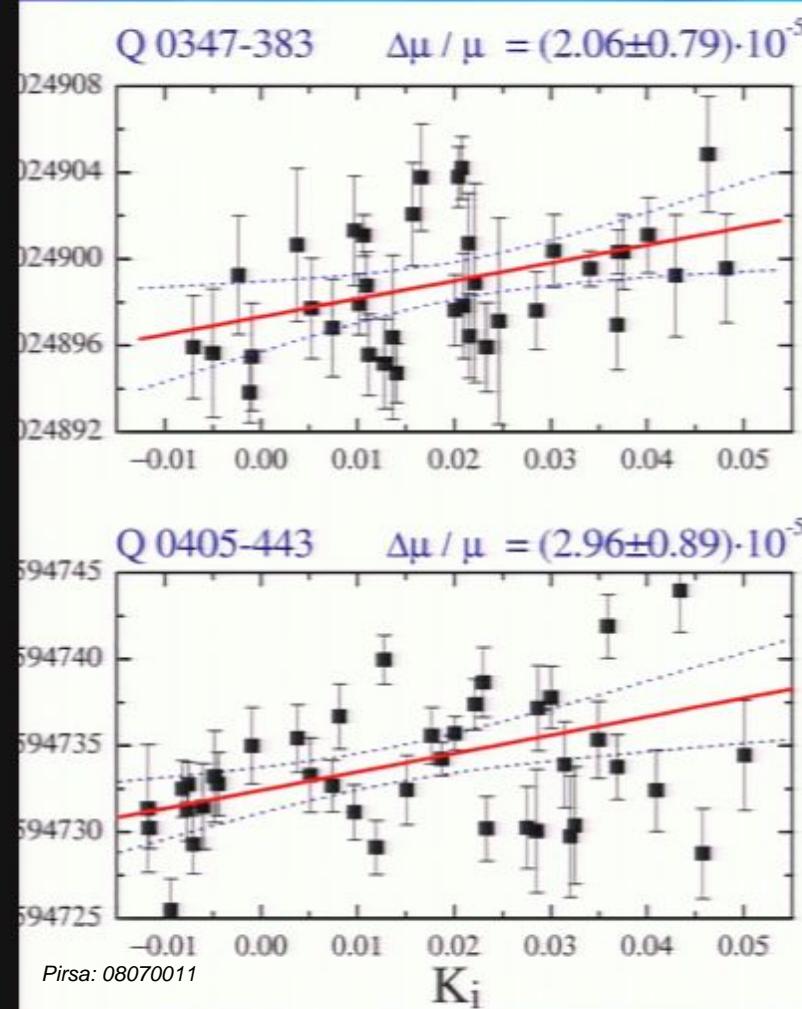
Possible systematic effects

Any effect that could produce a line center shift increasing nonotonously with wavelengths (air-vacuum wavelength conversion, Th-Ar calibration, atmospheric dispersion effect etc.)

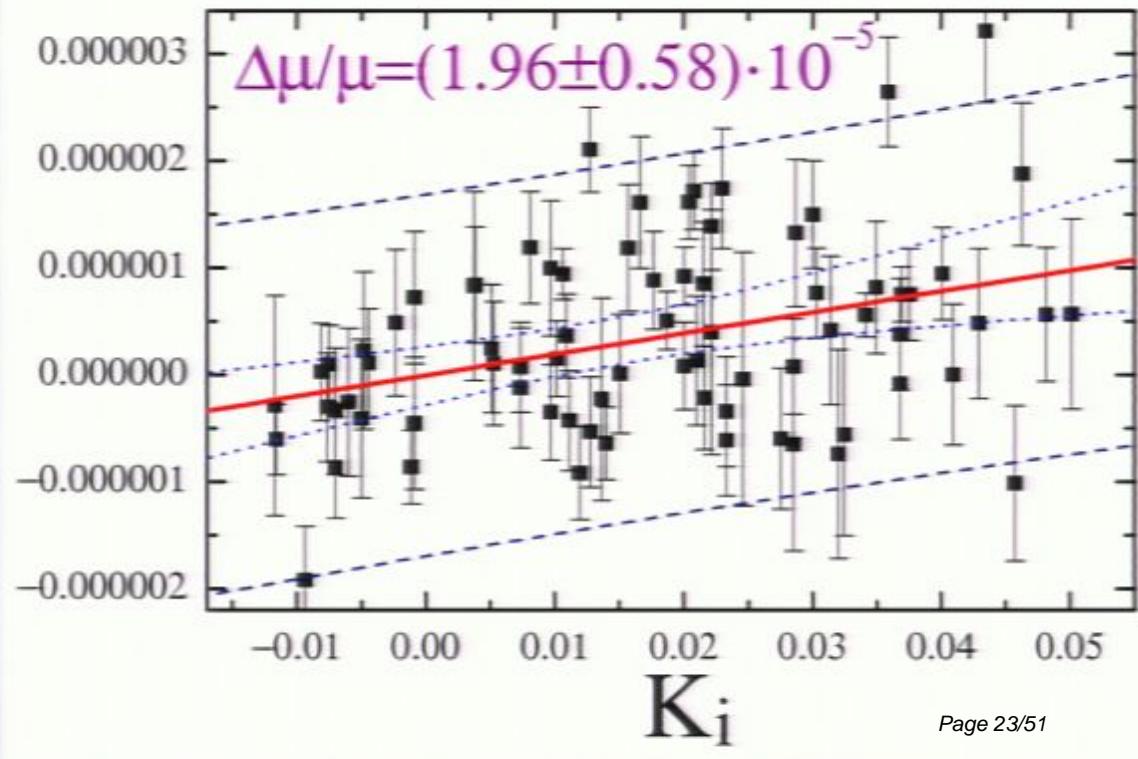
$$\frac{\lambda_i^{obs}}{\lambda_i^{lab}} = (1 + z_{abs}) \left(1 + K_i \frac{\Delta\mu}{\mu} \right)$$



Regression analysis



$$\zeta_i = \frac{z_i - z_{abs}}{1 + z_{abs}}$$



3.

