Title: Quark Gluon Plasma at RHIC (and in QCD and String Theory)

Date: Jun 04, 2008 09:00 AM

URL: http://pirsa.org/08060040

Abstract:

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Quark-Gluon Plasma at RHIC (and in QCD and String Theory)

W.A. Zajc
PASCOS 2008
Perimeter Institute

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×Uninteresting question:

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 - What happens when I crash two gold nuclei together?

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- \$ Compelling question:

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- ×Uninteresting question:
 - What happens when I crash two gold nuclei together?
- ✓Interesting question:
 - →Are there new states of matter at the highest temperatures and densities?
- \$ Compelling question:
 - What fundamental thermal properties of our gauge theories of nature can be investigated experimentally?

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- 10 GeV / fm³
 ~10¹⁶ gm / cm³
- T~170 MeV
 ~2 x 10¹² K

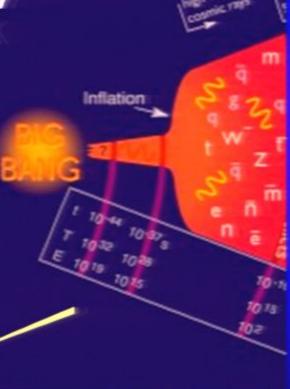
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• 10 GeV / fm³ History of the ~ 10¹⁶ gm / cm³

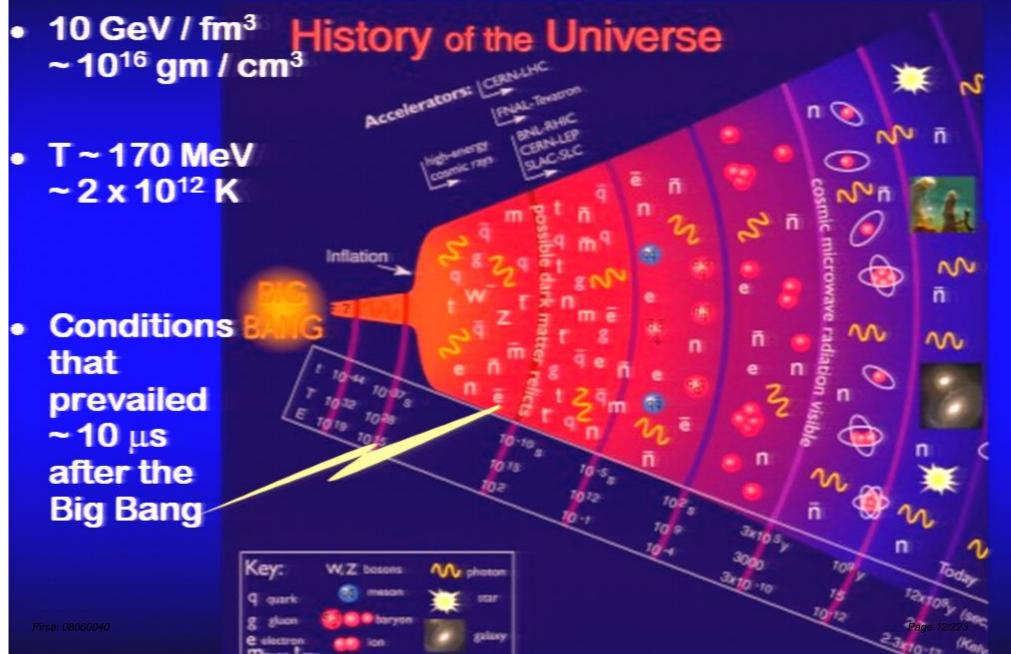
T~170 MeV ~2 x 10¹² K

• Conditions that prevailed ~ 10 μs after the Big Bang

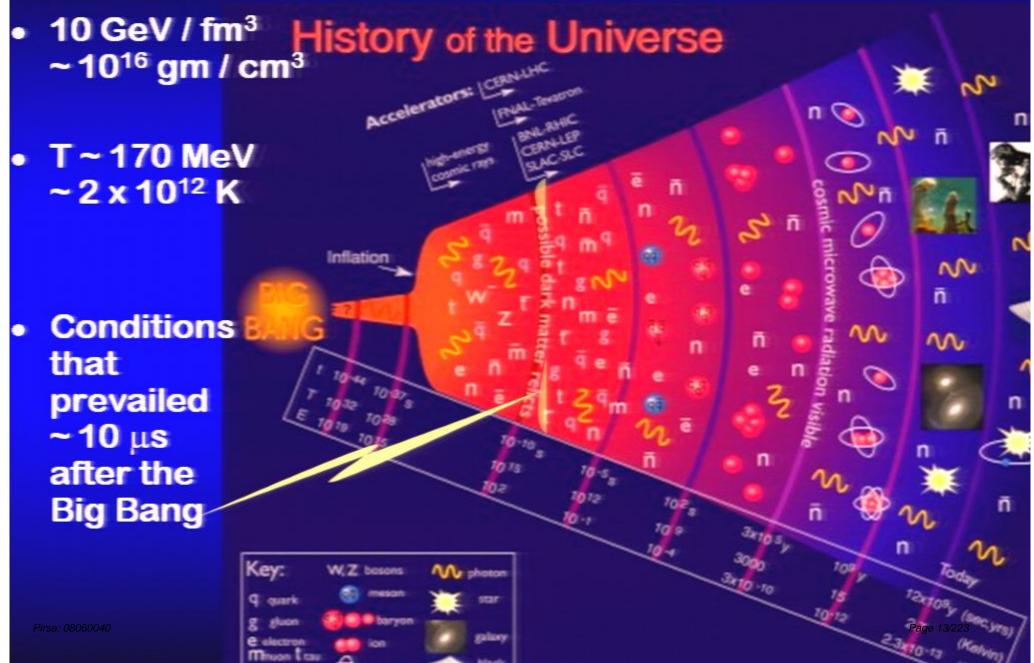




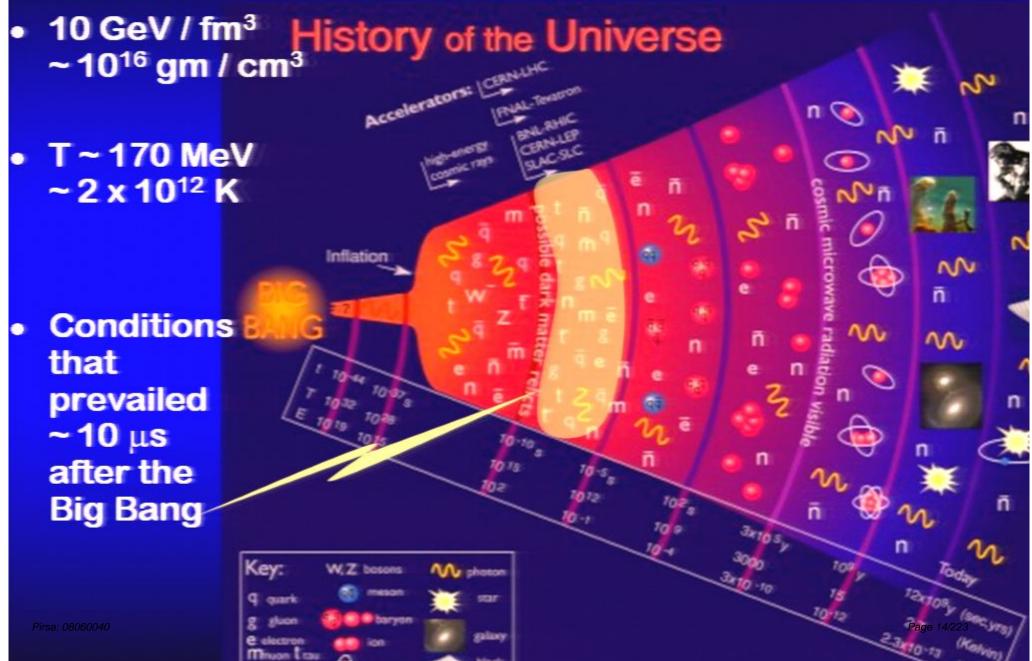














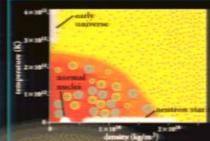
Expectations circa 2000

Nuclear Science



Expansion of the Universe



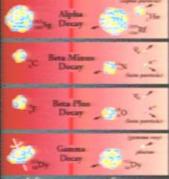


Phases of

Unstable Nucle



Radioactivity



Nucleus

neutron field proton quark

> electromagnetic field

Nuclear Energy



Chart of the Nuclides

Applications



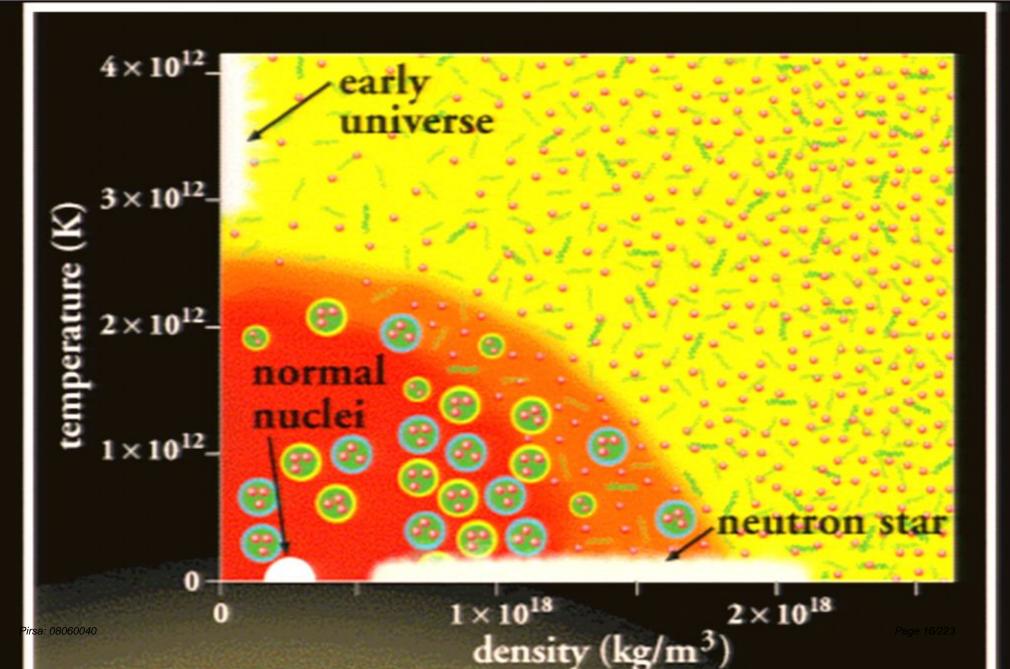






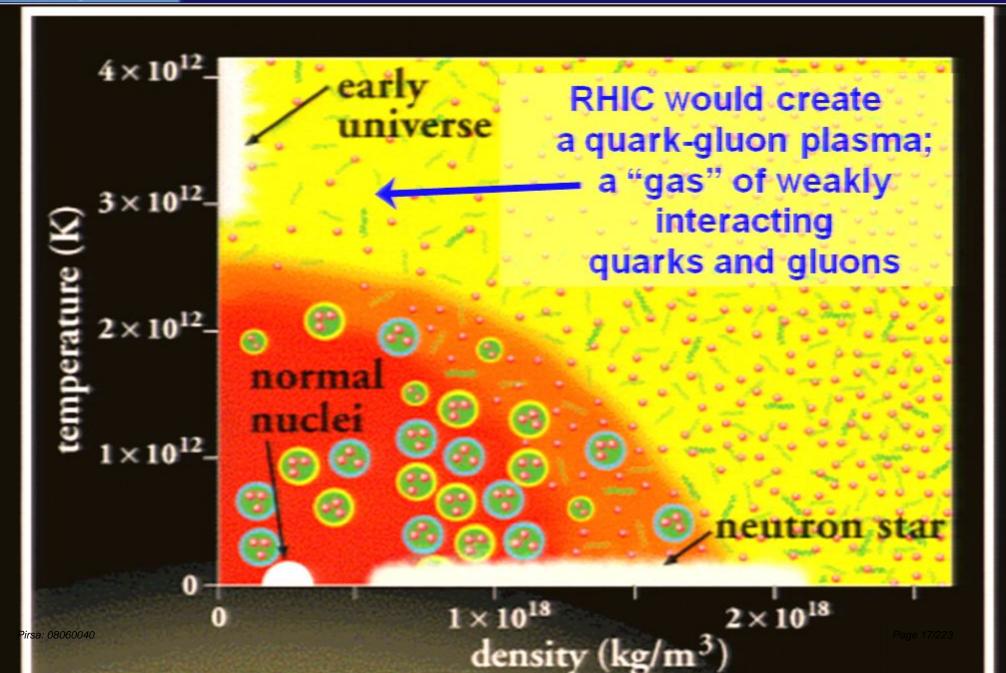


Expectations circa 2000





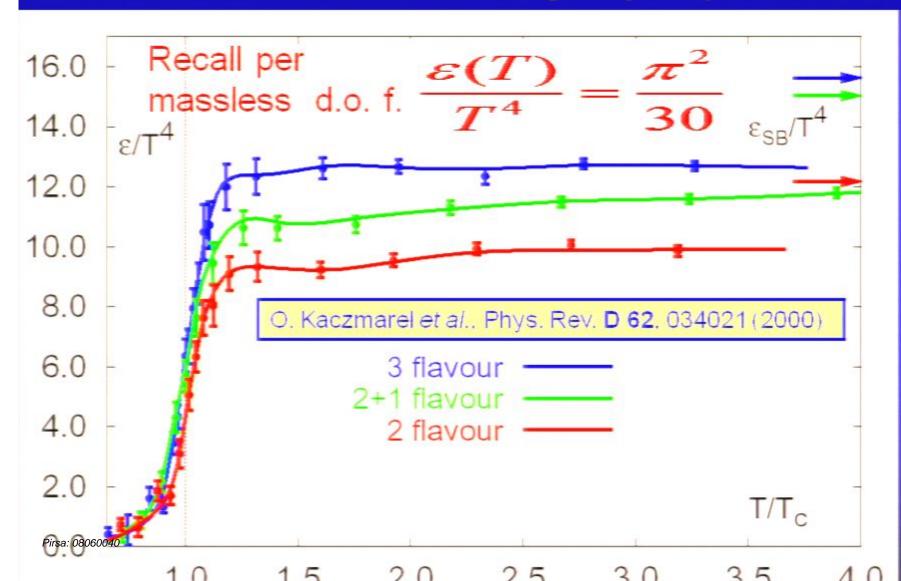
Expectations circa 2000





The Prevailing Prejudice

- Lattice within ~80% of "ideal gas" Stefan-Boltzmann limit
 - Quark-Gluon Plasma is a weakly coupled 'plasma'

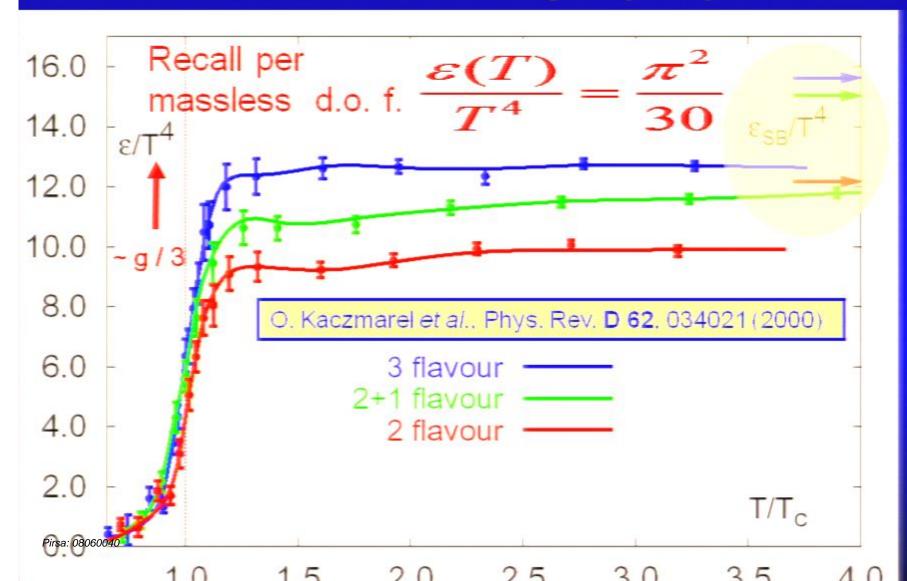


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The Prevailing Prejudice

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- Use RHIC's unprecedented capabilities
 - □ Large √s ⇒
 - Access to reliable pQCD probes
 - Clear separation of valence baryon number and glue
 - ◆ To provide definitive experimental evidence for/against Quark Gluon Plasma (QGP)
 - □ Polarized p+p collisions

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◆ Potential for upgrades in response to discoveries



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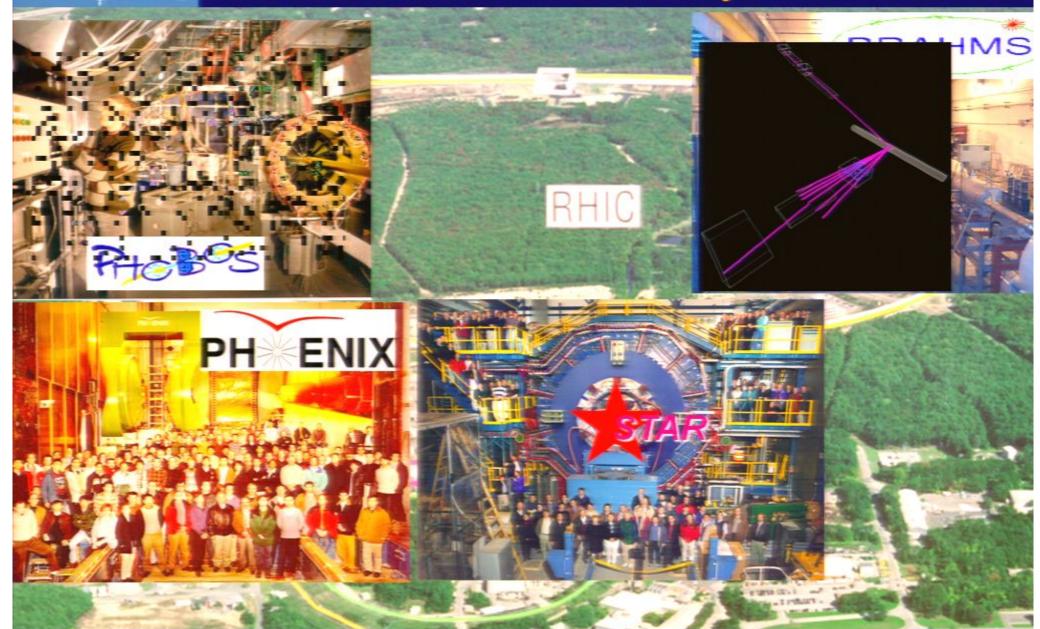




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Since Then...

