

Title: SICkening Axioms for Quantum Mechanics

Date: Jun 25, 2008 02:30 PM

URL: <http://pirsa.org/08060197>

Abstract:

SIC World (1999)

"symmetric informationally complete"

For $|\psi_i\rangle \in \mathcal{H}_d$ (d -dimensional)

find d^2 projections $\pi_i = |\psi_i\rangle\langle\psi_i|$

such that

$$\text{tr } \pi_i \pi_j = \text{constant}$$

for $i \neq j$.

Caves
Żauner

SIC World (1999)

"symmetric informationally complete"

For $|\psi_i\rangle \in \mathcal{H}_d$ (d -dimensional)

find d^2 projections $\pi_i = |\psi_i\rangle\langle\psi_i|$

such that

$$\text{tr } \pi_i \pi_j = \text{constant}$$

for $i \neq j$.

Caves
Zauner

Inequivalent SIC Sets

Let $d=3$, $\omega = e^{\frac{2\pi i}{3}}$.

Set 1

$$\begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ 0 \\ -1 \end{bmatrix}$$

$$\begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 \\ \omega \\ \omega^2 \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ -\omega \\ \omega^2 \end{bmatrix}$$

$$\begin{bmatrix} -1 \\ \omega \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 \\ \omega^2 \\ \omega \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ -\omega^2 \\ \omega \end{bmatrix}$$

$$\begin{bmatrix} -1 \\ \omega^2 \\ 0 \end{bmatrix}$$

Inequivalent SIC Sets

Let $d=3$, $\omega = e^{\frac{2\pi i}{3}}$.

Set 1

$\begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$	$\begin{bmatrix} 0 \\ \omega \\ \omega^2 \end{bmatrix}$	$\begin{bmatrix} 0 \\ \omega^2 \\ \omega \end{bmatrix}$
$\begin{bmatrix} 1 \\ 0 \\ -1 \end{bmatrix}$	$\begin{bmatrix} 1 \\ \omega \\ \omega^2 \end{bmatrix}$	$\begin{bmatrix} 1 \\ \omega^2 \\ \omega \end{bmatrix}$
$\begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix}$	$\begin{bmatrix} -1 \\ \omega \\ 0 \end{bmatrix}$	$\begin{bmatrix} -1 \\ \omega^2 \\ 0 \end{bmatrix}$

Set 2

$\begin{bmatrix} -2 \\ 1 \\ 1 \end{bmatrix}$	$\begin{bmatrix} -2 \\ \omega \\ \omega^2 \end{bmatrix}$	$\begin{bmatrix} -2 \\ \omega^2 \\ \omega \end{bmatrix}$
$\begin{bmatrix} 1 \\ -2 \\ 1 \end{bmatrix}$	$\begin{bmatrix} 1 \\ -2\omega \\ \omega \end{bmatrix}$	$\begin{bmatrix} 1 \\ -2\omega^2 \\ \omega \end{bmatrix}$
$\begin{bmatrix} 1 \\ 1 \\ -2 \end{bmatrix}$	$\begin{bmatrix} 1 \\ \omega \\ -2\omega^2 \end{bmatrix}$	$\begin{bmatrix} 1 \\ \omega^2 \\ -2\omega \end{bmatrix}$

Evidence for Existence

Analytical Constructions

$$d = 2 - 13, 15, 19$$

Numerical ($\epsilon \leq 10^{-11}$)

$$d = 2 - 47$$

SIC World (1999)

"symmetric informationally complete"

For $|\psi_i\rangle \in \mathcal{H}_d$ (d -dimensional)

find d^2 projections $\pi_i = |\psi_i\rangle\langle\psi_i|$

such that

$$\text{tr } \pi_i \pi_j = \text{constant}$$

for $i \neq j$.

Caves
Zauner

Existence

- Klappenecker et al, quant-ph/0503239
On approximately symmetric informationally ...
- Klappenecker et al, quant-ph/0502031
Mutually unbiased bases are complex projective...
- Appleby, quant-ph/0412001
SIC-POVMs and the extended Clifford group
- Grassl, quant-ph/0406175
On SIC-POVMs and MUBs in Dimension 6
- Wootters, quant-ph/0406032
Quantum measurements and finite geometry
- Renes et al, quant-ph/0310075
Symmetric informationally complete quantum ...
-
-
-
- Lemmens + Seidel, J. Alg. 24, p. 494 (1973)
Equiangular Lines

Activity at + through PI

In-house: Blume-Kohout
Flammia
Fuchs

Recent Visitors: Appleby
Hughston
Larsson

Previous Visitors: Barnum
Caves
Roetteler
Scott
Wootters
:

SIC World (1999)

"symmetric informationally
complete"

For $|\psi_i\rangle \in \mathcal{H}_d$ (d -dimensional)

