

Title: Entangling Power of an Expanding Universe

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Abstract: Quantum fields in the Minkowski vacuum are entangled with respect to local field modes. This entanglement can be swapped to spatially separated quantum systems using standard local couplings. A single, inertial field detector in the exponentially expanding (de Sitter) vacuum responds as if it were bathed in thermal radiation in a Minkowski universe. Using two inertial detectors, interactions with the field in the thermal case will entangle certain detector pairs that would not become entangled in the corresponding de Sitter case. The two universes can thus be distinguished by their entangling power.

Entangling Power of an Expanding Universe

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Information in Quantum Field Theory

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 - Hawking radiation

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Operational Approach

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$$10 \gamma_{\frac{1}{2}} \otimes 10 \gamma_{\frac{1}{2}}$$

Operational Approach

$$|0\rangle_i \otimes |0\rangle_{ii}$$

$$\phi(x, t)$$

$$|0\rangle_i \otimes |0\rangle_n$$

$$\int dx dt \phi(x, t)$$

$$|0\rangle_{\vec{r}} \otimes |0\rangle_{\vec{r}'}$$

$$\int dx dt \Phi(x, t)$$

$$|0\rangle_{\tilde{p}} \otimes |0\rangle_{\tilde{p}'}$$

$$\int dx dt \phi(x, t)$$

$$\tilde{\phi}(p, E)$$

$$|0\rangle_{\vec{r}} \langle 0|_{\vec{r}}$$

$$\int dx dt \phi(x, t)$$

$$\tilde{\phi}(p, E) = \int dx dt \phi(x, t) e^{-i p x - i E t}$$

Quantum fields have entangling power

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 - Local modes are entangled in the vacuum state
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 - Reznik, *et al.*

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- Entangling power therefore *measures properties of the field and of spacetime*

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 - But maybe...

Question

Can the universes be distinguished
by their *entangling power*?

Entangling resource: quantum field

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- Equation of motion ($m = 0$):

$$\left[\square_x + m^2 + \frac{1}{6}R(x) \right] \phi(x) = 0$$

Two cases

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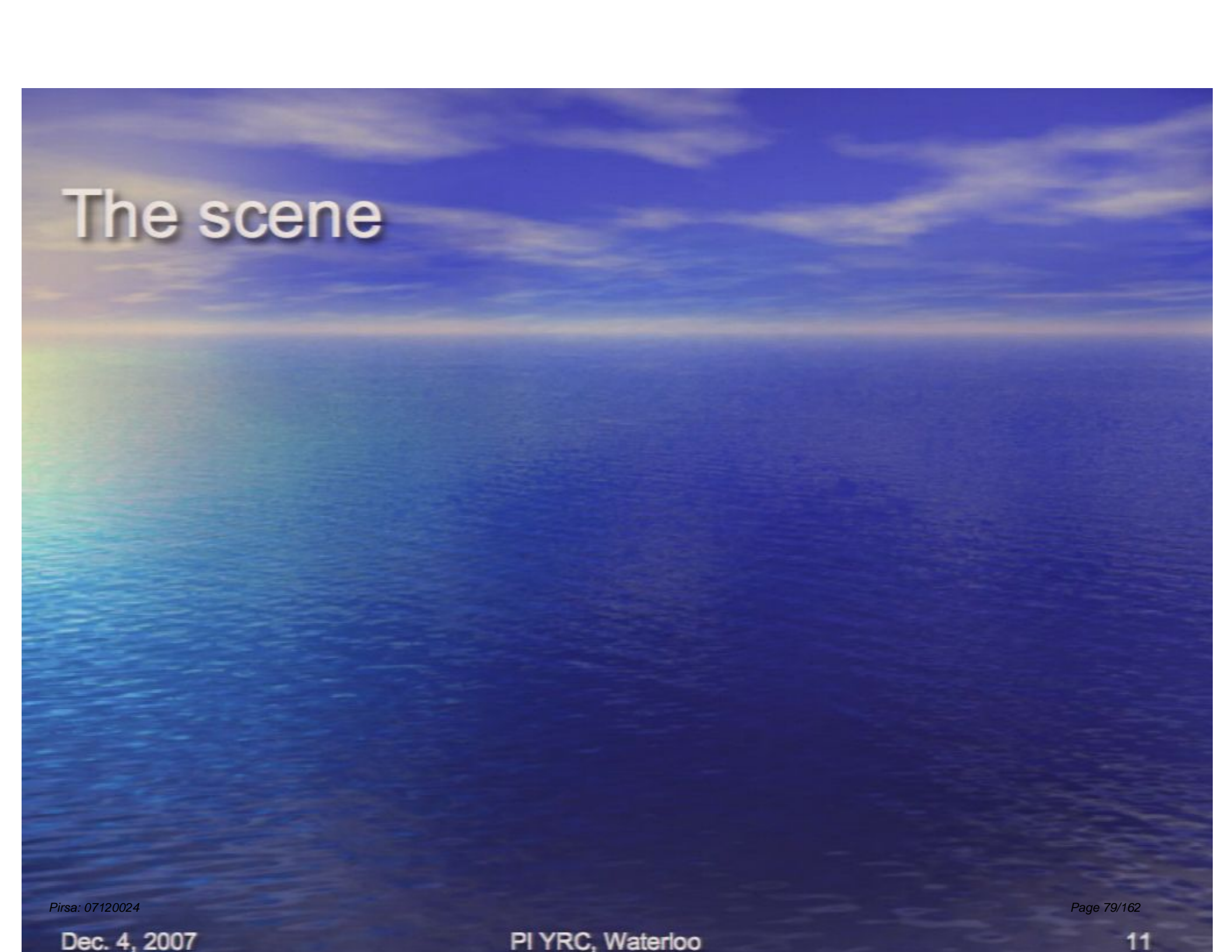
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 - $T_{\text{GH}} \sim (\text{CMB temp}) = 2.7 \text{ K} \rightarrow \kappa^{-1} \sim 0.1 \text{ mm}$

