

Title: Virtual extra dimensions from unparticle fields

Date: Jun 03, 2008 03:15 PM

URL: <http://pirsa.org/08060138>

Abstract: The phenomenology of TeV-scale physics can be modified by the existence of scale-invariant 'unparticle' fields that can couple to the standard model sector. In particular, it has been shown that unparticles can alter gravitational interactions in a similar fashion to extra dimensions. Observable results from this mechanism -- most notably mini-black hole formation -- will be discussed, and methods of differentiating unparticles from extra dimensional models will be addressed.

“Virtual” Extra Dimensions from Unparticles (and LHC black holes!)



J. R. Mureika
Department of Physics
Loyola Marymount University

Phys. Lett. B **660**, 561-566 (2008)
arXiv: 0712.1786



Black Holes in $4+n$ dimensions

Black Holes in 4+n dimensions

- (4+n)-D Schwarzschild metric (below R_{compact}):

$$ds^2 = \left(1 - h(r)^{n+1}\right) dt^2 - \frac{dr^2}{1 - h(r)^{n+1}} - r^2 d\Omega_{n-2}^2 \quad h(r) = \frac{r_H}{r} .$$

Black Holes in 4+n dimensions

- (4+n)-D Schwarzschild metric (below R_{compact}):

$$ds^2 = \left(1 - h(r)^{n+1}\right) dt^2 - \frac{dr^2}{1 - h(r)^{n+1}} - r^2 d\Omega_{n-2}^2 \quad h(r) = \frac{r_H}{r} .$$

- Modified horizon:

$$r_H = \frac{1}{\sqrt{\pi} M_*} \left(\frac{M_{BH}}{M_*}\right)^{\frac{1}{n+1}} \left(\frac{8\Gamma\left(\frac{n+3}{2}\right)}{n+2}\right)^{\frac{1}{n+1}}$$

Black Holes in 4+n dimensions

- (4+n)-D Schwarzschild metric (below R_{compact}):

$$ds^2 = \left(1 - h(r)^{n+1}\right) dt^2 - \frac{dr^2}{1 - h(r)^{n+1}} - r^2 d\Omega_{n-2}^2 \quad h(r) = \frac{r_H}{r} .$$

- Modified horizon:

$$r_H = \frac{1}{\sqrt{\pi} M_*} \left(\frac{M_{BH}}{M_*}\right)^{\frac{1}{n+1}} \left(\frac{8\Gamma\left(\frac{n+3}{2}\right)}{n+2}\right)^{\frac{1}{n+1}}$$

- $M_{BH} > M_* \sim 1 \text{ TeV}$, $r_H \sim \text{sub-fm scale}$
- Geometric cross-section: $\sigma = \pi r_H^2 \sim 100 \text{ pb}$
- Max. luminosity @ LHC: 1 event / second

Unparticle Physics Primer

Unparticle Physics Primer

Scale-invariant Banks-Zaks field (dim. d_{BZ}), weakly interacting with SM via massive exchange particle M_U :

$$\mathcal{L} = \frac{1}{M_U^k} \mathcal{O}_{SM} \mathcal{O}_{BZ}$$

Unparticle Physics Primer

Scale-invariant Banks-Zaks field (dim. d_{BZ}), weakly interacting with SM via massive exchange particle M_u :

$$\mathcal{L} = \frac{1}{M_u^k} \mathcal{O}_{SM} \mathcal{O}_{BZ}$$

Below some scale, undergoes dimensional transmutation $\mathcal{O}_{BZ} \rightarrow C_u \lambda^{d_{BZ}-d_U} \mathcal{O}_u$ to become an “unparticle” field \mathcal{O}_u (scalar, vector, tensor, spinor) at energy scale Λ_u :

$$\mathcal{L} = \frac{\kappa}{\Lambda_u^k} \mathcal{O}_{SM} \mathcal{O}_u \quad , \quad \kappa = C_u \left(\frac{\Lambda_u}{M_{BZ}} \right)^k$$

Unparticle Physics Primer

*Georgi, PRL 98, 221601 (2007);
PLB 650, 275 (2007)*

Unparticle Physics Primer

- Propagator:

*Georgi, PRL 98, 221601 (2007);
PLB 650, 275 (2007)*

$$\langle 0 | \mathcal{O}_u(x) \mathcal{O}_u^\dagger(0) | 0 \rangle = \int \frac{d^4 p}{(4\pi)^4} e^{iPx} |\langle 0 | \mathcal{O}_u(0) | P \rangle|^2 \rho(P^2)$$

Unparticle Physics Primer

- Propagator:

*Georgi, PRL 98, 221601 (2007);
PLB 650, 275 (2007)*

$$\langle 0 | \mathcal{O}_U(x) \mathcal{O}_U^\dagger(0) | 0 \rangle = \int \frac{d^4 p}{(4\pi)^4} e^{iPx} |\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2)$$

- Unparticle Phase Space:

$$|\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2) = A_{d_U} \theta(P^0) \theta(P^2) (P^2)^{d_U - 2}$$

Unparticle Physics Primer

- Propagator:

*Georgi, PRL 98, 221601 (2007);
PLB 650, 275 (2007)*

$$\langle 0 | \mathcal{O}_U(x) \mathcal{O}_U^\dagger(0) | 0 \rangle = \int \frac{d^4 p}{(4\pi)^4} e^{iPx} |\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2)$$

- Unparticle Phase Space:

$$|\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2) = A_{d_U} \theta(P^0) \theta(P^2) (P^2)^{d_U - 2}$$

- SM Particle Phase Space:

$$A_n \theta(P^0) \theta(P^2) (P^2)^{n-2} ,$$

The Unparticle Interpretation

$$|\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2) = A_{d_U} \theta(P^0) \theta(P^2) (P^2)^{d_U - 2}$$

The Unparticle Interpretation

$$|\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2) = A_{d_U} \theta(P^0) \theta(P^2) (P^2)^{d_U - 2}$$

**Unparticle stuff “looks like” d_U (non-integer)
indivisible, massless particles**

[Georgi, PRL 98, 221601 (2007)]

The Unparticle Interpretation

$$|\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2) = A_{d_U} \theta(P^0) \theta(P^2) (P^2)^{d_U - 2}$$

**Unparticle stuff “looks like” d_U (non-integer)
indivisible, massless particles**

[Georgi, PRL 98, 221601 (2007)]

- How does one quantify and/or measure this??