find d^2 projections $\pi_i = |\psi_i\rangle\langle\psi_i|$

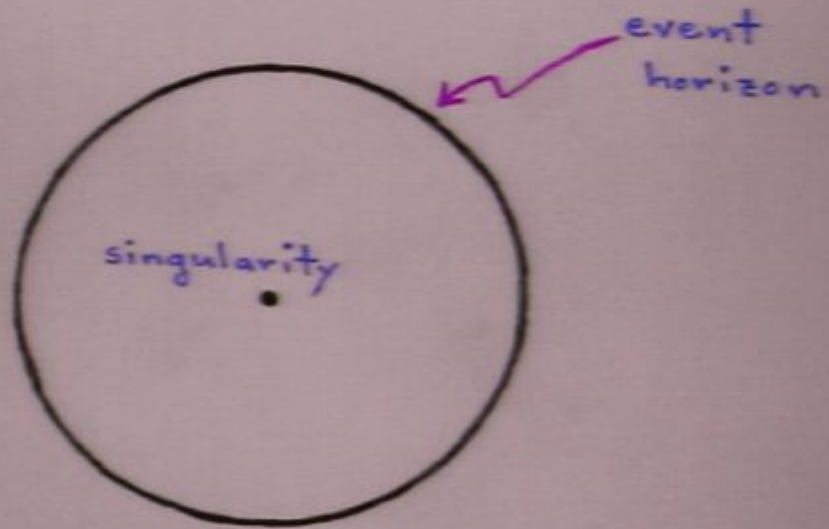
such that

$$\text{tr } \pi_i \pi_j = \text{constant}$$

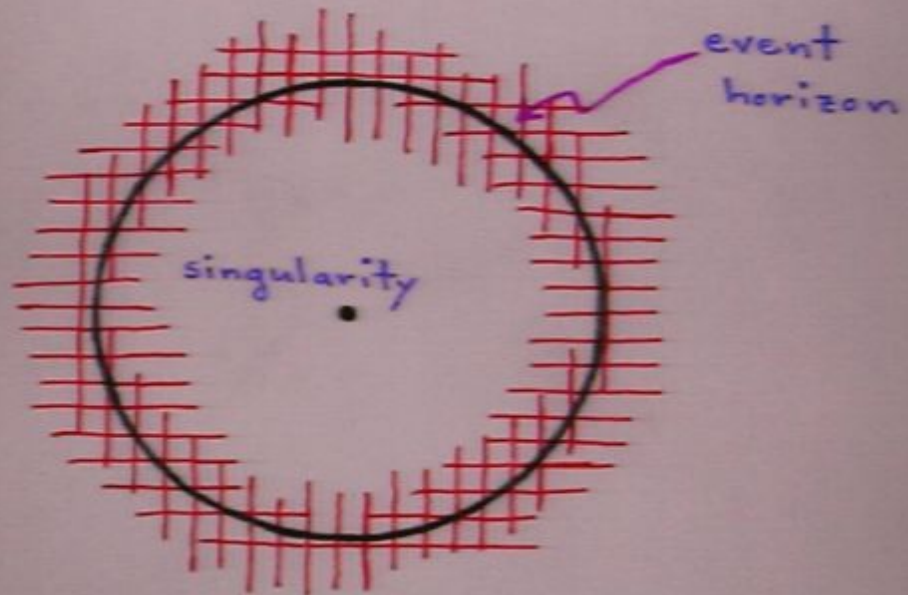
for $i \neq j$.

Caves
Zauner

Black Hole



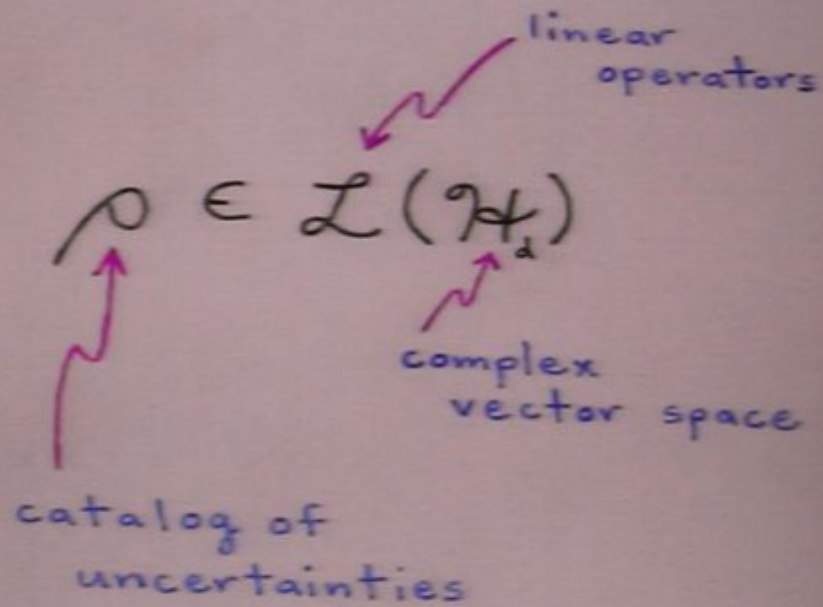
Black Hole



Eddington-
Finkelstein

124 >

Density Operators



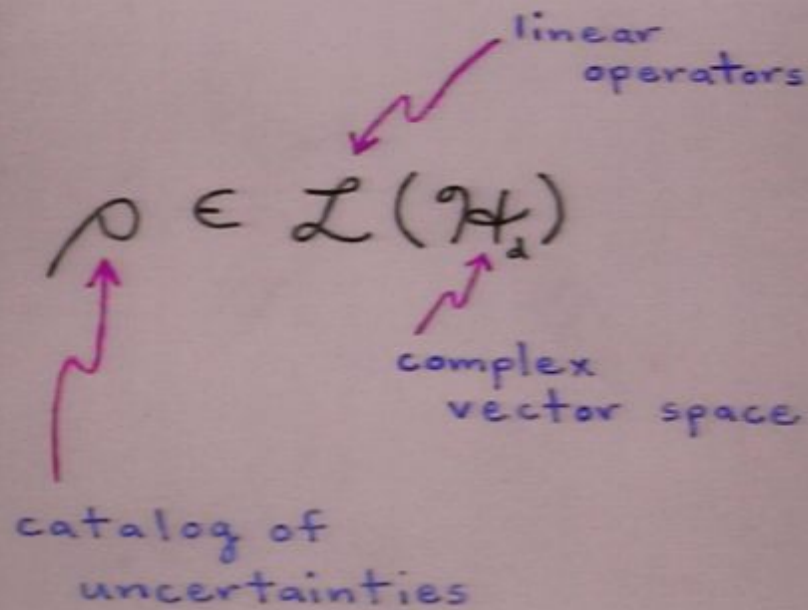
1) $\rho^\dagger = \rho$

2) $\text{tr } \rho = 1$

3) $\lambda_i(\rho) \geq 0$

convex hull of the set $\{ |\psi\rangle\langle\psi| : |\psi\rangle \in \mathcal{H}_d \}$

