

Title: Gravitino Dark Matter and Cosmological Constraints with a stau NLSP

Date: May 31, 2008 02:00 PM

URL: <http://pirsa.org/08050055>

Abstract:

Outline

- Thermal Gravitino Production
- Upper Bound on T_R from Thermal Production
- Gravitino Production from Decays
- Constraints from Big Bang Nucleosynthesis
- Constraints from Catalyzed Big Bang Nucleosynthesis
- Late-time Entropy Production
- Conclusion

From SUSY to SUGRA

QED:

$$S_0 = \int d^4x \bar{\psi}(x) i\gamma^\mu \partial_\mu \psi(x)$$

local gauge transformation:

$$\psi(x) \rightarrow e^{i\alpha(x)} \psi(x)$$

action not invariant:

$$\delta S_0 \sim \int d^4x J^\mu(x) \partial_\mu \alpha(x)$$

with $J^\mu = \bar{\psi}(x) \gamma^\mu \psi(x)$.

photon $A_\mu(x)$: $\delta A_\mu(x) \sim \partial_\mu \alpha(x)$

$$S = S_0 + \int d^4x J^\mu(x) A_\mu(x)$$

invariant under *local* U(1) transf.!

SUGRA:

$S_0 =$ invariant under global SUSY

local SUSY transformation:

$$\Phi(y, \theta, \bar{\theta}) \rightarrow e^{i(\bar{\varepsilon}(x)Q + \bar{Q}\varepsilon(x))} \Phi(y, \theta, \bar{\theta}) e^{-i(\dots)}$$

action not invariant:

$$\delta S_0 \sim \int d^4x \bar{J}_r^\mu(x) \partial_\mu \varepsilon_r(x)$$

where J_r^μ is the supercurrent.

gravitino $\psi_r^\mu(x)$: $\delta \psi_r^\mu(x) \sim M_{\text{Pl}} \partial_\mu \varepsilon_r(x)$

$$S = S_0 + \frac{1}{2M_{\text{Pl}}} \int d^4x \bar{J}^\mu(x) \psi_\mu(x) + \dots$$

invariant under *local* SUSY transf.!

Properties of the Gravitino \tilde{G}

- \tilde{G} is the gauge field of *local* SUSY (=SUGRA) transformations
- superpartner of graviton, spin 3/2 Majorana field

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 - goldstino becomes helicity $\pm 1/2$ components of \tilde{G}
 - depending on breaking: $10 \text{ eV} \lesssim m_{\tilde{G}} \lesssim 100 \text{ TeV}$
 - softly broken global SUSY (e.g. MSSM) + \tilde{G} interactions

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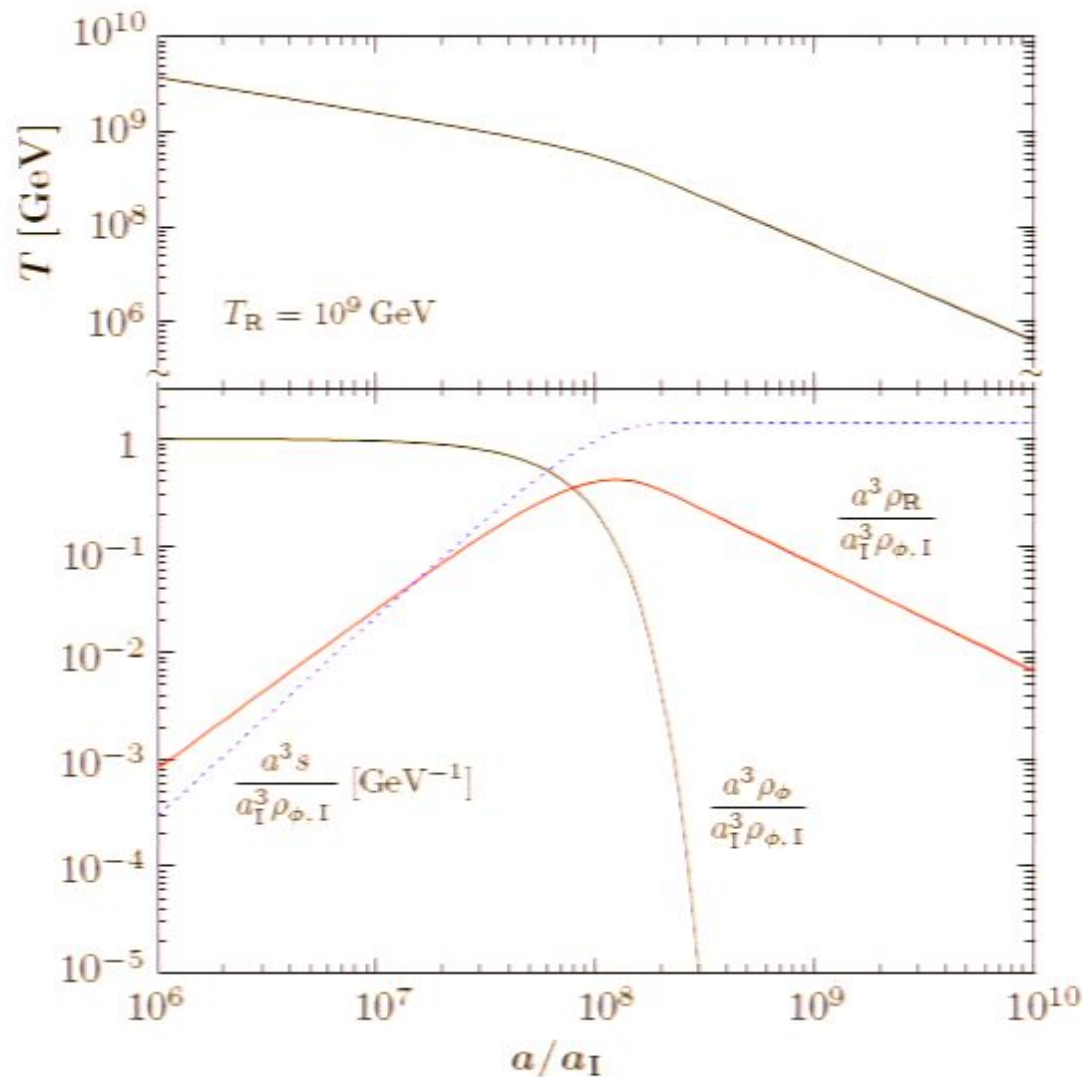
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→ couplings fixed by the symmetry, e.g.

$$-\frac{i}{4M_P} \delta_{ab} [\not{P}, \gamma^\rho] \gamma^\mu$$

$$-\frac{ig_\alpha}{\sqrt{2}M_P} T_{a..ji}^{(\alpha)} P_R \gamma^\rho \gamma^\mu$$

Reheating



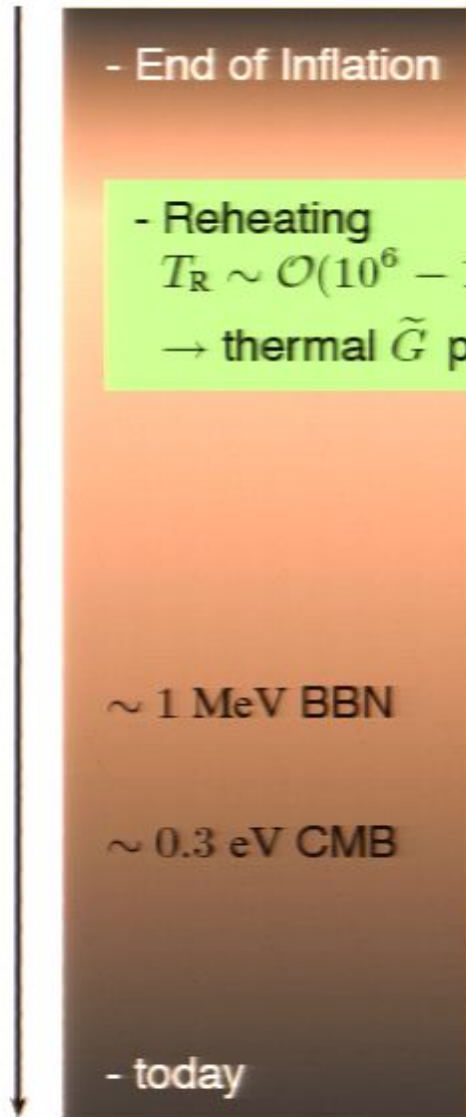
- Inflaton decays

$$\frac{d\rho_{\phi}}{dt} + 3H\rho_{\phi} = -\Gamma_{\phi}\rho_{\phi}$$

- Universe enters radiation dominated epoch

$$\frac{d\rho_{\text{rad}}}{dt} + 4H\rho_{\text{rad}} = \Gamma_{\phi}\rho_{\phi}$$

Thermal Gravitino Production - Framework



- time evolution of \tilde{G} -density governed by the Boltzmann equation

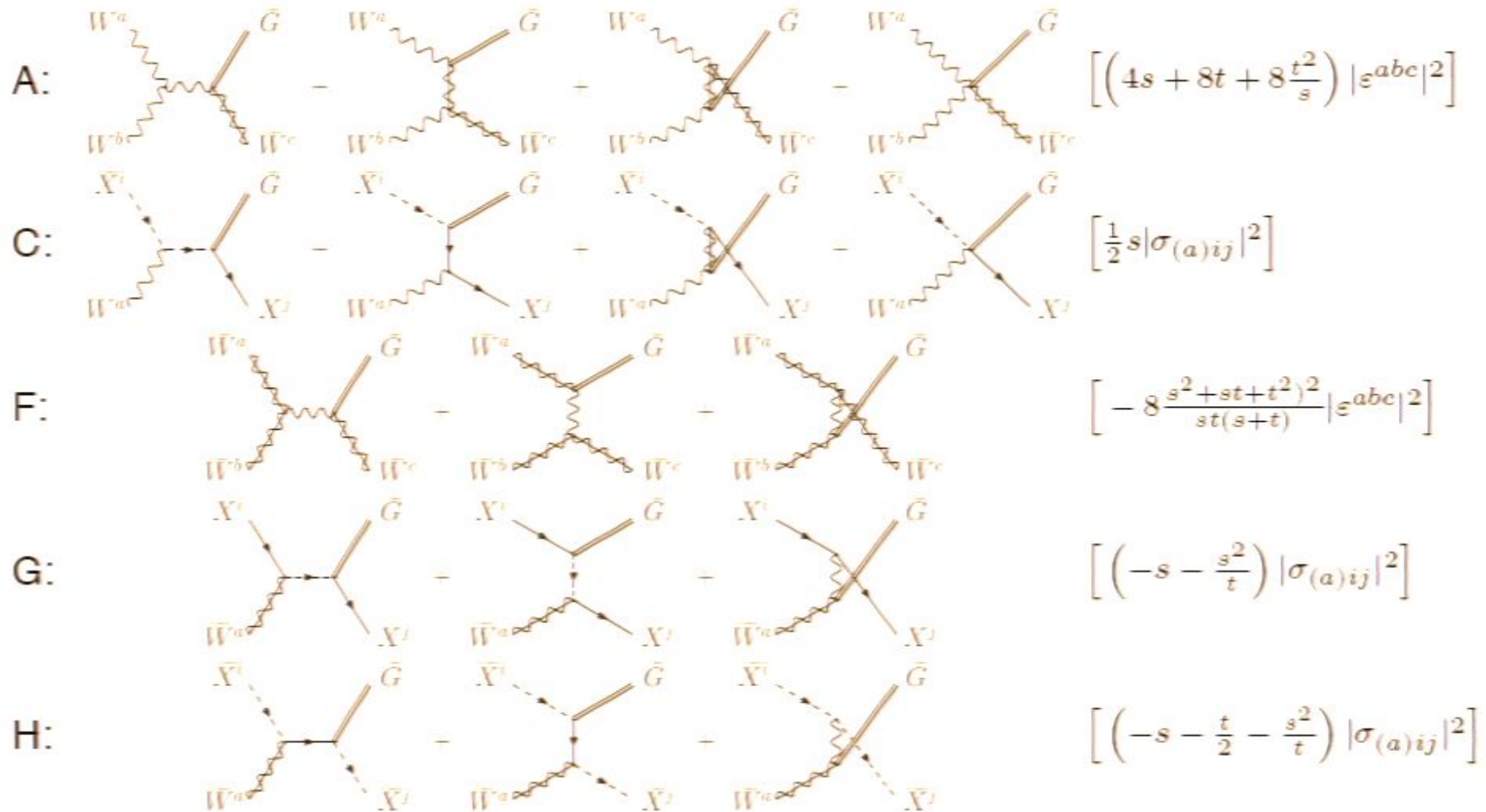
$$\frac{dn_{\tilde{G}}}{dt} + 3Hn_{\tilde{G}} = C_{\tilde{G}}$$

- $C_{\tilde{G}}$: $2 \rightarrow 2$ scatterings + finite T approach



$SU(2)_L$ Contributions

$$|\mathcal{M}_i|^2 = \frac{g_2^2}{M_P^2} \left(1 + \frac{M_2^2}{3m_{\tilde{G}}^2} \right) [\dots]$$



+ crossings \rightarrow in total 10 processes

$SU(2)_L$ Contributions

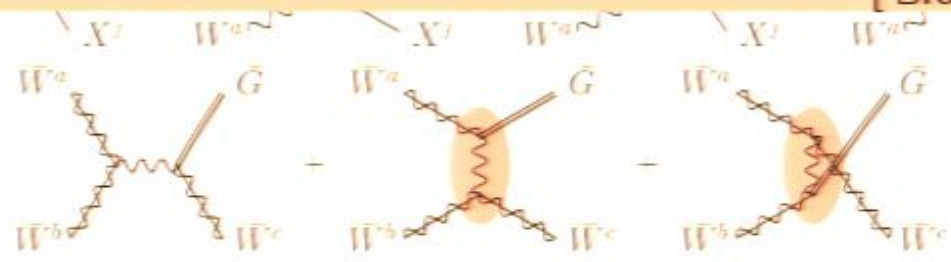
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A: logarithmic singularity $\left(\frac{t^2}{s} \right) |\epsilon^{abc}|^2$


→ regularized screening effects of the plasma

C: → Hard Thermal Loop (HTL) resummed propagator

[Braaten, Pisarski, 1990]

F:  $\left[-8 \frac{s^2 + st + t^2}{s t (s+t)} |\epsilon^{abc}|^2 \right]$

G:  $\left[\left(-s - \frac{s^2}{t} \right) |\sigma_{(a)ij}|^2 \right]$

H:  $\left[\left(-s - \frac{t}{2} - \frac{s^2}{t} \right) |\sigma_{(a)ij}|^2 \right]$

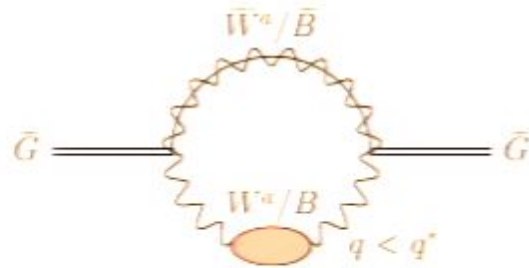
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seperate scales

$$\text{soft} \leftarrow gT \ll q^* \ll T \rightarrow \text{hard}$$

seperate scales
 soft $\leftarrow gT \ll q^* \ll T \rightarrow$ hard

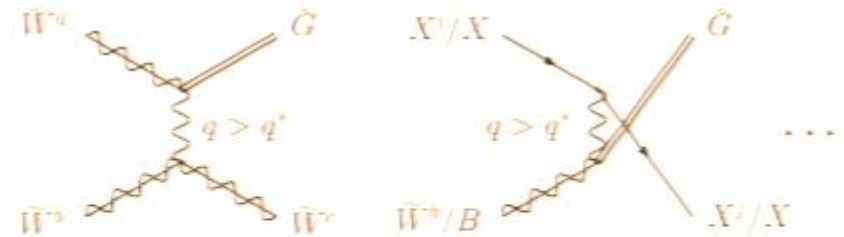
- Soft Part:



$$\Gamma_{\tilde{G}}^{\text{soft}}(E) = - \frac{\text{Im}\Sigma_{\tilde{G}}(E + i\varepsilon, \mathbf{p})}{E} \Big|_{|\mathbf{p}_1 - \mathbf{p}_3| < q^*}$$

$$= A_{\text{soft}} + B \ln \left(\frac{q^*}{m_{W/B}^{\text{therm}}} \right)$$

- Hard Part:

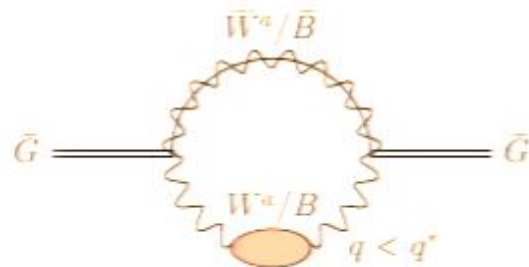


$$\Gamma_{\tilde{G}}^{\text{hard}}(E) = \int \dots |\mathcal{M}|^2 \theta(|\mathbf{p}_1 - \mathbf{p}_3| - q^*)$$

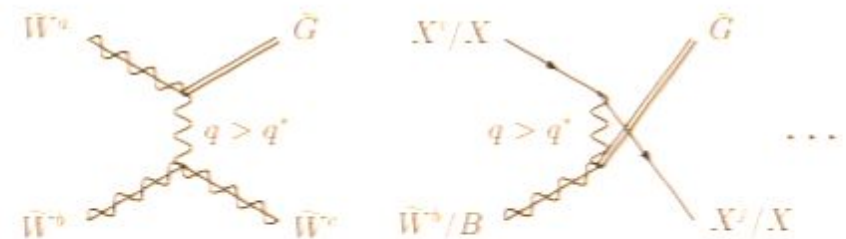
$$= A_{\text{hard}} + B \ln \left(\frac{T}{q^*} \right)$$

