

Title: The Exact Renormalization Group - Lecture 1: Wilsonian Renormalization

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Abstract: In this lecture, I will discuss Wilson's picture of renormalization and its relation to the Exact Renormalization Group (ERG). In particular, I will focus on how one can understand, in a physically intuitive way, what it is for a quantum field theory to be nonperturbatively renormalizable.

The Exact Renormalization Group: Introduction & Applications

Lecture 1: Wilsonian Renormalization

Oliver J. Rosten

Dublin Institute for Advanced Studies

April / May 2008

What is the Exact Renormalization Group?

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Physics at Different Scales

- The way we describe physics changes with scale
- The relevant degrees of freedom change with scale
- A few microscopic laws give rich macroscopic behaviour

From Microscopic to Macroscopic

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- **Lecture 1: Wilsonian Renormalization**
- Lecture 2: Exact Renormalization Group Equations
- Lecture 3: The Derivative Expansion
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Useful Literature

- K. Wilson and J. Kogut, "The Renormalization group and the epsilon expansion," *Phys. Rept.* **12** (1974) 75.
- F. Wegner , "The Critical State, General Aspects," in: C. Domb and M. S. Green (Eds.), *Phase Transitions and Critical Phenomena*, Vol VI, Acad. Press, N.-Y., 1976 p. 7.
- M. E. Fisher, "Renormalization group theory: Its basis and formulation in statistical physics," *Rev. Mod. Phys.* **70** (1998) 653.
- J. Cardy, "Scaling and Renormalization in Statistical Physics," CUP, 1996.

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Continued...

- T. R. Morris, “Elements of the continuous renormalization group,” *Prog. Theor. Phys.* **131** (1998) 395, [hep-th/9802039](#).
- C. Bagnuls and C. Bervillier, “Exact renormalization group equations: An introductory review,” *Phys. Rept.* **348** (2001) 91, [hep-th/0002034](#).
- J. Berges, N. Tetradis, and C. Wetterich, “Non-perturbative renormalization flow in quantum field theory and statistical physics,” *Phys. Rept.* **363** (2002) 223, [hep-ph/0005122](#).

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- 1 Broad Uses of the ERG
- 2 Qualitative Aspects of the ERG
 - The Basic Ideas
 - Relevance, Irrelevance and all that
- 3 Renormalizability
 - Continuum Limits
 - Examples
 - Sources of Misunderstanding
- 4 Recap

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The ERG in Statistical Physics

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Example: The Ferromagnet



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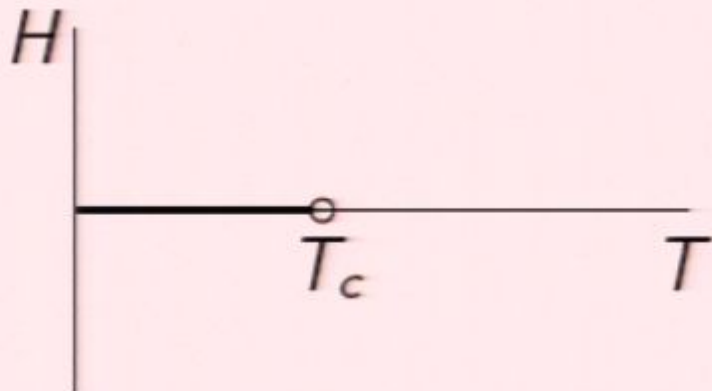
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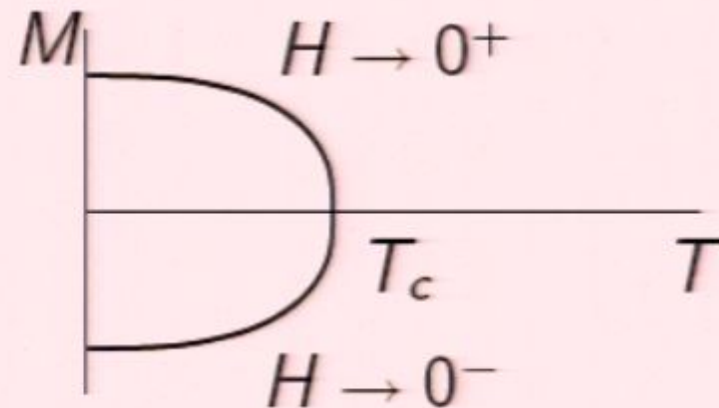
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Phase Diagram



M vs T

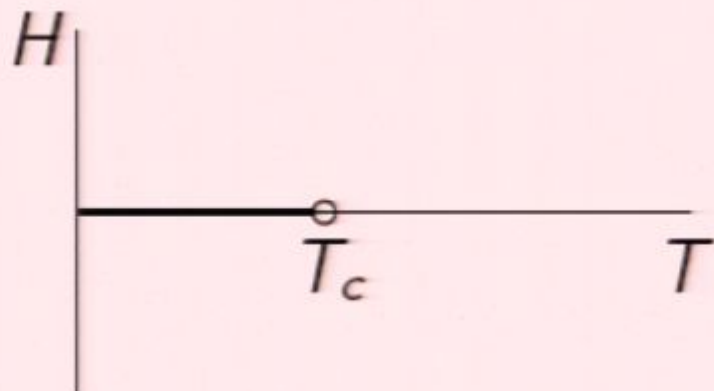


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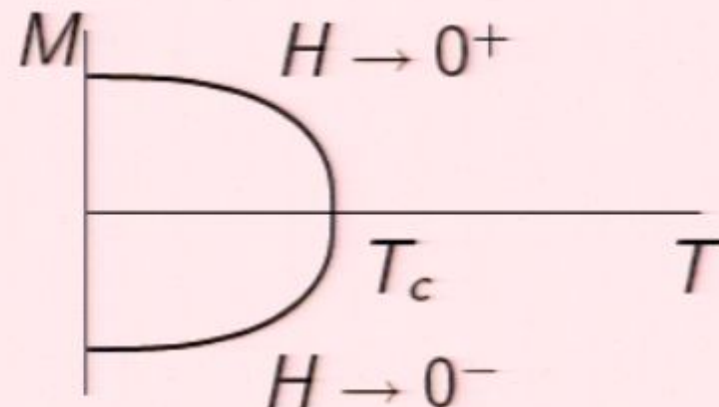
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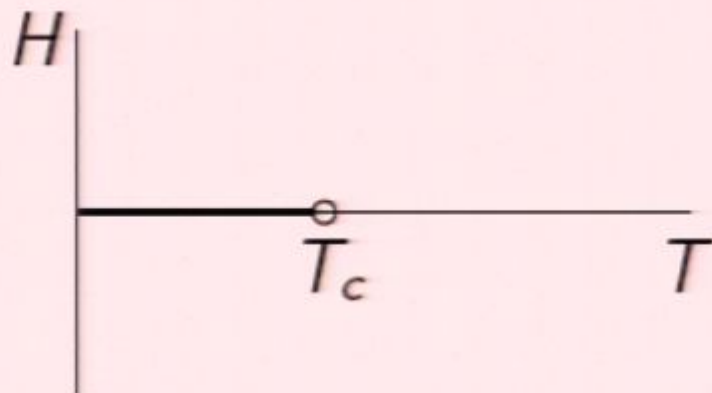
Understanding Critical Phenomena

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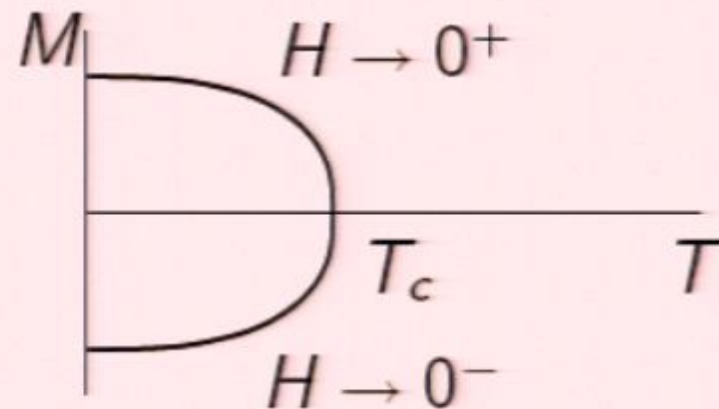
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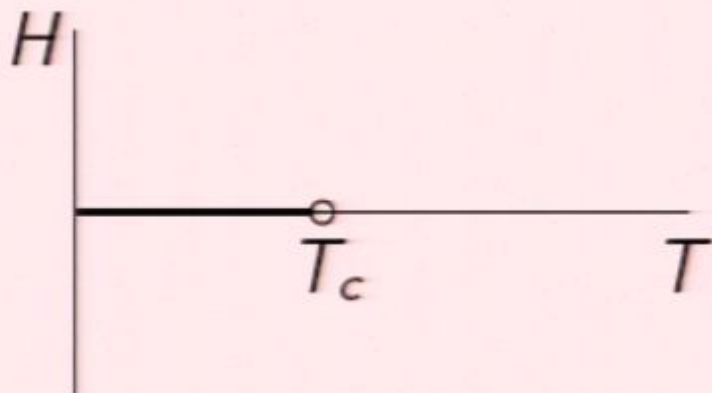
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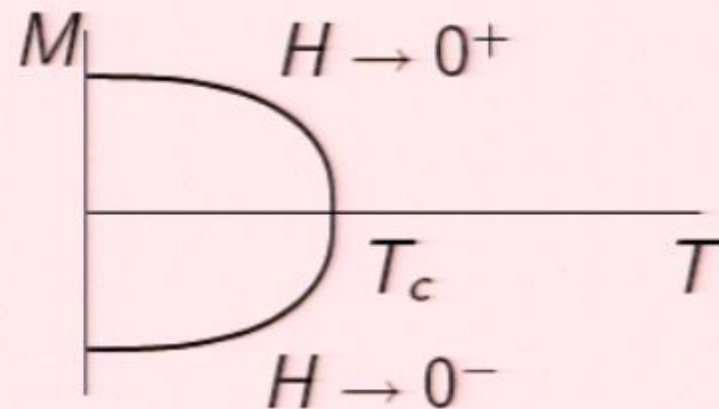
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Understanding Critical Phenomena

- Computation of critical exponents eg $M \sim (T_c - T)^\beta$
- Universality: why do apparently very different systems have the same critical exponents?

The ERG in QFT

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Qualitative Aspects

- Understanding the physical basis of renormalization
 - Asymptotic freedom
 - Asymptotic safety
 - Triviality
- Why perturbative renormalizability does not necessarily imply nonperturbative renormalizability

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- Calculations can be performed directly in terms of renormalized variables
- The ERG admits nonperturbative approximation schemes which preserve renormalizability

