

Title: Efficient quantum algorithms for an additive approximation of the Tutte polynomial and the q -state Potts model.

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Abstract: I will present an efficient quantum algorithm for an additive approximation of the famous Tutte polynomial of any planar graph at any point. The Tutte polynomial captures an extremely wide range of interesting combinatorial properties of graphs, including the partition function of the q -state Potts model. This provides a new class of quantum complete problems.

Our methods generalize the recent AJL algorithm for the approximation of the Jones polynomial; instead of using unitary representations, we allow non-unitarity, which seems counter intuitive in the quantum world. Significant contribution of this is a proof that non-unitary operators can be used for universal quantum computation.

An efficient quantum algorithm for approximating the Tutte polynomial and the q -state Potts model

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Previous quantum algorithms

“Shor type algorithms”

- Factoring (Shor, 1994)
- Group Order (Watrous, 2001)
- Legendre Symbol (van Dam, Hallgren, 2000)
- Pell’s Equation (Hallgren, 2002)
- Dihedral HSP (Kuperberg, 2003)
- HSP in various groups
- Hidden quadratic structures (Childs, Schulman, Vazirani, 2007)

Random-walk based algorithms

Quantum random walks (Childs, Cleve, Deotto, Farhi, Gutmann, Spielman, 2002). Provides an exponential speed-up over the classical case

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The AJL algorithm - approximating the Jones polynomial

The Jones polynomial

A knot-invariant polynomial, discovered by Jones in the mid 80's. Appears in topological field theory, statistical physics, biology, and many more



The AJL algorithm

An efficient quantum algorithm to approximate the Jones polynomial of a knot by Aharonov, Jones & Landau (2005). Based on work of Kitaev, Freedman, Larsen & Wang on Topological Quantum Computation (2002).

- Not based on quantum Fourier transform;
- Uses a mapping of the combinatorial problem into an **algebra**, and then uses a **unitary representation** of the algebra that can be implemented on a quantum computer.
- Provides an **additive** approximation of the Jones Polynomial.
- Is **quantum-complete**.

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
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This talk: extending the AJL approach for the Tutte polynomial

Main result

- An efficient approximation of the Tutte polynomial (a generalization of the Jones polynomial) at any point.
- For many points the approximation is quantum-hard.
- Many new quantum-complete problems.

Contributions

- **Breaking the unitarity barrier**
 - Quantum algorithm that uses non-unitary representations.
 - Quantum universality with non-unitary operators. 
- **The Potts model:** possible implications to the long-standing problem of approximating the partition function of the q -state Potts model.

Outline

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- 1 Background
- 2 The quantum algorithm
- 3 Universality
- 4 Summary and Conclusions



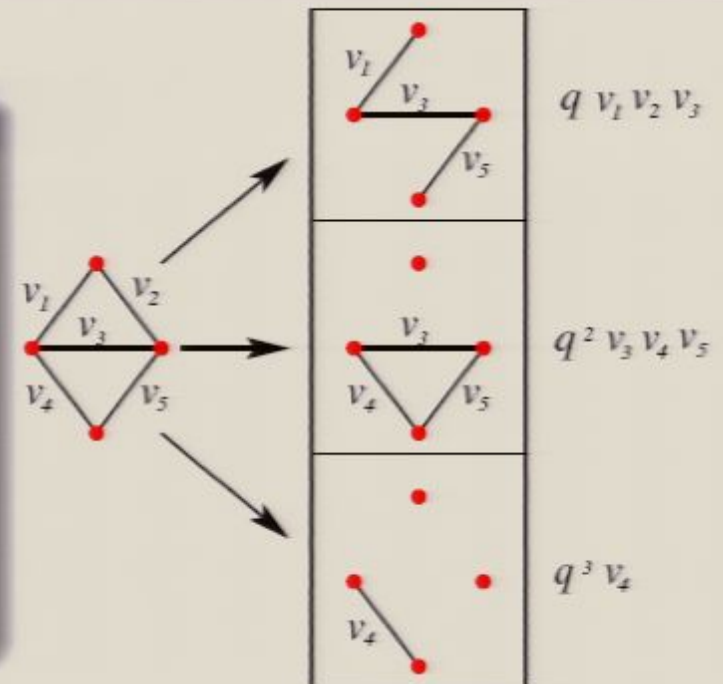
The (multivariate) Tutte polynomial

Definition

A defined for any graph $G = (V, E)$ with edge weights $\mathbf{v} \stackrel{\text{def}}{=} \{v_e | e \in E\}$

$$Z_G(q, \mathbf{v}) \stackrel{\text{def}}{=} \sum_{A \subseteq E} q^{k(A)} \prod_{e \in E} v_e$$

$k(A)$ - number of connected components in the subgraph $G = (V, A)$.



Some properties

- Encodes the reliability of a network, the number of spanning trees, chromatic polynomial and much more...
- Important special case: **the Jones polynomial**.
- For integer q and weights $v_e > -1$: the partition function of the q -state Potts model.

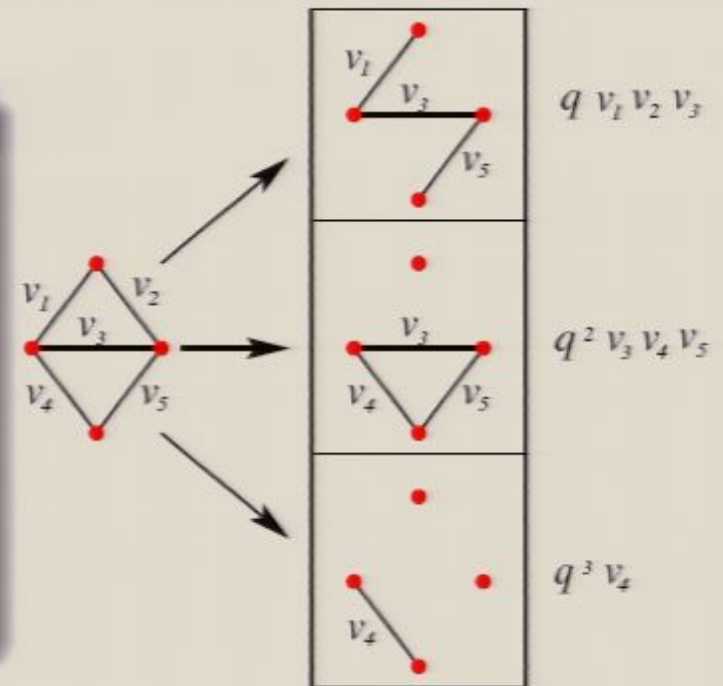
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Known complexity results (standard Tutte polynomial)

Exact evaluation

#P-hard for all (x, y) points but few trivial ones (Jaeger et al, 1990).

Multiplicative approximations (FPRAS)

- Exists for $q = 2$ with weights $v_e \geq 0$ (the Ferromagnetic Ising model) (Jerrum & Sinclair, 1993).
- NP-hard for $\sim 3/4$ of the (x, y) plane (Goldberg & Jerrum, 2006).
- For many points - including the ferromagnetic Potts model - complexity is unknown.



Can quantum computation do better?

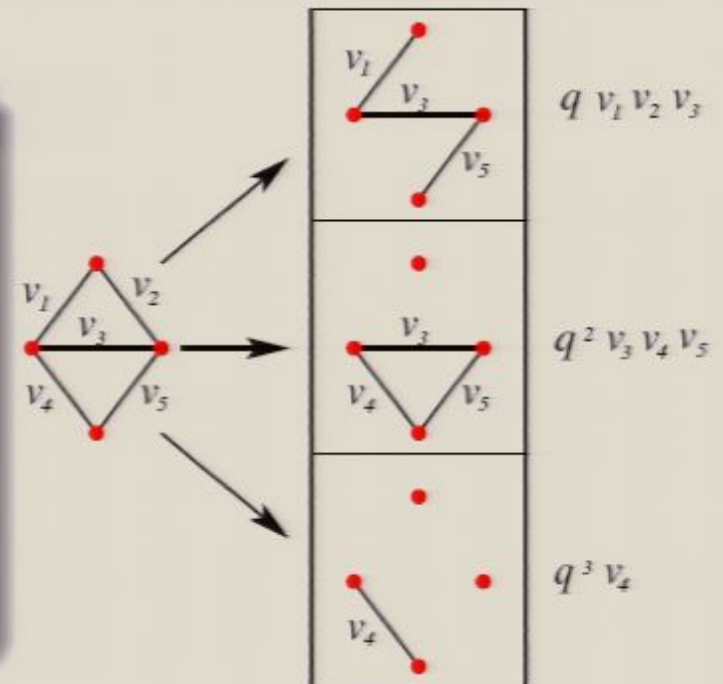
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Additive approximations - a relaxed notion of approximation

Suppose we wish to approximate the function $f(x)$ for a given x .

Multiplicative approximation

Find a y such that

$$f(x)(1 - \epsilon) \leq y \leq f(x)(1 + \epsilon), \quad \epsilon = O\left(\frac{1}{\text{poly}(n)}\right)$$

Additive approximation

Find a y such that

$$f(x) - \epsilon\Delta_x \leq y \leq f(x) + \epsilon\Delta_x, \quad \epsilon = O\left(\frac{1}{\text{poly}(n)}\right)$$

Δ_x - The approximation scale of the problem (easy to calculate).

Δ_x is crucial: the problem might become trivial if $\Delta_x \gg f(x)$

The algorithm: rough statement of results

Theorem I (rough version)

There exists an efficient quantum algorithm that provides an additive approximation for the multivariate Tutte polynomial $Z_G(q, \mathbf{v})$, i.e., outputs a number Y

$$Z_G(q, \mathbf{v}) - \frac{\Delta}{\text{poly}(n)} \leq Y \leq Z_G(q, \mathbf{v}) + \frac{\Delta}{\text{poly}(n)}$$

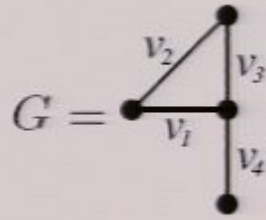
for **any** planar graph $G = (V, E)$ with **any** set of weights $\mathbf{v} = \{v_e\}$ and any complex q .



The approximation scale Δ depends on (G, q, \mathbf{v}) and on the way in which G is embedded in the plane.

Outline of the quantum algorithm

Reminder: our goal



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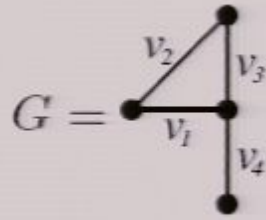
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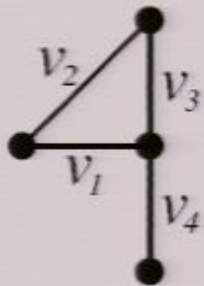
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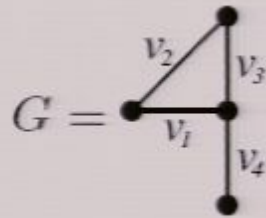
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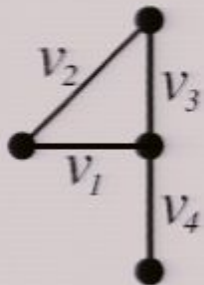
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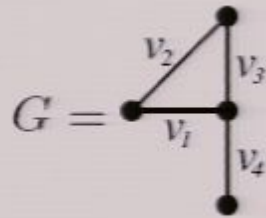
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Weighted medial
graph

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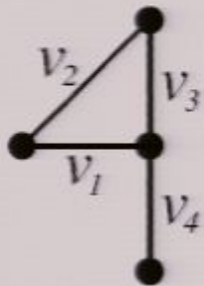
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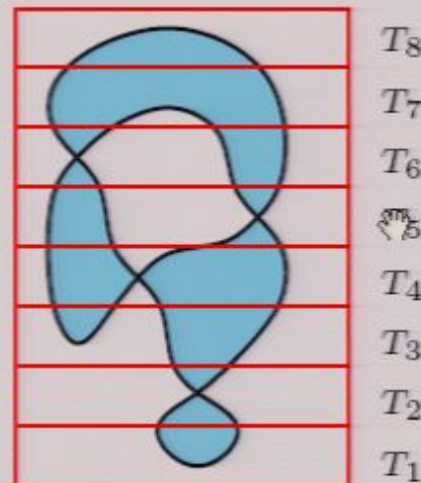
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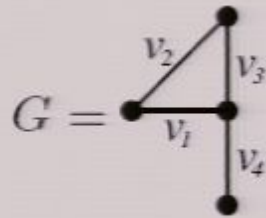
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$GTL(d)$ algebra
 $T_8 \cdots T_1$

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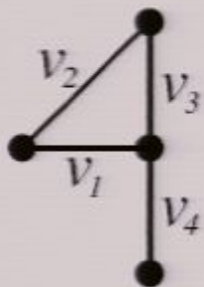
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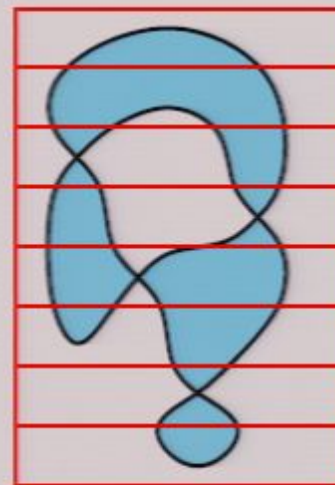
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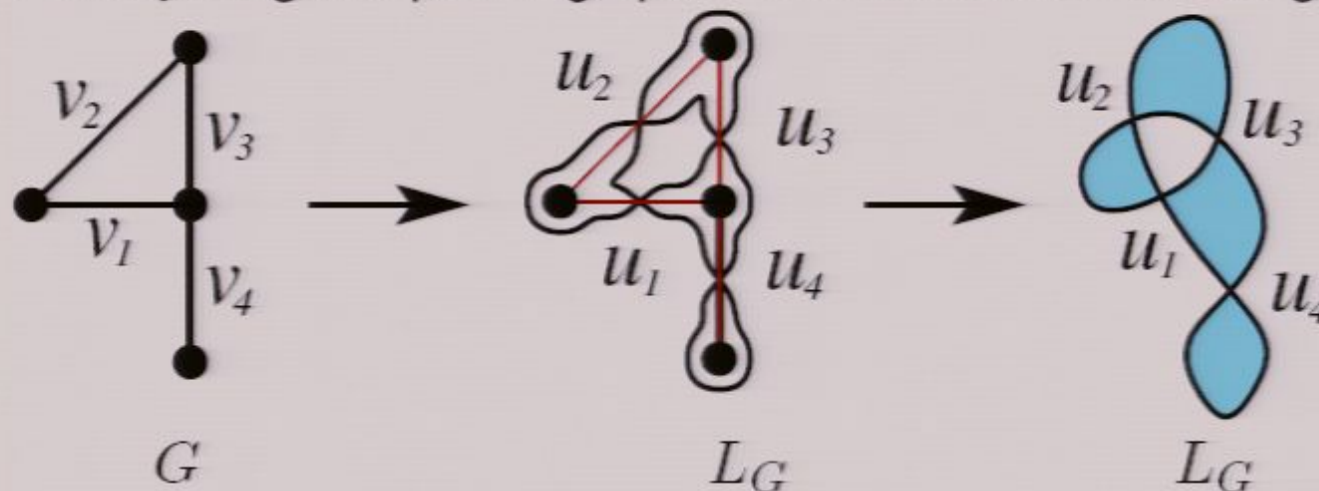
T_8	\longrightarrow	$\rho(T_8)$
T_7	\longrightarrow	$\rho(T_7)$
T_6	\longrightarrow	$\rho(T_6)$
T_5	\longrightarrow	$\rho(T_5)$
T_4	\longrightarrow	$\rho(T_4)$
T_3	\longrightarrow	$\rho(T_3)$
T_2	\longrightarrow	$\rho(T_2)$
T_1	\longrightarrow	$\rho(T_1)$

Matrix
Representation

From graphs to medial graphs

The Medial graph L_G

For any weighted planar graph G we can define a medial graph L_G .

