

Title: Mathematical structures of Quantum Mechanics and connections to operational principles

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Abstract:

*Mathematical structures for
Quantum Mechanics
and connections to operational principles*

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*Operational Quantum Physics and the
Quantum-Classical Contrast*

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On experimental science

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Knowledge on such state will allow us to **predict the results of forthcoming experiments** on a (similar) object in a similar situation.

Since necessarily we work with only partial prior knowledge of both system and experimental apparatus, the rules for the experiment must be given in a **probabilistic setting**.

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The interaction between object and apparatus produces one of a **set of possible transformations** of the object, each one occurring with some probability.

Postulates

- **Postulate 1 (Independent systems)** *There exist **independent** systems.*
- **Postulate 2 (Symmetric faithful state)** *For every composite system made of two identical physical systems there exists a **symmetric joint state** that is both **dynamically** and **preparationally faithful**.*
- **Postulate 3 (Local observability principle)** *For composite systems **local informationally complete observables** provide **global informationally complete observables**.*
- **Postulate 4 (Info-complete discriminating observable)** *For every system there exists a minimal info-complete observable that can be achieved using a joint **discriminating observable** on system+ ancilla.*

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Actions and outcomes

Experiment (or “action”): every experiment is described by a set $\mathbb{A} \equiv \{\mathcal{A}_j\}$ of possible transformations \mathcal{A}_j having overall unit probability, with the apparatus signaling the outcome j labeling which transformation actually occurred.

States

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State: A state ω for a physical system is a rule which provides the probability for any possible transformation within an experiment, namely:

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Identity transformation:
$$\omega(\mathcal{I}) = 1$$

Convex structure of states

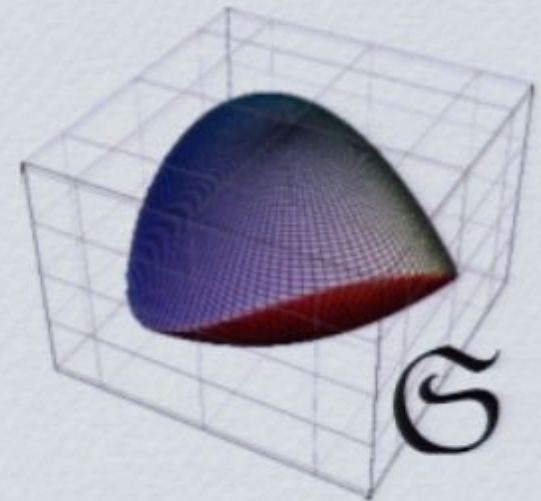
The possible states of a physical system make a convex set \mathfrak{S}

ω_1, ω_2 any two states:

$$\omega = \lambda\omega_1 + (1 - \lambda)\omega_2, \quad 0 \leq \lambda \leq 1,$$

corresponding to the probability rule

$$\omega(\mathcal{A}) = \lambda\omega_1(\mathcal{A}) + (1 - \lambda)\omega_2(\mathcal{A})$$



Monoid \mathcal{T} of transformations

Transformations make a monoid: the composition $\mathcal{A} \circ \mathcal{B}$ of two transformations \mathcal{A} and \mathcal{B} is itself a transformation. Consistency of composition of transformations requires associativity, namely

$$\mathcal{C} \circ (\mathcal{B} \circ \mathcal{A}) = (\mathcal{C} \circ \mathcal{B}) \circ \mathcal{A}$$

There exists the identical transformation \mathcal{I} which leaves the physical system invariant, and which for every transformation \mathcal{A} satisfies the composition rule

$$\mathcal{I} \circ \mathcal{A} = \mathcal{A} \circ \mathcal{I} = \mathcal{A}$$

Independent systems and local transformations

Independent systems and local experiments: two physical systems are “independent” if on each system it is possible to perform “local experiments” for which on every joint state one has the commutativity of the pertaining transformations

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Multipartite system: a collection of independent systems

Local state

For a multipartite system we define the local state $\omega|_n$ of the n -th system the state that gives the probability of any local transformation \mathcal{A} on the n -th system with all other systems untouched, namely

$$\omega|_n(\mathcal{A}) \doteq \Omega(\mathcal{I}, \dots, \mathcal{I}, \underbrace{\mathcal{A}}_{nth}, \mathcal{I}, \dots)$$

Conditional state

When composing two transformations \mathcal{A} and \mathcal{B} the probability that \mathcal{B} occurs conditioned that \mathcal{A} occurred before is given by the **Bayes rule**

$$p(\mathcal{B}|\mathcal{A}) = \frac{\omega(\mathcal{B} \circ \mathcal{A})}{\omega(\mathcal{A})}$$

Conditional state: the conditional state $\omega_{\mathcal{A}}$ gives the probability that a transformation \mathcal{B} occurs on the physical system in the state ω after the transformation \mathcal{A} occurred, namely

$$\omega_{\mathcal{A}}(\mathcal{B}) \doteq \frac{\omega(\mathcal{B} \circ \mathcal{A})}{\omega(\mathcal{A})}$$

No-signaling from the future

[Ozawa] The definition of conditional state needs to assume that

$$\sum_{\mathcal{B}_j \in \mathbb{B}} \omega(\mathcal{B}_j \circ \mathcal{A}) = \omega(\mathcal{A}), \quad \forall \mathbb{B}, \forall \mathcal{A}.$$

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Weights and Operations

Weight: un-normalized state

$$\omega = \frac{\tilde{\omega}}{\tilde{\omega}(\mathcal{I})}$$

$$0 \leq \tilde{\omega}(\mathcal{A}) \leq \tilde{\omega}(\mathcal{I}) < +\infty$$

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$$\tilde{\omega}_{\mathcal{A}}(\mathcal{B}) = \omega(\mathcal{B} \circ \mathcal{A})$$

Action of a transformation over a state (“Schrödinger picture”):

$$\mathcal{A} \omega := \text{Op}_{\mathcal{A}} \omega$$

$$(\mathcal{A} \omega)(\mathcal{B}) := \omega(\mathcal{B} \circ \mathcal{A})$$

Evolution as conditioning

Axioms

Theorems

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Evolution as conditioning

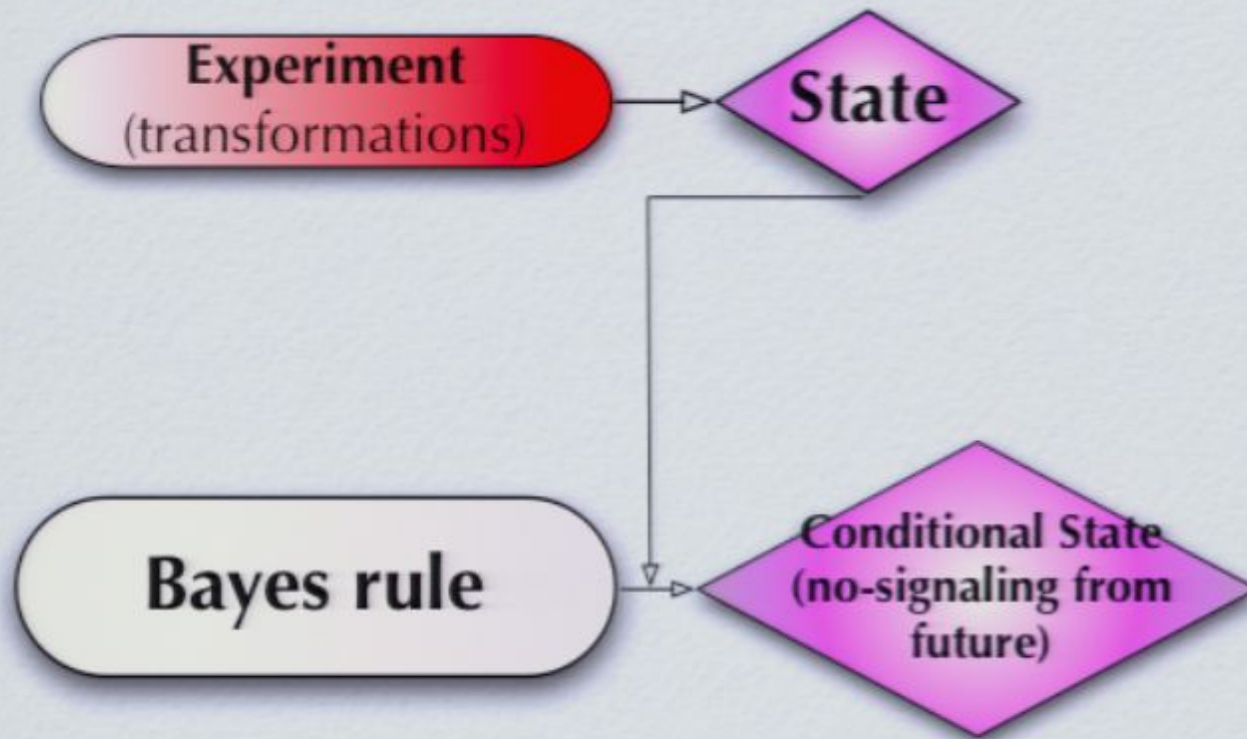
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Bayes rule

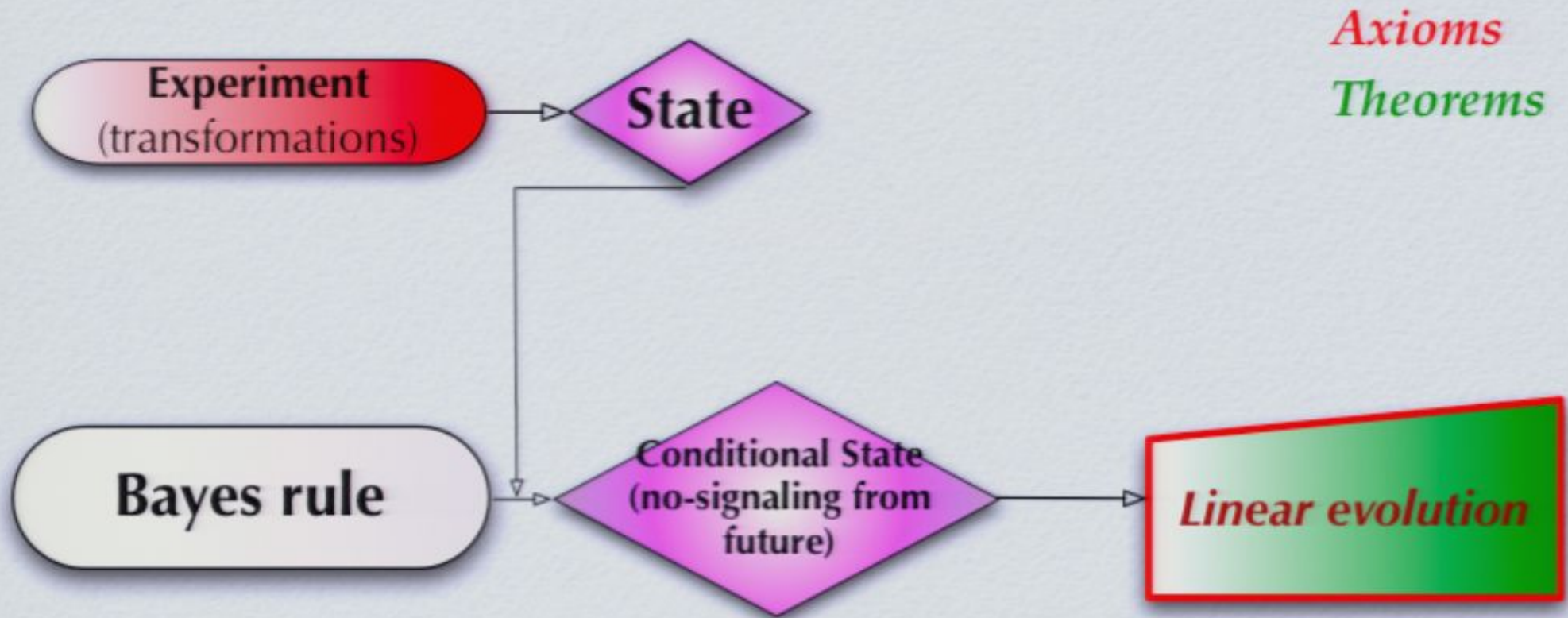
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Dynamical equivalence of transformations: two transformations \mathcal{A} and \mathcal{B} are dynamically equivalent if

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Informational equivalence of transformations: two transformations \mathcal{A} and \mathcal{B} are informationally equivalent if

$$\omega(\mathcal{A}) = \omega(\mathcal{B}) \quad \forall \omega \in \mathcal{G}$$

