

Title: Warped extra dimensions and partial compositeness

Date: Dec 03, 2007 04:00 PM

URL: <http://pirsa.org/07120020>

Abstract: Strong gauge dynamics can be given a holographic description in terms of a warped extra dimension. In particular, Randall-Sundrum models with bulk fields are dual to Standard Model partial compositeness. We identify a holographic basis of 4D fields that allows for a quantitative description of the elementary/composite mixing in these theories.



Warped Extra Dimensions and Partial Compositeness *

Brian Batell

University of Minnesota

*with Tony Gherghetta

- arXiv:0706.0890
- arXiv:0710.1838

I Perimeter YRC 11/29/07

5D Warped Dimension = 4D Strong Dynamics

- AdS/CFT duality: Extra dimension is a calculational tool
- Randall-Sundrum models \iff Standard Model partial compositeness
- How to quantify elementary/composite mixing?
 - Understand structure and phenomenology of 4D dual theory
- Answer: The Holographic Basis:

$$\Phi(x, y) = \varphi^s(x)g^s(y) + \sum_{n=1}^{\infty} \varphi_{CFT}^n(x)g^n(y)$$



Outline

- Context: The hierarchy problem
- Randall-Sundrum models and geometrical hierachies
- AdS/CFT and Holography
- The Kaluza-Klein Basis
- The Holographic Basis
- Elementary/composite content of SM fields
- Application: Flavor Changing Neutral Current suppression



Big Question : How do particles get mass?

- We think there is a Higgs boson waiting for us at the LHC
- but ... Hierarchy Problem - Higgs mass sensitive to quantum effects

$$m_{Higgs}^2 = m_0^2 + c\Lambda^2$$

- Suppose no new physics until Planck scale, $\Lambda \sim 10^{19}$ GeV:

$$(10^2 \text{GeV})^2 = m_0^2 + (10^{19} \text{GeV})^2$$

- \implies bare mass m_0 finely tuned

Ugly?



A "radical" solution

Warped extra dimension

Randall, Sundrum '99

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$$

- Warped geometry \Rightarrow
Energy scales depend on location
- Planck/Weak scale hierarchy:

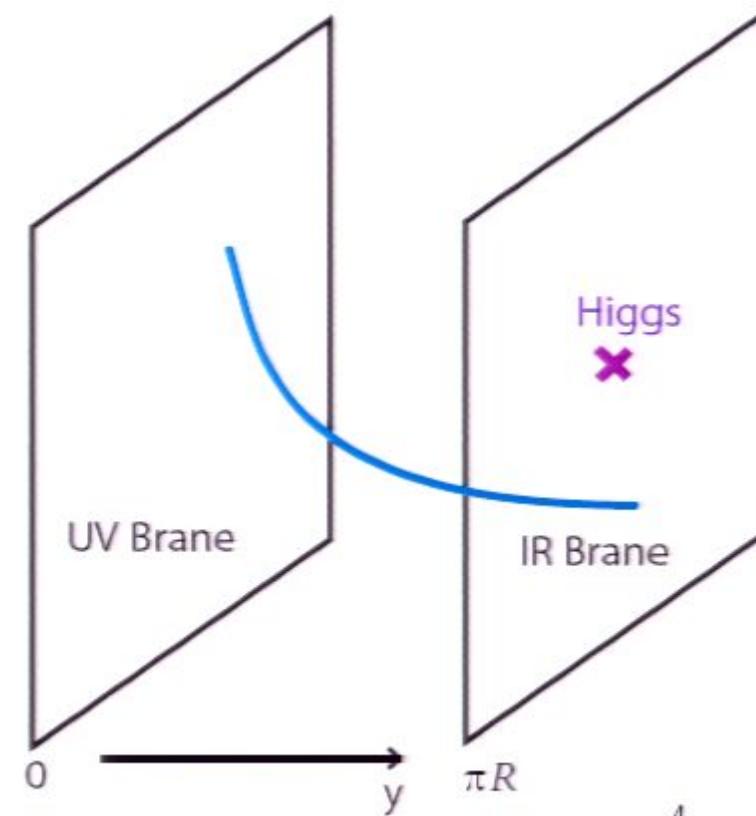
$$\Lambda_{weak} \sim M_P e^{-\pi k R}$$

$$k \sim \mathcal{O}(M_P), \quad \pi k R \sim \mathcal{O}(30)$$

- R can be naturally stabilized

Goldberger, Wise '99

Perimeter YRC 11/29/07



Standard Model in the bulk

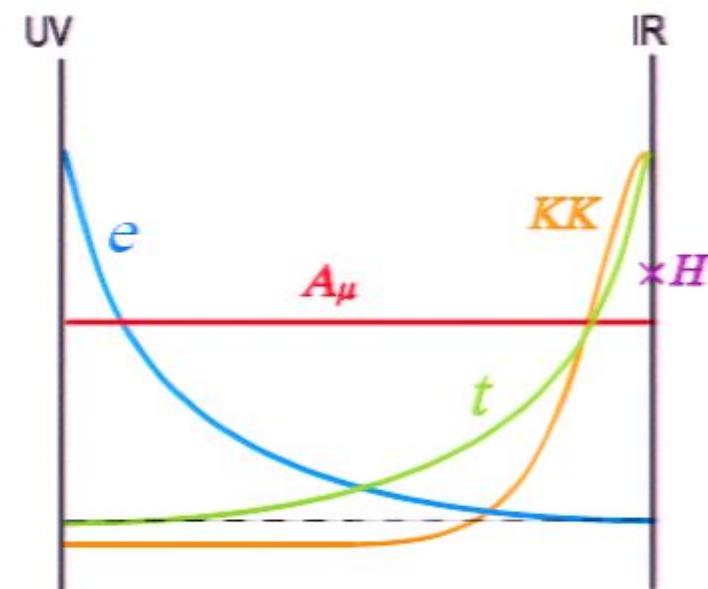
Davoudiasl, Hewett, Rizzo '99; Pomarol '99

Grossman, Neubert '99

Chang et al. '99; Gherghetta, Pomarol '00

Theory of Flavor!

- Natural Yukawa hierarchies
- FCNC suppressed



Bulk fields

Scalar field with tuned bulk and boundary masses

$$S = \int d^5x \sqrt{-g} \left[-\frac{1}{2}(\partial_M \Phi)^2 - \frac{1}{2}ak^2\Phi^2 - bk\Phi^2(\delta(y) - \delta(y - \pi R)) \right]$$

Tuning:

$$b = 2 \pm \sqrt{4 + a}$$

Why consider this toy model?

- Tuning allows for a localized zero mode: $\tilde{f}^0(y) \sim e^{(b-1)ky}$
 \Rightarrow Holographic interpretation depends on b
- special values for b mimic bulk graviton and gauge boson





1



6



69.5%



Find

Standard Model in the bulk

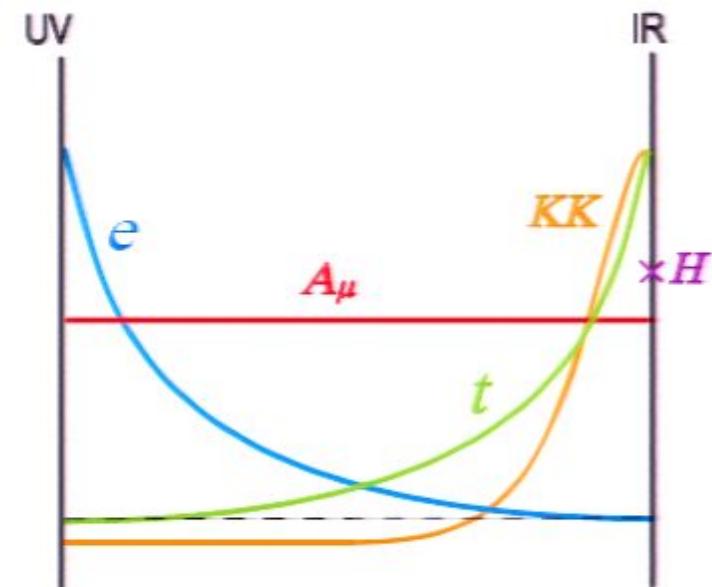
Davoudiasl, Hewett, Rizzo '99; Pomarol '99

Grossman, Neubert '99

Chang et al. '99; Gherghetta, Pomarol '00

Theory of Flavor!