Density Operators



1) $\rho^\dagger = \rho$

2) $\text{tr } \rho = 1$

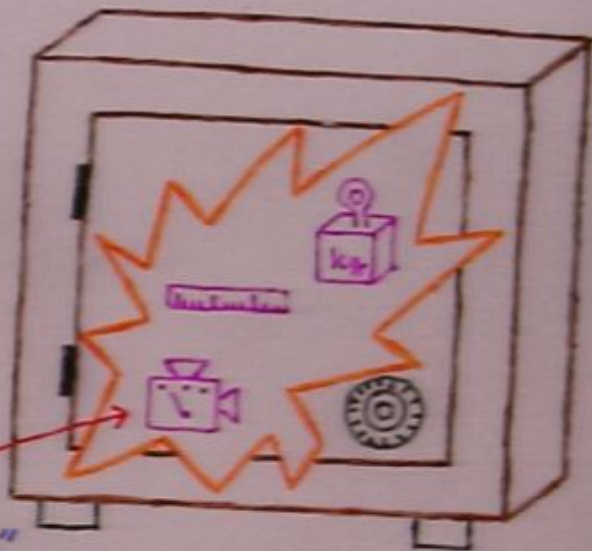
3) $\lambda_i(\rho) \geq 0$

eigenvalues

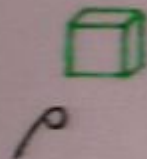
convex hull of the set $\{|\psi\rangle\langle\psi| : |\psi\rangle \in \mathcal{H}_d\}$

$\rho \longleftrightarrow \rho(h)$

Bureau of Standards



the
"standard"
quantum measurement



Quantum Probability

Given a state ρ ,
and an observable

$$H = \sum_i \alpha_i \Pi_i \quad (\text{eigendecomp}),$$

the probability of outcomes
is

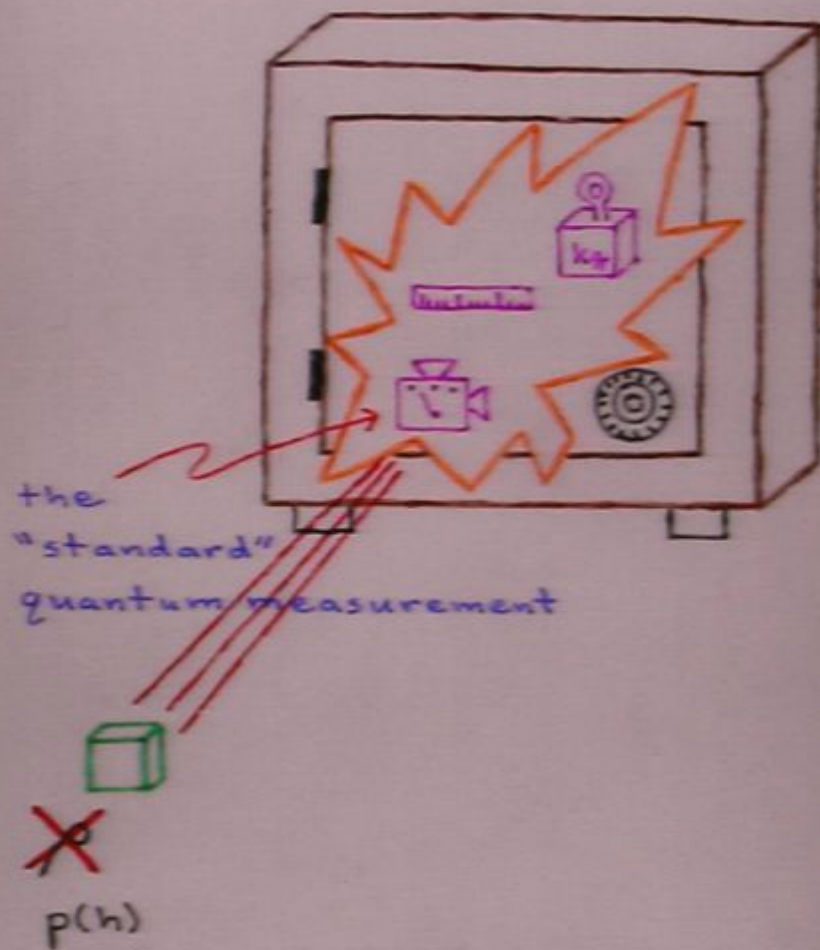
$$p_i = \text{tr} \rho \Pi_i .$$



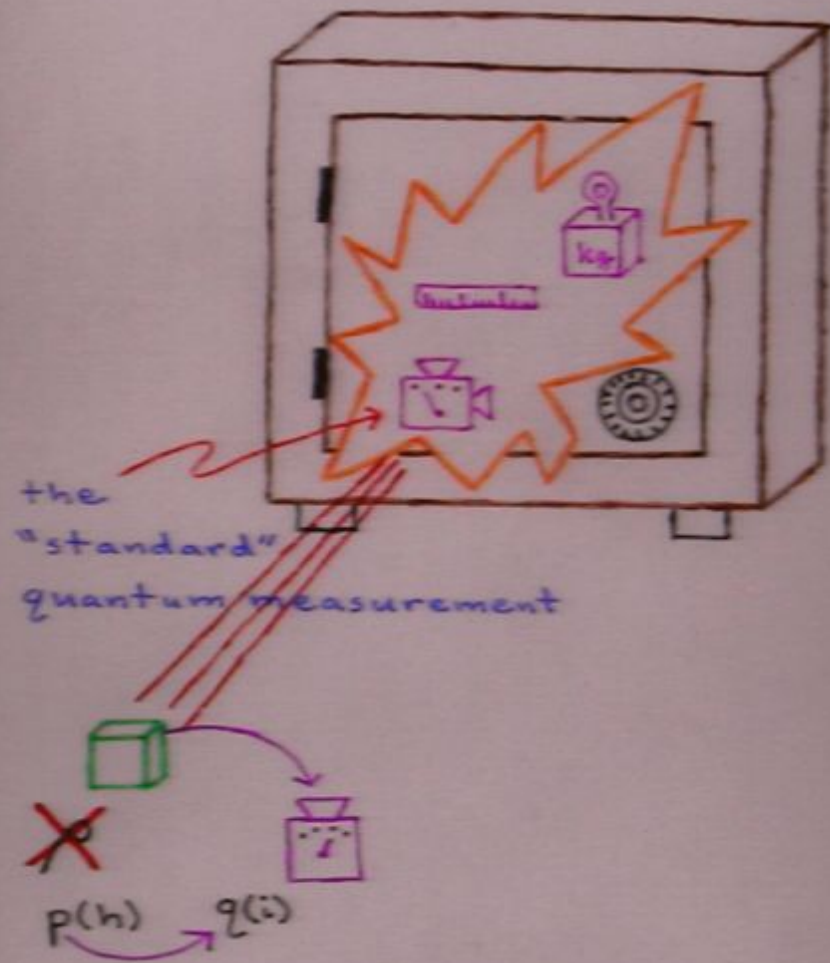
Why this rule and
not some other ?

