

Title: Cosmology present status and future prospects

Date: Apr 02, 2008 02:00 PM

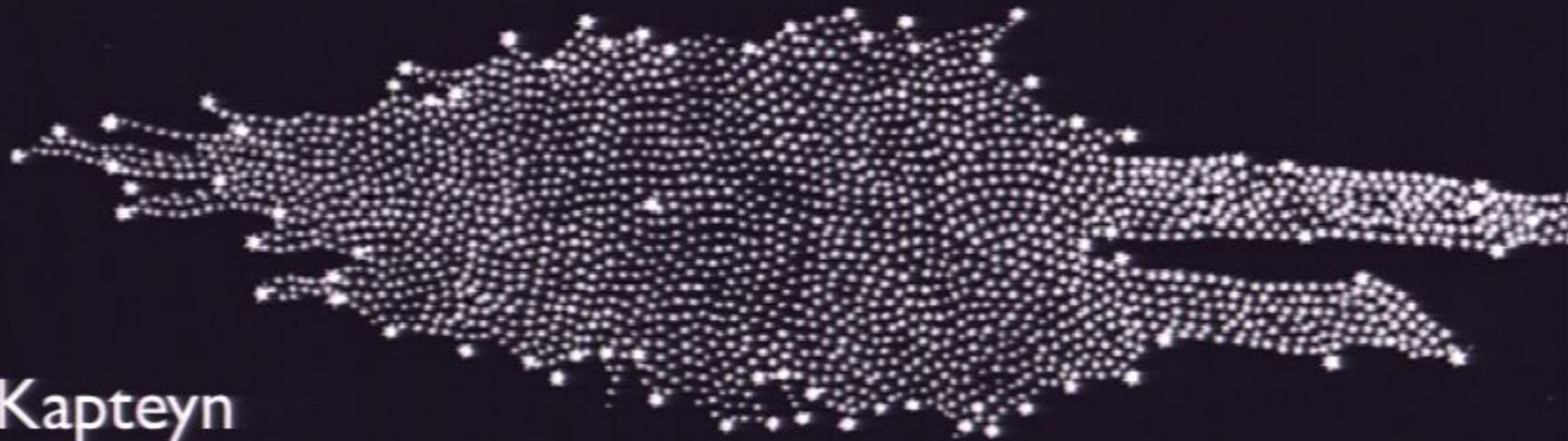
URL: <http://pirsa.org/08040002>

Abstract: I will summarize current observational constraints in cosmology with emphasis on what we have learned about the properties of the primordial density perturbations. I will describe future directions including observations of high redshift neutral hydrogen through is 21 cm line.

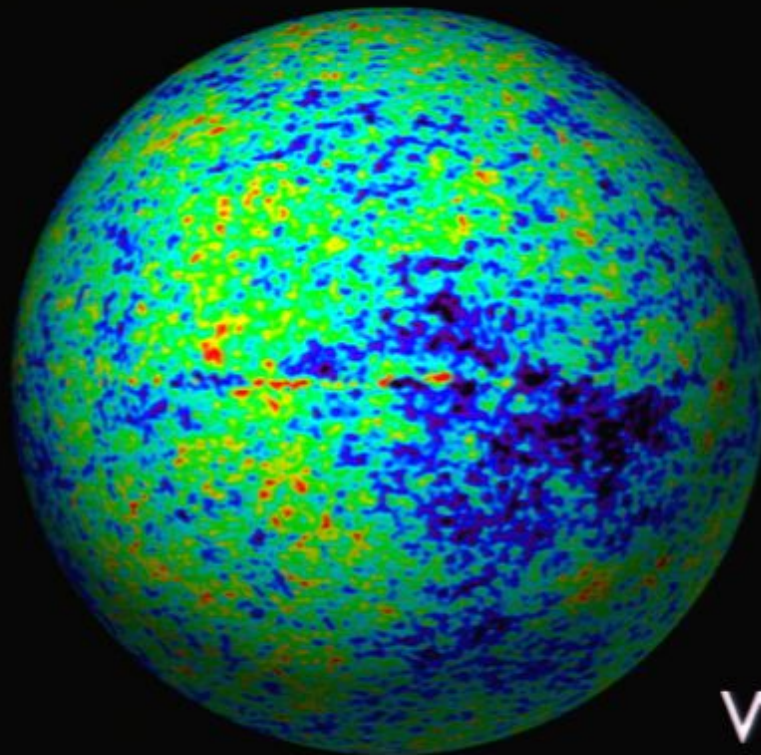
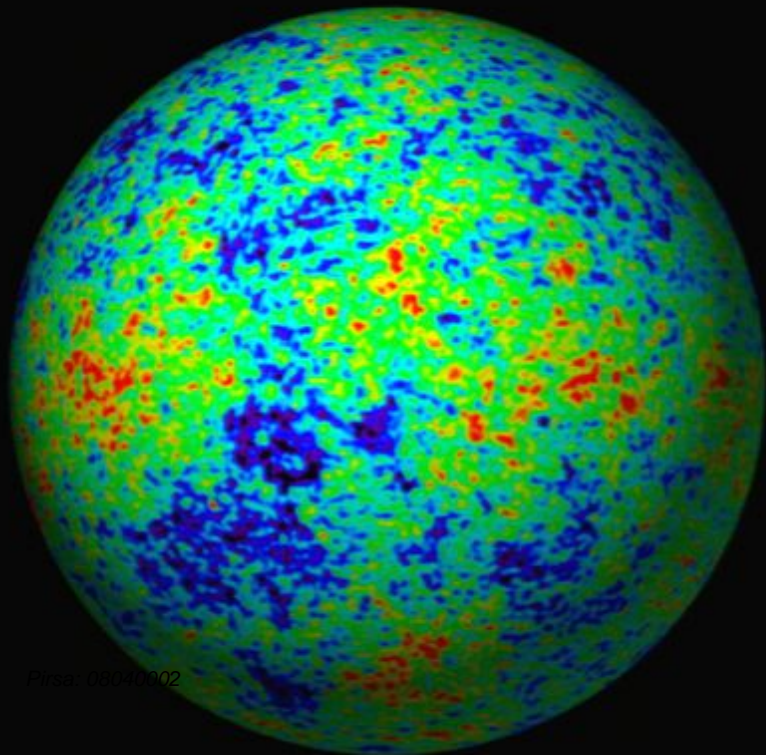
# Cosmology: Present status and future prospects

Matias Zaldarriaga

A century of progress



Kapteyn



WMAP

# The Era of Precision Cosmology

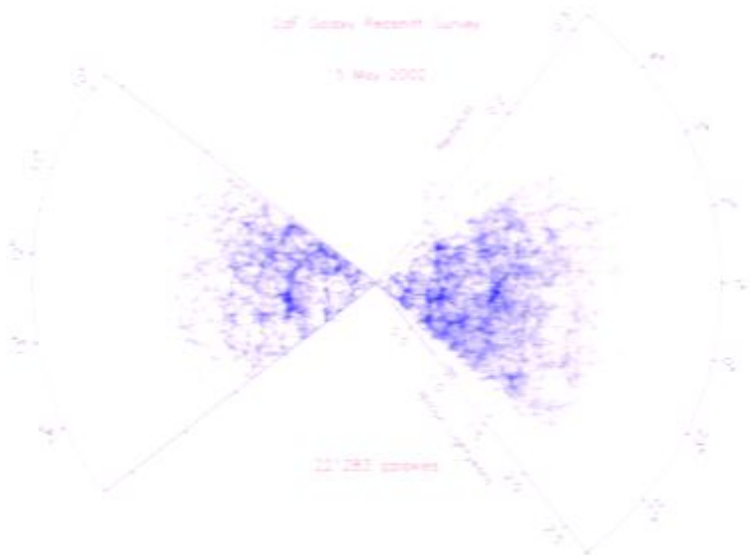
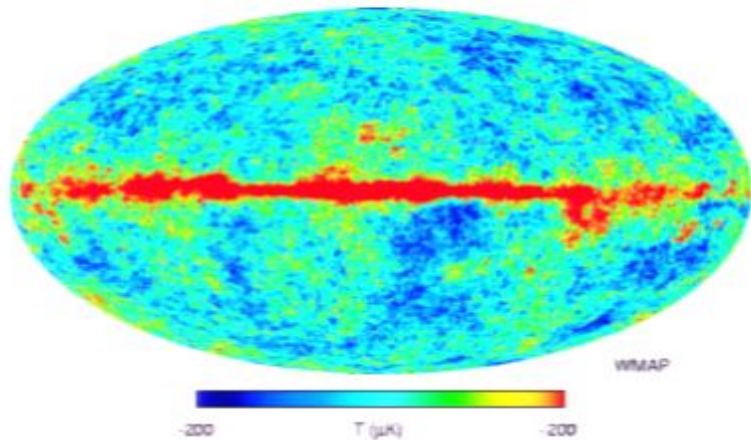


Table 6. Cosmological Parameter Summary

Description	Symbol	WMAP-only	WMAP+BAO+SN
Parameters for Standard $\Lambda$ CDM Model <sup>a</sup>			
Age of universe	$t_0$	$13.69 \pm 0.11$ Gyr	$13.73 \pm 0.12$ Gyr
Hubble constant	$H_0$	$71.9^{+0.6}_{-0.7}$ km/s/Mpc	$70.1 \pm 1.3$ km/s/Mpc
Baryon density	$\Omega_b$	$0.0441 \pm 0.0020$	$0.0462 \pm 0.0015$
Physical baryon density	$\Omega_b h^2$	$0.02271 \pm 0.00062$	$0.02265 \pm 0.00059$
Dark matter density	$\Omega_c$	$0.214 \pm 0.027$	$0.233 \pm 0.013$
Physical dark matter density	$\Omega_c h^2$	$0.1099 \pm 0.0062$	$0.1143 \pm 0.0034$
Dark energy density	$\Omega_\Lambda$	$0.742 \pm 0.030$	$0.721 \pm 0.015$
Curvature fluctuation amplitude, $k_0 = 0.002 \text{ Mpc}^{-1} h$	$\Delta_{\zeta}^2$	$(2.41 \pm 0.11) \times 10^{-9}$	$(2.48^{+0.092}_{-0.093}) \times 10^{-9}$
Fluctuation amplitude at $8h^{-1} \text{ Mpc}$	$\sigma_8$	$0.796 \pm 0.036$	$0.817 \pm 0.026$
$10^{10} C_{\ell}^{TT}(\ell=2)$	$C_{\ell=20}$	$5756 \pm 42 \mu\text{K}^2$	$5748 \pm 41 \mu\text{K}^2$
Scalar spectral index	$n_s$	$0.963^{+0.014}_{-0.015}$	$0.960^{+0.014}_{-0.013}$
Redshift of matter-radiation equality	$z_{\text{eq}}$	$1176^{+151}_{-150}$	$1280^{+88}_{-89}$
Angular diameter distance to matter-radiation eq <sup>c</sup>	$d_A(z_{\text{eq}})$	$14279^{+188}_{-189}$ Mpc	$14173^{+141}_{-139}$ Mpc
Redshift of decoupling	$z_*$	$1090.51 \pm 0.95$	$1091.00^{+0.77}_{-0.73}$
Age at decoupling	$t_*$	$380081^{+5642}_{-5641}$ yr	$375908^{+3148}_{-3115}$ yr
Angular diameter distance to decoupling <sup>c,d</sup>	$d_A(z_*)$	$14115^{+188}_{-191}$ Mpc	$14006^{+142}_{-141}$ Mpc
Sound horizon at decoupling <sup>d</sup>	$r_s(z_*)$	$146.8 \pm 1.8$ Mpc	$145.6 \pm 1.2$ Mpc
Acoustic scale at decoupling <sup>d</sup>	$l_*(z_*)$	$302.08^{+0.82}_{-0.84}$	$302.11^{+0.84}_{-0.82}$
Recombination optical depth	$\tau$	$0.087 \pm 0.017$	$0.084 \pm 0.016$
Redshift of recombination	$z_{\text{reco}}$	$11.0 \pm 1.4$	$10.8 \pm 1.4$
Age at recombination	$t_{\text{reco}}$	$427^{+88}_{-85}$ Myr	$432^{+90}_{-87}$ Myr
Parameters for Extended Models <sup>a</sup>			
Total density <sup>e</sup>	$\Omega_{\text{tot}}$	$1.009^{+0.100}_{-0.095}$	$1.0051 \pm 0.0064$
Equation of state <sup>g</sup>	$w$	$-1.05^{+0.41}_{-0.42}$	$-0.971^{+0.061}_{-0.060}$
Tensor to scalar ratio, $k_0 = 0.002 \text{ Mpc}^{-1} h, h$	$r$	$< 0.42$ (95% CL)	$< 0.20$ (95% CL)
Running of spectral index, $k_0 = 0.002 \text{ Mpc}^{-1} h, h$	$dn_s/d \ln k$	$-0.037 \pm 0.038$	$-0.032^{+0.021}_{-0.030}$
Neutrino density <sup>f</sup>	$\Omega_\nu h^2$	$< 0.014$ (95% CL)	$< 0.0065$ (95% CL)
Neutrino mass <sup>f</sup>	$\sum m_\nu$	$< 1.3 \text{ eV}$ (95% CL)	$< 0.61 \text{ eV}$ (95% CL)
Number of light neutrino families <sup>h</sup>	$N_{\text{eff}}$	$> 1.1$ (95% CL)	$4.4 \pm 1.5$

<sup>a</sup>The parameters reported in the first section assume the 6 parameter  $\Lambda$ CDM model, first using WMAP data only, then using WMAP+BAO+SN data.

<sup>b</sup> $h = 0.002 \text{ Mpc}^{-1}$  —  $h_{\text{ref}} \approx 30$ .

<sup>c</sup>Comoving angular diameter distance

"Cosmologists are often in error, though never in doubt." -- Lev Landau



## Not everything is yet settled:

- **Matter budget:** Dark matter/Dark energy
- **Early universe:** How were the initial seeds created?
- **Non-linear physics:** Galaxy formation and evolution, reionization, the first stars

## The matter content

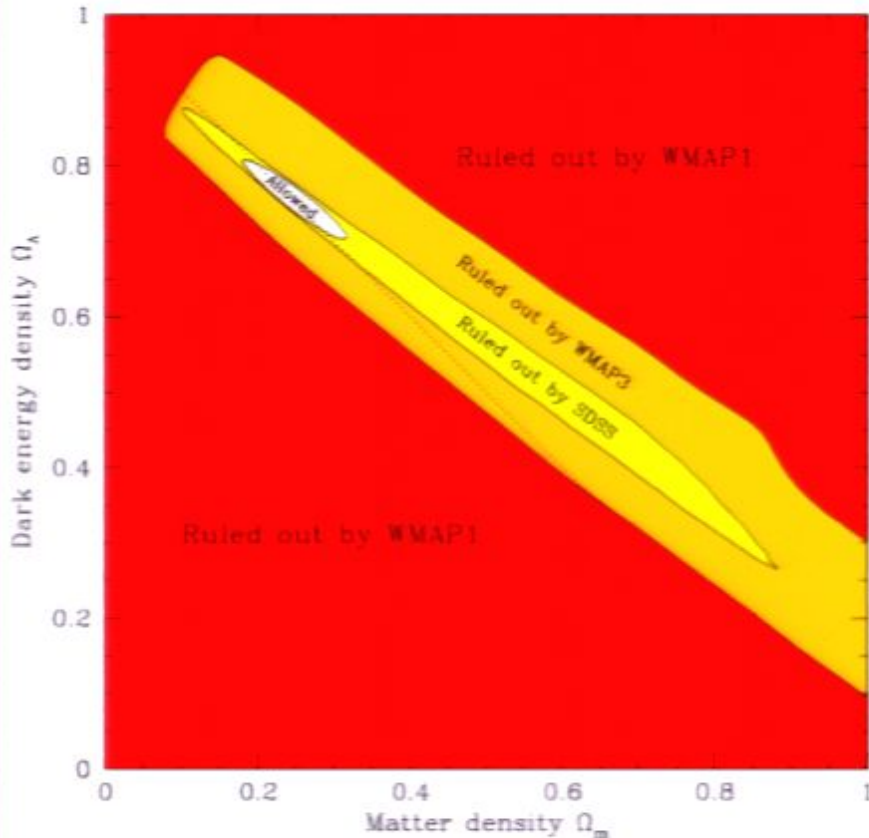


FIG. 15: 95% constraints in the  $(\Omega_m, \Omega_\Lambda)$  plane. The large shaded regions are ruled out by WMAP1 (red/dark grey) and WMAP3 (orange/grey) when spatial curvature is added to the 6 vanilla parameters, illustrating the well-known geometric degeneracy between models that all give the same acoustic peak locations in the CMB power spectrum. The yellow/light grey region is ruled out when adding SDSS LRG information, breaking the degeneracy mainly by measuring the acoustic peak locations in the galaxy power spectrum. Models on the diagonal dotted line are flat, those below are open and those above are closed. Here the yellow banana has been cut off from below by an  $h \gtrsim 0.4$  prior in the CosmoMC software.

The Universe appears to be flat with the energy density dominated by a fluid with equation of state

$$w = -1.004 \pm 0.089$$

0705.3323 (Percival et al.) using BAOs

Dark matter is required to explain the shape of the CMB peaks as well as LSS.

The baryon fraction inferred from the CMB is consistent with BBN.

## Windows into the inflationary epoch:

We now know for a fact that perturbations had to be created during inflation. Which properties of the perturbations can be used to test the model?

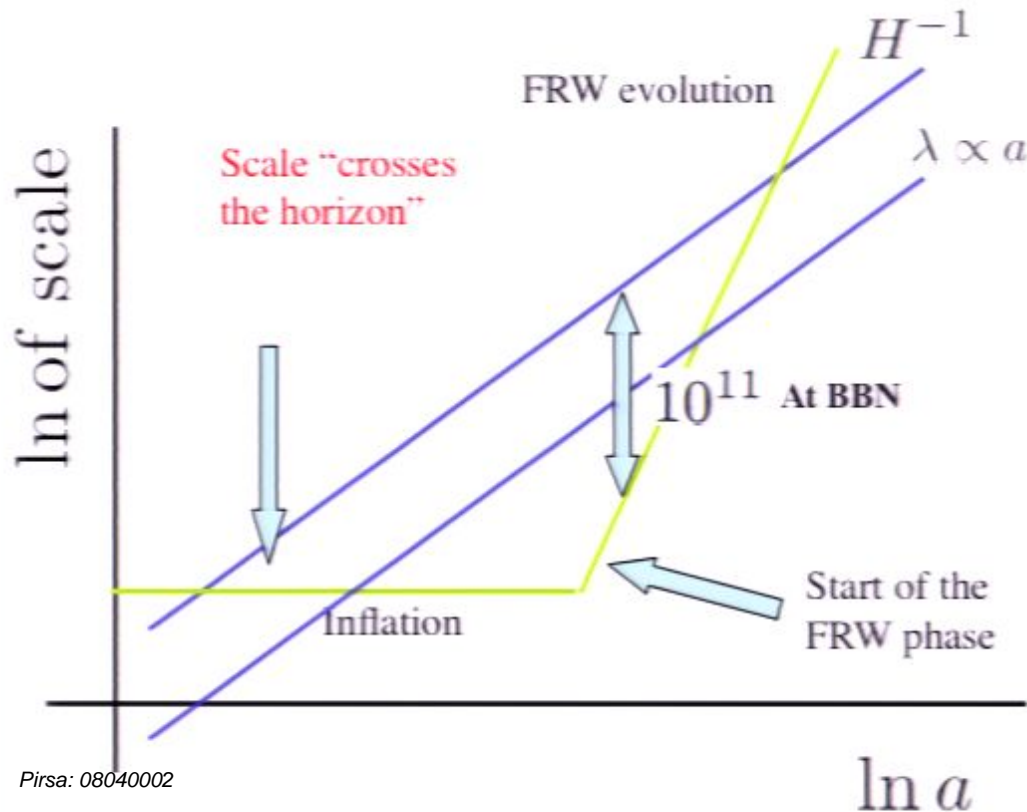
1. Departures from scale invariance of the density perturbations
2. Gravity waves
3. Departures from Gaussianity
4. Isocurvature fluctuations
5. Relics such as cosmic strings

- CMB experiments: WMAP, Planck, CMBPOL, SPT, ACT, etc
- Smaller scales from Lyman alpha forest, lensing, SZ, etc.