The Unparticle Interpretation

$$|\langle 0 | \mathcal{O}_U(0) | P \rangle|^2 \rho(P^2) = A_{d_U} \theta(P^0) \theta(P^2) (P^2)^{d_U - 2}$$

**Unparticle stuff “looks like” d_U (non-integer)
indivisible, massless particles**

[Georgi, PRL 98, 221601 (2007)]

- How does one quantify and/or measure this??
- More recently: “Sommerfield” model: massless fermions coupled to massive vector field in infrared limit

[Georgi and Kats, 0805.3593]

The Unparticle Interpretation

The Unparticle Interpretation

Unparticle stuff is a composite BZ-particle with a continuum mass spectrum

[Krasnikov, 0707.1419; Nikolic, 0801.4471; McDonald, 0805.1888]

The Unparticle Interpretation

Unparticle stuff is a composite BZ-particle with a continuum mass spectrum

[Krasnikov, 0707.1419; Nikolic, 0801.4471; McDonald, 0805.1888]

- Unparticle operator is “deconstructed” to sum of massive BZ fields with mass $M_n^2 = \Delta^2 n$

$$O_{\mathcal{U}} \rightarrow \sum_{n=0}^{\infty} F_n \phi_n \quad , \quad F_n^2 = \left| \langle 0 | O_{\mathcal{U}}(0) | \phi_n \rangle \right|^2 = \frac{A_{d_{\mathcal{U}}}}{2\pi} \Delta^2 (M_n^2)^{d_{\mathcal{U}}-2}$$

- In limit $\Delta \rightarrow 0$, mass spectrum, SM $\rightarrow \mathcal{U}$ decay rates, etc... become continuous and finite (infinite # of kinematically allowed states)
- \mathcal{U} to SM decay rates vanish (unparticles are stable)

“Ungravity”

$$T^{\mu\nu} + T_{\mathcal{U}}^{\mu\nu}, \quad T_{\mathcal{U}}^{\mu\nu} \sim \sqrt{|g|} T^{\alpha\beta} \mathcal{O}_{\alpha\beta}^{\mathcal{U}} g_{\mu\nu}.$$

Goldberg and Nath, PRL 100, 031803 (2008)

“Ungravity”

Introduce a tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$T^{\mu\nu} + T_{\mathcal{U}}^{\mu\nu}, \quad T_{\mathcal{U}}^{\mu\nu} \sim \sqrt{|g|} T^{\alpha\beta} \mathcal{O}_{\alpha\beta}^{\mathcal{U}} g_{\mu\nu}.$$

Goldberg and Nath, PRL 100, 031803 (2008)

“Ungravity”

- Introduce a tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$T^{\mu\nu} + T_{\mathcal{U}}^{\mu\nu}, \quad T_{\mathcal{U}}^{\mu\nu} \sim \sqrt{|g|} T^{\alpha\beta} \mathcal{O}_{\alpha\beta}^{\mathcal{U}} g_{\mu\nu}.$$

Goldberg and Nath, PRL 100, 031803 (2008)

- “Cosmological constant” term (but there could be other couplings as well)

“Ungravity”

- Introduce a tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$T^{\mu\nu} + T_{\mathcal{U}}^{\mu\nu}, \quad T_{\mathcal{U}}^{\mu\nu} \sim \sqrt{|g|} T^{\alpha\beta} \mathcal{O}_{\alpha\beta}^{\mathcal{U}} g_{\mu\nu},$$

Goldberg and Nath, PRL 100, 031803 (2008)

- “Cosmological constant” term (but there could be other couplings as well)
- Operator $\mathcal{O}_{\mu\nu}$ represents a spin-2 particle (ungraviton)

“Ungravity”

- Introduce a tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$T^{\mu\nu} + T_{\mathcal{U}}^{\mu\nu}, \quad T_{\mathcal{U}}^{\mu\nu} \sim \sqrt{|g|} T^{\alpha\beta} \mathcal{O}_{\alpha\beta}^{\mathcal{U}} g_{\mu\nu}.$$

Goldberg and Nath, PRL 100, 031803 (2008)

- “Cosmological constant” term (but there could be other couplings as well)
- Operator $\mathcal{O}_{\mu\nu}$ represents a spin-2 particle (ungraviton)
- Note that vector unparticle operator \mathcal{O}_{μ} can couple to baryon currents as $B^{\mu} \mathcal{O}_{\mu}$ (repulsive)

Newtonian limit of “ungravity”

Newtonian limit of “ungravity”

Do some quantum field theory on propagator; obtain non-relativistic limit of tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$V(r) = V_N(r) \left[1 + \frac{2}{\pi^{2d_U-1}} \frac{\Gamma(d_U + \frac{1}{2})\Gamma(d_U - \frac{1}{2})}{\Gamma(2d_U)} \left(\frac{R_*}{r}\right)^{2d_U-2} \right] = V_N(r) \left[1 + \Gamma_{d_U} \left(\frac{R_*}{r}\right)^{2d_U-2} \right],$$

Newtonian limit of “ungravity”

Do some quantum field theory on propagator; obtain non-relativistic limit of tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$V(r) = V_N(r) \left[1 + \frac{2}{\pi^{2d_U-1}} \frac{\Gamma(d_U + \frac{1}{2})\Gamma(d_U - \frac{1}{2})}{\Gamma(2d_U)} \left(\frac{R_*}{r}\right)^{2d_U-2} \right] = V_N(r) \left[1 + \Gamma_{d_U} \left(\frac{R_*}{r}\right)^{2d_U-2} \right],$$

Ungravity
interaction scale $R_* = \Lambda_U^{-1} \left(\frac{M_{Pl}}{\Lambda_U}\right)^{\frac{1}{d_U-1}} \left(\frac{\Lambda_U}{M_U}\right)^{\frac{d_{RZ}}{d_U-1}}$

Newtonian limit of “ungravity”

Do some quantum field theory on propagator; obtain non-relativistic limit of tensor unparticle operator $\mathcal{O}_{\mu\nu}$ coupling to stress-energy $T_{\mu\nu}$

$$V(r) = V_N(r) \left[1 + \frac{2}{\pi^{2d_U-1}} \frac{\Gamma(d_U + \frac{1}{2})\Gamma(d_U - \frac{1}{2})}{\Gamma(2d_U)} \left(\frac{R_*}{r}\right)^{2d_U-2} \right] = V_N(r) \left[1 + \Gamma_{d_U} \left(\frac{R_*}{r}\right)^{2d_U-2} \right],$$

$$\text{Ungravity interaction scale } R_* = \Lambda_U^{-1} \left(\frac{M_{Pl}}{\Lambda_U}\right)^{\frac{1}{d_U-1}} \left(\frac{\Lambda_U}{M_U}\right)^{\frac{d_{BZ}}{d_U-1}}$$

Interactions will depend on mass scales $M_{\mathcal{U}} > \Lambda_{\mathcal{U}}$ and mass dimensions d_{BZ} , $d_{\mathcal{U}}$.