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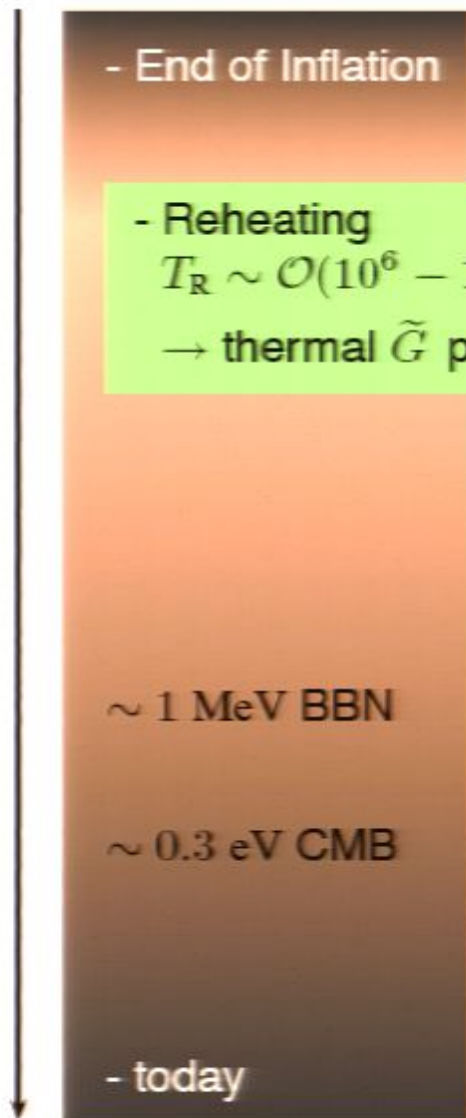


- Hard Part:



$$\begin{aligned}
 C_{\tilde{G}}(T) &= \int \frac{d^3\mathbf{p}}{(2\pi)^3} f_F(E) \left[\Gamma_{\tilde{G}}^{\text{soft}}(E) + \Gamma_{\tilde{G}}^{\text{hard}}(E) \right] \\
 &= \int \frac{d^3\mathbf{p}}{(2\pi)^3} f_F(E) \left[A_{\text{soft}} + A_{\text{hard}} + B \ln \left(\frac{1}{g} \right) \right]
 \end{aligned}$$

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- $C_{\tilde{G}}$: 2 \rightarrow 2 scatterings + finite T approach



$$C_{\tilde{G}} = \sum_{i=1}^3 \frac{3\zeta(3)T^6}{16\pi^3 M_{\text{Pl}}^2} \left(1 + \frac{M_i^2}{3m_{\tilde{G}}^2} \right) c_i g_i^2 \ln \left(\frac{k_i}{g_i} \right)$$

[Bolz, Brandenburg, Buchmüller, 2001]

[JP, Steffen, 2006]

(see also [Rychkov, Strumia, 2007])

Gravitino Dark Matter

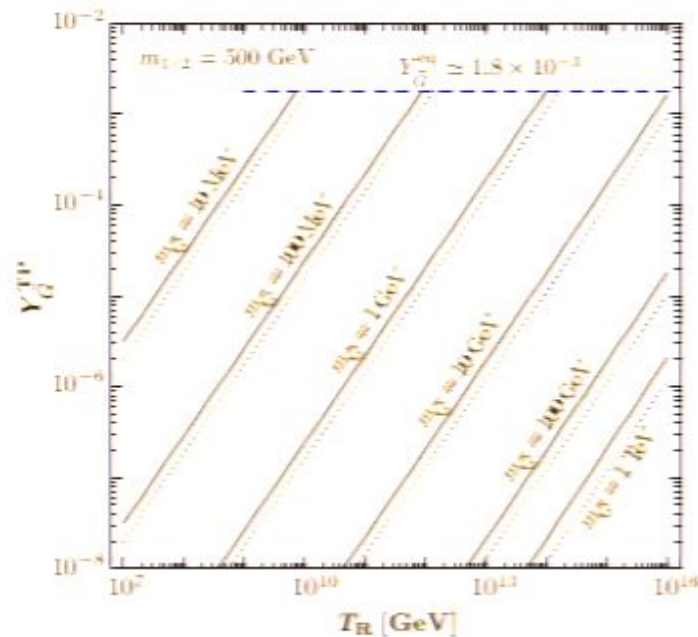
→ solve Boltzmann equation in terms of the Yield: $Y_{\tilde{G}}^{\text{TP}} \equiv \frac{n_{\tilde{G}}}{s}$

$$Y_{\tilde{G}}^{\text{TP}} = \sum_{i=1}^3 y_i g_i^2(T_{\text{R}}) \left(1 + \frac{M_i^2(T_{\text{R}})}{3m_{\tilde{G}}^2} \right) \ln \left(\frac{k_i}{g_i(T_{\text{R}})} \right) \left(\frac{T_{\text{R}}}{10^{10} \text{ GeV}} \right)$$

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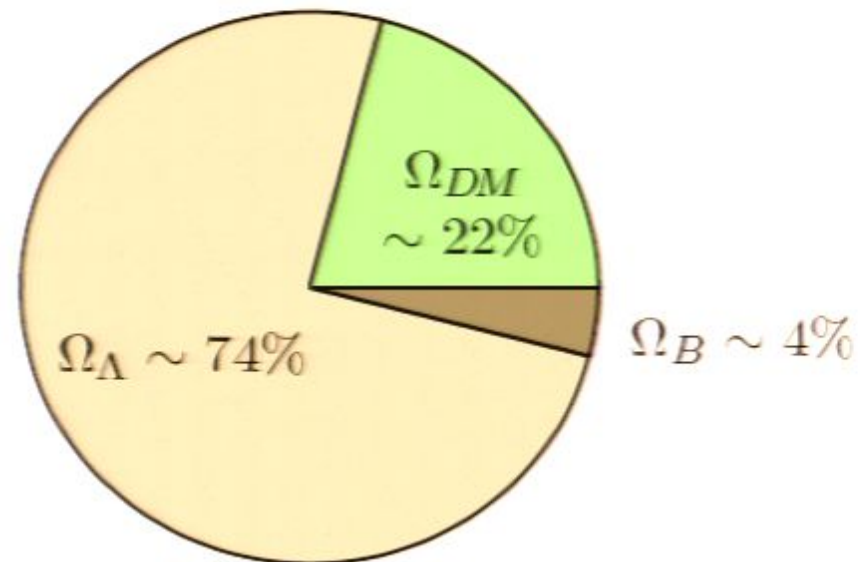
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Assume

- \tilde{G} lightest SUSY particle (LSP)
- R -Parity conservation

→ \tilde{G} stable and can be dark matter



Upper Limit on T_R from Thermal Production

relic \tilde{G} density:

$$\Omega_{\tilde{G}}^{\text{TP}} = m_{\tilde{G}} Y_{\tilde{G}}^{\text{TP}} s / \rho_c$$

observed DM abundance:

$$\Omega_{\text{dm}}^{3\sigma} h^2 = 0.105_{-0.030}^{+0.021}$$

→ upper limit on T_R :

$$\Omega_{\tilde{G}}^{\text{TP}} h^2 \leq 0.126$$

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T_R ... reheating temperature

$m_{\tilde{G}}$... gravitino mass

M_i ... gaugino masses

y_i ... $\mathcal{O}(10^{-12})$

$i = 3, 2, 1 \dots \text{SU}(3)_c, \text{SU}(2)_L, \text{U}(1)_Y$

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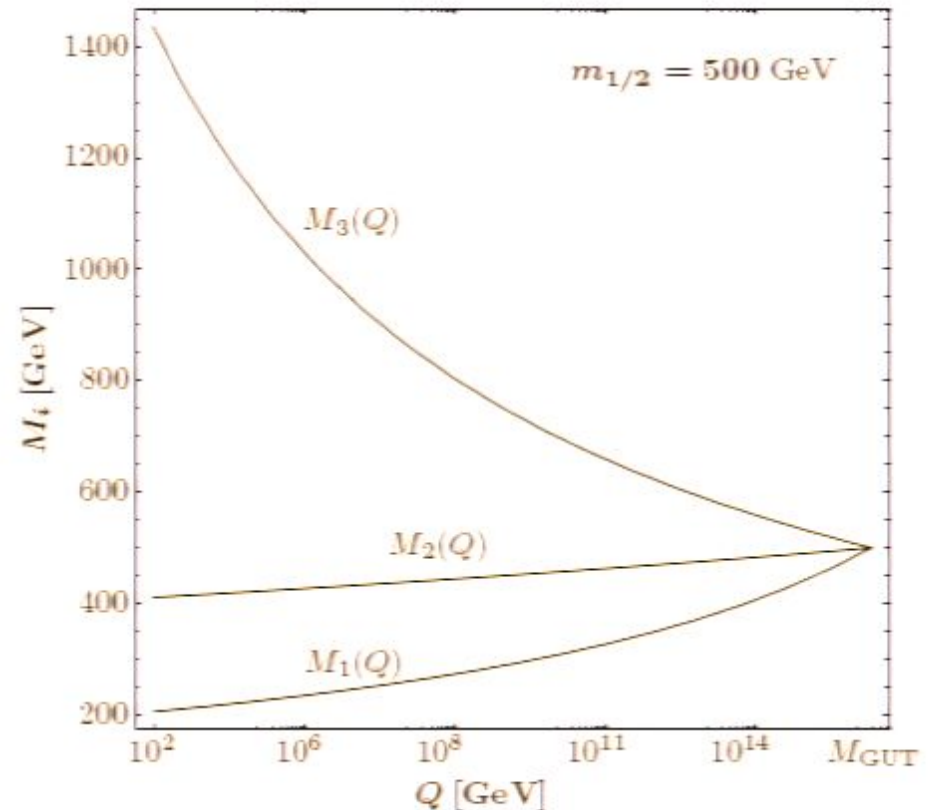
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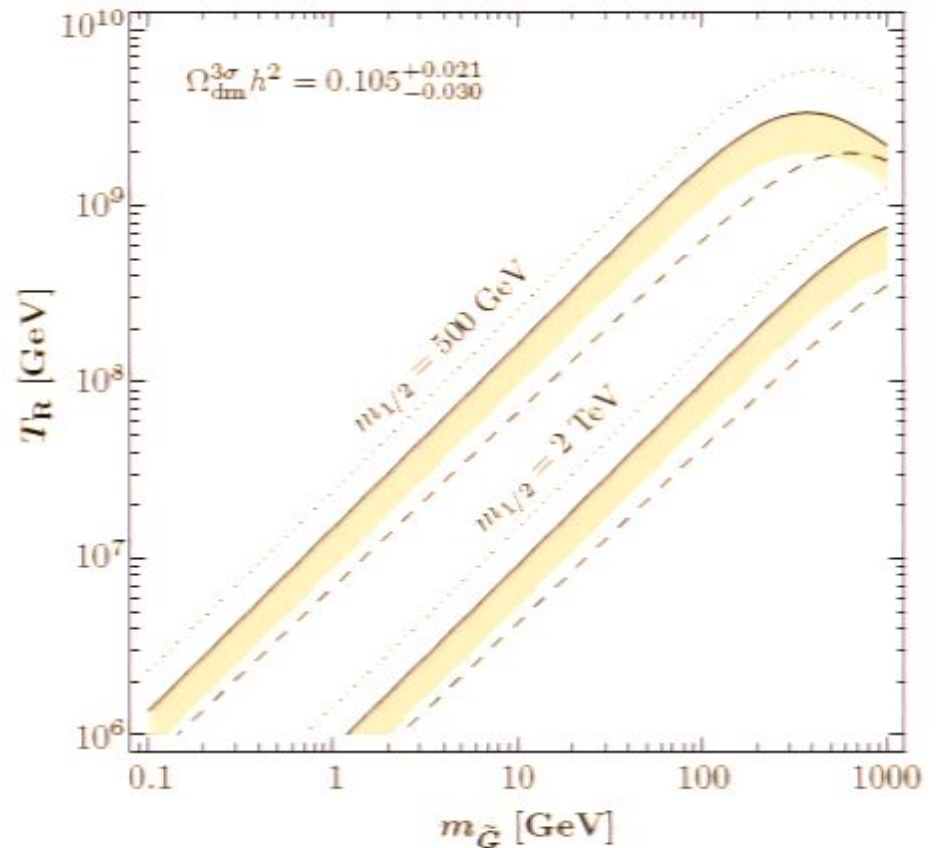
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Gravitino Production from Decays: NLSP $\rightarrow \tilde{G} + \text{SM}$

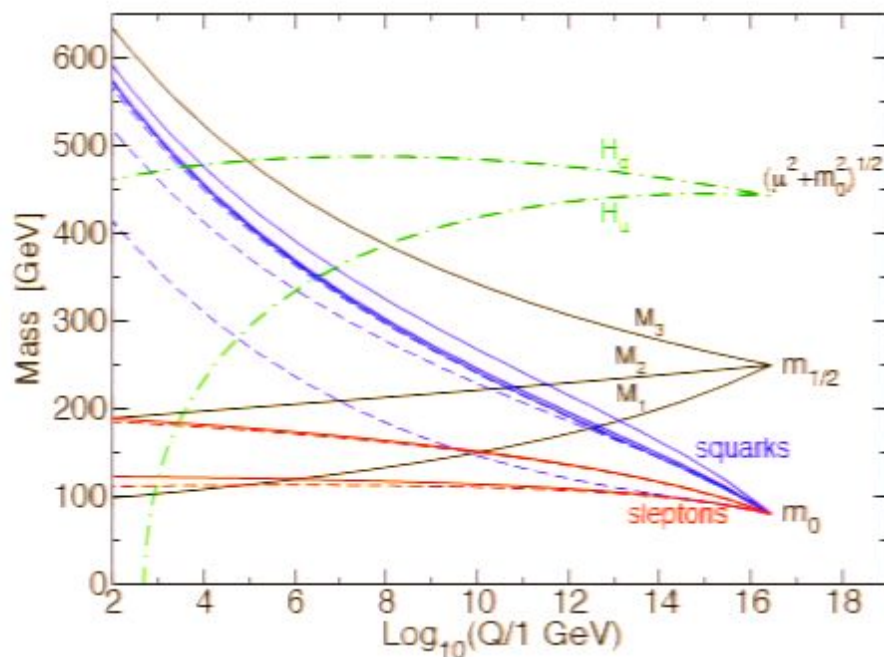
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systematic study of Ω_{NLSP}

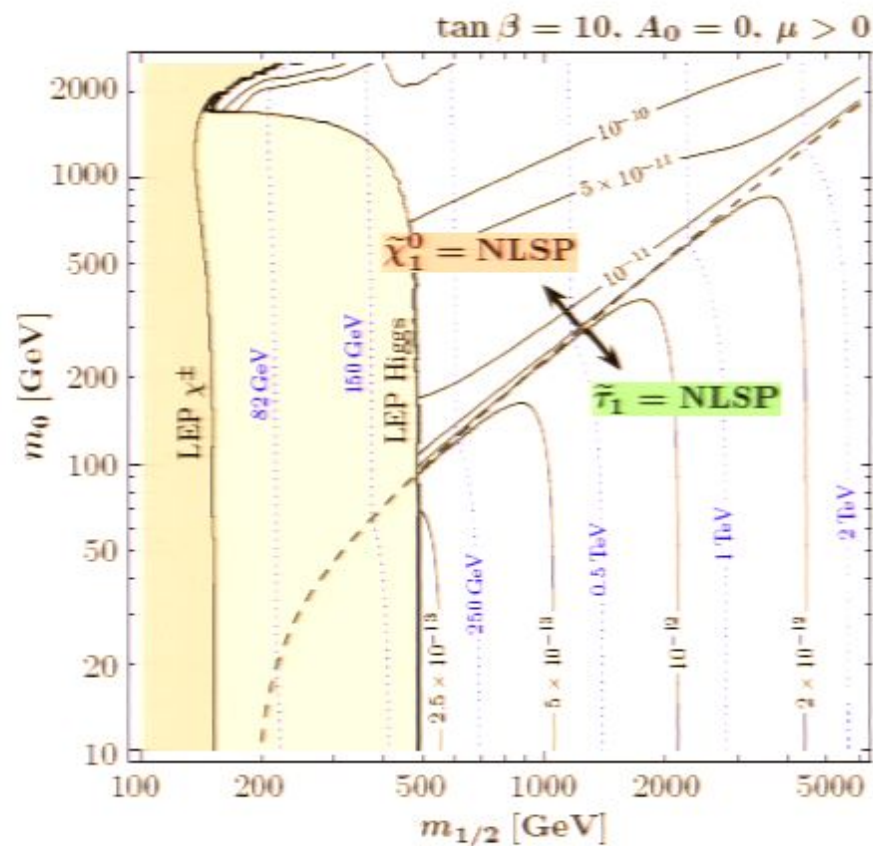
\rightarrow Constrained-MSSM

- $m_{1/2}$... universal gaugino mass
- m_0 ... universal scalar mass
- A_0 ... universal trilinear scalar interaction
- $\tan \beta$... ratio of the Higgs vev's
- $\text{sgn } \mu$... sign of the higgsino parameter

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Constraints on T_R in the CMSSM

$$0.075 \leq (\Omega_{\tilde{G}}^{\text{TP}} + \Omega_{\tilde{G}}^{\text{NTP}})h^2 \leq 0.126$$

