

Title: Universal resources for approximate and stochastic measurement-based quantum computation

Date: May 02, 2008 11:00 AM

URL: <http://pirsa.org/08050021>

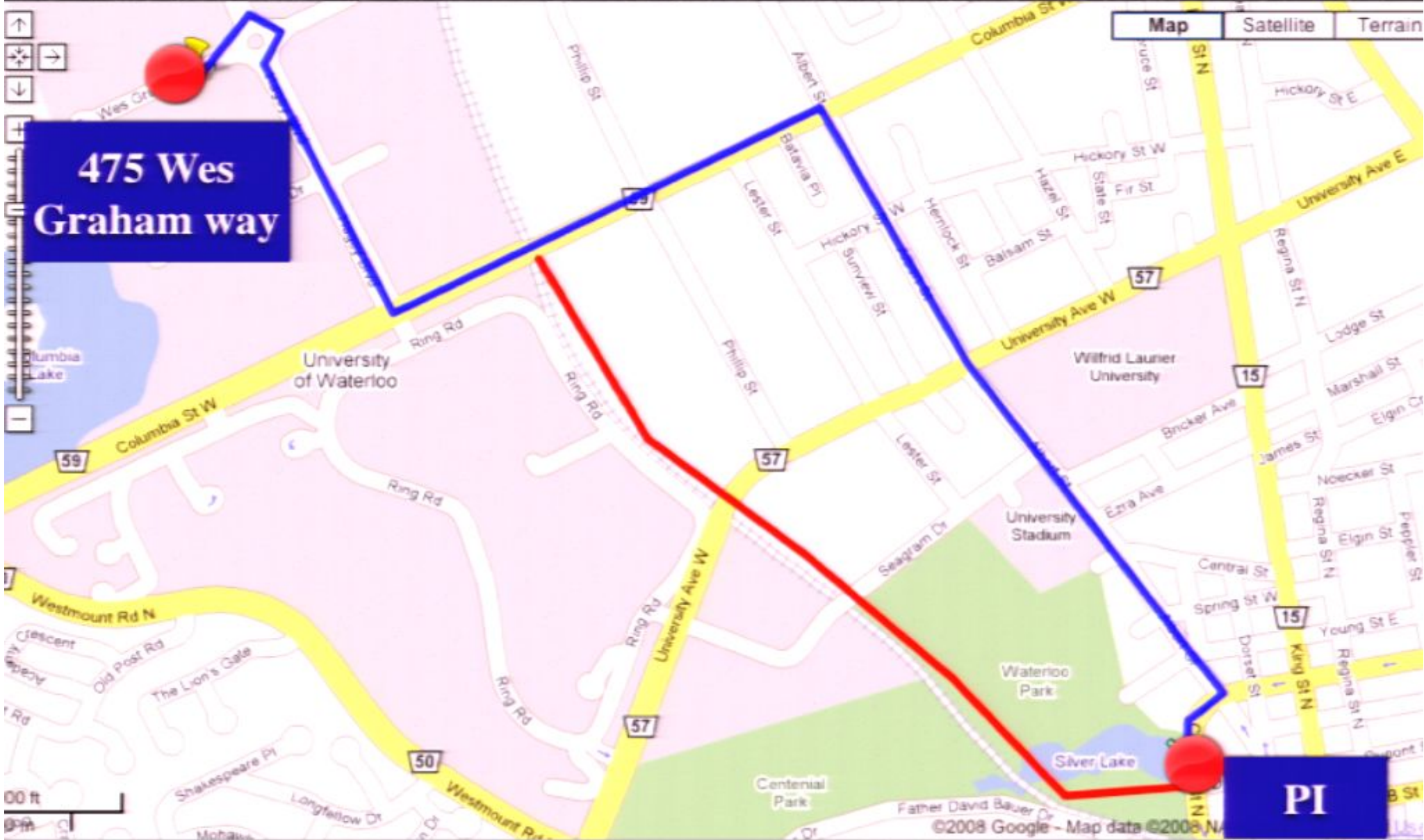
Abstract: We investigate which families of quantum states can be used as resources for approximate and/or stochastic universal measurement-based quantum computation, in the sense that single-qubit operations and classical communication are sufficient to prepare (with some fixed precision and/or probability) any quantum state from the initial resource. We find entanglement-based criteria for non-universality in the approximate and/or stochastic case. By applying them, we are able to discard some families of states as not universal also in this weaker sense. Finally, we show that any family Σ of states that is 'close' to an (approximate and/or stochastic) universal family Γ is approximate and stochastic universal, and we prove that if Γ was efficiently universal then also Σ is.

*Universal resources
for
approximate and stochastic
measurement-based quantum
computing*

~~*Katerina-E. Mora*~~
(IQC)

*Work in collaboration with
M. Piani, M. Van den Nest,
A. Miyake, W. Dür, H. Briegel*

Welcome barbecue this afternoon!



In the new temporary IOC building

Outline

- ◎ *Measurement-based computation*
 - ▶ *Universal exact “state preparators”*
 - ▶ *Entanglement-based criteria*
 - ▶ *Examples*

Outline

- ◎ *Measurement-based computation*
 - ▶ *Universal exact “state preparators”*
 - ▶ *Entanglement-based criteria*
 - ▶ *Examples*

- ◎ *Approximate/stochastic universality*
 - ▶ *Basic idea and tools: ϵ -measures of entanglement*
 - ▶ *Criteria in the approximate/stochastic case*
 - ▶ *Examples and “Go” results*

Outline

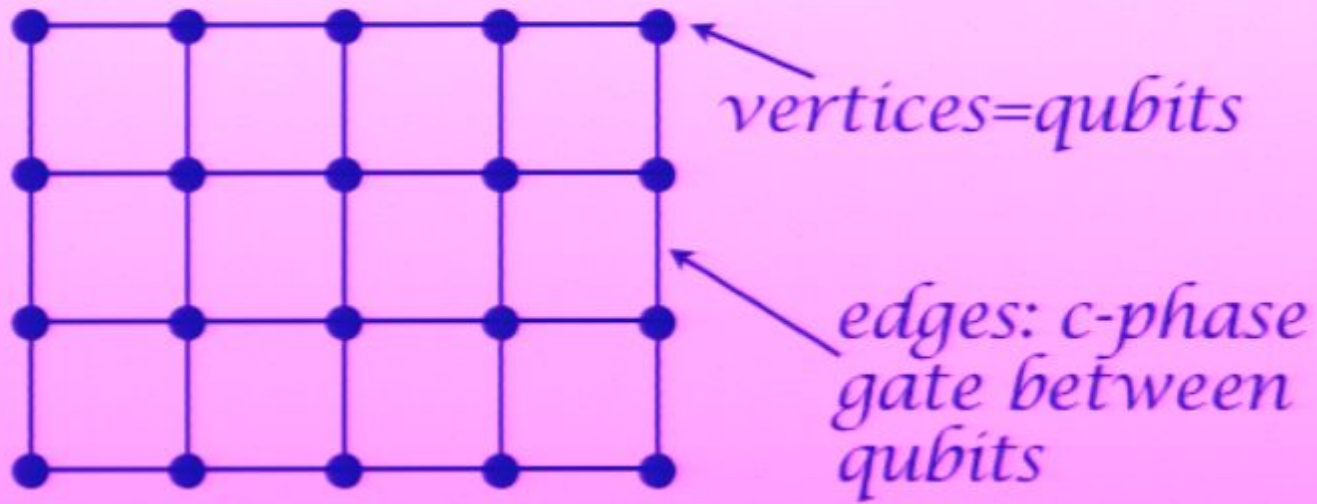
- ◎ *Measurement-based computation*
 - ▶ *Universal exact “state preparators”*
 - ▶ *Entanglement-based criteria*
 - ▶ *Examples*

- ◎ *Approximate/stochastic universality*
 - ▶ *Basic idea and tools: ϵ -measures of entanglement*
 - ▶ *Criteria in the approximate/stochastic case*
 - ▶ *Examples and “Go” results*

- ◎ *Conclusions*

The one-way model

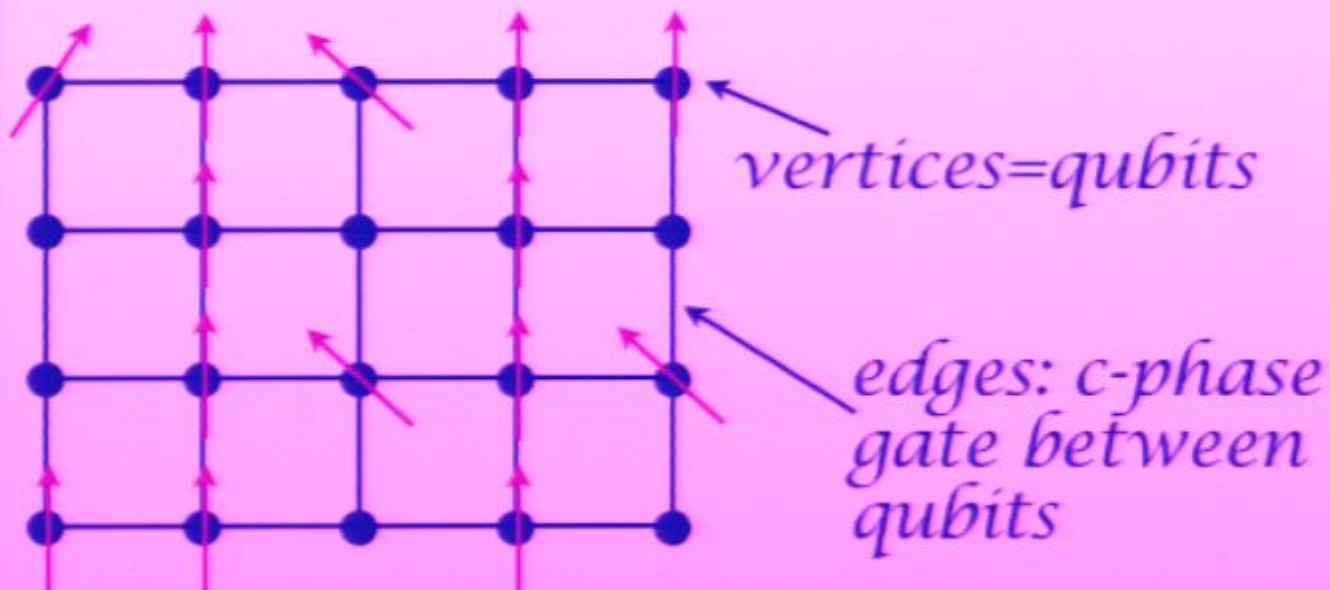
Cluster state



Remember
Dan Browne's talk!