Observable effects include orbital modifications and Cavendish anomalies

Unparticle-Enhanced Black Holes

Mureika, PLB 660 (2008) 561-566

Unparticle-Enhanced Black Holes

$f \Lambda_u \sim 1$ TeV, new EW physics without extra dimensions!

Unparticle-Enhanced Black Holes

f $\Lambda_u \sim 1$ TeV, new EW physics without extra dimensions!

- Metric (weak field and strong field matching):

$$ds^2 = \left[1 - \frac{2GM}{r} \left(1 + \Gamma_{d_u} \left(\frac{R_*}{r} \right)^{2d_u-2} \right) \right] dt^2 + \frac{dr^2}{1 - \frac{2GM}{r} \left(1 + \Gamma_{d_u} \left(\frac{R_*}{r} \right)^{2d_u-2} \right)} + r^2 d\Omega^2 .$$

Unparticle-Enhanced Black Holes

f $\Lambda_{\mathcal{U}} \sim 1$ TeV, new EW physics without extra dimensions!

- Metric (weak field and strong field matching):

$$ds^2 = \left[1 - \frac{2GM}{r} \left(1 + \Gamma_{d_{\mathcal{U}}} \left(\frac{R_*}{r} \right)^{2d_{\mathcal{U}}-2} \right) \right] dt^2 + \frac{dr^2}{1 - \frac{2GM}{r} \left(1 + \Gamma_{d_{\mathcal{U}}} \left(\frac{R_*}{r} \right)^{2d_{\mathcal{U}}-2} \right)} + r^2 d\Omega^2 .$$

- If $R_{int} \ll R_*$, acts like $(2d_{\mathcal{U}}-2)$ extra dimensions!

$$\Phi(r) \sim \frac{GM_{BH} \Gamma_{d_{\mathcal{U}}}}{r} \left(\frac{R_*}{r} \right)^{2d_{\mathcal{U}}-2} .$$

Unparticle-Enhanced Black Holes

f $\Lambda_U \sim 1$ TeV, new EW physics without extra dimensions!

- Metric (weak field and strong field matching):

$$ds^2 = \left[1 - \frac{2GM}{r} \left(1 + \Gamma_{d_U} \left(\frac{R_*}{r} \right)^{2d_U-2} \right) \right] dt^2 + \frac{dr^2}{1 - \frac{2GM}{r} \left(1 + \Gamma_{d_U} \left(\frac{R_*}{r} \right)^{2d_U-2} \right)} + r^2 d\Omega^2 .$$

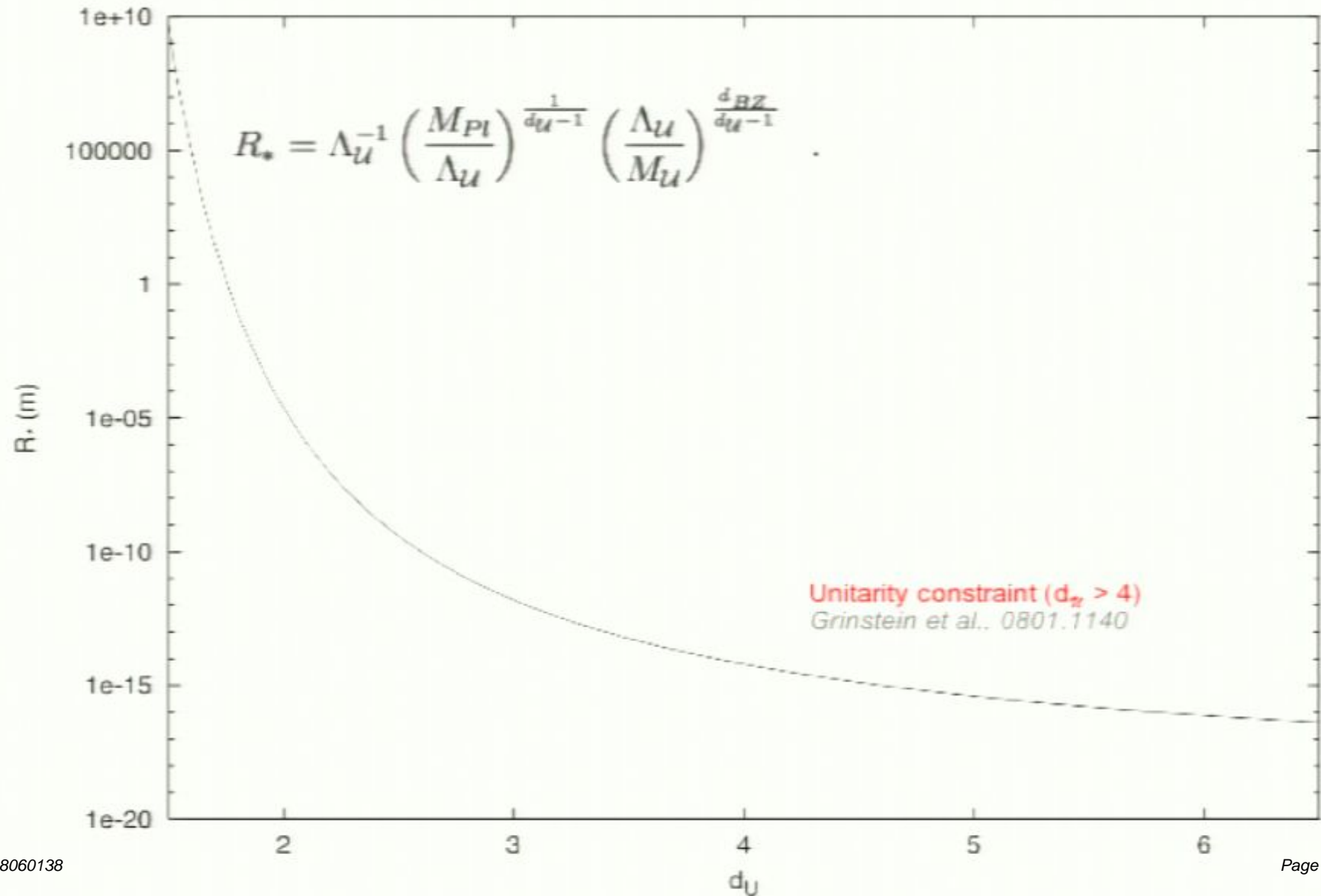
- If $R_{int} \ll R_*$, acts like $(2d_U - 2)$ extra dimensions!