The scene



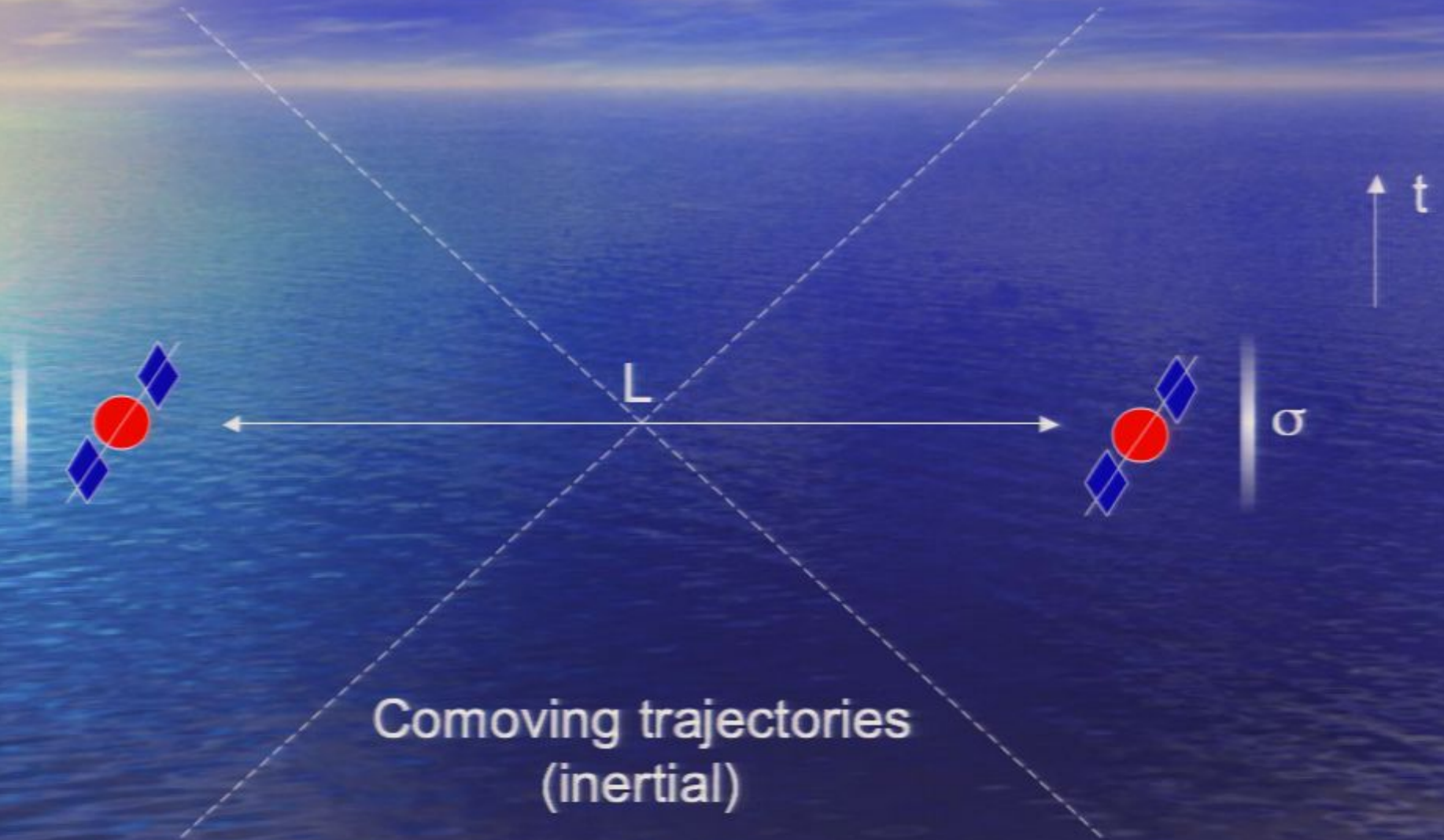
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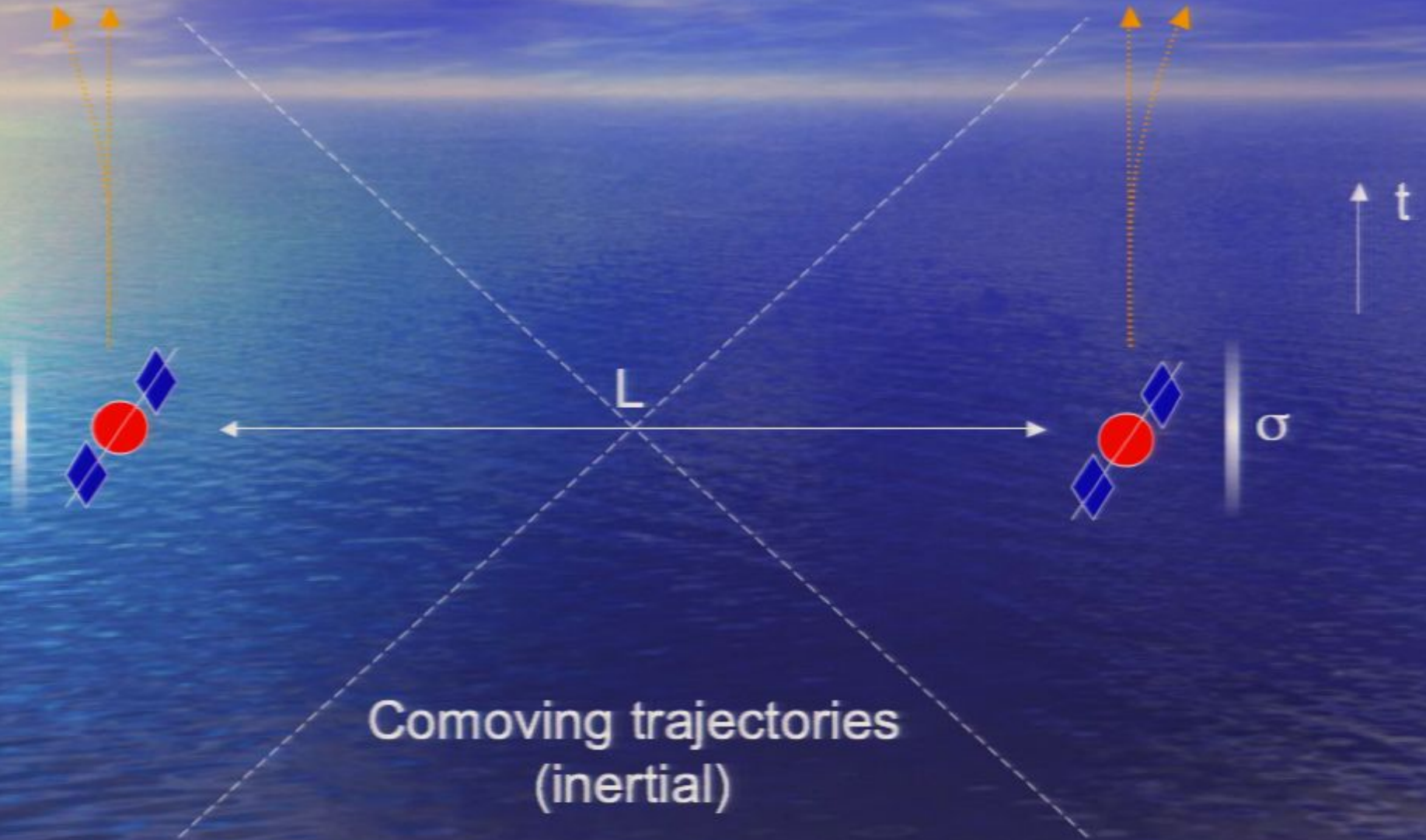


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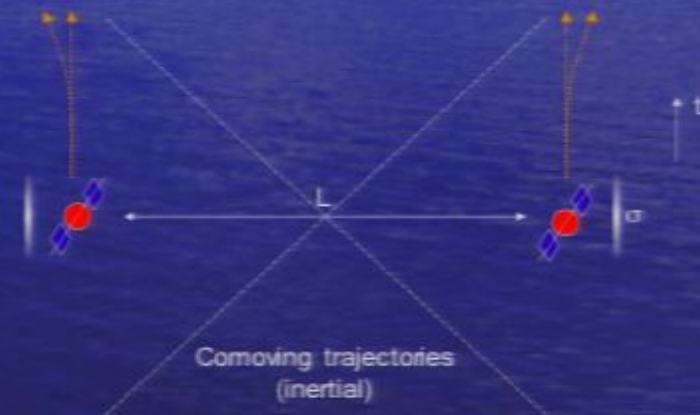


Comoving trajectories
(inertial)

The scene

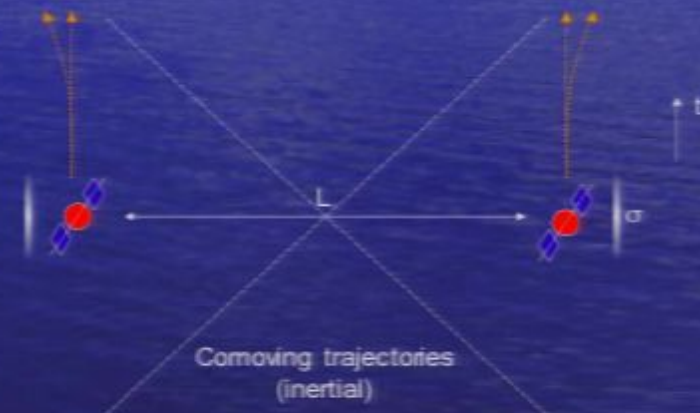


Two detectors



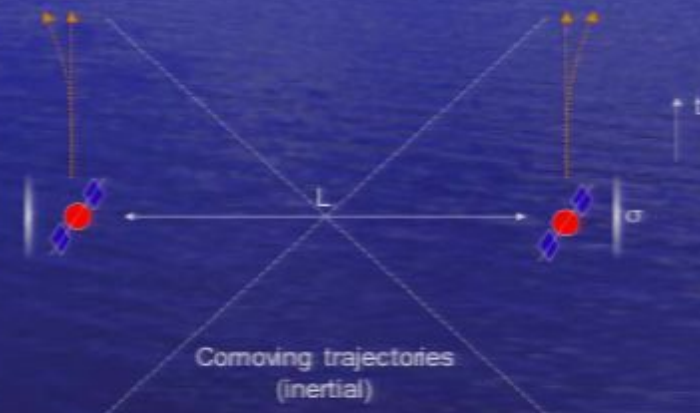
Two detectors

- Qubits with tunable energy gap (Ω)



