

Title: Experimentalist View of the LHC

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Abstract:

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<http://cmsdoc.cern.ch/~smaria/perimeter07/>

Maria Spiropulu

CERN-PH

Perimeter Institute, Summer School on Particle Physics, Cosmology and
Strings.



August 9, 2007

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2. We are lucky that is LHC is happening now (fine fun)

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1. The LHC will determine the future of particle physics -

2. We are lucky that is LHC is happening now (fine tuning?)

3. It will not be a walk in the park

Outline

24 years and counting

Mission Statement

The "Machine"

The Large Hadron Collider

Status of the LHC Experiments

The LHC Physics Besides the Standard Model

SUSY@LHC program of work

ATLAS+CMS SUSY search Highlights

Engineering the Discovery Plan

Wisdom from the Past



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lhc_atlas.swf

realPlayer

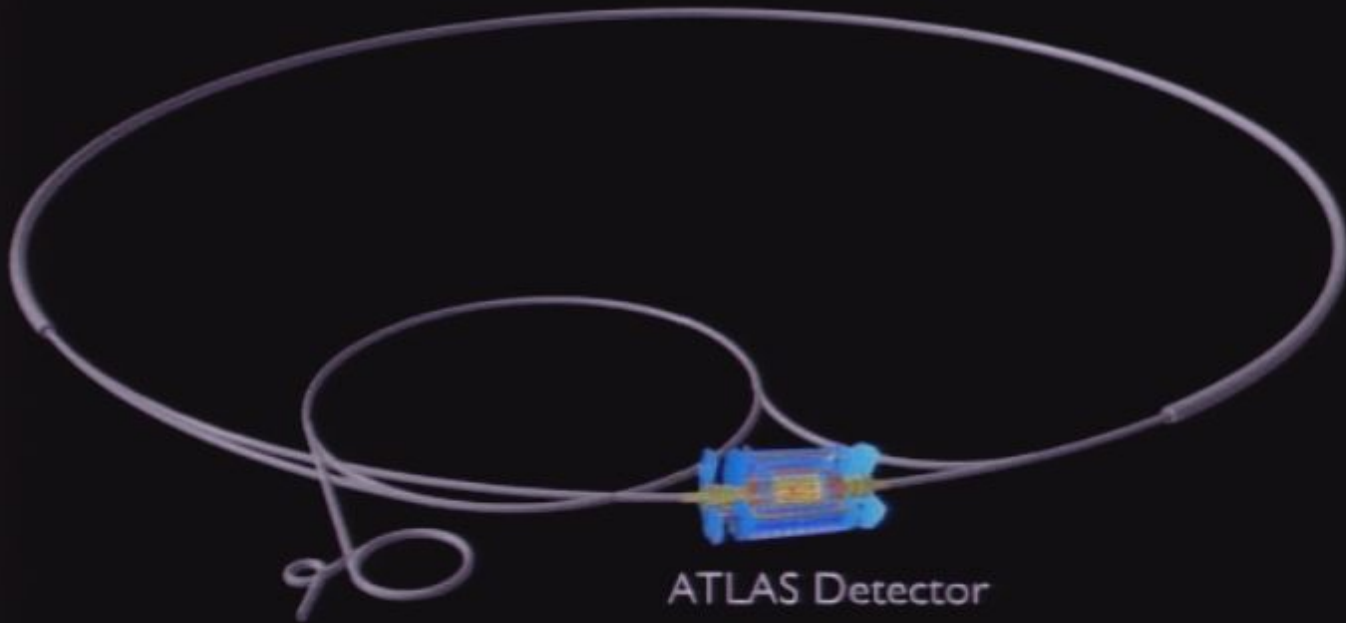
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PLAY ▶

Large Hadron Collider

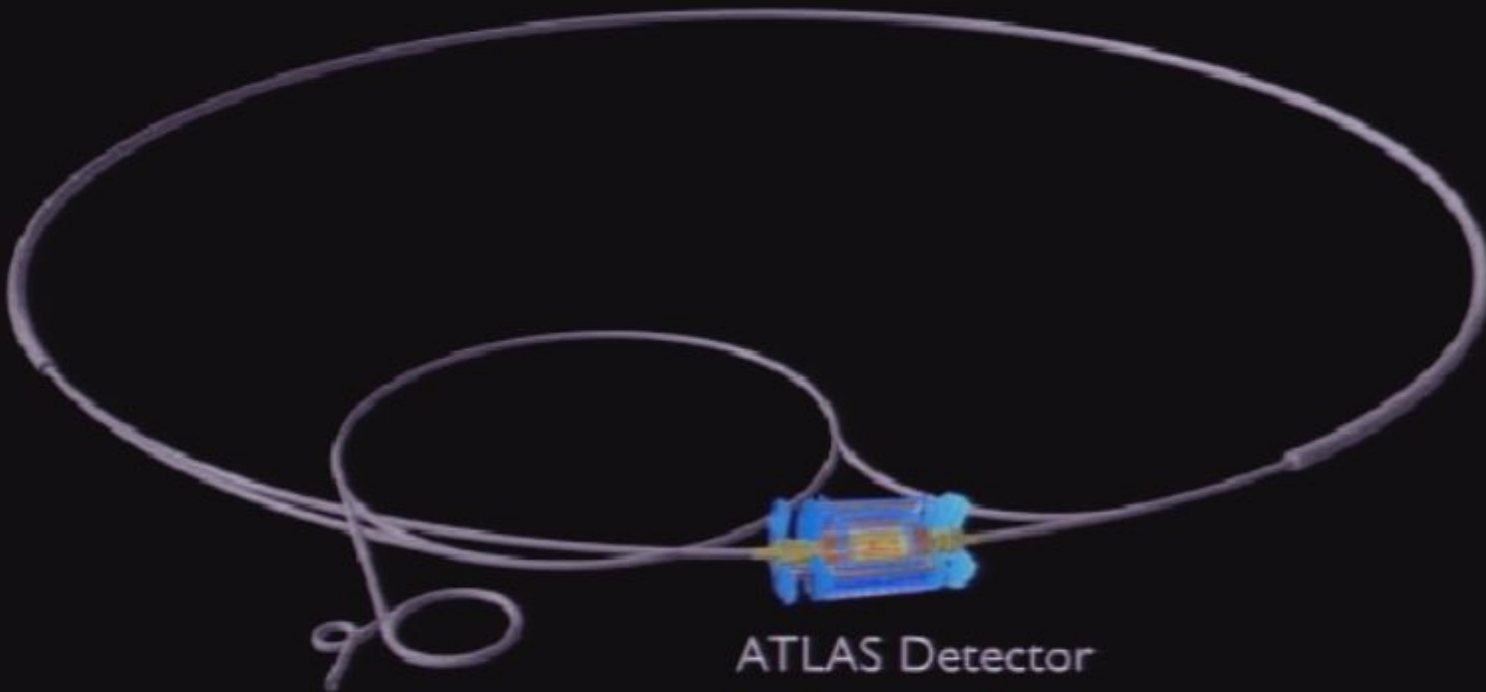


ATLAS Detector



PLAY ▶

Large Hadron Collider



ATLAS Detector

lhc_atlas.swf

Playing

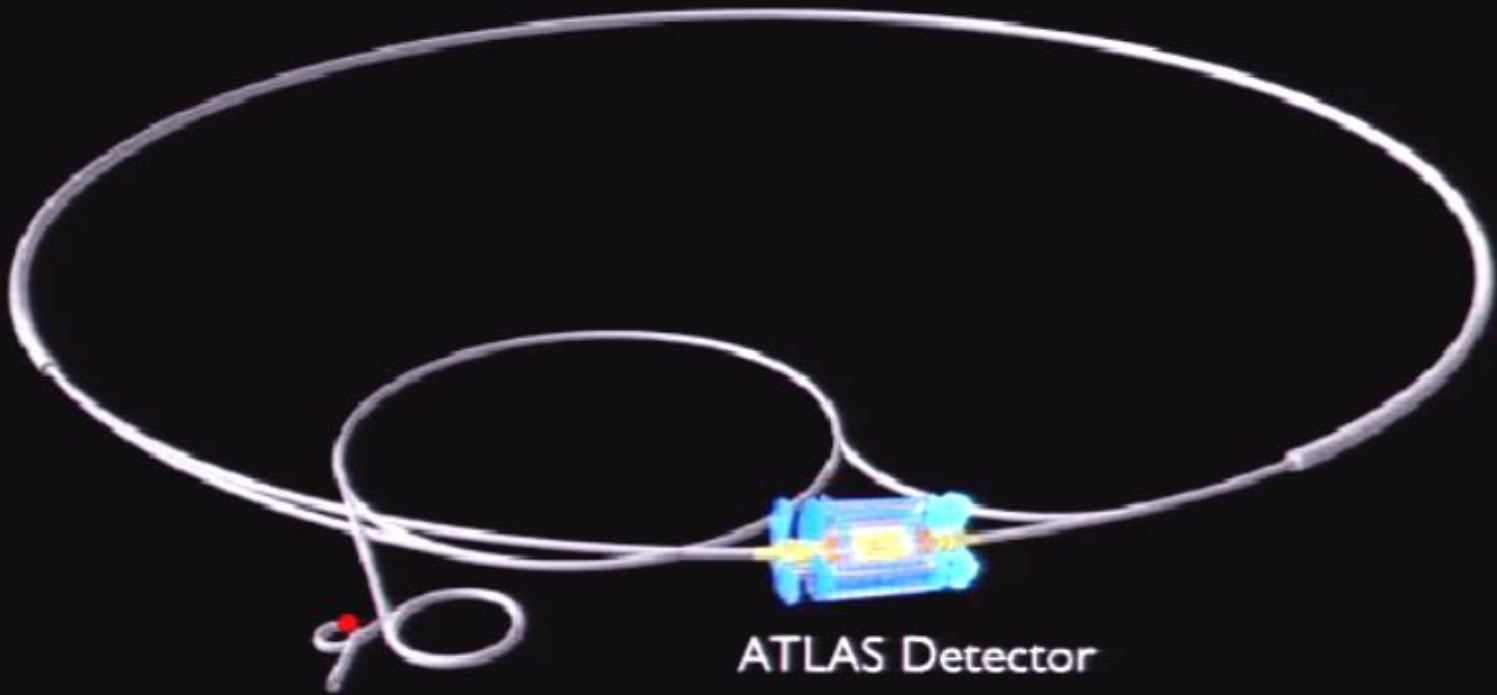
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Large Hadron Collider



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lhc_atlas.swf

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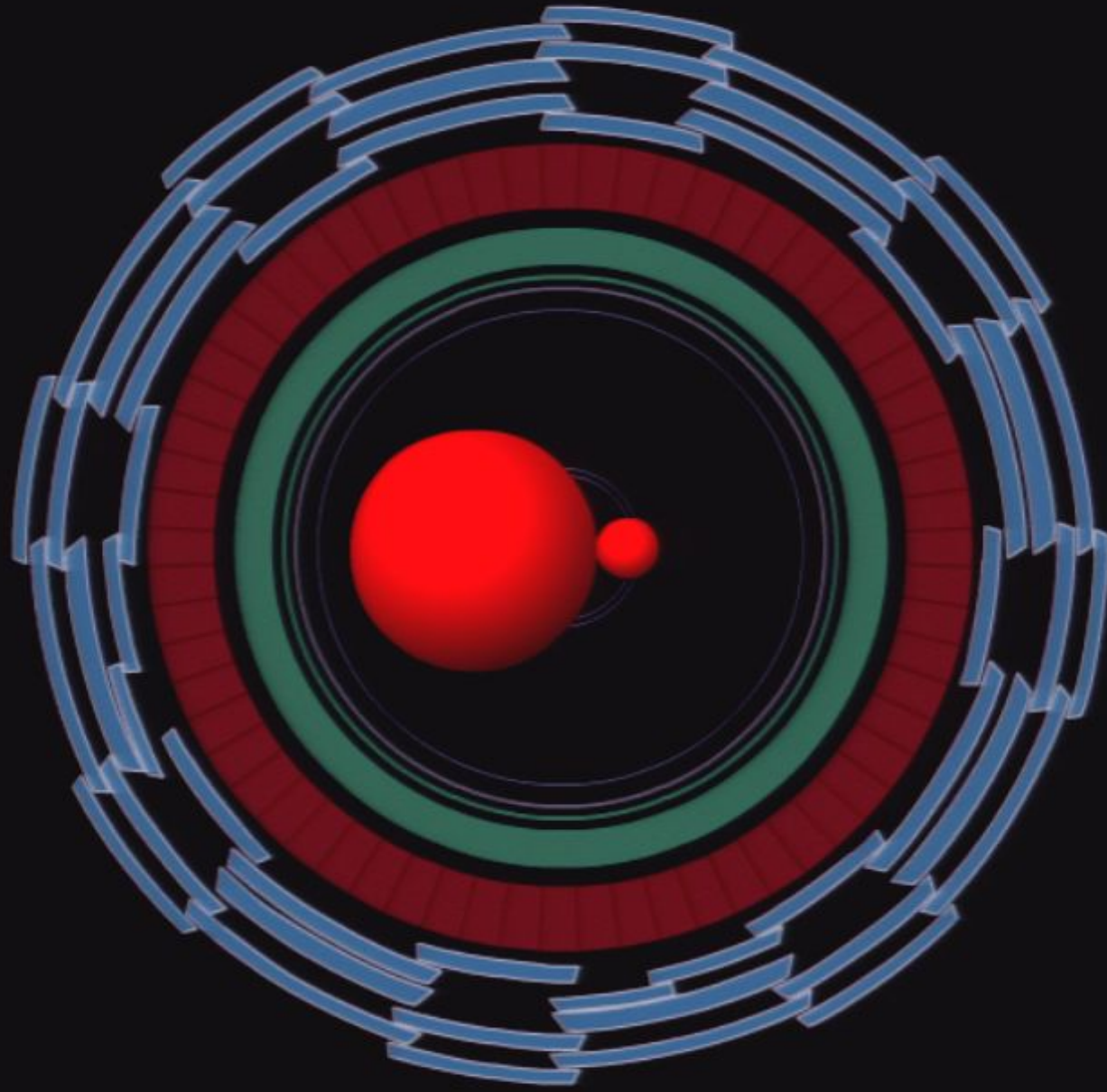
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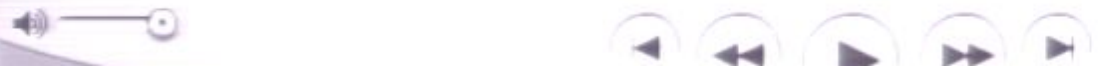




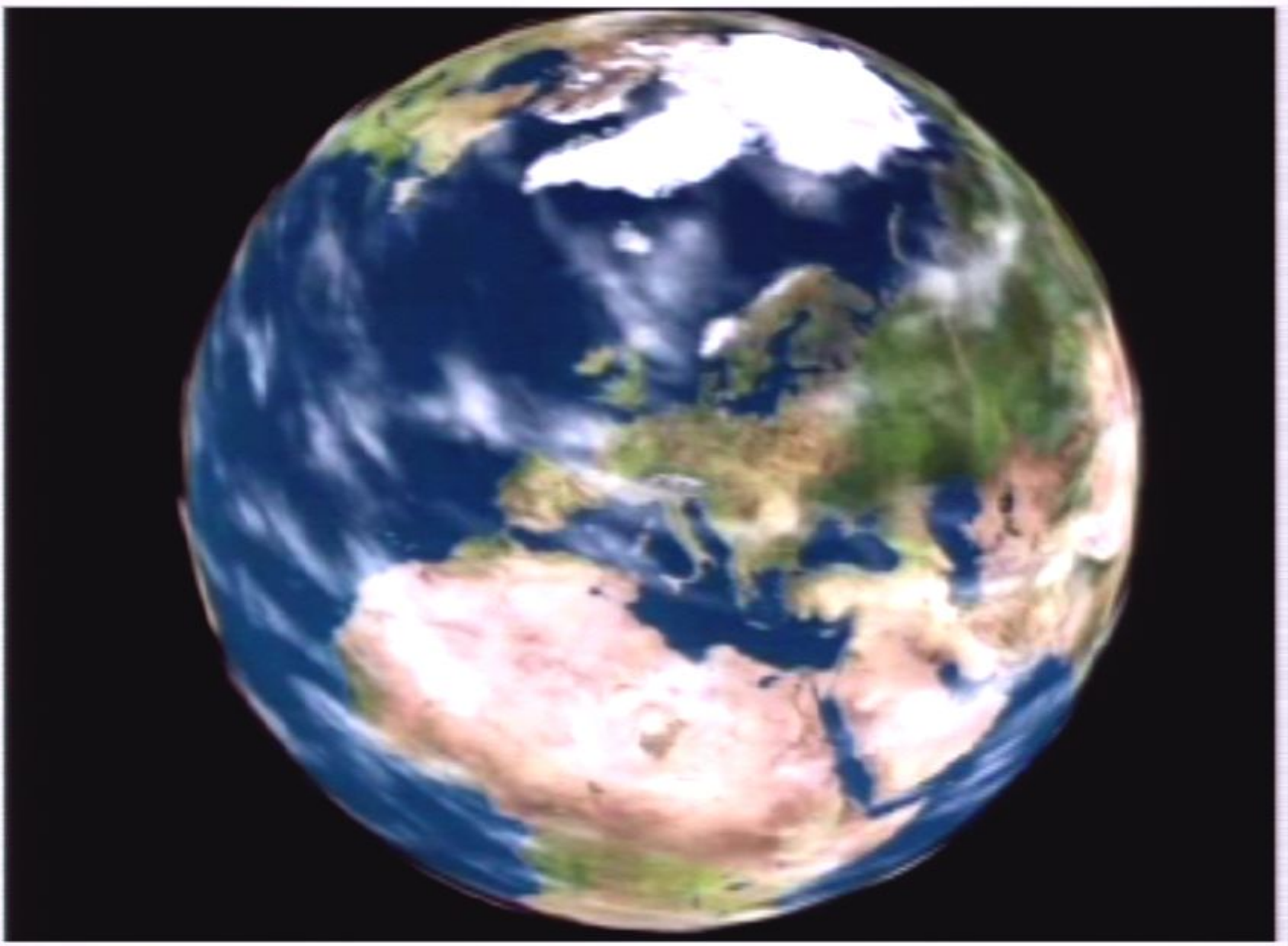
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Previous



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previous



00:00:00 [Progress Bar]

[Volume Control] [Play/Pause] [Previous] [Next] [Full Screen]

Navigation controls: Previous, Next, Page (3), Back/Forward, Zoom In, Zoom Out, Tool Mode.

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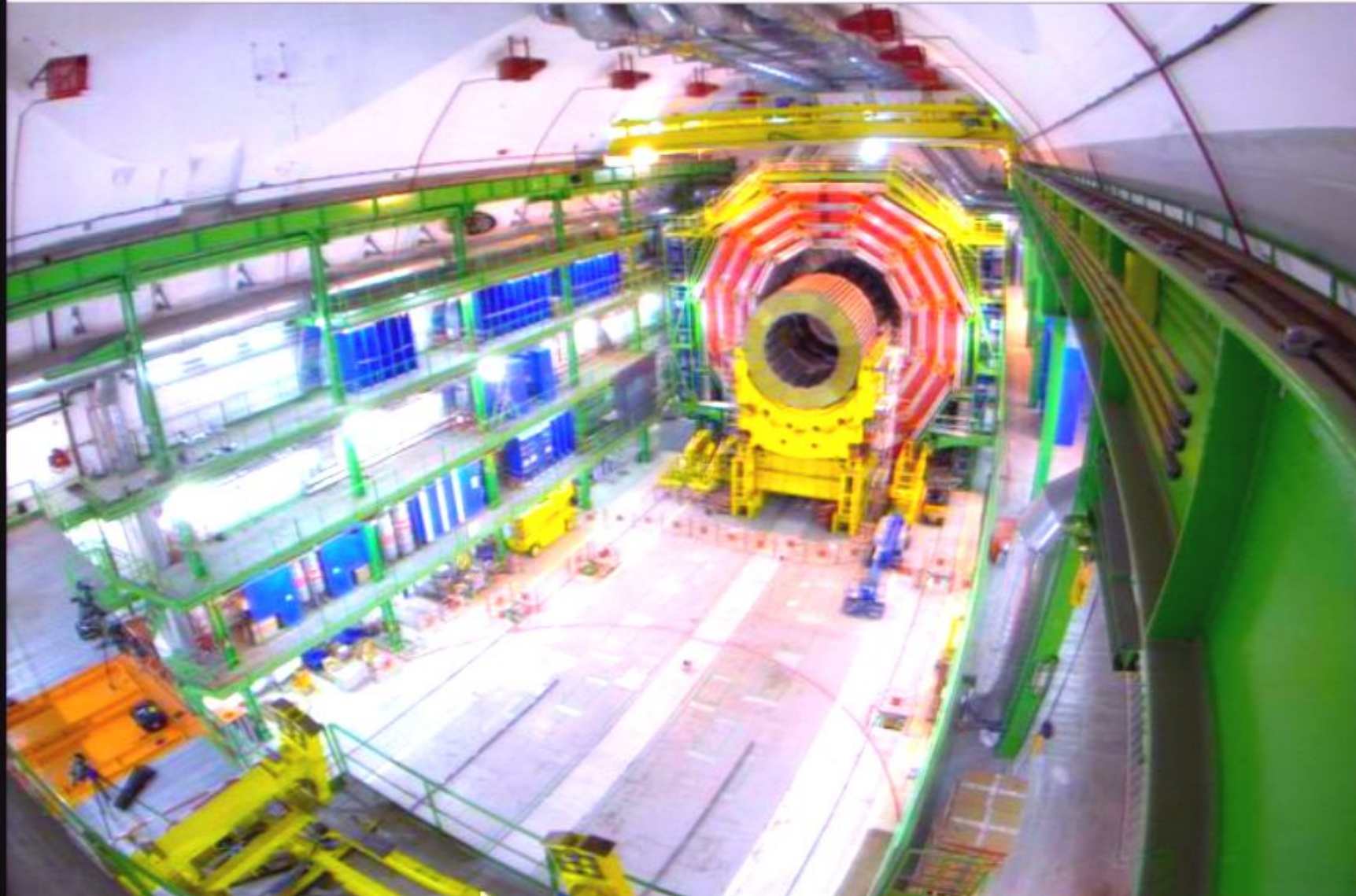
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CMSeye::YB0 Lowering on 28 February 2007

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LHC machine ...oto gallery Apple (146) Yahoo! News (1270) Indico [CMS meetings] CMSSW LHC Magnets



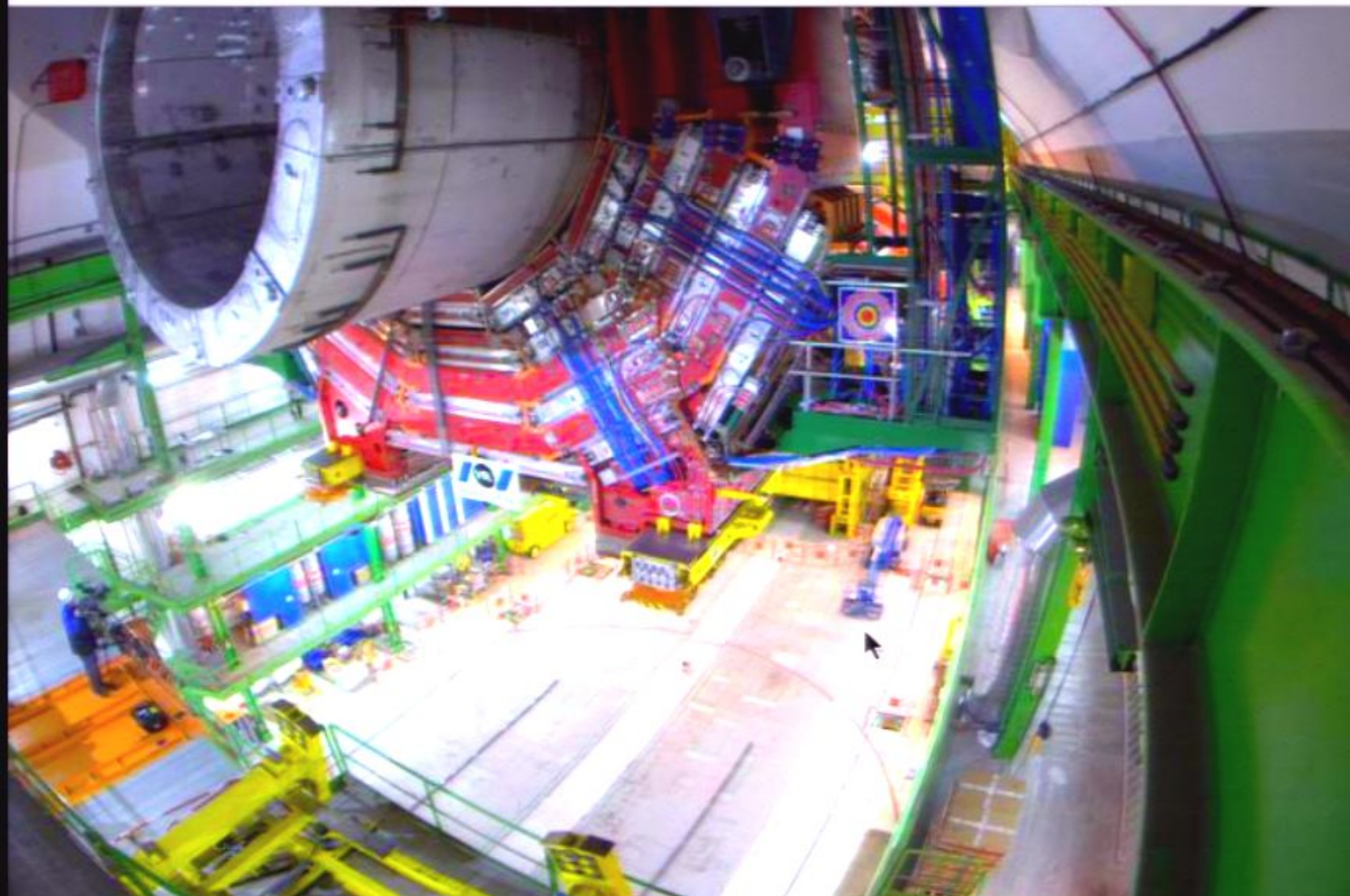
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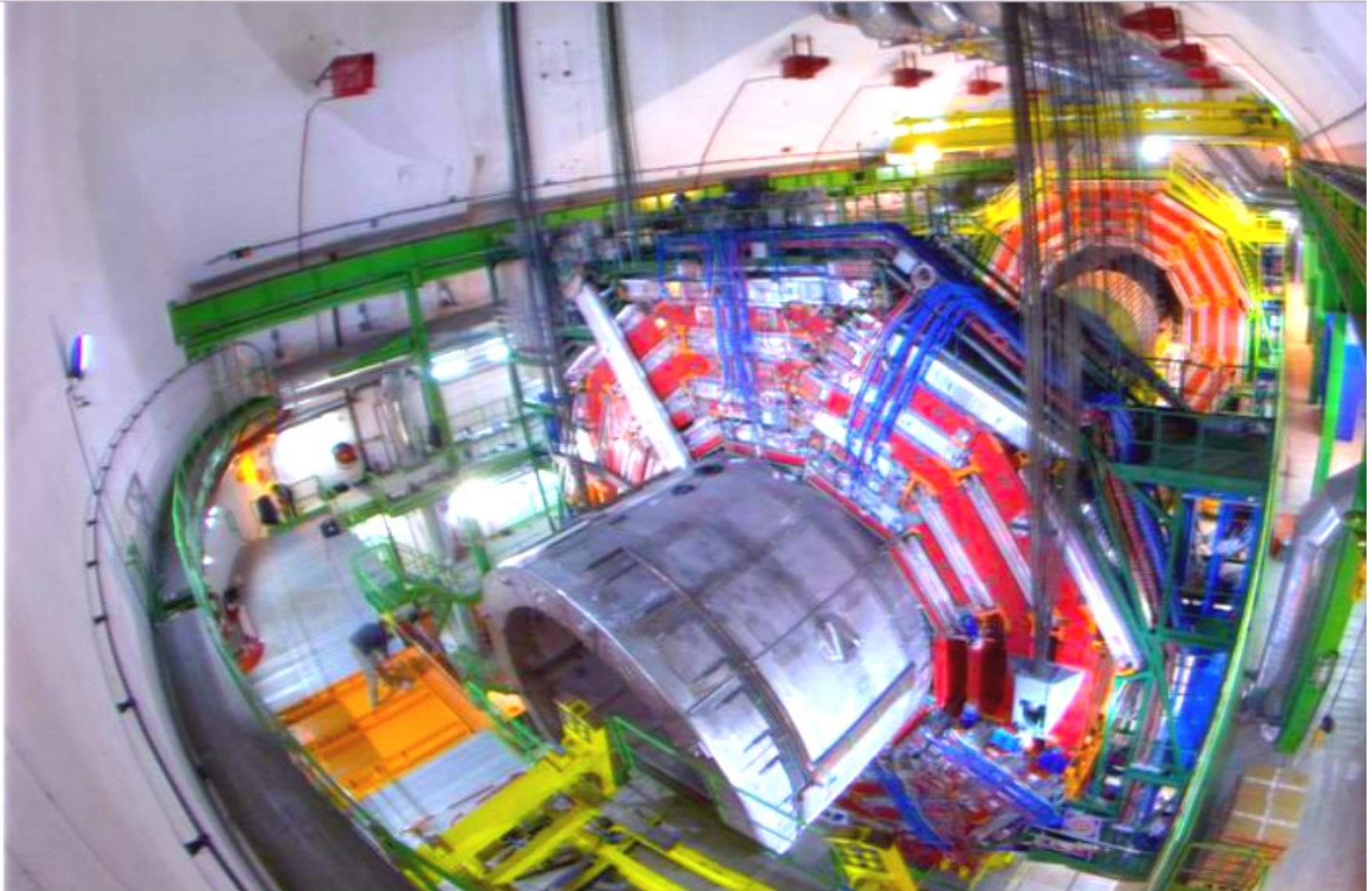
Start Stop Reset Fwd 1 Frame Back 1 Frame

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Navigation toolbar with icons for Previous, Next, Page (4), Back/Forward, Zoom In, Zoom Out, and Tool Mode.

24++

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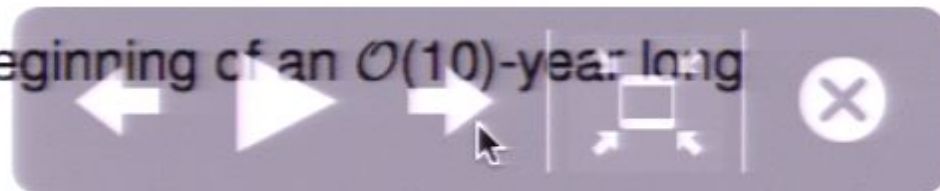
CMS cavern completed

2007-8

Experiments ready for beam

July 2008

14 TeV collision run - the beginning of an $\mathcal{O}(10)$ -year long experimental program.



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Mission Statement of the LHC and its experiments

From the Presentation to CERN Council in Dec 1993

[...] the Large Hadron Collider (LHC) [is] **the right machine for further significant advance in the field of high energy physics research** and for the future of CERN [...] The accelerator will produce not only higher energy but also a higher luminosity [...] than has been achieved before and it will reveal the behavior of fundamental particles of matter which has never been studied. [from Council report]



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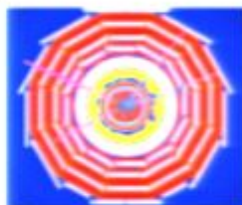


The major physics scope of the LHC

- Find and characterize the new particles that compose the dark matter of the universe



- Find the Higgs particle



- Find new particles, forces, extra dimensions of space



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The Large Hadron Collider



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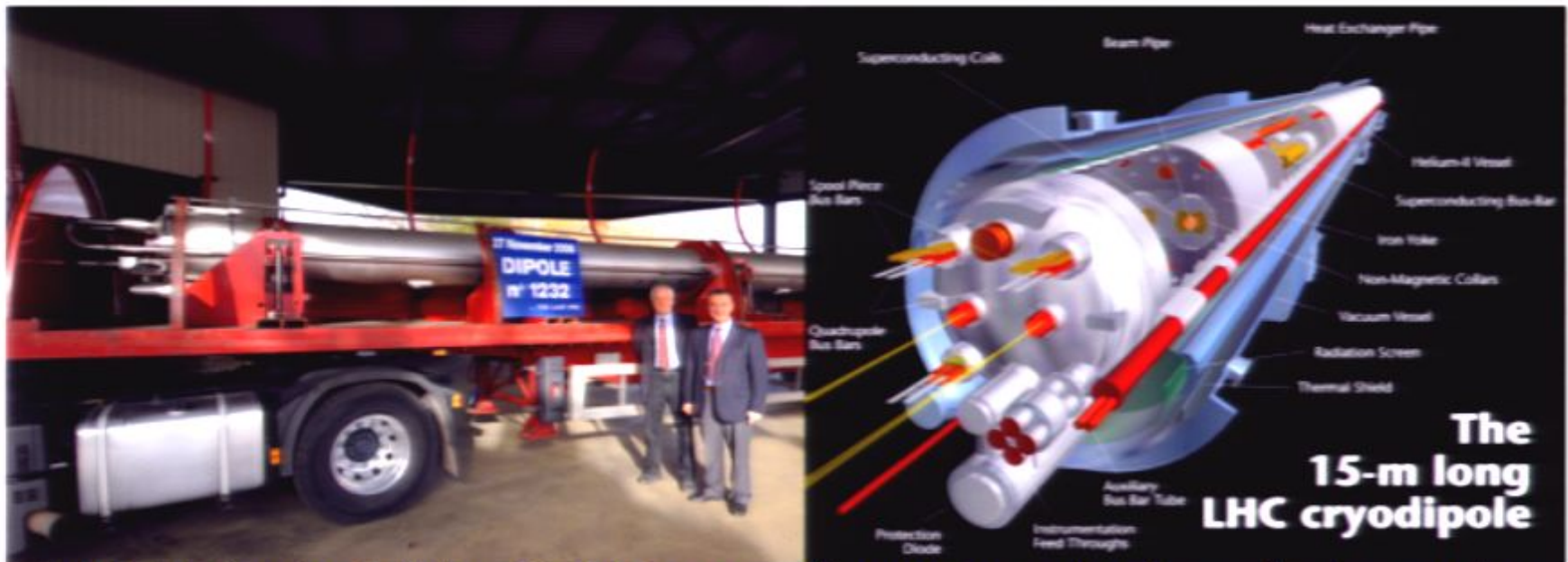
The machine



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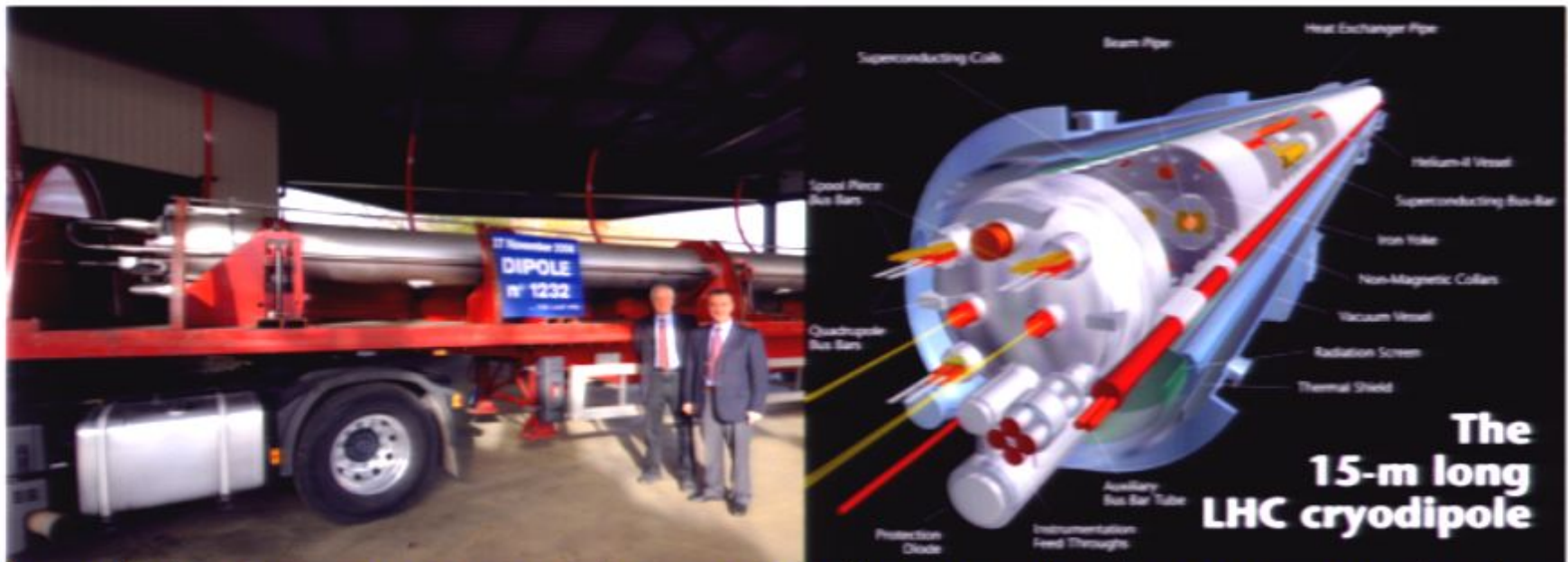
1232 Cryodipoles: The bending force



- 8.33 T dipole fields; 31,000 tons of superconducting dipoles to be operated at 1.9°K [currents of >15,000 amps] → 90 tons of liquid helium and > 1,000 tons of liquid nitrogen for cryo purposes.