$$\Phi(r) \sim \frac{GM_{BH} \Gamma_{d_U}}{r} \left(\frac{R_*}{r} \right)^{2d_U-2} .$$

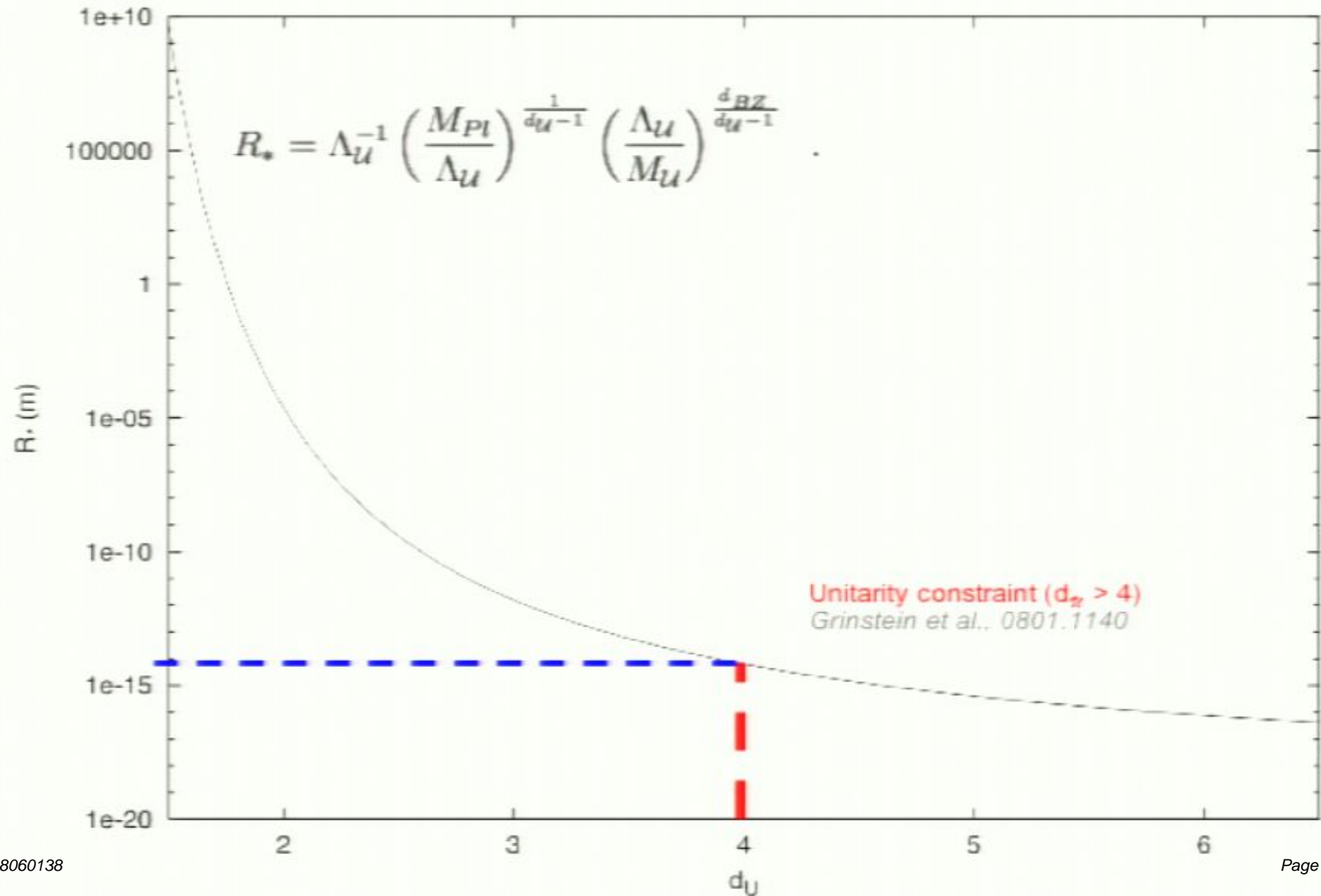
- Modified horizon (M_{Pl} independent):

$$r_H \approx \left(\frac{2M_{BH} \Gamma_{d_U}}{M_U^2 \Lambda_U^{-1}} \right)^{\frac{1}{2d_U-1}} \Lambda_U^{-1}$$