Addition of transformations

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$$\forall \omega \in \mathfrak{S} \quad \omega(\mathcal{A}_1 + \mathcal{A}_2) = \omega(\mathcal{A}_1) + \omega(\mathcal{A}_2) \quad (\text{info-class})$$

$$\forall \omega \in \mathfrak{S} \quad \omega_{\mathcal{A}_1 + \mathcal{A}_2} = \frac{\omega(\mathcal{A}_1)}{\omega(\mathcal{A}_1 + \mathcal{A}_2)} \omega_{\mathcal{A}_1} + \frac{\omega(\mathcal{A}_2)}{\omega(\mathcal{A}_1 + \mathcal{A}_2)} \omega_{\mathcal{A}_2} \quad (\text{dyn-class})$$

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$$(\mathcal{A}_1 + \mathcal{A}_2)\omega = \mathcal{A}_1\omega + \mathcal{A}_2\omega$$

Rescaling of transformations

Multiplication by a scalar: for each transformation \mathcal{A} the transformation $\lambda\mathcal{A}$ for $0 \leq \lambda \leq 1$ is defined as the transformation which is dynamically equivalent to \mathcal{A} but occurs with probability $\omega(\lambda\mathcal{A}) = \lambda\omega(\mathcal{A})$

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Convex structure for transformations \mathcal{T} and for actions

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“Heisenberg picture”:
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However a local action $\mathbb{A} \equiv \{\mathcal{A}_j\}$ on system 2 does not affect the local state on system 1, more precisely:

acausality of local actions: any local action on a system is equivalent to the identity transformation on another independent system.

$$\mathbb{A} \equiv \mathcal{I}(\mathbb{A}) := \sum_{\mathcal{A}_j \in \mathbb{A}} \mathcal{A}_j$$

$$\forall \Omega \in \mathfrak{S}^{\times 2}, \forall \mathbb{A}, \quad \Omega_{\mathbb{A}, \mathcal{I}}|_2 = \Omega|_2$$

No-signaling

Theorem 1 (No-signaling) *Any local action on a system does not affect another independent system. More precisely, any local action on a system is equivalent to the identity transformation when viewed from another independent system. In equations one has*

$$\forall \Omega \in \mathfrak{S}^{\times 2}, \forall \mathbb{A}, \quad \Omega_{\mathbb{A}, \mathcal{I}}|_2 = \Omega|_2. \quad (1)$$

Proof. Since the two systems are dynamically independent, for every two local transformations one has $\mathcal{A}^{(1)} \circ \mathcal{A}^{(2)} = \mathcal{A}^{(2)} \circ \mathcal{A}^{(1)}$, which implies that $\Omega(\mathcal{A}^{(1)} \circ \mathcal{A}^{(2)}) = \Omega(\mathcal{A}^{(2)} \circ \mathcal{A}^{(1)}) \equiv \Omega(\underline{\mathcal{A}}^{(1)}, \underline{\mathcal{A}}^{(2)})$. By definition, for $\mathcal{B} \in \mathfrak{T}$ one has $\Omega|_2(\mathcal{B}) = \Omega(\mathcal{I}, \mathcal{B})$, and using the addition rule for transformations and reminding the identification $\mathbb{A} \equiv \sum_j \mathcal{A}_j$, one has

$$\Omega(\mathbb{A}, \mathcal{B}) = \Omega(\underline{\mathbb{A}}, \underline{\mathcal{B}}) = \Omega(\underline{\mathcal{I}}, \underline{\mathcal{B}}) =: \Omega|_2(\mathcal{B}). \quad (2)$$

On the other hand, we have

$$\Omega_{\mathbb{A}, \mathcal{I}}|_2(\mathcal{B}) = \Omega((\mathcal{I}, \mathcal{B}) \circ (\mathbb{A}, \mathcal{I})) = \Omega(\mathbb{A}, \mathcal{B}), \quad (3)$$

No-signaling from dynamical independence

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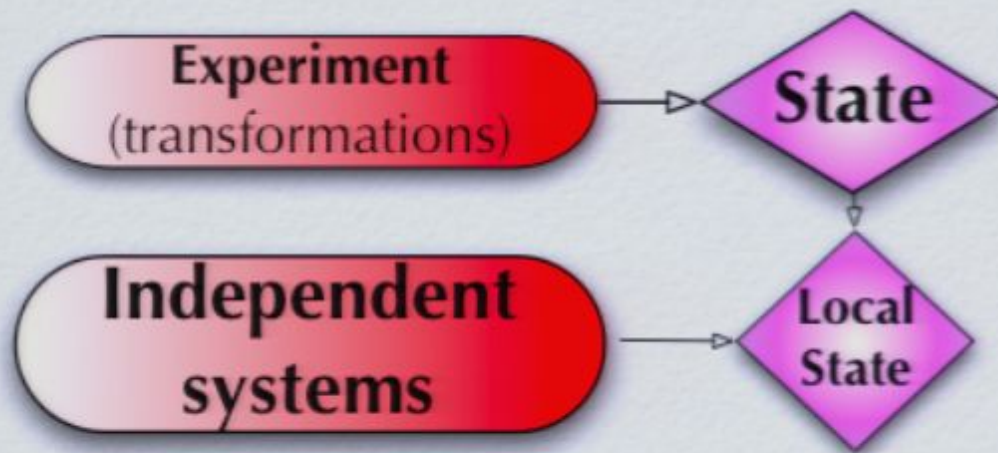
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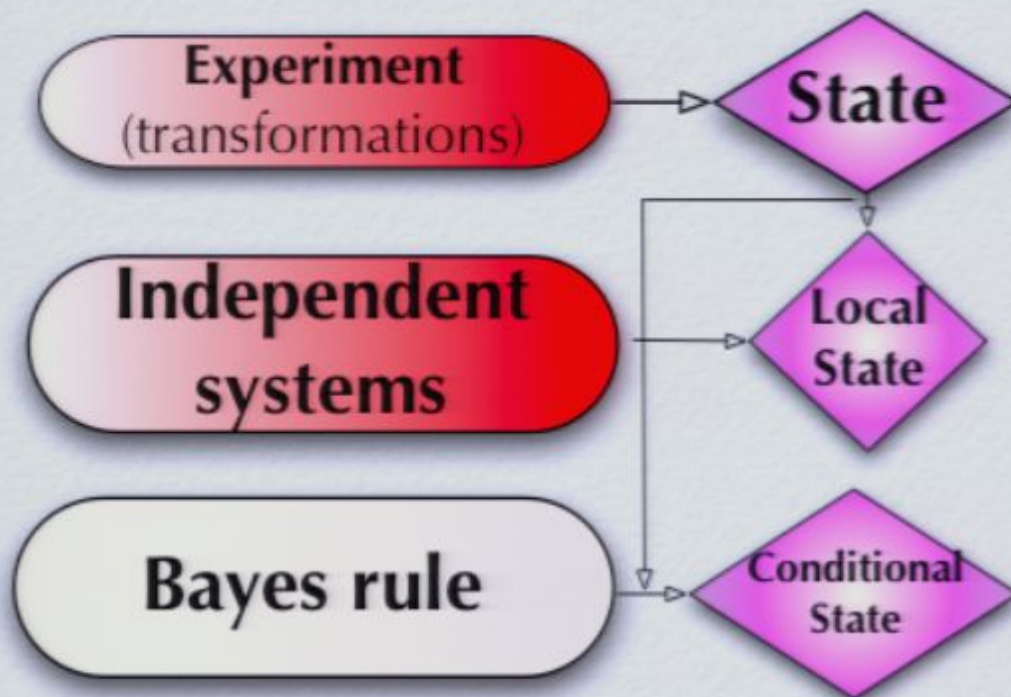
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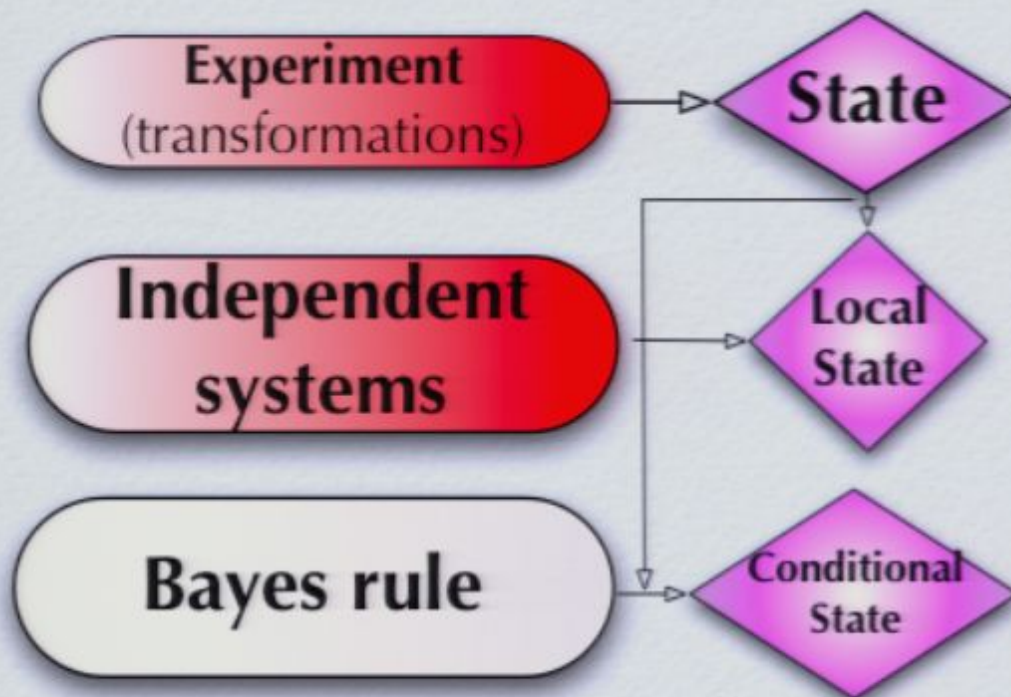
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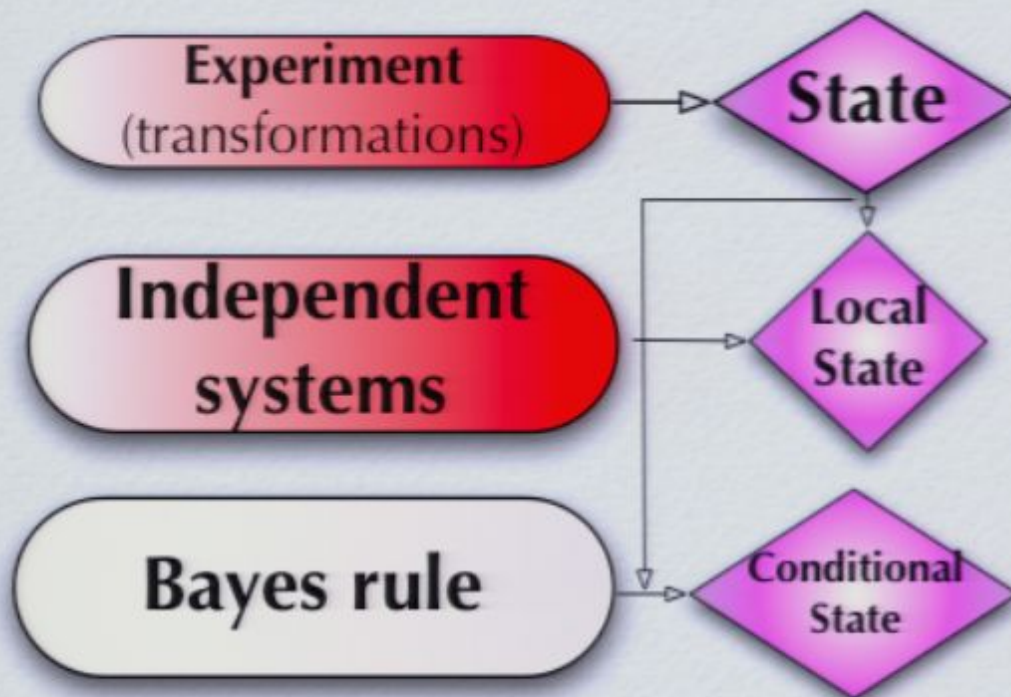
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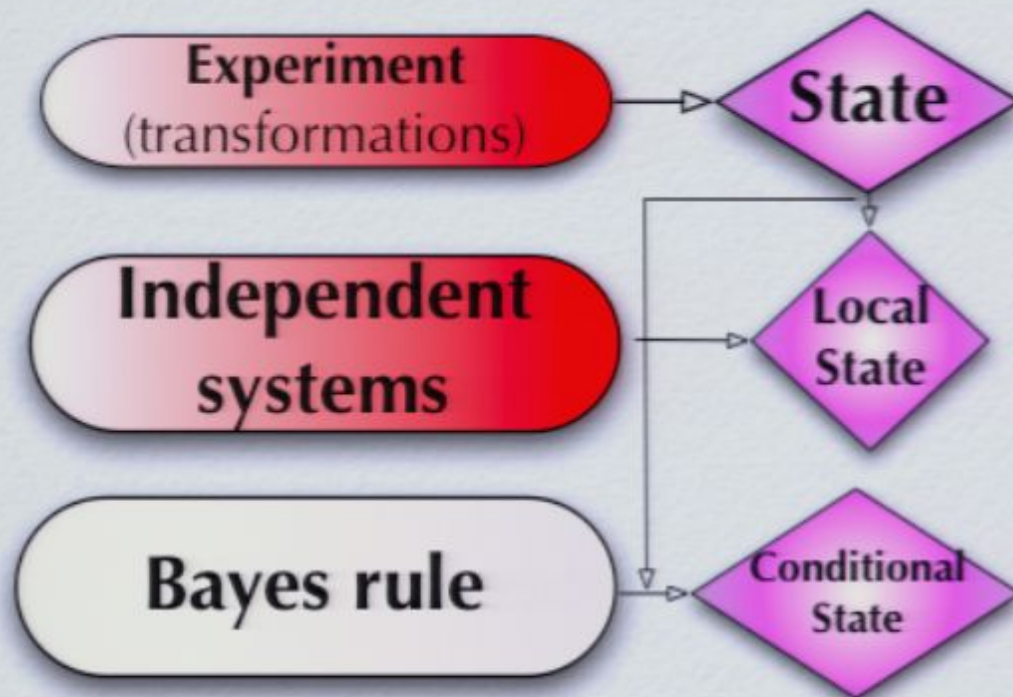
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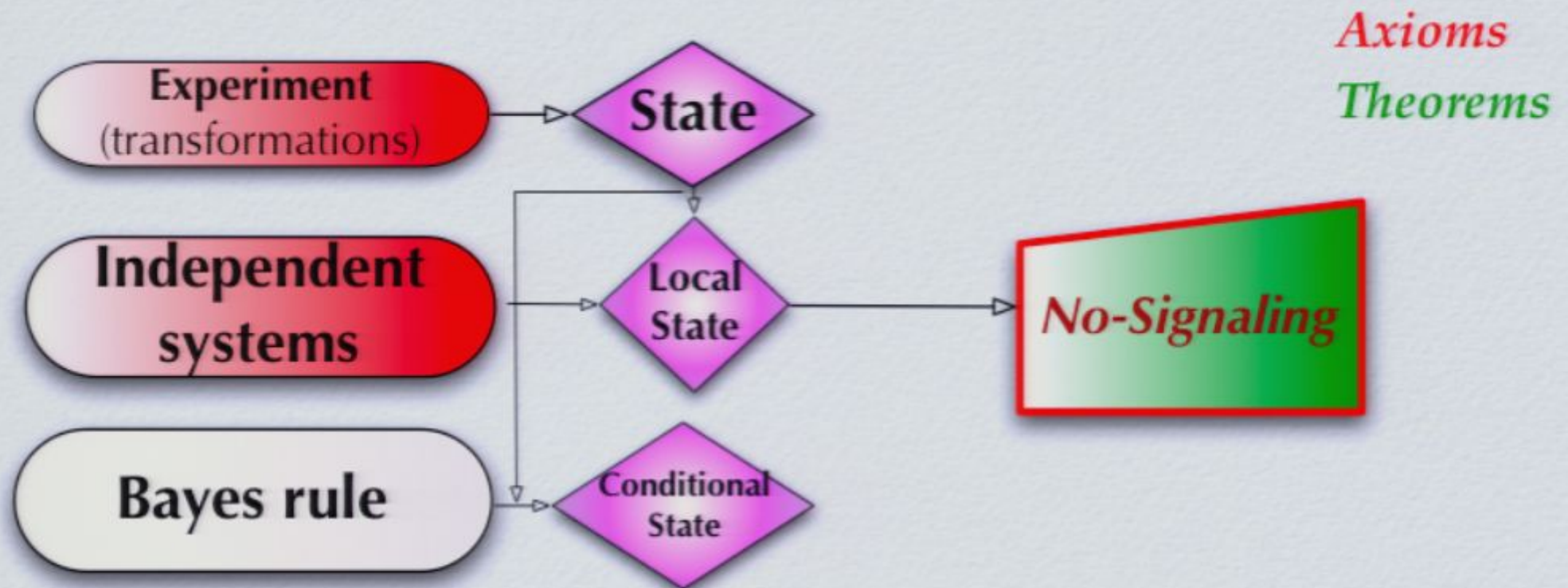
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Generalized weights, transformations, and effects

