

Title: Imprints of primordial non-gaussianity on large-scale structure

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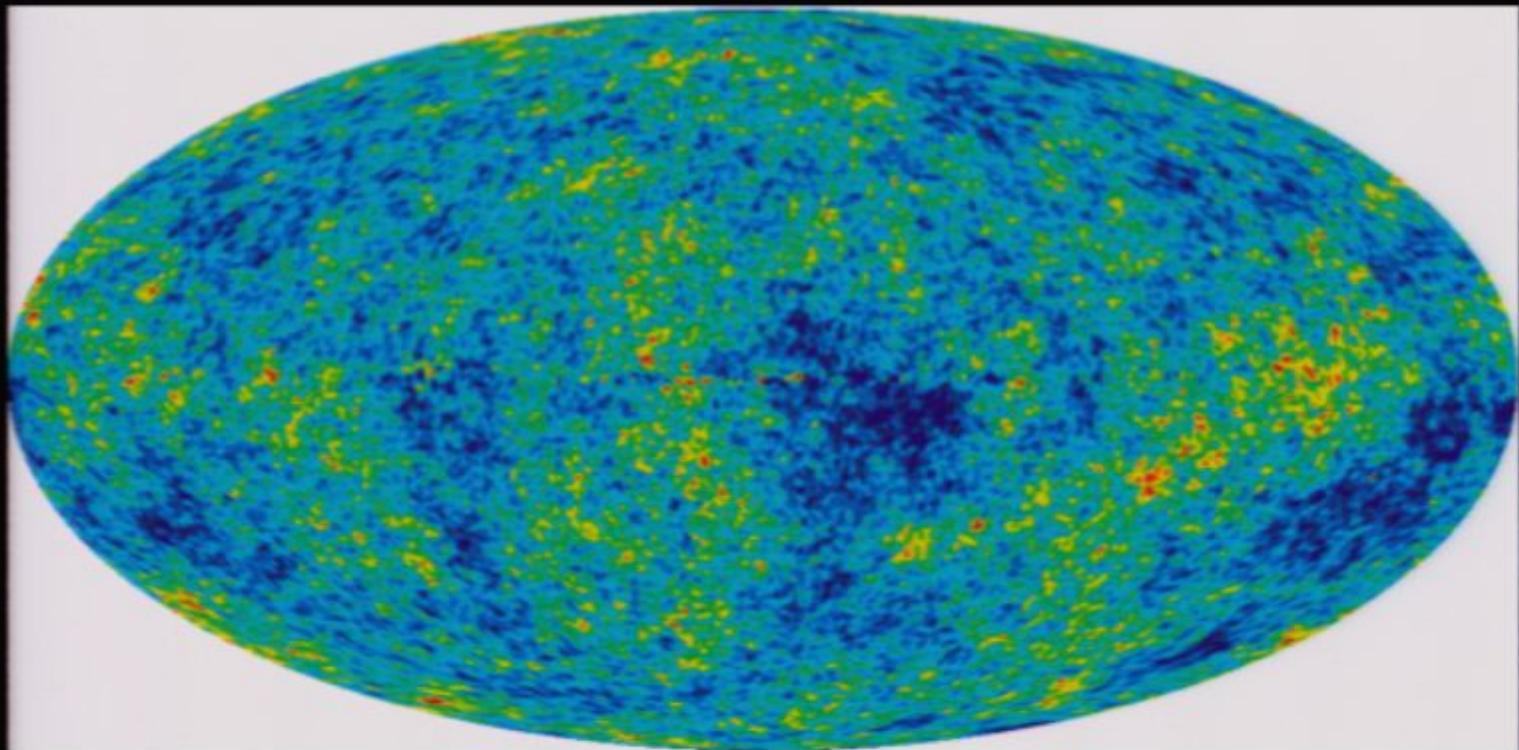
Abstract:

# Imprints of Primordial Nongaussianity on Large-scale Structure

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(University of Michigan)

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Olivier Doré (CITA)  
Alex Shirokov (CITA)

# Initial conditions in our universe



## Generic inflationary predictions:

- Nearly scale-invariant spectrum of density perturbations
- Background of gravity waves
- (Very nearly) gaussian initial conditions

# Inflation generically predicts (very nearly) gaussian random fluctuations

- Nongaussianity is proportional to slow-roll parameters,  $V'/V$  and  $V''/V$
- Reasonable and commonly used approximation
- If  $\Phi$  is Gaussian, then  $f_{NL} \approx 0(0.1)$ , which is basically too small to ever measure
- More exotic inflationary models can produce observable NG, however

$$\Phi = \phi + f_{NL}(\phi^2 - \langle \phi^2 \rangle)$$

Salopek & Bond 1990; Maldacena 2003;  
Seery & Lidsey 2005;  
Chen, Easther & Lim 2008

# Brief history of NG measurements: 1990's

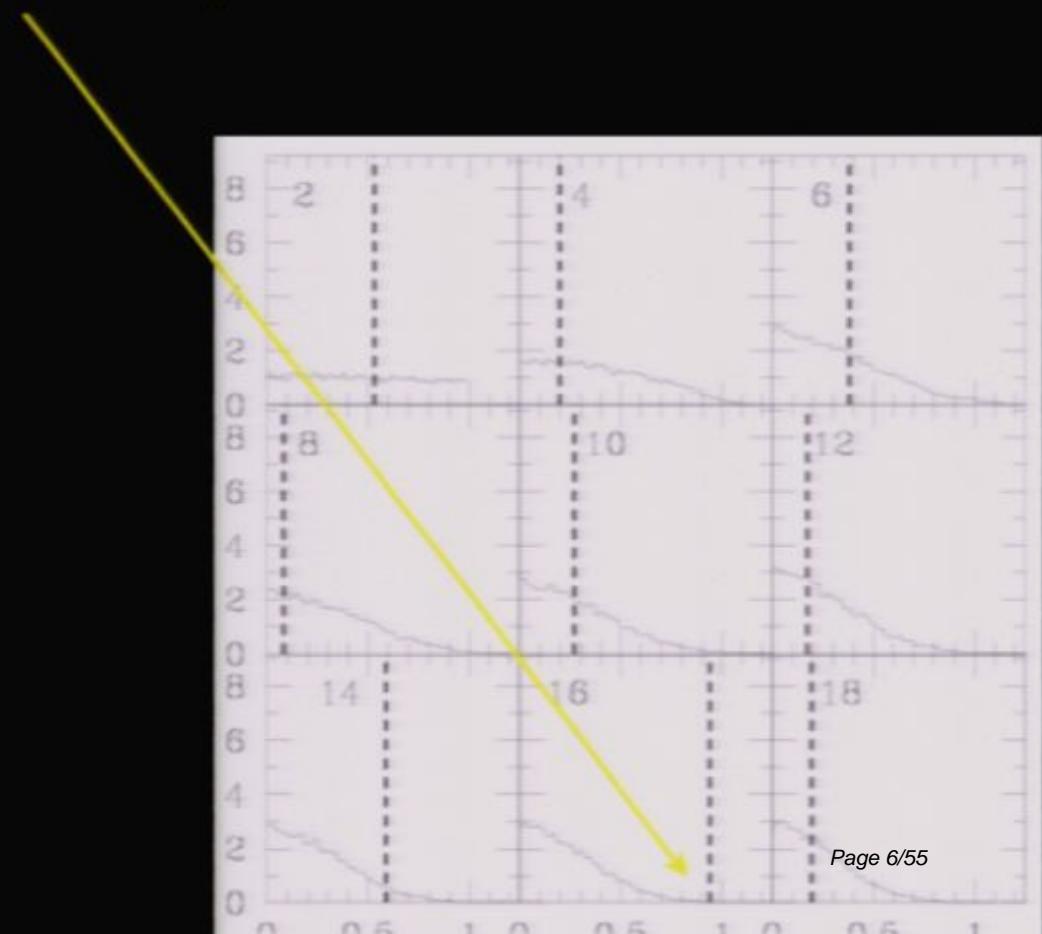
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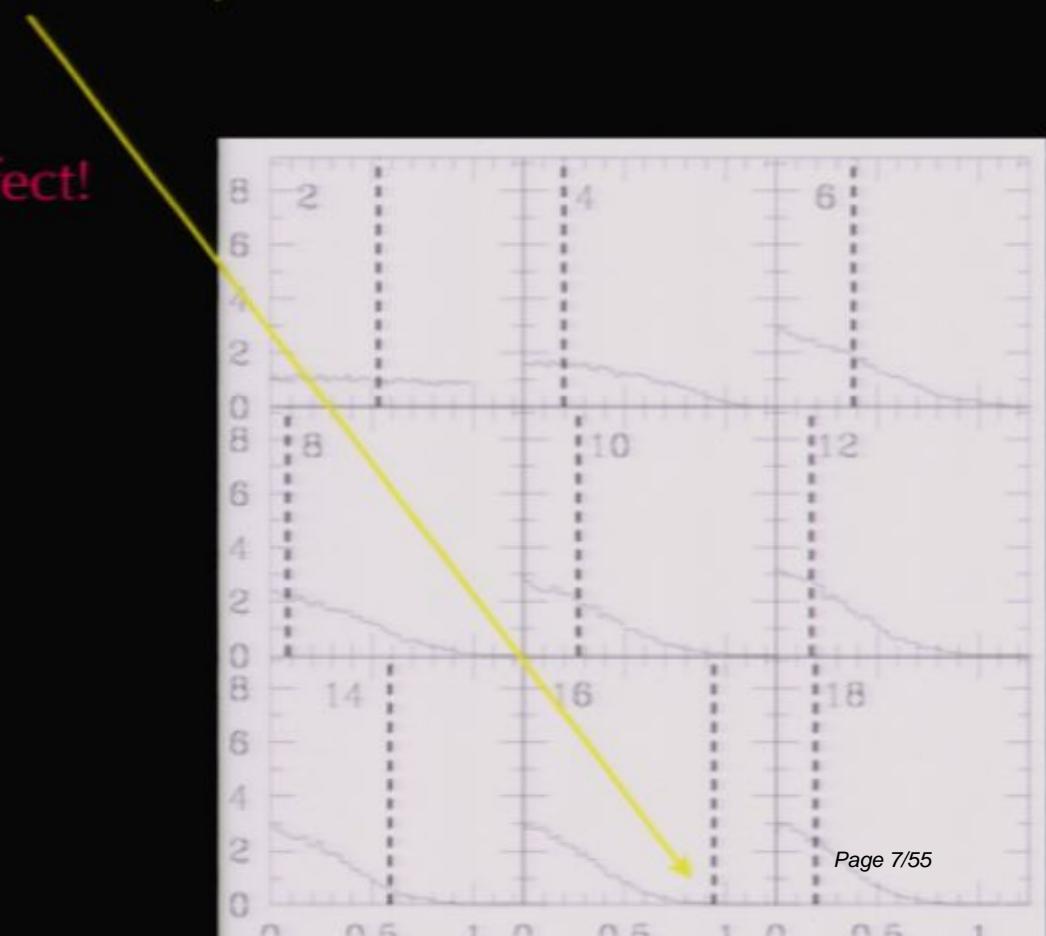


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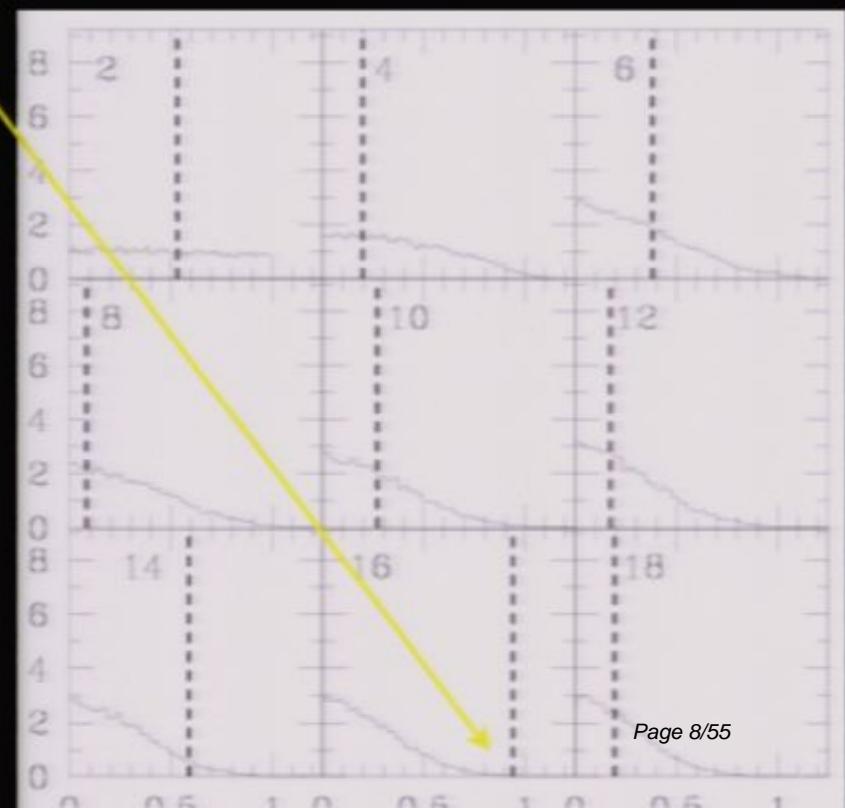
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and anyway isn't unexpected given all  
bispectrum configurations you can measure;  
Komatsu 2002)



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Pre-WMAP CMB: all is gaussian (e.g. MAXIMA; Wu et al 2001)

WMAP pre-2008: all is gaussian

(Komatsu et al. 2003; Creminelli, Senatore, Zaldarriaga & Tegmark 2007)