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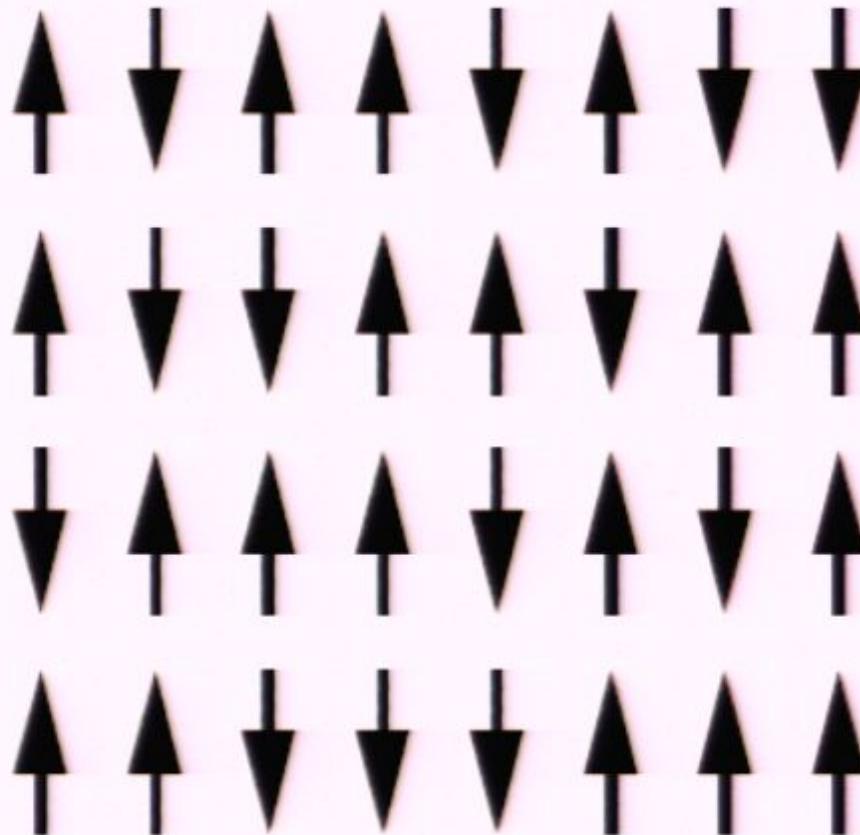
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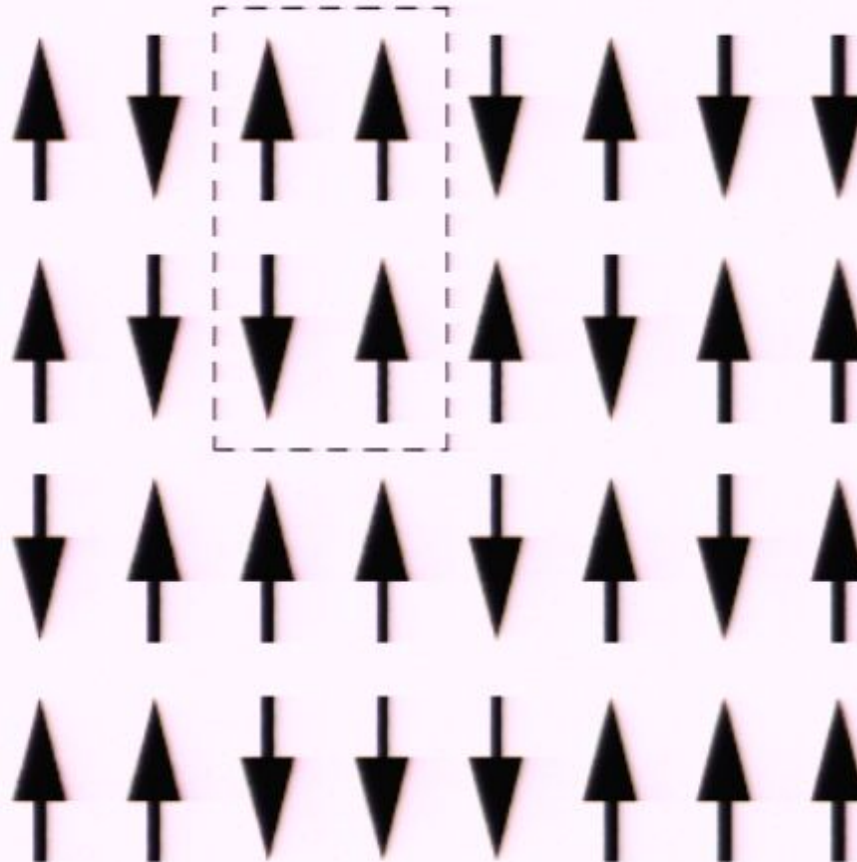
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- To go from micro to macro, **average** over groups of spins
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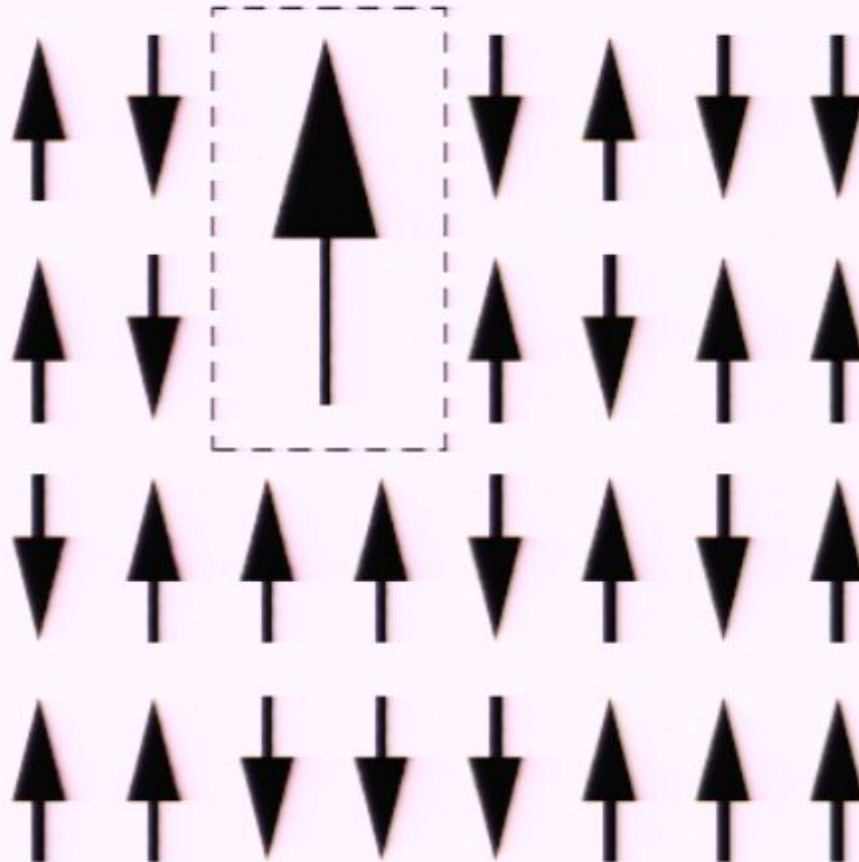
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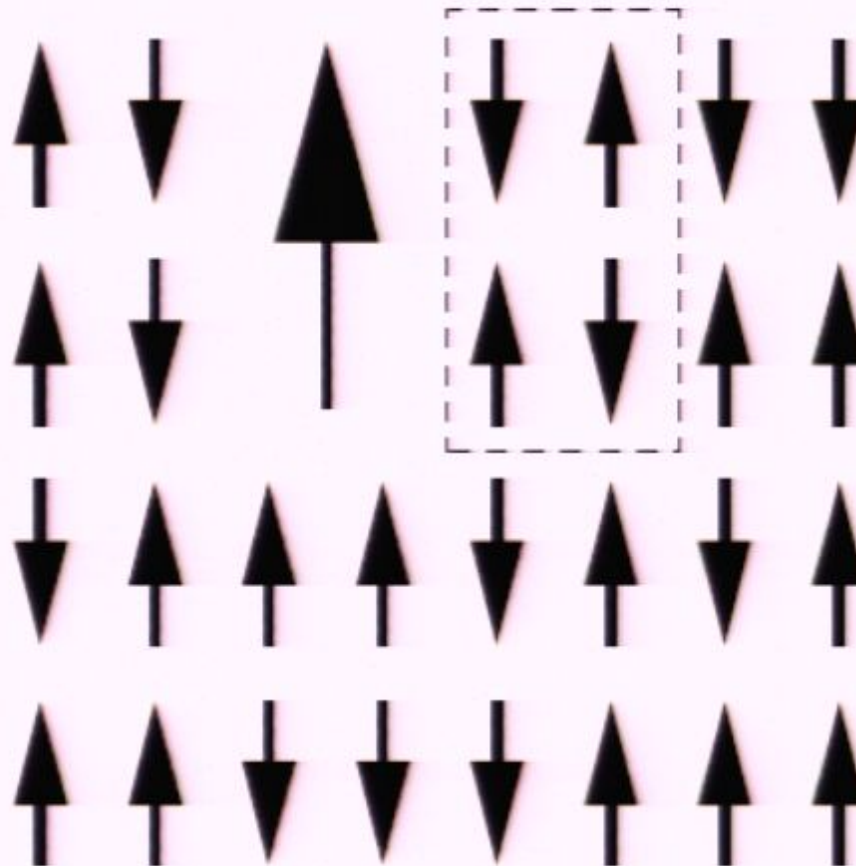
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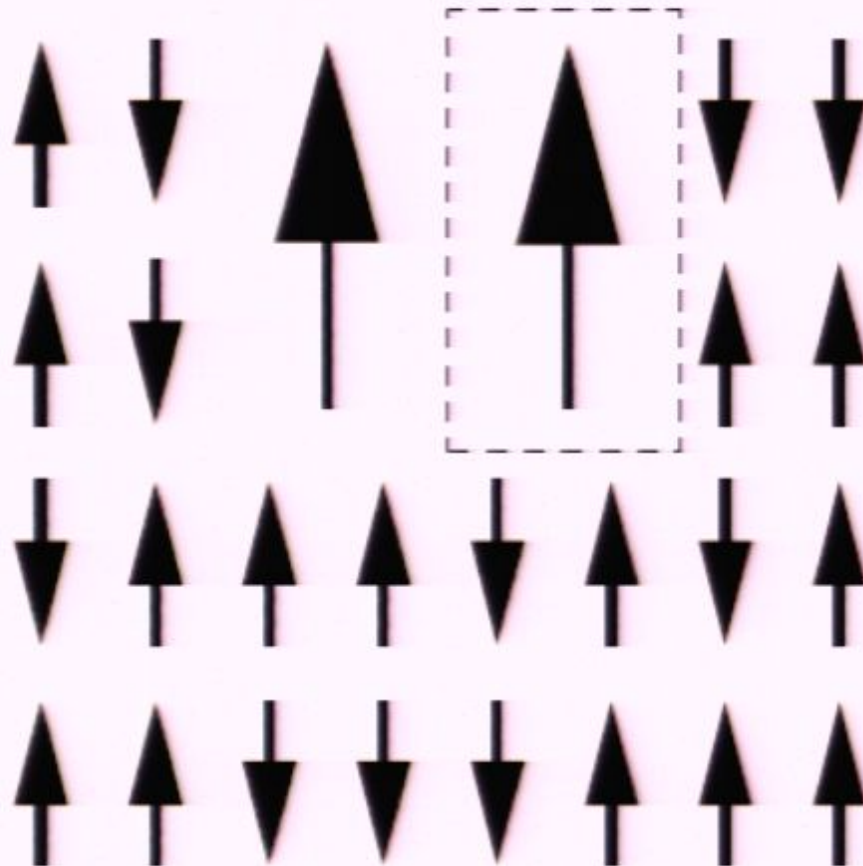
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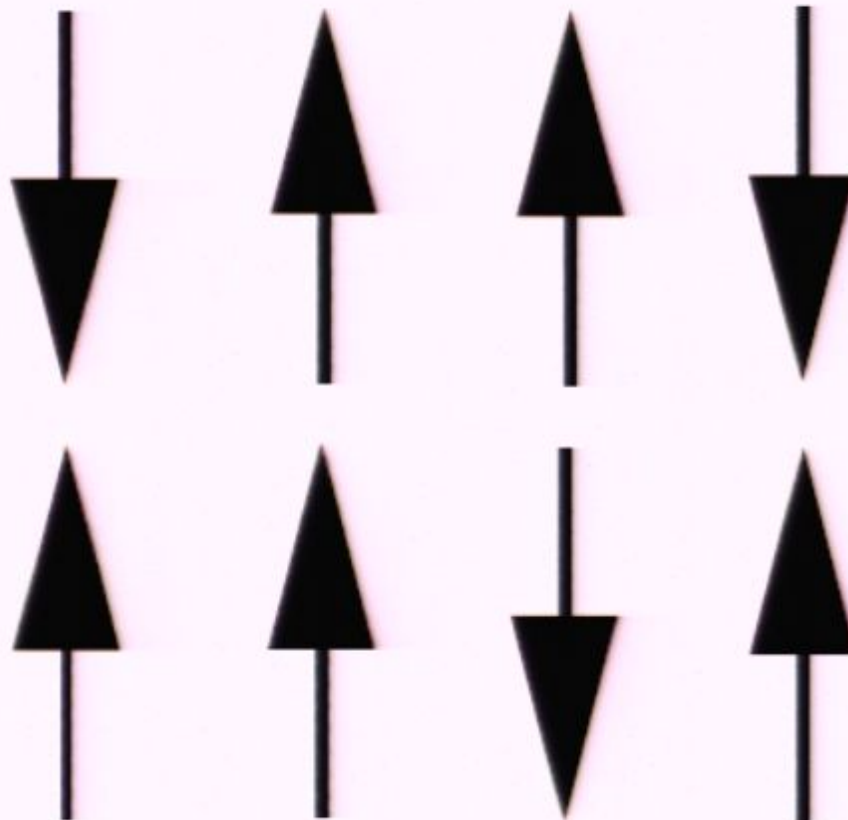
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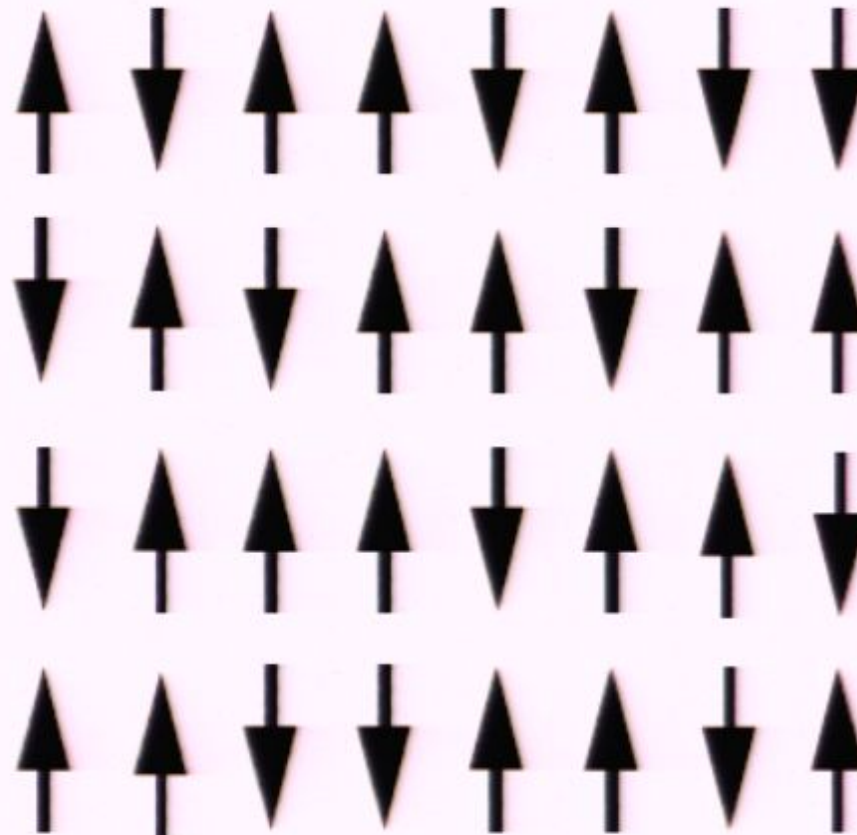
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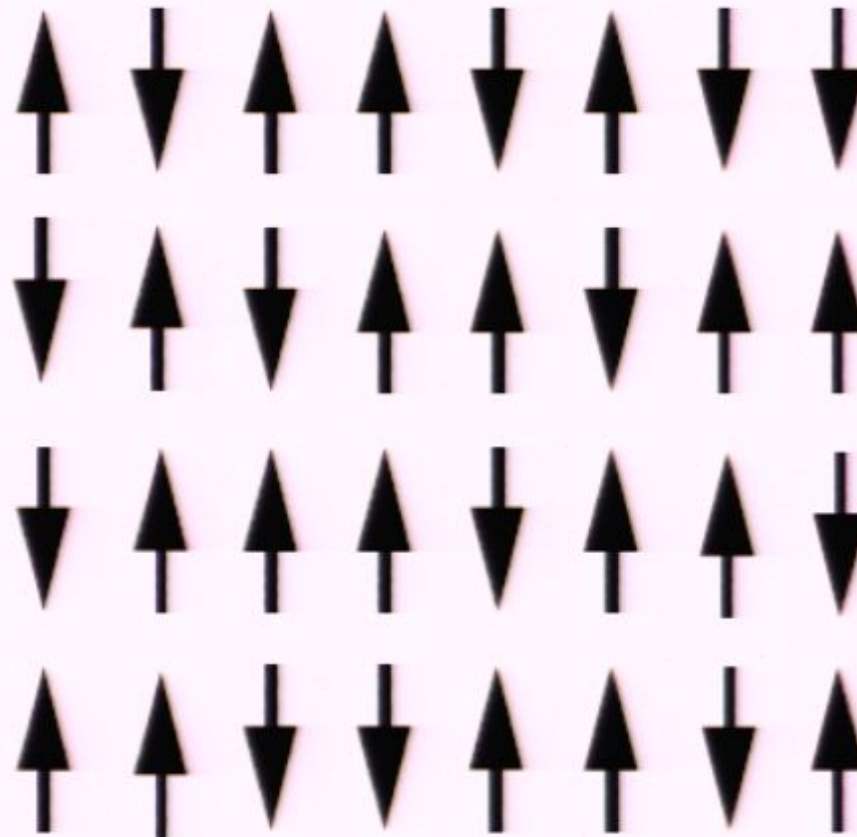
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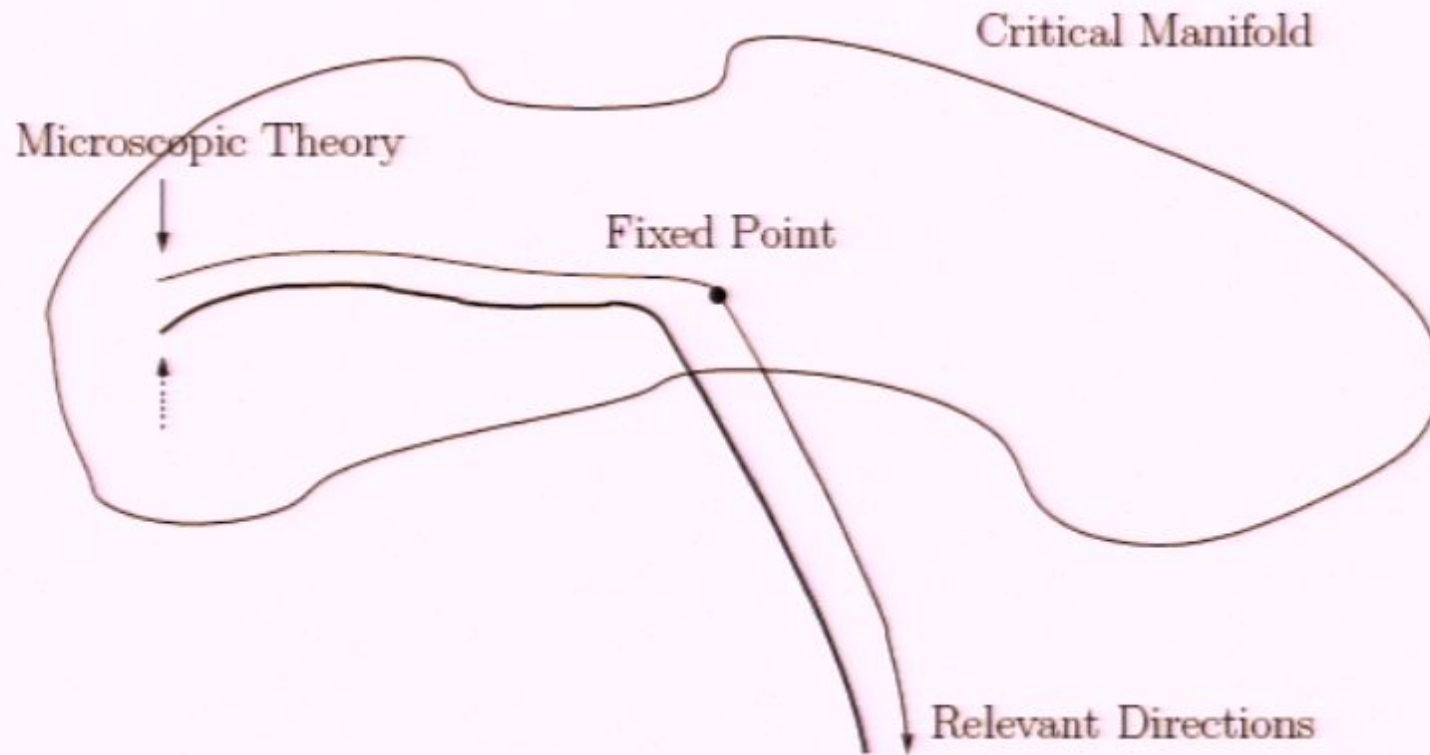
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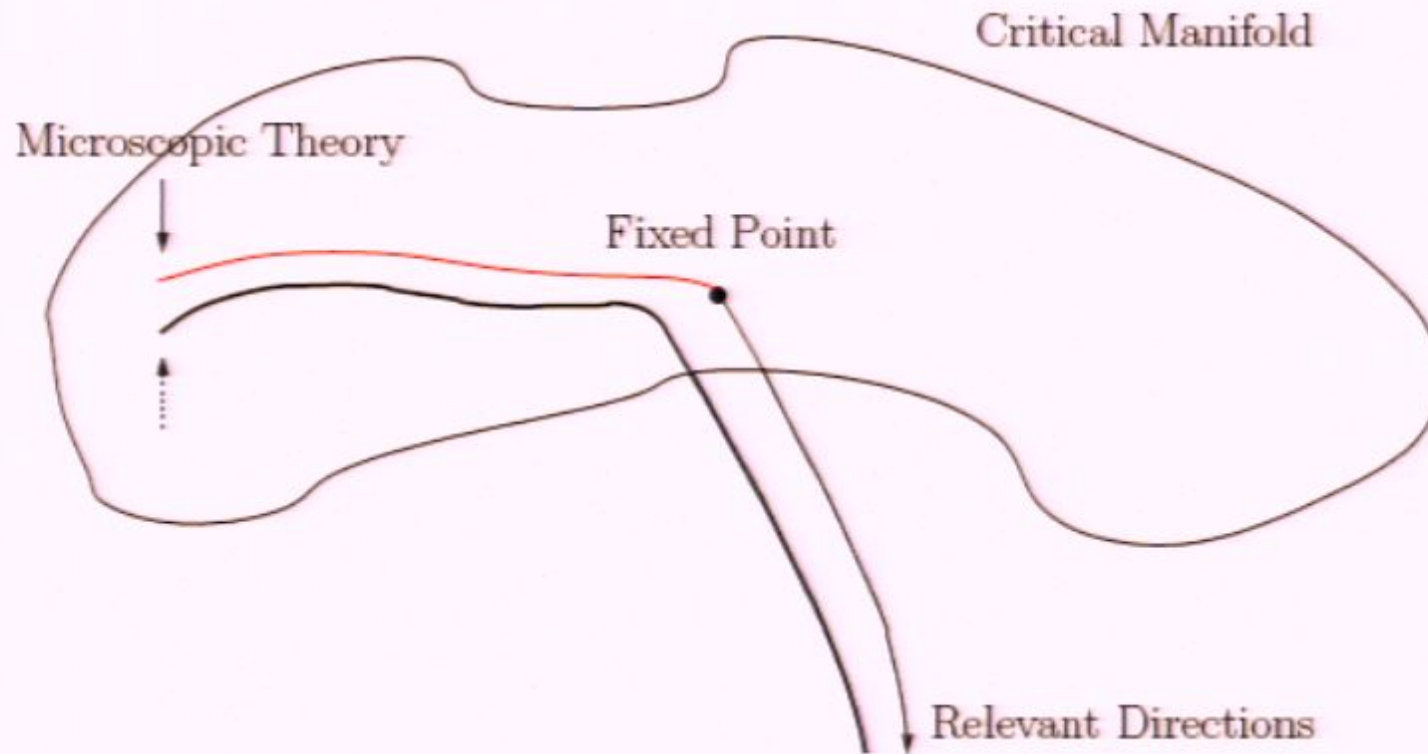
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- The transformation can have fixed points