Accelerator complex

Experiments:

Science

Future

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Since Then...

- Accelerator complex
 - □ Routine operation at 2-4 x design luminosity (Au+Au)
 - □ Extraor

- Experiments:
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Future

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- Accelerator complex
 - □ Routine operation at 2-4 x design luminosity (Au+Au)
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Pirsa: 08060040 Page 38/223



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 - ◆ Species: Au+Au, d+Au, Cu+Cu, p↑+p↑
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Pirsa: 08060040



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 - Demonstrated ability to upgrade
 - □ Key science

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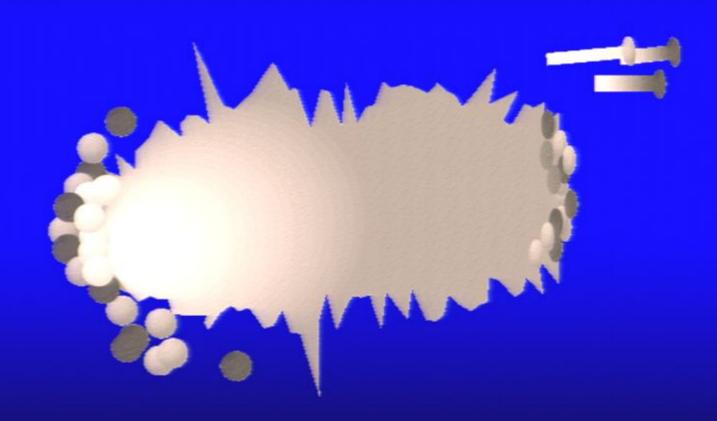
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 - Demonstrated ability to upgrade
 - Key science questions identified
 - □ Accelerator and experimental upgrade program underway to perform that science



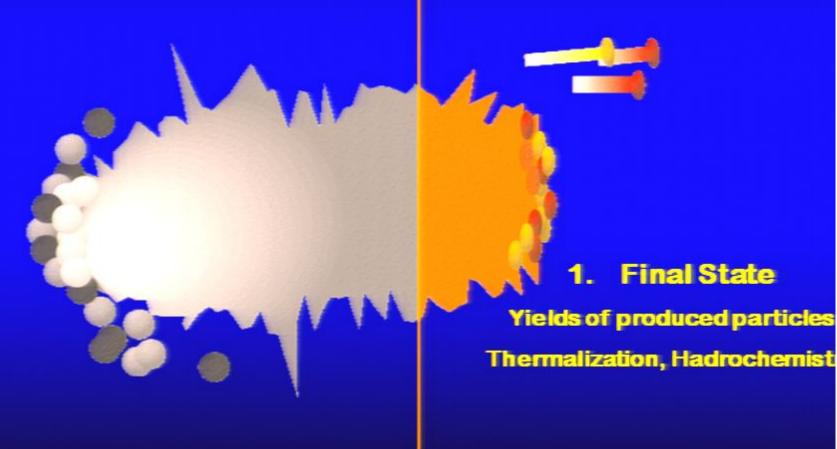
Will present *sample* of results from various points of the collision process:



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Will present *sample* of results from various points of the collision process:

2. Initial State

Hydrodynamic flow from in itial spatial asymmetries

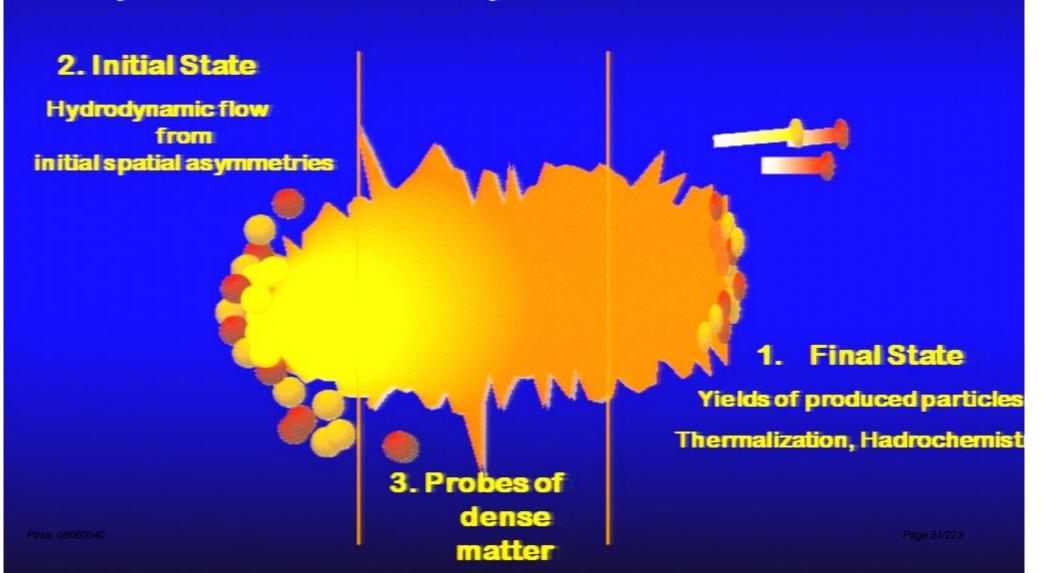
1. Final State

Yields of produced particles

Thermalization, Hadrochemist



Will present *sample* of results from various points of the collision process:



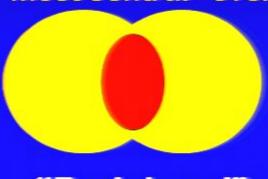


Assertion

 In these complicated events, we have (a posteriori) control over the event geometry:

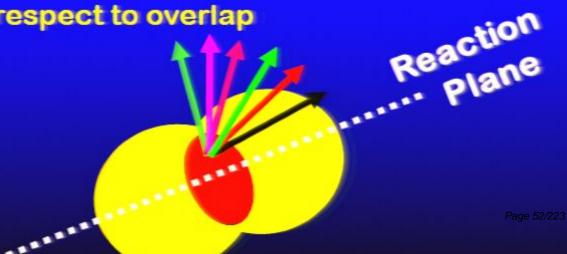
□ Degree of overlap- e.g. [0-10%] "most central" events





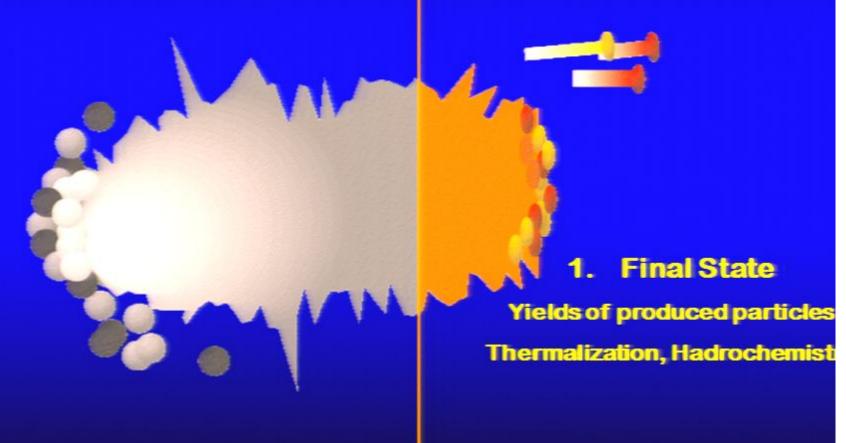
"Peripheral"

Orientation with respect to overlap



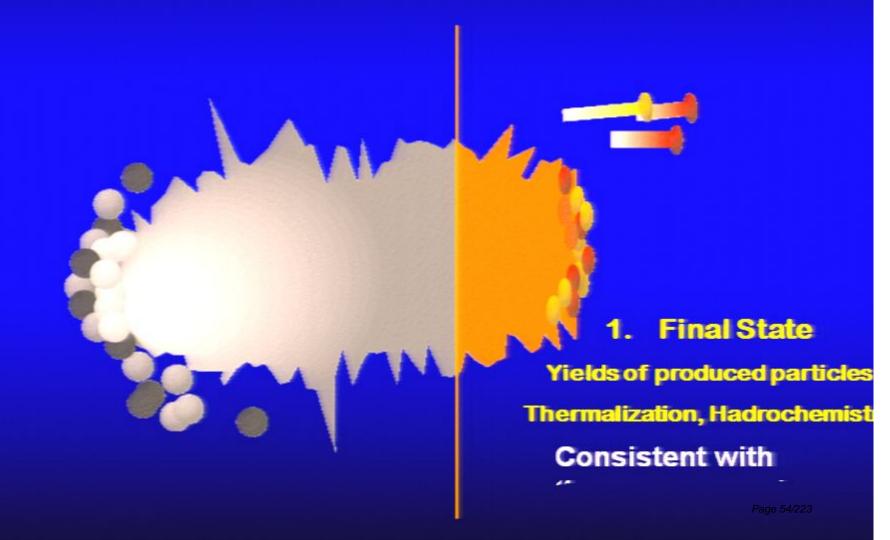


Does the huge abundance of final state particles reflect a *thermal* distribution?:



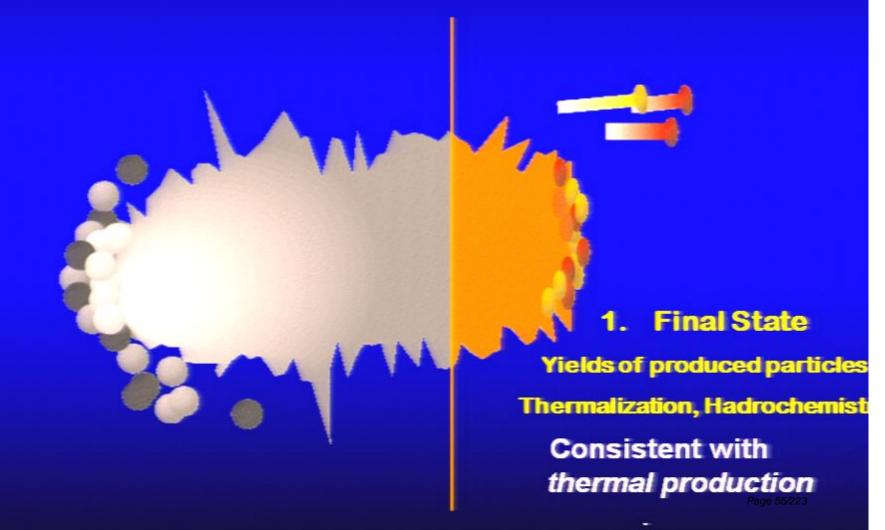


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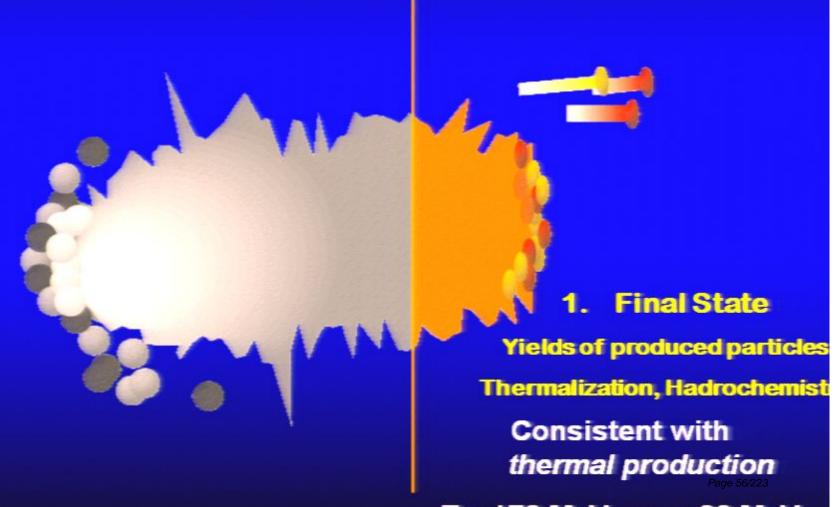


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T~ 170 MeV Up ~ 30 MeV



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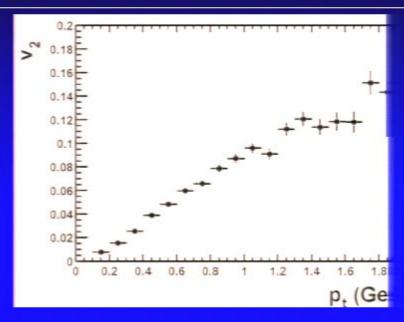


- Discovery of strong "elliptic" flow:
 - □ Elliptic flow in Au + Au collisions at √s_{NN}= 130 GeV,

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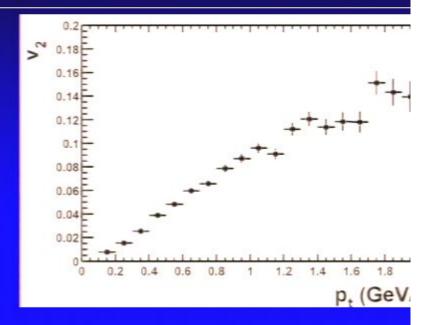
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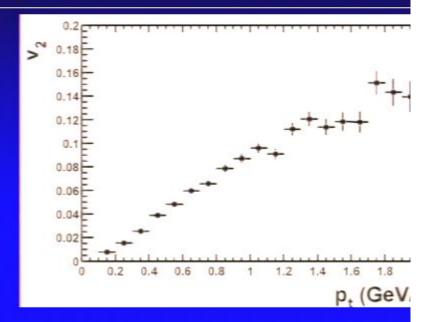
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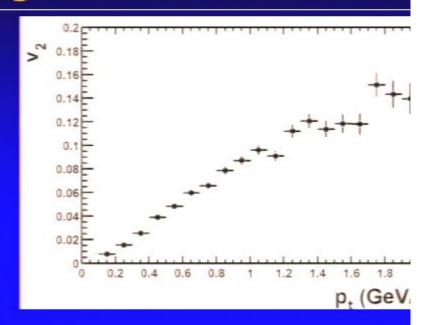


Discovery of

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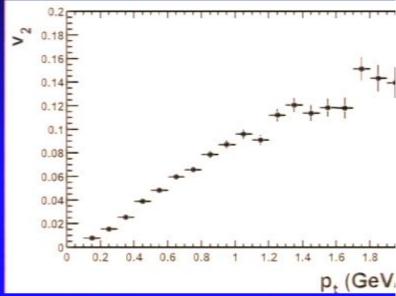


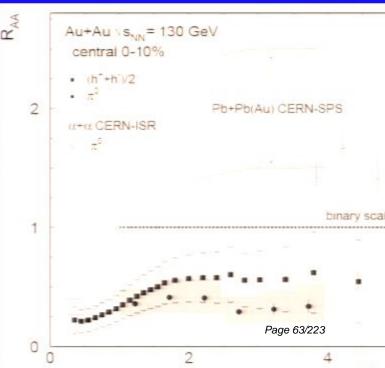
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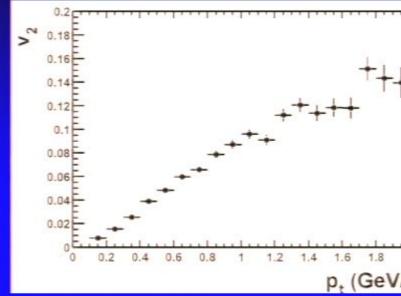
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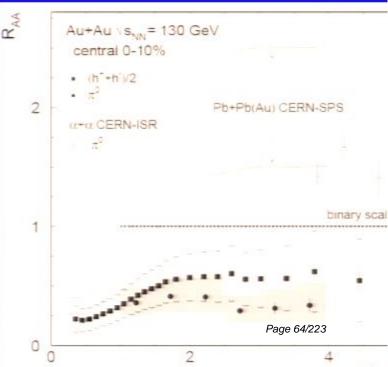






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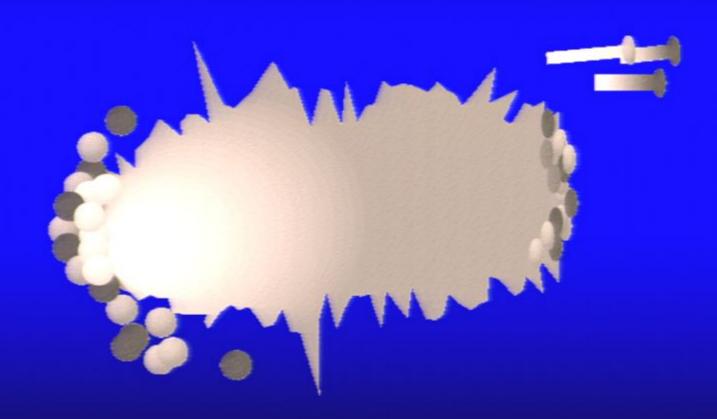






Initial State

How are the initial state densities and asymmetries imprinted on the detected distributions?



