Title: The Everett interpretation

Date: Sep 21, 2007 10:00 AM

URL: http://pirsa.org/07090069

Abstract: I shall present an overview of quantum mechanics in the Everett interpretation, that emphasises its structural characteristics, as a theory of what exists. In this respect it shares common ground with other fundamental theories in physics. As such its appeal is conservative; it makes do with the purely unitary equations of quantum mechanics as exceptionless and universal. It also makes do with standard methods for extracting \high level\' or \empty emergent\' ontology, the furniture of macroscopic worlds, from largish molecules on up. It would appeal all the more if it made do with standard epistemological principles too - for example, in the context of inductive statistical confirmation, with standard Bayesian epistemology. But this links to the question of the interpretation of probability in the Everett interpretation, and here the theory seems anything but conservative. It is a common complaint that the approach leaves no room at all for talk of uncertainty. I shall argue, again on conservative interpretative practises, that this claim is incorrect. Chance events are, indeed, revealed in a surprising light - as quantum branchings - but they are the more perspicuous, and their properties and quantitative measure better explained, in light of that.

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The Everett interpretation:

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Oxford University

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The Everett interpretation:

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The Everett interpretation: a user's manual

Simon Saunders

Oxford University

http://users.ox.ac.uk/~everett/

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To solve the measurement problem

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To solve the measurement problem

The problem of how to relate the (unitary, deterministic) equations of motion of quantum mechanics to observed phenomena

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The conventional method (measurement postulates):

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To solve the measurement problem

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The conventional method (measurement postulates):

 (i) Expand the state in a given basis Ψ=Σc_kφ_k depending on what is measured

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To solve the measurement problem

The problem of how to relate the (unitary, deterministic) equations of motion of quantum mechanics to observed phenomena

- The conventional method (measurement postulates):
- (i) Expand the state in a given basis Ψ=Σc_kφ_k depending on what is measured
- (ii) The state collapses to φ_k on measurement with probability | Ω₁₄₃|²

To solve the measurement problem

The problem of how to relate the (unitary, deterministic) equations of motion of quantum mechanics to observed phenomena

The theory is to apply universally

- The conventional method (measurement postulates):
- (i) Expand the state in a given basis Ψ=Σc_kφ_k depending on what is measured
- (ii) The state collapses to φ_k on measurement with probability | γ₁₄₃|²

To solve the measurement problem

The problem of how to relate the (unitary, deterministic) equations of motion of quantum mechanics to observed phenomena

- The theory is to apply universally
- Without any special mention of 'the observer' or 'measurement'

The conventional method (measurement postulates):

- Expand the state in a given basis Ψ=Σc_kφ_k depending on what is measured
- (ii) The state collapses to φ_k on measurement with probability | Ω₁₄₁₄₃|²

To solve the measurement problem

The problem of how to relate the (unitary, deterministic) equations of motion of quantum mechanics to observed phenomena

- The theory is to apply universally
- Without any special mention of 'the observer' or 'measurement'
- And without any special interpretative assumptions or additional equations

The conventional method (measurement postulates):

- Expand the state in a given basis Ψ=Σc_kφ_k depending on what is measured
- (ii) The state collapses to φ_k on measurement with probability | G₁₄₃|²



Why don't we see the earth move?



Why don't we see the earth move?

 if the earth moved, there should be a great wind



Why don't we see the earth move?

- if the earth moved, there should be a great wind
- if the earth moved, falling bodies would fall behind in the direction

of its motion



Why don't we see the earth move?

Answer: the earth and everything on its surface and in its atmosphere are subject to the same motion



Why don't we see the earth move?

Answer: the earth and everything on its surface and in its atmosphere are subject to the same motion

Only relative motions are

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Why don't we see superpositions?



Why don't we see superpositions?

 if there are only superpositions, we would all turn into blacmange



Why don't we see superpositions?

- if there are only superpositions, we would all turn into blacmange
- if there are only superpositions, there would be no meaning to probability



Why don't we see superpositions?

Answer: when macroscopic superpositions arise, everything enters into the superposition, including the apparatus and the environment



Why don't we see superpositions?

Answer: when macroscopic superpositions arise, everything enters into the superposition, including the apparatus and the environment

Only correlations are observable



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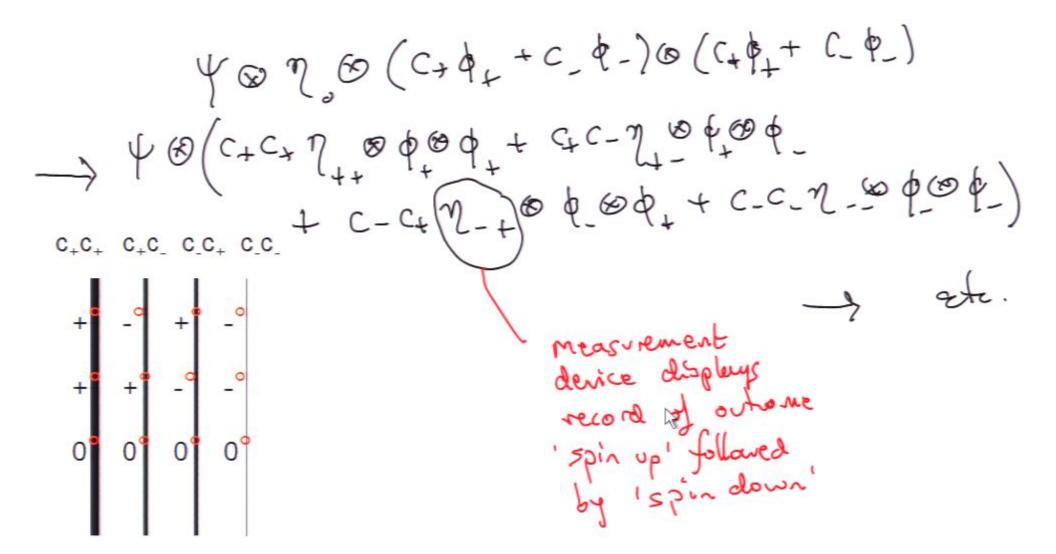
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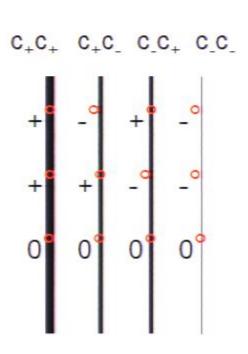
$$+ C_{-}C_{+}\eta_{-} + \otimes \varphi_{-}$$

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To solve the measurement problem...

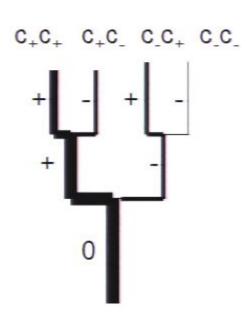
Jniversality • no mention of measurement • no special interpretative assumption or additional eqs

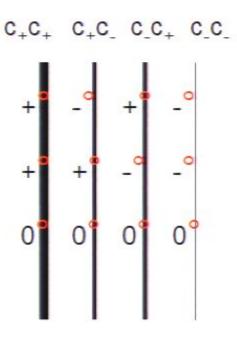


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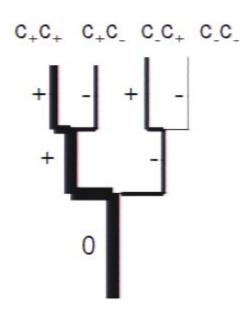


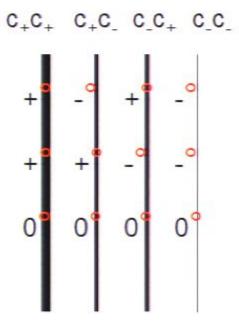


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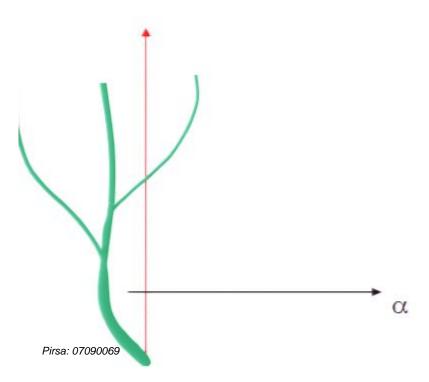




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To solve the measurement problem...

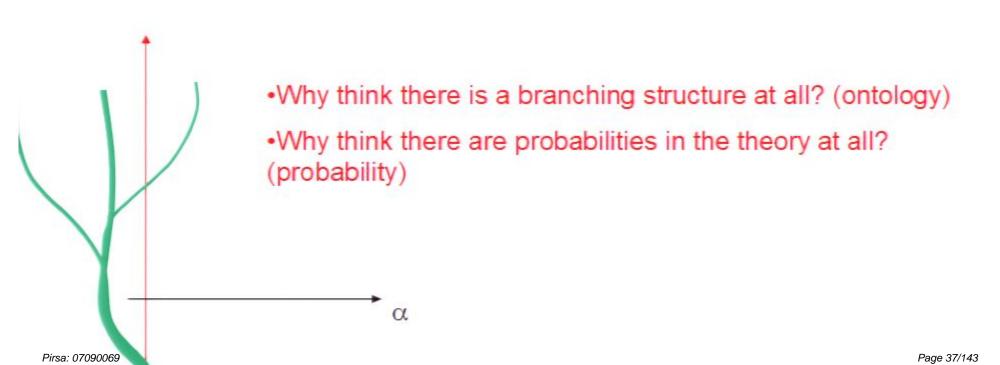
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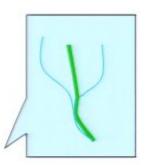


Why many worlds?

To solve the measurement problem...

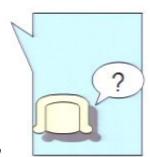
Jniversality • no mention of measurement • no special interpretative assumption or additional eqs

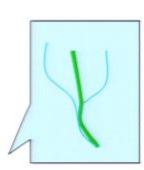




Ontology

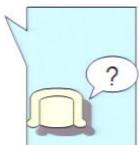
Probability



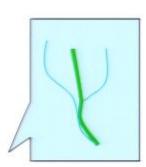


Ontology

Probability



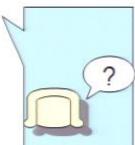
Incoherence problem



Ontology

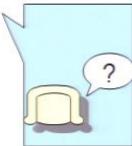
Probability

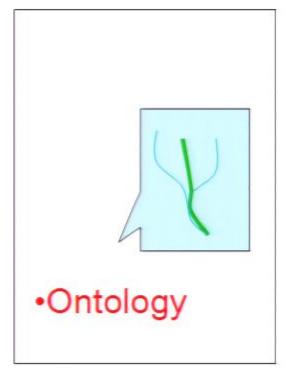
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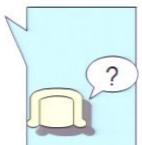


Incoherence problem

quantitative problem



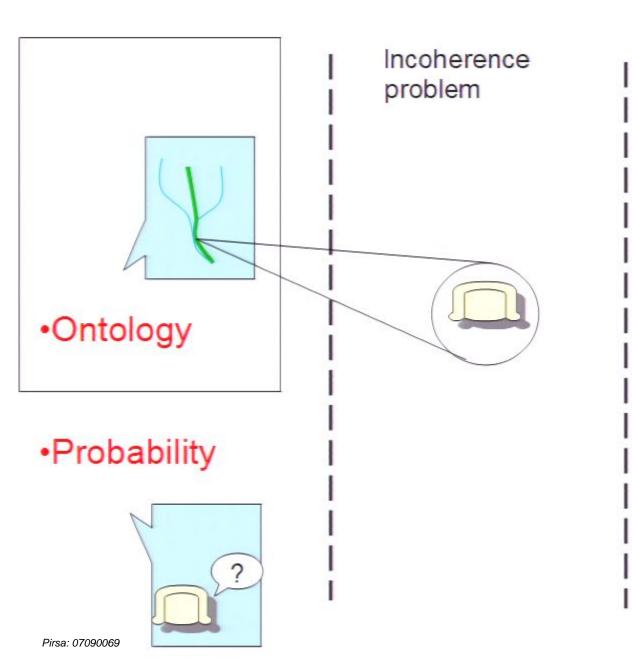




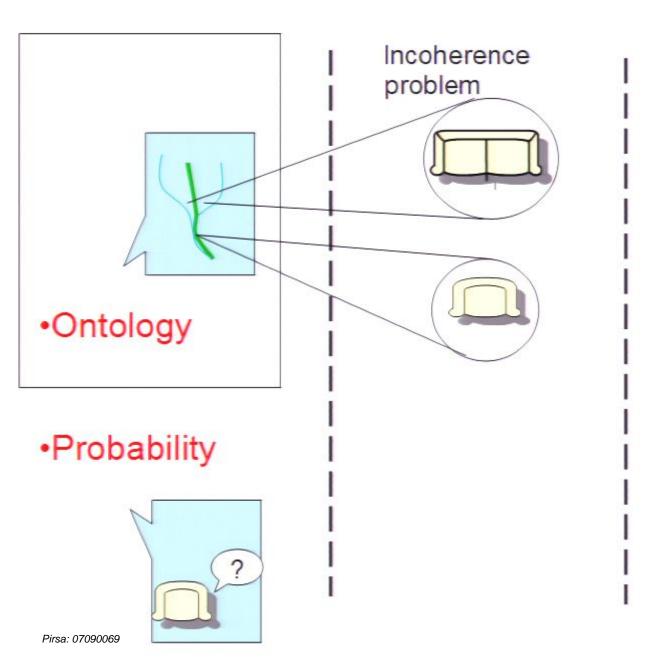
Incoherence problem

quantitative problem

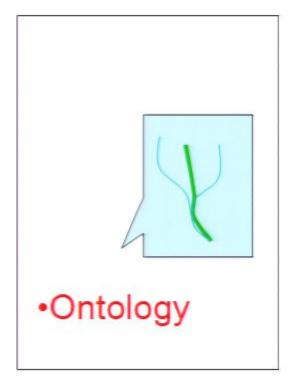


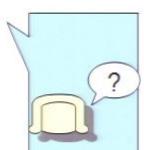


quantitative problem

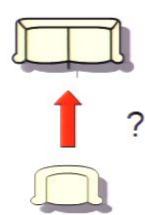


quantitative problem



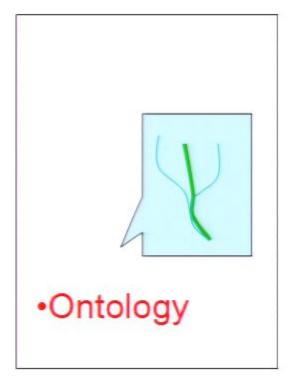


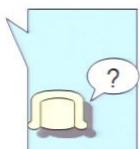
Incoherence problem



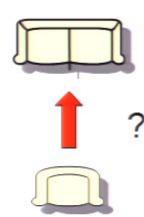
quantitative problem





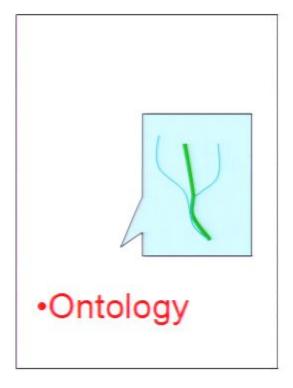


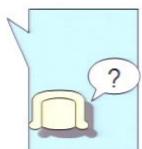
Incoherence problem



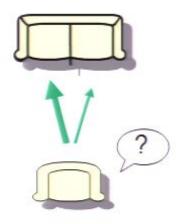
the problem arises with any materialistic theory of the self

quantitative problem

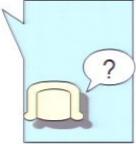




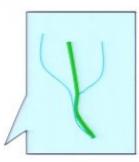
Incoherence problem



quantitative problem

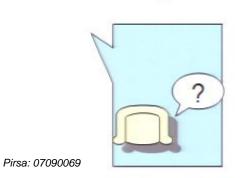


suppose solved

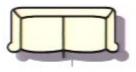


Ontology

Probability



Incoherence problem

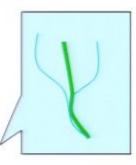




quantitative problem

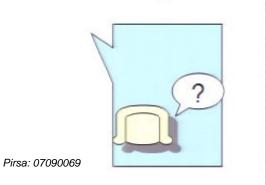
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suppose solved

