Title: Foundations of Quantum Mechanics #4

Date: Jan 17, 2008 06:30 PM

URL: http://pirsa.org/08010022

Abstract: Interferometry, measurement and interpretation. Beyond the quanta.

Pirsa: 08010022 Page 1/197

Interpretation, Reformulation or Replacement?



- Is it enough to attempt to just interpret quantum theory?
- · Reformulate it?
- Replace it?

Pirsa: 08010022

ERIC POISSON

< SPM TUESDAY

ERIC POISSON PHYSICS DEPT.

< SPM TUESDAY



Interpretation, Reformulation or Replacement?



- Is it enough to attempt to just interpret quantum theory?
- · Reformulate it?
- Replace it?

Interpretation, Reformulation or Replacement?

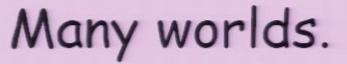


- Is it enough to attempt to just interpret quantum theory?
- · Reformulate it?
- Replace it?

Many worlds.



 Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.



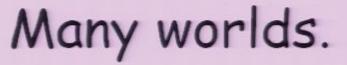


- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.



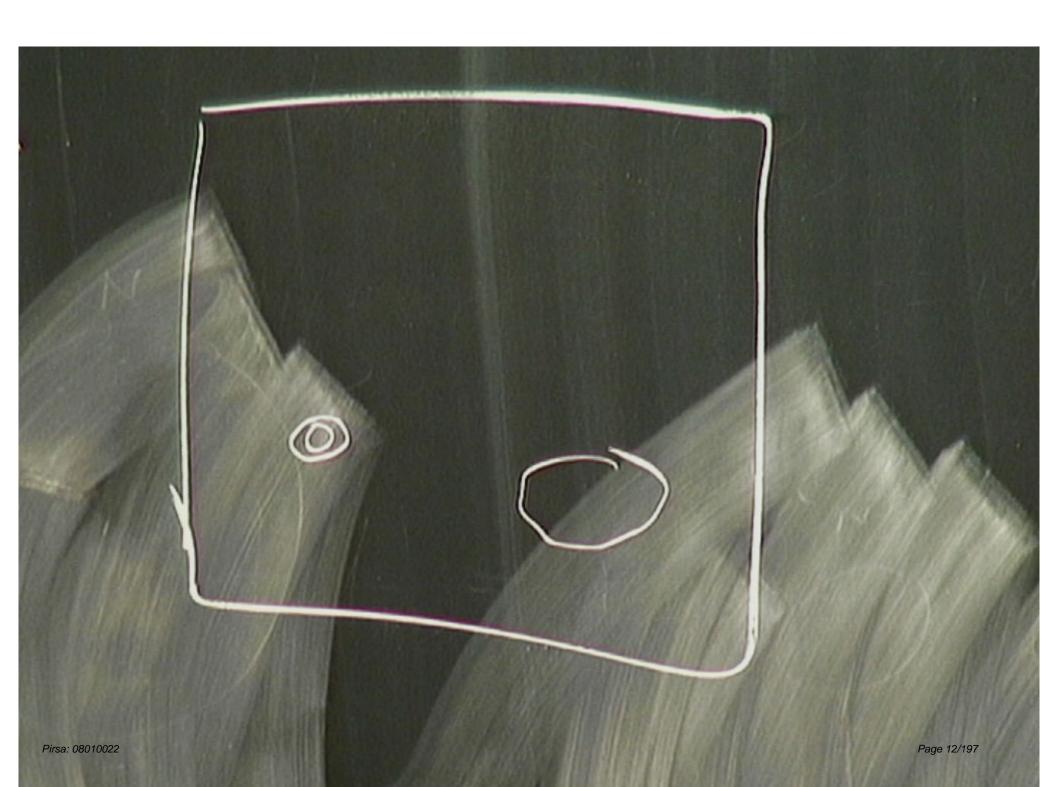


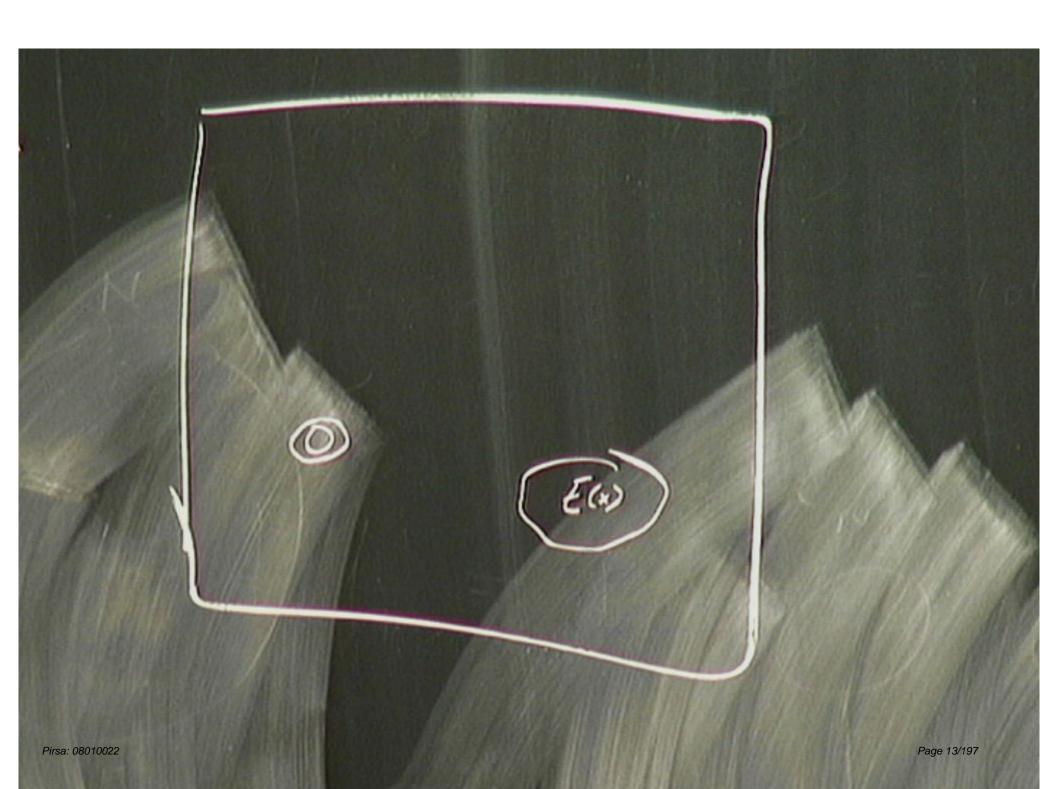
- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.
- Locality. It has been argued that branching universes do not require any nonlocality. This makes reconciliation between quantum theory and general relativity easier to achieve.

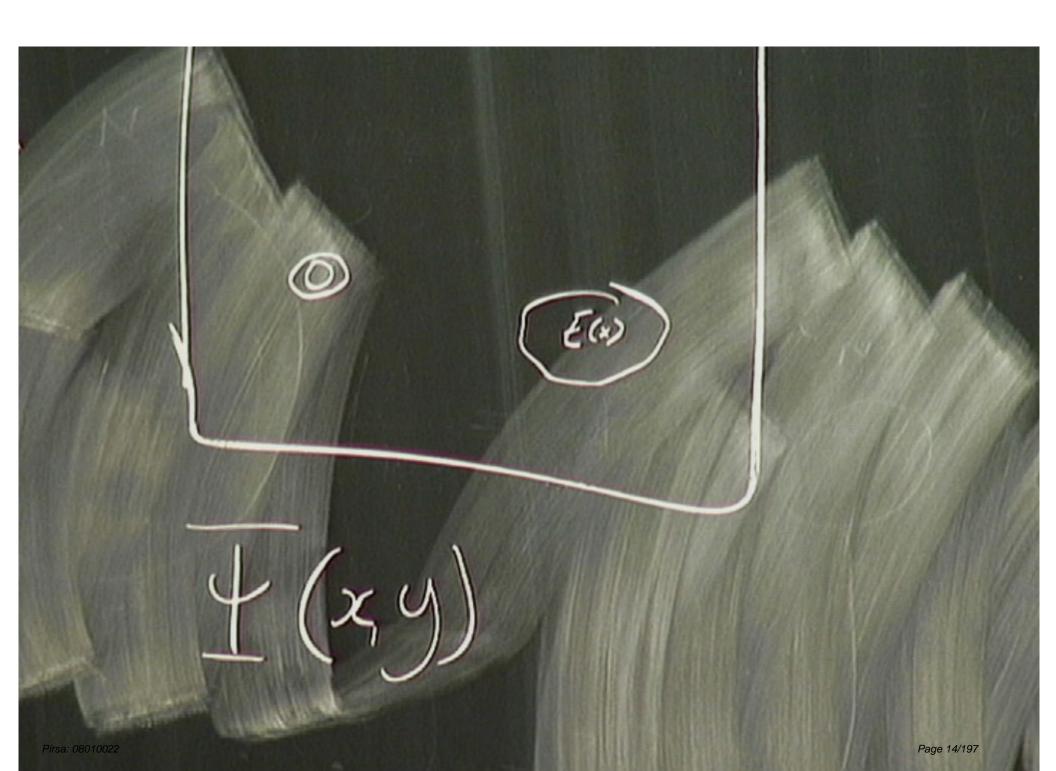


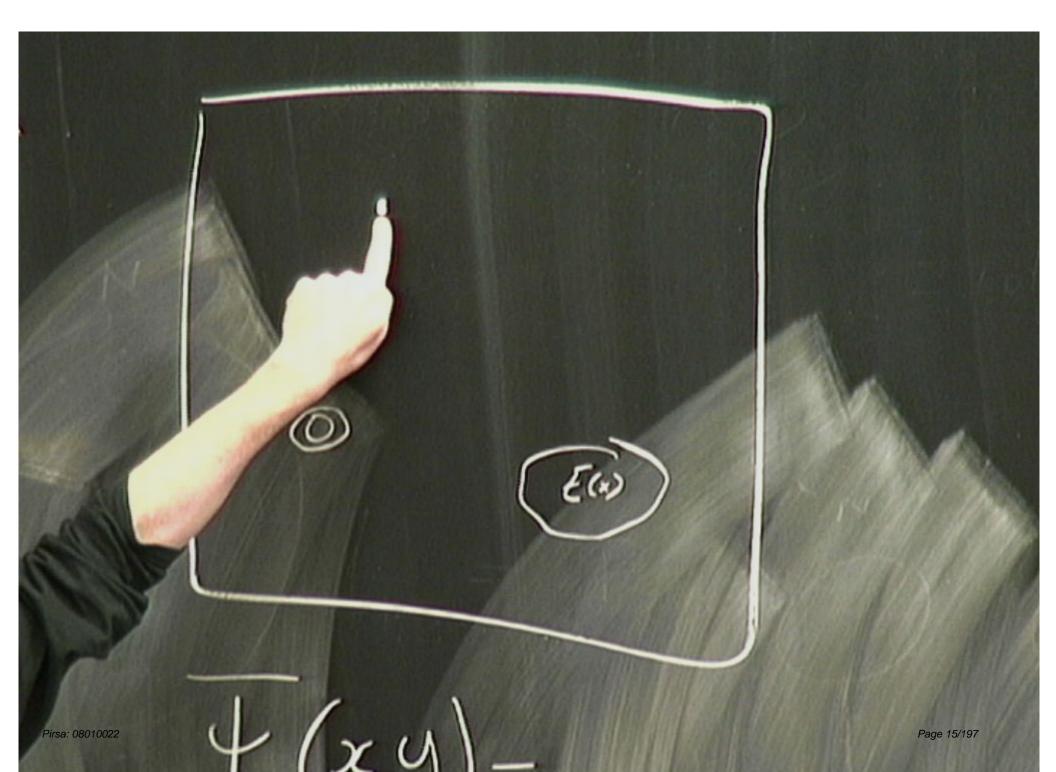


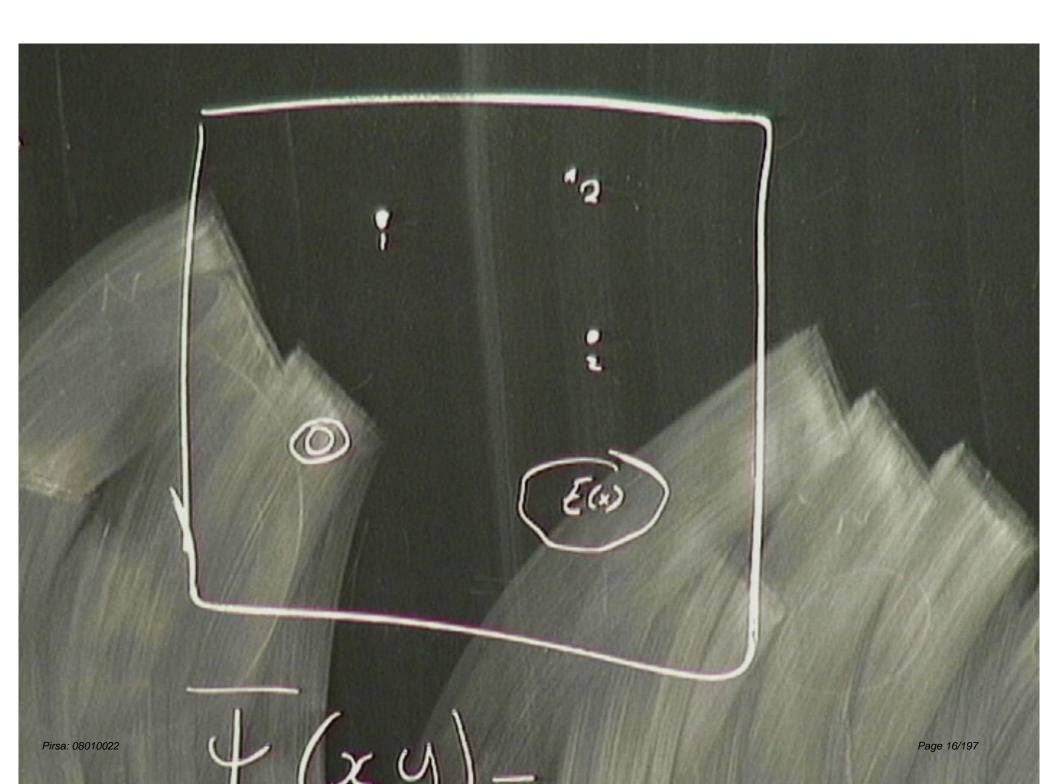
- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.
- Locality. It has been argued that branching universes do not require any nonlocality. This makes reconciliation between quantum theory and general relativity easier to achieve.
 - Is the wavefunction a local object?