Or simply, how to understand
it?

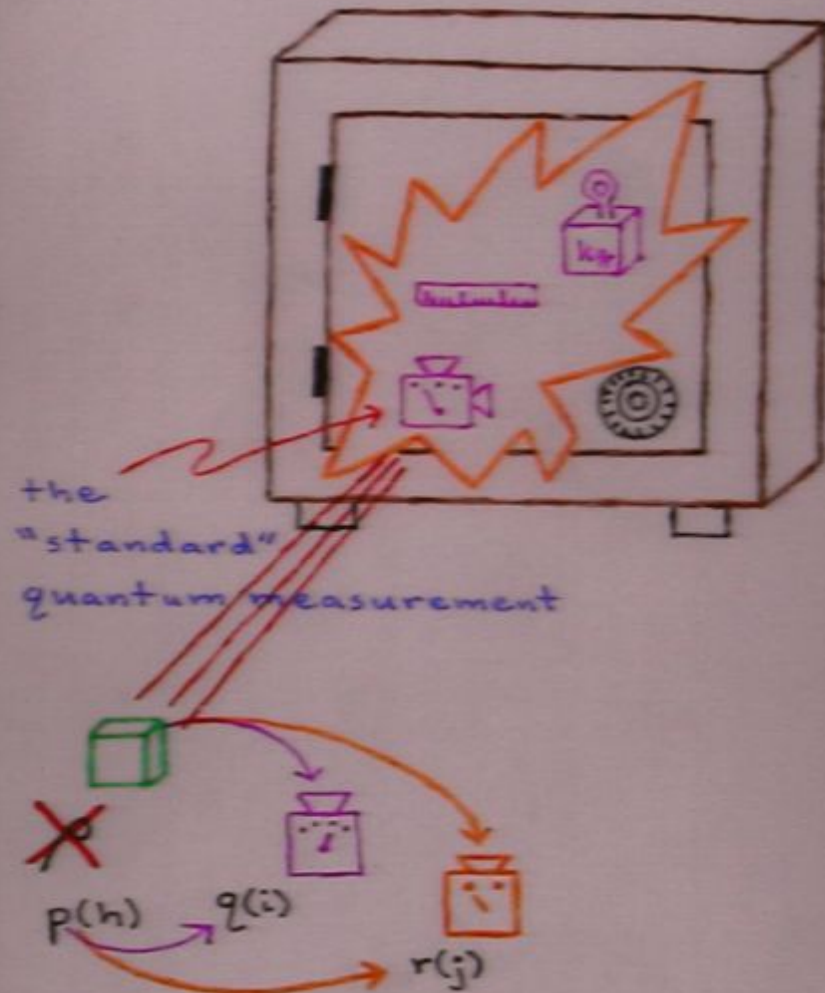
Bureau of Standards



Bureau of Standards



Bureau of Standards



Informational Completeness

quantum states

$\rho \in \mathcal{L}(\mathcal{H}_D)$ — D^2 -dimensional
vector space

Choose POVM $\{E_k\}_{k=1, \dots, D^2}$
with E_k all linearly independent.
(Can be done.)

D^2 numbers $p(k) = \text{tr} \rho E_k$ determine ρ .
↑ projection of ρ onto E_k
Success probability is an inner product

Any $\{E_k\}$ can be the
standard quantum measurement.

Informational Completeness

quantum states

$\rho \in \mathcal{L}(\mathcal{H}_D)$ — D^2 -dimensional
vector space

Choose POVM $\{E_h\}$, $h=1, \dots, D^2$,
with E_h all linearly independent.
(Can be done.)

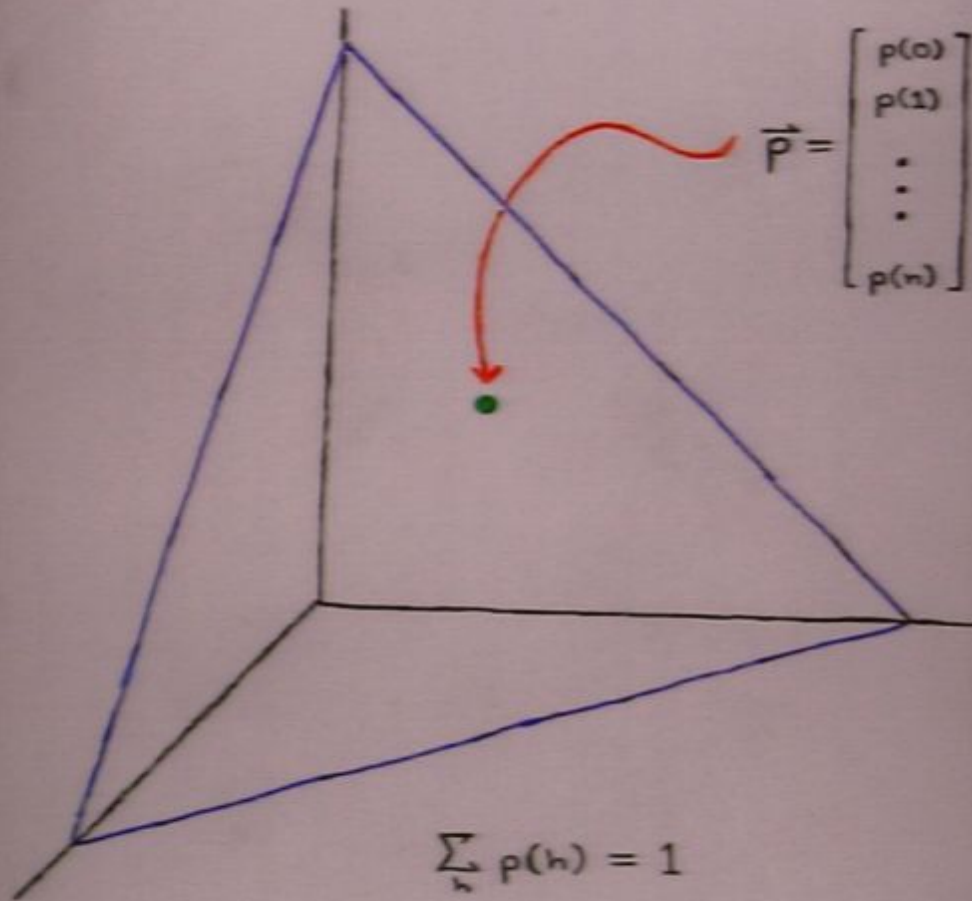
D^2 numbers $p(h) = \text{tr} \rho E_h$ determine ρ .