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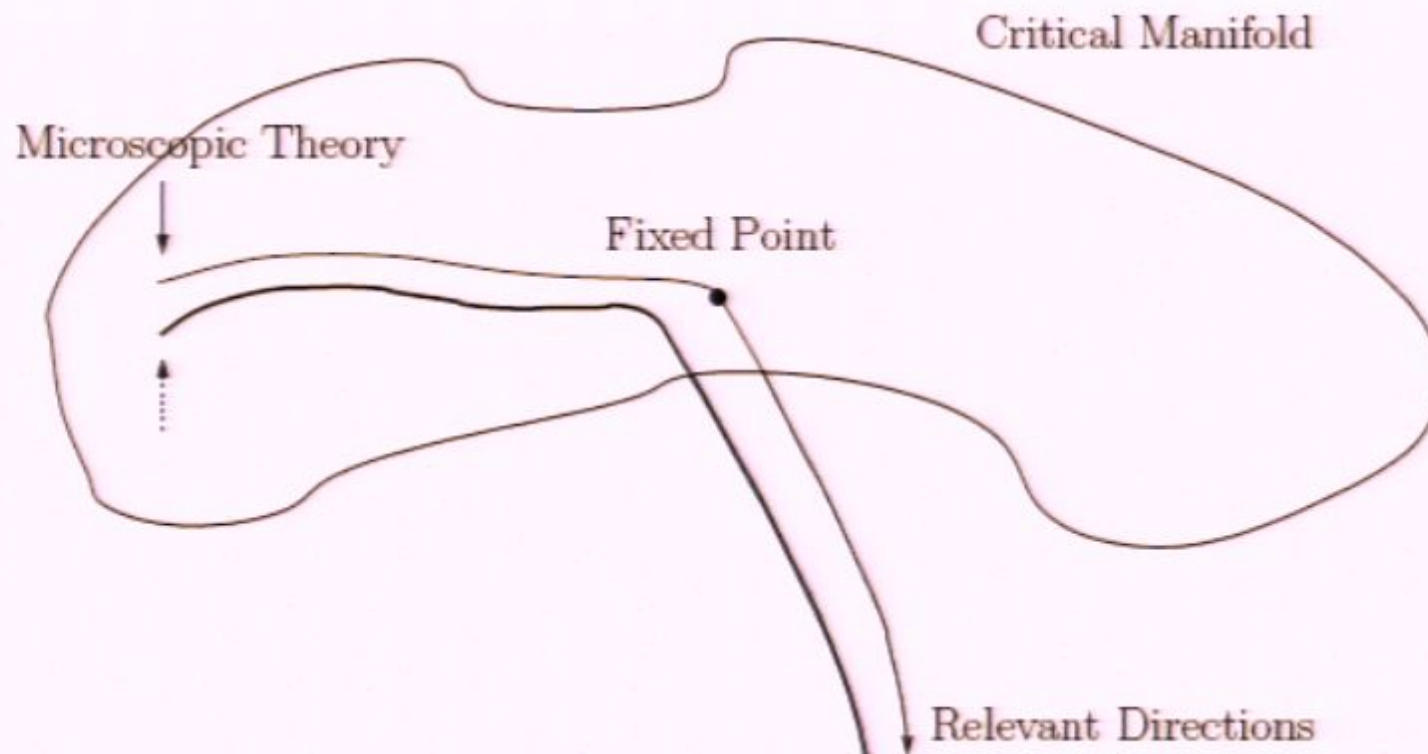


