

Title: Cosmological Constraints from Gravitational Lensing

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Abstract:

# Cosmological Constraints from Gravitational Lensing

James Taylor

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University of Waterloo

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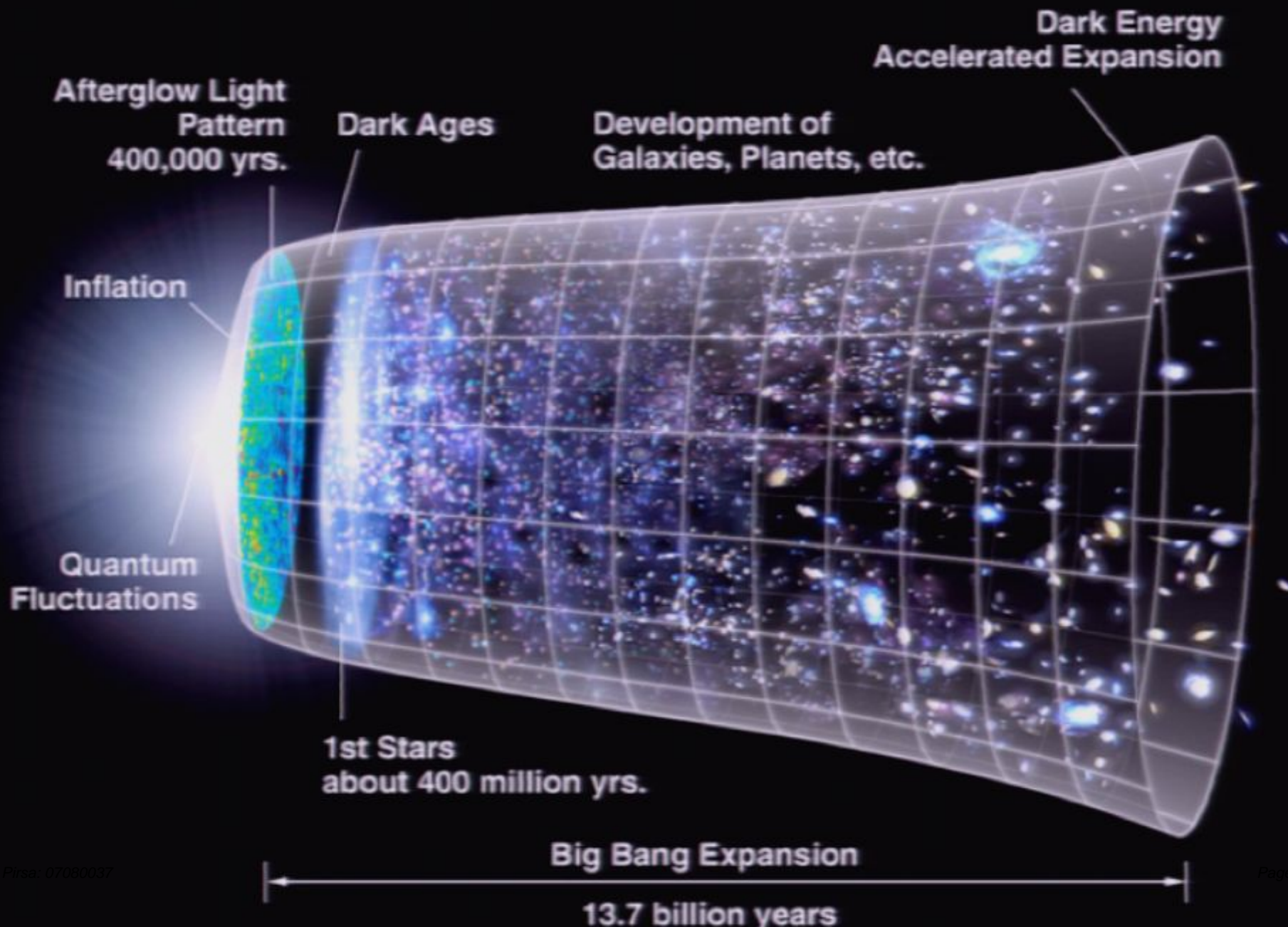
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## Outline

- 1) Introduction: what is observational cosmology and what does it tell us?
- 2) Lensing Basics, and forms of lensing
- 3) Weak Lensing results in cosmology
- 4) Results from a few recent surveys

# Overview of Cosmic History



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<u>Era</u>	<u>Redshift</u>	<u>time</u>	<u>T (K)</u>	<u>E</u>
Inflation, Topology		$10^{-35}$ s	$10^{27}$	$10^{14}$ GeV
The High-Energy Universe		$10^{-6}$ s	$10^{13}$	1 GeV
Last Scattering	1100	$3 \times 10^5$ yr	3300	1 eV
The Dark Ages	1000-70	1 Myr	1000	
The First Stars	70-50	50 Myr	100	
Reionization	8-15	.2-.6 Gyr	30	
Formation of Massive Galaxies & AGN	2-5	1-3 Gyr		
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# What can we learn from observational cosmology?

Basically two things:

\* Current expansion rate and expansion history of the smoothed background

← net contents of the universe (matter, radiation, vacuum energy, etc.)

\* Fluctuations in the smooth background

← production method (inflation) and expansion history of the early universe

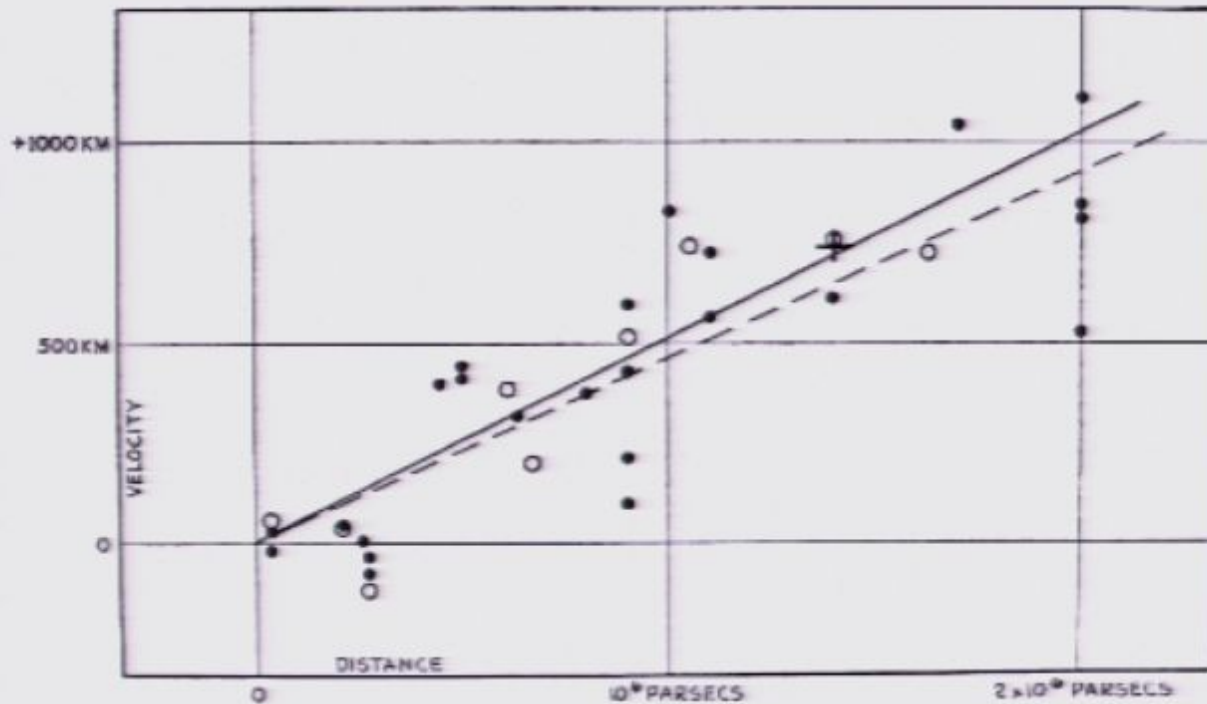
← physical properties (velocity, temperature, pressure, interactions) of the different matter and energy components



# What can we learn from observational cosmology?

a) The current expansion rate: the Hubble constant

$$\langle v \rangle \sim (\lambda - \lambda_0 / \lambda_0) = H_0 d; H_0 \sim 73 \pm 5 \text{ km/s/Mpc}$$



all observed distances, densities etc. scale with this constant, often written as

$$H_0 = 100 h \text{ km/s/Mpc where } h \sim 0.73$$

## What can we learn from observational cosmology?

b) The past expansion rate at all times,  $H(a) = \dot{a}/a$

N.B. easiest to measure the scale the expansion has reached:

scale factor  $a = 1/(1+z)$

$$H^2(a) = H^2(0) [ \Omega_R(0) a^{-4} + \Omega_m(0) a^{-3} + \Omega_k(0) a^{-2} + \Omega_\Lambda(0) a^0 ]$$

(Friedmann equation)

$$\text{here } \Omega_m(0) = \rho_m(0) / \rho_{\text{crit}}(0)$$

for dark energy replace  $\Omega_\Lambda a^0$  with  $\Omega_X a^{-n}$  where  $n = 3(1+w)$

$$w = p/\rho$$

equation of state parameter; can in principle be  $w(a)$

## What can we learn from observational cosmology?

c) Various integrals of the expansion and/or geometric terms:

“Luminosity distance”

(+ in principle also time)

$$d_L = (1+z) S_k(r) = a \int_0^z dz/H(z)$$

$$\text{where } S_k(r) = \begin{cases} R_0 \sin(r/R_0) & k = +1 \\ r & k = 0 \\ R_0 \sinh(r/R_0) & k = -1 \end{cases}$$

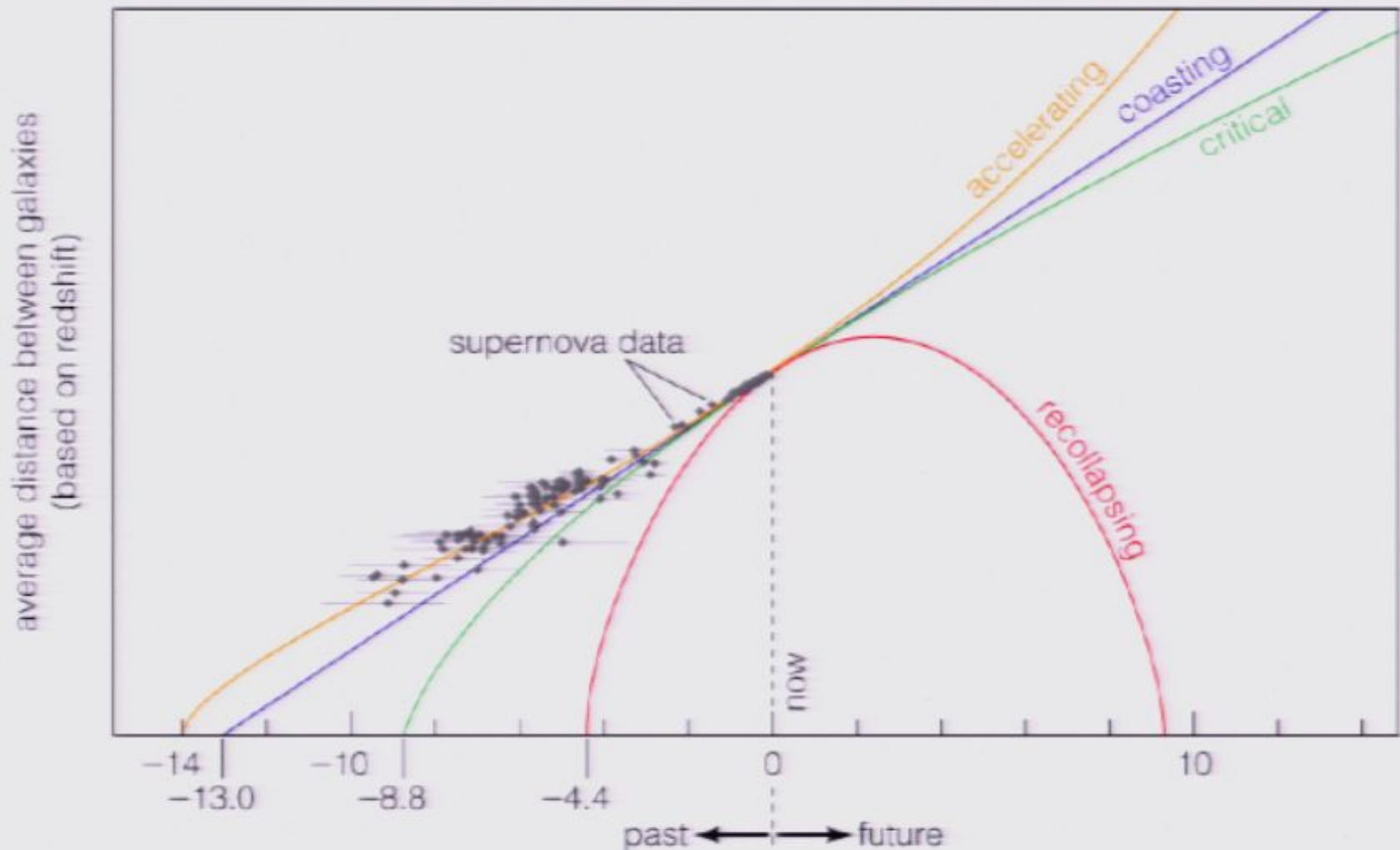
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 $\Rightarrow$  equation of state

“Angular diameter distance”

