Title: Physics Beyond the Horizon

Date: Jan 25, 2008 03:00 PM

URL: http://pirsa.org/08010043

Abstract: The history of human knowledge is often highlighted by our efforts to explore beyond our apparent horizon. In this talk, I will describe how this challenge has now evolved into our quest to understand the physics at/beyond the cosmological horizon, some twenty orders of magnitude above ColumbusÂ's original goal. I also argue why inflationary paradigm predicts the existence of non-trivial physics beyond the cosmological horizon, and how we can use the Integrated Sachs-Wolfe effect in the Cosmic Microwave Background to probe this physics, including the nature of gravity and primordial non-gaussianity on the horizon scale.

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PHYSICS BEYOND THE HORIZON

Niayesh Afshordi



PERIMETER P

NSTITUTE FOR THEORETICAL PHYSICS

Perimeter Institute Observatory



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Ptolemy's 150 AD World Map, 109 cm

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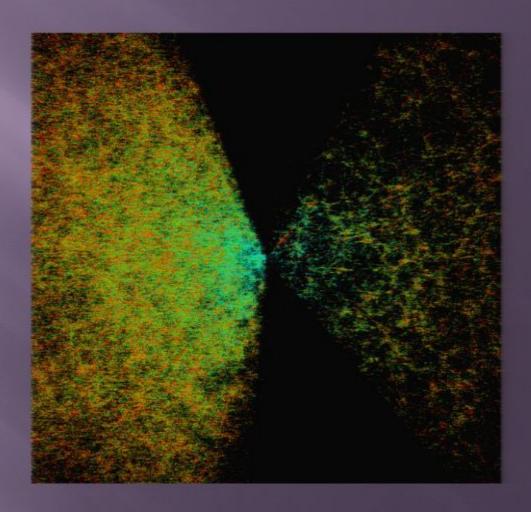
Ptolemy's 150 AD World Map, 109 cm

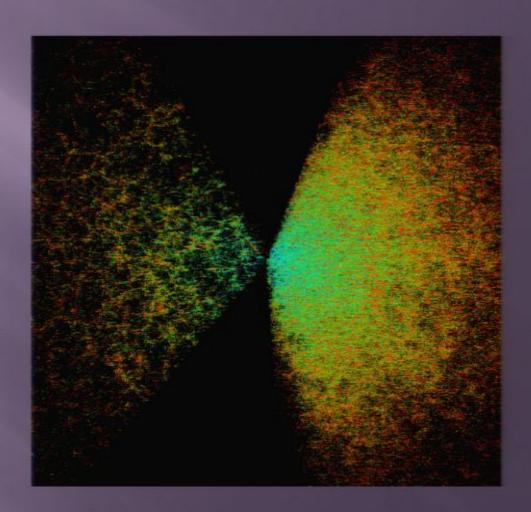
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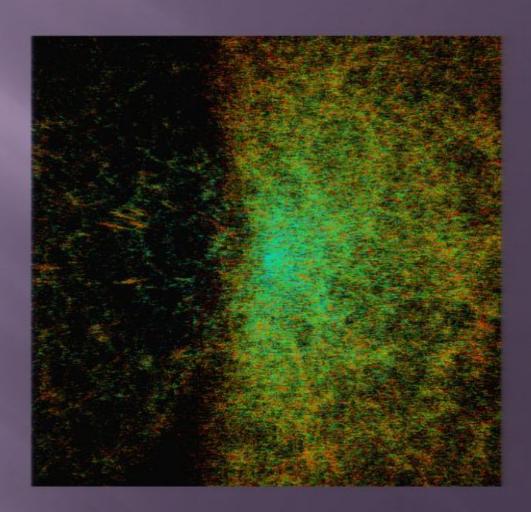


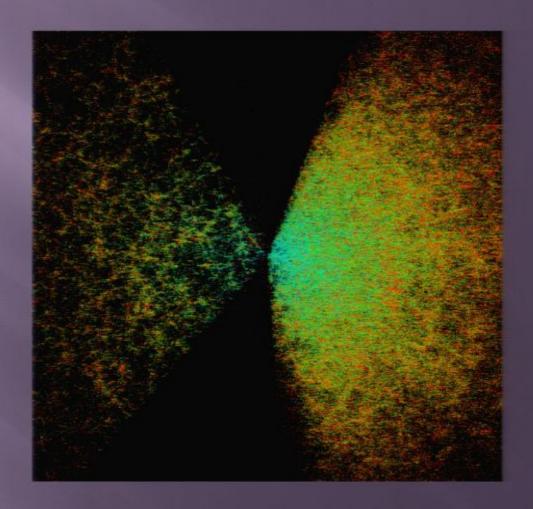
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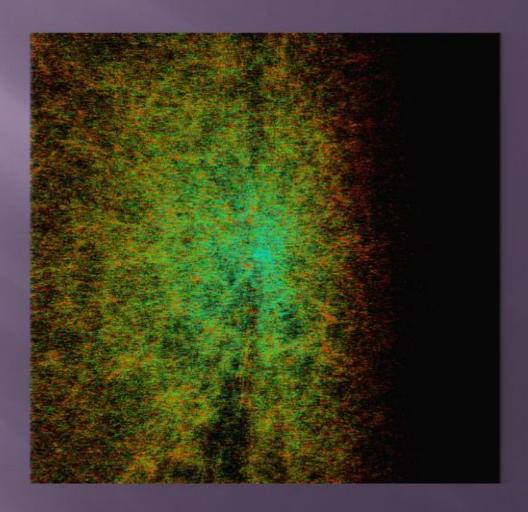
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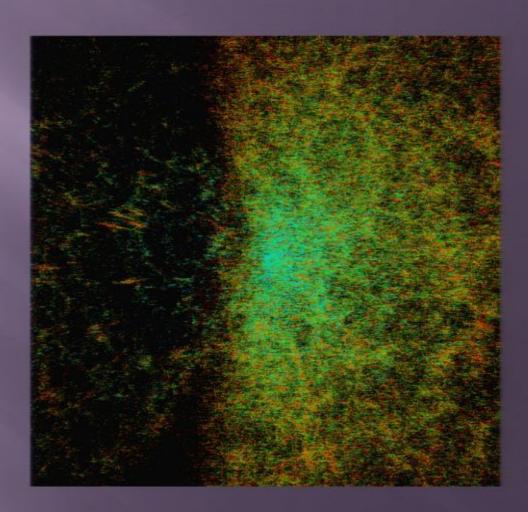