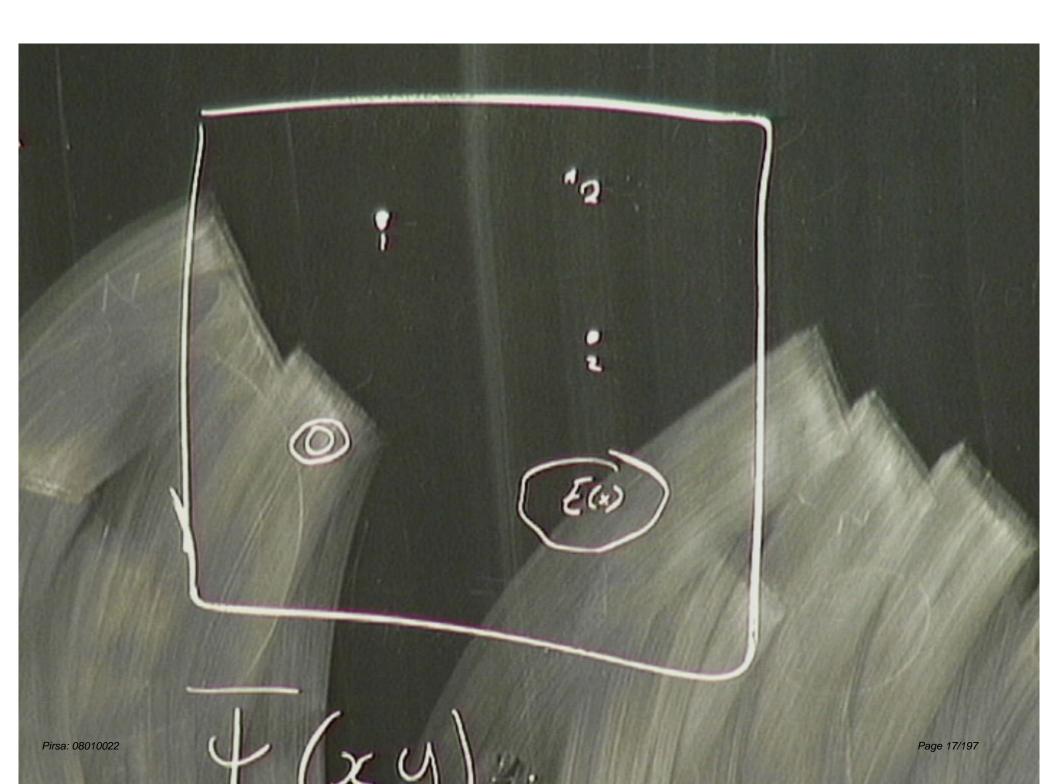
















- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.
- Locality. It has been argued that branching universes do not require any nonlocality. This makes reconciliation between quantum theory and general relativity easier to achieve.
 - Is the wavefunction a local object?

Many worlds.



- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.
- Locality. It has been argued that branching universes do not require any nonlocality. This makes reconciliation between quantum theory and general relativity easier to achieve.
 - Is the wavefunction a local object?
 - Deutsch-Hayden argument: against the Schrodinger picture



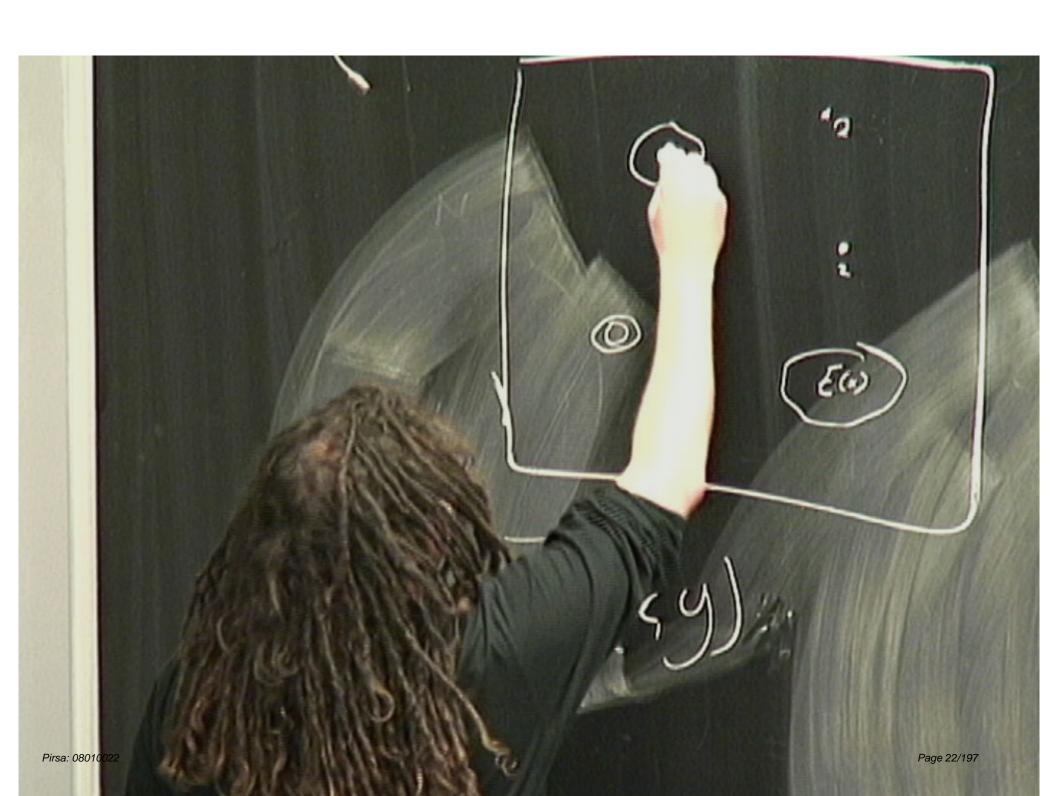


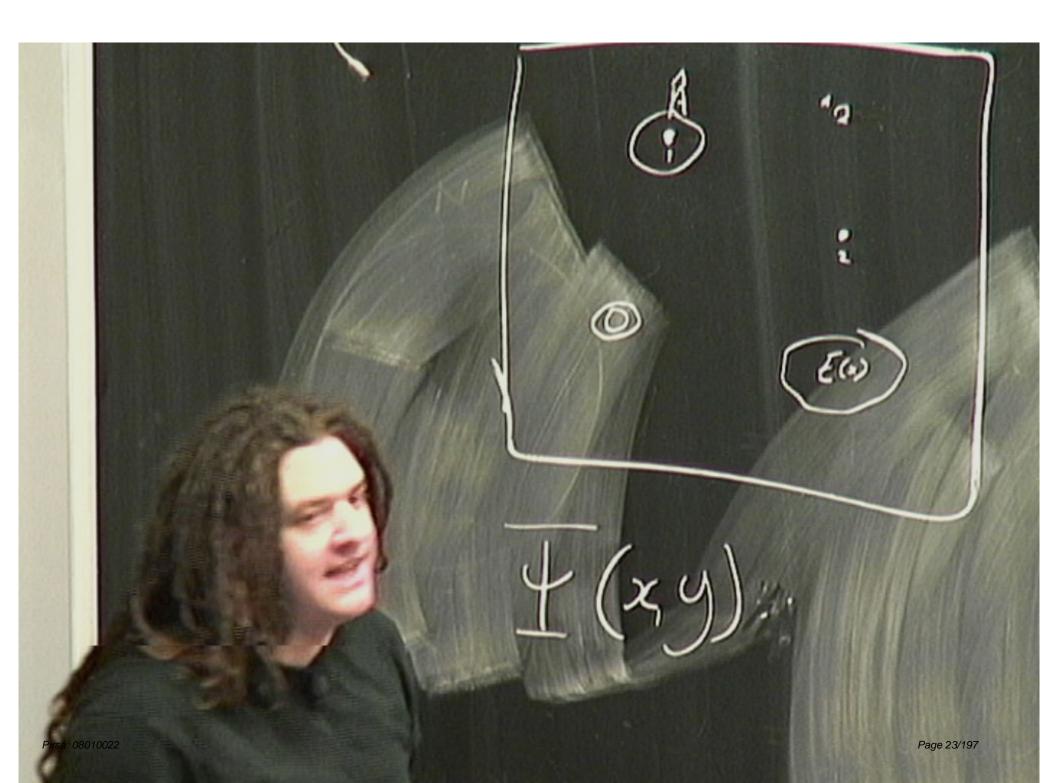
- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.
- Locality. It has been argued that branching universes do not require any nonlocality. This makes reconciliation between quantum theory and general relativity easier to achieve.
 - Is the wavefunction a local object?
 - Deutsch-Hayden argument: against the Schrodinger picture
 - · Operator Realism vs. Wavefunction Realism

Many worlds.



- Preferred basis. Hilbert spaces do not prefer any particular basis, yet for the Everettian interpretation to succeed, our perceptions must divide in a particular basis. Generally considered to be solved by decoherence.
- Probability. How to make sense of normal probabilistic assertions in a universe in which all possible outcomes do actually occur? Recent work has suggested a resolution, but it is still controversial.
- Locality. It has been argued that branching universes do not require any nonlocality. This makes reconciliation between quantum theory and general relativity easier to achieve.
 - Is the wavefunction a local object?
 - Deutsch-Hayden argument: against the Schrodinger picture
 - Operator Realism vs. Wavefunction Realism







- Original quantum theory! De Broglie, 1924-1927.
 - Wave or particle?
 - Wave and particle!

$$\psi(x,t) = |\psi(x,t)| e^{iS(x,t)}$$

For a plane wave Ae^{ikx} so $p=\hbar k=\hbar \nabla S(x,t)$

$$p = \hbar k = \hbar \nabla S(x, t)$$

$$m\dot{x} = \hbar \nabla S(x,t)$$
 $P(x|t=t_0) = |\psi(x,t_0)|^2$

Conservation equation $\frac{\partial P(x,t)}{\partial t} + \nabla J = 0$

$$J(x,t) = \frac{\psi^*(x,t)\hbar\nabla\psi(x,t) - \psi(x,t)\hbar\nabla\psi^*(x,t)}{2\,i\,m} = \frac{P(x,t)\hbar\nabla S(x,t)}{m} = P(x,t)\dot{x}$$



$$\psi(x,t) = |\psi(x,t)| e^{iS(x,t)}$$

$$m\dot{x} = \hbar \nabla S(x,t)$$

$$m\dot{x} = \hbar \nabla S(x,t)$$
 $P(x|t=t_0) = |\psi(x,t_0)|^2$

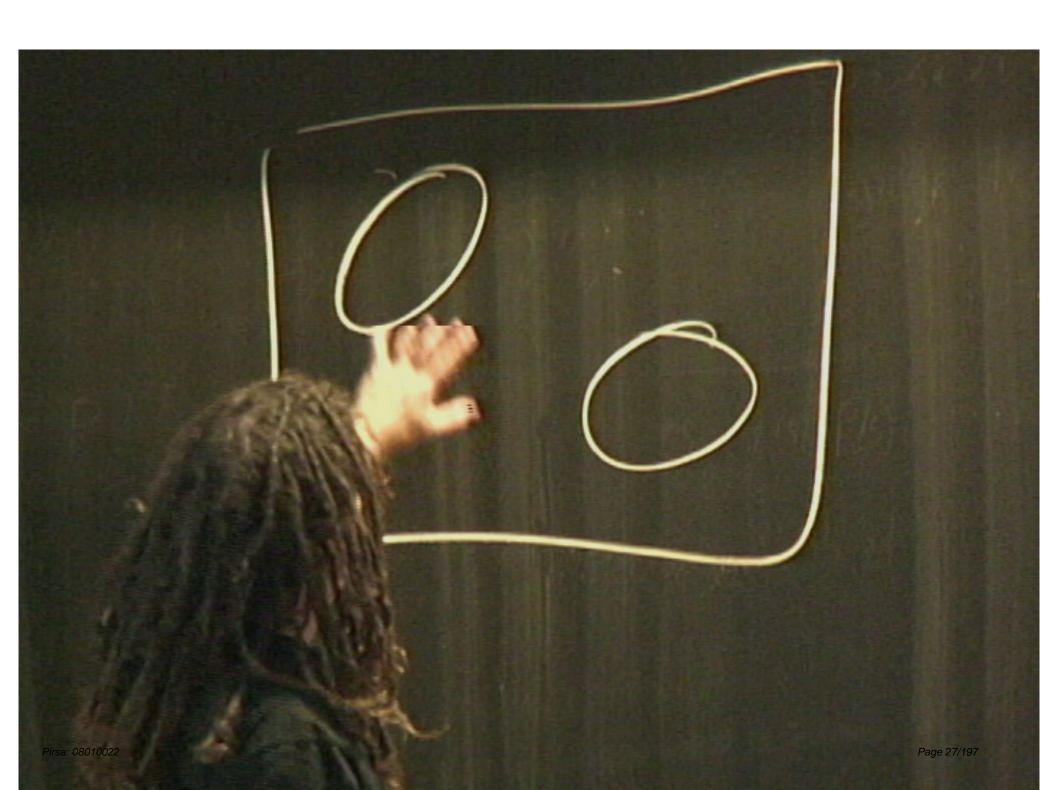
$$\psi(x,t) = \frac{1}{\sqrt{2}} \left(|\psi_u(x,t)| e^{iS_u(x,t)} + |\psi_d(x,t)| e^{iS_d(x,t)} \right) = |\psi(x,t)| e^{iS(x,t)}$$

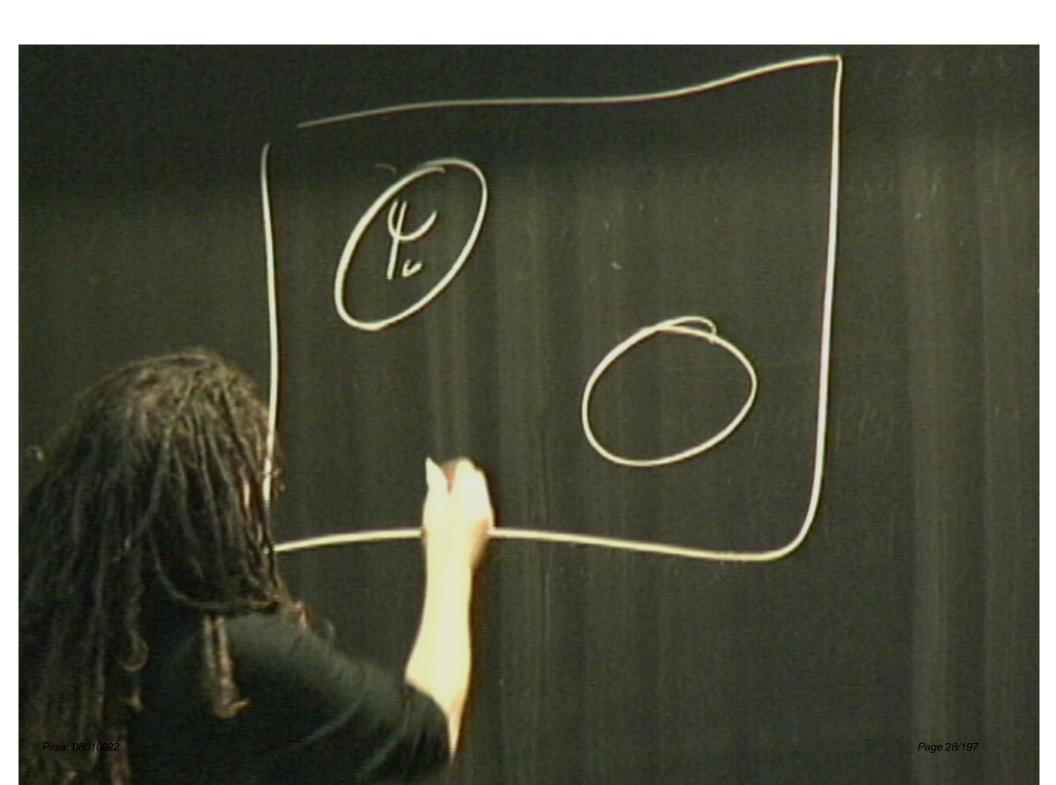
$$|\psi(x,t)| = \sqrt{\frac{|\psi_u(x,t)|^2 + |\psi_d(x,t)|^2 + 2|\psi_u(x,t)||\psi_d(x,t)|\cos(S_u(x,t) - S_d(x,t))}{S(x,t) = \arctan\left(\frac{|\psi_u(x,t)|\sin S_u(x,t) + |\psi_d(x,t)|\sin S_d(x,t)}{|\psi_u(x,t)|\cos S_u(x,t) + |\psi_d(x,t)|\cos S_d(x,t)}\right)}$$