- Natural Yukawa hierarchies
- FCNC suppressed





Bulk fields

Scalar field with tuned bulk and boundary masses

$$S = \int d^5x \sqrt{-g} \left[-\frac{1}{2}(\partial_M \Phi)^2 - \frac{1}{2}ak^2\Phi^2 - bk\Phi^2(\delta(y) - \delta(y - \pi R)) \right]$$

Tuning:

$$b = 2 \pm \sqrt{4 + a}$$

Why consider this toy model?

- Tuning allows for a localized zero mode: $\tilde{f}^0(y) \sim e^{(b-1)ky}$
 \Rightarrow Holographic interpretation depends on b
- special values for b mimic bulk graviton and gauge boson



Standard Model in the bulk

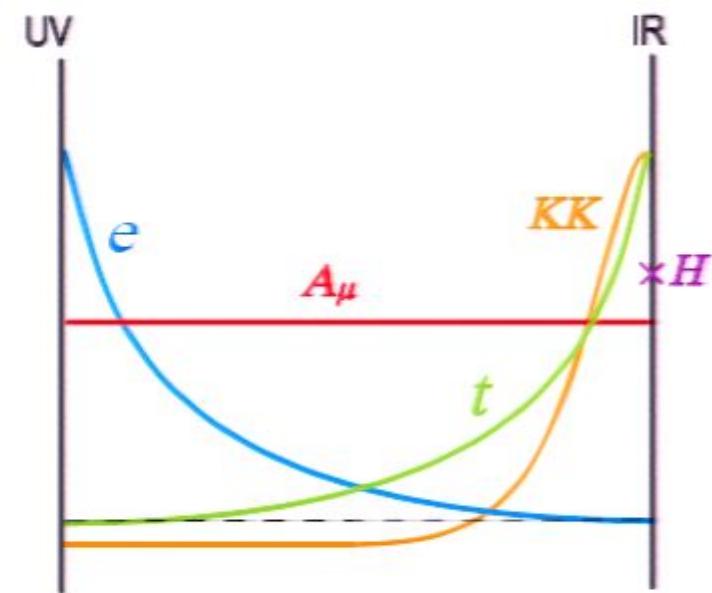
Davoudiasl, Hewett, Rizzo '99; Pomarol '99

Grossman, Neubert '99

Chang et al. '99; Gherghetta, Pomarol '00

Theory of Flavor!

- Natural Yukawa hierarchies
- FCNC suppressed



Bulk fields

Scalar field with tuned bulk and boundary masses

$$S = \int d^5x \sqrt{-g} \left[-\frac{1}{2}(\partial_M \Phi)^2 - \frac{1}{2}ak^2\Phi^2 - bk\Phi^2(\delta(y) - \delta(y - \pi R)) \right]$$

Tuning:

$$b = 2 \pm \sqrt{4 + a}$$

Why consider this toy model?

- Tuning allows for a localized zero mode: $\tilde{f}^0(y) \sim e^{(b-1)ky}$
 \Rightarrow Holographic interpretation depends on b
- special values for b mimic bulk graviton and gauge boson



AdS/CFT duality

Maldacena '97

is for our purposes . . .

Weakly coupled gravity in warped 5D	dual	Strongly coupled gauge theory (CFT) in 4D*
--	------	---



* Large N_c gauge theory

$$\left\langle \exp \left(- \int \varphi_0 \mathcal{O} \right) \right\rangle_{\text{CFT}} = \exp [- \Gamma(\varphi_0)]$$

Gubser, Klebanov, Polyakov '97; Witten '97



“Dictionary”

<u>5D</u>	<u>4D</u>
bulk field $\Phi(x, y)$	\iff CFT operator $\mathcal{O}(x)$
BC $\Phi(x, y_0) = \varphi_0(x)$	\iff source: $\varphi_0(x)\mathcal{O}(x)$
bulk mass	\iff dimension Δ of \mathcal{O}

e.g.

Bulk gauge field

global symmetry current

$$A_\mu(x, y) \iff J_\mu^{CFT}$$

$$m_A^2 = 0 \iff \Delta_J = 3$$



AdS/CFT duality

Maldacena '97

is for our purposes . . .

Weakly coupled gravity in warped 5D	dual	Strongly coupled gauge theory (CFT) in 4D*
	\iff	

* Large N_c gauge theory

$$\left\langle \exp \left(- \int \varphi_0 \mathcal{O} \right) \right\rangle_{\text{CFT}} = \exp [- \Gamma(\varphi_0)]$$

Gubser, Klebanov, Polyakov '97; Witten '97



“Dictionary”

<u>5D</u>	<u>4D</u>
bulk field $\Phi(x, y)$	\iff CFT operator $\mathcal{O}(x)$
BC $\Phi(x, y_0) = \varphi_0(x)$	\iff source: $\varphi_0(x)\mathcal{O}(x)$
bulk mass	\iff dimension Δ of \mathcal{O}

e.g.

Bulk gauge field

global symmetry current

$$A_\mu(x, y) \iff J_\mu^{CFT}$$

$$m_A^2 = 0 \iff \Delta_J = 3$$

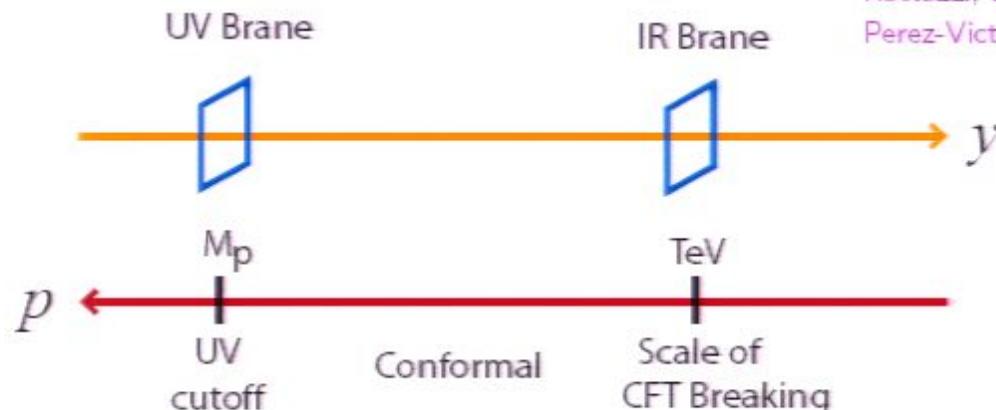


Holography for RS1

Arkani-Hamed, Poratti, Randall '00

Rattazzi, Zaffaroni '00

Perez-Victoria '00



zero mode \sim source field (elementary)

KK modes \sim CFT bound states (composites)

but wait ...