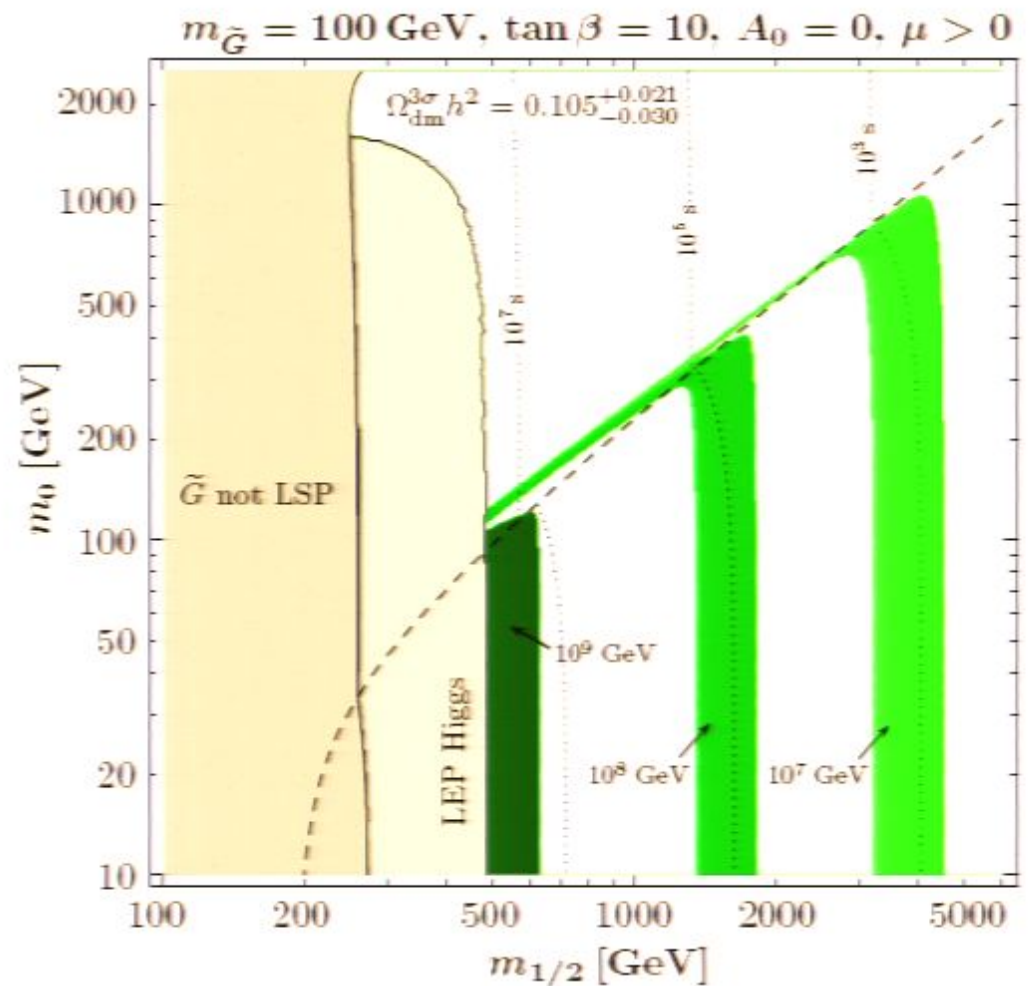
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$$m_{\tilde{G}} = 100 \text{ GeV}$$

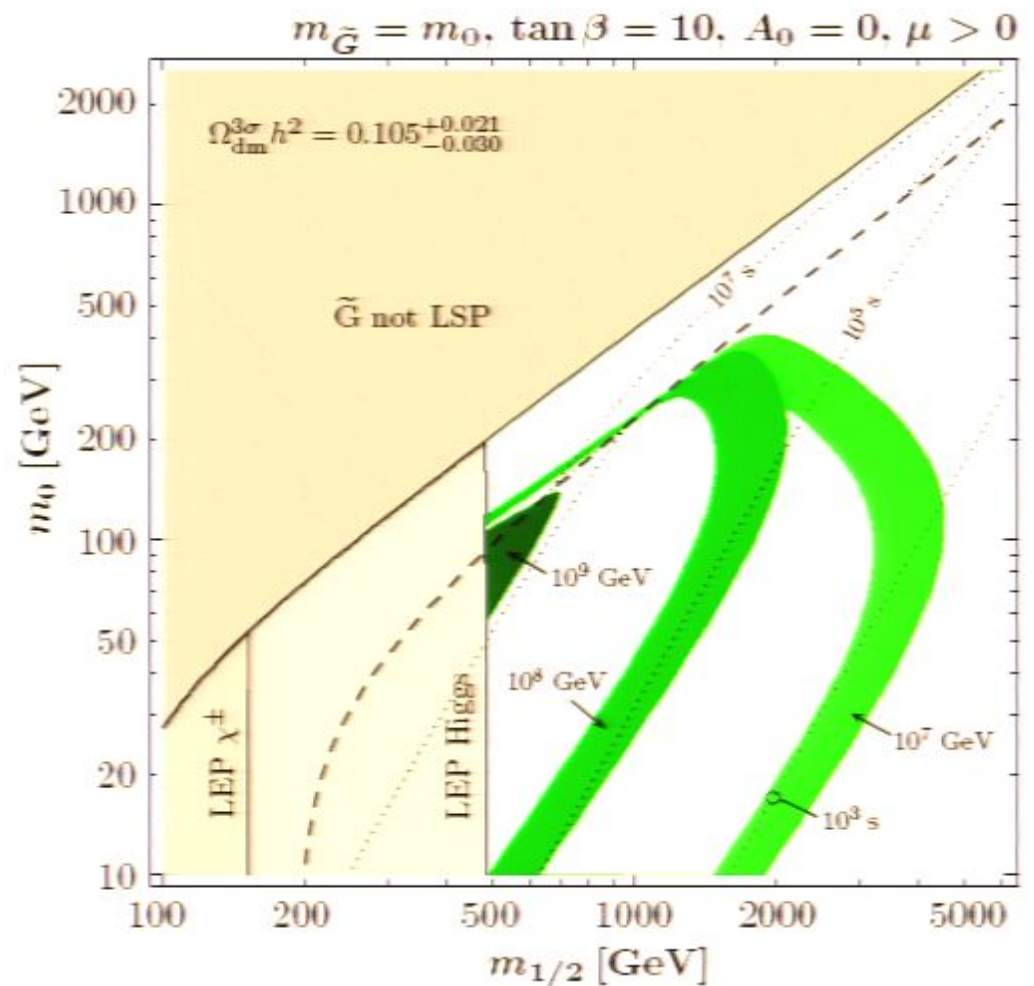


Constraints on T_R in the CMSSM

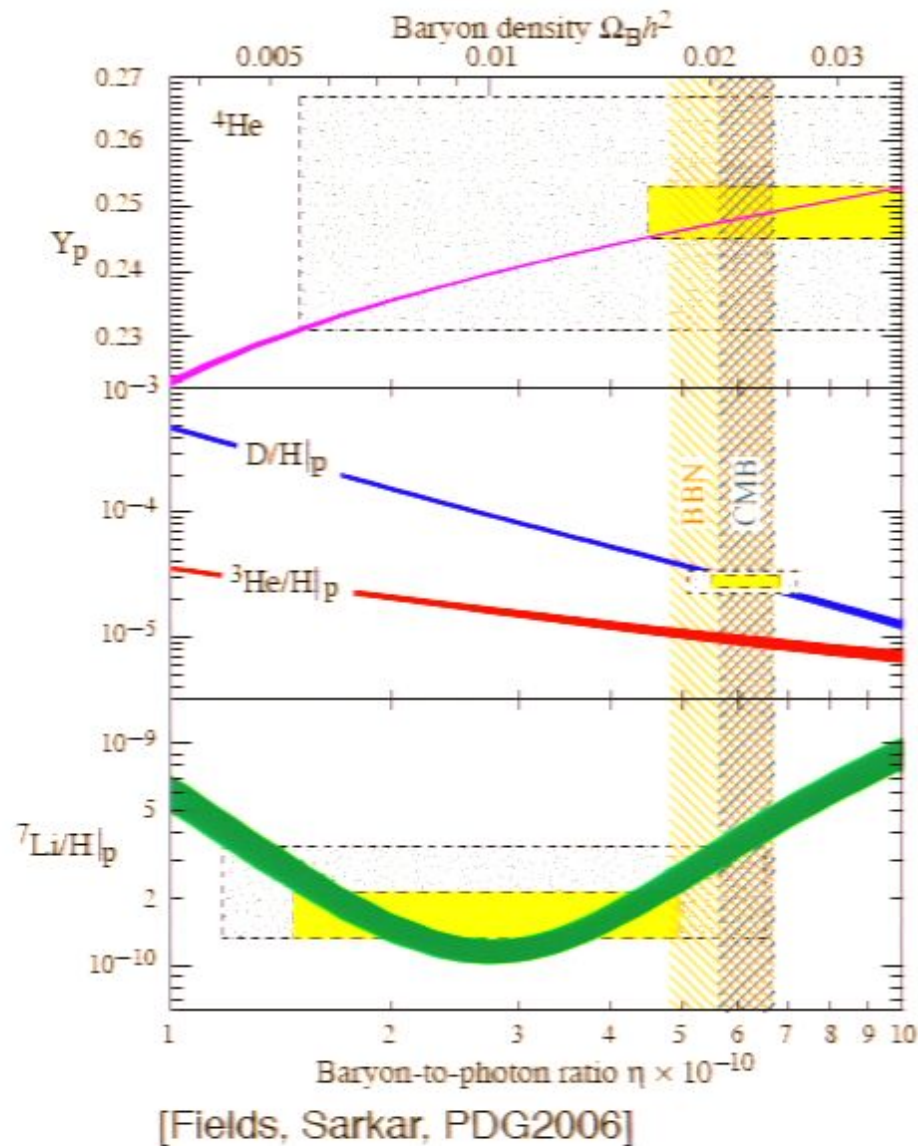
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$$m_{\tilde{G}} = m_0$$



Big Bang Nucleosynthesis (BBN)



Beyond the Standard Model:

Timing “extra neutrino species”

Non-thermal nuclear reactions

NLSP decay during BBN

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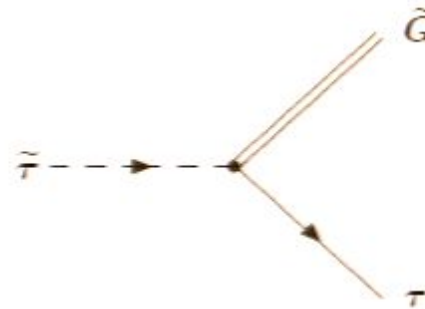
→ constraints on $m_{\tilde{G}}$, m_{NLSP} and Y_{NLSP}

- $\tilde{\chi}_1^0$ NLSP disfavoured in CMSSM
 $\tilde{\tau}_1$ region: BBN bounds important but much less severe
e.g. [Ellis et al., 2004; Feng et al., 2004; Jedamzik et al., 2006; Steffen, 2006; Cyburt et al., 2006]
- → in the following $\tilde{\tau}_1 = \text{NLSP}$

NLSP decay during BBN

- $\tilde{\tau}_1$ can decay during/after BBN, leading to ...

- electromagnetic energy release

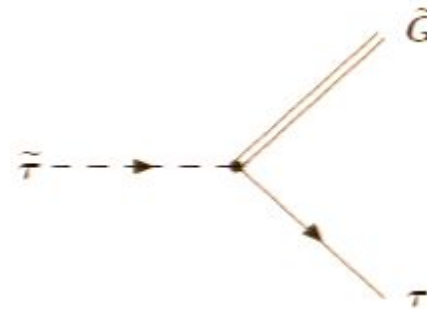


$$\tau_{\tilde{\tau}_1} \simeq \Gamma^{-1}(\tilde{\tau}_1 \rightarrow \tilde{G}\tau) = \frac{48\pi m_{\tilde{G}}^2 M_{\text{P}}^2}{m_{\tilde{\tau}_1}^5} \left(1 - \frac{m_{\tilde{G}}^2}{m_{\tilde{\tau}_1}^2}\right)^{-4}$$

NLSP decay during BBN

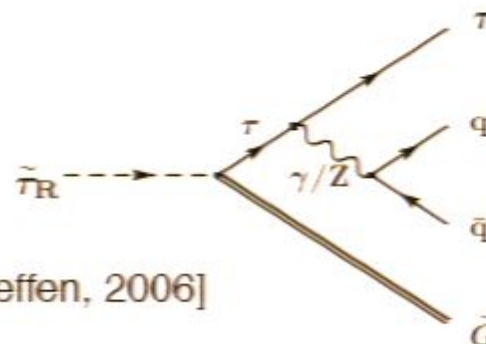
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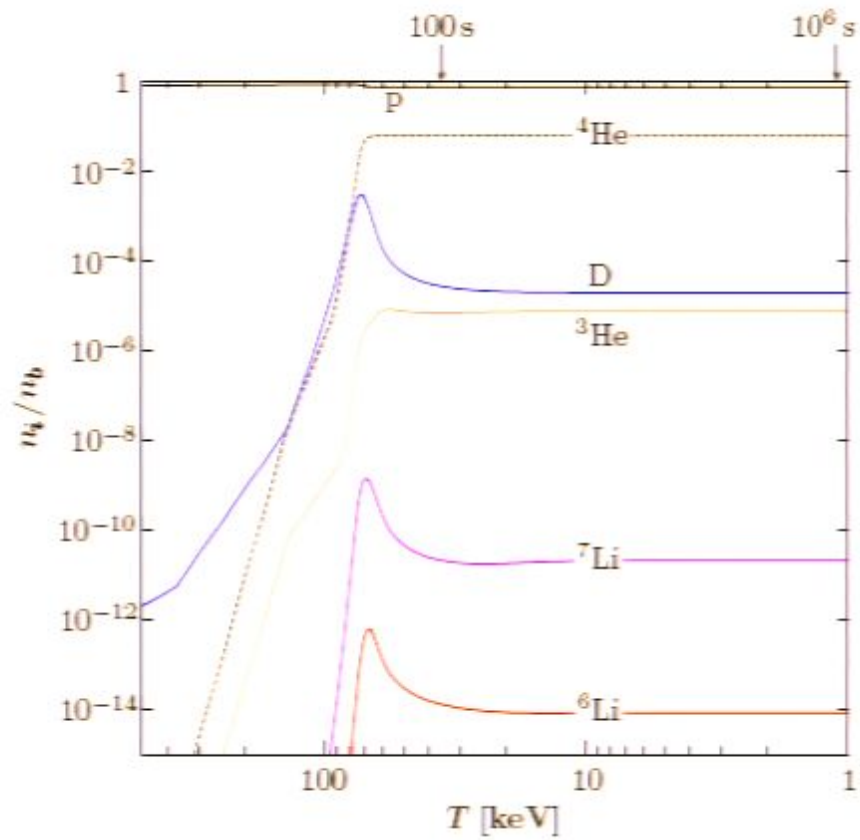
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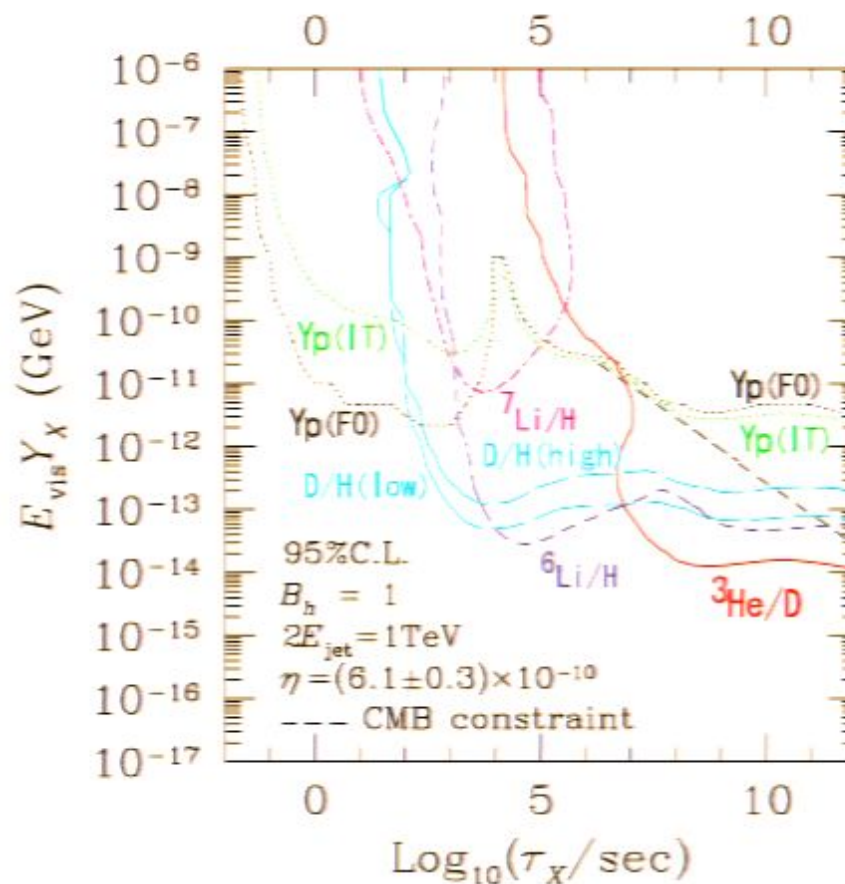
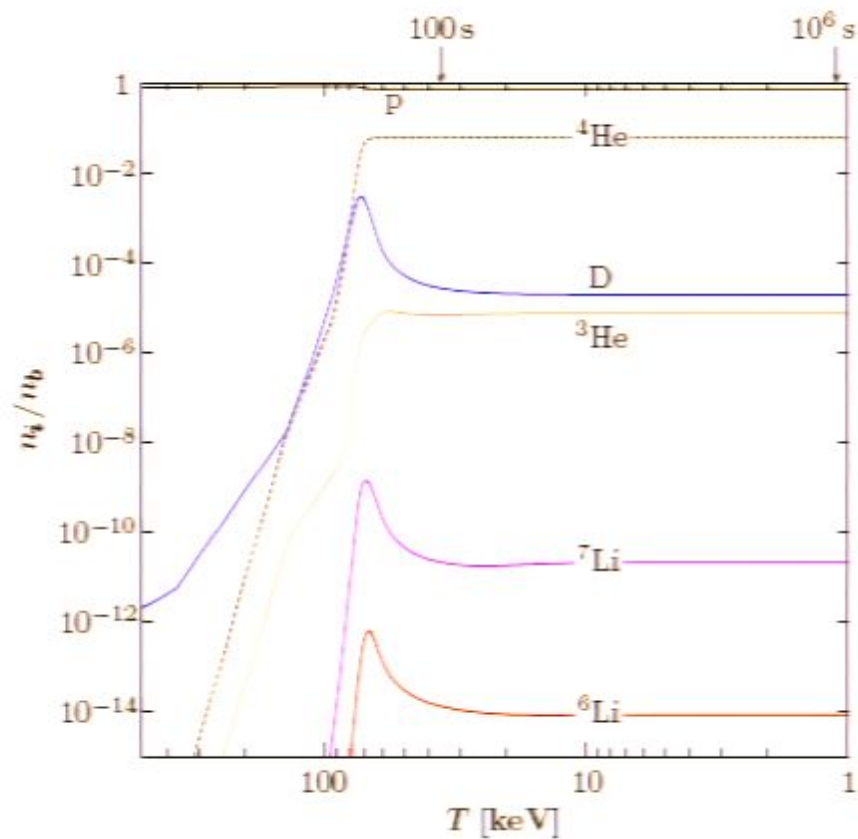


$$B_{\text{h}} \lesssim 3 \times 10^{-3} (m_{\tilde{\tau}_1} < 2 \text{ TeV}) \text{ [Steffen, 2006]}$$