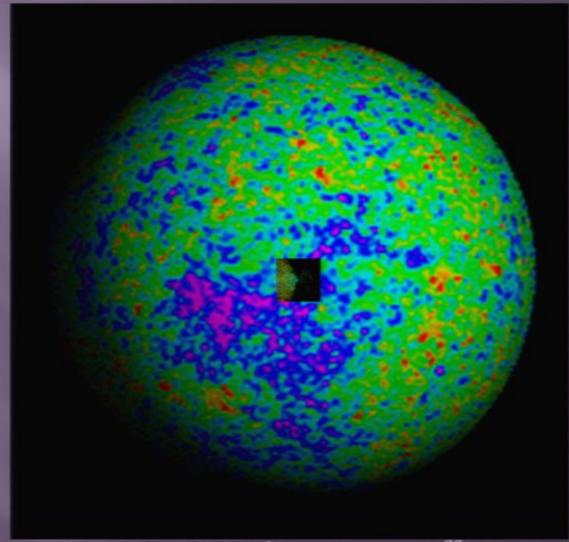












Wilkinson Microwave Anisotropy Probe: 10 Gpc ~ 10²⁹ cm (due to Max Tegmark)

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■ 4 x 10° cm (200 BC): Eratosthenes measured the earth circumference to 1% → (Columbus, Magellan 15th-16th century)

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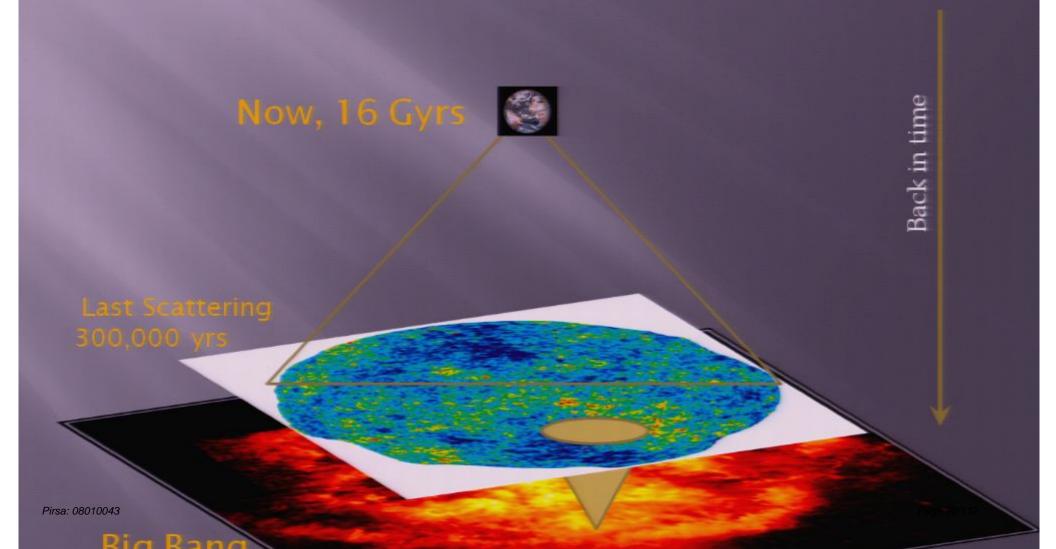
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Horizon Problem

At last scattering, causal horizon is much smaller than horizon today.

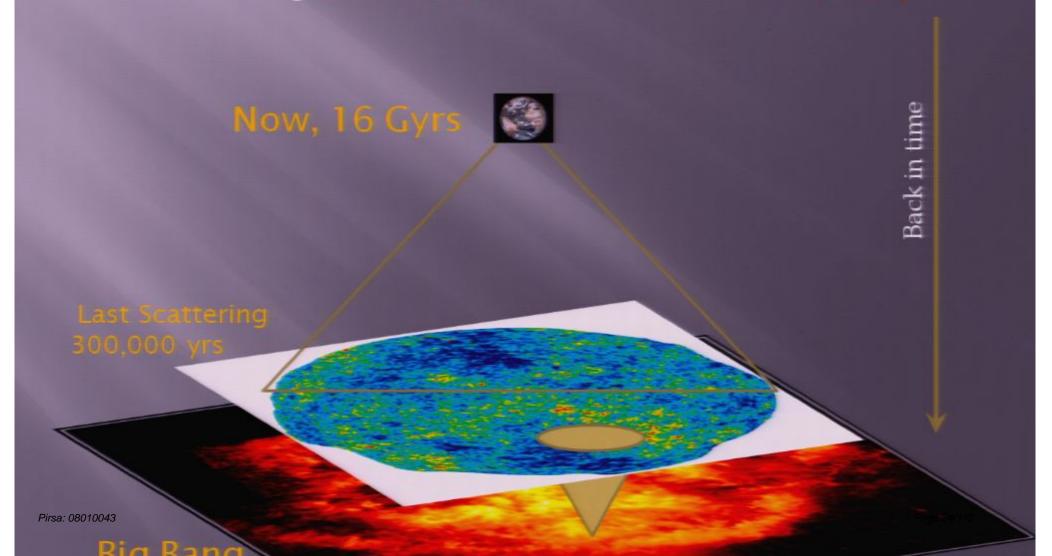


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- >10³² cm (2003): WMAP observations within the inflationary paradigm (Castro, Douspis, & Ferreira 2003)

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Horizon Problem

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- So, why is CMB isotropic to 1 part in 10⁵?