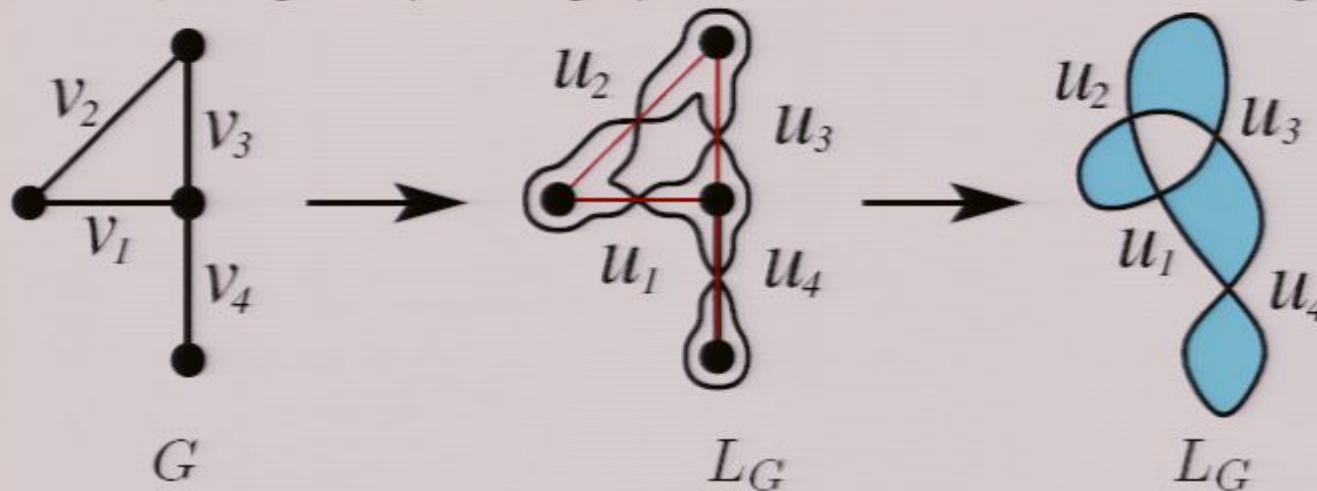


- L_G is a 4-regular planar graph.
- Edges in G are transformed into vertices of L_G , also called **crossings**.
- Edges weights v_i are transformed to crossing weights u_i .
- **Blue-and-White coloring:**
 - 1 Paint the external facet in white.
 - 2 Proceed by painting adjacent facets in opposite colors.
 - 3 \implies Blue regions correspond to vertices of G .

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The weighted Kauffman bracket (equivalent to the Tutte polynomial)

$L_G \longrightarrow$ **Kauffman bracket** $\langle L_G \rangle(d, \mathbf{u})$.

The Kauffman Bracket

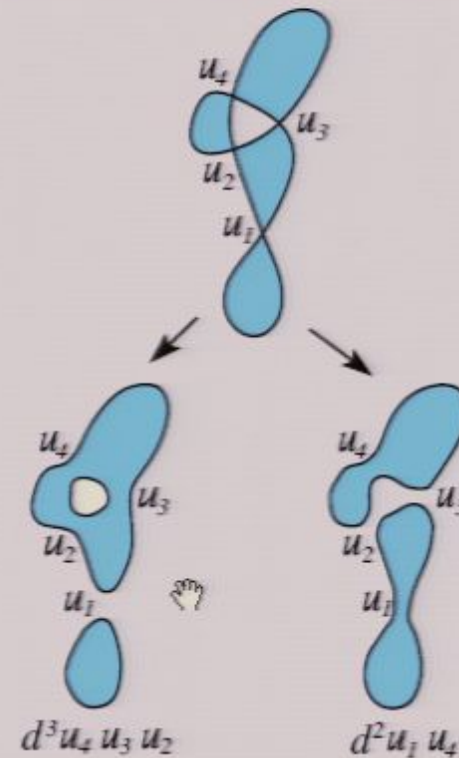
A polynomial that can be defined for any weighted medial graph L_G by **opening the crossings** in two ways:



$$\langle L_G \rangle(d, \mathbf{u}) \stackrel{\text{def}}{=} \sum_{\sigma} d^{\#\ell} \prod_{e \in \sigma} u_e$$

- d - a complex number
- σ An opening of the crossings.
 $e \in \sigma \Leftrightarrow$ shaded area connected.
- $\#\ell$ - Number of resulting loops

Opening a crossing



Relation to the Tutte polynomial

$$\langle L_G \rangle(d, \mathbf{u}) = d^{-|V|} Z_G(d^2, d\mathbf{u})$$

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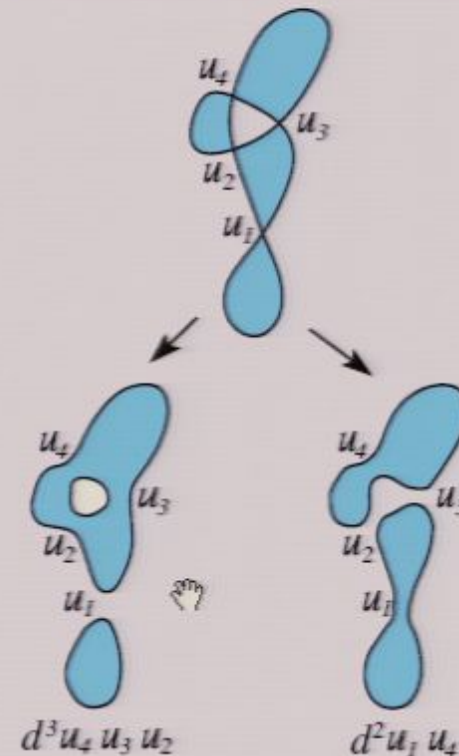
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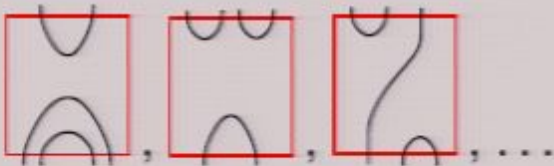
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Moving to the language of algebras: the (generalized) Temperley-Lieb algebra

Definition: the $GTL(d)$ algebra

The $GTL(d)$ is generated by tangle diagrams, with no crossings and no loops. **Number of in-pegs may be different from number of out-pegs.**



- A general element is the weighted sum of such elements:

$$\alpha \left[\begin{array}{c} \cup \\ \cap \end{array} \right] + \beta \left[\begin{array}{c} \cap \\ \cup \end{array} \right] + \dots$$

- **Product Rule:** simply place one on top of the other (when no. of strands match).

- Loops are taken out as d factors ($d = \text{loop value}$).

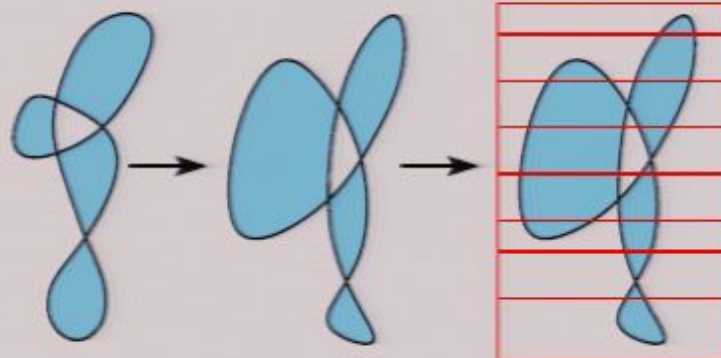
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From the medial graph to the $GTL(d)$ algebra

$$\Psi : L_G \rightarrow GTL(d)$$

We present L_G as a product of elementary diagrams:

- A Minima
- A Maxima
- A Crossing



All the operators are local!

Embedding an elementary diagram in $GTL(d)$

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- A Maxima:
- A Crossing: $\rightarrow u +$

Corollary



$$\Psi(L_G) = \langle L_G \rangle \mathcal{I}$$

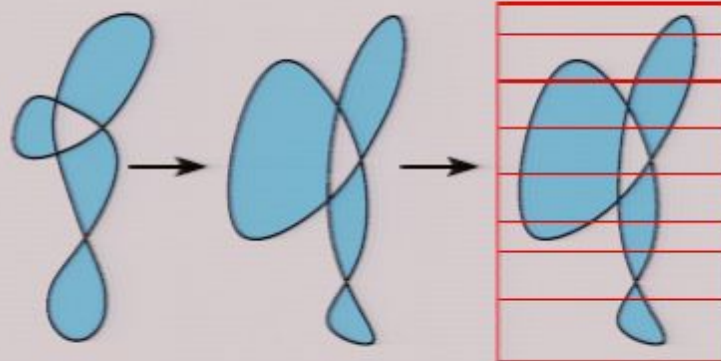
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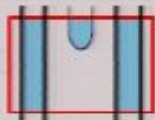
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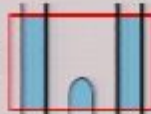
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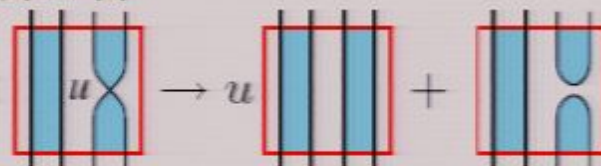
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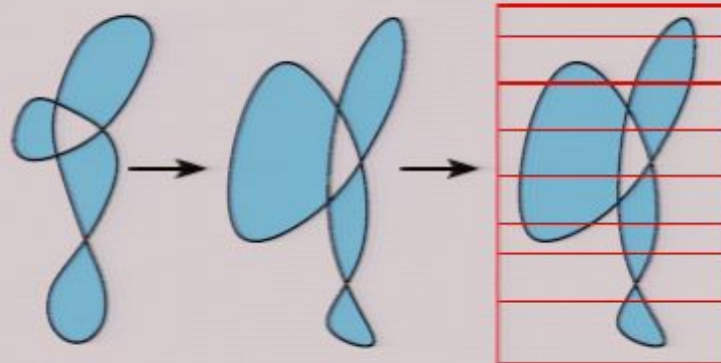
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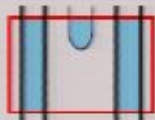
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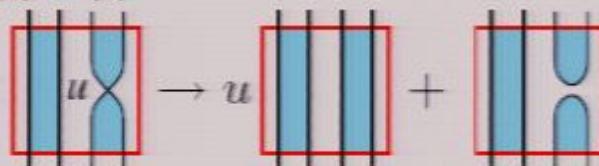
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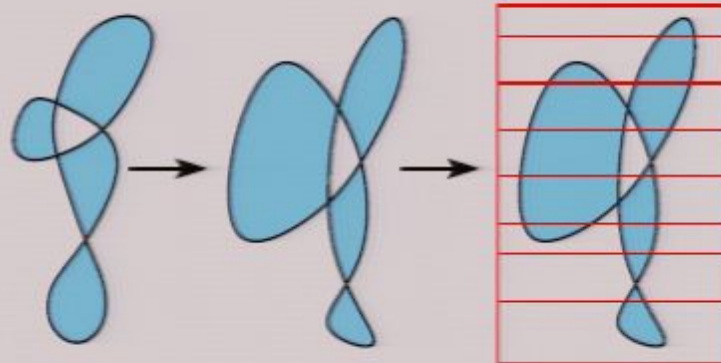
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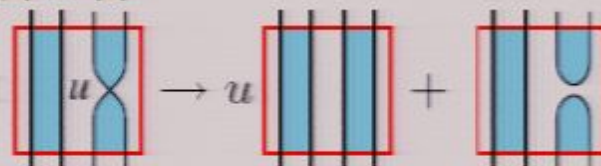
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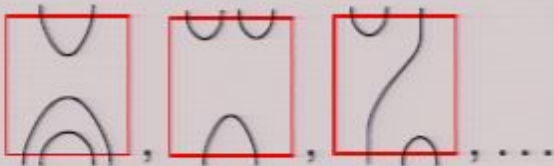
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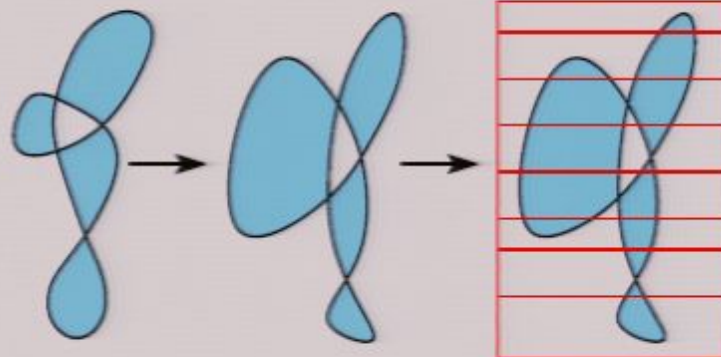
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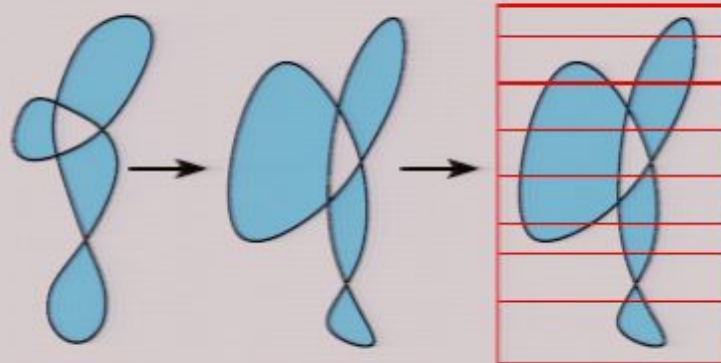
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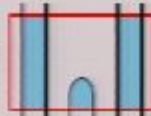
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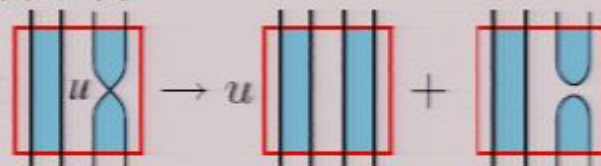
- A Minima:



- A Maxima:



- A Crossing:



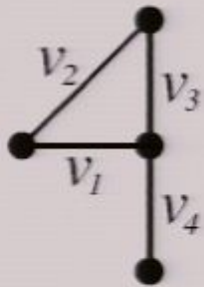
Corollary

$$\Psi(L_G) = \langle L_G \rangle \mathcal{I}$$

We want to find the prefactor before the identity!