Push to smaller scales to have more modes and measure polarization to look for Gravity Waves

# When were perturbations created?

The data definitively shows that they were created during Inflation



$$H = \frac{\dot{a}}{a}$$

$$R_H \sim H^{-1}$$

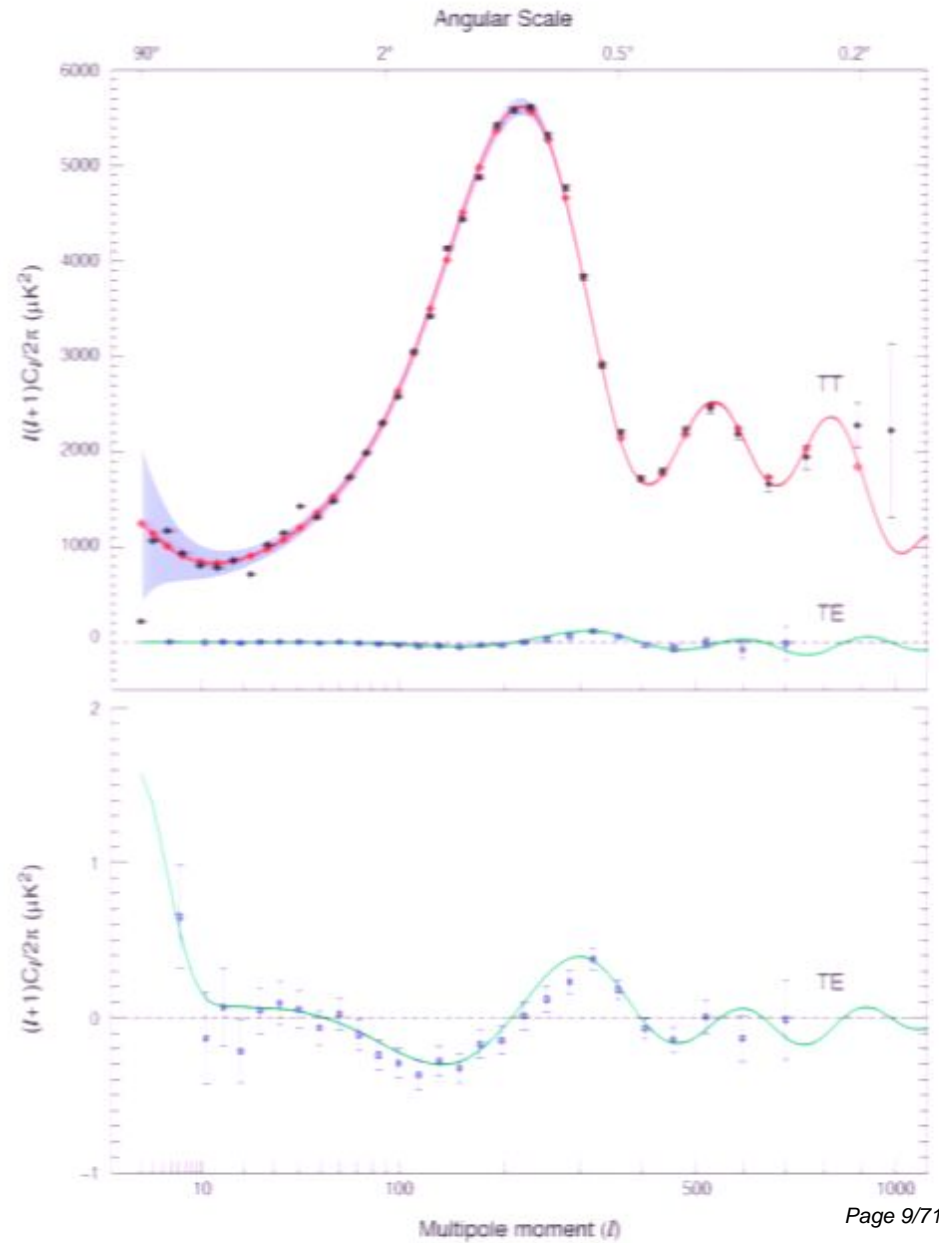
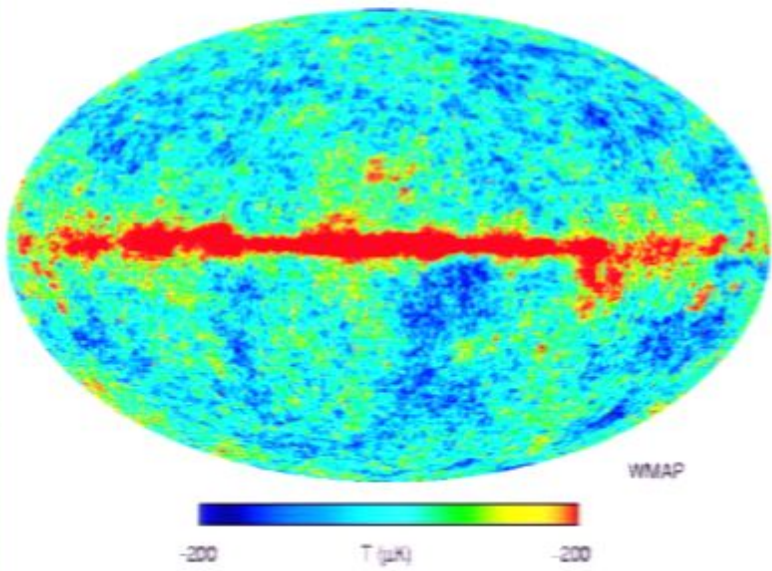
$$\lambda = \lambda_0 a$$

$$\frac{\lambda}{R_H} = \lambda_0 a H \propto \dot{a}$$

Gravity is attractive so  $\dot{a}$  usually decreases

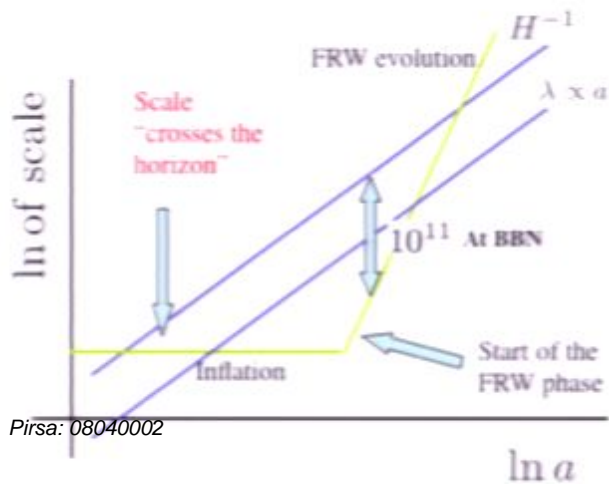
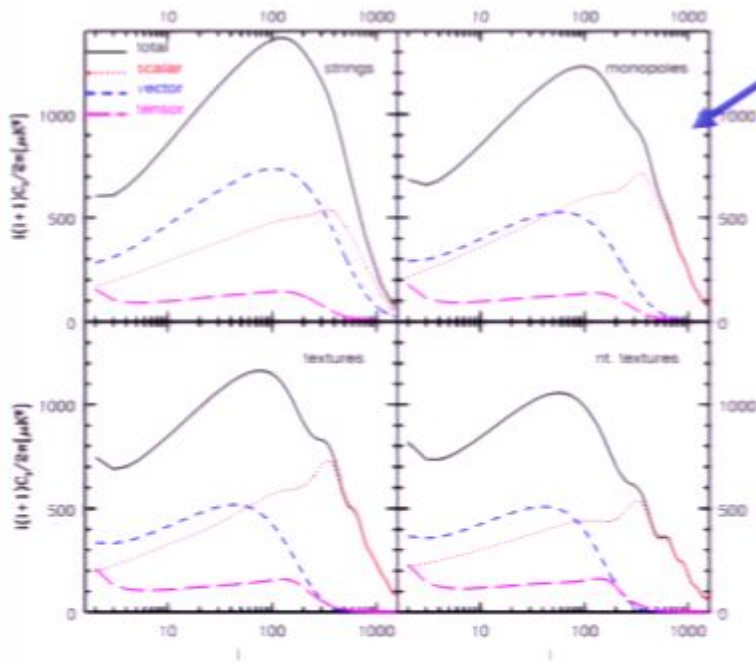


# Power Spectrum



# When were perturbations created?

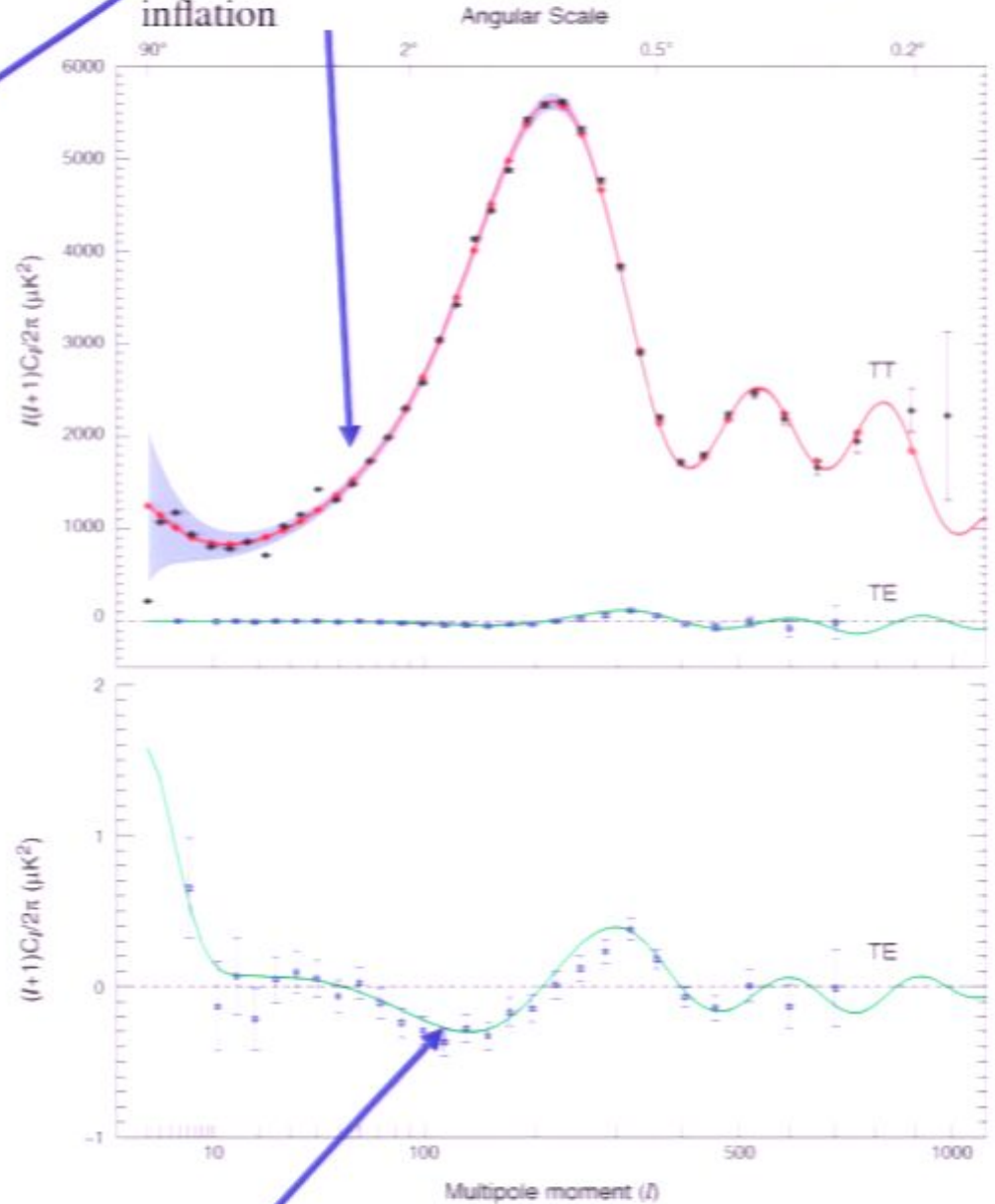
## Causal Seeds Pen. Seljak & Turok (1997)



Pirsa: 08040002

Sharp acoustic peaks are difficult to create without inflation

## WMAP 3yrs



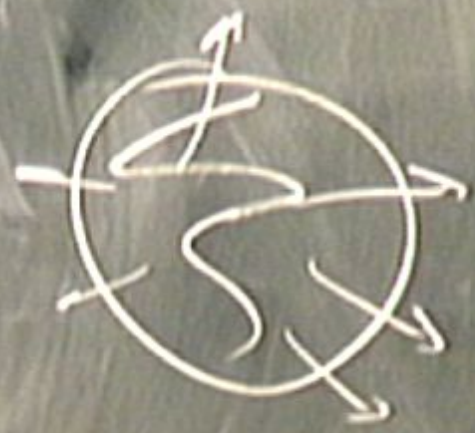
Negative peak imply fluctuations come from outside horizon

Hu & White (1996)  
 Spergel & MZ (1997)  
 Pieris et al WMAP (2003)

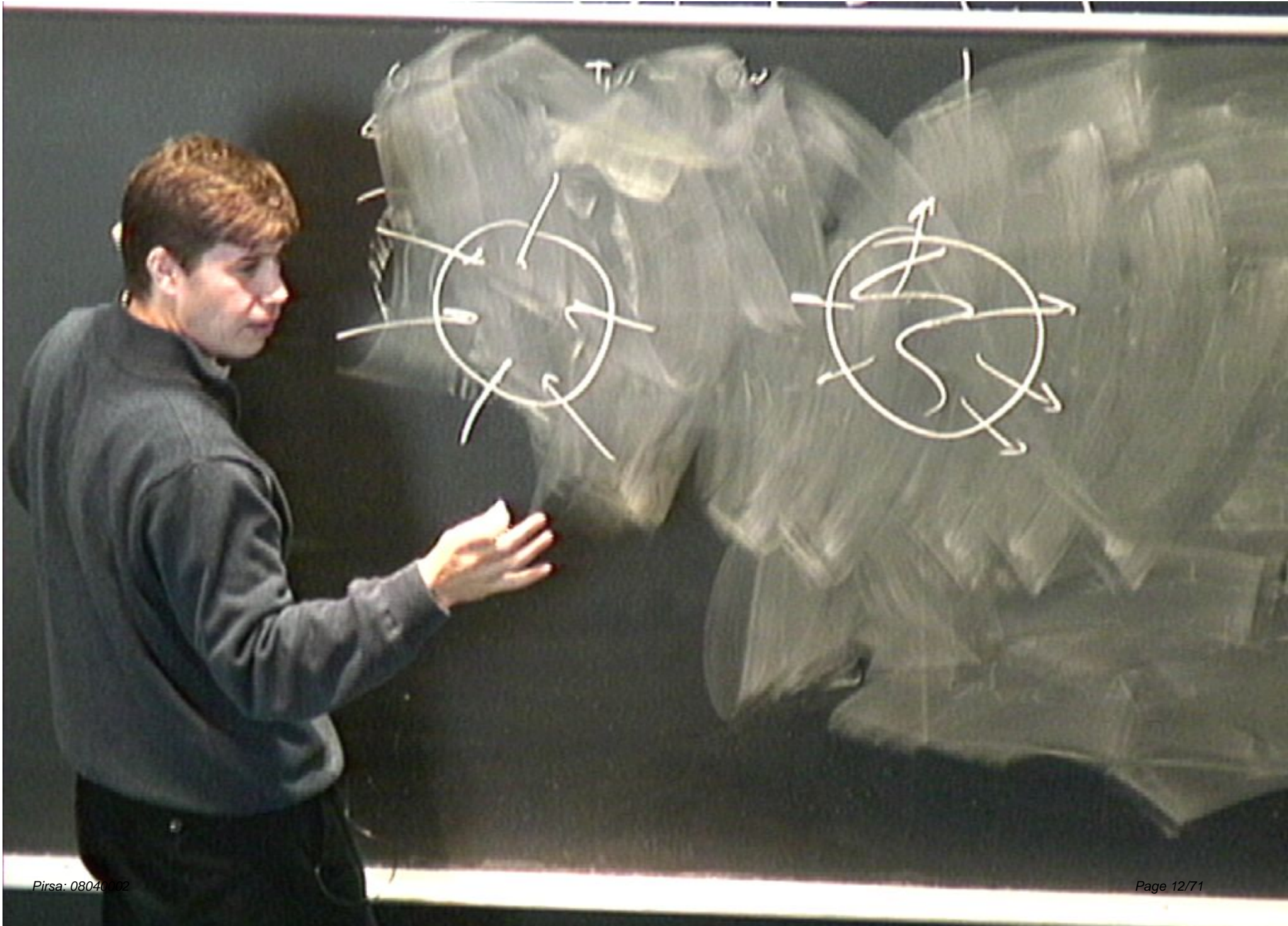


$$Q \rightarrow F_{\mu} \rightarrow T_{\mu} \rightarrow G_{\mu}$$

$$F \rightarrow Q$$










# The shape of the power spectrum

Standard inflation predicts departures from scale invariance, although small

$$n - 1 \sim 1/N_{efold} \sim 0.02$$

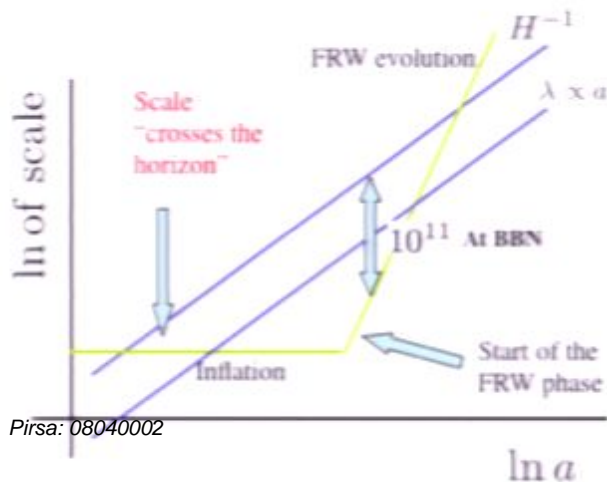
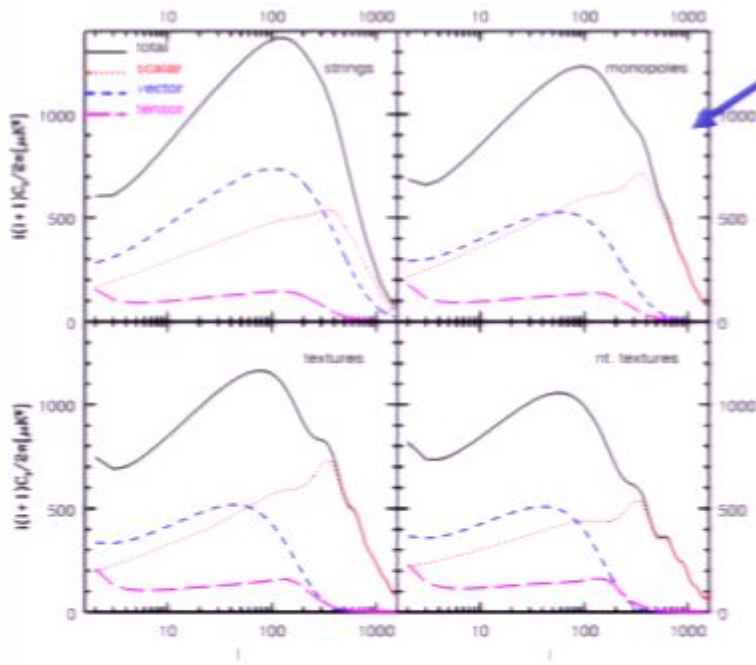
ratio of the size of fluctuations at two different scales, with wave-numbers  $k_1$  and  $k_2$