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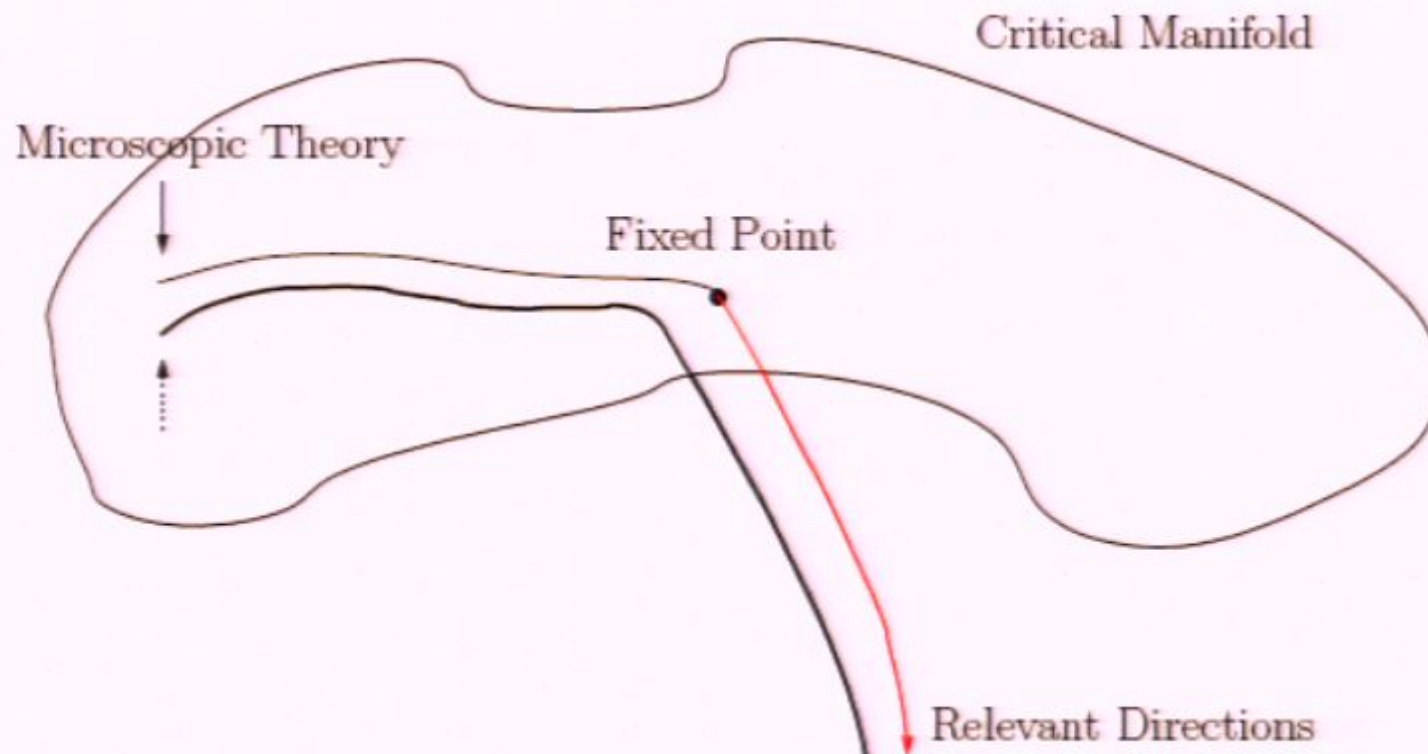
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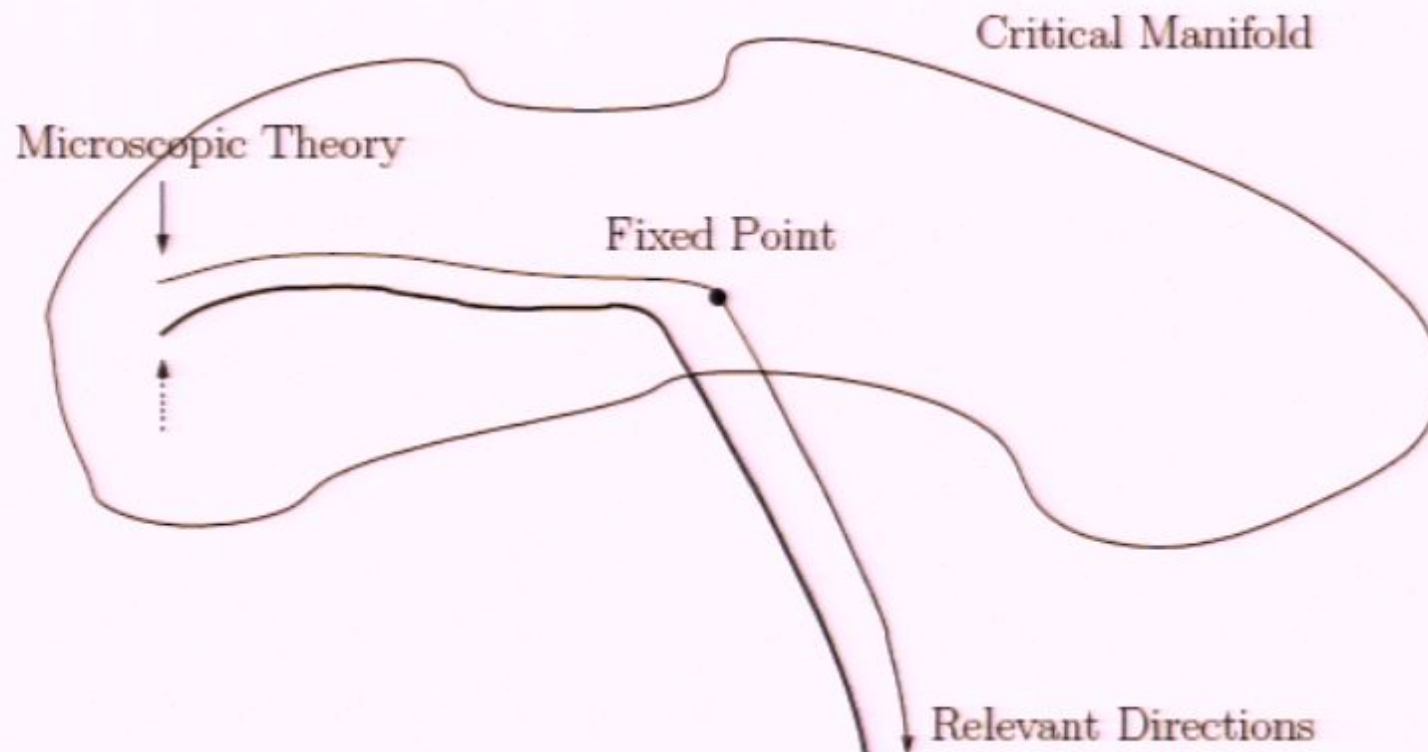
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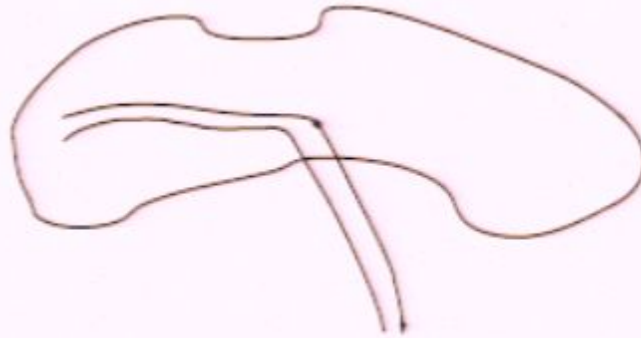
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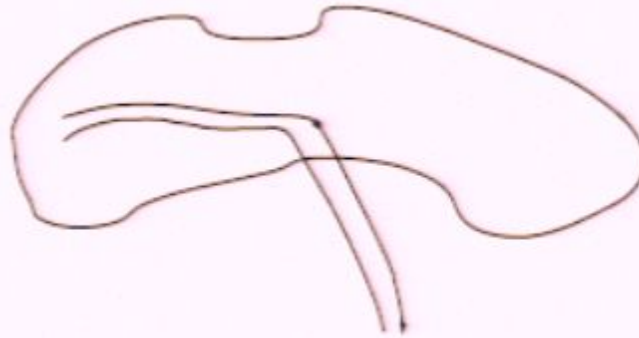


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- Flows along the relevant directions leave the critical surface
- If there are n relevant directions, then we must tune n quantities to get on to the critical surface

Universality

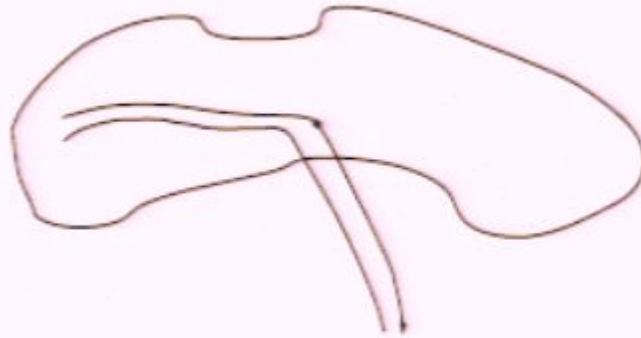


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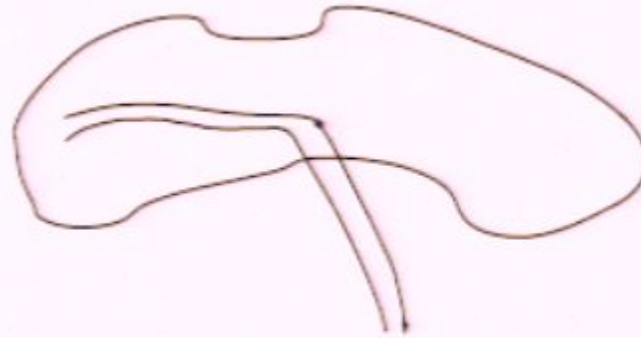
- All theories on the critical surface flow towards the fixed point
- Therefore the IR dynamics of these theories is the same
- The theories are in the same universality class
- To determine the critical exponents, move slightly away from the critical surface of $M \sim (T_c - T)^\beta$
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- The critical exponents are obtained by linearizing the flow around the fixed point
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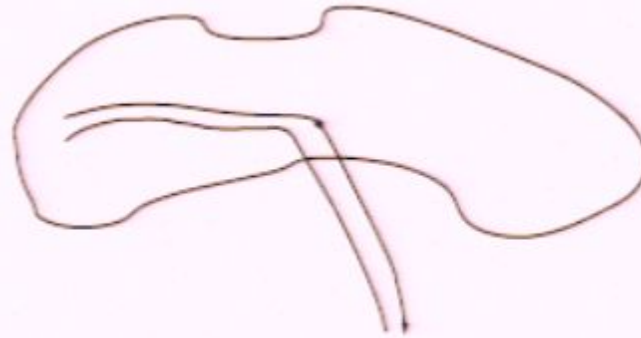
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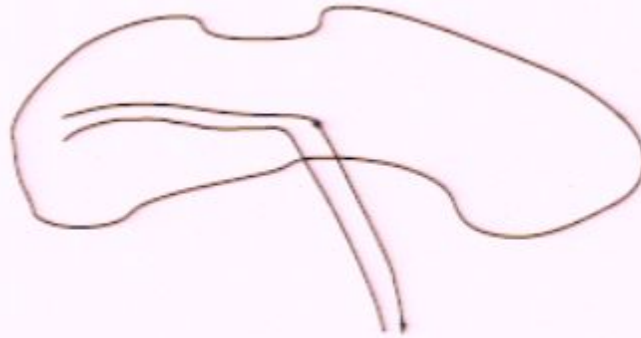
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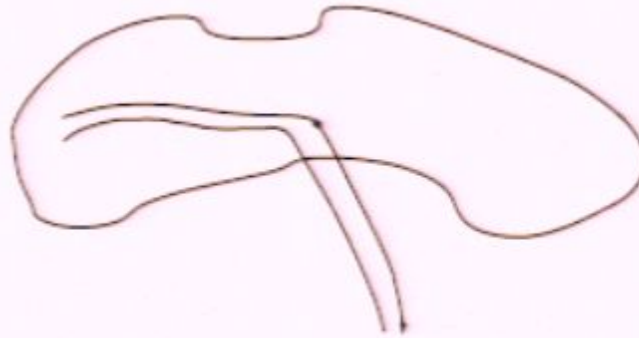
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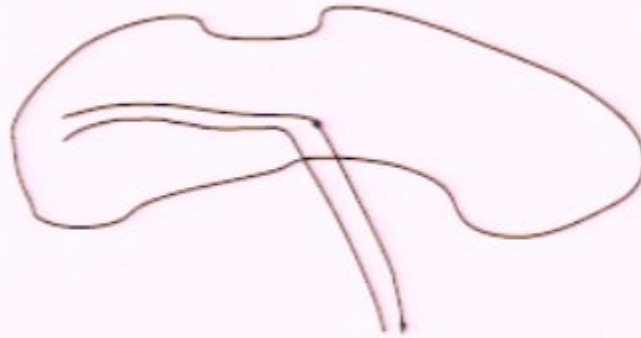
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$$Z = \int_{\Lambda_0} \mathcal{D}\Phi e^{-S_{\Lambda_0}[\Phi]} = \int_{\Lambda} \mathcal{D}\Phi e^{-S_{\Lambda}[\Phi]}$$

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- The bare (classical) action
- Integrate out modes between the bare scale and an intermediate scale, Λ

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The Wilsonian Effective Action

Start with the partition function

$$Z = \int_{\Lambda_0} \mathcal{D}\Phi e^{-S_{\Lambda_0}[\Phi]} = \int_{\Lambda} \mathcal{D}\Phi e^{-S_{\Lambda}[\Phi]}$$

- The bare scale
 - High energy (short distance) scale
 - Modes above this scale are cut off (regularized)
- The bare (classical) action
- Integrate out modes between the bare scale and an intermediate scale, Λ
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Very General ERGs

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Formulation

$$-\Lambda \partial_\Lambda e^{-S[\varphi]} = \int_x \frac{\delta}{\delta \varphi(x)} \left(\Psi_x[\varphi] e^{-S[\varphi]} \right)$$

- effective scale
- set of fields
- Wilsonian effective action
- partition function, $\int \mathcal{D}\varphi e^{-S[\varphi]}$, invariant under the flow
- defines our ERG

• parametrizes blocking procedure

• huge freedom in precise form—adjust to suit our needs

Flow Equation

$$-\Lambda \partial_\Lambda S = \int_x \frac{\delta S}{\delta \varphi(x)} \Psi_x - \int_x \frac{\delta \Psi_x}{\delta \varphi(x)}$$

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Ingredients of ERG Transformation

- Blocking (coarse-graining)
- Rescaling

Implementing Rescaling

What we need for this talk

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- Remember to take account of anomalous dimensions!
- i.e. $X \rightarrow X\Lambda^{\text{full scaling dimension}}$
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Linearized ERG Theory

Linearized ERG Theory

- At a fixed point we have $\partial_t S_* = 0$
- Near a fixed point, separate the ERG operator into linear and quadratic parts:

$$\partial_t(S_* + \Delta S) = \mathcal{L}\Delta S + \mathcal{Q}\Delta S$$

- Separate variables in the linearized equation:

$$\Delta S[\varphi] = \sum_i \alpha_i e^{\lambda_i t} \mathcal{O}_i[\varphi]$$

- Quantization of the λ_i will be discussed in Lecture 3
- Finally, $S_t = S_* + \sum_i \alpha_i e^{\lambda_i t} \mathcal{O}_i$

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Physical Eigenvalues: Scaling Operators

- $\lambda_i > 0$: \mathcal{O}_i is relevant
- $\lambda_i < 0$: \mathcal{O}_i is irrelevant
- $\lambda_i = 0$: \mathcal{O}_i is marginal;
 - $S_* = \mathcal{O}_i = 0$ is a fixed point
 - This might not be true beyond leading order
 - Eg the four-point coupling in $D=4$ scalar field theory is marginally irrelevant

Unphysical Eigenvalues: Redundant Operators

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$$S_t = S_* + \sum_i \alpha_i e^{\lambda_i t} \mathcal{O}_i$$

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- Physics is independent of these eigenvalues and their operators
- More in Lecture 3...