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gen. effects $\mathfrak{P}_{\mathbb{R}}$: $\|\underline{\mathcal{A}}\| := \sup_{\omega \in \mathfrak{S}} |\omega(\underline{\mathcal{A}})|$

gen. weights $\mathfrak{W}_{\mathbb{R}}$: $\|\tilde{\omega}\| := \sup_{\mathfrak{P}_{\mathbb{R}} \ni \|\underline{\mathcal{A}}\| \leq 1} |\tilde{\omega}(\underline{\mathcal{A}})|$

gen. transformations $\mathfrak{T}_{\mathbb{R}}$: $\|\underline{\mathcal{A}}\| := \sup_{\mathfrak{P}_{\mathbb{R}} \ni \|\underline{\mathcal{B}}\| \leq 1} \|\underline{\mathcal{B}} \circ \underline{\mathcal{A}}\|$

$\mathfrak{W}_{\mathbb{R}} \ \mathfrak{P}_{\mathbb{R}}$ *dual Banach pair* under the pairing

$$l_{\underline{\mathcal{A}}}(\omega) \doteq \omega(\underline{\mathcal{A}})$$

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Banach-space structures

Axioms

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Banach-space structures

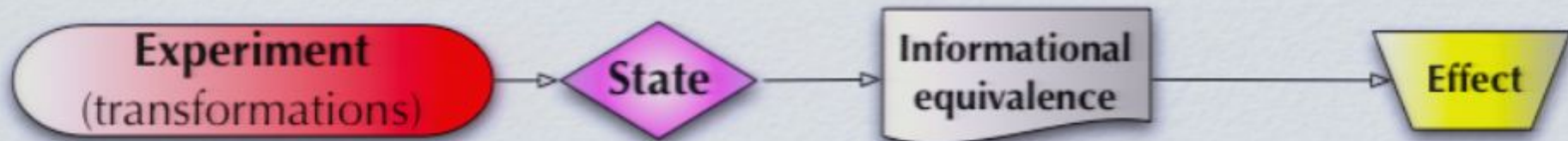
Axioms

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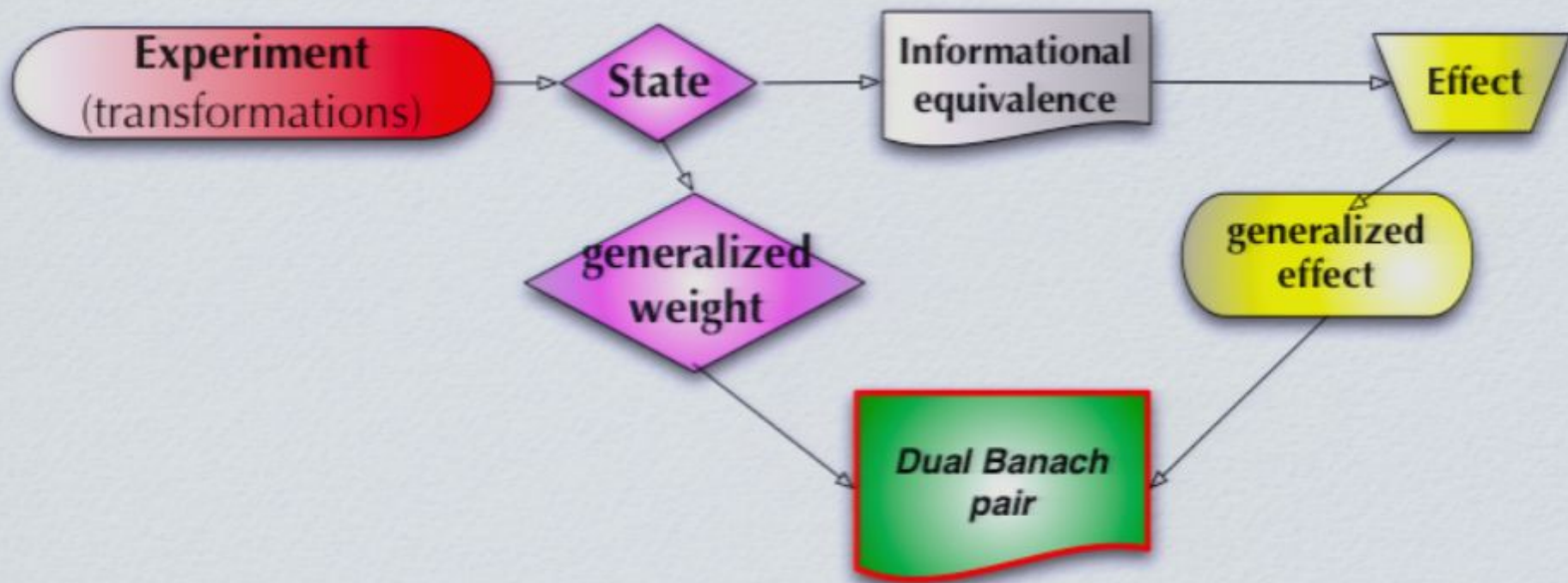
Theorems



Banach-space structures

Axioms

Theorems



Observable

Observable: a complete set of effects $\mathbb{L} = \{l_i\}$

$$\sum_j l_j = \underline{\mathcal{I}}$$

Informationally complete observable

Informationally complete observable: an observable $\mathbb{L} = \{l_i\}$ is informationally complete if any effect l can be written as linear combination of elements of \mathbb{L} , namely there exist coefficients $c_i(l)$ such that

$$l = \sum_{i=1}^{|\mathbb{L}|} c_i(l) l_i$$

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Bloch representation

$$l_{\underline{\mathcal{A}}} = \sum_j m_j(\underline{\mathcal{A}}) n_j \quad l_{\underline{\mathcal{A}}}(\omega) = \boldsymbol{m}(\underline{\mathcal{A}}) \cdot \boldsymbol{n}(\omega) + q(\underline{\mathcal{A}})$$

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**Conditioning:
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transformation**

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$$M_{ij}(\mathcal{A}) = \begin{pmatrix} q(\underline{\mathcal{A}}) & \mathbf{m}(\underline{\mathcal{A}}) \\ \mathbf{k}(\mathcal{A}) & \mathbf{M}(\mathcal{A}) \end{pmatrix}$$

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Theorem: *there always exists a minimal informationally complete observable.*

Proof. By definition $\mathfrak{P}_{\mathbb{R}} = \text{Span}_{\mathbb{R}}(\mathfrak{P})$, whence there must exist a spanning set for $\mathfrak{P}_{\mathbb{R}}$ that is contained in \mathfrak{P} . The maximal number of elements of this set that are linearly independent will constitute a *basis*, which we suppose has finite cardinality $\dim(\mathfrak{P}_{\mathbb{R}})$. It remains to be shown that it is possible to have a basis with sum of elements equal to \mathcal{I} , and that such basis is obtained operationally starting from the available observables from which we constructed \mathfrak{P} .