Now, 16 Gyrs 😂



Last Scattering

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Ria Rana

 A period of exponential expansion can inflate the causal horizon

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 - Quantum prturbations → nearly Gaussian linear fluctuations

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- A period of exponential expansion can inflate the causal horizon
- 60 e-foldings of inflation solves the horizon problem
- Also (correctly) predicts :
 - flatness
 - Quantum prturbations → nearly Gaussian linear fluctuations
- Other (model-dependent) features:
 - Non-gaussianity (secondary field, non-minimal kinetic term)
 - Gravitational waves (high energy scale inflation)

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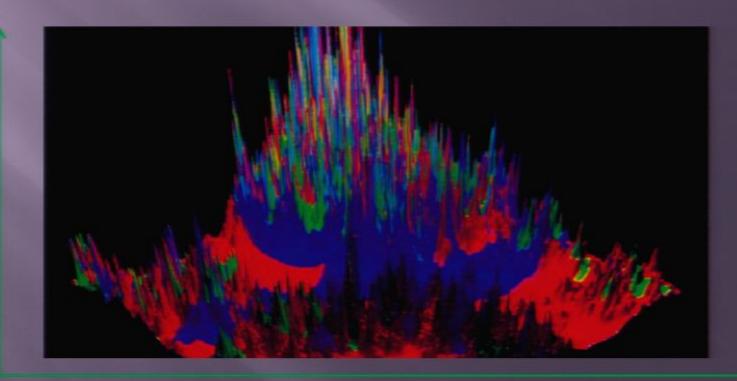
Inflation generically lasts much longer than 60 e-folds

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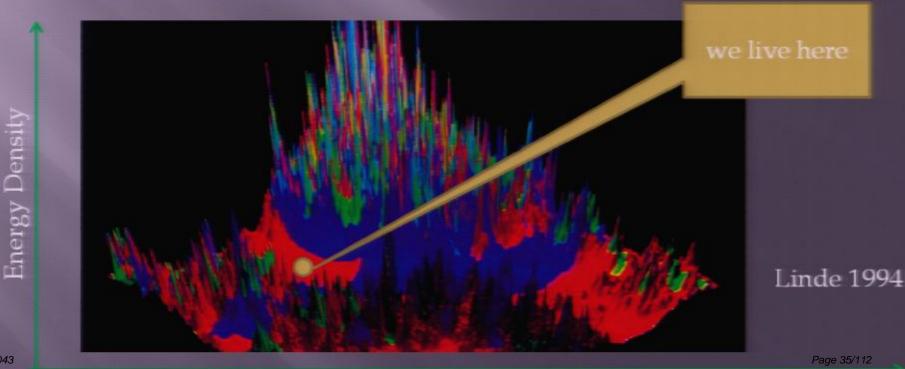
Linde 1994

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Energy Density

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Different Horizons

Inflationary horizon: how far light travels since Big Bang

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Different Horizons

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- Hubble horizon $\sim c/H$: important for linear perturbations

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Different Horizons

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- Hubble horizon ~ c/H: important for linear perturbations
- Today's Hubble horizon $\sim c/H_0$: scale of Dark Energy

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Different Horizons

- Inflationary horizon: how far light travels since Big Bang
- Hubble horizon ~ c/H: important for linear perturbations
- □ Today's Hubble horizon ~ c/H₀: scale of Dark Energy

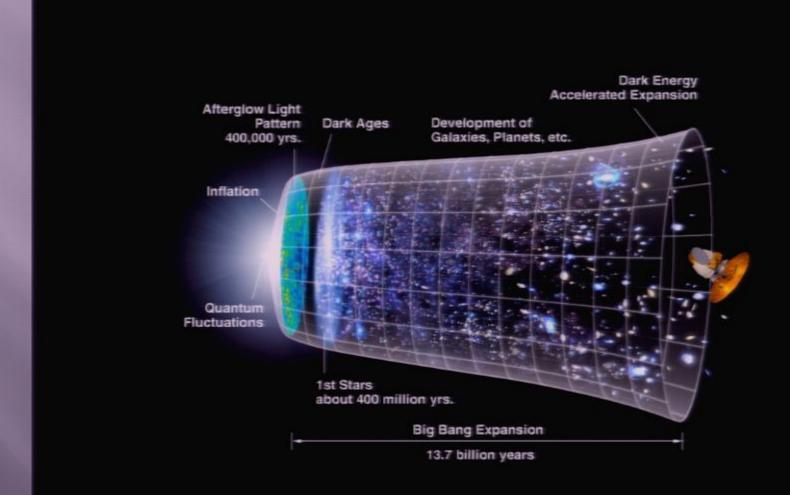
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Physics beyond the Horizon:

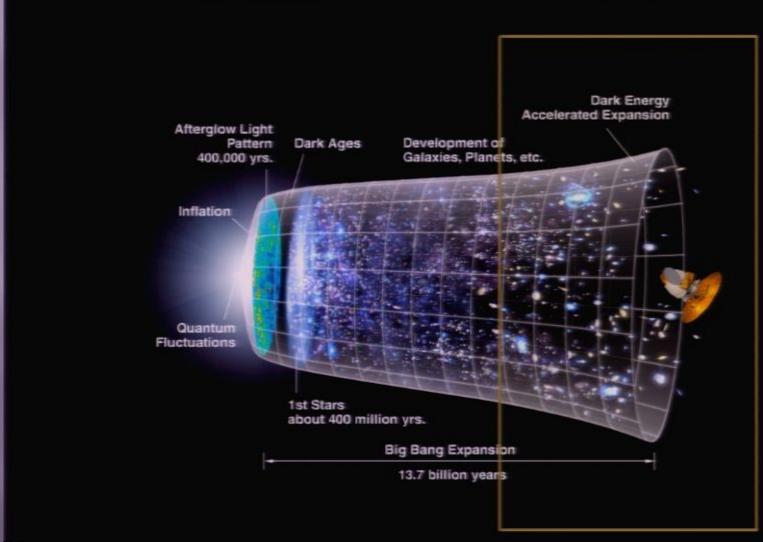
- Cosmic acceleration and the ISW effect
- Gravity on Horizon scale
- Statistics on Horizon scale

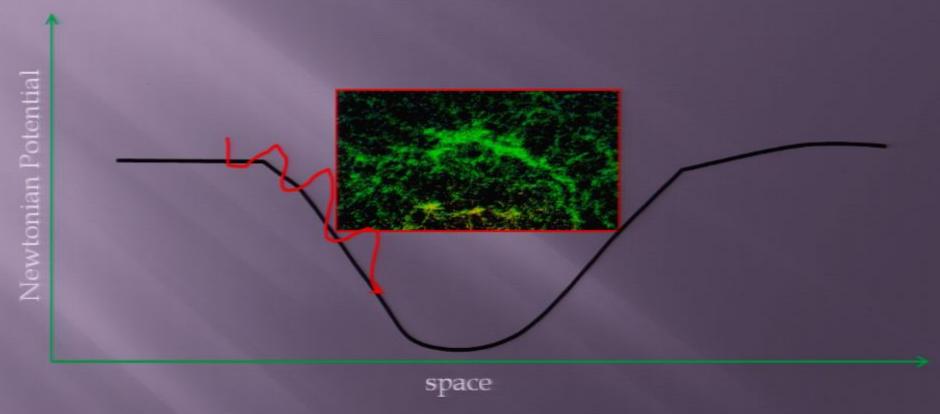
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Cosmic Timeline: Dark Energy Domination