Now *ab initio* nonadiabatic calculations of the wavelengths λ_i of the individual lines of the Lyman and Werner series of H_2 and corresponding sensitivity coefficients K_i (with accuracy better than 1%) have been performed.

V. Meshkov, A. Stolyarov, A. Ivanchik, and D. Varshalovich

EPL Letters 83, 2006

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So, we have all for the analyses:

$$\lambda_i^{\text{obs}}, \lambda_i^{\text{lab}}, K_i$$

$$\frac{\lambda_i^{\text{obs}}}{\lambda_i^{\text{lab}}} = (1 + z_{abs}) \left(1 + K_i \frac{\Delta \mu}{\mu} \right)$$

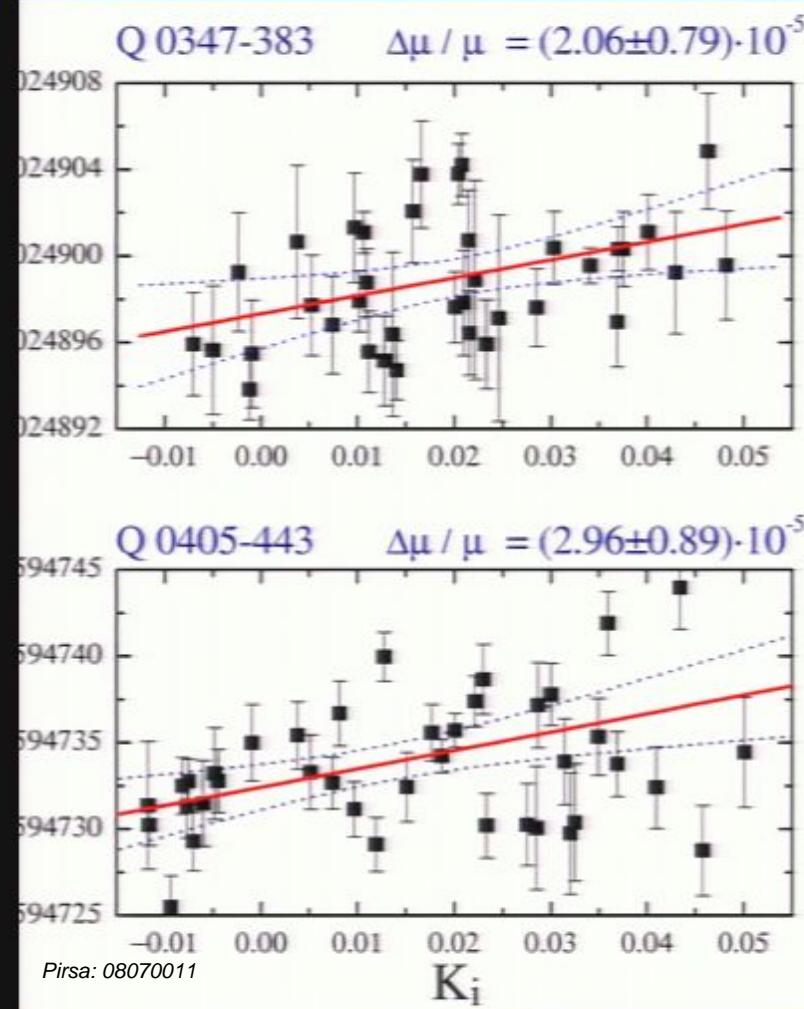
| Transition | $\lambda_i(H_2)$ | $\lambda_i(D_2)$ | $\lambda_i(T_2)$ | K_i |
|------------|------------------|------------------|------------------|----------|
| 0-0 R(1) | 1108.633 | 1103.351 | 1101.021 | -0.00719 |
| 1-0 R(0) | 1092.194 | 1091.765 | 1091.565 | -0.00055 |
| 2-0 R(0) | 1077.137 | 1080.882 | 1082.584 | +0.00508 |
| 3-0 P(1) | 1064.605 | 1071.312 | 1074.498 | +0.00865 |

So, we have all for the analyses:

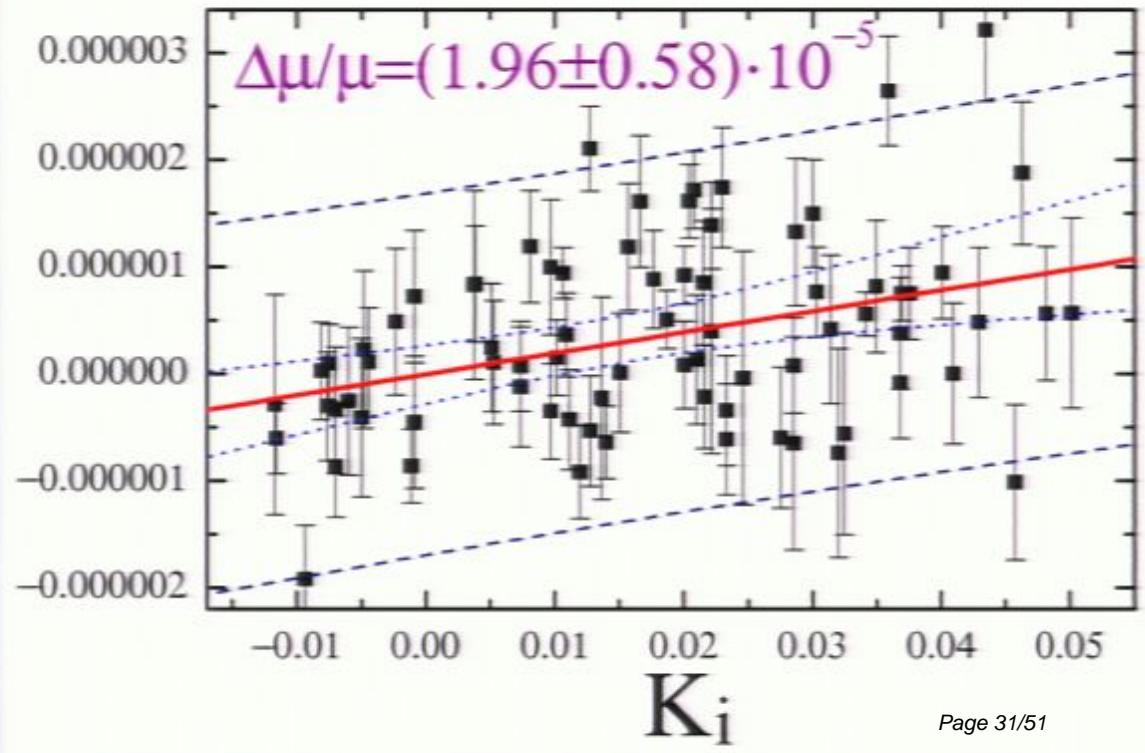
$$\lambda_i^{\text{obs}}, \quad \lambda_i^{\text{lab}}, \quad K_i$$

