

Title: Wilson Loops as D-branes

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Abstract:

## PLAN

WILSON LOOPS - INTRODUCTION  $U(1)$ ,  $SU(N)$

AdS/CFT - INTRODUCTION

-  $\mathcal{N}=4$  SYM

-  $\frac{1}{2}$  BPS WILSON LOOPS

HOLOGRAPHIC W.L. - D5, D3 PROBES

- LOW ENERGY  
FIELD THEORY

## WILSON LOOPS: NON LOCAL AND GAUGE INVARIANT OPERATORS

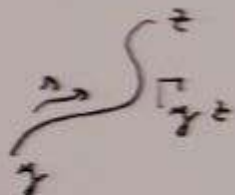
U(1) GAUGE THEORY  $[A_\mu(x)]$

GAUGE TRANSFORMATION:  $A_\mu(x) \xrightarrow{G.T.} A_\mu(x) + \partial_\mu \alpha(x)$

PHASE FACTOR

$$U[\Gamma_{yz}] = \text{EXP} \left[ iq \int_{\Gamma_{yz}} dx^\mu A_\mu(x) \right]$$

$$U[\Gamma_{yz}] \xrightarrow{G.T.} e^{-iq\alpha(y)} U[\Gamma_{yz}] e^{iq\alpha(z)}$$



PARALLEL  
TRANSPORTER

$$\text{IF } \Gamma = C$$



A "LOOP"

$$\longrightarrow W_q(C) = \text{EXP} \left[ i q \int_C dx^\mu A_\mu(x) \right]$$

GAUGE INVARIANT OPERATOR

- $q$  CHARGE [REPRESENTATION OF  $U(1)$ ]
- $C$  CLOSED CURVE

DESCRIBE THE COUPLING OF THE  $U(1)$   
GAUGE THEORY WITH A NON DYNAMICAL  
CHARGED PARTICLE.

## U(N) YANG-MILLS

$$W_R(C) = \text{Tr}_R \mathcal{P} \exp \left[ i \int_C ds A_\mu(x) \dot{x}^\mu \right]$$

- $\mathcal{P}$  PATH ORDERING ( $A_\mu = A_\mu^a \tau^a$ )
- $\text{Tr}$  TRACE ( $\tau^a$  GENERATORS OF  $U(N)$ )

$W-L$  IS CHARACTERIZED BY

$C$  A CLOSED CURVE

$R$  REPRESENTATION OF  $U(N)$

$R$  IS SUMMARIZED BY A YOUNG TABLEAU

$$R = (m_1, m_2, \dots, m_i, \dots, m_N)$$

$$N \left\{ \begin{array}{cccc} 1 & 2 & \dots & m_1 \\ 1 & 2 & \dots & m_2 \\ \dots & \dots & \dots & \dots \\ m_N & & & \end{array} \right.$$

BOXES IN THE  $i$ -TH ROW

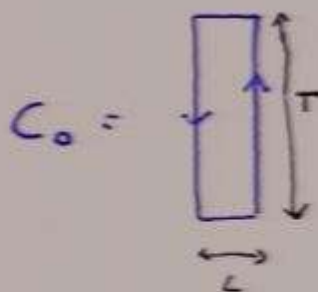
$$m_1 \geq m_2 \geq \dots \geq m_N$$

$W_R(\mathcal{C})$  DESCRIBE THE COUPLING OF  $U(N)$  YANG-MILLS WITH AN EXTERNAL NON DYNAMICAL PARTICLE WITH CHARGE  $R$  5

E.G. THE EXTERNAL PARTICLE IS A QUARK

$\rightarrow R = \square$  FUNDAMENTAL

$\langle W_{\square}(\mathcal{C}_0) \rangle = A(L) e^{-T E(L)}$  QUARK-ANTIQUARK POTENTIAL



$T \rightarrow +\infty$



$W-L$  : ORDER PARAMETER FOR CONFINEMENT

# WHY HOLOGRAPHIC?

6

## AdS/CFT DUALITY

STATEMENT:

TYPE IIB STRING  
ON  $AdS_5 \times S^5$   
WITH  $\int_{S^5} F_5 = N$

EQUIVALENT

$\Leftrightarrow$   $\mathcal{N} = 4$  SYM  
ON  $\mathbb{R}^{3,1}$

PARAMETERS:

RADIUS:  $L^4 = 4\pi\alpha' N^2$

$$\frac{L^2}{\alpha'} \sim N$$

$$g_s = g_{YM}^2$$

Y.M. COUPLING:  $g_{YM}$

$$\text{t'HOOFT: } \lambda = g_{YM}^2 N$$

$N \rightarrow +\infty$  LIMIT

NON INTERACTING  
STRING THEORY

PLANAR GRAPHS

$\lambda \gg 1$  NO  $\alpha'$   
CORRECTIONS

$\lambda \ll 1$  PERTURBATIVE  
REGIME

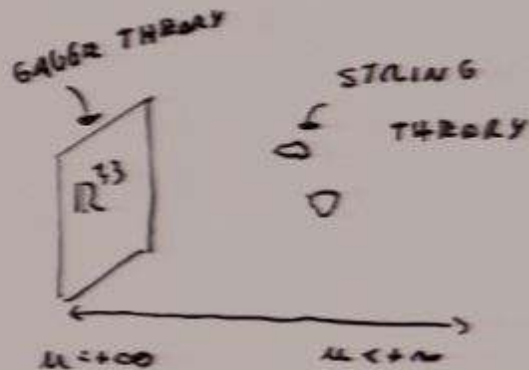


## HOLOGRAPHY

$$d^2_{\text{SAdS}_5\text{S}^5} = L^2 \left( \mu^2 \gamma_{\mu\nu} dx^\mu dx^\nu + \frac{d\mu^2}{\mu^2} \right) + L^2 (d\theta^2 + \sin^2\theta \rho(\mathbb{R}_4^2))$$

$\mu \rightarrow +\infty$  BOUNDARY OF THE SPACE  
 $\mathbb{R}^{1,3}$

THE PHYSICS OF THE BULK RESULTS  
 FROM A HOLOGRAPHIC IMAGE  
 OF THE PHYSICS ON THE BOUNDARY



# GRAVITATIONAL DESCRIPTION OF WILSON LOOPS

REY, YEE  
MALDACENA



$W_{\square}(C)$  IS ASSOCIATED TO A FUNDAMENTAL STRING IN THE BULK ENDING ON C AT THE BOUNDARY OF  $AdS$

→ STRING ACTION, PROPORTIONAL TO THE SURFACE OF THE WORLD-SHEET

$$\langle W_{\square}(C) \rangle \sim e^{-S}$$

$$\langle W_{\square}(C) \rangle = \int \mathcal{D}X e^{-\int \sqrt{-g} S(X)}$$

FOR NON-CONFORMAL THEORY

$$E(L) \propto L$$

CONFINEMENT

$\mathcal{N}=4$  SYM  $[A_\mu, \phi^I, \psi]$

W.L. DESCRIBE COUPLING OF A  
SUPERPARTICLE TO SYM

→ BOSONIC WILSON LOOPS

$$C \equiv (X^\mu(2), \gamma^I(1))$$

$$W_R(C) = \text{Tr}_R \mathcal{P} \exp \left[ i \int_C ds (A_\mu \dot{x}^\mu + \phi_I \dot{\gamma}^I) \right]$$

$\mathcal{N}=4$  SYM INVARIANT UNDER:

POINCARRE SUSY:  $\delta_{\epsilon_1} A_\mu = i \bar{\epsilon}_1 \gamma_\mu \psi$        $\delta_{\epsilon_1} \phi_I = i \bar{\epsilon}_1 \chi_I$

SUPER CONFORMAL:  $\delta_{\epsilon_2} A_\mu = i \bar{\epsilon}_2 X^\nu \gamma_\nu \gamma_\mu \psi$        $\delta_{\epsilon_2} \phi_I = i \bar{\epsilon}_2 X^\nu \chi_{\nu I}$