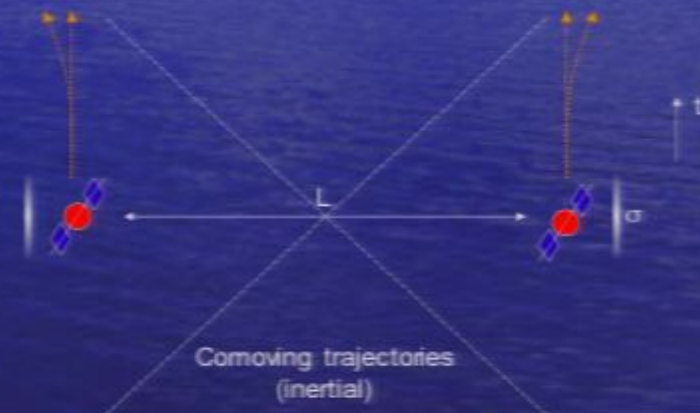
Two detectors

- Qubits with tunable energy gap (Ω)
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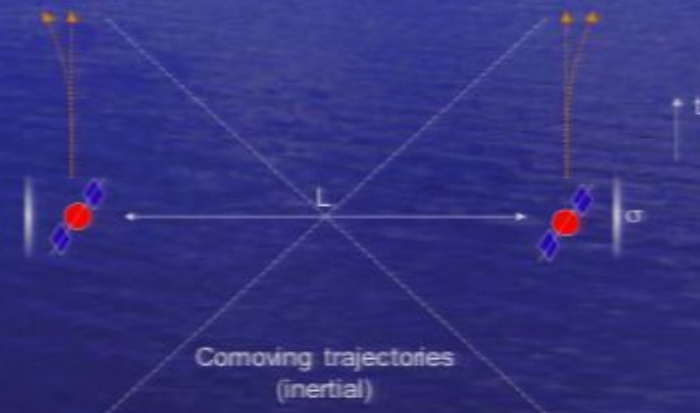
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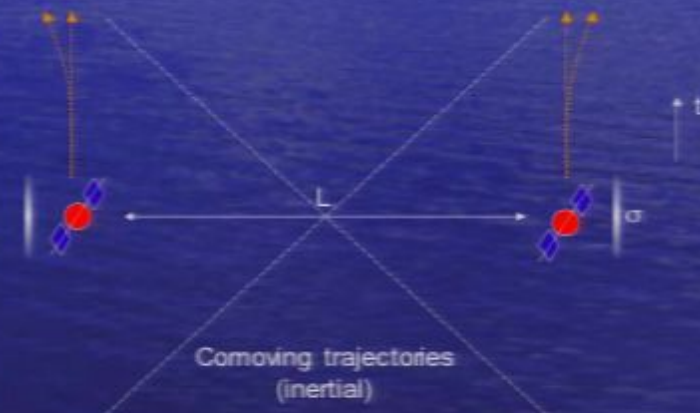
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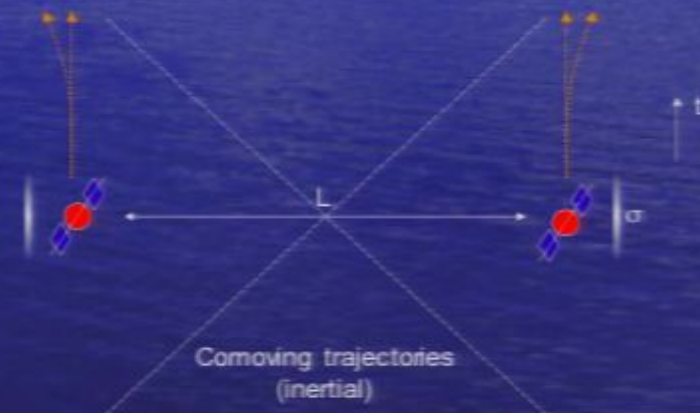
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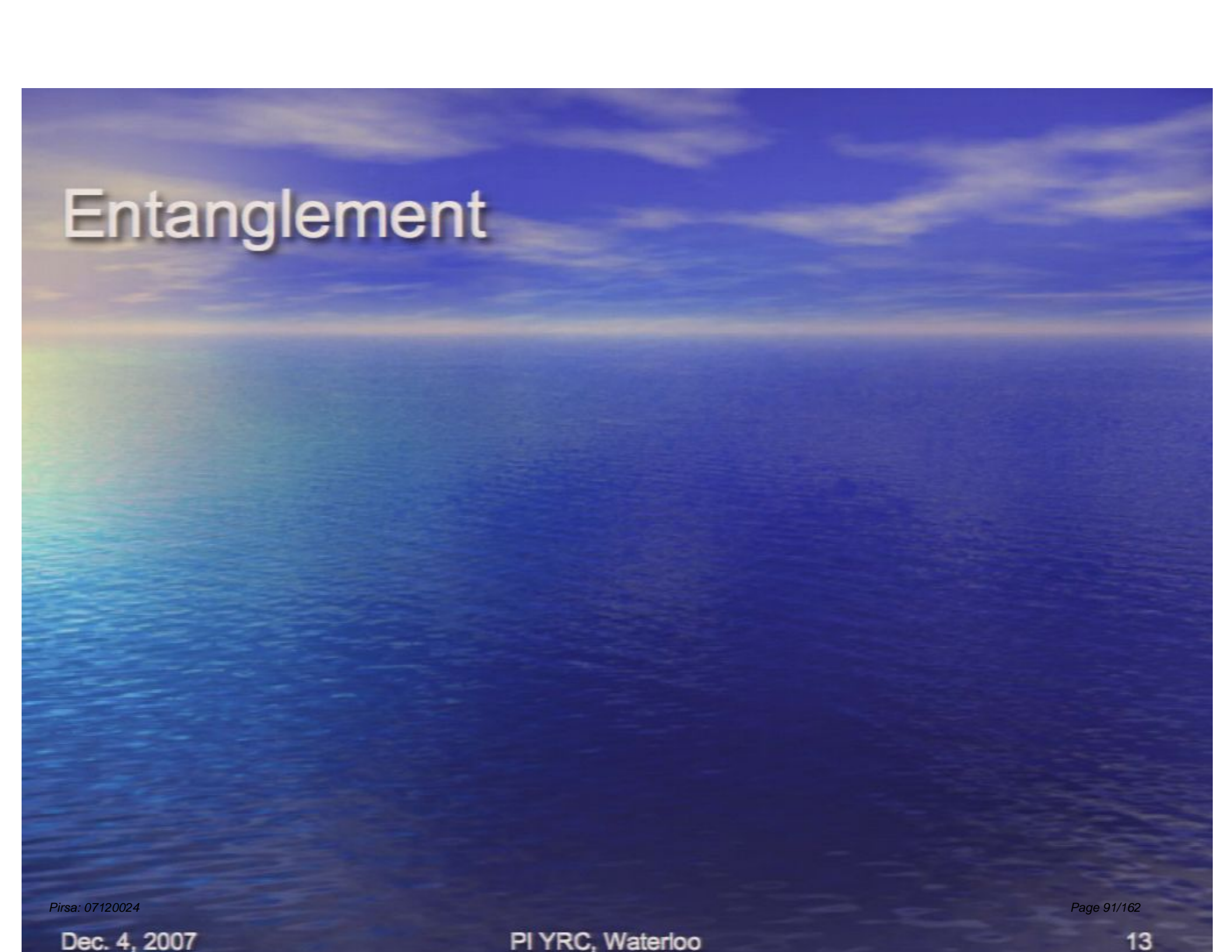
Two detectors