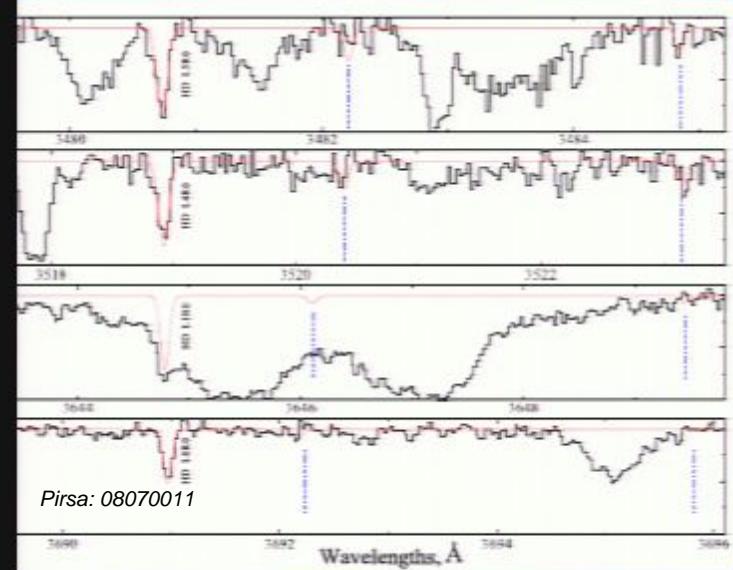
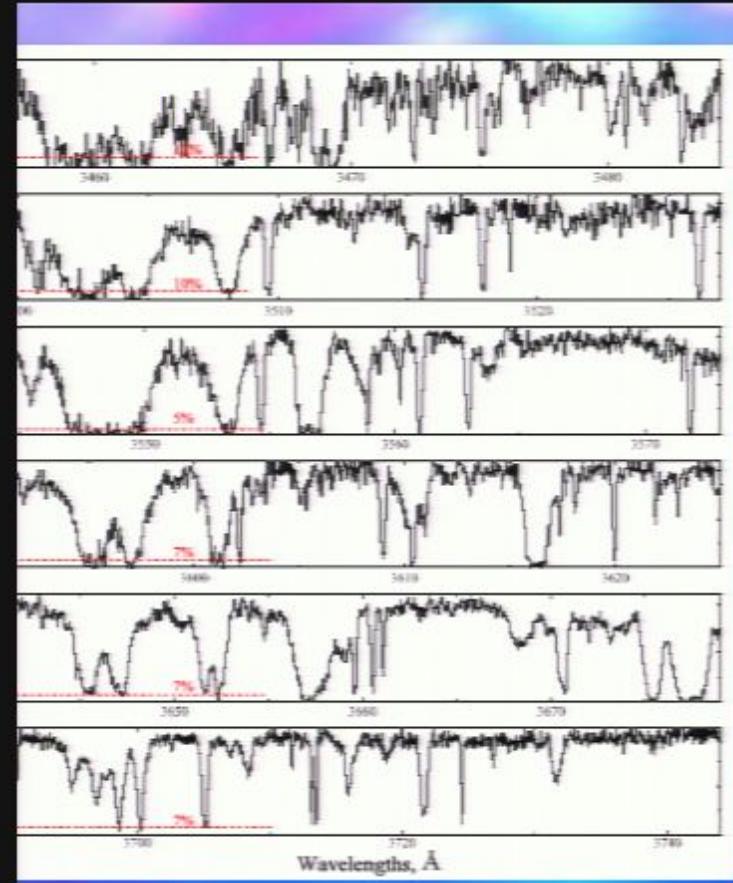
$$\frac{\lambda_i^{\text{obs}}}{\lambda_i^{\text{lab}}} = (1 + z_{abs}) \left(1 + K_i \frac{\Delta \mu}{\mu} \right)$$