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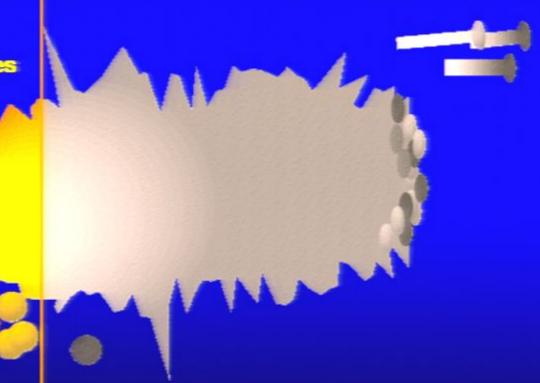


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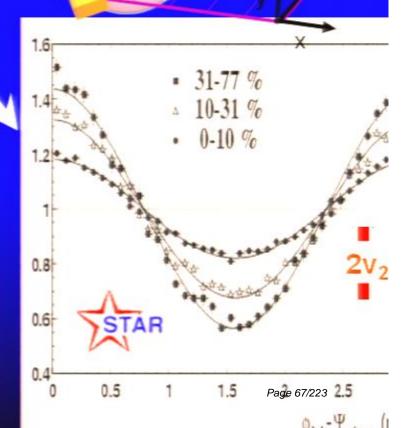
Hydrodynamic flow from in itial spatial asymmetries





- When does thermalization occur?
 - Strong evidence that final state bulk behavior reflects the initial state geometry

Because the initial azimuthal asymmetry persists in the final state dn/dφ ~ 1 + 2 v₂(p₁) cos (2 φ) + ...

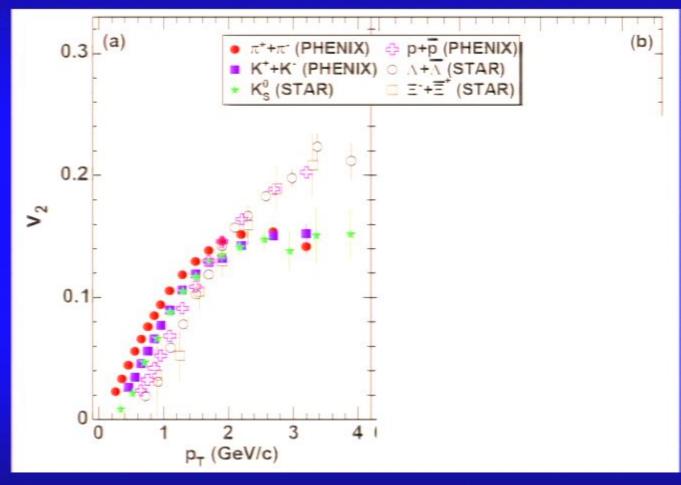




The "Flow" Is Perfect

 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

hydrodynamics

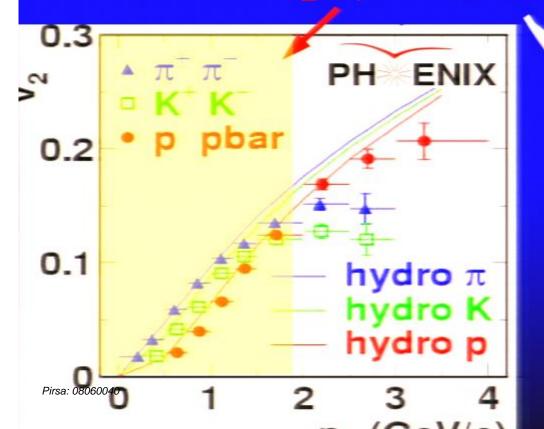


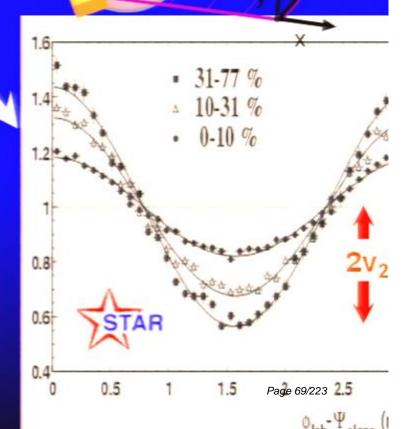
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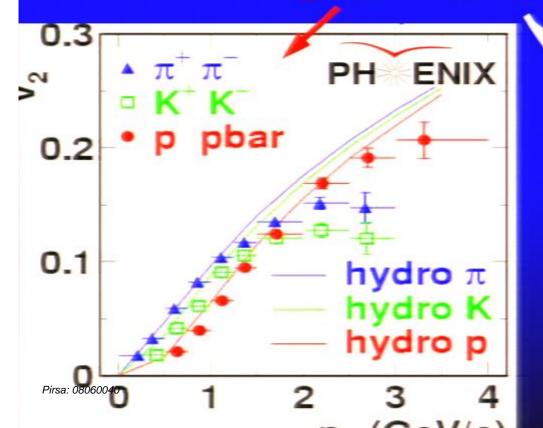


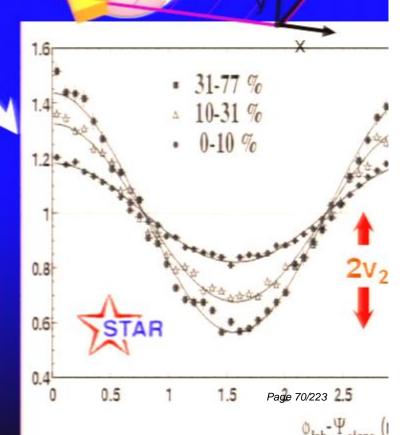




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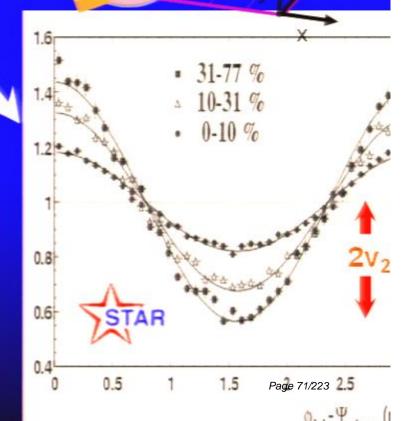






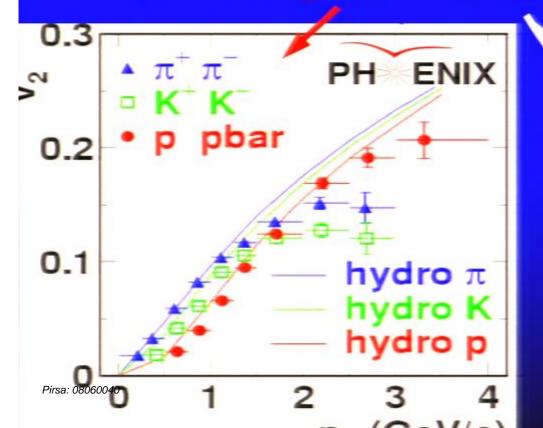
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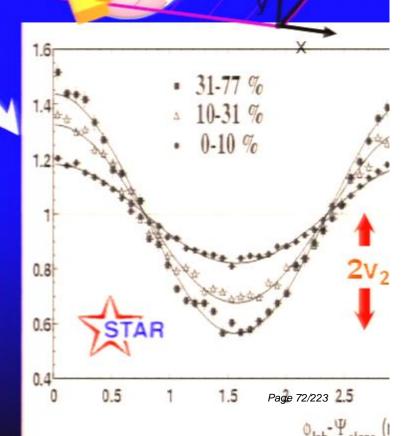
Because the initial azimuthal asymmetry persists in the final state dn/dφ ~ 1 + 2 v₂(p₁) cos (2 φ) + ...





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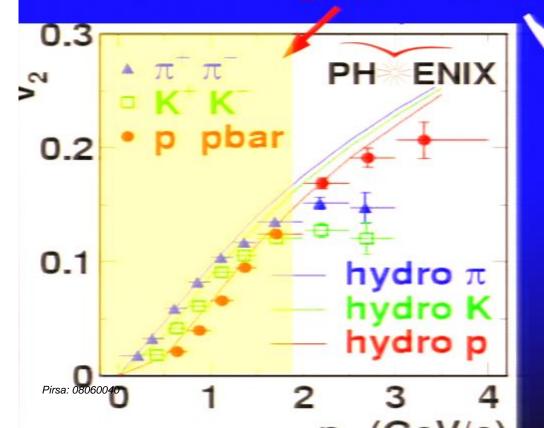


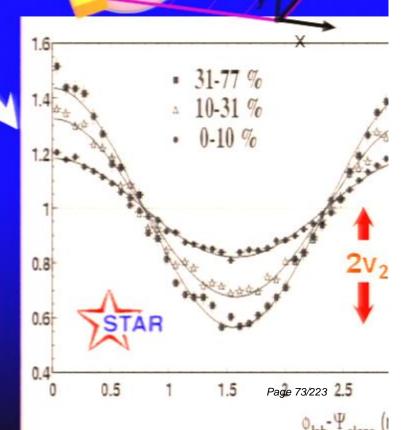


Motion Is Hydrodynamic

- When does thermalization occur?
 - Strong evidence that final state bulk behavior reflects the initial state geometry

Because the initial azimuthal asymmetry
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dn/dφ ~ 1 + 2 v₂(p₁) cos (2 φ) + ...

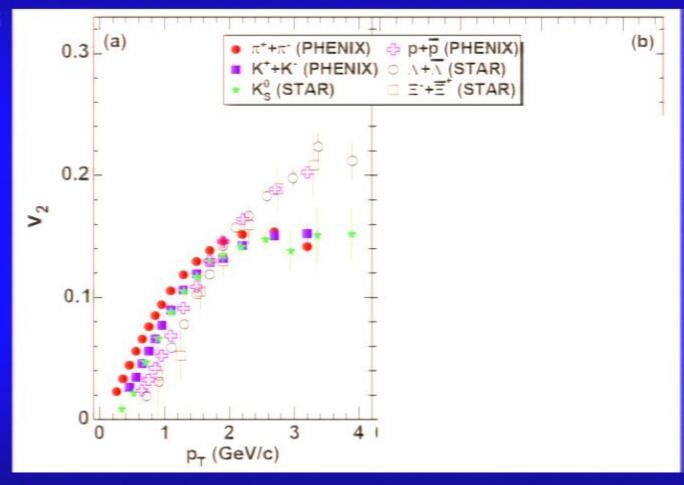






 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

hydrodynamics

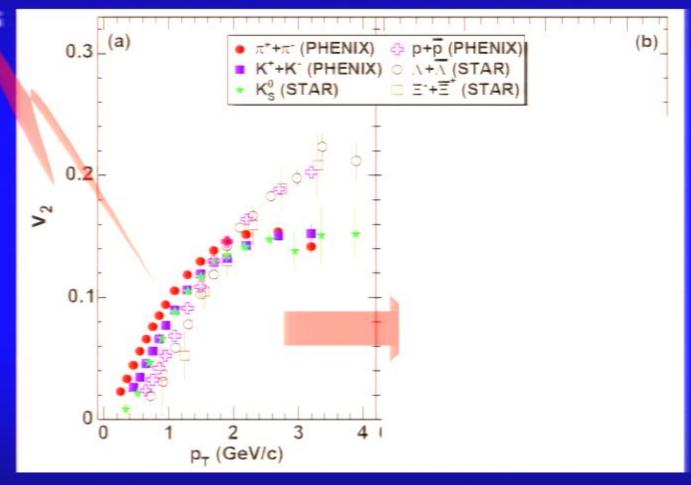


Pirsa: 08060040 Page 74/223



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hydrodynamics



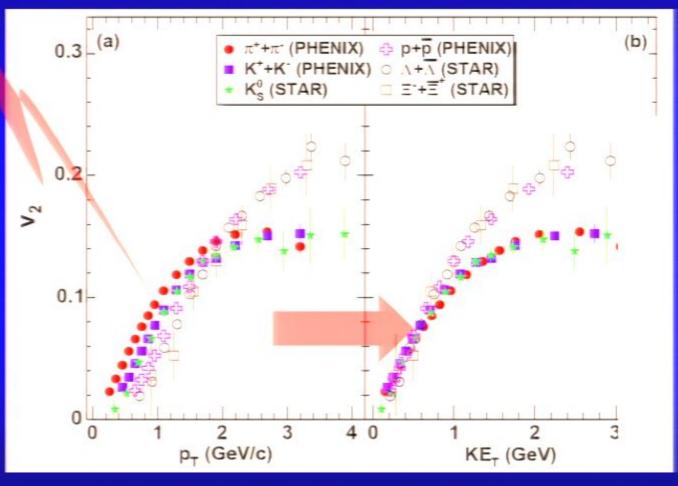
Pirsa: 08060040 Page 75/223



 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

hydrodynamics

$$KE_T \equiv \sqrt{m^2 + I}$$



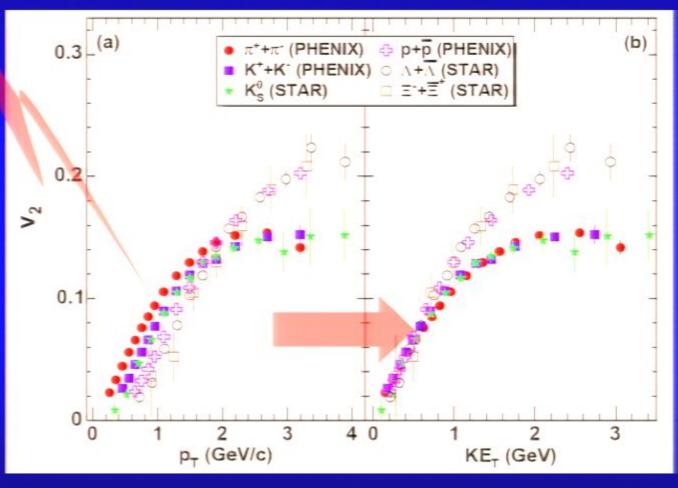
Pirea: 08060040



 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

hydrodynamics

$$KE_T \equiv \sqrt{m^2 + p_T^2}$$



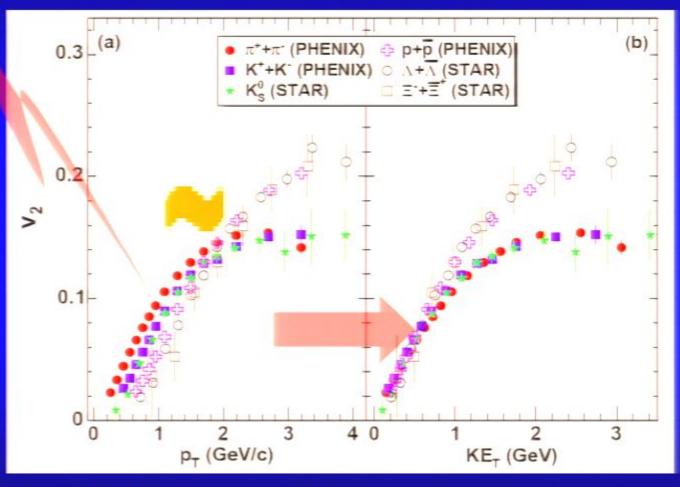
Pirsa: 08060040 Page 77/223



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hydrodynamics

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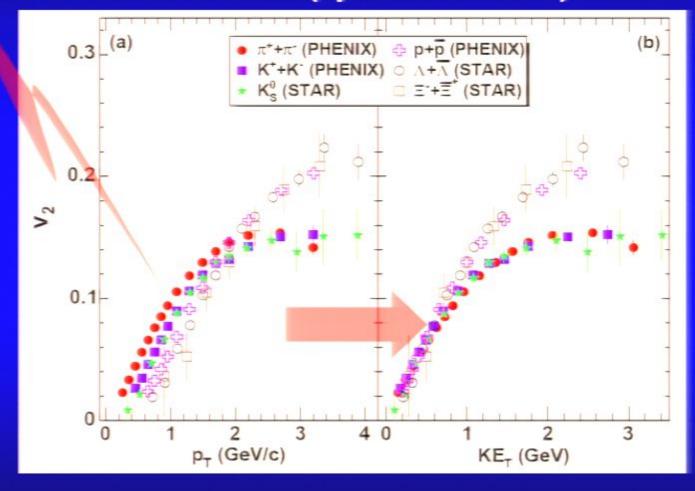
• Roughly: $\partial_{\nu} T^{\mu\nu} = 0 \rightarrow$ Work-energy theorem

$$\rightarrow \int \nabla P \, d(vol) = \Delta E_{\kappa} \simeq m_{T} - m_{0} = \Delta \kappa E_{T}$$