The one-way model

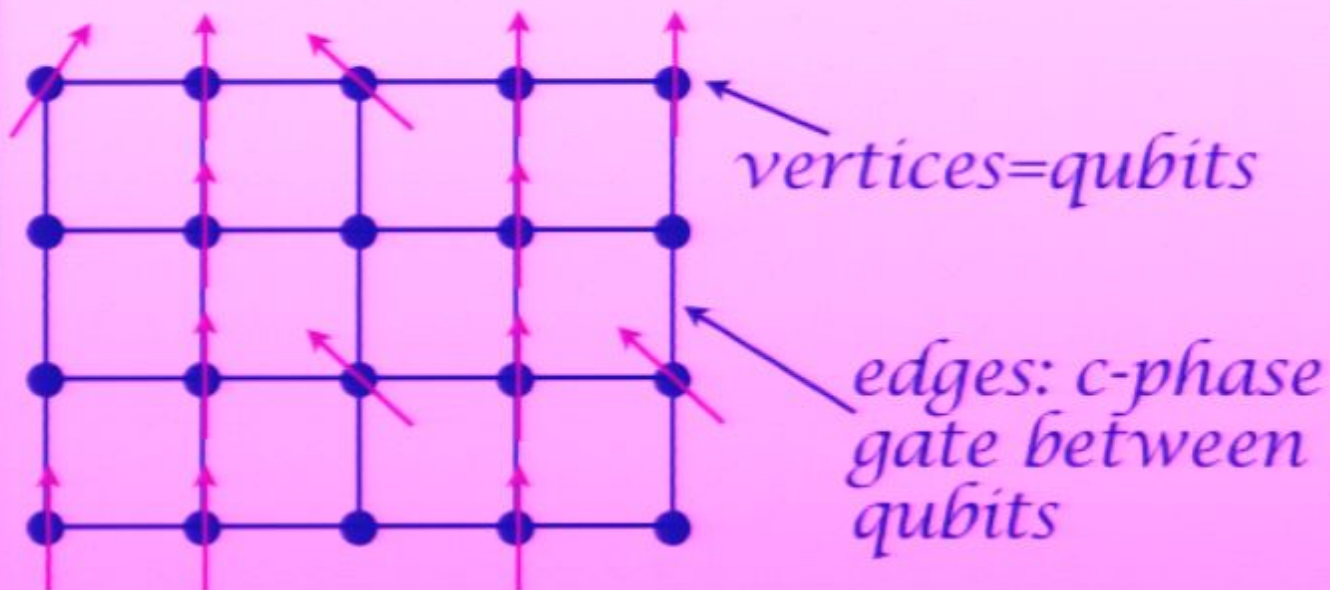
Cluster state + single-qubit measurements



Remember
Dan Browne's talk!

The one-way model

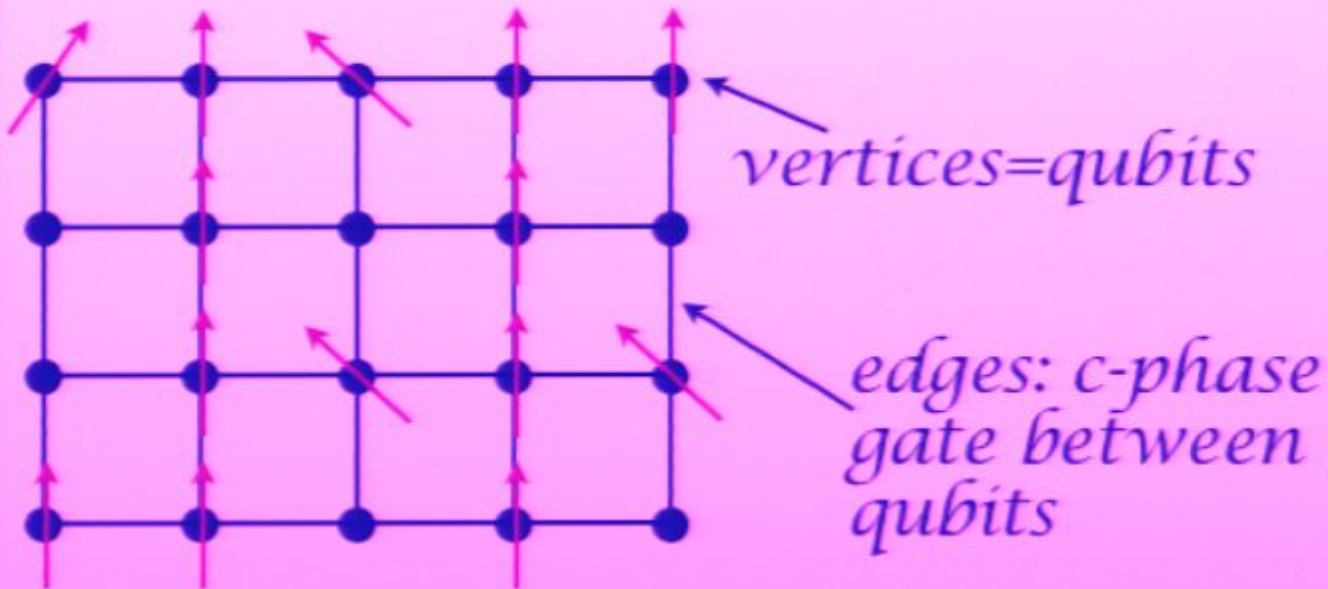
Cluster state + single-qubit measurements + classical communication



Remember
Dan Browne's talk!

The one-way model

Cluster state + single-qubit measurements + classical communication

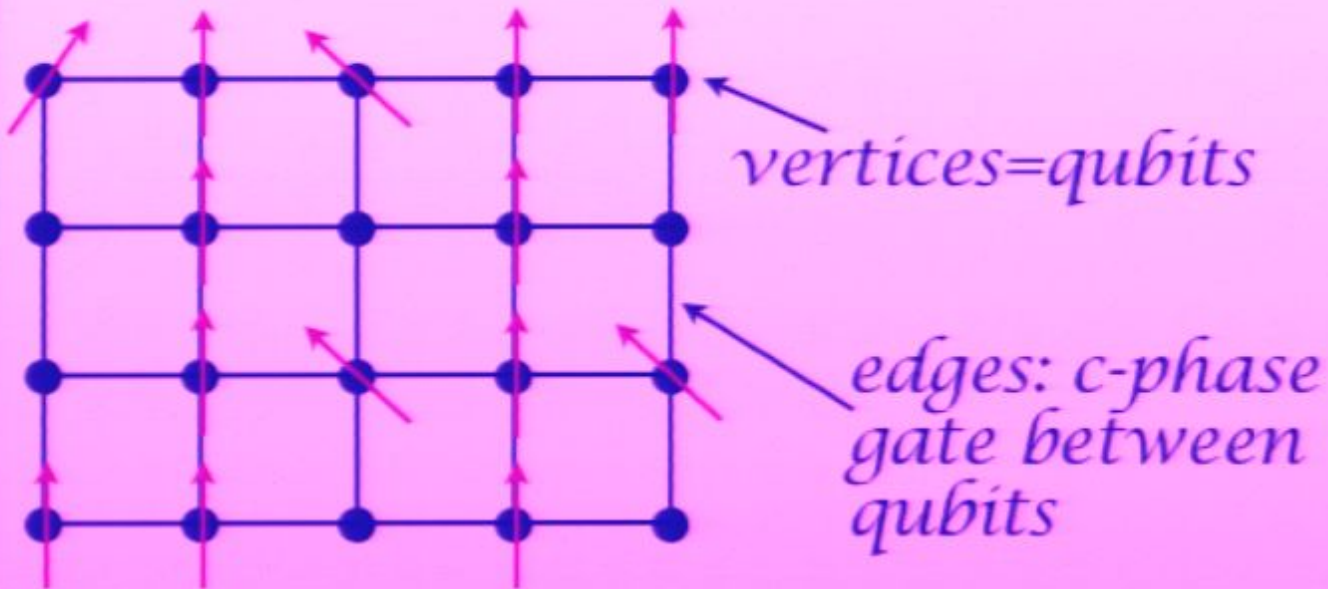


↓
quantum
computation

Remember
Dan Browne's talk!

The one-way model

Cluster state + single-qubit measurements + classical communication



↓
quantum
computation

➔ Equivalent to the circuit model

Remember
Dan Browne's talk!