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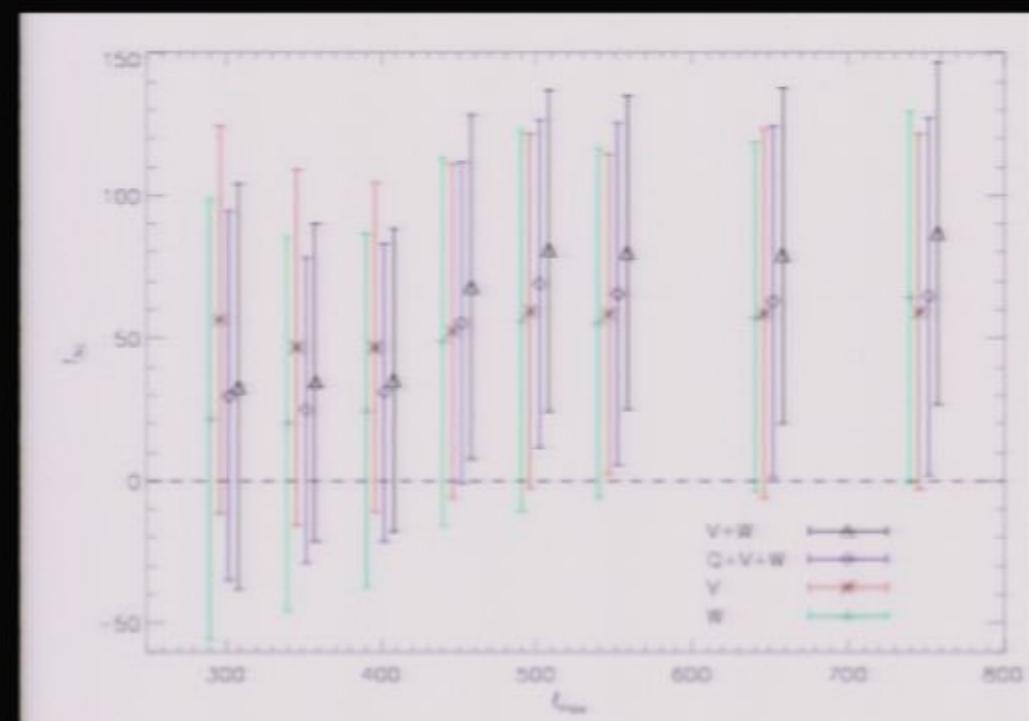
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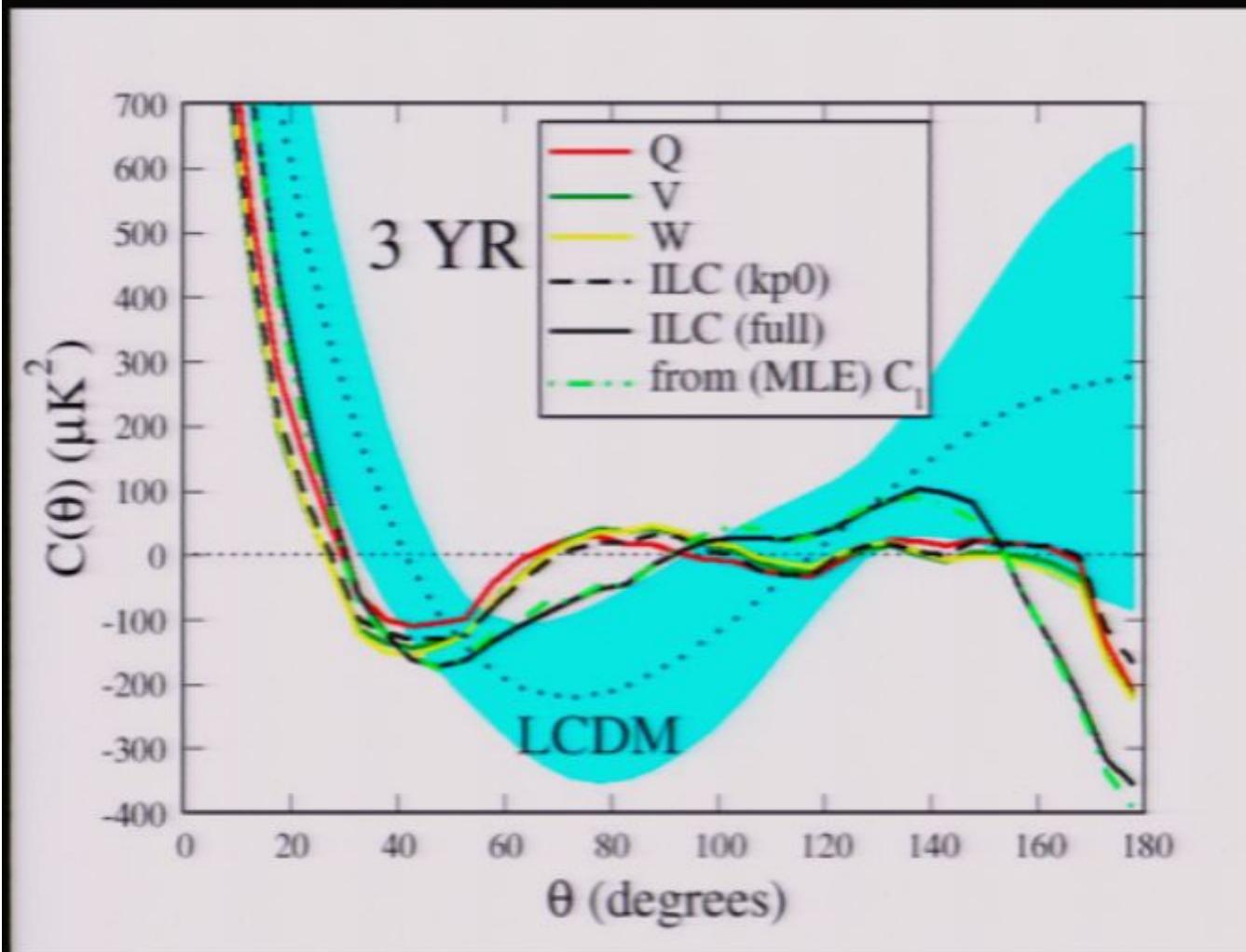
Dec 2007, claim of NG in WMAP

(Yadav & Wandelt arXiv:0712.1148)

$$27 < f_{NL} < 147 \quad (95\% \text{ CL})$$



... and also “large-scale anomalies”



e.g. lack of power  
at >60 deg;  
significant at 99.96%

Hinshaw et al. 1996

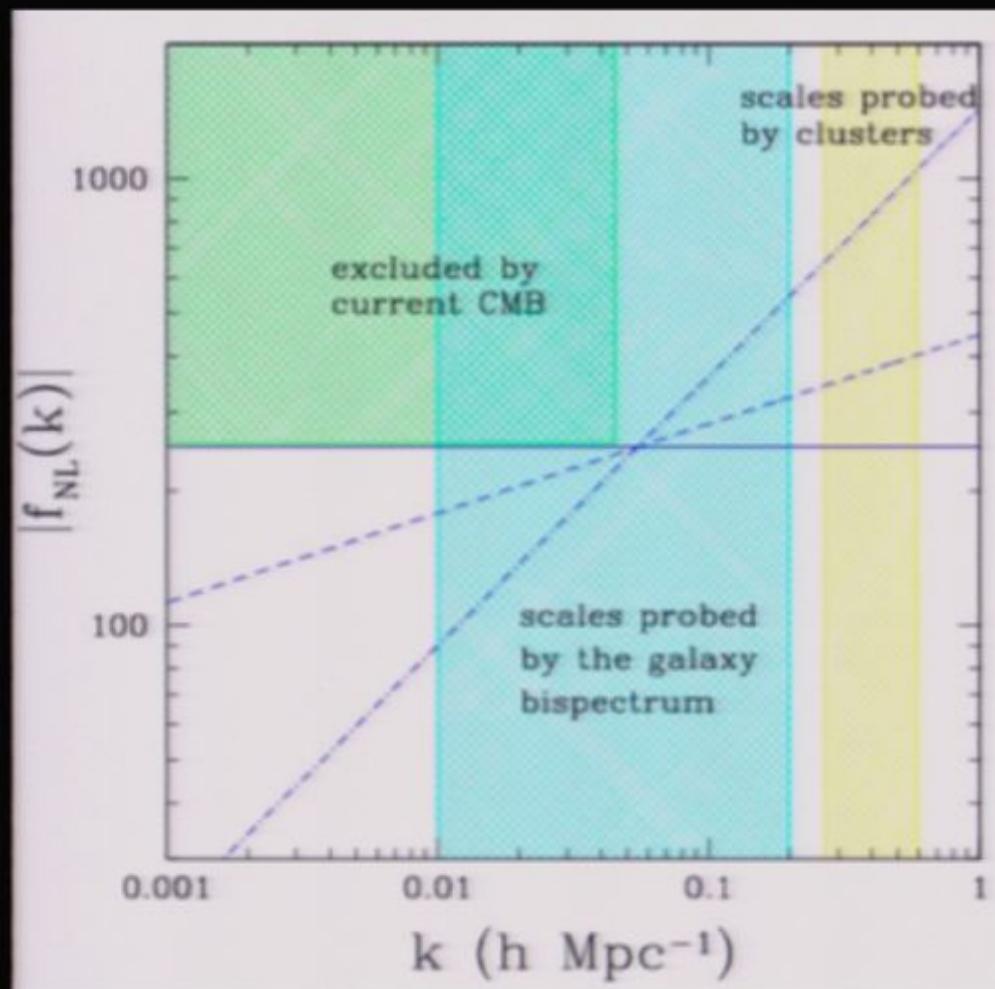
(COBE)

Spergel et al. 2003

(WMAP 1)

Copi, Huterer, Schwarz & Starkman 2007 (WMAP 3)

# Constraints from future LSS surveys



Sefusatti, Vale, Kadota & Frieman, 2006

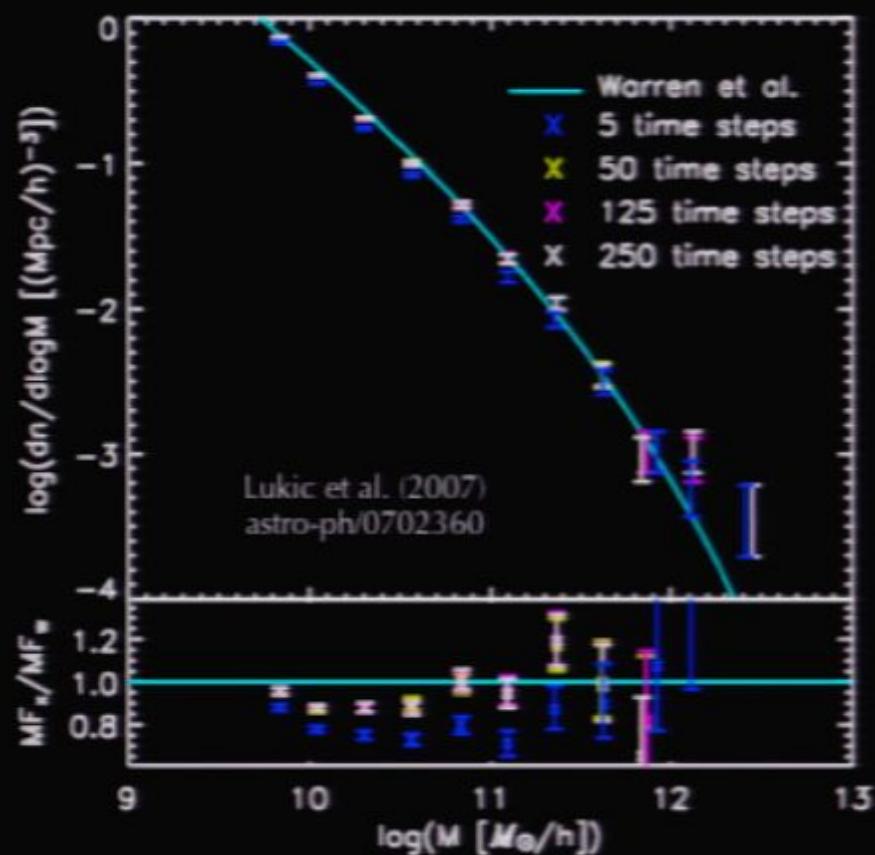
LoVerde, Miller, Shandera & Verde, arXiv:0711.4126

# Abundance of halos: the mass function

Lots of interest in using halo counts as a cosmological probe.

- Mass function can be computed precisely (~5%) and robustly for standard cosmology (Jenkins et al. 01, Warren et al. 03)
- $dN/dM$  appears universal — i.e.  $f(\sigma)$  — for standard cosmologies

$$\sigma^2(M, z) = \frac{1}{2\pi^2} \int_0^\infty k^2 P(k) W^2(k, M) dk$$



# Mass function, usual analytic approach

Press & Schechter 1974:

$$\frac{dn}{dM} dM = \frac{\rho_M}{M} \left| \frac{dF}{dM} \right| dM \quad F(> M) = 2 \int_{\delta_c/\sigma(M)}^{\infty} P_G(\nu) d\nu$$

therefore

$$\left( \frac{dn}{d \ln M} \right)_{\text{PS}} = 2 \frac{\rho_M}{M} \frac{\delta_c}{\sigma} \left| \frac{d \ln \sigma}{d \ln M} \right| P_G(\delta/\sigma)$$

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follow EPS, then expand  $P_{NG}$  in terms of skewness, do the integral

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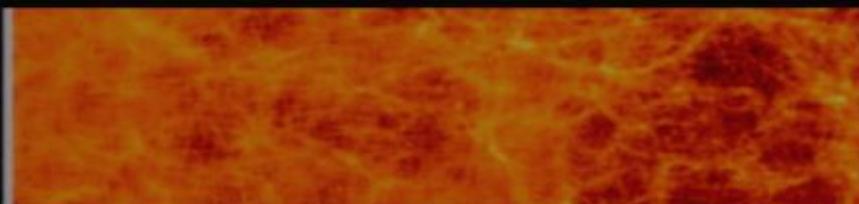
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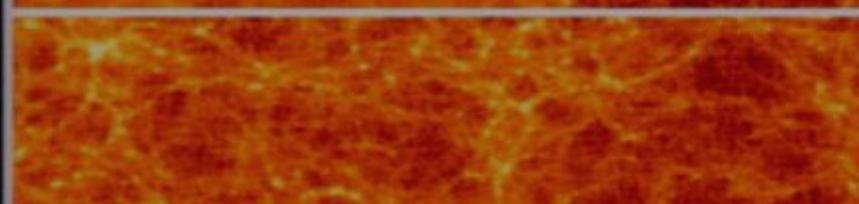
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Need to check these formulae with simulations