$$\frac{P(k_1)}{P(k_2)} = \left(\frac{k_1}{k_2}\right)^{n-1}$$

$$(400)^{n-1} \sim 1.3 \quad (1500)^{n-1} \sim 1.4 \quad (10^4)^{n-1} \sim 1.6$$

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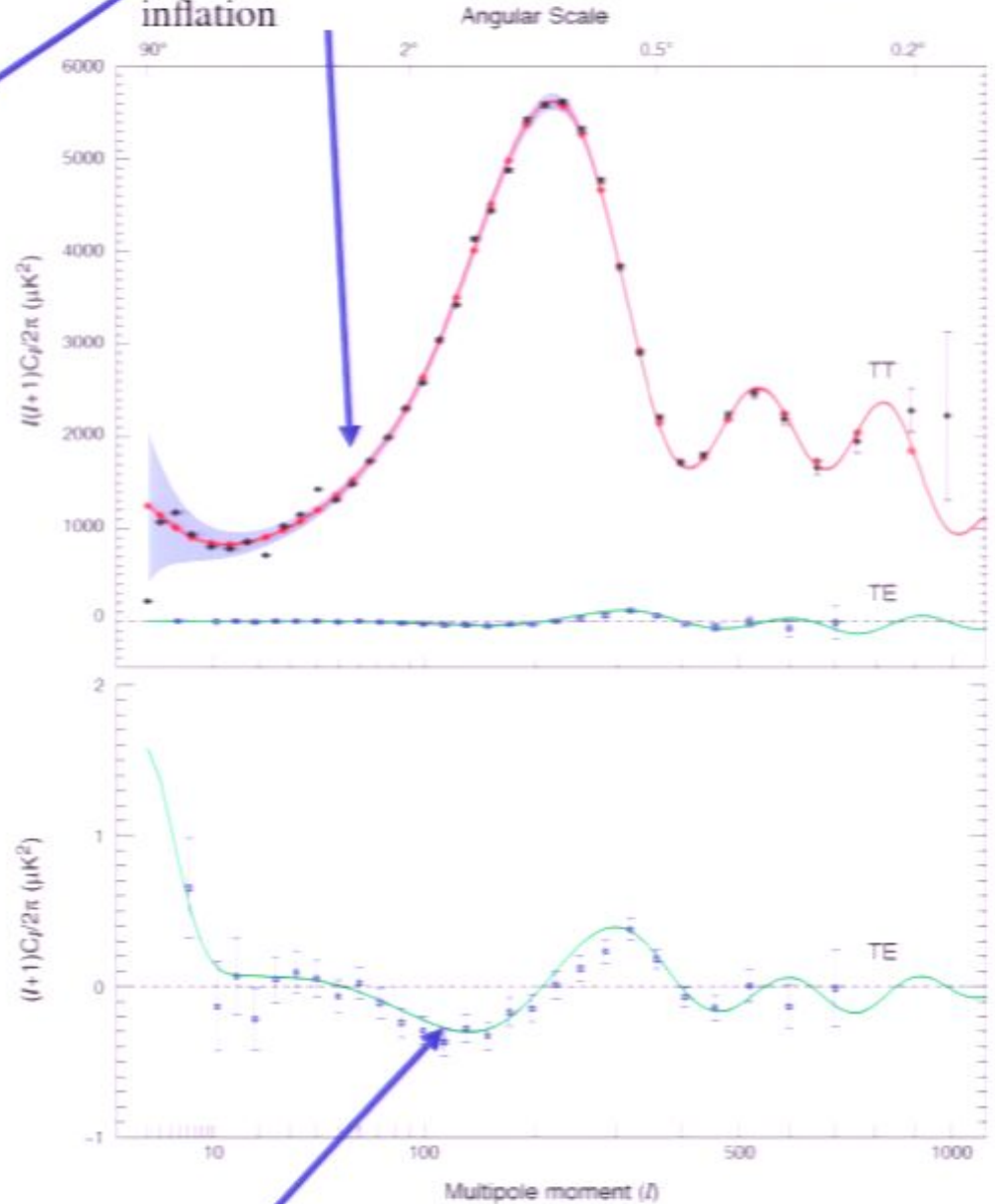
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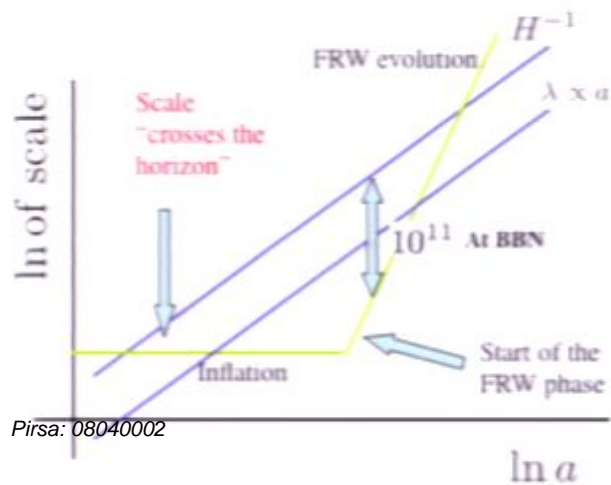
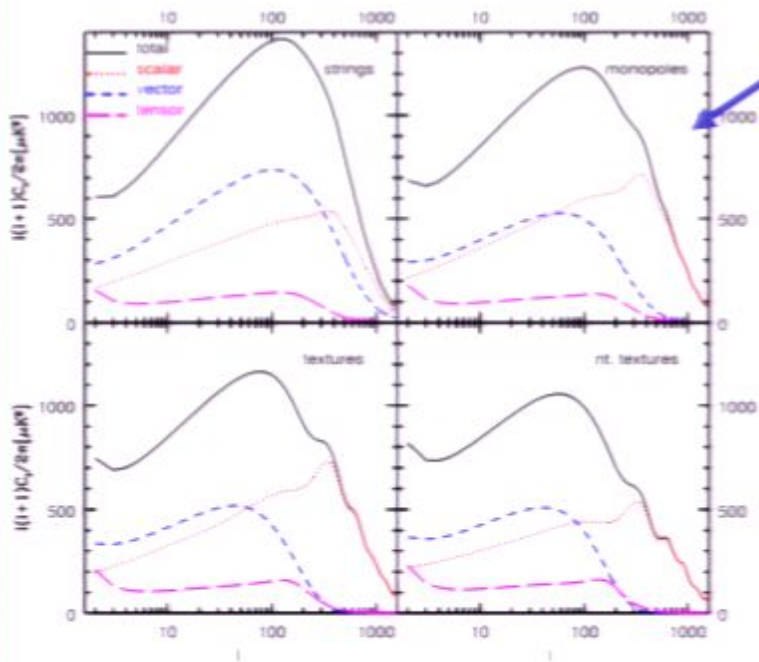
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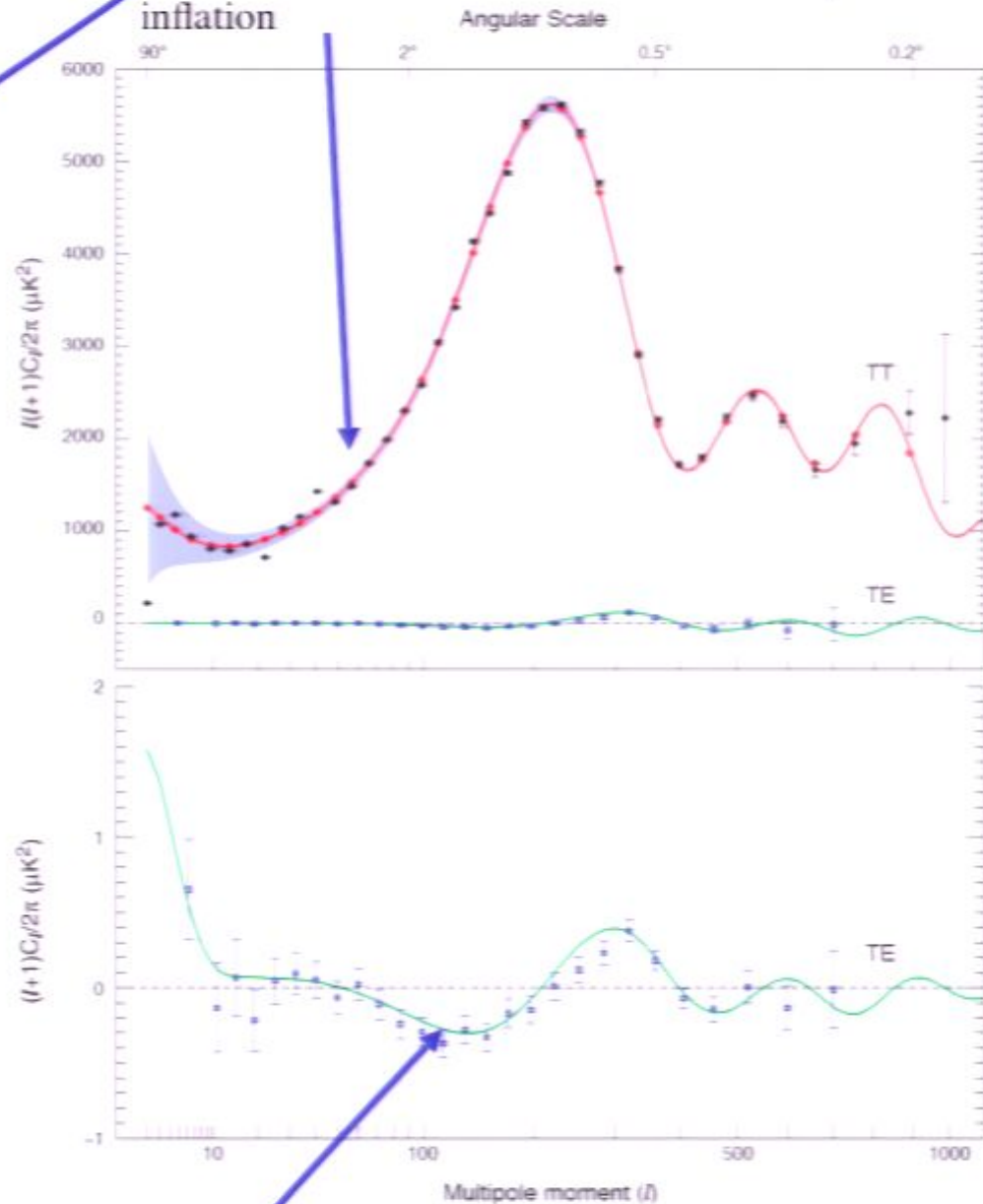
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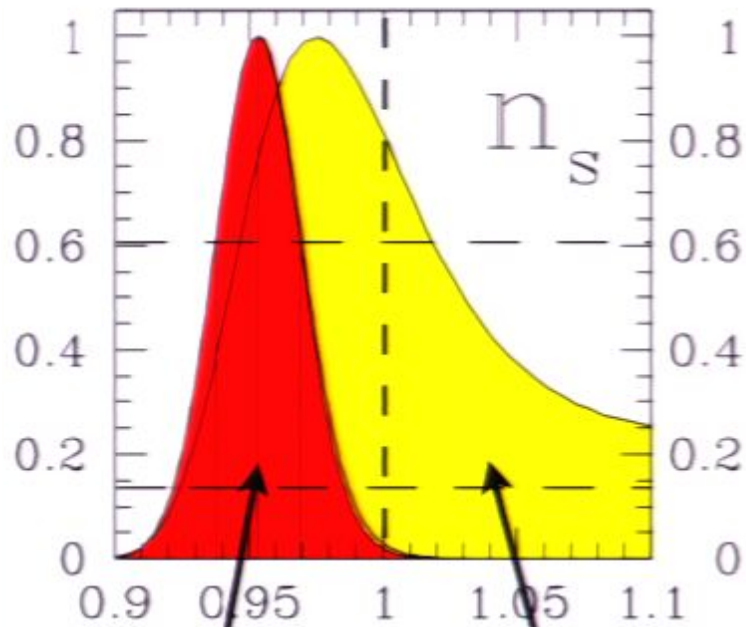
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## Departures from scale invariance after 3 yrs of WMAP

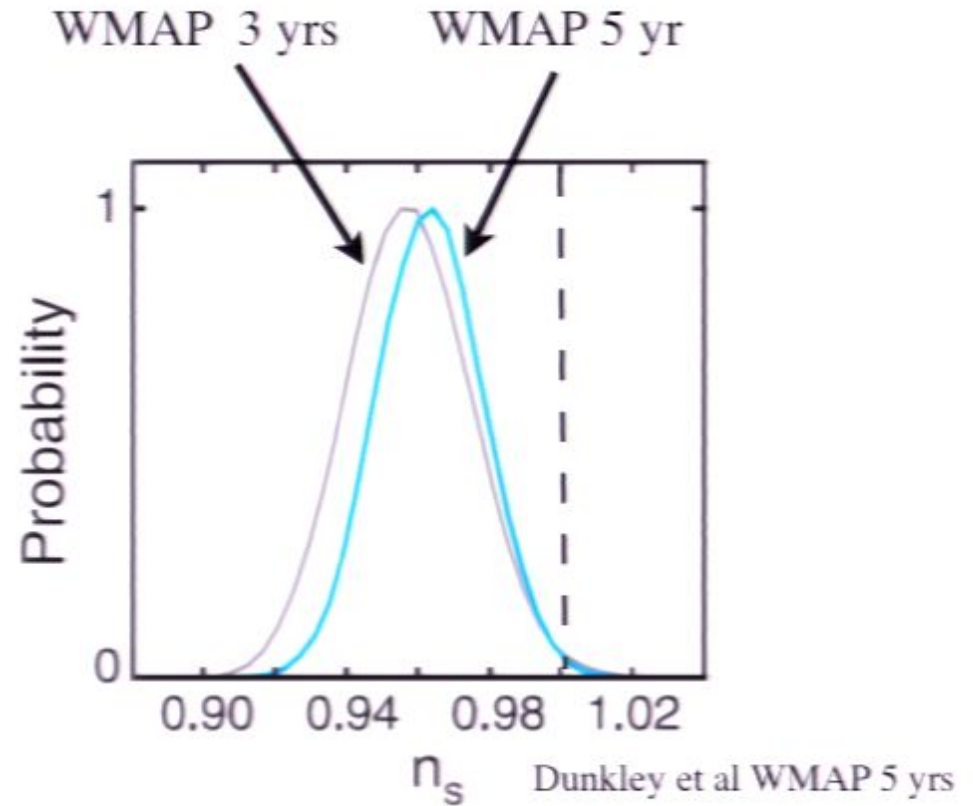


WMAP 3 yrs  
WMAP 1 yr

$n_s$	WMAP	$0.954^{+0.017}_{-0.016}$
	+SDSS	$0.953^{+0.016}_{-0.016}$

Pirsa: 08040002

Tegmark et al astro-ph/0608632



Several technical issues remain unsolved that appear to change the significance of this result (eg. different treatments of point sources can make  $n < 1$  be only a 1.8  $\sigma$  result 0710.1873).

There are tensions between WMAP other probes such as the Lyman  $\alpha$  forest, etc.

# The future is even brighter: Ground base experiments, Planck

## PLANCK

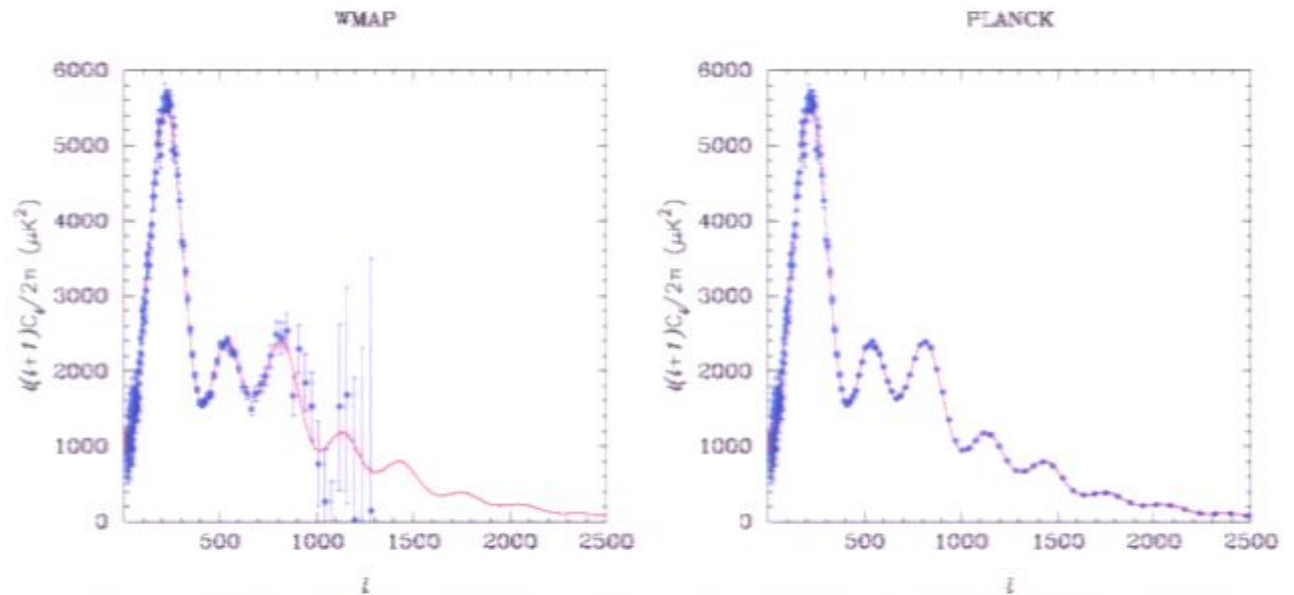
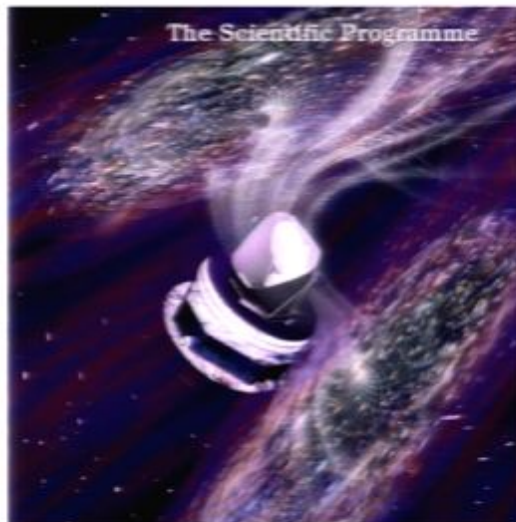
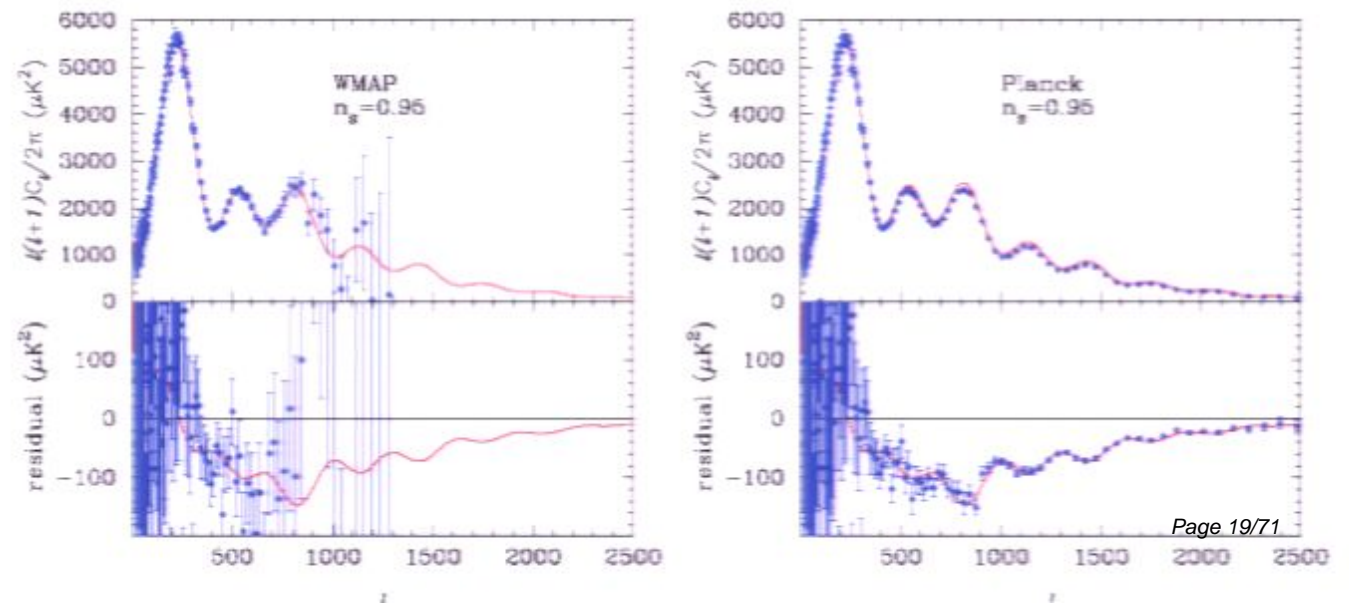


FIG 2.8.—The left panel shows a realisation of the CMB power spectrum of the concordance  $\Lambda$ CDM model (red line) after 4 years of WMAP observations. The right panel shows the same realisation observed with the sensitivity and angular resolution of *Planck*.





## The energy scale of inflation: gravitational waves

Inflation predicts the presence of a stochastic  
Background of Gravitational Waves

$$h_{ij} \sim \frac{H}{M_{PL}}$$

If  $\frac{H}{M_{PL}} \sim 10^{-5}$   Comparable to  
density perturbations

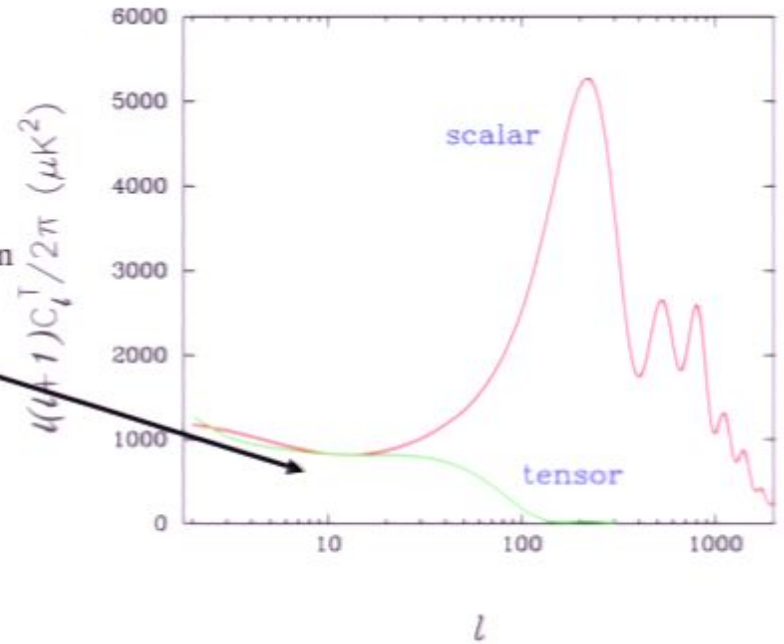
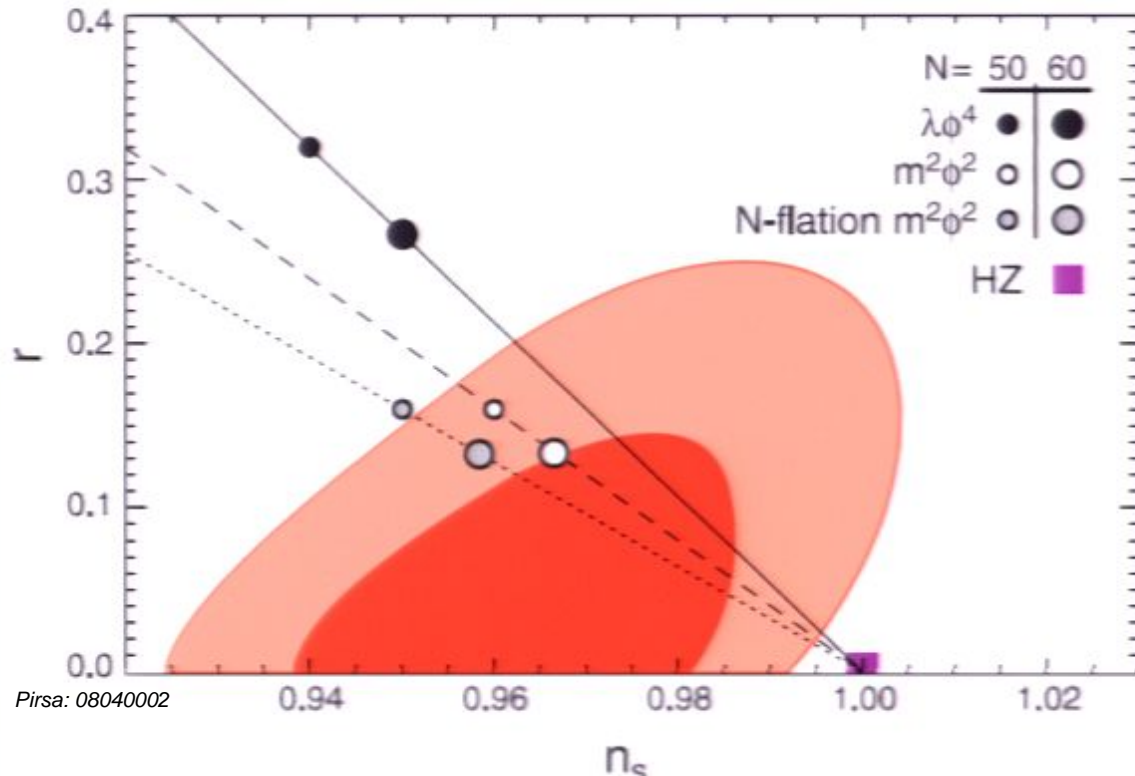
Directly measure the expansion rate during Inflation.