 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

hydrodynamics

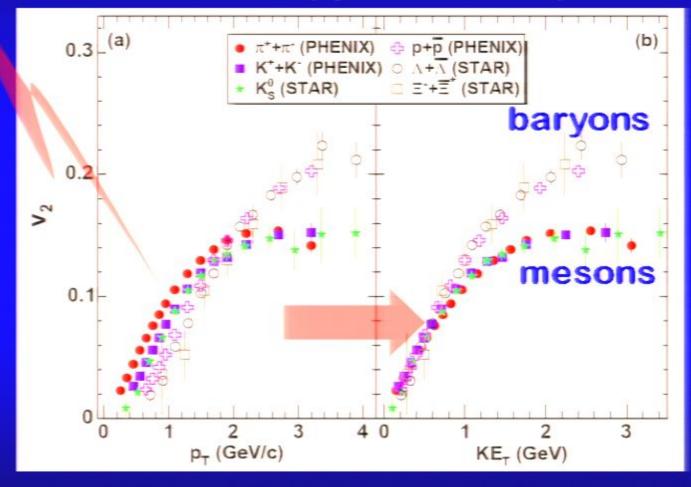


Pirsa: 08060040 Page 79/223



 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

hydrodynamics

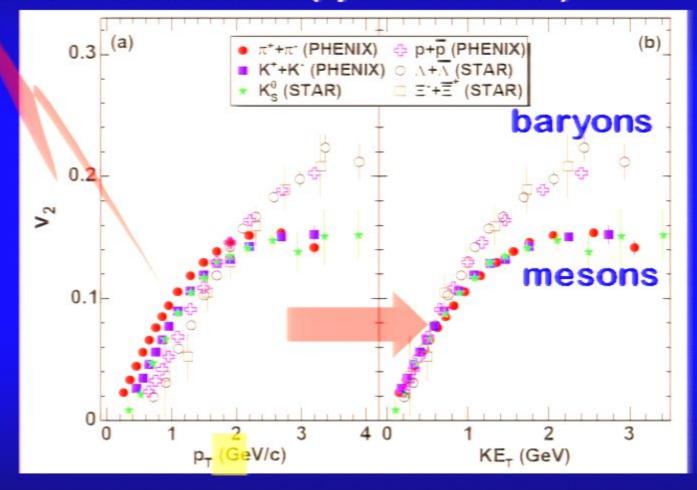


Pirsa: 08060040 Page 80/223



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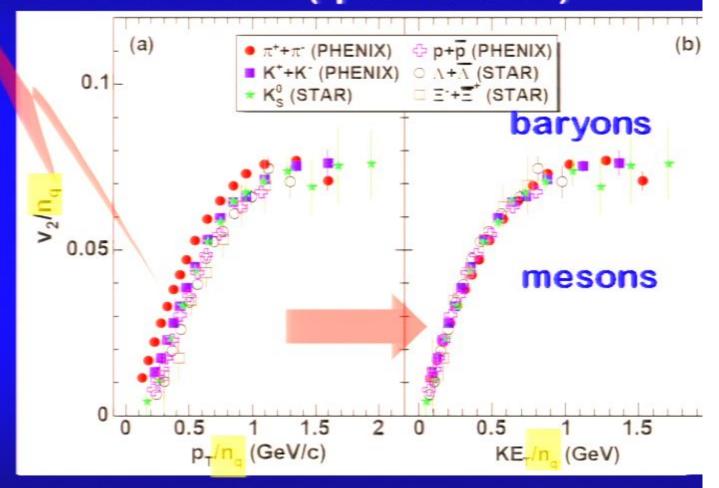


Scaling flow parameters by quark content n_q resolves
 Pire 1780 @ Scaling flow parameters by quark content n_q resolves
 Pire 1780 @ Scaling flow parameters by quark content n_q resolves



 The "fine structure" v₂(p_T) for different mass particles shows good agreement with ideal ("perfect fluid")

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Scaling flow parameters by quark content n_q resolves
 Ping 1990 Scaling flow parameters by quark content n_q resolves



"Perfect fluid" (and/or "ideal hydrodynamics")

defined as "zero viscosity" ⇔ "zero dissipation"

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"Perfect fluid" (and/or "ideal hydrodynamics")

defined as "zero viscosity" ⇔ "zero dissipation"

$$\eta_{QGP}\sim 1$$

Pirsa: 08060040 Page 84/22-



"Perfect fluid" (and/or "ideal hydrodynamics")

defined as "zero viscosity" ⇔ "zero dissipation"

$$\eta_{QGP} \sim 2 \times 10^{11}$$

Pirsa: 08060040 Page 85/223



"Perfect fluid" (and/or "ideal hydrodynamics")

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$$\eta_{QGP} \sim 2 \times 10^{11} \, \text{Pa} \cdot \text{s}$$

Pirsa: 08060040 Page 86/223



"Perfect fluid" (and/or "ideal hydrodynamics")

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$$\eta_{OGP} \sim 2 \times 10^{11} \, \text{Pa} \cdot \text{s}$$

$$\eta_{H_2O} \sim 1 \times$$

Pirsa: 08060040 Page 87/223



"Perfect fluid" (and/or "ideal hydrodynamics")

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$$\eta_{QGP} \sim 2 \times 10^{11} \text{ Pa} \cdot \text{s}$$

$$\eta_{H,o} \sim 1 \times 10^{-3} \,\mathrm{Pa} \cdot \mathrm{s}$$

Pirsa: 08060040 Page 88/223



"Perfect fluid" (and/or "ideal hydrodynamics")

defined as "zero viscosity" ⇔ "zero dissipation"

$$\eta_{QGP} \sim 2 \times 10^{11} \text{Pa} \cdot \text{s}$$
 $\Rightarrow \frac{\eta_{QGP}}{\eta_{H_2O}} \sim 2$
 $\eta_{H_2O} \sim 1 \times 10^{-3} \text{Pa} \cdot \text{s}$

Pirsa: 08060040 Page 89/223



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$$\eta_{QGP} \sim 2 \times 10^{11} \text{Pa} \cdot \text{s} \\
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η

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$$\eta_{Pitch} \sim 2.3$$

Pirsa: 08060040 Page 91/223



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$$\eta_{Pitch} \sim 2.3 \times 10^8 \, \mathrm{P}$$

Pirsa: 08060040 Page 92/223



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Pirsa: 08060040 Page 93/223



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$$\eta_{Pitch} \sim 2.3 \times 10^8 \, \text{Pa·s} \quad \eta_{Glass(A.P.)}$$

Pirsa: 08060040 Page 94/223



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$$\eta_{Pitch} \sim 2.3 \times 10^8 \, \text{Pa} \cdot \text{s}$$
 $\eta_{Glass(A.P.)} \sim 10^{12} \, \text{Pa}$

Pirsa: 08060040 Page 95/223



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Pirsa: 08060040 Page 96/223



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Pirsa: 08060040 Page 97/223



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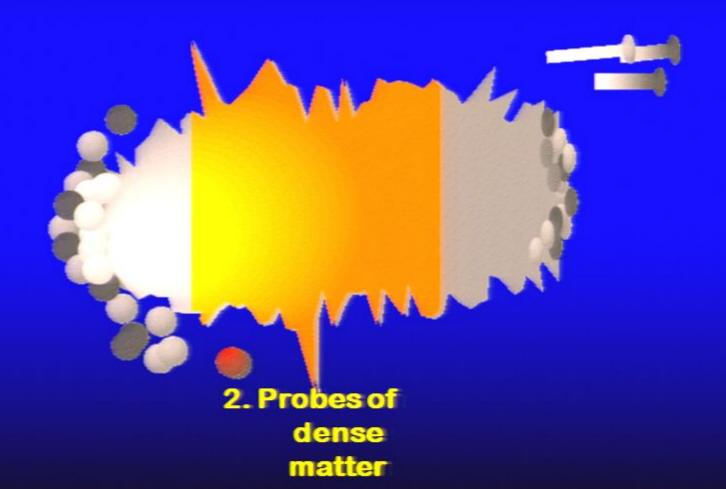
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Probes of Dense Matter

Q. How dense is the matter?

A. Do pQCD Rutherford scattering on deep interior using "auto-generated" probes:





Baseline p+p Measurements with pQCD

- Consider measurement of π^0 's in p+p collisions at RHIC.
- Compare to pQCD calculation

$$d\sigma = f_{a/A}(x_a, \mu^2) \otimes f_{b/B}(x_b, \mu^2)$$

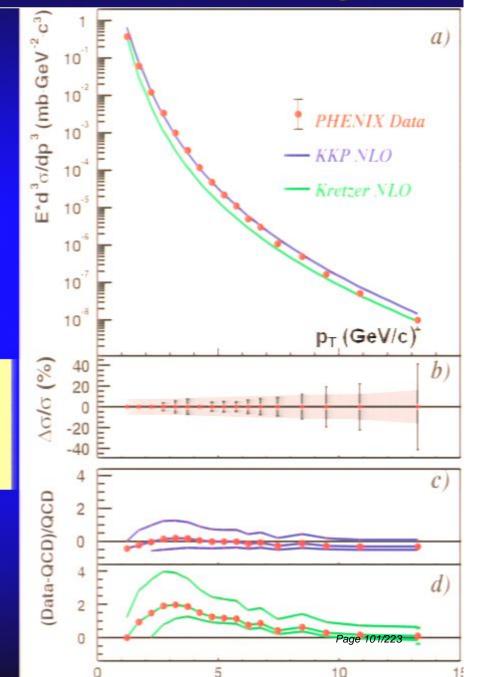
parton distribution functions,
for partons a and b
measured in DIS, universality

$\otimes d \overset{\wedge}{\sigma} (a+b \to c+d)$

- •perturbative cross-section (NLO)
- requires hard scale
- factorization between pdf and cross section

$$\otimes D_{h/c}(z_h,\mu^2)$$

- •fragmentation function
- ·measured in e+e-



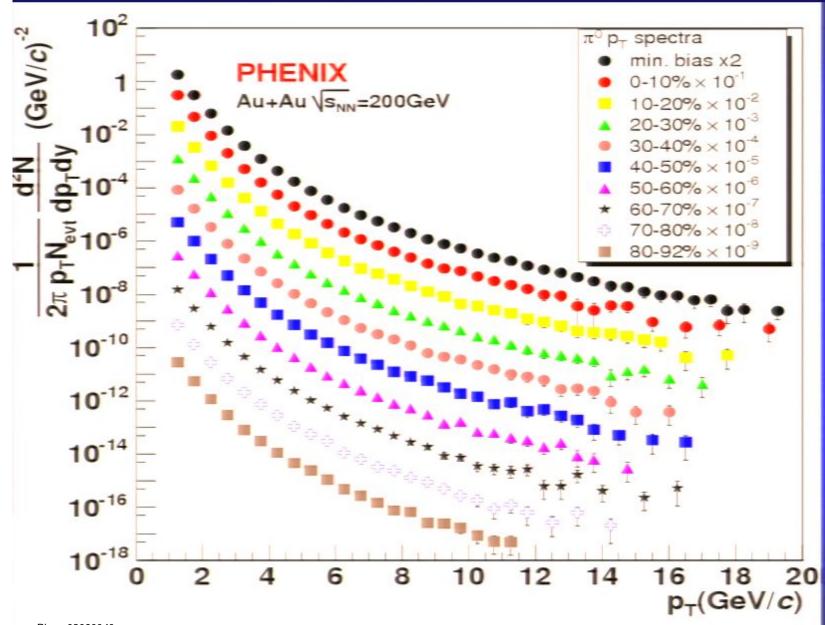
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Phys. Rev. Lett. 91, 241803 (2003)

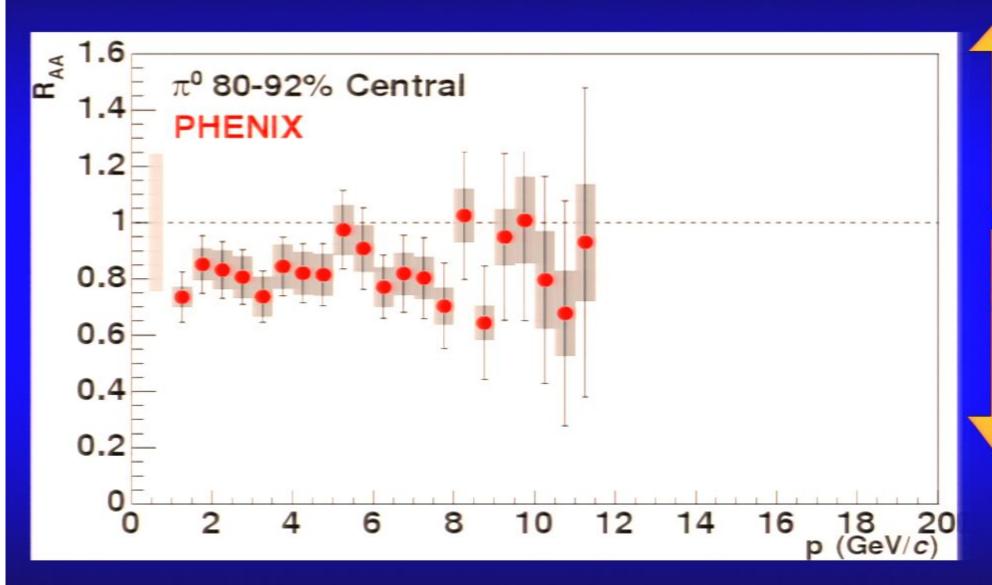


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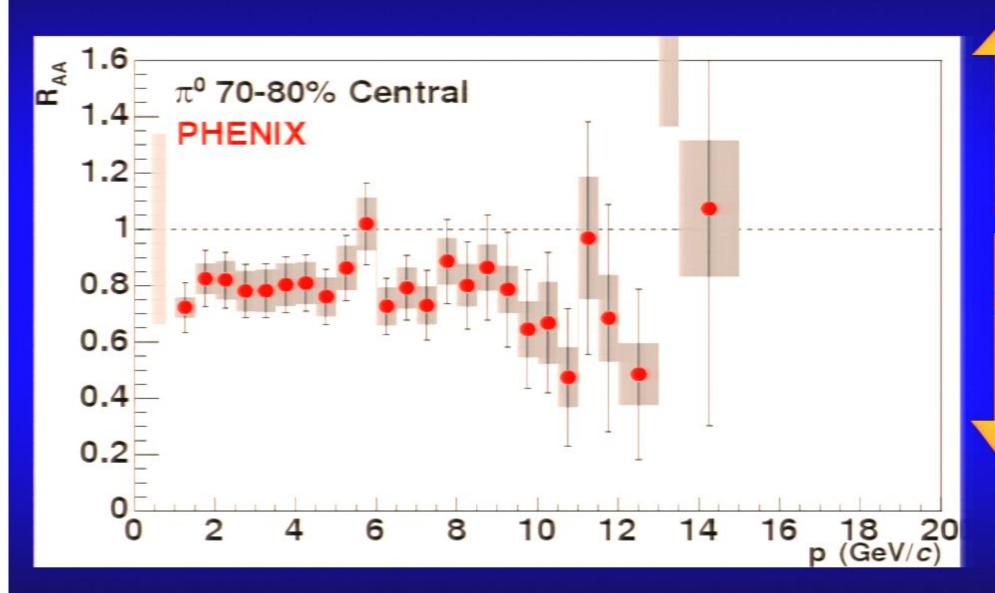






