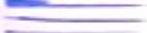


Title: K-Essence, superluminal propagation, causality and emergent geometry

Date: Sep 14, 2007 10:00 AM

URL: <http://pirsa.org/07090029>

Abstract: The k-essence theories admit the superluminal propagation of the perturbations on classical nontrivial backgrounds. In this talk I will review our arguments from arXiv:0708.0561v1 and show that in spite of the superluminal propagation the causal paradoxes do not arise in these theories and in this respect they are not less safe than General Relativity.

 
 **κ -essence,
SUPERLUMINAL
PROPAGATION,
CAUSALITY**

*and
emergent geometry*

A. Vikman $\left\{ \begin{array}{l} NYU-CCPP \\ LMU-ASC \end{array} \right\}$

based on arXiv:0708.0561
in collaboration with : V. Mukhanov
E. Babichev

● GENERAL formalism

S_C

Action: $S_\phi = \int d^4x \sqrt{-g} \tilde{P}(\phi, X)$

where $X = \frac{1}{2} \nabla_\mu \phi \nabla^\mu \phi$ (+, sign - ---)

EMT

$$T_{\mu\nu} = P_{,X} \nabla_\mu \phi \nabla_\nu \phi - g_{\mu\nu} P$$

$$NEC: T_{\mu\nu} n^\mu n^\nu \geq 0 \leftrightarrow P_{,X} \geq 0$$

Hydrodynamics: (if $X > 0$)

$$P = P(\phi, X), \quad \text{pressure}$$

$$\epsilon = 2X P_{,X} - P, \quad \text{energy dens.}$$

$$U_\mu = \frac{\nabla_\mu \phi}{\sqrt{2X}}, \quad \text{4 velocity}$$

$$C_s^2 = \frac{P_{,X}}{\epsilon_{,X}}, \quad \text{sound speed}$$

[Mukhanov
Garriga
99]

● GENERAL formalism

Sca.

Action: $S_\phi = \int d^4x \sqrt{-g} \tilde{P}(\phi, X)$

where $X = \frac{1}{2} \nabla_\mu \phi \nabla^\mu \phi$ (+, - ---)

EMT

$$T_{\mu\nu} = P_{,X} \nabla_\mu \phi \nabla_\nu \phi - g_{\mu\nu} P$$

NEC: $T_{\mu\nu} n^\mu n^\nu \geq 0 \leftrightarrow P_{,X} \geq 0$

Hydrodynamics: (if $X > 0$)

$$P = P(\phi, X), \quad \text{pressure}$$

$$\epsilon = 2X P_{,X} - P, \quad \text{energy dens.}$$

$$U_\mu = \frac{\nabla_\mu \phi}{\sqrt{2X}}, \quad \text{4 velocity}$$

$$C_s^2 = \frac{P_{,X}}{\epsilon_{,X}}, \quad \text{sound speed}$$

[Mukhanov
Garriga
99]

EoM:

$$6 \quad v_{\mu} v^{\nu} + \frac{1}{P_{,x}} = 0$$

where

$$G^{\mu\nu} = g^{\mu\nu} + \frac{P_{,xx}}{P_{,x}} \nabla^{\mu} \phi \nabla^{\nu} \phi$$