# Simulations with nongaussianity ( $f_{NL}$ )

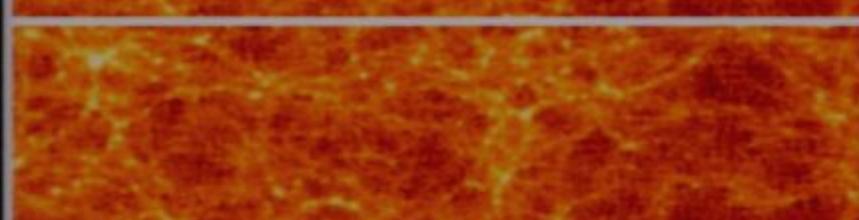
$f_{NL} = -5000$



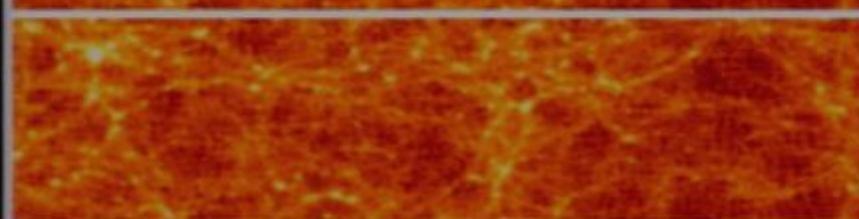
$f_{NL} = -500$



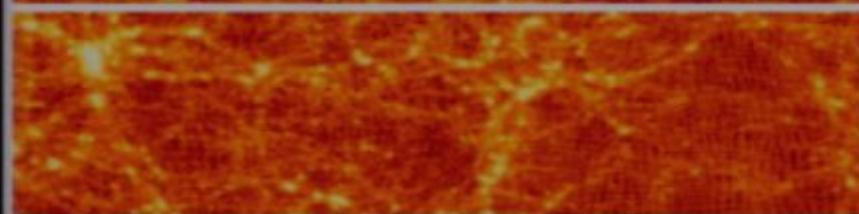
$f_{NL} = 0$



$f_{NL} = +500$



$f_{NL} = +5000$



■ Under-dense region evolution decrease with  $f_{NL}$

■ Over-dense region evolution increase with  $f_{NL}$

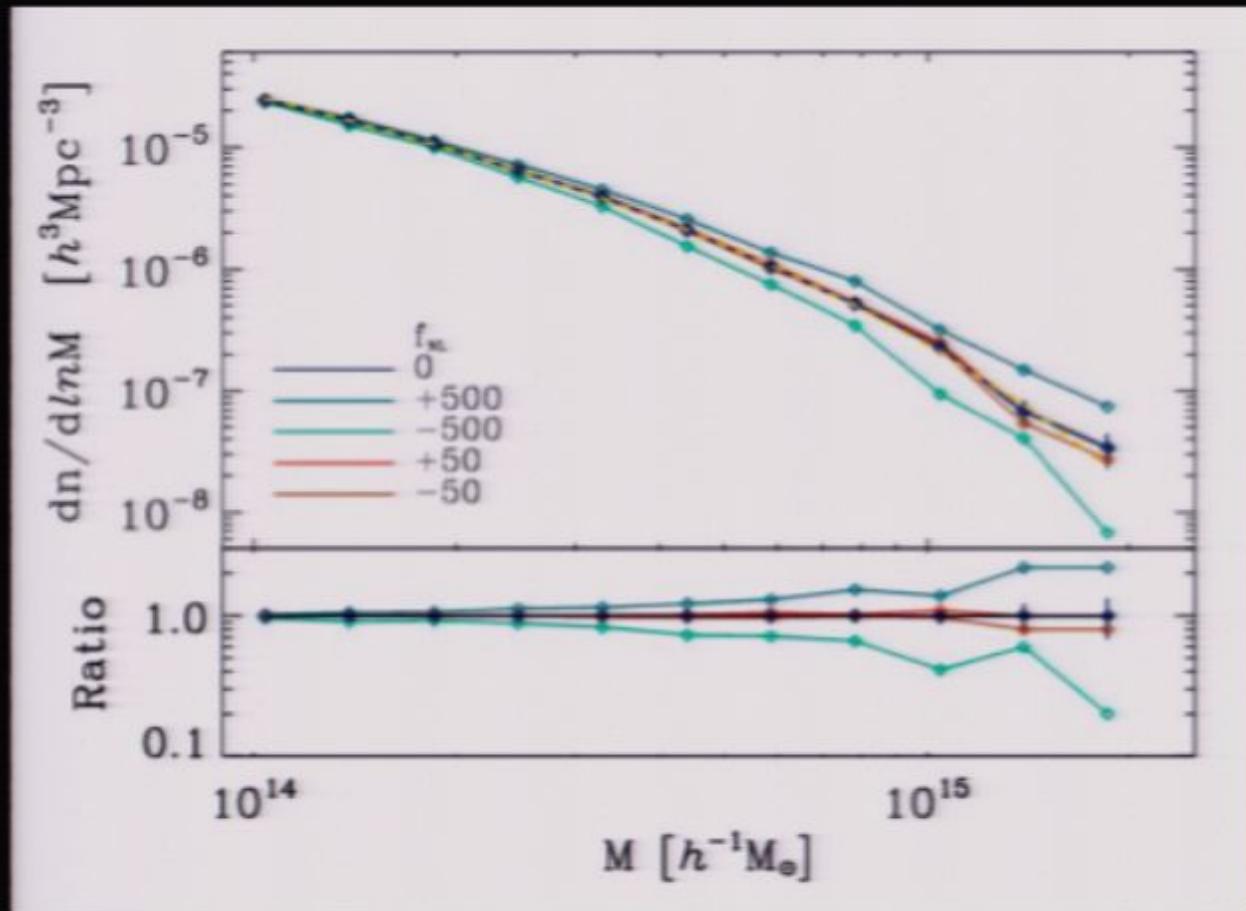
80 Mpc/h

375 Mpc/h

■ Same initial conditions, different  $f_{NL}$

■ Slice through a box in a simulation  $N_{part}=512^3$ ,  $L=800$  Mpc/h

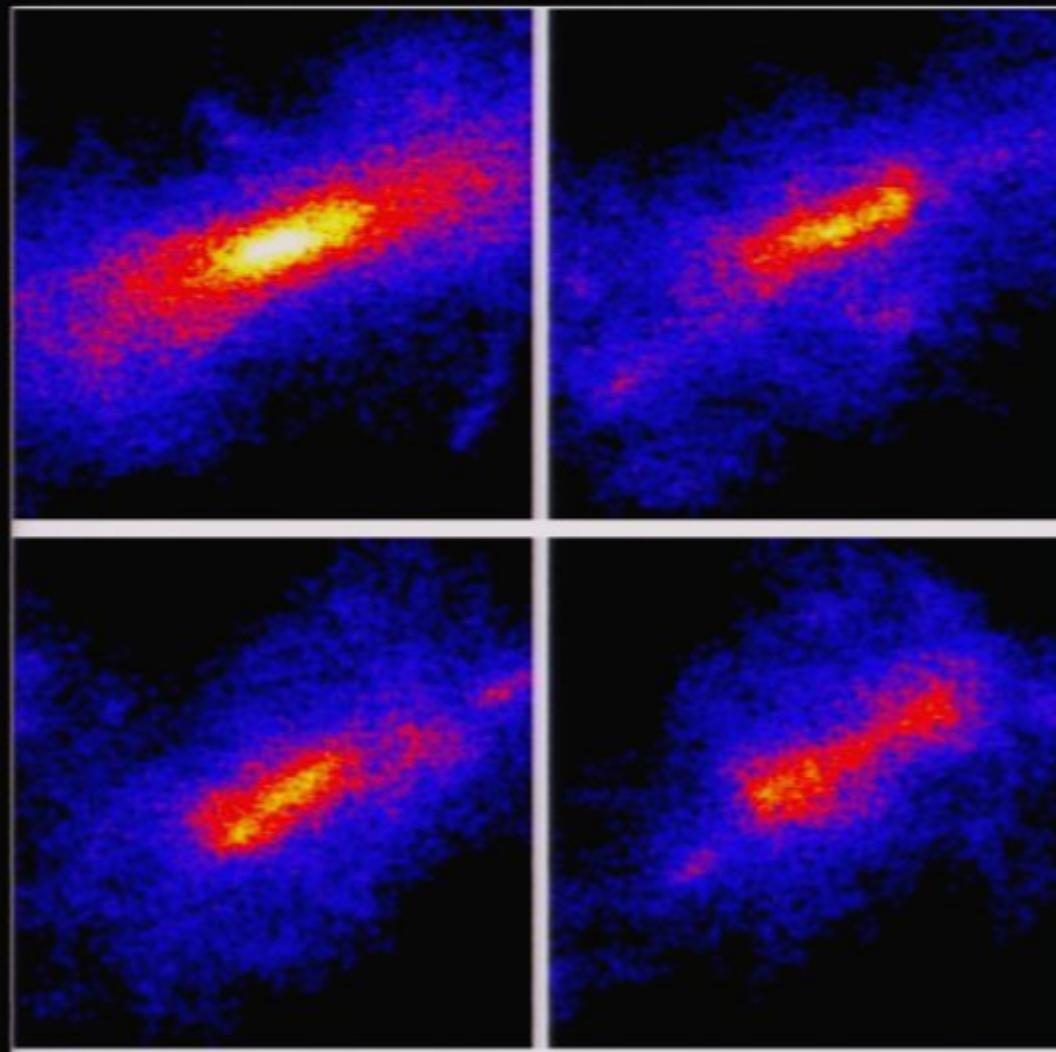
# The measured halo mass function



- $512^3$  ( $1024^3$ ) particle simulations with box size 800 (1600)  $\text{Mpc}/h$
- Gracos code ([www.gracos.com](http://www.gracos.com)); add quadratic Phi term in real space; apply transfer function in Fourier space

# Looking at one individual cluster

$f_{NL}=+5000$   
 $M=1.2 \cdot 10^{16} M_\odot$



$f_{NL}=+500$   
 $M=5.9 \cdot 10^{15} M_\odot$

$f_{NL}=0$   
 $M=5.1 \cdot 10^{15} M_\odot$

$f_{NL}=-500$   
 $M=4.3 \cdot 10^{15} M_\odot$

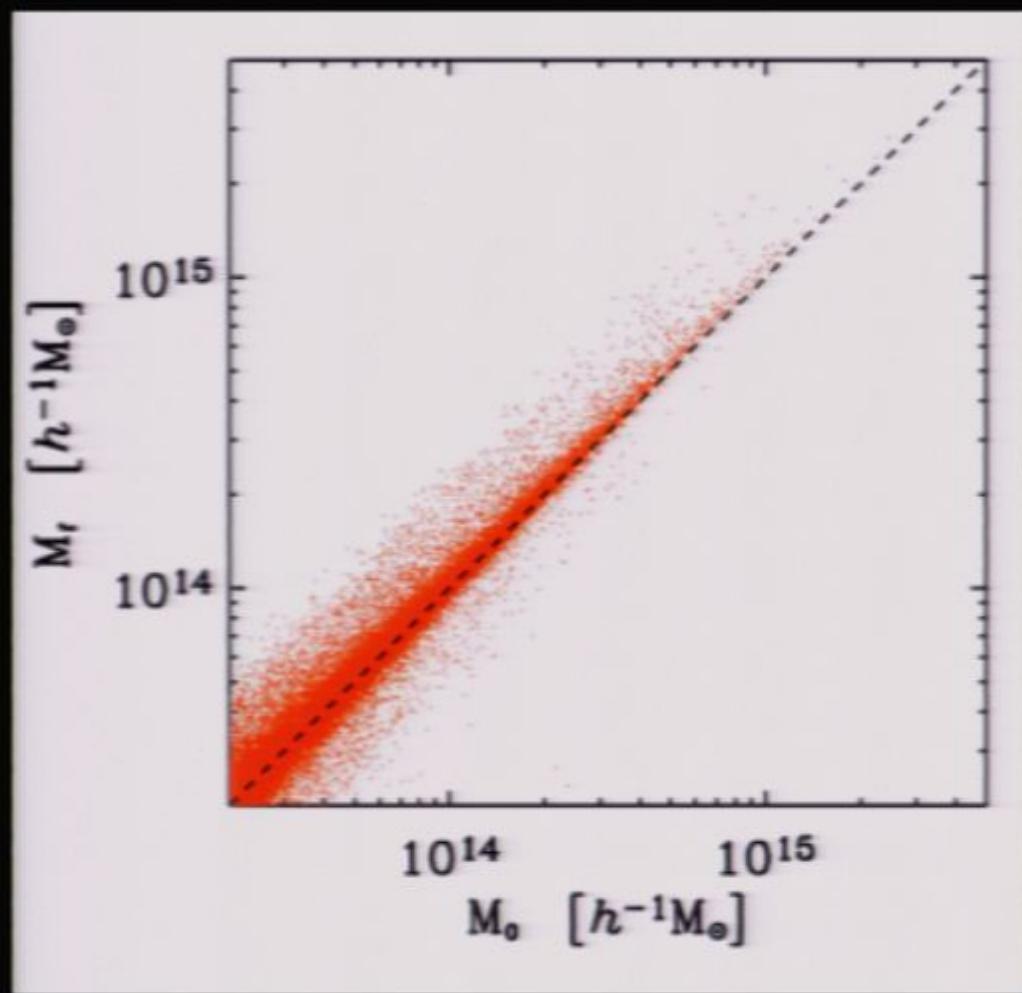
► Most massive cluster in our simulation

► For small enough  $f_{NL}$ , same peaks arise, with different heights (implying different masses)

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► Can we extend to any cluster?

# Building the $P(M_f|M_0)$ distribution



$f_{NL} = 500$

- Idea: identify the *same* cluster for different  $f_{NL}$ , keep track how its mass changed!
- Significantly saves computational expenses

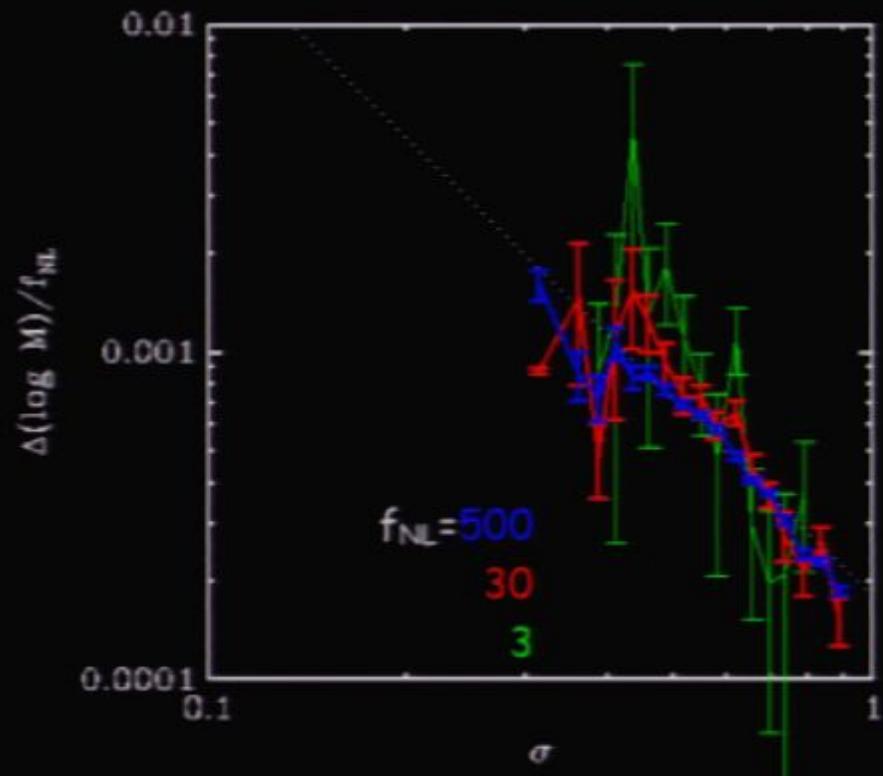
# Towards a fitting function

- If the mapping  $M_0 \rightarrow M_f$  is described by a PDF  $dP/dM_f(M_0)$ , then the non-gaussian mass function is a convolution over the (known) gaussian mass function

$$\frac{dN}{dM} = \int \frac{dP(M_f|M_0)}{dM_f} \frac{dN}{dM_0} dM_0$$

(e.g. Jenkins)

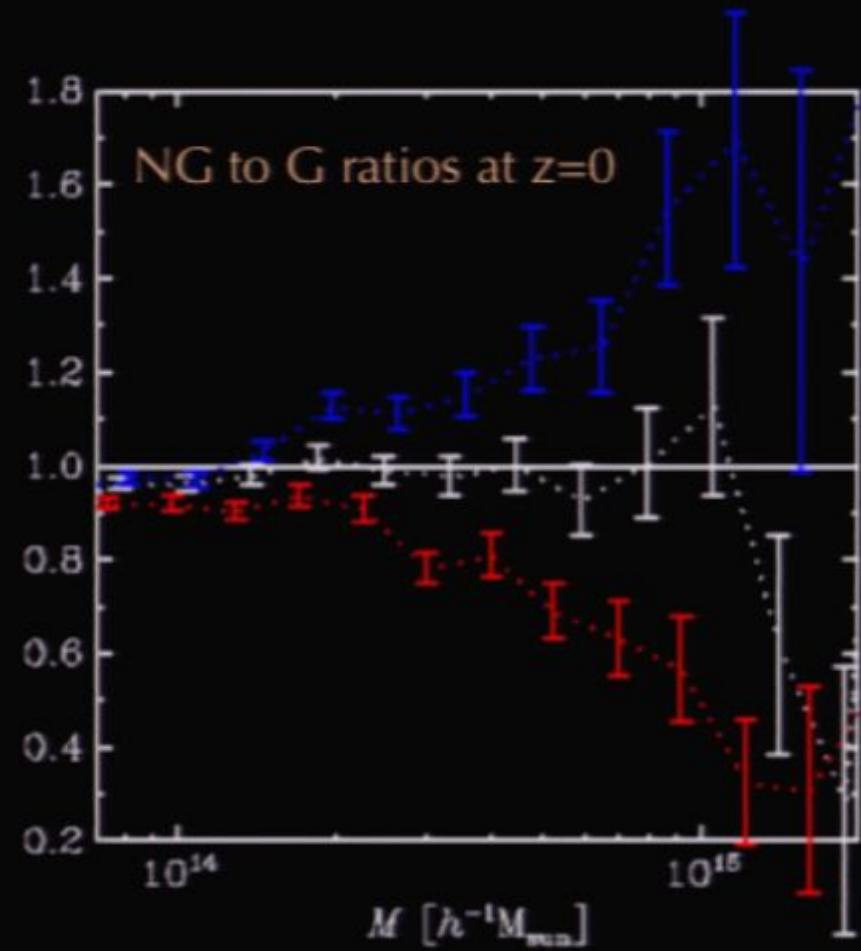
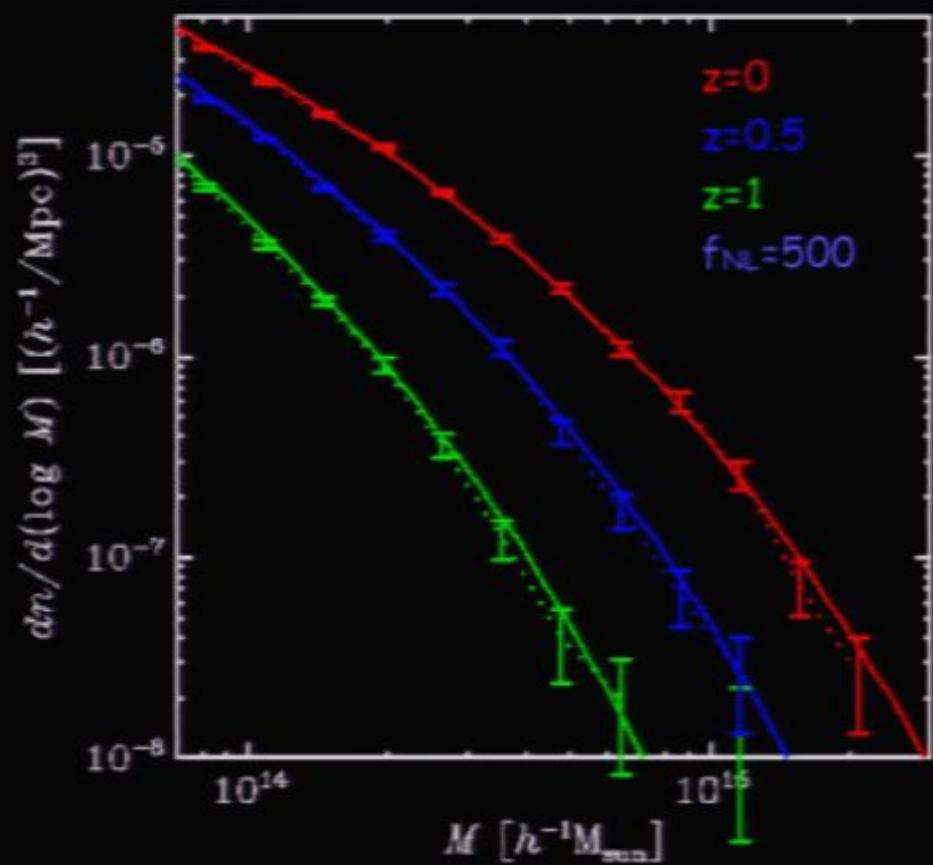
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- We'd expect the mean of the PDF to be shifted by  $\Delta(\log M) = f_{NL}$
- We find that a good fit is given by