This measurement has taken a greater significance now that it  
appears that GW might not be observable in “string-inflation”.



## Current constraints from WMAP 5 yrs

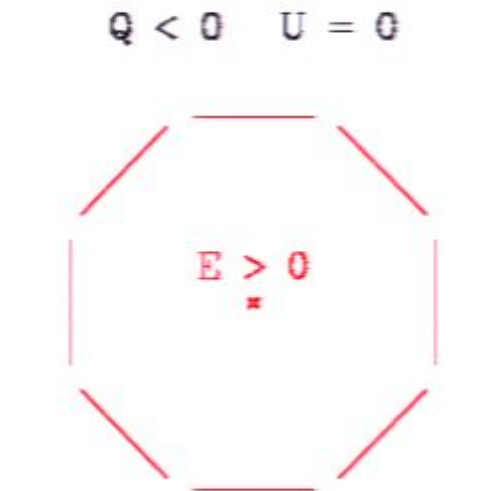
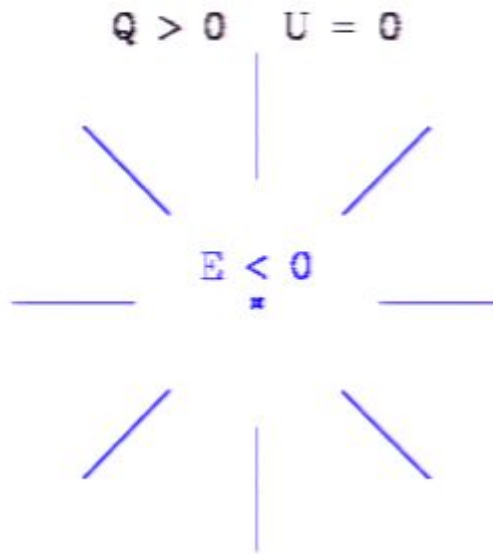
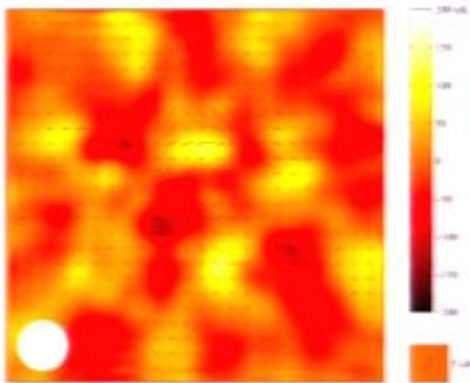
Current constraints come from fitting the shape of the temperature power spectrum.



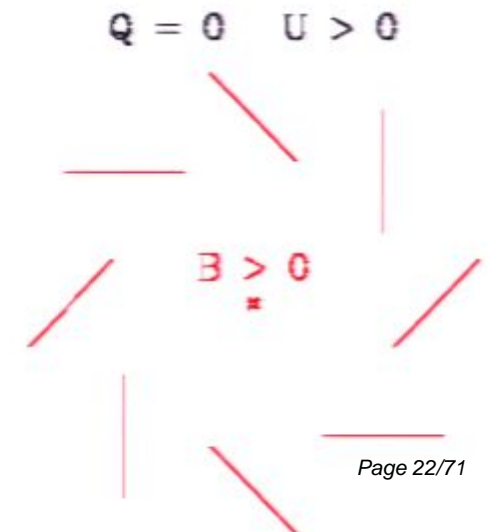
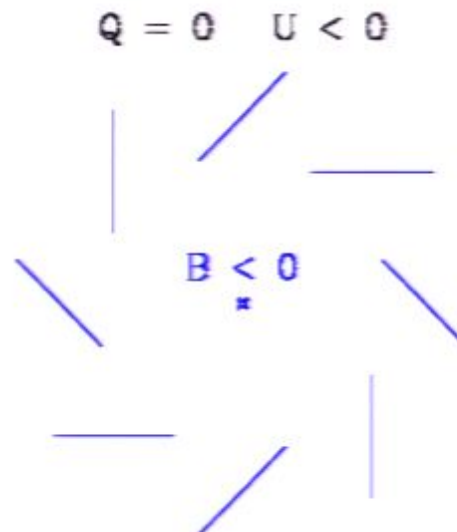
CMB probes a narrow range of energy scales

# The future of GW constraints: Polarization

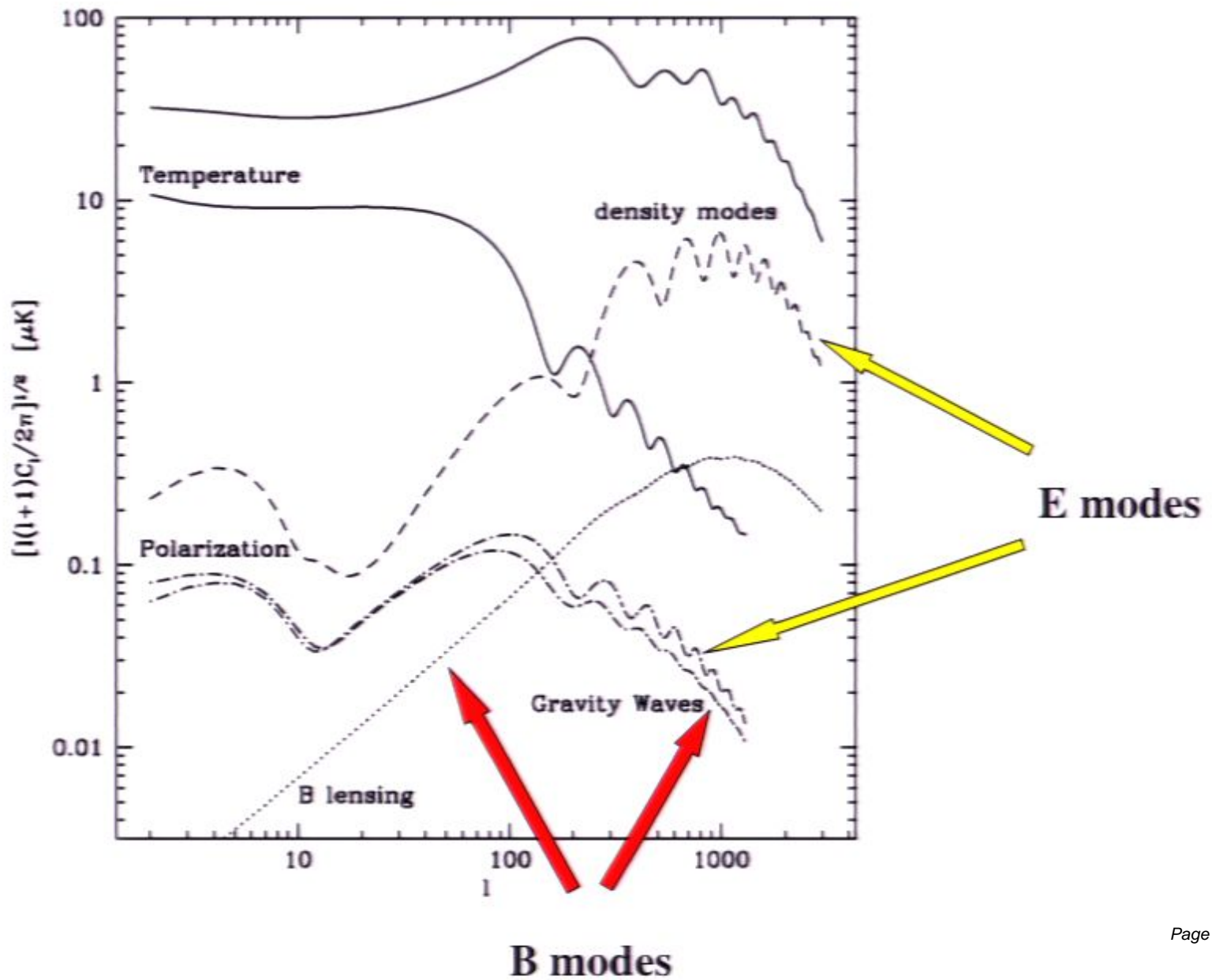
Density pert.  
&  
Gravity Waves



Gravity  
Waves



## Power spectra of E & B





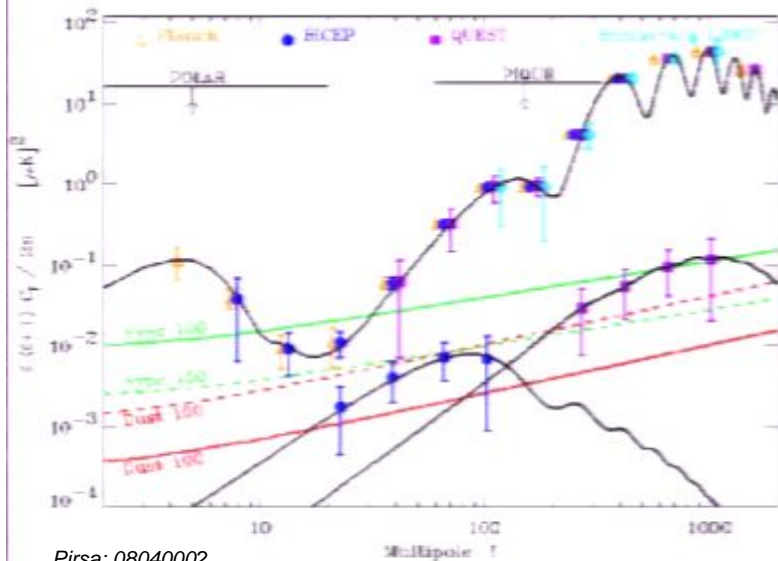
# Which experiment will first constrain Gravity waves using B modes?

Examples

Running

Future

Bicep



Pirsa: 08040002

1500 sq deg., 42 arcmin res., 2 Freq., 100, 150 GHz

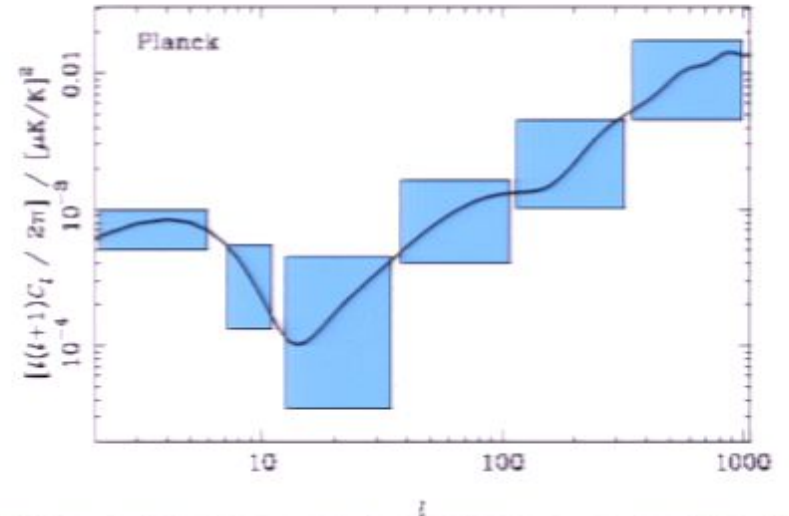
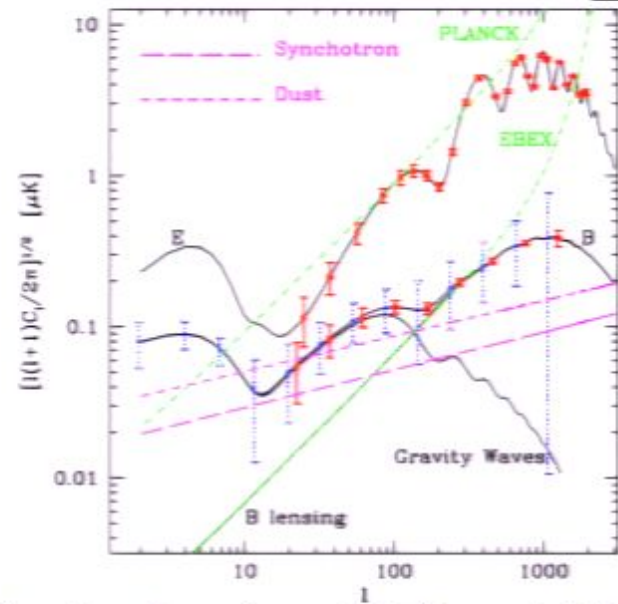


FIG 2.17.—Forecasts for the  $\pm 1\sigma$  errors on the  $B$ -mode polarization power spectrum  $C_l^B$  from Planck (for  $r = 0.1$  and  $\tau = 0.17$ ). Above  $l \sim 150$  the primary spectrum is swamped by weak gravitational lensing of the  $E$ -polarization produced by the dominant scalar perturbations. The cosmological model, and the assumptions about instrument characterization, are the same as in Figure 2.13.

EBEX

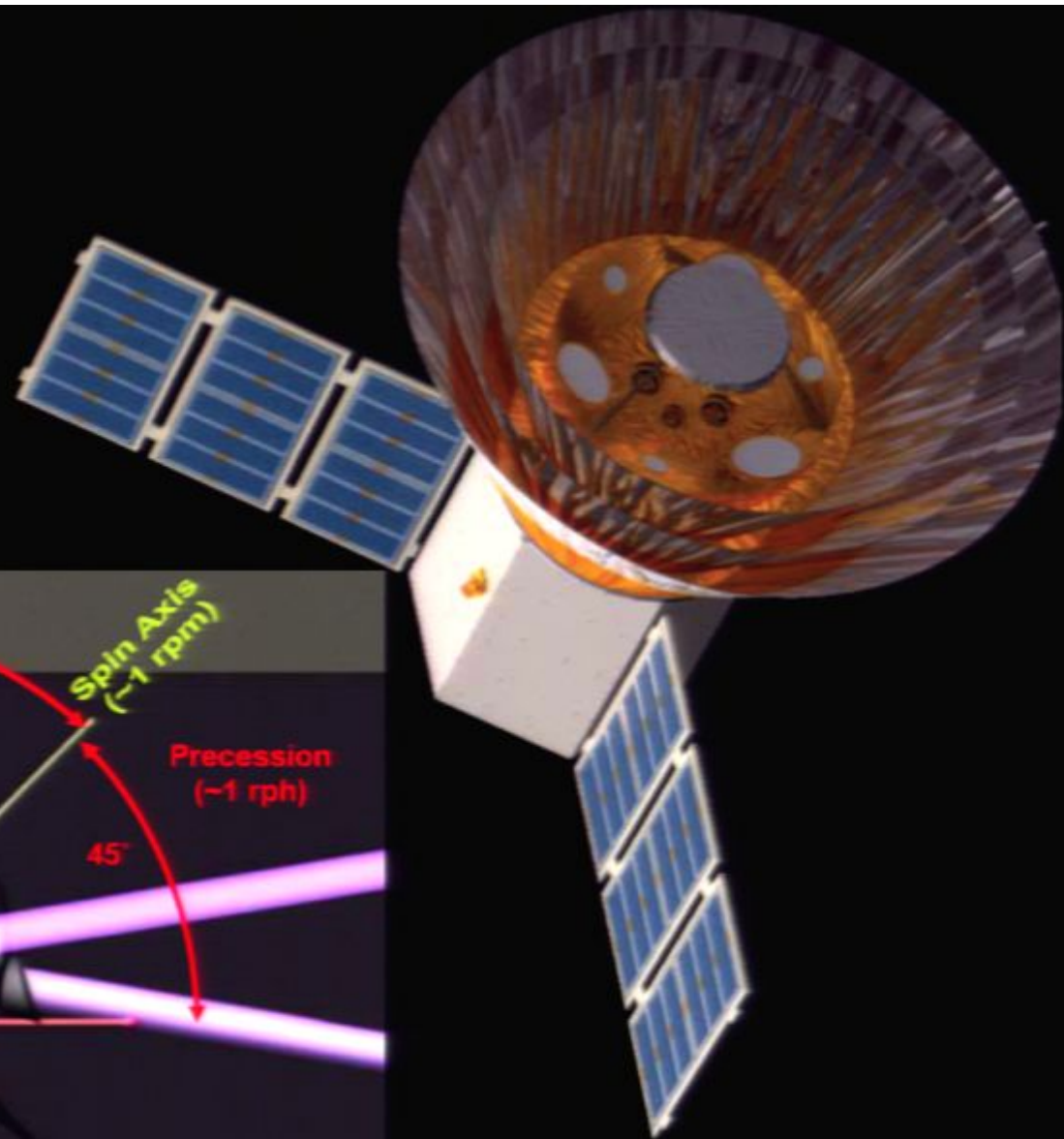


350 sq deg., 8 arcmin res., 0.7  $\mu$ K per pixel, 3 Freq., 150,

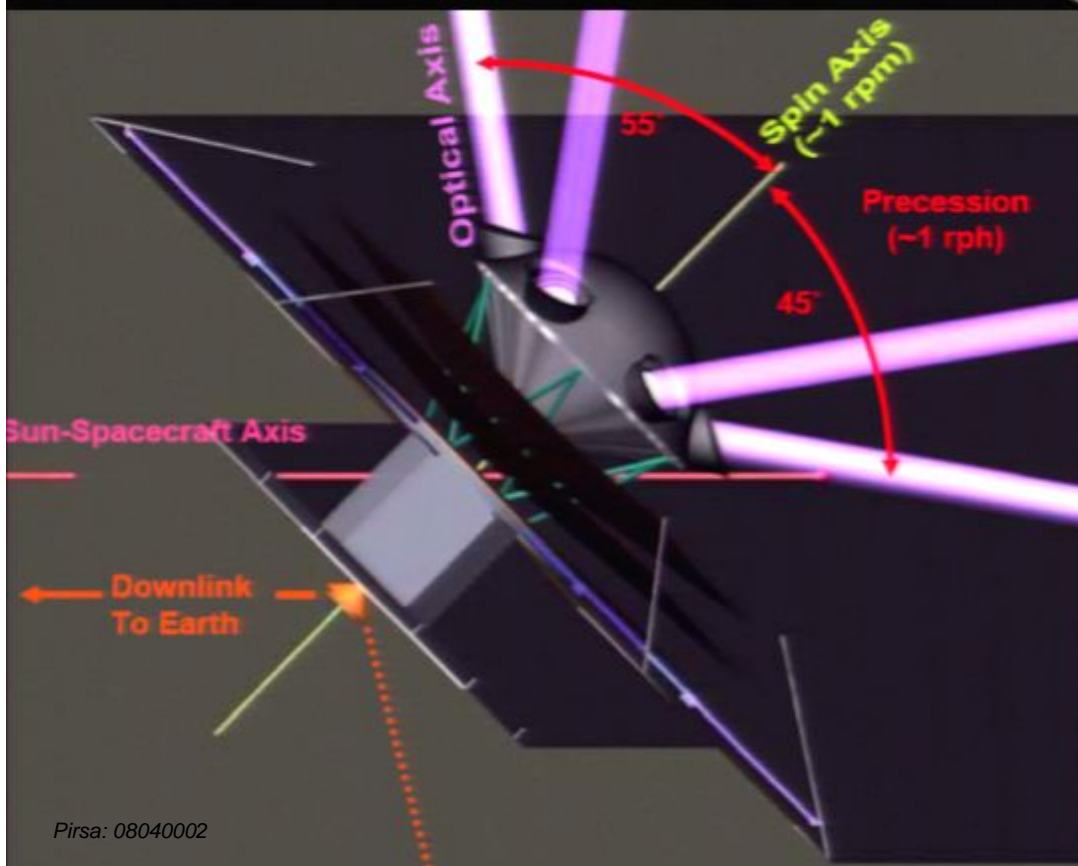
250 & 420 GHz

# Future

## Epic



# CMBPOL



# The statistics of the seeds

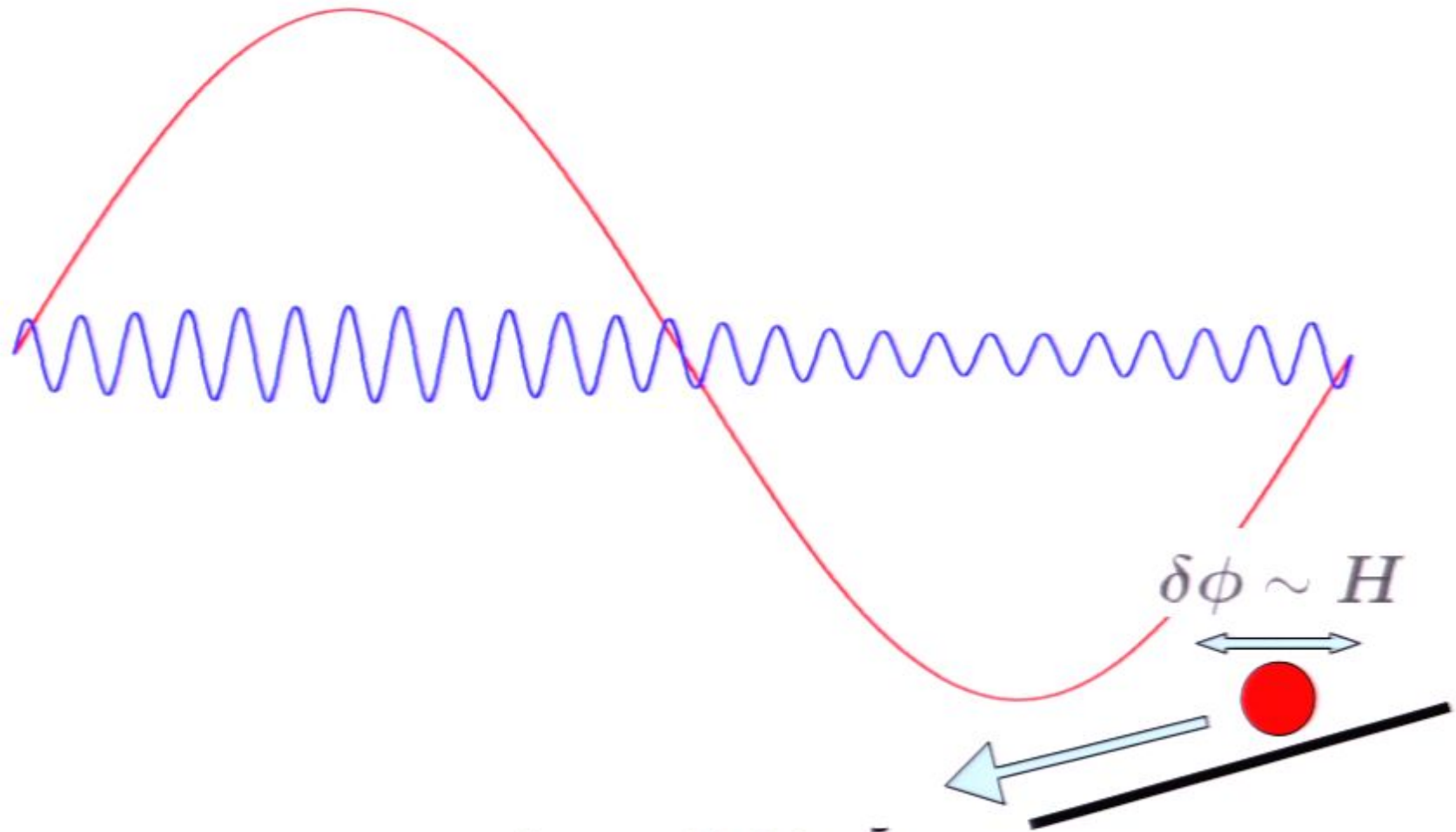
The simplest inflationary models predict very small (unobservable) departures from Gaussianity

Secondary anisotropies and residual foregrounds also lead to departures from Gaussianity.