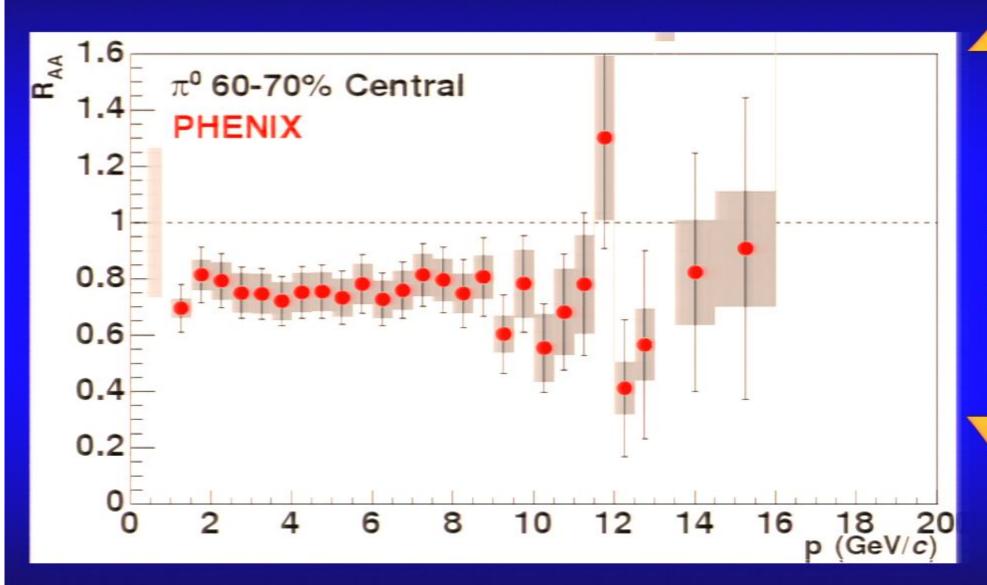
Pirsa: 08060040 Page 104/223





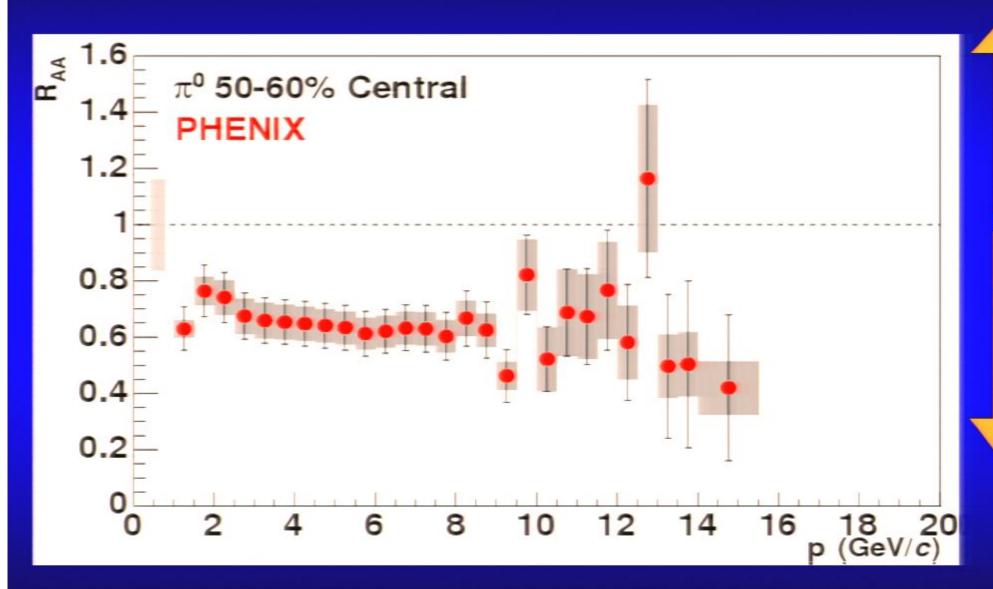
Pirsa: 08060040 Page 105/223





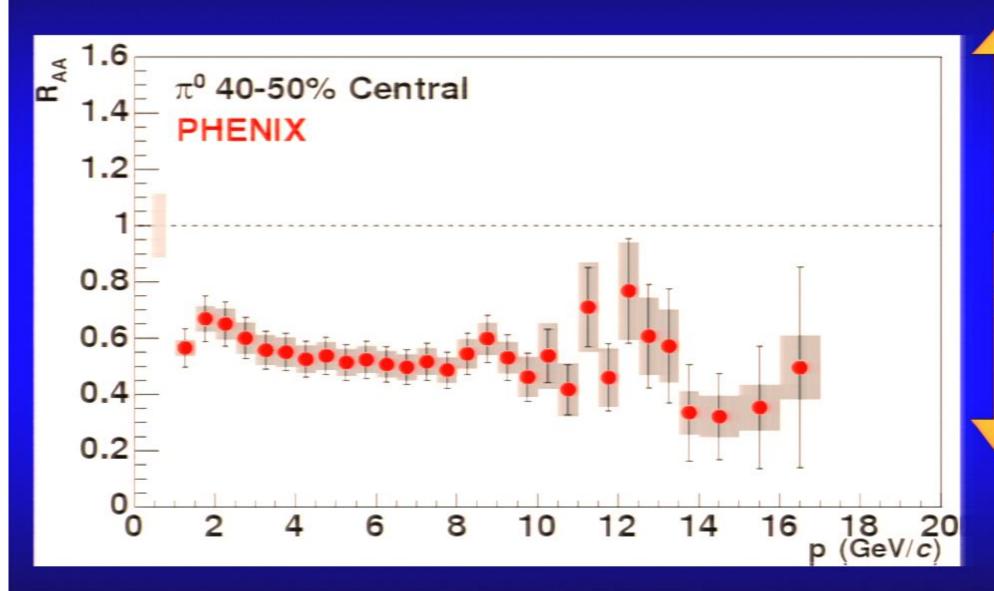
Pirsa: 08060040 Page 106/22-





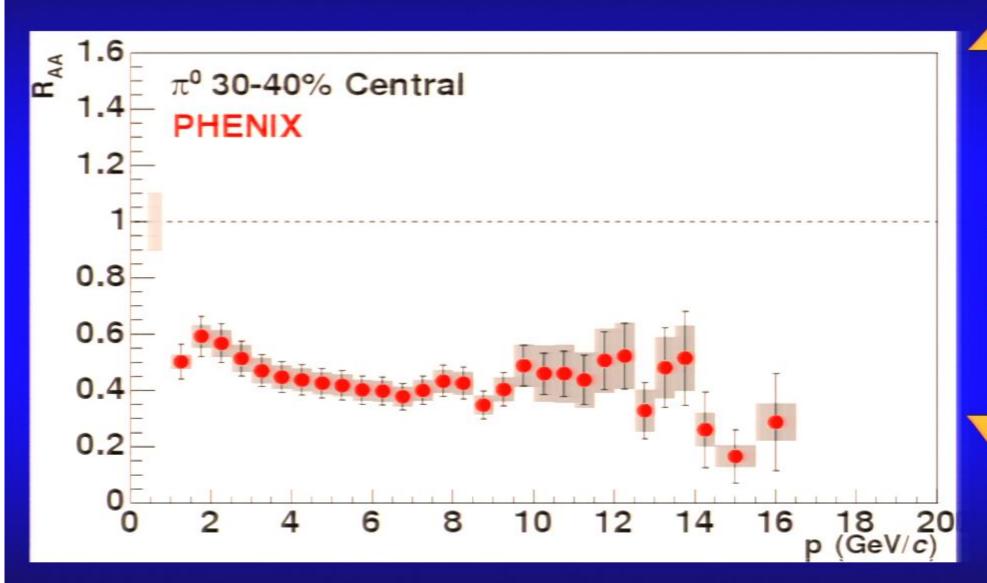
Pirsa: 08060040 Page 107/22-





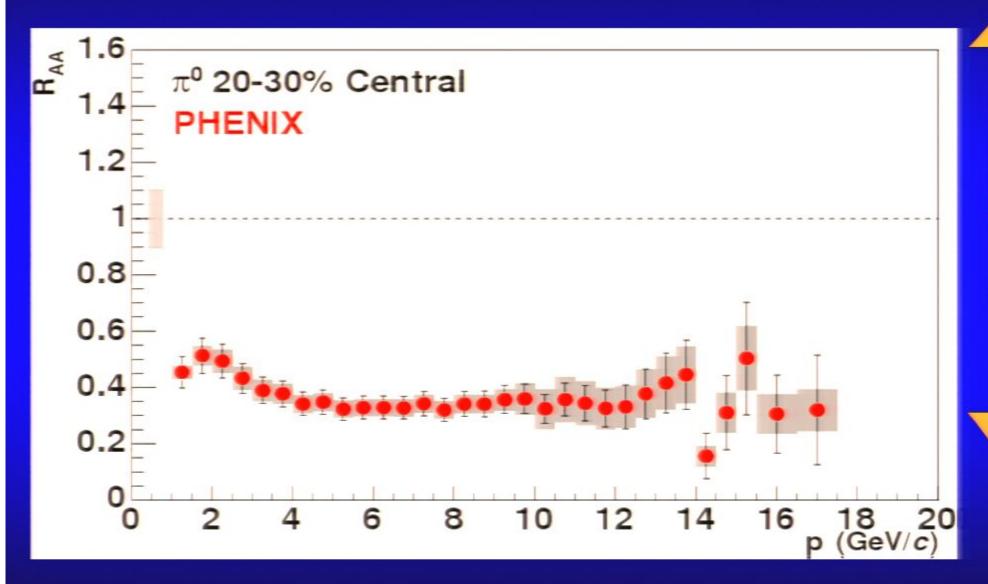
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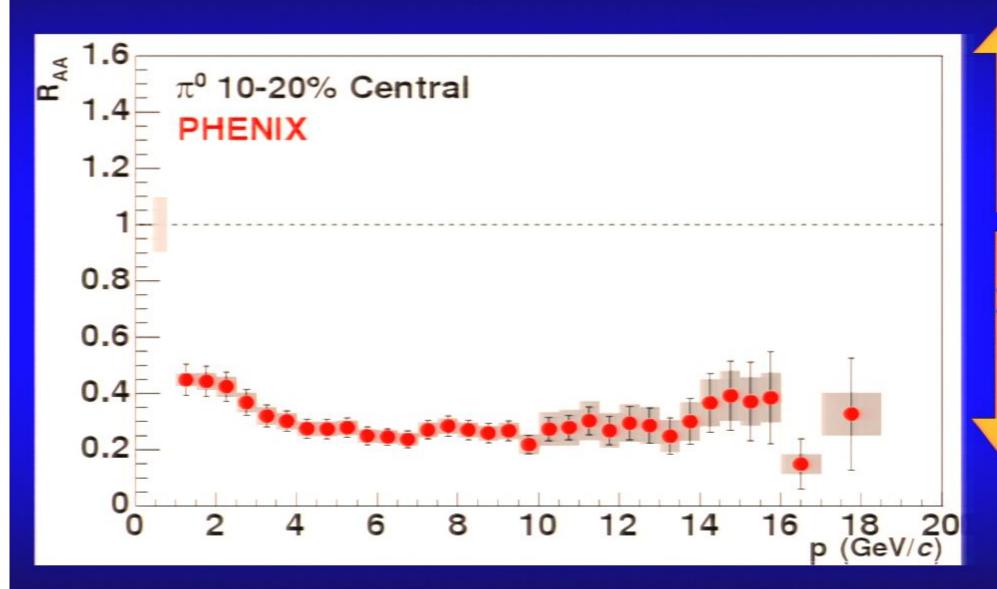
Pirsa: 08060040 Page 109/223





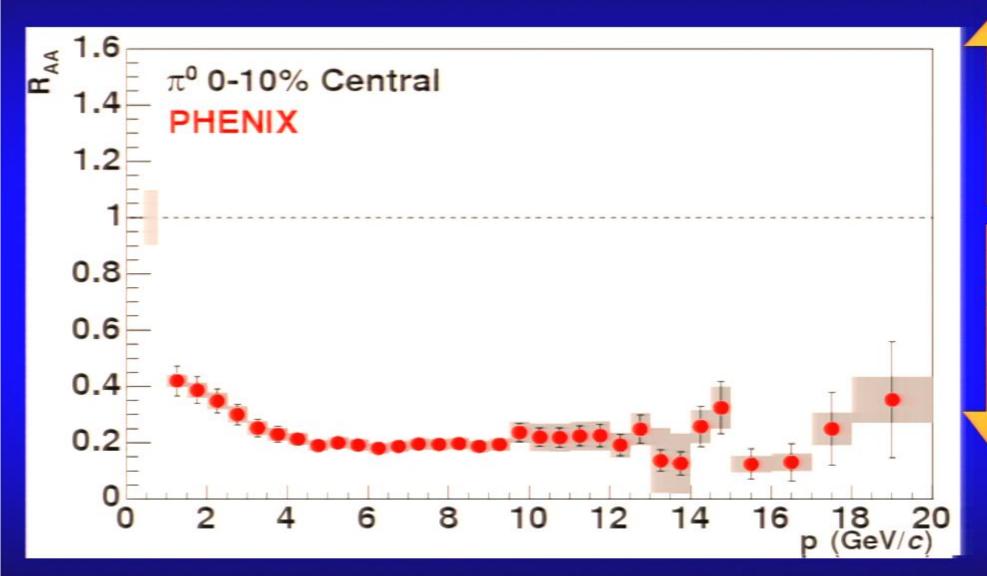
Pirsa: 08060040 Page 110/223





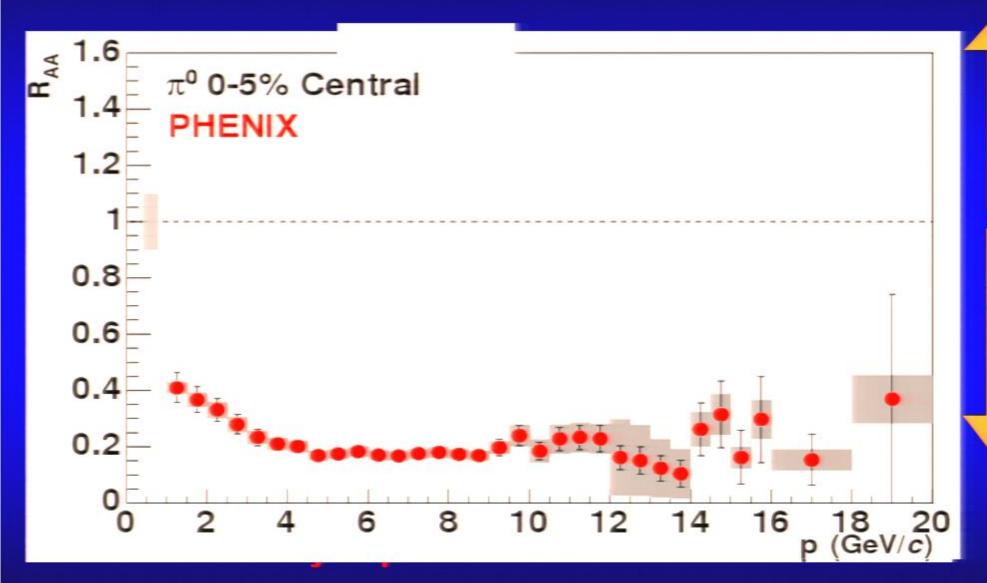
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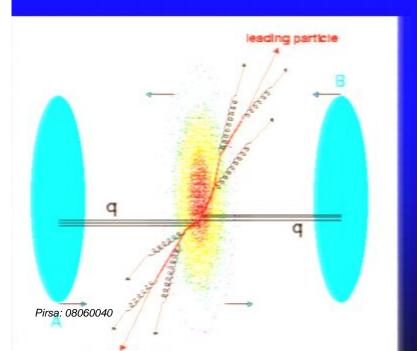


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The Matter is Opaque

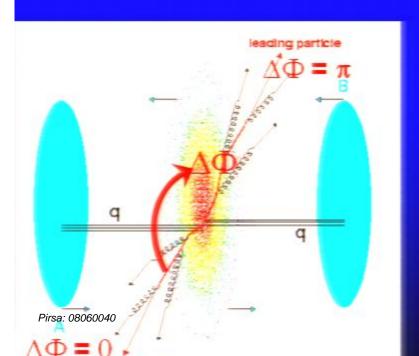
 STAR azimuthal correlation function shows
 complete absence of "away-side" jet





The Matter is Opaque

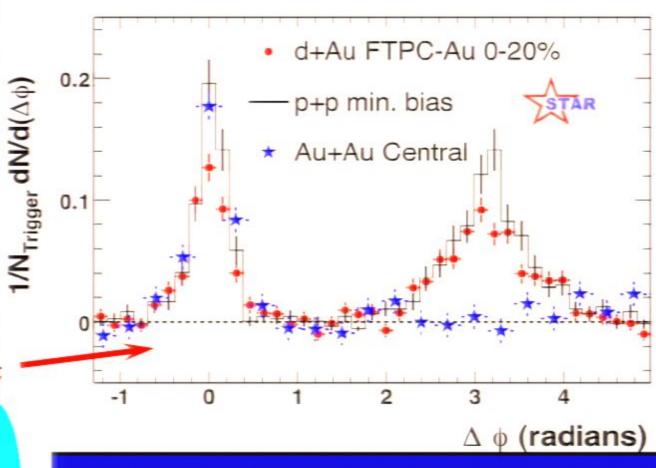
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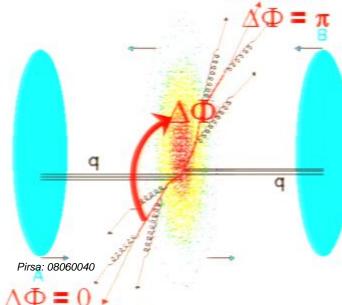