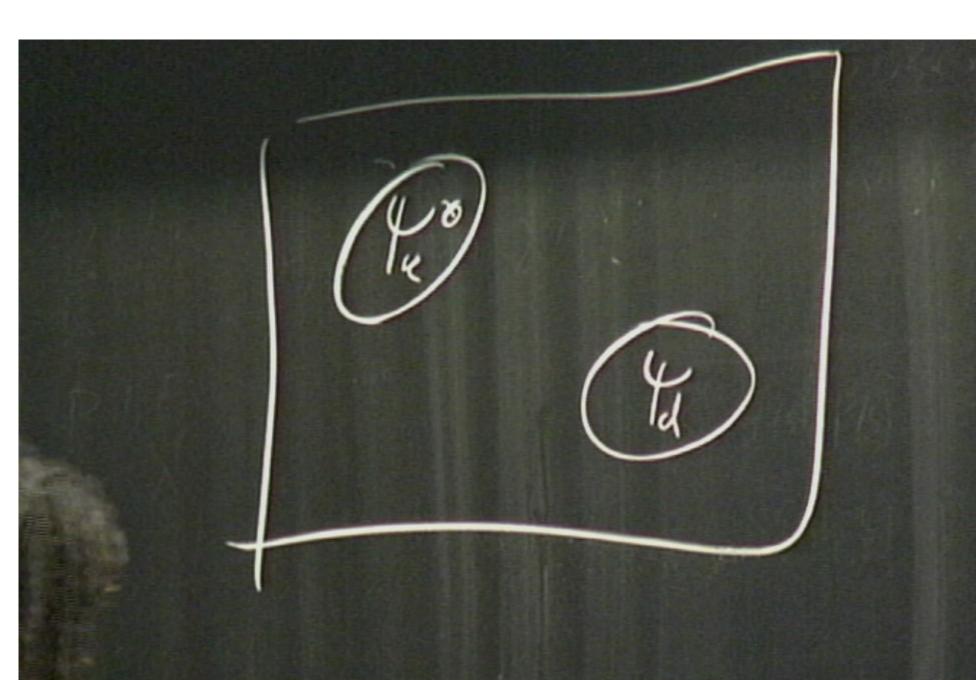
If, at some position $x' = |\psi_n(x',t)| \approx 0$ or $|\psi_n(x',t)| \approx 0$

then
$$\frac{|\psi\left(x^{\,\prime},\,t\right)|\approx|\psi_{u}\left(x^{\,\prime},\,t\right)|}{S\left(x^{\,\prime},\,t\right)\approx S_{u}\left(x^{\,\prime},\,t\right)} \quad \text{or} \quad \frac{|\psi\left(x^{\,\prime},\,t\right)|\approx|\psi_{d}\left(x^{\,\prime},\,t\right)|}{S\left(x^{\,\prime},\,t\right)\approx S_{d}\left(x^{\,\prime},\,t\right)}$$

Page 26/1972









$$\psi(x,t) = |\psi(x,t)| e^{iS(x,t)}$$

$$m\dot{x} = \hbar \nabla S(x,t)$$

$$m\dot{x}=\hbar\nabla S(x,t)$$
 $P(x|t=t_0)=|\psi(x,t_0)|^2$

$$\psi(x,t) = \frac{1}{\sqrt{2}} \left(|\psi_u(x,t)| e^{iS_u(x,t)} + |\psi_d(x,t)| e^{iS_d(x,t)} \right) = |\psi(x,t)| e^{iS(x,t)}$$

$$|\psi(x,t)| = \sqrt{\frac{|\psi_u(x,t)|^2 + |\psi_d(x,t)|^2 + 2|\psi_u(x,t)||\psi_d(x,t)|\cos(S_u(x,t) - S_d(x,t))}{S(x,t) = \arctan\left(\frac{|\psi_u(x,t)|\sin S_u(x,t) + |\psi_d(x,t)|\sin S_d(x,t)}{|\psi_u(x,t)|\cos S_u(x,t) + |\psi_d(x,t)|\cos S_d(x,t)}\right)}$$

If, at some position $x' = |\psi_n(x',t)| \approx 0$ or $|\psi_n(x',t)| \approx 0$

then
$$\frac{|\psi\left(x^{\,\prime},\,t\right)|\approx|\psi_{u}\left(x^{\,\prime},\,t\right)|}{S\left(x^{\,\prime},\,t\right)\approx S_{u}\left(x^{\,\prime},\,t\right)} \quad \text{or} \quad \frac{|\psi\left(x^{\,\prime},\,t\right)|\approx|\psi_{d}\left(x^{\,\prime},\,t\right)|}{S\left(x^{\,\prime},\,t\right)\approx S_{d}\left(x^{\,\prime},\,t\right)}$$



$$\psi(x,t) = |\psi(x,t)| e^{iS(x,t)}$$

$$m\dot{x} = \hbar \nabla S(x,t)$$

$$m\dot{x}=\hbar\nabla S(x,t)$$
 $P(x|t=t_0)=|\psi(x,t_0)|^2$

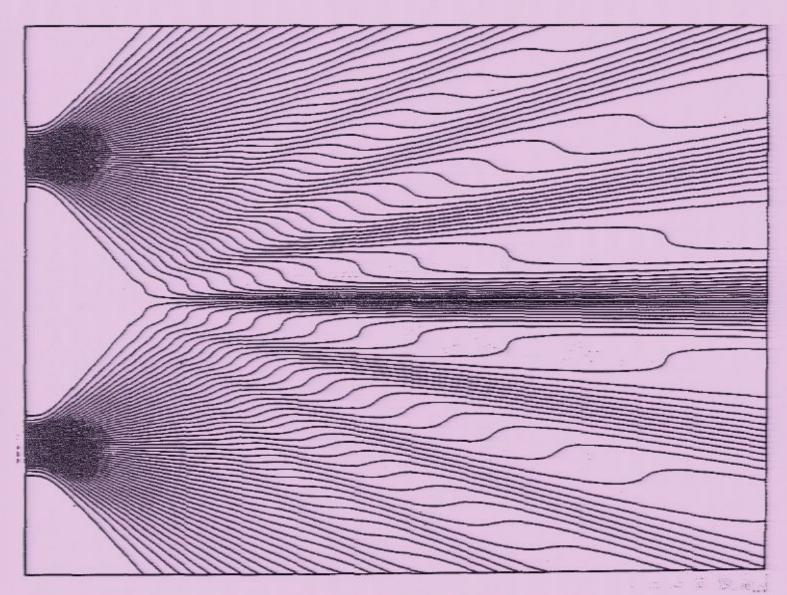
$$\psi(x,t) = \frac{1}{\sqrt{2}} \left(|\psi_u(x,t)| e^{iS_u(x,t)} + |\psi_d(x,t)| e^{iS_d(x,t)} \right) = |\psi(x,t)| e^{iS(x,t)}$$

$$|\psi(x,t)| = \sqrt{\frac{|\psi_u(x,t)|^2 + |\psi_d(x,t)|^2 + 2|\psi_u(x,t)||\psi_d(x,t)|\cos(S_u(x,t) - S_d(x,t))}{2}}{S(x,t) = \arctan\left(\frac{|\psi_u(x,t)|\sin S_u(x,t) + |\psi_d(x,t)|\sin S_d(x,t)}{|\psi_u(x,t)|\cos S_u(x,t) + |\psi_d(x,t)|\cos S_d(x,t)}\right)}$$

If, at some position $x' = |\psi_n(x',t)| \approx 0$ or $|\psi_n(x',t)| \approx 0$

then
$$\frac{|\psi(x',t)|\approx|\psi_u(x',t)|}{S(x',t)\approx S_u(x',t)} \quad \text{or} \quad \frac{|\psi(x',t)|\approx|\psi_d(x',t)|}{S(x',t)\approx S_d(x',t)}$$





Arrisan 1900 00022 ction to Quantum Foundations Source: Bohm, Hiley
Lecture 4: Interpretation, Reformulation or Replaced Addivided Universe pg. 33

New Horpage 32/1971 Fundamental Physics



$$\Psi(x,y,t) = |\Psi(x,y,t)| e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$



$$\Psi(x,y,t) = |\Psi(x,y,t)| e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \nabla \cdot \underline{J} = 0$$



$$\Psi(x,y,t)=|\Psi(x,y,t)|e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \underline{\nabla} \cdot \underline{J} = 0$$

produces two coupled equations:



$$\Psi(x,y,t)=|\Psi(x,y,t)|e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \nabla \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t)$$



$$\Psi(x,y,t)=|\Psi(x,y,t)|e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \nabla \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x, y, t)$$



With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

 $\Psi(x,y,t)=|\Psi(x,y,t)|e^{iS(x,y,t)}$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \underline{\nabla} \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x, y, t)$$



$$\Psi(x,y,t) = |\Psi(x,y,t)| e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \nabla \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x, y, t)$$

But if
$$\Psi(x, y, t) = \psi(x, t)\Phi(y, t)$$



 $\Psi(x,y,t) = |\Psi(x,y,t)| e^{iS(x,y,t)}$ With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \nabla \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x, y, t)$$

But if
$$\Psi(x,y,t)=\psi(x,t)\Phi(y,t)$$
 then $S(x,y,t)=S_{\psi}(x,t)+S_{\Phi}(y,t)$



With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

 $\Psi(x, y, t) = |\Psi(x, y, t)| e^{iS(x, y, t)}$

the probability current:

$$\frac{\partial P(x,y,t)}{\partial t} + \nabla \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x, y, t)$$

But if
$$\Psi(x,y,t) = \psi(x,t)\Phi(y,t)$$
 then $S(x,y,t) = S_{\psi}(x,t) + S_{\Phi}(y,t)$
$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S_{\psi}(x,t)$$



$$\Psi(x,y,t)=|\Psi(x,y,t)|e^{iS(x,y,t)}$$

With two degrees of freedom:

$$P(x, y|t=t_0)=|\Psi(x, y, t_0)|^2$$

the probability current:
$$\frac{\partial P(x,y,t)}{\partial t} + \underline{\nabla} \cdot \underline{J} = 0$$

produces two coupled equations:

$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x, y, t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x, y, t)$$

But if
$$\Psi(x,y,t) = \psi(x,t)\Phi(y,t)$$
 then $S(x,y,t) = S_{\psi}(x,t) + S_{\Phi}(y,t)$
$$\frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S_{\psi}(x,t) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S_{\Phi}(y,t)$$

$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$





$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$



$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$

$$|\Psi(x,y)| = \sqrt{\frac{|\Psi_{u}(x,y)|^{2} + |\Psi_{d}(x,y)|^{2} + 2|\Psi_{u}(x,y)||\Psi_{d}(x,y)|\cos(S_{u}(x,y) - S_{d}(x,y))}{2}}$$



$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$

$$|\Psi(x,y)| = \sqrt{\frac{|\Psi_{u}(x,y)|^{2} + |\Psi_{d}(x,y)|^{2} + 2|\Psi_{u}(x,y)||\Psi_{d}(x,y)|\cos(S_{u}(x,y) - S_{d}(x,y))}{2}}$$



$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$

$$|\Psi(x,y)| = \sqrt{\frac{|\Psi_{u}(x,y)|^{2} + |\Psi_{d}(x,y)|^{2} + 2|\Psi_{u}(x,y)||\Psi_{d}(x,y)|\cos(S_{u}(x,y) - S_{d}(x,y))}{2}}$$

$$\begin{split} & \Psi_u(x,y) = \psi_u(x) \, \Phi_u(y) \\ & \Psi_d(x,y) = \psi_d(x) \, \Phi_u(y) \end{split} \qquad S(x,y) = \arctan \left(\frac{|\Psi_u(x,y)| \sin S_u(x,y) + |\Psi_d(x,y)| \sin S_d(x,y)}{|\Psi_u(x,y)| \cos S_u(x,y) + |\Psi_d(x,y)| \cos S_d(x,y)} \right) \end{split}$$



$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$

$$|\Psi(x,y)| = \sqrt{\frac{|\Psi_u(x,y)|^2 + |\Psi_d(x,y)|^2 + 2|\Psi_u(x,y)||\Psi_d(x,y)|\cos(S_u(x,y) - S_d(x,y))}{2}}$$

$$\begin{split} & \Psi_u(x,y) = \psi_u(x) \Phi_u(y) \\ & \Psi_d(x,y) = \psi_d(x) \Phi_u(y) \end{split} \qquad S(x,y) = \arctan \left(\frac{|\Psi_u(x,y)| \sin S_u(x,y) + |\Psi_d(x,y)| \sin S_d(x,y)}{|\Psi_u(x,y)| \cos S_u(x,y) + |\Psi_d(x,y)| \cos S_d(x,y)} \right) \end{split}$$

If, at some position $y' \quad |\Phi_u(y')| \approx 0 \quad \text{or} \quad |\Phi_d(y')| \approx 0$



$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$

$$|\Psi(x,y)| = \sqrt{\frac{|\Psi_u(x,y)|^2 + |\Psi_d(x,y)|^2 + 2|\Psi_u(x,y)||\Psi_d(x,y)|\cos(S_u(x,y) - S_d(x,y))}{2}}$$

$$\begin{split} & \Psi_u(x,y) = \psi_u(x) \, \Phi_u(y) \\ & \Psi_d(x,y) = \psi_d(x) \, \Phi_u(y) \end{split} \qquad S(x,y) = \arctan \left(\frac{|\Psi_u(x,y)| \sin S_u(x,y) + |\Psi_d(x,y)| \sin S_d(x,y)}{|\Psi_u(x,y)| \cos S_u(x,y) + |\Psi_d(x,y)| \cos S_d(x,y)} \right) \end{split}$$

If, at some position y' $|\Phi_u(y')| \approx 0$ or $|\Phi_d(y')| \approx 0$

then
$$\frac{\Psi(x,y')\!\approx\!\psi_u(x)\Phi_u(y')}{\frac{dx}{dt}\!\approx\!\frac{\hbar}{m}\frac{\partial}{\partial x}S_{\psi_u}(x)}$$



$$\Psi(x,y) = \frac{1}{\sqrt{2}} \left(\psi_u(x) \Phi_u(y) + \psi_d(x) \Phi_d(y) \right)$$

$$\Psi(x,y) = |\Psi(x,y)| e^{iS(x,y)} \qquad \frac{dx}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial x} S(x,y) \qquad \frac{dy}{dt} = \frac{\hbar}{m} \frac{\partial}{\partial y} S(x,y)$$

$$|\Psi(x,y)| = \sqrt{\frac{|\Psi_{u}(x,y)|^{2} + |\Psi_{d}(x,y)|^{2} + 2|\Psi_{u}(x,y)||\Psi_{d}(x,y)|\cos(S_{u}(x,y) - S_{d}(x,y))}{2}}$$

$$\begin{aligned} & \Psi_u(x,y) = \psi_u(x) \, \Phi_u(y) \\ & \Psi_d(x,y) = \psi_d(x) \, \Phi_u(y) \end{aligned} \qquad S(x,y) = \arctan \left(\frac{|\Psi_u(x,y)| \sin S_u(x,y) + |\Psi_d(x,y)| \sin S_d(x,y)}{|\Psi_u(x,y)| \cos S_u(x,y) + |\Psi_d(x,y)| \cos S_d(x,y)} \right) \end{aligned}$$

If, at some position y' $|\Phi_u(y')| \approx 0$ or $|\Phi_d(y')| \approx 0$

then
$$\frac{\Psi(x,y') \approx \psi_u(x) \Phi_u(y')}{\frac{dx}{dt} \approx \frac{\hbar}{m} \frac{\partial}{\partial x} S_{\psi_u}(x)} \qquad \text{or} \qquad \frac{\Psi(x,y') \approx \psi_d(x) \Phi_d(y')}{\frac{dx}{dt} \approx \frac{\hbar}{m} \frac{\partial}{\partial x} S_{\psi_d}(x)}$$

ArPiism 49000022ction to Quantum Foundations
Lecture 4: Interpretation, Reformulation or Replacement?

New HorPage 506197n Fundamental Physics



Quantum field theory. Although Bohm presented a field ontology for the electromagnetic
field in 1952, most work has been on non-relativistic particle theories. Recent work has
shown how de Broglie-Bohm hidden variables can be constructed for general interacting
field theories.