## The Size of non-Gaussianities :

There is a direct connection between the departures from scale invariance and the three point function



$$\text{Non-Gaussianities} \sim (n - 1)10^{-5}$$

## Level of non-G in standard inflation

- $\frac{\langle \delta^3 \rangle}{\langle \delta^2 \rangle^{3/2}} \sim (n - 1)\delta \sim 10^{-2} \times 10^{-5}$

- The ultimate level using CMB is around  $10^{-4}$

Komatsu & Spergel

- Current constraint is around  $10^{-2}$

WMAP first year data

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WMAP first year data



In models of inflation where the density perturbations we see were produced by fluctuations in a field other than the inflaton, and get converted into density fluctuations at reheating or later, non-Gaussianities of the local form are common.

$$\frac{\delta\rho}{\rho} = g + f_{NL}(g^2 - \langle g^2 \rangle) + \dots$$

This is generic to all models where there is more than one field (or clock) so that it makes sense to talk about evolution outside the horizon (curvaton, multi-field models, variable decay rate).

$$f_{NL} \sim 10$$

## Expected level

$$\frac{\langle \delta^3 \rangle}{\langle \delta^2 \rangle^{3/2}} \sim \frac{V^{(3)}}{H}$$

$$\frac{(\partial\phi)^4}{M^4} \rightarrow \frac{\dot{\phi}H^3}{HM^4} \sim f\delta \quad f \sim \frac{\dot{\phi}^2}{M^4}$$

$f$  can be  $\gg (n-1)$  however if  $f$  is close to one the higher derivative operators will modify the homogeneous solution

$$\mathcal{L} = F\left[\frac{(\partial\phi)^2}{M^4}\right]$$

For detectable non-G the full function  $F$  is important for the dynamics.

Strongly coupled dynamics not traditional slow roll.

There are just two different things to look for, two distinct shapes of the three point function.

Multiple fields:  
evolution outside the  
horizon is meaningful

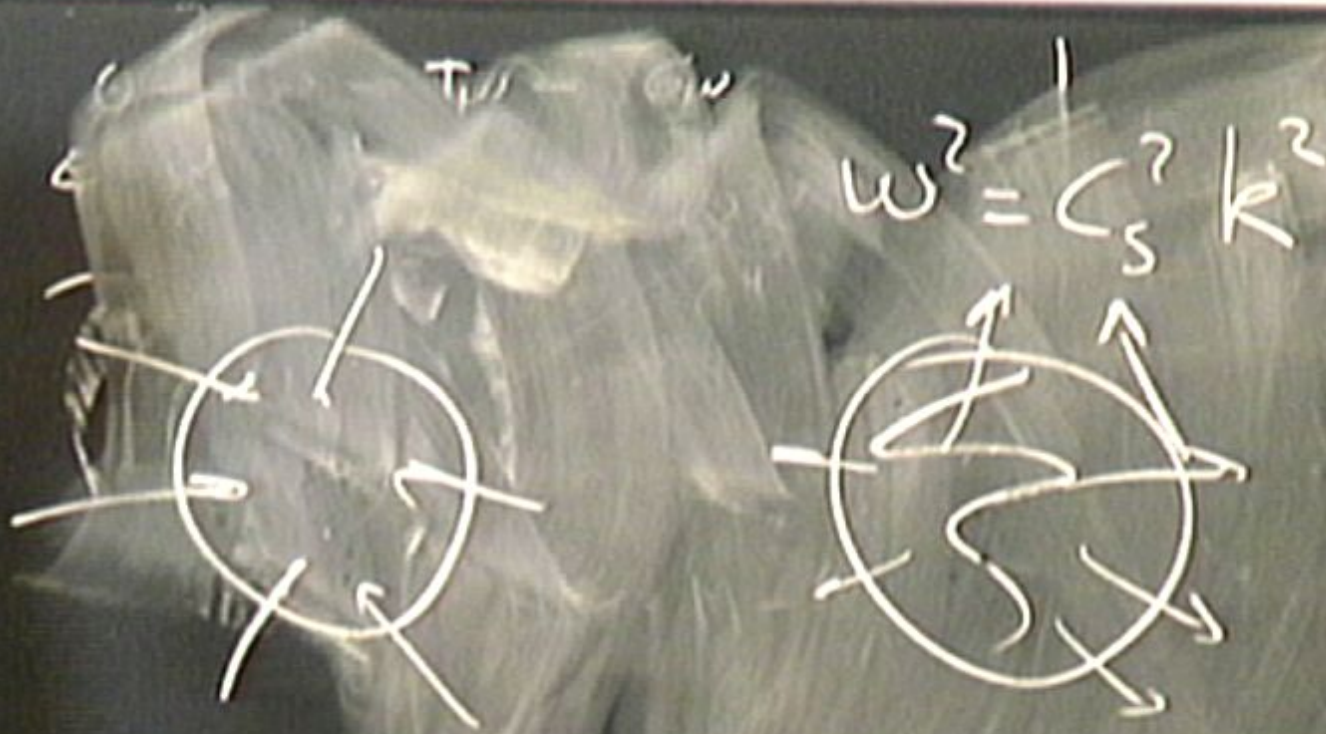
Correlates large  
and small scales

Only one clock but strongly  
coupled dynamics

Only correlates  
similar scales

In general we expect (= “can manage to get”) departures from  
Gaussianity that are on the small side for WMAP but could be seen by  
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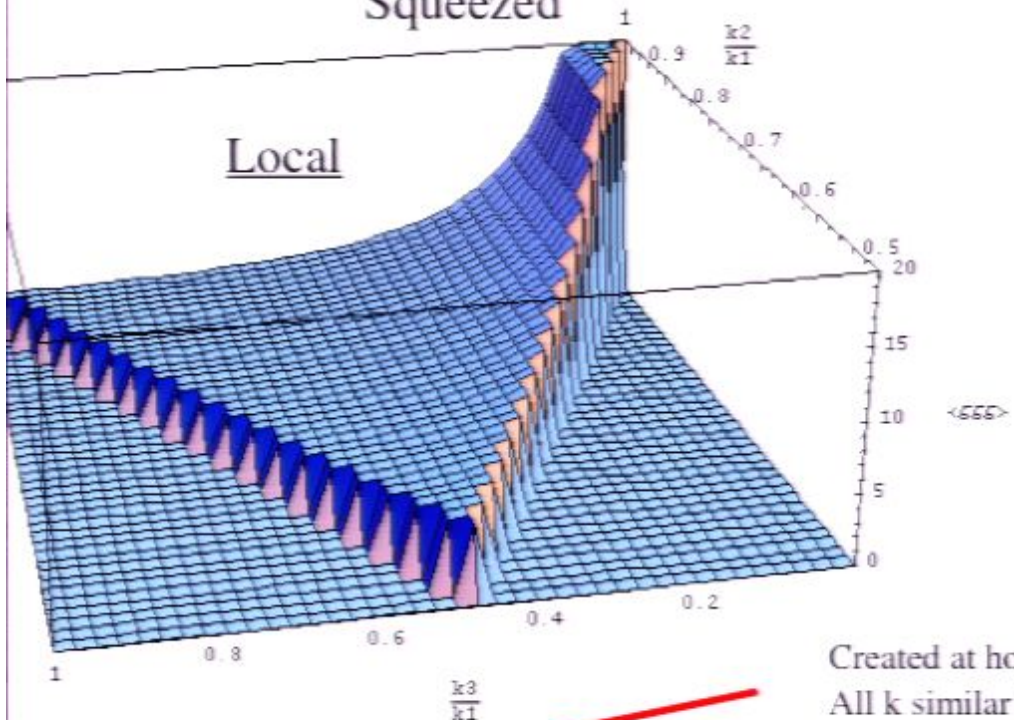
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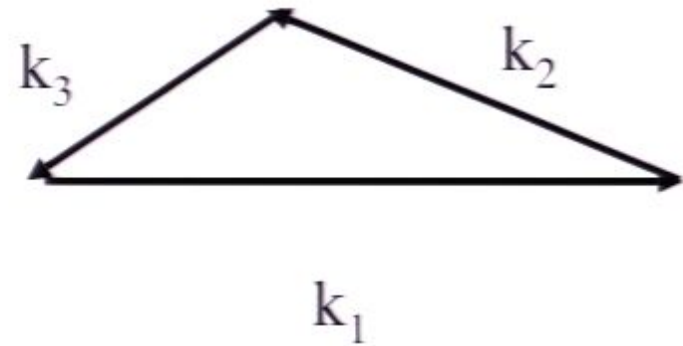


Squeezed

Local



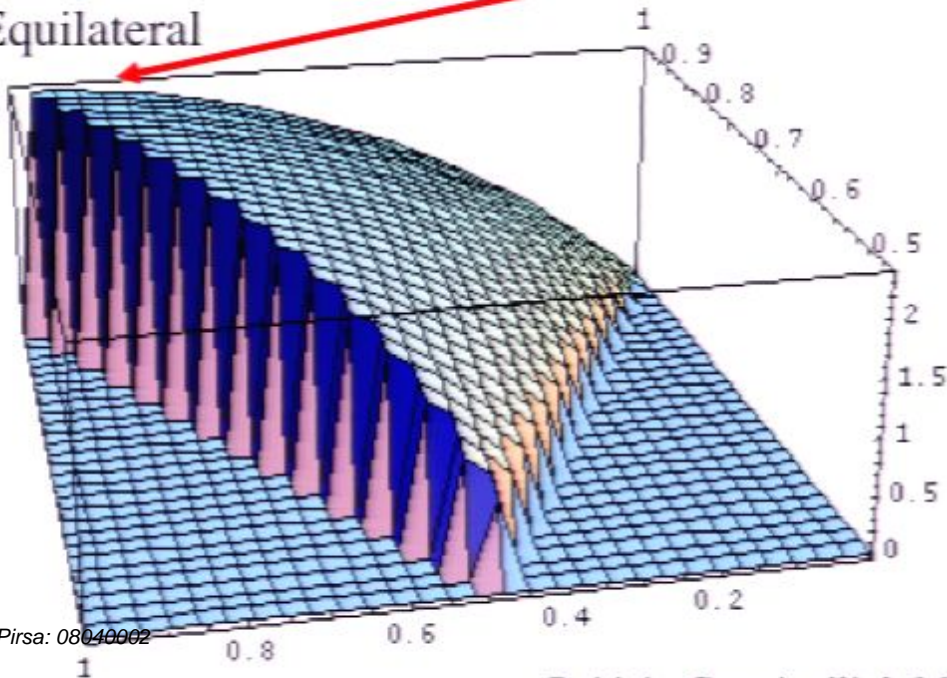
The 3-point function



Created at horizon crossing  
All k similar in order to contribute

$$\left\langle \frac{\delta\rho}{\rho}(k_1) \frac{\delta\rho}{\rho}(k_2) \frac{\delta\rho}{\rho}(k_3) \right\rangle$$

Equilateral



Higher derivative operators



## Searching for the signal

The best data set is the one with the largest number of high signal to noise measurements (pixels, Fourier modes). Constraints go as

$$N_{\text{pix}}^{-1/2}$$

WMAP and future all sky CMB experiments are the most promising.

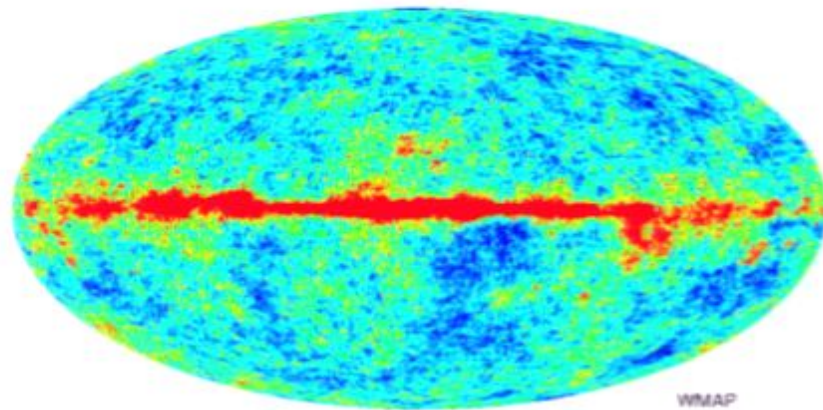
Surveys of hydrogen at high redshift using its 21 cm line could potentially do better.

## The basics of CMB Anisotropies

$$\frac{\delta T}{T} = \phi + \frac{\delta\gamma}{4} + \frac{v_r}{c}$$

All 3 effects have the same origin

$$P \sim \lambda_T \nabla v$$



Tight Coupling

14 Gpc

Free Streaming

Observer Today

## Summary so far

Experiments are beginning to constrain the characteristics of the primordial fluctuations at levels that are interesting for early universe physics.

- Small departures from scale invariance
- Stochastic background of GW
- Departures from Gaussianity

The best way to characterize the departures from Gaussianity is the three point function.

Departures come in just two types depending on the physical origin of the effect

The constraining power of a data set is just given by  $N_{\text{pix}}^{1/2}$

The data is consistent with Gaussianity.

Upper limits are close to plausible levels.



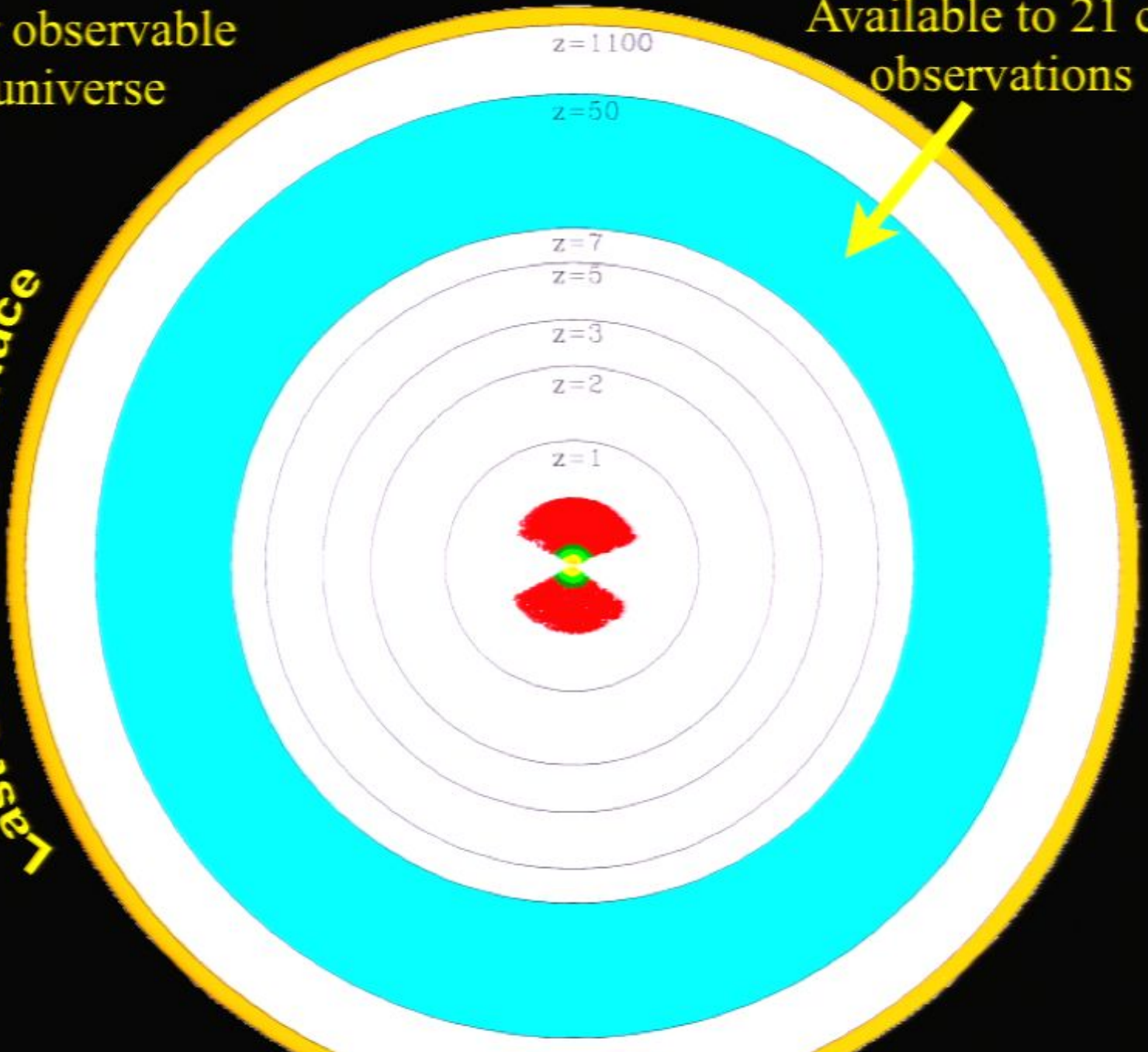
# 21 cm fluctuations: a new window for Cosmology

- 
- The diagram consists of two arrows pointing towards a central list of bullet points. A black arrow labeled 'First Step' points from the top right towards the first bullet point, 'Epoch of Reionization Science'. A red arrow labeled 'Future' points from the bottom right towards the third bullet point, 'Gaussianity'.
- Epoch of Reionization Science
  - Fluctuations on small scales
  - Gaussianity

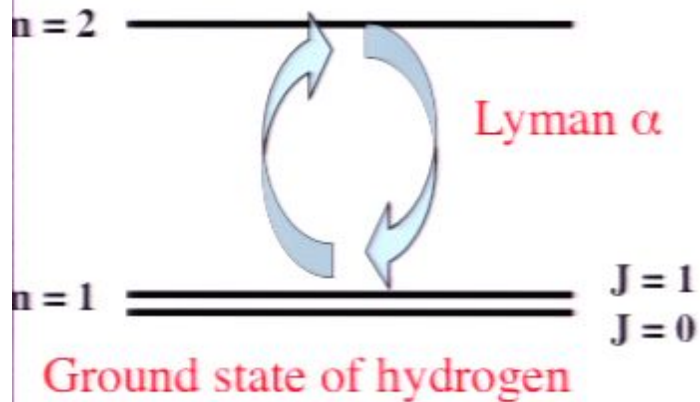
Our observable universe

Available to 21 cm observations

Last scattering surface



## 21 cm transition: a new window



$$\frac{n_1}{n_2} = \frac{g_1}{g_2} e^{-T_*/T_s}$$

$$T_* = 0.068 \text{ K}$$

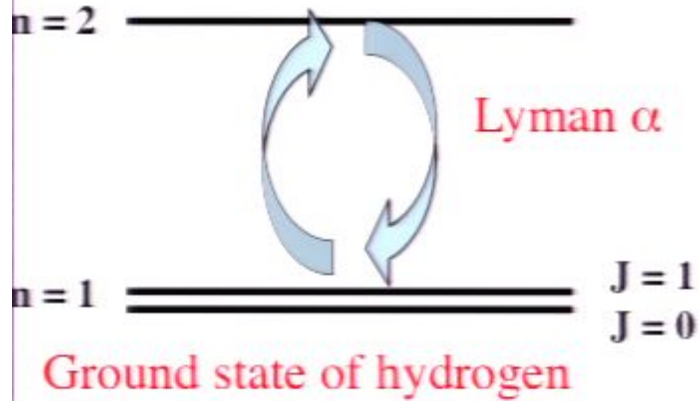
Before reionization we can observe hydrogen using the 21 cm line