- Qubits with tunable energy gap (Ω)
- Simple, local field coupling
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- Interaction Hamiltonian:

$$H_I(\tau) = \eta(\tau)\phi(x(\tau)) \left(e^{+i\Omega\tau} \sigma^+ + e^{-i\Omega\tau} \sigma^- \right)$$



Entanglement



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$$\rho = \begin{pmatrix} 1 - 2A - C & 0 & 0 & X \\ 0 & A & B & 0 \\ 0 & B^* & A & 0 \\ X^* & 0 & 0 & C \end{pmatrix}$$

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iff *negativity* > 0

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- $A = \textit{probability}$ for single detector to get excited
- Use perturbation theory to calculate these

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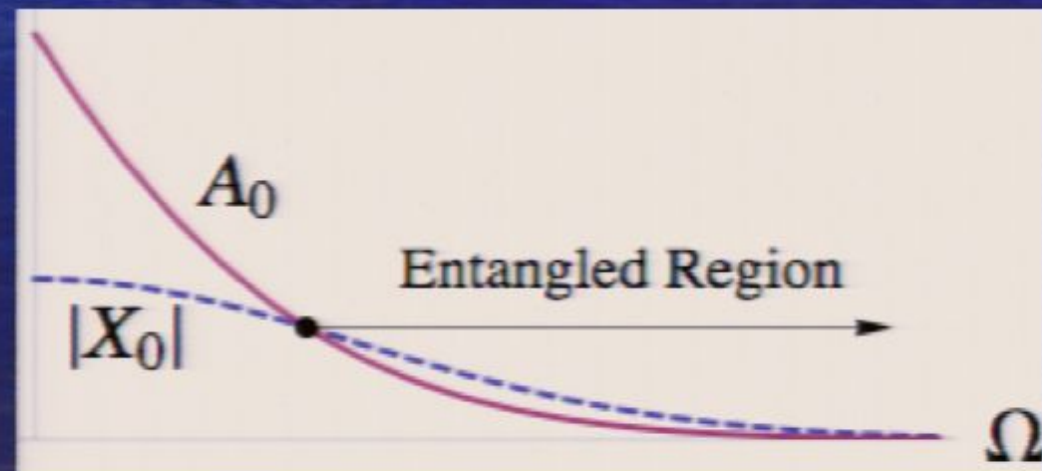
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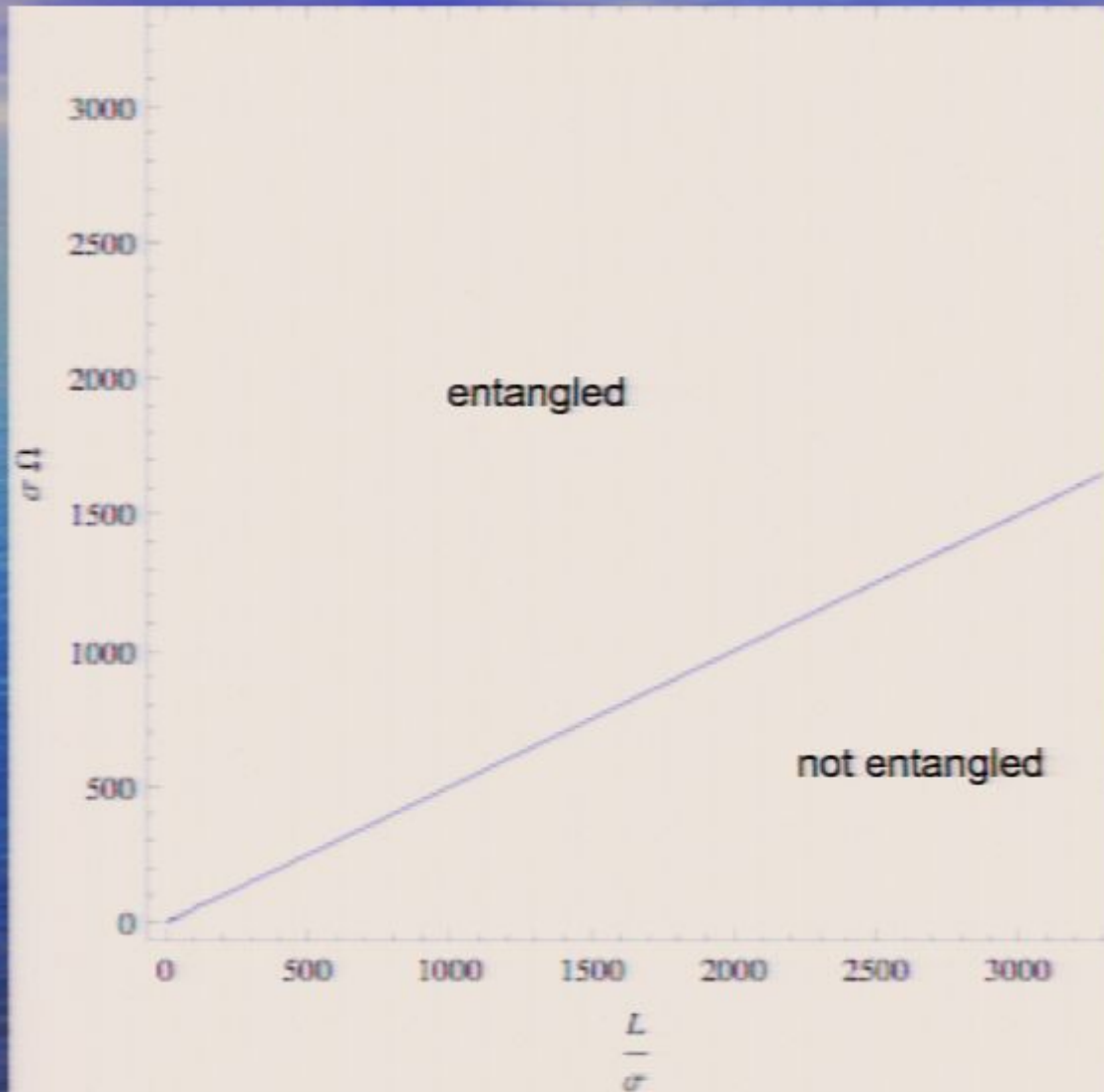
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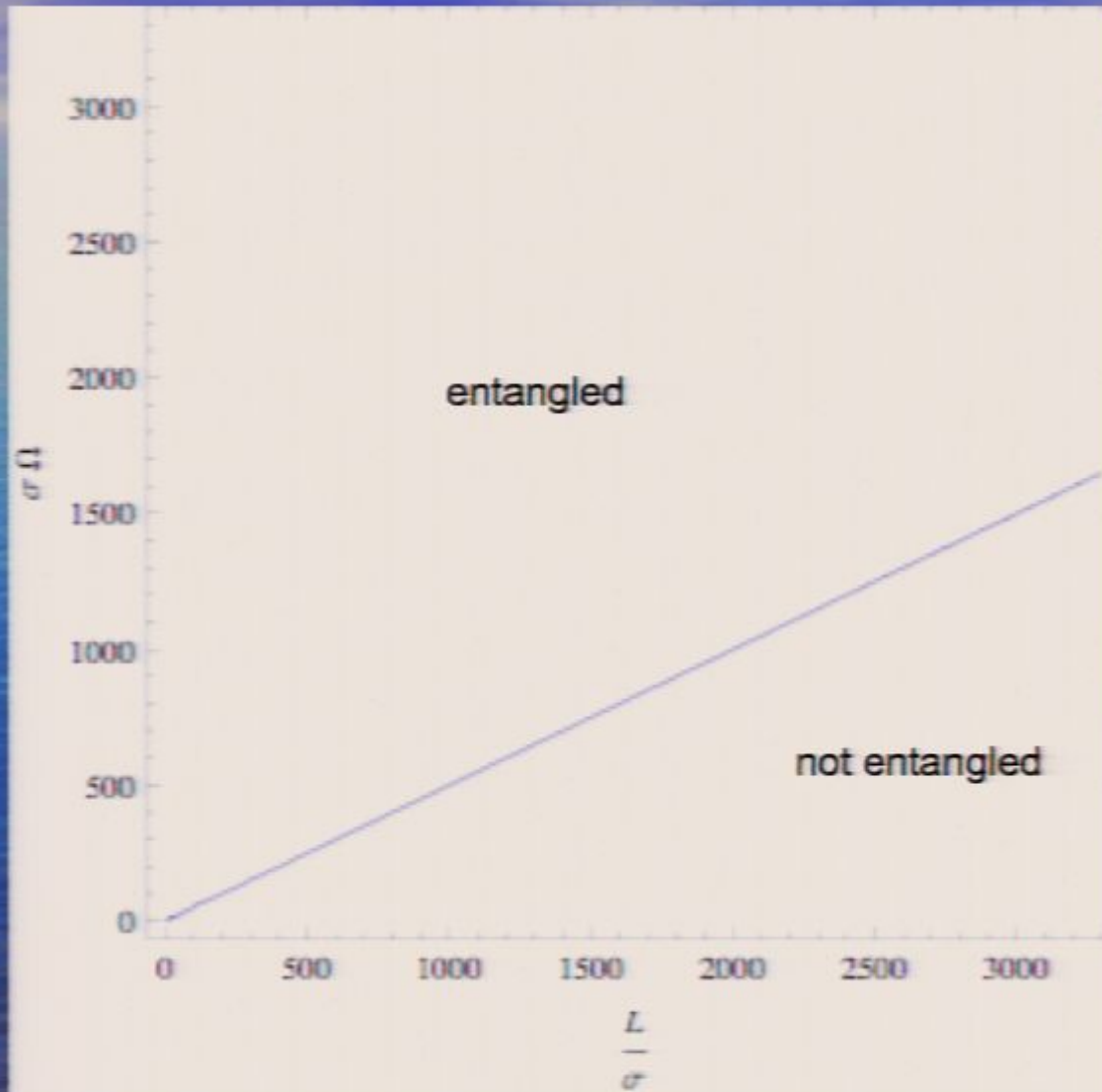
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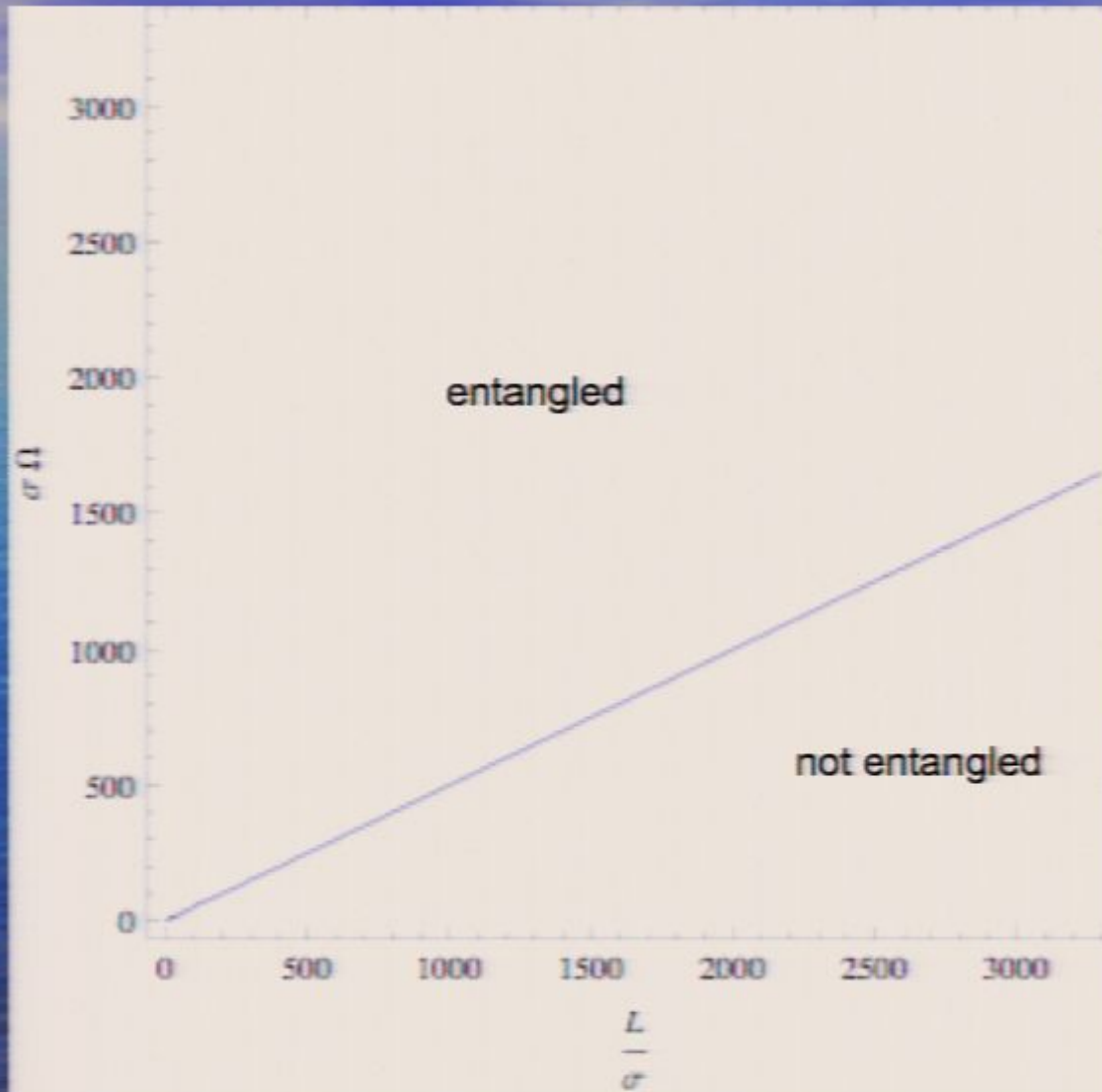


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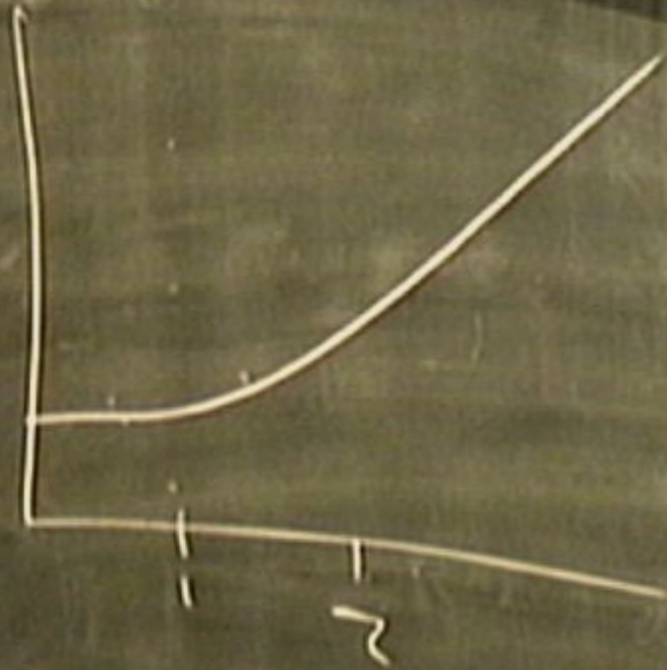