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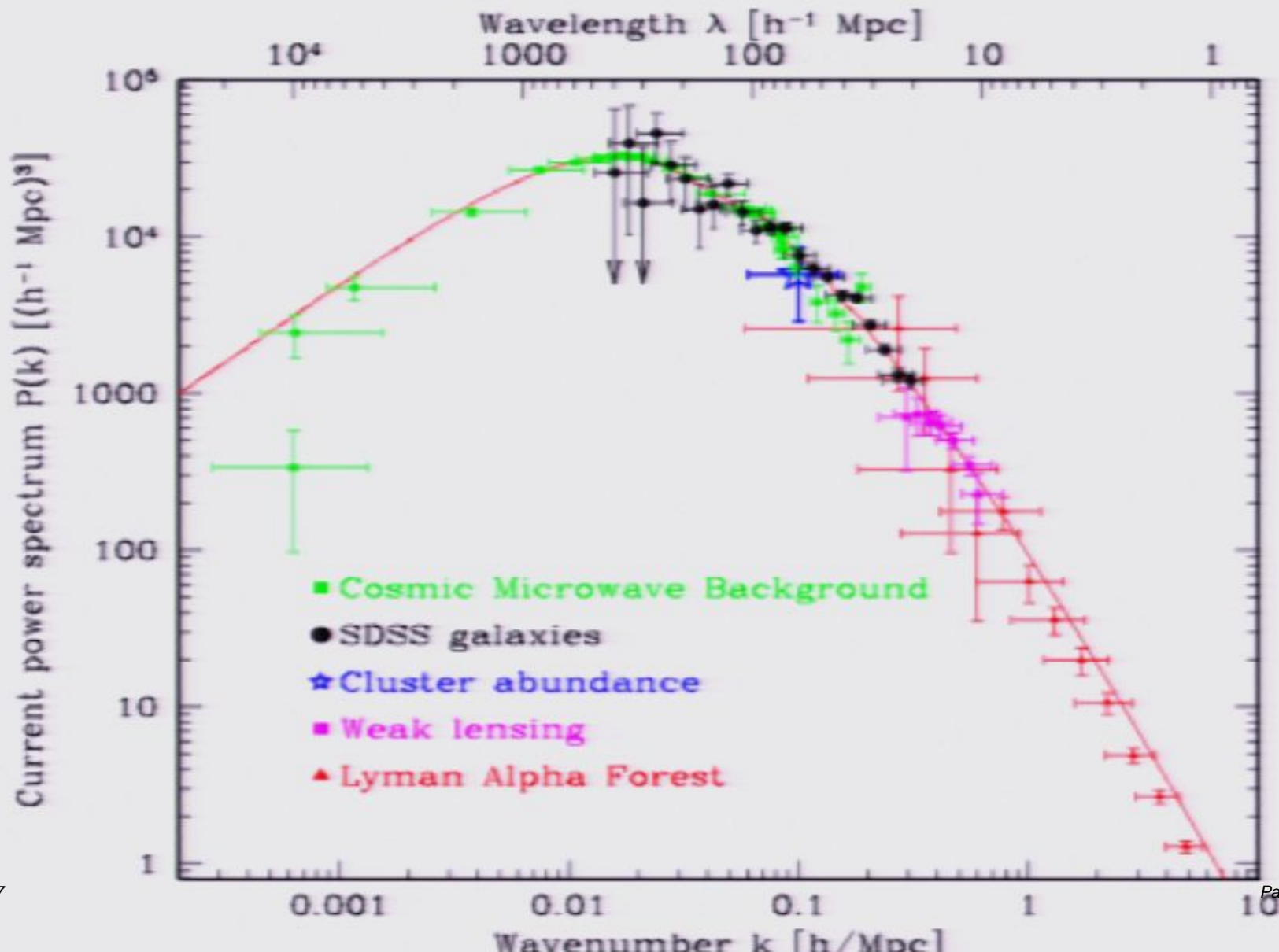
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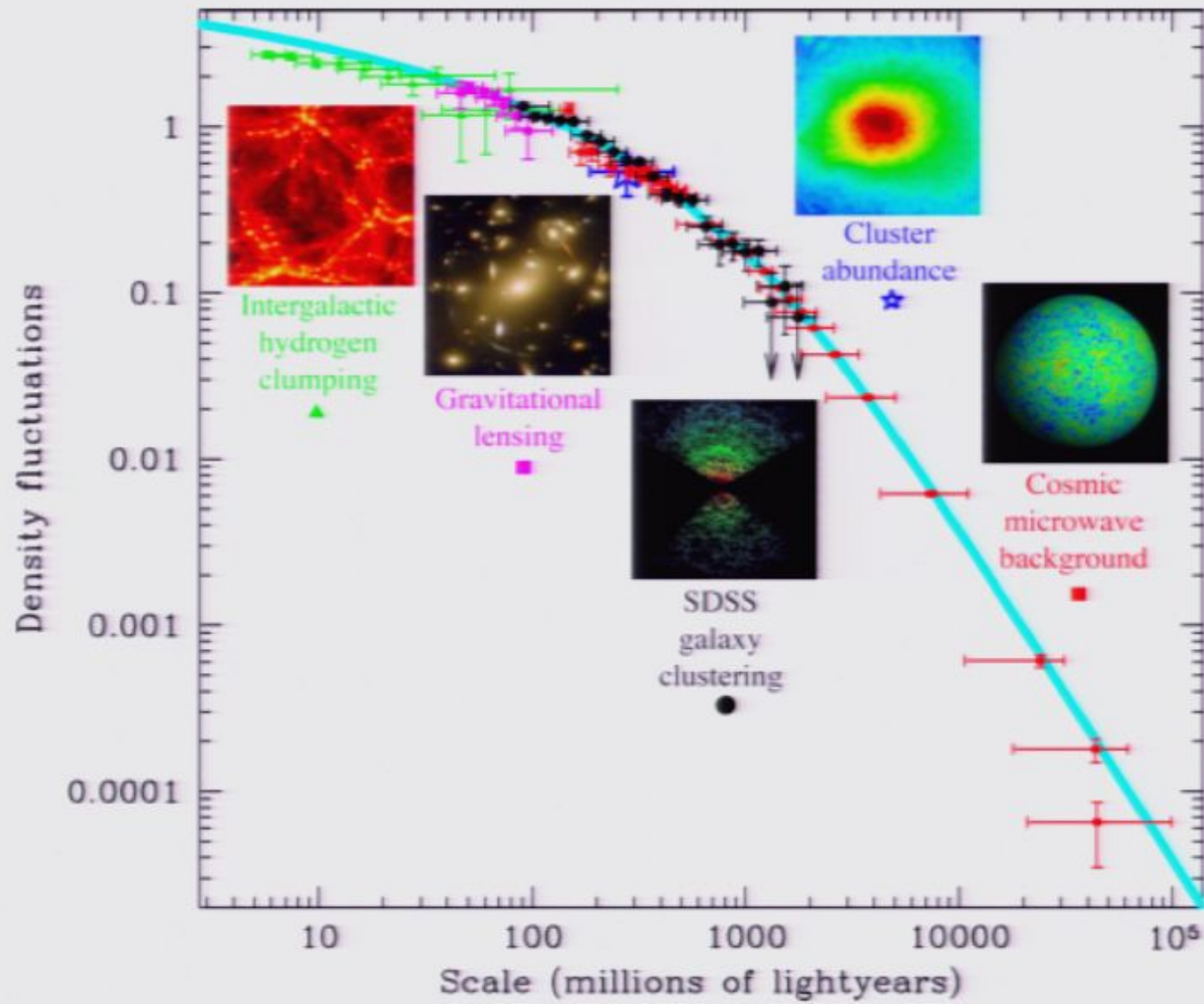
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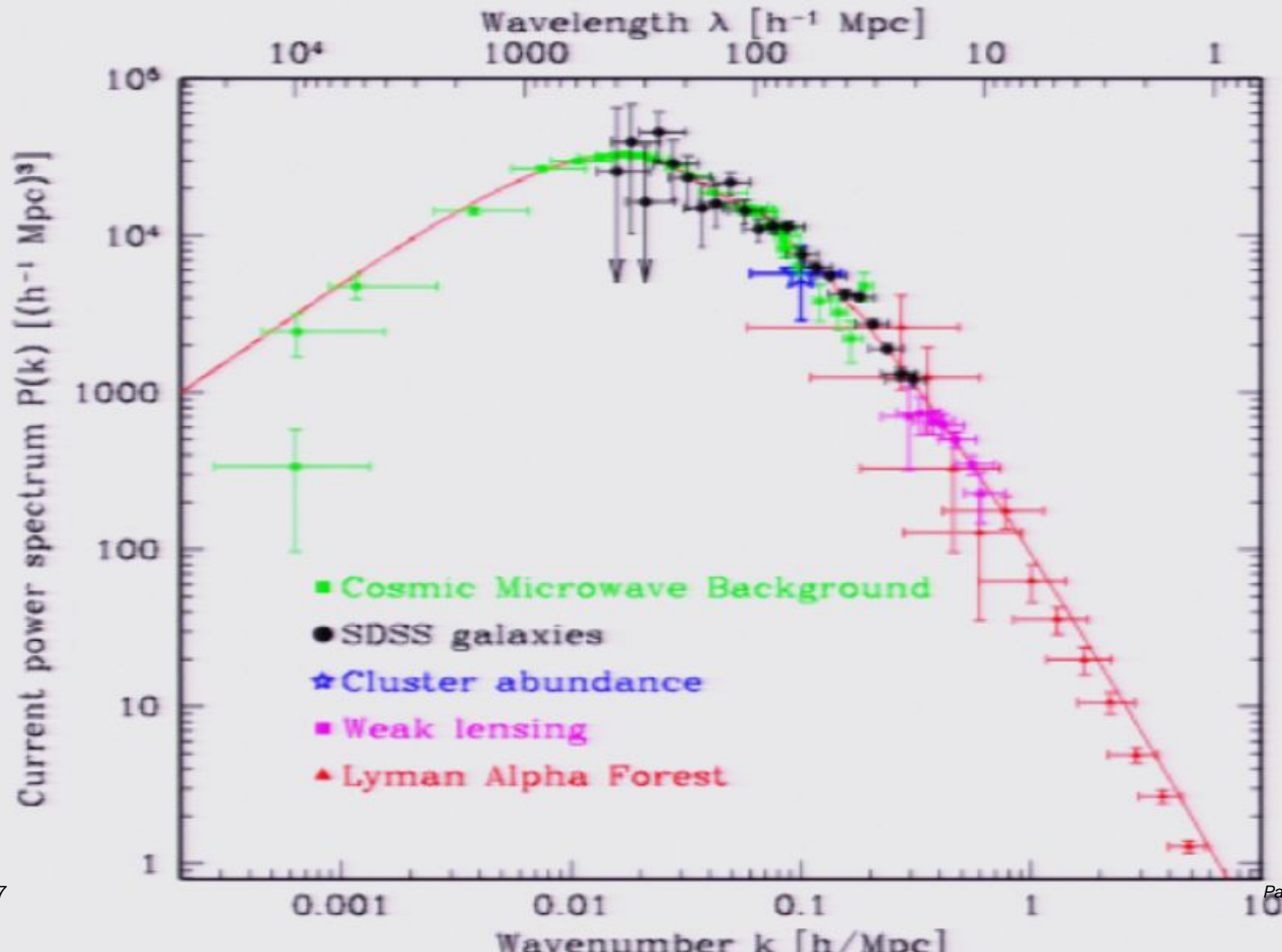
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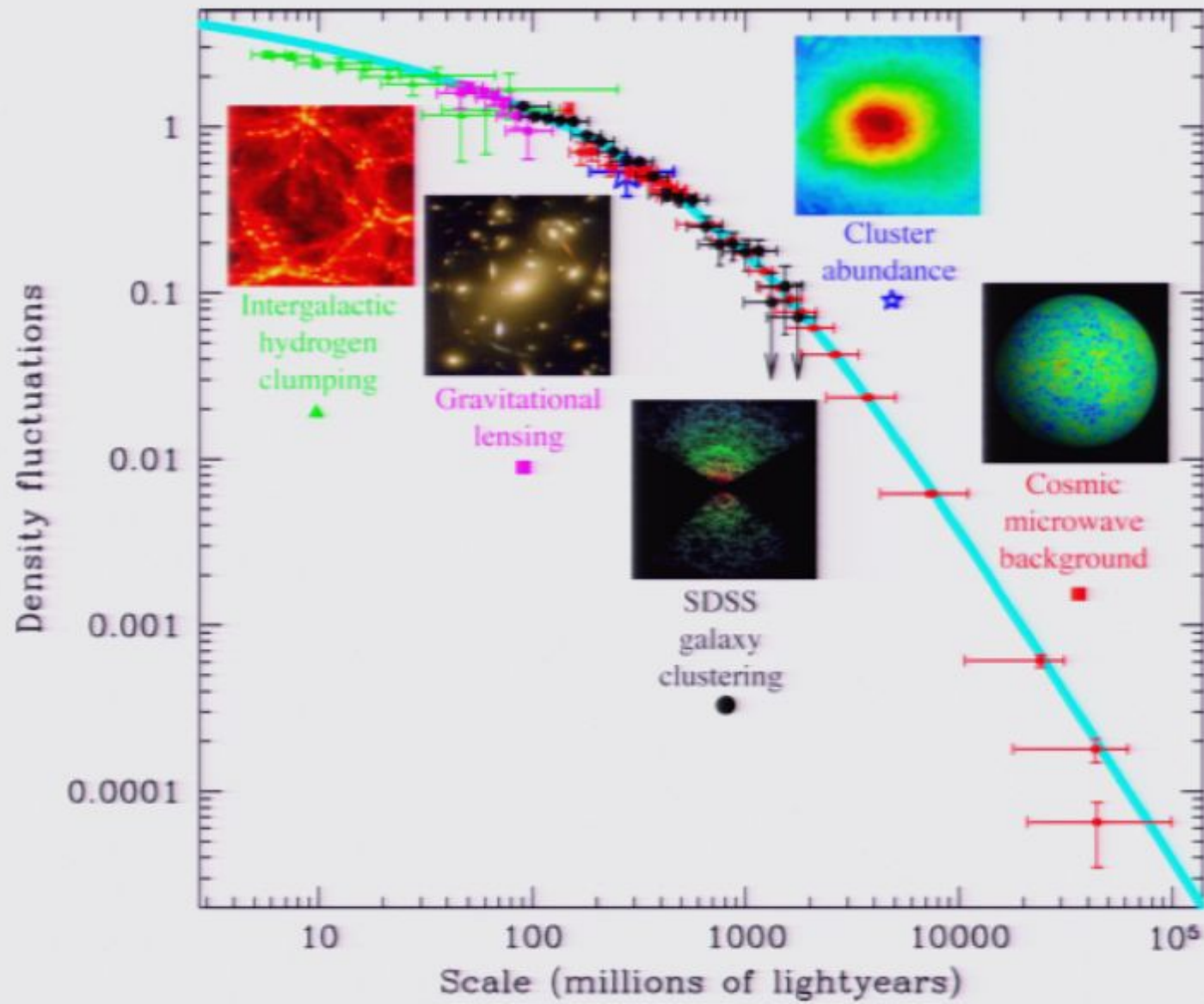
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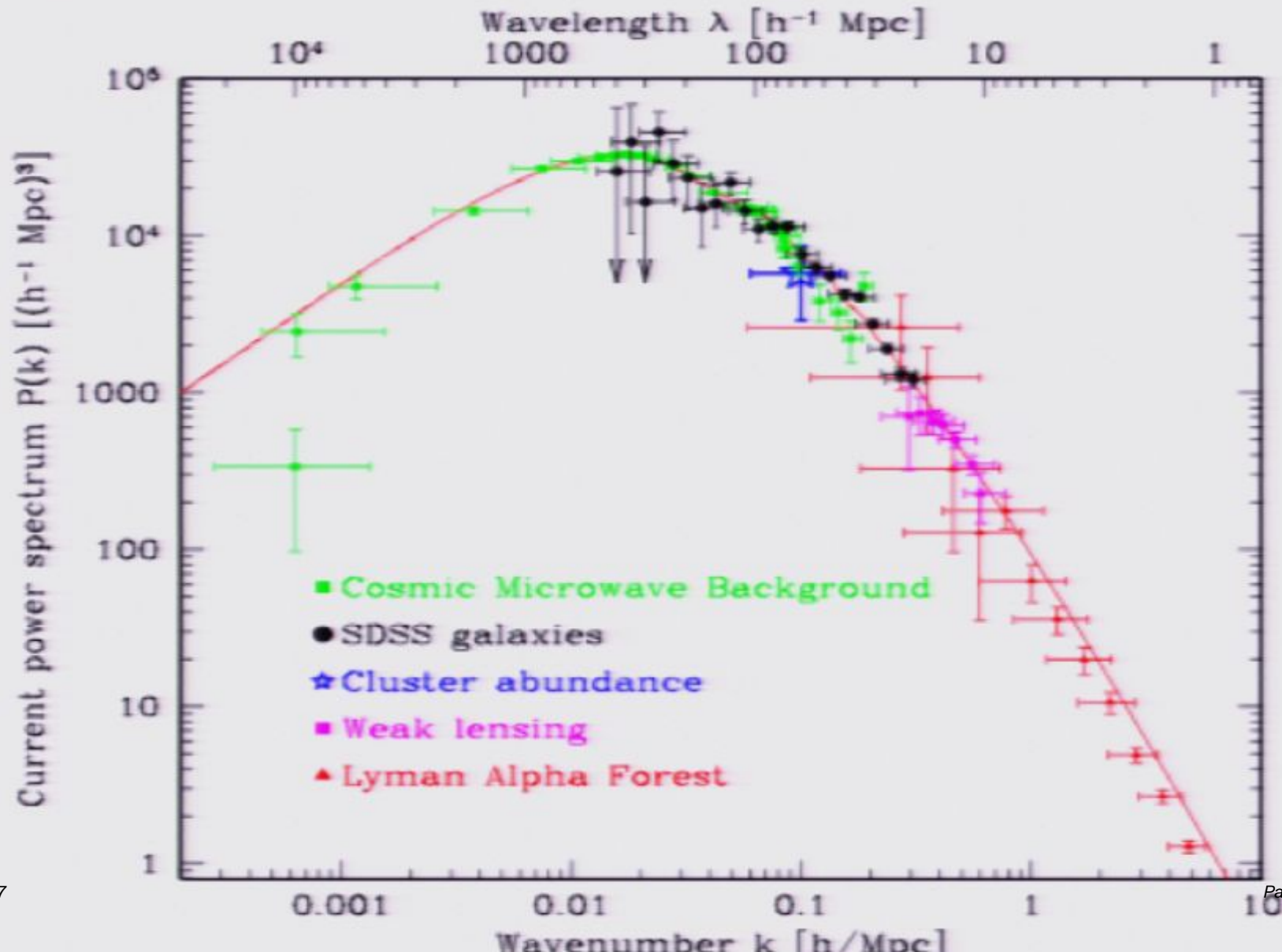
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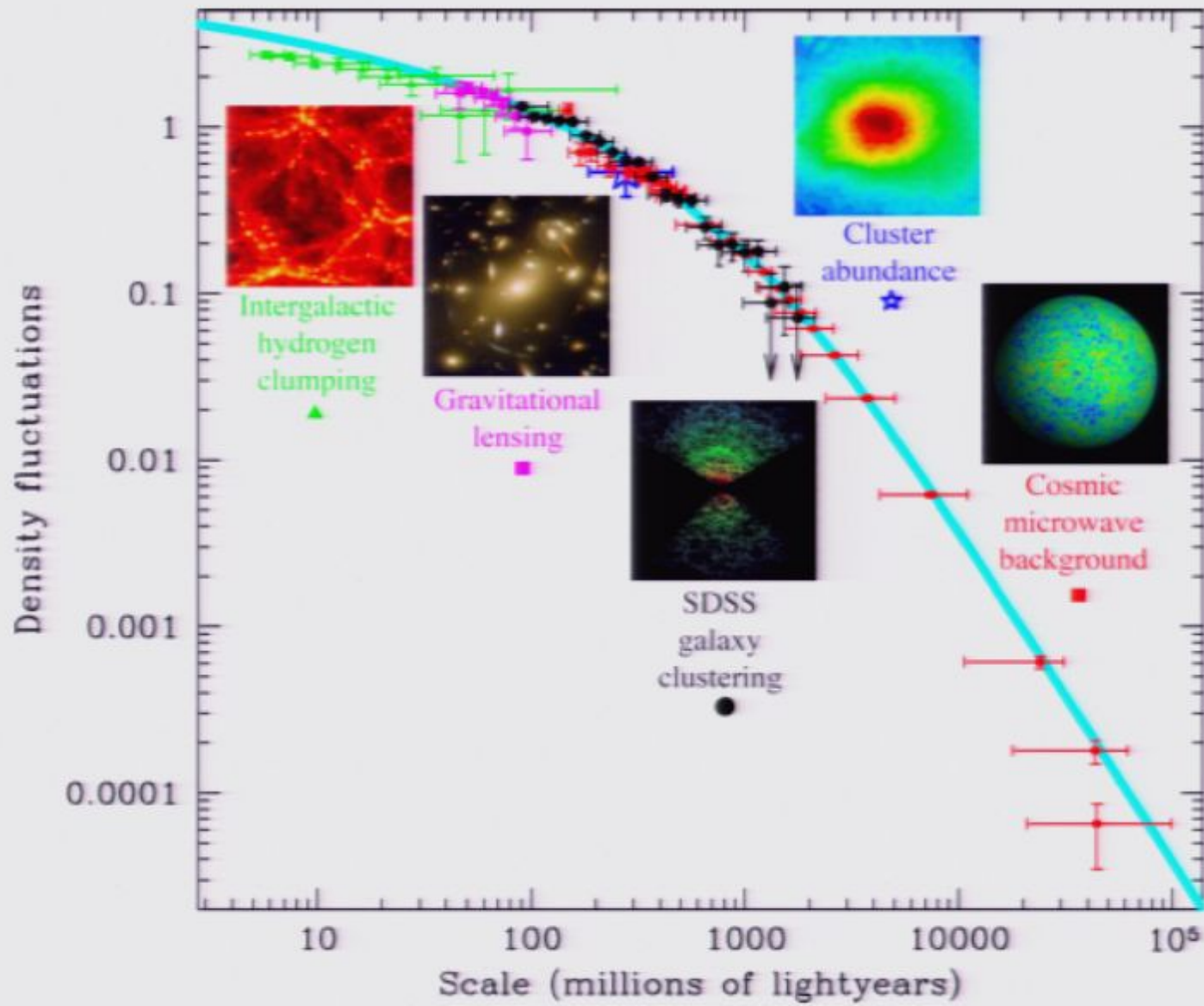
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# What can we learn from observational cosmology?

e) The physics of the main components, e.g. :

Dark Matter annihilation (e.g. Neutralino annihilation to photons/pions)

Dark Matter scattering cross-section

Dark Matter decays

other dark matter properties?

neutrino mass

neutrino oscillations

dark energy properties?

variations in the fundamental constants

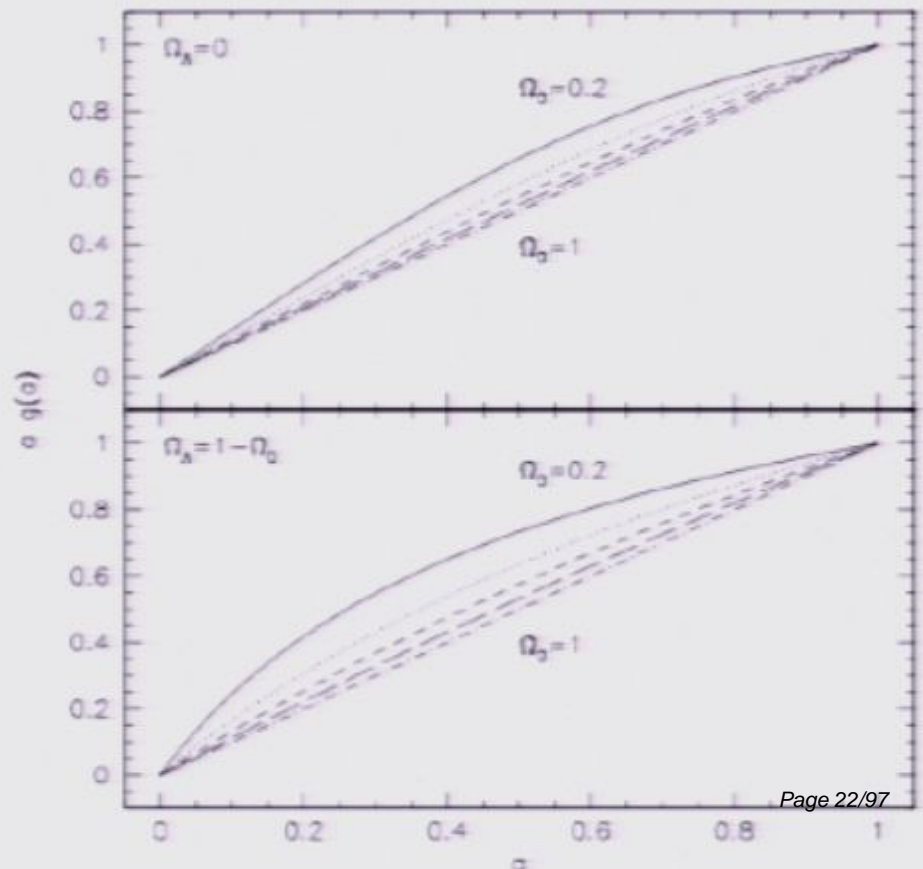
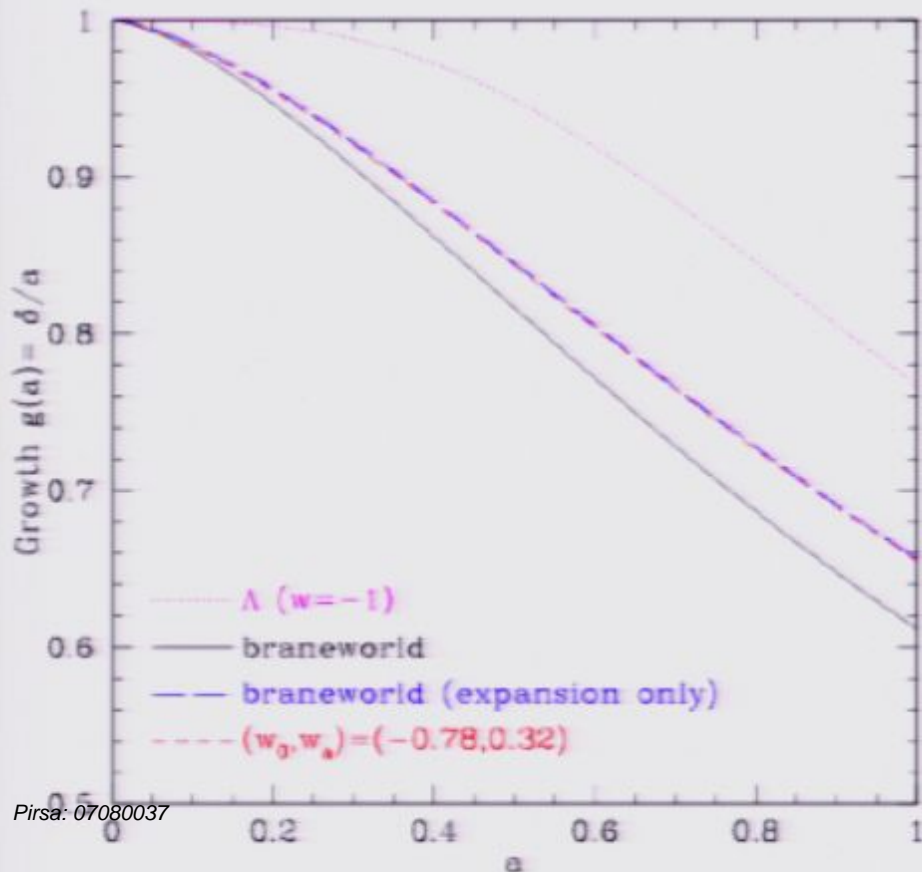
tests of strong and weak gravity

## N.B. the value of multiple constraints (cf. Linder 2005)

Can modify the Friedman equation  $H(a)$  either by changing the equation of state of dark energy or by modifying gravity; either could in principle produce exactly the same expansion history.

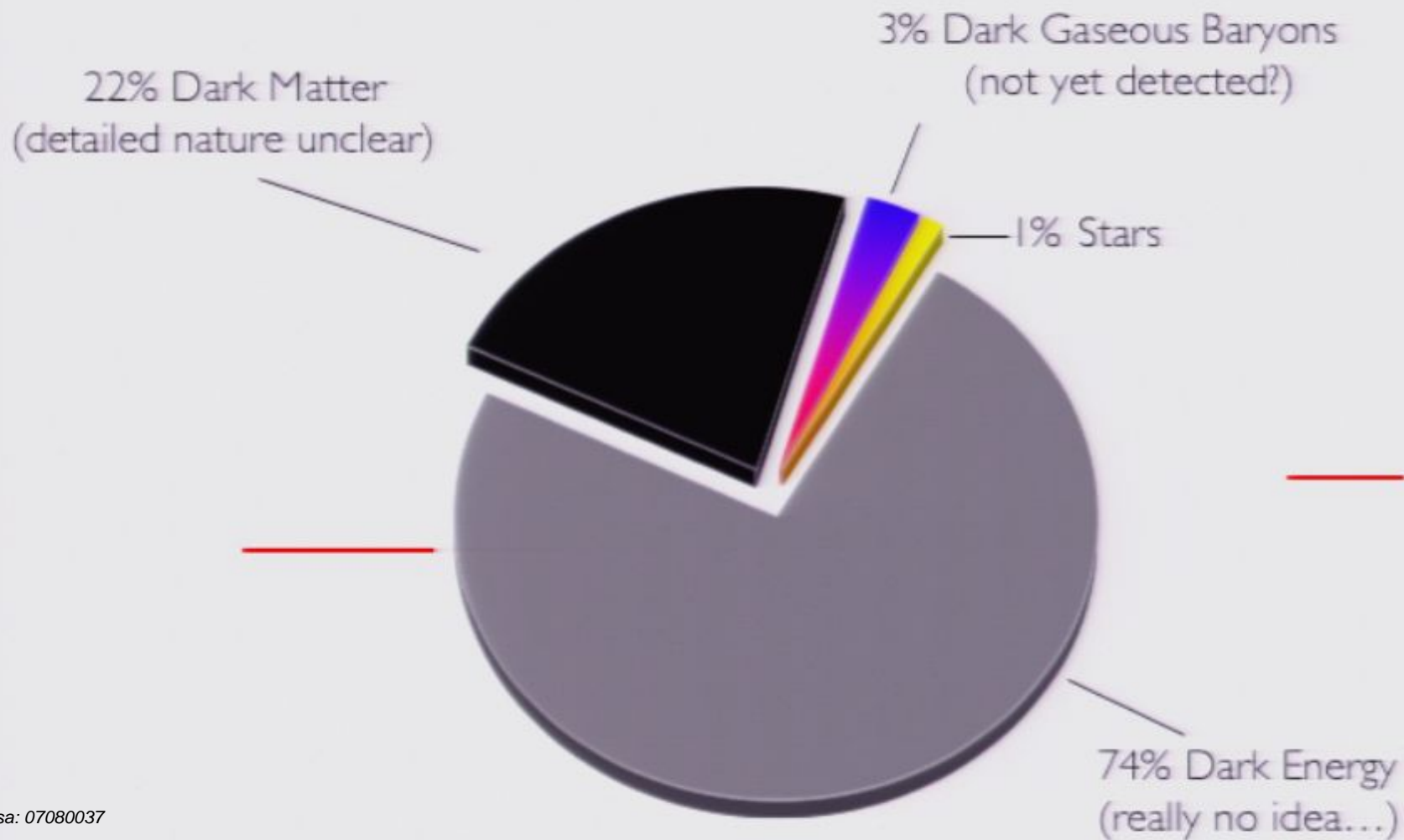
Linear growth factor describing the growth of perturbations can distinguish between the two (local vs. global effect), e.g. DGP braneworld gravity

$$g \equiv \frac{D}{a} = \frac{5\Omega_m}{2} \frac{H(a)}{a} \int_0^a \frac{da'}{a'^3 H(a')^3}$$



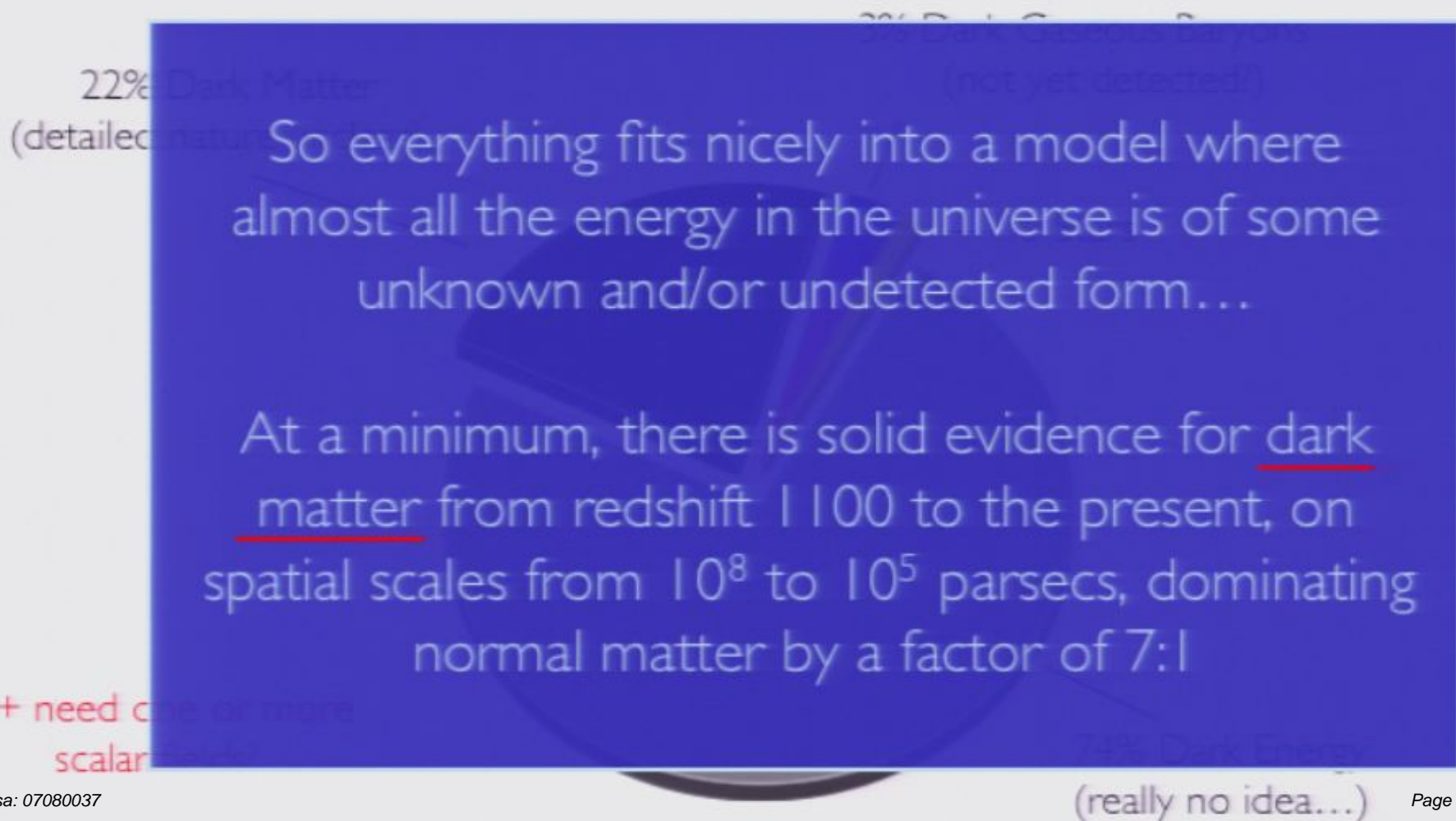
## Summary:

# The Composition of the Universe (third quarter 2007)



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# Gravitational Lensing

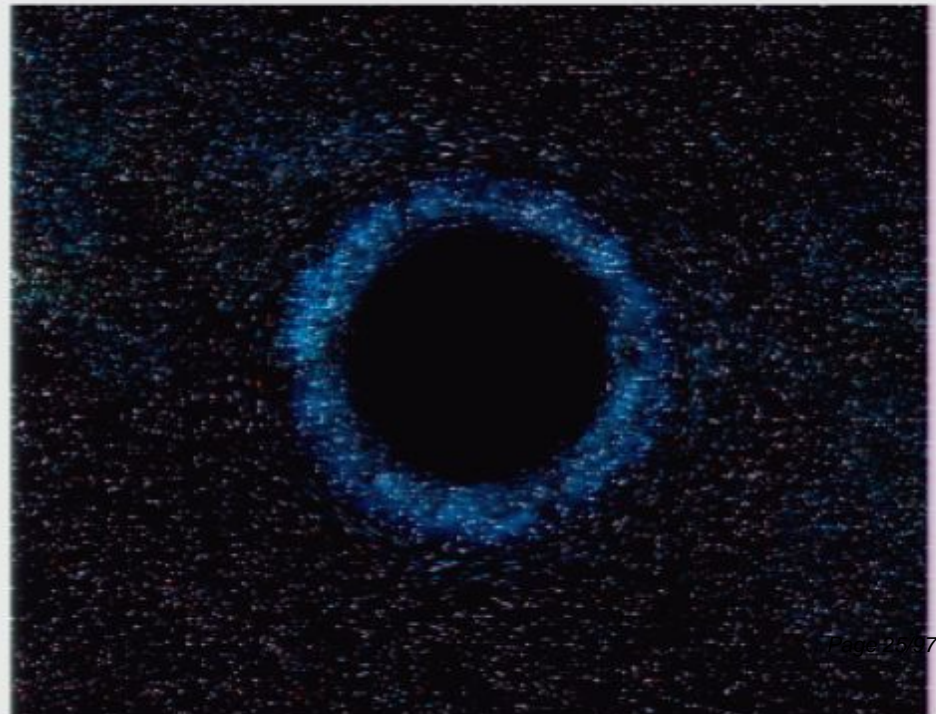
(cf. reviews by Schneider on astro-ph)

## Basic Idea:

point mass  $M$  deflects light; using the Schwarzschild solution  
calculate deflection angle  $\hat{\alpha}$  :

$$\hat{\alpha} = \frac{4GM}{c^2 \xi} = \frac{2R_s}{\xi}$$

where  $\xi$  is the impact parameter



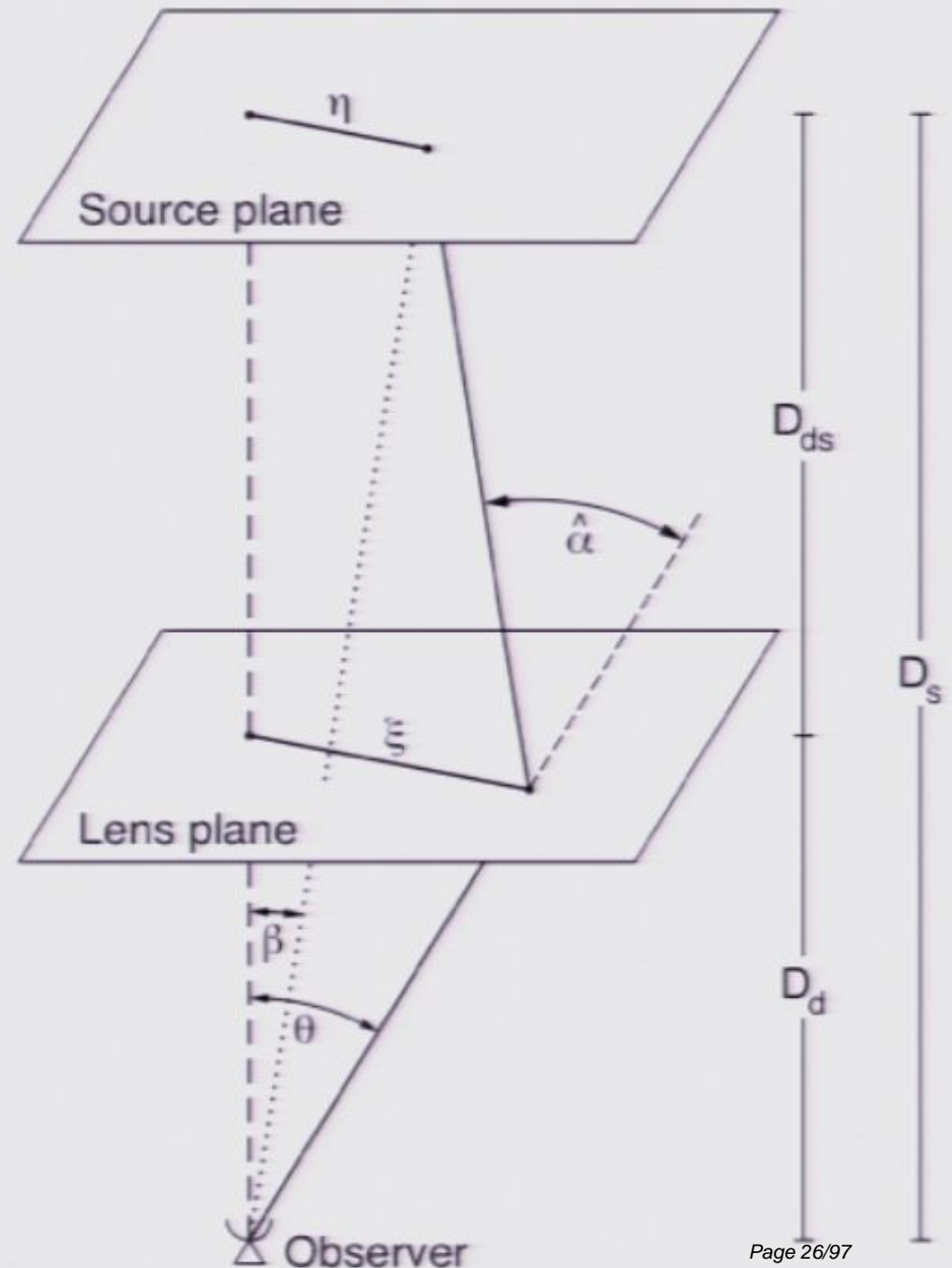
# Gravitational Lensing

More generally, for 2-D distribution of  $\Sigma$  deflection of ray from object at (2-D vector position  $\eta$  should be at  $\beta$ , instead seen

$$\theta = \hat{\alpha}(\theta) + \beta$$

where deflection angle  $\alpha$  is given by:

$$\hat{\alpha}(\xi) = \frac{4G}{c^2} \int d^2\xi' \Sigma(\xi') \frac{\xi - \xi'}{|\xi - \xi'|^2}$$



# Gravitational Lensing

Deflection angle can be written in terms of the convergence kappa:

$$\alpha(\boldsymbol{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} d^2\theta' \kappa(\boldsymbol{\theta}') \frac{\boldsymbol{\theta} - \boldsymbol{\theta}'}{|\boldsymbol{\theta} - \boldsymbol{\theta}'|^2}$$

where:

$$\kappa(\boldsymbol{\theta}) := \frac{\Sigma(D_d \boldsymbol{\theta})}{\Sigma_{\text{cr}}} \quad \text{with} \quad \Sigma_{\text{cr}} = \frac{c^2}{4\pi G} \frac{D_s}{D_d D_{\text{ds}}}$$

also note that we can define a potential such that:

$$\alpha = \nabla \psi, \quad \text{with} \quad \psi(\boldsymbol{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} d^2\theta' \kappa(\boldsymbol{\theta}') \ln |\boldsymbol{\theta} - \boldsymbol{\theta}'|$$

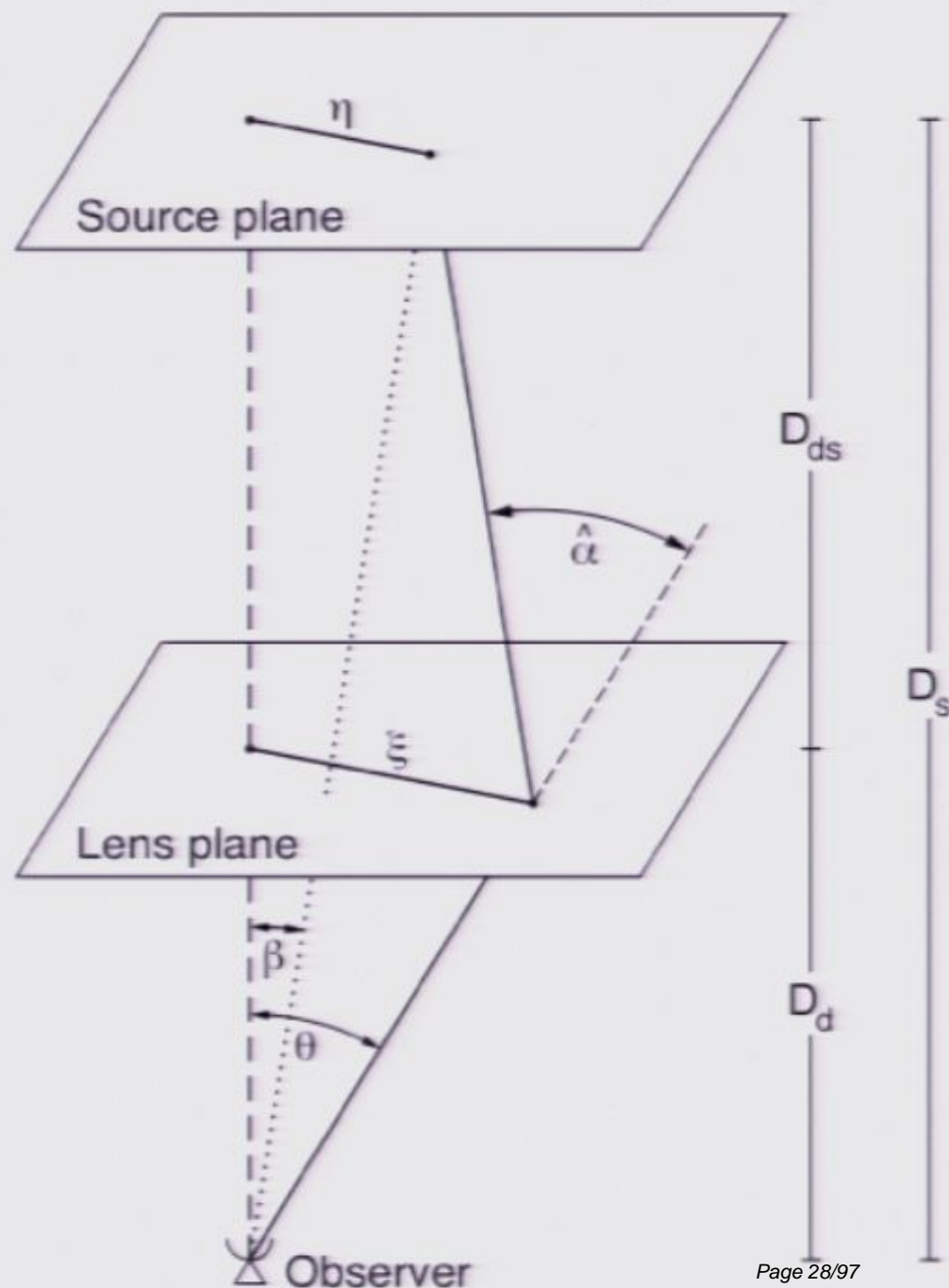
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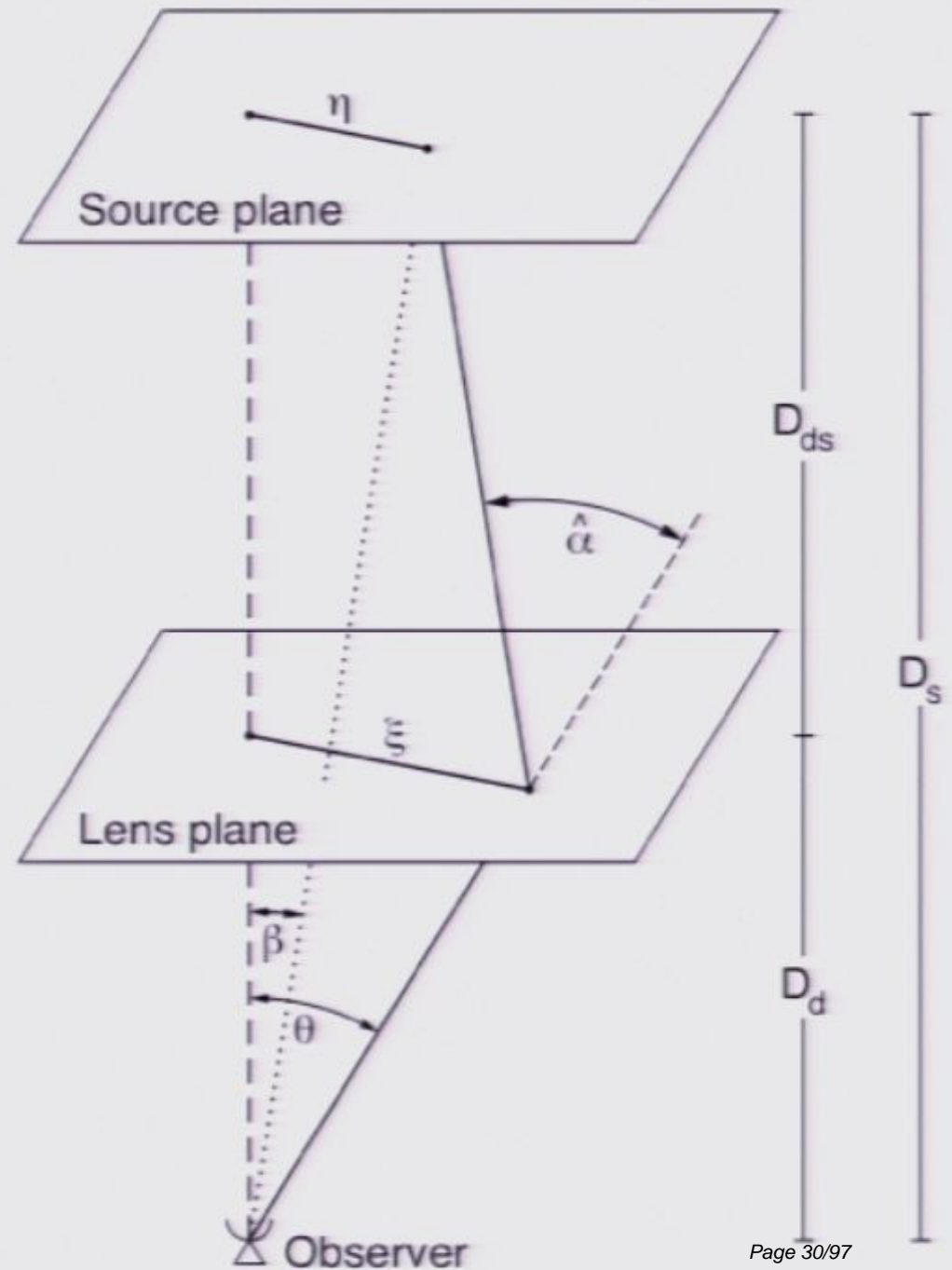
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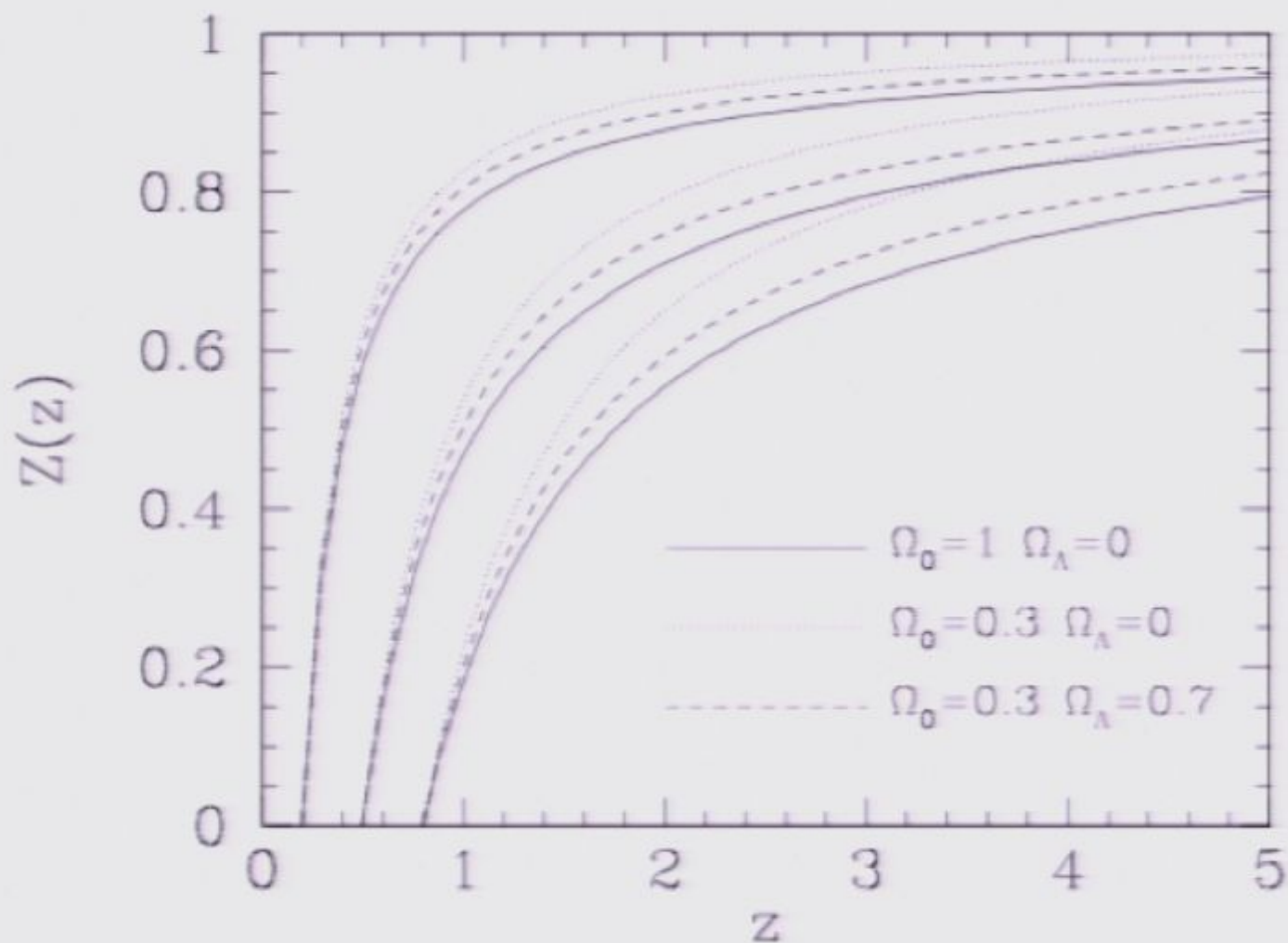
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# Gravitational Lensing

Effect of the geometric term is to make lenses most efficient when they are roughly halfway between the observer and the source:





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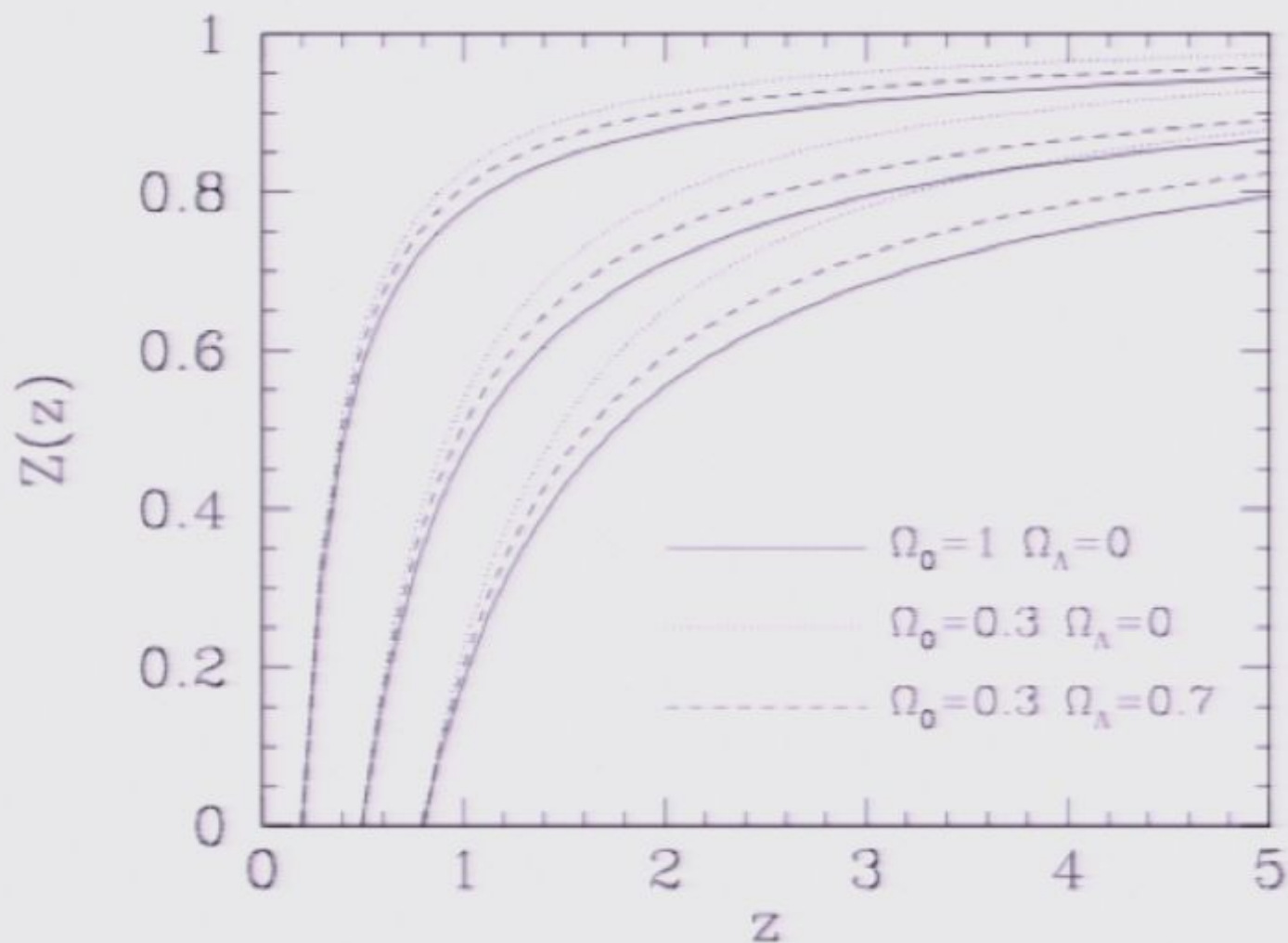
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## Shear and Weak Gravitational Lensing

Shear: consider a bundle of rays propagating from an extended source to the observer; if the mass distribution varies across the field different rays will be deflected by different amounts. Define a tensor describing the local deformation of objects:

$$\mathcal{A}(\boldsymbol{\theta}) = \frac{\partial \boldsymbol{\beta}}{\partial \boldsymbol{\theta}} = \left( \delta_{ij} - \frac{\partial^2 \psi(\boldsymbol{\theta})}{\partial \theta_i \partial \theta_j} \right) = \begin{pmatrix} 1 - \kappa - \gamma_1 & -\gamma_2 \\ -\gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix}$$

kappa: convergence (brightens/dims)

gamma: shear (distorts)