BBN Constraints from em./had. energy injection



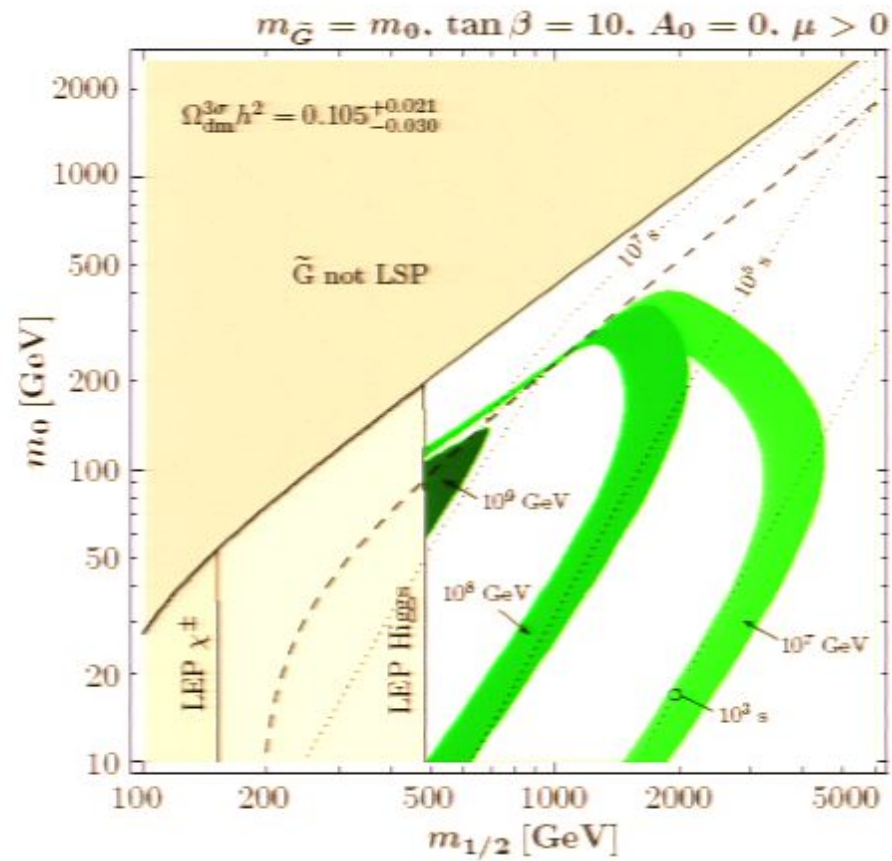
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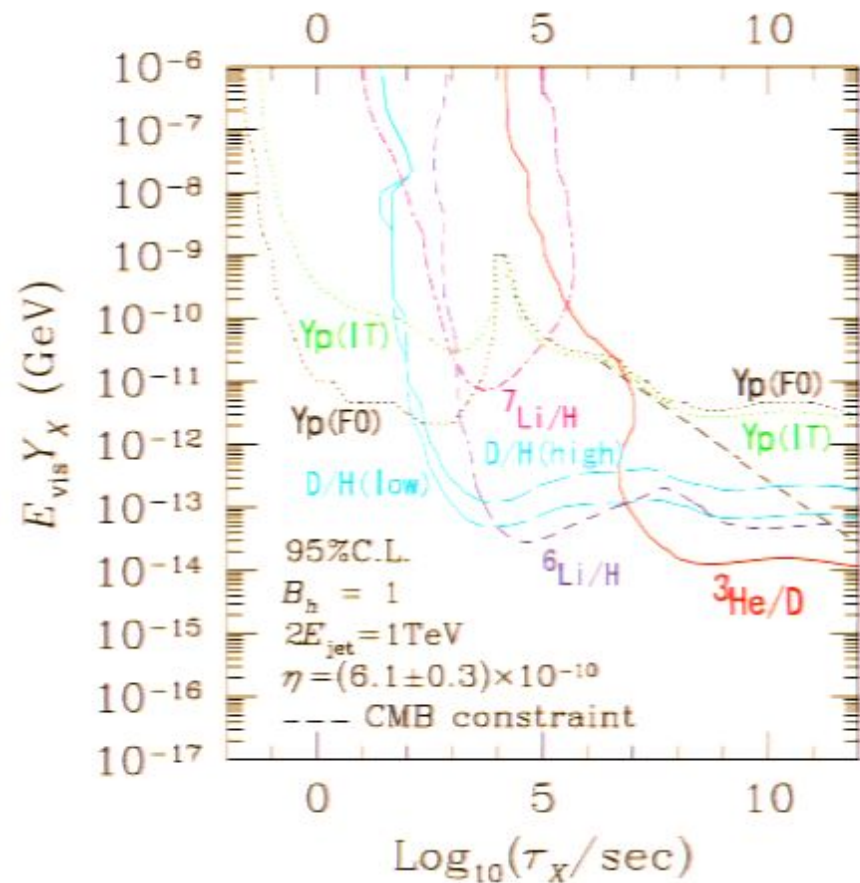
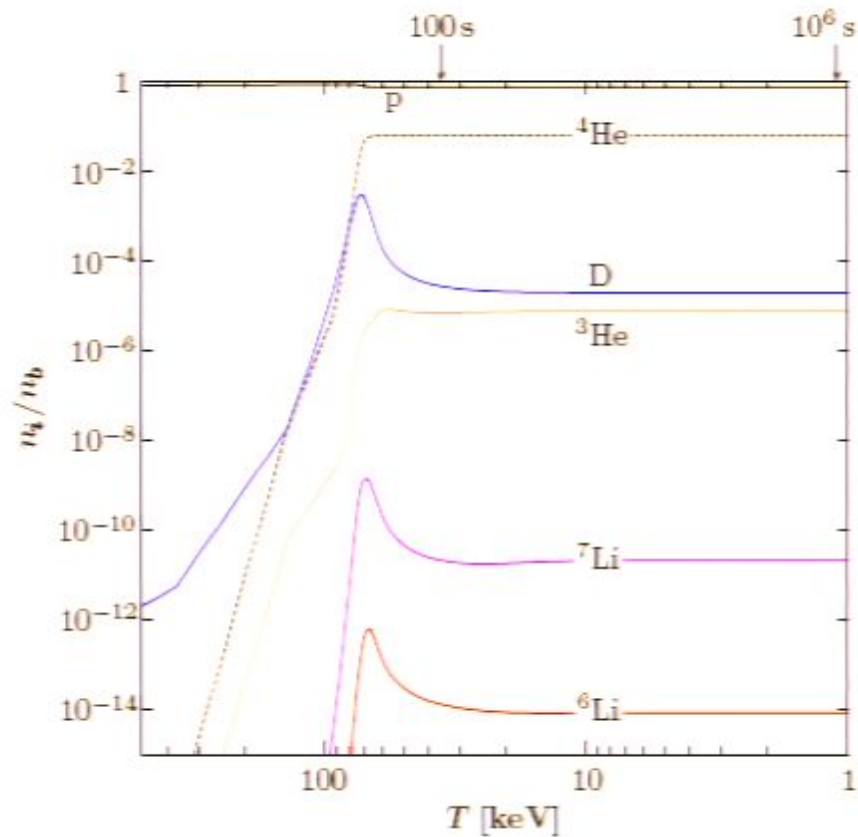
$m_X = 1\text{TeV}$

[Kawasaki, Kohri, Moroi, 2005]

Constraints on T_R from catalyzed BBN



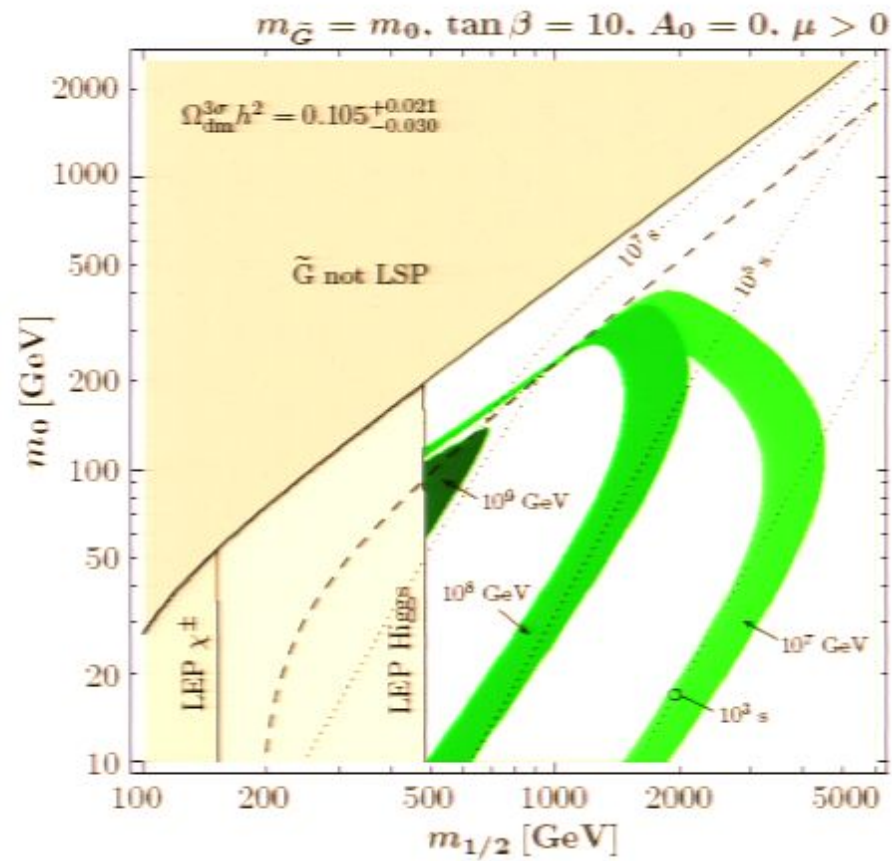
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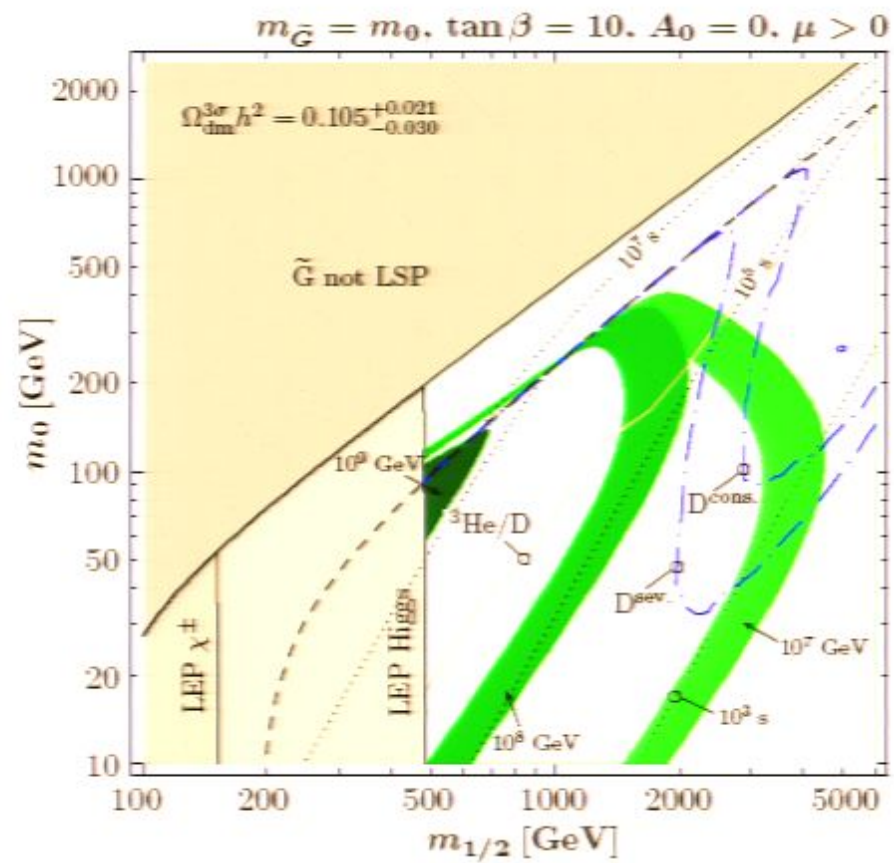
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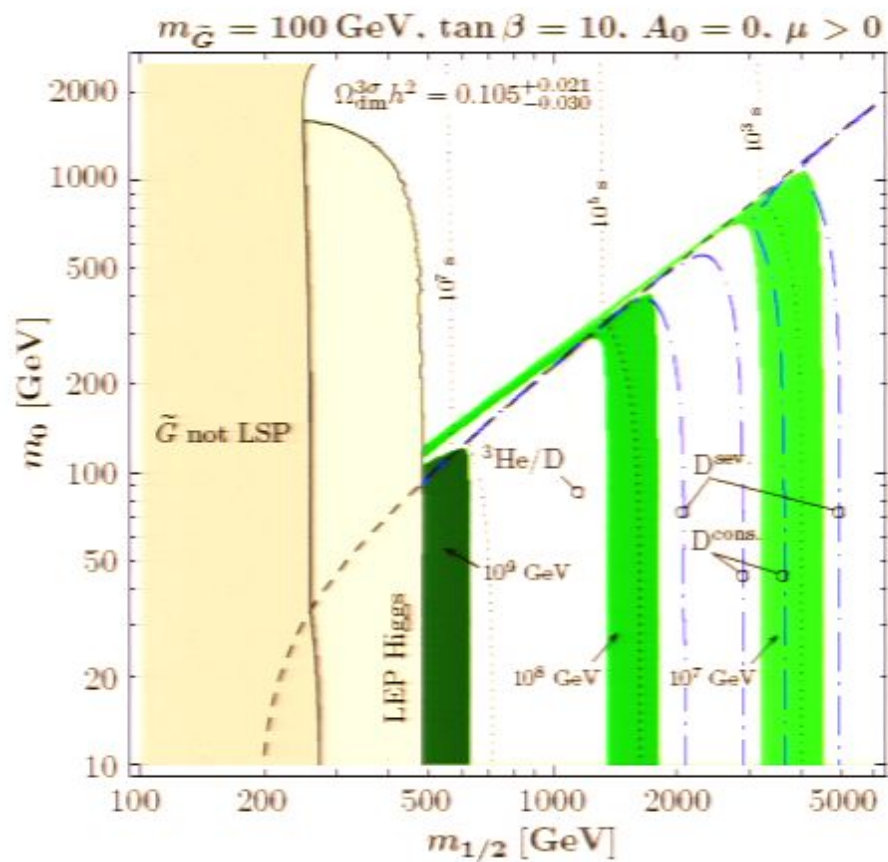
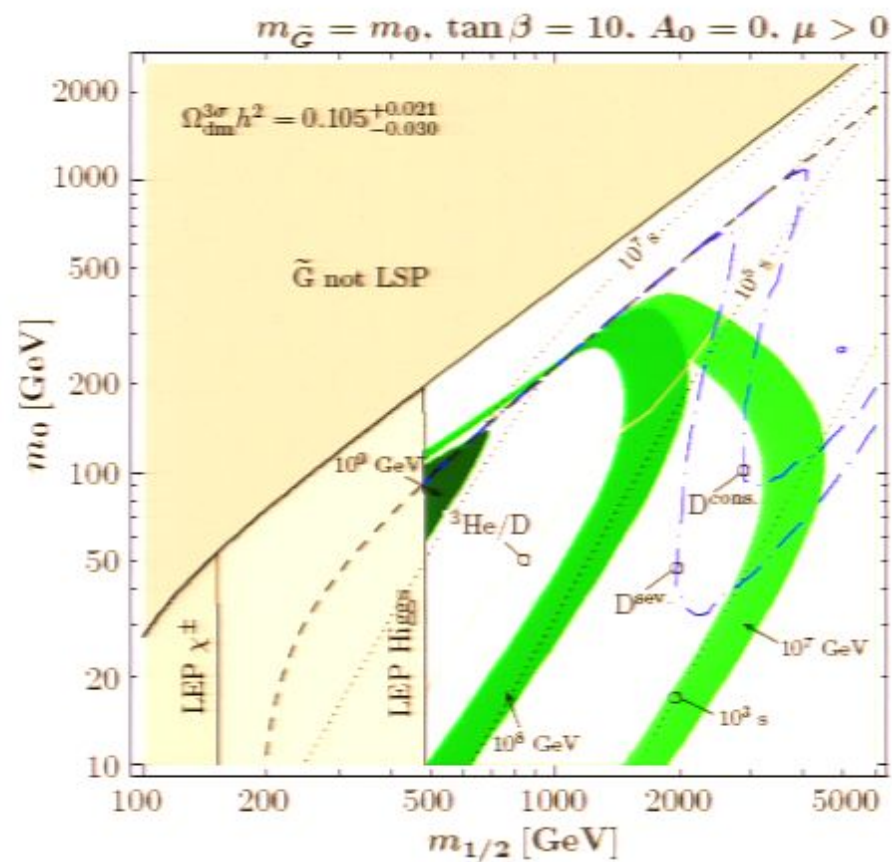
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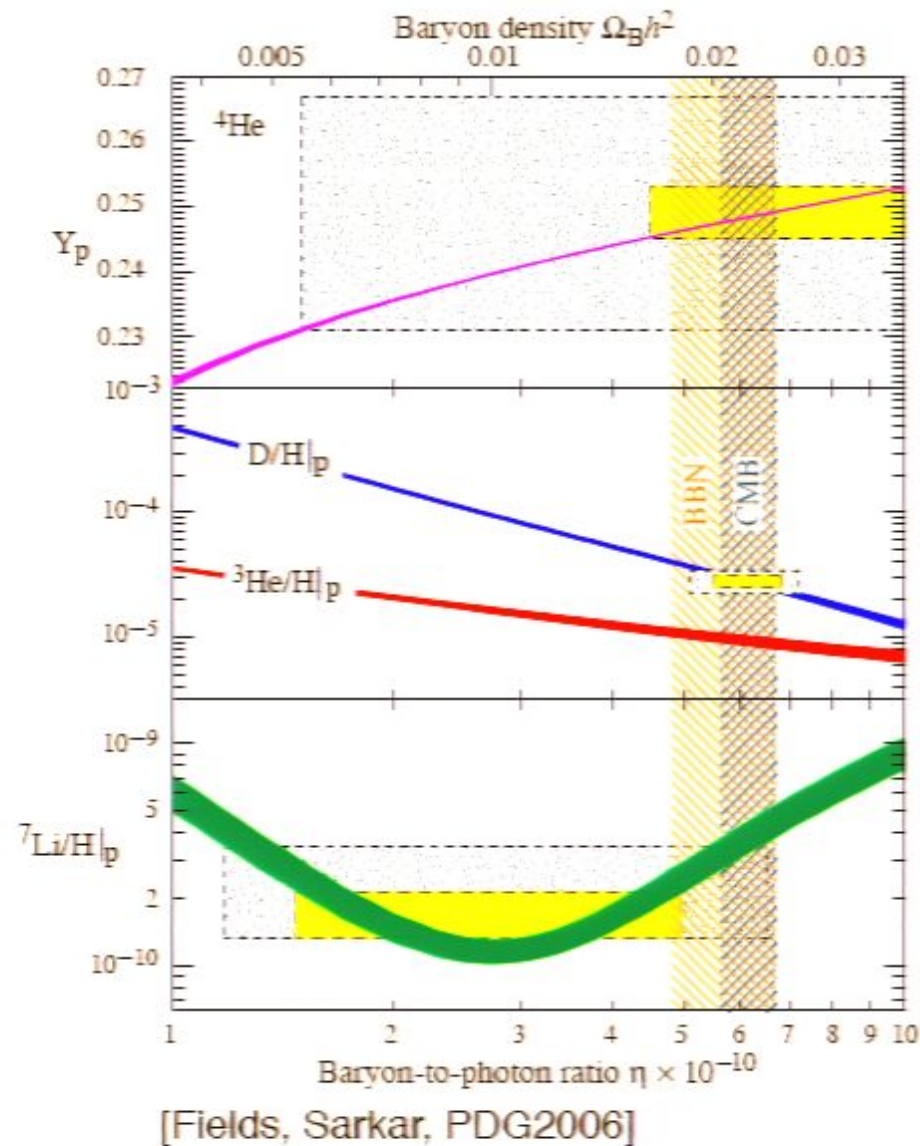
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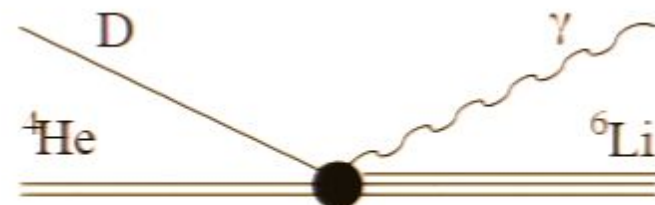
Catalyzed BBN