$$\left[ \frac{\bar{M}_f}{M_0} \right] - 1 = 6 \cdot 10^{-5} f_{NL} \sigma_8 \sigma(M_0, z)^{-2}$$

$$\sigma \left( \left[ \frac{\bar{M}_f}{M_0} \right] - 1 \right) = 0.012 (f_{NL} \sigma_8)^{0.4} \sigma(M_0, z)^{-0.5}$$

# Mass function from N-body simulation and our fitting formula



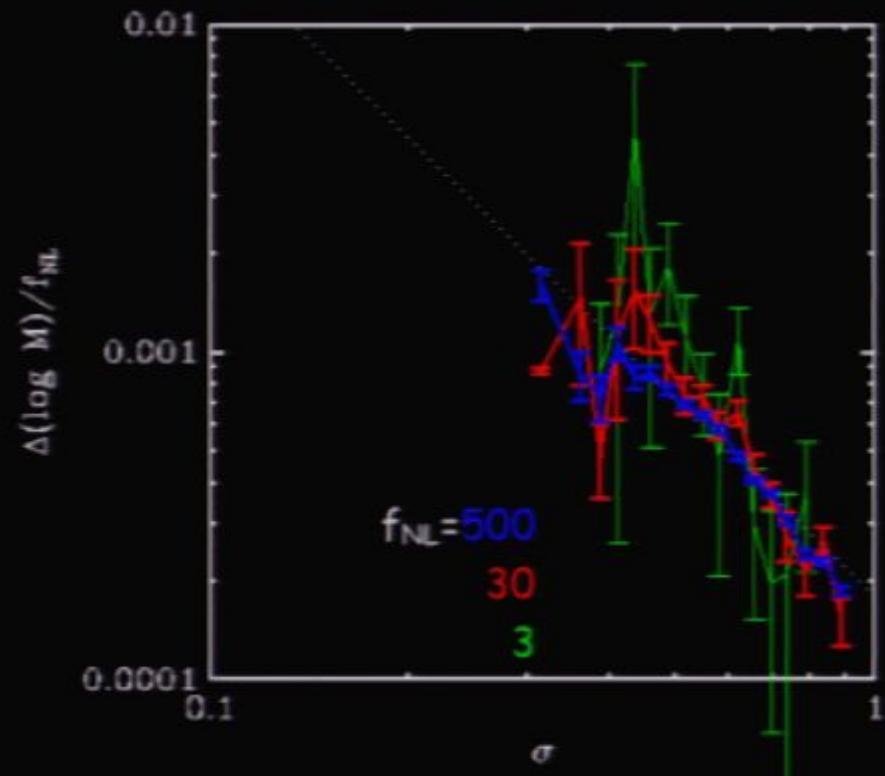
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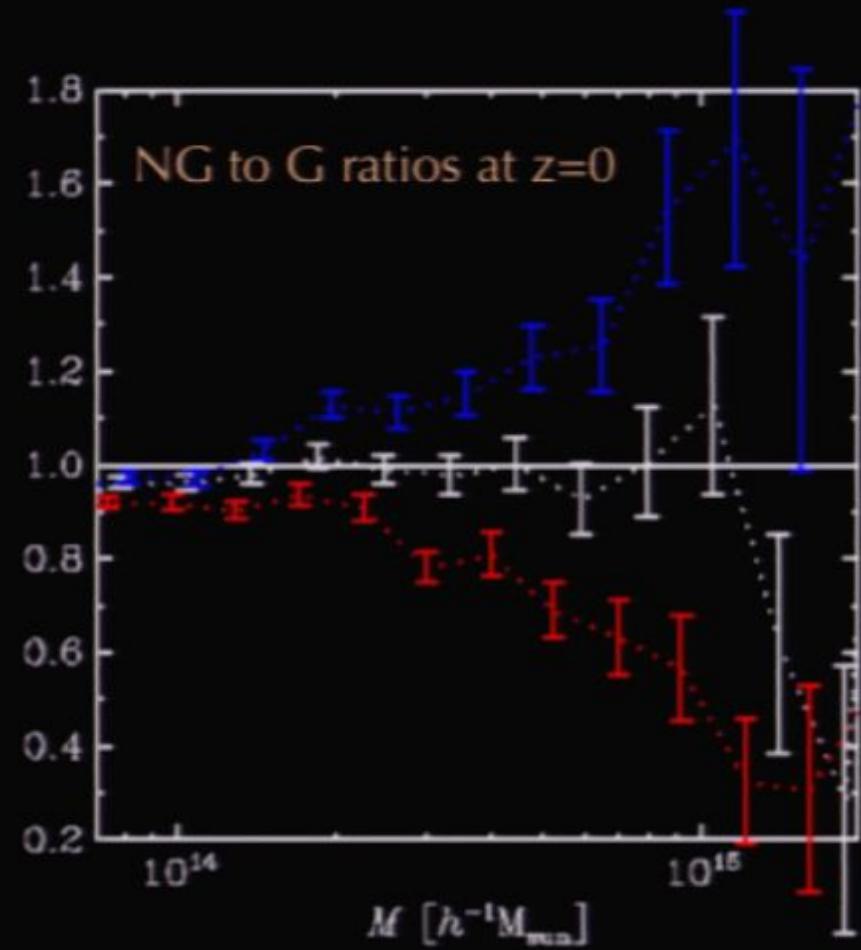
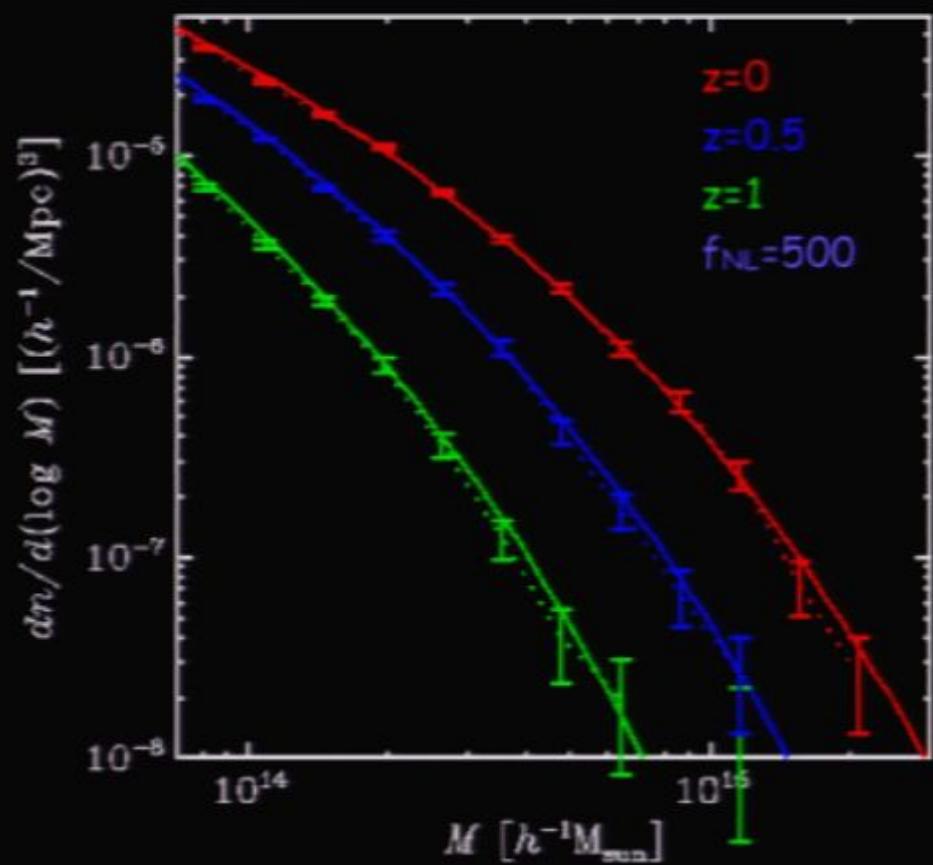
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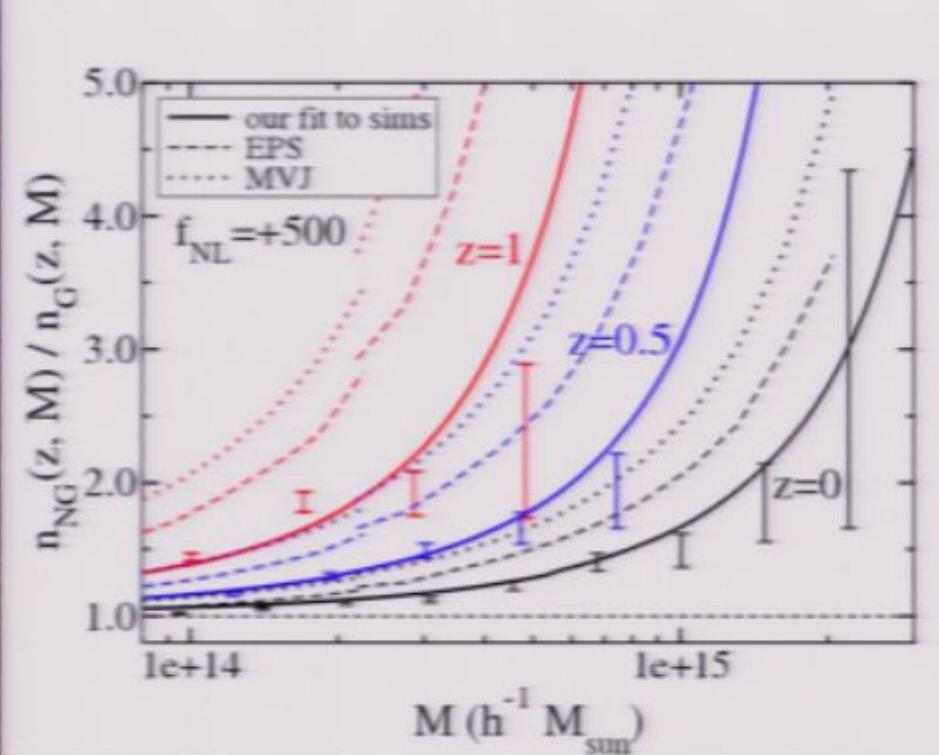
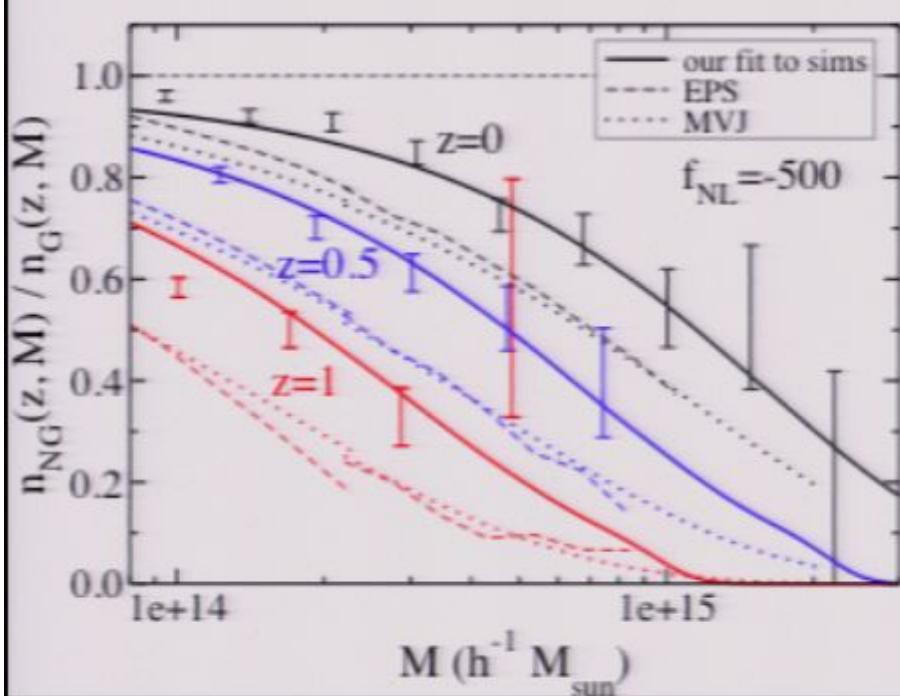
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off by O(100%) wrt truth



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than evaluate Extended Press-Schechter  $n(M)$ !