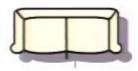


Ontology

Probability



Incoherence problem



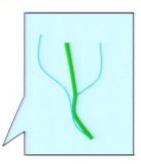


Uncertainty intelligible-'chance'

Uncertainty unintelligible-'caring measure' quantitative problem

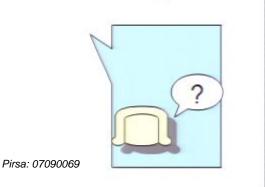
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suppose solved

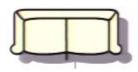


Ontology

Probability



Incoherence problem



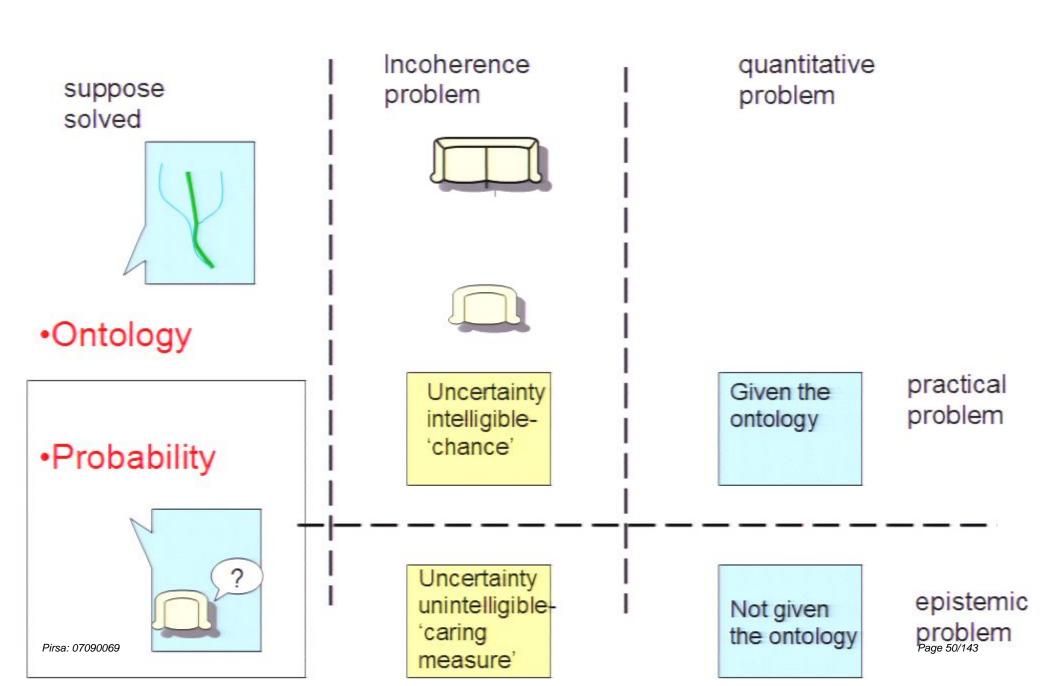


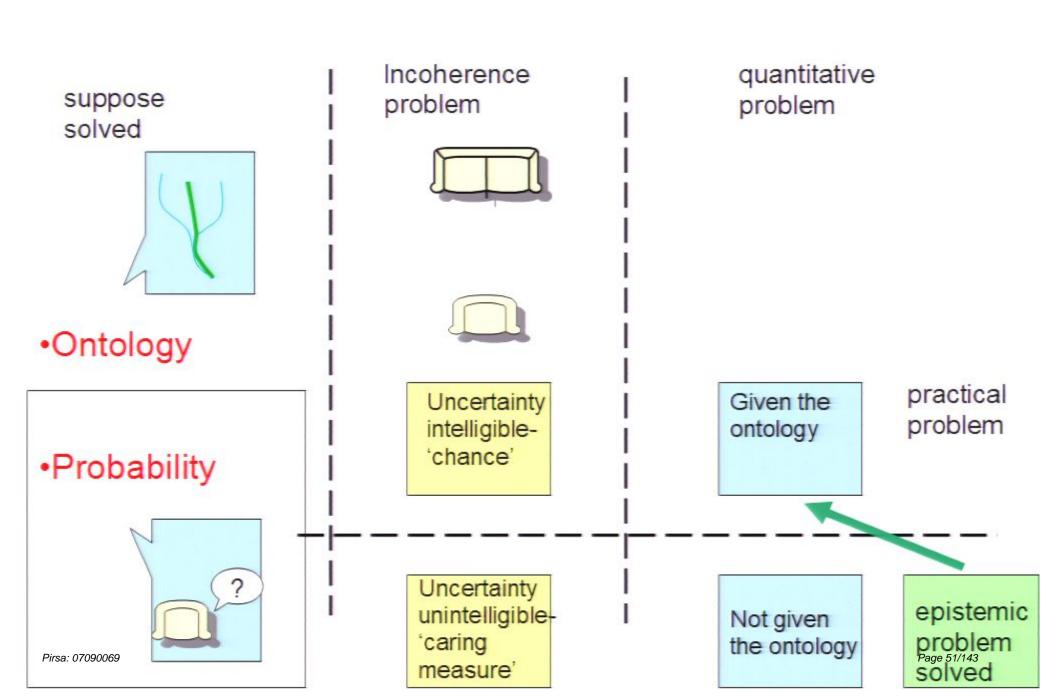
Uncertainty intelligible-'chance'

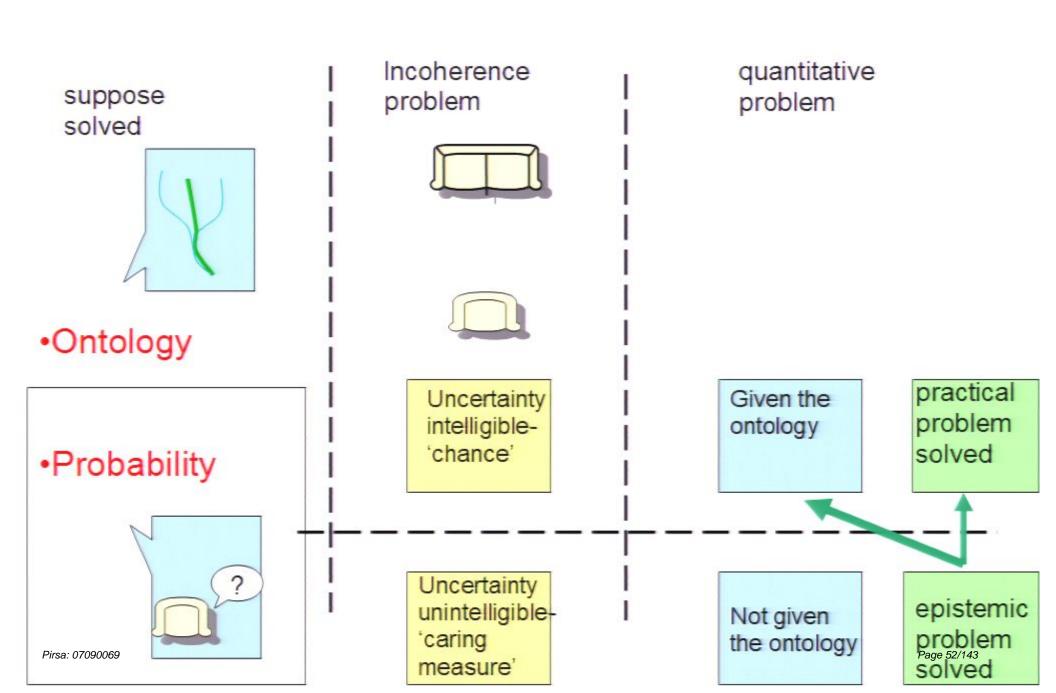
Uncertainty unintelligible-'caring measure' quantitative problem

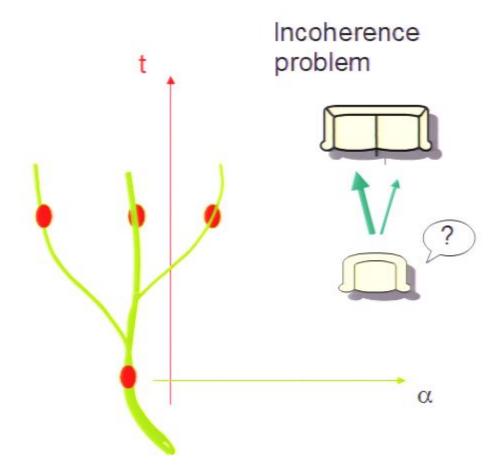
Given the ontology

Not given the ontology

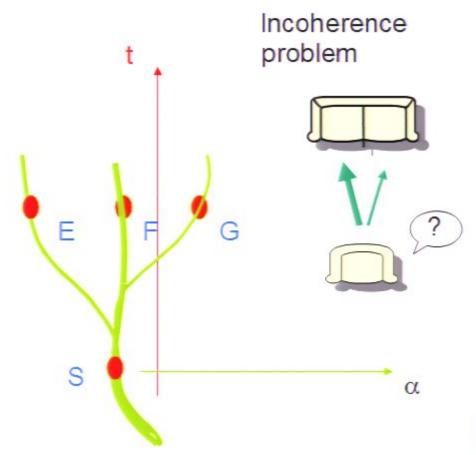






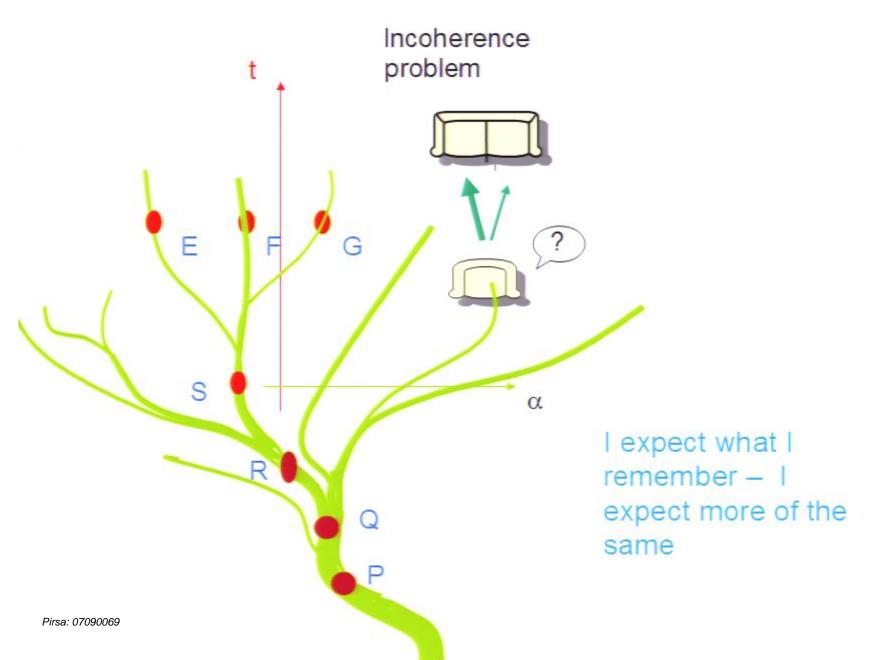


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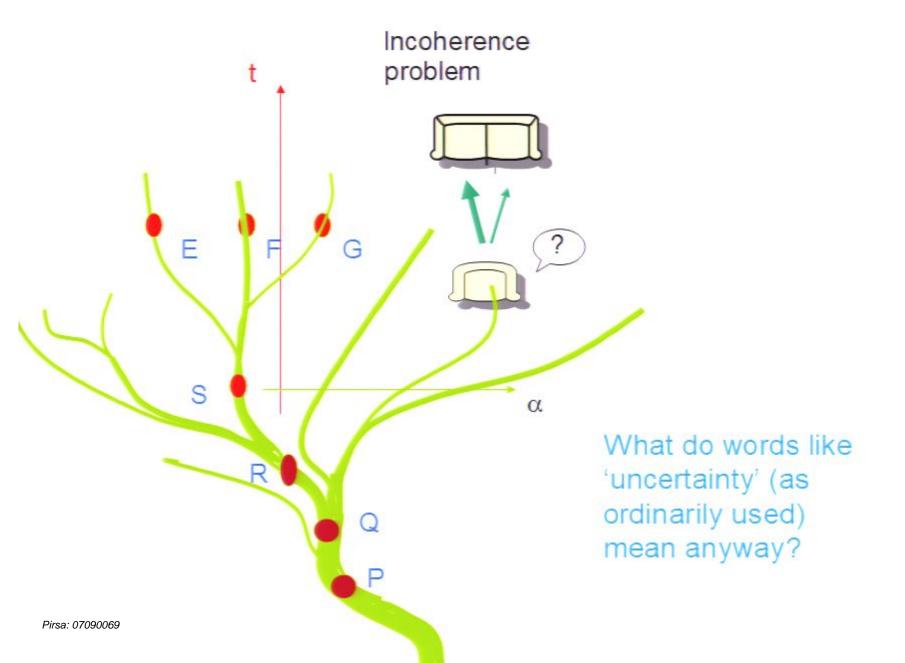