If all observables are *uninformative* (i. e. with all effects proportional to \mathcal{I}), then $\mathfrak{P}_{\mathbb{R}} = \text{Span}(\mathcal{I})$, \mathcal{I} is a minimal infocomplete observable, and the statement of the theorem is proved. Otherwise, there exists at least an observable $\mathbb{E} = \{l_i\}$ with $n \geq 2$ linearly independent effects. If this is the only observable, again the theorem is proved. Otherwise, take a new binary observable $\mathbb{E}_2 = \{x, y\}$ from the set of available ones (you can take different binary observables out of a given observable with more than two outcomes by summing up effects to yes-no observables). If $x \in \text{Span}(\mathbb{E})$ discard it. If $x \notin \text{Span}(\mathbb{E})$, then necessarily also $y \notin \text{Span}(\mathbb{E})$ [since if there exists coefficients λ_i such that $y = \sum_i \lambda_i l_i$, then $x = \sum_i (1 - \lambda_i) l_i$]. Now, consider the observable

$$\mathbb{E}' = \left\{ \frac{1}{2}y, \frac{1}{2}(l_1 + x), \frac{1}{2}l_2, \dots, l_n \right\} \quad (1)$$

(which operationally corresponds to the random choice between the observables \mathbb{E} and \mathbb{E}_2 with probability $\frac{1}{2}$, and with the events corresponding to x and l_1 made indistinguishable). This new observable has now $|\mathbb{E}'| = n + 1$ linearly independent effects (since y is linearly independent on the l_i and one has $y = \sum_{i=1}^n l_i - x = \sum_{i=2}^n l_i + l_1 - x$). By iterating the above procedure we reach $|\mathbb{E}'| = \dim(\mathfrak{P}_{\mathbb{R}})$, and we have so realized an apparatus that measures a minimal informationally complete observable. ■

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$$m(\omega|\omega^*) = \frac{\Delta F(\omega^*) (m(\omega) + k_1(\omega^*))}{m(\omega^*) + n(\omega^*) + q(\omega^*)}$$

$$M_{\omega^*}(\omega) := \begin{pmatrix} q(\omega^*) & m(\omega^*) \\ k_1(\omega^*) & \Delta F(\omega^*) \end{pmatrix}$$

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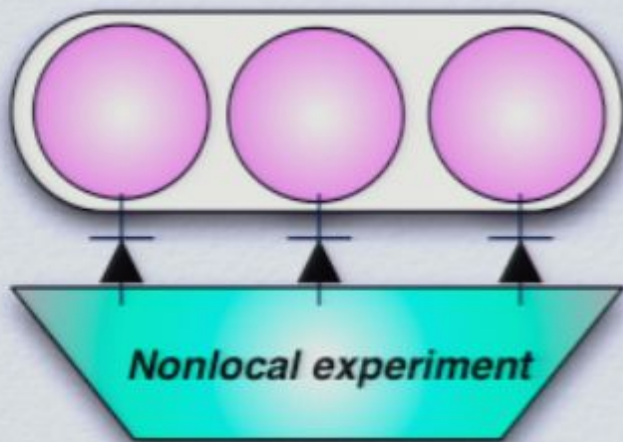
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Postulate 3: Local observability principle

For composite systems local info-complete observables provide global info-complete observables.

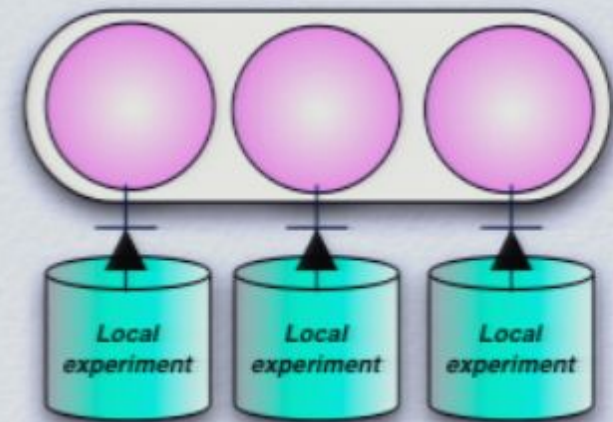
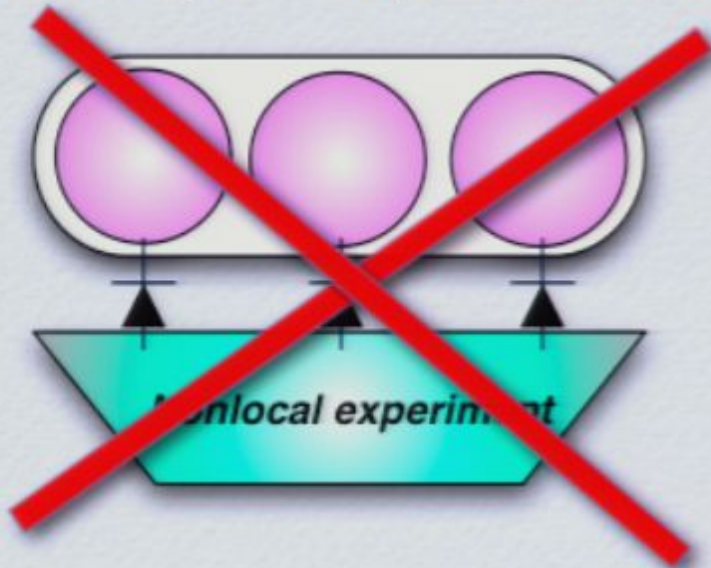
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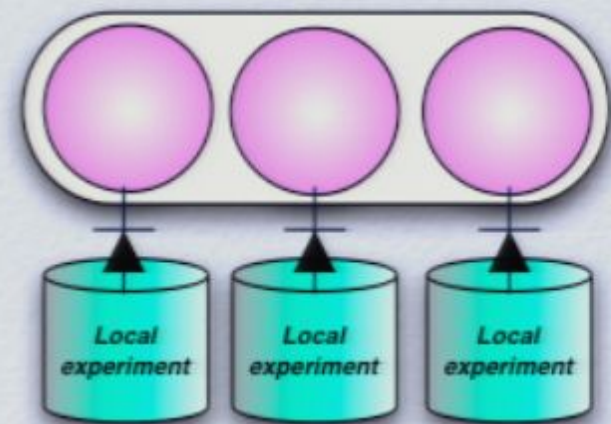
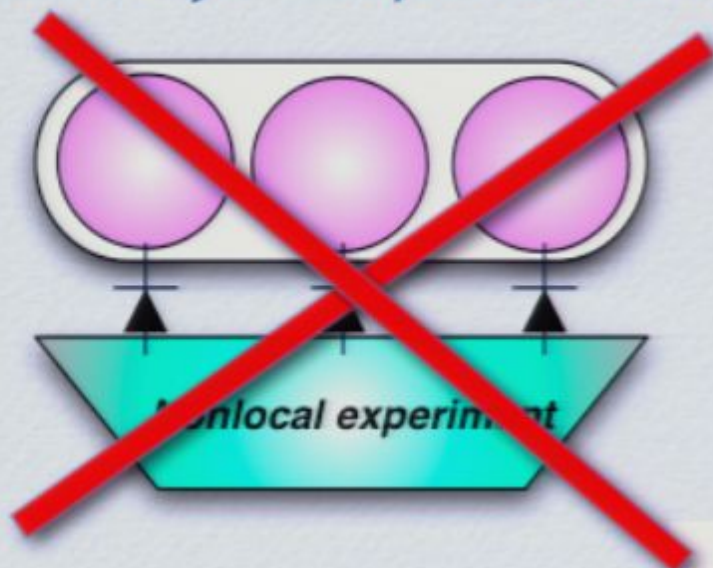
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Holism



Reductionism

identity for the affine dimension of composite systems

$$\dim(\mathfrak{S}_{1,2}) = \dim(\mathfrak{S}_1) \dim(\mathfrak{S}_2) + \dim(\mathfrak{S}_1) + \dim(\mathfrak{S}_2)$$

Postulate 3: Local observability principle

identity for the affine dimension of composite systems

$$\dim(\mathfrak{S}_{12}) = \dim(\mathfrak{S}_1) \dim(\mathfrak{S}_2) + \dim(\mathfrak{S}_1) + \dim(\mathfrak{S}_2)$$

Proof. We first prove that the left side is a lower bound for the right side. Indeed, the number of outcomes of a minimal informationally complete observable is $\dim(\mathfrak{S}) + 1$, since it equals the dimension of the affine space embedding the convex set of states \mathfrak{S} plus an additional dimension for normalization. Now, consider a global informationally complete measurement made of two local minimal informationally complete observables measured jointly. It has number of outcomes $[\dim(\mathfrak{S}_1) + 1][\dim(\mathfrak{S}_2) + 1]$. However, we are not guaranteed that the joint observable is itself minimal, whence the bound.

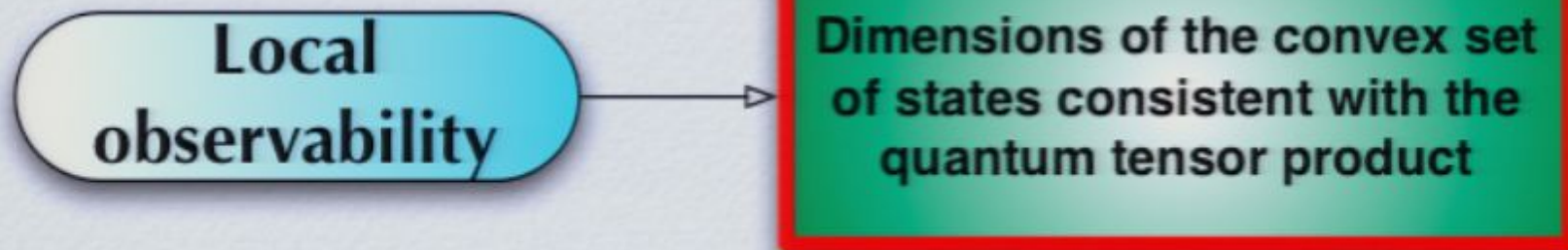
The opposite inequality can be easily proved by considering that a global informationally incomplete measurement made of minimal local informationally complete measurements should belong to the linear span of a minimal global informationally complete measurement. ■

Postulate 3: Local observability principle

Postulates

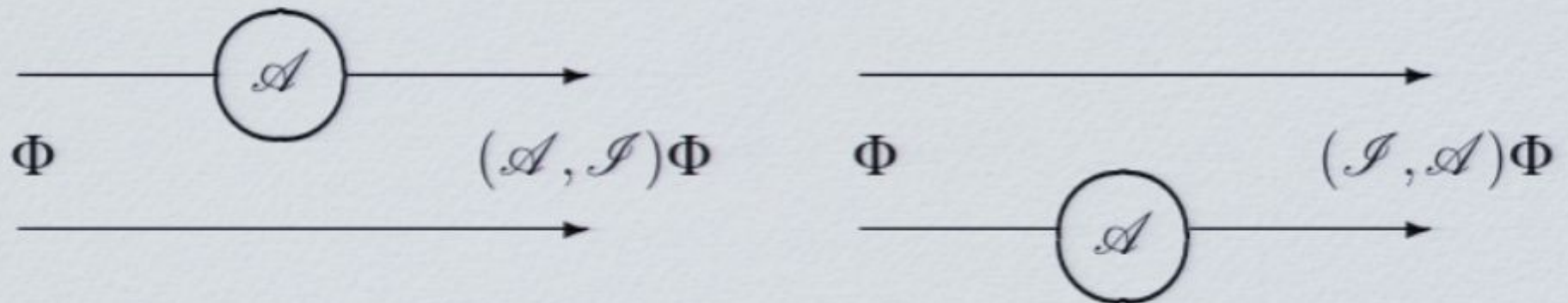
Axioms

Theorems



Faithful states

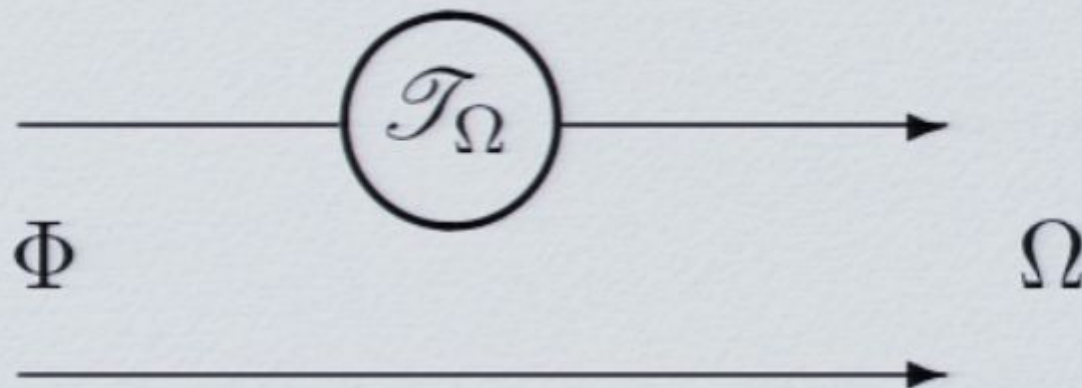
Dynamically faithful state: we say that a state Φ of a bipartite system is dynamically faithful if when acting on it with a local transformation \mathcal{A} on one system the output conditioned weight $(\mathcal{A}, \mathcal{I})\Phi$ is in 1-to-1 correspondence with the transformation \mathcal{A}



$$(\mathcal{A}, \mathcal{I})\Phi = 0 \implies \mathcal{A} = 0, \quad \forall \mathcal{A} \in \mathcal{I}_{\mathbb{R}}$$

Faithful states

Preparationally faithful state: we say that a state Φ of a bipartite system is preparationally faithful if every joint state Ω can be achieved by a suitable local transformation \mathcal{T}_Ω on one system occurring with nonzero probability



Faithful states

Symmetric bipartite state: we call a joint state Φ of a bipartite system symmetric if

$$\Phi(\mathcal{A}, \mathcal{B}) = \Phi(\mathcal{B}, \mathcal{A})$$

Perfectly discriminating observable

Perfectly discriminable observable/states:

an observable $\mathbb{L} = \{l_i\}$ such that there exist states $\{\omega_j\}$ satisfying

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Informational dimension $\dim_{\#}(\mathfrak{S})$: maximal number of perfectly discriminable states