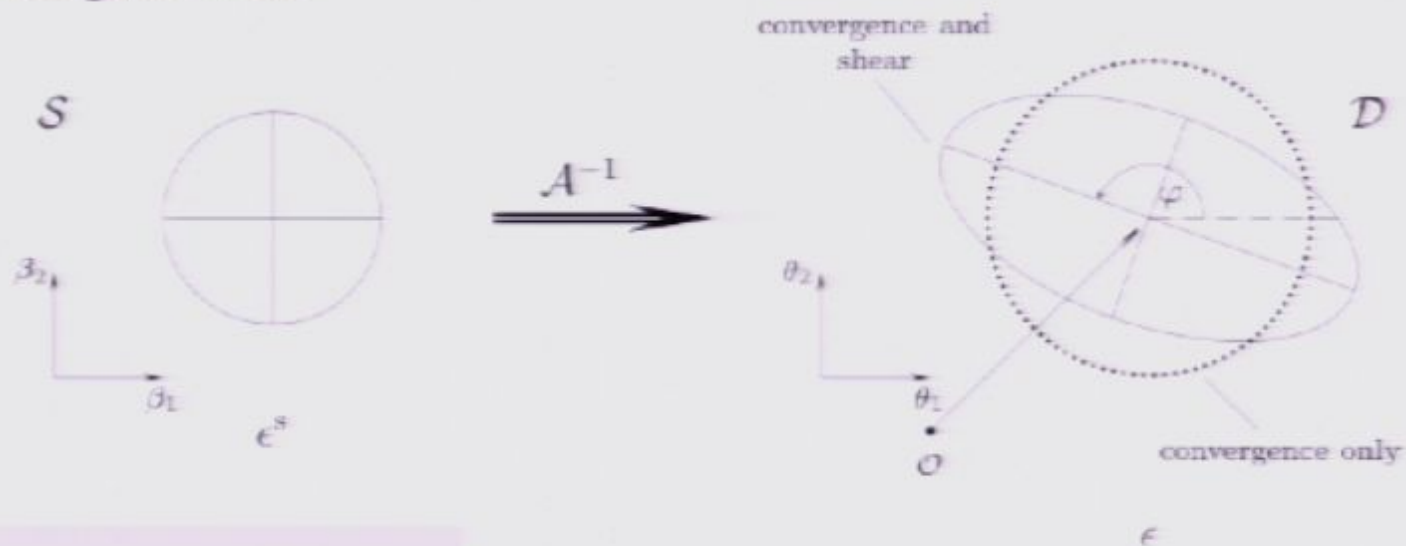
for kappa  $\ll 1$ , weak lensing regime; can make approximation:

$$\boxed{\mathcal{A}(\boldsymbol{\theta}) = (1 - \kappa) \begin{pmatrix} 1 - g_1 & -g_2 \\ -g_2 & 1 + g_1 \end{pmatrix}}, \quad \boxed{g_i = \frac{\gamma_i}{(1 - \kappa)}} : \text{reduced shear}$$

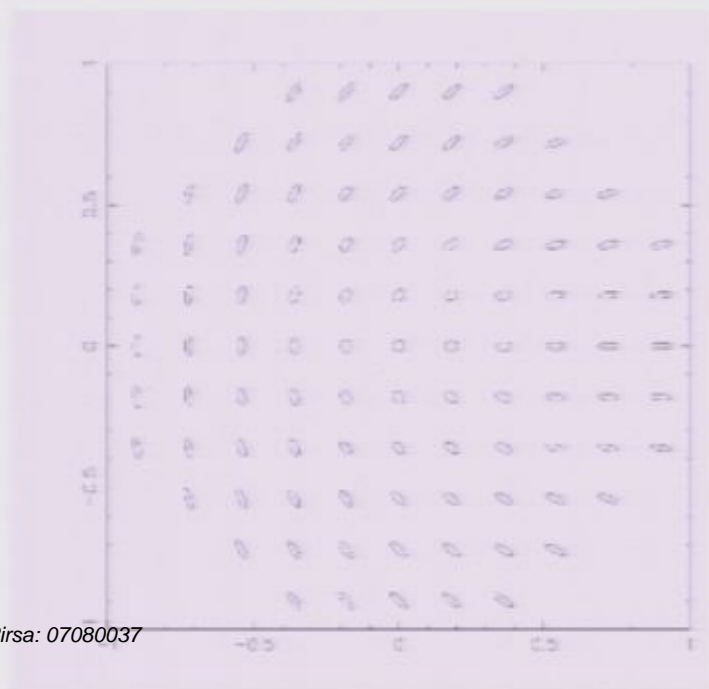
now distortion and amplification separated out

# Shear and Weak Gravitational Lensing

Shear versus magnification:

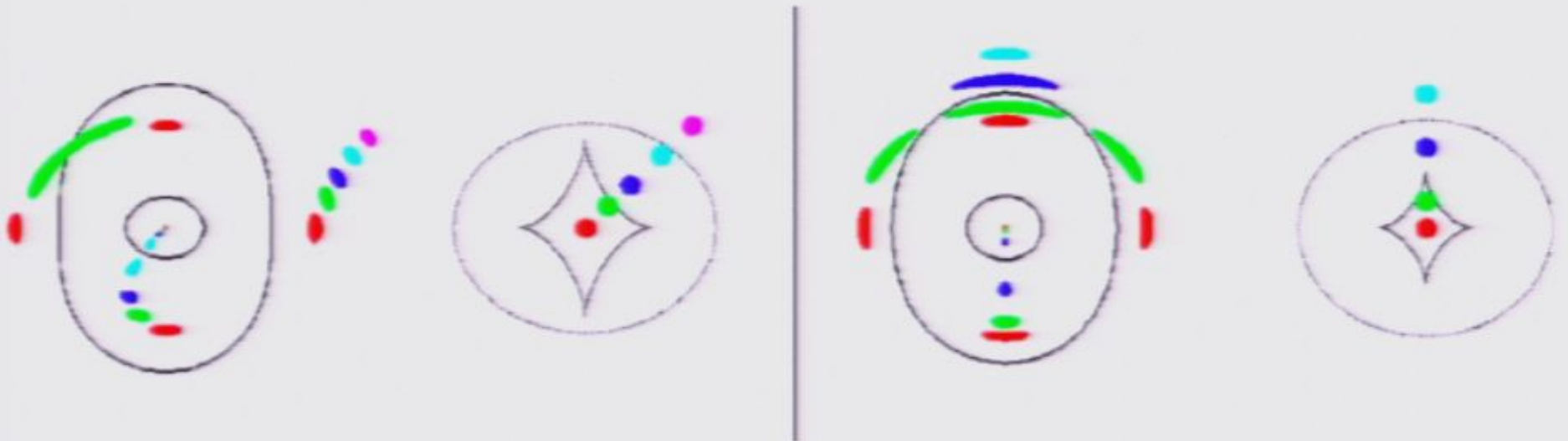


NB rather unintuitive description of shear in terms of two components, due to symmetry under 180 degree rotations



## N.B.: Strong Gravitational Lensing

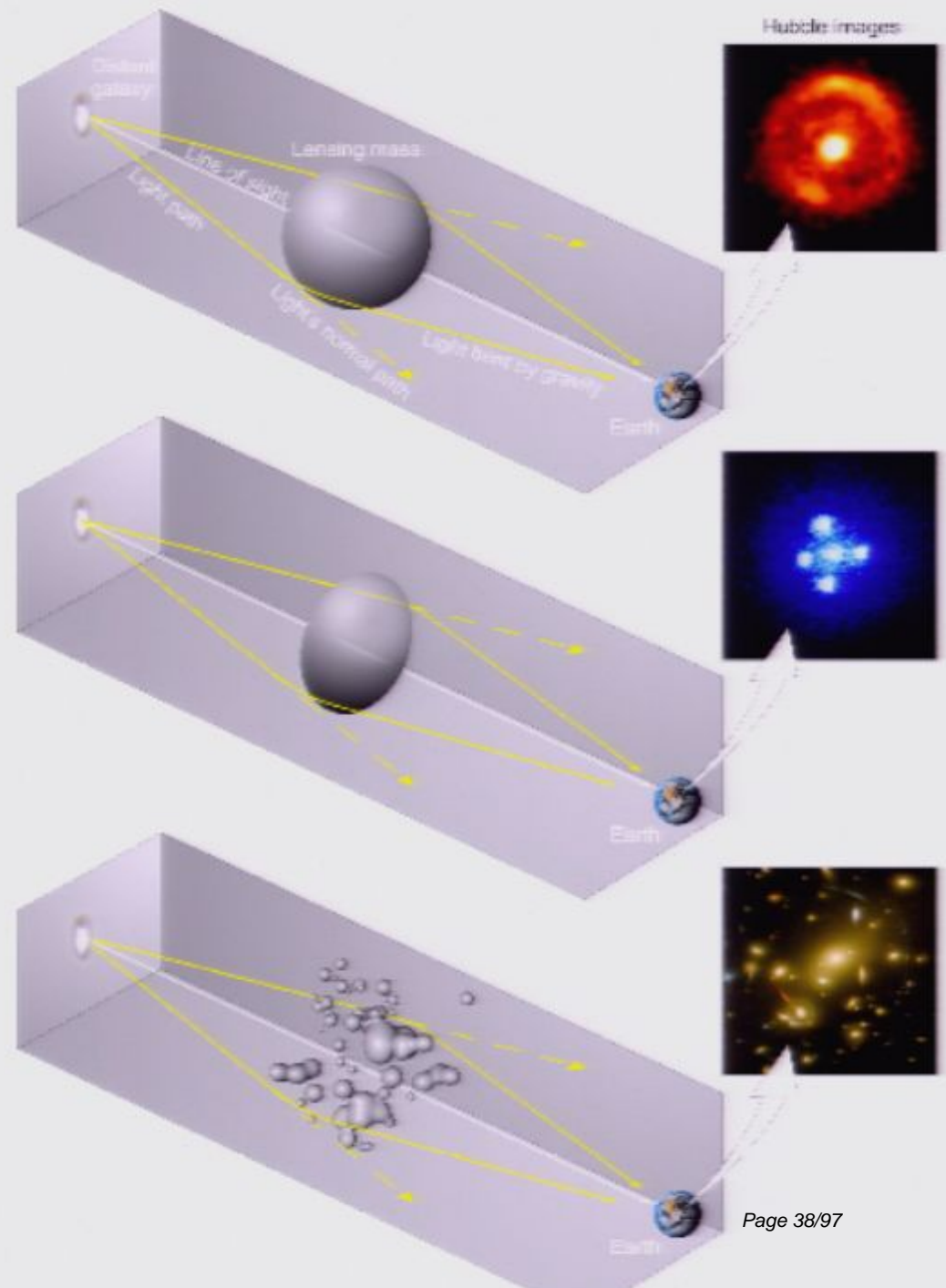
When  $\kappa$  exceeds 1, can get multiple images, as rays passing around either side of an object are bent towards the observer:



# N.B.: Strong Gravitational Lensing

Useful probe of galaxies -> measure of mass profiles, Hubble constant, searches for small scale-dark matter structure

Also very useful probe of clusters -> internal structure, dark matter substructure, searches for distant source galaxies



Strong  
Gravitational  
Lensing in  
Abell 1689



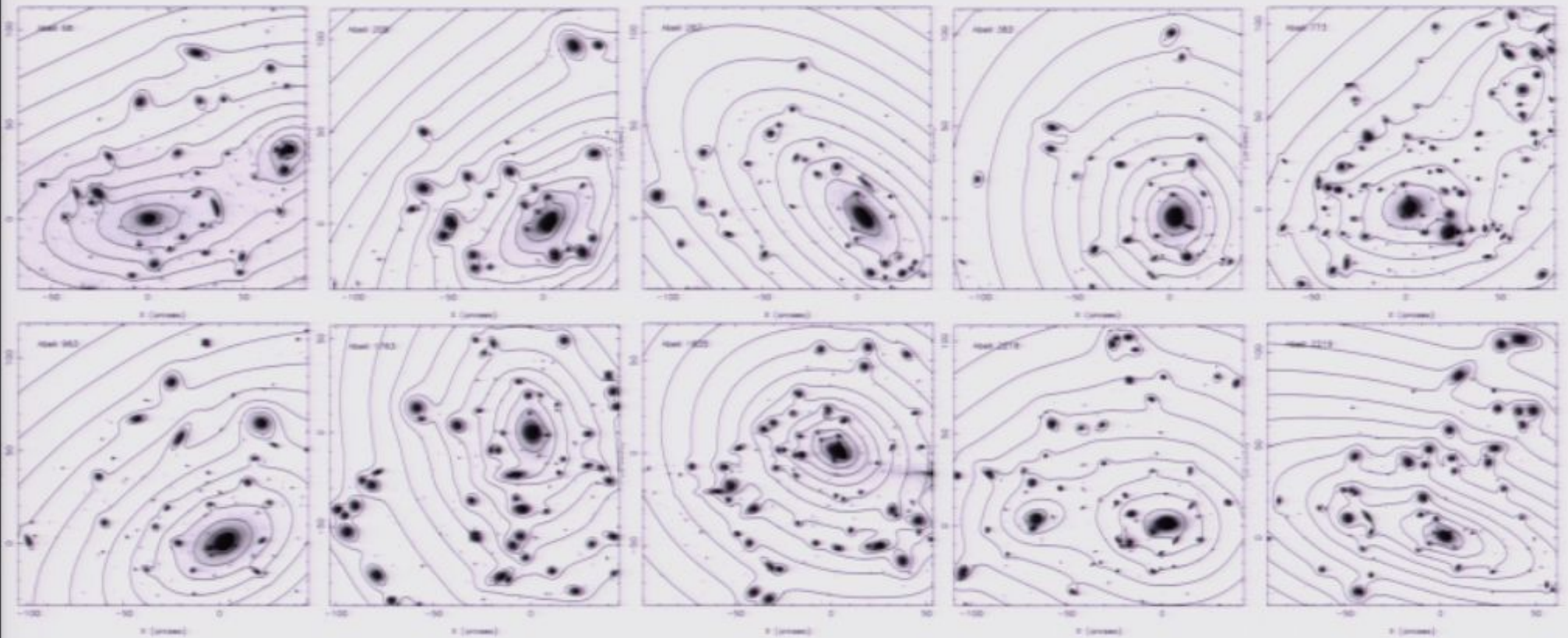
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Use strong + weak lensing to map cluster structure and substructure in detail

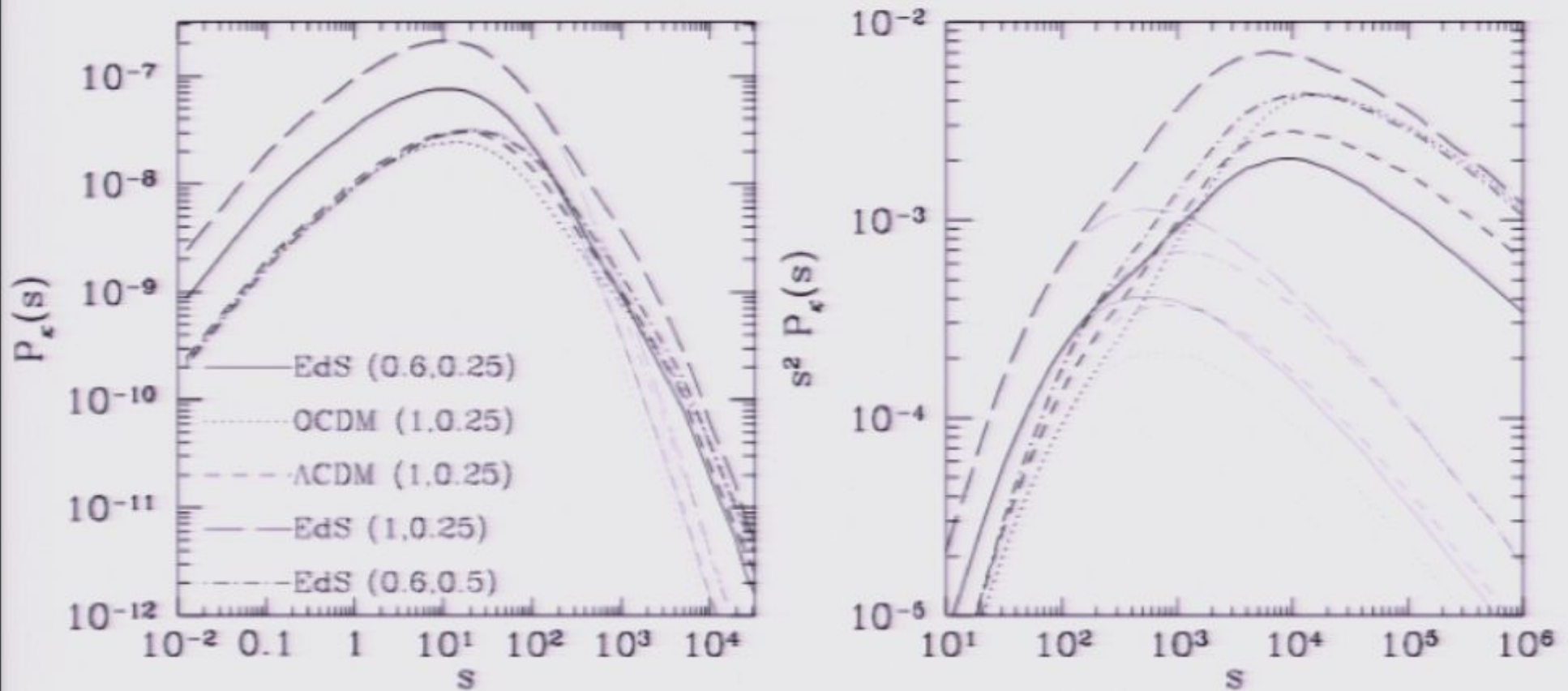
(Smith et al. 2005)



# Applications of Weak Gravitational Lensing

## Projected (Convergence/Surface Density) Power Spectrum:

In principle, calculate kappa at each point, take Fourier transform and there you are:

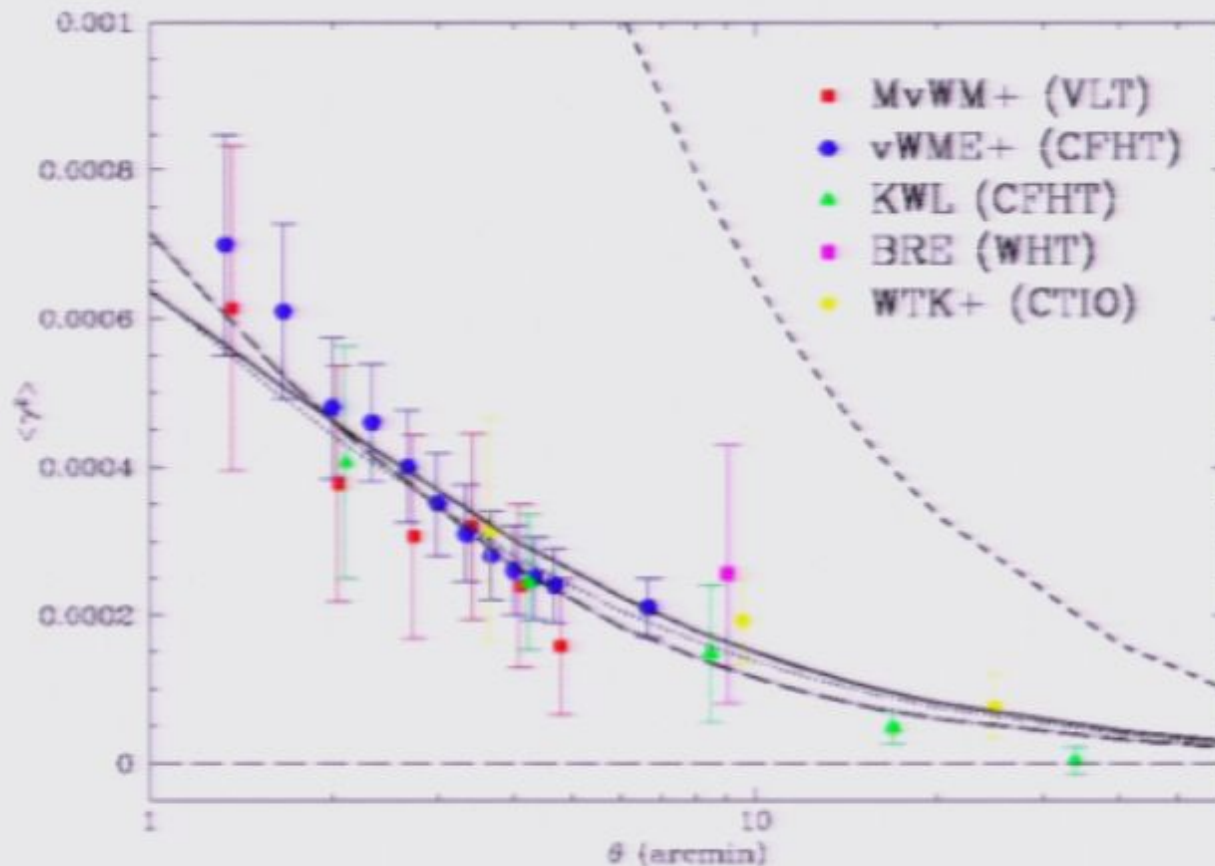


# Applications of Weak Gravitational Lensing

In practice, several other more practical measurements:

- 2-point shear correlation function
- mean shear dispersion
- aperture mass dispersion

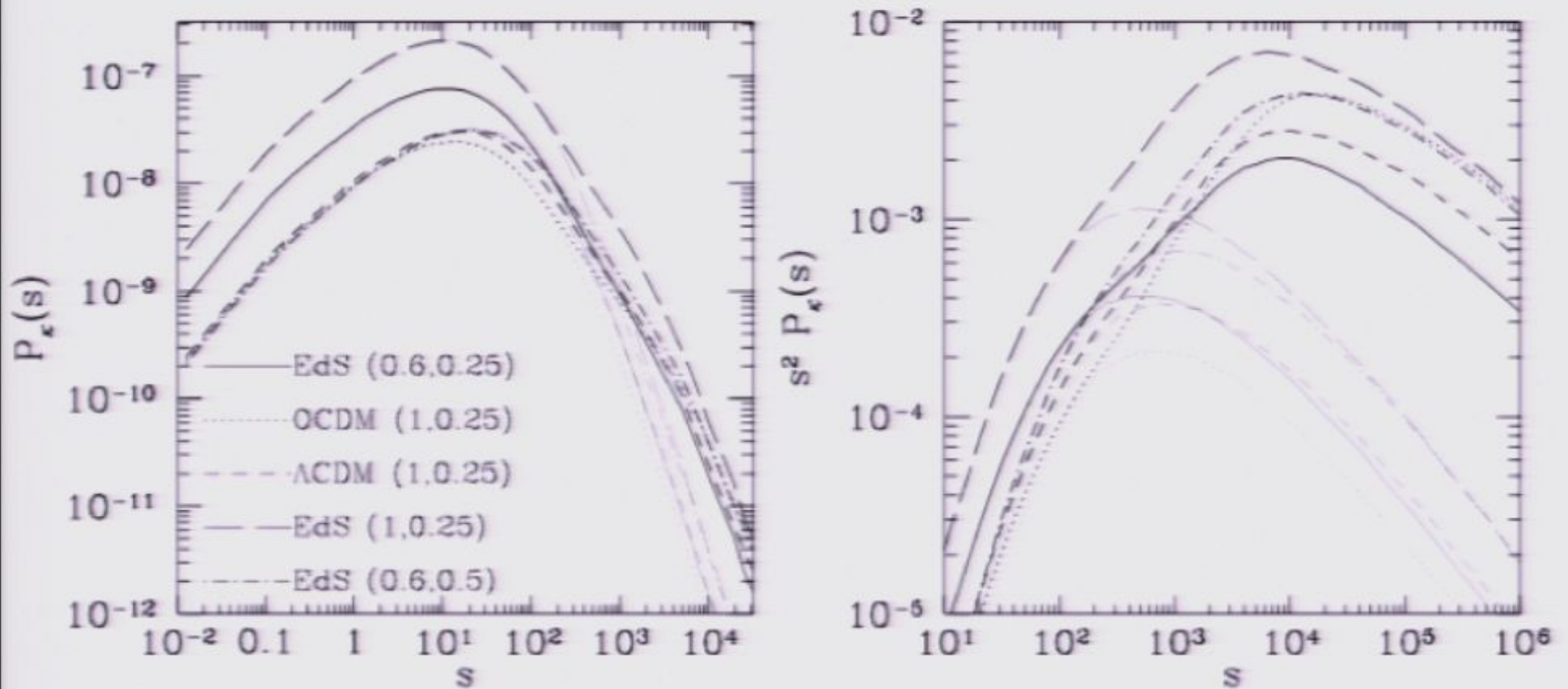
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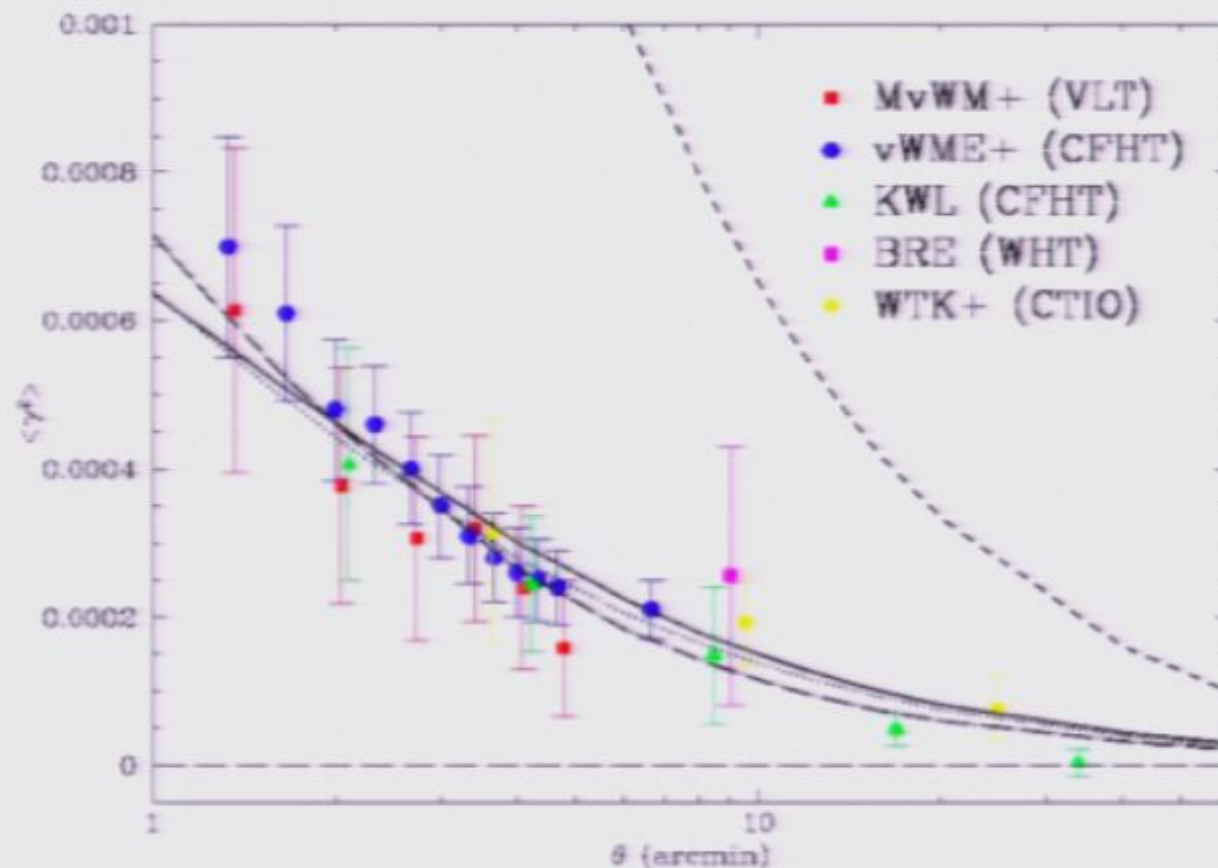