## $\frac{1}{2}$ BPS WILSON LOOPS

W.L. INVARIANT UNDER  $\frac{1}{2}$  POINCARÉ SUSY  
AND  $\frac{1}{2}$  SUPERCONFORMAL

→ SUSY SELECT A PREFERRED  
CURVE:

$$C = (x^0 = t, x^i = 0, A^I = m^I) \quad (m^I)^2 = 1$$

OPERATOR UNDER EXAM

$$W_R = \text{Tr}_R \mathcal{P} \exp \left[ i \int d\tau (A_0 + \phi) \right]$$

SYMMETRIES

SYM

$\frac{1}{2}$  BPS W.L.

$$SU(2,2) \simeq SO(2,4) \longrightarrow SU(2,1) \times SU(2)$$

$$SU(4) \simeq SO(6) \longrightarrow SO(5)$$

$\epsilon_1, \epsilon_2$

$$\longrightarrow \begin{cases} \gamma_{01} m^I \epsilon_1 = \epsilon_1 \\ \gamma_{02} m^I \epsilon_2 = -\epsilon_2 \end{cases}$$

$PSU(2,2|4)$

$$\longrightarrow \mathcal{O}_{\text{sup}} [4^2|4]$$

## $\frac{1}{2}$ BPS WILSON LOOPS

W.L. INVARIANT UNDER  $\frac{1}{2}$  POINCARÉ SUSY  
AND  $\frac{1}{2}$  SUPERCONFORMAL

→ SUSY SELECT A PREFERRED  
CURVE:

$$C = (x^0 = t, x^i = 0, y^I = n^I) \quad (n^I)^2 = 1$$

OPERATOR UNDER EXAM

$$W_R = \text{Tr}_R \mathcal{P} \exp \left[ i \int dt (A_0 + \phi) \right]$$

SYMMETRIES

SYM  $\frac{1}{2}$  BPS W.L.

$$SU(2,2) \simeq SO(2,4) \longrightarrow SU(1,1) \times SU(2)$$

$$SU(4) \simeq SO(6) \longrightarrow SO(5)$$

$$\epsilon_1, \epsilon_2 \longrightarrow \begin{cases} \gamma_{01} n^I \epsilon_1 = \epsilon_1 \\ \gamma_{01} n^I \epsilon_2 = -\epsilon_2 \end{cases}$$

$$PSU(2,2|4) \longrightarrow \text{osp}(4^*|4)$$

## GRAVITATIONAL DESCRIPTION OF $\frac{1}{2}$ BPS W.L.

WE LOOK FOR PROBES THAT PRESERVE  
THE  $O_{3p}(4^2|4)$  SYMMETRY, I.E. THE  
SAME AS  $\frac{1}{2}$  BPS W.L.

$R = \square$  FUNDAMENTAL REP.

→ FUNDAMENTAL STRING DESCRIBED  
BY  $X^M = X^M(\tau, \sigma^1)$

$$x^0 = \tau \quad \mu = \sigma^1 \quad x^i = 0 \quad x^I = \alpha^I$$

SPANS  $AdS_2$

SITS AT  $x^i = 0$

SITS AT A  
POINT IN  $S^5$

$$\rightarrow SU(1,1) \otimes SU(2) \otimes SO(5)$$

→ SAME CONDITION ON  $E_1, E_2$  AS W.L.

$O_{3p}(4^2|4)$  SYMMETRY

→ AT THE BOUNDARY THE STRING  
ENDS ON  $x^0 = t$

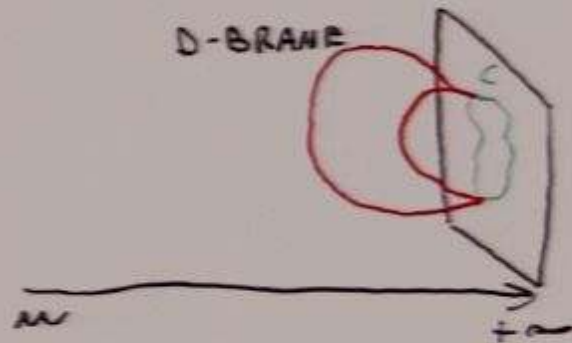
HIGHER REPRESENTATIONS



MULTIPLE  $(k)$  COINCIDENT FUNDAMENTAL STRINGS



DESCRIPTION IN TERMS OF A SINGLE D-BRANE WITH  $(k)$  UNIT OF FUNDAMENTAL STRING CHARGE



$DS_k$ -BRANE

GIANT WILSON LOOP

16

$$X^M = X^M(\sigma^0, \sigma^1, \dots, \sigma^5) \quad \text{EMBEDDING IN } AdS_5 \times S^5$$