Regression analysis



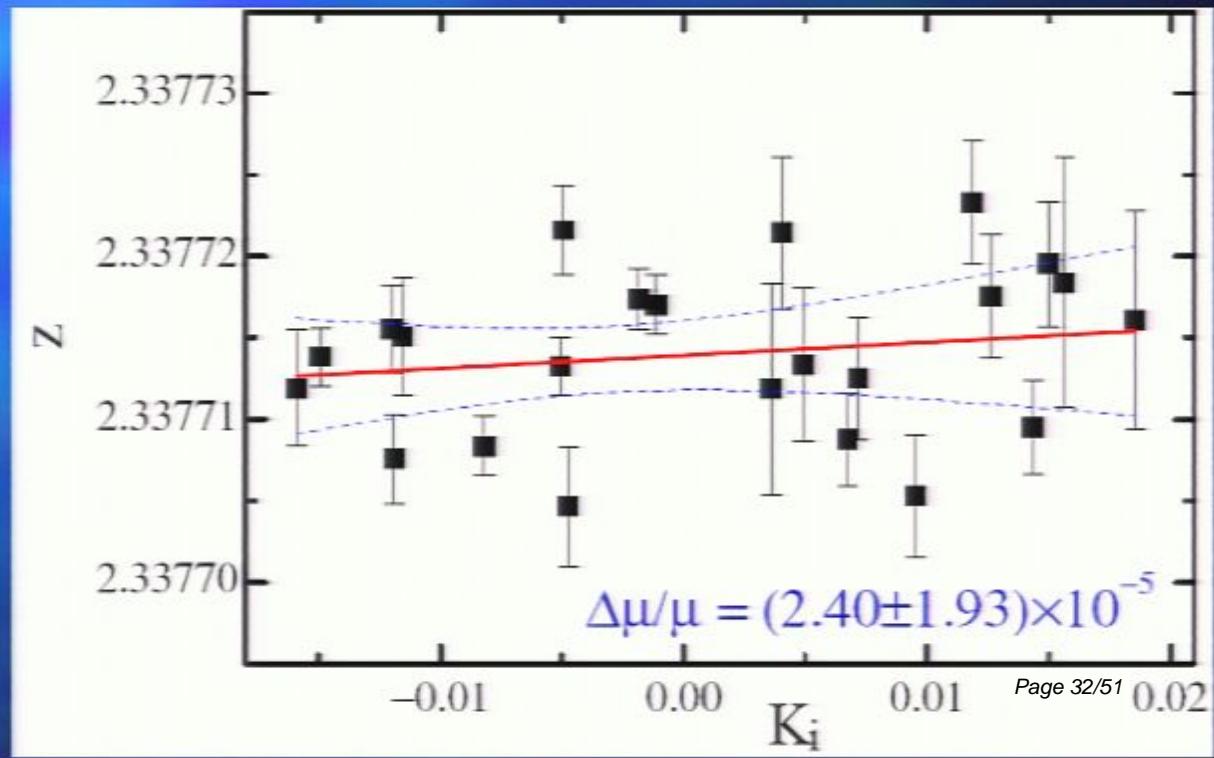
$$\zeta_i = \frac{z_i - z_{abs}}{1 + z_{abs}}$$