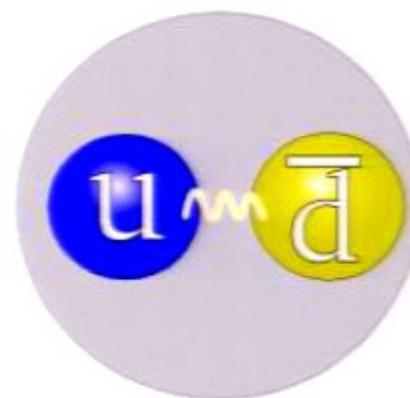
Mixing through operator $\varphi_0(x)\mathcal{O}(x)$

⇒ Mass eigenstates are elementary/composite mixtures

A “conservative” solution

Compositeness (Strong gauge dynamics)

- Quantum Chromodynamics:
 - $SU(3)_{color}$; quarks and gluons
 - Confinement at $\Lambda_{QCD} \sim 200\text{MeV}$
 - We observe composite particles,
e.g. Pion



Could the Higgs (or something like it) be composite, with $\Lambda \sim \text{TeV}$

- Technicolor Weinberg; Susskind '70s
- Composite Higgs Kaplan, Georgi '80s



The “Holographic Recipe”

Step 1: Evaluate bulk action for arbitrary boundary condition

$$\Phi(x, y_0) = \varphi_0(x) \text{ to obtain } \Gamma(\varphi_0)$$

Step 2: Take functional derivatives to compute correlation functions of CFT operators

$$\begin{aligned}\langle \mathcal{O} \mathcal{O} \rangle(p) &= \frac{\delta^2}{\delta \varphi_0^2} \left\langle \exp \left(- \int \varphi_0 \mathcal{O} \right) \right\rangle_{\text{CFT}} = \frac{\delta^2}{\delta \varphi_0^2} \exp[-\Gamma(\varphi_0)] \\ &= \mp ip \frac{J_{b-1} \left(\frac{ip}{k} \right) Y_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right) - Y_{b-1} \left(\frac{ip}{k} \right) J_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right)}{J_{b-2} \left(\frac{ip}{k} \right) Y_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right) - Y_{b-2} \left(\frac{ip}{k} \right) J_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right)}\end{aligned}$$

Step 3: Interpret $\langle \mathcal{O} \mathcal{O} \rangle(p)$



Operator Dimension

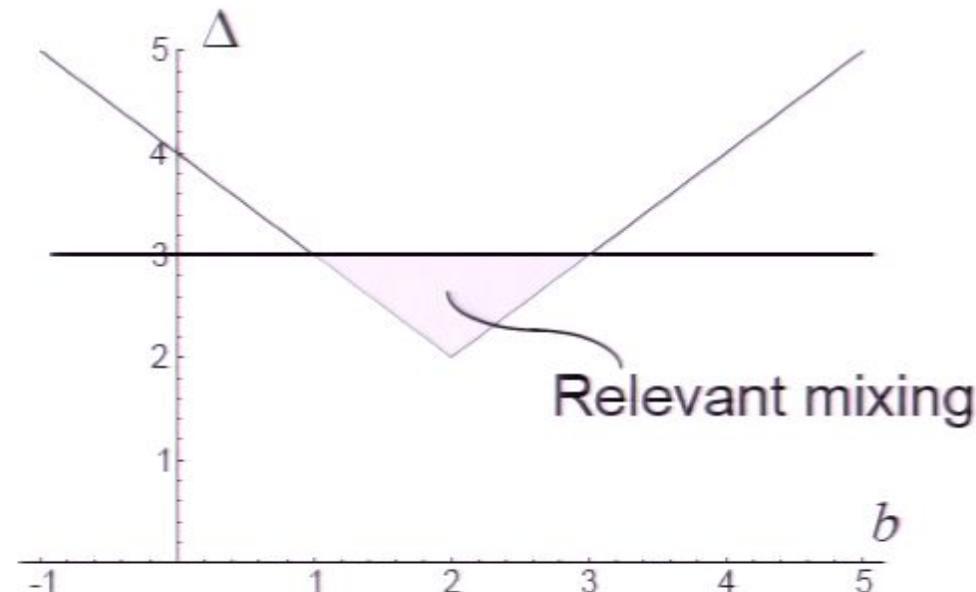
$$\Delta = 2 + \sqrt{4 + a} = 2 + |b - 2|$$

$$\varphi_0(x) \quad \widehat{\mathcal{O}(x)}$$

$1 \qquad \qquad 2 + |b - 2|$

$b < 1$ or $b > 3$
⇒ irrelevant mixing

$1 < b < 3$
⇒ relevant mixing



The “Holographic Recipe”

Step 1: Evaluate bulk action for arbitrary boundary condition

$$\Phi(x, y_0) = \varphi_0(x) \text{ to obtain } \Gamma(\varphi_0)$$

Step 2: Take functional derivatives to compute correlation functions of CFT operators

$$\begin{aligned}\langle \mathcal{O} \mathcal{O} \rangle(p) &= \frac{\delta^2}{\delta \varphi_0^2} \left\langle \exp \left(- \int \varphi_0 \mathcal{O} \right) \right\rangle_{\text{CFT}} = \frac{\delta^2}{\delta \varphi_0^2} \exp[-\Gamma(\varphi_0)] \\ &= \mp ip \frac{J_{b-1} \left(\frac{ip}{k} \right) Y_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right) - Y_{b-1} \left(\frac{ip}{k} \right) J_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right)}{J_{b-2} \left(\frac{ip}{k} \right) Y_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right) - Y_{b-2} \left(\frac{ip}{k} \right) J_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right)}\end{aligned}$$

Step 3: Interpret $\langle \mathcal{O} \mathcal{O} \rangle(p)$



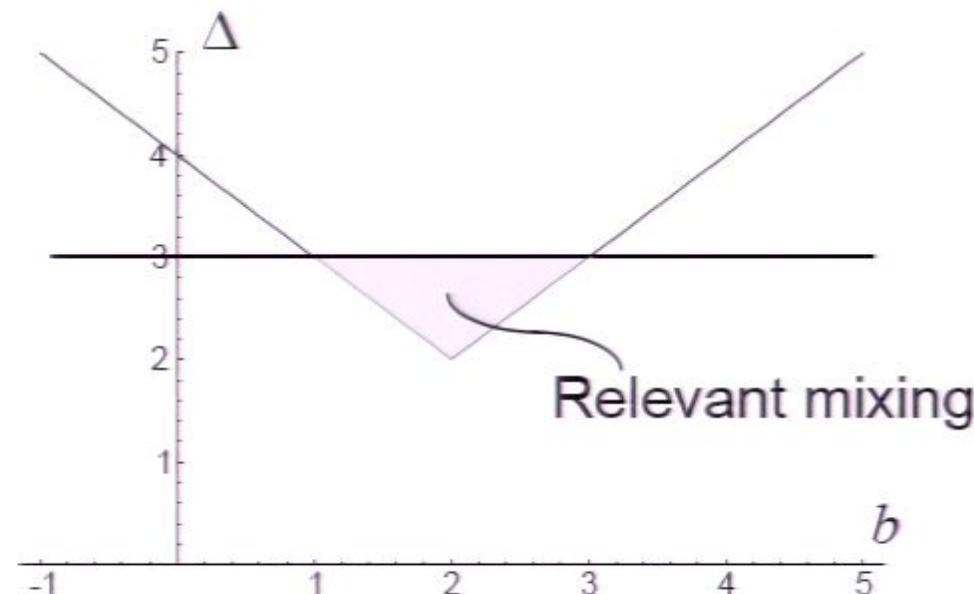
Operator Dimension

$$\Delta = 2 + \sqrt{4 + a} = 2 + |b - 2|$$

$$\varphi_0(x) \quad \overset{1}{\overbrace{\quad}} \quad 2 + |b - 2|$$
$$\mathcal{O}(x) \quad \overset{2 + |b - 2|}{\overbrace{\quad}}$$

$b < 1$ or $b > 3$
⇒ irrelevant mixing

$1 < b < 3$
⇒ relevant mixing





Two branches in dual theory

$$\Delta = 2 + |b - 2|$$

- $b < 2$:
 - source field $\varphi_0(x)$ massless
 - zero mode primarily elementary
 - Nearly all RS phenomenological examples are described by $b < 2$ (fermions too!)