State preparation and exact universality (I)

For any state $|\psi\rangle$ there are:

- ✓ a cluster state $|C_{N \times M}\rangle$ large enough
- ✓ a 1-qubit measurement pattern \mathcal{M}

State preparation and exact universality (I)

For any state $|\psi\rangle$ there are:

- ✓ a cluster state $|C_{N \times M}\rangle$ large enough
- ✓ a 1-qubit measurement pattern \mathcal{M}

such that

$$|C_N\rangle \xrightarrow{\mathcal{M}} |\psi\rangle$$

State preparation and exact universality (I)

For any state $|\psi\rangle$ there are:

- ✓ a cluster state $|C_{N \times M}\rangle$ large enough
- ✓ a 1-qubit measurement pattern \mathcal{M}

such that

$$|C_N\rangle \xrightarrow{\mathcal{M}} |\psi\rangle$$

exact state preparation

The family $\Psi_C = \{|C_{N \times M}\rangle\}$ of cluster states is universal for exact and deterministic quantum computation (state preparation)

State preparation and exact universality (II)

Just how “special” is the cluster state?

Cluster state + single-qubit measurements + classical communication

State preparation and exact universality (II)

Just how “special” is the cluster state?

Resource state + single-qubit measurements + classical communication

State preparation and exact universality (II)

Just how “special” is the cluster state?

Resource state + Local operations and classical communication

State preparation and exact universality (II)

Just how “special” is the cluster state?

Resource state + Local operations and classical communication



Which (families of) states still allow for universal quantum computation?

State preparation and exact universality (II)

Just how “special” is the cluster state?

Resource state + Local operations and classical communication



Which (families of) states still allow for universal quantum computation?

$$\Phi = \{|\varphi_k\rangle\}_k$$

family of quantum states

State preparation and exact universality (II)

Just how “special” is the cluster state?

Resource state + Local operations and classical communication



Which (families of) states still allow for universal quantum computation?

$\Phi = \{|\varphi_k\rangle\}_k$
family of
quantum states

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$
such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

State preparation and exact universality (II)

Just how “special” is the cluster state?

Resource state + Local operations and classical communication



Which (families of) states still allow for universal quantum computation?

Universal quantum computation

=

Universal state preparation

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

LOCC operations cannot increase entanglement!

→ $E(|\varphi\rangle) \geq E(|\psi\rangle)$ for any entanglement measure E

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

LOCC operations cannot increase entanglement!

➔ $\mathcal{E}(|\varphi\rangle) \geq \mathcal{E}(|\psi\rangle)$ for any entanglement measure \mathcal{E}

If Φ is universal, we can choose any $|\psi\rangle$

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

LOCC operations cannot increase entanglement!

→ $\mathcal{E}(|\varphi\rangle) \geq \mathcal{E}(|\psi\rangle)$ for any entanglement measure \mathcal{E}

If Φ is universal, we can choose any $|\psi\rangle$

$$\mathcal{E}^* = \sup_{|\psi\rangle} \mathcal{E}(|\psi\rangle) \quad (\mathcal{E}^* = \infty \text{ in many cases})$$

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

LOCC operations cannot increase entanglement!

→ $\mathcal{E}(|\varphi\rangle) \geq \mathcal{E}(|\psi\rangle)$ for any entanglement measure \mathcal{E}

If Φ is universal, we can choose any $|\psi\rangle$

$$\mathcal{E}^* = \sup_{|\psi\rangle} \mathcal{E}(|\psi\rangle) \quad (\mathcal{E}^* = \infty \text{ in many cases})$$

→ If Φ is universal, $\mathcal{E}(\Phi) = \mathcal{E}^*$ for any measure \mathcal{E}

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

LOCC operations cannot increase entanglement!

→ $\mathcal{E}(|\varphi\rangle) \geq \mathcal{E}(|\psi\rangle)$ for any entanglement measure \mathcal{E}

If Φ is universal, we can choose any $|\psi\rangle$

$$\mathcal{E}^* = \sup_{|\psi\rangle} \mathcal{E}(|\psi\rangle) \quad (\mathcal{E}^* = \infty \text{ in many cases})$$

→ If Φ is universal, $\mathcal{E}(\Phi) = \mathcal{E}^*$ for any measure \mathcal{E}

$$\mathcal{E}(\Phi) = \sup_{|\varphi\rangle \in \Phi} \mathcal{E}(|\varphi\rangle)$$

Entanglement and universality (The exact deterministic case)

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

LOCC operations cannot increase entanglement!

→ $E(|\varphi\rangle) \geq E(|\psi\rangle)$ for any entanglement measure E

If Φ is universal, we can choose any $|\psi\rangle$

$$E^* = \sup_{|\psi\rangle} E(|\psi\rangle) \quad (E^* = \infty \text{ in many cases})$$

→ If Φ is universal, $E(\Phi) = E^*$ for any entanglement measure E

$$E(\Phi) = \sup_{|\varphi\rangle \in \Phi} E(|\varphi\rangle)$$

Computing
 $E(\Phi)$ and E^* we prove
that some families are not
universal for MBQC

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N=2}^{\infty} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N=2}^{\infty} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $E(|\psi\rangle) = 1 - \sup \langle \varphi_{sep} | \psi \rangle$

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) = 1 - \sup \langle \varphi_{\text{sep}} | \psi \rangle$



$$\mathcal{E}^* = 1 - 2^{-n}$$

$$\mathcal{E}(\Phi_W) = 1 - 1/e$$

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N=3} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) =$



$$\mathcal{E}^* = 1 - 2^{-n} > \mathcal{E}(\Phi_W) = 1 - 1/e$$

W states
are not
universal!

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) =$



$$\mathcal{E}^* = 1 - 2^{-n} > \mathcal{E}(\Phi_W) = 1 - 1/e$$

W states
are not
universal!

Some graph states
(graphs with bounded rank)

Ground states
(of non-critical 1-D systems)

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) =$



$$\mathcal{E}^* = 1 - 2^{-n} > \mathcal{E}(\Phi_W) = 1 - 1/e$$

W states
are not
universal!

Some graph states

(graphs with bounded rank)

Ground states

(of non-critical 1-D systems)

Use Schmidt-rank width: multipartite entanglement measure associated to Schmidt rank

$$\mathcal{E}^* = \infty$$

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) = \max_{|a\rangle} |\langle a|\psi\rangle|^2$



$$\mathcal{E}^* = 1 - 2^{-n} > \mathcal{E}(\Phi_W) = 1 - 1/e$$

W states
are not
universal!

Some graph states

(graphs with bounded rank)

Ground states

(of non-critical 1-D systems)

Use Schmidt-rank width: multipartite entanglement measure associated to Schmidt rank

$$\mathcal{E}^* = \infty > \mathcal{E} \text{ bounded in these cases}$$

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 3} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) =$



$$\mathcal{E}^* = 1 - 2^{-n} > \mathcal{E}(\Phi_W) = 1 - 1/e$$

W states
are not
universal!

Some graph states

(graphs with bounded rank)

Ground states

(of non-critical 1-D systems)

Use Schmidt-rank width: multipartite entanglement

Not associated to Schmidt rank

Not
universal
either

$$\mathcal{E}^* = \infty > \mathcal{E} \text{ bounded in these cases}$$

Some examples (I)

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 3} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Consider the geometric measure: $\mathcal{E}(|\psi\rangle) =$



$$\mathcal{E}^* = 1 - 2^{-n} > \mathcal{E}(\Phi_W) = 1 - 1/e$$

W states
are not
universal!

Some graph states

(graphs with bounded rank)

Ground states

(of non-critical 1-D systems)

Use Schmidt-rank width: multipartite entanglement associated to Schmidt rank

Not
universal
either

$$\mathcal{E}^* = \infty > \mathcal{E} \text{ bounded in these cases}$$

Proof:

Sorry! not
enough space in
this small margin

Approximate and stochastic universality (I)

Exact universality: for all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that $|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$

Approximate and stochastic universality (I)

Exact universality: for all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that $|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$

Obs: Often exact preparation is not necessary (or possible)

Approximate and stochastic universality (I)

Exact universality: for all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that $|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$

Obs: Often exact preparation is not necessary (or possible)

➔ *Approximate/stochastic* universality

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that “most of the time” the output is “close enough” to $|\psi\rangle$.

Approximate and stochastic universality (II)

→ *Approximate/stochastic universality*

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that “most of the time” the output is “close enough” to $|\psi\rangle$.

Approximate and stochastic universality (II)

→ *Approximate/stochastic universality*

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that “most of the time” the output is “close enough” to $|\psi\rangle$.

More precisely:

$|\varphi\rangle$

Approximate and stochastic universality (II)

→ *Approximate/stochastic* universality

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that "most of the time" the output is "close enough" to $|\psi\rangle$.

More precisely:

$|\varphi\rangle$

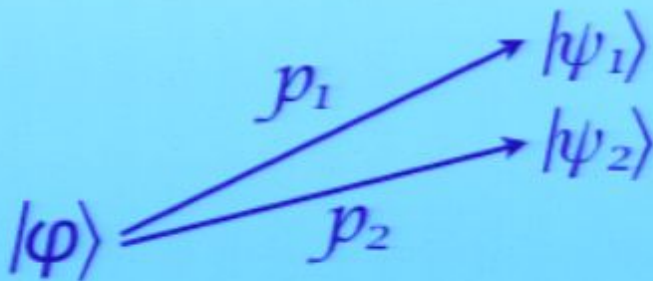
LOCC

Approximate and stochastic universality (II)

➔ *Approximate/stochastic* universality

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that "most of the time" the output is "close enough" to $|\psi\rangle$.