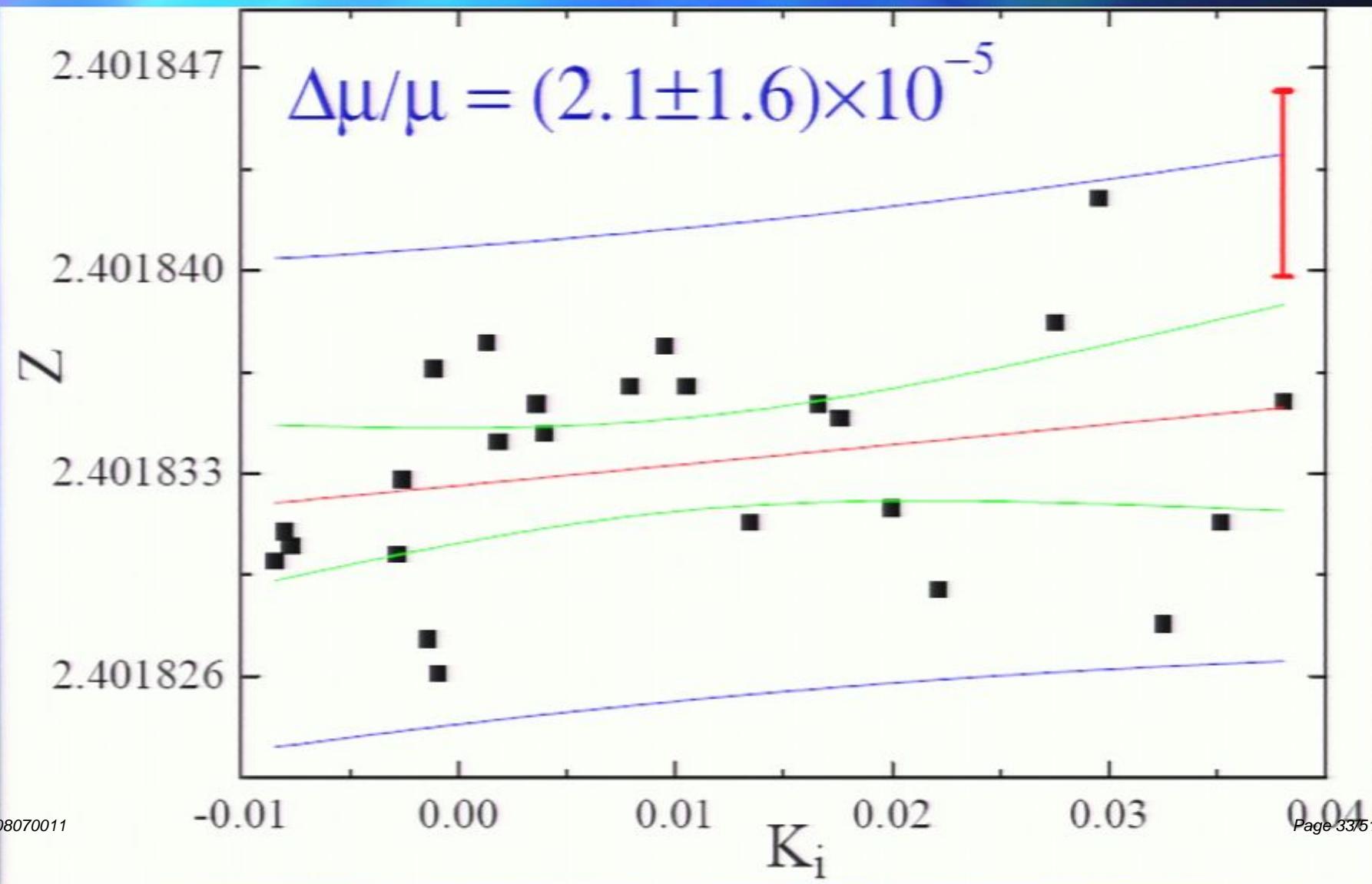


Q1232+082

H₂ & HD

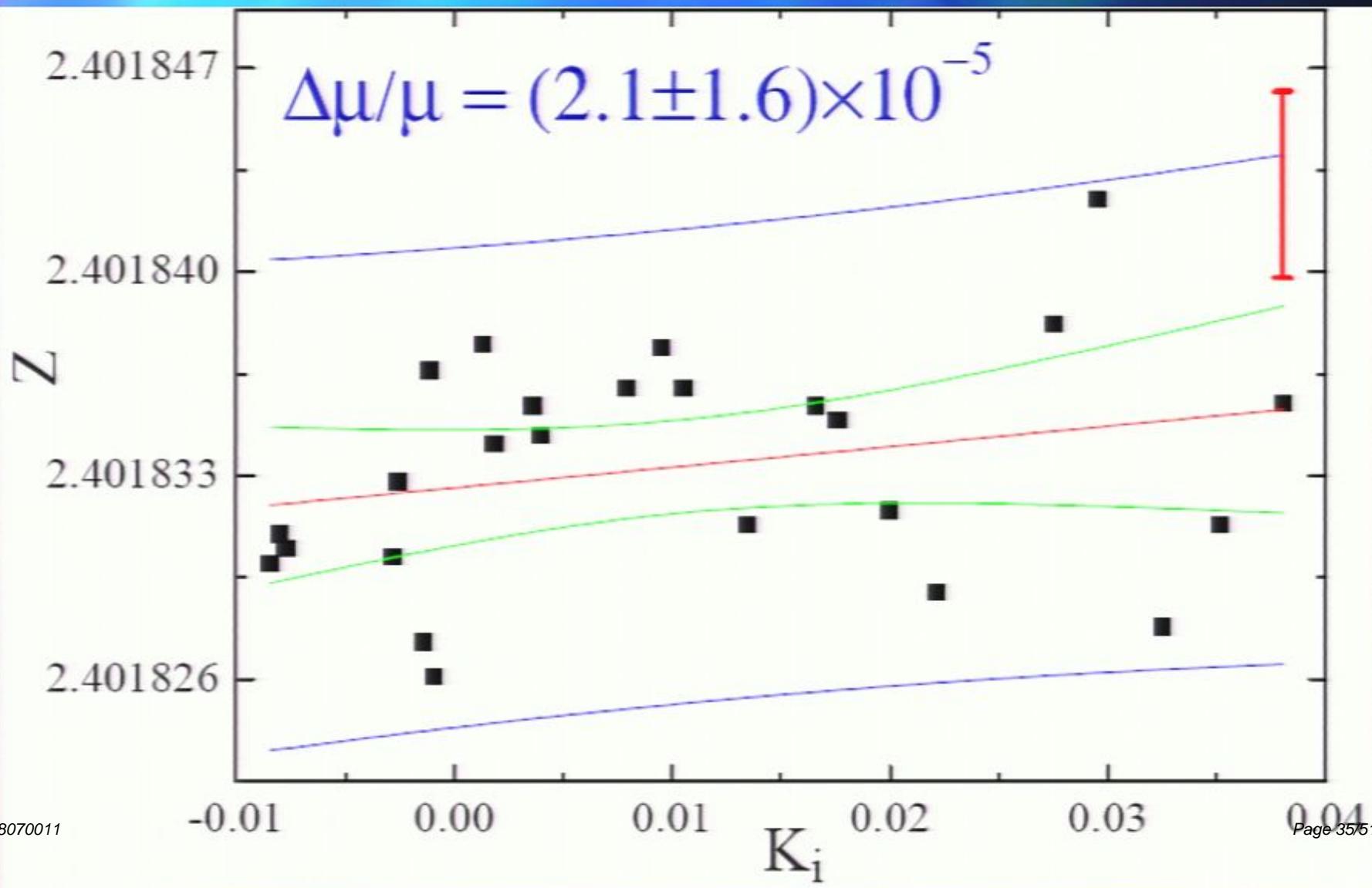


HE 0027-183



The main question:
Is it μ -variation
or
unknown systematic effect
?

HE 0027-183



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or
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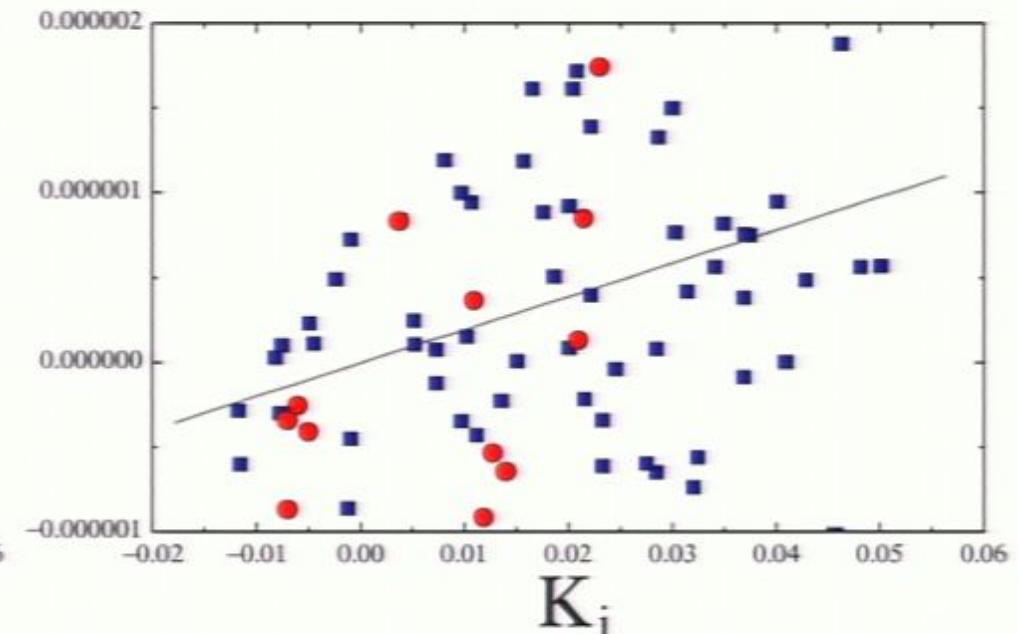
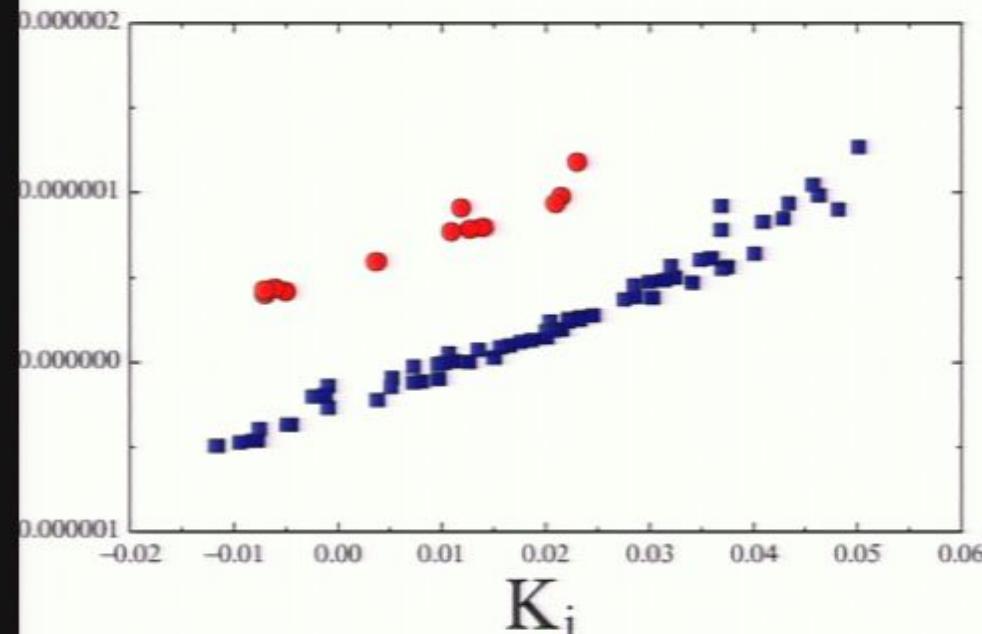
The main question:
Is it μ -variation
or
unknown systematic effect
?