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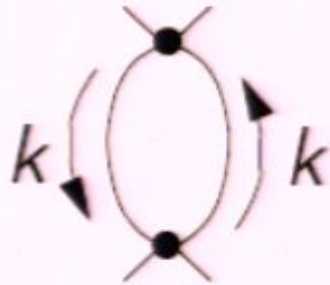
UV Divergences in QFT

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UV Divergences in QFT

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A Feynman diagram showing a loop with two external lines and two vertices. The loop is represented by two curved lines forming a circle. Two external lines, each with a black dot at its vertex, enter and exit the loop. The momentum of the loop is labeled k on both sides with arrows indicating the direction of flow.

$$\sim \int \frac{d^4 k}{k^4} \sim \ln \Lambda_0$$

UV Divergences in QFT

- Loop diagrams in quantum field theory yield UV divergences



A Feynman diagram showing a loop with two external lines and two vertices. The loop is formed by two curved lines, each with an arrow pointing clockwise. The two external lines are straight lines meeting the loop at two vertices, each marked with a black dot. The momentum of the loop is labeled as k on both sides.

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- If all divergences can be absorbed into a finite number of couplings, the theory is renormalizable
- Perturbative renormalizability by no means guarantees nonperturbative renormalizability!
- With its built in UV cutoff, the ERG is a natural tool to study renormalizability**

Continuum Limits I

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The Question

Are there effective actions $S_{\Lambda, \Lambda_0}[\varphi]$ for which we can safely send $\Lambda_0 \rightarrow \infty$?

The Simplest Answer

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Are there effective actions $S_{\Lambda, \Lambda_0}[\varphi]$ for which we can safely send $\Lambda_0 \rightarrow \infty$?

The Simplest Answer

- Rescale all quantities, using Λ
- Only dimensionless variables appear
- Fixed points of the ERG correspond to continuum limits!

$$S_*(\varphi) = 0$$

- S_* is independent of all scales, including Λ_0
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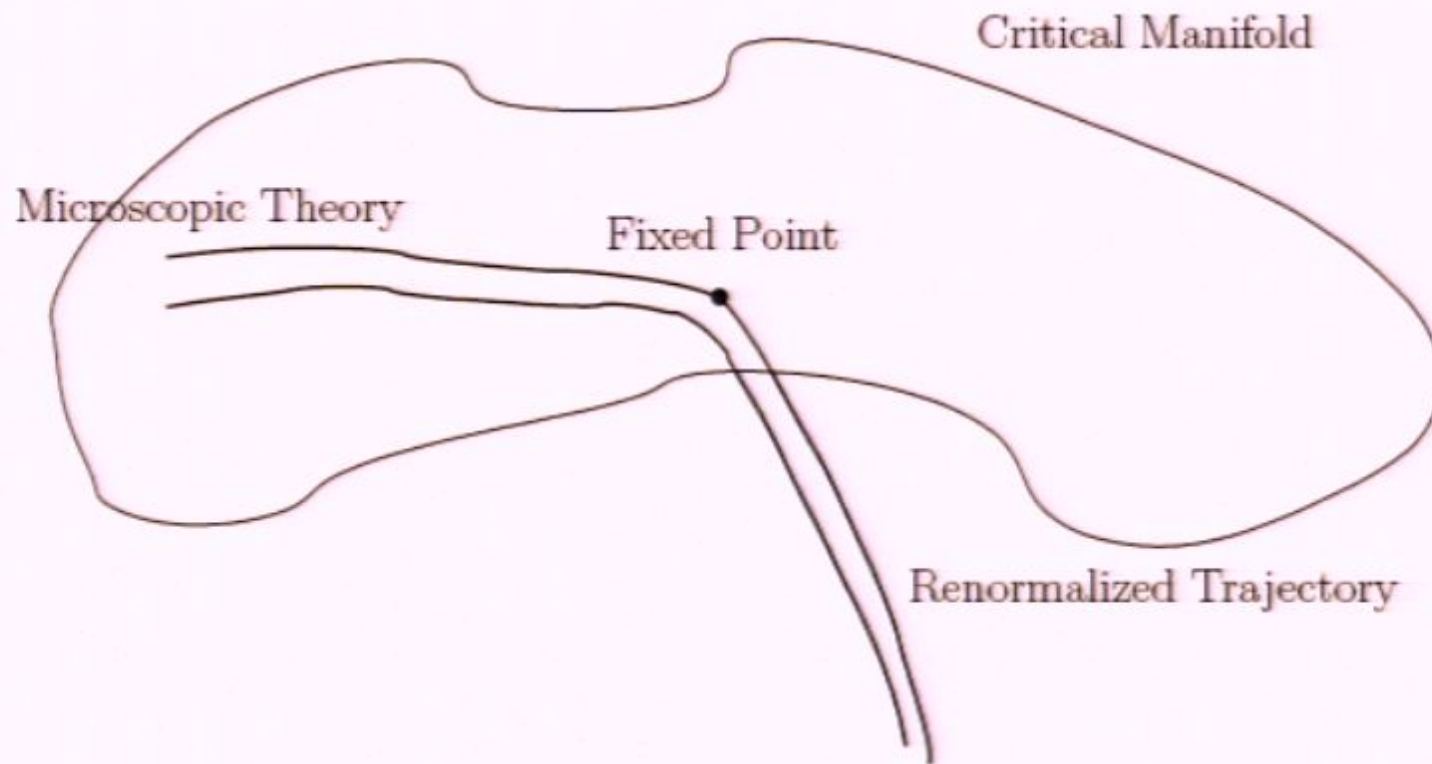
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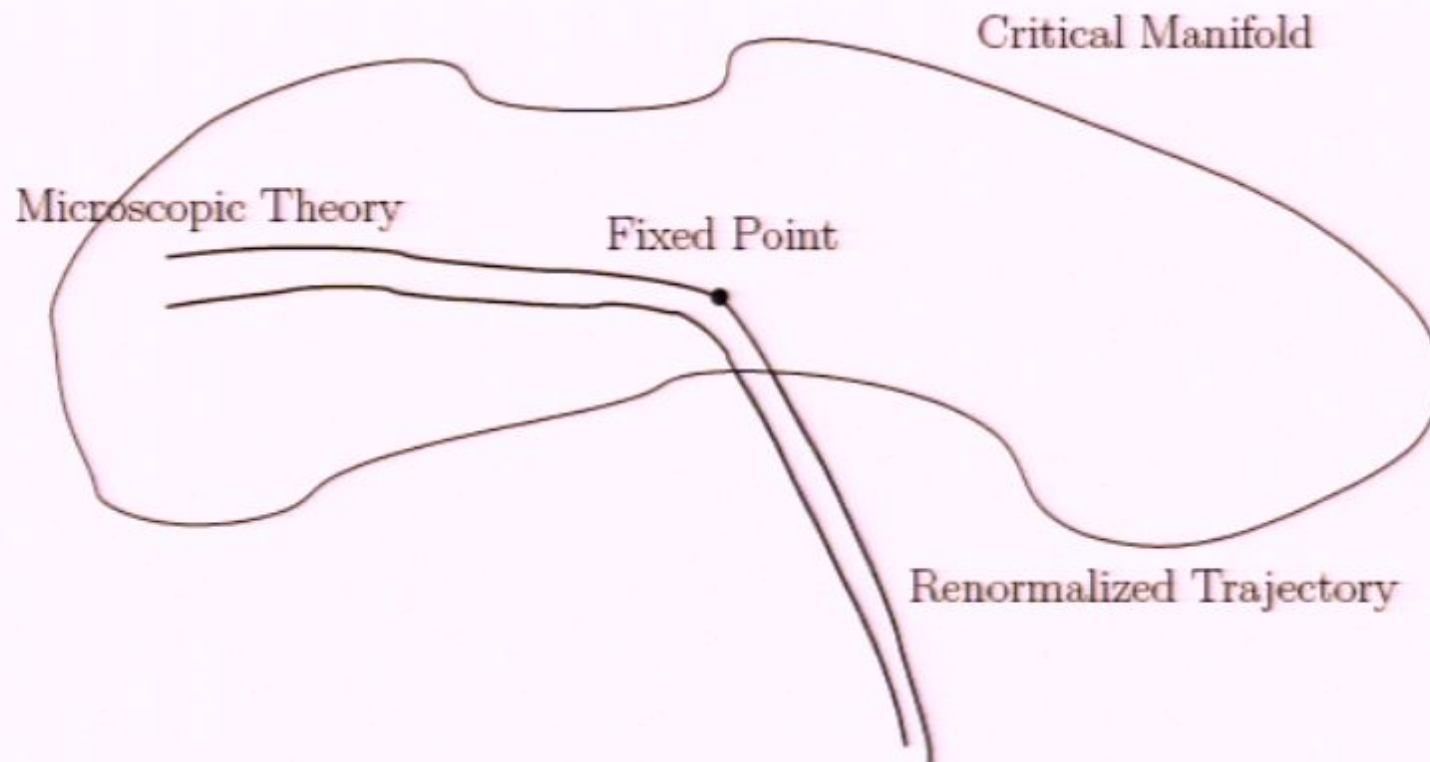
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Continuum Limits II

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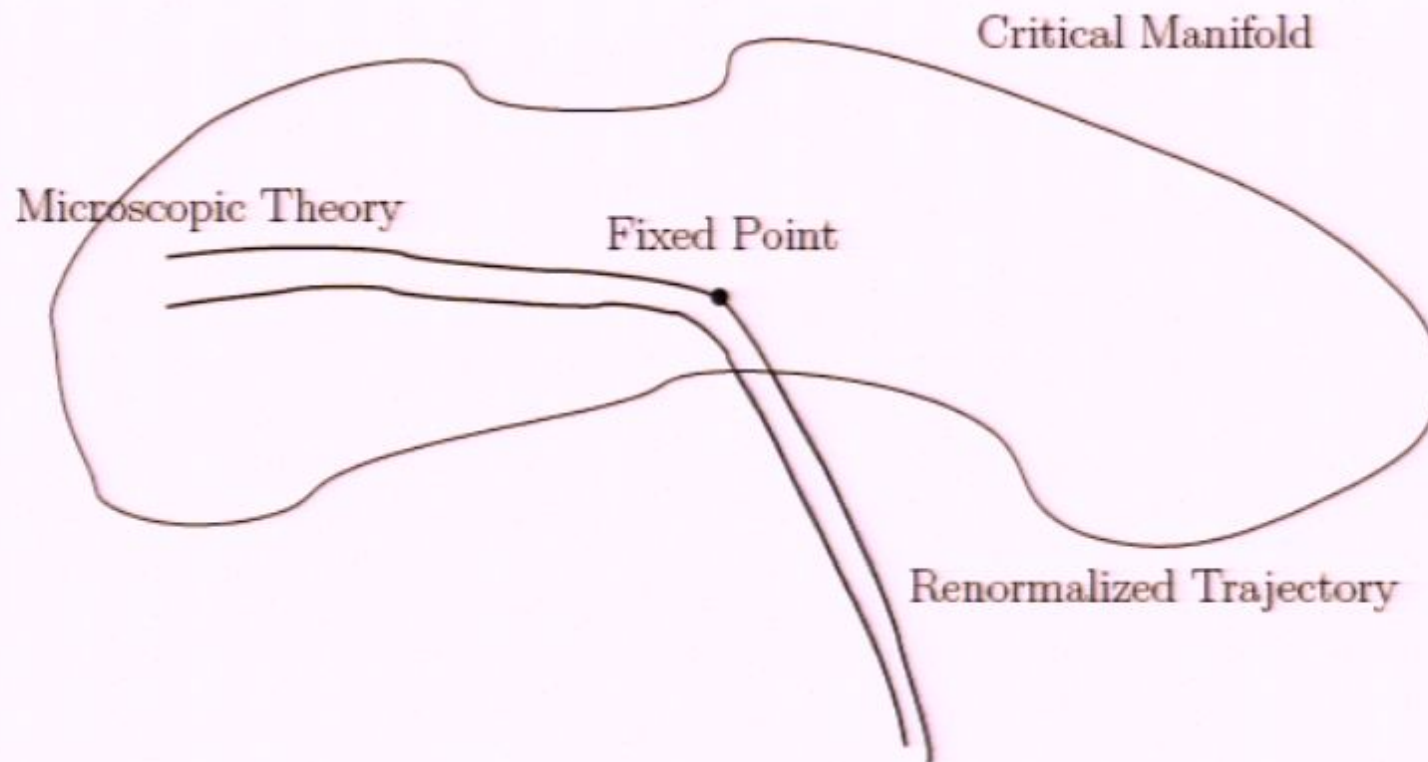