$$\begin{array}{ccc} AdS_2 & \oplus & S_4 \\ \uparrow & & \uparrow \\ AdS_5 & \otimes & S^5 \end{array} \quad \text{EMBEDDING}$$

$$\rightarrow SU(2, 1) \times SU(2) \times SO(5) \quad \text{SYMMETRY}$$

$$\rightarrow \text{SAME } \frac{1}{2} \text{ BPS SUSY}$$

EMBEDDING IN  $S^5$



$$\rightarrow \theta_k \quad \text{LATITUDE } \in [0, \pi]$$

$$\rightarrow \theta_0 = 0 \quad \theta_N = \pi$$

$\rightarrow$  INCREASE WITH  $K$

$K \leq N$  MAX NUMBER OF FUNDAMENTAL CHARGES

$$DS_k \leftrightarrow W \left. \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right\} k$$

$D3_K$  - BRANE      DUAL GIANT WILSON LOOP <sup>15</sup>

$$x^M = x^M(\sigma^0, \sigma^1, \sigma^2, \sigma^3) \quad \text{EMBEDDING}$$

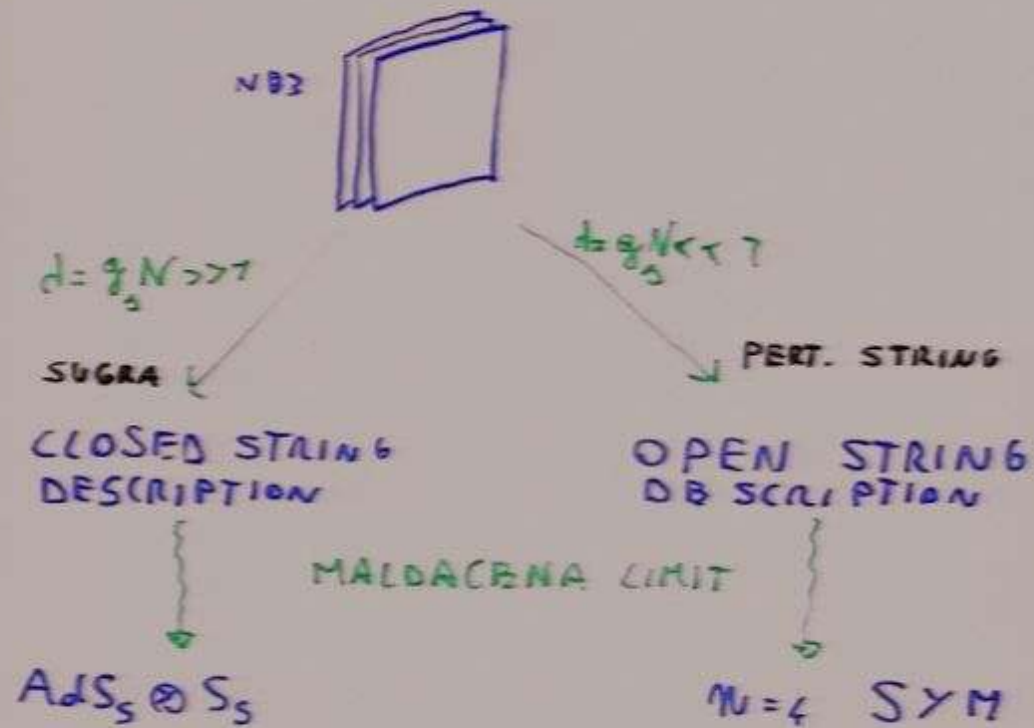
$AdS_2 \otimes S_2$	POINT
$\uparrow$	$\uparrow$
$AdS_5$	$S^5$