Today, we do not know it
because ...

Possible systematic effects

Any effect that could produce a line center shift increasing monotonously with wavelengths (air-vacuum wavelength conversion, Th-Ar calibration, atmospheric dispersion effect etc.)

$$\frac{\lambda_i^{obs}}{\lambda_i^{lab}} = (1 + z_{abs}) \left(1 + K_i \frac{\Delta\mu}{\mu} \right)$$



*But we know what we have
to do . . .*

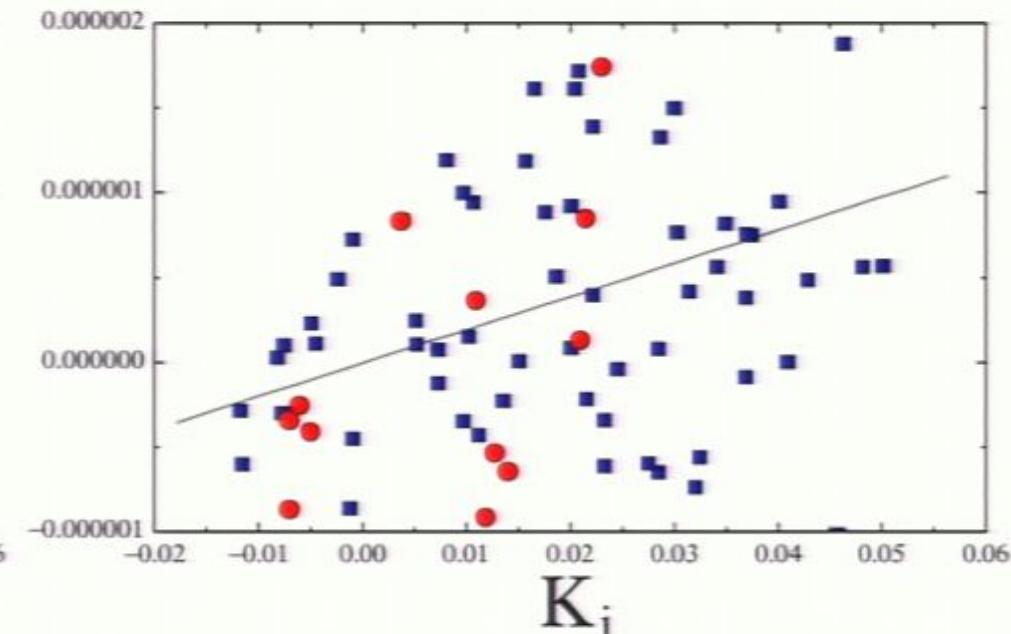
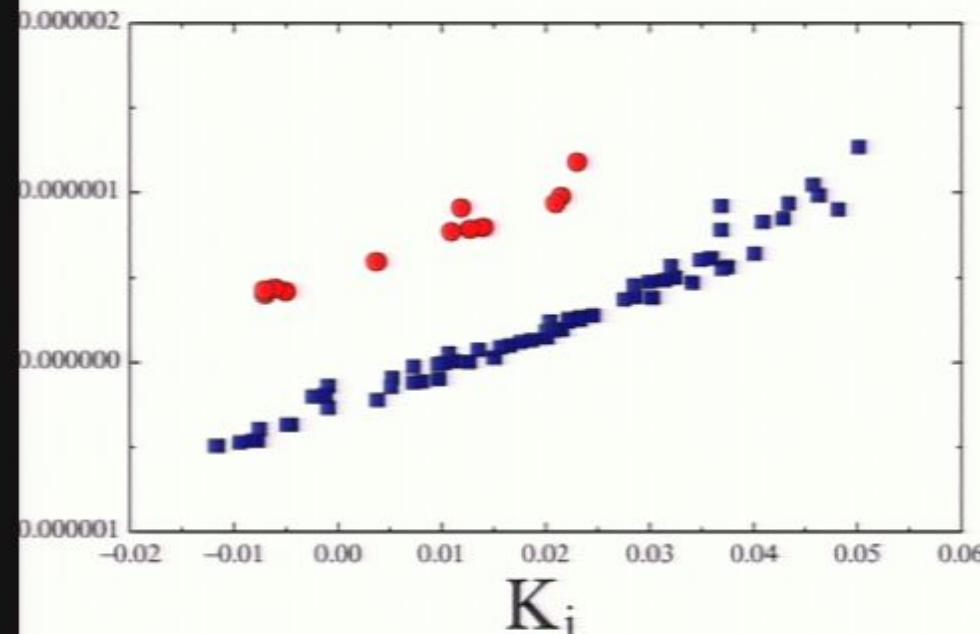
To distinguish systematic from real effect
we have to decrease the dispersion at least
two times

we hope that we will be able to answer
onto this question in near future 5-10 yr
(non-astronomical time scale 12 Gyr)