Intermediate summary

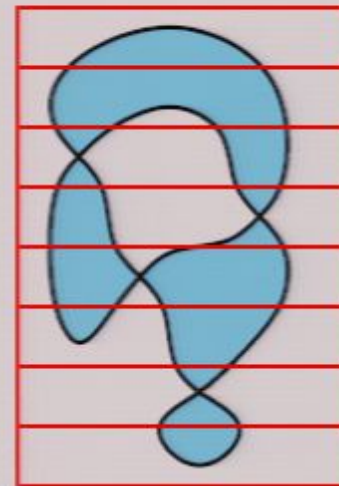
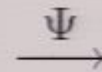
Outline of the quantum algorithm so far



Weighted planar graph $G = (V, E)$



Weighted medial graph



$GTL(d)$ algebra
 $\Psi(L_G) = \langle L_G \rangle \mathcal{I}$.

T_8	\longrightarrow	$\rho(T_8)$
T_7	\longrightarrow	$\rho(T_7)$
T_6	\longrightarrow	$\rho(T_6)$
T_5	\longrightarrow	$\rho(T_5)$
T_4	\longrightarrow	$\rho(T_4)$
T_3	\longrightarrow	$\rho(T_3)$
T_2	\longrightarrow	$\rho(T_2)$
T_1	\longrightarrow	$\rho(T_1)$

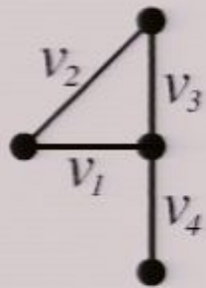
Move to matrix representation of $GTL(d)$

We need a representation that can be applied efficiently. We then apply

$$\rho(T_8) \cdot \dots \cdot \rho(T_1) = \rho(T_8 \cdot \dots \cdot T_1) = \rho(\langle L_G \rangle \mathcal{I}) = \langle L_G \rangle \mathbf{1}$$

Intermediate summary

Outline of the quantum algorithm so far

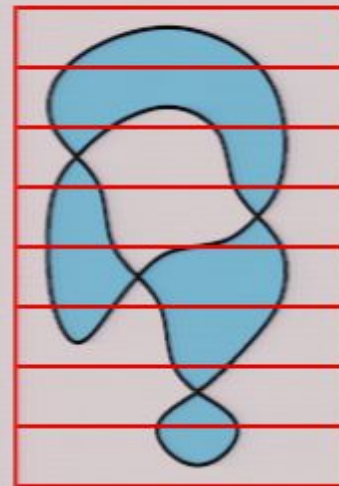


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$\Psi \rightarrow$



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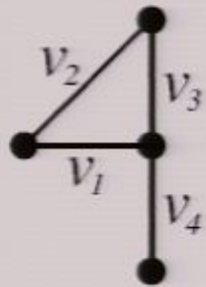
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Intermediate summary

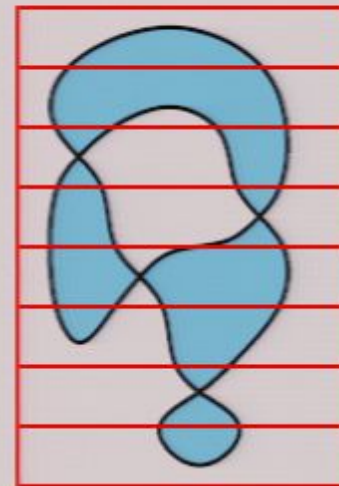
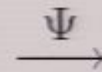
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Move to matrix
representation
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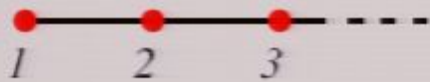
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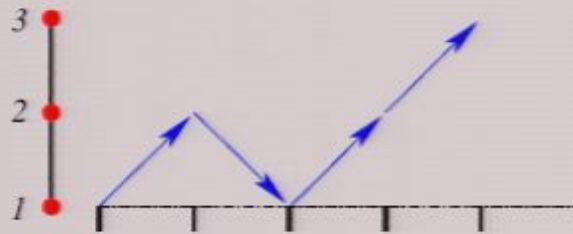
The Path-Model representation

The $\{H_n\}$ spaces

Spanned by n -steps paths on the line, starting from 1.



$|12123\rangle \in H_4$ is:



A diagram with 0 lower strands and 0 upper strands is simply a linear transformation from $H_0 = \text{span}\{|1\rangle\}$ into itself - a simple scalar multiplication.

How does a diagram operate on paths ?

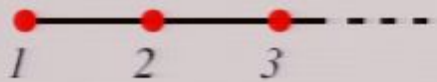
- Paint the **regions** of the diagram.
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- Example:

$$|12323\rangle = \alpha |1212323\rangle + \beta |1212343\rangle$$

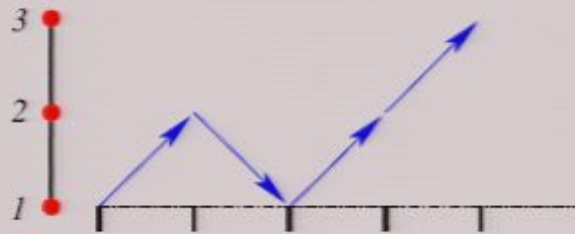
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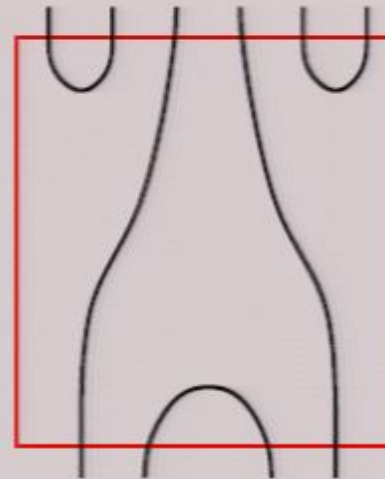


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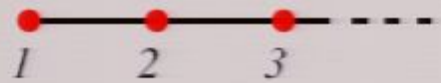
• Example:

$$|123\rangle \rightarrow \frac{1}{2} |12123\rangle + \frac{1}{2} |132\rangle$$

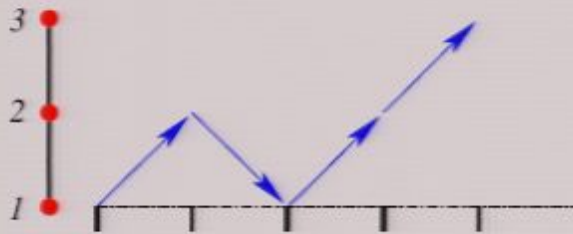
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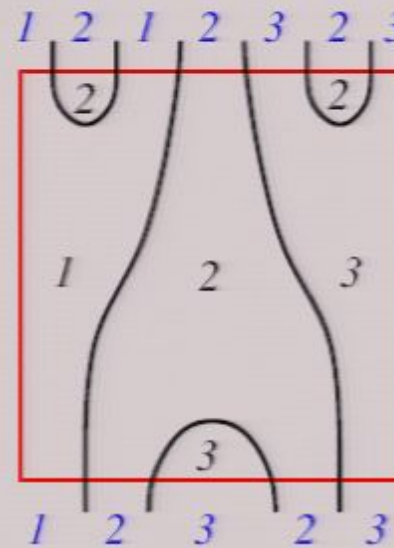


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 \uparrow
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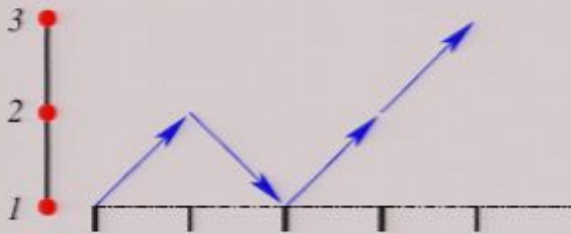
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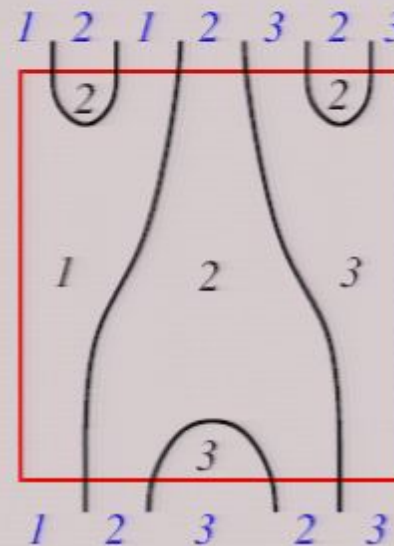


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Calculating $\langle L_G \rangle$ with the path-model representation

Recall

$$\rho(T_N) \cdot \dots \cdot \rho(T_1) = \rho(T_N \cdot \dots \cdot T_1) = \langle L_G \rangle \mathbf{1}$$

Calculating $\langle L_G \rangle$ with matrices



$$|1\rangle \in H_0$$

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$$|\psi_1\rangle = \rho(T_1) |1\rangle \in H_2$$

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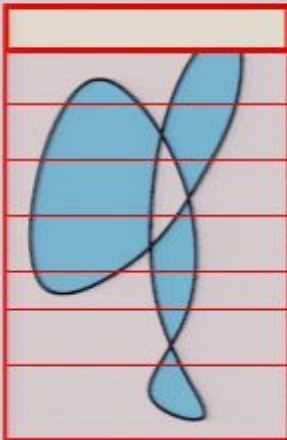
$$|\psi_2\rangle = \rho(T_2)\rho(T_1) |1\rangle \in H_2$$

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Calculating $\langle L_G \rangle$ with matrices



$$|\psi_{N-1}\rangle = \rho(T_{N-1}) \cdots \rho(T_1) |1\rangle$$

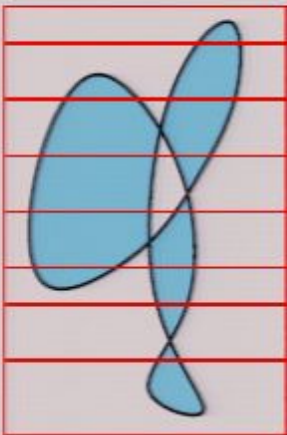


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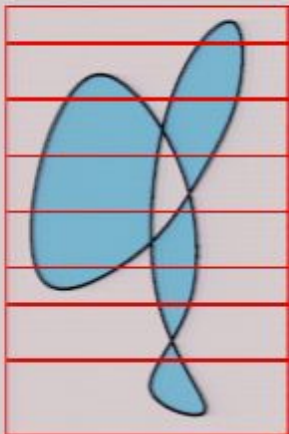
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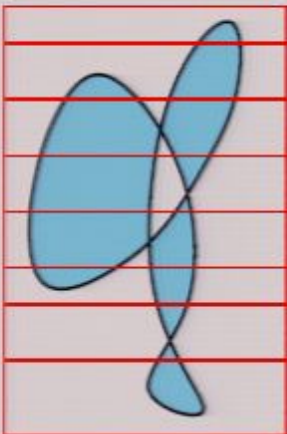
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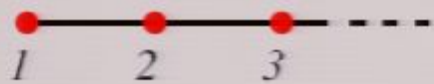
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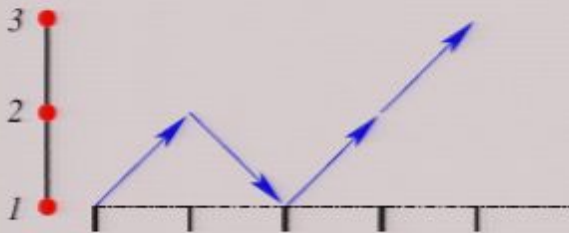
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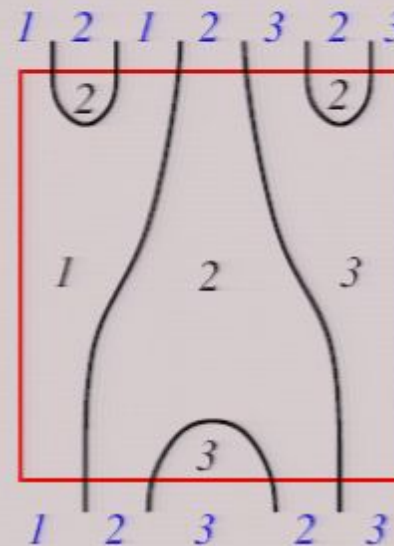


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How does a diagram operate on paths ?