$\rightarrow O_{sp}(4^*/4)$  SYMMETRY

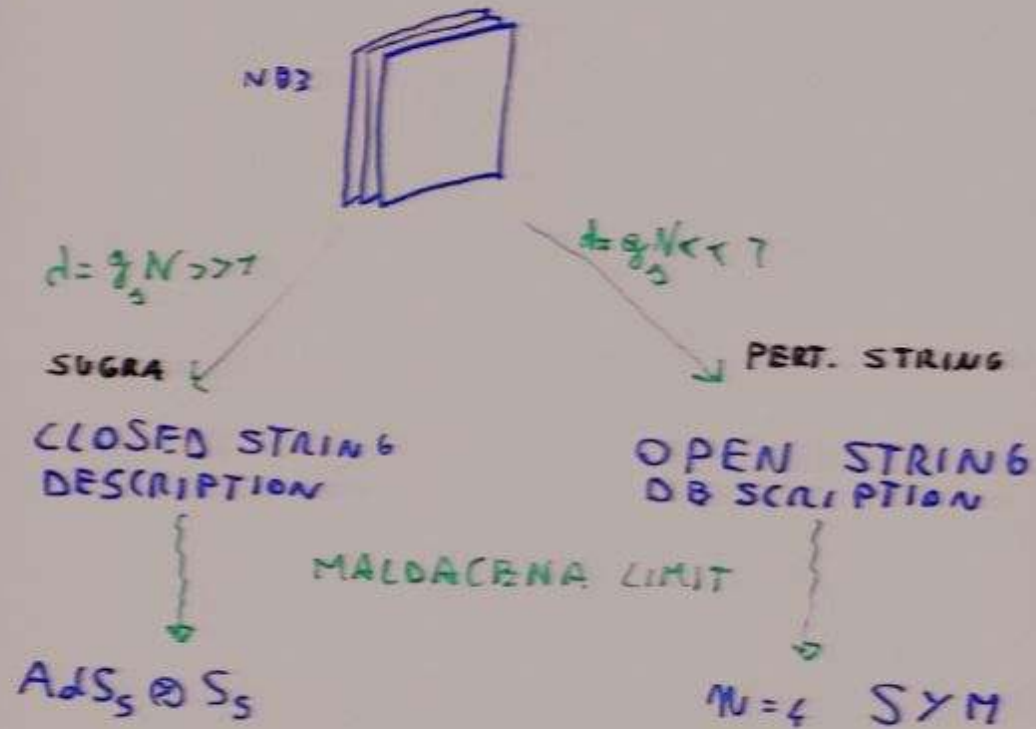
ARBITRARY NUMBER (K) OF FUNDAMENTAL CHARGES



	0	1	2	3	4	5	6	7	8	9	SPATIAL COORDINATES
N DB	x	x	x	x							



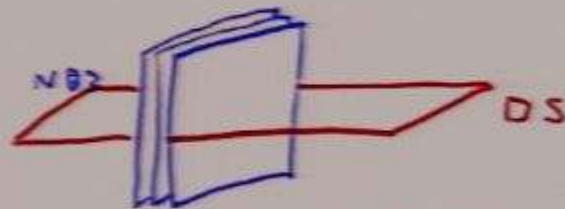
$0 \ 1 \ 2 \ 3 \ 4 \ 5 \ 6 \ 7 \ 8 \ 9$  SPATIAL  
 N D3 X X X X COORDINATES



# DS<sub>k</sub> GIANT WILSON LOOP

16  
7

	0	1	2	3	4	5	6	7	8	9	SPATIAL COORDINATES
N DB	x	x	x	x							
DS					x	x	x	x	x		



$$d = \frac{9}{8} N \gg 7$$

SUGRA

CLOSED STRING DESCRIPTION

AdS<sub>5</sub> ⊗ S<sub>5</sub>

→ DS<sub>k</sub> GIANT W.L.  
PROBE EMBEDDED  
IN AdS<sub>5</sub> × S<sub>5</sub>

$$d = \frac{9}{8} N \ll 7$$

PERT. STRING

OPEN STRING DESCRIPTION

N=4 SYM

→ DEFECT FIELD THEORY  
(NS-D, D-MS)

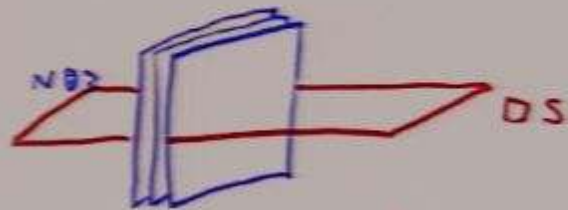
→ NON DYNAMICAL  
S-D.O.F.