Because $(A, B) = \text{tr} A^* B$ is an inner product.

↑ projection of ρ onto E_h

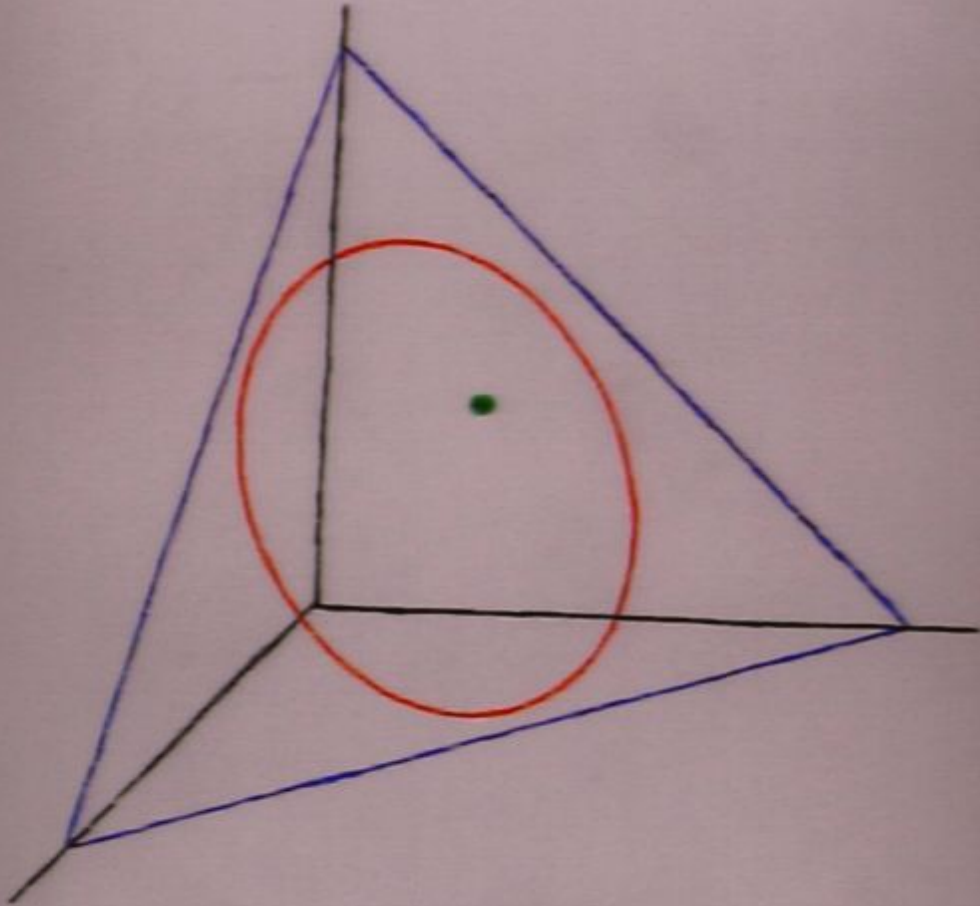
Any ^{such} $\{E_h\}$ can be the
standard quantum measurement.

Probability Simplex



$$\vec{p} = \begin{bmatrix} p(0) \\ p(1) \\ \vdots \\ p(n) \end{bmatrix}$$

$$\sum_h p(h) = 1$$
$$p(h) \geq 0 \quad \forall h$$



Path Back to Density Ops

Suppose $\{E_j\}$, $j=1, \dots, d^2$, is ICP.

Then $p(j)$ determines ρ .

But also $\rho = \sum_j \alpha_j E_j$ for some α_j 's.

Thus

$$p(j) = \text{tr} \rho E_j = \sum_k \alpha_k \text{tr} E_j E_k$$

i.e.

$$\vec{p} = M \vec{\alpha} \quad \text{where } M = [\text{tr} E_j E_k]$$

and so

↖ nonnegative matrix

$$\vec{\alpha} = M^{-1} \vec{p}$$

↙

Prettiest when $M_{jk} = a + b \delta_{jk}$.

Path Back to Density Ops

Suppose $\{E_j\}$, $j=1, \dots, d^2$, is ICP.

Then $p(j)$ determines ρ .

But also $\rho = \sum_j \alpha_j E_j$ for some α_j 's.

Thus

$$p(j) = \text{tr} \rho E_j = \sum_k \alpha_k \text{tr} E_j E_k$$

i.e.

$$\vec{p} = M \vec{\alpha} \quad \text{where } M = [\text{tr} E_j E_k]$$

and so

nonnegative matrix

$$\vec{\alpha} = M^{-1} \vec{p}$$

Prettiest when $M_{jk} = a + b \delta_{jk}$.

Example of Pretty Ones

Let $H(\mathcal{H}_d)$ be the real vector space of Hermitian operators.

Consider $d^2 - 1$ dimensional subspace $T(\mathcal{H}_d)$ of trace-free operators.

Choose $B_i \in T(\mathcal{H}_d)$, $i = 1, \dots, d^2$, so that

- 1) they lie on unit sphere
 $\text{tr } B_i^2 = 1$,
- 2) they form a regular simplex
 $\text{tr } B_i B_j = -1/(d^2 - 1)$.

Then can choose $\epsilon > 0$ so that

$$E_i = \frac{1}{2}(\mathbf{I} + \epsilon B_i)$$

are all positive semi-definite.

It follows that

- 1) E_i are ICP,
- 2) $M_{jk} = a + b \delta_{jk}$.

A Very Fundamental M.M.T.?

Caves, 1999

Suppose d^2 projectors $\Pi_i = |\psi_i\rangle\langle\psi_i|$
satisfying

$$\text{tr } \Pi_i \Pi_j = \frac{1}{d+1}, \quad i \neq j$$

exist.

Can prove:

- 1) the Π_i linearly independent
- 2) $\sum_i \frac{1}{d} \Pi_i = \mathbb{I}$

So good for Bureau of Standards.