- bound-state formation of X^- with light elements opens up new reaction channels

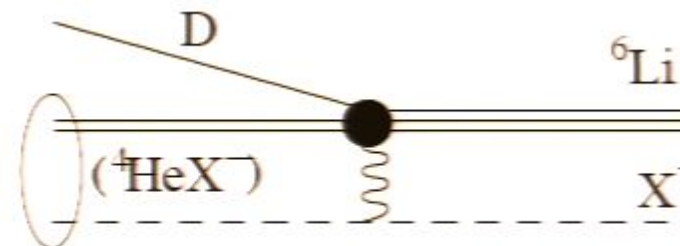
Catalyzed BBN

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- most significant for production of ${}^6\text{Li}$ [Pospelov, 2006]

standard BBN:
 $\rightarrow \langle \sigma_S v \rangle$



catalyzed BBN:
 $\rightarrow \langle \sigma_C v \rangle$



cross-section enhanced by 7 orders of magnitude

Numerical Study of CBBN

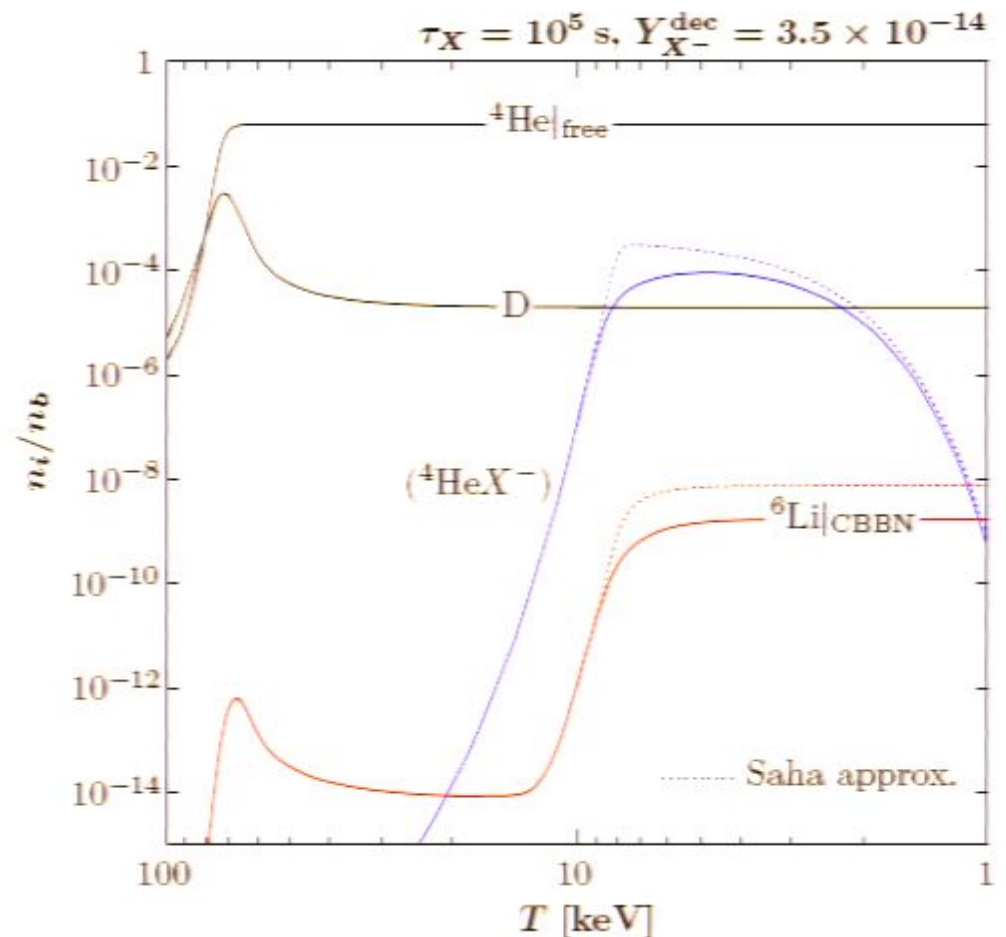
- Bound state (BS) formation becomes efficient at $T \sim 10$ keV after Standard BBN

$$\frac{dY_{6\text{Li}}}{dt} = \langle \sigma_C v \rangle s Y_{\text{BS}} Y_{\text{D}}$$

$$\frac{dY_{\text{BS}}}{dt} = \langle \sigma_r v \rangle s Y_{\delta} - \Gamma_{X^-} Y_{\text{BS}} - \langle \sigma_C v \rangle s Y_{\text{BS}} Y_{\text{D}}$$

etc ...

- $\langle \sigma_r v \rangle Y_{\delta}$ parameterizes competition between photo-dissociation and recombination of $({}^4\text{He}X^-)$
- $\langle \sigma_C v \rangle \dots$ [Kamimura et al., 2007]



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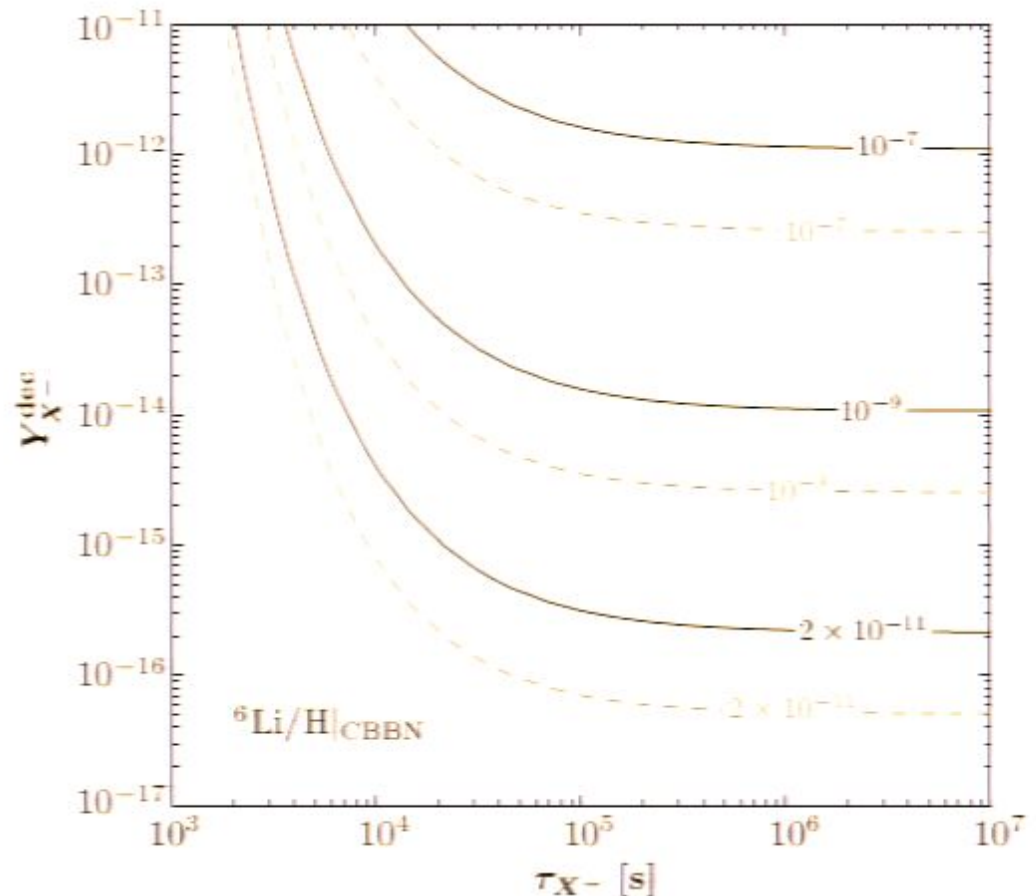
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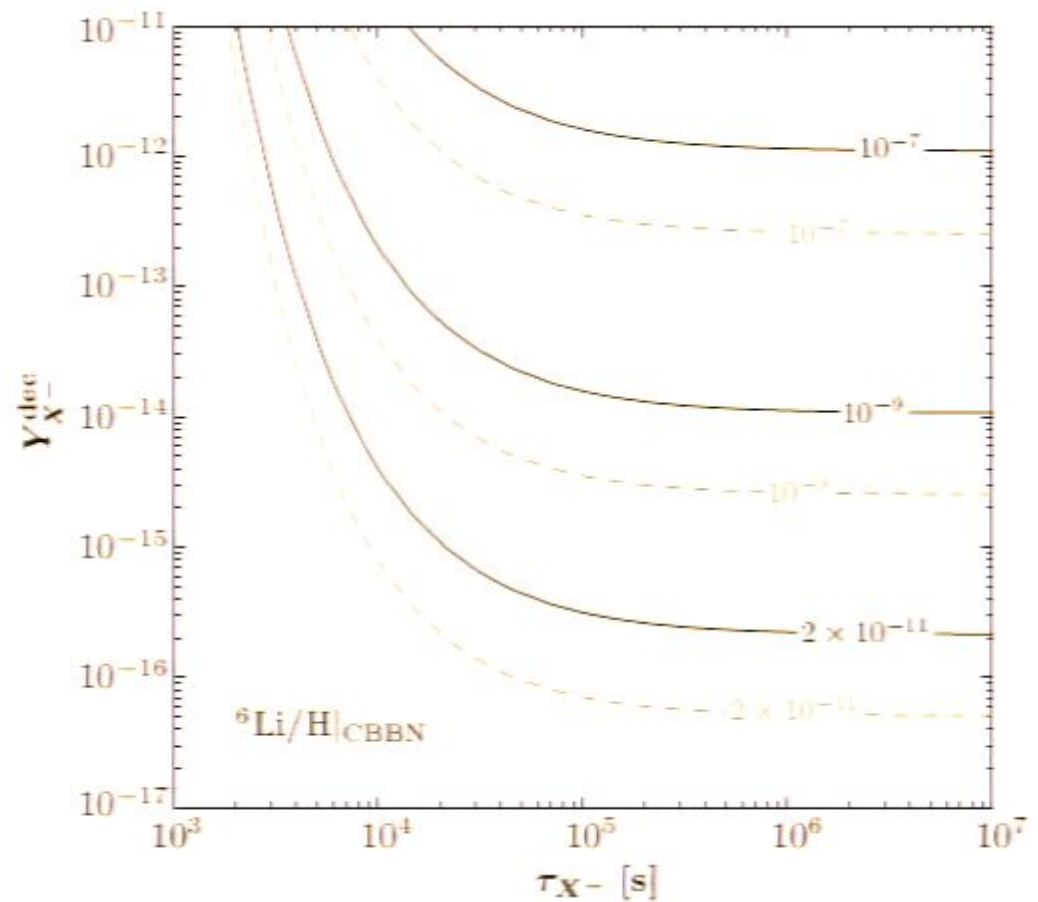
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- observationally inferred primordial abundance:

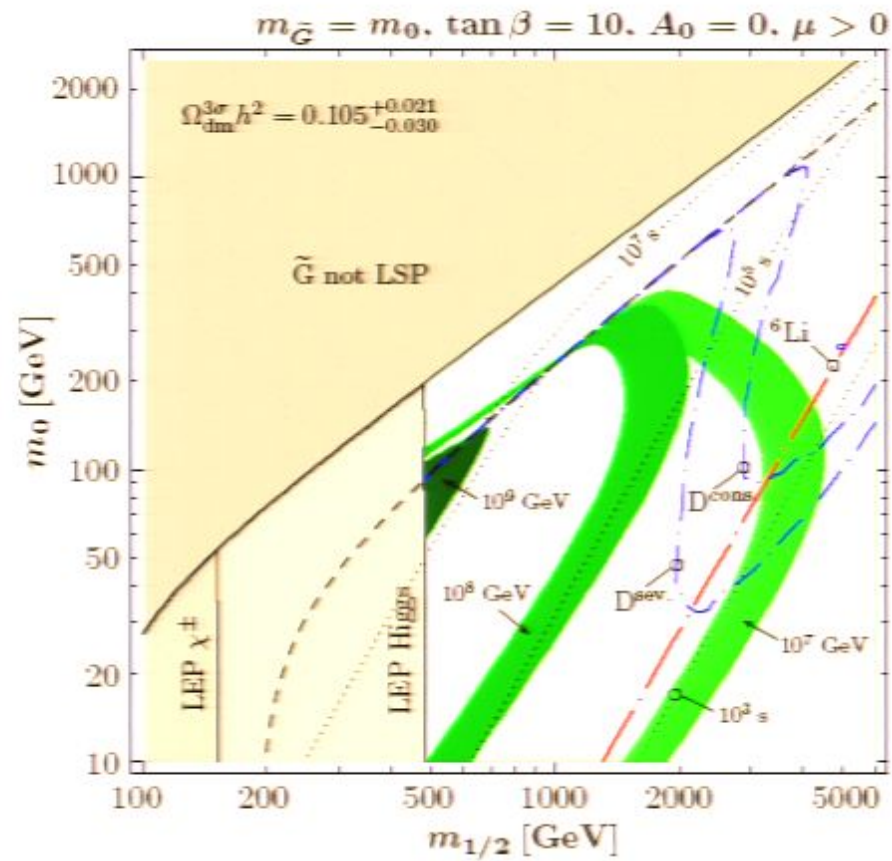
[Cyburt et al., 2002]

$$({}^6\text{Li}/\text{H})_{\text{p}} \lesssim 2 \times 10^{-11}$$

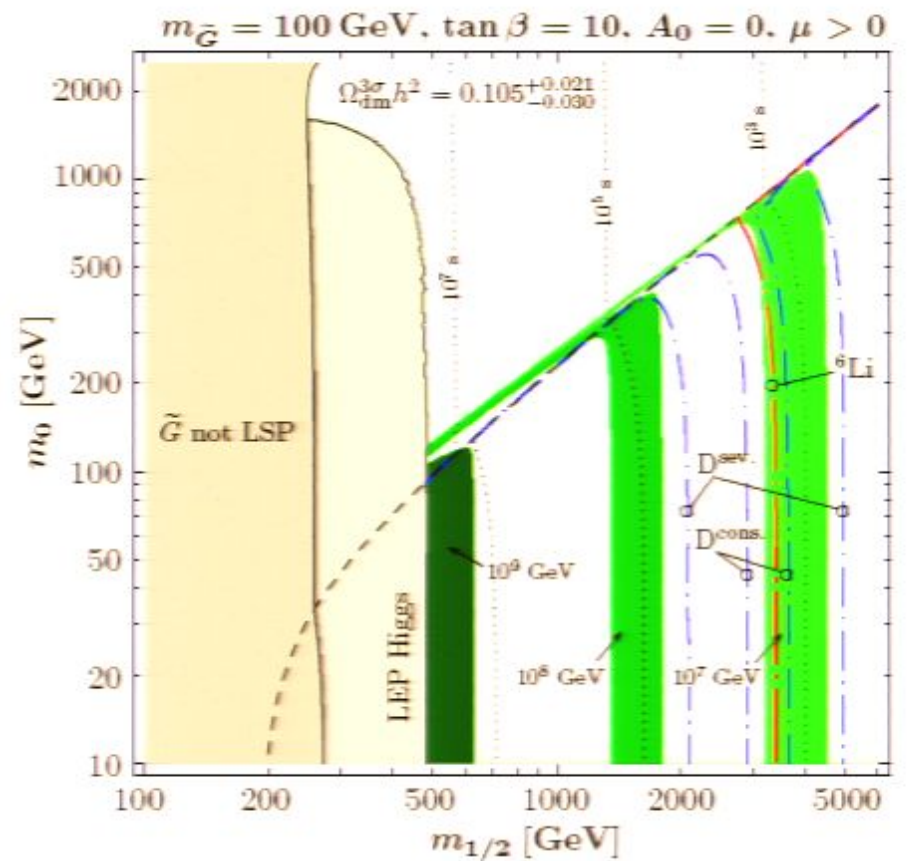
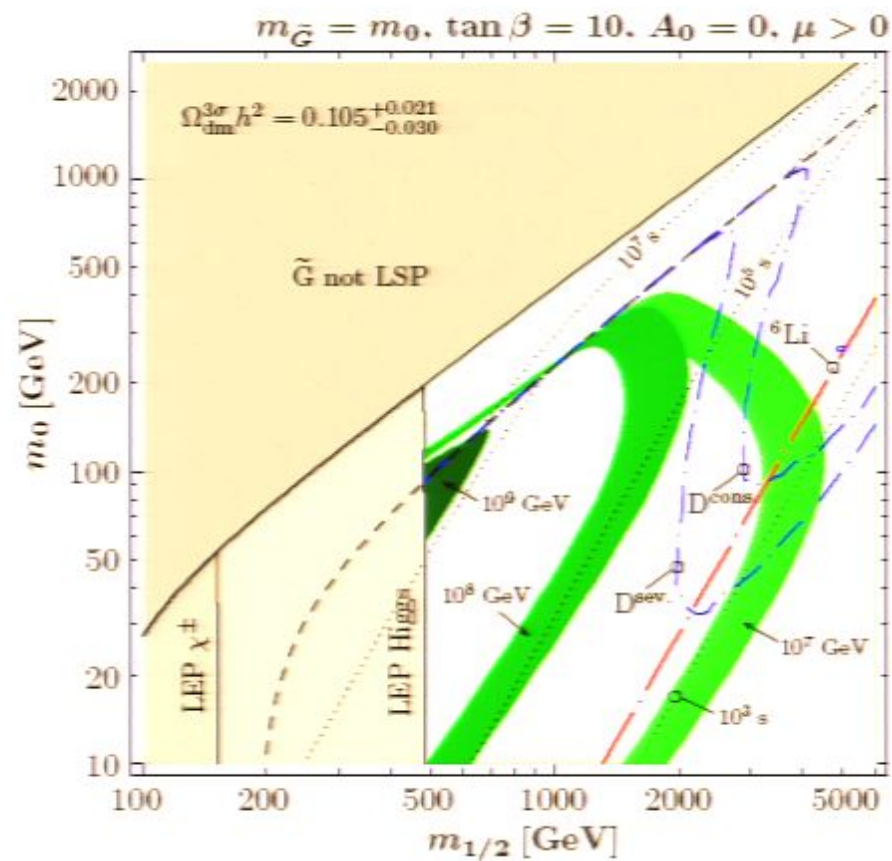
→ upper bound on
 $Y_{X^-} = Y_{\tilde{\tau}_1}/2$



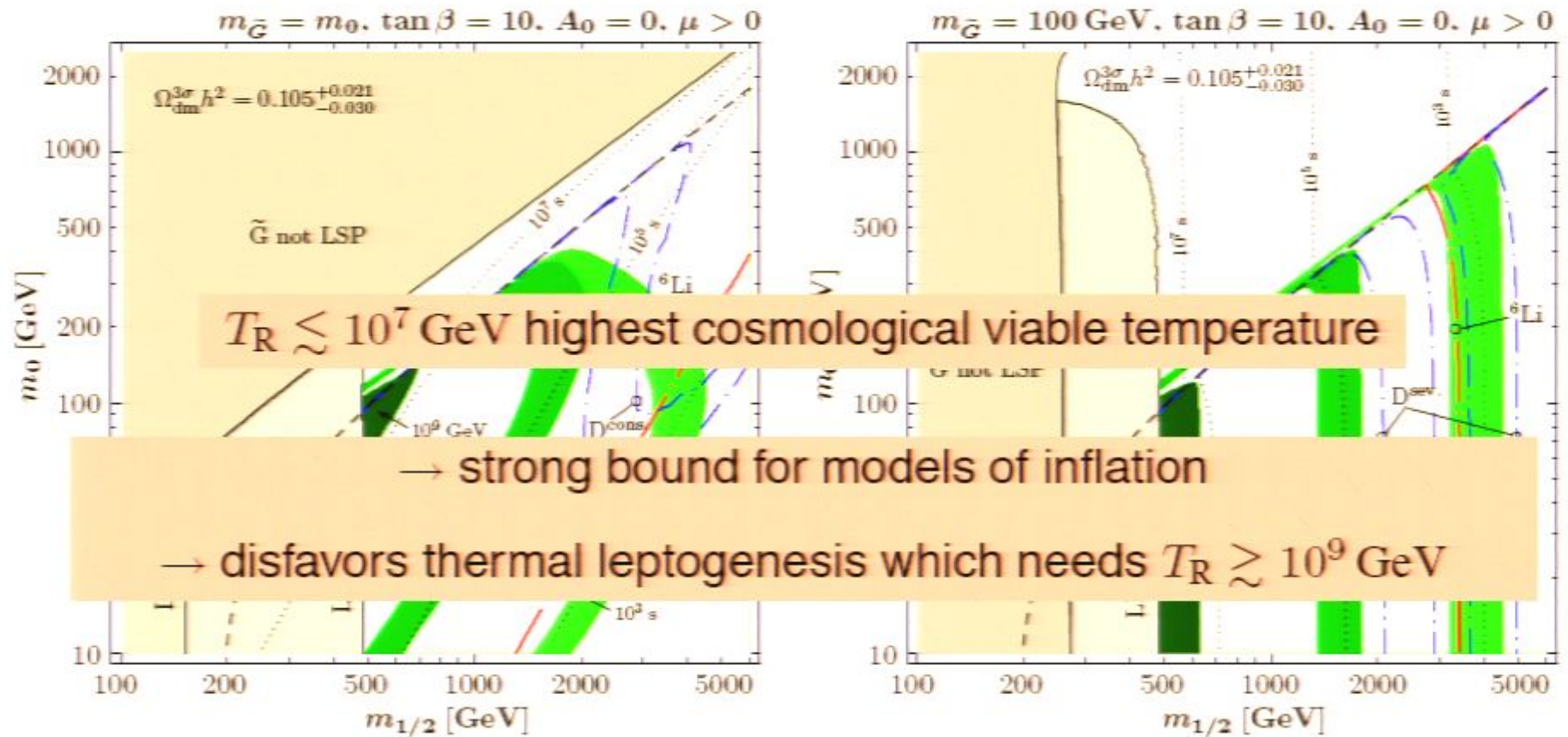
Constraints on T_R from catalyzed BBN



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- For typical $Y_{\text{NLSP}} \gtrsim 7 \times 10^{-14}$ the stau lifetime is constrained
 $\tau_{\tilde{\tau}_1} \lesssim 5 \times 10^3 \text{ s}$
- Implies a lower limit on the mass splitting between $m_{\tilde{\tau}_1}$ and $m_{\tilde{G}}$
- Scan over the CMSSM parameter space ($m_{\tilde{\tau}_1} \leftrightarrow m_{1/2}$):

$$\begin{aligned}m_{1/2} &= 0.1 - 6 \text{ TeV}, \\ \tan \beta &= 2 - 60, \\ \text{sgn } \mu &= \pm 1, \\ -4m_0 &< A_0 < 4m_0\end{aligned}$$

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- SUSY mass spectrum scales with $m_{1/2}$ in the $\tilde{\tau}_1$ NLSP region!

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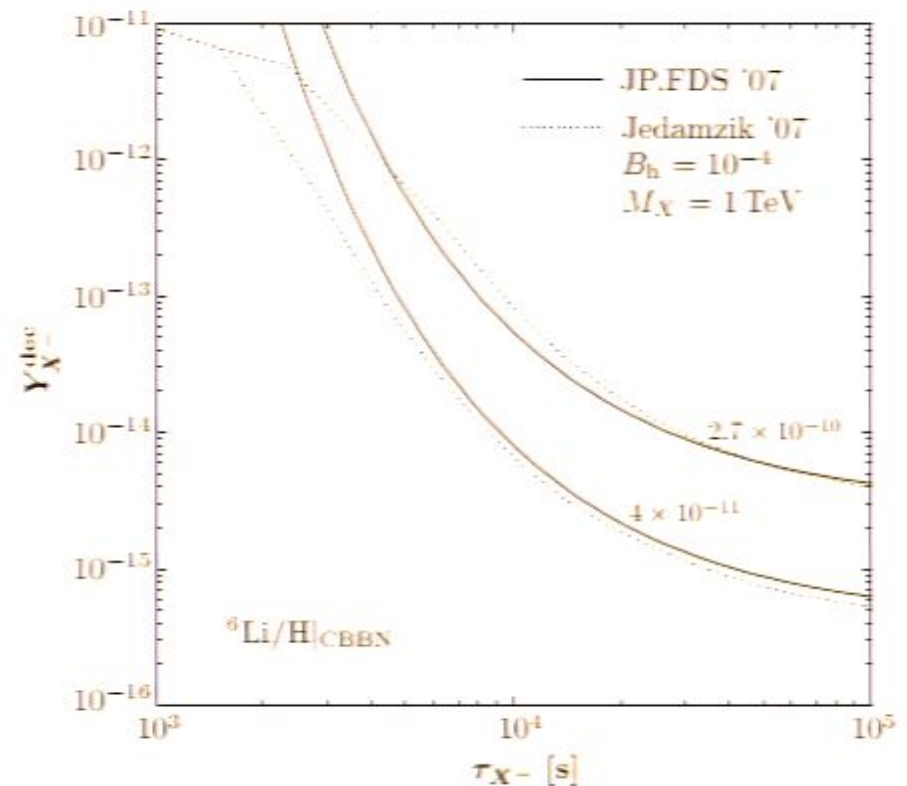
Adopting for

$$(^6\text{Li}/\text{H})_p \lesssim 4 \times 10^{-11} \quad (2.7 \times 10^{-10})$$

[Jedamzik, 2007]

$$m_{1/2} : \quad 0.87 \quad (0.78)$$

$$T_R : \quad 5.3 \times 10^7 \quad (6.5 \times 10^7)$$

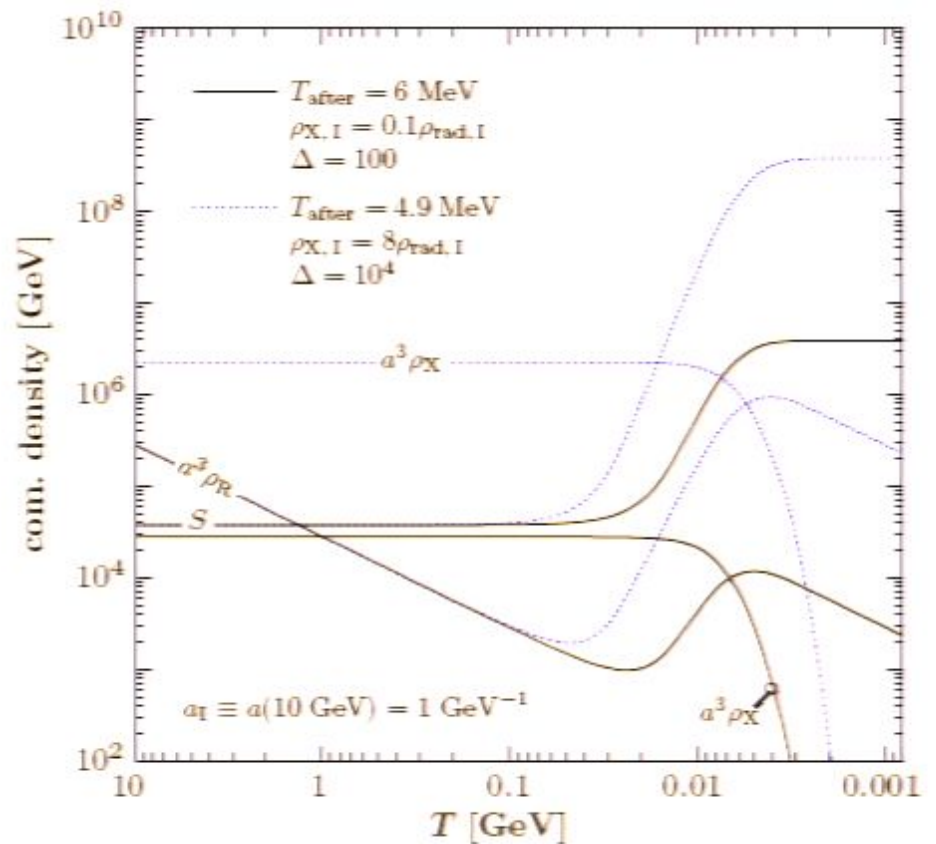


Late-time Entropy Production

out-of-eq. decay of heavy particle X:

$$\frac{dS}{dt} = \frac{dQ}{T} = \frac{\Gamma_X \rho_X a^3}{T}$$

see also [Buchmüller et al., 2006]