MALDACENA LIMIT

# DS<sub>k</sub> GIANT WILSON LOOP

16  
7

	0	1	2	3	4	5	6	7	8	9	SPATIAL COORDINATES
N DB	x	x	x	x							
DS					x	x	x	x	x	x	



$$d = \frac{9}{8} N \gg 7$$

SUGRA

CLOSED STRING DESCRIPTION

$$AdS_5 \otimes S^5$$

→ DS<sub>k</sub> GIANT W.L.  
PROBE EMBEDDED  
IN AdS<sub>5</sub> × S<sub>5</sub>

$$d = \frac{9}{8} N \ll 7$$

PERT. STRING

OPEN STRING DESCRIPTION

$$N=4 \text{ SYM}$$

→ DEFECT FIELD THEORY  
(NS-D, D-MS)  
→ NON DYNAMICAL  
S-S D.O.F.

MALDACENA LIMIT

## DEFECT FIELD THEORY

$$S = S_{N=4} + \int dt \left( i \chi^\dagger \partial_t \chi + \chi^\dagger (A_0 + \phi) \chi \right)$$

: THE FOUND. REP. OF  $U(N)$

$$+ S_{\text{EXTRA}}$$

NEW DYNAMICAL D.O.F. ( $\tilde{A}$ )

$$S_{\text{EXTRA}} = \int dt \left( \chi^\dagger \tilde{A}_0 \chi - K \tilde{A}_0 \right)$$

ENCODE THE CHOICE OF  $K$

WE NOW INTEGRATE OUT THE D.O.F. ASSOCIATED TO  $DS_K$ .

GAUGE CHOICE :  $A_0 + \phi = \text{DIAG}(\omega_1, \dots, \omega_N)$

E.O.M. FOR  $\chi$ 's :  $(i \partial_t + \omega_i) \chi_i = 0$

$$Z = \int [D\chi] [D\chi^\dagger] [D\tilde{A}_0] e^{i(S + S_{\text{EXTRA}})}$$

INTEGRATING ONLY THE  $\chi, \chi^\dagger$   $0 < t \leq \beta$  AND IGNORING  $S_{\text{EXTRA}}$

$$Z^* = e^{iS_{N=4}} \prod_{i=1}^N (1 + x_i) \quad x_i = e^{-\beta \omega_i}$$

EXPANDING

$$\prod_{i=1}^N (1+x_i) = \sum_{\ell=0}^N E_{\ell}(x_1, \dots, x_N)$$

$$E_{\ell}(x_1, \dots, x_N) = \sum_{i_1 < i_2 < \dots < i_{\ell}} x_{i_1} \dots x_{i_{\ell}}$$

WHERE

$$E_{\ell} = \text{Tr}_{\mathbb{R}} \left\{ \mathbb{P} \exp \left[ i \int dA (A_0 + \phi) \right] \right\}_{\ell}$$

$$= W_{\mathbb{R}} \left\{ \mathbb{R} \right\}_{\ell}$$

THUS:

$$Z^{\mathbb{R}} = \sum_{\ell=0}^N e^{iS_{N=4}} W_{\mathbb{R}} \left\{ \mathbb{R} \right\}_{\ell}$$