EoM is hyperbolic if

$$1 + \frac{P_{,xx}}{P_{,x}} 2X > 0$$

$$C_s^2 > 0$$

Mukhanov, Garriga

$$\left[C_s^2 = \left[1 + \frac{P_{,xx}}{P_{,x}} 2X \right]^{-1} \right]$$

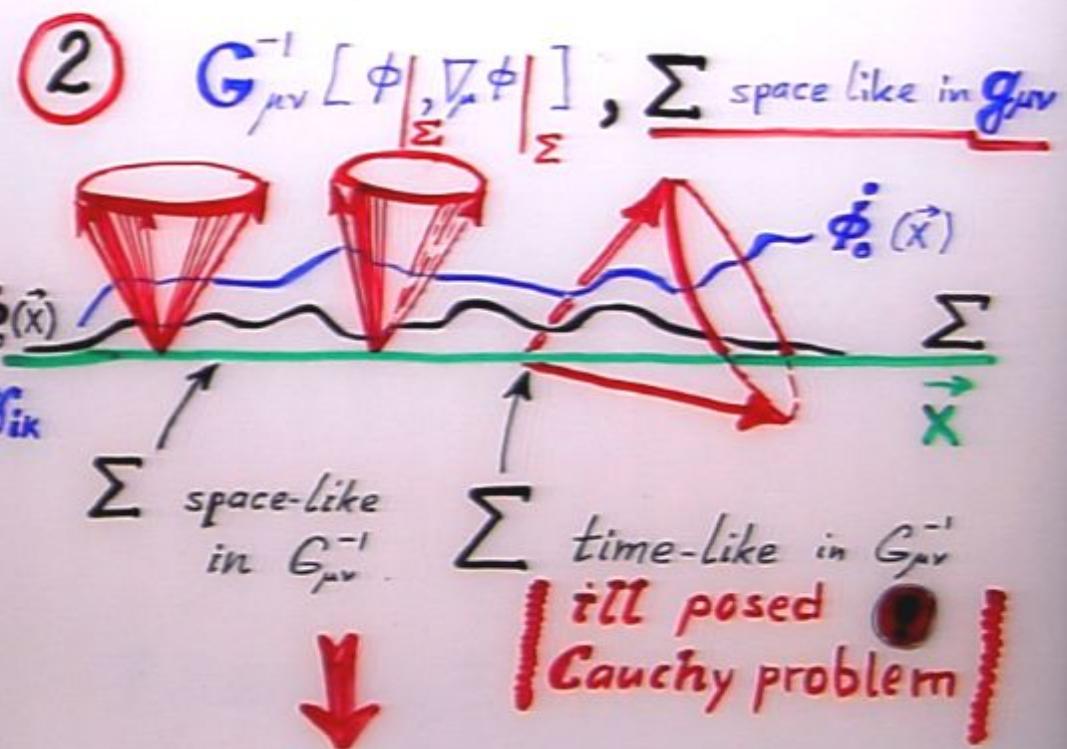
Aharanov, Komar,
Susskind, 69

Armendaric-Picon, Lim
05

Rendall 05

I VOL all initial data are allowed for well posed Cauchy problem.

① $1 + 2X \frac{P_{,XX}}{P_{,X}} > 0$ Hyperbolicity



allowed $(\Phi_0(\vec{x}), \dot{\Phi}_0(\vec{x})) : \sum$
space-like in $G_{\mu\nu}^{-1}$ and $g_{\mu\nu}$.

"Good initial data"

" = "

Well posed Cauchy problem

" = " !

" Small " derivatives (+ Hyperbol.)

$$1 + c_s^2 (\vec{\nabla} \phi_0)^2 \frac{L_{,XX}}{L_{,X}} > 0$$

$$(\vec{\nabla} \phi_0)^2 = g^{ik} \partial_i \phi_0 \partial_k \phi_0 \quad || \quad \bullet$$

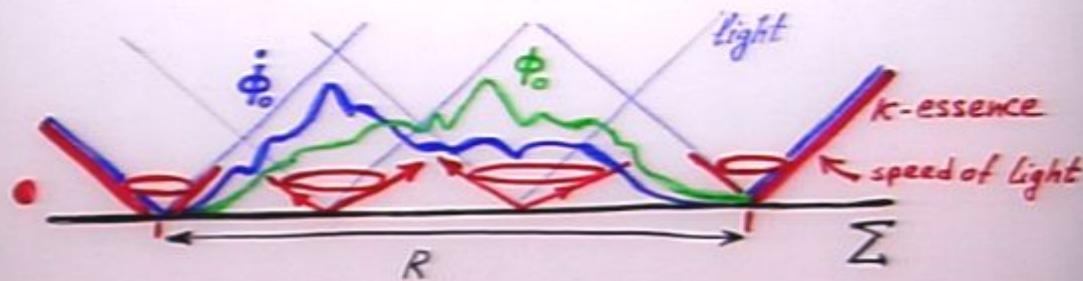
$$c_s^2 \left(1 + \dot{\phi}_0^2 \frac{L_{,XX}}{L_{,X}} \right) > 0$$