Flat thermal state (Minkowski, $T>0$)

Flat vacuum result



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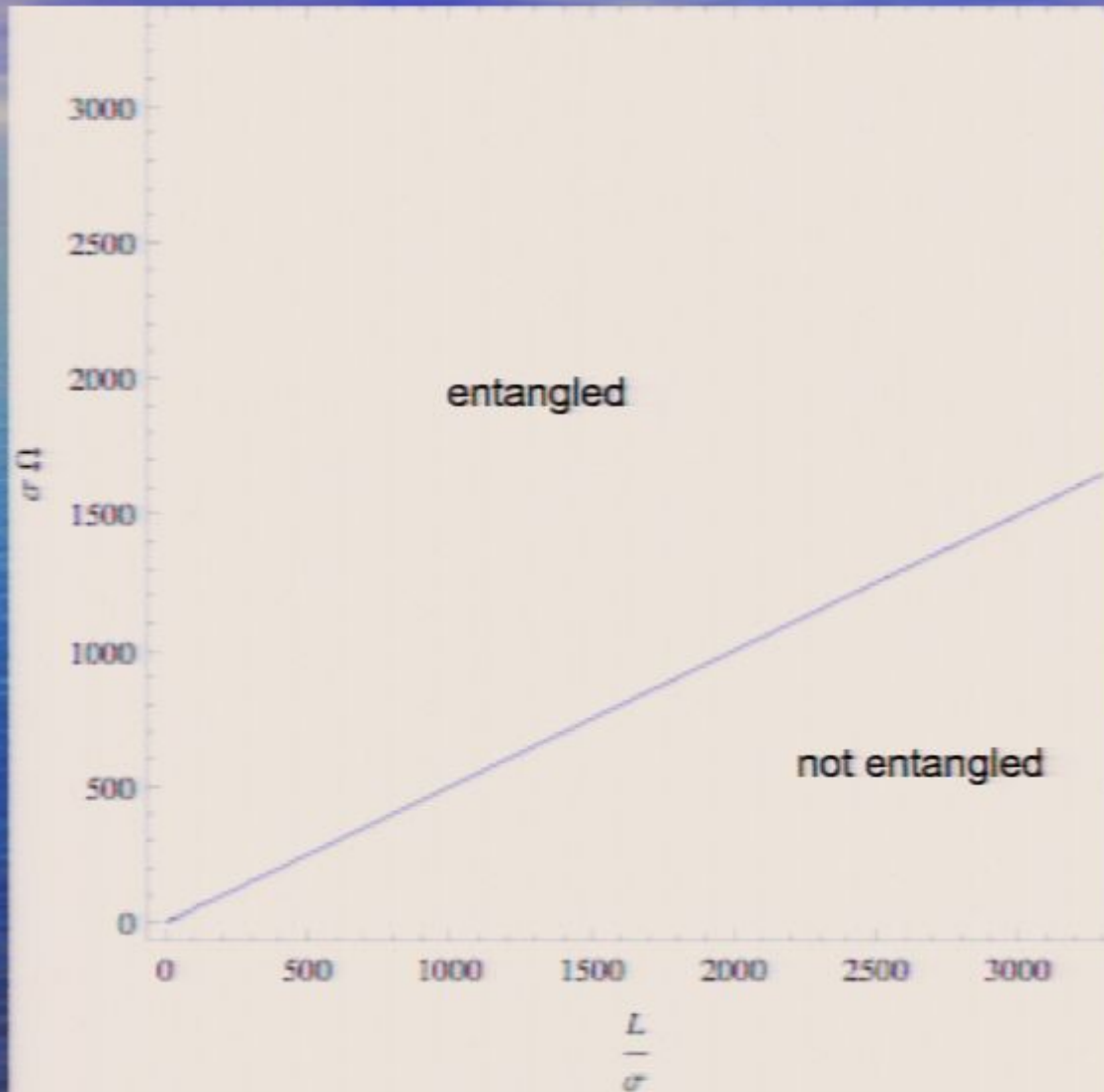


$$|0\rangle_i \otimes |0\rangle_{ii}$$

$$\int dx dt \phi(x)$$

$$\int$$

Flat vacuum result



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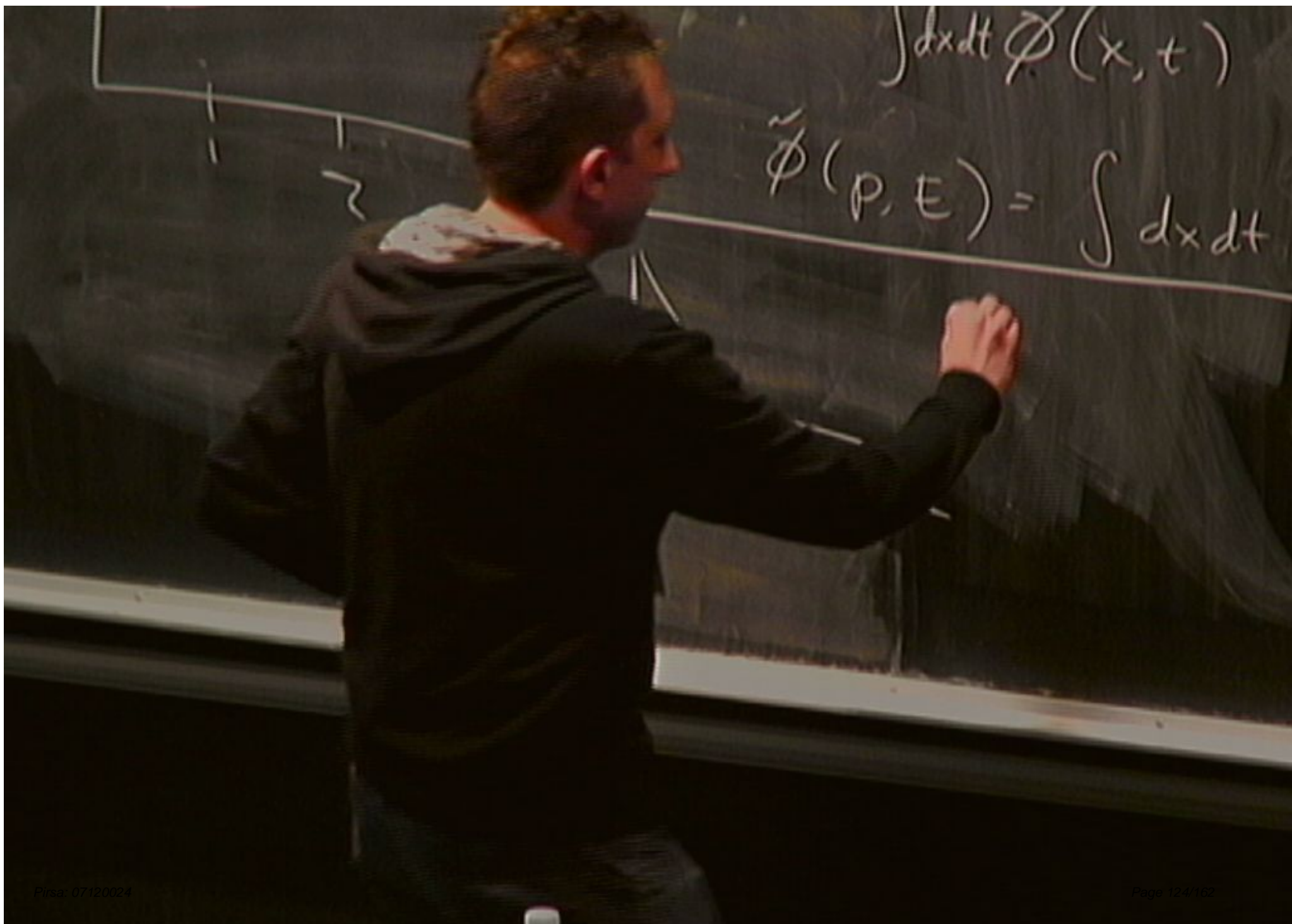
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 - If $T > 0$, then $\min_{\Omega} f(\Omega) > 0$