H_6
↑
 H_4

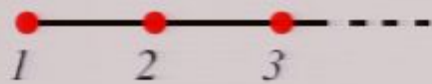
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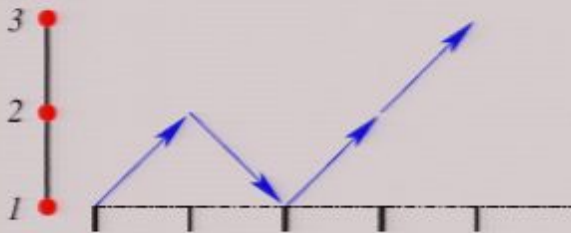
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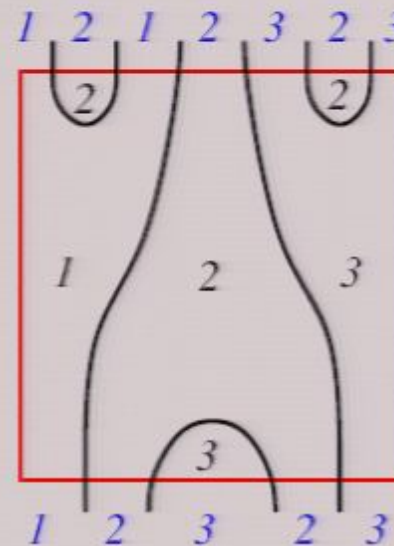


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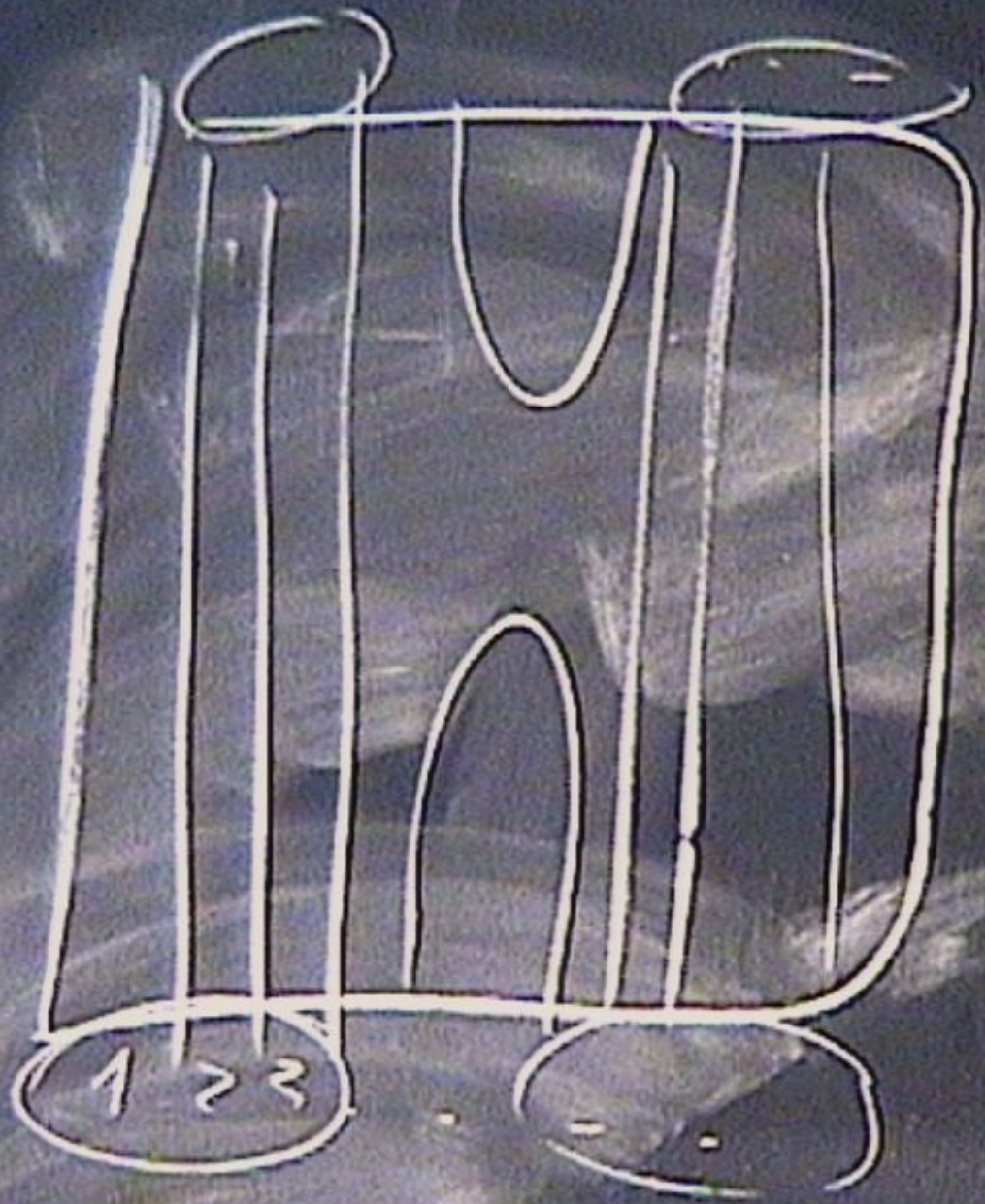
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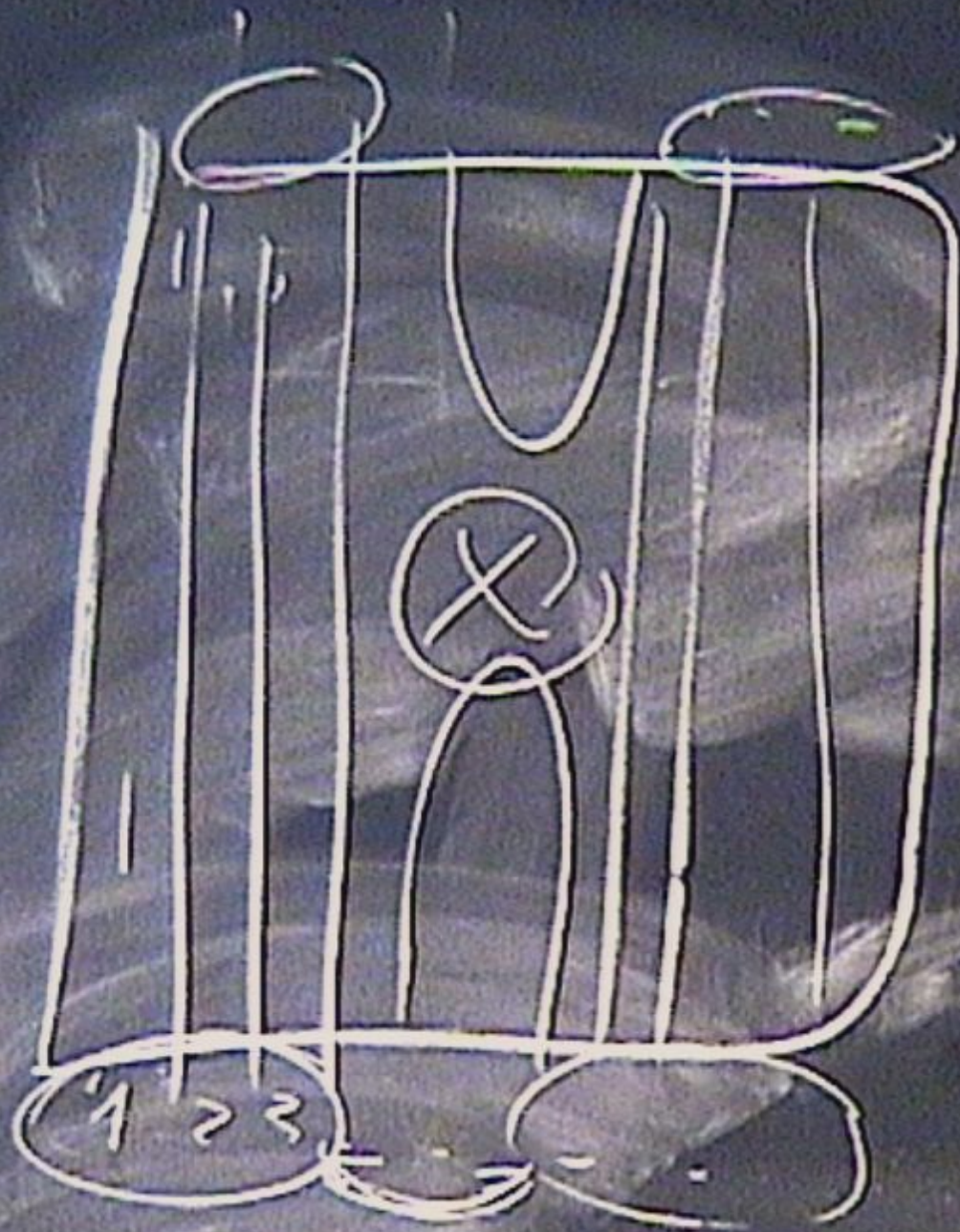
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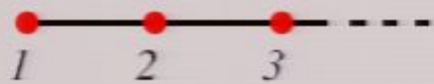
ϕ 2
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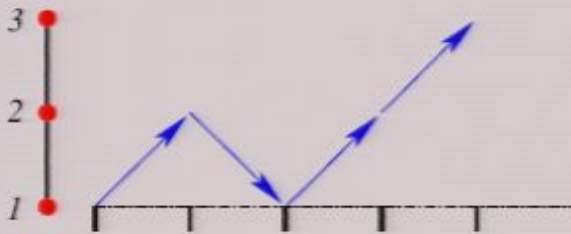
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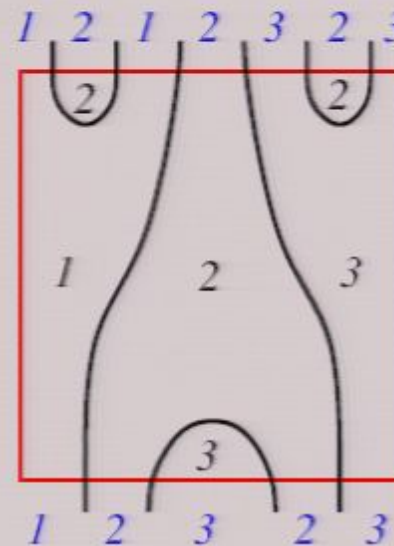


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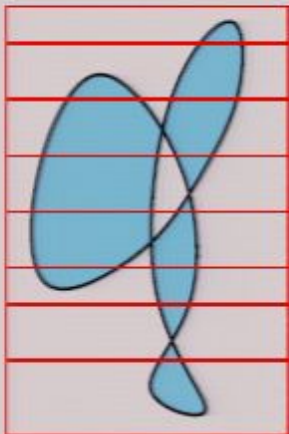
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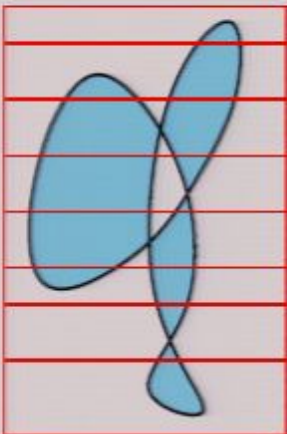
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Implementing non-unitary elementary operators

Crossing operator

- Reminder: The **Crossing** operator is a $(u + v) \times (u + v)$ (not necessarily unitary) square matrix $M : H_n \rightarrow H_n$.

See also the [Feynman Decomposition](#) slide.



Minima & Maxima operators

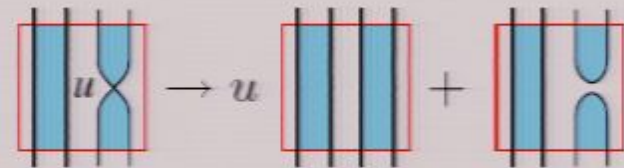
- A **Minima** ($H_n \rightarrow H_{n+2}$) is implemented by adding qubits.
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We use the **Polar Decomposition** $M = PU$



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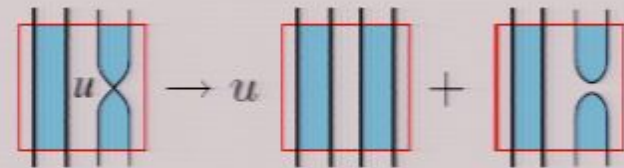
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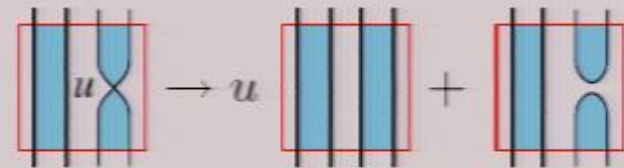
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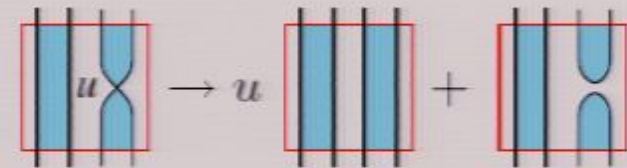
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$\Rightarrow U$ is unitary - easy to implement (because of locality).



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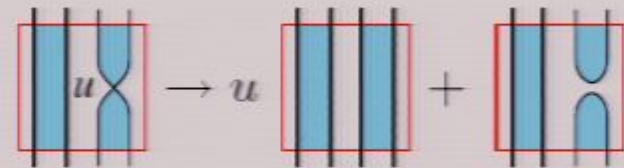
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$$\Rightarrow P = \text{diag}(r_1, r_2, \dots,)$$



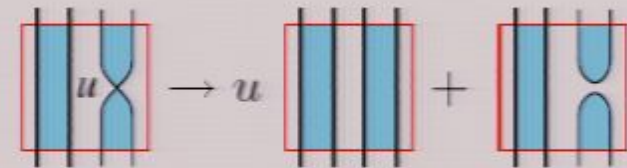
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$$\Rightarrow P = \text{diag}(r_1, r_2, \dots)$$

We implement $\frac{P}{\|P\|} = \text{diag}(1, r_2/r_1, r_3/r_1, \dots)$ by pushing the weights to an **ancilla qubit**

$$|i\rangle \otimes |0\rangle \rightarrow \frac{r_i}{r_1} |i\rangle \otimes |0\rangle + \sqrt{1 - (r_i/r_1)^2} |r_i\rangle \otimes |1\rangle$$



and book-keeping the norm $\|P\| = \|M\|$.