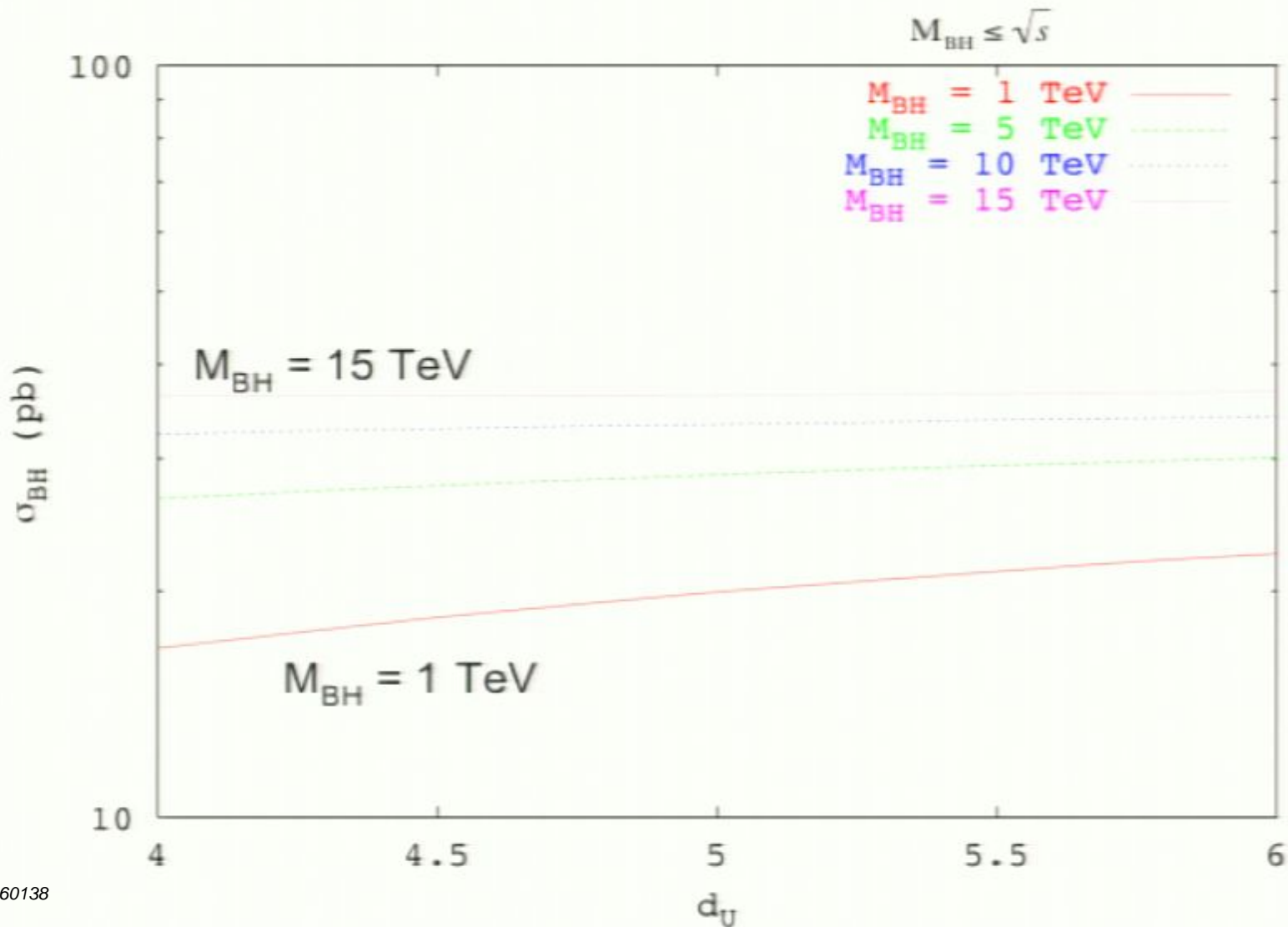
Unparticle interaction scale R_* as a function of operator mass dimension d_U



Unparticle interaction scale R_* as a function of operator mass dimension d_U



Geometric cross-section (pb) for unparticle-enhanced black hole formation, $\Lambda_u = 1 \text{ TeV}$



Differentiating “virtual” from real extra dims

Differentiating “virtual” from real extra dims

- Hawking temperature:

$$T_H = \frac{1}{4\pi r_H} = \frac{\Lambda_{\mathcal{U}}}{4\pi} \left(\frac{2M_{BH} \Gamma_{d_{\mathcal{U}}}}{M_{\mathcal{U}}^2 \Lambda_{\mathcal{U}}^{-1}} \right)^{-\frac{1}{2d_{\mathcal{U}}-1}}$$

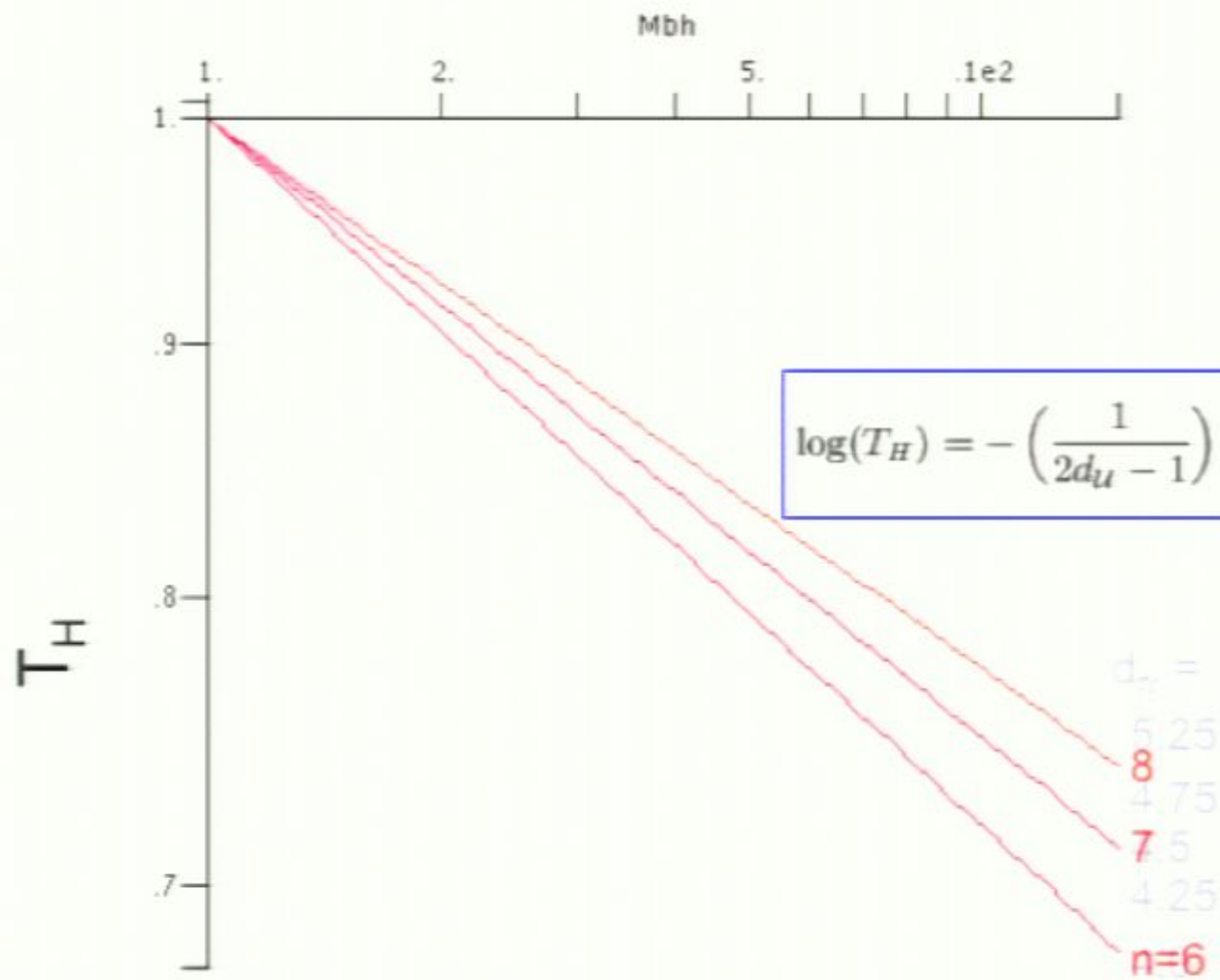
Differentiating “virtual” from real extra dims