1232 Cryodipoles: The bending force



To collide two counter-rotating proton beams, the beams must be in separate vacuum chambers (in the bending sections) with opposite B field direction: There are TWO LHCs: The twin-bore design was proposed by John Blewett, BNL in 1971.



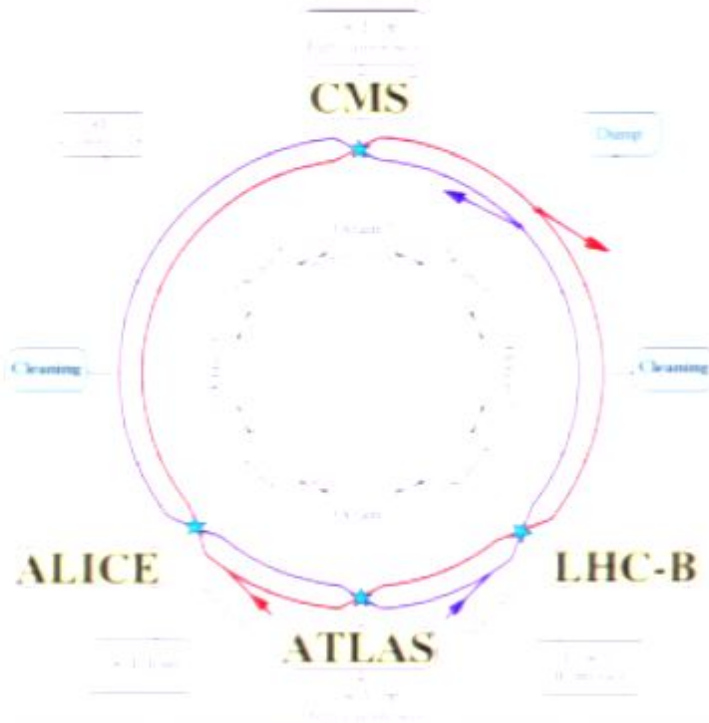
The machine



Navigation controls: back, play, forward, zoom in, zoom out, close.



Schematic View of the LHC



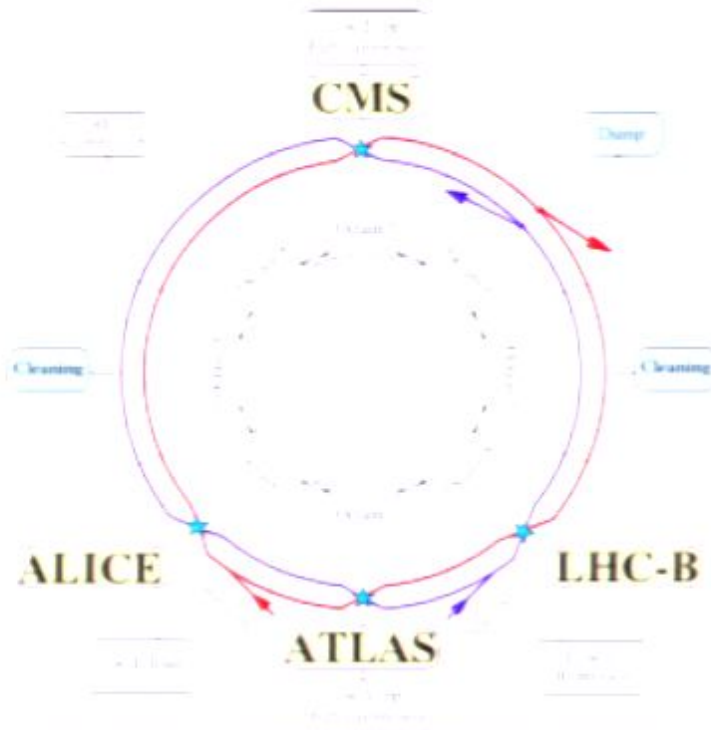
- IP1=ATLAS (general purpose)
- IP5=CMS (general purpose)
- IP2=ALICE (ion-ion, p-ion)
- IP8=LHCb (pp, B-physics, CP-violation)

• Design: 2808 × 2 bunches of protons, 10^{11} protons/bunch accelerated to 7 TeV; The total energy of each beam: $2808 \times 1.15 \times 10^{11}$ protons per bunch \times 7 TeV = $0.185 M_{\text{Planck}} = 362$ MJoules \rightarrow

Absolutely NO accidents allowed : if the beam is steered wrongly a surface of 1/10 mm will be subjected to 4 Terawatts of proton beam power \sim detonation of 86 Kg of TNT!



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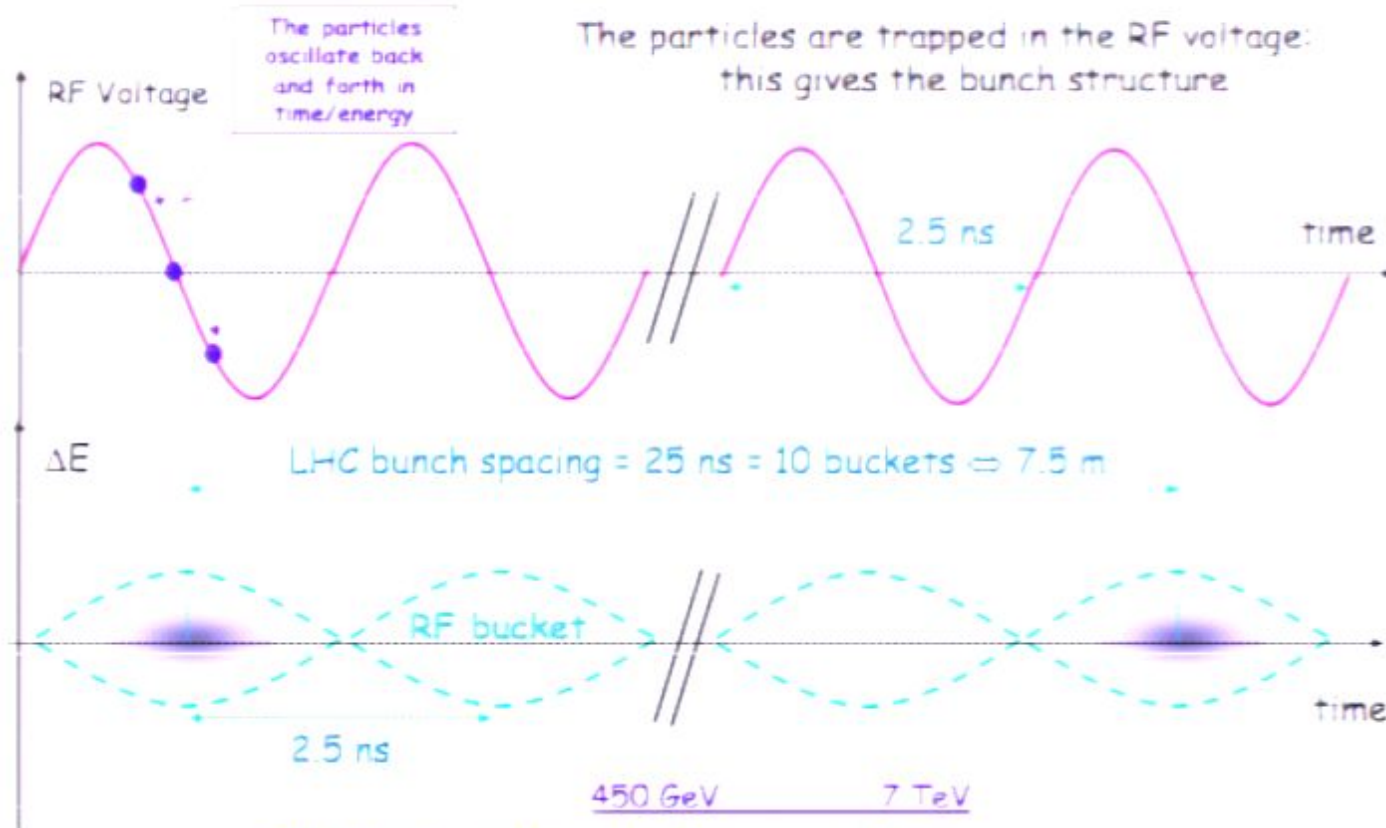
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The beam

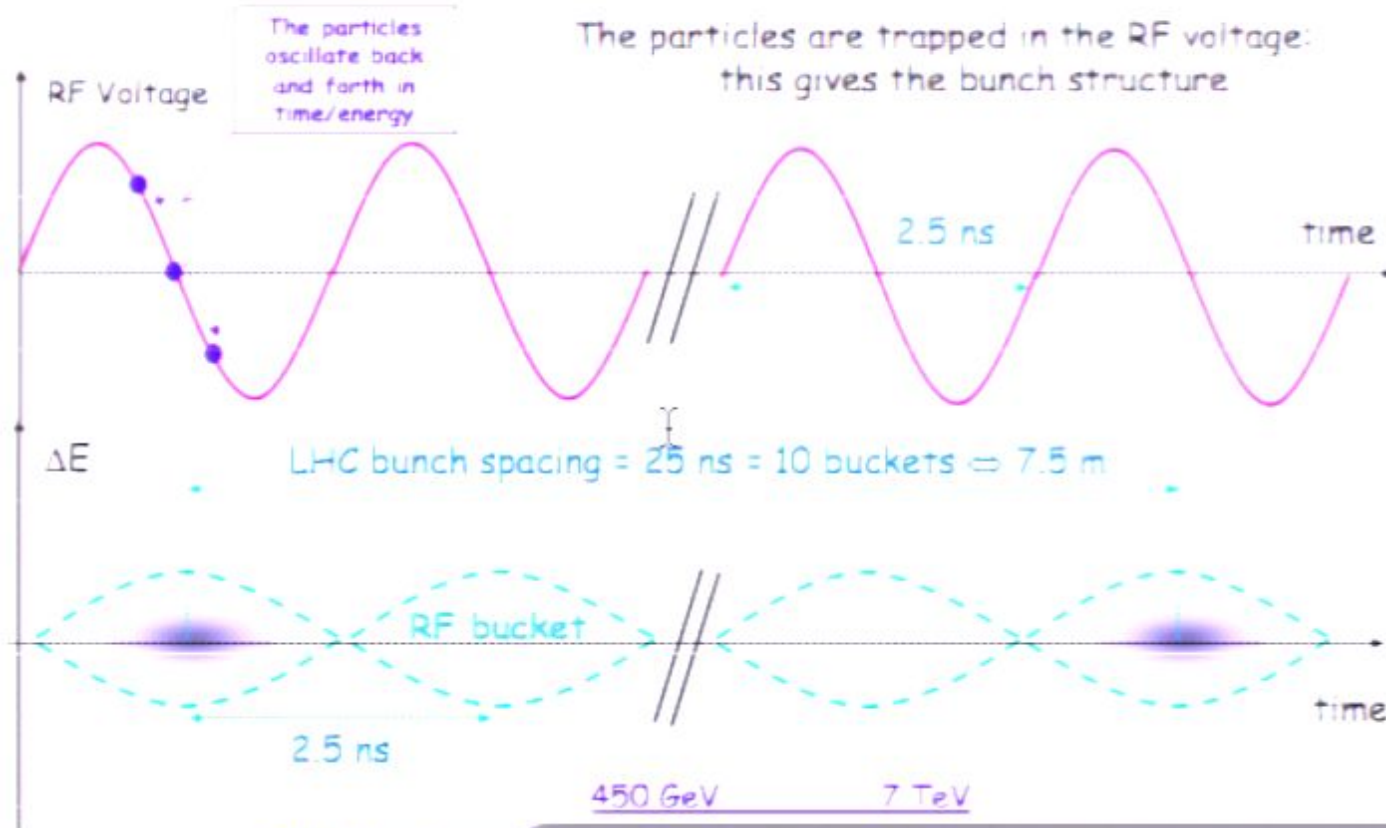
RF buckets and bunches



	450 GeV	7 TeV
RMS bunch length	11.2 cm	7.6 cm
RMS energy spread	0.031%	0.011%

The beam

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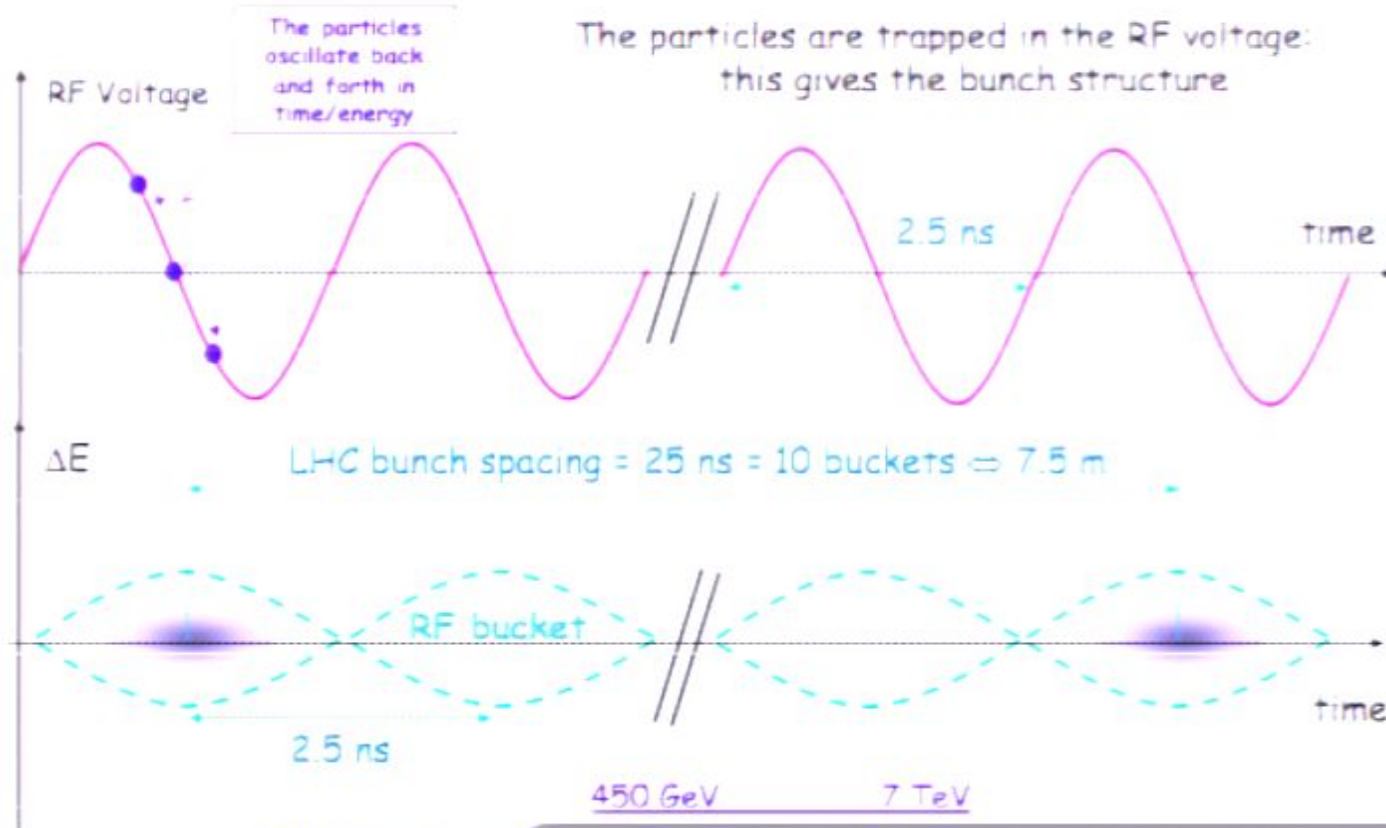
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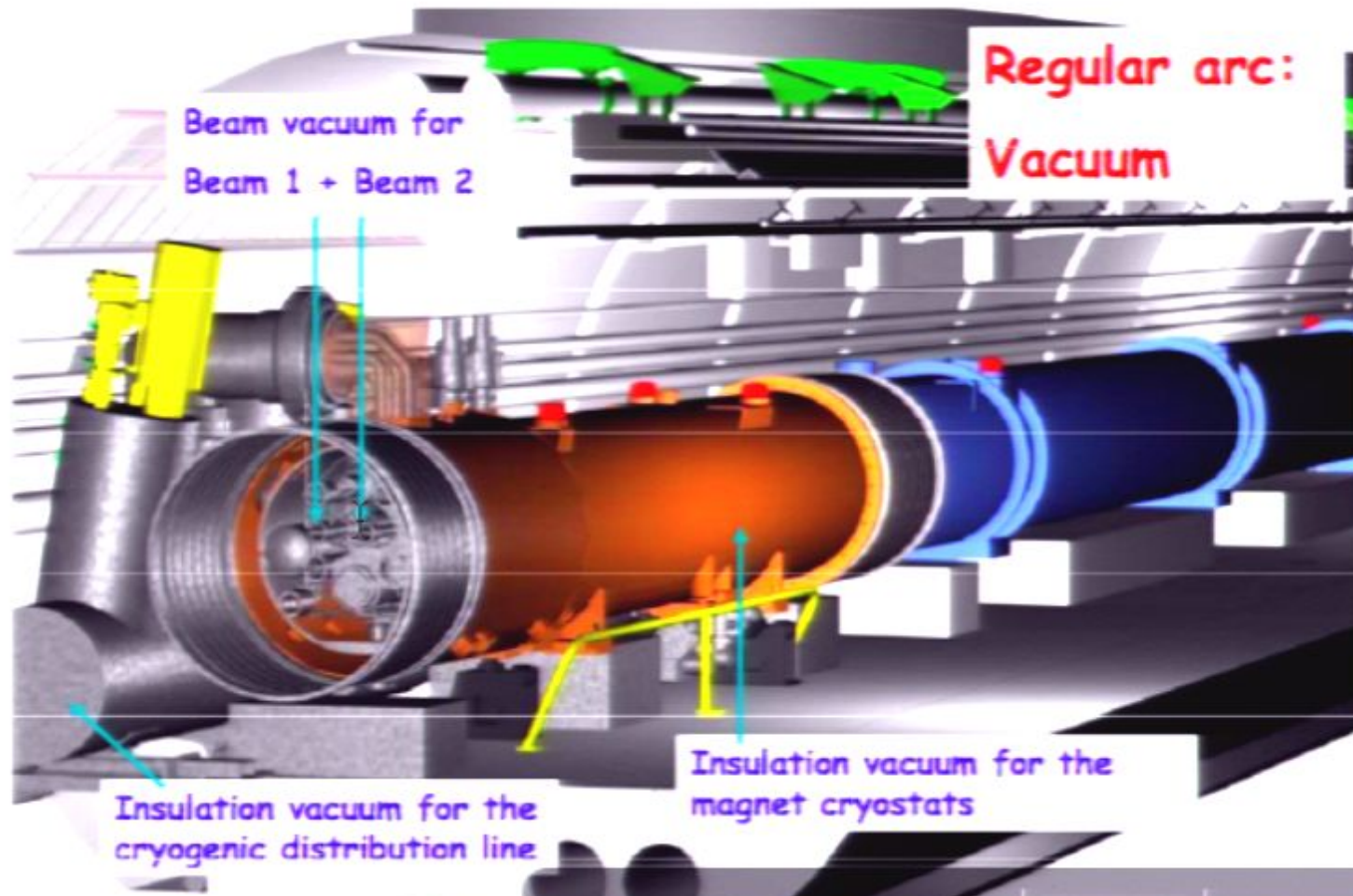


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Tunnel Instrumentation



Accelerator Complex



Schedule

Lyn Evans
SPC 18-June-2007

General LHC Schedule

- Engineering run originally foreseen at end 2007 now precluded by delays in installation and equipment commissioning.
- 450 GeV operation now part of normal setting up procedure for beam commissioning to high-energy
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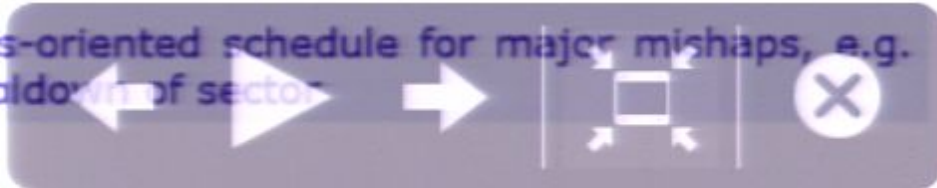
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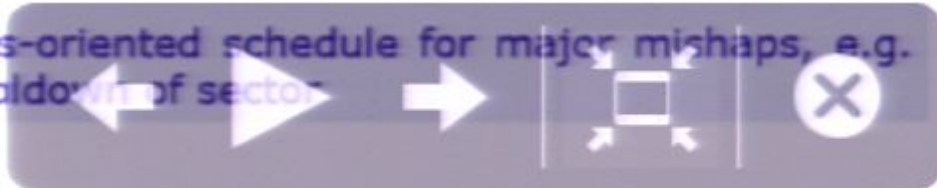


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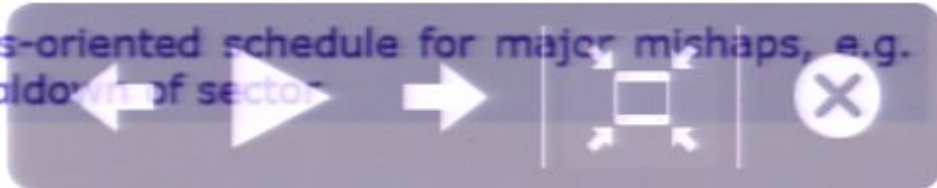


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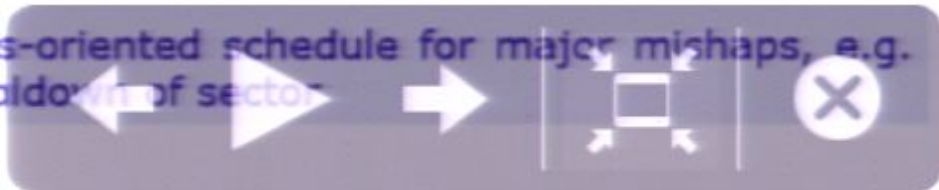


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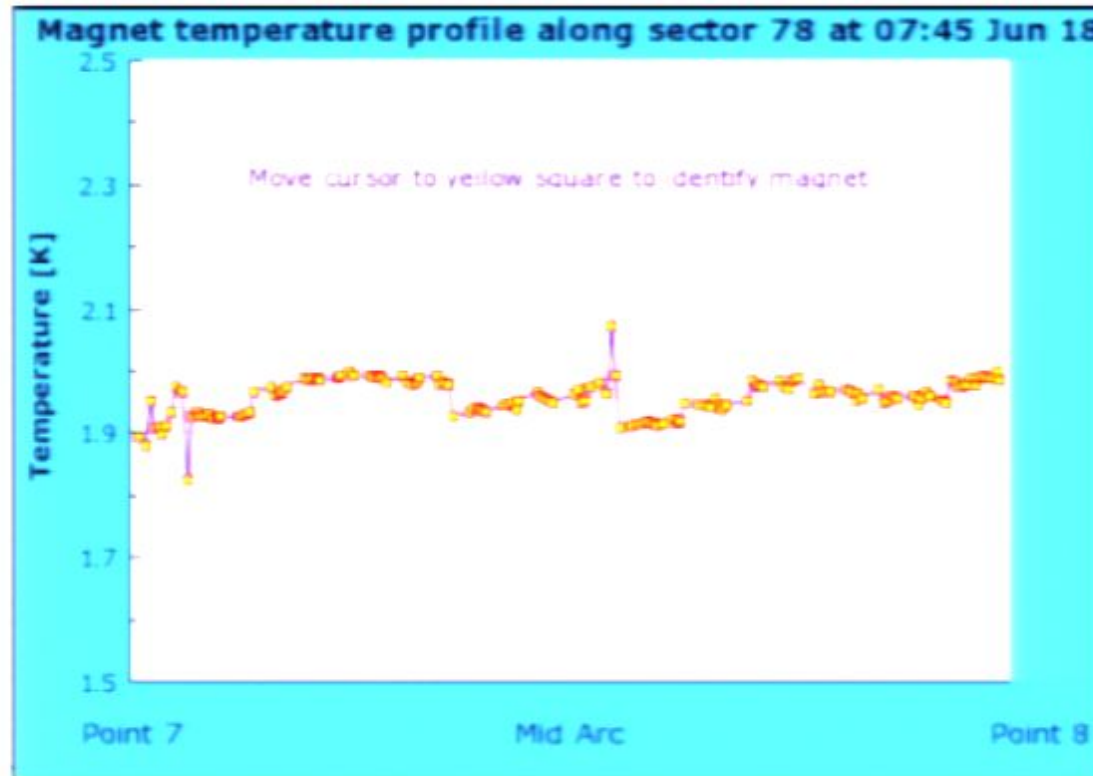
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Cool

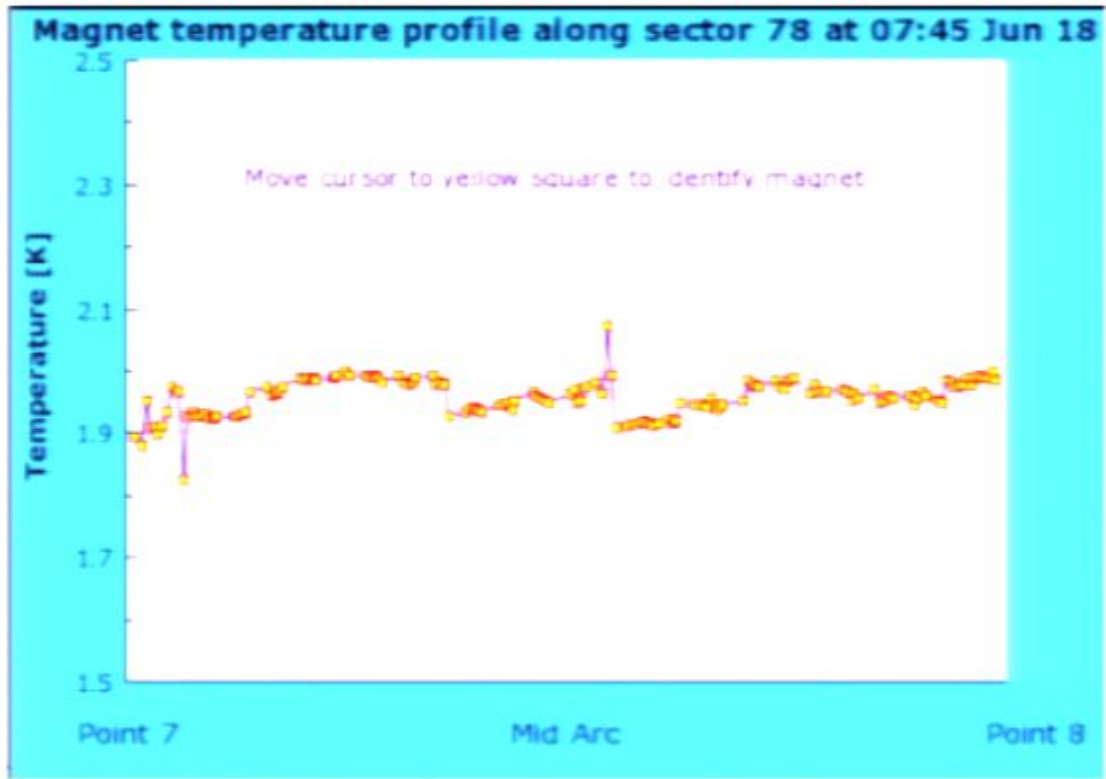


1 sector=3.3 Km, 154 dipoles. The first sector was cooled-down in Feb 07.

Timely cool-down of all sectors determines the schedule.



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Commissioning stages

The steps toward the design luminosity:

- number of bunches and bunch intensity
- crossing angle
- less focusing at the collision point (larger β^*)

Parameter	Phase A	Phase B	Phase C	Nominal
k / no. bunches	43-156	936	2808	2808
Bunch spacing (ns)	2021-566	75	25	25
N (10^{11} protons)	0.4-0.9	0.4-0.9	0.5	1.15
Crossing angle (μrad)	0	250	280	280
$\sqrt{\beta^*/\beta^*_{\text{nom}}}$	2	$\sqrt{2}$	1	1
σ^* (μm , IR1&5)	32	22	16	16
L ($\text{cm}^{-2}\text{s}^{-1}$)	$6 \times 10^{30} - 10^{32}$	$10^{32} - 10^{33}$	$(1-2) \times 10^{33}$	10^{34}



The Luminosity

Luminosity

$$\mathcal{L} = \frac{n_b N_L N_R f_{rev}}{A_T^{eff}} \sim 10^{34} \text{ cm}^{-2} \text{ s}^{-1}, \text{ where}$$

$$f_{rev} = c/27 \text{ km} \sim 10^4 \text{ Hz and}$$

$A_T^{eff} = 4\pi\sigma_b^2$, the effective transverse area of the proton beam
with $\sigma_b = 16$ microns

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cross section

total inelastic [non-diffractive] cross section 60 millibarns. [1
 barn= 10^{-24} cm^2]

- long range inelastic strong interactions modelled by meson exchange as in the original Yukawa theory [back of the envelope calculation, $\sigma = \pi \text{range}^2$] will give this answer as well!

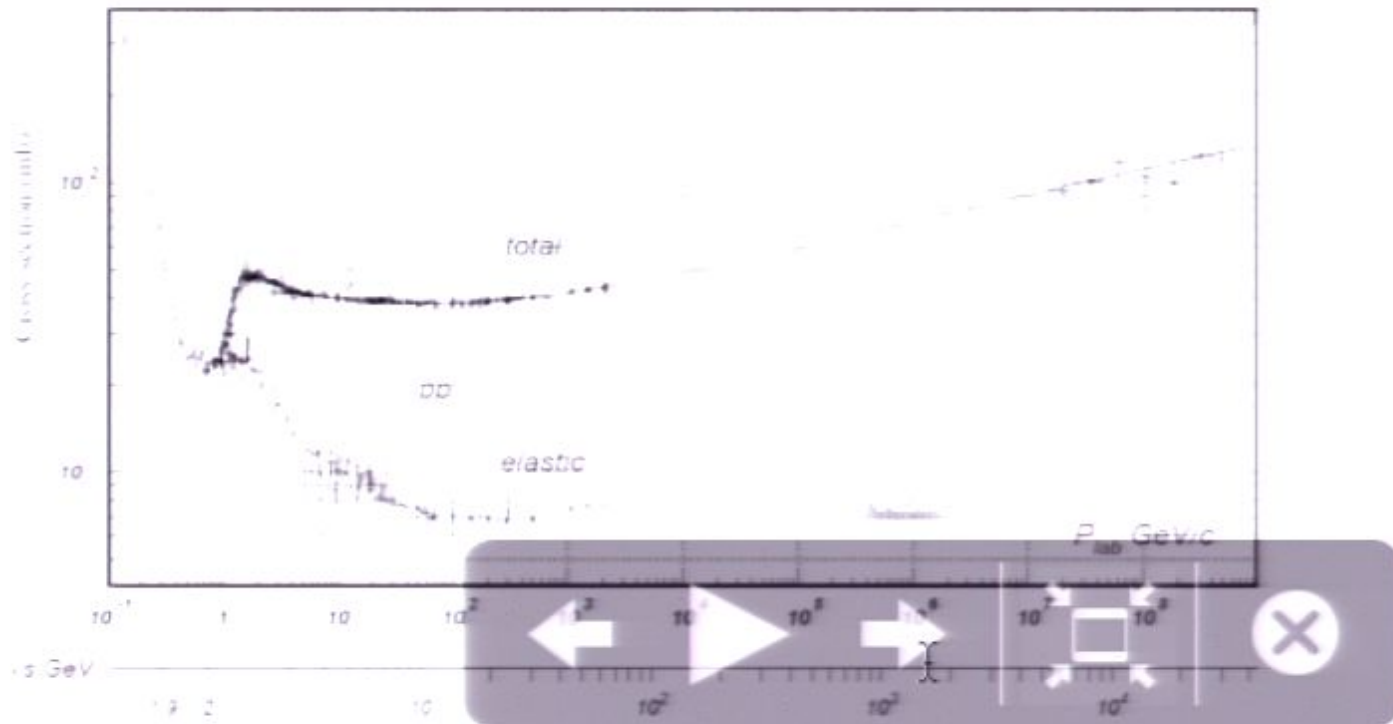
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$A_T^{eff} = 4\pi\sigma_b^2$, the effective transverse area of the proton beam
with $\sigma_b = 16$ microns

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of collisions per second = $\mathcal{L} \times \sigma \sim 10^9$ Hz

- detect, record and analyze this: a billion per second...