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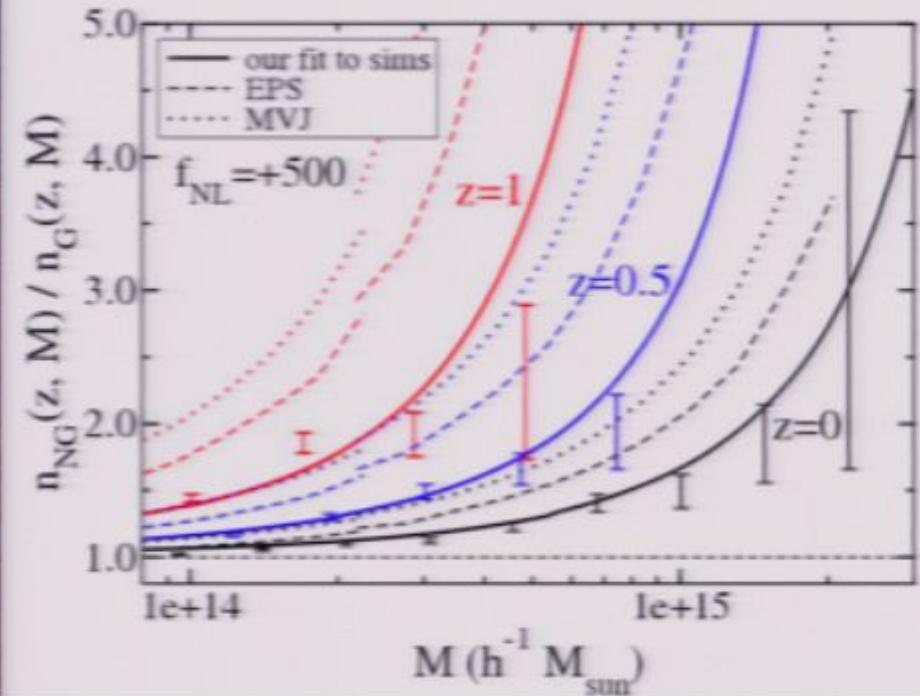
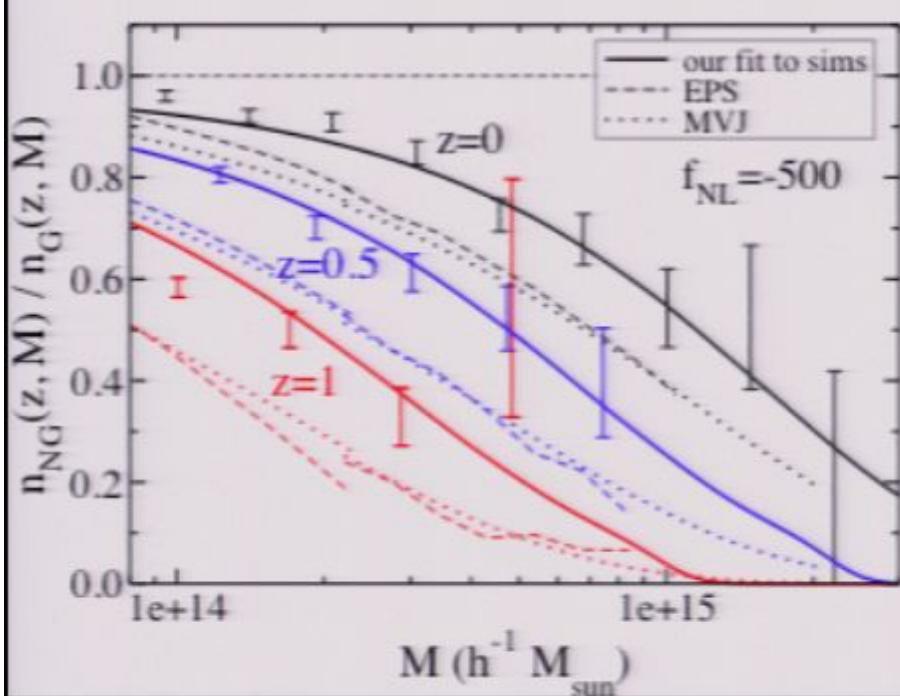
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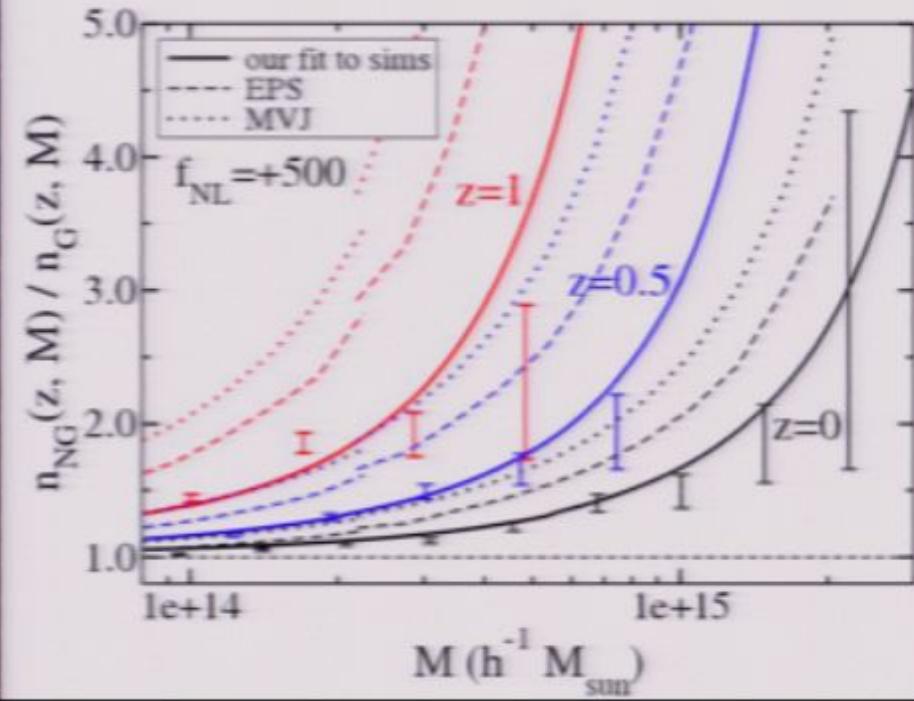
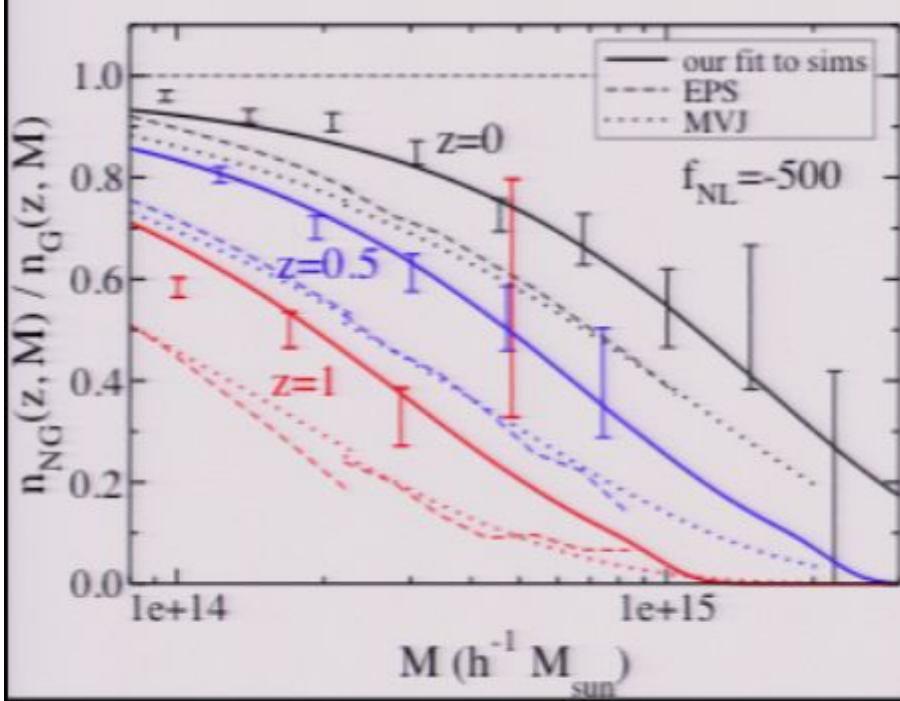
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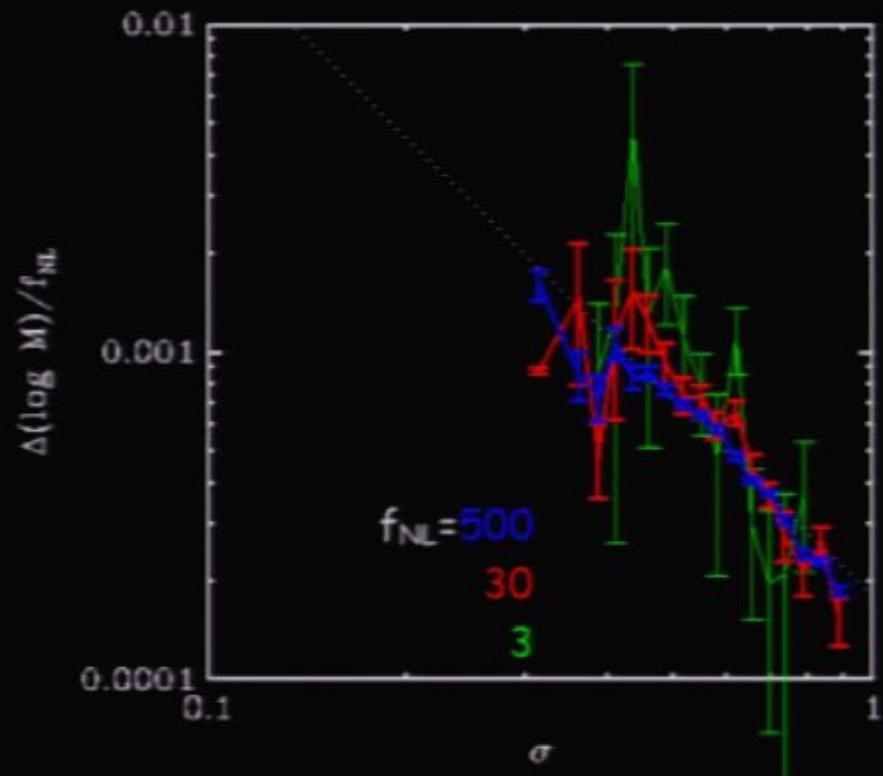
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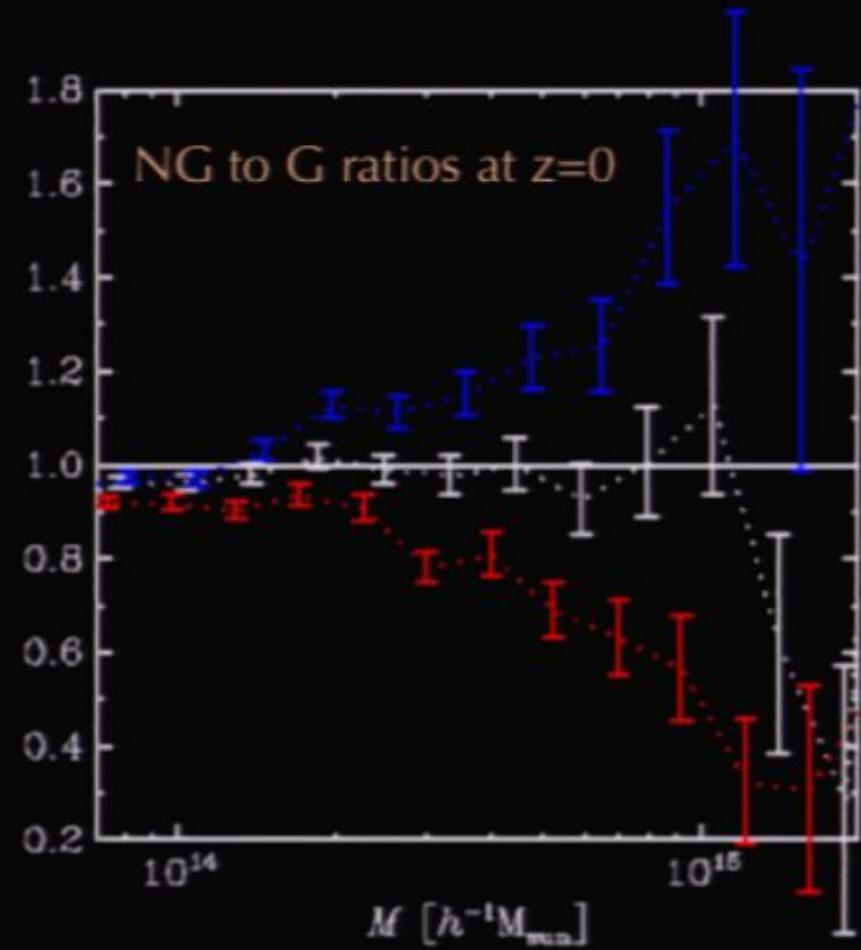
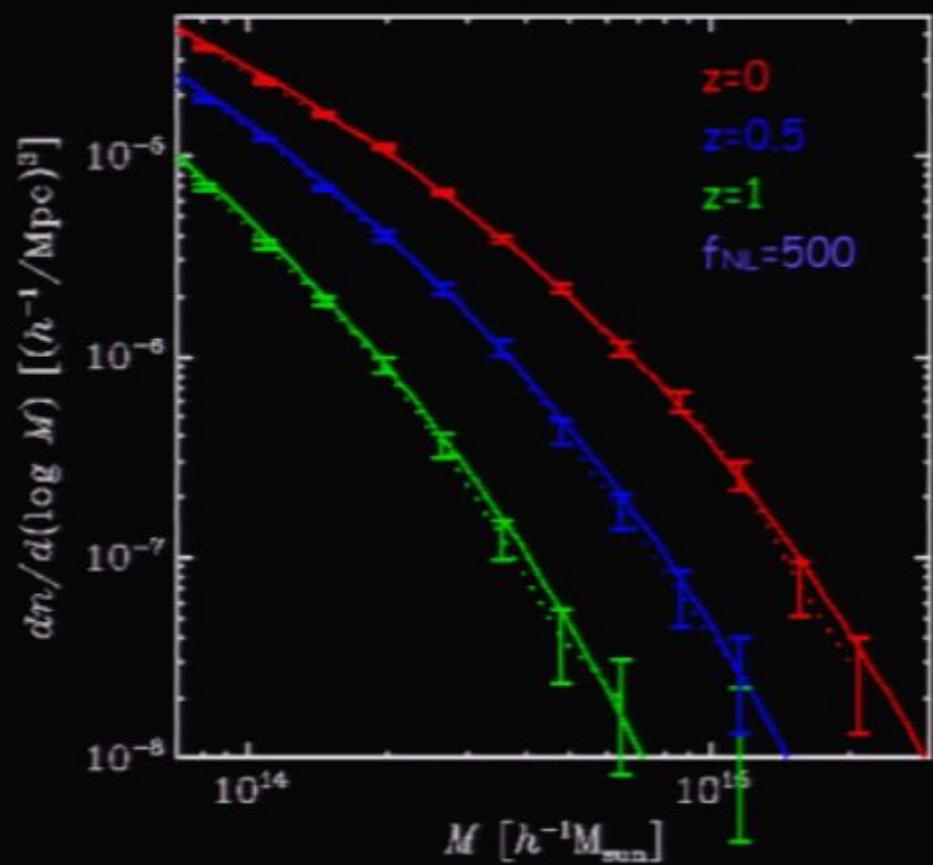
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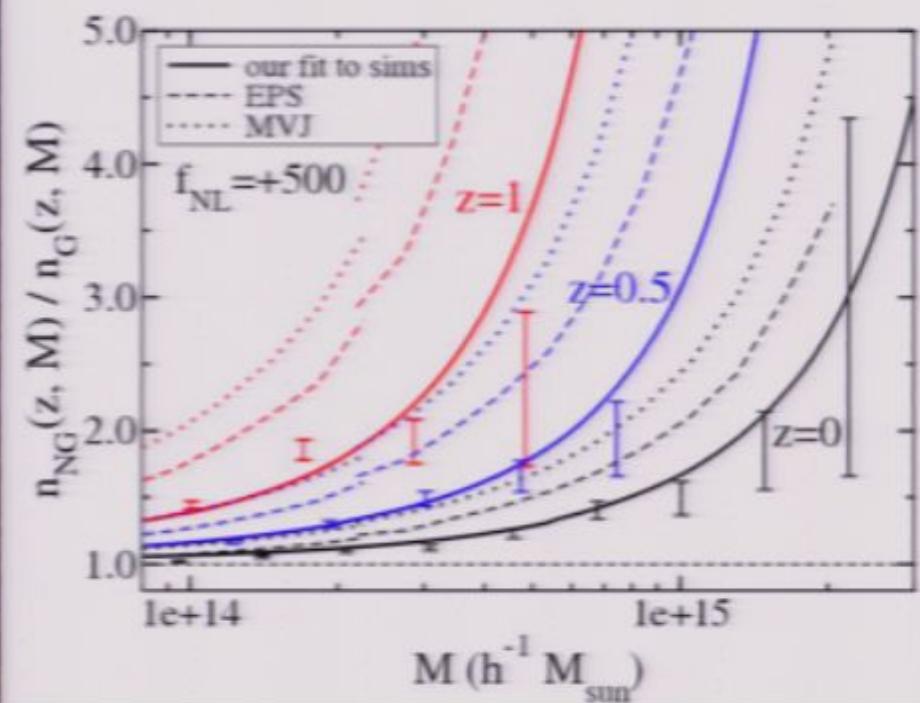
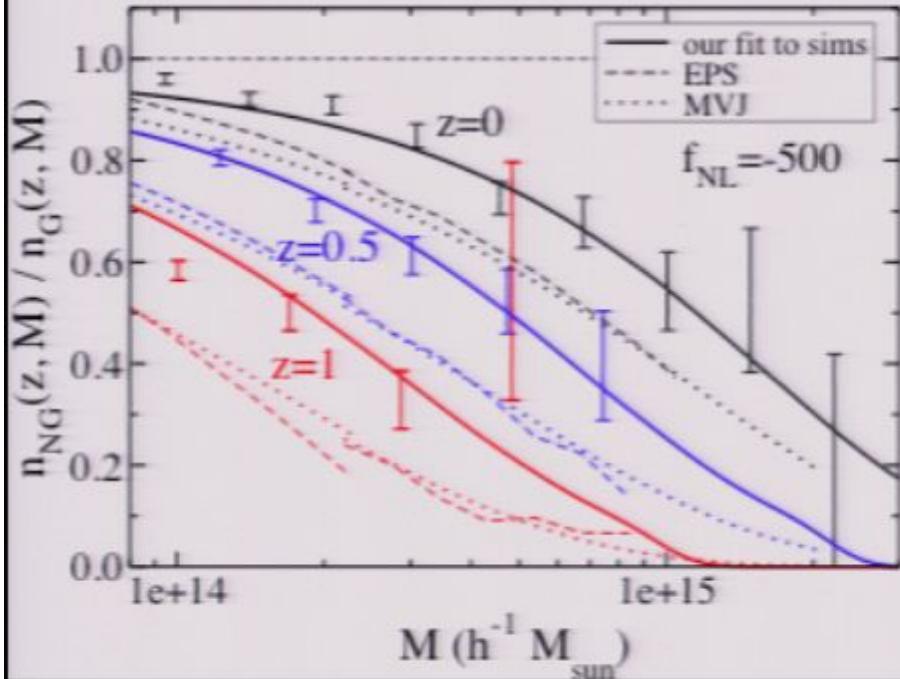
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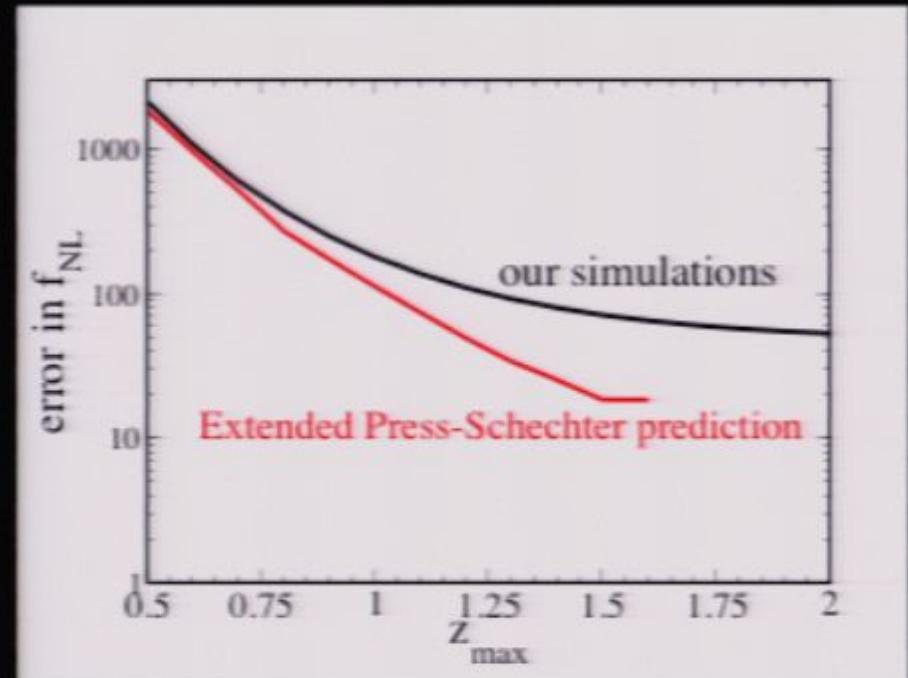
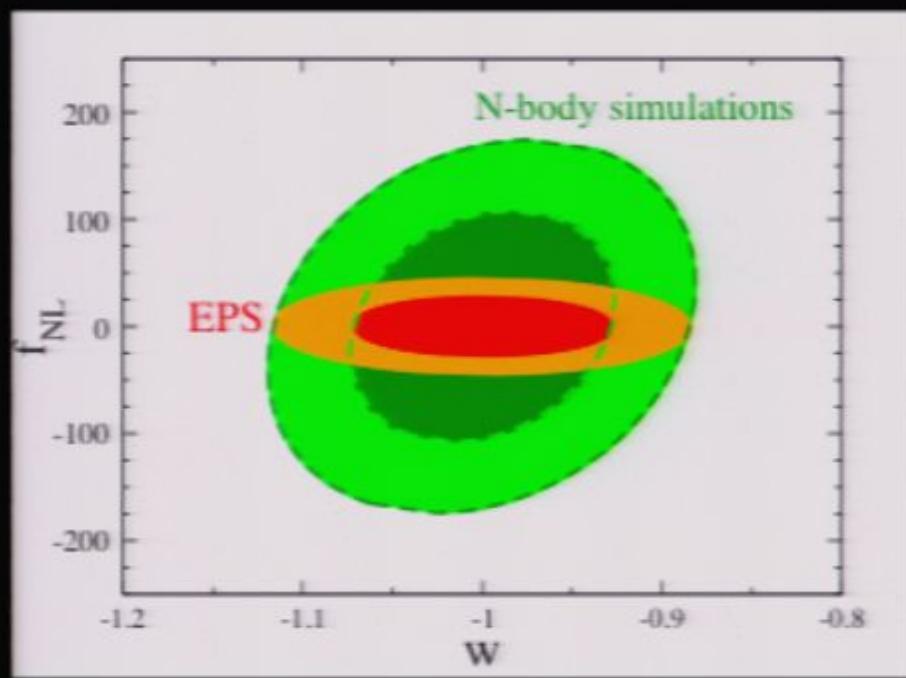