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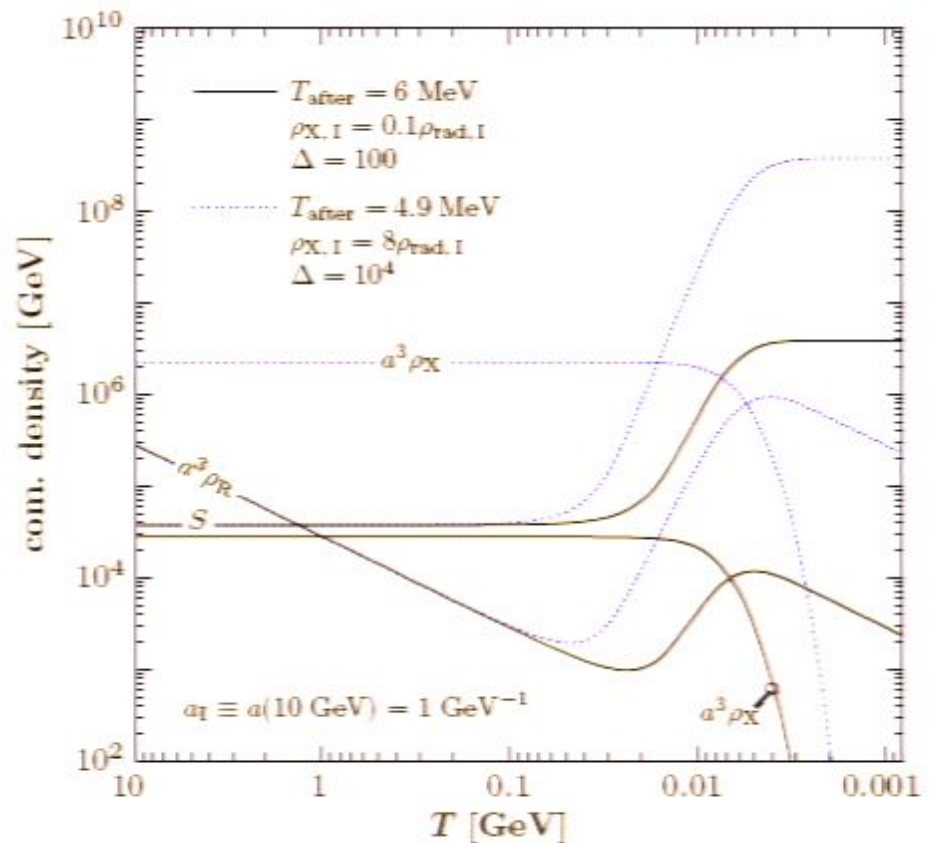
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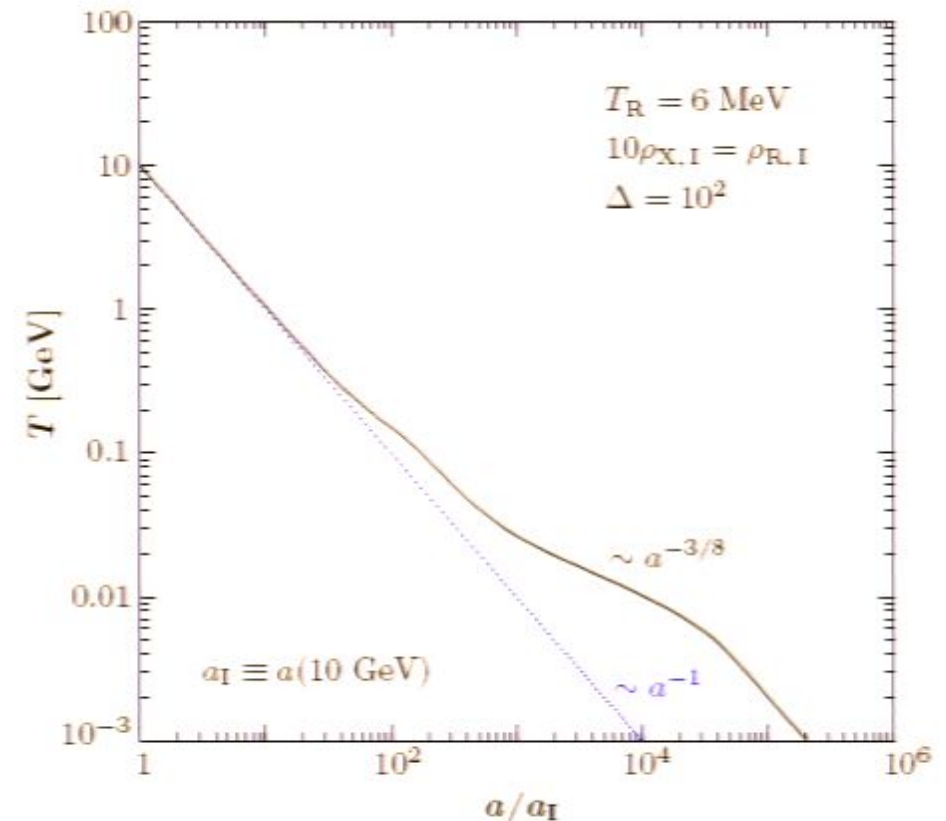
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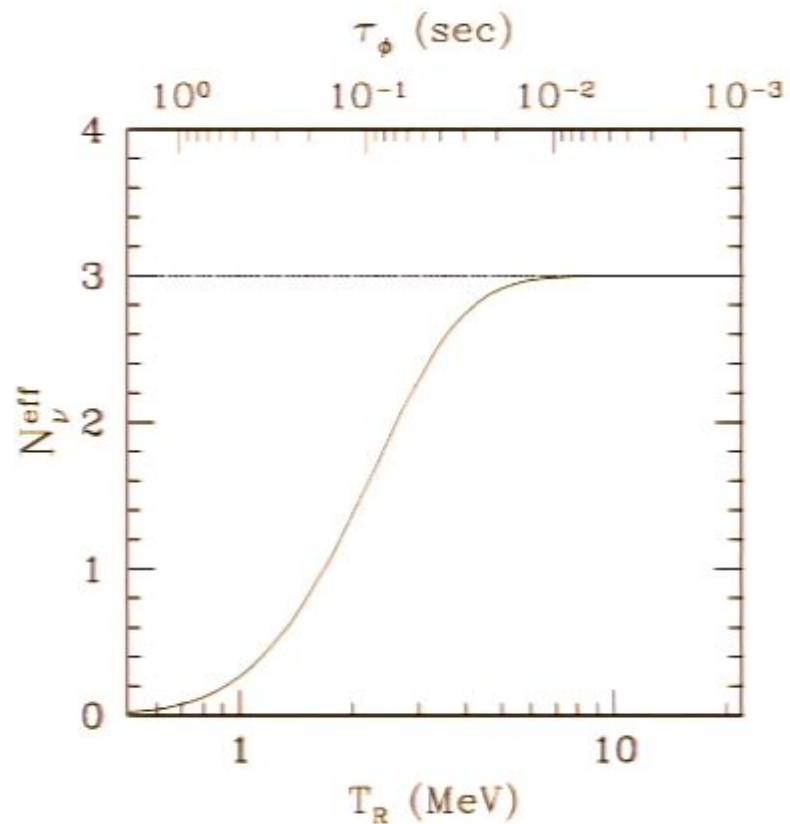
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[Kawasaki, Kohri, Sugiyama, 1999]

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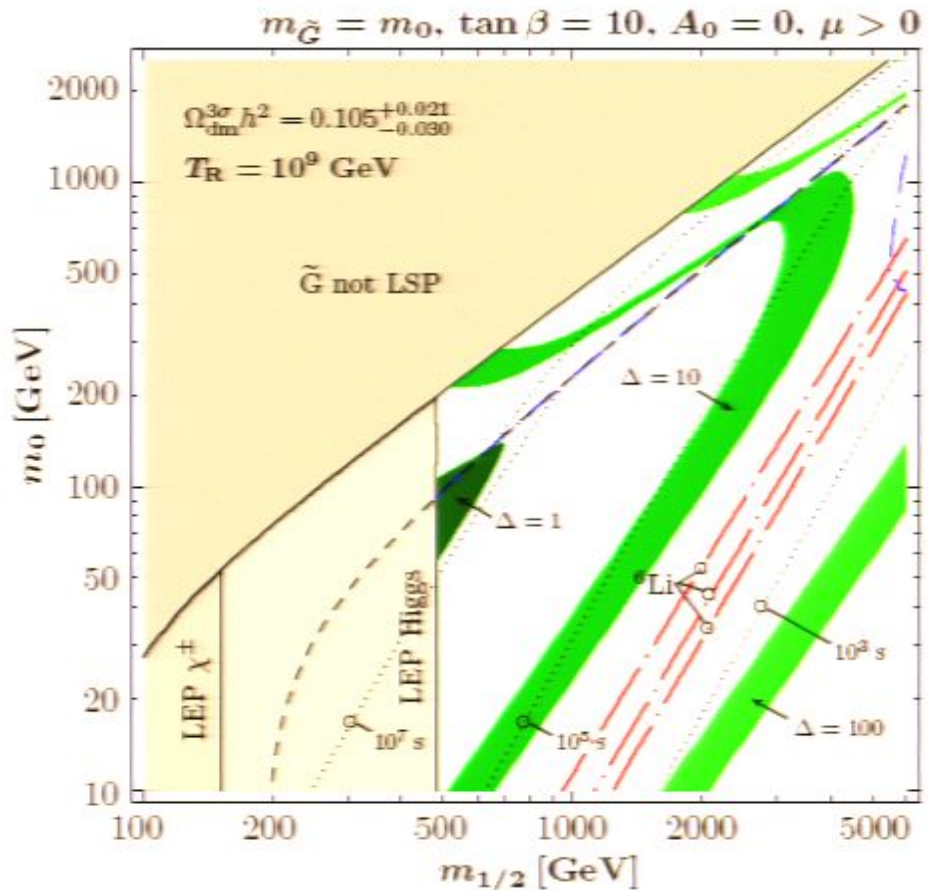
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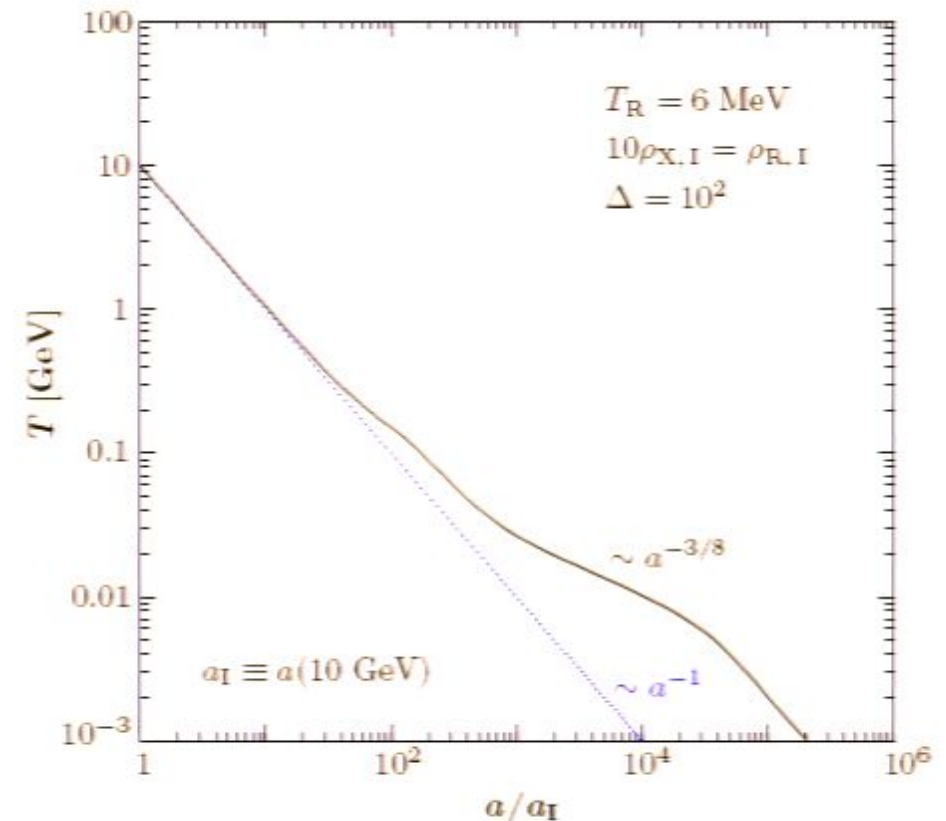
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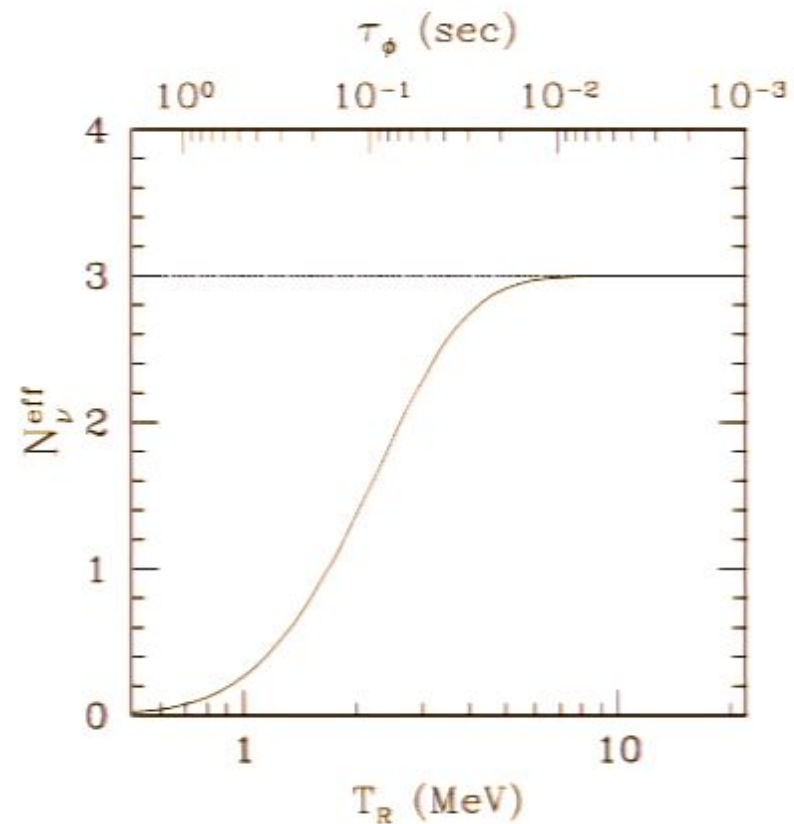
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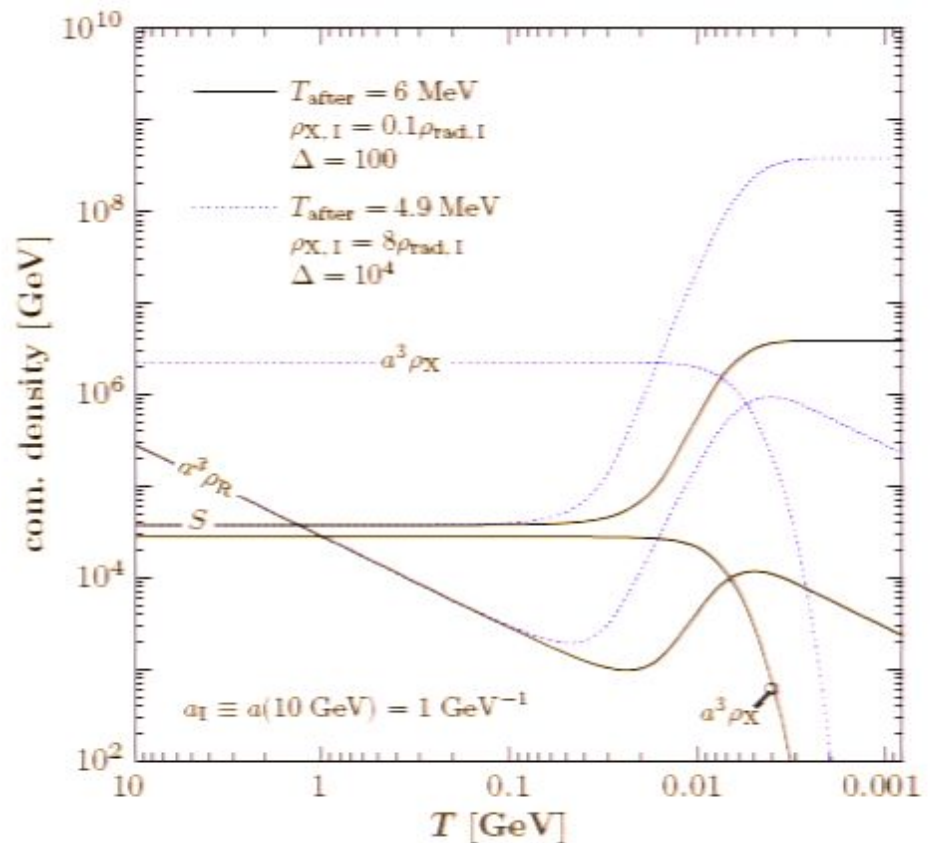
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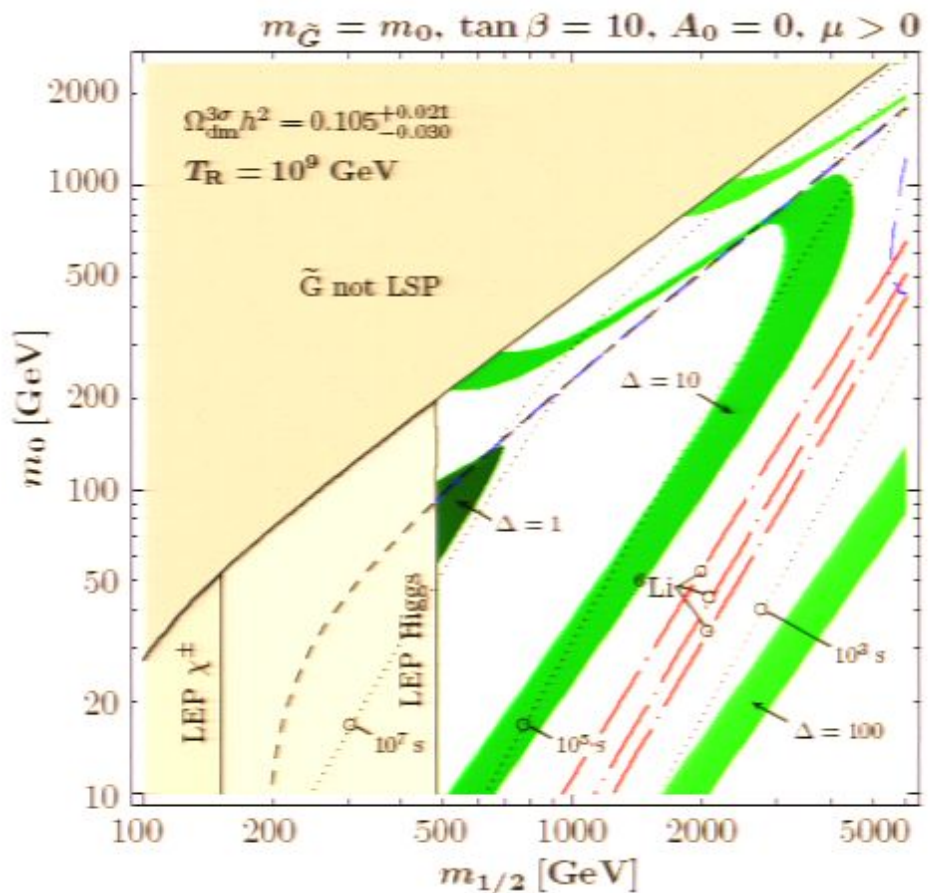
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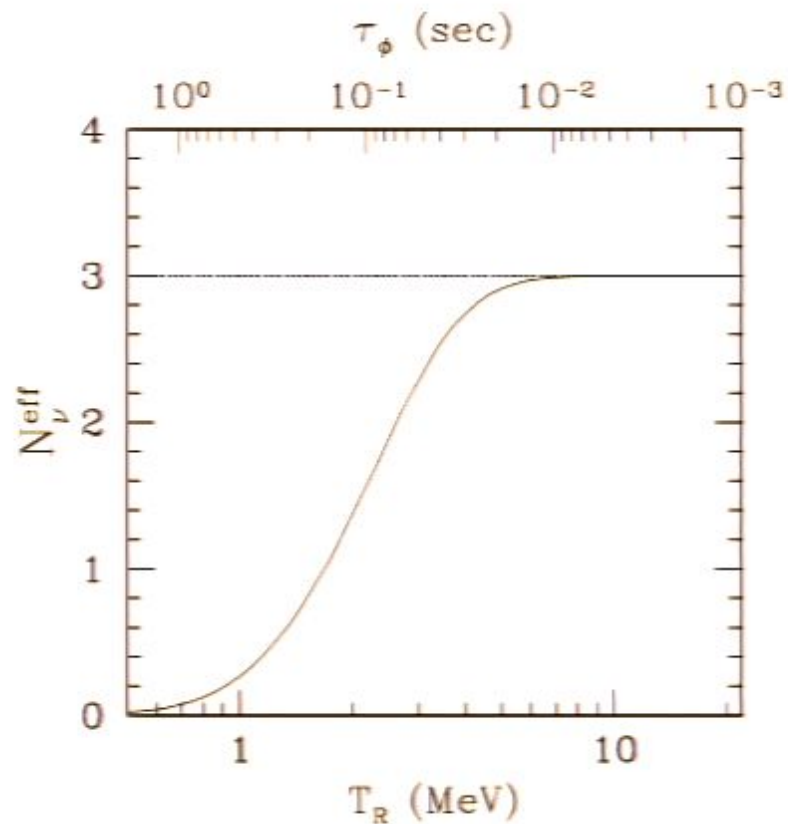
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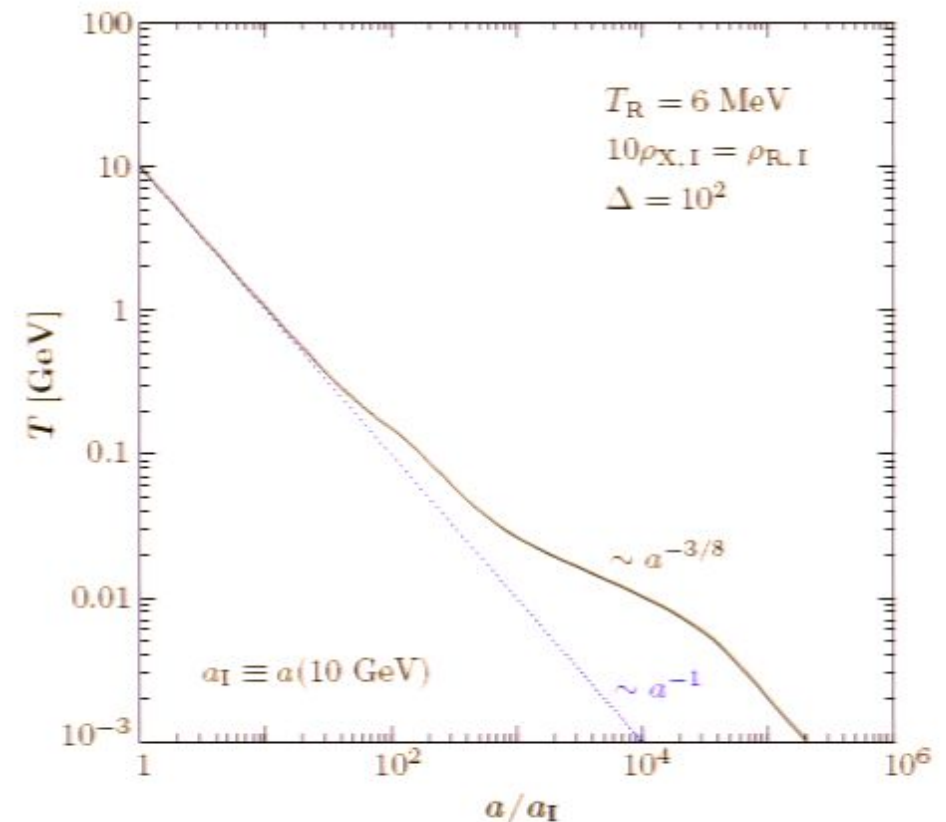
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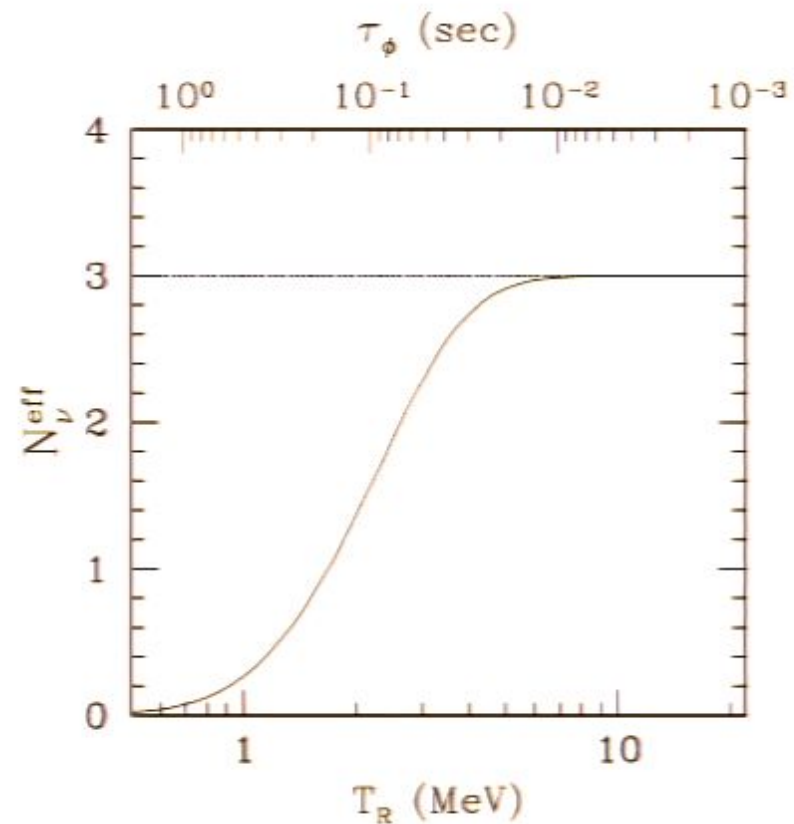
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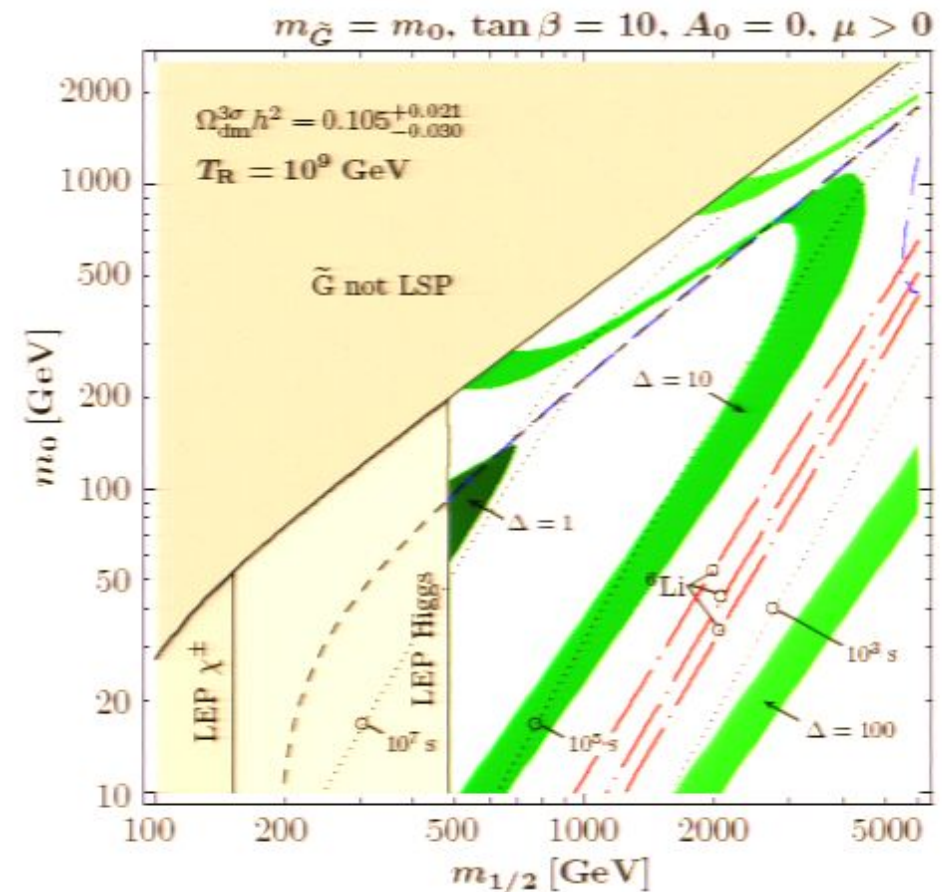
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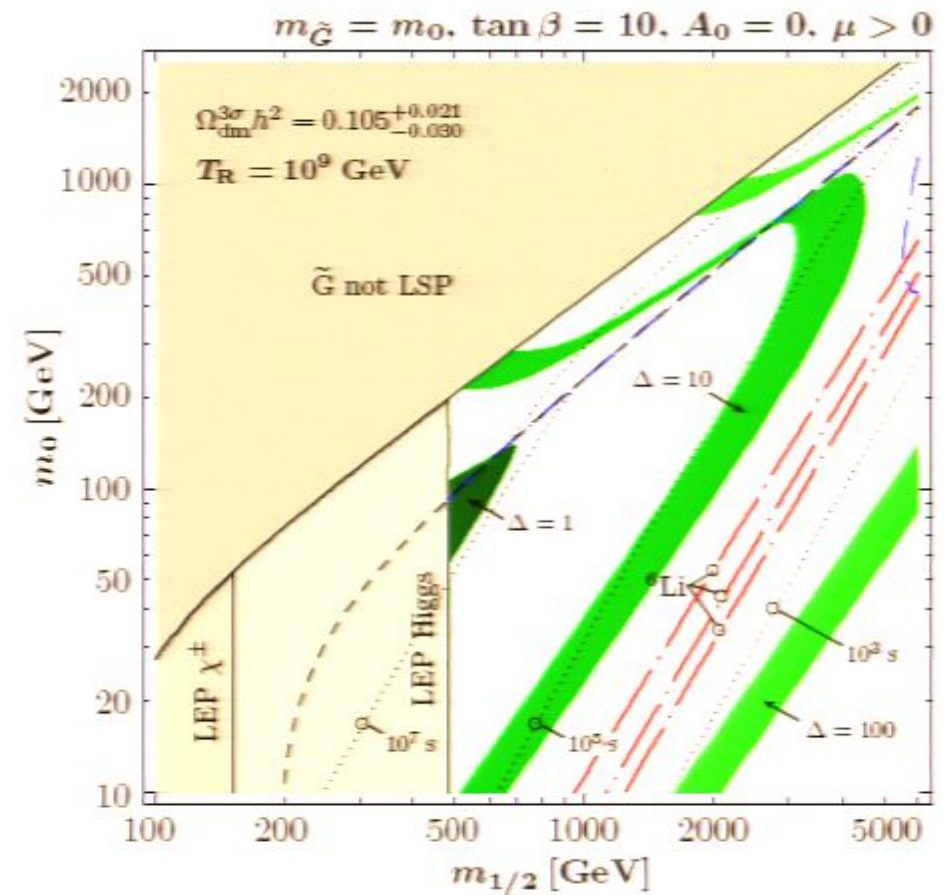
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Late-time Entropy Production