Also

$$p(i) = \frac{1}{d} \text{tr } \rho \Pi_i$$

$$\rho = \sum_i \left[(d+1)p(i) - \frac{1}{d} \right] \Pi_i$$

Analogy

$$|\psi\rangle = \sum_{i=1}^d \alpha_i |i\rangle$$

$$\rho = \sum_{i=1}^{d^2} \alpha_i A_i$$

operators

What properties should we demand of these?

orthonormal? $(A, B) = \text{tr} A^\dagger B$

density operators themselves?

measurement outcomes?

Measure of Orthonormality

ADF, 0707.2071

Suppose A_i , $i=1, \dots, d^2$
positive semi-definite.

And $\text{tr} A_i = 1$.

"Orthonormality"

$$K = \sum_{i \neq j} (\text{tr} A_i A_j)^2$$

smaller
the better

Can prove

$$K \geq \frac{d^2(d-1)}{d+1} \quad \text{with } = \text{ iff}$$

$$\text{tr} A_i A_j = \frac{1}{d+1} \quad \forall i \neq j$$

$$A_i - \text{rank-1}$$

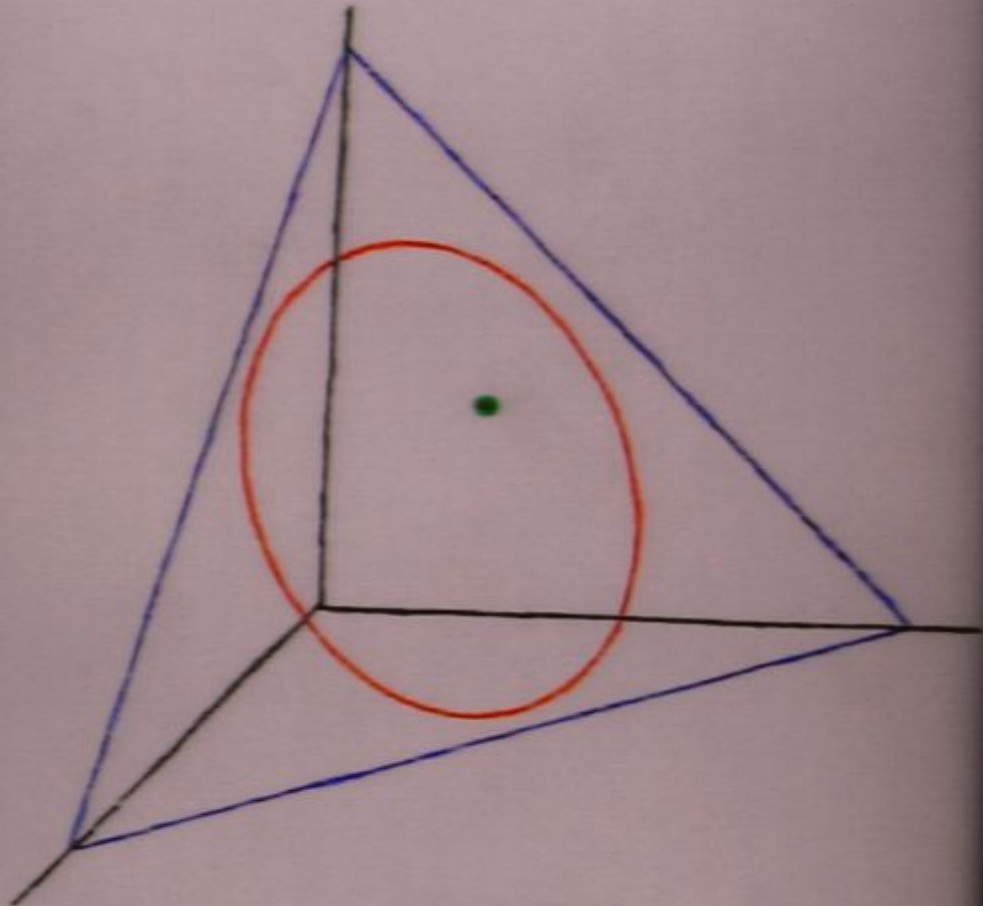
If They Exist ...

- 1) the $|\psi_i\rangle$ form a set of states maximally sensitive to eavesdropping in quantum crypto settings
CAF, quant-ph/0404122
- 2) are optimal for some natural cases of quantum tomography
A. J. Scott, quant-ph/0604049
- 3) in prime d form "minimum uncertainty" states, in analogy to coherent states, for complete sets of mutually unbiased bases
Appleby, Dang, CAF, 0707.2071

But do such sets of states

EXIST

?



Extreme Points

Characterization 1:

$$\rho = |\psi\rangle\langle\psi|$$

Characterization 2:

ρ is hermitian

$$\rho^2 = \rho$$

$$\text{tr } \rho = 1$$

Characterization 3:

ρ is positive semi-definite

$$\text{tr } \rho = 1$$

$$\text{tr } \rho^2 = 1$$

Remarkable Theorem

Jones & Linden, PRA 71 (2005)
Flammia, (unpub, 2004)

Let A be Hermitian, $A^\dagger = A$.

Then, $A = |\psi\rangle\langle\psi|$ if and only if

$$\text{tr } A^2 = \text{tr } A^3 = 1.$$

Proof:

a_i — eigenvalues of A

$$\operatorname{tr} A^2 = \sum_i a_i^2 = 1 \quad \Rightarrow \quad |a_i| \leq 1 \\ 1 - a_i \geq 0$$

$$0 = \operatorname{tr} A^2 - \operatorname{tr} A^3 = \sum_i a_i^2(1 - a_i)$$

$$\Rightarrow a_i = 0 \text{ or } 1 - a_i = 0$$

$$\operatorname{tr} A^2 = 1 \quad \Rightarrow \quad a_i = 1 \text{ for one and only one } i.$$

QED

Pure States in SIC Language

Conditions

$$\rho^\dagger = \rho \quad , \quad \text{tr } \rho^2 = \text{tr } \rho^3 = 1$$

translate to

$$\sum_i p(i)^2 = \frac{2}{d(d+1)}$$

and

$$\sum_{jkl} c_{jkl} p(j)p(k)p(l) = \frac{d+7}{(d+2)^2}$$

where

$$c_{jkl} = \text{Re } \text{tr } \Pi_j \Pi_k \Pi_l$$



Could these be independently
motivatable physical constants?