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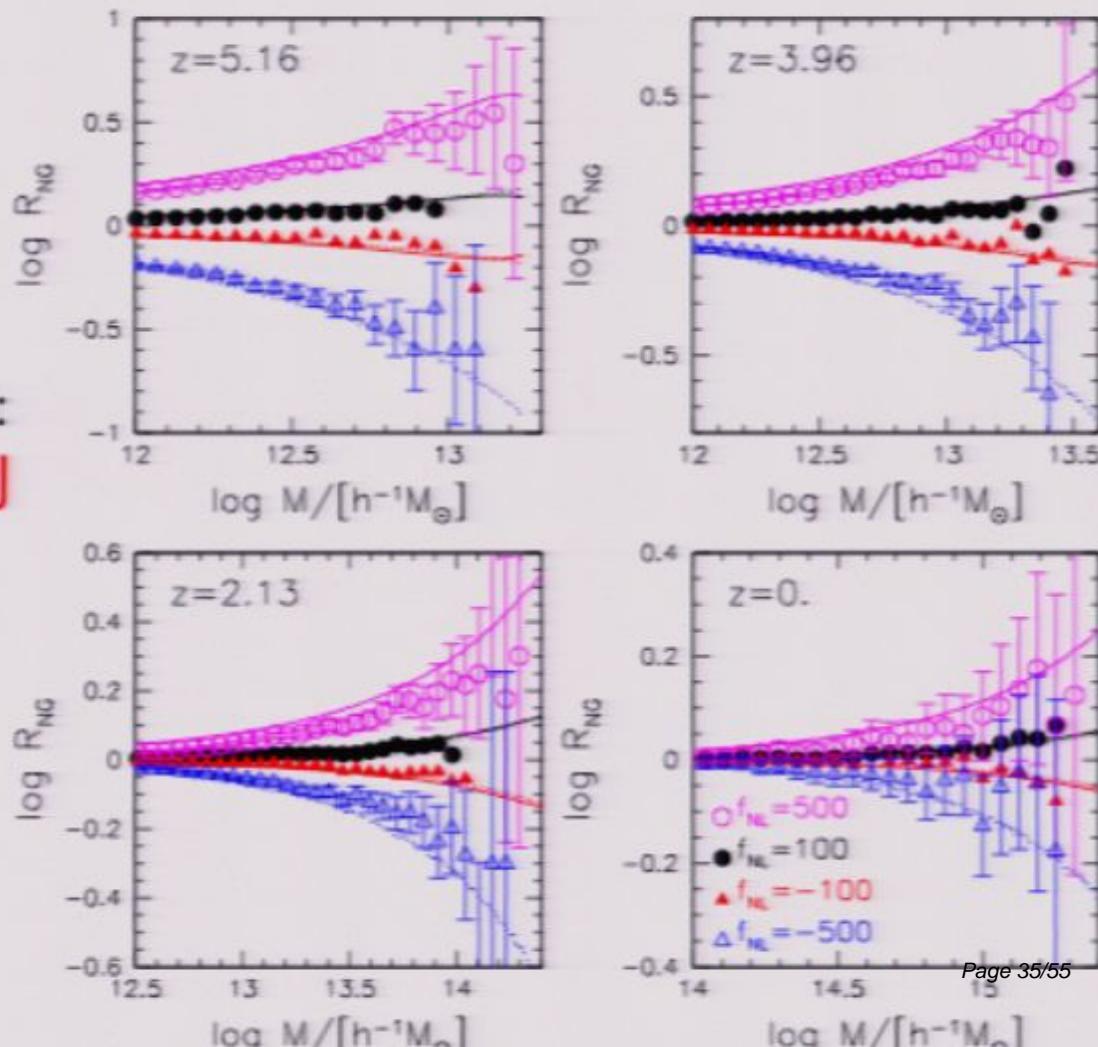
# Cosmological constraints - dark energy and NG



SPT-type survey, ~7,000 clusters, 4000 sq.deg.,  $0.1 < z < 1.5$   
Planck prior

# Comparison to other (numerical) work

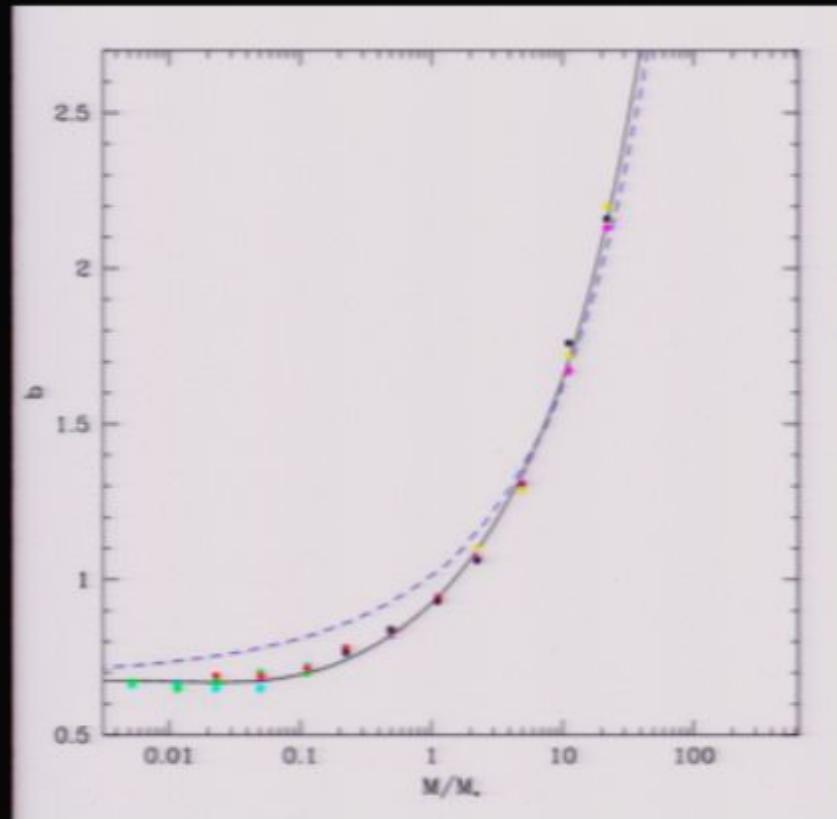
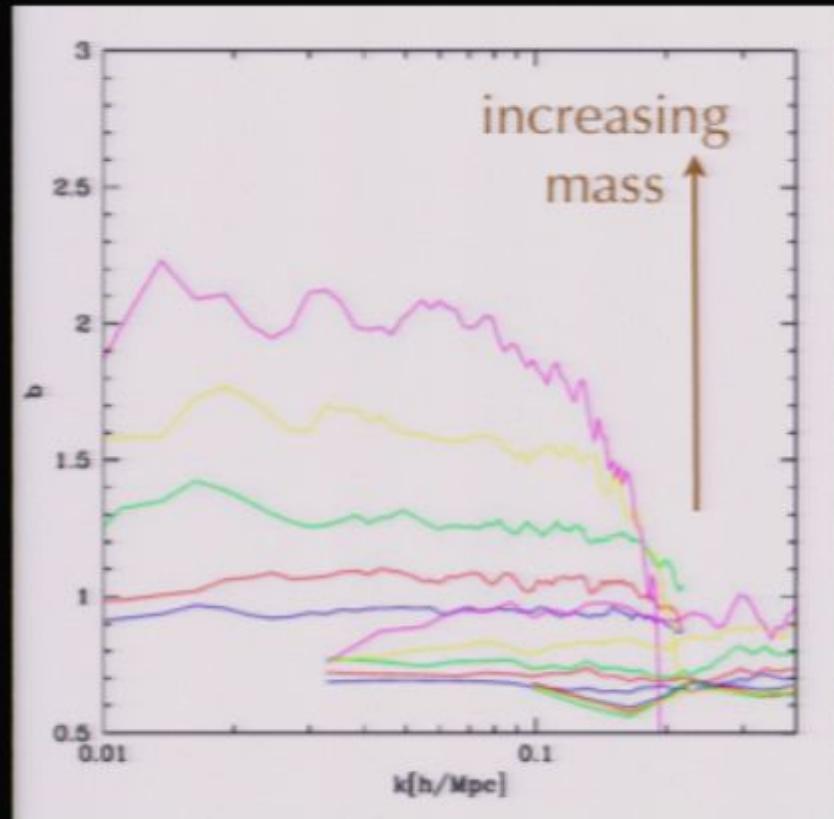
- I) Kang, Norberg & Silk (astro-ph/0701131):  
claim much *bigger* discrepancy with MVJ,  
but: their simulations are  $128^3$  (insufficient, as they note)