### Coupling mechanisms

- 21 cm processes  $T_{CMB}$
- Collisions  $T_K$
- Lyman  $\alpha$   $T_K$

$T_s > T_{CMB}$  Emission  
 $T_s < T_{CMB}$  Absorption

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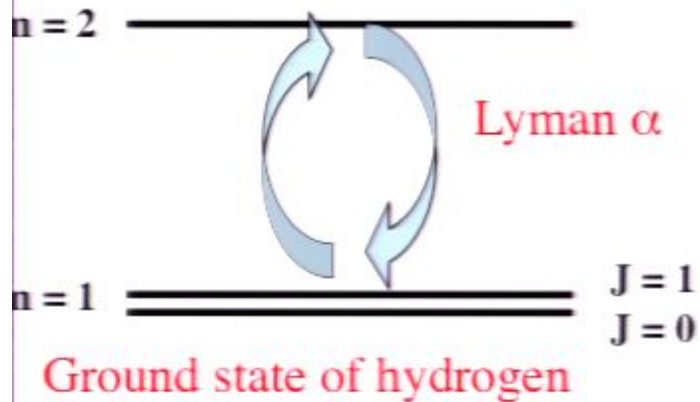
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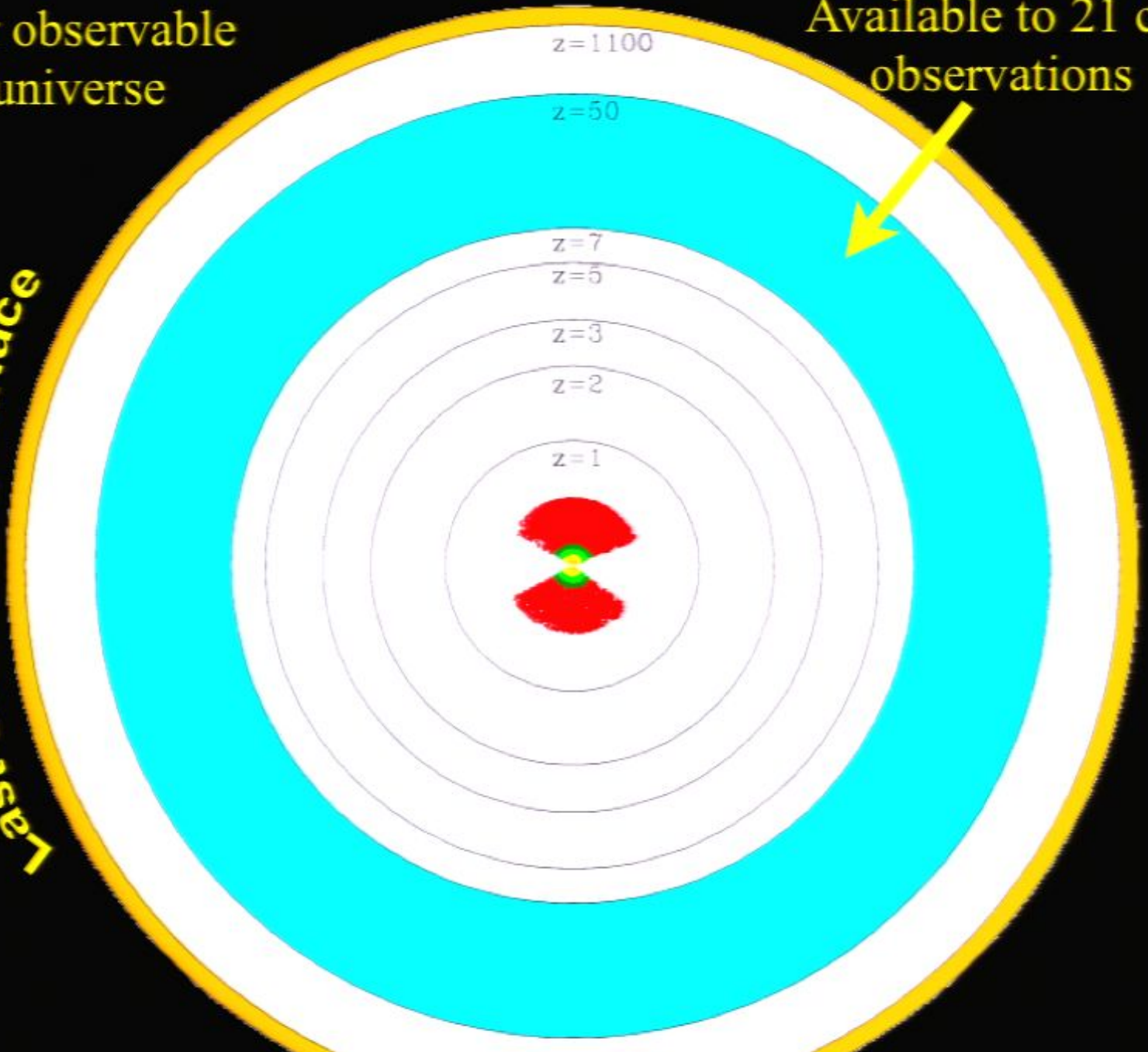
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Our observable universe

Available to 21 cm observations

Last scattering surface



## Brightness fluctuations

Fluctuation in the Neutral fraction

$$\psi = x_H(1 + \delta) \left( \frac{T_S - T_{CMB}}{T_S} \right)$$

Fluctuation in the density

Fluctuation in the temperature:

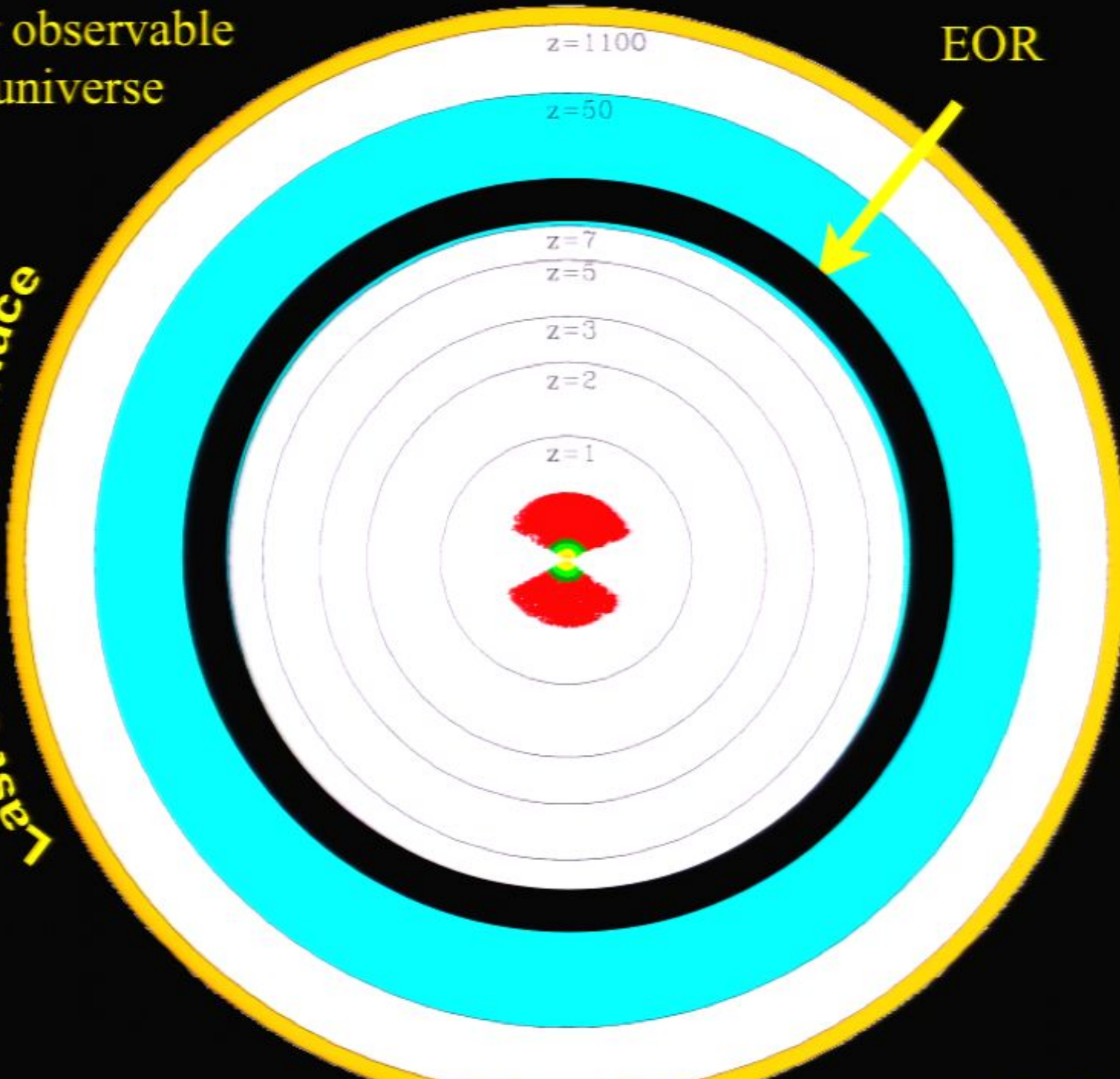
Unimportant for  $T_s \gg T_{cmb}$  (reionization)