Cosmic Timeline: Dark Energy Domination

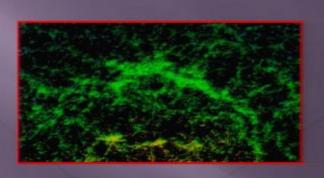




Accelerated Expansion results in decay of Newtonian potential
 ISW effect: decaying Newtonian potential causes secondary anisotropy

$$\delta T = 2T \int \dot{\Phi} dt$$

Newtonian Potential

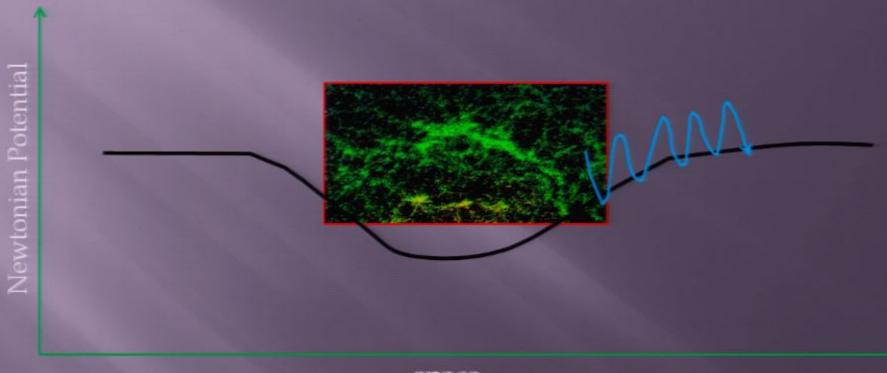


space

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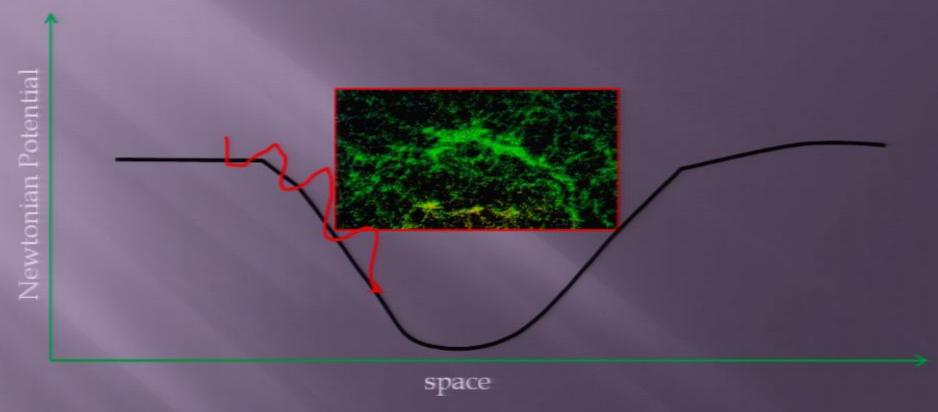
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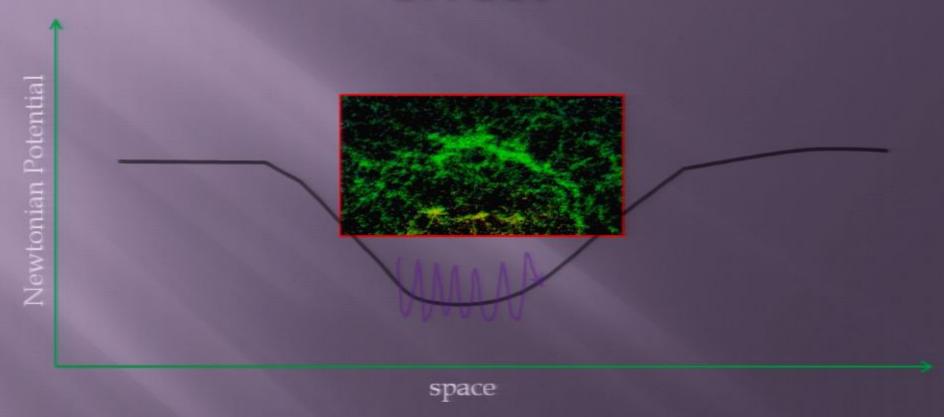


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in the CMB temperature:

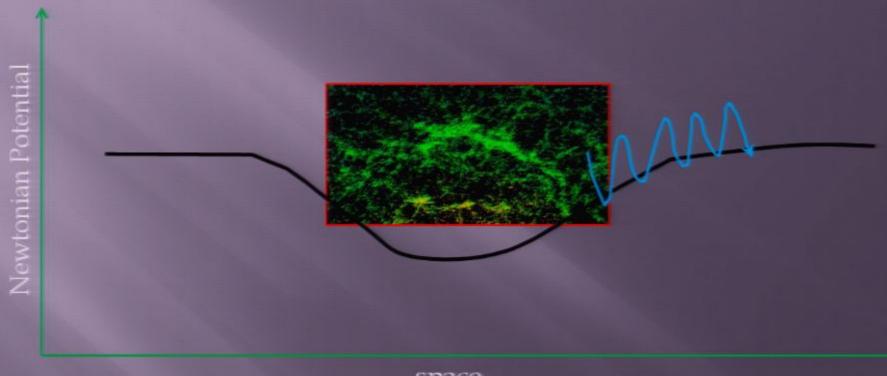
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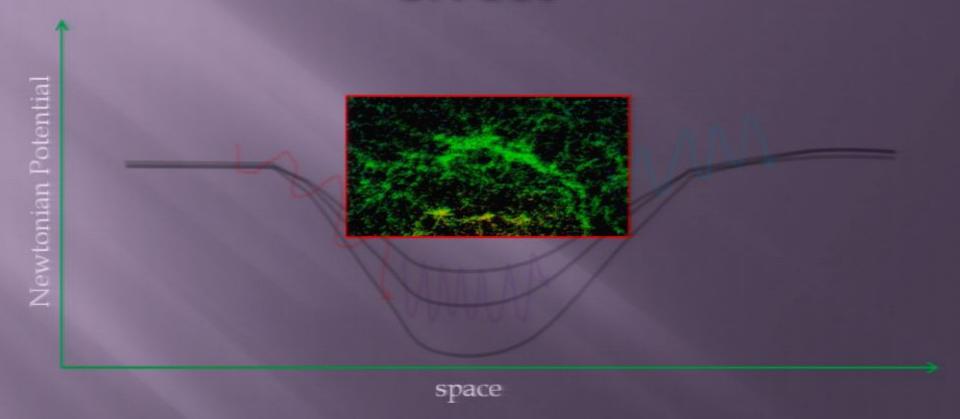
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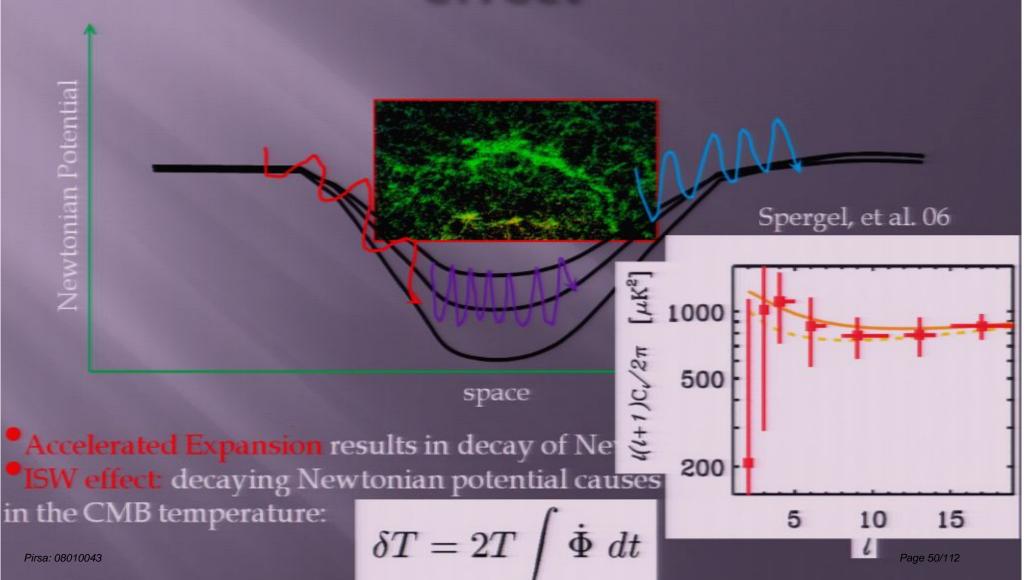
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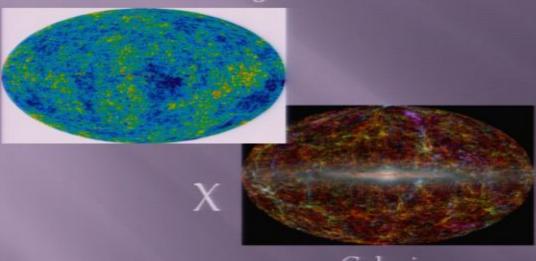


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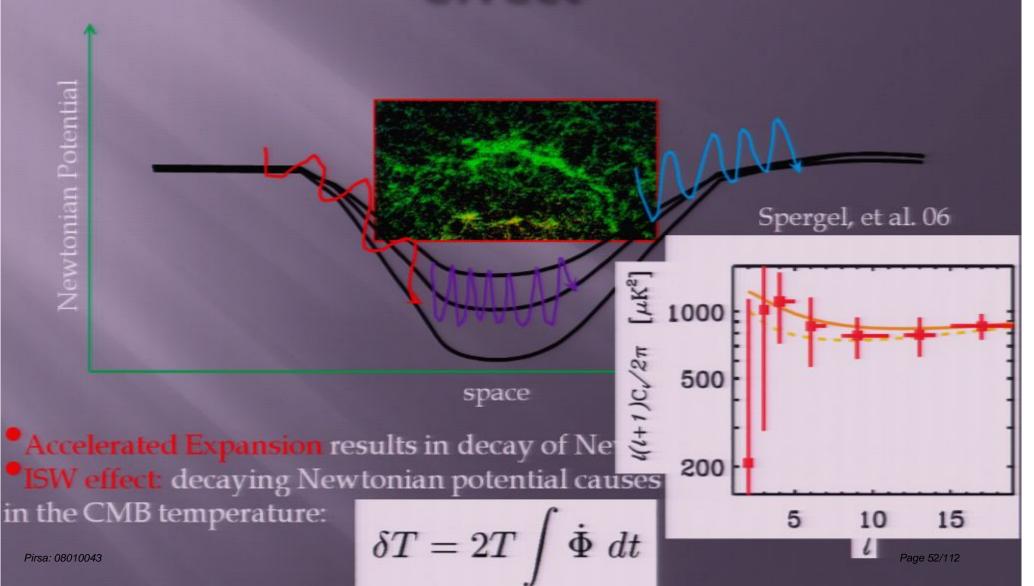
Cosmic Microwave Background



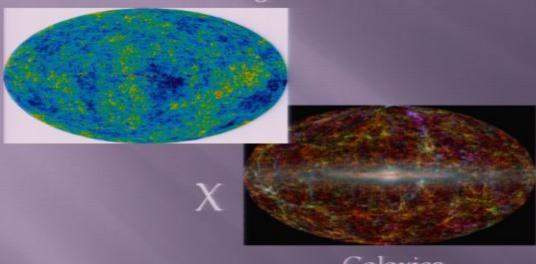
Galaxies

 Cross-correlating CMB with Galaxy distribution extracts the ISW signal from primary anisotropies

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Cosmic Microwave Background

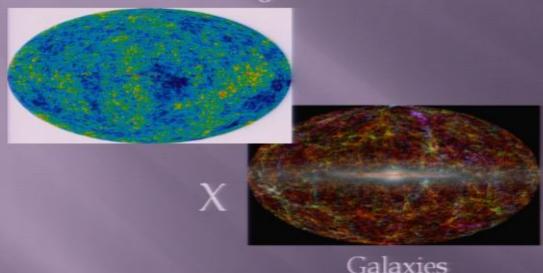


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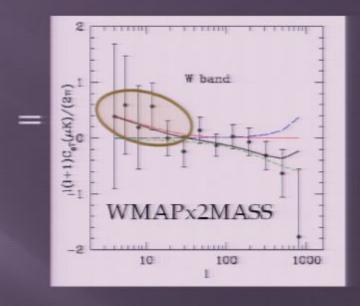
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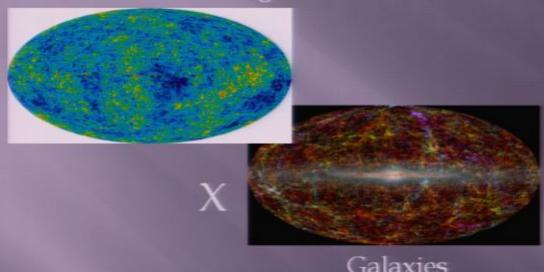
NA, Loh, & Strauss 2004