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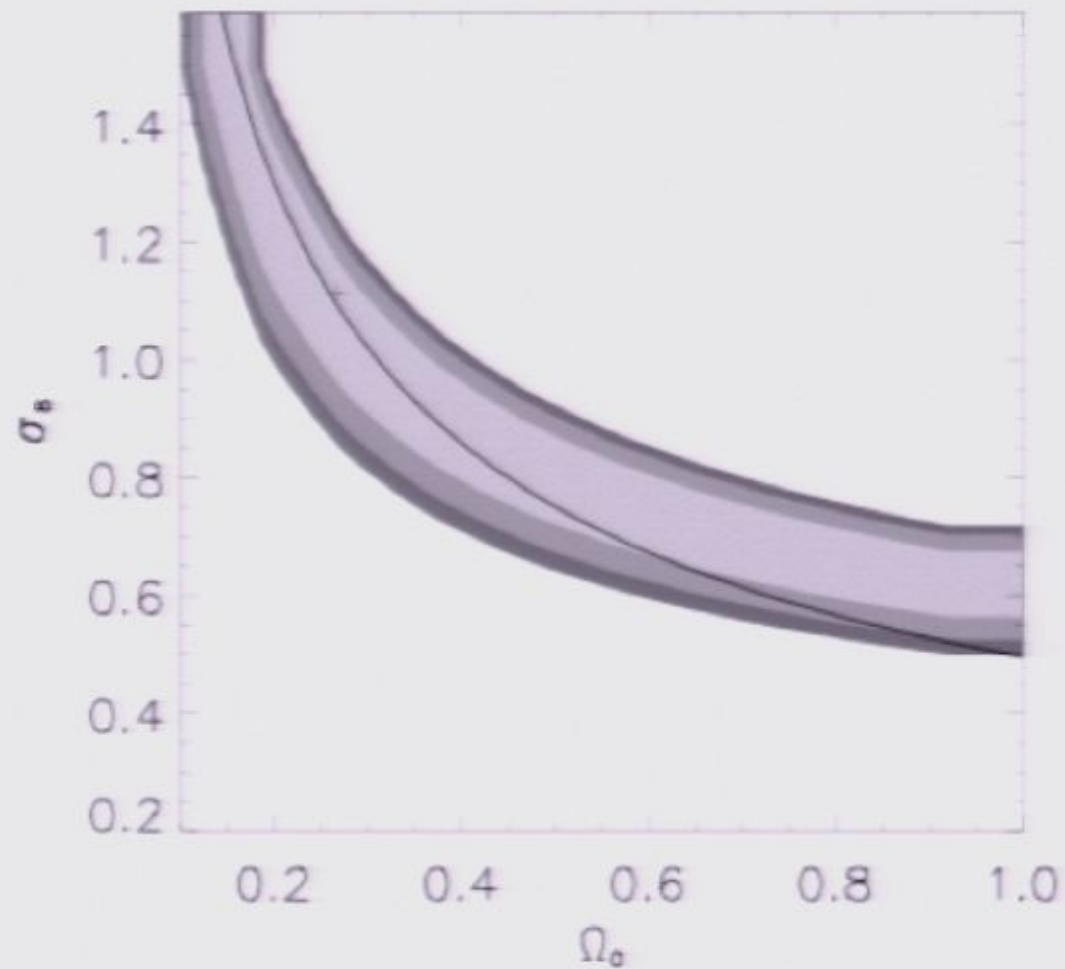
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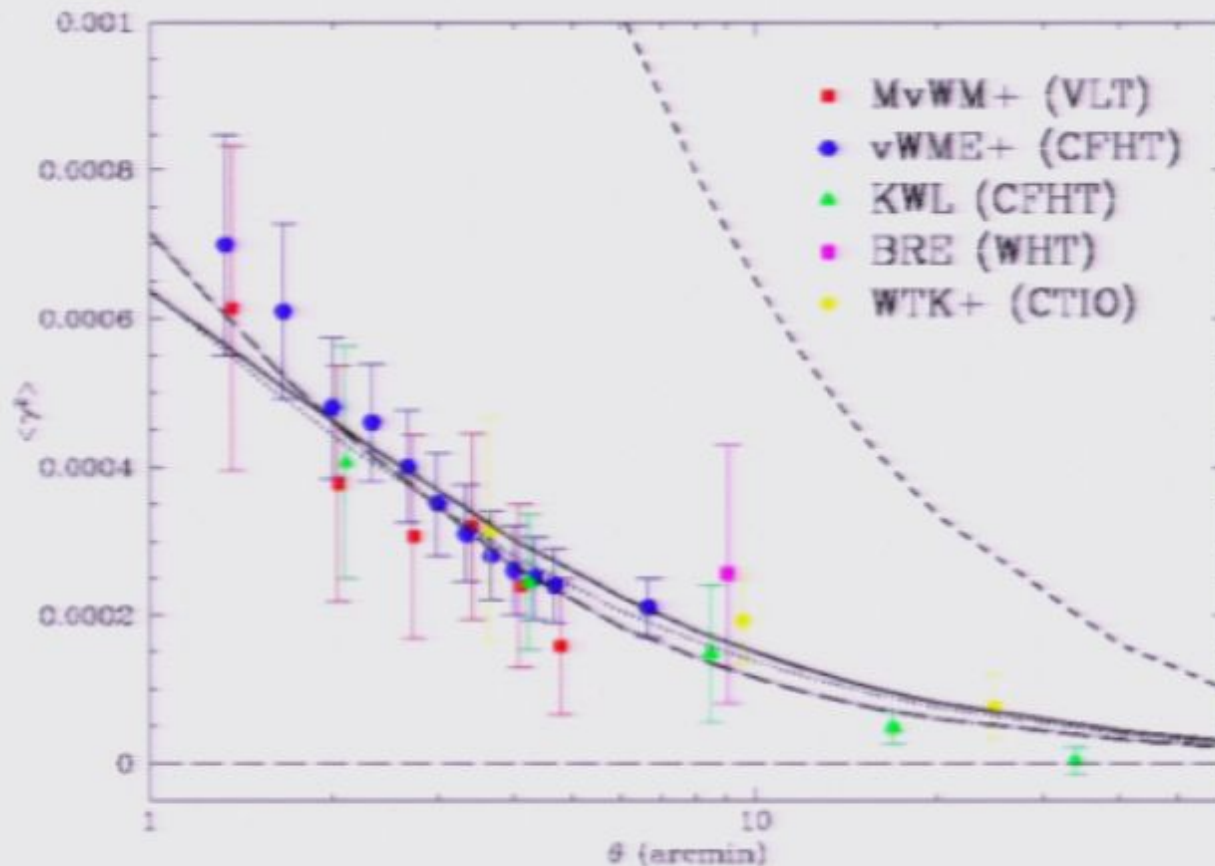


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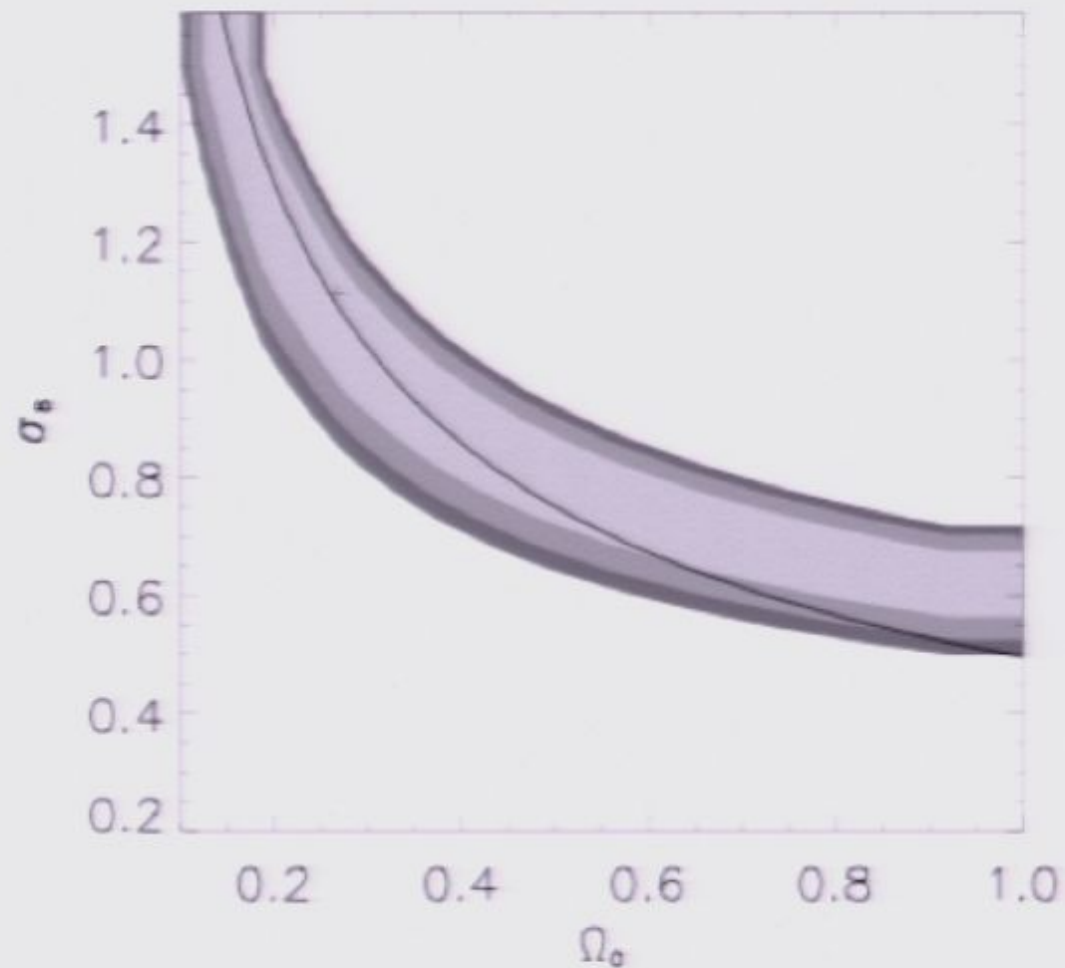
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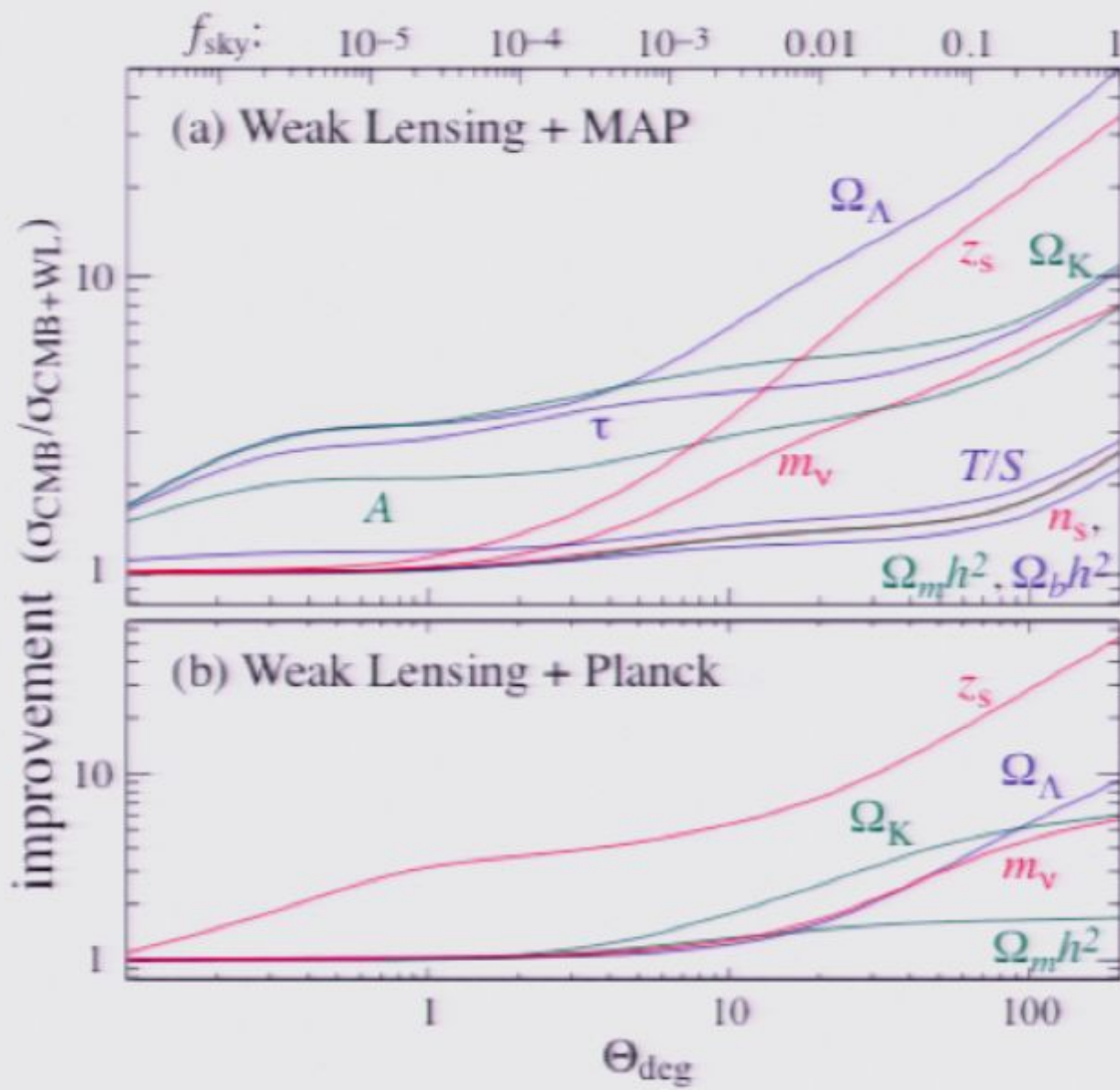
# Applications of Weak Gravitational Lensing

Why not just use the CMB?

- weak lensing measure the power spectrum at low redshifts and on small scales, so very different from the CMB
- low redshifts => sensitivity to dark energy
- small scales => (only) direct probe of dark matter distribution
- growth factor versus expansion history tests modifications to gravity
- even in conventional cosmologies, benefit of long lever arm

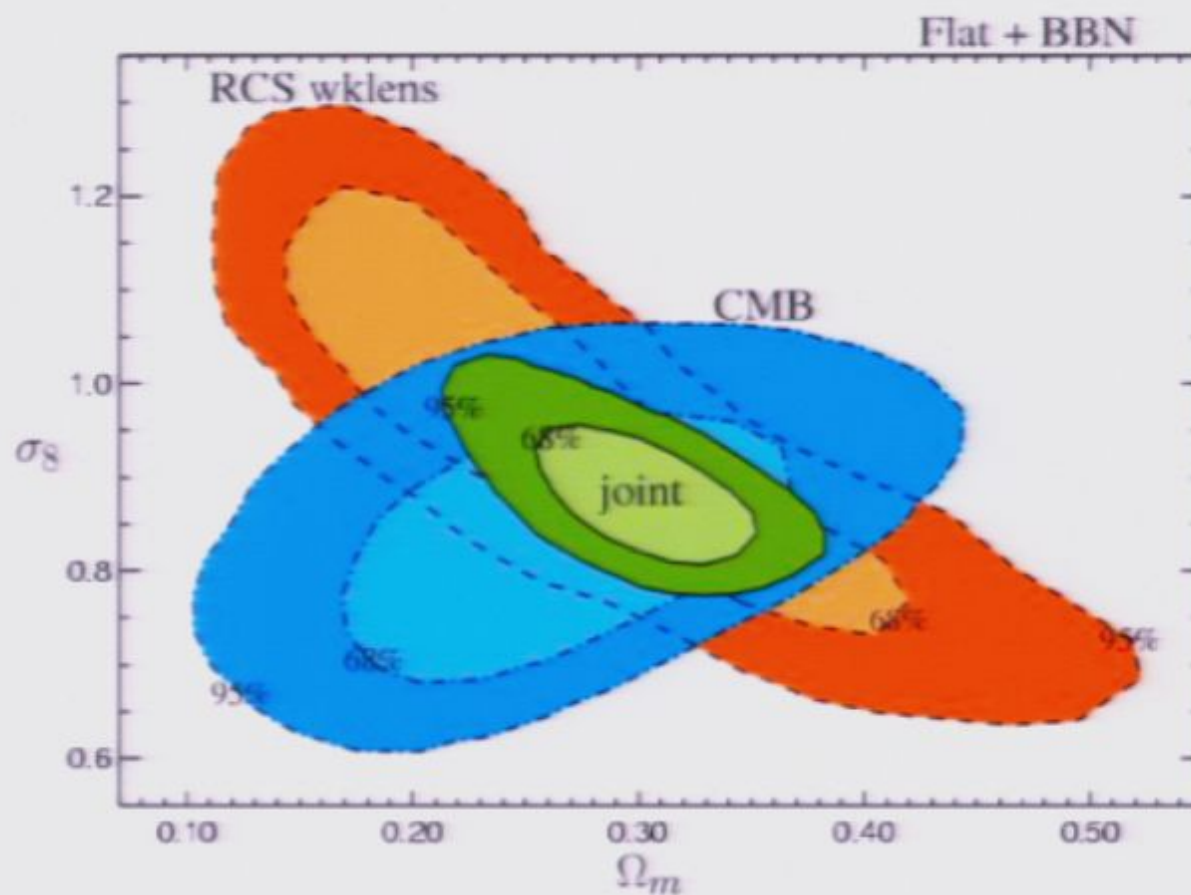
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# Applications of Weak Gravitational Lensing

Why not just use the CMB?



1E 0657-56 (The 'bullet' cluster)  
– a smoking gun?















But see Angus et al. 2007 (astro-ph/069125);  
Brownstein & Moffat 2007 (astro-ph/072146)

# Probes of structure growth using gravitational lensing:

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### 1) Cosmic shear from weak lensing

← linear and non-linear fluctuations on a range of scales  
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⇐ abundance of high peaks in the density field  
(degenerate in  $\sigma_8$  and growth)

3) Differential growth of non-linear structure

(interesting, but may be hard to calibrate?)

# Cosmic Shear, etc.

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Mass overdensity produces tangential distortion of light rays from background objects; underdensity produces radial distortion.



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overdense



underdense



B-mode  
(systematic)



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overdense



underdense



B-mode  
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Effect always small, so need to average over 100s of objects (i.e. galaxies)  
⇒ limit to scales probed, intrinsic shape noise.

# Cosmic Shear, etc.

Mass overdensity produces tangential distortion of light rays from background objects; underdensity produces radial distortion.

overdense



underdense



B-mode  
(systematic)



Effect always small, so need to average over 100s of objects (i.e. galaxies)  
⇒ limit to scales probed, intrinsic shape noise.

\* Can look at resulting signal statistically: shear correlation measuring correlation of overdense/underdense regions as a function of angular scale.

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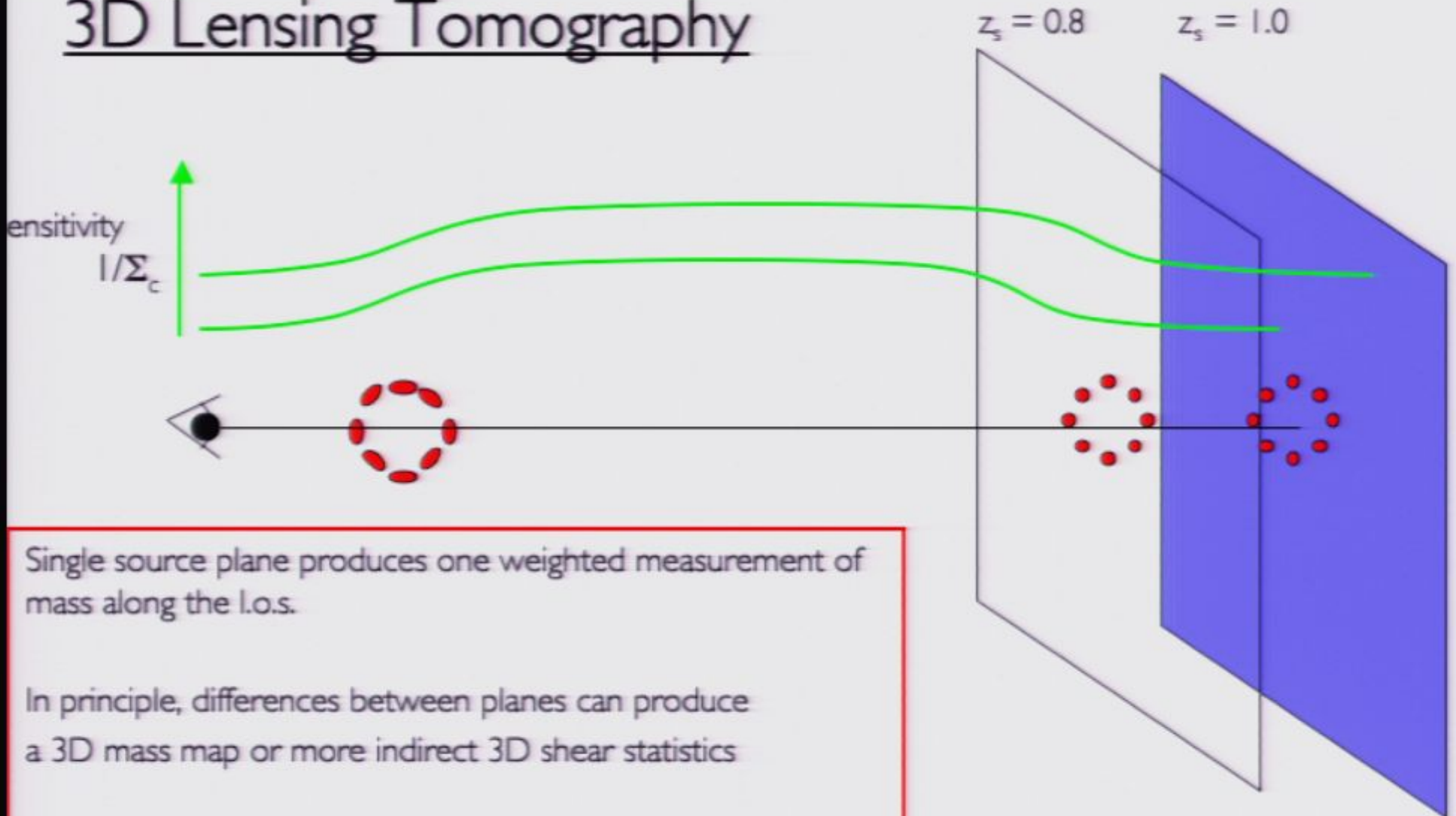
\* Can look at resulting signal statistically: shear correlation measuring correlation of overdense/underdense regions as a function of angular scale.