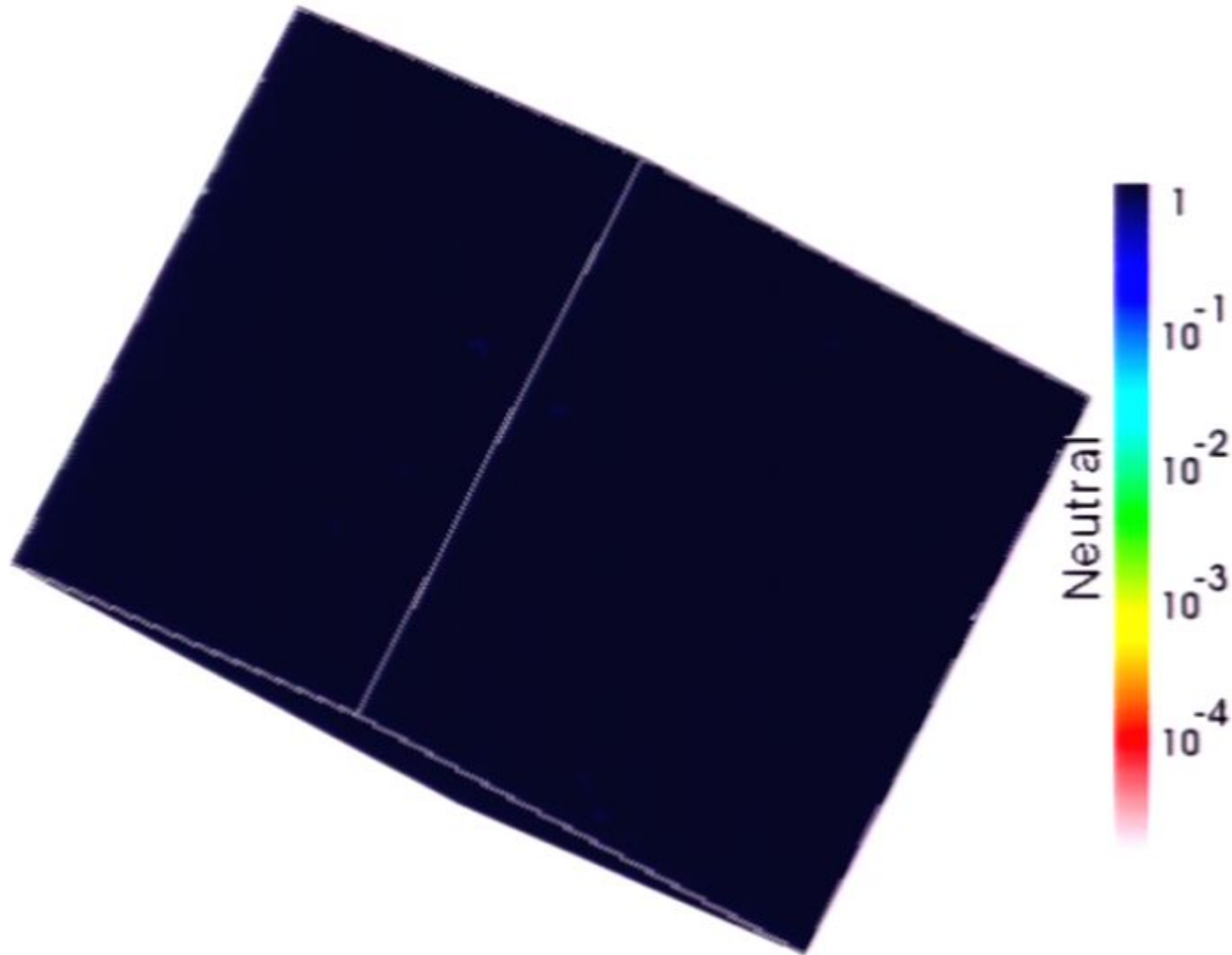
Our observable universe

EOR

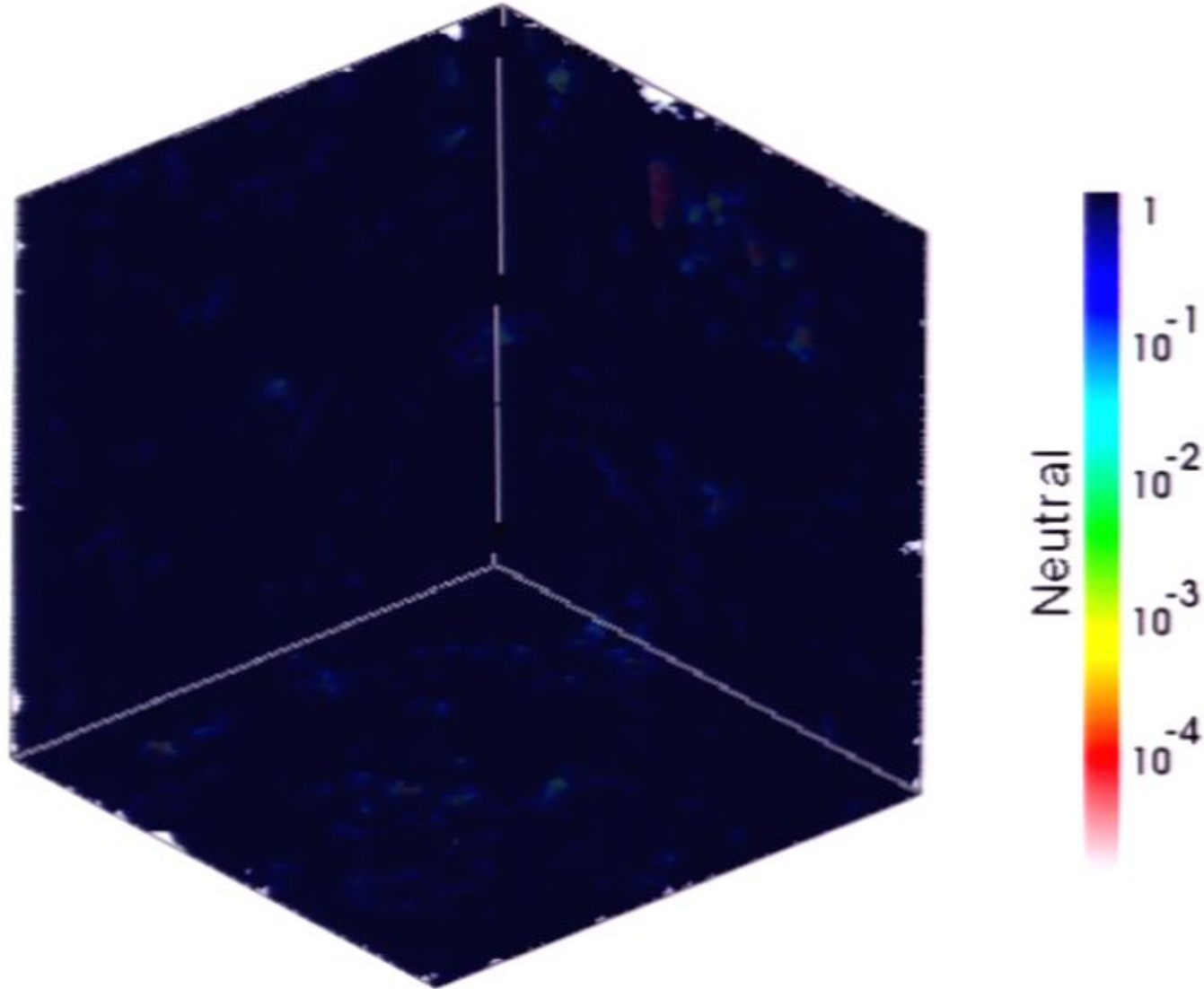
Last scattering surface



## The epoch of reionization

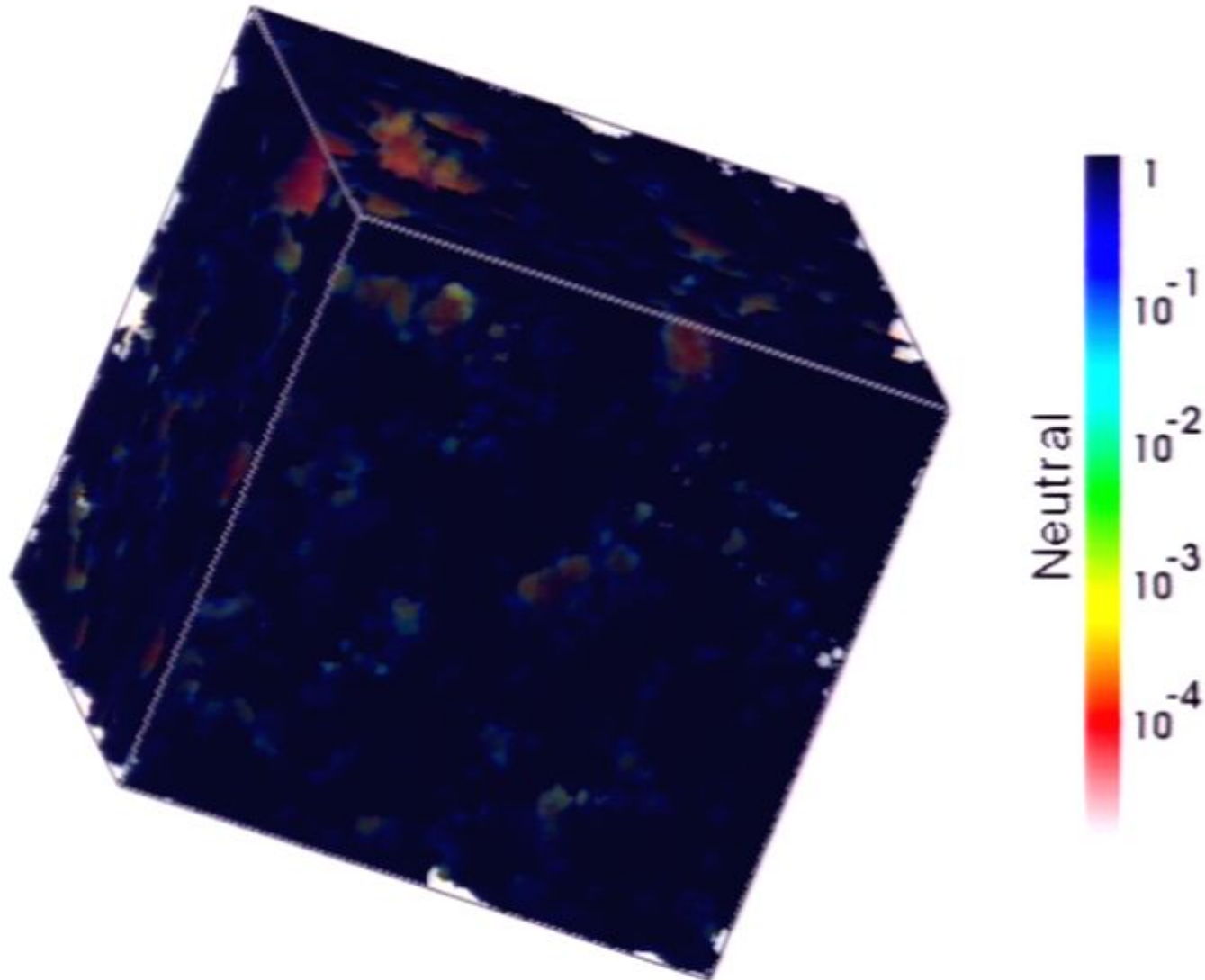


## The epoch of reionization

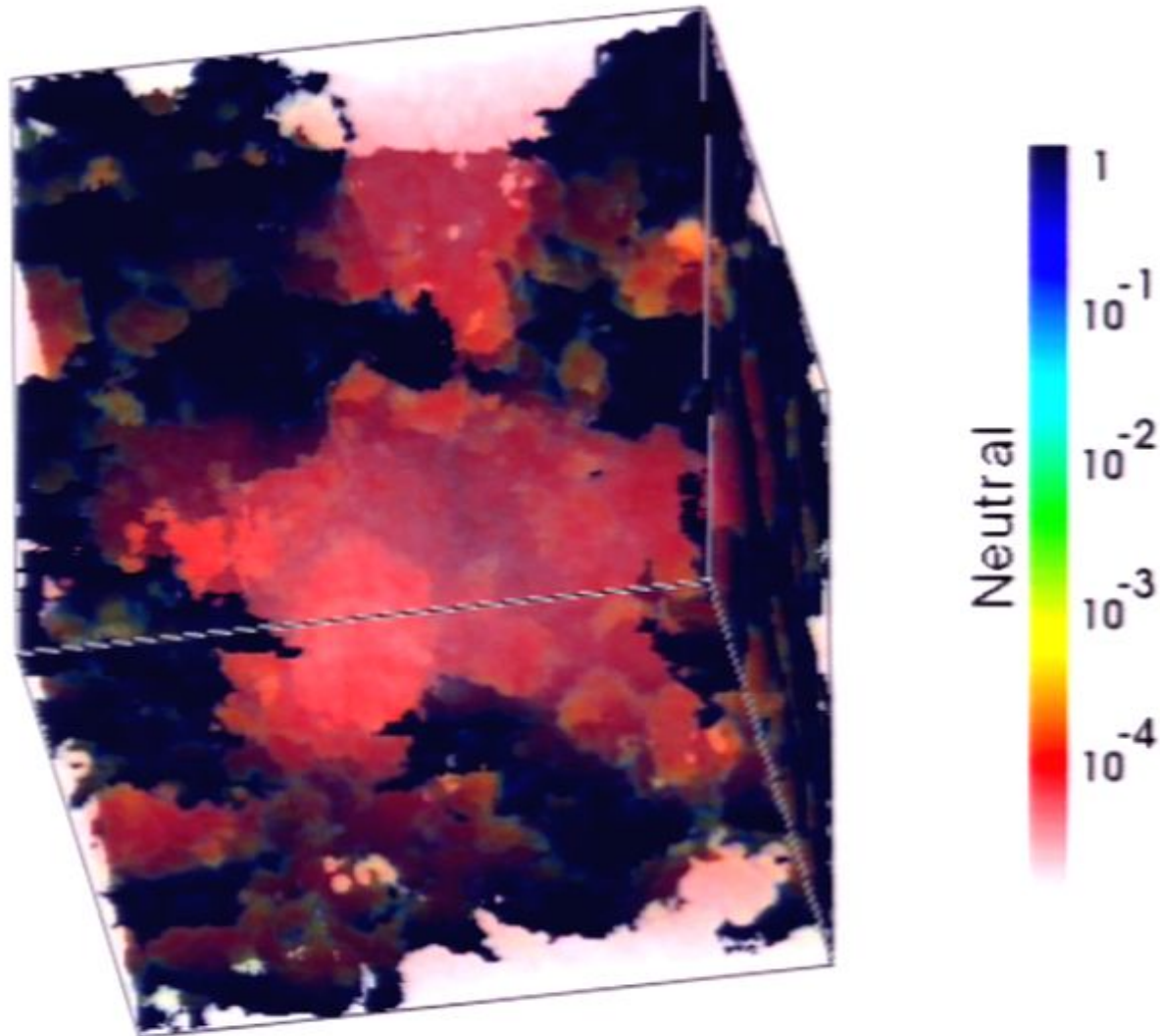




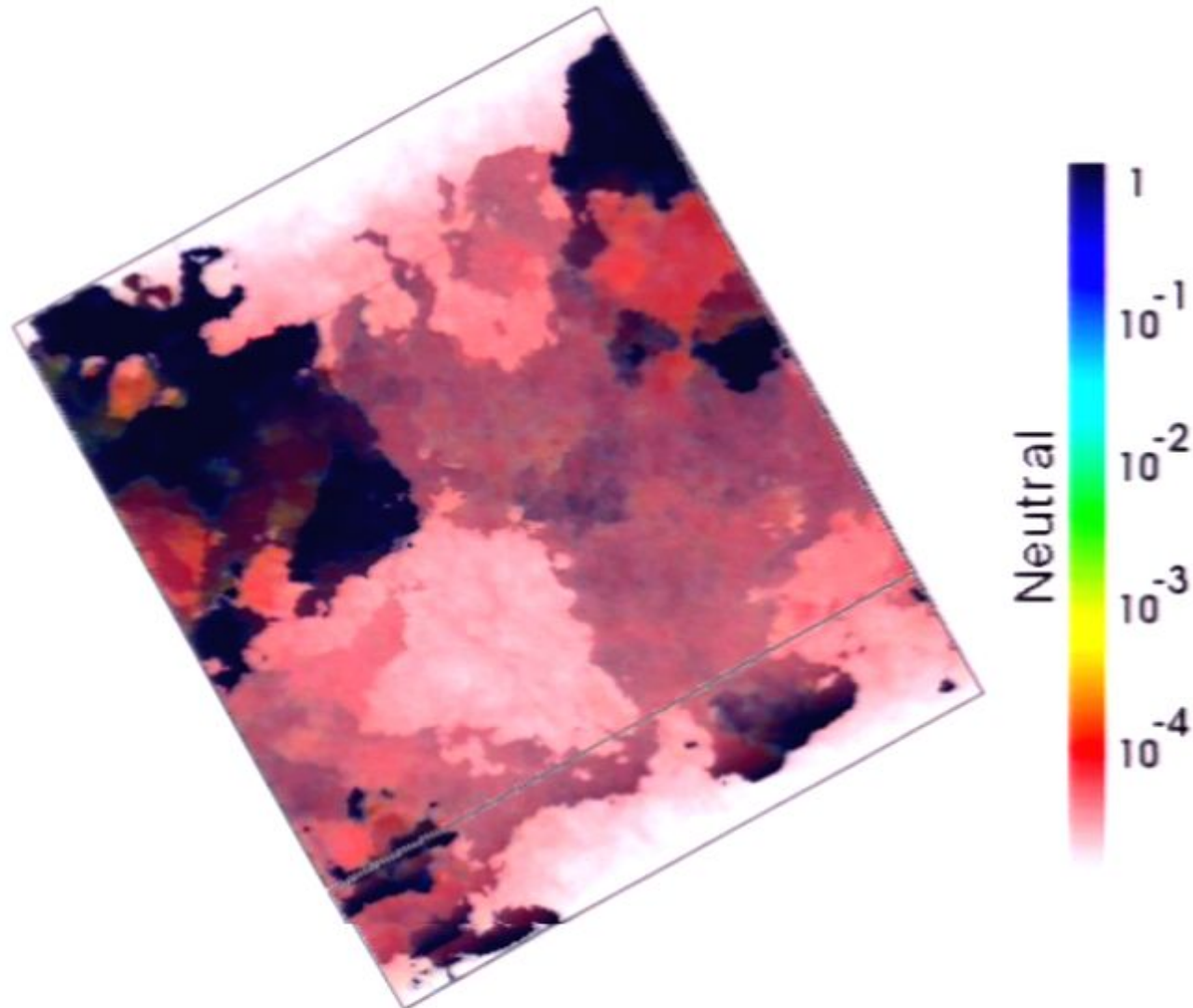
## The epoch of reionization



## The epoch of reionization

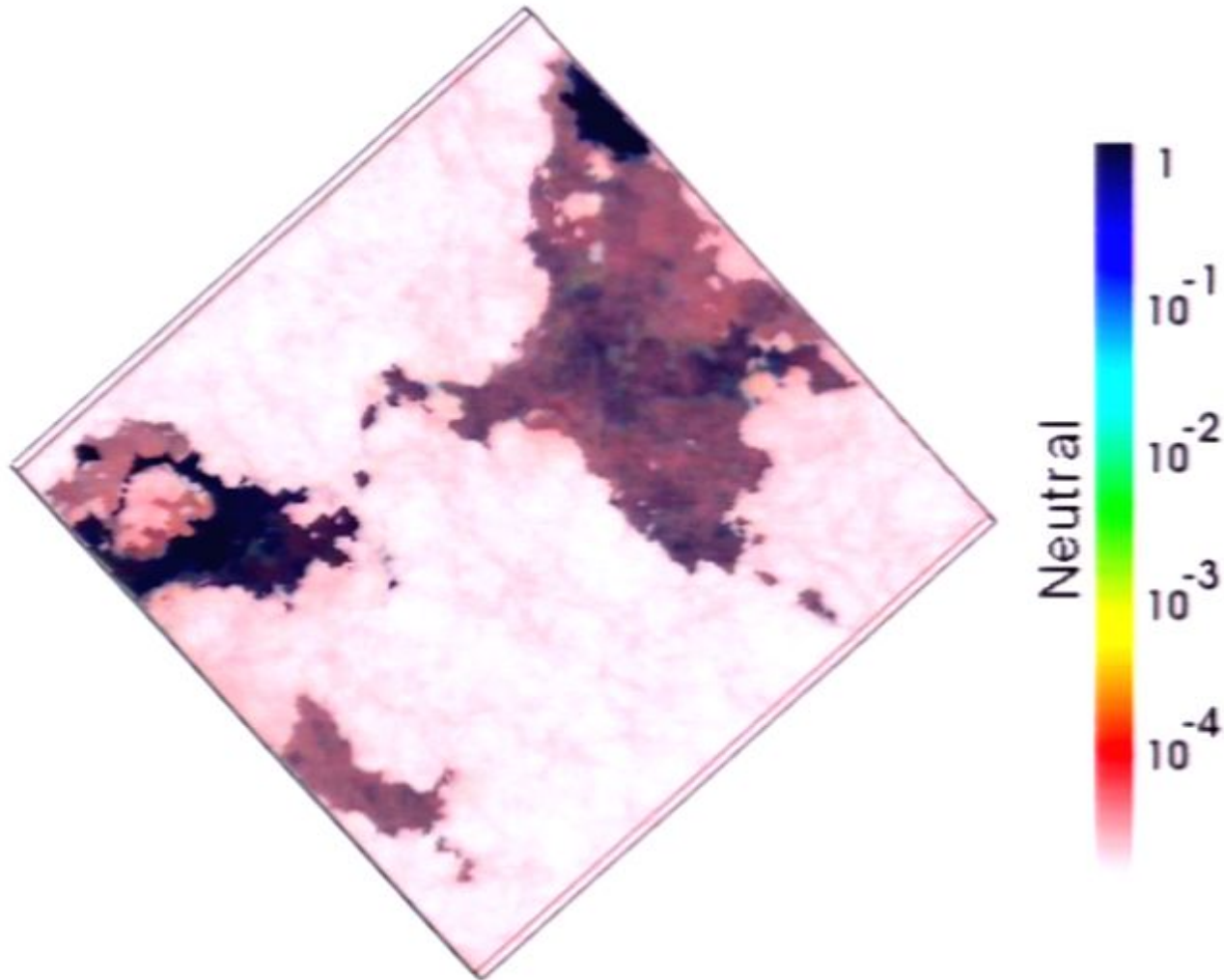


## The epoch of reionization





## The epoch of reionization



The Early window  
before structure  
formation

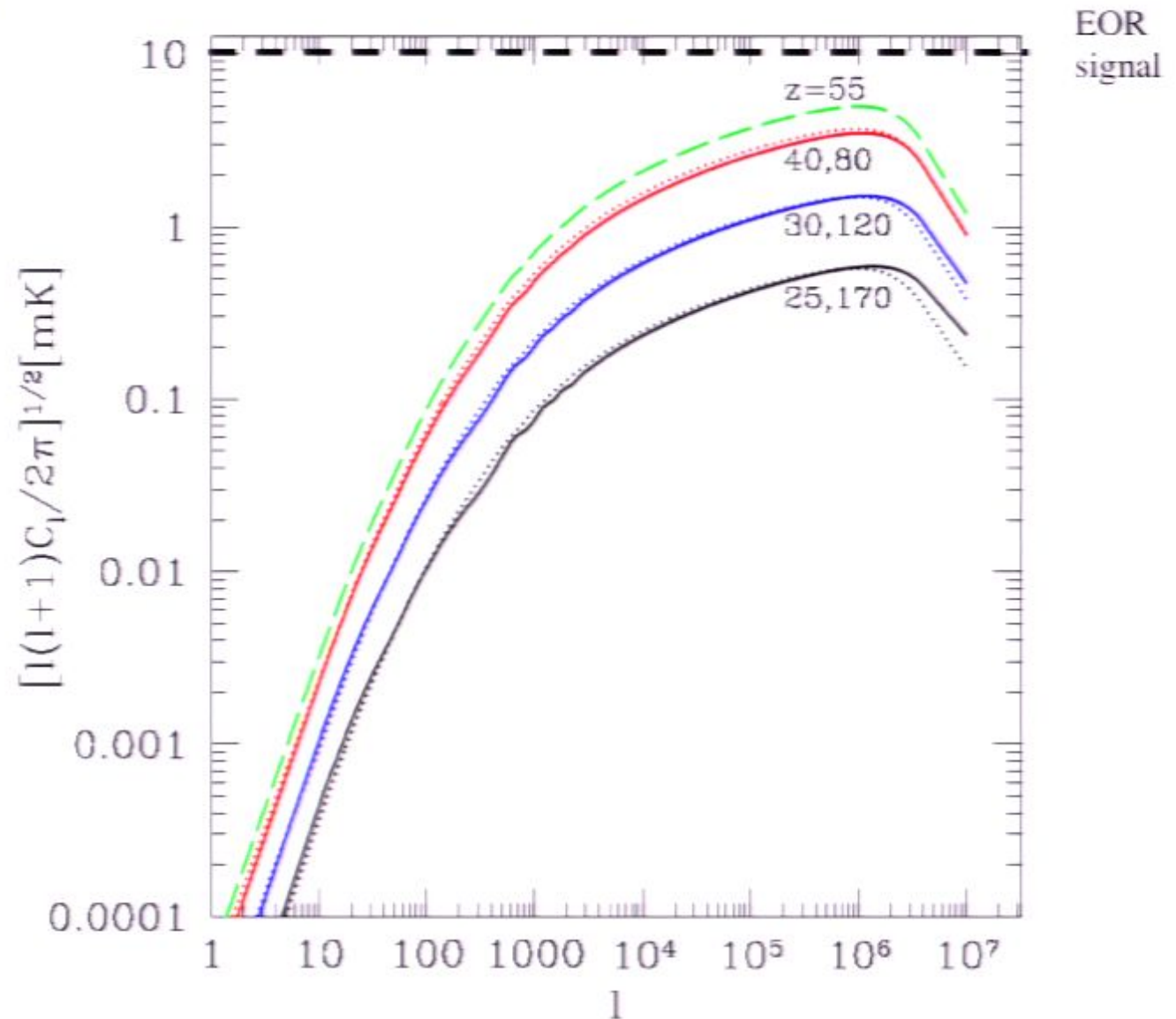


FIG. 2. Angular power spectrum of 21cm anisotropies on the sky at various redshifts. From top to bottom,  $z = 55, 40, 80, 30, 120, 25, 170$ .

## Experimental challenge

### CMB:

Background: 3K CMB

Temperature Fluctuations: 30  $\mu\text{K}$  ( $10^{-5}$ )

E mode polarization: 3  $\mu\text{K}$  ( $10^{-6}$ )

B mode polarization: 0.1  $\mu\text{K}$  ( $3 \times 10^{-8}$ )

Wavelength: 0.3 cm

Frequency: 100 GHz

### 21 cm Fluctuations:

Background: 200 K Galaxy

Temperature Fluctuations: 20 mK ( $10^{-4}$ )

Wavelength: 2 meters

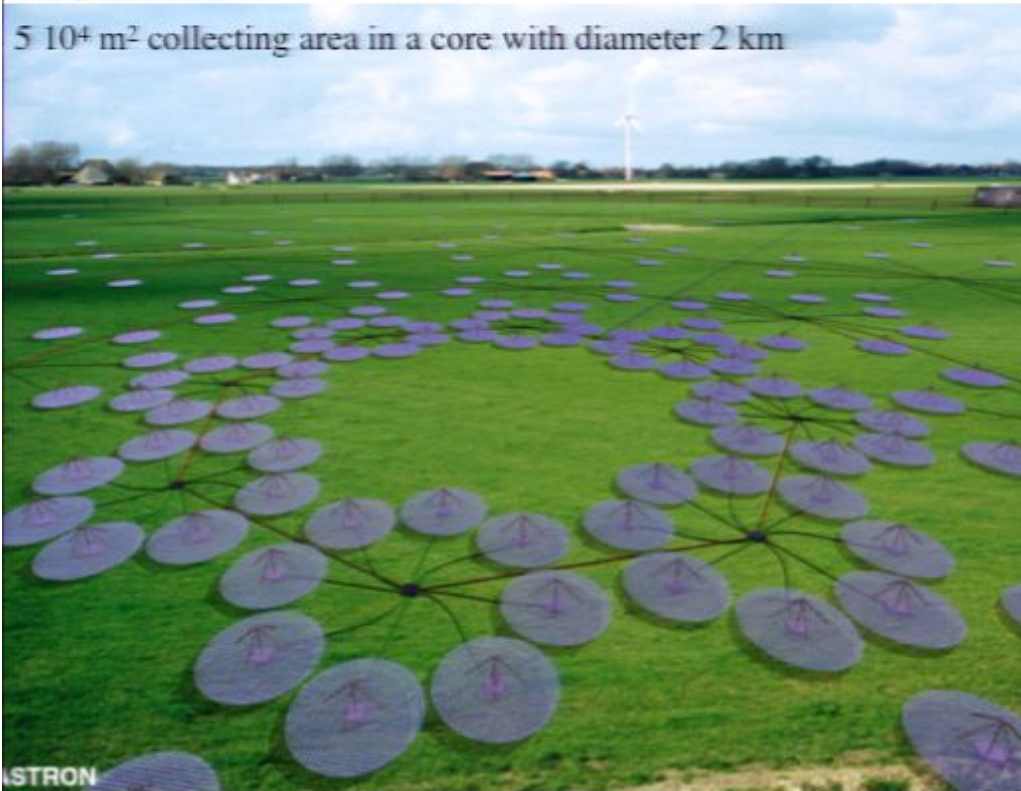
Frequency: 140 MHz

$$\Delta T \propto \frac{\lambda^2 T_{sys}}{A \sqrt{\Delta \nu t_0}}$$



# LOFAR

5 10<sup>4</sup> m<sup>2</sup> collecting area in a core with diameter 2 km



## 21CMA (PAST)

Prototype, Ulatai, West China

Table 1: Specifications of Baseline PAST instrument.

Parameter	Goal (at 100 MHz)
Total antennae	1000
Effective Elements	10
Effective area	9000 m <sup>2</sup>
Element effective receiver temperature	10 K
Sky brightness temperature	170 K
Imaging field of view	1 sq deg
Angular resolution	3 arc minutes

Pirsa: 08040002





# Mileura Wide-field Array Low Frequency Demonstrator (MWA-LFD)

## EOR Science Collaboration

MIT: Jacqueline Hewitt, Colin Lonsdale, Miguel Morales

Melbourne: Rachel Webster, Stuart Wyithe, David Barnes

CfA: Avi Loeb, Lincoln Greenhill, Lars Hernquist, Matias Zaldarriaga

Other Institutions: Steve Furlanetto (UCLA), Chris Carilli (NRAO), Frank Briggs (ANU)

## Characteristics of the Array

- 500 16-element antennas in a compact 1.5 km array
- Wide 80–300 MHz frequency range
- Digital receiver & filter chain, 32 MHz at 16 KHz resolution
- Full cross-correlation of all 500 antennas, 20–30° field of view





# Western Australia



Booldardy in the Murchinson area (MWA), 0.01 humans/km<sup>2</sup>

# Early Deployment

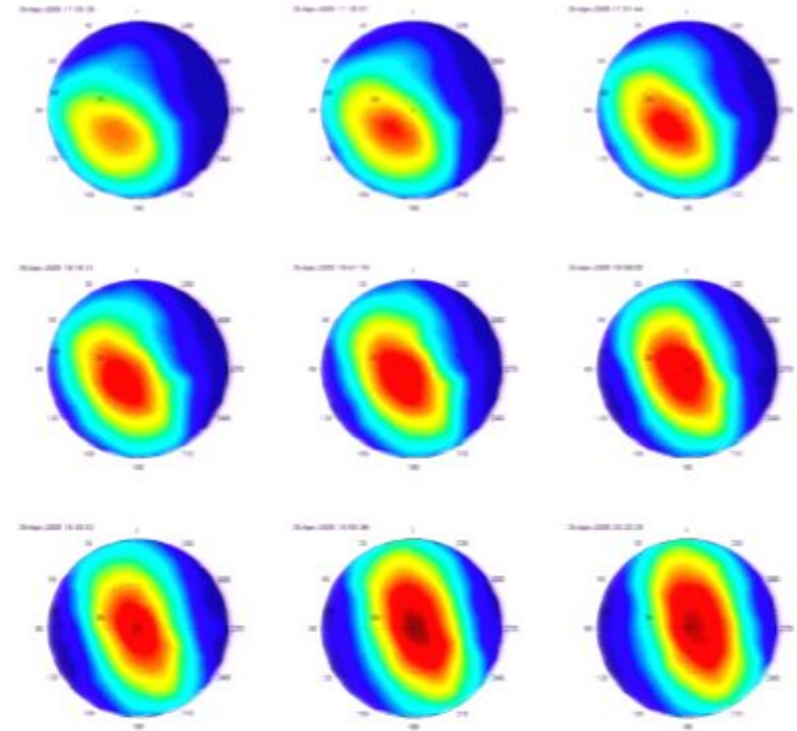


Figure 6. Image sequence showing the galactic center passing overhead at Mileura. The observing frequency is 108 MHz, and the images were made by steering the single-tile beam to 40 grid locations in rapid succession. The zenith is at the center of the images, and the edges area at 30 degrees elevation. The colors cover a brightness range of 4 to 1, and the tile beamwidth is about 34 degrees. The data have not been corrected for the single-dipole power pattern, which peaks at the zenith.