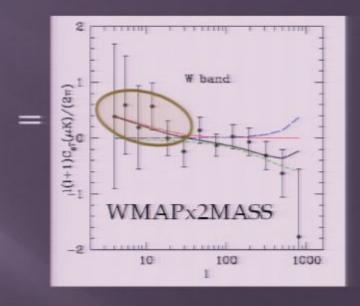
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Cosmic Microwave Background



NA, Loh, & Strauss 2004



- Cross-correlating CMB with Galaxy distribution extracts the ISW signal from primary anisotropies
- Independent Evidence for Dark Energy

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Cosmic Convergence

- Science's No. 1 breakthrough of the year 2003
 - → Consistency of different observations as evidence of Dark Energy



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 Not the best probe of conventional DE properties (w,w',...)

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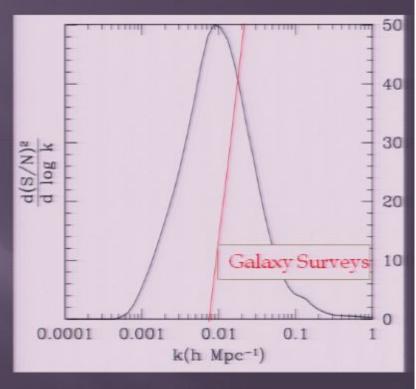


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- Not the best probe of conventional DE properties (w,w',...)
- Good probe of new physics close to the horizon scale

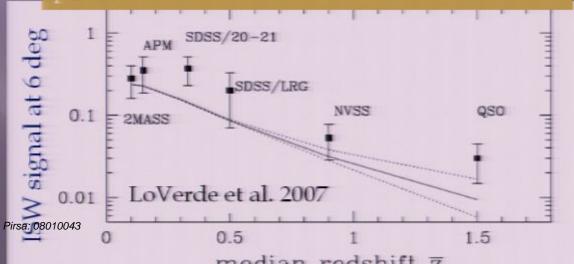


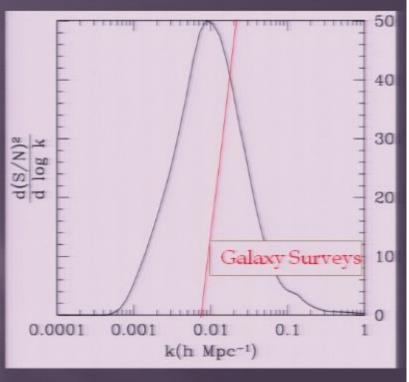
NA 2004

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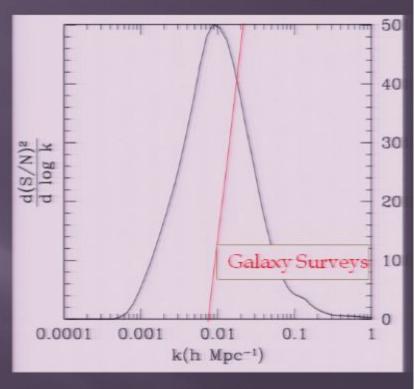
Why are all the observations larger than predicted?!





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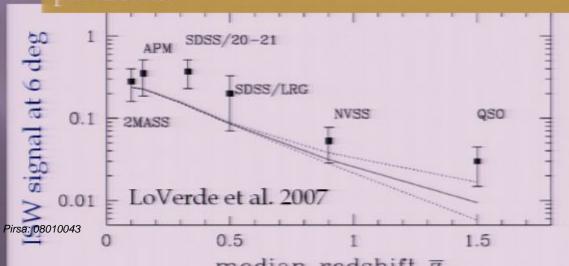


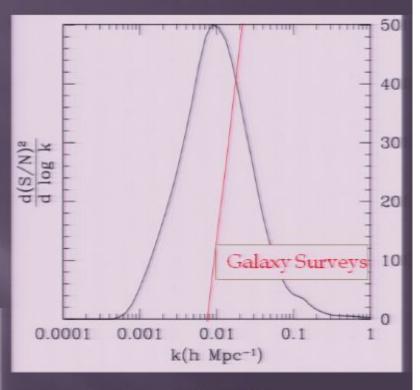
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Friedmann Equation

$$\frac{\ddot{a}}{a} = -\frac{4\pi G_F}{3}\rho$$

Newtonian Gravity

$$F = \frac{G_N m_1 m_2}{r^2}$$

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- Can Planck mass run on the cosmological horizon scale?
 - Scalar field with exponential potential (Ferreira & Joyce 1998)
 - □ Lorentz-violating vector field (Carroll & Lim 2004)
 - Cuscuton (Incompressible DE; Geshnizjani, Chung, & NA 2007)