I expect what I remember – I expect more of the same

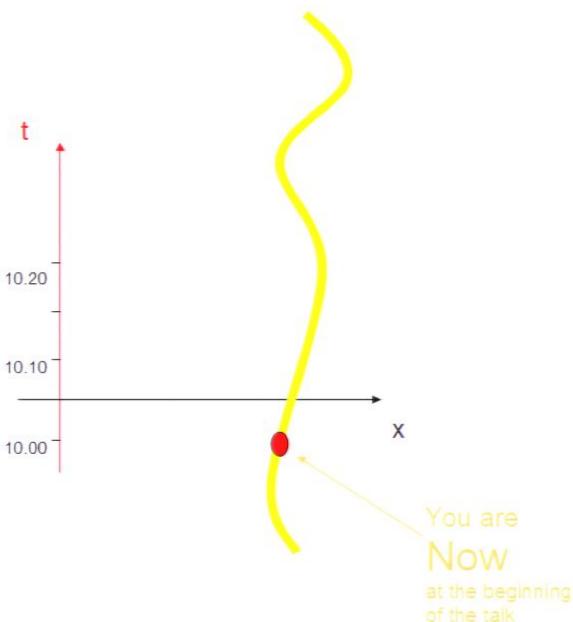
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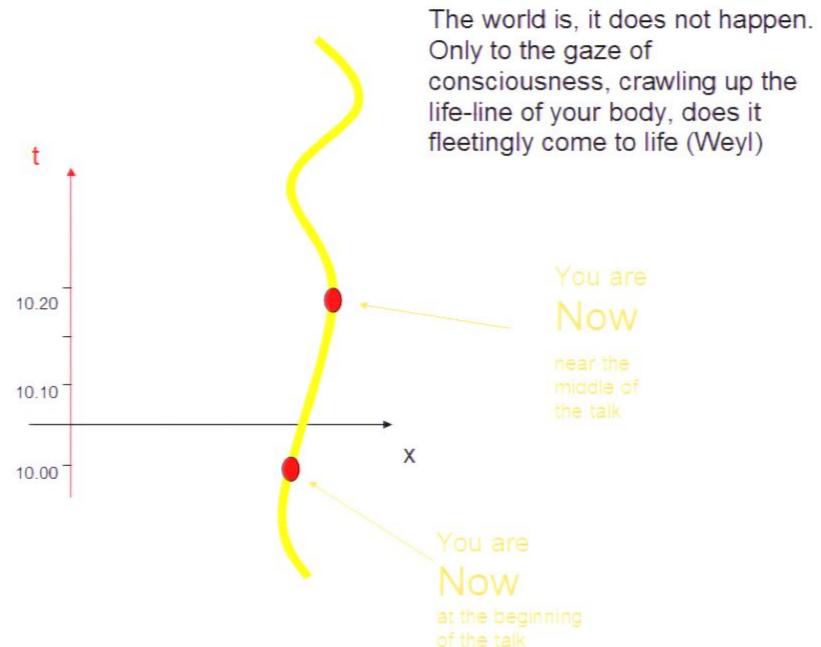
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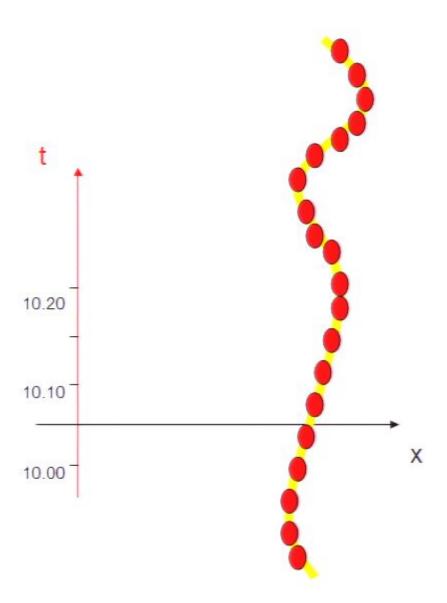
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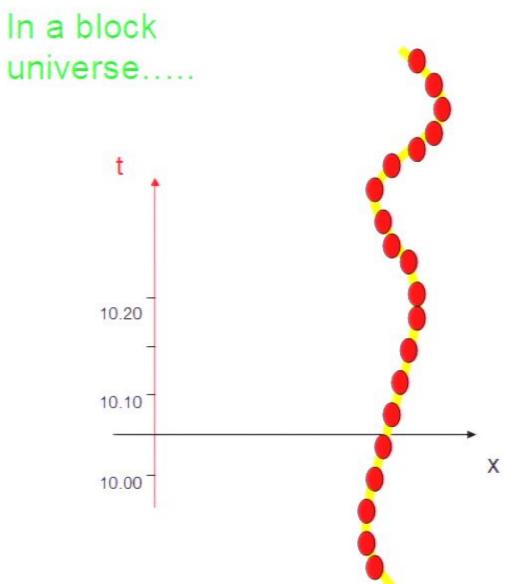
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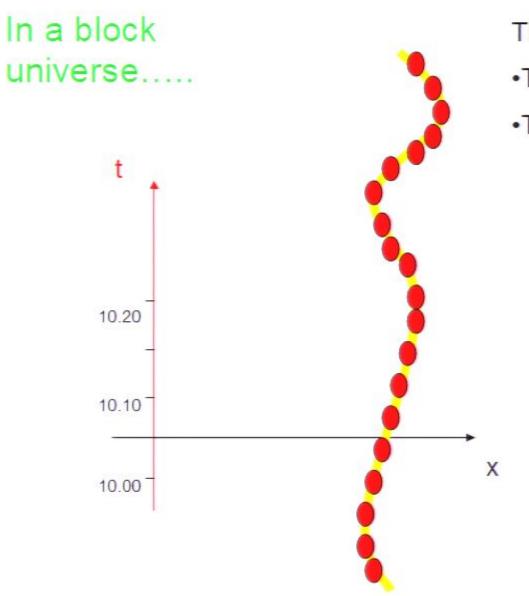


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There is no change

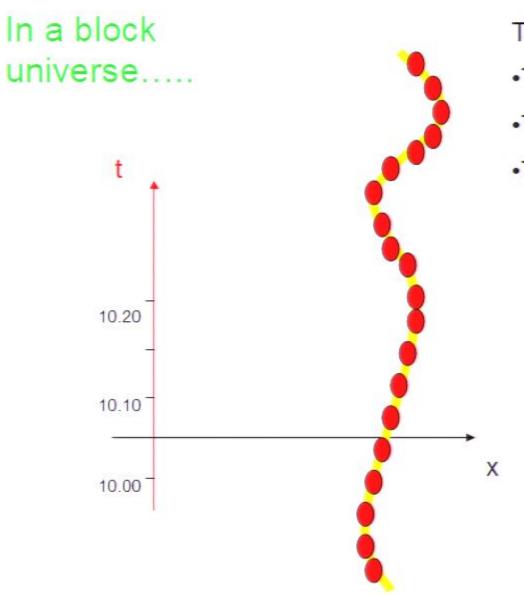
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There is no change

- •Time does not flow
- There is no becoming

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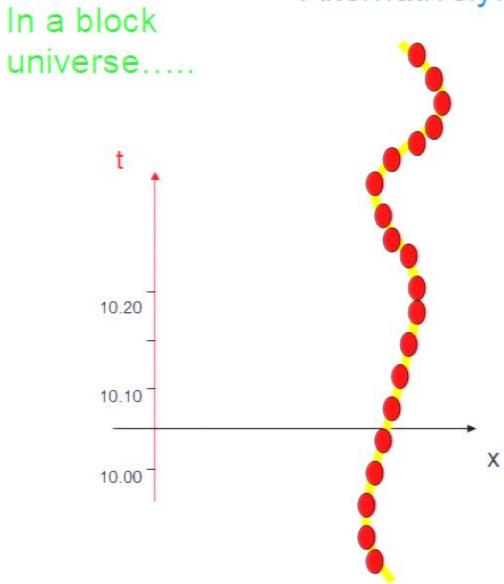


There is no change

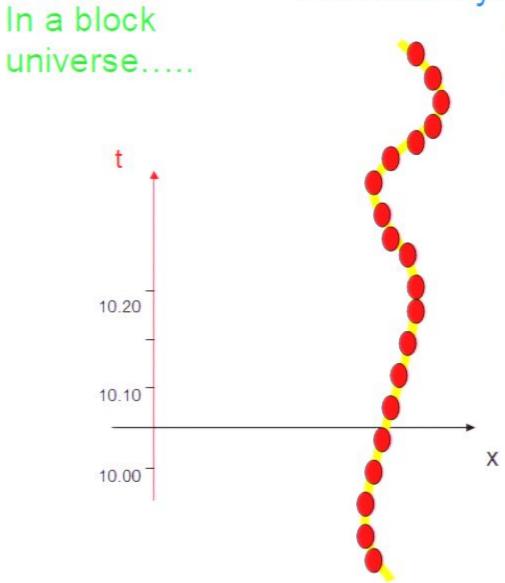
- Time does not flow
- There is no becoming
- •There is no free-will

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Alternatively:



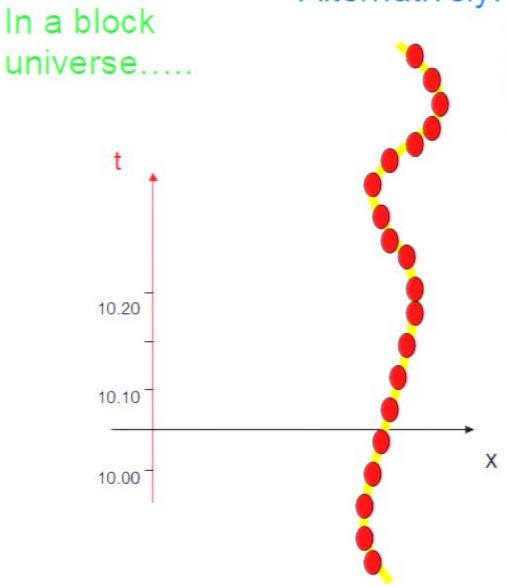




Change resides in the relations among events throughout your life-line

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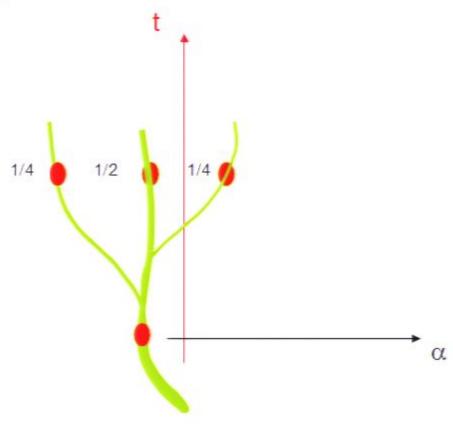




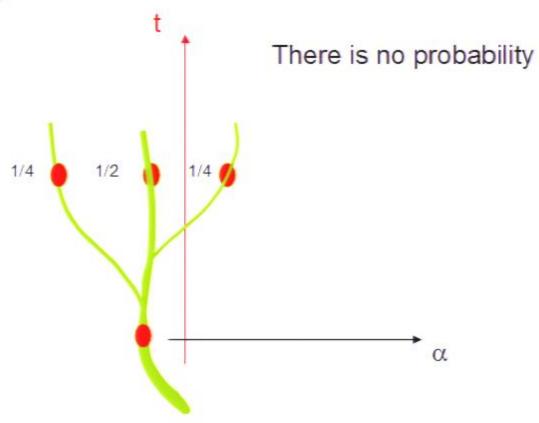
Change resides in the relations among events throughout your life-line

- The flow of time is....
- ·Becoming is...
- Free-will is...

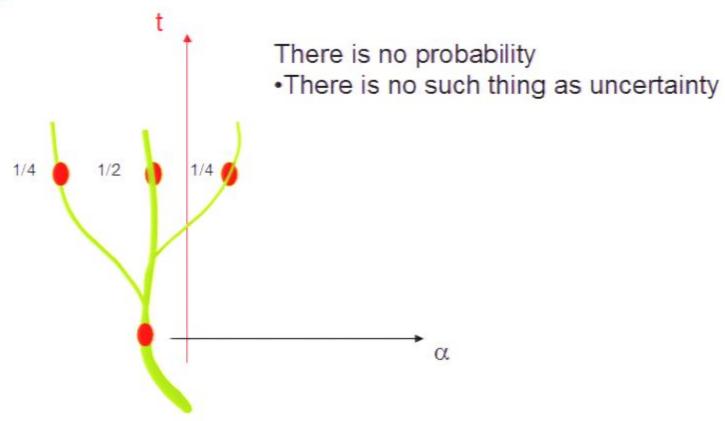
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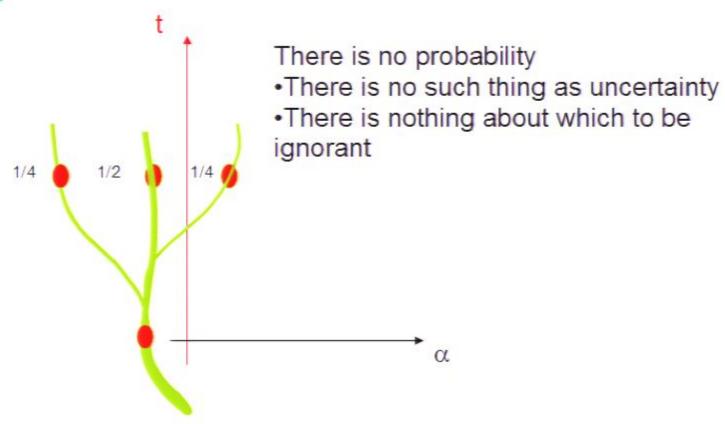
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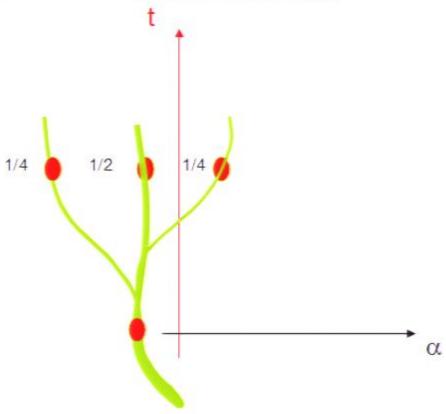


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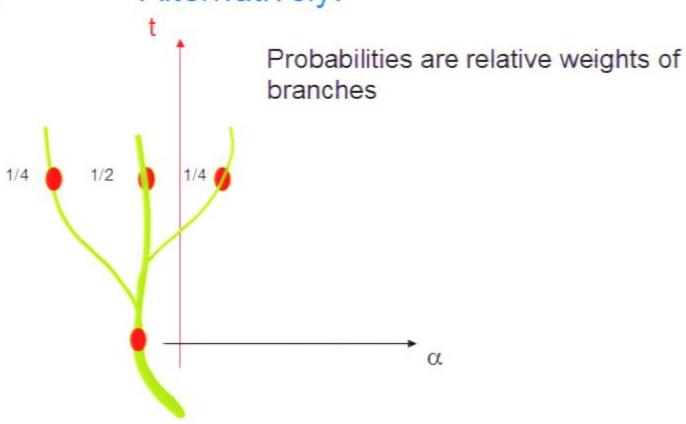
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Alternatively:



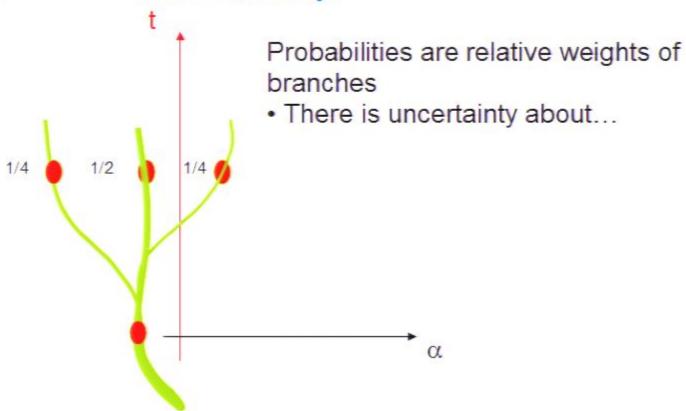
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Alternatively:



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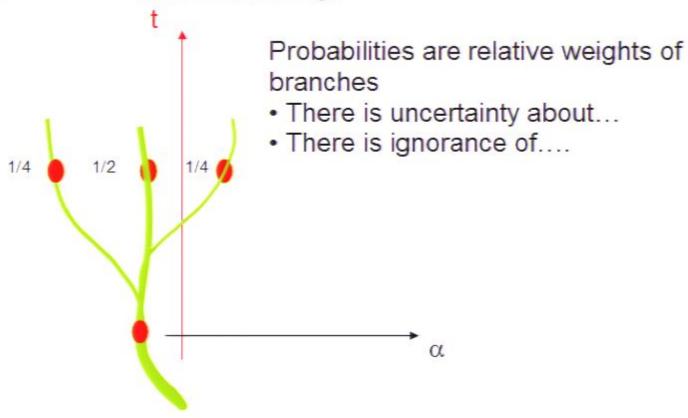
Alternatively:



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In a branching universe....