$\Delta = 100 \rightarrow T_R \simeq 10^9 \text{ GeV}$ OK



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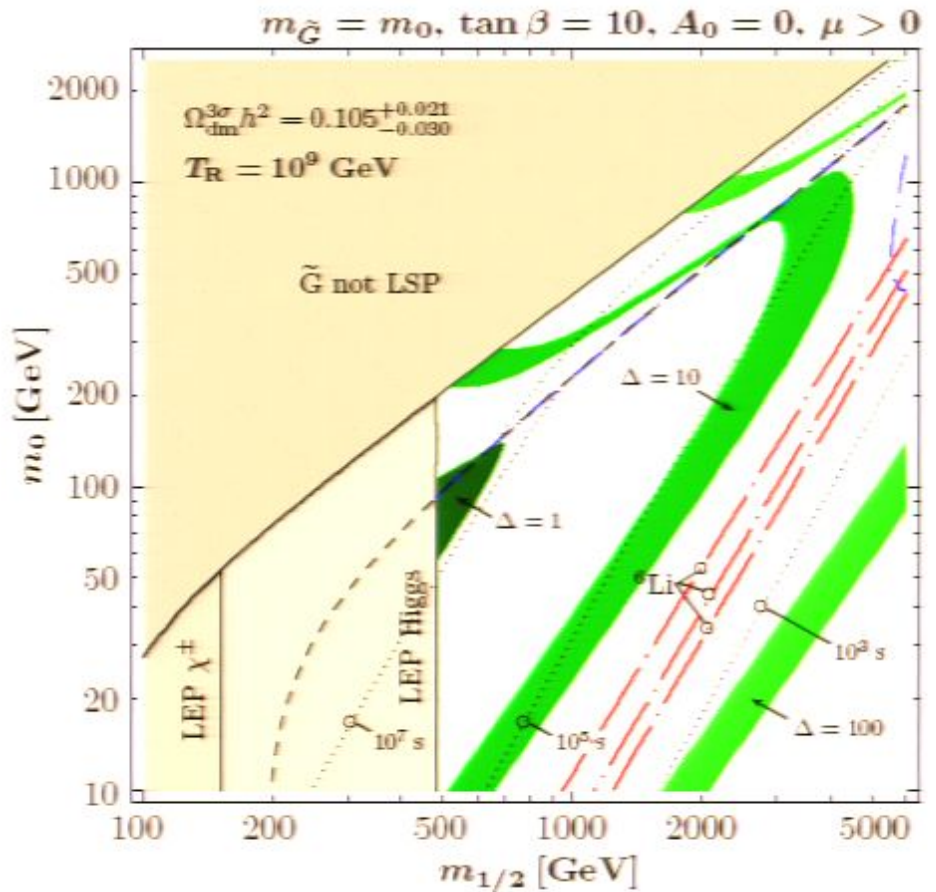
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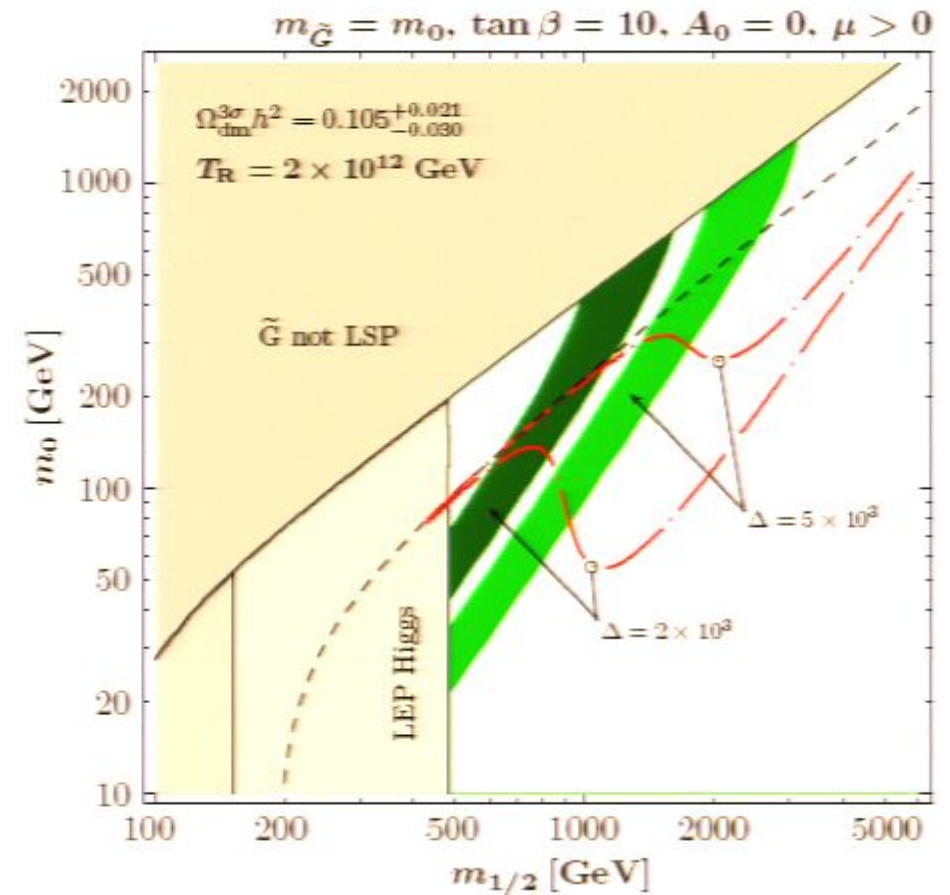
Late-time Entropy Production

$$\Delta \sim 10^3, T_R \sim 10^{12} \text{ GeV}$$

thermal leptogenesis OK

$$M_{R1} \sim T_R \sim 10^{12} \text{ GeV}$$

[Buchmüller, Di Bari, Plümacher, 2002]



Conclusions

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This talk summarizes

JP and F. D. Steffen, *Implications of Catalyzed BBN in the CMSSM with Gravitino Dark Matter*, arXiv:0710.2213 [hep-ph].

JP and F. D. Steffen, *Constraints on the reheating temperature in gravitino dark matter scenarios*, Phys. Lett. B **648** (2007) 224.

JP and F. D. Steffen, *Thermal gravitino production and collider tests of leptogenesis*, Phys. Rev. D **75** (2007) 023509.

No Signal

VGA-1