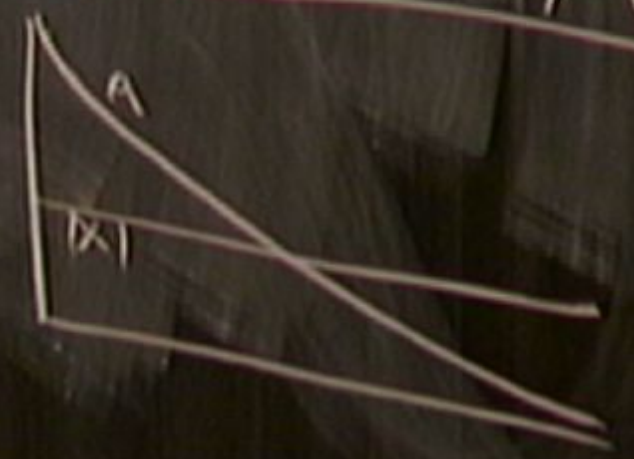
Flat thermal state (Minkowski, $T > 0$)

- $A = \exp(-\sigma^2 \Omega^2) f(\Omega)$
- $X = \exp(-\sigma^2 \Omega^2) g(L)$
- At fixed L ...
 - Entanglement occurs iff $f(\Omega) < (\text{const.})$
 - If $T > 0$, then $\min_{\Omega} f(\Omega) > 0$
 - Since $g(L)$ decreases monotonically with L , entanglement is restricted to $L < L_{\text{max}}$



$$\int dx dt \phi(x, t)$$

$$\tilde{\phi}(p, E) = \int dx dt \phi(x, t)$$



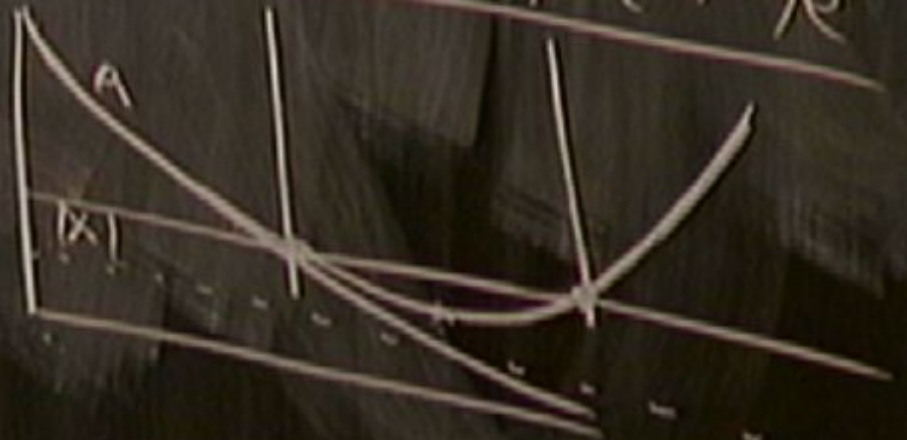
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$$\tilde{\phi}(p, E) = \int dx dt \phi(x, t) e^{-i p x - i E t}$$