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The rate that the 2808 proton bunches "cross" each other at each interaction region = $2808 c/27000 \sim 3 \times 10^7$ Hz

- each bunch crossing produces >20 proton-proton scatterings!
 - bunch crossings every 25 ns; in 25 ns light travels 7.5m, much less than the size of the detectors! No electronic system can read the data from the entire detector and decide to do anything with it before the next event comes about!
- Readout is intergated over many bunch crossings**



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integrated luminosity

$\int \mathcal{L} dt = L$; units cm^{-2} or fb^{-1} . One year is about $\pi \times 10^7$ s. One
 "collider running year" [or Snowmass year] is $1/\pi$ this. So the
 integrated luminosity over a year at the LHC is $10^{41} \text{ cm}^{-2} \rightarrow$
 100 fb^{-1} . **Compare with the total integrated luminosity of the
 Tevatron after 10 years of running which is 2 fb^{-1} !**

- Caveat: Design luminosity at the LHC is a subject matter of 2010
- Experiments are preparing for the first 2008 1 fb^{-1}



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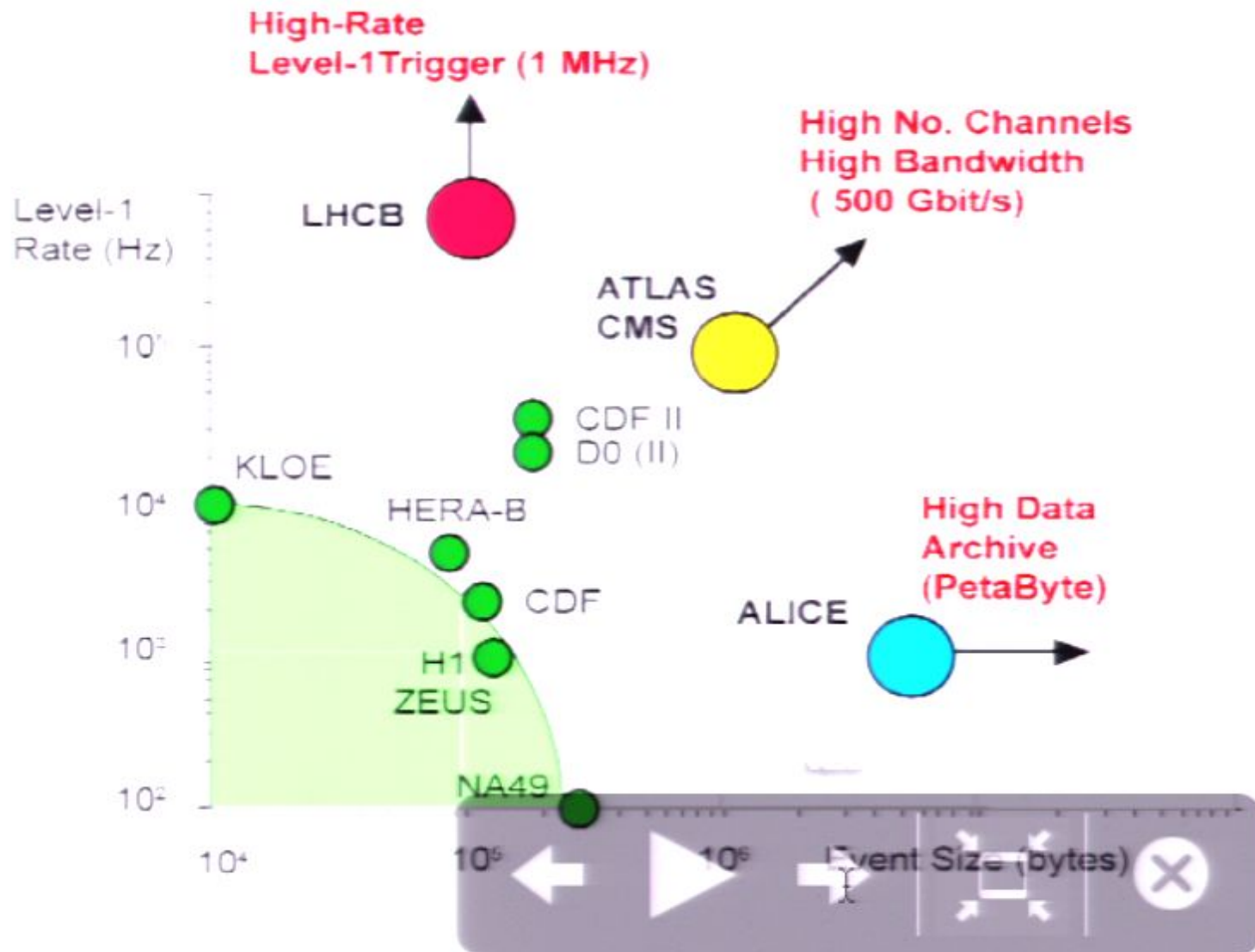
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the harvest [example]

process	cross section in fb
pair of 500 GeV jets	10,000
SUSY with 1 TeV superpartners	3,000
light Higgs	2,500
heavy Higgs	1,000
3 TeV Z'	2

- N.B. this is before taking account the trigger, detector etc efficiencies

A lot of data



The Answer is the Trigger/DAQ Architecture

- The triggers pipeline the data coming out of the detector subsystems (to buy a little time), then decide quickly which events are sufficiently interesting to pass along to the next layer of filtering.
- There are enough layers of triggers to reduce the original 40 MHz event rate down to 100-200 Hz of interesting events, which are then permanently "written to tape" for later analysis (analysis "offline")
- a good trigger is a fast trigger that keeps all the interesting events and is reprogrammable to respond to detector understanding and possible change of the physics targets

Nota Bene

- events that do not pass the trigger are rejected; so for example heavy charged slow particles can be missed in the trigger;
- Unlike the Tevatron the L1 at the LHC is all based on calorimetry and muon triggers.



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And the Grid

- The vast amount of data cannot be stored at any one single place
- the Grid architecture with the distributed analysis model is a huge global enterprise.
- The data from the LHC experiments will be distributed around the globe, according to a four-tiered model. A primary backup will be recorded on tape at CERN, the "Tier-0" centre of LCG. After initial processing, this data will be distributed to a series of "Tier-1" centres, large computer centres with sufficient storage capacity for a large fraction of the data, and with round-the-clock support for the Grid.
- The "Tier-1" centres will make data available to Tier-2 centres, each consisting of one or several collaborating computing facilities. Individual scientists will access these facilities through "Tier-3" computing resources, which can consist of local clusters in a University Department or even individual PCs, and which may be allocated to LCG on a regular basis.



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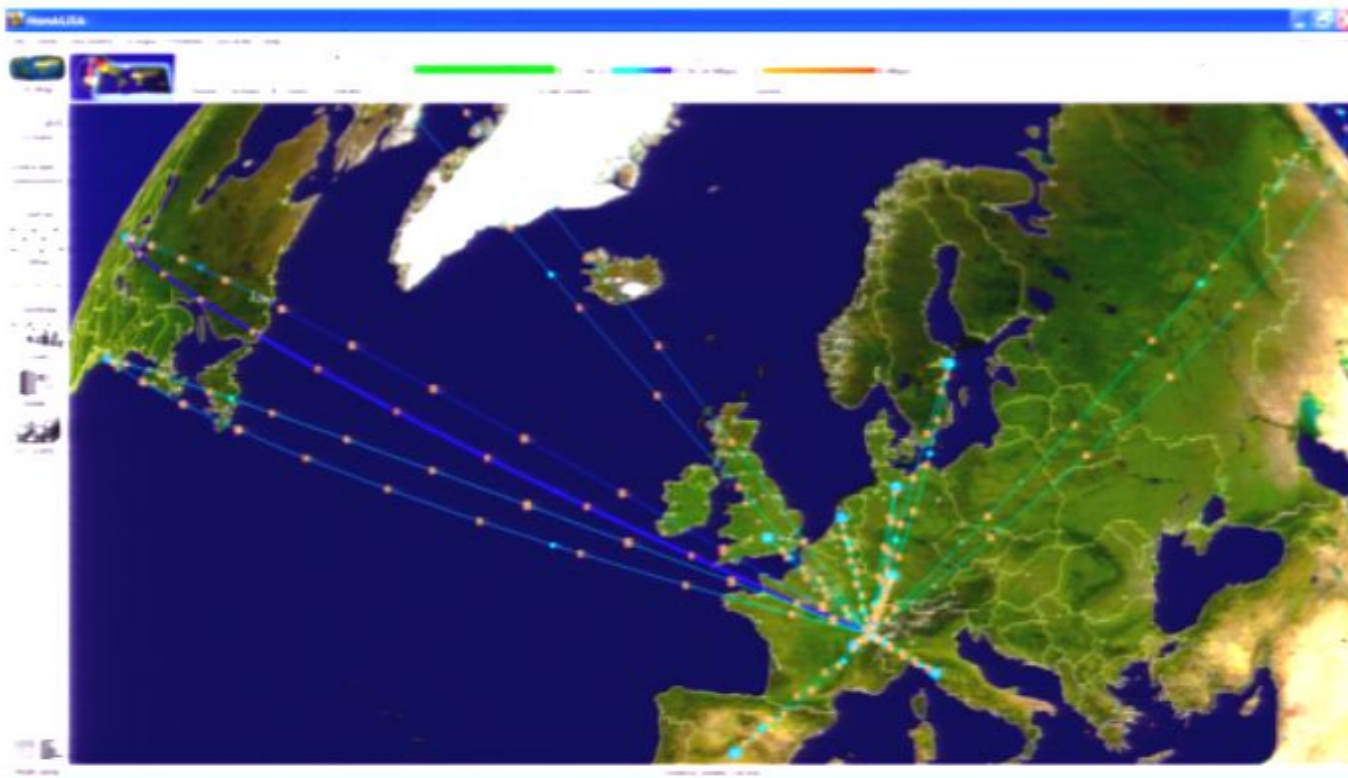


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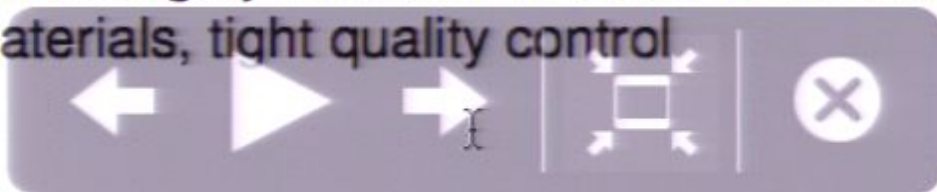
The GRID



- monitoring $>40,000$ computers, 120 computing centers, 40 countries.

A lot of data: Impact on Detector Design

- The large luminosity of the LHC is required such that we get a chance to capture the data from smaller than the standard model cross section processes.
- Even before we go the two general puprose detectors, we can see the impact on the detector design:
 - The detector subsystems must have fast repsonse [20-50 ns] → very challenging readout electronics
 - The segmentation should be fine to minimize the chance of overalap of say a photon from the Higgs event with a photon from a min-bias event from a neighboring bunch crossing → very large number of readout channels and high cost
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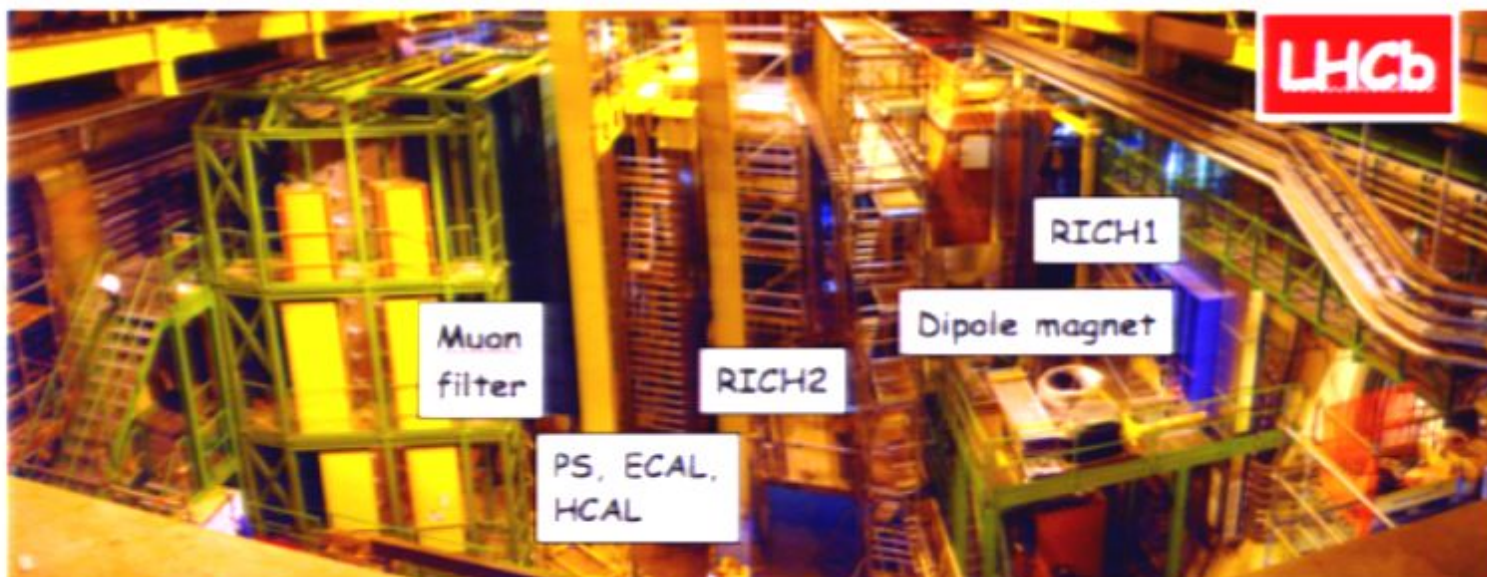
Outline

Status of the LHC Experiments

I



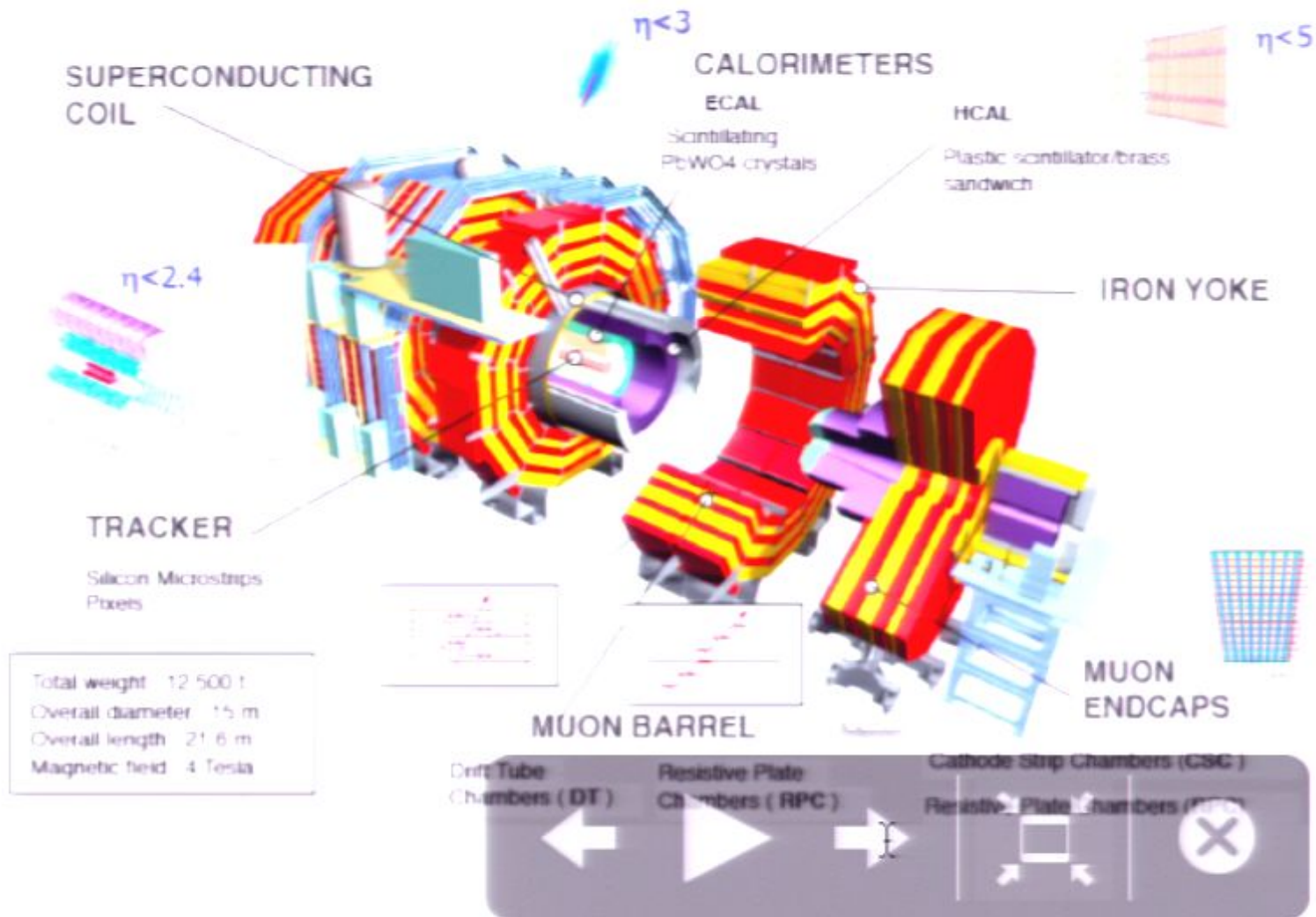
Status of ATLAS and CMS



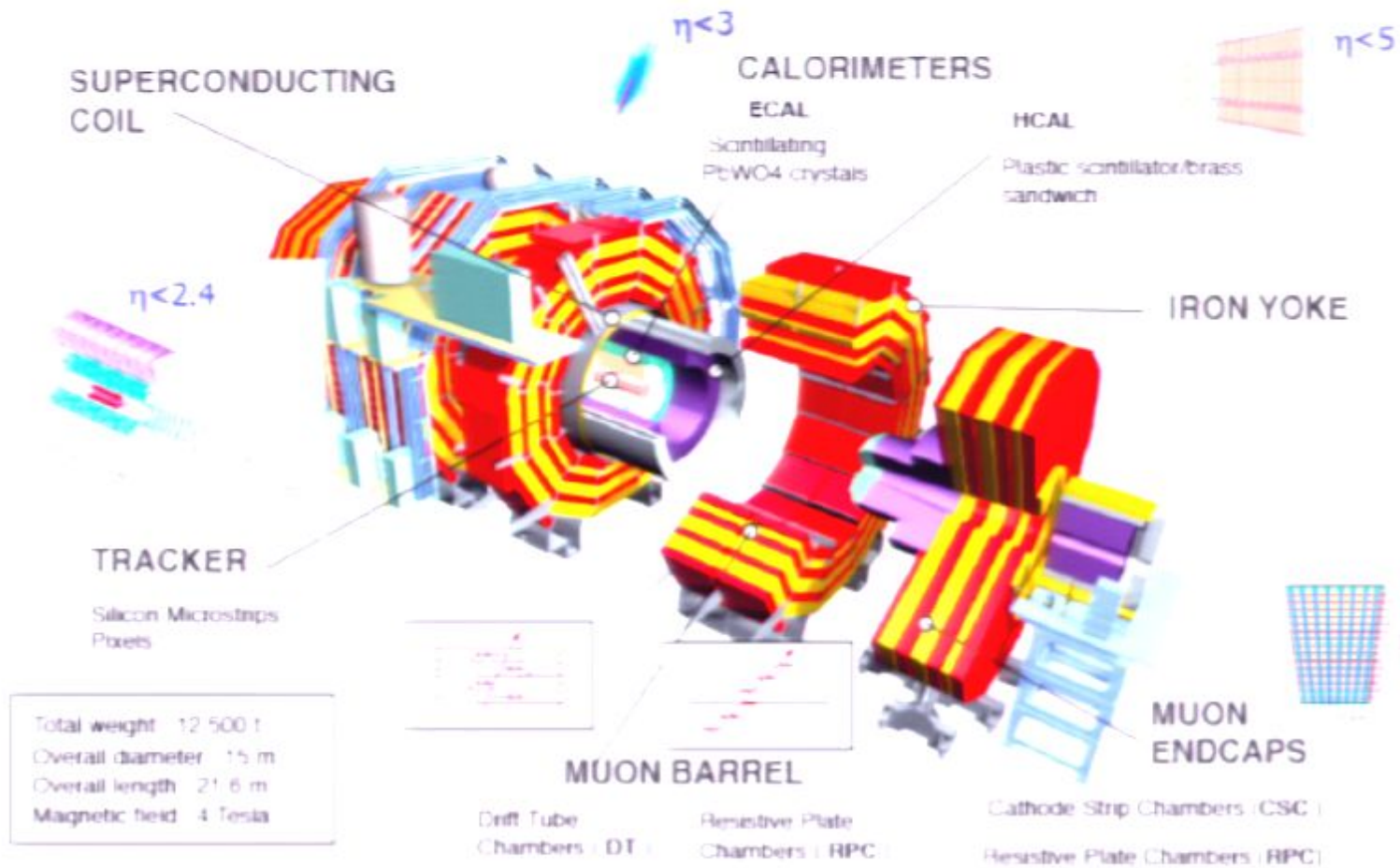
- LHCb and ALICE are on track
- here highlights from CMS and ATLAS



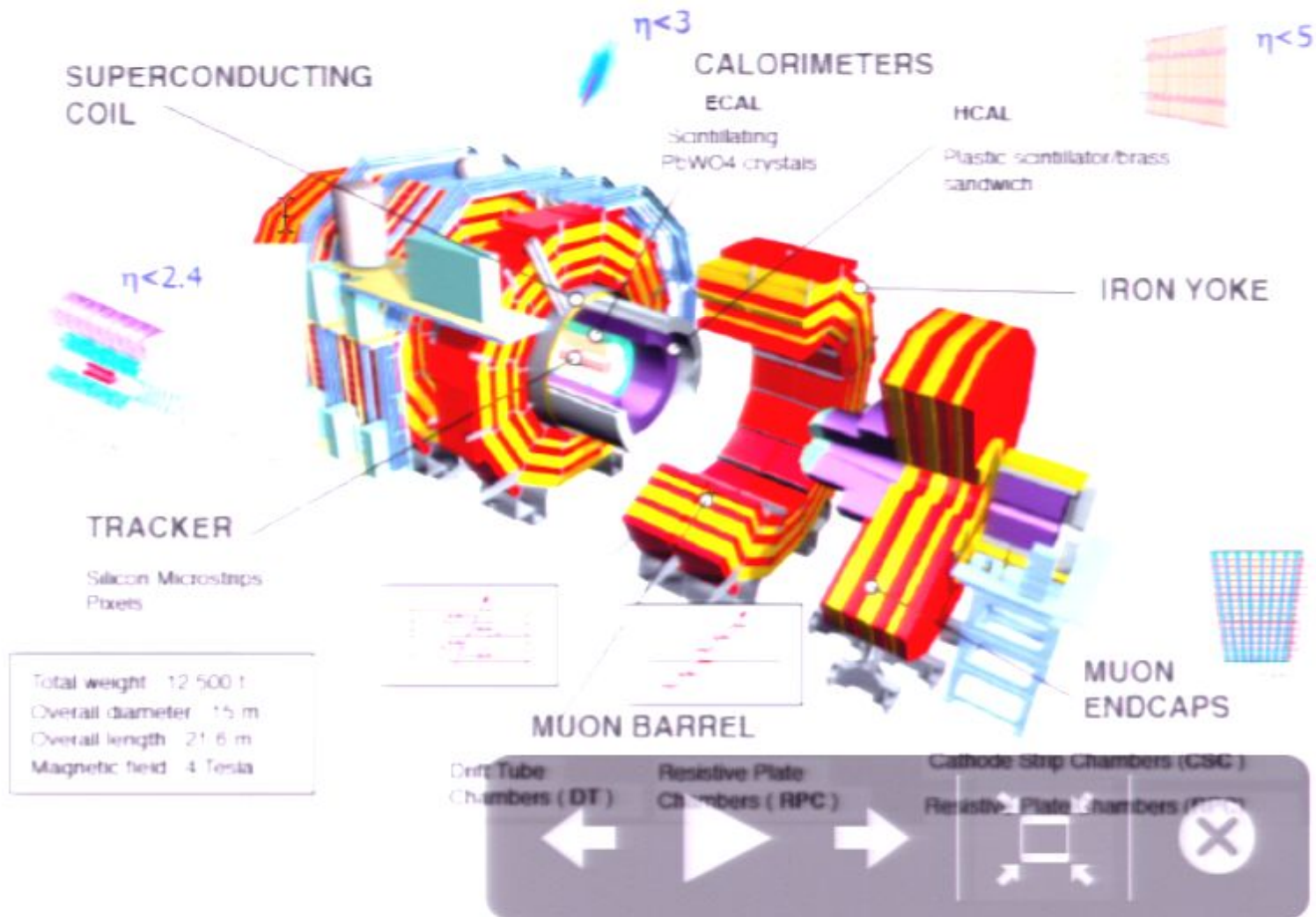
CMS: Compact Muon Solenoid detector at LHC



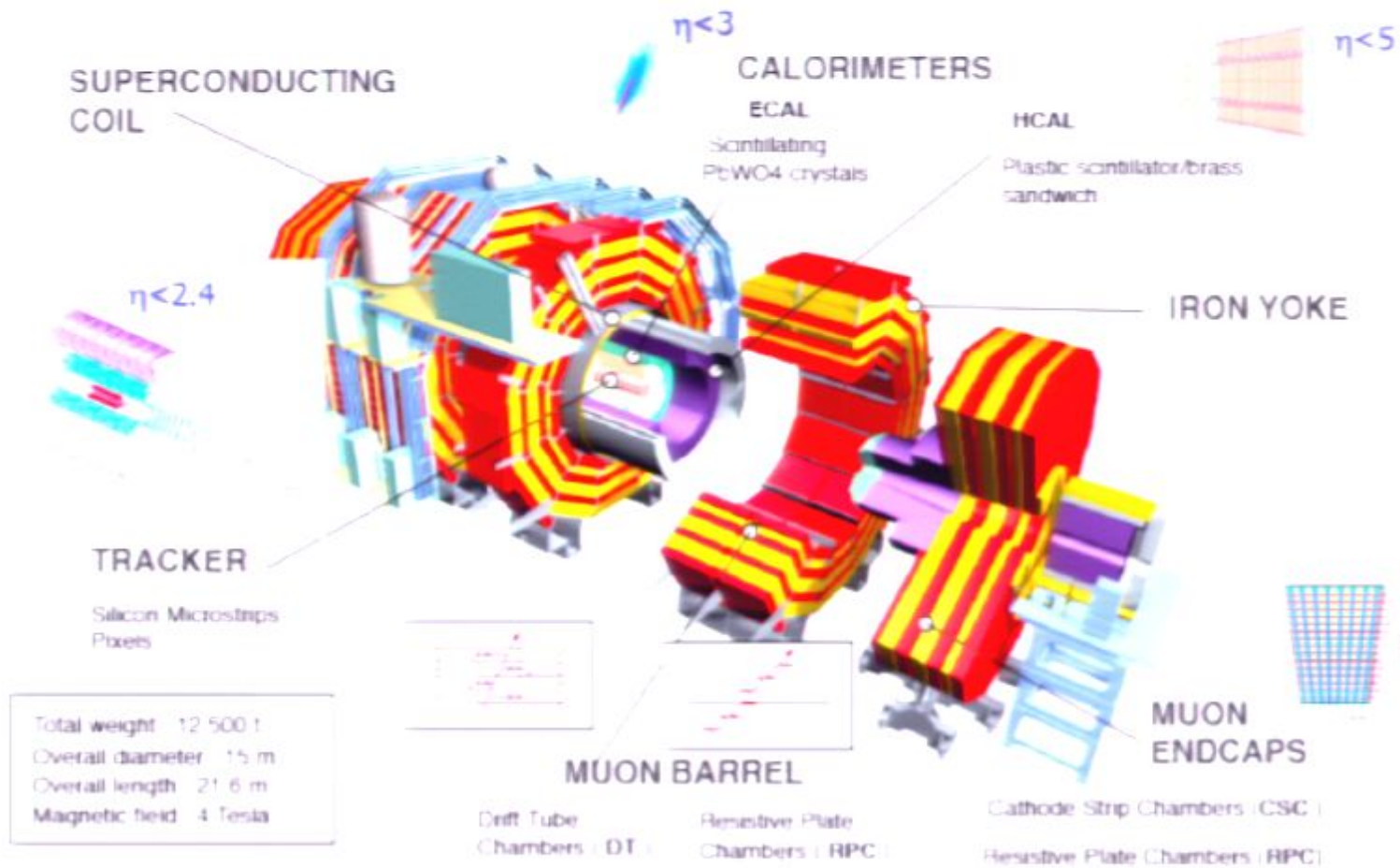
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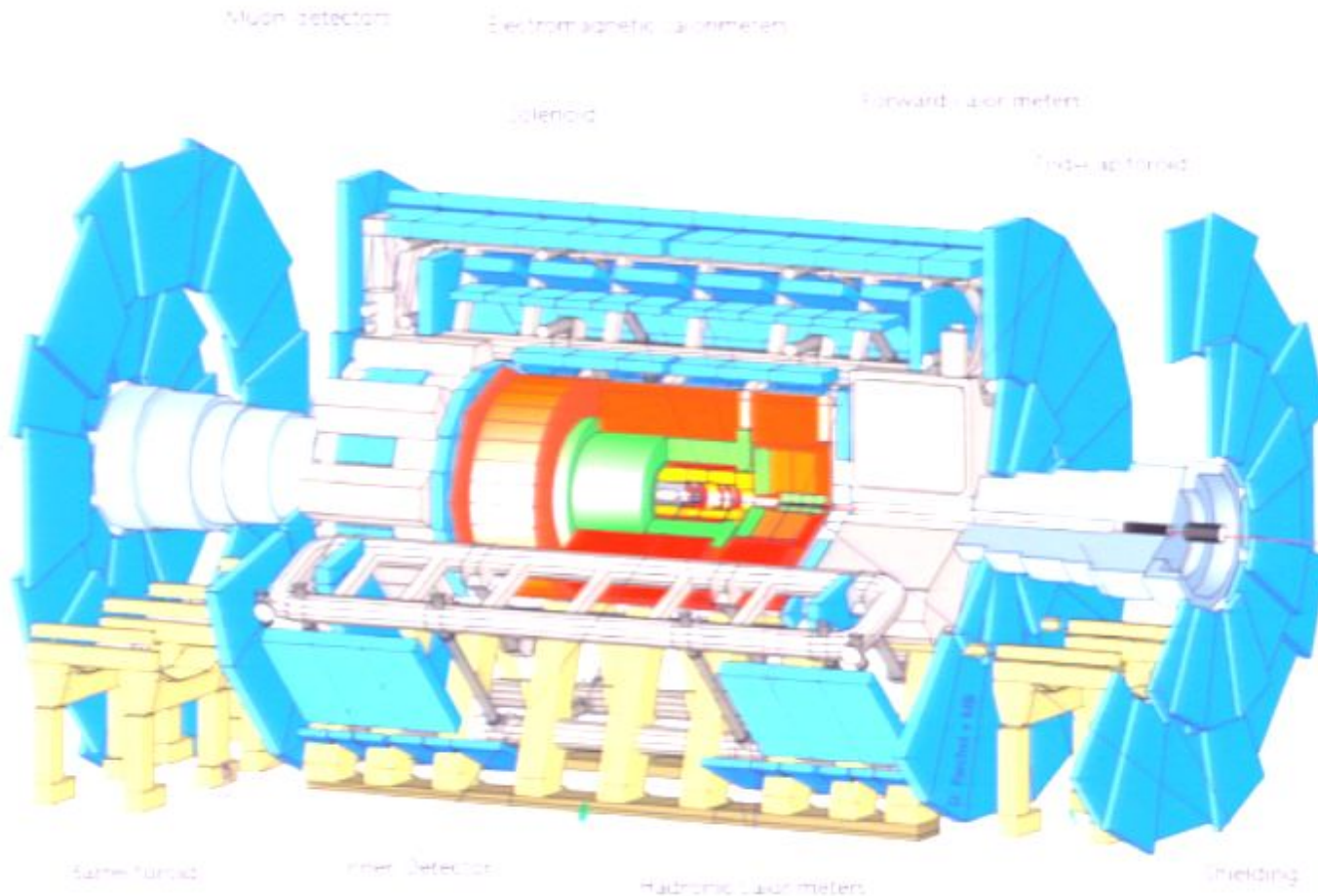
CMS: Compact Muon Solenoid detector at LHC



CMS: Compact Muon Solenoid detector at LHC



ATLAS: A Toroidal LHC Apparatus



• Tracking Si pixels and strips; Transition Radiation Detector •

Calorimetry EM : Pb-LAr with Accordion shape, HAD: Fe/scintillator

Cu/W-L Ar : • Muon Spectrometer : air-core toroids with muon



	ATLAS	CMS
weight (tons)	7,000	12,500
diameter (m)	22	15
length (m)	46	20
magnetic field for tracking (T)	2	4
solid angle coverage $\Delta\eta \times \Delta\phi$	$2\pi \times 5.0$	$2\pi \times 5.0$
cost in M CHF	550	550

Object	Weight (tons)
Boeing 747 [fully loaded]	200
Endeavor space shuttle	368
ATLAS	7,000
Eiffel Tower	7,300
USS John McCain	8,300
CMS	12,500



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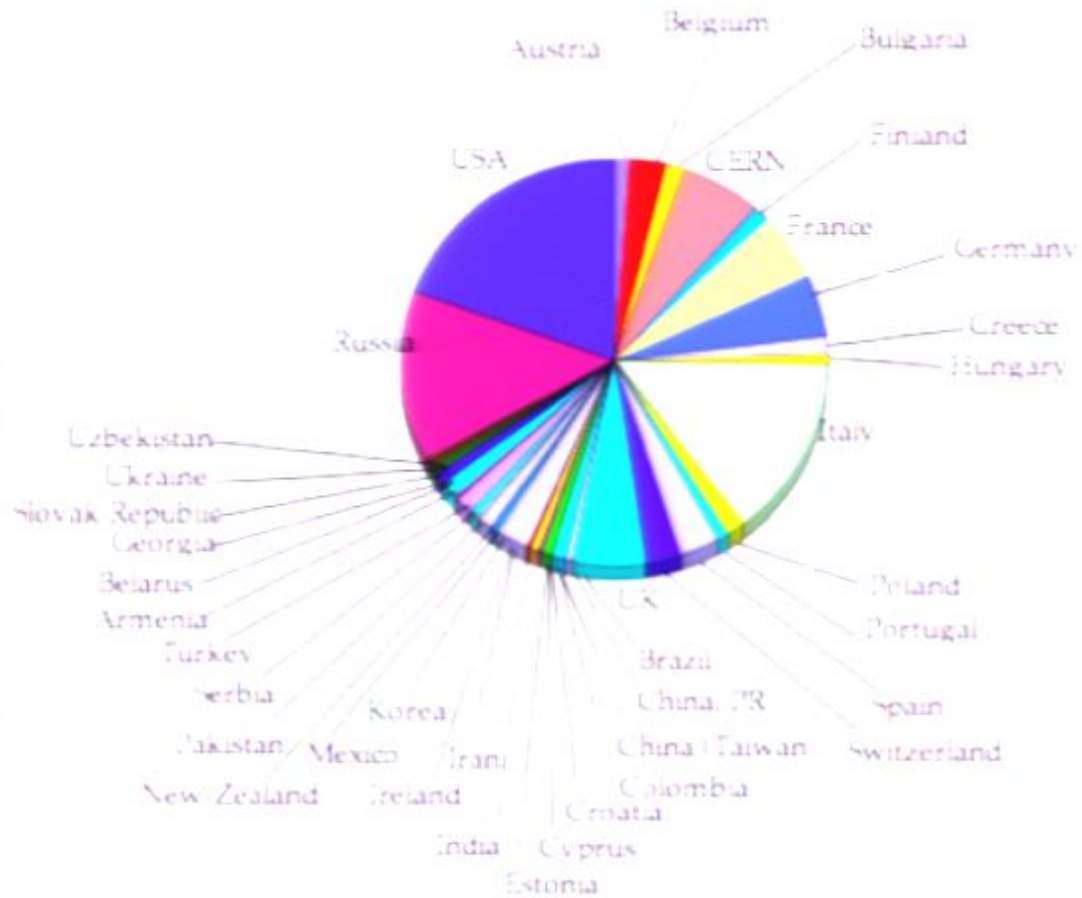
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	Number of Laboratories
Member States	61
Non-Member States	64
USA	49
Total	174

	Nr of Scientific Authors
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Non-Member States	428
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Total	2030

Associated Institutes	
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Number of Laboratories	8

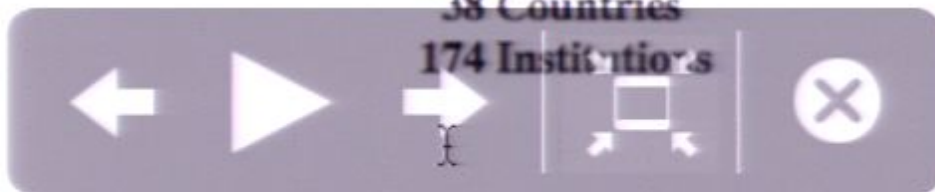


2030 Scientific Authors

38 Countries

174 Institutions

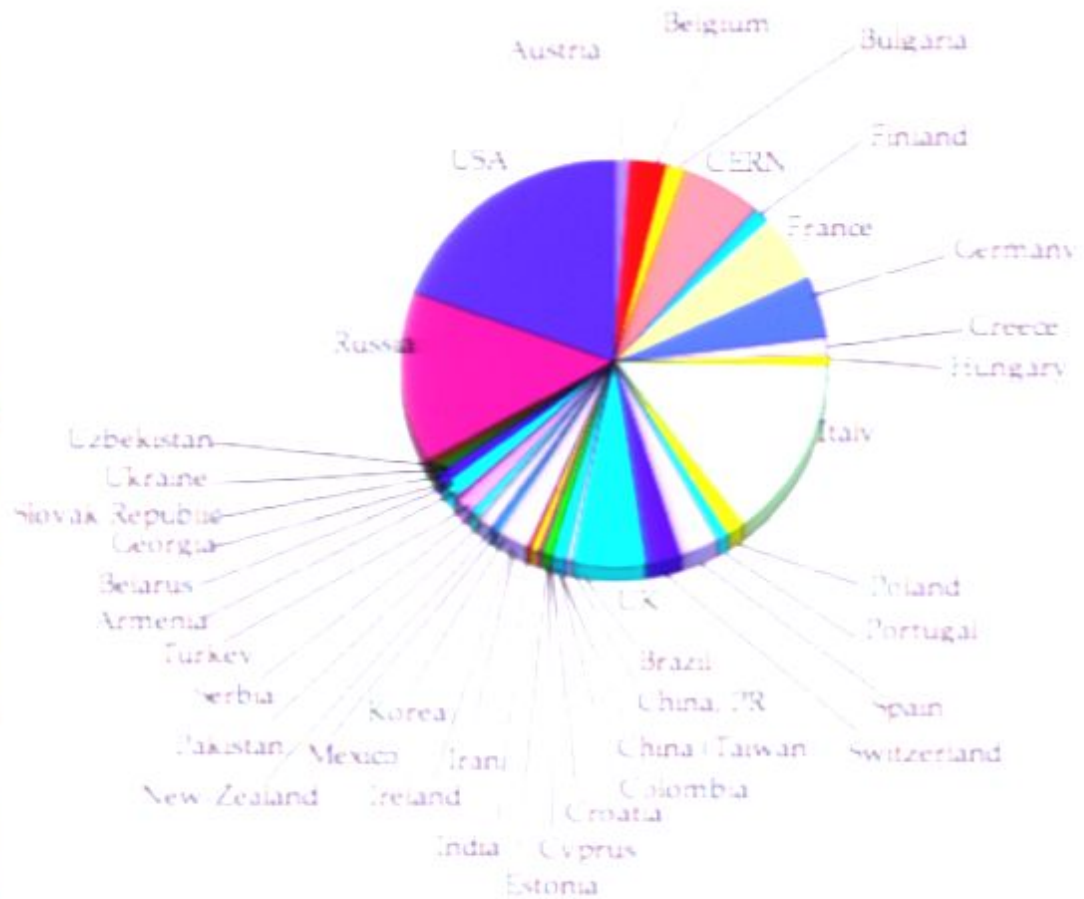
May 14 2006 pm
<http://www.cern.ch/press/rel/050606/050606.html>