Continuum Limits II



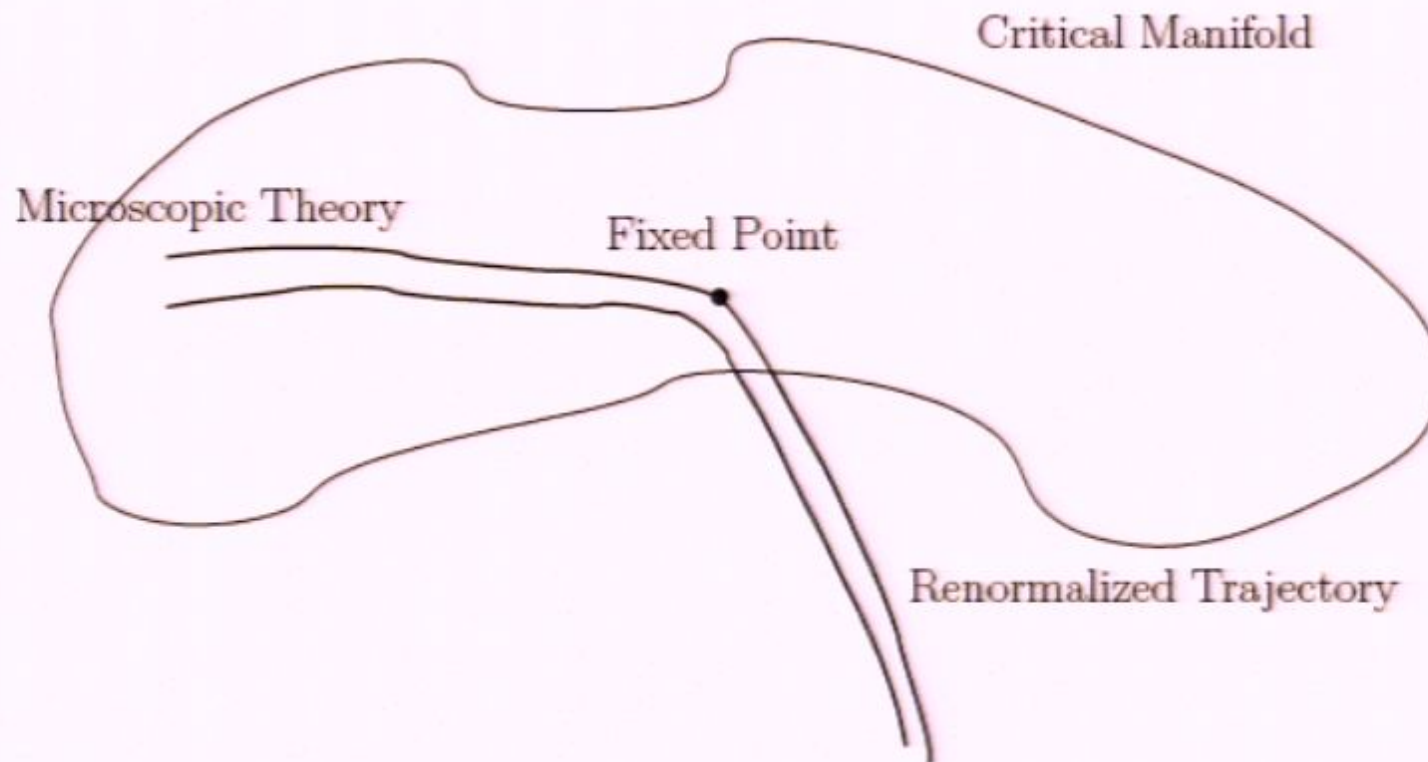
- Tune the trajectory towards the critical surface, as $\Lambda_0 \rightarrow \infty$

Continuum Limits II



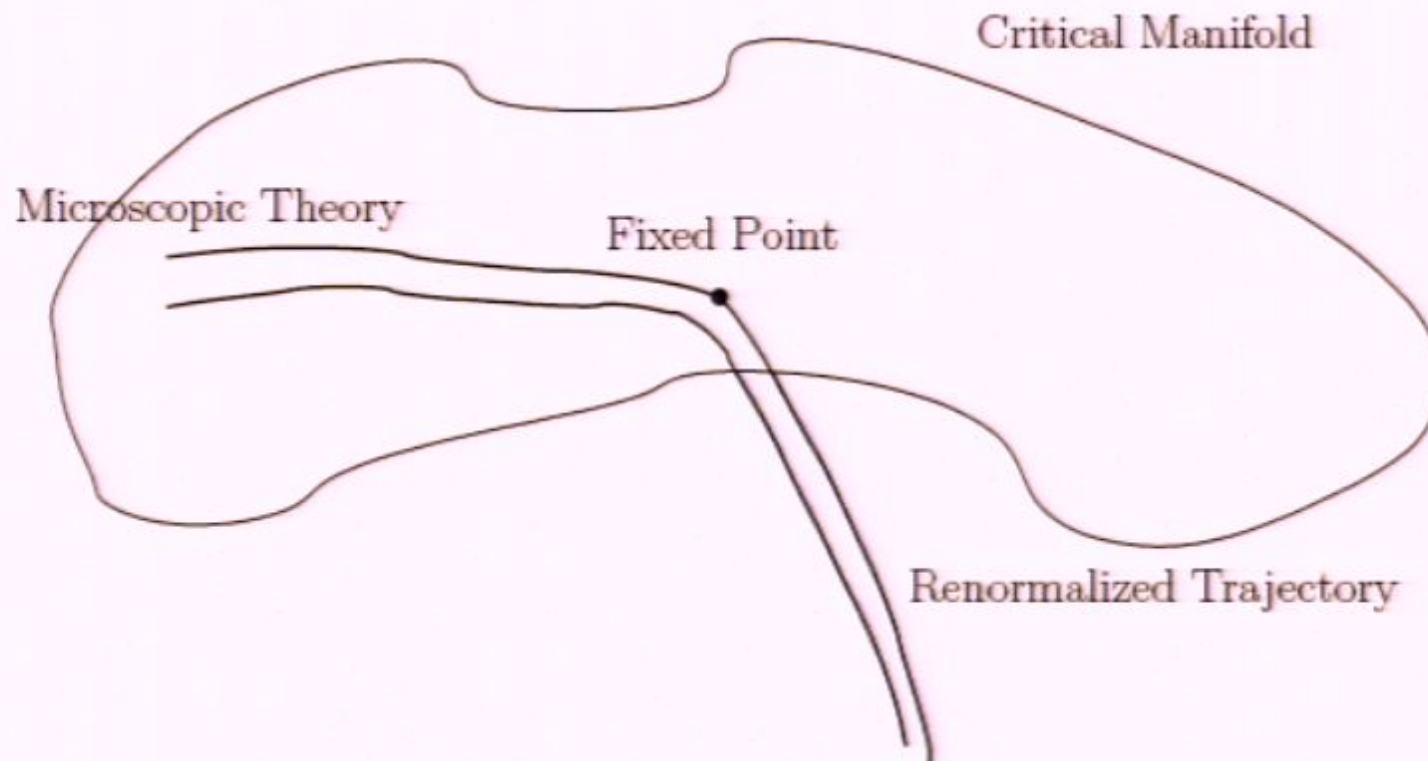
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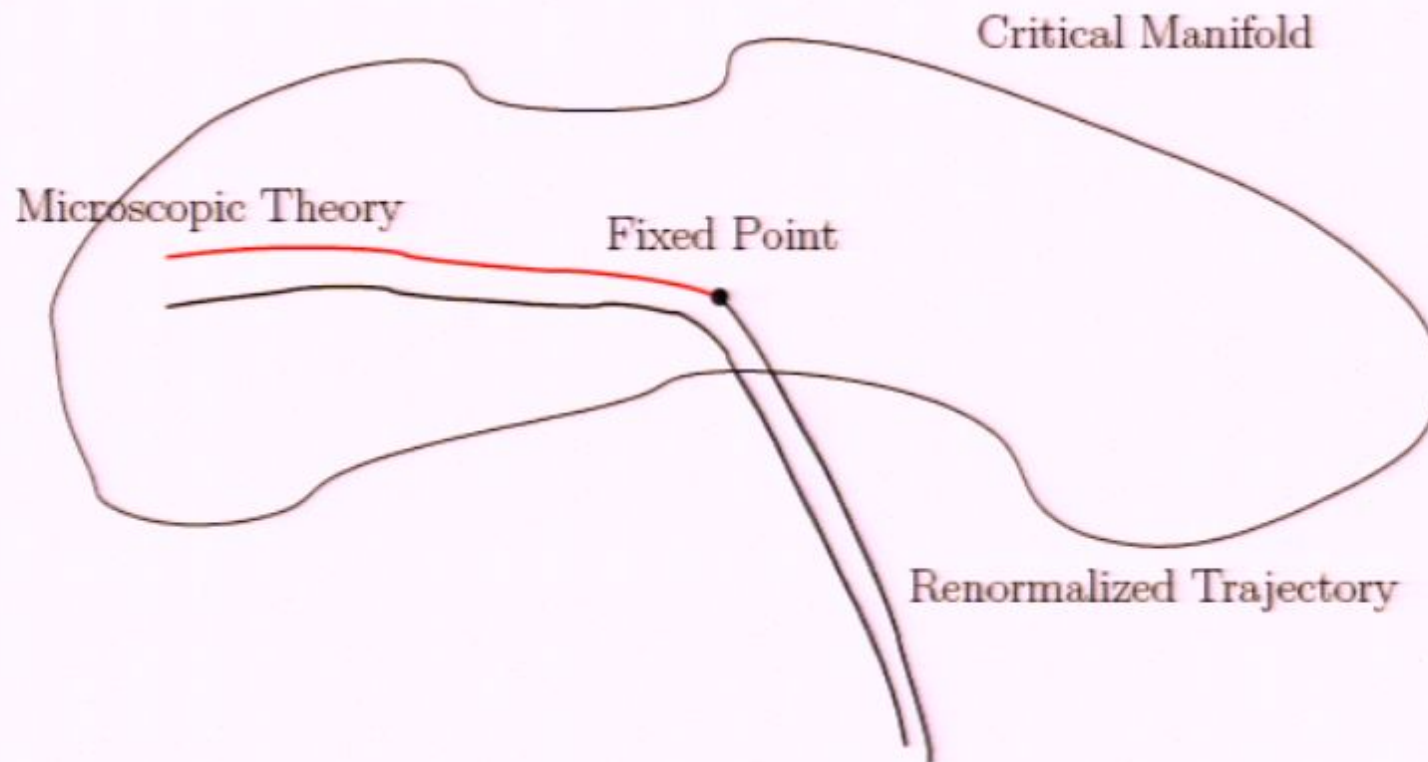
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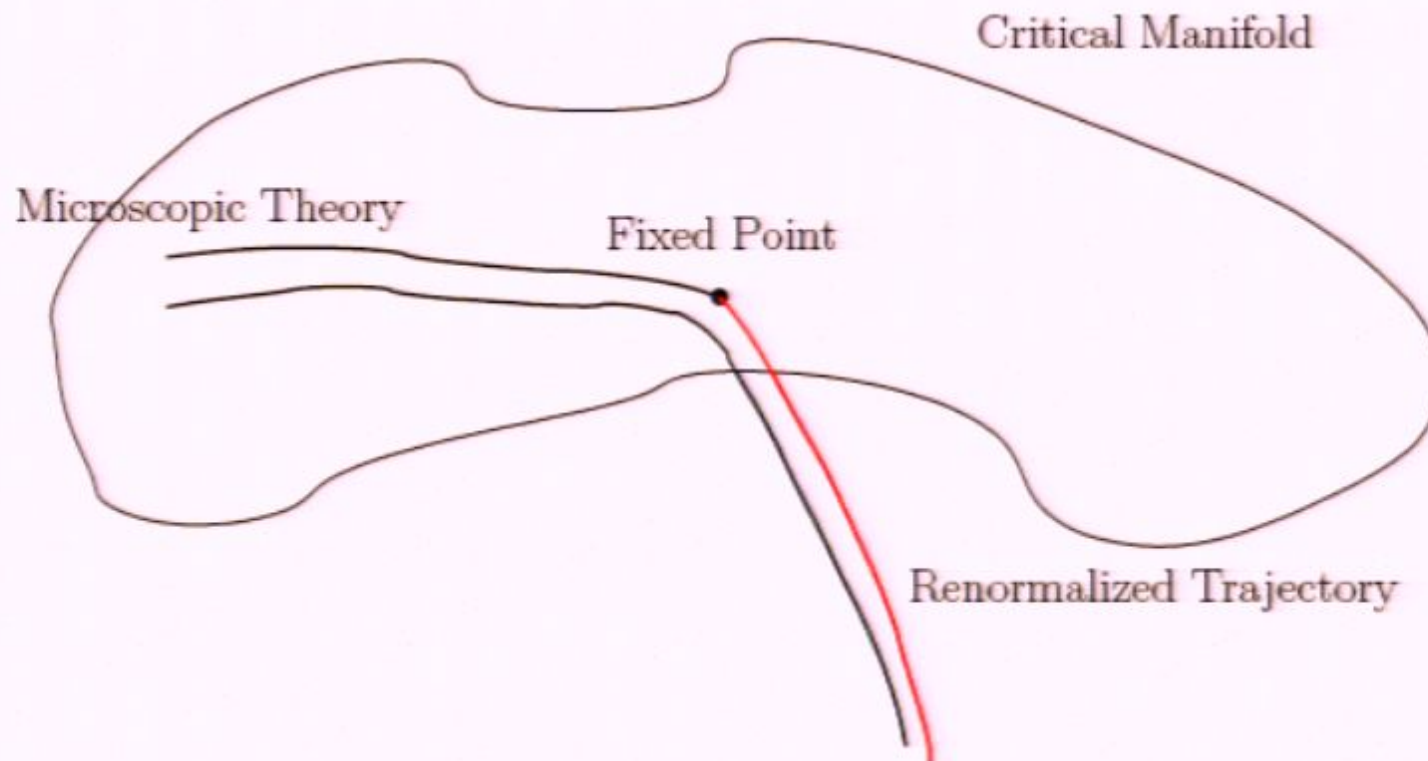
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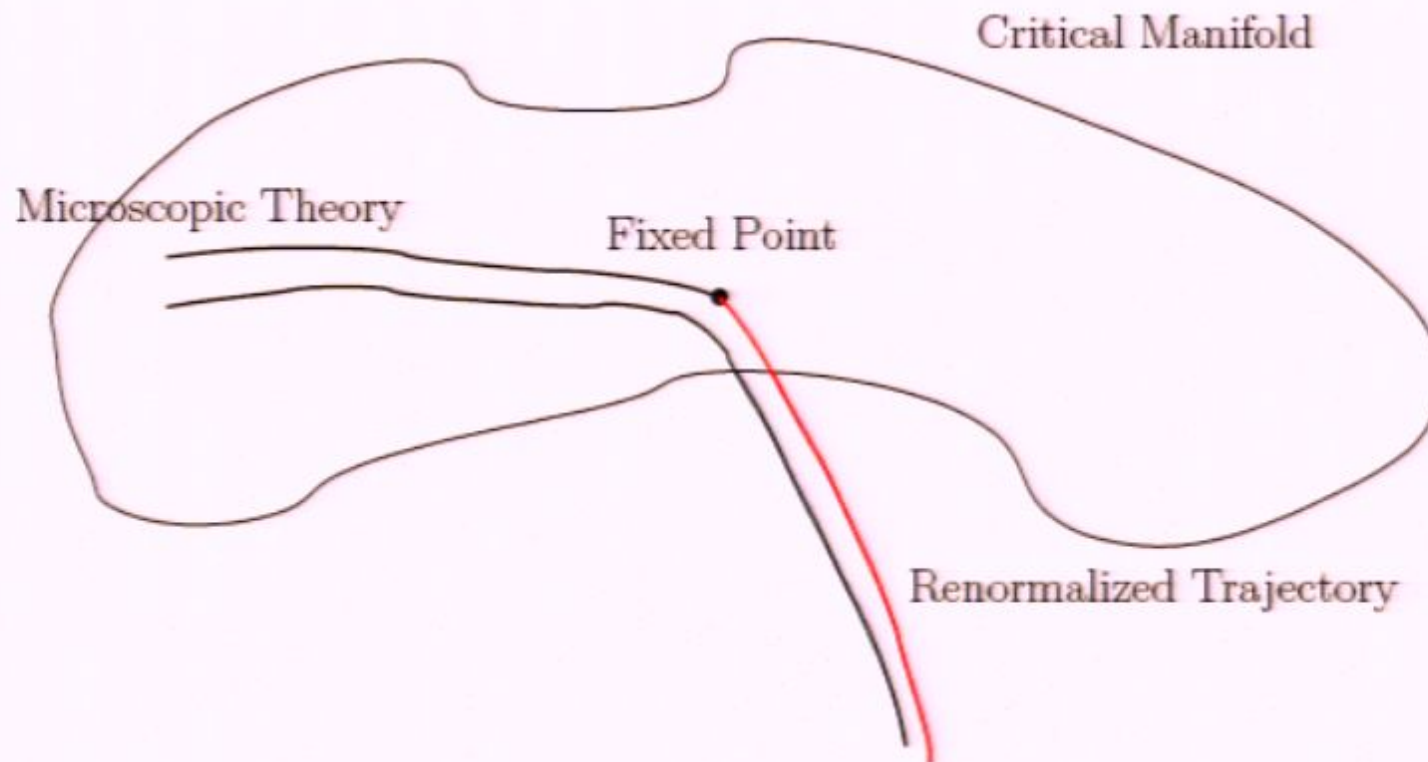
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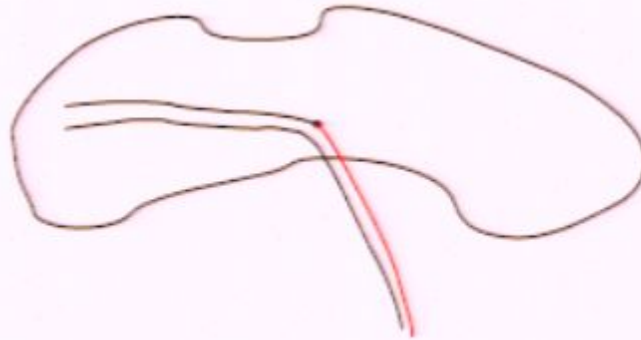
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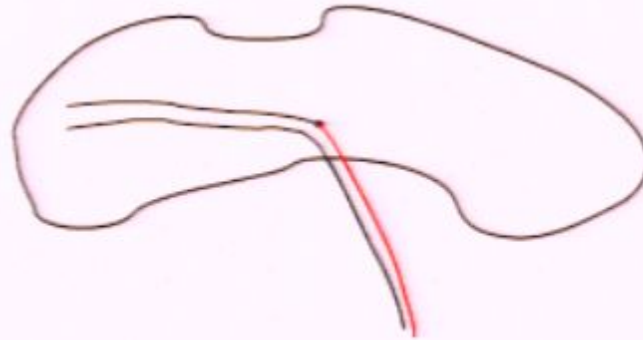
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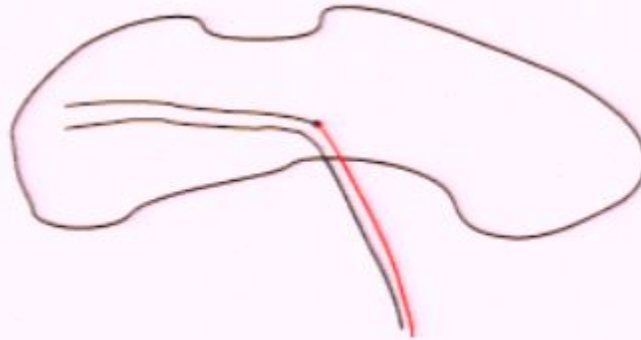