- Quantum field theory. Although Bohm presented a field ontology for the electromagnetic field in 1952, most work has been on non-relativistic particle theories. Recent work has shown how de Broglie-Bohm hidden variables can be constructed for general interacting field theories.
- Different choices of hidden variable. Particle and field configuration hidden variables
 present intuitively clear routes to distinct outcomes. Alternative hidden variables spin,
 orientation, momentum, matrix valued, grassman number valued may or may not be
 feasible.



- Quantum field theory. Although Bohm presented a field ontology for the electromagnetic field in 1952, most work has been on non-relativistic particle theories. Recent work has shown how de Broglie-Bohm hidden variables can be constructed for general interacting field theories.
- Different choices of hidden variable. Particle and field configuration hidden variables
 present intuitively clear routes to distinct outcomes. Alternative hidden variables spin,
 orientation, momentum, matrix valued, grassman number valued may or may not be
 feasible.
- Empty waves. In de Broglie-Bohm theories, the portions of the wavefunction that
 correspond to the outcomes that did not occur, still exist. It can be argued that these
 outcomes are just as real and that hidden variable theorists are Everettians "in a chronic
 state of denial".



- Quantum field theory. Although Bohm presented a field ontology for the electromagnetic field in 1952, most work has been on non-relativistic particle theories. Recent work has shown how de Broglie-Bohm hidden variables can be constructed for general interacting field theories.
- Different choices of hidden variable. Particle and field configuration hidden variables
 present intuitively clear routes to distinct outcomes. Alternative hidden variables spin,
 orientation, momentum, matrix valued, grassman number valued may or may not be
 feasible.
- Empty waves. In de Broglie-Bohm theories, the portions of the wavefunction that
 correspond to the outcomes that did not occur, still exist. It can be argued that these
 outcomes are just as real and that hidden variable theorists are Everettians "in a chronic
 state of denial".
- Non-equilibrium. Hidden variable theories reproduce quantum mechanics for particular probability distributions over the hidden variable state. This distribution is often referred to as "quantum equilibrium", as it's justifications is similar to thermal equilibrium. The possibility of systems with non-equilibrium distributions would lead to novel experimental results and possibilities.

ArPisen 490000022ction to Quantum Foundations
Lecture 4: Interpretation, Reformulation or Replacement?

New Horfage 544197n Fundamental Physics



$$\Psi = \frac{1}{\sqrt{2}} \left(\psi_u \Phi_u + \psi_d \Phi_d \right)$$



$$\Psi = \frac{1}{\sqrt{2}} \left(\psi_u \Phi_u + \psi_d \Phi_d \right)$$

 The wavefunction is real and does represent the state of a physical object.



$$\Psi = \frac{1}{\sqrt{2}} \left(\psi_u \Phi_u + \psi_d \Phi_d \right)$$

- The wavefunction is real and does represent the state of a physical object.
- Linear evolution is not right



$$\Psi = \frac{1}{\sqrt{2}} \left(\psi_u \Phi_u + \psi_d \Phi_d \right)$$

- The wavefunction is real and does represent the state of a physical object.
- Linear evolution is not right
- The state of the world is actually



$$\Psi = \frac{1}{\sqrt{2}} \left(\psi_u \Phi_u + \psi_d \Phi_d \right)$$

- The wavefunction is real and does represent the state of a physical object.
- Linear evolution is not right
- The state of the world is actually

$$\psi_u \Phi_u$$
 or $\psi_d \Phi_d$



Linear evolution changes at some point. When?



- Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?





- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.



- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.
- Ghirardi, Rimini, Weber (1986)
 - For each degree of freedom, there is a random chance of a spontaneous collapse taking place, with a probability: $1/\tau$ per unit time.



- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.
- · Ghirardi, Rimini, Weber (1986)
 - For each degree of freedom, there is a random chance of a spontaneous collapse taking place, with a probability: $1/\tau$ per unit time.
 - When a collapse takes place the wavefunction changes from: $\Psi_o(x,y,z...)$



- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.
- · Ghirardi, Rimini, Weber (1986)
 - For each degree of freedom, there is a random chance of a spontaneous collapse taking place, with a probability: $1/\tau$ per unit time.
 - When a collapse takes place the wavefunction changes from: $\,\Psi_{o}(x\,,y\,,z...)\,$

to
$$\Psi_{N}(x, y, z...) = \frac{f(x'-x)}{N(x)} \Psi_{o}(x', y, z...)$$



- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.
- · Ghirardi, Rimini, Weber (1986)
 - For each degree of freedom, there is a random chance of a spontaneous collapse taking place, with a probability: $1/\tau$ per unit time.
 - When a collapse takes place the wavefunction changes from: $\Psi_o(x,y,z...)$

to
$$\Psi_N(x, y, z...) = \frac{f(x'-x)}{N(x)} \Psi_o(x', y, z...)$$

Where: f(x'-x) is a function sharply peaked around (x'=x)



- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.
- · Ghirardi, Rimini, Weber (1986)
 - For each degree of freedom, there is a random chance of a spontaneous collapse taking place, with a probability: $1/\tau$ per unit time.
 - When a collapse takes place the wavefunction changes from: $\Psi_o(x,y,z...)$

to
$$\Psi_{\scriptscriptstyle N}(x\,,\,y\,,\,z...) = \frac{f\,(\,x\,'-x\,)}{N(\,x\,)} \Psi_{\scriptscriptstyle o}(x\,'\,,\,y\,,\,z...)$$

Where: f(x'-x) is a function sharply peaked around (x'=x)

Normalisation requires:
$$N(x) = \sqrt{\int_{x',y,z...} \left| f\left(x'-x\right) \Psi_o(x',y,z...) \right|^2 dx' \, dy dz..}$$



- · Linear evolution changes at some point. When?
 - Macroscopic, classical, observed, irreversible, conscious?
 - As there is a difference between the empirical predictions of the uncollapsed and collapsed wavefunctions, any unambiguous model for how it happens leads to definite empirical predictions.
- · Ghirardi, Rimini, Weber (1986)
 - For each degree of freedom, there is a random chance of a spontaneous collapse taking place, with a probability: $1/\tau$ per unit time.
 - When a collapse takes place the wavefunction changes from: $\Psi_o(x,y,z...)$

to
$$\Psi_N(x, y, z...) = \frac{f(x'-x)}{N(x)} \Psi_o(x', y, z...)$$

Where: f(x'-x) is a function sharply peaked around (x'=x)

Normalisation requires:
$$N(x) = \sqrt{\int_{x',y,z...} \left| f\left(x'-x\right) \Psi_o(x',y,z...) \right|^2 dx' \, dy dz..}$$

and x' is selected at random with probability:

$$P(x') = |N(x')|^2 dx'$$

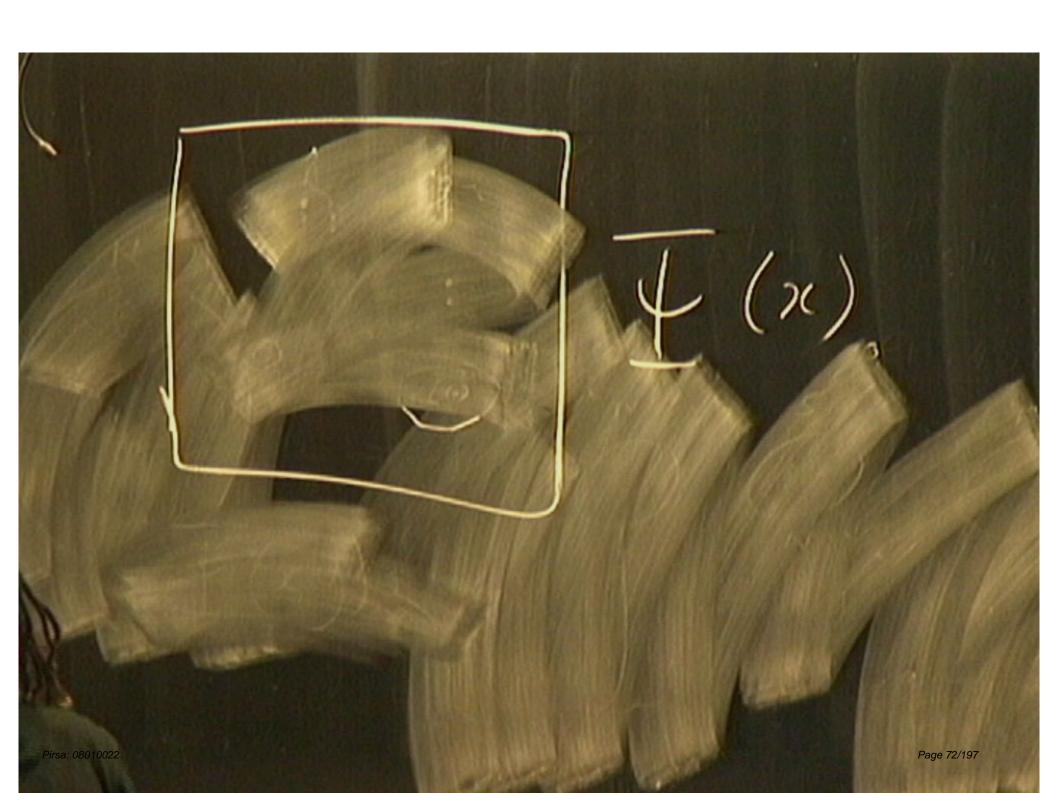
ArPisantention to Quantum Foundations
Lecture 4: Interpretation, Reformulation or Replacement?

New Horfage 695197n Fundamental Physics

$$\Psi_o(x) \rightarrow \Psi_N(x) = \frac{f(x'-x)}{N(x)} \Psi_o(x')$$



Pirsa: 08010022 Page 71/197



Page 73/197

$$\Psi_o(x) \rightarrow \Psi_N(x) = \frac{f(x'-x)}{N(x)} \Psi_o(x')$$





$$\boldsymbol{\varPsi}_{o}(\boldsymbol{x}) \!\rightarrow\! \boldsymbol{\varPsi}_{N}(\boldsymbol{x}) \!=\! \frac{f\left(\boldsymbol{x}^{\,\prime} \!-\! \boldsymbol{x}\right)}{N\left(\boldsymbol{x}\right)} \boldsymbol{\varPsi}_{o}(\boldsymbol{x}^{\,\prime})$$

 $\tau \approx 10^{15} s$ In 1 second probability of a collapse in x is $p \approx 10^{-15}$

Average time for collapse: $\tau_x \approx 10^{15} s$

$$\Psi_o(x) = \frac{\psi_u(x) + \psi_d(x)}{\sqrt{2}} \quad \text{collapses to:} \quad \Psi_N(x) = \frac{f\left(x' - x\right)\psi_u(x') + f\left(x' - x\right)\psi_d(x')}{N(x)\sqrt{2}}$$

Provided $\psi_u(x)\psi_d(x)\approx 0$ for all values of x and $f(x'-x)\approx 0$ for $x'\neq x$



$$\boldsymbol{\varPsi}_{o}(\boldsymbol{x}) \!\rightarrow\! \boldsymbol{\varPsi}_{N}(\boldsymbol{x}) \!=\! \! \frac{f\left(\boldsymbol{x}^{\,\prime} \!-\! \boldsymbol{x}\right)}{N\left(\boldsymbol{x}\right)} \boldsymbol{\varPsi}_{o}(\boldsymbol{x}^{\,\prime})$$

 $\tau \approx 10^{15} s$ In 1 second probability of a collapse in x is $p \approx 10^{-15}$

Average time for collapse: $\tau_x \approx 10^{15} s$

$$\Psi_o(x) = \frac{\psi_u(x) + \psi_d(x)}{\sqrt{2}} \quad \text{collapses to:} \quad \Psi_N(x) = \frac{f\left(x' - x\right)\psi_u(x') + f\left(x' - x\right)\psi_d(x')}{N(x)\sqrt{2}}$$

Provided $\psi_u(x)\psi_d(x) \approx 0$ for all values of x and $f(x'-x) \approx 0$ for $x' \neq x$



$$\boldsymbol{\varPsi}_{o}(\boldsymbol{x}) \!\rightarrow\! \boldsymbol{\varPsi}_{N}(\boldsymbol{x}) \!=\! \! \frac{f\left(\boldsymbol{x}^{\,\prime} \!-\! \boldsymbol{x}\right)}{N\left(\boldsymbol{x}\right)} \boldsymbol{\varPsi}_{o}(\boldsymbol{x}^{\,\prime})$$

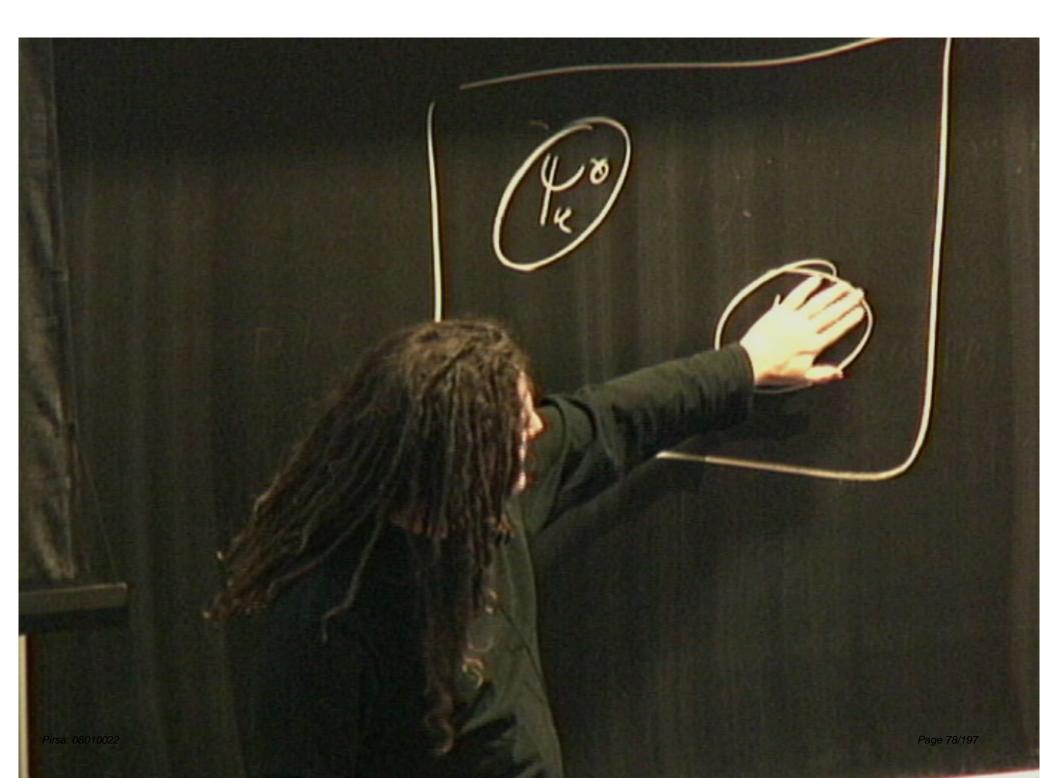
 $\tau \approx 10^{15} s$ In 1 second probability of a collapse in x is $p \approx 10^{-15}$