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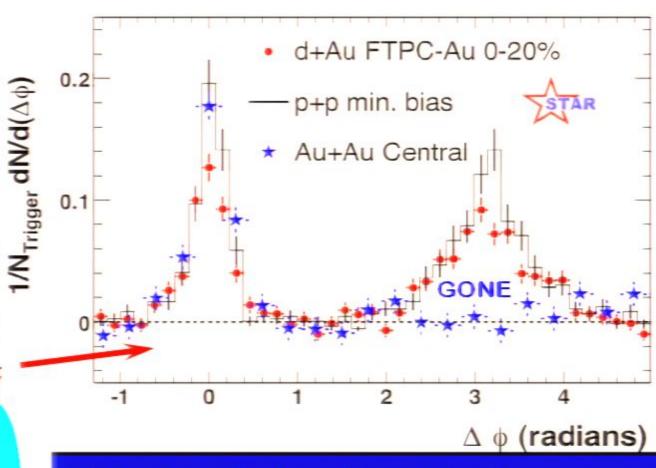




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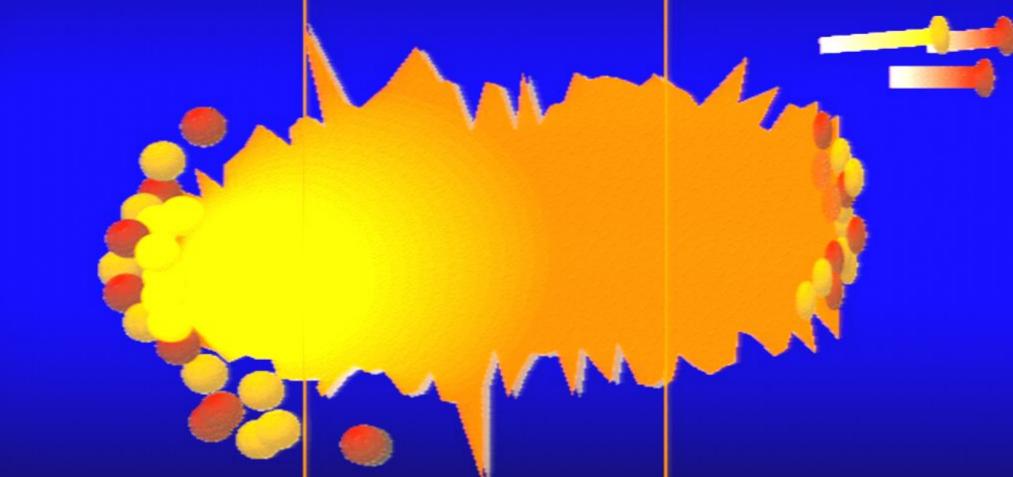


Partner in hard scatter is completely absorbed in the dense medium



Schematically (Partons)

Scattered partons on the "near side" lose energy, but emerge;



those on "far side" are totally absorbed in the *dense* mediu



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Scattered partons on the "near side" lose energy, but emerge;

those on "far side<mark>" are totally absorbed</mark> in the *dense* medium



- Any engineer will tell you
 - □ Kinematic viscosity η / ρ ~ [Velocity] x [Length] is what matters

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- Any engineer will tell you
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Any relativist will tell you

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- Any engineer will tell you
 - Kinematic viscosity η / ρ ~ [Velocity] x [Length] is what matters

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 - $\square \rho \rightarrow \epsilon + p$

Pirea: 08060040



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 - $\Box \epsilon + p = Ts$ (at $\mu_B = 0$)
- So
 - $\Box \eta/\rho \rightarrow \eta/(\epsilon + p) \rightarrow (\eta/sT) = (\eta/s)$

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• The initial discoveries at RHIC

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The initial discoveries at RHIC clearly demonstrated:

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The initial discoveries at RHIC clearly demonstrated:

**Essentially perfect fluid behavior

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The initial discoveries at RHIC clearly demonstrated:

Essentially perfect fluid behavior of extraordinarily dense matter

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The initial discoveries at RHIC clearly demonstrated:

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Matter about as far away from

Pirsa: 08060040 Page 134/223



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Essentially perfect fluid behavior of extraordinarily dense matter

Matter about as far away from "asymptotically free gas" as possible

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The initial discoveries at RHIC clearly demonstrated:

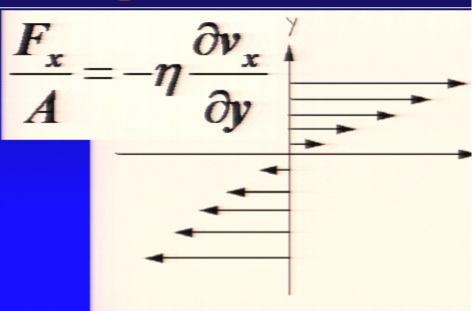
Essentially perfect fluid behavior of extraordinarily dense matter

Matter about as far away from "asymptotically free gas" as possible

How to quantify these statements?



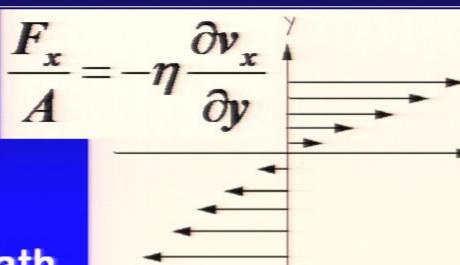
 Parameterizes momentum transport between "layers"



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 Parameterizes momentum transport between "layers"

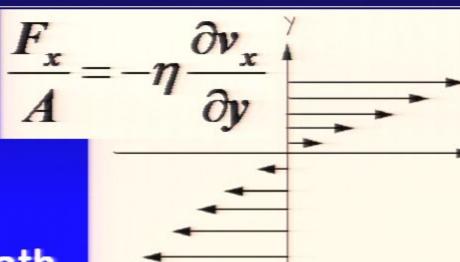


⇒Viscosity ~ mean free path

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 Parameterizes momentum transport between "layers"

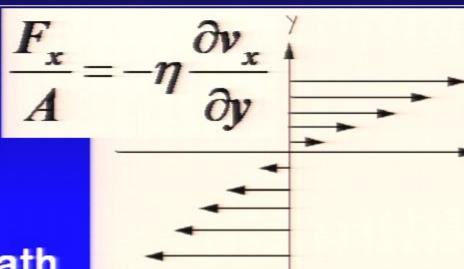


⇒Viscosity ~ mean free path

$$\eta \sim n \ \overline{p} \ \lambda_{mfp}$$



 Parameterizes momentum transport between "layers"



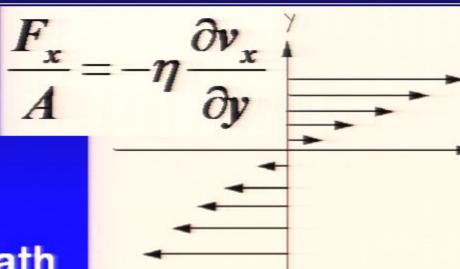
⇒Viscosity ~ mean free path

$$\eta \sim n \ \overline{p} \ \lambda_{mfp}$$

□Small viscosity → Small λ_{mfp}

Pirea: 08060040

 Parameterizes momentum transport between "layers"



⇒Viscosity ~ mean free path

$$\eta \sim n \ \overline{p} \ \lambda_{mfp}$$

- □Small viscosity → Small λ_{mfp}
- \square "Ideal hydro" $\rightarrow \lambda_{mfp} \rightarrow 0 \rightarrow \eta \rightarrow 0$



KSS Bound

"Strong coupling" ⇒ "small viscosity" - but:



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KSS Bound

- "Strong coupling" ⇒ "small viscosity" but:
- "A Viscosity Bound Conjecture",
 P. Kovtun, D.T. Son, A.O. Starinets, hep-th/0405231

$$\frac{\eta}{s} \geq$$





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 "small viscosity" but:
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 P. Kovtun, D.T. Son, A.O. Starinets, hep-th/0405231

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 P. Kovtun, D.T. Son, A.O. Starinets, hep-th/0405231

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Where do "ordinary" fluids sit wrt this limit?



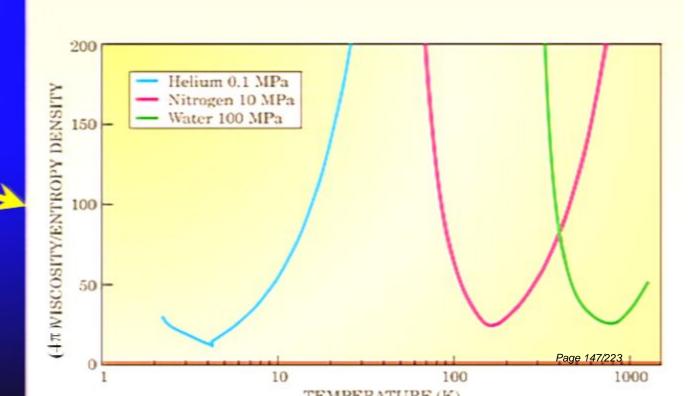


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Where do "ordinary" fluids sit wrt this limit?

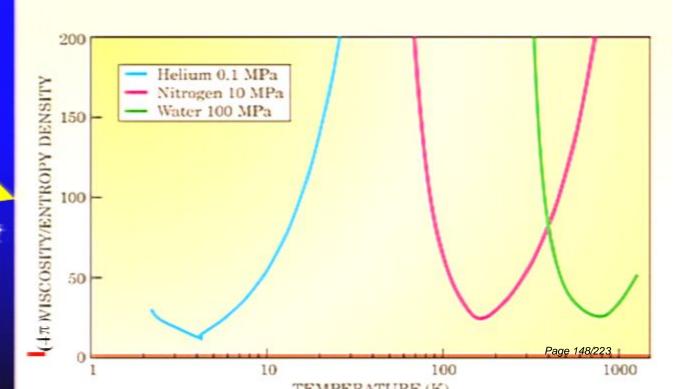




- "Strong coupling" ⇒ "small viscosity" but:
- "A Viscosity Bound Conjecture",
 P. Kovtun, D.T. Son, A.O. Starinets, hep-th/0405231

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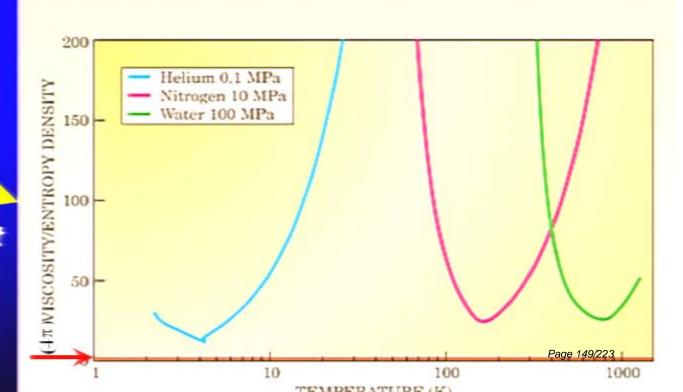




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- Exploit Maldacena's
 - "D-dimensional strongly coupled gauge theory ⇔ (D+1)-dimensional stringy gravity"
- Thermalize with massive black brane
- Calculate viscosity η = "Area"/16πG

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Exploit Maldacena's

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Thermalize with massive black brane

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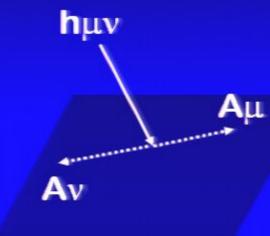


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Exploit Maldacena's

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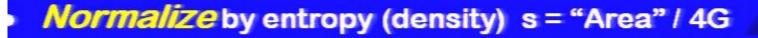


Exploit Maldacena's

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Thermalize with massive black brane

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Dividing out the infinite "areas":

 $\frac{\eta}{s}$:

hμν Αμ Av Infinite "Area"!



Exploit Maldacena's

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Calculate viscosity η = "Area"/16πG



Dividing out the infinite "areas":



$$\frac{\eta}{s} = (\frac{\hbar}{k}) \frac{1}{4k}$$



- **Exploit** Maldacena's
 - "D-dimensional strongly coupled gauge theory ⇔ (D+1)-dimensional stringy gravity"
- Thermalize with massive black brane
- Calculate viscosity η = "Area"/16πG
- Normalize by entropy (density) s = "Area" / 4G
 - Dividing out the infinite "areas":

See next talk: K. Rajagopal



$$\frac{\eta}{s} = (\frac{\hbar}{k}) \frac{1}{4\pi}$$

- Conjectured to be a lower bound "for all relativistic quantum field theories at finite temperature and zero chemical potential".
- See Viscosity in strongly interacting quantum field theories from black hole physics, P. Kovtun, D.T. Son, A.O. Starinets, Phys.Rev.Lett.94:111601, 2005, hep-th/0405231



Damping (flow, fluctuations, heavy quark motion) ~ η/S

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- Damping (flow, fluctuations, heavy quark motion) $\sim \eta/S$
 - □ FLOW: Has the QCD Critical Point Been Signaled by Observations at RHIC?, R. Lacey et al., Phys.Rev.Lett.98:092301,2007 (nucl-ex/0609025)
 - □ The Centrality dependence of Elliptic flow, the Hydrodynamic Limit, and the Viscosity of Hot QCD, H.-J. Drescher et al., (arXiv:0704.3553)

$$\frac{\eta}{s}$$

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$$\frac{\eta}{s} = (1.1 \pm 0.2 \pm 1.2) - \frac{\eta}{2}$$

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$$\frac{\eta}{s} = (1.9 - 2.5)$$

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 Phys.Rev.Lett.97:162302,2006
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$$\frac{\eta}{s} = (1.1 \pm 0.2 \pm 1.2) - \frac{\eta}{s}$$

$$\frac{\eta}{s} = (1.9 - 2.5) \frac{\eta}{4}$$

$$\frac{\eta}{s} = 0$$

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$$\frac{\eta}{s} = (1.0 - 3.$$

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C

Measuring (Estimating) η/s

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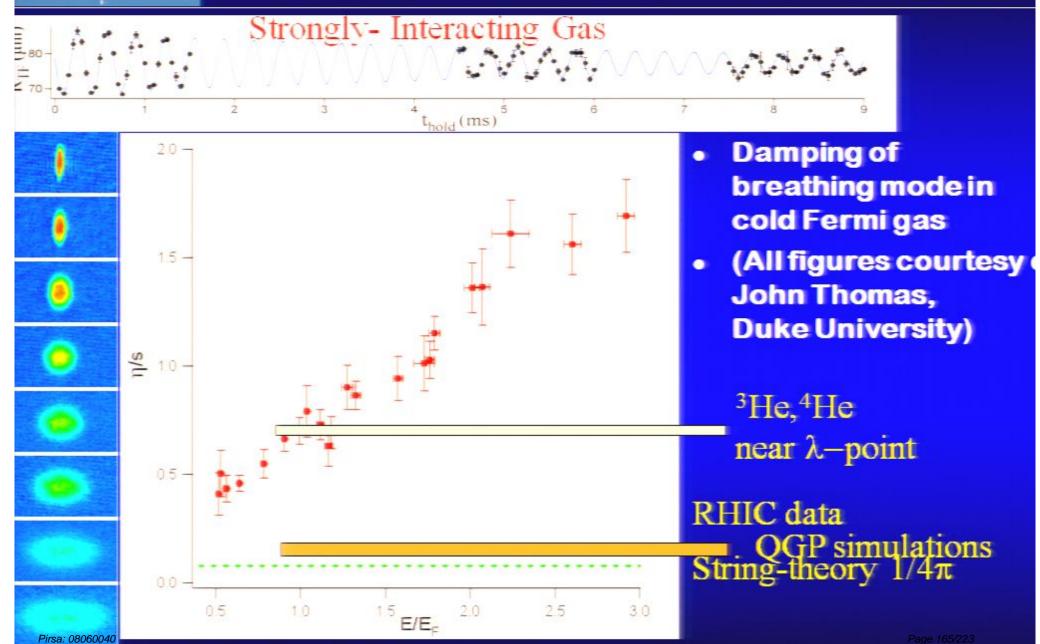
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$$\frac{\eta}{s} = (1.3 - 2.$$



Compare to the Competition





Non-relativistic: Damping given by

$$\dot{E} = -\frac{1}{2} \int d^3x \, \eta(x) \left(\partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right)^2$$

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Non-relativistic: Damping given by

$$\dot{E} = -\frac{1}{2} \int d^3x \, \eta(x) \, \left(\partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right)^2$$

Relativistic:



Non-relativistic: Damping given by

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- Relativistic: Causal, second-order expansion:
 - Relativistic Fluid Dynamics:
 Physics for



Non-relativistic: Damping given by

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Pirea: 08060040



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Working out the divergence of the entropy current, and making use of the equations of motion, we arriv

$$\nabla_{\mu}s^{\mu} = -\frac{1}{T}\tau \left[\nabla_{\mu}u^{\mu} + \beta_{0}u^{\mu}\nabla_{\mu}\tau - \alpha_{0}\nabla_{\mu}q^{\mu} - \gamma_{0}Tq^{\mu}\nabla_{\mu}\left(\frac{\alpha_{0}}{T}\right) + \frac{\tau T}{2}\nabla_{\mu}\left(\frac{\beta_{0}u^{\mu}}{T}\right)\right]$$

$$-\frac{1}{T}q^{\mu}\left[\frac{1}{T}\nabla_{\mu}T + u^{\nu}\nabla_{\nu}u_{\mu} + \beta_{1}u^{\nu}\nabla_{\nu}q_{\mu} - \alpha_{0}\nabla_{\mu}\tau - \alpha_{1}\nabla_{\nu}\tau^{\nu}_{\mu}\right]$$

$$+\frac{T}{2}q_{\mu}\nabla_{\nu}\left(\frac{\beta_{1}u^{\nu}}{T}\right) - (1 - \gamma_{0})\tau T\nabla_{\mu}\left(\frac{\alpha_{0}}{T}\right) - (1 - \gamma_{1})T\tau^{\nu}_{\mu}\nabla_{\nu}\left(\frac{\alpha_{1}}{T}\right)\right]$$

$$-\frac{1}{T}\tau^{\mu\nu}\left[\nabla_{\mu}u_{\nu} + \beta_{2}u^{\alpha}\nabla_{\alpha}\tau_{\mu\nu} - \alpha_{1}\nabla_{\mu}q_{\nu} + \frac{T}{2}\tau_{\mu\nu}\nabla_{\alpha}\left(\frac{\beta_{2}u^{\alpha}}{T}\right) - \gamma_{1}Tq_{\mu}\nabla_{\nu}\left(\frac{\alpha_{1}}{T}\right)\right].$$

In this expression it should be noted that we have introduced (following Lindblom and Hiscock) two furtiparameters, γ_0 and γ_1 . They are needed because without additional assumptions it is not clear how the "mixed" quadratic term should be distributed. A natural way to fix these parameters is to appeal to the Onsager symmetry principle [58.9], which leads to the mixed terms being distributed "equally" and hen $\gamma_0 = \gamma_1 = 1/2$.