$$C_s^2 = \left(1 + 2X \frac{L_{,XX}}{L_{,X}} \right)^{-1} > 0 \bullet$$

Hyperb.

J_S $g_{\mu\nu}$

- $\mathcal{L}(X, \phi) \simeq V(\phi) + K_1(\phi)X + K_2(\phi)X^2 + \dots$
for $X \approx 0$



- if $K_1 = 0 \rightarrow c_s^2 = \frac{1}{1+2(n-1)} < 1$

- For $\mathcal{L}(X, \phi) = \mathcal{L}(X)$

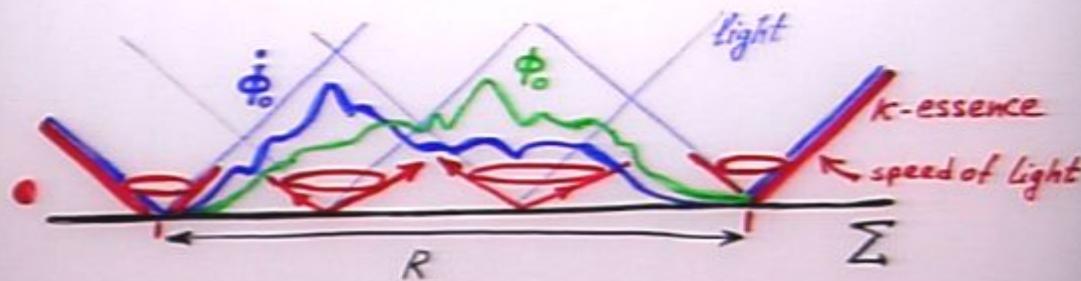
$$\exists \underline{\phi = f(x \pm t)}$$



Causal Limit for
Localized configurations
is given by $g_{\mu\nu}$!

$$\underline{JS} \quad \underline{g_{\mu\nu}}$$

- $\mathcal{L}(X, \phi) \simeq V(\phi) + K_1(\phi)X + K_2(\phi)X^2 + \dots$
for $X \approx 0$



- if $K_1 = 0 \rightarrow c_s^2 = \frac{1}{1+2(n-1)} < 1$
- For $\mathcal{L}(X, \phi) = \mathcal{L}(X)$

$$\exists \quad \underline{\underline{\phi = f(x \pm t)}}$$



Causal Limit for
Localized configurations
is given by $g_{\mu\nu}$!

• π is the variables of

and emergent
geometry.

- In which metric do π 's around $\phi(x)$ "live"?

- EoM for π 's in $\tilde{G}_{\mu\nu}^{-1}$ (test field appr.)
 $\tilde{G}_{\mu\nu}^{-1}$ generally covariant
- $\tilde{G}_{\mu\nu}^{-1}$ can be only conf. transf. of $G_{\mu\nu}^{-1}$

$$\tilde{G}_{\mu\nu}^{-1} = \frac{\lambda_{,X}}{c_s} \left[g^{\mu\nu} - c_s^2 \frac{\lambda_{,XX}}{\lambda_{,X}} \nabla_\mu \phi \nabla_\nu \phi \right]$$

$$\tilde{G}^{\mu\nu} D_\mu D_\nu \pi + M_{\text{eff}}^2 \pi = 0$$

$$D_\mu \tilde{G}_{\alpha\beta}^{-1} = 0 \quad \pi \text{ causal in } \tilde{G}^{\mu\nu}$$

$$M_{\text{eff}}^2 = \frac{c_s}{\lambda_{,X}} \left(E_{,\phi\phi} + (\lambda_{,X} G^{\mu\nu})_{,\phi} \nabla_\mu \nabla_\nu \phi_0 \right)$$

Legendre transform and emergent geometry.