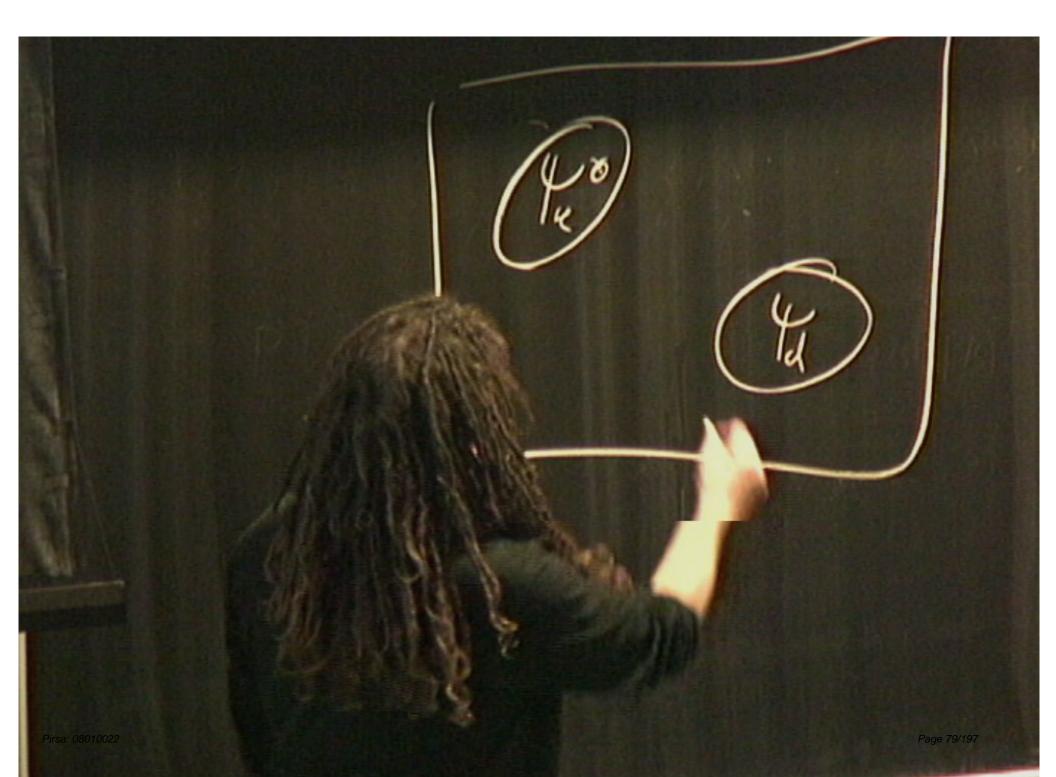
Average time for collapse: $\tau_x \approx 10^{15} s$

$$\Psi_o(x) = \frac{\psi_u(x) + \psi_d(x)}{\sqrt{2}} \quad \text{collapses to:} \quad \Psi_N(x) = \frac{f(x' - x)\psi_u(x') + f(x' - x)\psi_d(x')}{N(x)\sqrt{2}}$$

Provided $\psi_u(x)\psi_d(x) \approx 0$ for all values of x and $f(x'-x) \approx 0$ for $x' \neq x$

either:
$$f(x'-x)\psi_u(x) \approx 0$$
 or: $f(x'-x)\psi_d(x) \approx 0$







Now take:
$$\Psi_o(x,y,z...) = \frac{\psi_u(x) \varPhi_u(y,z,...) + \psi_d(x) \varPhi_d(y,z,...)}{\sqrt{2}}$$



Now take:
$$\Psi_{o}(x,y,z...) = \frac{\psi_{u}(x) \varPhi_{u}(y,z,...) + \psi_{d}(x) \varPhi_{d}(y,z,...)}{\sqrt{2}}$$



Now take:
$$\Psi_o(x,y,z...) = \frac{\psi_u(x) \varPhi_u(y,z,...) + \psi_d(x) \varPhi_d(y,z,...)}{\sqrt{2}}$$

Provided



Now take:
$$\Psi_o(x,y,z...) = \frac{\psi_u(x) \Phi_u(y,z,...) + \psi_d(x) \Phi_d(y,z,...)}{\sqrt{2}}$$

Provided

$$\Phi_u(y,z,...)\Phi_d(y,z,...)\approx 0$$
 for all values of y and $f(y'-y)\approx 0$ $y'\neq y$

$$f(y'-y)\Phi_u(y,z,...)\approx 0$$
 or: $f(y'-y)\Phi_d(y,z,...)\approx 0$

$$\text{SO:} \quad \Psi_{N}(x,y,z...) = \frac{\psi_{u}(x) \, f(\,y\,'-\,y) \Phi_{u}(\,y\,',z\,,...) + \psi_{d}(x) \, f(\,y\,'-\,y) \Phi_{d}(\,y\,',z\,,...)}{N(\,y) \sqrt{2}}$$



Now take:
$$\Psi_o(x, y, z...) = \frac{\psi_u(x)\Phi_u(y, z, ...) + \psi_d(x)\Phi_d(y, z, ...)}{\sqrt{2}}$$

Provided

$$\Phi_u(y,z,...)\Phi_d(y,z,...)\approx 0$$
 for all values of y and $f(y'-y)\approx 0$ $y'\neq y$

$$f(y'-y)\Phi_u(y,z,...)\approx 0$$
 or: $f(y'-y)\Phi_d(y,z,...)\approx 0$

$$\text{SO:} \quad \Psi_{N}(x,y,z...) = \frac{\psi_{u}(x) \, f\left(\, y\,' - y\right) \varPhi_{u}(\, y\,',z\,,...) + \psi_{d}(x) \, f\left(\, y\,' - y\right) \varPhi_{d}(\, y\,',z\,,...)}{N(\, y) \sqrt{2}}$$

$$\Psi_N(x,y,z...) \approx \frac{\psi_u(x) \, f\left(y^{\,\prime}-y\right) \Phi_u(y^{\,\prime},z\,,...)}{N(y) \sqrt{2}} \qquad \text{or:} \qquad \Psi_N(x,y,z...) \approx \frac{\psi_d(x) \, f\left(y^{\,\prime}-y\right) \Phi_d(y^{\,\prime},z\,,...)}{N(y) \sqrt{2}}$$



Now take:
$$\Psi_o(x, y, z...) = \frac{\psi_u(x)\Phi_u(y, z, ...) + \psi_d(x)\Phi_d(y, z, ...)}{\sqrt{2}}$$

Provided

$$\Phi_u(y,z,...)\Phi_d(y,z,...)\approx 0$$
 for all values of y and $f(y'-y)\approx 0$ $y'\neq y$

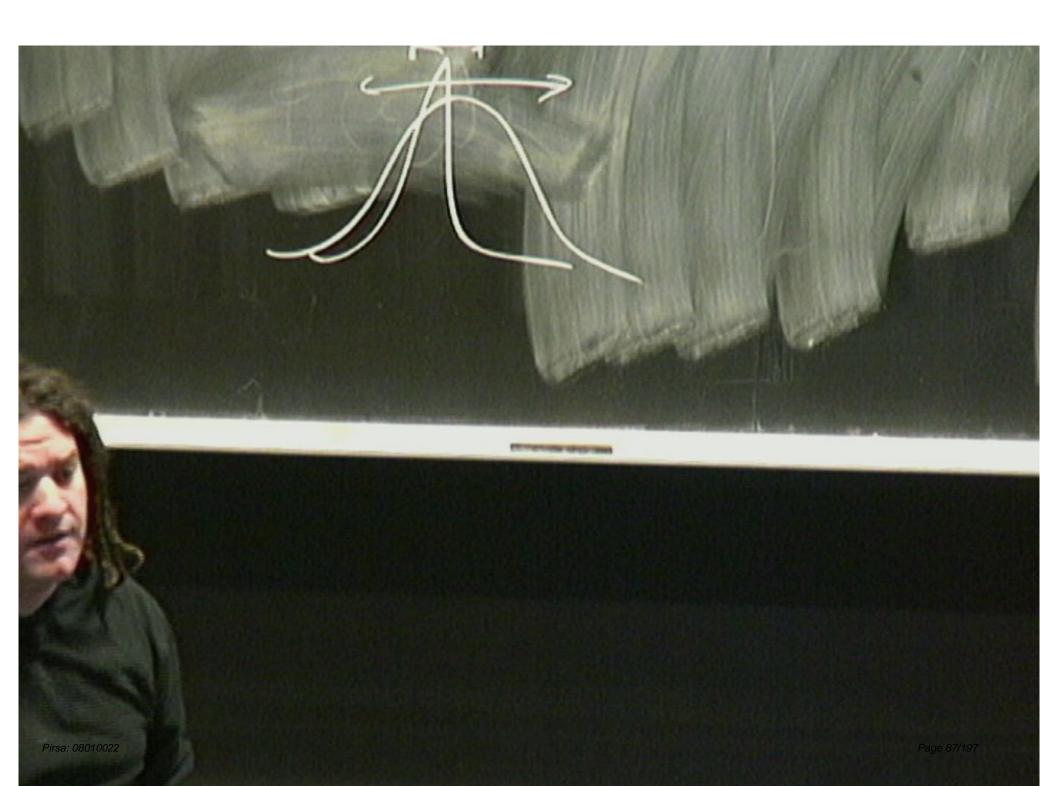
$$f(y'-y)\Phi_u(y,z,...)\approx 0$$
 or: $f(y'-y)\Phi_d(y,z,...)\approx 0$

$$\text{SO:} \quad \Psi_{N}(x,y,z...) = \frac{\psi_{u}(x) \, f(\,y\,'-\,y) \varPhi_{u}(\,y\,',\,z\,,...) + \psi_{d}(x) \, f(\,y\,'-\,y) \varPhi_{d}(\,y\,',\,z\,,...)}{N(\,y) \sqrt{2}}$$

$$\begin{split} \Psi_{N}(x,y,z...) \approx & \frac{\psi_{u}(x) f(y'-y) \Phi_{u}(y',z,...)}{N(y)\sqrt{2}} \quad \text{or:} \quad \Psi_{N}(x,y,z...) \approx & \frac{\psi_{d}(x) f(y'-y) \Phi_{d}(y',z,...)}{N(y)\sqrt{2}} \\ & \tau_{y,z...} \approx & \frac{10^{15} s}{10^{23}} = & 10^{-8} s \end{split}$$



 Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.









 Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.



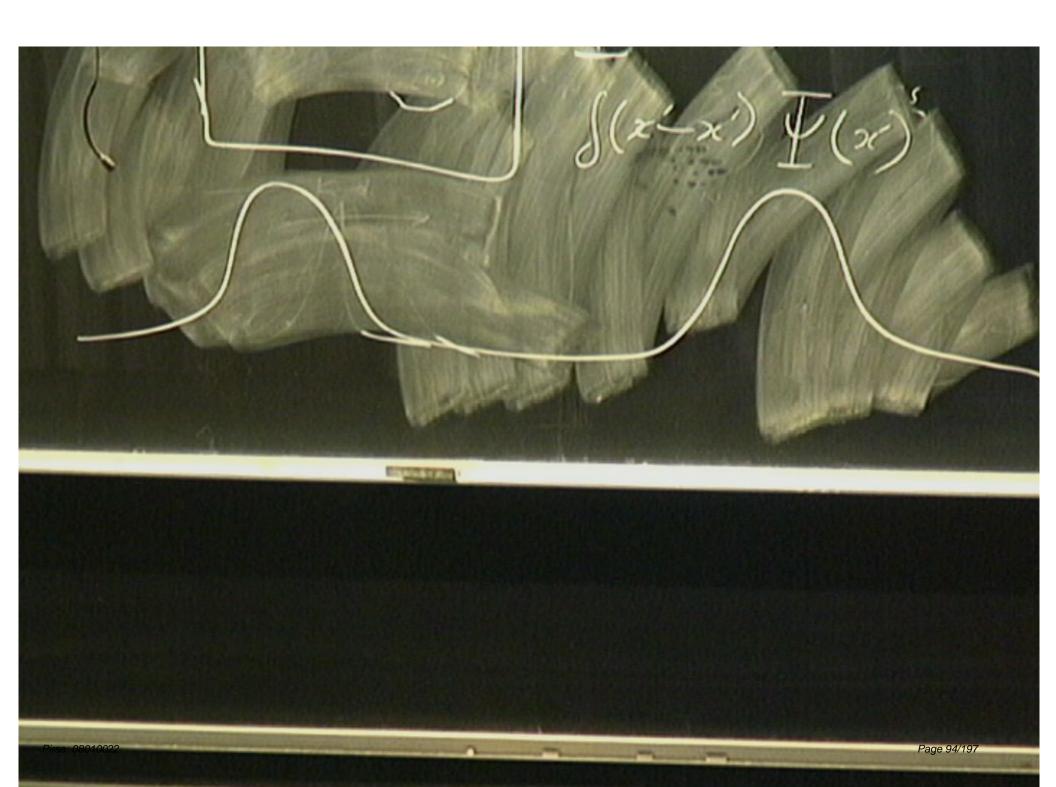
- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.

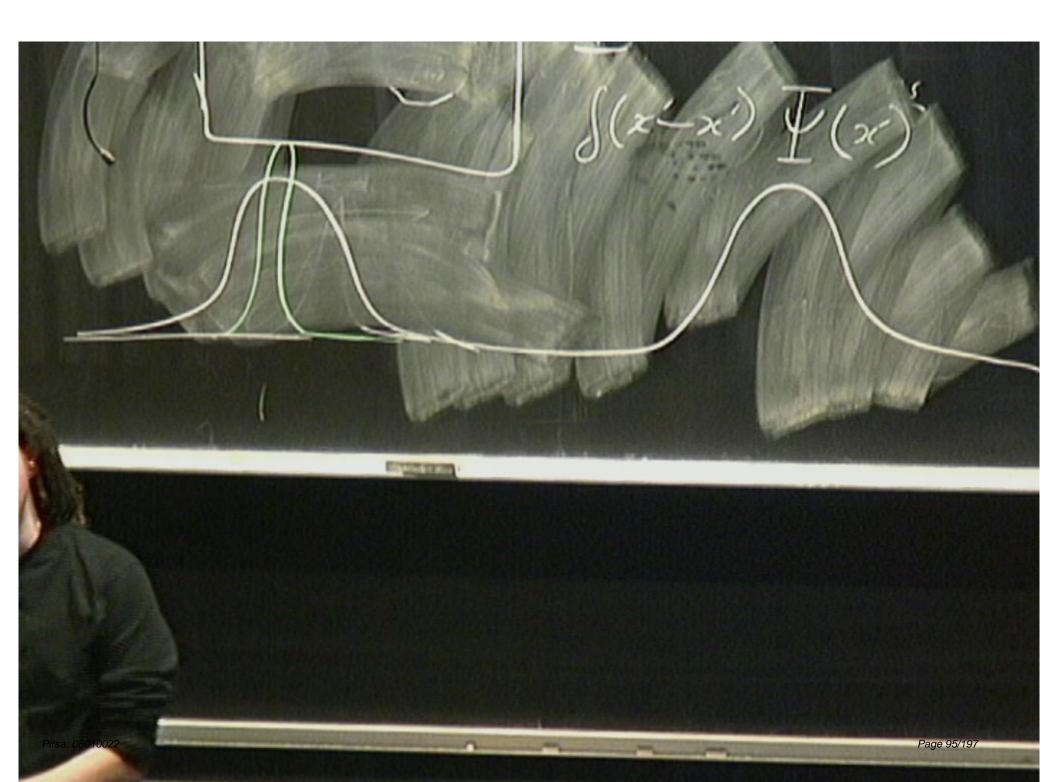


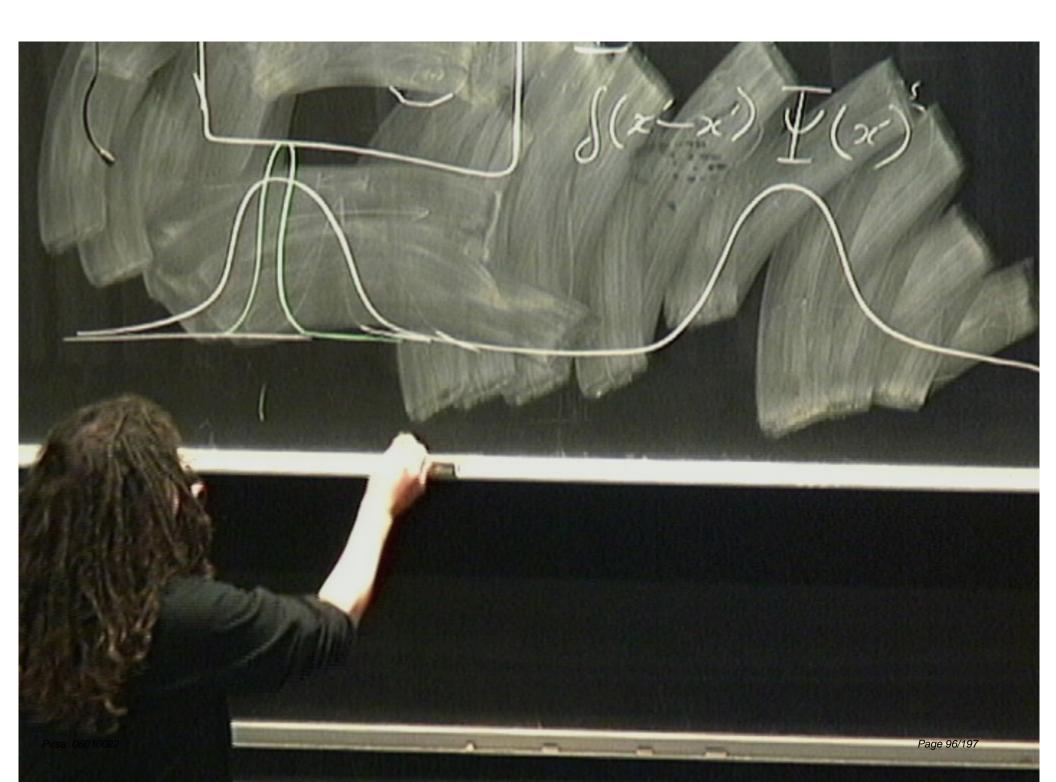
- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.