Denoting the comoving derivative by a dot i.e. using $u^{\mu}\nabla_{\mu}\tau = \tau$ etc. we see that the second law



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$$+\frac{T}{2}q_{\mu}\nabla_{\nu}\left(\frac{\beta_{1}u^{\nu}}{T}\right) - (1 - \gamma_{0})\tau T\nabla_{\mu}\left(\frac{\alpha_{0}}{T}\right) - (1 - \gamma_{1})T\tau^{\nu}_{\mu}\nabla_{\nu}\left(\frac{\alpha_{1}}{T}\right)\right]$$

$$-\frac{1}{T}\tau^{\mu\nu}\left[\nabla_{\mu}u_{\nu} + \beta_{2}u^{\alpha}\nabla_{\alpha}\tau_{\mu\nu} - \alpha_{1}\nabla_{\mu}q_{\nu} + \frac{T}{2}\tau_{\mu\nu}\nabla_{\alpha}\left(\frac{\beta_{2}u^{\alpha}}{T}\right) - \gamma_{1}Tq_{\mu}\nabla_{\nu}\left(\frac{\alpha_{1}}{T}\right)\right].$$

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Denoting the comoving derivative by a dot, i.e. using $u^{\mu}\nabla_{\mu}\tau=\dot{\tau}$ etc. we see that the second law thermodynamics is satisfied if we choose

$$\tau = -\zeta \left[\nabla_{\mu} u^{\mu} + \beta_0 \dot{\tau} - \alpha_0 \nabla_{\mu} q^{\mu} - \gamma_0 T q^{\mu} \nabla_{\mu} \left(\frac{\alpha_0}{T} \right) + \frac{\tau T}{2} \nabla_{\mu} \left(\frac{\beta_0 u^{\mu}}{T} \right) \right],$$

$$q^{\mu} = -\kappa T \perp^{\mu\nu} \left[\frac{1}{T} \nabla_{\nu} T + \dot{u}_{\nu} + \beta_1 \dot{q}_{\nu} - \alpha_0 \nabla_{\nu} \tau - \alpha_1 \nabla_{\alpha} \tau^{\alpha}_{\nu} + \frac{T}{2} q_{\nu} \nabla_{\alpha} \left(\frac{\beta_1 u^{\alpha}}{T} \right) \right],$$

$$-(1 - \gamma_0) \tau T \nabla_{\nu} \left(\frac{\alpha_0}{T} \right) - (1 - \gamma_1) T \tau^{\alpha}_{\nu} \nabla_{\alpha} \left(\frac{\alpha_1}{T} \right) + \gamma_2 \nabla_{[\nu} u_{\alpha]} q^{\alpha} \right],$$

$$\tau_{\mu\nu} = -2\eta \left[\beta_2 \dot{\tau}_{\mu\nu} + \frac{T}{2} \tau_{\mu\nu} \nabla_{\alpha} \left(\frac{\beta_2 u^{\alpha}}{T} \right) + \left\langle \nabla_{\mu} u_{\nu} - \alpha_1 \nabla_{\mu} q_{\nu} - \gamma_1 T q_{\mu} \nabla_{\nu} \left(\frac{\alpha_1}{T} \right) + \gamma_3 \nabla_{[\mu} u_{\alpha]} \tau_{\nu}^{\alpha} \right\rangle \right],$$

where the angular brackets denote symmetrization as before. In these expression we have added yet another two terms, representing the coupling to vorticity. These bring further "fpage 171/223ers" and is easy to see that we are allowed to add these terms since they do not affect the entropy production. It fact, a large number of similar terms may, in principle, be considered (see note added in proof in [53.9]). The presence of coupling terms of the particular form that we have introduced is suggested by kinetic.



Non-relativistic: Damping given by

$$\dot{E} = -\frac{1}{2} \int d^3x \, \eta(x) \, \left(\partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right)^2$$

- Relativistic: Causal, second-order expansion:
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 Physics for
 Many Different Scales
- Neglect various terms at your own risk:
 - Natsuume and Okamura,
 Comment on
 "Viscous hydrodynamics

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In this expression it should be noted that we have introduced (following Lindblom and Hiscock) two furtiparameters, 70 and 71. They are needed because without additional assumptions it is not clear how the "mixed" quadratic term should be distributed. A natural way to fix these parameters is to appeal to the Onsager symmetry principle [58.9], which leads to the mixed terms being distributed "equally" and hen $\gamma_0 = \gamma_1 = 1/2$.

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In this expression it should be noted that we have introduced (following Lindblom and Hiscock) two furtiparameters, γ_0 and γ_1 . They are needed because without additional assumptions it is not clear how the "mixed" quadratic term should be distributed. A natural way to fix these parameters is to appeal to the Onsager symmetry principle [58.9], which leads to the mixed terms being distributed "equally" and hen $\gamma_0 = \gamma_1 = 1/2$.

Denoting the comoving derivative by a dot, i.e. using $u^{\mu}\nabla_{\mu}\tau=\dot{\tau}$ etc. we see that the second law thermodynamics is satisfied if we choose

$$\begin{split} \tau &= -\zeta \left[\nabla_{\mu} u^{\mu} + \beta_{0} \dot{\tau} - \alpha_{0} \nabla_{\mu} q^{\mu} - \gamma_{0} T q^{\mu} \nabla_{\mu} \left(\frac{\alpha_{0}}{T} \right) + \frac{\tau T}{2} \nabla_{\mu} \left(\frac{\beta_{0} u^{\mu}}{T} \right) \right], \\ q^{\mu} &= -\kappa T \perp^{\mu\nu} \left[\frac{1}{T} \nabla_{\nu} T + \dot{u}_{\nu} + \beta_{1} \dot{q}_{\nu} - \alpha_{0} \nabla_{\nu} \tau - \alpha_{1} \nabla_{\alpha} \tau^{\alpha}_{\nu} + \frac{T}{2} q_{\nu} \nabla_{\alpha} \left(\frac{\beta_{1} u^{\alpha}}{T} \right) \right. \\ & \left. - (1 - \gamma_{0}) \tau T \nabla_{\nu} \left(\frac{\alpha_{0}}{T} \right) - (1 - \gamma_{1}) T \tau^{\alpha}_{\nu} \nabla_{\alpha} \left(\frac{\alpha_{1}}{T} \right) + \gamma_{2} \nabla_{[\nu} u_{\alpha]} q^{\alpha} \right], \\ \tau_{\mu\nu} &= -2\eta \left[\beta_{2} \dot{\tau}_{\mu\nu} + \frac{T}{2} \tau_{\mu\nu} \nabla_{\alpha} \left(\frac{\beta_{2} u^{\alpha}}{T} \right) + \left\langle \nabla_{\mu} u_{\nu} - \alpha_{1} \nabla_{\mu} q_{\nu} - \gamma_{1} T q_{\mu} \nabla_{\nu} \left(\frac{\alpha_{1}}{T} \right) + \gamma_{3} \nabla_{[\mu} u_{\alpha]} \tau_{\nu}^{\alpha} \right] \right]. \end{split}$$

where the angular brackets denote symmetrization as before. In these expression we have added yet another two terms, representing the coupling to vorticity. These bring further 'fipage 173/223 ers' 2 and is easy to see that we are allowed to add these terms since they do not affect the entropy production. It fact, a large number of similar terms may, in principle, be considered (see note added in proof in [53.9]. The presence of coupling terms of the particular form that we have introduced is suggested by kinetic.



Viscosity Information from Relativistic Nuclear Collisions: How Perfect is the Fluid Observed at RHIC?, P. Romatschke and U. Romatschke, Phys. Rev. Lett. 99:172301, 2007

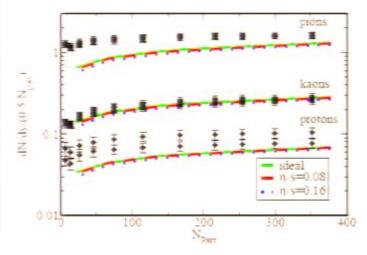
- Signature: dN/dy, v₂, <p_T>
- Calculation: 2nd order causal viscous hydro:

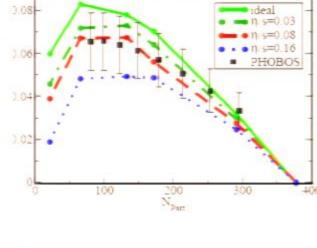
$$\begin{split} (\epsilon + p)Du^{\mu} &= \nabla^{\mu}p - \Delta^{\mu}_{\alpha}d_{\beta}\Pi^{\alpha\beta}, \\ D\epsilon &= -(\epsilon + p)\nabla_{\mu}u^{\mu} + \frac{1}{2}\Pi^{\mu\nu}\nabla_{\nu}u_{\mu}, \\ \Delta^{\mu}_{\alpha}\Delta^{\nu}_{\beta}D\Pi^{\alpha\beta} &= -\frac{\Pi^{\mu\nu}}{\gamma_{\Pi}} + \frac{\eta}{\gamma_{\Pi}}\nabla^{\mu}u^{\nu} - 2\Pi^{\alpha(\mu}\omega^{\nu)}_{\alpha} \\ &+ \frac{1}{2}\Pi^{\mu\nu}\left[5D\ln T - \nabla_{\alpha}u^{\alpha}\right]. \end{split} \tag{2}$$

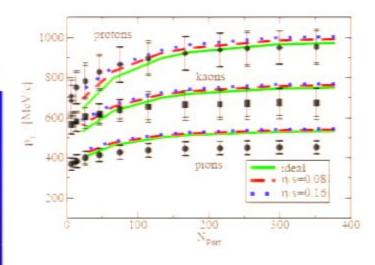
(Glauber IC's)

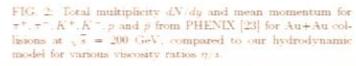
Payoff Plots:

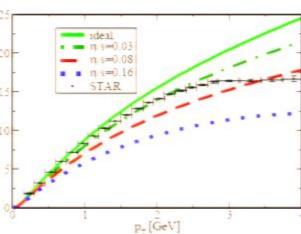
$$\Rightarrow \frac{\eta}{s} = (0 - 2.0) \frac{1}{4\pi}$$











7. 3. PHOBOS [24] data on ρ_T integrated ρ₂ and STAR data on minimum bias ρ₂, for charged particles in Au+Au issons at √s = 200 GeV, compared to our hydrodynamic del for various viscosity ratios η_cs. Error bars for PHO-S data show 90% confidence level systematic errors while STAR only statistical errors are shown.



- So-RHIC flow data within factors 1-4
 of the bound on η/s calculated via AdS/CFT
- Can AdS/CFT also be applied to the strong energy loss in the fluid?

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$$\hat{\boldsymbol{q}} \sim \frac{\langle \boldsymbol{p_T}^2 \rangle}{\lambda}$$

$$\rightarrow \Delta E_{rad} \approx \alpha_s \dot{q}$$



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 - Dainese, Loizides, Paic: hep-ph/04062

Pirsa: 08060040 Page 178/22

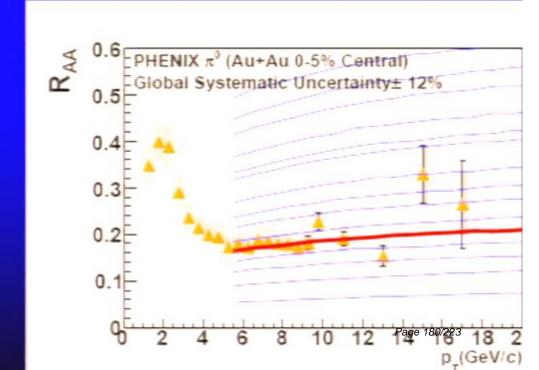


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- Detailed statistical analysis of energy loss systematics
 - ☐ Dainese, Loizides, Paic: hep-ph/0406201
 - PHENIX: arXiv:0801.166

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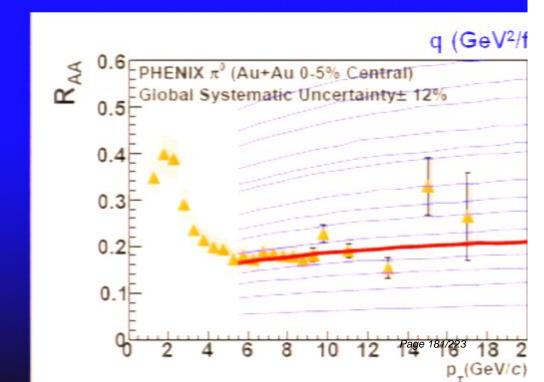
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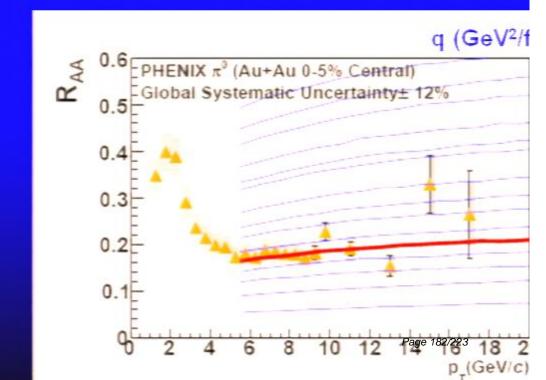
$$\hat{q} = 13.2_{-21}^{+2.3} \text{ G}$$





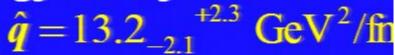
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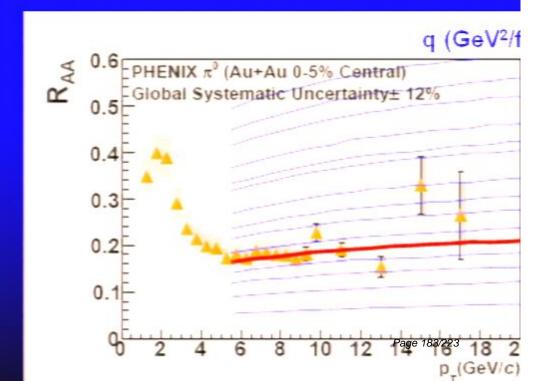
$$\hat{q} = 13.2_{-2.1}^{+2.3} \text{ GeV}^2/\text{fin}$$





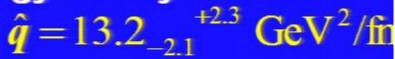
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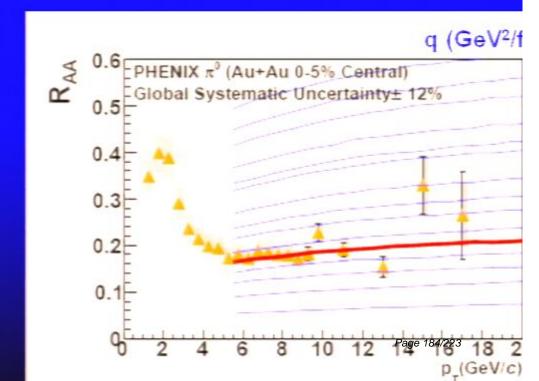






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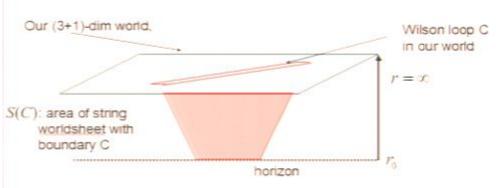
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AdS/CFT?:

Wilson loop from AdS/CFT

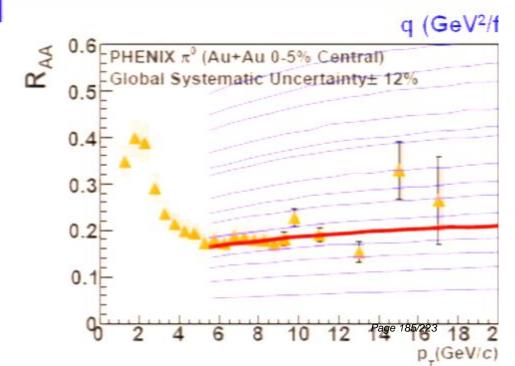
Maidacena (1998): Rey and Yee (1998)

Recipe:
$$\langle W(\mathcal{C}) \rangle = \exp[iS(\mathcal{C})]$$



Pirsa: 08060040 Black hole in AdS spacetime:

- radial coordinate r.