Pure States in SIC Language

Conditions

$$\rho^\dagger = \rho \quad , \quad \text{tr } \rho^2 = \text{tr } \rho^3 = 1$$

translate to

$$\sum_i p(i)^2 = \frac{2}{d(d+1)}$$

and

$$\sum_{jkl} c_{jkl} p(j)p(k)p(l) = \frac{d+7}{(d+1)^2}$$

where

$$c_{jkl} = \text{Re tr } \Pi_j \Pi_k \Pi_l$$

Could there be a more natural
motivation for this definition?

Building Hilbert Space

Given c_{ijk} , prove representation theorem!

For instance, can define a Jordan product:

$$\pi_i \bullet \pi_j \equiv \sum_k \alpha_{ijk} \pi_k$$

with

$$\alpha_{ijk} = \frac{1}{d(d+1)} \left[c_{ijk} - \frac{d\delta_{ij} + 1}{(d+1)^2} \right] .$$

Unitarity

$$\begin{array}{ccc} \rho & \longrightarrow & U\rho U^\dagger \\ \downarrow & & \downarrow \\ p(i) & \longrightarrow & q(j) \end{array}$$

$$q(j) = \sum_i \langle j | U | i \rangle \rho_{ii} \langle i | U^\dagger | j \rangle$$

unitary transformation

Then

$$q(j) = \sum_i \langle j | U | i \rangle \rho_{ii} \langle i | U^\dagger | j \rangle$$

could maybe be simpler

Unitarity

$$\begin{array}{ccc} \rho & \longrightarrow & U\rho U^\dagger \\ \downarrow & & \downarrow \\ p(i) & \longrightarrow & q(j) \end{array}$$

Unitarity

$$\begin{array}{ccc} \rho & \longrightarrow & U\rho U^\dagger \\ \downarrow & & \downarrow \\ p(i) & \longrightarrow & q(j) \end{array}$$

Define $t(j|i) = \frac{1}{d} \text{tr} U \pi_i U^\dagger \pi_j$

\uparrow
doubly stochastic
matrix

Then

$$q(j) = (d+1) \sum_i p(i) t(j|i) - \frac{1}{d}$$

Could hardly be simpler.

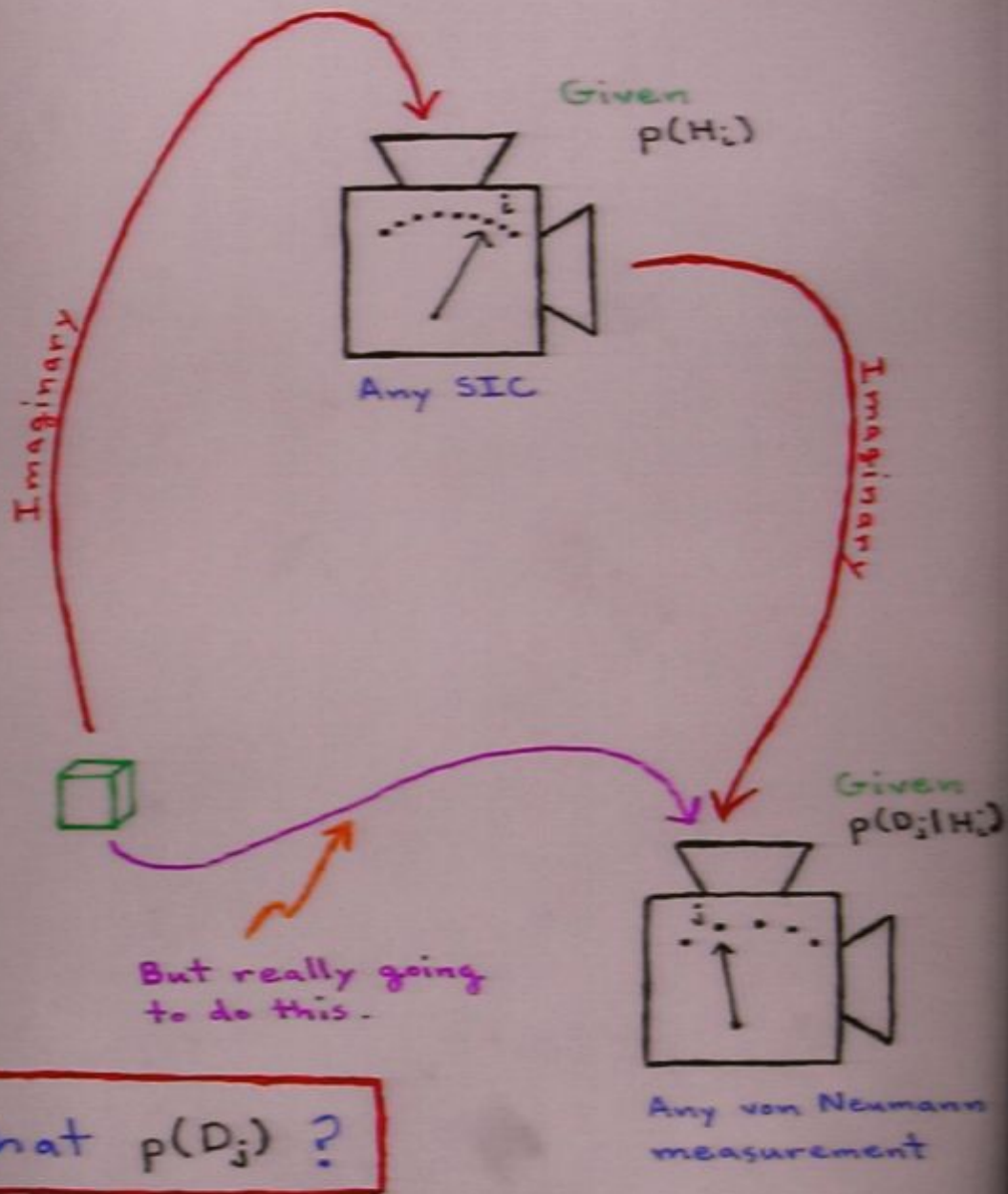
Role of the Born Rule

quantum state ρ
measurement $\{E_n\}$

$$q(h) = \text{tr } \rho E_n$$

To transform or relate
probabilities!