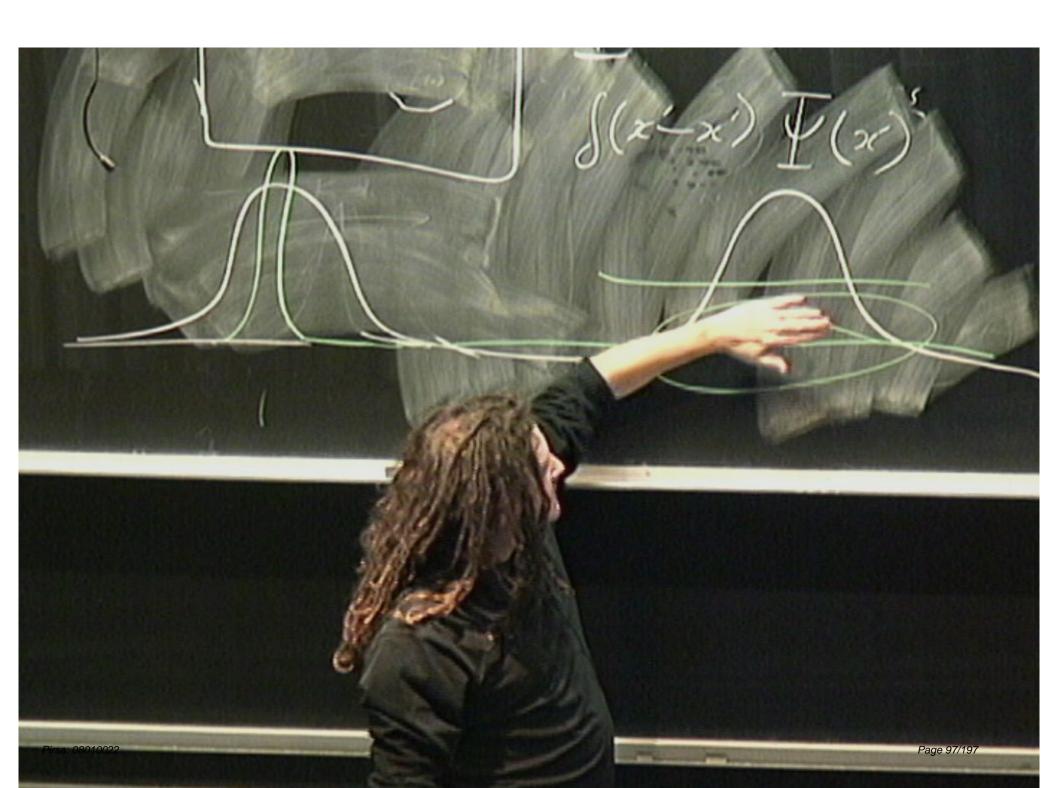


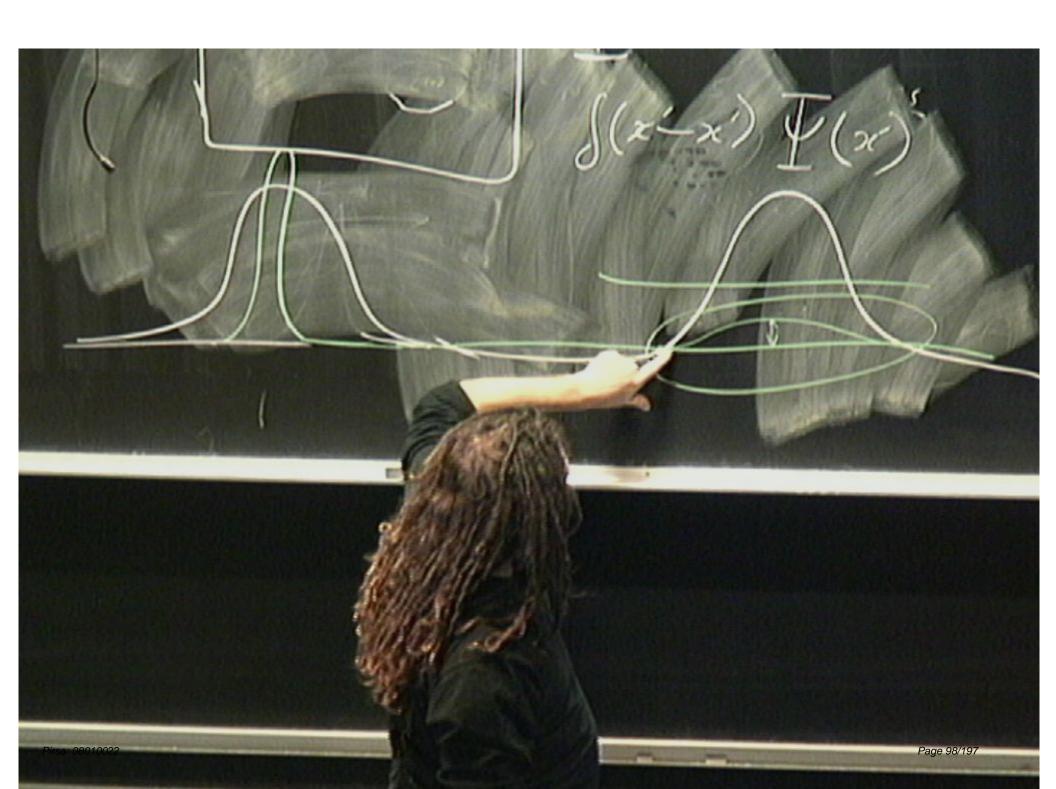
- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.

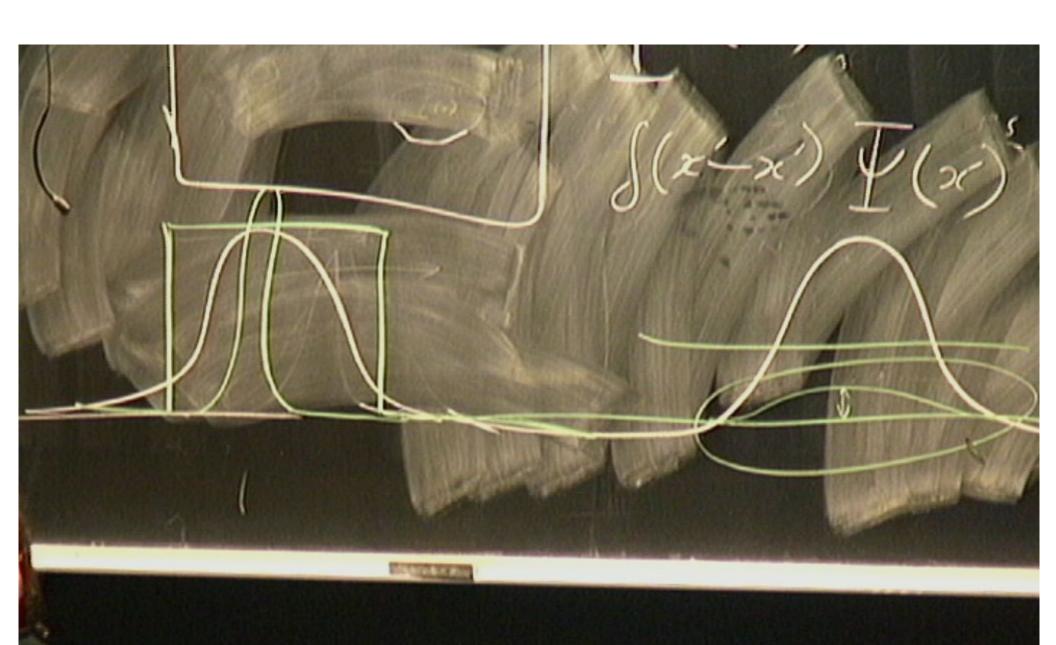


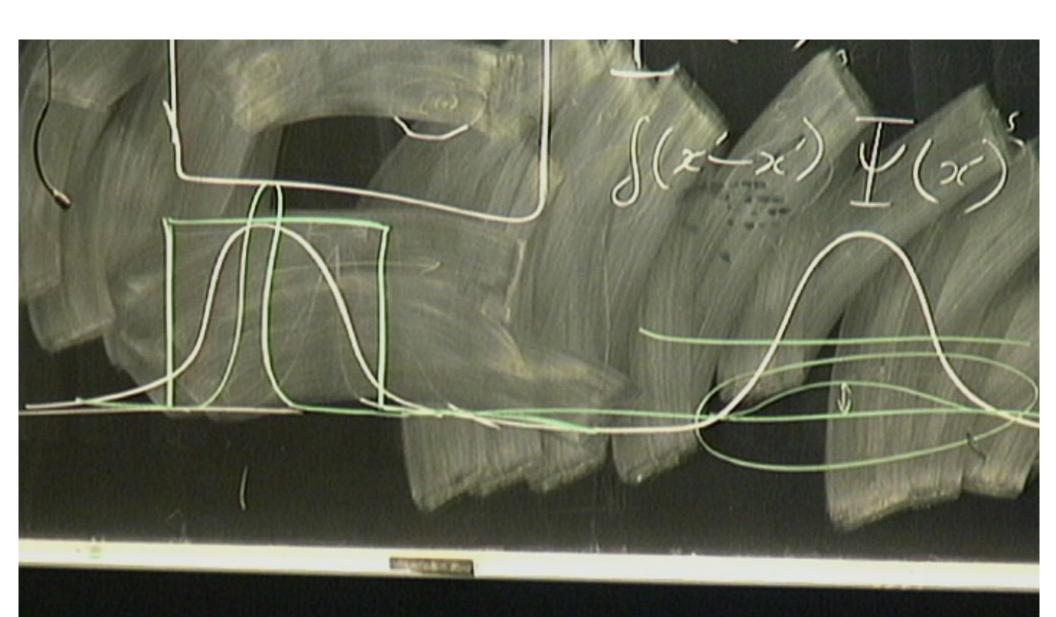














- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?



- Non-conservation of energy. Objective collapse models, that produce localised states, generically
 violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?

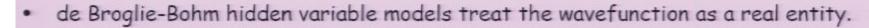


- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?
 - Flash ontology: moments of experience? What is real is only the event at the moment of the collapse.



- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?
 - Flash ontology: moments of experience? What is real is only the event at the moment of the collapse.
 - Matter field: extra structure to the world in addition to the wavefunction

Ψ epistemic theories







- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?
 - Flash ontology: moments of experience? What is real is only the event at the moment of the collapse.
 - Matter field: extra structure to the world in addition to the wavefunction



- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?
 - Flash ontology: moments of experience? What is real is only the event at the moment of the collapse.



Now take:
$$\Psi_{a}(x, y, z...) = \frac{\psi_{u}(x)\Phi_{u}(y, z, ...) + \psi_{d}(x)\Phi_{d}(y, z, ...)}{\sqrt{2}}$$

Provided

$$\Phi_u(y,z,...)\Phi_d(y,z,...)\approx 0$$
 for all values of y and $f(y'-y)\approx 0$ $y'\neq y$

$$f(y'-y)\Phi_u(y,z,...)\approx 0$$
 or: $f(y'-y)\Phi_d(y,z,...)\approx 0$

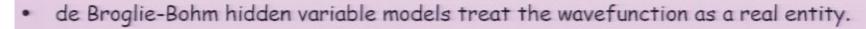
$$\text{SO:} \quad \Psi_{N}(x,y,z...) = \frac{\psi_{u}(x) \, f(\,y\,'-\,y) \Phi_{u}(\,y\,',z\,,...) + \psi_{d}(x) \, f(\,y\,'-\,y) \Phi_{d}(\,y\,',z\,,...)}{N(\,y) \sqrt{2}}$$

$$\begin{split} \Psi_N(x\,,\,y\,,\,z...) \approx & \frac{\psi_u(x)\,f\,(y\,'-y)\,\Phi_u(\,y\,',\,z\,,...)}{N\,(\,y\,)\sqrt{2}} \qquad \text{or:} \qquad \Psi_N(x\,,\,y\,,\,z...) \approx & \frac{\psi_d(x)\,f\,(\,y\,'-y)\,\Phi_d(\,y\,',\,z\,,...)}{N\,(\,y\,)\sqrt{2}} \\ & \tau_{y,z...} \approx & \frac{10^{15}\,s}{10^{23}} = 10^{-8}\,s \end{split}$$

Collapse models



- Non-conservation of energy. Objective collapse models, that produce localised states, generically violate the conservation of energy, even on average.
- Relativistic invariance. Explicit collapse models have generally been non-relativistic. Making a
 relativistically invariant collapse model presents difficulties. Recent progress has been made for noninteracting systems.
- Empirical predictions. Wavefunction collapse models generically produce situations where different predictions can be made to quantum theory. Experimentally probing these situations are hard, as environmental decoherence must be excluded.
- Tails and signalling Collapse models leave "tails" that include uncollapsed traces of unobserved outcomes. Eliminating these traces leads to violations of no-signalling.
- What is the ontology?
 - Bare theory : Everettian?
 - Flash ontology: moments of experience? What is real is only the event at the moment of the collapse.
 - Matter field: extra structure to the world in addition to the wavefunction







- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.



- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.
 - Is it possible to construct models in which the wavefunction is only a probability distribution over a microscopic reality?



- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.
 - Is it possible to construct models in which the wavefunction is only a probability distribution over a microscopic reality?
 - Wavefunction collapse might then correspond to an epistemic "updating" of probabilities. (Although measurement would still need to be invasive)



- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.
 - Is it possible to construct models in which the wavefunction is only a probability distribution over a microscopic reality?
 - Wavefunction collapse might then correspond to an epistemic "updating" of probabilities. (Although measurement would still need to be invasive)
 - · Empty waves would not be real

$$P(O_i|\psi) = \sum_{X} P(O_i|X) P_{\psi}(X)$$



- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.
 - Is it possible to construct models in which the wavefunction is only a probability distribution over a microscopic reality?
 - Wavefunction collapse might then correspond to an epistemic "updating" of probabilities. (Although measurement would still need to be invasive)
 - Empty waves would not be real

$$P(O_i|\psi) = \sum_{x} P(O_i|X) P_{\psi}(X)$$

For
$$\Psi$$
 ontic $X = (\Psi, x)$ $X \Rightarrow \Psi$



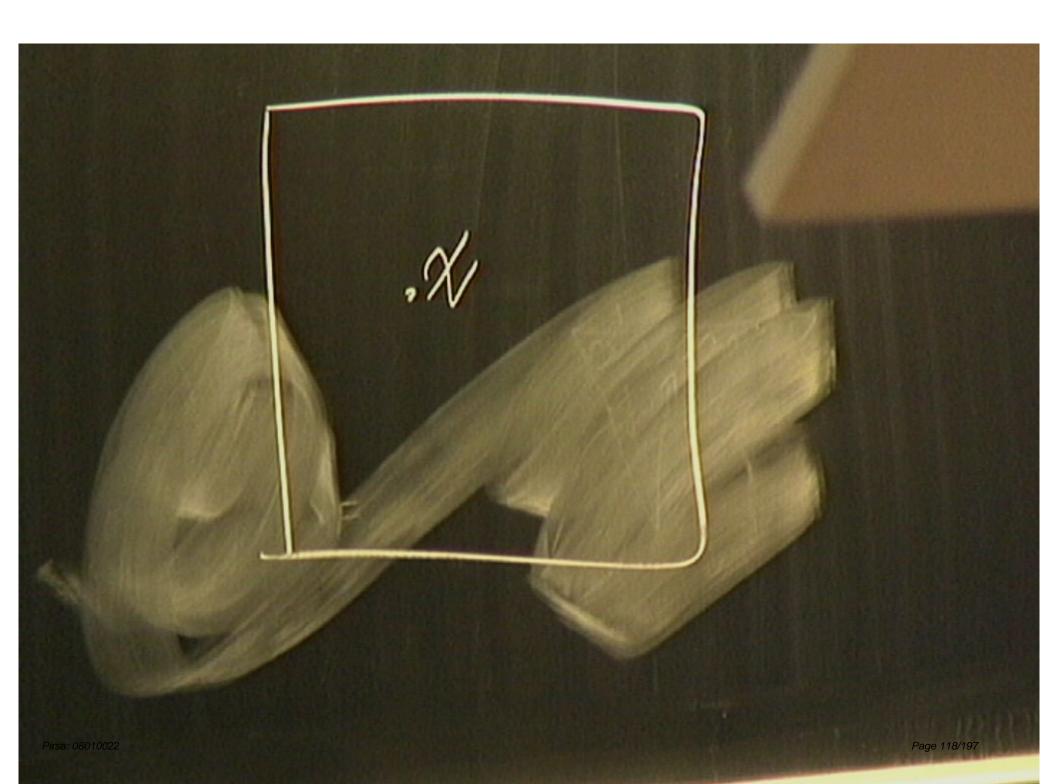
- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.
 - Is it possible to construct models in which the wavefunction is only a probability distribution over a microscopic reality?
 - Wavefunction collapse might then correspond to an epistemic "updating" of probabilities. (Although measurement would still need to be invasive)
 - Empty waves would not be real

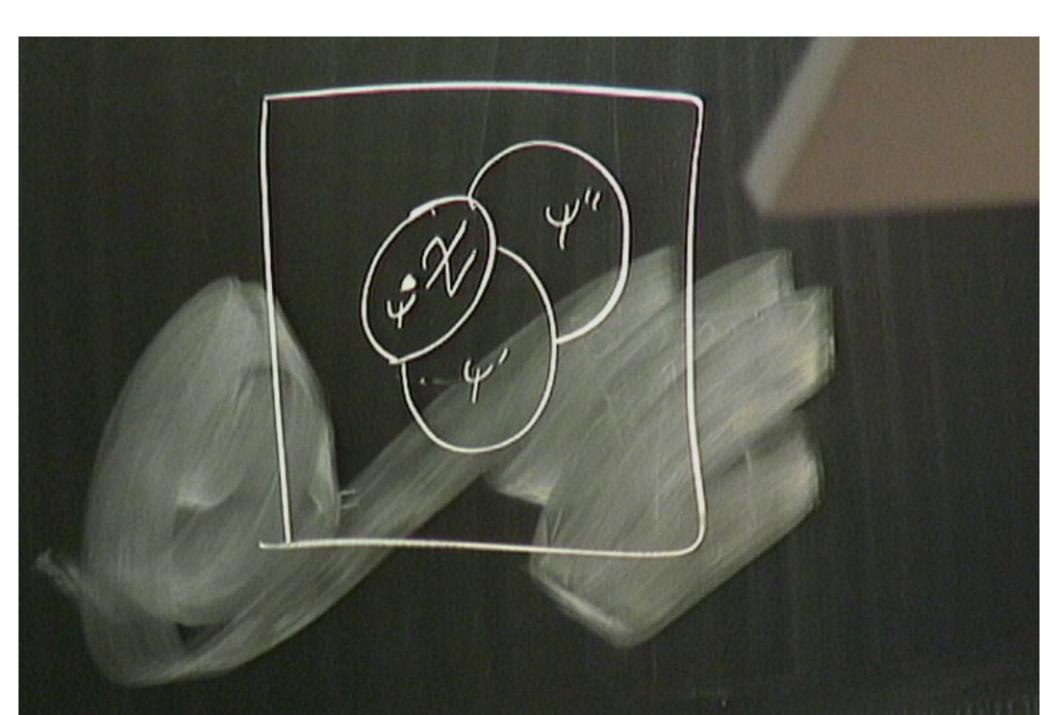
$$P(O_i|\psi) = \sum_{x} P(O_i|X) P_{\psi}(X)$$

For
$$\Psi$$
 ontic $X = (\Psi, x)$ $X \Rightarrow \Psi$

If:
$$P_{\psi}(X)P_{\psi}(X)=\delta_{\psi',\psi}(P_{\psi}(X))^2 \quad \forall X$$

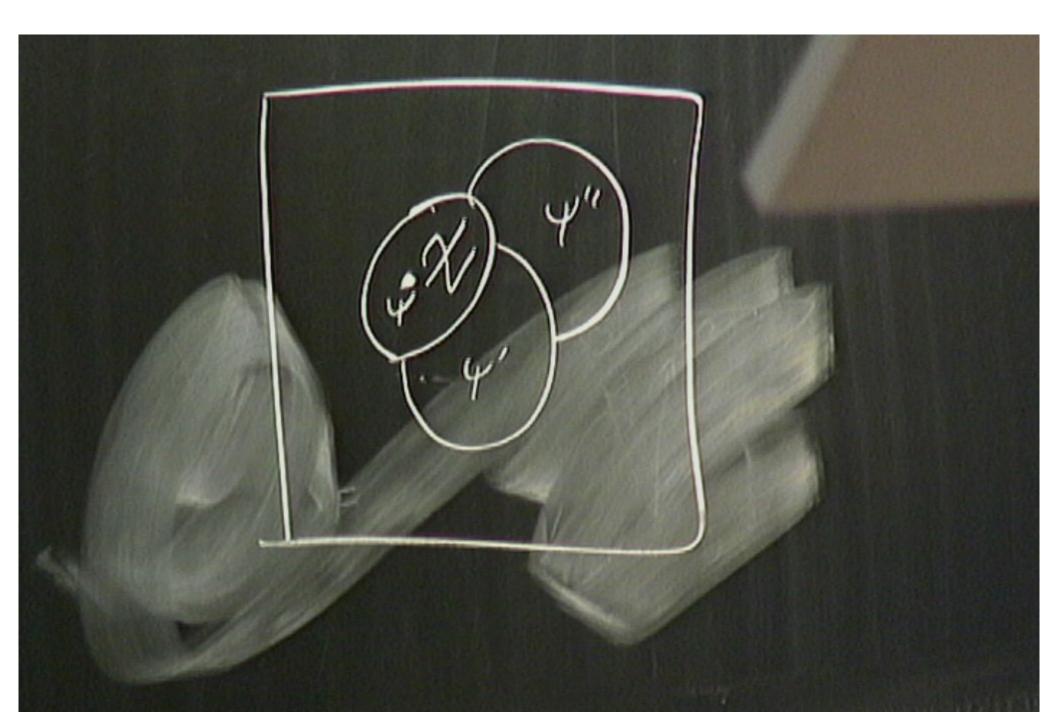
Pirsa: 08010022





Pirsa: 08010022

Page 119/197



Pirsa: 08010022



- de Broglie-Bohm hidden variable models treat the wavefunction as a real entity.
 - The wavefunction shares some properties with classical probability distributions.
 - Is it possible to construct models in which the wavefunction is only a probability distribution over a microscopic reality?
 - Wavefunction collapse might then correspond to an epistemic "updating" of probabilities. (Although measurement would still need to be invasive)
 - Empty waves would not be real

$$P(O_i|\psi) = \sum_{X} P(O_i|X) P_{\psi}(X)$$

For Ψ ontic $X = (\Psi, x)$ $X \Rightarrow \Psi$

If:
$$P_{\psi}(X)P_{\psi}(X) = \delta_{\psi',\psi}(P_{\psi}(X))^2 \quad \forall X$$

then: $\chi \Rightarrow \Psi$



$$\left[\phi_a,\phi_b,\phi_c\right]$$



$$\left[\phi_a,\phi_b,\phi_c\right]$$



$$\left[\phi_a,\phi_b,\phi_c\right]$$
 $V(\phi_a)=0,1$



$$\left[\phi_a,\phi_b,\phi_c\right]$$
 $V(\phi_a)=0,1$ $V(\phi_a)+V(\phi_b)+V(\phi_c)=1$



$$\left[\phi_a, \phi_b, \phi_c \right] \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_b) + V(\phi_c) = 1$$

$$\left[\phi_a, \phi_d, \phi_e \right]$$



$$\begin{bmatrix} \phi_a, \phi_b, \phi_c \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_b) + V(\phi_c) = 1$$

$$\begin{bmatrix} \phi_a, \phi_d, \phi_e \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_d) + V(\phi_e) = 1$$



$$\begin{bmatrix} \phi_a, \phi_b, \phi_c \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_b) + V(\phi_c) = 1$$

$$\begin{bmatrix} \phi_a, \phi_d, \phi_e \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_d) + V(\phi_e) = 1$$

- Value Definiteness
 - · All observables defined for a QM system have definite values at all times



$$\begin{bmatrix} \phi_a, \phi_b, \phi_c \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_b) + V(\phi_c) = 1$$

$$\begin{bmatrix} \phi_a, \phi_d, \phi_e \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_d) + V(\phi_e) = 1$$

- Value Definiteness
 - · All observables defined for a QM system have definite values at all times
- Non Contextuality
 - If a QM system possesses a property (value of an observable), then it does so
 independently of any measurement context, i.e. independently of how that value is
 eventually measured.

$$V(\phi_a)_{[\phi_a,\phi_b,\phi_c]} = V(\phi_a)_{[\phi_a,\phi_d,\phi_e]}$$



$$\begin{bmatrix} \phi_a, \phi_b, \phi_c \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_b) + V(\phi_c) = 1$$

$$\begin{bmatrix} \phi_a, \phi_d, \phi_e \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_d) + V(\phi_e) = 1$$

- Value Definiteness
 - · All observables defined for a QM system have definite values at all times
- Non Contextuality
 - If a QM system possesses a property (value of an observable), then it does so
 independently of any measurement context, i.e. independently of how that value is
 eventually measured.

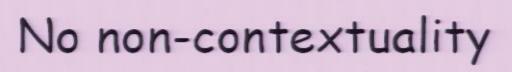


$$\begin{bmatrix} \phi_a, \phi_b, \phi_c \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_b) + V(\phi_c) = 1$$

$$\begin{bmatrix} \phi_a, \phi_d, \phi_e \end{bmatrix} \qquad V(\phi_a) = 0,1 \qquad V(\phi_a) + V(\phi_d) + V(\phi_e) = 1$$

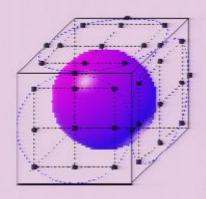
- Value Definiteness
 - · All observables defined for a QM system have definite values at all times
- Non Contextuality
 - If a QM system possesses a property (value of an observable), then it does so independently of any measurement context, i.e. independently of how that value is eventually measured.

$$V(\phi_a)_{[\phi_a,\phi_b,\phi_c]} = V(\phi_a)_{[\phi_a,\phi_d,\phi_e]}$$

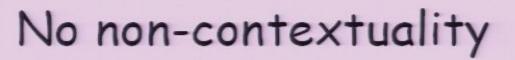




Example by Peres

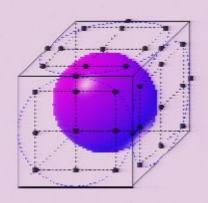


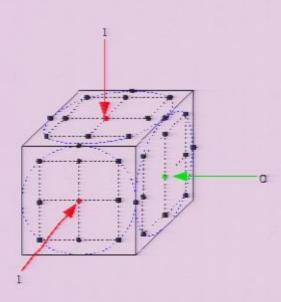
http://www.cs.auckland.ac.nz/~jas/one/freewill-theorem.html





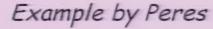
Example by Peres

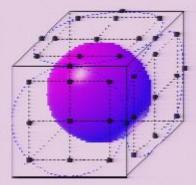


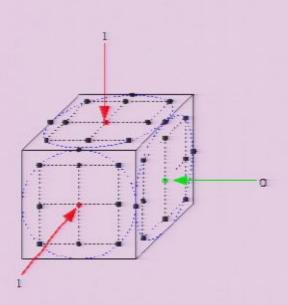


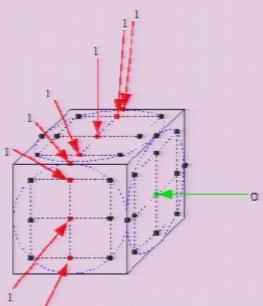
http://www.cs.auckland.ac.nz/~jas/one/freewill-theorem.html











http://www.cs.auckland.ac.nz/~jas/one/freewill-theorem.html

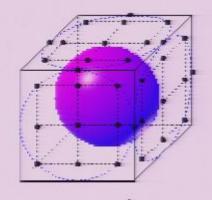
ArPirsa 08010022ction to Quantum Foundations

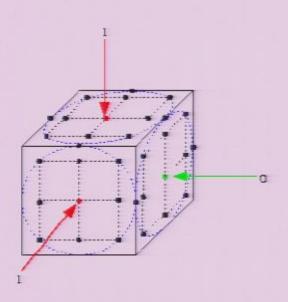
Lecture 4: Interpretation, Reformulation or Replacement?

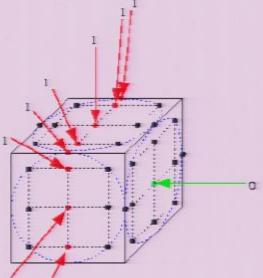
New Hopage 184197n Fundamental Physics

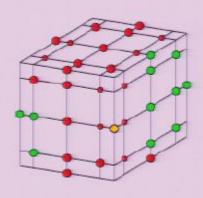


Example by Peres









http://www.cs.auckland.ac.nz/~jas/one/freewill-theorem.html

ArPisa 08010022ction to Quantum Foundations

Lecture 4: Interpretation, Reformulation or Replacement?

New Hopage 185497n Fundamental Physics





 Can quantum mechanics be better understood as a theory of information?



- Can quantum mechanics be better understood as a theory of information?
 - Quantum theory has always seemed to present restrictions upon what can be known about a system.



- Can quantum mechanics be better understood as a theory of information?
 - Quantum theory has always seemed to present restrictions upon what can be known about a system.
 - Perhaps quantum theory is simply about information itself, not information about something - the process of acquiring information creating the very information that is acquired.



- Can quantum mechanics be better understood as a theory of information?
 - Quantum theory has always seemed to present restrictions upon what can be known about a system.
 - Perhaps quantum theory is simply about information itself, not information about something - the process of acquiring information creating the very information that is acquired.
 - Restrictions on how much information may be known, or how much information may be stored, means that new information acquisition must invalidate old information.