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Results

We find a correlation between z_i and K_i that could be interpreted as a variation of μ over 12 Gyr

$$\Delta\mu / \mu = (1.96 \pm 0.58) \times 10^{-5}, \quad 3\sigma$$

! But today, we can not distinguish between real μ -variation and unknown systematic effect. So, the value can be treat as the most stringent limit on possible cosmological μ -variation (at $z=3$, $t=12$ Gyr)

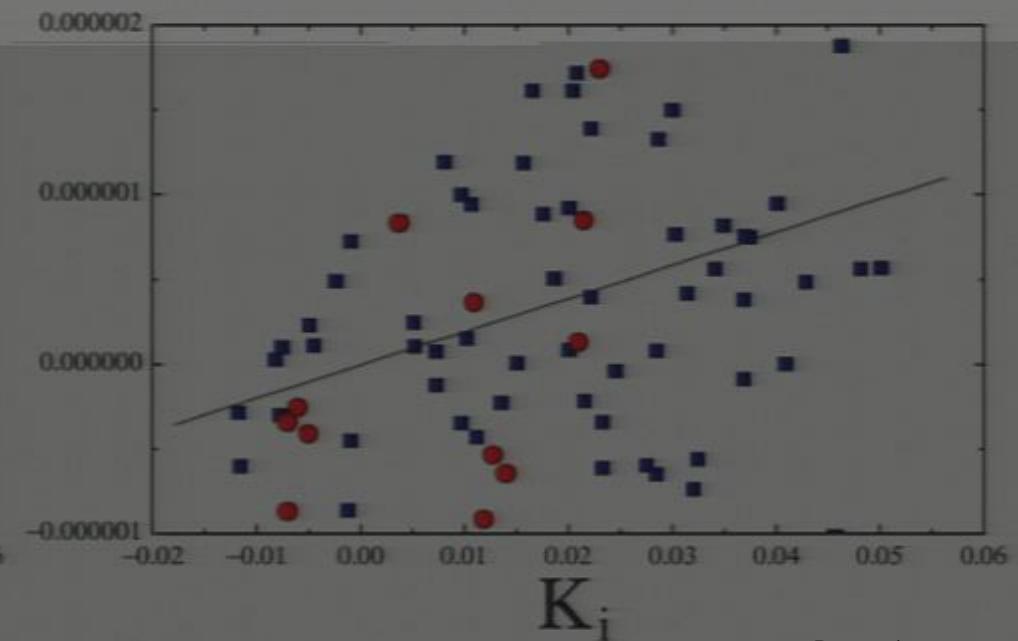
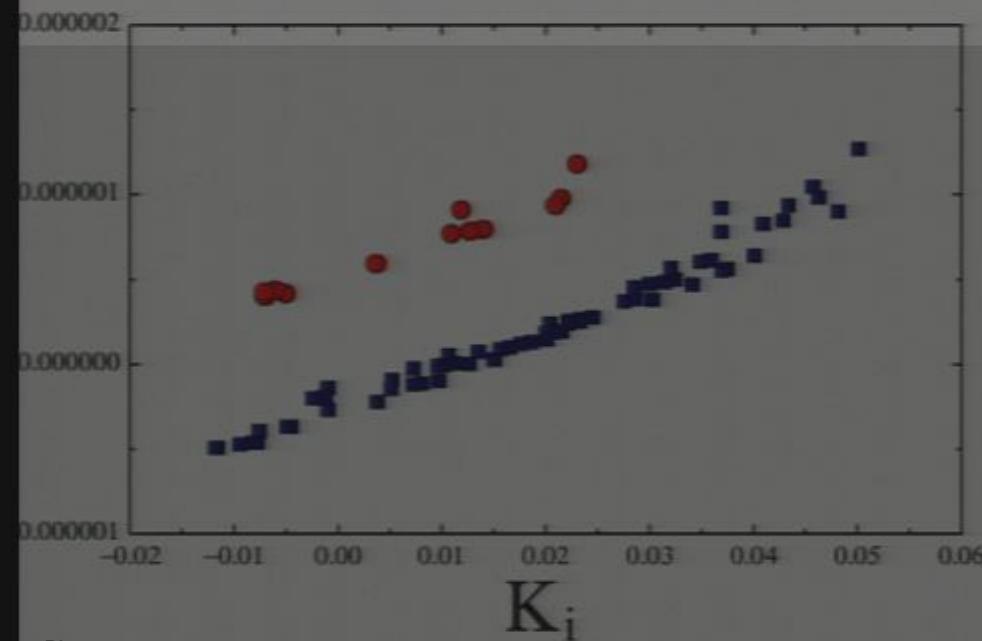
Thanks a lot
for your attention