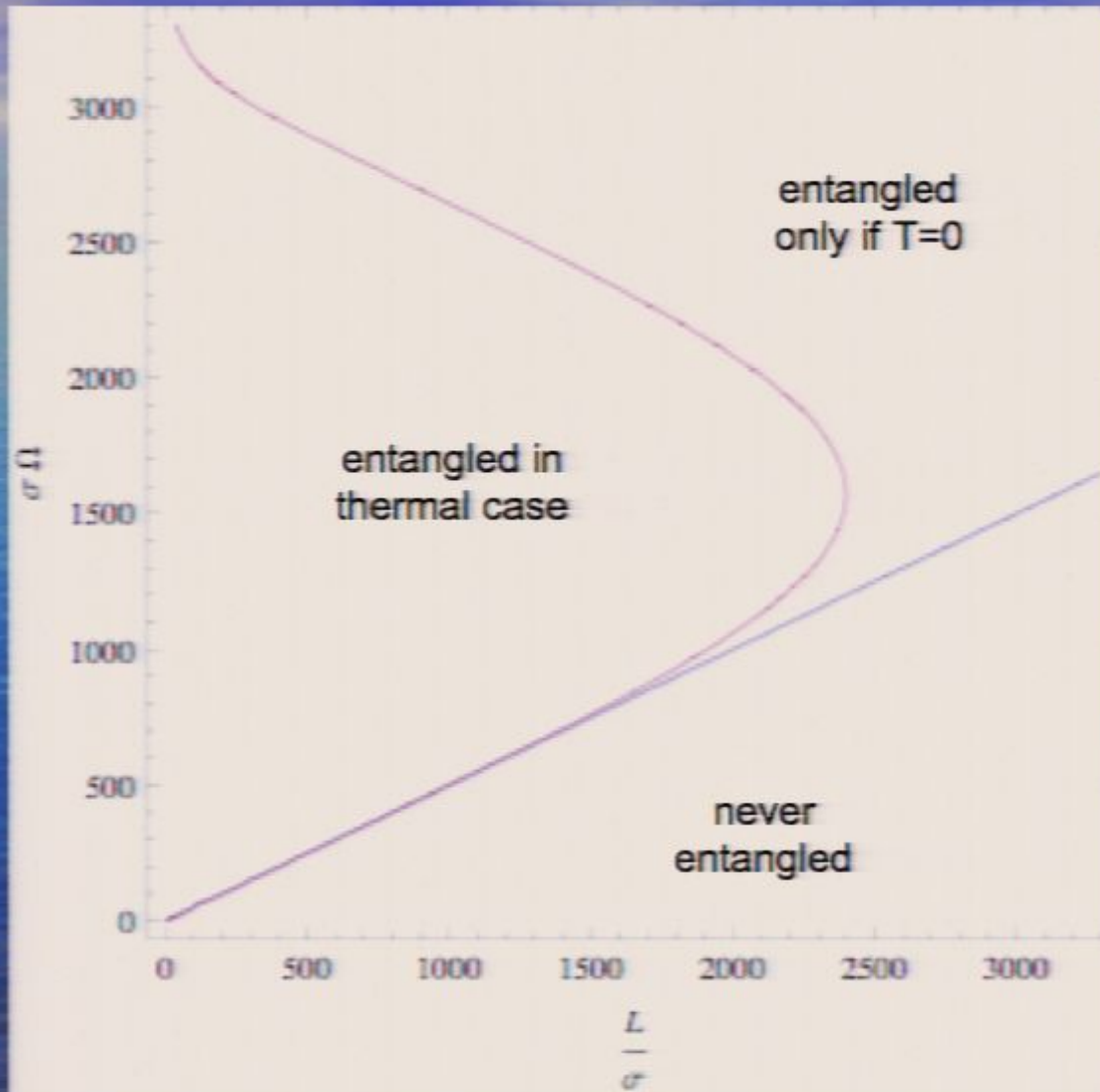


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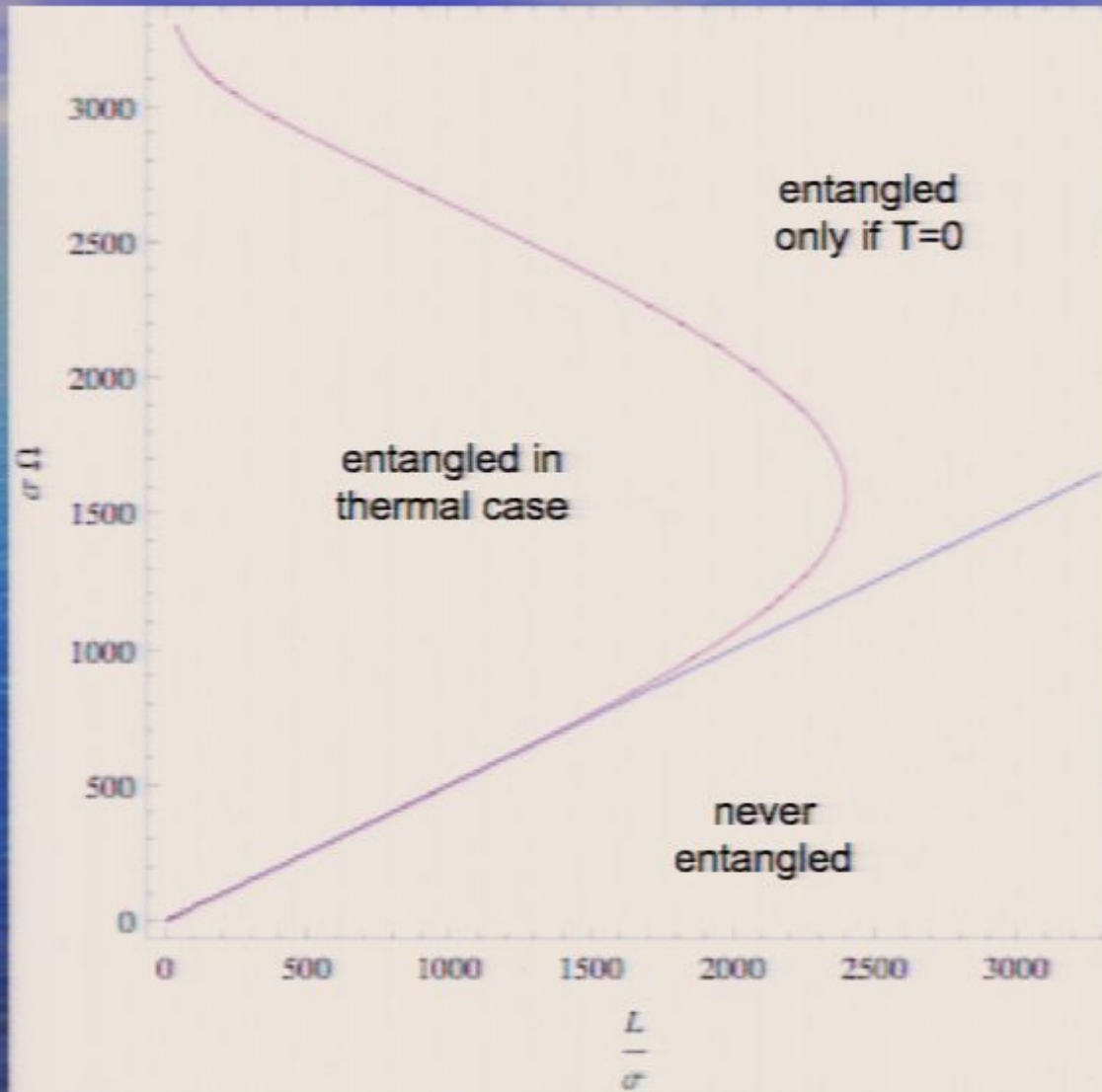


Flat spacetime



$$(2\pi\sigma T = .001)$$

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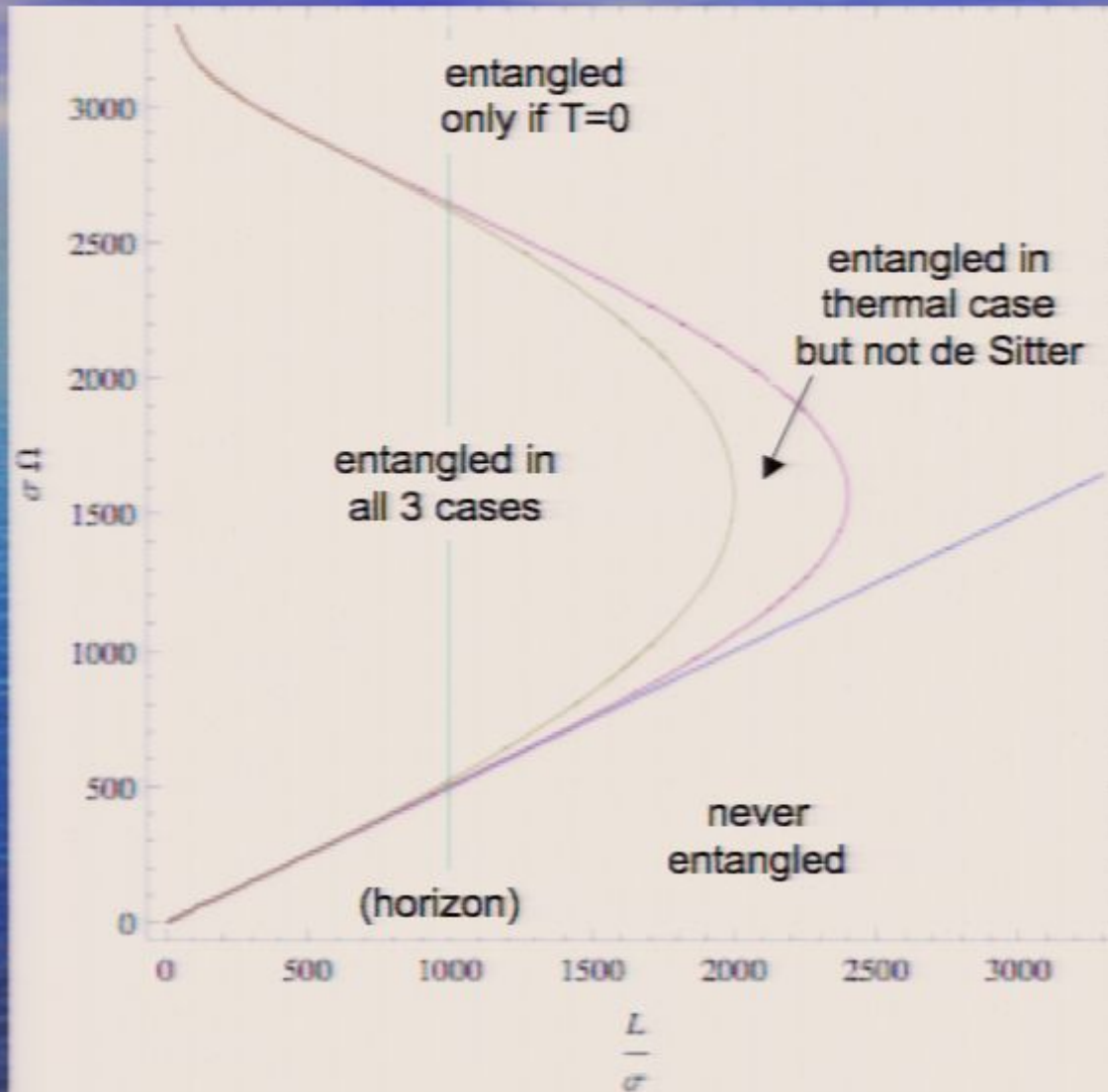
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- Can be distinguished by *entanglement*
- de Sitter vacuum: *less* entangling power than Minkowski thermal state ($T = T_{\text{GH}}$)
- Both have less power than the Minkowski vacuum (which has the least local noise)

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 - Other types of curvature

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- People
 - John Preskill
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 - Gerard Milburn
 - Carl Caves
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 - NSF
 - DoD

Thank you!