More precisely:



LOCC

Approximate and stochastic universality (II)

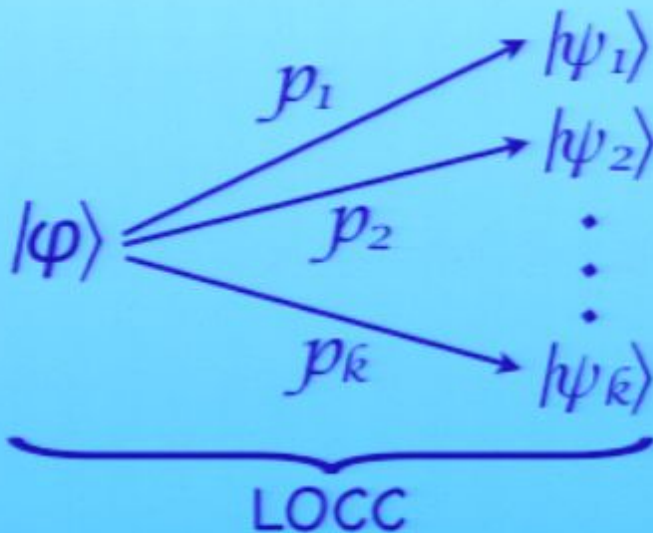
→ *Approximate/stochastic* universality

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that "most of the time" the output is "close enough" to $|\psi\rangle$.

More precisely:

ϵ -approximate:

$$D(\psi, \psi_j) \leq \epsilon$$

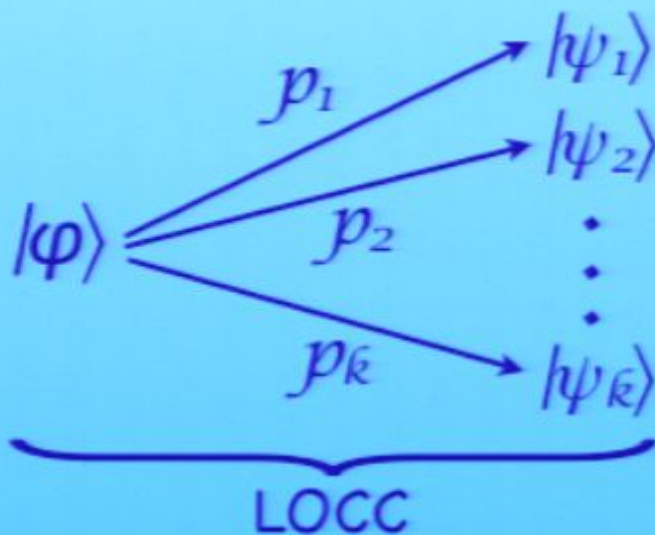


Approximate and stochastic universality (II)

→ *Approximate/stochastic universality*

For all $|\psi\rangle$ there is $|\varphi\rangle \in \Phi$ such that "most of the time" the output is "close enough" to $|\psi\rangle$.

More precisely:



ϵ -approximate:

$$D(\psi, \psi_j) \leq \epsilon$$

δ -stochastic:

$$\sum_{j:\epsilon\text{-close}} p_j \geq 1 - \delta$$

Entanglement and universality (The approximate and stochastic case)

Entanglement gave us criteria in the exact deterministic case.

Entanglement and universality (The approximate and stochastic case)

Entanglement gave us criteria in the exact deterministic case.

Recall!
 $\mathcal{E}(\Phi) = \mathcal{E}^*$
for all
universal families

Entanglement and universality (The approximate and stochastic case)

Entanglement gave us criteria in the exact deterministic case.




And in the ε -approximate and δ -stochastic case?


Recall!
 $\mathcal{E}(\Phi) = \mathcal{E}^*$

for all
universal families

Entanglement and universality (The approximate and stochastic case)

Entanglement gave us criteria in the exact deterministic case.


 And in the ε -approximate and δ -stochastic case?


 We can still find good criteria, but we need a new “type” of entanglement measure

Recall!
 $\mathcal{E}(\Phi) = \mathcal{E}^*$
for all
universal families

Entanglement and universality (The approximate and stochastic case)

Entanglement gave us criteria in the exact deterministic case.

 And in the ε -approximate and δ -stochastic case?

 We can still find good criteria, but we need a new “type” of entanglement measure

Idea: we are happy to go close to the state: we don't need to reach maximal entanglement

Recall!
 $\mathcal{E}(\Phi) = \mathcal{E}^*$
for all
universal families

Entanglement and universality (ϵ -measures of entanglement)

Consider a state ρ

• ρ

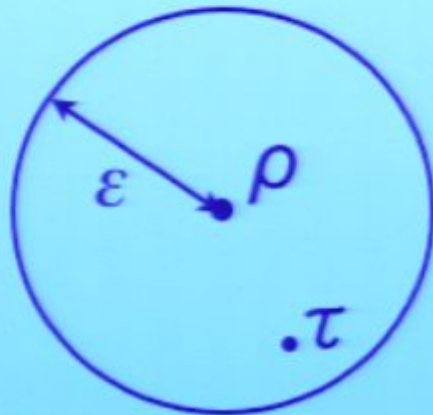
Entanglement and universality (ϵ -measures of entanglement)

Consider a state ρ and a ball of radius ϵ around it



Entanglement and universality (ε -measures of entanglement)

Consider a state ρ and a ball of radius ε around it

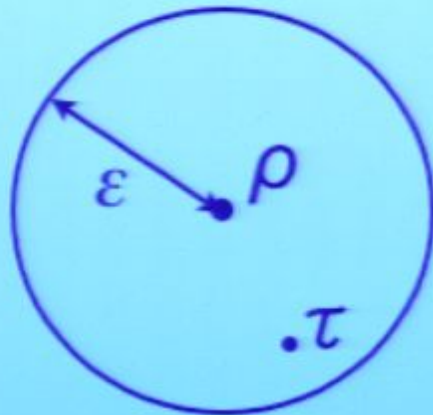


Given an entanglement measure \mathcal{E} , consider:

$$\mathcal{E}_\varepsilon(\rho) = \min\{E(\tau) \text{ s.t. } D(\rho, \tau) \leq \varepsilon\}$$

Entanglement and universality (ε -measures of entanglement)

Consider a state ρ and a ball of radius ε around it



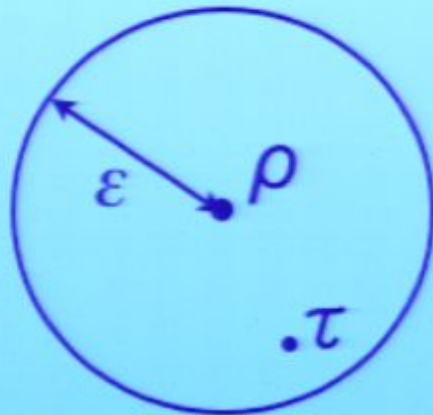
Given an entanglement measure \mathcal{E} , consider:

$$\mathcal{E}_\varepsilon(\rho) = \min\{E(\tau) \text{ s.t. } D(\rho, \tau) \leq \varepsilon\}$$

➔ The ε -version of a measure is a good entanglement measure:

Entanglement and universality (ε -measures of entanglement)

Consider a state ρ and a ball of radius ε around it



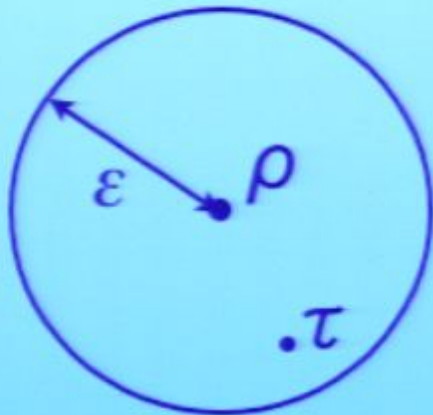
Given an entanglement measure \mathcal{E} , consider:

$$\mathcal{E}_\varepsilon(\rho) = \min\{E(\tau) \text{ s.t. } D(\rho, \tau) \leq \varepsilon\}$$

➔ The ε -version of a measure is a good entanglement measure: \mathcal{E}_ε is zero on separable states

Entanglement and universality (ε -measures of entanglement)

Consider a state ρ and a ball of radius ε around it



Given an entanglement measure \mathcal{E} , consider:

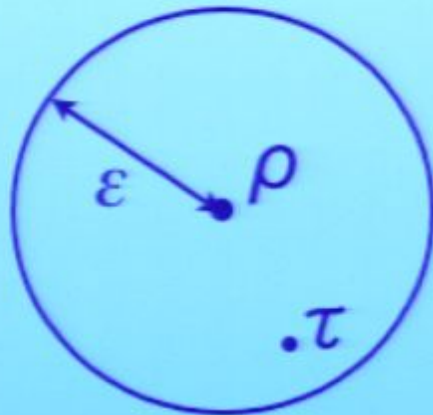
$$\mathcal{E}_\varepsilon(\rho) = \min\{E(\tau) \text{ s.t. } D(\rho, \tau) \leq \varepsilon\}$$

➔ The ε -version of a measure is a good entanglement measure: \mathcal{E}_ε is zero on separable states

\mathcal{E}_ε does not increase under LOCC

Entanglement and universality (ε -measures of entanglement)

Consider a state ρ and a ball of radius ε around it



Given an entanglement measure \mathcal{E} , consider:

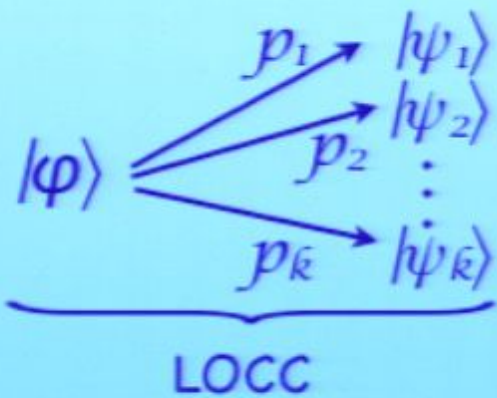
$$\mathcal{E}_\varepsilon(\rho) = \min\{E(\tau) \text{ s.t. } D(\rho, \tau) \leq \varepsilon\}$$