- Hawking temperature:

$$T_H = \frac{1}{4\pi r_H} = \frac{\Lambda_{\mathcal{U}}}{4\pi} \left(\frac{2M_{BH} \Gamma_{d_{\mathcal{U}}}}{M_{\mathcal{U}}^2 \Lambda_{\mathcal{U}}^{-1}} \right)^{-\frac{1}{2d_{\mathcal{U}}-1}}$$

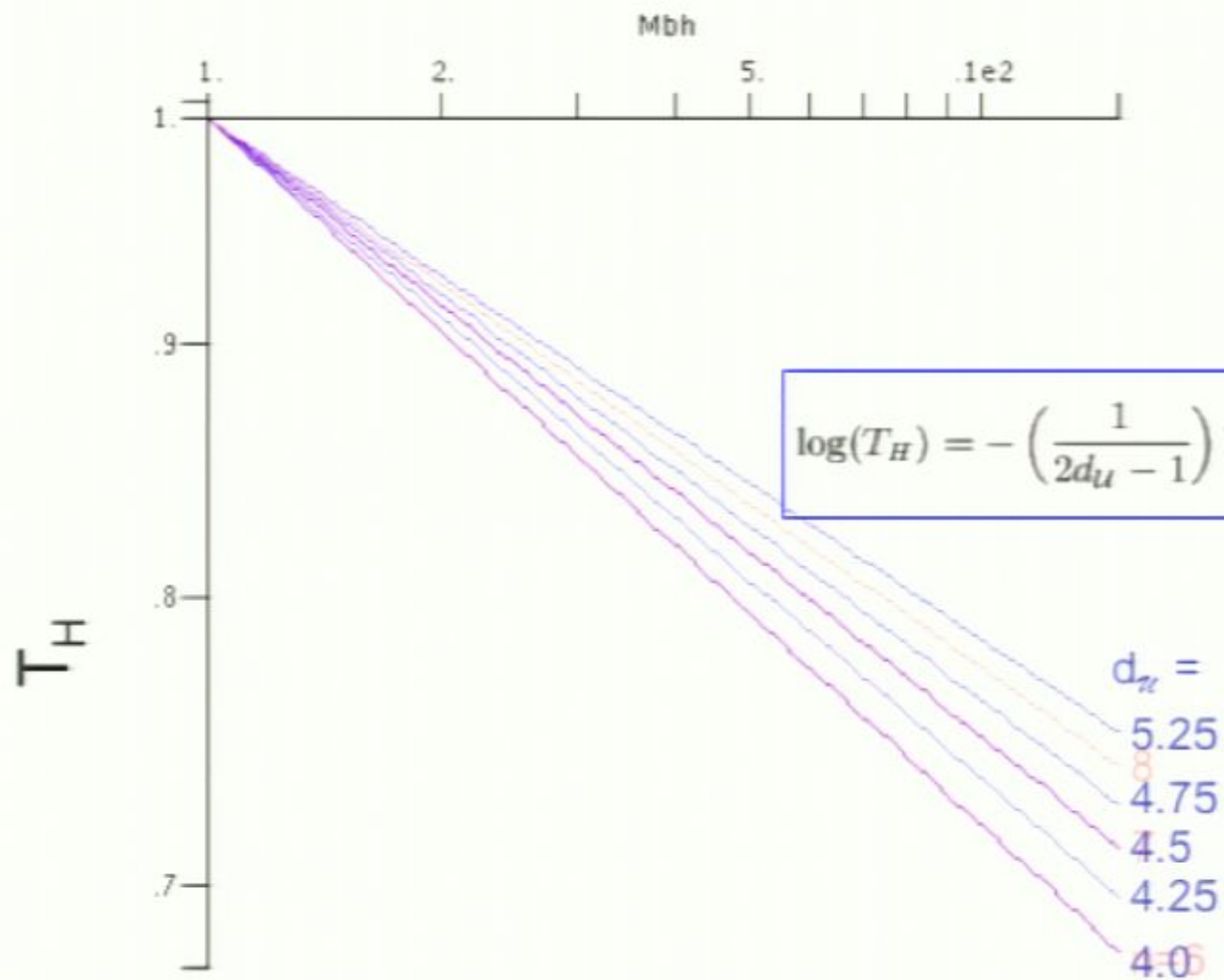
- Unparticle black holes:

$$\log(T_H) = - \left(\frac{1}{2d_{\mathcal{U}} - 1} \right) \log M_{BH} + (\text{constant})$$



$$\log(T_H) = - \left(\frac{1}{2d_U - 1} \right) \log M_{BH} + (\text{constant})$$

$$\log(T_H) = - \left(\frac{1}{n + 1} \right) \log M_{BH} + (\text{constant})$$



Differentiating “virtual” from real extra dims

- Hawking temperature:

$$T_H = \frac{1}{4\pi r_H} = \frac{\Lambda_{\mathcal{U}}}{4\pi} \left(\frac{2M_{BH} \Gamma_{d_{\mathcal{U}}}}{M_{\mathcal{U}}^2 \Lambda_{\mathcal{U}}^{-1}} \right)^{-1/2d_{\mathcal{U}}-1}$$