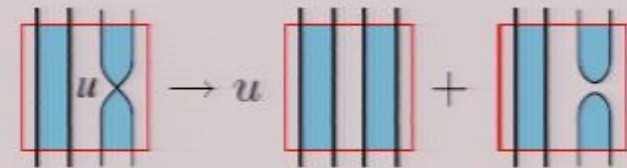
Minima & Maxima operators

- A **Minima** ($H_n \rightarrow H_{n+2}$) is implemented by adding qubits.
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Implementing non-unitary elementary operators

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Summary of quantum algorithm

- The computation starts with the state $|1\rangle$ (a path with 0 steps).
- Apply the elementary operators: minimas, maximas, crossings.
- In the end, we get

$$|\psi_N\rangle = \frac{\langle L_G \rangle}{\|\rho(T_1)\| \cdot \dots \cdot \|\rho(T_N)\|} |1\rangle \otimes \overbrace{|0\rangle \otimes \dots \otimes |0\rangle}^{\text{ancilla qubits}} + \dots$$

- Estimate the inner product $\langle 1 \otimes 0 \otimes \dots \otimes 0 | \psi_N \rangle$ using the **Hadamard test**, Obtain:

$$\frac{\langle L_G \rangle}{\|\rho(T_1)\| \cdot \dots \cdot \|\rho(T_N)\|} \pm 1/\text{poly}(n)$$

The approximation scale

$$\Delta_{alg} = \|\rho(T_1)\| \cdot \dots \cdot \|\rho(T_N)\|$$

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Conclusion of algorithmic part

Theorem 1

There exists an efficient quantum algorithm that provides an additive approximation for the multivariate Tutte polynomial $Z_G(q, \mathbf{v})$ of **any** planar graph $G = (V, E)$ with **any** set of weights $\mathbf{v} = \{v_e\}$ and any complex q .

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- We prove non-triviality by proving BQP-hardness.
- This is already known for the special case of the Jones polynomial \sim the unitary case.
- **Main difficulty:** how to prove universality for the non-unitary case?
- We generalize the universality proof of the Jones polynomial of Aharonov & Arad (2006), following the work of Freedman, Kitaev, Larsen & Wang.
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Welcome to the dark side of the talk...



Quantum Universality

Universality

Given a quantum circuit $U = U_L \cdot \dots \cdot U_1$ over n qubits, we need to decide whether $|\langle 0^{\otimes n} | U | 0^{\otimes n} \rangle|^2 \leq 1/3$ or $|\langle 0^{\otimes n} | U | 0^{\otimes n} \rangle|^2 \geq 2/3$.

We will construct a weighted graph $G = (V, E)$ such that

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Outline of the universality proof

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Finding G

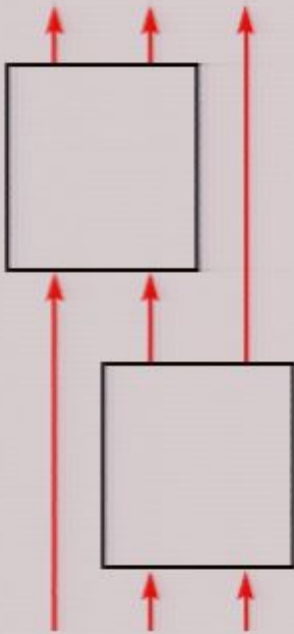


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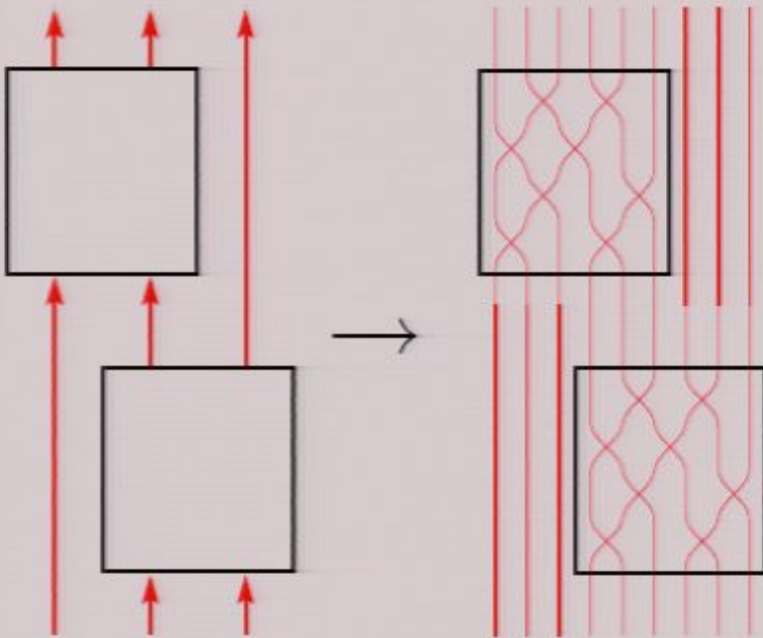
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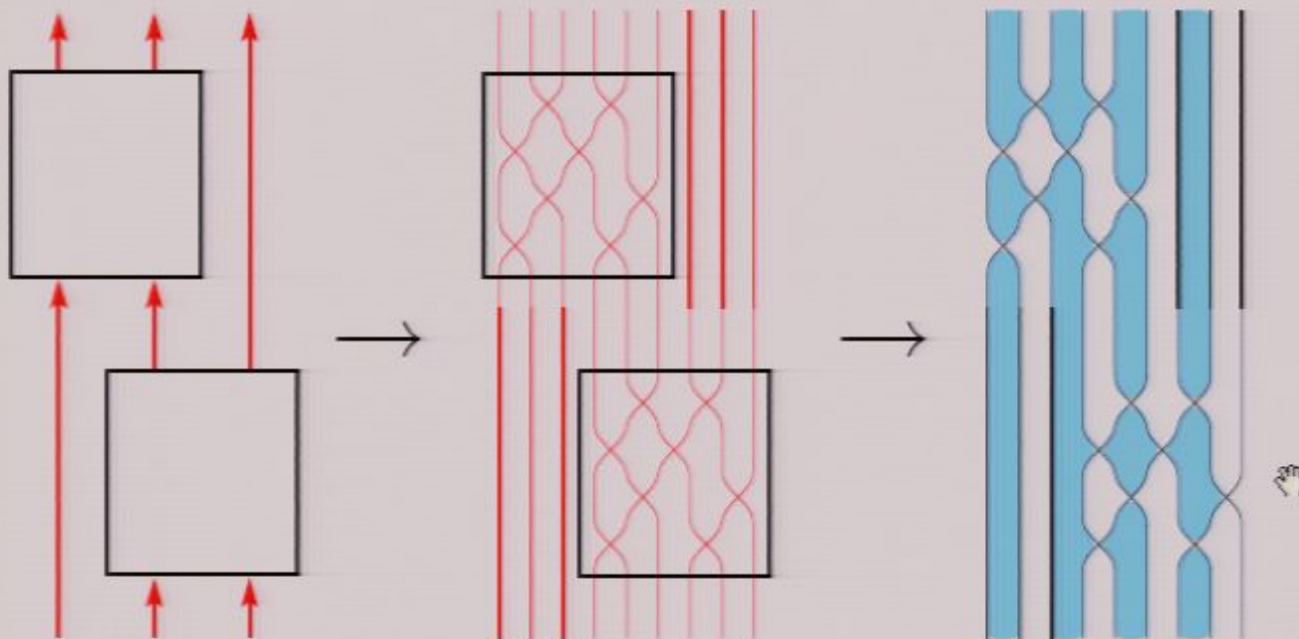
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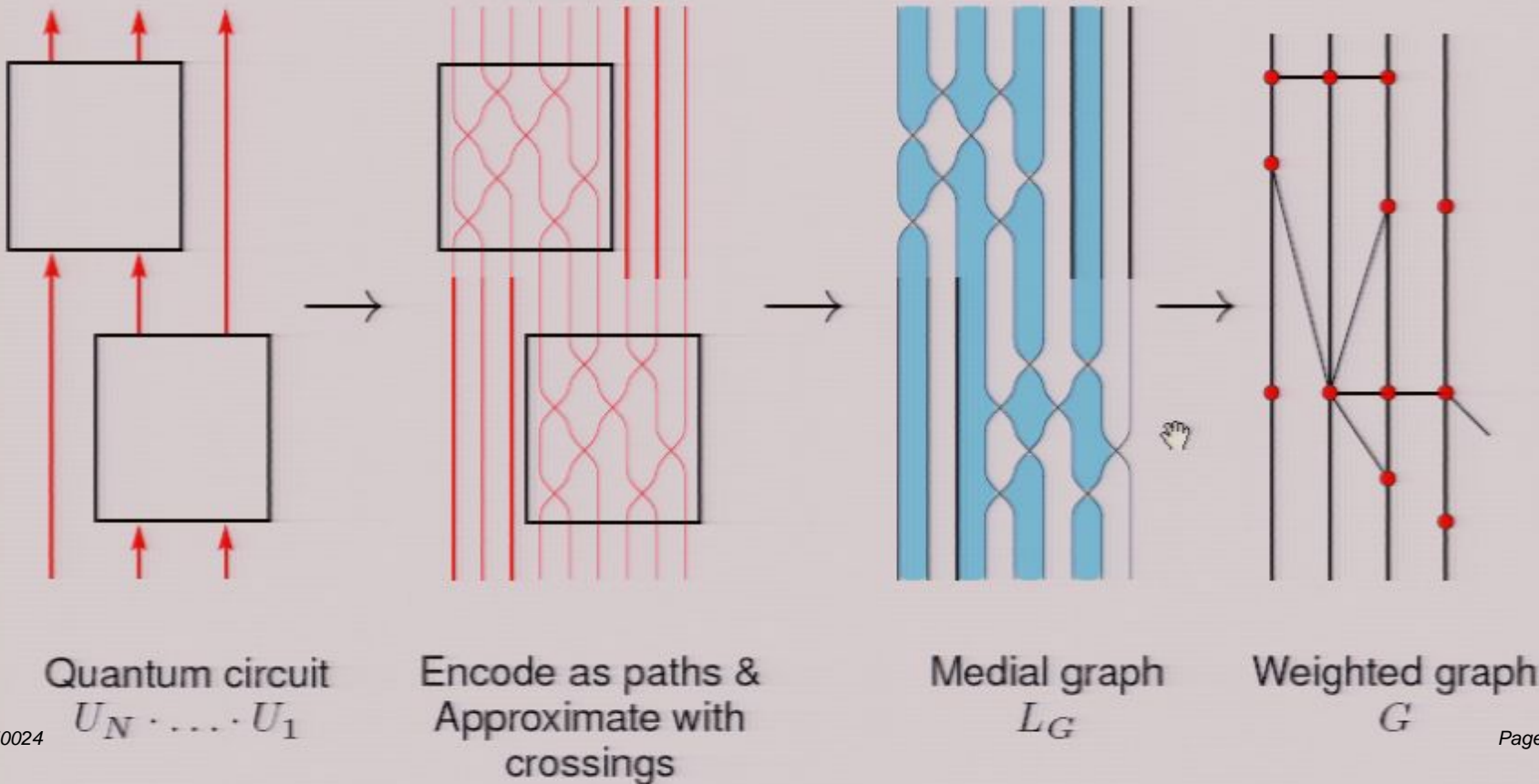
Medial graph
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The 4-steps encoding (Kitaev, Wocjan & Yard)

Encoding qubits

Every qubit is mapped into a 4-steps path that starts at 1 and ends at 1:

$$\begin{aligned}
 |\underline{0}\rangle &\stackrel{\text{def}}{=} |12121\rangle, & \begin{array}{c} \text{---} \\ \diagup \quad \diagdown \\ \text{---} \end{array} & \begin{array}{c} \text{---} \\ \diagdown \quad \diagup \\ \text{---} \end{array} \\
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String of n qubits is mapped to a $4n$ -path in the subspace $L_{4n} \subseteq H_{4n}$.