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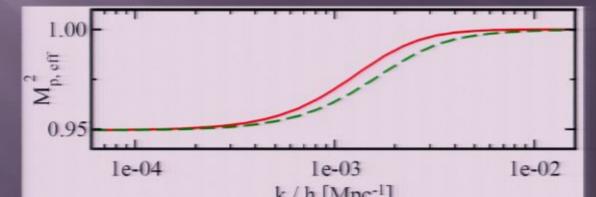
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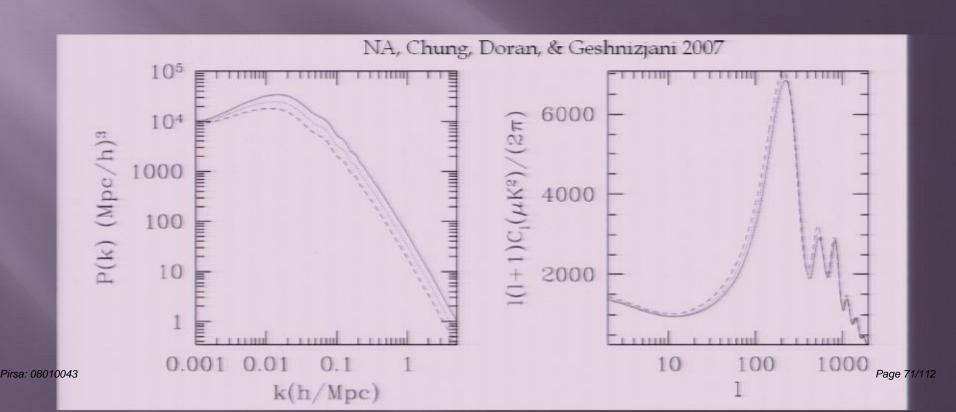
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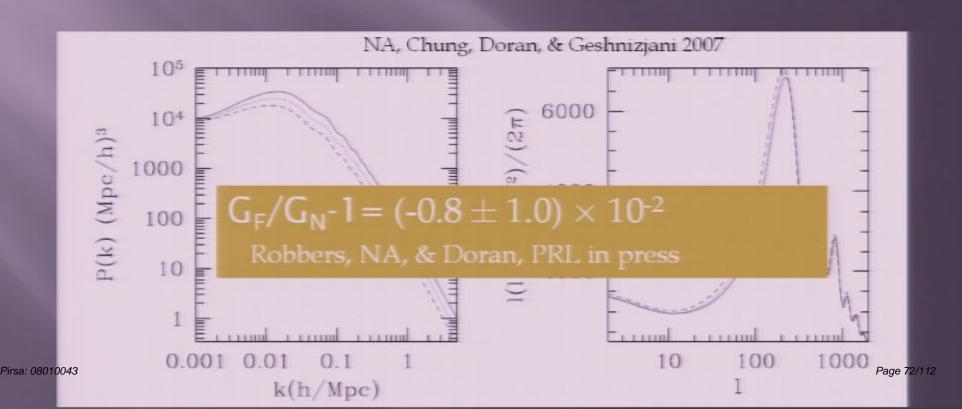
Does **Planck mass** run on the Horizon scale?

- If $G_F > G_N \rightarrow \Phi$ decays during matter era \rightarrow
 - Early ISW → boosts CMB spectrum
 - Structure formation suppressed (e.g. Lyman-α forest)



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Statistics on Horizon scale

Inflation typically generates a non-gaussian modulation of the metric perturbations:

$$\Phi = \Phi_G + f_{NL} \Phi_G^2$$

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- This introduces mode-mode coupling
- \rightarrow A large scale mode can modulate the statistics on small scales: $\Phi^s = \Phi^s_G (1 + 2f_{NL}\Phi_G^L)$

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ISW and non-Gaussianity

- On large scales, galaxy distribution follows the Newtonian potential (rather than density)
 Dalal et al. 2007
- This yields much bigger power on large scales

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- → Statistics on small scales (e.g. galaxies) could probe modes on the horizon scale

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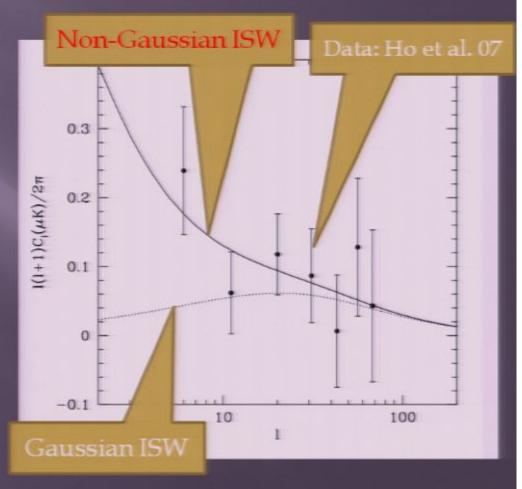
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Cross-power spectrum NVSSxCMB NA & Tolley, in prep. Page 78/112

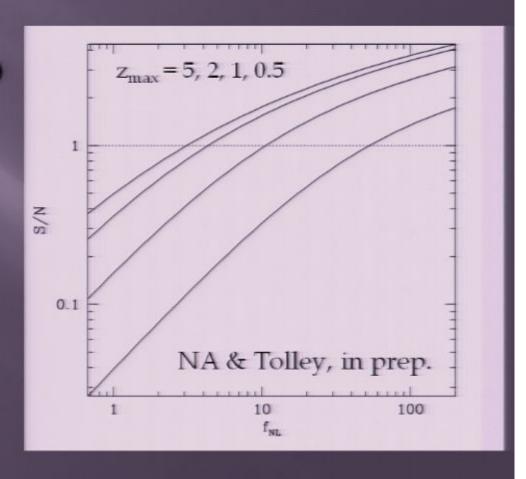
Primordial non-Gaussianity, on the horizon

- □ $f_{NL} = 87 \pm 30$ (Yadav & Wandelt 07; WMAP3, $l_{max} = 750$)
- $\triangle f_{NL} \sim 5$ (Planck satellite, in 3-5 years)

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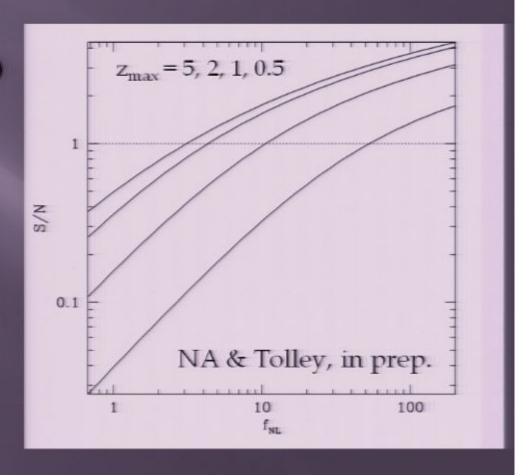
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- →insensitive to systematics



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- $\triangle f_{NL} \sim 3$ (Best accuracy for non-Gaussian ISW detection)
- →insensitive to systematics
- Similar accuracy for upcoming large scale surveys (in lieu of systematics)??