$$q(h) = \sum_i [(d+1)\rho(i) - \frac{1}{d}] \text{tr } \pi_i E_n$$



Laws of Probability

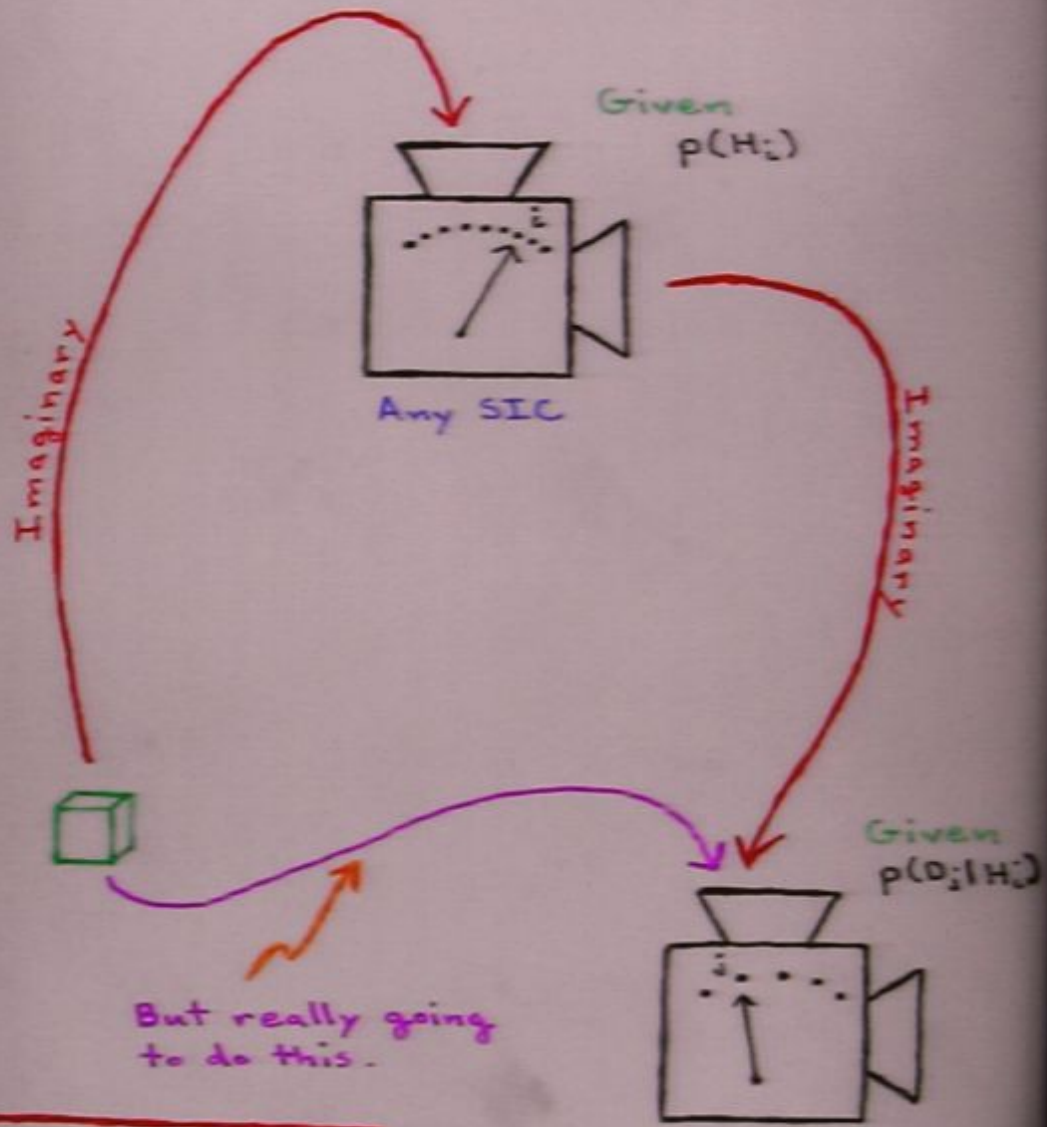
H_i - various hypotheses one might have

D_j - data values one might gather

Given: $p(D_j | H_i)$ ← expectations for data given hypothesis
 $p(H_i)$ ← expectations for hypotheses themselves

Question: What expectations should one have for the D_j ?

Answer: $P(D_j) = \sum_i p(H_i) p(D_j | H_i)$



What $p(D_j)$?

$$p(D_j) = (d+1) \sum_i p(H_i) p(D_j | H_i) - 1$$

Quantum (under $p(D_j)$)

(Usual) Bayesian (under $\sum_i p(H_i) p(D_j | H_i)$)

Magic!

"Measurement"

Does it reveal a pre-existing,
but unknown, value?

or

Does it in some sense go toward
creating the very value?

Laws of Probability

H_i - various hypotheses one might have

D_j - data values one might gather

Given: $p(D_j | H_i)$ ← expectations for data given hypothesis
 $p(H_i)$ ← expectations for hypotheses themselves

Question: What expectations should one have for the D_j ?

Answer: $P(D_j) = \sum_i p(H_i) p(D_j | H_i)$

$$p(D_j) = (d+1) \sum_i p(H_i) p(D_j | H_i) - 1$$

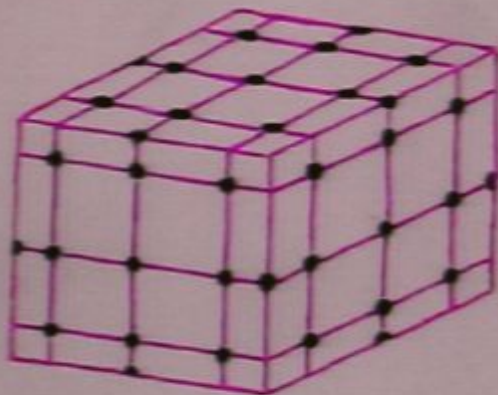
Quantum

(Usual) Bayesian

Magic!

Kochen - Specker

Cannot be colored:



33 rays, Peres

(when completed into full triads, consists of 40 triads made from 57 rays)

$$p(D_j) = (d+1) \sum_i p(H_i) p(D_j | H_i) - 1$$

Quantum

(Usual) Bayesian

Magic!

Think SIC thoughts!

... and maybe by way of it
we'll come to understand
quantum mechanics a
little better.