Postulate 4: Informationally complete discriminating observable

For every system there exists a minimal info-complete observable that can be achieved by means of a joint discriminating observable on system+ancilla[†]

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$$\dim(\mathfrak{S}) = \dim_{\#}(\mathfrak{S}^{\times 2}) - 1$$

Dimensionality identities

\Rightarrow		
state-effect duality	$\dim(\mathfrak{P}_{\mathbb{R}}) = \dim(\mathfrak{S}) + 1$	(D1)
P3 (loc. observability)	$\dim(\mathfrak{S}_{12}) = \dim(\mathfrak{S}_1) \dim(\mathfrak{S}_2) + \dim(\mathfrak{S}_1) + \dim(\mathfrak{S}_2)$	(D2)
P4 (infoc. as joint discr.)	$\dim(\mathfrak{S}) = \dim_{\#}(\mathfrak{S}^{\times 2}) - 1$	(D4)
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$$\dim(\mathfrak{P}_{\mathbb{R}}) = \dim_{\#}(\mathfrak{S})^2$$

Positive bilinear form

Positive form over generalized effects: from Φ real symmetric form over effects obtain the positive form (for finite dimensions)

$$|\Phi| := \Phi_+ - \Phi_-$$

$$|\Phi|(\underline{\mathcal{A}}, \underline{\mathcal{B}}) = \Phi(\underline{\mathcal{A}}, \varsigma(\underline{\mathcal{B}})), \quad \varsigma(\underline{\mathcal{A}}) = (\mathcal{P}_+ - \mathcal{P}_-)(\underline{\mathcal{A}}) \\ \varsigma^2 = \mathcal{I}$$

The complex Hilbert space formulation



For finite dimensions the real Hilbert space $\mathfrak{P}_{\mathbb{R}}$ is isomorphic to the real Hilbert space of Hermitian complex matrices representing selfadjoint operators over a complex Hilbert space \mathbf{H} of dimensions $\dim(\mathbf{H}) = \dim_{\#}(\mathfrak{S})$.

**This is the Hilbert space formulation
of Quantum Mechanics**

The complex Hilbert space formulation

If the state is also preparationally faithful then one can make every state correspond to an effect



Then one can write the probability rule in terms of a real scalar product pairing between states and effects, with the convex cones of effects and states corresponding to the convex cone of positive matrices.

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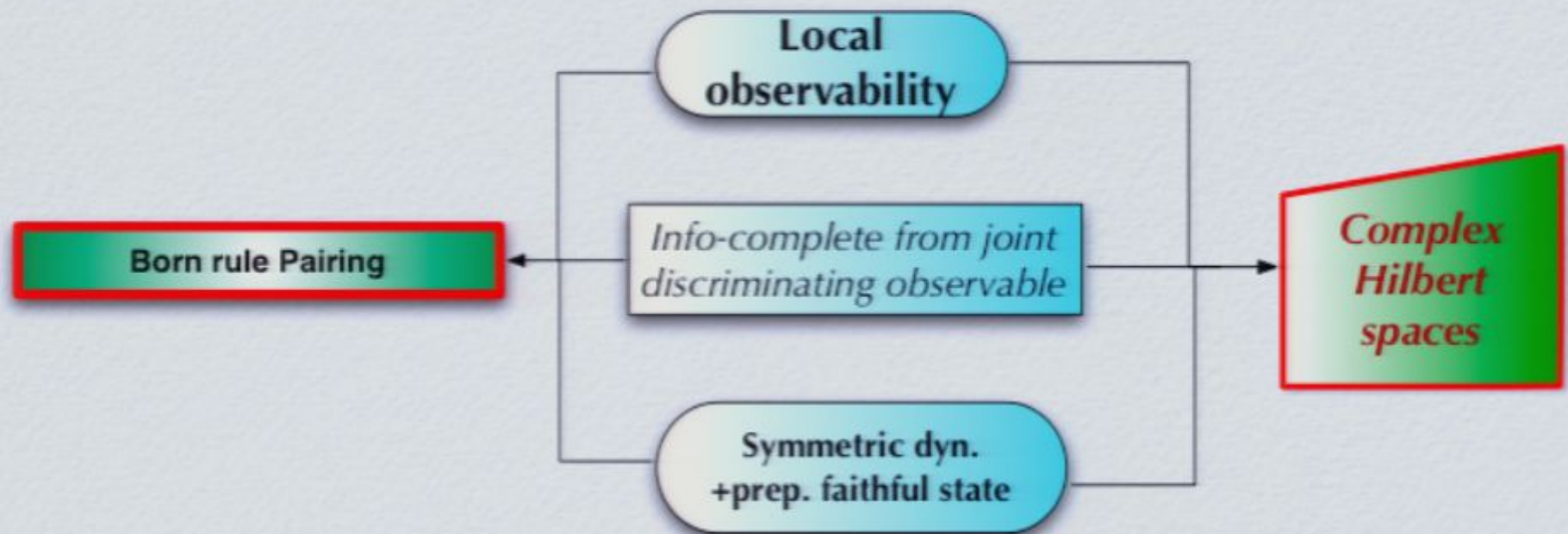
Since Φ is *preparationally* faithful, then for every state ω there exists a suitable transformation \mathcal{T}_ω such that $\omega = \Phi_{\mathcal{I}, \mathcal{T}_\omega}|_1$ with probability $\Phi(\mathcal{I}, \mathcal{T}_\omega) > 0$

Then we can write the probability rule in terms of the pairing between states and effects:

$$\omega(\underline{\mathcal{C}}) = \Phi_{\mathcal{I}, \mathcal{T}_\omega}|_1(\underline{\mathcal{C}}) = |\Phi|(\underline{\mathcal{C}}, \widetilde{\mathcal{T}}_\omega),$$

$$\widetilde{\mathcal{T}}_\omega = \frac{\varsigma(\underline{\mathcal{T}}_\omega)}{\Phi(\underline{\mathcal{I}}, \underline{\mathcal{T}}_\omega)}$$

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Rest of the construction:

- *construct complex operators by complex linear combination of effects*

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• ...

Construction of a C^* -algebra of transformations

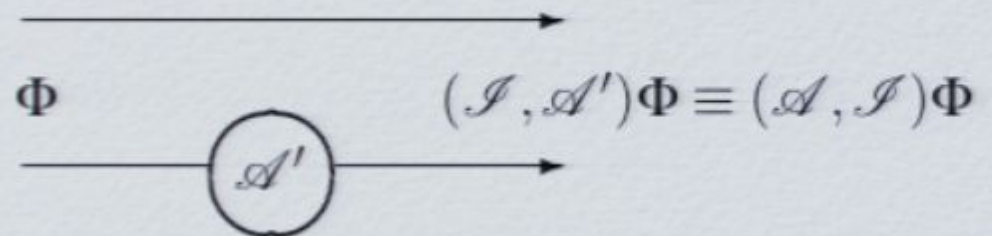
Operational definition of *transposed*

Existence of symmetric faithful states



“transposition” over the real algebra \mathcal{A} of (generalized) transformations

$$\mathcal{A} \longleftrightarrow \mathcal{A}'$$



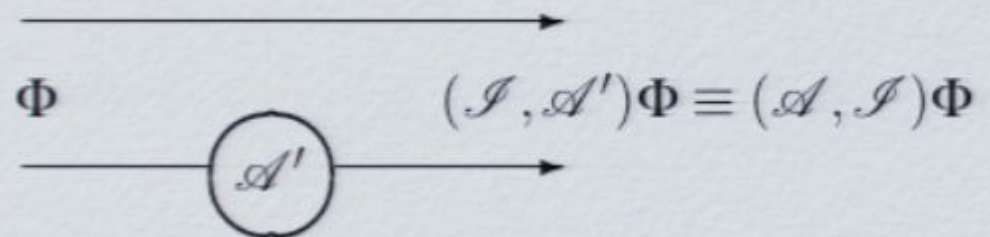
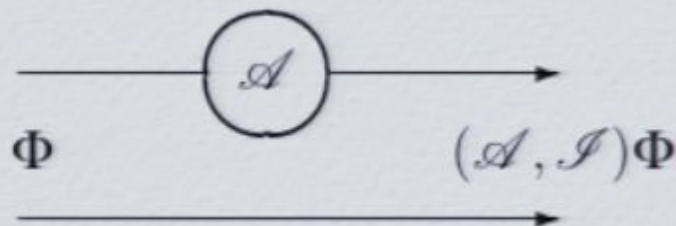
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$$\Phi(\mathcal{B} \circ \mathcal{A}, \mathcal{C}) = \Phi(\mathcal{B}, \mathcal{C} \circ \mathcal{A}')$$

The complex conjugation

The involution ς corresponds to a generalized transformation

$$\varsigma(\underline{\mathcal{A}}) = \underline{\mathcal{A}} \circ \mathcal{Z}$$

Correspondingly the involution over transformations reads

$$\varsigma(\mathcal{A}) = \mathcal{Z} \circ \mathcal{A} \circ \mathcal{Z}$$

which is composition preserving, namely

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The adjoint

Scalar product over $\mathfrak{P}_{\mathbb{R}}$:

$$\Phi \langle \underline{\mathcal{B}} | \underline{\mathcal{A}} \rangle_{\Phi} := \Phi(\varsigma(\underline{\mathcal{B}}'), \underline{\mathcal{A}}') = \Phi|_1(\mathcal{B}^{\dagger} \circ \mathcal{A})$$

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ζ works as a complex-conjugation in the sense that

$\mathcal{A}^{\dagger} := \zeta(\mathcal{A}')$ works as an adjoint, namely

$$\Phi \langle \mathcal{C}^{\dagger} \circ \underline{\mathcal{A}} | \underline{\mathcal{B}} \rangle_{\Phi} = \Phi \langle \underline{\mathcal{A}} | \mathcal{C} \circ \underline{\mathcal{B}} \rangle_{\Phi}$$

GNS construction for representing transformations

Representations π_Φ of transformations $\mathcal{A} \in \mathcal{A}$ over effects \mathcal{A}/\mathcal{I}

$$\pi_\Phi(\mathcal{A})|\underline{\mathcal{B}}\rangle_\Phi \doteq |\underline{\mathcal{A} \circ \mathcal{B}}\rangle_\Phi$$

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$$\omega(\underline{\mathcal{A}}) = {}_\Phi \langle \underline{\mathcal{A}}^\dagger | \varrho \rangle_\Phi$$

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C^* -algebra of transformations

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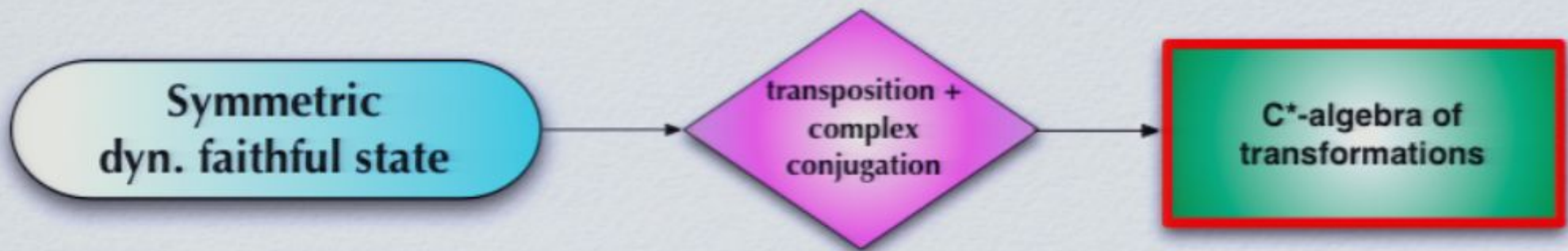


C^* -algebra of transformations

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C*-algebra of transformations

state-effect duality	$\dim(\mathfrak{P}) = \dim(\mathfrak{S}) + 1$	(D1)
P2 (prep. faith.)	$\dim(\mathfrak{T}) = \dim(\mathfrak{S}^{\times 2}) + 1$	(T)
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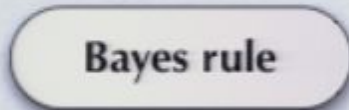
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Summary