➔ The ε -version of a measure is a good entanglement measure: \mathcal{E}_ε is zero on separable states

\mathcal{E}_ε does not increase under LOCC

✓ Measures the entanglement contained in a *state that is known partially* (as for imperfect preparation)

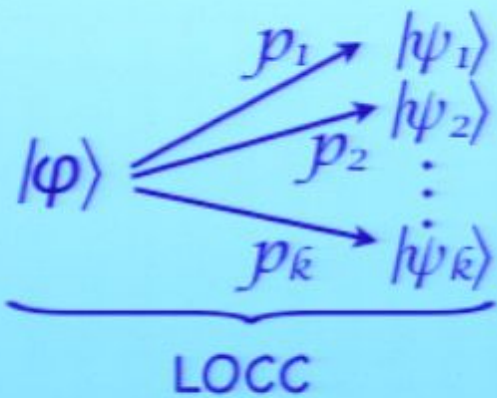
Entanglement and universality (The approximate/stochastic case)



1) $\delta=0$: all $|\psi_j\rangle$ must be "close enough" to $|\psi\rangle$

$$\sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

Entanglement and universality (The approximate/stochastic case)



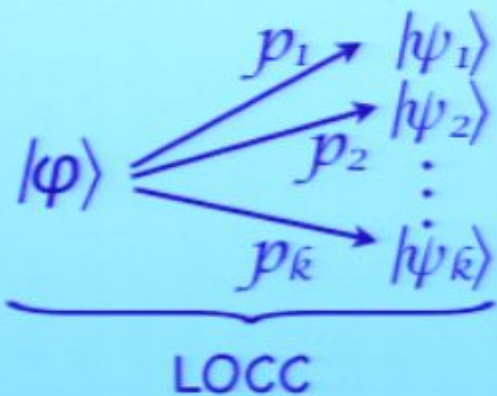
$$\sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

1) $\delta=0$: all $|\psi_j\rangle$ must be "close enough" to $|\psi\rangle$



We must be able to "reach" the minimum entanglement contained in a state close to $|\psi\rangle$

Entanglement and universality (The approximate/stochastic case)



$$\sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

1) $\delta=0$: all $|\psi_j\rangle$ must be "close enough" to $|\psi\rangle$



We must be able to "reach" the minimum entanglement contained in a state close to $|\psi\rangle$

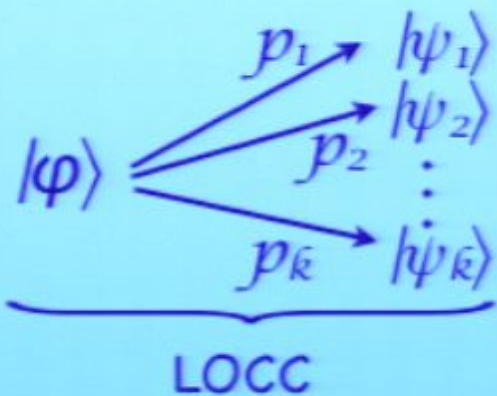
$$\Rightarrow \mathcal{E}(\Phi) \geq \mathcal{E}_\epsilon^*$$

Exact case:

$$\mathcal{E}(\Phi) \geq \mathcal{E}^*$$

with $\mathcal{E}^* \geq \mathcal{E}_\epsilon^*$

Entanglement and universality (The approximate/stochastic case)



$$\sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

1) $\delta=0$: all $|\psi_j\rangle$ must be "close enough" to $|\psi\rangle$



We must be able to "reach" the minimum entanglement contained in a state close to $|\psi\rangle$

$$\Rightarrow \mathcal{E}(\Phi) \geq \mathcal{E}_\epsilon^*$$

Exact case:

$$\mathcal{E}(\Phi) \geq \mathcal{E}^*$$

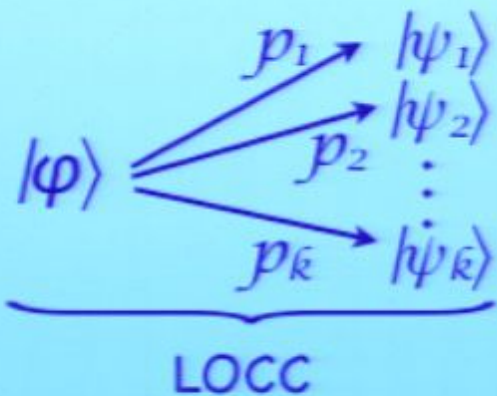
with $\mathcal{E}^* \geq \mathcal{E}_\epsilon^*$

General case: if $\delta > 0$, less entanglement is sufficient (some of the outputs can be thrown away)

More examples

Computing $\mathcal{E}_\varepsilon^$ is hard in general, but can be done!*

Entanglement and universality (The approximate/stochastic case)



$$\sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

1) $\delta=0$: all $|\psi_j\rangle$ must be "close enough" to $|\psi\rangle$



We must be able to "reach" the minimum entanglement contained in a state close to $|\psi\rangle$

$$\Rightarrow \mathcal{E}(\Phi) \geq \mathcal{E}_\epsilon^*$$

Exact case:

$$\mathcal{E}(\Phi) \geq \mathcal{E}^*$$

with $\mathcal{E}^* \geq \mathcal{E}_\epsilon^*$

General case: if $\delta > 0$, less entanglement is sufficient (some of the outputs can be thrown away)

$$\Rightarrow \mathcal{E}(\Phi) \geq (1 - \delta) \mathcal{E}_\epsilon^*$$

More examples

Computing $\mathcal{E}_\varepsilon^$ is hard in general, but can be done!*

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Again the geometric measure

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N=3} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Again the geometric measure

$$\mathcal{E}_\varepsilon^*(|\psi\rangle) \geq 1 - f(\varepsilon)$$

$$\mathcal{E}(\Phi_W) = 1 - 1/e$$

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Again the geometric measure

$$\mathcal{E}_\varepsilon^*(|\psi\rangle) \geq 1 - f(\varepsilon)$$

$$\mathcal{E}(\Phi_W) = 1 - 1/e$$

Reasonable function
 $f(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Again the geometric measure

$$\mathcal{E}_\varepsilon^*(|\psi\rangle) \geq 1 - f(\varepsilon)$$

Reasonable function
 $f(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$

>

$$\mathcal{E}(\Phi_W) = 1 - 1/e$$

for ε small

W states
are not
approximate
universal!
for ε small

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 1} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Again the geometric measure

$$\mathcal{E}_\varepsilon^*(|\psi\rangle) \geq 1 - f(\varepsilon)$$

Reasonable function
 $f(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$

>

$\mathcal{E}(\Phi_W) = 1 - 1/c$ In other cases, compute $\mathcal{E}_\varepsilon^*$ for $\varepsilon \rightarrow 0$

for ε small

W states are not approximate universal! for ε small

More examples

Computing $\mathcal{E}_\varepsilon^*$ is hard in general, but can be done!

W states

$$\Phi_W = \{ |W_N\rangle \}_{N \geq 3} \quad |W_3\rangle \propto |100\rangle + |010\rangle + |001\rangle$$

Again the geometric measure

$$\mathcal{E}_\varepsilon^*(|\psi\rangle) \geq 1 - f(\varepsilon)$$

Reasonable function
 $f(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$

>

$$\mathcal{E}(\Phi_W) = 1 - 1/c$$

In other cases, compute $\mathcal{E}_\varepsilon^*$ for $\varepsilon \rightarrow 0$

for ε small

W states are not approximate universal! for ε small

Previous examples not ε -approximate universal for some values of ε

Efficiency and universality (I)

Different questions:

- 1) Can we prepare any state?
- 2) Is it very expensive?

Efficiency and universality (I)

Different questions:

1) Can we prepare any state?

2) Is it very expensive?

time

resources

Efficiency and universality (I)

Different questions:

1) Can we prepare any state? (universality)

2) Is it very expensive?

time

resources

Efficiency and universality (I)

Different questions:

1) Can we prepare any state? (universality)

2) Is it very expensive? (efficiency)

time

resources

Efficiency and universality (I)

Different questions:

1) Can we prepare any state? (universality)

2) Is it very expensive? (efficiency)

time

resources

Φ is *efficiently universal* if, for all $|\psi\rangle \in (\mathbb{C}^2)^n$ there is $|\varphi\rangle \in \Phi$ such that

Efficiency and universality (I)

Different questions:

1) Can we prepare any state? (universality)

2) Is it very expensive? (efficiency)

time resources

Φ is *efficiently universal* if, for all $|\psi\rangle \in (\mathbb{C}^2)^n$ there is $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$$

Efficiency and universality (I)

Different questions:

1) Can we prepare any state? (universality)

2) Is it very expensive? (efficiency)

time resources

Φ is *efficiently universal* if, for all $|\psi\rangle \in (\mathbb{C}^2)^n$ there is $|\varphi\rangle \in \Phi$ such that

$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$ with $\left\{ \begin{array}{l} \text{polynomial} \\ \text{resources} \end{array} \right. \quad |\varphi\rangle \in (\mathbb{C}^2)^{O(\text{poly}(n))}$

Efficiency and universality (I)

Different questions:

1) Can we prepare any state? (universality)

2) Is it very expensive? (efficiency)

time resources

Φ is *efficiently universal* if, for all $|\psi\rangle \in (\mathbb{C}^2)^n$ there is $|\varphi\rangle \in \Phi$ such that

$|\varphi\rangle \xrightarrow{\text{LOCC}} |\psi\rangle$ with $\left\{ \begin{array}{l} \text{polynomial} \\ \text{resources} \end{array} \right. \quad |\varphi\rangle \in (\mathbb{C}^2)^{O(\text{poly}(n))}$

$\left\{ \begin{array}{l} \text{polynomial} \\ \text{time} \end{array} \right. \quad \text{At most } \text{poly}(n) \text{ LOCC steps}$

Efficiency and universality (II)

Can we prepare the states?