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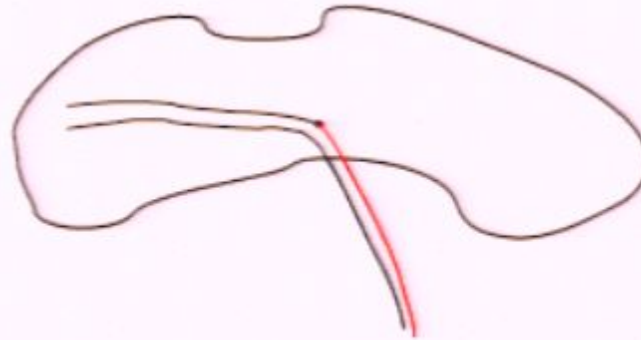
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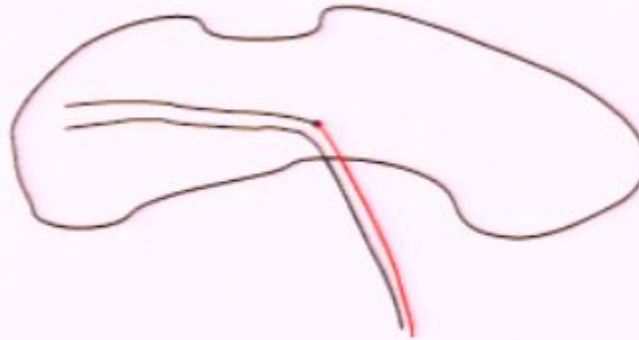
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$$\lim_{t \rightarrow -\infty} S_t[\varphi] = S_*[\varphi] + \sum_{i=1}^n \alpha_i e^{\lambda_i t} \mathcal{O}_i$$

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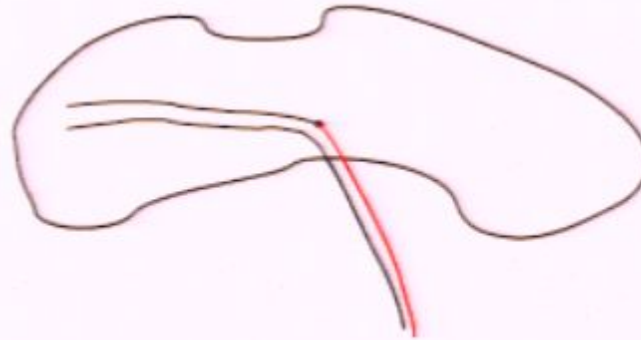


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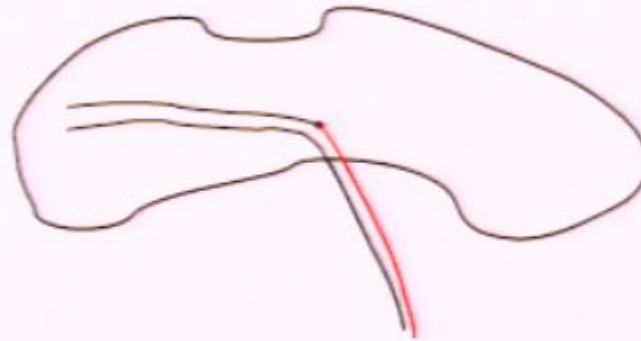


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- If there are marginally relevant terms, include additional terms which sink logarithmically into the fixed point

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$$\lim_{\Lambda \rightarrow \infty} S_\Lambda[\varphi] = S_\star[\varphi] + \sum_{i=1}^n \alpha_i (\mu/\Lambda)^{\lambda_i} \mathcal{O}_i$$

- $S_\Lambda[\varphi] = S_\Lambda[\varphi](\alpha_1, \dots, \alpha_n)$
- Renormalization Conditions:

- From the functions $g_i(\Lambda), \gamma(\Lambda)$

- Therefore, we can trade Λ and α_i for $g(\Lambda)$ and $\gamma(\Lambda)$

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- Only the Gaussian FP exists
- The mass is relevant
- The four point coupling is marginally irrelevant
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M. Beneke, "Renormalons," *Phys. Rept.* **317** (1999) 1, [hep-th/9807443](#)



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- Then there exist numbers K_N such that

$$\left| R(\alpha) - \sum_{n=0}^N r_n \alpha^n \right| < K_{N+1} \alpha^{N+1}$$

for all α in \mathcal{C}

- For asymptotic series, the truncation error decreases at first, but then increases
- If this happens after a large number of terms, N_* , then the truncation error is exponentially small

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M. Beneke, "Renormalons," *Phys. Rept.* 317 (1999) 1, hep-th/9807443

- Consider the divergent series $R \sim \sum_n r_n \alpha^n$
- Divergent series can be useful if they are asymptotic to R in a region, \mathcal{C} , of the complex α -plane
- Then there exist numbers K_N such that

$$\left| R(\alpha) - \sum_{n=0}^N r_n \alpha^n \right| < K_{N+1} \alpha^{N+1}$$

for all α in \mathcal{C}

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$$\bar{R} = \int_0^{\infty} dt e^{-t/\alpha} B[R](t)$$

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$$\lambda_{\text{bare}} = f(\lambda_n) + \cancel{\lambda_0}$$

Massive, Trivial



$$\lambda_{\text{Gure}} = \underbrace{f(\lambda_n)}_{\text{circled}} + \cancel{\lambda/\lambda_0}$$

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Scalar Field Theory: Four Dimensions

Nonperturbative

- Only the Gaussian FP exists
- The mass is relevant
- The four point coupling is marginally irrelevant
- All other couplings are irrelevant
- The only nonperturbatively renormalizable scalar field theories in four dimensions are trivial!

Perturbative

- Order by order in perturbation theory, $\lambda\varphi^4$ is renormalizable
- $\lambda_{\text{bare}} = \sum a_n \lambda_{\text{renorm}}^n + f(\Lambda_{\text{renorm}}/\Lambda_0)$
- Formally, $\lim_{\Lambda_0 \rightarrow \infty} f(\Lambda_{\text{renorm}}/\Lambda_0) = 0$



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Gaussian Fixed Point

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Wilson-Fisher Fixed Point



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Scalar Field Theory: Three Dimensions

Example of a Continuum Limit in $D=3$

Wilson-Fisher FP



Gaussian FP



Asymptotic Freedom etc.



Asymptotic Freedom etc.

Triviality

Asymptotic Freedom

Asymptotic Safety

GFP



no
interacting
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NT FP



renormalizability
determined
in UV

FP

Theory appears
non renormalizable
in IR



What is the Bare Action?