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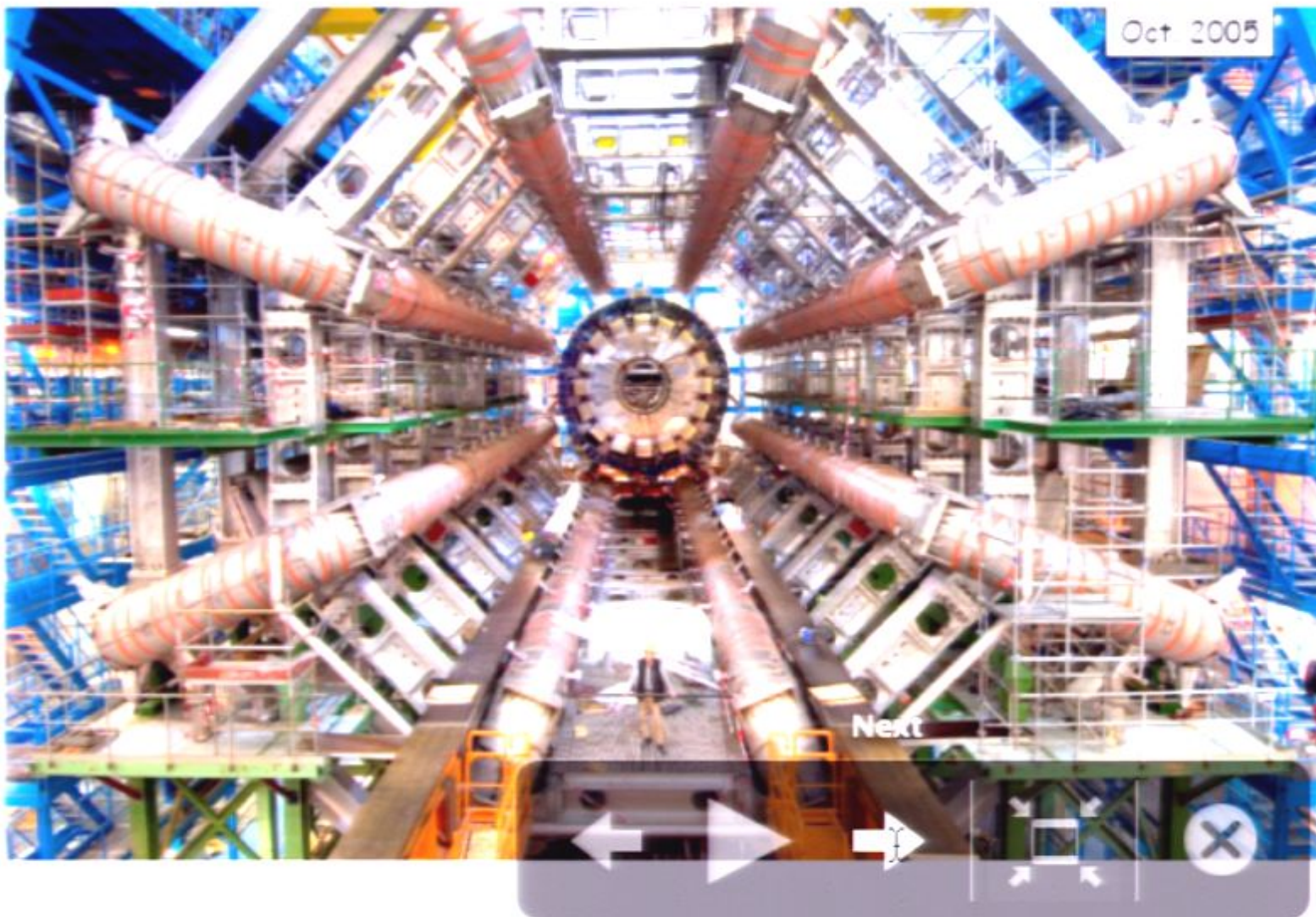


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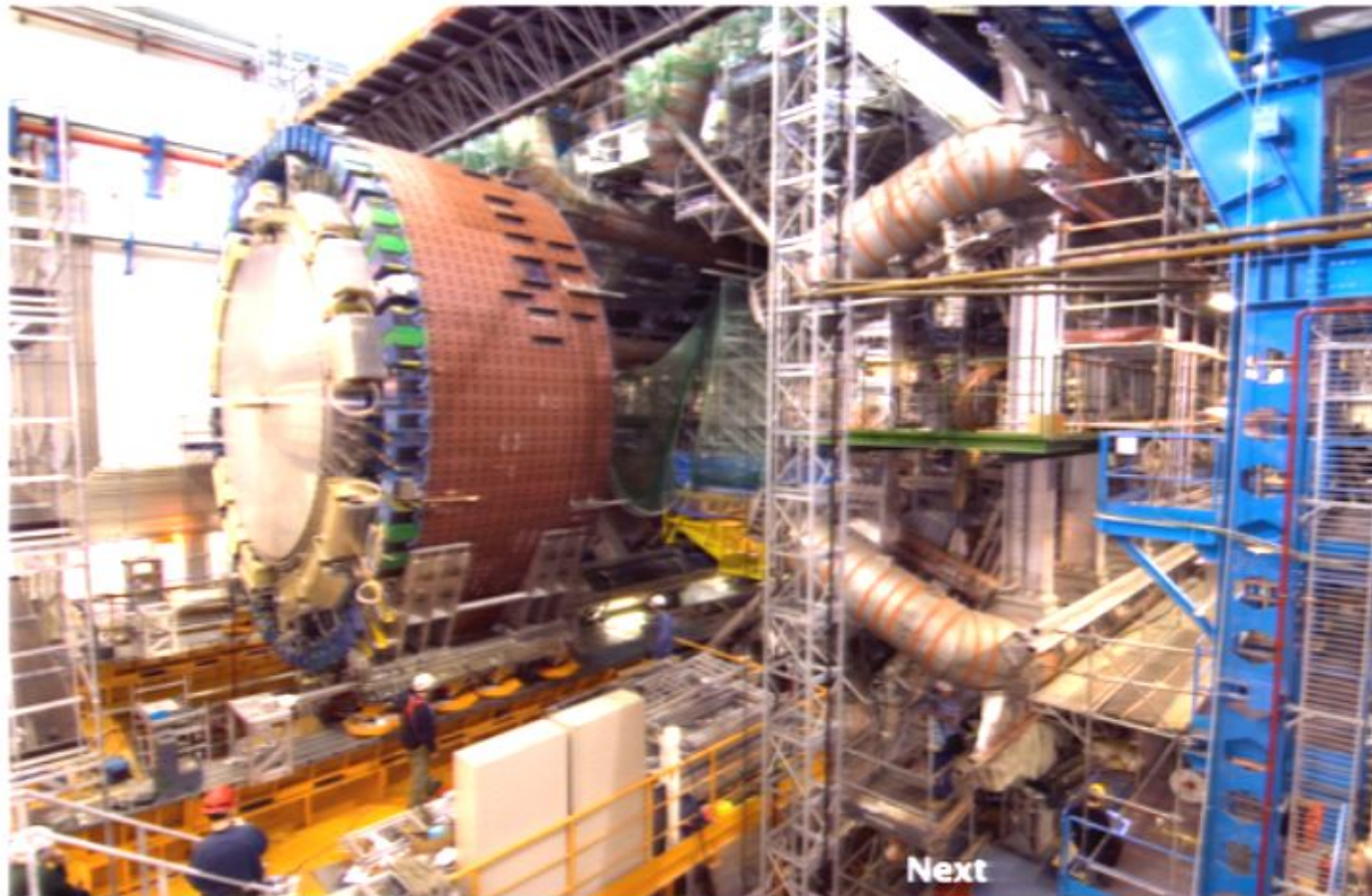
May 14 2006 gm
<http://lhc.cern.ch/press/rel/2006/05/14/06051401.html>



ATLAS in cavern-2005



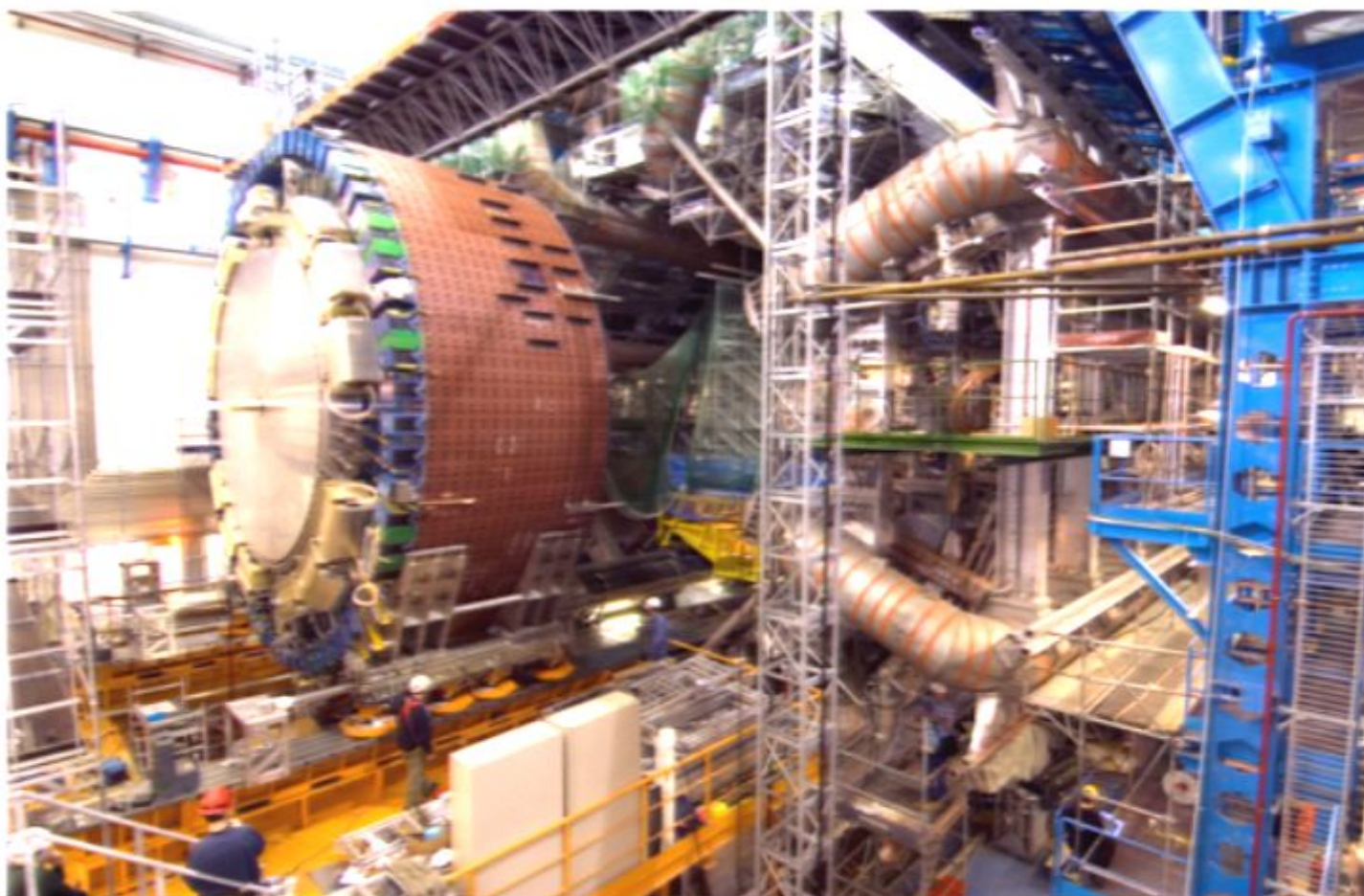
ATLAS EndCap Calorimeter



- One end-cap calorimeter (LAr EM, LAr HAD, LAr Forward Tilecal) being moved inside the barrel toroid



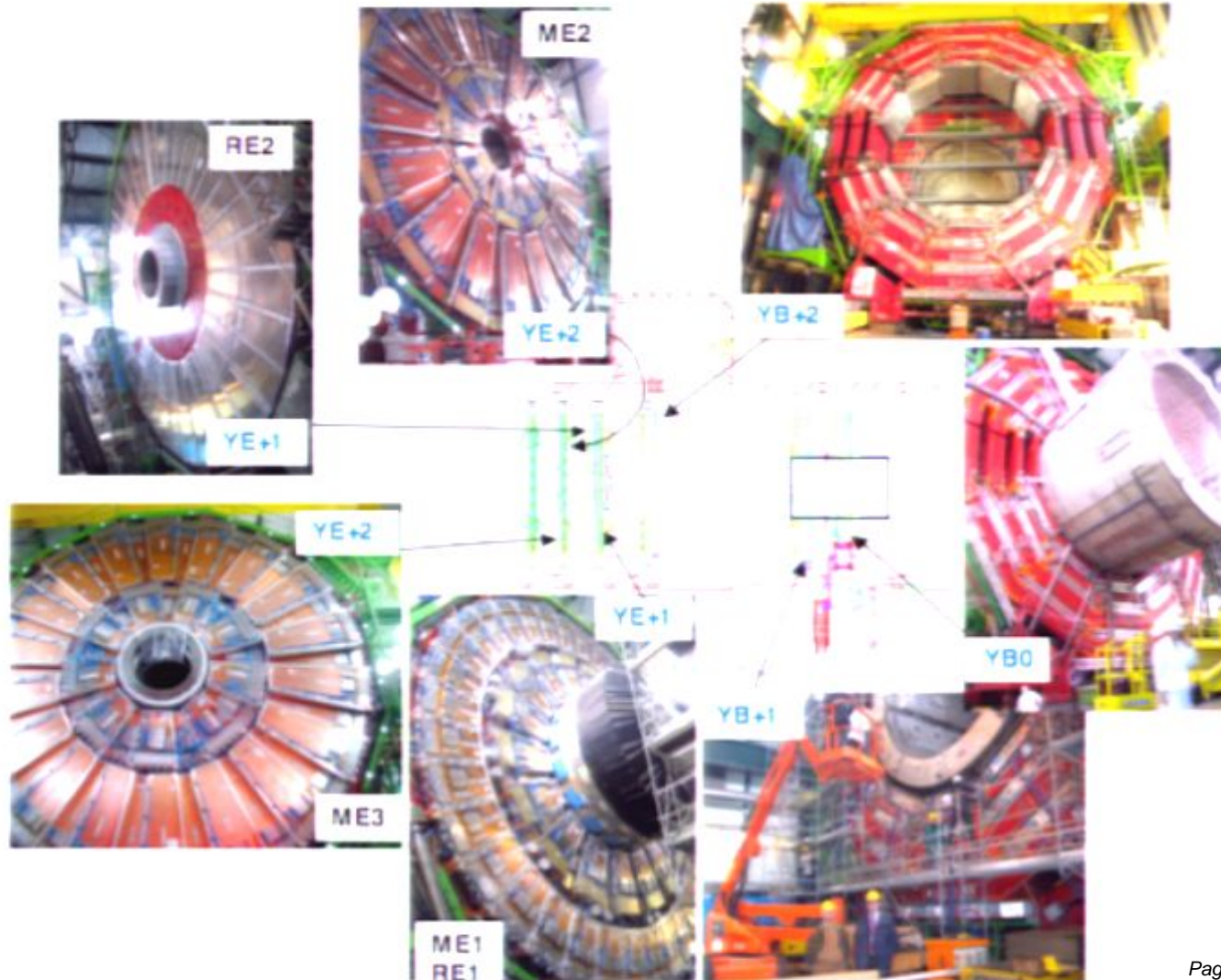
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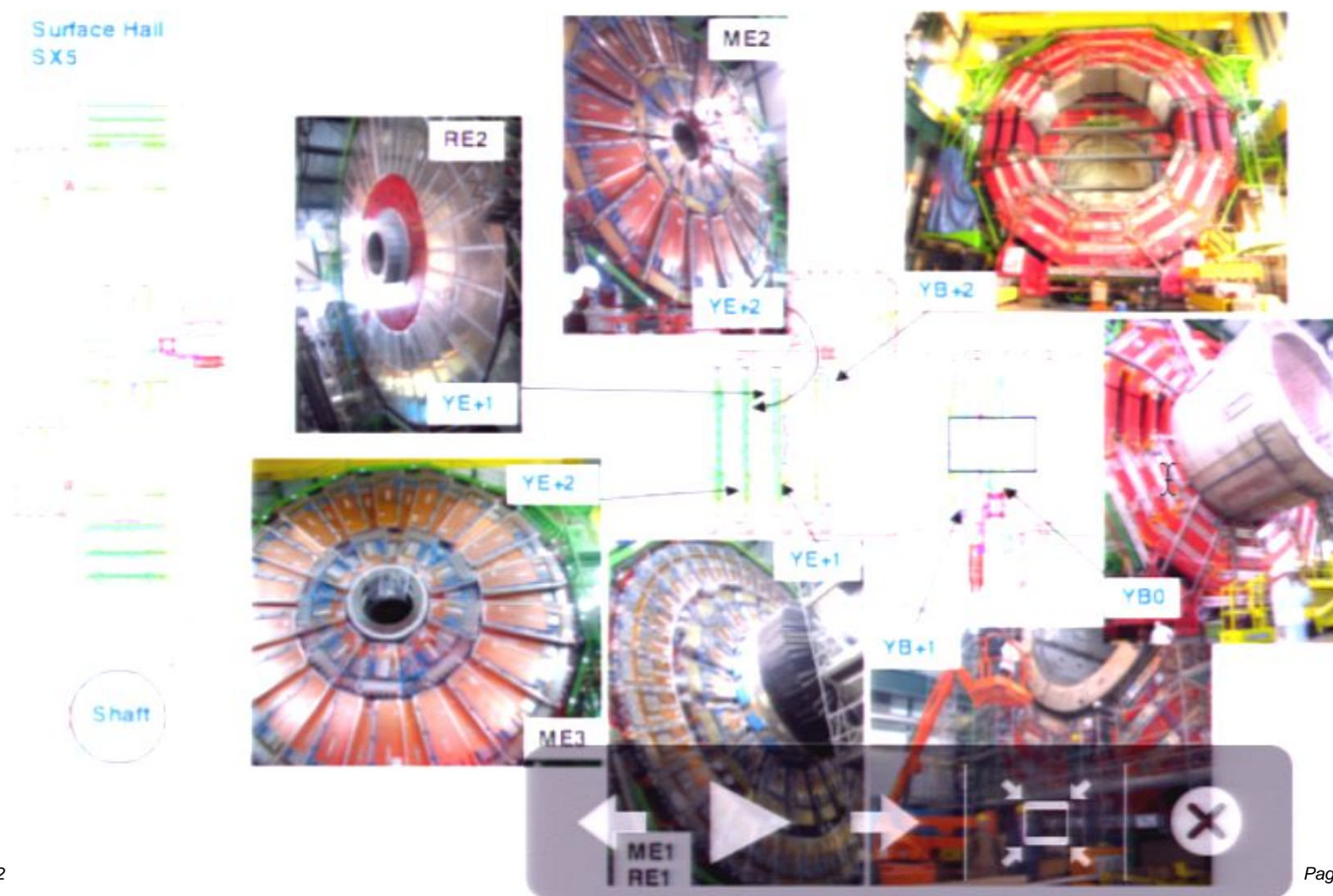
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CMS surface assembly



CMS surface assembly



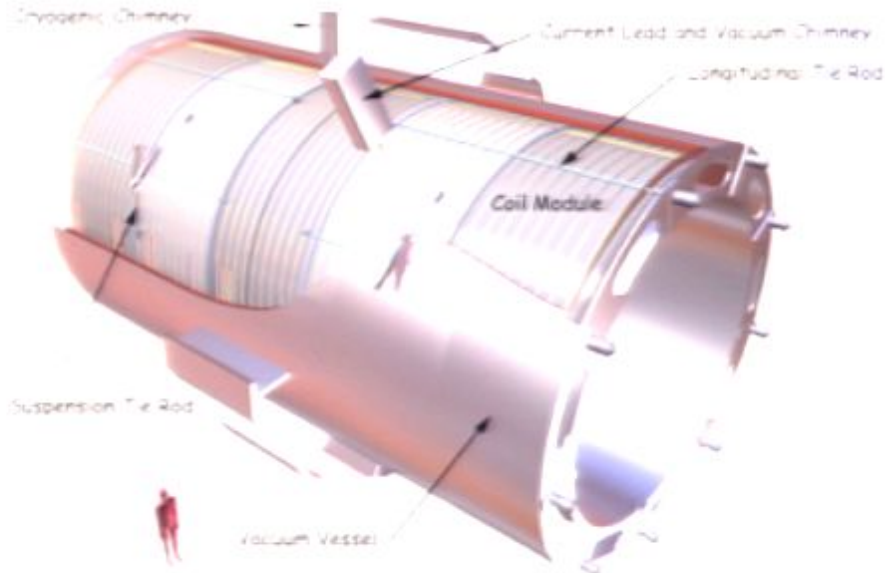
1. The LHC will determine the future of particle physics -

2. We are lucky that is LHC is happening now (fine tuning?)

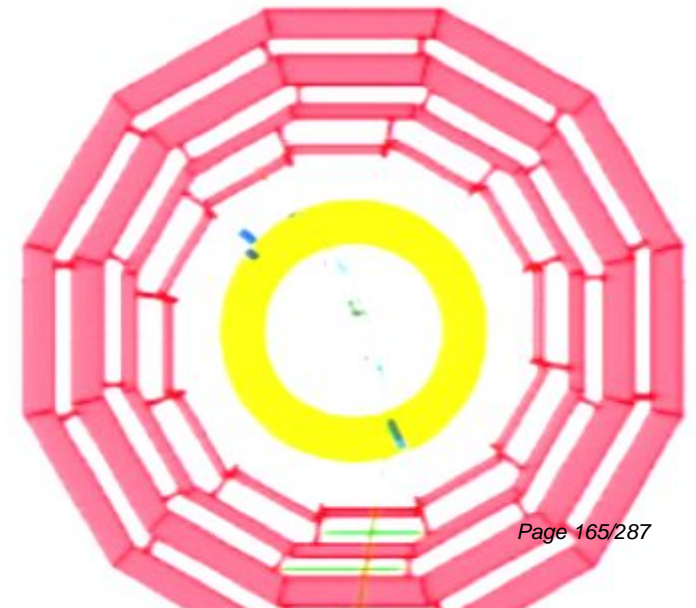
3. It will not be a walk in the park



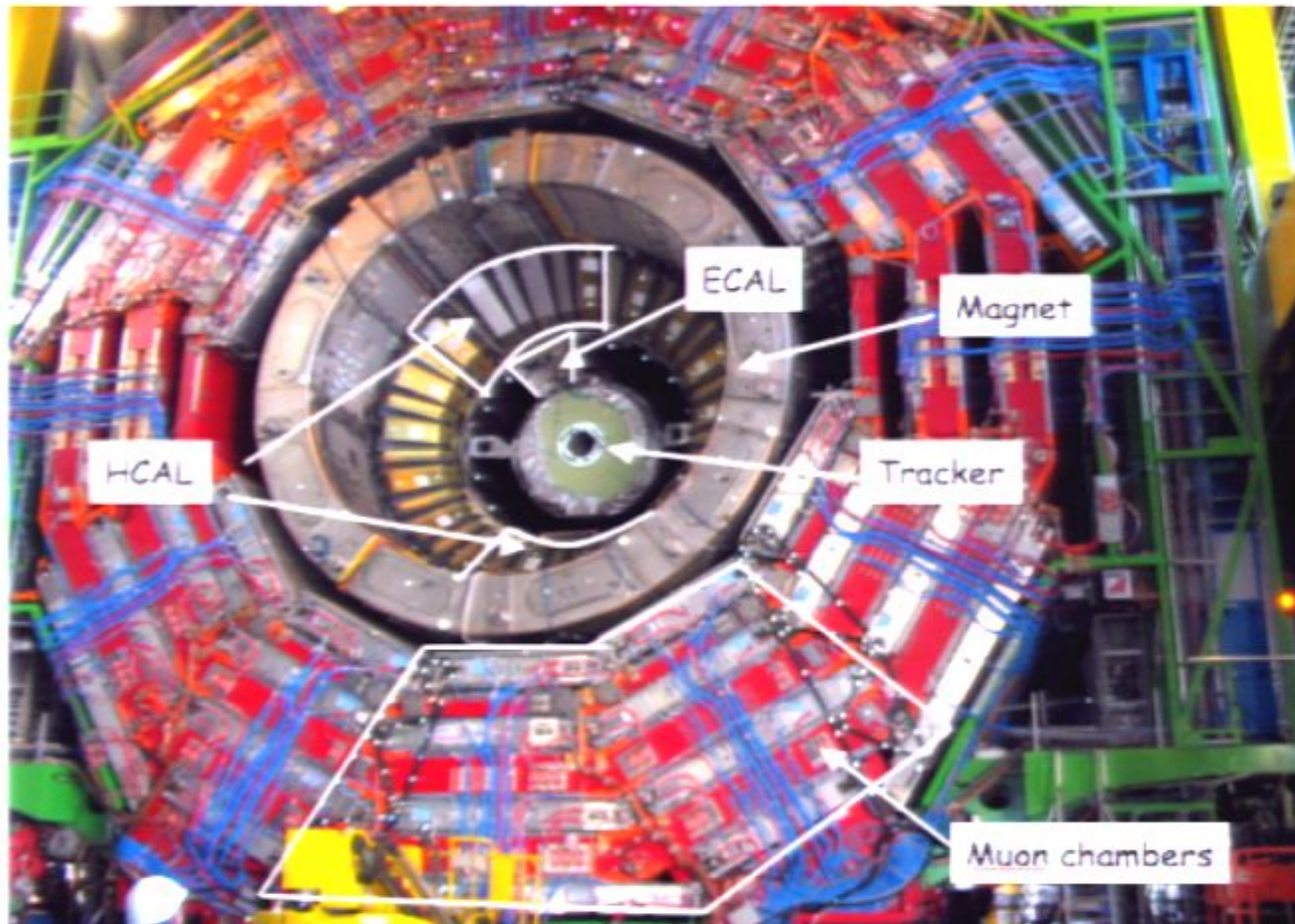
CMS superconducting solenoid



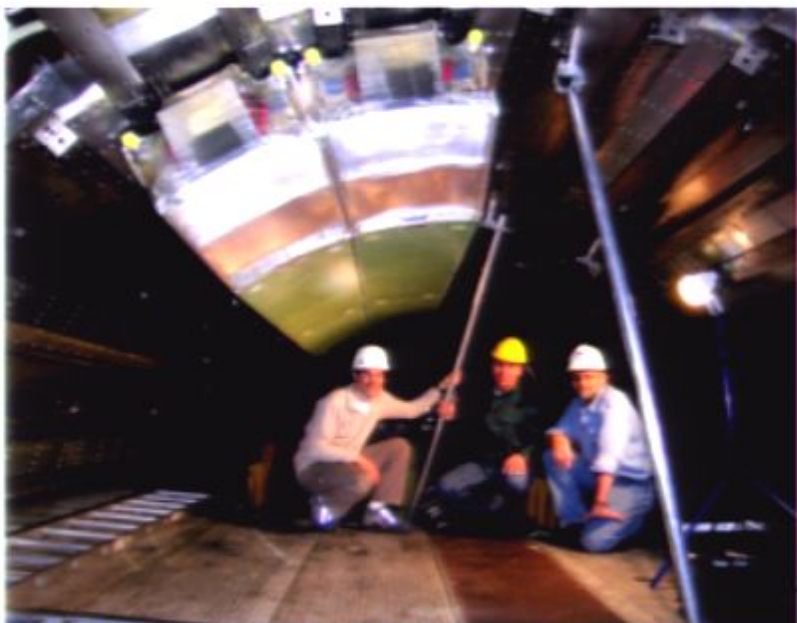
- Magnet length 12.5 m; Diameter 6 m; Magnetic field 4 T; Nominal current 20 kA; Stored energy 2.7 GJ! ; the most energetic magnet ever built+you can drive a truck into



CMS surface configurations for the Magnet and Cosmic Slice Test



CMS ECAL



ALL 61.5k PbWO₄ barrel crystals have been delivered. Endcap crystal delivery from Russia and China is ongoing and foreseen to end in March 2008. First endcap Dee will be ready for insertion end of Feb. 2008. Construction of second endcap will be finished by early summer 2008.



CMS continues detailed tests of Gravity



CMS continues detailed tests of Gravity



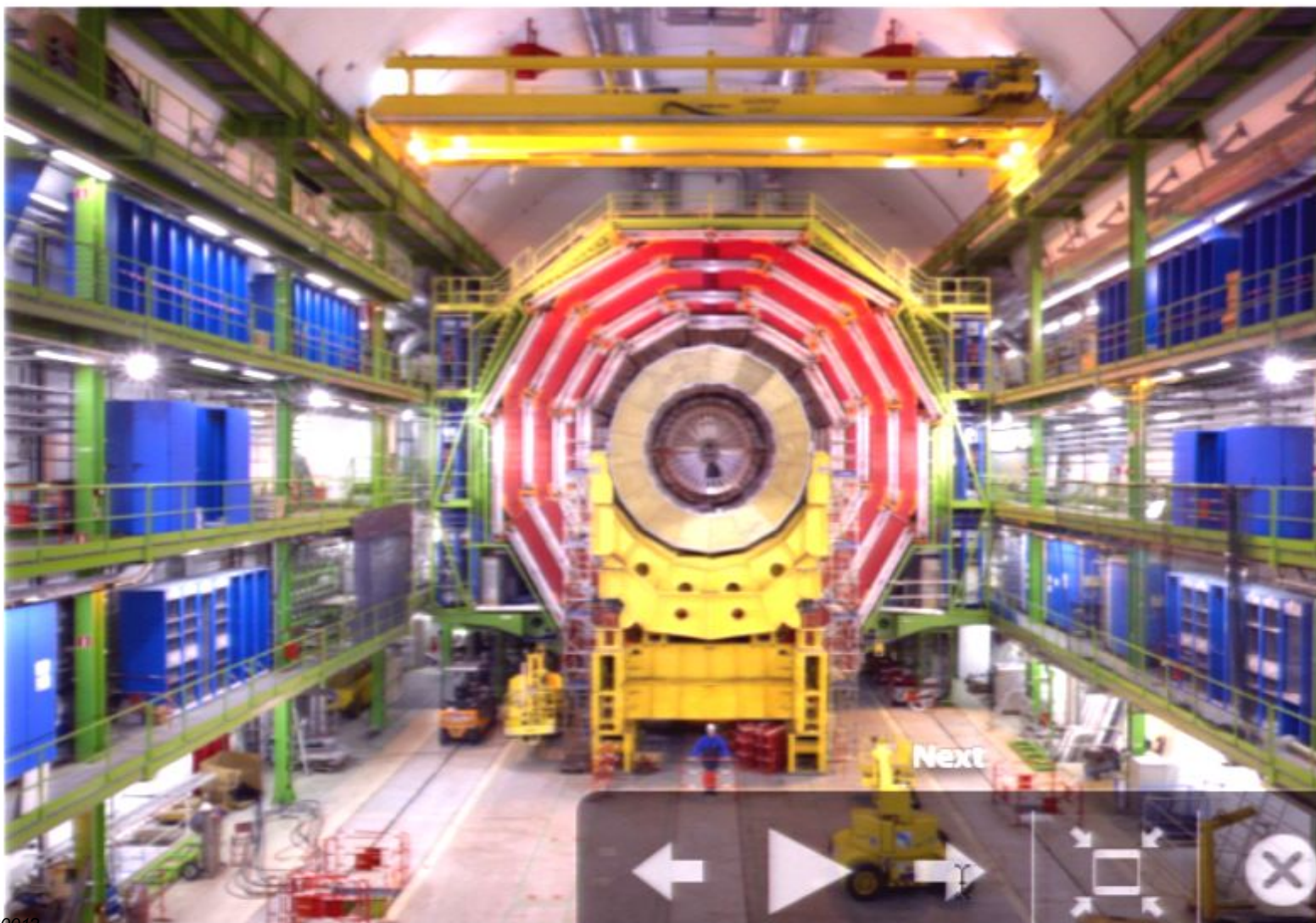
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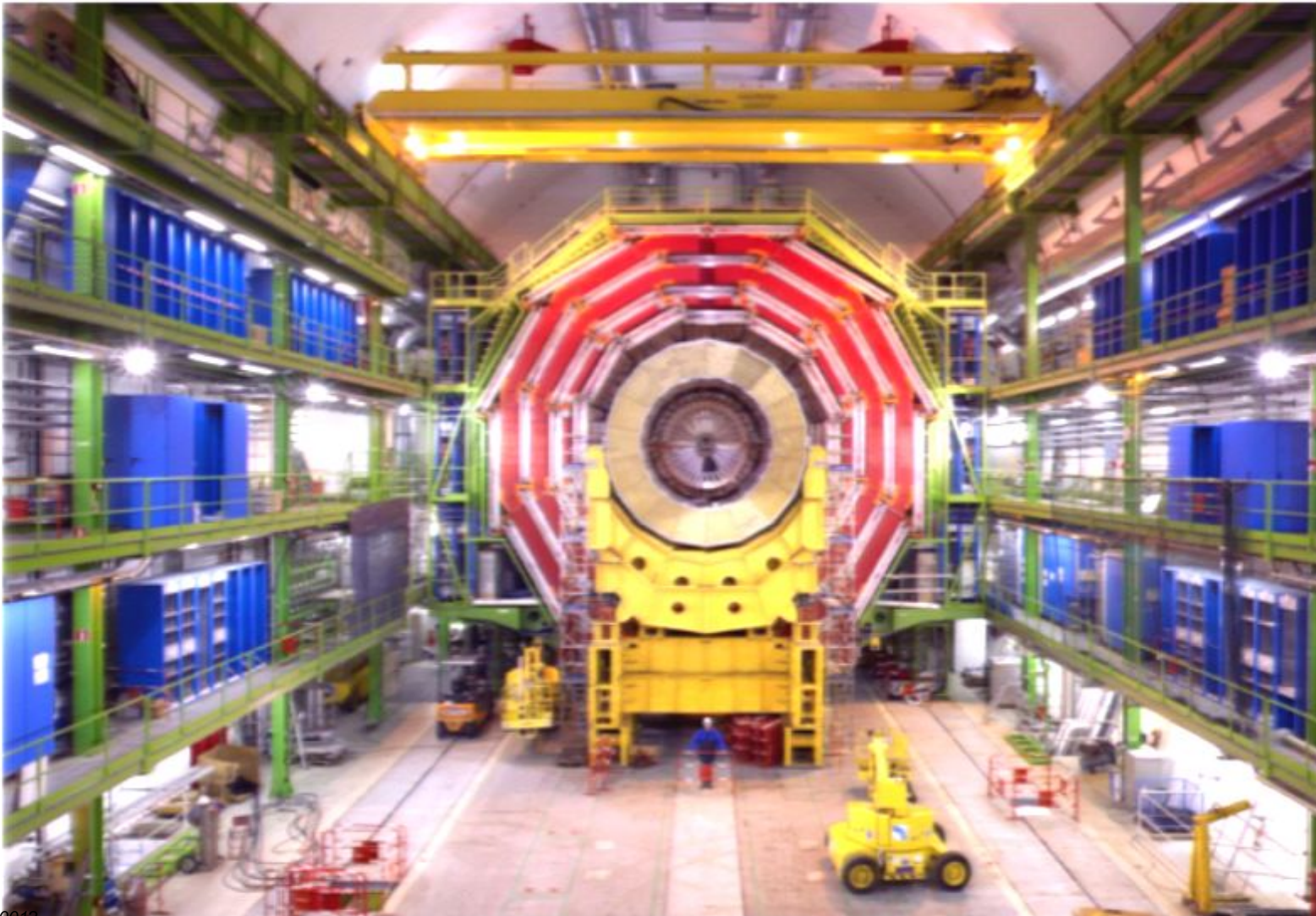
Next