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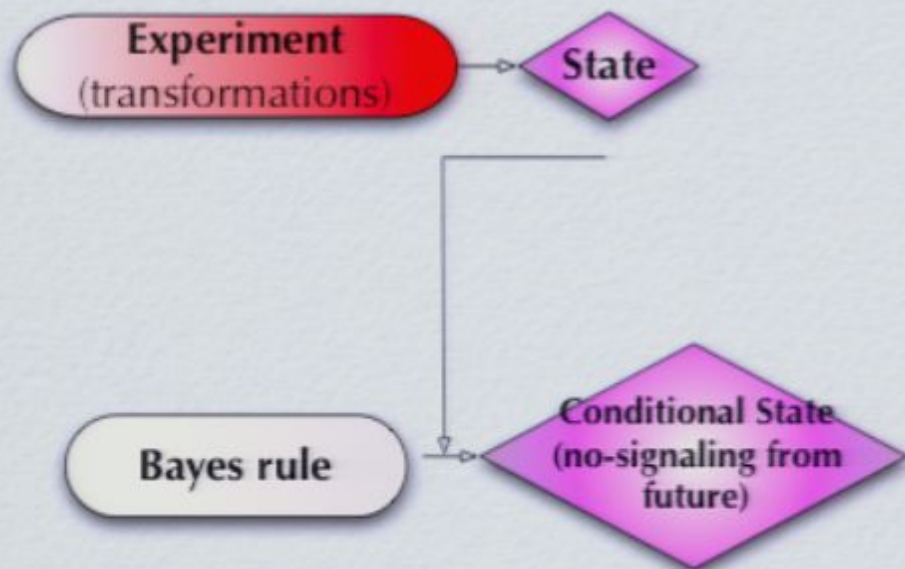
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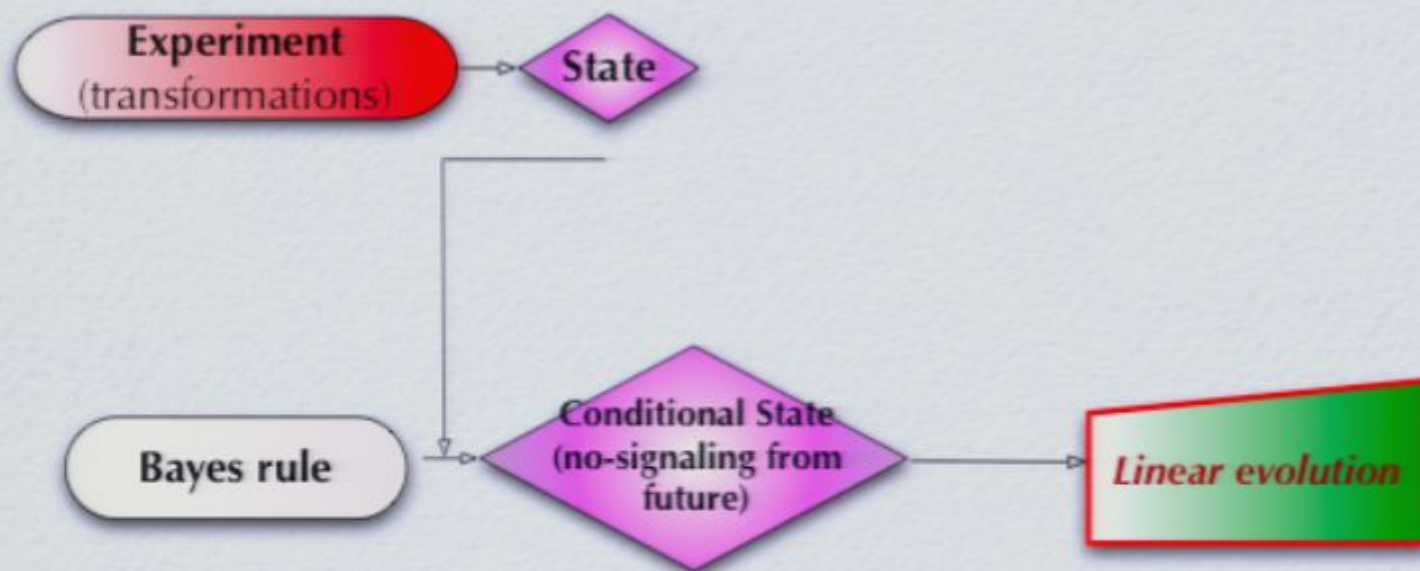
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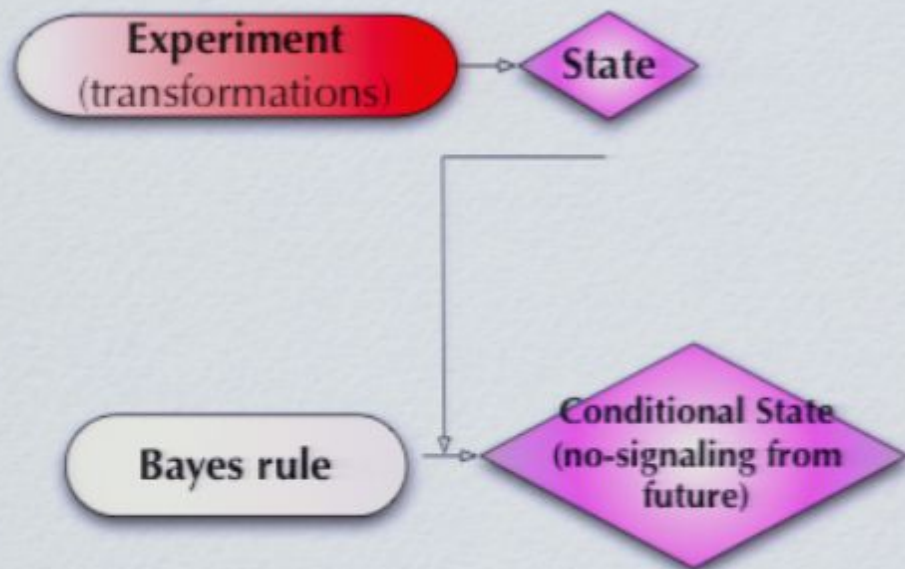
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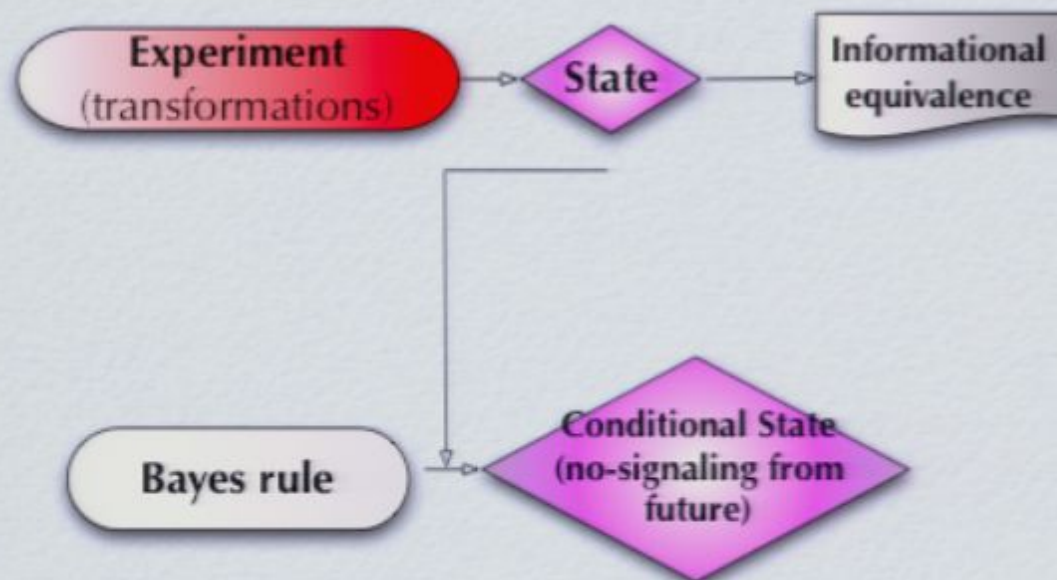
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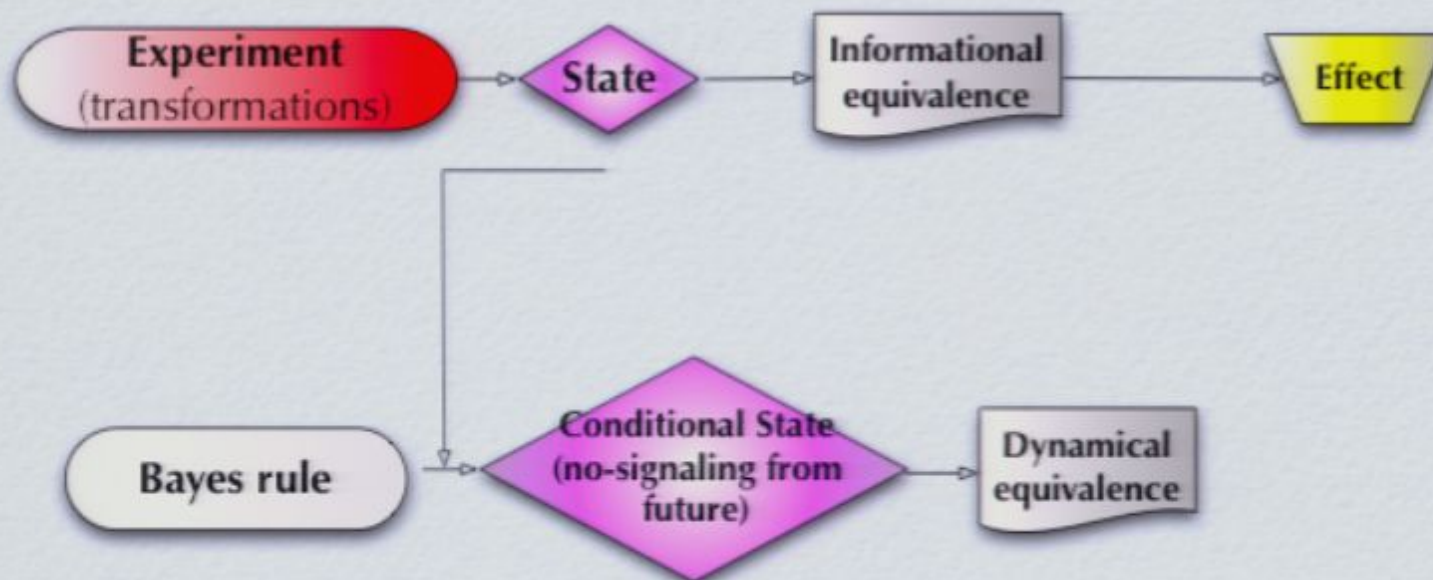
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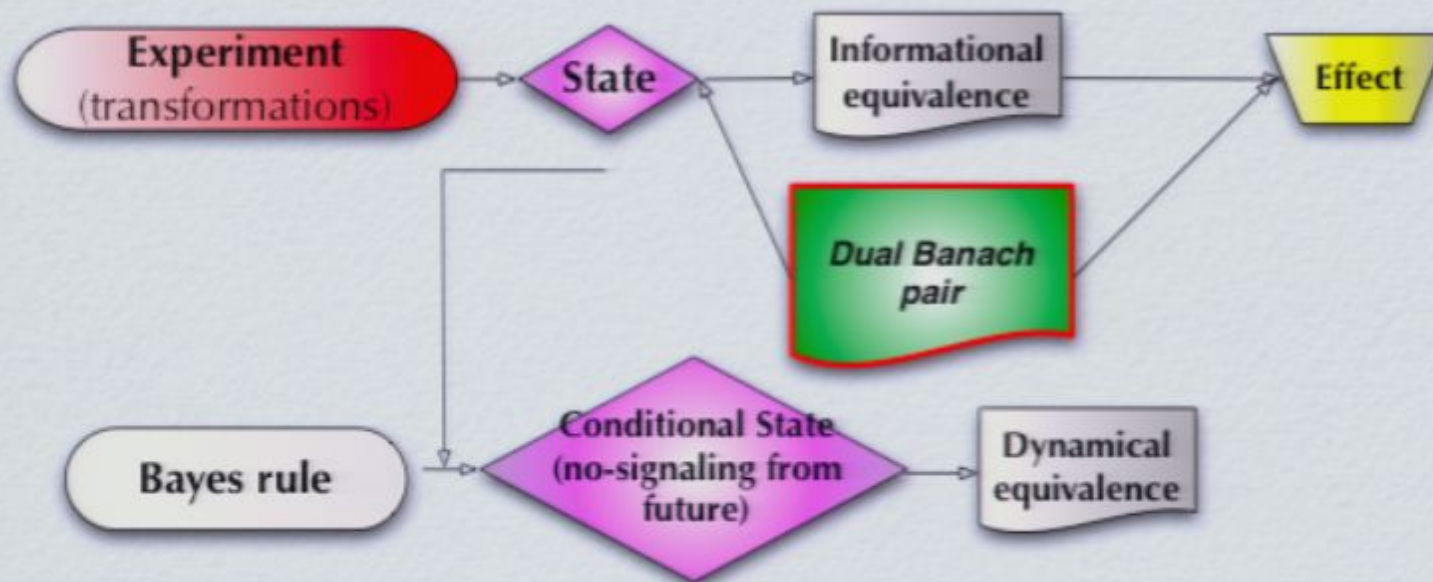
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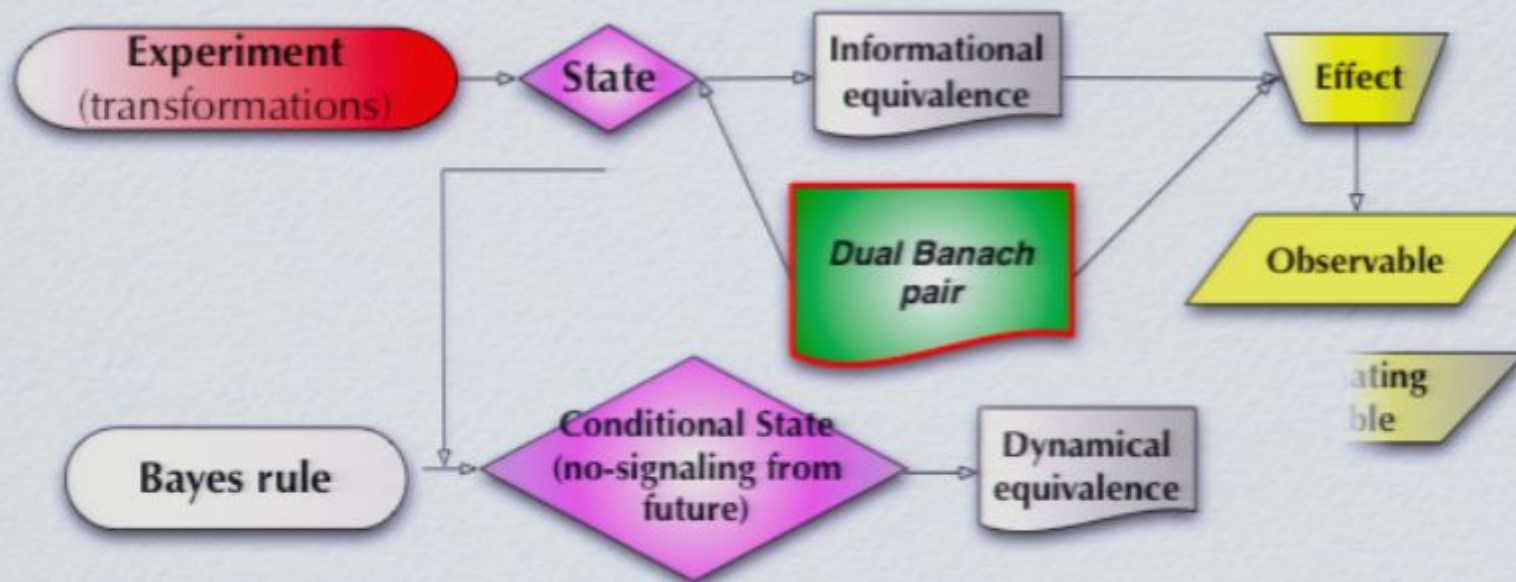
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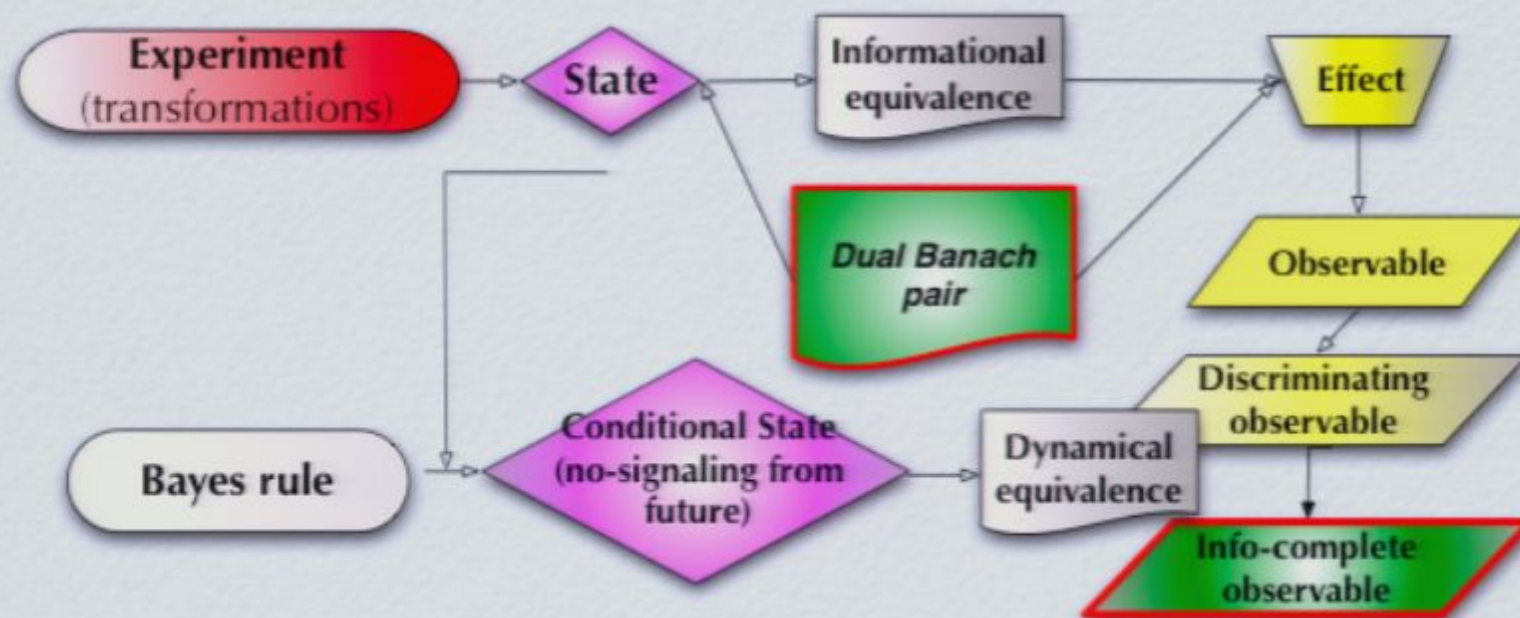
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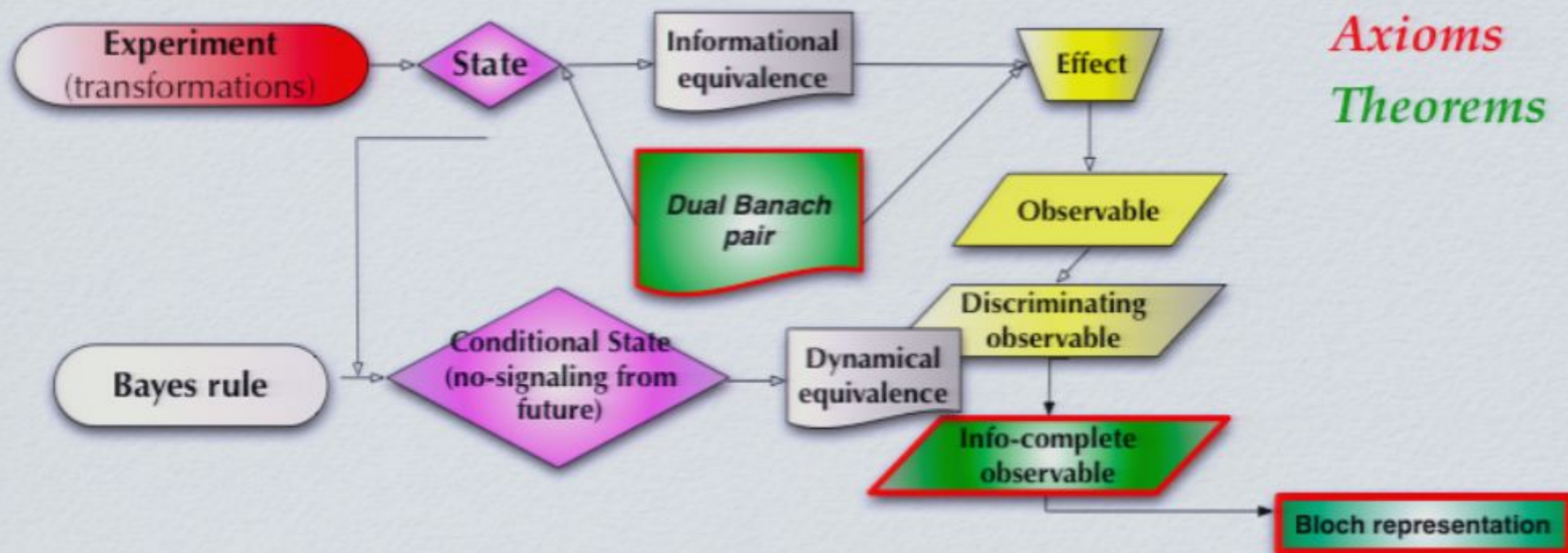
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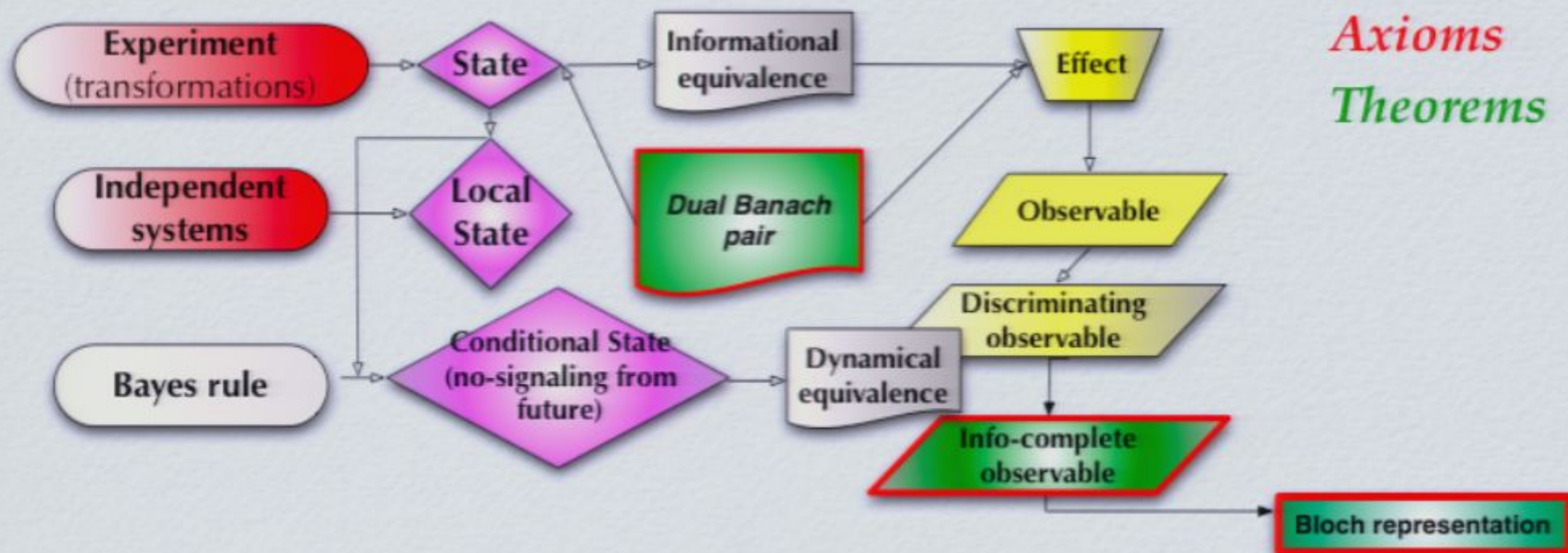