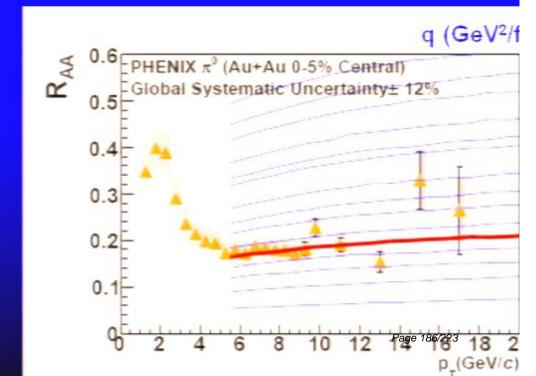


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- AdS/CFT?:
 - Liu, Rajagopal, Wiedemann: hep-ph/0605178

$$\hat{q} = \sqrt{\lambda} \, \pi^{3/2} \frac{\Gamma(3/4)}{\Gamma(5/4)} T^3$$

 $\sim 4.5 \pm ? \text{ GeV}^2/\text{fm}$

 $\hat{q} = 13.2_{-2.1}^{+2.3} \text{ GeV}^2/\text{fin}$



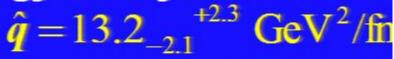


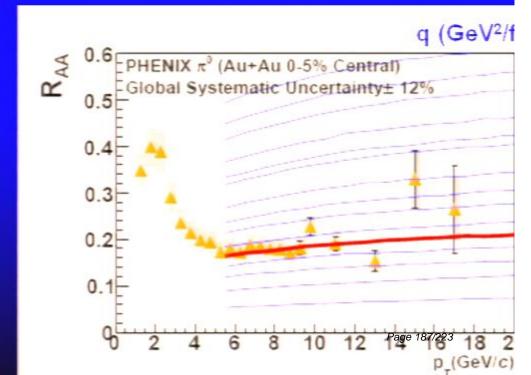
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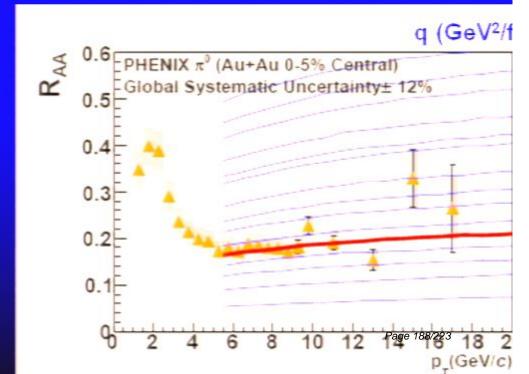


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Factor of ~3 difference:



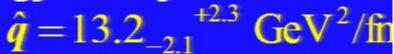


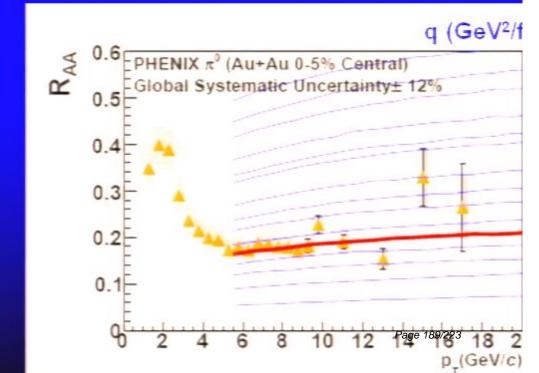
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- Factor of ~3 difference:
 - □ AdS/CFT excluded at 3σ?







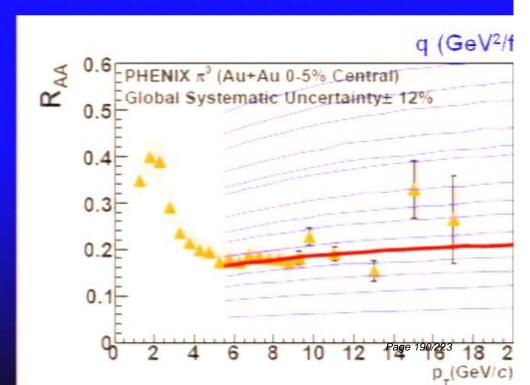
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Suc





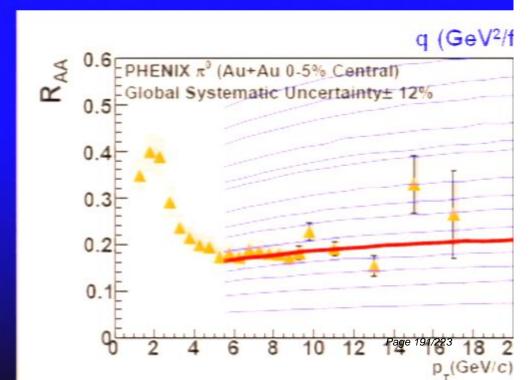
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□ Success?





While the (conjectured) bound $\frac{\eta}{2} > \frac{\hbar}{2}$

$$\frac{\eta}{s} \ge \frac{\hbar}{4\pi}$$

is a purely quantum mechanical result . . .

The state of the s the Anti-de Sitter space / Conformal Field Theory correspondence

- Weak form:
 - "Four-dimensional N=4 supersymmetric SU(N_c) gauge theory is equivalent to type IIB string theory with AdS₅ x S⁵ boundary conditions." (The Large N limit of superconformal field theories and supergravity, J. Maldacena, Adv. Theor. Math. Phys. 2, 231, 1998 hep-th/9711200)



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- Strongest form: Only with QCD can we explore experimentally these fascinating connections over the full range of the coupling constant to study QGP = Quantum Gauge Phluid
- Uber-Strongest form: "Susskind says that by studying heavy



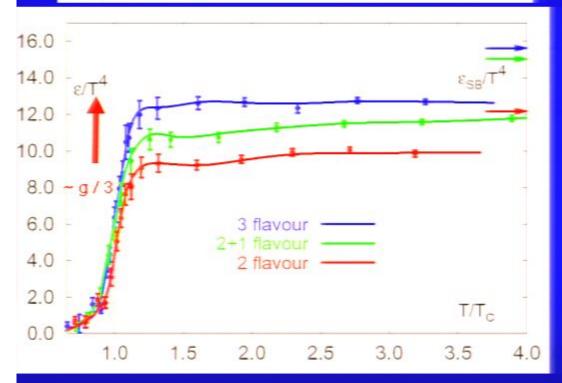
- While the (conjectured) bound $\frac{\eta}{s} \ge \frac{\hbar}{4\pi}$
 - is a purely quantum mechanical result . . .
- It was derived in and motivated by the Anti-de Sitter space / Conformal Field Theory correspondence
- Weak form:
 - "Four-dimensional N=4 supersymmetric SU(N_c) gauge theory is equivalent to type IIB string theory with AdS₅ x S⁵ boundary conditions." (The Large N limit of superconformal field theories and supergravity, J. Maldacena, Adv. Theor. Math. Phys. 2, 231, 1998 hep-th/9711200)
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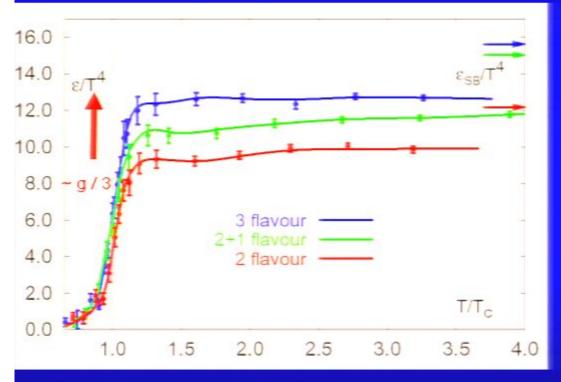
$$\frac{\mathcal{E}(T > T_C)}{\mathcal{E}_{Stefan-Boltzmann}} = 80\% \Rightarrow \text{almost}$$
 asymptotic ally free



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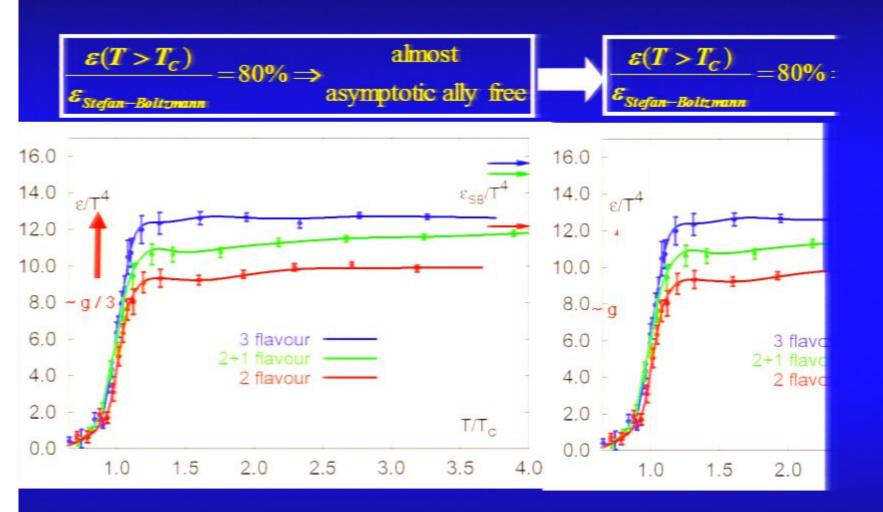






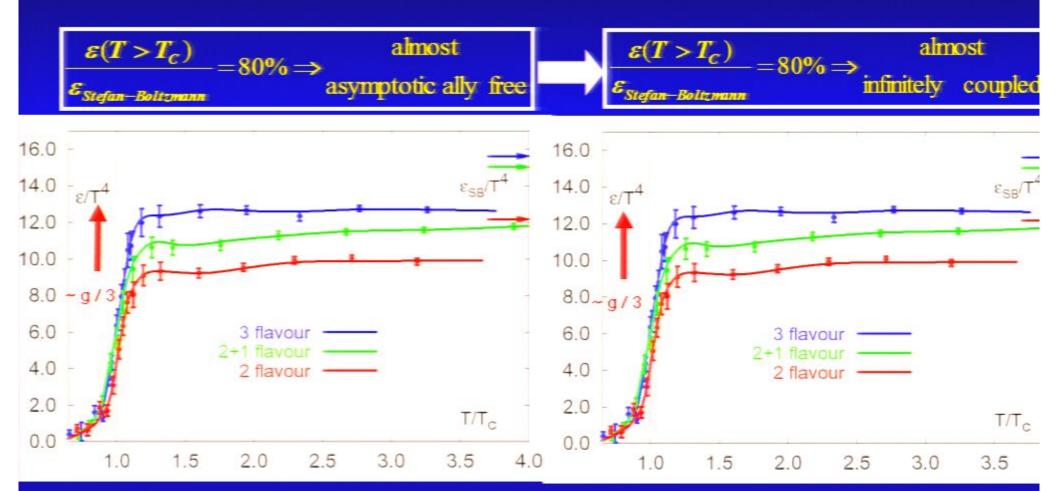
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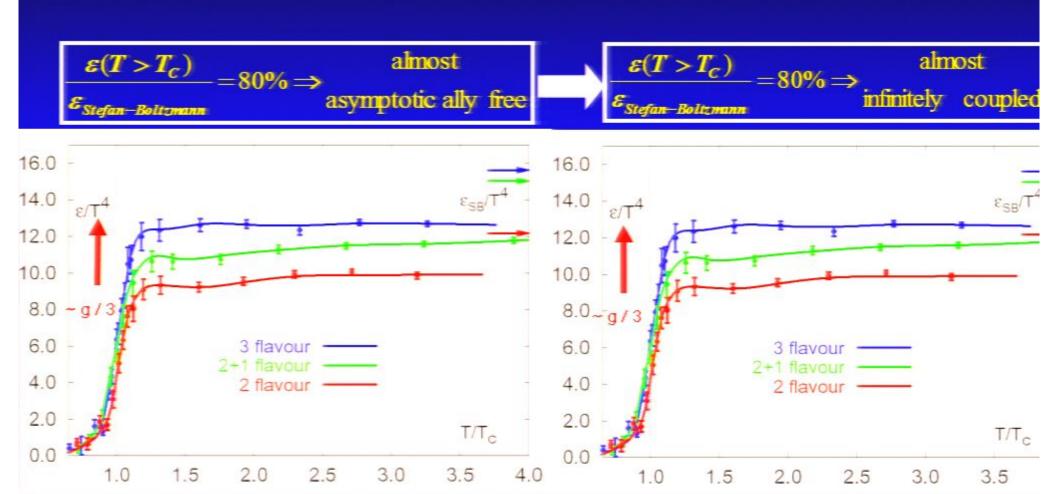
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'Strongly' motivated by famous Gubser, Klebanov, Peet result

$$\frac{s(\lambda=\infty)}{s(\lambda=0)}$$



$$\begin{array}{c} \boldsymbol{\varepsilon}(T > T_C) \\ \boldsymbol{\varepsilon}_{Stefan-Boltzmann} \end{array} = 80\% \Rightarrow \begin{array}{c} \text{almost} \\ \text{asymptotic ally free} \end{array} \qquad \begin{array}{c} \boldsymbol{\varepsilon}(T > T_C) \\ \boldsymbol{\varepsilon}_{Stefan-Boltzmann} \end{array} = 80\% \Rightarrow \begin{array}{c} \text{almost} \\ \text{infinitely coupled} \end{array}$$

'Strongly' motivated by famous Gubser, Klebanov, Peet result

3.5

$$\frac{s(\lambda = \infty)}{s(\lambda = 0)} = \frac{3}{4} = 75\%$$

0.0

1.0

1.5

3.0

3.5

2.5

2.0

1.0

1.5

2.0

2.5

3.0

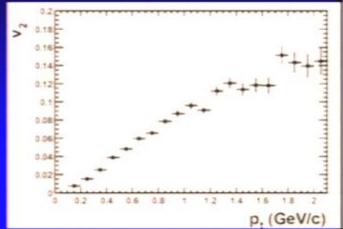


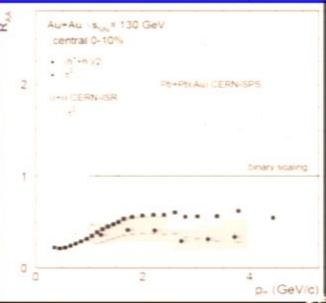
Visual Summary

EACH

RHIC's Two Major Discoveries

- Discovery of strong "elliptic" flow:
 - □ Elliptic flow in Au + Au collisions at √SNA = 130 GeV,
 STAR Collaboration, (K.H.
 Ackermann et al.).
 Phys.Rev.Lett.86:402-407,2001
 - □ 345 citations
- Discovery of "jet quenching"
 - Suppression of hadrons with large transverse momentum in central Au+Au collisions at √s_{NM} = 130 GeV, PHENIX Collaboration (<u>K. Adcox et</u> al.), Phys.Rev.Lett.88:022301,2002
 - 429 citations





□ Indicative of strongl coupled "QGP"



Visual Summary

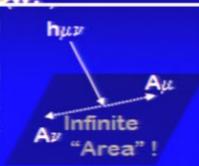
EACH

RHIC's Two Major Discoveries

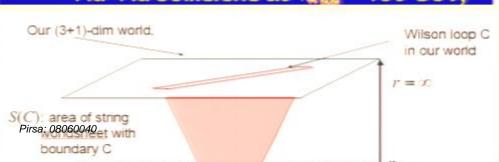
Discovery of strong "elliptic" flow:

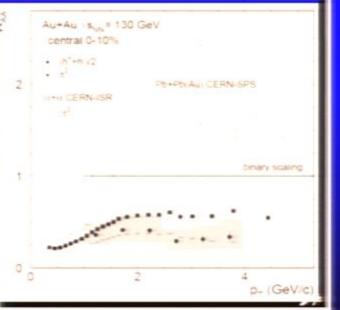
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Indicative of strong!
 coupled "QGP"

□ Consister
to within
factors of
2-3 with
AdS/CFT
results

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- The near-simultaneous development of
 - □ RHIC Discoveries
 - □ Theoretical Synthesis
 - □ Overthrow of "classical" QGP
 - □ Emergence of AdS/CFT connections

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- The near-simultaneous development of
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- The near-simultaneous development of
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represents a

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- The near-simultaneous development of
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represents a true para

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- The near-simultaneous development of
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represents a true paradigm &

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- The near-simultaneous development of
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in (

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- The near-simultaneous development of
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represents a true paradigm shift

in our understanding of

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represents a true paradigm shift

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See next talk by Krishna Rajagopal for the theoretical details

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No Signal VGA-1

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