- 2) Grossi et al (arXiv:0707.2516):  
claim perfect agreement with MVJ

# Bias of dark matter halos - Gaussian case

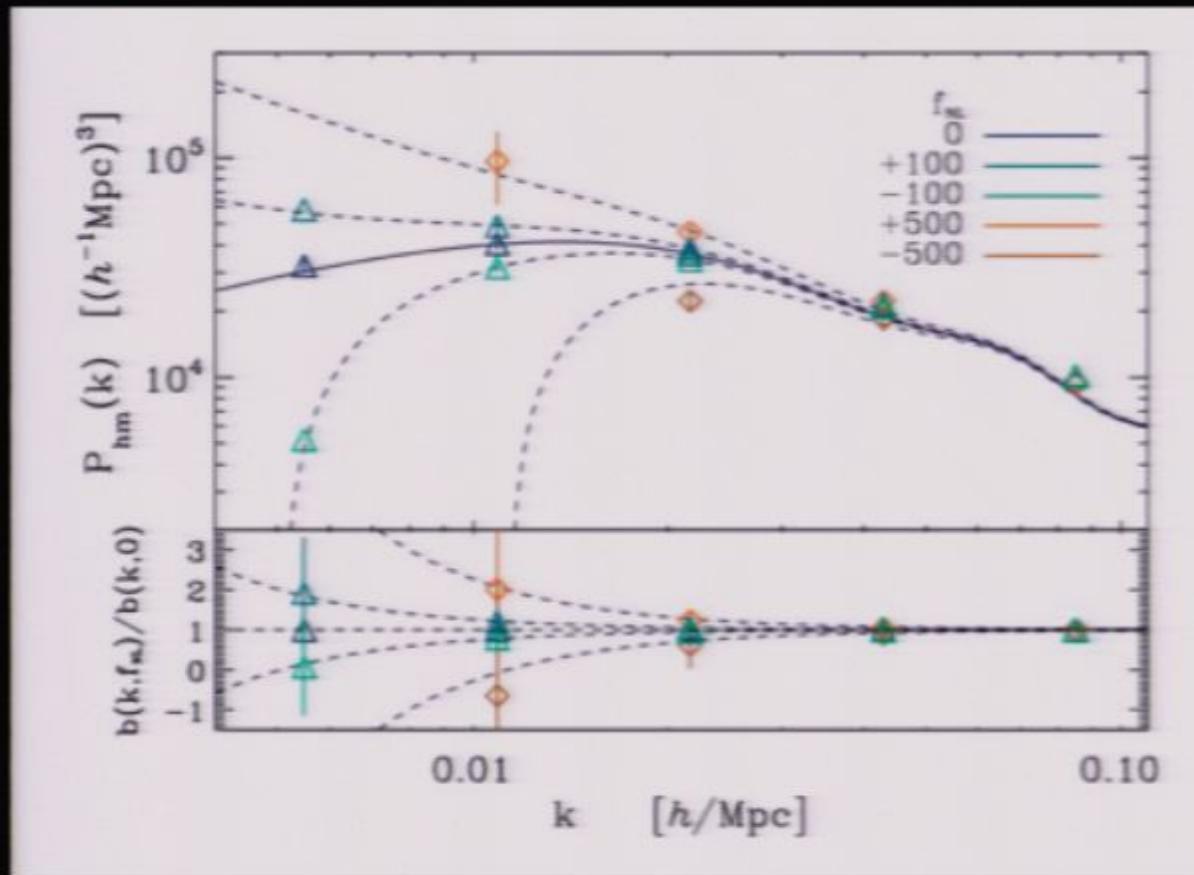
$$b \equiv \delta_h / \delta_{\text{DM}}$$



Seljak & Warren 2006

Simulations and theory both say: large-scale bias is scale-independent  
(theorem if halo abundance is function of local density)

# Scale dependence of NG halo bias!



- Strong scale dependence in linear regime
- Good agreement between theory and simulation
- $512^3$  ( $1024^3$ ) particle simulations with box size 800 (1600)  $\text{Mpc}/h$

## Halo clustering with NG: Analytic estimates

$$\Phi_{\text{NG}} = \phi + f_{\text{NL}}(\phi^2 - \langle \phi^2 \rangle)$$

Then

$$\nabla^2 \Phi_{\text{NG}} = \nabla^2 \phi + 2f_{\text{NL}} (\phi \nabla^2 \phi + |\nabla \phi|^2)$$

We know the statistics of all terms, so we can compute anything, e.g.

Skewness

$$S_3 = \frac{\langle \delta_{\text{NG}}^3 \rangle}{\langle \delta_{\text{NG}}^2 \rangle^2} = 6f_{\text{NL}} \frac{\langle \phi \delta \rangle}{\sigma_\delta^2}$$

And in particular

$$\delta_{\text{NG}} = \delta(1 + 2f_{\text{NL}}\phi)$$

## Halo clustering with NG: Analytic estimates

Definition of bias:  $\delta_h = b_L \delta$

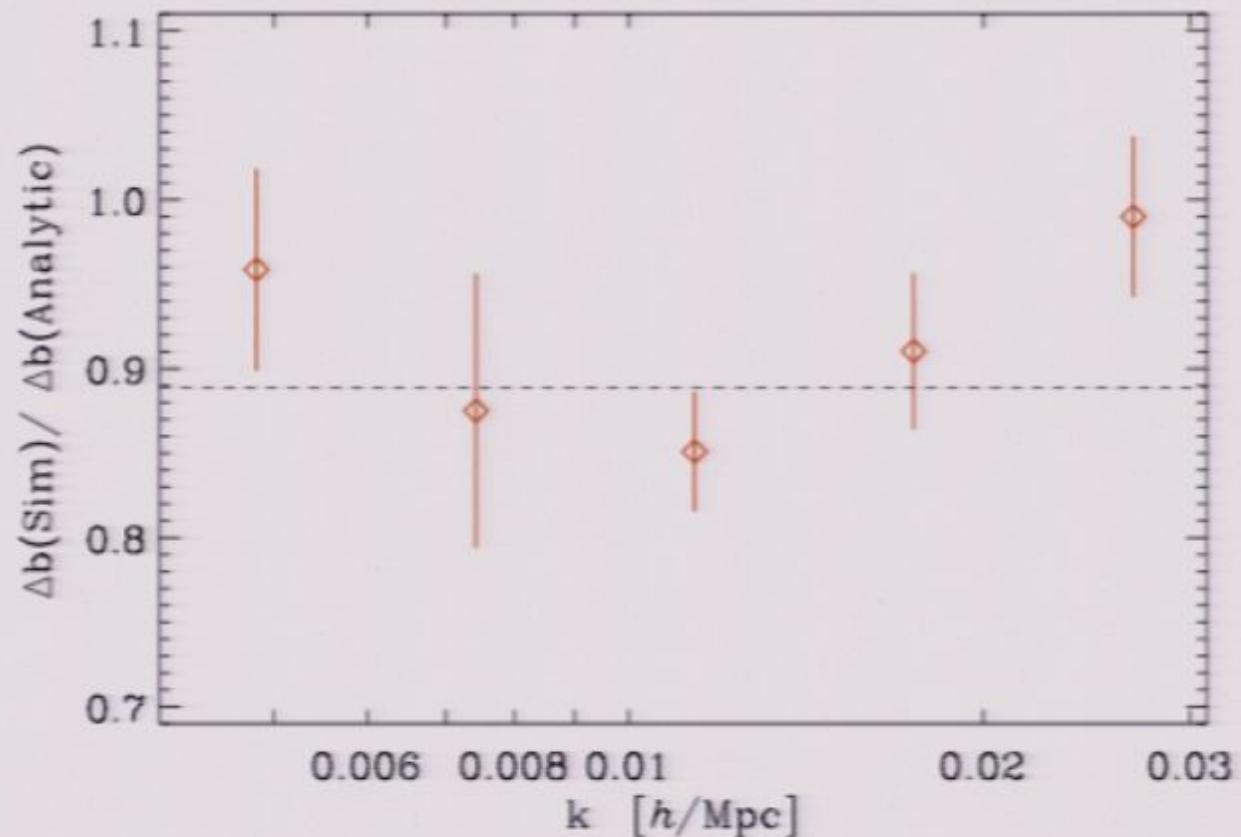
With NG, for peaks:  $\delta \rightarrow \delta + 2f_{\text{NL}}\phi_p\delta_c$

Assuming  $\delta_h \rightarrow (b_L + \Delta b(k))\delta$

and using Poisson equation it follows that

$$\Delta b(k) = 2b_L f_{\text{NL}} \delta_{\text{crit}} \frac{3\Omega_M}{2ar_H^2 k^2}$$

## Analytic and numerical results agree



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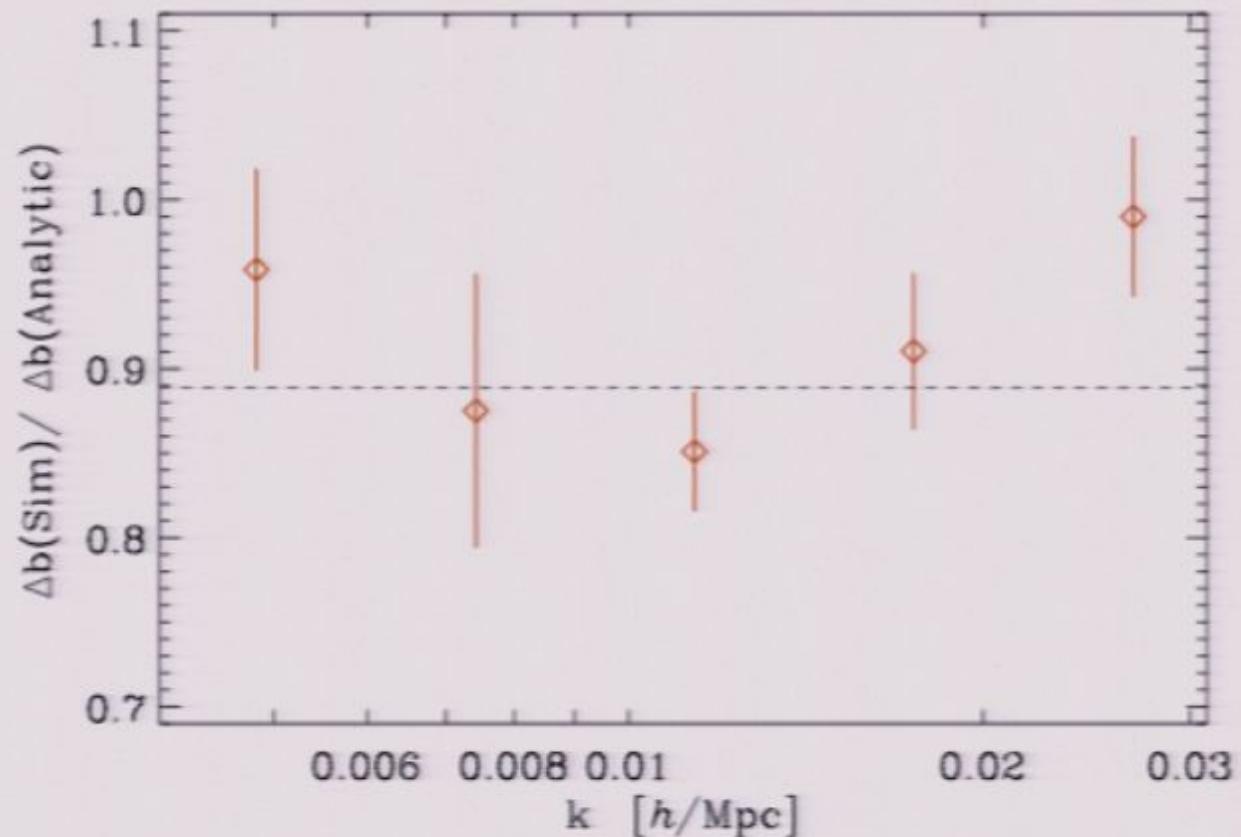
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## NG from measurements from the b(k) aspect of large-scale structure

- Numerous cosmological probes, such as the baryon acoustic oscillations (BAO) or probes of Integrated Sachs-Wolfe effect (galaxy-CMB cross-corr) can be used to measure  $b(k)$
- The effect (going as  $k^{-2}$ ) provides a fairly unique signature and an easy target
- Back-of-envelope calculation shows  $\sigma(f_{NL}) \sim 10$  for a typical LRG survey out to  $z \sim 0.7$  and  $1/4$  sky
- Current ISW observations should reach  $\sigma(f_{NL}) \sim 100$ ; and future ones a factor of  $\sim 3$  better (see also Afshordi talk)
- Need to estimate these numbers more accurately

# Conclusions

- Cosmological models with (local) primordial NG lead to significant scale dependence of halo bias; theory and simulations appear to be in remarkable agreement on this
- Therefore, LSS probes (baryon oscillations, galaxy-CMB cross-correlations, etc) are likely to lead to constraints on NG stronger than previously thought
- Back-of-envelope calculation shows  $\sigma(fNL) \sim 10$ , more accurate estimates needed; watch out for upcoming constraints using current data (Afshordi & Tolley; also Padmanabhan & Slosar)
- We also provided accurate, easy-to-use fitting formula for cluster abundance with NG
- What about other NG models?

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# Measurements from the scale structure

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more accurately

## Halo clustering with NG: Analytic estimates

$$\Phi_{\text{NG}} = \phi + f_{\text{NL}}(\phi^2 - \langle \phi^2 \rangle)$$

Then

$$\nabla^2 \Phi_{\text{NG}} = \nabla^2 \phi + 2f_{\text{NL}} (\phi \nabla^2 \phi + |\nabla \phi|^2)$$

We know the statistics of all terms, so we can compute anything, e.g.

Skewness

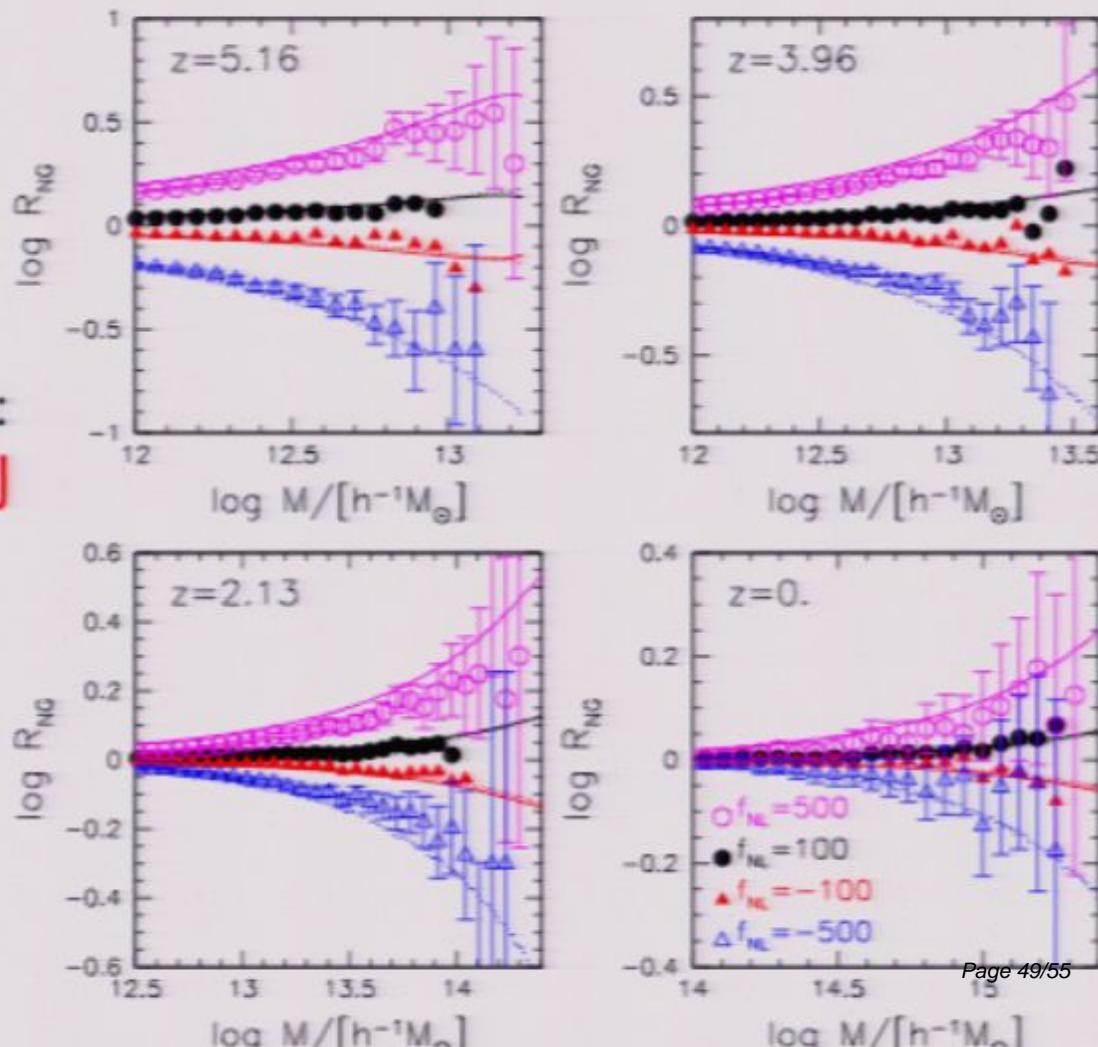
$$S_3 = \frac{\langle \delta_{\text{NG}}^3 \rangle}{\langle \delta_{\text{NG}}^2 \rangle^2} = 6f_{\text{NL}} \frac{\langle \phi \delta \rangle}{\sigma_\delta^2}$$