Efficiency and universality (II)

Can we prepare the states?



For any $|\psi\rangle \in (\mathbb{C}^2)^n$ is there $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} \{p_k, |\psi_k\rangle\}_k ; \sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

Efficiency and universality (II)

Can we prepare the states efficiently?



For any $|\psi\rangle \in (\mathbb{C}^2)^n$ is there $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} \{p_k, |\psi_k\rangle\}_k ; \quad \sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

Efficiency and universality (II)

Can we prepare the states efficiently?



For any $|\psi\rangle \in (\mathbb{C}^2)^n$ is there $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} \{p_k, |\psi_k\rangle\}_k ; \quad \sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

polynomial resources

$$|\varphi\rangle \in (\mathbb{C}^2)^{O(\text{poly}(n, 1/\delta, 1/\epsilon))}$$

Efficiency and universality (II)

Can we prepare the states efficiently?



For any $|\psi\rangle \in (\mathbb{C}^2)^n$ is there $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} \{p_k, |\psi_k\rangle\}_k ; \quad \sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

polynomial resources

$$|\varphi\rangle \in (\mathbb{C}^2)^{O(\text{poly}(n, 1/\delta, 1/\epsilon))}$$

polynomial time

$$O(\text{poly}(n, 1/\delta, 1/\epsilon)) \\ \text{LOCC steps}$$

Efficiency and universality (II)

Can we prepare the states efficiently?



For any $|\psi\rangle \in (\mathbb{C}^2)^n$ is there $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} \{p_k, |\psi_k\rangle\}_k ; \quad \sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

polynomial resources

$$|\varphi\rangle \in (\mathbb{C}^2)^{O(\text{poly}(n, 1/\delta, 1/\epsilon))}$$

polynomial time

$$O(\text{poly}(n, 1/\delta, 1/\epsilon)) \\ \text{LOCC steps}$$



In the circuit model: $[O(\text{poly}(n \log 1/\epsilon))]$

Efficiency and universality (II)

Can we prepare the states efficiently?



For any $|\psi\rangle \in (\mathbb{C}^2)^n$ is there $|\varphi\rangle \in \Phi$ such that

$$|\varphi\rangle \xrightarrow{\text{LOCC}} \{p_k, |\psi_k\rangle\}_k ; \quad \sum_{j: D(\psi, \psi_j) \leq \epsilon} p_j \geq 1 - \delta$$

polynomial resources

$$|\varphi\rangle \in (\mathbb{C}^2)^{O(\text{poly}(n, 1/\delta, 1/\epsilon))}$$

polynomial time

$$O(\text{poly}(n, 1/\delta, 1/\epsilon)) \\ \text{LOCC steps}$$



In the circuit model: $[O(\text{poly}(n \log 1/\epsilon))]$

*Problem: want to treat δ and ϵ in the same way
no Solovay-Kitaev theorem for δ*

Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

- 1) The family must have enough entanglement

Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

- 1) The family must have enough entanglement
- 2) The entanglement must grow “quickly enough” with the size of the states

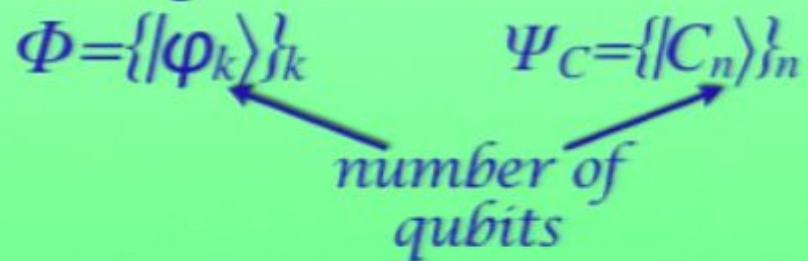
Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

- 1) The family must have enough entanglement
- 2) The entanglement must grow “quickly enough” with the size of the states

$$\Phi = \{|\varphi_k\rangle\}_{k} \quad \Psi_C = \{|\mathcal{C}_n\rangle\}_n$$

number of qubits



Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

- 1) The family must have enough entanglement
- 2) The entanglement must grow "quickly enough" with the size of the states

$$\Phi = \{|\varphi_k\rangle\}_k$$

$$\Psi_C = \{|C_n\rangle\}_n$$

number of
qubits

$$\mathcal{E}(|C_n\rangle) = \mathcal{E}(|\varphi_{k=O(\text{poly}(n))}\rangle)$$

Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

- 1) The family must have enough entanglement
- 2) The entanglement must grow "quickly enough" with the size of the states

$$\Phi = \{|\varphi_k\rangle\}_{k=1}^k$$

$$\Psi_C = \{|C_n\rangle\}_{n=1}^n$$

number of
qubits

$$\mathcal{E}(|C_n\rangle) = \mathcal{E}(|\varphi_{k=O(\text{poly}(n))}\rangle)$$

2') In the approximate/stochastic case:

Similar, but \mathcal{E}_ε is enough: $\mathcal{E}_\varepsilon(|C_n\rangle) = \mathcal{E}(|\varphi_{k=O(\text{poly}(n))}\rangle)$

Entanglement and efficiency

The efficiency requirement gives us stronger criteria:

- 1) The family must have enough entanglement
- 2) The entanglement must grow "quickly enough" with the size of the states

If Φ is not approximate efficient universal (with $\delta=0$) then $\delta>0$ doesn't help

$\Psi_C = \{ |C_n\rangle \}_n$
of

$$\mathcal{E}(|C_n\rangle) = \mathcal{E}(|\varphi_{k=O(\text{poly}(n))}\rangle)$$

approximate/stochastic case:

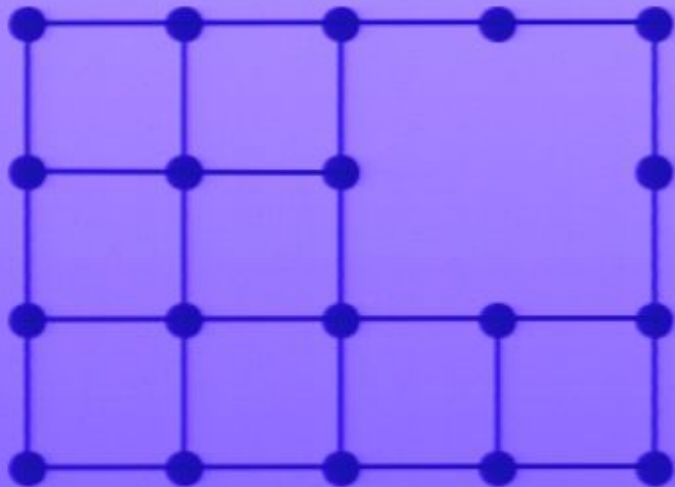
is enough: $\mathcal{E}_\varepsilon(|C_n\rangle) = \mathcal{E}(|\varphi_{k=O(\text{poly}(n))}\rangle)$

$$E(\Phi) \geq (1-\delta)E_{\varepsilon}^*$$



Good news (something works I) (Cluster state with holes)

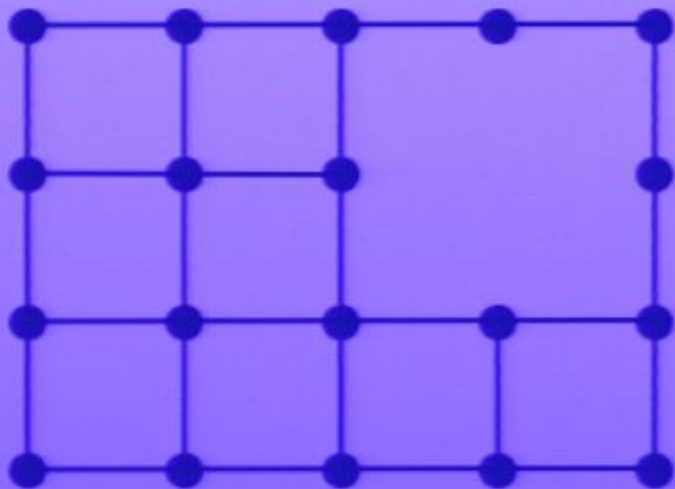
Suppose we want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



Problem: not every site got its qubit

Good news (something works I) (Cluster state with holes)

Suppose we want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



Problem: not every site got its qubit

Good news (something works I) (Cluster state with holes)

Suppose we want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$

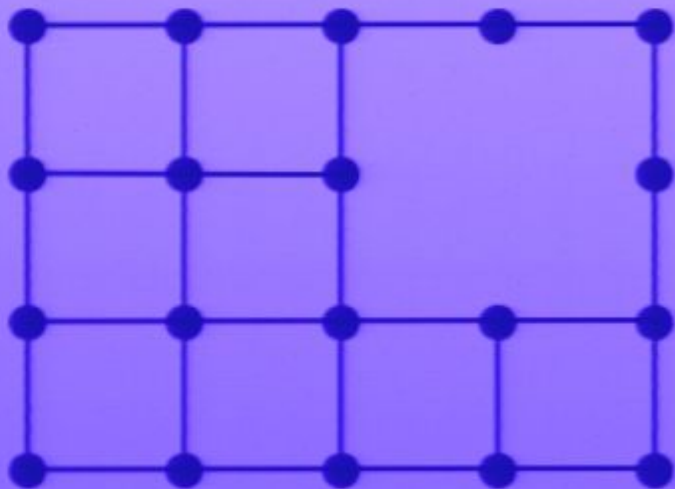


Problem: not every site got its qubit

Assume: we know where the holes are

Good news (something works I) (Cluster state with holes)