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Redshift dependence of effect is very slow, producing a very broad kernel in redshift space:

$$\frac{4\pi G}{c^2} \int \frac{D_L D_{LS}}{D_S} \rho dl$$

# 3D Lensing Tomography



Single source plane produces one weighted measurement of mass along the l.o.s.

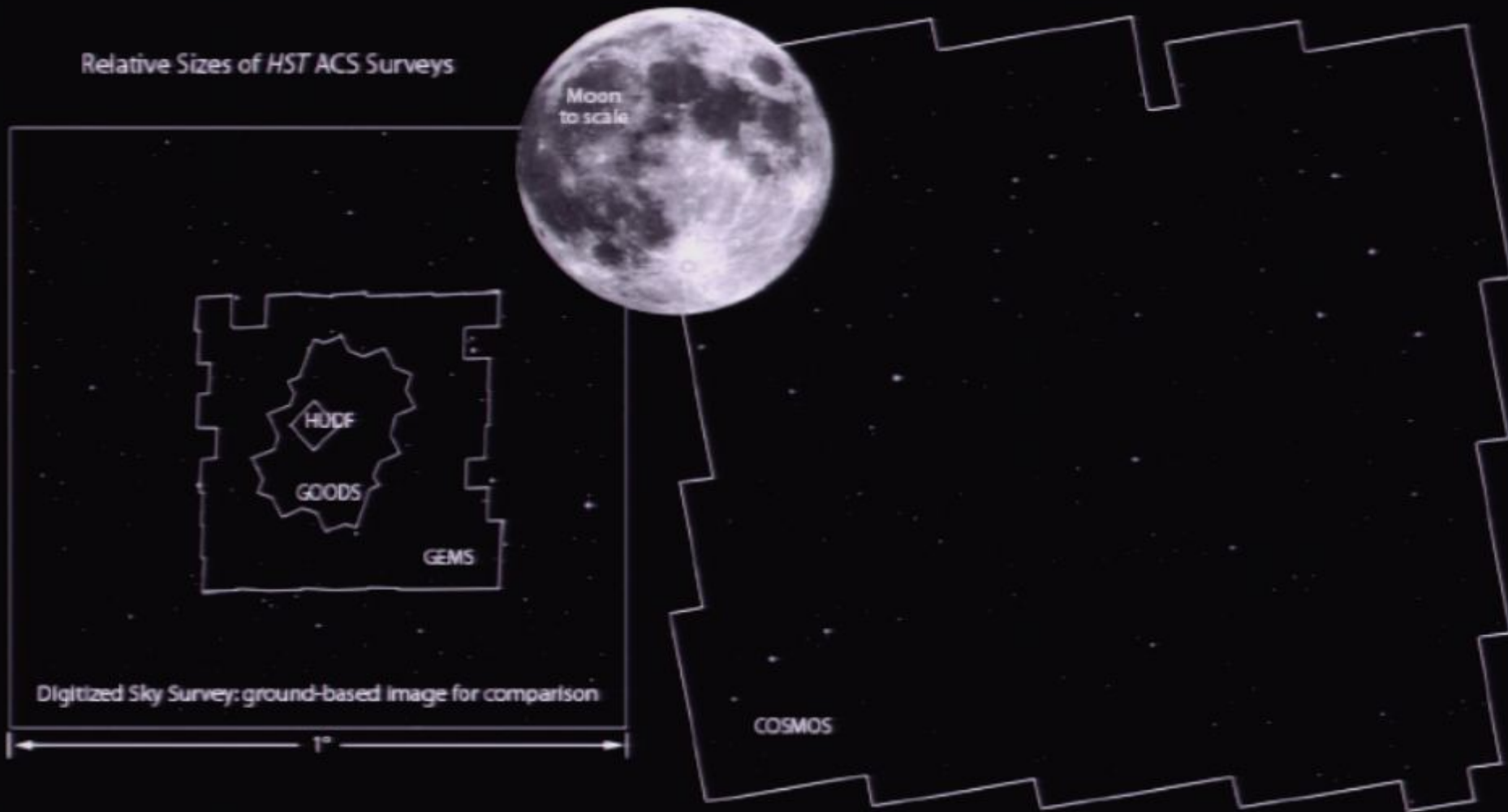
In principle, differences between planes can produce a 3D mass map or more indirect 3D shear statistics

Some problems from extent of redshift kernel, inaccuracies in photo-zs

# The COSMOS Survey

P.I. Nick Scoville

Relative Sizes of HST ACS Surveys



# The COSMOS Survey

- 2 square degree ACS mosaic
- current lensing results from 1.64 square degrees
- 2-3 million galaxies down to  $F814W_{AB} = 26.6$
- 15-band photometry, photo-zs with  $dz \sim 0.03(1+z)$  to  $z = 1.4$  and  $I_{F814W} = 24$
- follow-up in X-ray, radio, IR, UV

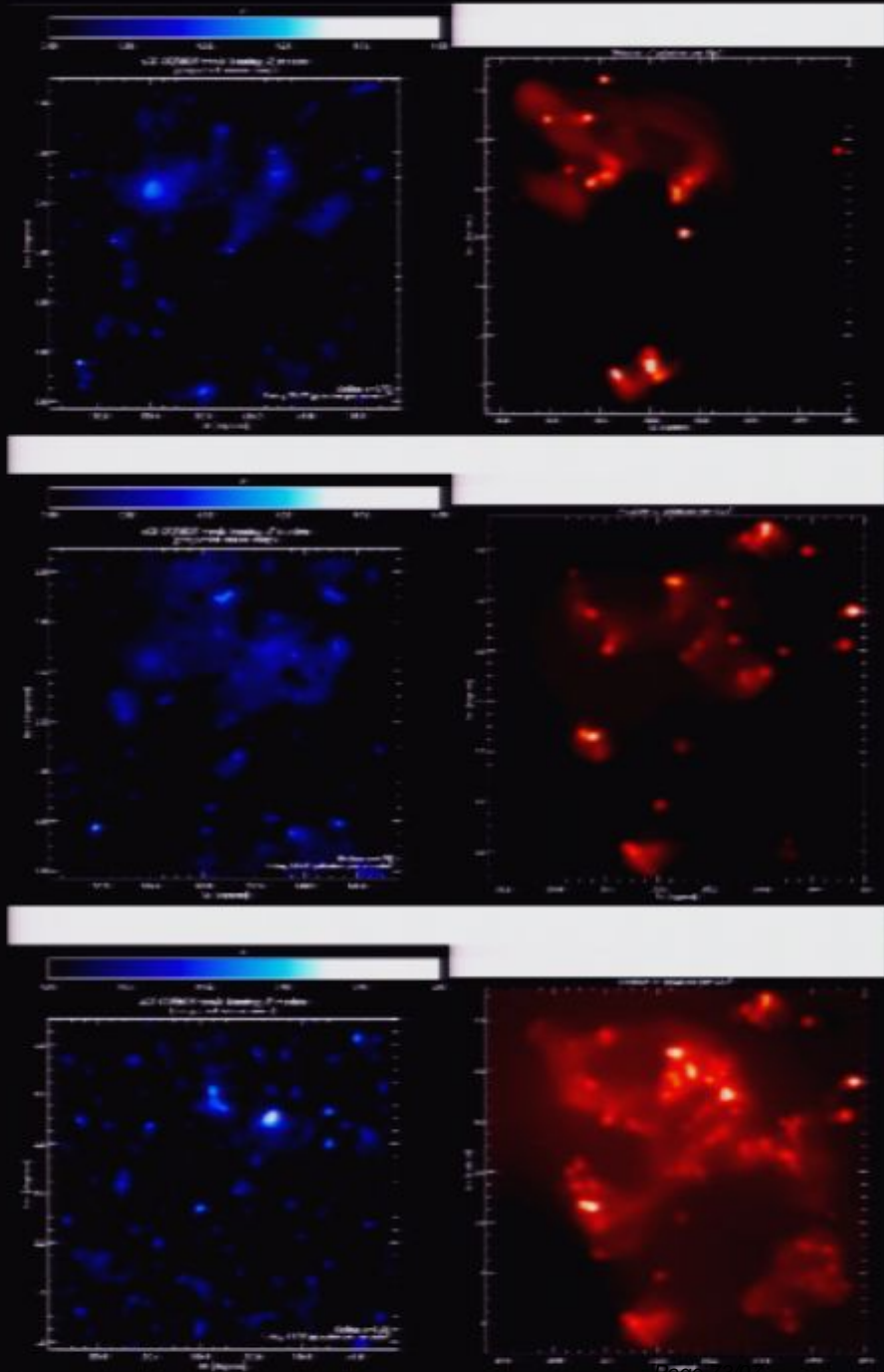




# WL Convergence Maps

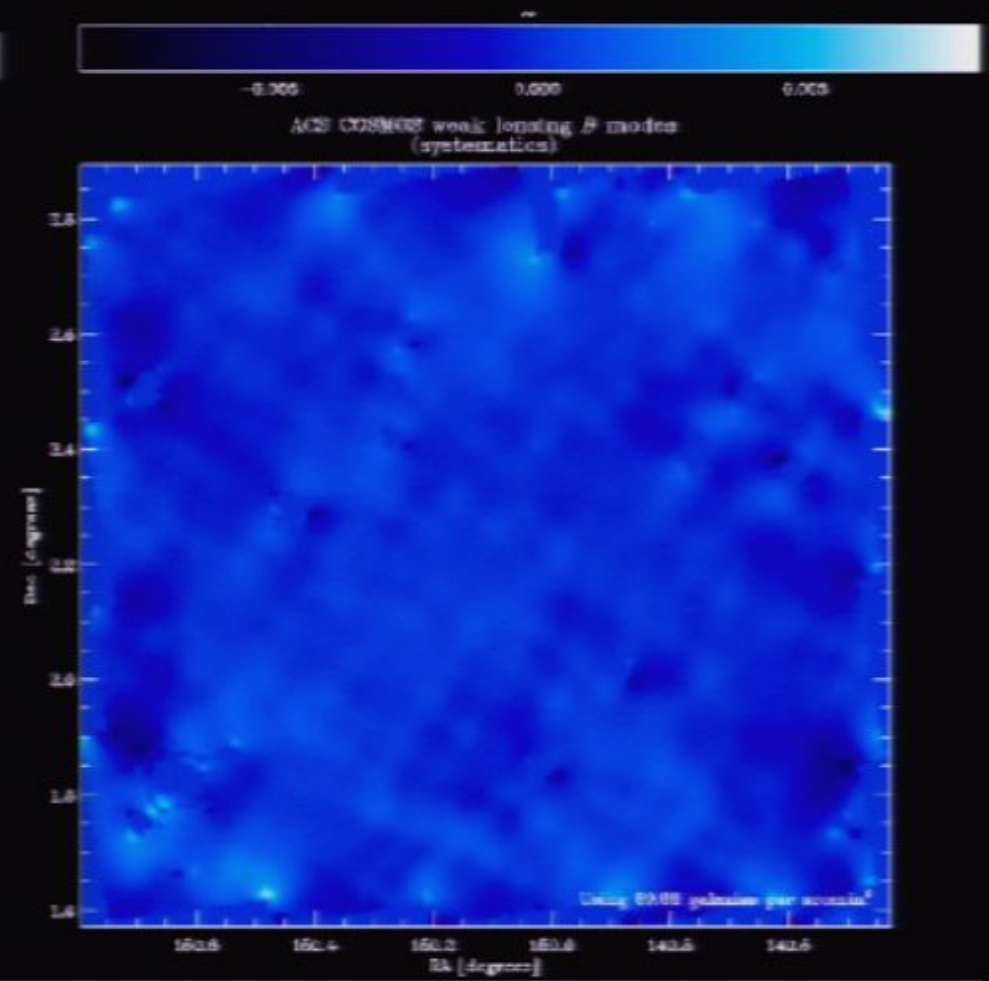
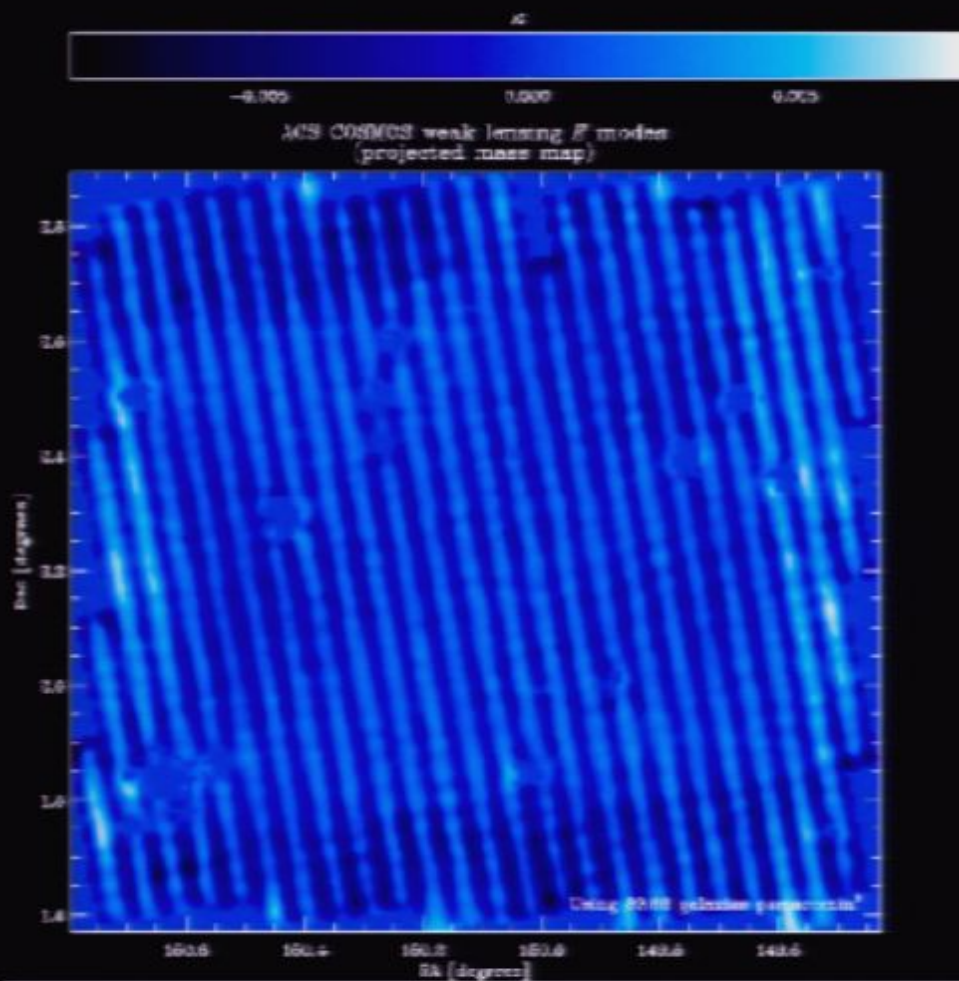
Massey, Rhodes, Leauthaud  
Capak, Koekemoer, Scoville, Refregier

- cut catalogue down to 40 galaxies/arcmin<sup>2</sup> to remove bad zs
- correct for PSF variations, CTE
- Get lensing maps, low-resolution 3D maps, various measures of power in 2D and restricted 3D
- results compare well with baryonic distributions (e.g. galaxy distribution)

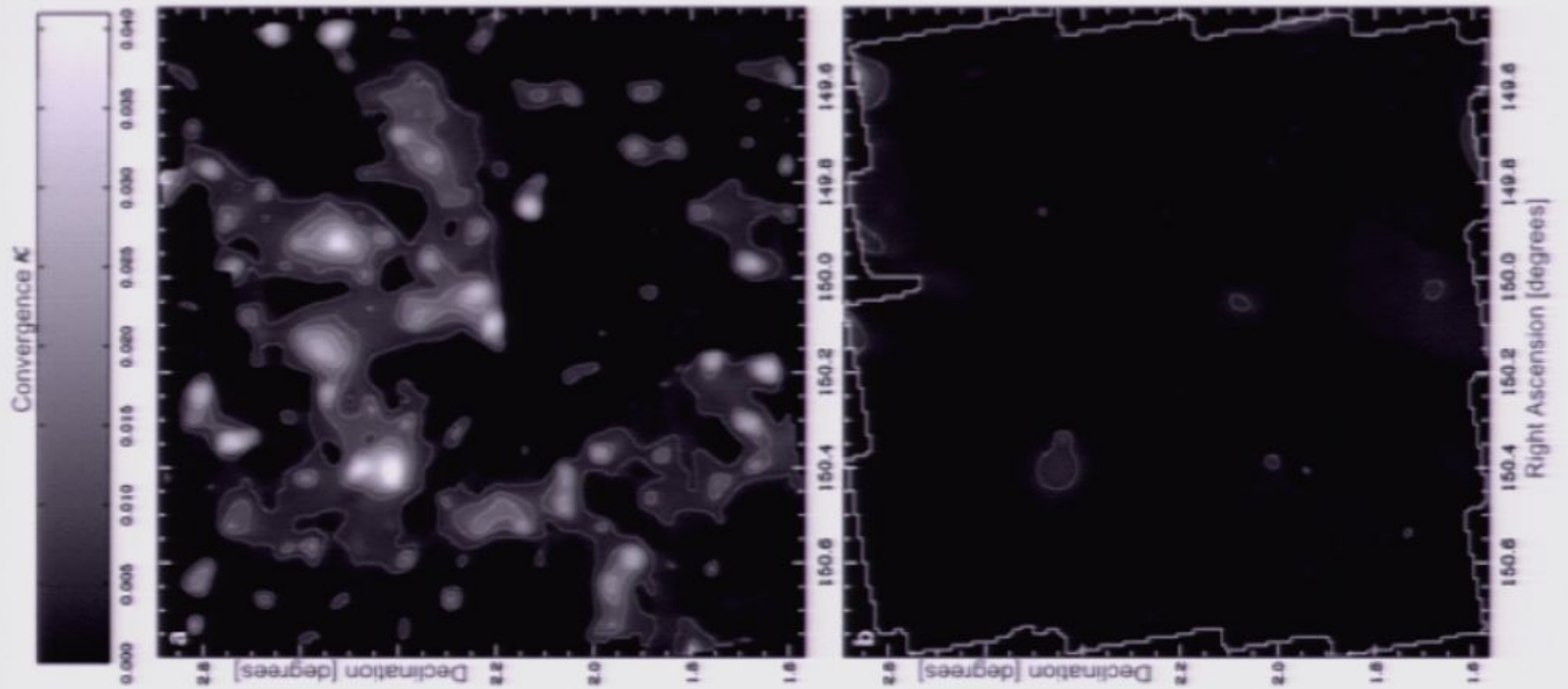


# Systematics

instrumental systematics a major and unanticipated headache  
(PSF variations, CTE)



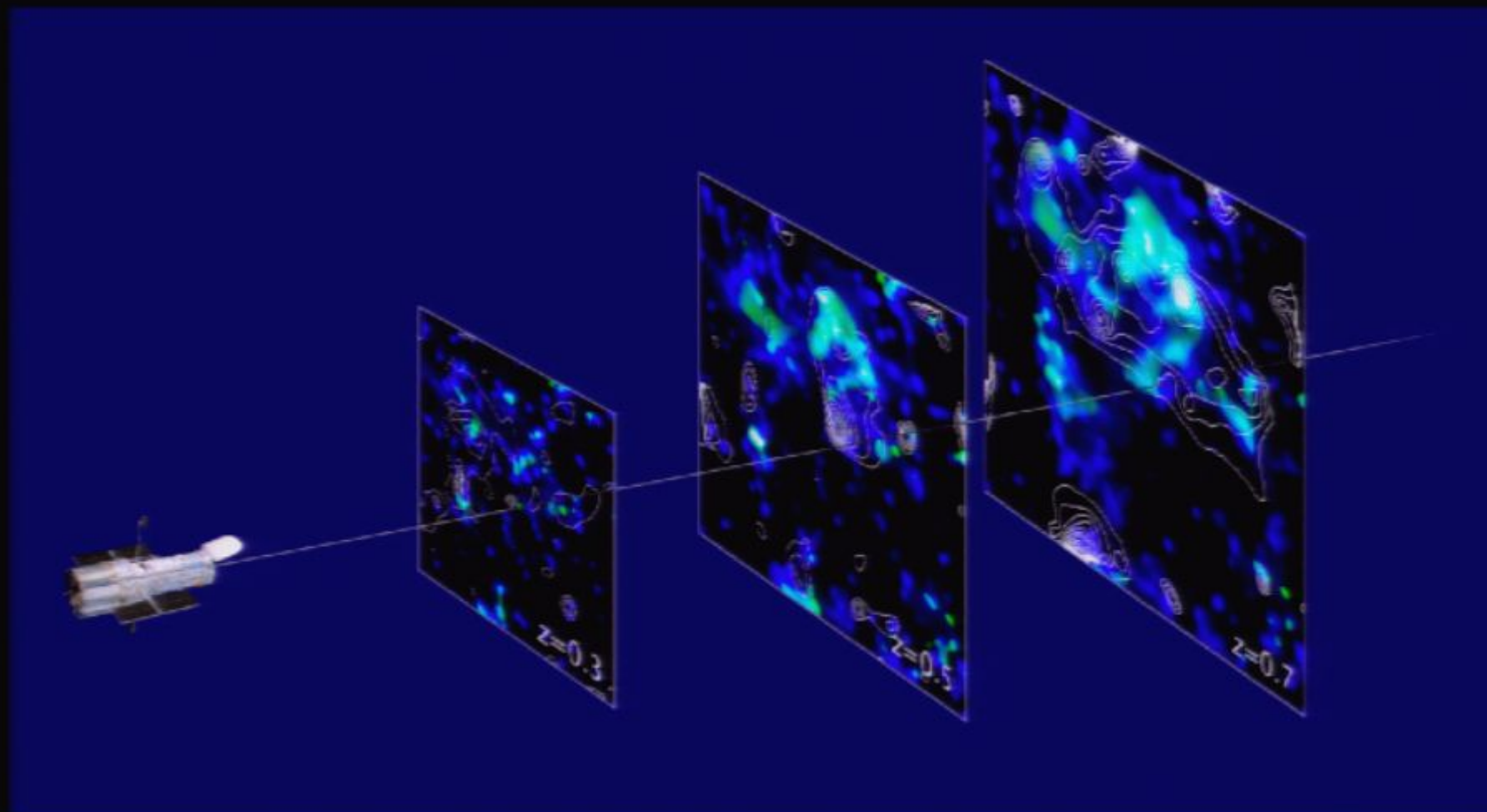
# The Final Result



E-modes (left) versus B-modes (right)