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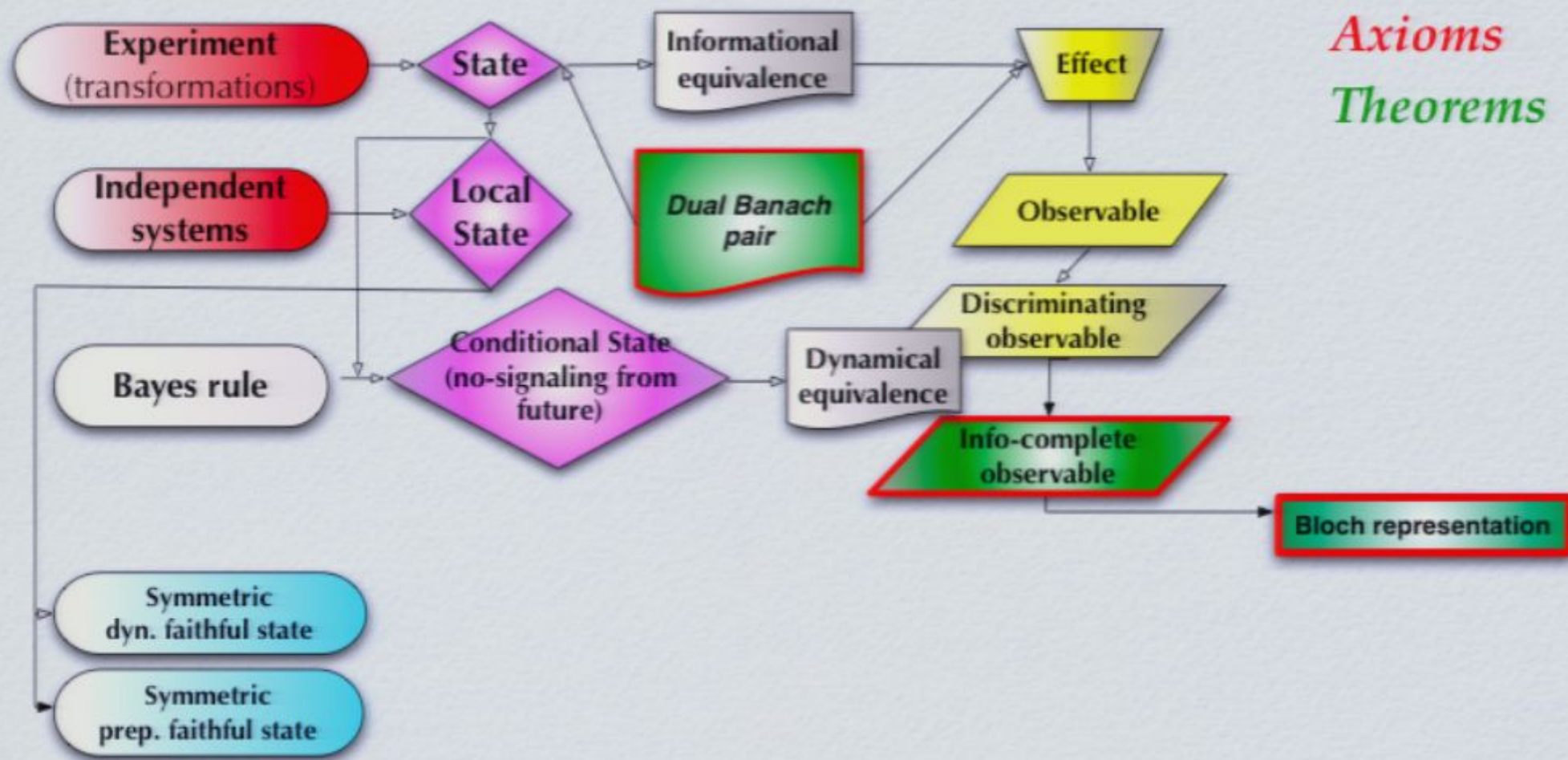