- In which metric do Π 's around $\Phi(x^\mu)$ "live"?

"(test field appr.)

- EoM for Π 's generally covariant

$$\tilde{G}_{\mu\nu}^{-1}$$

- $\tilde{G}_{\mu\nu}^{-1}$ can be only conf. transf.
of $G_{\mu\nu}^{-1}$

$$\tilde{G}_{\mu\nu}^{-1} = \frac{Z_{,X}}{C_s} \left[g^{\mu\nu} - C_s^2 \frac{Z_{,XX}}{Z_{,X}} \nabla_\mu \phi \nabla_\nu \phi \right]$$

$$\tilde{G}^{\mu\nu} D_\mu D_\nu \Pi + M_{\text{eff}}^2 \Pi = 0$$

$$D_\mu \tilde{G}_{\alpha\beta}^{-1} = 0 \quad \Pi \text{ causal in } \tilde{G}^{\mu\nu}$$

$$M_{\text{eff}}^2 = \frac{C_s}{Z_{,X}} \left(\mathcal{E}_{,\phi\phi} + (Z_{,X} G^{\mu\nu})_{,\phi} \nabla_\mu \nabla_\nu \phi_0 \right)$$

$$S^2 = \frac{1}{2} \int d^4x \sqrt{\tilde{G}} \left[\tilde{G}^{\mu\nu} \partial_\mu \pi \partial_\nu \pi - M_{\text{eff}}^2 \pi^2 \right]$$

also for general cosmological perturbations

different M_{eff}^2 but the same $\tilde{G}^{\mu\nu}$
for $\pi \equiv \delta\phi + \frac{\dot{\phi}}{H} \psi$

in cosmology (flat Friedmann Universe)

$$\nu \equiv \sqrt{\epsilon_{,X}} a \pi$$

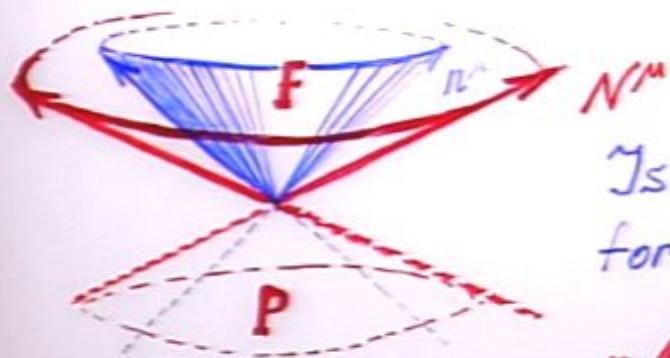
$$z = \frac{\dot{\phi}}{H} \sqrt{\epsilon_{,X}} a$$

$$S^2 = \frac{1}{2} \int d\eta d^3X \left(\nu'^2 - c_s^2 (\vec{\nabla} \nu)^2 + \frac{z''}{z} \right)$$

(Garriga, Mukhanov
1999)

$S^2 \rightarrow$ Quantization

Chronology Protection



Is this a problem
for causality?