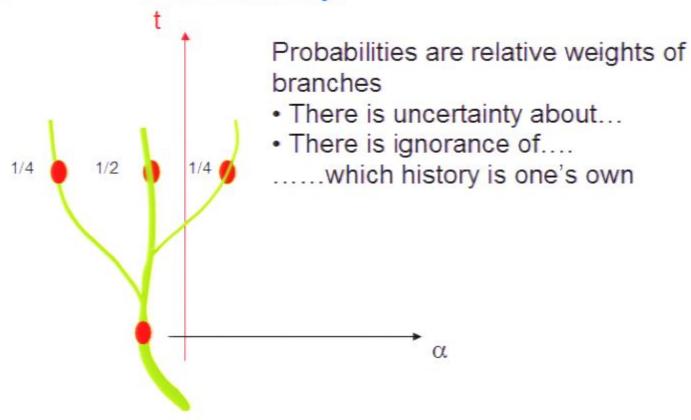
Alternatively:



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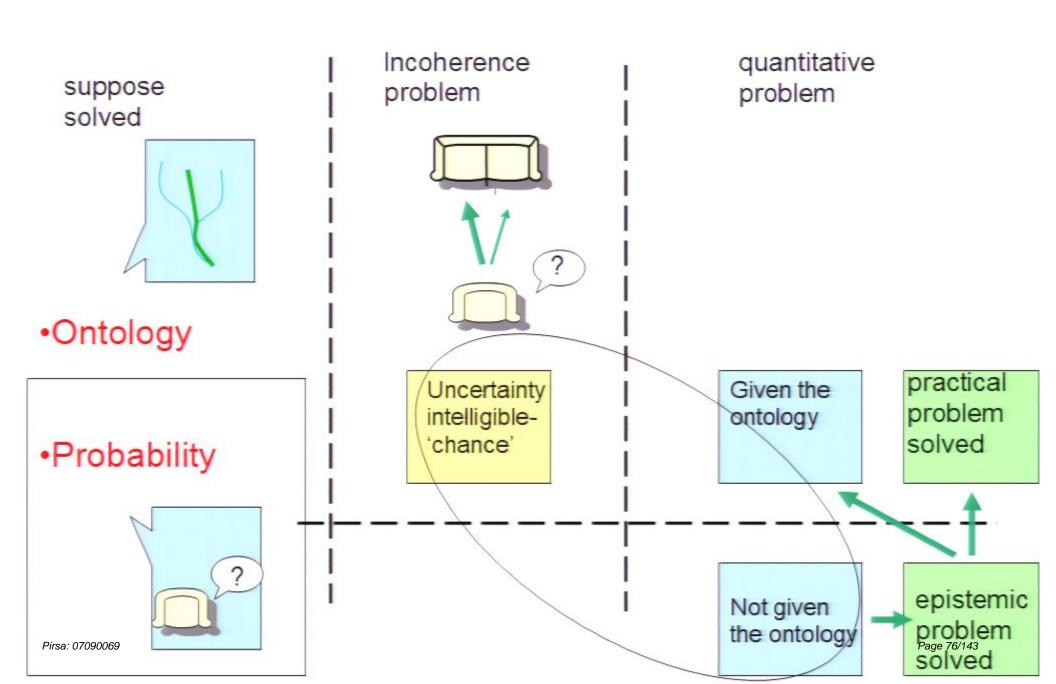
In a branching universe....

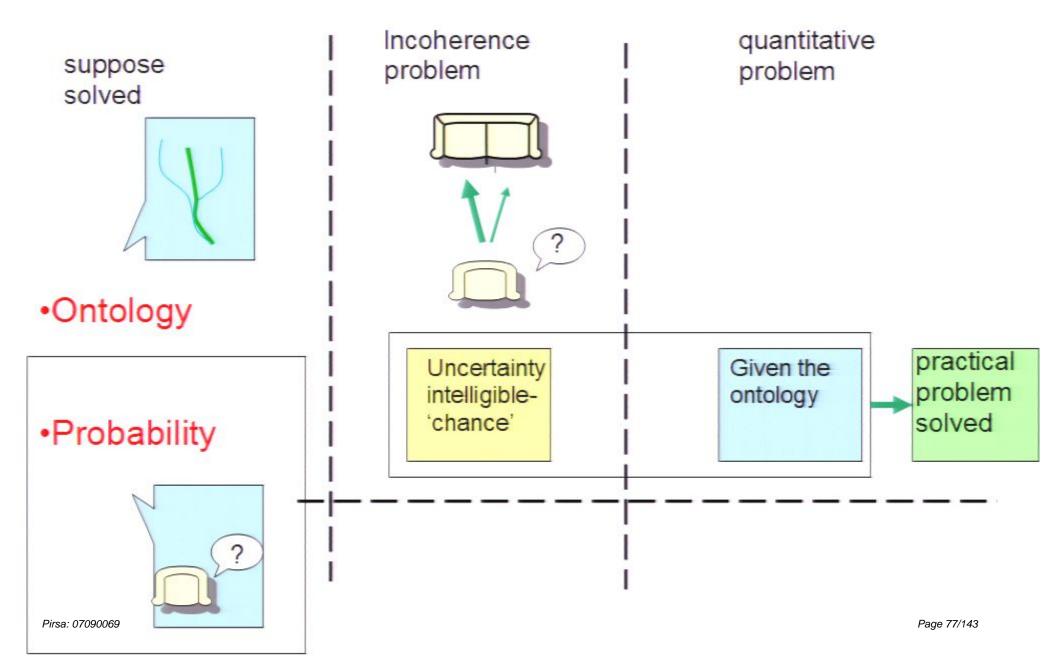
Alternatively:

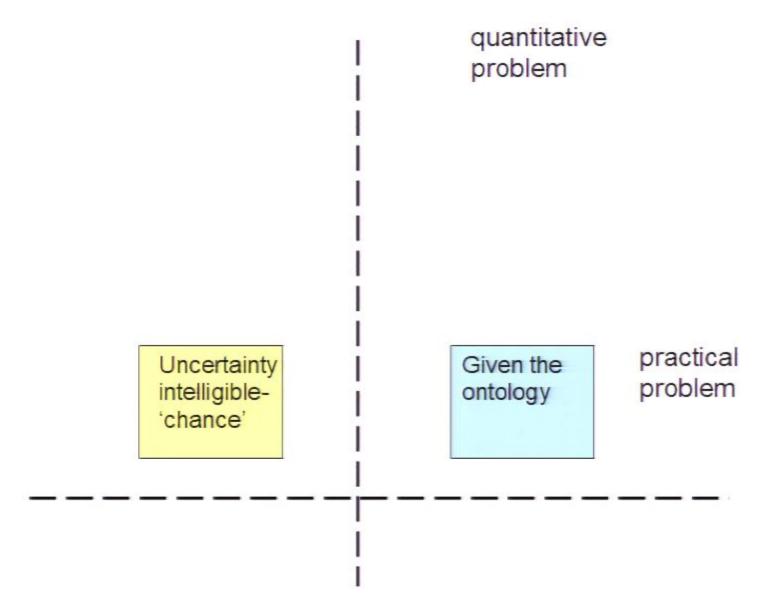


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Incoherence quantitative suppose problem problem solved Ontology Uncertainty intelligible-'chance' Probability Page 75/143 Pirsa: 07090069







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 Introduce axioms of rationality, to be satisfied in the rankings of preferences of a rational actor among quantum games, in accordance with his utilities

quantitative problem

Uncertainty intelligible-'chance' Given the ontology

practical problem

 Introduce axioms of rationality, to be satisfied in the rankings of preferences quantitative of a rational actor among problem quantum games, in accordance with his utilities Prove that there exists an essentially unique utliity function and probability distribution over outcomes of games, under which the preferences of the actor are the same as those which practical Uncertainty Given the maximize the expected problem intelligibleontology utilities 'chance'

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 And prove (this is the essential point) that that probability measure is the quantitative same as the Born-rule problem measure of quantum mechanics, independent of the agent's utility function, if only it agrees with it in the case of EQUALLY weighted branches (i.e. given the 'equivalence rule') - the Deutsch-Wallace representation theorem practical Uncertainty Given the problem intelligibleontology 'chance'

Pirsa: 07090069 Page 81/143

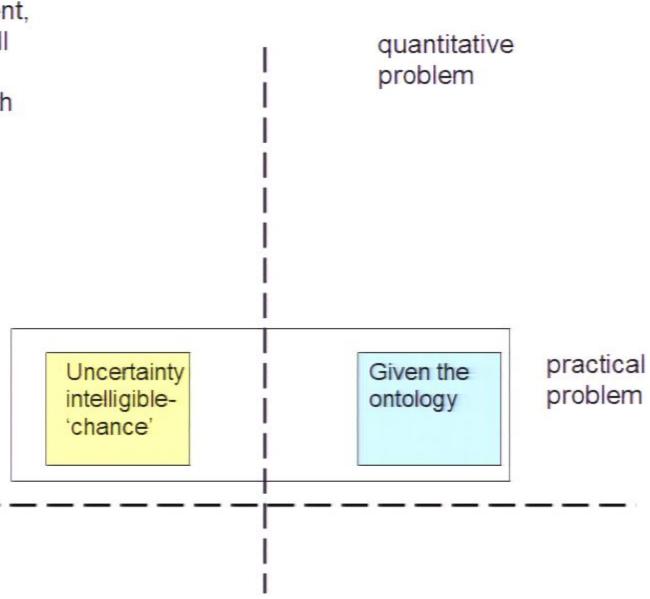
Conclude: if he obeys the equivalence rule, the agent, on pain of irrationality, will maximize his expected utilities in accordance with the Born rule.

quantitative problem

Uncertainty intelligible-'chance' Given the ontology

practical problem

Conclude: if he obeys the equivalence rule, the agent, on pain of irrationality, will maximize his expected utilities in accordance with the Born rule.



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Conclude: if he obeys the equivalence rule, the agent, on pain of irrationality, will maximize his expected utilities in accordance with the Born rule.

quantitative problem

Given the equivalence rule

Uncertainty intelligible-'chance' Given the ontology

practical problem solved

Why the equivalence rule?

Consider a Stern-Gerlach experiment, in which an atom is prepared in the superposion:

$$\phi = \frac{1}{\sqrt{2}} \left(+ \phi_{-} \right)$$

Game 1: The agent receives the payoff iff the result is spin up Game 2: The agent receives the payoff iff the results is spin downaw

inel:
$$Y \otimes Z \otimes \phi \rightarrow Y \otimes \frac{1}{72} (N_{+} \otimes \phi_{+} + N_{-} \otimes \phi_{-})$$

Some 2: $Y \otimes N \otimes \phi \rightarrow Y' \otimes \frac{1}{52} (N_{+}^{7} \otimes \phi_{+} + N_{-} \otimes \phi_{-})$
where Y' , Y'' differ in inessential ways that the agent doesn't care about

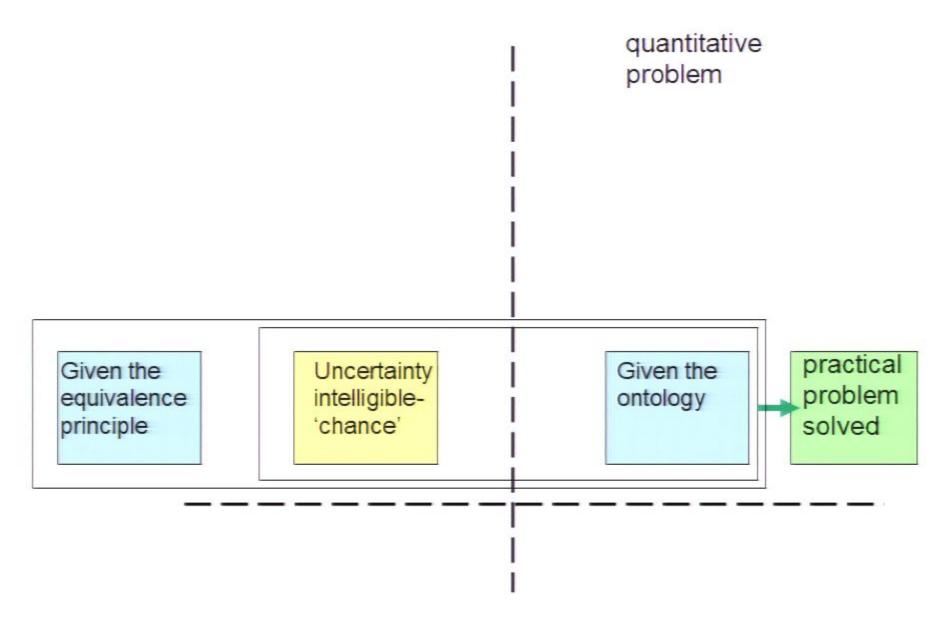
Why the equivalence rule?

$$n^{r} = |X\rangle |Reward\rangle$$
 $n^{r} = |B\rangle |Reward\rangle$
 $n^{r} = |X\rangle |No Reward\rangle$
 $n^{r} = |B\rangle |Reward\rangle$
 $n^{r} = |B\rangle |No Reward\rangle$

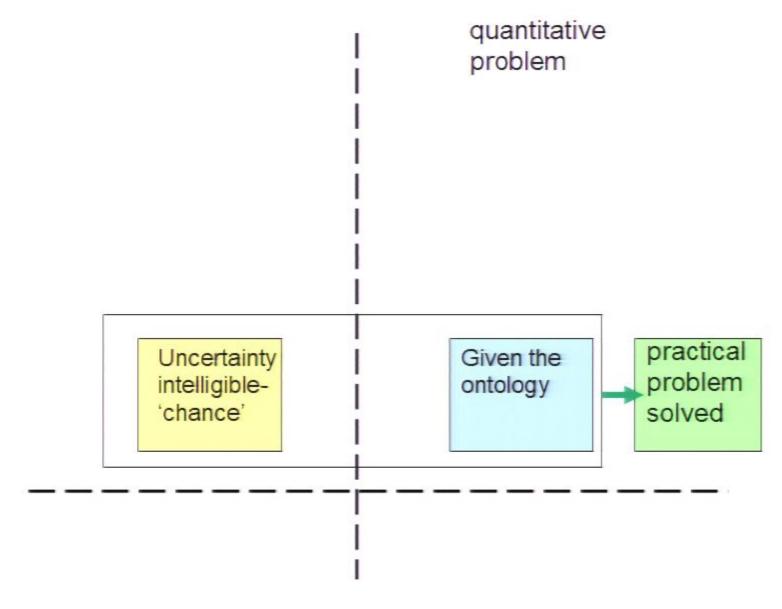
So the choice & between

Why the equivalence rule?