The path space of 2 qubits - L_8

$$|\underline{00}\rangle \quad |\underline{01}\rangle \quad |\underline{10}\rangle \quad |\underline{11}\rangle$$

A 2-qubit gate is translated to a unitary operator over L_8

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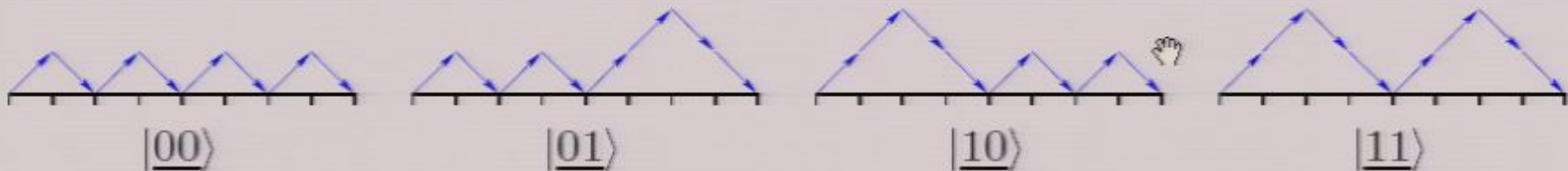
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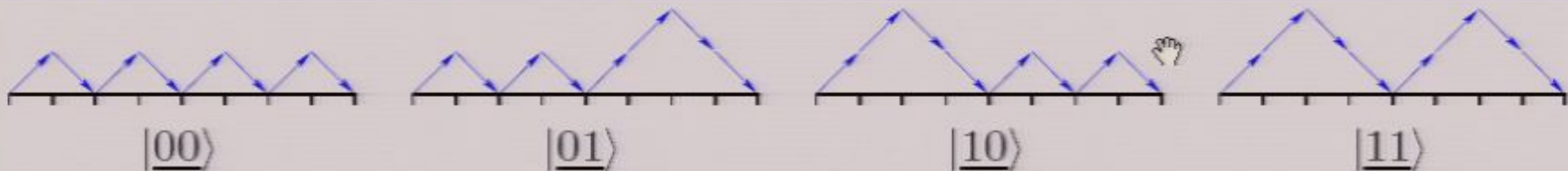
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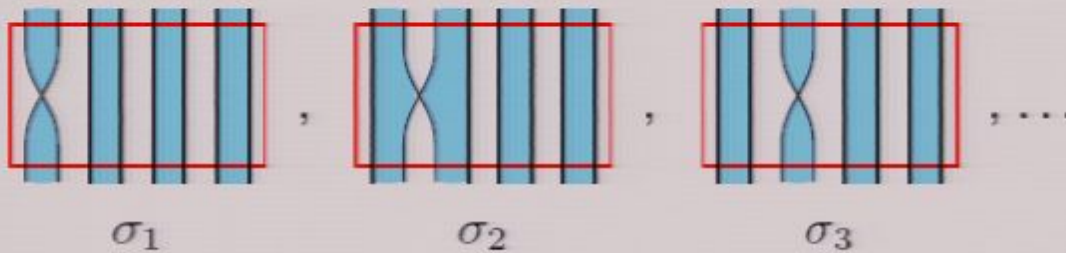


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The 8-steps subspace

The crossing operators on L_8

7 crossing operators $\{\sigma_1, \sigma_2, \dots, \sigma_7\}$ on 8 strands

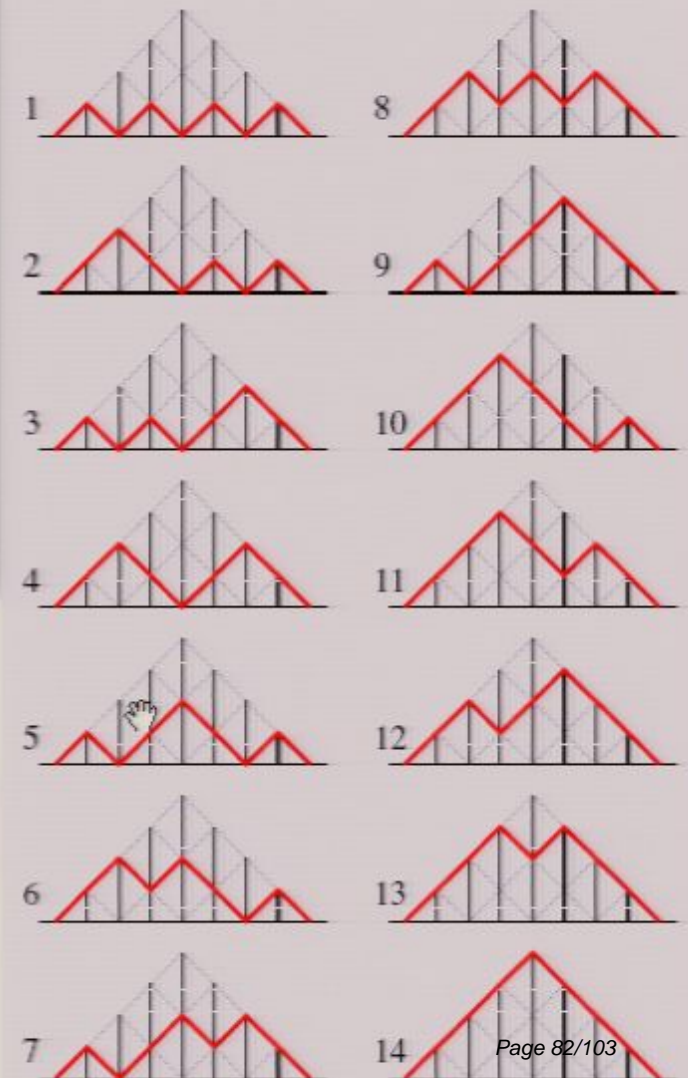


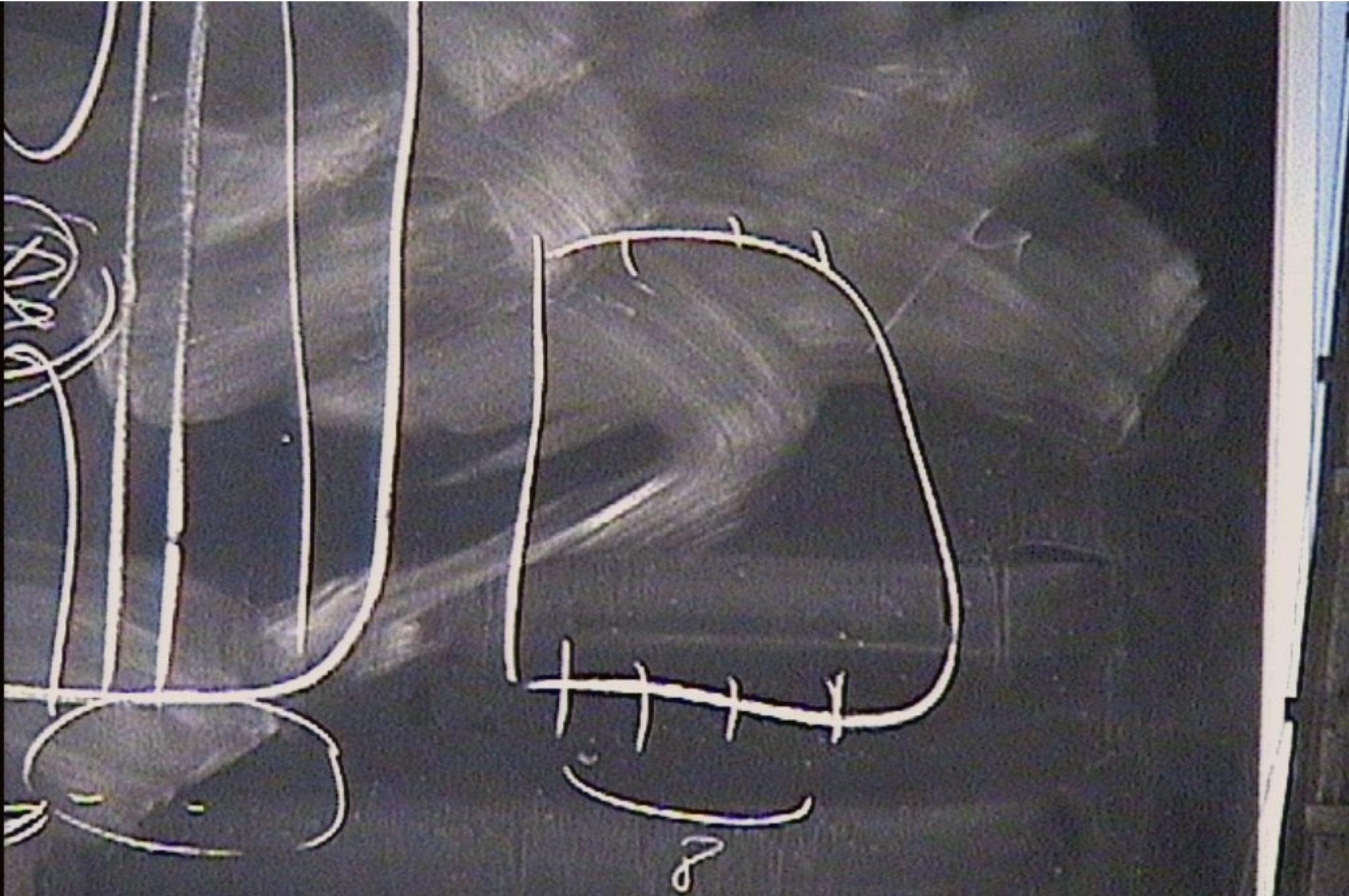
The smallest subspace that includes L_8 and is invariant to $\sigma_1, \dots, \sigma_7$ is

$$H_{8,1 \rightarrow 1} = \text{all 8-steps paths from } 1 \rightarrow 1$$

We need to show density in $SU(14)$ or $SL(14)$

$$\dim(H_{8,1 \rightarrow 1}) = 14$$

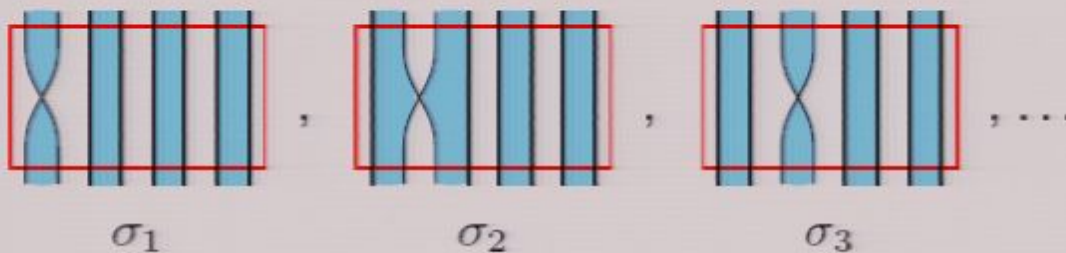




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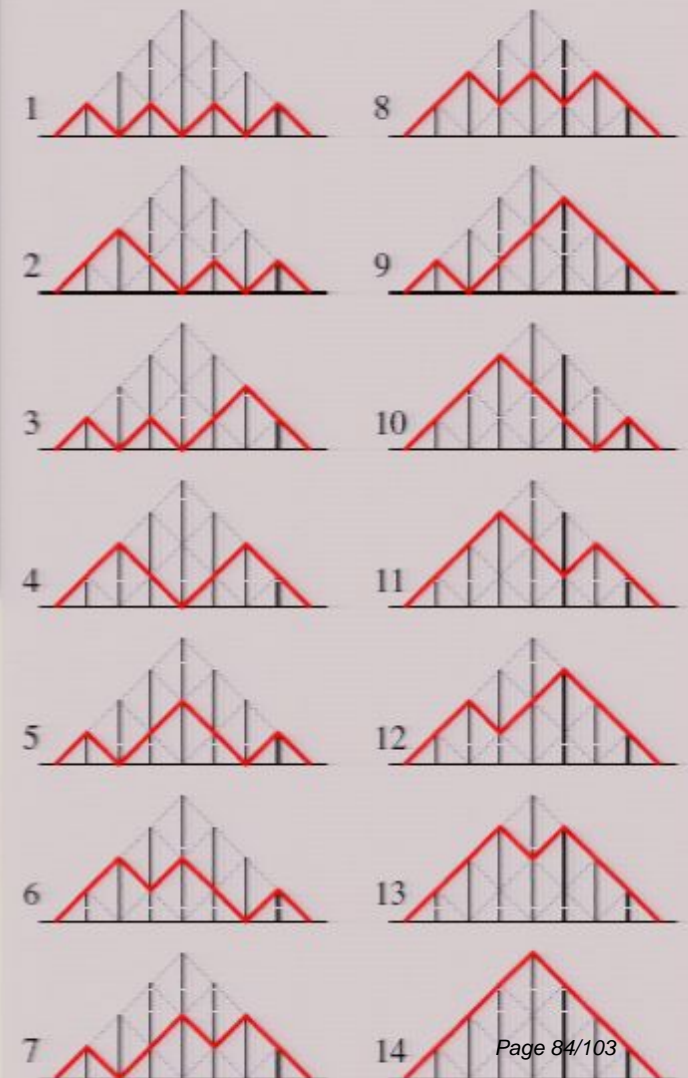


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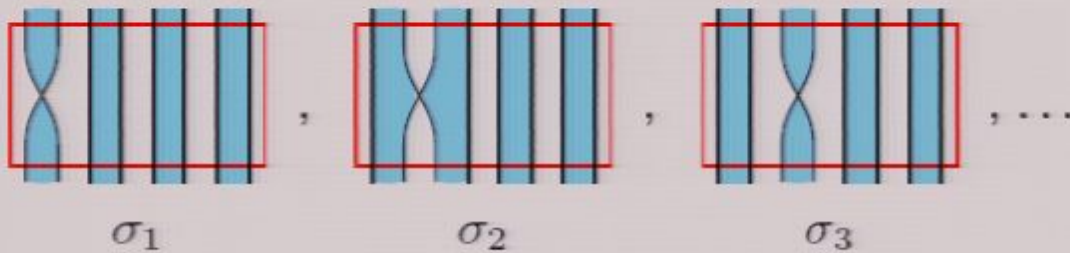
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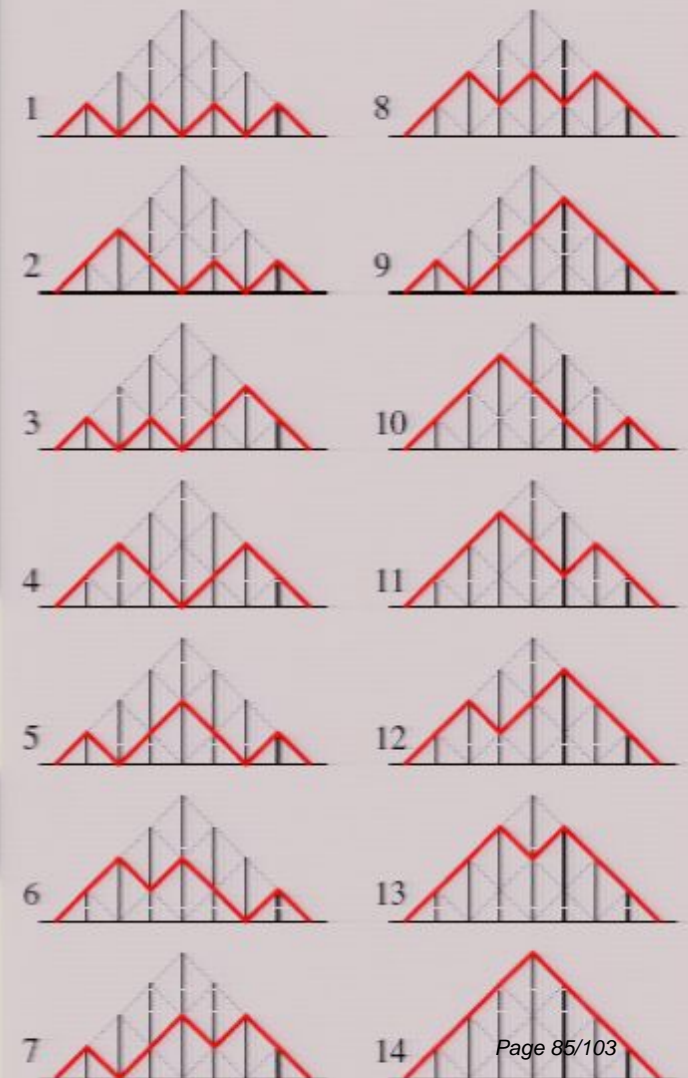