- $b > 2$:
 - source field $\varphi_0(x)$ massive $M_0 \sim k$
 - zero mode primarily composite
 - Higgs; perhaps t_R in some models



Operator Dimension

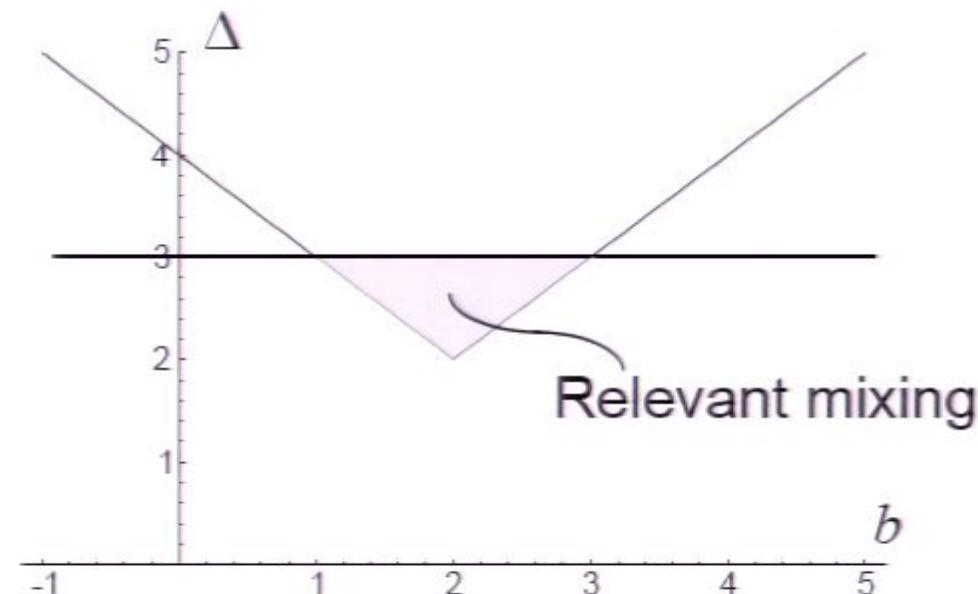
$$\Delta = 2 + \sqrt{4 + a} = 2 + |b - 2|$$

$$\varphi_0(x) \quad \widehat{\mathcal{O}(x)}$$

$1 \qquad \qquad 2 + |b - 2|$

$b < 1$ or $b > 3$
⇒ irrelevant mixing

$1 < b < 3$
⇒ relevant mixing



Two branches in dual theory

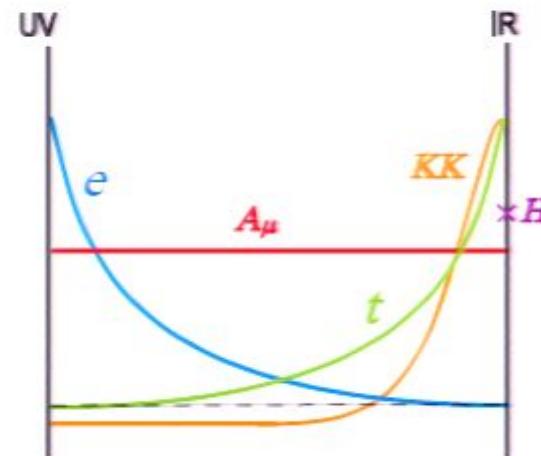
$$\Delta = 2 + |b - 2|$$

- $b < 2$:
 - source field $\varphi_0(x)$ massless
 - zero mode primarily elementary
 - Nearly all RS phenomenological examples are described by $b < 2$ (fermions too!)

- $b > 2$:
 - source field $\varphi_0(x)$ massive $M_0 \sim k$
 - zero mode primarily composite
 - Higgs; perhaps t_R in some models



Partial compositeness of SM fields



- UV localized \iff mostly elementary
- IR localized \iff mostly composite

Can we quantify source/CFT (elementary/composite) mixing?

Kaluza-Klein mass eigenbasis

KK decomposition:

$$\Phi(x, y) = \sum_{n=0}^{\infty} \phi^n(x) f^n(y), \quad \text{BC : } (++)$$
$$(\partial_5 - bk) f^n(y) \Big|_{0, \pi R} = 0$$

Localized massless mode:

$$\tilde{f}^0(y) \sim e^{(b-1)ky}, \quad -\infty < b < \infty$$

The fields $\phi^n(x)$ are the mass eigenstates

- Spectrum:

$$J_{b-1} \left(\frac{m_n}{k} \right) Y_{b-1} \left(\frac{m_n e^{\pi k R}}{k} \right) - Y_{b-1} \left(\frac{m_n}{k} \right) J_{b-1} \left(\frac{m_n e^{\pi k R}}{k} \right) = 0$$



Two branches in dual theory

$$\Delta = 2 + |b - 2|$$

- $b < 2$:
 - source field $\varphi_0(x)$ massless
 - zero mode primarily elementary
 - Nearly all RS phenomenological examples are described by $b < 2$ (fermions too!)

- $b > 2$:
 - source field $\varphi_0(x)$ massive $M_0 \sim k$
 - zero mode primarily composite
 - Higgs; perhaps t_R in some models



The “Holographic Recipe”

Step 1: Evaluate bulk action for arbitrary boundary condition

$$\Phi(x, y_0) = \varphi_0(x) \text{ to obtain } \Gamma(\varphi_0)$$

Step 2: Take functional derivatives to compute correlation functions of CFT operators

$$\begin{aligned}\langle \mathcal{O} \mathcal{O} \rangle(p) &= \frac{\delta^2}{\delta \varphi_0^2} \left\langle \exp \left(- \int \varphi_0 \mathcal{O} \right) \right\rangle_{\text{CFT}} = \frac{\delta^2}{\delta \varphi_0^2} \exp[-\Gamma(\varphi_0)] \\ &= \mp ip \frac{J_{b-1} \left(\frac{ip}{k} \right) Y_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right) - Y_{b-1} \left(\frac{ip}{k} \right) J_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right)}{J_{b-2} \left(\frac{ip}{k} \right) Y_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right) - Y_{b-2} \left(\frac{ip}{k} \right) J_{b-1} \left(\frac{ipe^{\pi k R}}{k} \right)}\end{aligned}$$

Step 3: Interpret $\langle \mathcal{O} \mathcal{O} \rangle(p)$



Kaluza-Klein mass eigenbasis

KK decomposition:

$$\Phi(x, y) = \sum_{n=0}^{\infty} \phi^n(x) f^n(y), \quad \text{BC : } (++)$$
$$(\partial_5 - bk) f^n(y) \Big|_{0, \pi R} = 0$$

Localized massless mode:

$$\tilde{f}^0(y) \sim e^{(b-1)ky}, \quad -\infty < b < \infty$$

The fields $\phi^n(x)$ are the mass eigenstates

- Spectrum:

$$J_{b-1} \left(\frac{m_n}{k} \right) Y_{b-1} \left(\frac{m_n e^{\pi k R}}{k} \right) - Y_{b-1} \left(\frac{m_n}{k} \right) J_{b-1} \left(\frac{m_n e^{\pi k R}}{k} \right) = 0$$



Holographic basis

Basic idea:

Expand the bulk field directly in terms of a source field $\varphi^s(x)$ and composite CFT states $\varphi_{CFT}^n(x)$:

$$\Phi(x, y) = \varphi^s(x)g^s(y) + \sum_{n=1}^{\infty} \varphi_{CFT}^n(x)g^n(y)$$

- Leads to kinetic and mass mixing in 4D effective theory
- Mass eigenstates will be a mixture of $\varphi^s(x)$ and $\varphi_{CFT}^n(x)$



Source profile $g^s(y)$

$g^s(y)$ can be determined from mass of source

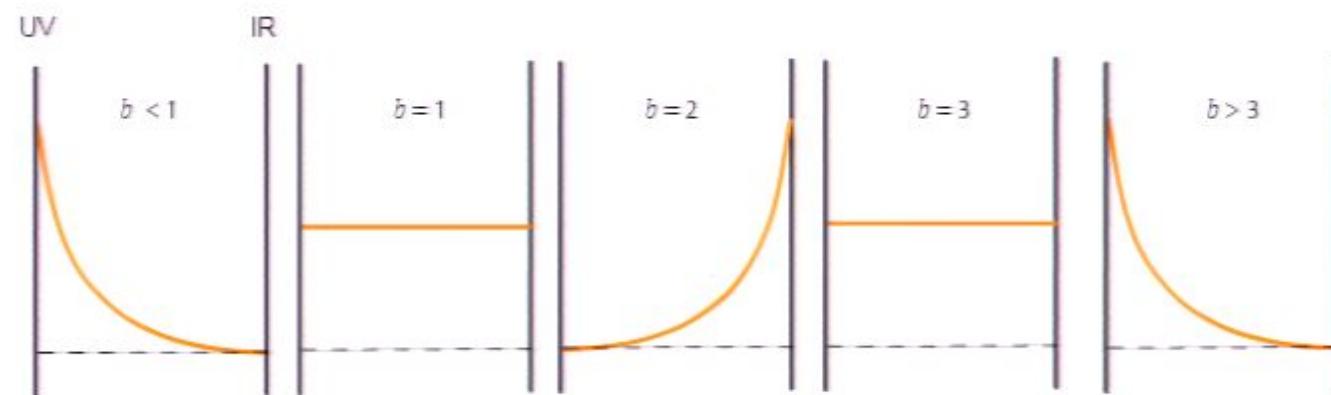
$$M_s^2 = \begin{cases} 0 & \text{for } b < 2 \\ 4(b-2)(b-3)k^2 & \text{for } b > 2 \end{cases}$$

$$\implies \tilde{g}^s(y) \sim e^{-ky} e^{(4-\Delta)ky} = \begin{cases} e^{(b-1)ky} & \text{for } b < 2 \\ e^{(3-b)ky} & \text{for } b > 2 \end{cases}$$