# Seeing the Growth of Structure Directly

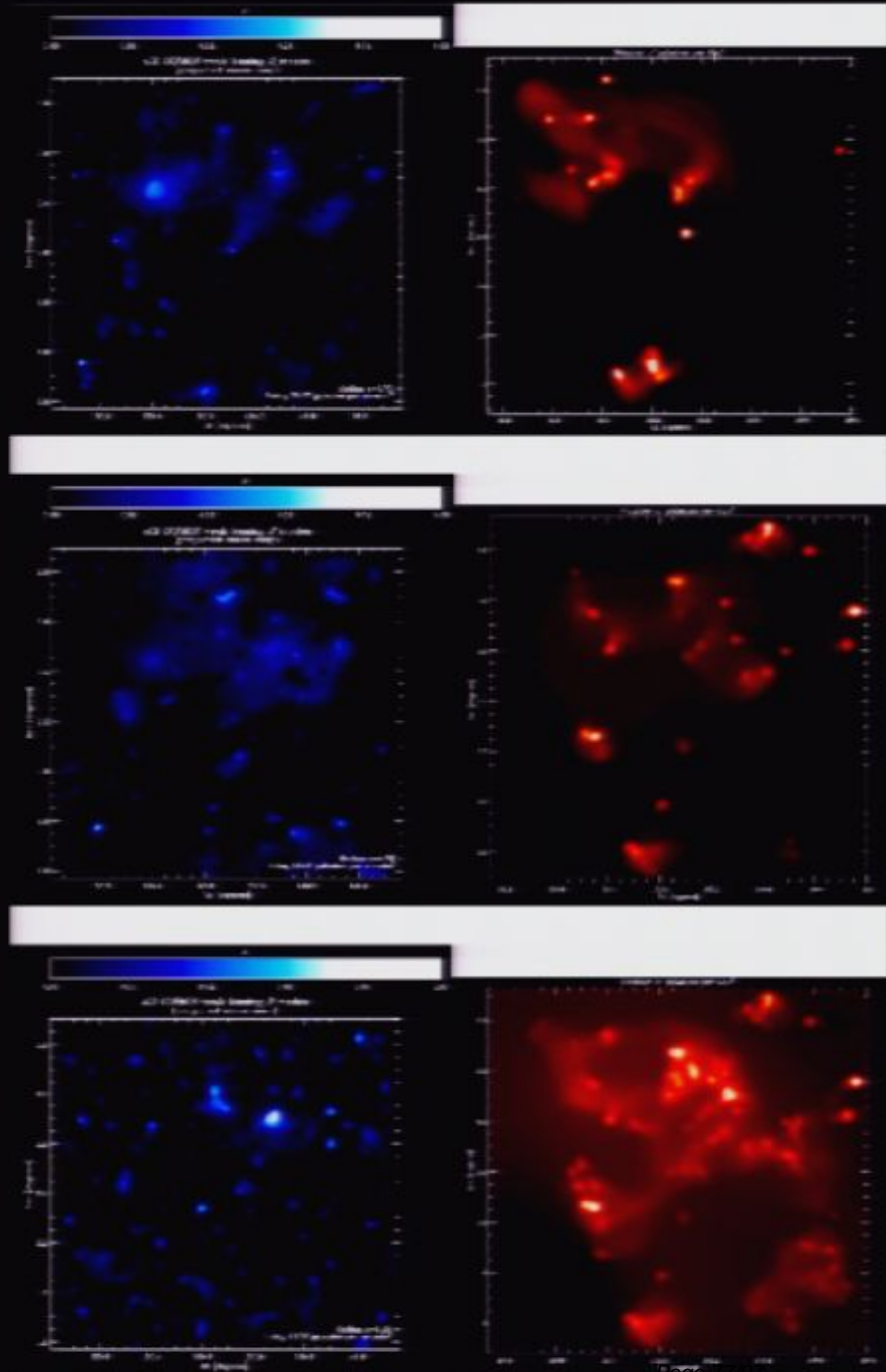
actually quite complicated to show the growth of structure going forward in time



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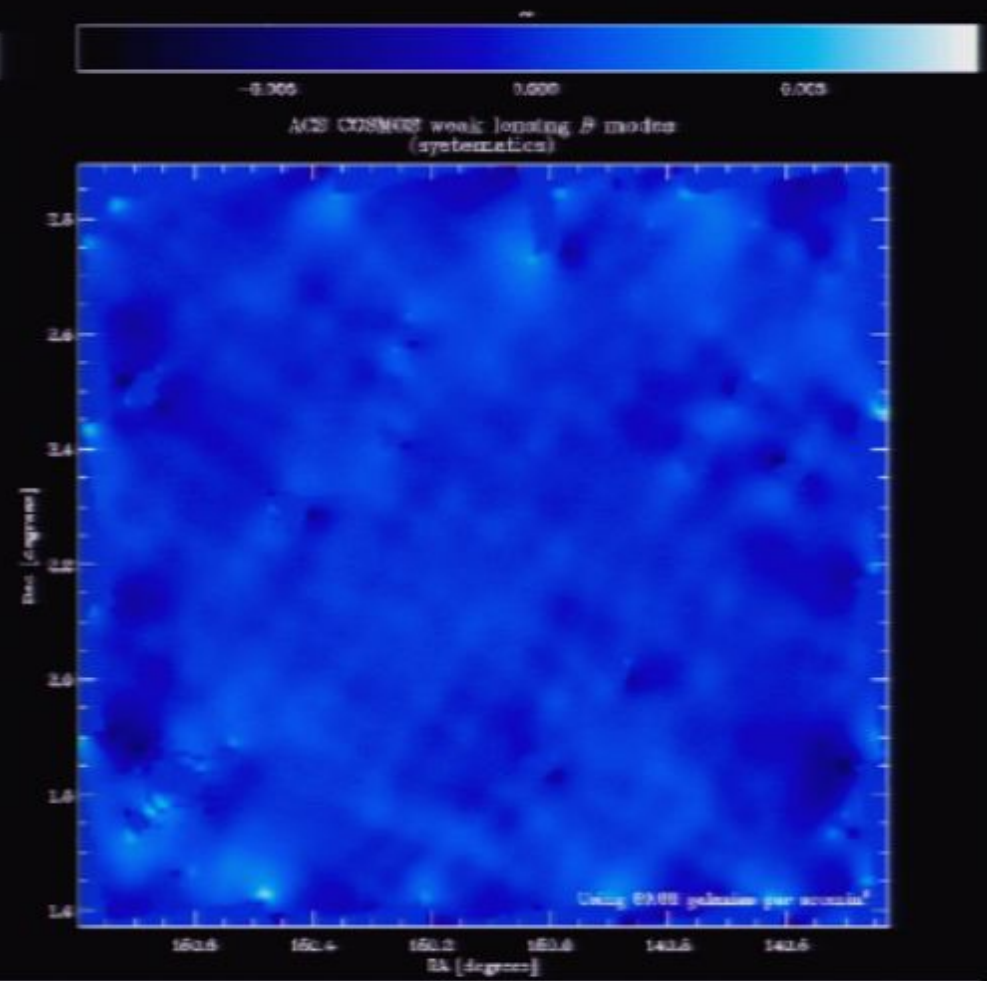
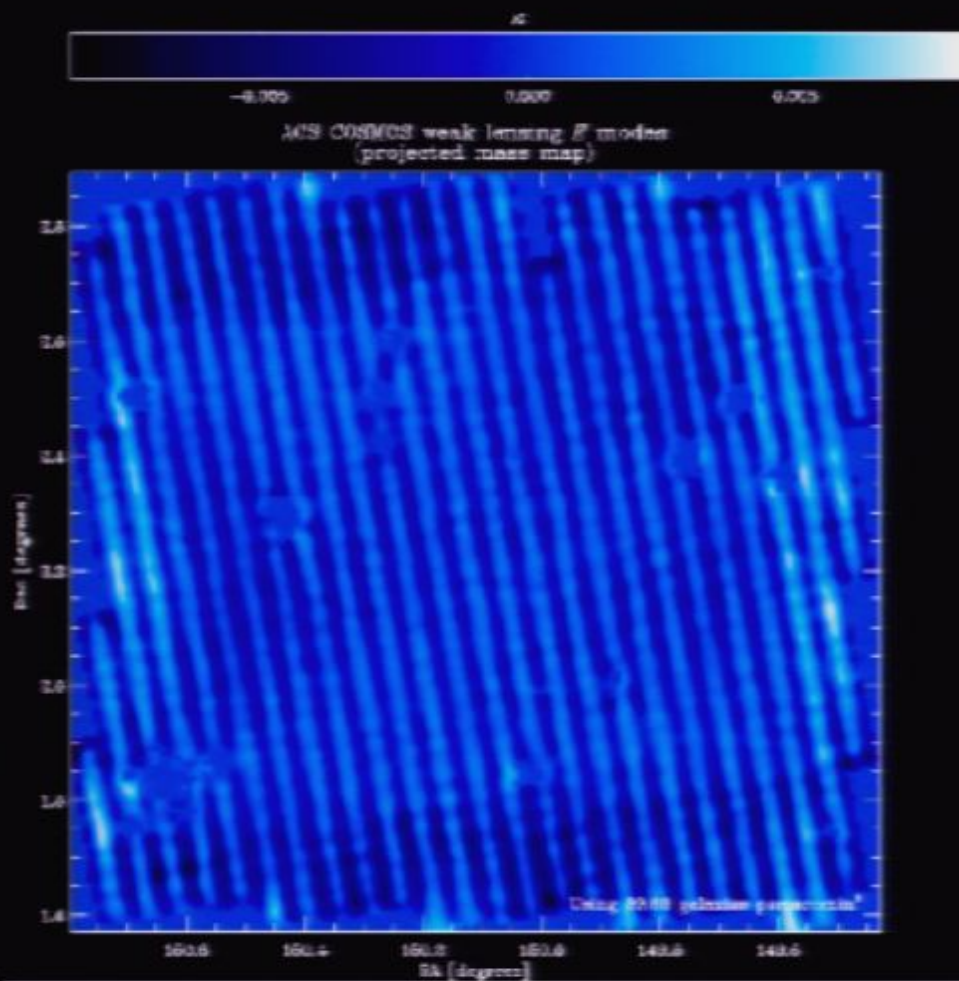
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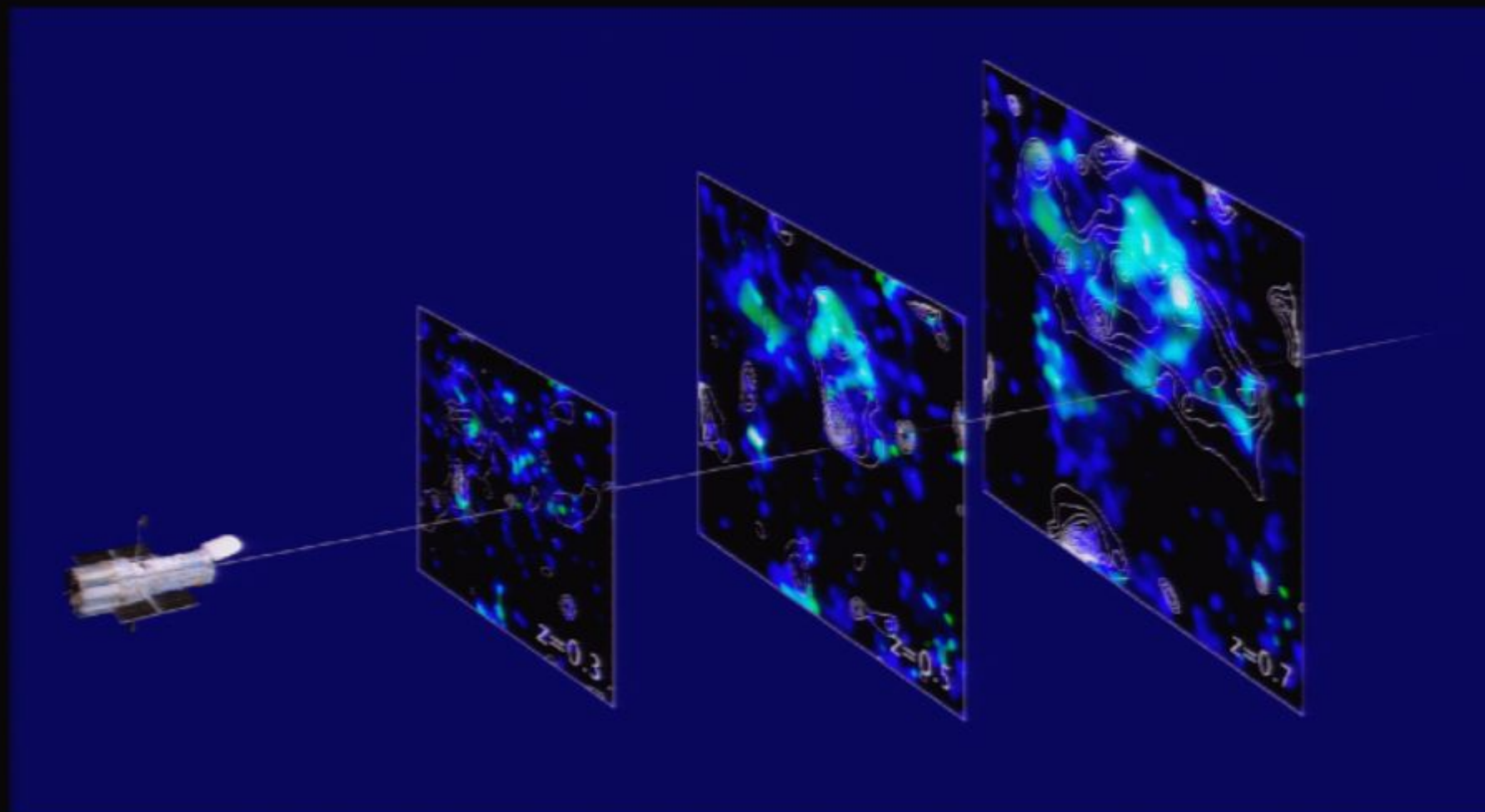
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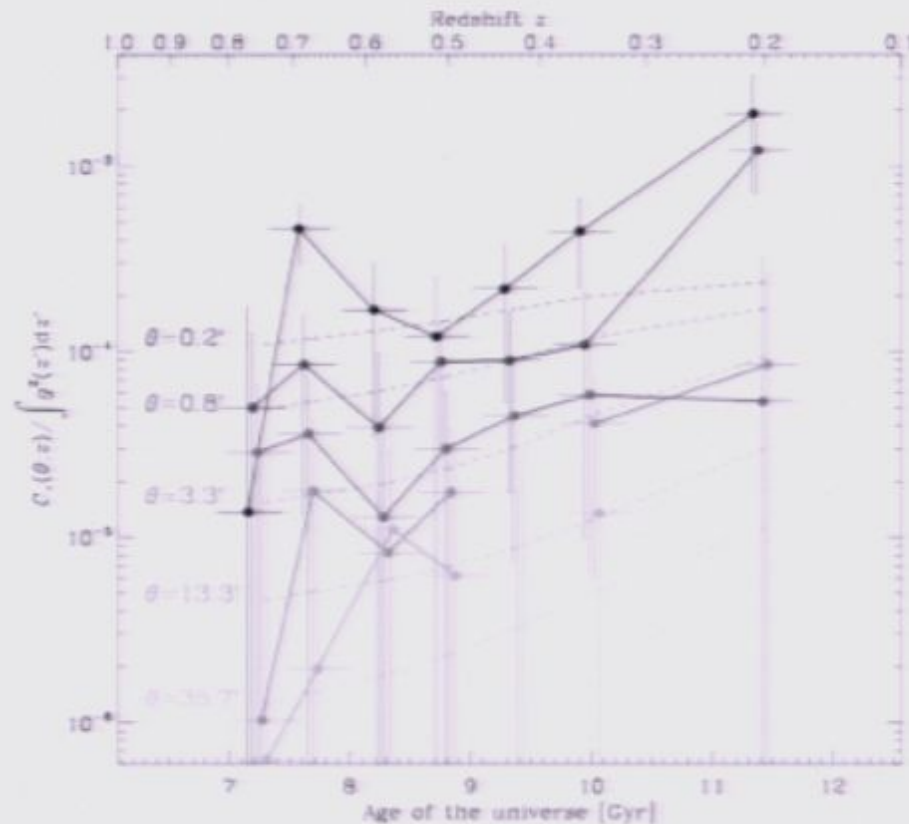
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# Measuring the Growth of Structure

Calculate the shear correlation function divided by the integrated lensing sensitivity to a given slice, and rescale to a fixed physical scale, assuming the signal is coming from the redshift of peak sensitivity.

Get (rather noisy) measure of the growth of structure:

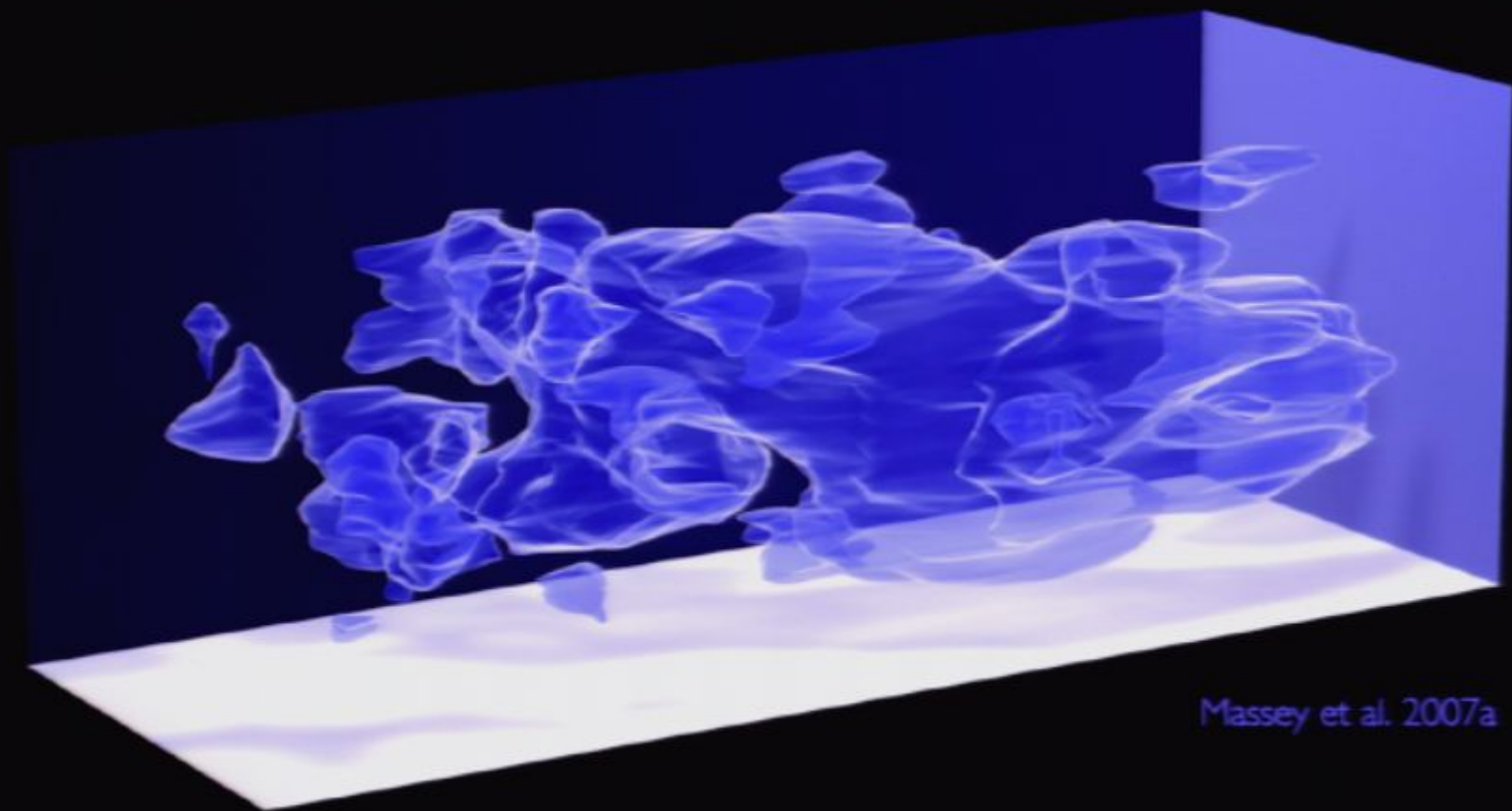


Massey et al. 2007b



## 3D Mass Distribution

Final result: first 3D map of the mass (or potential) distribution in a large volume.  
Note this is only 1.6 square degrees on the sky.