„No!“ / Global properties

Example: Friedmann universe

ϕ - cosmological scalar : $\phi(t)$

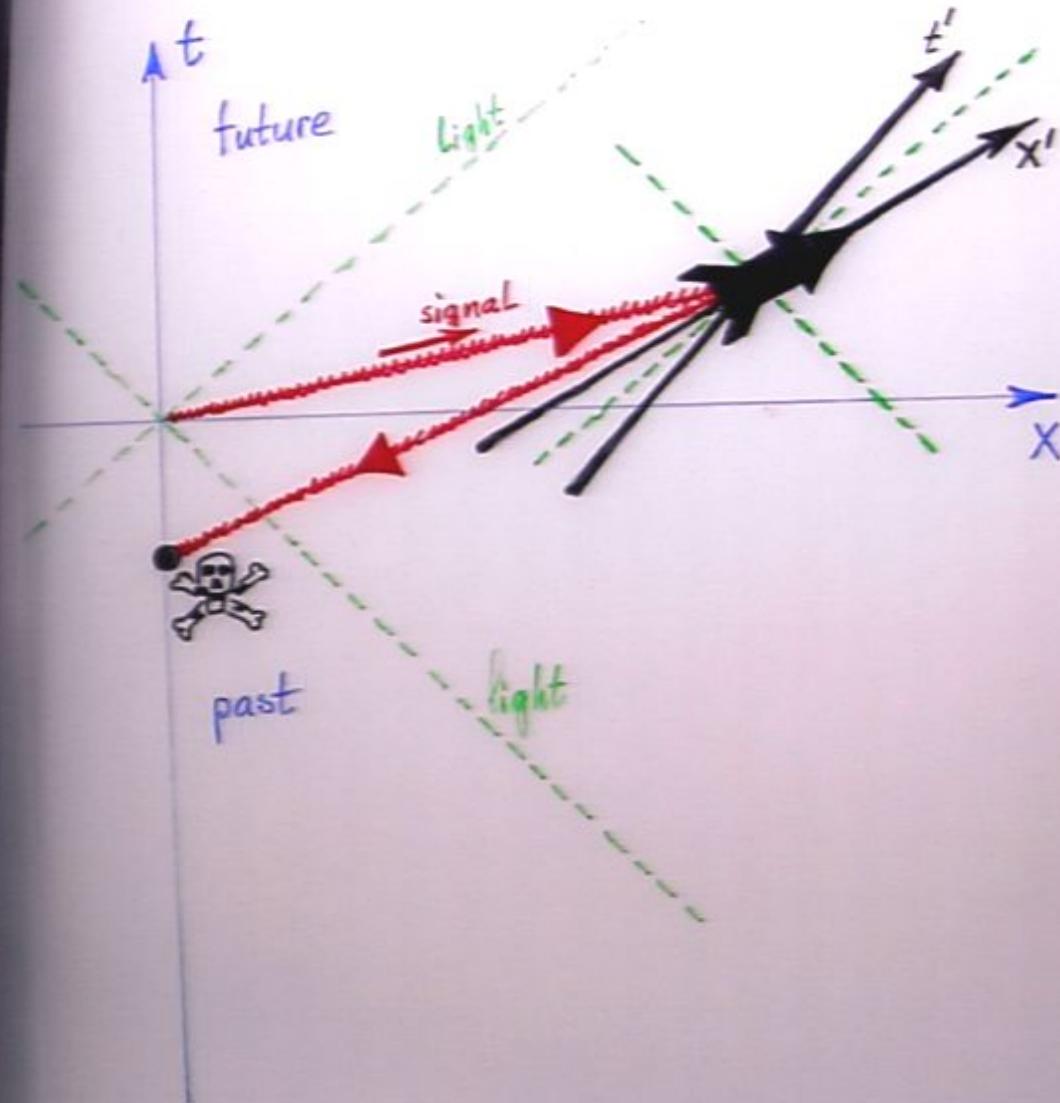
$$ds^2 = dt^2 - \alpha^2(t) d\vec{x}^2 = g_{\mu\nu} dx^\mu dx^\nu$$