Source profiles mimic operator dimensions:

$$\Delta = 2 + |2 - b|$$



- Indicates when mixing is relevant, marginal, or irrelevant



Source profile $g^s(y)$

$g^s(y)$ can be determined from mass of source

$$M_s^2 = \begin{cases} 0 & \text{for } b < 2 \\ 4(b-2)(b-3)k^2 & \text{for } b > 2 \end{cases}$$

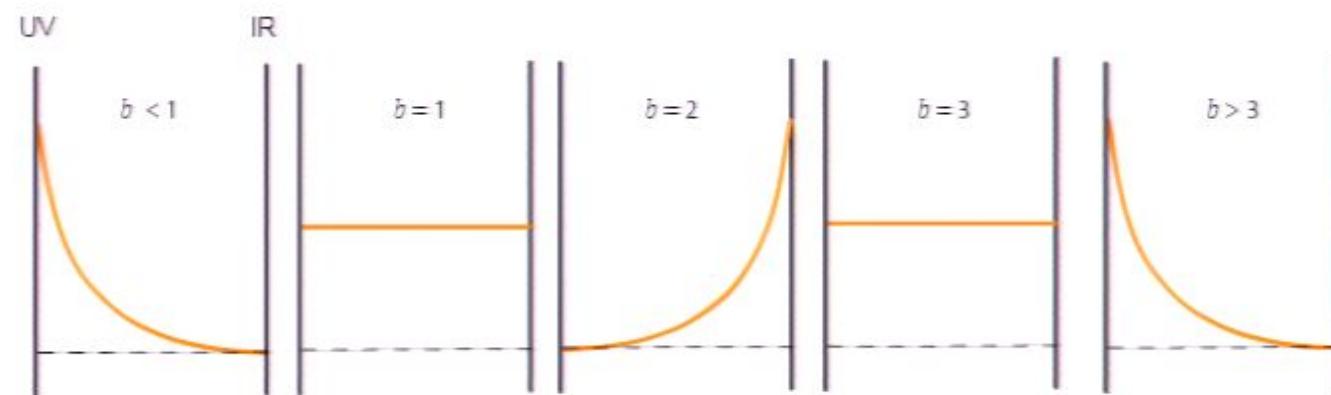
$$\implies \tilde{g}^s(y) \sim e^{-ky} e^{(4-\Delta)ky} = \begin{cases} e^{(b-1)ky} & \text{for } b < 2 \\ e^{(3-b)ky} & \text{for } b > 2 \end{cases}$$





Source profiles mimic operator dimensions:

$$\Delta = 2 + |2 - b|$$



- Indicates when mixing is relevant, marginal, or irrelevant



CFT composite profiles $g^n(y)$

CFT spectrum obtained from poles in 2-point function:

$$J_{b-2} \left(\frac{M_n}{k} \right) Y_{b-1} \left(\frac{M_n e^{\pi k R}}{k} \right) - Y_{b-2} \left(\frac{M_n}{k} \right) J_{b-1} \left(\frac{M_n e^{\pi k R}}{k} \right) = 0$$

Note different from KK spectrum!

Identical to the spectrum obtained with the following BC for $g^n(y)$:

BC : $(-+)$

$$g^n(y) \Big|_0 = 0$$

$$(\partial_5 - bk) g^n(y) \Big|_{\pi R} = 0$$



Effective 4D Lagrangian in the holographic basis

$$\mathcal{L} = \frac{1}{2} \vec{\varphi}^T \mathbf{Z} \square \vec{\varphi} - \frac{1}{2} \vec{\varphi}^T \mathbf{M}^2 \vec{\varphi},$$

where $\vec{\varphi}^T = (\varphi^s, \varphi_{CFT}^1, \varphi_{CFT}^2, \dots)$

$$\mathbf{Z} = \begin{pmatrix} 1 & z_1 & z_2 & z_3 & \cdots \\ z_1 & 1 & 0 & 0 & \cdots \\ z_2 & 0 & 1 & 0 & \cdots \\ z_3 & 0 & 0 & 1 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}, \quad \mathbf{M}^2 = \begin{pmatrix} M_s^2 & \mu_1^2 & \mu_2^2 & \mu_3^2 & \cdots \\ \mu_1^2 & M_1^2 & 0 & 0 & \cdots \\ \mu_2^2 & 0 & M_2^2 & 0 & \cdots \\ \mu_3^2 & 0 & 0 & M_3^2 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

Notice kinetic mixing \Rightarrow nonorthogonal basis

z_n and μ_n^2 computed from wavefunction overlap integrals

Diagonalization leads to KK basis





69.5%



Graviton $h_{\mu\nu}$

$$\tilde{f}^0(y) \sim e^{-ky}$$

$b = 0; \Delta = 4 \implies$ irrelevant mixing

$$\begin{pmatrix} h^0 \\ h^1 \\ \vdots \end{pmatrix} = \begin{pmatrix} 1 & \sim e^{-\pi kR} & \dots \\ 0 & \sim -1 & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} h^s \\ h^{1(CFT)} \\ \vdots \end{pmatrix}$$

- 4D graviton $h_{\mu\nu}^0(x) \sim$ elementary source ; compositeness negligible
- KK modes are purely composite





Gauge field A_μ

$$\tilde{f}^0(y) = \frac{1}{\sqrt{\pi R}}$$

$b = 1; \Delta = 3 \implies$ marginal mixing

$$\begin{pmatrix} A_\mu^0 \\ A_\mu^1 \\ A_\mu^2 \\ \vdots \end{pmatrix} = \begin{pmatrix} 1 & -0.19 & 0.13 & \dots \\ 0 & -0.98 & -0.03 & \dots \\ 0 & 0.01 & -0.99 & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} A_\mu^s \\ A_\mu^{1(CFT)} \\ A_\mu^{1(CFT)} \\ \vdots \end{pmatrix}$$

- massless eigenstate $A_\mu^0(x)$ is primarily elementary
- KK modes are purely composite