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- Inflationary paradigm predicts a rich structure beyond the horizon
- ISW effect
 - probes new physics close to horizon
 - severely constrains the running of Planck mass (<1%)
 - in correlation with galaxy surveys, probes primordial non-gaussianity (tentative evidence at 2σ level)

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- Large Scale Structure surveys: LSST, Pan Starrs, JDEM, SKA, ...
 - ISW detection, Non-gaussianity

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 - ISW detection, Non-gaussianity
- CMB surveys: WMAP, Planck
 - Better constraints on running of Planck mass

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- CMB surveys: WMAP, Planck
 - Better constraints on running of Planck mass
- SZ cluster surveys: ACT, SPT, Planck, ...
 - Large scale statistics of clusters > primordial nongaussianity

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- Our knowledge/intelect has often thrived through exploring beyond our horizons: Cosmology is no exception
- Inflationary paradigm predicts a rich structure beyond the horizon
- ISW effect
 - probes new physics close to horizon
 - severely constrains the running of Planck mass (<1%)

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- Our knowledge/intelect has often thrived through exploring beyond our horizons: Cosmology is no exception
- Inflationary paradigm predicts a rich structure beyond the horizon

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Statistics on Horizon scale

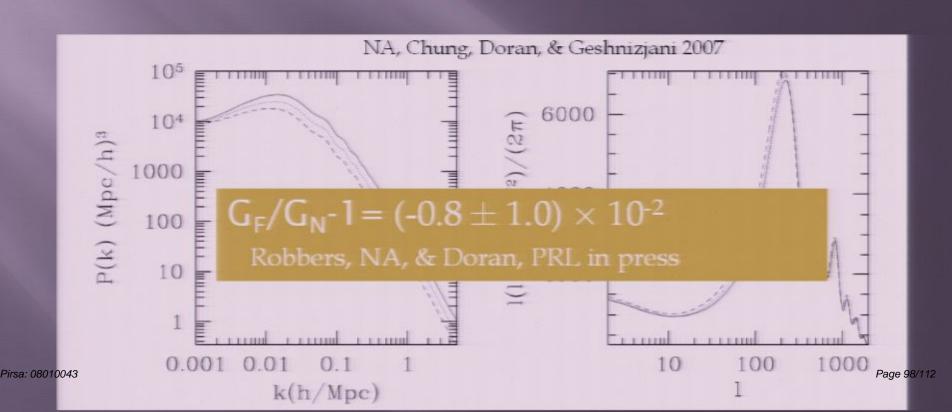
Inflation typically generates a non-gaussian modulation of the metric perturbations:

$$\Phi = \Phi_G + f_{NL} \Phi_G^2$$

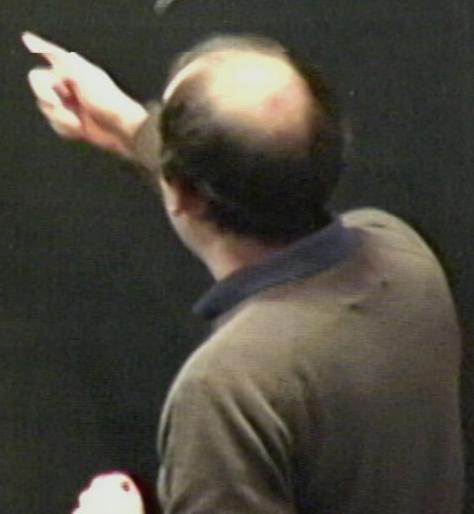
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Does **Planck mass** run on the Horizon scale?

- If $G_F > G_N \rightarrow \Phi$ decays during matter era \rightarrow
 - Early ISW → boosts CMB spectrum
 - Structure formation suppressed (e.g. Lyman-a forest)



5 + 2 H S = 4 TTGP 8



5 + 2 H S = 4 11 GP 8 H'= 4 TGF

5+2H8=4πGP 8 Sxt³/₂1ς, H'= 4πGF P

Sx t33

1 3 (GF-1)



$$|\ddot{S} + 2H\dot{S} = 4\pi G \rho \, \delta \, \delta \times t^{3/3}$$

$$H' = 4\pi G \rho \, \rho \, \delta \, \delta \times t^{3/3}$$

$$\ddot{G} + 2H\dot{G} = 3G_N H^2 \delta \, \delta \, \delta \times t^{3/3}$$

$$\ddot{G} + 2H\dot{G} = 3G_N H^2 \delta \, \delta \, \delta \times t^{3/3}$$

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$$|\ddot{\delta} + 2H\dot{\delta} = 4\pi G \rho \delta \delta t^{3}$$

$$H' = 4\pi G \rho \rho + \frac{2}{3} (\frac{2}{6} - 1)$$

$$H' = 3\frac{1}{3} (\frac{2}{6} - 1)$$

$$(\ddot{\delta} - 2H\dot{\delta} - \frac{3}{3} \frac{G_{h}H'\delta}{G_{F}})$$

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$$\frac{|\ddot{5} + 2H\dot{8} = 4\pi GP}{|\ddot{5} + 2H\dot{8} = 4\pi GP} = \frac{1}{3} \frac{1}{3}$$

$$\frac{|\ddot{5} + 2H\dot{6} = 4\pi G\rho \delta}{|\ddot{4}| \frac{G_{5}}{G_{5}} - 1}$$

$$\frac{|\ddot{5} + 2H\dot{6} = 4\pi G\rho \delta}{|\ddot{4}| \frac{G_{5}}{G_{5}} - 1}$$

$$\frac{|\ddot{5} + 2H\dot{6} = 4\pi G\rho \delta}{|\ddot{5}| \frac{1}{3} \frac{1}{3} \frac{G_{5}}{G_{5}} - 1}$$

$$\frac{|\ddot{5} + 2H\dot{6} = 4\pi G\rho \delta}{|\ddot{5}| \frac{1}{3} \frac{1}{3} \frac{G_{5}}{G_{5}} - 1}$$

$$\frac{|\ddot{5} + 2H\dot{6} = 4\pi G\rho \delta}{|\ddot{5}| \frac{1}{3} \frac{1}{3} \frac{G_{5}}{G_{5}} - 1}$$

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$$\frac{|\ddot{5} + 2H\dot{6} = 4\pi G\rho \delta}{|\ddot{5}| \frac{1}{3} \frac{1}{3} \frac{G_{5}}{G_{5}} - 1}$$

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V=- 6; V Φ= «ΠGP =- V Φ- Ην

PDE 4 TCN Page 110/112

Gravity on Horizon scale

Friedmann Equation

$$\frac{\ddot{a}}{a} = -\frac{4\pi G_F}{3}\rho$$

Newtonian Gravity

$$F = \frac{G_N m_1 m_2}{r^2}$$

- Can Planck mass run on the cosmological horizon scale?
 - Scalar field with exponential potential (Ferreira & Joyce 1998)
 - Lorentz-violating vector field (Carroll & Lim 2004)
 - Cuscuton (Incompressible DE; Geshnizjani, Chung, & NA 2007)