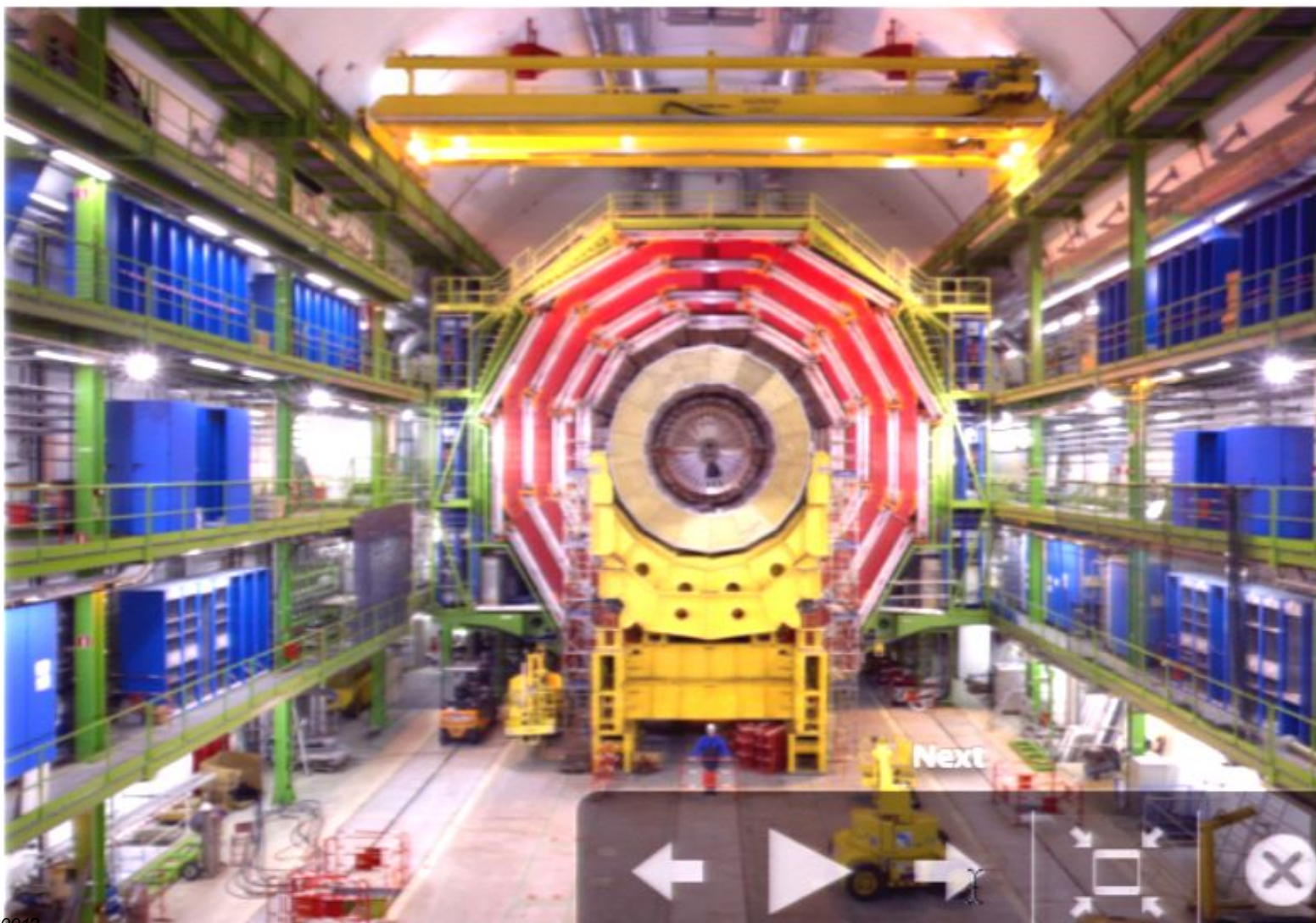
CMS at the pit now



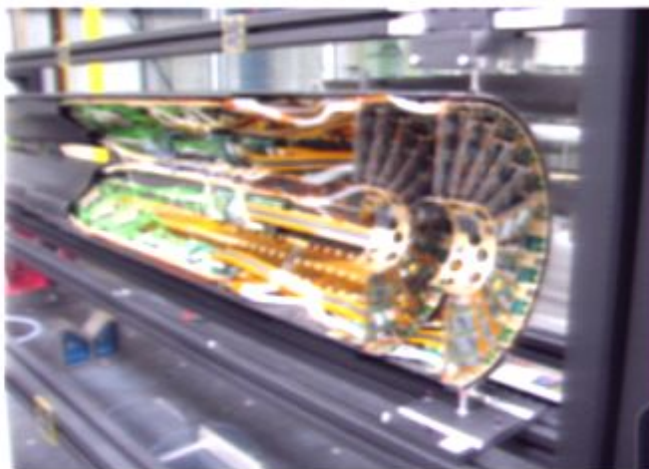
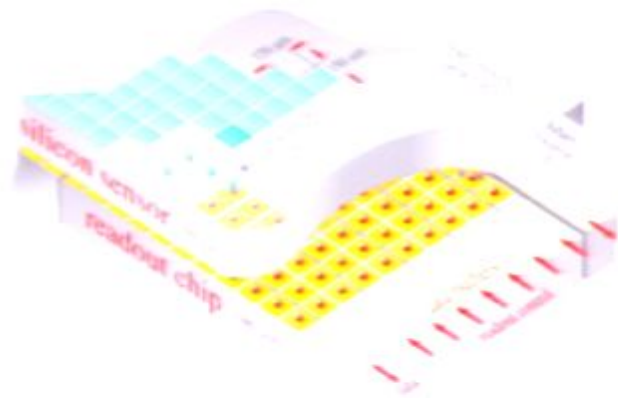
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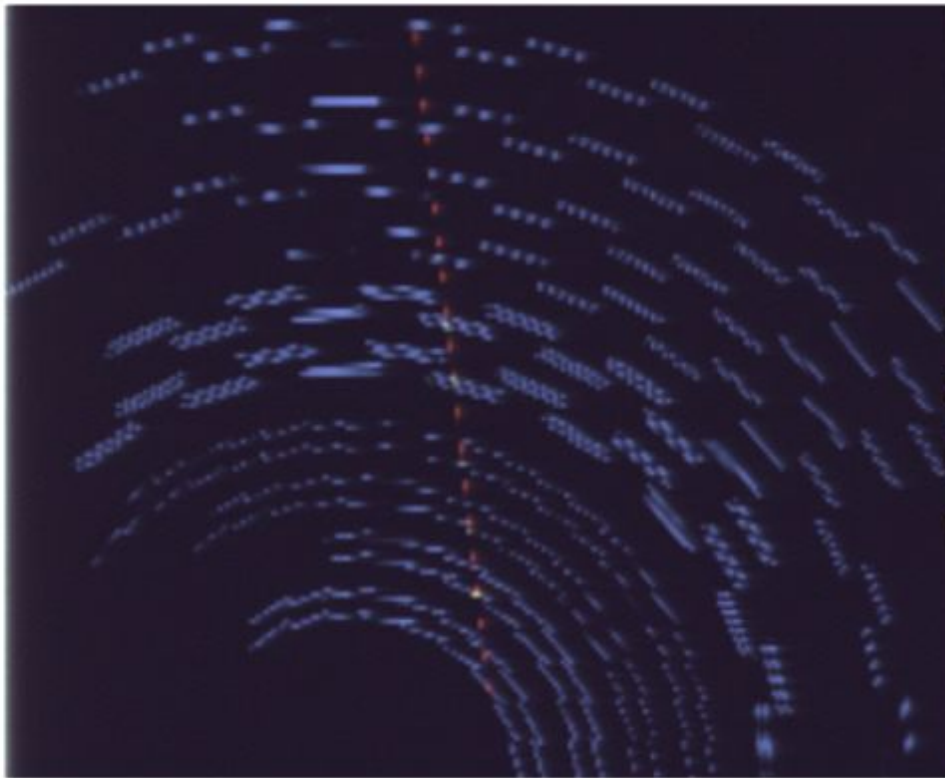


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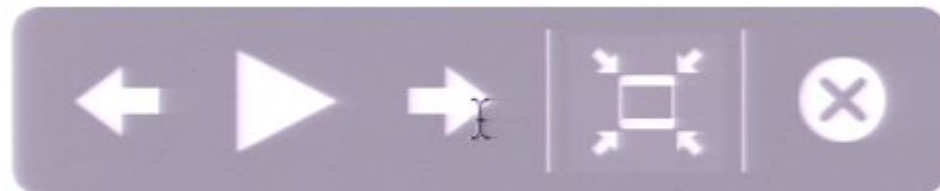


Silicon (pixels (L), SST (R))





alignment goal: 30μ 04.08
SST installation: starting
SST cosmic commissioning:
 $\mathcal{O}(M)$ recorded cosmic events
Pixel installation: <summer 08



Schedule

Both experiments are in status of commissioning and integration in preparation for data in summery 2008.

May 07 and few days every end of month	commissioning/global runs/ cosmics/test data
Sep. 07	CSA07
Oct-Dec 07	cosmics/integration
Feb 08	beam pipe close
++	install pixels/ECAL EE (one)
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Apr 08	close CMS for 14 TeV run
Jun 08	14 TeV physics run



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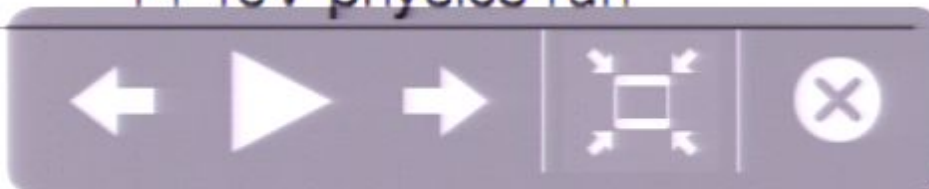
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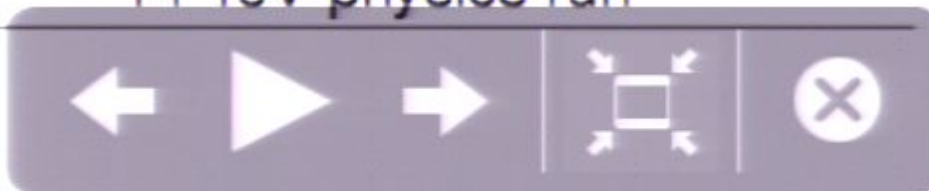
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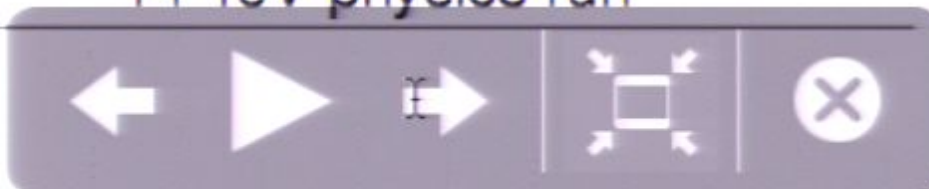
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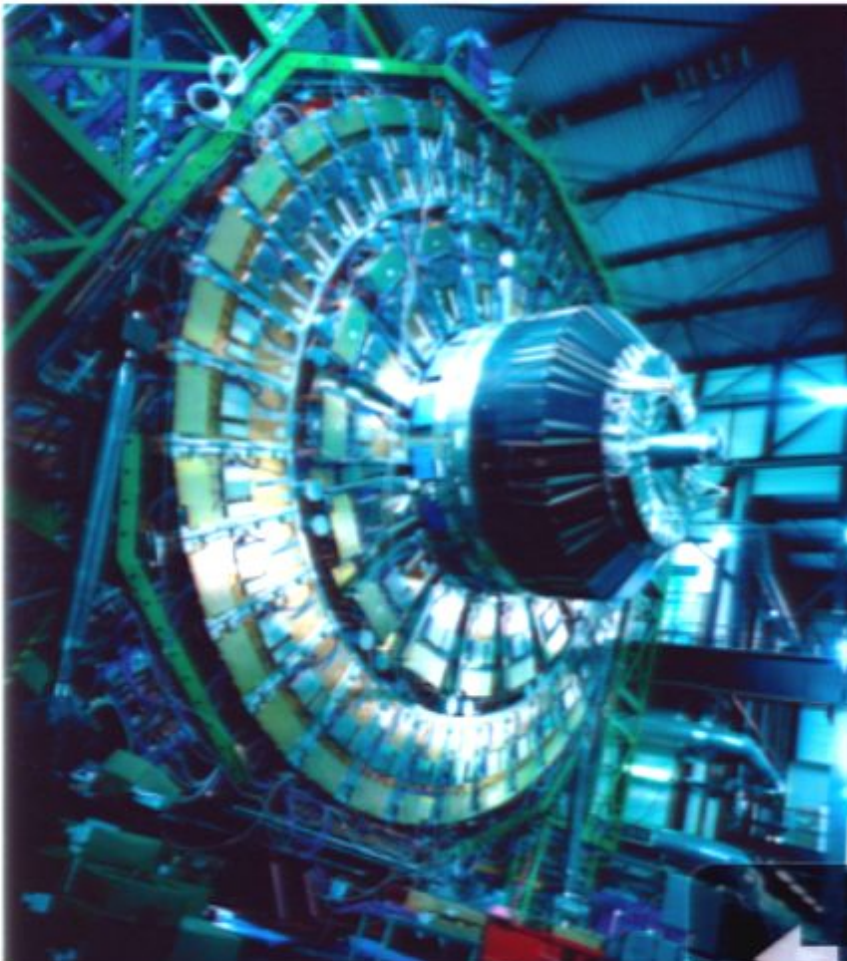
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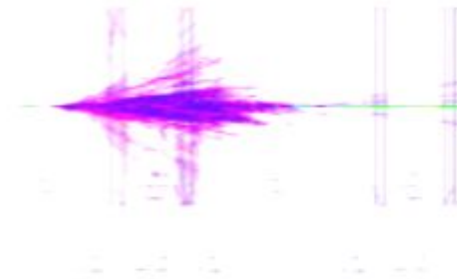


big and challenging

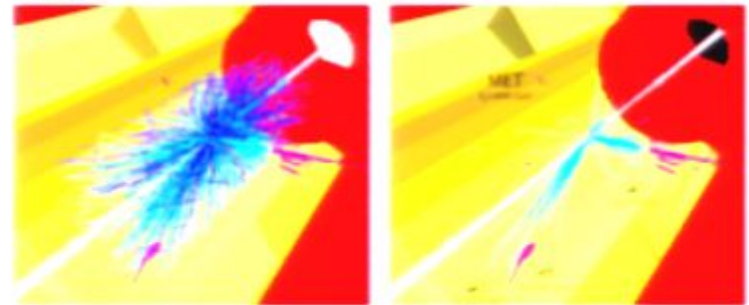


- high energy

Test it out with a "catastrophic" energy loss of 22 GeV



- high luminosity

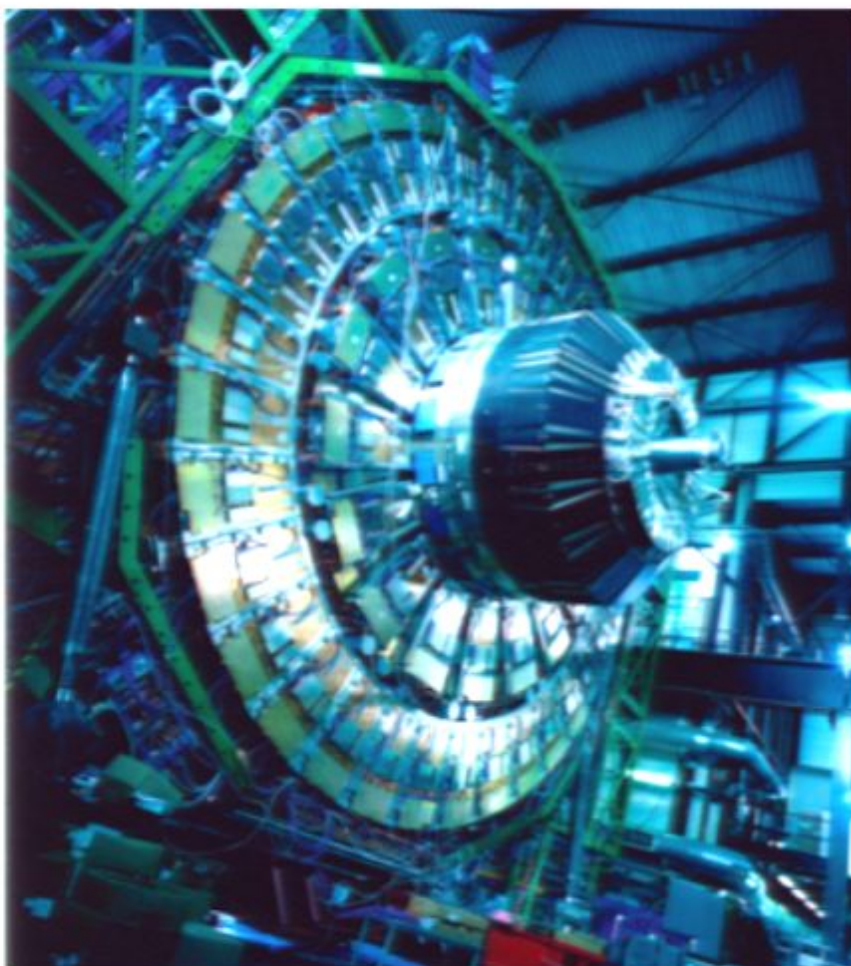


- high data rates (**trigger**, **GRID**)

- ...not a walk in the park

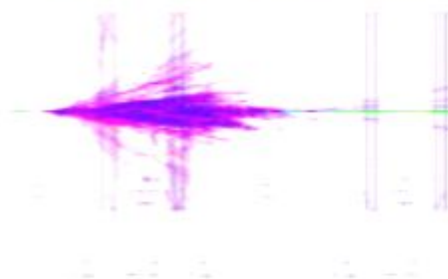


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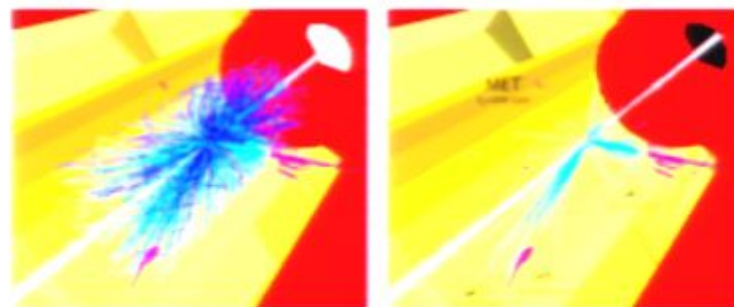


- high energy

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- high luminosity



- high data rates (trigger, GRID)
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Detector Concept

- The detectors designed to identify most of the very energetic particles emerging from the proton-proton collisions, and to measure as efficiently and precisely as feasible their trajectories and momenta.
- The interesting particles are produced over a wide range of energies (from a few hundred MeV to a few TeV) and over the full solid angle. They therefore need to be detected down to small polar angles (θ) with respect to the incoming beams (a fraction of a degree corresponding to pseudorapidities η up to 5, where $\eta = -\log[\tan(\theta/2)]$).
- No detectable particle should escape unseen.
- The design is "cylindrical onion"



The homework assignment is then...

Why ATLAS and CMS are so differently designed?



you can say they are "dual"

Clue1: Benchmark design requirement

Measure the momentum of 1 TeV muon to no worse than 10% precision

you can say they are "dual"

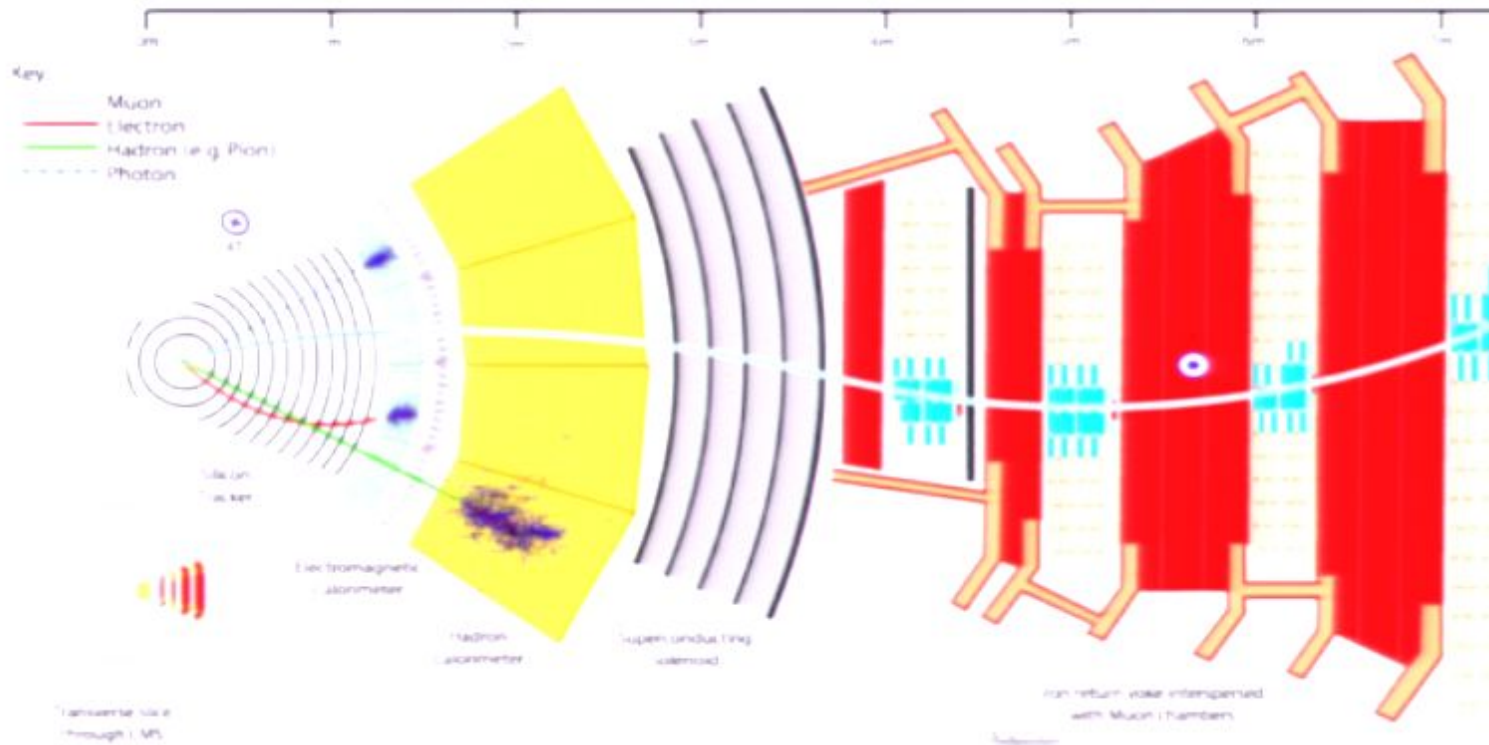
Clue1: Benchmark design requirement

Measure the momentum of 1 TeV muon to no worse than 10% precision

Clue2: Benchmark design requirement has at least 2 solutions!

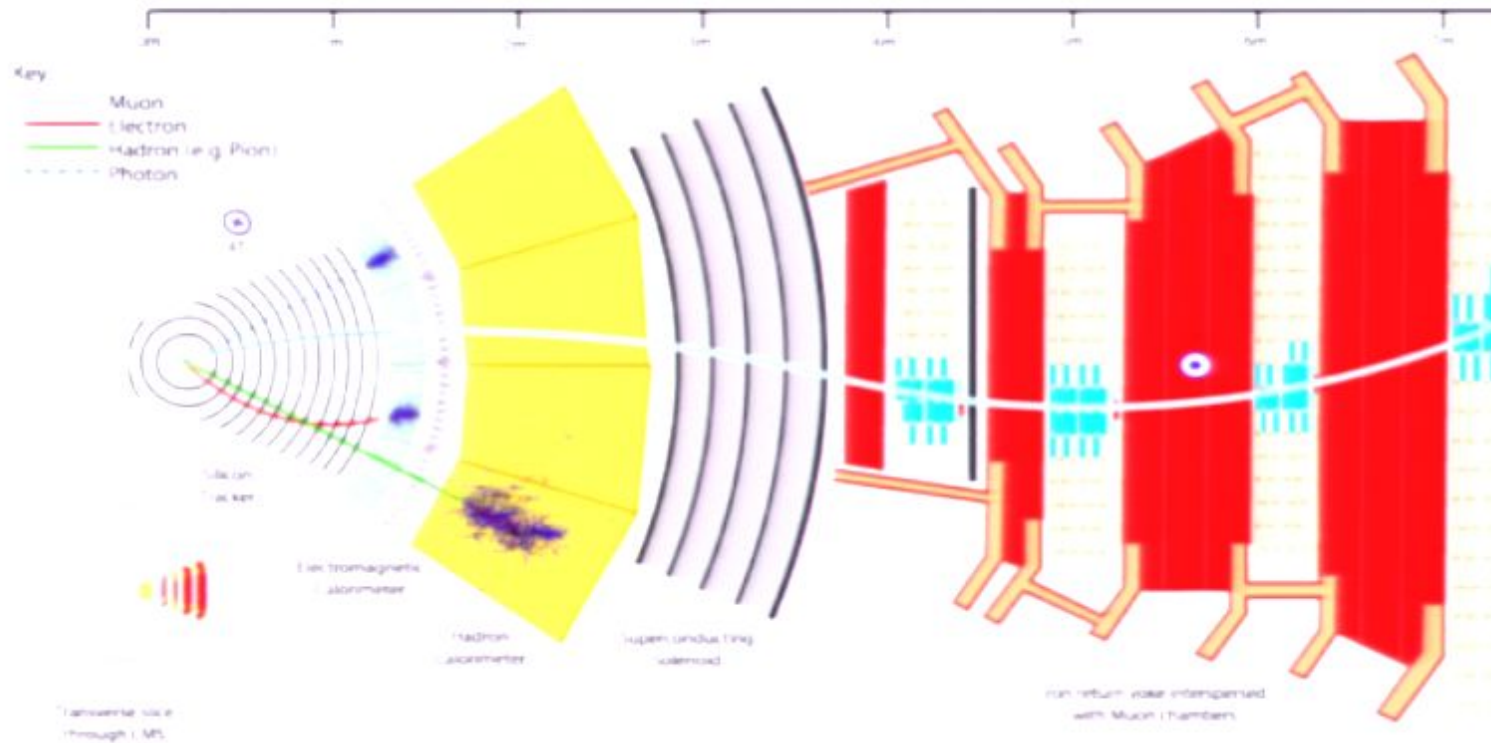
Homework

Because $\frac{\Delta p}{p} \sim \frac{1}{BL^2}$ where B is the magnetic field strength and L is the distance over which the bending of the muon takes place.



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Physics objects

- jets (calorimetry/tracking)
- missing E_T (calorimetry)
- electrons/photons (calorimetry/tracking)
- muons (tracking, muon chambers)
- hadronic taus (narrow jets)
- b's (jets with secondary vertex)
- top, W , Z (as reconstruction/calibration examples)



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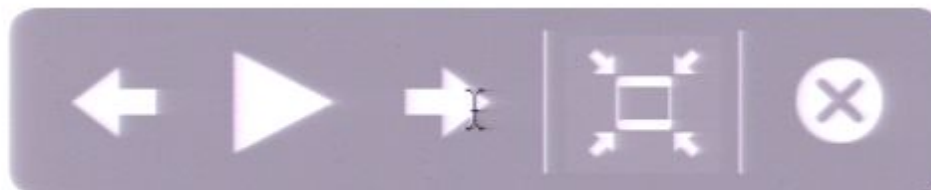


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Physics Objects @14 TeV

- jets from a highly boosted top are on top of each-other
- 44% of the energy of a $P_T=35$ GeV electron in the barrel is lost in brems (algorithms to recover ok)
- TeV muons shower and brem too
- material before calorimetry in ATLAS and CMS ~dark matter of the universe (interactions with matter complicate reconstruction)



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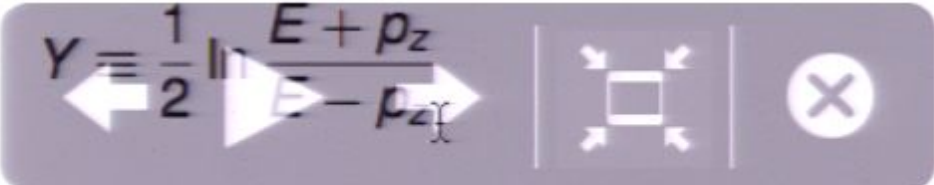


Coordinates and solid angle segmentation (η, ϕ)

- The origin $x = y = z = 0$ is at the nominal collision point in the geometrical center of the detector
- The z direction is the direction of the proton beam; the detector solid angle segmentation is designed to be invariant under boost along the z -direction. ϕ is the azimuthal angle about the z -axis η (*pseudorapidity*) is related to the polar angle θ by the relation:

$$\eta \equiv -\ln \tan \frac{\theta}{2} \quad (1)$$

- The pseudorapidity, η , is equivalent to the *rapidity*, of a particle in the limit of $p \gg m$, where p is the momentum of the particle and m its mass. The rapidity Y , is defined by

$$Y \equiv \frac{1}{2} \ln \frac{E + p_z}{E - p_z} \quad (2)$$




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If Y is expanded in terms of η and $\alpha = m/P_T$)

$$Y = \frac{1}{2} \ln\left(\frac{\sqrt{\cosh^2 \eta + \alpha^2} + \sinh \eta}{\sqrt{\cosh^2 \eta + \alpha^2} - \sinh \eta}\right)$$

$$\sim \eta - .5\alpha^2 \tanh \eta + O(\alpha^3) \eta(3)$$

- The rapidity Y of a particle transform linearly under a boost, $Y \rightarrow Y + \text{constant}$. Taking the derivative shows that segmentation in rapidity is Lorentz invariant, $dY \rightarrow dY$. The rapidity is a function of the particle's mass and polar angle and is not used to define the angular segmentation. Pseudorapidity depends only on the polar angle and is approximately Lorentz invariant under z -boosts for high P_T particles. It is used to define angular segmentation. Being transverse to the z -direction the ϕ is also invariant under z -boosts and is the orthogonal solid angle variable.



Jets

- Hadronic jets: final state of outgoing partons (subject to fragmentation) in a hard scattering
- Jet reconstruction algorithms: Clumping "clusters" of final state particles (towers, hadrons, tracks) into "jets" so that their kinematic properties (momenta) are related to the original parton
- Jet measurements: tests of QCD, mass reconstruction, search for new physics: (test of QCD predictions [PDFs x-section 'jet constituent'] is deposited by a massless particle striking the tower center and by summing the 4-momenta of all the towers(i) that by clustering construct the jet: .

$$E = \sum_i (E_i^{em} + E_i^{had})$$

$$p_x = \sum_i (E_i^{em} \sin \theta_i^{em} + E_i^{had} \sin \theta_i^{had}) \cos \phi_i$$



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Jets: Assignment

Show that jets are approximately circular in $\eta\phi$



Jets Algorithms: Requirements

- Detector Independence: There should be no dependence on cell type, numbers, or size.
- Theoretically Well Behaved: IR/Collinear Safe, no ad-hoc clustering parameters. [IR safety: soft radiation between two jets may cause a merging when above seed threshold. Collinear safety: jet not reconstructed if the seed energy is below threshold being split among several detector towers; number of jets depending on seed E_T ordering]
- Order Independence: Equivalence at the parton, particle, and detector levels.
- Fully Specified: Energy, angles, and underlying event should be clearly and completely defined. If necessary, pre-clustering, merging, and splitting algorithms must be completely described.
- Boost Invariant.



missing transverse energy (\cancel{E}_T)

Neutrinos do not interact in the detector. From transverse momentum conservation, the presence of undetected neutrinos results in transverse energy imbalance in the detector which is proportional to the neutrino momentum and it is called missing E_T or \cancel{E}_T or MET or P_T^{miss} . [Since the longitudinal component of the colliding partons momentum is unknown, only the transverse component of the neutrino momentum is measured.]

The raw \cancel{E}_T , is defined as the negative vector sum of the transverse energy in the calorimeter: At the online level, $\vec{\cancel{E}}_T$ is the vector sum of all calorimeter towers, both electromagnetic and hadronic, and is defined as

$$\vec{\cancel{E}}_T = \sum_{towers} (E_i \sin \theta_i) \hat{n}_i \quad (5)$$

where E_i is the energy of the i th tower, and \hat{n}_i is a transverse unit vector pointing to the center of each tower. θ_i is the polar angle of the line pointing from $z = 0$ to the i th tower. The sum extends to $|\eta_d| < 5$



Outline

The LHC Physics Besides the Standard Model

SUSY@LHC program of work

ATLAS+CMS SUSY search Highlights



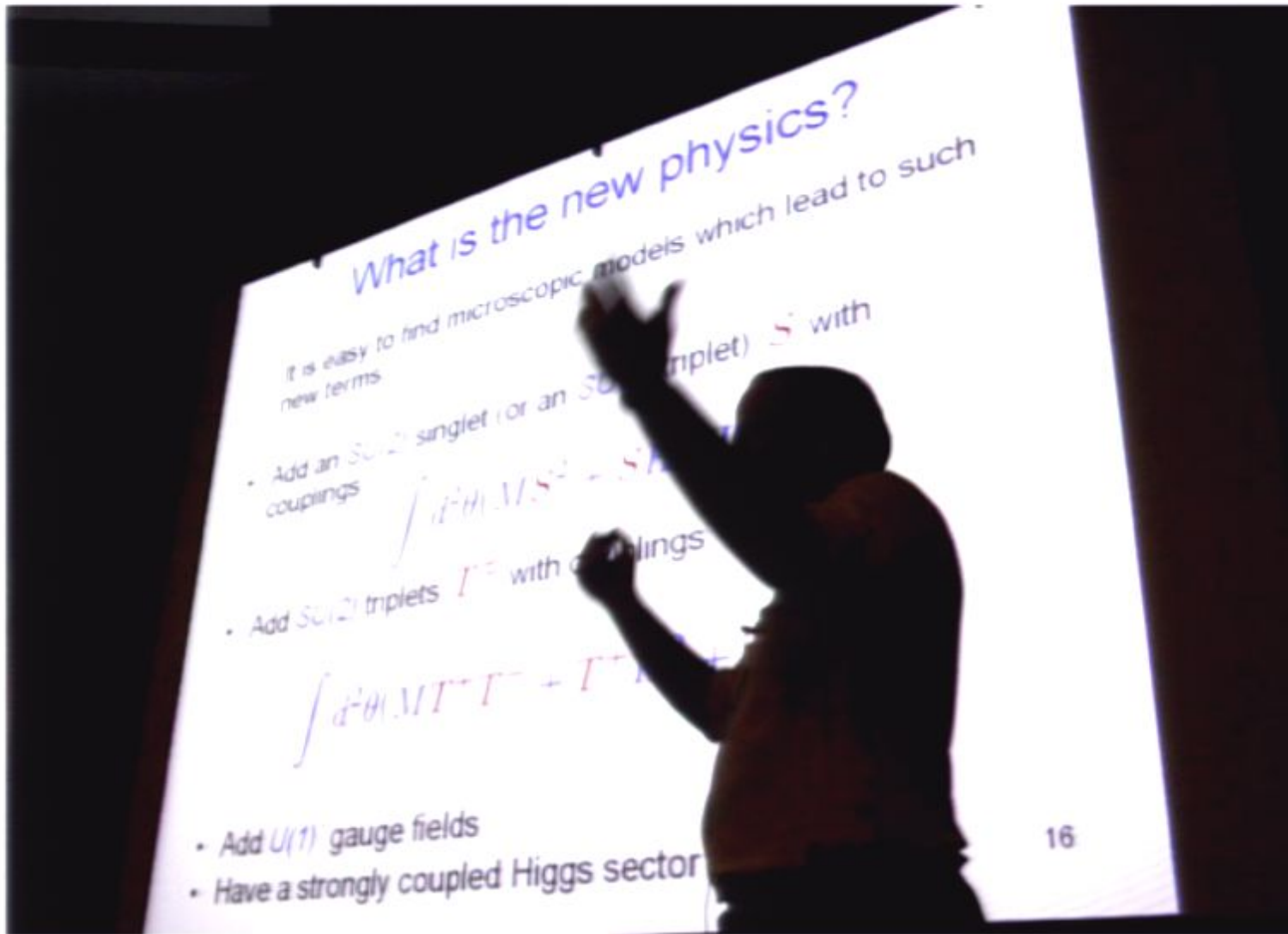
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Shelly Glashow, Stockholm 1979



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Strings 07. (Nati assumes) The LHC will discover some of the particles in the MSSM. These include some or all of the 5 massive Higgs particles of the MSSM. No particle outside the MSSM will be discovered (and concludes) There could be measurable deviations from MSSM relations at the LHC. These could point to new higher energy physics.