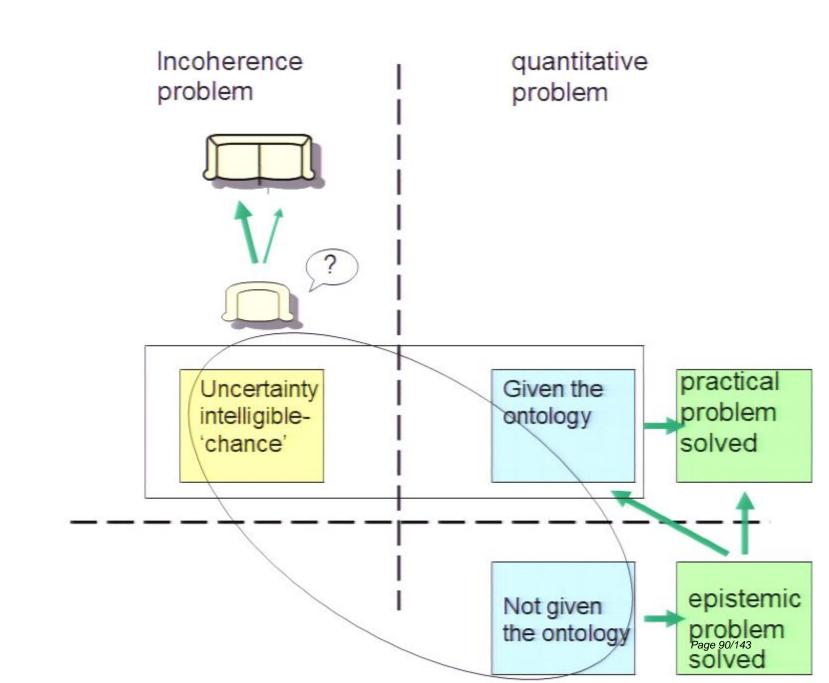
But he agent doesn't care about (d) and (B)! So let the numerical value displayed be crossed The result is: I (1 crosed) (remard) + | ersed) (no reword) Jame 1: (levased > (No revered > + levered) (Revered >) Game 2:

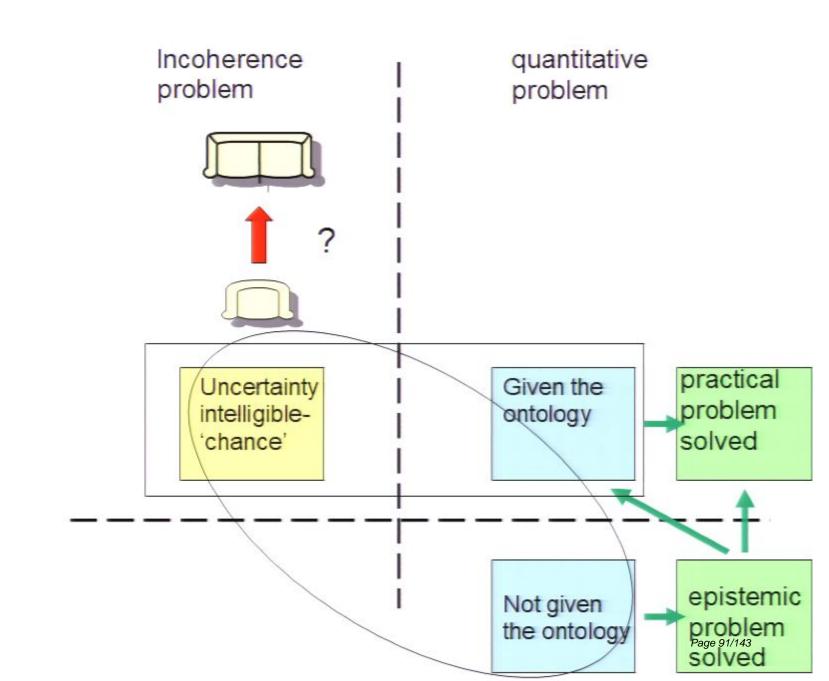


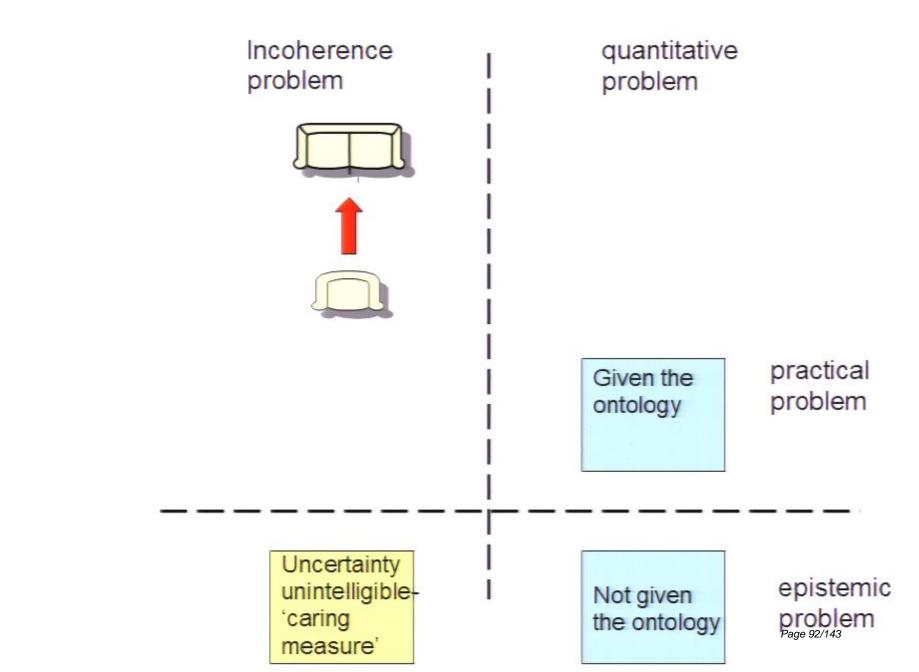
Pirsa: 07090069 Page 88/143

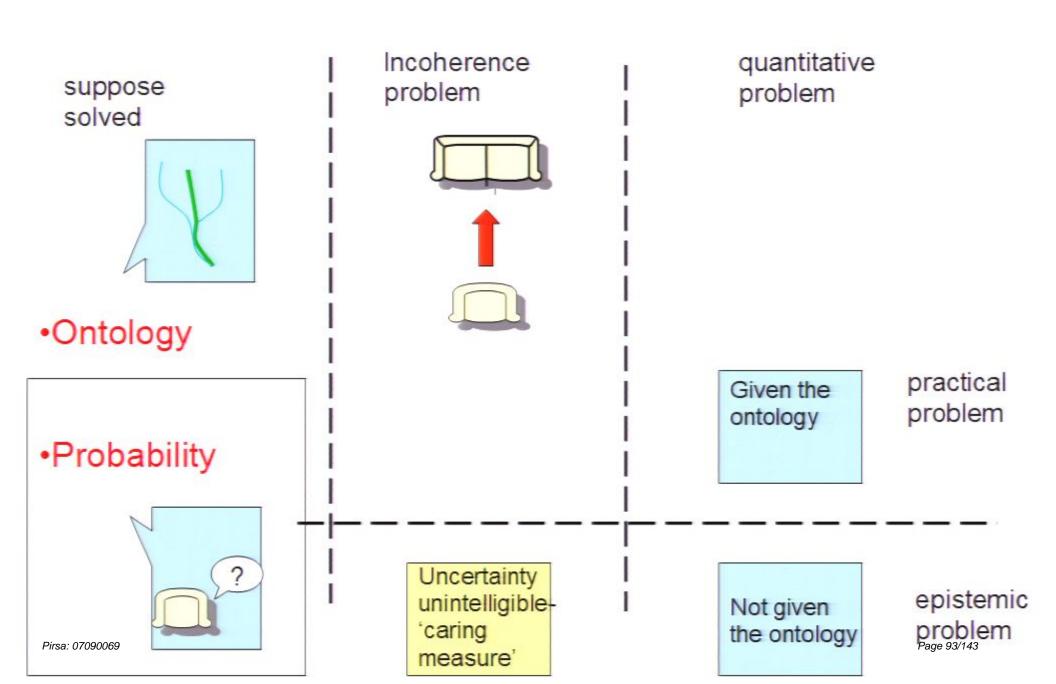


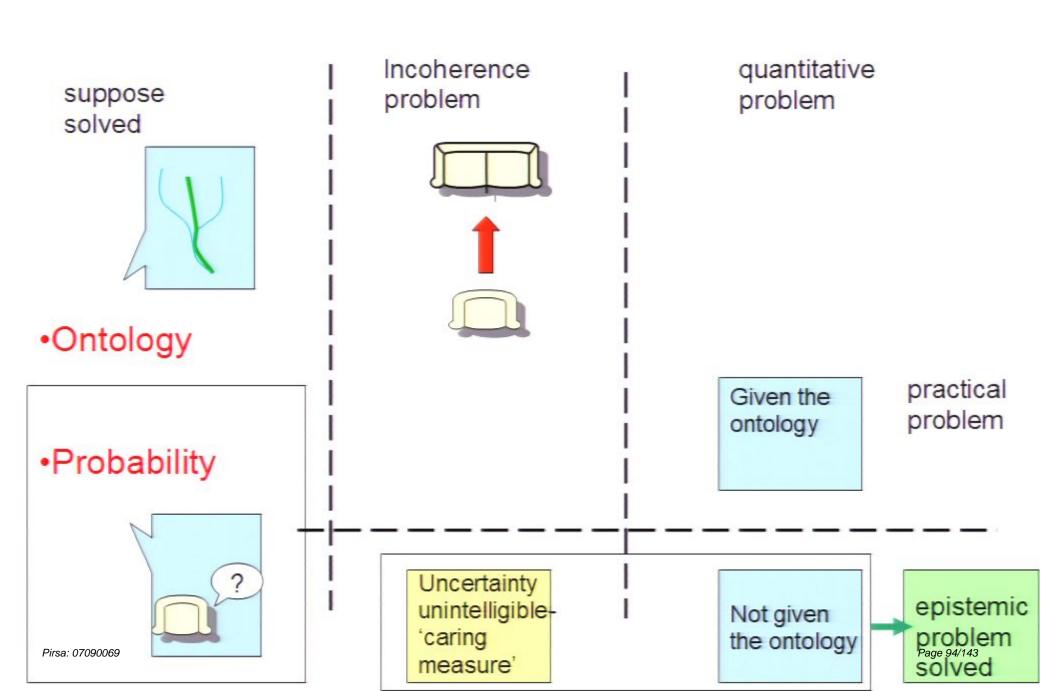
Pirsa: 07090069 Page 89/143

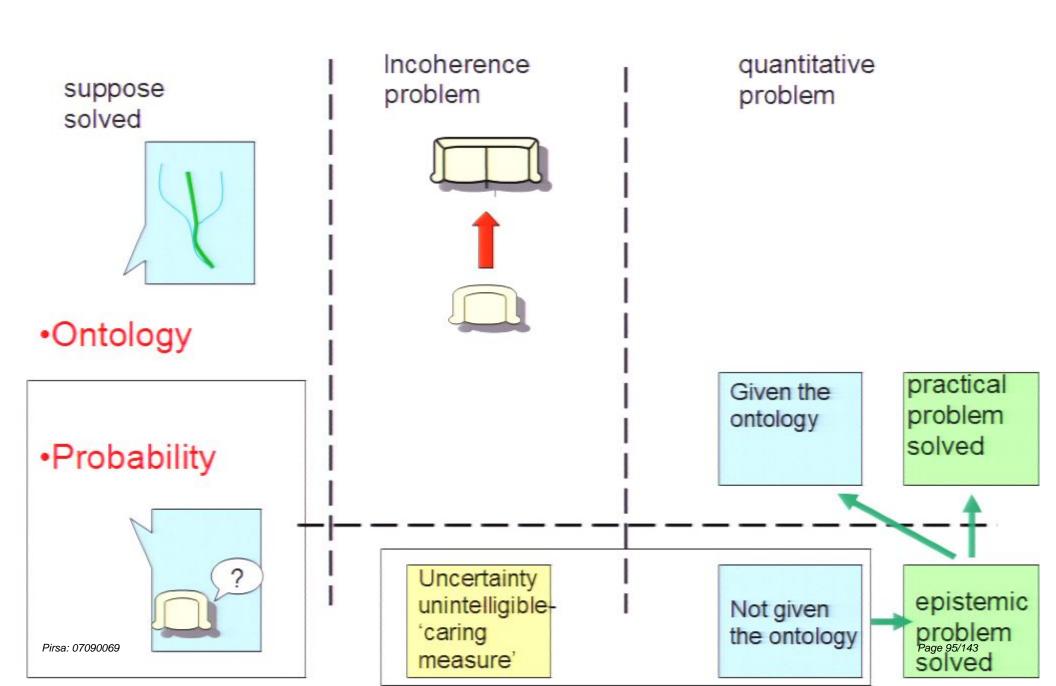


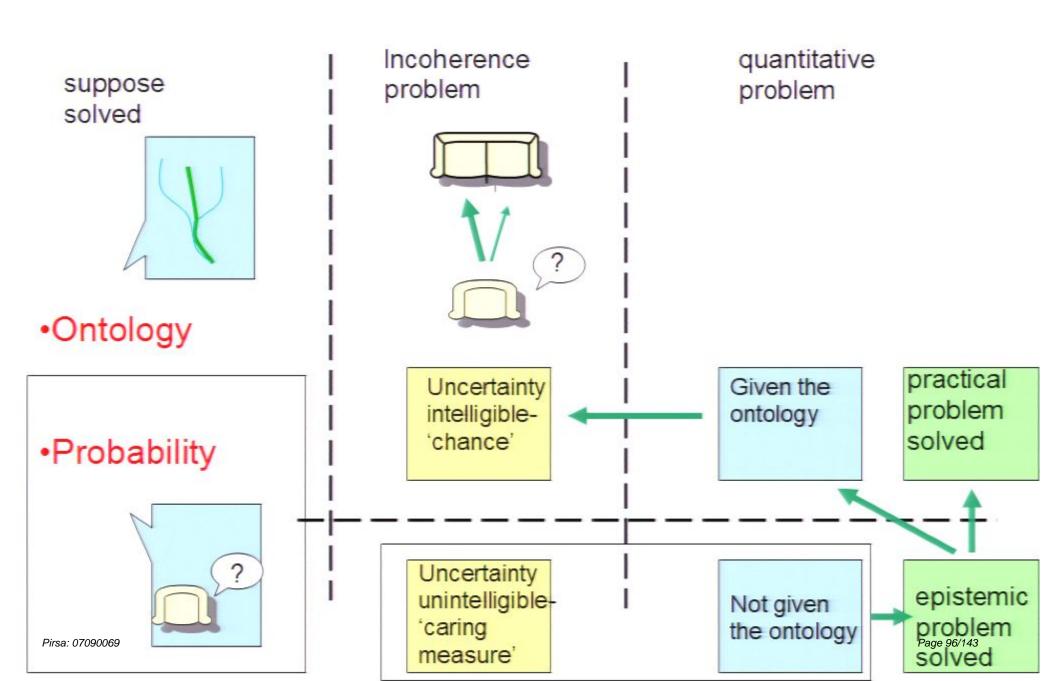




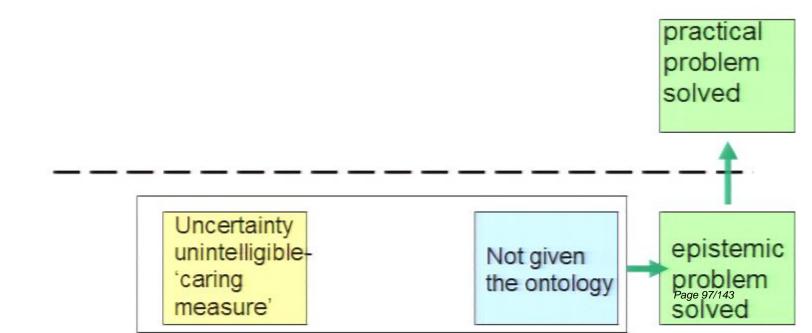






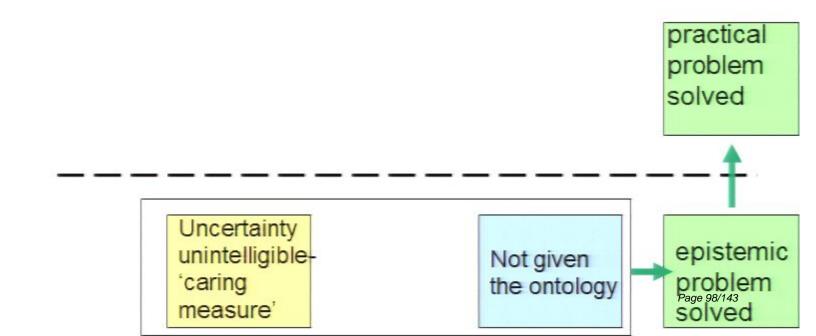


quantitative problem



 Introduce axioms of rationality as before

 to be satisfied in the rankings of preferences of a rational actor among quantum games quantitative problem



 Introduce axioms of rationality (Savage)

 to be satisfied in the rankings of preferences of a rational actor among quantum games.

•Prove there exists an essentially unique utility function and probability distribution over outcomes of games, under which the preferences of the actor are the same as those which maximize the expected utilities.

quantitative problem

practical problem solved

Uncertainty unintelligible-'caring measure'

Not given the ontology

epistemic problem

- Introduce axioms of rationality (Savage)
- to be satisfied in the rankings of preferences of a rational actor among quantum games.
- in accordance with which she acts as if employing standard Bayesian updating of priors, on given evidence.

quantitative problem

practical problem solved

Uncertainty unintelligible-'caring measure'

Not given the ontology

epistemic problem Page 100/143 solved Introduce axioms of rationality (Savage)

 to be satisfied in the rankings of preferences of a rational actor among quantum games.

 in accordance with which she acts as if employing standard Bayesian updating of priors, on given evidence.

•and (and this is the central point): this reasoning applies — equally to branching as to non-branching universes

quantitative problem

practical problem solved

†

Uncertainty unintelligible-'caring measure'

Not given the ontology problem Page 101/143 solved Introduce axioms of rationality (Savage)

 to be satisfied in the rankings of preferences of a rational actor among quantum games.

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quantitative problem

Given the ontology

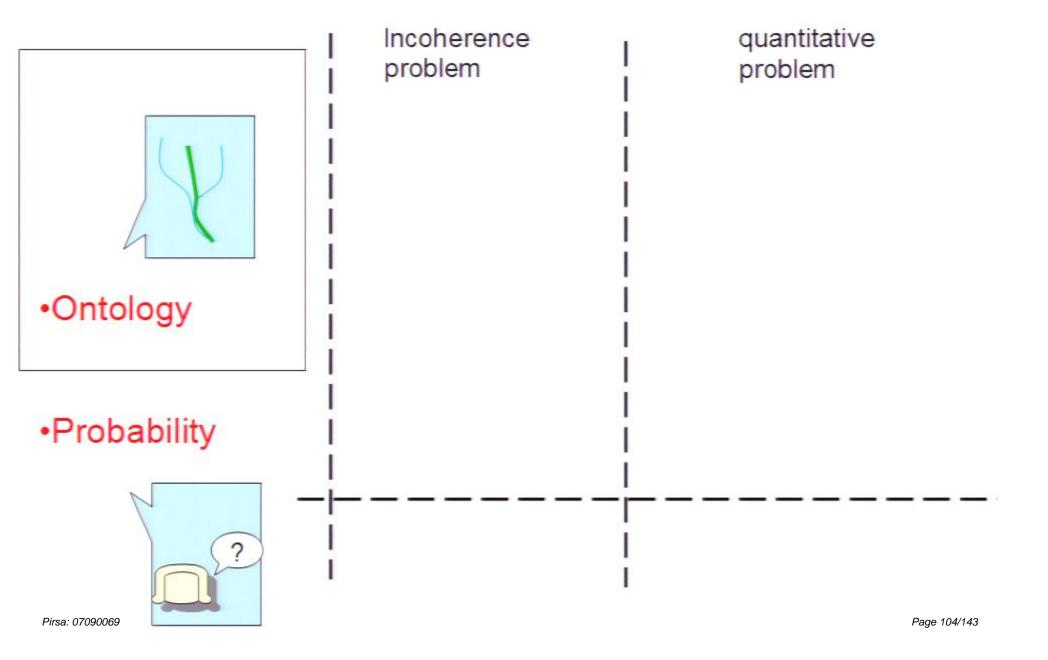
practical problem solved

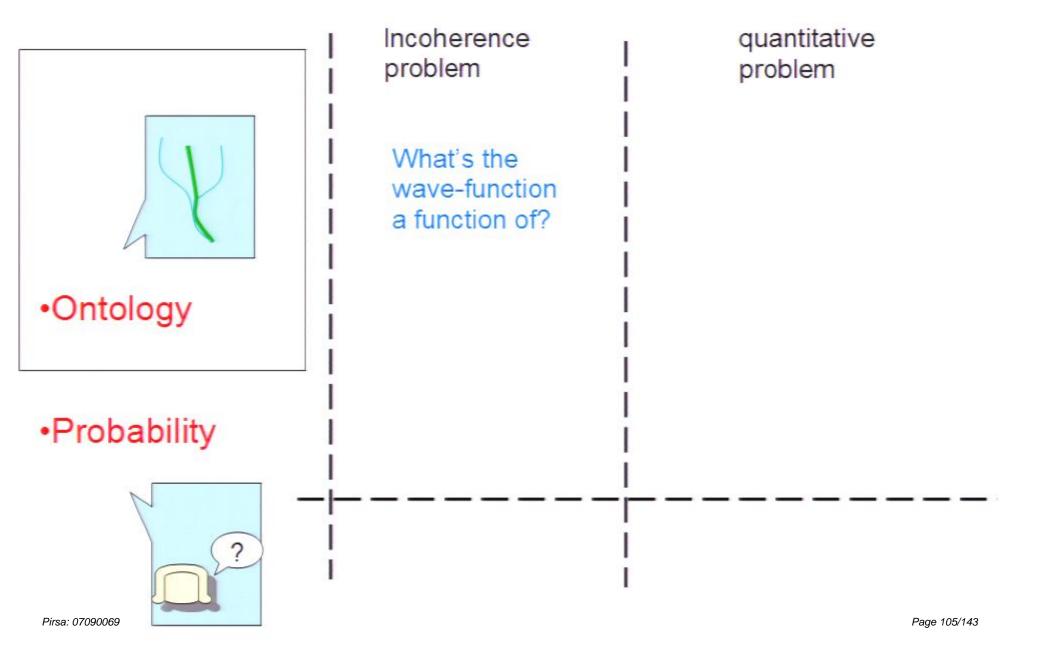
Uncertainty unintelligible-'caring measure'

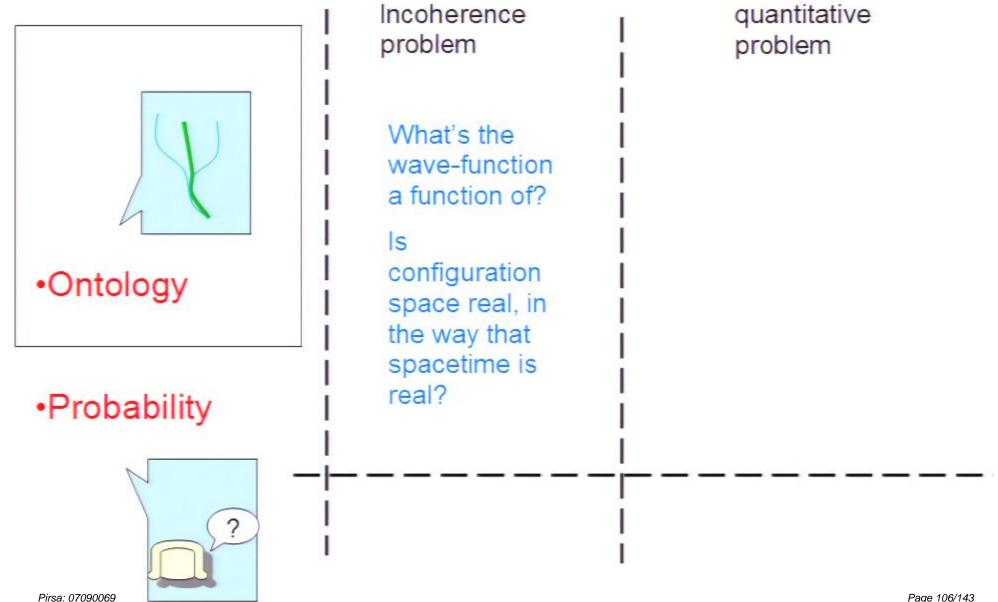
Not given the ontology

epistemic problem Page 102/143

Incoherence quantitative suppose problem problem solved Ontology Probability



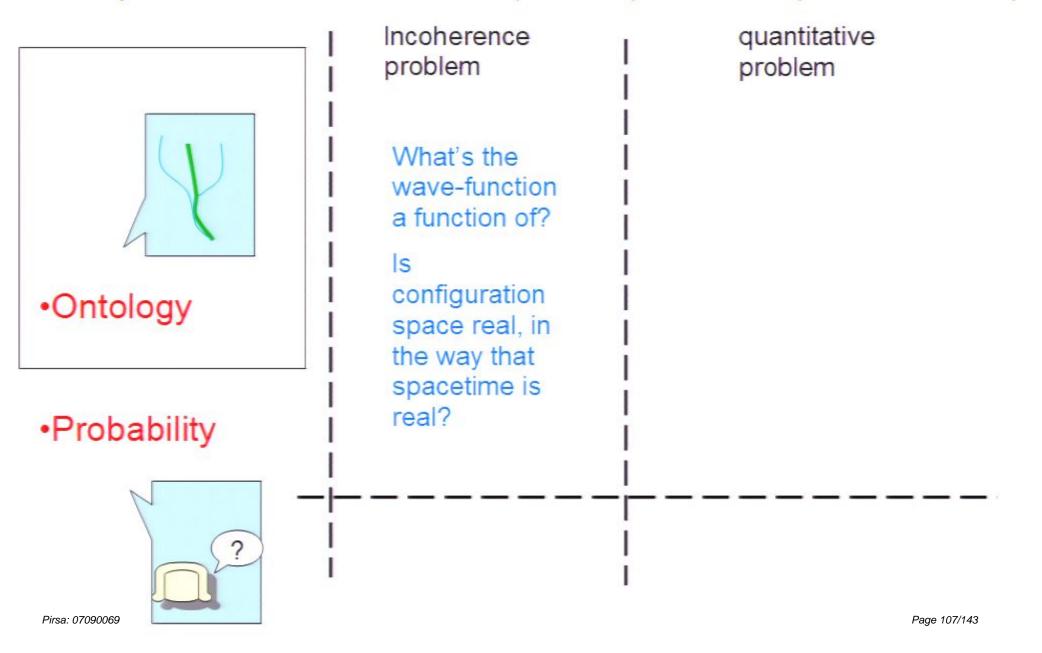




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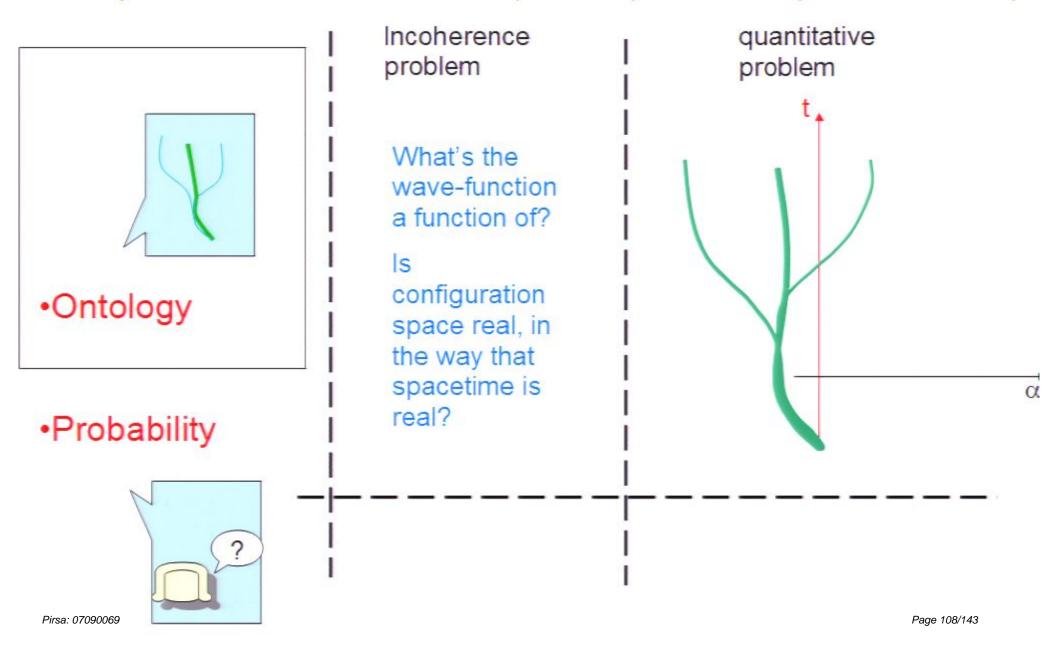
To solve the measurement problem...