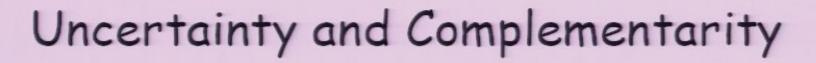
- Can quantum mechanics be better understood as a theory of information?
 - Quantum theory has always seemed to present restrictions upon what can be known about a system.
 - Perhaps quantum theory is simply about information itself, not information about something - the process of acquiring information creating the very information that is acquired.
 - Restrictions on how much information may be known, or how much information may be stored, means that new information acquisition must invalidate old information.
 - Physics as a sequence of yes-no questions (with limited storage)



- Can quantum mechanics be better understood as a theory of information?
 - Quantum theory has always seemed to present restrictions upon what can be known about a system.
 - Perhaps quantum theory is simply about information itself, not information about something - the process of acquiring information creating the very information that is acquired.
 - Restrictions on how much information may be known, or how much information may be stored, means that new information acquisition must invalidate old information.
 - Physics as a sequence of yes-no questions (with limited storage)
 - · Quantum state as a disposition to provide particular answers to questions



- Can quantum mechanics be better understood as a theory of information?
 - Quantum theory has always seemed to present restrictions upon what can be known about a system.
 - Perhaps quantum theory is simply about information itself, not information about something - the process of acquiring information creating the very information that is acquired.
 - Restrictions on how much information may be known, or how much information may be stored, means that new information acquisition must invalidate old information.
 - Physics as a sequence of yes-no questions (with limited storage)
 - Quantum state as a disposition to provide particular answers to questions
- "Asking a question" replaces "Making a measurement" but does this help?







- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)



 Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?
 - Simultaneous measurability of position and momentum by a single experimental arrangement?



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?
 - Simultaneous measurability of position and momentum by a single experimental arrangement?
- Reformulate as precise mathematical notions. Positive Operator Value measurements ("unsharp" measurements)



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?
 - Simultaneous measurability of position and momentum by a single experimental arrangement?
- Reformulate as precise mathematical notions. Positive Operator Value measurements ("unsharp" measurements)
 - Informationally Complete Measurements (and others)



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?
 - Simultaneous measurability of position and momentum by a single experimental arrangement?
- Reformulate as precise mathematical notions. Positive Operator Value measurements ("unsharp" measurements)
 - Informationally Complete Measurements (and others)
- What of other complementary observables?



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?
 - Simultaneous measurability of position and momentum by a single experimental arrangement?
- Reformulate as precise mathematical notions. Positive Operator Value measurements ("unsharp" measurements)
 - Informationally Complete Measurements (and others)
- What of other complementary observables?
 - Energy-Time uncertainty relationships



- Heisenbergy Uncertainty Relations can be derived as a mathematical consequence of Hilbert space structure. But what do they mean?
 - Disturbance of a momentum value by making a position measurement (and vice versa?)
 - Spread in variance of position and momentum values when position and momentum measurements are made on equivalently prepared systems?
 - Simultaneous measurability of position and momentum by a single experimental arrangement?
- Reformulate as precise mathematical notions. Positive Operator Value measurements ("unsharp" measurements)
 - Informationally Complete Measurements (and others)
- What of other complementary observables?
 - Energy-Time uncertainty relationships
 - Photon Phase-Number statistics





Schrodinger Picture
$$i\hbar \frac{\partial \psi}{\partial t} = \hat{H}\psi$$



Schrodinger Picture
$$i\hbar \frac{\partial \psi}{\partial t} = \hat{H}\psi$$
 Heisenberg Picture $\frac{d\hat{A}}{dt} = \frac{1}{i\hbar}[\hat{H},\hat{A}] + \frac{\partial \hat{A}}{\partial t}$



Schrodinger Picture
$$i\hbar\frac{\partial\psi}{\partial t}=\hat{H}\psi$$
 Heisenberg Picture $\frac{d\hat{A}}{dt}=\frac{1}{i\hbar}[\hat{H},\hat{A}]+\frac{\partial\hat{A}}{\partial t}$ Path Integral $\psi(x,t)=\int_{x'}\int_{q(t)}e^{\frac{i}{\hbar}\int_{t'}^{t}S(q(t),t)dt}D(q(t))\psi(x',t')dx'$



Perhaps the mathematical formulation of quantum theory is wrong?

Schrodinger Picture
$$i\hbar\frac{\partial\psi}{\partial t} = \hat{H}\psi$$
 Heisenberg Picture $\frac{d\hat{A}}{dt} = \frac{1}{i\hbar}[\hat{H},\hat{A}] + \frac{\partial\hat{A}}{\partial t}$

Path Integral
$$\psi(x,t) = \int_{x'} \int_{q(t)} e^{\frac{i}{\hbar} \int_{t'}^{t} S(q(t),t) dt} D(q(t)) \psi(x',t') dx'$$

Histories Formalism:

A history is a sequence of events $H_a = E_i^0 E_i^1 \dots E_k^{n-1} E_i^n$

$$H_{\alpha} = E_i^0 E_j^1 \dots E_k^{n-1} E_l^n$$



Perhaps the mathematical formulation of quantum theory is wrong?

Schrodinger Picture
$$i\hbar\frac{\partial\psi}{\partial t} = \hat{H}\psi$$
 Heisenberg Picture $\frac{d\hat{A}}{dt} = \frac{1}{i\hbar}[\hat{H},\hat{A}] + \frac{\partial\hat{A}}{\partial t}$

Path Integral
$$\psi(x,t) = \int_{x'} \int_{q(t)} e^{\frac{i}{\hbar} \int_{t'}^{t} S(q(t),t) dt} D(q(t)) \psi(x',t') dx'$$

Histories Formalism:

A history is a sequence of events

$$H_a = E_i^0 E_j^1 \dots E_k^{n-1} E_l^n$$

"Weight" of a history is:

$$W(H_{\alpha}) = \text{Tr}[H_{\alpha}H_{\alpha}^{\dagger}]$$



· Perhaps the mathematical formulation of quantum theory is wrong?

Schrodinger Picture
$$i\hbar\frac{\partial\psi}{\partial t} = \hat{H}\psi$$
 Heisenberg Picture $\frac{d\hat{A}}{dt} = \frac{1}{i\hbar}[\hat{H},\hat{A}] + \frac{\partial\hat{A}}{\partial t}$

Path Integral
$$\psi(x,t) = \int_{x'} \int_{q(t)} e^{\frac{i}{\hbar} \int_{t'}^{t} S(q(t),t) dt} D(q(t)) \psi(x',t') dx'$$

Histories Formalism:

A history is a sequence of events

"Weight" of a history is:

Probability of a history is:

$$H_{\alpha} = E_i^0 E_j^1 \dots E_k^{n-1} E_l^n$$

$$W(H_{\alpha}) = \text{Tr}[H_{\alpha}H_{\alpha}^{\dagger}]$$

$$P(H_{\alpha}) = \frac{W(H_{\alpha})}{\sum_{\beta} W(H_{\beta})}$$



Perhaps the mathematical formulation of quantum theory is wrong?

Schrodinger Picture
$$i\hbar\frac{\partial\psi}{\partial t} = \hat{H}\psi$$
 Heisenberg Picture $\frac{d\hat{A}}{dt} = \frac{1}{i\hbar}[\hat{H},\hat{A}] + \frac{\partial\hat{A}}{\partial t}$

Path Integral
$$\psi(x,t) = \int_{x'} \int_{q(t)} e^{\frac{i}{\hbar} \int_{t'}^{t} S(q(t),t) dt} D(q(t)) \psi(x',t') dx'$$

Histories Formalism:

A history is a sequence of events

$$H_{\alpha} = E_i^0 E_j^1 \dots E_k^{n-1} E_l^n$$

"Weight" of a history is:

$$W(H_{\alpha}) = \text{Tr}[H_{\alpha}H_{\alpha}^{\dagger}]$$

Probability of a history is:

$$P(H_{\alpha}) = \frac{W(H_{\alpha})}{\sum_{\beta} W(H_{\beta})}$$

If a consistency condition is met, all normal probability rules can be used.

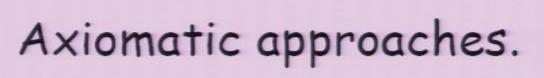
$$\operatorname{Re}\left\{\operatorname{Tr}\left[H_{\alpha}H_{\beta}^{\dagger}\right]\right\} = 0, \forall \alpha, \beta, \alpha \neq \beta$$

ArPiisa 08000022ction to Quantum Foundations

Lecture 4: Interpretation, Reformulation or Replacement?

New - Page 162/197 Fundamental Physics







· What is different to classical theories?



- · What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.



- What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.
 - · Clifton Bub Halvorson.
 - "No cloning. No signalling. No bit commitment" to construct a principle quantum theory



- What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.
 - Clifton Bub Halvorson.
 - "No cloning. No signalling. No bit commitment" to construct a principle quantum theory
 - Generalised probability.



- What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.
 - · Clifton Bub Halvorson.
 - "No cloning. No signalling. No bit commitment" to construct a principle quantum theory



- What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.
 - Clifton Bub Halvorson.
 - "No cloning. No signalling. No bit commitment" to construct a principle quantum theory
 - Generalised probability.



- What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.
 - · Clifton Bub Halvorson.
 - "No cloning. No signalling. No bit commitment" to construct a principle quantum theory
 - Generalised probability.
 - Weaken Kolmogorov axioms.



- What is different to classical theories?
 - Principle theories, such as relativity and thermodynamics, constructed from prohibitions.
 - · Clifton Bub Halvorson.
 - "No cloning. No signalling. No bit commitment" to construct a principle quantum theory
 - Generalised probability.
 - Weaken Kolmogorov axioms.
 - · Generalised logic.
 - Propositions about a system do not form a Boolean lattice.

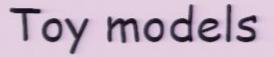
Toy models



Toy models

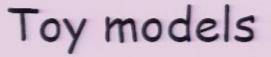


 Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.



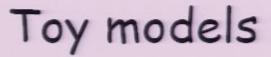


- Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.
 - Local hidden variable models which can simulate teleportation or dense coding, despite the fact that quantum theory requires entanglement to do so.



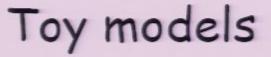


- Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.
 - Local hidden variable models which can simulate teleportation or dense coding, despite the fact that quantum theory requires entanglement to do so.
 - Theories which do not violate Bell inequalities, but do violate higher order inequalities.



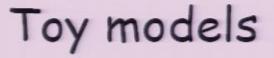


- Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.
 - Local hidden variable models which can simulate teleportation or dense coding, despite the fact that quantum theory requires entanglement to do so.
 - Theories which do not violate Bell inequalities, but do violate higher order inequalities.
 - Possible theories which violate Bell inequalities but do not violate higher order inequalities.



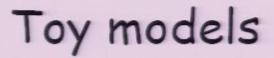


- Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.
 - Local hidden variable models which can simulate teleportation or dense coding, despite the fact that quantum theory requires entanglement to do so.
 - Theories which do not violate Bell inequalities, but do violate higher order inequalities.
 - Possible theories which violate Bell inequalities but do not violate higher order inequalities.
- Constructs which do things quantum mechanics cannot do.





- Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.
 - Local hidden variable models which can simulate teleportation or dense coding, despite the fact that quantum theory requires entanglement to do so.
 - Theories which do not violate Bell inequalities, but do violate higher order inequalities.
 - Possible theories which violate Bell inequalities but do not violate higher order inequalities.
- Constructs which do things quantum mechanics cannot do.
 - Popescu-Rohrlich non-local boxes, which are more non-local than quantum theory, although still do not permit signalling.





- Constructs which do not reproduce all of quantum theory but which can reproduce some characteristic quantum effects.
 - Local hidden variable models which can simulate teleportation or dense coding, despite the fact that quantum theory requires entanglement to do so.
 - Theories which do not violate Bell inequalities, but do violate higher order inequalities.
 - Possible theories which violate Bell inequalities but do not violate higher order inequalities.
- Constructs which do things quantum mechanics cannot do.
 - Popescu-Rohrlich non-local boxes, which are more non-local than quantum theory, although still do not permit signalling.
 - Inequalities not violated by QT but violated by other possible theories.

And more ...





- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?



- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian



- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian
- Relativistic Mechanics
 - Relationship between quantum theory and relativity



- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian
- Relativistic Mechanics
 - Relationship between quantum theory and relativity
 - Many body relativity, the "no-interaction theorem", Wheeler-Feynman electrodynamics



- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian
- Relativistic Mechanics
 - Relationship between quantum theory and relativity
 - Many body relativity, the "no-interaction theorem", Wheeler-Feynman electrodynamics
 - Hegerfeldt box paradox, localisability
 - Quantum Field theory



- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian
- Relativistic Mechanics
 - Relationship between quantum theory and relativity
 - Many body relativity, the "no-interaction theorem", Wheeler-Feynman electrodynamics
 - Hegerfeldt box paradox, localisability
 - Quantum Field theory
 - Haag's Theorem, Dyson's Theorem, Spin-statistics and identical particles



Emergent classicality

- If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian

Relativistic Mechanics

- Relationship between quantum theory and relativity
 - Many body relativity, the "no-interaction theorem", Wheeler-Feynman electrodynamics
 - Hegerfeldt box paradox, localisability
- Quantum Field theory
 - Haag's Theorem, Dyson's Theorem, Spin-statistics and identical particles
 - Axiomatic and Algebraic Quantum Field Theory



- Emergent classicality
 - If the world is fundamentally quantum in behaviour, why does the everyday world behave so classically?
 - Ehrenfest theorem, quantum chaos, decoherence, restrictions on observables and hamiltonian
- Relativistic Mechanics
 - Relationship between quantum theory and relativity
 - Many body relativity, the "no-interaction theorem", Wheeler-Feynman electrodynamics
 - Hegerfeldt box paradox, localisability
 - Quantum Field theory
 - Haag's Theorem, Dyson's Theorem, Spin-statistics and identical particles
 - Axiomatic and Algebraic Quantum Field Theory
- Experimental tests
 - precision testing of macroscopic interference, collapse models, non-equilibrium hidden variables...

ArPisatee000022ction to Quantum Foundations
Lecture 4: Interpretation, Reformulation or Replacement?

New HoPages1884197n Fundamental Physics





 Quantum phenomena can be understood, but there seems no easy or "right" way to do so.



- Quantum phenomena can be understood, but there seems no easy or "right" way to do so.
 - All approaches involve something that seems, intuitively, unreasonable.



- Quantum phenomena can be understood, but there seems no easy or "right" way to do so.
 - All approaches involve something that seems, intuitively, unreasonable.
 - Progress can be (and has been) made in understanding what combination of assumptions are tenable (and what are not).



- Quantum phenomena can be understood, but there seems no easy or "right" way to do so.
 - All approaches involve something that seems, intuitively, unreasonable.
 - Progress can be (and has been) made in understanding what combination of assumptions are tenable (and what are not).
 - Experiments can be (and have been) performed relevant to this.



- Quantum phenomena can be understood, but there seems no easy or "right" way to do so.
 - All approaches involve something that seems, intuitively, unreasonable.
 - Progress can be (and has been) made in understanding what combination of assumptions are tenable (and what are not).
 - Experiments can be (and have been) performed relevant to this.
- In any attempt to understand quantum phenomena



- Quantum phenomena can be understood, but there seems no easy or "right" way to do so.
 - All approaches involve something that seems, intuitively, unreasonable.
 - Progress can be (and has been) made in understanding what combination of assumptions are tenable (and what are not).
 - Experiments can be (and have been) performed relevant to this.
- In any attempt to understand quantum phenomena
 - Be logically consistent
 - Be conceptually clear and coherent



- Quantum phenomena can be understood, but there seems no easy or "right" way to do so.
 - All approaches involve something that seems, intuitively, unreasonable.
 - Progress can be (and has been) made in understanding what combination of assumptions are tenable (and what are not).
 - Experiments can be (and have been) performed relevant to this.
- In any attempt to understand quantum phenomena
 - Be logically consistent
 - Be conceptually clear and coherent
 - Be compatible with observed phenomena



"One way or another, God has played us a nasty trick.

The voice of nature has always been faint,
but in this case it speaks in riddles and mumbles as well."

T. Maudlin "Quantum Non-locality and Relativity"