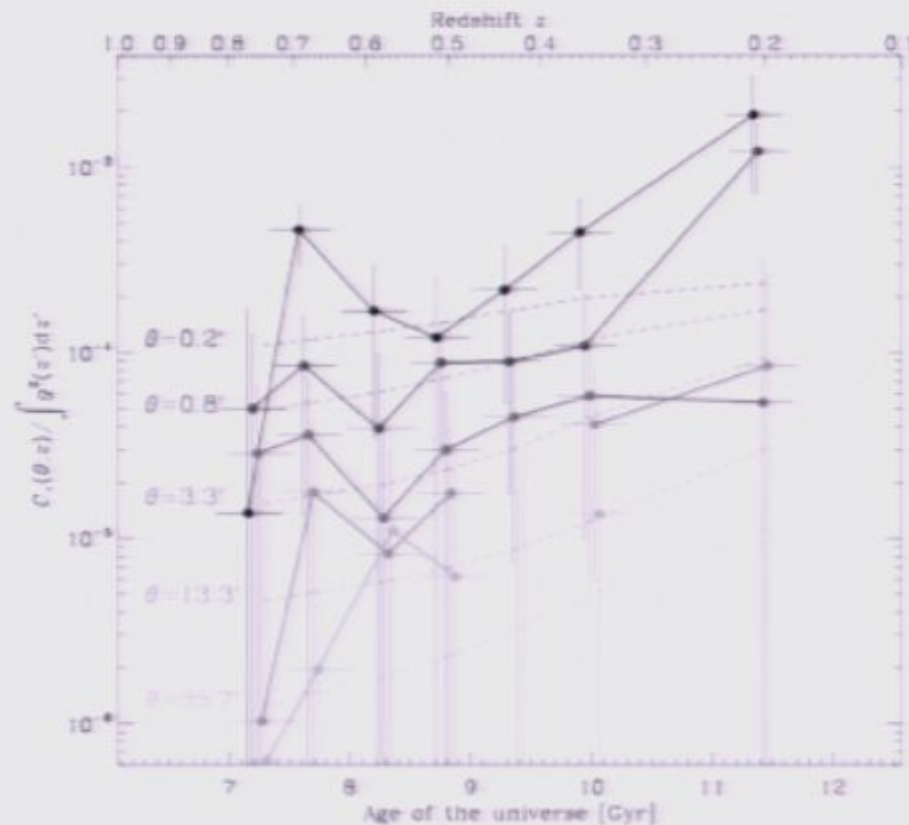


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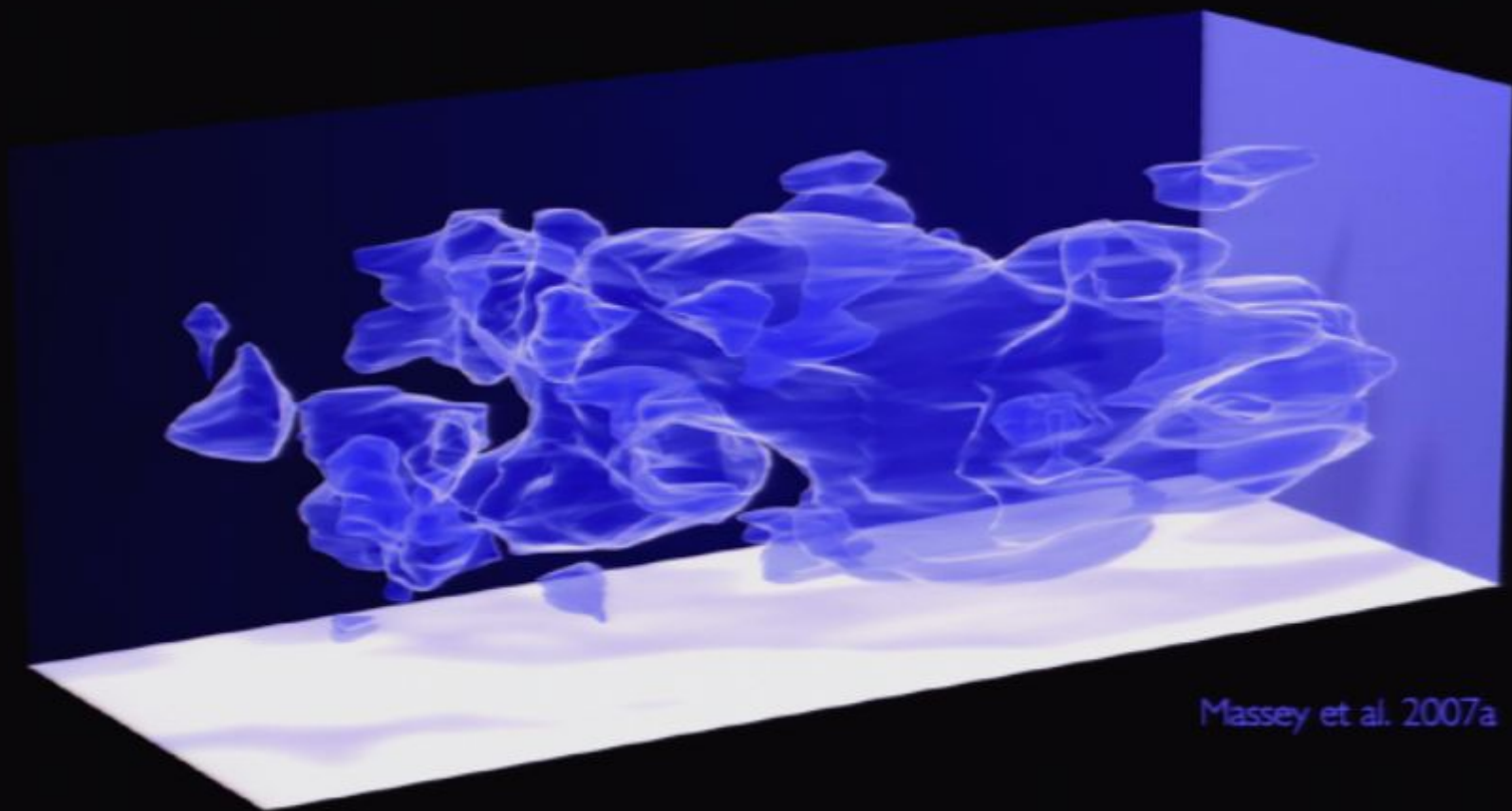
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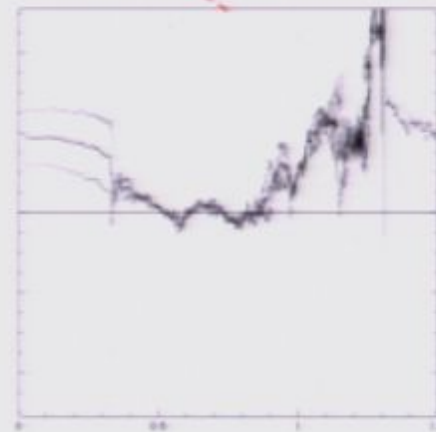
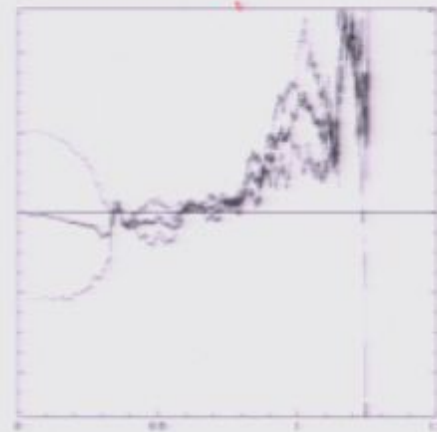
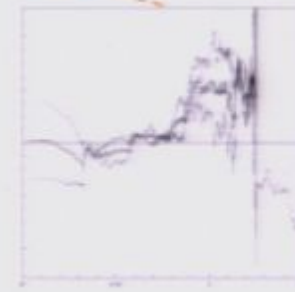
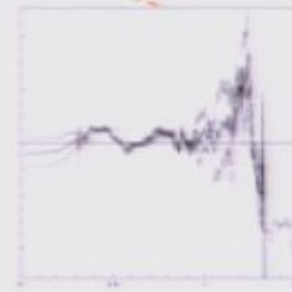
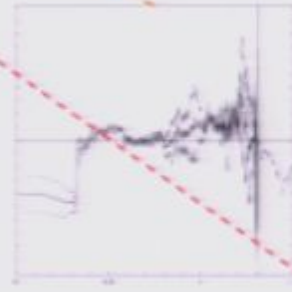
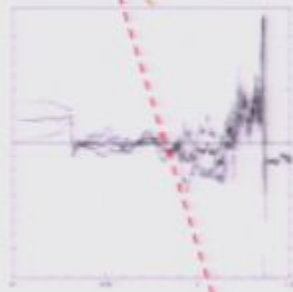
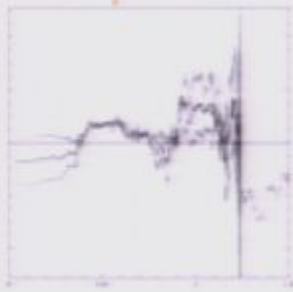
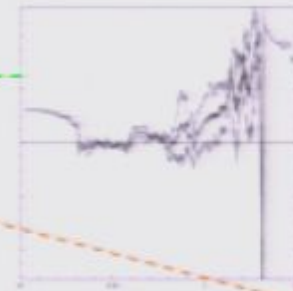
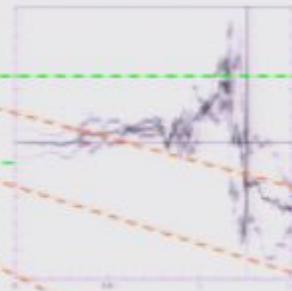
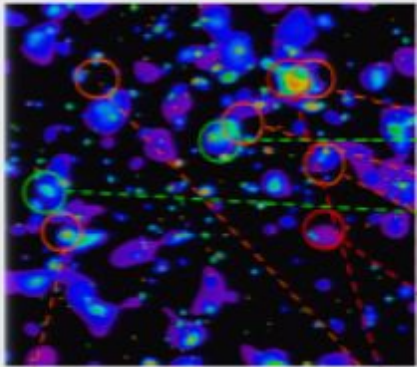
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Massey et al. 2007a

# Shear vs. photo-z around peaks, along promising lines of sight



## Summary

Can get constraints on the amplitude of fluctuations:

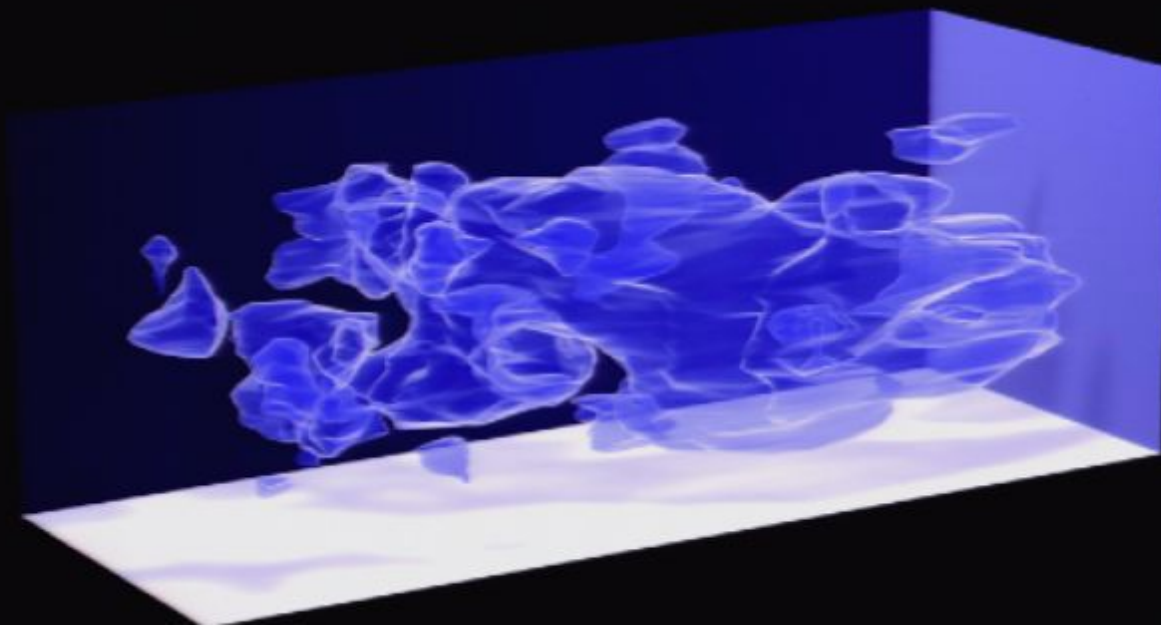
$$2D \quad \sigma_8 \left( \frac{\Omega_m}{0.3} \right)^{0.48} = 0.81 \pm 0.17$$

$$3D \quad \sigma_8 \left( \frac{\Omega_m}{0.3} \right)^{0.44} = 0.866^{+0.085}_{-0.068}$$

Need a larger volume to get constraints on the equation of state  
(current estimate:  $w = -1 \pm 1$  with this volume)

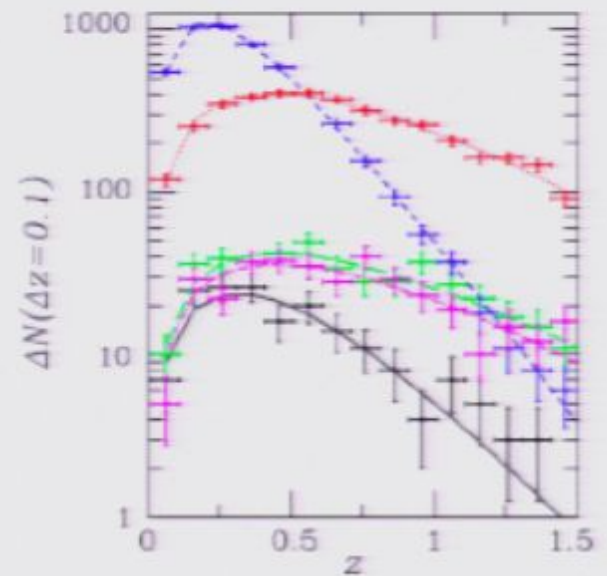
Working from space does not automatically make systematics vanish.

Physically, projection effects and photo-z errors a major issue.



## Second Method: Cluster Number Counts

- Count the number of high peaks in the universe
- In principle a very sensitive measure
- In practice, mass calibrations a problem
- Also constrained to do this at fairly low redshift (esp. for x-ray) and fairly large mass



Weller & Battye 2003

N.B. Press-Schechter: 
$$\frac{dn}{dM} dM = \frac{\rho_0}{M} \exp\left[-\frac{\nu^2}{2}\right] \frac{d\nu}{dM} dM$$

where

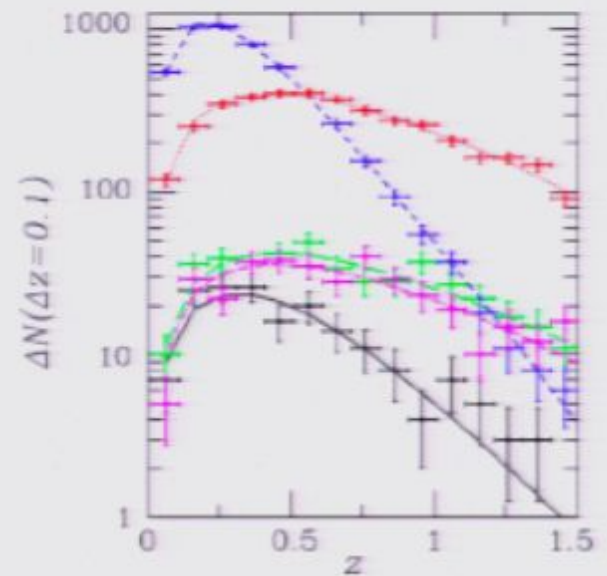
$$\nu = \frac{\delta_c}{D\sigma}$$

How to break sigma-growth factor degeneracy?

## Third Method: Differential Growth

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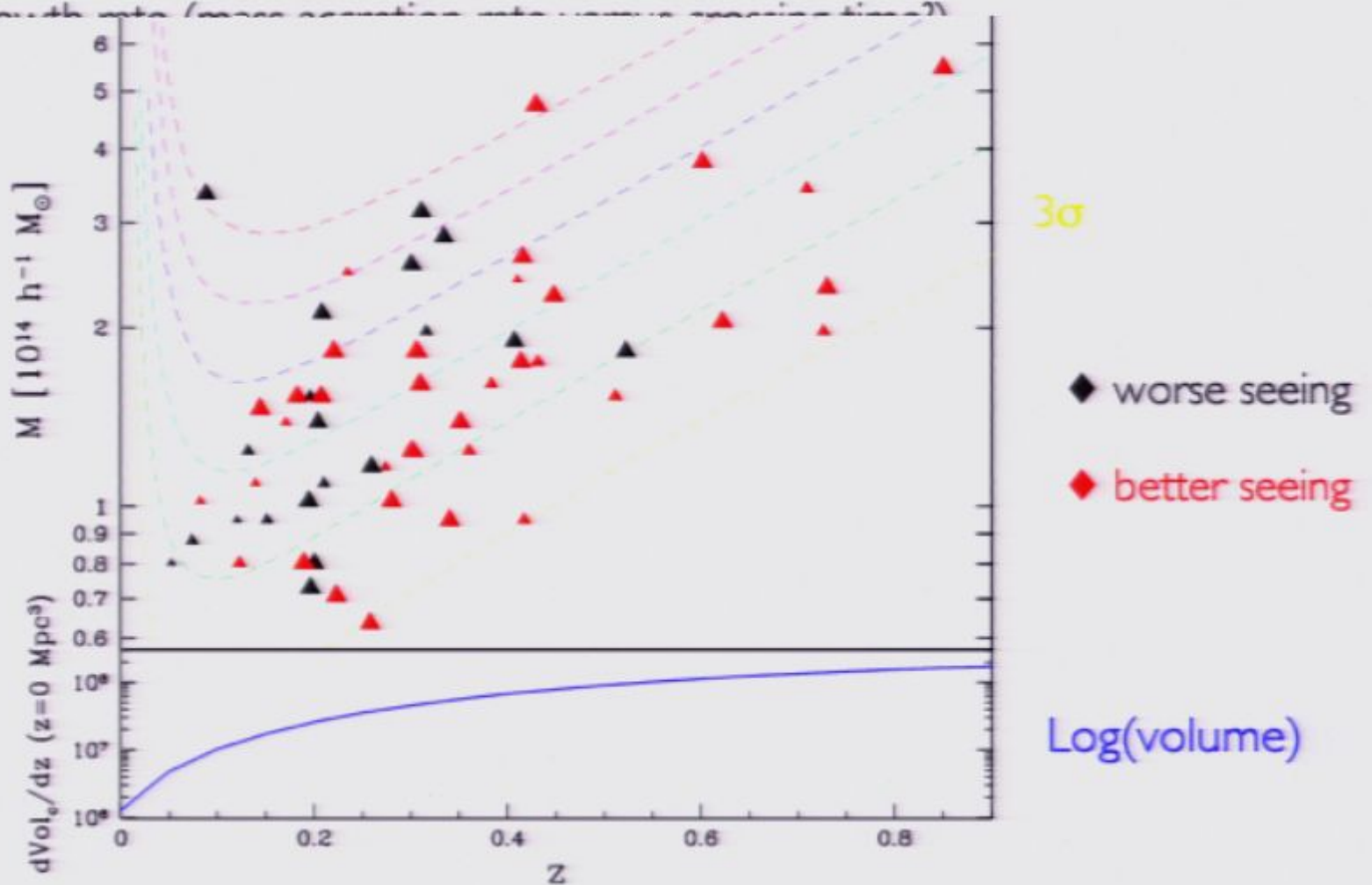
- Dark matter halos achieve a different degree of relaxation depending on their growth rate (mass accretion rate versus crossing time?)
- The change between the mixed and unmixed regions corresponds to a change in the slope of the density profile:  $r^{-1}$  to  $r^{-3}$

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- The change between the mixed and unmixed regions corresponds to a change in the slope of the density profile:  $r^{-1}$  to  $r^{-3}$
- This is measured as halo concentration:  $c \equiv r_s / r_{\text{vir}}$ ;  $r_s$  where slope is  $r^{-2}$

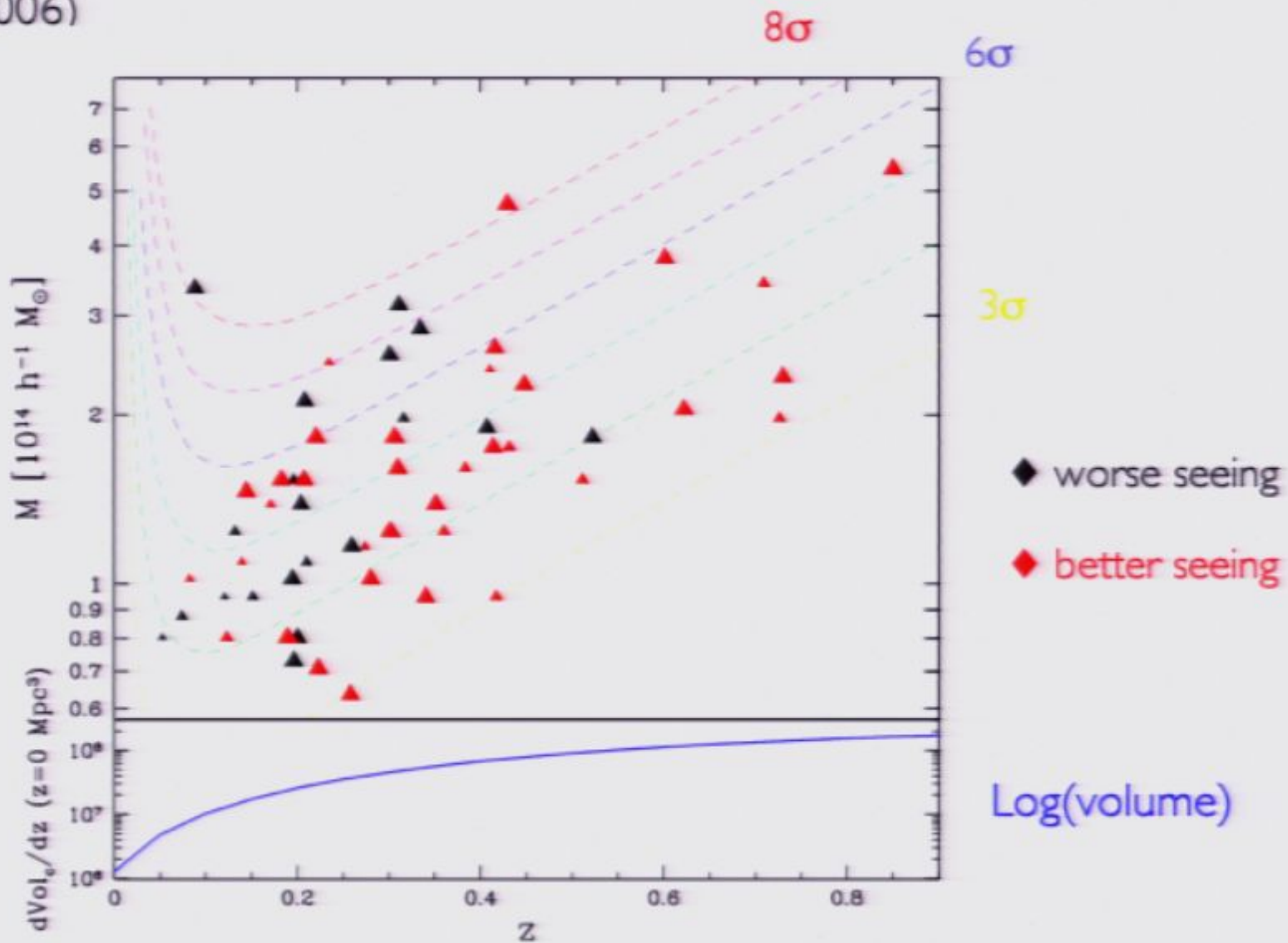
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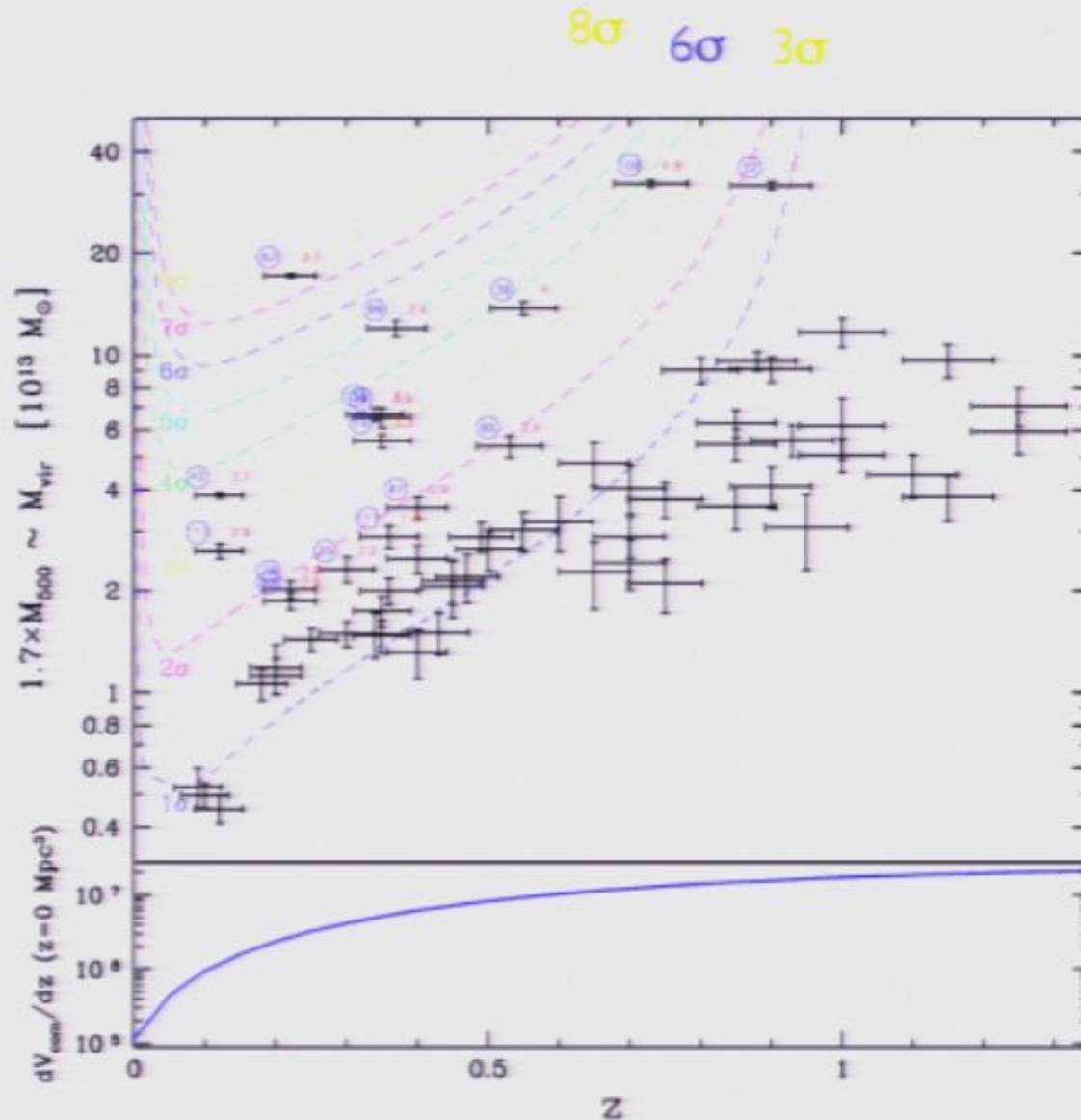
# Recent Suprime22 Results

(Green et al. 2006)



# COSMOS X-ray-lensing comparison

(Finoguenov et al. 2006, Taylor et al in prep.)



## Summary

Gravitational lensing in general and weak gravitational lensing in particular provide the most direct route to measuring the mass distribution in the universe.

Most measurements boil down to the expansion history and the linear growth factor, although lensing also tests the physical properties of dark matter (thermal velocity, scattering cross-section).

These measurements are and will remain complementary too CMB measurements and results from other methods.

Field is currently limited by available facilities => hope of future space missions?

- [-] Cosmological Constraints from Gravitational Lensing
  - James Taylor
  - Department of Physics and Astronomy
  - University of Waterloo
- [-] Outline
  - 1) Introduction: what is observational cosmology and what does it tell us?
  - 2) Lensing Basics, and forms of lensing
  - 3) Weak Lensing results in cosmology
  - 4) Results from a few recent surveys
- [-] Overview of Cosmic History
- [-] Overview of Cosmic History Era
 

Era	Redshift	time T (K)	E
Inflation, Topology	$10^{-35}$	s	$10^{27}$
	$10^{14}$	GeV	
The High-Energy Universe	$10^{-6}$	s	$10^{13}$
	1	GeV	
Last Scattering	1100	$3 \times 10^5$	yr 3300

**Cosmological Constraints from Gravitational Lensing**  
 James Taylor  
 Department of Physics and Astronomy  
 University of Waterloo

Click to add notes

Modified	S
ay, 12:10 AM	
y, 12:17 AM	
y, 12:40 AM	6
y, 12:43 AM	10
y, 12:44 AM	3
y, 12:46 AM	11
y, 12:48 AM	6
y, 12:50 AM	13
y, 12:53 AM	3
y, 12:55 AM	3
y, 1:05 PM	3
y, 2:03 AM	10
erday, 6:32 PM	21
erday, 6:36 PM	12
erday, 11:27 PM	
erday, 6:30 PM	13
y, 1:35 PM	16.
erday, 6:30 PM	16
erday, 6:34 PM	1.



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