And in particular

$$\delta_{\text{NG}} = \delta(1 + 2f_{\text{NL}}\phi)$$

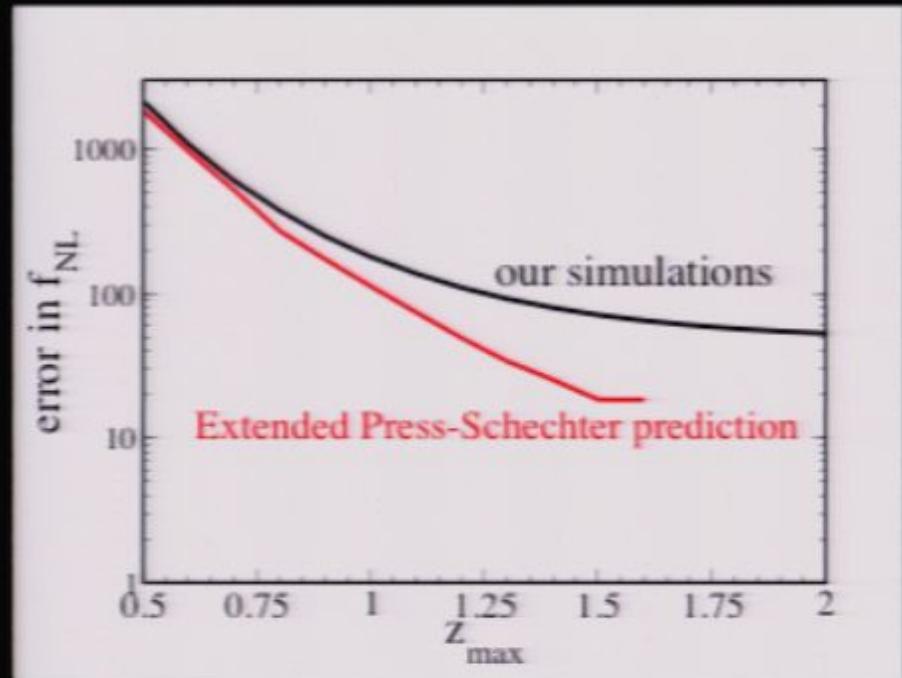
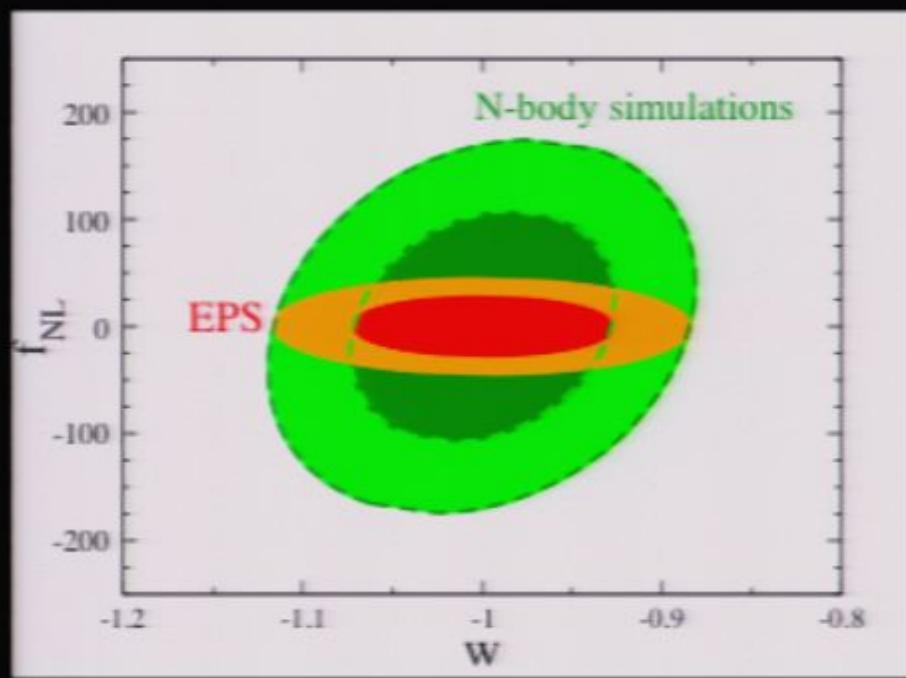
# Comparison to other (numerical) work

- 1) Kang, Norberg & Silk (astro-ph/0701131):  
claim much *bigger* discrepancy with MVJ,  
but: their simulations are  $128^3$  (insufficient, as they note)



- 2) Grossi et al (arXiv:0707.2516):  
claim perfect agreement with MVJ

# Cosmological constraints - dark energy and NG



SPT-type survey, ~7,000 clusters, 4000 sq.deg.,  $0.1 < z < 1.5$   
Planck prior

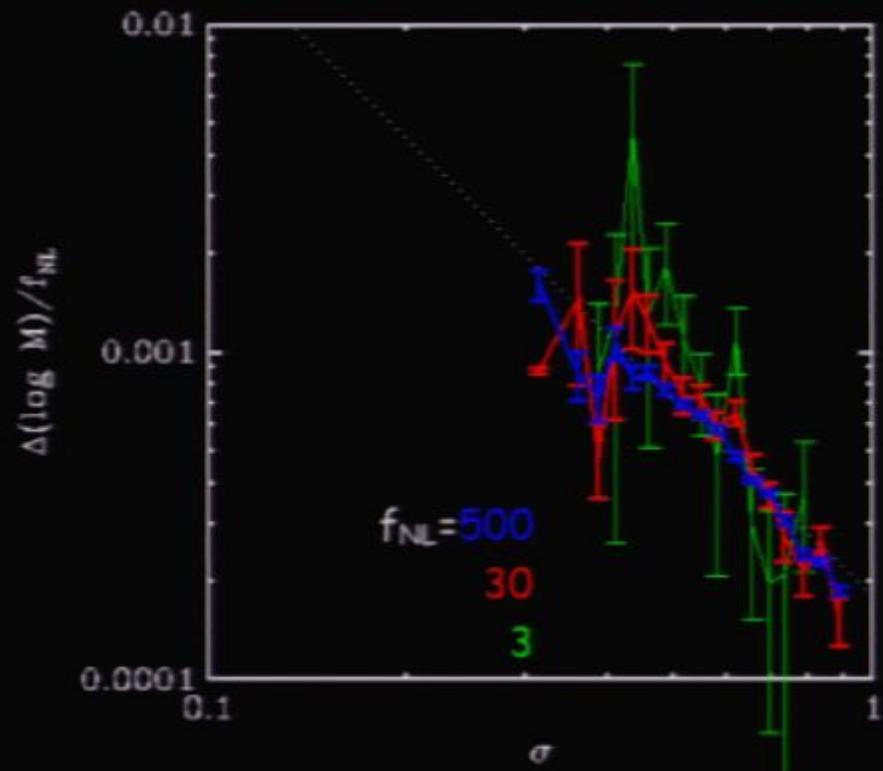
Recall, this is just from the cluster counts;  
CMB provides stronger constraints

# Towards a fitting function

- If the mapping  $M_0 \rightarrow M_f$  is described by a PDF  $dP/dM_f(M_0)$ , then the non-gaussian mass function is a convolution over the (known) gaussian mass function

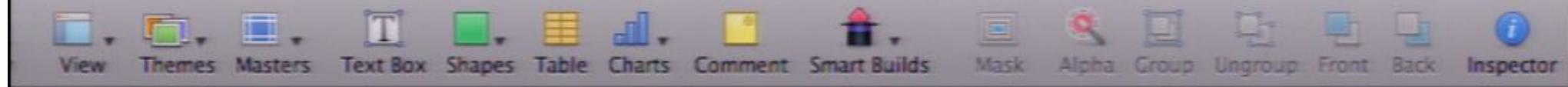
$$\frac{dN}{dM} = \int \frac{dP(M_f|M_0)}{dM_f} \frac{dN}{dM_0} dM_0$$

(e.g. Jenkins)



- We thus aim at fitting the mean and rms of  $\Delta(\log M)(z)$
- The simplest thing to do is to consider a... Gaussian...
- We'd expect the mean of the PDF to be shifted by  $\Delta(\log M) = f_{NL}$
- We find that a good fit is given by

$$\left[ \frac{\bar{M}_f}{M_0} \right] - 1 = 6 \cdot 10^{-5} f_{NL} \sigma_8 \sigma(M_0, z)^{-2}$$
$$\sigma \left( \left[ \frac{\bar{M}_f}{M_0} \right] - 1 \right) = 0.012 (f_{NL} \sigma_8)^{0.4} \sigma(M_0, z)^{-0.5}$$

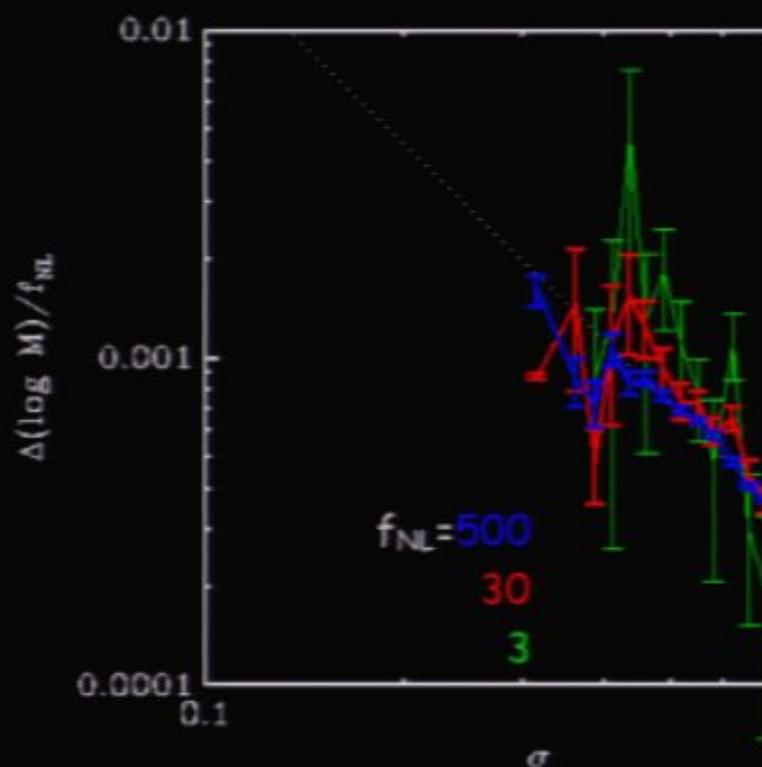


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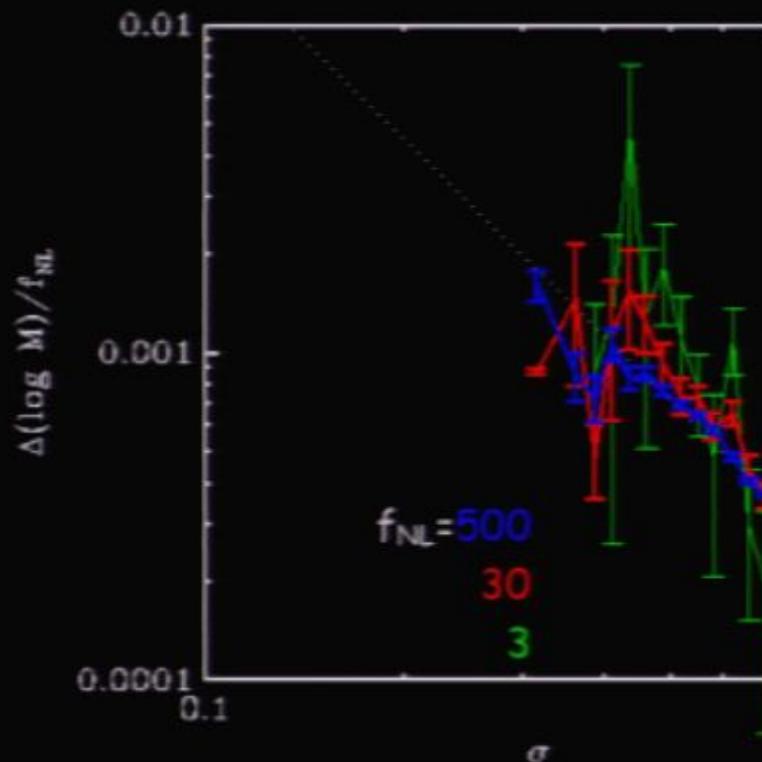
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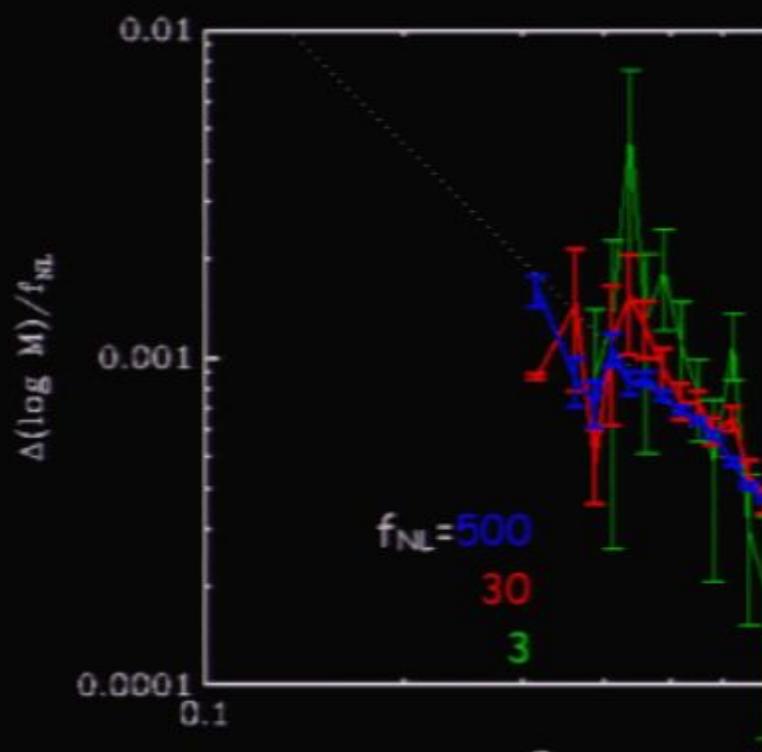
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