Top 2 Experimental Manifestations of New Physical Mechanisms (the known unknowns)

- the W and Z boson measured masses
- the dark matter

The expectation is then that the LHC will discover a new sector of particles associated with electroweak symmetry breaking and dark matter: →

Look for SUSY or "SUSY-mutation" in multijets+missing energy final state and 1xx GeV WIMP/neutralino dark matter candidate(s).

Top 1 Current Experimental Hint for Low Mass SUSY

$m_t = 171.4$ GeV and going down (currently 170.9)



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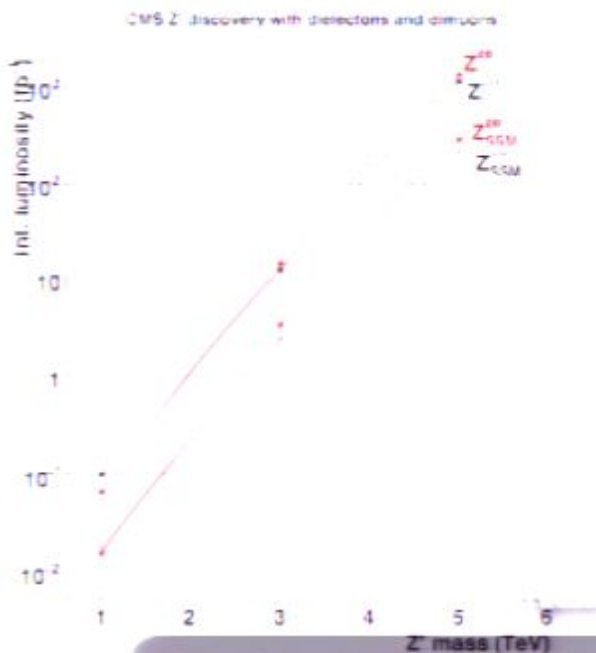
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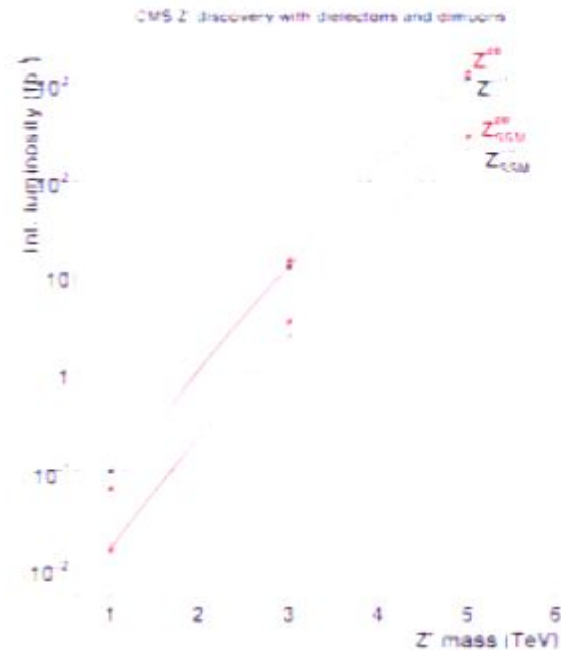
I will skip new signals from strings

i.e. **Z-primes** because they are easy to be discovered fast if they are 1 TeV heavy



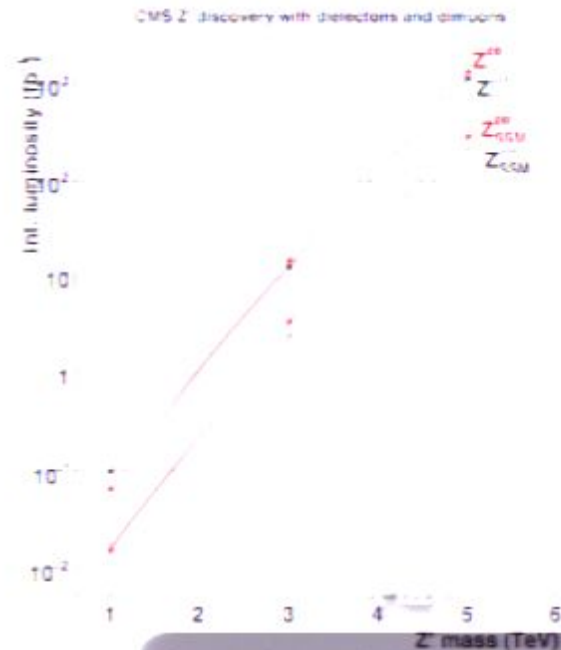
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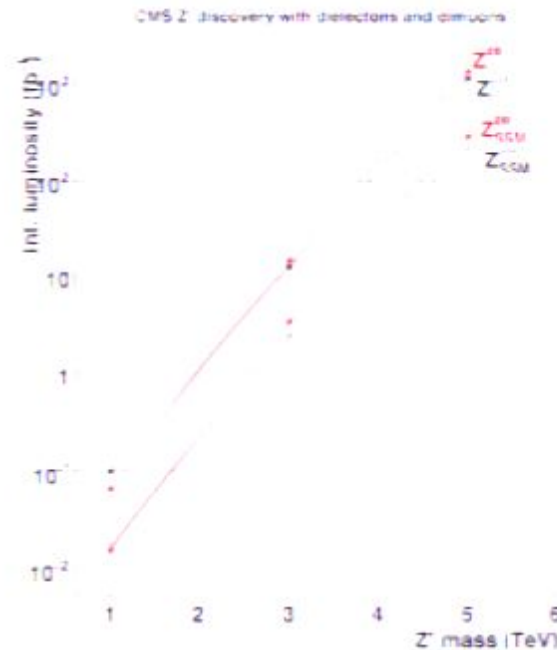
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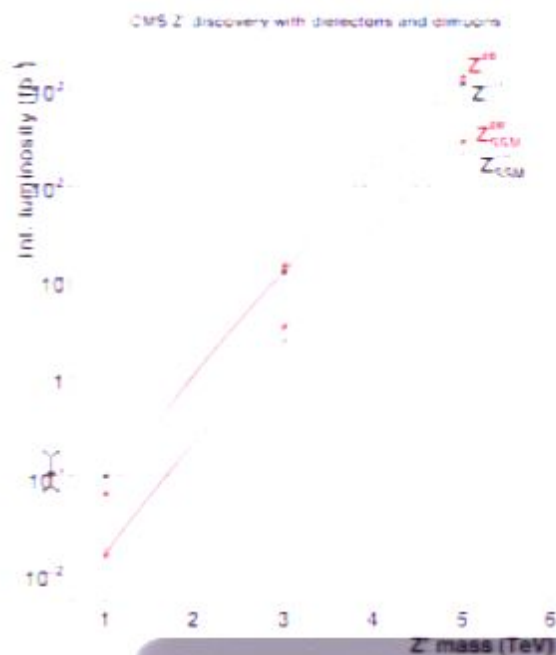
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ATLAS/CMS program of work

ATLAS CSC SUSY Notes

- SUSY1 : Data-driven Estimation of Z/W backgrounds to SUSY
- SUSY2 : Data-driven Estimation of top Backgrounds to SUSY
- SUSY3 : Data-driven Estimation of QCD Backgrounds to SUSY
- SUSY4 : Estimation of Heavy Flavor backgrounds and associated systematic
- SUSY5 : Searches and inclusive studies for SUSY events
- SUSY6 : Exclusive measurements for SUSY events
- SUSY7 : Gaugino direct productions
- SUSY8 : Studies for Gauge mediated SUSY

ATLAS CSC Exotics Notes

- Black Holes
- Dibosons
- Lepton+jets
- Dileptons
- Leptons+ Emiss



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ATLAS/CMS program of work

- **Leptonic searches (SUSY) [WG1]**
 - Search for SUSY in ≥ 1 lepton + E_T^{miss} + jets at 14 TeV in the electron and muon channels ($\mathcal{O}(100 \text{ pb}^{-1})$).
 - in dilepton pairs + E_T^{miss} + jets at 14 TeV in the electron and muon channels ($\mathcal{O}(100 \text{ pb}^{-1})$).
 - Search for SUSY in trileptons + jets at 14 TeV (1 fb^{-1}).
- **Hadronic searches (SUSY) [WG2]**
 - Search for SUSY in 0 lepton + E_T^{miss} + jets at 14 TeV ($\mathcal{O}(100 \text{ pb}^{-1})$).
 - in $b\bar{b}$ + E_T^{miss} + jets at 14 TeV ($\mathcal{O}(100 \text{ pb}^{-1})$).
 - in di-tau + E_T^{miss} at 14 TeV ($\mathcal{O}(100 \text{ pb}^{-1})$).
- **HSCP and photonic searches (GMSB) [WG3]**
 - Search and reconstruction of heavy stable charged particles at 14 TeV using TOF and dE/dx (500 pb^{-1} , model dependent).
 - Search for GMSB using prompt photons at 14 TeV (500 pb^{-1}).
- **High Energy Pair Searches (U(1)'/ED/other) [WG4]**
 - Search for TeV mass resonances in diEM events at 14 TeV (100 pb^{-1} $\mathcal{O}(\text{fb}^{-1})$)
 - in dimuon events at 14 TeV (100 pb^{-1})
 - in dijet events at 14 TeV ($\mathcal{O}(\text{fb}^{-1})$)



Alerts and Emphases in the Past Year

- Trigger aware analyses
- Data-driven background estimates, use of control samples
- Understanding of Standard Model QCD associated production of anything at $\sqrt{s} = 14$ TeV (especially W/Z/top, see talk by Michelangelo Mangano)
- Engineering of "first data" strategies – i.e. shifting gears to the pb^{-1} from the fb^{-1}
- Engineering of faster, simpler navigation/orientation compass in the vast parameter space



First Data and SUSY

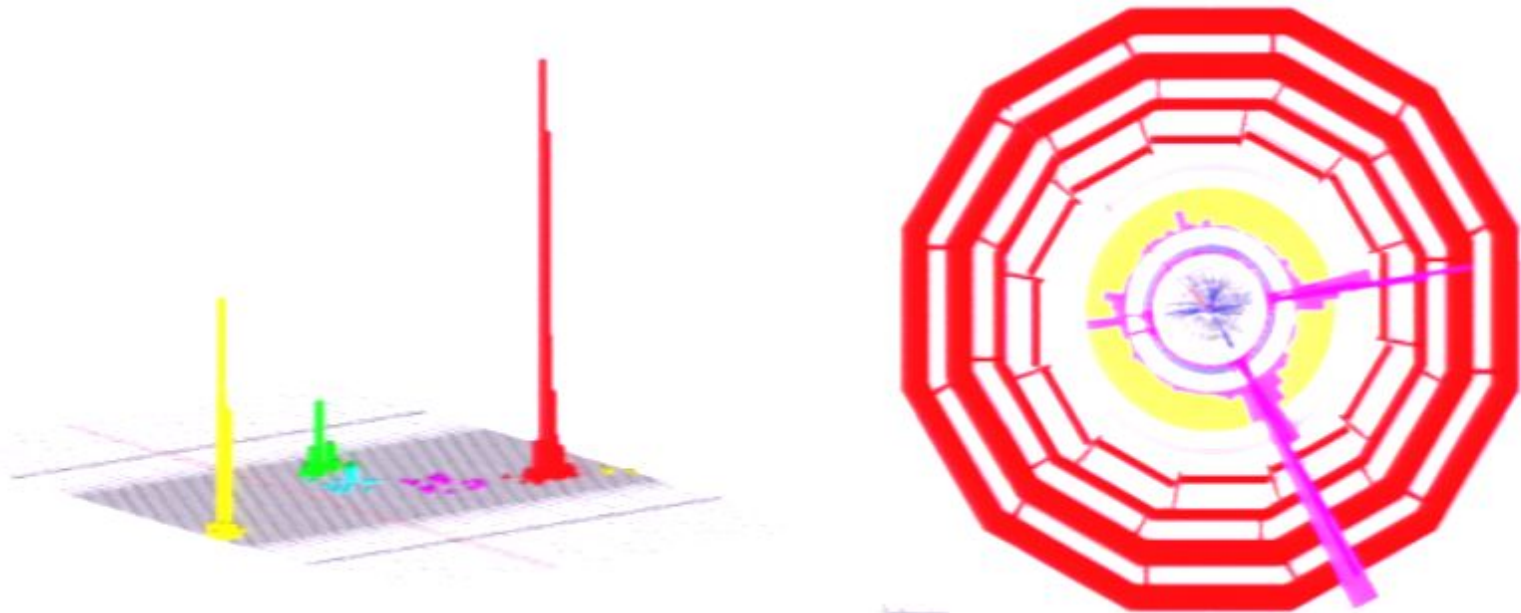
General Strategy (SUSY06, Giacomo P. /Maria S. for ATLAS/CMS)

- Choose signatures identifying well defined decay chains
- Extract constraints on masses, couplings, spin from decay kinematics/rates
- try to match emerging pattern to tentative template models
- having adjusted template models to measurements, try to find additional signatures to discriminate different options



Canonical Dark Matter Searches Using Missing Energy

$$E_T^{miss} = 360 \text{ GeV}, E_T(1) = 330 \text{ GeV}, E_T(2) = 140 \text{ GeV}, E_T(3) = 60 \text{ GeV}$$

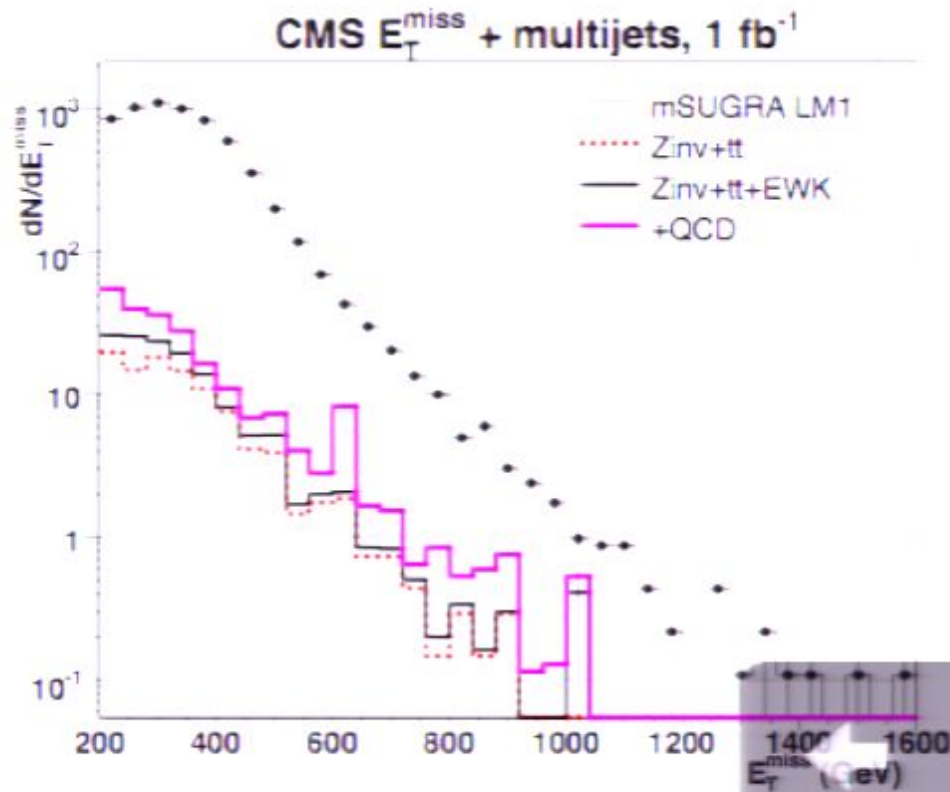


LHC can discover such events **fast**. Their cross section is **huge** :
10,000 to 1,000,000 events per year...



Inclusive Signatures with Missing Energy

E_T^{miss} +jets, gluino=600 GeV, neutralino=100 GeV



- fast-track to discovery of “low mass” SUSY $\mathcal{O}(10) \text{ pb}^{-1}$ b/c of signal cross section – control of systematics using SM processes (e.g. Z+jets, top)
- the time between $\mathcal{O}(10)$ and $\mathcal{O}(100) \text{ pb}^{-1}$ of well understood data will be critical for the discovery and characterization of SUSY



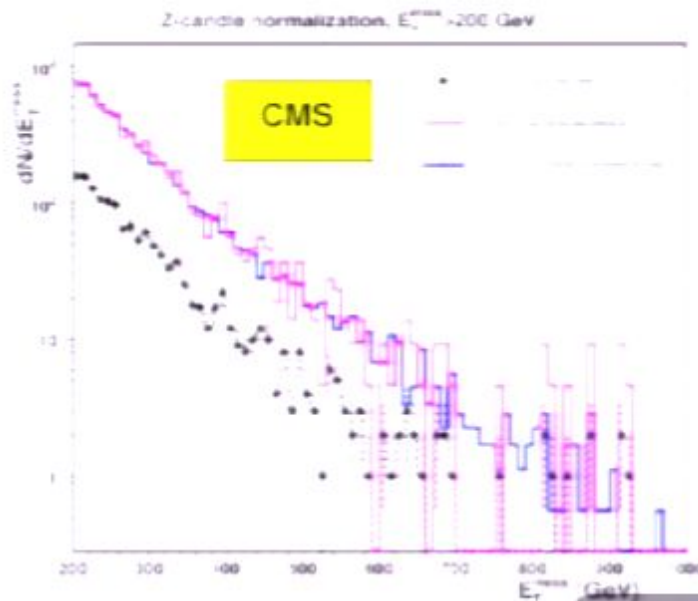
Standard Model Candles and Handles

Example

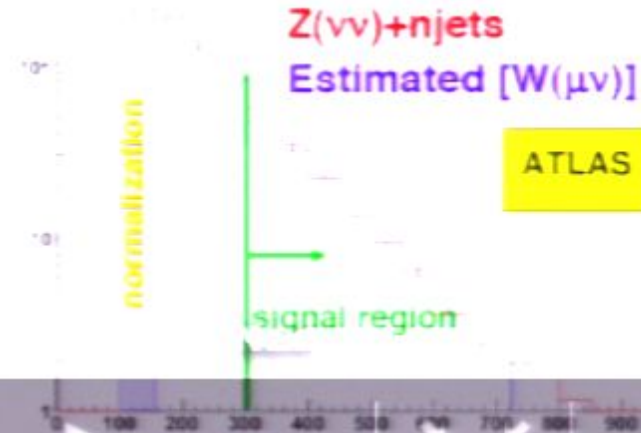
The Z candle



W/Z+jets Background



- MC to data normalization avoids systematics due to QCD scale, PDFs, ISR, FSR, jet energy scale ...



- CMS : Systematic uncertainty dominated by luminosity, measured ratio R and p ; 5% precision expected with ~ 1.5 fb⁻¹



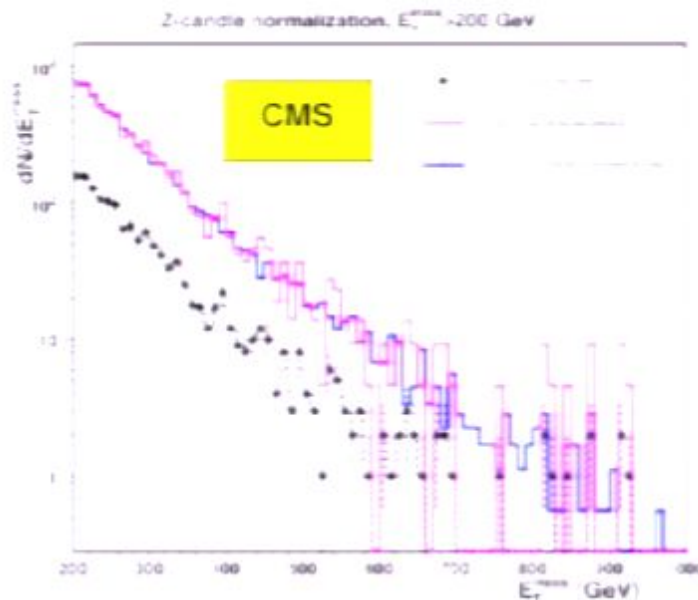
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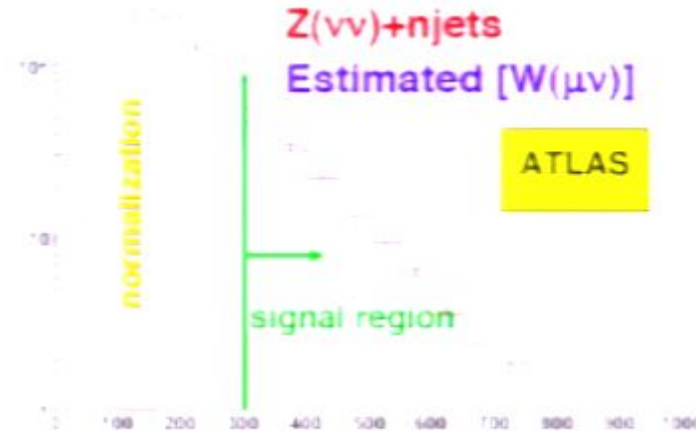
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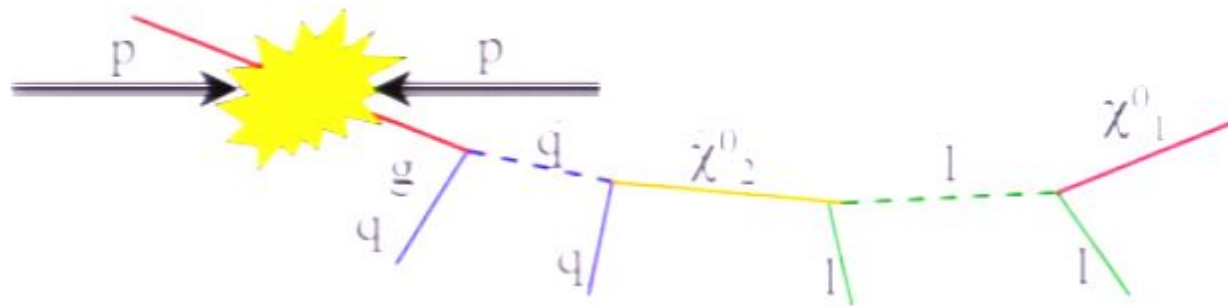
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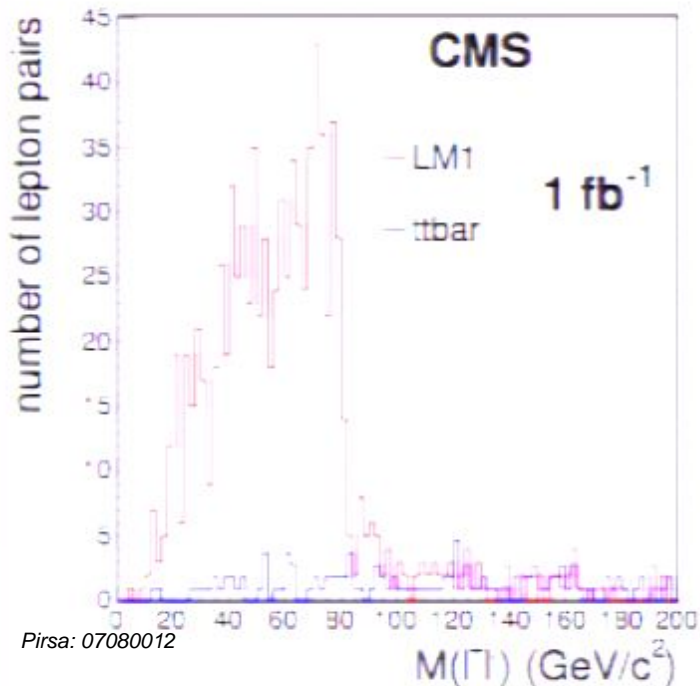
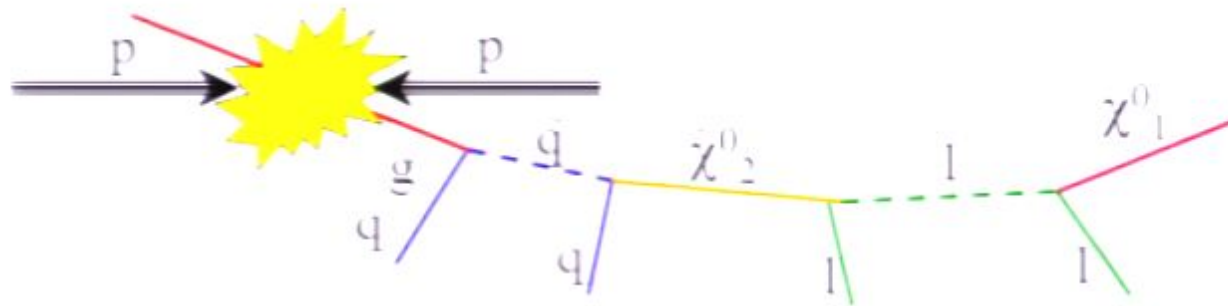
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First/Fast Mass Clues (dileptons)

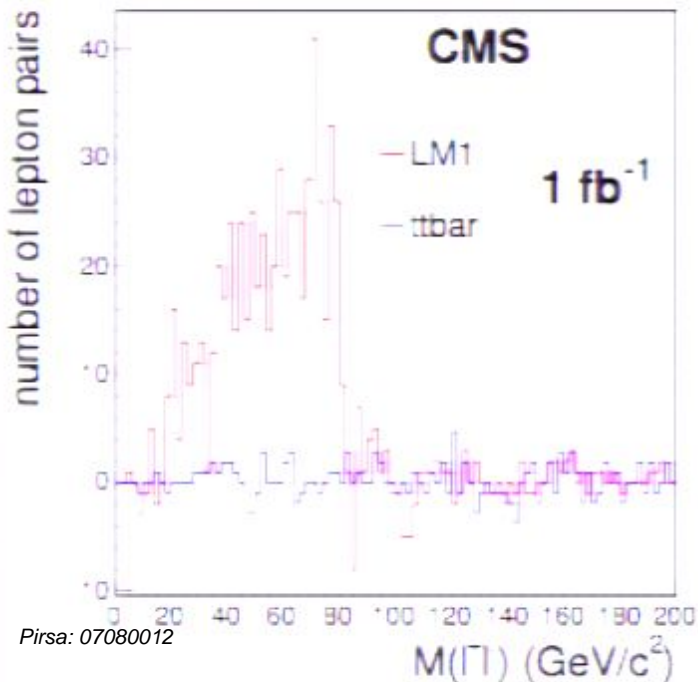
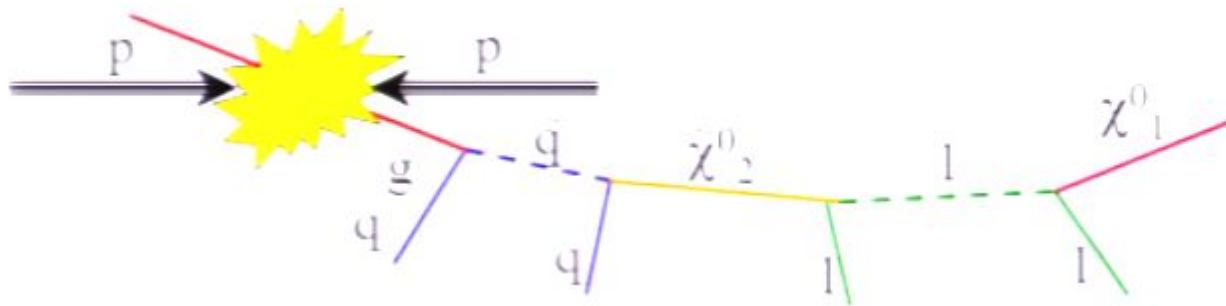


First/Fast Mass Clues (dileptons)



- SFOS dilepton+jets+ E_T^{miss}
- $t\bar{t}$: $WW+j$: $Z+j$: other $\sim 6:1:1:1$
- flavor subtraction ($e^- \mu^+ + e^- \mu^-$) to suppress chargino, W , $t\bar{t}$, WW , "other"
- L1+HLT trigger path required
- overall systematic on the background 20% (JES dominated)
- 5σ discovery with $\sim 20 \text{ pb}^{-1}$ (of data understood as expected with 1 fb^{-1}).

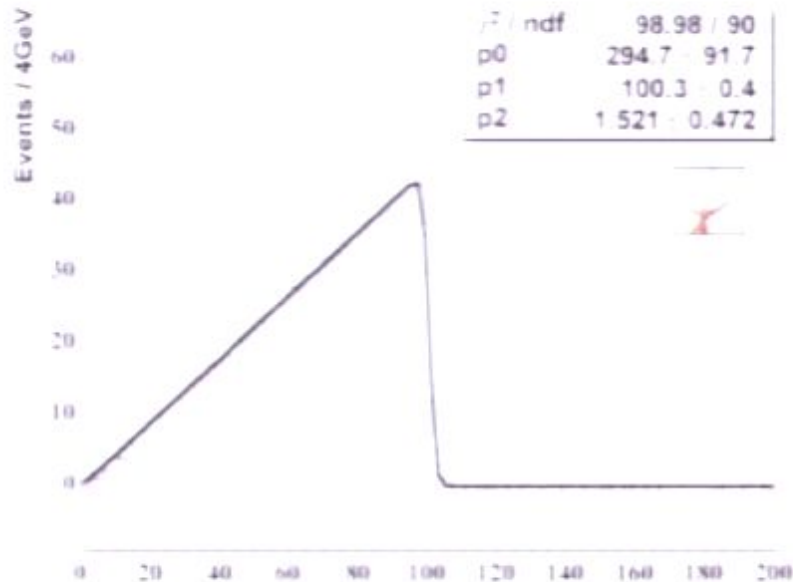
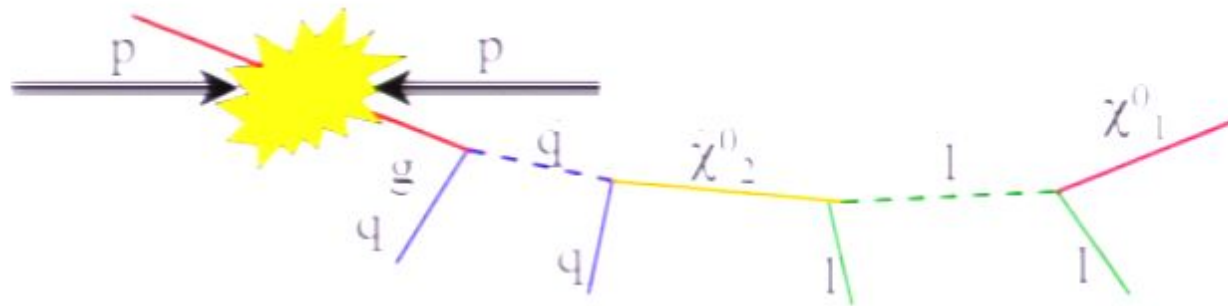
First/Fast Mass Clues (dileptons)



- $M_{\ell\ell}^{\max} = M(\tilde{\chi}_2^0) \sqrt{1 - \frac{M^2(\tilde{\ell}_R)}{M^2(\tilde{\chi}_2^0)}} \sqrt{1 - \frac{M^2(\tilde{\chi}_1^0)}{M^2(\tilde{\ell}_R)}}$
- $M_{\ell\ell}^{\max}(\text{meas}) = 80.42 \pm 0.48 \text{ GeV}/c^2$, cfr with
- expected $M_{\ell\ell}^{\max} = 81 \text{ GeV}/c^2$ [given $M(\tilde{\chi}_1^0) = 95$, $M(\tilde{\chi}_2^0) = 180$ and $M(\tilde{\ell}_R) = 119 \text{ GeV}/c^2$]



First/Fast Mass Clues (dileptons)



- $m_0 = 100 \text{ GeV}$, $m_{1/2} = 300 \text{ GeV}$, $A_0 = -300 \text{ GeV}$, $\tan \beta = 6$ $\text{sgn}(\mu) = +1$ [ATLAS SU3 test-point]
- $M_{\ell\ell}^{\text{max}}(\text{meas}) = 100.3 \pm 0.4 \text{ GeV}/c^2$ with 4.20 fb^{-1} (Geant-4 based simulation, no systematics)
- ATLAS-preliminary

towards SUSY reconstruction

Lepton+Jet Endpoint



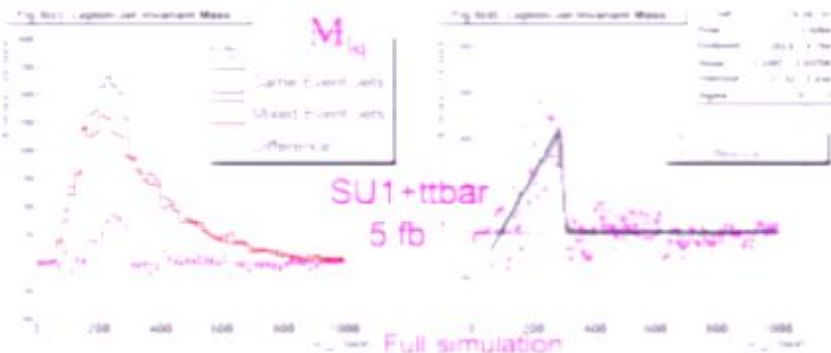
$$\bar{q}_L \rightarrow q\tilde{\chi}_1^0 \rightarrow ql'\tilde{\nu}_l \rightarrow ql'\nu_l\tilde{\chi}_1^0$$

$$M_{\nu_l}^{max} = \sqrt{\frac{(M_{\tilde{\nu}_l}^2 - M_{\tilde{q}}^2)(M_{\tilde{q}}^2 - M_{\tilde{\chi}_1^0}^2)}{M_{\tilde{\chi}_1^0}^2}}$$

Event selection:

- one lepton (electron / muon) $p_T > 20$ GeV $|\eta| < 2.5$ $chi2/ndof < 10$ for muon
- EMlikehood > 0.95 lepton isolation cut $E_{miss} < 10$ GeV in 0.45 cone
- Cone7 (R=0.7) jets, $\Delta R(\text{lepton}, \text{jet}) > 0.7$
- to reduce lepton+jets and dilepton ttbar background
 - leading and second leading jet with $E_T > 200$ GeV
 - transverse mass $M_T < 60$ GeV or $M_T < 100$ GeV
 - $E_{miss} > 250$ GeV

- Use **mixed event technique** to subtract combinatorial jet background (randomly pair jets from a different event (satisfying same event selection) with the lepton)
- Assumptions of the technique: The jet from signal decay chain and the jet from the decay of the other squark have **similar kinematic distributions**. And the squarks are produced at rest so the event is **roughly spherical**. Both are valid for SU1 but not for ttbar
- Subtract mixed-event-jet distributions from same-event-jets distribution to obtain an inferred "correct jet" distribution (normalization correction applied)
- Statistical significance of discovery is 5 σ with 5 fb
- Expected endpoint: 284 GeV
- Fitted endpoint: 283.6 \pm 4.8 GeV