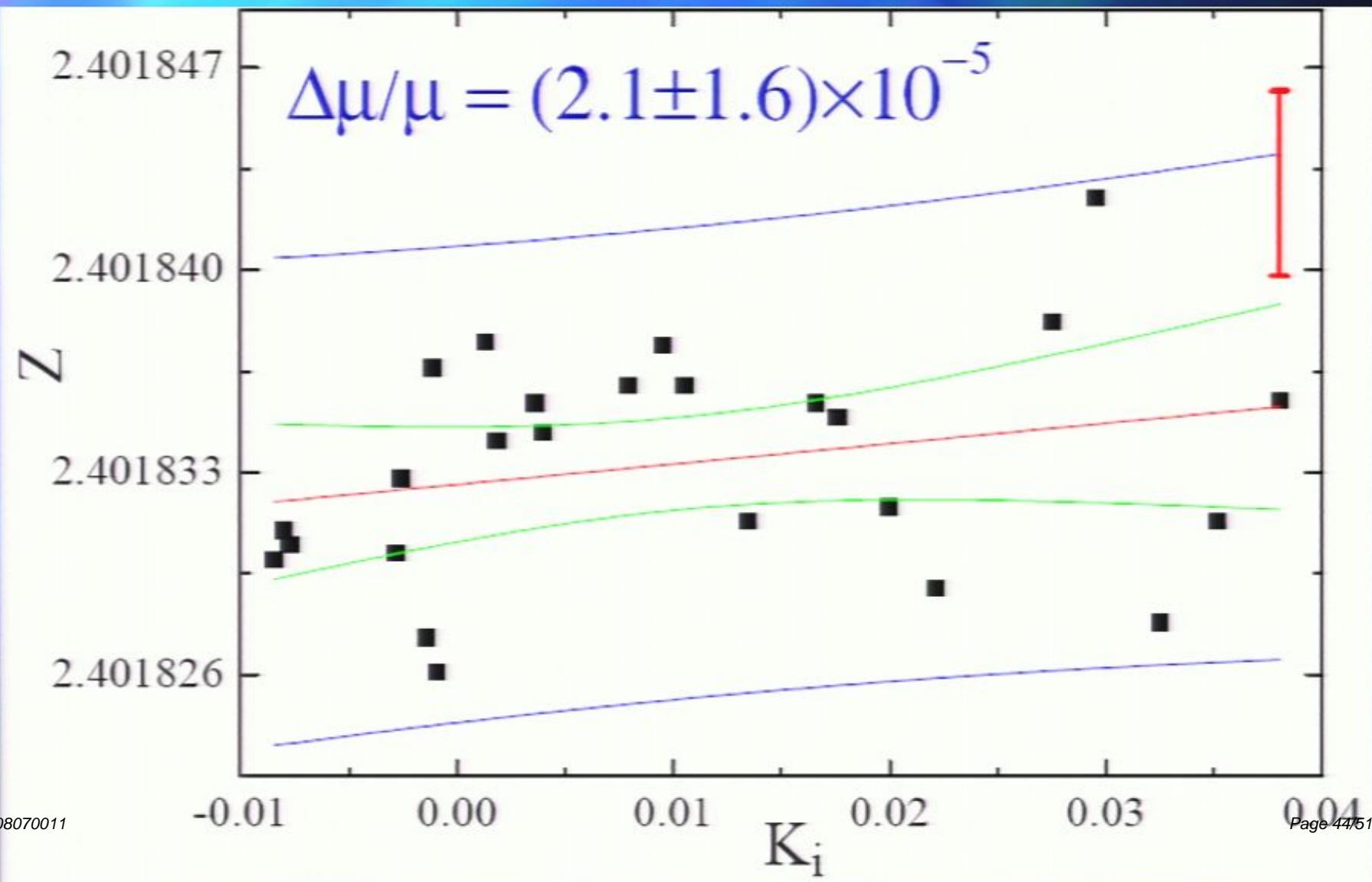
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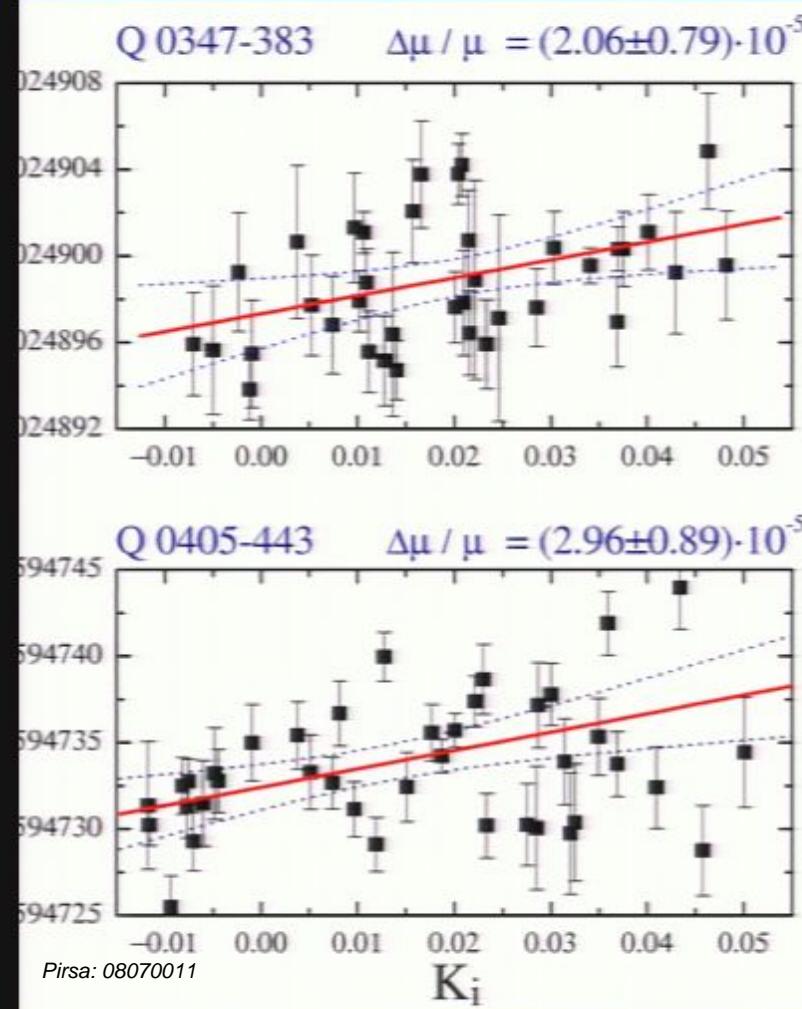
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HE 0027-183



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