Jniversality • no mention of measurement • no special interpretative assumption or additional eqs



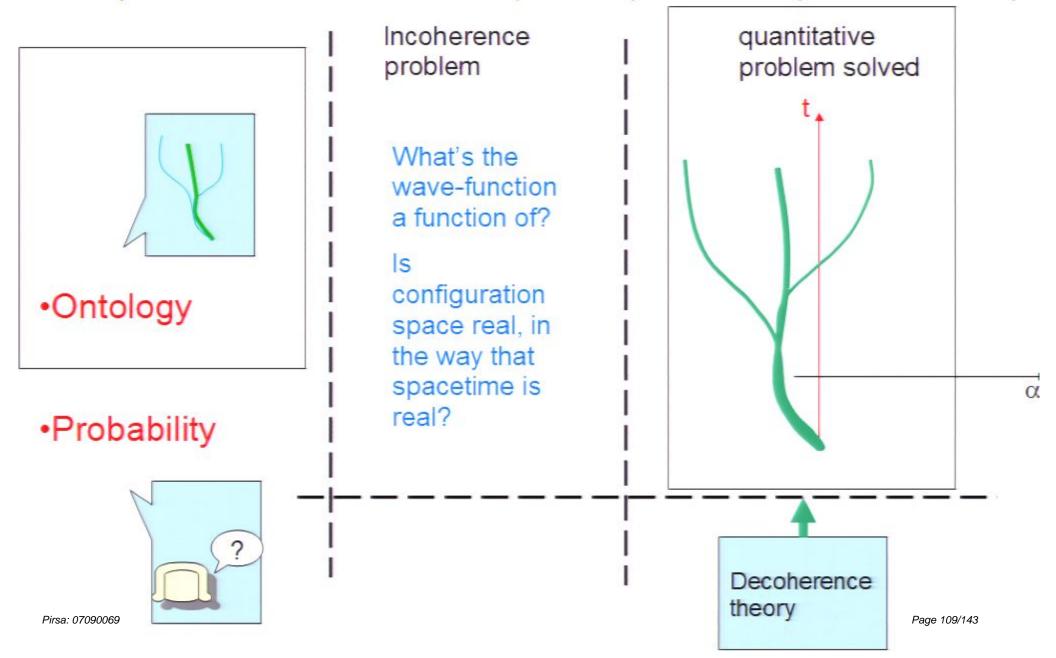
To solve the measurement problem...

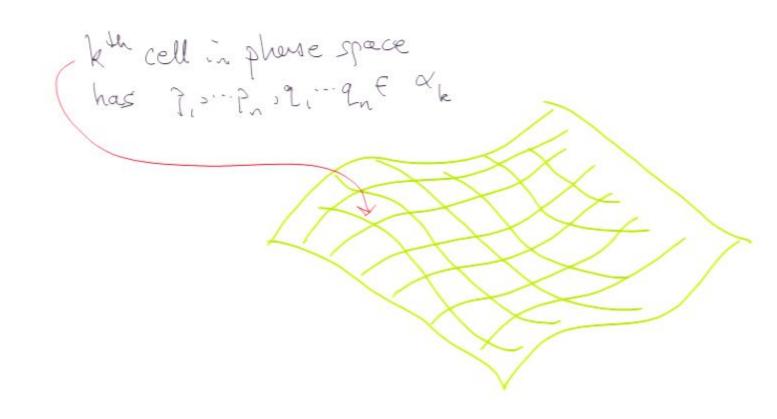
Jniversality • no mention of measurement • no special interpretative assumption or additional eqs



To solve the measurement problem...

Jniversality • no mention of measurement • no special interpretative assumption or additional eqs





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has 7, sup, 2, 2, 2, 2, 6 0 k Defre Heisenberg pultre PX(E) = EiHEP eiHE A 'history' is a sequence of such projection $C_{x} = P_{x}(t_{n}) \dots P_{x}(t_{n})$

Absolute probability
of a history x in state 4 is

Vy(x) = | Ca Y||²

A 'history' is a sequence of such projections $C_{\alpha} = P_{\alpha}(t_{\alpha}) \dots P_{\alpha}(t_{i})$

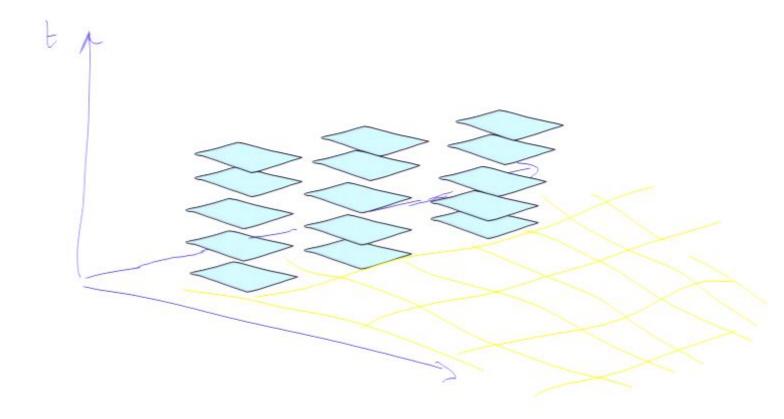
Conditional or relative probability

Wy (X/B) = 11 CzCR Y12

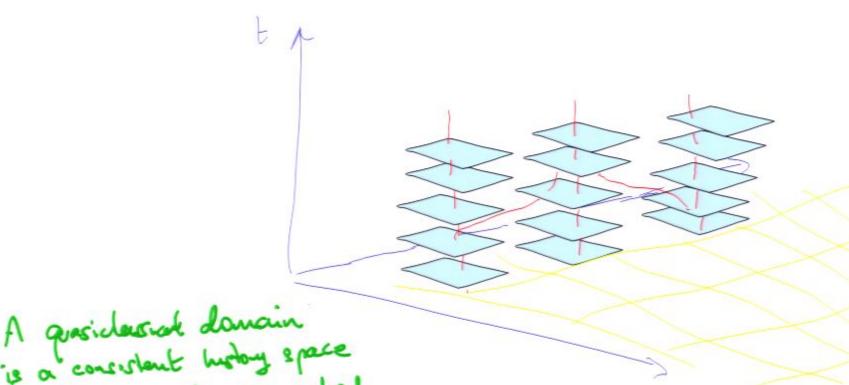
11 CBY12

Considerey as relations of the form: $\leq W(xididn)$ W(Xoxk) = satisfied if { Cx +} are offungered (Cx4, Cx4) = 0 Two stat experiments: W(K, df) + W(xix, af)

A consistent or decoherent history space requires a choice of projections at each time to smining to mity satisfying the confidency concertion



H consistent or decoherent history space requires a choice of projections at each time to smining to vivity satisfying the contribution



is a consistent history space and a state sharply peaked on histories degring approximately classial

A formulation of the expression $< C_{\alpha} \psi$, $C_{\beta} \psi >$ when the projections are coarse-grainings in q at time t:

$$\int (q'(t), q(t)) = \delta(q'_{5} \cdot l_{5}) \exp \left[-\frac{1}{L}(S(q) - q'_{2})\right] p(q'_{0}, q_{0})$$

$$\int (\Delta \alpha', \Delta \lambda) = \int S q' \int dq$$

$$q' \in \Delta \lambda' \quad q \in \Delta \lambda$$

$$= \int S q' \int S q \quad \delta(q'_{5} - q'_{5}) \exp \left[-\frac{1}{L}(S_{free}(q') - S_{free}(q) + M_{q'_{1}}, q'_{2})\right]$$
where
$$N(q', q) = -M_{Y} \int dt \quad (q', q' - q'_{2} - q'_{2} - q'_{2} - q'_{2} - q'_{2})$$

$$+ i_{2}MY_{k}T \int dt \quad [q'(t) - q'_{2}, q'_{2}]^{2}$$

1) ecoherence time ton I for small molecules thermal me ~ small molecules - J Sq' J Sq S (9'_5-9'_5) exp [in (Sfree 9') - Sfree 2)4 - Mg f dt (2'9'-92-92-12) + 12Mg kt fat [2'(1)-2(1)]

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evangles:

Brown an motion (density matrix approach)

Joss + Zeh 1985

(quantum state defesion)

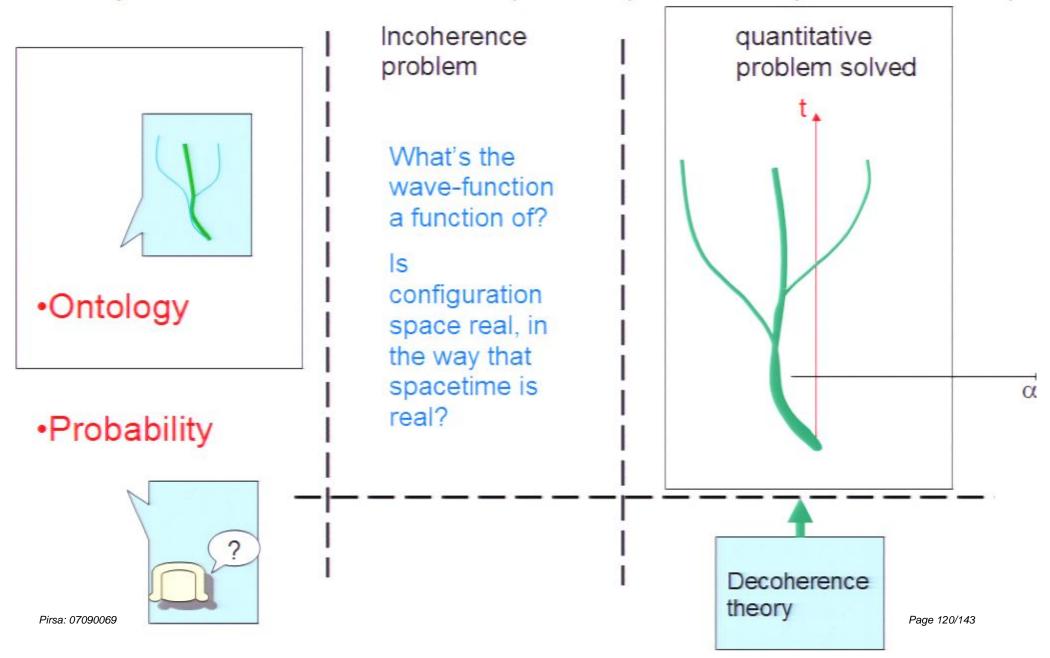
Zurek 1994, 2001,

Hallivell + Zapas 1997.

Difrsion in fluide + gases (hydrodynamic venilles) Hallivell 1998.

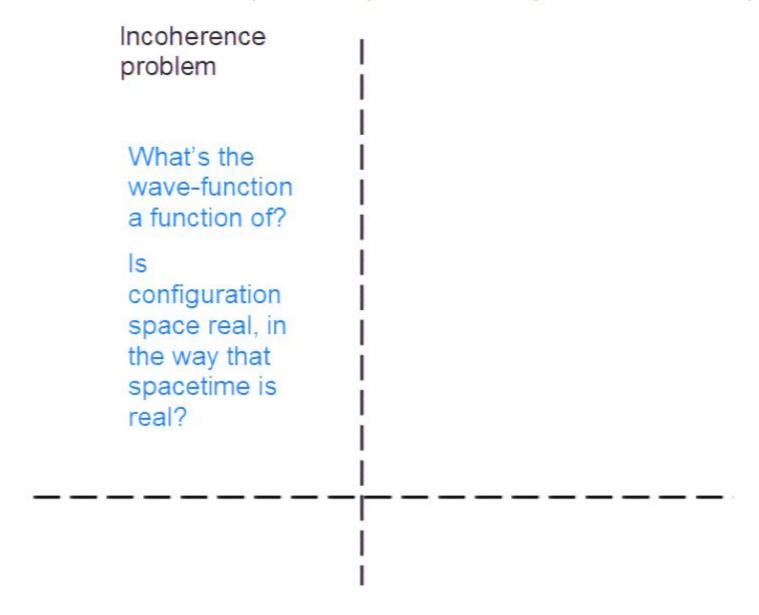
To solve the measurement problem...

Jniversality • no mention of measurement • no special interpretative assumption or additional eqs



To solve the measurement problem...

Jniversality • no mention of measurement • no special interpretative assumption or additional eqs



Conclitional or relative probability

Wy (X/B) = 11 CzCR Y12

11 CpY1/2

Conditional or relative correlations

Wy (X/B) = || C_2C_B Y||^2

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$$\psi \otimes \eta_{,} \otimes (C_{+} \varphi_{+} + C_{-} \varphi_{-})$$

$$\longrightarrow \psi \otimes (C_{+} \eta_{+} \otimes \varphi_{+} + C_{-} \eta_{-} \otimes \varphi_{-})$$

$$\longrightarrow C_{+} \psi_{+} \otimes \eta_{+} \otimes \varphi_{+} + C_{-} \psi_{-} \otimes \varphi_{-}$$

$$\uparrow \text{ Relative state' af } \eta_{+} \text{ is } \varphi_{+}$$

$$\uparrow \eta_{-} \text{ is } \varphi_{-}$$

Conditional or relative correlations W (X/B) = 1 CLCB 41/2 11 CBY 1/2 Define relation Detp us: Defp(xn/xi) - y (du/di) = 1 "Configuration Xx is value definite relative to di

Pirsa: 07090069

in state p"

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Conclinates

(as 7, 5.1. 2, 1.