Right-handed top t_R

$$\tilde{f}^0(y) = e^{(\frac{1}{2}-c)ky} \quad m_\psi = ck$$

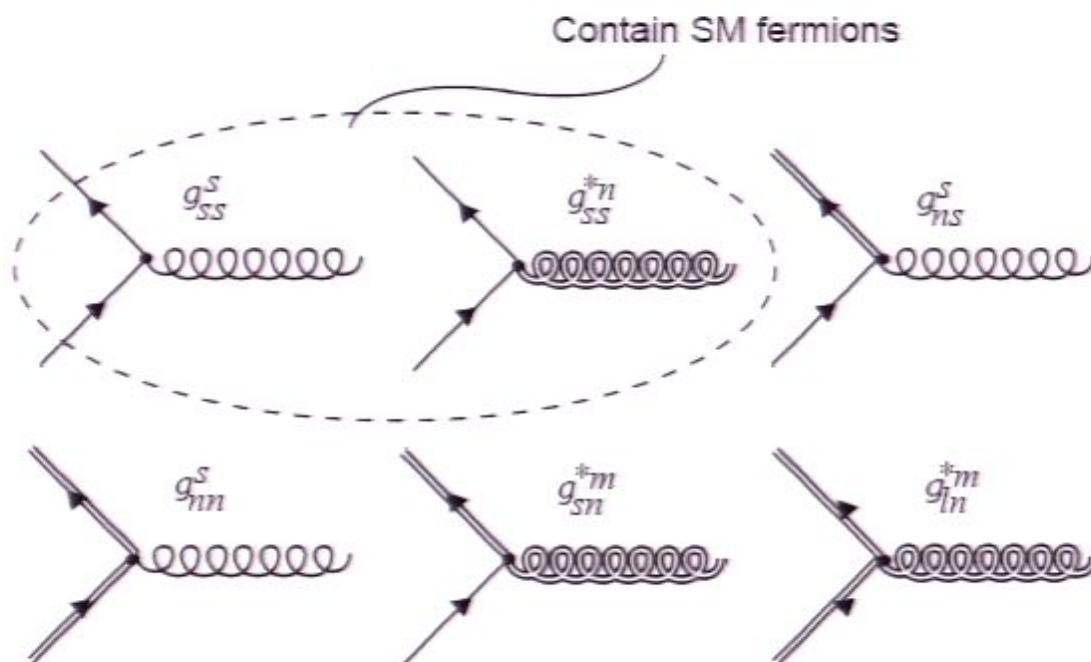
Take e.g. $c = -0.7$; $\Delta = 1.7 \Rightarrow$ relevant mixing

$$\begin{pmatrix} t_R^{(0)} \\ t_R^{(1)} \\ t_R^{(2)} \\ t_R^{(3)} \\ \vdots \end{pmatrix} = \begin{pmatrix} 0.9796 & \sim -1 & \sim 0 & \sim 0 & \dots \\ -0.1816 & \sim 0 & \sim -1 & \sim 0 & \dots \\ 0.0514 & \sim 0 & \sim 0 & \sim -1 & \dots \\ 0.0471 & \sim 0 & \sim 0 & \sim 0 & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} t_R^s \\ t_R^{CFT(1)} \\ t_R^{CFT(2)} \\ t_R^{CFT(3)} \\ \vdots \end{pmatrix}$$

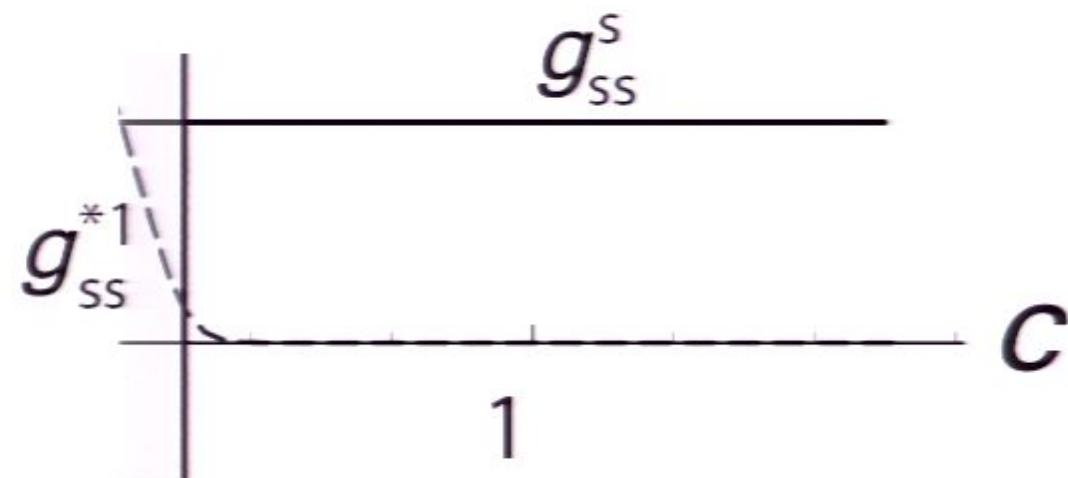
- massless eigenstate $t_R^0(x)$ roughly equal mixture of source/CFT
- KK modes contain elementary component