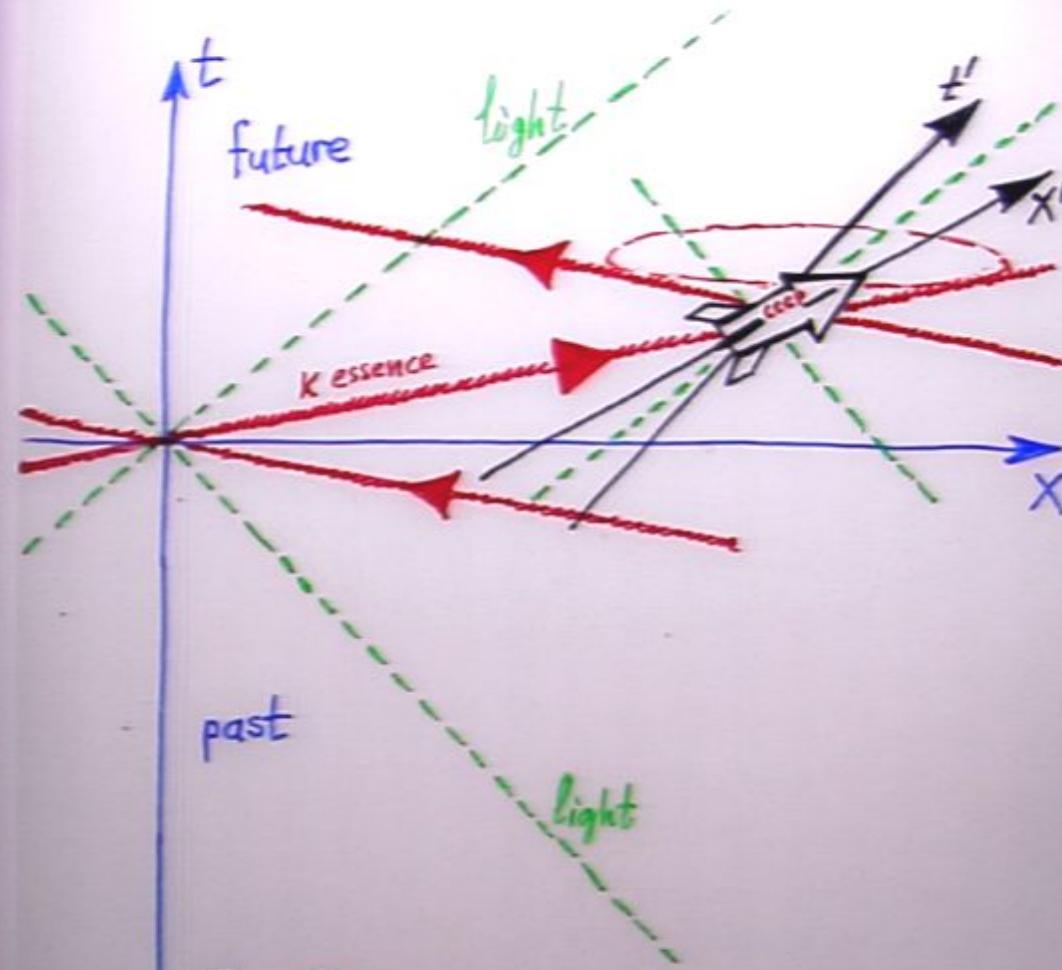
$$d\ell^2 = G_{\mu\nu} dx^\mu dx^\nu = ds^2 - \\ - c_s^2 \frac{p_{,xx}}{p_{,x}} 2X dt^2 \rightarrow$$

$$d\ell^2 = c_s^2(t) dt^2 - \alpha^2(t) d\vec{x}^2$$

$$dt = c_s dt \rightarrow d\ell^2 \text{ - „Friedmann“}$$

$$d\ell^2 = d\tau^2 - \alpha^2(\tau) d\vec{x}^2$$





exists Global time $f \rightarrow$
stable causality

exists diff. f : $\nabla^M f$ is future directed
 timelike vector field. \rightarrow No CCC in
 "the most interesting" backgrounds.

- Conclusion

Superluminal propagation
does not lead to causal
paradoxes in K-essence.

Thanks for attention!

- Conclusion

Superluminal propagation
does not lead to causal
paradoxes in κ -essence.

Thanks for attention!