## The raw constraining power of 21 cm tomography:

For a 2D survey

$$N_{2D} = l_{max}^2 = 9 \times 10^6 \left( \frac{l_{max}}{3000} \right)^2$$

In the case of 21 cm, there is extra radial information

$$N_{21cm} = 3 \times 10^{16} \left( \frac{l_{max}}{10^6} \right)^3 \left( \frac{\Delta\nu}{\nu} \right) \left( \frac{z}{100} \right)^{-1/2} \propto l_{max}^3$$

Constraints on departures from scale invariance and on primordial non Gaussianity scale as

$$\frac{1}{\sqrt{N_{obs}}}$$

For CMB

$$\frac{1}{\sqrt{N_{obs}}} \sim \frac{1}{l_{max}} \sim 10^{-3}$$

For inflationary models we expect to see  $(n-1)$  but not running  $(n-1)^2$

$$NG \sim (n - 1) \frac{\delta\rho}{\rho} \quad \text{Too small in simple models}$$

6.2 The John M. Hedgepeth Concept

Getting away from RFI and increasing collecting area

Lunar orbit

Earth Trailing orbit

Wire mesh telescope:  
The John M.  
Hedgepeth Concept  
(1965)

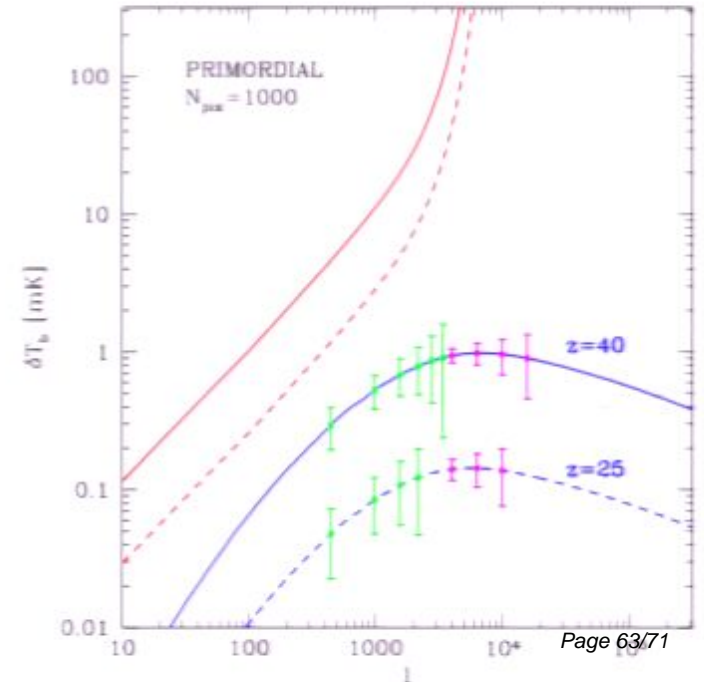
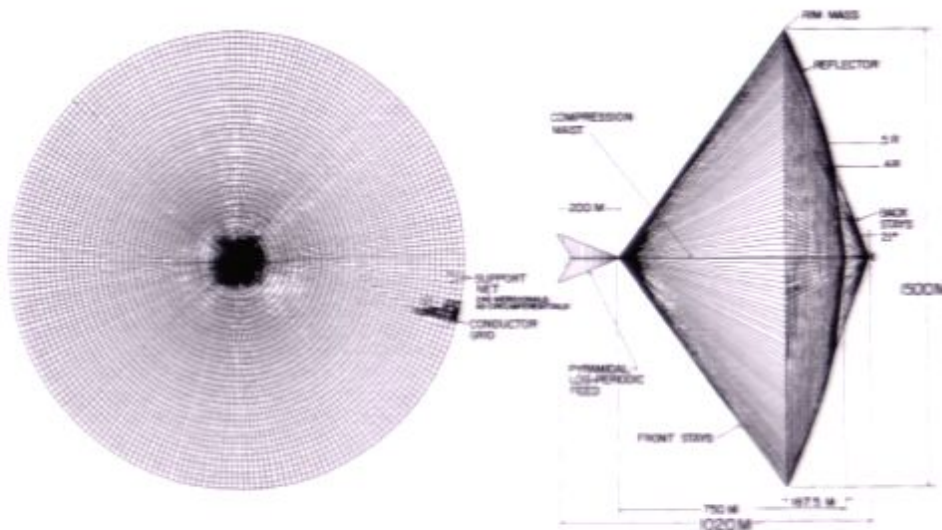
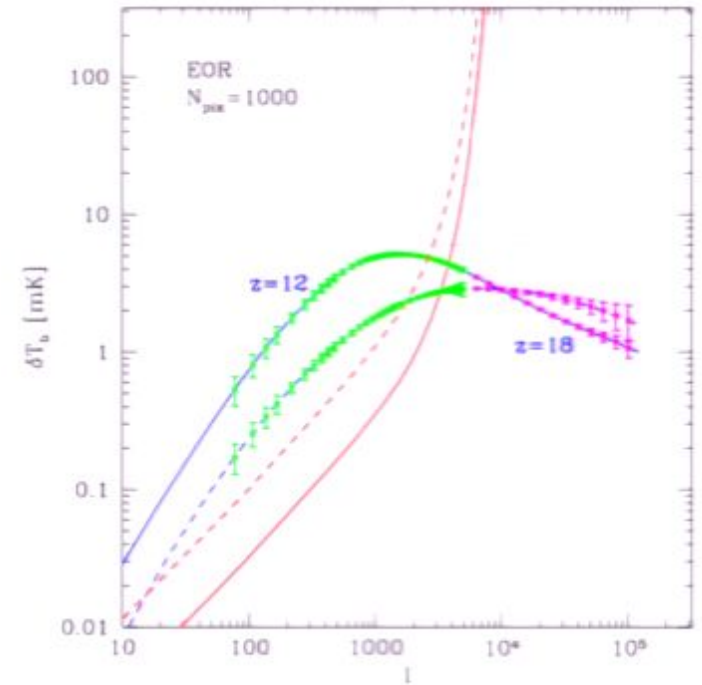
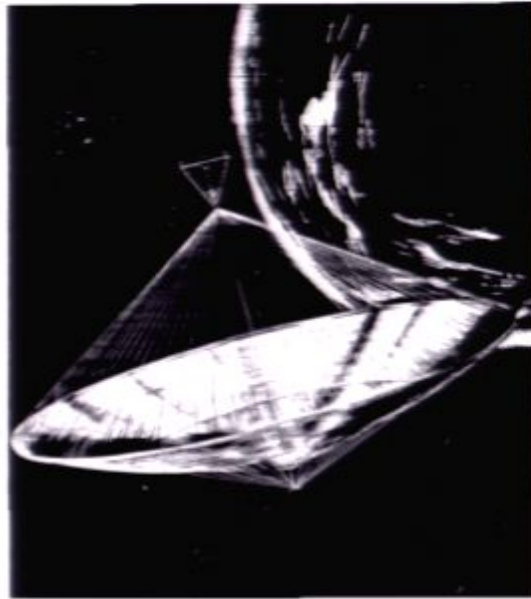


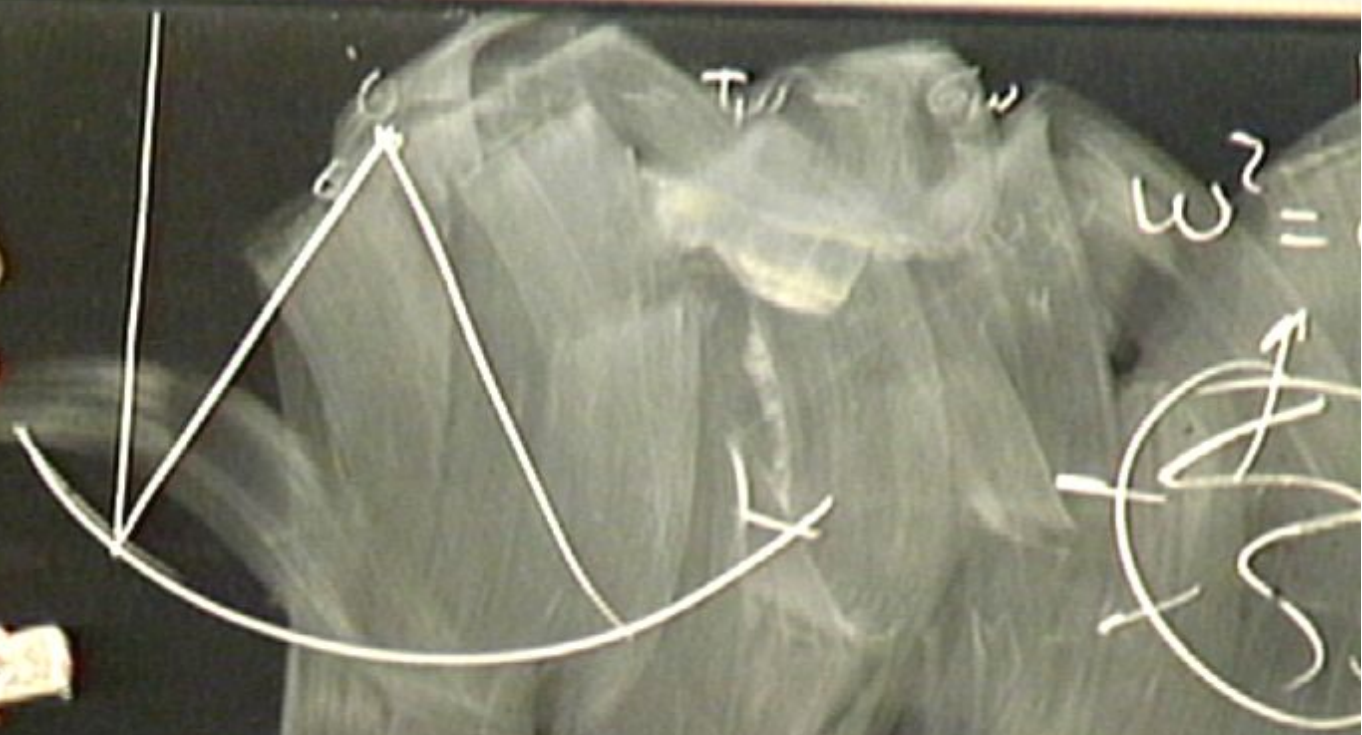
Figure 10. Near Parabolic John M Hedgepeth Concept

# Summary:

- Upcoming 21 cm observations have the sensitivity to detect fluctuations in the background which will constrain both cosmological parameters and the physics of reionization, opening a new window for Cosmology.
- In principle this technique might enable us to map the largest fraction of our Hubble volume, resulting in the a huge “number of numbers” measured.
- Will we ever be able to use a substantial fraction of the information that is out there in the 21 cm signal? We will have a much better idea of the answer to this question in a few years when the first observatories come on line.



$$\| \text{set } T \| = (K) T \quad \text{yours}$$





$$\mathbf{p} \cdot \mathbf{t} = (\mathbf{K}) \cdot \mathbf{t}$$

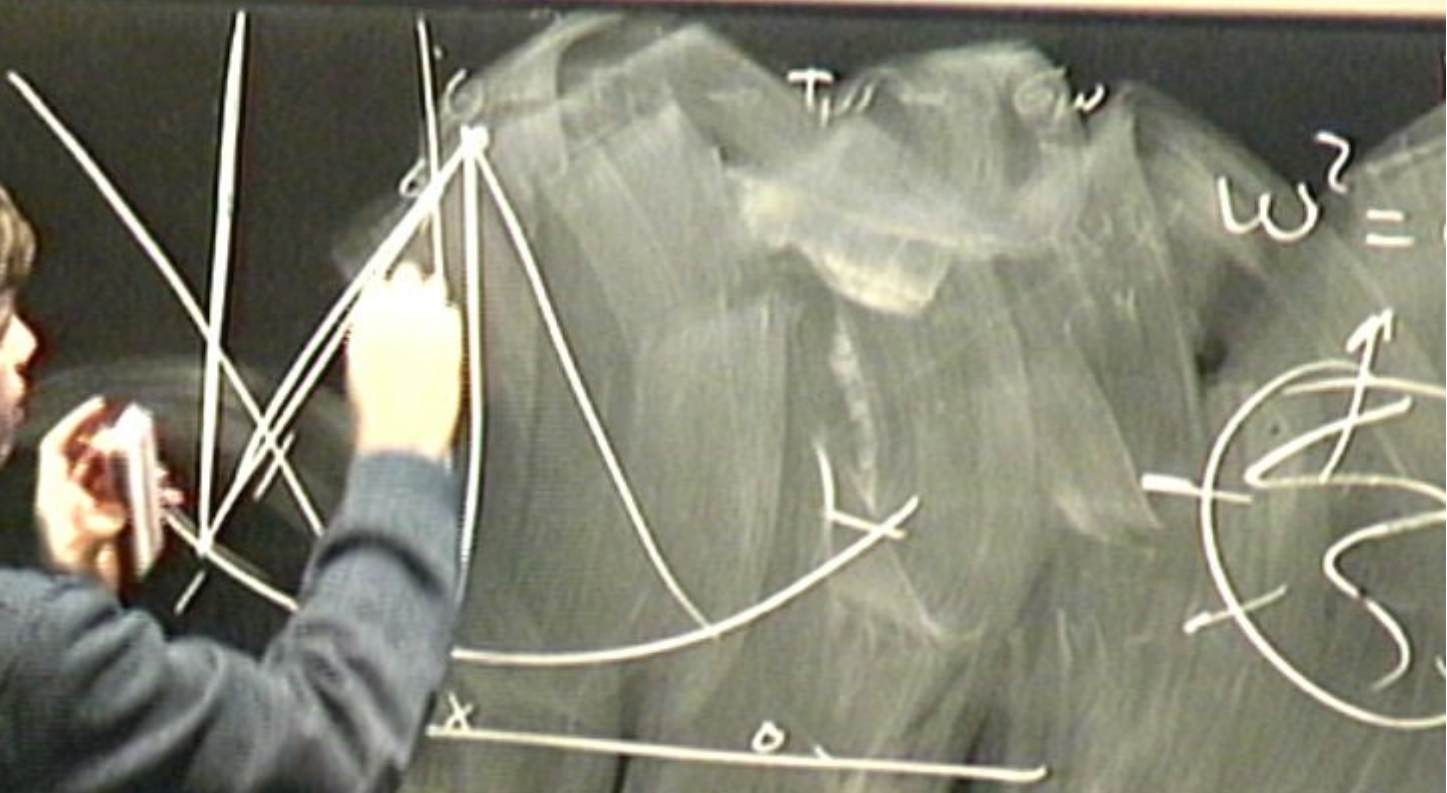


$$\omega^2 =$$





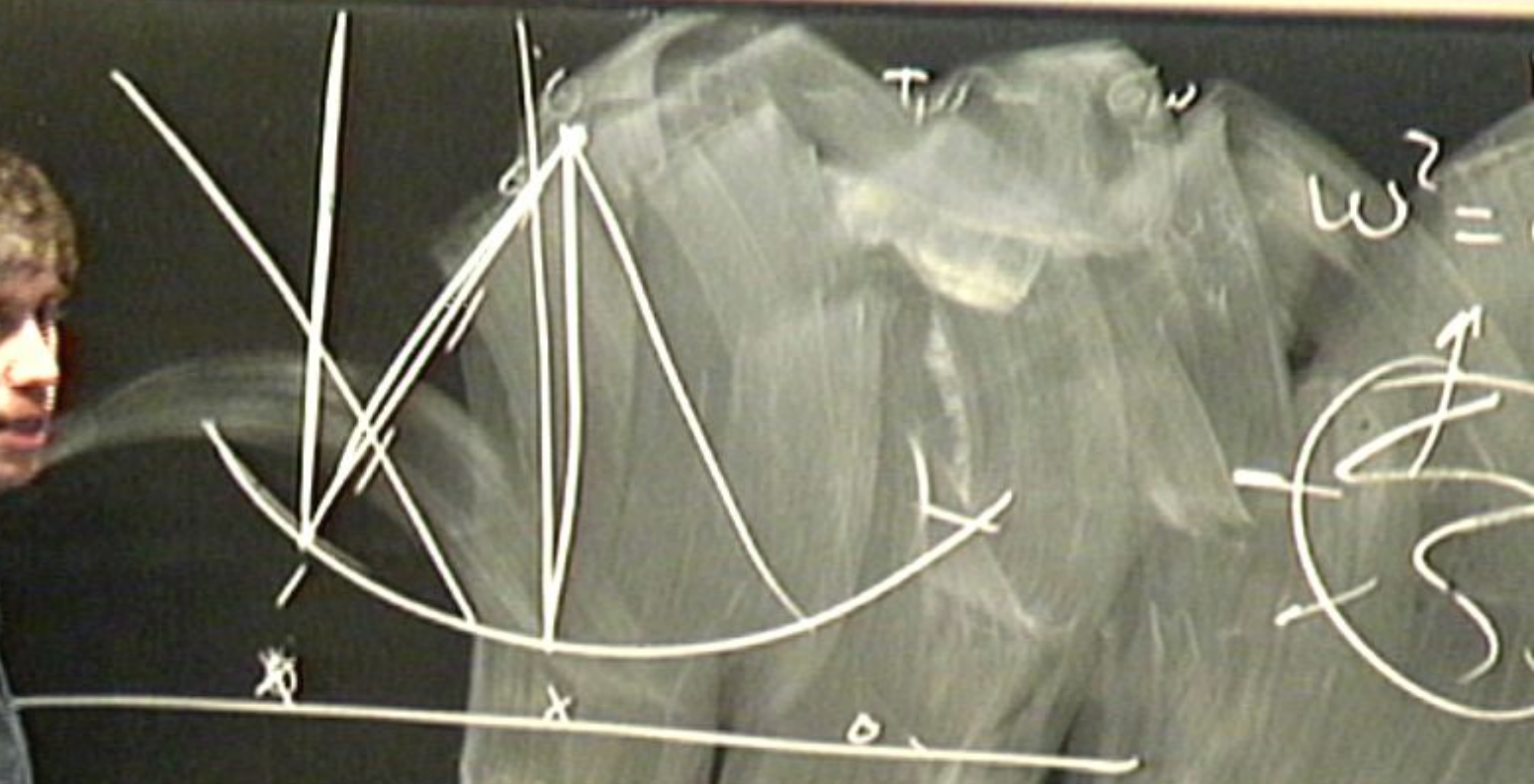
$$\mathbf{H} \mathbf{z} + \mathbf{t} = (\mathbf{K}) \mathbf{T} \quad \text{yours}$$



$$\omega^2 =$$



$$\text{Tr } \rho T = \text{Tr } (T \rho) = \text{Tr } \rho T$$



$$\omega^2 =$$

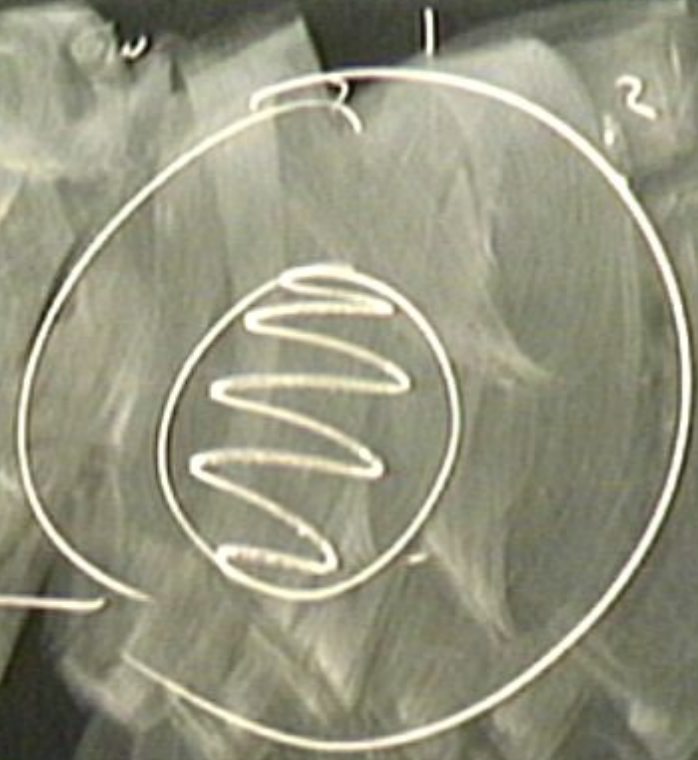




$$= e^{k \cdot v}$$

$$H^2_{sp} + 1 \equiv H(k) + \dots$$

quadratic eqn in non-affine





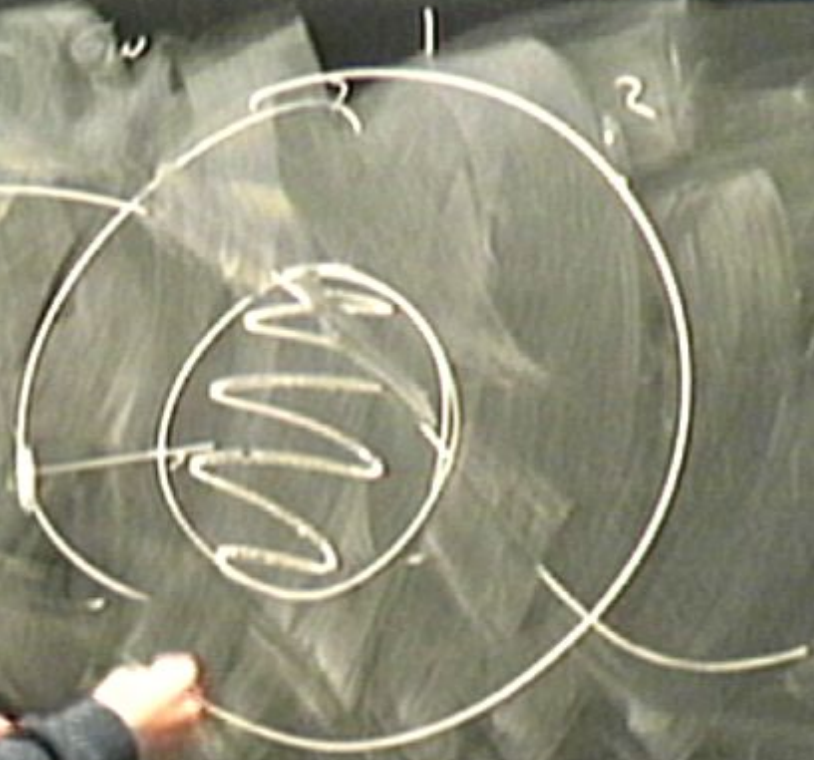
$$= e^{k \cdot v}$$

$$H^2_{ip} + \dots = H(k) + \dots$$

quadratic map in non-affine



$T_{11} = \dots$





$$= e^{kx}$$

$$H^2_{sp} + \mu = H(k) + \dots$$

quadratic eqn in non-affine