Tuesday 17/2/2009

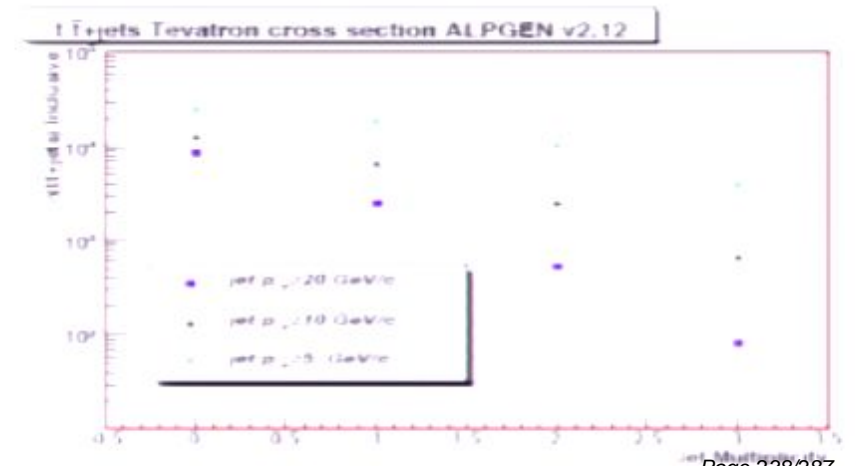
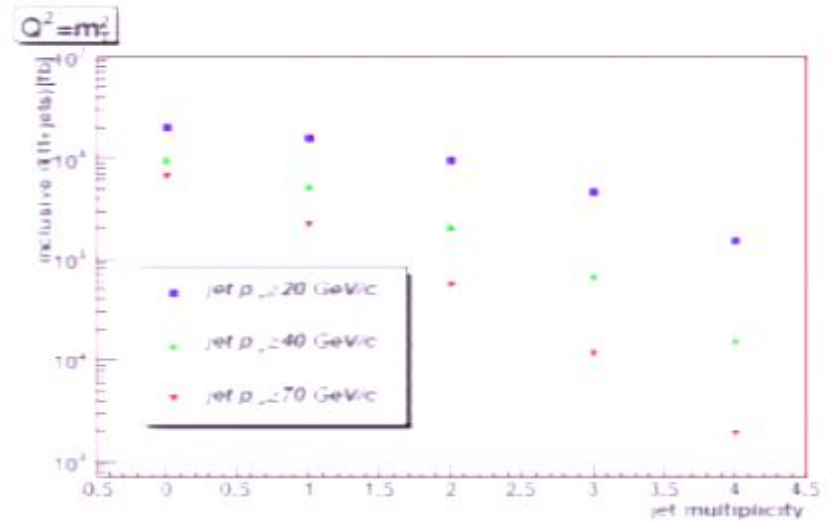
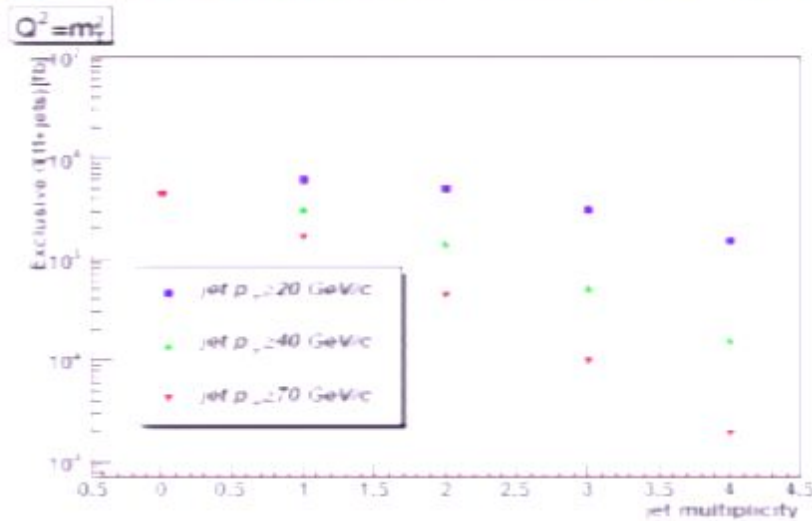
ATLAS note: Cooke et al ATLAS-COM-PHYS-2007-002 14th Feb 2007 11/20

Spin measurements, SFITTER, FITTINO results (usually require large luminosity)



Top QCD Matters @LHC: non trivial

MS, M.Pierini ALPGEN, tt+Njets... (see also MLM)



$$m^2 \psi^* \psi$$

$$MCFMn^2 \psi^* \chi \psi$$

ME PJ match \rightarrow

$$m^2 \psi^* \chi \psi$$



$$m^2 \psi^* \psi \leftarrow m^2 \psi^* \chi \psi$$

$$m \chi \chi - m \chi \sigma^{ij} F_{ij}$$

$$+ F^2 \chi_i \sigma^{ij} \chi_j$$

$$\bar{\chi}_i \sigma^{ij} \chi_j$$

$$+ \frac{F}{M} m^2 \psi^* \chi \psi - \frac{F}{M} m \chi \sigma^{ij} F_{ij}$$

$$m^2 \psi^* \psi + m c F m^2 \psi^* \chi \psi$$

δ $m^2 \psi^*$ ψ ME p, match $\sum X_i = \sum X_i^*$



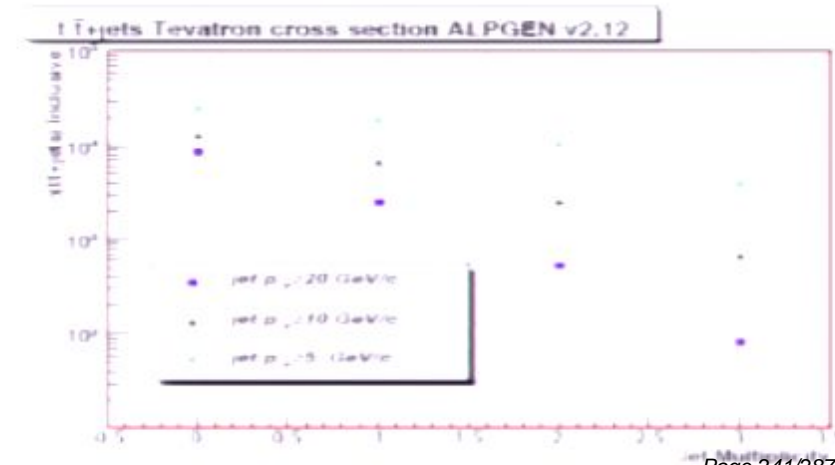
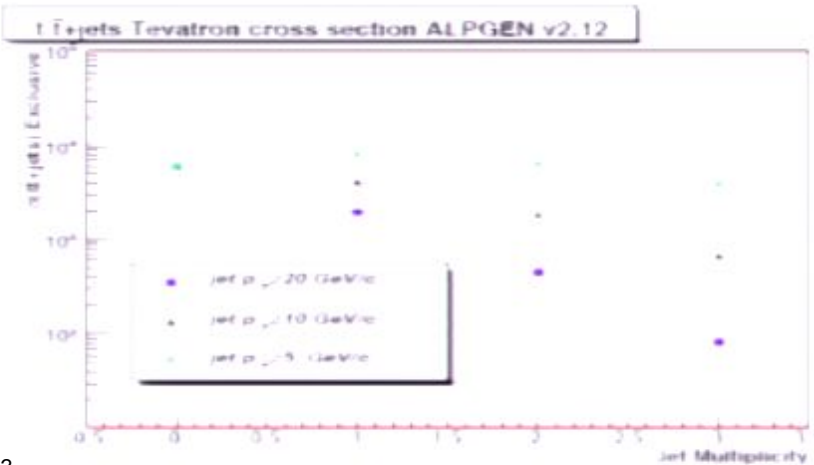
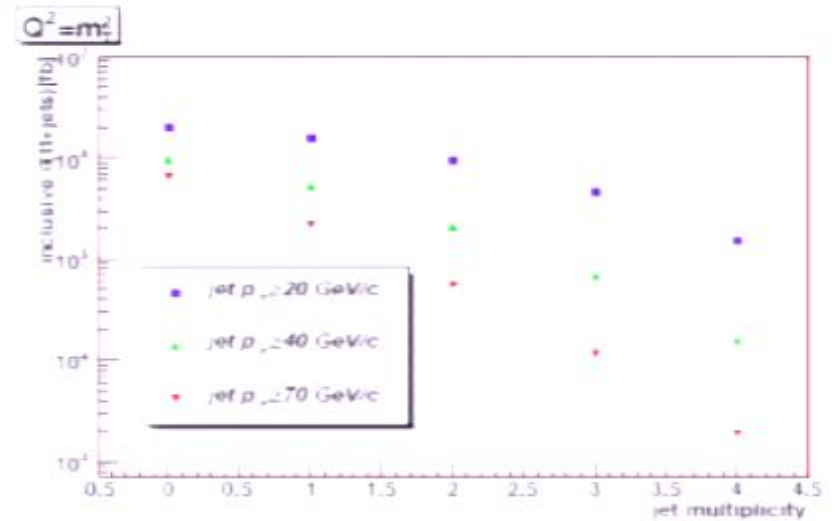
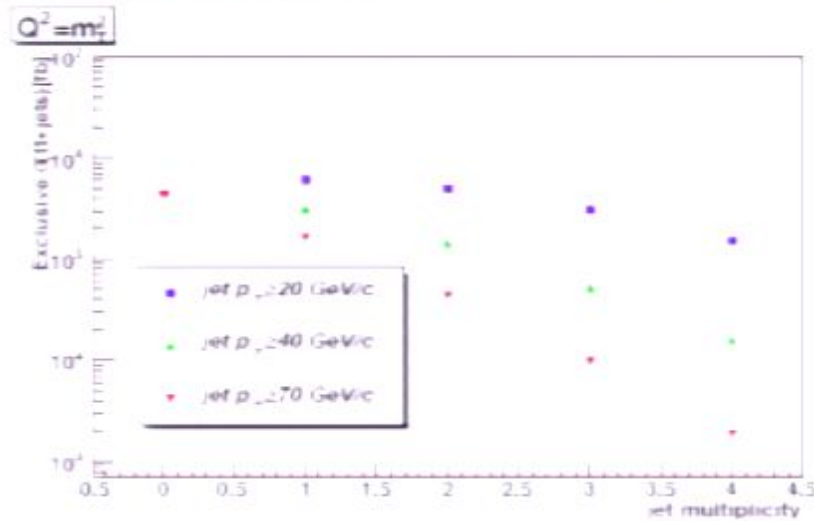
$$m^2 \psi^* \psi$$

$$m^2 \psi^* \psi + m^2 \psi^* \chi \psi + m \chi \chi - m \chi \psi m F_{ij}$$

$$+ F^4 \chi_i \sigma_i \chi + \chi_i \sigma_i \chi + m^2 \psi^* \psi + m \chi \chi - m \chi \psi m F_{ij}$$

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$$m^2 \varphi^* \varphi$$

$$MCFM m^2 \varphi^* \chi \chi$$

δ

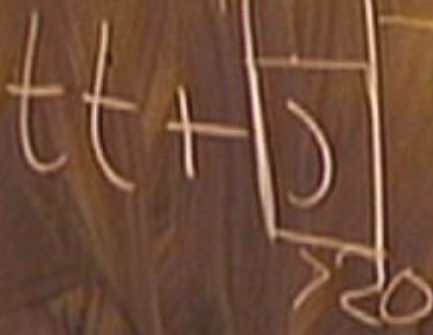
$$m^2 \varphi^*$$

χ

ME

ps match

$$m^2 \varphi^* \chi \chi$$



$$m^2 \varphi^* \varphi \leftarrow m^2 \varphi^* \chi \chi$$

$$+ m \chi \chi - m \chi \sigma \chi$$

$$+ F^4 \bar{\chi}_i \sigma \chi$$

$$\bar{\chi}_i \sigma \chi$$

$$+ \frac{m^2}{F} \varphi^* \chi \chi - m \chi \sigma \chi$$

$$m^2 \varphi^* \varphi$$

$$MCFM m^2 \varphi^* \chi \psi$$

δ

$$m^2 \varphi^* \chi$$

ME ps match



$$m^2 \varphi^* \chi \psi$$

$$m^2 \varphi^* \varphi \leftarrow m^2 \varphi^* \chi \psi$$

$$+ m \chi \chi - m \chi \psi$$

$$+ F^4 \bar{\chi}_i \sigma \partial \chi$$

$$\bar{\chi}_i \sigma \partial \chi$$

$$+ \frac{m^2}{F} \varphi^* \chi \psi - \frac{m}{F} \chi \psi$$

$$m^2 \varphi^* \varphi$$

$$MCFM m^2 \varphi^* \chi \chi$$

δ

$$m^2 \varphi^* \chi$$

ME ps match



$$m^2 \varphi^* \chi \chi$$

$$m^2 \varphi^* \varphi \leftarrow m^2 \varphi^* \chi \chi$$

$$+ m \chi \chi - m \chi \chi$$

$$+ F^4 \bar{\chi}_i \sigma \chi$$

$$\bar{\chi}_i \sigma \chi$$

$$+ \frac{m^2}{F} \varphi^* \chi \chi - \frac{m}{F} \chi \chi$$

$$m \varphi^* \varphi + M C E M m^2 \varphi^* \chi \psi$$



$$m^2 \varphi^* \varphi - m^2 \varphi^* \chi \psi$$

$$+ m \chi \chi - m \chi \sigma \chi$$

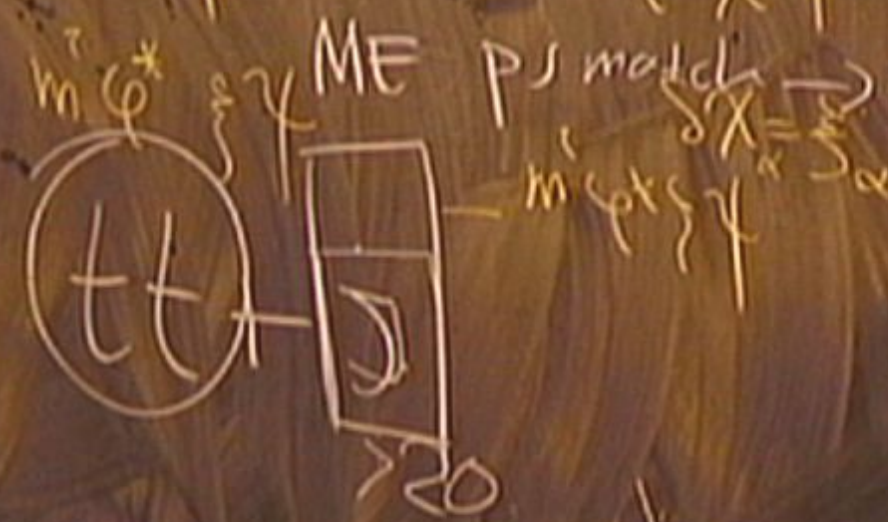
$$+ F^4 \bar{\chi}_i \sigma \chi$$

$$\bar{\chi}_i \sigma \chi$$

$$+ \frac{m^2}{F} \varphi^* \chi \psi - \frac{m}{F} \chi \sigma \chi$$

$$m \varphi^* \varphi$$

$$MCFM m^2 \varphi^* \chi \psi$$



$$m^2 \varphi^* \varphi - m^2 \varphi^* \chi \psi$$

$$+ m \chi \psi - m \chi \psi$$

$$+ F^4 - \bar{\chi} i \sigma \partial \chi$$

$$\bar{\chi} i \sigma \partial \chi$$

$$+ \frac{m^2}{F} \varphi^* \chi \psi - m \chi \psi$$

$$m^2 \varphi^* \varphi$$

$$MCFM m^2 \varphi^* \chi \chi$$

$$m^2 \varphi^* \chi \text{ ME } p, s \text{ match}$$

$$m^2 \varphi^* \chi \chi$$



$$m^2 \varphi^* \varphi - m^2 \varphi^* \chi \chi$$

$$+ m \chi \chi - m \chi \sigma \chi$$

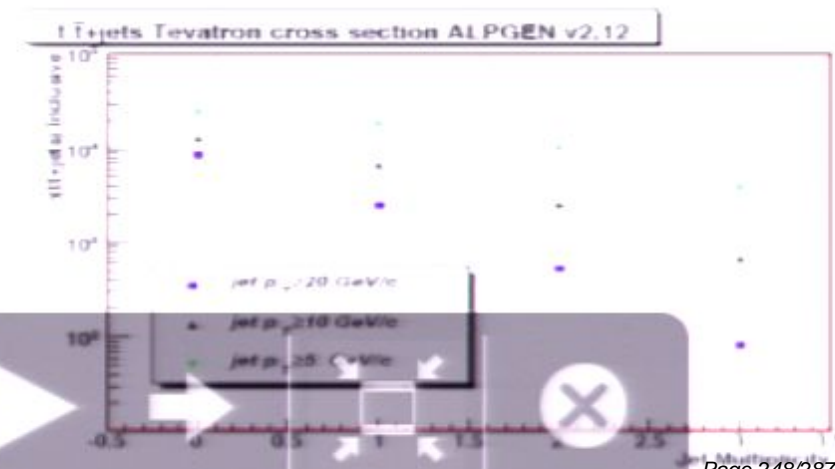
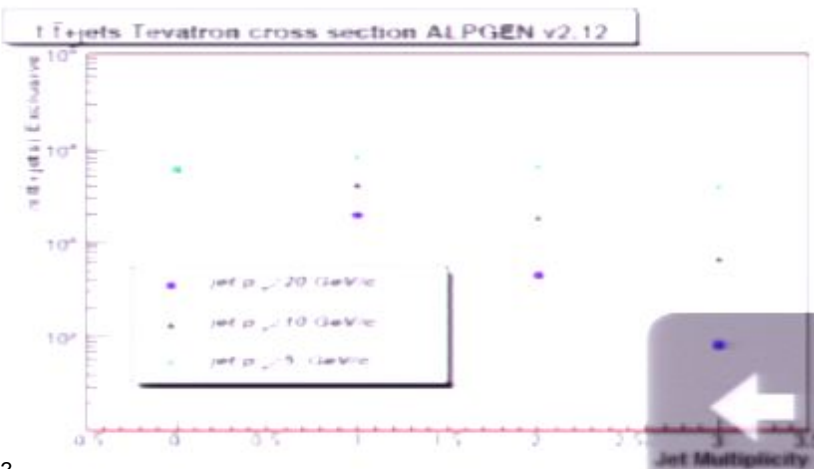
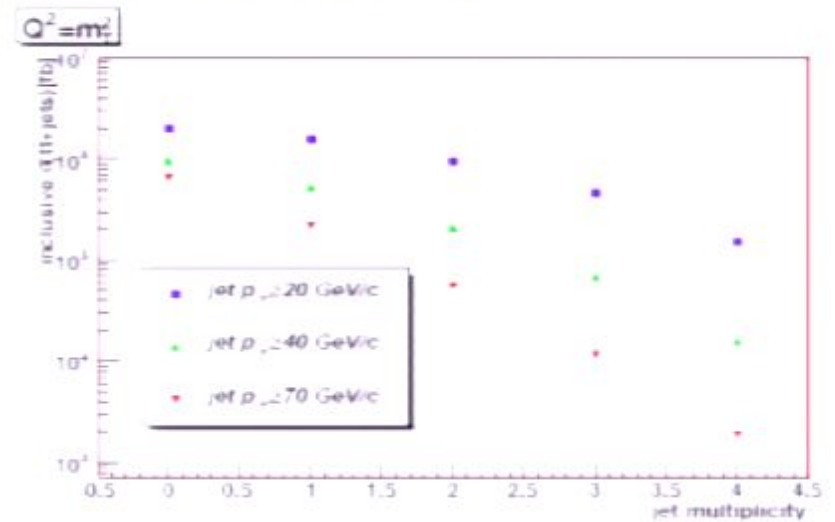
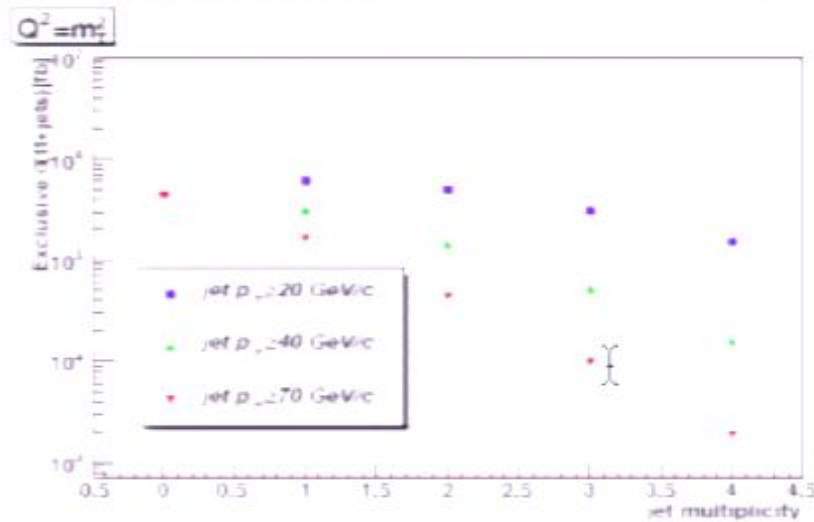
$$+ F^4 - \bar{\chi}_i \sigma \chi$$

$$\bar{\chi}_i \sigma \chi$$

$$+ \frac{m^2}{F} \varphi^* \chi \chi - \frac{m}{F} \chi \sigma \chi$$

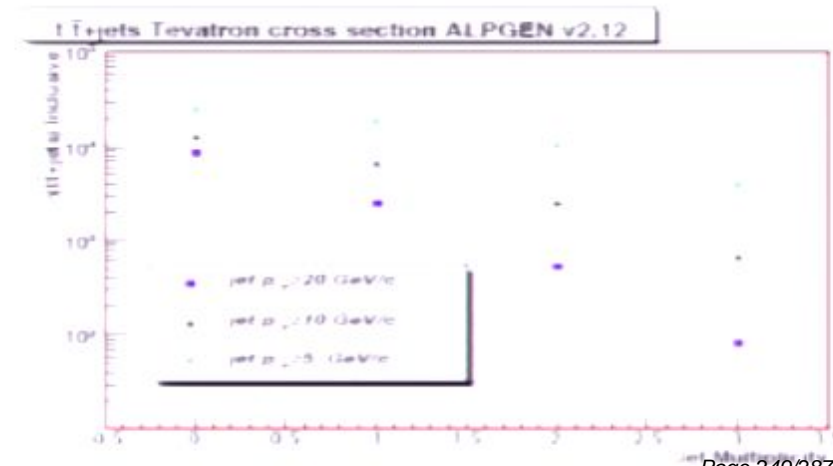
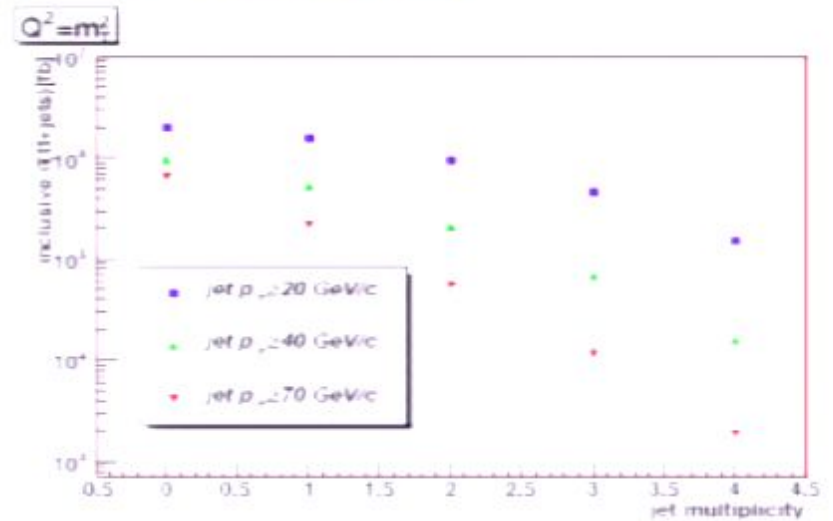
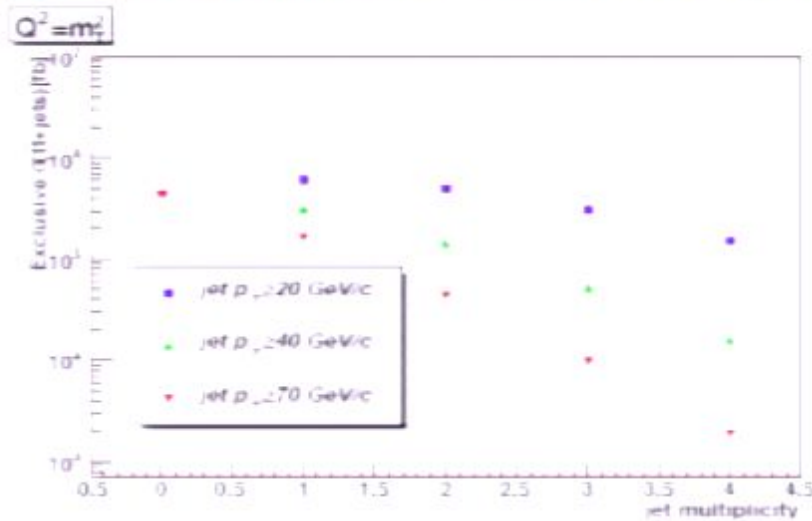
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bus

$$m \psi^* \psi$$

$$MCEM m^2 \psi^* \chi \psi$$

$$m \psi^* \chi$$

ME ps match



$$m \psi^* \chi \psi$$

$$m^2 \psi^* \psi \leftarrow m^2$$

$$+ m \chi \chi - m$$

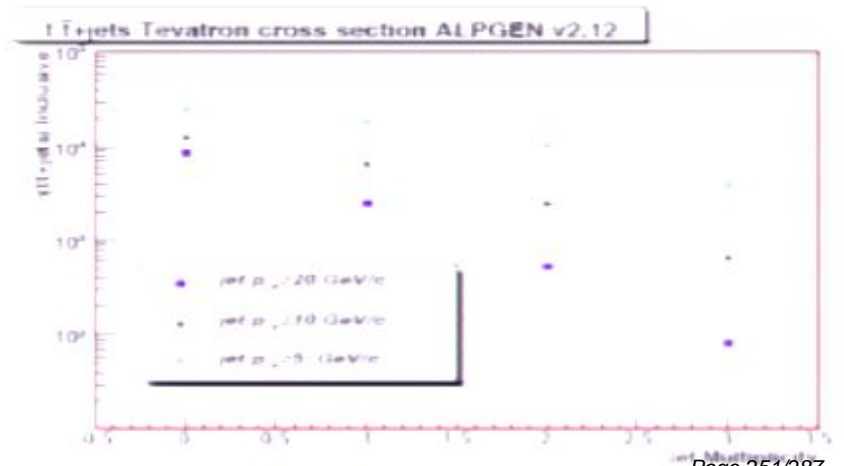
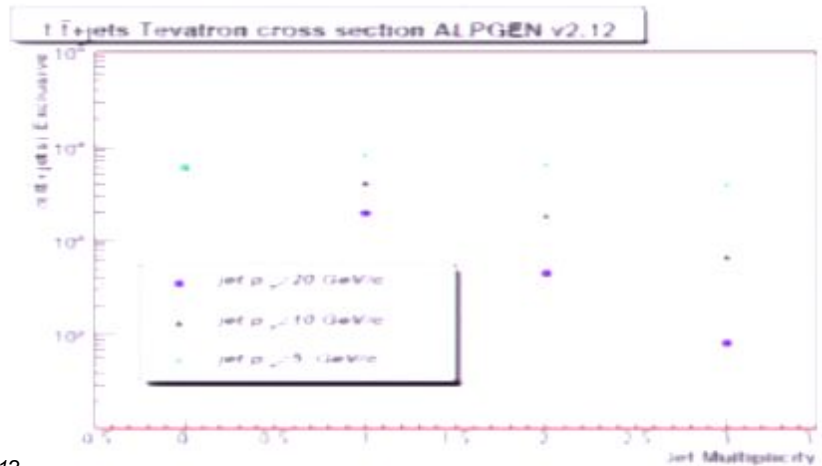
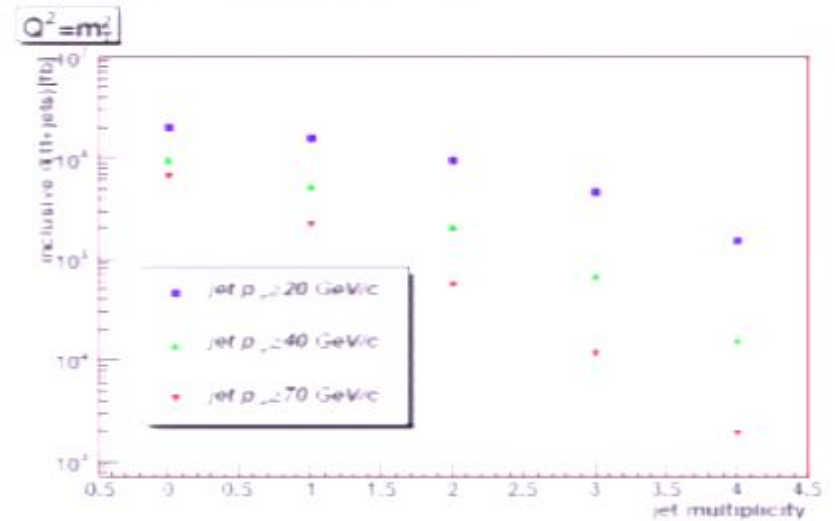
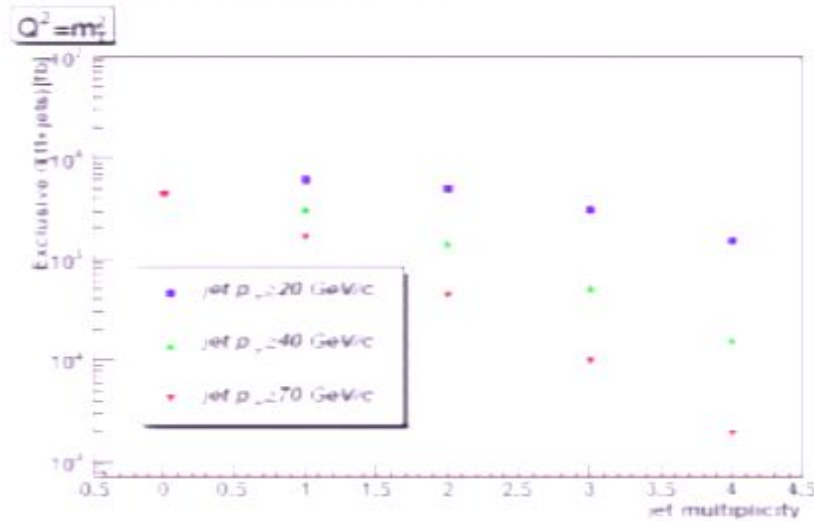
$$+ F^4 \chi_i$$

$$\bar{\chi}_i \psi \chi_i$$

$$+ \frac{F}{M^2}$$

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CMS SUSYBSM discovery plan

- Search for SUSY (Evidence for excess) in >1 lepton + E_T^{miss} + jets at 14 TeV in the electron and muon channels (100 pb^{-1}).
- Search for SUSY (Evidence for excess) in opposite sign dilepton pairs + E_T^{miss} + jets at 14 TeV in the electron and muon channels (20 pb^{-1})
- Search for SUSY (Evidence for excess) in same-sign dilepton pairs + E_T^{miss} + jets at 14 TeV in the electron and muon channels (200 pb^{-1})
- Search for SUSY (Evidence for excess) in Z^0 leptonic decays + E_T^{miss} + jets at 14 TeV in the electron and muon channels (100 pb^{-1})
- Search for LVF SUSY (Evidence for excess) in $e - \mu$ final state at 14 TeV (500 pb^{-1})
- Search for SUSY (Evidence for excess) in trileptons + jets at 14 TeV. ($\sim \text{fb}^{-1}$)
- Search for SUSY (Evidence for excess) in $b\bar{b}$ + 1 lepton at 14 TeV.
- Search for SUSY (Evidence for excess) in 0 lepton + E_T^{miss} + jets at 14 TeV (10 pb^{-1})
- Search for SUSY (Evidence for excess) in $b\bar{b}$ + E_T^{miss} + jets at 14 TeV (100 pb^{-1})
- Search for SUSY (Evidence for excess) in top hadronic decays + E_T^{miss} at 14 TeV (200 pb^{-1})
- Search for SUSY (Evidence for excess) in opposite-sign ditau + E_T^{miss} at 14 TeV (200 pb^{-1})
- Search for GMSB (Evidence for excess) in prompt photon final states at 14 TeV (500 pb^{-1})
- Search for GMSB (Evidence for excess) in non-pointing photons at 14 TeV (1 fb^{-1})
- Search and reconstruction of heavy stable charged particles at 14 TeV using TOF and dE/dx (500 pb^{-1})
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-



Some Discovery signatures

- canonical inclusive
 - jets+ E_T^{miss} (*) includes strategies for beam halo/noise, first data
 - jets+ ℓ + E_T^{miss}
 - same-sign dimuon + E_T^{miss}
 - opposite-sign same flavor dielectron and dimuon + E_T^{miss}
 - opposite-sign same flavor hadronic ditau + E_T^{miss}
- higher reco object inclusive
 - Z + E_T^{miss}
 - t hadronic + E_T^{miss}
 - $h^0(b\bar{b})$ + E_T^{miss}



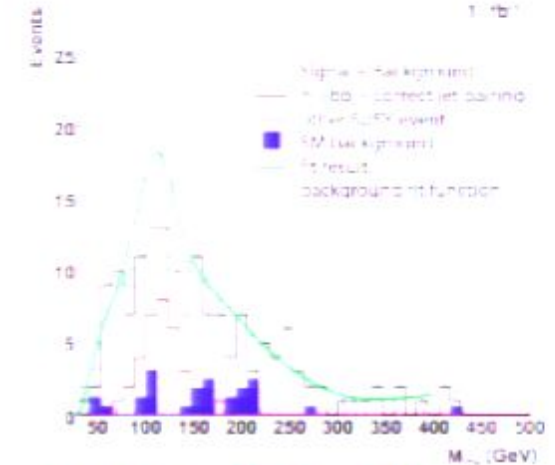
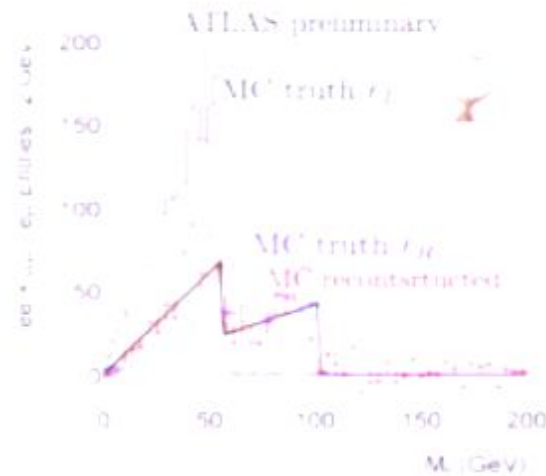
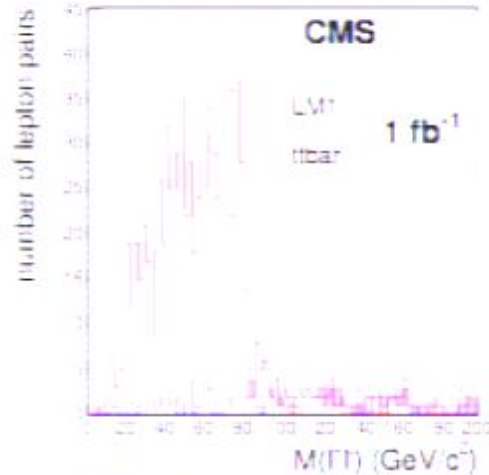
is it an amoeba or is it Sharon Stone? decoding the compositions of primary constituents



- Just like decoding DNA we have to decode the signals we will observe. And we do expect more similarities than differences, so fast discrimination will require smart and simple measurements.



Features directly from the kinematics



- But can we also tell the twins apart? Just as DNA is **“what”** we are but not **“who”** we are, we have to find additional traits and put the pieces together in order to orient and navigate in the vast parameter spaces of the models.