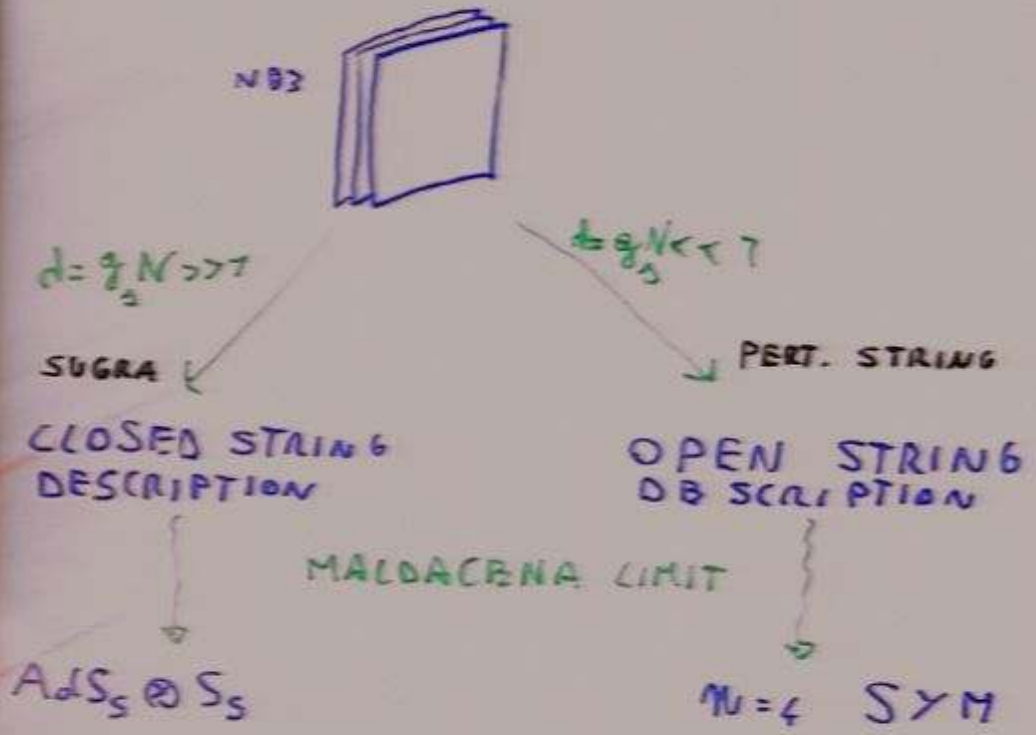
S<sub>EXTRA</sub>

INTEGRATING  $\tilde{A}_0 \leadsto \sum_{i=1}^N x_i^{\dagger} x_i = K$

$$\rightarrow \ell = K$$

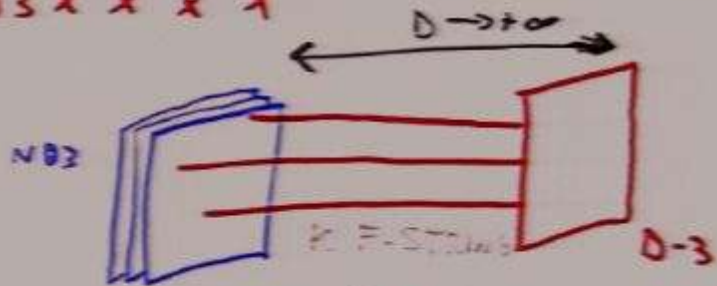
$$DS_K \longleftrightarrow Z = e^{iS_{N=4}} W_{\mathbb{R}} \left\{ \mathbb{R} \right\}_K$$

	0	1	2	3	4	5	6	7	8	9	SPATIAL COORDINATES
N DB	X	X	X	X							



# D3<sub>N</sub> DUAL GIANT WILSON LOOP <sup>7</sup><sub>7</sub><sup>19</sup>

	0	1	2	3	4	5	6	7	8	9	SPATIAL COORDINATES
N D3	x	x	x	x							
D3	x	x	x	x							



$$d = \frac{2}{3} N \gg 7$$

SUGRA

CLOSED STRING DESCRIPTION

$$AdS_5 \otimes S^5$$

$$\Rightarrow AdS_2 \otimes S^2 \subset AdS_5$$

PROBE D-3

$$d = \frac{2}{3} N \ll 7$$

PERT. STRING

OPEN STRING DB DESCRIPTION

$$N=4 \text{ SYM}$$

NON RELATIVISTIC W-BOSONS

W-BOSONS INSERTIONS IN THE PATH INTEGRAL

MALDACENA LIMIT

# LOW ENERGY FIELD THEORY D3

CONSIDER  $U(N+1)$  SYM IN THE COULOMB BRANCH

$$\langle \phi \rangle = \begin{vmatrix} 0 & 0 \\ 0 & D \end{vmatrix}$$

$\hookrightarrow$  DISTANCE BETWEEN TWO STACKS OF BRANA

$$U(N+1) \rightarrow U(N) \times U(1)$$

$$\hat{A}_\mu = \begin{vmatrix} A_\mu & w_\mu \\ w_\mu^\dagger & \tilde{A}_\mu \end{vmatrix}$$