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Building up the density - a tale of 3 lemmas

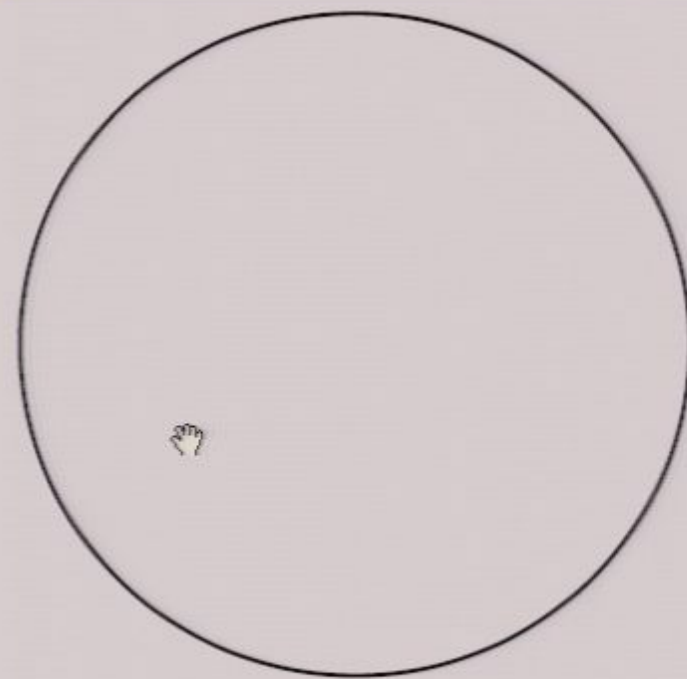
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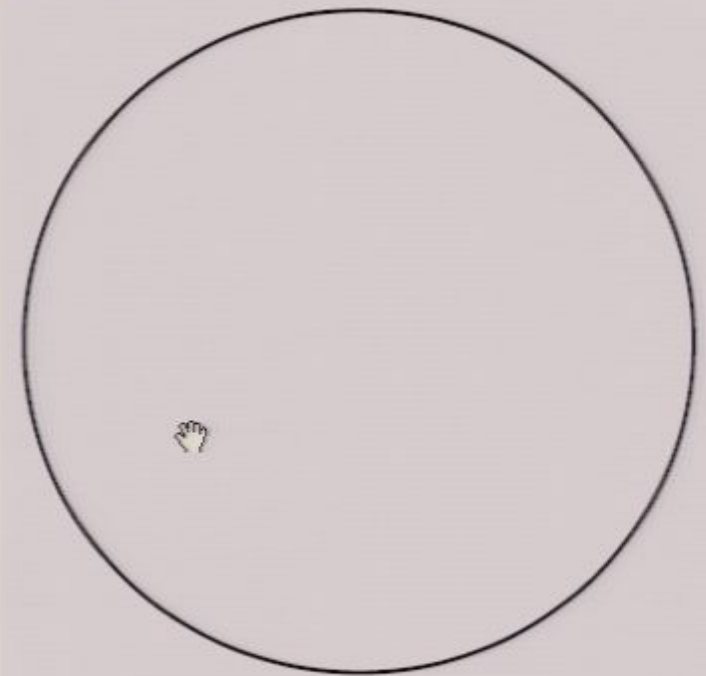
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Proof: Classification of finite groups in $SO(3)$
- **Non-unitary case:** $SL(2)$ is non-compact; **density does not happen most of the times.**
Proof: Möbius transformations + Jørgensen inequality

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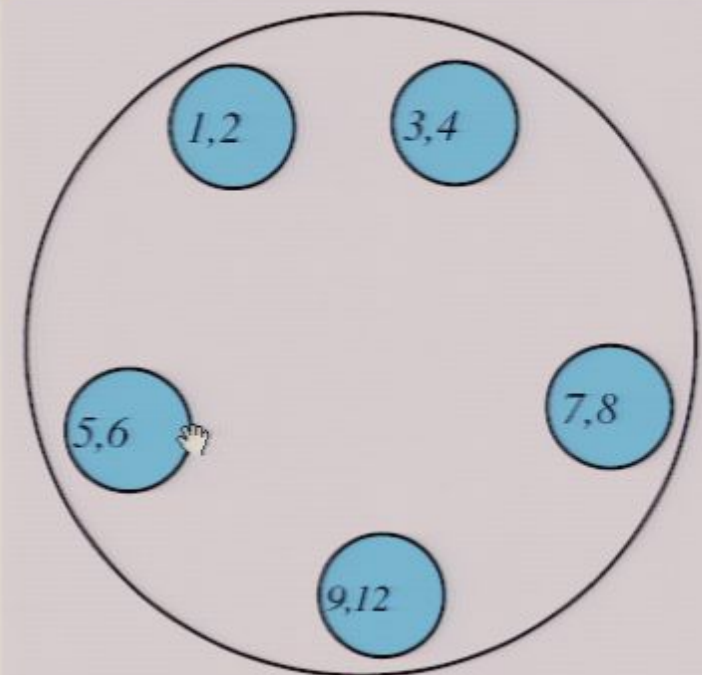
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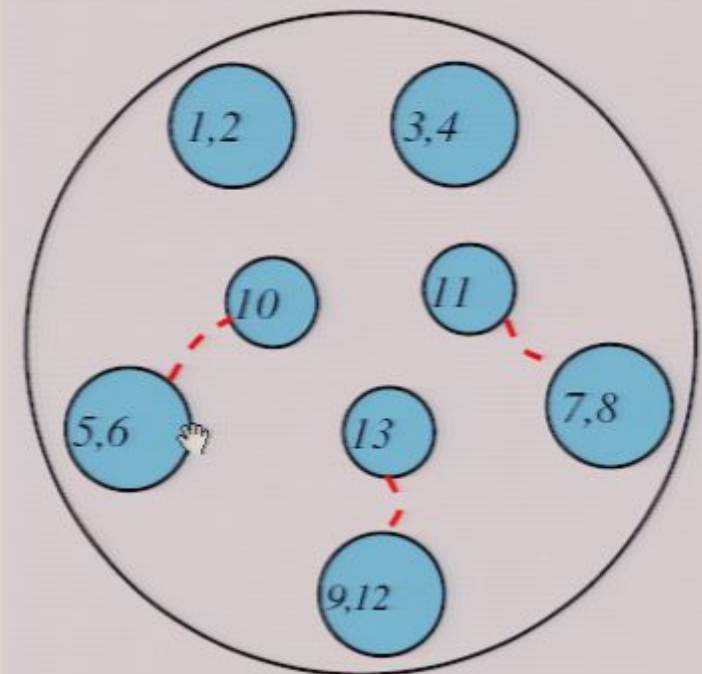
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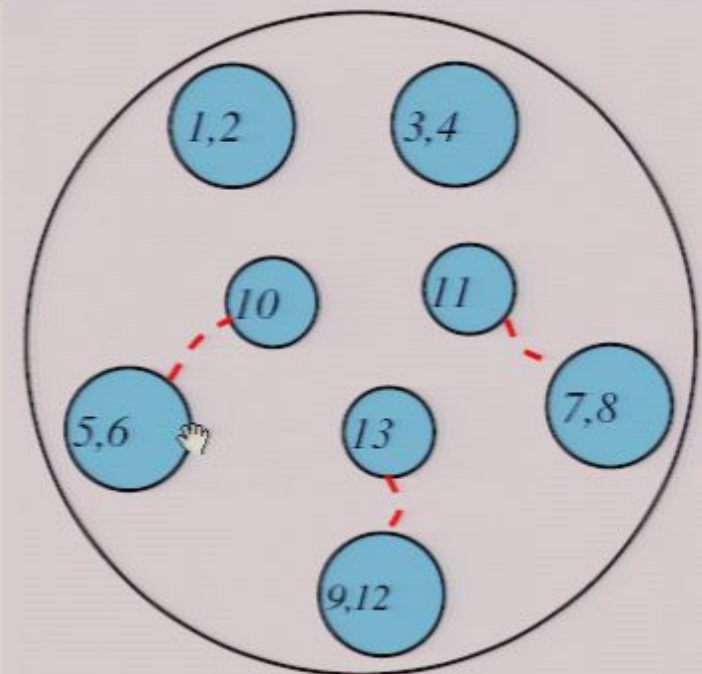
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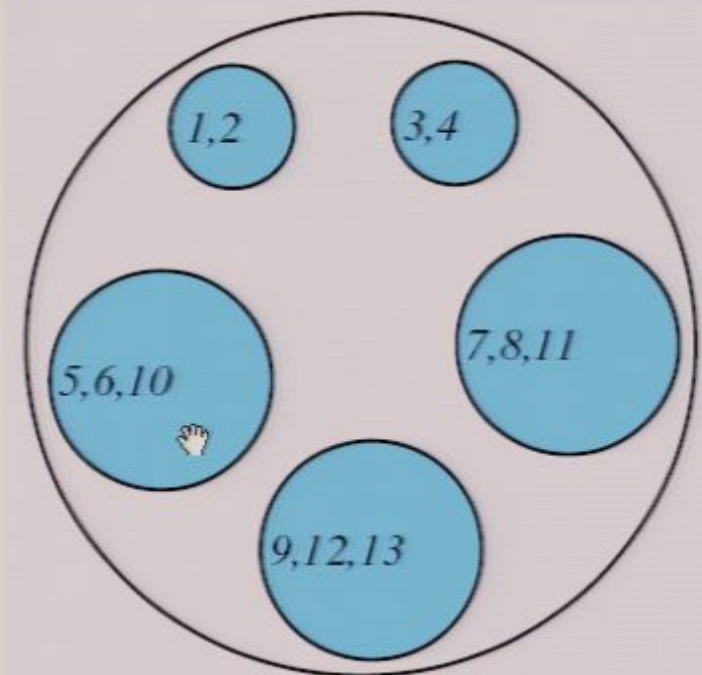
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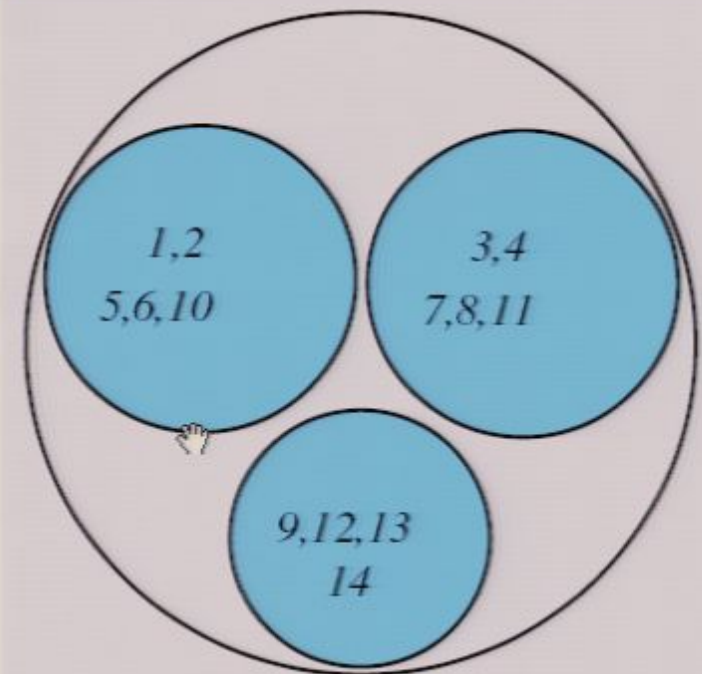
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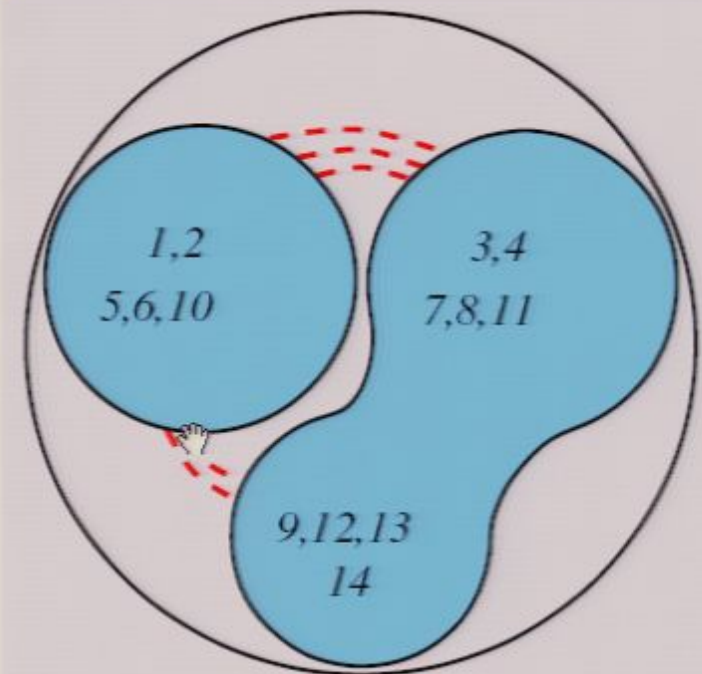
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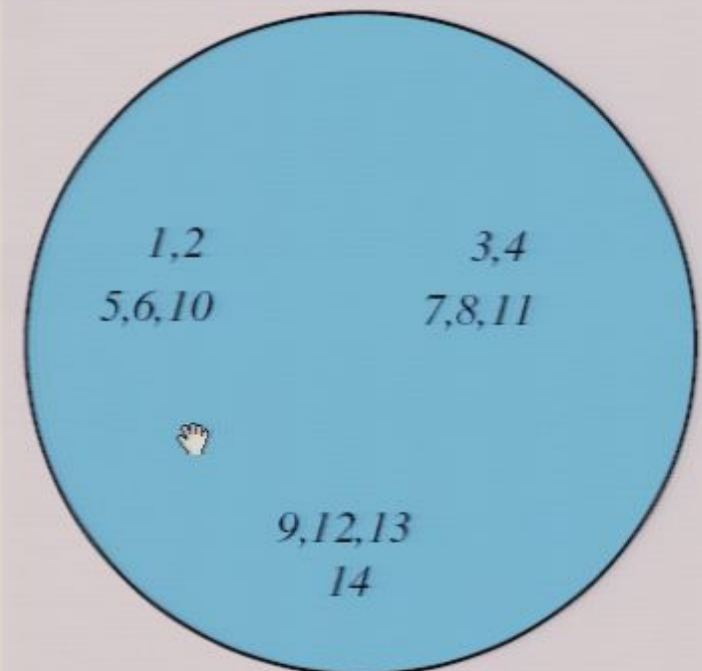
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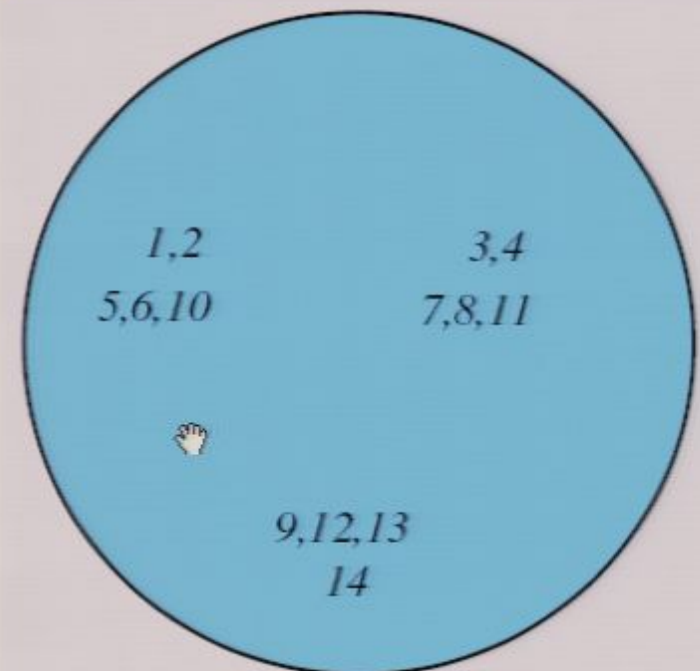
Density \implies efficiency