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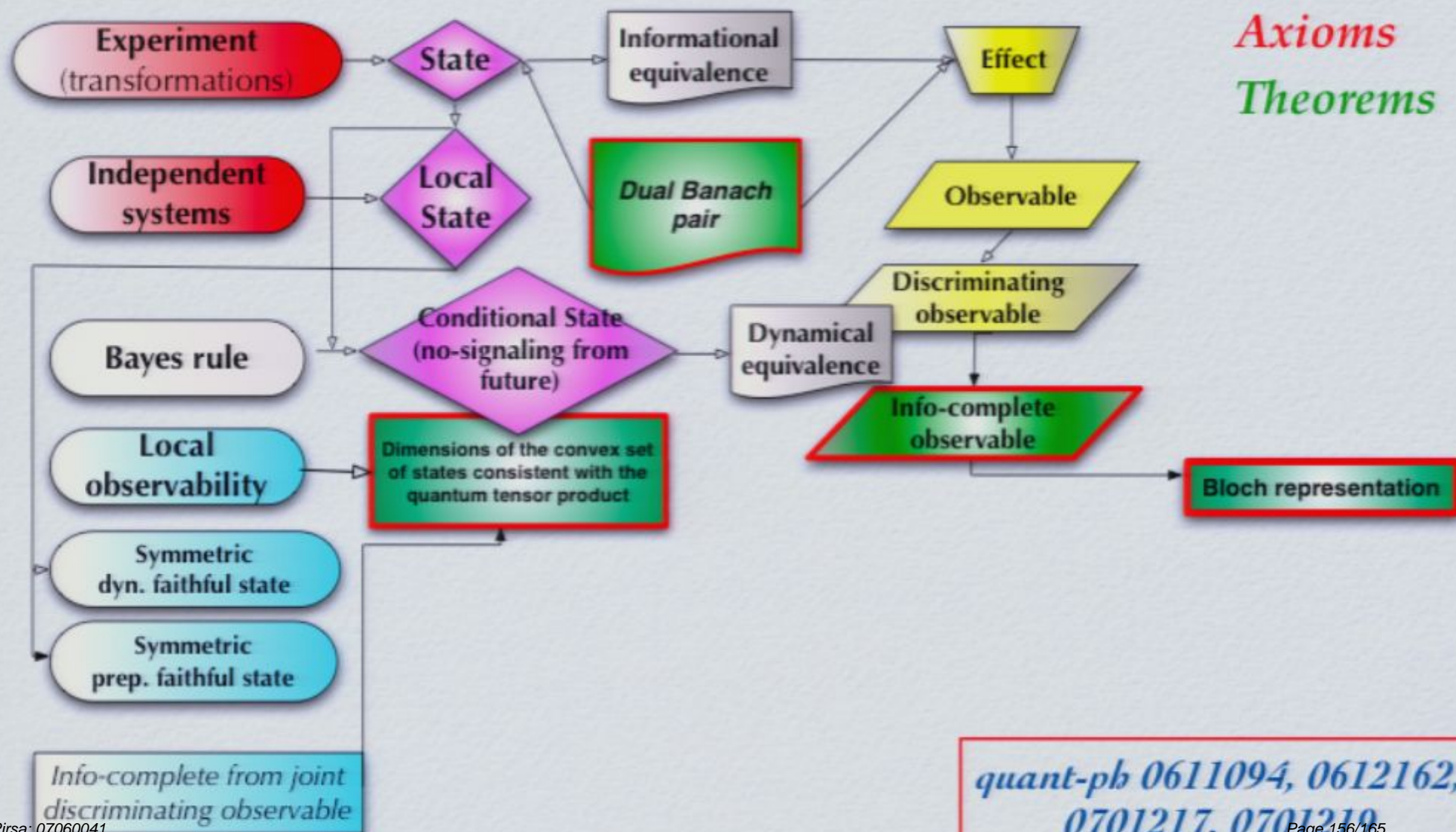


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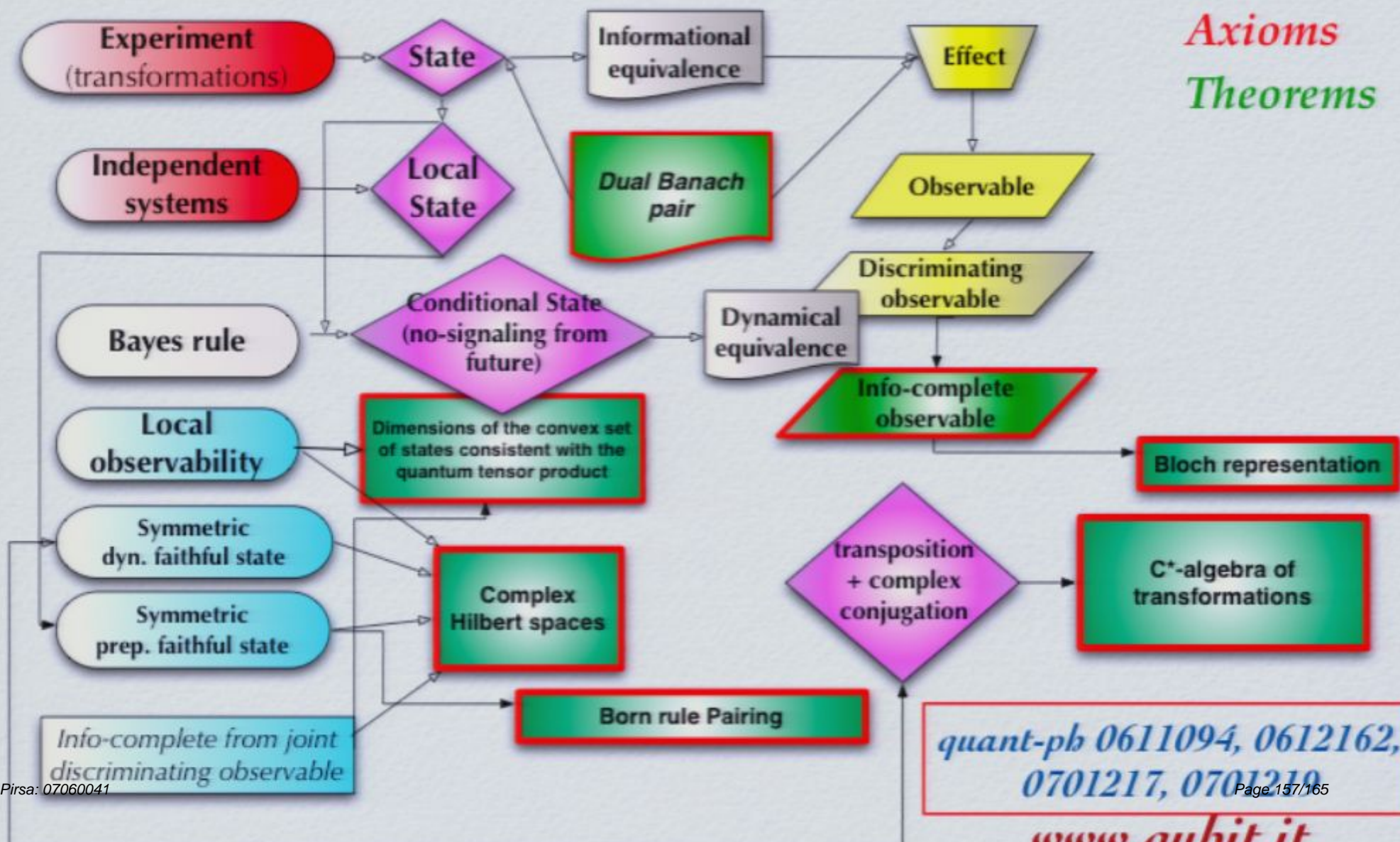


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Open problems

$\dim(\mathfrak{P}) = \infty$ Existence of ς (i.e. existence of the decomposition of the Banach space $\mathfrak{P}_{\mathbb{R}}$ into positive and negative parts for the symmetric real form Φ)

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Exploit purity of Φ

Postulates (in progress)

- **Postulate 1 (Independent systems)** *There exist **independent** systems.*
- **Postulate 2 (Symmetric faithful state)** *For every composite system made of two identical physical systems there exists a **symmetric joint state** that is both **dynamically** and **preparationally faithful**.*
- **Postulate 3 (Pure symmetric faithful state)** *If there exists a pure symmetric faithful state then we have Quantum Mechanics*

The complex conjugation

The involution ς corresponds to a generalized transformation

$$\varsigma(\underline{\mathcal{A}}) = \underline{\mathcal{A}} \circ \mathcal{Z}$$

Correspondingly the involution over transformations reads

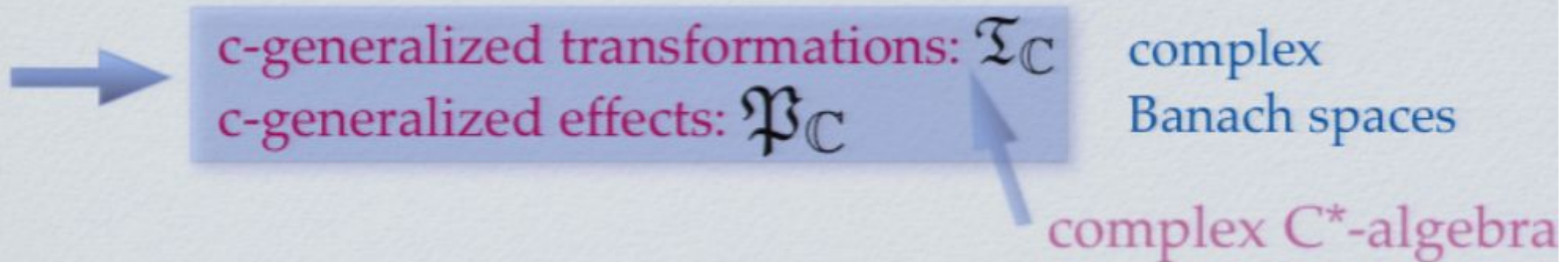
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The C^* -algebra of generalized transformations

Take complex linear combinations of generalized transformations and define $\varsigma(c\mathcal{A}) = c^* \varsigma(\mathcal{A})$ for $c \in \mathbb{C}$



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