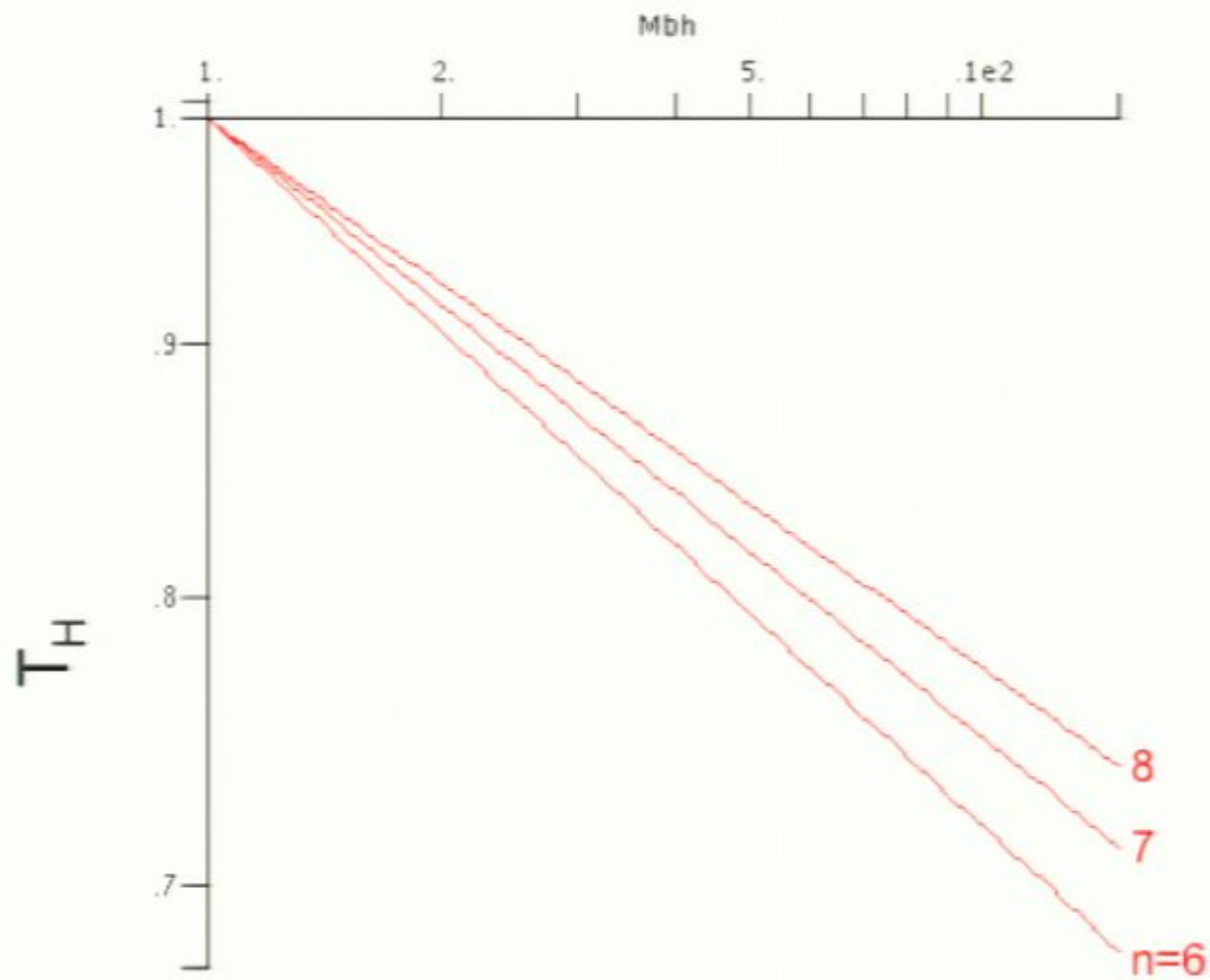
- Unparticle black holes:

$$\log(T_H) = - \left(\frac{1}{2d_{\mathcal{U}} - 1} \right) \log M_{BH} + (\text{constant})$$

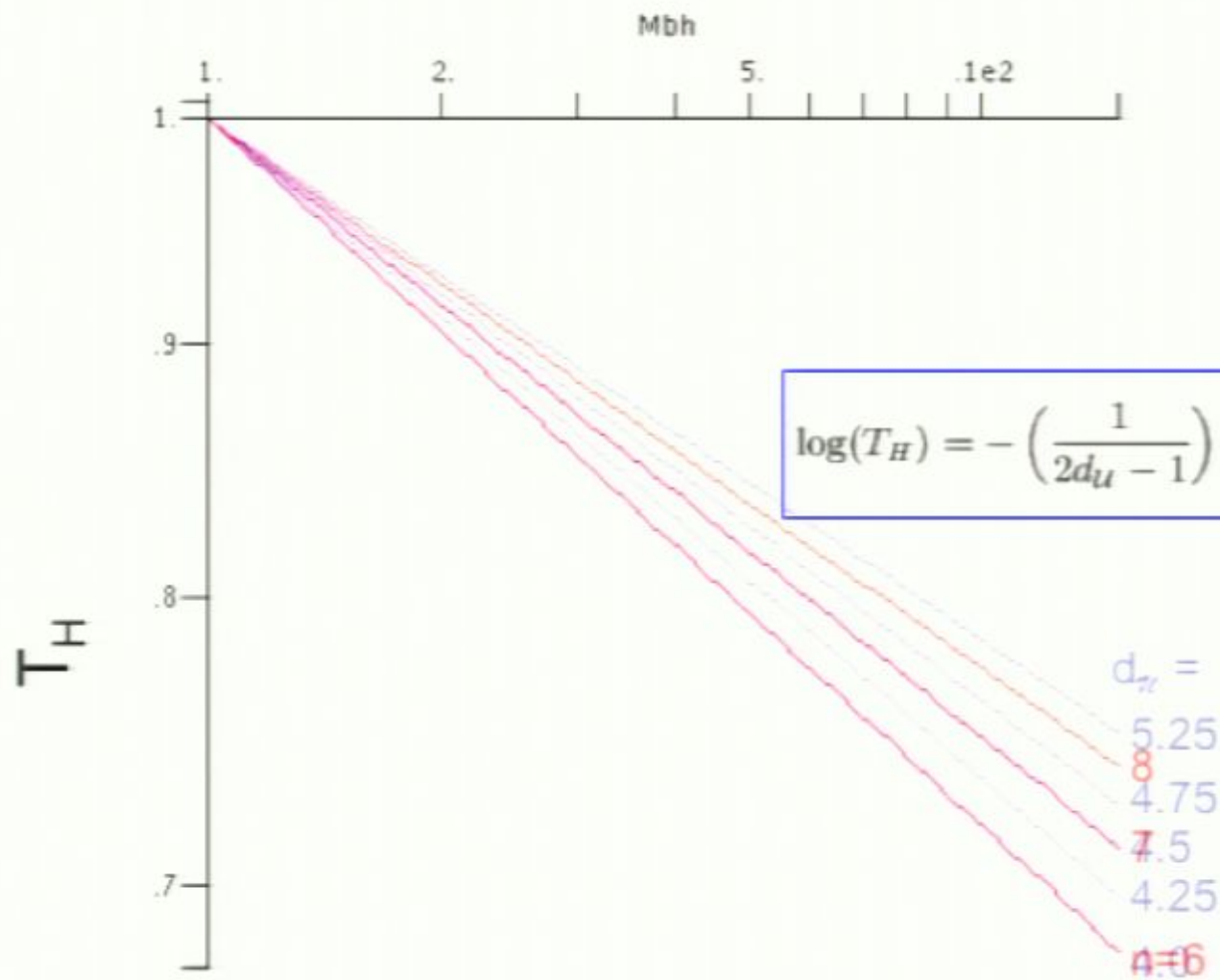
- ED black holes:

$$\log(T_H) = - \left(\frac{1}{n + 1} \right) \log M_{BH} + (\text{constant})$$

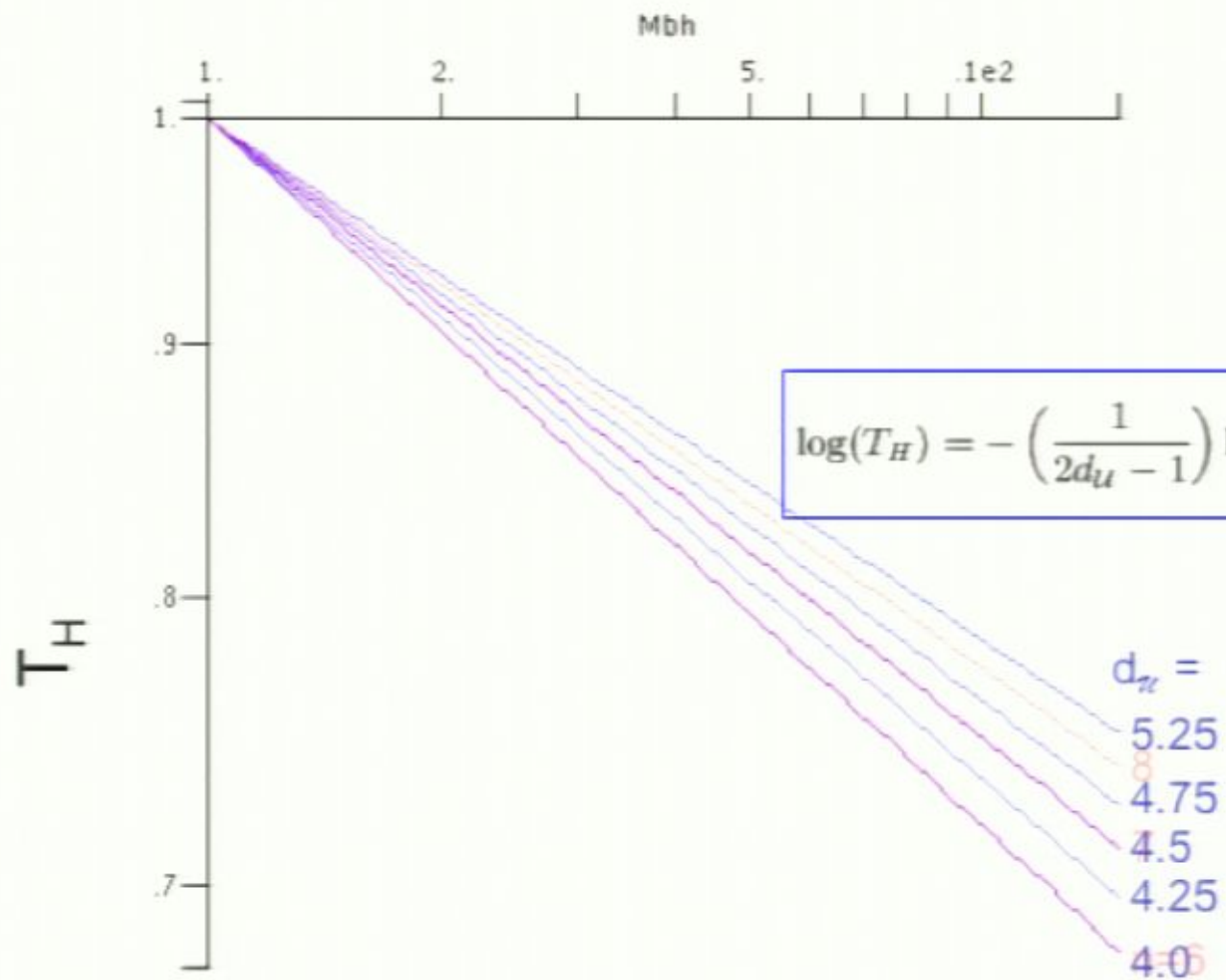
- $d_{\mathcal{U}} \geq 4 \rightarrow n \geq 6$



$$\log(T_H) = - \left(\frac{1}{n+1} \right) \log M_{BH} + (\text{constant})$$



$$\log(T_H) = -\left(\frac{1}{n + 1}\right) \log M_{BH} + (\text{constant})$$



$$\log(T_H) = - \left(\frac{1}{n + 1} \right) \log M_{BH} + (\text{constant})$$

Conclusions and future work

Conclusions and future work

- Ungravity provides a source of TeV-scale black holes without need for extra dimensions

Conclusions and future work

- Ungravity provides a source of TeV-scale black holes without need for extra dimensions
- Parton distributions - unique spectrum too?

Conclusions and future work

- Ungravity provides a source of TeV-scale black holes without need for extra dimensions
- Parton distributions - unique spectrum too?
- Properties of “ungravity metric” (cosmological implications?)
 - Vector ungravity: *repulsive* correction (Reissner-Nordström / Kerr type solutions without charge or a.m.)

Conclusions and future work

- Ungravity provides a source of TeV-scale black holes without need for extra dimensions
- Parton distributions - unique spectrum too?
- Properties of “ungravity metric” (cosmological implications?)
 - Vector ungravity: *repulsive* correction (Reissner-Nordström / Kerr type solutions without charge or a.m.)
- CFT/unitarity considerations?

Conclusions and future work

- Ungravity provides a source of TeV-scale black holes without need for extra dimensions
- Parton distributions - unique spectrum too?
- Properties of “ungravity metric” (cosmological implications?)
 - Vector ungravity: *repulsive* correction (Reissner-Nordström / Kerr type solutions without charge or a.m.)
- CFT/unitarity considerations?
- Beyond LHC (\gg TeV): interesting stuff?

Thank you!

Email: jmureika@lmu.edu

Mureika, Phys. Lett. B. **660**, 561-566 (2008)
arXiv: 0712.1786