$\hookrightarrow$  W-BOSONS, TRANSFORM IN THE FUNDAMENTAL OF  $U(N)$

$D$  IS THE MASS OF W-BOSONS

$D \rightarrow +\infty$  IS A NON RELATIVISTIC LIMIT

$$S = S_{m=4} + \int i \chi^\dagger \partial_t \chi + \chi^\dagger (A_0 + \phi) \chi$$

$$w = \frac{1}{\sqrt{D}} e^{-i\epsilon D} \chi$$

# LOW ENERGY FIELD THEORY D3<sup>20</sup>

CONSIDER  $U(N+1)$  SYM IN THE  
COULOMB BRANCH

$$\langle \phi \rangle = \begin{vmatrix} 0 & 0 \\ 0 & D \end{vmatrix}$$

↳ DISTANCE BETWEEN TWO  
STACKS OF BRANE

$$U(N+1) \rightarrow U(N) \times U(1)$$

$$\hat{A}_\mu = \begin{vmatrix} A_\mu & W_\mu \\ W_\mu^\dagger & \tilde{A}_\mu \end{vmatrix}$$

↳  $W$ -BOSONS, TRANSFORM IN THE  
FUNDAMENTAL OF  $U(N)$

$D$  IS THE MASS OF  $W$ -BOSONS

$D \rightarrow +\infty$  IS A NON RELATIVISTIC  
LIMIT

$$S = S_{m=4} + \int i \bar{\chi} \not{\partial} \chi + \bar{\chi} (A_0 + \phi) \chi$$

$$W = \frac{1}{\sqrt{D}} e^{-i\phi D} \chi$$

K STRETCHED STRINGS

$$= \begin{cases} K \text{ W-BOSON CREATION AT } t=-\infty \\ K \text{ W-BOSON ANNIHILATION AT } t=+\infty \end{cases}$$

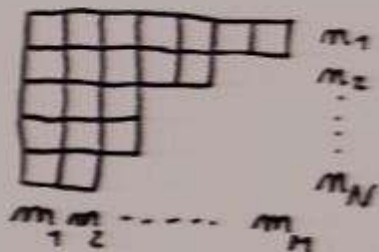
$$Z = e^{iS_{M=4}} \int [D\chi][D\chi^\dagger] e^{iS_\chi}$$

$$= \frac{1}{K!} \sum_{i_1, \dots, i_K} \chi_{i_1}(\infty) \dots \chi_{i_K}(\infty) \chi_{i_1}^\dagger(-\infty) \dots \chi_{i_K}^\dagger(-\infty)$$

$$= e^{iS_{M=4}} \underbrace{W}_K$$

$$D3_K \longleftrightarrow \underbrace{W}_K$$

$\frac{1}{2}$  BPS WILSON LOOP IN  
A GENERIC REPRESENTATION



IS ASSOCIATED TO

$$(D5_{m_1}, D5_{m_2}, \dots, D5_{m_M})$$

OR

$$(D3_{m_1}, D3_{m_2}, \dots, D3_{m_M})$$

W. L. EXPECTATION VALUE

ANTISYMMETRIC REPRESENTATION

$$\langle W_{\vec{k}} \rangle = e^{-S_0 S_{\vec{k}}}$$

CONSIDERING  $C = C_0$

$$\rightsquigarrow E_{\vec{k}} = \frac{2N}{3\pi} \sin^3 \theta_{\vec{k}} \cdot E_0(\vec{k})$$

$$K = N \Rightarrow \theta_N = \pi \Rightarrow E_{\left. \begin{matrix} \vec{k} \\ \vdots \\ \vec{k} \end{matrix} \right\} N} = 0$$

SINGULET STATES  
DO NOT INTERACT

## CONCLUSION

- ANY  $\frac{1}{2}$  BPS WILSON LOOPS HAS A DUAL DESCRIPTION IN TERMS OF D5 OR ALTERNATIVELY IN TERMS OF D3
- COMPUTATION OF WILSON LOOPS EXPECTATION VALUE USING BRANES