Suppose we want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



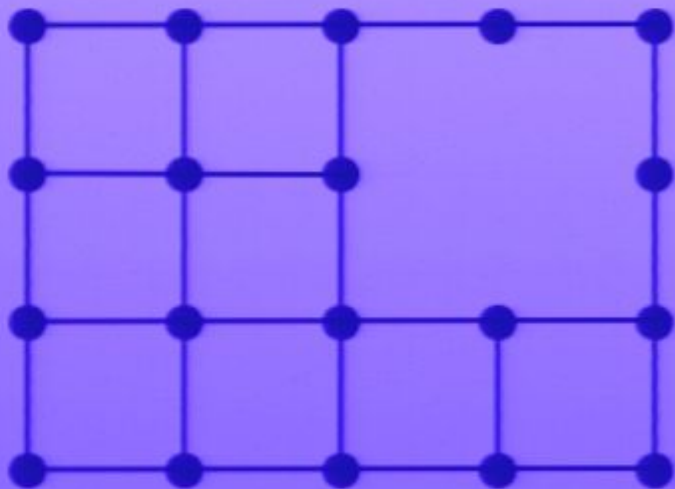
Problem: not every site got its qubit

Assume: we know where the holes are

Exact deterministic computation is not possible

Good news (something works I) (Cluster state with holes)

Suppose we want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



Problem: not every site got its qubit

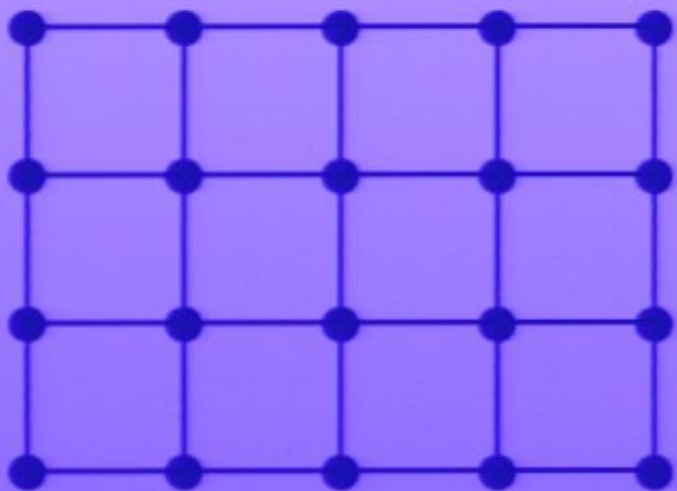
Assume: we know where the holes are

Exact deterministic computation is not possible

Still useful! Universal resource for exact δ -stochastic computation *for all $\delta > 0$*

Good news (something works II) (Stability of universal resources)

Suppose we again want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



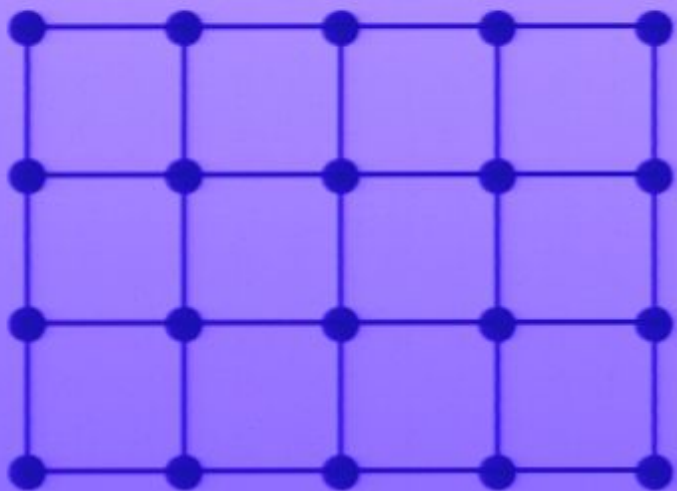
Assume $\mathcal{D}(\sigma_n, C_n) \leq \eta$

The protocol does not increase the damage too much
(LOCC operations don't increase \mathcal{D})

\Rightarrow Most branches will still end up "close enough" to
the desired output

Good news (something works II) (Stability of universal resources)

Suppose we again want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



But we are in the real world!

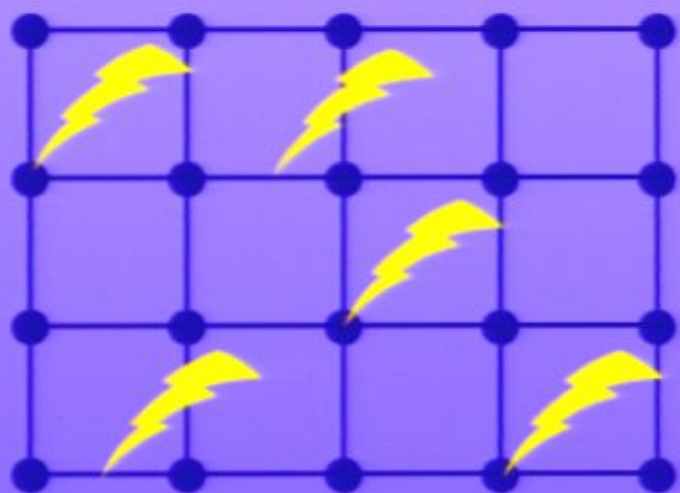
Assume $\mathcal{D}(\sigma_n, C_n) \leq \eta$

The protocol does not increase the damage too much
(LOCC operations don't increase \mathcal{D})

\Rightarrow Most branches will still end up "close enough" to
the desired output

Good news (something works II) (Stability of universal resources)

Suppose we again want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



But we are in the real world!

Errors!!!!!!

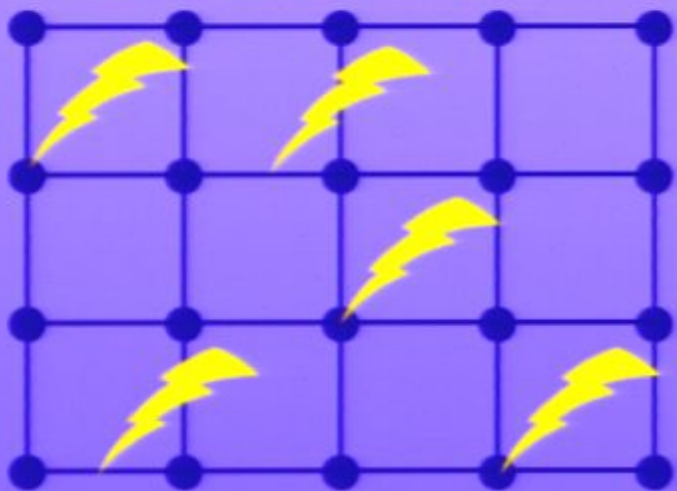
Assume $\mathcal{D}(\sigma_n, C_n) \leq \eta$

The protocol does not increase the damage too much
(LOCC operations don't increase \mathcal{D})

\Rightarrow Most branches will still end up "close enough" to
the desired output

Good news (something works II) (Stability of universal resources)

Suppose we again want to use cluster states: $\Psi_C = \{|C_n\rangle\}_n$



But we are in the real world!

Errors!!!!!!



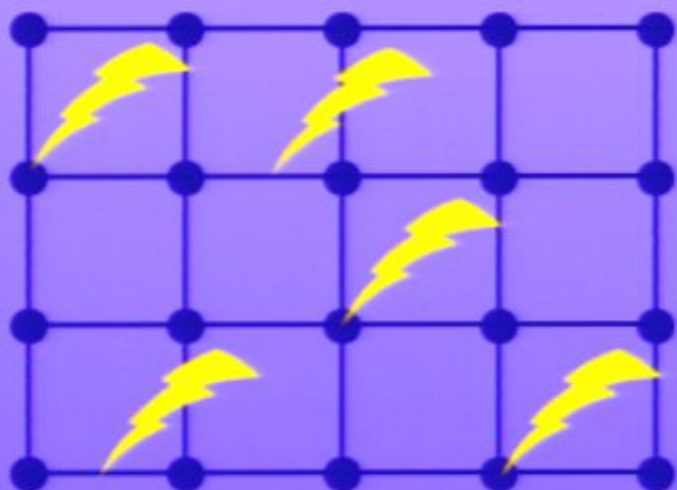
We actually have $\Sigma_C = \{|\sigma_n\rangle\}_n$

Assume $\mathcal{D}(\sigma_n, C_n) \leq \eta$

The protocol does not increase the damage too much
(LOCC operations don't increase \mathcal{D})