To finish up use the (non-unitary) Solovay-Kitaev theorem

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$SU(14)$ or $SL(14)$



Summary of universality proof

Universality

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Finding G

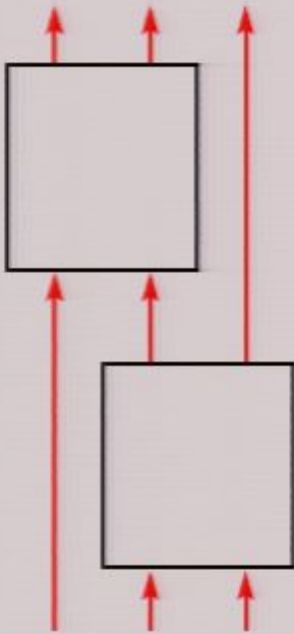


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Finding G



Quantum circuit

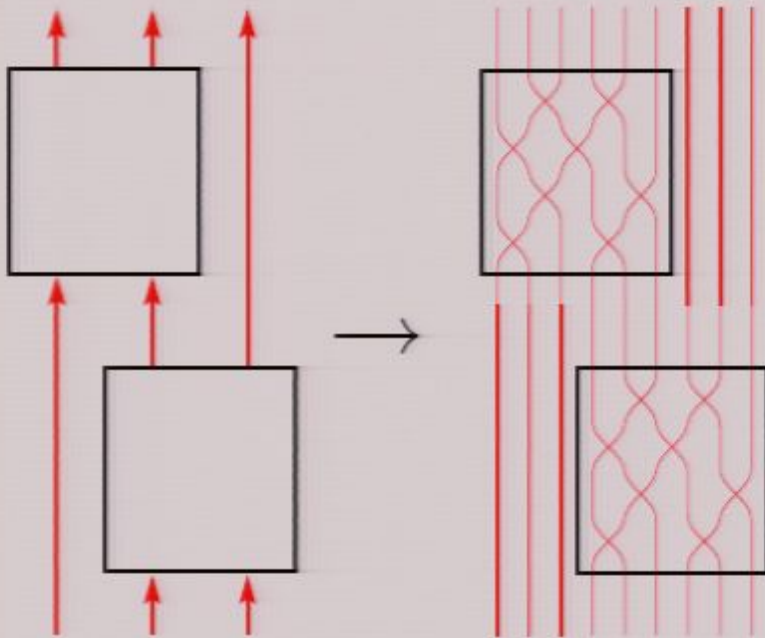
$$U_N \cdots U_1$$

Summary of universality proof

Universality

$$Z_G(q, \mathbf{v}) \simeq \Delta_{hard} \langle 0^{\otimes n} | U | 0^{\otimes n} \rangle$$

Finding G



Quantum circuit
 $U_N \cdots U_1$

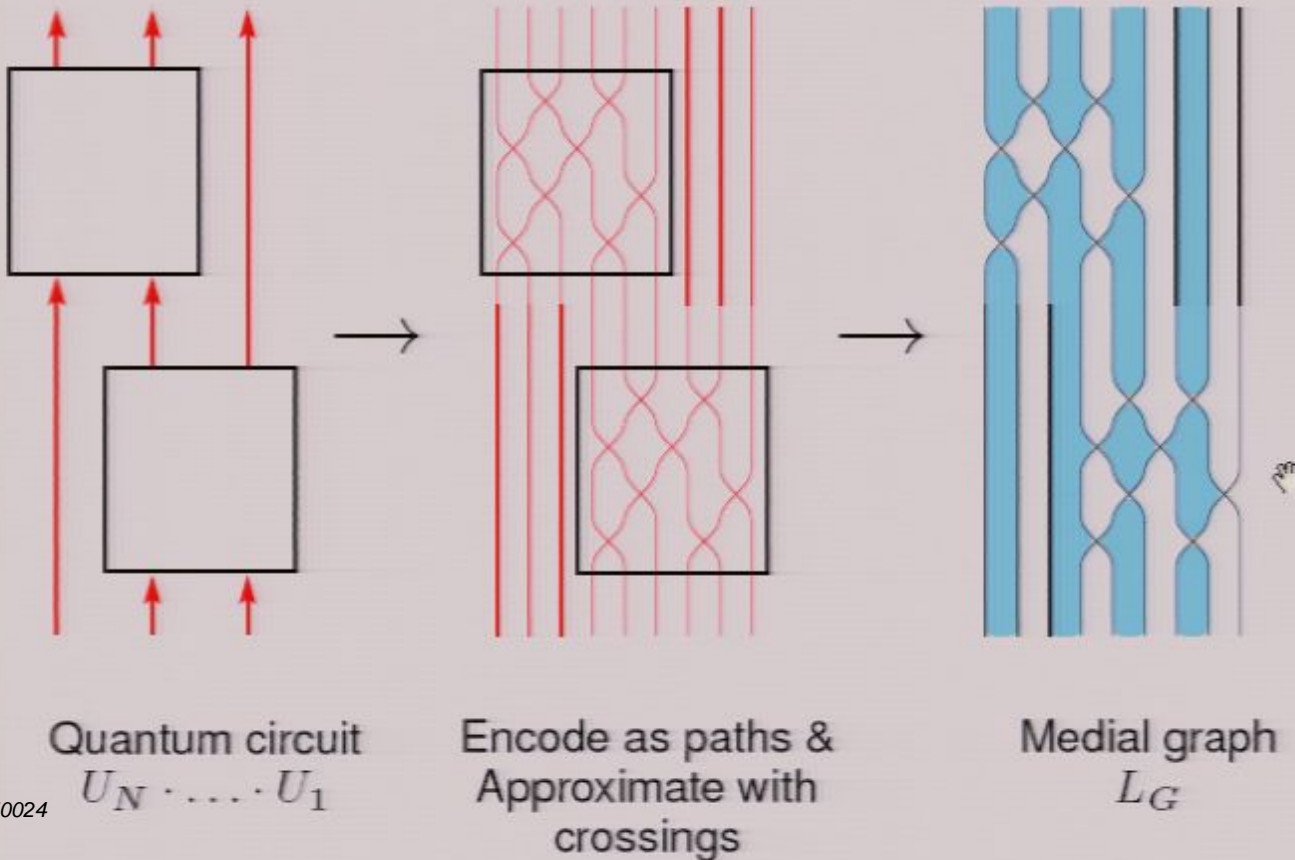
Encode as paths &
 Approximate with
 crossings

Summary of universality proof

Universality

$$Z_G(q, \mathbf{v}) \simeq \Delta_{hard} \langle 0^{\otimes n} | U | 0^{\otimes n} \rangle$$

Finding G

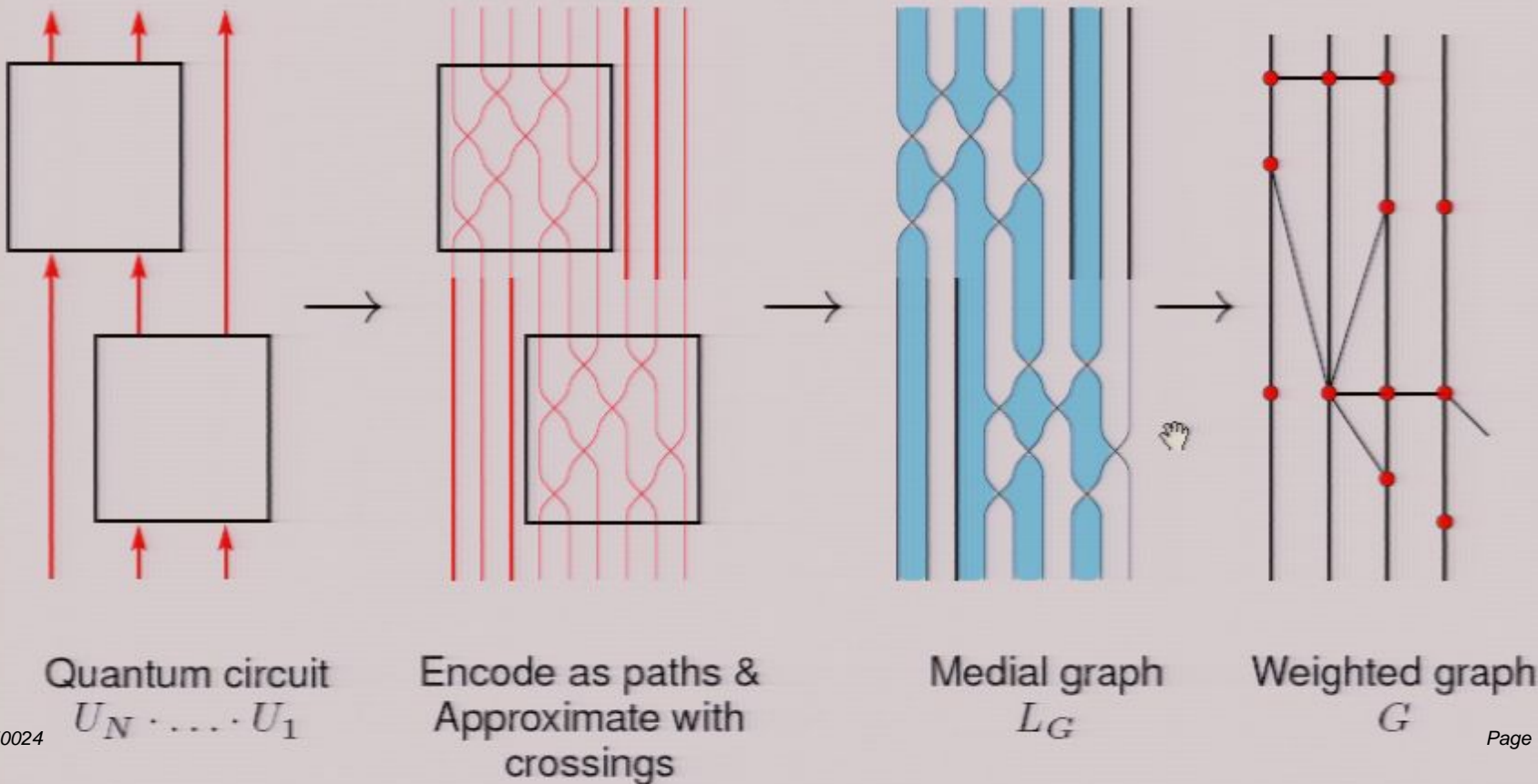


Summary of universality proof

Universality

$$Z_G(q, \mathbf{v}) \simeq \Delta_{hard} \langle 0^{\otimes n} | U | 0^{\otimes n} \rangle$$

Finding G



Summary and open questions

Summary

- An efficient quantum algorithm to approximate the Tutte polynomial for **any** planar graph and **any** set of weights. Uses a local mapping of the combinatorial problem into an **algebra** and implementing on a quantum computer using a **representation**. Many of the cases quantum-complete.
- **Breaking the unitarity barrier**
 - Quantum algorithm that uses non-unitary representations.
 - Quantum universality with non-unitary operators.
- **The Potts model:** possible implications to the long-standing problem of approximating the partition function of the q -state Potts model.

Open questions

- Universality does not apply for the Potts model; what is the complexity of the problem?
- Can it be generalized to non-planar graphs?
- Other combinatorial problems encoded by other algebras?
- Non-unitary universality for quantum: does it have any interesting implications?

$$|Y - f(x)| \leq \epsilon \cdot \Delta_x$$

$$\epsilon \in \frac{1}{\text{Poly}(n)}$$

$$\|T - f(x)\| \leq \epsilon \in (\Delta)$$

$$\epsilon \in \text{Poly}(n)$$

Cóculo $\rightarrow \approx \frac{1}{2}$