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Renormalized Trajectories

- Start from a fixed point ($\Lambda = \infty$)
- Choose the integration constants associated with the relevant directions
- Solve for the RT action at all other scales
- Actions along an RT are sometimes called 'Perfect Actions'
- Bare action is sometimes referred to as
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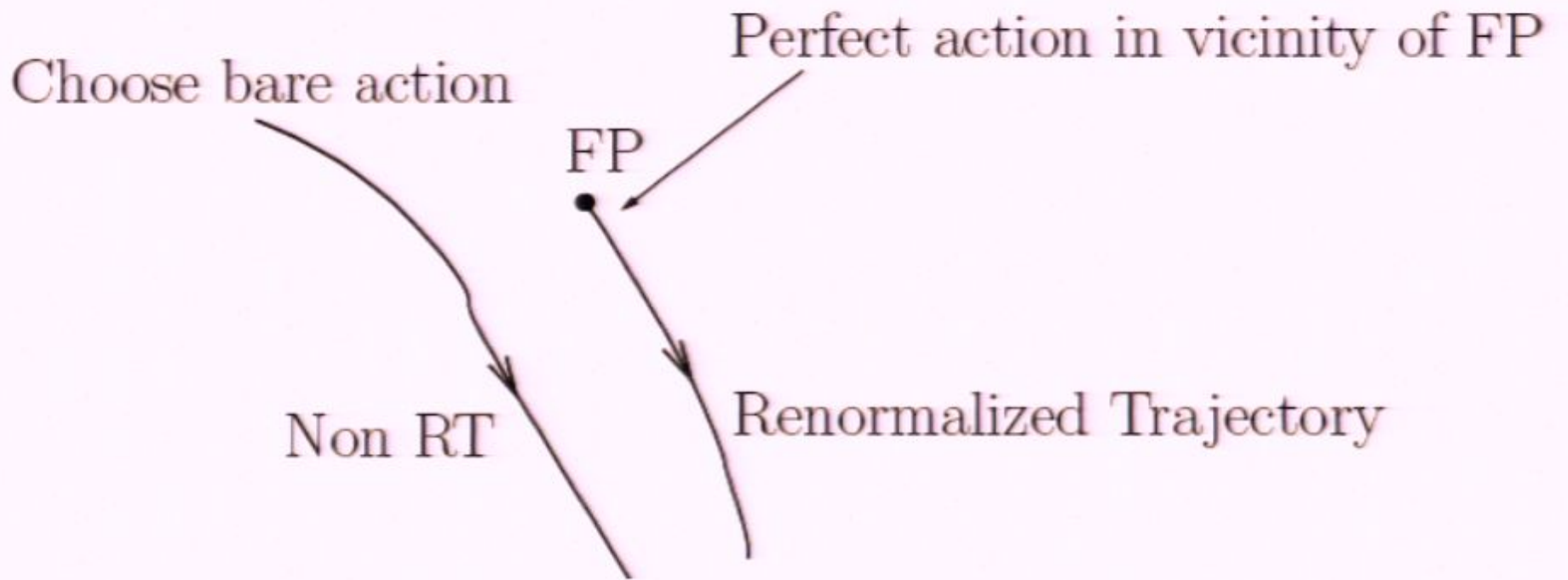
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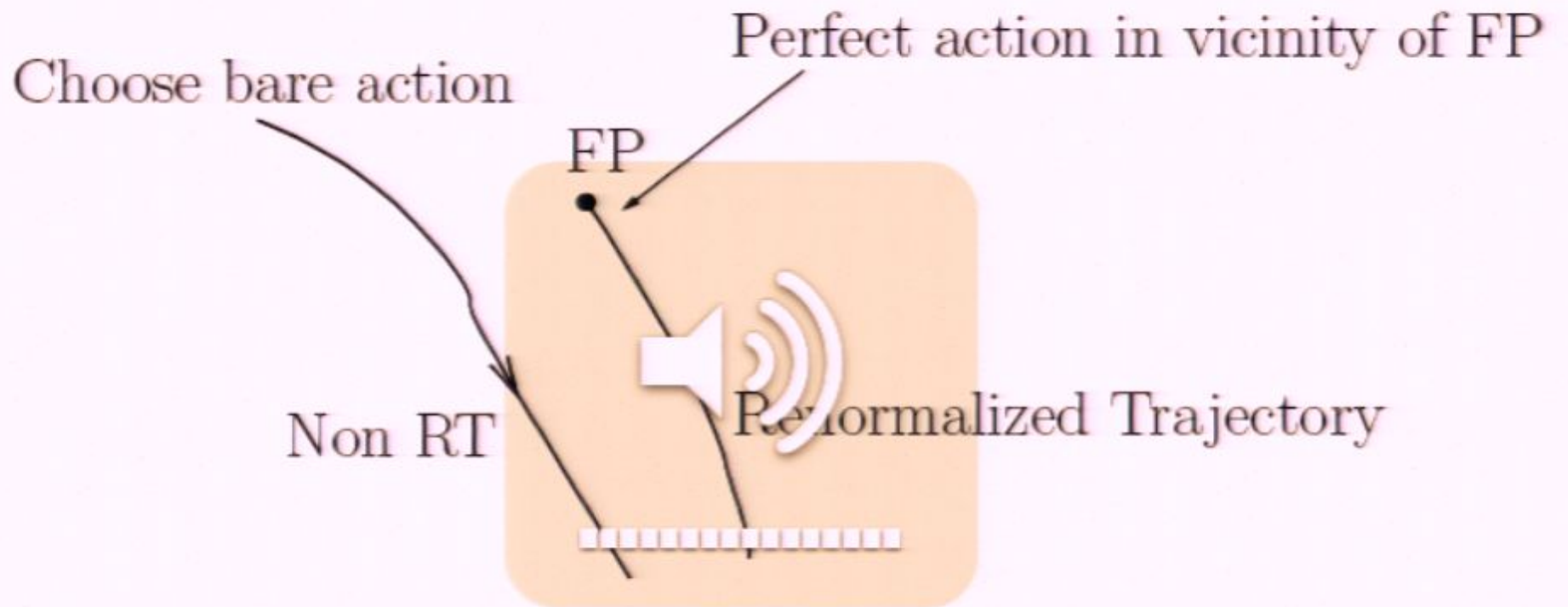
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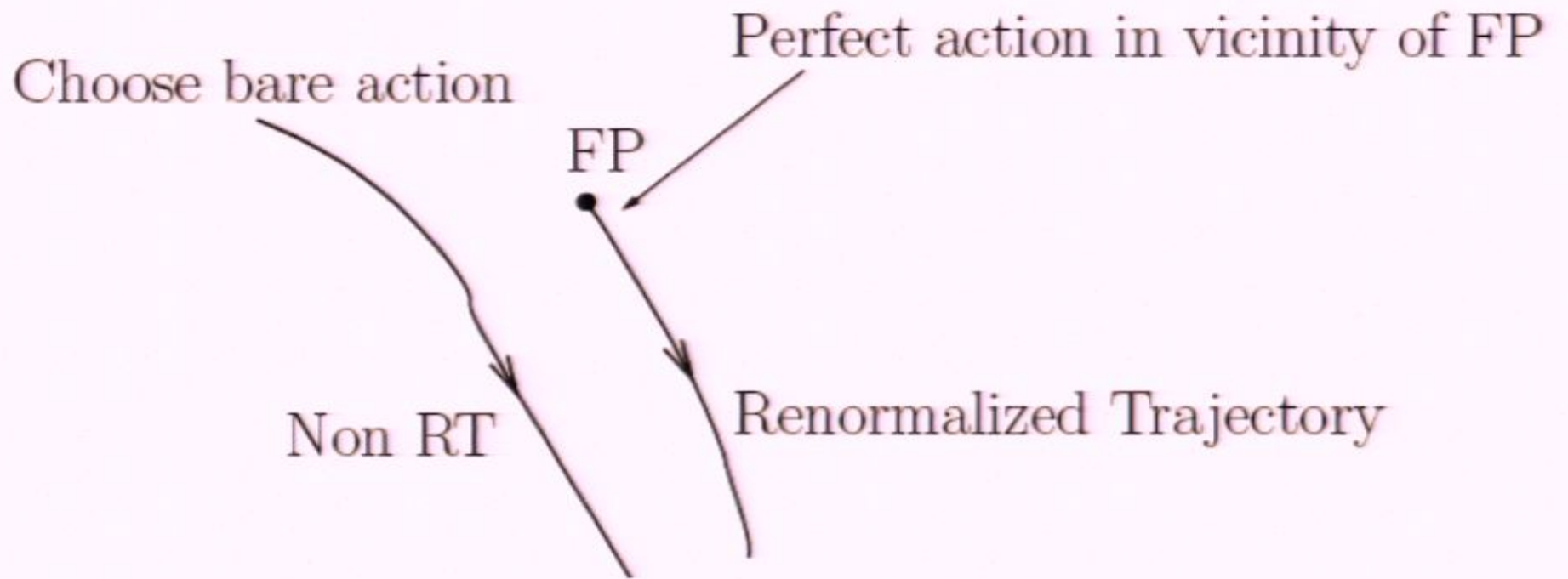
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How Much do the Relevant Directions Tell Us?



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- Along an RT

$$S_\Lambda[\varphi] = S[\varphi](g_1(\Lambda), \dots, g_n(\Lambda), \gamma(\Lambda))$$

- All scale dependence occurs through $g_i(\Lambda)$ and $\gamma(\Lambda)$
- But this does not mean that eg the non-trivial RT leaving the Gaussian fixed point in $D = 3$ does not develop $c^{(6)}\varphi^6$, $c^{(8)}\varphi^8 \dots$ interactions
- It just means that $c^{(6)}(\Lambda) = c^{(6)}(g_i, \gamma)$, with $c^{(6)}(\Lambda = \infty) = 0$
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$D = 3$ scalar field theory



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- Classification of the relevant / irrelevant directions of the Gaussian Fixed Point
- Computing the approximate scale dependence of the action in the vicinity of the Gaussian fixed point

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Critical Point versus Fixed Point



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- Consider a fixed point with n relevant directions then
 - To tune a model on to the critical surface, an experimentalist must 'adjust n knobs'
 - Eg for a ferromagnet, we adjust the external magnetic field (to zero) and the temperature
 - Tuning these knobs moves our model around in parameter space
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 - Under ERG transformations all critical point theories in the same universality class flow into the same fixed point



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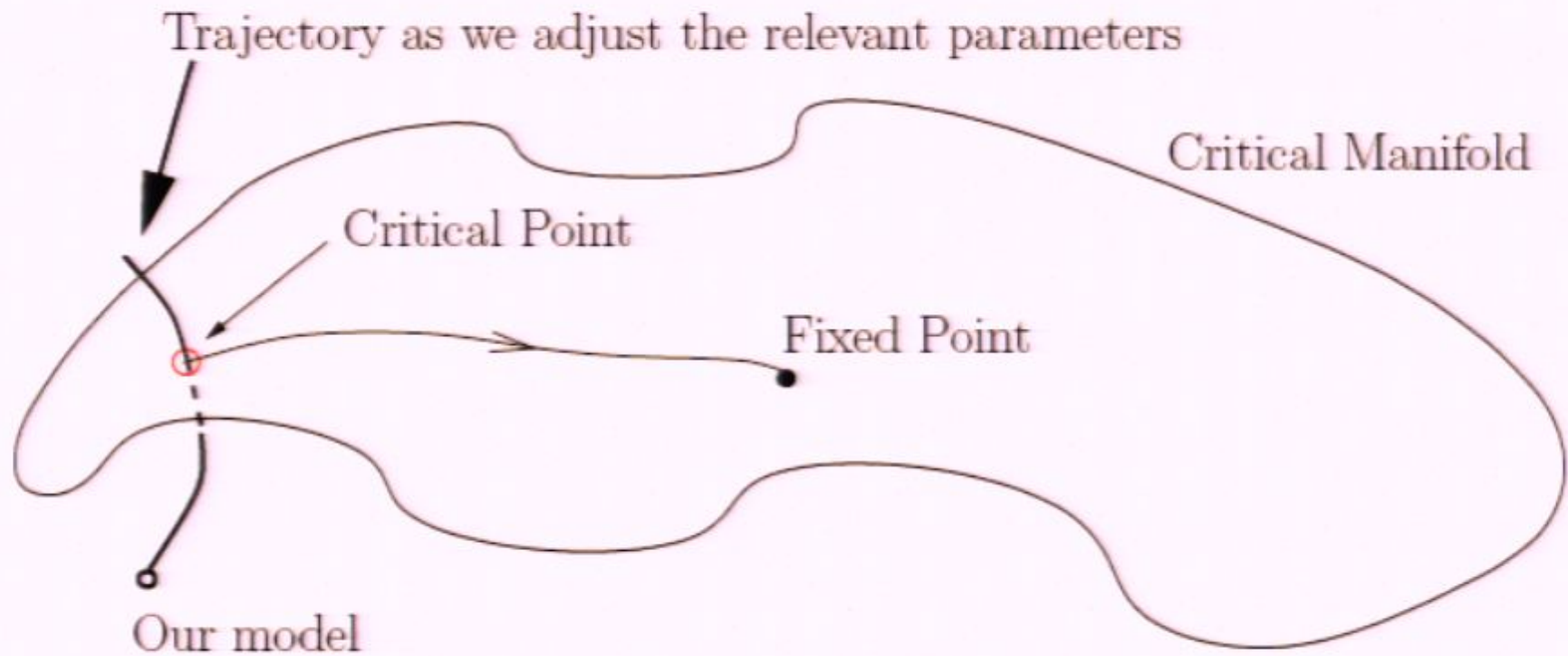


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Recap



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- The ERG relates physics at different scales via a **coarse-graining procedure**
- Under an ERG transformation, the effective action evolves
- Including a rescaling step, we find fixed points,
- By linearizing the ERG equation, we define the scaling fields which can be relevant, irrelevant, marginal (or redundant)
- The relevant (and marginally relevant) directions can be used to construct scale dependent continuum limits
- These 'Renormalized Trajectories' (RTs) are self-similar
- We can understand, intuitively, what it is for a QFT to be nonperturbatively renormalizable



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Scalar Field Theory: Three Dimensions

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Gaussian Fixed Point

- The mass term is relevant
- The four-point coupling is relevant
- The six-point coupling is marginally irrelevant
- Non-trivial renormalizable theories exist along the $\lambda\phi^4$ direction!

Wilson-Fisher Fixed Point

- In addition to the Gaussian FP, there is a non-trivial FP
- The W-F FP possesses a single relevant direction
- This can also be used to construct an RT

↓ NT Relevant direction

$$T_{\text{time}} = f(n) + \cancel{n \log n}$$



Massive, Trimed