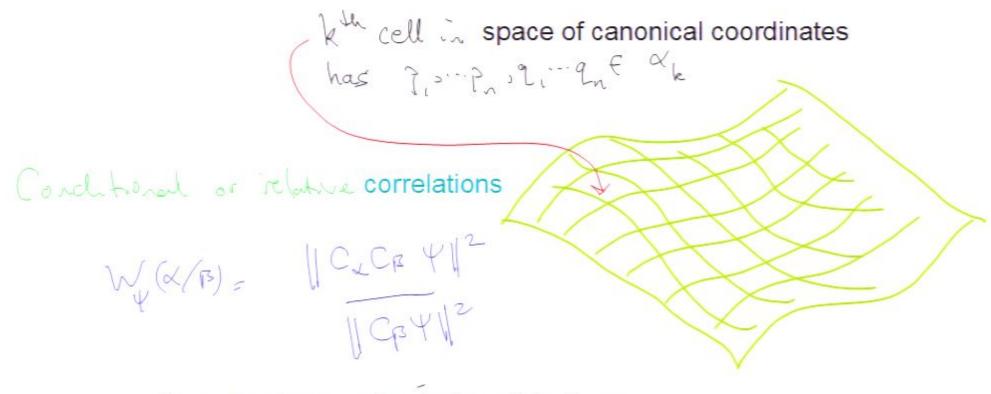
 Def_{ψ} is transitive for spacelike projectors and for consistent timelke projectors

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Conditional or relative correlations

Wy (X/B) = ||C_2C_B Y||^2

In a quantum mechanical particle theory:



In a quantum mechanical particle theory:

 All correlations among patterns of particles (relative positions and momenta) exist

Conditional or relative correlations

Wy (X/F) = ||C_2C_F Y||^2

||C_3Y||^2

In a quantum mechanical particle theory:

 More generally: all correlations among patterns of polarizations on n-particle phase space exist

Conditional or relative correlations

We (all in space of canonical coordinates has 7,5...p., 2,...2nf alk

Conditional or relative correlations

We (all in space of canonical coordinates

has 7,5...p., 2,...2nf alk

I Conditional or relative correlations

In a quantum mechanical particle theory:

- More rigorously: all correlations among patterns of polarizations on n-particle phase space exist
- as given by invariant relations among projections

Conditional or relative correlations

We (x/B) = ||C_2C_B Y||^2

In a quantum mechanical particle theory:

- More rigorously: all correlations among patterns of polarizations on n-particle phase space exist
- as given by invariant relations among projections
- but only some of these correlations among patterns are stable under perturbations (typically coarsegrainings), yielding dynamically decoupled (branching) histories.

Conditional or relative correlations

Wy (x/B) = ||C_2C_B Y||^2

In a quantum mechanical field theory:

- More rigorously: all correlations among field configurations (polarizations of phase space) exist
- as given by invariant relations among projections
- but only some of these correlations of field configurations are stable under perturbations (typically coarse-grainings), yielding dynamically decoupled (branching) histories.

Conditional or relative correlations

Wy (x/B) = || CyCBY||^2

In classical general relativity:

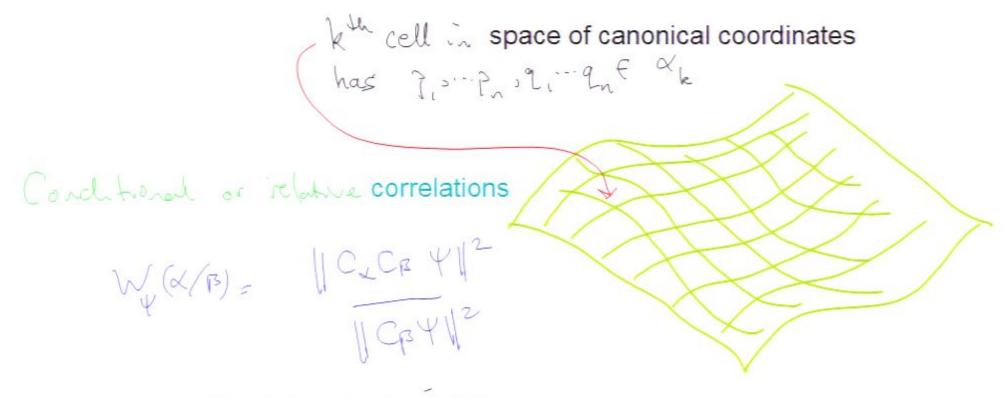
- All correlations among values of fields on the manifold exist
- as given by diffeomorphic-invariant relations among them
- but only some of these correlations of field configurations are stable under perturbations (typically coarse-grainings), yielding dynamically decoupled (but non branching) histories

Conditional or relative correlations

We (x/B) = ||C_LC_B Y||^2

In classical general relativity:

·...like galaxies, stars, planets



In classical general relativity:

 ...patterns of values of fields can be referred to any foliation, but some foliations are more perspicuous than others (simpler, better 'adapted', etc.)

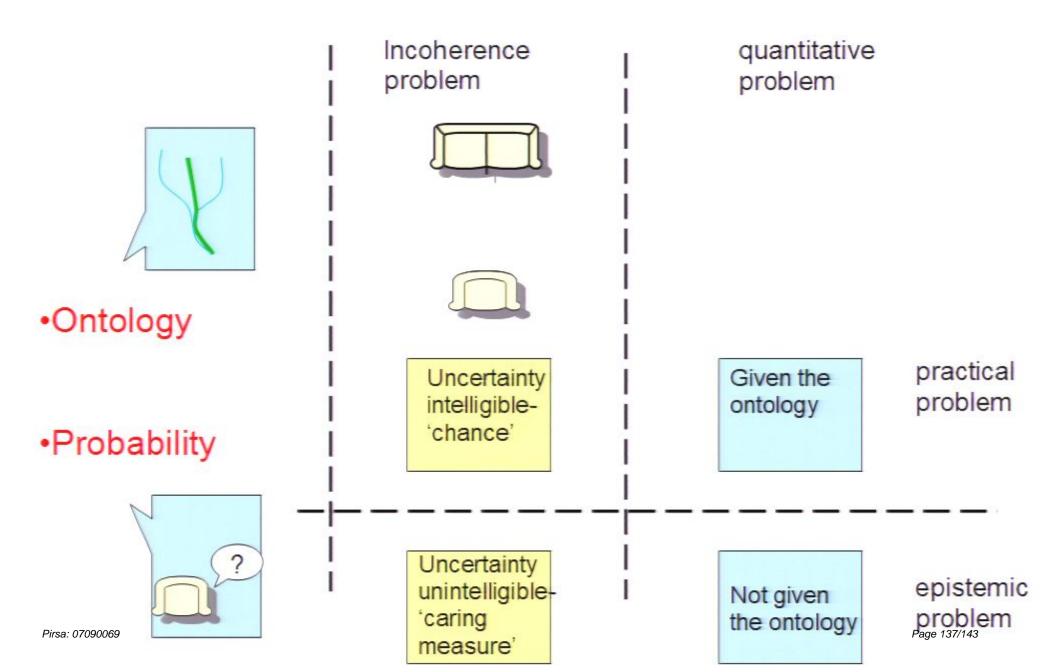
Pirsa: 07090069 Page 135/143

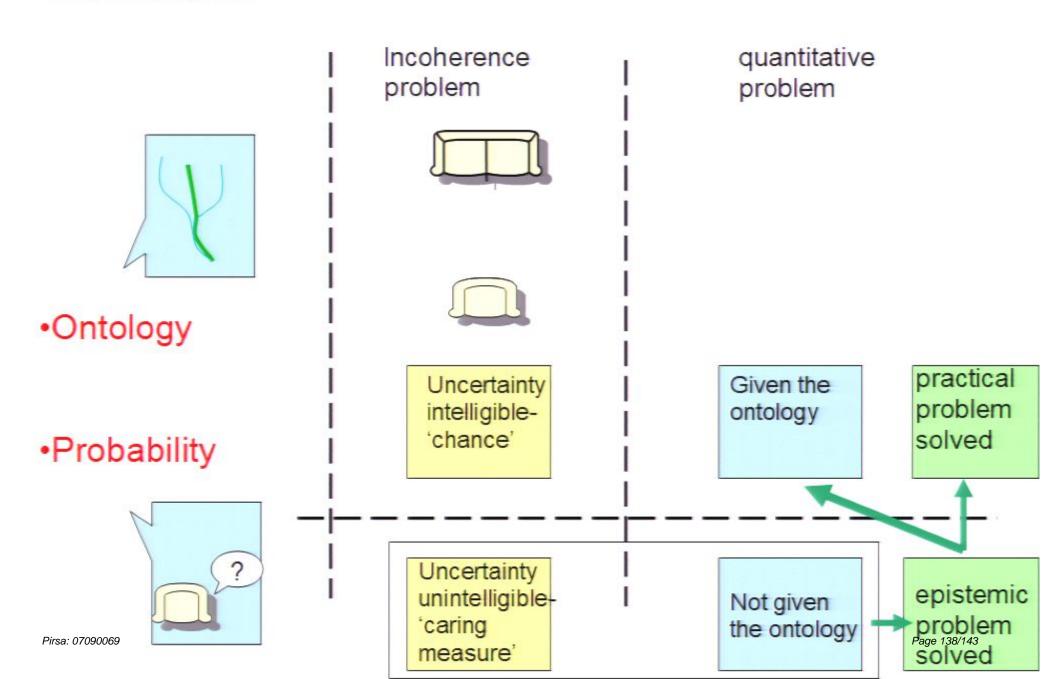
Conclutional or relative correlations

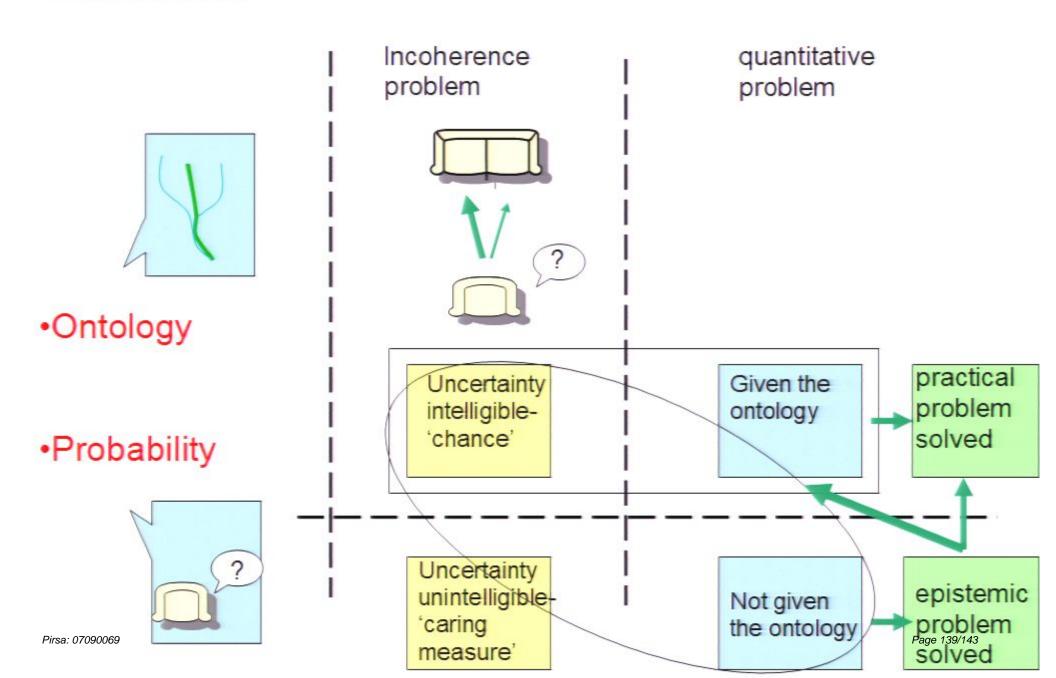
Wy (X/B) = || C_2C_B Y||^2

In classical general relativity:

- ...patterns of values of fields can be referred to any foliation, but some foliations are more perspicuous than others (simpler, better 'adapted', etc.)
- but only if ALL (correlations among field-values on) leaves of a foliation are real.

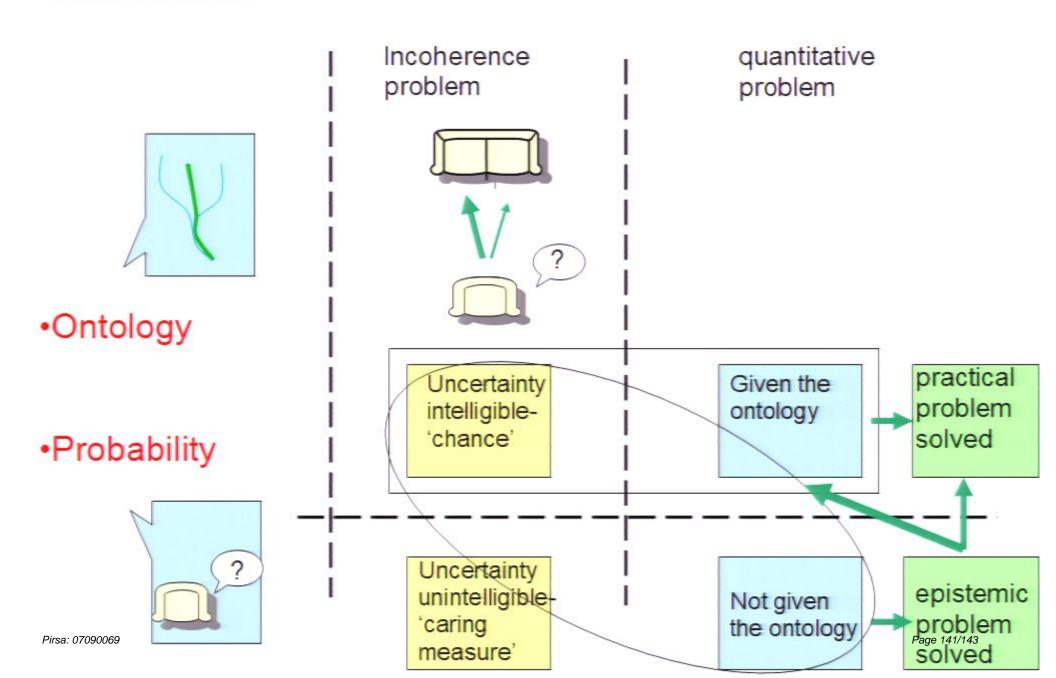






End of slide show, click to exit.

Pirsa: 07090069 Page 140/14.



Conditional or relative correlations

Wy (x/F) = ||C_2C_F Y||^2

In a quantum mechanical particle theory:

- More rigorously: all correlations among patterns of polarizations on n-particle phase space exist
- as given by invariant relations among projections
- but only some of these correlations among patterns are stable under perturbations (typically coarsegrainings), yielding dynamically decoupled (branching) histories.

Conditional or relative correlations

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