Gauge interactions $g_{\psi\psi}^A$

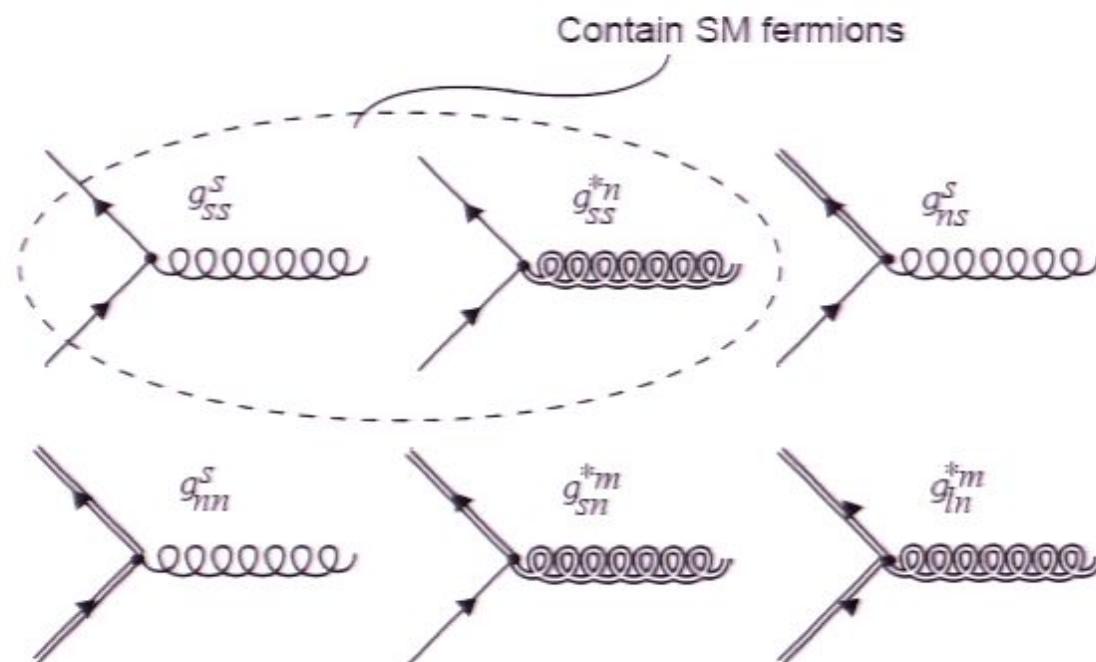


Key point: composites contain no zero mode

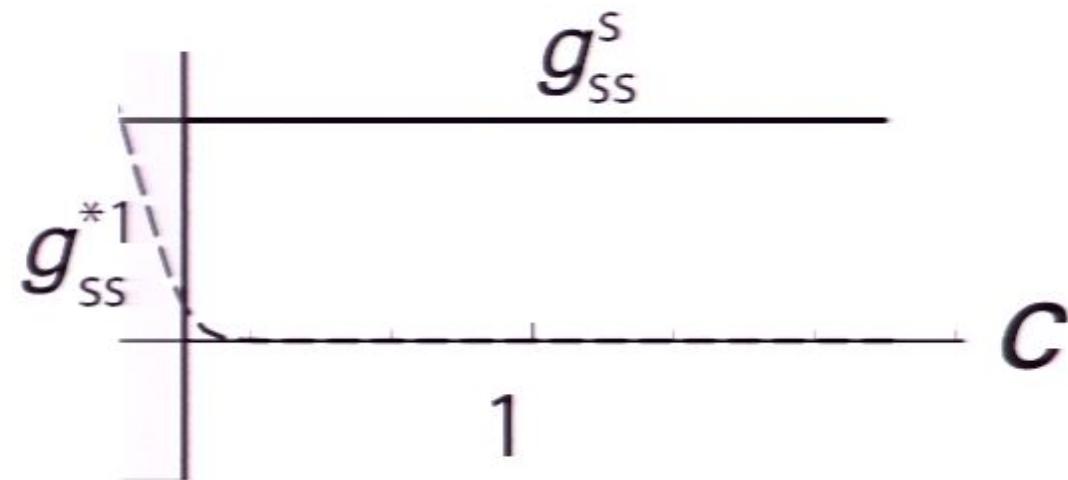


- 3-source vertex dominates

Gauge interactions $g_{\psi\psi}^A$



Key point: composites contain no zero mode

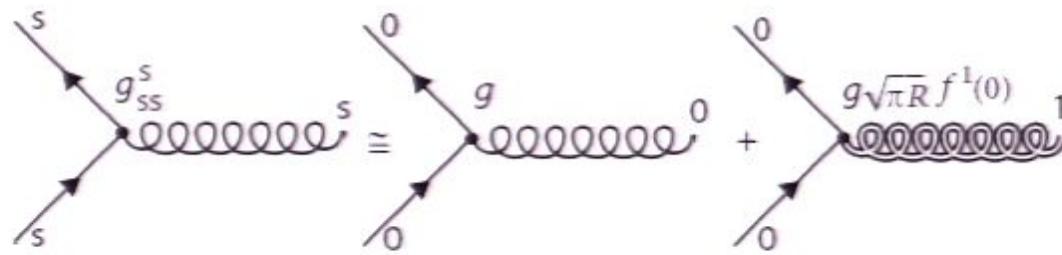


- 3-source vertex dominates



RS GIM mechanism

For light fermions, $c > 1/2$, 3-source vertex dominates:

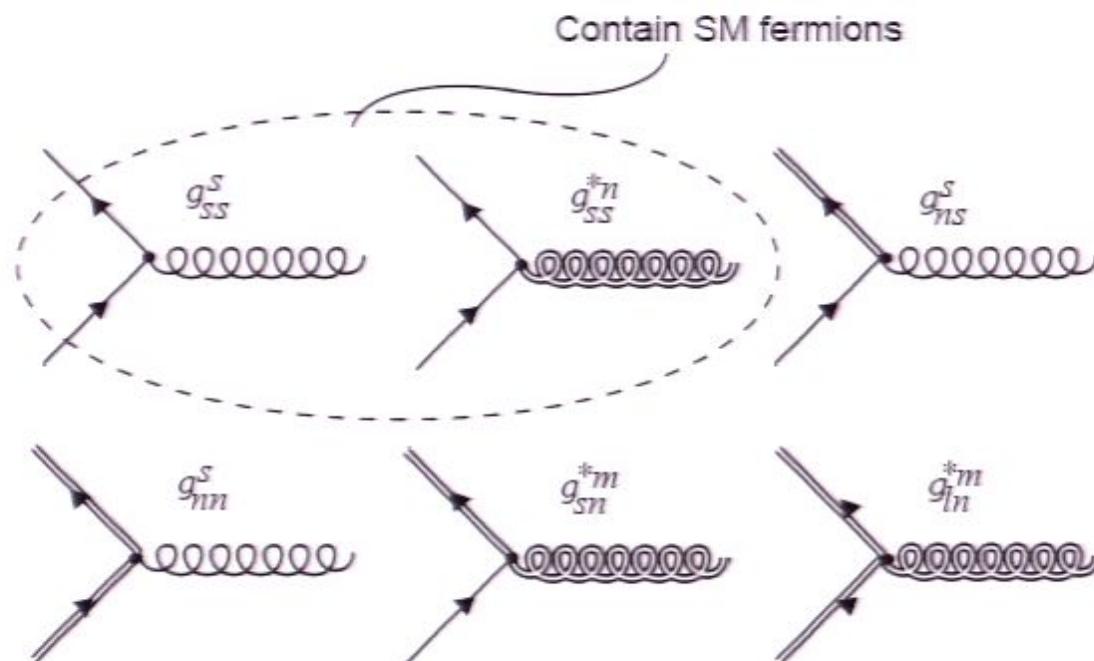


KK gauge boson couplings are approx. universal for light fermions

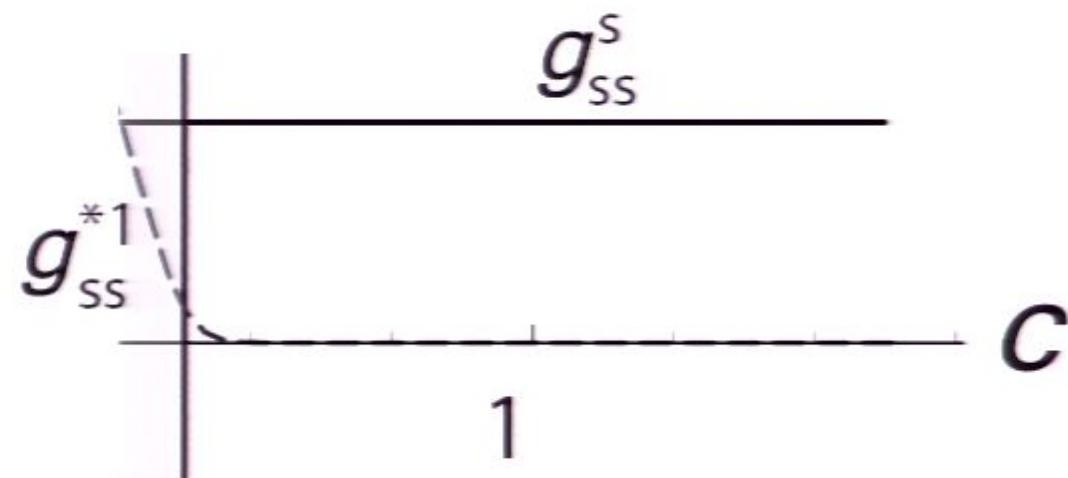
⇒ FCNCs suppressed



Gauge interactions $g_{\psi\psi}^A$



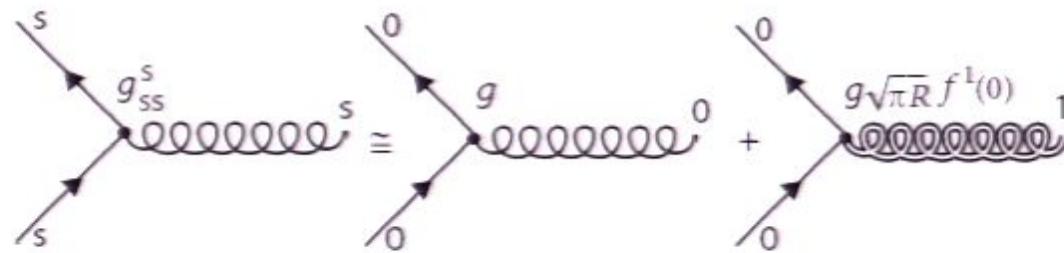
Key point: composites contain no zero mode



- 3-source vertex dominates

RS GIM mechanism

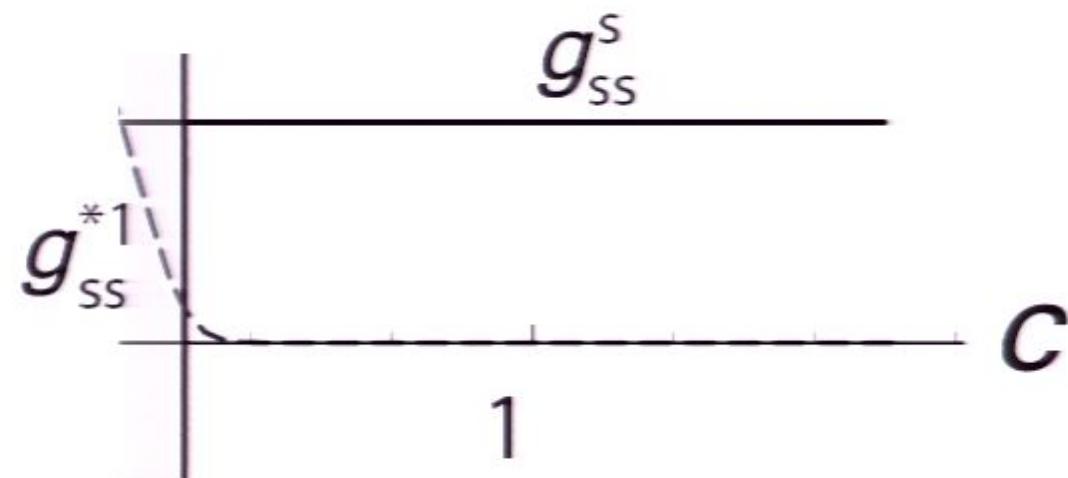
For light fermions, $c > 1/2$, 3-source vertex dominates:



KK gauge boson couplings are approx. universal for light fermions

⇒ FCNCs suppressed

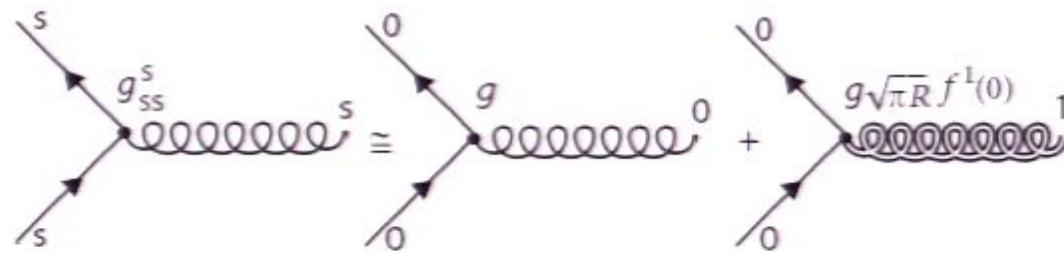




- 3-source vertex dominates

RS GIM mechanism

For light fermions, $c > 1/2$, 3-source vertex dominates:



KK gauge boson couplings are approx. universal for light fermions

⇒ FCNCs suppressed



Right-handed top t_R

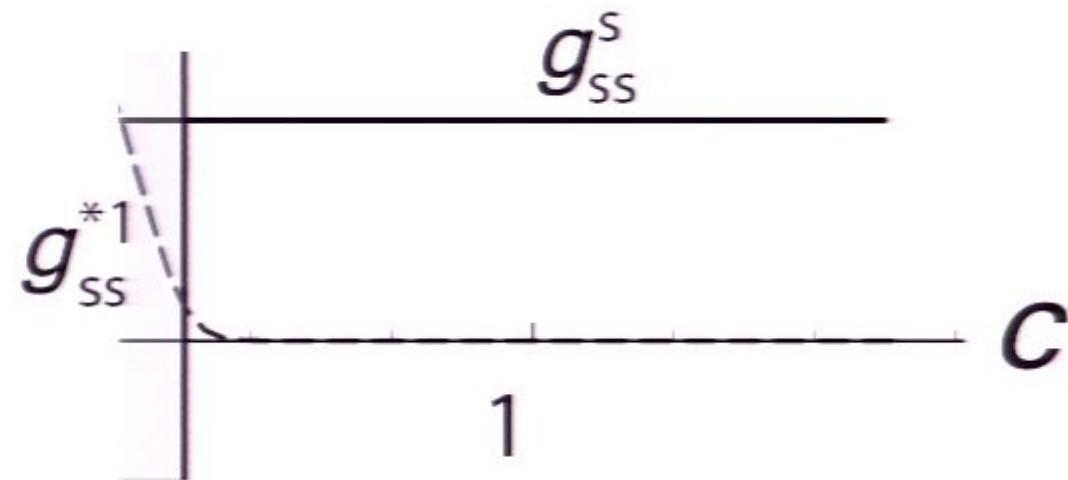
$$\tilde{f}^0(y) = e^{(\frac{1}{2}-c)ky} \quad m_\psi = ck$$

Take e.g. $c = -0.7$; $\Delta = 1.7 \Rightarrow$ relevant mixing

$$\begin{pmatrix} t_R^{(0)} \\ t_R^{(1)} \\ t_R^{(2)} \\ t_R^{(3)} \\ \vdots \end{pmatrix} = \begin{pmatrix} 0.9796 & \sim -1 & \sim 0 & \sim 0 & \dots \\ -0.1816 & \sim 0 & \sim -1 & \sim 0 & \dots \\ 0.0514 & \sim 0 & \sim 0 & \sim -1 & \dots \\ 0.0471 & \sim 0 & \sim 0 & \sim 0 & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} t_R^s \\ t_R^{CFT(1)} \\ t_R^{CFT(2)} \\ t_R^{CFT(3)} \\ \vdots \end{pmatrix}$$

- massless eigenstate $t_R^0(x)$ roughly equal mixture of source/CFT
- KK modes contain elementary component

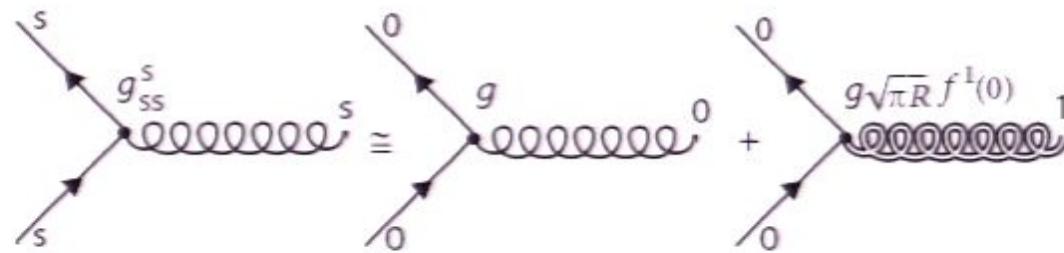




- 3-source vertex dominates

RS GIM mechanism

For light fermions, $c > 1/2$, 3-source vertex dominates:



KK gauge boson couplings are approx. universal for light fermions

⇒ FCNCs suppressed



Conclusions

- Holographic basis: bulk field expanded in source and CFT resonances
- Quantitatively describe elementary/composite mixing in warped duals
- Explain warped physics in terms of strong gauge dynamics
- Things to do:
 - Other applications: Higgsless models, warped SUSY, Gauge-Higgs models (QCD?)
 - Loop diagrams
 - important for EWPT, gauge coupling unification etc.
 - Brane localized kinetic terms - could modify composite content
 - More general geometries?



Holographic basis

Basic idea:

Expand the bulk field directly in terms of a source field $\varphi^s(x)$ and composite CFT states $\varphi_{CFT}^n(x)$:

$$\Phi(x, y) = \varphi^s(x)g^s(y) + \sum_{n=1}^{\infty} \varphi_{CFT}^n(x)g^n(y)$$

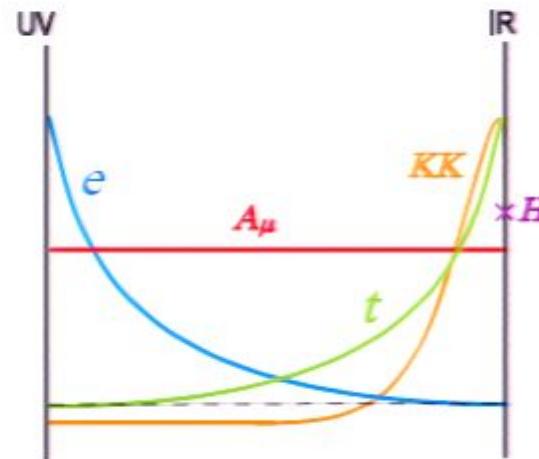
The fields $\phi^n(x)$ are the mass eigenstates

- Spectrum:

$$J_{b-1}\left(\frac{m_n}{k}\right)Y_{b-1}\left(\frac{m_n e^{\pi k R}}{k}\right) - Y_{b-1}\left(\frac{m_n}{k}\right)J_{b-1}\left(\frac{m_n e^{\pi k R}}{k}\right) = 0$$



Partial compositeness of SM fields



- UV localized \iff mostly elementary
- IR localized \iff mostly composite

Can we quantify source/CFT (elementary/composite) mixing?