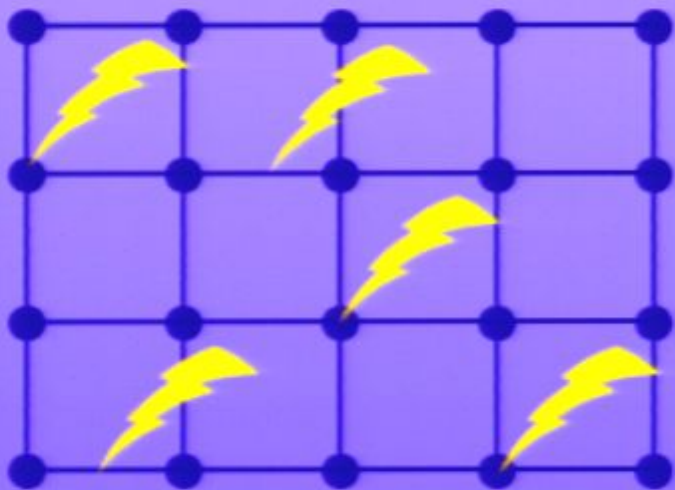
⇒ Most branches will still end up “close enough” to
the desired output

Good news (something works IIB)
(Stability of universal resources)



*Cluster not so special: argument
holds in general*

Good news (something works IIb) (Stability of universal resources)

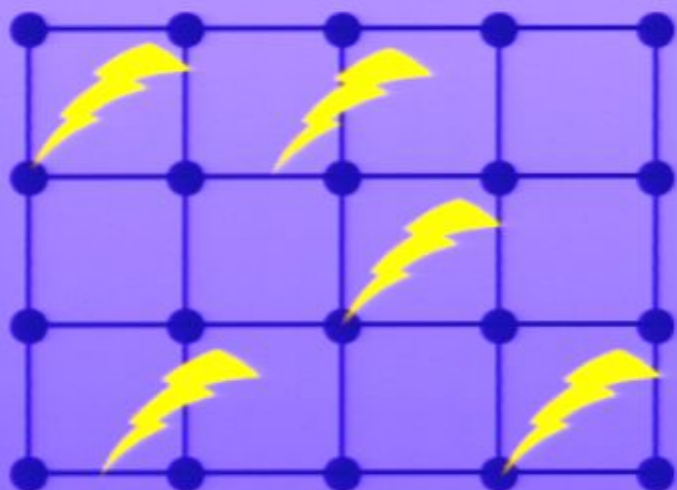


Cluster not so special: argument holds in general

$$\Phi = \{|\varphi_k\rangle\}_k$$

↑
 ε -approximate
 δ -stochastic

Good news (something works IIB) (Stability of universal resources)

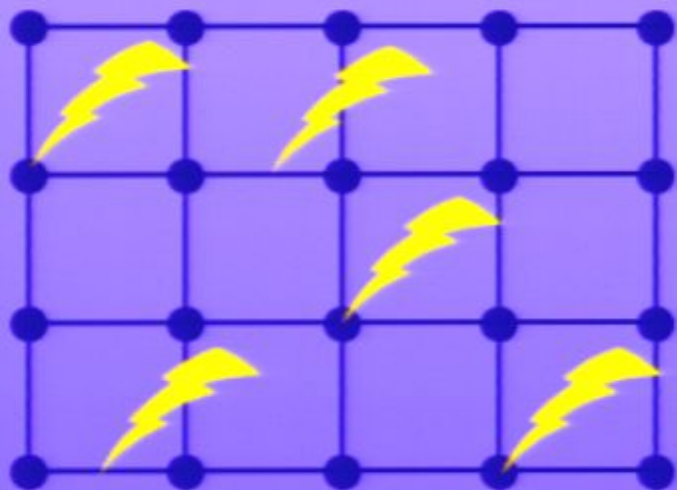


Cluster not so special: argument holds in general

$$\Phi = \{|\varphi_k\rangle\}_{k \rightarrow} \begin{cases} \Sigma = \{\sigma_k\}_{k \rightarrow} \\ \mathcal{D}(\sigma_k, \varphi_k) \leq \eta \end{cases}$$

↑
 ε -approximate
 δ -stochastic

Good news (something works IIB) (Stability of universal resources)



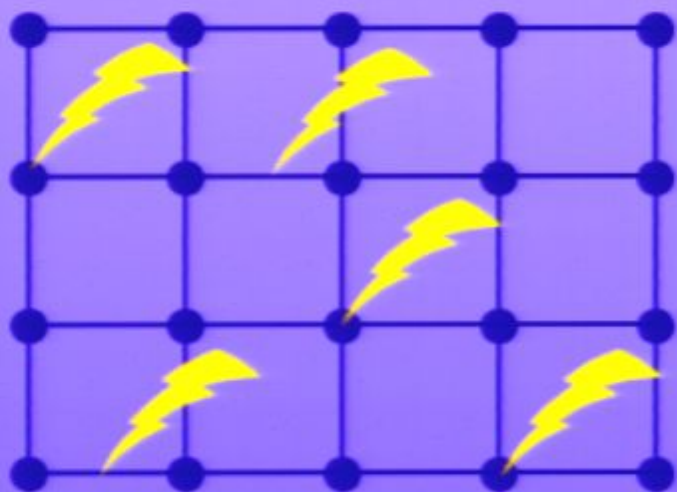
Cluster not so special: argument holds in general

$$\Phi = \{|\varphi_k\rangle\}_{k \rightarrow} \begin{cases} \Sigma = \{\sigma_k\}_{k \rightarrow} \\ \mathcal{D}(\sigma_k, \varphi_k) \leq \eta \end{cases}$$

↑
 ε -approximate
 δ -stochastic

$\Rightarrow \Sigma$ is still approximate stochastic universal

Good news (something works IIb) (Stability of universal resources)



Cluster not so special: argument holds in general

$$\Phi = \{|\varphi_k\rangle\}_{k \rightarrow} \begin{cases} \Sigma = \{\sigma_k\}_{k \rightarrow} \\ \mathcal{D}(\sigma_k, \varphi_k) \leq \eta \end{cases}$$

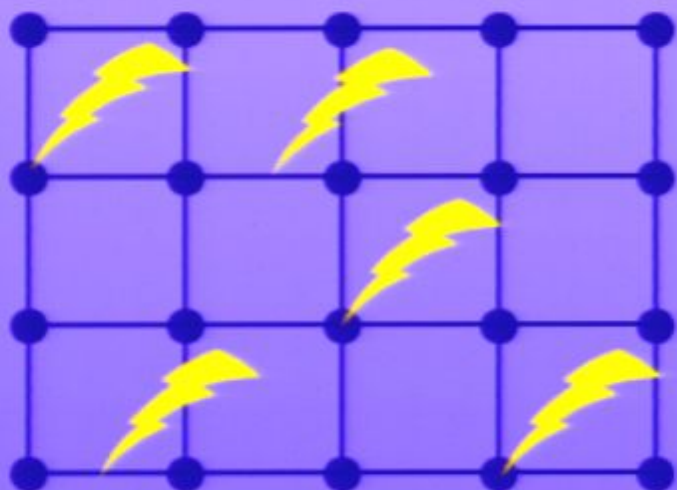
↑
 ε -approximate
 δ -stochastic

$\Rightarrow \Sigma$ is still approximate stochastic universal

New parameters δ', ε' such that $\delta'\varepsilon' \geq \delta + \varepsilon + \eta$

Good news (something works IIB)

(Stability of universal resources)



Cluster not so special: argument holds in general

$$\Phi = \{ |\varphi_k\rangle \}_{k} \xrightarrow{\text{lightning bolt}} \begin{cases} \Sigma = \{ \sigma_k \}_{k} \\ \mathcal{D}(\sigma) \end{cases}$$

↑
 ε -approximate
 δ -stochastic

Universal resources are stable under perturbations

$\Rightarrow \Sigma$ is still approximate stochastic universal

New parameters δ', ε' such that $\delta'\varepsilon' \geq \delta + \varepsilon + \eta$

Conclusions

- ◎ *Defined universality for quantum computation*
 - ▶ *Universal exact “state preparators”*
 - ▶ *Approximate and stochastic universality*
 - ▶ *Efficiency and universality*
- ◎ *Constructed entanglement-based criteria for URs*
 - ▶ *Criteria in the exact and deterministic case*
 - ▶ *ε -measures and generalized criteria*
 - ▶ *Examples of applications*
- ◎ *“Go” results for universal families*
 - ▶ *Deformed cluster*
 - ▶ *Approximate resources*

Exact case

*M. Van den Nest, A. Miyake,
W. Duer, H.-J. Briegel,
New J. Phys 9, 204 (2007)*

Approximate/stochastic

*C.M., M. Piani, and all above
On the arXiv (sooner or later)*

Exact case

*M. Van den Nest, A. Miyake,
W. Duer, H.-J. Briegel,*

New J. Phys 9, 204 (2007)

Approximate/stochastic

*C.M., M. Piani, and all above
On the arXiv (sooner or later)*

Thank you

Welcome barbecue this afternoon!



In the new temporary IOC building

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal

VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal
VGA-1

No Signal
VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal

VGA-1

No Signal
VGA-1

No Signal

VGA-1

No Signal
VGA-1





Support Center



MyDVD LE



Try WordPerfect



Owner's Manual



Windows Media Player



setupfiles



misc



08



fglrpc



Octave



My Documents



Burn CDs & DVDs with...



Recycle Bin

Options

- Introduction
- Graph Quantum Co
- Finite Geometry Qu
- Subsystem LDPC C

Two Approaches to Sparse Graph Quantum Codes

Pradeep Kiran Sarvepalli

Department of Computer Science
Texas A&M University

Joint work with Andreas Klappenecker and Martin Rötteler

Quantum Information and Graph Theory: Emerging Connections
Perimeter Institute for Theoretical Physics
May 2, 2008

Two Approaches to Sparse Graph Quantum Codes

Pradeep Kiran Sarvepalli

Department of Computer Science
Texas A&M University

Joint work with Andreas Klappenecker and Martin Rötteler

Quantum Information and Graph Theory: Emerging Connections
Perimeter Institute for Theoretical Physics
May 2, 2008