Towards Systematic and Deeper Mining of the Event

Examples of the simple questions to ask and answer (fast) and recent examples of works:

- excess of SS dileptons \rightarrow ?
- $++ / -- = 2 \rightarrow ?$
- excess if OS dileptons $\rightarrow ?$
- triangle in dilepton invariant mass $\rightarrow ?$
- double triangle $\rightarrow ?$
- no triangle $\rightarrow ?$
- Z^0 and no triangle $\rightarrow ?$
- Z^0 and triangle $\rightarrow ?$
- ratios: $1\ell/2\ell/3\ell/4\ell$, jets
- ... more "footprint" homeworks



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Arkani-Hamed, Toro, Shuster et al Very clever "marmoset" solutions to the navigation problem: see Olympics and black boxes challenges

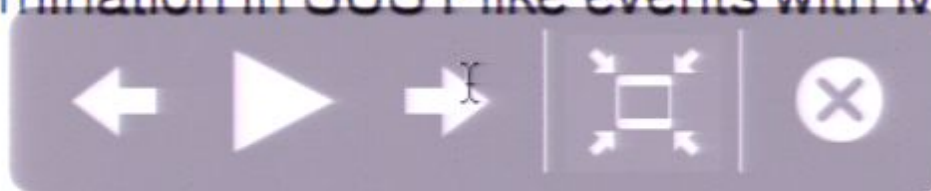


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Gunion et al. Mass Determination in SUSY-like events with Missing Energy



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- ... more "footprint" homeworks

A number of post-modern multivariate very sophisticated analyses with DTs, NNs, fancy statistics etc: with care and caveats



N.B. Prerequisite path of Discovery Work

- understand detectors and SM backgrounds
- control/understand: trigger, initial calibrations, scales, resolutions, efficiencies...
- minimize poorly estimated standard model backgrounds
- use SM "candle"/control samples (W/Z/top) to estimate backgrounds as possible
- use ratios as much as possible to get rid of luminosity dependence and other cross section related systematics
- adapt methods for background extraction as a function of luminosity
- have in place MC tools and reasonable "significance" calculations

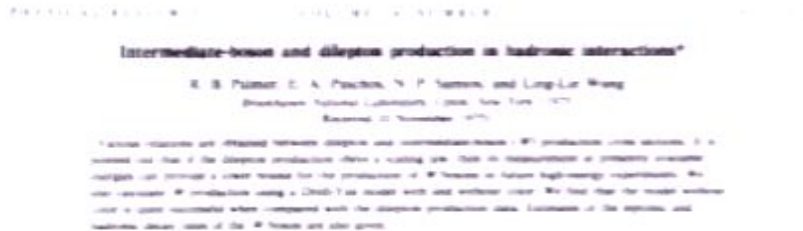


Outline

Wisdom from the Past



W-s and Z's: from 1976 to ALPGEN



Extrapolating from 1975

- BNL 400 GeV collidtr proposal
- s/M_W^2 Tevatron ~ 500
 $\rightarrow \sigma = 16$ nb
- s/M_W^2 LHC ~ 30000
 $\rightarrow \sigma = 12$ nb
- The Tevatron W-inclusive cross section is (measured) 8 nb
- The LHC is ~ 20 nb

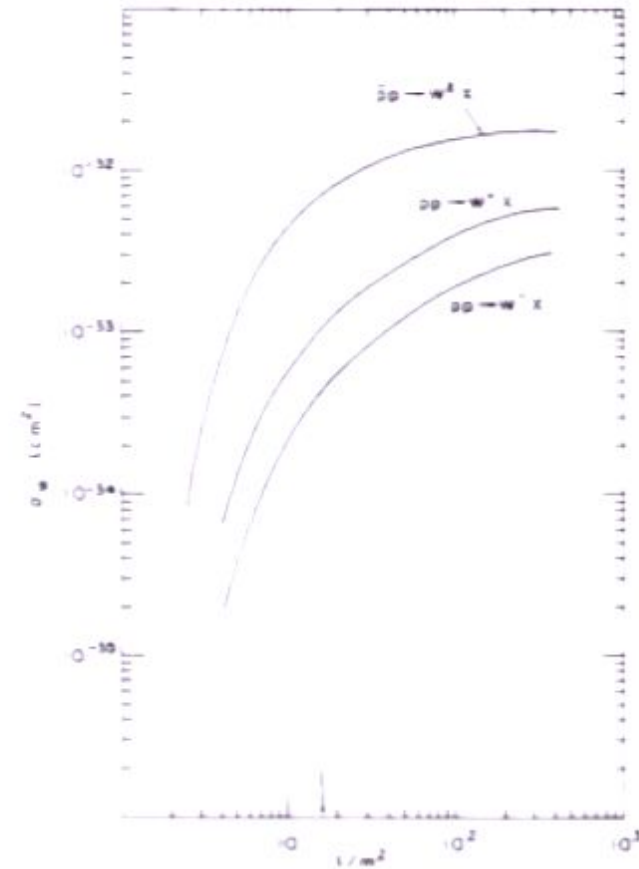
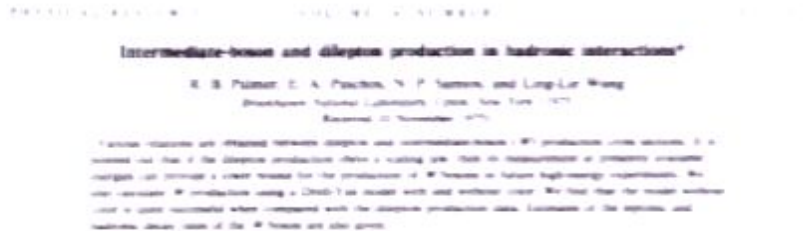


FIG. 6. Production cross section for $pp \rightarrow W^+ X$ and $\bar{p}p \rightarrow W^+ X$, calculated from the model given in Sec. III. The arrow points to the s/M_W^2 value corresponding to the production of a 100-GeV W meson by a 200-GeV/c-on-200-GeV storage machine.



W-s and Z's: from 1976 to ALPGEN



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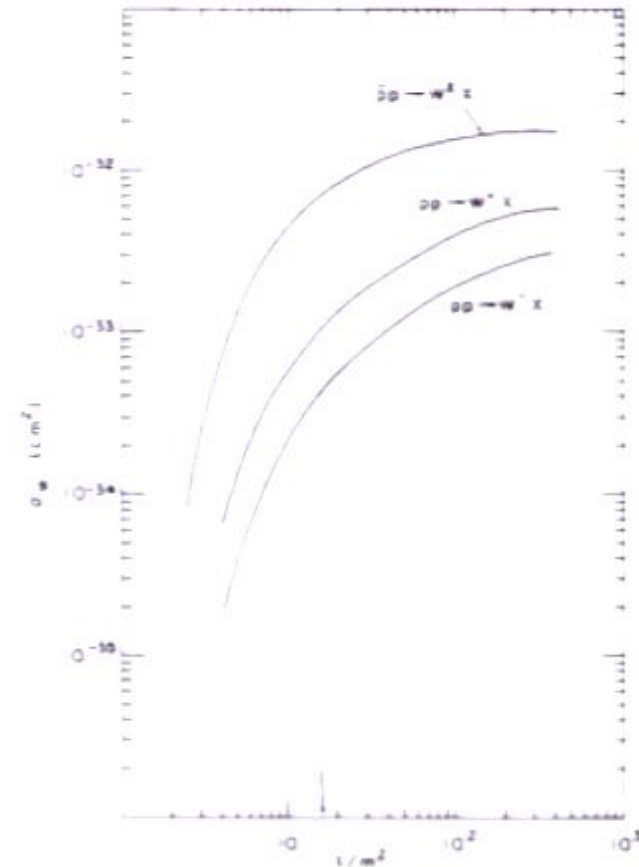
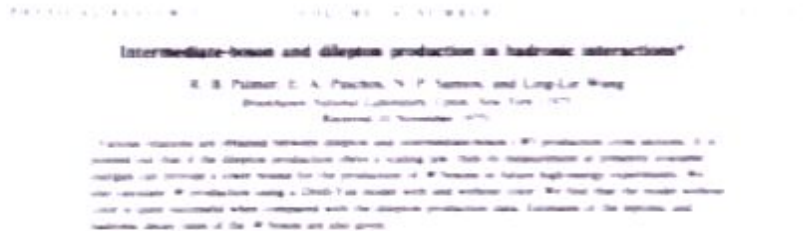


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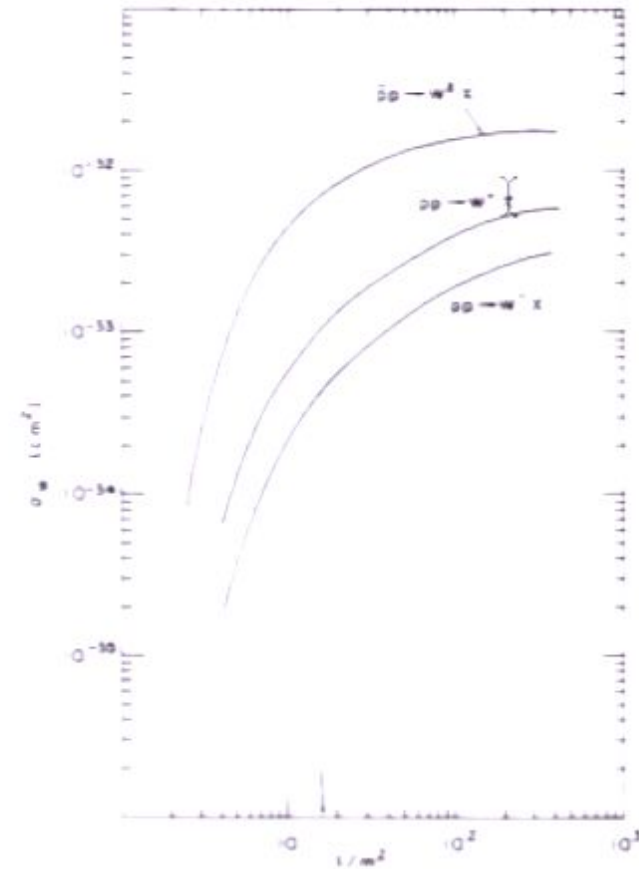
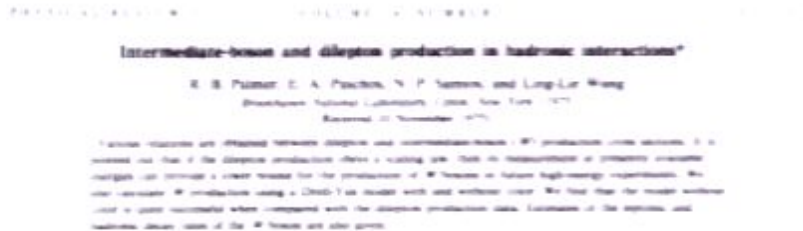


FIG. 6. Production cross section for $pp \rightarrow W^+ X$ and $\bar{p}p \rightarrow W^- X$, calculated from the model given in Sec. III. The arrow points to the s/M_W^2 value corresponding to the production of a 100-GeV W meson by a 200-GeV/c-on-200-GeV storage machine.



W-s and Z's: from 1976 to ALPGEN



Extrapolating from 1975

- BNL 400 GeV collidr proposal
- s/M_W^2 Tevatron ~ 500
 $\rightarrow \sigma = 16$ nb
- s/M_W^2 LHC ~ 30000
 $\rightarrow \sigma = 12$ nb
- The Tevatron W-inclusive cross section is (measured) 8 nb
- The LHC is ~ 20 nb

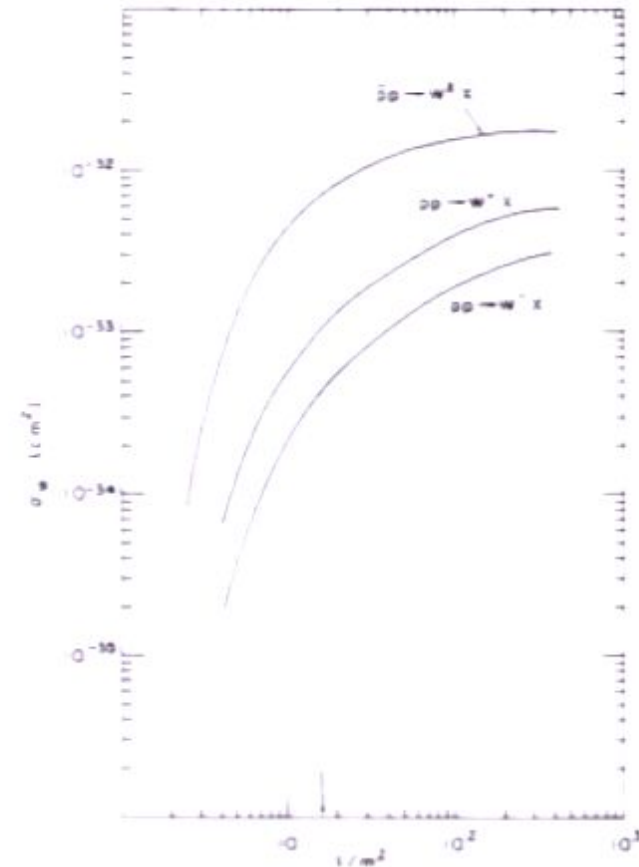


FIG. 6. Production cross section for $pp \rightarrow W^+ \lambda$ and $\bar{p}p \rightarrow W^+ \lambda$, calculated from the model given in Sec. III. The arrow points to the s/M^2 value corresponding to the production of a 100-GeV W meson by a 200-GeV/c-on-200-GeV/c storage machine.



Resonances, Reflections: a non-trivial discovery example from 30 years ago (MLM digged this)



$K\pi$ and $K3\pi$ spectra 2.08.76 PRL 37 (76) 255

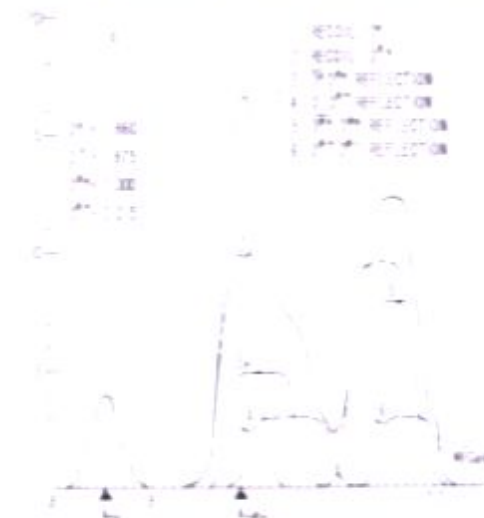
Experiment (hadron) $\sqrt{s} = 1.96$ MeV and $\sqrt{s} = 1.00$ GeV. Using our technique of fits we tabulate available energies \sqrt{s} and branching ratios B for J^P decays as follows:

$J^P = 0^{-+}$ at 1.1 MeV $B = 0.97$
 $J^P = 0^{-+}$ at 1.1 MeV $B = 0.79$
 $J^P = 0^{-+}$ at 1.90 MeV $B = 0.13$

$J^P = 0^{-+}$ at 1.0 MeV $B = 0.80$
 $J^P = 0^{-+}$ at 1.0 MeV $B = 0.71$
 $J^P = 0^{-+}$ at 1.40 MeV $B = 0.14$

Although none of the masses we use are precisely determined, the width Γ (near Γ_{total}) is roughly a neutral J^P meson (mass) Γ_{total} (width) ratio.

To the extent that $\Gamma_{\text{total}} \approx \Gamma_{\text{total}} \approx \Gamma_{\text{total}}$ (width) ratio, the result of neutral J^P is expected to be seven times greater than the total $\Gamma_{\text{total}} \approx \Gamma_{\text{total}} \approx \Gamma_{\text{total}}$ (width) ratio.



Resonances, Reflections: a non-trivial discovery example from 30 years ago (MLM digged this)



Experiment (radiative) $J/\psi \rightarrow \gamma e^+e^-$ (36 MeV) and $J/\psi \rightarrow \gamma \mu^+\mu^-$ (100 GeV) using our estimate of error we tabulate available energies \sqrt{s} and branching ratios B for J/ψ decays as follows:

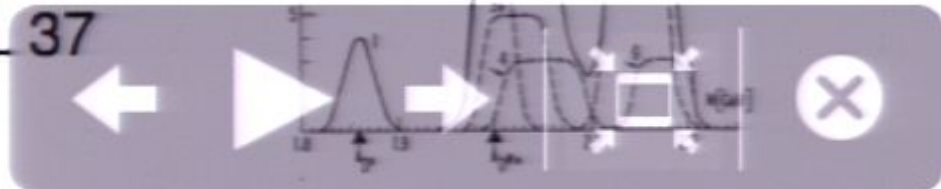
$J/\psi \rightarrow \gamma e^+e^-$ 31 MeV $B = 0.40\%$
 $J/\psi \rightarrow \gamma \mu^+\mu^-$ 100 MeV $B = 0.17\%$
 $J/\psi \rightarrow \gamma \mu^+\mu^-$ 140 MeV $B = 0.13\%$

$J/\psi \rightarrow \gamma \mu^+\mu^-$ 190 MeV $B = 0.00\%$
 $J/\psi \rightarrow \gamma \mu^+\mu^-$ 300 MeV $B = 0.00\%$
 $J/\psi \rightarrow \gamma \mu^+\mu^-$ 340 MeV $B = 0.00\%$

Although none of the masses we use are precisely determined, the width of each resonance is assumed to be equal to the mass of the resonance. To the extent that $B(\psi \rightarrow \mu^+\mu^-) \approx B(\psi \rightarrow e^+e^-)$ holds, the predicted J/ψ ratios, the result of neutral J/ψ is expected to be seven times greater than the field of $\mu^+\mu^-$ at J/ψ energy.



$K\pi$ and $K3\pi$ spectra 2.08.76 PRL 37 (76) 255



Resonances, Reflections: a non-trivial discovery example from 30 years ago (MLM digged this)



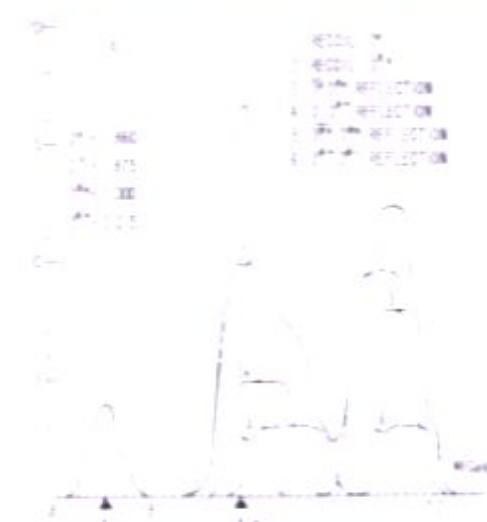
$K\pi$ and $K3\pi$ spectra 2.08.76 PRL 37 (76) 255

Experiment (indicated) $\sqrt{s} = 2.08$ GeV and $\sqrt{s} = 2.10$ GeV. Using our estimate of error we tabulate available energies \sqrt{s} and branching ratios β for D^0 decays as follows:

- $\sqrt{s} = 2.08$ GeV $\beta = 0.40$
- $\sqrt{s} = 2.09$ GeV $\beta = 0.47$
- $\sqrt{s} = 2.10$ GeV $\beta = 0.53$
- $\sqrt{s} = 2.11$ GeV $\beta = 0.60$
- $\sqrt{s} = 2.12$ GeV $\beta = 0.67$
- $\sqrt{s} = 2.14$ GeV $\beta = 0.77$

Although some of the classes we use are precisely determined, the most accurate features of branching in several D^0 channels are still to be determined.

To the extent that $\beta = 0.40$ to 0.77 applies, our plot shows the predicted D^0 ratios. The trend of several D^0 is expected to be better than greater than the trend of D^0 to D^+ in higher energy.



...excerpts

Observation in e^+e^- Annihilation of a Narrow State at 1865 MeV c^2 Decaying to $K\pi$ and $K\pi\pi\pi$

G. Goldhaber,* F. M. Pierre,† G. S. Abrams, M. S. Alam, A. M. Boyarski,‡ M. Breidenbach,
W. C. Carithers, W. Chouhowsky, S. C. Cooper, R. G. DeVoe, § M. Dothan, G. I. Feldman,
C. E. Friedberg, D. Fryberger, G. Hanson, J. Jaros, A. D. Johnson, J. A. Kadys,
H. R. Lahren, D. Lake,§ V. Lüth, H. L. Lynch, R. J. Madaras, C. C. Morehouse, †
H. K. Nguyen,** J. M. Paterson, M. L. Perl, † Peruzzi,** M. Piccolo,††
T. P. Pun, P. Rapidis, B. Richter, B. Sadoulet, R. B. Schindler,
R. F. Schmitters, † Siegrist, W. Tanenbaum, G. H. Trilling,
D. Vannucci,‡‡ J. S. Whitaker, and J. E. Wiss

*Lawrence Berkeley Laboratory and Department of Physics, University of California, Berkeley, California 94720,
and Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305
Received 14 June 1976*

We present evidence, from a study of multihadronic final states produced in e^+e^- annihilation at center-of-mass energies between 3.9 and 4.0 GeV, for the production of a new neutral state with mass 1865 ± 13 MeV c^2 and decay width less than 40 MeV c^2 that decays to $K^+\pi^-$ and $K^+\pi^-\pi^+\pi^-$. The recoil-mass spectrum for this state suggests that it is produced only in association with systems of comparable or larger mass.

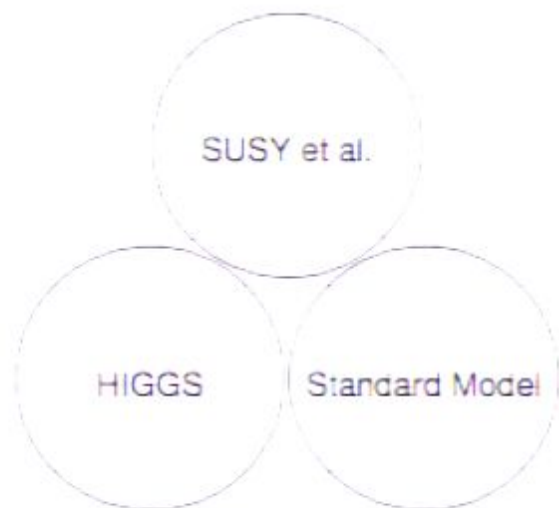
is charm found? De Rujula, Georgi, Glashow, PRL 37,(76) 398

In summary, it may be fruitful to search for charmed states—both mesons and baryons—at e^+e^- energies of 4–6 GeV. Peaks in invariant masses of several final hadrons are expected, accompanied by rich and energy-dependent structure in the recoil-mass spectra. Dozens of new hadrons await discovery.

We wish to thank Professor Gerson Goldhaber for discussing the data with us prior to publication.



SUSY is marking the engineering of the discovery path of beyond the standard model signatures[which consist of more than SUSY and SUSY-like frameworks and duals]:



- SUSY
- Heavy Stable Charged Particles
- $U(1)'$
- Technicolor
- Compositeness
- Flavor violation
- Little higgs
- Extra dimensions

Navigation controls including a search bar with the text "Black holes", a play/pause button, a zoom in/out button, and a close button.



Flavor for SUSY and beyond @LHC

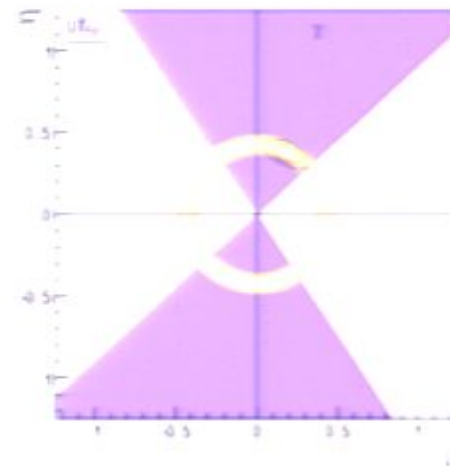
flavor has caused SUSY a lot of suffering...

In order to disentangle new physics effects, we should first determine CKM parameters by "tree-level" processes.

$$|V_{ub}|, f_{B\gamma} \quad V_{ub} = V_{ub} e^{i\theta_3}$$

- + Bd mixing and CP asymmetries
- + Bs mixing and CP asymmetries
- + ϵ_K and $B(K \rightarrow \pi\nu\nu)$

Fit from tree-level processes



We know (or constrain) which sector is affected by new physics.
Improvement of $f_{B\gamma}$ is essential.

More on the essential measurements for new physics searches at the B-factories and LHCb at the flavor session.

with the discovery of any new physics @LHC the role of flavor will be revised and upgraded into tool for navigation and characterization of the new physics



Theory-Experiment crossing

Since the time of Fermi, it took LHC to bring theory and experiments back close together

Das Nullte Theorem der Wissenschaftsgeschichte

lautet dass eine Entdeckung (Regel, Gesetzmässigkeit, Einsicht), die nach einer Person bennant ist, nicht von dieser Person herrührt

–E.P. Fischer, Fremde Federn, Im Gegenteil, July 24 2006, Die Zeit (from J. D. Jackson's "the first and the famous" collection) my message: Work and make the discoveries with care and respect to the data: the famous are often not the first and the first are often not famous but this is irrelevant in particular at a monumental time like the LHC one. Getting it right is important.



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Theory-Experiment crossing



Theory-Experiment crossing

Since the time of Fermi, it took LHC to bring theory and experiments back close together



Theory-Experiment crossing

Since the time of Fermi, it took the threat of the LHC to bring theory on shell



Theory-Experiment crossing

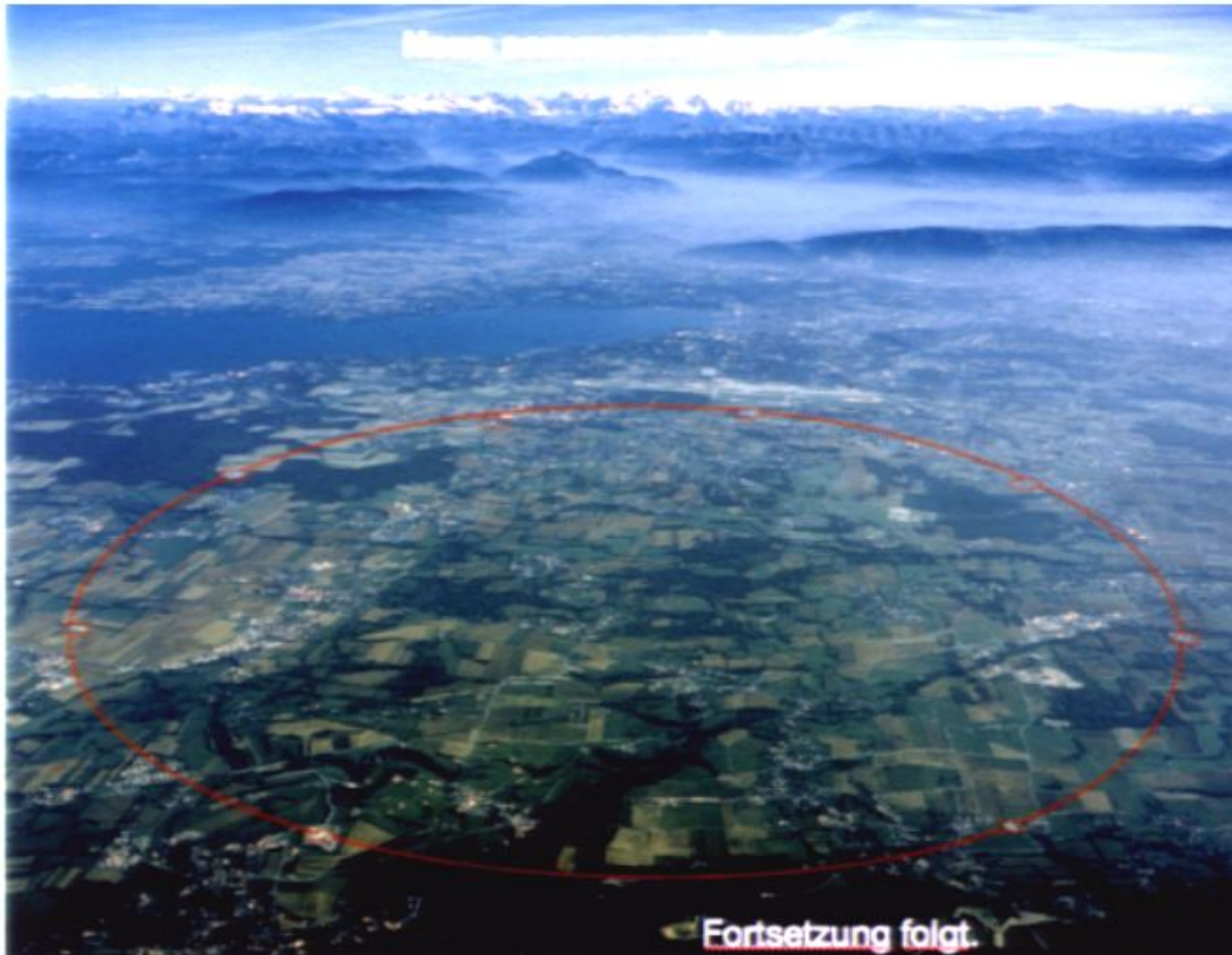
Since the time of Fermi, it took the threat of the LHC to bring theory on shell



Theory-Experiment crossing

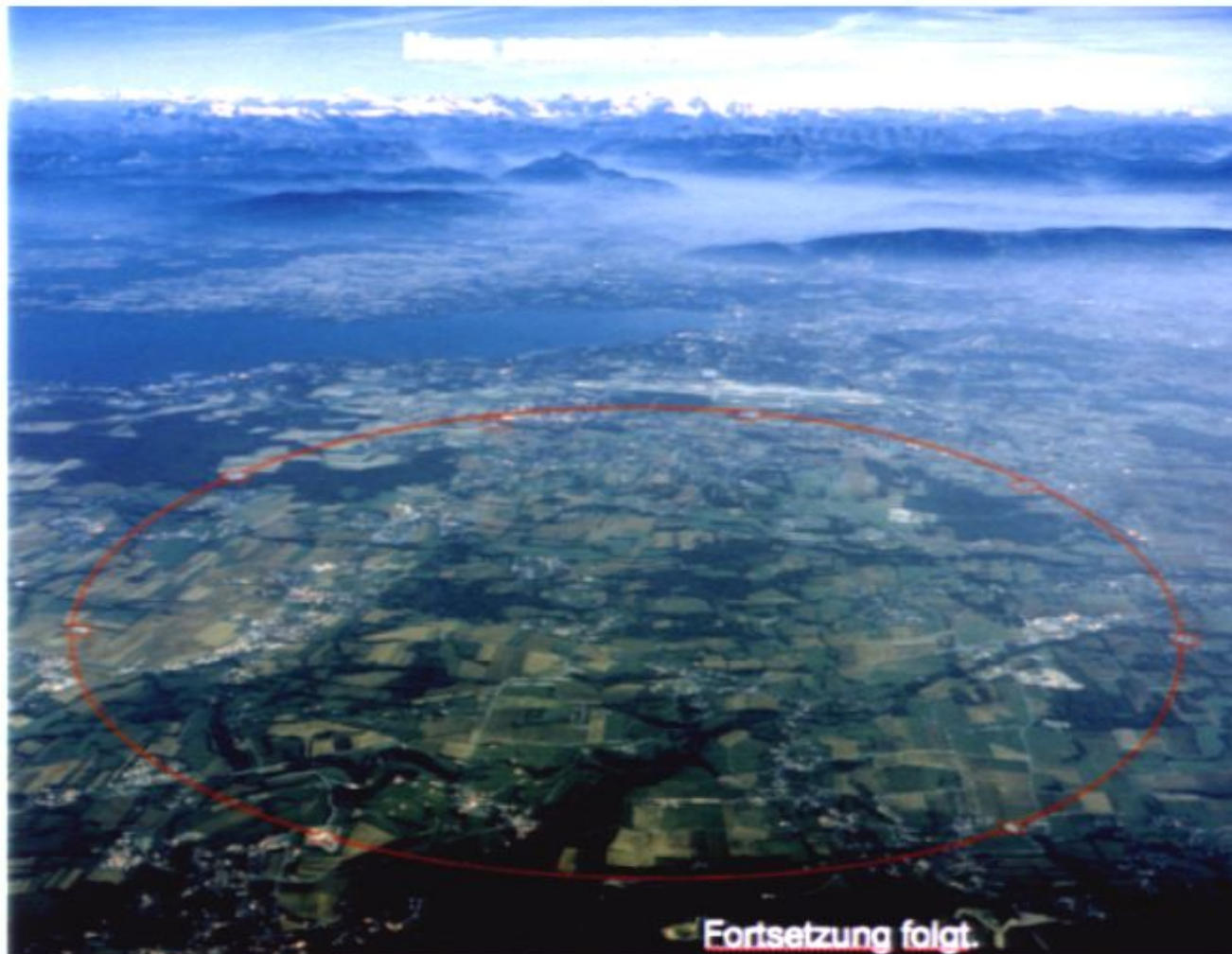
getting finally to experimental discoveries: PRICELESS





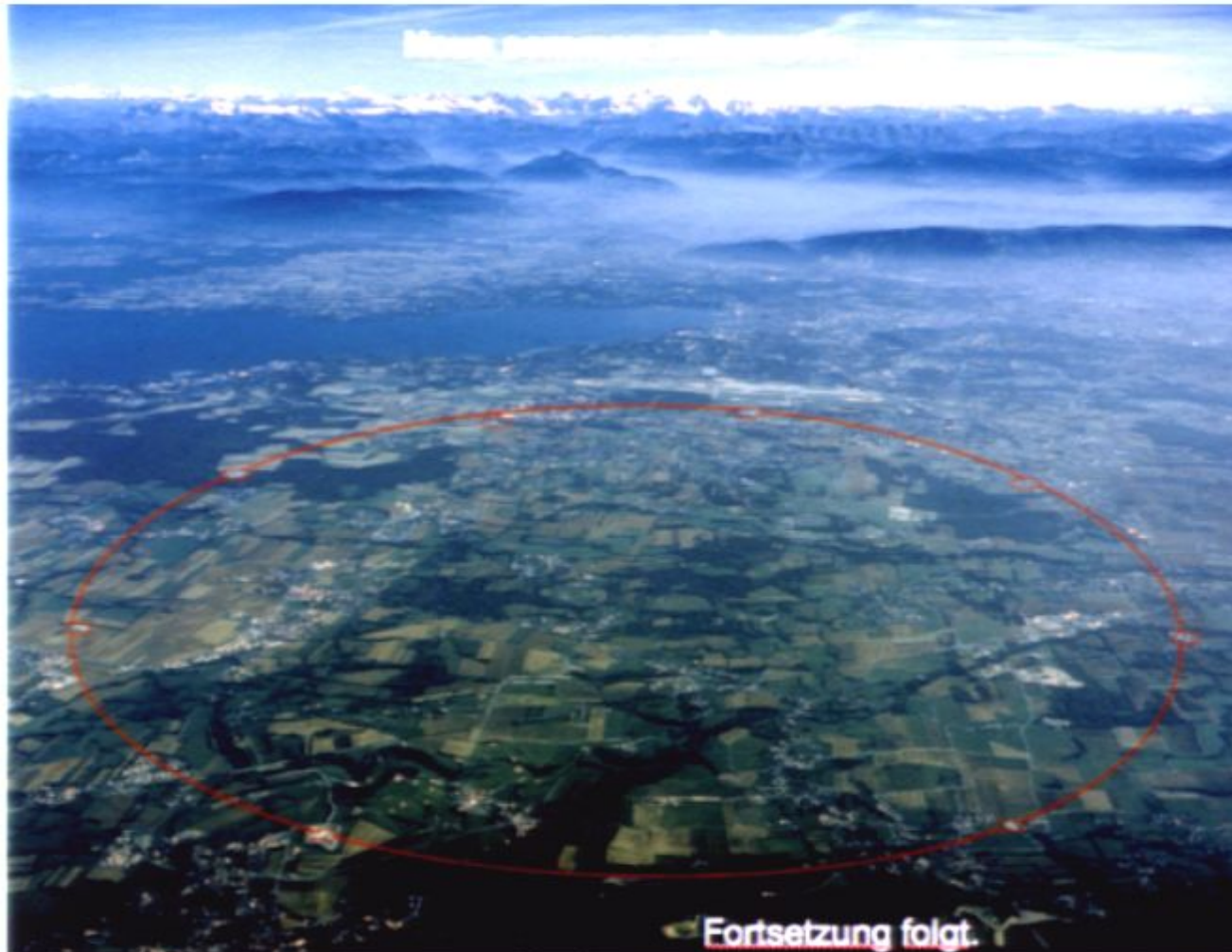
"faster, better, simpler, smarter" with the LHC data: The field needs the LHC input at the $\mathcal{O}(\text{pb}^{-1})$ time scale. Move to the data without strong pre-conceived ideas (i.e. escape the tyranny of MSSMs)





"faster, better, simpler, smarter" with the LHC data: The field needs the LHC input at the $\mathcal{O}(\text{pb}^{-1})$ time scale. Move to the data without strong pre-conceived ideas (i.e. escape the tyranny of MSSMs)





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<http://cmsdoc.cern.ch/~smaria/perimeter07/>

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CERN-PH

Perimeter Institute, Summer School on Particle Physics, Cosmology and Strings.



August 9, 2007